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### **Title**

ACDOS2: AN IMPROVED NEUTRON-INDUCED DOSE RATE CODE

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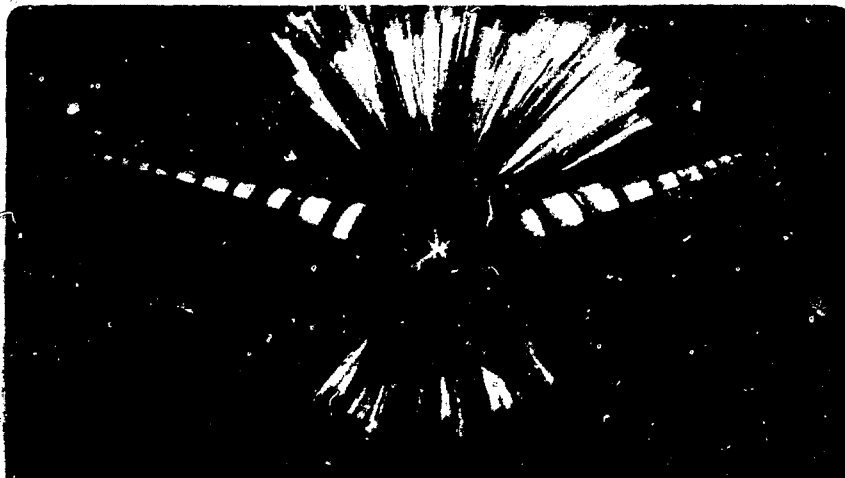
## Accelerator & Fusion Research Division

**MASTER**

ACDOS2: AN IMPROVED NEUTRON-INDUCED  
DOSE RATE CODE

Jean-Charles Lagache  
(M.S. thesis)

June 1981



Prepared for the U.S. Department of Energy under Contract W-7405-ENG-48

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ACDOS2: AN IMPROVED NEUTRON-INDUCED DOSE RATE CODE

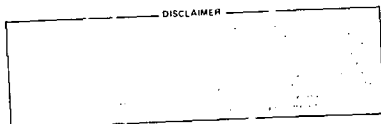
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## I. Introduction

Although it is generally recognized that activation problems will result from several of the large fusion devices now planned or under construction, it appears that such problems may have to be faced even before such devices operate. This consequence follows from the fact that neutral-beam design must necessarily precede applications by several years. A forthcoming upgrade program at LBL will require prolonged periods of testing of a deuterium 170-kev neutral beam, at 65A, with a 10% duty factor. Such a source can produce appreciable dose rates for personnel who must do maintenance work following a prolonged period of testing. To calculate the expected dose rate as a function of geometry, composition, and time after shutdown a computer code, ACDOS2, was written, which utilizes up-to-date libraries of cross-sections and radio-isotope decay data. ACDOS2 is in ANSI FORTRAN IV, in order to make it readily adaptable elsewhere.

The methodology for solving the overall problem is broken down into several steps. The first is that of determining the production of neutrons in the case of a neutral-beam injector. In practice, neutron production is not uniform but occurs during a succession of evenly spaced pulses of short time duration. During the pauses, no neutrons are produced. An effective steady neutron source term is generated by utilizing an appropriate

duty factor over the time duration of the test periods. The neutron production rate is then calculated on the basis of the energy and current of the accelerated beam.

With the unity-normalized average group fluxes assumed to be available from a previous calculation, the second step involves calculating the activities due to the activation of a target or of the walls, by the neutrons. First the flux weighted group cross-sections have to be computed by making use of a library containing microscopic cross-sections. Prerequisite to this calculation is the determination of a weighting flux for use in averaging cross-section data over the appropriate energy groups. This is accomplished by assuming a room-temperature Maxwell-Boltzmann distribution for the thermal region, a  $1/E$  distribution for intermediate energies, and an exponential function for the source-neutron groups. Arbitrary constants in the weighting fluxes are determined by integral and/or boundary conditions on the known flux for the particular energy group in question. With the weighting flux a determinable function of energy, the required flux-weighted group cross-sections can be evaluated. The microscopic cross-section library contains up to ten different neutron reactions per target nuclide, at neutron energies up to 20 MeV. The calculation of the resultant activities given the neutron production rate, the flux weighted group cross-sections, the target nuclide mass and type, and the specific times after shutdown at which induced activities are desired, can then be performed. This is done by solving the pertinent ordinary

differential multigroup equations that describes the system undergoing activation and then algebraically modifying the standard solution to make it amenable to the injector test schedule.

The final step is a calculation of the dose rates as a function of time and geometry, given the previously calculated activities. This routine interrogates a decay library and extracts the necessary gamma ray energy and intensity information. The program treats radioactive daughters accurately. The user specifies the type of geometry to be used in the actual calculations (point, sphere, cylinder-on axis, cavity) and the strategy (non-absorbing source, absorbing source), in addition to the distance from the activated component, at which the dose rate is desired.

The code can handle from one to fifty neutron groups with up to thirty target nuclides. It has an option which allows the user to substitute a neutron source term other than that from a neutral-beam injector, and a further option which allows for the evaluation of the approximate fluxes in the walls, and thereby the dose-rate from the walls. The end result is then a code that is tailored to solve injector activation problems, but contains enough versatility to be useful in solving a wide variety of general activation problems produced by neutrons of energy less than 20 MeV.

This report constitutes the user manual of ACDOS2, which is an improved version of ACDOS1.<sup>1</sup> Basically, ACDOS2 uses the same framework as ACDOS1 to calculate neutron induced activation and



corresponding dose rates as a function of geometry, composition and time after shutdown. Improvements have been made in the physical models and the versatility has been increased by the introduction of new options. Specifically, the radioactive daughters are now accurately treated, a fourth possible geometry (cavity) is supplied, both absorption and build-up in volumetric sources are taken into account, the dose due to the walls can be computed, and an option to indirectly couple a slowing-down code is provided.

The main part of the report deals with the different models and equations used by ACDOS2. Some sections (§ II, IV-a) are reproduced directly from the report on ACDOS1 <sup>1</sup>. A new user\* will find enough information to run the code by just reading the Appendices which completely describe the input required for ACDOS2.

\*ACDOS2 is available on magnetic tape. Address inquiries to: Head, Neutral Beam Development Group, Bldg. 4, Lawrence Berkeley Laboratory, Berkeley CA 94720.

## II. Neutron Production Model

The first major step in solving the overall activation problem is to supply the average neutron source strength. Alternatively, ACDOS2 makes provision for the user to directly enter the average neutron source strength or to determine from data supplied by the user, the production of neutrons in the specific case of neutral-beam injectors.

In neutral-beam testing, neutron production is not uniform but occurs during an arbitrary succession of evenly spaced injector pulses of short time duration. The neutron production coincides with the injector pulses. Figure II-1 represents the duty factor (i.e. the fraction of time in which the injector is actually on) during the test periods. INSNPS is the instantaneous number of neutrons produced per second. T1 is the length of the test periods--typically a few hours. T2 is the length of time between test periods. N is the number of test periods, following the last of which, the activities are calculated. It is assumed throughout the activation calculations that T1 and T2 do not vary during operation of the injector.

The neutrons originate from two sources within the injector system: gaseous deuterium that leaks from the neutralizer and the ion-source, adsorbs on the surfaces of the neutral-beam calorimeter and the deflected-beam dump. A (d, n) reaction occurs as D<sup>0</sup> and D<sup>+</sup> impinge upon these surfaces. The reaction is:

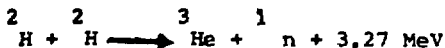
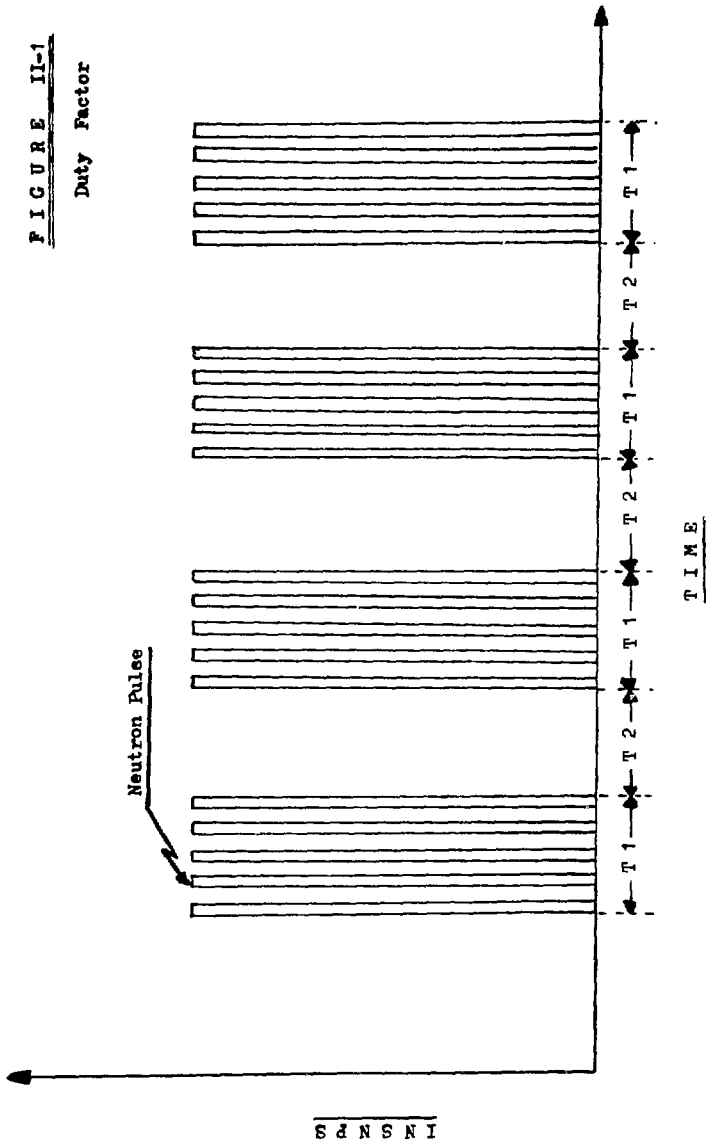


FIGURE II-1  
Duty Factor



INSNPS

TIME

The calculation that determines the instantaneous number of neutrons produced per seconds is empirical in nature and based on data taken under accelerator test conditions.<sup>2</sup> The expression used is:

$$\text{INSNPS} = 8.64 \times 10^4 \left( \frac{\text{neutron}}{\mu\text{coulomb}} \right) \times A \times 10^6 \left( \frac{\mu\text{coulomb}}{\text{second}} \right) \times F \times CF \times DF$$

INSNPS: instantaneous number of neutrons produced per second.

A: current in Amperes

DF: duty factor during the test periods

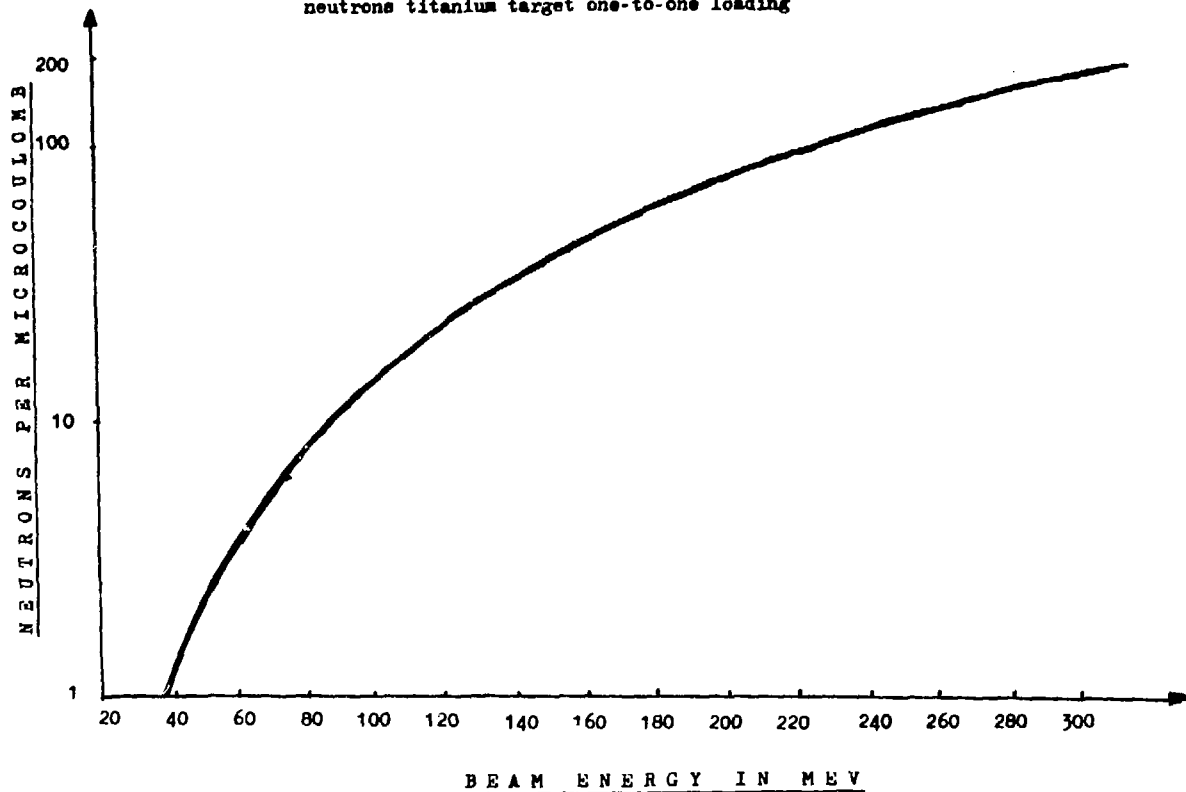
F: fraction of the beam which is monoatomic

CF: yield correction factor for voltages different from 150 kV.

The yield correction factor for voltages different than 150 kV, is taken from the thick-target yield curve<sup>3</sup> reproduced in Figure II-2. To facilitate the calculation of INSNPS during execution of the program, a power fit of the form  $y=a^x$  was applied to the points read from the curve to determine an analytical expression for CF. Three different fits were used to insure a high coefficient of determination. These expressions are incorporated into ACOS2 to calculate INSNPS.

FIGURE II-2

Theoretical thick target yield of  
neutrons titanium target one-to-one loading



### III. Fluxes

The mono-energetic source of neutrons produces a continuous spectrum by slowing-down in the shielding walls. The group fluxes in the test-cell are not calculated in ACDOS2, but rather are supplied from a previous calculation, with, for example, a Monte-Carlo code like MORSE or a transport code like ANISN. The output of such a calculation is coupled to ACDOS2 either from a tape or from the input data deck.

The coupling of ACDOS2 with a slowing-down code can then be done in two steps. First, the unity-normalized average group fluxes are computed by using the slowing down code, based on a unity source placed in the center of the test-cell. At the end of the slowing-down calculation, the space-independent fluxes inside the test cell are stored on a tape, starting with the higher energy-group flux and according to the format of ACDOS2 (E 10.3, 2X). Then, ACDOS2 is executed by reading the fluxes from the tape.

ACDOS2 provides for the calculation of the dose due to the activation of an object placed inside the test-cell or the dose due to the activation of the walls of the test-cell, which is approximated to be a hollow sphere. In this case, the fluxes in the walls are assumed to be the same as the fluxes in the test-cell, but corrected by a shape factor. This factor was calculated on the basis of one-group diffusion theory.

According to this model, the flux resulting from a point source placed in the center of a hollow sphere is constant inside a hollow sphere. This is due to the fact that the net current is zero at the inner radius. In the walls, the flux satisfies:

$$\nabla^2 \phi + K^2 \phi = 0$$

where:  $K^2 = 1/\lambda^2 = \Sigma_a/D$

and also satisfies the boundary condition:

$$\phi(R + H) = 0$$

where R is the inner radius and H the thickness of the hollow sphere. The solution to this equation is:

$$\phi(r) = A \sinh[K(R + H - r)]/Kr$$

where A is a constant.

The average flux in the wall is then:

$$\langle \phi \rangle = \frac{\int_R^{R+H} \frac{A \sinh K(R+H-r)}{Kr} 4\pi r^2 dr}{4/3\pi[(R+H)^3 - R^3]}$$

or after integration:

$$\langle \phi \rangle = \frac{3A}{[K(R+H)]^3 - [KR]^3} [\sinh(KD) + KR \cosh(KD) - (KR+KD)]$$

Defining the shape factor FF by the ratio of the flux inside the hollow sphere to the average flux in the walls, we have:

$$FF = \frac{\langle \phi \rangle}{\phi(r=R)}$$

or

$$FF = \frac{3KR}{(KR+KH)^3 - (KR)^3} \left( 1 - \frac{KR+KH}{\sinh(KH)} + \frac{KR \cosh(KH)}{\sinh(KH)} \right)$$

In ACOS2, the dose due to the walls, can be computed approximately by multiplying the fluxes inside the test-cell, available from a previous calculation, by the shape factor FF.



#### IV. Activation

##### A. Weighting flux and flux weighted group cross-sections

In order to average cross-section data over the appropriate energy groups to acquire group cross-sections, a weighting flux is required. The flux-weighted group cross section can be evaluated from the expression:

$$\bar{\sigma} = \frac{\int_{\Delta E} \sigma(E) \phi(E) dE}{\int_{\Delta E} \phi(E) dE}$$

The weighting fluxes were determined by assuming that the thermal region could be represented by a Maxwell-Boltzmann distribution with  $kT = 0.025$  eV, the intermediate energies by a function proportional to  $1/E$ , and the fast groups by an exponential function. All energy dependences are thus determined, except for arbitrary constants.

The decision to use these energy dependences came in part from the predictions of slowing-down theory plus the results of a Monte-Carlo calculation involving a neutral beam injector surrounded by thick concrete walls. The group fluxes calculated<sup>4</sup> for this case are given in Table IV-1. To verify the above assumptions, column 3 (neutrons/cm<sup>2</sup>.eV.source neutron) was plotted as a function of the arithmetic average of the group boundaries found in column 2 for groups 1 through 19 (figure IV-1).

Inspection of the curve suggests a  $1/E$  type of behavior for groups 4 through 19 since a quantity that varies as a constant

T A B L E IV-1

Group Fluxes Inside the Test Cell

GROUP	ENERGY INTERVAL (eV)	NEUTRONS	NEUTRONS
		$\frac{2}{\text{cm}^2 \cdot \text{eV} \cdot \text{SOURCE NEUTRON}}$	$\frac{2}{\text{cm}^2 \cdot \text{SOURCE NEUTRON}}$
1	2.385E+6 2.307E+6	1.013E-10	7.901E-6
2	2.307E+6 1.827E+6	2.879E-13	1.382E-7
3	1.827E+6 1.108E+6	1.959E-13	8.333E-8
4	1.108E+6 5.502E+5	1.896E-13	1.058E-7
5	5.502E+5 1.576E+5	3.510E-13	1.378E-7
6	1.576E+5 1.111E+5	6.111E-13	2.842E-8
7	1.111E+5 5.248E+4	8.172E-13	4.790E-8
8	5.248E+4 2.479E+4	1.394E-12	3.860E-8
9	2.479E+4 2.188E+4	2.550E-12	7.420E-9
10	2.188E+4 1.033E+4	3.049E-12	3.522E-8
11	1.033E+4 3.355E+3	6.552E-12	4.570E-8
12	3.355E+3 1.234E+3	1.896E-11	4.021E-8
13	1.234E+3 5.829E+2	4.673E-11	3.043E-8
14	5.829E+2 1.013E+2	1.502E-11	7.234E-8
15	1.013E+2 2.902E+1	6.481E-10	4.684E-8
16	2.902E+1 1.068E+1	1.642E-09	3.011E-8
17	1.068E+1 3.059E+0	5.184E-09	3.951E-8
18	3.059E+0 1.125E+0	1.787E-08	3.456E-8
19	1.125E+0 4.140E-1	4.327E-08	3.076E-8
20	4.140E-1 1.000E-5	2.002E-06	8.288E-7

FIGURE IV-1  
Neutrons/cm<sup>2</sup>·eV vs Neutron energy

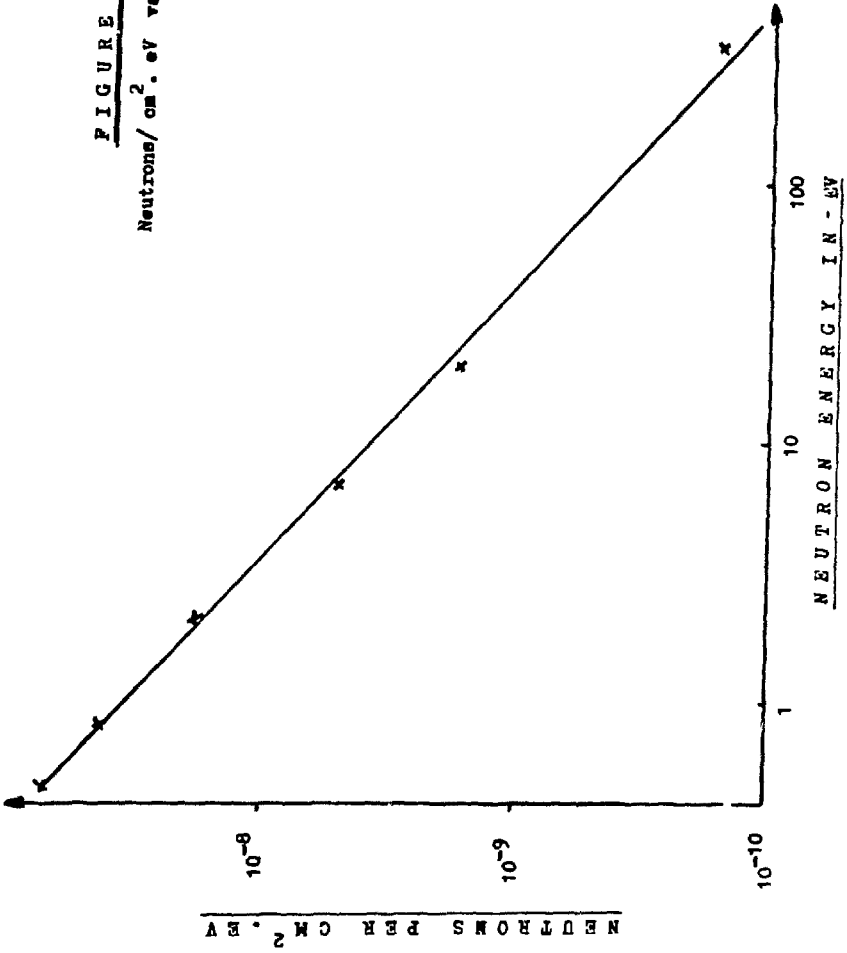


FIGURE IV-1 (Cont.)

Neutrons/cm<sup>2</sup>.eV vs Neutron energy

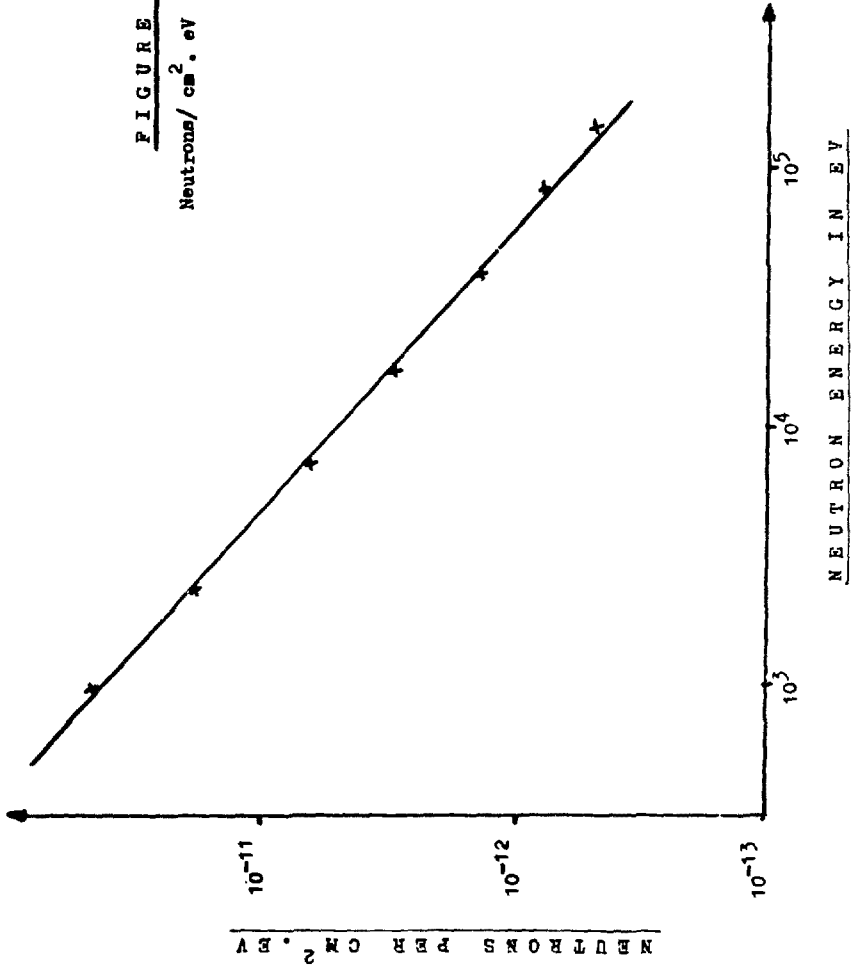
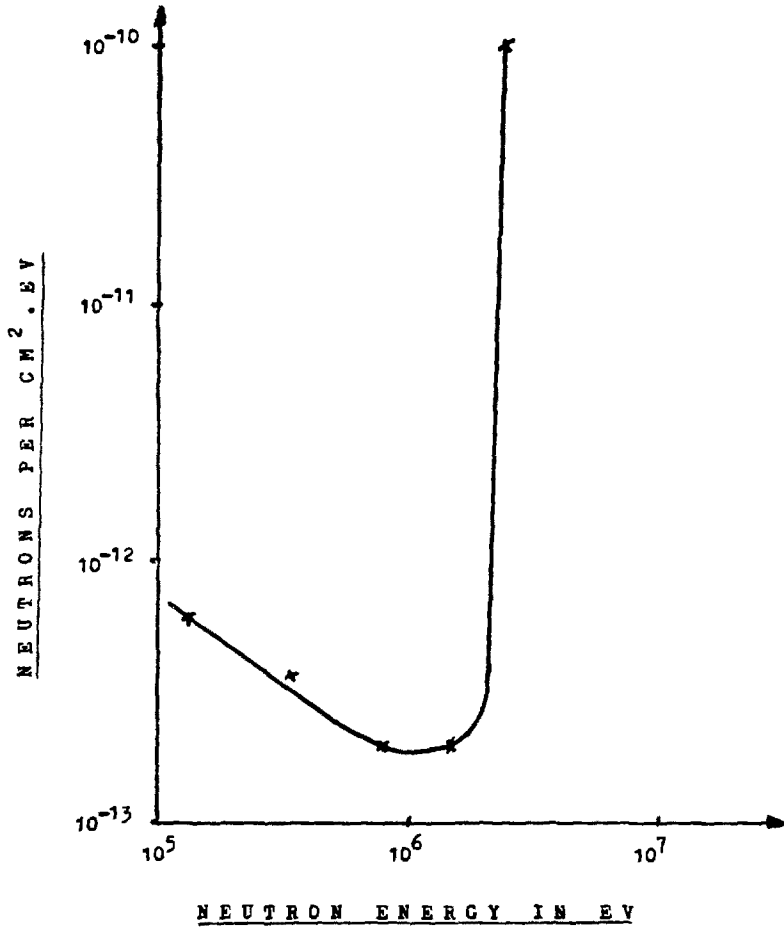


FIGURE IV-1 (Cont.)  
Neutrons/cm<sup>2</sup>.eV vs Neutron energy



over  $E$  will plot as a straight line on log-log paper and will have a slope of  $-1$ . Calculation of the slope at various points along the line did reveal a slope close to  $-1$ . At higher energies, above  $2 \times 10^4$  eV, the curve departs somewhat more from  $1/E$  behavior and finally reaches a region just below fast energies where the curve is neither  $1/E$  nor exponential. Even though this region just below the exponential part of the curve deviates from  $1/E$  behavior, it is assumed, for ease of calculations, that it does behave as  $1/E$ .

