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Recent Work

Title

Time Dependent Focusing

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Time Dependent Focussing

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VNL Physics Meeting

Background

Focussing Concept

NDCX II

Fusion Driver

Background

The VNL program now hinges on beam manipulations similar to those of light ion fusion:

neutralized drift compression
neutralized final focus

To be attractive WDM and fusion system concepts have evolved:

low ion kinetic energy
low ion mass
high line charge density
modular accelerators
few (~20) drift lines
high longitudinal compression
highly stripped ions

→ Perveance $> 10^3$ doesn't allow quadrupole transport in vacuum

Perveance may actually exceed unity in final focus!

Solenoidal transport in vacuum can handle much larger permeance, but there is a space charge limit:

$$\lambda < \left(\frac{10 \mu\text{C}}{\text{m}}\right) \left(\frac{B}{10\text{T}}\right)^2 \left(\frac{q}{10\text{cm}}\right)^2 \left(\frac{133}{A/\text{s}}\right)$$

+ high space charge potential:

$$\Delta\phi_{\text{center-to-edge}} = (90\text{kV}) \left(\frac{1}{10 \mu\text{C}/\text{m}}\right)$$

• Use electron neutralization in solenoids - from plasma or co streaming

Not yet clear this really works when electrons must cross field lines

{ If a design did not use neutralization, Grant would improve it by increasing λ so it did require neutralization

(4)

Some consequences of neutralization:

Momentum tilt is not removed by space charge prior to final focus

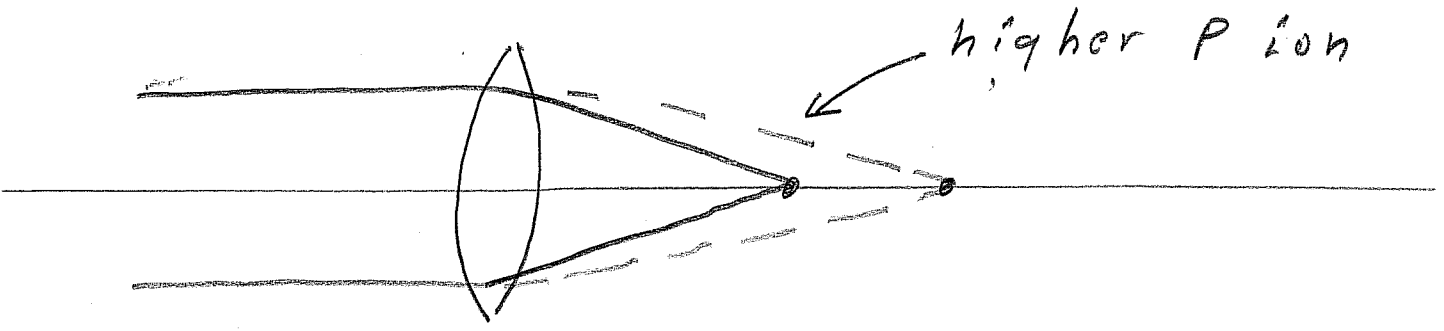
Strong quadrupoles probably can't be used - problem for electron flow

Solenoids may be "ok" - what happens to electrons when beam does not follow field lines?

Large tilts ($\Delta p/p_0 \sim \pm .05$) for large compression and short drift lengths

A static final focus system has a severe chromatic aberration ("second order")

The focal spot position depends on momentum



parallel-to-point focus

(We are actually considering thick solenoidal lenses)

Random momentum spread blows up spot radius - cant fix with SLC type scheme using sextupoles

