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## Title

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An Algorithm for the Generation of Nuclear Spin Species and Nuclear Spin Statistical Weights

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## Abstract

This paper develops a set of algorithms for the computer generation of nuclear spin species and nuclear spin statistical weights potentially useful in molecular spectroscopy. These algorithms generate the nuclear spin species from group structures known as generalized character cycle indices (GCCI's). Thus the required input for these algorithms is just the set of all GCCI's for the symmetry group of the molecule which can be easily computed from the character table. The algorithms are executed and illustrated with examples.

In recent years chemical applications of non-numerical computational techniques are becoming important in several areas ${ }^{1-4}$. Generation of discrete structures such as graphs and combinatorial structures has been made possible with the advent and rapid growth of algorithms and computational techniques. Randic ${ }^{1-3}$ has pioneered some such chemical applications of discrete non-numerical computational techniques. In this paper we consider another important application of combinatorial algorithrns to a problem in molecular spectroscopy that has been of considerable interest for several decades.

An experiment or theory in molecular spectroscopy typically concerns with possible spectral lines and the intensities of such possible lines. The intensities of allowed inter-rovibronic transitions depend upon the nuclear spin statistical weights of the rovibronic levels. 5-6 The conventional technique for obtaining the nuclear spin statistical weights is to find the character of the representation spanned by all the possible nuclear spin functions of the molecule and then break it down into irreducible components. For a molecule containing $b_{1}$ nuclei of the same kind with $a_{1}$ spin states, $b_{2}$ nuclei of the same kind with $a_{2}$ spin states etc., there are $a_{1}{ }^{b 1} a_{2}{ }^{b} 2 \ldots$ nuclear spin functions. For ${ }^{12} \mathrm{C}$ triphenylene this number is 4096 . Thus the technique mentioned above is quite difficult for polyatomics. The objective of the present paper is to develop very efficient algorithms that would generate nuclear spin species and statistical weights from a minimal input.

The statistical weights of rotational levels in the rotational subgroup have been discussed by Placzek and Teller ${ }^{7}$, Wilson ${ }^{8-9}$ and Mizushima ${ }^{10}$. Hougen ${ }^{11}$ correlated these to the point groups of molecules.

Galbraith ${ }^{12}$ obtained the nuclear spin statistical weights using the unitary group approach and Schur's theorem. Recently, Weber 13-14 discussed the nuclear spin statistics of spherical top molecules belonging to the point groups $D_{n d}$ or $D_{n h}(n \leqslant 6)$. The present author ${ }^{15}$ recently developed a general method for the nuclear spin statistics of molecules belonging to any symmetry. In Sec. 2 we briefly review this method. The present author ${ }^{16-24}$ has been interested in group theoretical methods for problems in chemistry. Even though this paper is written with the intent of making it self-contained, a more extensive account. of preliminaries and definitions can be found in text books. 16 -19 This paper uses a theorem of Williamson 20 for characters of one dimensional representations which was recently generalized by Merris ${ }^{21}$ to characters of higher dimensional representations.

In Sec. 2 we outline the theoretical methods, and in Sec. 3 we describe the set of algorithms for nuclear spin statistical weights.

## 2. Theory

Let $D$ be the set of nuclei of the same kind in the molecule and let $R$ be the set of possible spin states of the nuclei in $D$. In this paper we shall treat each kind of nuclei (such as $\mathrm{H},{ }^{13} \mathrm{C},{ }^{19} \mathrm{~F}$, etc.) separately, obtain the nuclear spin species of each kind of nuclei and eventually the overall nuclear spin species are obtained as direct products of each kind of nuclear spin species. Let F be the set of all maps from $D$ to $R$. The image of $E$ is the set of spin functions. For example, for a set of spin $1 / 2$ nuclei, $R$ would be a set consisting of 2 elements which can be denoted by $\alpha$ (spin up) and $\beta$ (spin down). A map
$\mathrm{E}_{1}$ from D to R is shown below.

