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Superconductivity in Dilute Hydrides of Ammonia under Pressure

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Abstract

In the last decade, there has been great progress in predicting and synthesizing polyhydrides that exhibit superconductivity when squeezed. Dopants allow these compounds to become metals at pressures lower than those required to metallize elemental hydrogen. Here, we show that by combining the fundamental planetary building blocks of molecular hydrogen and ammonia, conventional superconducting compounds can be formed at high pressure. Through extensive theoretical calculations we predict metallic metastable structures with NH_n ($n = 10, 11, 24$) stoichiometries that are based on NH_4^+ superalkali cations and complex hydrogenic lattices. The hydrogen atoms in the molecular cation contribute to the superconducting mechanism, and the estimated superconducting critical temperatures, T_c s, are comparable to the highest values computed for the alkali metal polyhydrides. The largest calculated (isotropic Eliashberg) T_c is 179 K for $Pnma\text{-NH}_{10}$ at 300 GPa. Our results suggest that other molecular cations can be mixed with hydrogen under pressure yielding superconducting compounds.

The long-hypothesized phase of metallic hydrogen¹ tantalizes with exotic properties including high-temperature superconductivity² and superfluidity.³ Though numerous groups have attempted to create this sought-after phase, reports of metallicity in statically compressed samples remain controversial,⁴ and superconductivity has not been measured. However, a route towards lowering the metallization pressure, based upon the addition of an impurity element that can “chemically precompress” the hydrogen lattice^{5,6} has set the stage for the prediction and discovery of a series of hydrogen-based conventional superconductors at pressures found in Earth’s core. Most of the studied compounds are covalent metals that contain hydrogen in combination with an electropositive metal or a *p*-block element.⁷ In this paper, we theoretically investigate the structure and superconducting properties of NH_n ($n = 7 - 24$) polyhydrides that are based on hydrogen and ammonia molecules, which are building blocks of giant planets.

Ab initio crystal structure prediction (CSP) studies have proposed a plethora of stable or low-lying metastable phases formed from doping hydrogen by a Group-I element under pressure.⁸ Numerous exotic structural motifs such as stretched $\text{H}_2^{\delta-}$ molecules (*e.g.* in LiH_6 ⁹), hydridic H^- ions (*e.g.* in NaH_9 ¹⁰), linear or ever-so-slightly bent H_3^- units (*e.g.* in KH_5 ,¹¹ RbH_5 ¹² and CsH_3 ¹³) and $(\text{H}_3^-)_\infty$ polymeric chains (*e.g.* in RbH_6 ¹²) are present in the CSP-found phases. The trends in their computed superconducting critical temperatures, $T_{c,s}$, are related to their structural features, with compounds containing H^- and H_3^- units typically being poor candidates for superconductivity, those accommodating $\text{H}_2^{\delta-}$ units exhibiting intermediate $T_{c,s}$, and phases with non-molecular hydrogenic motifs showing the most promise.¹⁴ The computed $T_{c,s}$ of polyhydrides of the alkaline earths are typically higher than those of the alkali metals.⁷ Nonetheless, various superconducting Group-I high hydrides have been predicted including: $R\bar{3}m$ - LiH_6 ($T_c = 38$ K at 150 GPa and 82 K at 300 GPa,¹⁵ or 130-160 K at 300 GPa¹⁶), $I422$ - LiH_8 (31 K at 100 GPa¹⁵), $C2/c$ - KH_6 (~ 70 K at 166 GPa¹⁷), $Immm$ - KH_{10} (140 K at 150 GPa),¹⁸ and $C2/m$ - RbH_{12} (133 K at 150 GPa),¹⁸ $Pm\bar{3}m$ - NaH_6 (260 K at 100 GPa) and $P6_3/mmc$ - NaH_9 (252 K at 500 GPa).¹⁶

Encouragingly, lithium polyhydrides have been synthesized under pressure (though their structure has not been determined),¹⁹ both NaH_7 and NaH_3 have been made below a megabar,^{20,21} and

recently the synthesis of cesium and rubidium polyhydrides has been reported.²² Density functional theory (DFT) calculations predict the stabilization pressure of these systems to decrease with decreasing ionization potential (IP) of the metal, so whereas the enthalpy of formation, ΔH_F , of the lithium polyhydrides becomes negative around 120 GPa,²³ the cesium polyhydrides are stable by a mere 2 GPa.¹³ Quasi-spherical molecular complexes, such as $\text{Li}(\text{NH}_3)_4$, possessing frontier molecular orbitals with the same symmetry as those of the Group-I elements, but extending outside the component atoms, can be thought of as superalkali atoms.²⁴ We therefore wondered if other molecular species, with superatom characteristics resembling the Group-I elements, could be combined with hydrogen under pressure to access polyhydrides with unique chemical compositions? And, would the hydrogen atoms comprising the molecular complex help enhance their superconducting properties?

