

5.14 The Stable Isotopic Composition of Atmospheric O₂

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5.14.1 Introduction

During the 1930s, Harold Urey and his collaborators demonstrated that ratios between heavy to light isotopes of the same element varied among various compounds that occur in nature and in the laboratory. They were successful in developing a theory that explained many of the observed variations from differences in the vibrational frequency of molecules (Urey, 1947) and thus obtained convincing evidence that isotopes of the same element differed not only in atomic mass but also in chemical properties. Using this theory, Urey and Greiff (1935) calculated equilibrium constants for exchange reactions involving light elements, such as carbon, oxygen, and sulfur. From the evaluations of the effect of the isotopic exchange, they could predict appreciable variations of atomic weights. Indeed, shortly after the publication of Urey and Greiff's paper, two independent laboratories discovered that the atomic weight of atmospheric oxygen was considerably larger than the atomic weight of oxygen in water (Dole, 1935; Morita, 1935).

For the 30 years that followed this discovery, Dole and his collaborators, almost alone, worked on the global cycle of oxygen and its isotopes. In a review in 1965, Dole concluded that the primary cause for the difference in atomic weight of oxygen between air and water is isotopic fractionation (to be defined in the following section) taking place during respiration by various organisms. The difference in the isotopic composition between atmospheric O₂ and ocean water is named the 'Dole effect' (DE) after Dole who discovered it and devoted many years of his career to studying much about its global controlling

mechanisms. Yet, much remained to be discovered and understood about this effect after Dole had left the scene, and an updated review on the DE is given in the succeeding section.

Oxygen has three naturally occurring stable isotopes of atomic mass numbers 16, 17, and 18. The difference between the atomic weights of oxygen in water and air (DE) is caused mainly by differences in the ratio ¹⁸O/¹⁶O, and most studies on stable oxygen isotopes deal with variations in the ¹⁸O/¹⁶O ratio. This is so because until 1973 it was thought that all natural processes causing variations in stable isotopic ratios are mass-dependent. In this case, relative variations in the ¹⁷O/¹⁶O ratio are expected to be only about one-half those in the ¹⁸O/¹⁶O ratio. The discovery that in meteorites, variations in ¹⁷O/¹⁶O are about the same as in ¹⁸O/¹⁶O rather than one-half (Clayton et al., 1973) changed that view and showed that mass-independent processes occur in nature. After this discovery, Thiemens and his collaborators started in 1983 a successful line of research and demonstrated that mass-independent fractionations are quite common and important both in the laboratory and in field observations (Thiemens, 1999). Motivated by these studies, Luz et al. (1999) studied atmospheric O₂ and showed that it has a deficit of ¹⁷O due to photochemical reactions among O₂, ozone, and CO₂ in the stratosphere. Among other factors, the extent of this deficiency depends on the rate of photosynthesis and is thus useful in studies involving marine and global production of O₂ by plants (e.g., Blunier et al., 2002; Luz and Barkan, 2000, 2009). Variations in both the ¹⁸O/¹⁶O and ¹⁷O/¹⁶O ratio in molecular O₂ and their significance in geochemistry are reviewed in the succeeding dedicated section.

5.14.2 Methodology and Terminology

5.14.2.1 Abundance and Fractionation of Oxygen Isotopes

As defined in classical chemistry, elements that differ only in numbers of neutrons are called *isotopes*. Many elements have two, or more, naturally occurring isotopes. In the case of oxygen, there are three stable (nonradioactive) isotopes ¹⁶O, ¹⁷O, and ¹⁸O. Their nominal abundances are about 99.76, 0.04, and 0.20%, respectively. The superscripts (16, 17, and 18) refer to their atomic masses.

In isotope geochemistry, abundances of rare isotopes are expressed by the ratio of rare/abundant isotope (*X). For example, for H₂O, this ratio is defined as

$${}^*X(\text{H}_2\text{O}) = \frac{({}^*\text{O in water})}{({}^{16}\text{O in water})} \quad [1]$$

where ‘*’ stands for 17 or 18.

The ratios *X can be measured with very high precision with mass spectrometers (for more details, see [Mook, 2000](#)). These isotope ratios are usually not measured on the pure elements, but instead on a gas containing the element under consideration. For most purposes, oxygen isotope ratios are measured on O₂ or CO₂. The common procedure is to compare the isotope ratio of a sample relative to that of a standard, which is assumed to have a known isotopic ratio. In this way, variations due to instrumental drifts and instabilities are considerably reduced because they occur in both sample and standard and cancel out. Using international standards has an additional advantage because it allows comparison of isotopic ratios among laboratories.

Because the measurements are done with respect to standards, it is customary to report the results as deviations (δ) of the isotope ratio of a sample from that of a standard:

$${}^*\delta = \frac{{}^*X_{\text{sample}} - {}^*X_{\text{standard}}}{{}^*X_{\text{standard}}} = \frac{{}^*X_{\text{sample}}}{{}^*X_{\text{standard}}} - 1 \quad [2]$$

As can be seen, the δ value is dimensionless. For natural substances, the δ values are small, and thus, they are usually multiplied by 10³ and given in ‰ (per mil).

For a specific isotope (e.g., ¹⁸O in O₂), eqn [2] is usually presented in the form

$$\delta^{18}\text{O} = \frac{({}^{18}\text{O}/{}^{16}\text{O})_{\text{sample}}}{({}^{18}\text{O}/{}^{16}\text{O})_{\text{standard}}} - 1 \quad [3]$$

The choice of standard depends on the material being analyzed. In the case of O₂ gas, measurements are performed directly on the material of interest, and atmospheric O₂ is an ideal standard because atmospheric mixing is very rapid compared to the processes altering the isotopic composition of oxygen and, therefore, δ¹⁸O and δ¹⁷O of atmospheric O₂ are very uniform.

In contrast, in the case of water, most measurements are not done on the original material. Instead, they are performed on gaseous CO₂ or O₂ reacted with or obtained from samples and standards, and the Vienna Standard Mean Ocean Water (VSMOW) standard is used for water. These standards have been discussed in detail in various publications (e.g., [Coplen, 1994](#); [Coplen et al., 1983](#); [Gonfiantini, 1978](#)), so here, only a short summary is given.

Standard, VSMOW is a well-preserved batch of distilled water prepared mainly from ocean water. In 1976, it was decided by the International Atomic Energy Agency (IAEA) to use VSMOW for fixing the zero point of the δ¹⁸O water scale (δ¹⁸O_{VSMOW} ≡ 0). Analyses of ¹⁸O¹⁶O are performed by equilibrating a water sample with CO₂ of a known isotopic composition, followed by mass spectrometric analysis of equilibrated CO₂. Usually, water samples consisting of unknown samples are run together with the standard or some laboratory reference water well calibrated relative to VSMOW. Samples and standards are run under the same conditions:

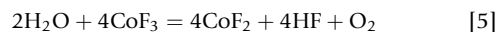
$$\begin{aligned} \delta^{18}\text{O}(\text{sample vs. VSMOW}) \\ = \delta^{18}\text{O}(\text{sample CO}_2 \text{ vs. VSMOW-CO}_2) \end{aligned} \quad [4]$$

The relationship between VSMOW and air O₂ is known from [Barkan and Luz \(2011\)](#) as

$$\delta^{18}\text{O}_{\text{gas}}(\text{vs. VSMOW}) = 0.97668\delta^{18}\text{O}_{\text{gas}}(\text{vs. air O}_2) - 23.324$$

Measurements of δ¹⁷O are difficult, first, because ¹⁷O is less abundant than ¹⁸O by a factor of about five and, second, because precise measurements are not possible with CO₂ gas. The reason for this is the masking effect of the much more abundant ¹³C on the isotopically substituted CO₂ at the same molecular mass of 45 (¹³C¹⁶O¹⁶O and ¹²C¹⁷O¹⁶O). It is noted that even with the recent elegant development of [Assonov and Brenninkmeijer \(2001\)](#), the measurement precision of δ¹⁷O in CO₂ gas was only ±0.3‰. Such precision is insufficient for most geochemical applications.

Measurements of δ¹⁷O are straightforward on pure O₂ gas and with proper correction also in O₂+Ar mixture ([Barkan and Luz, 2003](#)). In the case of water, it can be quantitatively converted to O₂ gas by fluorination using CoF₃:



[Barkan and Luz \(2005\)](#) applied this method and demonstrated very high precision (0.01–0.03‰) in measurements of both δ¹⁷O and δ¹⁸O.

5.14.2.1.1 Isotopic fractionation

The chemical and physical properties of compounds containing rare isotopes are very similar to those of the same compound containing major isotopes. Nonetheless, there are small differences in the physical properties (e.g., density, vapor pressure, boiling, and melting points) mainly due to the higher vibration energy of the lighter isotope and higher binding energies in the heavier molecules ([Bigeleisen, 1952](#); [Bigeleisen and Wolfsberg, 1958](#); [Urey, 1947](#)). Consequently, as a rule, lighter molecules have higher zero-point energies and react and move faster. This results, for example, in smaller evaporation pressure of H₂¹⁸O than H₂¹⁶O and smaller respiration rate of ¹⁸O¹⁶O than ¹⁶O¹⁶O.

The above differences lead to separation of isotopic compounds during processes such as evaporation or condensation, melting or crystallization, diffusion in different media, and isotopic exchange reactions in chemical equilibrium (e.g., dissolved CO₂ and bicarbonate) or in physical reactions (e.g., water vapor and liquid water). Such separation is called *isotope fractionation*. The extent of this fractionation depends on mass differences among various isotopes and on

temperature. For example, fractionation is about twice as large for ¹⁸O/¹⁶O/¹⁶O/¹⁶O as for ¹⁷O/¹⁶O/¹⁶O/¹⁶O, or in other words, variations in δ¹⁷O are about one-half those in δ¹⁸O.

The isotope fractionation is described mathematically by comparing the isotope ratios of two compounds in equilibrium or of the compounds before and after physical or chemical process (e.g., Mook, 2000). A quantity known as the *fractionation factor* (α) is defined as

$$\alpha_{B/A} = \frac{X_B}{X_A} \quad [6a]$$

or

$$\alpha_{A/B} = \frac{X_A}{X_B} \quad [6b]$$

Equations [6a] and [6b] express the isotope ratio in the phase or compound B relative to that in A or vice versa. For example, X_A may represent ¹⁸O/¹⁶O (or ¹⁸O¹⁶O/¹⁶O¹⁶O) in O₂ reservoir and X_B ¹⁸O/¹⁶O (or ¹⁸O¹⁶O/¹⁶O¹⁶O) in O₂ removed by respiration. In this case, α is known as the *respiratory fractionation factor*.

Typically, fractionation factors are quantities slightly larger [6b] or smaller [6a] than unity (e.g., α_{B/A} = 0.9800 and α_{A/B} = 1.0204). Often, it is convenient to express fractionation as isotope effect (ε), and the relation between α and ε is

$$\varepsilon = \alpha - 1 \quad [7]$$

Because magnitudes of ε are always small, they are multiplied by 1000 and reported in ‰. For example, ε_{B/A} for ordinary respiration is about -20‰ (or ε_{A/B} = 20.4‰).

Comparing eqns [6] and [7] clearly shows that while fractionation factors are multiplicative (α_{total} = α₁α₂...α_n), ε is approximately additive (ε_{total} ≈ ε₁ + ε₂ + ... + ε_n). This explains why in precise calculations of fractionations that result from more than one process, it is necessary to use α and not ε.

Isotope effects can be measured directly in cases where the two phases or compounds are available. For example, the isotope effect of photosynthesis can be measured by comparison of δ¹⁸O of O₂ formed by this process from water to δ¹⁸O of the water substrate. However, for certain processes, such as respiration, it is either impossible or not practical to measure the starting and end materials. In such cases, the isotope effect (ε) is determined from fractionation experiments in which a starting amount of material is lost by a process whose isotope effect is unknown (e.g., evaporation or respiration) and from the change in isotopic ratio X (¹⁸O/¹⁶O or ¹⁷O/¹⁶O) as the amount of the material (N) goes down (such experiments are known as Rayleigh fractionation experiments in the stable isotope literature):

$$\ln \frac{X}{X_0} = \left(\frac{1}{\alpha - 1} \right) \ln \left(\frac{N}{N_0} \right) \quad [8]$$

where the subscript 0 stands for initial X and N (see Mook (2000) for further detail).

An example of the derivation of soil respiratory fractionation from a Rayleigh experiment is shown in Figure 1. Soil was placed in a closed and dark container and the fraction of remaining oxygen and its δ¹⁸O were recorded. The results are shown in Figure 1. In this particular experiment, the consumption of ¹⁶O was about 1.5% faster than ¹⁸O consumption. Because of this difference in consumption rates, the ratio

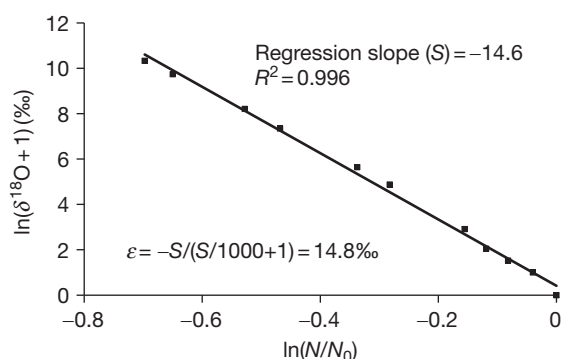


Figure 1 ln(δ¹⁸O + 1) versus ln(N/N₀) in Rayleigh fractionation.

¹⁸O/¹⁶O increased as the remaining fraction of O₂ went down. The isotope effect (ε) was calculated from the regression slope in Figure 1 as 14.8‰.

5.14.2.2 Relationships between Fractionations of ¹⁷O/¹⁶O and ¹⁸O/¹⁶O

In certain geochemical studies, such as the isotopic composition of atmospheric O₂, it is of interest to know the relationships between fractionations of ¹⁷O/¹⁶O and ¹⁸O/¹⁶O. For most terrestrial processes, this relation is mass-dependent and can be derived from theoretical considerations (e.g., Matsuhisa et al., 1978; Weston, 1999; Young et al., 2002) showing that

$${}^{m+1}\alpha_{A/B} = ({}^{m+2}\alpha_{A/B})^\lambda \quad \text{or} \quad \ln({}^{m+1}\alpha_{A/B}) = \lambda \ln({}^{m+2}\alpha_{A/B}) \quad [9]$$

where m is the mass of the abundant isotope (16 in the case of oxygen).

