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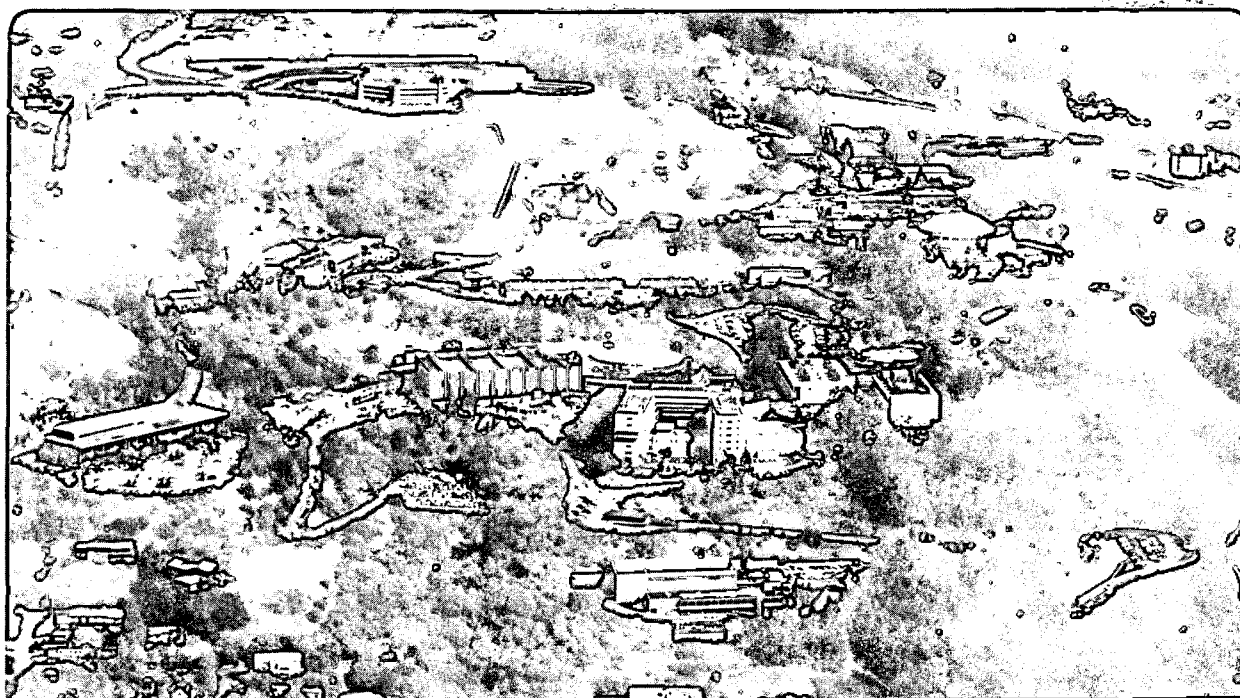
A COMPUTATIONAL MODEL FOR CRITICAL FLOW THROUGH
INTERGRANULAR STRESS CORROSION CRACKS

V.E. Schrock, S.T. Revankar, S.Y. Lee,
and C.-H. Wang

August 1986

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A Computational Model for
Critical Flow Through Intergranular Stress
Corrosion Cracks

by

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ABSTRACT

The presence of intergranular stress corrosion cracks (IGSCC) in thermal stressed zones in stainless steel piping and associated components is of much concern in reactor safety. The prediction of leak rates through the cracks is important in assessing the plant reliability. An analytical model has been developed to predict flow rates of initially subcooled or saturated water through these cracks. The model assumes the flow in the crack to be homogeneous and in thermal equilibrium. The crack geometry was idealized as a convergent straight slit of constant gap thickness. The fluid is assumed to enter the crack without separation. The one dimensional model accounts for the changing cross sectional area in the flow direction. The effects of wall friction, expansions/contractions and tortuosity of the actual flow path are lumped into an equivalent friction. The numerical scheme developed for the model solution has been programmed into a Fortran computer code called SOURCE. A companion subroutine STEAM provides the saturated fluid properties. Inputs to SOURCE are the upstream stagnation pressure and temperature, the crack geometry specification, and the equivalent friction factor.

SOURCE has been assessed against the experimental data obtained in the Battelle Columbus Laboratories (BCL) study using actual crack specimens. From a parametric study of the BCL data using SOURCE, a subcooling correction factor was developed to modify source predictions. A general methodology was developed for predicting a lumped friction factor (includes effects of bends, contractions and expansions) for use in SOURCE. Corrected SOURCE predictions using the general friction factor methodology agree well with BCL data. This method is recommended for estimating leak rates and assessing the leak-before-break criterion.

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NOMENCLATURE

Symbol	Description	Dimension
A	Cross-sectional area of crack	L^2
a	Half of entrance crack width	L
C	Subcooling Correction Factor	-
c	Half of exit crack width	L
D	Equivalent hydraulic diameter	L
E	Modulus of elasticity	$ML^{-3}t^{-2}$
f	Friction factor	-
g_c	Gravitational conversion factor	Lt^{-2}
G	Mass flux	$ML^{-2}t^{-1}$
h	Specific enthalpy	Lt^{-2}
K	Pressure loss coefficients	
L	Depth of crack	L
ℓ	Surface crack length	L
M	Mach number	
\dot{m}	Mass flow rate	Mt^{-1}
P	Pressure	$ML^{-3}t^{-2}$
P_m	Wetted perimeter	L
T	Temperature	T
t	Equivalent slit thickness	L
v	Specific volume	L^3M^{-1}
V	Fluid velocity	Lt^{-1}
w	Local width of crack	L
x	Thermodynamic quality	
z	Length variable along crack channel	L

Greek Letters

δ	Average crack thickness (Figure 3.2)	L
Δ	Difference	
ϵ	Surface roughness of crack wall	L
θ	Percentage tolerance	
ρ	Density	$L^{-3}M$
Σ	Summation	-
σ	Tangential stress	$ML^{-3}t^{-2}$
σ_y	Yield stress	$ML^{-3}t^{-2}$
τ_w	Wall shear stress	$ML^{-3}t^{-2}$

Subscripts

o	Stagnation
l	Entrance
a	Sound speed
c	Choking condition
e	Crack exit
f	Saturated liquid
fg	Saturated liquid-vapor difference
fl	Evaluated at flashing plane
g	Saturated vapor
i	Grid point
l	Liquid or flashing point
s	Evaluated at constant entropy
T	Total
sat	Evaluated at saturation condition corresponding at the local pressure
sub	Subcooling

Superscripts

-	Average
---	---------

1. INTRODUCTION

The presence of intergranular stress corrosion cracks (IGSCC) in the weld heat affected zones of types 304 and 316 stainless steel piping and associated components of commercial boiling water reactors (BWR) and steam generator tubes in pressurized water reactors (PWR) has attracted considerable attention over the past several years. [1,2]. Because of economic and safety considerations, it is highly desirable to determine if the failure of the piping system will occur in a 'leak-before-break' mode. Leak-before-break is demonstrated by establishing that postulated cracks in a pipe will be detected by leak detection methods before such cracks reach a critical size to cause unstable fracture. The ability to predict the leak rates through the cracks in piping and generator tubes before the crack size becomes critical is vital to demonstrate leak-before-break.

Most reports on the critical two-phase flow are related to flow in pipes, nozzles and orifices and there is little direct literature available on a flow rate of two-phase flow through tight cracks (crack width of less than 1 mm). Agostinelli et al. [3] studied flows of flashing water and steam through smooth annular passages of hydraulic diameters with constant area in the range of 0.15 mm to 0.43 mm. Test data were obtained with stagnation conditions of pressure ranging from 3.50 to 20.51 MPa and fluid subcoolings from 9.3°C to 67°C. Hendricks et al. [4] made a qualitative study of radially inward flow through 0.076 mm annular clearance between two glass plates using liquid nitrogen. In their experiments for inlet conditions with 10 K or larger subcoolings, flashing was seen to occur near the end of the 0.72 cm flow passage. Simoneau [5] carried out an experimental study of two-phase nitrogen flow through a rectangular slit. The test section used was 2.54 cm in length and width, with an opening dimension which was nominally 0.292 mm. He concluded that a uniform two phase flow pattern existed in most of the test runs and that vaporization was occurring at or near the exit plane, for the tests carried out with stagnation pressures up to 6.8 MPa and for four different subcoolings in the range $0.84 < T_R < 1.03$, $T_R = T_O/T_C$. Amos and Schrock [6]

carried out experiments on constant area rectangular slits 20 mm in width and opening dimension range nominally from 0.127 to 0.318 mm. Stagnation pressure ranged between 4.1 and 16.2 MPa and water subcoolings from zero to 65°C. The length-to-hydraulic diameter ratios (L/D) of slits were between 83 and 400, with flow direction dimension fixed at 6.35 cm. Their study strongly suggested that the frictional effects are predominant in two-phase critical flows with large L/D.

More recently an experimental and analytical research program to study IGSCC associated with BWR piping was carried out at Battelle Columbus Laboratories (BCL) by Collier et al. [7] under the joint sponsorship of the Electric Power Research Institute (EPRI) and the BWR Owners Group. The experiments were carried out in two phases. In the first phase, simulated cracks were used. For simulated cracks, the ratio L/D and the surface roughness could be controlled. In the second phase, actual

IGSCC in stainless steel pipes were utilized. An analytical model to predict two-phase flow through cracks was developed [8] based on the non-equilibrium model suggested by Henry [9]. The model was developed by extending the Henry model to account for wall friction, flow area change and bends in the flow path. Further modifications were made to this model by Abdollahian and Chexal [10] to improve its agreement with the data. Both versions of this model, coded into programs LEAK and LEAK 01 respectively, assumed that flashing always begins at an L/D of 12 and that the quality varies linearly with distance along the flow path. Quality was evaluated assuming an isentropic process in LEAK and an isenthalpic process in LEAK 01. The calculations were done by separately calculating channel pressure drop due to momentum and friction based upon length averaged conditions rather than solving the equations in a marching method to obtain the distribution of pressure and quality along the crack length (flow direction).

The present study, aimed at developing a computational scheme for steady state two-phase critical flow through IGSCC, takes account of more realistic assumptions in modelling the flow through cracks. From the theoretical and experimental study by Amos and Schrock [6] it was observed that the predictions of critical two-phase flow through simulated cracks of large L/D using the homogeneous equilibrium model (HEM) agreed fairly well when compared with their experimental measurements (with deviations less than 20%). As the IGSCC have large L/D ratios, in the present study the HEM (with account of wall friction, area change and tortuosity) is used with a correction factor dependent upon upstream subcooling. The computational scheme is coded into a Fortran program called SOURCE. The inputs to the program require stagnation conditions and specifications of crack geometry. A properties subroutine, STEAM, is provided to the SOURCE code to calculate the fluid thermodynamic properties. The STEAM subroutine was developed using the equations for water given by Ishimoto et al. [15]. In the following sections the development of the computational scheme and the prediction of BCL tests are presented. A method of estimating the crack equivalent friction factor is also presented along with predicted results using these friction factors. The SOURCE flow chart and listing are presented in Appendices B and C, respectively.

2. CHARACTERISTICS OF FLOW THROUGH IGSCC

The BCL tests [7] on actual IGSCC were carried out on two stainless steel pipes which contained approximately 90 percent through-the-wall circumferential cracks. The crack depths were not actually uniform and an electric potential technique was used to locate the deepest section. At this location a portion of the pipe surface was machined away to expose the tip of the IGSCC. Progressively deeper cuts were made to obtain various L/D ratios. Figure 2.1 shows a sketch of the crack used by BCL. The exit to entrance area ratio for cracks was determined by inspection of the flow surface convergence following the experiments. The measured leak rates and the test parameters are given in Table 2.1.

From the stress analysis [11] the crack cross sectional area, A, can be related to the surface crack length, ℓ , as

$$A = \frac{\sigma}{2E} \ell^2 F\left(\frac{\sigma}{\sigma_y}\right) \quad (2.1)$$

Here the function F is of order unity. All the symbols are defined in the nomenclature.

For a typical stress of $\sigma = 8 \times 10^3$ psi and $E = 26 \times 10^6$ psi, the area A is given as $A \approx 5 \times 10^{-4} \ell^2$. Hence the equivalent slit thickness t Figure 2.1 is given as

$$t = 5 \times 10^{-4} \ell \quad (2.2)$$

Equation (2.2) gives an idea of crack width to length ratio. As indicated earlier it is important to define the margins for "leak-before-break" (LBB) i.e., the leak rate through a crack before failure of the pipe with events such as burst or break. With an assumed typical critical mass flux $G \sim 3900 \text{ lb}_m/\text{ft}^2\text{sec.}$, the estimations [12] of the crack length that will give the observable flow rate of 5 GPM and the critical crack length before break are given in Table 2.2 for different pipe sizes.

Usually the non circular flow cross sections are modelled by defining an equivalent diameter and using the correlations developed for round tubes. The crack width in BCL experiments was between 0.02 mm and 0.2 mm, the ratio of the crack opening length to width was between 30 and 150 and the L/D varied from 50 to 475. When the pipes were destructively tested after the BCL phase II tests, the flow path was found to have a number of turns.

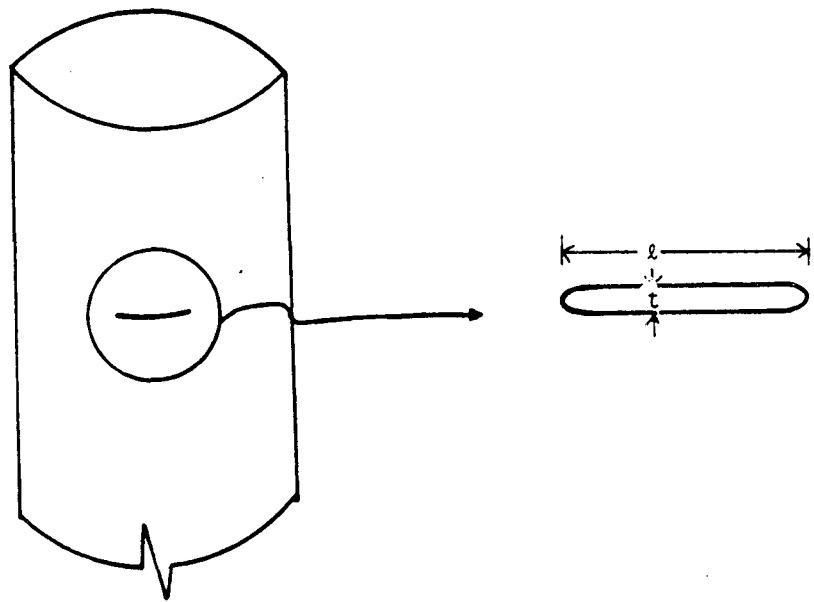
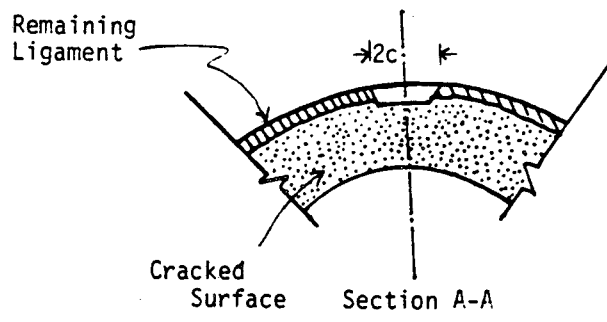
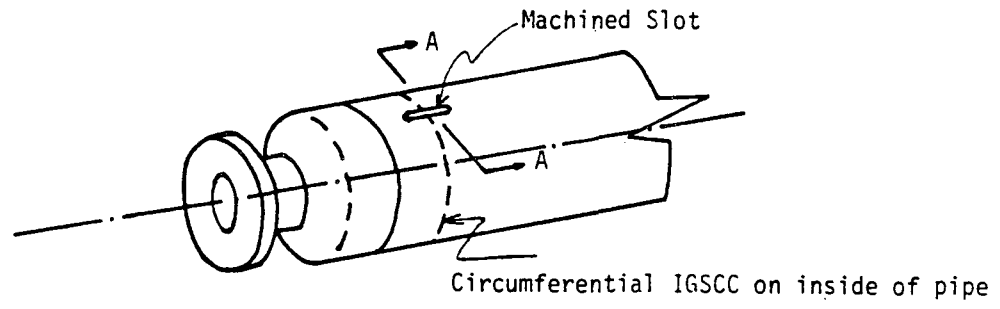


Figure 2.1 Illustration of BCL IGSCC Test Sections

Table 2.1 Test Parameters for BCL Phase II Experiments

Test No.	Stagnation Conditions			Crack Geometry					Measured
	P ₀ (MPa)	T ₀ (°C)	ΔT Subcooling (°C)	Depth L (mm)	Width δ (mm)	Length 2c (mm)	Area Ratio AR	Opening Area A (mm ²)	Leak Rate W (kg/s)
1	7.316	272.8	18.4	19.27	0.074	3.63	0.1	0.268	1.02 x 10 ⁻³
2	7.316	272.8	18.4	"	"	"	"	"	"
3	8.667	267.2	34.2	"	"	"	"	"	1.10 x 10 ⁻³
4	9.412	260.6	46.8	"	"	"	"	"	"
5	9.412	260.6	46.8	"	"	"	"	"	"
6	9.322	267.8	38.8	"	"	"	"	"	"
7	9.301	284.4	22.1	18.63	0.0220	0.74	0.04	0.0162	5.49 x 10 ⁻⁵
8	9.446	281.7	25.9	"	0.0218	"	"	0.0161	6.66 x 10 ⁻⁵
9	9.467	278.3	29.4	"	0.0216	"	"	0.0159	3.08 x 10 ⁻⁵
10	5.702	273.3	.1	"	0.0208	"	"	0.0153	3.08 x 10 ⁻⁵
11	5.971	262.8	13.6	"	0.0201	"	"	0.0148	1.34 x 10 ⁻⁵
12	5.868	260.0	15.2	"	0.0199	"	"	0.0146	1.09 x 10 ⁻⁴
13	5.854	268.3	6.7	"	0.0205	"	"	0.0151	1.51 x 10 ⁻⁴
14	5.868	272.8	2.4	"	0.0208	"	"	0.0153	1.01 x 10 ⁻⁴
15	3.523	251.7	-7.1	"	0.0190	"	"	0.0140	9.30 x 10 ⁻⁵
16	3.379	241.1	1.1	"	0.0183	"	"	0.0135	1.05 x 10 ⁻⁴
17	5.641	278.9	-5.7	"	0.0212	"	"	0.0156	1.01 x 10 ⁻⁴
18	7.081	260.0	27.6	"	--	"	"	--	6.21 x 10 ⁻⁵
19	7.309	273.9	15.8	19.27	0.108	9.53	0.13	1.026	3.04 x 10 ⁻²
20	7.309	282.2	7.4	"	"	"	"	"	2.90 x 10 ⁻²
21	9.005	280.0	24.2	"	"	"	"	"	4.85 x 10 ⁻²
22	9.005	273.3	30.8	"	"	"	"	"	4.62 x 10 ⁻²
23	8.964	256.7	39.3	"	"	"	"	"	4.52 x 10 ⁻²
24	8.888	256.7	36.3	"	"	"	"	"	4.44 x 10 ⁻²
25	7.164	256.1	32.2	"	"	"	"	"	3.67 x 10 ⁻²
26	7.164	256.7	31.7	"	"	"	"	"	3.59 x 10 ⁻²
27	5.571	260.6	11.4	"	"	"	"	"	2.51 x 10 ⁻²
28	5.626	267.8	4.8	"	"	"	"	"	2.35 x 10 ⁻²
29	5.592	241.7	30.5	"	"	"	"	"	3.01 x 10 ⁻²
30	5.592	243.9	28.3	"	"	"	"	"	2.98 x 10 ⁻²
31	7.219	242.8	46.1	"	"	"	"	"	3.96 x 10 ⁻²
32	7.219	238.3	46.6	"	"	"	"	"	3.94 x 10 ⁻²
33	8.812	241.7	60.9	"	"	"	"	"	4.51 x 10 ⁻²
34	8.812	235.0	67.6	"	"	"	"	"	4.73 x 10 ⁻²
35	7.129	226.1	61.9	"	"	"	"	"	3.84 x 10 ⁻²
36	7.129	222.8	65.2	"	"	"	"	"	3.92 x 10 ⁻²
37	5.592	220.0	52.2	"	"	"	"	"	3.84 x 10 ⁻²
38	5.44	223.2	47.1	"	"	"	"	"	2.97 x 10 ⁻²
39	3.937	230.6	20.3	"	"	"	"	"	2.62 x 10 ⁻²
40	3.937	232.8	18.1	"	"	"	"	"	1.57 x 10 ⁻²

Table 2.1 (continued)

Test No.	Stagnation Conditions			Crack Geometry					Measured
	P_0 (MPa)	T_0 (°C)	ΔT Subcooling (°C)	Depth L (mm)	Width δ (mm)	Length 2c (mm)	Area Ratio AR	Opening Area A_c (mm ²)	Leak Rate \dot{W} (kg/s)
41	7.309	278.9	10.8	18.63	0.0535	1.59	0.09	0.0849	1.81 x 10 ⁻³
42	8.060	277.2	19.2	"	0.0534	"	"	0.0848	1.98 x 10 ⁻³
43	8.722	276.7	25.2	"	0.0535	"	"	0.0850	2.19 x 10 ⁻³
44	9.377	274.4	32.6	"	0.0533	"	"	0.0847	2.37 x 10 ⁻³
45	7.35	278.3	11.9	"	0.0534	"	"	0.0848	3.62 x 10 ⁻⁴
46	8.067	275.0	21.4	"	0.0530	"	"	0.0842	5.26 x 10 ⁻⁴
47	8.757	275.0	27.2	"	0.0533	"	"	0.0845	3.91 x 10 ⁻³
48	9.432	272.2	35.3	"	0.0530	"	"	0.0841	4.37 x 10 ⁻³
49	7.35	276.1	13.9	"	0.0530	"	"	0.0841	3.22 x 10 ⁻³
50	8.067	274.4	20.3	"	0.0529	"	"	0.0840	3.59 x 10 ⁻³
51	5.213	258.3	9.4	"	0.0490	"	"	0.0778	2.51 x 10 ⁻³
52	6.923	253.9	32.1	"	0.0487	"	"	0.0773	3.54 x 10 ⁻³
53	8.646	252.8	48.5	"	0.0491	"	"	0.0779	4.11 x 10 ⁻³
54	5.213	256.7	11.1	"	0.0487	"	"	0.0773	2.20 x 10 ⁻³
55	6.902	254.4	31.4	"	0.0488	"	"	0.0775	2.61 x 10 ⁻³
56	8.646	253.3	47.9	"	0.0492	"	"	0.0781	3.17 x 10 ⁻³
57	3.413	233.9	8.9	"	0.0438	"	"	0.0696	7.57 x 10 ⁻⁴
58	5.171	232.8	34.5	"	0.0442	"	"	0.0711	9.03 x 10 ⁻⁴
59	6.923	231.1	54.9	"	0.0445	"	"	0.0706	1.50 x 10 ⁻³
60	8.646	228.9	72.4	"	0.0446	"	"	0.0708	1.85 x 10 ⁻³
61	8.646	231.7	69.6	"	0.0451	"	"	0.0717	1.75 x 10 ⁻³
62	7.033	228.3	58.0	"	0.0440	"	"	0.0698	1.76 x 10 ⁻³
63	5.206	225.0	42.7	"	0.0428	"	"	0.0679	1.53 x 10 ⁻³
64	3.448	222.8	20.6	"	0.0415	"	"	0.0663	1.33 x 10 ⁻³
65	3.461	235.0	8.6	"	0.0441	"	"	0.0700	1.84 x 10 ⁻³
66	5.185	232.2	35.3	"	0.0441	"	"	0.0700	2.32 x 10 ⁻³
67	6.895	230.0	55.8	"	0.0443	"	"	0.0703	2.69 x 10 ⁻³
68	8.646	228.3	71.8	"	0.0445	"	"	0.0707	2.69 x 10 ⁻³
69	8.626	241.7	59.4	19.27	0.235	27.89	0.21	6.547	1.51 x 10 ⁻¹
70	6.936	236.8	49.5	"	0.227	"	"	6.338	1.39 x 10 ⁻¹
71	6.861	236.7	48.8	"	0.227	"	"	6.334	1.57 x 10 ⁻¹
72	8.536	241.1	59.3	"	0.234	"	"	6.528	1.75 x 10 ⁻¹
73	5.109	252.8	13.7	"	0.178	"	"	4.952	1.41 x 10 ⁻¹
74	6.861	253.9	31.6	"	0.243	"	"	6.787	1.69 x 10 ⁻¹
75	8.605	254.4	46.5	"	0.247	"	"	6.882	1.91 x 10 ⁻¹
76	8.598	250.6	50.3	"	0.243	"	"	6.779	2.00 x 10 ⁻¹
77	6.861	248.3	37.1	"	0.238	"	"	6.641	1.78 x 10 ⁻¹
78	5.192	246.1	21.4	"	0.173	"	"	4.828	1.51 x 10 ⁻¹
79	3.937	240.6	10.3	"	0.144	"	"	4.024	1.28 x 10 ⁻¹
80	5.089	236.1	30.2	"	0.166	"	"	4.633	1.55 x 10 ⁻¹
81	6.771	237.2	47.3	"	0.228	"	"	6.345	1.76 x 10 ⁻¹
82	8.577	234.4	66.3	"	0.228	"	"	6.355	1.99 x 10 ⁻¹

In Figure 2.2 photomicrographs of typical IGSCC in weld sensitized type 304 stainless steel are shown. These pictures reveal the tortuous and irregular nature of the channel with many expansions and contractions along the passage.

Table 2.2 Leak Before Failure Criterion

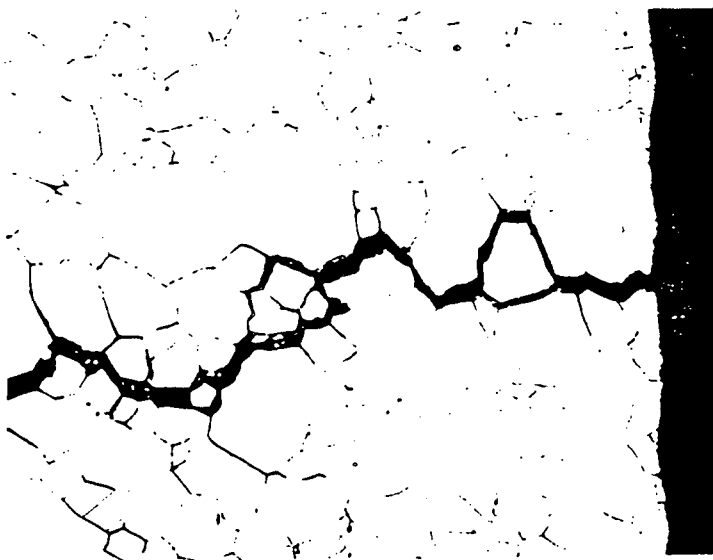
$$G_c = 3900 \text{ lb}_m/\text{ft}^2 \text{ sec. Assumed.}$$

Pipe ID in.	Pipe Wall Thickness, in.	$l_{5\text{GPM}}$ in.	l_{crit} , in.	$\frac{l_{5\text{GPM}}}{l_{\text{crit}}}$
4	0.3	4.5	6.5	0.69
10	0.7	6.9	16	0.43
24	1.0	7.0	36	0.20

It should be noted that the BCL method of measuring the average crack thickness δ was not documented. Inconsistencies may be noted between some mass flowrate data and crack opening areas given in Table 2.1. It seems likely that this may be due in part to the difficulty in measuring δ . Another source of discrepancies may be the non-uniformity of δ and possible plugging of some cracks by small solid particles. These difficulties are real and impact the predictive capability of any model.



IGSCC IN WELD HEAT AFFECTED ZONE - 8X



HIGHER MAGNIFICATION - 125X

Figure 2.2 Photomicrograph of Typical IGSCC in Weld Sensitized Type 304 Stainless Steel

3. MODELLING

In the models developed by Collier [7] and by Abdollahian and Chexal [10] it was assumed, similar to Henry's model [9], that the flow was homogeneous and the non-equilibrium effects were taken into account through the empirical parameter, N , which is a function of equilibrium quality and the flow path length to diameter ratio L/D . In addition to these assumptions the critical flow model developed and coded in the program LEAK included friction and acceleration pressure drops within the flow path and the pressure drop at entrance (sharp-edged entrance was assumed). In the code LEAK 00 the quality in the crack duct was assumed to vary linearly along the flow direction and was evaluated with assumption of isentropic flow. The friction pressure drop was calculated with the mass flux evaluated as the average of inlet and outlet mass flux. The surface roughness was taken as $\epsilon = 0.00178$ mm and six 45 degree turns in the flow path were assumed to determine the friction pressure drop. In a subsequent code, LEAK 01, the quality was evaluated from an isenthalpic assumption, which is more nearly correct thermodynamically. The average mass flux was evaluated with the assumption of linearly varying cross-sectional area of the crack duct. For the friction pressure drop evaluation, a surface roughness of $\epsilon = 0.0051$ mm and twenty 45-degree turns in the flow path were assumed.

In actual pipe IGSCCs the entrance to the crack duct does not represent sharp-edged geometry. The crack L/D ratios are large (>50) hence the homogeneous equilibrium assumption with a rounded entrance seems more appropriate. The values for L/D beyond which homogeneous equilibrium assumptions can be made have been reported in the literature. These values vary from 1.5 for reactor scale pipe breaks [13] to 25 for 4 mm diameter tubes with sharp-edged entrance geometry. In the model of the LEAK code the flashing of the liquid in the crack channel is assumed to occur at $Z = 12D$ irrespective of crack dimensions and thermodynamic properties of the upstream fluid. In the present work, the more realistic assumptions are made in modelling the critical flow through IGSCC as described in the next section.

3.1 Assumptions of the Present Model

The actual flow channel in an IGSCC (as shown in Figure 2.1) is approximated with a one dimensional convergent slit as shown in Figure 3.1. The convergence of the channel is determined by the specifications of inlet to exit area ratio and crack depth. Thus the flow area is assumed to vary linearly along the z -direction with the flow channel, while the slit gap is taken as uniform in transverse and in flow directions. As mentioned earlier, the fluid is considered homogeneous and in thermodynamic equilibrium. The entrance region is assumed to be smooth, hence frictionless. The gravitational effects

3-2

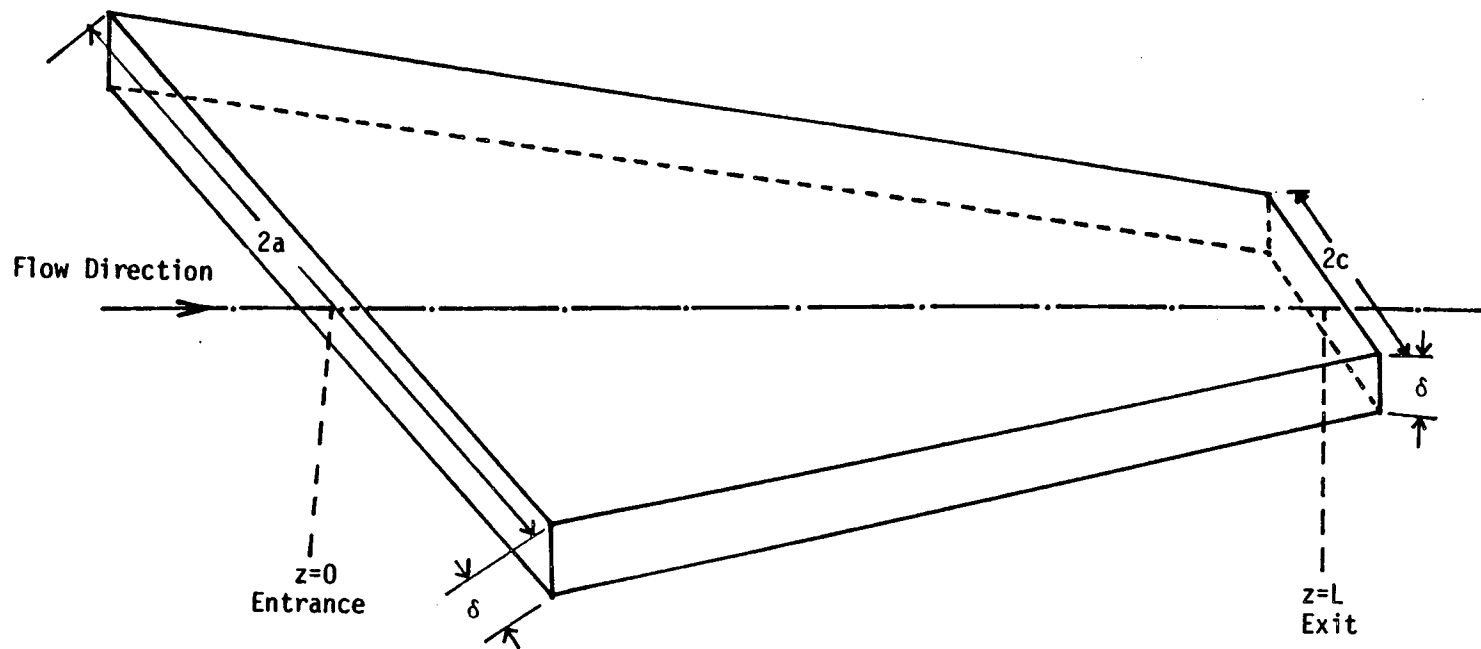


Figure 3.1 Idealized Geometry of the Crack

are neglected and the control surface of the geometry considered is assumed to be adiabatic. The flow is considered in steady state.