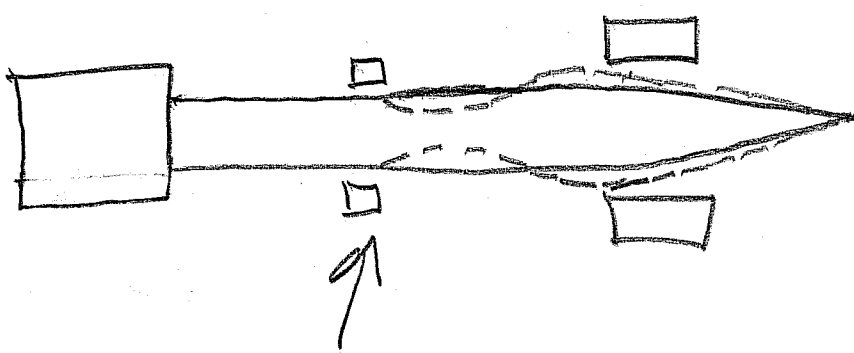
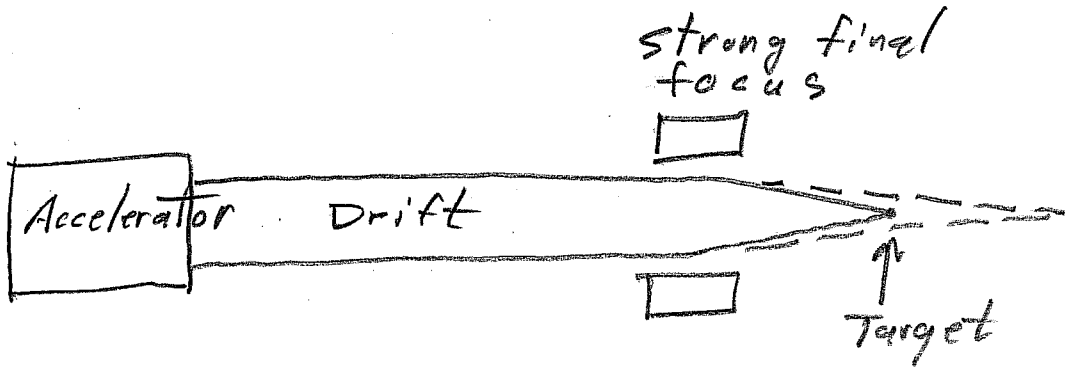
→ Must have sufficiently cold beam before compression

Systematic tilt moves focal spot - blows up the mean spot radius at fixed position

But systematic tilt can be compensated by time dependant upstream lenses (TDL)

$$(\text{spot motion from tilt}) + (\text{spot motion from TDL}) = 0$$

Focussing Concept (Drawing by GL)



time dependent
lens for finite ΔP

Make ions with $\Delta P > 0$ enter
final focus with a finite
converging angle

~ One period of envelope
oscillation from pulsed lens
to final focus

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There is a connection between beam radius in drift and drift distance from pulsed lens to final focus !!

$$\frac{d^2 a}{dz^2} = - \left[\frac{B_d}{z(B\rho)} \right]^2 a + \frac{E^2}{a^3}$$

B_d = drift field

$(B\rho) = \beta \gamma \frac{Mc}{qe} = \text{ion rigidity}$

$a = \text{beam radius}$

equilibrium $a = \sqrt{\frac{z(B\rho) E}{B_d}} \sim \sqrt{\frac{M E_n}{\beta \gamma B_d}}$

envelope oscillation:

$$\frac{d^2 \delta a}{dz^2} = - \left[\frac{B_d}{z(B\rho)} \right]^2 \delta a - \frac{3E^2}{a^4} \delta a$$

$$\delta a \sim \cos(kz)$$

$$k^2 = 4 \left[\frac{B_d}{z(B\rho)} \right]^2 \quad (= 4 k_p^2)$$

set $k = 2\pi / L_{\text{drift}}$ Pulse Tail

$$\rightarrow B_{\text{drift}} = k(B\rho) = \frac{2\pi(B\rho)}{L_{\text{drift}}}$$

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Summary of concept:

Pulse head goes straight through

Pulse tail gets inward envelope
kick + goes through one
envelope oscillation

In between parts of the
pulse get kicked and have
less than one oscillation

For a given final strong
lens find out how hard
each part of pulse must
be kicked to compensate
focal position motion

