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High-quality SnSe₂ single crystal: Electronic and thermoelectric properties

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ABSTRACT

The high-quality SnSe₂ single crystal was successfully synthesized using the temperature gradient method. N-type characteristic and strong anisotropic transport properties of SnSe₂ single crystal were exhibited between ab-plane and c-axis. At 673 K, the power factor (PF) value along ab-plane is $3.43 \mu W cm^{-1} K^{-2}$ while it along c-axis is $0.92 \mu W cm^{-1} K^{-2}$. The ratio between thermal conductivities along ab-plane (κ_{ab}) and along c-axis (κ_c) is 7.6 order at 300 K, while as at 673 K, this value about 5.6. The thermoelectric figure of merit ZT along c-axis (0.15) is higher than that (0.1) along the ab-plane, according to the ultra-low out-of-plane thermal conductivity. The electronic band structure results, which were examined by angle-resolved photoemission spectroscopy (ARPES) predicted the potential of improving thermoelectric performance of SnSe₂ single crystal by electron doping.

Keywords: *Thermoelectric material, single crystal, SnSe₂, ultra-low thermal conductivity, ARPES.*

I. INTRODUCTION

Recently, thermoelectric materials, which directly convert waste heat into electricity, have increasingly attracted attention. The thermoelectric performance is evaluated by the dimensionless thermoelectric figure of merit, $ZT = \frac{S^2\sigma T}{\kappa}$, where S , σ , T , and κ represent the Seebeck coefficient, electrical conductivity, absolute temperature, and thermal conductivity, respectively^{1,2}. To reach excellent performance as a commercial product, S and σ should stay at high value, while as κ remains low. Since these three transport coefficients are inter-dependent, it is difficult to optimize ZT value. There are typical ways to improve ZT value such as enhancing the power factor $S^2\sigma$ via optimizing the carrier concentration n and mobility μ or suppressing the lattice thermal conductivity κ_L by introducing the scattering centers. As a result, there are currently only a few materials used for commercial applications.

In layered structure materials, the difference in bonding mechanism between in-plane and out-of-plane leads to unique electronic and thermal transport properties. Among them, layered chalcogenide crystals possessing strong intra-layer covalent bonding and noticeable weak inter-layer van der Waals bonding, such as Bi_2Te_3 and Sb_2Te_3 , are by far recognized as good thermoelectric materials at around room temperature. In addition, layered SnSe with the weak cross-plane covalent bonds, has been reported with in-plane $ZT = 2.6$ at 923 K in p-type SnSe^3 , and $ZT = 2.2$ at 733 K in Bi-doped n-type SnSe^4 ; out-of-plane $ZT = 2.8$ at 773 K in Br-doped n-type SnSe^5 .

In the layered structure materials family, Tin diselenide (SnSe_2) possesses the CdI_2 -type hexagonal structure with space group $P\bar{3}m1$, which is formed by a van der Waals bonding between Sn-Se-Sn slabs along out-of-plane direction (c-axis). Lattice parameters of SnSe_2 are $a = 3.811 \text{ \AA}$ and $c = 6.137 \text{ \AA}^6$. SnSe_2 has yet confirmed belonging to n-type semiconductor with indirect band gap of 0.97 eV⁷, which was considered as applications such as phase-change memory, optoelectronic devices, field-effect transistors (FETs), and gas sensor devices⁸⁻¹². Recently, SnSe_2 has attracted interest from scientists due to its ultralow thermal conductivity. Some theoretical calculations have predicted that thermoelectric properties of SnSe_2

1 are significantly influenced by carrier concentration. ZT value of p-type SnSe₂ monolayer can be obtained
2 up to 0.94 at 600 K with carrier concentration around 10¹⁹ cm⁻³, while that of n-type SnSe₂ is around 0.8
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4 with carrier concentration around 10²⁰ cm⁻³ ¹³. Ding *et al.* predicted that at 800 K, ZT values of n-type
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6 SnSe₂ are 0.01, 0.2, 1.1, and 2.95 with carrier concentrations of 10¹⁷, 10¹⁸, 10¹⁹, and 10²⁰ cm⁻³,
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8 respectively¹⁴. By doping Cl, Xu *et al.* achieved $ZT = 0.4$ at 673 K in SnSe_{1.88}Cl_{0.12} polycrystals¹⁵. By
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10 doping Ag, Li *et al.* also obtained $ZT = 0.4$ in Sn_{0.99}Ag_{0.01}Se₂ at 773 K¹⁶. By substituted Br into Se sites of
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12 SnSe₂, Wu *et al.* improved ZT to 0.6 at 750 K¹⁷.

