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THE CRYSTAL STRUCTURE OF  $\text{Xe}_2\text{F}_n+\text{AuF}_6^-$  AND THE RAMAN SPECTRUM OF  $\text{Xe}_2\text{F}_{11}+$

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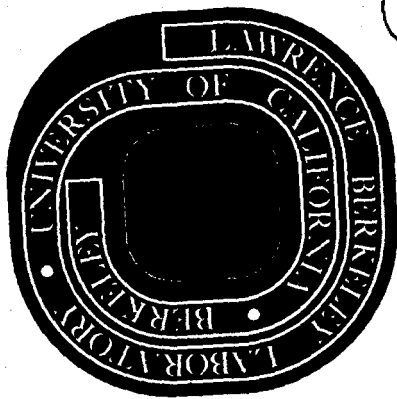
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THE CRYSTAL STRUCTURE OF  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  AND THE RAMAN SPECTRUM OF  $\text{Xe}_2\text{F}_{11}^+$

by

Kevin Leary, Allan Zalkin and Neil Bartlett

ABSTRACT

Crystals of  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  are orthorhombic with  $\underline{a} = 9.115(6)$ ,  $\underline{b} = 8.542(25)$ ,  $\underline{c} = 15.726(20)$  Å,  $\underline{V} = 1224$  Å<sup>3</sup>,  $\underline{z} = 4$ ,  $\underline{D}_c = 4.24$  g cm<sup>-3</sup>, space group  $\underline{Pnma}$ . A structure determination using three dimensional Mo K $\alpha$  X-ray data resulted in a conventional  $\underline{R}$  factor of 0.036 for 862 independent reflections for which  $I \geq 3\sigma(I)$  ( $\underline{R} = 0.052$  for the 1140 independent data, including zero weight data). The anion is essentially octahedral, with an average Au-F bond distance of 1.87(1) Å. The  $\text{Xe}_2\text{F}_{11}^+$  consists of two crystallographically independent  $\text{XeF}_5$  groups bridged by a common fluorine atom, with a bridge-angle of 169.2(2)°; the bridge bond lengths are 2.23 Å (av.) versus 1.84 (av.) for the other Xe-F distances. Each  $\text{XeF}_5$  group departs significantly from the ideal  $\underline{C}_{4v}$  symmetry of the  $\text{XeF}_5^+$  cation. However, the  $\text{F}_{ax}\text{-Xe-F}_{eq}$  angles are ~ 80° for both  $\text{XeF}_5^+$  and the  $\text{XeF}_5$  groups in  $\text{Xe}_2\text{F}_{11}^+$ . The cis angle furthest from the bridging fluorine atom is larger than the others indicating that the bridging F atom may be deflecting the non-bonding Xe(VI) valence electron pair from its ideal position in each pseudo-octahedral,  $\text{XeF}_5^+$ -like component of  $\text{Xe}_2\text{F}_{11}^+$ .

Raman data indicate that the complex cation behaves vibrationally like two weakly coupled  $\text{XeF}_5^+$  species with a 'bridge stretch' at  $\approx 360 \text{ cm}^{-1}$ . This and the structural data indicate that  $\text{F}_5\text{Xe}^+ \text{F}^- \text{XeF}_5^+$  must be a major canonical form in the resonance hybrid description of the cation.

### INTRODUCTION

Recently we set out to synthesize  $\text{AuF}_6^-$  and obtained<sup>1</sup> our first salt of this anion in the form of the complex cation salt  $\text{Xe}_2\text{F}_{11}^+ \text{AuF}_6^-$ . Since both ions were novel and of structural interest, we were fortunate that our synthetic method yielded suitable single crystals for an X-ray structural analysis.

Bartlett and his coworkers<sup>2</sup> had prepared a salt of empirical formula  $\text{F}_{17}\text{PtXe}_2$ , at the time they characterized the salt  $\text{XeF}_5^+ \text{PtF}_6^-$ , and considered it likely to be  $\text{Xe}_2\text{F}_{11}^+ \text{PtF}_6^-$ . The composition of the latter, which can also be expressed as  $2\text{XeF}_6 \cdot \text{PtF}_5$  adduct, suggested that the compound  $2\text{XeF}_6 \cdot \text{SbF}_5$ , described even earlier<sup>3</sup> by Gard and Cady, was also probably an  $\text{Xe}_2\text{F}_{11}^+$  salt.

On the basis of Raman data<sup>4</sup> for  $2\text{XeF}_6 \cdot \text{AsF}_5$  and  $\text{XeF}_6 \cdot \text{AsF}_5$  and the crystal structure of  $\text{XeF}_5^+ \text{AsF}_6^-$  (ref. 5) Bartlett and Wechsberg<sup>4</sup> concluded that the former complex was  $\text{Xe}_2\text{F}_{11}^+ \text{AsF}_6^-$ . Although Bartlett and Wechsberg were able to obtain single crystals of the arsenic complex<sup>4</sup>, all showed disorder or gross twinning features and were unsuitable for an X-ray structure determination. Sladky and Bartlett<sup>6</sup> had similar difficulties in preparing single crystals of  $\text{Xe}_2\text{F}_{11}^+ \text{RuF}_6^-$ , which was considered to be the best platinum-metal-pentafluoride case (because of the lower atomic number of ruthenium) for a structural characterization of the cation.