The NH_4 radical is a Rydberg molecule, whose singly occupied molecular orbital (SOMO) is extremely diffuse. The MOs of this tetrahedral molecule are reminiscent of an alkali metal atom's, with the nearly spherical $1a_1$ SOMO resembling the s -orbital, and the higher lying unoccupied triply degenerate $1t_2$ orbitals being of p -type symmetry (Figure S2). Comparison of the radial distribution functions of these orbitals shows that the profile computed for NH_4 resembles most closely that of K or Rb (Figure S3). The experimentally measured IP of NH_4 , 4.62 eV,²⁵ falls in-between the IPs of Na and K (5.14 and 4.34 eV, respectively²⁶), and is somewhat higher than the IP of Rb (4.18 eV²⁶). Thus, based upon its IP and the radial extent of its SOMO, NH_4 is analogous to Rb or K.

In the gas phase NH_4 is a short-lived metastable species that readily decomposes into NH_3 and a hydrogen radical. In condensed phases and above pressures of ~ 90 GPa computations predict that NH_3 self-ionizes into phases containing NH_4^+ and NH_2^- ions, driven by a volume decrease.²⁷ Moreover, CSP searches predict that the ammonium cation is found in a number of stable high pressure phases including $\text{NH}_4^+ \cdots \text{OH}^-$ by 5 GPa²⁸ and $(\text{NH}_4^+)_2 \cdots \text{O}^{2-}$ by 65 GPa,²⁹ and in NH_3 -HF mixtures of varying composition.³⁰ Support for the existence of NH_4^+ in a subset of these systems has subsequently been obtained in high-pressure experiments.^{31,32} Especially relevant for

the work presented here, CSP studies on NH_n ($n = 1 - 9$) found that NH_7 has the most negative ΔH_{F} (from H_2 and NH_3) between 50-200 GPa.³³ Below 60 GPa the most stable NH_7 phase (of $R\bar{3}m$ symmetry) contained $\text{NH}_3\text{-H-NH}_3$, as well as H_2 and H^- units, and at higher pressures a $P4_12_12$ structure, which remained insulating to at least 100 GPa, with tetrahedral NH_4^+ , H_2 and H^- units was preferred instead.

Given the relative ease of formation of the ammonium ion under pressure and its similarity with that of an alkali-metal cation, we hypothesized that superconducting ammonium polyhydrides, with NH_n ($n \geq 5$) stoichiometries, might be (meta)stable under pressure. Therefore, CSP searches were performed (using the XTALOPT evolutionary algorithm version 12^{34,35}) for $n = 3, 5 - 12, 16, 20$ at 100, 200 and 300 GPa with the projector-augmented wave³⁶ DFT³⁷ framework (see Section S1 for further details). The many predicted low enthalpy structures were filtered, and only those that were metallic, high-symmetry, as well as frozen-phonon-dynamically and either meta- or thermodynamically stable were considered for further analysis and T_{c} calculations. Thermodynamically stable structures lay on the (0 K) convex hull generated using NH_3 and H_2 as the endpoints, and metastable ones were defined as those within 50 meV/atom of the hull (Figure S1E-F).

An intriguing phase that emerged in our CSP searches, $Pnma\text{-NH}_{10}$ (Figure 1A), was dynamically stable above 280 GPa (Figure S4A). Though it was unstable with respect to decomposition into H_2 and NH_3 ($\Delta H_{\text{F}} \sim 23$ meV/atom at 300 GPa within the PBE functional³⁸), inclusion of dispersion³⁹ lowered the magnitude of the instability ($\Delta H_{\text{F}} \sim 7.5$ meV/atom), and the zero-point energy (ZPE) corrections were further stabilizing, resulting in an overall ΔH_{F} of ~ 6.5 meV/atom. This phase can be described as sheets of NH_4^+ molecules that lie in the bc plane and are separated by one-dimensional (1D) zig-zag hydrogenic chains that run along the c -axis. H_2 units (labelled as H(5) and H(6)) whose bond lengths are slightly stretched relative to those in the most stable H_2 phases at 1 atm ($P6_3/m$) or 300 GPa ($Cmca\text{-H}_2$), and whose Bader charges assign them as being nearly neutral, and hydridic hydrogens (H(4)), with a Bader charge of -0.15 comprise these chains.