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18 Based on potential applications of SnSe₂ crystal, the synthesis of high-quality SnSe₂ single crystal is
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20 beneficial in studying the intrinsic thermoelectric properties. In this article, high crystalline quality SnSe₂
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22 has been grown by simple temperature gradient method and has been studied anisotropic thermoelectric
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24 transport properties. The thermoelectric performance of SnSe₂ single crystal shows better than the results
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26 on polycrystalline. Besides, the ARPES results show the potential for improving thermoelectric
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28 performance by electron doping.
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31 **II. EXPERIMENT**

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34 SnSe₂ single crystal was grown by the temperature gradient method¹⁸. First, high purity (5N) Se and Sn
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36 powders with the mole ratio of 2:1 were weighed and sealed in quartz ampoule under vacuum (<10⁻⁴ Torr).
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38 Then, the ampoule was heated to 750 °C (60 °C h⁻¹) and soaked for 10 hours. Finally, the temperature was
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40 gradually cooled with a rate of 1 °C h⁻¹ to 550 °C, then rapidly dropped to room temperature (50 °C h⁻¹).
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42 A high-quality SnSe₂ crystal with sizes of 14 mm diameter × 30 mm length was obtained. The crystal
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44 structure was studied by means of X-ray diffraction (XRD). Field Emission Scanning Electron
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46 Microscopy (FE-SEM) and Energy Dispersive X-ray Spectroscopy (EDS) were used to observe the layer
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48 morphology and stoichiometry of the samples, respectively.
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53 To study the thermoelectric properties in both directions, parallel (ab-plane) and perpendicular (c-
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55 axis) to the layered plane, bar-shaped samples used to measure the Seebeck coefficient and electrical
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conductivity, and disk-shaped samples used to measure thermal diffusivity were cut into a precise size $3 \times 3 \times 10 \text{ mm}^3$ and $13 \times 13 \times 1 \text{ mm}^3$, respectively. Seebeck coefficient and electrical conductivity were determined simultaneously in an Argon environment to prevent oxidation and evaporation. Hall measurement was studied in various temperatures from 300 to 673 K via the van der Pauw method under a reversible magnetic field of 0.7 T. The laser flash diffusivity method (LFA-457, NETZSCH, Germany) was used to evaluate the thermal diffusivity D . The heat capacity was determined from the measured values of Wiedemeier *et al*¹⁹. by $C_p = 73.39 + 1.15 \times 10^{-2}T - 1.92 \times 10^{-5}T^{-2} \text{ (J K}^{-1} \text{ mol}^{-1}\text{)}$ with the uncertainty of 1% for SnSe₂. The total thermal conductivity was obtained following the formula $\kappa_{\text{tot}} = D.C_p.\rho$, with ρ is the mass density measured by the Archimedes' principle at 300 K. The lattice thermal conductivity κ_L was obtained by subtracting the electronic contribution ($\kappa_e = LT/\rho$) from the total thermal conductivity κ_{tot} , where the Lorenz number²⁰ ($L \sim [1.5 + \exp(-\frac{|S|}{116})] \times 10^{-8} \text{ W } \Omega \text{ K}^{-2}$) was estimated from the Seebeck coefficient data. Transport property measurements (S , σ , and κ) were conducted from room temperature to 673 K, with the uncertainty of each being $\sim 5\%$. Angle-resolved photoemission spectroscopy (ARPES) was performed at the Beamline 10.0.1, Advanced Light Source, using photon energy 55 eV. The measurements were made under base pressure below 4×10^{-11} Torr at 30 K.

III. RESULTS AND DISCUSSION

By using temperature gradient method, we have been successfully grown high-quality SnSe₂ single crystals. Figure 1a describes the real image of the grown samples with cleaved shiny surfaces. EDS measurement indicates the proportion of Se and Sn to be 2:1, as shown in Figure 1b. The surface of cleaved SnSe₂ single crystal was observed by FE-SEM. The results exhibited a clear surface and a lamellar microstructure with an average thickness of a few tens nanometer, as shown in Figure 1c. For STM (Scanning Tunneling Microscopy) studies, the sample was cleaved in-situ to obtain clean surfaces of

