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# Taking Care of Business in a Flash †: Constraining the Timescale for Low-Mass Satellite Quenching with ELVIS

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## ABSTRACT

The vast majority of dwarf satellites orbiting the Milky Way and M31 are quenched, while comparable galaxies in the field are gas-rich and star-forming. Assuming that this dichotomy is driven by environmental quenching, we use the ELVIS suite of  $N$ -body simulations to constrain the characteristic timescale upon which satellites must quench following infall into the virial volumes of their hosts. The high satellite quenched fraction observed in the Local Group demands an extremely short quenching timescale ( $\sim 2$  Gyr) for dwarf satellites in the mass range  $M_\star \sim 10^6 - 10^8 M_\odot$ . This quenching timescale is significantly shorter than that required to explain the quenched fraction of more massive satellites ( $\sim 8$  Gyr), both in the Local Group and in more massive host halos, suggesting a dramatic change in the dominant satellite quenching mechanism at  $M_\star \lesssim 10^8 M_\odot$ . Combining our work with the results of complementary analyses in the literature, we conclude that the suppression of star formation in massive satellites ( $M_\star \sim 10^8 - 10^{11} M_\odot$ ) is broadly consistent with being driven by starvation, such that the satellite quenching timescale corresponds to the cold gas depletion time. Below a critical stellar mass scale of  $\sim 10^8 M_\odot$ , however, the required quenching times are much shorter than the expected cold gas depletion times. Instead, quenching must act on a timescale comparable to the dynamical time of the host halo. We posit that ram-pressure stripping can naturally explain this behavior, with the critical mass (of  $M_\star \sim 10^8 M_\odot$ ) corresponding to halos with gravitational restoring forces that are too weak to overcome the drag force encountered when moving through an extended, hot circumgalactic medium.

**Key words:** Local Group – galaxies: formation – galaxies: evolution – galaxies: dwarf – galaxies: star formation

## 1 INTRODUCTION

Foremost among the results of galaxy surveys over the last decade has been the realization that the galaxy population at  $z \lesssim 2$  is bimodal in nature (e.g. Strateva et al. 2001; Baldry et al. 2004; Bell et al. 2004; Cooper et al. 2006). That is, galaxies both locally and out to intermediate redshift are effectively described as one of two distinct types: red, early-type galaxies lacking significant star formation and blue, late-type galaxies with active star formation. In color-magnitude space, the red galaxies populate a tight relation (often called the red sequence), while the distribution of blue galaxies is more scattered (sometimes referred to as the blue cloud). While the red and blue populations

comprise approximately equal portions of the cosmic stellar mass budget at  $z \sim 1$ , galaxies on the red sequence dominate today, following a growth in stellar mass within the red population of roughly a factor of 2 over the past 7 Gyr (Bell et al. 2004; Bundy et al. 2006; Faber et al. 2007; Brown et al. 2007). Despite uncertainty regarding the particular physical process(es) at play, the suppression (or quenching) of star formation in blue galaxies, thereby making them red, is one of the principal drivers of this dramatic growth in the number density of quiescent systems at late cosmic time.

At both low and intermediate redshift, the local environment of a galaxy is known to be well-correlated with the suppression of star formation, such that passive or quiescent galaxies preferentially live in higher-density environments (Balogh et al. 2004; Kauffmann et al. 2004; Blanton et al. 2005; Cooper et al. 2006, 2007, 2010a). While the stel-

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lar mass or surface mass density of a galaxy may be more closely connected to quenching for massive systems (Peng et al. 2010; Cheung et al. 2012; Woo et al. 2013), recent work has shown that environment is likely the dominant driver of quenching at the lowest mass scales ( $M_* < 10^9 M_\odot$ , Geha et al. 2012). For example, studies comparing satellite galaxies to isolated field systems of similar stellar mass in the local Universe find that satellites tend to exhibit lower star formation rates, more bulge-dominated morphologies, as well as older and more metal-rich stellar populations (Baldry et al. 2006; van den Bosch et al. 2008; Cooper et al. 2010b; Pasquali et al. 2010; Tollerud et al. 2011; Phillips et al. 2014). This observed suppression of star formation in satellite galaxies is commonly referred to as “environmental quenching”. At present, galaxy formation models fail to reproduce this environment-dependent effect in detail, with both semi-analytic and hydrodynamical models overpredicting the number of quenched satellites at low masses (Kimm et al. 2009; Weinmann et al. 2011, 2012; Hirschmann et al. 2014; Wang et al. 2014).

Understanding the physics of environmental quenching is complicated by the wide range of physical processes that could be responsible for suppressing star formation in low-mass satellites. For example, various quenching mechanisms preferentially operate on satellite galaxies within overdense environments, including [i] “starvation” or “strangulation” – whereby gas accretion onto a satellite galaxy is halted following infall, thus eventually eliminating the fuel for star formation (Larson, Tinsley & Caldwell 1980; Kawata & Mulchaey 2008), [ii] “harassment” – by which close encounters between densely packed cluster or group members strip gas from around the interacting galaxies (Moore et al. 1996), and [iii] “ram-pressure stripping” – where the cold dense gas at the center of a satellite (or the hot halo surrounding it) is removed from the galaxy as a result of a high-speed interaction with the hot gas halo of the host (Gunn & Gott 1972; Bekki 2009).

To further complicate matters, each of these processes may be important in different host halos and at different satellite mass scales. In an attempt to differentiate between these various physical mechanisms, we aim to measure the timescale upon which satellite quenching occurs across a broad range of satellite stellar mass. Throughout this work, we define this quenching timescale ( $\tau_{\text{quench}}$ ) relative to the time at which a satellite was accreted onto its current host system (i.e. relative to the infall time of the satellite). While quenching via starvation will proceed according to the available cold gas reservoir for star formation within an accreted satellite, ram-pressure stripping will largely act on a timescale set by the density distribution of the host’s hot halo and infall velocity of the satellite. As such, constraining the timescale for satellite quenching may therefore serve as a critical step towards differentiating between various physical mechanisms.

By comparing satellite samples at low and intermediate redshift to galaxy and dark matter-only simulations, several recent analyses of satellites at high stellar mass ( $> 5 \times 10^9 M_\odot$ ) point towards relatively long quenching timescales of several Gyr or more (Wetzell et al. 2013; De Lucia et al. 2012). Remarkably, when pushing this analysis to lower masses ( $\sim 10^{8.5} - 10^{9.5} M_\odot$ ), Wheeler et al. (2014) find that quenching remains surprisingly inefficient, such that the

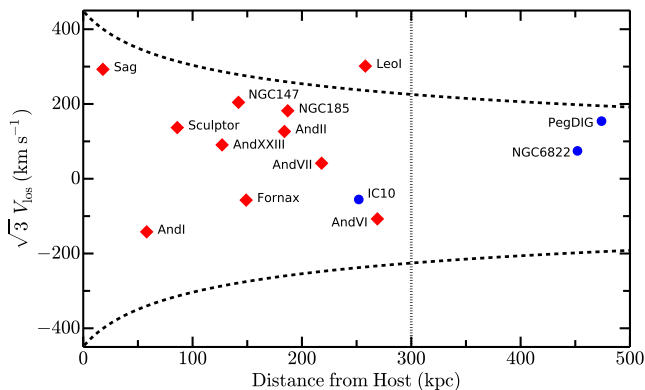
average quenching timescale increases to  $> 7$  Gyr. Here, we extend these previous analyses of quenching timescales to yet lower masses by studying the satellite galaxies of the Local Group. Through comparison to simulations, we aim to establish a coherent picture of the timescale upon which satellite quenching acts, spanning more than five orders of magnitude in satellite mass. In §2, we detail the population of Local Group satellites along with an estimation of the observed quenched fraction. In §3 and §4, we describe the relevant  $N$ -body simulations and physically-motivated quenching models, respectively, detailing the relevant constraints on the derived quenching timescale at low masses. Finally, in §5 and §6, we discuss the implications of our analysis, including a “back of the envelope” explanation of the underlying physics operating on low-mass satellites and a summary of our main results.