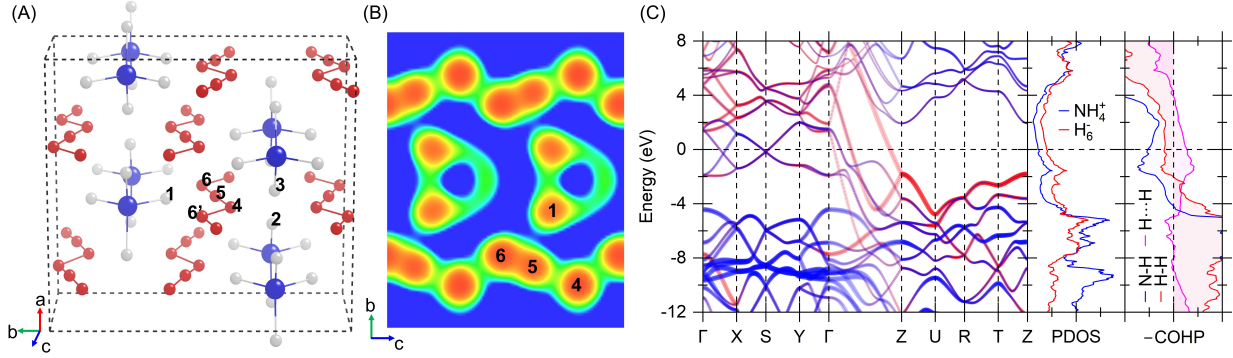


Figure 1: $Pnma$ - NH_{10} at 300 GPa: (A) Unit cell. N atoms are blue, H atoms in NH_4^+ are white and the H atoms comprising the 1D chains are red. (B) Plot of the Electron Localization Function (ELF) using an isosurface value of 0.5, with the contours colored from 0.5 (blue) to 1.0 (red). (C) Atom projected band structure and density of states showing contributions from atoms within the NH_4^+ units (blue) and atoms within the 1D chains (red). $-\text{COHP}$ averaged over N-H bonds within NH_4^+ units (blue), H-H bonds within the 1-D chain (red), and between the H atoms within NH_4^+ and those in the chain (pink).

The crystal orbital Hamilton population integrated to the Fermi level, E_F , ($-\text{ICOHP}$),^{40,41} which gives a measure of the covalent bond strength, was smaller for the dihydrogen units within $Pnma$ - NH_{10} than for the H-H bonds within $P6_3/m$ and $Cmca$ - H_2 (Table 1). Visualization of the electron localization function (ELF) (Figure 1B) revealed an even weaker covalent $\text{H-H}\cdots\text{H}^-$ interaction. Nonetheless, these three hydrogens can be thought to comprise an asymmetric H_3^- molecule, whose bond lengths are more equalized than in the gas phase local minimum, predicted to possess one short (0.75 Å) and one long (2.84 Å) H-H distance via *ab initio* calculations.⁴² The gas phase species possesses a double well minimum, with a symmetric H_3^- transition state, corresponding to the simplest example of a three centered four electron bond. Previous CSP calculations revealed such H_3^- motifs, which were symmetrized under pressure, within predicted RbH_3 , RbH_5 ¹² and KH_5 ¹¹ phases. Within $Pnma$ - NH_{10} the nearest neighbor intermolecular H_3^- distances are only somewhat longer than the longest intramolecular distance (*cf.* H(4)-H(6') vs. H(4)-H(5)), suggesting that it may be more appropriate to view these hydrogenic motifs as a series of parallel ${}^1_\infty[\text{H}_2\cdots\text{H}^-]$ chains, resembling those predicted in $Imma$ - RbH_6 ¹² and $R\bar{3}m$ - SrH_6 ⁴³ phases near 250 GPa.

Because the NH_4 molecule can be thought of as a superalkali metal atom whose properties

Table 1: H-H distances in the NH_n ($n = 10, 11, 24$) phases, along with the corresponding crystal orbital Hamiltonian populations integrated to the Fermi level (-ICOHP) at 300 GPa. Values are also provided for solid H_2 at 1 atm ($P6_3/m$) and 300 GPa ($Cmca$). The hydrogen atom numbering scheme is provided in Figures 1, 2 and 3. H_{mol} , H_{chain} and H_{ammon} label hydrogens comprising molecular H_2 , 1D-chain, or ammonium cation units, respectively.

bond type	distance (Å)	-ICOHP (eV/bond)
<i>P6₃/m</i> -H ₂ (1 atm)		
H _{mol} -H _{mol}	0.750	5.51
H _{mol} -H _{mol}	0.751	5.51
<i>Cmca</i> -H ₂ (300 GPa)		
H _{mol} -H _{mol}	0.756	6.15
H _{mol} -H _{mol}	0.775	5.78
H _{mol} -H _{mol}	1.100	1.45
H _{mol} -H _{mol}	1.131	1.23
H _{mol} -H _{mol}	1.146	1.18
<i>Pnma</i> -NH ₁₀ (300 GPa)		
H5 _{chain} -H6 _{chain}	0.812	4.71
H4 _{chain} -H5 _{chain}	0.928	3.10
H4 _{chain} -H6' _{chain}	0.965	2.90
H1 _{ammon} -H5 _{chain}	1.047	1.92
H3 _{ammon} -H4 _{chain}	1.136	1.38
H2 _{ammon} -H4 _{chain}	1.173	1.19
<i>Cmc2₁</i> -NH ₁₁ (300 GPa)		
H5 _{chain} -H6 _{chain}	0.785	5.63
H3 _{chain} -H4 _{chain}	0.935	3.27
H4 _{chain} -H5 _{chain}	0.955	3.04
H1 _{ammon} -H3 _{chain}	1.067	1.51
H4 _{chain} -H6' _{chain}	1.100	1.60
H1 _{ammon} -H3' _{chain}	1.146	1.00
H2 _{ammon} -H6 _{chain}	1.155	1.22
<i>C2</i> -NH ₂₄ (300 GPa)		
H9 _{mol} -H9' _{mol}	0.733	6.86
H10 _{mol} -H11 _{mol}	0.744	6.38
H7 _{mol} -H8 _{mol}	0.764	5.94
H5 _{chain} -H6 _{chain}	0.797	5.14
H2 _{chain} -H3 _{chain}	0.842	4.60
H2 _{chain} -H4 _{chain}	0.961	2.87
H3 _{chain} -H4 _{chain}	0.969	2.83
H4 _{chain} -H5 _{chain}	0.975	2.68
H6 _{chain} -H8 _{mol}	1.029	1.96
H1 _{ammon} -H5 _{chain}	1.114	1.13
H3 _{chain} -H7 _{mol}	1.148	1.14
H1 _{ammon} -H6 _{chain}	1.154	0.93