1 SnSe₂ single crystal. High resolution STM image in Figure 1d proved the appearance of hexagonal
 2 structure of Se on the surface of SnSe₂.
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 6 The XRD result of cleaved SnSe₂ single crystal without any impure peaks has been shown in Figure
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 8 2a. Compared with JCPDS PDF 89-3197 data, all diffraction peaks belong to the hexagonal SnSe₂
 9 structure and the patterns showed only (00*l*) peak group. Lattice parameters of SnSe₂ single crystal were
 10 calculated from powder XRD data (as shown in Figure 2b) with a *CuKα*₁ radiation; a = 3.808 Å and c =
 11 6.129 Å are comparable with earlier reported. The calculated mass density of the grown SnSe₂ single
 12 crystal by lattice parameters is 5.966 g/cm³, which agreed well with that of the directly measured one,
 13 5.937 ± 0.11 g/cm³.
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22 Figure 3a represents the temperature dependence of carrier concentration $n(T)$ which is obtained from
 23 Hall measurement by using formula $\frac{V_H}{I} = \frac{1}{ned}H$, where V_H is Hall voltage, I is the current, n is the number
 24 of carriers, e is electrical charge, H is magnetic field, and d is sample thickness. As increasing temperature,
 25 n gradually increases from $2.26 \times 10^{18} \text{ cm}^{-3}$ at 300 K to $3.37 \times 10^{18} \text{ cm}^{-3}$ at 473 K and then
 26 dramatically increases up to $3.05 \times 10^{19} \text{ cm}^{-3}$ at 673 K. The number of carrier at 300 K is one order of
 27 magnitude higher than that of the previous reported single crystal by Julien *et al.* prepared by the Bridgman
 28 technique, while the electron mobility along ab-plane, $\mu = 31.6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ (Figure 3b), is quite
 29 similar⁶. The mobility along both directions decreases sharply with temperature and obeys the trend of
 30 $T^{-3/2}$, pointing out that the acoustic phonon scattering is the major contribution. These results are
 31 reasonable compared with previous works on single crystals and polycrystals of SnSe₂. However, the
 32 mobility value in single crystal is much higher than those in polycrystalline samples^{6,21–23}. Thermoelectric
 33 properties of SnSe₂ single crystal have been measured along cleaved surface (ab-plane) and perpendicular
 34 cleaved surface (c-axis) from 300 to 673 K, as shown in Figure 4. Electrical conductivity showed
 35 semiconducting behavior. However, measurement results along cleaved surface and perpendicular cleaved
 36 surface are different due to anisotropic transport properties of this material. Especially at above 500 K,
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1 electrical conductivity along the cleaved plane (σ_{ab}) increases up to 35.97 S cm^{-1} at 673 K caused by the
 2 rapid rise of carrier concentration. Seebeck coefficient along and perpendicular to cleaved surface follows
 3 a similar trend, as shown in Figure 4b. From room temperature to 500 K , SnSe_2 single crystal exhibited
 4 the metallic transport behavior. This transport behavior can also be observed in Figure 3b when the carrier
 5 mobility forms a sharp degradation, in contrast, the carrier concentration slightly increases along with the
 6 temperature. As generally expected, Seebeck coefficients slowly rise with the temperature, which
 7 corresponds to the fall of electrical conductivity. At above 500 K , sample shows thermally activated
 8 transport behavior. So, both electrical conductivity and Seebeck coefficient show an opposite trend
 9 compared to those in the lower temperatures. Figure 4c shows the thermoelectric power factor ($PF = S^2\sigma$
 10) along both directions. Maximum PF along cleaved plane ($3.47 \mu\text{W cm}^{-1} \text{ K}^{-2}$) is larger than that along c-
 11 axis direction ($0.97 \mu\text{W cm}^{-1} \text{ K}^{-2}$) owing to the relatively big difference in electrical conductivities between
 12 the two directions. Furthermore, its anisotropy becomes stronger with temperature. Because of the
 13 limitation of material thermal stability, it is impractical to carry out the experiment at higher temperatures.
 14 Therefore, peak electrical conductivities are not obtained in this temperature range, and the highest value
 15 is around 35.97 S cm^{-1} at 673 K .

