## Lawrence Berkeley National Laboratory

**Recent Work** 

## Title

THE REACTION DYNAMICS OF ELECTRONICALLY EXCITED ALKALI ATOMS WITH SIMPLE MOLECULES

**Permalink** https://escholarship.org/uc/item/1x4990gr

## Author

Weiss, P.S.

**Publication Date** 

1986-03-01



UC-4

B1-212



Prepared for the U.S. Department of Energy under Contract DE-AC03-76SF00098

#### DISCLAIMER

This document was prepared as an account of work sponsored by the United States Government. While this document is believed to contain correct information, neither the United States Government nor any agency thereof, nor the Regents of the University of California, nor any of their employees, makes any warranty, express or implied, or assumes any legal responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by its trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof, or the Regents of the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof or the Regents of the University of California.

## THE REACTION DYNAMICS OF ELECTRONICALLY EXCITED ALKALI ATOMS WITH SIMPLE MOLECULES

### Paul Storch Weiss

#### Materials and Molecular Research Division Lawrence Berkeley Laboratory

and Department of Chemistry University of California Berkeley, California 94720

The reactions of electronically excited sodium have been studied using the crossed molecular beams method. The apparatus and optical pumping are described in chapter I.

Chapter II describes the reactions of Na with the hydrogen halides. The reactions of Na(3S,3P,4D,5S) + HCl have been studied in detail. A large increase in the reactive cross section with electronic energy is observed. A change in the reaction mechanism leads to very different product scattering distributions for Na(3P) vs. Na(4D,5S) scattering. Laser polarization dependences show that this is due to long range electron transfer in Na(4D) + HCl that does not occur for Na(3P) + HCl. No reaction was observed for Na(3S,3P,4D) + HF up to collision energies of 13 kcal/mole. Na(3S,3P) + HBr behave much like the equivalent HCl reactions, except that the reduced endothermicity leads to more ground state reaction for HBr. Chapter III describes the reaction of Na with  $0_2$ . No reaction is seen for Na(3S,3P,4S,4P,5S) +  $0_2$ . Reaction to form NaO + 0 is observed for Na(4D) +  $0_2$  above the threshold translational energy of 15±1 kcal/mole. The reaction cross section is small,  $\sigma_R=0.9 \text{ Å}^2$  at 18 kcal/mole collision energy. The reaction is direct, does not occur via long range electron transfer, and a near collinear approach geometry is required for reaction. Some information about the transition state symmetry is derived from laser polarization dependences.

The reactions of  $Na(3S, 3P, 4D) + CH_3Br$  are described in chapter IV. As in all previously studied alkali plus alkyl monohalide reactions, the alkali halide product is predominantly back scattered due to steric restrictions. With increasing electronic excitation, a small relaxation of the restriction is observed, leading to NaBr scattered to lower center-of-mass angles than for the ground state reaction.

The reactions of Na(3S,3P) + Cl<sub>2</sub> are described in chapter V. At a collision energy of 6 kcal/mole, the reactive cross section increases 60% on electronic excitation, while at a higher collision energy, 19 kcal/mole, there is an increase of only 16%. A stripping mechanism explains the very similar scattering distributions of each state.

Yuan T. Lee

#### ACKNOWLEDGEMENTS

i

Many people have contributed to the work presented here, both in and out of the laboratory. There is room to mention only a few of them below, but none are forgotten, and I thank them all for their help.

On the experimental team, I have had the good fortune to work with a number of talented and bright individuals. First and foremost among these is Professor Yuan Lee. His proficiency at both molecular beam experiments and their interpretation, and his love for his work are awe-inspiring. Matt Vernon taught me how to use the apparatus, and set up much of what has produced the results described here. He also taught me his sound method of getting things done, which I have occasionally called into service. Dr. James Joseph Anthony Michael O'Brien helped set up the experiment before it came to fruition. The visits of Dr. Hartmut Schmidt and Dr. Jean Michel Mestdagh set the stage for several quantum leaps forward on these studies. Also working on these experiments was Michael "Bud" Covinsky, who was a constant source of entertainment. It has been great fun to work with Barbara Balko, who has shown me unending kindness in the last one and a half years. I have no regrets for keeping her up at lab for 36 hours at a time just for her company. The most recent addition to the A machine is Isabelle Duborg who helped with the final  $0_2$  experiments. I also thank her for her herculean effort in generating a great many of the figures that follow.

Several other members of the Lee group have supported me during these last few difficult months -- Gil Nathanson, Distinguished Miller Fellow, Lisa Yeh, and Babs Balko. I have learned something from everyone in the Lee group, but of particular note, conversations with Gary Robinson, Bob Continetti, Gil Nathanson, Carl Hayden, and Howard Nathel have been helpful in my work and in getting me through the days. and nights here. Ann Weightman has gone to much trouble to shield us from many of the harsh realities of bureaucracy, and for this and her help in the preparation of this thesis, I thank her.

A number of the LBL and UC campus support staff have been instrumental in keeping our experiments going. Among these, the ones that saved us time and time again are: Will Lawrence, Charlie Taylor, Tony Moscarelli, Fred Wolfe, and Fred Vogelsberg. I enjoyed my association with Fred Vogelsberg and Jacques Millaud on the multichannel scaler and the Timer-Gater-Scaler projects. A lot of good has come from their skilled and diligent work.

The San Francisco Laser Center, and its director, Dr. Andy Kung have come to the rescue many times when disaster struck or we just wanted to try something new. Steve Bittenson and Tim Ling have been helpful in keeping our lasers running well.

I first got started in chemistry at MIT, and it was Professor Bob Field who provided me with the most stimulating atmosphere I have ever encountered. He was patient, kind, and encouraging, and through his efforts I learned much more than I otherwise could have. All of what I

ii

learned from him and from two of his group members, Rick Gottscho and Ron Marks, has served me well at Berkeley.

This work was supported by the Director, Office of Energy Research, Office of Basic Energy Sciences, Chemical Sciences Division of the U.S. Department of Energy under Contract No. DE-AC03-76SF00098.

Others, outside of science have given me the support I needed to carry out this work, and much more. The O'Connor family, and Harry Pasternack and Liza Cramer helped me through critical times in Berkeley. Kendall's parents, Walter and Judy Munk, and her grandmother, Mrs. Winter Horton have made me feel very special, and this has given me an extra gear to work harder. All my life, my parents have given me their love, confidence, and support, and none of this would have been possible without them.

My wife Kendall and son Walter have waited patiently for me to finish. Kendall has been ever so understanding and encouraging with my odd schedules from the very beginning, and it is to Kendall that this thesis and I are dedicated.

iii

# THE REACTION DYNAMICS OF ELECTRONICALLY EXCITED ALKALI ATOMS WITH SIMPLE MOLECULES

Ta	<b>h</b> '	le	of	Con	tents
	<b>U</b>		01	0011	661163

Abstract						
Ackno	wled	gements	i			
Table	of	Contents	iv			
I.	EXPERIMENTAL					
	Α.	The Experimental Apparatus	1			
	B.	Atomic and Molecular Beam Conditions	20			
	С.	Data Acquisition	32			
	D.	Optical Pumping	39			
		1. Introduction	39			
		2. The Na D <sub>2</sub> Transition	43			
		3. Subsequent Excitation to $Na(4^2D_5/2)$	57			
		4. Subsequent Excitation to $Na(5^2S_{1/2})$	62.			
	E.	References	65			
II.	THE REACTIONS OF GROUND AND EXCITED ALKALI ATOMS WITH HYDROGEN HALIDE MOLECULES					
	Α.	Introduction	68			
	B.	Results	81			
		1. Na(4D,5S) + HCl	81			
		2. Na + HF	97			
		3. Na(3S,3P) + HBr	97			

ίv

			v
·		C. Analysis of Experimental Results	101
		D. Discussion	117
		<ol> <li>The Mechanism and Distribution of Product Energy in the Na(4D,5S) + HCl Reaction</li> </ol>	117
		2. The Lack of Reaction for Na + HF	129
		E. Conclusions	130
		F. References	132
	III.	THE STATE SELECTIVE REACTION OF ELECTRONICALLY EXCITED SODIUM ATOMS WITH MOLECULAR OXYGEN	
		A. Introduction	137
		B. Results	144
		C. Analysis of Experimental Results	179
		D. Discussion	194
		E. Conclusions	202
		F. References	204
	IV.	THE REACTIONS OF GROUND AND EXCITED STATE ALKALI ATOMS WITH METHYL HALIDE MOLECULES	
• .		A. Introduction	207
		B. Results	212
		C. Analysis of Experimental Results	218
		D. Discussion	229
• • •		1. Observations on the Experimental Results	229
		2. Other Possible Processes	230.
		E. Conclusions	232
·		F. References	233
		·	

· . .

۷.	THE REACTIONS OF GROUND AND EXCITED STATE ALKALI ATOMS WITH HALOGEN MOLECULES						
	Α.	Intro	duction	237			
	B.	Resul	ts	254			
	C.	Analy	sis of Experimental Results	261			
	D.	Discu	ssion	272			
	D.	Concl	usions	276			
	Ε.	Refer	ences	278			
VI.	CON	CONCLUSIONS					
	Α.	Concl Excite	usions on the Reaction Dynamics of Electronically ed Alkali Atoms	282			
VII.	APP	APPENDICES					
	Α.	Data Acquisition for Angular Distributions, Polarization Dependences, Machine Condition Optimization: The Data Acquisition System and Program SANG					
	÷	1. Iı	ntroduction to the Data Acquisition System	284			
		2. In	ntroduction to Program SANG	286			
		3. L <sup>.</sup>	isting of Program SANG	294			
	Β.	Data / The Ne	Acquisition for Time of Flight Measurements: ew Multichannel Scaler and Program TUF	399			
		1. TH	ne Purpose of the Multichannel Scaler	399			
		2. C/	AMAC Commands for the New Multichannel Scaler	401			
		3. Mu	ultichannel Scaler Hardware Specifications	406			
		4. Mi De	ultichannel Scaler Circuit Logic and Signal	409			

vi

	5.	Introduction to the Program TUF	417
	6.	Listing of Program TUF	419
C.	Ref	erences	474

vii

# To Kendall

#### I. EXPERIMENTAL

A. The Experimental Apparatus

The experimental measurements described below were carried out on a crossed molecular beams machine modified to allow laser excitation at the crossing of the atomic and molecular beams. The basic apparatus has been described previously, 1-5 but will be reviewed here.

An atomic sodium beam, a molecular beam, and one or two dye laser beams cross orthogonally in a vacuum chamber under single collision conditions as shown schematically in figure 1. A mass spectrometric detector rotates about the collision region in the plane defined by the atomic and molecular beams. Three types of experiments are performed: (1) product angular distributions are measured by rotating the detector, (2) product velocities for fixed detector angles are measured by modulating the product directly or the reaction itself by modulating the laser, and (3) the effect of rotating the laser polarization and thus the excited orbital alignment upon the reactive product at a fixed detector angle is measured.

The sodium atomic beam source was designed by Dr. H. Schmidt and is a hybrid of the sodium source design used in the I. V. Hertel group at the Freie Universitaet Berlin, and the lithium source used for earlier measurements of Li + HF, HCl.<sup>3</sup> Figure 2 is a schematic of

1.



XBL 855-2592 B

Fig. 1. Schematic of the crossed molecular beams apparatus.

Fig. 2. Schematic of the seeded supersonic sodium atomic beam source (adapted from a similar figure in reference 2, with permission):

3

(1) Primary source chamber, (2) Sodium reservoir, (3) Thermal coaxial heaters for reservoir, (4) Bottom oven thermocouple, (5) Bottom oven radiation shields, (6) Nozzle tube, (7) Thermal coaxial heaters for nozzle, (8) Nozzle disk, (9) Nozzle radiation shields, (10) Nozzle thermocouple, (11) Oven support rods, (12) Insulating ceramic for support rods, (13) Thin wire to thermally insulate support rods from (14) Copper blocks attached to (15) Water-cooled copper structure to hold oven rigidly aligned, (16) Upper reservoir, (17) Radiative heater for upper reservoir, (18) Top oven radiation shields, (19) Upper reservoir thermocouple, (20) Gas inlet tube, (21) Flow constrictor, (22) Heated Skimmer, (23) Thermal coaxial heaters for skimmer, (24) Heated pre-skimmer shield, (25) Thermal coaxial heaters for pre-skimmer, and (26) Rectangular beam-defining aperture.



XBL 834-9475A



the sodium source chamber (1). The stainless steel oven chamber is divided into three areas. The bottom of the reservoir (2) contains the sodium metal that is heated by thermal coaxial cable brazed to the outside of the chamber (3). The temperature, typically 500°-600°C, of this chamber controls the vapor pressure of the sodium, and is measured by a thermocouple attached to the bottom of the oven (4). Two radiation shields surround the oven (5). This temperature is stabilized to ±0.5°C using a thermocouple comparator circuit to regulate the heater supply. A 3" long, 0.25" diameter tube (6) protrudes from near the top of the bottom reservoir region. This tube is heated separately to 750°C by thermal coaxial cable (7) vacuum brazed to it in order to prevent nozzle clogging, to prevent condensation of the sodium atoms during the supersonic expansion, and to determine the terminal velocity of the atoms. The nozzle temperature is also regulated to  $\pm 0.5$ °C using a thermocouple comparator circuit. A stainless steel disk 0.20" thick is welded to the end of the tube (8). The nozzle diameter in the center of this disk ranges from  $100\mu$  to  $250\mu$ . The nozzle tube is wrapped in three layers of radiation shielding -- thin sheets of stainless steel (9). Underneath the radiation shields approximately 10 mm. from the nozzle, a thermocouple (10) measures the temperature of the tube. The thermocouple is kept this far from the nozzle because by attaching the thermocouple directly to the nozzle enough heat is conducted away from the nozzle to cool and regularly clog the nozzle. When the ovens were heated under identical conditions and the sodium beam had the same velocity and profile, the thermocouple further from

the nozzle measured 120°C lower than the thermocouple at the nozzle. At the same height as the nozzle tube, three 1/4" stainless steel rods protrude from the oven (11). A 5 mm. long, loosely fitting ceramic tube (12) is placed on each rod to prevent the top reservoir heater from being in electrical contact with the oven. A thin stainless steel wire (13) is wrapped on the end of these rods to minimize the thermal contact with the water cooled copper blocks (14) that are clamped down on the rods. These blocks are connected to a larger copper structure that attaches rigidly to the front of the source chamber (15) and is used to align the nozzle before placing the source chamber into the molecular beams machine. Since the copper blocks are water cooled and do not change temperature substantially, the oven and nozzle are held in the correct aligned position. The upper reservoir chamber (16) has a conflat-like knife edge seal with a 0.005" nickel gasket. This port is used for loading sodium into the oven. This region is heated radiatively by a tungsten ribbon wrapped around eight alumina posts (17), held in place by slotted beads, and shielded by three layers of radiation shielding -- thin sheets of tantalum (18). A thermocouple (19) is attached to the top of the upper mini-conflat, and the temperature here is kept 25°C hotter than that of the bottom reservoir to prevent any condensation in this region. On the upper mini-conflat a 1/4" tube (20) extends up for the introduction of seeding gas into the oven chamber. Inside this tube is a 1/4" long molybdenum flow constrictor with a 0.005" clearance from the inside of the tube (21). This flow constrictor is held in place 1/4" from the opening of the

inlet tube into the reservoir chamber by a 0.018" wire spot welded several inches back along the gas inlet path. It is essential that this flow constrictor be pulled back from the opening since when it is nearer the opening it frequently clogs the gas inlet line. This flow constrictor prevents back diffusion of the sodium into the gas inlet path by increasing the local velocity of the seed gas around it. Only ultrahigh purity (0.99999) helium, neon, argon, and mixtures of the above are used for seeding the sodium beam.

The beam conditions are optimized for reactive signal and beam quality (high speed ratio). The sodium vapor pressure at 500°-600°C (the bottom reservoir temperature) varies from 3 to 25 torr. The beam velocity distributions for the various seed gases are given in section B below.

The Na rare gas mixture is expanded through the nozzle into a vacuum chamber maintained at  $0.5-2 \times 10^{-4}$  torr. 0.25" to 0.30" in front of the nozzle is a stainless steel skimmer (22) with a diameter of 0.045". This skimmer is heated conductively by thermal coaxial cable (23) to  $500^{\circ}-600^{\circ}$ C to prevent clogging with condensed sodium. A heated stainless steel shield (24), termed a pre-skimmer, prevents sodium condensation on the inner walls of the water-cooled source chamber, which would otherwise form a bridge to the skimmer, and clog it. The pre-skimmer is made from two disks sandwiching thermal coaxial cable (25), all vacuum brazed together. It is also heated to  $500^{\circ}-600^{\circ}$ C.

Beyond the skimmer is a differential pumping region with a typical pressure of 2-10 x  $10^{-6}$  torr. This region is separated from the main chamber in which the scattering takes place by collimating slits. These slits are fashioned from 4 razor blades spot welded together to give a 0.53 mm. by 1.90 mm. rectangular aperture (26). This determines the beam dimensions at the scattering center to be 1.1 mm. high and 3.0 mm. wide.

The collision region is in the main chamber with a liquid nitrogen cold shield, and is typically at  $1-2 \times 10^{-7}$  torr. A triply differentially pumped ultrahigh vacuum quadrupole mass spectrometer equipped with an electron bombardment ionizer and a Daly ion detector rotates about the collision region in the horizontal (scattering) plane. Its angular range is -6° to 98 5° (where 0° has the detector pointing into the atomic beam, and 90° has the detector pointing into the molecular beam). For a detailed description of the detector, see reference 5.

The (secondary) molecular beam source is a heated stainless steel tube with a 70 $\mu$  platinum nozzle on the end as shown in figure 3. This nozzle is an electron microscope aperture held in place with a small copper gasket tightened on by screwing down a stainless steel cap (1) on the threaded end of the tube. Heating tape (2) is wrapped around the tube up to near a keying device (3). Two copper block heaters (4,5) with thermal coaxial cable running through them are clamped on the tube, one in back of the keying device, and one in front, flush with the nozzle tip. There are thermocouples (6) attached to the nozzle's copper block and underneath the heating tape. This source is

Fig. 3. Schematic of the supersonic molecular beam source and the associated chamber (adapted from a similar figure in reference 2, with permission):

(1) Nozzle cap, (2) Heating tape for the gas inlet tube,

(3) Keying device, (4) Copper block nozzle heater,

(5) Nozzle pre-heater, (6) Thermocouple, and (7) Skimmer.



used for all reactant molecular beams. The temperature of the nozzle's copper block is stabilized to  $\pm 0.5$ °C with a thermocouple comparator circuit referenced to an ice water bath. Neat or seeded supersonic beams are made by passing gas (neat or premixed) or by bubbling a gas through a volatile liquid held in a bubbler. The beam is expanded into a 0.5-2 x  $10^{-4}$  torr region. The beam is skimmed by a stainless steel skimmer with a diameter of 0.023" (7). Beyond the skimmer is a differential pressure chamber at 1-5 x  $10^{-6}$  torr which contains a 150 Hz tuning fork chopper (not shown) for modulating the secondary beam. This is mounted on a water cooled copper block mounted on the wall of the source chamber.

A collimating aperture 2.1 mm. high by 1.8 mm. wide similar to the one described above separates the differential region from the main chamber, and defines the secondary beam to be 2.7 mm. high by 2.1 mm. wide at the collision region.

The modifications to the crossed molecular beams machine all have to do with the introduction of the exciting lasers and the monitoring of the atomic fluorescence.