## 2 LOCAL GROUP SATELLITES

In assembling our sample of Milky Way and M31 satellites, we utilize the compilation of local dwarf galaxy properties presented by McConnachie (2012). To minimize the potential impact of incompleteness at low masses, we limit our analysis to those dwarfs with measured stellar masses in the range  $10^6 - 10^8 M_\odot$ ; however, as discussed in more detail below, including systems at lower masses has very little impact on our qualitative or quantitative results. From this population of local dwarfs, satellite systems are selected to be within 300 kpc of the Milky Way or M31 and bound with respect to their host – i.e.  $\sqrt{3} \cdot V_{\text{los}} < V_{\text{esc}}(r)$ , where  $V_{\text{esc}}(r)$  is the escape velocity at a radial distance of  $r$  for an NFW halo with a concentration ( $c$ ) of 8 and a total halo mass of  $2 \times 10^{12} M_\odot$ .<sup>1</sup> To account for potential tangential motion, we multiply all observed line-of-sight velocities from McConnachie (2012) by a factor of  $\sqrt{3}$  in evaluating whether each system is bound. Within the stellar mass range of  $10^6 M_\odot$  to  $10^8 M_\odot$ , Figure 1 shows the position and velocity of all Local Group dwarfs relative to their host (or nearest massive neighbor in the case of field systems such as the Pegasus dIrr).

We identify 12 satellite systems meeting our selection criteria (see Table 1), with the final set of satellites and corresponding quenched fraction relatively independent of the adopted selection criteria. For instance, field dwarfs such as NGC 6822, Leo T, and Phoenix, which are nearly in our stellar mass limits, are located significantly beyond our radial selection limit (i.e.  $> 400$  kpc from the Milky Way). Moreover, while the observed line-of-sight velocity of Leo I suggests that it may be unbound (see Fig. 1), recent proper motion observations conclude that the system is very likely bound to the Milky Way even at a distance of more than 250 kpc (Sohn et al. 2013; Boylan-Kolchin et al. 2013). Finally, given the high metallicity and lack of overdensity in RR Lyrae stars, we categorize Canis Major as an overdensity in the Milky Way and not a satellite (Momany et al. 2004; Butler et al. 2007; Mateu et al. 2009, but see also Martínez-Delgado et al. 2005).

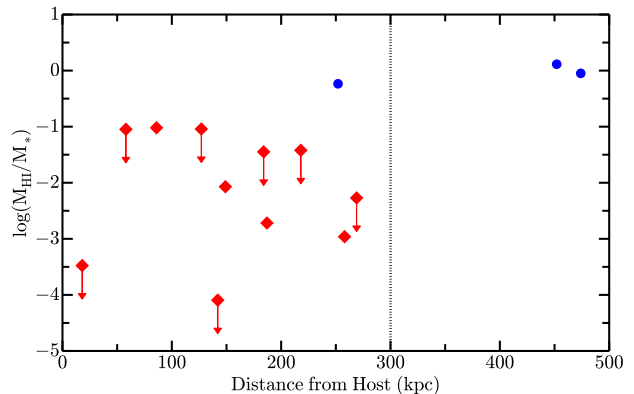
<sup>1</sup> Varying the adopted host halo mass and concentration (and thus our satellite selection limits) does not have a significant effect on our final sample of Local Group satellites.



**Figure 1.** Host-centric velocity versus distance for all known Local Group galaxies in the stellar mass range of  $10^6 - 10^8 M_{\odot}$ . To account for unknown tangential velocities, the observed line-of-sight velocities have been multiplied by a factor of  $\sqrt{3}$ . The dotted vertical line at 300 kpc corresponds to the expected virial radius of the Milky Way and M31 and represents the radial limit for our satellite selection. The dashed lines show the escape velocity for an NFW halo with a concentration ( $c$ ) of 8 and a total mass of  $2 \times 10^{12} M_{\odot}$ . Satellites are color-coded as star-forming (blue circles) or quenched (red diamonds) according to their observed HI gas fraction (see Fig. 2).

To determine whether a given satellite is quenched, we use the observed HI gas fraction as a proxy for the current star formation rate (SFR), defining dwarfs with atomic gas fractions less than 10% to be quenched (see Fig. 2). This criterion leads to only one satellite in our sample, IC 10, categorized as currently star-forming and yields a quenched fraction of  $> 90\%$  for satellites bound to the Milky Way and M31 in the selected mass range. Alternative definitions of quenched versus star-forming – e.g. using star-formation histories derived from photometric observations of spatially-resolved stellar populations (Weisz et al. 2014a, 2015) or star-formation rates inferred from H $\alpha$  emission (Kaisin & Karachentsev 2013; Karachentsev & Kaisina 2013) yield similar results. Moreover, our measurement of the satellite quenched fraction is not strongly dependent on the adopted mass range, as the quenched fraction is largely unchanged if we extend the lower mass limit down to  $3 \times 10^5 M_{\odot}$  so as to account for potential uncertainties or scatter in the stellar mass measurements. Over this expanded stellar mass range, the sample would then also include Leo II, And XXI, And XXV, and LGS 3. While HI measurements suggest that LGS 3 is likely star-forming, a variety of observations, including star formation rates inferred from H $\alpha$  narrow-band imaging, indicate that Leo II, And XXI, and And XXV are quiescent (Grcevich & Putman 2009; Kaisin, Karachentsev & Ravindranath 2012; Kaisin & Karachentsev 2013; Spekkens et al. 2014; Weisz et al. 2014a,b). When including these four lower-mass systems in our sample, the resulting quenched fraction remains quite high ( $\sim 87\%$ ); furthermore, pushing to yet lower masses will similarly yield little variation in the quenched fraction, as the lowest-mass satellites in the Local Group are universally quenched (e.g. Okamoto et al. 2008; Sand et al. 2009, 2010; Weisz et al. 2014a; Brown et al. 2014).

While our census of satellites in the Local Group is likely incomplete at low masses due to obscuration by or confu-



**Figure 2.** Observed HI gas fraction as a function of distance from the nearest host galaxy for local dwarfs included in our mass range of  $10^6 - 10^8 M_{\odot}$ . Systems with an atomic gas fraction less than 10% are classified as quenched (red diamonds), with more gas-rich dwarfs identified as star-forming (blue circles). At these low masses, roughly 90% of satellites are quenched in the Local Group.

sion with the disk of both the Milky Way and M31, the focus of our analysis is the relative fraction of star-forming versus quenched systems. Assuming that neither quenched nor star-forming systems preferentially reside in or behind the disk, the quenched fraction should be independent of any such incompleteness. From an analysis of the Palomar and UK Schmidt photographic plates, which cover more than 20,000 square degrees of sky at high Galactic latitude ( $|b| > 20^\circ$ ), Irwin (1994) finds that the known satellite population for the Milky Way is complete down to below our adopted mass limit ( $\gtrsim 3 \times 10^5 M_{\odot}$ ). This conclusion is supported by deeper imaging data from the Sloan Digital Sky Survey (SDSS, York et al. 2000), Pan-STARRS (Kaiser et al. 2010), and Dark Energy Survey (DES, Diehl et al. 2014), which cover a significant portion of the northern and southern sky and are complete to beyond the host virial radius (i.e.  $> 300$  kpc) at these stellar masses (Tollerud et al. 2008; Koposov et al. 2008; Laevens et al. 2014; Koposov et al. 2015; Bechtol et al. 2015).