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 36 Temperature-dependent total thermal conductivities (κ_{tot}) along and perpendicular cleaved surface are
 37 shown in Figure 4d. The data revealed the lattice thermal conductivity accounts for 98 % of κ_{tot} ; phonon
 38 transport contributed to most total heat transfer in SnSe_2 . From the inset, the lattice thermal conductivities
 39 (in-plane & out-of-plane) decrease as T^{-1} along with temperature, this could result from the anharmonic
 40 phonon-phonon interactions. In addition, the obtained in-plane thermal conductivity (κ_{ab}) is significantly
 41 larger than that out of plane one (κ_c) at all temperatures. At 300 K , we found $\kappa_{ab} = 6.9 \text{ W m}^{-1} \text{ K}^{-1}$ which is
 42 nearly eight times higher than $\kappa_c = 0.9 \text{ W m}^{-1} \text{ K}^{-1}$. The ratio $\kappa_{ab}/\kappa_c = 7.6$ indicates strong anisotropy in
 43 thermal conductivity of SnSe_2 . Moreover, at 673 K , the out of plane κ_c even drops to an ultralow value of
 44 $0.43 \text{ W m}^{-1} \text{ K}^{-1}$. The large anisotropy in thermal transport properties is well explained from crystal
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1 structure and bonding characters of SnSe₂. This is a consequence of ultraweak van der Waals forces among
 2 layers compared with the intralayer covalent bonding. These values agree well with both previous
 3 experimental²¹ and theoretical studies^{22,23}.
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8 *ZT* values along and perpendicular cleaved surface have been shown in Figure 5. *ZT* values of both
 9 directions increase with temperature. Ultra-low out-of-plane thermal conductivity is the cause of higher
 10 *ZT_c* value compared to that along ab-plane *ZT_{ab}*. We determine the maximum *ZT* value of SnSe₂ single
 11 crystal along c-axis direction is 0.15 at 673 K, while that along ab-plane is 0.1. Note that, in previous
 12 experiment reports, *ZT* value of un-doped SnSe₂ polycrystal is usually smaller 0.05^{9,10}. In two theoretical
 13 calculations by Li *et al.* and Ding *et al.*, *ZT* value of monolayer and single crystal of SnSe₂ are strongly
 14 dependent on carrier concentration. In n-type SnSe₂ monolayer, *ZT* value at 900 K with electron concentration
 15 of 10¹⁸ cm⁻³ is 0.1 and it increases up to 1.1 with *n* of 10²⁰ cm⁻³. In the SnSe₂ single crystal, *ZT* value at 800
 16 K is 0.15 with *n* ≈ 10¹⁸ cm⁻³ and increases up to 2.95 with *n* ≈ 10²⁰ cm⁻³ ^{14,27}.
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29 To study the relation between *ZT* and electronic band structures, we analyze the electronic band
 30 structure of SnSe₂ single crystals using ARPES. Overall band structure obtained from ARPES experiments
 31 is consistent with the reported theoretical results^{28,29}. The position of the valence band maximum (VBM)
 32 of pristine SnSe₂ is 1.36 eV below Fermi energy (*E_F*) with no photoemission intensity near *E_F*, implying
 33 that SnSe₂ is an intrinsic *n*-type semiconductor (Figure 6a, b). The constant energy contour at VBM shows
 34 a flower-like shape, indicating three-fold rotational symmetry of SnSe₂. By potassium (K) doping, SnSe₂
 35 is electron-doped so that conduction bands are clearly revealed at the M point (Figure 6c), which is
 36 comparable with the earlier calculation result²⁸. The extracted effective masses from the conduction band
 37 of SnSe₂ are $m^*_{\Gamma-M} = (2.63 \pm 0.12)m_e$ and $m^*_{K-M} = (1.56 \pm 0.10)m_e$ (Figure 6d), which are
 38 similar to the theoretical value³⁰. Note that the ARPES measurement is performed at 30 K to illustrate the
 39 band structure clearly, while most of the discussion in literatures are performed at 300 – 800 K where the
 40 thermoelectric applications are most relevant²⁶. Therefore, it is better to focus on providing useful
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1 information on how the added electrons would populate conduction band by potassium doping on the
 2 surface rather than looking for direct signature of conduction band in our low temperature ARPES results.
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 6 To achieve high ZT of a thermoelectric material, band engineering is proposed as an effective
 7 strategy, but the resulting enhancement is often subtle due to the competition among the factors
 8 constituting the ZT value. ZT is determined by electrical conductivity, Seebeck coefficient, and thermal
 9 conductivity which depend on charge carrier concentration (n) in a conflicting manner suggesting that
 10 optimizing n is a significant strategy for enhancing ZT . For SnSe_2 , ZT maximizes when n is around 10^{20}
 11 cm^{-3} which is higher than the n of pristine SnSe_2 ($\sim 10^{18} \text{ cm}^{-3}$)²⁶. Our ARPES data clearly shows that
 12 electron doping fills up the conduction bands rather than forming the in-gap states resulting in the metallic
 13 band character (Figure 6c, d)), which is in contrast to the semiconducting band character of pristine SnSe_2
 14 (Figure 6a, b). This proposes that thermoelectric properties of tin diselenide can be improved by enhanced
 15 n yielding the changes in the band character from semiconducting with $n \sim 10^{18} \text{ cm}^{-3}$ to metallic with n
 16 $\sim 10^{20} \text{ cm}^{-3}$, induced by electron doping.
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31 According to the calculations of the partial and total density of state (DOS) for SnSe_2 ³¹, the Sn $4p$
 32 orbitals mainly contribute to the lowest conduction bands, in which the contribution of Se $4p_z$ is less than
 33 that of Se $4p_x$ and $4p_y$. This leads to the superiority of in-plane electrical conductivity compared to the
 34 out-of-plane direction. However, with increasing carrier concentration by heavy donor doping, the carriers
 35 will fill Se $4p_z$ partially when Se $4p_x$ and $4p_y$ are already filled, leading to a steady increase in out-of-plane
 36 electrical conductivity. Thus, theoretical calculations support our higher ZT_c of 0.15 along the c -axis,
 37 compared to ZT_{ab} of 0.1 in the ab -plane, indicating better out-of-plane thermoelectric properties. We
 38 expect the significant boost of thermoelectric performance in out-of-plane ZT_c , when SnSe_2 can be heavily
 39 electron-doped.
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51 52 **IV. CONCLUSIONS** 53 54 55 56 57 58 59 60

