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Surface state reconstruction in ion-damaged SmB₆N. Wakeham,¹ Y. Q. Wang,¹ Z. Fisk,² F. Ronning,¹ and J. D. Thompson¹¹*Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA*²*Department of Physics and Astronomy, University of California, Irvine, California 92697-4575, USA*

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We have used ion irradiation to damage the (001) surfaces of SmB₆ single crystals to varying depths, and have measured the resistivity as a function of temperature for each depth of damage. We observe a reduction in the residual resistivity with increasing depth of damage. Our data are consistent with a model in which the surface state is not destroyed by the ion irradiation, but instead the damaged layer is poorly conducting and the initial surface state is reconstructed below the damage. This behavior is consistent with a surface state that is topologically protected.

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I. INTRODUCTION

Topological insulators (TIs) have a set of bulk wave functions that are topologically distinct from the vacuum. One consequence of this is that at the interface between the topologically nontrivial insulator and the trivial vacuum the energy gap must close [1]. This gives rise to a gapless surface state that in a time-reversal invariant system is said to be topologically protected because it is robust against time-reversal invariant perturbations. TIs have been intensively studied because of their ability to further our understanding of the importance of the topology of wave functions, but also because properties of the surface state, for example, spin-momentum locking, have practical applications in areas such as spintronics [2].

SmB₆ has long been established as a Kondo insulator, in which the hybridization of the *f* electrons with the conduction electrons opens a small energy gap [3]. This energy gap causes the resistivity to diverge with decreasing temperature down to 5 K. However, below 5 K, the resistivity saturates. Recently, SmB₆ was predicted to be topologically nontrivial, and therefore a three-dimensional topological Kondo insulator with a topologically protected surface state [4]. Experiments by several groups have established strong evidence that the resistivity saturation is indeed due to a conductive and intrinsic surface state [5–7]. Furthermore, transport measurements have shown weak antilocalization [8] and the sensitivity of the surface state to time-reversal breaking perturbations [9] that are consistent with the surface state being topological. Angle resolved photoemission spectroscopy (ARPES) experiments have also shown evidence for a topological surface state through observation of in-gap surface states and spin-momentum locking [10,11]. However, the topological nature of the surface state still needs to be conclusively proven. There is evidence from scanning tunneling spectroscopy measurements for surface reconstructions [12], and x-ray measurements have indicated the physical surface of SmB₆ crystals has a different average valence of Sm ions compared with the bulk [13,14]. These factors may give rise to a nontopological conducting surface state, perhaps as well as a topological one. Hence, the properties of this surface state are still an area of interest.

Here, we report on the nature of the intrinsic surface state in SmB₆ through nonmagnetic ion irradiation of the surface

of SmB₆ single crystals. We have measured the temperature dependent resistivity of two samples as a function of the depth to which the surface is damaged, and find that the intrinsic surface state is not destroyed by this time-reversal invariant damage. Instead, the damaged area becomes more metallic than bulk SmB₆, and the intrinsic surface state is reconstructed below this damaged layer. Such insensitivity to this time-reversal invariant perturbation is consistent with the surface state in SmB₆ being topologically protected.

II. EXPERIMENTAL TECHNIQUES

Single crystals of SmB₆ were grown using aluminum flux and then polished to give platelike crystals in the (001) direction with approximate dimensions of the length, width, and thickness of 300, 180, and 35 μm , respectively. The resistance *R* of the samples as a function of temperature *T* was measured using a four-probe low frequency ac resistance bridge with spot-welded contacts of 25 μm platinum wires attached to the top (001) face. These contacts remained in place during all of the measurements and ion irradiation. Temperature was controlled within a Quantum Design physical property measurement system (PPMS). Note that the platelike geometry of the samples means that the sides of the samples are not a significant conducting channel compared with the larger faces, and are therefore ignored in our discussions. The (001) surfaces of the samples were damaged by ion irradiation. The ion used, acceleration energy, ion fluence, and exposure time to produce a peak damage of at least one displacement per atom (DPA) are shown in Table I. The level of damage as a function of depth for each exposure was calculated using the SRIM Monte Carlo code in the full cascade mode, and is shown in the inset to Fig. 1 [15]. The damage depth is defined as the depth at which there is half of the maximum damage.

III. RESULTS

The temperature dependent resistance of the samples was measured after being damaged to the stated depth on the top face, and then again after being damaged on the bottom face. *R*(*T*) as the depth of damage was varied on each face is shown in Fig. 1.