For subcooled stagnation conditions, the flow in the channel is divided into two regions, single phase and two-phase. The two phase region is assumed to begin at the plane where the liquid reaches the saturation pressure at T_0 . Both single phase and two phase frictional pressure losses are calculated by assuming a single equivalent friction factor. Evaluation of the equivalent friction factor, which accounts for the tortuosity (bends, contractions, expansions) of the channel, is a key element in the modelling. In section 4.1 we present a parametric study of the BCL data to obtain optimized friction factors for each test section. In section 4.2 we present a generalized approach for estimating friction factors for IGSCC. This method is applied to the BCL test sections in Appendix A and the resulting friction factors used in an independent prediction of the BCL critical flow data.

3.2 Governing Equations

The continuity equation is

$$\frac{d}{dz} (\rho VA) = 0 \quad (3.1)$$

where $\dot{m} = \rho VA$.

The momentum equation is

$$-\frac{dP}{dz} = \rho V \frac{dV}{dz} + \frac{\tau_w P}{A} \quad (3.2)$$

The first term on right hand side of momentum equation is the acceleration pressure drop and the second term is the friction pressure drop.

The energy equation is

$$\frac{d}{dz} \left(h + \frac{V^2}{2} \right) = 0 \quad (3.3)$$

The state equations are

i) for two-phase mixture

$$v = (1 - x)v_f(P) + x v_g(P) = v_f(P) + x[v_g(P) - v_f(P)] \quad (3.4)$$

$$h = (1 - x)h_f(P) + x h_g(P) = h_f(P) + x [h_g(P) - h_f(P)] \quad (3.5)$$

ii) for single-phase liquid

$$v_l \approx \text{constant} = v_f(T_0) . \quad (3.6)$$

$$h_l \approx h_f(T_0) . \quad (3.7)$$

Critical Flow Criteria

Homogeneous equilibrium sound speed criterion was used to determine when the flow was critical. This is given as

$$v = v_a = \left\{ \left(\frac{\partial P}{\partial \rho} \right)_s \right\}^{1/2} \quad (3.8)$$

at the channel exit.

Using equations (3.2) to (3.5), equation (3.8) becomes

$$v_a = v \left[- \left\{ (1-x) \frac{dv_f}{dP} + x \frac{dv_g}{dP} + (v_g - v_f) \left(\frac{\partial x}{\partial P} \right)_s \right\} \right]^{1/2} \quad (3.9)$$

where

$$\left(\frac{\partial x}{\partial P} \right)_s = \frac{\left(\frac{ds_f}{dP} \right) (1-x) + x \left(\frac{ds_g}{dP} \right)}{(s_g - s_f)} \quad (3.10)$$

3.3 Development of Finite Difference Forms of Governing Equations

To achieve a numerical solution, the above equations are reduced to finite difference forms using a forward differencing scheme, which is numerically stable. The continuity and the energy equations, (3.1) and (3.3), are given in finite difference form as

$$\frac{\Delta}{\Delta z_i} (\rho AV)_i = \frac{\Delta \dot{m}_i}{\Delta z_i} = 0 \quad (3.11)$$

$$\frac{\Delta h_i}{\Delta z_i} + \frac{\Delta}{\Delta z_i} \left(\frac{1}{2} v^2 \right)_i = 0 . \quad (3.12)$$

Here $\Delta z_i = z_{i+1} - z_i$, $\Delta h_i = h_{i+1} - h_i$,

$$\text{and } \Delta \left(\frac{1}{2} v^2 \right)_i = \frac{1}{2} v_{i+1}^2 - \frac{1}{2} v_i^2 .$$

The momentum equation (3.2) can be written as

$$- \left(\frac{dP}{dz} \right) = \frac{\dot{m}^2}{A^2} \left(\frac{dv}{dz} \right) - \frac{\dot{m}^2 v}{A^3} \left(\frac{dA}{dz} \right) + f \left(\frac{P}{A} \right) \frac{\dot{m}^2 v}{2A^2} \quad (3.13)$$

where f is an equivalent friction factor including effects of bends, contractions and expansions.

In equation (3.13) the first term on the right hand side is the pressure gradient due to phase change, the second term is the pressure gradient due to area change and the third term is pressure gradient due to friction in the crack channel. Equation (3.13) is integrated along the flow path to evaluate the overall pressure drop across the crack. The channel region is divided into two regions: a) entrance region; and b) main channel region.

a) Entrance region

i) Single phase entrance. (High stagnation subcooling)

With ideal flow assumed, the single phase pressure drop for flow between stagnation region and entrance region is obtained by using the Bernoulli equation (special case of equation (3.13))

$$\frac{P_0}{\rho_0} = \frac{P_1}{\rho_1} + \frac{v_1^2}{2}$$

or

$$P_{0-1} = \frac{1}{2} \rho_0 v_1^2 = \frac{\dot{m}^2 v_0^2}{2A_1^2} . \quad (3.14)$$

This case corresponds to $P_1 > P_{\text{sat}}(T_0)$. For the case with $P_1 = P_{\text{sat}}(T_0)$, there will be flashing at the entry. The entrance pressure drop is then

$$\Delta P_{0-1} = P_0 - P_{\text{sat}}(T_0) . \quad (3.15)$$

ii) Two phase entrance

In the case of very low stagnation subcooling the liquid may start to flash before entering the channel. In such cases the entrance pressure drop is calculated by the use of thermodynamic relationships for isentropic flow, i.e., both entropy and stagnation enthalpy remain constant. Thus we have

$$s_o = s_1 = s_{f1} + x_1 s_{fg1} \quad (3.16)$$

$$h_o = h_{f1} + x_1 h_{fg1} + \frac{v_1^2}{2} \quad (3.17)$$

and continuity

$$\dot{m} = \frac{A_1 v_1}{v_1} \quad (3.18)$$

where

$$v_1 = v_{f1} + x_1 v_{fg1}$$

With known values of A_1 , \dot{m} , s_o , and h_o equations (3.16), (3.17) and (3.18) determine P_1 , x_1 and v_1 .

b) Main channel region.

i) Single phase flow

The single phase region ends at the point where the flashing occurs. The pressure drop between the entrance plane and the plane of flashing (where $P = P_{sat}$) is obtained from equation (3.13) neglecting the compressibility of the liquid, as

$$P_1 - P_{sat} = \frac{\dot{m}^2}{2} v_o \left(1 + \frac{\delta f}{\eta}\right) \left(\frac{1}{A_{f\ell}^2} - \frac{1}{A_1^2}\right) + \frac{\dot{m}^2 f v_o}{\delta \eta} \left(\frac{1}{A_{f\ell}} - \frac{1}{A_1}\right) \quad (3.19)$$

where $\eta \doteq \frac{2\delta}{L} (a-c)$.

Hence the flashing location is obtained in terms of the flashing area A_f

$$A_{f\ell} = \frac{1 + \left[1 + (\delta^2 \eta^2 + \delta^3 \eta f) / (\dot{m}^2 f^2 v_o)\right] C_1}{(\delta \eta / (\dot{m}^2 f v_o)) C_1} \quad (3.20)$$

where

$$C_1 \doteq 2(P_1 - P_{\text{sat}}) + \frac{\dot{m}^2 v_0}{A_1^2} \left(1 + \frac{\delta f}{\eta}\right) + \frac{2\dot{m}^2 f v_0}{\delta \eta A_1} .$$

ii) Two-phase region

Now the pressure drop between two local points z_{i+1} and z_i in the channel can be written as

$$\Delta P_{(i+1-i)} = P_i - P_{i+1} = - \int_{z_i}^{z_{i+1}} \left(\frac{dP}{dz} \right) dz \quad (3.21)$$

From equations (3.13) and (3.21) we have the pressure drop between two local points z_{i+1} and z_i as

$$\Delta P_{i+1-i} = \dot{m}^2 \int_{v_i}^{v_{i+1}} \frac{dv}{d^2} - \dot{m}^2 \int_{A_i}^{A_{i+1}} \frac{v}{A^3} dA + \frac{\dot{m}^2}{2} \int_{z_i}^{z_{i+1}} f \left(\frac{P_m}{A} \right) \left(\frac{v}{A^2} \right) dz \quad (3.22)$$

Here

$$v_i = v(z = z_i)$$

$$v_{i+1} = v(z = z_{i+1})$$

$$A_i = A(z = z_i)$$

$$A_{i+1} = A(z = z_{i+1})$$

Now, the first integral in equation (3.22) is evaluated as follows

$$\dot{m}^2 \int_{v_i}^{v_{i+1}} \frac{dv}{A^2} = \frac{\dot{m}^2}{\bar{A}_i^2} (v_{i+1} - v_i) \quad (3.23)$$

where

$$\bar{A}_i^2 = \frac{1}{2} (A_i^2 + A_{i+1}^2)$$

and the second integral is approximated by

$$\dot{m}^2 \int_{A_i}^{A_{i+1}} \frac{v}{A^3} (-dA) = \frac{\dot{m}^2 \bar{v}_i}{2} \left(\frac{L}{A_{i+1}^2} - \frac{1}{A_i^2} \right) \quad (3.24)$$

where

$$\bar{v}_i = \left(\frac{v_i + v_{i+1}}{2} \right).$$

In order to evaluate third integral we have to evaluate (P/A) for the convergent crack geometry. A linear change in area of the crack channel is assumed in the present model. The local area $A(z)$ is given as

$$A(z) = A_1 - \eta z \quad (3.25)$$

where

$$A_1 = 2a\delta \quad \text{and} \quad \eta = \frac{2\delta}{L} (a-c).$$

Hence

$$(P_m/A) = \frac{2}{A} \left(\frac{A}{\delta} + \delta \right) = 2 \left(\frac{1}{\delta} + \frac{\delta}{A_1 - \eta z} \right) \quad (3.26)$$

Using equations (3.25) and (3.26) we have after simplifications

$$\begin{aligned} \frac{\dot{m}^2}{2} \int_{z_i}^{z_{i+1}} f \left(\frac{P_m}{A} \right) \left(\frac{v}{A^2} \right) dz &= \frac{\dot{m}^2 f \bar{v}_i}{\delta \eta} \left[\left\{ \frac{1}{(A_1 - \eta z_{i+1})} - \frac{1}{(A_1 - \eta z_i)} \right\} \right. \\ &\quad \left. + \frac{\delta^2}{2} \left\{ \frac{1}{(A_1 - \eta z_{i+1})^2} - \frac{1}{(A_1 - \eta z_i)^2} \right\} \right] \quad (3.27) \end{aligned}$$

Using equations (3.22), (3.23), (3.24) and (3.27) we have after some rearrangement

$$\begin{aligned} (\Delta P)_{i+1-i} &= \frac{\dot{m}^2}{A_i} (v_{i+1} - v_i) + \frac{\dot{m}^2 v_i}{2} \left(1 + \frac{\delta f}{\eta} \right) \\ &\times \left\{ \frac{1}{(A_1 - \eta z_{i+1})^2} - \frac{1}{(A_1 - \eta z_i)^2} \right\} + \frac{\dot{m}^2 f \bar{v}_i}{\delta \eta} \left\{ \frac{1}{(A_1 - \eta z_{i+1})} - \frac{1}{(A_1 - \eta z_i)} \right\} \quad (3.28) \end{aligned}$$

The total pressure drop through the crack is then given as

$$(\Delta P)_T = (\Delta P)_{0-1} + (\Delta P)_{1-e} \quad (3.29)$$

The solution must simultaneously satisfy the continuity, energy and state equations.

The state equations (3.4) to (3.7) are given in difference form as

i) for two-phase homogeneous mixture

$$v_i = v_f(P_i) + x_i [v_g(P_i) - v_f(P_i)] \quad (3.30)$$

$$h_i = h_f(P_i) + x_i [h_g(P_i) - h_f(P_i)] \quad (3.31)$$

ii) for single-phase liquid

$$v_l = v_l(T_0) \approx \text{constant} \quad (3.32)$$

$$h_l = h_l(T_0) \quad (3.33)$$

Using equations (3.30) to (3.31) the energy equation (3.12) becomes

$$h_0 = h_f(P_i) + x_i [h_g(P_i) - h_f(P_i)] + \frac{\dot{m}^2}{2A_i^2} [v_f(P_i) + x_i \{v_g(P_i) - v_f(P_i)\}]^2 \quad (3.34)$$

Equation (3.34) can be rearranged to give the quality at mode i as

$$x_i = -C_2 + \left[C_2^2 - \frac{2A_i^2(h_g - h_f)}{\dot{m}^2(v_g - v_f)} + \left(\frac{v_f}{v_g - v_f} \right)^2 \right]^{1/2} \quad (3.35)$$

where

$$C_2 = \left[\frac{v_f}{v_g - v_f} + \frac{2A_i^2(h_g - h_f)}{\dot{m}^2(v_g - v_f)} \right] \quad \text{and } h \text{ and } v \text{ are evaluated}$$

at P_i .

3.4 Computer Code (SOURCE) Logic

The difference equations presented in section 3.3 were programmed into a Fortran code called SOURCE. The logic of the code method is described in a SOURCE flow chart given in Appendix B. The SOURCE program listing is given in Appendix C. A brief description of the key features follows.

The governing steady equations are solved iteratively marching downstream from the upstream stagnation region using an assumed mass flowrate. This develops the location of incipient flashing and the distribution of pressure and quality in the flow direction. At each node point the converged value of fluid velocity is compared with the HEM sound speed calculated for the local state. The marching procedure continues until either, a. the exit plane is reached at a subcritical flow state or b. a critical flow state is reached at a position upstream of the exit. In the first instance the assumed flowrate was too low, in second instance it was too high. The mass flowrate is adjusted systematically until a critical flow state is found at the exit plane within a small tolerance.

In principle any first guess flowrate will do, however, to reduce the number of iterations a simple approximate method of estimating the flowrate was built into SOURCE. The method uses the IHEM nozzle results tabulated by Hall [17] with a multiplier deduced from the BCL data. The multiplier is a linear function of L/D with subcooling as a parameter. As programmed this estimate may be made for stagnation temperature $480K \leq T_0 \leq 550K$ and stagnation pressure $3 \leq P_0 \leq 10 \text{ MPa}$.

The location of incipient flashing is found precisely using equation (3.20) for the case of incipient flashing within the channel. When flashing is initiated upstream of the channel equations (3.16) - (3.18) give the pressure and quality at the entrance. These results identify the region of two phase channel flow. Because the pressure gradient, in friction dominated critical flow, increases rapidly as the critical location is approached, a non-uniform axial discretization was chosen. Each successive distance $Dz(i) = (z_{i+1} - z_i)$ is just half the preceding one such that $Dz(i) = (DzT)2^{N-i}$. With the total two phase length given by $z_{TP} = \sum_{i=1}^N Dz(i)$ the number of nodes N is fixed by the requirement that $DzT \leq 10^{-7} \text{ m}$.

The first node in the two phase region used a subroutine GUESS for an estimate of the pressure gradient, which changes stepwise at incipient flashing in the HEM model. GUESS uses the single phase pressure gradient to estimate the pressure P_{i+1} by forward extrapolation. This pressure is in turn used to evaluate an approximate quality x_{i+1} from equation (3.34). This quality is then used to evaluate the two phase

friction multiplier with the aid of a linear fit to the result of Martinelli and Nelson [18].

$$\phi_{10}^2 = 1.20 (v_{fg}/v_f)(1.42 - 1.67 \times 10^{-4}P) x^{0.9} + 1. \quad (3.36)$$

A second estimate of P_{i+1} is then obtained by extrapolating the two phase pressure gradient. The second estimate of x_{i+1} is then the first guess for iterative solution of equations (3.28) through (3.35).

For subsequent axial nodes the last evaluated pressure gradient is extrapolated forward to obtain a first estimate of P_{i+1} . Use of the energy equation (3.34) then gives the first guess x_{i+1} for iterative solution of equations (3.28) through (3.35).

After downstream marching is terminated by reaching a critical flow state or the channel exit, the flowrate is then corrected accordingly. When the exit is reached at a Mach number $M < 1.0$, the flowrate is corrected by

$\dot{m}_{new} = \dot{m}_{old} M^{-0.05}$ and a new solution is generated. If again the flow is subcritical at the exit the same correction is applied. This procedure is repeated until the new flowrate results in a critical flow state upstream of the exit. The desired solution is then bracketed by the last subcritical flow \dot{m}_1 and the last supercritical flow \dot{m}_2 . The flowrate correction

$\dot{m} = \frac{\dot{m}_1 + \dot{m}_2}{2}$ is then used. If this flowrate gives a critical flow state upstream of the exit, it becomes the new \dot{m}_2 . If instead it gives a subcritical flow at the exit it becomes the new \dot{m}_1 . Successive corrections then converge rapidly and the solution is accepted when, for $M_e < 1$, $(1 - M_e) \leq 10^{-3}$ or when a critical flow state is reached at a position $\leq 10^{-7}m$ upstream of the exit.

Thermodynamic properties of saturated liquid and vapor are evaluated using a subroutine STEAM which is described in Appendix D. This subroutine is very fast running and contributes to the efficiency of the numerical calculations.

The code was found to be very fast running. A typical computation required about 22 flowrate corrections and was completed in about 3 seconds of CPU time on an IBM 3081.

4. RESULTS

4.1 Parametric Study of BCL Data

Table 2.1 gives data of BCL phase II experiments carried out on IGSCC. In order to establish appropriate friction factors for the actual cracks all the BCL data except 21 were studied with SOURCE. These 21 tests (numbers 7, 8, 9, 10, 11, 13, 18, 41, 42, 43, 57, 58, 59, 60, 61, 62, 63, 73, 78, 79 and 80) were disqualified for the parametric study, because they showed inconsistent measured mass flow rates in relation to the stagnation conditions, when compared with the main body of experimental data for their respective test sections. For initial SOURCE runs, certain values of friction factor were tried for each crack size so that the predicted mass flowrates compared well with the measured data. This process required running SOURCE least three friction factor values for each experimental test condition. Then an optimum value of friction factor was obtained for each crack size. This involved a graphical interpolation procedure. The optimum values of friction factor for each test section are given in Table 4.1. From the table we find that for the test section used for runs 1-6, the friction factor so obtained is quite high. For these test numbers the mass flow rate is in the range of 10^{-3} kg/s. The opening area of the crack was 0.268 mm which is quite high compared with the test section used for test numbers 41 to 68, for which the measured mass flow rates were in the range of 10^{-3} kg/s as well. Comparing the mass flow rates and crack geometry for tests 1-6 with other test sections, it is concluded that either the crack channel was blocked by particles during experimental measurements, as observed in destructive tests of pipes following BCL experiments, or the geometry of the crack is not properly given.

Once the optimum equivalent friction factor was obtained it was used again to run SOURCE for each test condition. When these results were compared it was apparent that a systematic deviation existed between the prediction and the measured data that was dependent upon the stagnation subcooling. It was found that the higher value of friction factor was necessary for smaller subcooling than at higher subcooling. This is physically plausible since at lower subcooling the two-phase region is longer than that with higher subcooling, and hence the effective friction factor for the case with longer two-phase region should be greater than for the case with smaller two-phase region. The subcooling for the BCL tests varied from 0.1C to 78C. A correction factor was developed to account for the subcooling effects. The subcooling correction factor is given by

$$\left. \begin{aligned} C &= 1.3015 - 5.3075 \times 10^{-3} \Delta T_{\text{sub}} \text{ for } \Delta T_{\text{sub}} < 60^{\circ}\text{C} \\ &= 1.0 \text{ for } \Delta T_{\text{sub}} > 60^{\circ}\text{C} \end{aligned} \right\} (4.1)$$

Table 4.1 BCL Friction Factors from Parametric Study

Test Numbers	δ mm Average	L mm	A_e mm ² Average	\dot{m} Kg/s Range	f
1 - 6	0.074	19.27	0.268	10^{-3}	9.0
7 - 18	0.0205	18.63	0.0151	10^{-4}	0.8
19 - 40	0.108	19.27	1.026	10^{-2}	0.07
41 - 68	0.0488	18.63	0.0777	10^{-3}	0.02
69 - 82	0.220	19.27	6.133	10^{-1}	0.30

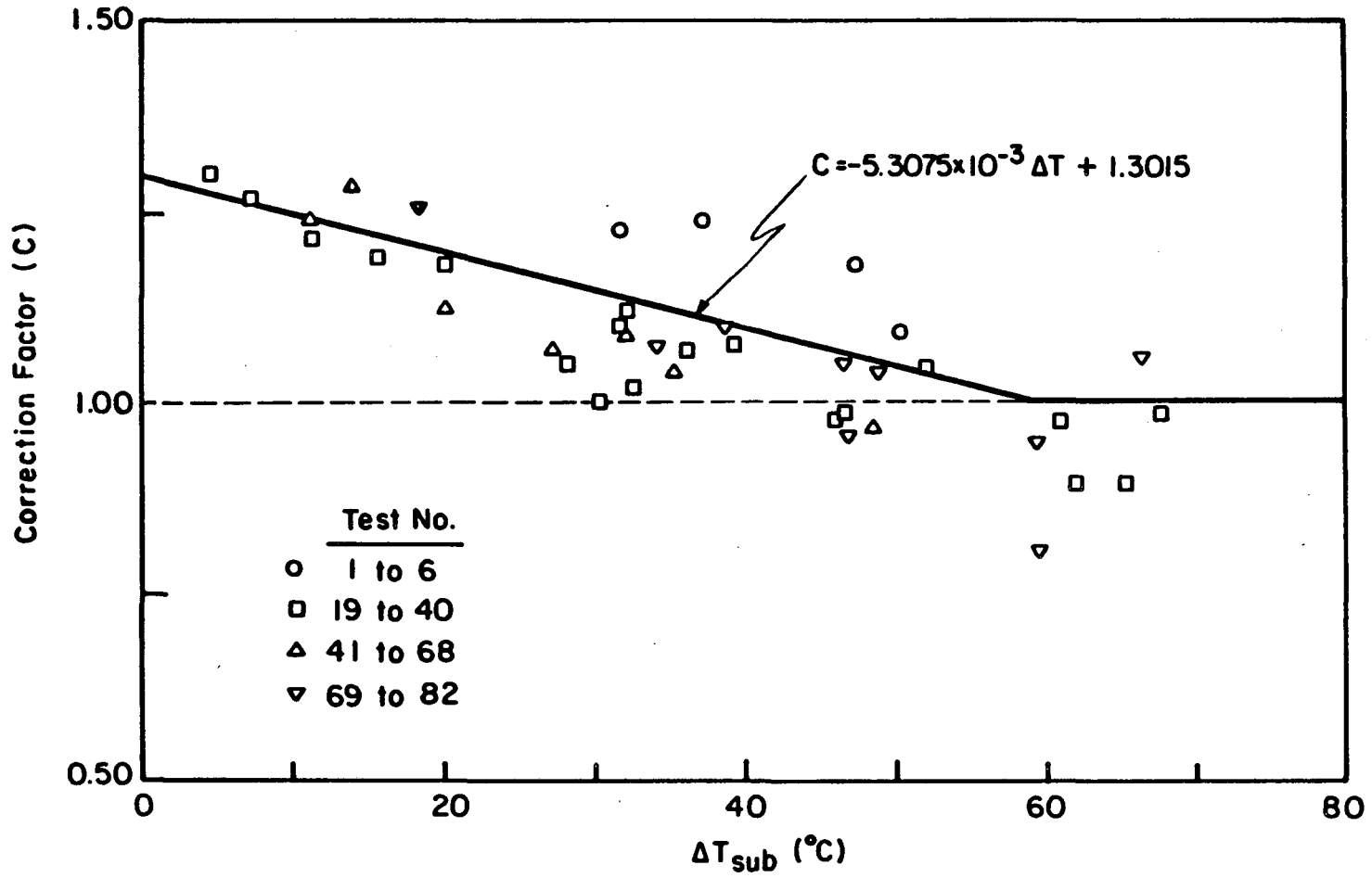


Figure 4.1 Subcooling Correction for SOURCE

In Table 4.2, the calculated mass flow rates, pressure and quality at the exit plane using SOURCE are presented for all the BCL phase II tests. These calculated values of mass flow rates when multiplied by the correction factor C, to account for the subcooling effect, then become the predicted mass flow rates. The predicted mass flow rates for 61 test numbers have been compared with BCL measured data in Figure 4.2. The comparison is remarkably good with standard deviation of 1.35%. When all the total 82 test number predictions were compared with BCL data, the standard deviation was 6.45%. Considering the evident lack of coherence in the basic BCL data, these comparisons show that the predictive method developed here offers an excellent capability to predict crack flows if the geometry is known.

4.2 Generalized Prediction of Equivalent Friction Factors

Choice of the optimum friction factor used to predict the results given in Table 4.2 was based on the parametric study of BCL data. For general use of SOURCE a methodology to estimate the equivalent friction factor from global features of IGSCCs is required. In this section we present a methodology which has been tested against the results of the parametric study of the BCL data in Appendix A.

The friction pressure drop in equation (3.13) was presented by an equivalent friction factor in the form

$$\left(\frac{dp}{dz}\right)_F = f \left(\frac{P_m}{A}\right) \cdot \frac{\dot{m}^2 v}{2A^2} \quad (4.2)$$

and

$$f = f_p + D \sum_i n_i K_i \quad (4.3)$$

where f_p represents "pipe" friction and the K_i represent the loss factors for specific effects such as bends, contractions, expansions and tortuosity and the n_i represents the number per unit length in the flow direction of the type i .

Assigning $i = 1$, for bends, $i = 2$ for contraction and $i = 3$ for expansions. The respective loss coefficients per unit length are given by [14].

$$\text{For bends } n_i K_i = (n_i - \frac{1}{L})(0.25 f_p \cdot \pi \cdot (\frac{r}{D}) + 0.5 K') + K'/L \quad (4.4)$$

when n_i is number of bends per unit length, r is the radius of the bend D is the hydraulic diameter, f_p is the tube friction factor, and

Table 4.2 SOURCE Predictions Using Friction Factors from BCL Parametric Study

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Predictions		
	P_0 , MPa	T_0 , C		\dot{m} Kg/sec	x_e	P_e MPa
1	7.316	272.8	1.02×10^{-3}	8.129×10^{-4}	2.297×10^{-1}	0.762
2	7.316	272.8	1.02×10^{-3}	8.129×10^{-4}	2.297×10^{-1}	0.762
3	8.667	267.2	1.10×10^{-3}	1.025×10^{-3}	1.984×10^{-1}	1.005
4	9.412	260.6	1.10×10^{-3}	1.153×10^{-3}	1.840×10^{-1}	0.952
5	9.412	260.6	1.10×10^{-3}	1.153×10^{-3}	1.840×10^{-1}	0.952
6	9.322	267.8	1.10×10^{-3}	1.000×10^{-3}	1.806×10^{-1}	1.278
7	9.301	284.4	5.49×10^{-5}	1.636×10^{-4}	2.495×10^{-1}	0.837
8	9.446	281.7	6.66×10^{-5}	1.676×10^{-4}	1.542×10^{-1}	2.371
9	9.467	278.3	3.08×10^{-5}	1.706×10^{-4}	1.420×10^{-1}	2.433
10	5.702	273.3	3.08×10^{-5}	8.809×10^{-4}	1.735×10^{-1}	1.604
11	5.971	262.8	1.34×10^{-4}	1.052×10^{-4}	1.415×10^{-1}	1.674
12	5.868	260.0	1.09×10^{-4}	1.043×10^{-4}	1.426×10^{-1}	1.525
13	5.854	268.3	1.51×10^{-4}	6.286×10^{-5}	4.817×10^{-1}	1.327
14	5.868	272.8	1.01×10^{-4}	9.394×10^{-5}	1.799×10^{-1}	1.445
15	3.523	251.7	9.30×10^{-5}	6.226×10^{-5}	1.560×10^{-1}	1.051
16	3.379	241.1	1.05×10^{-4}	5.483×10^{-5}	1.498×10^{-1}	0.858
17	5.861	278.9	1.01×10^{-4}	6.736×10^{-5}	3.503×10^{-1}	5.321
18	7.081	260.0	6.21×10^{-5}	-	-	-

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Table 4.2 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Predictions		
	P_0 , MPa	T_0 , C		\dot{m} Kg/sec	x_e	P_e MPa
19	7.309	273.9	3.04×10^{-2}	2.506×10^{-2}	2.985×10^{-2}	4.998
20	7.309	282.2	2.90×10^{-2}	2.252×10^{-2}	6.109×10^{-2}	4.830
21	9.005	280.0	4.85×10^{-2}	3.120×10^{-2}	1.202×10^{-2}	6.070
22	9.005	273.3	4.62×10^{-2}	3.479×10^{-2}	0.0	5.796
23	8.964	256.7	4.52×10^{-2}	4.200×10^{-2}	0.0	4.448
24	8.888	256.7	4.44×10^{-2}	4.165×10^{-2}	0.0	4.448
25	7.164	256.1	3.67×10^{-2}	3.285×10^{-2}	0.0	4.404
26	7.164	256.7	3.59×10^{-2}	3.255×10^{-2}	0.0	4.448
27	5.571	260.6	2.51×10^{-2}	2.155×10^{-2}	2.011×10^{-2}	3.967
28	5.626	267.8	2.35×10^{-2}	1.805×10^{-2}	5.524×10^{-2}	3.852
29	5.592	241.7	3.01×10^{-2}	2.931×10^{-2}	0.0	3.449
30	5.592	243.9	2.98×10^{-2}	2.833×10^{-2}	0.0	3.584
31	7.219	242.8	3.96×10^{-2}	3.850×10^{-2}	0.0	3.516
32	7.219	238.3	3.94×10^{-2}	4.003×10^{-2}	0.0	3.249
33	8.812	241.7	4.51×10^{-2}	4.640×10^{-2}	0.0	3.449
34	8.812	235.0	4.73×10^{-2}	4.819×10^{-2}	0.0	3.074
35	7.129	226.1	3.84×10^{-2}	4.317×10^{-2}	0.0	2.603
36	7.129	222.8	3.92×10^{-2}	4.402×10^{-2}	0.0	2.447
37	5.592	220.0	3.84×10^{-2}	3.686×10^{-2}	0.0	2.320
38	5.440	223.2	2.97×10^{-2}	3.506×10^{-2}	0.0	2.465

Table 4.2 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Predictions		
	P_0 , MP _a	T_0 , °C		\dot{m} Kg/sec	x_e	P_e MPa
39	3.937	230.6	2.62×10^{-2}	2.224×10^{-2}	0.0	2.827
40	3.937	232.8	1.57×10^{-2}	2.010×10^{-2}	0.0	2.944
41	7.309	278.9	1.81×10^{-3}	2.344×10^{-3}	2.439×10^{-2}	5.586
42	8.060	277.2	1.98×10^{-3}	3.049×10^{-3}	0.0	6.153
43	8.722	276.7	2.19×10^{-3}	3.553×10^{-3}	0.0	6.103
44	9.377	274.4	2.37×10^{-3}	2.334×10^{-3}	0.0	5.895
45	7.350	278.3	3.62×10^{-4}	2.405×10^{-3}	1.912×10^{-2}	5.676
46	8.067	275.0	5.26×10^{-4}	3.161×10^{-3}	0.0	5.950
47	8.757	275.0	3.91×10^{-3}	3.658×10^{-3}	0.0	5.950
48	9.432	272.2	4.37×10^{-3}	4.207×10^{-3}	0.0	5.950
49	7.350	276.1	3.22×10^{-3}	2.510×10^{-3}	0.0	6.051
50	8.067	274.4	3.59×10^{-3}	3.197×10^{-3}	0.0	5.895
51	5.213	258.3	2.51×10^{-3}	1.689×10^{-3}	0.0	4.566
52	6.923	253.9	3.54×10^{-3}	3.257×10^{-3}	0.0	4.246
53	8.646	252.8	4.11×10^{-3}	4.260×10^{-3}	0.0	4.169
54	5.213	256.7	2.20×10^{-3}	1.777×10^{-3}	9.473×10^{-3}	4.203
55	6.902	254.4	2.61×10^{-3}	3.229×10^{-3}	0.0	4.282
56	8.646	253.3	3.17×10^{-3}	4.256×10^{-3}	0.0	4.204
57	3.413	233.9	7.57×10^{-4}	1.166×10^{-3}	9.002×10^{-3}	2.817
58	5.171	232.8	9.03×10^{-4}	2.706×10^{-3}	0.0	2.944
59	6.923	231.1	1.50×10^{-3}	3.644×10^{-3}	0.0	2.855