The controversy concerning the nature of the bonding in  $\text{XeF}_6$ <sup>7,8,9</sup> and the role of the 'non-bonding' Xe valence electron pair in determining the shape of the molecule, gives added interest to the geometry of the  $\text{Xe}_2\text{F}_{11}^+$  cation. It seemed even at the outset, however, (see ref. 4) that the complex cation would be a symmetrical fluorine bridged  $\text{F}_5\text{Xe}\dots\text{F}\dots\text{XeF}_5$  species and, in particular, a relationship to the crystalline  $\text{XeF}_6$  structure<sup>10</sup> was anticipated.

Although the structure determination of an alkali fluoroaurate would have been more satisfactory for the description of the  $\text{AuF}_6^-$  ion, single crystals of  $\text{MAuF}_6$  (M = Cs, Rb, K, etc.) have not yet been obtained. Nevertheless, with allowance for the perturbing influence of the unsymmetrical cation, an adequate description of  $\text{AuF}_6^-$  has been provided by the structure of  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$ .

#### EXPERIMENTAL

Crystal Preparation.-  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  was prepared as previously described<sup>1</sup>. Crystals were grown by placing  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  (1.28 mmoles) and  $\text{XeF}_2$  (5.91 mmoles), prepared as previously described<sup>11</sup>, in a Monel autoclave bomb. Fluorine gas (70 mmoles) was added by condensing with liquid nitrogen. The bomb was heated at 400° for 48 hours under a fluorine pressure of 1000 psi. It was then cooled slowly to room temperature overnight and the excess  $\text{F}_2$  and  $\text{XeF}_6$  were removed under vacuum. The bomb was opened in the dry nitrogen atmosphere of a Vacuum Atmospheres Corporation Dri Lab. The  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  lay in the bottom of the bomb as a mass of small yellow-green plates, whose crystal habit was orthorhombic. Crystals were wedged into small quartz capillaries with

Pyrex push rods, then sealed temporarily with Kel-F grease. On removal from the Dri Lab the capillaries were immediately sealed in a small flame. Because  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  is extremely water sensitive, the utmost precautions were taken to exclude water from all apparatus and materials.

Crystal Data.-  $\text{F}_{17}\text{AuXe}_2$  (mol. wt. 782.5) is orthorhombic with  $a = 9.115(6)$ ,  $b = 8.542(25)$  and  $c = 15.726(20)$  Å,  $V = 1224.3$  Å<sup>3</sup>,  $z = 4$ ,  $D_c = 4.24$  g cm<sup>-3</sup>,  $\mu$  (Mo K $\alpha$ ) = 182 cm<sup>-1</sup>, and  $F(000) = 1345.72$ . The rather large estimated standard deviations of the cell constants reflect changes during data collection, presumably as a consequence of some decomposition. A powder photograph of the bulk material was indexed using the single crystal unit cell dimensions. The unit cell volume satisfies Zachariason's criterion<sup>12</sup> for close packed fluoride lattices, since the effective volume per fluorine atom is 18.0 Å<sup>3</sup>. Single crystal Weissenberg photographs indicated that the space group was either  $\text{Pnma}$  (#62) or  $\text{Pn}2_1b$  (#33 in a non-standard setting). The structure was successfully refined in the centrosymmetric group  $\text{Pnma}$ . Due to the extreme reactivity of the material no attempt was made to obtain an experimental density.

X-Ray Measurements.- A Picker automatic four-circle diffractometer equipped with a fine focus Mo anode tube ( $\lambda$  Mo K $\alpha_1 = 0.70926$  Å), and a graphite monochromator was used for data collection. Accurate cell dimensions were obtained from a least-squares refinement of the orientation matrix and of the cell parameters based on the four angle settings ( $2\theta$ ,  $\omega$ ,  $\psi$  and  $\phi$ ) of 12 high angle ( $45 \leq 2\theta < 50^\circ$ ) reflections. Intensity data were collected by the  $\theta$ - $2\theta$  scan technique, at a scan rate of 2°/min. The scan width was 1.4°. Background counts were offset from the scan limits by 0.8°, and each count lasted 4 seconds. Three standards were checked every 50

reflections. The temperature during data collection was  $24 \pm 1^\circ$ . Because  $\text{Xe}_2\text{F}_{11}^+$  decomposes slowly in X-rays, we had to use four crystals in the data collection. All four were flat plates, elongated along the b axis, and each was mounted with the b axis along the  $\phi$  axis of the diffractometer. The crystals had the following dimensions for the h, k, and l directions, respectively: #1, 0.255, 0.345, 0.047; #2, 0.282, 0.351, 0.060; #3, 0.204, 0.462, 0.096; #4, 0.231, 0.600, 0.072 mm (the precision of these measurements is probably no better than  $\pm 0.005$  mm). Each of the crystals was bounded by the six planes of the forms [100], [010] and [001]. The first three crystals provided a complete set of +h, +k, +l data to  $2\theta \leq 50^\circ$  (1162 reflections). Crystal #4 yielded a complete set of +h, +k, +l and +h, -k, +l data to  $2\theta \leq 40^\circ$  (1204 reflections). A crystal was discarded when an  $\omega$  scan half width of any standard reflection became  $\geq 0.25^\circ$ . Intensity decay of the standards was no greater than 20.0% in any one crystal. Corrections for decay were made.