Since neutrons are "born" with a narrow group of discrete energies in the injector, as contrasted with a fission spectrum in a reactor, one would intuitively expect the fast region to appear more as a "spike" than as a continuous distribution. The graph definitely illustrates this behavior and for this reason the assumption of an exponential form for the fast groups appears justified. As stated previously, the Maxwell-Boltzmann distribution is used for describing the thermal region.

At this point, the forms of the equations for describing the weighting flux are known except for arbitrary constants. To evaluate the constants, each analytic expression for the weighting flux is integrated over the appropriate energy group and set equal to the numerical value of the corresponding group quantities in column 4 of table IV-1, which are normalized total group fluxes. For example, for the thermal region, group 20:

$$\int_{10^{-5}}^{4.14 \times 10^{-1}} C \sqrt{E} e^{-\frac{E}{0.025}} dE = 8.288 \times 10^{-7}$$

therefore:

$$C = 2.36 \times 10^{-4}$$

For an intermediate neutron group, group 17:

$$\int_{3.059}^{10.68} C/E dE = 3.951 \times 10^{-8}$$

therefore:

$$C = 3.16 \times 10^{-8}$$

The determination of the constants associated with fast groups, however, is somewhat more involved. Since the assumed form for the fast exponential is

$$\phi(E) = A e^{\frac{E_1 - E_2}{B}}$$

where  $E_1$  and  $E_2$  are known and  $E_2$  is the highest energy of that group, two equations are required to evaluate the two constants. One equation results from the usual integral condition that

$$\int_{E_1}^{E_2} A e^{\frac{E_1 - E_2}{B}} dE = \text{total flux for that group}$$

while the other condition results from a continuity of flux requirement at the boundary of the two groups, namely:

$$C/E_1 = A e^{\frac{E_1 - E_2}{B}}$$

where  $C$  is known constant from the integral condition placed on the last  $1/E$  group. The analysis is then always performed, starting with the lowest energy group. With two equations and two unknowns, the constants can be determined by solving the resultant transcendental equations. If there is another fast

group, the same method is applied. One equation results from the required integral condition on that group, and a second equation results from a continuity of flux requirement at the boundary with the preceding group.

The result of this method is that the first exponential group is always matched to the last 1/E group, and any additional fast groups are always matched to the preceding fast groups. In the case of only two groups, one thermal and one fast, a match is required at the boundary between the groups to allow evaluation of the two constants in the assumed exponential function following the Maxwellian.

Only one equation is required to evaluate 1/E groups as there is only one constant to be evaluated. Since 1/E constants are calculated solely from the integral condition placed on the equation for that group, discontinuities sometimes result in the flux at the group boundary between two adjacent 1/E or between the thermal and the first 1/E group. This is not a serious problem because the discontinuities are not significantly large. The important consideration is that the integral of the weighting-flux function over the group be equal to the total flux for that group so that the correct averaging of the flux over that same interval is maintained.

Table IV-2 shows the results of this approach to the weighting flux function determination. The integral of these functions over their proper energy range always results in the numerical value of the total flux for that group, as it should. A plot of the data in Table IV-2 (groups 1-19) is shown in Figure IV-2.



T A B L E IV-2Weighting Fluxes

<u>GROUP</u>	<u>ENERGY INTERVAL</u> (eV)	<u>EQUATION <math>\phi(X)</math></u> <u>(neutrons/cm<sup>2</sup> .source neutron)</u>
1	2.385E+6 2.307E+6	$7.07E-10 * EXP((X-2.385E6)/11184)$
2	2.307E+6 1.827E+6	$6.62E-13 * EXP((X-2.307E6)/2.42E6)$
3	1.827E+6 1.108E+6	$1.666E-7 * (1/X)$
4	1.108E+6 5.502E+5	$1.511E-7 * (1/X)$
5	5.502E+5 1.576E+5	$1.102E-7 * (1/X)$
6	1.576E+5 1.111E+5	$8.129E-8 * (1/X)$
7	1.111E+5 5.248E+4	$6.387E-8 * (1/X)$
8	5.248E+4 2.479E+4	$5.147E-8 * (1/X)$
9	2.479E+4 2.188E+4	$5.942E-8 * (1/X)$
10	2.188E+4 1.033E+4	$4.693E-8 * (1/X)$
11	1.033E+4 3.355E+3	$4.064E-8 * (1/X)$
12	3.355E+3 1.234E+3	$4.020E-8 * (1/X)$
13	1.234E+3 5.829E+2	$4.057E-8 * (1/X)$
14	5.829E+2 1.013E+2	$4.134E-8 * (1/X)$
15	1.013E+2 2.902E+1	$3.747E-8 * (1/X)$
16	2.902E+1 1.068E+1	$3.012E-8 * (1/X)$
17	1.068E+1 3.059E+0	$3.160E-8 * (1/X)$
18	3.059E+0 1.125E+0	$3.455E-8 * (1/X)$
19	1.125E+0 4.140E-1	$3.077E-8 * (1/X)$
20	4.140E-1 1.00E-5	$2.36E-4 * (X)^{1/2} * EXP(-X/0.025)$

FIGURE IV-2

Weighting Flux

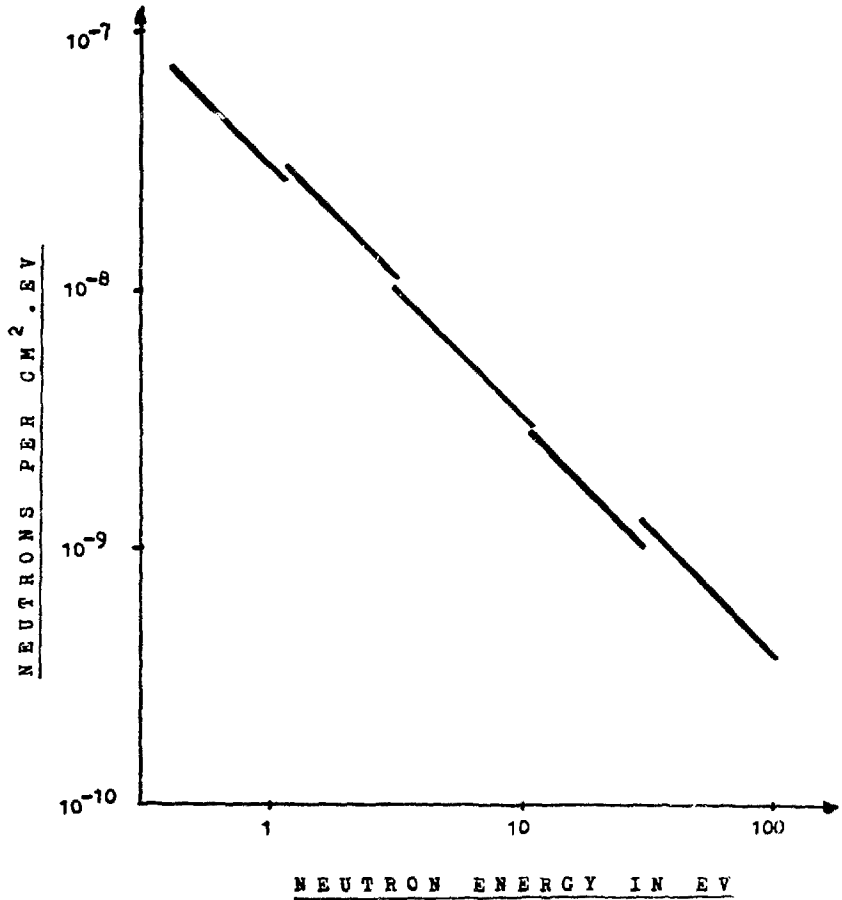


FIGURE IV-2 (Cont.)

Weighting Flux

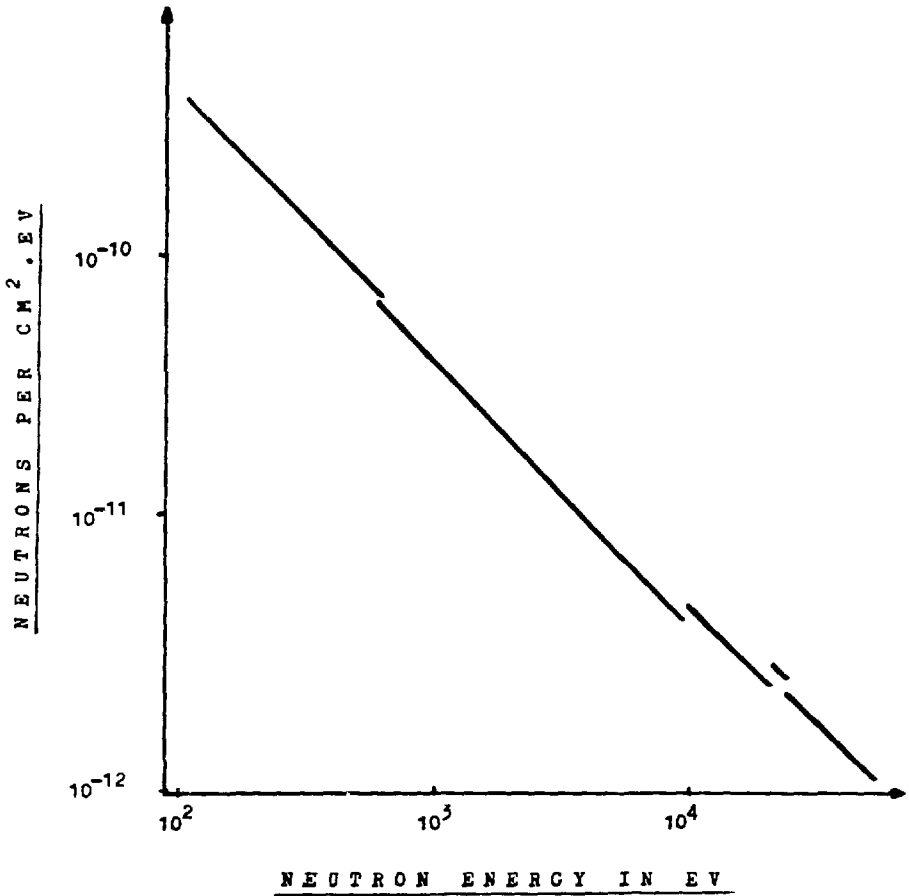
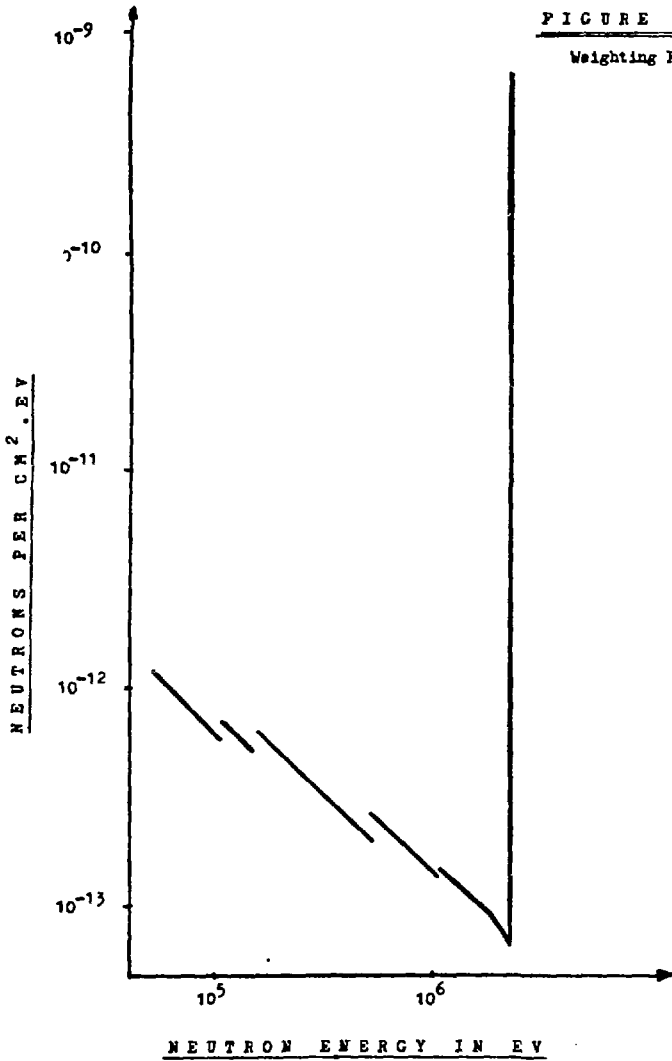


FIGURE IV-2 (Cont.)

Weighting Flux



As previously mentioned, slight flux discontinuities at boundaries between 1/E group exist. Note that the discontinuities are greater in the region just below the two fast groups. This is not surprising since that region deviates somewhat from 1/E but was assumed to be 1/E in weighting flux calculations. Comparison of Figure IV-1 with Figure IV-2 shows a close resemblance as is desired.

With the weighting flux function known for each group, the flux weighted group cross-section can be calculated from the expression:

$$\bar{\sigma} = \frac{\int_{\Delta E} \sigma(E) \phi(E) dE}{\int_{\Delta E} \phi(E) dE}$$

where  $\phi(E)$  is the previously determined weighting flux function. In the calculation of the activation, the cross-section information for the target nuclide is read from the ACTLMFE library which is a subset (without fission cross-sections) of the ACTL library.<sup>5</sup> The data consist of two numbers, an energy in MeV and an associated cross-section for that energy in barns.

To calculate  $\bar{\sigma}$ , ACOS2 first merges the group boundary values into the energy-cross-section pair data and performs a linear interpolation to calculate the cross-section at the group-boundaries.<sup>6</sup> The result is two arrays, one with energy values and the other with cross-section values. There is always a one to one correspondence. If GP(J) is an arbitrary group boundary which initially had no cross-section value associated with it. The value is obtained by a linear interpolation using the first value on each side of the group boundary assuming that

cross-sections are linearly interpolable in energy. For two arbitrary cross-sections,  $\sigma(I)$  and  $\sigma(I+1)$  whose values lie within the group boundaries  $GP(J)$  and  $GP(J+1)$ , an analytical expression is found for the line connecting them, using the common two point formula:

$$\sigma(E) - \sigma(I) = \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)]$$

so that  $\sigma(E)$  for  $E(I) \leq E \leq E(I+1)$

$$\sigma(E) = \sigma(I) + \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)]$$

Now that an analytical expression exists for  $\sigma(E)$  and  $\phi(E)$ , the product of these two functions is numerically integrated between  $E(I)$  and  $E(I+1)$  and the value kept as a running sum. This process is continued, point by point, until the upper boundary is reached. The particular form of  $\phi(E)$  used, depends upon the energy range in which the specific group boundaries lie. For example, in the Monte-Carlo calculation involving a neutral beam injector surrounded by thick concrete walls, for the thermal region, where

$$GP(J) = 10^{-5} \text{ eV} \quad \text{and} \quad GP(J+1) = 0.414 \text{ eV}$$

from Table IV-2, we have:

$$\phi(E) = 2.36 \times 10^{-4} \times \sqrt{E} \times e^{-\frac{E}{0.025}}$$

so that the expression for  $\sigma(E) \times \phi(E)$  is:

$$\left\{ \sigma(I) + \frac{\sigma(I+1) - \sigma(I)}{E(I+1) - E(I)} [E - E(I)] \right\} \times \left\{ 2.36 \times 10^{-4} \times \sqrt{E} \times e^{-\frac{E}{0.025}} \right\}$$

where  $E$  is in eV.

It is this expression that is numerically integrated between each thermal  $E(I)$  and  $E(I+1)$ .

When all numerical integrations are completed for a specific group, the running sum is divided by the total flux for that group. The quotient is the desired quantity, the flux weighted group cross section for that group. This entire process is repeated until all energy groups have been addressed. Several test cases have been run to acquire results for comparison with another averaging program which originated at LLNL.<sup>6</sup> The values obtained by this method are remarkably close to those obtained by the LLNL code for fast and intermediate neutron groups. For the thermal region, however, the above-mentioned procedure produced better results than the LLNL code.

Two further comments are in order considering the method of averaging cross-sections. First, for the case of nuclear reactions which are threshold oriented, a zero is substituted for each group cross-section whose energy range is below the threshold energy. Calculations begin only when the energy at which a cross-section value was measured, equals or exceeds the threshold energy. Typically, that first cross section value is located somewhere within the group boundaries at which the calculations begin. The contribution to that particular group cross section comes solely from evaluation of the pertinent quantities for energies greater than or equal to the first cross-section value and less than or equal to the first encountered group boundary. This is shown in figure IV-3. In this example, the reaction has a threshold located between the arbitrary group boundaries, GP(J+3) and GP(J+4). The group cross section for group (J+3)

results strictly from the evaluation of the pertinent quantities from the threshold value,  $T_v$ , to the first encountered group boundary,  $GP(J+4)$ . Even if there is only one cross-section value located in group  $(J+3)$ , the calculation proceeds since there is a cross section value associated with  $GP(J+4)$  as a result of the initial merging of group boundary energies into the energy-cross-section data. In any case, however, zeros are substituted for group cross-sections in group  $J$  through  $J+2$  in accordance with the above discussion.

The second comment concerns the calculation of the thermal group cross section. Since no provision is made for interpolating cross section values at the first group boundary, averaging calculations begin at the first encountered cross section value above the energy of the first group boundary. As a result a very small portion of the Maxwellian, generally below 0.0001 eV, is not accounted for. However, the resultant error is insignificant, since the contribution from this part of the Maxwellian is exceedingly small.

#### B. Activation

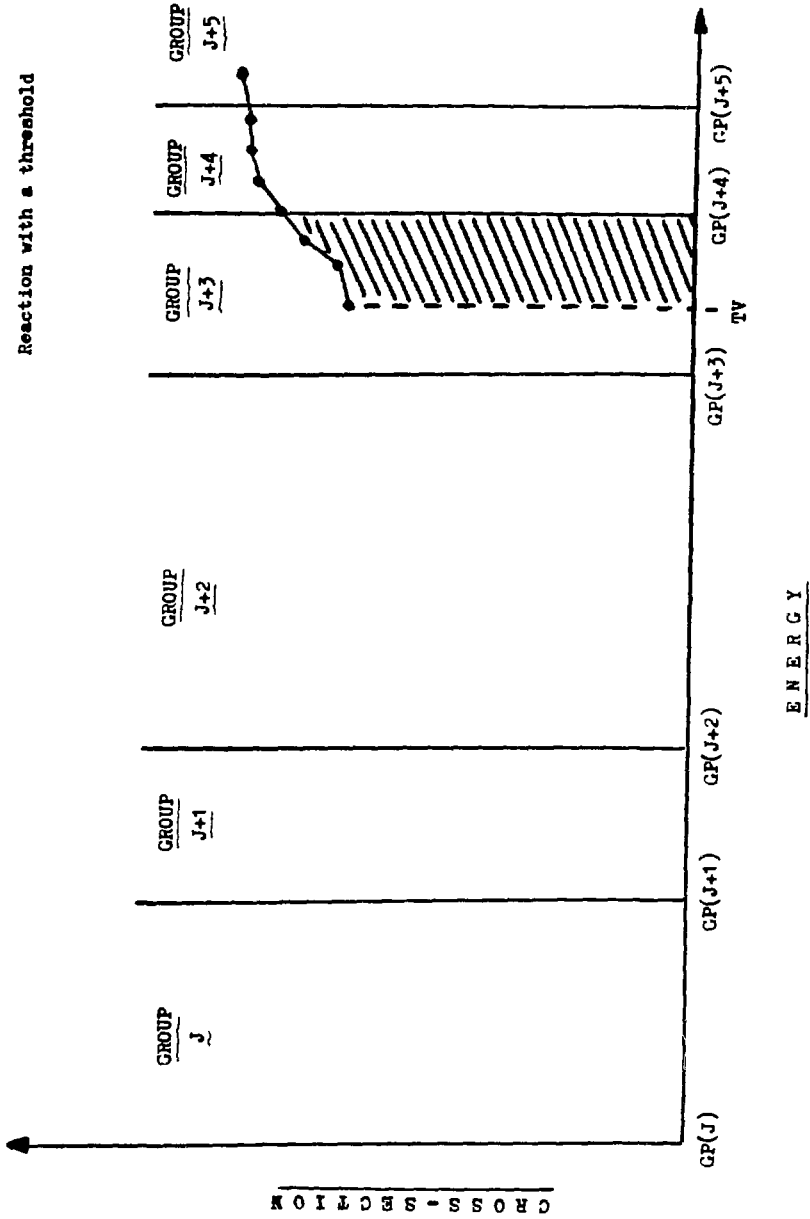
The differential equations describing the activation of target nuclides  $T$  giving radioactive product nuclides  $P$  which decay into daughter nuclides  $D$  are:<sup>10</sup>

$$\frac{dP(t)}{dt} = \sigma_T \phi T - \lambda_P P(t)$$

$$\frac{dD(t)}{dt} = \lambda_P P(t) - \lambda_D D(t)$$



FIGURE IV-3  
Reaction with a threshold



where:

- \*T(t) is the number of atoms of the target nuclide
- \*P(t) is the number of atoms of the product nuclide
- \*D(t) is the number of atoms of the daughter nuclide
- \*λ<sub>P</sub> is the decay constant of nuclide P
- \*λ<sub>D</sub> is the decay constant of nuclide D
- \*σ<sub>T</sub> is the group activation cross section for nuclide T
- \*φ is the group neutron flux.

The solutions are:<sup>10</sup>

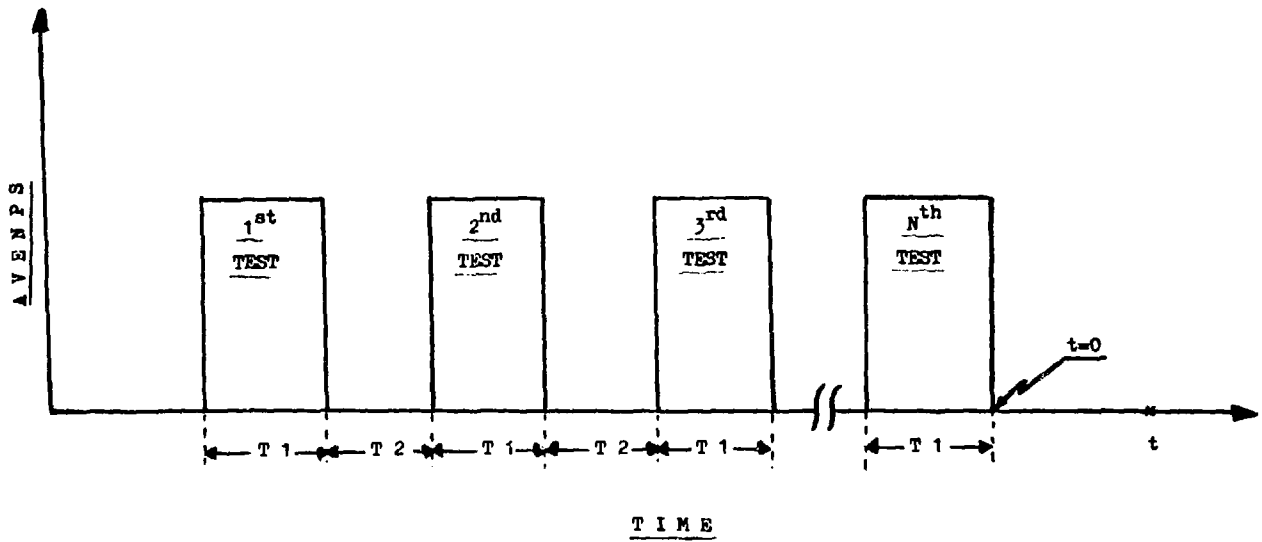
$$P(t) = \frac{\sigma_T \phi T}{\lambda_P} \left( 1 - e^{-\lambda_P t} \right)$$

$$D(t) = \frac{\sigma_T \phi T}{\lambda_D} \left( 1 - \frac{\lambda_D e^{-\lambda_P t} - \lambda_P e^{-\lambda_D t}}{\lambda_D - \lambda_P} \right)$$

These are the basic solutions that are used in subsequent activation calculations. These basic forms, however, must be algebraically modified to take into account pauses of length T<sub>2</sub> in the injector test schedule as shown in Figure IV-4.

According to this figure, let T<sub>1</sub> be the time length of one test, T<sub>2</sub> the time length of one pause, N the total number of tests, t=0 the time at which shutdown occurs and t, an arbitrarily selected time after shutdown where activities are required. It is desired to calculate the amount of an activated nuclide at t=0 due to an arbitrary neutron testing history. The method of approach is to calculate the amount at t=0 due to each test prior t=0. The total amount at t=0 will then be the sum of the contributions from each test.

FIGURE IV-4  
Modified Neutron Pulses



The contribution at  $t = 0$  of the  $i$ 'th test will be:

$$\underline{P}_i = \frac{\sigma_T \phi T}{\lambda_P} \cdot \left( 1 - e^{-\lambda_P T_1} \right) \cdot e^{-(N-i)(T_1+T_2)\lambda_P}$$

and

$$D_i = \frac{\sigma_T \phi T}{\lambda_D} \cdot \left( 1 - \frac{\lambda_D e^{-\lambda_P T_1} - \lambda_P e^{-\lambda_D T_1}}{\lambda_D - \lambda_P} \right) \cdot e^{-(N-i)(T_1+T_2)\lambda_P}$$

Then for some arbitrary time  $t$  after shutdown, the amount of nuclide  $\underline{P}$  will be given by:

$$\underline{P}(t) = \frac{\sigma_T \phi T}{\lambda_P} \cdot \left( 1 - e^{-\lambda_P T_1} \right) \cdot \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_P} \right) e^{-\lambda_P t}$$

and the amount of nuclide  $D$  by:

$$D(t) = \frac{\sigma_T \phi T}{\lambda_D} \cdot \left( 1 - \frac{\lambda_D e^{-\lambda_P T_1} - \lambda_P e^{-\lambda_D T_1}}{\lambda_D - \lambda_P} \right) \cdot \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_D} \right) e^{-\lambda_P t}$$

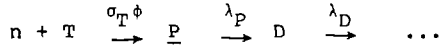
These expressions, however, represent only one group of neutrons. To generalize the relations for up to energy groups, the above calculations of the amounts of nuclides must be done for each group, that is:

$$\underline{P}(t) = \left( \sum_{k=1}^n \frac{\sigma_{Tk} \phi_k T}{\lambda_P} \right) \left( 1 - e^{-\lambda_P T_1} \right) \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_P} \right) e^{-\lambda_P t}$$

and

$$D(t) = \left( \sum_{k=1}^n \frac{\sigma_{Tk} \phi_k^T}{\lambda_D} \right) \left( 1 - \frac{\lambda_D e^{-\lambda_P T_1} - \lambda_P e^{-\lambda_D T_1}}{\lambda_D - \lambda_P} \right) \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_D} \right) e^{-\lambda_P t}$$

For the general case of the reaction:



the specific activities of the product nuclide and the daughter D in becquerels, t seconds after shutdown will be:

$$\lambda_P \underline{P}(t) = \left( \sum_{k=1}^n \sigma_{Tk} \phi_k^T \right) \left( 1 - e^{-\lambda_P T_1} \right) \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_P} \right) e^{-\lambda_P t} \quad (1)$$

and:

$$\lambda_D^D(t) = \left( \sum_{k=1}^n \sigma_{Tk} \phi_k^T \right) \left( 1 - \frac{\lambda_D e^{-\lambda_P T_1} - \lambda_P e^{-\lambda_D T_1}}{\lambda_D - \lambda_P} \right) \left( \sum_{i=1}^N e^{-(N-i)(T_1+T_2)\lambda_D} \right) e^{-t\lambda_D} \quad (2)$$

Equation is used in ACDOS2 for calculating the activity due to the product nuclides. Then, the activity of the daughters, if any, is computed from Equation 2.

## V. Dose-rate calculation

The final step is to calculate dose rates as a function of geometry and time after shutdown.