For mass-dependent fractionations, the value of factor λ is slightly larger than 0.5. For example, in the case of respiratory fractionation of ¹⁷O/¹⁶O and ¹⁸O/¹⁶O, it can be approximated as 0.515 = ln(16/17)/ln(16/18) (e.g., Young et al., 2002).

Experimental studies showed that indeed in most natural processes the fractionation for the heaviest isotope is about twice as large as that for the lighter isotope (Barkan and Luz, 2007; Helman et al., 2005; Miller et al., 2002). It was also shown that the factor λ varies slightly (0.500–0.529) depending on the isotope fractionation processes (kinetic or steady state) and among the various biological processes. When results of a Rayleigh process (eqn [10]) are displayed in ln(δ¹⁷O + 1) versus ln(δ¹⁸O + 1) plots, they fall on straight lines. An example of such a process is given in Figure 2 for the case of O₂ consumption by respiration. As can be seen, the ln(¹⁷α)/ln(¹⁸α) slope (λ) in this experiment equals 0.518, which is nearly as expected from theory.

Whereas mass-dependent fractionation is the rule for most terrestrial processes, there are exceptions with mass-independent fractionation in which λ is much greater than 0.52 (e.g., Clayton et al., 1973; Thiemens and Heidenrich, 1983). For example, the isotopic composition of atmospheric O₂ is affected by photochemical reactions in the stratosphere (e.g., Luz et al., 1999; Thiemens, 1999). Due to these reactions, ozone and CO₂ in the stratosphere have anomalously high δ¹⁷O, and atmospheric O₂ is anomalously low in δ¹⁷O in comparison with O₂ produced from water by photosynthesis (Luz et al., 1999).

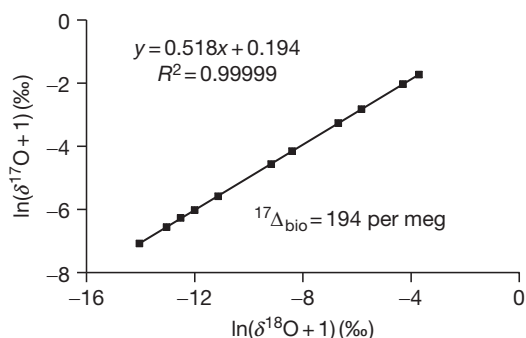


Figure 2 $\ln(\delta^{17}\text{O}+1)$ versus $\ln(\delta^{18}\text{O}+1)$ in Rayleigh respiratory fractionation.

Deviations from mass-dependent behavior have been described in the literature, and in many cases, the anomaly of ^{17}O ($\Delta^{17}\text{O}$) was calculated as $\Delta^{17}\text{O} = \delta^{17}\text{O} - 0.52\delta^{18}\text{O}$ (e.g., Thiemens, 1999). Yet, as discussed by Miller (2002), differences among calculated values of $\Delta^{17}\text{O}$ of various materials depend on the choice of isotopic references. To overcome this problem and in order to facilitate direct comparison of results among laboratories, he proposed using natural log transformed deltas (e.g., $\ln(\delta^{17}\text{O}+1)$) instead of $\delta^{17}\text{O}$). Following Miller (2002), Luz and Barkan (2005) defined a useful parameter ($^{17}\Delta$) for studies of atmospheric and marine dissolved O₂:

$$^{17}\Delta = \ln(\delta^{17}\text{O}+1) - \lambda \ln(\delta^{18}\text{O}+1) \quad [10]$$

The $^{17}\Delta$ parameter indicates an excess or deficit of ^{17}O in a sample with respect to the atmospheric O₂. Because the magnitudes of $^{17}\Delta$ values calculated with eqn [10] are very small, they are multiplied by 10^6 and the reported values are in per meg.

From eqn [10], it is clear that the magnitude of a calculated $^{17}\Delta$ depends on the chosen value of λ and that the absolute value of this magnitude will increase when $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ of measured samples are large. Luz and Barkan (2005) discussed such $\delta^{17}\text{O}$, $\delta^{18}\text{O}$, and $^{17}\Delta$ in studies of dissolved marine O₂. In such studies, it is advantageous to use atmospheric O₂ as a standard (rather than VSMOW) because measured $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ values are generally close to those of atmospheric oxygen. In addition, the reference slope λ should be selected such that it optimally fits the system under consideration, and in this respect, a value of 0.518 is optimal for studies of marine dissolved O₂.

It is evident that a Rayleigh process will result in perfectly linear trends in plots of $\ln(\delta^{17}\text{O}+1)$ versus $\ln(\delta^{18}\text{O}+1)$, and Figure 2 is a good illustration of such linearity. However, when two O₂ volumes having different isotopic composition are mixed, the resulting mixtures will have isotopic compositions that fall on curves in $\ln(\delta^{17}\text{O}+1)$ versus $\ln(\delta^{18}\text{O}+1)$ plots. Such mixing, for example, takes place when photosynthetic O₂ is added to an existing reservoir of dissolved O₂ (Figure 3).

5.14.3 $^{18}\text{O}/^{16}\text{O}$ Ratios in Atmospheric O₂

5.14.3.1 The Dole Effect and Its Magnitude

On timescales of thousands of years, the concentration of atmospheric O₂ is held nearly constant by equal rates of photosynthetic production and respiratory consumption (e.g., Broecker, 1970). Because the substrate from which

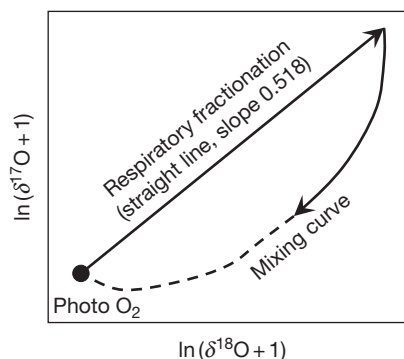


Figure 3 $\ln(\delta^{17}\text{O}+1)$ versus $\ln(\delta^{18}\text{O}+1)$ showing that respiratory fractionation results in isotopic composition changing along a straight line. Photosynthesis adds new O₂ that mixes with existing O₂ along a curved line.

photosynthesis produces O₂ is water, and because the ocean is, by far, the largest reservoir of water on Earth, the $\delta^{18}\text{O}$ of atmospheric O₂ must be strongly linked to that of seawater. Nevertheless, as demonstrated independently by Dole (1935) and Morita (1935), the isotopic composition of H₂O and air O₂ is not identical, and the latter is enriched in ^{18}O . The deviation of $\delta^{18}\text{O}$ of air O₂ from that of seawater oxygen is known as the DE.

An accurate determination of the value of the DE was first published by Kroopnick and Craig (1972) as $23.5 \pm 0.3\%$. They applied the CO₂ equilibration method for measuring $\delta^{18}\text{O}$ of the SMOW standard (which represents seawater), while CO₂ produced by the oxidation of graphite with oxygen gas was used for the measurement of $\delta^{18}\text{O}$ of atmospheric O₂. The value determined in this case depends on the fractionation factor for the equilibrium of CO₂ with water. This factor varies in the range 1.0407–1.0417 (e.g., Brenninkmeijer et al., 1983) and is a main source of uncertainty in the DE magnitude. To circumvent this problem, Horibe et al. (1973) converted atmospheric O₂ to water (by oxidation of H₂ gas) and then used the CO₂ equilibration method for measuring $\delta^{18}\text{O}$ of this water and also that of the seawater standard SMOW. The magnitude in this case was $23.8 \pm 0.3\%$. By avoiding the problem of CO₂–H₂O equilibration in an alternative way, Barkan and Luz (2005) obtained a value of $23.88 \pm 0.02\%$ by direct comparison of the isotopic composition of atmospheric oxygen with O₂ produced by fluorination of the VSMOW standard.

5.14.3.2 Processes Influencing the Dole Effect

5.14.3.2.1 Biological O₂ consumption

The major causes of the DE are biological processes that preferentially consume O₂ molecules containing only light stable isotopes ($^{16}\text{O}^{16}\text{O}$). The result of this consumption is that O₂ molecules containing heavier isotopes ($^{17}\text{O}^{16}\text{O}$ and $^{18}\text{O}^{16}\text{O}$) become enriched in the remaining gas. The magnitude of such fractionations has been studied primarily in small-scale Rayleigh-type experimental systems (see Section 5.14.2.1) and also in large-scale field experiments.

Rabinowitch (1945) was the first to suggest that the cause of the DE might be isotopic fractionation during respiration. Following his suggestion, Lane and Dole (1956) conducted simple respiration experiments and found preferential

consumption of ¹⁶O over ¹⁸O (¹⁸ε) of 7–25‰. Similar simple experiments with natural plankton or cultures of marine organisms were done by Kroopnick (1975), Kiddon et al. (1993), Quay et al. (1995), and Luz et al. (2002). Kiddon et al. (1993) determined the discrimination by main marine eukaryotes as 21‰, but weaker by bacteria (19‰). Kroopnick (1975), in experiments with surface ocean communities, obtained an average respiratory fractionation of 20.8‰, while Luz et al. (2002) reported a value of 21.6‰ for Lake Kinneret in which the plankton community is dominated by phytoplankton. Weaker fractionation (17.6‰) was obtained by Quay et al. (1995) for the Amazon River where O₂ uptake is dominated by bacteria.

In experiments with plants and phytoplankton, Guy et al. (1989, 1992) showed that ¹⁸ε in ordinary respiration through the cytochrome oxidase (COX) pathway was 17.4–19.9‰, and much larger ¹⁸ε (24.1–26.2‰) was associated with the alternative oxidase (AOX), which is resistant to cyanide. They also pointed out that in large objects, O₂ supply is limited by slow diffusion and suggested that this was the reason for very weak fractionation in intact carrot and potato (8–10‰) in the early experiments of Lane and Dole (1956). In an opposite way, this may also be the reason for the larger ¹⁸ε values (~21‰ for COX and ~32‰ for AOX) reported from measurements of isolated mitochondria by Ribas-Carbo et al. (1995). The same authors also showed that ¹⁸ε values due to AOX were larger in green tissues than in nongreen tissues (30–32‰ and 24–26‰, respectively).

All the experiments described above were done in darkness in order to avoid interference from O₂ addition by photosynthesis, but, by their design, such dark experiments are not useful for measuring O₂ consumption by mechanisms that are engaged only under illumination. One of these mechanisms is photorespiration, which involves two enzymes, rubisco oxygenase and glycolate oxidase, and is engaged when CO₂ supply is below optimum for plant growth. It is estimated that some 30% of the global photosynthetic production of O₂ is consumed by photorespiration (Bender et al., 1994). Additional O₂-consuming processes that require light are Mehler reaction (Mehler, 1951), chlororespiration (Bennoun, 1994), and the plastoquinol terminal oxidase (Bailey et al., 2008).

Guy et al. (1993) designed special experimental setups for *in vitro* measurement of ¹⁸O/¹⁶O isotope effects (¹⁸ε) of Mehler

reaction, rubisco oxygenase, and glycolate oxidase in the absence of O₂ production by photosynthesis. Later on, Helman et al. (2005) followed Guy et al.'s (1993) design and ran similar experiments in which they measured, in addition to ¹⁸ε, the ¹⁷O/¹⁶O isotope effects (¹⁷ε). The ¹⁸ε results from all these experiments range from about 11‰ for Mehler reaction to 22‰ for glycolate oxidase (see Table 1 for more details). The magnitude of ¹⁸ε associated with O₂ consumption by chlororespiration or the plastoquinol terminal oxidase is unknown.

From the results of the experimental work reviewed earlier, it is clear that biological consumption of O₂ is an important factor in the DE. However, the range of ¹⁸ε values in various organisms and biological mechanisms is very broad (11–32‰), and there is no a priori way to accurately obtain a weighted average ¹⁸ε of relevance for global calculations of the DE. Fortunately, certain globally important systems lend themselves to large-scale field experiments that are useful for estimating overall biological fractionations relevant for understanding the DE. These systems include soils and the ocean below the surface photic zone.

Excluding light-dependent mechanisms, such as photorespiration and Mehler reaction, soils are the major O₂-consuming systems in the terrestrial environment. Soils contain variable amounts of air in passages among soil particles. Because air diffusion into soils is always limited to a certain extent and because roots and bacteria respire, O₂ concentration in soil air is generally significantly lower than that in the overlaying atmosphere. Likewise, due to respiration, δ¹⁸O of soil air is greater than that of atmospheric O₂. These differences make it possible to calculate effective ¹⁸ε (¹⁸ε_{effective}) of soil O₂ uptake from measurements of the ratio O₂/Ar and δ¹⁸O in soil air. The effects of slow diffusion on ¹⁸ε_{effective} of soils have been studied extensively by Angert and Luz (2001) and Angert et al. (2001, 2003a).

