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ABSTRACT

Analytic formulas suitable for numerical computation are derived for two-dimensional air coil systems. The computer code, FIELDS, is described, and evaluation of a sextupole magnet is included as a sample problem. A complete listing of the program coding is also provided in an appendix.

I. INTRODUCTION

A method of calculating two-dimensional magnetic fields in vacuum will be developed consisting of evaluation of analytical complex formulas. A description of the resulting computer code, FIELDS, is included, and user instructions are provided including a sample sextupole magnet problem. The method treats infinite straight conductors with arbitrary cross section and uniform current density, j , in each conductor, and no permeable material may be present. Field calculations may be obtained both inside and outside the coil area without special treatment. In addition a multipole expansion of the fields is provided.

II. MAGNETIC FIELD EQUATIONS FOR TWO-DIMENSIONAL COILS

Complex algebra will be used in the following field derivations, and MKS units will be used throughout. Magnetic fields can be described by the complex quantity $H = H_x + iH_y$. Since only the complex conjugate of H , namely $H^* = H_x - iH_y$, is an analytic function of $z = x + iy$, this is the quantity used.

For a current filament with positive current I (positive in the $\vec{x} \times \vec{y}$ direction, or directed out of the page) shown in Fig. 1 at point z , the magnetic field at a point z_0 has a strength equal to $I/(2\pi |z - z_0|)$. The field direction (perpendicular to the vector $z_0 - z$) may be given by $i(z_0 - z)/|z_0 - z|$.

In equation form,

$$H = \frac{I}{2\pi |z - z_0|} \frac{i(z_0 - z)}{|z_0 - z|}$$

This may be rewritten as

$$H = \frac{I}{2\pi} \frac{i(z_0 - z)}{(z_0 - z)(z_0^* - z^*)} = \frac{Ii}{2\pi(z_0^* - z^*)} \quad (2)$$

where z^* and z_0^* are complex conjugates of z and z_0 respectively. Introducing H^* in (2)

$$H^* = \frac{Ii}{2\pi(z - z_0)} \quad (3)$$

For several current filaments the fields at z_0 may be obtained by superposition. Extending this to a conductor of arbitrary cross section with uniform current density j , the problem becomes one of integration over the cross section.

$$H^* = \frac{ji}{2\pi} \int \frac{1}{z - z_0} dx dy \quad (4)$$

The integral in (4) is of the form $\int F(z) dx dy$, where $F(z)$ is analytic and single valued over the integration region, and may have one singularity at $z = z_0$. (If no field values are to be evaluated inside the coil, z_0 is outside the integration region, and $F(z)$ has no singularities). For a more general $F(z)$ (with "n" first order singularities inside the integration region) it is shown in Appendix A that the integral of (4) may be represented by

$$\frac{1}{2i} \left(\oint F(z) z^* dz - \sum_{\text{Sum over "n" poles}} \oint F(z) z_n^* dz \right)$$

The contour integrations in the sum are performed for infinitesimal circles about each pole.

In (4) $F(z) = 1/(z - z_0)$, and H^* may be written as

$$H^* = \frac{j}{4\pi} \oint \frac{z^* - z_0^*}{z - z_0} dz \quad (5)$$

Equation (5) is valid for field points both inside and outside the coil region. This formulation differs slightly from that developed by R. A. Beth⁽¹⁾ in that he considers interior and exterior points separately. Equation (5) is more convenient to use for numerical computation since no checks are required to determine the type of point. Another advantage of (5) is that the absolute value of the integrand is always unity. For field points near the coil Beth's technique requires taking differences between large numbers over part of the integration contour. Use of (5) avoids both problems, and is used for calculating all field values in the computer code. For completeness, we derive (5) in the appendix with a procedure that differs from Beth's, and allows extension to other problems.

III. MULTIPOLE COEFFICIENTS

Multipole coefficients are obtained by expanding the fields in a power series over a region excluding all currents. In such a region the complex potential $G(z) = A + iV$ (where A and V are, respectively, the vector and scalar potentials) is analytic and $H^*(z)$ can be obtained from $G(z)$ through $H^*(z) = iG'(z)$. Therefore $G(z)$ may be expanded in a converging Taylor Series of the complex variable z_0 .

$$G(z_0) = a_0 + a_1 z_0 + a_2 z_0^2 + \dots + a_n z_0^n \quad (7)$$

$$G(z_0) = \sum_{n=0}^{\infty} a_n z_0^n$$

Similarly

$$H^*(z_0) = i \sum_{n=1}^{\infty} a_n z_0^{n-1} = \sum_{n=1}^{\infty} b_n z_0^{n-1} \quad (8)$$

where a_n and b_n are in general complex constants.

This Taylor Series has a convergence radius equal to the distance from the expansion origin to the closest singularity (the closest coil dimension, in this case). The series converges, then, only within a circle extending to the nearest coil boundary. In such a region (5) may be written as

$$H^* = \frac{j}{4\pi} \oint \frac{z^*}{z - z_0} dz = \frac{j}{4\pi} \oint \frac{z^*}{z(1 - \frac{z_0}{z})} dz \quad (9)$$

and for $|z_0/z| < 1$ this becomes

$$H^* = \frac{j}{4\pi} \oint \frac{z^*}{z} \sum_{n=1}^{\infty} \frac{z_0^{n-1}}{z^{n-1}} dz, \quad (9a)$$

Equating the two representations for H^* in this circular region, namely (8) and (9a)

$$\sum_{n=1}^{\infty} b_n z_0^{n-1} = \frac{j}{4\pi} \oint \frac{z^*}{z} \sum_{n=1}^{\infty} \frac{z_0^{n-1}}{z^{n-1}} dz, \quad (10)$$

and solving (10) for all values of b_n , one obtains

$$b_n = \frac{j}{4\pi} \oint \frac{z^*}{z^n} dz \quad (11)$$

Although certain coil symmetries produce only real or only imaginary coefficients, b_n in general may be complex coefficients.