For M31, our knowledge of satellites is largely based on imaging surveys such as the Pan Andromeda Archaeological Survey (PAndAS, McConnachie et al. 2009, see also Ferguson et al. 2002; Ibata et al. 2007; McConnachie et al. 2008). As shown by Brasseur et al. (2011), PAndAS is complete down to our mass limit within a survey footprint that covers a radial distance of roughly 150 kpc surrounding M31. While recent analysis of imaging data, including that from Pan-STARRS, has likely uncovered several new satellites of M31 at radial distances of  $\gtrsim 150$  kpc (And XXX-XXXIV, Conn et al. 2012; Martin et al. 2013b,a, 2014), we exclude these systems from our study given that distances (relative to M31) are poorly constrained and potentially greater than 300 kpc. Early analysis of the stellar populations in Andromeda XXXI and XXXII, however, do suggest passive populations devoid of young stars (Martin et al. 2013b). As likely quenched satellites, we note that inclusion of And XXX-XXXIV would have little impact on the observed quenched fraction in the Local Group and thus on the results of our analysis. Moreover, while the M31 satel-

Galaxy Name	Morphology (1)	$M_{\star}$ ( $10^6 M_{\odot}$ ) (2)	$M_{\text{HI}}$ ( $10^6 M_{\odot}$ ) (3)	$D_{\text{host}}$ (kpc) (4)	$V_{\text{los}}$ (km s $^{-1}$ ) (5)	Quenched (6)
Milky Way						
Sagittarius	dSph	21	< 0.007	18	169	Yes
Sculptor	dSph	2.3	0.22	86	79	Yes
Fornax	dSph	20	0.17	149	-33	Yes
Leo I	dSph	5.5	< 0.006	258	174	Yes
M31						
And I	dSph	3.9	< 0.35	58	-82	Yes
And XXIII	dSph	1.1	< 0.1	127	52	Yes
NGC 147	dE/dSph	62	< 0.003	142	118	Yes
And II	dSph	7.6	< 0.27	184	73	Yes
NGC 185	dE/dSph	68	0.13	187	105	Yes
And VII	dSph	9.5	< 0.36	218	24	Yes
IC 10	dIrr	86		252	-32	No
And VI	dSph	2.8	< 0.015	269	-62	Yes

**Table 1.** Observed properties of satellite galaxies in the Local Group with measured stellar masses in the range  $10^6 - 10^8 M_{\odot}$  and located within 300 kpc of either the Milky Way or M31 (separated according to host system): (1) morphological classification and (2) stellar mass from McConnachie (2012); (3) atomic gas mass measurements from McConnachie (2012) and Grcevich & Putman (2009), plus upper limits from the literature (Burton & Lockman 1999; Grcevich & Putman 2009; Huang et al. 2012b; Giovanelli et al. 2013); (4) distance from the nearest host as given by McConnachie (2012); (5) line-of-sight velocity with respect to the host from McConnachie (2012), with the exception of And XXIII (Kirby et al. 2014); (6) identification as quenched versus star-forming, according to observed atomic gas fraction (see Fig. 2).

lite population may be incomplete at radial distances of 150 – 300 kpc from M31, this should have minimal impact on the observed quenched fraction in the Local Group (even assuming that star-forming satellites are possibly biased towards large radial distances, as our results would imply). In §4, we discuss this potential selection effect in more detail.

### 3 SIMULATIONS

To model the evolution of satellite galaxies in the Local Group, we utilize the Exploring the Local Volume In Simulations (ELVIS) suite of 48 high-resolution, dissipationless simulations of Milky Way-like halos (Garrison-Kimmel et al. 2014a). The suite includes 24 isolated halos as well as 12 mass-matched Local Group-like pairs, simulated within high-resolution uncontaminated volumes spanning 2-5 Mpc in size using a particle mass of  $1.9 \times 10^5 M_{\odot}$  and a Plummer-equivalent force softening of  $\epsilon = 141$  physical parsecs. Within the high-resolution volumes, the halo catalogs are complete down to  $M_{\text{halo}} > 2 \times 10^7 M_{\odot}$ ,  $V_{\text{max}} > 8 \text{ km s}^{-1}$ ,  $M_{\text{peak}} > 6 \times 10^7 M_{\odot}$ , and  $V_{\text{peak}} > 12 \text{ km s}^{-1}$  — thus more than sufficient to track the evolution of halos hosting Local Group dwarfs with stellar masses of  $> 10^6 M_{\odot}$ . ELVIS adopts a cosmological model based on Wilkinson Microwave Anisotropy Probe 7 (Komatsu et al. 2011), with the following  $\Lambda$ CDM parameters:  $\sigma_8 = 0.801$ ,  $\Omega_m = 0.266$ ,  $\Omega_{\Lambda} = 0.734$ ,  $n_s = 0.963$ , and  $h = 0.71$ .

As hosts for the Local Group dwarfs in the stellar mass range of  $10^6 - 10^8 M_{\odot}$ , we select halos in the ELVIS simulations with masses of  $M_{\text{peak}} = 5 \times 10^9 - 6 \times 10^{10} M_{\odot}$ , following the abundance matching prescription of Garrison-Kimmel et al. (2014a). Abundance matching, a common technique for populating simulated dark matter distributions with galaxies, assumes a one-to-one relation between

a galaxy’s stellar mass and the mass of its parent dark matter (sub)halo (Behroozi et al. 2013; Moster, Naab & White 2013). While this simple empirical approach to modeling galaxy formation yields great success for massive galaxies, matching a wide range of clustering statistics as a function of cosmic time (e.g. Berrier et al. 2006; Conroy, Wechsler & Kravtsov 2006), it potentially breaks down at low masses, where the most massive subhalos of Milky Way-like simulations are inconsistent with the observed dynamics of nearby dwarf galaxies (Boylan-Kolchin, Bullock & Kaplinghat 2011, 2012; Garrison-Kimmel et al. 2014b; Kirby et al. 2014; Tollerud, Boylan-Kolchin & Bullock 2014).

While uncertainties remain in the potential efficacy or robustness of abundance matching prescriptions at low masses, our results are largely unaffected by changes in the assumed stellar mass-halo mass relation. As our fiducial model, we adopt the abundance matching relation of Garrison-Kimmel et al. (2014a), which in our stellar mass range ( $10^6 - 10^8 M_{\odot}$ ) yields 9 subhalos as candidate satellite hosts within a typical ELVIS parent halo (after removing stripped or disrupted subhalos). When applying the shallower stellar mass-halo mass relation from Behroozi et al. (2013), the number of potential subhalo hosts increases accordingly, but the quantitative and qualitative results with regard to satellite quenching remain unchanged. To account for potential increased scatter in the abundance matching relation at low masses, we also mimic scatter in the Garrison-Kimmel et al. (2014a) relation, by broadening the range of halo masses associated with the observed stellar mass range, randomly selecting 9 subhalos within  $2 \times 10^9 < M_{\text{peak}}/M_{\odot} < 2 \times 10^{11}$  as satellite hosts. Again, while this biases the subhalo population to lower masses, our results with regard to satellite quenching remain unaffected.