When O₂ is respired in roots or microorganisms in soil aggregates, ¹⁸ε_{effective} is influenced by ¹⁸ε of respiration (¹⁸ε_r) and by ¹⁸ε of O₂ diffusion (¹⁸ε_{diff}) in air or water. In soils, several steps of diffusion and respiration occur in series, but the resulting effective fractionation cannot be simply calculated by addition of the effects in each of these steps. The total effect of a series of fractionating steps has been treated theoretically by Farquhar et al. (1982). Following them in simple case of O₂ consumption in an airtight chamber by plant roots or soil aggregates (Figure 4), ε_{effective} is expressed by

Table 1 Isotopic effects in biological O₂ uptake processes

	¹⁸ ε (‰)	¹⁷ ε/ ¹⁸ ε
Ordinary respiration by cytochrome oxidase (COX)	17.4–19.9 (Guy et al., 1989, 1992)	0.518 (Angert et al., 2003a; Luz and Barkan, 2005)
Cyanide resistant respiration by alternative oxidase (AOX)	24.1–26.2 (Guy et al., 1989, 1992)	0.518 (Angert et al., 2003a; Luz and Barkan, 2005)
Mehler reaction – reduction of O ₂ to produce water when photosynthetic production of O ₂ is high	10.9–15.3 (Guy et al., 1993; Helman et al., 2005)	0.526 (Helman et al., 2005)
Photorespiration – O ₂ consumption when CO ₂ supply is limited		
Occurs in two steps:		
1. By rubisco oxygenase	21.3–21.8 (Guy et al., 1993; Helman et al., 2005)	0.517 (Helman et al., 2005)
2. By glycolate oxidase	22.0–22.7 (Guy et al., 1993; Helman et al., 2005)	0.501 (Helman et al., 2005)

$$^{18}\epsilon_{\text{effective}} = ^{18}\epsilon_{\text{diff}} + \frac{(^{18}\epsilon_r - ^{18}\epsilon_{\text{diff}})[\text{O}_2]_i}{[\text{O}_2]_a} \quad [11]$$

where $[\text{O}_2]_i$ and $[\text{O}_2]_a$ stand for O₂ concentration inside the root or aggregate and in the chamber, respectively. For root walls and often for soil aggregates, O₂ diffusion is through water in which $^{18}\epsilon_{\text{diff}}$ is very small (Farquhar and Lloyd, 1993). Taking it as zero, the equation simplifies to $^{18}\epsilon_{\text{effective}} = (^{18}\epsilon_r [\text{O}_2]_i) / [\text{O}_2]_a$. In turn, the ratio $[\text{O}_2]_i / [\text{O}_2]_a$ is unity when O₂ is free to move back and forth from the root to the chamber free space and vice versa and becomes smaller when resistance to O₂ exchange becomes stronger. Accordingly, $^{18}\epsilon_{\text{effective}}$ will equal $^{18}\epsilon_r$ or become much smaller, respectively. Such simple relationships have been demonstrated in experiments (Table 2) that showed that $^{18}\epsilon_{\text{effective}}$ of soil aggregates decreased by wetting them with water. The added water increased resistance to O₂ exchange and consequently $[\text{O}_2]_i / [\text{O}_2]_a$ must have been also decreased. As a result, $^{18}\epsilon_{\text{effective}}$ became smaller. An opposite tendency was demonstrated by disaggregation and suspension of small soil particles in water. There was little resistance to O₂ exchange by these small particles, $[\text{O}_2]_i / [\text{O}_2]_a$ must have been close to unity, and $^{18}\epsilon_{\text{effective}}$ increased significantly.

In soils, the situation is more complex. They can be pictured as a mixture of small soil aggregates, roots, and free passages among them. The free space of the passages may be occupied by air, water, or both. Respiration and associated isotope effects take place in roots and aggregates, and in addition, there are also isotope effects of diffusion on a micro scale, and the latter may be quite large (e.g., Aggarwal et al., 2004). Yet, there is no

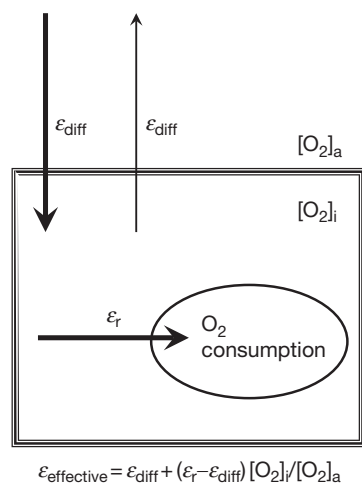


Figure 4 Schematic illustration of effective and respiratory isotope effects.

Table 2 Effects of diffusion on $^{18}\epsilon_{\text{effective}}$

Incubation experiment	$^{18}\epsilon_{\text{effective}}$ (‰)
Soil aggregates	13.6
Soil aggregates + water	8.5
Soil disaggregated in water	20.5

Source: Angert A, Luz B, and Yakir D (2001) Fractionation of oxygen isotopes by respiration and diffusion in soils: Implications for the isotopic composition of atmospheric O₂. *Global Biogeochemical Cycles* 14: 871–881.

practical way for estimating the isotope effects of each of these processes in individual aggregates or roots. It is possible, on the other hand, to assess the overall effective fractionation of the entire soil complex. In doing so, it is assumed that diffusive fractionation in the larger air-filled passages is 14.3‰ as for binary diffusion of O₂ in air (e.g., Mason and Marrero, 1970).

By applying a model and using measurements of $\delta^{18}\text{O}$ and $\delta\text{O}_2/\text{Ar}$ of soil air, Angert et al. (2003a) estimated $^{18}\epsilon_{\text{effective}}$ of soils from various climatic regions. These were 10.8, 17.8, and 22.5‰ for tropical, temperate, and boreal soils. From these data, Landais et al. (2007a) calculated a global weighted average of soil fractionation as 15.8‰. The small fractionation in tropical soils may be related to the moisture content blocking free movement of O₂, thus limiting supply to consumption sites. In turn, the larger fractionation in temperate and boreal soils than in tropical soils is likely the result of engagement of respiration through the AOX, which strongly discriminates against ^{18}O . Heat production by AOX is probably advantageous at low soil temperatures.

The ocean below the photic zone (usually below ~ 150 m) is another system where large-scale fractionations can be studied with relative ease. Variations in the concentration of O₂ below the photic zone are controlled primarily by bacterial consumption and horizontal and vertical mixing. Typically, an O₂ minimum zone exists in mid depth of all oceanic basins where bacterial consumption dominates. This minimum zone was targeted by Dole and coworkers for studying $^{18}\epsilon_{\text{effective}}$ of the deep sea (Rakestraw et al., 1951) because there they expected to find large concentration and isotopic variations. Indeed, they found good correlation between high $^{18}\text{O}/^{16}\text{O}$ ratios and low O₂ concentration, and the strong positive correlation led them to the conclusion that there is deep metabolism in the ocean. Assuming that mixing played only a minor role, they used a closed-system Rayleigh calculation to derive a weak isotope effect of only several per mil. Following them, Kroopnick and Craig (1976) measured O₂ concentration and $\delta^{18}\text{O}$ in several stations in the world ocean. They applied a more sophisticated model in which vertical diffusive mixing was included in addition to respiration and derived an optimal respiratory isotope effect of 10‰, a weak fractionation similar to that reported by Rakestraw et al. (1951). More recently, similar weak fractionation for the O₂ minimum zone has been reported by Bender (1990) and Levine et al. (2009).

Whereas, due to large concentration and isotopic variations, the O₂ minimum zone seems ideal for studying marine isotope effects, the fractionation there is considerably smaller than in laboratory experiments with marine organisms. This raises the question whether these experiments indicate relevant isotope effects for natural marine systems. One possibility is that most of the respiration in the O₂ minimum takes place inside sinking particles into which gas diffusion is slow with steep exterior–interior O₂ gradients. In this case, similarly to roots and soil aggregates, $^{18}\epsilon_{\text{effective}}$ would be small. And a further question is whether similar fractionation, regardless of its mechanistic explanation, would apply to overall marine respiration.

It is well known that most of the sinking organic matter produced by photosynthesis in the euphotic zone is oxidized at a relatively shallow depth of the subphotic zone (e.g., Jenkins and Goldman, 1985). Away from the equatorial region, there are significant seasonal variations in O₂ concentration in the

subphotic zone, starting with high concentration after the late winter overturn followed by gradual lowering of O₂ levels in the spring–summer stratified period (Figure 5). Because of the seasonal variations in this zone, estimation of the respiratory fractionation is difficult as it requires careful δ¹⁸O measurements of many samples. So such estimation has been done in only one study in the North Atlantic near Bermuda from δ¹⁸O measurements taken from the same samples shown in Figure 5(a) (Luz and Barkan, 2011). As O₂ decreased between March and September, δ¹⁸O of dissolved O₂ increased. Assuming water at 250 m as a closed system, ¹⁸ε_{effective} for the subphotic zone can be calculated as 19.7‰ (Figure 5(b)). This calculation is in good agreement with Levine et al.’s (2009) estimate of 19‰ derived from a wide ocean budget for the mesopelagic zone (excluding the equatorial O₂ minimum zone mentioned earlier). Both estimates are in agreement with isotope effects derived from small-scale experiments with bacteria.

An important question is to what extent the O₂ minimum zone impacts the DE. To answer this question, it is instructive to compare respiration rates there to such rates in the subphotic zone. For instance, the respiration rate in the subtropical Atlantic is about 27 μmol kg⁻¹ year⁻¹ (Figure 5(a)), and rates in the more fertile parts of the ocean must be greater. In contrast, the rate of respiration in the O₂ minimum is very low, and Levine et al. (2009) estimated only ~2.9 μmol kg⁻¹ year⁻¹ in the 300–1000 m layer of the equatorial Atlantic where [O₂] is minimal. So the respiration rate in the O₂ minimum zone is two orders of magnitude less than the rate in the subphotic zone and the global effect of the small respiratory fractionation in this zone is negligible. Therefore, an overall ε_{effective} of 19.4 ± 0.5‰ can be taken as a representative value for in situ respiration in the entire ocean below the photic zone.

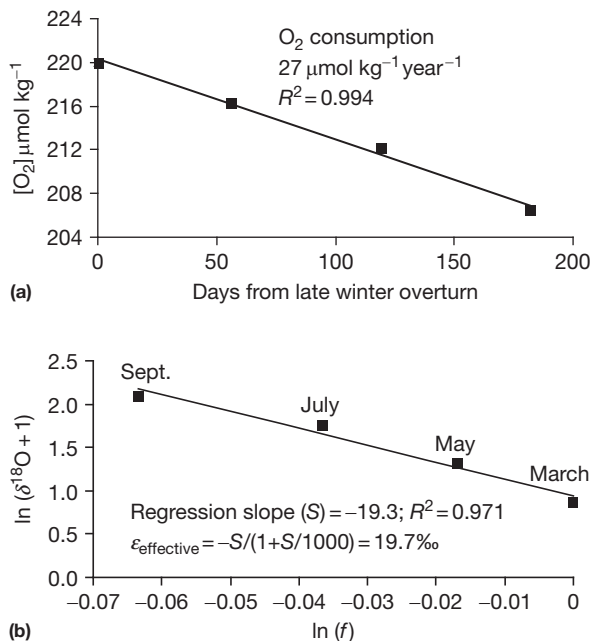


Figure 5 (a) Rate of O₂ consumption in the oceanic thermocline. (b) ¹⁸ε_{effective} in the oceanic thermocline.

In summary, the ¹⁸ε_{effective} of the dark parts of the ocean of 19.4‰ is larger than that of global soils (~15.8‰). These are the only components of the global O₂ system for which it is possible to derive direct integrated estimates.

In the terrestrial biosphere, aboveground O₂ uptake is probably dominated by photorespiration with combined isotope effects of both rubisco oxygenase and glycolate oxidase of ~21.9‰. O₂ uptake by Mehler reaction (¹⁸ε ~ 11‰) in leaves is considered to be less than 10% of total O₂ uptake (e.g., Bader et al., 2000), and ordinary and cyanide-resistant respiration is assumed to account for about 16% (Landais et al., 2007a). Some details on the proportions of these O₂ uptake pathways and associated fractionations are given in Table 3. The overall O₂ uptake isotope effect in leaves is ~19.2‰, but there is substantial uncertainty in this value.

In addition to the biological O₂ consumption discussed earlier, it is known that a considerable fraction of the global oxygen uptake takes place in the ocean surface. However, its ¹⁸ε_{effective} cannot be estimated without first considering isotopic effects of photosynthesis.

5.14.3.2.2 Photosynthetic O₂ production

With the exception of negligible O₂ production by the photolysis of atmospheric water vapor, all O₂ on Earth is produced by photosynthesis. Knowing this, Dole and Jenks (1944) ran photosynthesis experiments and obtained O₂ gas produced by plants and algae and compared its ¹⁸O/¹⁶O ratio to that of the substrate water. Expressed in the present conventional δ¹⁸O terms, the O₂ gas was enriched in ¹⁸O by about 5‰ with respect to the substrate water. However, the validity of this pioneering observation has been repeatedly challenged by several studies (Guy et al., 1993; Helman et al., 2005; Luz and Barkan, 2005; Vinogradov et al., 1959; Yakir et al., 1994). In particular, Guy et al. (1993) demonstrated similar δ¹⁸O of water and photosynthetic O₂ produced by spinach thylakoids, cyanobacteria, and diatoms. Helman et al. (2005) confirmed this observation in experiments with cyanobacteria. Moreover, from mass balance calculations, Luz and Barkan (2005) confirmed the similarity of δ¹⁸O of water and O₂ produced by *Phylodendron*. All these studies led to the understanding that the enrichment reported by Dole and Jenks (1944) does not represent the original isotopic composition of photosynthetic O₂ and was caused by their failure to eliminate partial consumption by respiration and/or atmospheric contamination in their experimental setup.

Table 3 Calculated ¹⁸ε_{effective} of O₂ uptake in leaves

Process and fraction of terrestrial gross O ₂ production	¹⁸ ε _{effective} (‰)
Photorespiration (30%)	21.9 ± 1.0
Mehler reaction (10%)	10.9 ± 0.2
COX and AOX respiration (16%)	19.4 ± 1.0
Total leaf O ₂ uptake (56%)	19.2 ± 2.0

Source: Landais A, Lathiere J, Barkan E, and Luz B (2007a) Reconsidering the change in global biosphere productivity between the Last Glacial Maximum and present day from the triple oxygen isotopic composition of air trapped in ice cores. *Global Biogeochemical Cycles* 21: GB1025, <http://dx.doi.org/10.1029/2006GB002739>.

The notion that there is no fractionation of oxygen isotopes during photosynthesis has been deeply entrenched and has guided the interpretation of experimental results until very recently. Yet, failure of global budgets, such as in [Bender et al. \(1994\)](#), to account for the magnitude of the DE has led [Eisenstadt et al. \(2010\)](#) to revisit the problem of possible fractionation in photosynthesis, and they ran photosynthetic experiments in which O₂ was produced by various cultures of marine phytoplankton. These experiments complement a small number of well-controlled ones done by [Guy et al. \(1993\)](#), [Luz and Barkan \(2005\)](#), and [Helman et al. \(2005\)](#).