Two optical flats were placed on the main flange of the machine (the main door) facing the secondary molecular beam source -- one directly in the path of the molecular beam, and one 10" below that. One or two dye laser beams enter through the lower window and are reflected up into the crossing region perpendicular to the scattering plane (thus avoiding a Doppler broadening from all but the transverse velocity components of the sodium beam, and also keeping the laser

polarization in the scattering plane). The upper window is used for the collection of atomic fluorescence (particularly when products are predominantly scattered near the sodium beam, and the detector would block the view from the window located on the opposite side of the sodium beam, which is more commonly used). It can also be used for the laser to enter in the scattering plane, but still perpendicular to the sodium beam in order to study the difference between in plane and out of plane atomic orbital alignment.

A large flange in the line of the atomic beam was replaced by one with seven conflat ports. Three of these are equipped with conflat viewports that directly view the collision region. One is directly in the path of the sodium beam, and the image of the sodium beam through the rectangular collimating slits is apparent on it after running for a few hours. Another viewport is in the scattering plane and in line with the detector when the detector is placed at  $-12.5^{\circ}$ . This port is normally used for the fluorescence monitor described below. All the fluorescence measurements shown were recorded with the fluorescence monitor in this position. The detector blocks the view of this port if it is at an angle below 16°.

Two other optical ports are on the machine on the two sides onto which the beam sources are mounted. Each makes a 45° angle with the atomic and molecular beam axes at their crossing point in the scattering plane so as to give a clear line of sight that passes through the crossing region. There is a Brewster window and a set of three sharp blackened apertures at each port to spatially filter the laser

and stray light. The laser is passed through these two ports in order to measure the Doppler profile of the Na  $D_2$  line and thus the velocity profile of the sodium beam. The port near the sodium beam (the laser exit flange for the 45° crossing) is equipped with a simple system to measure two different colors of fluorescence using an interference filter and two photomultipliers.

The optical pumping setup is shown in figure 4. The  $Na(3^{2}P_{3/2} + 3^{2}S_{1/2}) D_{2}$  transition is pumped by a Coherent Radiation Model 599-21 cw single-frequency actively stabilized linear dye laser using Rhodamine 6G, pumped by 2.7 W at 5145Å using an argon ion laser. This laser outputs 70-100 mW at the 5892Å (yellow-orange) Na  $D_2$  transition wavelength. A second transition, either the  $Na(5^2S_{1/2} + 3^2P_{3/2})$ transition at 6162Å (red), or the Na( $4^{2}D_{5/2} \epsilon^{3^{2}P}_{3/2}$ ) transition at 5690Å (green), is pumped by a Coherent Radiation Model 699-29 or 699-21 cw single-frequency actively stabilized ring dye laser using Rhodamine 6G pumped with 6.0 W at 5145Å with an argon ion laser. At the 6162Å (red) transition it put out 200-500 mW, and at the 5690Å (green) transition it put out 400-800 mW. In some cases the laser frequencies were dithered ±2-5 MHz at 510 Hz and 23 Hz, respectively, and locked to the peak of the fluorescence (of the upper transition if two lasers were in use) using lock-in amplifiers. Neutral density filters are inserted into the laser path in order to ensure that saturation effects are not causing a reduction in the fraction of Na atoms optically pumped. This is discussed further in section D.

Fig. 4. Schematic of the optical pumping system used for the sodium atomic excitation. The symbols for optional components are in parentheses. The components are:

(BS) Beam Splitter, (M) Mirror, (GT) Glan-Thompson Prism,  $(\lambda/4)$  Quarter-Wave Plate, (FP) Fabry-Perot Etalon, (S) Shutter, (F) Filter, (PMT) Photomultiplier and/or Fluorescence Monitor, (Na) Sodium Source, (2°) Secondary Molecular Beam Source, and (MS) Quadrupole Mass Spectrometer.



XBL 862-631

Fig. 4

As shown in figure 4 the two lasers are combined by having one laser beam nearly graze a mirror (M') off the edge of which the other beam is reflected. The two beams are thus very nearly parallel and their overlap is maximized in the collision region (after a path length of greater than 3 m.) by maximizing the laser produced reactive signal. The two laser polarizations are kept parallel at all times. Each laser is monitored with a spectrum analyzer (FP) with a 1.5 GHz free spectral range. An iodine cell is used to determine the frequency of the laser using the iodine atlas of Gerstenkorn and Luc. $^{6}$ Approximately once an hour, a computer controlled shutter (S) sends a fraction of the 5890Å light to the 45° Doppler crossing path in order to measure sodium beam velocity profiles as described above. The lasers can be circularly polarized before being combined using quarter-wave plate  $(\lambda/4)$  for the 5890Å laser, and another guarter-wave plate or a Soleil-Babinet Compensator for the other laser. After being combined and before entering the main chamber the two lasers can be passed through a cross-correlation time-of-flight wheel. The wheel has a pseudo-random sequence of 255 on/off bits and can be spun up to 400 Hz in a vacuum chamber pumped down to  $10^{-2}$  torr to give a time resolution of 10.0 usec.<sup>7</sup> One or both of the lasers can be modulated at 3 Hz using a stepping motor driven beam flag (not shown) run synchronously with the 150 Hz tuning fork secondary beam chopper. Also the linear laser polarizations can be rotated in tandem using a computer controlled stepping motor driven double fresnel rhomb polarization rotator. As described above, the lasers enter the vacuum

chamber parallel to and below the secondary beam, and are reflected up into the collision region from a mirror on the bottom of the machine. This mirror is mounted on a heated platform to keep diffusion pump oil from condensing on it, and to prevent movement as the liquid nitrogen cold shield is cooled. Above the mirror is an aperture to spatially filter the laser light. This aperture is attached to the rotating detector, but the aperture is located on the detector's rotating axis, and is thus centered vertically below the collision region so that its position does not change as the detector is rotated. A photodiode attached to the differential wall of the machine is directly above the collision region and is used for the rough alignment of the laser position. The fine adjustment of position is done so as to optimize the reactive signal due to the laser excitation as seen by the mass spectrometer.

The machine has three possible positions for fluorescence monitors. The most commonly used positions are in line with the detector positions  $-12.5^{\circ}$  and  $90^{\circ}$  (relative to the sodium beam). These fluorescence monitors are shown schematically in figure 5. A Nikon 300 mm. f/4.5 lens is placed near the machine window and is attached to a Nikon PB-6 bellows. If two lasers are in use, an interference filter or a color filter is placed behind the bellows so that the monitor only views the fluorescence of the upper transition. At the end of another bellows, a focal point is formed by the front lens at which an aperture is placed to spatially filter out stray light. This is followed by a reversed Nikon 28 mm. f/2.8 lens which is attached to yet another



XBL 861-10007

Fig. 5. Schematic of the fluorescence monitor used for the magnification and measurement of the sodium atomic fluorescence.

bellows, which is attached to a Nikon F3 35 mm. single lens reflex camera with a hole in its back and a photomultiplier tube (RCA 1P28) attached to it. This set up gives a magnification of 10x, making the 1 mm. x 3 mm. crossing region appear 10 mm. x 30 mm. on the photomultiplier tube. The fluorescence can be viewed visually through a 6X magnifying viewfinder (Nikon DW-4). A fluorescence monitor on the exit of the 45° laser crossing consists of a tilted interference filter which reflects one wavelength and transmits another so that two photomultipliers (RCA 1P28) record each fluorescence wavelength simultaneously. A shutter is closed when a Doppler velocity measurement is being done so as to avoid damaging the photomultipliers, and a different fluorescence monitor is used to record the velocity measurement.

The experiment is controlled by an LSI-11/23 computer. Angular distributions and polarization dependences are recorded, and machine conditions are optimized with a program entitled SANG which is described in Appendix A. Time-of-flight measurements are made with a program named TUF, described in Appendix B. The details of the data acquisition are given in section C, below.

B. Atomic and Molecular Beam Conditions

As discussed above, three sodium atomic beam energies are most commonly used. The rare gas backing pressures used are: 200-1000 torr for helium, 100-400 torr for neon, and 100-400 torr for argon. The large ranges are due predominantly to the various nozzle sizes at which measurements were made. It was found necessary to seed the atomic beam with gases of research grade purity (0.99999, Scientific Gases), as impurities in lower grades of rare gases clogged the sodium beam too quickly for measurements to be made. No additional benefit was found when the backing gases were passed through a liquid nitrogen cold trap. The gas feed line to the reservoir was all metal, and this line must be leak tight.

Figure 6 shows three velocity measurements taken by scanning the dye laser across the  $Na(3^{2}P_{3/2} + 3^{2}S_{1/2}) D_{2}$  transition with the laser crossing the Na beam at both 45° and 90°. The recorded fluorescence shows (from right to left -- red to blue) the Doppler broadened and shifted 45° peaks and then the sharp unbroadened, and unshifted 90° peaks. There are two peaks of each type because the two ground state hyperfine levels (F=1,2) are separated by 1.77 GHz, and are thus spectrally distinct. Below each fluorescence trace is the simultaneously recorded output of a spectrum analyzer (Fabry-Perot etalon) used as a relative frequency standard. This etalon (Tropel Model 210) has a free spectral range of 1.5 GHz and a finesse of 200. In the cases of sodium

Fig. 6. Velocity measurement of sodium atomic beams produced with three different seed gases: a) Helium, b) Neon, and c) Argon.





Fig. 6(a)






Fig. 6(c)

seeded in neon and argon, the weak (corresponding to F''=1) 45° peak is buried under the strong (corresponding to F''=2) 90° peak. In each case, to convert the peak separations into velocities the distance between the strong 45° and 90° peaks and between the Fabry-Perot peaks, and the full width half maximum of the strong 45° peak are measured either by hand or if the computer was used to record the Doppler scan, it is also capable of making these measurements automatically. The 90° to 45° peak separation and the 45° width are converted into frequencies and the following Doppler formula is used:

$$v = \frac{c \Delta v}{v \sin 45^{\circ}}$$
(1)

where v is the velocity being measured in cm/sec, c is the speed of light, in cm/sec,  $\Delta v$  is the measured frequency difference between the two velocity components of interest in GHz, and v is the sodium D<sub>2</sub> transition frequency,  $5.092 \times 10^5$  GHz (16973 cm<sup>-1</sup>). The speed ratio (S = v/ $\Delta v$ ) is just the ratio of the measured 45° to 90° peak separation to the measured width of the 45° peak. The values of the peak velocity and speed ratio for the three backing gases commonly used are shown in table I. Intermediate velocities are obtained by mixing combinations of helium and neon. A Matheson rotameter is used to keep the flow of each gas constant to 5%. In this case, the sodium beam velocity is measured as above.

Seed Gas	Peak Sodium Velocity ( /10 <sup>4</sup> cm/sec)	Speed Ratio
Helium	30.0	6
Neon	16.0	5
Argon	10.8	5

Table	Ι.	Measured	sodium	beam	velocities	and	speed	ratios	for	the	three
		commonly	used se	ed ga	ases.						

The secondary beams' conditions used for the measurements discussed in the following chapters are shown in table II. When gas mixtures are shown, these are premixed either commercially or in our lab. The beam velocities were measured using the time-of-flight method on another crossed molecular beam machine using flight length and detection delays calibrated with rare gas beams of helium and argon. A 17.8 cm. diameter wheel with eight .75 mm. slots was used to modulate the beam. After modulation, the neutral flight length was approximately 20 cm. The distributions were recorded using the program  $\mathsf{TUF}^\mathsf{8}$  with dwell times of 500-1500 nsec per channel. The program KELVIN<sup>1</sup> was modified (to accept the output of TUF) and used to fit the measured distributions to peak velocities and speed ratios. The measured beam time-of-flight distributions are shown for neat beams of HC1 and  $0_2$  and  $0_2$  seeded in helium in figures 7, 8a, and 8b, respectively. These measurements were not always made for beams for which quantitative conversions of laboratory data to the center-of-mass frame of reference were not to be



Fig. 7. Measured time-of-flight distribution for a neat HCl molecular beam under the conditions used for the experiments in chapter II. The solid line is the fit to the distribution using program KELVIN.

Fig. 8. Measured time-of-flight distributions for the  $0_2$  molecular beams used for the experiments in chapter III. The solid line is the fit to the distribution using program KELVIN.<sup>2</sup> The distributions are: a) neat  $0_2$ , and b)  $0_2$  seeded in helium.



Fig. 8(a)



XBL 863-825

Fig. 8(b)

made. The velocities of the unmeasured beams were estimated from nozzle temperatures assuming essentially complete rotational relaxation (as in reference 1).

Table II.	Velocities and speed ratios for molecular beams used in	
	reactive scattering studies.	

Reactant Beam	Seed Gas	Nozzle Temperature (°C)	Beam Velocity ( /10 <sup>4</sup> cm/sec)	Speed Ratio	Fig.	Chap.
CH <sub>3</sub> Br	none	200	5.8	4		4
5% CH <sub>3</sub> Br	He	200	12.6	4		4
C1 <sub>2</sub>	none	179	6.1	5		5
HBr	none	× 200	5.8 (Est.)	4		2
HC1	none	179	8.6	4.5	7	2
HF	none	250	12.3 (Est.)	4.6		2
0 <sub>2</sub>	none	200	9.8	8	8a	3
5% 0 <sub>2</sub>	He	200	15.0	12	8b	. <b>3</b>

## C. Data Acquisition

The data acquisition is automated using a Digital Equipment Corporation LSI-11/23 computer and a CAMAC crate (IEEE Standard 583-1975). The programs used for data acquisition are described in appendices A and B, but their specific usage is described below.

On each new reactive system the first experiment performed is the measurement of one or more product angular distributions. Six channels of data are recorded simultaneously from two sources. These channels are defined in table III. The signal from the quadrupole mass spectrometer (labeled detector in table III) is taken from the Daly detector's photomultiplier (RCA 8850) and put into a quad discriminator (LeCroy 821), the fanout of which is put into the four inputs of a CAMAC 125 MHz guad scaler with external gating (Joerger S1-Ind). The other source of signal is from the fluorescence monitor (labeled fluorescence in table III). This signal is run through a current amplifier (Keithley 427) into a gated integrator (Evans 4130). the overflow of which is converted to NIM pulses and is input into another channel of the quad discriminator for the fanout capabilities, and then is recorded using another (Joerger) guad scaler. The signal is gated using the output of a CAMAC powered Laser Gate Generator (LBL, Lawrence Berkeley Laboratory). This gate generator takes as input the output of a NIM Timer-Gater module (LBL 13X-3050-II) which converts the signal

	tions a various	nd polarization signals.	n measurements, and the	origins of the
Channe 1		Laser(s') Status	Molecular Beam Status	Signal Source
la 1b		On On	On Off	Detector Detector

On Off

Independent Independent

Off

0ff

0n

Off

from the 150 Hz tuning fork chopper that modulates the secondary
(molecular) beam into gate information. The phase and gate width are
variable on this device. The laser gate generator also is used with a
CAMAC powered Motor Sync Scaler (LBL 13X-3441-P1-A) which outputs a
pulse every 25 periods of the tuning fork chopper in order to change
the state of the laser at a frequency of 150/25, or 6 changes per
second, giving a total modulation frequency of 3 Hz. After every 24
periods of the tuning fork chopper, all counting is gated off for 1
period and a stepping motor is used to change the state of the laser
beam flag. If only one laser is in use, it is blocked and unblocked.
If two lasers are in use, one or both can be blocked and unblocked.
In two laser experiments, the reactions of four electronic states are
measured. First, the Na(4D) vs. Na(3P) distributions are measured by

Table III. The definitions of channels recorded for angular distribu-

 $\sim$ 

2a

2b

3a

3b

Detector Detector

Fluorescence Fluorescence

having the yellow-orange (Na(3P < 3S) transition) laser always on, with the green (Na(4D)<Na(3P) transition) laser blocked and unblocked. Then, the Na(5S) vs. Na(3S) distributions are measured by blocking and unblocking the yellow-orange (Na(3P < 3S) transition) and the red (Na(5S < 3P) transition) lasers with the beam flag.

The scattering signals are derived as shown below. The laser off signal (OFF) is:

$$OFF = (2a - 2b) / TIME,$$
 (2)

where 2a and 2b are the number of counts in scaler channels 2a and 2b, respectively, from table III, and TIME is the actual counting time for each channel. The counting time for a single measurement is:

$$TIME = \frac{1}{2} \frac{24}{25}$$
 (Gate Width) (Countdown Time), (3)

where the gate width and countdown time are Timer-Gater settings, in seconds. The laser on signal (ON) is:

$$ON = (1a - 1b) / TIME.$$
 (4)

The signal due to the excited state (EXC) can be easily derived from the above:

$$EXC = \frac{ON - [(1 - FE)OFF]}{FE},$$
 (5)

where FE is the fraction excited. Note that the contribution to ON due to the fraction remaining in the lower state (1 - FE) is removed.

"Loop mode" of program SANG (described in appendix A) is used to record the angular distributions. An angular range and an angular interval are chosen for the scan. For example, the range might be 20° to 80°, with data recorded every 2°. Each angle is measured for a set countdown time, typically 60 seconds in real time. An angle is chosen for the normalization of the data (to counteract the effects of long term drifts in machine conditions); this angle is usually chosen to be the peak of the reactive signal. After every ten measurements, this angle is measured two, four, or more times. The time at which each measurement is taken is recorded, and the data is normalized later by taking a linear interpolation <u>in time</u> between successive sets of normalization measurements.

Polarization measurements are recorded with program SANG in "polarization rotation mode" with this same setup, but with the addition of the broadband polarization rotator (Spectra Physics 310A). This polarization rotator is mounted and attached to a stepping motor (Slo-Syn M062-FC09E) driven by a CAMAC pulse generator (Kinetic Systems 3360). The polarization angle resolution (minimum step size) is 0.8°, where 180° is a full cycle. After over 1000 cycles of the polarization rotator no deviation from its expected angular position is observed. Eight steps are typically chosen per cycle, and, as for angular distribution measurements, the measurement is usually made for 60 seconds in real time after which the rotator is automatically advanced to its next position. Note that each polarization rotation measurement is made

with the mass spectrometer held at a fixed angle, and different distributions are seen for different detector angles. For the experiments discussed in subsequent chapters the laser polarization was only rotated in the scattering plane, however by using a different laser entry window, the laser polarization can be rotated out of the scattering plane also. This has been done in recent measurements of Na + HC1.<sup>9</sup>

The favored polarization angle  $(\phi)$  and the amplitude of the polarization dependence (2A) are determined by fitting a cosine curve of the form:

Normalized Signal =  $A\cos(2(\Theta - \phi)) + (1 - A)$  (6) to the reactive signal. The same is done for the fluorescence signal for the signal from a fixed fluorescence monitor in the plane of the scattering. This fluorescence amplitude provides a direct measure of the fraction of atoms aligned as discussed in section D. The fit of equation (6) to the fluorescence data (shown in chapters II and III) gives the amplitudes as 2A=0.35 for the 3P excitation, and 2A=0.32 for the 4D excitation.

The time-of-flight measurements reported here are made by modulating the laser with a pseudo-random binary sequence of 255 bits photoetched as slots on a wheel that is spun up to 392 Hz to give a time resolution of 10 usec over the 20.8 cm. flight path of the reaction products. As discussed in reference 1, this has the disadvantage of modulating any signal due to ground state in the opposite direction, and this data must be either measured or accurately

synthesized in order to be taken into account in the subsequent data analysis. An auxiliary slot in the cross-correlation wheel is monitored by having a light-emitting diode on one side of the wheel and a photodiode on the other. The brief signal from the wheel is converted into a TTL pulse by triggering a pulse generator (Systron Donner 101) the negative output of which goes to the trigger input of a CAMAC 4096 channel scaler (LBL, described in appendix B). The data is the mass spectrometer signal taken from the fanout of the quad discriminator (described above) which is converted to TTL pulses, then is sent to the data input of the multichannel scaler. Program TUF (described in appendix B) is used to record the time-of-flight data.