### A. Source

Using the previously calculated total activities as part of the input information, ACDOS2 first reads a decay library, a shortened version of LEVDEC, which is a sublet of ENSL<sup>7</sup> (See Appendix 1), for the energy and multiplicity associated with each gamma ray produced by a specific activated nuclide. A search for radioactive daughters is also made. During the interrogation of the library, ACDOS2 checks the mode of decay of the radioactive products and daughters. If the mode of decay is  $\gamma$ ,  $\beta^+$  or  $\beta^-$  decay, the reaction is considered. If not, the reaction is skipped and interrogation continues. Secondly, when the decay is a  $\beta^+$  desintegration, two annihilation  $\gamma$  of 0.511 MeV each are automatically included in the dose rate calculation, with the appropriate multiplicities. The case of the same product nuclide created from two different target nuclides is allowed for in the LEVDEC-read algorithm (see Appendix 1). Any daughter of a radioactive daughter is assumed to be stable.

Next, the code calculates an effective photon flux that yields an absorbed dose of 2.5 mrem/h in soft tissue for any

gamma ray or annihilation quantum read from the decay library. The latter calculation was accomplished by first fitting thirteen curves to the data<sup>8</sup> shown in Table V-1. Twelve of the curves are of the form  $ax^b$  and one of the form  $ae^{bx}$ . A large number of curves were chosen so as to accurately reproduce the data over a wide energy range. As a result, all coefficients of determination arising from the curve fitting process are in excess of 0.99. Table V-2 shows the result of the curve fittings. The appropriate equation for calculating the effective particle flux, FLUXE, is chosen according to the value of E, the energy of a particular gamma ray, as read from the decay library.

The source strength divided by the flux per unit dose rate can then be computed in  $\text{cm}^2 \cdot \text{mrem/h}$  as follow

$$S_0 = \frac{\text{MULT} \times 2.5 \times \text{ACT}}{\text{FLUXE}}$$

where MULT is the multiplicity of the gamma ray, FLUXE is the effective particle flux corresponding to a dose rate of 2.5 mrem/h in soft tissue, ACT is the previously calculated activity in Bq .

For a volumetric source of volume VOLU, the volumetric gamma ray source strength divided by the flux per unit dose rate is then:

$$\text{SOU} = S_0 = \frac{\text{MULT} \times 2.5 \times \text{ACT}}{\text{VOLU} \times \text{FLUXE} \times \text{VOLU}}$$

<u>E(MeV)</u>	<u>DOSE RATE OF 2.5 MREM/H</u> <u>2</u> <u>CORRESPONDS TO # QUANTA/ CM .SEC</u>
0.010	956
0.015	2310
0.020	4320
0.030	9980
0.040	17400
0.050	23100
0.060	25200
0.080	23200
0.100	18600
0.150	10800
0.300	5020
0.400	3660
0.500	2920
0.600	2440
0.800	1880
1.000	1550
1.250	1330
1.500	1130
2.000	912
3.000	686
4.000	559
5.000	460
6.000	420
8.000	339
10.000	284

TABLE V-1

Effective Photon  
Fluxes



T A B L E V-2

Curve Fittings of FLUXE .

<u>ENERGY INTERVAL</u> (MeV)	<u>PARTICLE FLUX EQUATION</u> FLUXE(X)
0.010 0.020	21495330.64*(X**2.17594470)
0.020 0.040	11460233.22*(X**2.0136022)
0.040 0.050	1036918.975*(X**1.2698653)
0.050 0.060	96497.38032*(X**0.4772399)
0.060 0.080	11225.12221*(X**(-0.2874410))
0.080 0.150	55482.78326*EXP(X*(-10.9132069))
0.150 0.400	1330.569205*(X**(-1.1036085))
0.400 0.600	1462.198966*(X**(-1.0004513))
0.600 1.000	1546.997798*(X**(-0.8891058))
1.000 1.500	1549.091320*(X**(-0.7797676))
1.500 3.000	1508.205757*(X**(-0.7188784))
3.000 6.000	1488.362190*(X**(-0.7051622))
6.000 10.000	1656.852874*(X**(-0.7649757))

B. Geometry

The second consideration is that of geometry. ACDOS2 provides four possible geometries (point source, sphere, cylinder, hollow sphere) and two possible strategies (non absorbing source, absorbing source) using the following formulae.

Non-absorbing Source

Point source

The point source approximation gives the following dose rate:

$$DSR = \frac{S}{4 D^2} = \frac{MULT \times 2.5 \times ACT}{FLUXE \times 4 \times D^2}$$

where the variables are the same as previously defined and D is the radial distance to the point where the dose rate is desired.

Spherical source

Outside a non-absorbing sphere the dose rate is given by<sup>8</sup>:

$$DSR = \frac{SOU}{(R+D)} \left( (2R(R+D) - (D - 2RD) \ln\left(\frac{2R+D}{D}\right)) \right)$$

where R is the radius of the sphere, D is the distance from the surface of the sphere to the observer.

Cylindrical source

On the axis of a cylinder of radius R and height H, at a distance D from one end, the dose rate is:<sup>8</sup>

$$DSR = \frac{SOU}{4} \left( (D+H) \left( \ln\left(1 + \frac{R^2}{(1+D)^2}\right) + \frac{2R}{H+D} \arctan\left(\frac{H+D}{R}\right) - D \left( \ln\left(1 + \frac{R^2}{D^2}\right) + \frac{2R}{D} \arctan\left(\frac{D}{R}\right) \right) \right) \right)$$

Hollow sphere

At the center of a hollow sphere of inner radius R and thickness D the dose rate is:

$$DSR = SOU \times H$$

• Self-absorbing source

Analytical formulae independent of the nature of the nuclide were developed for computing the constants needed in self-absorption calculations. Therefore, self-absorption for every gamma ray energy, read from the decay library, can be calculated without having to know the emitter nuclide species .

The mass attenuation factor AMU can be computed as a function of the gamma ray energy E in MeV, by using:

$$AMU = 0.0488 E^{-0.4633}$$

Table V-3 shows that this formula is conservative, since it gives a somewhat smaller value than the experimental values found for any nuclide.

The attenuation factor MU in  $cm^{-1}$  is then given by:

$$MU = \frac{AMU \times TMASS}{VOLUME}$$

where TMASS is the total mass of the target and VOLUME its volume, assuming that these values are consistent with the density of the material.

The build-up factor B can be calculated by using a Taylor expansion

$$B = A e^{-\alpha_1 \mu t} + (1-A) e^{-\alpha_2 \mu t}$$

which later will enable an integration of this factor over the geometry.

TABLE V-3  
Mass Attenuation Coefficient (cm<sup>2</sup>/g)

	<u>Gamma ray Energy (MeV)</u>							
	<u>0.15</u>	<u>0.3</u>	<u>0.5</u>	<u>1.0</u>	<u>1.5</u>	<u>2.0</u>	<u>5.0</u>	<u>10.0</u>
<u>Analytical</u> <u>function</u>	.1175	.0852	.0673	.0488	.0404	.0354	.0232	.0168
<u>H</u>	.2650	.2120	.1730	.1260	.1030	.0876	.0510	.0321
<u>Be</u>	.1190	.0945	.0773	.0565	.0459	.0394	.0234	.0161
<u>C</u>	.1340	.1060	.0870	.0636	.0518	.0444	.0270	.0194
<u>Al</u>	.1340	.1030	.0840	.0614	.0500	.0432	.0282	.0229
<u>Cu</u>	.2060	.1080	.0820	.0585	.0476	.0418	.0316	.0305
<u>Pb</u>	1.840	.3560	.1450	.0684	.0512	.0457	.0426	.0489
<u>Concrete</u>	.1390	.1070	.0870	.0635	.0517	.0445	.0287	.0229

The coefficient A is estimated by using:

$$\begin{array}{ll} A = 44.1 - 10.2 \times E & \text{if } E \leq 3 \text{ MeV} \\ A = 16.2 - 0.9 \times E & \text{if } 3 < E \leq 8 \text{ MeV} \\ A = 29 - 2.5 \times E & \text{if } 8 < E \leq 10 \text{ MeV} \\ A = 4 & \text{if } E > 10 \text{ MeV} \end{array}$$

Table V-4 shows that this formulation is conservative since the calculated A is larger than the experimental values for all the nuclides.

The coefficient  $\alpha_1$  is computed according to the following formula:

$$\begin{array}{ll} \alpha_1 = 0.063 \times E - 0.182 & \text{if } E \leq 2 \text{ MeV} \\ \alpha_1 = -0.02 \times E - 0.015 & \text{if } 2 < E \leq 8 \text{ MeV} \\ \alpha_1 = -0.0085 \times E - 0.107 & \text{if } E < 8 \text{ MeV} \end{array}$$

This expression is conservative, since the calculated values of  $-\alpha_1$  are larger than the experiment ones (Table V-5).

The coefficient  $\alpha_2$  is calculated by using:

$$\begin{array}{ll} \alpha_2 = -0.032 \times E + 0.048 & \text{if } E \leq 1 \text{ MeV} \\ \alpha_2 = -0.001 \times E + 0.017 & \text{if } 1 < E \leq 2 \text{ MeV} \\ \alpha_2 = 0.012 \times E - 0.009 & \text{if } 2 < E \leq 4 \text{ MeV} \\ \alpha_2 = 0.00125 \times E + 0.034 & \text{if } 4 < E \leq 8 \text{ MeV} \\ \alpha_2 = 0.012 \times E - 0.052 & \text{if } E > 8 \text{ MeV} \end{array}$$

Table V-6 shows that this analytical form can be applied conservatively to a great number of nuclides, since the calculated values are bigger than the experimental values.

T A B L E V-4  
Parameter A of the Taylor  
Expansion of the Build-up

Factor

	<u>Gamma Ray Energy (MeV)</u>							
	0.5	1	2	3	4	6	8	10
<u>Analytical</u> <u>function</u>	39.000	33.900	23.700	13.500	12.600	10.800	9.000	4.000
<u>Al</u>	38.911	28.782	16.981	10.583	7.526	5.713	4.716	3.999
<u>Fe</u>	31.379	24.957	17.622	13.218	9.624	5.867	3.243	1.747
<u>Sn</u>	11.440	11.426	8.783	5.400	3.496	2.005	1.100	0.708
<u>Pb</u>	1.677	2.984	5.421	5.580	3.897	0.926	0.368	0.311
<u>Concrete</u>	30.225	25.507	18.089	13.640	11.460	10.780	8.972	4.015

T A B L E V-5

Parameter  $-\alpha_1$  of the Taylor expansion

of the Build-up factor

	<u>Gamma Ray Energy (MeV)</u>							
	<u>0.5</u>	<u>1</u>	<u>2</u>	<u>3</u>	<u>4</u>	<u>6</u>	<u>8</u>	<u>10</u>
<u>Analytical</u>	0.1505	0.1190	0.0560	0.075	0.095	0.1350	0.1750	0.1920
<u>Function</u>								
<u>Al</u>	0.1002	0.0682	0.0459	0.0407	0.0397	0.0393	0.0384	0.0390
<u>Fe</u>	0.0684	0.0609	0.0463	0.0443	0.0470	0.0615	0.0750	0.0990
<u>Sn</u>	0.0180	0.0427	0.0535	0.0744	0.0952	0.1373	0.1729	0.1920
<u>Pb</u>	0.0308	0.0350	0.0348	0.0542	0.0847	0.1786	0.2370	0.2402
<u>Concrete</u>	0.1482	0.0723	0.0425	0.0320	0.0260	0.0152	0.0130	0.0288

T A B L E V-6

Parameter of the Taylor expansion

of the Build-up Factor

	<u>Gamma Ray Energy (MeV)</u>							
	0.5	1	2	3	4	6	8	10
<u>Analytical</u>	0.0320	0.0160	0.0150	0.0270	0.0390	0.0415	0.0440	0.0680
<u>function</u>								
<u>Al</u>	-0.0631	-0.0297	0.0027	0.0251	0.0386	0.0435	0.0443	0.0413
<u>Fe</u>	-0.0374	-0.0246	-0.0053	-0.0009	0.0018	-0.0019	0.0212	0.0663
<u>Sn</u>	0.0319	0.0160	0.0150	0.0208	0.0260	-0.0150	-0.0179	0.0155
<u>Pb</u>	0.3094	0.1349	0.0438	0.0061	-0.0238	-0.0464	-0.0586	-0.0278
<u>Concrete</u>	-0.1058	-0.0184	0.0085	0.0202	0.0245	0.0293	0.0298	0.0684



The next task concerns the geometry and the calculation of the self-absorption over the entire volume of the source.

Spherical source.

The dose rate at a distance D outside a spherical volume source of radius R can be represented by the expression:<sup>8</sup>

$$DSR = \frac{SOU \times B \times R}{\pi} \cdot G(p, MU \times R)$$

where  $p = \frac{D + R}{R}$  and the function  $G(p, MU \times R)$  can be approximated by the following analytical formula:

$$G(p, MU \times R) = \exp - ((0.342 \times MU \times R + 2.07322) \times \ln p + \text{Min} ( 60 \times MU \times R - 0.144 , 46.37 \times MU \times R - 0.035 , 28.78 \times MU \times R + 0.3168 , 14.55 \times MU \times R - 0.144 , , 8.17 \times MU \times R + 1.3966 , 4.72 \times MU \times R + 1.742 ))$$

where D and R are in centimeters.

Table V-7 shows the differences between this analytical form and an exact computer calculation of G.

Using the "buildup factor method"<sup>8</sup>, the Taylor expansion of B:

$$B = A e^{-\alpha_1 \mu t} + (1-A) e^{-\alpha_2 \mu t}$$

can be combined with G to give:

$$DSR = \frac{SOU \times R}{\pi} ( A \times G(p, (1 + \alpha_1) \times MU \times R) + (1 - A) \times G(p, (1 + \alpha_2) \times MU \times R)$$

Cylindrical source.

In the case of the calculation of the dose rate on the axis on the end of a cylinder, such a function as G for the

TABLE V-7  
8

Comparison of G tabulated and

G analytical

		<u>MU x R</u>						
		<u>0.4</u>	<u>0.8</u>	<u>2</u>	<u>4</u>	<u>8</u>	<u>10</u>	<u>15</u>
<u>p=1.25</u>	<u>Analytical</u>	5.719E-1	4.497E-1	2.576E-1	1.446E-1	8.055E-2	6.828E-2	5.371E-2
	<u>Tabulated</u>	6.011E-1	4.732E-1	2.718E-1	1.528E-1	8.526E-2	7.244E-2	5.725E-2
<u>p=1.5</u>	<u>Analytical</u>	3.918E-1	3.080E-1	1.763E-1	9.884E-2	5.492E-2	4.650E-2	3.646E-2
	<u>Tabulated</u>	3.923E-1	3.082E-1	1.763E-1	9.869E-2	5.446E-2	4.598E-2	3.580E-2
<u>p=2.0</u>	<u>Analytical</u>	2.158E-1	1.695E-1	9.689E-2	5.423E-2	3.001E-2	2.536E-2	1.978E-2
	<u>Tabulated</u>	2.094E-1	1.642E-1	9.366E-2	5.255E-2	2.922E-2	2.474E-2	1.930E-2
<u>p=3.0</u>	<u>Analytical</u>	9.306E-2	7.304E-2	4.169E-2	2.327E-2	1.280E-2	1.079E-2	8.358E-3
	<u>Tabulated</u>	9.000E-2	7.045E-2	4.000E-2	2.233E-2	1.230E-2	1.036E-2	7.992E-3
<u>p=5.0</u>	<u>Analytical</u>	3.226E-2	2.529E-2	1.441E-2	8.011E-3	4.377E-3	3.679E-3	2.823E-3
	<u>Tabulated</u>	3.190E-2	2.495E-2	1.414E-2	7.857E-3	4.306E-3	3.625E-3	2.813E-3
<u>p=10.0</u>	<u>Analytical</u>	7.661E-3	5.999E-3	3.407E-3	1.886E-3	1.020E-3	8.540E-4	6.471E-4
	<u>Tabulated</u>	7.924E-3	6.197E-3	3.511E-3	1.947E-3	1.054E-3	8.807E-4	6.695E-4

spherical case, exists.<sup>8</sup> But since this function depends on four parameters, no simple analytical form could be found.

An approximate solution in the axial direction was obtained by substituting for the cylindrical volume source a truncated cone of the same height as the cylinder, H, and the apex of which is at the point where the dose rate is calculated.<sup>9</sup> The equations for the attenuation of radiation for a source in the form of a truncated cone can then be used. The upper limit of the dose rate will be found for a truncated cone of angle, at the apex

$$\theta = \theta_1 = \arctan \left( \frac{R}{D} \right)$$

(The small base of the cone is then the closest base of the cylinder) and the lower limit dose will be found for

$$\theta = \theta_2 = \arctan \left( \frac{R}{D+H} \right)$$

(The large base of the cone is then the farthest base of the cylinder).

The most conservative attitude would be to use the expression yielding the upper limit of the dose rate. But such an attitude results in a gross overestimate, especially when comparing the dose rate from a sphere and a cylinder geometrically similar. As a result of such considerations, it was decided that a truncated cone should be used, which has an angle at the apex equal to:

$$\theta = \arctan \left( \frac{R}{D+H/3} \right)$$

If the height of the cylinder satisfies the relation:

$$H \leq 3/MU$$

then the dose rate is:<sup>9</sup>

$$DSR = \frac{B \times SOU}{2 \times MU} \left( 1 - \cos\theta - E_2(MU \times H) + \cos\theta \times E_2\left(\frac{MU \times H}{\cos\theta}\right) \right)$$

where  $E_2$  is the exponential integral function of the second kind and can be approximated by:<sup>10</sup>

$$E_2(X) = \exp(-X) \times (1/(X+2) + 1/(X+2)^3)$$

Using the Taylor expansion of B in the "build-up factor method"<sup>9</sup> gives:

$$DSR = (SOU/2) \times (A/(1+\alpha_1) \times MU) \times (1 - \cos\theta - E_2((1+\alpha_1) \times MU \times H)) + \cos\theta \times E_2((1+\alpha_1) \times MU \times H / \cos\theta) + (1-A)/(1+\alpha_2) \times MU \times (1 - \cos\theta - E_2((1+\alpha_2) \times MU \times H) + \cos\theta \times E_2((1+\alpha_2) \times MU \times H / \cos\theta))$$

When the height of the cylinder satisfies the relation:

$$H > 3/MU$$

the contribution of radiation only from a partial cylinder of height:

$$H' = 3/MU$$

is taken into account, assuming that only the first three mean-free paths of cylinder material will contribute appreciably to the dose rate which then is:<sup>9</sup>

$$DSR = \frac{SOU \times B}{2 \times MU} (1 - \cos\theta)$$

or using the "buildup factor method":

$$DSR = SOU \times (1 - \cos\theta) \times (A/2 \times MU \times (1+\alpha_1)) + (1-A)/(2 \times MU \times (1+\alpha_2))$$

Hollow sphere

The dose rate at the center of an absorbing sphere of radius R is:<sup>9</sup>

For an absorbing cavity of inner radius R and thickness H, the dose rate at the center is then:

$$DSR = \frac{B \times SOU}{MU} \left( 1 - e^{-MUxH} \right)$$

Using the "buildup factor" method this expression becomes:

$$DSR = SOU \times \left( \frac{A(1 - e^{-MUx(1+\alpha_1)xH}}{(1+\alpha_1)xMU} \right) + \frac{(1-A)x(1 - e^{-MUx(1+\alpha_2)xH}}{(1+\alpha_2)xMU} \right)$$

## VI. Conclusion

The first application of ACDOS2 was to the dose-rate produced by the target calorimeter of a TFTR-upgraded neutral-beam injector (see Appendix 2). About 55 mrem/h, one hour after shutdown, were found at 10 cm from the surface of a calorimeter made of copper, and around 2 mrem/h for a calorimeter made of molybdenum. From the point of view of dose-rates to personnel, molybdenum is thus preferable to copper. Even though ACDOS2 is designed to be specific to the case of injectors, it can also be used for the calculation of the entire test cell as described by Asmiller et al.<sup>11</sup>

Some improvements are possible in the models used in ACDOS2. In particular, the temperature of the calorimeter should be taken into account in the calculation of the neutron-source strength. Secondly, tabulated functions, that are approximated by analytical functions, could be more accurately defined. Such changes, however, would not drastically improve the code since the actual tabulated data have been computed by analytic approximations to within 5%.

For some applications, hand calculations were performed to verify the accuracy of the results given by ACDOS2. For example, the dose at the center of the injector test cell due to the activation of the concrete walls (Appendix 2) comes almost entirely from one product-  $^{28}\text{Al}$ . The hand calculation gave, at the time of shutdown for  $^{28}\text{Al}$ , an activity of  $1.02 \times 10^{10}$  Bq to be compared

with  $1.21 \times 10^{10}$  Bq given by the code, and a dose-rate of 4.07 mrem/h by hand to be compared with 4.05 mrem/h by the code. In the case of a thin disk of copper, which is a special case of a cylinder, the dose-rate found by a hand computation also matched the one found by running the code. Additionally, the accuracy of ACDOS2 should be evaluated by running some benchmark problems whose results could be compared with other codes. In the future, dose-rate measurements obtained from the test facility at LBL will allow an experimental verification of ACDOS2 predictions.

#### Acknowledgments

I am indebted to Robert J. Howerton, of LLNL, who provided the ACTLMFE and LEVDEC libraries, and to Gregory S. Keney, who is the author of the ACDOS1 code.

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Appendix 1

Code description

- Subroutine description
- Subroutine level chart
- Input data description
- Output description
- Tape description
  - General structure
  - ACTLMFE description
  - LEVDEC description



Subroutine description

ACDOS2 is written in a structured form, using one executing program to call up a number of subroutines. Variably dimensioned arrays are used, where possible, to conserve memory. The name and purpose of each subroutine is discussed below.

- 1) INPUT: for entering all variable input data into computer memory. All data is printed out for user confirmation. Error tests are provided to check input data.
- 2) ARRAYIN: for entering all array data. Array data is printed for user confirmation.
- 3) SOURCE: given the current, voltage, beam fraction, and duty factor, calculates the instantaneous and average number of source neutrons produced per second during a test. This subroutine is bypassed if the user wishes to use his own source term.
- 4) GRFLUX: uses the calculated values of the average number of source neutrons produced per second during a test, or the user supplied source term, and calculated the average flux for each group by taking the product of the unity normalized total group flux and source neutron term.
- 5) WTFLUX: determines the constants for the assumed weighting flux functions by applying integral and/or group

boundary constraints to the pertinent equations. Also prints out the integrals of the weighting flux functions over their appropriate energy intervals for user verification.

6) ACTVAT: takes the following parameters and calculates activities according to the equation of section IV.

- number of tests
- time length of the tests
- time length of the pauses
- number of kilograms of a particular nuclide
- specific times after shutdown
- average group fluxes
- flux weighted group cross sections

The result of each activation calculation is printed so the user can determine what reaction is the most significant for a particular target nuclide. Moreover, two running sums are maintained in order to print out the total activity from the activation of a particular target nuclide and the total system activity.

ACTVAT calls the following subroutines.

\*AVRAGE: determines an analytical expression for  $\sigma(E)$  over a specific  $\Delta E$  and numerically integrates the product of  $\sigma(E)$  and  $\phi(E)$ , the weighting function, over the group energy interval. It then takes the sum

of the integrated products over the unity normalized total flux for the same group to calculate the flux weighted group cross sections.

\*SMOOTH: first determines if the energy associated with the first energy-cross-section pair read from the ACTLMFE library is below the greatest group boundary. If not, the particular reaction is skipped and the next one considered. If so, the routine then merges the group boundary energies into the energy cross-section data read from the ACTLMFE library and linearly interpolates to find the value of the cross section at the group boundary.

\*POSITN: positions the file marker in the ACTLMFE library.

\*POSIT2: Same as POSITN

7) DOSRTE: using the previously calculated total activities, computes dose rates as a function of geometry and time after shutdown. The result of each dose-rate calculation is printed and two running sums of dose rate values are printed, one for a particular target nuclide and one for the system as a whole.

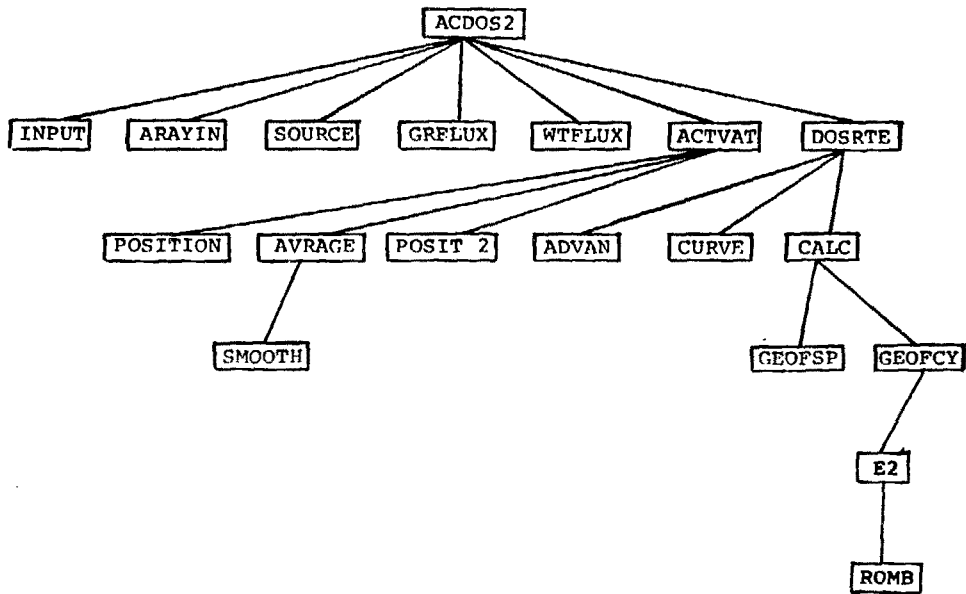
Since the CDC compilers do not permit the BACKSPACE command with formatted tapes, the interrogation of the decay library performed in DOSRTE is somewhat complicated. In order to cope with the case of the same product nuclide created from two or

more different target nuclides, and with the search for radioactive daughters, alternative solutions to the use of BACKSPACE had to be found. The first case is handled by putting the product nuclides in increasing order and by checking, during the interrogation of the library, to see if the same product appears twice or more. For the search for radioactive daughters, the tape containing the library is rewound and the interrogation starts again from the top. Possible radioactive daughters are detected by checking the previously stored information on the parents. No grand-daughter's search is done. DOSRTE calls the following subroutines:

- \*CURVE: given the energy E of the gamma ray, calculates the effective particle flux that yields a unit absorbed dose in soft tissue, FLUXE.
- \*ADVAN: for positioning the file marker in the LEVDEC library
- \*CALC: calculated dose rates using one of the four geometries, given: FLUXE, activity in Bq, multiplicity of the gamma ray or annihilation quantum, and proper dimensions.

CALC can call the following subroutines:

- GEOFSP: calculates the geometrical correction in the case of an absorbing sphere.
- GEOFCY: calculates the geometrical correction in the case of an absorbing cylinder by calling the subroutine E2 which estimates the exponential integral function of the second species using subroutine ROMB.



