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Invariance in ecological pattern

Steven A. Frank 1 and Jordi Bascompte 2

The abundance of different species in a community often follows the log series distribution. Other ecological patterns also have simple forms. Why does the complexity and variability of ecological systems reduce to such simplicity? Common answers include maximum entropy, neutrality, and convergent outcome from different underlying biological processes. This article proposes a more general answer based on the concept of invariance, the property by which a pattern remains the same after transformation. Invariance has a long tradition in physics. For example, general relativity emphasizes the need for the equations describing the laws of physics to have the same form in all frames of reference. By bringing this unifying invariance approach into ecology, one can see that the log series pattern of species abundances dominates when the consequences of density dependent processes are invariant to addition or multiplication. The lognormal pattern dominates when the total of the density independent processes obeys rotational invariance (symmetry) with respect to the summing up of the individual processes acting on abundance. Recognizing how these invariances connect pattern to process leads to a synthesis of previous approaches. First, invariance provides a simpler and more fundamental maximum entropy derivation of the log series distribution. Second, invariance provides a simple derivation of the key result from neutral theory: the log series at the metacommunity scale and a clearer form of the skewed lognormal at the local community scale. The invariance expressions of those patterns are easy to understand because they uniquely describe the basic underlying components that shape pattern.

"It was Einstein who radically changed the way people thought about nature, moving away from the mechanical viewpoint of the nineteenth century toward the elegant contemplation of the underlying symmetry [invariance] principles of the laws of physics in the twentieth century" (ref. [1,](#page-14-0) p. 153).

Introduction

Ecologists have been interested in species abundance distributions (SADs) since the classic papers by Fisher 2 2 and Preston $^3.$ $^3.$ $^3.$ Two major patterns have been identified depending on the size of the community. In a large community, abundances often follow the log series distribution 4 4 . Specifically, the probability that a species has a population size of n individuals follows p^n/n . Communities differ only in their average population size, described by the parameter, p . At smaller spatial scales, the species abundance pattern often follows a skewed lognormal (a random variable is lognormally distributed when its logarithm is nor-mally distributed)^{[5,](#page-14-4)[6](#page-14-5)}.

It is intriguing that the species abundance distribution follows these simple patterns irrespective of the particular group (birds, insects, mammals) and region considered. Other ecological patterns also fol-low simple probability distributions^{[7](#page-14-6)-9}. Those patterns have attracted a lot of attention. Why does the variability and complexity of biology reduce to such a small range of simple distributions? How can we understand the relations between complex processes and simple patterns?

Approaches such as Harte's maximum entropy formalism 9 9 and Hubbell's neutral theory 5 5 have attempted to explain the generality of the log series and skewed lognormal patterns in species abundance distributions. Maximum entropy describes probability distributions that are maximally random subject to satisfying certain constraints^{[10](#page-14-8)[–12](#page-15-0)}. This approach has a long tradition in physics, both in statistical me-

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chanics and information theory. An early maximum entropy approach in ecology derived the biomass pat-tern of populations^{[13](#page-15-1)}.

Neutral theory derives probability distributions by assuming that all individuals are equivalent^{[14](#page-15-2)}. Variation arises by random processes acting on the mechanistically identical individuals. Put another way, the mechanistic processes are "neutral" apart from random processes. Both maximum entropy and neutral theory have been shown to provide a good fit to the empirical patterns of species abundance distributions. In this article, we subsume these two different ways of understanding the log series and skewed lognormal patterns with a more general perspective based on the concept of invariance 15 .

Invariance can be defined as the property by which a system remains unchanged under some transformation. For example, a circle is the same (invariant) before and after rotation (Fig. [1a](#page-3-0)). In ecology, pattern often depends on the ways in which form remains invariant to changes in measurement. Some patterns retain the same form after uniformly stretching or shrinking the scale of measurement (Fig. [2b](#page-4-0)). Measures of length provide a common example of stretch invariance. One can measure lengths equivalently in millimeters or centimeters without loss of information. As we will see, that kind of invariance often determines the form of observed pattern.

To give another example, consider the common and widely familiar pattern of the normal distribution. By the central limit theorem, when independent random variables are added, their properly normalized sum tends toward a normal distribution even when the component variables themselves are not normally distributed. The central limit theorem and the normal distribution are often considered as unique aspects of pattern that stand apart from other commonly observed patterns.

The invariance perspective that we promote shows how the normal distribution is in fact a specific example of a wider framework in which to understand the commonly observed patterns of nature. In particular, the normal distribution arises from the rota-tional invariance of the circle^{[16](#page-15-4)}. For two independent variables, x and y , with a given squared length, the same radius, r , lie along the circumference of a $2^2 + y^2 = r^2$, all combinations of the variables with circle (Fig. [1a](#page-3-0)). When each combination is equally likely, the rotationally invariant radius is sufficient to describe the probability pattern.

It is this rotational invariance that gives the particular mathematical form of the normal distribution, in which the average squared radius sets the variance of the distribution. By this perspective, the mathematical forms of all commonly observed distributional patterns express their unique invariances 16 16 16 .

The perspective of invariance was the basis for most of the great conceptual advances of physics in the twentieth century $^1\!\!$ $^1\!\!$ $^1\!\!$. For example, Gell-Mann's pioneering theoretical work on the fundamental particles of nature derived from invariance (symmetry) properties that unified understanding of known particles and predicted new particles such as quarks, which were subsequently observed. By contrast, general aspects of invariance have not been used consistently as the fundamental basis for understanding patterns in ecology. One exception concerns scale invariance, which is often discussed in ecology $17-19$ $17-19$. But scale invariance is typically limited to special kinds of patterns rather than forming a unified approach to diverse patterns.

The point of this paper is that invariance is the most general way in which to understand commonly observed patterns. Species abundance distributions provide an excellent illustration. For species abundances, we show that maximum entropy and neutral models can succeed in certain cases because they derive from invariance principles. However, maximum entropy and neutrality are often difficult to interpret because they hide their underlying basis in invariance.

Our unifying invariance analysis clarifies why seemingly different conceptual approaches have the same consequences for pattern. Similarly, seemingly different biological processes may often lead to the same observed pattern, because those different processes share a common basis in invariance. That deeper understanding suggests a more insightful way to think about alternative mechanistic models. It also suggests the kinds of empirical tests that may differentiate between alternative causal processes.