Because of the large absorption coefficient ( $182 \text{ cm}^{-1}$ ) and the fact that all the crystals were much larger than the optimum size, the data were corrected for absorption using a program developed by P. Coppens, L. Leisenwitz, and D. Rabinovich<sup>13</sup>, modified by D. Cahen and J. Ibers<sup>14,15</sup>, and adapted for local use. This method incorporates numerical integration using a Gaussian grid. The data were treated and weights assigned as previously described<sup>16</sup> with the exception that a q factor (used to decrease the weights of large intensities) of zero was used. Scattering factors<sup>17</sup> for neutral gold, xenon and fluorine were used. Values for anomalous dispersion,  $\Delta f'$  and  $\Delta f''$ , were taken from Cromer and Liberman<sup>18</sup>.



Structure Refinement.- Initially, the structure was solved using the data from Crystal #4 only ( $2\theta \leq 40^\circ$ ). This set of data yielded an averaged set of 633 unique reflections, of which 529 satisfying the condition  $I \geq 3\sigma(I)$  were used in the least-squares refinement.

A Patterson function yielded the positions of the two Xe and one Au atoms and refined to  $R = 0.21$ . The Patterson map confirmed the choice of the centric space group since all the Harker sections had either  $y = 0$  or  $y = 1/2$ . A difference Fourier based on the set of phases generated by the least-squares refinement gave the positions of all of the fluorine atoms. Least-squares refinement incorporating anisotropic temperature factors for all atoms gave a final  $R$  factor of 0.04, using the data from Crystal #4.

The data obtained from the first three crystals (+h, +k, +l,  $20 \leq 50^\circ$ ) were then scaled and averaged with the set obtained from crystal #4. Of the 1162 total independent data, the 874 which satisfied the condition  $I \geq 3\sigma(I)$  were used for the least-squares refinement. A least-squares refinement using the combined data gave  $R = 4.48\%$ , and weighted  $R_2 = 3.98\%$ .

It was noted that the higher angle data ( $40^\circ \leq 2\theta \leq 50^\circ$ ) had large weighted discrepancies,  $w(\Delta F)^2$ . These data had been measured only once whereas the data below  $2\theta = 40^\circ$  were measured at least three times. For multiply measured data, the estimated standard deviation is based on the larger of either their counting statistics or their scatter, and thus the standard deviations of the high angle data were consequently smaller. To correct this situation a minimum value corresponding to what was observed for the weaker lower angle data was then applied as

a lower limit to the standard deviations of the higher angle data. The four sets of data were rescaled. All of the data with  $\sin\theta/\lambda \leq 0.15$ , (a total of 22 reflections) were arbitrarily deleted because of excessively large discrepancies; this is no doubt due to the inadequacies of the absorption correction. Examination of the data showed no extinction effects.

The final least-squares refinement, with all atoms anisotropic, gave an  $R$  factor of 0.052 for all 1140 reflections, and 0.036 for the 862 ( $I \geq 3\sigma(I)$ ) non-zero weighted reflections. The weighted  $R_2$  was 0.025. The standard deviation of an observation of unit weight was 1.36. The largest shift of any parameter, divided by its estimated standard deviation on the last cycle of least-squares was  $\leq 0.0012$ .

A final difference Fourier showed that the largest residual electron density was 1.91 electrons/ $\text{\AA}^3$  near the gold atom. Table I gives the positional and thermal parameters from the final refinement. Observed structure factors, standard deviations and differences are given in Table II which will appear in the microfilm version of this paper<sup>19</sup>. Table III gives chemically significant distances and angles.

Raman Spectra.- Microcrystalline samples of  $[\text{XeF}_5^+]_2\text{PdF}_6^{2-}$ ,  $[\text{Xe}_2\text{F}_{11}^+]_2\text{PdF}_6^{2-}$ ,  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  and  $\text{XeF}_5^+\text{AuF}_6^-$ , enclosed in 1 mm o.d. quartz capillaries, were excited at 6328  $\text{\AA}$ , using a 100 mW He/Ne ion laser and spectra were recorded from a Spex Model 1400 double monochromator. Spectra were also obtained on a Cary 83 spectrometer equipped with a 100 mW  $\text{Ar}^+$  ion (4880  $\text{\AA}$ ) laser. The spectra are tabulated in Table IV and the  $\text{AuF}_6^-$  salt spectra are given in the microfilm version of this paper<sup>19</sup> as Figure 4.

DESCRIPTION OF THE STRUCTURE

As may be seen from Figure 1, and Table III, the structure analysis clearly defines an  $\text{AuF}_6$  group and a  $\text{Xe}_2\text{F}_{11}$  group. The latter consists of two similar  $\text{XeF}_5$  groups linked by an additional common F atom. All of these groups ( $\text{AuF}_6$ ,  $\text{Xe}_2\text{F}_{11}$  and its  $\text{XeF}_5$  components) possess mirror symmetry.