Subroutine level chart

Input data description

<u>card #</u>	<u>Variable description</u>	<u>columns</u>	<u>format</u>
1	<p><u>OPTION:</u>                      -if 0.000 ACDCS2 will calculate the average number of neutrons                      -any other value will be used as the average number of neutrons</p>	1 to 14	E14.7
2	<p><u>A</u> current in ampere (only if <u>OPTION=0.000</u>)</p> <p><u>V</u> voltage in kilovolts (only if <u>OPTION=0.000</u>)</p> <p><u>F</u> fraction of the beam which is monoatomic (only if <u>OPTION=0.000</u>)</p> <p><u>DF</u> duty factor (only if <u>OPTION=0.000</u>)</p> <p><u>T1</u> time length of test in hours</p> <p><u>T2</u> time length of pauses in hours</p> <p><u>N</u> number of tests</p>	<p>1 to 10</p> <p>13 to 22</p> <p>25 to 34</p> <p>37 to 46</p> <p>49 to 58 (if <u>OPTION=0</u>) 1 to 10 (if <u>OPTION≠0</u>)</p> <p>61 to 70 (if <u>OPTION=0</u>) 13 to 22 (if <u>OPTION≠0</u>) 72 to 75 (if <u>OPTION=0</u>) 25 to 27 (if <u>OPTION≠0</u>)</p>	<p>E10.3</p> <p>E10.3</p> <p>E10.3</p> <p>E10.3</p> <p>E10.3</p> <p>E10.3</p> <p>I 3</p>
3	<u>NOEGPS</u> number of energy groups (up to 50)	1 to 2	I2
4	<u>NOPAS</u> number of points after shutdown (up to 12)	1 to 2	I2
5	<p><u>R</u> radius of the sphere or of the cylinder or inner radius of the cavity (in m)</p> <p><u>D</u> distance from point source or surface of the sphere or the cylinder. Zero for inside a cavity (in m)</p>	<p>1 to 10</p> <p>13 to 22</p>	<p>E10.3</p> <p>E10.3</p>

	H height of the cylinder or thickness of the cavity (in m)	25 to 34	E10.3
	<u>IGEOM</u> 1 for point source 2 for outside a sphere 3 for a cylinder on the axis 4 for at the center of a cavity	37	I1
	<u>ISAB</u> 0 for non absorbing source 1 for absorbing source	38	I1
6	<u>NONUCL</u> number of target nuclides (up to 30)	1 to 2	I2
7	<u>NF</u> number of fast neutron groups	1 to 2	I2
8 & 9 if needed	<u>STAFS</u> (I) (I from 1 to NOEPAS) Times after shutdowns in hours (6 per card).	1 to 72	NOEPAS x (E10.3, 2x)
Above to as needed (NONUCL cards: one per target nuclide)	<u>ZNAM</u> 1(I) first part of the name of the Ith nuclide	1 to 10	A10
	<u>ZNAM</u> 2(I) second part of the name of the Ith nuclide	11 to 20	A10
	<u>IDNO</u> (I) ID number (1000Z+A) of the Ith nuclide	22 to 26	I5
	<u>MASS</u> (I) number of kilograms of the Ith nuclide	27 to 36	E10.3
Above to as needed	<u>GP</u> (I) (I from 1 to NOEGPS + L) group boundaries in MeV in ascending order	1 to 72	(NOEGPS+1) x(E10.3, 2x)
One Above	<u>IFLUX</u> : option on the fluxes 0 fluxes read from input data deck 1 fluxes read from tape 8	1 to 2	I2
	<u>IWALL</u> option on the place of the activation 0 activation of an object inside the test cell. 1 activation of the wall	4 to 5	I2
	<u>KWALL</u> (Only if IWALL=1) inverse of the relaxation length of the wall material in cm <sup>-1</sup> .	6 to 15	E10.3

<u>Above to as needed</u>	BFLUX (I) (only if IFLUX=0) (I from 1 to NOEGPS), Unity normalized group fluxes in $\text{cm}^{-2} \text{s}^{-1}$ .	1 to 72	NOEGPS x (E10.3, 2X)
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Output description

Prior to the execution of activation calculations, all input informations are printed out for user verification or recheck in case of abortion of the code. The following heading is used:

"THE FOLLOWING DATA HAVE BEEN ENTERED INTO MEMORY".

The results are then printed out under the following heading:

"THE FOLLOWING HAS BEEN CALCULATED BY ACDOS2".

The instantaneous and average number of neutrons produced per second are printed out. The integrals of the weighting flux functions over their appropriate energy intervals are printed out for user verification.

Calculated activities are printed out for each target nuclide under the two different headings:

TIME (H)	-	-
TARGET PRODUCT	ACT (T1)	... ACT(T12)
-	-	-

and:

TIME (H)	-	-
TARGET	ACT (T1)	ACT (T12)
-	-	-

Calculated dose-rates are printed out in the same way including daughters. The first heading is:

TIME (H)		-	-
TARGET	PRODUCT	DSR (T1)	DSR (T12)
-	-	-	-
-	-	-	-
PRODUCT	DAUGHTER	DSR (T1)	DSR (T12)

The second heading giving the dose due to each target nuclide is:

TIME (H)	-	-
TARGET	DSR (T1)	DSR (T12)

The final heading announces the total dose:

TIME (H)	-	-
SYSTEM	DSR(T1)	DSR (T12)
	-	-

## Tape description

### General structure

ACDOS2 was put on a nine-track magnetic tape comprising two files:

1. /NAME/ **ACDOS2/ LIBS**, XXXXX. containing the ACTLMFE and the condensed LEVDEC libraries.

2. /NAME/ACDOS2/SOURCE, XXXXX. containing the source listing (see Appendix 3)

- NAME is the name of the owner of the tape (optional when reading)
- XXXXX is the registration number of the tape

### ACTLMFE description

The ACTLMFE library contains 20,177 lines. The format of the file is:

<u>Record</u>	<u>Column</u>	<u>Variable</u>	<u>Format</u>
1	1-6	ZA (1000Z+A)	I6
	7-13	IGNORE	7X
	14-24	TARGET MASS (amu)	E11.4
	25-35	IGNORE	11X
	36-46	LEVEL OF TARGET (MeV)	E11.4
	47	IGNORE	1X
	48-58	TARGET HALF-LIFE (S)	E11.4
2	1 -2	REACTION ID NUMBER	I2
	3-8	IGNORE	6X

	9-20	Q VALUE (MeV)	E12.4	
	21-32	ZA OF PRODUCT	I6	
	33-44	LEVEL OF PRODUCT (MeV)	E12.4	
	45-56	PRODUCT HALF-LIFE (s)	E12.4	
	3	1-3	NUMBER OF ENERGY- CROSS-SECTION PAIRS (NP)	I3
	3	1-66	EN(K), CS(K) (K=1,NP)	6(E11.4
et seq				
After the last line of EN, CS pairs	72	REACTION SEPARATOR SENTINEL	71X, I1	

The above pattern is repeated for each reaction. The last reaction is for  $^{240}\text{V} (n, \gamma)$ .

LEVDEC Description

LEVDEC is a decay-mode library based on the more detailed library ENSL. The format for each set is:

<u>record</u>	<u>column</u>	<u>variable</u>	<u>format</u>
first of each set	1-6	AZ (1000Z+A)	I6
	7-17	LEVEL	E11.4
	18-21	PARITY	F4.1
	22-26	SPIN	F .1
	27-37	HALF-LIFE (s)	E11.4
	38-41	NUMBER of DECAY MODE TO FOLLOW (NMODE)	I3
2nd and seq. records to NMODE	1-38	IGNORE	38X
	39-40	MODE OF DECAY	I2

41-47	ZA of DAUGHTER	I7
48-58	LEVEL of DAUGHTER	E11.4
59-69	PROBABILITY OF DECAY TO THAT LEVEL	E12.4

Twelve modes of decay are possible for each nuclide:

<u>Decay identifier</u>	<u>Mode of decay</u>
1	neutron
2	proton
3	deuteron
4	triton
5	He <sup>3</sup>
6	$\alpha$
7	$\gamma$
8	$\beta^+$
9	$\beta^-$
10	EC
18	unresolved EC + $\beta^+$
99	No decay-stable ground state

A shortened version of LEVDEC was obtained by deleting the stable nuclides (half-life bigger than  $10^{50}$  s) and the nuclides with a very short half-life (smaller than 1 second) which are not taken into account in ACDOS2. For example, for hydrogen, LEVDEC gives:

1001	0.	1.0	.5	1.000E+SI	1		
					99	0.0	0
1002	0.	1.0	.0	1.000E+51	1		
					99	00	0
1003	0.	1.0	.5	.388E+09	1		
					9	20030	.1E+01

when the new version gives only:

1003 0.	1.0	.5	.388E+09	1		
				9	2003 0.	.1E+01

Such a procedure allows for condensing LEVDEC from three files to one, which reduces the running time of ACDOS2.

Appendix 2

Sample Problems

1. Calorimeter of a TFTR-upgraded neutral beam injector: copper vs molybdenum.
2. TFTR neutral beam injector shielding walls.

Calorimeter of a TFTR-upgraded neutral beam injector

Copper

Assuming a spherical geometry the calorimeter is made of 372 kg of copper. From the "chart of nuclides", 9th edition, 1966:

<u>Isotope</u>	<u>Atomic %</u>	<u>Weight%</u>	<u>Mass (kg)</u>
Cu <sup>63</sup>	69.1	68.45	117.36
Cu <sup>65</sup>	30.9	31.55	254.64

Since the density of copper is 8.96 g/cm<sup>3</sup> the volume of the sphere is:

$$VOLU = \frac{372000}{8.96} = 4.15 \times 10^5 \text{ cm}^3$$

and its radius

$$R = \left( \frac{3 \times VOLU}{4 \times \pi} \right)^{1/3} = 21.48 \text{ cm}$$

The observer is standing at a distance D= 10 cm from the sphere. The fluxes, read from a tape, were provided by a Monte-Carlo calculation (MORSE) involving a neutral beam injector surrounded by thick concrete walls (Table IV-1).

The following parameters are used:

A	current (A)	65
V	voltage (kV)	170
F	beam fraction	0.5
DF	duty factor	0.1
T1	length of test (h)	8
T2	length of test (h)	16
N	number of test	7



NOEGPS	number of energy groups	20
NOPAS	number of points after shutdown	12
IGEOM	geometry designator	2
ISAB	strategy designator	1
NONUCL	number of nuclides	2
NF	number of fast groups	2
IFLUX	option on fluxes	1
IWALL	option on walls	0

Table A shows the input data deck and table B the output listing.

Molybdenum

Assuming the same conditions as for the copper case, the only parameters to be changed concern the sphere. Using 554 kg of molybdenum from the "chart of nuclides" 9th edition, 1966, we have:

<u>Isotope</u>	<u>Atomic %</u>	<u>Weight %</u>	<u>Mass (kg)</u>
Mo <sup>92</sup>	14.8	14.24	78.91
Mo <sup>94</sup>	9.1	8.93	49.45
Mo <sup>95</sup>	15.9	15.75	87.24
Mo <sup>95</sup>	16.7	16.75	92.80
Mo <sup>97</sup>	9.5	9.63	53.34
Mo <sup>98</sup>	24.4	24.67	136.69
Mo <sup>100</sup>	9.6	10.03	55.57

Since the density of Molybdenum is  $10.2 \text{ g/cm}^3$ , the volume of the sphere is:

TABLE A  
Input Data Deck for Copper

0.0000000E+00						
20 .650E+02	.170E+03	.100E+01	.010E+01	.000E+01	.100E+02	7
12 .215E+00	.100E+00	.040E+00	21			
2						
2						
0.000E+1	0.100E+1	0.200E+1	0.300E+1	0.400E+1	0.500E+1	
0.600E+1	0.700E+1	0.800E+1	0.900E+1	1.000E+1	2.000E+1	
	COPPER-65 29065	1.1735E+02				
	COPPER-63 29063	2.5465E+02				
1.000E-11	4.140E-07	1.125E-06	3.059E-06	1.008E-05	2.902E-05	
1.013E-04	5.829E-04	1.234E-03	3.355E-03	1.033E-02	2.198E-02	
2.479E-02	5.248E-02	1.111E-01	1.576E-01	5.502E-01	1.108E+00	
1.827E+00	2.307E+00	2.385E+00				
1	J					

T A B L E B

Output Listing for Copper

THE FOLLOWING DATA HAVE BEEN FILTERED INTO MEMORY  
 .....

AMPERES	KILCVOLTS	EFFICIENCY	DUTYFACTOR	I1	I2	N
.650E+02	.170E+03	.100E+01	.100E+00	.000E+01	.160E+02	7

NUMBER OF ENERGY GROUPS-NCEGCS 20  
 NUMBER OF TARGET NUCLIDES-MNNUCL 2  
 NUMBER OF POINTS AFTER SHUTDOWN-NCPAS 12

ISAR=1

IGSON=2

YOU ARE USING AN ABSORBING SOURCE

THIS SOURCE IS A SPHERE  
 ITS RADIUS IS RW .015E+00 METERS  
 THE OBSERVER IS STANDING AT A DISTANCE DT .100E+00 METERS FROM THE SURFACE OF THIS SPHERE

NUMBER OF GROUPS BECAUSIES-NG	21
NUMBER OF FAST NEUTRON GROUPS-NF	2

SPECIFIC TIMES AFTER SHUTDOWN TIME

TIME 1	0.
TIME 2	.100E+01
TIME 3	.200E+01
TIME 4	.300E+01
TIME 5	.400E+01
TIME 6	.500E+01
TIME 7	.600E+01
TIME 8	.700E+01
TIME 9	.800E+01
TIME 10	.900E+01
TIME 11	.100E+02
TIME 12	.200E+02

OPTION ON THE FLUXES-FLUX= 1  
 THE FLUXES WILL BE REPORTED TO TAPE A

OPTION ON THE PART OF THE ROOM WHERE ACTIVATIONS,  
 TAPES PLACE-INWALL= 0  
 ACTIVATIONS ON AN OBJECT PLACED INSIDE THE ROOM

UNIT NORMALIZED FLUXES INSIDE THE ROOM

GROUP 1	.750E-05
GROUP 2	.110E-05
GROUP 3	.430E-07
GROUP 4	.100E-06

GROUP 5           .11378F-06  
 GROUP 6           .2442F-07  
 GROUP 7           .4750F-07  
 GROUP 8           .7820F-07  
 GROUP 9           .9427F-07  
 GROUP 10          .3177F-07  
 GROUP 11          .4570F-07  
 GROUP 12          .4011F-07  
 GROUP 13          .3247F-07  
 GROUP 14          .7754F-07  
 GROUP 15          .4444F-07  
 GROUP 16          .2011F-07  
 GROUP 17          .7541F-07  
 GROUP 18          .3457F-07  
 GROUP 19          .1078F-07  
 GROUP 20          .8777F-06

GROUP BOUNDARIES (MEMS)

BOUNDARY 1       .2785F+01  
 BOUNDARY 2       .2707F+01  
 BOUNDARY 3       .1177F+01  
 BOUNDARY 4       .1109F+01  
 BOUNDARY 5       .5502F+00  
 BOUNDARY 6       .1576F+00  
 BOUNDARY 7       .1111F+00  
 BOUNDARY 8       .2748F-01  
 BOUNDARY 9       .7485F-01  
 BOUNDARY 10      .2188F-01  
 BOUNDARY 11      .1031F-01  
 BOUNDARY 12      .2344F-02  
 BOUNDARY 13      .1774F-02  
 BOUNDARY 14      .5929F-02  
 BOUNDARY 15      .1013F-02  
 BOUNDARY 16      .2502F-04  
 BOUNDARY 17      .1049F-04  
 BOUNDARY 18      .1045F-04  
 BOUNDARY 19      .2127F-04  
 BOUNDARY 20      .6140F-06  
 BOUNDARY 21      .1000F-10

NAME	IE-NUMBER	CLASSIFICATION
CCORP-47	29077	.11378F-07
CCORP-45	29077	.11378F-07

ALOPS IS NOW READY TO RUN  
 .....

THE FOLLOWING HAS BEEN CALCULATED BY ACOS3  
 .....

INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND     .7645E+13  
 AVERAGE NUMBER OF ACTINONS PRODUCED PER SECOND         .7765E+12

INTEGRALS OF THE FITTING FUNCTION OVER THE GROUP INTERVALS  
 GROUP LOWER BOUNDARY UPPER BOUNDARY REQUIRED INTEGRAL CALCULATED INTEGRAL

20	.100E+04	.414E+00	.829E-06	.829E-06
19	.434E+00	.117E+01	.129E-07	.130E-07
18	.117E+01	.306E+01	.544E-07	.344E-07
17	.306E+01	.107E+02	.317E-07	.195E-07
16	.107E+02	.290E+02	.217E-07	.101E-07
15	.290E+02	.101E+03	.468E-07	.468E-07
14	.101E+03	.577E+03	.727E-07	.727E-07
13	.577E+03	.127E+04	.766E-07	.766E-07
12	.127E+04	.576E+04	.407E-07	.407E-07
11	.576E+04	.102E+05	.417E-07	.417E-07
10	.102E+05	.235E+05	.352E-07	.352E-07
9	.235E+05	.248E+05	.742E-07	.742E-07
8	.248E+05	.525E+05	.785E-07	.380E-07
7	.525E+05	.111E+06	.479E-07	.479E-07
6	.111E+06	.198E+06	.286E-07	.286E-07
5	.198E+06	.570E+06	.158E-06	.158E-06
4	.570E+06	.111E+07	.108E-06	.108E-06
3	.111E+07	.183E+07	.833E-07	.833E-07
2	.183E+07	.221E+07	.139E-06	.139E-06
1	.221E+07	.229E+07	.705E-05	.740E-05

THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVITIES IN MG

TIME (HR)	UO	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)
29062	24027	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29064	29062	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29067	27059	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29065	29067	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04	.174E+04
29043	27060	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29067	27060	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29065	27060	.427E+10	.404E+10	.393E+10	.377E+10	.344E+10	.326E+10	.301E+10	.292E+10	.277E+10	.262E+10	.248E+10

TIME (HR)	UO	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)
18251	18251	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29063	18251	.437E+10	.404E+10	.393E+10	.377E+10	.344E+10	.326E+10	.301E+10	.292E+10	.277E+10	.262E+10	.248E+10

TIME (HR)	UO	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)
29065	29064	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29066	29064	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29065	29064	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C
29065	29061	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C	.C

27065	28065	.C	.0	.0	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
27065	27067	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
27065	27062	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
27065	27066	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
27065	27066	.151E+10	.435E+06	.105E+02	.740E-01	.104E-04	.799E-08	.860E-12	.248E-15	.713E-19	.205E-22	.591E-26	.231E-61						

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
1845E1	ACT(F11)	ACT(F12)	ACT(F13)	ACT(F14)	ACT(F15)	ACT(F16)	ACT(F17)	ACT(F18)	ACT(F19)	ACT(F10)	ACT(F11)	ACT(F12)
27065	.551E+10	.435E+06	.105E+02	.740E-01	.104E-04	.799E-08	.860E-12	.248E-15	.713E-19	.205E-22	.591E-26	.231E-61

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
SYSTEM	ACT(F11)	ACT(F12)	ACT(F13)	ACT(F14)	ACT(F15)	ACT(F16)	ACT(F17)	ACT(F18)	ACT(F19)	ACT(F10)	ACT(F11)	ACT(F12)
	.591E+10	.404E+10	.203E+10	.362E+10	.244E+10	.326E+10	.308E+10	.292E+10	.277E+10	.262E+10	.240E+10	.145E+10

THE FOLLOWING OUTPUT DATA ARE CALCULATED DOSE RATES IN MREM/H

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27063	27065	.C	.C	.0	.0	.0	.0	.0	.C	.0	.C	.0
27063	27060	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27063	27060	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27063	28062	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27063	27062	.C	.0	.0	.0	.0	.3	.0	.0	.0	.0	.C
27064	27062	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27063	27064	.551E+10	.422E+02	.479E+02	.444E+02	.421E+02	.398E+02	.377E+02	.358E+02	.339E+02	.321E+02	.187E+02
PRODUCED DUGHTER	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27054	30064	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
1845E1	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27063	.551E+10	.422E+02	.455E+02	.469E+02	.444E+02	.421E+02	.398E+02	.377E+02	.358E+02	.339E+02	.321E+02	.187E+02

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
1845E1	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27065	27062	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27062	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27062	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27064	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27065	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27064	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27064	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27065	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27065	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
27065	27065	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
PRODUCED DUGHTER	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27055	33064	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C

TIME(H)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	DS(F11)	DS(F12)	DS(F13)	DS(F14)	DS(F15)	DS(F16)	DS(F17)	DS(F18)	DS(F19)	DS(F10)	DS(F11)	DS(F12)
27065	.804E+01	.271E-02	.766E-06	.152E-09	.552E-11	.190E-16	.458E-20	.137E-23	.379E-27	.109E-30	.214E-34	.123E-66

$$\text{VOLUME} = 554000 / 10.2 = 5.43 \times 10^5 \text{ cm}^3$$

and its radius

$$R = (3 \times \text{VOLUME} / 4 \times \pi)^{1/3} = 23.49 \text{ cm}$$

The only parameters changed compared with the previous calculation are:

NONUCL number of nuclides 7

R radius of the sphere 23.49 cm

and of course the target nuclides ID and mass.

Table C lists the input data deck and Table D shows the output listing.

### Discussion of results

Thermal-neutron activation was found in both cases to produce the dominant dose rates. In the case of the copper beam dump, the activation of  $^{63}\text{Cu}$  to  $^{64}\text{Cu}$  (12.7 h), which undergoes a  $\beta^+$  decay producing two 0.511 MeV  $\gamma$  rays, constitutes the main dose. In the case of the molybdenum beam dump, after an hour to allow the  $^{101}\text{Mo}$  (14.6 mn) and its daughter  $^{101}\text{Tc}$  to decay,  $^{99}\text{Mo}$  (66.7 h) as a product of  $^{98}\text{Mo}$ , becomes the principal dose-rate producer. However, the dose rate from the Mo dump is only about 5% of that from the Cu dump for the first 10 hours, even though the former is 30% larger volumetrically





TABLE C  
Input Data Deck for Molybdenum

0,000000E+00						
.650E+02	.170E+03	.100E+01	.010E+01	.900E+01	.169E+02	7
20						
12	.235E+00	.100E+00	.000E+00	21		
7						
2						
0.000E+1	0.100E+1	0.200E+1	0.300E+1	0.400E+1	0.500E+1	
0.600E+1	0.700E+1	0.800E+1	0.900E+1	1.000E+1	2.000E+1	
MOLYBONUM-97	42097	53.340E+00				
MOLYBONUM-95	42095	87.240E+00				
MOLYBONUM-92	42092	78.900E+00				
MOLYBONUM-94	42094	49.450E+00				
MOLYBONUM-98	42098	13.669E+01				
MOLYBONUM-96	42096	92.600E+00				
MOLYBONUM-100	42100	55.570E+00				
1.000E-11	4.140E-07	1.125E-06	3.059E-06	1.068E-05	2.902E-05	
1.013E-04	5.823E-04	1.234E-03	3.355E-03	1.033E-02	2.188E-02	
2.479E-02	5.249E-02	1.111E-01	1.576E-01	5.562E-01	1.109E+00	
1.827E+00	2.307E+00	2.385E+00				
1 0						

TABLE D

Output listing for Molybdenum

THE FOLLOWING DATA HAVE BEEN ENTERED INTO MEMORY  
\*\*\*\*\*

AMPERES	KILO-VOLTS	BEAM-RADIUS	COLLECTOR	T1	T2	A
.150E+02	.170E+03	.100E+01	.100E+00	.800E+01	.140E+02	7

NUMBER OF ENERGY GROUPS-NGRPS 20

NUMBER OF TARGET NUCLEI-NRNUCL 7

NUMBER OF POINTS AFTER SHUTDOWN-NRPTS 12

ISAD=1

IGEDN=2

YOU ARE USING AN ABSORBING SOURCE

THIS SOURCE IS A SPHERE

ITS RADIUS IS R= .250E+00 METERS

THE OBSERVER IS STANDING AT A DISTANCE D= .100E+00 METERS FROM THE SURFACE OF THIS SPHERE

NUMBER OF GROUPS BOUNDARIES-NG 21

NUMBER OF FAST NEUTRON GROUPS-NF 2

SPECIFIC TIMES AFTER SHUTDOWN (MS)

TIME 1	0.
TIME 2	.100E+01
TIME 3	.200E+01
TIME 4	.300E+01
TIME 5	.400E+01
TIME 6	.500E+01
TIME 7	.600E+01
TIME 8	.700E+01
TIME 9	.800E+01
TIME 10	.900E+01
TIME 11	.100E+02
TIME 12	.200E+02

THE POSITION OF THE REACTOR-RELOC= 1  
THE GROUPS WILL BE ORDERED FROM TOP TO

BOTTOM BY THE PART OF THE BEAM WHICH ACTIVATES,  
TARGET, PLATE-COLL= 1  
ACTIVATION OF AN OBJECT OCCURS INSIDE THE BEAM

UNIT NORMALIZED FLUXES INSIDE THE BEAM

GROUP 1	.150E-05
GROUP 2	.110E-06
GROUP 3	.800E-07
GROUP 4	.100E-07

GROUP 5	.1178F-06
GROUP 6	.2942F-07
GROUP 7	.4706F-07
GROUP 8	.6470F-07
GROUP 9	.8234F-07
GROUP 10	.1000F-07
GROUP 11	.2764F-07
GROUP 12	.4528F-07
GROUP 13	.6292F-07
GROUP 14	.8056F-07
GROUP 15	.9820F-07
GROUP 16	.1584F-07
GROUP 17	.3348F-07
GROUP 18	.5112F-07
GROUP 19	.6876F-07
GROUP 20	.8640F-06

GROUP BOUNDARIES (MEV)

BOUNDARY 1	.2185E+01
BOUNDARY 2	.3367E+01
BOUNDARY 3	.4549E+01
BOUNDARY 4	.5731E+01
BOUNDARY 5	.6913E+00
BOUNDARY 6	.8095E+00
BOUNDARY 7	.9277E+00
BOUNDARY 8	.10459E-01
BOUNDARY 9	.22281E-01
BOUNDARY 10	.34103E-01
BOUNDARY 11	.45925E-01
BOUNDARY 12	.57747E-02
BOUNDARY 13	.69569E-02
BOUNDARY 14	.81391E-02
BOUNDARY 15	.93213E-02
BOUNDARY 16	.10503E-04
BOUNDARY 17	.12285E-04
BOUNDARY 18	.14067E-05
BOUNDARY 19	.15849E-05
BOUNDARY 20	.17631E-06
BOUNDARY 21	.19413E-10

NAME	IC-NUMBER	MASSICG1
PLUTONIUM-239	42092	.787E+02
PLUTONIUM-240	42094	.675E+02
PLUTONIUM-241	42096	.563E+02
PLUTONIUM-242	42098	.451E+02
PLUTONIUM-243	42100	.339E+02
PLUTONIUM-244	42098	.227E+02
PLUTONIUM-245	42100	.115E+02

ALOPS IS NOW READY TO RUN  
\*\*\*\*\*

THE FOLLOWING HAS BEEN CALCULATED BY ACTOS?  
 \*\*\*\*\*

INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND .7645E+13

AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND .7649E+12

INTERVALS OF THE FITTING FUNCTION OVER THE GROUP INTERVALS  
 GROUP LOWER BOUNDARY UPPER BOUNDARY REQUIRED INTEGRAL CALCULATED INTEGRAL

20	.107E+04	.414E+00	.879E-06	.802E-06
19	.414E+00	.113E+01	.709E-07	.308E-07
18	.113E+01	.306E+01	.246E-07	.346E-07
17	.306E+01	.107E+02	.255E-07	.799E-07
16	.107E+02	.270E+02	.201E-07	.301E-07
15	.270E+02	.101E+05	.449E-07	.448E-07
14	.101E+05	.583E+03	.727E-07	.773E-07
13	.583E+03	.122E+04	.204E-07	.704E-07
12	.122E+04	.236E+04	.402E-07	.402E-07
11	.236E+04	.107E+05	.471E-07	.457E-07
10	.107E+05	.219E+04	.272E-07	.252E-07
9	.219E+05	.248E+05	.742E-07	.742E-07
8	.248E+05	.527E+05	.204E-07	.386E-07
7	.527E+05	.111E+06	.479E-07	.479E-07
6	.111E+06	.178E+06	.294E-07	.294E-07
5	.178E+06	.540E+04	.178E-06	.138E-06
4	.540E+05	.111E+07	.176E-06	.136E-06
3	.111E+07	.192E+07	.833E-07	.433E-07
2	.192E+07	.221E+07	.138E-06	.138E-06
1	.221E+07	.229E+07	.790E-05	.790E-05

THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVITIES IN PG

TIME(H)	1	2	3	4	5	6	7	8	9	10	11	12	13	14
180SET PRODUCT	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)	ACT(13)	ACT(14)
42392	42091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42092	41052	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42093	41091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42094	40095	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42095	42093	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05

TIME(H)	1	2	3	4	5	6	7	8	9	10	11	12	13	14
150 SET	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)	ACT(13)	ACT(14)
42392	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05	.148E+05

TIME(H)	1	2	3	4	5	6	7	8	9	10	11	12	13	14
180SET PRODUCT	ACT(1)	ACT(2)	ACT(3)	ACT(4)	ACT(5)	ACT(6)	ACT(7)	ACT(8)	ACT(9)	ACT(10)	ACT(11)	ACT(12)	ACT(13)	ACT(14)
42096	42063	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42097	42092	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42098	42092	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42099	41054	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42099	42092	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
42099	40091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

42094	42095	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42094	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42095	42094	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42095	42093	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42095	41095	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42095	41094	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42095	41092	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42095	42096	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+02	.200E+02	.300E+02	.400E+02	.500E+02	.600E+02	.700E+02	.800E+02	.900E+02	.100E+03	.200E+03		
TARGET	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42095	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42096	42095	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42096	42094	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42096	41356	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42096	43093	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42096	42097	.C	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42096	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42097	42094	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42097	42095	.0	.C	.C	.0	.0	.0	.0	.0	.0	.0	.0		.C
42097	41097	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42097	43094	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42097	42099	.0	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42097	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
TIME(HH)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02		
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)		
42098	42097	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42098	42096	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42098	41398	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42098	43095	.C	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0		.C
42098	42099	.104E+00	.104E+00	.100E+00	.101E+00	.101E+00	.101E+00	.101E+00	.101E+00	.101E+00	.101E+00	.101E+00		.C

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
199PT	ACTE111	ACTE121	ACTE131	ACTE141	ACTE151	ACTE161	ACTE171	ACTE181	ACTE191	ACTE1101	ACTE1111	ACTE1121
4209A	.176E+05	.104E+09	.107E+07	.122E+09	.107E+09	.130E+05	.991E+09	.691E+09	.171E+08	.961E+08	.451E+08	.457E+08

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
TARGET PROJECT	ACTE111	ACTE121	ACTE131	ACTE141	ACTE151	ACTE161	ACTE171	ACTE181	ACTE191	ACTE1101	ACTE1111	ACTE1121
4210J 4209A	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4210J 4209A	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4210J 4210C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4210J 4009A	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4210J 4210J	.110E+05	.435E+07	.769E+06	.213E+05	.124E+04	.717E+02	.455E+01	.241E+00	.140E-01	.909E-03	.449E-04	.201E-16

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
TARGET	ACTE111	ACTE121	ACTE131	ACTE141	ACTE151	ACTE161	ACTE171	ACTE181	ACTE191	ACTE1101	ACTE1111	ACTE1121
4210J	.110E+05	.435E+07	.769E+06	.213E+05	.124E+04	.717E+02	.455E+01	.241E+00	.140E-01	.805E-03	.449E-04	.201E-16

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
SYSTEM	ACTE111	ACTE121	ACTE131	ACTE141	ACTE151	ACTE161	ACTE171	ACTE181	ACTE191	ACTE1101	ACTE1111	ACTE1121
	.215E+05	.111E+05	.104E+09	.102E+09	.101E+09	.100E+09	.991E+09	.981E+08	.971E+08	.961E+08	.951E+08	.947E+08

THE FOLLOWING OUTPUT DATA ARE CALCULATED POS. VALUES IN PERCENT

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
TARGET PROJECT	DSR1111	DSR1121	DSR1131	DSR1141	DSR1151	DSR1161	DSR1171	DSR1181	DSR1191	DSR11101	DSR11111	DSR11121
4209J 4209A	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209J 41091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209J 41092	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209J 42091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209J 42093	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
PROJECTS (DSR1111)	DSR1111	DSR1121	DSR1131	DSR1141	DSR1151	DSR1161	DSR1171	DSR1181	DSR1191	DSR11101	DSR11111	DSR11121

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
100E1	DSR1111	DSR1121	DSR1131	DSR1141	DSR1151	DSR1161	DSR1171	DSR1181	DSR1191	DSR11101	DSR11111	DSR11121
4209J	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	W	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+C2	200E+02
TARGET PROJECT	DSR1111	DSR1121	DSR1131	DSR1141	DSR1151	DSR1161	DSR1171	DSR1181	DSR1191	DSR11101	DSR11111	DSR11121
4209A 42091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209A 41091	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209A 4209A	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209A 4209C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209A 4209J	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
4209A 4209J	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

42094	42385	0	0	0	0	0	0	0	0	0	0	0	0	0	0
PRODUCT	DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42094	42385	0	0	0	0	0	0	0	0	0	0	0	0		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42095	40092	0	0	0	0	0	0	0	0	0	0	0	0		
42095	41094	0	0	0	0	0	0	0	0	0	0	0	0		
42095	41095	0	0	0	0	0	0	0	0	0	0	0	0		
42095	42095	0	0	0	0	0	0	0	0	0	0	0	0		
42095	42094	0	0	0	0	0	0	0	0	0	0	0	0		
42095	42096	0	0	0	0	0	0	0	0	0	0	0	0		
PRODUCT	DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42095	42095	0	0	0	0	0	0	0	0	0	0	0	0		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42096	40093	0	0	0	0	0	0	0	0	0	0	0	0		
42096	41394	0	0	0	0	0	0	0	0	0	0	0	0		
42096	42094	0	0	0	0	0	0	0	0	0	0	0	0		
42096	42095	0	0	0	0	0	0	0	0	0	0	0	0		
42096	42097	0	0	0	0	0	0	0	0	0	0	0	0		
PRODUCT	DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42097	42097	0	0	0	0	0	0	0	0	0	0	0	0		
42097	41093	0	0	0	0	0	0	0	0	0	0	0	0		
42097	42095	0	0	0	0	0	0	0	0	0	0	0	0		
42097	42096	0	0	0	0	0	0	0	0	0	0	0	0		
PRODUCT	DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
TIME(H)		0	100E+01	200E+01	300E+01	400E+01	500E+01	600E+01	700E+01	800E+01	900E+01	100E+02	200E+02		
TARGET	PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(10)	DSR(11)	DSR(12)		
42097	40094	0	0	0	0	0	0	0	0	0	0	0	0		

I  
C  
N  
I





TFTR neutral beam injector shielding walls

Assuming the test cell a spherical cavity, the inner radius can be determined by conserving the interior surface of the parallelepipedic real cell (6.2 x 7.4 x 20.2 meters).

$$\begin{aligned} A &= \text{inside surface} \\ &= 2 \times 7.4 \times 20.2 + 20.2 \times 6.8 + 6.8 \times 7.4 \\ &= 674.32 \text{ m}^2 \end{aligned}$$

hence we have:

$$R = ( A / (4 \times \pi) )^{1/2} = 7.325 \text{ m}$$

The thickness used is the same as the one of the real cell:

$$H = 0.30 \text{ m}$$

Assuming a concrete density of 2.6 g/cm<sup>3</sup> the total mass is

$$\text{TMASS} = 2.6 \times \text{VOLU}$$

$$\text{where VOLU} = \frac{4}{3} \cdot \pi \cdot ((R+H)^3 - R^3) = 2.11 \times 10^8 \text{ m}^3$$

$$\text{then TMASS} = 547 \text{ 753 kg}$$

distributed over the different elements as listed in Table E.

This time, the space-independent group-fluxes inside the test cell are inputed from the input data deck. The average group-fluxes in the walls are computed by setting: IWALL = 1 with KWALL = 0.1111 since for concrete  $\lambda \approx 9$  cm. The input data deck is shown in Table F. Table G is the output listing. The dose is almost entirely due to the 1.78 MeV gamma rays emitted by <sup>28</sup>Al, a product of <sup>27</sup>Al. Since the half-life of this product is short (134 s), after one hour, the dose-rate due to the activation of the walls is

T A B L E E

Concrete Composition

<u>ELEMENT</u>	<u>% IN CONCRETE</u>	<u>ISOTOPES</u>	<u>ATOMIC %</u>	<u>WEIGHT %</u>	<u>MASS (Kg)</u>
H	16.3	1 H	99.985	99.97	89257.22
		2 H	0.015	0.03	26.78
O	54.5	16 O	99.759	99.73	297723.64
		17 O	0.037	0.04	116.44
		18 O	0.204	0.23	684.92
Al	2.1	27 Al	100.	100.	11503.00
Si	23.9	28 Si	92.27	91.92	120344.50
		29 Si	4.68	4.82	6307.01
		30 Si	3.05	3.26	4261.49
Ca	3.2	40 Ca	97.13	96.89	16981.50
		42 Ca	0.64	0.67	117.50
		43 Ca	0.15	0.16	28.20
		44 Ca	2.06	2.26	396.20
		48 Ca	0.02	0.02	4.20

T A B L E F

Input Data Deck for Concrete

0.000000E+00	0.650E+02	.170E+03	.100E+01	.010E+01	.800E+01	.160E+02	7
20							
12	7.325E+00	0.000E+00	0.300E+00	41			
14							
2	0.000E+1	0.100E+1	0.200E+1	0.300E+1	0.400E+1	0.500E+1	
	0.600E+1	0.700E+1	0.800E+1	0.900E+1	1.000E+1	2.000E+1	
	HYDROGEN-1	1001	89.257E+03				
	HYDROGEN-2	1002	26.780E+00				
	OXYGEN-16	8016	29.772E+04				
	OXYGEN-17	8017	11.644E+01				
	OXYGEN-18	8018	68.492E+01				
	ALUMINIUM-27	13027	11.503E+03				
	SILICON-28	14028	12.034E+04				
	SILICON-29	14029	63.070E+02				
	SILICON-30	14030	42.615E+02				
	CALCIUM-40	20040	16.982E+03				
	CALCIUM-42	20042	11.750E+01				
	CALCIUM-43	20043	28.200E+00				
	CALCIUM-44	20044	39.620E+01				
	CALCIUM-48	20048	4.200E+00				
	1.000E-11	4.140E-07	1.125E-06	3.059E-06	1.068E-05	2.902E-05	
	1.013E-04	5.829E-04	1.234E-03	3.355E-03	1.033E-02	2.188E-02	
	2.479E-02	5.248E-02	1.111E-01	1.576E-01	5.502E-01	1.106E+00	
	1.827E+00	2.307E+00	2.385E+00				
	0.1111E-01						
	8.288E-7	3.076E-8	3.456E-8	3.951E-8	3.011E-8	4.684E-8	
	7.234E-8	3.043E-8	4.021E-8	4.570E-8	3.522E-8	7.420E-9	
	3.860E-8	4.790E-8	2.842E-8	1.378E-7	1.058E-7	8.333E-8	
	1.382E-7	7.901E-6					

T A B L E G

Output Listing for Concrete

THE FOLLOWING DATA HAVE BEEN PRINTED INTO MEMORY  
 .....

AMPELES	KILC-VOLTS	BEAM-EFFECT	DENSITY	T1	T2	N
.650E+02	.170E+03	.107E+01	.100E+00	.400E+01	.160E+02	7

NUMBER OF ENERGY GROUPS-NRNGPS 20

NUMBER OF TARGET NUCLIDES-NRNUCL 14

NUMBER OF POINTS AFTER SHUTDOWN-NRPS 12

ISAD=1

IGCM=4

YOU ARE USING AN ABSORBING SOURCE

THIS SOURCE IS A SPHERICAL CAVITY  
 ITS RADIUS IS R= .1E+01 METERS  
 ITS WIDTH IS W= .100E+00 METERS  
 THE OBSERVER IS AT THE CENTER OF THE CAVITY

NUMBER OF GROUPS BEYOND-NRNG 21

NUMBER OF FAST NEUTRON GROUPS-NF 7

SPECIFIC TIMES AFTER SHUTDOWN (H)

TIME 1	C.
TIME 2	.100E+01
TIME 3	.200E+01
TIME 4	.300E+01
TIME 5	.400E+01
TIME 6	.500E+01
TIME 7	.600E+01
TIME 8	.700E+01
TIME 9	.800E+01
TIME 10	.900E+01
TIME 11	.100E+02
TIME 12	.200E+02

OPTION IN THE FLUXES=1BLL= 0  
 THE FLUXES WILL BE READ FROM THE INPUT

POSITION ON THE MAP OF THE ROOM-WALLS ACTIVATION,  
 TAPES PLACE=1BLL= 1  
 ACTIVATION OF THE WALLS OF THE ROOM

AVERAGE UNIT NORMALIZED FLUXES IN THE WALLS (WALL= .100E+00)

GROUP 1	.214E-05
GROUP 2	.214E-07
GROUP 3	.221E-07

GROUP 4	.286E+07
GROUP 5	.2737E+07
GROUP 6	.2765E+08
GROUP 7	.125E+07
GROUP 8	.164E+07
GROUP 9	.201E+08
GROUP 10	.654E+08
GROUP 11	.1277E+07
GROUP 12	.168E+07
GROUP 13	.244E+08
GROUP 14	.197E+07
GROUP 15	.124E+07
GROUP 16	.817E+08
GROUP 17	.167E+07
GROUP 18	.636E+08
GROUP 19	.633E+08
GROUP 20	.274E+08

GROUP BOUNDARIES (MEV)

BOUNDARY 1	.278E+01
BOUNDARY 2	.2767E+01
BOUNDARY 3	.1827E+01
BOUNDARY 4	.116E+01
BOUNDARY 5	.550E+00
BOUNDARY 6	.153E+00
BOUNDARY 7	.111E+00
BOUNDARY 8	.224E+01
BOUNDARY 9	.243E+01
BOUNDARY 10	.217E+01
BOUNDARY 11	.1637E+01
BOUNDARY 12	.234E+02
BOUNDARY 13	.1274E+02
BOUNDARY 14	.582E+02
BOUNDARY 15	.1982E+02
BOUNDARY 16	.2502E+04
BOUNDARY 17	.104E+04
BOUNDARY 18	.205E+05
BOUNDARY 19	.112E+05
BOUNDARY 20	.414E+10
BOUNDARY 21	.1ECCF+10

NAME	ID-NUMBER	MASS (MEV)
HYDROGEN-1	1001	.938E+09
HELIUM-4	1002	.372E+09
HYDROGEN-16	9016	.938E+09
DEUTERIUM-2	9022	.201E+09
DEUTERIUM-3	9028	.332E+09
ALUMINUM-27	13027	.227E+09
SILICON-28	14028	.280E+09
SILICON-29	14029	.283E+09
SILICON-30	14030	.287E+09
URANIUM-238	92038	.395E+09
CALIFORNIA-47	98047	.470E+09
CALIFORNIA-49	98049	.490E+09
CALIFORNIA-50	98050	.500E+09
CALIFORNIA-51	98051	.510E+09
CALIFORNIA-52	98052	.520E+09
CALIFORNIA-53	98053	.530E+09
CALIFORNIA-54	98054	.540E+09
CALIFORNIA-55	98055	.550E+09
CALIFORNIA-56	98056	.560E+09
CALIFORNIA-57	98057	.570E+09
CALIFORNIA-58	98058	.580E+09
CALIFORNIA-59	98059	.590E+09
CALIFORNIA-60	98060	.600E+09
CALIFORNIA-61	98061	.610E+09
CALIFORNIA-62	98062	.620E+09
CALIFORNIA-63	98063	.630E+09
CALIFORNIA-64	98064	.640E+09
CALIFORNIA-65	98065	.650E+09
CALIFORNIA-66	98066	.660E+09
CALIFORNIA-67	98067	.670E+09
CALIFORNIA-68	98068	.680E+09
CALIFORNIA-69	98069	.690E+09
CALIFORNIA-70	98070	.700E+09
CALIFORNIA-71	98071	.710E+09
CALIFORNIA-72	98072	.720E+09
CALIFORNIA-73	98073	.730E+09
CALIFORNIA-74	98074	.740E+09
CALIFORNIA-75	98075	.750E+09
CALIFORNIA-76	98076	.760E+09
CALIFORNIA-77	98077	.770E+09
CALIFORNIA-78	98078	.780E+09
CALIFORNIA-79	98079	.790E+09
CALIFORNIA-80	98080	.800E+09
CALIFORNIA-81	98081	.810E+09
CALIFORNIA-82	98082	.820E+09
CALIFORNIA-83	98083	.830E+09
CALIFORNIA-84	98084	.840E+09
CALIFORNIA-85	98085	.850E+09
CALIFORNIA-86	98086	.860E+09
CALIFORNIA-87	98087	.870E+09
CALIFORNIA-88	98088	.880E+09
CALIFORNIA-89	98089	.890E+09
CALIFORNIA-90	98090	.900E+09
CALIFORNIA-91	98091	.910E+09
CALIFORNIA-92	98092	.920E+09
CALIFORNIA-93	98093	.930E+09
CALIFORNIA-94	98094	.940E+09
CALIFORNIA-95	98095	.950E+09
CALIFORNIA-96	98096	.960E+09
CALIFORNIA-97	98097	.970E+09
CALIFORNIA-98	98098	.980E+09
CALIFORNIA-99	98099	.990E+09

ACCTS IS NOW READY TO SUP.  
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THE FOLLOWING HAS BEEN CALCULATED BY ACCESS

INSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SECOND \* 7459143

AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND

INTEGRALS OF THE RATING FUNCTION OVER THE CORE IN INTERVALS  
GROUP NUMBER LOWER NUMBER INDICATE INITIAL CALCULATED INTEGRAL

20	.403E+24	1.17E+01	1.17E+01	1.17E+01
19	4.47E+02	1.17E+01	1.17E+01	1.17E+01
18	1.17E+01	1.17E+01	1.17E+01	1.17E+01
17	.378E+01	1.17E+01	1.17E+01	1.17E+01
16	1.17E+01	1.17E+01	1.17E+01	1.17E+01
15	.257E+07	1.17E+01	1.17E+01	1.17E+01
14	1.01E+03	1.17E+01	1.17E+01	1.17E+01
13	4.63E+03	1.17E+01	1.17E+01	1.17E+01
12	1.17E+01	1.17E+01	1.17E+01	1.17E+01
11	1.17E+04	1.17E+01	1.17E+01	1.17E+01
10	1.17E+04	1.17E+01	1.17E+01	1.17E+01
9	1.17E+03	1.17E+01	1.17E+01	1.17E+01
8	1.17E+01	1.17E+01	1.17E+01	1.17E+01
7	1.17E+07	1.17E+01	1.17E+01	1.17E+01
6	1.17E+03	1.17E+01	1.17E+01	1.17E+01
5	1.17E+05	1.17E+01	1.17E+01	1.17E+01
4	1.17E+05	1.17E+01	1.17E+01	1.17E+01
3	1.17E+07	1.17E+01	1.17E+01	1.17E+01
2	1.17E+07	1.17E+01	1.17E+01	1.17E+01
1	1.17E+07	1.17E+01	1.17E+01	1.17E+01

THE FOLLOWING GROUP DATA ARE CALCULATED ACTIVITIES IN BQ

(TIME)	1001	1002	1003	1004
1001	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1002	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1003	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1004	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1005	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1006	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1007	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1008	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1009	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1010	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1011	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1012	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1013	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1014	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1015	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1016	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1017	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1018	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1019	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1020	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1021	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1022	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1023	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1024	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1025	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1026	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1027	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1028	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1029	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1030	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1031	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1032	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1033	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1034	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1035	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1036	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1037	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1038	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1039	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1040	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1041	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1042	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1043	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1044	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1045	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1046	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1047	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1048	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1049	1.00E+01	1.00E+01	1.00E+01	1.00E+01
1050	1.00E+01	1.00E+01	1.00E+01	1.00E+01

TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
8015	8015	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
8016	7015	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
8016	8012	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
8016	7016	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
8015	7015	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
8016	8017	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
TIME(4)	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
8015	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C

ATTENTION-TARGET NUCLEIDE NUMBER 8017 DOES NOT EXIST IN DATA LIBRARY-THEREFORE-NO CONTRIBUTION FROM THIS NUCLEIDE IN CASE CALS.

ATTENTION-TARGET NUCLEIDE NUMBER 8018 DOES NOT EXIST IN DATA LIBRARY-THEREFORE-NO CONTRIBUTION FROM THIS NUCLEIDE IN CASE CALS.

TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
13027	13026	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
13027	12027	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
13027	12026	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
13027	12025	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
13027	11024	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
13027	13029	.120E+11	.187E+03	.279E-05	.425E-13	.644E-21	.987E-29	.150E-36	.229E-44	.349E-52	.531E-60	.809E-68

TIME(4)	TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
13027	.110E+11	.187E+03	.279E-05	.425E-13	.644E-21	.987E-29	.150E-36	.229E-44	.349E-52	.531E-60	.809E-68	.949-146	

TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
14029	14027	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	13029	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	12024	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	13029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	13025	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	14026	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C

TIME(4)	TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
14029	.110E+11	.187E+03	.279E-05	.425E-13	.644E-21	.987E-29	.150E-36	.229E-44	.349E-52	.531E-60	.809E-68	.949-146	

TARGET PRODUCT	ACT(11)	ACT(12)	ACT(13)	ACT(14)	ACT(15)	ACT(16)	ACT(17)	ACT(18)	ACT(19)	ACT(10)	ACT(11)	ACT(12)
14029	14029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	13029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	13026	.C	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C
14029	12025	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.C



14029	13029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
14029	12076	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
14029	14030	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
14029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02	
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)	
14030	14029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	
14030	14029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	
14030	13029	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	
14030	13028	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	
14030	12027	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	
14030	14031	.165E+10	.129E+10	.900E+09	.780E+09	.583E+09	.448E+09	.344E+09	.264E+09	.202E+09	.155E+09	.119E+09	.847E+07

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
14030	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
14030	.165E+10	.129E+10	.900E+09	.780E+09	.583E+09	.448E+09	.344E+09	.264E+09	.202E+09	.155E+09	.119E+09	.847E+07

TIME(H)	.0	.100E+01	.200E+01	.900E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
20040	20039	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	19039	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	19038	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	17040	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00	.113E+00
20040	19039	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	19037	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09
20040	20041	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04	.113E+04

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
20040	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09	.792E+09

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
20042	20041	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	19041	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	18038	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	17042	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	19036	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	ACT(111)	ACT(112)	ACT(113)	ACT(114)	ACT(115)	ACT(116)	ACT(117)	ACT(118)	ACT(119)	ACT(110)	ACT(111)	ACT(112)
20042	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	.0	.100E+01	.700E+01	.500E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
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TARGET PRODUCT	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20043 20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043 18079	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043 17043	.122E+06	.119E+06	.115E+06	.112E+06	.109E+06	.105E+06	.102E+06	.986E+05	.946E+05	.926E+05	.898E+05	.859E+05
20043 17043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043 17043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043 20044	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
1AF611	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20043	.122E+06	.119E+06	.115E+06	.112E+06	.109E+06	.105E+06	.102E+06	.986E+05	.946E+05	.926E+05	.898E+05	.859E+05

TARGET PRODUCT	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20044 20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20044 17044	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20044 17043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20044 18043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20044 20044	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
1AF611	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20044	.122E+06	.119E+06	.115E+06	.112E+06	.109E+06	.105E+06	.102E+06	.986E+05	.946E+05	.926E+05	.898E+05	.859E+05

TARGET PRODUCT	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20048 20047	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20048 18048	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20048 20049	.743E+C7	.559E+05	.594E+03	.518E+01	.440E-01	.409E-03	.362E-C5	.211E-D7	.285E-09	.253E-11	.224E-13	.676E-16

TIME(H)	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
1AF611	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20048	.743E+C7	.559E+05	.594E+03	.518E+01	.440E-01	.409E-03	.362E-C5	.211E-D7	.285E-09	.253E-11	.224E-13	.676E-16

TIME(H)	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
1AF611	.0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
20048	.743E+C7	.559E+05	.594E+03	.518E+01	.440E-01	.409E-03	.362E-C5	.211E-D7	.285E-09	.253E-11	.224E-13	.676E-16

THE FOLLOWING OUTPUT FILE ARE CALCULATED THESE VALUES IN 8006/4

TIME(H)	ACT1111	ACT1121	ACT1131	ACT1141	ACT1151	ACT1161	ACT1171	ACT1181	ACT1191	ACT1101	ACT1111	ACT1121
TARGET PRODUCT	CSR1111	CSR1121	CSR1131	CSR1141	CSR1151	CSR1161	CSR1171	CSR1181	CSR1191	CSR1101	CSR1111	CSR1121
1001 1002	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
PRODUCT DAJGH1E	CSR1111	CSR1121	CSR1131	CSR1141	CSR1151	CSR1161	CSR1171	CSR1181	CSR1191	CSR1101	CSR1111	CSR1121

Table with columns: TIMEFHM, TARGET PRODUCT, VALUE, and TIMEFHM. The table contains multiple rows of data, many of which are zero values. The data is organized in columns, with some rows grouped under specific labels like 'TIMEFHM', 'TARGET PRODUCT', 'VALUE', and 'TIMEFHM'.



TIME(H)	FC	.100F+01	.200F+01	.300F+01	.400F+01	.500F+01	.600F+01	.700F+01	.800F+01	.900F+01	.100E+02	.200E+02
TARGET	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
14033	.642E-C3	.716E-C3	.243E-C3	.194E-C3	.147E-C3	.110E-C3	.242E-C4	.644E-04	.496E-04	.181E-04	.292E-04	.228E-05

TIME(H)	CD	.100F+01	.200F+01	.300F+01	.400F+01	.500F+01	.600F+01	.700F+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20042	18036	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20040	18037	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043	17038	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20045	17039	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043	17040	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20045	20035	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043	20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
PRODUCT DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)

TIME(H)	CD	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	CD	.100F+01	.200F+01	.300F+01	.400F+01	.500F+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20042	18038	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	18035	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	19041	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	19042	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	20041	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
PRODUCT DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)

TIME(H)	FC	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20042	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	CD	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20043	19035	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043	17043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	17047	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20043	17043	.395E-04	.377E-04	.357E-04	.337E-04	.316E-04	.295E-04	.272E-04	.249E-04	.225E-04	.202E-04	.179E-04
20043	20047	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
20042	20044	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0
PRODUCT DAUGHTER	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
17043	20043	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0	.0

TIME(H)	FC	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)
20043	.239E-C4	.277E-C4	.312E-C4	.347E-C4	.380E-C4	.412E-C4	.443E-C4	.473E-C4	.501E-C4	.528E-C4	.554E-C4	.579E-C4

TIME(H)	CD	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02
TARGET PRODUCT	DSR(11)	DSR(12)	DSR(13)	DSR(14)	DSR(15)	DSR(16)	DSR(17)	DSR(18)	DSR(19)	DSR(110)	DSR(111)	DSR(112)

20044	18041	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20044	17041	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20044	16044	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20044	20043	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20044	20045	.154E-C7	.154E-07	.153E-07	.153E-07	.153E-07	.153E-07	.153E-07	.153E-C7	.152E-07	.153E-07	.152E-07	.153E-07	.153E-07	.153E-07
PRODUCT	NAUGHTER	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141
20045	21045	0	0	0	0	0	0	0	0	0	0	0	0	0	0

TIME(H)	0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02	.300E+02	.400E+02	.500E+02
TARGET	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151
20044	.154E-C7	.154E-07	.153E-07	.153E-07	.153E-07	.153E-07	.153E-C7	.152E-07	.153E-07	.152E-07	.153E-07	.152E-07	.153E-07	.152E-07	.153E-07

TIME(H)	0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02	.300E+02	.400E+02	.500E+02
TARGET	PRODUCT	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141
20048	19048	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20048	20047	0	0	0	0	0	0	0	0	0	0	0	0	0	0
20048	20045	.329E-C2	.292E-04	.259E-06	.220E-04	.204E-10	.181E-12	.160E-14	.142E-16	.126E-18	.112E-20	.993E-22	.900E-24	.820E-26	.740E-28
PRODUCT	DAUGHTER	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141
20048	21049	.151E-C5	.726E-06	.390E-08	.160E-06	.414E-07	.393E-07	.180E-C7	.913E-08	.440E-08	.212E-08	.102E-08	.694E-12	.390E-12	.220E-12

TIME(H)	0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02	.300E+02	.400E+02	.500E+02
TARGET	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151
20040	.125E-C1	.255E-04	.439E-06	.172E-06	.814E-07	.397E-07	.149E-C7	.912E-08	.440E-08	.212E-08	.102E-08	.695E-12	.390E-12	.220E-12	.150E-12

TIME(H)	0	.100E+01	.200E+01	.300E+01	.400E+01	.500E+01	.600E+01	.700E+01	.800E+01	.900E+01	.100E+02	.200E+02	.300E+02	.400E+02	.500E+02
SFSTEP	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151	D5P1161	D5P1171	D5P1181	D5P1191	D5P1101	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151
	.405E-C1	.253E-01	.297E-03	.221E-03	.173E-03	.143E-07	.118E-C2	.557E-C4	.798E-04	.672E-04	D5P1111	D5P1121	D5P1131	D5P1141	D5P1151

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insignificant. In fact a more accurate calculation should take into account heavy elements which, even in small quantities in concrete, could significantly increase the dose.