This manuscript is organized as follows. First, we highlight new theoretical results for species abundance distributions. Second, we review how invariance defines probability patterns in a general

(a) (b) (c) (d)

Figure 1: Rotational and asymptotic invariance. (a) Transforming a circle by rotation leaves the circle unchanged (invariant), with an invariant radial distance at all points along the circumference. (b) Rotating regular polygons changes pattern. However, as more rotated polygons are added, the form converges asymptotically to a rotationally invariant circle, in which adding another rotated polygon does not change the pattern. Many common patterns of nature are asymptotically invariant. In this case, aggregation causes loss of all information except invariant radial distance. (c) The normal distribution is asymptotically invariant. The left curve describes an arbitrary probability pattern. The second curve expresses the sum of two randomly chosen values from the first curve. The height is the relative probability of the summed values. The third, fourth, and fifth curves express the sum of 4, 8, and 16 randomly chosen values from the first curve. Each curve width is shrunk to match the first curve. In this case, aggregation smooths the curve, causing loss of all information except the average squared distance from the center (the variance), which is equivalent to the average squared radial distance of rotationally invariant circles. (d) Extreme value distributions are asymptotically invariant. The left curve is an arbitrarily chosen probability pattern. The second curve expresses the probability of the largest value in a sample of two randomly chosen values from the first curve. The third, fourth, and fifth curves show the probability of the largest value of 4, 8, and 16 randomly chosen values. The asymptotically invariant curve on the right expresses exponential scaling at small values and linear scaling at large values, labeled in green and blue. Commonly observed probability distributions often express simple combinations of linear, logarithmic, and exponential scaling. Panels (a-c) modified from ref. [15.](#page-15-3)

way^{[16](#page-15-4)[,20,](#page-15-7)[21](#page-15-8)}. The log series distribution^{[22](#page-15-9)} and the new gamma-lognormal distribution for species abundances follow directly from the universal invariance expression of probability patterns. Third, we show that maximum entropy and neutrality can easily be understood as special examples of invariance principles. Finally, we discuss the broad role of invariance in the analysis of ecological pattern.

New results

This article develops two new theoretical results. We highlight those results before starting on the general overview of invariance and pattern.

First, we present a simple maximum entropy derivation of the log series pattern. We show that constraining the average abundance per species is sufficient when analyzing randomness and entropy on the proper demographic scale.

The simplicity of our maximum entropy derivation contrasts with Harte's more complicated maximum entropy model^{[9](#page-14-7)[,23](#page-15-10)}. Harte had to assume an additional unnecessary constraint on energy usage. He required that unnecessary constraint because he evaluated randomness on the scale of measured abundances rather than on the scale of demographic process. This will be made explicit below.

We use this new result to demonstrate that max-

Figure 2: Shift and stretch invariance of the exponential distribution. (a) The left panel shows $e^{-(x+a)}$ for $a = 0, -1, \ldots, -4$.
Decreasing values of a shift the curve to the right, which is equivalent to shifting the x Decreasing values of a shift the curve to the right, which is equivalent to shifting the x axis by resetting the zero point. For probability patterns, the total probability must be normalized to one, which means that all curves must have the same area under the curve for values of x between 0 and ∞. To normalize the curves, the right panel plots $k_a e^{-(x+a)}$ with $k = e^a$. Thus all curves become e^{-x} invariantly with respect to different shift values, a. (b) The le $k_a = e^a$. Thus, all curves become e^{-x} invariantly with respect to different shift values, a. (b) The left panel shows e^{-bx}
for $b = 2^0$, 2^{-1} and 2^{-4} . Decreasing values of b stratch the x axis by a factor of 2 f for $b = 2^0, 2^{-1}, \ldots, 2^{-4}$. Decreasing values of b stretch the x axis by a factor of 2 for each halving of b. To normalize the system of each probability curve to be the same, the right panel shows $e^{-\lambda_b bx}$ for $\lambda_b = 1/b$ the average value of each probability curve to be the same, the right panel shows $e^{-\lambda_b bx}$ for $\lambda_b = 1/b$. Thus, all curves become e^{-x} invariantly with respect to different stratch values h become e^{-x} invariantly with respect to different stretch values, *b*.

imum entropy is the outcome of deeper underlying principles of invariance and pattern. By working at the deeper level of invariance, one obtains a simpler and more powerful understanding of pattern.

The second new result shows that Hubbell's $⁵$ $⁵$ $⁵$ neu-</sup> tral model is the simple expression of three basic invariances. Hubbell's full range of log series and skewed lognormal (zero sum multinomial) results follows immediately from those three underlying invariances.

The three invariances correspond to a maximum entropy model that constrains the average abundance of species and the average and variance of the demographic processes influencing abundance. The three invariances lead to a simple gamma-lognormal distribution that matches the neutral theory pattern for species abundances. The gamma-lognormal is a product of the standard gamma and lognormal distributions.

Invariance

This section reviews how invariance considerations lead to the log series distribution^{[22](#page-15-9)}. We delay discussion of the gamma-lognormal until the later section on Hubbell's neutral model.

Canonical form of probability distributions

We can rewrite almost any probability distribution as

$$
q_z = k e^{-\lambda T_z}, \tag{1}
$$

in which $T(z) \equiv T_z$ is a function of the variable, z, and k and λ are constants. For example, Student's t-distribution, usually written as

$$
q_z = k \left(1 + z^2 / v \right)^{-(\nu + 1)/2}
$$

can be written in the form of eqn [1](#page-4-1) with $\lambda = (\nu + 1)/2$ and $T_z = \log(1 + z^2/v)$.
The probability patts

The probability pattern, q_z , is invariant to a constant shift, $T_z \mapsto a + T_z$, because we can write the transformed probability pattern in eqn [1](#page-4-1) as

$$
q_z = k_a e^{-\lambda(a+T_z)} = k e^{-\lambda T_z},
$$

with $k = k_a e^{-\lambda a}$ (Fig. [2a](#page-4-0)). We express k in this way
because k adjusts to satisfy the constraint that the because k adjusts to satisfy the constraint that the total probability be one. In other words, conserved total probability implies that the probability pattern is shift invariant with respect to T_z (see ref. [16\)](#page-15-4).

Now consider the consequences if the average of some value over the distribution q_z is conserved. For example, the average of z is the mean, μ , and the average of $(z - \mu)^2$ is the variance. A constraint causes
the probability pattern to be invariant to a multiplica the probability pattern to be invariant to a multiplicative stretching (or shrinking), $T_z \mapsto bT_z$, because

$$
q_z = k e^{-\lambda_b b T_z} = k e^{-\lambda T_z},
$$

with $\lambda = \lambda_b b$ (Fig. [2b](#page-4-0)). We specify λ in this way because λ adjusts to satisfy the constraint of conserved average value. Thus, invariant average value implies that the probability pattern is stretch invariant with respect to T_z .

Conserved total probability and conserved average value cause the probability pattern to be invariant to an affine transformation of the T_z scale, $T_z \mapsto a + bT_z$, in which "affine" means both shift and stretch.

The affine invariance of probability patterns with respect to T_z induces significant structure on the form of T_z and the associated form of probability patterns. Understanding that structure provides insight into probability patterns and the processes that generate them 16,20,21 16,20,21 16,20,21 16,20,21 16,20,21 .