The  $\text{AuF}_6$  group is approximately octahedral, with only one Au-F interatomic distance ( $\text{Au-F8} = 1.90(1) \text{ \AA}$ ) departing significantly from the average value of  $1.86(1) \text{ \AA}$ . The cis F-Au-F angles are close to  $90^\circ$ ; the greatest deviation being for  $\text{F8-Au-F12} = 88.0(5)^\circ$ .

Each  $\text{XeF}_5$  group of the  $\text{Xe}_2\text{F}_{11}$  species approximates to a square-based pyramid, with the xenon atom placed below the base. The  $\text{F}_{\text{ax}}\text{-Xe-F}_{\text{eq}}$  angles are  $\approx 80^\circ$  in both  $\text{XeF}_5$  groups. On the other hand, each of the groups departs significantly from  $\text{C}_{4v}$  symmetry and the cis  $\text{F}_{\text{eq}}\text{-Xe-F}_{\text{eq}}$  angles in each group are not equivalent. The greatest cis angle of each  $\text{XeF}_5$  equatorial set, is that furthest from the atom (F7) which links the  $\text{XeF}_5$  groups to define the  $\text{Xe}_2\text{F}_{11}$  species. Coincidentally, the greatest cis equatorial angles in each  $\text{XeF}_5$  group are also associated with the longest Xe-F distances within the group.

Although the interatomic distances  $\text{Xe1-F7}$  and  $\text{Xe2-F7}$  are sufficiently short [ $2.21(1)$  and  $2.26(1)$  respectively] to warrant the identification of an  $\text{Xe}_2\text{F}_{11}$  group, all other inter-group contacts are sufficiently long that they may be classified as van der Waals contacts. The  $\text{Xe1-F7-Xe2}$  angle is not quite linear [ $169.2(2)^\circ$ ]. The  $\text{F7-Xe}$  distances are not significantly different and indeed the entire  $\text{Xe}_2\text{F}_{11}$  group has essentially  $\text{C}_{2v}$  symmetry.

The bridging F atom (F7) is not the only F atom of interest in relationship to the  $\text{XeF}_5$  groups. It is seen that each Xe atom of each  $\text{XeF}_5$  is approached by three other F atoms (of neighboring  $\text{AuF}_6$  groups) as well as by F7. Thus Xe1 is associated with F7, F9, F9' and F8 and Xe2 with F7, F9, F9' and F12 [see Figure 1(c)]. These sets of four F atoms are arranged about the base of the approximately square-pyramidal  $\text{XeF}_5$  groups, such that, together with the F atoms of the  $\text{XeF}_5$ , they form a distorted capped archimedean antiprism arrangement. The arrangement is illustrated, for the Xe1 case, in Figure 2.

#### DISCUSSION

Prior to the structure determination, chemical, vibrational, spectroscopic and Mössbauer evidence<sup>1</sup> had indicated the formulation  $\text{Xe}_2\text{F}_{11}^+$   $\text{AuF}_6^-$  for the  $\text{F}_{17}\text{Xe}_2\text{Au}$  material. The X-ray structure is fully compatible with that formulation.

The geometry of the  $\text{AuF}_6^-$  ion is defined for the first time in this structure. A low spin  $d_{-t2g}^6$  Au(V) electron configuration is anticipated to be akin to the configurations of Pt(IV) and Pd(IV) and like them to favor a regular octahedral  $\text{MF}_6$  species. Any departures from octahedral symmetry, in this structure, can be excused on the basis of possible distorting influences of the  $\text{Xe}_2\text{F}_{11}^+$  cation which is far from octahedral symmetry itself. The average Au-F anion interatomic distance of 1.86(1) Å compares with the average Pt-F distance<sup>2</sup> of 1.89(5) Å for  $\text{PtF}_6^-$  in  $\text{XeF}_5^+\text{PtF}_6^-$ . The greater nuclear charge of Au relative to Pt is anticipated<sup>20</sup> to result in a shorter M-F bond in the Au case.

This structure is at least of as much interest for its cation as

for its anion. To a first approximation the complex cation has the form anticipated for an assembly of two  $\text{XeF}_5^+$  ions and a common  $\text{F}^-$ . Thus the geometry of each  $\text{XeF}_5$  component resembles that of  $\text{XeF}_5^+$  (ref. 21) and the coordination of each Xe atom (represented in Figure 2) closely resembles that for Xe in  $\text{XeF}_5^+\text{RuF}_6^-$ . A comparison of the geometry of the  $\text{XeF}_5$  unit in  $\text{Xe}_2\text{F}_{11}^+$  with that of the  $\text{XeF}_5^+$  cation in  $\text{XeF}_5^+\text{RuF}_6^-$  is given in Figure 3. The close approach of F7 to each Xe atom and the departure of the Xe-F-Xe angle from linearity, suggest that a measure of covalency should be incorporated into any bonding description. Nevertheless, the ionic model  $\text{XeF}_5^+\text{F}^-\text{XeF}_5^+$  accounts for a number of the observed structural features of  $\text{Xe}_2\text{F}_{11}^+$ . In previous papers<sup>21,22</sup> we have argued for steric activity of the non-bonding Xe(VI) electron pair in the pseudo-octahedral  $\text{XeF}_5^+$  species. We can allow that the close approach of  $\text{F}^-$  to  $\text{XeF}_5^+$  would deflect the non-bonding valence electron pair from its axial position towards the bisectors of the F3-Xe1-F3' and the F5-Xe2-F5' angles. Increase in the non-bonding pair repulsive interactions with the Xe1-F3 and Xe2-F5 bonds, consequent upon such deflection of the electron pairs, could also account for the lengthening of the Xe1-F3 and Xe2-F5 interatomic distances. Similarly, the deflection of the Xe(VI) 'pairs' could account for shortening of the Xe1-F2 and Xe2-F6 distances. In  $\text{XeF}_5^+$  the Xe-F equatorial distances<sup>21</sup> are 1.84 Å.