$$
\begin{aligned}
f_{1}(1) & =\underset{\sim}{\beta} \\
f_{1}(2) & =\underset{\sim}{\alpha} \\
f_{1}(3) & =\underset{\sim}{\beta} \\
f_{1}(4) & =\alpha
\end{aligned}
$$

The spin function thus generated by $f_{1}$ is $\beta \alpha \beta \alpha$. The group $G$ is the point group or more generally, the permutation inversion group of the molecule. G acts on the elements of F in accordance to the following recipe.

$$
g(f(i))=f\left(g^{-1}\right) \text { for every } i \varepsilon D
$$

To illustrate if $g$ is the permutation $(12)(34)$, then $\mathrm{g}^{-1}=(12)(34)$. Consequently, $f_{1}$ shown above is transformed as

$$
\begin{aligned}
& g f_{1}(1)=f_{1}\left(g^{-1} 1\right)=f_{1}(2)=\underset{\sim}{\alpha} \\
& g f_{1}(2)=f_{1}\left(g^{-1} 2\right)=f_{1}(1)=\underset{\sim}{\underset{\sim}{\sim}} \\
& g f_{1}(3)=f_{1}\left(g^{-1} 3\right)=f_{1}(4)=\underset{\sim}{\alpha} \\
& g f_{1}(4)=f_{1}\left(g^{-1} 4\right)={\underset{\sim}{1}}^{\sim}(3)=\underset{\sim}{\sim}
\end{aligned}
$$

Thus, the spin function $\underset{\sim}{\beta \alpha \beta \alpha}$ is transformed into $\underset{\sim \sim \sim}{\alpha \beta \alpha \beta}$ by the action of $g=(12)(34)$. Define the cycle representation of a permutation $g \varepsilon G$, containing $b_{1}$ cycles of length $1_{1} b_{2}$ cycles of length 2 , etc, as $x_{1} b_{1} b_{2} \ldots$ For example, the cycle representation of the permutation (12)(34), is $x_{2}^{2}$. The cycle index of $G$, denoted as $P_{G}$ is defined as the sum of all cycle representations of $g \varepsilon G$ divided by $|G|$, the number of elements in G. In symbols,

$$
P_{G}=\frac{1}{|G|} \sum_{g \in G} x_{1}^{b_{1} x_{2} b_{2}} \ldots \ldots
$$

In order to book-keep the number of various spin states in a spin function, let us introduce the concept of weight of a spin function. With
each $r \varepsilon R$, let us associate a weight $w(x)$, which is just a formal symbol that differentiates the various spin states in the set $R$. Then the weight of any function fer is defined as the product of the weights of the corresponding images. Symbolically,

$$
W(f)=\pi w(f(d)) .
$$

To illustrate, the map $f_{1}$ defined above which maps 4 nuclei into the spin function $\underset{\sim \sim}{\beta \alpha \beta \alpha}$ would have the weight $\alpha^{2} \beta^{2}$ if the weight of the spin states $\underset{\sim}{\alpha}$ and $\underset{\sim}{\beta}$ are $\alpha$ and $\beta$, respectively.

Let $g \rightarrow \chi(g)$ be the character of an irreducible representation $I$ in $G$. Williamson treated $x$ as the character of one dimensional representations and Merris generalized to characters of higher dimensional represenations. Define the generalized character cycle index (CCCI) corresponding to the irreducible representation $\Gamma$ whose character is $\times$ as follows.

$$
P_{G}^{x}=\frac{1}{|G|} \sum_{\varepsilon G} x(g) x_{1}^{b_{1}} x_{2}^{b_{2}} \ldots \ldots
$$

Then using the theorems of Williamson and Merris, it can be shown that the generating function for nuclear spin species which corresponds to the irreducible representation $\Gamma$ whose character is $x$, is given by:

$$
\text { Theorem 1: G.F. }=\mathrm{P}_{\mathrm{G}}^{\mathrm{X}}\left(\mathrm{x}_{\mathrm{k}} \rightarrow \sum_{r \in \mathrm{R}}(w(\mathrm{r}))^{\mathrm{k}}\right)
$$