resemble those of K or Rb, the formula of $Pnma\text{-NH}_{10}$ can be written as $(\text{NH}_4)\text{H}_6$, in further analogy with the previously predicted RbH_6 and SrH_6 phases. However, in stark contrast to the bonding that is observed within these “simple” hexahydrides, weak covalent bonds are formed between the hydrogen atoms in the NH_4^+ superalkali ion, and the hydrogens within the hydrogenic chains. As Table 1 reveals these interactions (H(1)-H(5), H(2)-H(4) and H(3)-H(4)), though small, are not negligible, and are of similar magnitude as those between the second nearest neighbor hydrogens within $Cmca\text{-H}_2$ at 300 GPa.

RbH_6 , SrH_6 and $Pnma\text{-NH}_{10}$ are all good metals with a high density of states (DOS) at E_F . Their metallicity is in part due to the formation of 1D hydrogenic chains, whose antibonding levels are partially filled via donation from the electropositive alkali or superalkali atom. However, the weak interactions formed between the hydrogen atoms within NH_4^+ and those comprising the ${}^1_{\infty}[\text{H}_2 \cdots \text{H}^-]$ chains in $Pnma\text{-NH}_{10}$ ensures that character from hydrogens in both the anionic and cationic sublattices are found at E_F , as is evident in the projected DOS (Figure 1C). Moreover, plots of the -ICOHP averaged over the N-H bonds within NH_4^+ (blue curve) and over all H-H pairs within the 1D zig-zag chains (red curve) show that all of the bonding levels are filled, and the antibonding levels are partially occupied at E_F . The bonding states between the hydrogens comprising NH_4^+ and the hydrogen chain (pink curve), however, are not completely filled. These weak but numerous covalent interactions may prevent the melting of the 1D chain. Our molecular dynamics (MD) simulations confirmed the thermal stability of $Pnma\text{-NH}_{10}$ at 150 K and 300 GPa. (Figure S6A) Although at 200 K some protons started to diffuse to nearby empty lattice sites, the chains do not melt until 450 K, greatly exceeding the T_c calculated for this compound, whereas the ${}^1_{\infty}[\text{H}_3^-]$ chains within RbH_6 exhibited liquid-like behavior at comparable temperatures.¹²

Another CSP-found phase that satisfied our criteria for further analysis, $Cmc2_1\text{-NH}_{11}$, was computed to be slightly preferred over its neutral molecular constituents ($\Delta H_F = -7.8$ meV/atom at 300 GPa, with dispersion and ZPE corrections), and it was dynamically stable above 280 GPa (Figure S4B). The NH_4^+ superalkali cation was also a key constituent of this phase, with neighboring cations stacked along the c -axis (Figure 2A). Along the a and b axes the NH_4^+ units were

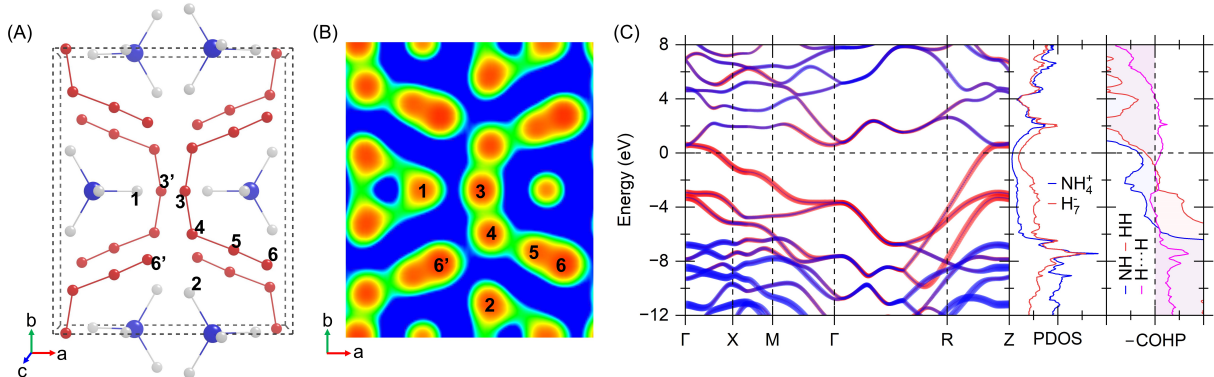


Figure 2: $Cmc2_1$ - NH_{11} at 300 GPa: (A) Standard conventional unit cell. (B) ELF profile in a plane that cuts through the H_7^- molecular fragments, colored using the same scheme as in Figure 1. (C) Atom projected band structure and density of states showing contributions from atoms within the NH_4^+ units (blue) and atoms comprising the H_7^- motifs (red). $-COHP$ averaged over N-H bonds within NH_4^+ (blue), H-H bonds within the H_7^- units (red), and weak H1-H3, H4-H6', H1-H3', and H2-H6 interactions as in Table 1 (pink).