As may be seen from the Raman data given in Table IV, the complex  $4\text{XeF}_6 \cdot \text{PdF}_4$  (ref. 23) contains essentially the same cation as  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  and may, therefore, be confidently formulated as  $[\text{Xe}_2\text{F}_{11}^+]_2\text{PdF}_6^{2-}$ . The F bridging of two  $\text{XeF}_5$  groups in the  $\text{Xe}_2\text{F}_{11}^+$  cation appears to be characterized by a 'bridge stretch' at  $\sim 360 \text{ cm}^{-1}$ ; but, in keeping with the observed structure, the complex cation otherwise behaves vibrationally

like two weakly coupled  $\text{XeF}_5^+$  species.

Since crystalline  $\text{XeF}_6$  can be described<sup>10</sup> as  $\text{XeF}_5^+\text{F}^-$  (clustered either in tetramers or hexamers) it is not surprising that  $\text{Xe}_2\text{F}_{11}^+$  looks like a fragment of an  $\text{XeF}_6$  tetramer or hexamer. The resemblance is closest to the tetramer. In a bonding description of  $\text{Xe}_2\text{F}_{11}^+$  and crystalline  $\text{XeF}_6$ ,  $\text{XeF}_5^+$  and  $\text{F}^-$  are clearly important contributing canonical forms. Perhaps the best description of  $\text{Xe}_2\text{F}_{11}^+$  is as a resonance hybrid of  $(\text{XeF}_5^+\text{F}^-\text{XeF}_5^+)$ ,  $(\text{XeF}_6\text{XeF}_5^+)$  and  $(\text{XeF}_5^+\text{XeF}_6)$  with the first canonical form dominant.

It seems probable that all  $\text{XeF}_6$  complexes with fluoride ion acceptors will prove to be either  $\text{XeF}_5^+$  or  $\text{Xe}_2\text{F}_{11}^+$  salts. In particular, the  $2\text{XeF}_6 \cdot \text{MF}_5$  and  $4\text{XeF}_6 \cdot \text{MF}_4$  complexes reported by Cady and his coworkers<sup>3,24</sup> are very probably  $\text{Xe}_2\text{F}_{11}^+$  salts.

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TABLE I

Positional and Thermal Parameters for  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$ 

ATCM	X	Y	Z	B11	B22	B33	B12	B13	B23
AU	.15834(6)	.250*	.02276(5)	2.42(3)	3.32(3)	3.95(3)	0*	-.44(2)	0*
XE(1)	.29072(9)	.250*	.66069(7)	2.80(4)	4.73(6)	3.03(5)	0*	.32(3)	0*
XE(2)	.2013(1)	.250*	.38245(7)	3.08(4)	5.06(6)	3.36(5)	0*	-.65(3)	0*
F(1)	.234(1)	.250*	.7726(6)	7.0(5)	7.9(7)	2.3(4)	0*	1.0(4)	0*
F(4)	.077(1)	.250*	.2933(7)	6.5(6)	11.1(8)	5.3(6)	0*	-3.7(4)	0*
F(7)	.2238(8)	.250*	.5254(6)	4.4(4)	4.8(5)	3.7(5)	0*	-.2(4)	0*
F(8)	.041(1)	.250*	-.0771(6)	4.1(4)	8.3(6)	5.4(5)	0*	-1.6(4)	0*
F(10)	.270(1)	.250*	.1212(7)	5.6(5)	7.2(6)	5.9(7)	0*	-3.2(4)	0*
F(11)	.3300(9)	.250*	-.0407(7)	3.1(4)	9.3(8)	6.9(7)	0*	1.0(4)	0*
F(12)	-.0143(9)	.250*	.0857(6)	3.1(4)	9.9(7)	5.5(5)	0*	.6(3)	0*
F(2)	.1448(5)	.1026(9)	.6578(4)	4.3(3)	6.2(4)	5.3(4)	-1.9(3)	1.0(2)	.2(3)
F(3)	.4114(7)	.0949(8)	.7046(4)	5.0(3)	6.3(4)	5.1(4)	1.5(3)	.0(2)	1.3(3)
F(5)	.2914(7)	.096(1)	.3168(5)	5.9(4)	11.3(7)	7.0(5)	-1.7(4)	-.0(3)	-4.5(5)
F(6)	.0708(6)	.1013(8)	.4160(5)	3.6(3)	5.4(4)	8.1(5)	-1.4(2)	-.9(2)	.1(3)
F(9)	.1587(8)	.033(1)	.0236(6)	8.2(4)	4.3(4)	9.6(6)	-.2(3)	-3.0(4)	.3(4)