The coefficient of a typical term $w_{1} b_{2} b_{2} \ldots$ in the G.F. gives the frequency of occurrence of the irreducible representation $r$ in the set of spin functions with the weight ${ }_{w}{ }_{1}{ }_{w} b_{2} \ldots$. The beauty of this method is that it did not require the character of all spin functions of dimension $a_{1}^{b_{1}} b_{2} \ldots$. to decompose into irreducible components. All that we need is the set of GCCI's of the group which can be obtained immediately by examining
the character table.
If one is interested only in finding the irreducible representations occurring in $r^{\text {spin }}$, the reducible representation spanned by all spin functions, then all one needs to do is to set all the weights to unity. Then theorem 1 reduces to corollary 1.

Corollary 1: The number of times the irreducible representation $\Gamma$ occurs in $r^{\text {spin }}$ is given by

$$
P_{G}^{X}\left(x_{k} \rightarrow|R|\right)
$$

where $|R|$ is the number of elements in the set $R$.

## 3. Algorithms

Theorem 1 has simplified the problem of obtaining nuclear spin species with great elegance. Yet one needs to expand several multinomials, combine them, and collect the total coefficients corresponding to all possible weights. For example, if one looks at an aggregate of bosons with spin 2 arranged in $\mathrm{O}_{\mathrm{h}}$ symmetry, the GCCI of $\mathrm{T}_{1 \mathrm{~g}}$ representation and the corresponding G.F. are shown below.

$$
\begin{aligned}
& \operatorname{GCCI}\left(T_{1 g}\right)=\frac{1}{48}\left(3 x_{1}{ }^{6}-3 x_{2}^{3}+6 x_{1}{ }^{2} x_{4}-9 x_{1}{ }^{2} x_{2}^{2}+6 x_{2} x_{4}-3 x_{1}{ }^{4} x_{2}\right) \\
& \text { G.F. }=\frac{1}{48} 3(\alpha+\beta+\gamma+\delta+\varepsilon)^{6}-3\left(\alpha^{2}+\beta^{2}+\gamma^{2}+\delta^{2}+\varepsilon^{2}\right)^{3} \\
& +6(\alpha+\beta+\gamma+\delta+\varepsilon)^{2}\left(\alpha+\alpha^{4}+\gamma^{4}+\delta^{4}\right) \\
& -9(\alpha+\beta+\gamma+\delta+\varepsilon)^{2}\left(\alpha^{2}+\beta^{2}+\gamma^{2}+\delta^{2}+\varepsilon^{2}\right)^{2}+6\left(\alpha^{2}+\beta^{2}+\gamma^{2}+\delta^{2}+\varepsilon^{2}\right) \\
& \left(\alpha^{4}+\beta^{4}+\gamma^{4}+\delta^{4}+\varepsilon^{4}\right) \\
& -3(\alpha+\beta+\gamma+\delta+\varepsilon)^{4}\left(\alpha^{2}+\beta^{2}+\gamma^{2}+\delta^{2}+\varepsilon^{2}\right)
\end{aligned}
$$

where $\alpha, \beta, \gamma, \delta$ and $\varepsilon$ are the weights of the 5 spin states. There are 210 terms in the above G.F. One has to collect the coefficients of all the 210 terms in order to obtain the nuclear spin species. Thus, one

$$
-7-
$$

can see the necessity of an algorithm for computer generation of G.E.'s and, subsequently, nuclear spin species. Let us describe a set of algorithms which will take GCCI's as inputs and generate nuclear spin species.