This is not a linear
dependence on Δp

→ Find how the pulsed field
must be ramped & see
if it is possible

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NDCX II Example

$$\text{Ion} = \text{Na}^{+9} \quad (A = 23)$$

$$\begin{aligned} E_{\text{head}} &= 20 \text{ MeV} \\ E_{\text{tail}} &= 24 \text{ MeV} \end{aligned} \quad \left. \vphantom{\begin{aligned} E_{\text{head}} \\ E_{\text{tail}} \end{aligned}} \right\} \text{tilt} \approx .095$$

$$\epsilon_n = 2.3 \times 10^{-6} \text{ m-r}$$

$$L_{\text{drift}} = 4.0 \text{ m}$$

$$B_{\text{drift}} = .59 \text{ T}$$

$$d_{\text{drift}} = .0079 \text{ m}$$

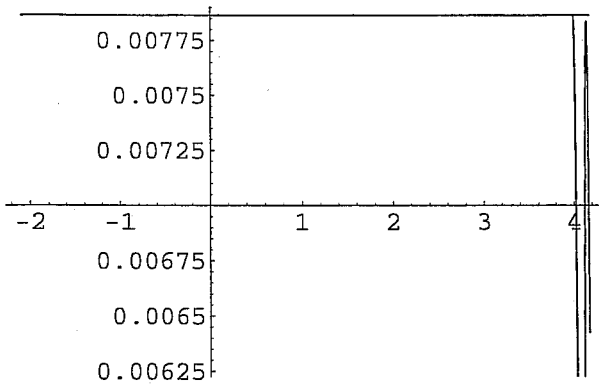
$$B_{\text{final focus}} = 15 \text{ T}$$

$$\text{Pulse Magnet length} = .2 \text{ m}$$

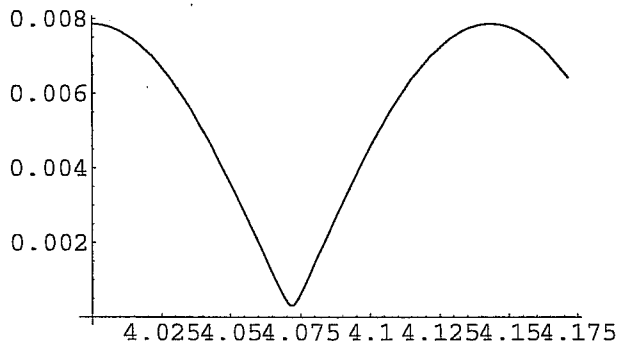
$$\text{Spot radius of head} = .00031 \text{ m}$$

$$\text{Cone Half angle of FF} \approx 150 \text{ mrad}$$

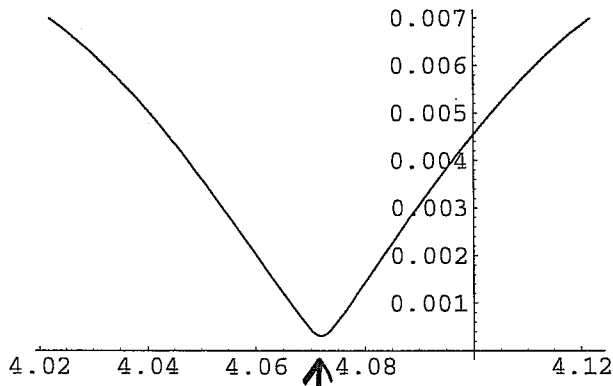
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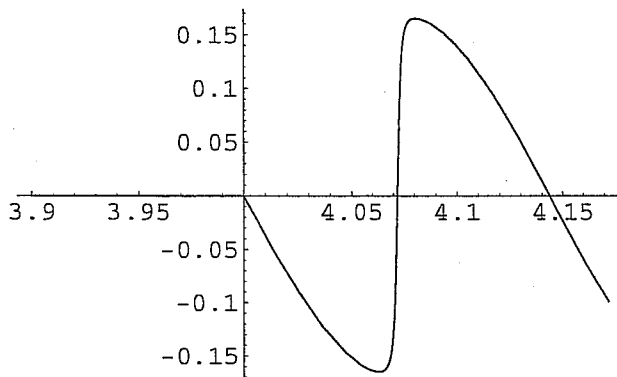
envelope radius