1 In summary, we synthesized a high-quality SnSe₂ single crystal using the temperature gradient method.
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3 The bulk material showed strong anisotropy in both electrical and thermal conductivity. An ultralow
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5 thermal conductivity $\kappa_c = 0.43 \text{ W m}^{-1} \text{ K}^{-1}$ was obtained at 673 K, which creates a favorable condition to
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7 enhance thermoelectric performance. The out-of-plane thermoelectric figure of merit (ZT_c) value is 0.15
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9 at 673 K, which is higher than that in-plane because of the lower out-of-plane thermal conductivity of
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11 SnSe₂ single crystal. ARPES data predicts that the power factor can be significantly improved via
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13 electronic doping.
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20 **Author contributions**

21 The manuscript was written through contributions of all authors. All authors have given approval to the
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23 final version of the manuscript.
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Figure captions

Figure 1. (a) Photo of samples, (b) EDS measurement data, (c) FE-SEM image, and (d) STM image of SnSe₂ single crystal.

Figure 2. Room temperature XRD patterns of (a) cleaved and (b) powdered SnSe₂.

Figure 3. Temperature dependence of (a) carrier concentrations and (b) carrier mobility of SnSe₂ single crystal.

Figure 4. Temperature-dependent thermoelectric transport properties along both directions. (a) Electrical resistivity $\sigma(T)$, (b) Seebeck coefficient $S(T)$, (c) power factor $S^2\sigma$, and (d) thermal conductivity $\kappa(T)$.

Figure 5. Temperature-dependent dimensionless figure of merit ZT of SnSe₂ along both directions.

Figure 6. The electronic band structure of pristine and potassium (K) doped SnSe₂ single crystal. (a) Constant energy contour of pristine SnSe₂ for the binding energy of -1.36 eV. (b) ARPES spectra of pristine SnSe₂ along Γ -M direction with VBM at 1.36 eV below E_F . The blue dotted lines are the calculated band structures (DFT) adapted from ref. 26, Copyright © 2019 AIP Publishing. (c) Fermi surface of K doped SnSe₂ with conduction bands at M points. (d) Electronic structures along the momentum directions, A and B on (c). Red dots are the peaks of EDCs fitted by Gaussian function and orange curves indicate the fitted parabolic curves near the band bottom to extract the effective mass.

FIG. 1, Pham et al.