In particular, Frank & Smith^{[20](#page-15-7)} showed that the invariance of probability patterns to affine transformation, $T_z \mapsto a + bT_z$, implies that T_z satisfies the differential equation

$$
\frac{\mathrm{d}T_z}{\mathrm{d}w} = \alpha + \beta T_z,
$$

in which $w(z)$ is a function of the variable z. The solution of this differential equation expresses the scaling tion of this differential equation expresses the scaling of probability patterns in the generic form

$$
T_z = \frac{1}{\beta} \Big(e^{\beta w} - 1 \Big), \tag{2}
$$

in which, because of the affine invariance of T_z , we have added and multiplied by constants to obtain a have added and multiplied by constants to obtain a convenient form, with $T_z \rightarrow w$ as $\beta \rightarrow 0$. With this expression for T_z , we may write probability patterns generically as

$$
q_z = k e^{-\lambda (e^{\beta w} - 1)/\beta}.
$$
 (3)

Proportional processes and species abundances

To understand the log series, we must consider the relation $n = e^r$ between the observed pattern of abun-
danger r and the processes r . Here, r represents dances, n , and the processes, r . Here, r represents the total of all proportional processes acting on abun- $dance²²$ $dance²²$ $dance²²$.

A proportional process simply means that the number of individuals or entities affected by the process increases in proportion to the number currently present, n. Demographic processes, such as birth and death, act proportionally.

The sum of all of the proportional processes on abundance over some period of time is

$$
r=\int_0^\tau m(t)\mathrm{d}t.
$$

Here, $m(t)$ is a proportional process acting at time t to change abundance. Birth and death typically occur as proportional processes. The value of $r = \log n$ is the total of the *m* values over the total time, τ . For simplicity, we assume $n_0 = 1$.

The log series follows as a special case of the generic probability pattern in eqn [3.](#page-5-0) To analyze abundance, focus on the process scale by letting the variable of interest be $z \equiv r$, with the key scaling simply the process variable itself, $w(r) = r$. Then eqn [3](#page-5-0) becomes

$$
q_r dr = k e^{-\lambda (e^{\beta r} - 1)/\beta} dr, \qquad (4)
$$

in which $q_r dr$ is the probability of a process value, r, in the interval $r + dr$. From the relation between abundance and process, $n = e^r$, we can change from
the process scale, r, to the abundance scale, n, by the the process scale, r , to the abundance scale, n , by the substitutions $r \mapsto \log n$ and $dr \mapsto n^{-1}dn$, yielding the identical probability pattern expressed on the abund identical probability pattern expressed on the abundance scale

$$
q_n \mathrm{d}n = k n^{-1} e^{-\lambda \left(n^{\beta} - 1\right)/\beta} \mathrm{d}n. \tag{5}
$$

The value of k always adjusts to satisfy the constraint of invariant total probability, and the value of λ always adjusts to satisfy the constraint of invariant average value.

For $\beta = 1$, we obtain the log series distribution^{[22](#page-15-9)}

$$
q_n = kn^{-1}e^{-\lambda n},\tag{6}
$$

replacing $n - 1$ by n in the exponential term which, because of affine invariance, describe the same probability pattern. The log series is often written with serves discrete values $n = 1, 2, \ldots$. See Frank^{[22](#page-15-9)} for $-\lambda = p$, and thus $q_n = kp^n/n$. One typically ob-
organization values $n = 1, 2$. See Frank²² for

a discussion of the general relation between discrete and continuous distributions. The continuous analysis here is sufficient to understand pattern.

We can also write the log series on the process scale, r, from eqn [4,](#page-5-1) as 22 22 22

$$
q_r = k e^{-\lambda e^r}.
$$
 (7)

This form shows that the log series is the simplest expression of generic probability patterns in eqn [3.](#page-5-0) The log series arises from $\beta = 1$, associated with $n = e^r$,
and from the base scale as $w = r$ for proportional and from the base scale as $w \equiv r$ for proportional processes, r.

Invariances of the log series

This subsection summarizes a few technical points about invariance. These technical points provide background for our simpler and more general derivation in the following section of maximum entropy models for species abundances. Those prior models focused only on abundances, n , without considering the underlying process scale, r.

We begin with invariance on the process scale, r . On that scale, the log series in eqn [7](#page-6-0) is the pure expression of additive shift invariance to r and lack of multiplicative stretch invariance to r . For example, note in eqn [7](#page-6-0) that an additive change, $r \mapsto r + a$, is compensated by a change in λ to maintain the overall invariance, whereas a multiplicative change, $r \mapsto br$, cannot be compensated by a change in one of the constants. For example, if r is net reproductive rate, then an improvement in the environment that adds a constant to everyone's reproductive rate does not alter the log series pattern. By contrast, multiplying reproductive rates by a constant does alter pattern.

To understand the parameter, β , from eqn [2,](#page-5-2) consider that

$$
T = \frac{1}{\beta} \left(e^{\beta r} - 1 \right) = \frac{1}{\beta} \left(n^{\beta} - 1 \right),
$$

in which $β$ is the relative curvature of the measurement scale for abundance, n , with respect to the scale for process, *r*. The relative curvature is $\beta = T''/T'$,
with the primes depoting differentiation with respect with the primes denoting differentiation with respect to r.

For the log series, the curvature of $\beta = 1$ describes the amount of bending of the abundance scale, $n = e^r$,

with respect to multiplying the process scale, r , by a constant—the departure from stretch invariance.

The simple invariances with respect to process, r , become distorted and more difficult to interpret when we focus only on the observed scale for abundance, n , associated with the log series in eqn [6.](#page-5-3) In that form of the distribution, the canonical scale is

$$
T = \frac{1}{\lambda} \log n + n. \tag{8}
$$

In this expression, purely in terms of abundances, the logarithmic term dominates when n is small, and the linear term dominates when n is large. Thus, the scale changes from stretch but not shift invariant at small magnitudes to both shift and stretch invariant at large magnitudes^{[22](#page-15-9)}.

Without the simple insight provided by the process scale, r, we are left with a complicated and nonintuitive pattern that is separated from its simple cause. That difficulty has led to unnecessary complications in maximum entropy theories of pattern.

Maximum entropy

Maximum entropy describes probability distributions that are maximally random subject to satisfying certain constraints 10^{-12} 10^{-12} . In eqn [1,](#page-4-1) with the generic description for distributions as

$$
q_z \mathrm{d} z = k e^{-\lambda T_z} \mathrm{d} z,
$$

maximum entropy interprets this form as the expression of maximum randomness with respect to the scale z, subject to the constraint that the average of T_z is fixed^{[21](#page-15-8)}.
This section

This section begins with a maximum entropy derivation for the log series based on our separation between the scales of process, r , and observed abundance, n.

We then discuss Harte's^{[9,](#page-14-7)[23](#page-15-10)} alternative maximum entropy derivation of the log series. Harte's derivation emphasizes mechanistic aspects of energy constraints rather than our emphasis on the different scales of process and abundance.