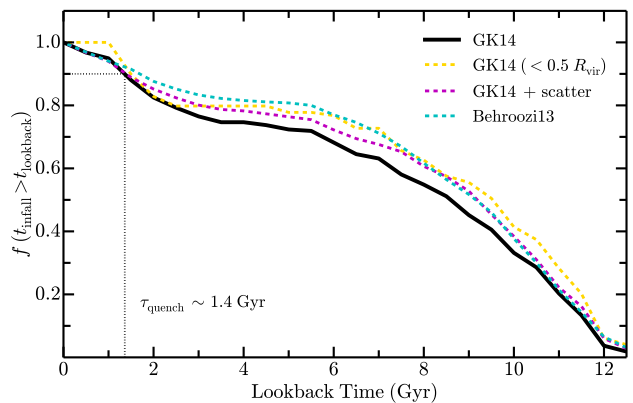
Independent of the particular abundance matching pre-

scription, we account for satellites that have been tidally-stripped or disrupted — and thus not included in our observational sample — by restricting our analysis to subhalos for which  $V_{\max}/V_{\text{peak}} > 0.3$ , where  $V_{\max}$  is the current ( $z = 0$ ) maximum circular velocity for the halo and  $V_{\text{peak}}$  is the value of  $V_{\max}$  when the halo mass is at its maximum. While a halo’s maximum halo mass generally occurs prior to infall (typically at  $> 1.5 R_{\text{vir}}$ ),  $V_{\text{peak}}$  is a good approximation for  $V_{\text{infall}}$  since a halo’s mass, in general, does not change appreciably prior to infall (Behroozi et al. 2014). By restricting our analysis to present-day halos with  $V_{\max}/V_{\text{peak}} > 0.3$ , we eliminate all subhalos (i.e. satellites) that have lost greater than 90% of their mass since infall, thereby selecting those subhalos that are within the desired mass range and are highly likely to currently host dwarf galaxies (i.e. not disrupted or significantly stripped). Finally, we note that our results are not strongly dependent upon the adopted  $V_{\max}/V_{\text{peak}}$  cut; applying a more conservative selection criterion (e.g.  $V_{\max}/V_{\text{peak}} > 0.5$ ) does not significantly change our results.

The ELVIS merger trees include 75 snapshots from  $z = 125$  to  $z = 0$ , providing time resolution of roughly 200 Myr. As discussed in more detail in §4, we are interested in using ELVIS to explore models where satellite quenching occurs at a fixed physical radius from the host (e.g. within 50 kpc radially). The crossing time on such scales ( $\sim 100$  kpc), however, is typically greater than the median time resolution of the snapshots, such that subhalos would potentially pass into and back out of the quenching radius between two successive snapshots. To more precisely determine subhalo orbits within ELVIS, we map the spatial position of each subhalo (relative to its host) at 20 Myr intervals by spline interpolating the position of each subhalo and corresponding host halo across each of the 75 snapshots. This interpolation also allows us to determine the time at which each subhalo was accreted onto its host halo (i.e. the infall time or  $t_{\text{infall}}$ ) at greater precision. Due to the linear spacing of the ELVIS snapshots with scale factor (i.e. non-linear spacing with lookback time), the difference in the measured infall time between our interpolated data and that inferred from the standard ELVIS merger trees depends upon the lookback time at which a subhalo is accreted. For systems with  $t_{\text{infall}} < 3$  Gyr, which place the strongest constraints on the satellite quenching timescale, the difference in  $t_{\text{infall}}$  is relatively modest, with a median offset of less than 0.5 Gyr for our fiducial sample of subhalos.

## 4 QUENCHING MODELS

As highlighted in §1, observations of low-mass galaxies in the local Universe suggest that all isolated (or “field”) systems at  $\lesssim 10^9 M_{\odot}$  are star-forming. Thus, all quenching at these low masses is likely driven by environmental effects associated with dwarf galaxies being accreted onto more massive halos, such as that of the Milky Way or M31. Two likely quenching mechanisms operating on satellite systems are starvation, where future gas accretion onto the satellite is halted following infall into a more massive host halo, and ram-pressure stripping, where the cold gas in a satellite galaxy is violently removed via interaction with the hot gas halo of the host. To mimic these respective processes, we



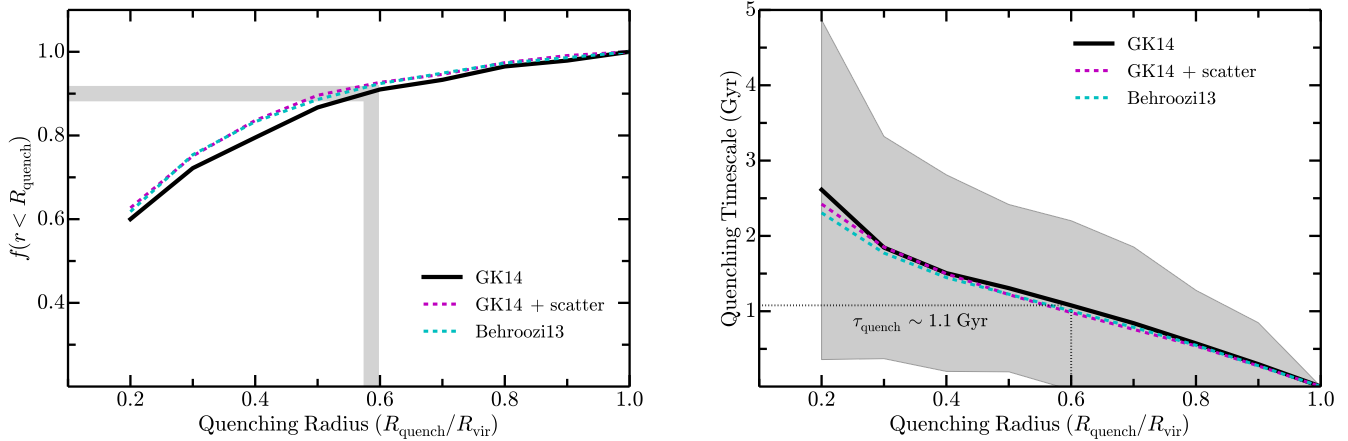
**Figure 3.** The cumulative distribution of infall times ( $t_{\text{infall}}$ ) as a function of lookback time for those subhalos in the ELVIS simulations identified as hosts of the Local Group satellites at  $10^6 < M_*/M_{\odot} < 10^8$ . The solid black line corresponds to our default selection criteria, where subhalos are selected according to the abundance matching relation of Garrison-Kimmel et al. (2014a), restricting to those subhalos with  $V_{\max}/V_{\text{peak}} > 0.3$  and within the host virial radius at  $z = 0$ . The dashed cyan and magenta lines show the corresponding distribution of infall times when applying alternative abundance matching prescriptions (see §3), while the dashed gold line traces the accretion history of only those subhalos within  $0.5 R_{\text{vir}}$  at  $z = 0$  following the Garrison-Kimmel et al. (2014a) abundance matching relation. A quenched fraction of 90% today, as observed for the Local Group at these masses, corresponds to a quenching timescale ( $\tau_{\text{quench}}$ ) of  $\sim 1.4$  Gyr.

implement two quenching models in the ELVIS simulations: (i) a “starvation” model in which quenching depends only on time since infall into the host halo and (ii) a “stripping” model in which quenching occurs once a satellite reaches a fixed physical distance from the host.