To yield meaningful results, it is necessary to completely eliminate any respiratory consumption of the newly produced O₂ or alternatively to correct for respiratory consumption. The latter is possible only if accurate rates of gross O₂ production, respiratory consumption, and its isotope effect are known. Accordingly, three experimental setups have been designed:

1. O₂ was produced *in vitro* by thylakoids (photosynthetic membranes) separated from fresh market spinach. These membranes do not respire, and O₂ produced by photosynthesis had δ¹⁸O identical (±0.2‰) to that of the substrate water (see [Guy et al. \(1993\)](#) for details).
2. An airtight terrarium experiment was set up under controlled conditions in which δ¹⁸O leaf water was known, photorespiration was prevented by high CO₂ in the terrarium air, and practically all O₂ uptake was by soil respiration, whose ¹⁸ε_{effective} was known from dark experiments. Rates of gross (G) and net (N) O₂ production and thus light O₂ uptake (U=G-N) were also known from a separate experiment with the same terrarium using ¹⁸O-spiked O₂ gas. More details on the experiment are given in [Luz and Barkan \(2005\)](#). Using their data, it is estimated that ¹⁸O of the photosynthetic O₂ was only slightly enriched (~0.7‰) with respect to the leaf water.
3. Experiments for O₂ production by marine cyanobacteria and algae have been set up by [Guy et al. \(1993\)](#), [Helman et al. \(2005\)](#), and [Eisenstadt et al. \(2010\)](#). Special attention was paid to prevent δ¹⁸O alteration by partial O₂ consumption or by atmospheric contamination. In order to allow optimal growth conditions and prevent photorespiration, cultured algal and cyanobacteria cells were suspended in an adequate growth medium supplemented with at least 5 mM NaHCO₃ and a buffer for maintaining a pH of 8.2 to prevent photorespiration. The light intensity was adjusted for maximal O₂ production. The medium was illuminated and O₂ produced was immediately removed by stirring and vigorously bubbling with 99.999% grade He. Results from these experiments are given in [Table 4](#) and show negligible (~0.4‰) to substantial (~6‰) ¹⁸O enrichment in the newly produced O₂ with respect to the substrate water.

Table 4 Photosynthetic enrichment in ¹⁸O in marine phytoplankton

	δ ¹⁸ O versus substrate water (‰)
Cyanobacteria (Helman et al., 2005)	0.47
Green algae (Eisenstadt et al., 2010)	2.85
Diatoms (Eisenstadt et al., 2010)	4.43
Coccolithophores (Eisenstadt et al., 2010)	5.81

Experiments of types 1 and 2 are the only ones giving information relevant to higher plants and thus to terrestrial O₂ production. The results show that O₂ originating from terrestrial photosynthesis is similar or slightly enriched in ¹⁸O with respect to leaf water. Experiments of type 3 are relevant to O₂ produced by marine photosynthesis and show that with the exception of cyanobacteria, O₂ originating from marine photosynthesis is enriched in ¹⁸O with respect to seawater. The extent of this enrichment in the ocean depends on the proportion of gross production rates by various phytoplankton groups (but not necessarily their abundances). So the overall ¹⁸O enrichment in marine systems cannot be easily assessed from a priori consideration. Despite that, ocean-wide estimates of a combined isotope effect due to photosynthesis and O₂ uptake (¹⁸ε_{up}) in the surface ocean can be obtained.

The combined isotope effect in the surface ocean can be determined from the mass balance of ¹⁸O in dissolved O₂. Such mass balance should take into account O₂ gains and losses due to photosynthesis, uptake, air-sea gas exchange, and mixing of O₂ from the deeper layer below the surface mixed layer. In a complete mass balance, it is necessary to take account of the temporal changes in the depth of the mixed layer and variations in O₂ concentration and δ¹⁸O. A simplified mass balance is possible, if quasi-steady state is assumed in the mixed layer. [Figure 6](#) shows the major fluxes used in a simple steady-state calculation, neglecting O₂ exchange across the oceanic pycnocline. For such a simplified case, the following mass balance equation can be written (see derivation details in [Luz and Barkan, 2011](#)):

$${}^{18}\epsilon_{\text{up}} = \frac{[(C/C_{\text{eq}} - 1)R_w + N/G(R_{\text{eq}} - C/C_{\text{eq}}R_{\text{dis}})]}{[(C/C_{\text{eq}} - 1)(1 - N/G)R_{\text{dis}}]} - 1 \quad [12]$$

where *N* is the net O₂ production, *G* is the gross O₂ production, *U* is the O₂ uptake by various respiratory mechanisms, *C* and *C*_{eq} are measured and equilibrium O₂ concentrations, respectively, and *R*_w, *R*_{eq}, and *R*_{dis} stand for ¹⁸O¹⁶O/¹⁶O¹⁶O ratios in O₂ derived from complete conversion of water to O₂ gas, in O₂ at equilibrium with the atmosphere, and in O₂ dissolved in the ocean, respectively.

Such a calculation was first applied by [Quay et al. \(1993\)](#) who derived ¹⁸ε_{up} of 22 ± 6‰ for the subarctic Pacific. The relatively large error in their result was due to uncertainties in the magnitudes of the various fluxes. For a case of much larger biological fluxes (*G* and *N*) and small gas exchange in Lake Kinneret, the errors were smaller and [Luz et al. \(2002\)](#) reported ¹⁸ε_{up} ranging 22–25‰. Assuming that there was no fractionation in photosynthesis, they concluded that the larger estimates may be the result of engagement of the AOX pathway which is known to strongly discriminate against ¹⁸O (see [Section 5.14.3.2.1](#)).

[Hendricks et al. \(2004\)](#) proposed a method based on similar fluxes as in eqn [12], but with smaller uncertainties, by using measured δO₂/Ar, δ¹⁷O, and δ¹⁸O of dissolved O₂. In this method, *G* is estimated from δ¹⁷O and δ¹⁸O (see [Section 5.14.4.1](#)) and *N* from δO₂/Ar. [Hendricks et al. \(2004, 2005\)](#) applied this method in their studies of the Southern Ocean and the equatorial Pacific and estimated the combined respiratory

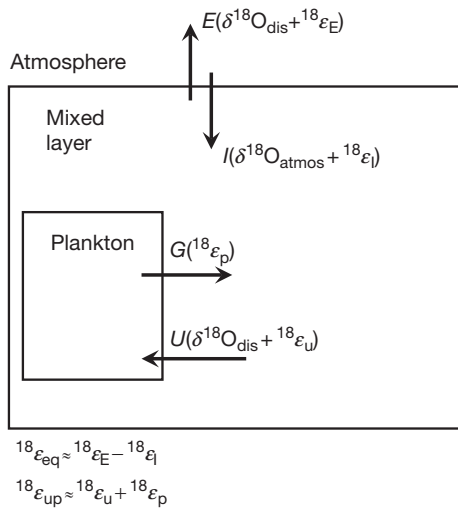


Figure 6 Schematic illustration of O₂ isotope effects in the oceanic mixed layer. ${}^{18}\epsilon_{\text{eq}}$ is the $\delta^{18}\text{O}$ difference of dissolved O₂ from atmospheric O₂ due to small equilibrium isotope effect (${}^{18}\epsilon_{\text{I}}$) in O₂ dissolution. ${}^{18}\epsilon_{\text{eq}}$ is a combined isotope effect that includes ${}^{18}\text{O}$ enrichment in photosynthetic production from water (${}^{18}\epsilon_{\text{P}}$) and further enrichment due to respiratory uptake (${}^{18}\epsilon_{\text{U}}$).

photosynthesis isotope effect as 22 ± 1 and $21 \pm 2\%$, respectively. In the same way, [Luz and Barkan \(2011\)](#) applied the method in various marine locations, including two which were sampled during spring blooms. In addition, they applied the method to the largest available database on oceanic $\delta\text{O}_2/\text{Ar}$, $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ ([Reuer et al., 2007](#)). Recently, [Prokopenko et al. \(2011\)](#) suggested an improved equation for calculating N/G . Using their equation, ${}^{18}\epsilon_{\text{up}}$ values were recalculated for various ocean locales ([Table 5](#)).

From [Table 5](#), it is clear that the magnitudes of ${}^{18}\epsilon_{\text{up}}$ are not randomly distributed in the ocean. Stronger fractionations were observed in oceanic sites which were at their spring blooms or just about to start blooming (S. Ocean and Celtic Sea). However, there is no correspondence between photosynthetic accumulation of O₂, as indicated by $\delta\text{O}_2/\text{Ar}$ in the Celtic Sea, and ${}^{18}\epsilon_{\text{up}}$. Likewise, there is no correspondence between iron addition and ${}^{18}\epsilon_{\text{up}}$ in the Eisenex experiment in the Southern Ocean. This suggests that before the start of the bloom, the phytoplankton communities were already primed for strong fractionation during O₂ uptake or photosynthesis.

The database of [Reuer et al. \(2007\)](#) that was used for deriving the ${}^{18}\epsilon_{\text{up}}$ for the Southern Ocean is by far the most extensive one available for the world ocean and thus deserves special attention. In [Figure 7](#), all the individual calculated ${}^{18}\epsilon_{\text{up}}$ values were plotted. As can be seen, there is a slight tendency for larger ${}^{18}\epsilon_{\text{up}}$ with increasing net O₂ production. This, in turn, is consistent with the observation of strong fractionation during blooms ([Table 5](#)). Because of the broad area covered by this database, the average ${}^{18}\epsilon_{\text{up}}$ of $25.2 \pm 1.9\%$ can be taken as a representative of the entire ocean. If ${}^{18}\epsilon_{\text{U}}$ of 19.7% for average marine respiration below the photic zone applies also to respiration in the surface ocean, then it can be calculated that ${}^{18}\epsilon_{\text{P}} \approx 5\%$ for mean oceanic photosynthetic enrichment with respect to seawater.

Table 5 Isotopic effects in the mixed layer

	ϵ_{up} (‰)
Southern Ocean average ^a	25.2 ± 1.9
Eisenex experiment ^b	
with Fe addition	26.9 ± 0.3
without Fe addition	27.1 ± 0.9
Celtic Sea ^b	
bloom	26.7 ± 1.3
prebloom	25.3 ± 0.6
N. Atlantic near Bermuda ^b	23.7 ± 1.8
Red Sea near Eilat ^b	21.9 ± 1.0

^aData from [Reuer et al. \(2007\)](#)

^bData from [Luz and Barkan, 2011](#); due to large computation errors, data points for which ${}^{17}\Delta$ and/or N/G were less than 20 per meg and 0.01, respectively, were not used in the calculation of ϵ_{up} .

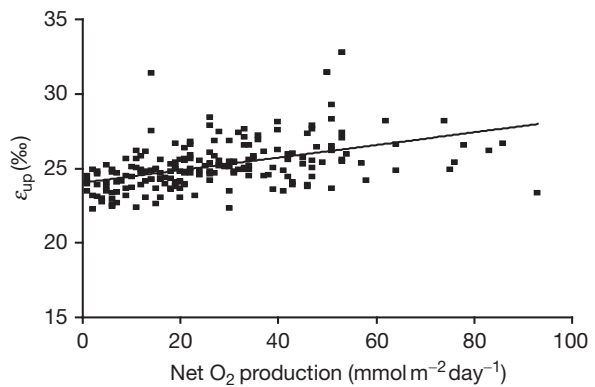


Figure 7 Combined isotopic effect of respiration and photosynthesis in the surface ocean.

5.14.3.2.3 Hydrologic processes

The substrate of O₂ produced by terrestrial photosynthesis is leaf water and therefore its $\delta^{18}\text{O}$ is an essential factor of the DE. In turn, $\delta^{18}\text{O}$ of leaf water is controlled by two main mechanisms. The first one involves processes that determine $\delta^{18}\text{O}$ of meteoric water, and the second involves processes affecting $\delta^{18}\text{O}$ of leaf water in plant transpiration.

Isotopic variations in meteoric water have been a target for much investigation starting from the very early days of stable isotope geochemistry. In general, meteoric waters have $\delta^{18}\text{O}$ close to that of seawater and become depleted at higher latitudes or high topographic elevations (e.g., [Bowen and Revenaugh, 2003](#); [Dansgaard, 1964](#)). Accordingly, water used by plants and trees has a wide range of $\delta^{18}\text{O}$ values. These, in turn, increase by leaf transpiration into dry air. The effect of air humidity on $\delta^{18}\text{O}$ of leaf water was first studied by [Gonfiantini et al. \(1965\)](#). The cause of the ${}^{18}\text{O}$ enrichment in leaf water is similar to that of evaporation from small water bodies. It depends on the liquid–water isotopic equilibrium and on fractionation during the diffusion of water vapor in air. The effect of the latter becomes stronger as air relative humidity goes low (e.g., [Gat, 1996](#)). [Dongmann et al. \(1972\)](#) explained that because leaf water is the substrate for terrestrial photosynthetic O₂, the transpiration enrichment mechanism enhances the

magnitude of the DE. Subsequently, Dongmann (1974) estimated the contribution of land photosynthesis to the DE as 8‰. This estimate, however, has been subject to much debate in subsequent publications.

While the general trends of $\delta^{18}\text{O}$ in meteoric water and the effects of transpiration on $\delta^{18}\text{O}$ of leaf water are well established, the detailed distribution of leaf water and its integrated effect on $\delta^{18}\text{O}$ of atmospheric O₂ are very complex. Yet, for the explanation of the isotopic composition of atmospheric O₂, it is the integrated effect of leaf water that is of much importance. An ingenious way for obtaining such an integrated estimate was proposed by Farquhar et al. (1993) from analyses of $\delta^{18}\text{O}$ of atmospheric CO₂. It is known that O₂ gas and oxygen in water vapor do not exchange isotopes. In turn, two main reservoirs of water influence $\delta^{18}\text{O}$ of atmospheric CO₂ by isotopic exchange, seawater and leaf water. Despite its very small size in comparison to the ocean, leaf water has a major effect. This is because CO₂-H₂O equilibration, through which the isotopic exchange occurs, is very rapid in leaves.