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Table 4.2 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Predictions		
	P_0 , MPa	T_0 , °C		\dot{m} Kg/sec	x_e	P_e MPa
60	8.646	228.9	1.85×10^{-3}	4.414×10^{-3}	0.0	2.742
61	8.646	231.7	1.75×10^{-3}	4.420×10^{-3}	0.0	2.886
62	7.033	228.3	1.76×10^{-3}	3.708×10^{-3}	0.0	2.712
63	5.206	225.0	1.53×10^{-3}	2.810×10^{-3}	0.0	2.550
64	3.448	222.8	1.33×10^{-3}	1.672×10^{-3}	0.0	2.447
65	3.461	235.0	1.84×10^{-3}	1.171×10^{-3}	1.050×10^{-2}	2.844
66	5.185	232.2	2.32×10^{-3}	2.688×10^{-3}	0.0	2.913
67	6.895	230.0	2.69×10^{-3}	3.640×10^{-3}	0.0	2.797
68	8.646	228.3	2.69×10^{-3}	4.418×10^{-3}	0.0	2.712
69	8.626	241.7	1.51×10^{-1}	1.797×10^{-1}	0.0	3.449
70	6.936	236.8	1.39×10^{-1}	1.455×10^{-1}	0.0	3.163
71	6.861	236.7	1.57×10^{-1}	1.440×10^{-1}	0.0	3.158
72	8.536	241.1	1.75×10^{-1}	1.741×10^{-1}	0.0	3.413
73	5.109	252.8	1.41×10^{-1}	5.656×10^{-2}	7.546×10^{-2}	2.447
74	6.861	253.9	1.69×10^{-1}	1.339×10^{-1}	2.205×10^{-2}	3.709
75	8.605	254.4	1.91×10^{-1}	1.731×10^{-1}	0.0	4.282
76	8.598	250.6	2.00×10^{-1}	1.747×10^{-1}	0.0	4.018
77	6.861	248.3	1.78×10^{-1}	1.371×10^{-1}	1.995×10^{-3}	3.818
78	5.192	246.1	1.51×10^{-1}	6.733×10^{-2}	5.056×10^{-2}	2.606
79	3.937	240.6	1.28×10^{-1}	3.762×10^{-2}	6.916×10^{-2}	1.985
80	5.089	236.1	1.55×10^{-1}	6.793×10^{-2}	2.330×10^{-2}	2.643

Table 4.2 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Predictions		
	P_0 , MP _a	T_0 , °C		\dot{m} Kg/sec	x_e	P_e MPa
81	6.771	237.2	1.76×10^{-1}	1.422×10^{-1}	0.0	3.186
82	8.577	234.4	1.99×10^{-1}	1.777×10^{-1}	0.0	3.030

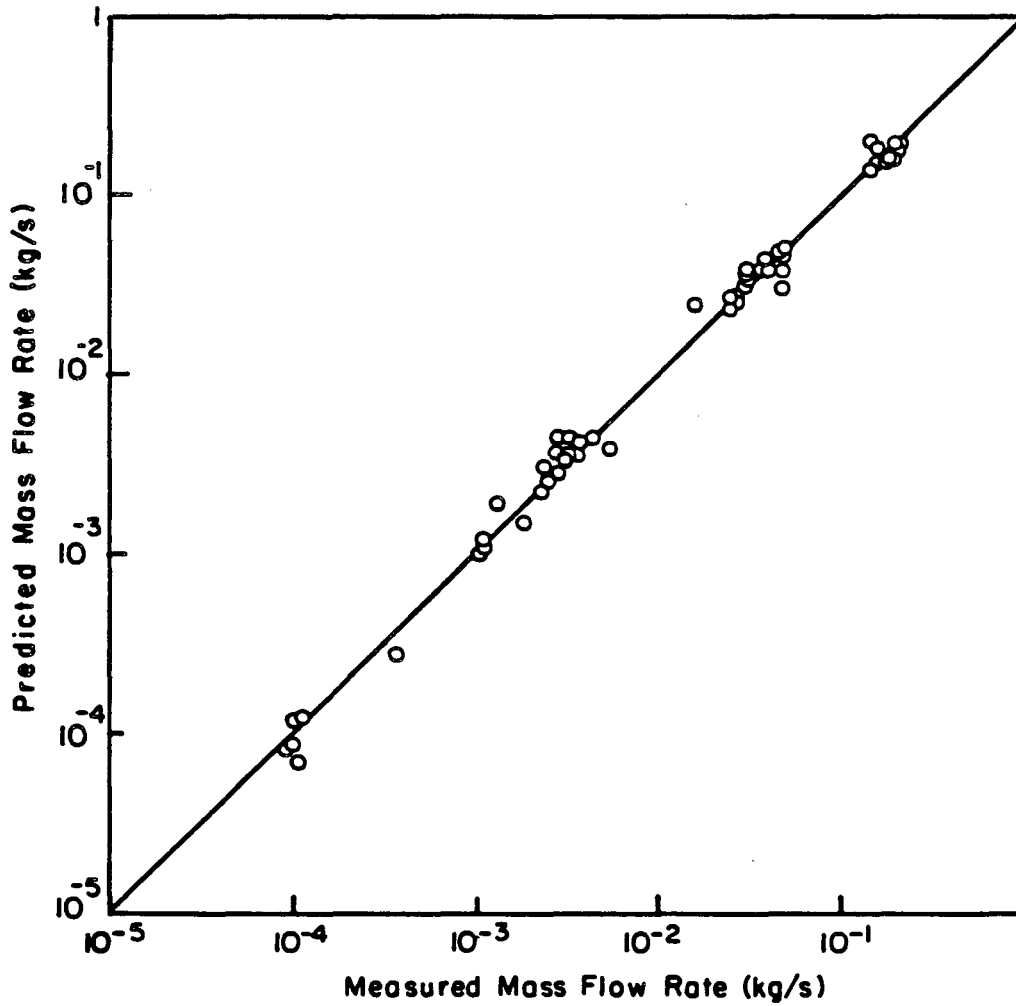


Figure 4.2 Corrected SOURCE Predictions (Based on Optimum Friction Factor) Compared with BCL Data

$K' \approx 15 f_p$ for 45° bends $K' \approx 30 f_p$ for 90° bends.

$$\text{For contractions } n_2 K_2 = (0.8 \sin \frac{\alpha}{2} (1-\beta^2)) n_2 \quad (4.5)$$

$$\text{For expansions } n_3 K_3 = (2.6 \sin \frac{\alpha}{2} (1-\beta^2)^2) \cdot n_3 \quad (4.6)$$

Here n_2, n_3 are respectively, number per unit length of contractions and expansions, α is the angle of convergence, and β is the ratio of lower to higher diameter.

Close observation of photomicrographs of the IGSCC reveals that the bends, contractions and expansions each appear in the crack channel at intervals of approximately 10δ . For evaluating the loss coefficients K_i , the bends in the crack were approximated on the average as 45° bends. The ratio r/D was approximated a constant 0.60. Then the loss coefficient K_1 for bends can be calculated as

$$K_1 = 8f_p + \frac{7f_p}{n_1 L_1} \quad (4.7)$$

To evaluate the loss coefficient due to contraction and expansions, the angle of convergence was approximated to be 45 degrees and $\beta = 0.5$ was used.

The number of bends per unit length was evaluated as the

$$n_1 = \frac{1}{10\delta} \quad (4.8)$$

The number of contractions or expansions per unit length was evaluated as

$$n_2 + n_3 = \frac{1}{10\delta} \quad \text{with } n_2 = n_3 \quad (4.9)$$

The method of evaluating the equivalent friction factor involves determining the friction f_p for the crack and then the loss coefficients are determined using equations (4.5) through (4.7). Then using the equation (4.3) the equivalent friction factor f is determined.

To determine the friction factor f_p , Reynolds number has to be evaluated. As the inlet and outlet area are different the Reynolds number was taken as the average value of the Reynolds numbers at inlet and

outlet plane of the crack. The surface roughness was taken as $\epsilon = 1.78\mu$. This value for surface roughness was used in the BCL LEAK 00 code. Depending on the relative surface roughness ϵ/D and Re the friction factor was determined from the standard charts.

The above methodology was applied for all the test cases of BCL test sections to evaluate the equivalent friction factor (see Appendix A). Using these values of equivalent friction factor the mass flow rate predictions were obtained for the 61 tests. SOURCE predictions using these friction factors, corrected for subcooling, are compared with BCL data in Table 4.3 and in Figure 4.3. Except for the prediction for Tests 1-6, the results compare well with BCL data, with a standard deviation of 15.90%. These predictions appear to be better than the prediction by Chexal et al. [16] with LEAK-01, which is shown in Figure 4.4.

As already noted, the test section used for Test 1-6 was evidently partially plugged probably during the machining. As a result the friction factor given by the generalized method does not represent that test section and the predicted flow rate is about 4 times the measured value. Although there is no conclusive evidence, it seems unlikely that such extensive plugging value, could occur in an actual IGSCC. If it does, the extent of the crack would be under-estimated by the observed leak rate. Further experimental investigation is needed using improved techniques to produce test cracks.

4.3 Quality and Pressure Profiles

SOURCE calculations develop predictions of the profile of fluid state along the flow path, including the state (critical state) at the point of choking, as well as the critical flow rate. In Appendix E we present graphs of the quality and pressure profiles predicted for 20 of the BCL test conditions. Negative quality characterizes the sub-cooled liquid region so that flashing starts at the point where $x = 0$.

The very steep slopes at the exit point up the need for very fine noding near the exit to avoid error in the numerical computations.

Table 4.3 SOURCE Predictions Using Friction Factors
from Generalized Methodology

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Prediction \dot{m} Kg/sec
	P_0 , MPa	T_0 , C		
1	7.316	272.8	1.02×10^{-3}	4.386×10^{-3}
2	7.316	272.8	1.02×10^{-3}	4.386×10^{-3}
3	8.667	267.2	1.10×10^{-3}	4.785×10^{-3}
4	9.412	260.6	1.10×10^{-3}	5.005×10^{-3}
5	9.412	260.6	1.10×10^{-3}	5.005×10^{-3}
6	9.322	267.8	1.10×10^{-3}	4.895×10^{-3}
12	5.868	260.0	1.09×10^{-4}	1.482×10^{-4}
14	5.868	272.8	1.01×10^{-4}	1.343×10^{-4}
15	3.523	251.7	9.30×10^{-5}	9.486×10^{-5}
16	3.379	241.1	1.05×10^{-4}	8.715×10^{-5}
17	5.681	278.9	1.01×10^{-4}	1.020×10^{-4}
19	7.309	273.9	3.04×10^{-2}	2.067×10^{-2}
20	7.309	282.2	2.90×10^{-2}	1.682×10^{-2}
21	9.005	280.0	4.85×10^{-2}	2.619×10^{-2}
22	9.005	273.3	4.62×10^{-2}	2.910×10^{-2}
23	8.694	256.7	4.52×10^{-2}	3.026×10^{-2}
24	8.888	256.7	4.44×10^{-2}	2.931×10^{-2}
25	7.164	256.1	3.67×10^{-2}	2.312×10^{-2}
26	7.164	256.7	3.59×10^{-2}	2.297×10^{-2}
27	5.571	260.6	2.51×10^{-2}	1.456×10^{-2}
28	5.626	267.8	2.35×10^{-2}	1.245×10^{-2}
29	5.592	241.7	3.01×10^{-2}	2.062×10^{-2}
30	5.592	243.9	2.98×10^{-2}	1.996×10^{-2}
31	7.219	242.8	3.96×10^{-2}	2.853×10^{-2}
32	7.219	238.3	3.94×10^{-2}	2.916×10^{-2}
33	8.812	241.7	4.51×10^{-2}	3.382×10^{-2}
34	8.812	235.0	4.73×10^{-2}	3.263×10^{-2}
35	7.129	226.1	3.84×10^{-2}	2.957×10^{-2}
36	7.129	222.8	3.92×10^{-2}	2.942×10^{-2}

Table 4.3 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Prediction \dot{m} Kg/sec
	P_0 , MPa	T_0 , C		
37	5.592	220.0	3.84×10^{-2}	2.573×10^{-2}
38	5.440	223.2	2.97×10^{-2}	2.435×10^{-2}
39	3.937	230.6	2.62×10^{-2}	1.546×10^{-2}
40	3.937	232.8	1.57×10^{-2}	1.316×10^{-2}
43	8.722	276.7	2.19×10^{-3}	1.664×10^{-3}
44	9.377	274.4	2.37×10^{-3}	1.848×10^{-3}
47	8.757	275.0	3.91×10^{-3}	1.603×10^{-3}
48	9.432	272.2	4.37×10^{-3}	1.857×10^{-3}
49	7.350	276.1	3.22×10^{-3}	1.385×10^{-3}
50	8.067	274.4	3.59×10^{-3}	1.436×10^{-3}
51	5.213	258.3	2.51×10^{-3}	0.978×10^{-3}
52	6.923	253.9	3.54×10^{-3}	1.451×10^{-3}
53	8.646	252.8	4.11×10^{-3}	1.891×10^{-3}
54	5.213	256.7	2.20×10^{-3}	0.910×10^{-3}
55	6.902	254.4	2.61×10^{-3}	1.409×10^{-3}
56	8.646	253.3	3.17×10^{-3}	1.933×10^{-3}
62	7.033	228.3	1.76×10^{-3}	1.487×10^{-3}
63	5.206	225.0	1.53×10^{-3}	1.117×10^{-3}
64	3.448	222.8	1.33×10^{-3}	0.731×10^{-3}
65	3.461	235.0	1.84×10^{-3}	0.717×10^{-3}
66	5.185	232.2	2.32×10^{-3}	1.206×10^{-3}
67	6.895	230.0	2.69×10^{-3}	1.614×10^{-3}
68	8.646	228.3	2.69×10^{-3}	2.313×10^{-3}
69	8.626	241.7	1.51×10^{-1}	2.219×10^{-1}
70	6.936	236.8	1.39×10^{-1}	1.876×10^{-1}
71	6.861	236.7	1.57×10^{-1}	1.931×10^{-1}
72	8.536	241.1	1.75×10^{-1}	2.275×10^{-1}
74	6.861	253.9	1.69×10^{-1}	1.690×10^{-1}
75	8.605	254.4	1.91×10^{-1}	2.421×10^{-1}
76	8.598	250.6	2.00×10^{-1}	2.344×10^{-1}

Table 4.3 (continued)

Test Number	Stagnation Conditions		Measured \dot{m} Kg/sec	Present Prediction \dot{m} Kg/sec
	P_o , MPa	T_o , C		
77	6.861	248.3	1.78×10^{-1}	2.344×10^{-1}
81	6.771	237.2	1.76×10^{-1}	1.901×10^{-1}
82	8.577	234.4	1.99×10^{-1}	2.328×10^{-1}

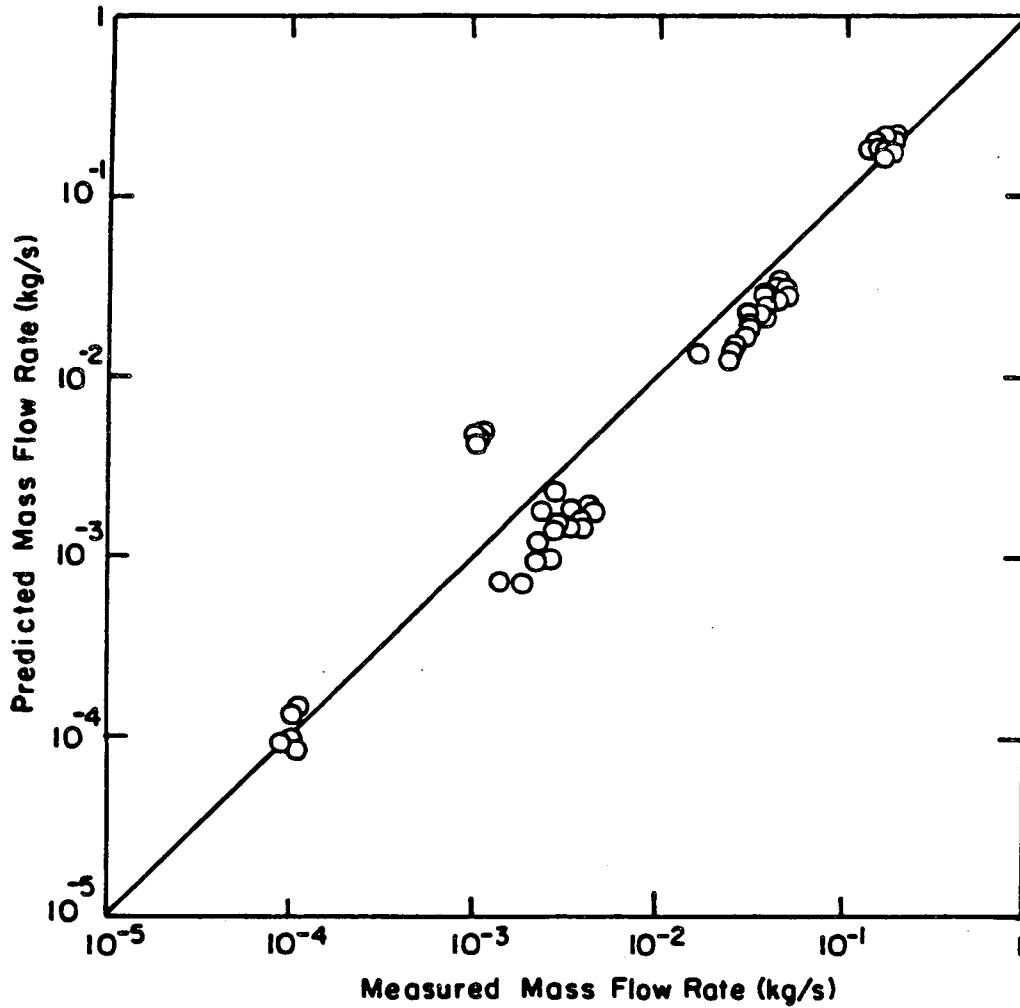


Figure 4.3 Corrected SOURCE Predictions (Based on Friction Factor from Generalized Method) Compared with BCL Data

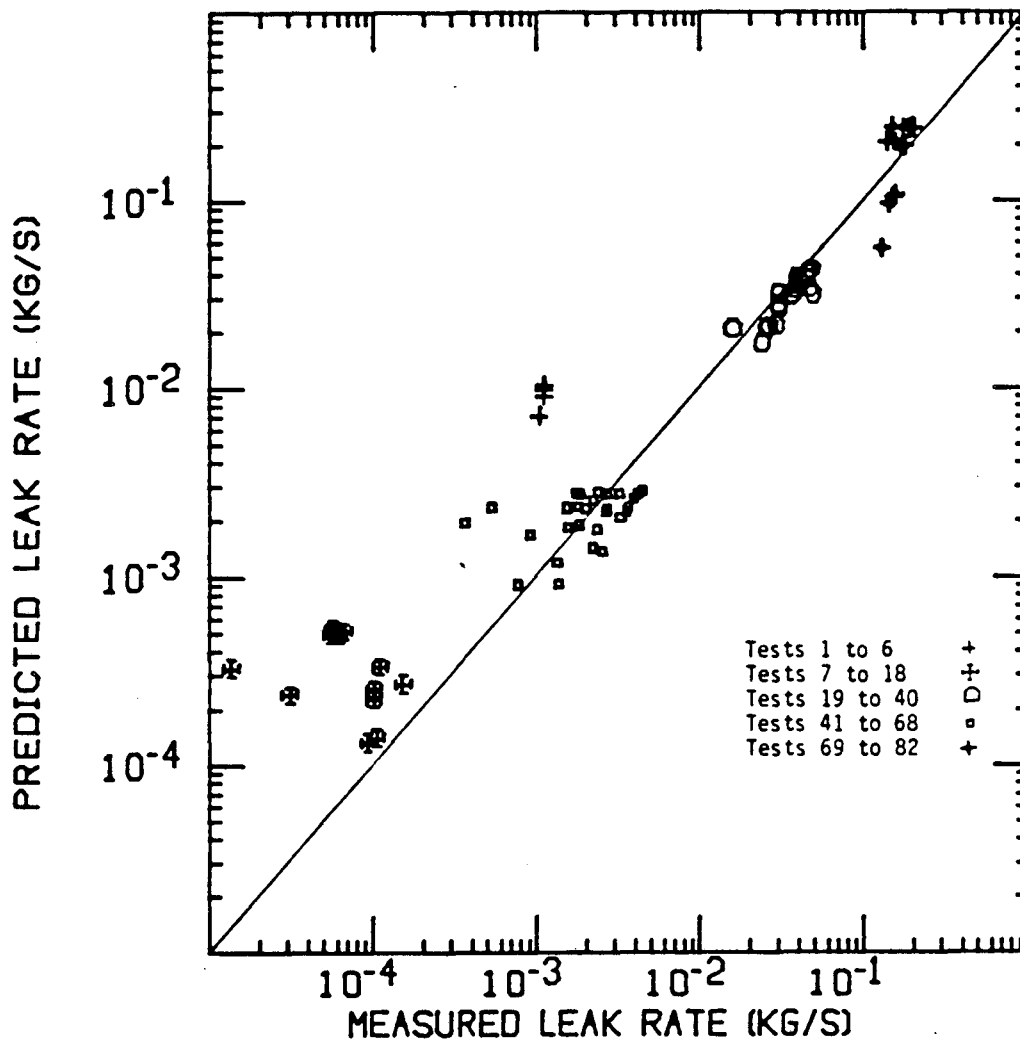


Figure 4.4 Chexal et al. [16] Predictions Compared with BCL Data

5. CONCLUSIONS

A computer code (SOURCE) has been developed to estimate leak rates through IGSCC cracks. The basis of the code is the steady state Homogeneous Equilibrium Model with friction and area change. A general method was developed to estimate the equivalent friction factor, lumping the effects of wall friction, tortuosity of the flow path and expansions and contractions, from physical features of IGSCC. Experimental results of Collier et al. (BCL) were used to develop a sub-cooling dependent correction factor to apply to SOURCE critical flow predictions. Using optimum friction factors deduced from BCL data for each test section, corrected SOURCE predictions agree with the BCL measured flowrates for 61 qualified tests with a standard deviation of 1.35% and for all 82 tests with a standard deviation of 6.5%. Using the independently estimated friction factor (from the generalized method) the predictions agree with qualified data with a standard deviation of 15.9%. This is an improvement as compared with Chexal's method.

The code SOURCE is very fast running and should be adaptable to large systems codes without significant sacrifice in cost. The method developed is recommended for estimating leak rates through IGSCC.

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APPENDIX A

Here we estimate the equivalent friction factor for each BCL test section using the generalized method described in Section 4.2.

i) Test section with $\delta = 0.074$ mm:

$$\text{the hydraulic diameter } D = \frac{4A}{P} = 0.145 \text{ mm.}$$

The grain sizes for thermally stressed steel pipes are around $\sqrt{2} \mu\text{m} = 2 \times 10^{-3}$ mm. This grain size is taken to yield surface roughness $\epsilon = 1.78 \mu = 1.78 \times 10^{-3}$ mm. Hence the relative roughness is

$$\epsilon/D = 1.78 \times 10^{-3} / 0.145 = 0.0122$$

The Reynolds number at the exit is 5.4×10^5 . Here we have evaluated the hydraulic diameter and the relative surface roughness at the exit. According to the present modelling of the crack geometry with constant δ , these parameters will be the same throughout the crack channel. However, the Reynolds number is smaller at the upstream of the crack exit. Hence an average Reynolds number is used. For relative surface roughness of 0.0122 the friction factor f is independent of Reynolds number and is given as

$$f_p = 8[2.46 \ln(D/\epsilon) + 3.22]^{-2}$$

Hence for $G/D = 0.0122$ we have $f = 0.048$.

Now we evaluate the loss coefficients $K_i n_i$ for various bends and expansion along the crack channel.

For this test section, $\delta = 0.074$, we have an average number of bends, 1.35 per mm and the average number of contraction and expansion, each of 0.67 per mm.

$$\begin{aligned} \text{Hence for } 45^\circ \text{ bends } K_1 n_1 &= 0.535/\text{mm} \\ \text{for contraction } K_2 n_2 &= 0.135/\text{mm} \text{ and} \\ \text{for expansion } K_3 n_3 &= 0.272/\text{mm} . \\ \therefore f &= f_p + D \sum_1 K_i n_i \\ &= 0.184 \end{aligned}$$

ii) Test section with $\delta = 0.022$ mm:

$$\text{Hydraulic diameter } D = 0.041 \text{ mm.}$$

$$\text{Relative roughness } \epsilon/D = 1.78 \times 10^{-3} / 0.041 = 0.043.$$

$$\text{Reynolds number range } 2 \times 10^5 \text{ to } 2 \times 10^4.$$

$$\therefore f_p = 0.07$$

Loss coefficients: Here $n_1 = 4.87$ $n_2 = n_3 = 2.43$.

$$K_1 n_1 = 2.75/\text{mm}$$

$$K_2 n_2 = 0.476/\text{mm}$$

$$K_3 n_3 = 0.989/\text{mm}$$

$$\begin{aligned}\therefore f &= f_p + D \sum_i K_i n_i \\ &= 0.243\end{aligned}$$

iii) Test section with $\delta = 0.108$

Hydraulic diameter $D = 0.212$

Relative roughness $\epsilon/D = 1.78 \times 10^{-2} / 0.212 = 0.0084$.

Exit Reynolds number range 8.7×10^6 to 2.7×10^6

Hence friction factor $f_p = 0.038$

Loss coefficients: Here $n_1 = 0.925$ $n_2 = n_3 = 0.46$

$$K_1 n_1 = 0.303/\text{mm}$$

$$K_2 n_2 = 0.098/\text{mm}$$

$$K_3 n_3 = 0.203/\text{mm}$$

$$\begin{aligned}\text{Hence } f &= f_p + D \sum_i K_i n_i \\ &= 0.166\end{aligned}$$

iv) Test section with $\delta = 0.048$ mm:

Hydraulic diameter $D = 0.092$

Relative roughness $\epsilon/D = 0.0193$

Exit Reynolds number range 5×10^6 to 1×10^6

Hence friction factor $f_p = 0.05$

Loss coefficients: Here $n_1 = 2.08$ $n_2 = n_3 = 1.04$

$$K_1 n_1 = 0.851/\text{mm}$$

$$K_2 n_2 = 0.203/\text{mm}$$

$$K_3 n_3 = 0.423/\text{mm}$$

$$\begin{aligned}\text{Hence } f &= f_p + D \sum_i K_i n_i \\ &= 0.186\end{aligned}$$

v) Test section with $\delta = 0.22$ mm:

Hydraulic diameter $D = 0.441$

Relative roughness $\epsilon/D = 0.00403$

Exit Reynolds number range: 1×10^7 to 2×10^7 .

Hence friction factor $f_p = 0.030$

Loss coefficients: Here $n_1 = 0.45$ $n_2 = n_3 = 0.227$

$$K_1 n_1 = 0.119/\text{mm}$$

$$K_2 n_2 = 0.045/\text{mm}$$

$$K_3 n_3 = 0.092/\text{mm}$$

$$\text{Hence } f = f_p + D \sum_i K_i n_i$$

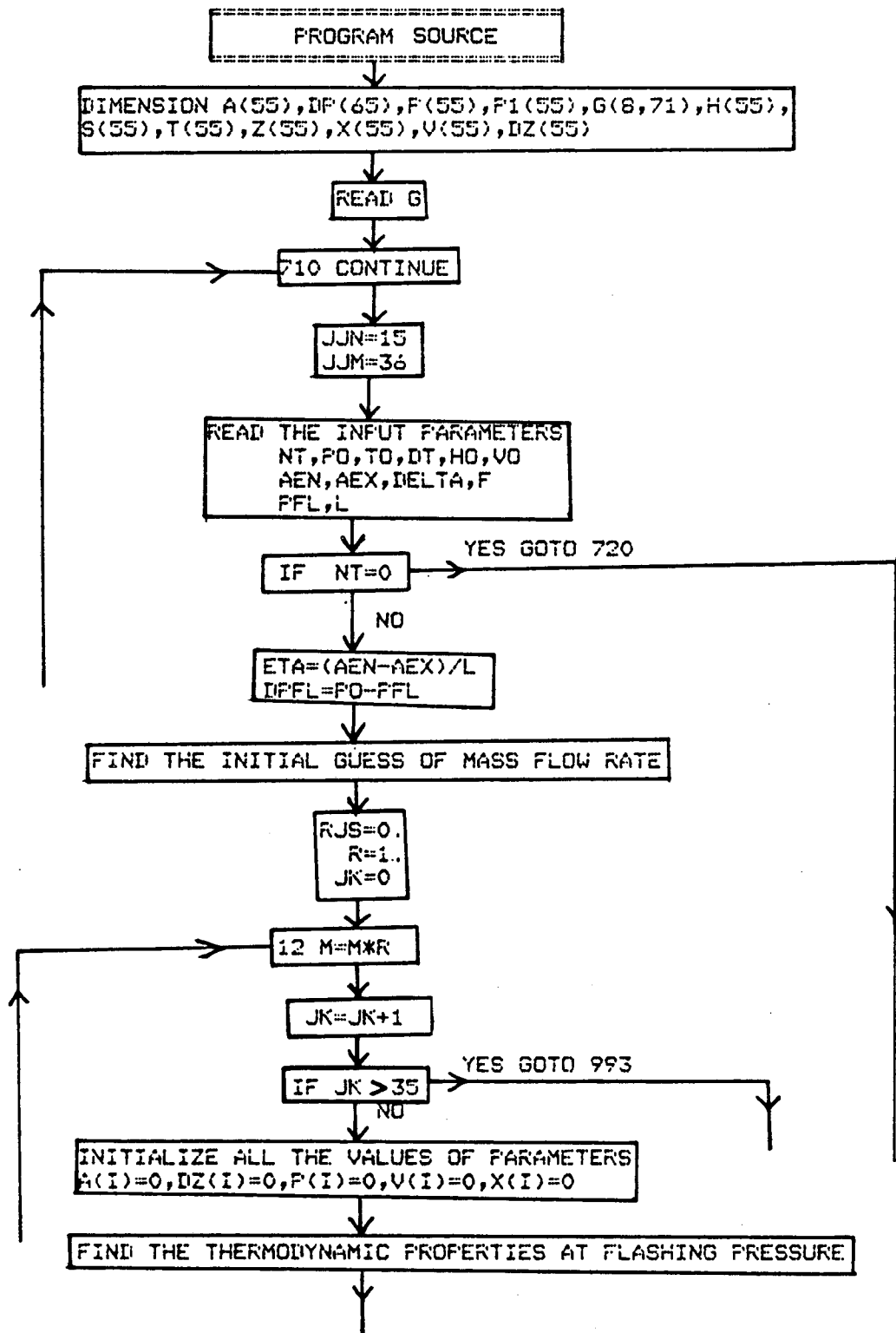
$$= 0.143$$

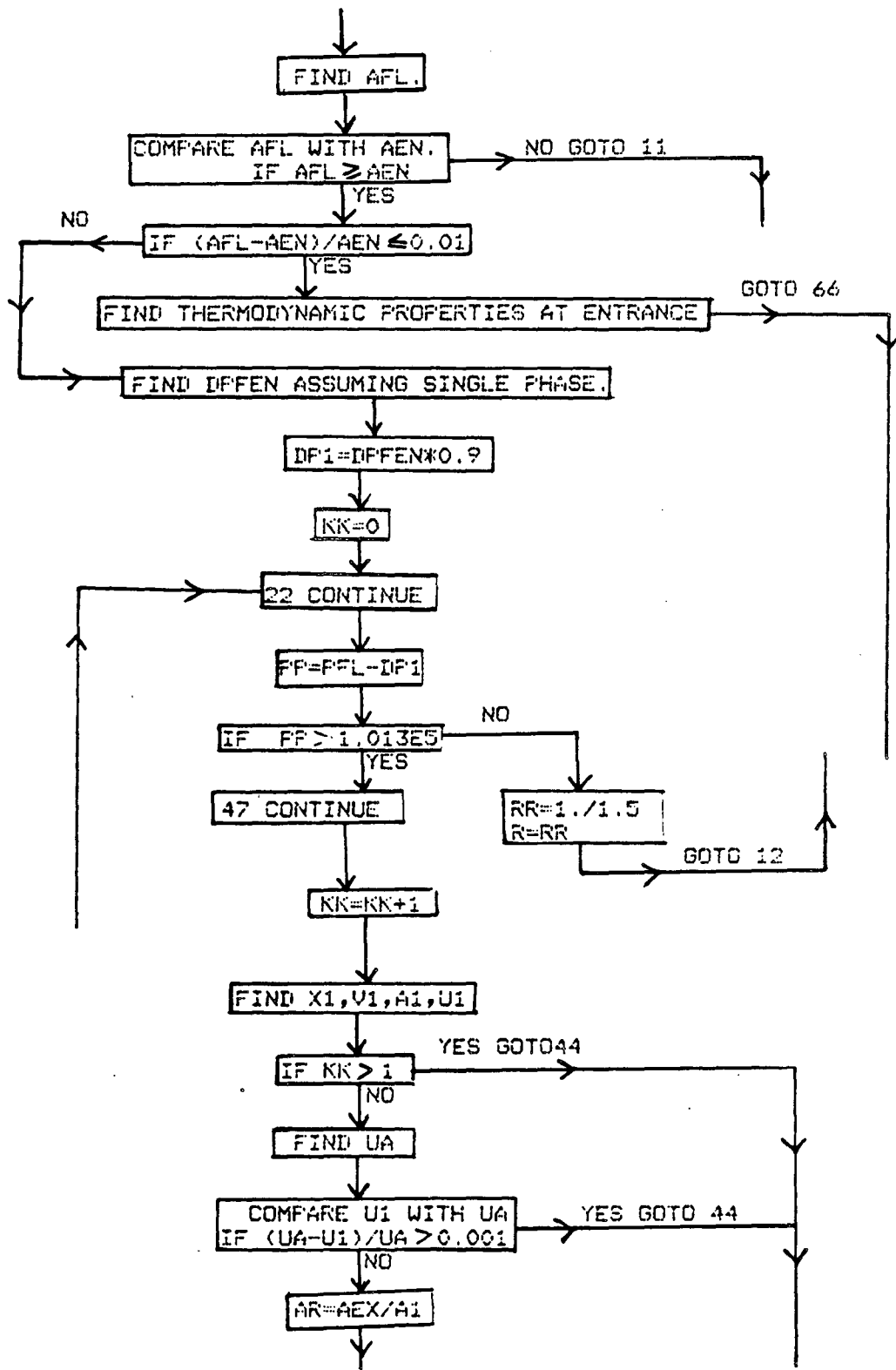
Table A.1 compares the friction factors obtained from the present generalized method with those obtained from parametric study.