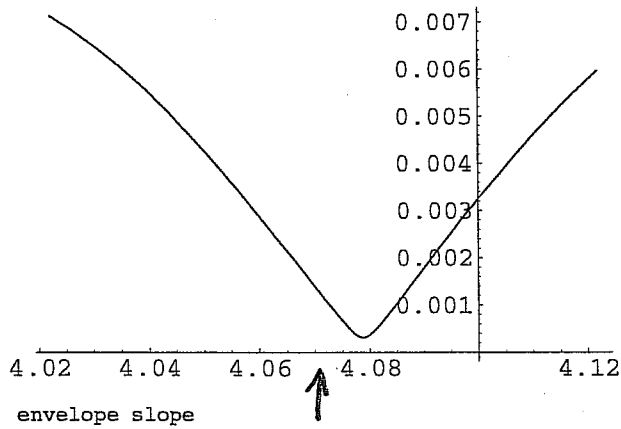
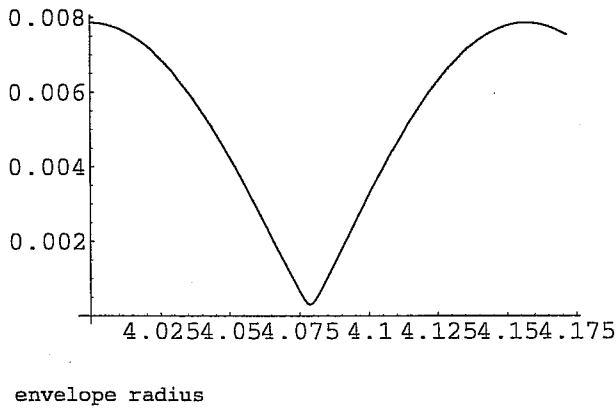
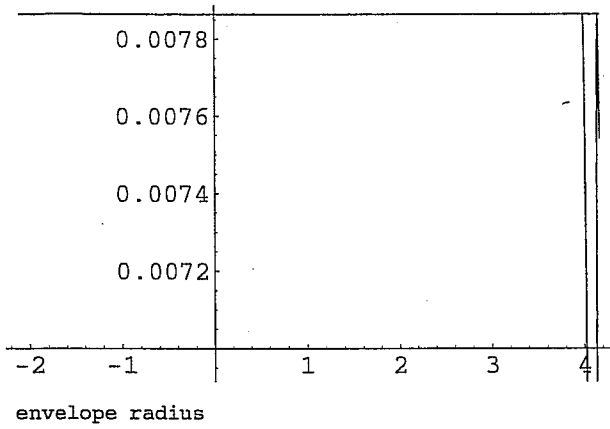
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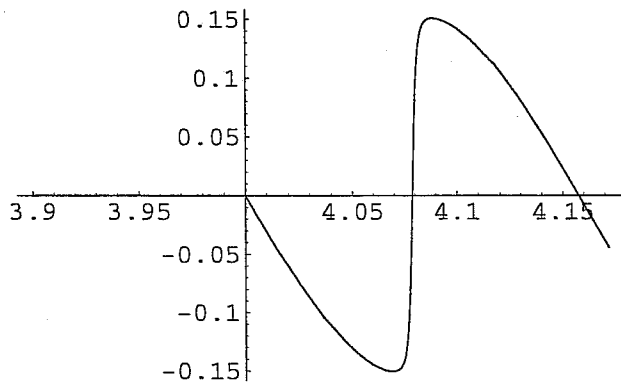
envelope slope



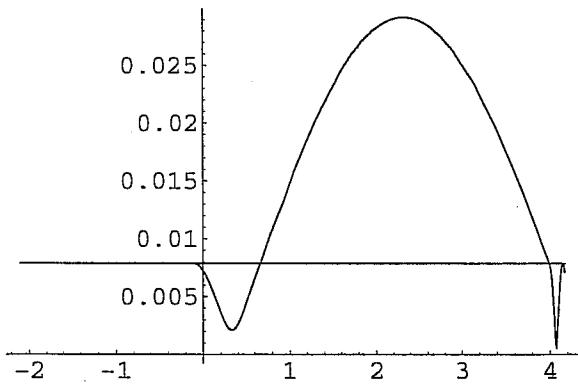
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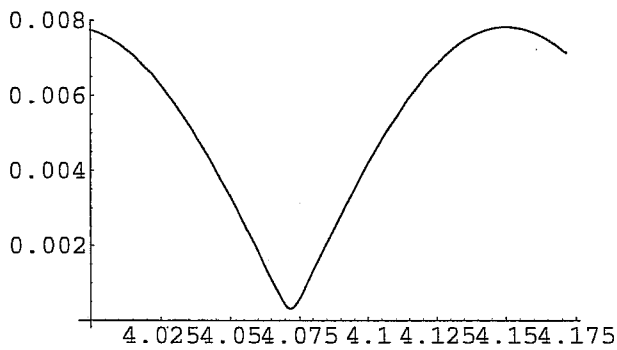
Uncompensated Envelope