separated from each other via H_7^- molecular fragments, which can be readily identified in a 2D plot of the ELF that passes through them (Figure 2B). This non-planar H_7^- C_s symmetry anion contained two tightly bonded H_2 units (H(5)-H(6)) with Bader charges of -0.10. The H(5)-H(4) and H(4)-H(3) distances were ~ 0.16 Å longer than in these strongly bonded dihydrogen units resulting in interactions that were 2.5 eV weaker, on average (Table 1). The Bader charge of H(4) was +0.08, and that of H(3) was -0.12. This charge distribution differs from the one computed for an H_7^- anion in the gas phase because of the numerous bonds present between the hydrogens comprising the superalkali cation and the C-shaped H_7^- anion. However, within $Cmc2_1$ - NH_{11} these interactions are even weaker than in $Pnma$ - NH_{10} , as evidenced by the longer bond distances, smaller $-ICOHPs$, as well as a decreased NH_4^+ projected DOS and $-COHP$ between the respective hydrogen atoms at E_F (Figure 2C). MD simulations revealed that $Cmc2_1$ - NH_{11} is thermally stable at 300 GPa and 100 K, but by 150 K the NH_4^+ molecules rotate as the N-H bonds vibrate (Figure S6B).

Finally, we describe the structural peculiarities of the most dilute hydride of ammonia that was analyzed, $C2$ - NH_{24} . Prior high pressure experiments have synthesized compounds with high weight percent H_2 content including $HI(H_2)_{13}$ (17.7%)⁴⁴ and $(H_2)_8CH_4$ (33.4%).⁴⁵ The phase pre-

dicted here contains an impressive 52.6 or 55.3 weight % H_2 (when NH_4 or NH_3 are treated as indivisible entities, respectively). It is dynamically stable from 200-380 GPa (Figure S4C), with a ΔH_F of 3.8 and 5.3 meV/atom at 200 and 300 GPa, respectively, when ZPE and dispersion corrections are included. MD simulations showed that $C2-NH_{24}$ is stable up to 150 K at 300 GPa (Figure S6C).