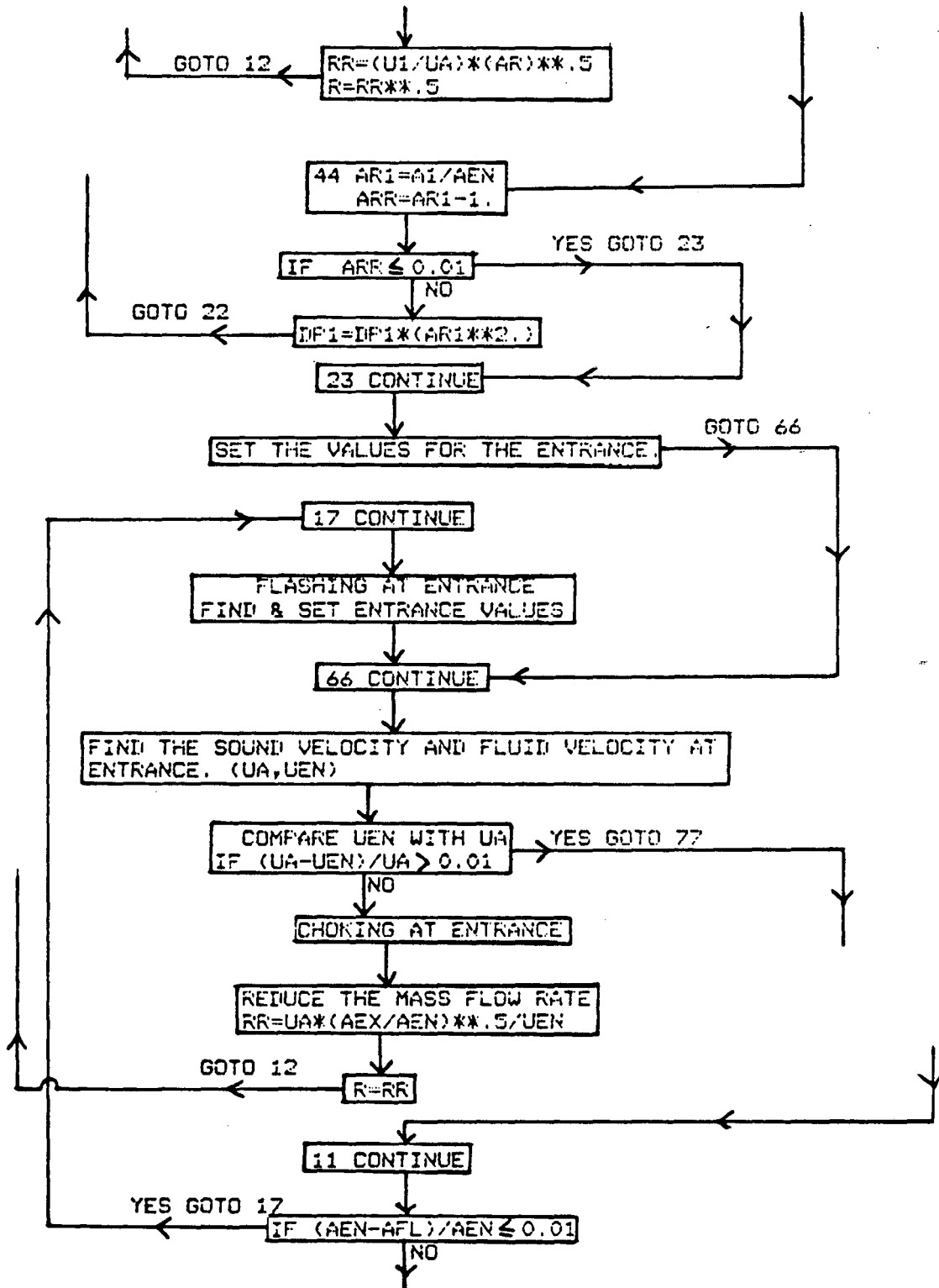
Table A.1

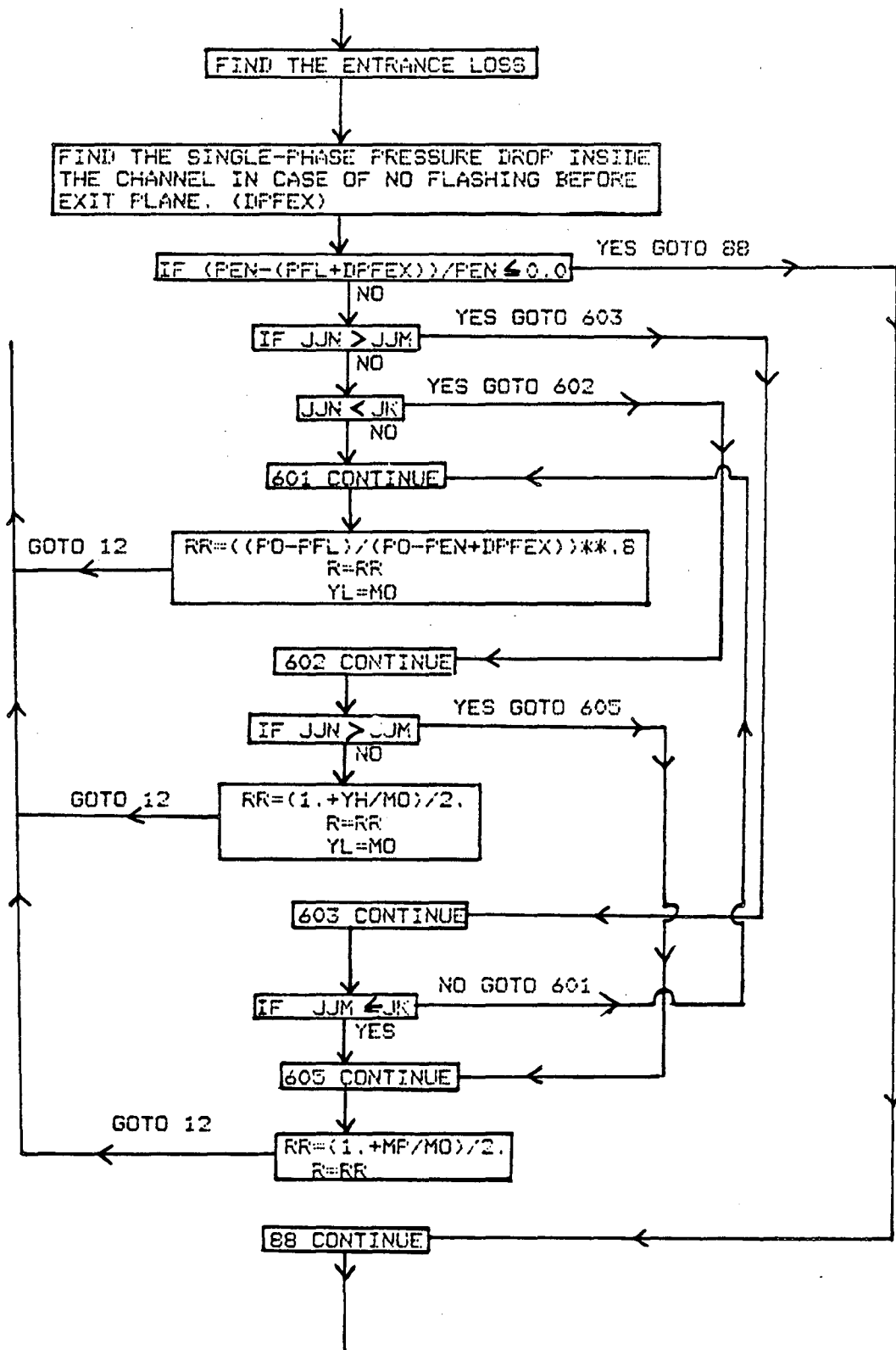
Comparison of friction factors calculated based on generalized methodology with those from parametric study

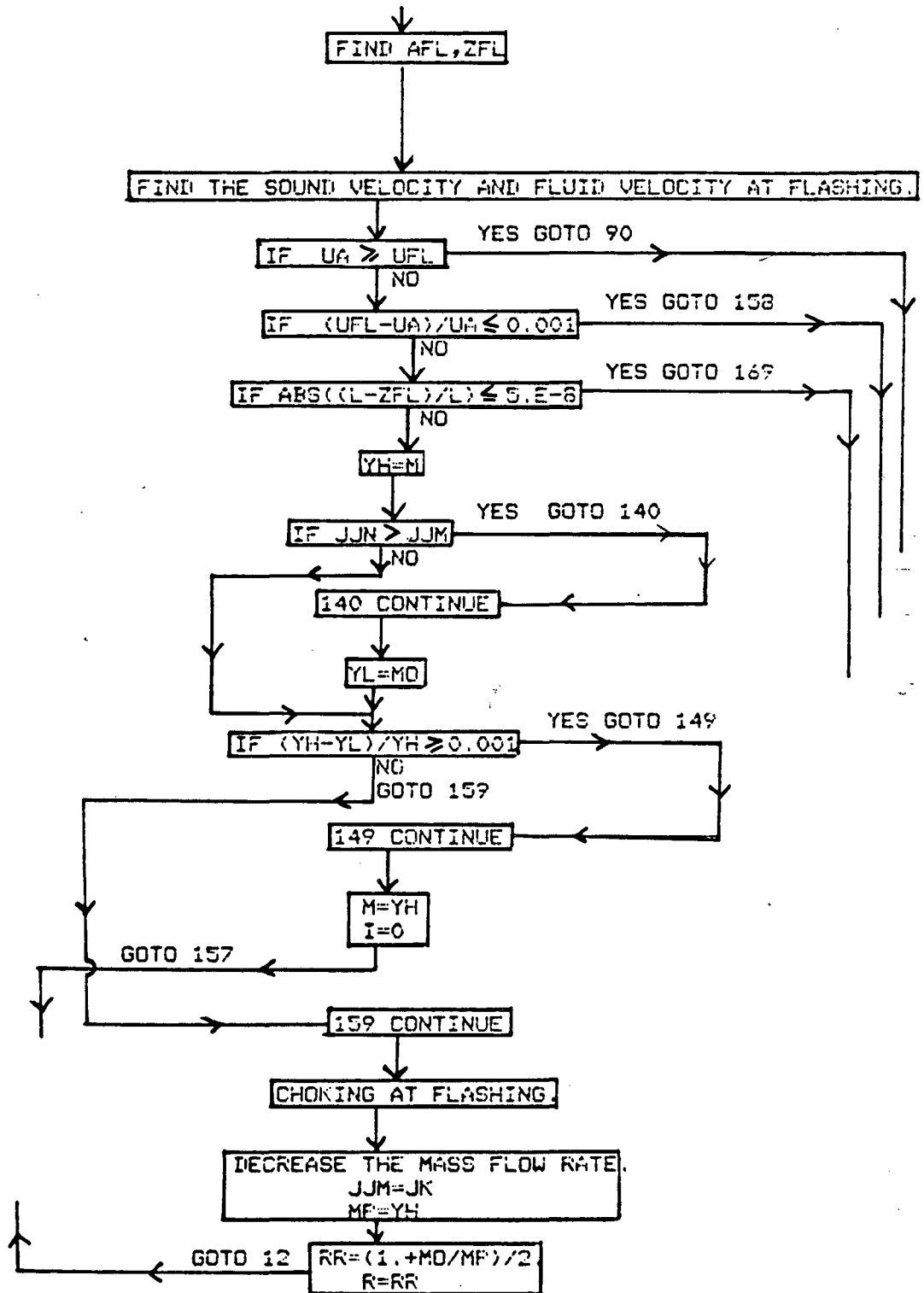
(mm)	D(mm)	Re (average)	f_p	$D \sum_i K_i n_i$	f	
					Generalized Methodology	Parametric Study
0.074	0.144	5.483×10^3	0.048	0.136	0.184	9.0
0.022	0.041	1.684×10^3	0.070	0.173	0.243	0.80
0.108	0.212	6.675×10^4	0.038	0.128	0.166	0.07
0.048	0.092	1.476×10^4	0.050	0.136	0.186	0.02
0.22	0.441	9.550×10^4	0.030	0.113	0.143	0.30

