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## **RF Gas Plasma Source Development for Heavy Ion Fusion\***

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### *Abstract*

Presently the Heavy Ion Fusion Virtual National Laboratory is researching ion sources and injector concepts to understand how to optimize beam brightness over a range of currents (50-2000 mA argon equivalent). One concept initially accelerates millimeter size, milliamp beamlets to 1 MeV before merging them into centimeter size, ampere beams. Computer simulations have shown the final brightness of the merged beams is dominated by the emittance growth of the merging process, as long as the beamlets' ion temperature is below a few eV. Thus, a RF multicusp source capable of high current density can produce beams with better brightness compared to ones extracted from a colder source with a large aperture and lower current density. As such, experiments have begun to develop a RF multicusp source capable of delivering one amp of extracted beam current. It is expected that it will require 10 kW of 13 MHz RF power delivered via a quartz shielded, one and half turn, four inch diameter antenna. Important considerations in the development of the source include the dependence of current density and beam ion temperature on consumed RF power and gas pressure. A fast rise time (~100 ns) for the extracted beam pulse must also be achieved. Progress on these experiments will be presented.

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### **THE BEAMLET ARCHITECTURE**

The source injector for a heavy ion fusion driver must deliver  $\sim 10^{16}$  ions to the accelerator [1]. Given current limits of sources and transport limits of accelerators, the injector

will almost assuredly have multiple beams. Most present conceptual designs have an injector energy of 1.6 MeV and current per beam of 0.5 A [2]. Presently, the Heavy Ion Fusion Virtual National Laboratory is pursuing a concept of creating the 0.5 A beams by initially extracting many,  $\sim 100$ , small beamlets and accelerating them to 1 MeV before they are allowed to merge together [3]. This concept promises to produce brighter, a key figure of merit, higher current beams and a more compact injector. Initial simulations have given encouraging results.

Another advantage of the beamlet architecture is the beam brightness of the merged beam is dominated by the emittance growth of the merging process and, is only a weak function of the ion temperature of the source. This allows the possibility of using sources with higher ion temperatures instead of surface ionization sources. The merged beams in the beamlet injector only achieve high brightness if the current density of the beamlet is high,  $\sim 100$  mA/cm<sup>2</sup>. In fact, current density becomes a more important criterion for the source than ion temperature.

### **THE RF PLASMA SOURCE**

For the beamlet concept to be practical, there must be a source that delivers high current density and easily accommodates the geometry of many millimeter size beams. A plasma source is such a source. Specifically, the HIF program has started development of a multicusp RF plasma source for HIF. These sources have been used to produce high current density beams and can form the beamlets by using a single RF bucket with a grid for the extraction plate. In fact, K.N. Leung's group at LBNL demonstrated over 200 mA/cm<sup>2</sup> for Ar<sup>+</sup> [4].

Another important requirement for the sources in HIF is the ability of a fast turn on. To minimize non-linear effects in the head of the beam, the rise time of the voltage waveform across the first gap must be  $\sim 100$  ns. Plasma formation times are usually of the order 10  $\mu$ s, implying the necessity for a fast high voltage pulser. Even with such a pulser, a fast rise time

is not guaranteed. The emission surface in a plasma source does change as the voltage pulse ramps up, so a fast beam current pulse may not arise with a fast high voltage pulse. K.N. Leung's group also demonstrated fast rise time,  $\sim 1 \mu\text{s}$ , with a RF plasma source, but with only a few microamps of beam current [5]. Thus, HIF source injector program recently conducted experiments aimed at simultaneously demonstrating fast rise time and high current density

### EXPERIMENTAL SETUP

The experiments were performed with the same 10 cm multicusp source used in reference [4]. The plasma chamber has an inner diameter of 10 cm with 20 SmCo magnet columns. A 1.5-turn, quartz antenna with a 5 cm diameter coil generates plasma inside the chamber. The power is delivered to the antenna from a RF amplifier via a fifty-ohm transmission line and an impedance matching network. The RF amplifier generates a 2 ms burst of 13.56 MHz RF signal at a rep rate of 10 Hz. The maximum peak RF output power of the amplifier is 5 kW. The gas, Argon, is introduced to the plasma chamber through a needle valve and the absolute pressure is measured by a capacitance manometer. A small tungsten starter filament is used to feed seed electrons and typically draw about 60 W of power in order to provide an adequate source of electrons.

The extraction system consists of a plasma electrode with an aperture of 3 mm. The gap to the ground electrode is 3 mm, which has a 4 mm diameter opening for the beam to travel through. A pulser, manufactured by Diversified Electronics, provides the extraction voltage. The system is capable of delivering a 50-kV, 20-microsecond pulse with a peak current of 25 A. A Faraday cup with a magnetic filter was used to measure the beam current extracted from the source.