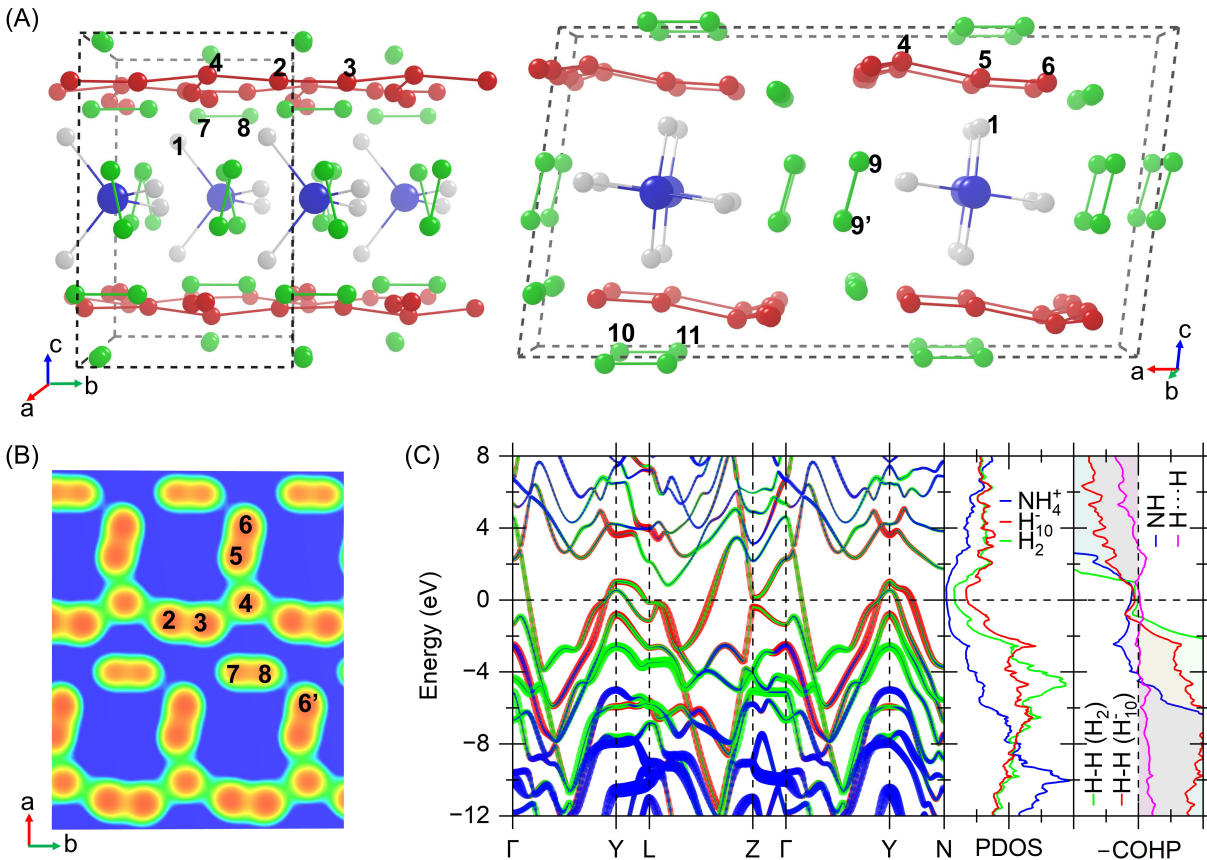


Figure 3: $C2-NH_{24}$ at 300 GPa: (A) Standard conventional unit cell from (left) front view and (right) side view. (B) ELF profile in a plane that cuts through the 1D chain, colored using the same scheme as in Figure 1. (C) Atom projected band structure and density of states showing contributions from atoms within the NH_4^+ units (blue), atoms within the H_2 units (green), and atoms within the 1D chains (red). $-COHP$ averaged over N-H bonds within NH_4^+ (blue), H-H bonds within the H_2 units (green), H-H bonds within the 1D chain (red), and between the H atoms within NH_4^+ or H_2 and those in the chain (H6-H8, H1-H5, H3-H7, and H1-H6 as in Table 1, pink).

The standard conventional cell of $C2-NH_{24}$ (Figure 3A) contains two NH_4^+ molecules stacked parallel to the b -axis. In the ab plane these rows of ammonium cations are separated from one

another by two 1D polymeric hydrogen motifs (H(2) to H(6)) that themselves are separated by H₂ molecules (H(10)-H(11)), and a further set of H₂ molecules separate the superalkali cations in the *bc* plane (H(9)-H(9')). Therefore, another way to write the formula of this phase is NH₄(H₁₀)(H₂)₅. Turning to the 1D chain, whose ELF profile is plotted in Figure 3B, reveals pairs of hydrogens that form shorter stronger bonds (H(5)-H(6) and H(2)-H(3)) surrounding a nearly neutral lone hydrogen atom (H(4)). The bond distances and strengths of the ensuing interactions are provided in Table 1. Similar to the previously discussed phases, weak but numerous interactions between the hydrogen atoms within NH₄⁺ and those within the anionic sublattice are found. However, *C2*-NH₂₄ is predicted to be dynamically stable within the harmonic approximation to pressures as low as 200 GPa, whereas *Pnma*-NH₁₀ required at least 280 GPa to prevent the onset of a dynamic instability. The lowering of the stabilization pressure might be attributed to the dilution of the negative charge on the extended hydrogenic framework from H₆⁻ in *Pnma*-NH₁₀ to H₁₀⁻ in *C2*-NH₂₄. Thus, we postulate that phases containing an even larger weight percent hydrogen, with extensive polymeric hydrogen chains, could be (meta)stable at lower pressures, in particular when anharmonic fluctuations, likely to be important for these hydrogen atoms, are considered.

As the band structure and DOS plots at 300 GPa show (Figure 3A), *C2*-NH₂₄ is metallic in the whole pressure range of its stability. Here, a *little* bit of nitrogen has really done much for hydrogen!⁹ The formal charge on the anionic lattice in *C2*-NH₂₄ (assuming full transfer of the unpaired electron in NH₄ to the hydrogenic lattice) is H₂₀⁻ as compared to H₆⁻ for NH₁₀. As a result, there are fewer antibonding states filled within NH₂₄, which will decrease further in the limit of infinite dilution.

Let us now explore the propensity for superconductivity in these superhydride phases and identify the vibrations that furnish the greatest contribution to their predicted T_c s. To do so, we calculated the Eliashberg spectral function, $\alpha^2F(\omega)$, and from it we obtained the electron-phonon-coupling (EPC) parameter, λ , as well as the logarithmic average of the phonon frequencies, ω_{\log} . The combined nuclear quantum and anharmonic effects, which could be significant in these systems, were neglected, the Coulomb parameter, μ^* , was set to 0.1, and T_c was estimated via the

Allen-Dynes modified McMillan (ADM) equation (including the strong coupling and shape dependence),⁴⁶ and within isotropic Eliashberg theory.⁴⁷

One way to identify the phonons most important for increasing the T_c in a conventional superconductor is by calculating its functional derivative with respect to the Eliashberg spectral function, $\delta T_c / \delta \alpha^2 F$.^{46,48} Such an analysis has previously been applied to compressed hydrides, revealing the important role of bending vibrations at low pressures, while at high pressures mixed bending and stretching vibrations spanning a wide range of frequencies contributed to the EPC.⁴⁹ Using the ADM, we calculated the functional derivative not only of T_c , but also of λ and ω_{\log} numerically using a frequency interval of 1 cm^{-1} . The results for the three studied superhydrides at 300 GPa are plotted in Figure 4 (red curves). An additional analysis was performed where $\alpha^2 F(\omega)$ was set to zero in a frequency range between ω to $\omega + \Delta\omega$, and the superconducting properties were re-computed. The change induced by this perturbation is plotted as $\Delta\lambda$, $\Delta\omega_{\log}$ and ΔT_c in the same figures (blue curves).