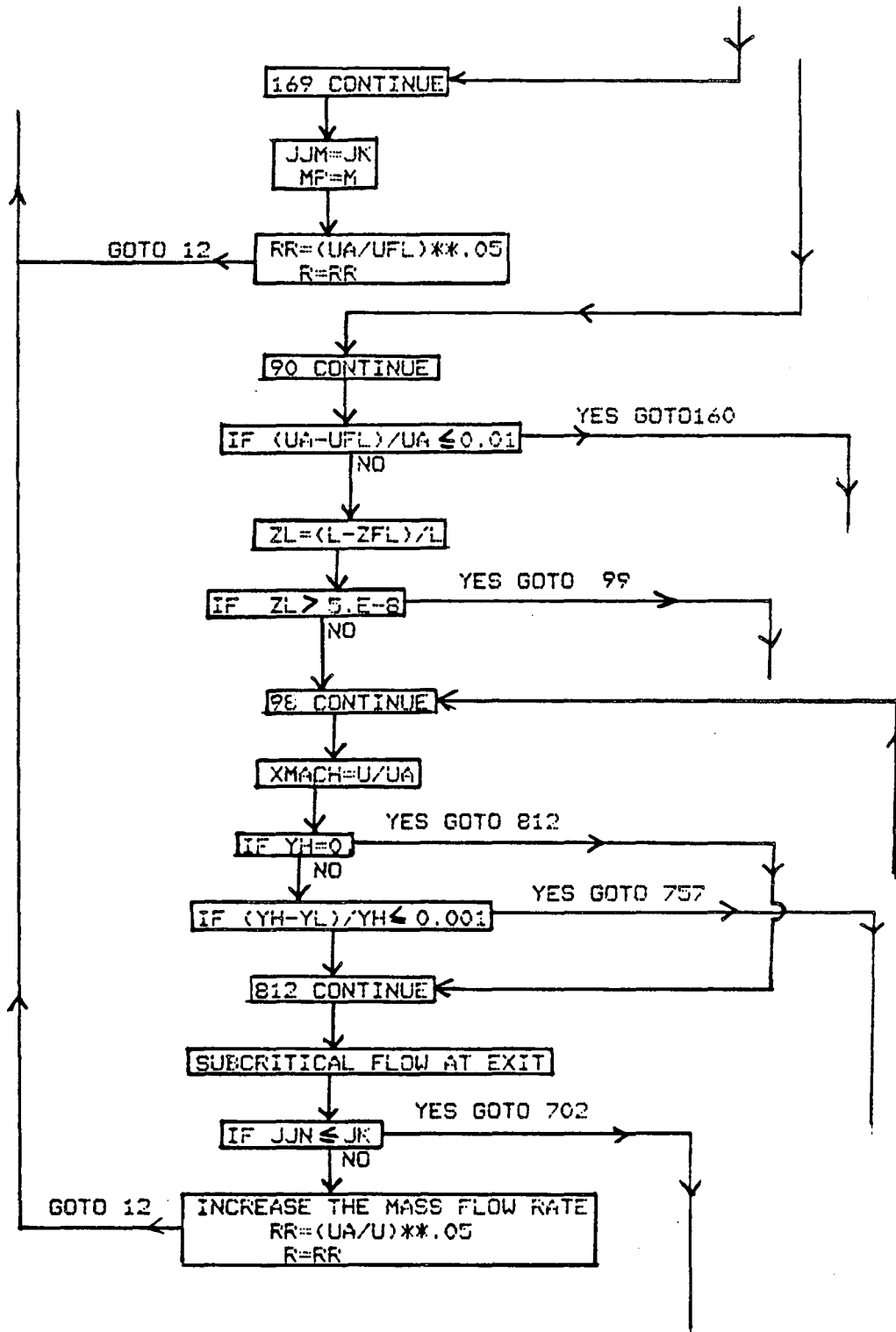


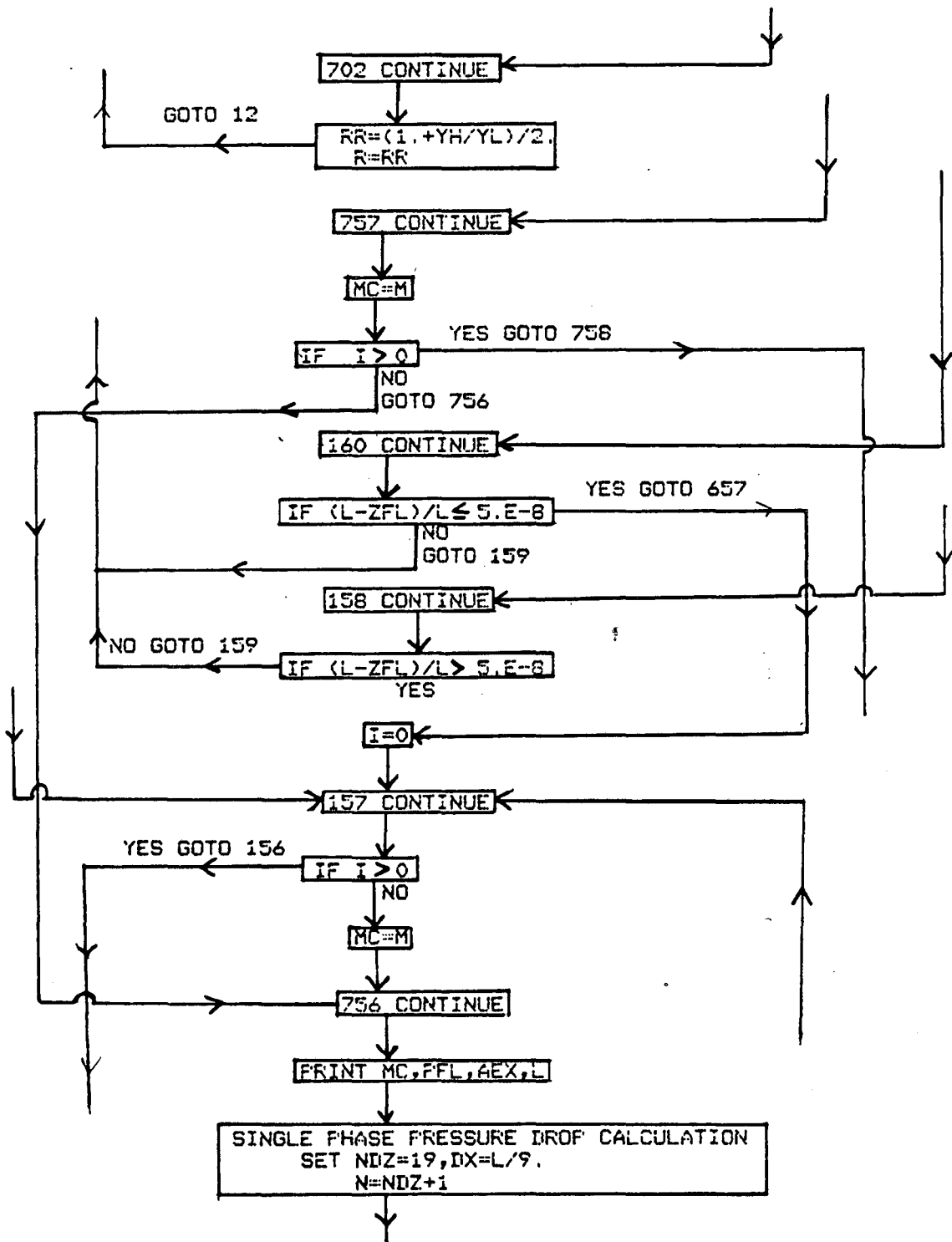


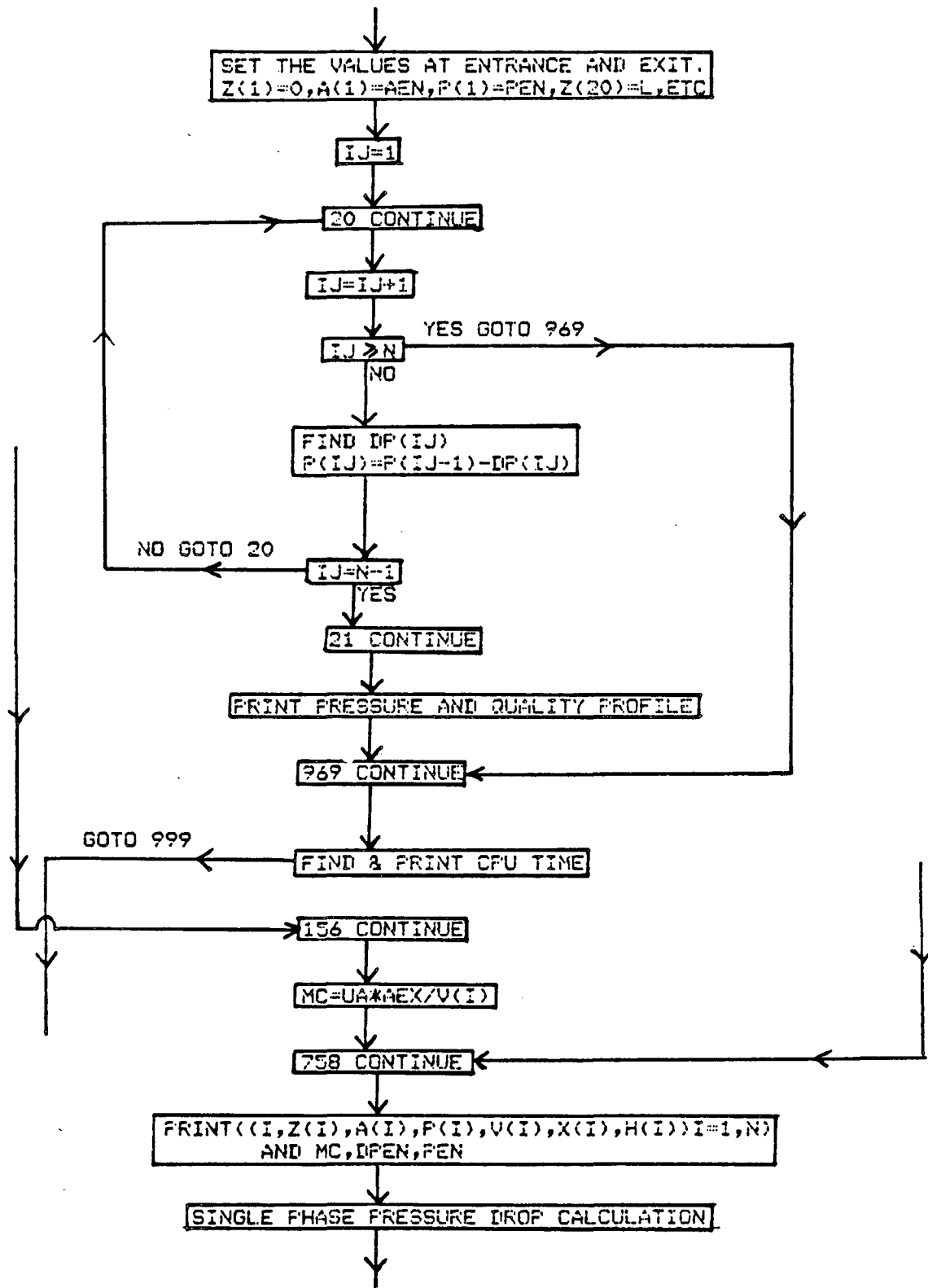


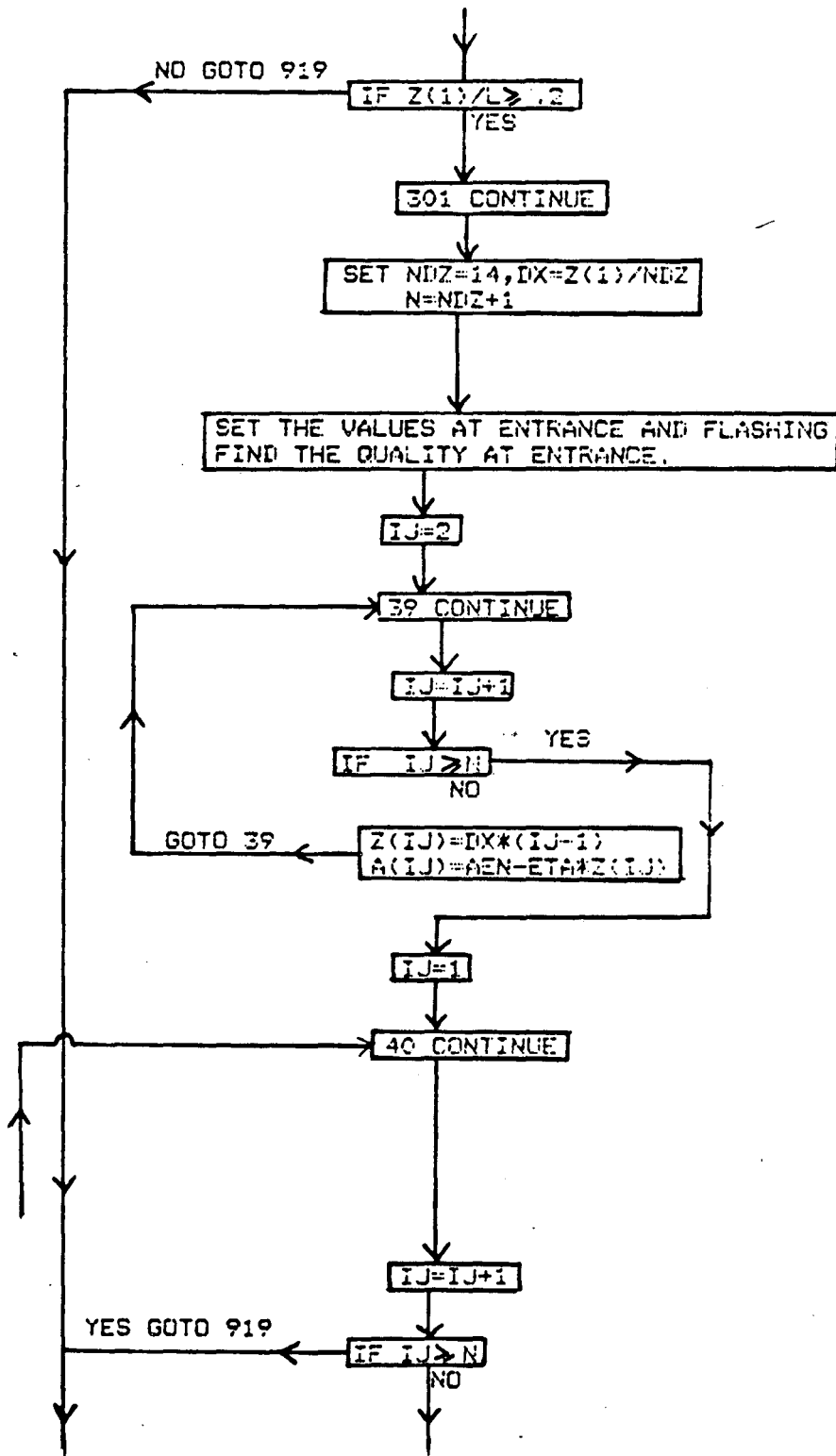


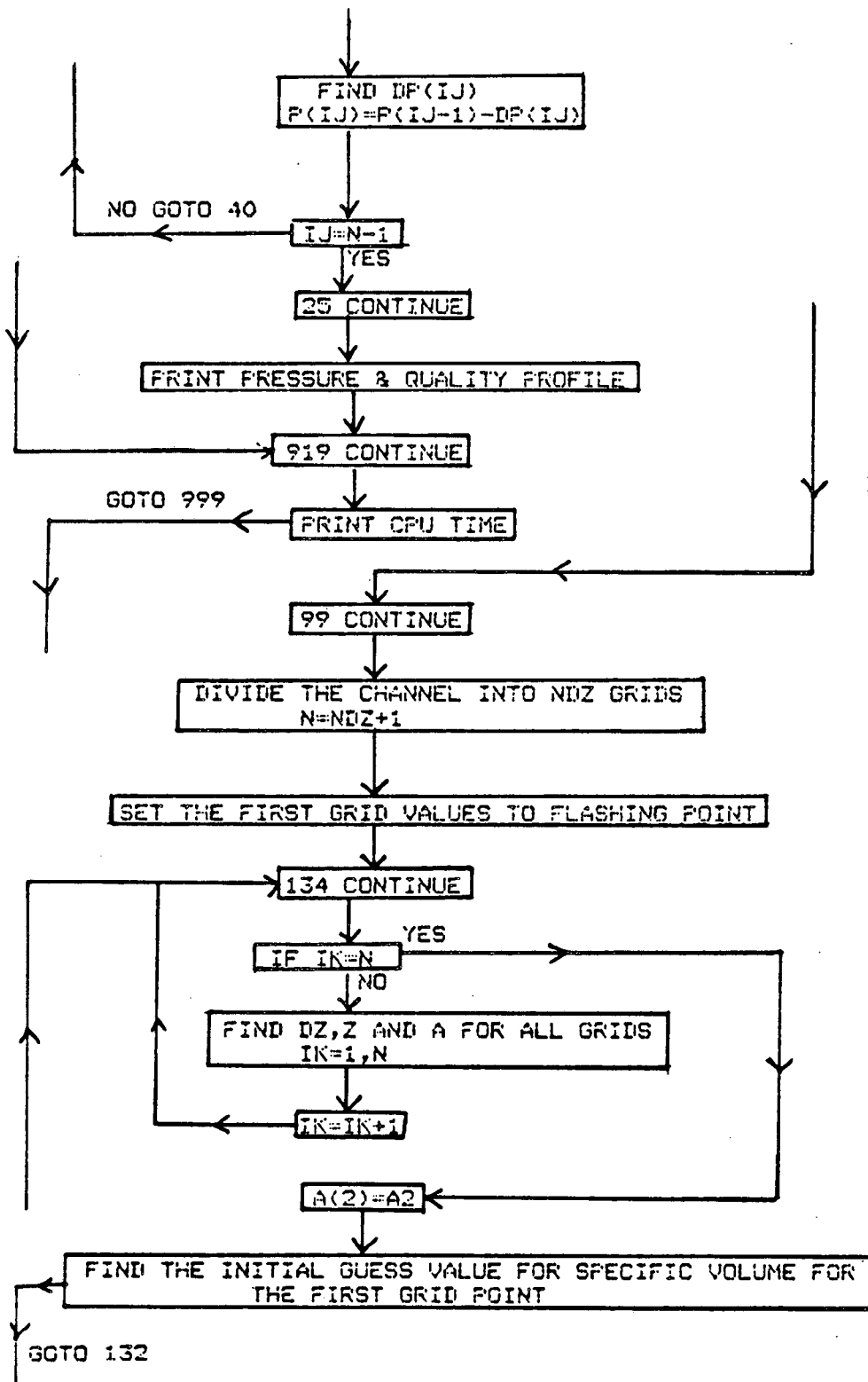


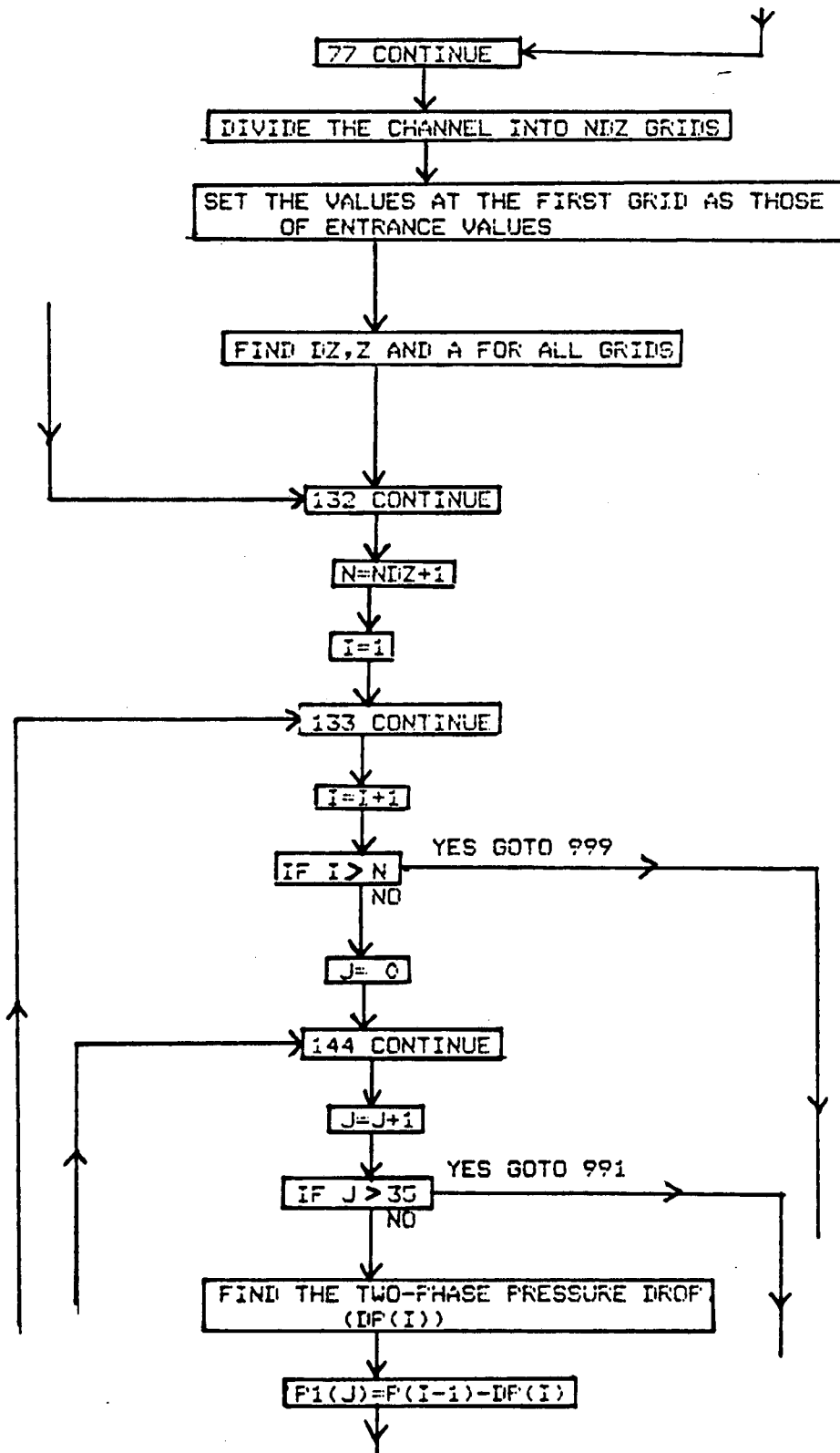


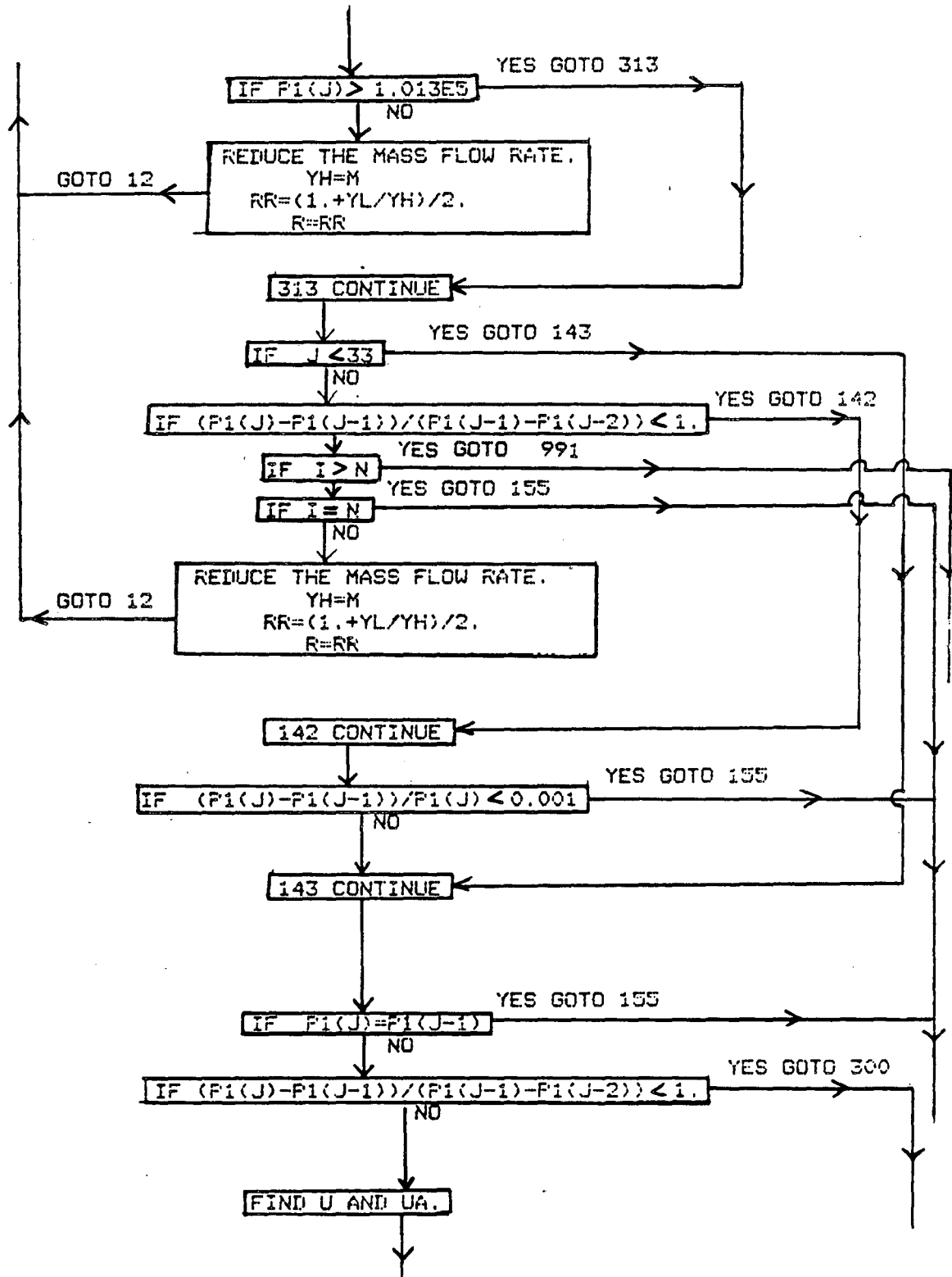


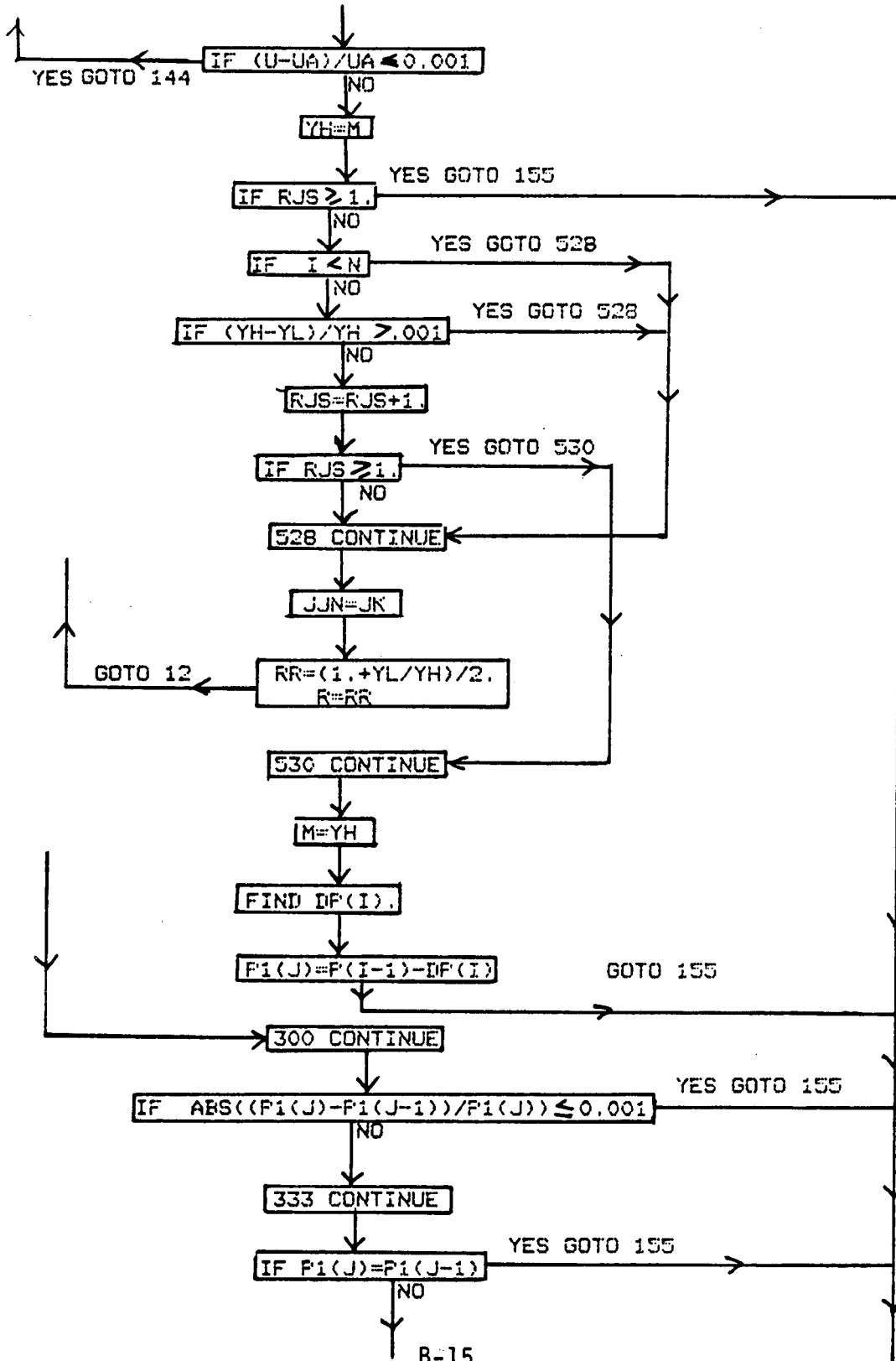


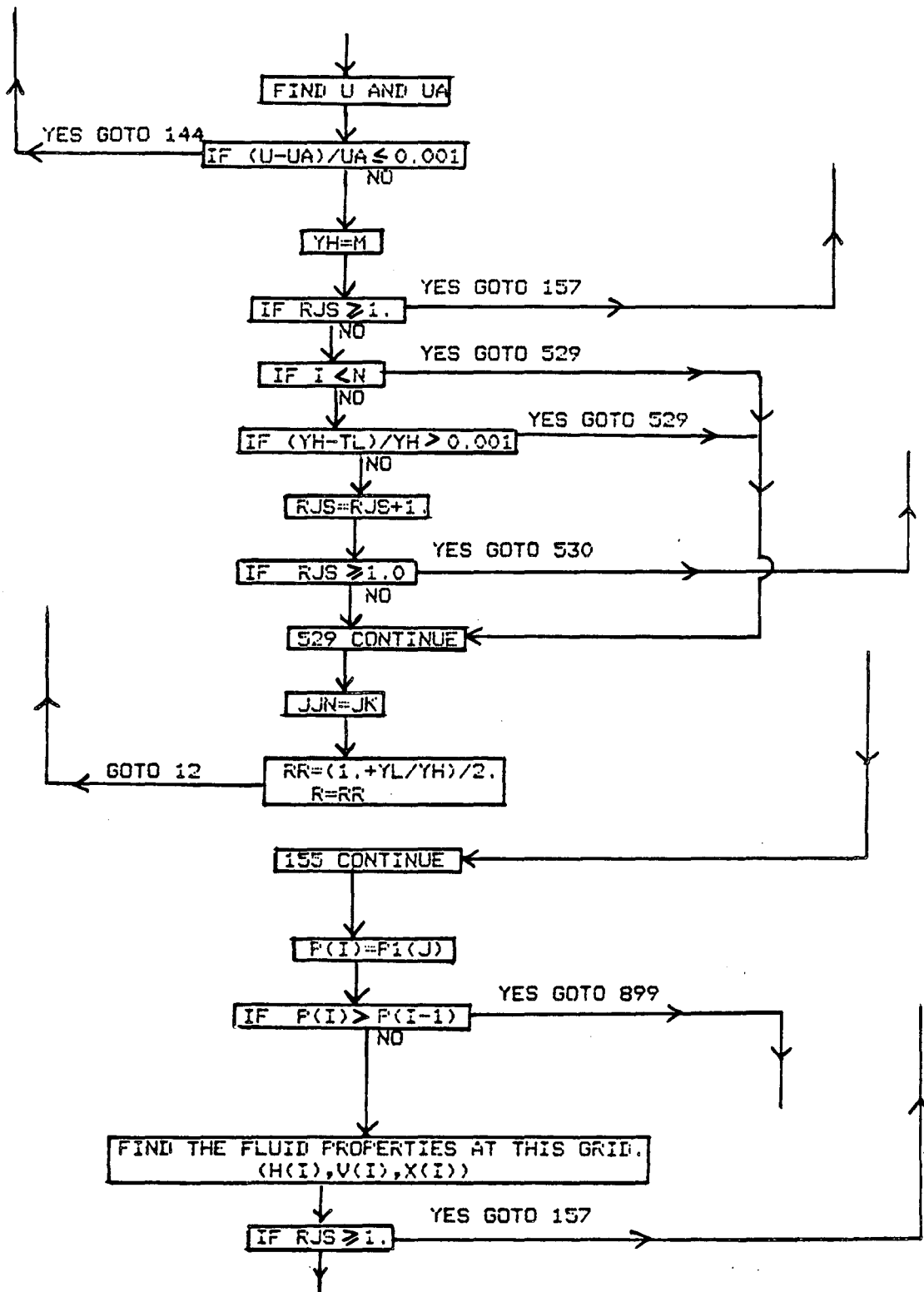


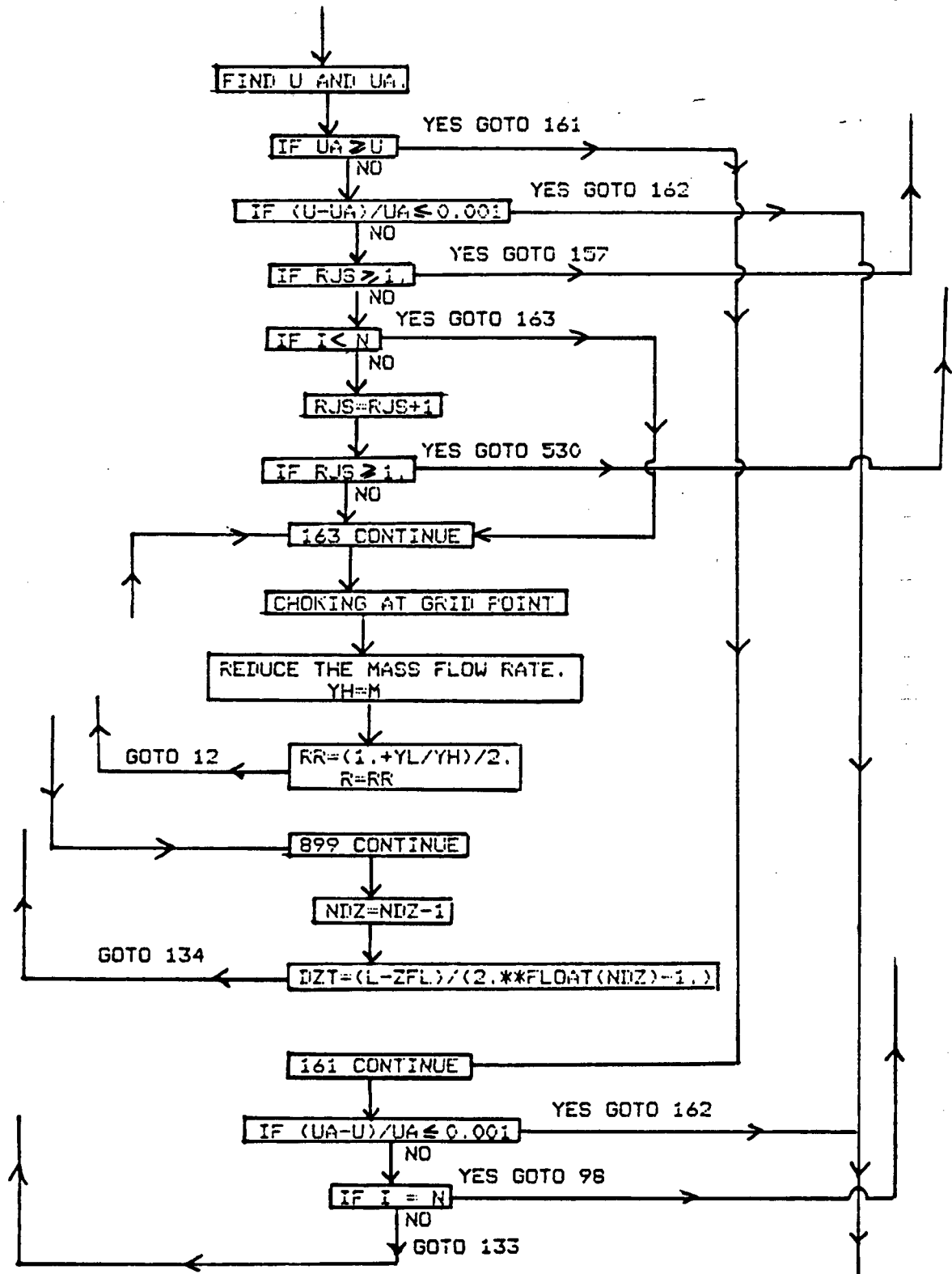


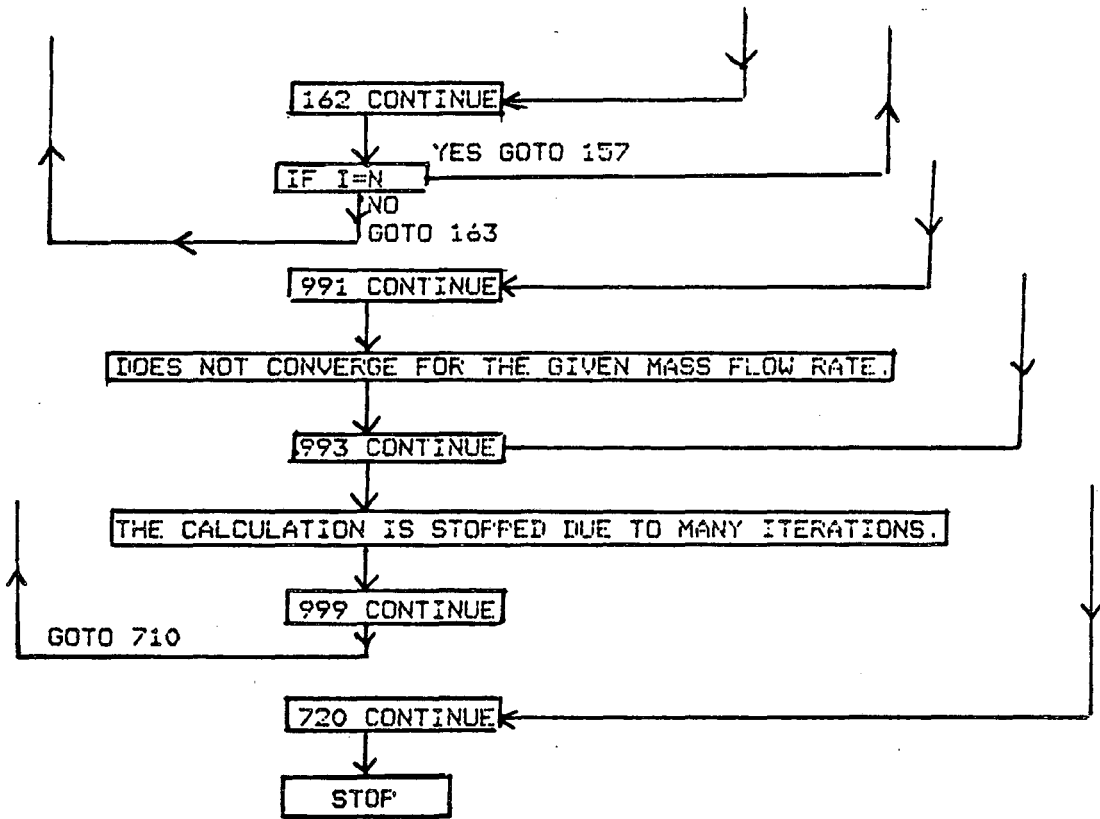












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C.....SOU00010
C
C      THIS IS A SOURCE PROGRAM      SOU00020
C
C      PROGRAM SOURCE                SOU00030
C
C                                     SOU00040
C                                     SOU00050
C                                     SOU00060
C-----SOU00070
C PURPOSE                            SOU00080
C THE PROGRAM CALCULATES THE CRITICAL MASS FLOW RATE THROUGH A INTER- SOU00090
C GRANULAR STRESS COROSION CRACKS ON PIPES SUCH AS FOUND IN THE SOU00100
C PRIMARY COOLANT SYSTEM OF NUCLEAR REACTOR PIPES. SOU00110
C-----SOU00120
C
C      IMPLICIT REAL*8(A,L,M)        SOU00130
C      REAL*4 ECPU, ETIME, ETCPU     SOU00140
C      CHARACTER*23 DATTIM           SOU00150
C      DIMENSION A(55),DP(55),P(55),P1(65),G(8,71),H(55),S(55),T(55), SOU00160
C      &DZ(55),X(55),V(55),Z(55)     SOU00170
C                                     SOU00180
C                                     SOU00190
C                                     SOU00200
C      COMMON/BLOCK1/VF,DVFD, VG,DVGD, HF,DHFD, HG,DHGD, SF,DSFD, SG, SOU00210
C      &DSGD                           SOU00220
C      COMMON/BLOCK2/M,VFL,ETA,DELTA,F,AFL,ZFL,NDZ,DZT,PFL,HFL SOU00230
C      OPEN(10,FILE='HEMD.DATA',STATUS='OLD',FORM='UNFORMATTED') SOU00240
C                                     SOU00250
C      VARIABLES USED IN THE PROGRAM SOU00260
C-----SOU00270
C VARIABLE                            DESCRIPTION SOU00280
C-----SOU00290
C A(I)    CROSS SECTIONAL AREA OF CRACK AT LOCATION Z(I) SOU00300
C AEN     ENTRANCE AREA SOU00310
C AEX     EXIT AREA SOU00320
C AAV     AVERAGE AREA BETWEEN TWO NODES SOU00330
C A1      DUMMY AREA BEFORE ENTRANCE SOU00340
C AR1     AREA RATIO AT TWO PHASE ENTRANCE A1/AEN SOU00350
C AFL     AREA AT FLASHING SOU00360
C AR      AREA RATIO AEX/A1 SOU00370
C                                     SOU00380
C C3      TOLARANCE FOR LENTGH SOU00390
C                                     SOU00400
C SA      HALF WIDTH OF CRACK CHANNEL AT ENTRANCE SOU00410
C SB      HALF WIDTH OF CRACK CHANNEL AT EXIT SOU00420
C                                     SOU00430
C DELTA   CRACK CHANNEL THICKNESS SOU00440
C DP      PRESSURE DROP BETWEEN TWO NODE POINTS SOU00450
C DPEN    ENTRANCE PRESSURE LOSS SOU00460
C DP1     DUMMY PRESSURE LOSS BETWEEN STAGN AND ENTRANCE SOU00470
C DH      HYDRAULIC DIAMETER OF CRACK SOU00480
C DPFEX   PRESSURE DROP FOR NO FLASH BEFORE EXIT SOU00490
C DPFL    PRESSURE DROP BETWEEN FLASH AND STAGN SOU00500
C DPFEN   PRESSURE DROP BETWEEN FLASH AND ENTRANCE TWO PHASE ENTRANCE SOU00510
C DPSP    SINGLE PHASE PRESSURE DROP AT ENTRANCE SOU00520
C DPSP1   SINGLE PHASE PRESSURE DROP BETWEEN FLASH AND FIRST GRID SOU00530
C DX      STEP SIZE IN SINGLE IN SINGLE-PHASE REGION SOU00540
C DZ(I)   STEP SIZE AT ITH GRID SOU00550
C DT      SUBCOOLING OF LIQUID SOU00560
C DZT     TERMINAL GRID SIZE SOU00570
C                                     SOU00580
C ETA     AREA SLOPE SOU00590
C                                     SOU00600

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C	F	FRICITION FACTOR	SOU00610
C			SOU00620
C	G	MASS FLUX	SOU00630
C			SOU00640
C	H(I)	ENTHALPY AT GRID I	SOU00650
C	HO	STAGNATION ENTHALPY	SOU00660
C	HF	LIQUID SATURATED ENTHALPY	SOU00670
C	HG	GAS SATURATED ENTHALPY	SOU00680
C	HEN	ENTRANCE LOCAL ENTHALPY	SOU00690
C	HFL	ENTHALPY AT FLASH	SOU00700
C			SOU00710
C	IJ	SPATIAL INDEX FOR SINGLE-PHASE REGION	SOU00720
C	I	SPATIAL INDEX FOR CHANNEL STEPS	SOU00730
C			SOU00740
C	J	INDEX FOR ITERATION ON MOMENTUM EQN	SOU00750
C			SOU00760
C	L	CHANNEL LENGTH	SOU00770
C			SOU00780
C	M	MASS FLOW RATE ASSUMED	SOU00790
C	MC	CRITICAL MASS FLOW RATE	SOU00800
C			SOU00810
C	NDZ	NUMBER OF SPATIAL STEPS	SOU00820
C	NT	TEST NUMBER	SOU00830
C			SOU00840
C	P(I)	LOCAL PRESSURE AT GRID I	SOU00850
C	PO	STAGN PRESSURE	SOU00860
C	PFL	PRESSURE AT FLASH	SOU00870
C	PEX	PRESSURE AT EXIT PLANE	SOU00880
C	PEN	PRESSURE AT ENTRANCE PLANE	SOU00890
C	PP	DUMMY LOCAL PRESSURE BEFORE ENTRANCE	SOU00900
C	PHI	MARTINNELI-NELSON FRICTION MULTIPLIER	SOU00910
C			SOU00920
C	R	CHANGE FACTOR ON MASS FLOW RATE	SOU00930
C	TR	INTENDED CHANGE FACTOR ON MASS FLOW RATE	SOU00940
C			SOU00950
C	S(I)	LOCAL ENTROPY AT GRID I	SOU00960
C	SO	STAGN ENTROPY	SOU00970
C	SFL	ENTROPY AT FLASH	SOU00980
C	SF	SATURATED LIQUID ENTROPY	SOU00990
C	SG	SATURATED GAS ENTROPY	SOU01000
C			SOU01010
C	T(I)	LOCAL SAT. TEMPERATURE	SOU01020
C	TFL	SATURATED TEMPERATURE	SOU01030
C	TO	STAGN TEMPERATURE	SOU01040
C			SOU01050
C	U	LOCAL FLUID VELOCITY	SOU01060
C	UJFL	FLUID VELOCITY AT FLASH	SOU01070
C	UA	SONIC VELOCITY	SOU01080
C			SOU01090
C	V(I)	LOCAL SPECIFIC VOLUME	SOU01100
C	VO	STAGN SP. VOLUME	SOU01110
C	VF	SATURATED SP. VOLUME OF LIQUID	SOU01120
C	VG	SATURATED GAS SPECIFIC VOLUME	SOU01130
C	VFL	SP VOLUME AT FLASH	SOU01140
C	VEN	SP VOLUME AT ENTRANCE	SOU01150
C	VEX	SP VOLUME AT EXIT	SOU01160
C			SOU01170
C	X(I)	QUALITY AT GRID I	SOU01180
C	XEN	ENTRANCE QUALITY	SOU01190
C	XEX	EXIT QUALITY	SOU01200

C		SOU01210
C	YH THE HIGHEST MASS FLOW RATE FOR ITERATION	SOU01220
C	YL THE LOWEST MASS FLOW RATE FOR ITERATION	SOU01230
C		SOU01240
C	Z(I) SPATIAL DISTANCE AT GRID I	SOU01250
C	ZFL FLASHING LOCATION	SOU01260
C		SOU01270
C	-----	SOU01280
C	ALL UNITS USED ARE IN STANDARD MKS AND TEMPERATURE IN DEG KELVIN	SOU01290
C	-----	SOU01300
C		SOU01310
C	READ MASS FLUX DATA	SOU01320
	READ(10) G	SOU01330
	710 CONTINUE	SOU01340
	JJN=15	SOU01350
	JJM=38	SOU01360
C	READ INPUT PARAMETERS	SOU01370
	READ(5,*)NT,PO,TO,DT,HO,VO	SOU01380
	READ(5,*)AEN,AEX,DELTA,F	SOU01390
	READ(5,*)PFL,L	SOU01400
	IF(NT.EQ.0)GOTO 720	SOU01410
	ETA=(AEN-AEX)/L	SOU01420
	DPFL=PO-PFL	SOU01430
C	PRINT INPUT PARAMETERS	SOU01440
	WRITE(6,100)NT	SOU01450
	WRITE(6,101)PO,TO,DT	SOU01460
	WRITE(6,102)AEX,AEN,L	SOU01470
	WRITE(6,103)F	SOU01480
	PRINT*, 'ETA=',ETA	SOU01490
	100 FORMAT(1H1,///T25,'TEST NUMBER = ',I5,/)	SOU01500
	101 FORMAT(//T5,'PO=',E15.7,2X,'TO=',E15.7,2X,'DT=',E15.7)	SOU01510
	102 FORMAT(1H+,T65,'AEX=',E15.7,2X,'AEN=',E15.7,2X,'L=',E15.7)	SOU01520
	103 FORMAT(/T30,' F =',E15.7)	SOU01530
C	TO FIND THE INITIAL GUESS VALUE OF M	SOU01540
	TRE=480.	SOU01550
	PRE=2.900E6	SOU01560
C	FIRST FIND THE G(I,J)	SOU01570
	S1=(TO-TRE)/10.+0.5	SOU01580
	K1=IFIX(S1)	SOU01590
	S2=(PO-PRE)/100000.+0.5	SOU01600
	K2=IFIX(S2)	SOU01610
	S3=G(K1,K2)	SOU01620
C	FIND L/DH AND THEN THE INITIAL GUESS M	SOU01630
	S4=L/(DELTA*2.)	SOU01640
	S5=DT/(TO-273.1)	SOU01650
	Q1=-0.0013506	SOU01660
	Q2=-0.0013663	SOU01670
	Q3=-0.0018824	SOU01680
	IF(S5.GT.0.2) THEN	SOU01690
	M=S3*AEX*(Q1*S4+0.52)	SOU01700
	ELSE IF(S5.GT.0.1) THEN	SOU01710
	M=S3*AEX*(Q2*S4+0.69)	SOU01720
	ELSE	SOU01730
	M=S3*AEX*(Q3*S4+0.97)	SOU01740
	END IF	SOU01750
	IF(M.GT.0.)GOTO 391	SOU01760
	IF(NT.LT.7) THEN	SOU01770
	M=1.0E-3	SOU01780
	ELSE IF(NT.LT.19) THEN	SOU01790
	M=1.0E-4	SOU01800