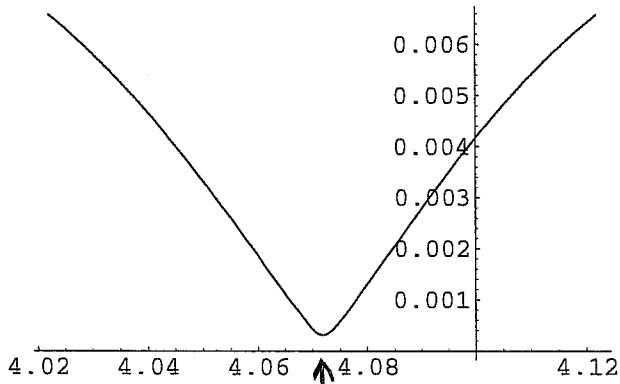
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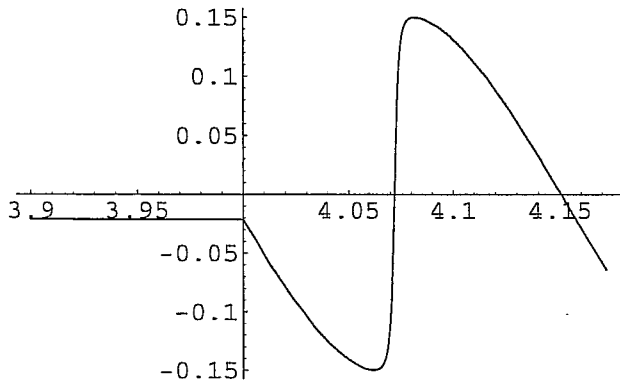
envelope radius



envelope radius



envelope slope



Compensated
Envelope

NDCX-II Pulsed Fields :

% AP	B _{pulsed}	spot
0	0 (T)	.000 309 (m)
20	.777	.000 369
40	1.166	.000 392
60	1.503	.000 387
80	1.882	.000 359
100	2.385	.000 311

We have compensated spot position

If we minimize spot radius at target we can get a smaller spot!

- An increase in beam radius entering final focus can decrease the spot more than the shift of focus increases it. —

Driver Example

Ion = K^{+19} (A=39)

$E_{head} = 611 \text{ MeV}$
 $E_{tail} = 739 \text{ MeV}$ } tilt $\approx .1$

$E_n = 5 \times 10^{-6} \text{ m-r}$

$L_{drift} = 400 \text{ m}$

$B_{drift} = 1.0203 \text{ T}$

$a_{drift} = .056 \text{ m}$

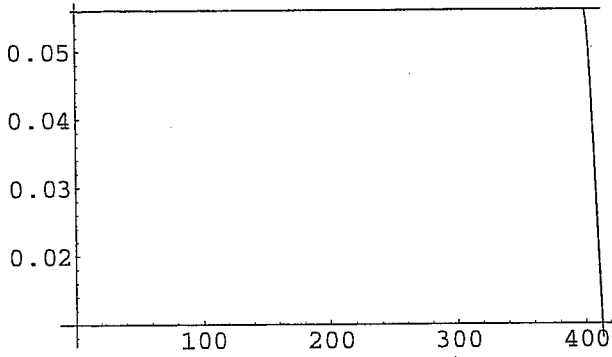
$B_{final \text{ focus}} = .3 \text{ T}$