IV. DESCRIPTION OF FIELDS PROGRAM

A Fortran IV computer code has been written to evaluate magnetic fields and multipole coefficients using equations (5) and (11). Input data requires only specification of coil boundary and current information, certain coil symmetry conditions, and a description of the desired grid for field evaluation points.

Coil areas are approximated by inscribed polygons to enable numerical treatment of the line integrals in (5) and (11). A "Trapezoid rule" technique is currently used for the integration. The user may change the number of polygon line segments at will, and also could easily implement additional integration methods (Simpson's rule among others). The symmetry order is treated automatically and only coordinate points for a single coil are actually specified.

The total current, I in amperes, per coil must be given, and a current density is obtained by dividing I by the calculated polygon area. Thus although the calculated model and physical current densities may differ slightly (since true coil area and polygon areas are not quite identical), the total current is identical for both.

A typical problem is shown in Fig. 2. The input data for this sextupole first must specify the total current per coil, I_1 . The coil boundary is next specified (in a counter clockwise direction for a positive current element) by giving end points $((r_1, \theta_1)$ and $(r_2, \theta_2))$ for each coil side. A code $K = 1$

or $K = 2$ is used to indicate whether the coil side is a straight line segment or an arc segment respectively. Finally a value for N is specified to indicate a desired division of the coil segment into N equal length polygon sides.

Actual data cards for the coil boundary shown would be:

2, r_1 , θ_1 , r_2 , θ_2 , 50

1, r_3 , θ_3 , 10

2, r_4 , θ_4 , 60

1, r_1 , θ_1 , 10

where K is either 1 or 2, units for r and θ are cm and radians respectively, and the desired division shown results in a 130 sided polygon approximation. Note that after the first boundary segment is described by first and last points, only the last point is required for each succeeding sector.

Two additional data cards are required to set up the rectangular edit grid. The first card specifies (in order) the number of field points in the x direction, the initial value for x , and the incremental distance between the points. The second card gives identical quantities for y .

All input data is read by subroutine DATA which allows complete field freedom in punching data on the cards. The user should become familiar with the use of this routine as described by Bill Benson's write up: "Subroutine Data - A Fortran 66 Format Free Input Routine", June 9, 1965. (Ref. 2)

The symmetry order of the coil system is specified through a Fortran data statement in the main program. The variable name NCOIL is used, and it represents the number of rotationally symmetrical coils (dipole = 2, sextupole = 6, etc.). In some cases, the user may wish to describe more than one coil in a symmetrical

set (provision for small tuning coils, variable current densities over portions of the total coil, etc.). The variable name IGRP represents the number of separate coils comprising a single symmetrical group or set. If IGRP \neq 1, then separate input cards must describe each of the coil boundaries in the set.

Multipole coefficients are obtained from (11) and multiplied by r_a^{n-1} , where r_a is the radius of the expansion aperture (or a value specified by the user). The thus normalized multipole coefficients describe, according to (8), the contribution of the individual multipoles to the field at the normalization radius. It can easily be shown that rotational symmetries will result in certain " b_n " values being exactly zero for particular "n" in the series of (11). If NCOIL represents the number of rotationally symmetrical coil sets, a general expression for the particular "n" orders that may exist in many coil types is

$$n = \frac{\text{NCOIL}}{2} (2M - 1)$$

where $M = 1, 2, 3, \dots$.

The user may wish to include expressions for additional coil types.

A complete listing of the program coding and the results from the sextupole sample problem are in Appendix B.

APPENDIX

A. Evaluation of $\int F(z) dx dy$

The function $F(z)$ must be analytic, single valued over the integration region, and must have no singularities higher than first order inside the boundary. It is convenient to represent $F(z)$ as the derivative of another function $G(z)$.

$$F(z) = G'(z)$$

$$I = \int F(z) dx dy = \int G'(z) dx dy \tag{A1}$$

Since the integration is over a total area, and since integration over an infinitesimal circular area containing a simple pole contributes only infinitesimally to this total, the region in Fig. A1 may be divided into strips, thus excluding the area inside of the infinitesimal circles around each pole. Integrating first over x , (A1) becomes

$$I = \sum_{\text{sum for all strips}} \int^{G(z, \text{right})} G(z, \text{left}) dy = \sum_{\text{strips}} \oint G(z) dy \tag{A2}$$

Once the contour is defined, $G(z)$ becomes a function of y alone since $z = x(y) + iy$ for points on the boundary. The line integral of (A2) is integrated by parts.

Let

$$u = G(z) \quad du = G'(z) (x'y + i) dy = G'(z) dz$$

$$dv = dy \quad v = y$$

then,
$$\oint G(z) dy = G(z) y \Big|_{\text{start}}^{\text{end}} - \oint G'(z) y dz \quad (\text{A3})$$

Since $G(z)$ is regular and single valued within each loop (strip), and since the path starts and ends at the same point, $G(z) y \Big|_{\text{start}}^{\text{end}} = 0$.

Replacing y by $\frac{1}{2i} (z^* - z)$ and $G'(z)$ by $F(z)$, (A3) becomes

$$\oint G(z) dy = \frac{1}{2i} \oint F(z) (z^* - z) dz \quad (\text{A4})$$

Since all poles have been excluded from the integration strips,

$$\frac{1}{2i} \oint F(z) z dz \text{ must be zero from Cauchy's integral formula.}$$

Therefore

$$I = \frac{1}{2i} \sum_{\text{sum over all strips}} \oint F(z) z^* dz = \frac{1}{2i} \oint_{\text{outer contour}} F(z) z^* dz - \frac{1}{2i} \sum_{\text{each pole inside contour}} \oint F(z) z^* dz \quad (\text{A5})$$

The integrals about each pole must be taken over infinitesimal circular areas.