ELSE IF (NT.LT.41) THEN	SOU01810
M=2.E-2	SOU01820
ELSE IF (NT.LT.69) THEN	SOU01830
M=1.E-3	SOU01840
ELSE	SOU01850
M=1.0E-1	SOU01860
END IF	SOU01870
391 CONTINUE	SOU01880
RJS=0.	SOU01890
R=1.	SOU01900
JK=0	SOU01910
12 M=M*R	SOU01920
C IF ITERATION ON MASS FLOW RATE ADJUSTMENT GOES MORE THAN N TIMES	SOU01930
C STOP	SOU01940
JK=JK+1	SOU01950
IF (JK.GT.35) GOTO 993	SOU01960
C FIND THE FLASHING AREA ASSUMING SINGLE PHASE LIQUID AND HENCE FIND	SOU01970
C UFL	SOU01980
C INITIALISE ALL THE ARRAYS	SOU01990
DO 5 I=1,55	SOU02000
A(I)=0.	SOU02010
DZ(I)=0.	SOU02020
P(I)=0.	SOU02030
P1(I)=0.	SOU02040
V(I)=0.	SOU02050
X(I)=0.	SOU02060
5 CONTINUE	SOU02070
C FIND THE THERODYNAMIC PROPERTIES AT FLASH PRESSURE	SOU02080
CALL STEAM(PFL)	SOU02090
C SET THE VALUES AT FLASH	SOU02100
HFL=HF	SOU02110
VFL=VF	SOU02120
SFL=SF	SOU02130
AFL=(M**2.*V0*0.5/DPFL)**0.5	SOU02140
C COMPARE AFL WITH AEN TO SEE IF SINGLE PHASE ENTRANCE IF YES GO TO 11	SOU02150
IF (AFL.LT.AEN) GOTO 11	SOU02160
C COMPARE THE TOLARENCE	SOU02170
IF ((AFL-AEN)/AEN.LE.0.01) GOTO 17	SOU02180
C FIND PRESSURE DROP BETWEEN FLASH AND ENTRANCE	SOU02190
DPFEN=(M**2.*VFL*0.5/AEN**2.)*(1.-(AEN/AFL)**2.)	SOU02200
DP1=DPFEN*0.9	SOU02210
C SET FOR ITERATION AND FIND THERMODYNAMIC PROPERTIES AT PP	SOU02220
KK=0	SOU02230
22 CONTINUE	SOU02240
PP=PFL-DP1	SOU02250
IF (PP.GT.1.013E5) GOTO 47	SOU02260
RR=1./1.5	SOU02270
R=RR	SOU02280
GOTO 12	SOU02290
47 CONTINUE	SOU02300
KK=KK+1	SOU02310
CALL STEAM(PP)	SOU02320
X1=(S0-SF)/(SG-SF)	SOU02330
V1=VF+X1*(VG-VF)	SOU02340
H1=HF+X1*(HG-HF)	SOU02350
U1=(2.*(HFL-H1)+UFL**2.)**0.5	SOU02360
A1=M*V1/U1	SOU02370
IF (KK-1) 42,42,44	SOU02380
42 CONTINUE	SOU02390
C FIND SOUND VELOCITY	SOU02400

	DX1DP=-((DSFDP)*(1.-X1)+X1*DSGDP)/(SG-SF)	SOU02410
	UA=V1*((-((1.-X1)*DVFDP+X1*DVGDP+(VG-VF)*DX1DP))**(-0.5))	SOU02420
C	COMPARE U1 AND UA	SOU02430
	IF((UA-U1)/UA.GT.0.001)GOTO 44	SOU02440
	AR=AEX/A1	SOU02450
C	REDUCE MASS FLOW RATE	SOU02480
	RR=U1*(AR**0.5)/UA	SOU02470
	R=RR**0.5	SOU02480
	GOTO 12	SOU02490
	44 CONTINUE	SOU02500
C	CHECK FOR THE CONVERGENCE FOR AREA AT ENTRANCE	SOU02510
	AR1=A1/AEN	SOU02520
	ARR=AR1-1.	SOU02530
	IF(ARR.LE.0.01)GOTO 23	SOU02540
	DP1=DP1*(AR1**2.0)	SOU02550
	GOTO 22	SOU02560
	23 CONTINUE	SOU02570
C	SET THE ENTRANCE VALUES FOR THEY ARE NOW DETERMINED	SOU02580
	PEN=PP	SOU02590
	HEN=H1	SOU02600
	VEN=V1	SOU02610
	UEN=U1	SOU02620
	XEN=X1	SOU02630
	SEN=SF+X1*(SG-SF)	SOU02640
	GOTO 66	SOU02650
C	THE FOLLOWING CORRESPOND TO THE CASE OF FLASH AT ENTRANCE PLANE	SOU02660
	17 CONTINUE	SOU02670
	CALL STEAM(PFL)	SOU02680
	PEN=PFL	SOU02690
	HEN=HFL	SOU02700
	VEN=VF	SOU02710
	SEN=SF	SOU02720
	UEN=M*VEN/AEN	SOU02730
	XEN=0.	SOU02740
	66 CONTINUE	SOU02750
C	FIND THE SOUND VELOCITY AT ENTRANCE PLANE	SOU02760
	DX1DP=-((DSFDP)*(1.-XEN)+XEN*DSGDP)/(SG-SF)	SOU02770
	UA=VEN*((-((1.-XEN)*DVFDP+XEN*DVGDP+(VG-VF)*DX1DP))**(-0.5))	SOU02780
C	COMPARE UEN WITH UA. IT'S TWO-PHASE ENTRANCE IF YES ON GOTO 77.	SOU02790
	IF((UA-UEN)/UA.GT.0.001)GOTO 77	SOU02800
C	REDUCE MASS FLOW RATE	SOU02810
	RR=UA*((AEX/AEN)**0.5)/UEN	SOU02820
	R=RR	SOU02830
	GOTO 12	SOU02840
	11 CONTINUE	SOU02850
	IF((AEN-AFL)/AEN.LE.0.01)GOTO 17	SOU02860
C	DESIGNATE SINGLE PHASE AS TP=0.	SOU02870
	TP=0.	SOU02880
C	FIND ENTRANCE LOSS	SOU02890
	DPEN=M**2.*V0*0.5/AEN**2.	SOU02900
	PEN=P0-DPEN	SOU02910
C	FIND PRESSURE DROP FOR NO FLASH BEFORE EXIT PLANE	SOU02920
	AEXS=AEN*AEN+AEX*AEX	SOU02930
	V1=2.*M*M*(VFL-V0)/(AEXS)	SOU02940
	V2=M*M*0.25*(VFL+V0)*(1.+F*DELTA/ETA)*(1./AEX**2.-1./AEN**2.)	SOU02950
	V3=M*M*0.5*F*(VFL+V0)*(1./AEX-1./AEN)/(ETA*DELTA)	SOU02960
	DPFEX=V1+V2+V3	SOU02970
C	CHECK FOR FLASH POSITION INSIDE CHANNEL IF YES GOTO 88	SOU02980
	IF((PEN-(PFL+DPFEX))/PEN.LE.0.000)GOTO 88	SOU02990
C	INCREASE MASS FLOW RATE	SOU03000

	MO=M	SOU03010
	IF (JJN.GT.JJM) GOTO 603	SOU03020
	IF (JJN-JK) 602,602,601	SOU03030
601	CONTINUE	SOU03040
	RR=((PO-PFL) / (PO-PEN+DPFEX)) ** .8	SOU03050
	R=RR	SOU03060
	YL=MO	SOU03070
	GOTO 12	SOU03080
602	CONTINUE	SOU03090
	IF (JJN.GT.JJM) GOTO 605	SOU03100
	RR=(1.+YH/MO) / (2.)	SOU03110
	R=RR	SOU03120
	YL=MO	SOU03130
	GOTO 12	SOU03140
603	CONTINUE	SOU03150
	IF (JJM-JK) 605,605,601	SOU03160
605	CONTINUE	SOU03170
	RR=(1.+MP/MO) / (2.)	SOU03180
	R=RR	SOU03190
	GOTO 12	SOU03200
88	CONTINUE	SOU03210
C	FLASH INSIDE THE CHANNEL	SOU03220
	AFF=2.*(PEN-PFL)+M*M*VFL*(1.+DELTA*F/ETA)/AEN**2.*(2.*M*M*F*VFL/	SOU03230
	&(ETA*DELTA*AEN))	SOU03240
	AFL=(1.+(1.+((ETA**2.*DELTA**2.+DELTA**3.*ETA*F)/(M*M*F*F*VFL)))*	SOU03250
	&AFF)**0.5)/((ETA*DELTA/(M*M*F*VFL))*AFF)	SOU03260
	ZFL=(AEN-AFL)/ETA	SOU03270
C	WRITE ZFL AND AFL AT PFL	SOU03280
C	FIND THE SOUND VELOCITY AND FLUID VELOCITY	SOU03290
	CALL STEAM(PFL)	SOU03300
	XFL=0.	SOU03310
	VFL=VF	SOU03320
	HFL=HF	SOU03330
	SFL=SF	SOU03340
	DXDP=- (DSFDP) / (SG-SF)	SOU03350
	UA=VFL*((- (DVFD* (VG-VF) *DXDP)) ** (-0.5))	SOU03360
	UFL=(M*VFL)/AFL	SOU03370
C	COMPARE THE SOUND VELOCITY AND THE FLUID VELOCITY	SOU03380
	IF (UA.GE.UFL) GOTO 90	SOU03390
	IF ((UFL-UA) / UA.LE.0.001) GOTO 158	SOU03400
	IF (ABS((L-ZFL) / L).LE.0.5E-7) GOTO 169	SOU03410
	YH=M	SOU03420
	IF (JJN.GT.JJM) GOTO 140	SOU03430
	GOTO 141	SOU03440
140	CONTINUE	SOU03450
	YL=MO	SOU03460
141	CONTINUE	SOU03470
	IF ((YH-YL) / YH.LE.0.001) GOTO 149	SOU03480
	GOTO 159	SOU03490
149	CONTINUE	SOU03500
	M=YH	SOU03510
	I=0	SOU03520
	GOTO 157	SOU03530
159	CONTINUE	SOU03540
C	DECREASE THE MASS FLOW RATE	SOU03550
	JJM=JK	SOU03560
	MP=YH	SOU03570
	RR=(1.+MO/MP) / (2.)	SOU03580
	R=RR	SOU03590
	GOTO 12	SOU03600

169	CONTINUE	SOU03610
	JJM=JK	SOU03620
	MP=M	SOU03630
	RR=(UA/UFL)**.05	SOU03640
	R=RR	SOU03650
	GOTO 12	SOU03660
90	CONTINUE	SOU03670
	IF((UA-UFL)/UA.LE.0.001)GOTO 160	SOU03680
C	CHECK FOR LENGTH	SOU03690
	ZL=(L-ZFL)/L	SOU03700
	IF(ZL.GT.0.5E-7)GOTO 99	SOU03710
	U=UFL	SOU03720
98	CONTINUE	SOU03730
	XMACH=U/UA	SOU03740
	IF(YH.EQ.0.)GOTO 812	SOU03750
	IF((YH-YL)/YH.LE.0.001)GOTO 757	SOU03760
812	CONTINUE	SOU03770
	YL=M	SOU03780
	XMACH=U/UA	SOU03790
C	INCREASE MASS FLOW RATE	SOU03800
	IF(JJN-JK)702,702,701	SOU03810
701	CONTINUE	SOU03820
	RR=(UA/U)**0.05	SOU03830
	R=RR	SOU03840
	GOTO 12	SOU03850
702	CONTINUE	SOU03860
	RR=(1.+YH/YL)/(2.)	SOU03870
	R=RR	SOU03880
	GOTO 12	SOU03890
757	CONTINUE	SOU03900
	MC=M	SOU03910
	IF(I.GT.0)GOTO 758	SOU03920
	GOTO 758	SOU03930
C	CONTINUE FOR STEP 158	SOU03940
160	CONTINUE	SOU03950
	IF((L-ZFL)/L.LE.50E-7)GOTO 657	SOU03960
	GOTO 159	SOU03970
158	CONTINUE	SOU03980
C	FIRST CHECK FOR LENGTH WHEATHER AT EXIT	SOU03990
	IF((L-ZFL)/L.GT.0.50E-7)GOTO 159	SOU04000
657	I=0	SOU04010
157	CONTINUE	SOU04020
C	WRITE CHOKE EXACTLY AT EXIT AND FIND CRITICAL MASS FLOW RATE	SOU04030
	IF(I.GT.0)GOTO 158	SOU04040
	MC=M	SOU04050
758	CONTINUE	SOU04060
	PRINT*, 'THE FINAL CRITICAL MASS FLOW RATE IS',MC	SOU04070
	PRINT*, '.....'	SOU04080
	WRITE(6,113)L,AEX,PFL,MC	SOU04090
113	FORMAT(///T25,'RIGHT CONDITIONS ATTAINED AT EXIT FLASHING',//T5	SOU04100
	&,'Z(I)=' ,E15.7,'AEX=' ,E15.7,'PFL=' ,E15.7,2X,'MC=' ,E15.7)	SOU04120
C	FIND THE PRESSURE DISTRIBUTION IN SINGLE-PHASE REGION.	SOU04130
	NDZ=19	SOU04140
	N=NDZ+1	SOU04150
	DX=L/19.	SOU04160
	Z(1)=0.0	SOU04170
	A(1)=AEN	SOU04180
	P(1)=PEN	SOU04190
	P(20)=PFL	SOU04200

Z(20)=L	SOU04210
A(20)=AEX	SOU04220
X(20)=0.0	SOU04230
DO 19 IJ=2,19	SOU04240
Z(IJ)=DX*(IJ-1)	SOU04250
A(IJ)=AEN-ETA*Z(IJ)	SOU04260
19 CONTINUE	SOU04270
CALL STEAM(PEN)	SOU04280
X(1)=(HO-HF)/(HG-HF)	SOU04290
M=MC	SOU04300
IJ=1	SOU04310
20 CONTINUE	SOU04320
IJ=IJ+1	SOU04330
IF(IJ.GE.N)GOTO 969	SOU04340
DP(IJ)=M*M*0.5*VO*(1.+DELTA*F/ETA)*(1./A(IJ)**2.-1./A(IJ-1)**2.)	SOU04350
&*M*M*F/(2.*DELTA*ETA)*(2.*VO)*(1./A(IJ)-1./A(IJ-1))	SOU04360
P(IJ)=P(IJ-1)-DP(IJ)	SOU04370
CALL STEAM(P(IJ))	SOU04380
X(IJ)=(HO-HF)/(HG-HF)	SOU04390
IF(IJ.EQ.(N-1))GOTO 21	SOU04400
GOTO 20	SOU04410
21 CONTINUE	SOU04420
WRITE(6,478)	SOU04430
478 FORMAT(///T6,'I',7X,'Z(I)',15X,'A(I)',15X,'P(I)',15X,'X(I)')	SOU04440
DO 248 IJ=1,N	SOU04450
248 WRITE(6,249)IJ,Z(IJ),A(IJ),P(IJ),X(IJ)	SOU04460
249 FORMAT(/6X,I2,4(2X,E15.7))	SOU04470
969 CONTINUE	SOU04480
CALL DATETM(DATTIM,23,VCPU,CTIME,TCPU)	SOU04490
WRITE(6,801)DATTIM,VCPU,CTIME,TCPU	SOU04500
GOTO 999	SOU04510
156 CONTINUE	SOU04520
MC=UA*AEX/V(I)	SOU04530
768 CONTINUE	SOU04540
WRITE(6,114)I,P(I),I,X(I),I,V(I),MC	SOU04550
114 FORMAT(//T25,'RIGHT CONDITIONS ATTAINED AT EXIT PLANE',//T6,'P(',	SOU04560
&I2,')=',E15.7,'X(',I2,')=',E15.7,'V(',I2,')=',E15.7,'MC =',E15.7)	SOU04570
PRINT*, 'THE FINAL CRITICAL MASS FLOW RATE IS',MC	SOU04580
PRINT*, '.....'	SOU04590
&***'	SOU04600
PRINT*, 'DPEN=',DPEN,'PEN=',PEN	SOU04610
WRITE(6,477)	SOU04620
477 FORMAT(///T6,'I',7X,'Z(I)',15X,'A(I)',15X,'P(I)',15X,'V(I)',12X,'	SOU04630
&X(I)',12X,'H(I)')	SOU04640
Z(1)=ZFL	SOU04650
A(1)=AFL	SOU04660
DO 662 KKK=2,N	SOU04670
DZ(KKK)=DZT*2.** (FLOAT(N-KKK))	SOU04680
Z(KKK)=Z(KKK-1)+DZ(KKK)	SOU04690
A(KKK)=AEN-ETA*Z(KKK)	SOU04700
662 CONTINUE	SOU04710
DO 246 I=1,N	SOU04720
246 WRITE(6,247) I,Z(I),A(I),P(I),V(I),X(I),H(I)	SOU04730
247 FORMAT(/6X,I2,6(2X,E15.7))	SOU04740
IF(Z(1)/L.GE.0.2)GOTO 301	SOU04750
GOTO 919	SOU04760
C FIND THE PRESSURE DISTRIBUTION IN SINGLE-PHASE REGION.	SOU04770
301 CONTINUE	SOU04780
NDZ=14	SOU04790
N=NDZ+1	SOU04800

DX=(Z(1))/14.	SOU04810
P(15)=P(1)	SOU04820
P(1)=PEN	SOU04830
A(15)=A(1)	SOU04840
A(1)=AEN	SOU04850
Z(15)=Z(1)	SOU04860
Z(1)=0.0	SOU04870
X(15)=0.0	SOU04880
CALL STEAM(PEN)	SOU04890
X(1)=(HO-HF)/(HG-HF)	SOU04900
DO 39 IJ=2,14	SOU04910
Z(IJ)=DX*(IJ-1)	SOU04920
A(IJ)=AEN-ETA*Z(IJ)	SOU04930
39 CONTINUE	SOU04940
M=MC	SOU04950
IJ=1	SOU04960
40 CONTINUE	SOU04970
IJ=IJ+1	SOU04980
IF(IJ.GE.N)GOTO 919	SOU04990
DP(IJ)=M*M*0.5*VO*(1.+DELTA*F/ETA)*(1./A(IJ)**2.-1./A(IJ-1)**2.)	SOU05000
&+M*M*F/(2.*DELTA*ETA)*(2.*VO)*(1./A(IJ)-1./A(IJ-1))	SOU05010
P(IJ)=P(IJ-1)-DP(IJ)	SOU05020
CALL STEAM(P(IJ))	SOU05030
X(IJ)=(HO-HF)/(HG-HF)	SOU05040
IF(IJ.EQ.(N-1))GOTO 25	SOU05050
GOTO 40	SOU05060
25 CONTINUE	SOU05070
PRINT*,'**THE FOLLOWING ARE THE PRESSURE AND QUALITY DISTRIBUTION	SOU05080
&IN SINGLE-PHASE REGION**'	SOU05090
WRITE(6,478)	SOU05100
DO 48 IJ=1,N	SOU05110
48 WRITE(6,249)IJ,Z(IJ),A(IJ),P(IJ),X(IJ)	SOU05120
919 CONTINUE	SOU05130
C WRITE CPU TIME FOR THE CALCULATIONS	SOU05140
CALL DATETM (DATTIM,23,VCPU,CTIME,TCPU)	SOU05150
WRITE(6,801)DATTIM,VCPU,CTIME,TCPU	SOU05160
801 FORMAT('DATE/TIME:',A23/,'VCPU=',F10.3,'SEC.', 'CONNECT TIME=',F	SOU05170
&10.3,'SEC.', 'TCPU=',F10.3,'SEC.')	SOU05180
GOTO 999	SOU05190
C CONTINUE FOR SINGLE PHASE CALCULATIONS FROM STEP 92 TO 99	SOU05200
99 CONTINUE	SOU05210
C DIVIDE THE CHANNEL IN GRIDS	SOU05220
C3=0.01E-6	SOU05230
IF((L-ZFL).GT.C3)GOTO 988	SOU05240
NDZ=1	SOU05250
DZT=(L-ZFL)	SOU05260
GOTO 98	SOU05270
988 CONTINUE	SOU05280
IF(ZL.GE.0.01)GOTO 511	SOU05290
NDZ=15	SOU05300
GOTO 398	SOU05310
511 CONTINUE	SOU05320
NDZ=25	SOU05330
398 CONTINUE	SOU05340
DZT=(L-ZFL)/(2.*(FLOAT(NDZ))-1.)	SOU05350
IF(DZT.GT.C3)GOTO 98	SOU05360
NDZ=NDZ-1	SOU05370
GOTO 398	SOU05380
98 CONTINUE	SOU05390
C SUBROUTINE GRID DOES THE DIVISION OF GRIDS INTO NDZ GRIDS WITH A(I)	SOU05400

C	Z(I) DETERMINED	SOU05410
	N=NDZ+1	SOU05420
C	SET THE FLASH AS THE ZERO POINT	SOU05430
	Z(1)=ZFL	SOU05440
	P(1)=PFL	SOU05450
	A(1)=AFL	SOU05460
	V(1)=VFL	SOU05470
	X(1)=0.	SOU05480
	H(1)=HFL	SOU05490
134	CONTINUE	SOU05500
	DO 6 IK=2,N	SOU05510
	DZ(IK)=DZT*2.*(FLOAT(N-IK))	SOU05520
	Z(IK)=Z(IK-1)+DZ(IK)	SOU05530
	A(IK)=AEN-ETA*Z(IK)	SOU05540
6	CONTINUE	SOU05550
	A2=A(2)	SOU05560
C	GUESS DOES THE INITIAL GUESS FOR SPECIFIC VOLUME AT FIRST GRID POINT	SOU05570
	CALL GUESS(A2,0,Q)	SOU05580
	V(2)=0	SOU05590
	X(2)=Q	SOU05600
	GOTO 132	SOU05610
C	THE STEP COMING FROM THE TWO PHASE ENTRANCE IS THE FOLLOWING IE 77	SOU05620
	77 CONTINUE	SOU05630
C	DIVIDE THE CHANNEL IN TO GRID POINTS	SOU05640
	NDZ=15	SOU05650
75	CONTINUE	SOU05660
	DZT=L/(2.*(FLOAT(NDZ))-1.)	SOU05670
	IF(DZT.GT.C3)GOTO 76	SOU05680
	NDZ=NDZ-1	SOU05690
	GOTO 75	SOU05700
78	CONTINUE	SOU05710
C	SET THE ENTRANCE AS THE ZERO TH GRID	SOU05720
	Z(1)=0.	SOU05730
	A(1)=AEN	SOU05740
	V(1)=VEN	SOU05750
	P(1)=PEN	SOU05760
	X(1)=XEN	SOU05770
	V(2)=V(1)	SOU05780
	DO 7 II=2,N	SOU05790
	DZ(II)=DZT*2.*(FLOAT(N-II))	SOU05800
	Z(II)=Z(II-1)+DZ(II)	SOU05810
	A(II)=AEN-ETA*Z(II)	SOU05820
7	CONTINUE	SOU05830
132	CONTINUE	SOU05840
	N=NDZ+1	SOU05850
	I=1	SOU05860
C	NOW START CALCULATING THE PRESSURE DROP FOR SUCESSIVE GRID POINT	SOU05870
133	CONTINUE	SOU05880
	I=I+1	SOU05890
C	PRINT*, 'I =', I	SOU05900
C	STOP IF END POINT IS REACHED	SOU05910
	IF(I.GT.N)GOTO 999	SOU05920
	J=0	SOU05930
C	LOOP FOR THE CONVERGENCE OF PRESSURE AT NEXT GRID POINT	SOU05940
144	CONTINUE	SOU05950
	J=J+1	SOU05960
C	IF ITERATION GOES MORE THAN 35 TIMES THEN STOP	SOU05970
	IF(J.GT.35)GOTO 991	SOU05980
C	FIND PRESSURE DROP USING TWO-PHASE MOMENTUM EQUATION	SOU05990
	DP(I)=2.*M*M/(A(I)**2.+A(I-1)**2.)*(V(I)-V(I-1))+M*M*0.25*(V(I-1)+SOU06000	SOU06000

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      &V(I))* (1.+DELTA*F/ETA)*(1./A(I)**2.-1./A(I-1)**2.)+M*M*F/(2.*DELTA)SOU06010
      &*ETA)*(V(I-1)+V(I))*(1./A(I)-1./A(I-1))SOU06020
C   PRINT*, 'DP(I) =', DP(I)SOU06030
C   FIND THE PRESSURE AT THE I TH GRID POINTSOU06040
      P1(J)=P(I-1)-DP(I)SOU06050
      IF(P1(J).GT.1.013E5)GOTO 313SOU06060
C   REDUCE THE MASS FLOW RATESOU06070
      YH=M
      RR=(1.+YL/YH)/(2.)SOU06080
      R=RRSOU06090
      GOTO 12SOU06100
313 CONTINUESOU06110
      IF(J.LT.33)GOTO 143SOU06120
C   CHEK FOR CONVERGENCE IF ITERATION IS MORE THAN 30 TIMES
      IF((P1(J)-P1(J-1))/(P1(J-1)-P1(J-2)).LT.1.)GOTO 142SOU06130
      IF(I.GT.N)GOTO 991SOU06140
      IF(I.EQ.N)GOTO 155SOU06150
      YH=M
      RR=(1.+YL/YH)/(2.)SOU06160
      R=RRSOU06170
      GOTO 12SOU06180
142 CONTINUESOU06190
C   SEE FOR CONVERGENCE
      IF((P1(J)-P1(J-1))/P1(J).LE.0.001)GOTO 155SOU06200
143 CONTINUESOU06210
      IF(J-3) 333,333,000SOU06220
800 CONTINUESOU06230
      IF(P1(J).EQ.P1(J-1))GOTO 155SOU06240
      IF((P1(J)-P1(J-1))/(P1(J-1)-P1(J-2)).LT.1.)GOTO 300SOU06250
      CALL STEAM(P1(J))SOU06260
      VC=A(I)*A(I)*(HO-HF)/(M*M*(VG-VF)**2.)SOU06270
      VA=VF/(VG-VF)SOU06280
      VB=A(I)*A(I)*(HG-HF)/(M*M*(VG-VF)**2.)SOU06290
      X(I)=- (VA+VB)+((VA+VB)**2.+2.*VC-VA**2.)**(.5)SOU06300
      V(I)=VF+X(I)*(VG-VF)SOU06310
      U=M*V(I)/A(I)SOU06320
      DXDP=- (DSFDP*(1.-X(I))+X(I)*DSGDP)/(SG-SF)SOU06330
      UA=V(I)*((-((1.-X(I))*DVFDP+X(I)*DVGDP+(VG-VF)*DXDP))**(-.5))SOU06340
      XMACH=U/UA
      IF((U-UA)/UA.LE.0.001)GOTO 144SOU06350
C   REDUCE THE MASS FLOW RATE.SOU06360
      YH=M
      IF(RJS.GE.1)GOTO 157SOU06370
      IF(I.LT.N)GOTO 528SOU06380
      IF((YH-YL)/YH.GT.0.001)GOTO 528SOU06390
      RJS=RJS+1.
      IF(RJS.GE.1.)GOTO 530SOU06400
528 CONTINUESOU06410
      JJN=JK
      RR=(1.+YL/YH)/(2.)SOU06420
      R=RRSOU06430
      GOTO 12SOU06440
530 CONTINUESOU06450
      M=YH
      DP(I)=2.*M*M/(A(I)**2.+A(I-1)**2.)*(V(I)-V(I-1))+M*M*0.25*(V(I-1)
      &+V(I))*(1.+DELTA*F/ETA)*(1./A(I)**2.-1./A(I-1)**2.)+M*M*F/(2.*
      &DELTA*ETA)*(V(I-1)+V(I))*(1./A(I)-1./A(I-1))SOU06460
      P1(J)=P(I-1)-DP(I)SOU06470
      GOTO 155SOU06480
300 CONTINUESOU06490

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IF(ABS((P1(J)-P1(J-1))/P1(J)).LE.0.001)GOTO 155	SOU06610
333 CONTINUE	SOU06620
IF(P1(J).EQ.P1(J-1))GOTO 155	SOU06630
C FIND THE VALUEES OF PROPERTIES AT NEW VALUE OF P1(J)	SOU06640
CALL STEAM(P1(J))	SOU06650
C FIND QUALITY USING ENERGY EQUATION	SOU06660
VC=A(I)*A(I)*(HO-HF)/(M*M*(VG-VF)**2.)	SOU06670
VA=VF/(VG-VF)	SOU06680
VB=A(I)*A(I)*(HG-HF)/(M*M*(VG-VF)**2.)	SOU06690
X(I)=- (VA+VB) + ((VA+VB)**2.+2.*VC-VA**2.)**(.5)	SOU06700
V(I)=VF+X(I)*(VG-VF)	SOU06710
H(I)=HF+X(I)*(HG-HF)	SOU06720
C FIND FLUID VELOCITY AND SOUND VELOCITY AT DRID POINTS	SOU06730
U=M*V(I)/A(I)	SOU06740
DXDP=- (DSFDP*(1.-X(I))+X(I)*DSGDP)/(SG-SF)	SOU06750
UA=V(I)*((-((1.-X(I))*DVFDP+X(I)*DVGDP+(VG-VF)*DXDP))**(-.5))	SOU06760
XMACH=U/UA	SOU06770
IF((U-UA)/UA.LE.0.001)GOTO 144	SOU06780
C REDUCE THE MASS FLOW RATE.	SOU06790
YH=M	SOU06800
IF(RJS.GE.1.)GOTO 157	SOU06810
IF(I.LT.N)GOTO 529	SOU06820
IF((YH-YL)/YH.GT.0.001)GOTO 529	SOU06830
RJS=RJS+1.	SOU06840
IF(RJS.GE.1.)GOTO 530	SOU06850
529 CONTINUE	SOU06860
JJN=JK	SOU06870
RR=(1.+YL/YH)/(2.)	SOU06880
R=RR	SOU06890
GOTO 12	SOU06900
155 CONTINUE	SOU06910
P(I)=P1(J)	SOU06920
IF(P(I).GT.P(I-1))GOTO 899	SOU06930
C WRITE THE CONVERGED VALUES AT THIS GRID POINT.	SOU06940
C FIND THE FLUID PROPERTIES	SOU06950
CALL STEAM(P(I))	SOU06960
C FIND THE QUALITY	SOU06970
VC=A(I)*A(I)*(HO-HF)/(M*M*(VG-VF)**2.)	SOU06980
VA=VF/(VG-VF)	SOU06990
VB=A(I)*A(I)*(HG-HF)/(M*M*(VG-VF)**2.)	SOU07000
X(I)=- (VA+VB) + ((VA+VB)**2.+2.*VC-VA**2.)**0.5	SOU07010
V(I)=VF+X(I)*(VG-VF)	SOU07020
H(I)=HF+X(I)*(HG-HF)	SOU07030
C WRITE THE FLUID STATE AT THE GRID I	SOU07040
IF(RJS.GE.1.)GOTO 157	SOU07050
C FIND SOUND VELOCITY AND FLUID VELOCITY	SOU07060
U=M*V(I)/A(I)	SOU07070
DXDP=- (DSFDP*(1.-X(I))+X(I)*DSGDP)/(SG-SF)	SOU07080
UA=V(I)*((-((1.-X(I))*DVFDP+X(I)*DVGDP+(VG-VF)*DXDP))**(-0.5))	SOU07090
XMACH=U/UA	SOU07100
C COMPARE THE SOUND AND FLUID VELOCITIES	SOU07110
IF(UA.GE.U)GOTO 161	SOU07120
IF((U-UA)/UA.LE.0.001)GOTO 162	SOU07130
IF(RJS.GE.1.)GOTO 157	SOU07140
IF(I.LT.N)GOTO 163	SOU07150
RJS=RJS+1.	SOU07160
IF(RJS.GE.1.)GOTO 530	SOU07170
163 CONTINUE	SOU07180
C DECREASE MASS FLOW RATE	SOU07190
YH=M	SOU07200