### 4.1 Starvation Model

When quenching occurs via starvation, the satellite galaxy is deprived of any additional supply of cold gas after infall, such that star formation is halted following the consumption of its current gas reservoir — i.e. quenching occurs within the gas depletion timescale ( $\tau_{\text{depl}} = M_{\text{gas}}/\text{SFR}$ ). While observations suggest that atomic and/or molecular depletion timescales may vary with galaxy mass (Leroy et al. 2008; Bigiel et al. 2011; Boselli et al. 2014; Somerville, Popping & Trager 2015), we adopt a fixed quenching timescale (following infall) over the limited stellar mass range probed in our analysis ( $10^6 - 10^8 M_{\odot}$ ).

In the models, we define the infall time as the lookback time at which a halo *first* crosses the virial radius,  $R_{\text{vir}}(z)$ , of the host. For the  $\sim 23\%$  of subhalos that cross the host virial radius more than once, we are likely overestimating the infall time. However, 90% of these subhalos are accreted for a final time  $\gtrsim 6$  Gyr ago, thereby contributing little to no effect on the final inferred quenching timescale. Figure 3 shows the cumulative distribution of infall times for subhalos selected as hosts of Local Group satellites in the simulations following various abundance matching prescriptions. Universally, we find that subhalos are accreted relatively uniformly over



**Figure 4.** *Left:* The dependence of the satellite quenched fraction on the assumed quenching radius,  $R_{\text{quench}}$ , for our ram-pressure stripping model (see §4.2). The solid black and dashed cyan and magenta lines correspond to subhalo populations selected according to different abundance matching prescriptions (see §3). In all cases, a quenching radius of  $0.6 R_{\text{vir}}$  is required to reproduce a quenched fraction of 90% as observed in the Local Group. *Right:* The dependence of the mean quenching timescale ( $\tau_{\text{quench}}$ ) on the assumed quenching radius ( $R_{\text{quench}}$ ) for the same set of abundance matching relations. For our fiducial subhalo sample, the grey shaded region illustrates the  $1\sigma$  scatter in the quenching timescale across the 48 ELVIS hosts. As for the starvation model, the high observed quenched fraction in the Local Group implies a short quenching timescale ( $\sim 1.1$  Gyr) in our ram-pressure stripping model. The typical  $R_{\text{vir}}$  (at  $z = 0$ ) for ELVIS hosts is  $\sim 300$  kpc.

cosmic time, with the vast majority infalling into their host halo  $> 2$  Gyr ago.

As shown in Figure 3, to match the observed satellite quenched fraction of  $\sim 90\%$  in the Local Group, the starvation model – independent of the adopted abundance matching prescription – favors a quenching timescale ( $\tau_{\text{quench}}$ ) of roughly  $1.4^{+0.9}_{-0.7}$  Gyr, including the error associated with the uncertainty in the observed Local Group quenched fraction (according to binomial statistics). Thus, on average, satellites of the Milky Way and M31 (at  $M_{\star} \sim 10^6 - 10^8 M_{\odot}$ ) must quench within 1.4 Gyr of infalling onto their respective host halo. Within the individual ELVIS simulations there is non-negligible scatter in the derived quenching timescale, given the stochasticity of subhalo accretion events; assuming a quenched fraction of 90%, the  $1\sigma$  scatter in  $\tau_{\text{quench}}$  across the 12 Local Group-like host pairs in ELVIS is approximately 0.7 Gyr.

A quenched fraction of 90% in the local group could be dependent on completeness of the dwarf sample in this mass range, as discussed in the introduction. This is particularly true around M31, where PAndAS imaging only surveys the volume within a radius of  $\sim 150$  kpc. To test the impact of this potential selection effect on our results, we also compute the distribution of infall times for those subhalos that reside within  $0.5 R_{\text{vir}}$  today. While simulated subhalos with earlier infall times are slightly biased towards smaller host-centric distances (e.g. Diemand & Kuhlen 2008; Rocha, Peter & Bullock 2012; Wetzel, Deason & Garrison-Kimmel 2015), the correlation between radial distance from the host and infall time exhibits significant scatter, such that the quenching timescale derived from subhalos within  $0.5 R_{\text{vir}}$  today is fully consistent with that based on the full sample of subhalos within  $1 R_{\text{vir}}$  (see Fig. 3).

In our starvation model, it is assumed that the quenching timescale (i.e. depletion timescale) is independent of cosmic time. While studies at higher redshift are limited to significantly more massive systems, current work suggests that

molecular depletion timescales are relatively weakly dependent on redshift (Geach et al. 2011; Magdis et al. 2012; Sain-tonge et al. 2013; Bauermeister et al. 2013), favoring shorter timescales at higher  $z$ . This conclusion is echoed in analysis of the satellite quenching timescale across a broad range of halo masses at intermediate redshift (McGee et al. 2011; McGee, Bower & Balogh 2014; Mok et al. 2014; Muzzin et al. 2014; Bahé & McCarthy 2015). Altogether, a modest variation in the quenching timescale with redshift would yield no significant impact on our results, as our analysis is focused on the cumulative quenched fraction today and not on the detailed quenching times of individual satellites.

## 4.2 Stripping Model

To mimic quenching via ram-pressure stripping, we implement a model in which satellites are quenched at a fixed radial distance ( $R_{\text{quench}}$ ) from their host. In a stripping scenario, this radial quenching scale would be set by the volume over which the host’s hot halo reaches a density capable of stripping an infalling dwarf. To account for growth of the hot halo over cosmic time, we define our quenching radius,  $R_{\text{quench}}$ , in terms of the host’s virial radius as a function of redshift (i.e.  $R_{\text{quench}} \propto R_{\text{vir}}(z)$ ) and require that a subhalo cross within the quenching radius at  $z < 3$ , allowing the host’s hot halo time to form (Birnboim & Dekel 2003; Kereš et al. 2005; Dekel & Birnboim 2006). For subhalos in ELVIS that have crossed within the quenching radius at  $z < 3$ , we classify the corresponding satellite as quenched; all other systems remain star-forming, independent of the time spent in the host halo (i.e. independent of infall time).

Figure 4a shows the dependence of the satellite quenched fraction on the adopted quenching radius (from  $0.2 R_{\text{vir}}$  to  $1 R_{\text{vir}}$ ). We find that to reproduce the observed quenched fraction in the Local Group (90%) requires a quenching radius of roughly  $0.6 R_{\text{vir}}$ , again with little dependence on the assumed abundance matching relation. Adopt-



ing this value for the quenching radius ( $R_{\text{quench}} = 0.6 R_{\text{vir}}$ ), we are able to determine the time at which each subhalo in the simulation (i.e. satellite) quenched following infall — and thus the quenching timescale (measured relative to infall). In Figure 4b, we show the mean quenching timescale as a function of the chosen quenching radius ( $R_{\text{quench}}$ ) within the 48 ELVIS simulations. For large values of  $R_{\text{quench}}$ , each satellite — largely independent of orbit — spends little time in the host halo prior to being quenched, such that the effective quenching timescale ( $\tau_{\text{quench}}$ ) is short. Adopting a quenching radius of  $R_{\text{quench}} = 0.6 R_{\text{vir}}$ , we find an average quenching timescale of 1.1 Gyr, in very good agreement with that derived from our starvation model (see §4.1). Finally, a complementary analysis by Slater & Bell (2014), in which quenching commences at pericenter passage, finds very similar results ( $\tau_{\text{quench}} \sim 1 - 2$  Gyr, see also Weisz et al. 2015; Wetzel, Tollerud & Weisz 2015).