The data is transformed from the laboratory frame of reference in which it is recorded to the center-of-mass frame of reference in which it is easily interpretable using a derivative of the program CMLAB<sup>10</sup> entitled GM which runs on the DEC VAX 8600 computers at the Lawrence Berkeley Laboratory computer center. This program assumes separable parameterized center-of-mass angular  $(T(\theta))$  and translational energy  $(P(E_T))$  distributions, and calculates laboratory angular and velocity distributions taking into account beam velocity spreads, and finite modulation width and ionizer length.<sup>9,10</sup> The parameters for the center-of-mass distributions are varied iteratively in order to minimize the differences of the measured and calculated laboratory angular and velocity distributions. In this calculation it is assumed that the reactant orientations average to give cylindrically symmetric

product distributions about the relative velocity vector. In the case of laser excited sodium atoms this is, in general, not correct (when the laser polarization is along the relative velocity vector it is correct). The effect of this asymmetry is not more than 15 percent as shown by the polarization dependences in chapters 2 and 3, and thus has been ignored for the analyses done here.

It should be noted that in the version of CMLAB given in reference 10, and the versions used for the last few years, the product translational energy distribution is normalized incorrectly, giving false values of the relative cross sections. As far as can be determined, no such values have been reported. The problem lies in the fact that the laboratory scattering data is synthesized from the center-of-mass information by using a grid of laboratory velocities. As pointed out in reference 10, this is nearly equivalent to using P(u), the probability distribution of product center-of-mass velocities (u), rather than  $P(E_T)$ , the probability distribution of recoil energies. Thus, since

 $E_{\rm T}^{\prime} = \frac{1}{2} \mu_{\rm Prod} u_{\rm rel}^2,$  (7)

where  $\mu_{Prod}$  is the product reduced mass, if P(u) is substituted for P(E) in the transformation Jacobian, there is a further transformation factor:

$$P(u) = P(E) \frac{dE}{du} = P(E) \mu_{Prod} u_{rel}^{o}$$
(8)

The factor u was correctly taken into account in the transformation, but the product reduced mass was not. This made any comparison of branching ratios where product channels have different reduced masses incorrect. Also, calibration of relative cross sections using elastic scattering data was not possible. Furthermore, P(E) was normalized so that all intensity information was carried in the computed scaling factors, and in  $T(\Theta)$ , the center-of-mass product angular distribution. However, since P(u) is used in the transformation, this is what needs to be normalized.

Both the problems of the product reduced mass and the normalization of P(u) have been corrected in all versions of GM dating from the beginning of 1986 onward.

## D. Optical Pumping

## 1. Introduction

In determining the reactivity of the excited state atoms it is of fundamental importance to know the fraction of the atoms populating these states. It is difficult in the apparatus described above to make a direct measurement of the optical pumping efficiency. It is certainly possible to optimize the optical pumping (laser frequencies, powers, and positions) on reactive signal due only to the excited state produced, but it remains difficult to know the actual fraction excited. One possible method of obtaining the fraction excited is to find a system in which the excited state gives no signal under the same circumstances that the ground state does. Then, the fraction of the ground state signal depleted on excitation is exactly equal to the fraction excited. Of course, unless it was certain that the excited state yielded no signal, the excited state fraction obtained in this manner would be only a lower limit. No good example of such a situation was found in the experiments conducted.<sup>11</sup>

The Na atomic energy levels are shown in figure 9. The fine and hyperfine structure are not shown in this figure. The three upward pointing arrows show the essential transitions optically pumped. The transition moments are quite well known for Na atoms. The Einstein A coefficients have been measured experimentally  $1^{12}$  and calculated  $1^{13}$ for all the states relevant to the present study. These values are given in table IV. To pump the  $Na(3^2P_{3/2})$  state a single laser (usually the Coherent 599-21 cw dye laser) is tuned to the 16973.379 cm<sup>-1</sup> (yellow-orange) Na( $3^{2}P_{3/2}$ , F=3  $\leftarrow$   $3^{2}S_{1/2}$ , F=2) transition, part of the Na D<sub>2</sub> line. In order to pump the  $5^2S_{1/2}$  state, a second laser (the Coherent 699-21 or 699-29 cw ring dye laser) is tuned to the 16227.317  $cm^{-1}$  (red) Na(5<sup>2</sup>S<sub>1/2</sub>, F=2  $\leftarrow$  3<sup>2</sup>P<sub>3/2</sub>, F=3) transition. To pump the  $4^2D_{5/2}$  state, the second laser (the Coherent 699 as for the 5S) is tuned to the 17575.410 cm<sup>-1</sup> (green) Na( $4^{2}D_{5/2}$ , F=4  $\leftarrow$   $3^{2}P_{3/2}$ , F=3) transition. Note that in pumping the  $5^2S_{1/2}$  and  $4^2D_{5/2}$  states, the 4S and 4P states are radiatively populated. This is discussed further in section D3, below.

Fig. 9.

The sodium atomic levels relevant to the optical pumping. The fine and hyperfine structure levels are not shown. The relative populations of the various sodium levels as calculated in reference 22 for optically pumping the  $Na(4^2D_{5/2})$  [Na(5<sup>2</sup>S<sub>1/2</sub>)] state.





XBL 861-80

Fig. 9

Another important question is what are the population distributions among the various hyperfine levels and sublevels excited? Also, what is the meaning of the hyperfine level and sublevel (F,  $m_F$ ) populations in terms of the excited electronic orbitals (L,  $m_L$ )? If the chemistry of the excited state is to be understood, it would be very useful to have a clear picture of the orbitals produced in the excitation. These pictures can be quantified in terms of the multipole moments of the distributions populated as discussed by Hertel and coworkers, <sup>14,15</sup> and reviewed below.

## 2. The Na D<sub>2</sub> Transition

Figure 10 shows the hyperfine levels of the states involved in pumping the Na  $D_2$  line. The quantum numbers L, S, and J have their usual meanings, but since  $^{23}$ Na has a nuclear spin I=3/2, the quantum number F, where

$$F = I + J \tag{9}$$

describes the total angular momentum. Fischer and Hertel have discussed the optical pumping of the Na  $D_2$  transition.<sup>14</sup> They came to the conclusion that in steady state 31% of the atoms could be in the excited state.<sup>16</sup> Since the 3P lifetime is only 16.9 nsec,<sup>13</sup> a cycling occurs in the optical pumping in which the atoms are excited and radiatively decay many times while in the laser field. In the experiments described here, the mean number of pumping cycles on the  $D_2$  transition ranges from 60 to 180, depending on the residence time



XBL 863-790

Fig. 10. Hyperfine structure of the  $Na(3^2P_{3/2} < 3^2S_{1/2})$  transition. Reproduced for reference 15 with permission.

Table IV.	Einstein A coefficients for the transitions	important in the
	optical pumping of the Na(3P,4D,5S) states.	The values of A
	given are in units of $10^6$ sec <sup>-1</sup> .	

Lower State ←	Upper State	Calculated <sup>13</sup> A	Experimental <sup>12</sup> A
3 <sup>2</sup> S <sub>1/2</sub>	3 <sup>2</sup> P1/2	58.9	61.8
3 <sup>2</sup> S <sub>1/2</sub>	3 <sup>2</sup> P3/2	59.1	62.2
3 <sup>2</sup> S <sub>1/2</sub>	4 <sup>2</sup> P1/2	2.85	2.81
3 <sup>2</sup> S <sub>1/2</sub>	4 <sup>2</sup> P3/2	2.85	2.81
3 <sup>2</sup> P1/2	4 <sup>2</sup> S <sub>1/2</sub>	8.26	8.9
3 <sup>2</sup> P1/2	3 <sup>2</sup> D3/2	41.3	45.3
3 <sup>2</sup> P1/2	5 <sup>2</sup> S <sub>1/2</sub>	2.29	2.6
3 <sup>2</sup> P3/2	3 <sup>2</sup> D3/2	8.26	9.0
3 <sup>2</sup> P3/2	3 <sup>2</sup> 05/2	49.5	54.
3 <sup>2</sup> P3/2	4 <sup>2</sup> S <sub>1/2</sub>	16.4	17.6
3 <sup>2</sup> P3/2	4 <sup>2</sup> D <sub>5/2</sub>	11.9	12.
3 <sup>2</sup> P3/2	5 <sup>2</sup> S <sub>1/2</sub>	4.58	5.2
3 <sup>2</sup> D3/2	4 <sup>2</sup> P1/2	0.15	-
3 <sup>2</sup> D3/2	4 <sup>2</sup> P3/2	0.016	-

Table IV, cont.

Lower State *	Upper State	Calculated <sup>13</sup>	Experimental <sup>12</sup>
3 <sup>2</sup> 0 <sub>5/2</sub>	4 <sup>2</sup> P3/2	0.14	
4 <sup>2</sup> S <sub>1/2</sub>	4 <sup>2</sup> P1/2	6.44	· _
4 <sup>2</sup> S <sub>1/2</sub>	4 <sup>2</sup> P3/2	6.47	-
4 <sup>2</sup> P1/2	5 <sup>2</sup> S <sub>1/2</sub>	1.70	
4 <sup>2</sup> 93/2	4 <sup>2</sup> D5/2	6.84	12.
4 <sup>2</sup> P3/2	5 <sup>2</sup> S <sub>1/2</sub>	3.38	-

of the Na atoms in the interaction region, and thus upon the Na beam velocity. Subsequent experiments by Hertel and coworkers have shown the optical pumping efficiency to be more on the order of 15-20%.<sup>17</sup> Hertel and Stoll used a rate equation approach to model the excitation process. Cohen-Tannoudji has treated a two level system more rigorously with a "dressed atom" model to take into account more fully the interaction of the atom with the field.<sup>18</sup> Cohen-Tannoudji has shown that for a two level system the rate equation method gives similar results. Hertel and Stoll have extended this to show that rate equations are useful in the case of pumping the Na D<sub>2</sub> transition.

Linearly or circularly polarized light can be used to optically pump this transition. For linearly polarized light, the selection rules have  $\Delta F = \pm 1, 0$ ,  $\Delta m_F = 0$  transitions allowed, except for the forbidden  $\Delta F = \Delta m_F = m_F = 0$  transition. For right (left) circularly polarized light only the  $\Delta F = \pm 1, 0$ ,  $\Delta m_F = \pm 1$  (-1) transitions are allowed. In fluorescence, the  $\Delta F = \pm 1, 0$ ,  $\Delta m_F = \pm 1, 0$  transitions are all allowed except for the forbidden  $\Delta F = \Delta m_F = m_F = 0$  transition.

The most stable situation for optically pumping to the 3P level is to pump only the F'=3  $\leftarrow$  F"=2 transition, because the selection rule  $\Delta F=\pm1,0$  forbids radiative decay into the unpumped F"=1 level. The ground state levels are spectrally distinct, but care must be taken not to excite into more than one hyperfine level of the excited state, since if the F'=2 level is excited, it will fluoresce down to the dark F"=1 level.

There are a number of factors contributing to the effective linewidth of the transition pumped. These are the natural linewidth of the transition, the power broadening due to the dynamic Stark effect, and the Doppler width under the conditions used.

Although Hertel and Stoll assumed that the transition was saturated, they made no effort to take power broadening into account. The power broadening can be determined by:<sup>19</sup>

 $\Delta v_{Power} = 1.381 \times 10^7 \langle \mu \rangle |S|^{1/2}$ (10) where  $\Delta v$  is the half width half maximum (HWHM) of the saturated line width (if  $\Delta v_{Power} \rangle$  the natural linewidth) in Hz,  $\mu$  is the transition dipole matrix element in Debye, and S is the Poynting vector (the laser power density in the interaction region) in W/cm<sup>2</sup>. The

Table V.	The transverse velocity and resultant Doppler broadening of
	the Na D2 transition for the seed gases used. All velo-
	cities are in cm/sec.

Seed Gas	<u>Na Beam Velocity</u>	Transverse Velocity	Doppler Broadened Linewidth
Helium	$3.0 \times 10^5$	$3.3 \times 10^4$	56 MHz
Neon	$1.6 \times 10^{5}$	$1.8 \times 10^4$	30 MHz
Argon	$1.1 \times 10^{5}$	$1.2 \times 10^4$	20 MHz

dipole transition matrix element can easily be derived from the Einstein A coefficient:<sup>20</sup>

$$|\mu|^2 = \frac{3hA}{64\pi^4 v^3},$$
 (11)

where  $\mu$  is the transition dipole matrix element in esu-cm ( =10<sup>18</sup> D), A is the Einstein A factor in sec<sup>-1</sup>, and  $\nu$  is the transition frequency in cm<sup>-1</sup>. From (11), for the D<sub>2</sub> line < $\mu$ >=6.2 D. For a typical power of 30 mW at the collision region focused to 3 mm. diameter, the power density is .33 W/cm<sup>2</sup>. This yields the power broadened line width of 50 MHz.

Since the transition is optically pumped perpendicular to the motion of the atomic beam only the transverse velocity of the beam is

important in determining the Doppler width. This is sometimes referred to as the second order Doppler effect. To reduce this effect to a minimum the height of the atomic beam at the collision region was reduced to its present value of 1.1 mm. The maximum transverse velocity is related to the beam velocity by a geometric factor, in our case this is just:

$$v_{\perp} = \pm (.011) v_{Na}$$
 (12)

The Doppler width (HWHM) of the transition is then:

$$\Delta v_{\text{Doppler}} = \frac{v_{\text{L}}}{c}.$$
 (13)

The values of  $v_{\perp}$  and  $\Delta v$  for the three commonly used Na beam velocities are given in table V. Note that since the transverse velocity is directly proportional to the distance from the center of the atomic beam, and since the laser comes up vertically into the interaction region, the Doppler shift is proportional to the height within the sodium beam. When the total broadening gets too large, it is possible to see (with the naked eye) the top and/or bottom of the sodium beam darken (less excitation results in less fluorescence). This has only been observed with the sodium seeded in helium, which gives the largest Doppler shifts.

The natural linewidth is only 5 MHz. The various contributions to the line width of the transition can be combined approximately as:

 $\Delta v_{\text{Total}} \cong \left[ \left( \Delta v_{\text{nat}} \right)^2 + \left( \Delta v_{\text{Power}} \right)^2 + \left( \Delta v_{\text{Doppler}} \right)^2 \right]^{1/2}.$  (14) This yields values from 54 to 75 MHz for the broadened linewidth. These values compare unfavorably to the separation of the F'=2 and F'=3 levels shown in figure 10. As Fischer and Hertel have pointed out, the deleterious effects of this broadening are reduced substantially by the selection rule forbidding  $\Delta F = \Delta m_F = m_F = 0$  transitions. That is the F'=2,  $m_F = 0$  transition is not allowed whereas the majority of the population in the case of linearly polarized light resides in  $m_F = 0$  (see below). Also, since the excited state fraction is optimized, and only (nearby) transitions to the red are to be avoided, the optimal laser frequency of the saturated transition would be to the blue (lower wavelength) of the center frequency of the transition, thus reducing the fraction pumped into levels other than F'=3.

It has been observed (particularly when the laser is focused) that reducing the power can increase both the fluorescence and the reactive signal due to the excited state. For each experiment the fluorescence and eventually the reactive signal are peaked on the dye laser power using neutral density filters. Under normal operating conditions, however, so long as the laser is not focused, no neutral density filters are required.

Fischer and Hertel have calculated the distribution of populations in the hyperfine sublevels ignoring whatever pumping to levels other than F'=3 occurs.<sup>14</sup> In the initial moments of pumping, the distribution reflects only the relative transition moments of the various up transitions. In the quasi-steady state situation the level populations have been redistributed over many pumping-fluorescence cycles. The distributions for linearly and right circularly polarized laser pumping of the transition are shown in figure 11, from reference 14 with

Fig. 11. Estimated relative populations in magnetic sublevels  $(m_F)$  of Na $({}^{2}P_{3/2}, F=3)$  on optical pumping with a) and b) a linearly polarized laser, and c) and d) a circularly polarized laser; a) and c) for inital moment pumping, and b) and d) for steady state pumping. Reproduced from reference 14 with permission.



XBL 863-789



permission. In the case of right (left) circularly polarized light, with the constant bias of  $\Delta m_F = +1$  (-1) with each pumping cycle, <u>all</u> the population ends up in  $m_F = (-)F$  for both the ground state level and the excited level pumped (the unpumped ground state level is of course unaffected).

From the hyperfine sublevel populations one can calculate the degree of alignment for linearly polarized light, or the degree of orientation for circularly polarized light. Following the notation of reference 14, the alignment can be quantified as:

$$a_0^{ph} = \sum_{m=-F}^{F} [3m^2 - (F(F^{+1}))] \sigma_m^{-1} / \sum_m \sigma_m^{-1}, \quad (15)$$

where  $a_0^{ph}$  is a measure of the alignment along the z axis in the photon frame of reference (along the electric vector of the laser), m is m<sub>F</sub>, and  $\sigma_m$  is the fraction in sublevel m. There is another measure of the alignment along the x axis with respect to the y axis which is 0 everywhere here. The parameter  $a_0^{ph}$  for F=3 varies between the limits of -12 and 15, where  $a_0^{ph}$ =-12 corresponds to the maximum alignment with all population in m<sub>F</sub>=0,  $a_0^{ph}$ =0 corresponds to equal population in each m<sub>F</sub> level, and  $a_0^{ph}$ =15 corresponds to minimum alignment with all population in m<sub>F</sub>=43 (and the maximum orientation if all of the population is in only one of these two levels). The orientation can be measured by:

$$\sigma_0^{\text{ph}} = \sum_{m=-F}^{F} m \sigma_m / \sum_m \sigma_m^{}.$$
(16)

The limits of parameter  $o_0^{ph}$  for F=3 are -3 and +3 which correspond to all population in m<sub>F</sub>=-3 and +3, respectively, and maximum orientation. The minimum orientation occurs if the sublevel population is symmetric about m<sub>F</sub>=0, then  $o_0^{ph}=0$ . This is the case any time linearly polarized light is used to excite the transition.

The alignment and orientation parameters can be converted into alignment and orientation parameters with respect to the electronic orbital angular momentum L which is a more chemically meaningful representation. The transformation for the alignment parameter is:

$$a_0^{ph}(L) = a_0^{ph} V_2(F,L) (2F+1)(2J+1) \begin{cases} F F 2 \\ J J I \end{cases} \begin{cases} J J 2 \\ L L S \end{cases} (-)^{F+2J+L+I+S}, (17)$$

where the bracketed quantities are 6-j coefficients, and

$$V_{2}(F,L) = (-)^{(2F+2L)} \left[ \frac{(2F-2)! (2L+3)!}{(2F+3)! (2L-2)!} \right]^{1/2} .$$
(18)

The transformation from F to L for the orientation parameter is:

$$o_0^{ph}(L) = o_0^{ph} V_1(F,L) (2F+1)(2J+1) \begin{cases} F F 1 \\ J J I \\ L L S \end{cases} \begin{pmatrix} J J 1 \\ L L S \\ - \end{pmatrix}^{F+2J+L+I+S}, (19)$$

where

$$V_{1}(F,L) = (-)^{(2F+2L)} \left[ \frac{(2F-1)! (2L+2)!}{(2F+2)! (2L-1)!} \right]^{1/2} .$$
(20)

From equation (15) for the  $3^{2p}_{3/2}$ , F=3 level, the conversion of the alignment parameter is  $a_0^{ph}(L) = a_0^{ph}/15$ . From equation (19), the conversion of the orientation parameter for this level is  $o_0^{ph}(L) = o_0^{ph}/3$ .

The fraction of the excited state population in the  $3p_z$  orbital (the orbital parallel to the electric vector of the laser) can be related to the  $a_0^{ph}(L)$  parameter in the case of linearly polarized laser pumping:

$$P(p_z) = \frac{1}{3} - a_0^{ph}(L)/3.$$
 (21)

If all the excited state population were in  $m_F=0$ , the fraction in the  $p_z$  orbital would be 0.6. The remainder of the excited state population is distributed equally between the  $p_x$  and  $p_y$  orbitals.