### RESULTS

First, it was verified that this source could produce  $100 \text{ mA/cm}^2$  of beam current density. The matched condition for the source configuration was determined by varying the

extraction voltage and measure the flattop beam current. Figure 1 shows a plot of current density versus extraction voltage. The data shows a typical voltage to the three-halves power behavior for low voltage. In this region, the Child-Langmuir limit for space charge limited ion extraction is lower than the emission limit from the plasma source. At around 15 kV, the two limits are near equal. Above, this voltage, the extracted current density plateaus, indicating emission limited beam extraction. For this example, the matched condition was determined to be 14.4 kV. This plot clearly shows that 100 mA/cm<sup>2</sup> can be produced by this source.

With the starter filament in place, 2 mT was found to be the minimum pressure for stable operation of the source. Below this pressure, the RF pulse did not always form plasma. Below about 1.0 mT, plasma was never ignited. The RF amplifier was set to 3 kW of peak power for the 2 ms pulse. The extraction pulser was set to fire 500  $\mu$ s after the RF pulse started in order to allow the plasma to stabilize before beam extraction. The extraction voltage was set to 14.4 kV.

Figure 2 shows the resulting extraction voltage waveform and extracted current density. The delay between the two is the time of flight of the argon ions to the Faraday cup. The measure current density is 93.3 mA/cm<sup>2</sup>. This, however, is not a fundamental limit of the setup, but simply the matched condition for this RF power setting and pressure. Figure 3 is the same plot with the time axis expanded around the front of the pulse. Analysis of the waveforms revealed a 2.2 +/- 0.1  $\mu$ s rise time for 10%-90% of the flattop. The extraction pulse also has a similar rise time, 2.3 +/- 0.1  $\mu$ s, indicating that the rise time is being limit by the pusler and not the plasma source itself.

The rise time of the pulser is dominated by the stray capacitance and inductance of the system. One way to significantly reduce the effect of the load is the remove starter filament and the isolation transformer needed to deliver power to it. This was removed from the system and the measurement repeated. Unfortunately, without the starter filament the minimum pressure for stable operation is 18 mT. To get a similar current density, the RF

power was reduced to 2 kW. For these settings, a matched condition was achieved at 15.7 kV. Figure 4 shows the resulting waveforms. For this configuration the rise time of the extraction pulse is  $1.4 \pm 0.1 \mu\text{s}$ . The current pulse again appears to follow the extraction pulse with a rise time of  $1.5 \pm 0.1 \mu\text{s}$ . Again, there is no evidence the plasma source itself is limiting the rise time. Unfortunately, further reduction in the rise time of the extraction pulse would require a different pulser or major changes to the test setup.

### **DISCUSSION**

These results demonstrate it is possible to achieve simultaneously high current density,  $\sim 100 \text{ mA/cm}^2$ , and fast rise time,  $\sim 1 \mu\text{s}$ , with a multicusp RF gas plasma source. The data indicate the rise time is limited by the extraction pulse and not the plasma source itself, though there may be a plasma source limit below a  $1 \mu\text{s}$ . This result is achieved by fast switching of the extraction voltage and not the RF power. A starter filament allows for operation at a lower pressure, 2 mT as compared to 18 mT without. Also only 3 kW of peak RF power is needed to generate current densities of  $100 \text{ mA/cm}^2$ . Even though more power would be needed for a source large enough to generate 100, 5 mA beamlets, a duty factor of less than 0.1 % is needed for the RF power, which implies a low average power system is needed.

### **FUTURE PLANS**

Based on these results, the HIF program has decided to continue development of the plasma source. A larger plasma source capable of delivering 100, 5 mA beamlets is under construction. Figure 5 shows a mechanical drawing of the source that is currently being assembled. The inner diameter of the plasma cavity is 26 cm and it has 38 magnets surrounding the outside. The antenna is a one and two-thirds turn, quartz antenna with a diameter of 10 cm. Based on the results of the previous source, it is estimated that  $\sim 10 \text{ kW}$  of peak RF power is need to create plasma with a high enough density to extract  $100 \text{ mA/cm}^2$ . This RF power will be provided by a pulsed RF system, which can be floated at high voltage.



The source will be tested on the source pulser system of the Recirculator [6]. This pulser system is capable of delivering up to 100 kV and rise times of ~300 ns when driving a 10 nF load. This system may be capable of even faster rise times for loads with less capacitance, as is expected when this source is connected to this platform. This pulser system should allow full scale testing for the first gap of the beamlet architecture. Experiments will be focused on exploring extracted current density and emittance versus RF power and gas pressure. Charge exchange in the extraction gap is an important concern and energy measurements will be performed to determine the magnitude of this effect. Also, charge state purity will be explored with time of flight measurements. A pulsed gas system may also be developed to minimize gas load and thus charge exchange in the extraction gap. Results from the experiments should be available by April of 2002.