This figure clearly shows a marked reduction in the low temperature residual resistance *R*₀ as the damage depth was

TABLE I. Parameters used in the ion irradiation of SmB_6 to produce a damage profile of approximately one displacement per atom down to the stated depth.

Depth (nm)	Ion	Energy (keV)	Ion fluence (cm^{-2})	Time (s)
12.9	Xe^+	20	5.0×10^{14}	800
108	Xe^{2+}	350	3.0×10^{14}	1034
237	Ar^{2+}	300	1.5×10^{15}	357

increased. A fit to the data (not shown) based on an activation-like resistance at high temperature and temperature independent surface resistance at low temperature, discussed by Kim *et al.* [7], indicates no significant change to the bulk energy gap of the system after the surfaces were damaged. Therefore,

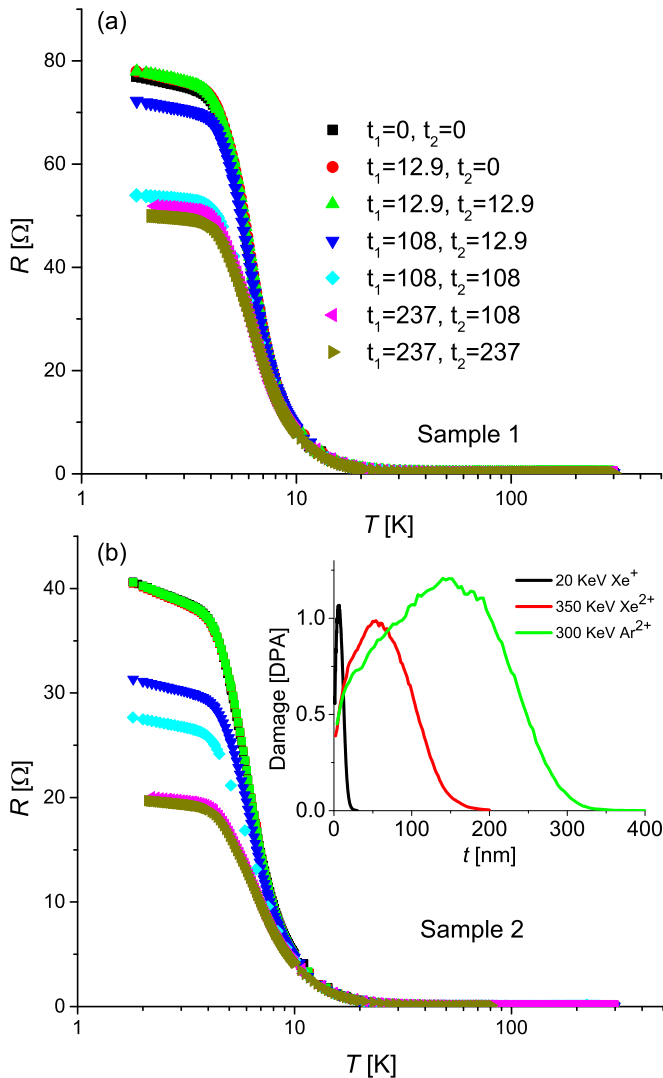


FIG. 1. (Color online) Temperature dependence of the resistance of SmB_6 after differing depths of ion-irradiation damage to each face. t_1 is the depth of damage to the top face of the sample (to which contacts were made), and t_2 is the depth of damage to the bottom face, both in units of nm. Shown for (a) sample 1 and (b) sample 2. The inset shows the calculated damage as a function of depth for each ion beam exposure described in Table I.

the reduction in R_0 suggests that either the damage increased the conductivity of the intrinsic surface state, or it introduced a new conducting channel. Karkin *et al.* irradiated SmB_6 with neutrons to produce damage of 10 DPA throughout the sample. This damage reduced the residual resistivity by three orders of magnitude, and gave a largely temperature independent resistivity [16]. Even modest damage to the bulk of SmB_6 by neutrons reduced the low temperature resistivity [17]. The shapes of the defects induced by low energy ion irradiation and neutrons are expected to be qualitatively similar [18]. We conclude that in our measurements one effect of the ion irradiation was to create a thin damaged layer of relatively low resistivity that reduced the overall resistance of the sample at low temperatures. In addition to this, we must also determine whether the intrinsic surface state present in the undamaged sample was destroyed by the damage or was reconstructed below it.