With the magnetic sublevel distributions calculated in reference 14 and shown in figure 11, the initial moment pumping with linearly polarized light gives  $a_0^{ph}=-7.2$  and  $a_0^{ph}(L)=-0.48$ , and steady state values of  $a_0^{ph}=-10$ ., and  $a_0^{ph}(L)=-0.67$ .<sup>14</sup> However, the  $a_0$  values can be related to the measured fluorescence dependence 2A from equation (6). This relationship is:

$$a_0^{ph} = (2A) / [g^{(2)}(1.5 - 2A)]$$
 (22)

where g<sup>(k)</sup> is:

$$g^{(k)} = (-)^{F'-F''} [F'(F'+1)]^{-1} \begin{cases} F' F' k \\ 1 & 1 & F'' \end{cases} / \begin{cases} F' F' k \\ 1 & 1 & F'' \end{cases}, \quad (23)$$

which for F"=2, F'=3 gives  $g^{(2)}=-1/15$ . Fischer and Hertel measured alignments up to  $a_0^{ph}=-9.4$ . In the experiments described here, the measured polarization dependence of the fluorescence for the 3S to 3P excitation was typically 2A=0.35. Assuming that fluorescence from only F'=3 is observed, equation (22) yields a value of  $a_0^{ph}=-4.6$ , or  $a_0^{ph}=-0.3$ . This implies that 0.43 of the excited state population is in the p<sub>z</sub> orbital vs. 0.28 in the p<sub>x</sub> and p<sub>y</sub> orbitals.

Fischer and Hertel show a sharp drop in alignment with increasing Na density above  $10^{11}$  atoms/cc.<sup>14</sup> Above  $10^{11}$  atoms/cc, radiation trapping occurs,<sup>15</sup> that is fluorescence photons are absorbed by other Na atoms, destroying the alignment produced by the laser field. The Na atom densities in the collision region in the experiments described here are in the range of  $10^{11}$ - $10^{12}$  atoms/cc, and this certainly affects the orbital alignment obtained.

No successful measurements were made of the orientation dependence of any of the reactions described in chapters II-V.<sup>21</sup> Also, no attempts were made to measure the degree of orientation. Fischer and Hertel predict that in steady state essentially all of the Na(3P) atoms can be pumped to  $F=m_F=3$ , giving  $o_0^{ph}=3$ , and  $a_0^{ph}=15$ . Experimentally, they found that the orientation can be as high as  $o_0^{ph}=2.9$ , and  $a_0^{ph}=14.7.^{14}$  Once again radiation trapping becomes significant above  $10^{11}$  atoms/cc, and the orientation drops off sharply above this density.

Thus, the atomic densities that were found necessary to conduct the reactive scattering experiments interfere with but do not preclude

the measurement of alignment and orientation effects upon the reactive differential cross sections of the Na(3P) state.

Since the signal due to the excited state is nearly inversely proportional to the fraction excited (see equation (5)), a larger assumed fraction excited gives a lower estimate of the differential cross section of the excited state. For this reason the fraction excited was assumed to take the value of its upper limit, 20%, 17 for the data analyzed in chapters II-V.

3. Subsequent Excitation to  $Na(4^2D_{5/2})$ 

As discussed above, a second laser can be used to excite the sodium atoms to the  $4^{2}D_{5/2}$  state. In analogy to the pumping of the Na( $3^{2}P_{3/2}$ , F=3  $\leftarrow 3^{2}S_{1/2}$ , F=2) transition, only the Na( $4^{2}D_{5/2}$ , F=4  $\leftarrow 3^{2}P_{3/2}$ , F=3) transition is pumped, and the other hyperfine levels are left unpumped. A schematic of the relevant energy levels is shown in figure 12.

Hertel and coworkers have performed quenching experiments with Na(4D,5S) atoms.<sup>22</sup> Faced with the same optical pumping problem, they opted for using circularly polarized light to effectively reduce this to a three level case. The three levels are of course  $m_F=F$  where F is 2, 3, and 4 for the 3S, 3P, 4D levels, respectively. This automatically isolates the F=4 level of the  $4^2D_{5/2}$  state from the lower F levels. However, there is nothing to prevent fluorescence to  $4^2P_{3/2}$ , F=3, instead of  $3^2P_{3/2}$ , F=3 (the branching ratio is 1.7:1 for 3P:4P). When this occurs, there is some scrambling of levels, as the 4P state



Fig. 12. Hyperfine structure of the  $Na(4^{2}D_{5/2} < 3^{2}P_{3/2})$  transition.

decays predominantly to the 4S state (the branching ratio is 2.3:1:0.06 for 4S:3S:3D), which then fluoresces to either  $3^{2}P_{3/2}$  or  $3^{2}P_{1/2}$  (the branching ratio is 2:1) before returning to the ground state (see figure 9). Of course once the  $4^{2}S_{1/2}$  or  $3^{2}P_{1/2}$  levels, or any hyperfine level that can lead to population of the unpumped  $3^{2}S_{1/2}$  level is reached, atoms will become unavailable for pumping. There are additional complications when other hyperfine levels of the  $4^{2}D_{5/2}$  state are pumped. Jamieson et al. have estimated the fraction in each state from the ratios of the Einstein A coefficients of Heavens (reference 13 and table V).<sup>22</sup> They assumed pumping with circularly polarized light. Their estimates are shown in figure 9.

Using linearly polarized light, the excited 4d orbital can be aligned much as the 3p orbital was. Using circularly polarized light, the 4d orbital can be oriented, but no orientation effects have been observed in the experiments described.<sup>21</sup> The interpretation of the alignment experiments is made difficult because in addition to the  $d_z^2$ orbital the  $d_{xz}$  and  $d_{yz}$  orbitals contribute to the alignment, and to the alignment parameter  $a_0^{ph}$ . No independent measure was made of the population of the  $d_{xz}$  and  $d_{yz}$  orbitals. There is no problem determining the multipole moments for the excitation as in the previous section. The maximum conceivable alignment occurs for all the excited state population in F=4,  $m_F=0$ . From (15) this gives  $a_0^{ph}=-20$ . For F=4, L=2, and J=5/2, equation (17) gives  $a_0^{ph}(L) = 3a_0^{ph}/14$ . For linearly polarized light the orientation is of course still zero;  $o_0^{ph}=0$ . The other limit of  $a_0$  is for  $m_F=\pm F$ , which yields  $a_0^{ph}=28$ . Once again if
all the population is in one or the other outermost  $m_F$  level, the maximum orientation is attained, and  $o_0^{ph} = \pm 4$ . This is attained with circularly polarized light, both exciting lasers must have the same sense. There is no confusion as to the meaning of the orientation due to the additional orbitals; when  $m_F = \pm 4$ ,  $m_I = \pm 4$ .

Cycling also occurs for the  $3P \Rightarrow 4D$  excitation, but the ratio of the lifetimes (derived from table V) shows that the cycling occurs more than three times slower for the upper transition (Rabi flopping via stimulated emission is not included in this since no redistribution of magnetic sublevel population takes place as a result of it). To a good approximation then, the distribution of magnetic sublevels in the  $4^{2}D_{5/2}$ , F=4 state should reflect those in the  $3^{2}P_{3/2}$ , F=3 state. Using the values for the steady state populations in the 3P state calculated by Hertel and Stoll<sup>15</sup> as the distribution in the  $4^{2}D_{5/2}$ , F=4 state gives an expected alignment of  $a_0^{ph} = -17.9$ , and  $a_0^{ph}(L) = -3.8$ . The measured polarization dependence of the fluorescence of the  $Na(4^2D_{5/2})$ state typically gave 2A=0.32. Using a narrow band pass filter only the green fluorescence of the Na( $3^{2}P_{3/2} \leftarrow 4^{2}D_{5/2}$ ) transition was monitored. For F'=4, and F"=3, from (23),  $g^{(2)}=-1/28$ , giving  $a_0^{ph}=-7.6$  from (22). This gives  $a_0^{ph}(L) = -1.62$ . As discussed above, the fraction in the  $d_z^2$ orbital cannot be determined from a<sub>0</sub>.

Experiments on reactive scattering have shown that having both lasers circularly polarized in the same sense yields an excited state fraction approximately 15% higher than having both lasers linearly polarized with their electric vectors parallel. If the two lasers are

circularly polarized with the opposite sense, the fraction excited is reduced by more than a factor of 3. Several of the experiments in chapters II-V on the reactive scattering of Na(4D) atoms were performed with both the linear and circular polarization optical pumping schemes, and no fundamental differences were observed between the cases. No polarization dependence of the reactive scattering at a single detector angle of greater than 17% was seen, so this should limit the effect on the angular distributions recorded.

For the reactive scattering experiments of the 4D state, it was not possible to separate the contributions of the radiatively populated 4S and 4P states. Thus, the angular distributions derived for the 4D state in general contain the contributions of these states. No experiments were done to directly excite the Na(4S,4P) states. For the distributions, it was assumed that (1) the excited state fraction in the 3P state remains unchanged when the second laser is added, (2) the ground state contribution to the reactive scattering of the ground state is negligible, and (3) the fraction excited to the 4D state is 80% that to the 3P state. Assumption (1) is based on the experimental observation that the fluorescence from the 3P does not change when the second laser excitation is modulated. Assumption (2) is an empirical observation. Assumption (3) is based on the estimates of Jamieson et al.<sup>22</sup> The scattering due to the Na(4D) state was determined as:

$$4D = (ON4D - OFF3P) / 0.8,$$
 (24)

where 4D is the signal due to the Na(4D) state, ON4D is the signal with

both lasers on, OFF3P is the signal with the green laser off and the yellow laser on (so that the 3P level is still excited). Compare this to (5).

4. Subsequent Excitation to  $Na(5^{2}S_{1/2})$ 

As in the 4D case, a second dye laser can be used to excite from the  $3^{2}P_{3/2}$  level to the  $5^{2}S_{1/2}$  level. The hyperfine levels involved in the excitation of the  $5^{2}S_{1/2}$  state are shown in figure 13. The  $5^{2}S_{1/2}$ level can fluoresce into both the fine structure states of the 3P and 4P states (the branching ratio from table V is 0.5:1:0.4:0.7 for  $3^{2}P_{1/2}:3^{2}P_{3/2}:4^{2}P_{1/2}:4^{2}P_{3/2}$ ). There is no simple way to avoid fluorescence into dark levels, but as pointed out by Jamieson et al.,<sup>22</sup> and observed here, there does not seem to be a large difference in the pumping efficiency of this level relative to pumping the 4D level.<sup>22</sup> Once again, there are no accurate determinations of the fraction excited. Linearly polarized lasers whose electric vectors are parallel are used to pump the transitions. Estimates by Jamieson et al. for the fraction in each state with excitation to the 5S state are shown as the bracketed quantities in figure 9.

As with the 4D state, it was not possible to separate the contributions to the reactive scattering by the radiatively populated 4S and 4P states. Thus, the distributions derived for the 5S state also contain the contributions of these states. The same assumptions regarding the excitation were made as for the excitation of the 4D state (see



Fig. 13. Hyperfine structure of the Na( $5^2S_{1/2} \leftarrow 3^2P_{3/2}$ ) transition.

above). It was assumed that the fraction excited to the 5S state is 80% that to the 3P state. With the replacement of "5S" for each "4D" equation (24) was used to determine the scattering of the 5S state.

## E. References

- M. F. Vernon, H. Schmidt, P. S. Weiss, M. H. Covinsky, and Y. T. Lee, <u>J. Chem. Phys.</u>, to be published.
- M. F. Vernon, Ph.D. Thesis, University of California, Berkeley (1983).
- 3. C. H. Becker, P. Casavecchia, P. W. Tiedemann, J. J. Valentini, and Y. T. Lee, <u>J. Chem. Phys. 73</u>, 2833 (1980).
- 4. Y. T. Lee, J. D. McDonald, P. R. LeBreton, and D. R. Herschbach, Rev. Sci. Inst. 40, 1402 (1969).
- P. E. Siska, J. M. Parson, T. P. Schaefer, and Y. T. Lee, <u>J. Chem.</u> Phys. 55, 5762 (1971).
- S. Gerstenkorn and P. Luc, <u>Atlas du Spectre d'Absorption de la</u> Molecule d'Iode, (Editions du CNRS, Paris, 1978).
- 7. For an explanation of the cross-correlation technique, see:
  G. Comsa, R. David, and B. J. Schumacher, <u>Rev. Sci. Instrum. 52</u>, 789 (1981).
- H. Schmidt, J. M. Mestdagh, M. H. Covinsky, B. A. Balko, and Y. T.
   Lee, unpublished results (1985).
- 9. T. T. Warnock and R. B. Bernstein, J. Chem. Phys. 49, 1878 (1968).
- R. J. Buss, Ph.D. Thesis, University of California, Berkeley, California (1979).

- 11. For Na(3P,4D) + HCl  $\Rightarrow$  <u>NaCl</u> + H at a collision energy of 3.4 kcal/mole a small decrease is observed in reactive scattering at large laboratory detector angles in going from 3P  $\Rightarrow$  4D, but this is the <u>only</u> case where a decrease has been observed on excitation.
- 12. W. L. Wiese, M. W. Smith, and B. M. Miles, <u>Atomic Transition</u> <u>Probabilities Volume II: Sodium Through Calcium</u>, Natl. Stand. Ref. Data Ser., Nat. Bur. Stand. (U. S.), Washington, DC (1969).
- 13. 0. S. Heavens, <u>J. Opt. Soc. 51</u>, 1058 (1961).
- 14. A. Fischer and I. V. Hertel, <u>Z. Phys. A 304</u>, 103 (1982).
- 15. I. V. Hertel and W. Stoll, <u>Adv. At. Mol. Phys. 13</u>, 113 (1978).
- 16. This is just one half of the fraction initially in F=2 of the ground state (5/8); so the excited state fraction would be 5/16.
- 17. H. Schmidt and I. V. Hertel, private communication.
- C. Cohen-Tannoudji, in <u>Atomic Physics 4</u>, 589, G. zu Putlitz, E. W.
   Weber, and A. Winnacker, eds. (Plenum, New York, 1975).
- J. C. Hemminger, R. Cavanagh, J. M. Lisy, and W. Klemperer,
   J. Chem. Phys. 67, 4952 (1977).
- 20. G. Herzberg, <u>Molecular Spectra and Molecular Structure I. Spectra</u> <u>of Diatomic Molecules</u>, 20 (Van Nostrand Reinhold Co., New York, 1950).
- 21. Recent measurements by H. Schmidt, J. M. Mestdagh, B. A. Balko, M. H. Covinsky, and Y. T. Lee on Na(4D) + NO<sub>2</sub> have shown an orientation dependence.

22. G. Jamieson, W. Reiland, C. P. Schulz, H.-U. Tittes, and I. V. Hertel, <u>J. Chem. Phys. 81</u>, 5805 (1984).

## II. THE REACTIONS OF GROUND AND EXCITED STATE ALKALI ATOMS WITH HYDROGEN HALIDE MOLECULES

## A. Introduction

The reactions of alkali atoms (M) with hydrogen halide molecules (HX) have been studied in great detail both experimentally<sup>1-12</sup> and theoretically.<sup>13-22</sup> In 1955, Taylor and Datz reported the first angular distribution measured for any reaction in a crossed molecular beams experiment for K + HBr.<sup>1</sup> Since then, much more quantitative experiments have been performed in order to determine the dynamics of these reactions as well as the effects of reagent translational,<sup>2-6</sup> vibrational,<sup>5-9</sup> and rotational<sup>8-10</sup> excitation on the reaction. Until recently, the effect of electronic excitation remained unexplored.<sup>11</sup>

An interesting feature of this family of reactions is that with the appropriate choice of alkali and halogen atoms, the reaction can be made exothermic, thermoneutral, or endothermic to allow tests of the effects of various forms of reagent energy in the threshold region. Of course, this does complicate attempts to make direct comparisons between results measured on different reactions (i.e. on different combinations of reactants). The  $\Delta H_0^0$  values for the M + HX  $\Rightarrow$  MX + H reactions for the various combinations of alkali and halogen atoms,

Table I	I Heats of reaction (ΔH <sub>0</sub> <sup>O</sup> ) in kcal/mole for M + HX ⇒ (M = Li, Na, K, Rb, Cs; X = F, Cl, Br, I) calculat the diatomic molecular dissociation constants give reference 23.					MX + H ed from n in
М:	Χ:	F	C1	·Br	I	
Li		-].	-8.8	-12.3	-10.6	
Na		13.	4.7	0.4	1.2	
К		18.4	2.1	-3.5	-5.9	
Rb		20.	3.	-3.3	-5.5	
Cs		17.3	-3.4	-9.5	-11.7	

calculated from spectroscopic and thermochemical dissociation constants,<sup>23</sup> are given in table I for reference.

At the Conference on Potential Energy Surfaces in Chemistry in 1971, Herschbach proposed that theoreticians calculate an <u>ab initio</u> potential energy surface for studying the dynamics of the prototypical Li + HF system.<sup>24</sup> A careful study in crossed molecular beams was performed for the reactions Li + HX  $\Rightarrow$  LiX + H (X = F, Cl) in order to provide a test for theoretical calculations.<sup>2</sup> Most of the theoretical studies of the reactions of alkali atoms with hydrogen halides have in fact concentrated on Li + HF because of its relative simplicity (i.e. Li, H, and F are all top row elements). The first calculation of the Li + HF system was performed by Lester and Krauss in 1970, who examined some of the features of the entrance channel, and

found a small potential well there.<sup>14</sup> A valence bond calculation by Balint-Kurti and Yardley (BKY) found a small entrance channel well and an exit channel barrier that was lowest for a bent configuration.<sup>15</sup> Included in BKY is a calculation of a few low-lying excited state surfaces including the one corresponding to Li(2P) + HF. Zeiri and Shapiro calculated semi-empirical potential energy surfaces (ZS), and found a small barrier that was between the entrance and exit channels for a collinear configuration, but moved into the entrance channel with an increasingly bent orientation. $^{16}$  They have since extended their calculations to provide potential energy surfaces for 14 of the reactions in the alkali + hydrogen halide family, and in a revised Li + HF surface (ZS II) found that the barrier was in the exit channel 17in agreement with BKY. They found that the barrier was lowest for a bent configuration in the M + HF potential energy surfaces, but was lowest for collinear and near collinear geometries in the M + HCl. HBr. and HI potential energy surfaces. They found a strong orientation dependence with a large barrier (>3 eV) for the collinear M-H-Xgeometry.<sup>17</sup> Chen and Schaefer have calculated an <u>ab initio</u> potential surface which shows an entrance channel well and an exit channel barrier which is Jowest for a bent transition state, much like that of BKY.<sup>18</sup>

Shapiro, Zeiri, and Pollack have done classical trajectory calculations on the ZS surface in which they found oscillatory behavior for the reaction probability as a function of collision energy.<sup>19</sup> This was attributed in part to the opening up of inelastic processes

such as HF vibrational excitation with increasing collision energy. NoorBatcha and Sathyamurthy have done quasiclassical trajectory calculations<sup>20-22</sup> on an analytic function fit to the Chen-Schaefer surface by Carter and Murrell,<sup>25</sup> and on the ZS surface. They looked at the effects of vibration, rotation, and orientation on reaction probability. They predicted a large ( $\cong$ 1 eV) vibrational threshold to reaction, which was <u>not</u> found in the earlier crossed molecular beam experiments (in which essentially all of the HF and HC1 molecules were in v=0).<sup>2</sup> They also predicted an initially decreasing then increasing dependence of reaction probability on rotational energy. This effect, although of a lower magnitude was seen in the experiments of Blackwell, Polanyi, and Sloan<sup>10</sup> described below.