Constraint of average abundance on process scale

The log series in eqn [7](#page-6-0) is

$$
q_r dr = k e^{-\lambda e^r} dr.
$$

Here, $T = e^r = n$. This distribution expresses max-
imum entropy with respect to the precess scale r . imum entropy with respect to the process scale, r . The constraint is the ecological limitation on average abundance

$$
\langle T \rangle_r = \langle e^r \rangle_r = \langle n \rangle_r, \tag{9}
$$

in which $\langle \cdot \rangle_r$ denotes average value with respect to the process scale r . the process scale, r.

In this case, process values, r , are maximally random, subject to the ecological constraint that limits abundance, n. Thus, maximizing entropy with respect to the process scale, r , subject to a constraint on the observed pattern scale, n , leads immediately to the log series.

Relating the process scale, r , to the scale of ecological constraint, n, often makes sense. Typically, environmental perturbations associate with changes in demographic variables, such as birth and death rates. Such demographic factors typically act proportionally on populations, consistent with our interpretation of r as the aggregate of proportionally acting processes. The perturbations, acting on demographic variables, associate the process scale with the scale of randomness.

In contrast with the process scale of perturbation and randomness for the demographic variables, the scale of constraint naturally arises with respect to a limit on the number of individuals, n . Thus, randomness happens on the r scale and constraint happens on the n scale.

It is, of course, possible to formulate alternative models in which randomness and constraint happen on scales that differ from our interpretation. Different formulations are not intrinsically correct or incorrect. Instead, they express different assumptions about the relations between process, randomness, and invariance. The next section considers an alternative formulation.

Harte's joint constraints of abundance and energy

Harte developed comprehensive maximum entropy models of ecological pattern. He tested those theories with the available data. His work synthesizes many aspects of ecological pattern^{[9](#page-14-7)}.

For species abundances, Harte^{[9,](#page-14-7)[23](#page-15-10)} analyzed maximum randomness with respect to the scale of abundance values, n. Maximum entropy derivations commonly evaluate randomness on the same scale as the observations. In this case, with observations for the probabilities of abundances, p_n , entropy on the same scale is the sum or integral of $-p_n \log p_n$.

However, there is no a priori reason to suppose that the scale of observation is the same as the scale of randomness. The fact that observation, randomness, and process may occur on different scales often makes maximum entropy models difficult to develop and difficult to interpret. For example, we may observe the probabilities of abundances, p_n , but randomness may be maximized on the scale of process, as the sum or integral of $-p_r \log p_r$.
In the final part of the

In the final part of this section, we argue that invariance provides a truer path to the natural scale of analysis and to the mechanistic processes that generate pattern than does maximum entropy. Before comparing invariance and maximum entropy, it is useful to sketch the details of Harte's maximum entropy model for species abundances.

The simplest maximum entropy model analyzes entropy with respect to abundance, n , subject to a constraint on the average abundance, $\langle n \rangle$. That analysis yields an exponential distribution

$$
q_n \mathrm{d}n = k e^{-\lambda n} \mathrm{d}n.
$$

The exponential pattern differs significantly from the observed log series pattern. Thus, maximizing entropy with respect to the scale of abundance, n , and constraining the average abundance is not sufficient.

From our invariance perspective, it is natural to think of the scale of randomness in terms of dr , the scale of proportional processes, rather than in terms of dn, the scale of abundance. Maximizing randomness with respect to dr leads directly to the log series, as shown in the previous section.

Harte did not consider the distinction between the exponential and log series patterns with respect to the scale of randomness. Instead, to go from the default exponential pattern of maximum entropy to the log series, his maximum entropy analysis required additional assumptions. He proceeded in the following way.

Suppose that the total quantity of some variable, ϵ , is constrained to be constant over all individuals of all species. The average value per individual is $\langle \epsilon \rangle$. It does not matter what the variable ϵ is. All that matters is that the constraint exists. Harte assumed that ϵ is energy, but that assumption is unnecessary with regard to the species abundance distribution.

The value ϵ is distributed over individuals independently of their species identity. Thus, the variable $\delta |n = n\epsilon$ is the total value in a species with *n* individuals, with average value $\langle \delta | n \rangle = n \langle \epsilon \rangle$.

The joint distribution of *n* and δ is

$$
q_{n,\delta}=q_nq_{\delta|n}.
$$

The explicit form of this joint distribution can be obtained by maximizing entropy subject to the constraints on the average abundance per species, $\langle n \rangle$, and the average total value in a species with n individuals, $\langle \delta | n \rangle$, yielding

$$
q_{n,\delta} = k e^{-\lambda n} e^{-\lambda' \delta}
$$

We obtain the form presented by Harte^{[23](#page-15-10)} using the equivalence $\delta = n\epsilon$, yielding

$$
q_{n,\epsilon}=ke^{-\lambda n}e^{-\lambda' n\epsilon}.
$$

The species abundance distribution is obtained by

$$
q_n = \int q_{n,\epsilon} d\epsilon = k e^{-\lambda n} \int e^{-\lambda' n \epsilon} d\epsilon.
$$

Noting that $\int e^{-\lambda' n \epsilon} = 1/\lambda' n$, and absorbing the con-
stant λ' into k, we obtain the log series for the species $\lim_{n \to \infty} \frac{1}{n}$ into k, we obtain the log series for the species abundance distribution

$$
q_n = kn^{-1}e^{-\lambda n}.
$$

Maximum entropy and invariance

Harte's maximum entropy derivation of the log series assumes joint constraints of abundance, n , and some auxiliary variable, ϵ , which he labeled as energy. He evaluated entropy on the scales of n and ϵ .

By contrast, our invariance derivation arises from a constraint on abundance plus evaluation of invariance or entropy on the scale $r = \log n$. On that scale, the log series arises in a simple and clear way. There is no need for constraint of a second auxiliary variable.

Without an invariance argument, nothing compels us to analyze with respect to the r scale. Harte, without focus on invariance, followed the most natural approach of using n as the scale for maximization of randomness and for constraint. That approach required the creation of an auxiliary constraint on a second scale to arrive at the log series.

Harte's approach was a major step in unifying the analysis of empirical pattern. But, in retrospect, his approach was unnecessarily complicated.

One might say that Harte's approach provided a richer theory because it led to predictions about both abundance and energy. However, the data on abundance patterns match very closely to the log series, whereas the data for different proxies of energy vary considerably^{[9](#page-14-7)}.

Our invariance approach strips away the unnecessary, vaguely conceived, and weakly supported auxiliary variable. The invariance theory therefore provides a much simpler way to derive and to understand abundance patterns.

Maximum entropy can be thought of purely as a basic invariance method of analysis. Maximum entropy distributions have the form in eqn [1](#page-4-1) as

$$
q_z \mathrm{d} z = k e^{-\lambda T_z} \mathrm{d} z,
$$

in which T_z is the affine-invariant scale that defines the probability pattern. Thus, the method of maximum entropy is simply a method for deriving the affine invariant expression, T_z . In practice, maximum entropy has three limitations.