Since the pioneering study of Mills and Urey (1940), it has been known that oxygen isotope exchange takes place between CO₂ and H₂O through the formation of carbonic acid H₂CO₃. However, the reaction of carbon dioxide hydration and dehydration is not instantaneous, and the time for oxygen isotope equilibration between CO₂ and liquid H₂O is on the order of minutes or longer. On the other hand, Roughton (1935) showed that hydration-dehydration is very rapid when it is catalyzed by the enzyme carbonic anhydrase (CA). CA is ubiquitous in leaves and can potentially bring about equilibrium between H₂O and CO₂ during the short transit time of carbon dioxide in photosynthesizing leaves. Indeed, Francey and Tans (1987) found a clear signature of leaf water $\delta^{18}\text{O}$ in atmospheric CO₂. Until their study, it was generally believed that the isotopic composition of atmospheric carbon dioxide was dominated by isotopic exchange with the ocean because $\delta^{18}\text{O}$ of CO₂ was close to the equilibrium value with seawater (Bottinga and Craig, 1969; Keeling, 1961). However, the observations of Francey and Tans (1987) showed that this could not be the case because $\delta^{18}\text{O}$ of CO₂ in high northern latitudes was too low to be explained by equilibration with cold oceanic water. They suggested that their observations were the result of rapid exchange with leaf water.

Farquhar et al. (1993) took the atmospheric observations a step further and constructed a model which included atmospheric mixing times and isotopic exchange between seawater, soil moisture, and leaf water with atmospheric CO₂. A key element in this model is the very rapid CA-catalyzed exchange in leaves in comparison to slow exchange with seawater and soil moisture. Because the isotopic exchange with leaves takes place during photosynthesis inside chloroplasts, $\delta^{18}\text{O}$ of chloroplast water affects both atmospheric CO₂ and newly produced O₂. From model calculations, they estimated that 4.4‰ of Earth's DE would result from ^{18}O enrichment of leaf water. This estimate fell several per mil short of the 8‰ that was needed to account for the DE in the global budgets of Dongmann (1974) and Bender et al. (1994).

An explanation for the discrepancy was suggested by Gillon and Yakir (2001), who conducted a survey of CA activity and CO₂ hydration-dehydration rates in species representing trees,

shrubs, herbs, and grasses. From these measurements, they calculated the degree of isotopic equilibration for the various plant groups and found that it was low for C₄ grasses. Overall, they reported that global CO₂-H₂O equilibration was not complete and estimated an average equilibration value of ~80%. From this estimate, they suggested that ^{18}O enrichment in leaf water must be larger than in Farquhar et al.'s (1993) calculation and further suggested that ^{18}O enrichment in leaf water accounts for 6–8‰ of the DE.

An alternative way for obtaining a global average of ^{18}O enrichment in leaf water is from direct observations on their $\delta^{18}\text{O}$ values. Yet, because of the high variability on $\delta^{18}\text{O}$ of leaf water on both temporal and geographical scales, obtaining such a global average is difficult and there is considerable uncertainty in the obtained value. Nevertheless, such averages have been derived and a recent one, based on global models, of 6.5‰ has been published by West et al. (2008). Considering the uncertainty in this estimate ($\pm 2.1\text{‰}$), it is similar to the $\delta^{18}\text{O}$ of leaf water suggested by Gillon and Yakir (2001).

5.14.3.2.4 Stratospheric photochemical reactions

Exchange of oxygen isotopes between atmospheric O₂ and CO₂ was suggested by Vinogradov et al. (1959) as a major mechanism for the explanation of the DE. Their argument started from a well-known fact that due to isotopic equilibration between CO₂ and H₂O, atmospheric CO₂ is enriched in ^{18}O by more than 40‰ with respect to seawater. They further suggested that in the stratosphere, UV radiation should cause some transfer of this ^{18}O enriched oxygen from CO₂ to O₂. Yet, they did not have any data to support their hypothesis.

Because of temperature stratification, stratospheric air does not readily exchange with the troposphere and it resides for several years above the tropopause (Appenzeller et al., 1996; Holton, 1990). Therefore, if Vinogradov et al. (1959) were right, stratospheric CO₂ would have been depleted in ^{18}O with respect to that gas in the troposphere. But, in fact, observations show that stratospheric CO₂ is highly enriched in ^{18}O (e.g., Gamo et al., 1989; Lammerzahl et al., 2002). Thus, the Vinogradov et al. (1959) explanation must be ruled out. However, the high ^{18}O enrichment of stratospheric CO₂ is interesting in its own right and will be discussed in Section 5.14.4.1. In short, it is caused by mass-independent transfer of ^{18}O and ^{17}O from O₂ to O₃ and then to CO₂ (e.g., Thieme, 1999).

Because the ultimate source of the heavy oxygen isotopes in stratospheric CO₂ is atmospheric O₂, it is expected that as CO₂ gets enriched, O₂ would become depleted in ^{17}O and ^{18}O . From stratospheric mass balance of oxygen isotopes, Bender et al. (1994) calculated a decrease of about 0.4‰ in both $\delta^{18}\text{O}$ and $\delta^{17}\text{O}$ of atmospheric O₂. Later on, this mass-independent depletion was experimentally confirmed by Luz et al. (1999), and from their data, about 0.3‰ lowering of $\delta^{18}\text{O}$ of atmospheric O₂ by stratospheric photochemistry was estimated. Clearly, the effect of stratospheric photochemistry on atmospheric $\delta^{18}\text{O}$ is negligible in comparison to the effect of respiration, photosynthesis, or hydrology. It is noted, however, that, as discussed in Section 5.14.4.1, this small mass-independent change in air O₂ can be used for estimating the ratio of gross to net O₂ production in the surface ocean. This ratio, in turn, is very useful for estimating oxygen isotope

effects of respiration and photosynthesis in the ocean (see Section 5.14.3.2.2).

5.14.3.3 Global Budgets of Processes Influencing the Dole Effect

Efforts to explain the DE started after its discovery in 1935 and experiments were set up to find out whether it could result from ¹⁸O enrichment in photosynthesis. The first results (Dole and Jenks, 1944) showed some 5‰ enrichment in photosynthetic O₂ with respect to the substrate water. These experiments were followed by the respiration experiments of Lane and Dole (1956) that showed substantial discrimination against ¹⁸O in O₂ consumed by respiration and ¹⁸O enrichment in the remaining O₂. Dole (1965) reviewed the available information and concluded that, similarly to the elemental O₂ cycle, there was a natural cycle of oxygen isotopes. In his view, when both cycles operated in steady state between production and consumption, atmospheric O₂ became somewhat enriched in ¹⁸O in the photosynthetic production step and further enrichment (~16–18‰) was attributed to respiratory consumption. Given that at the time Dole wrote his review there were no accurate data on the magnitude of the DE, these two enrichments combined seemed sufficient for explaining the measured effect. However, as discussed in Section 5.14.3.2.2, other researchers of his time rejected the idea that O₂ produced by photosynthesis was enriched in ¹⁸O with respect to the water substrate. This left the origin of the DE unexplained.

A solution was proposed by Dongmann and his colleagues (Dongmann, 1974; Dongmann et al., 1972) who suggested that the gap could be closed by terrestrial photosynthesis, the substrate of which is leaf water. Based on their measurements, Dongmann and coworkers proposed substantial 8‰ ¹⁸O enrichment in leaf water, sufficient for closing the gap.

The idea of Dongmann was further developed by Bender et al. (1994) in a more quantitative way. Bender's group utilized additional information on the biogeochemical mechanisms involved in the global oxygen cycle, as well as more precise experimental data on oxygen isotopic fractionation in various processes. They divided the global DE (DE_{glob}) into two main components: (1) terrestrial DE (DE_{terr}), which is defined as the effect that would result from O₂ exchange by photosynthesis and respiration on land alone, and (2) marine DE (DE_{mar}), which is defined analogously and represents the ocean alone. Following this concept, DE_{glob} is expressed as

$$DE_{\text{glob}} = DE_{\text{mar}} \times f_{\text{mar}} + DE_{\text{terr}} \times f_{\text{terr}} - \epsilon_{\text{strat}} \quad [13]$$

where f_{mar} and f_{terr} stand for the fractions of marine and terrestrial gross O₂ production, respectively, and ϵ_{strat} represents lowering of $\delta^{18}\text{O}$ of atmospheric O₂ by stratospheric photochemistry. The values of DE_{terr} and DE_{mar} were estimated as 22.4 and 18.9‰, respectively, and then, taking f_{mar} and f_{terr} as 0.37 and 0.63, respectively, the global DE was calculated as 20.8‰.

By means of the same logic, but with a 3D numerical model, Hoffmann et al. (2004) obtained the following: DE_{terr} = 25.9‰, DE_{mar} = 17.0‰, and DE_{glob} = 22.9‰. The greater DE_{terr} was due to their use of stronger ¹⁸O enrichment in leaf water (~6 instead of 4.4‰). The smaller DE_{mar} was the result of the larger fraction of marine respiration assigned to O₂ uptake at intermediate

oceanic depths (0.2 instead of 0.05 in Bender et al., 1994), where the overall fractionation is known to be weaker (Bender, 1990; Kroopnick and Craig, 1976; Levine et al., 2009; Rakestraw et al., 1951). In both cases, the calculated magnitude of DE_{glob} strongly depended on the sizes of f_{mar} and f_{terr} . From this dependence, it was expected that variations in the $f_{\text{mar}}/f_{\text{terr}}$ ratio would result in changing magnitudes of DE_{glob}.

In the preceding sections, the current understanding of the magnitude of ¹⁸O/¹⁶O fractionation of the major mechanisms that control DE_{glob} has been reviewed. Among these, there are two main observations that were not available to either Bender et al. (1994) or Hoffmann et al. (2004). These are significant fractionation in marine photosynthesis and attenuated fractionation in soil respiration. In the succeeding text, these are included in calculations of DE_{mar} and DE_{terr}.

Following Bender et al. (1994), the marine DE can be expressed as

$$DE_{\text{mar}} = f_{\text{surface}} DE_{\text{surface}} + f_{\text{deep}} DE_{\text{deep}} \quad [14]$$

where f_{surface} and f_{deep} are the fractions of O₂ uptake in the surface and deep ocean, respectively. DE_{surface} is the difference between $\delta^{18}\text{O}$ of seawater and atmospheric O₂ that would result in a hypothetical case in which the only components are the surface ocean and the atmosphere. DE_{deep} is defined similarly for a case where the deep ocean and the atmosphere are the only components.

DE_{surface} is simply calculated by subtraction of the isotopic effect of the equilibrium isotopic fractionation of dissolved O₂ with respect to atmospheric O₂ ($^{18}\epsilon_{\text{eq}} \approx 0.75\text{‰}$) from $^{18}\epsilon_{\text{up}}$ ($25.2 \pm 1.9\text{‰}$) as 24.4‰. On the other hand, the in situ fractionation in the deep isolated parts of the ocean (ϵ_{u}) is not fully expressed in the atmosphere, but DE_{deep} can be calculated from eqn [15] as (see derivation in Luz and Barkan, 2011)

$$DE_{\text{deep}} = \frac{10^3 [10^{-3}\epsilon_{\text{p}}(1-f) + f^{\alpha}(1+10^{-3}\epsilon_{\text{eq}}) - f - 10^{-3}\epsilon_{\text{eq}}]}{[1-f^{\alpha}]} \quad [15]$$

where $\alpha = 1/(10^{-3}\epsilon_{\text{u}} + 1)$ and f is the fraction of initial O₂ that remains after respiration took place in an isolated parcel of water in the deep sea. The rest of the parameters are as defined earlier.

From the ratio of respiration to the size of O₂ inventories in the water column just below the illuminated upper ocean where most marine respiration takes place, f can be estimated. At steady state, respiration is equal to the flux of organic aggregates produced in the photic zone, which, in turn, equals net production. Net production was estimated from $\delta\text{O}_2/\text{Ar}$ and wind speed measurements (e.g., Hendricks et al., 2004) from the vast Southern Ocean database of Reuer et al. (2007) as about 2.7 mol m⁻². This amounts to ~6% of the O₂ inventory in the 100–300 m interval (~44 mol m⁻²). A similar calculation based on data from the BATS station near Bermuda (Luz and Barkan, 2009) yields ~9%. Thus, f can be taken as 0.91–0.94. Then, substituting $\epsilon_{\text{p}} = 5\text{‰}$, $\alpha = 1/(19.7 \times 10^{-3} + 1)$, and $\epsilon_{\text{eq}} = 0.75\text{‰}$ into eqn [15], DE_{deep} is obtained as 23.1–23.4‰.

In order to calculate DE_{mar}, the ratio of respiration below the photic zone to respiration in the photic zone needs to be known. Bender et al. (1994) assumed the fraction consumed

below the photic zone as 5% of the total, but they did not argue their choice. In principle, the given ratio corresponds to the N/G ratio, because on several years scale total marine respiration equals gross production (G) and the deep respiration rate equals net production (N). As discussed in Luz and Barkan (2009), correct N/G values should be determined from in situ measurements of O₂/Ar ratio, ¹⁷O/¹⁶O, and ¹⁸O/¹⁶O of dissolved O₂ (see Section 5.14.3.2.2). Such N/G estimates were obtained for the subtropical Pacific (average of 0.1; Juranek and Quay, 2005), equatorial Pacific (average of 0.06; Hendricks et al., 2005), Southern Ocean (average of 0.13; Reuer et al., 2007), and the Atlantic Ocean near Bermuda (average of 0.14; Luz and Barkan, 2009). From these ratios, a representative value for N/G ratio in the ocean is estimated as 0.11. So based on ocean-wide information, the fraction of respiration in the deep sea is about 11%.

Finally, using the estimates of DE_{surface}, DE_{deep} and the fraction of respiration in the deep sea, DE_{mar} is calculated as 23.3‰ (Table 6). Obviously, there is considerable uncertainty in this value, and its main source is the scatter of DE_{surface} in the database ($\pm 1.9\%$).