C        DOSPRO(I)        DOSE DUE TO A PRODUCT AT SHUTDOWN  
C        DOSPRC(I)        DOSE DUE TO A GIVEN PRODUCT ALONE  
C        DOSSUM(I)        DOSE RATE DUE TO A PARTICULAR TARGET NUCLIDE  
C                        FOR UP TO 12 DIFFERENT TIMES  
C        DPMVAR         VALUE OF DEPENDENT VARIABLE IN NUMERICAL  
C                        INTEGRATIONS  
C        E                ENERGY OF GAMMA RAY  
C        EN(I)            ENERGY VALUE ASSOCIATED WITH A CROSS SECTION  
C        EI(I)            ENERGY VALUE ASSOCIATED WITH A CROSS SECTION  
C        F                BEAM FRACTION  
C        FCONST(I)        EXPONENTIAL CONSTANTS FOR FAST NEUTRON GROUPS  
C        FF                SHAPE FACTOR FOR THE FLUXES IN THE WALL  
C        FLUXBC            WEIGHTING FLUX AT A GROUP BOUNDARY  
C        FPRIME           EVALUATED VALUE OF THE DERIVATIVE OF THE  
C                        TRANSCENDENTAL EQUATION FOR B  
C        FXSUBI            EVALUATED VALUE OF THE TRANSCENDENTAL EQUATION  
C                        FOR B  
C                        ABSORBING CYLINDER  
C        GFC                GEOMETRIC CORRECTION IN THE CASE OF AN  
C        GFS                GEOMETRIC CORRECTION IN THE CASE OF AN  
C                        ABSORBING SPHERE  
C        GP                GROUP EQUADARIES  
C        GPFLUX(I)        FEW GROUP FLUXES-UP TO 50  
C        GPXSEC(I)        FEW GROUP CROSS SECTIONS-UP TO 50  
C        GUESSI            INITIAL GUESS FOR SOLVING TRANSCENDENTAL  
C                        EQUATIONS  
C        H                HEIGHT OF CYLINDER (M)  
C        ICOUNT            NEUTRON GROUP COUNTER  
C        ICPI             ICOUNT+1  
C        ICPE             ICOUNT+2  
C        IDNC(I)           5 DIGIT ID NUMBER OF TARGET NUCLIDE  
C        IDLM             DUMMY VARIABLE FOR ADVANCING DISK FILE RECORDS  
C        IEPH             COUNTS FOR INTERMEDIATE NEUTRON GROUPS  
C        IFLUX             OPTICN ON THE WAY TO INPUT THE FLUXES  
C        IGEOM             DESIGNATES PROBLEM GEOMETRY  
C        INSTANS           INSTANTANLCUS NUMBER OF NEUTRONS PRODUCED PER  
C                        SECOND  
C        IORDER(I)        HOLDING ARRAY FOR PRODUCT NUCLIDE ID NUMBERS  
C        IPRGD             PRODUCT NUCLIDE COUNTS  
C        ISAB             OPTICN ON THE ABSORPTION STRATEGY  
C        ITAPE             READ VARIABLE FOR DISK FILE DATA  
C        ITER             ITERATION COUNTER  
C        IWALL             OPTICN ON THE PLACE OF THE ACTIVATION  
C                        INVERSE OF THE RELAXATION LENGTH OF THE WALL  
C                        MATERIAL  
C        MASS(I)           NUMBER OF KILOGRAMS OF A PARTICULAR NUCLIDE  
C        MU                ATTENUATION COEFFICIENT  
C        MULT             PLTIPLICITY OF GAMMA RAY  
C        N                NUMBER OF TESTS  
C        NEPH             NUMBER OF INTERMEDIATE NEUTRON GROUPS  
C        NEI(I)            NUMBER OF E1 E1 PAIRS  
C        NF                NUMBER OF FAST NEUTRON GROUPS  
C        NG                NUMBER OF GROUP BOUNDARIES-UP TO 51  
C        NDEGFS            NUMBER OF ENERGY GROUPS-UP TO 50  
C        NONUCL            NUMBER OF TARGET NUCLIDES-UP TO 30 PER RUN  
C        NOPAS            NUMBER OF POINTS AFTER SHUTDOWN-UP TO 12  
C        NOTORC(I)        HOLDING ARRAY FOR PRODUCT NUCLIDE NUMBERS THAT  
C                        HAVE NOT BEEN RE-ARRANGED IN ASCENDING ORDER  
C        NP                NUMBER OF ENERGY-CROSS-SECTION PAIRS  
C        OPTICN            VARIABLE THAT DETERMINES WHETHER OR NOT AVEPNS  
C                        WILL BE CALCULATED BY ACOS (R READ FROM A  
C                        DATA CARD AS PRE-DETERMINED INPUT

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C      PARSUP          PARTIAL SLP IN NUMERICAL INTGRATIONS
C      PRODHL(I,J)    PRODUCT NUCLIDE HALF-LIFE (S)
C      PRODNUI(I,J)   5 DIGIT ID NUMBER OF PRODUCT NUCLIDE
C      R              RADIUS OF HOMOGENEOUS NON-ABSORBING SPHERE OR
C                   CYLINDER (M)
C      REALDUM        REAL DUPHY VARIABLE FOR ADVAAING DISK FILE
C                   RECORDS
C      SOU            DOSE RATE SOURCE STRENGTH DIVIDED BY THE
C                   UNIT DCSE RATE FLUX
C      STAFS(I)       SPECIFIC TIMES AFTER SHUTDOWN (M)
C      STOPHL(I)      HOLDING ARRAY FOR PRODUCT NUCLIDE ID NUMBERS
C      STOTAC(I,J)    HOLDING ARRAY FOR PRODUCT NUCLIDE ACTIVITIES
C      SUMACT(I)       SUM OF ACTIVITIES PRODUCED FROM A PARTICULAR
C                   TARGET NUCLIDE FOR UP TO 12 DIFFERENT TIMES
C      SUM1           EXPONENTIAL EFFECT OF TESTS AND PAUSES ON THE
C                   SHUTDOWN TIME ACTIVITY OF A SPECIFIC PRODUCT
C                   NUCLIDE
C      SUMGEL         RUNNING SUM OF AVERAGE CROSS SECTIONS
C      SUMINT         RESULT OF NUMERICAL INTEGRATIONS
C      SYSACT(I)      SYSTEM ACTIVITY FOR UP TO 12 DIFFERENT TIMES
C      SYSDOSE(I)     SYSTEM DOSE RATE FOR UP TO 12 DIFFERENT TIMES
C      TARMAS         ATOMIC WEIGHT OF TARGET NUCLIDE (AMU)
C      TEMPS         TEMPORARY STORAGE FOR DD LOOP SUMS
C      TIMACT(I,J,K) INDUCEC ACTIVITY AS A FUNCTION OF TARGET
C                   NUCLIDE,REACTION, AND TIME
C      TMASS         TOTAL NUMBER OF KILOGRAMS OF THE TARGET
C      T1            LENGTH OF TEST (H)
C      T2            LENGTH OF PAUSE (H)
C      V             VOLTAGE (KV)
C      VOLU          VOLUME OF THE SOURCE
C      WIDTH         DELTA E USED IN NUMERICAL INTEGRATIONS
C      XSBIP1        LATEST REFINED GUESS FOR B
C      XSUB1         INDEPENDENT VARIABLE OF THE TRANSCENDENTAL
C                   EQUATION FOR B
C      ZINCR         INCREMENT IN ENERGY USED IN NUMERICAL
C                   INTEGRATIONS
C      ZNAME         VARIABLE FOR NUCLICE NAMES
C      ZNUMER        NUMEFATOR IN THE TRANSCENDENTAL EQUATION FOR B

```

```

C
C      DIMENSION TIMACT(30,12,12),SYSACT(12),SUMACT(12),SYSDOSE(12),PPROHL
C      + (30,13),PRODNUI(30,13),DCSSUM(12),ZNAF1(30),STAFS(12),BFLUX(50),IDN
C      + 0(30),GFXSEC(50),PASS(30),GPFUX(50),GPI5(1),ZNAF2(130)
C      DIMENSION CONST(50),FCNST(50)
C      DIMENSION DOSNUC(13),DOSEGT(13),DOSPRO(13)
C      COMMON A,V,F,DF,T1,T2,N,R,D,OPTION,IGEOM,ISAB
C      REAL INSNFS,MASS
C      CALL INPUT(NDNUCL,NOEGFS,AOFAS,NG,NF)
C      CALL ARAY INIBFLUX,ZNAM1,IGND,PASS,STAFS,GP,NOEGFS,MONUCL,MOFAS,NG,
C      + TMASS,ZNAF2)
C      IF(DFTICH.EQ.0.D000000) GC TC 1
C      AVEAFS=OPTION
C      WRITE(6,20)
20  FORMAT(1H1,48H THE FOLLOWING DATA HAVE BEEN ENTERED IN'D MEMORY,/,
C      + 48H .....//)
C      GC TC 2
1  CALL SCURCE(AVEAFS)
2  CALL GRFLUX(AVEAFS,NOEGFS,GPFUX,BFLUX)
C      CALL W7FLUX(BFLUX,GP,CONST,FCNST,NOEGFS,NF,NG)
C      CALL ACTVAT(TIMACT,SYSACT,SLPACT,GPFUX,PRODNUI,PRODHL,GFXSEC,ICNC,

```



```

SUBROUTINE INPUT (NONUCL, NOEGPS, NOPAS, NG, NF)
C
C THIS PROGRAM IS USED FOR ENTERING "VARIABLE" DATA INTO MEMORY
C
C ARGUMENTS IN=NONE
C ARGUMENTS OUT=NONUCL, NOEGPS, NOPAS, NG, NF
C
COMMON A, V, F, DF, T1, T2, N, R, D, H, OPTION, IGEOM, ISAB
READ (5, 61) OPTION
61 FORMAT (I4, ?)
IF (OPTION.NE.0.000000E+00) GO TO 62
READ (5, 1) A, V, F, DF, T1, T2, N
1 FORMAT (6E10.3, 2X), I3)
3 READ (5, 2) NOEGPS
2 FORMAT (I2)
READ (5, 2) NOPAS
5 READ (5, 3) R, D, H, IGEOM, ISAB
FORMAT (3(E10.3, 2X), 2(I1))
READ (5, 2) NONUCL
READ (5, 2) NF
GO TO 19
62 READ (5, 63) T1, T2, N
63 FORMAT (2(E10.3, 2X), I3)
GO TO 3
19 WRITE (6, 20)
20 FORMAT (1H1, 40MTHE FOLLOWING DATA HAVE BEEN ENTERED INTO MEMORY, /,
+9M *****//)
WRITE (6, 21)
21 FORMAT (1H0, 85HAMPERES KILO-VOLTS BEAMFRACTION DUTYFA
+CTOR T1 T2 N)
IF (OPTION.EQ.0.000000) WRITE (6, 13) A, V, F, DF, T1, T2, N
13 FORMAT (1H0, E8.3, 4X, E8.3, 8X, E8.3, 11X, E8.3, 4X, E8.3, 3X, I3, //
+1)
IF (OPTION.NE.0.000000) WRITE (6, 65) T1, T2, N
65 FORMAT (1H0, 55H***** *****)
+***, 6X, 2(E8.3, 3X), I3, ///)
WRITE (6, 22) NOEGPS
22 FORMAT (1H, 39HNUMBER OF ENERGY GROUPS=NOEGPS, I2, /)
WRITE (6, 75) NONUCL
75 FORMAT (1H, 39HNUMBER OF TARGET NUCLIDES=NONUCL, I2, /)
WRITE (6, 23) NOPAS
23 FORMAT (1H, 39HNUMBER OF POINTS AFTER SHUTDOWN=NOPAS, I2, /)
WRITE (6, 80) ISAB, IGEOM
80 FORMAT (1X, 5HISAB=, I1, //, 7H IGEOM=, I1, //)
IF (ISAB.EQ.0) WRITE (6, 83)
83 FORMAT (1X, 36HYOU ARE USING A NON-ABSORBING SOURCE, /)
IF (ISAB.EC.1) WRITE (6, 84)
84 FORMAT (1X, 33HYOU ARE USING AN ABSORBING SOURCE, /)
IF (IGEOM.EQ.1) WRITE (6, 85)
85 FORMAT (1X, 25HTHE SOURCE IS A POINT-SOURCE, /, 42H THE OBSERVER IS ST
+ANDING AT A DISTANCE D=.E10.3, 23H METERS FROM THE SOURCE, /)
IF (IGEOM.EQ.2) WRITE (6, 86) R, D
86 FORMAT (1X, 23HTHE SOURCE IS A SPHERE, /, 17H ITS RADIUS IS R=.E10.3,
+7H METERS, /, 42H THE OBSERVER IS STANDING AT A DISTANCE D=.E10.3, 39
+0H METERS FROM THE SURFACE OF THE SPHERE, /)
IF (IGEOM.EQ.3) WRITE (6, 87) R, H, D
87 FORMAT (1X, 25HTHE SOURCE IS A CYLINDER, /, 17H ITS RADIUS IS R=.E10.
+3, 7H METERS, /, 17H ITS HEIGHT IS H=.E10.3, 7H METERS, /, 45H THE DESER
+VER FACES THE AXIS AT A DISTANCE D=.E10.3, 7H METERS, /)

```

```
IF(IGEOM.EQ.4)WRITE(6,99)R,M
88 FORMAT(1X,33HTHIS SOURCE IS A SPHERICAL CAVITY,/,17H ITS RADIUS IS
+ R=.E10.3,7H METERS,/,16H ITS WIDTH IS H=.E10.3,7H METERS,/,44H TH
+ E OBSERVER IS AT THE CENTER OF THE CAVITY,/)
NG=NDEGPS+1
WRITE(6,72) MG
72 FORMAT(1X,39HNUMBER OF GROUPS BOUNDARIES-MG           ,I2,/)
WRITE(6,76) MF
76 FORMAT(1X,39HNUMBER OF FAST NEUTRON GROUPS-MF          ,I2,/)
IF(DPTION.NE.0.000000) WRITE(6,66) OPTION
66 FORMAT(1X,46H AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND,4X,E1
+ 0,4,/)
IF(DPTION.NE.0.00000E+00)GO TO 64
IF(A.LE.0.0) GO TO 36
IF(V.LT.38.0.OR.V.GT.300.0) GC TO 37
IF(F.LE.0.0.OR.F.GT.1.0) GC TO 38
IF(DF.LE.0.0.OR.DF.GT.1.0) GC TO 39
64 IF(1.LE.0.0) GO TO 40
IF(2.LT.0.0) GO TO 41
DC 42 K=1,999
IF(M.EQ.K) GO TO 430
42 CONTINUE
GO TO 44
430 DO 45 K=1,50
IF(ACEGPS.EQ.M) GO TO 460
45 CONTINUE
GO TO 9
460 DO 47 K=1,1E
IF(NCPAS.EQ.K) GO TO 480
47 CONTINUE
GO TO 1E
480 IF(IGEOM.EQ.1)GO TO 60
IF(IGEOM.EQ.2)GO TO 8E
IF(IGEOM.EQ.3)GO TO 8E
IF(IGEOM.EQ.4)GC TO 80
GO TO 67
80 IF(ISAB.EQ.0)GO TO 7C
IF(ISAB.EQ.1)GO TO 7C
GC TO 81
70 IF(R.LT.0.0.OR.D.LT.0.0) GC TO 49
IF(H.LT.0.0) GC TO 77
QP(R/NDEG.EQ.K) GO TO 520
50 CONTINUE
GO TO 11
9 WRITE(6,1E)
1E FORMAT(1X,310HPROGRAM *ACDOS* ABORTED-NUMBER OF ENERGY GROUPS MUS
+ T BE INTEGER AND LESS THAN OR EQUAL TO 58-RECHECK THIS DATA)
STOP
10 WRITE(6,14)
14 FORMAT(1X,110HPROGRAM *ACDOS* ABORTED-NUMBER OF POINTS AFTER SHUT
+ DOWN MUST BE INTEGER AND LESS THAN OR EQUAL TO 12-RECHECK THIS DAT
+ A)
STOP
11 WRITE(6,17)
17 FORMAT(1X,90HPROGRAM *ACIOS* ABORTED-NUMBER OF MCLIDES IN SYSTEM
+ MUST BE INTEGER AND LESS THAN OR EQUAL TO 3C1)
STOP
36 WRITE(6,52)
52 FORMAT(1X,20HPROGRAM *ACICS* ABORTED-CURRENT MUST BE GREATER THAN
+ 0 AMPERES-RECHECK THIS DATA)
STOP
```

```
37 WRITE(6,53)
53 FORMAT(1H ,09HPROGRAM "ACDC" ABORTED-VOLTAGE MUST BE BETWEEN 30 A
+ND 300 KV INCLUSIVE-RECHECK THIS DATA)
STOP
38 WRITE(6,54)
38 FORMAT(1H ,107HPROGRAM "ACDS" ABORTED-BEAM FRACTION MUST BE GREATER
+R THAN 0 AND LESS THAN OR EQUAL TO 1.0-RECHECK THIS DATA)
STOP
39 WRITE(6,55)
55 FORMAT(1H ,107HPROGRAM "ACDS" ABORTED-DUTY FACTOR MUST BE GREATER
+THAN 0.0 AND LESS THAN OR EQUAL TO 1.0-RECHECK THIS DATA)
STOP
40 WRITE(6,56)
56 FORMAT(1H ,70HPROGRAM "ACDS" ABORTED-PULSE LENGTH MUST BE GREATER
+THAN 0.0-RECHECK THIS DATA)
STOP
41 WRITE(6,57)
57 FORMAT(1H ,90HPROGRAM "ACDS" ABORTED-PAUSE LENGTH MUST BE GREATER
+THAN OR EQUAL TO 0.0-RECHECK THIS DATA)
STOP
44 WRITE(6,58)
58 FORMAT(1H ,104HPROGRAM "ACDS" ABORTED-NUMBER OF PULSES MUST BE IN
+TEGER AND LESS THAN OR EQUAL TO 999-RECHECK THIS DATA)
STOP
49 WRITE(6,60)
60 FORMAT(1H ,134HPROGRAM "ACDS" ABORTED-R MUST BE GREATER THAN OR
+EQUAL TO 0.0 AND 0 GREATER THAN 0.0-RECHECK THIS DATA)
STOP
77 WRITE(6,70)
70 FORMAT(1H ,80HPROGRAM "ACDS" ABORTED-M MUST BE GREATER THAN OR EQ
+UAL TO 0.0-RECHECK THIS DATA)
STOP
67 WRITE(6,68)
68 FORMAT(1H ,664HPROGRAM "ACDS" ABORTED-IGEOM MUST BE 1,2,3, OR 4-RECH
+ECK THIS DATA)
STOP
81 WRITE(6,82)
82 FORMAT(1H ,61HPROGRAM "ACDS" ABORTED-ISAB MUST BE 0, OR 1-RECHECK
+THIS DATA)
STOP
510 RETURN
END
```

SUBROUTINE AFRAY (BFLUX,ZNAM1,IGNO,MASS,STAFS,CP,NOEGPS,NCNUCL,NDPAS,  
PAS,KG,THASS,ZNAM2)

C  
C  
C  
C  
C  
C  
C  
C  
C

THIS PROGRAM IS USED FOR ENTERING "ARRAY" DATA INTO MEMORY

ARGUMENTS IN=STAFS,NOEGPS,NCNUCL,NDPAS,NG  
ARGUMENTS OUT=BFLUX,ZNAM1,IGNO,MASS,GP

```

DIMENSION BFLUX(NDGPS),ZNAM1(NCNUCL),IGNO(NCNUCL),MASS(NCNUCL),ST
*AFS(NDGPS),GP(NG),ZNAM2(NDNUCL)
REAL MASS,KWALL
COMMON A,V,F,DF,T1,T2,A,F,C,M,OPTION,IGEOM,ISAE
R4CD(5,-1)(STAFS(I),I=1,NDPAS)
4 FORMAT (E10.3,2X)
DO 7 I=1,NCNUCL
  REAC(5,-1) ZNAM1(I),ZNAM2(I),IGNO(I),MASS(I)
  9 FORMAT (2A10,2X,E10.3)
7 CONTINUE
  REAC(5,-1)GP(I),I=1,NG)
  DO 80 I=1,NDNUCL-1
    IDMIN=IGNO(I)
    IMIN=I
    DO 90 J=I+1,NDNUCL
      IF (IDMIN.LE.IDN(J))GO TO 90
      IDMIN=IDN(J)
      IMIN=J
90 CONTINUE
    IF (I.EQ.IMIN) GO TO 80
    IDTEMP=IDN(I)
    IDN(I)=IDN(IMIN)
    IDN(IMIN)=IDTEMP
    TEMPM=MASS(I)
    MASS(I)=MASS(IMIN)
    MASS(IMIN)=TEMPM
    Z1TEMP=ZNAM1(I)
    ZNAM1(I)=ZNAM1(IMIN)
    ZNAM1(IMIN)=Z1TEMP
    Z2TEMP=ZNAM2(I)
    ZNAM2(I)=ZNAM2(IMIN)
    ZNAM2(IMIN)=Z2TEMP
80 CONTINUE
    WRITE (6,B4)
84 FORMAT (/,1H .33/SPECIFIC TIMES AFTER SHUTDOWN (1),/)
    DO 34 I=1,NDPAS
      WRITE (6,3E) I,STAFS(I)
35 FORMAT (1H .4HTIME,1Y,1Z,10A,E10.3)
36 CONTINUE
    REAC(5,-1)BFLUX,KNALL,KWALL
    10 FORMAT (2 (I2,1X),E10.3)
    WRITE (6,11)BFLUX
    11 FORMAT (/,28H OPTION ON THE FLUXES-IFLUX=,I2)
    IF (IFLUX.EQ.0)WRITE (6,12)
    22 FORMAT (9H THE FLUXES WILL BE READ FROM THE INPUT)
    IF (IFLUX.EQ.1)WRITE (6,13)
    13 FORMAT (36H THE FLUXES WILL BE READ FROM TAPE B)

```

```

IF((IFLUX.NE.1).AND.(IFLUX.NE.2))GO TO 40
WRITE(6,24)IWALL
14 FORMAT(//,49H OPTION ON THE PART OF THE ROOM WALLS ACTIVATION,/,19H
+ TAKES PLACE-IWALL=.22)
IF(IWALL.EQ.0)WRITE(6,15)
15 FORMAT(//7H ACTIVATION OF AN OBJECT PLACED INSIDE THE ROOM)
IF(IWALL.EQ.1)WRITE(6,16)
16 FORMAT(//3EH ACTIVATION OF THE WALLS OF THE ROOM)
IF((IWALL.NE.0).AND.(IWALL.NE.1))GO TO 41
GO TO(17,18)(IFLUX+1)
17 READ(5,4)(BFLUX(I),I=1,NCEGPS)
GO TO 19
18 READ(8,4)(BFLUX(I),I=1,NCEGPS)
19 GO TO(20,21)(IWALL+1)
21 R=R*KWALL*100.
DK=C*KWALL*100.
FF=3*R*(1-(R*K+K)/SINH(DK))+R*K/TANH(DK)/((R*K+K)**3-R*K**3)
DO 22 I=1,NCEGPS
BFLUX(I)=FF*BFLUX(I)
22 CONTINUE
WRITE(6,27)KWALL
23 FORMAT(//,52H AVERAGE UNIT NORMALIZED FLUXES IN THE WALLS (KWALL=
+.E10,3,1H),/)
GO TO 24
24 WRITE(6,27)
27 FORMAT(//,39H UNIT NORMALIZED FLUXES INSIDE THE ROOM,/)
DO 22 I=1,NCEGPS
I=NCEGPS+1-I
WRITE(6,28)I,BFLUX(I)
28 FORMAT(1H+GROUP,1X,I2,1GX,E11.4)
32 CONTINUE
WRITE(6,73)
73 FORMAT(//,1H+22H GROUP ECUNARIES (MEV),/)
WRITE(6,74)(I,GP(NG+1-I),I=1,NG)
74 FORMAT(1H+8H BOLDIFY,1X,I3,5X,E11.4)
WRITE(6,29)
29 FORMAT(1H+//,1H+10X,4HNAME,10X,9HID=NUMBER,4X,8HMASS(KG),/)
DO 30 I=1,NONUCL
WRITE(6,31)ZNAM1(I),ZNAM2(I),IDNO(I),MASS(I)
31 FORMAT(1X,2A10,5X,I5,5X,E8.3)
30 CONTINUE
THASS=C.
DO 30 I=1,NONUCL
THASS=THASS+MASS(I)
50 CONTINUE
WRITE(6,25)
25 FORMAT(//,55X,25H CODS IS NOW READY TO RUN,/,55X,25H*****
+******)
GO TO 26
40 WRITE(6,42)
42 FORMAT(//,63H PFCGRAM *ACDGS* ABORTED-IFLUX MUST BE 0 OR 1-RECHECK
+ THIS DATA)
STOP
41 WRITE(6,43)
43 FORMAT(//,63H PFCGRAM *ACDGS* ABORTED-IWALL MUST BE 0 OR 1-RECHECK
+ THIS DATA)
STOP
26 RETLBN
&P

```



```

          SUBROUTINE SOURCE (AVENPS)
C
C
C THIS PROGRAM CALCULATES THE SOURCE NEUTRON PRODUCTION RATE
C
C
C ARGUMENTS IN=NDNL
C ARGUMENTS OUT=AVENPS
C
      REAL INSAPS
      COMMON /4,V,F,DF,T1,T2,N,F,L,M,OPTION,IGEOM,ISAE
      IF (V.GE.30.0.AND.V.LE.80.0) GO TO 3
      IF (V.GT.80.0.AND.V.LE.150.0) GO TO 4
      IF (V.GT.180.0.AND.V.LE.300.0) GO TO 5
3      CF=(3.25917*V**2.83194)/38.5
      GO TO 7
4      CF=(11.51725*V**2.42306)/38.5
      GO TO 7
5      CF=(6.64661*V**2.14132)/38.5
7      INSAPS=(9.64E+6)* (4)*(1.0E+6)*F*(CF*1.0E-4)
      AVENPS = INSAPS*DF
      WRITE (6,1)
9      FORMAT (1/1M ,43HTHE FOLLOWING HAS BEEN CALCULATED BY ACOS2,44X
4      *.....//)
      WRITE (6,1) INSAPS
1      FORMAT (1/1M ,52HINSTANTANEOUS NUMBER OF NEUTRONS PRODUCED PER SEC
*END*4X,11.4)
      WRITE (6,2) AVENPS
2      FORMAT (1/1M ,46H AVERAGE NUMBER OF NEUTRONS PRODUCED PER SECOND*4X,
*E11.4)
      RETURN
      END

```

```
      SUBROUTINE GFFLUX(AVENFS,NDLGPS,GPFLUX,BFLUX)
C
C  THIS PROGRAM CALCULATES UP TO 50 GROUP FLUXES
C
C  ARGUMENTS IN-AVENFS,NOEGPS,BFLUX
C  ARGUMENTS OUT-GPFLUX
C
C
      DIMENSION GPFLUX(NOEGPS),BFLUX(NOEGPS)
      DO 1 I = 1,NOEGPS
      GPFLUX(I) = BFLUX(I)*AVENFS
1  CONTINUE
      RETURN
      END
```