It is easy to show that for this case

$$\oint_{\text{pole}} F(z) z_n^* dz = z_n^* \oint_{\text{pole}} F(z) dz, \quad (\text{A6})$$

If there is one or no singularity inside the outer contour, (A5) can be expressed as

$$I = \frac{1}{2i} \oint_{\text{outer contour}} F(z) (z^* - z_n^*) dz \quad (A7)$$

It becomes evident by inspection that if there are n singularities, not all of which have to be inside the contour, I becomes

$$I = \int_{\substack{\text{area} \\ \text{inside contour}}} F(z) dx dy = \frac{1}{2i} \oint_{\text{contour}} F(z) (z^* - \sum_n z_n^* R_n(z)) dz \quad (A8)$$

with

$$R_n(z) = \frac{P_n(z)}{P_n(z_n)}, \quad \text{and} \quad P_n(z) = \frac{\prod_{\gamma} (z - z_{\gamma})}{z - z_n}$$

PROGRAM FIELDS(INPUT,OUTPUT)

COMMON/BDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)
COMMON/GRID/X0,Y0,Z0,Z0S
COMMON/MPOL/RNORM
COMMON/XY/IXY

DATA IDIM/2000/
DIMENSION CUR(5), CURDEN(5)

C IGRP REPRESENTS NUMBER (UP TO 5) OF COILS PER GROUP OR SET
DATA IGRP/1/

C NCOIL REPRESENTS NUMBER OF ROTATIONALLY SYMMETRICAL SETS
C (DIPOLE = 2, SEXTUPOLE = 6, ETC)
DATA NCOIL/4/

C IALT = -1 IF SIGNS OF CURRENT ELEMENTS ALTERNATE, OTHERWISE +1
DATA IALT/-1/

C RN IS THE NORMALIZING RADIUS USED IN THE MULTIPOLE ANALYSIS.
C IF RN IS NEGATIVE THE POINT ON COIL NEAREST ORIGIN IS USED FOR
C NORMALIZATION.
DATA RN/-1./

C IF SENTINEL IXY EQUALS 1, THE USER MAY INPUT DATA IN (X,Y)
C INSTEAD OF (R,THETA) COORDS.
DATA IXY/1/

DATA (CUR(I), I=1,4)/5*0.0/
COMPLEX SUM, HSTAR, ROT
COMPLEX Z, ZS, Z0, Z0S

C FIRST DATA CARD CONTAINS CURRENTS -- C1, C2, ... S
C AN S IS REQUIRED IF LESS THAN 5 CURRENT ELEMENTS ARE DESCRIBED.
C DESCRIBE BOUNDARY OF COIL IN A COUNTER CLOCKWISE DIRECTION.
C ENCLOSE POSITIVE CURRENT CARRYING ELEMENT
C REGION MUST CLOSE -- FIRST PT AND LAST PT COORDS MUST BE IDENTICAL.
C FIRST CARD READ MUST CONTAIN (IN ORDER), THE KIND OF CURVE
C SEGMENT (1 FOR STRAIGHT LINE, 2 FOR CIRCULAR ARC), R1, PHI1,
C R2, PHI2, AND THE NUMBER OF INTERVALS DESIRED.
C FOLLOWING CARDS CONTAIN ONLY THE KIND, R2, PHI2, AND NUMBER
C OF INTERVALS.

PI = 3.1415926536898

C READ CURRENTS FOR COILS IN GROUP
CALL DATA(CUR,5)

DO 700 N=1,IGRP

PRINT 600, N
600 FORMAT(1H1* DATA CARDS FOR COIL BDRY *I3* ARE LISTED BELOW*///)

CALL DATA(KIND,1,R1,1,PHI1,1,R2,1,PHI2,1,NINT,1)

R0 = R1
PHI0 = PHI1

IF(N.EQ.1) GO TO 8

NPTS(N) = NPTS(N-1)

GO TO 9

8 NPTS(1)=1

9 M = NPTS(N)

X(M)=R1

Y(M)=PHI1

IF(IXY.EQ.1) GO TO 10

X(M) = R1 * COS(PHI1)

Y(M) = R1 * SIN(PHI1)

10 IF(KIND .EQ. 1) CALL LINE(R1,PHI1,R2,PHI2,NINT,N)

IF(KIND .EQ. 2) CALL ARC(R1,PHI1,R2,PHI2,NINT,N)

C HAS REGION BEEN CLOSED

20 IF((R2.EQ.R0).AND.(PHI2.EQ.PHI0)) GO TO 1000

R1 = R2

PHI1 = PHI2

CALL DATA(KIND,1,R2,1,PHI2,1,NINT,1)

IF(KIND .EQ. 1) CALL LINE(R1,PHI1,R2,PHI2,NINT,N)

IF(KIND .EQ. 2) CALL ARC(R1,PHI1,R2,PHI2,NINT,N)

IF(NPTS(N) .GT. IDIM) GO TO 25

GO TO 20

C DIMENSIONS HAVE BEEN EXCEEDED

25 PRINT 26, NPTS(N), IDIM

26 FORMAT(///# DIMENSIONS OF X AND Y EXCEEDED#/# NPTS = #15

#* , MAX ALLOWED IS#15 //)

STOP

1000 CONTINUE

700 CONTINUE

C CALCULATE COIL AREAS AND CURDENS

DO 18 I=1,IGRP

C FIRST AND LAST PTS ARE SAME == SUBTRACT ONE FROM NPTS

NPTS(I) = NPTS(I) - 1

CALL AREA(A,I)

CURDEN(I) = CUR(I)/A

PRINT 66, CURDEN(I)

66 FORMAT(/# CURDEN = #E15.6# A/CM2#)

18 CONTINUE

C CALCULATE Z AND ZSTAR FROM X,Y PAIRS

C ALSO OBTAIN NORMALIZING RADIUS FOR MULTIPOLES

RNORM = 10000.