First, maximum entropy is silent with respect to the proper choice for the scale on which entropy is maximized and the constraints that set the affine invariant expression, T_z . By contrast, focus on invariance led us to the shift invariance of the process scale, r . That scale provided a much simpler analysis, in which r is the incremental scale with respect to invariance and the measurement scale with respect to entropy.

In other words, maximum entropy is a blind application of the most basic invariance principles, without any guidance about the proper scales for invariance, randomness, and constraint. By contrast, an explicit invariance approach takes advantage of the insight provided by the analysis of invariance.

Second, by focusing on invariance, we naturally ob-

tain the full invariance (symmetry) group expression in eqn [3](#page-5-0) as the generic form of probability patterns

$$
q_z = k e^{-\lambda (e^{\beta w} - 1)/\beta}
$$

That generic expression leads us to a generalization of the log series in eqn [5](#page-5-4) as 22

$$
q_n \mathrm{d}n = k n^{-1} e^{-\lambda \left(n^{\beta}-1\right)/\beta} \mathrm{d}n,
$$

which is a two parameter distribution for abundances with respect to λ and β . The log series is a special case with $\beta = 1$.

Third, invariance leads to a deeper understanding of the relation between observed pattern and alternative mechanistic models of process. The following section provides an example.

Neutrality

Here, we analyze Hubbell's 5 5 neutral model of species abundances in the light of our invariance perspective. With that example in mind, we then discuss more generally how neutral models relate to invariance and maximum entropy.

Hubbell's neutral model

The strong recent interest in Hubbell's neutral model follows from the match of the theory to the contrasting patterns of species abundance distributions (SADs) that have been observed at different spatial scales. In the theory, many local island-like communities are connected by migration into a broader metacommunity. Sufficiently large metacommunities follow the log series pattern of species abundances. Each local community follows a distribution that Hubbell called the zero-sum multinomial^{[24](#page-15-11)}, which is similar to a skewed lognormal. As noted by Rosindell et al. 6, 6, 6, it is this flexibility of the classic neutral model to reconcile the log series and lognormal distributions that allows it to fit empirical data well^{[25](#page-15-12)}.

Invariance and the gamma-lognormal distribution

Broad consensus suggests that species abundances closely follow the log series pattern at large spatial scales. Extensive data support that conclusion^{[4](#page-14-3)}.

Observed pattern at small spatial scales differs from the log series. Consensus favors a skewed lognormal pattern. The data typically show an excess of rare species, causing a skew relative to the symmetry of the lognormal when plotted on a logarithmic scale.

At small spatial scales, most recent analyses focus on data from a single long-term study of tree species in Panama^{[5,](#page-14-4)[24](#page-15-11)}. Thus, some ambiguity remains about the form and consistency of the actual pattern at small scales.

The blue curve of Fig. [3](#page-10-0) shows Chisholm and Pacala's^{[24](#page-15-11)} fit of the neutral theory to the Panama tree data for species abundances at small spatial scales. The gold curve shows the close match to the neutral theory pattern by a simple probability distribution derived from the analysis of invariance.

To obtain the matching distribution derived by invariance, we begin with the canonical form for probability distributions in eqn [3.](#page-5-0) That canonical form expresses pattern in terms of the shift-invariant scale, w. Next, we need to find the specific form of the scale w that relates this canonical form for probability distributions to the neutral theory. Because the neutral theory derives abundance, n , as an outcome of demographic processes, r , the fundamental shift-invariant scale for neutral theory is expressed in terms of the demographic process variable as

$$
w = \log \left(e^r - \frac{a}{\lambda} r + \frac{\alpha}{\lambda} r^2 \right).
$$
 (10)

Below, we discuss why this is a natural shift-invariant scale for neutral theory. For now, we focus on the details of the mathematical expressions. Recall that $n = e^r$ relates measured abundances, *n*, to the demo-
graphic process scale, *r*. If we assume that $\beta = 1$ in graphic process scale, r. If we assume that $\beta = 1$ in eqn [3](#page-5-0) and use w from eqn [10,](#page-9-0) we obtain

$$
q_r = k e^{-\lambda e^r + ar - \alpha r^2}, \qquad (11)
$$

with parameters λ , a, and α . We can write this distribution equivalently on the n scale for abundance as

$$
q_n = kn^{\tilde{a}-1}e^{-\lambda n}e^{-\alpha(\log n - \mu)^2}.
$$
 (12)

In the second distribution, $\mu = (a - \tilde{a})/2a$. Thus, both distributions have the same three parametric degrees of freedom.

The right-hand exponential term of eqn [12](#page-9-1) is a

Figure 3: Match of the gamma-lognormal pattern in gold to the neutral theory fit in blue for Panama tree species abun-dances. The neutral theory fit to the data comes from Chisholm and Pacala's^{[24](#page-15-11)} analysis in their Fig. 1. They used Hubbell's neutral theory model with parameters $J = 21,060$, $m = 0.075$, and $\theta = 52.1$ in their eqn 3, originally from Alonso and McKane^{[26](#page-15-13)}. The gamma-lognormal model in eqn [11](#page-9-2) produces essentially the identical pattern with parameters $\lambda = 0.00205$, $a = 0.491$, and $\alpha = 0.0559$. The abundance scale can be expressed equivalently on the process scale, $\log_2 n = r/\log 2$. See the Data accessibility statement for the calculations used to produce this plot.

lognormal distribution with parameters μ and $\sigma^2 =$
1/2 α . The remaining terms are a gamma distribution $1/2\alpha$. The remaining terms are a gamma distribution with parameters \tilde{a} and λ . We call this product of the gamma and lognormal forms the gamma-lognormal distribution.

Figure [3](#page-10-0) showed that the gamma-lognormal distribution matches the neutral theory fit for the Panama tree data. Fig. [4](#page-11-0) shows that the shape of the gammalognormal matches the shape of the neutral theory predictions for various mechanistic parameters of the neutral theory.

In summary, the neutral theory distribution appears to be nearly identical to a gamma-lognormal distribution when compared over realistic parameter values. Both distributions have the same three parametric degrees of freedom.

Maximum entropy and the gamma-lognormal

The constraints on pattern can be seen most clearly by rewriting eqn [11](#page-9-2) as

$$
q_r = k e^{-\lambda T_r} = k e^{-\lambda e^r + \tilde{a}r - \alpha \tilde{r}^2}, \qquad (13)
$$

in which $\tilde{r}^2 = (r - \mu)^2$ is the squared deviation from μ in which μ is the average value of r. This express μ , in which μ is the average value of r. This expression remains a three-parameter distribution because, as noted above, $\mu = (a - \tilde{a})/2a$.