```

SUBROUTINE WFLX(BFLUX,GP,CONST,FCONST,NOEGPS,NF,NG)
C
C THIS PROGRAM DETERMINES THE CONSTANTS ASSOCIATED WITH THE ASSUMED
C WEIGHTING FUNCTIONS
C
C
C
C ARGUMENTS IN=BFLX,GP,NOEGPS,NF,NG
C ARGUMENTS OUT=CONST,FCONST
C
C
C DIMENSION BFLUX(NOEGPS),GP(NG),CONST(NOEGPS),FCONST(NOEGPS)
DO 10E JC=1,NG
GP(JC)=GF(JC)*1.000E5
10E CONTINUE
C CALCULATE THE THERMAL GROUP CONSTANT
ICOUNT=1
SUMINT=0.0
WIDTH=(GF(2)-GF(1))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(1)*ZINCR
DO 1 I=1,200
PARSUM=SQRT(DPVAR)*EXP(-(DPVAR/0.025)*WIDTH)
SUMINT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
1 CONTINUE
CONST(ICOUNT)=BFLX(ICOUNT)/SUMINT
IF(NOEGPS.EQ.1) GO TO 1000
IF(NGEQPS.EQ.2) GO TO 100
C CALCULATE THE 1/E GROUP CONSTANT(S)
ICOUNT=ICOUNT+1
IEPTH=0
2 SUMINT=0.0
WIDTH=(GF(ICOUNT+1)-GF(ICOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GF(ICOUNT)*ZINCR
DO 3 I=1,200
PARSUM=(1.0/DPVAR)*WIDTH
SUMINT=SUMINT+PARSUM
DPVAR=DPVAR+WIDTH
3 CONTINUE
CONST(ICOUNT)=BFLX(ICOUNT)/SUMINT
IEPTH=IEPTH+1
NEPTH=NOEGPS-1-NF
IF(IEPTH.EQ.NEPTH) GO TO 150
ICOUNT=ICOUNT+1
GO TO 2
100 ICOUNT=ICOUNT+1
C THE 1 GROUP FAST SPECTRUM WILL BE MATCHED TO THE MAXWELLIAN AT THE
CGROUP BOUNDAry
FLUXBO=CONST(1)*SQRT(GP(2))*EXP(-(GP(2)/0.025))
ZNUMH=GUESS1=BFLX(INOEGPS)/FLUXBO
ITER=1
E FXSLBI=ZNUMH/(EXP((GP(3)-GP(2))/GUESS1)-1)-GUESS1
XSLBI=GUESS1
FPRIME=ZNUMH*(GP(3)-GP(2))*EXP((GP(3)-GP(2))/GUESS1)/GUESS1**2*(
+EXP((GP(3)-GP(2))/GUESS1)-1)**2)-1
XSEIP1=(SUEI-FXSLBI/FPRIME)

```

```

DIFF=ABS(XSBIP1-GUESS1)
IF(DIFF.LT.0.1) GO TC 5
ITER=ITER+1
IF(ITER.GT.100) GO TO 75
GUESS1=XSBIP1
GC TC 6
75 WRITE(6,76)
76 FORMAT(1X,'.116H***PROGRAM ABORTED***THE TRANSCENDENTAL EQUATION
+USED TO MATCH THE THERMAL GROUP TO FAST GROUP IS NOT CONVERGING)
WRITE(6,27)
STOP
5 B=XSBIP1
A=BFLUX(INOEGPS)/(B*(1-EXP((GF(2)-GP(3))/B)))
CONST(ICOUNT)=A
FCOAST(ICOUNT)=B
GC TO 100
150 IF(INF.GT.2) GO TC 16J
CTHE 1 GROUP FAST SPECTRUM WILL BE MATCHED TO THE LAST 1/E GROUP
ITER=1
NGP1=NCEGPS+1
ZNUMER=GUESS1=(BFLUX(INOEGPS)*GP(INOEGPS)/CONST(ICOUNT))
91 FXSLBI=ZNUMER/(EXP((GF(NGP1)-GP(INOEGPS))/GUESS1)-1)-GUESS1
XSUBI=GUESS1
FPRIME=ZNUMER*(GF(NGP1)-GP(NCEGPS))*EXP((GF(NGP1)-GP(INOEGPS))/GUES
+S1)/GUESS1**2*(EXP((GF(NGP1)-GP(INOEGPS))/GUESS1)-1)**2)-1
XSBIP1=XSUBI-FXSLBI/FPRIME
DIFF=ABS(XSBIP1-GUESS1)
IF(DIFF.LT.0.1) GO TO 92
ITER=ITER+1
IF(ITER.GT.100) GO TO 25
GUESS1=XSBIP1
GC TO 91
92 B=XSBIP1
A=BFLUX(INOEGPS)/(B*(1-LXP((GF(INOEGPS)-GP(NSP1))/B)))
CONST(ICEGPS)=A
FCOAST(ICEGPS)=B
GC TO 100
16J ICP1=ICOUNT+1
CTHE FIRST OF 2 CP MORE FAST GROUPS WILL BE MATCHED TO THE LAST 1/E
CGROUP
ICP2=ICOUNT+2
ITER=1
ZNUMER=GUESS1=(BFLUX(ICP1)*GP(ICP1)/CONST(ICOUNT))
20 FXSLBI=ZNUMER/(EXP((GF(ICP2)-GP(ICP1))/GUESS1)-1)-GUESS1
XSUBI=GUESS1
FPRIME=ZNUMER*(GF(ICP2)-GP(ICP1))*EXP((GF(ICP2)-GP(ICP1))/GUESS1)/
+GUESS1**2*(EXP((GF(ICP2)-GP(ICP1))/GUESS1)-1)**2)-1
XSBIP1=XSUBI-FXSLBI/FPRIME
DIFF=ABS(XSBIP1-GUESS1)
IF(DIFF.LT.0.1) GO TO 30
ITER=ITER+1
IF(ITER.GT.100) GO TO 25
GUESS1=XSBIP1
GC TO 20
25 WRITE(6,26)
26 FORMAT(1X,'.116H***PROGRAM ABORTED***THE TRANSCENDENTAL EQUATIONS
+WHICH MATCH 1/E TO FAST OR FAST TO FAST GROUPS ARE NOT CONVERGING)
WRITE(6,27)
27 FORMAT(1X,'.75HMCST LIKELY PROBLEM IS UNREALISTIC OP UNSAFELY
+AND GROUP BOUNDARY DATA)
STOP
30 B=XSBIP1

```

```

A=BFLUX(ICF1)/(E*(1-EXP((GP(ICP1)-GF(ICP2))/E1)))
ICOUNT=ICOUNT+1
ICP1=ICOUNT+1
ICF2=ICCLAT+2
CONST(ICCOUNT)=A
FCNST(ICCOUNT)=B
IF (ICOUNT*.10.NOF GPS) GO TO 1000
ZNUPE=BFLUX(ICF1)/CONST(ICCOUNT)
ITER=1
GO TO 20

1000 WRITE (6,11)
11 FORMAT (//,1H ,63H      INTEGRALS OF THE FITTING FUNCTION OVER THE G
+FOUF INTERVALS)
WRITE (6,12)
12 FORMAT (1F ,7=HGROUP LOWER BOUNDARY UPPER BOUNDARY REQUIRED INTEGR
+AL CALCULATED INTEGRAL,/)
CCALCULATE THE THERMAL INTEGRAL
ICOUNT=1
SUMINT=0.0
WIDTH=(GP(2)-GP(1))/200.0
ZINCR=WIDTH/2.0
DPVAR=GP(1)+ZINCR
DO 21 I=1,200
PARSLM=CONST(1)*SCT(DPVAR)*EXP(-(DPVAR/0.025))*WIDTH
SUMINT=SUMINT+PARSLM
DPVAR=DPVAR+WIDTH
21 CONTINUE
IF (ICOUNT.EQ.NOE GPS) GO TO 1009
IF (ICOUNT.EQ.2) GO TO 46
JJ=NG-ICOUNT
WRITE (6,22) JJ,GP(ICOUNT),GP(ICOUNT+1),BFLUX(ICOUNT),SUMINT
22 FORMAT (1F ,2),I2,5),E10.3,4X,E10.3,5X,E10.3,9X,E10.3)
ICOUNT=ICOUNT+1
CCALCULATE THE 1/2 INTEGRAL(S)
45 SUMINT=0.0
WIDTH=(GF(ICOUNT+1)-GF(ICOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GF(ICOUNT)+ZINCR
DO 23 I=1,200
PARSLM=(CONST(ICCOUNT)/DPVAR)*WIDTH
SUMINT=SUMINT+PARSLM
DPVAR=DPVAR+WIDTH
23 CONTINUE
JJ=NG-ICOUNT
WRITE (6,22) JJ,GP(ICOUNT),GP(ICOUNT+1),BFLUX(ICOUNT),SUMINT
ICOUNT=ICOUNT+1
IF (ICOUNT.LE.NEFTM+1) GO TO 45
GO TO 16

46 JJ=2
WRITE (6,22) JJ,GP(ICOUNT),GP(ICOUNT+1),BFLUX(ICOUNT),SUMINT
IF (NG.EQ.2) ICOUNT=ICOUNT+1
15 SUMINT=0.0
CCALCULATE THE FAST INTEGRAL(S)
WIDTH=(GF(ICOUNT+1)-GF(ICCOUNT))/200.0
ZINCR=WIDTH/2.0
DPVAR=GF(ICOUNT)+ZINCR
DO 24 I=1,200
PARSLM=CONST(ICCOUNT)*EXP((DPVAR-GP(ICOUNT+1))/FCNST(ICCOUNT))*WIDT
44 SUMINT=SUMINT+PARSLM
DPVAR=DPVAR+WIDTH
24 CONTINUE

```

```
JJ=NG-ICOUNT
WRITE (6,22) JJ,GF(ICOUNT),GP(ICOUNT+1),BFLUX(ICOUNT),SUMINT
ICOUNT=ICOUNT+1
IF(ICOUNT.LE.NCGPS) GO TO 16
GO TO 15
109 JJ=1
WRITE (6,22) JJ,GF(ICOUNT),GP(ICOUNT+1),BFLUX(ICOUNT),SUMINT
15 DO 105 JC=1,NG
GP(JC)=GF(JC)/1.005E6
105 CONTINUE
RETURN
END
```

```

SUBROUTINE ACTVAT(TI*ACT,SYS*ACT,SUM*ACT,GP*FLUX,F*CDNU,PR*CDHL,G*FXSEC
+ ,ID*AC,P*ASS,S*TAFS,GP,CONST,F*CONST,B*FLUX,ND*NUCL,A*DEG*PS,ND*PAS,NG,N*F)
C
C THIS PROGRAM CALCULATES ACTIVITIES DUE TO NEUTRON INDUCED REACTIONS
C
C ARGUMENTS IN=GP*FLUX,GP*YSEC,ID*NO,P*ASS,S*TAFS,GP,CONST,F*CONST,B*FLUX
C ARGUMENTS IN=ND*NUCL,A*DEG*PS,ND*PAS,NG,N*F
C ARGUMENTS OUT=TI*ACT,SYS*ACT,PR*CDNU,PR*CDHL
C
C DIMENSION IDNO(A*NUCL),PASS(A*NUCL),STAFS(N*CPAS)
C DIMENSION TI*ACT(N*NUCL,13,N*CPAS),SYS*ACT(N*CPAS),SUM*ACT(N*CPAS),GP*FL
+ UY(A*DEG*PS),PR*CDHL(A*NUCL,13),PR*CDHL(A*NUCL,13),GP(N*G),GP*XSEC(A*DEG*
+ S),CONST(A*DEG*PS),F*CONST(A*DEG*PS),B*FLUX(A*DEG*PS)
C COMMON A,V,F,DF,T1,T2,M,N,F,E,M,OPTION,I*GEOM,IS*AE
REAL PASS
WRITE(E,102)
102 FORMAT(///,1P,'57THE FOLLOWING OUTPUT DATA ARE CALCULATED ACTIVIT
+ IES IN 3C)
I=1
J=0
C INITIALIZE FLAGS AND ARRAYS
DO 40 M=1,ND*NUCL
DO 40 MM=1,13
PR*CDHL(M,MM)=0.000000E+00
PR*CDHL(M,MM)=0.000000E+00
40 CONTINUE
DO 41 M=1,ND*NUCL
DO 41 MM=1,13
DO 41 P*MM=1,N*CPAS
TI*ACT(M,P*MM)=0.0000E+00
41 CONTINUE
DO 22 J*1,N*CPAS
SYS*ACT(J*1)=SUM*ACT(J*1)=0.0
22 CONTINUE
K*FLAG=LF*FLAG=C
C PRINT HEADINGS
WRITE(E,37)1STAFS(K),K=1,N*CPAS)
37 FORMAT(1P,1P,4X,7HTIME(H),4X,3D(E8.3,1X),2X,2(E9.3,2Y))
WRITE(E,41)
1 FORMAT(1P,125HTARGET PR*CDHL ACT(I1) ACT(I2) ACT(I3) ACT(I4)
+ ACT(I5) ACT(I6) ACT(I7) ACT(I8) ACT(I9) ACT(I10) ACT(I11)
+ ACT(I12))
2 READ(17,25) I*TAPE,T*ARMAS
25 FORMAT(2E,7X,2E,11,4)
IF(I*TAPE.EQ.ID*NO(I)) GO TO 7
IF(I*TAPE.GT.ID*NO(I)) GO TO 44
CALL PCS*ITCN
GO TO 2=
C THE ABOVE STATEMENTS SL*FROM THE TAPE FOR A MATCHING TARGET ID NUMBER
4= WRITE(E,45) ID*NO(I)
45 FORMAT(///,32H ATTENTION-TARGET NUC*LUIDE NUMBER,1E,9H DOES NOT EXI
+ ST IN DATA LIBRARY-THEREFORE=NO CONTRIBUTION FROM THIS NUC*LUIDE IN
+ D*CASE CALLS,///)
IF(I*TAPE.EQ.ID*NO(I+1)) GO TO 46
I=I+1

```

```

      CALL PCSICN
CPRINT HEADINGS
WRITE (6,37) (STAFS(K),K=1,NOPAS)
WRITE (6,1)
GO TO 24
46 I=I+1
CPRINT HEADINGS
WRITE (6,37) (STAFS(K),K=1,NCFAS)
WRITE (6,1)
7 J=J+1
FEAC(7,5) PRCDNU(I,J)*FRCDML(I,J)
FEAC(7,95) NP
IF (ITAPE.EQ.IDNC(ACNUCL)) LFLAG=1
99 FOPMAT(3)
IF (FRCDML(I,J).EQ.0.100DE+50) GO TO 50
8 FORMAT (2D3,1E12.4,1E12.4)
CALL AVRAGE(GPXSEC,CNST,FCCNST,BFLUX,GP,NDEGPS,NG,NF,NP)
IF (GPXSEC(1).GT.0.50DE+49) GO TO 103
SUM1A=SUM2A=C.C
DO 11 J=1,N
TEMPS=EXP(-(0.693*(N-JL)*((T1*3600.0)+(T2*3600.0)))/PRCDML(I,J)
*)
SUM1A=SUM1A+TEMPS
11 CONTINUE
DO 12 L=1,NDEGPS
TEMPS=SUM1*(1.0E-24)*GPYSEC(L)*GPF LUX(L)
SUM2A=SUM2A+TEMPS
12 CONTINUE
ACT=(MASS(I)*1000.0/TARFAS)*E.Q.23E+23*SUM2A*(1-EXP(-0.693*3600.0*T
+1/PROCDML(I,J)))
DO 13 K=1,NCFAS
TIMACT(I,J,K)=ACT*L.VD(-(STAFS(K)*C.693*3600.0/FRCDML(I,J)))
SUMACT(K)=SUMACT(K)+TIMACT(I,J,K)
13 CONTINUE
103 WRITE (6,20) IDNC(I),FIX(PRCDNU(I,J)),(TIMACT(I,J,K),K=1,NCFAS)
20 FORMAT (1X,1F15.2X,1F5.2X,10(E8.3,1X),2X,2(E8.3,2X))
READ(7,25) ITAPE,TARFAS
IF (EOF(7)) NE.G) KFLAG=1
IF (KFLAG.EQ.1) GO TO 27
IF (LFLAG.EQ.1.AND.ITAPE.NE.IDNC(NDNUCL)) GO TO 27
IF (ITAPE.EQ.IDNC(ACNUCL)) LFLAG=1
IF (ITAPE.NE.IDNC(I)) GO TO 27
GO TO 7
CPRINT HEADINGS
27 WRITE (6,37) (STAFS(K),K=1,NCFAS)
WRITE (6,14)
14 FORMAT (1F12EM TARGET ACT(T1) ACT(T2) ACT(T3) ACT(T4)
+ ACT(T5) ACT(T6) ACT(T7) ACT(T8) ACT(T9) ACT(T10) ACT(T11)
+ACT(T12))
WRITE (6,15) IDNC(I),(SUMACT(K),K=1,NCFAS)
15 FORMAT (1X,5X,1F5.2X,10(E8.3,1X),2X,2(E8.3,2X),///)
DO 10 K=1,NCFAS
SYSACT(K)=SYSACT(K)+SUMACT(K)
10 CONTINUE
IF (ITAPE.EQ.IDNC(I+1)) GO TO 32
IF (LFLAG.EQ.1) GO TO 16
32 DO 23 JN=1,NCFAS
SUMACT(JN)=0.0
23 CONTINUE
IF (ITAPE.EQ.IDNC(I+1)) GO TO 31
IF (ITAPE.GT.IDNC(I+1)) GO TO 47
I=I+1

```



```
CALL PCSITON
CPRINT HEADINGS
WRITE (6,37)(STAFS(I),K=1,NOPAS)
WRITE (6,1)
J=0
GO TO 24
47 WRITE (6,45) IDNO (I+1)
CALL PCSITON
I=I+2
IF (ITAPE.GT.IDNC(I)) WRITE (6,45) IDNC (I)
IF (ITAPE.GT.IDNO (I)) I=I+1
CPRINT HEADINGS
WRITE (6,37)(STAFS(K),K=1,NOPAS)
WRITE (6,1)
J=0
GO TO 24
CPRINT HEADINGS
31 WRITE (6,37)(STAFS(K),K=1,NOPAS)
WRITE (6,1)
I=I+1
J=0
GO TO 7
CPRINT HEADINGS
16 WRITE (6,544)
54- FORMAT ('/')
WRITE (6,37)(STAFS(K),K=1,NOPAS)
WRITE (6,21)
21 FORMAT (1X,125H SYSTEM ACT(T1) ACT(T2) ACT(T3) ACT(T4)
+ ACT(T5) ACT(T6) ACT(T7) ACT(T8) ACT(T9) ACT(T10) ACT(T11)
+ACT(T12))
WRITE (6,51)(SYSACT(K),K=1,NOPAS)
51 FORMAT (1X,15Y,10(E3.3,1X),2X,2(E3.2X))
IF (K.LAG.50.) GO TO 42
RETLRN
42 CALL PCSITON
43 READ(7,25) ITAPE,ITAFMAS
IF (EQ(' ',N2.0)) RETLRN
CALL PCSITON
GO TO 43
53 CALL PCSITON(2(NP))
GO TO 103
END
```

```

SUBROUTINE AVERAGE (GPXSEC,CONST,FCNST,BFLUX,GP,NCEGPS,NG,NF,NP)
C
C
C THIS PROGRAM CALCULATES FLUX WEIGHTED GROUP CROSS SECTIONS BY AVERAGING
C MICROSCOPIC CROSS SECTION DATA
C
C
C ARGUMENTS IN=CONST,FCNST,BFLUX,GP,NCEGPS,NG,NF,NP
C ARGUMENTS = GPXSEC
C
C
C DIMENSION GPXSEC (NCEGPS),CONST (NCEGPS),FCNST (NCEGPS),BFLUX (NCEGPS
+1,CFIN)
C DIMENSION CS (150),EN (150),CS1 (201),E1 (201)
C INITIALIZE ARRAY GPXSEC
DO 41 J=1,NCEGPS
GPXSEC (J)=0.0
41 CONTINUE
C LOAD ENERGY AND CROSS SECTION DATA FROM THE ACTLME LIBRARY
REAC (7,3) (EN (I),CS (I)),I=1,NP
3 FORPAT (E (E1 (4)))
C DETERMINE IF THE ENERGY OF THE FIRST ENERGY-CROSS SECTION PAIR READ
CFROM THE ACTLME LIBRARY IS BELOW THE GREATEST GROUP BOUNDARY
IF (CN (1)-GP (NG)) 125,102,102
105 CALL SMCCT (GP,NG,EN,CS,NP,E1,CS1,NE1)
I=1
J=1
SUMINT=0.0
C DETERMINE IF THE ENERGY AT WHICH THE CROSS SECTION WAS MEASURED IS
CBELOW THE NEXT GROUP BOUNDARY
50 IF (E1 (2).LT.GP (J+1)) GO TO 51
GPXSEC (J)=0.0
J=J+1
GO TO 50
52 IF (E1 (I).LT.GP (J+1)) GO TO 51
C CALCULATE THE GROUP CROSS SECTION
GPXSEC (J)=SUMINT/BFLUX (J)
IF (J.EQ.NCEGPS) GO TO 106
J=J+1
SUMINT=0.0
GO TO 52
51 IF (J.EQ.1) GO TO 1003
IF (NCEGPS.GE.3.AND.J.GT.1.AND.J.LE.(NCEGPS-NF)) GO TO 2000
GO TO 3000
1003 SUMCEL=J.0
C AVERAGE MICROSCOPIC CROSS SECTION DATA OVER THERMAL ENERGIES
WIDTH=(E1 (I+1)-E1 (I))/200.0
ZINCF=WIDTH/2.0
DPVAR=E1 (I)+ZINCF
DO 1001 KK=1,200
FAESUM=(CONS (J)*SOFT (DPVAR*1.0E 6)*EXP (- (DPVAR*1.0E 6/0.025)))*(CS1
+ (I+1)-CS1 (I))* (DPVAR-E1 (I))/(E1 (I+1)-E1 (I))+CS1 (I))*WIDTH*1.0E 6
SUMCEL=SUMDEL+FAESUM
DPVAR=DPVAR+WIDTH
1001 CONTINUE
SUMINT=SUMINT+SUMDEL
I=I+1
GO TO 52

```

```

2000 SUMDEL=0.0
CAVERAGE MICROSCOPIC CROSS SECTION DATA OVER INTERMEDIATE ENERGIES
WIDTH=(E1(I+1)-E1(I))/200.0
ZINCR=WIDTH/2.0
DPVAR=E1(I)+ZINCR
DO 2001 KK=1,200
PARSLP=(CONST(J)/(DPVAR*1.0E6))*WIDTH*1.0E6*(CS1(I+1)-CS1(I))*OP
+VAF-E1(I)/(E1(I+1)-E1(I))+CS1(I))
SUMDEL=SUMDEL+PARSLP
DPVAR=DPVAR+WIDTH
2001 CONTINUE
SUMINT=SUMINT+SUMDEL
I=I+1
GO TO 52
3000 SUMDEL=0.0
CAVERAGE MICROSCOPIC CROSS SECTION DATA OVER FAST ENERGIES
WIDTH=(E1(I+1)-E1(I))/200.0
ZINCR=WIDTH/2.0
DPVAR=E1(I)+ZINCR
DO 3001 KK=1,200
PARSLP=CONST(J)*EXP((DPVAR-OP(J+1))*1.0E6/CONST(J))*WIDTH*1.0E6*(
+CS1(I+1)-CS1(I))*DPVAR-E1(I)/(E1(I+1)-E1(I))+CS1(I))
SUMDEL=SUMDEL+PARSLP
DPVAR=DPVAR+WIDTH
3001 CONTINUE
SUMINT=SUMINT+SUMDEL
I=I+1
GO TO 52
102 GPYSEC(1)=1.0E+50
106 RFAD(7,5) 13(M)
5  FORMAT(7Y,1:1)
RETLFN
END

```

SUBROUTINE SMCOTR(X1,NX1,X2,Y2,NX2,X3,Y3,NX3)

C  
C  
C  
C  
C  
C  
C  
C  
C  
C  
C  
C

THIS PROGRAM CALCULATES CROSS SECTIONS AT GROUP BOUNDARIES BY LINEAR  
INTERPOLATION.

ARGUMENTS IN-GP,NG,EN,CS,NP  
ARGUMENTS OUT-E1,CS1,NE1

DIMENSION X1(1),Y2(1),Y2(1),X3(1),Y3(1)  
K=1  
12 IF (X2(K)-Y1(I)) 20,11,13  
20 K=K+1  
GO TO 12  
11 X3(1)=X2(K)  
Y3(1)=Y2(K)  
K=K+1  
L=2  
I1=2  
GO TO 14  
13 I1=2  
IF (X2(K)-LT.>1(I1)) GO TO 11  
24 I1=I1+1  
IF (Y2(K)-LT.>1(I1)) GO TO 25  
GO TO 24  
25 Y3(1)=X2(K)  
Y3(1)=Y2(K)  
K=K+1  
L=2  
14 DO 1 I=I1,NX1  
5 IF (X2(K)-Y1(I)) 2,3,4  
2 X3(L)=X2(K)  
Y3(L)=Y2(K)  
L=L+1  
K=K+1  
IF (K=NX2) 5,5,10  
3 X3(L)=X2(K)  
Y3(L)=Y2(K)  
L=L+1  
K=K+1  
IF (K=NX2) 1,1,10  
4 X3(L)=X1(I)  
Y3(L)=Y3(L-1)+((Y2(K)-Y3(L-1))/(X2(K)-X3(L-1)))\*X3(L)-X3(L-1)  
L=L+1  
1 CONTINUE  
10 NX3=L-1  
RETURN  
END

```
      SUBROUTINE ADVAN(NOMODE)
C
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE LEVDEC
CLIBRARY
C
C
C   ARGUMENTS IN-NOMODE
C   ARGUMENTS OUT-NONE
C
C
      DD 305 JB=1,NOMODE
      READ(7,2(2) IDUWE
202  FORMAT(36X,I2)
305  CONTINUE
      RETURN
      END
```

```
      SUBROUTINE POSITON
C
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE ACTLEMFC
CLIBRARY
C
C
C ARGUMENTS IN-NONE
C ARGUMENTS OUT-NONE
C
C
      READ(7,30) NP
30  FORMAT(/,I3)
      NP1=INT(FLOAT(NP)/3.000+0.000)
      DO 52 ICDUNT=1,NP1
      READ(7,53) REALDUM
53  FORMAT(E11.4)
52  CONTINUE
      READ(7,5) IDUMB
5   FORMAT(71X,I1)
      RETURN
      END
```

```
      SUBROUTINE PCSIT2(API)
C
C
C THIS PROGRAM IS USED FOR POSITIONING THE FILE MARKER IN THE ACTLEMPF
CLIBRARY
C
C
C   ARGUMENTS IN-NP
C   ARGUMENTS OUT-NCNE
C
C
      NP1=INT(FLCAT(API)/3.000+0.000)
      CC 52 ICCUNT=1+NP1
      REAC(7,53) REALDLM
53  FORMAT (11.4)
52  CONTINUE
      REAC(7,5) IDLMB
5   FORMAT (7IX,11)
      RETURN
      END
```