IST = NPTS(IGRP)

DO 19 K= 1,IST

R = SQRT(X(K)*X(K) + Y(K)*Y(K))

IF(R.LT.RNORM) RNORM = R

Z(K) = CMPLX(X(K), Y(K))

ZS(K) = CMPLX(X(K), -Y(K))

19 CONTINUE

IF(RN.GT.0.) RNORM = RN

C SET UP EDIT GRID

PRINT 602

602 FORMAT(///# DATA CARDS READ TO SET UP GRID ARE LISTED BELOW#///)

CALL DATA(NX,1,X01,1,DELX0,1)

CALL DATA(NY,1,Y01,1,DELY0,1)

XP = X01

YP = Y01

C CALCULATE POWERS FOR MULTIPOLES

ROT = CMPLX(0.,0.)

PRINT 301

301 FORMAT(1H1 * MULTIPOLE ANALYSIS*//)

PRINT 302, RNORM

302 FORMAT(2X,*THE NORMALIZING RADIUS IS*, F10.4, *CM*//)

PRINT 303

303 FORMAT(47H BSTAR = SUM OVER N (CN*((Z0/R) ** (N-1))) //)

DO 11 M=1,10

C ANY EXPRESSION FOR MULTIPOLE ANALYSIS (IPOWER) MAY BE SURSTITUTED
C FOR THE GENERAL ONE USED BELOW

IPOWER = (NCOIL/2)*(2*M-1)

HSTAR = CMPLX(0.,0.)

DO 899 I=1,IGRP

CALL SUMMM(SUM, I, IPOWER)

HSTAR = HSTAR - CURDEN(I) * 0.1*ROT*SUM

899 CONTINUE

HSTAR = FLOAT(NCOIL) * HSTAR * RNORM

PRINT 300, IPOWER, HSTAR

300 FORMAT(/* N = *I3, 10X, * CN = * 2E15.6)

11 CONTINUE

PRINT 201

201 FORMAT(1H1,14X,*X(CM)*,5X,*Y(CM)*,12X,*BX(G)*13X,*BY(G)*,13X,*BT(G
X)*//)

DO 500 J=1,NY

DO 501 JJ=1,NX

HSTAR = CMPLX(0.,0.)

DO 499 JI = 1,NCOIL

KNT = JI + 1

ISN = (IALT)**KNT

ANG = FLOAT(JI-1)*2.*PI/FLOAT(NCOIL)

ROT = CMPLX(0.,-ANG)

Z0 = CMPLX(XP,YP)

Z0 = Z0 * CEXP(ROT)

Z0S = CMPLX(XP,-YP)

Z0S = Z0S * CEXP(ROT)

X0 = REAL(Z0)

Y0 = AIMAG(Z0)

DO 550 I=1,IGRP

CALL SUMS(SUM,I)

IF(CABS(SUM) .EQ. 0.) GO TO 77

SUM = SUM * CURDEN(I) * CEXP(ROT) * 0.1

HSTAR = HSTAR + SUM*FLOAT(ISN)

550 CONTINUE

499 CONTINUE

RX = REAL(HSTAR)
BY = -AIMAG(HSTAR)
B = SQRT(BX*BX + BY*BY)

PRINT 200, XP, YP, BX, BY, B
200 FORMAT(/10X, 2F10.4, 5X, 3(E15.6, 3X))

GO TO 78
77 PRINT 202, XP, YP
202 FORMAT(/10X, 2F10.4, 10X* EDIT PT IS ON COIL BDRY*)
78 XP = XP + DELX0
501 CONTINUE

XP = X01
YP = YP + DELY0
500 CONTINUE

END

SUBROUTINE AREA(A,I)

COMMON/BDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)
COMMON/GRID/X0,Y0,Z0,Z0S

IF(I.EQ.1) GO TO 1
IS = NPTS(I-1) + 1
GO TO 2

1 IS = 1
2 IE = NPTS(I)

AS1 = X(IS)*(Y(IS + 1) - Y(IE))
AS2 = X(IE)*(Y(IS) - Y(IE - 1))
A = AS1 + AS2

JS = IS + 1
JE = IE - 1

DO 3 J= JS, JE

AS = X(J) * (Y(J+1) - Y(J-1))
A = A + AS

3 CONTINUE

A = A * 0.5

RETURN
END

SUBROUTINE SUMS(SUM,I)

COMMON/BDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)

COMMON/GRID/X0,Y0,Z0,Z0S

COMPLEX Z, ZS, Z0,Z0S

COMPLEX SUM, AS

SUM = CMPLX(0.,0.)

AS = CMPLX(0.,0.)

IF(I.EQ.1) GO TO 1

IS = NPTS(I-1) + 1

GO TO 2

1 IS = 1

2 IE = NPTS(I)

IF(CABS(Z(IS) - Z0) .EQ. 0. .OR. CABS(Z(IE)-Z0) .EQ. 0.) GO

* TO 8

SUM = ((ZS(IS) - Z0S)/(Z(IS)-Z0))*(Z(IS+1)-Z(IE))

AS = ((ZS(IE)-Z0S)/(Z(IE)-Z0))*(Z(IS)-Z(IE-1))

SUM = SUM+AS

JS = IS + 1

JE = IE - 1

DO 3 J = JS, JE

IF(CABS(Z(J) - Z0) .EQ. 0.) GO TO 8

AS = ((ZS(J)-Z0S)/(Z(J)-Z0))*(Z(J+1)-Z(J-1))

SUM = SUM + AS

3 CONTINUE

SUM = SUM * 0.5

RETURN

8 SUM = CMPLX(0.,0.)

RETURN

END

SUBROUTINE SUMMM(SUM,I,IPOWER)

COMMON/BDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)

COMMON/GRID/X0,Y0,Z0,Z0S

COMMON/MPOL/R

COMPLEX Z, ZS, Z0,Z0S

COMPLEX SUM, AS

SUM = CMLX(0.,0.)

AS = CMLX(0.,0.)