The effect of varying translational energy was studied in the experiments of Becker et al. on Li + HF, HC1.<sup>2</sup> At a collision energy of 3 kcal/mole the Li + HF forms a long-lived complex (relative to a rotational period of the complex), because the bottom of the potential energy surface is sampled at this energy. At higher collision energies, such as 9 kcal/mole, this complex is not formed, and the reaction proceeds directly to products. The Li + HF reaction proceeds with a low cross section,  $\sigma_R < 1$  Å<sup>2</sup>, at these collision energies, yielding an average opacity of ~0.1. In Li + HC1, at low collision energy (3 kcal/mole) the LiC1 product is scattered backwards relative to the initial Li velocity, whereas at higher collision energy (9 kcal/mole) the LiC1 is scattered backwards-sideways. The Li + HC1

had higher reactive cross sections,  $\sigma_R > 25 \text{ Å}^2$  for the collision energies studied, yielding an average opacity of 1.

Brooks and coworkers have looked at the effect of translational energy on the K + HCl reaction.<sup>4,5</sup> They have found an increase in reaction probability as the collision energy is increased from 3 to 10 kcal/mole, and then a fall-off as the collision energy is further increased to 18 kcal/mole. They have concluded that translational excitation can never be as effective in promoting reaction as the excitation of a single quantum of vibration in HCl.

Heismann and Loesch have looked at the effect of translational energy on the K + HX(v=0,1)  $\Rightarrow$  KX + H (X = F, C1).<sup>6</sup> The drop-off of reaction probability for K + HC1(v=0) seen by Geiss et al. was not seen by Heismann and Loesch, who saw reaction probability rise steadily until reaching a constant value for collision energies from 21 kcal/mole up to 48 kcal/mole. This is shown in figure 1 which is adapted from reference 6 with permission. The open circles are the results of Brooks and coworkers, while the squares are the results of Heismann and Loesch. For the substantially endothermic K + HF(v=0)  $\Rightarrow$  KF + H, the reaction probability rises steadily for collision energies between 18 and 41 kcal/mole.

Lacmann and Herschbach have studied K + HCl at collision energies of 1-20 eV, however they only looked for the production of K\*(4P) and  $K^+$ .<sup>3</sup> They only saw K<sup>+</sup>, which was produced only above the threshold for the appearance of both K<sup>+</sup> and Cl<sup>-</sup>. Balint-Kurti and Yardley have

Fig. 1. Reaction probability vs. reagent translational energy from the reaction K + HCl(v=0) > KCl + H including the data of Brooks and coworkers,  $^{4,5}$  as open circles, and that of Heismann and Loesch,  $^{6}$  as triangles. This figure is adapted from figure 5 of reference 6, with permission.



Fig. 1

XBL 861-79

suggested that this is due to KCl molecules formed in very high vibrational levels which then dissociate to ion pairs.<sup>15</sup>

Odiorne, Brooks, and Kasper have studied the effect of reagent vibrational excitation on the reaction K + HCl(v=0,1)  $\Rightarrow$  KCl + H by exciting the HCl with a cw chemical laser.<sup>7</sup> They found a rise in the estimated reaction cross section from  $\sigma_R(v=0)=0.15 \ \text{A}^2$  to  $\sigma_R(v=1)=20 \ \text{A}^2$ , for a collision energy of 3-4 kcal/mole. This shows that vibrational excitation can effectively lift the M + HX reaction out of the threshold region, and "turn the reaction on."

Heismann and Loesch found similar results for K + HCl(v=0,1) at low collision energies (such as the one studied by Odiorne et al.), but saw a decrease in the vibrational enhancement, measured as:

$$\epsilon = \frac{\sigma_{\mathsf{R}}(\mathsf{v}=1)}{\sigma_{\mathsf{R}}(\mathsf{v}=0)},$$

with higher collision energy.<sup>6</sup> Their vibrational excitation was thermally produced by heating the hydrogen halide molecules in an oven prior to expansion in a molecular beam. Similar behavior was found for the reaction K +  $HF(v=0,1) \Rightarrow KF + H$ , but in this case the much larger endothermicity pushes the threshold of this process out to higher collision energies in the absence of vibrational excitation. It was not possible to reach collision energies at which  $\varepsilon$  fell below 70 (this was reached when  $E_{coll}=48.4$  kcal/mole). In both cases, the dropping off of the vibrational enhancement,  $\varepsilon$ , is due more to the increase in

(1)

 $\sigma_{\rm R}(v=0)$  with increasing collision energy, than to a large drop in  $\sigma_{\rm R}(v=1)$ , although there is a gentle decrease in  $\sigma_{\rm R}(v=1)$  at higher collision energies in the cases of both HCl and HF.

Dispert, Geis, and Brooks have studied the effect of rotation on the K + HCl(v=1,J)  $\Rightarrow$  KCl + H reaction by exciting specific rotational levels of the HCl(v=1) level with a cw chemical laser.<sup>9</sup> They found a factor of 2 decrease in the reactive cross section in going from J=1 to J=4. This is the anticipated effect, since in going from very low to moderate rotational excitation, the favored orientation for alignment is destroyed.

Along similar lines, Hoffmeister, Potthast, and Loesch have looked at the effect of rotation on the K + HF(v=1,J)  $\Rightarrow$  KF + H reaction for J = 5, 6, and 7.<sup>8</sup> They observed an overall increase in cross section with increasing rotational level. This effect is also as expected; moving from moderate to higher rotational levels there should be a gradual change from the troublesome destruction of favored orientation to a blurring of orientations due to fast rotation. Alternatively, this could be attributed to favorable excitation of bending in the transition state which could help overcome the barrier to reaction. A continued increase in cross section with further rotational excitation (and thus bending excitation) would be expected in an exact analogue to vibrational enhancement of reactivity.

Blackwell, Polanyi, and Sloan looked at the reactions Na + HX  $\Rightarrow$  NaX + H for HX = HF(v=1-4, J=0-14), and HX = HCl(v=1-4, J=0-19) in a

fascinating series of experiments that monitored the depletion of the infrared chemiluminescence of chemically produced HF and HCl by reaction with Na atoms. Thus, Blackwell et al. were able to compare the reactivities of the various rotational levels within each vibrational level. They saw an initial decrease in reactivity with increasing J, and then a subsequent increase. The minimum  $\sigma(J)$  occured at J $\cong$ 7 for HF, and J $\cong$ 11 for HCl for each vibrational level studied. Examples of the dependences they observed are shown in figure 2, reproduced with permission from reference 10. The explanation given for this behavior is, as above, a destruction of the proper orientation for reaction with increasing J at first, until after a certain point, the incoming atom sees a spatially averaged potential. The rotational energies achieved in this study are quite high. For HF,  $B_{a}=20.6$  $cm^{-1}$ , so that for J=14, the rotational energy is approximately 0.5 eV, or about one-fourth the energy of the first electronic excitation, or about the energy of one vibrational quantum. While there is an effect of rotation, it is certainly not "energy efficient" to try to . promote reaction in this way. Also, in general it is very difficult to rotationally excite molecules to such high rotational energy by any method other than this clever chemical production technique, as so many h of angular momentum are needed to reach the requisite J level. It should be noted that for Na + HF, a system discussed in section C2, no depletion, and thus no reaction, was observed for HF(v=1) which with a collision energy of >1.7 kcal/mole is energetically allowed to react to



XBL 861-78

- Fig. 2.
- Reaction probability vs. rotational level for the reaction Na + HCl(v=1-4) measured by the chemiluminescence depletion method by Blackwell, Polanyi, and Sloan, and reproduced with permission from reference 10.

NaF + H. This also gives credence to the assumption that depletion of the upper vibrational levels is not due to vibrational relaxation.

The first reaction of electronically excited sodium atoms was studied in our laboratory.<sup>11</sup> This was:

Na + HCl > NaCl + H  $\Delta H_0^0 = +4.7$  kcal/mole, (2) for which the reactive scattering of the ground  $(3^2S_{1/2})$  and first excited  $(3^{2}P_{3/2})$  states were studied. At a collision energy of 5.4 kcal/mole the ground state barely reacts due to the endothermicity of the process. The reaction of the Na(3P) gives the product NaCl center of mass contour map as shown in figure 3, adapted from reference 11. The product NaCl is predominantly backwards scattered in the center of mass relative to the incoming Na atoms with some forward-sideways This implies a direct reaction with collinear or nearpeaking. collinear Na-Cl-H approach geometries required for reaction. The product translational energy distribution was found to be broad and flat, remaining approximately constant from nearly zero translational energy up to all of the available energy going into translation. At high collision energy, 19.4 kcal/mole, there is still an increase in reaction with electronic excitation, but the enhancement is clearly lower. This effect is certainly the same as that seen by Heismann and Loesch for vibrational excitation, where enhancement is largest in (or below) the threshold region.



XBL 861-77

Fig. 3. NaCl center-of-mass contour map from the reaction of Na(3P) + HCl at a collision energy of 5.4 kcal/mole, as determined in reference 11.

The intention in performing the experiments described in section B was to enable a comparison between the effects of reagent electronic excitation on these reactions with the effects of other forms of energy already investigated. Specifically, how do the cross sections vary with electronic energy, and how are the reaction dynamics affected by varying electronic energy and symmetry?

## B. Results

Three systems from this family were chosen for a study of the effects of electronic excitation on these reactions. The reactions studied were Na with HF, HC1, and HBr. As discussed in chapter I, the collision energies were varied by changing the seed gases of the sodium beam from argon to neon to helium, and the Na(3P,4D,5S) states were optically pumped with one or two single-frequency cw dye lasers. The experimental apparatus and the treatment of data are described in detail in chapter I.

1. Na(4D, 5S) + HC1

When excited to the Na( $4^{2}D_{5/2}$ ,  $5^{2}S_{1/2}$ ) states and reacted with HC1 in crossed molecular beams, the NaCl product angular and velocity distributions change substantially from what was observed for the Na( $3^{2}P_{3/2}$ ) + HCl reaction. Angular distributions for Na(3S, 3P, 4D, 5S) + HCl > NaCl + H were measured at collision energies of 3.4, 5.6, and

16.3 kcal/mole. Newton velocity vector diagrams detailing the kinematics of the system at these three collision energies are shown in figures 4-6. The angular distributions themselves are shown in figures 7-9. The distributions shown are the raw data -- the 4D and 5S curves include the contributions of 3S and 3P as well as the radiatively populated 4S and 4P states.

Two dramatic effects are immediately apparent: the angular distributions have narrowed and the signals have increased in going from Na(3P) to Na(4D,5S). The unexpected narrowing of the NaCl angular distributions with increasing electronic energy is illustrated in figure 10, in which the angular distribution due to Na(3P) is blown up to the same scale as that due to Na(4D) for a collision energy of 5.6 kcal/mole. Time-of-flight measurements for Na(4D) + HCl at collision energies of 3.4 and 5.6 kcal/mole are shown for various detector angles in figures 11 and 12, respectively. Similar time-of-flight measurements for Na(5S) + HCl are shown in figures 13 and 14. The time-offlight distributions include the contribution of Na(3P) also, as well as the negative contribution due to the depletion of the ground state as discussed in reference 11. These measurements confirm that there is actually less translational energy in the NaCl due to the reaction of Na(4D,5S) than that due to Na(3P).

At a collision energy of 5.6 kcal/mole for Na(4D) + HCl, the dependence of the reactive signal at several detector angles upon the laser polarization was measured as described in chapter I. The derived distributions are shown in figures 15a-e and summarized in table II.

10 20 0 30 40 50 **5**2 -60 \4D 55 v Na -70 ЗP 1.0x10<sup>5</sup> cm/sec -80 -90 V<sub>HC1</sub>

8.6x10<sup>4</sup> cm/sec

------ 5.0x10<sup>4</sup> cm/sec

Fig. 4. Newton velocity vector diagram for Na(3S,3P,4D,5S) +  $HC1 \Rightarrow$ <u>NaC1</u> + H at a collision energy of 3.4 kcal/mole. The sodium beam is seeded in argon, the HC1 beam is neat.

83

E<sub>coll</sub>=3 kcal/mole

E<sub>coll</sub>=6 kcal/mole



└──── 5.0x10<sup>4</sup> cm/sec

XBL 863-858

Fig. 5. Newton velocity vector diagram for Na(3S,3P,4D,5S) + HCl  $\Rightarrow$ <u>NaCl</u> + H at a collision energy of 5.6 kcal/mole. The sodium beam is seeded in neon, the HCl beam is neat.



└──── 5.0x10<sup>4</sup> cm/sec

XBL 863-859

Fig. 6. Newton velocity vector diagram for Na(3S,3P,4D,5S) + HC1 → <u>NaCl</u> + H at a collision energy of 16.3 kcal/mole. The sodium beam is seeded in helium, the HCl beam is neat.



Fig. 7. NaCl angular distributions for Na(3S,3P,4D,5S) + HCl at a collision energy of 3.4 kcal/mole.







Fig. 9. NaCl angular distributions for Na(3S,3P,4D) + HCl at a collision energy of 16.3 kcal/mole.



Fig. 10. NaCl angular distributions for Na(3P,4D) + HCl at a collision energy of 5.6 kcal/mole from figure 8. The distribution arising from the reaction of Na(3P) is blown up so as to be able to compare the angular widths of the two distributions.



Fig. 11. NaCl product time-of-flight measurements for the reaction of Na(4D) + HCl at a collision energy of 3.4 kcal/mole for the laboratory detector angles shown on each frame.



Fig. 12. NaCl product time-of-flight measurements for the reaction of Na(4D) + HCl at a collision energy of 5.6 kcal/mole for the laboratory detector angles shown on each frame.



Fig. 13. NaCl product time-of-flight measurements for the reaction of Na(5S) + HCl at a collision energy of 3.4 kcal/mole for the laboratory detector angles shown on each frame.



Fig. 14. NaCl product time-of-flight measurements for the reaction of Na(5S) + HCl at a collision energy of 5.6 kcal/mole for the laboratory detector angles shown on each frame.

Fig. 15. Laser polarization dependences of the NaCl reactive signal due to the Na(4D) + HCl reaction at a collision energy of 5.6 kcal/mole for the angles shown in each frame. The solid lines are the best fits to the data as discussed in chapter I.



Polarization Angle (degrees)

Signal (arb. units)

XBL 863-895

Fig. 15
Laboratory Detector Angle	Laboratory Peak of Polarization Dependence (ø <sub>LAB</sub> )	Center of Mass Peak of Polarization Dependence (ø <sub>CM</sub> )	Amplitude (2A)
35°	-17°	13°	.084
40°	-22°	8 <b>°</b>	.062
44°	-25°	5°	.089
50 <b>°</b>	-22°	8°	.077
56 <b>°</b>	-34°	_4°	.089

Table II. A summary of polarization dependences found for the reaction  $Na(4D) + HCl \Rightarrow NaCl + H$  at a collision energy of 5.6 kcal/mole. These results are tabulated from figures 15a-e.

As seen in table II, at all detector angles measured the favored polarization angles for reaction are within a few degrees of the relative velocity vector of the system.

Simultaneously recorded polarization dependences for Na(3P) + HC1show no polarization dependence except for the farthest backwards scattering (at a laboratory detector angle of 56°) for which a 3% dependence (2A = 0.03) is observed favoring reaction when the laser polarization is parallel to the relative velocity vector. In previous measurements, no effect was seen for Na(3P) + HC1, but the minimum detectable effect for those measurements was 10–15%, so that it is not surprising that nothing was seen.<sup>11</sup> The upper limit of any polarization dependences for the other angles (30°-50°) for which none were observed could safely be put at 2%. 2. Na + HF

At collision energies up to 13.4 kcal/mole <u>no</u> reaction was observed for:

 $Na(3S,3P,4D,5S) + HF \Rightarrow NaF + H$   $\Delta H_0^0 = +13.$  kcal/mole. (3) Extensive signal averaging yielded no detectable NaF within the angular limits defined by its kinematics at m/e=42 (NaF<sup>+</sup>) nor at m/e=23 (Na<sup>+</sup>). As mentioned above, reaction of ground state sodium atoms with vibrationally excited HF has been observed in chemiluminescence experiments.<sup>10</sup>

3. Na(3S, 3P) + HBr

Angular distributions have been measured for the reaction:

 $Na(3S,3P) + HBr \Rightarrow NaBr + H$   $\Delta H_0^0 = +0.4 \text{ kcal/mole}$  (4) at collision energies of 3.0, 5.3, and 20.0 kcal/mole, as shown in figures 16, 17, and 18, respectively. Being very nearly thermoneutral in the absence of electronic and collision energy, a substantial amount of ground state reaction is observed for each collision energy studied. Still, there is a significant increase in the reactive cross section for the reaction when the Na atoms are excited to the 3P state.



XBL 863-879

Fig. 16. NaBr angular distributions for Na(3S,3P) + HBr at a collision energy of 3.0 kcal/mole.

98<sup>.</sup>



XBL 863-880





Fig. 18. NaBr angular distributions for Na(3S,3P) + HBr at a collision energy of 20. kcal/mole.

C. Analysis of Experimental Results

As described in chapter I, experimental angular and time-of-flight distributions were fitted using independent center of mass scattering angular and translational energy distributions using the program GM (a derivative of CMLAB). Detailed transformations of Na(4D,5S) + HCl to the center of mass frame were carried out in this way for the 5.6 kcal/mole data. The best fits to the measured angular distributions for Na(4D,5S) + HCl are shown in figure 19. The fits of the same center of mass distributions to the time-of-flight data are shown in figures 20 and 21 for Na(4D) + HCl and Na(5S) + HCl, respectively.

A comparison of the best fit translational energy distributions  $(P(E_T))$  distributions) at collision energies of 5.4 kcal/mole for Na(3P) in figure 22, and 5.6 kcal/mole for Na(4D) and Na(5S) in figure 23 points up this large change in the dynamics of the reaction. Note that with increasing electronic energy, the peaks and the mean energies in the P(E) distributions move to lower energy, with  $E_T^{peak} = 16$ , 7, and 4 kcal/mole and  $\langle E_T^{-} \rangle = 20$ , 7, and 7 kcal/mole for the reactions of the 3P, 5S, and 4D, respectively. The tails of the recoil energy distributions in all cases go out to the maximum possible energy available.

The calculated center of mass angular distributions  $(T(\Theta))$  shown in figure 24 for the product NaCl due to Na(5S) and Na(4D), at 5.6



XBL 863-761

Fig. 19. Angular distributions for Na(4D,5S) + HCl at a collision energy of 5.6 kcal/mole. The solid lines are the best fit to the data generated from center-of-mass angular and translational energy distributions.

Fig. 20. Time-of-flight distributions for Na(4D) + HCl at a collision energy of 5.6 kcal/mole. The solid lines are the best fit to the data generated from center-of-mass angular and translational energy distributions. The distributions are for the detector angles shown in each frame.



Fig. 20

XBL 863-756



XBL 863-754

Fig. 21. Same as figure 20, but for Na(5S) + HC1.



XBL 863-851

Fig. 22. Best fit center-of-mass product translational energy distributions for Na(3P) + HCl at a collision energy of 5.4 kcal/mole, from reference 11.



Fig. 23. Best fit center-of-mass product translational energy distributions for Na(4D,5S) + HCl, at a collision energy of 5.6 kcal/mole.



XBL 863-762

Fig. 24. Best fit product NaCl center-of-mass scattering angular distributions at a collision energy of 5.6 kcal/mole for Na(4D,5S) + HCl.

kcal/mole show strong backwards peaking with slight forward sideways peaking. This is similar to the  $T(\Theta)$  found for the reaction of the Na(3P) atoms.<sup>11</sup> The features of the two upper state  $T(\Theta)$  distributions are sharper and exhibit a sharp rise at approximately 90° for each.