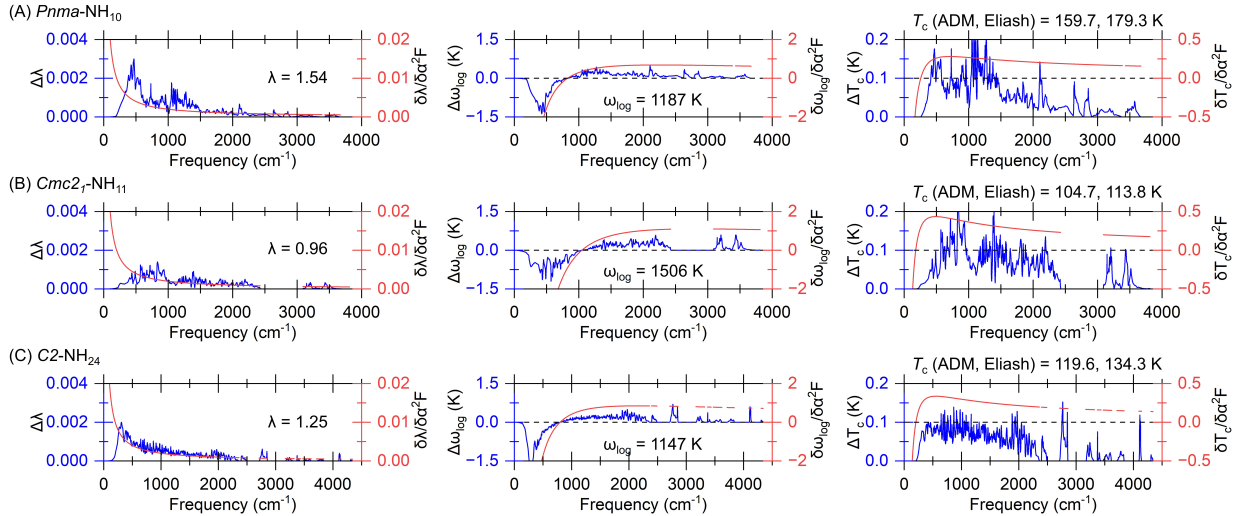


Figure 4: Superconducting properties of (a) $Pnma\text{-NH}_{10}$, (b) $Cmc2_1\text{-NH}_{11}$, and (c) $C2\text{-NH}_{24}$ at 300 GPa. Contribution of the vibrational frequencies to the electron-phonon-coupling ($\Delta\lambda$), logarithmic average frequency ($\Delta\omega_{\log}$) and superconducting critical temperature (ΔT_c) as a function of frequency (blue curves). The functional derivatives with respect to the Eliashberg spectral function, $\delta\lambda/\delta\alpha^2 F$, $\delta\omega_{\log}/\delta\alpha^2 F$ and $\delta T_c/\delta\alpha^2 F$, are also plotted (red curves). The T_c predicted with the Allen-Dynes modified McMillan (ADM) and Eliashberg equations and $\mu^* = 0.1$ is provided.

The plot of $\Delta\lambda$ vs. ω for $Pnma\text{-NH}_{10}$ showed that modes whose frequencies were below

490 cm^{-1} contributed most to λ . Further analysis revealed that the mode with the strongest EPC, *i.e.* largest $\lambda_{q\nu}$, is found at the first phonon branch at the X point with a frequency of 180 cm^{-1} . Visualization of this mode showed that it is an H(4)-H(5)-H(6) stretch coupled with an NH_4^+ rotation. The former, which corresponds to vibrations of the 1D chain, could result in the formation of H_2/H^- units with the concomitant opening of a gap near E_F along the Γ - Z high symmetry line (Figure 1). The bands that cross E_F along Γ - X - S - Y - Γ exhibit character from both the anionic and cationic sublattices, as evident in a plot of the partial charge density calculated within ± 0.15 eV of the Fermi level. This plot showed the H_6^- chains enveloped by tubes connected via charge density arising from the ammonium cations, by way of a rhombus formed from H(2), H(3), and two H(4)s from two neighbouring H_6^- chains (Figure S7). Rotation of the NH_4^+ elongates the distances within the H_4 rhombus, and perturbs the numerous weak interactions between the hydrogens comprising the anionic and cationic lattices, resulting in the EPC.