\* Fixed Parameter



TABLE III

Interatomic Distances( $\text{\AA}$ ) and Angles(DEG.) for  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$ 

(Estimated Standard Deviations are in Parentheses)

<u>Anion</u>		<u>Cation</u>			
F9-Au-F9'	179.1(3) $^\circ$	F1-Xe1-F2	79.6(4) $^\circ$	F4-Xe2-F6	79.4(4) $^\circ$
F8	90.4(4)	F3	79.0(4)	F5	81.1(5)
F8-Au-F11	91.6(6)	F2-Xe1-F2	86.8(3)	F6-Xe2-F6	88.6(3)
12	88.0(5)	F3	87.1(4)	F5	87.4(6)
F10-Au-F11	89.2(5)	F3-Xe1-F3	91.1(3)	F5-Xe2-F5	89.9(3)
F12	91.2(6)	Xe2-F7-Xe1	169.2(2)		

Interionic Distances and Angles

Xe1-F8	2.64(1) $\text{\AA}$	F1-Xe1-F7	147.6(6) $^\circ$	F4-Xe2-F7	146.3(7) $^\circ$
F9	3.27(1)	F8	136.3(6)	F12	139.8(7)
Xe2-F12	2.64(1)	F9	132.2(3)	F9	135.9(2)
F9	3.52(1)				

TABLE IV

Raman Spectra of  $\text{Xe}_2\text{F}_{11}^+$  Salts and Related Species (shifts in  $\text{cm}^{-1}$ )

$\text{IF}_5^{(a)}$	$(\text{NO}^+)_2\text{PdF}_6^{2-}$	$(\text{XeF}_5^+)_2\text{PdF}_6^{2-}$	$(\text{Xe}_2\text{F}_{11}^+)_2\text{PdF}_6^{2-}$	$(\text{Xe}_2\text{F}_{11}^+)\text{AuF}_6^-$	$\text{Cs}^+\text{AuF}_6^-$
710( $\underline{\nu}_1$ ) vs	.....	653 s	.....	[ 668 w 651 vs	..... 661 s
631( $\underline{\nu}_7$ ) sh	.....	634 vw	.....	630 m	..... 640 w
614( $\underline{\nu}_2$ ) vs	.....	[ 616 vw 606 w	.....	[ 615 w 610 m 606 w	..... [ 626 w 600 s
602( $\underline{\nu}_4$ ) sh	.....	590 ms	.....	[ 591 s 583 m	..... 595 vs ..... 590 vs
	573( $\underline{\nu}_1$ ) s	.....	558 ms	.....	568 s
	554( $\underline{\nu}_2$ ) ms	.....	535 s	.....	546 w
					530( $\underline{\nu}_2$ ) vw
370( $\underline{\nu}_8$ ) w	.....	[ 425 vw 396 vw	.....	[ 412 vw 396 vw	..... 400 w
				375 w	356 w
318( $\underline{\nu}_3$ ) m	.....	309 w	.....	296 w (b)	..... 290 w
274( $\underline{\nu}_6$ ) w	.....	269 w	.....	270 w (b)	---
	243( $\underline{\nu}_5$ ) ms	.....	245 w	.....	245 w (b)
					221 m
					224( $\underline{\nu}_5$ ) s

(a) L. E. Alexander and I. R. Beattie, J. Chem. Soc.(A),  
1971, 3091.