IV. DISCUSSION

To interpret the change in low temperature resistance, it is useful to convert the residual resistance to a sheet resistance R_{S0} , given by $R_{S0} = R_0 2w/l$, where w is the width of the sample and l is the distance between voltage contacts. Figure 2(a) shows R_{S0} as a function of the total depth of surface damage d , where d is the sum of the depth of damage of the top and bottom faces.

The inset of Fig. 2(a) shows R_{S0} as a function of $1/d$ and demonstrates there is not a $1/d$ dependence. The residual sheet conductance of the sample is given by $G_{S0} = 1/R_{S0}$, and is plotted against d in Fig. 2(b). The solid curves in Figs. 2(a) and 2(b) show that the d dependence can be described by a model of reconstruction of the surface state represented by two parallel resistors.

A schematic diagram of the parallel resistor scenario is shown in Fig. 3. The first resistor represents the intrinsic surface state, and is assumed to have a constant sheet resistance R_{SS} , as a function of d . This constant is given by R_{S0} in the undamaged sample. The second resistor represents the damaged layer, and therefore the sheet resistance is given by $2\rho_{DL}/d$, where ρ_{DL} is the resistivity of the damaged layer, which is a fitted parameter. The bulk is assumed to be perfectly insulating. The total sheet resistance R_{S0} is therefore given by

$$\frac{1}{R_{S0}} = G_{S0} = \frac{1}{R_{SS}} + \frac{d}{2\rho_{DL}}. \quad (1)$$

The fitted parameter ρ_{DL} is $1.0 \pm 0.2 \text{ m}\Omega \text{ cm}$ and $2.4 \pm 0.4 \text{ m}\Omega \text{ cm}$ in sample 1 and 2, respectively. This is in reasonable agreement with the resistivity of $1.5 \text{ m}\Omega \text{ cm}$ found in the heavily neutron-damaged sample [16], as is expected since in both cases the irradiated region is sufficiently damaged to be considered approximately amorphous.

The origin of the reduction in magnitude of dR/dT with increasing depths of damage at temperatures below 4 K shown in Fig. 1 cannot be ascertained from our data. It may arise from a temperature dependent resistivity of the damaged layer, or a change in the temperature dependence of the intrinsic surface state. The data shown in Fig. 2 demonstrate deviations from the fitted line that are seemingly beyond the expected statistical error associated with the uncertainty in the depth of

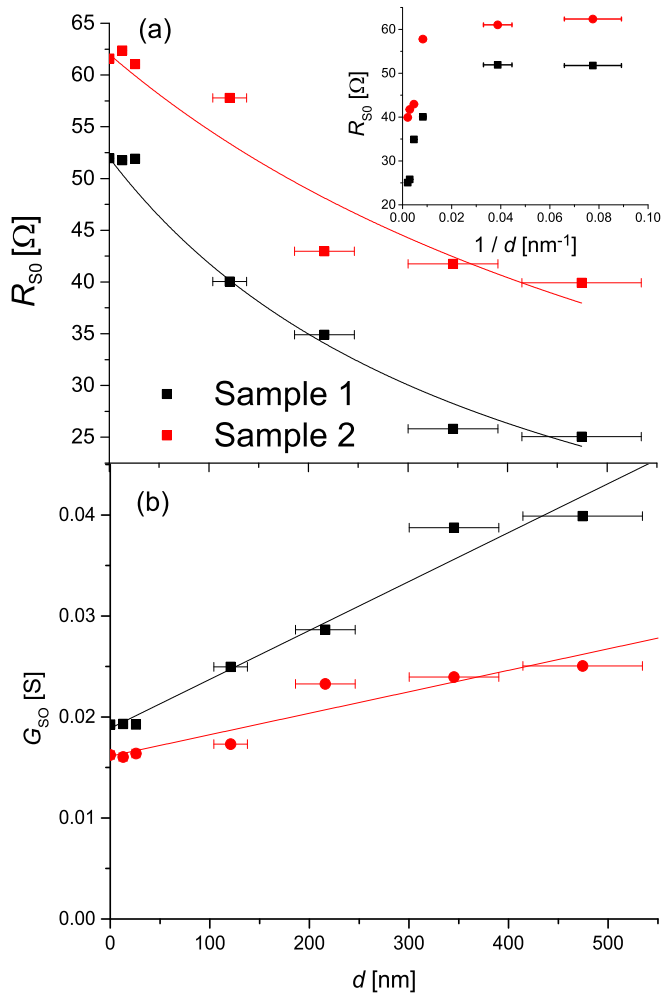


FIG. 2. (Color online) (a) Residual sheet resistance R_{S0} and (b) residual sheet conductance G_{S0} of two samples of SmB_6 as a function of the effective depth of ion-radiation damage d , given by the sum of the damage depth to the top and bottom faces. Solid lines are fits to the formula $1/R_{S0} = G_{S0} = 1/R_{SS} + d/2\rho_{DL}$, where ρ_{DL} is the resistivity of the damaged layer. The inset shows R_{S0} as a function of $1/d$ to emphasize the lack of proportionality between the two.