With this set of parameters, the affine-invariant scale is

$$
T_r = e^r - \frac{\tilde{a}}{\lambda}r + \frac{\alpha}{\lambda}\tilde{r}^2.
$$
 (14)

git • preprint @ preprint-1.0-1::f0222b9-2019-06-17 (2019-06-18 00:54Z) • safrank **0 2 4 6 8 10 12** Note that T and w are related by eqn [2.](#page-5-2) We are using
w from eqn 10 and $\beta = 1$, as noted below eqn 10. We w from eqn [10](#page-9-0) and $\beta = 1$, as noted below eqn [10.](#page-9-0) We ignore the extra -1 term in T of eqn [2,](#page-5-2) because the canonical form of probability distributions is invariant to adding a constant to T . The tilde parameters of the distribution in eqn [13](#page-10-1) are interchangeable with the nontilde parameters of the identical distribution in eqn [12.](#page-9-1) The tilde expressions focus on the invariances that will help us to interpret ecological pattern. The nontilde expressions describe pattern in terms of the classic forms for the gamma and lognormal distributions.

By the standard theory of maximum entropy, q_r maximizes entropy on the incremental scale dr subject to a constraint on the average value of the defining affine-invariant scale, $\langle T \rangle_r$. That constraint is the linear combination of three constraints: the average linear combination of three constraints: the average abundance on the process scale, $\langle n = e^r \rangle_r$, the aver-
age demographic process value $\langle r \rangle$ and the variance age demographic process value, $\langle r \rangle$, and the variance
in the demographic process values $\langle \tilde{x}^2 \rangle$ in the demographic process values, $\langle \tilde{r}^2 \rangle$.
By maximum optropy, all of the inf

By maximum entropy, all of the information in Hubbell's mechanistic process theory of neutrality

Figure 4: Match of Hubbell's neutral theory to the gamma-lognormal distribution. The blue curve for the neutral theory and the gold curve for the gamma-lognormal are calculated as described in Fig. [3.](#page-10-0) The parameters for the neutral theory are the same as in Fig. [3,](#page-10-0) except as shown in each panel. I fit the parameters for the gamma-lognormal to each neutral theory curve, with values for each panel: (a) $\lambda = 0.01115$, $a = 0.4452$, and $\alpha = 0.03660$; (b) $\lambda = 0.0004209$, $a = 0.4622$, and $\alpha = 0.05014$; (c) $\lambda = 0.002765$, $a = 0.3182$, and $\alpha = 0$; (d) $\lambda = 0.001777$, $a = 0.2217$, and $\alpha = 0.03576$; (e) $\lambda = 0.0001509$, $a = 0.3851$, and $\alpha = 0.03667$; (f) $\lambda = 0.02519$, $a = 0.3726$, and $\alpha = 0.006900$. See the Data accessibility statement for the calculations used to produce these plots.

and the matching gamma-lognormal pattern reduces to maximum randomness subject to these three constraints.

However, it is very unlikely that we would have derived the correct form by maximum entropy without knowing the answer in advance. This limitation emphasizes that maximum entropy provides deep insight into process and pattern, but often we need an external theory to guide our choice among various possible maximum entropy formulations.

Put another way, maximum entropy and process oriented theories, such as Hubbell's model, often work together synergistically to provide deeper insight than either approach alone.

Invariance, information and scale

Before turning to invariance and the gammalognormal pattern of neutral theory, it is useful to consider some basic properties of invariance and in-formation^{[27,](#page-15-14)[28](#page-15-15)}. In particular, this subsection develops our claim that the affine-invariant scale provides the deepest insights into the relations between pattern and process.

We start by noting that, in the general expression

for probability distributions

$$
q_z = k e^{-\lambda T_z},
$$

the affine-invariant scale, T_z , is equivalent to a common expression for the information content in a measurement, z, as

$$
I_z=-\log q_z.
$$

This expression follows from assuming that: information depends on the probability, q_z , of observing the measured value and not on the value itself; rarely observed values provide more information than commonly observed values; and the information in two independent measurements is the sum of the information in each measurement.

From the general expression for probability distributions

$$
I_z = -\log q_z = -\log k + \lambda T_z.
$$

Thus, an incremental change in information is equal to an incremental change in the affine-invariant scale

$$
dI_z = -d \log q_z = \lambda dT_z.
$$

Equivalently, the change in information with respect

ke<sup>-
$$
\lambda
$$
T</sup>
(a)
 λ
 τ
(b)
 $\pm \sqrt{T}$
 $\pm \sqrt{T}$

Figure 5: Continuous probability distributions can often be expressed as exponential or normal distributions with respect to the affine-invariant scale. A continuous distribution typically can be written as $q_z = ke^{-\lambda T_z}$, from eqn [1.](#page-4-1) In the figure, $T = T$, (a) A parametric plot of $a_y \text{ or } T$ is exponential. All of differences between probabili $T \equiv T_z$. (a) A parametric plot of q_z vs T_z is exponential. All of differences between probability distributions are contained in the form of the affine-invariant scale, T_z . The change in information for each increment of the affine-invariant scale is λ ,
as in ean 15, (b) A parametric plot of a , $y = \pm \sqrt{T}$ is pormally distributed when desc as in eqn [15.](#page-12-0) (b) A parametric plot of q_z vs $\pm \sqrt{T_z}$ is normally distributed when describing the deviations from a unimodal
peak of a. The average of the deviations on the affine-invariant scale $\langle T \rangle$ relative to me peak of q_z . The average of the deviations on the affine-invariant scale, $\langle T \rangle$, relative to measurements on the square root реак от q_z . тне
of that scale, $\sqrt{ }$ $\overline{T_z}$, is the variance, σ^2 . For the normal distribution, we can think of a deviation from the central location variant scale $T = R^2$ as the squared radial deviations along the circumference of a circle with radius on the affine-invariant scale, $T_z = R_z^2$, as the squared radial deviations along the circumference of a circle with radius R_z , describing the squared vector length for an aggregation of variables. The variance is the av describing the squared vector length for an aggregation of variables. The variance is the average of the squared radial describing the squared vector length for an aggregat
deviations relative to the scale of radial measures, $\sqrt{ }$ $\overline{T_z} = R_z$. Most continuous unimodal distributions are, in this way, respect to the square root of the affine invariant measure. See refs. equivalent to a normal distribution when scaled with respect to the square root of the affine-invariant measure. See refs. [15](#page-15-3)[,16](#page-15-4) for details.

to a change in the affine-invariant scale,

$$
\frac{\mathrm{d}I_z}{\mathrm{d}T_z} = \lambda,\tag{15}
$$

is constant at all magnitudes of the measurement, z. Every measured increment on the T_z scale provides the same amount of information about pattern. Constancy of information at all magnitudes is the ideal for a measurement scale. Thus, affine invariance provides the ideal scale on which to evaluate the pat-tern in measurements^{[21](#page-15-8)}. Figure [5](#page-12-1) illustrates some key properties of the affine-invariant scale.

Information is sometimes thought of as a primary concept. However, it is important to understand that, in this context, information and affine invariance are the same thing. Neither is intrinsically primary.

We prefer to emphasize invariance, because it is an explicit description of the properties that pattern and process must obey^{[15](#page-15-3)[,29](#page-15-16)[,30](#page-15-17)}. Further analysis of invariant properties leads to deeper insight. By contrast, "information" is just a vague word that associates with underlying invariances. Further analysis of information requires unwinding the definitions to return to the basic invariances.