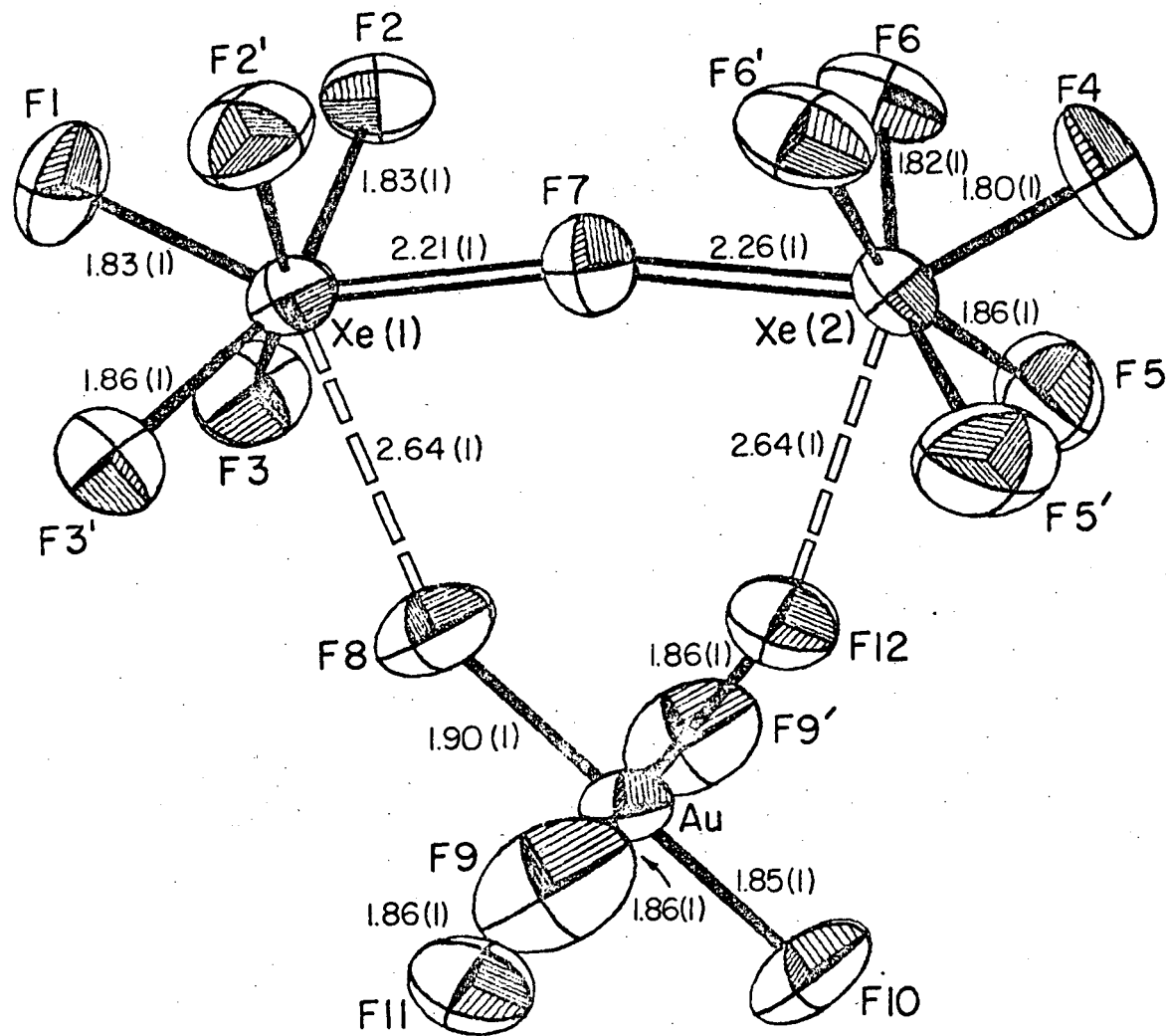
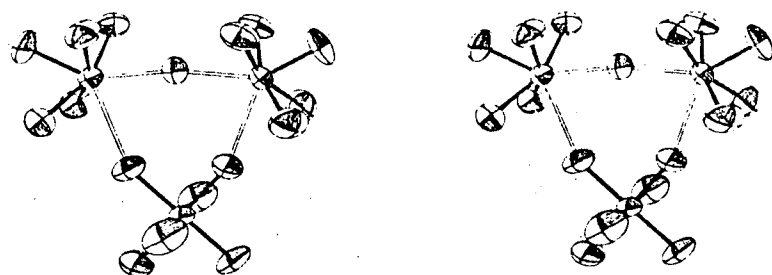
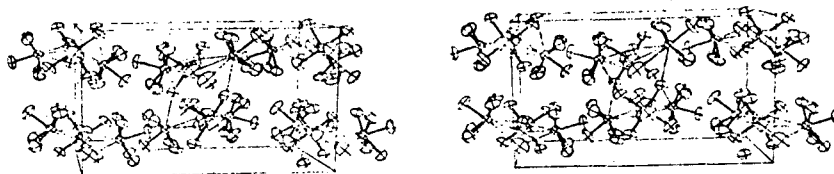


Figure 1a: The structural unit  $\text{Xe}_2\text{F}_{11}\text{AuF}_6^+$



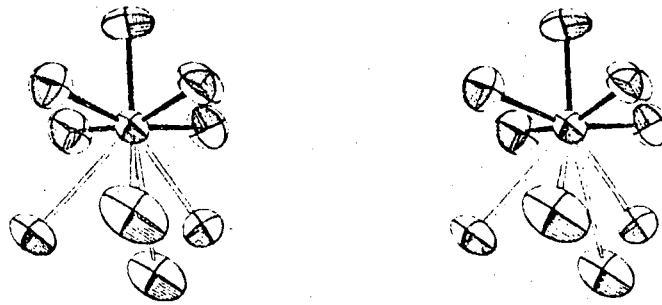
XBL 7210-7124

Figure 1b: Stereogram of the  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  Structure Unit.



XBL 7210-7127

Figure 1c: Stereogram showing the arrangement of the  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$  units in the unit cell - view along the b axis.



XBL 7210-7125

Figure 2: Stereogram showing the typical Xe coordination in F atoms (exemplified by Xe(1) coordination).