Visualization of a subset of (randomly chosen) phonon branches shows they resemble the mode with the largest $\lambda_{q\nu}$, therefore one might expect that the phonon linewidth ($\gamma_{q\nu}$) for most of the phonon modes would be similar, as confirmed by a plot of the phonon band structure decorated by the $\gamma_{q\nu}$ (Figure S8). Despite their large and positive contribution towards the overall λ , the low frequency modes below ~ 500 cm^{-1} decrease the overall ω_{\log} (and are thus shown with a negative sign in the middle panel in Figure 4A), and their overall contribution to T_c (right panel) is small. In fact setting the $\alpha^2F(\omega)$ values below 490 cm^{-1} to zero decreases the T_c by only 16 (24) K in the ADM (Eliashberg) approximation. A similar analysis for the phonon branches above 2330 cm^{-1} , which are linear combinations of N-H and H-H stretching modes, yields a drop of the ADM (Eliashberg) T_c by only 23 (28) K. Therefore, similar to Ref.,⁴⁹ we find that the intermediate frequency regime from 490 to 2300 cm^{-1} contributes most to the superconductivity. The maximum of the functional derivative, $\delta T_c / \delta \alpha^2 F$, which defines the optimal frequency (ω_{opt}), can be used to estimate the T_c as $\omega_{\text{opt}} \sim 7k_B T_c$,⁴⁸ where k_B is the Boltzmann constant, yielding ~ 169 K, which is close to our Allen-Dynes value of 160 K, and slightly smaller than our Eliashberg value of 180 K.

$Cmc2_1$ - NH_{11} has a smaller DOS at E_F (0.069 states/Ry/ \AA^3) as compared to $Pnma$ - NH_{10}

(0.108 states/Ry/Å³), and a plot of the band structure decorated by the $\gamma_{q\nu}$ (Figure S9) illustrates that only a few of the phonon modes have large linewidths, in contrast to the decahydride. Indeed, Figure 4B illustrates that $Cmc2_1$ -NH₁₁ possesses fewer low frequency modes with a high EPC, resulting in a smaller total λ , but a larger ω_{\log} . The ADM and Eliashberg predicted T_c values are similar (105 vs. 114 K). The lower critical temperature of NH₁₁ (as compared to NH₁₀) is not surprising as it possesses a molecular hydrogenic sublattice, in agreement with the trends in T_c that have been observed for other high pressure polyhydrides with electropositive elements.¹⁴ Finally, $C2$ -NH₂₄ with both molecular and 1D-periodic hydrogenic motifs falls somewhere between NH₁₀ and NH₁₁ in terms of the three functional derivatives, the way in which certain frequencies contribute to the integrated superconducting parameters, and the DOS at E_F (0.064 states/Ry/Å³). Therefore, it should not be a surprise that its estimated T_c is also intermediate, with a 120 K ADM (134 K Eliashberg) value.

In conclusion, evolutionary structure searches predicted a number of metallic, high-symmetry NH_{*n*} phases that were metastable under pressure. These compounds, containing ammonium cations and a wide variety of hydrogenic lattices, including 1D hydrogenic chains, dihydrogen, and more complex hydrogenic molecular units, were predicted to be conventional superconductors. Because the quasi-spherical NH₄⁺ molecule can be thought of as a superalkali atom, with a radius and ionization potential resembling that of K or Rb, these compounds can be seen as extensions of the alkali metal polyhydrides. The studied NH₁₀, NH₁₁ and NH₂₄ compounds, whose formulae could be written as (NH₄)(H₆), (NH₄)(H₇) and NH₄(H₂)₅H₁₀, afforded weak bonds between the hydrogens comprising the NH₄⁺ cation and those in the anionic lattice, with H-H distances of \sim 1-1.2 Å. Analysis of the electron phonon coupling mechanism showed that rotations of the NH₄⁺ molecules modify these weak interactions, thereby perturbing the hydrogen atoms in the anionic sublattice and contributing to the microscopic mechanism of superconductivity. The predicted (Eliashberg) T_c s, of 179, 114, and 134 K for NH₁₀, NH₁₁ and NH₂₄, respectively, at 300 GPa, are comparable to some of the higher T_c s computed for the alkali polyhydrides.¹⁶ Our study suggests that molecular cations can serve the role of the electropositive element in superconducting hydrides, including

ones that may be constituents of the interiors of gas giant planets.

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Supporting Information Available

The Supporting Information is available free of charge on the ACS Publication website. It includes full computational details, 2D convex hulls, molecular orbitals of the NH_4 radical, phonon band structures and density of states, electronic structure analysis, trajectories of molecular dynamics runs, projected phonon linewidths and Eliashberg spectral functions, and structural parameters.

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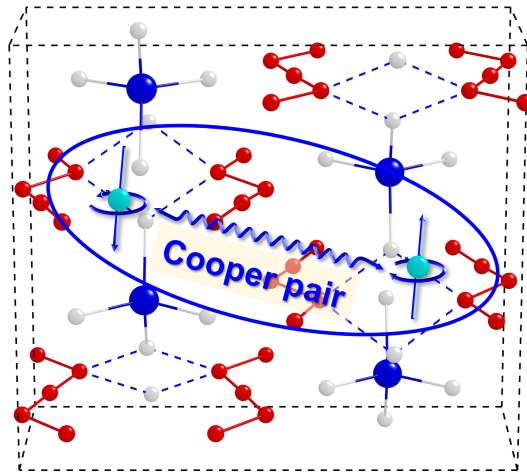
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Table of Contents Graphic



$Pnma-NH_{10}, T_c = 180 \text{ K at } 300 \text{ GPa}$