Invariance interpretation of the gammalognormal

We turn now to the neutral theory model for abundances at local spatial scales. We showed that all of the information about pattern and process in the neutral theory is captured by the gamma-lognormal pattern in eqn [13](#page-10-1) as

$$
q_r = k e^{-\lambda T_r} = k e^{-\lambda e^r + \tilde{a}r - \alpha \tilde{r}^2},
$$

which defines the affine-invariant scale in eqn [14](#page-10-2) as

$$
\lambda T_r = \lambda e^r - \tilde{a}r + \alpha \tilde{r}^2. \tag{16}
$$

On this scale, changes in r provide the same amount of information about pattern at all magnitudes. Shifting the scale by a constant does not change the information about pattern in measurements. In other words, it does not matter where we set the zero point for T_r . Similarly, uniformly stretching or shrinking
the scale T_r does not change the information in mea the scale, T_r , does not change the information in measurements of r surements of r.

We can parse the terms of eqn [16](#page-12-2) with respect to constraint and invariance. When r is large, the term $\lambda e^{r} = \lambda n$ dominates the shape of the distribution in the upper tail, which decays as

$$
q_r dr = k e^{-\lambda e^r} dr = k e^{-\lambda n} dr
$$

for sufficiently large $e^r = n$. The smaller the value
of λ relative to \tilde{a} and α the greater e^r must be for of λ relative to \tilde{a} and α , the greater e^r must be for
this pattern to dominate. When λ is relatively large this pattern to dominate. When λ is relatively large compared with \tilde{a} and α , this pattern dominates at all magnitudes and leads to the log series.

With respect to constraint, for large values of abundance, n, the constraint on average abundances dominates the way in which altered process influences pattern. In other words, density dependence dominates.

With respect to invariance, a process that additively shifts or multiplicatively stretches the $e^r = n$ values
does not alter the pattern in the upper tail. Similarly does not alter the pattern in the upper tail. Similarly, pattern is invariant to a process that additively shifts process values, r , but processes that multiplicatively change r alter pattern. Thus, we can evaluate the role of particular processes by considering how they change n or r .

The pattern at small and intermediate values of r depends on the relative sizes of the parameters. If the ar term dominates, then the constraint, $\langle r \rangle$, on the average process value is most important. For example, rare species may be most strongly influenced by density-independent limitations on opportunities for growth or density-independent death processes.

With respect to invariance when ar dominates, a process that additively shifts or multiplicatively stretches the r values does not alter the pattern in the lower tail. That lower tail is a rising exponential shape, $e^{\tilde{a}r}$, as in Fig. [4c](#page-11-0).

When the $\alpha \tilde{r}^2$ term is negligible at all magnitudes,
a combination of the dominance by density index the combination of the dominance by density independent limitations in the lower tail, ar , and the dominance by the density dependent limitations in the upper tail, λe^r , yields the gamma distribution pattern
on the abundance scale, n on the abundance scale, n.

Finally, for magnitudes of r at which the $\alpha \tilde{r}^2 = (r - \mu)^2$ term dominates, the constraint σ^2 $\alpha (r - \mu)^2$ term dominates, the constraint, $\sigma^2 =$
 $\left(r - \mu^2\right)$ on the variance in process values is most im- $\langle r - \mu^2 \rangle$, on the variance in process values is most im-
portant. In this case, pattern follows a pormal distri portant. In this case, pattern follows a normal distribution, $e^{-\alpha(r-\mu)^2}$, on the *r* scale, which is a lognormal distribution on the abundance scale, *n*.

When combining numerous process values to obtain an overall net r value, approximate rotational invariance is sufficient for the pattern to be very close

to a perfect normal curve (see Introduction). When measuring net squared deviations from the mean, which is the squared radial distance, the pattern is invariant to shift and stretch of the squared radial measures, $(r - \mu)^2$.
In practice

In practice, the lognormal pattern of abundance dominates when a density independent constraint on process, r, dominates, and net process values obey rotational invariance (symmetry) with respect to the summing up of the individual processes acting on abundance.

Any theory of process that leads to those three basic invariances will follow the gamma-lognormal pattern.

Conclusions

Invariance does not replace other approaches. Rather, invariance clarifies why different approaches and seemingly different assumptions often lead to the same pattern. Common pattern arises because of simple shared invariances. Thus, as noted in the introduction, the normal distribution arises when the squared vector length of non-normal random variables distributes uniformly over the rotationally invariant circumference of the circle, with the average squared radius as the variance of the distribution 16 16 16 .

Invariance also improves the derivation and understanding of key results. For example, maximum entropy is an invariance method. However, maximum entropy hides its underlying basis in invariance, often making it difficult to apply to particular problems and to understand the results.

Another example is neutrality. Neutrality analyzes fundamental processes, such as birth, death, and migration. However, neutrality by itself often misleads, because it hides the simple invariances that link process to pattern. Combining neutrality and invariance provides deeper understanding and simpler analytical results.

Different biological processes transform inputs to outputs in different ways. But ultimately those different processes must produce the same simple invariances if they lead to the same simple patterns $24,31-39$ $24,31-39$ $24,31-39$.

Invariance therefore provides a method to compare different models of process. In a model that combines many different component processes, how do the various transformations interact to produce the invariances that define commonly observed pattern? Among the many complex component processes that may occur in a model, which truly matter? In other words, which component processes shape the defining invariances and which are irrelevant?

Suppose we wish to compare two different models that predict the same pattern. Invariance suggests two questions. First, for each model, which components define the same key invariance? Second, for the focal component processes of each model, which empirical tests would tell us the more likely match to natural processes?

The apparent simplicity of invariance can mislead about its ultimate power. For example, probability patterns express a shift and stretch invariant scaling. That affine invariant scaling provides a constant measure of information at all magnitudes.

Shift and stretch invariance seem almost trivially simple. Yet, by analyzing how repeated transformations of shift and stretch retain invariance, we obtain the most general form that expresses various affineinvariant scales (eqn [2\)](#page-5-2). That affine symmetry group defines the simple, general structure of probability patterns and their uniform measurement scales.

Knowing the general invariant form of probability patterns reveals the relations between different approaches. Invariance provides powerful methods to analyze pattern and process.

To sum up, our invariance approach is not just another one among various alternatives. Rather, it is the only way to relate process to pattern, because the essence of pattern is invariance. Only by understanding what pattern actually is and how it generally arises can one begin to formulate testable hypotheses about mechanism.

Put another way, pattern is always the interaction between, on the one hand, the generic aspects of invariance and scale that arise in all cases and, on the other hand, the particular aspects of biology that operate in each case. Without a clear view of that duality between the generic and the particular, it is easy to mistakenly attribute generic aspects of observed pattern to particular causes. To properly understand the role of specific mechanistic aspects in shaping pattern, one must evaluate pattern simultaneously from the perspectives of the generic and the particular.