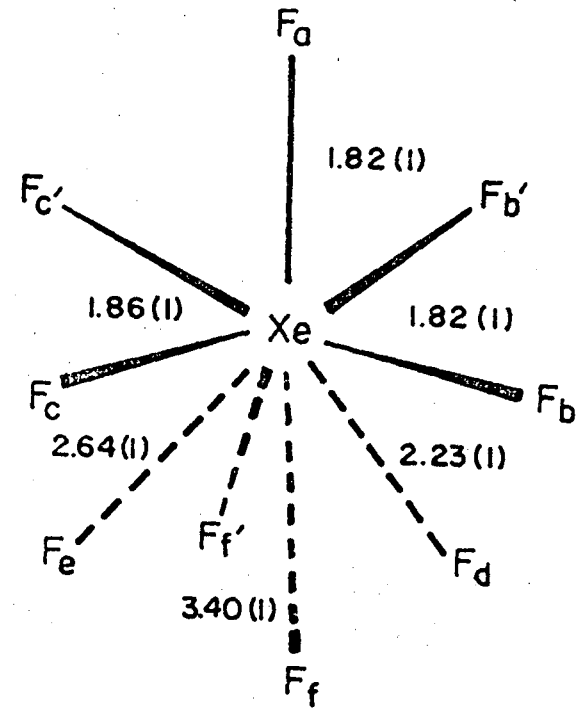
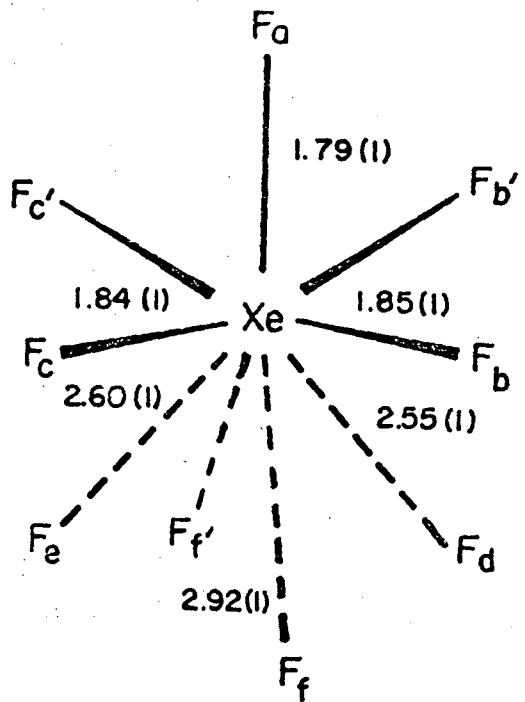
Table for Figure 3  
 (insert figure here as instructed)

$\text{XeF}_5^+$  in  $\text{XeF}_5^+\text{RuF}_6^-$  (a)       $\text{XeF}_5$  unit (av.) in  $\text{Xe}_2\text{F}_{11}^+\text{AuF}_6^-$

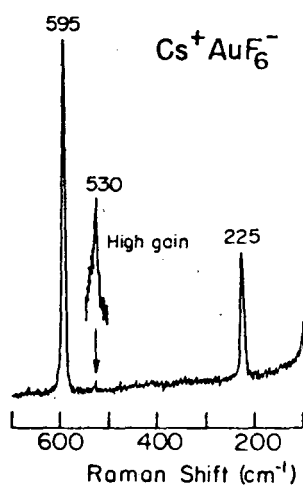
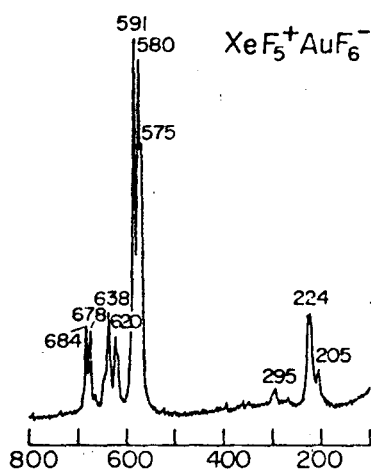
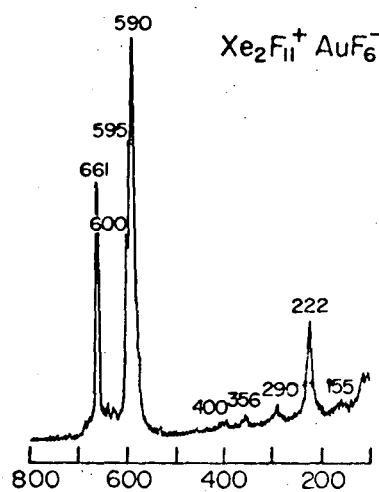
$F_{\text{eq}}-\text{Xe}-F_{\text{ax}}$	79.0(2)°	79.9(2)°
$F_{\text{b}}-\text{Xe}-F_{\text{b}}'$	87.8(2)	87.7(2)
$F_{\text{c}}-\text{Xe}-F_{\text{c}}'$	87.0(2)	90.5(2)
$F_{\text{b}}-\text{Xe}-F_{\text{c}}$	88.4(4)	87.2(3)
$F_{\text{d}}-\text{Xe}-F_{\text{a}}$	142.3(7)	147.0(5)
$F_{\text{e}}-\text{Xe}-F_{\text{a}}$	140.6(6)	138.0(5)
$F_{\text{f}}-\text{Xe}-F_{\text{a}}$	129.6(3)	134.0(2)

(a) Ref. 22

Figure 3



XBL 733-5892



XBL 7210-7074

Figure 4: Raman Spectra of some  $\text{AuF}_6^-$  salts.

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