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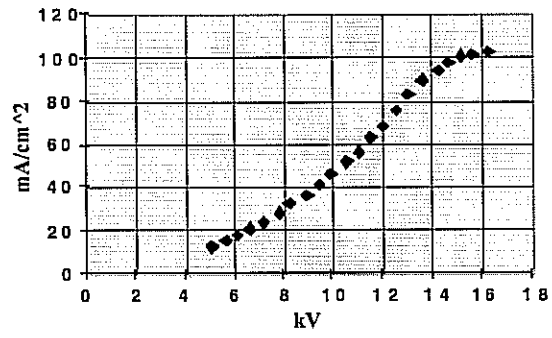


Figure 1: Current density versus extraction voltage.

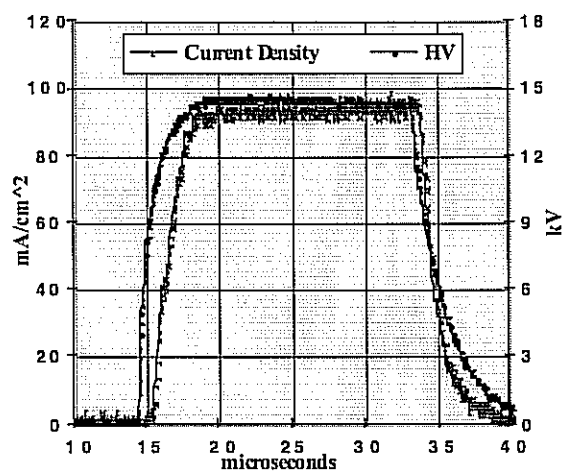


Figure 2: High voltage pulse and extracted current pulse for plasma source with starter filament.

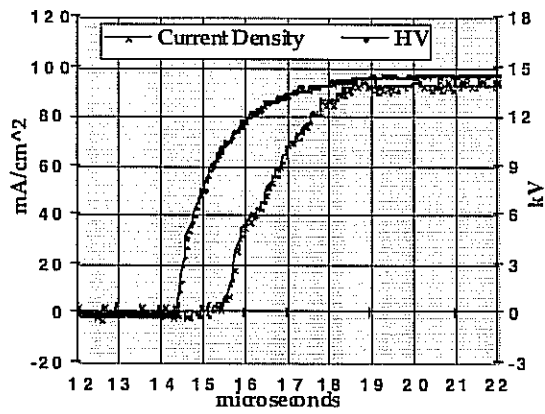


Figure 3: Same as figure 2, only with the time scale expanded around start of pulse.

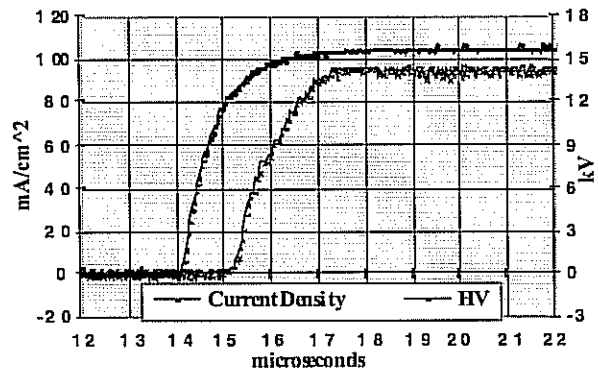


Figure 4: Extraction voltage and current density versus time for plasma source without the starter filament.

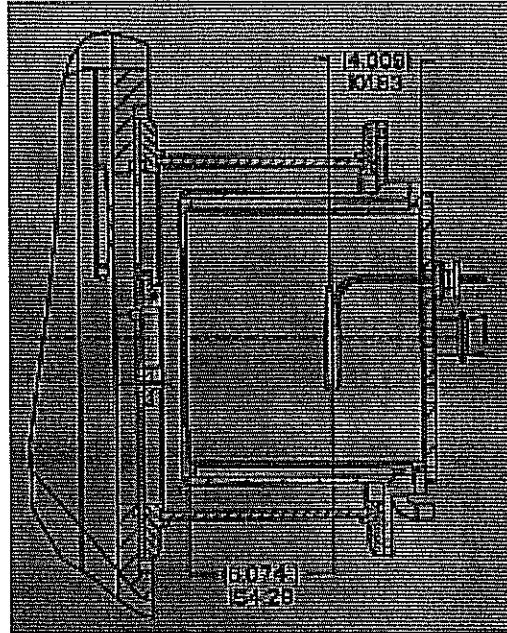


Figure 5: Drawing of plasma source for next series of experiments.

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