damage. This may arise from inhomogeneity of the damage to the sample, particularly as a function of d . In addition, there may also be changes to the conductivity of the intrinsic

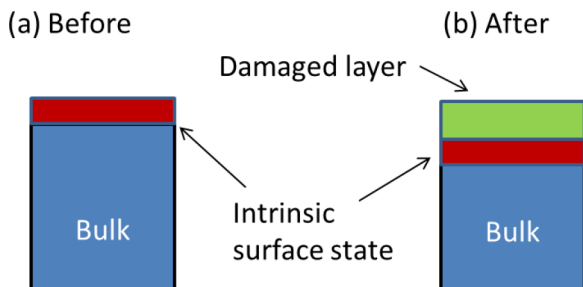


FIG. 3. (Color online) Schematic diagram of the top surface of SmB_6 (a) before and (b) after ion-irradiation damage. This shows the model of reconstruction of the intrinsic surface state (red) underneath a conductive damaged layer (green).

surface state not considered in our model. Note that this would be particularly significant for small d .

Let us now consider scenarios for the effect of the damage that are alternatives to the surface state reconstruction discussed above. First, note that the inset of Fig. 2 shows that the sheet resistance does not follow a $1/d$ dependence that would be expected if the intrinsic surface state depth were very shallow (<10 nm), and this were destroyed by the ion irradiation and replaced with the conducting damaged layer. We therefore exclude this possibility.

The depth of the surface state ξ may be crudely estimated from measurements of the Fermi velocity v_F of the surface state using $\xi = \hbar v_F / \Delta$, where Δ is the magnitude of the Kondo energy gap, taken to be $38k_B$ J, and k_B is the Boltzmann constant [7]. Using this equation, the value of $v_F \approx 6$ eV \AA obtained from de Haas–van Alphen (dHvA) measurements implies a metallic surface depth of ~ 150 nm [19]. Therefore, we must consider another scenario, that the intrinsic surface state has a depth of around 100 nm, and is destroyed by the ion damage. In this picture, ion-irradiation damage to 12.9 nm may have little effect on the sample resistance. However, once damaged to 108 nm, the surface state could be destroyed, and then at larger d the resistance as a function of $1/d$ would follow a linear dependence. The calculated resistivity of the damaged layer obtained by fitting R_{S0} for $d > 108$ nm with a $1/d$ dependence is of order $100 \mu\Omega$ cm. This resistivity value is more than an order of magnitude smaller than that measured for the neutron-damaged SmB_6 [16]. In addition, the room temperature bulk resistivity of the measured crystals was $\sim 150 \mu\Omega$ cm. It therefore seems unlikely that the heavily damaged material could support such a low resistivity. Work by others also argues against the surface state having a thickness of ~ 100 nm. For instance, the values of $v_F \approx 300$ meV \AA obtained from ARPES measurements are much smaller than obtained through dHvA measurements [10,20]. This smaller value implies the penetration of the surface state is $\xi \approx 10$ nm. We also note that measurements of 100 nm thick films of SmB_6 show a temperature dependence similar to that of bulk crystals, and is argued to reflect the development of a Kondo-derived gap, and a saturated resistivity associated with a conductive surface state [21]. Because some fraction of the film's volume hosts the Kondo gap, the penetration depth of the surface state must be well below 50 nm. Theoretical estimates of the depth of the topological surface state are 1–25 lattice spacings or 0.4–10 nm [22–24]. Based on all of these factors, we conclude that a surface state of 100 nm and a highly conducting damaged layer are unlikely to be the explanation of our results.