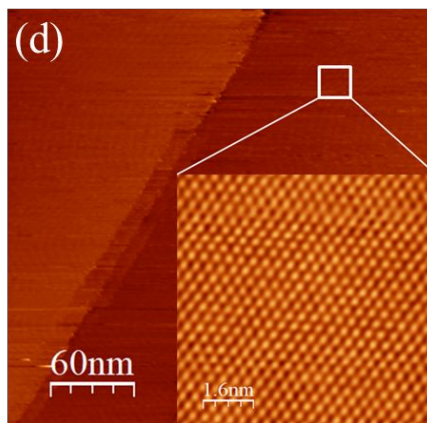
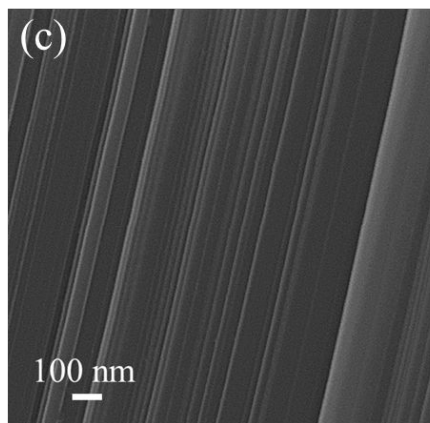
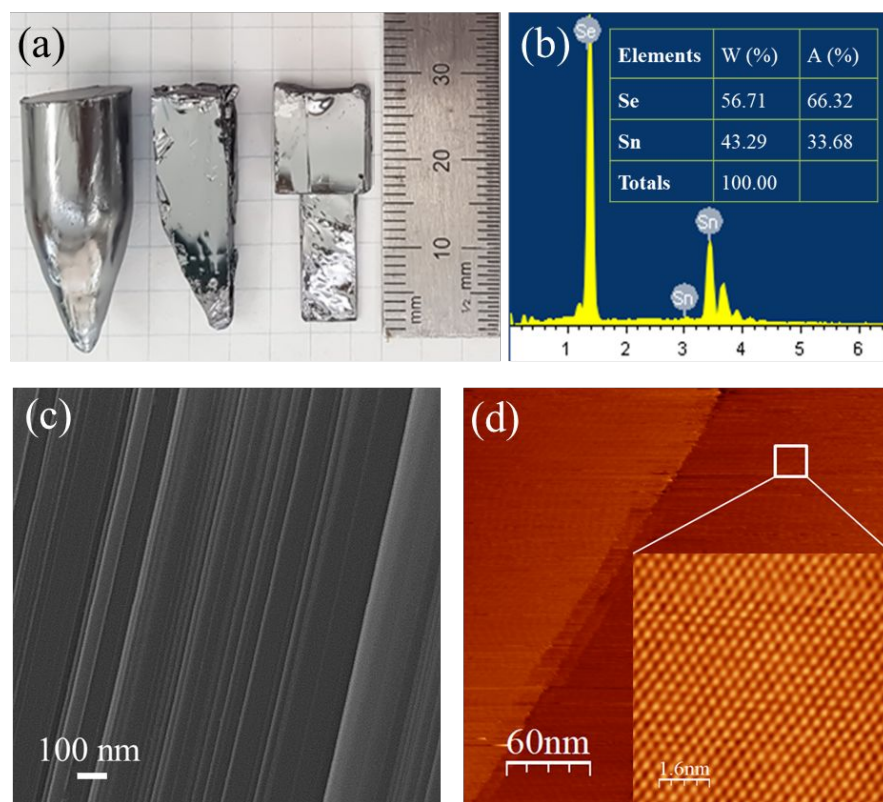
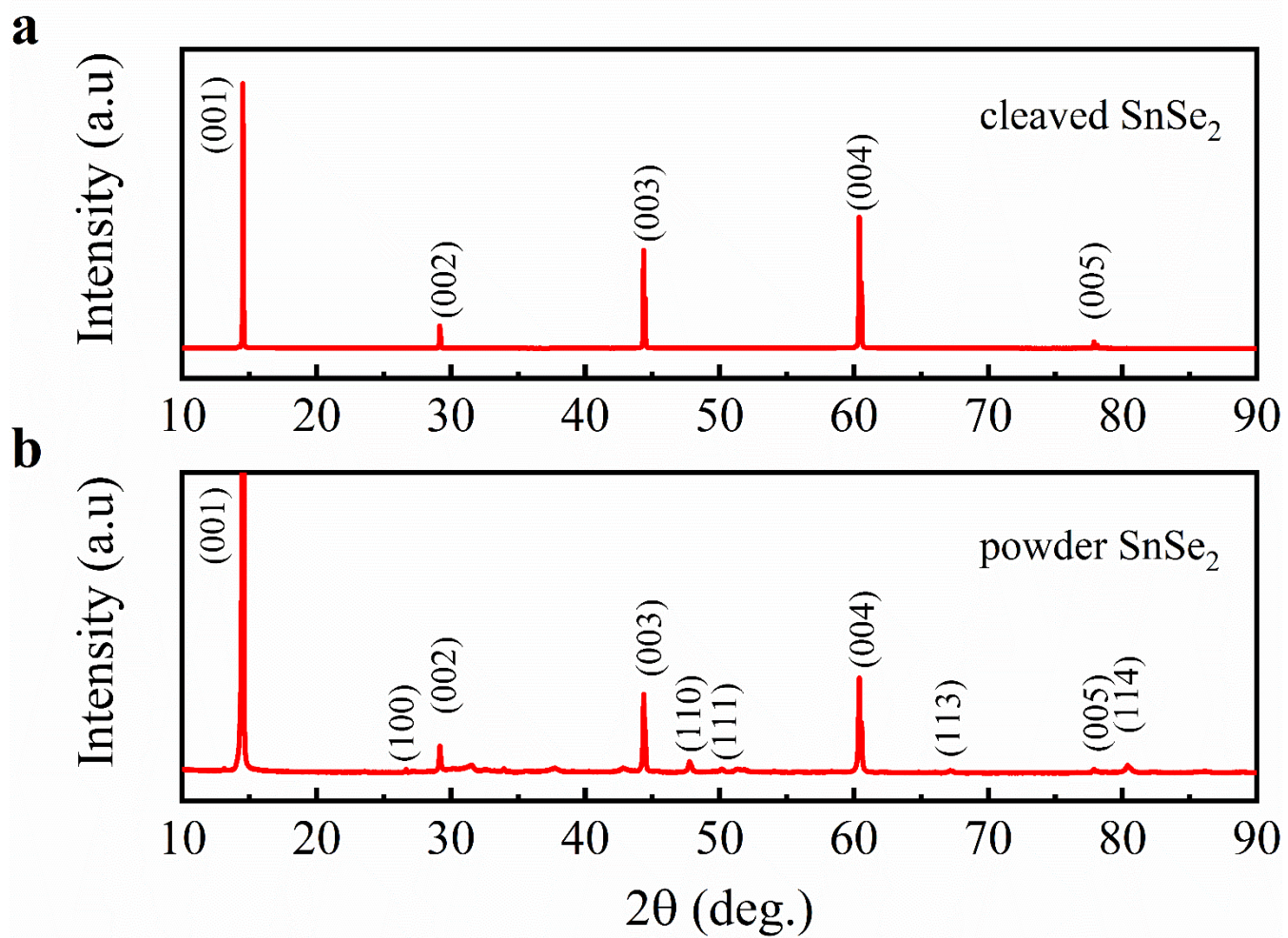


FIG. 2, Pham et al.