RR=(1.+YL/YH)/(2.)	SOU07210
R=RR	SOU07220
GOTO 12	SOU07230
899 CONTINUE	SOU07240
NDZ=NDZ-1	SOU07250
N=NDZ+1	SOU07260
DZT=(L-ZFL)/(2.** (FLOAT(NDZ))-1.)	SOU07270
GOTO 134	SOU07280
161 CONTINUE	SOU07290
IF((UA-U)/UA.LE.0.001)GOTO 162	SOU07300
C CHECK FOR THE LENGTH	SOU07310
IF(I.EQ.N)GOTO 98	SOU07320
GOTO 133	SOU07330
C CONTINUE FOR STEP 162	SOU07340
162 CONTINUE	SOU07350
C FIRST CHECK FOR LENGTH WHEATHER AT EXIT	SOU07360
IF(I.EQ.N)GOTO 157	SOU07370
GOTO 163	SOU07380
991 CONTINUE	SOU07390
WRITE(8,119)P1(J),P1(J-1),P1(J-2),P1(J-3),P1(J-4)	SOU07400
119 FORMAT(//T10,'DOES NOT CONVERGE FOR THE GIVEN MASS FLOW RATE',//,	SOU07410
&5(2X,E15.7))	SOU07420
993 CONTINUE	SOU07430
PRINT*, 'THE CALCULATION IS STOPPED DUE TO MANY ITERATIONS'	SOU07440
999 CONTINUE	SOU07450
GOTO 710	SOU07460
720 CONTINUE	SOU07470
STOP	SOU07480
END	SOU07490
C	SOU07500
C.....	SOU07510
C	SOU07520
C-----	SOU07530
C	SOU07540
C THE FOLLOWING ARE SUBROUTINES	SOU07550
C-----	SOU07560
SUBROUTINE GUESS(AA,C,D)	SOU07570
C THIS SUBROUTINE DETERMINES FIRST GUESS VALUE FOR SP VOLUME AT FIRST	SOU07580
C GRID POINT AFTER THE FLASHING POINT	SOU07590
IMPLICIT REAL*8(A,L,M)	SOU07600
COMMON/BLOCK1/VF,DVFD, VG,DVGDP, HF,DHFD, HG,DHGD, SF,DSFSP,SG,	SOU07610
&DSGDP	SOU07620
COMMON/BLOCK2/M,VFL,ETA,DELTA,F,AFL,ZFL,NDZ,DZT,PFL,HFL	SOU07630
C FIND SINGLE PHASE PRESSURE DROP BETWEEN FLASH AND FIRST GRID	SOU07640
DPSP1=M*M*0.5*VFL*(1.+DELTA*F/ETA)*(1./AA**2.-1./AFL**2.)+M*M*F	SOU07650
&*VFL/ETA/DELTA*(1./AA-1./AFL)	SOU07660
C FIND PRESSURE AT FIRST GRID	SOU07670
PS1=PFL-DPSP1	SOU07680
C FIND PROPERTIES OF FLUID AT PS1	SOU07690
IF(PS1.LT.100.)GOTO 10	SOU07700
CALL STEAM(PS1)	SOU07710
C FIND QUALITY AT FIRST GRID	SOU07720
X1=(HFL-HF)/(HG-HF)	SOU07730
C FIND MARTINELLI-NELSON FRICTION MULTIPLIER	SOU07740
C	SOU07750
PS1=PS1*1.46041E-4	SOU07760
C	SOU07770
PHI=EXP((1.85E-10*PS1**4.-8.9E-7*PS1**3.+1.79E-3*PS1**2.-1.63*PS1+	SOU07780
&6.02E2)*X1)	SOU07790
C FIND TWO-PHASE PRESSURE DROP	SOU07800

P2P=PHI*DPSP1	SOU07810
PP2=PFL-P2P	SOU07820
C FIND FLUID PROPERTIES WITH THIS VALUE OF PP2	SOU07830
CALL STEAM(PP2)	SOU07840
D=(HFL-HF)/(HG-HF)	SOU07850
C=VF*D*(VG-VF)	SOU07860
C	SOU07870
10 CONTINUE	SOU07880
RETURN	SOU07890
END	SOU07900
C	SOU07910
C-----	SOU07920
SUBROUTINE STEAM(P)	SOU07930
COMMON/BLOCK1/VF,DVFD, VG,DVGD, HF,DHFD, HG,DHGD, SF,DSFD, SG,	SOU07940
&DSGSP	SOU07950
C	SOU07960
C	SOU07970
C THIS IS A SUBROUTINE TO CALCULATE THE THERMODYNAMIC PROPERTIES OF	SOU07980
C SATURATED WATER AND STEAM AND THE DERIVATIVES OF SPECIFIC VOLUME V,	SOU07990
C ENTROPY,S,ENTHALPY H,ETC AS REQUIRED.	SOU08000
C THE UNITS WE ARE USING ARE FOR P KG/CM**2.,FOR T DEG. K , FOR V,	SOU08010
C CM**3/GM,FOR ENTHALPY H KCAL/KG AND FOR S KCAL/KG K.	SOU08020
C THE INPUT AND OUTPUT FROM THIS SUBROUTINE GIVE THE CALCULATED FLUID	SOU08030
C PROPERTIES IN MKS UNITS	SOU08040
C-----	SOU08050
C	SOU08060
C DEFINE THE CRITICAL PRESSURE,TEMPERATURE,SPECIFIC VOLUME,ENTHALPY	SOU08070
C AND THE ENTROPY WITH THE UNITS ALREADY SAID ABOVE	SOU08080
C	SOU08090
PC=2.2119258D7	SOU08100
TC=847.27	SOU08110
VC=3.1698D-3	SOU08120
HC=2.107209D6	SOU08130
SC=4.4427222D3	SOU08140
	SOU08150
Z=-ALOG(P/PC)	SOU08160
C CONSTANTS FOR SATURATION TEMPERATURE	SOU08170
A1=0.1295426849D0	SOU08180
A2=0.4513602436D0	SOU08190
A3=0.2122378436D0	SOU08200
A4=-0.3581437096D-2	SOU08210
A5=0.9090316283D-4	SOU08220
A6=0.3338530622D1	SOU08230
A7=0.1426558985D1	SOU08240
C EQUATION FOR SATURATION TEMPERATURE	SOU08250
A8=1.+A6*Z+A7*Z*Z	SOU08260
A9=A1*Z+A2*Z*Z+A3*(Z**3.)+A4*(Z**4.)+A5*(Z**5.)	SOU08270
TA=TC*(A8/(A8+A9))	SOU08280
THETA=TA/TC	SOU08290
C DERIVATIVE DT/DP	SOU08300
A10=A8+2.*A7*Z	SOU08310
A11=A10+A1+2.*A2*Z+3.*A3*Z*Z+4.*A4*(Z**3.)+5.*A5*(Z**4.)	SOU08320
A12=A10/A8	SOU08330
A13=A11/(A8+A9)	SOU08340
DTDP=-TA*(A12-A13)/P	SOU08350
C CONSTANTS FOR SPECIFIC VOLUME OF WATER VF	SOU08360
B1=0.1030837042D5	SOU08370
B2=0.5538426569D7	SOU08380
B3=0.1799772193D9	SOU08390
B4=-0.4743843407D9	SOU08400

	B5=0.2019749779D10	SOU08410
	B6=-0.6258564009D10	SOU08420
	B7=0.1230937510D11	SOU08430
	B8=-0.1351160480D11	SOU08440
	B9=0.6353769714D10	SOU08450
	B10=0.1163268984D5	SOU08460
	B11=.7686775534D7	SOU08470
	B12=0.3694089380D9	SOU08480
C		SOU08490
	Y=1.-THETA	SOU08500
C	EQUATION FOR SPECIFIC VOLUME OF WATER VF	SOU08510
	B13=1.+B1*Y+B2*Y*Y+B3*(Y**3.)+B4*(Y**4.)	SOU08520
	B14=B5*(Y**5.)+B6*(Y**6.)+B7*(Y**7.)+B8*(Y**8.)+B9*(Y**9.)	SOU08530
	B15=B13+B14	SOU08540
	B16=1.+B10*Y+B11*Y*Y+B12*Y**3.	SOU08550
C		SOU08560
	VF=VC*B15/B16	SOU08570
C	DERIVATIVE DVF/DP : DVFDP	SOU08580
	B17=B1+2.*B2*Y+3.*B3*Y*Y+4.*B4*Y**3.	SOU08590
	B18=5.*B5*Y**4.+6.*B6*Y**5.+7.*B7*Y**6.+8.*B8*Y**7.+9.*B9*Y**8.	SOU08600
	B19=B17+B18	SOU08610
	B20=B19/B15	SOU08620
	B21=B10+2.*B11*Y+3.*B12*Y*Y	SOU08630
	B22=B21/B16	SOU08640
C		SOU08650
	DVFDP=-VF*DTDP*(B20-B22)/TC	SOU08660
C		SOU08670
C		SOU08680
C	CONSTANTS FOR THE SPECIFIC VOLUME OF STEAM VG	SOU08690
	C1=0.4116876614D4	SOU08700
	C2=0.8624691596D6	SOU08710
	C3=0.4332657100D7	SOU08720
	C4=-0.2835427189D8	SOU08730
	C5=0.3658875474D8	SOU08740
	C6=0.1299835973D7	SOU08750
	C7=-0.1729498423D8	SOU08760
	C8=0.4780391119D4	SOU08770
	C9=0.1436266862D7	SOU08780
	C10=0.2732195974D8	SOU08790
	C11=-0.8611392578D8	SOU08800
	C12=0.8981196903D8	SOU08810
C	EQUATION FOR THE SPECIFIC VOLUME OF STEAM VG	SOU08820
	C13=1.+C8*Y+C9*Y*Y+C10*Y**3.+C11*Y**4.+C12*Y**5.	SOU08830
	C14=1.+C1*Y+C2*Y*Y+C3*Y**3.+C4*Y**4.	SOU08840
	C15=C5*Y**5.+C6*Y**6.+C7*Y**7.	SOU08850
	C16=C14+C15	SOU08860
	C17=C16*THETA**6.	SOU08870
C		SOU08880
	VG=VC*C13/C17	SOU08890
C		SOU08900
C	DERIVATIVE DVGDP	SOU08910
	C18=C8+2.*C9*Y+3.*C10*Y*Y+4.*C11*Y**3.+5.*C12*Y**4.	SOU08920
	C19=C18/C13	SOU08930
	C20=C1+2.*C2*Y+3.*C3*Y*Y+4.*C4*Y**3.	SOU08940
	C21=5.*C5*Y**4.+6.*C6*Y**5.+7.*C7*Y**6.	SOU08950
	C22=C20+C21	SOU08960
	C23=C22/C16	SOU08970
C		SOU08980
	DVGDP=VG*DTDP*(-6./THETA-C19-C23)/TC	SOU08990
C		SOU09000

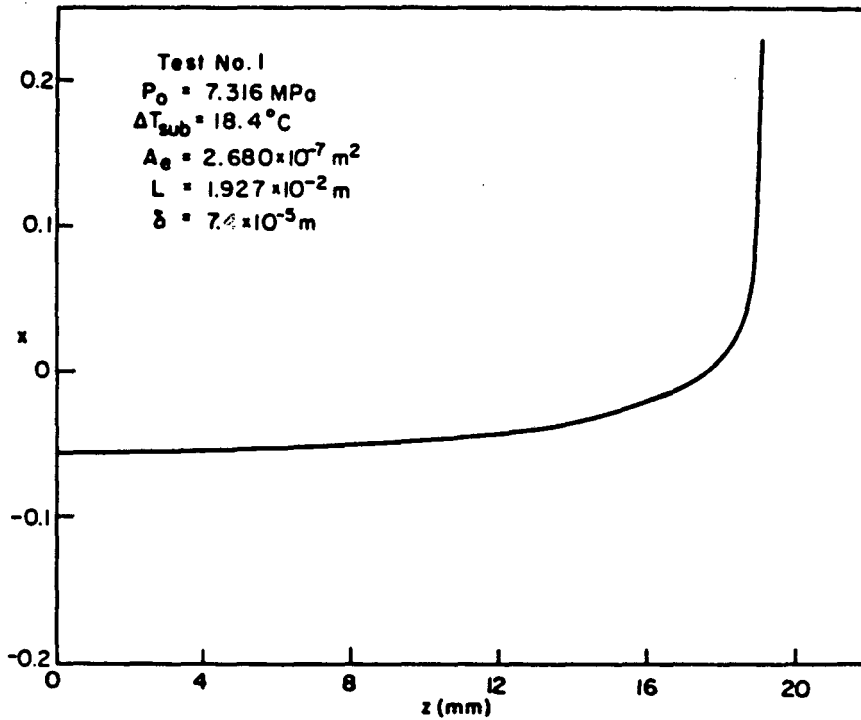
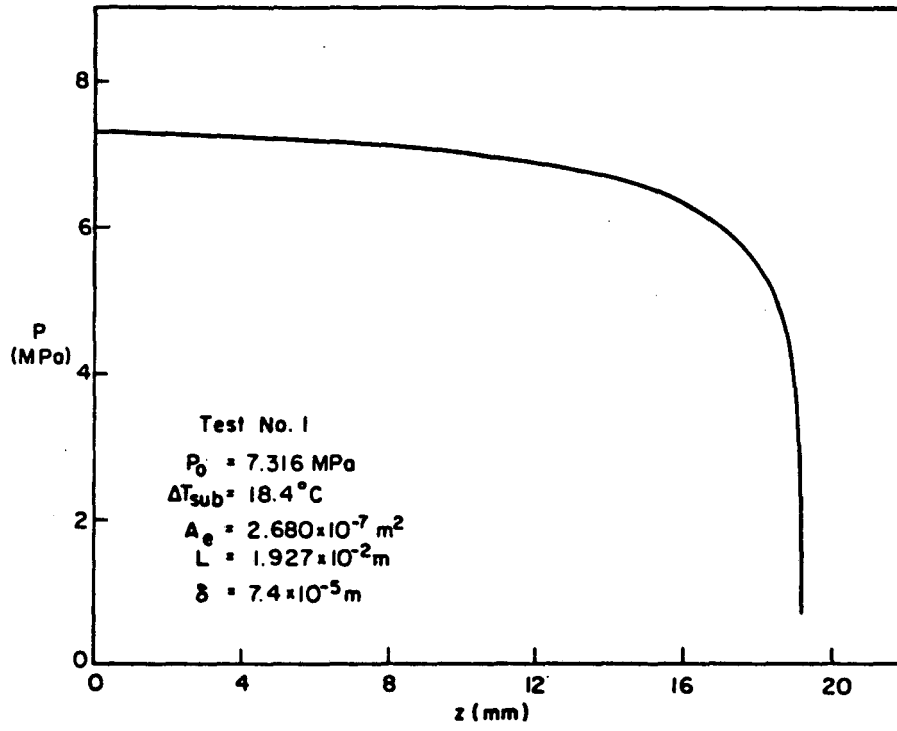
C		SOU09010
C	CONSTANTS FOR SPECIFIC ENTHALPY OF WATER HF	SOU09020
	D1=0.503588841804	SOU09030
	D2=0.852988031006	SOU09040
	D3=0.397566129407	SOU09050
	D4=-0.132121217408	SOU09060
	D5=0.747256735807	SOU09070
	D6=-0.169003475307	SOU09080
	D7=0.528873337704	SOU09090
	D8=0.982421289106	SOU09100
	D9=0.782590628807	SOU09110
	D10=-0.462022817707	SOU09120
C		SOU09130
C	EQUATION FOR SPECIFIC ENTHALPY OF WATER HF	SOU09140
	D11=1.+D1*Y+D2*Y*Y+D3*Y**3.+D4*Y**4.+D5*Y**5.+D6*Y**6.	SOU09150
	D12=1.+D7*Y+D8*Y*Y+D9*Y**3.+D10*Y**4.	SOU09160
	HF=HC*D11/D12	SOU09170
C	DERIVATIVE DHF/DP ; DHFDP	SOU09180
	D13=D1+2.*D2*Y+3.*D3*Y*Y+4.*D4*Y**3.+5.*D5*Y**4.+6.*D6*Y**5.	SOU09190
	D14=D7+2.*D8*Y+3.*D9*Y*Y+4.*D10*Y**3.	SOU09200
	D15=D13/D11	SOU09210
	D16=D14/D12	SOU09220
	DHFDP=-HF*(D15-D16)*DTDP/TC	SOU09230
C	CONSTANTS FOR SPECIFIC ENTHALPY OF STEAM HG	SOU09240
	E1=0.498130313704	SOU09250
	E2=0.155796861107	SOU09260
	E3=0.321871068908	SOU09270
	E4=0.958496511407	SOU09280
	E5=-0.439895952308	SOU09290
	E6=0.272498390808	SOU09300
	E7=0.480296564104	SOU09310
	E8=0.138587610407	SOU09320
	E9=0.232356620508	SOU09330
	E10=0.813454111006	SOU09340
C	EQUATION FOR SPECIFIC ENTHALPY OF STEAM HG	SOU09350
	E11=1.+E1*Y+E2*Y*Y+E3*Y**3.+E4*Y**4.+E5*Y**5.+E6*Y**6.	SOU09360
	E12=1.+E7*Y+E8*Y*Y+E9*Y**3.+E10*Y**4.	SOU09370
	HG=HC*E11/E12	SOU09380
C	DERIVATIVE DHGDP	SOU09390
	E13=E1+2.*E2*Y+3.*E3*Y*Y+4.*E4*Y**3.+5.*E5*Y**4.+6.*E6*Y**5.	SOU09400
	E14=E7+2.*E8*Y+3.*E9*Y*Y+4.*E10*Y**3.	SOU09410
	E15=E13/E11	SOU09420
	E16=E14/E12	SOU09430
	DHGDP=-HG*(E15-E16)*DTDP/TC	SOU09440
C	CONSTANTS FOR SPECIFIC ENTROPY OF WATER SF	SOU09450
	F1=0.500361430504	SOU09460
	F2=0.827385789806	SOU09470
	F3=0.389385240707	SOU09480
	F4=-0.115968636408	SOU09490
	F5=0.336964036507	SOU09500
	F6=0.122650602107	SOU09510
	F7=0.518729581304	SOU09520
	F8=0.916803401906	SOU09530
	F9=0.631959188407	SOU09540
	F10=-0.67626451707	SOU09550
C		SOU09560
C	EQUATION FOR THE SPECIFIC ENTROPY OF WATER SF	SOU09570
	F11=1.+F1*Y+F2*Y*Y+F3*Y**3.+F4*Y**4.+F5*Y**5.+F6*Y**6.	SOU09580
	F12=1.+F7*Y+F8*Y*Y+F9*Y**3.+F10*Y**4.	SOU09590
	SF=SC*F11/F12	SOU09600

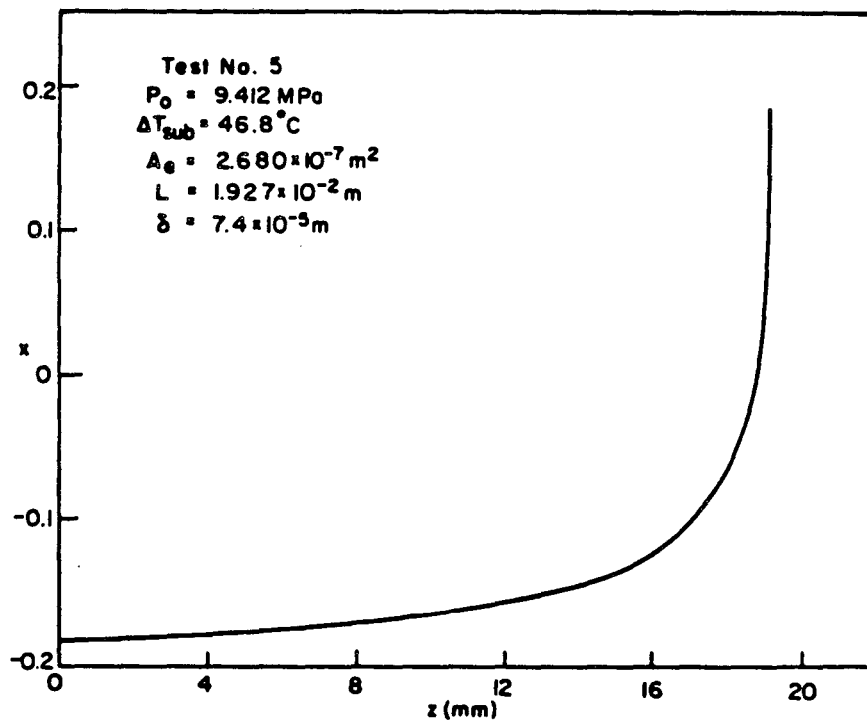
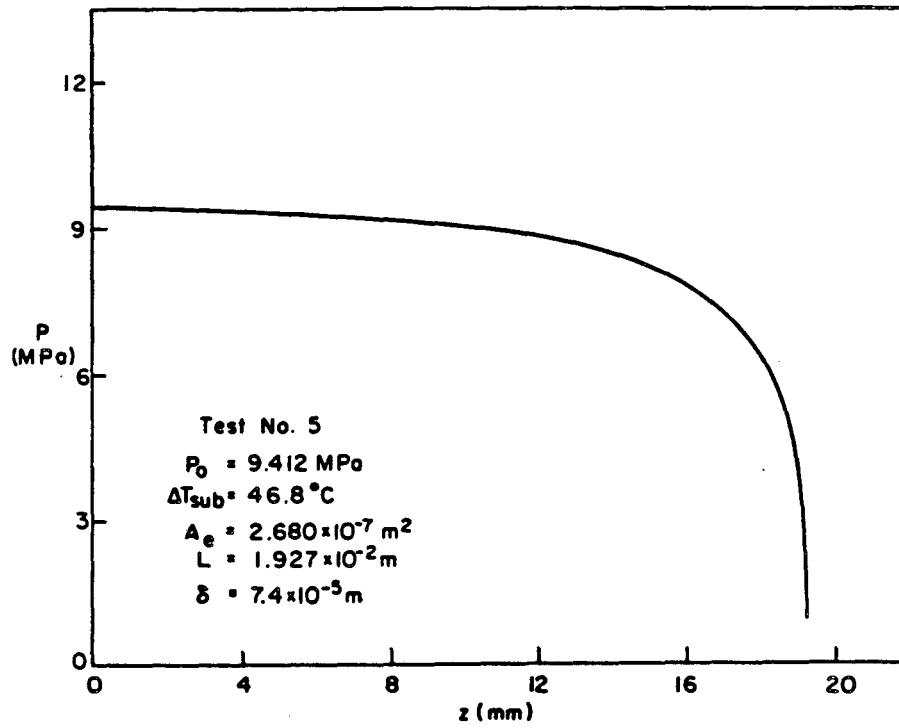
C DERIVATIVE DSF/DP ; DSFDP	SOU09610
F13=F1+2.*F2*Y+3.*F3*Y*Y+4.*F4*Y**3.+5.*F5*Y**4.+6.*F6*Y**5.	SOU09620
F14=F7+2.*F8*Y+3.*F9*Y*Y+4.*F10*Y**3.	SOU09630
F15=F13/F11	SOU09640
F16=F14/F12	SOU09650
DSFDP=-SF*(F15-F16)*DTDP/TC	SOU09660
C CONSTANTS FOR SPECIFIC ENTROPY OF STEAM SG	SOU09670
G1=0.5922502679D4	SOU09680
G2=0.2121506453D7	SOU09690
G3=0.4367159918D8	SOU09700
G4=-0.1755438026D8	SOU09710
G5=-0.3706247525D8	SOU09720
G6=0.4303459784D8	SOU09730
G7=0.5782905801D4	SOU09740
G8=0.1950840259D7	SOU09750
G9=0.3351957025D8	SOU09760
G10=-0.3604330231D8	SOU09770
C EQUATION FOR THE SPECIFIC ENTROPY OF STEAM SG	SOU09780
G11=1.+G1*Y+G2*Y*Y+G3*Y**3.+G4*Y**4.+G5*Y**5.+G6*Y**6.	SOU09790
G12=1.+G7*Y+G8*Y*Y+G9*Y**3.+G10*Y**4.	SOU09800
SG=SC*G11/G12	SOU09810
C DERIVATIVE DSG/DP ; DSGDP	SOU09820
G13=G1+2.*G2*Y+3.*G3*Y*Y+4.*G4*Y**3.+5.*G5*Y**4.+6.*G6*Y**5.	SOU09830
G14=G7+2.*G8*Y+3.*G9*Y*Y+4.*G10*Y**3.	SOU09840
G15=G13/G11	SOU09850
G16=G14/G12	SOU09860
DSGDP=-SG*(G15-G16)*DTDP/TC	SOU09870
RETURN	SOU09880
END	SOU09890

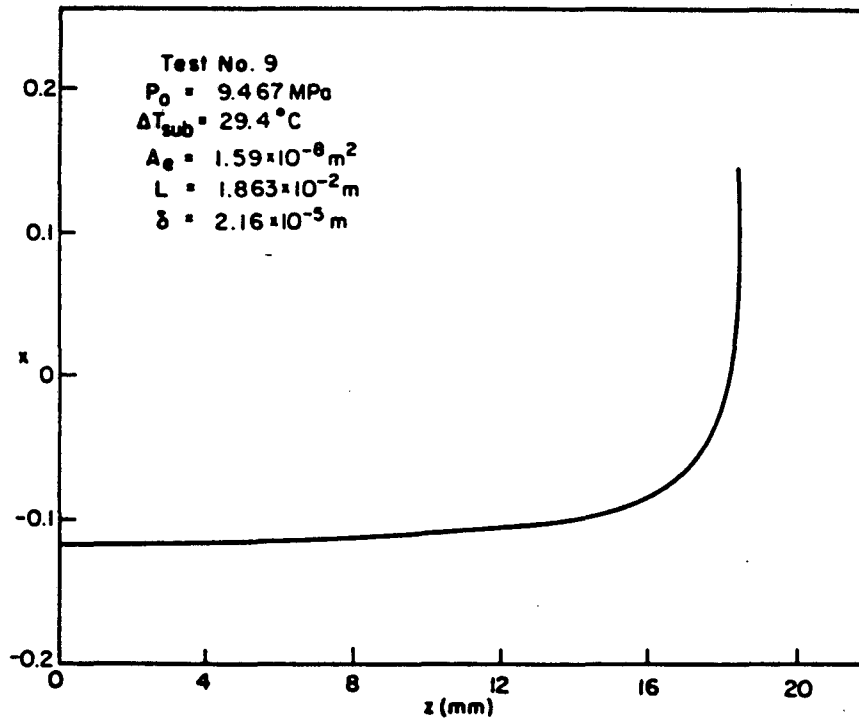
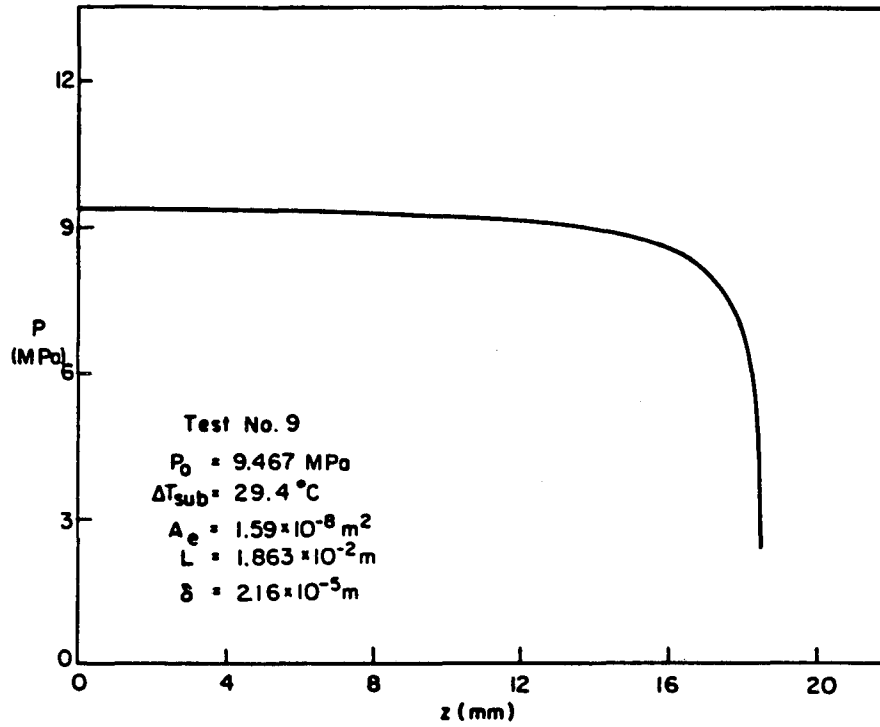
APPENDIX D

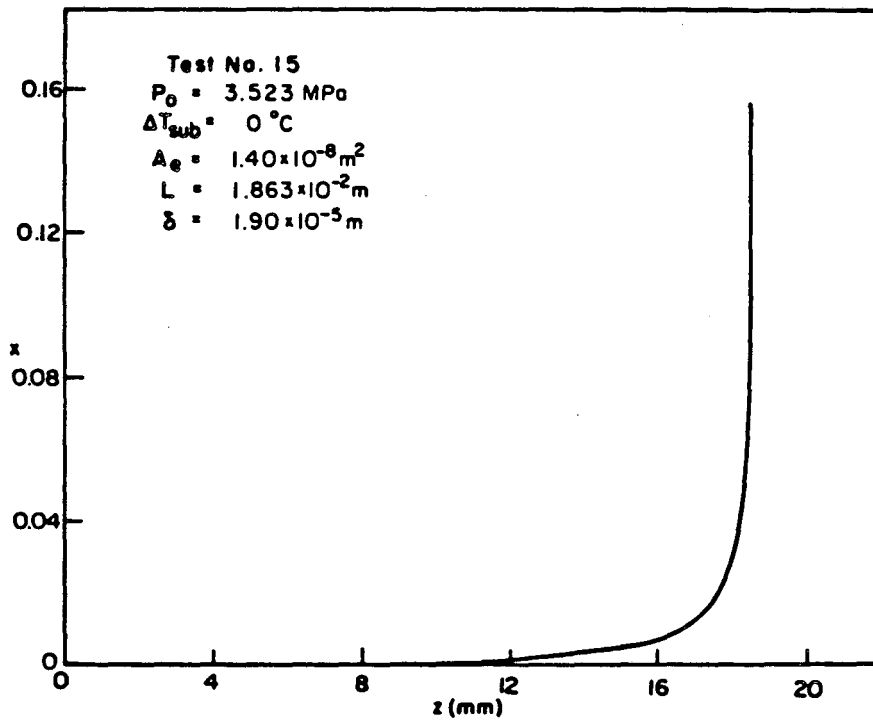
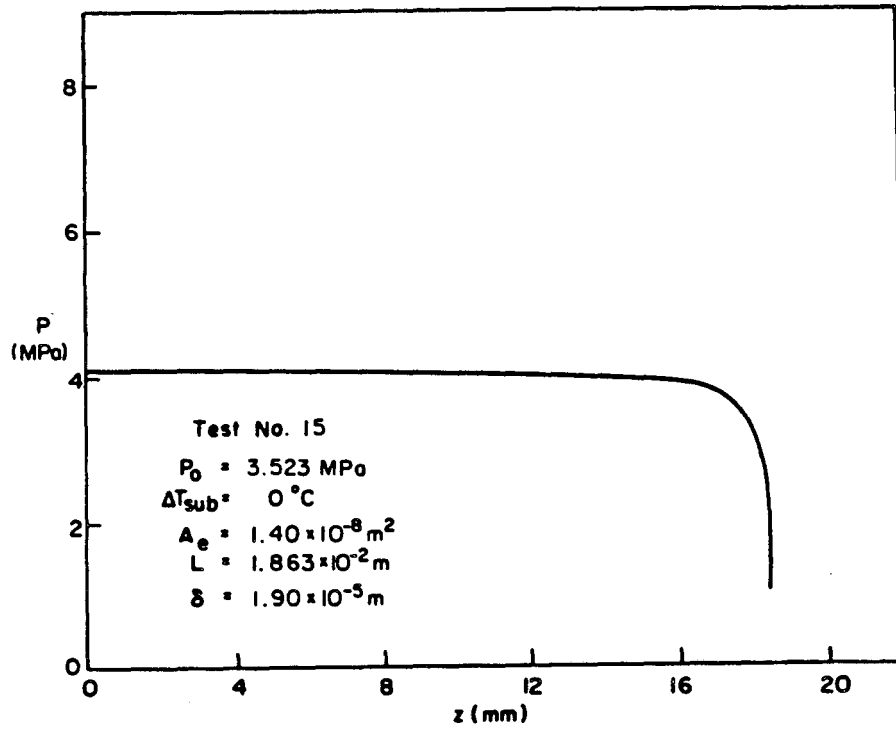
Pressure and Quality Profiles

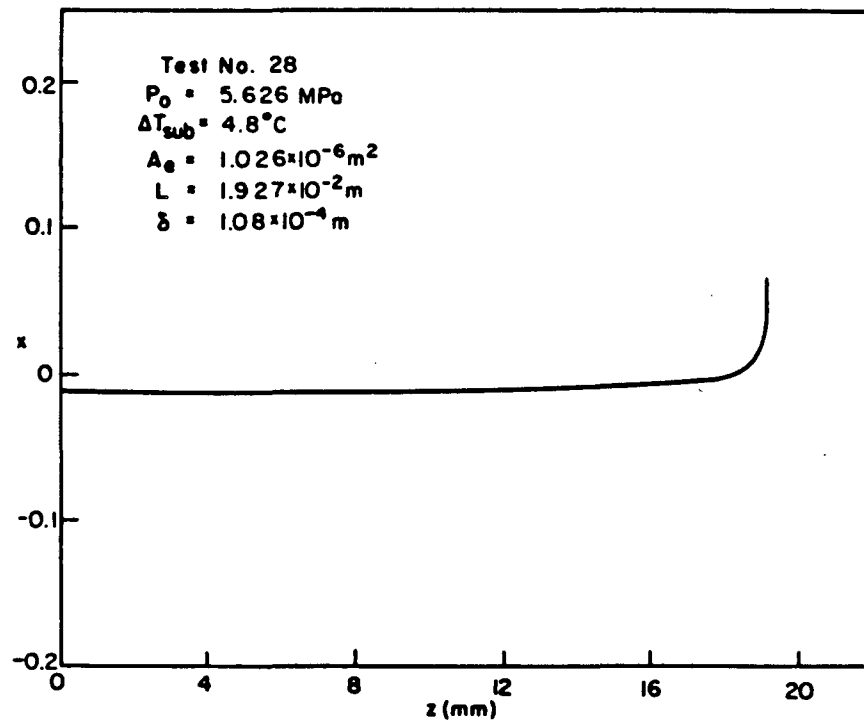
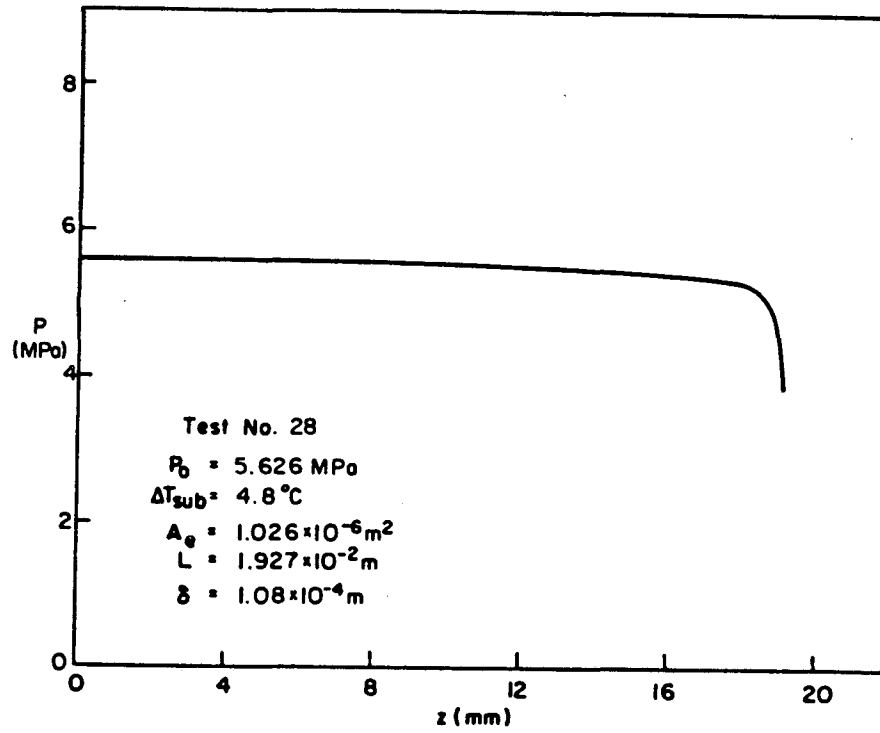
In this appendix the pressure and quality profiles obtained with SOURCE for some BCL tests are given. The choice of test numbers is based on highest and lowest subcoolings for each of the five BCL crack test sections. Hence total of 10 test numbers are chosen for presentation. The test parameters associated with each test are illustrated in each figure.