A final scenario to consider in the explanation of the residual resistance as a function of damage depth is that the undamaged sample had a topologically trivial insulating surface layer greater than 200 nm deep, with the intrinsic surface state below that. The layer would most likely be the oxide that forms on SmB_6 crystals in air, and is known to be an insulator [25]. The presence of such a deep layer would mean that the irradiation damage never went deep enough to affect the intrinsic surface state, it merely converted an insulating surface layer into a poor metallic layer, and therefore produced a reduced residual resistance. However, such an oxide layer would be expected to form to a depth of order 10 unit cells or approximately

4 nm, and indeed has recently been reported to be 2 nm deep in SmB₆ [26]. This is much less than the 200 nm required for this scenario. It is also unlikely that the highly insulating oxide layer would be made conducting by ion irradiation. Hence, we rule out the oxide layer as the explanation of our results.

V. CONCLUSION

Given all the considerations discussed above, we conclude that ion irradiation of the surface of SmB₆ produced a damaged layer with a resistivity of around 1.5 mΩ cm, but that the intrinsic conductive surface state present in the undamaged sample was reconstructed below the damaged layer. Theoretical predictions for the interface between a topological insulator and a conducting amorphous material are not available. However, it was shown theoretically that the interface between a crystalline metal and TIs will retain the Dirac spectrum expected at the interface to a vacuum, if the orbital overlap between the metal and the TI is small relative to the hopping parameter of the metal [27]. In our case, the large disorder in the metallic damaged layer means the resistivity is large and close to the Anderson localization limit, and therefore there is likely to be little orbital overlap with the TI. One further consideration is that ion irradiation causes a distribution of disorder as a function of depth into the sample. Therefore, the surface state is likely to have formed within a partially disordered region of the sample. Roy *et al.*

have calculated that this may shift the Dirac cone of the surface state relative to the valence band, but will not destroy the surface state [22]. Hence, we conclude that our deduction of a reconstructed surface state is consistent with the theoretical predictions for a topological surface state in SmB₆.

Our results may be relevant to the study of heterostructures of SmB₆ combined with insulators, superconductors, and normal metals that have been proposed [28–30]. We show that although the surface state will likely be present in such structures, the surface state may reside deeper within the SmB₆ layer than anticipated. In the future it will be interesting to compare our results with the effect of implanting magnetic ions into the surface of SmB₆. One might expect this to break the time-reversal symmetry at the surface and therefore destroy the intrinsic surface state, as was shown by the introduction of Gd ions into the bulk of SmB₆ [9], and the deposition of iron onto the surface of Bi₂Se₃ [31].