IF(I.EQ.1) GO TO 1

IS = NPTS(I-1) + 1

GO TO 2

1 IS = 1

2 IE = NPTS(I)

SUM = ((ZS(IS)/R)/((Z(IS)/R)**IPOWER))* (Z(IS+1)/R-Z(IE)/R)

AS = ((ZS(IE)/R)/((Z(IE)/R)**IPOWER))* (Z(IS)/R-Z(IE-1)/R)

SUM = SUM+AS

JS = IS + 1

JE = IE - 1

DO 3 J = JS, JE

AS = ((ZS(J)/R)/((Z(J)/R)**IPOWER))* (Z(J+1)/R-Z(J-1)/R)

SUM = SUM + AS

3 CONTINUE

SUM = SUM * 0.5

RETURN

END

```
SUBROUTINE LINE(R1,PHI1,R2,PHI2,NINT,N)  
COMMON/RDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)  
COMMON/GRID/X0,Y0,Z0,Z0S  
COMMON/XY/IXY
```

```
IEND = NPTS(N) + NINT  
X(IEND) = R2  
Y(IEND) = PHI2  
IF(IXY.EQ.1) GO TO 10  
X(IEND)=R2*COS(PHI2)      $ Y(IEND)=R2*SIN(PHI2)
```

```
10 NN = NPTS(N)  
XDIFF = (X(IEND) - X(NN) )/FLOAT(NINT)  
YDIFF = (Y(IEND) - Y(NN) )/FLOAT(NINT)  
NSTART = NPTS(N) + 1  
IK = IEND - 1  
DO 1 I=NSTART,IK
```

```
X(I) = X(I-1) + XDIFF  
Y(I) = Y(I-1) + YDIFF
```

```
1 CONTINUE
```

```
NPTS(N) = IEND  
RETURN  
END
```

```

SUBROUTINE ARC(R1,PHI1,R2,PHI2,NINT,N)
COMMON/RDRY/X(2000),Y(2000),NPTS(5),Z(2000),ZS(2000)
COMMON/GRID/X0,Y0,Z0,Z0S
COMMON/XY/IXY

```

```

IF(IXY.EQ.1) GO TO 200

```

```

IF(R1.EQ.R2) GO TO 2

```

```

PTS NOT ON ARC

```

```

PRINT 100, R1,R2

```

C

```

100 FORMAT(///# FIRST AND LAST PTS OF SEGMENT DONT HAVE SAME RADIUS#

```

```

** R1 = *F10.4,* R2 = *F10.4///)

```

```

STOP

```

```

200 PRINT 201

```

```

201 FORMAT(///# MUST USE (R,THETA) COORD SYSTEM TO DESCRIBE ARCS, (I

```

```

*XY MAY NOT BE 1)*//)

```

```

STOP

```

```

2 CONTINUE

```

```

IEND = NPTS(N) + NINT

```

```

X(IEND)=R2*COS(PHI2) $ Y(IEND)=R2*SIN(PHI2)

```

```

PHIDIF = (PHI2 - PHI1)/FLOAT(NINT)

```

```

NSTART = NPTS(N) + 1

```

```

PHI = PHI1

```

```

IK = IEND - 1

```

```

DO 1 I=NSTART,IK

```

```

PHI = PHI+PHIDIF

```

```

X(I)=R2*COS(PHI)

```

```

Y(I)=R2*SIN(PHI)

```

```

1 CONTINUE

```

```

NPTS(N) = IEND

```

```

RETURN

```

```

END

```


SUBROUTINE DATA(X1,N1,X2,N2,X3,N3,X4,N4,X5,N5,X6,N6,X7,N7)

C 53

COMMON/REDMAN/FLTING,TYPX,PRINX,VALUED,CFLAG,REPEAT,RC,VALUE,

* STAR,K,R,DOLLAR,S,C
* COLCT,COLMAX,TYPTAB(7),DECIDE(10),XPA(4),XPB(10),XPC(10),
* BLANK,BITS,MINUS,A,ALLBIT,TOOBIG,CARDCT

COMMON /CCNAP/ I
INTEGER TYPX,PRINX,DELTAJ,VALU,GOTX,CFLAG,REPEAT,RC,STAR
INTEGER K,R,DOLLAR,S,C,CARDCT,COLCT

\$\$\$\$\$

INTEGER COLMAX

\$\$\$\$\$

DATA STAR,K,R,DOLLAR,S,C,COLCT,COLMAX/0000047,0000013,0000022,

* 0000053,0000023,0000003,82,81/

\$\$\$\$\$* 0000053,0000023,0000003,74,73/

DATA TYPTAB/0005050705040505050505050705050505050707A,

* 05050505050505010101B,01010101010101030307B,

* 07070707070607020000B,0,0/

DATA DECIDE/005030212120114130000,005041515151515130000,

* 006151515151515130000,00613131313131313130000,

* 005061707171717130000,006171707171717130000,

* 011171013130713130000,011171313131013130000,

* 011171717171717130000,021212112122121130000/

DATA XPA/1.0,1.0,1.0/

DATA XPB/1.0,1.0E10,1.0E20,1.0E30,1.0E40,1.0E50,1.0E60,1.0E70,

* 1.0E80,1.0E90/

DATA XPC/1.0,1.0E1,1.0E2,1.0E3,1.0E4,1.0E5,1.0E6,1.0E7,1.0E8,

* 1.0E9/

DATA BLANK,BITS,MINUS,A,ALLBIT,TOOBIG/0000055,07700000000000000,

* 0000046,0000001,03777000000000000000,00001000000000000000/

DATA CARDCT/0/

CFLAG=0

REPEAT=0

COLCT=COLMAX+1

\$\$\$\$\$ COLCT=74

CALL CCNARG

IF(I.LE. 1)GO TO 103

CALL DATAI(X1,N1)

IF(I.LE. 3)GO TO 103

CALL DATAI(X2,N2)

IF(I.LE. 5)GO TO 103

CALL DATAI(X3,N3)

IF(I.LE. 7)GO TO 103

CALL DATAI(X4,N4)