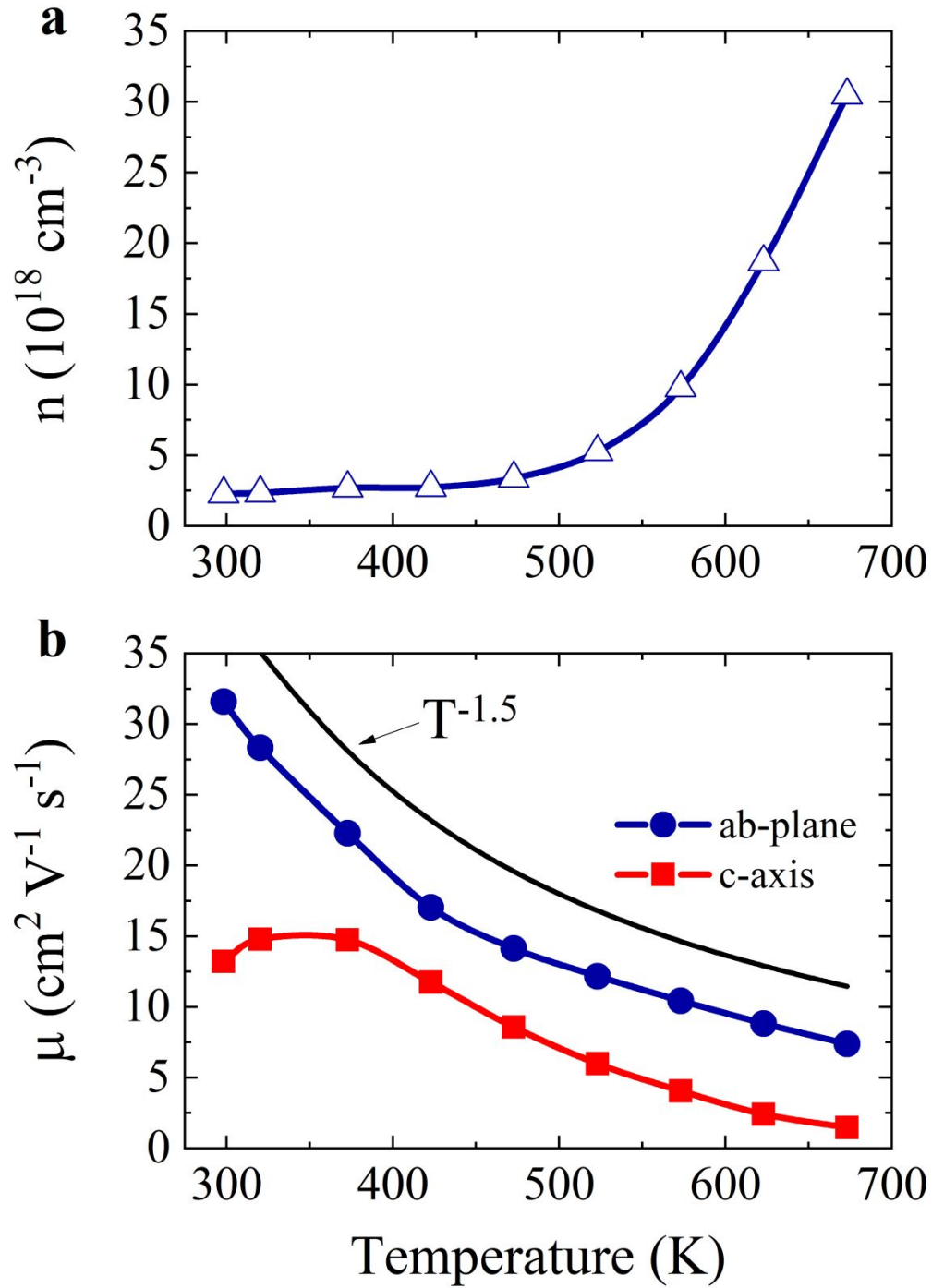
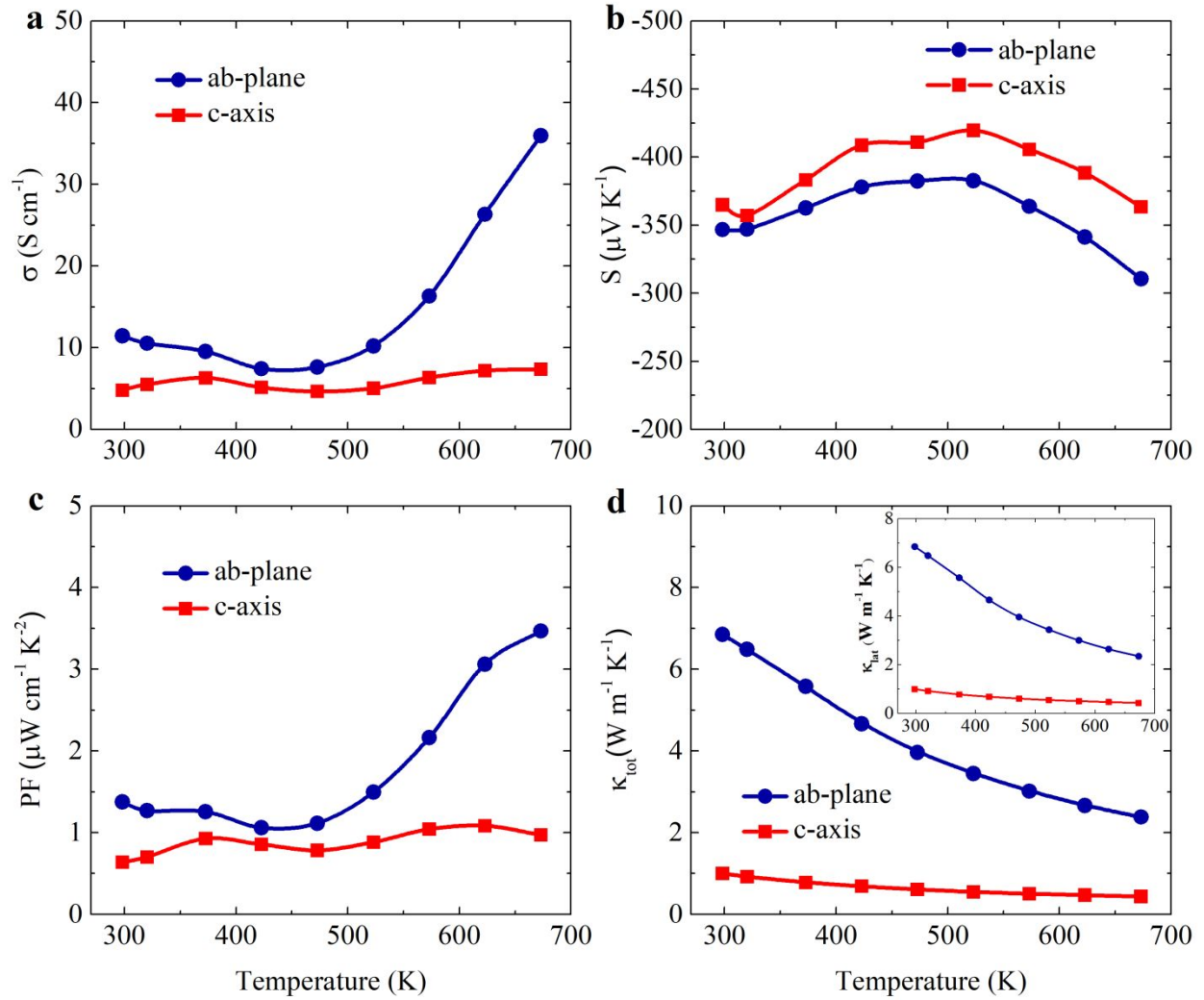