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- [1] Y. Ando, *J. Phys. Soc. Jpn.* **82**, 102001 (2013).
 - [2] T. Yokoyama and S. Murakami, *Physica E* **55**, 1 (2014).
 - [3] T. Kasuya, *Europhys. Lett.* **26**, 277 (1994).
 - [4] M. Dzero, K. Sun, V. Galitski, and P. Coleman, *Phys. Rev. Lett.* **104**, 106408 (2010).
 - [5] P. Syers, D. Kim, M. S. Fuhrer, and J. Paglione, [arXiv:1408.3402](https://arxiv.org/abs/1408.3402).
 - [6] S. Wolgast, C. Kurdak, K. Sun, J. W. Allen, D.-J. Kim, and Z. Fisk, *Phys. Rev. B* **88**, 180405 (2013).
 - [7] D. J. Kim, S. Thomas, T. Grant, J. Botimer, Z. Fisk, and J. Xia, *Sci. Rep.* **3**, 3150 (2013).
 - [8] S. Thomas, D. J. Kim, S. B. Chung, T. Grant, Z. Fisk, and J. Xia, [arXiv:1307.4133](https://arxiv.org/abs/1307.4133).
 - [9] D. J. Kim, J. Xia, and Z. Fisk, *Nat. Mater.* **13**, 466 (2014).
 - [10] M. Neupane, N. Alidoust, S.-Y. Xu, T. Kondo, Y. Ishida, D. J. Kim, C. Liu, I. Belopolski, Y. J. Jo, T.-R. Chang, H.-T. Jeng, T. Durakiewicz, L. Balicas, H. Lin, A. Bansil, S. Shin, Z. Fisk, and M. Z. Hasan, *Nat. Commun.* **4**, 2991 (2013).
 - [11] N. Xu, P. K. Biswas, J. H. Dil, R. S. Dhaka, G. Landolt, S. Muff, C. E. Matt, X. Shi, N. C. Plumb, M. Radović, E. Pomjakushina, K. Conder, A. Amato, S. V. Borisenko, R. Yu, H.-M. Weng, Z. Fang, X. Dai, J. Mesot, H. Ding, and M. Shi, *Nat. Commun.* **5**, 4566 (2014).
 - [12] S. Röbner, T.-H. Jang, D.-J. Kim, L. H. Tjeng, Z. Fisk, F. Steglich, and S. Wirth, *Proc. Natl. Acad. Sci. USA* **111**, 4798 (2014).
 - [13] M. Aono, R. Nishitani, C. Oshima, T. Tanaka, E. Bannai, and S. Kawai, *Surf. Sci.* **86**, 631 (1979).
 - [14] W. A. Phelan, S. M. Koohpayeh, P. Cottingham, J. W. Freeland, J. C. Leiner, C. L. Broholm, and T. M. McQueen, *Phys. Rev. X* **4**, 031012 (2014).
 - [15] J. F. Ziegler, M. D. Ziegler, and J. P. Biersack, *Nucl. Instrum. Methods Phys. Res., Sect. B* **268**, 1818 (2010).
 - [16] A. Karkin, Y. Akshentsev, and B. Goshchitskii, *Physica C* **460-462**, 811 (2007).
 - [17] J. Morillo, C.-H. de Novion, and J. Jun, *Solid State Commun.* **48**, 315 (1983).
 - [18] G. Was, Z. Jiao, E. Getto, K. Sun, A. M. Monterrosa, S. Maloy, O. Anderoglu, B. Sencer, and M. Hackett, *Scr. Mater.* **88**, 33 (2014).
 - [19] G. Li, Z. Xiang, F. Yu, T. Asaba, B. Lawson, P. Cai, C. Tinsman, A. Berkley, S. Wolgast, Y. S. Eo, D.-j. Kim, C. Kurdak, J. W. Allen, K. Sun, X. H. Chen, Y. Y. Wang, Z. Fisk, and L. Li, *Science* **346**, 1208 (2014).
 - [20] J. Jiang, S. Li, T. Zhang, Z. Sun, F. Chen, Z. R. Ye, M. Xu, Q. Q. Ge, S. Y. Tan, X. H. Niu, M. Xia, B. P. Xie, Y. F. Li, X. H. Chen, H. H. Wen, and D. L. Feng, *Nat. Commun.* **4**, 3010 (2013).
 - [21] J. Yong, Y. Jiang, D. Usanmaz, S. Curtarolo, X. Zhang, L. Li, X. Pan, J. Shin, I. Takeuchi, and R. L. Greene, *Appl. Phys. Lett.* **105**, 222403 (2014).
 - [22] B. Roy, J. D. Sau, M. Dzero, and V. Galitski, *Phys. Rev. B* **90**, 155314 (2014).
 - [23] V. Alexandrov, M. Dzero, and P. Coleman, *Phys. Rev. Lett.* **111**, 226403 (2013).
 - [24] V. Alexandrov, P. Coleman, and O. Erten, [arXiv:1501.03031](https://arxiv.org/abs/1501.03031)
 - [25] T. Tanaka, R. Nishitani, C. Oshima, E. Bannai, and S. Kawai, *J. Appl. Phys.* **51**, 3877 (1980).
 - [26] Y. S. Eo, S. Wolgast, T. Ozturk, G. Li, Z. Xiang, C. Tinsman, T. Asaba, F. Yu, B. Lawson, J. W. Allen, K. Sun, L. Li, C. Kurdak, D.-J. Kim, and Z. Fisk, [arXiv:1410.7430](https://arxiv.org/abs/1410.7430).

- [27] E. Zhao, C. Zhang, and M. Lababidi, *Phys. Rev. B* **82**, 205331 (2010).
- [28] C.-Y. Hou, K. Shtengel, G. Refael, and P. M. Goldbart, *New J. Phys.* **14**, 105005 (2012).
- [29] I. Affleck and D. Giuliano, *J. Stat. Mech.: Theory Exp.* (2013) P06011.
- [30] J. R. Williams, A. J. Bestwick, P. Gallagher, S. S. Hong, Y. Cui, A. S. Bleich, J. G. Analytis, I. R. Fisher, and D. Goldhaber-Gordon, *Phys. Rev. Lett.* **109**, 056803 (2012).
- [31] L. A. Wray, S.-Y. Xu, Y. Xia, D. Hsieh, A. V. Fedorov, Y. S. Hor, R. J. Cava, A. Bansil, H. Lin, and M. Z. Hasan, *Nat. Phys.* **7**, 32 (2010).