## A. Composition Algorithm

Let n and k be two positive integers. Then a composition of n into $k$ parts is an ordered sum

$$
n=t_{1}+t_{2}+\ldots+t_{k}
$$

where $t_{i} \geqslant 0$ 。
To illustrate, all the 10 compositions of 3 into 3 parts are shown below.

$$
\begin{aligned}
& 3+0+0,0+3+0, ~ 0+0+3,2+1+0,2+0+1, \\
& 1+2+0,1+0+2,0+1+2,0+2+1, \text { and } 1+1+1 \text {. }
\end{aligned}
$$

The algorithm described here is the one given in Nijenhuis and Wilf 22 which is based on a method which actually constructs such an arrangement. Suppose one wants to distribute n indistinguishable balls into k labelled cells such that any cell contains $0,1,2 \ldots$ up to $n$ balls, then each such distribution yields a composition. This can be accomplished by constructing the boundaries of cells as follows: Consider a big cell with two walls containing $n+k-1$ available spaces, the two walls occupying the 2 spaces in the extremities as shown below.

$$
\frac{1}{1} \quad \ldots \ldots . \overline{n+k} \frac{1}{n+k+1}
$$

The number of ways of distributing $n$ balls in the available $n+k-1$ slots can be seen to be

$$
\binom{n+k-1}{n}
$$

After distributing these balls, one can construct the cell boundaries in the rest $k-1$ available spaces. Such an arrangement contains exactly $k$ cells among which we have distributed the $n$ indistinguishable balls.

Thus, we have obtained a composition of $n$ into $k$ parts. Therefore, the number of such compositions $\mathrm{C}(\mathrm{n}, \mathrm{k})$ is given by

$$
C(n, k)=\binom{n+k-1}{n}=\binom{n+k-1}{k-1}
$$

The process described above is equivalent to finding all the possible $k-1$ - subsets from a $n+k-1$ subset. The following algorithm describes a way of generating k - subsets from a n - set in a lexicographic order. For example, we list below the 2 - subsets from a 3 - set in a lexicographic order.

12
13
23
Let $a_{1}, a_{2}, \ldots, a_{k}$ be a given $k$ - subset. Then all the $k$ - subsets are generated in a lexicographic order by the following algorithm.

Algorithm 1:
Al. 1: (Initial set)
$a_{j} \leqslant j(j=1, k) ;$ go to A1.2.
A1.2: (Subsequent sets) $h+\min \left\{j / a_{k+1-j} \neq n+1-j, j=1, k\right\}$ 。
$m_{1} \leqslant a_{k+1-h} ; \quad a_{k+1-h} \leqslant m_{1}+j \quad(j=1, h)$,
if $a_{1}=n-k+1$, stop.
Let us now illustrate the above algorithm with example 1.

Example 1: All 3-subsets from a 5-set.
A1.1: Produces $a_{1}=1, a_{2}=2 ; a_{3}=3$.
A1.2: $\operatorname{step} 1: h=1 ; m_{1}=a_{3+1-1}=3 ; \quad a_{3}=4$ 。
step 2: $h=1 ; m_{1}=4 ; a_{3}=5$
step 3: $h=2 ; m_{1}=2 ; a_{2}=3, a_{3}=4$, etc.
step 9: $h=3 ; m_{1}=2 ; a_{1}=3 ; a_{2}=4 ; a_{3}=5$
since $a_{1}=n-k+1=3$. stop.

Now let us proceed to generating compositions in a lexicographic order using the algorithm 1 by finding all the possible $\mathrm{k}-1$ - subsets from a $n+k-1$ - set. When one translates the language of subsets into the language of composition, Al modifies to A2.

Algorithm 2:
A2.1: (Initial entry): $t_{1}+n_{;} t_{i}+0($ for $i=2, k) ;$ go to $A 2.2$
A2.2: (Subsequent entries): $h \leftarrow \min \{i / t i \neq 0\}: r \leftarrow t_{h} ; t_{h} \leqslant 0 ; t_{1} \leqslant r-1$; $t_{h+1} \leftarrow t_{h+1}+l_{\text {; }}$ if $t_{k}=n_{\text {, }}$ stop.