The relative integrated cross sections for the reaction at 5.6 kcal/mole substantiate the idea that the cross section increases with increasing electronic energy. The ratios of integrated cross sections for the excited electronic states are:

3.5 (4D): 2.4(5S): 1.0 (3P) at 5.6 kcal/mole. (5) A relative optical pumping efficiency of 0.8 for the upper state (4D or 5S) as compared to the 3P state is assumed. That is, the fraction of atoms in the Na(3P) state is assumed unchanged when the second laser is turned on, and the fraction of atoms in the Na(4D) or Na(5S) state is assumed to be 80% of the fraction in the Na(3P) state as discussed in chapter I.

Assuming 20% efficiency in optically pumping the Na(3P) state from the ground state, the following ratios of reactive cross sections are obtained:

90 (3P): 1 (3S) at 5.4 kcal/mole, and (6)

6 (3P): 1 (3S) at 16.3 kcal/mole. (7)

The 20% pumping efficiency estimate is quite liberal, particularly in the most difficult to pump case of Na seeded in He (16.3 kcal/mole data) as discussed in chapter I. If the pumping efficiency were lower, it would imply larger enhancements than those reported above. Also, by

lowering the collision energy below the endothermicity of the reaction (turning off the ground state reaction), the enhancement can be made arbitrarily large as long as the Na(3P) still reacts. At 3.4 kcal/mole, this ratio is in fact infinite because of the lack of reaction of Na(3S).

The absolute reactive cross sections can be estimated by comparison with small angle elastic scattering.<sup>27</sup> This method eliminates the need to know the reactant beam number densities at the interaction region -- difficult quantities to determine accurately. Also, since both elastically scattered Na and reactively scattered NaCl are observed at m/e=23 (as Na<sup>+</sup>), only the relative probability of ionization of the two to Na<sup>+</sup> needs to be known, rather than the exact ionization efficiencies and the transmission functions of the quadrupole mass spectrometer at two different masses. In order to make the comparison between elastic and reactive contributions to the m/e=23 angular distributions, the absolute differential elastic cross section must be calculated. Then, one or more small angles in the measured ground state angular distribution at which only elastic scattering is observed are used to calibrate the angular distributions for each electronic state at that collision energy and the integrated reactive cross sections derived from them.

If large impact parameter collisions do not lead to chemical reaction, the absolute differential cross section for small angle

elastic scattering can be approximated for a spherically symmetric potential:

$$\left(\frac{d\sigma_{el}}{d\omega}\right)_{abs} = 0.239 \left(\frac{C_6}{E}\right)^{1/3} e^{-7/3}, \qquad (8)$$

where  $C_6$  is the coefficient for  $r^{-6}$  in a van der Waals potential in kcal/mole  $A^6$ , E is the collision energy in kcal/mole, and  $\Theta$  is the center of mass scattering angle in radians.<sup>28</sup> Equation (8) is the small angle scattering limit of the classical scattering due to a van der Waals interaction  $V(r) = C_6 r^{-6} \cdot {}^{28}$  This approximation is valid where sin $\Theta \simeq \Theta$ . The  $C_6$  parameter is not known for Na-HCl but it can be estimated. The  $C_6$  parameter will have two main contributions -- a dispersion portion and an induction portion -- with

$$C_6 = C_{6, \text{disp}} + C_{6, \text{ind.}}$$
 (9)

The dispersive portion can be estimated by the Slater-Kirkwood approximation:<sup>27,29-30,2</sup>

$$C_{6,disp} = (241.55 \text{ kcal/mole } \mathbb{A}^{3/2} \text{ a.u.}) \left( \frac{3}{2} \frac{\alpha(\text{Na})\alpha(\text{HC1})}{[\alpha(\text{Na})]^{1/2} + [\alpha(\text{HC1})/N(\text{HC1})]^{1/2}} \right)$$
(10)

where  $C_{6,disp}$  is in kcal/mole Å<sup>6</sup>,  $\alpha$  is the atomic or molecular polarizability in Å<sup>3</sup>, and N is the effective number of electrons. The polarizability of ground state Na is 24.5 Å<sup>3</sup>.<sup>31</sup> The polarizability of HCl is 2.63 Å<sup>3</sup>.<sup>32</sup> The effective numbers of electrons are: N(Na)=1, and N(HCl)=8. The value of  $C_{6,disp}$  obtained in this way is expected to be accurate to  $\pm 50\%$ .<sup>33</sup> Equation (10) yields a value of C<sub>6</sub>, disp = 4230 kcal/mole Å<sup>6</sup>. This compares favorably with the C<sub>6</sub> parameter measured for the isoelectronic Na-Ar ground state potential -- C<sub>6</sub>(Na-Ar) = 5500 kcal/mole Å<sup>6</sup>.<sup>34</sup> From (10), and the atomic polarizability of Ar  $\alpha(Ar)=1.6$  Å<sup>3</sup>,<sup>35</sup> the estimated C<sub>6</sub> would be C<sub>6</sub> = 2630 kcal/mole Å<sup>6</sup>, which is off by more than a factor of 2.

The inductive portion of  $C_6$ , the dipole induced dipole contribution, can be estimated from the Debye equation:<sup>36</sup>

 $C_{6,ind} = (14.394 \text{ Å}^3 \text{ kcal/mole}-D^2) \alpha(\text{Na}) \mu^2(\text{HCl}),$  (11)

where  $\mu$ (HCl) is the permanent dipole moment of HCl, 1.08D,<sup>37</sup> and once again  $\alpha$  is in Å<sup>3</sup>. Equation (11) yields a value of C<sub>6,ind</sub> = 411 kcal/mole Å<sup>6</sup>.

Thus from equations (9), (10), and (11),  $C_6 = 4640 \text{ kcal/mole } \text{Å}^6$ . The value obtained is expected to be limited by the accuracy of its dominant component  $C_{6,\text{disp}}$  to  $\pm 50\%$ .<sup>33</sup> However, the absolute differential cross section is slowly varying with  $C_6$  (as  $C_6^{1/3}$ ) so that the error induced in  $d\sigma_{el}$  by  $C_6$  is then only 20%.

The lowest laboratory scattering angle measured at a collision energy of 5.6 kcal/mole for (ground state) Na(3S) + HCl was 15°. Using the peak beam velocities, the angles and velocities were determined to carry out the center of mass to laboratory frame transformation as described below. Note that a substantial error is introduced by only using one Newton diagram in this manner. The center of mass scattering angle corresponding to  $\theta_{1ab}$ =15.0° is  $\theta_{cm} = 25.1° = 0.438$  radians. In order to convert from center of mass product flux  $(I(\Theta_{cm}))$  to laboratory frame number density  $(N(\Theta_{lab}))$ , the transformation:

$$N(\Theta_{ab}, v) = \frac{v}{u^2 |\cos \delta|} I(\Theta_{cm}, u)$$
(12)

is used.<sup>38</sup> The parameters are: the laboratory frame elastically scattered Na velocity v=1.791 km/sec, the center of mass frame elastically scattered Na velocity u=1.108 km/sec, and the detector viewing angle (the angle between the vectors u and v)  $\delta$ =18.2°.<sup>38</sup> Of course for elastic scattering only one center of mass velocity u is possible for a given collision energy. Thus all elastically scattered Na atoms fall precisely on the Newton sphere, and the laboratory velocity is determined by the scattering angle. Using equation (8) at a center of mass angle of 25.1°, the estimated value of the differential cross section is 15.5 Å<sup>2</sup>/steradian. Equation (8) is used to calculate the center of mass angular distribution in this manner. Then, a least squares fit to the low angle elastic scattering is carried out to determine the conversion between the relative and absolute differential cross sections.

The same calibration can be made for the elastic scattering of the Na(3P) atoms by HCl. The experimental value for the polarizability of Na( $3^2P_{3/2}$ ) is  $\alpha$ =53.6 Å<sup>3</sup>.<sup>39,40</sup> From (10), C<sub>6,disp</sub> = 6470 kcal/mole Å<sup>6</sup>, and from (11) C<sub>6,ind</sub> = 900 kcal/mole Å<sup>6</sup>. This gives a total C<sub>6</sub> = 7370 kcal/mole Å<sup>6</sup>. All angles and velocities are the same as for the ground state atoms, so that at 15° in the laboratory frame, with the

same assumptions as for the ground state scattering, from (8) the estimated value of the differential scattering is 18.1  $Å^2$ /steradian.

It is then necessary to know the relative detection efficiency of Na atoms and NaCl atoms with the mass spectrometer set at m/e=23 (Na<sup>+</sup>). This is only dependent upon the ionization cross section of Na bombarded by 200 eV electrons, and by the cross section of the fragmentation and ionization of NaCl to Na<sup>+</sup> when bombarded by 200 eV electrons ( $\sigma$ '). The ionization cross section of Na atoms has been measured as 2.1 Å<sup>2</sup> for 200 eV electrons.<sup>41</sup> The fragmentation ratio of NaCl upon ionization by 200 eV electrons has also been measured as:

 $\frac{[Na^+]}{[NaC1^+]} = 1.38$ , or 58% fragmentation.<sup>42</sup> Note that for much of the

vibrationally excited NaCl<sup>‡</sup> produced (from Na<sup>\*</sup> + HCl), this ratio is significantly higher.<sup>11</sup> However, the ionization cross section of NaCl has not been measured. Following the method described in reference 2, the cross section for ionization of NaCl by 200 eV electrons can be estimated by assuming that the ionization cross section of NaCl varies with electron energy as does that of argon atoms. The ratio of their ionization cross sections at a given electron energy is taken to be equal to the ratio of their polarizabilities. The polarizability of argon is  $\alpha(Ar)=1.6$  Å<sup>3</sup>.<sup>35</sup> The polarizability of NaCl can be estimated from the polarizabilities of Na<sup>\*</sup> and Cl<sup>-</sup> by:

$$\alpha(\text{NaC1}) \approx \alpha(\text{Na}^+) + \alpha(\text{C1}^-).^{36}$$
(13)

The polarizabilities of Na<sup>+</sup> and Cl<sup>-</sup> within the NaCl molecule at the equilibrium atomic separation have been calculated by Brumer and Karplus to be  $\alpha(Na^+)=0.239$  Å<sup>3</sup> and  $\alpha(Cl^-)=2.507$  Å<sup>3</sup>.<sup>43</sup> Thus from equation (13), the NaCl polarizability is  $\alpha(NaCl)=2.746$  Å<sup>3</sup>. The ionization cross section of argon for 200 eV electrons is 2.46 Å<sup>2</sup>,<sup>44</sup> and the ionization cross section of NaCl is then estimated to be 2.46x2.746/1.6 = 4.22 Å<sup>2</sup>. From the fragmentation ratio above, the ionization cross section of NaCl to Na<sup>+</sup> is:  $\sigma' = 4.22x0.58 = 2.5$  Å<sup>2</sup>. For vibrationally excited NaCl<sup>+</sup>, the fragmentation ratio is assumed to be 90%, and the ionization cross section of NaCl<sup>+</sup> to Na<sup>+</sup> becomes 4.22x0.9 = 3.8 Å<sup>2</sup>.<sup>45</sup> In any case the values of the ionization cross section of NaCl to Na<sup>+</sup> are probably only good to within a factor of 2, and they limit the accuracy of the value of the total cross sections derived below. The relative detection efficiency (Eff) of NaCl and Na at m/e=23 is the ratio of their ionization cross sections to Na<sup>+</sup>, or:

$$\frac{\text{Eff(NaC1)}}{\text{Eff(Na)}} = \frac{\sigma_{+}^{+}(NaC1)}{Na}.$$
(14)
$$\sigma_{+}(Na)$$
Na

For vibrationally cold molecules this is estimated to be 1.2, while for vibrationally hot molecules this is estimated to be 1.8.

The ratio of absolute to relative differential center of mass flux of elastically scattered Na is determined from the elastic signal count rate at small angles, and the transformation Jacobian in equation (12) for both ground state and Na(3P) scattering. Then by taking into account the relative detection efficiencies of Na and NaCl at m/e=23,

Table III. Estimated total reaction cross sections for Na(3S,3P,4D,5S) + HCl at a collision energy of 5.6 kcal/mole from the elastic scattering of Na(3S,3P) atoms and estimates of the absolute cross sections from Van der Waals attraction (VdW, at small angles), and hard sphere collisions (HS, at large angles).

Electronic State	Reactive Cross Sections: <u>from Na(3S) VdW</u> from Na(3P) VdW from Na(3S) HS			
3S	0.072 Å <sup>2</sup>	0.034 Å <sup>2</sup>	0.081 Å <sup>2</sup>	
3P	4.5 Å <sup>2</sup>	2.1 Å <sup>2</sup>	5.0 Å <sup>2</sup>	
4D	16 Å <sup>2</sup>	7 Å <sup>2</sup>	17 Å <sup>2</sup>	
5S	11 Å <sup>2</sup>	5 Å <sup>2</sup>	12 Å <sup>2</sup>	

the factor converting the integrated relative reactive cross sections for each electronic state measured into absolute total reactive cross sections ( $\sigma_R$ ) is determined. The cross sections determined have an accuracy of no better than a factor of two. The cross sections for the reaction at a collision energy of 5.6 kcal/mole are given in table III. The ground state cross section is so small that there should be no significant depletion in the elastic scattering measured at small angles. The elastic scattering of the Na(3P) atoms could be depleted by reaction. This would give an underestimate of the reactive cross sections. No calibration could be made at a collision energy of 16.3 kcal/mole, because the kinematics at this collision energy move the reactive scattering to lower angles, and thus even at the lowest angle recorded no elastic scattering is observed. Because there is a significant increase in background at angles near the sodium beam, it is impossible to measure closer to the beam.

The absolute cross section can also be estimated by comparison with large angle elastic scattering which is attributable to low impact parameter collisions. A hard sphere collision is assumed which will lead to isotropically scattered products if the impact parameter of the collision is smaller than the sum of the spheres' radii. The sum of the radii is equivalent to the classical turning point of the Na-HCl potential curve, which by analogy to  $He^*(2^1S) + Ar^{46}$  is estimated to be 3 Å at a collision energy of 5 kcal/mole. This gives a hard sphere cross section of 28 Å<sup>2</sup>. The values of the absoluted reactive cross section obtained in this way are shown in the right hand column of table III, and agree well with the values obtained from the small angle elastic scattering of ground state sodium atoms.

## D. Discussion

The Mechanism and Distribution of Product Energy in the Na(4D,5S) +
 HCl Reaction

The dominant features of the Na<sup>\*</sup> + HCl scattering are the decreasing product translational energy and the increasing reactive cross section with increasing electronic energy. Most excess energy must necessarily go into NaCl vibration and rotation, since there are no available electronically excited NaCl states and relatively little was found in product translation. Conservation of energy and angular momentum arguments can be used to determine where the internal energy is deposited.

While an electron jump mechanism has never before adequately explained the dynamics of an alkali-hydrogen halide reaction, <sup>10</sup> the very low ionization potential of the highly excited alkali atoms could increase the effect of such a process. HCl<sup>--</sup> is known to be dissociative. <sup>47</sup> HCl is dissociated by slow electrons, and dissociative electron attachment measurements yield an estimated (negative) vertical electron affinity of -0.815 eV.<sup>47</sup> Thus if an electron transfers from the sodium atom to the HCl at long range, the HCl<sup>-</sup> could quickly dissociate. Since the H atom is so much lighter than the Cl<sup>-</sup> ion, the H atom will rapidly depart, and

$$v_{\rm H} = -35 v_{\rm Cl} -$$
(15)

if the effect of the Na<sup>+</sup> is ignored. The Na<sup>+</sup> and the Cl<sup>-</sup> ions are then left at a separation approximately equal to the electron jump radius. The NaCl molecule formed is confined to travel in the direction opposite to the direction of departure of the H atom in the center of mass coordinate, because momentum must be conserved. Also, the velocity of the NaCl is severely restricted by conservation of momentum since the momenta of NaCl and of H must be equal and opposite in the center of mass, giving:

$$I_{\rm H} = -58v_{\rm NaCl}^{\circ}$$
(16)

This early departure of the H atom is able to explain the low product recoil energy found for the NaCl product.

If an electron jump mechanism is assumed, then the largest impact parameter leading to reaction would be the same as the position of the crossing of the neutral and ionic potential energy curves. By neglecting the long range neutral attractions, this can be estimated by:

$$IP - EA = \frac{e^2}{r_c}, \qquad (17)$$

where IP is the ionization potential of the appropriate sodium level, EA is the electron affinity of the HCl molecule, and  $r_c$  is the crossing radius of the neutral and ionic curves.<sup>48</sup> Equation (17) can be rewritten:

$$r_{c} = \frac{14.35}{IP - EA}$$
, (18)

where  $r_c$  is in Å, and the ionization potential and electron affinity are in eV. For the calculations below, the electron affinity of HCl is taken to be -0.815 eV.<sup>47</sup> The well-known Na atomic energies<sup>49</sup> are used to calculate the ionization potentials of the particular energy levels excited. For the four electronic states of sodium studied, this yields the values of  $r_c$  shown in table IV.

As pointed out in reference 11, the electron transfer radii for the Na(3S,3P) states are too small to be able to separate the breakage of the HCl bond (by electron transfer) from the formation of the NaCl bond. By the time the electron transfers from the alkali atom, the H atom is already feeling the repulsion of the closed shell NaCl molecule that is being formed. With the larger electron transfer radii of the

Na(nL)	IP[Na(nL)]	r <sub>c</sub>	۳r <sub>c</sub> <sup>2</sup>
3S	5.138 eV	2.41 Å	18 Å <sup>2</sup>
3P	3.033 eV	3.72 Å	47 Å <sup>2</sup>
4D	0.854 eV	8.60 Å	232 Å <sup>2</sup>
· 5S	1.021 eV	7.82 Å	192 Å <sup>2</sup>

Table IV.	Neutral-ionic	curve crossing radii	for various Na electronic
	states in the	Na + HCl reaction.	·

Na(4D,5S) states, this is no longer the case. The H atom feels only the repulsion of the closed shell  $Cl^-$  atom.

The electron transfer cross section ( $\sigma_e$ ) is expected to be:

$$\sigma_{\rm e} = \pi r_{\rm c}^2 . \tag{19}$$

The values of  $\sigma_{p}$  are given in the right-hand column of table IV.

The center of mass NaCl angular distributions derived for the reactions of Na(4D,5S) with HCl lead to the conclusion that the opacity function (the reaction probability as a function of impact parameter) is not P(b)=1 for all  $b < r_c$ . That is, not all encounters close enough for electron transfer lead to chemical reaction. The strong backwards scattering implies that the necessary configuration for reaction (of Na(4D,5S)) is Na-Cl-H, although the wider center of mass angular range in the backward hemisphere shows that the configuration is less constrained to be nearly collinear than was the reaction configuration of Na(3P) + HCl.

For Na(3P) + HCl the NaCl formed is due to an impulsive interaction which could impart a significant amount of momentum to the departing H atom.<sup>11</sup> This puts an upper limit on the impact parameter that leads to reaction which is certainly less than the 3.7 Å predicted by the neutral-ionic curve crossings. Also, the clear dependence of the reaction upon the relative orientation of the three atoms in the reaction of the Na(3P) atoms limits the value of  $\sigma_R^{3P}$  which is in agreement with the values of 2-5 Å<sup>2</sup> found (see table III). The values of  $\sigma_R^{3P}$ ,  $\sigma_R^{4D}$ , and  $\sigma_R^{5S}$  shown in table III are all significantly lower than the electron transfer cross sections given in table IV. Also, from the relative cross sections it is clear that the cross sections do not go up as the square of the crossing radii. This is indicative of what is already known from the center of mass angular distributions the opacity function P(b) is significantly less than 1 for b < r<sub>c</sub>.