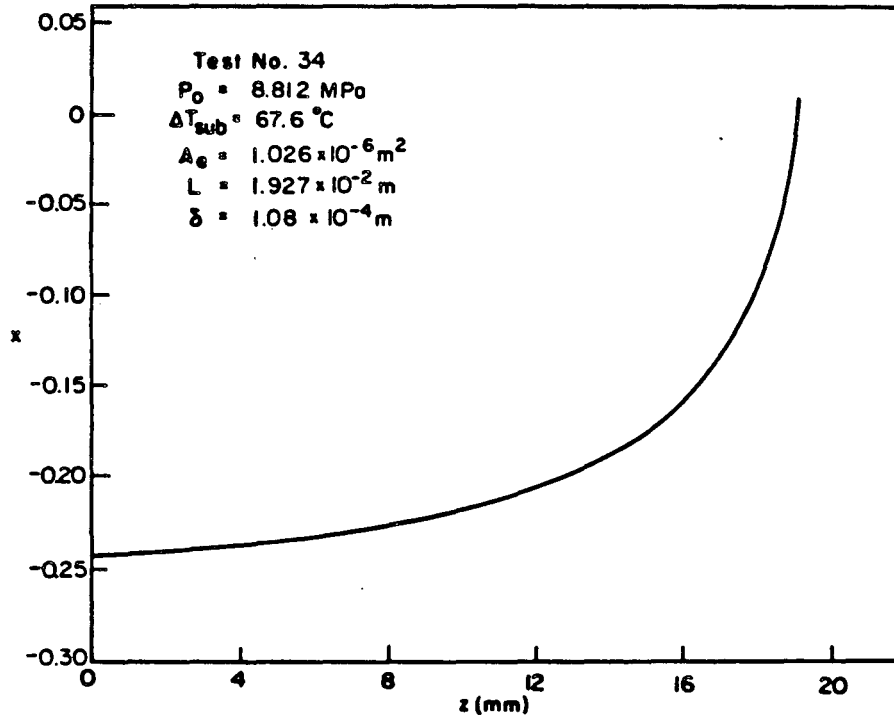
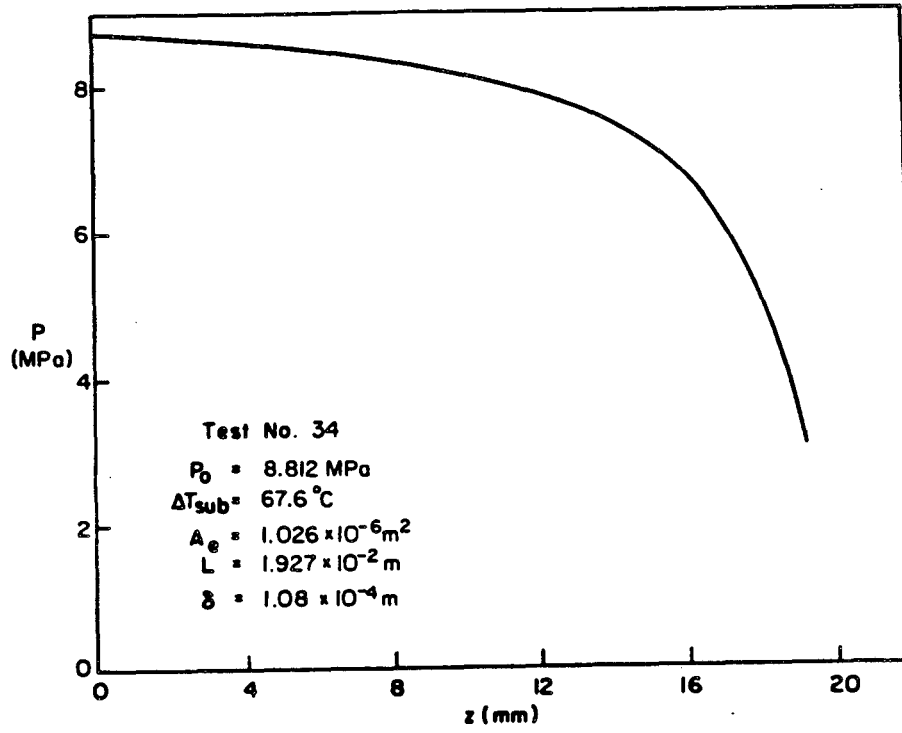


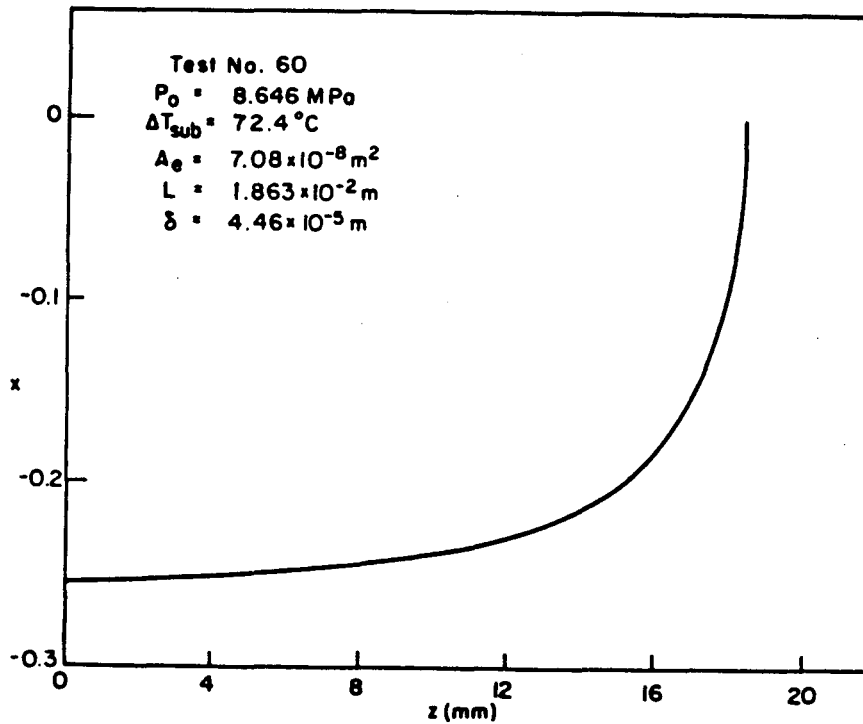
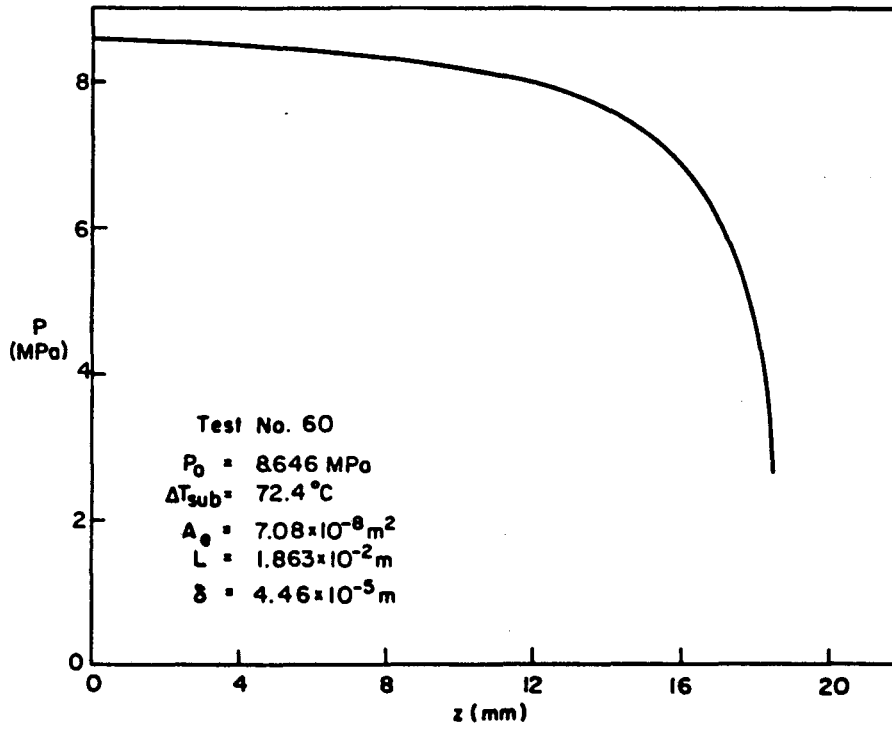


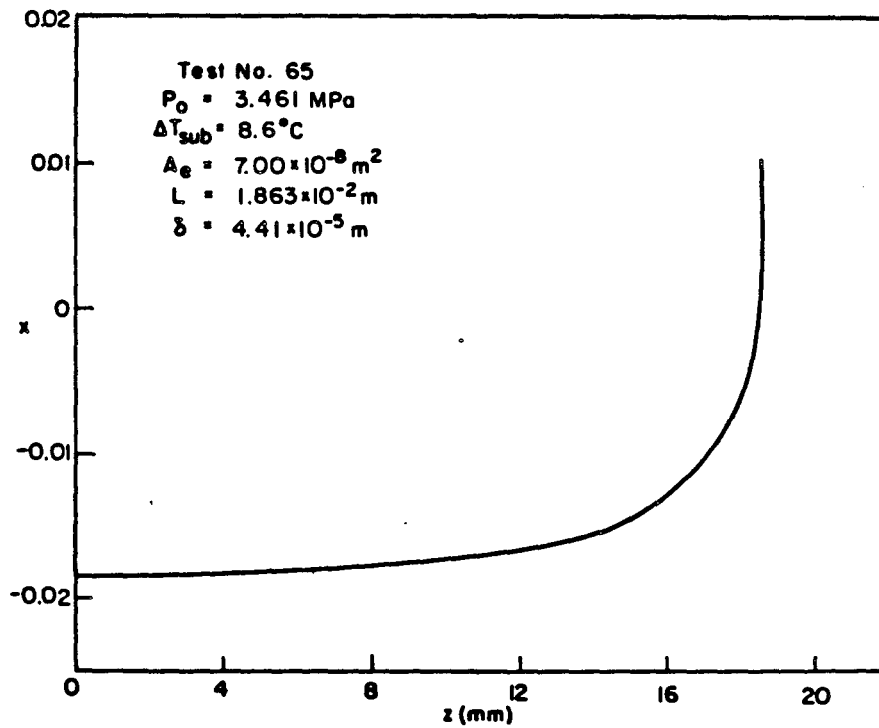
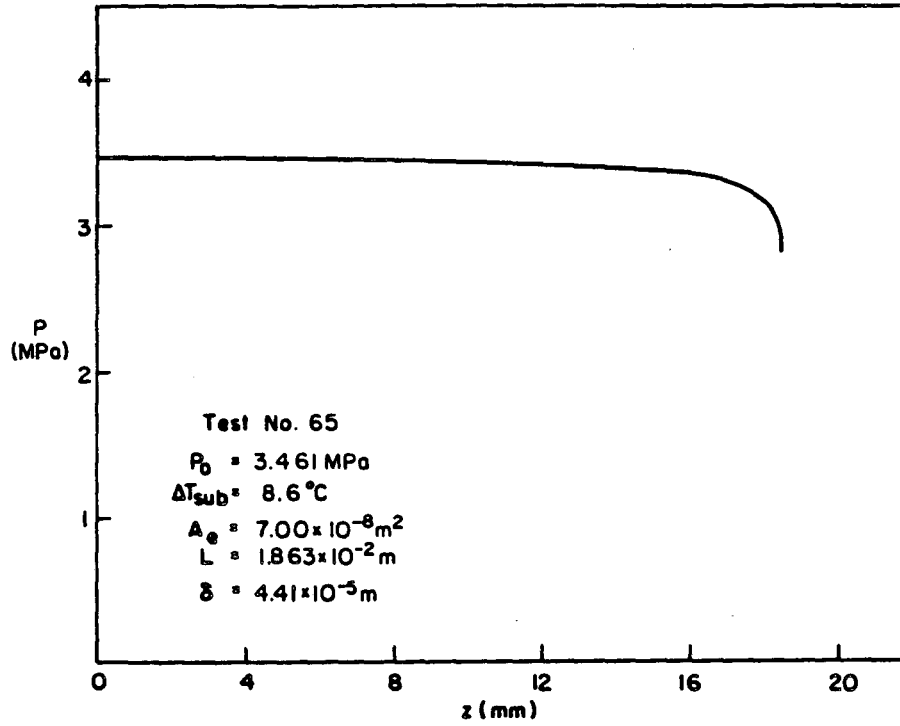


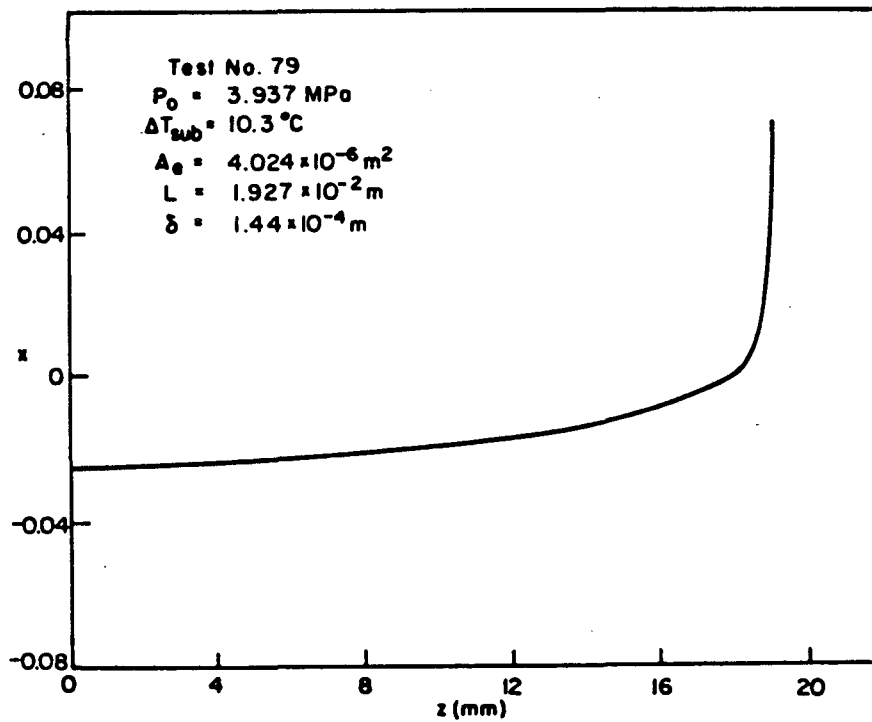
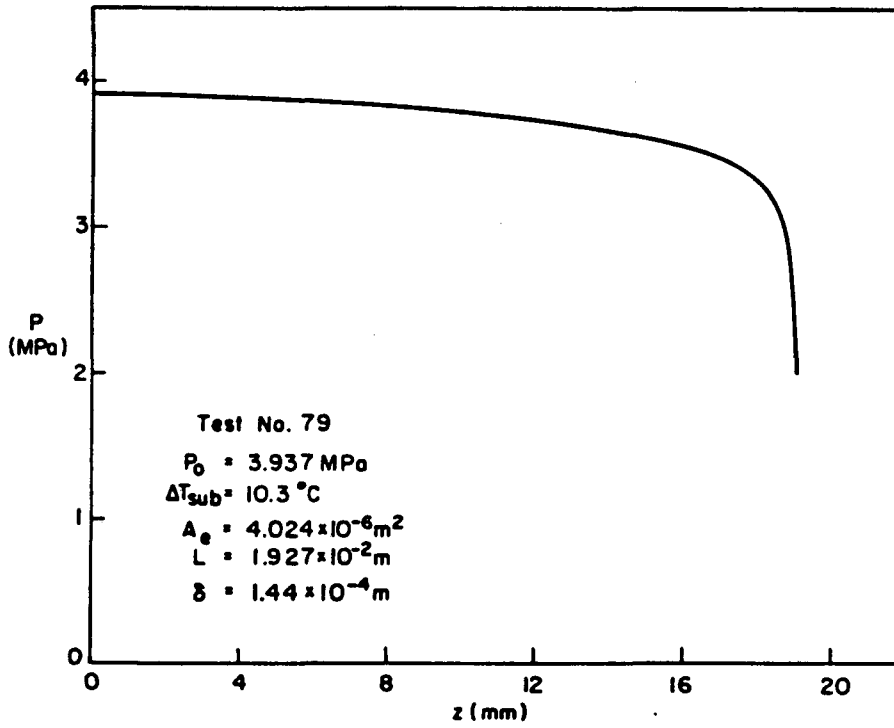


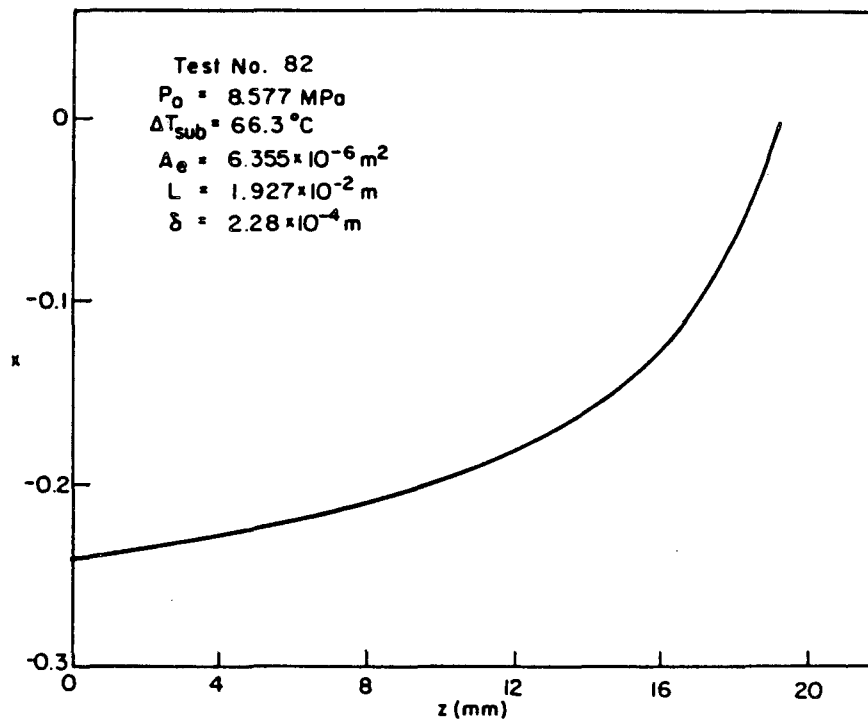
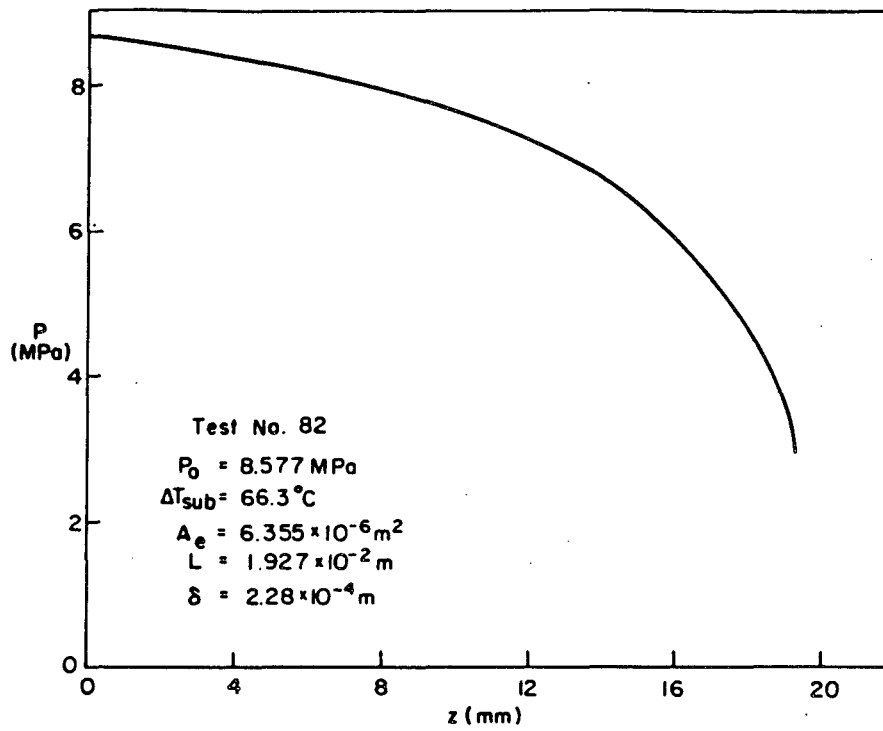












APPENDIX E

Thermodynamic Properties of Saturated Water and Steam

The code SOURCE uses the subroutine STEAM to calculate the thermodynamic properties of saturated water and steam. The input parameter to STEAM is pressure. For given value of pressure, STEAM first calculates the corresponding saturation temperature, and then uses this value of saturation temperature to calculate the other parameters, h_f , h_g , v_f , v_g , s_f and s_g . STEAM also calculates the derivatives of the saturation properties, enthalpy, specific volume and the entropy, with respect to pressure. For this calculation, the derivatives of each property with respect to the pressure are expressed in the following form, for example, for specific volume of water,

$$\begin{aligned} dv_f/dp &= v_c \cdot \frac{d}{dp} (v_f/v_c) \\ &= v_c \cdot \frac{d}{dT} (v_f/v_c) \cdot \frac{dT}{dp} \end{aligned}$$

The listing of the subroutine STEAM is included in the main program listing SOURCE. When subroutine STEAM is called upon in SOURCE with a pressure value, it returns the saturated values of specific volume, enthalpy and entropy and their derivation with respect to pressure for the input pressure value.

The thermodynamic properties equations are due to Ishimoto et al. [15] and are presented below.

1. Saturation temperature:

$$\frac{T_c}{T} = \left(1 + \sum_{i=1}^5 A_i x^i \right) / \left(1 + \sum_{i=6}^7 A_i x^{i-5} \right)$$

2. Specific volume of saturated water:

$$v_f/v_c = \left(1 + \sum_{i=1}^9 B_i \gamma^i \right) / \left(1 + \sum_{i=10}^{12} B_i \gamma^{i-9} \right)$$

3. Specific volume of saturated steam:

$$v_c/v_g = \left(\frac{T}{T_c} \right)^6 \left(1 + \sum_{i=1}^7 C_i \gamma^i \right) / \left(1 + \sum_{i=8}^{12} C_i \gamma^{i-7} \right)$$

4. Specific enthalpy of saturated water:

$$h_f/h_c = \left(1 + \sum_{i=1}^6 D_i \gamma^i \right) / \left(1 + \sum_{i=7}^{10} D_i \gamma^{i-6} \right)$$

5. Specific enthalpy of saturated steam:

$$h_g/h_c = \left(1 + \sum_{i=1}^6 E_i \gamma^i \right) / \left(1 + \sum_{i=7}^{10} E_i \gamma^{i-6} \right)$$

6. Specific entropy of saturated water:

$$s_f/s_c = \left(1 + \sum_{i=1}^6 F_i \gamma^i \right) / \left(1 + \sum_{i=7}^{10} F_i \gamma^{i-6} \right)$$

7. Specific entropy of saturated steam:

$$s_g/s_c = \left(1 + \sum_{i=1}^6 G_i \gamma^i \right) / \left(1 + \sum_{i=7}^{10} G_i \gamma^{i-6} \right)$$

Critical Properties for Water:

$$h_c = 2.107209 \times 10^6 \text{ J/kg}$$

$$P_c = 2.2119256 \times 10^7 \text{ N/m}^2$$

$$s_c = 4.4427222 \times 10^3 \text{ J/kg, K}$$

$$T_c = 647.27 \text{ K}$$

$$v_c = 3.1698 \times 10^{-3} \text{ m}^3/\text{kg}$$

Symbols:

$$\beta = \text{reduced pressure } P/P_c$$

$$X = \ln(1/\beta)$$

$$Y = 1 - \theta$$

Constants:

$$A_1 = 0.129\ 542\ 684\ 9 \times 10^0$$

$$A_2 = 0.451\ 360\ 243\ 6 \times 10^0$$

$$A_3 = 0.212\ 237\ 843\ 6 \times 10^0$$

$$A_4 = -0.358\ 143\ 709\ 6 \times 10^{-2}$$

$$A_5 = 0.909\ 031\ 628\ 3 \times 10^{-4}$$

$$A_6 = 0.333\ 853\ 062\ 2 \times 10^1$$

$$A_7 = 0.142\ 655\ 898\ 5 \times 10^1$$

Constants (continued)

$$B_1 = 0.103\ 083\ 704\ 2 \times 10^5$$

$$B_2 = 0.553\ 842\ 656\ 9 \times 10^7$$

$$B_3 = 0.179\ 977\ 219\ 3 \times 10^9$$

$$B_4 = -0.474\ 384\ 340\ 7 \times 10^9$$

$$B_5 = 0.201\ 974\ 977\ 9 \times 10^{10}$$

$$B_6 = -0.625\ 856\ 400\ 9 \times 10^{10}$$

$$B_7 = 0.123\ 093\ 751\ 0 \times 10^{10}$$

$$B_8 = -0.135\ 116\ 048\ 0 \times 10^{11}$$

$$B_9 = 0.635\ 376\ 971\ 4 \times 10^{10}$$

$$B_{10} = 0.116\ 326\ 898\ 4 \times 10^5$$

$$B_{11} = 0.768\ 677\ 553\ 4 \times 10^7$$

$$B_{12} = 0.369\ 408\ 938\ 0 \times 10^9$$

$$C_1 = 0.411\ 687\ 661\ 4 \times 10^4$$

$$C_2 = 0.882\ 469\ 159\ 6 \times 10^6$$

$$C_3 = 0.433\ 265\ 710\ 0 \times 10^7$$

$$C_4 = -0.283\ 542\ 718\ 9 \times 10^8$$

$$C_5 = 0.365\ 887\ 547\ 4 \times 10^8$$

$$C_6 = 0.129\ 983\ 597\ 3 \times 10^7$$

$$C_7 = -0.172\ 949\ 642\ 3 \times 10^8$$

$$C_8 = 0.478\ 039\ 111\ 9 \times 10^4$$

$$C_9 = 0.143\ 626\ 686\ 2 \times 10^7$$

$$C_{10} = 0.273\ 219\ 597\ 4 \times 10^8$$

$$C_{11} = -0.861\ 139\ 257\ 8 \times 10^8$$

$$C_{12} = 0.898\ 119\ 690\ 3 \times 10^8$$

$$D_1 = 0.503\ 588\ 841\ 8 \times 10^4$$

$$D_2 = 0.852\ 988\ 031\ 0 \times 10^6$$

$$D_3 = 0.397\ 566\ 129\ 4 \times 10^7$$

$$D_4 = -0.132\ 121\ 217\ 4 \times 10^8$$

$$D_5 = 0.747\ 256\ 735\ 8 \times 10^7$$

$$D_6 = -0.169\ 003\ 475\ 3 \times 10^7$$

$$D_7 = 0.528\ 873\ 227\ 7 \times 10^4$$

$$D_8 = 0.982\ 421\ 289\ 1 \times 10^6$$

$$D_9 = 0.782\ 590\ 628\ 8 \times 10^7$$

$$D_{10} = -0.462\ 022\ 817\ 7 \times 10^7$$

$$E_1 = 0.498\ 130\ 313\ 7 \times 10^4$$

$$E_2 = 0.155\ 796\ 661\ 1 \times 10^7$$

$$E_3 = 0.321\ 871\ 068\ 9 \times 10^8$$

$$E_4 = 0.958\ 496\ 511\ 4 \times 10^7$$

$$E_5 = -0.439\ 895\ 952\ 3 \times 10^8$$

$$E_6 = 0.272\ 498\ 390\ 8 \times 10^8$$

$$E_7 = 0.480\ 296\ 564\ 1 \times 10^4$$

$$E_8 = 0.138\ 587\ 610\ 4 \times 10^7$$

$$E_9 = 0.232\ 356\ 620\ 5 \times 10^8$$

$$E_{10} = 0.813\ 454\ 111\ 0 \times 10^6$$

Constants (continued)

$$F_1 = 0.500\ 361\ 430\ 5 \times 10^4$$

$$F_2 = 0.827\ 385\ 789\ 8 \times 10^6$$

$$F_3 = 0.389\ 385\ 240\ 7 \times 10^7$$

$$F_4 = -0.115\ 966\ 636\ 4 \times 10^8$$

$$F_5 = 0.336\ 964\ 036\ 5 \times 10^7$$

$$F_6 = 0.122\ 650\ 602\ 1 \times 10^7$$

$$F_7 = 0.518\ 729\ 581\ 3 \times 10^4$$

$$F_8 = 0.916\ 803\ 401\ 9 \times 10^6$$

$$F_9 = 0.631\ 959\ 188\ 4 \times 10^7$$

$$F_{10} = -0.676\ 264\ 517\ 9 \times 10^7$$

$$G_1 = 0.592\ 250\ 267\ 9 \times 10^4$$

$$G_2 = 0.212\ 150\ 645\ 3 \times 10^7$$

$$G_3 = 0.436\ 715\ 991\ 8 \times 10^8$$

$$G_4 = -0.175\ 543\ 802\ 6 \times 10^8$$

$$G_5 = -0.370\ 624\ 752\ 5 \times 10^8$$

$$G_6 = 0.430\ 345\ 978\ 4 \times 10^8$$

$$G_7 = 0.578\ 290\ 580\ 1 \times 10^4$$

$$G_8 = 0.195\ 084\ 025\ 9 \times 10^7$$

$$G_9 = 0.335\ 195\ 702\ 5 \times 10^8$$

$$G_{10} = -0.360\ 433\ 023\ 1 \times 10^8$$

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