IF(I.LE. 9)GO TO 103

CALL DATAI(X5,N5)

IF(I.LE.11)GO TO 103

CALL DATAI(X6,N6)

IF(I.LE.13)GO TO 103

CALL DATAI(X7,N7)

103 PRINT 301,CARDCT,COLCT

301 FORMAT(1H 85X23HDATA POSITIONED AT CARDI4,8H, COLUMNI3)

\$\$\$\$\$ FORMAT(24H DATA POSITIONED AT CARDI4,8H, COLUMNI3)

RETURN

END

ASCENTF SUBROUTINE IDBYTE(I,J)

C 54

RDBYTE

PS		000000000000
PS		000000000000
PS		000000000000
PS		000000000000
PS		000000000000
SA2=B2	.X2=J	56220
SX3=1	.X3=1	7130 000001
IX4=X2-X3	.X4=J-1	37423
SX5=10	.X5=10	7150 000012
PX7=B0,X4	.X7=FLOAT(J-1)	27704
PX0=B0,X5	.X0=10.	27005
NX0=B0,X0		24000
FX7=X7/X0	.X7=(J-1)/10	44770
UX7=B7,X7		26777
LX7=B7,X7		22777
PX7=B0,X7		27707
UX1=B0,X7	.X1=N	26107
SA2=X1+B1	.X2=N+I=I(N+1)=II	53211
SA3=B2	.X3=J	56320
PX7=B0,X5	.X7=10.	27705
PX0=B0,X1	.X0=FLOAT(N)	27001
DX7=X7*X0	.X7=10*N	42770
UX7=B0,X7		26707
IX4=X3-X7	.X4=J-10*N	37437
SX1=6	.X1=6	7110 000006
PX7=B0,X1	.X7=6.	27701
PX0=B0,X4	.X0=FLOAT(J-10*N)	27004
DX7=X7*X0	.X7=6*(J-10*N)	42770
UX6=B0,X7	.X6=MKK	26607
SB3=X6	.B3=MKK	63360
SX3=77B	.X3=A MASK	7130 000077
LX6=B3,X2	.SHIFT II MKK PLACES	22632
BX6=X6*X3	.MASK THE LOWER 6 BITS	11663
JP RDRYTE	.RETURN	0200 L00001
END		

SUBROUTINE DATAI(X,N)

C 55

COMMON/REDMAN/FLTNG,TYPX,PRINX,VALUED,CFLAG,REPEAT,RC,VALUE,

* STAR,K,R,DOLLAR,S,C,

* COLCT,COLMAX,TYPTAB(7),DECIDE(10),XPA(4),XPB(10),XPC(10),

* BLANK,BITS,MINUS,A,ALLBIT,TOOBIG,CARDCT

INTEGER TYPX,PRINX,DELTAJ,VALU,GOTX,CFLAG,REPEAT,RC,STAR

INTEGER K,R,DOLLAR,S,C,CARDCT,COLCT

DIMENSION X(1)

EQUIVALENCE (VALUE,VALU)

IF(N.LE.0)GO TO 101

J=1

DELTAJ=1

DO 100 I=1,N

1 IF(CFLAG.EQ.0)GO TO 2

CFLAG=0

GO TO 10

2 IF(REPEAT.EQ.0)GO TO 3

RC=RC-1

IF(RC.GT.0)GO TO 10

REPEAT=0

3 CALL RDNUM(VALUE)

IF(TYPX)3,4,10

4 GOTX=7

IF(VALU.EQ.STAR)GOTX=1

IF(VALU.EQ.K)GOTX=2

IF(VALU.EQ.R)GOTX=3

IF(VALU.EQ.DOLLAR)GOTX=4

IF(VALU.EQ.S)GOTX=5

IF(VALU.EQ.C)GOTX=6

GO TO (11,12,13,14,101,16,3),GOTX

11 CALL RDNUM(J)

IF(TYPX.EQ.1.AND.FLTNG.EQ.0.0)GO TO 3

GO TO 17

12 CALL RDNUM(DELTAJ)

IF(TYPX.EQ.1.AND.FLTNG.EQ.0.0)GO TO 3

GO TO 17

13 CALL RDNUM(RC)

VALUE=VALUED

REPEAT=1

IF(TYPX.EQ.1.AND.FLTNG.EQ.0.0.AND.RC.GT.0)GO TO 2

GO TO 17

14 PRINX=1

CALL RDNUM(VALUE)

IF(VALU=DOLLAR)14,1,14

16 VALU=I-1

CFLAG=1

GO TO 101

C ERROR***

17 CONTINUE

PRINT 302,CARDCT,COLCT

302 FORMAT(15H ERROR AT CARD 14,5H,COL 12)

C ERROR***

10 X(J)=VALUE

VALUED=VALUE

J=J+DELTAJ

100 CONTINUE

101 RETURN

END

SUBROUTINE RDNUM(ISIT)

C 56

```

COMMON/REDMAN/FLTNG,TYPX,PRINX,DUMMY(11),
* COLCT,COLMAX,TYPTAB(7),DECIDE(10),XPA(4),XPR(10),XPC(10),
* BLANK,BITS,MINUS,A,ALLBIT,TOOBIG,CARDCT
INTEGER TYPX,PRINX,COLCT,COLMAX,BLANK,BITS,MINUS,A,ALLBIT
INTEGER TOOBIG,CARDCT
INTEGER VALUE,STAGE,CHARAC,CHTYPE,OLDCH,      WORD,EXPON
INTEGER DIGITS,ESIGN,NSIGN
DIMENSION BUFFER(8)
EQUIVALENCE(X,I)

```

1

STAGE=1

N=0

DIGITS=0

EXPON=0

WORD=0

TESDC=0

NSIGN=0

ESIGN=0

FLTNG=0.