Evidence that reaction takes place at distances larger than the impact parameter of the reactive collisions comes from the polarization dependences. All favored laboratory polarization angles ( $\phi_{LAB}$ ) are near the angle of the relative velocity vector as shown in figure 15 and table II. This is made clear in the  $\phi_{CM}$  values, which are the angles between the favored laboratory polarization angles and the angle of the relative velocity vector of the system in the laboratory frame of reference. This is in sharp contrast to the lack of polarization dependences seen for Na(3P) + HCl. In the case of Na(3P) + HCl at a collision energy of 5.4 kcal/mole the only polarization dependence observed was for the farthest backwards scattered NaCl, corresponding

to the most collinear Na-Cl-H geometries, and this was only a (2A = ) 3% effect. All smaller laboratory angles had polarization dependences of less than the measurable limit of 2%. In the limiting case of  $C_{\infty V}$  (collinear) geometry the required symmetry for reaction to form the ionic NaCl is  $^{2}\Sigma$  since the ion pair derives this symmetry from:

$$Na^{+}(^{1}S) + HC1^{-}(^{2}\Sigma)$$
 (20)

Because of the ground state  $HC1(X^2\Sigma^+)$  symmetry, all the electronic angular momentum of the Na(4D) must be along the Na-C1-H axis -requiring the  $4d_z^2$  orbital to be along the Na-C1-H axis. For zero impact parameter collisions, this requires the  $4d_z^2$  orbital to be aligned along the relative velocity vector of the system.

If this collinear geometry and  ${}^{2}\Sigma$  symmetry were a strict requirement for reaction, then as the impact parameter was raised the Na-Cl-H orientation (axis) required for reaction would tilt with respect to the relative velocity vector. In the impulsive reaction limit, this would lead to lower center of mass (NaCl) scattering angles and thus lower laboratory scattering angles. Also, this would imply that the favored polarization for reaction would rotate with respect to the relative velocity vector to match the Na-Cl-H axis. Of course, since different impact parameters would lead to different reactive orientations and different scattering angles, the polarization dependence would vary as a function of laboratory scattering angle. This is <u>not</u> observed. Therefore, the assumption of collinearity and/or the impulsive approximation that the scattering angle depends upon the impact parameter are incorrect if larger impact parameters lead to reaction.

From the center of mass angular distributions the relative orientation of the Na-C1-H was shown to be important. Since the reaction probability depends strongly on the Na-C1-H angle, similar arguments to those above can be used to show that as the impact parameter is increased, the range of reactive orientations will change with respect to the relative velocity vector, as well as the NaCl scattering angles. By considering the united atom (diatomic-like) approximation in which the HCl molecule shrinks to an Argon  $(^{1}S)$ atom, the correct symmetry for electron transfer remains  ${}^{2}\Sigma$ . Electron transfer from the Na(4D) atom to the HCl requires only that the 4d\_2 orbital be aligned toward the (reduced) HCl molecule (i.e. toward the Cl atom which is nearly at the center of mass of the molecule). Once again the polarization dependence would vary as a function of the laboratory scattering angle. Since the variation is not observed, the assumption that larger impact parameter collisions lead to reaction must be false.

If the Na(4D) + HCl reaction only takes place for small impact parameter collisions, two questions remain. First, how small does the impact parameter need to be in order to observe no variation of the polarization dependence with laboratory scattering angle? Secondly, if only small impact parameter collisions lead to reaction, how is the scattering angle determined?

To answer the first question, it is again useful to consider the united atom approximation. Here the center of mass frame favored polarization angle for electron transfer can be described:

$$\phi_{\rm cm} = \sin^{-1} \frac{b}{r_{\rm c}} , \qquad (21)$$

where b is the impact parameter of the collision and  $r_c$  is the crossing radius of the neutral and ionic potential curves -- the distance at which electron transfer is expected to occur. To obtain  $\phi_{\rm Cm} \cong 0^\circ$ , it is necessary to have  $b << r_c$ . Therefore this is the restriction on the impact parameters for reactive collisions -- b must be small with respect to the distance at which electron transfer takes place, 8.6 Å for Na(4D) + HC1.

To answer the second question, recall that HCl is dissociated by slow electrons. The (dissociative)  $HCl^{(-)}$  orientation at the time of electron transfer is then the dominant determinant of the NaCl product scattering direction (due to conservation of momentum).

Further evidence that only low impact parameters lead to reaction comes from the analysis below. An electron transfer can occur at each crossing point of the neutral and ionic potential curves. While the earliest crossing is at 8.6Å for Na(4D) + HCl, the final crossing is also at 8.6Å -- after the Na atom has gone by the molecule. With the opposite orientation of HCl to what normally leads to reaction -- that is initially Na-H-Cl -- after the atom and molecule pass by each other, the orientation becomes Na-Cl-H, perfect for reaction by the arguments above, assuming the excitation has survived to this point. This would lead to forward scattered NaCl as the H atom would be ejected in the backwards direction. This forward scattering is <u>not</u> observed. Therefore one or both of the two underlying assumptions in this argument are

incorrect: the excitation does <u>not</u> survive the collision and/or impact parameters allowing such a "fly by" cannot lead to reaction. Since a large increase in elastic scattering is observed for the increasingly excited electronic states, it seems probable that a significant fraction of the excited atoms survive the collision, and thus the conclusion is that only low impact parameter collisions lead to reaction.

Two features of the reaction imply that it proceeds through long range electron transfer and early H atom departure, even though only in a limited set of orientations and impact parameters. These are the very low product recoil energy, and the polarization dependences for the reaction of Na(4D) versus the lack of a polarization dependence for the reaction of Na(3P). That the reaction occurs through early departure of the H atom, but only for small impact parameter collisions, indicates that a coupling of nuclear and electronic motions is required for reaction. Otherwise, any trajectory for which electron transfer is energetically possible (i.e. any b <  $r_c$ ) would lead to reaction.

The dynamics of the Na(4D,5S) + HCl systems (the early H atom departure driven by the dissociative HCl<sup>-</sup>) appear to qualify these reactions for the type of calculation carried out by Kuntz, Polanyi, and coworkers, and termed DIPR for "Direct Interaction with Product Repulsion."  $^{50,51}$  As pointed out by Harris and Herschbach, the problem with the DIPR model is that it assumes a constant product repulsion, when in fact a distribution of product repulsions would be

more appropriate. 52,53 The approach of Harris and Herschbach was to assume a distribution of products as in photodissociation (in the DIPR-DIP model), since the negative ionic and electronically excited states of many closed shell systems have very similar dissociative potential energy surfaces.<sup>54</sup> For HCl<sup>-</sup>, an approximate potential has been derived for the dissociative ground state by Bardsley and Wadehra from their dissociative electron attachment experiments.<sup>47</sup> Since a potential is available for HC1<sup>-</sup>, there is no reason to use an excited neutral HCl potential curve, and a new variation on the DIPR theme would be more appropriate, perhaps to be called "Direct Interaction with Product Repulsion -- Distributed as In Electron-induced Dissociation," or "DIPR-DIED." The procedure used in a DIPR-DIED calculation would be exactly the same as for a DIPR-DIP calculation in which a Franck-Condon transition from the ground state potential to the repulsive ionic potential curve using the ground state vibrational wave function is assumed. The only real difference between DIPR-DIED and its previous incarnation (DIPR-DIP) is that DIPR-DIED will be using what is presumably a more accurate potential energy curve. The same limitations apply to the two models, namely that if the incoming alkali atom distorts the molecular reactant bond length, then the vibrational wave function is also distorted, and the product energies calculated will be incorrect. It would be most interesting to examine the geometric dependence of the reaction probability in this way.

Three other processes are energetically possible under the conditions used in the collisions of electronically excited Na atoms with HCl. It is possible to have the reaction proceed to form NaH:

Na + HCl  $\Rightarrow$  NaH + Cl  $\Delta H_0^0 = +58.8$  kcal/mole. (22) It is also possible to collisionally dissociate HCl:

Na + HCl  $\Rightarrow$  Na + Cl + H  $\Delta H_0^0 = \pm 102.2 \text{ kcal/mole.}$  (23) The third set of energetically possible processes involves quenching to other electronic states of Na:

$$Na(nL) + HC1 \rightarrow Na(n'L') + HC1(v',J').$$
(24)

The electronic energies available for reaction are 98.8 kcal/mole and 94.9 kcal/mole for the 4D and 5S states of Na, respectively. Both these energies are more than sufficient for the production of NaH (process (22)). However, no evidence for the production of NaH is observed. This is not surprising; if electron transfer occurs at relatively long range, the H atom is quickly ejected, before encountering the alkali atom. Another interpretation has the alkali atom acting as an alkali ion when close to the hydrogen halide. This is helpful in explaining the alkali favoring of the halide end of the HX molecule on approach, as seen in the strong variation of the barrier to (ground state) reaction as a function of M-X-H angle calculated by Shapiro and Zeiri; the barrier rises several eV as the orientation approaches the M-H-X collinear configuration.<sup>17</sup> The long range dispersive forces alone would have the alkali atom favoring the hydrogen end of the molecule.<sup>11,55</sup>

With collision energies of at least 3.5 kcal/mole and 7.0 kcal/mole, respectively, the 4D and 5S states of Na have enough energy to dissociate HCl (process (23)). This could only occur for the smallest impact parameter collisions, since if a significant amount of energy goes into product NaCl rotation, there will not be enough left to overcome the endothermicity of the dissociation. This is necessary in any case, since the maximum possible impulse would be required to induce dissociation. Another way of looking at collisional dissociation is as reaction to form NaCl, except that NaCl is formed above the dissociation limit. Although m/e=23 (Na<sup>+</sup>) product is detected. this is not due to collisional dissociation of HCl. Very similar product distributions are seen for collision energies above and below the minimum at which HCl dissociation can occur. The threshold for dissociation to  $Na^+ + Cl^- + H$  is at 137.4 kcal/mole, so this process, analogous to that seen by Lacmann and Herschbach, 3,15 is not yet energetically allowed for the conditions studied.

Quenching (process (24)) certainly occurs, but these measurements are particularly insensitive to it. Therefore, quenching processes will not be discussed. Hertel and coworkers, among others, have performed many experiments on the quenching of excited sodium atoms by various diatomic molecules (although no hydrogen halides have been reported to date), and the quenching cross sections are typically quite large,  $\sim 0(10-100 \text{ A}^2)$ .<sup>56</sup>

## Na(3S, 3P) + HBr

The Na + HBr reaction is nearly thermoneutral and there is a substantial cross section for the ground state reaction. However, the reactive cross section still increases upon excitation of the Na atoms to the 3P level. As with Na + HCl, the ground state cross section increases more with increasing collision energy than does the cross section of the Na(3P) state. Preliminary analysis of this data has shown it to be very similar in nature to Na(3P) + HCl, where product NaX is mostly backwards scattered, but there is some forward sideways peaking, and the product recoil energy distributions are broad and flat, and extend out to the total energy available.

2. The Lack of Reaction for Na + HF.

No reaction was observed for Na(3S,3P,4D,5S) + HF (3). Apparently, electronic energy was unable to make up for the substantial endothermicity and exit channel barrier (see reference 17) of the reaction of Na + HF. For comparison, recall that it took two vibrational quanta of HF excitation for Blackwell et al. to observe evidence of reaction to form NaF. There are no close resonances for electronic to vibrational energy transfer of HF that might preclude reaction in the manner predicted by Zeiri et al. for translational to vibrational energy transfer in Li + HF.<sup>19</sup> The vertical electron affinity of HF is approximately -4 eV and like HCl<sup>-</sup>, HF<sup>-</sup> is dissociative.<sup>26</sup> This puts

the crossing point of the neutral and ionic curves, and thus electron transfer at under 3 Å even for the 4D and 5S states.

In Li + HF, reaction was observed,<sup>2</sup> but the energetics were quite a bit different. From table I, the Li + HF reaction is slightly exothermic, whereas the Na + HF reaction is endothermic by 13 kcal/mole. Also, the barrier to reaction calculated by Zeiri and Shapiro in the ground state surface is 29 kcal/mole (vs. 12 kcal/mole for Li + HF) not including the effect of the zero point energy which will reduce it by 5 kcal/mole. This says nothing, however, about the electronically excited surfaces on which the reaction was attempted.

## E. Conclusions

The electronic excitation of the reagent alkali atom can indeed promote reaction. As seen in Na(4D,5S) + HCl, a new mechanism comes to dominate the reaction at high electronic energy, although the reactive cross sections do not reach the limit in which all collisions for which electron transfer is possible lead to reaction. This is due to the restrictive geometric constraints necessary for reaction. Low impact parameter collisions with Na approaching the Cl end of the HCl molecule are almost exclusively responsible for production of NaCl, since a coupling of the electronic and nuclear motion is necessary for reaction.

To some extent electronic excitation is not directly comparable to other forms of energy since the potential energy surface on which the

reaction takes place is different for each electronic state -- with different topography and different possible symmetries. Nevertheless it is certainly interesting to ask how electronic excitation compares to other forms of energy in promoting reaction. Like electronic energy, exciting the first (or second) vibration is capable of turning on the reaction when the collision energy is insufficient to overcome the endothermicity. This was shown for K + HF, HC1<sup>6</sup> (and for Na + HF(v=2)<sup>10</sup>). As discussed in section B and reference 6, with increasing translational energy the first vibration in K + HCl no longer enhances the reaction cross section, whereas there remains over an order of magnitude enhancement in K + HF. Rotational excitation has only been studied in vibrationally excited HCl, HF + Na, K, and seems to adversely affect the reaction cross section or have little effect. $^{8-10}$  Translational energy overcomes the threshold to reaction in K + HCl, HF and continues to enhance reaction until a saturated value of the reaction cross section is reached.<sup>6</sup> As shown in the nonreaction of  $Na^* + HF$ , electronic excitation is not necessarily sufficient for reaction. However, in Na + HCl, HBr the initial electronic excitation to Na(3P) enhances the reactivity strongly whether or not there is sufficient collision energy for the ground state reaction. In Na + HCl, there is a small further increase (less than a factor of two from  $3P \rightarrow 4D$ , 5S) with further increased electronic energy.
# F. References

1.	E. H. Taylor and S. Datz, <u>J. Chem. Phys. 23</u> , 1711 (1955).
2.	C. H. Becker, P. Casavecchia, P. W. Tiedemann, J. J. Valentini,
	and Y. T. Lee, <u>J. Chem. Phys. 73</u> , 2833 (1980).
3.	K. Lacmann and D. R. Herschbach, <u>Chem. Phys. Lett. 6</u> , 106 (1970).
4.	J. G. Pruett, F. R. Grabiner, and P. R. Brooks, <u>J. Chem. Phys. 63</u> ,
	1173 (1975).
5.	M. W. Geis, H. Dispert, T. L. Budzynski, and P. R. Brooks, in
	State to State Chemistry, P. R. Brooks and E. F. Hayes, eds., ACS
	Symposium Series No. 56, 103, (American Chemical Society,
	Washington, DC, 1977).
6.	F. Heismann and H. J. Loesch, <u>Chemical Physics 64</u> , 43 (1982).
7.	T. J. Odiorne, P. R. Brooks., and J. V. V. Kasper, <u>J. Chem. Phys.</u>
	<u>55</u> , 1980 (1971).
8.	M. Hoffmeister, L. Potthast, and H. J. Loesch, <u>Chemical Physics</u>
	<u>78</u> , 369 (1983).
9.	H. H. Dispert, M. W. Geis, and P. R. Brooks, <u>J. Chem. Phys. 70</u> ,
	5317 (1979)
10.	B. A. Blackwell, J. C. Polanyi, and J. J. Sloan, Chemical Physics
	<u>30</u> , 299 (1978).
11.	M. F. Vernon, H. Schmidt, P. S. Weiss, M. H. Covinsky, and Y. T.
	Lee, <u>J. Chem. Phys.</u> , to be published.

- 12. R. R. Herm, in <u>Alkali Halide Vapors: Structure, Spectra, and</u> <u>Reaction Dynamics</u>, 189 (Academic Press, New York, 1979), and references therein.
- J. L. Kinsey, <u>MTP Int. Rev. Sci., Phys. Chem. Ser. One 9</u>, 173 (1972).
- 14. W. A. Lester, Jr. and M. Krauss, J. Chem. Phys. 52, 4775 (1970).
- 15. G. G. Balint-Kurti and R. N. Yardley, <u>Faraday Disc. Chem. Soc. 62</u>, 77 (1977).
- 16. Y. Zeiri and M. Shapiro, <u>Chem. Phys.</u> 31, 217 (1978).
- 17. M. Shapiro and Y. Zeiri, J. Chem. Phys. 70, 5264 (1979).
- M. M. L. Chen and H. F. Schaefer III, <u>J. Chem. Phys. 72</u>, 4376 (1980).
- 19. Y. Zeiri, M. Shapiro, and E. Pollack, Chem. Phys. 60, 239 (1981).
- 20. I. NoorBatcha and N. Sathyamurthy, J. Chem. Phys. 76, 6447 (1982).
- 21. I. NoorBatcha and N. Sathyamurthy, <u>J. Am. Chem. Soc. 104</u>, 1766 (1982).
- 22. I. NoorBatcha and N. Sathyamurthy, <u>Chem. Phys. Lett. 93</u>, 432 (1982).
- 23. K. P. Huber and G. Herzberg, <u>Molecular Spectra and Molecular</u> <u>Structure, IV. Constants of Diatomic Molecules</u> (Van Nostrand Reinhold, Co., New York, 1979).
- 24. D. R. Herschbach, <u>Proceedings of the Conference on Potential</u> <u>Energy Surfaces in Chemistry</u>, 44, W. A. Lester, Jr., ed., IBM Research Laboratory, San Jose, California (1971).
- 25. S. Carter and J.N. Murell, Mol. Phys. 41, 567 (1980),

- 26. D. C. Frost and C. A. McDowell, <u>J. Chem. Phys. 29</u>, 503 (1958).
- 27. J. H. Birely, R. R. Herm, K. R. Wilson, and D. R. Herschbach, <u>J.</u> <u>Chem. Phys. 47</u>, 993 (1967).
- 28. E. A. Mason, J. T. Vanderslice, and C. J. G. Raw, <u>J. Chem. Phys.</u> <u>40</u>, 2153 (1964).
- 29. J. C. Slater and J. G. Kirkwood, Phys. Rev. 37, 682 (1931).
- 30. K. S. Pitzer, Adv. Chem. Phys. 2, 59 (1959).
- 31. H.-J. Werner and W. Meyer, <u>Phys. Rev. A 13</u>, 13 (1976).
- 32. A. A. Maryott and F. Buckley, National Bureau of Standards (U. S.) Circ. 537 (1953).
- 33. see for example the comparison of  $C_{6}$  values in ref. 30.
- 34. G. York, R. Scheps, and A. Gallagher, <u>J. Chem. Phys. 63</u>, 1052 (1975).
- 35. J. S. Cohen and R. T. Pack, J. Chem. Phys. 61, 2372 (1974).
- 36. J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, <u>Molecular</u> <u>Theory of Gases and Liquids</u> (John Wiley and Sons, Inc., New York, 1964).
- 37. R. D. Nelson, Jr., D. R. Lide, Jr., and A. A. Maryott, <u>National</u> <u>Standard Reference Data Series 10</u>, National Bureau of Standards (U. S.) (1967).
- 38. For further explanation see R. B. Bernstein, <u>Chemical Dynamics via</u> <u>Molecular Beam and Laser Techniques</u>, pp. 86-7 (Oxford University Press, Oxford, 1982).
- P. Hannaford, W. R. MacGillivray, and M. C. Standage, <u>J. Phys. B:</u> Atom Molec. Phys. 12, 4033 (1979).