Pulsed magnet = 1.0 m
length

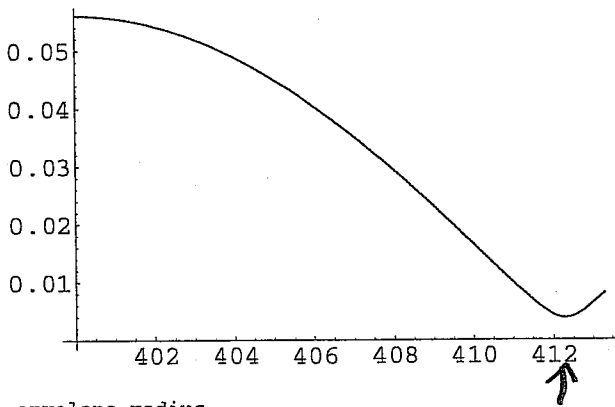
spot radius
of head = .00379 m

Cone Half
angle of FF $\approx .006$ radians

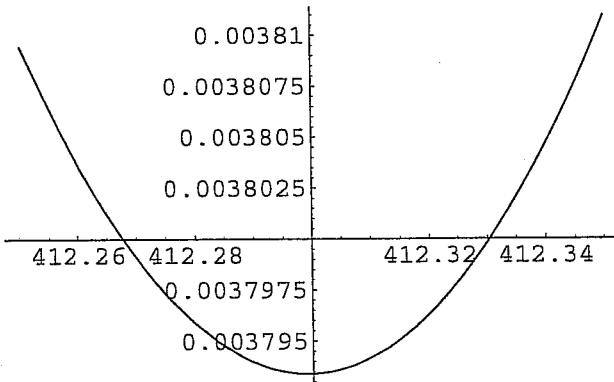
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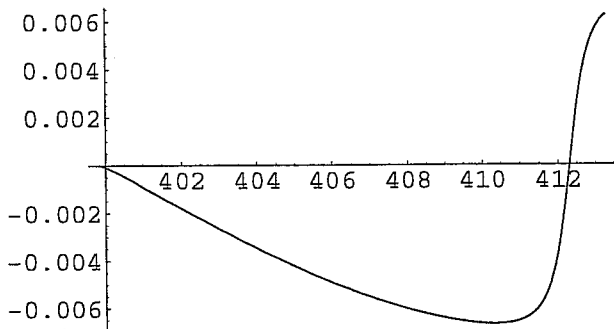
envelope radius



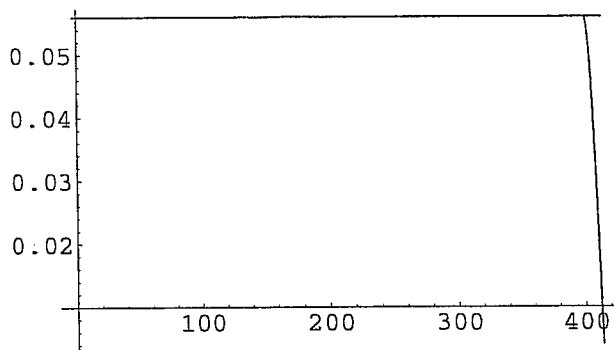
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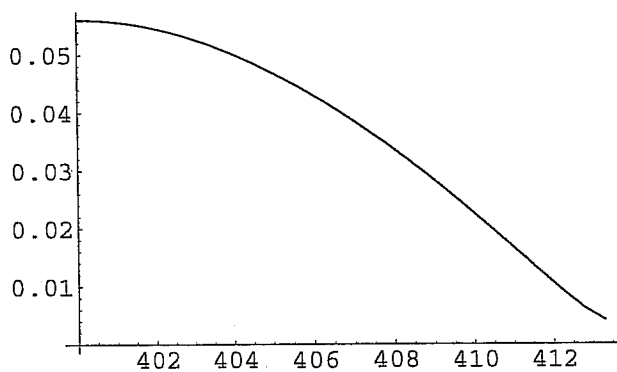
envelope slope