11

OLDCH=CHARAC

IF (COLCT-COLMAX) 3,6,2

6

CHARAC=BLANK

GO TO 7

2

READ 500,BUFFER

500

FORMAT(8A10)

\$\$\$\$\$

PRINT 510,BUFFER

510

FORMAT(1H 8A10)

\$\$\$\$\$

CARDCT=CARDCT+1

5

COLCT=1

3

CHARAC=IDBYTE (BUFFER,COLCT)

7

COLCT=COLCT+1

CHTYPE=IDBYTE (TYPTAB,CHARAC+1)

IF (CHTYPE.EQ.0)CHTYPE=8

130

STAGE=IDBYTE (DECIDE (STAGE),CHTYPE)

GO TO(11,12,11,11,15,16,17,18,19,20,101,102,103,104,105,106,107),

* STAGE

12

IF (CHARAC.EQ.MINUS)NSIGN=1

GO TO 11

15

IF (CHARAC.LT.27.OR.CHARAC.GT.36)GO TO 25

24

I=N+N

I=I+I+N

I=I+I+(CHARAC-27)

IF (I.GT.TOOBIG)GO TO 25

26

N=I

GO TO 11

25

DIGITS=DIGITS+1

GO TO 11

16

DIGITS=DIGITS-1

C

DECIMAL POINT ENCOUNTERED. SET FOR FLOATING POINT.

FLTNG=1.

GO TO 15

C

E ENCOUNTERED. SET FOR FLOATING POINT.

17

FLTNG=1.

GO TO 11

18

IF (CHARAC.EQ.MINUS)ESIGN=1

GO TO 11

19

EXPON=EXPON+10+(CHARAC-27)

GO TO 11

```

20 I=WORD.AND.BITS
   IF(I)11,28,11
28 WORD=64*WORD
   WORD=WORD.OR.CHARAC
   GO TO 11
C   SCAN COMPLETE.PREPARE FOR EXIT.
101 I=COLCT-1
   IF(PRINX.EQ.0)PRINT 502,I,CHARAC,CARDCT
   GO TO 1
502 FORMAT(19H INPUT ERROR COLUMN,I3,11H CHARACTER ,02.6H, CARD,I4)
103 COLCT=COLCT-1
   CHARAC=OLDCH
102 ISIT=CHARAC
   TYPX=0
114 RETURN
105 COLCT=COLCT-1
104 DIGITS=(DIGITS+(1-FSIGN-ESIGN)*EXPON
   TYPX=1
   IF(NSIGN.NE.0)N=-N
108 IF(FLTING)124,109,124
109 ISIT=N
   IF(DIGITS)126,114,112
112 DO 120 I=1,DIGITS
   N=ISIT+ISIT
   N=N+N+ISIT
120 ISIT=N+N
   IF(ISIT.LT.TOOBIG)GO TO 114
   IF(PRINX.EQ.0)PRINT 503
503 FORMAT(18H INTEGER TOO LARGE)
   ISIT=ALLBIT
   GO TO 114
126 DIGITS=-DIGITS
   DO 125 I=1,DIGITS
125 ISIT=ISIT/10
   GO TO 114
124 X=N
   K=DIGITS
   IF(DIGITS.LT.0)K=-DIGITS
   J1=K/100
   K=K-100*J1
   J2=K/10
   K=K-J2*10
   Y=XPA(J1+1)*XPB(J2+1)*XPC(K+1)
   IF(DIGITS)115,116,117
115 X=X/Y
116 ISIT=I
   GO TO 114
117 X=X*Y
   GO TO 116
107 COLCT=COLCT-1
106 TYPX=-1
   ISIT=WORD
   GO TO 114
C   DIGIT=NUMBER OF DIGITS TO RIGHT OF THE DECIMAL POINT.
C   N=INTEGER VALUE.
C   EXPON =EXPONENT VALUE.
END

```

Sample Sextupole Magnet
See Fig. 2

DATA CARDS FOR COIL BDRY 1 ARE LISTED BELOW

2	6.	.8726646	6.	.1745329	30
1	8.	.1745329	10.		
2	8.	.8726646	30		
1	6.	.8726646	10		

CURDEN = 1.023231E+02 A/CM2

DATA CARDS READ TO SET UP GRID ARE LISTED BELOW

8	0.	1.
3	0.	1.

MULTIPOLE ANALYSIS

THE NORMALIZING RADIUS IS 6.0000CM

BSTAR = SUM OVER N (CN * ((Z0/R) ** (N-1)))

N = 3 CN = 8.175091E-06 -1.064536E+02

N = 9 CN = 4.784895E-12 -7.582603E-08

N = 15 CN = -2.960939E-06 7.711295E+00

N = 21 CN = -2.712846E-06 5.046552E+00

N = 27 CN = 5.418791E-13 -5.314642E-08

N = 33 CN = 3.503158E-06 -4.147010E+00

N = 39 CN = 4.286323E-06 -4.293479E+00

N = 45 CN = -5.827757E-13 9.476247E-08

N = 51 CN = -6.424398E-06 4.920979E+00

N = 57 CN = -7.745811E-06 5.308615E+00

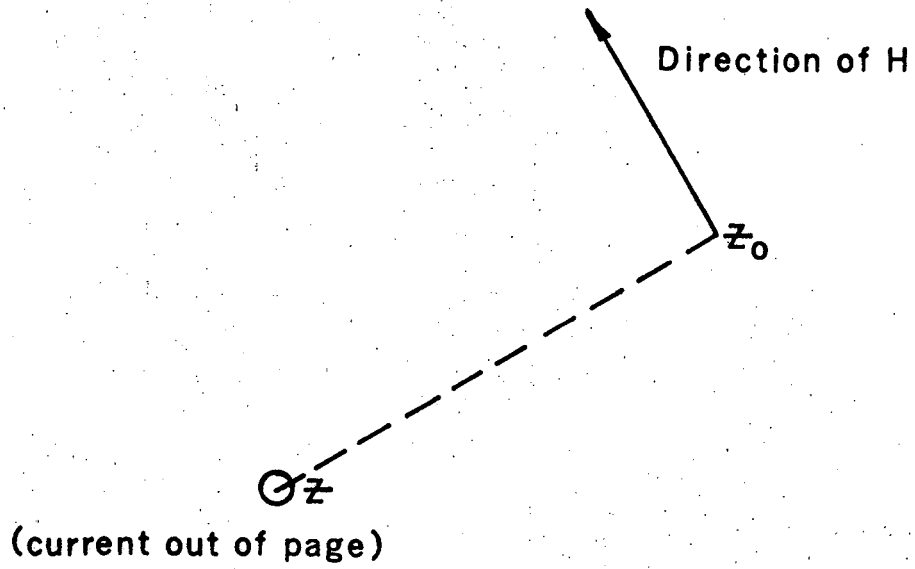
Note: The sextupole geometry described in this example (namely with each coil area equal to twice the area between successive coils) should contain no field contribution for N = 9, 27, 45, --- etc.