The terrestrial DE can be expressed as

$$DE_{\text{terr}} = \delta^{18}\text{O}_{\text{leaf-water}} + \varepsilon_{\text{u-leaves+soil}} - \varepsilon_{\text{eq}} \quad [16]$$

As discussed in Section 5.14.3.2.3, the global weighted average of $\delta^{18}\text{O}$ of leaf water is 6.5‰. From Section 5.14.3.2.1, the isotopic effects in leaf O₂ uptake and in soil respiration are 19.2 and 15.8‰, respectively, and the weighted average of the overall terrestrial respiratory fractionation ($\varepsilon_{\text{u-leaves+soil}}$) is 17.7‰. Substituting these values in eqn [16] and subtracting 0.75‰ to account for solubility equilibration between air O₂ and oxygen dissolved in leaf water, DE_{terr} can be estimated as 23.5 \pm 2.3‰ (Table 6).

By using eqn [16] with values of DE_{mar}, DE_{terr}, $\varepsilon_{\text{strat}}$, f_{mar} , and f_{terr} from Table 6, a global estimate of the DE can be calculated as 23.5 \pm 2.5‰. Given the calculation uncertainty, this value is equal to the measured DE (23.88‰) and also equal to both calculations of DE_{mar} (24.3 \pm 1.1‰) and DE_{terr} (23.5 \pm 2.3‰). Clearly, despite uncertainties in the sizes of the

fractions of terrestrial and marine productivities in Table 6, the calculated DE_{glob} is not sensitive to the ratio of f_{mar} to f_{terr} . Therefore, variations in $f_{\text{mar}}/f_{\text{terr}}$ cannot explain the record of past variations in the DE, and the cause has to be sought elsewhere as will be discussed in the next section.

5.14.3.4 Temporal Variations in the Dole Effect

The past history of atmospheric gases can be reconstructed from air bubbles trapped in ice cores, revealing useful information about biogeochemistry and climate. The pioneering ice core study of the isotopes of O₂ by Michael Bender (Bender et al., 1985) showed an enriched value of $\delta^{18}\text{O}_{\text{atm}}$ in the past glacial period, roughly as expected from the +1‰ higher seawater $\delta^{18}\text{O}$ at that time. Past seawater $\delta^{18}\text{O}$ is inferred from the $\delta^{18}\text{O}$ of benthic foraminifera calcite with appropriate corrections for temperature effects (Waelbroeck et al., 2002), and past values of the DE can be estimated from the difference between the measured ice core $\delta^{18}\text{O}_{\text{atm}}$ and the reconstructed seawater $\delta^{18}\text{O}$.

Subsequently, refinements to the precision of the ice core gas technique and longer ice core records showed that the DE also changes over time (Bender et al., 1994; Malaize et al., 1999). These changes seemed to have a strong precessional (~ 23 ky) periodicity and were thought to be partly due to low-latitude hydrologic cycle variations, stemming both from the isotopic composition of the rainfall and the relative humidity (which controls the amount of evaporative enrichment of chloroplast water; Farquhar et al., 2007). Heavy convective rains and high humidity both tend to decrease the $\delta^{18}\text{O}$ of photosynthetic substrate water. Mélières et al. (1997) pointed out that a minimum in the DE occurred at 175 ka coincident with a strong North African monsoon and strong precession-driven June insolation. This particular case was instructive because it occurred during a glacial period when boreal and temperate ecosystems were quiescent (Masson et al., 2000).

It now appears likely that the DE is primarily modulated over time by variations in the strength of the low-latitude northern hemisphere monsoon systems (Landais et al., 2010),

Table 6 Calculation of the Dole effect

	$\delta^{18}\text{O}$ (vs. VSMOW), $^{18}\varepsilon$ or DE _x (all in ‰)	Fraction of global O ₂ production	References
Marine Dole effect			
DE _{surface}	25.2 \pm 1.9	0.33	This chapter
DE _{deep}	23.2 \pm 0.5	0.04	Luz and Barkan (2011)
DE _{mar}	24.3 ^a \pm 2.0	0.37	
Terrestrial Dole effect			
$\delta^{18}\text{O}$ of photosynthetic O ₂	6.5 \pm 2.1		West et al. (2008)
Total respiratory fractionation	17.7 \pm 1.0		Landais et al. (2007a)
ε_{eq} of O ₂ leaf water relative to air	0.75 ^b \pm 0.05		Benson and Krause (1984)
DE _{terr}	23.5 ^c \pm 2.3	0.63	
Stratospheric isotopic effect ($\varepsilon_{\text{strat}}$)	0.3 ^d \pm 0.1		Luz et al. (1999)
Calculated global Dole effect	23.5 ^e \pm 2.5		
Measured global Dole effect	23.88 \pm 0.02		Barkan and Luz (2005)

^a24.3 = (0.33*24.2 + 0.04*23.2)/0.37.

^b0.75 is an average for possible temperature-dependent values ranging from 0.7 to 0.8‰.

^c23.5 = 6.5 + 17.7 - 0.75.

^dCalculated from eqn [10] and $^{17}\Delta$ of 0.166‰ (Section 5.14.2) and assuming in the first approximation $\delta^{17}\text{O} = \delta^{18}\text{O}$, the authors obtain 0.166 = $\delta^{18}\text{O} - 0.516 * \delta^{18}\text{O}$, and $\delta^{18}\text{O} \sim 0.3\%$.

^e23.5 = 24.3*0.37 + 23.5*0.63 - 0.3.

with strong summer (wet phase) monsoons causing a smaller DE (Figure 8). During the Holocene, further support for this monsoon hypothesis is provided by the observed excellent correlation ($R=0.95$) of the Chinese cave calcite $\delta^{18}\text{O}$ with the inferred fractionation that gives rise to the DE (Severinghaus et al., 2009). As the monsoon strength has long been known to be affected by northern hemisphere low-latitude summer insolation (Kutzbach, 1981), it comes as no surprise that the DE has a strong inverse correlation to 30° N summer solstice insolation.

Perhaps more surprising is that the DE also varies with the abrupt millennial-scale oscillations known as Heinrich and Dansgaard–Oeschger events (Figure 9; Landais et al., 2007b; Severinghaus et al., 2009), as do the Chinese cave records that are thought to be related to monsoon strength (Cheng et al., 2009; Wang et al., 2001). The abrupt monsoon variations may be related to rapid southward shifts of tropical rainfall belts and the Intertropical Convergence Zone, in response to abrupt increases in northern high-latitude winter sea ice cover (Cheng et al., 2009; Chiang and Bitz, 2005; Wang et al., 2007). However, some caution is required as the exact cause of the Chinese cave isotopic variations is still not understood, and it is possible that they do not solely record summer monsoon rainfall intensity (Clemens et al., 2010).

A third factor that probably contributes to the link between the DE and monsoons, more recently identified

(Angert et al., 2003a; Section 5.14.3.2.1), is the weak respiratory fractionation characteristic of wet tropical soils. This study also found strong respiratory fractionation in boreal soils, which by itself should cause a positive contribution to the DE. However, boreal precipitation is quite depleted in ^{18}O , causing the opposite effect.

In summary, the low $\delta^{18}\text{O}$ of precipitation in heavy summer monsoon rains, the weak evaporative enrichment of leaf water in conditions of high relative humidity, and the weak respiratory fractionation probably all contribute to the observed monsoon–DE linkage, although it is difficult to quantify the relative importance of each factor. The constructive interference of respiratory and hydrologic fractionation at low latitudes, versus their destructive interference at high latitudes, may operate as a sort of ‘monsoon rectifier effect’ that causes low-latitude hydrology to dominate the variability in the DE (Luz and Barkan, 2011; Severinghaus et al., 2009).

On longer timescales (using the full 800 ky record now available from ice cores; Figure 8), there appears to be a broad 100 ky peak in the power spectrum of DE variations, and the DE reaches a maximum every time that a glacial termination occurs (Landais et al., 2010). This likely stems from the fact that the Asian and African monsoons are extremely weak during the times when the northern ice sheets are melting, probably due to winter sea ice cover on the North Atlantic Ocean (Cheng et al., 2009). On even longer timescales, there

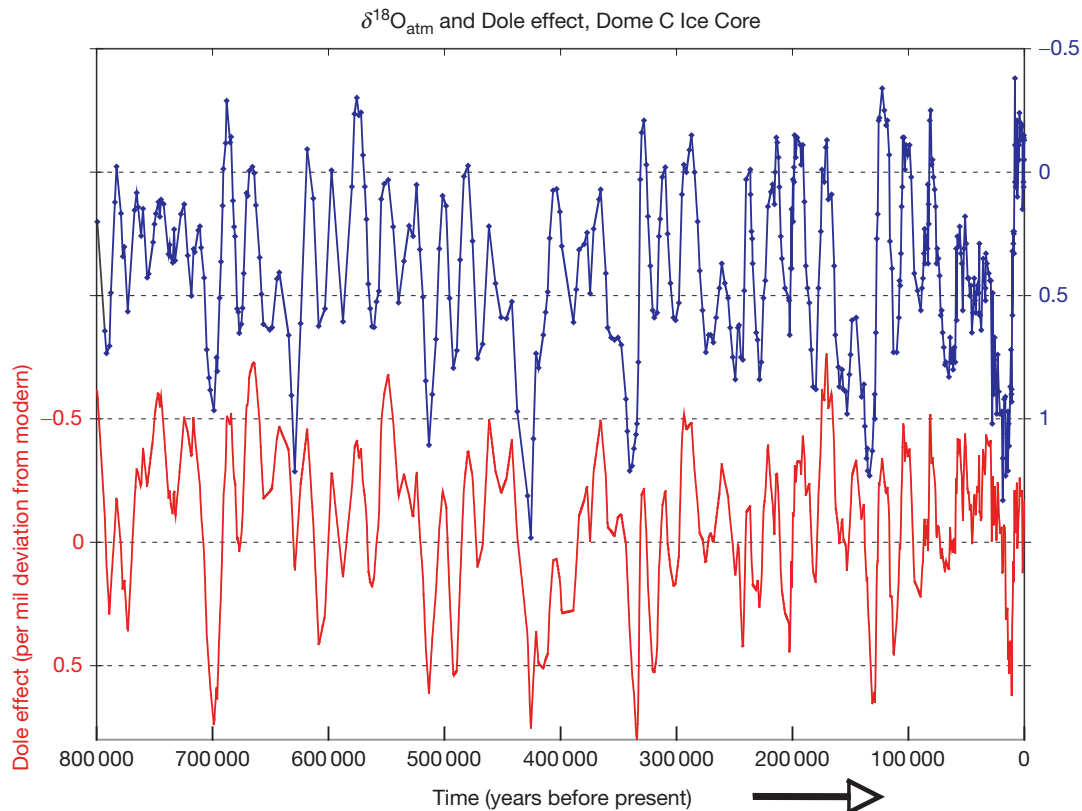


Figure 8 Past variations in the Dole effect (DE) reconstructed from measurements of $\delta^{18}\text{O}_{\text{atm}}$ in the EPICA Dome C ice core from Antarctica (Landais et al., 2010). Values are plotted using the modern atmosphere as a reference, with inverted axes (following convention) and with time proceeding toward the right. Note the strong precessional (~ 23 ky) periodicity in the record.

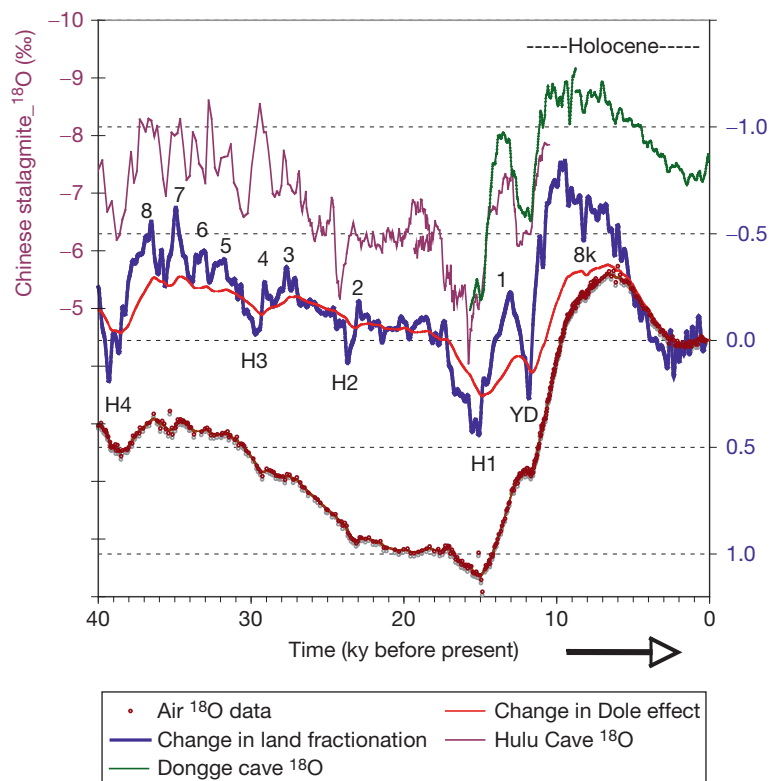


Figure 9 Observed variations in the Dole effect (DE) over the past 40 000 years, from the Siple Dome ice core, Antarctica (Severinghaus et al., 2009). The derivative of the observed DE curve was used to infer the change in land fractionation ($\Delta\epsilon_{\text{LAND}}$), effectively the main driver of changes in the DE (shown in the purple curve). Note the strong similarity with the Chinese cave records, which are thought to reflect monsoon variations (Dykoski et al., 2005; Wang et al., 2001, 2005). Abrupt climate events including Dansgaard–Oeschger events (numbered 1–8), Heinrich events (H1–H4), Younger Dryas (YD), and 8.2 ka event (8k) are shown.

seems to be a tendency for the DE to be larger during times of small eccentricity, such as the interval between 350 and 450 ka (Landais et al., 2010). Interestingly, power spectra show that there is virtually no signal of obliquity in the DE; rather, it is dominated by the precession signal (Landais et al., 2010).