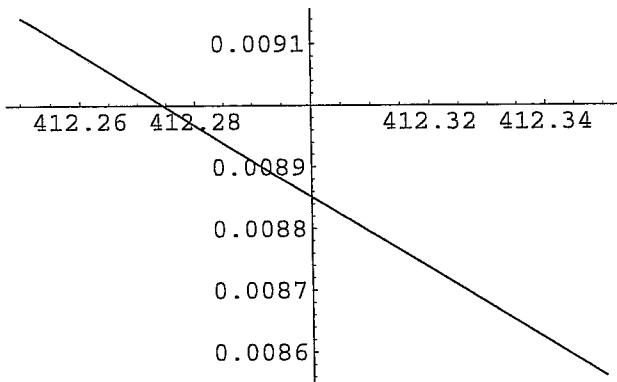
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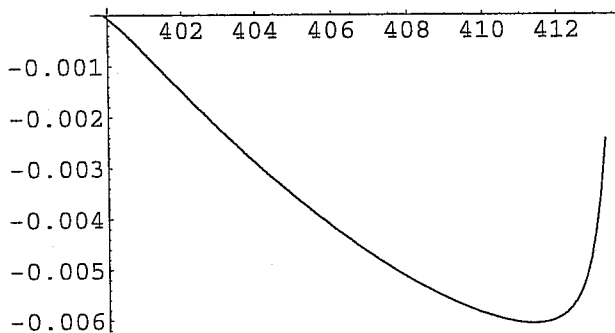
envelope radius



envelope radius

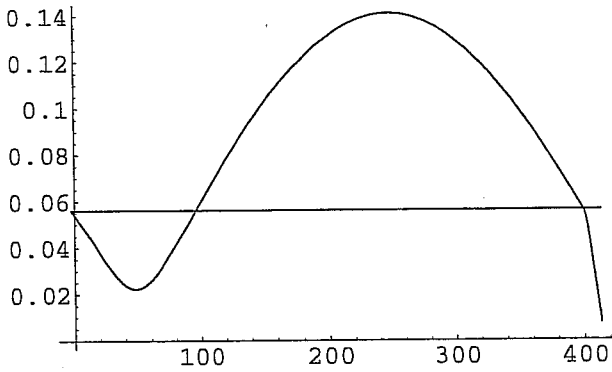


envelope slope

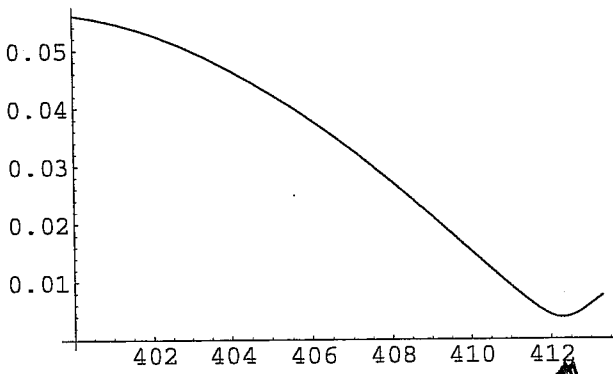


Uncompensated Envelope

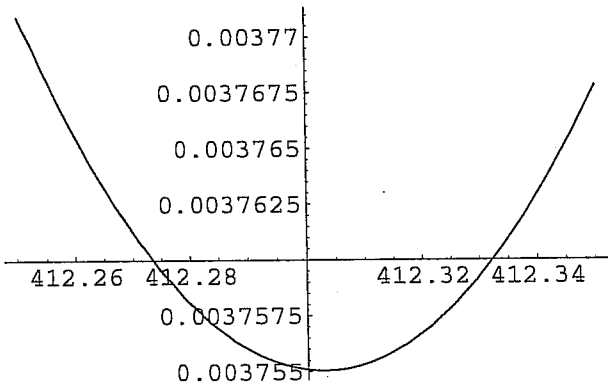
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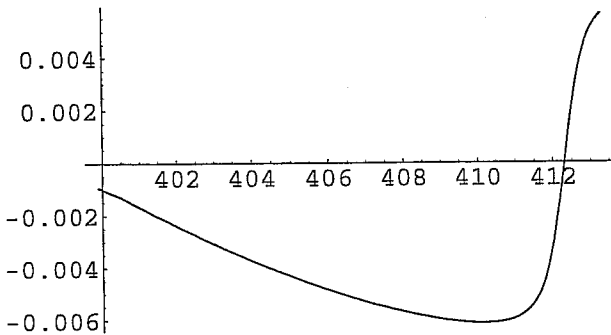
envelope radius



envelope radius



envelope slope



Compensated
Envelope

Note We need small cone angle to prevent very large amplitude envelope oscillations

Driver Pulsed Fields:

% AP	B_{pulsed}	d_{spot}
0	0 (T)	.00379(m)
20	.132	.00426
40	.183	.00443
60	.224	.00438
80	.266	.00415
100	.314	.00375