Also, all components should be purely imaginary numbers. These features may be noted in the above calculated results.

X (CM)	Y (CM)	BX (G)	BY (G)	BT (G)
0.	0.	-3.231719E-09	1.136868E-13	3.231719E-09
1.0000	0.	2.956433E+00	1.058579E-03	2.956433E+00
2.0000	0.	1.182577E+01	4.167218E-03	1.182577E+01
3.0000	0.	2.660847E+01	7.712695E-03	2.660847E+01
4.0000	0.	4.728758E+01	-4.854259E-03	4.728758E+01
5.0000	0.	7.330725E+01	-1.693231E-01	7.330745E+01
6.0000	0.	9.779448E+01	-3.102370E-01	9.779497E+01
7.0000	0.	1.003171E+02	4.013926E-01	1.003179E+02
0.	1.0000	-2.957656E+00	1.059564E-03	2.957656E+00
1.0000	1.0000	-1.219125E-03	-5.911968E+00	5.911968E+00
2.0000	1.0000	8.867920E+00	-1.182276E+01	1.477896E+01
3.0000	1.0000	2.364744E+01	-1.773296E+01	2.955773E+01
4.0000	1.0000	4.435790E+01	-2.366153E+01	5.027416E+01
5.0000	1.0000	7.158070E+01	-2.944211E+01	7.739919E+01
6.0000	1.0000	1.162422E+02	-2.408271E+01	1.187107E+02
7.0000	1.0000	1.145786E+02	1.647872E+01	1.157575E+02
0.	2.0000	-1.183066E+01	4.304832E-03	1.183066E+01
1.0000	2.0000	-8.874152E+00	-1.182307E+01	1.478295E+01
2.0000	2.0000	-3.825049E-03	-2.364743E+01	2.364743E+01
3.0000	2.0000	1.477795E+01	-3.545759E+01	3.841391E+01
4.0000	2.0000	3.534466E+01	-4.724030E+01	5.899907E+01
5.0000	2.0000	6.107473E+01	-5.983869E+01	8.550317E+01
6.0000	2.0000	7.690362E+01	-4.029403E+01	8.682036E+01
7.0000	2.0000	6.581025E+01	3.264225E+01	7.346091E+01

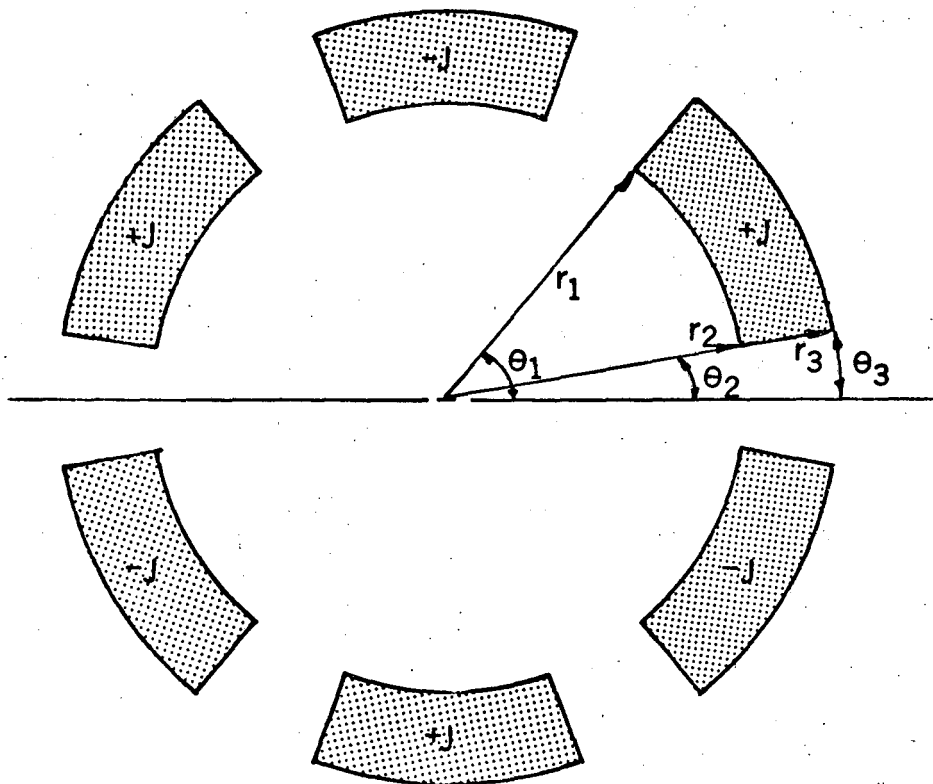
*This work was done under the auspices of the U. S. Atomic Energy Commission.

1. R. A. Beth, "Some Extensions of Complex Methods For Two-Dimensional Fields"; Proceedings of the Sixth International Conference on High Energy Accelerators, Cambridge Electron Accelerator, 1967.
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Infinite straight current filament
Fig. 1

XBL 688 4902



Model sextupole Geometry
Fig. 2

XBL 688 4903

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