```

SUBROUTINE DOSRTE (TIMACT,PRODNJ,NO NUCL,NOPAS,DOSSUM,SYSDOS,IDNO,ST
*AFS,PROJHL,DOSNUC,THASS,DOSEGT,DOSPRO)
C
C
C THIS PROGRAM CALCULATES DOSE RATES AS A FUNCTION OF COMPONENT COMPOSIT-
C ION, GEOMETRY, AND TIME AFTER SHUTDOWN
C
C
C
DIMENSION TI*ACT (NO NUCL,13,NOPAS),PRODNJ (NO NUCL,13),DOSSUM (NOPAS),
*SYSDOS (NOPAS),IDNO (NO NUCL),STAFS (NOPAS),DOSEGT (NOPAS),DOSPRO (NOPAS
*)
DIMENSION PROJHL (NO NUCL,13),DOSNUC (NOPAS),DOSDAU (390),DMA (390)
DIMENSION NT (390),NP (390),AC (390),MA (390),MD (390),DOSPRO (390)
DIMENSION IDLKM (30),IDAU (30),DLEVL (30),PRJB (30)
COMMON A,V,F,DF,T1,T2,N,F,C,H,OPTION,IGEOM,ISAB
REAL MULT
WRITE (6,100)
100 FORMAT (///,1H ,61THE FOLLOWING OUTPUT DATA ARE CALCULATED DOSE
*RATES IN MFEM/H,/)
DO 1 M=1,NOPAS
DOSPRO (M)=0.0
DOSEGT (M)=0.0
SYSDOS (M)=0.0
DOSNUC (M)=0.0
DOSSUM (M)=0.0
1 CONTINUE
DO 101 I=1,390
NT (I)=0
NP (I)=0
MD (I)=0
AC (I)=0.0
DOSPRO (I)=0.0
MA (I)=0.0
DMA (I)=0.0
DOSDAU (I)=0.0
101 CONTINUE
DO 2 I=1,NO NUCL
DO 2 J=1,13
K=(I-1)*13+J
NP (K)=IF IX (PRODNJ (I,J))
MA (K)=PRCOHL (I,J)
AC (K)=TI*ACT (I,J,1)
NT (K)=IDNO (I)
2 CONTINUE
NQ=NO NUCL*13-1
DO 3 J=1,NQ
NTE MP1=NP (J+1)
NTE MP2=NT (J+1)
ZTE MP3=AC (J+1)
ZTE MP4=MA (J+1)
DO 4 K=1,J
I=J+1-K
IF (NTE MP1.GE.NP (I)) GO TO 5
NP (I+1)=NP (I)
NT (I+1)=NT (I)
AC (I+1)=AC (I)

```



```

      HA(I+1)=HA(I)
4    CONTINUE
      I=0
5    NP(I+1)=NTEMP1
      NT(I+1)=NTEMP2
      AC(I+1)=ZTEMP3
      HA(I+1)=ZTEMP4
3    CONTINUE
      L=1
7    IF(NP(L).NE.0) GO TO 205
      L=L+1
      GO TO 7
205  IBASE=L
CBEGIN INTERROGATION OF THE LEVDEC LIBRARY
200  READ(7,16) ITAPE,HALFLF,NDMODE
      16  FORMAT(16,20X,E11.4,I3)
      IF(IOP(7).NE.0) GO TO 200
201  IF(ITAPE.EQ.NP(L)) GO TO 17
      IF(ITAPE.GT.NP(L)) GO TO 28
      CALL ADVAN(NDMODE)
      GO TO 200
26  IF(AC(L).EQ.0.0) GO TO 38
      WRITE(6,29) NP(L)
29  FORMAT(/,33M ATTENTION-PRODUCT NUCLIDE NUMBER,16,90M DOES NOT EXIS
      +T IN DECAY LIBRARY-THEREFORE NO CONTRIBUTION FROM THIS NUCLIDE IN
      +DOSE CALCS)
      WRITE(6,37) NT(L)
37  FORMAT(LN,20M THE TARGET ATOM WAS ,IS,///)
38  L=L+1
      IF(NP(L).EQ.0) GO TO 813
      GO TO 280
C/CHECK TO SEE IF NP(L) EQUALS NP(L+1)
17  IF(NP(L).EQ.NP(L+1)) GO TO 800
CCOMPAR E MALFLIVES
313  IF(HA(L)/HALFLF.LT.D.0.0R.HA(L)/HALFLF.GT.1.2) GO TO 19
      IF(HALFLF.EQ.1.0E+50.0R.AC(L).EQ.0.0) GO TO 19
      IMODE=0
302  IMODE=IMODE+1
CBEGIN INTERROGATION OF THE LEVDEC LIBRARY
      REAC(7,300) IDEKMO,IDAUG,OLEVEL,MULT
300  FORMAT(30X,I2,I7,E11.4,E12.4)
      IF(IDEKMC.EQ.9) ND(L)=IDAUGH
      IF(IDEKMO.EQ.7.0R.IDEKMC.EQ.9) GO TO 301
      IF(IDEKMO.EQ.8) GO TO 310
303  IF(IMODE.EQ.NDMODE1) GO TO 304
      GO TO 302
301  IF(OLEV=L.EQ.0.0) GO TO 303
      CALL CURVE(OLEVEL,FLUXE)
      CALL CALC(OLEVEL,AC(L),DCSPRO(L),MULT,FLUXE,TMASS)
      GO TO 303
310  IF(OLEVEL.EQ.0.0) GO TO 311
      CALL CURVE(OLEVEL,FLUXE)
      CALL CALC(OLEVEL,AC(L),DCSPRO(L),MULT,FLUXE,TMASS)
CIN ANY CASE ADD IN THE 0.511 MEV ANNIHILATION GAMMAS
311  DD 312 IBETA=1,2
      CALL CURVE(0.511,FLUXE)
      CALL CALC(0.511,AC(L),DCSPRO(L),MULT,FLUXE,TMASS)
312  CONTINUE
      GO TO 303
      19  CALL ADVAN(NDMODE)
C/CREAD THE NEXT LEVDEC NUCLIDE
30~  READ(7,16) ITAPE,HALFLF,NDMODE

```

```

IF (EOF(7).NE.0) GO TO 304
IF (ITAPE.EQ.NP(L)) GO TO 313
IF (ITAPE.EQ.NP(L+1)) GO TO 21
IF (NP(L+1).EQ.0) GO TO 813
L=L+1
GO TO 266
21 L=L+1
GO TO 17
800 DO 801 IL=1,NDMODE
READ(7,36) IDEKM(IL),IDAUG(IL),DLEVL(IL),PRCB(IL)
801 CONTINUE
IBEGIN=L
802 L=L+1
IF (NP(L).EQ.NP(L+1)) GO TO 802
IEND=L
810 DO 805 IP=IBEGIN,IEND
IF (HA(IP)/HALFLF.GT.0.8.AND.HA(IP)/HALFLF.LT.1.2.AND.HALFLF.NE.1.0
+E*50.AND.AC(IP).NE.0.0) GO TO 834
GO TO 835
804 DO 820 IL=1,NDMODE
IF (IDEKM(IL).EQ.9) NO(IP)=IDAUG(IL)
IF (IDEKM(IL).EQ.7.OR.IDEKM(IL).EQ.9) GO TO 801
IF (IDEKM(IL).EQ.8) GO TO 510
GO TO 820
501 IF (DLEVL(IL).EQ.6.0) GO TO 820
CALL CURVE(DLEVL(IL),FLUXE)
CALL CALC(DLEVL(IL),AC(IP),DOSPRD(IP),PROB(IL),FLUXE,TMASS)
GO TO 820
510 IF (DLEVL(IL).EQ.0.0) GO TO 511
CALL CURVE(DLEVL(IL),FLUXE)
CALL CALC(DLEVL(IL),AC(IP),DOSPRD(IP),PROB(IL),FLUXE,TMASS)
511 DO 821 I=BETA-1,2
CALL CURVE(D,511,FLJXE)
CALL CALC(D,511,AC(IP),DCSPRO(IP),PROB(IL),FLUXE,TMASS)
821 CONTINUE
820 CONTINUE
805 CONTINUE
906 READ(7,16) ITAPE,HALFLF,NDMODE
IF (EOF(7).NE.0) GO TO 806
IF (ITAPE.EQ.NP(IBEGIN)) GO TO 808
GO TO 807
808 DO 809 IL=1,NDMODE
READ(7,300) IDEKM(IL),IDAUG(IL),DLEVL(IL),PROB(IL)
809 CONTINUE
GO TO 810
807 IF (ITAPE.EQ.NP(IEND+1)) GO TO 811
L=IEND+1
IF (ITAPE.GT.NP(IEND+1)) GO TO 28
GO TO 290
911 L=IEND+1
GO TO 17
812 WRITE(6,29) NP(L)
WRITE(6,37) NT(L)
L=L+1
IF (ITAPE.GT.NP(L)) GO TO 812
GO TO 200
813 REMIND ?
MOVE THE FILE MARKER THROJGH THE ACTLME LIBRARY AND POSITION AT THE
CBEGINNING OF LEVDEC
49 READ(7,47) ITAPE
47 FORMAT (6)
IF (EOF(7).NE.0) GO TO 700

```

```

      CALL POSITION
      GO TO 49
700  ISTOP=L
      L=IBASE
CBEGIN INTERROGATION OF THE LEVDEC LIBRARY FOR POSSIBLE RADIOACTIVE
CDAUGHTERS=REAC UNTIL THE ID NUMBER IS MATCHED
701  IF(ND(L).NE.0) GO TO 702
      L=L+1
      IF(L.EQ.ISTOP+1) GO TO 913
      GO TO 701
702  READ(7,16) ITAPE,HALFLF,NOMODE
      IF(IEDF(7).NE.0) GO TO 702
      IF(ITAPE.EQ.ND(L)) GO TO 703
      IF(ITAPE.GT.ND(L)) GO TO 926
      CALL ADVAN(NOMODE)
      GO TO 702
926  L=L+1
      IF(L.EQ.ISTOP+1) GO TO 913
      IF(ITAPE.EQ.ND(L)) GO TO 703
      CALL ADVAN(NOMODE)
      GO TO 701
C CHECK THE HALF LIFE OF THE DAUGHTER
703  DHA(L)=HALFLF
      IF(HALFLF.GT.1.0E+08) GO TO 704
      GO TO 705
704  CALL ADVAN(NOMODE)
904  L=L+1
      IF(L.EQ.ISTOP+1) GO TO 913
      GO TO 701
705  IMODE=0
902  IMODE=IMODE+1
C BEGIN INTERROGATION OF THE LEVDEC LIBRARY FOR DAUGHTER DECAY INFO
707  FORMAT(30X,I2.7X,E12.4,E12.4)
      SUMA1=0.0
      SUMA2=0.0
      DO 60 I=1,N
      SUMA1=SUMA1+EXP(-IN-I)*(T1+T2)*(0.693*3600./HA(L))
      SUMA2=SUMA2+EXP(-IN-I)*(T1+T2)*(0.693*3600./DHA(L))
60  CONTINUE
      ACC=AC(L)*SUMA2*(1-ID.693/DHA(L))*EXP(-0.693*3600.*T1/HA(L))-0.693/
      +HA(L)*EXP(-0.693*3600.*T1/DHA(L))/(0.693/DHA(L)-0.693/HA(L))/SUM
      +A1/(1.-EXP(-0.693*3600.*T1/HA(L)))
      IF(IDEKMC.EQ.7.OR.IDEKMC.EQ.9) GO TO 901
      IF(IDEKMC.EQ.8) GO TO 910
903  IF(IMODE.EQ.NOMODE) GO TO 904
      GO TO 902
901  IF(DLEVEL.EQ.0.0) GO TO 903
      CALL CURVE(DLEVEL,FLUXE)
      CALL CALC(DLEVEL,ACC,DCSCAU(L),MULT,FLUXE,YMASS)
      GO TO 903
910  IF(DLEVEL.EQ.0.0) GO TO 911
      CALL CURVE(DLEVEL,FLUXE)
      CALL CALC(DLEVEL,ACC,DCSCAU(L),MULT,FLUXE,YMASS)
C IN ANY CASE ADD IN THE 0.511 MEV ANNIHILATION GAMMAS
911  DO 912 IBETA=1,2
      CALL CURVE(0.511,FLUXE)
      CALL CALC(0.5110,ACC,DCSCAU(L),MULT,FLUXE,YMASS)
912  CONTINUE
      GO TO 903
913  RLEIND ?
      NO=NDNUCL*13

```

```

DD 815 I=1,NONUCL
WRITE (6, 9) (STAF5(K), K=1, NOPAS)
3  FORMAT (//, 1H, 4X, 7HTIME(H), 4X, 10 (E0.3, 1X), 2X, 2 (E0.3, 2X))
WRITE (6, 40B)
*00 FORMAT (1P, 125HTARGET PRODUCT DSR (T1) DSR (T2) DSR (T3) DSR (T4)
+ DSR (T5) DSR (T6) DSR (T7) DSR (T8) DSR (T9) DSR (T10) DSR (T11)
+ DSR (T12))
DD 818 L=IBASE,NO
IF (NT(L).EQ.IDNO(I)) GO TO 817
GO TO 818
817 DO 819 K=1,NOPAS
DCSPRO(K)=DOSPRD(L)*EXP(-(STAF5(K)*0.693*3600./HA(L)))
IF (CHA(L).EQ.0.0) DOSDGT(K)=0.0
IF (DMA(L).NE.0.0) DOSDGT(K)=DOSDAU(L)*EXP(-(STAF5(K)*0.693*3600./O
+HA(L)))
DOSNUC(K)=DOSPRO(K)+DOSDGT(K)
DOSSUM(K)=DOSSUM(K)+DOSNUC(K)
SYSDOS(K)=SYSDOS(K)+DOSNUC(K)
819 CONTINUE
WRITE (6, 20) IDNO(I), NP(L), (DCSPRO(K), K=1, NOPAS)
20  FORMAT (1X, 1X, 15, 2X, 15, 2X, 10 (E0.3, 1X), 2X, 2 (E0.3, 2X))
818 CONTINUE
WRITE (6, 11)
11  FORMAT (126HPRODUCT DAUGHTER DSR (T1) DSR (T2) DSR (T3) DSR (T4) DS
+R (T5) DSR (T6) DSR (T7) DSR (T8) DSR (T9) DSR (T10) DSR (T11) DSR
+ (T12))
DD 618 L=IBASE,NG
IF (NT(L).NE.IDNO(I)) GO TO 619
IF (ND(L).NE.C.0) GO TO 619
IF (L.NE.NO) GO TO 618
GO TO 618
619 WRITE (6, 20) NP(L), ND(L), (DOSDGT(K), K=1, NOPAS)
618 CONTINUE
WRITE (6, 8) (STAF5(K), K=1, NOPAS)
WRITE (6, 10)
10  FORMAT (1P, 125H TARGET DSR (T1) DSR (T2) DSR (T3) DSR (T4)
+ DSR (T5) DSR (T6) DSR (T7) DSR (T8) DSR (T9) DSR (T10) DSR (T11)
+ DSR (T12))
WRITE (6, 23) IDNO(I), (DOSSUM(K), K=1, NOPAS)
23  FORMAT (1X, 5X, 15, 5X, 10 (E0.3, 1X), 2X, 2 (E0.3, 2X), ///)
DD 822 K=1,NOPAS
DOSSUM(K)=0.0
922 CONTINUE
815 CONTINUE
WRITE (6, 414)
414  FORMAT (//)
WRITE (6, 8) (STAF5(K), K=1, NOPAS)
WRITE (6, 26)
26  FORMAT (1P, 125H SYSTEM DSR (T1) DSR (T2) DSR (T3) DSR (T4)
+ DSR (T5) DSR (T6) DSR (T7) DSR (T8) DSR (T9) DSR (T10) DSR (T11)
+ DSR (T12))
WRITE (6, 51) (SYSDOS(K), K=1, NOPAS)
51  FORMAT (1X, 15X, 10 (E0.3, 1X), 2X, 2 (E0.3, 2X))
RETURN
END

```

```

SUBROUTINE CALC(CIE,ACT,DSFRD,MULT,FLUXE,THASS)
C
C
CTHIS PROGRAM CALCULATES DOSERATE FOR 4 DIFFERENT GEOMETRIES
C
C
C
C
REAL MULT
REAL MU
COMMON A,V,F,DF,T1,T2,N,F,D,H,OPTION,IQEOM,ISAE
GO TO(11,12)(ISAB+1)
11 GO TO(31,32,33,34)IQEOM
31 DSFRAT=(ACT*MULT*2.5/(4.0*3.1416*FLUXE*(D*100.0)**2))
DSFRD=DSFRD+DOSFRAT
RETURN
CTHE ABOVE DOSERATE CALCULATION IS FOR A POINT SOURCE
32 VOLU=4*3.141E*(R*100.)**3)/3
SDU=(ACT*MULT*2.5)/(VOLU*FLUXE*4.)
DOSFRAT=SDU*(2*R*(R+D)*10000.-D*(2*R+D)*16000.*(ALOG((2*R+D)/D)))/(
+R*D)*100.)
DSFRD=DSFRD+DOSFRAT
RETURN
CTHE ABOVE DOSERATE CALCULATION IS FOR A HOMOGENEOUS NON-ABSORBING
CSPHERE
33 VOLU=3.1416*((R*100.)**2)*H*100.
SDU=(ACT*MULT*2.5)/(VOLU*FLUXE*4.)
DOSFRAT=SDU*(H+D)*100.*(ALOG(1+(R**2)/(H+D)**2))+2*E*ATAN(H+D)/R
+1/(H+D))-D*16000.*(ALOG(1+(R**2)/(D**2))+2*R*ATAN(C/F)/D)
DSFRD=DSFRD+DOSFRAT
RETURN
CTHE ABOVE CALCULATION IS FOR A HOMOGENEOUS NON-ABSORBING CYLINDER WHERE
CTHE OBSERVER FACES THE AXIS
34 VOLU=4.*3.1416*((R*100.)**3)-(R*100.)**3)/3.
SDU=(ACT*MULT*2.5)/(VOLU*FLUXE)
DOSFRAT=SDU*H*100.
DSFRD=DSFRD+DOSFRAT
RETURN
C ABOVE CALCULATION IS FOR A HOMOGENEOUS SPHERICAL CAVITY
C WITHOUT ABSORPTION WHERE THE OBSERVER IS AT THE CENTER
12 GO TO(41,42,42,42)IQEOM
41 WRITE(6,50)
50 FORMAT(10X,36HNO SELF-ABSORPTION FOR A POINT SOURCE)
GO TO 31
C THE ABOVE STATEMENTS RETURN TO NON-ABSORBING POINT SOURCE
42 ANJ=0.0488*E**(-0.4633)
C CALCULATION OF THE ATTENUATION FACTOR
IF(E.LE.3.)BFA=44.1-(10.2*E)
IF(E.GT.(3.0)).AND.(E.LE.(8.0))BFA=-0.9*E+16.2
IF(E.GT.(8.0))BFA=AMAX1(4.0,(-2.5*E+29.))
C CALCULATION OF THE BUILD-UP FACTOR COEFFICIENT A
IF(E.LE.2.)BALPH1=0.063*E-0.182
IF(E.GT.(2.0)).AND.(E.LE.(8.0))BALPH1=-0.02*E-C.015
IF(E.GT.(8.0))BALPH1=-0.0085*E-0.107
C CALCULATE THE BUILD-UP FACTOR COEFFICIENT ALPHA1
IF(E.LE.(1.0))BALPH2=-0.032*E+0.048
IF(E.GT.(1.0)).AND.(E.LE.(2.0))BALPH2=-0.001*E+0.017
IF(E.GT.(2.0)).AND.(E.LE.(4.0))BALPH2=0.012*E-0.009

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      IF (E.GT.(6.0)).AND.(E.LE.(8.0))$BALPH2=0.00125*E+0.034
      IF (E.GT.(8.0))$BALPH2=0.012*E-0.052
C   CALCULATE THE BUILD-UP FACTOR COEFFICIENT ALPHA2
      GO TO (4,51,52,53)$IGEOM
51  VOLU=4.*3.1416*((R*100.)**3)/3
      MU=1000.*AMU*TMASS/VOLU
      SOJ=(ACT*MULT*2.5)/(VOLU*FLUXE)
      CALL GEOPSP(R,D,MU,BALPH1,GFS)
      DOSRAT=SOU**100.*GFS*BFA/3.1416
      CALL GEOPSP(P,D,MU,BALPH2,GFS)
      DOSRAT=DOSRAT+SOL*R*100.*(1.-BFA)*GFS/3.1416
      DSPRD=DSFRD+DOSRAT
      RETURN
C   CALCULATE THE DOSE RATE AT A PCINT OUT OF AN ABSORBING SPHERE
52  VOLU=3.1416*((R*100.)**2)*H*100.
      MU=1000.*AMU*TMASS/VOLU
      SDU=(ACT*MULT*2.5)/(VOLU*FLUXE)
      TETA=ATAN(R/(D+H/3.))
      IF (H*100.).GT.(13./MU))GO TO 60
      CALL GEOPCY(BALPH1,MU,TETA,H,GFC)
      DOSRAT=SDU*(1.-COS(TETA)+GFC)*3FA/(2.*MU*(1.+BALPH1))
      CALL GEOPCY(BALPH2,MU,TETA,H,GFC)
      DOSRAT=SDU*(1.-COS(TETA)+GFC)*(1.-BFA)/(2.*MU*(1.+BALPH2))+DOSPAT
      DSPRD=DSFRD+DOSRAT
      RETURN
50  DOSRAT=SOU*(1.-COS(TETA))*(BFA/((1.+BALPH1)*MU)+(1.-BFA)/((1.+BALP
      +H2)*MU))/2.
      DSPRD=DSFRD+DOSRAT
      RETURN
C   CALCULATE THE DOSE RATE FOR A HOMOGENECUS ABSORBING CYLINDER
C   WHERE THE OBSERVER FACES THE AXIS
53  VOLU=4.*3.1416*((((?+H)*100.)**3)-(R*100.)**3)/3.
      MU=1000.*AMU*TMASS/VOLU
      SOJ=(ACT*MULT*2.5)/(VOLU*FLUXE)
      DOSRAT=(SOU*BFA*(1.-EXP(-MU*(1.+BALPH1)*H*100.))/MU/(1.+BALPH1))+
      +(SOJ*(1.-BFA*(1.-EXP(-MU*(1.+BALPH2)*H*100.)))/(MU*(1.+BALPH2)))
      DSPRD=DSFRD+DOSPAT
      RETURN
C   CALCULATE THE DOSE RATE AT THE CENTER OF A SPHERICAL
C   ABSORBING CAVITY
      END

```

SUBROUTINE GE0FSF(X,Y,Z,T,GFS)

C  
C  
C  
C  
C  
C  
C

GIVEN THE RADIUS IN X, THE DISTANCE FROM THE SURFACE IN Y, THE  
ATTENUATION FACTOR IN Z AND THE BUILD-UP FACTOR COEFFICIENTS IN T  
CALCULATES THE GEOMETRIC FUNCTION FOR ABSORBING SPHERE USING AN  
ANALYTICAL APPROXIMATION OF THE G FUNCTION

GFS=0.  
PSP=(X+Y)/X  
GFS=EXP(-((0.342\*X\*Z\*(1.+T)+2.07322)\*ALOG(PSP)+AMIN1((60.\*X\*Z\*(1.+  
T)-0.144),(46.37\*X\*Z\*(1.+T)-0.035),(28.78\*X\*Z\*(1.+T)+0.3168),(14.5  
+5\*X\*Z\*(1.+T)+0.8861),(8.17\*X\*Z\*(1.+T)+1.3966),(4.72\*X\*Z\*(1.+T)+1.7  
+42))))  
RETURN  
END

SUBROUTINE GE0FCY(X,Y,T,Z,GFC)

C  
C  
C  
C  
C  
C  
C

GIVEN THE BUILD-UP FACTOR COEFFICIENTS IN X, THE ATTENUATION FACTOR  
IN Y, THE HEIGHT IN Z, THE ANGLE IN T, CALCULATES THE GEOMETRIC  
FUNCTION FOR ABSORBING CYLINDER USING AN APPROXIMATION OF THE E2  
FUNCTION

REAL VAR,VARI  
VAR=0.  
VARI=0.  
GFC=0.  
VAR=(1+X)\*Y\*Z\*100.  
VARI=VAR/COS(T)  
CALL E2(VARI,RES)  
GFC=COS(T)\*RES  
CALL E2(VAR,FES)  
GFC=GFC-RES  
RETURN  
END



SUBROUTINE E2(X,FES)

C  
C DESTINATE E2(X) THE SECOND SPECIES EXPONENTIAL INTEGRAL FUNCTION  
C

RES=0.  
IF(X.GE.0.)GO TO 33  
TOL=0.0001  
TOTAL=0.0  
A=X

B=5.\*X  
DO 10 I=1,10  
CALL RCMB(A,B,TOL,RESULT)  
TOTAL=TOTAL+RESULT

A=3  
B=B\*5.  
10 CONTINUE  
RES=X\*TOTAL

33 RETURN  
RES=EXP(-X)\*(1/(2.\*X)+2./((2.\*X)\*\*3))  
RETURN  
END

SUBROUTINE ROME(A,B,TOL,RESULT)

C  
C  
C  
C

```
DIMENSION TRAP(11,11)
H=(B-A)/10.
SUM=(EXP(-A))/(A**2)+(EXP(-B))/(B**2)
X=A
DO 10 I=2,10
  X=X+H
  SUM=SUM+((EXP(-X))/(X**2))*2.
10 CONTINUE
TRAP(1,1)=H/2.*SUM
DO 20 I=1,10
  H=H/2
  X=A+H
  K=10*2**I
  DO 30 J=2,K,2
    SUM=SUM+((EXP(-X))/(X**2))*2
  X=K+H+H
30 CONTINUE
  TRAP(I,I+1)=H/2.*SUM
  DO 40 L=1,I
    TRAP(L,I+1)=TRAP(L,I+1)+1./(-.**L-1.)*(TRAP(L,I+1)-TRAP(L,I))
40 CONTINUE
20 CONTINUE
IF (ABS(TRAP(I+1,I+1)-TRAP(I,I+1))-TOL) > .50, .50, 20
50 RESULT=TRAP(I,I)
RETURN
END
```

## SUBROUTINE CURVE (E, FLUXE)

C  
C  
C GIVEN THE ENERGY OF THE GAMMA RAY, THIS PROGRAM CALCULATES THE PARTICLE  
C FLUX (QUANTA/CM\*\*2\*SEC) THAT CORRESPONDS TO A DOSE RATE OF 2.5 MR/HR --  
C FLUXE

C  
C  
C ARGUMENTS IN-DLE VEL  
C ARGUMENTS OUT-FLUXE  
C

```

IF (E.GE.0.010.AND.E.LE.0.020) GO TO 1
IF (E.GT.0.020.AND.E.LE.0.040) GO TO 2
IF (E.GT.0.040.AND.E.LE.0.050) GO TO 3
IF (E.GT.0.050.AND.E.LE.0.060) GO TO 4
IF (E.GT.0.060.AND.E.LE.0.080) GO TO 5
IF (E.GT.0.080.AND.E.LE.0.100) GO TO 6
IF (E.GT.0.100.AND.E.LE.0.150) GO TO 7
IF (E.GT.0.150.AND.E.LE.0.200) GO TO 8
IF (E.GT.0.200.AND.E.LE.0.300) GO TO 9
IF (E.GT.0.300.AND.E.LE.0.400) GO TO 10
IF (E.GT.0.400.AND.E.LE.0.500) GO TO 11
IF (E.GT.0.500.AND.E.LE.0.700) GO TO 12
IF (E.GT.0.700.AND.E.LE.1.000) GO TO 13
FLUXE=1.0E+99
RETURN
1 FLUXE=21.495330.64*E**2.1759447
RETURN
2 FLUXE=11.460233.22*E**2.0136022
RETURN
3 FLUXE=1036918.975*E**1.2698653
RETURN
4 FLUXE=96497.38032*E**0.4772395
RETURN
5 FLUXE=11225.12221*E**(-0.2074410)
RETURN
6 FLUXE=55482.78326*E**E*(-10.9132069)
RETURN
7 FLUXE=1330.569205*E**(-1.1036005)
RETURN
8 FLUXE=2.462.298566*E**(-1.0604513)
RETURN
9 FLUXE=1546.957795*E**(-0.0091050)
RETURN
10 FLUXE=1549.091320*E**(-0.7797676)
RETURN
11 FLUXE=1508.205757*E**(-0.7196794)
RETURN
12 FLUXE=1488.362190*E**(-0.7051622)
RETURN
13 FLUXE=1656.052874*E**(-0.7649757)
RETURN
END

```

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