FIG. 4, Pham et al.



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FIG. 5, Pham et al.

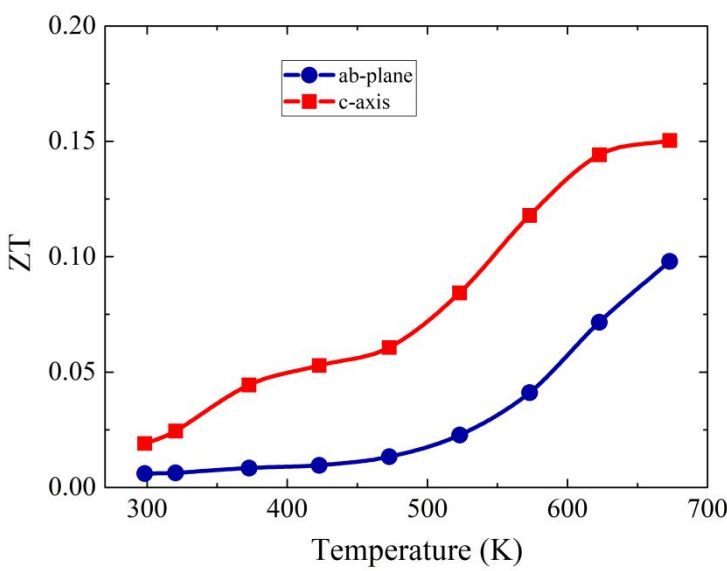


FIG. 6, Pham et al.