The algorithm 2 is illustrated with example 2.

Example 2: Compositions of 3 into 3 parts.
A2. 1: $t_{1}=3, t_{2}=0, t_{3}=0$. Produces the ordered vector ( $3,0,0$ )
A2.2: step 1: $h=1 ; r=t_{1}=3 ; t_{1}=2 . t_{2}=1$. Produces $(2,1,0)$
step 2: $h=1 ; r=2 . t_{1}=1 . t_{2}=2$. Produces $(1,2,0)$
step 3: $h=1 ; x=1$. $t_{1}=0$. $t_{2}=3$. Produces $(0,3,0)$
step 4: $h=2 ; r=3, t_{2}=0, t_{1}=2, t_{3}=1$. Produces $(2,0,1)$, etc.
step 9: $h=2 ; r=1 . t_{2}=0, t_{1}=0 . t_{3}=3$. Produces $(0,0,3)$.
The algorithm A2 produces all possible terms in the multinomial expansion $\left(x_{1}+x_{2}+\ldots+x_{k}\right)^{n}$. A typical composition vector $\left(a_{1}, a_{2}, \ldots, a_{k}\right)$ such that $\sum_{i=1}^{k} a_{i}=n$, corresponds to the term $x_{1}{ }_{1} x_{2} \ldots x_{k} a_{k}$ in this multinomial expansion.
B. VECIOR Algorithm

The algorithm which we call "VECTOR" essentially picks up the compositions generated by A2 and "multiplies" them and collects the coeffients and the overall vector in an expansion of the form

$$
\left(a_{1}^{b_{1}}+a_{2}^{b_{1}}+\ldots+a_{k}^{b_{1}}\right)^{n_{1}}\left(a_{1}^{b_{2}}+a_{2}^{b_{2}}+\ldots+a_{k}^{b_{2}}\right)^{n_{2}} \ldots
$$

Let max $n_{1}, n_{2}, \ldots$ be $m$. Then the number of terms in the above expansion
is given by the number of compositions of $m$ into $k$ parts. This algorithm calls A2 for each term in the product with $n=n_{i}$, obtains the composition vector for each term, the coefficient for each term and the overall coeffient and the vector. Let the number of terms in the above product be t. Algorithm VECIOR generates the overall vector array vec $(i, j)$, $i=1, k$ and $j=1, \operatorname{lin}_{i=1}$ and the corresponding coefficients $c(j)$ in the expansion. The vector array is the total vector generated as a result of multiplication of all the terms in the expansion

$$
{ }_{i=1}^{t_{j}}\left(j_{1}^{k} a_{j}^{b_{i}}\right)^{n_{i}}
$$

For example, a typical vector ( $l_{1}, l_{2}, \ldots$ ) stands for the term $a_{1} l_{1} a_{2}{ }_{2} \ldots$ in the above expansion. The resulting coefficient is a product of several multinomial coefficients.

## C. Poly Algorithm

Poly is the main algorithm which calls the VECIOR algorithm for each term in the GCCI. This algorithm obtains the vector for each term in the GCCI and matches it with the initial vector generated by VECTOR for the initial term $\left(a_{1}+a_{2}+\ldots+a_{k}\right)^{n}$. When there is a match it adds the product of the corresponding coefficient in the GCCI and the appropriate coefficient generated by vector to the already existing inventory of coefficients. When Poly exhausts all the terms in the GCCI, the resulting vector and the corresponding coefficients give the generating function with the vector $(1,12 \ldots)$ representing the term $a_{1}^{1} a_{2} 1_{2} \ldots$ Then it calls an algorithm SPIN described below which generates the spin species from the generating function.