## 5 DISCUSSION

Given that nearly all field dwarfs are star-forming at stellar masses of  $\lesssim 10^9 M_{\odot}$  (Grebel 1997; Mateo 1998; Geha et al. 2012), environment-dependent processes are the dominant driver of quenching at low masses and must act on a timescale so as to avoid over- or under-producing the number of star-forming satellite galaxies in the Local Group. Independent of the particular physical mechanism(s) at play, the high quenched fraction ( $\sim 90\%$ ) observed for low-mass Local Group satellites thereby dictates a remarkably short quenching timescale, measured relative to infall. Through a detailed comparison of the observed Milky Way and M31 satellites to the ELVIS suite of simulations, we show that this quenching timescale ( $\tau_{\text{quench}}$ ) is  $\sim 2$  Gyr at stellar masses of  $10^6 - 10^8 M_{\odot}$  in the Local Group.

A very similar analysis by Wheeler et al. (2014), in which the properties of nearby dwarfs in the SDSS are compared to the Millennium-II simulation (Boylan-Kolchin et al. 2009), finds that the quenching timescale must be much longer ( $> 7.5$  Gyr) at slightly higher stellar masses ( $\sim 10^9 M_{\odot}$ ). This longer quenching timescale is dictated by the relatively low quenched fractions observed in nearby groups (Geha et al. 2012). While Wheeler et al. (2014) preferentially study higher-mass host halos than that of the Milky Way and M31, analysis of the Local Group is in broad agreement. At  $M_{\star} \sim 10^8 - 10^{10} M_{\odot}$ , the satellite quenched fraction in the Local Group is only  $\sim 40\%$  (Wheeler et al. 2014; Phillips et al. 2015; Slater & Bell 2014), which requires a quenching timescale of  $\sim 7.8$  Gyr when comparing the infall times of ELVIS subhalos selected to have  $M_{\star} \sim 10^{8.5} - 10^{9.5} M_{\odot}$  according to the abundance matching relation of Garrison-Kimmel et al. (2014a).

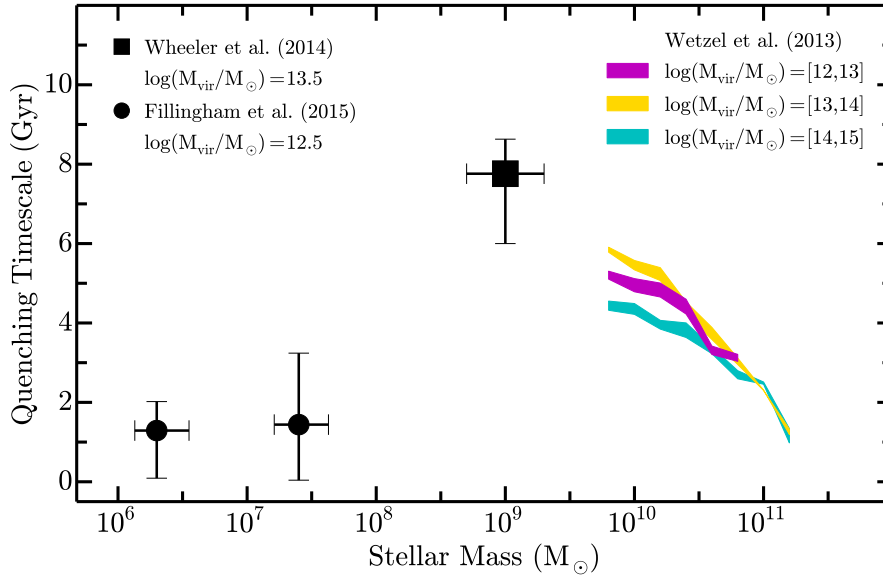
In an effort to study quenching for the most massive satellite systems, Wetzel et al. (2013) and De Lucia et al. (2012) compare group and cluster populations from the SDSS (Yang et al. 2007) to  $N$ -body and semi-analytic models, respectively. While analysis at these masses is complicated by the fact that some isolated or field systems are quenched independent of any environment-related effects, these independent analyses conclude that the quenching timescale at  $\sim 10^{10} - 10^{11} M_{\odot}$  increases with decreasing satellite stellar mass, reaching as long as  $\sim 5 - 6$  Gyr (see

also Hirschmann et al. 2014). Interestingly, for these most massive satellites, no evidence is found for variation in the quenching timescale with host halo mass (Wetzel et al. 2013, but see also Phillips et al. 2015).

In Figure 5, we combine our estimate of the quenching timescale for low-mass satellites in the Local Group ( $M_{\star} \sim 10^7 M_{\odot}$ ) with the complementary constraints from Wheeler et al. (2014) and Wetzel et al. (2013) at higher stellar masses. Due to differences in the adopted definition of infall time (see discussion in Wheeler et al. 2014), we assume a quenching timescale of 8 Gyr at  $10^9 M_{\odot}$ . This value for  $\tau_{\text{quench}}$  is fully consistent with the results from Wheeler et al. (2014,  $\tau_{\text{quench}} \sim 7.5 - 9$  Gyr); as previously highlighted, analysis of higher-mass subhalos in the ELVIS simulations — systems corresponding to  $10^{8.5} - 10^{9.5} M_{\odot}$  following the Garrison-Kimmel et al. (2014a) abundance matching relation — yields a very similar result ( $\tau_{\text{quench}} \sim 7.8$  Gyr; see also De Lucia et al. 2012). Altogether, we find that the quenching timescale increases with decreasing mass at  $\gtrsim 10^8 M_{\odot}$ , such that environmental processes are relatively inefficient in suppressing star formation at  $\sim 10^9 M_{\odot}$ . At lower stellar masses ( $\lesssim 10^8 M_{\odot}$ ), however, the quenching timescale decreases dramatically, indicating a stark increase in the satellite quenching efficiency for low-mass systems. Given that all known satellites of the Milky Way and M31 are quenched below stellar masses of  $\sim \text{a few} \times 10^5 M_{\odot}$ , the quenching timescale is expected to remain short to even lower masses. For the lowest-mass systems (i.e. ultra-faint dwarfs) environmental quenching is likely irrelevant, as the smallest dark matter halos may have their star formation suppressed significantly by non-environmental mechanisms like cosmic reionization (e.g. Rees 1986; Efsthathiou 1992; Wyithe & Loeb 2006; Oñorbe et al. 2015).

For an infalling satellite of any mass, assumed to be isolated from subsequent gas accretion, an upper limit on the longevity of any star formation activity is set by the gas supply within the satellite upon infall — including the satellite’s hot halo, from which gas may cool. Processes such as ram-pressure stripping, tidal effects, and harassment would then work to diminish this gas supply, so as to accelerate the quenching process. The long quenching timescales observed for higher-mass satellites ( $\sim 10^8 - 10^{10} M_{\odot}$ ) thus favor starvation as the dominant quenching mechanism. Ignoring future accretion, recycling, or gas cooling, the duration of star formation should follow the gas depletion timescale ( $\tau_{\text{depl}} = M_{\text{gas}}/\text{SFR}$ , where  $M_{\text{gas}}$  is the mass of a satellite’s cold gas reservoir upon infall). While studies of nearby star-forming galaxies suggest that the surface density of ongoing star formation is most closely tied to the dense  $\text{H}_2$  gas (Bigiel et al. 2008, 2011; Leroy et al. 2008), observed molecular depletion timescales are much shorter than the inferred quenching timescales. For relatively massive galaxies ( $M_{\star} \gtrsim 10^{10} M_{\odot}$ ), spanning a broad range of redshift, current observations point towards molecular depletion timescales of roughly  $1 - 2$  Gyr (Bigiel et al. 2011; Leroy et al. 2013; Saintonge et al. 2011; Tacconi et al. 2010, 2013; Genzel et al. 2010, 2015). Moreover, recent work targeting larger samples of lower-mass galaxies conclude that the molecular depletion timescale decreases with decreasing stellar mass, such that infalling satellites at  $10^9 M_{\odot}$  should, on average, exhaust their molecular reservoirs in less than 1 Gyr (Saintonge et al. 2011; Boselli et al. 2014).