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References

- 1. Lederman, L. M. & Hill, C. T. *Symmetry and the Beautiful Universe* (Prometheus Books, Amherst, N.Y., 2004). URL [http://www.loc.gov/catdir/](http://www.loc.gov/catdir/toc/ecip0417/2004008558.html) [toc/ecip0417/2004008558.html](http://www.loc.gov/catdir/toc/ecip0417/2004008558.html).
- 2. Fisher, R. A., Corbet, A. S. & Williams, C. B. The relation between the number of species and the number of individuals in a random sample of an animal population. *Journal of Animal Ecology* **12**, 42–58 (1943).
- 3. Preston, F. W. The commonness, and rarity, of species. *Ecology* **29**, 2554–283 (1948).
- 4. Baldridge, E., Harris, D. J., Xiao, X. & White, E. P. An extensive comparison of species-abundance distribution models. *PeerJ* **4**, e2823 (2016). URL <https://doi.org/10.7717/peerj.2823>.
- 5. Hubbell, S. P. *The Unified Neutral Theory of Biodiversity and Biogeography* (Princeton University Press, Princeton, NJ, 2001).
- 6. Rosindell, J., Hubbell, S. P. & Etienne, R. S. The unified neutral theory of biodiversity and biogeography at age ten. *Trends in Ecology and Evolution* **26**, 340–348 (2011).
- 7. Brown, J. H. *Macroecology* (University of Chicago Press, Chicado, 1995).
- 8. Gaston, K. J. & Blackburn, T. M. *Pattern and Process in Macroecology* (Blackwell Science, Oxford, UK, 2000).
- 9. Harte, J. *Maximum Entropy and Ecology: A Theory of Abundance, Distribution, and Energetics* (Oxford University Press, New York, 2011).
- 10. Jaynes, E. T. Information theory and statistical mechanics. II. *Phys. Rev.* **108**, 171–190 (1957).
- 11. Jaynes, E. T. Information theory and statistical mechanics. *Phys. Rev.* **106**, 620–630 (1957).
- 12. Jaynes, E. T. *Probability Theory: The Logic of Science* (Cambridge University Press, New York, 2003).
- 13. Wagensberg, J. & Valls, J. The extended maximum entropy formalism and the statistical structure of ecosystems. *Bulletin of Mathematical Biology* **49**, 531–538 (1987).
- 14. Alonso, D., Ostling, A. & Etienne, R. S. The implicit assumption of symmetry and the species abundance distribution. *Ecology Letters* **11**, 93– 105 (2008).
- 15. Frank, S. A. The invariances of power law size distributions. *F1000Research* **5**, 2074 (2016).
- 16. Frank, S. A. Common probability patterns arise from simple invariances. *Entropy* **18**, 192 (2016). URL [http://www.mdpi.com/](http://www.mdpi.com/1099-4300/18/5/192) [1099-4300/18/5/192](http://www.mdpi.com/1099-4300/18/5/192).
- 17. Solé, R. V., Manrubia, S. C., Benton, M., Kauffman, S. & Bak, P. Criticality and scaling in evolutionary ecology. *Trends in Ecology & Evolution* **14**, 156–160 (1999). URL [https://doi.org/](https://doi.org/10.1016/S0169-5347(98)01518-3) [10.1016/S0169-5347\(98\)01518-3](https://doi.org/10.1016/S0169-5347(98)01518-3).
- 18. Jordano, P., Bascompte, J. & Olesen, J. M. Invariant properties in coevolutionary networks of plant-animal interactions. *Ecology Letters* **6**, 69– 81 (2003).
- 19. Stouffer, D. B., Camacho, J., Guimerà, R., Ng, C. A. & Amaral, L. A. N. Quantitative patterns in the structure of model and empirical food webs. *Ecology* **86**, 1301–1311 (2005).
- 20. Frank, S. A. & Smith, E. A simple derivation and classification of common probability distributions based on information symmetry and measurement scale. *Journal of Evolutionary Biology* **24**, 469–484 (2011).
- 21. Frank, S. A. How to read probability distributions as statements about process. *Entropy* **16**, 6059– 6098 (2014).
- 22. Frank, S. A. The common patterns of abundance: the log series and Zipf's law. *F1000Research* **8**, 334 (2019).
- 23. Harte, J. Maximum entropy and the statevariable approach to macroecology. *Ecology* **89**, 2700–2711 (2008).
- 24. Chisholm, R. A. & Pacala, S. W. Niche and neutral models predict asymptotically equivalent species abundance distributions in high-diversity

ecological communities. *Proceeding of the National Academy of Sciences USA* **107**, 15821– 15825 (2010).

- 25. Etienne, R. S. & Rosindell, J. The spatial limitations of the current neutral models of biodiversity. *PLoS ONE* **6**, e14717 (2011).
- 26. Alonso, D. & McKane, A. J. Sampling hubbell's neutral theory of biodiversity. *Ecology Letters* **7**, 901–910 (2004).
- 27. Tribus, M. *Thermostatics and Thermodynamics: An Introduction to Energy, Information and States of Matter, with Engineering Applications* (Van Nostrand, New York, 1961).
- 28. Cover, T. M. & Thomas, J. A. *Elements of Information Theory* (Wiley, New York, 1991).
- 29. Frank, S. A. Invariant death. *F1000Research* **5**, 2076 (2016).
- 30. Frank, S. A. Measurement invariance explains the universal law of generalization for psychological perception. *Proceedings of National Academy of Sciences USA* **115**, 9803–9806 (2018).
- 31. Cohen, J. E. Alternative derivations of a species abundance relation. *American Naturalist* **152**, 165–172 (1968).
- 32. Chave, J., Alonso, D. & Etienne, R. S. Comparing models of species abundance. *Nature* **441**, E1–E2 (2006).
- 33. Adler, P. B., Hillerislambers, J. & Levine, J. M. A niche for neutrality. *Ecology Letters* **10**, 95–104 (2007).
- 34. Zillo, T. & Condit, R. The impact of neutrality, niche differentiation and species input on diversity and abundance distributions. *Oikos* **116**, 931–940 (2007).
- 35. Etienne, R. S., Alonso, D. & McKane, A. J. The zero-sum assumption in neutral biodiversity theory. *Journal of Theoretical Biology* **248**, 522–536 (2007).
- 36. Frank, S. A. The common patterns of nature. *Journal of Evolutionary Biology* **22**, 1563–1585 (2009).
- 37. McGill, B. J. Towards a unification of unified theories of biodiversity. *Ecology Letters* **13** (2010).
- 38. Matthews, T. J. & Whittaker, R. J. Neutral theory and the species abundance distribution: recent developments and prospects for unifying niche and neutral perspectives. *Ecology and Evolution*

4, 2263–2277 (2014).

39. Matthews, T. J. & Whittaker, R. J. Fitting and comparing competing models of the species abundance distribution: assessment and prospect. *Frontiers of Biogeography* **6**, 2 (2014).