## D. SPIN Algorithm

SPIN algorithm takes the generating function and the value of the nuclear spin quantum number for the states whose weights are in the
generating function and gives the nuclear spin species. Let the spin values of the states whose weights are $a_{1}, a_{2}, \ldots a_{k}$ be $m_{1}, m_{2} \ldots \ldots, m_{k}$. Then the total spin quantum number of a term $a_{1} b_{1} b_{2} \ldots a_{k} b_{k}$ is $m_{1} b_{1}+m_{2} b_{2}+\ldots+m_{k} b_{k}$. The corresponding coefficient in the generating function gives the frequency of occurrance of the irreducible representation $\Gamma$ in the set of spin functions with the total spin quantum number $m_{1} b_{1}+m_{2} b_{2}+\ldots+m_{k} b_{k}$. The spin algorithm first sorts the coefficients in the generating function in accordance to their total spin quantum number. Let $C_{j}$ be the array of such coefficients sorted in accordance to the total spin quantum number. Let $\lambda_{i}$ be the multiplicity of the $i$ th species and let $d_{i}$ be the corresponding frequency of occurrence. Then the algorithm described below generates the nuclear spin species from the sorted coefficients. Let n be the total number of nuclei in the molecule whose spin species are under consideration. D1. Algorithm Spin

For Bosons
$\left[\begin{array}{l}\begin{array}{l}\lambda \text { for } i=1, n, d o \\ \lambda_{i}+2 i-1 \\ j\end{array} \\ \quad \text { if } i=1, d_{j}+2-i \\ \text { final exit }\end{array}\right.$
for Fermions
if n is odd, $\mathrm{f} \leftarrow \mathrm{n}+1$.
$\left[\begin{array}{l}\text { for } i=2, f, 2, \text { do } \\ \lambda_{i} \leftarrow i \\ j \leftarrow n+3-i \\ \text { if } i=2, d_{j}+c_{j} ; \text { else } d_{j}+c_{j}-c_{j+2}\end{array}\right.$

```
final exit
if n is even, f & n+l
for i=1, f, 2,do
\lambda
j}~n+2-
if i=l, dj * C cj; else d}\mp@subsup{d}{j}{}\leqslant\mp@subsup{c}{j}{-}\mp@subsup{C}{j}{}+
final exit.
```

E. A complete example: A1 proton spin species of rigid butane.

Algorithm 'POLY' with the help of VECTOR and A2 generate the total generating function corresponding to the Al representation of the $\mathrm{C}_{2 \mathrm{v}}$ group which is shown below.

| Coefficient | $\frac{\text { Vector }}{1}$ |
| ---: | ---: |
| 3 | $(10,0)$ |
| 15 | $(8,1)$ |
| 32 | $(7,3)$ |
| 60 | $(6,4)$ |
| 66 | $(5,5)$ |
| 60 | $(4,6)$ |
| 32 | $(3,7)$ |
| 15 | $(2,8)$ |
| 3 | $(1,9)$ |
| 1 | $(0,10)$ |

where, for example, the vector $(7,3)$ stands for $\alpha^{7} \beta^{3}$ ( $\alpha$ and $\beta$ being the weights of the spin states $\underset{\sim}{\alpha}$ and $\underset{\sim}{\beta}$ of protons) and the corresponding coefficients in the generating function is 32. The algorithm spin now sorts the coefficients according to their total spin value. In this
in the G.F. and the array $\mathrm{C}_{\mathrm{j}}$. Now algorithm Dl is called for ferminons with $n=6$, Dl sets $\lambda_{1}=1, \lambda_{2}=3, \ldots, \lambda_{6}=11$. The corresponding $d_{j}$ 's are $d_{1}=6$, $d_{2}=28, d_{3}=17, d_{4}=12, d_{5}=2$ and $d_{6}=1$. Thus the proton species which correspond to $A_{1}$ representation are

$$
1_{A_{1}}(6),{ }^{3} A_{1}(28),{ }^{5} A_{A_{1}}(17),{ }^{7}{A_{1}}(12),{ }^{9} A_{A_{1}}(2), \text { and }{ }^{11}{ }_{A_{1}}(1)
$$

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