One potential clue about the origin of DE variations comes from direct measurements over the past several decades of the $\delta^{18}\text{O}$ of atmospheric CO₂ (Welp et al., 2011). This tracer shares in common with the DE a strong dependence on the $\delta^{18}\text{O}$ of chloroplast water because atmospheric CO₂ diffuses rapidly in and out of the stomata of terrestrial plants, rapidly exchanging its isotopes with chloroplast water thanks to the enzyme CA (Farquhar et al., 2007). However, CO₂ has an atmospheric turnover time of the order of 0.5 year in contrast to the ~ 1000 year turnover time of O₂, making CO₂ a much more sensitive tracer of chloroplast water variations. The CO₂ isotope record shows clearly that during Indonesian drought events, such as accompany El Niño events, chloroplast water becomes isotopically heavy. Thus, the implication is that a weakening of the Indonesian Low, or an increase in the frequency of El Niños, should cause an increase in the DE (Welp et al., 2011). The paucity of El Niños in the early Holocene is consistent with the weak DE during that time. Unfortunately, the CO₂ $\delta^{18}\text{O}$ is not well preserved in ice cores due to isotopic exchange with the ice, so it is not possible to reconstruct this tracer back in time.

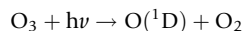
5.14.4 Oxygen-17 and Oxygen-18 in Atmospheric O₂

5.14.4.1 Mass-Independent Fractionation and Biological Normalization

A new field in stable isotope geochemistry was born when Thiemens and Heidenrich (1983) discovered mass-independent isotopic fractionation in the formation of ozone from O₂ gas by electrical discharge. In their experiment, ozone became equally enriched in both ¹⁷O and ¹⁸O with respect to the starting O₂, while in the remaining O₂, both isotopes became equally depleted. This experimental result was followed by experiments in which mass-independent fractionation was demonstrated with UV radiation (e.g., Thiemens and Jackson, 1987). In addition, observations on $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ of stratospheric ozone and CO₂ showed that both were enriched in ¹⁷O and ¹⁸O in a mass-independent way (e.g., Boering et al., 2004; Lammerzahl et al., 2002; Mauersberger, 1987). Efforts to explain the chemical origin of mass-independent fractionation started in 1983, but it has taken some two decades before a rigorous theory started to develop (e.g., Hathorn and Marcus, 2001; Marcus, 2004). The theoretical aspects of mass-independent fractionation are beyond the scope of this chapter, and interested readers are referred to a review by Thiemens (2006).

The current understanding of the stratospheric observations is based on the photochemical model of Yung et al. (1991, 1997), where the anomalous ozone enrichment is transferred to CO₂. In this model, ultraviolet photolysis of ozone in the

stratosphere generates an electronically excited oxygen atom (O(¹D)), which can undergo isotopic exchange with CO₂ through a transition state (CO₃^{*}):



Bender et al. (1994) suggested that because the ultimate source of the oxygen in O₃ and O(¹D) in the stratosphere is the O₂ reservoir, this reservoir must become anomalously depleted as CO₂ becomes anomalously enriched. The stratospheric enrichment in CO₂ is rapidly lost at the Earth's surface by isotope exchange with liquid water in leaves and also in the ocean. In contrast, O₂ does not exchange isotopes with water directly, and the depletion disappears only through the consumption of O₂ by respiration and its replacement by photosynthesis. Yet, the respiratory and photosynthetic fluxes are relatively small in comparison to the stratospheric production, and Bender et al. (1994) suggested that the anomalous depletion should accumulate to a measurable level of about 200 per meg over the 1200 year residence time of atmospheric O₂. Among other factors discussed in the succeeding text, the extent of this anomalous depletion depends on the intensity of UV radiation in the stratosphere, the concentration of atmospheric CO₂, and the rate of global gross O₂ production. Assuming the former factors are known, it should be possible to estimate global gross O₂ production from measurements of the ¹⁷O anomaly in air occluded in polar ice cores, which will be discussed in Section 5.14.4.2.

Luz et al. (1999) set up an experiment for testing Bender's hypothesis, but before this experiment can be discussed, it is necessary to give some background on definitions of mass-independent anomalies, and for details on triple isotope systematic, the reader is referred to Section 5.14.2.2 in this chapter. In studies of mass-independent fractionation in ozone and CO₂ in both stratospheric observations and laboratory experiments, the isotopic standards used are VSMOW or CO₂ in equilibrium with VSMOW, and ¹⁷O anomalies are calculated, assuming that a δ¹⁷O/δ¹⁸O ratio (or λ; see Section 5.14.2.2) of about 0.52 represents normal mass-dependent fractionation. In such studies, the anomaly is calculated as Δ¹⁷O = δ¹⁷O - λ*δ¹⁸O (λ = ~0.52) and its magnitude varies over ranges of several per mil or more. For such large variations, most Earth materials, excluding certain stratospheric gases, can be considered as being affected by mass-dependent fractionations and are thus defined 'normal.' However, the magnitude of the negative mass-independent fractionation in atmospheric O₂ or of oceanic dissolved O₂ is very small and must be treated more rigorously, and both standards and values of λ should be chosen such that they optimally fit the case being studied.

The standard of choice for studies of O₂ gas is present atmospheric O₂ (Luz and Barkan, 2005). Yet, as discussed earlier, this standard is not 'normal' in that it is deficient in ¹⁷O. So ¹⁷Δ as defined in eqn [10] (Section 5.14.2) is zero for atmospheric O₂ and has a positive value for O₂ produced by photosynthesis from natural water. As for values of λ, the best choice is for those of dark respiration which is the most ubiquitous O₂ uptake mechanism on Earth. Depending on whether

the application is for estimation of marine gross O₂ production (see succeeding text) or global gross O₂ production, the optimal values of λ are 0.518 and 0.516, respectively (Luz and Barkan, 2005).

The experiment of Luz et al. (1999) was run in an airtight terrarium that contained *Philodendron* plant, soil, and natural meteoric water. The terrarium was illuminated by artificial light with no UV radiation. Inside the terrarium, O₂ production and consumption occurred by the plant and by soil microorganisms. The experiment started with atmospheric air in the terrarium. As the experiment progressed, the initial atmospheric O₂ was partially consumed by respiration and gradually replaced by newly produced photosynthetic O₂. The isotopic composition of O₂ in the terrarium was monitored and ¹⁷Δ was calculated from δ¹⁷O and δ¹⁸O measurements. The results from one of the experiments are shown in Figure 10 and show an initial gradual increase in ¹⁷Δ, and then the values level off indicating that O₂ in the terrarium was completely replaced with new oxygen of biological origin alone with no effects of UV radiation. As can be seen, the ¹⁷Δ of this bio-O₂ has reached a maximum value designated ¹⁷Δ_{bio} (¹⁷Δ_{max} in various previous publications). In the experiment shown in Figure 10, ¹⁷Δ_{bio} is about 150 per meg. Bio-O₂ is considered 'normal' because it has lost the anomalous photochemical signature that atmospheric O₂ contains.

Depending on the experimental conditions and the water substrate from which photosynthesis makes O₂, there is a range of various 'normal' bio-O₂s. For example, experiments similar to that described earlier have been carried out with marine organisms in aquaria and the average value of ¹⁷Δ_{bio} in these latter experiments was 249 per meg (Barkan and Luz, 2011; Luz and Barkan, 2000). In Figure 11, it is shown schematically how ¹⁷Δ_{bio} is attained in such experiments.

In the open ocean, a simple steady-state situation where respiration and photosynthesis produce bio-O₂ and gas exchange tends to bring ¹⁷Δ of dissolved O₂ (¹⁷Δ_{dis}) to equilibrium with the atmosphere (¹⁷Δ_{eq}) may be considered. In this case, measurements of ¹⁷Δ_{dis} indicate the ratio of gross O₂ production with isotopic composition of ¹⁷Δ_{bio} to air-sea gas exchange introducing dissolved O₂ with isotopic composition of ¹⁷Δ_{eq}. Luz and Barkan (2000) introduced a method in which gross O₂ production could be estimated from measurements of ¹⁷Δ_{dis} and the rate of air-gas exchange. This method has been applied all over the ocean (e.g., Hendricks et al., 2004; Quay et al., 2010; Reuer et al., 2007; Sarma et al., 2008). Recently, Prokopenko et al. (2011) improved the method by obtaining a

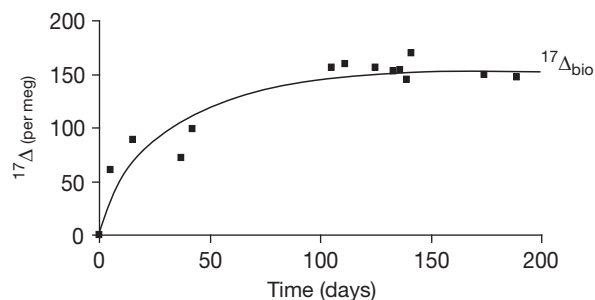


Figure 10 Removal of the mass-independent signature from atmospheric O₂ in a terrarium experiment.

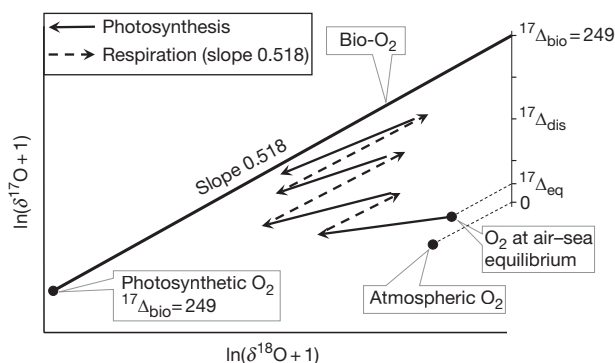


Figure 11 Schematic illustration (not to scale) showing how photosynthesis and respiration affect $\delta^{18}\text{O}$, $\delta^{17}\text{O}$, and $^{17}\Delta$ of dissolved O₂ ($^{17}\Delta_{\text{dis}}$). Newly produced photosynthetic O₂ mixes with the existing O₂, and depending on the extent of photosynthesis, the mixture isotopic composition will move toward that of seawater (solid arrows, which are slightly curved as in [Figure 3](#)) and $^{17}\Delta_{\text{dis}}$ increases. Respiration (dashed arrows) increases both $\delta^{18}\text{O}$ and $\delta^{17}\text{O}$ but does not change $^{17}\Delta_{\text{dis}}$ because $\ln(\delta^{17}\text{O} + 1)/\ln(\delta^{18}\text{O} + 1) = 0.518$. With sufficient photosynthesis and respiration cycles, $^{17}\Delta_{\text{dis}}$ will approach $^{17}\Delta_{\text{bio}}$. From this point on, photosynthesis will decrease both $\delta^{18}\text{O}$ and $\delta^{17}\text{O}$ and respiration will increase them but $^{17}\Delta_{\text{dis}}$ will not change and its value will remain equal to $^{17}\Delta_{\text{bio}}$.

rigorous equation for the steady-state case. This equation has been applied by [Luz and Barkan \(2011b\)](#) for the calculation of $^{18}\epsilon_{\text{up}}$ as described in [Section 5.14.3.2.2](#).

Returning to measurements of $^{17}\Delta_{\text{bio}}$, the different values in aquarium experiments with seawater and terrarium experiments with freshwater cannot be explained from interplay between mass-independent fractionation and mass-dependent effects because in both types of experiments there was no UV radiation. Instead, the cause of the difference in the values of $^{17}\Delta_{\text{bio}}$ is that several mass-dependent processes having somewhat different $^{17}\epsilon/^{18}\epsilon$ ratios come into play. The main controlling factors in aquarium experiments with seawater are the isotopic composition of the water and ordinary respiration with only a small effect of fractionation in photosynthesis ([Barkan and Luz, 2011](#)). On the other hand, in the terrarium experiments, the substrate water came from natural meteoric water that was affected by hydrologic processes of evaporation and precipitation. An extensive survey and discussion of these processes is found in [Luz and Barkan \(2010\)](#). Briefly, the hydrologic cycle starts from oceanic evaporation with an excess of ^{17}O ($^{17}\text{O}_{\text{excess}}$) of about 33 per meg with respect to seawater, but the $^{17}\epsilon/^{18}\epsilon$ ratio in vapor liquid equilibrium, which governs precipitation, is larger than in respiration. Consequently, $^{17}\Delta_{\text{bio}}$ in the terrarium experiment of [Luz and Barkan \(2005\)](#) was smaller (194 per meg) than in seawater (249 per meg). It is noted that there is no one universal value for all meteoric waters, and in general, if terrarium experiments were run with water having lower $\delta^{18}\text{O}$ than in the experiment of [Luz and Barkan \(2005\)](#), the resulting $^{17}\Delta_{\text{bio}}$ would have been smaller than 194 per meg.

In the natural terrestrial biosphere, the situation is more complex because in contrast to the terrarium, air relative humidity is less than 100% and due to transpiration into low humidity air, leaf water becomes enriched in both ^{17}O and ^{18}O in

comparison to meteoric water. [Landais et al. \(2006\)](#) studied the isotopic effects of transpiration and their impact on $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ of atmospheric O₂. These studies showed that the $\ln(^{17}\alpha)/\ln(^{18}\alpha)$ ratio in transpiration is considerably smaller than in global meteoric waters and despite the enrichment in ^{18}O in this process that significantly increases $\delta^{18}\text{O}$ of global leaf water (see [Section 5.14.3.2.3](#)), its $^{17}\text{O}_{\text{excess}}$ goes down with respect to seawater. But it is noted that in calculations of $^{17}\text{O}_{\text{excess}}$ in hydrology, λ is 0.528, while in global calculations of $^{17}\Delta$, it is 0.516, which is quite similar to λ in global transpiration (0.517, [Landais et al., 2007a](#)). Therefore, in O₂ emission from the terrestrial biosphere, $^{17}\Delta_{\text{bio}}$ becomes smaller and smaller as $\delta^{18}\text{O}$ of meteoric water goes low, while transpiration only slightly increases its value. The effects of the hydrologic cycle and leaf transpiration are shown schematically in [Figure 12](#).

A further complication in terrestrial systems arises from CO₂ limitation and the resulting consumption of O₂ by photorespiration. This was demonstrated by [Angert et al. \(2003b\)](#) who showed that $^{17}\Delta_{\text{bio}}$ in a terrarium experiment became about 100 per meg when experiments were carried out at continuous illumination that caused CO₂ concentration to go below 500 ppmv. In contrast, experiments with the same terrarium, plant, soil, and water, but with changing light-dark cycles and CO₂ remaining always above 2000 ppmv, yielded $^{17}\Delta_{\text{bio}}$ of 194 per meg ([Luz and Barkan, 2005](#)). Indeed, direct measurements by *in vitro* experiments of [Helman et al. \(2005; summarized in Table 1\)](#) demonstrated that $^{17}\epsilon/^{18}\epsilon$ ratio of O₂ uptake by glycolate oxidase, an important reaction in photorespiration, was 0.501 and significantly smaller than in dark respiration (0.518).