**Figure 5.** The dependence of the satellite quenching timescale ( $\tau_{\text{quench}}$ ) on satellite stellar mass at  $M_{\star} \sim 10^6 - 10^{11} M_{\odot}$ . The inset legend lists the typical host halo mass for each data set. The magenta, gold, and cyan colored bands show the constraints for satellites in different mass host halos derived from analysis of galaxy groups and clusters in the SDSS by Wetzel et al. (2013). The black square denotes the typical quenching timescale for slightly lower-mass satellites ( $\sim 10^9 M_{\odot}$ ) from Wheeler et al. (2014), with the horizontal error bars denoting the 25 – 75% range in stellar mass probed by that work and the vertical error bars giving the variation in  $\tau_{\text{quench}}$  corresponding to satellite quenched fractions of 25 – 55% (as derived from analysis of subhalo populations in ELVIS, see §5). Our estimate for the low-mass satellite quenching timescale in the Local Group is given by the black circles ( $\tau_{\text{quench}} \sim 1.4$  Gyr, see §4.1), where our sample of subhalos is divided into two stellar mass bins of  $10^6 - 10^7 M_{\odot}$  and  $10^7 - 10^8 M_{\odot}$ . Each point gives the quenching timescale for subhalos in that mass bin assuming a satellite quenched fraction of 90%, while the vertical error bars illustrate the variation in  $\tau_{\text{quench}}$  corresponding to satellite quenched fractions of 80 – 100%. Our results are largely independent of mass, with increased scatter at higher masses due to the smaller number (and thus increased stochasticity in the infall times) of massive subhalos. The horizontal error bars show the quartiles of the subhalo mass distribution within each bin. At high masses, the quenching timescale increases with decreasing stellar mass. Below  $M_{\star} \sim 10^8 M_{\odot}$ , we find that satellite quenching becomes dramatically more efficient, suggesting a likely change in the physical mechanisms at play.

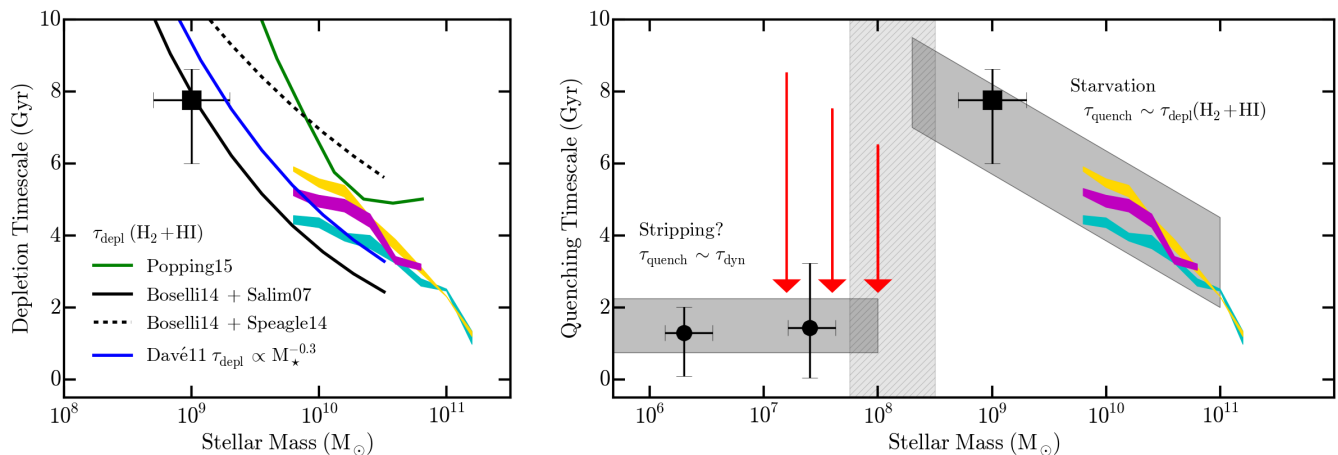
These relatively short molecular depletion timescales are in direct conflict with the derived quenching timescales for satellites at  $\sim 10^9 - 10^{10} M_{\odot}$ , which indicate that star formation typically continues for  $> 4$  Gyr after infall. Given that strong variation in star-formation efficiency for  $\text{H}_2$  is disfavored by both observations of the Kennicutt-Schmidt relation (Bigiel et al. 2008; Leroy et al. 2008; Krumholz & McKee 2005) and the bimodality of rest-frame galaxy colors (or specific star formation rates, Balogh et al. 2009; Weinmann et al. 2010),<sup>2</sup> the extended star formation in satellite galaxies seemingly must be driven by a fuel supply beyond the molecular reservoir. While gas cooling from a satellite’s hot halo could replenish the fuel supply for star formation, there is little observational evidence for significant hot gas coronae around star-forming galaxies, even at the high-mass regime and especially when restricting to gas with a sufficiently short cooling time (Sun et al. 2007; Rasmussen et al. 2009; Crain et al. 2010; Miller & Bregman 2015).

On the other hand, if atomic gas is included as a potential fuel source for star formation, there may be better agreement between the expected gas depletion and satellite quenching timescales. As previously stated, star formation

is observed to follow molecular gas density more closely than atomic gas density, although the weaker correlation with HI surface density could be largely due to atomic gas spanning a broader range of densities, thus on average tracing dense star-forming clouds more poorly (Glover & Clark 2012a,b; Clark et al. 2012). As shown in Figure 6a, when including the atomic gas component, the inferred quenching timescales at high stellar masses ( $M_{\star} \gtrsim 10^8 M_{\odot}$ ) are in broad agreement with the measured and predicted (atomic plus molecular) gas depletion timescales and thus consistent with a picture where satellite quenching is driven by starvation at high masses.

In the low-mass regime ( $\lesssim 10^8 M_{\odot}$ ), the atomic component becomes increasingly important, with observations of nearby star-forming systems finding that the typical HI gas fraction increases steadily towards lower mass (Skillman, Côté & Miller 2003; Geha et al. 2006; Leroy et al. 2008; Schiminovich et al. 2010; Catinella et al. 2010; Huang et al. 2012a). More importantly, this relative growth in atomic gas outpaces the corresponding increase in the average specific star formation rate (e.g. Noeske et al. 2007; Salim et al. 2007; Speagle et al. 2014), such that the gas depletion timescale is generally long ( $\gtrsim 10$  Gyr) for low-mass star-forming galaxies — a result that is fully consistent with the dearth of quiescent field galaxies at  $< 10^9 M_{\odot}$  (Geha et al. 2012; Oñorbe et al. 2015). In comparison to the typical quench-

<sup>2</sup> While a reduction in the star formation efficiency for satellites would yield longer depletion (and thus quenching) timescales, it would also lead to overproducing galaxies in the “green valley”.