[Landais et al. \(2007a\)](#) used the information on photorespiration and fractionation in the hydrologic cycle and in transpiration and calculated $^{17}\Delta_{\text{bio}}$ of the terrestrial biosphere as 110 ± 35 per meg with respect to atmospheric O₂. At the time of their publication, information on $^{17}\text{O}_{\text{excess}}$ of meteoric water was not available. It is now known that meteoric waters have $^{17}\text{O}_{\text{excess}}$ of 33 per meg ([Luz and Barkan, 2010](#)). On the other hand, $^{17}\Delta$ of seawater with respect to atmospheric O₂ is 223 and not 249 per meg, which is the value of $^{17}\Delta$ produced by photosynthesis from seawater ([Barkan and Luz, 2011](#)). These effects are included in [Figure 12](#), but an important point is that the $^{17}\text{O}_{\text{excess}}$ of meteoric water is almost entirely compensated by the lower $^{17}\Delta$ of seawater (223 per meg) than that of photosynthetic O₂ (249 per meg).

In [Figure 12](#), it is shown graphically how the various isotopic effects of hydrology and biology are used for deriving the terrestrial and global $^{17}\Delta_{\text{bio}}$ as 117 ± 35 and 166 ± 51 per meg, respectively. Evidently, there is considerable uncertainty in these calculated values, but despite that, they are useful for gaining insight on the increase of global gross O₂ production since the Last Glacial Maximum (LGM).

5.14.4.2 Temporal Variations in $^{17}\text{O}/^{16}\text{O}$ and $^{18}\text{O}/^{16}\text{O}$

The number of studies on past variations of $^{17}\Delta$ in atmospheric O₂ is small. The first measurements were done by [Luz et al. \(1999\)](#) on samples of air extracted from the GISP-2 ice core that had been drilled in Summit, Greenland. This record spans the last 82000 years. A more detailed study on the same core was

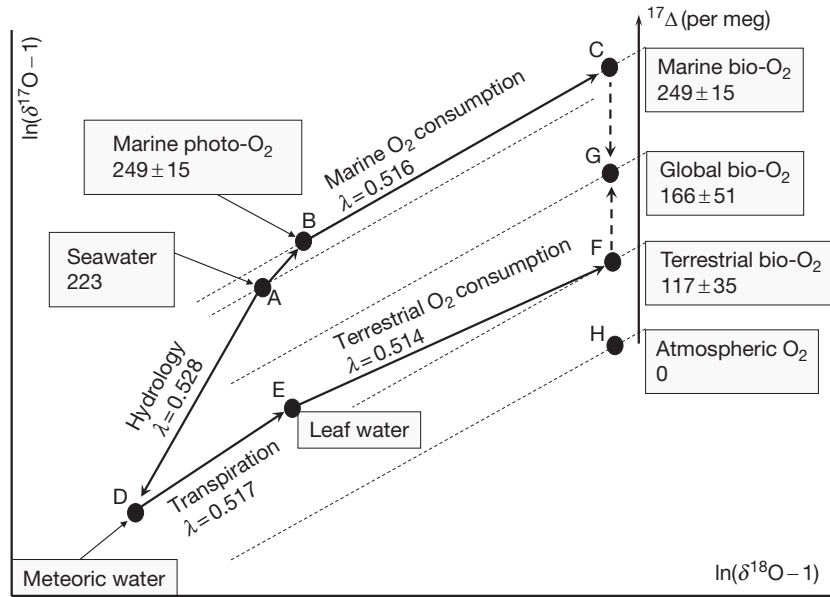


Figure 12 Schematic illustration of processes affecting the $^{17}\Delta$ of atmospheric O₂. In order to show the variations in line slopes and $^{17}\Delta$ values, the $\ln(\delta^{17}\text{O}+1)$ scale is highly exaggerated and is not uniform. Both axes are not drawn to scale. $^{17}\Delta$ values are constant along the dashed lines with slope of 0.516. The slope of arrow AB is 0.521.

published by Blunier et al. (2002) and showed the same trends in $^{17}\Delta$. In general, during the LGM, atmospheric O₂ was about 40 per meg higher in $^{17}\Delta$ than the present atmosphere. During the rest of the last glacial, $^{17}\Delta$ varied around an average value of about 30 per meg. In a much more indirect way, Bao et al. (2008) suggested that $\delta^{17}\text{O}$ and $\delta^{18}\text{O}$ of sulfate in barites and evaporites from the early Cambrian reflect, in part, the isotopic composition of the atmosphere of that time. Cast in terms of $^{17}\Delta$ as defined in this chapter, the maximum deviations they reported were about -700 per meg with respect to the present atmosphere. Because, as they proposed, only 5–20% of the sulfate oxygen was of atmospheric origin, it is calculated that $^{17}\Delta$ of the early Cambrian air reached an extremely low level of $-14\,000$ per meg or -14% with respect to today's air.

Before explaining the origin of these past variations in $^{17}\Delta$, it should be recognized that in addition to biological and photochemical controls, abundances of isotopes and elements of air occluded in ice cores are affected by gravitational and thermal diffusion fractionations (e.g., Severinghaus et al., 1998). The records presented by Luz et al. (1999) and Blunier et al. (2002) were corrected for gravitational fractionations by subtraction of measured $\delta^{15}\text{N}$ of N₂ in the same samples from the measured $\delta^{17}\text{O}$ and by subtraction of $2(\delta^{15}\text{N})$ from the measured $\delta^{18}\text{O}$. The size of the correction is on the order of 10–15 per meg in $^{17}\Delta$.

Following Luz et al. (1999), $^{17}\Delta$ of oxygen in the atmosphere, $^{17}\Delta_{\text{atm}}$, can be expressed as

$$F_{\text{bio}} \times (^{17}\Delta_{\text{bio}} - ^{17}\Delta_{\text{atm}}) = F_{\text{strat}} \times (^{17}\Delta_{\text{strat}} - ^{17}\Delta_{\text{atm}}) \quad [17]$$

where $^{17}\Delta_{\text{strat}}$ and $^{17}\Delta_{\text{bio}}$ are the $^{17}\Delta$ of stratospheric O₂ flux (F_{strat}) and the biological O₂ flux (F_{bio}), respectively. The ratio of the global gross O₂ production between the LGM and the present is calculated from eqn [18] as (Landrais et al., 2007a)

$$\frac{F_{\text{bio,LGM}}}{F_{\text{bio,PST}}} = \frac{F_{\text{strat,LGM}} \times (^{17}\Delta_{\text{strat,LGM}} - ^{17}\Delta_{\text{atm,LGM}})}{F_{\text{strat,PST}} \times (^{17}\Delta_{\text{strat,PST}} - ^{17}\Delta_{\text{atm,PST}})} \times \frac{(^{17}\Delta_{\text{bio,PST}} - ^{17}\Delta_{\text{atm,PST}})}{(^{17}\Delta_{\text{bio,LGM}} - ^{17}\Delta_{\text{atm,LGM}})} \quad [18]$$

By definition, $^{17}\Delta_{\text{atm,PST}} = 0$, and $^{17}\Delta_{\text{atm,LGM}}$ was determined by Blunier et al. (2002) as $+43$ per meg (note that the original value of Blunier et al. (2002) is $+38$ per meg since it was calculated with λ of 0.521 instead of 0.516, which is the optimal value for global calculations).

Luz et al. (1999) and Blunier et al. (2002) assumed that the ratio of the production rates of anomalously depleted O₂ in the stratosphere, $F_{\text{strat}} \times (^{17}\Delta_{\text{strat}} - ^{17}\Delta_{\text{atm}})$, between the LGM and the present is proportional to the ratio of atmospheric CO₂ concentrations between the LGM and the preindustrial Holocene. This assumption implies that the photochemical reactions involved in the production of CO₂ high in $^{17}\Delta$ and O₂ low in $^{17}\Delta$ are first order with respect to CO₂ concentration. Bao et al. (2008) dealt with this assumption and showed that it is reasonable at least for CO₂ concentration up to about 2000 ppm. With this assumption and using CO₂ concentrations of 280 ppmv for the preindustrial Holocene and 190 ppmv for the LGM (Barnola et al., 1987), eqn [18] becomes

$$\frac{F_{\text{bio,LGM}}}{F_{\text{bio,PST}}} = \frac{190}{280} \times \frac{^{17}\Delta_{\text{bio,PST}}}{(^{17}\Delta_{\text{bio,LGM}} - 43)} \quad [19]$$

Clearly, the correction for the different CO₂ concentration between the two periods is large, and indeed, the main driving force for the glacial to postglacial change in $^{17}\Delta$ was the rise in concentration of atmospheric CO₂. Figure 13 shows the $^{17}\Delta$ of atmospheric O₂ versus CO₂ from the data of Luz et al.

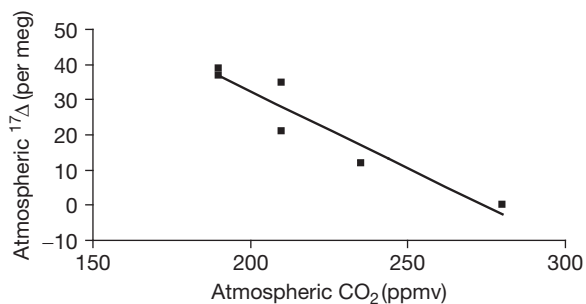


Figure 13 $^{17}\Delta$ versus CO₂ concentration in air extracted from GISP-2 ice core.

(1999) and a similar relation can be shown for Blunier et al. (2002) results.

Before continuing with the ice core record, it is noted that Bao et al. (2008) suggested that extremely high CO₂ concentration (up to 25 000 ppmv) can explain the huge negative $^{17}\Delta$ they measured in early Cambrian sulfates. The same logic was followed by Gehler and Pack (2010) who proposed high atmospheric CO₂, although far less dramatic, in the middle Eocene.

Application of eqn [19] requires that $^{17}\Delta_{\text{bio}}$ be known for both the PST and the LGM. Applying, as did Luz et al. (1999), 155 per meg for both, $F_{\text{bio,LGM}}/F_{\text{bio,PST}}$ is calculated as 0.94. If 249 per meg is applied instead, $F_{\text{bio,LGM}}/F_{\text{bio,PST}}$ becomes 0.82, and if this ratio is calculated with $^{17}\Delta_{\text{bio,LGM}}$ value equal to 214 and the PST value equal to 166 per meg, a ratio of 0.66 is obtained. These simple calculations demonstrate the sensitivity of $F_{\text{bio,LGM}}/F_{\text{bio,PST}}$ to uncertainty in the value of $^{17}\Delta_{\text{bio}}$.

Understanding the importance of reliable $^{17}\Delta_{\text{bio}}$ values for estimates of global O₂ production motivated Landais et al. (2006, 2007a) to study the factors controlling them at present and in the last glacial. Their main results on the present $^{17}\Delta_{\text{bio}}$ and the differences between land and sea have been reviewed in the previous section, and a global $^{17}\Delta_{\text{bio}}$ value of 166 per meg has been derived. However, this value cannot be used for the LGM in which $^{17}\Delta_{\text{bio}}$ must have been larger for three main reasons: (1) in the LGM, $^{17}\Delta$ of seawater was larger, (2) $^{17}\Delta$ of meteoric water used by plants at that time was also higher, and (3) smaller contribution of photorespiration led to the emission of O₂ with larger $^{17}\Delta$ by the glacial terrestrial biosphere.

At the LGM, $\delta^{18}\text{O}$ of seawater was about 1‰ larger than today (Waelbroeck et al., 2002), and from the understanding of isotope hydrology, there must have been parallel change in $\ln(\delta^{17}\text{O} + 1)$ and it can be calculated as $0.528 \cdot \ln(\delta^{18}\text{O} + 1)$. Because for global considerations $^{17}\Delta$ is calculated with λ equal 0.516 (see lines of constant $^{17}\Delta$ in Figure 13), it can be calculated that this isotopic change in seawater involved a 12 ($=[(0.528 - 0.516) \cdot \ln(1.0/1000 + 1)] \cdot 10^6$) per meg increase in $^{17}\Delta_{\text{bio}}$ of the glacial ocean.

In a similar way, due to the widespread glaciers, vegetation in the high latitudes of the northern hemisphere was greatly reduced. Landais et al. (2007a) estimated that, on average, $\delta^{18}\text{O}$ of meteoric water used by the LGM vegetation was about 1.2‰ higher than at present. Calculating in the same way as earlier, such a shift in $\delta^{18}\text{O}$ should have been associated with a 15 per meg increase in LGM leaf water.

Finally, global O₂ uptake via photorespiration was reduced due to a large reduction in the abundance of C₃ vegetation and the increase of the relative abundance of C₄ plants (it is known that C₄ vegetation concentrates CO₂ and therefore suppresses photorespiration). At present, the $\ln(\delta^{17}\text{O} + 1)/\ln(\delta^{18}\text{O} + 1)$ in terrestrial O₂ uptake is reduced from 0.516 to 0.514 due to photorespiration. Landais et al. (2007a) calculated that this ratio at the LGM was larger and consequently $^{17}\Delta_{\text{bio}}$ was further increased by 15 per meg.

Considering the above, Landais et al. (2007a) estimated that in the LGM the terrestrial $^{17}\Delta_{\text{bio}}$ was 30 per meg larger than today and then calculated $^{17}\Delta_{\text{bio}}$ estimates for the present as 124–189 per meg and 156–234 per meg for the LGM. Finally, they concluded that $F_{\text{bio,LGM}}$ was only 60–75% of $F_{\text{bio,PST}}$. Despite the uncertainty, this calculation shows that global gross production at the LGM was substantially smaller than in previous estimates.

The above analysis shows that $^{17}\Delta$ of air from ice cores is useful for estimating past changes in global photosynthetic rates. However, derivation of meaningful conclusions requires careful consideration of hydrology and changes in terrestrial vegetation and their effects on $^{17}\Delta_{\text{bio}}$.

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