**Figure 6.** *Left:* The dependence of the satellite quenching timescale ( $\tau_{\text{quench}}$ ) and gas depletion timescale ( $\tau_{\text{depl}}$ ) on satellite stellar mass at  $M_{\star} \gtrsim 10^8 M_{\odot}$ . The magenta, gold, and cyan colored bands show the constraints for satellites in different mass host halos derived from analysis of galaxy groups and clusters in the SDSS by Wetzel et al. (2013). The black square gives the typical quenching timescale for lower-mass satellites ( $\sim 10^{8.5} - 10^{9.5} M_{\odot}$ ) from Wheeler et al. (2014). The solid and dashed black lines show the mass dependence of the gas depletion timescale ( $\text{H}_2 + \text{HI}$ ) derived from combining the observed gas fractions of Boselli et al. (2014) with the star formation rate-stellar mass relation of Salim et al. (2007) and Speagle et al. (2014), respectively. The solid blue and solid green lines give the corresponding predictions from the hydrodynamical and semi-empirical models of Davé, Oppenheimer & Finlator (2011) and Popping, Behroozi & Peebles (2015), respectively. The inferred satellite quenching timescales at  $\gtrsim 10^8 M_{\odot}$  show broad agreement with the observed gas depletion timescales for field systems, suggesting that starvation is the main driver of satellite quenching at these masses. *Right:* The dependence of the satellite quenching timescale on satellite stellar mass at  $10^6 < M_{\star}/M_{\odot} < 10^{11}$ , including the estimates of  $\tau_{\text{quench}}$  from Fig. 5. The light grey shaded regions highlight the expected dominant quenching mechanism and projected mass dependence of  $\tau_{\text{quench}}$ , while the vertical hashed grey region demarcates the critical mass scale below which a more efficient physical process (e.g. ram-pressure stripping) drives relatively rapid quenching in the Local Group. At high masses ( $\gtrsim 10^8 M_{\odot}$ ), satellite quenching is consistent with being driven by starvation; below  $\sim 10^8 M_{\odot}$ , however, we posit that ram-pressure stripping becomes increasingly effective, such that depletion (and thus quenching) timescales are significantly shorter.

ing timescales at low masses, however, the overabundance of atomic gas in these systems suggests that starvation alone cannot drive satellite quenching at  $\lesssim 10^8 M_{\odot}$ . Instead, a secondary process, such as ram-pressure stripping, is needed to remove cool gas from infalling systems, thereby decreasing the gas depletion timescale significantly.

As observations of systems such as Leo I suggest (Sohn et al. 2013), we would expect stripping to occur as a satellite infalls from  $R_{\text{vir}}$  towards pericenter, such that quenching at low masses would proceed according to the dynamical time ( $\tau_{\text{dyn}}$ ), which is roughly 1–2 Gyr for a Milky Way-like halo today (e.g. Stewart et al. 2009). While a detailed analysis of the efficacy of ram-pressure stripping (or other quenching mechanisms such as tidal forces) is beyond the scope of this work (see Fillingham et al. in prep), Figure 6b presents a qualitative depiction of the potential physical processes at play, where satellite quenching is primarily driven by starvation at higher masses ( $\gtrsim 10^8 M_{\odot}$  for the Local Group) while ram-pressure stripping becomes increasingly effective below a critical mass scale of  $\sim 10^8 M_{\odot}$ , yielding shorter quenching timescales and high observed quenched fractions. In this qualitative picture, we expect that ram-pressure stripping (or some other efficient quenching mechanism) plays a critical role in suppressing star formation down to the mass scales at which reionization inhibits gas accretion and therefore star formation (perhaps,  $M_{\star} \sim 10^4 M_{\odot}$ , Brown et al. 2014; Oñorbe et al. 2015; Wheeler et al. 2015). Applying this qualitative picture to higher-mass host halos (e.g. rich groups and clusters), the critical scale at which ram-pressure stripping becomes effective should move to higher satellite

masses due to the increased density of the host’s hot halo in concert with higher infall velocities for the satellites.

## 6 SUMMARY

Comparing observations of the Local Group satellite population to corresponding high-resolution  $N$ -body simulations from the ELVIS suite, we investigate the typical timescale (following infall) upon which star formation is suppressed in low-mass ( $< 10^8 M_{\odot}$ ) satellite galaxies. When combining our work with complementary analyses of higher-mass satellite populations, we present a comprehensive picture of satellite quenching spanning roughly five orders of magnitude in stellar mass. Our principal results are as follows:

- To reproduce the high fraction of quenched low-mass satellites in the Local Group, the typical quenching timescale at  $M_{\star} \sim 10^6 - 10^8 M_{\odot}$  must be relatively short ( $\tau_{\text{quench}} \sim 2$  Gyr).
- The longer quenching timescales inferred for higher-mass satellites are roughly consistent with the observed cold gas ( $\text{H}_2$  plus  $\text{HI}$ ) depletion timescales in corresponding field systems; this suggests that satellite quenching is largely driven by starvation at  $M_{\star} \gtrsim 10^8 M_{\odot}$ .
- At low masses ( $M_{\star} \lesssim 10^8 M_{\odot}$ ), the much shorter satellite quenching timescales are potentially set by ram-pressure stripping or some other process that removes the cold gas

reservoirs in infalling systems, such that nearly all low-mass satellites in the Local Group (or more massive host halos) are quenched.

- If ram-pressure stripping is responsible for increasing the satellite quenching efficiency at low masses ( $M_{\star} \lesssim 10^8 M_{\odot}$ ) within the Local Group, we expect that this critical quenching scale will shift to higher satellite masses in higher-mass host halos (e.g. rich groups and clusters) due to the increased density of the host's hot halo in concert with higher infall velocities for the satellites.

We have suggested here that the transition to rapid quenching in the smallest satellites ( $M_{\star} \lesssim 10^8 M_{\odot}$ ) is potentially explained as a ram-pressure scale (see Figure 6 and Fillingham et al. in prep). Above this scale, satellite quenching occurs on a gas depletion time, as accreted galaxies are starved of fresh fuel for star formation. Below this scale, a more rapid, drag-induced stripping becomes dominant, such that the gravitational restoring force of lower-mass dark matter halos can no longer retain cool gas against the fluid pressure they experience while in orbit around a larger host.

The idea that Local Group dwarfs are quenched via ram pressure has been around for some time (e.g. Einasto et al. 1974; Lin & Faber 1983; Blitz & Robishaw 2000) but the evidence for very rapid quenching provided here adds further argument in its favor, and identifies a characteristic mass scale where it may become dominant. Underlying this hypothesis is the notion that the Milky Way and M31 both harbor extended ( $\sim 150$  kpc) reservoirs of hot baryons with sufficient density to strip small galaxies of their gas. There is evidence that M31 does indeed host an extended circumgalactic medium (Lehner, Howk & Wakker 2015), and X-ray studies are consistent with this possibility around the Milky Way (Fang, Bullock & Boylan-Kolchin 2013; Miller & Bregman 2015; Fang et al. 2015). If extended low-density coronae of this kind are common around  $\sim L^*$  galaxies, then they may prove important for understanding the baryon cycle and the global census of baryons in the Universe.

Of course, one shortcoming of our analysis is that we have restricted ourselves to satellites of M31 and the Milky Way. Given completeness issues, the Local Group is a reasonable starting point, but it remains possible that satellites in our vicinity are unusually quenched, perhaps owing to an uncommonly dense distribution of hot baryons in our vicinity, or atypically early infall times for the Local Group satellite population. Ongoing efforts to discover and characterize dwarf satellite systems around other massive hosts will be important for solidifying the existence of a sharp transition scale at  $M_{\star} \lesssim 10^8 M_{\odot}$ , where satellite quenching appears to become extremely efficient in the Local Group. If future studies reveal that the Local Group is indeed typical, then pinpointing the origin of this quenching scale and testing the ram-pressure hypothesis will be of utmost importance.

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