- 40. H. T. Duong and J.-L. Picque, <u>J. Physique 33</u>, 513 (1972).
- 41. R. H. McFarland and J. D. Kinney, Phys. Rev. 137, A1058 (1965).
- 42. P. Mohazzabi and A. W. Searcy, <u>Intl. J. Mass Spect. Ion Phys. 24</u>, 469 (1977).
- 43. P. Brumer and M. Karplus, J. Chem Phys. 58, 3903 (1973).
- 44. S. C. Brown, <u>Basic Data of Plasma Physics</u>, 2<sup>nd</sup> Ed. (MIT Press, Cambridge, Massachusetts, 1967).
- 45. The fragmentation ratio of 0.9 is used for the NaCl produced from the reaction of the electronically excited Na atoms.
- 46. C. H. Chen, H. Haberland, and Y. T. Lee, <u>J. Chem. Phys. 61</u>, 3095 (1974).
- 47. J. N. Bardsley and J. M. Wadehra, J. Chem. Phys. 78, 7227 (1983).
- 48. J. L. Magee, <u>J. Chem. Phys. 8</u>, 687 (1940), or R. D. Levine and R.
  B. Bernstein, <u>Molecular Reaction Dynamics</u>, 86 (Oxford University Press, New York, 1974).
- 49. W. L. Wiese, M. W. Smith, and B. M. Miles, <u>Atomic Transition</u> <u>Probabilities Volume II: Sodium Through Calcium</u>, Natl. Stand. Ref. Data Ser., Nat. Bur. Stand. (U. S.), Washington, DC (1969).
- 50. P. J. Kuntz, E. M. Nemeth, and J. C. Polanyi, <u>J. Chem. Phys. 50</u>, 4607 (1969).
- 51. P. J. Kuntz, M. H. Mok, and J. C. Polanyi, ibid., 4623 (1969).
- 52. R. M. Harris, Ph. D. Thesis, Harvard University, Cambridge, Massachusetts (1970).
- 53. D. R. Herschbach, Faraday Disc. Chem. Soc. 55, 233 (1973).
- 54. This is discussed in chapter IV, or see reference 31.

- 55. H. Schmidt, P. S. Weiss, J. M. Mestdagh, M. H. Covinsky, M. F. Vernon, and Y. T. Lee, in preparation.
- 56. I. V. Hertel and W. Stoll, <u>Adv. At. Mol. Phys. 13</u>, 113 (1977), and references therein.

# III. THE STATE-SELECTIVE REACTION OF ELECTRONICALLY EXCITED SODIUM ATOMS WITH MOLECULAR OXYGEN

### A. Introduction

The Na\* +  $0_2$  reaction has been chosen for study since the energetics preclude reaction in the ground state. The large endothermicity of the reaction:

 $Na(3^{2}S) + O_{2}(X^{3}\Sigma_{g}) \Rightarrow NaO(X^{2}\pi) + O(^{3}P) \Delta H_{O}^{0} = +58 \text{ kcal/mole}^{1}$  (1) is lower than the electronic energy of Na atoms in the 4S, 4P, 4D, and 5S states.<sup>3</sup> With at least 10 kcal/mole of translational energy Na atoms in the 3P state also have more than enough energy to overcome this endothermicity. Another point of interest is that the positive electron affinity of  $O_{2}$ , EA=0.44 eV,<sup>2</sup> undoubtedly leads to an ionic intermediate, Na<sup>+</sup>O<sub>2</sub>, which must play an important role in the collision between Na\* and  $O_{2}$ .

The NaO<sub>2</sub> molecule has been studied both theoretically and spectroscopically.<sup>4-7</sup> The theoretically determined lowest energy state of the NaO<sub>2</sub> system, which has  $C_{2v}$  symmetry, has the Na atom at the apex of an isoceles triangle, and is 1.6 eV more stable than ground state reactants (Na(3<sup>2</sup>S) + O<sub>2</sub>(X<sup>3</sup> $\Sigma_{g}^{-}$ )).<sup>4</sup> Infrared, Raman, and ESR spectroscopies of NaO<sub>2</sub> trapped in low temperature rare gas matrices have also shown the triatomic to be an isoceles triangle.<sup>5,6,7</sup>

Calculations on the NaO molecule indicate that it is ionic, Na<sup>+</sup>O<sup>-</sup>.<sup>8</sup> The ground state is the NaO(X<sup>2</sup> $\pi$ ) state, and there is a very low-lying first excited state at ~1500 cm<sup>-1</sup>, the NaO(A<sup>2</sup> $\Sigma$ <sup>+</sup>) state.<sup>8</sup> Certainly many other states exist for this open shell molecule, but nearly nothing is known of them. The first step to amend this was made by Pfeiffer and Gole, who reported seeing chemiluminescence arising from NaO from 10000-15000 cm<sup>-1</sup>, but left the observed spectra unassigned.<sup>9</sup> An attempt is being made to measure and assign the electronic spectra of NaO and other alkali monoxides by the Gole and Schulz groups at Georgia Tech.<sup>10</sup>

A compilation of symmetries of important reactant and intermediate states for Na +  $0_2$  is given in table I for reference in the following discussion.

Kempter and coworkers have studied the excitation of the Na(3P) level in 1-17 eV collisions with  $0_2$  by monitoring the Na D-line fluorescence.<sup>11</sup> In the process:

$$Na(3S) + 0_2 \Rightarrow Na^*(3P) + 0_2$$
 (2)

they find the threshold at 2.1 eV, the Na(3P) excitation energy. At a collision energy of 5 eV, they find the cross section for this process to be 26  $\mathbb{A}^2$ , and see a further increase of approximately an order of magnitude at a collision energy of 8 eV. At collision energies of 8 eV and above, 20% of the fluorescence observed appears to be due to states other than Na(3P), implying higher excitation takes place. The ground and excited states are strongly coupled via the ion-pair state (see table I). Because the difference between the Na ionization potential

System	Symmetry		<u>N</u>	<u>Multiplicity</u>
	C <sub>∞v</sub>	C <sub>s</sub>	C <sub>2v</sub>	
$Na(ns^{2}S)+0_{2}(X^{3}\Sigma_{g}^{-})$	.Σ <sup></sup>	Α"	<sup>8</sup> 1	2,4
$Na(np^2P)+O_2(X^3\Sigma_q)$	Σ	Α"	A <sub>2</sub>	2,4
5	+ 11	A۱	A <sub>1</sub>	2,4
	π	Α"	B <sub>1</sub>	2,4
Na(nd <sup>2</sup> D)+0 <sub>2</sub> (X <sup>3</sup> Σ <sup>-</sup> g)	Σ	Α"	<sup>B</sup> 1	2,4
	+ π	Α'	A <sub>1</sub>	2,4
	π	Α "	A2	2,4
	Δ+	Α'	<sup>8</sup> 2	2,4
	Δ.	Α"	<sup>B</sup> 1	2,4
NaO <sub>2</sub> (lowest 2 states)	π	A۳	A <sub>2</sub>	. 2
	+ 1î	Α'	<sup>B</sup> 2	2
$Na^{+}(^{1}S)+O_{2}^{-}(X^{2}\pi_{g})$	+ π	A۱	<sup>B</sup> 2	2
	π	Α"	A <sub>2</sub>	2
$Na^{+}(^{1}S)+O_{2}^{-}(a^{4}\Sigma_{u}^{-})$	Σ	A"	A <sub>2</sub>	4
$Na^{+}(^{1}S)+O_{2}^{-}(^{2}\Delta)$	Δ	A'	A <sub>1</sub>	2
- 	۵-	Α".	B <sub>1</sub>	2

Table I. Symmetries of important reactant and intermediate states for Na +  $0_2$  collisions.

and the electron affinity is larger than the Na(3P) excitation energy, excitation is expected to be more important than collisional ionization.<sup>12</sup>

A merging beams study of Na +  $0_2$  was carried out by Neynaber, Myers and Trujillo at collision energies ranging from 4.7 to 25 eV.<sup>13</sup> Only processes producing ions were measured, these were:

- $Na + 0_{2} \Rightarrow Na^{+} + 0_{2}^{-} \qquad \Delta H_{0}^{0} = \pm 108.3 \text{ kcal/mole}^{1} \qquad (3)$   $Na + 0_{2} \Rightarrow Na^{+} + 0_{2} + e^{-} \qquad \Delta H_{0}^{0} = \pm 118.5 \text{ kcal/mole}^{1} \qquad (4)$ 
  - $Na + 0_2 \rightarrow Na^+ + 0 + 0^- \Delta H_0^0 = +202.6 \text{ kcal/mole}^1$  (5)
  - Na +  $0_2 \gg Na0^+ + 0^- \Delta H_0^0 = +184.2 \text{ kcal/mole}^1$  (6)
- $Na + 0_2 \Rightarrow Na0^+ + 0 + e^-. \Delta H_0^0 = +218.0 \text{ kcal/mole}^1$  (7)

Ions were observed at collision energies of 6 eV and above, but not at 4.7 eV (the endothermicity of process (3)). The maximum cross section for the production of Na<sup>+</sup> (the sum of the cross sections of processes (3-5)) was measured to be  $\sigma_{Na}=0.05 \text{ Å}^2$ . The maximum cross section for the production of Na0<sup>+</sup> (the sum of the cross sections of processes (6 and 7)) was measured to be  $\sigma_{Na0}=0.004 \text{ Å}^2$ . The positive ions produced in the charge transfer processes are observed in the forward center-of-mass direction, as expected. There was no attempt to observe neutral products such as NaO or Na<sup>\*</sup>.

Kleyn, Los, and coworkers have made measurements of charge exchange processes (3, 6, and 7) at collision energies of 4-2000 eV.<sup>14,15</sup> Cross sections for production of  $0_2^-$  (process (3)) appear to be as high as 5 Å<sup>2</sup>, while cross sections for production of  $0^-$  (processes (6 and 7)) are always less than 0.3 Å<sup>2</sup>. The

discrepancy between the cross sections of Neynaber et al. $^{13}$  and those of Kleyn et al.<sup>13</sup> remains a subject of controversy with conflicting measurements also reported for analogous  $M + O_2$  (M = alkali atom) systems by others.<sup>14,16</sup> In that the calibration of the cross sections in the merging beams apparatus is more difficult, the values of Kleyn et al. are perhaps more dependable. These values are still significantly lower than the excitation cross sections of Kempter et al.,<sup>11</sup> bearing out Kempter's predictions as to the relative importance of collisional excitation and ionization.<sup>12</sup> Kleyn et al. have shown that at the high collision energies studied the change of the phase of the  $0_2^{(-)}$  vibration during the collision is very important in determining whether ions are produced. This has to do with the relative positions of the neutral and ionic curves after the collision, due to vibrational excitation from the collision; this effect is termed "vibronic excitation." This can only be observed in their studies because the high collision energies lead to very short collision times, ~O(10<sup>-14</sup> sec), which are on the same time scale as the  $O_2$  vibrational period. As such the crossing radius of the neutral and ionic curves,  $r_c$ , depends on the phase of the  $0_2^{(-)}$  vibration, and formula (18) of chapter II must be rewritten:

$$r_{c} = \frac{e^{2}}{IP - EA(r_{0_{2}})}, \qquad (8)$$

where IP is the ionization potential of the alkali atom, and  $EA(r_{0_2})$  is the electron affinity of  $0_2$ , which varies as a function of the  $0_2$ 

interatomic distance. Some of the known  $0_2$  and  $0_2^-$  potential energy curves are shown in figure 1, which is reproduced from reference 17 with permission. Note that the vertical electron affinity of  $0_2$  is approximately zero.

The collisions studied here are at much lower translational energies than those of Kleyn, and the direct effects of vibronic excitation cannot be observed. However, since the collision times are much longer here, it is important to remember that many  $0_2$  vibrations will be taking place, effectively making the energy separation of the neutral and ionic curves oscillate during the collision. Additionally, if the approaching Na atoms allow the  $0_2$  molecule to stretch, then the electron affinity can take on its adiabatic value of 0.44 eV.

Many measurements of charge exchange processes for other alkali atoms with  $O_2$  (analogues to processes (3-7)) have been made.<sup>14,16</sup> The excitation process is stronger for Na than for heavier alkali atoms<sup>11</sup> due in part to its higher ionization potential.

A crude study of the translational energy dependence of all ionproducing processes for  $Na(3S, 3P) + 0_2$  has been performed by Roethe, Mather, and Beck.<sup>18</sup> They have found that there are similar cross sections for ion-producing processes for each state with a lower threshold energy for Na(3P). No values of the threshold, nor of the cross sections were reported. This measurement shows, however, that ionization does not gain appreciable probability when the ionization potential of the alkali is effectively lowered by excitation.



Fig. 1. Potential energy curves for  $0_2$  and  $0_2^-$ , reproduced from reference 17 with permission.

In the experiments described here, electronic excitation is used in an attempt to overcome the very substantial endothermicity of the neutral reaction (1).

#### B. Results

The apparatus used to make the measurements reported below is described in chapter I. Laboratory angular distributions of products were measured for each of the three optically pumped Na levels (3P, 4D, 5S) at three center-of-mass collision energies: 8, 16, and 18 kcal/mole. The Newton velocity vector diagrams corresponding to these are shown in figures 2, 3, and 4, while the angular distributions themselves are shown in figures 5, 6, and 7, respectively. Also, angular distributions were recorded for Na(3S,4D) +  $0_2$  at mean collision energies of 11, 12, 13, and 22 kcal/mole. The Newton diagrams for these energies are shown in figure 8, for 11-13 kcal/mole, in which the Na velocity was varied by using mixtures of He and Ne as backing gases, and in figure 9 for 22 kcal/mole. The corresponding angular distributions are shown in figures 10a-c, and 11, respectively. At the lowest collision energy (8 kcal/mole) nothing other than elastic scattering was observed for any of the Na levels as shown in the angular distributions in figure 5. At all other collision energies, a peak was observed only when the sodium atoms were optically pumped to the Na(4D) state. Nothing other than elastic scattering was

Fig. 2. Newton velocity vector diagram for Na(3P,4D,5S) +  $0_2 \Rightarrow \underline{NaO}$ + 0 at a collision energy of 8 kcal/mole. The sodium beam is seeded in neon, and the molecular beam is 5%  $0_2$  in helium.



Fig. 2

XBL 863-856

Fig. 3. Newton velocity vector diagram for Na(3P,4D,5S) +  $0_2 \rightarrow Na0 + 0$ at a collision energy of 16 kcal/mole. The sodium beam is seeded in helium, and the molecular beam is neat  $0_2$ .



E<sub>coll</sub>=16 kcal/mole

Fig. 3

XBL 863-853

Newton velocity vector diagram for Na(3P,4D,5S) +  $0_2 \Rightarrow \underline{Na0} + 0$ Fig. 4. at a collision energy of 18 kcal/mole. The sodium beam is seeded in helium, and the molecular beam is 5%  $\boldsymbol{0}_2$  in helium.

E<sub>coll</sub>=18 kcal/mole



Fig. 4

XBL 863-855

Fig. 5. NaO product angular distribution for Na(3S,4D) + O<sub>2</sub> at a collision energy of 8 kcal/mole. The low angle rise is due to elastically and inelastically scattered Na atoms. Note that no reactive scattering is observed for either state.





Fig. 6. NaO product angular distribution for Na(3P,4D) + 0<sub>2</sub> at a collision energy of 16 kcal/mole. The low angle rise is due to elastically and inelastically scattered Na atoms.



Fig. 7. NaO product angular distribution for Na(4D,5S) + O<sub>2</sub> at a collision energy of 18 kcal/mole. The low angle rise is due to elastically and inelastically scattered Na atoms.



Fig. 8. Newton velocity vector diagram for Na +  $0_2 \Rightarrow \underline{Na0}$  + 0 at collision energies of 11, 12, and 13 kcal/mole. The sodium beam is seeded in a mixture of helium and neon, and the molecular beam is 5%  $0_2$  in helium. The helium-neon mixture is varied to produce the different Na beam velocities shown.

155

13 E<sub>coll</sub>= 12 kcal/mole 11

⊣ 5.0x10<sup>4</sup> cm/sec

E<sub>coll</sub>=22 kcal/mole



----- 5.0x10<sup>4</sup> cm/sec

XBL 863-857

Fig. 9. Newton velocity vector diagram for Na +  $0_2 \Rightarrow Na0$  + 0 at a collision energy of 22 kcal/mole. The sodium beam is seeded in helium, and the molecular beam is 5%  $0_2$  in helium.









observed for  $Na(5S) + O_2$ . In addition to the distributions shown, an effort was made to see scattering from Na(5S) at collision energies of 8 and 16 kcal/mole, although only the 18 kcal/mole data is shown. Since the 5S level radiatively populates the 4S and 4P levels, there is no reaction for these levels as well. It is clear that that the signal is indeed due only to the 4D state. All these distributions were measured with the mass spectrometer tuned to a mass to charge ratio (m/q) of 23  $(Na^{+})$  because the product NaO would fragment to  $Na^{+}$  in the electron bombardment ionizer. Angular distributions were also taken for the 18 kcal/mole collision energy with the quadrupole mass spectrometer set to  $m/q=16 (0^+ \text{ and } 0^{++}_2)$  and  $m/q=39 (Na0^+)$ , and are shown in figures 12 and 13, respectively, however the large background at each of these mass settings frustrated all attempts to observe a true mass-dependent signal. Product velocities were measured by the cross-correlation time-of-flight method at a number of angles for each energy, 16 and 18 kcal/mole, for m/q=23 (Na<sup>+</sup>), and these are shown in figures 14 and 15, respectively. Time-of-flight distributions were also measured at a collision energy of 18 kcal/mole for  $m/q=16 (0^+)$ and  $0_2^{++}$ ) and m/q=39 (NaO<sup>+</sup>), shown in figures 16 and 17, respectively.

It was discovered while making these measurements that signal could be observed with the detector ionizer and the majority of the ion optics turned off. In fact, only the -30 kV ion target and photomultiplier need be powered in order to see some signal. Reactive product angular distributions were recorded with the electron







Fig. 13. Product angular distribution for  $Na(4D) + O_2$  at a collision energy of 18 kcal/mole with the detector set to m/q=39(Na0<sup>+</sup>).



Fig. 14. NaO product time-of-flight measurements at various detector angles (shown in each frame) at a collision energy of 16 kcal/mole for m/q=23 (Na<sup>+</sup>).







Fig. 16. Product time-of-flight measurements at various detector angles at a collision energy of 18 kcal/mole with the detector set to m/q=16 (0<sup>+</sup>, 0<sup>+</sup><sub>2</sub>).

Fig. 17. Product time-of-flight measurement for 40° at a collision energy of 18 kcal/mole with the detector set to m/q=39(Na0<sup>+</sup>).



bombardment ionizer filament emission off, but all other conditions the This subsequent ionization is what is responsible for the tails same. on the product velocity distributions in figures 14-17. An example of this effect is shown in the two time-of-flight distributions shown in figure 18. The only difference between the two traces is that for the dotted line the ionizer filament is emitting electrons, and for the solid line, no electrons are emitted. This signal is due to ionization at a number of points in the detector. This has been shown by turning off the detector electronics sequentially and seeing time-of-flight peaks disappear one by one. Of course when this is done the peaks change position since the ion flight time is changed as the various accelerating optics are turned off. Time-of-flight measurements with the ionizer emission off, thus giving no electron bombardment, are shown for the quadrupole spectrometer set to m/q=16 (0<sup>+</sup> and 0<sup>+</sup><sub>2</sub>), m/g=23 (Na<sup>+</sup>), and m/g=39 (NaO<sup>+</sup>) for a collision energy of 18 kcal/mole and for a number of detector angles in figures 19, 20, and 21, respectively.

The ratio of the intensity of the first TOF peak, due to electron bombardment ionization, to that of the final peak for which only the doorknob and the photomultiplier need be powered, is 10:1 for the 18 kcal/mole data. This ratio remains constant accross the various lab angles measured. This is also seen in that the angular distributions have the same shape with the detector on and off.

Where signal is small, better statistics are found in the emission off distributions. These are shown for collision energies 11, 12, and
Fig. 18. Time-of-flight distributions for 40° at a collision energy of 18 kcal/mole with the ionizer emission on (dotted line) and off (solid line).

169 XBL 863-869 s: ; ...... 40°

Fig. 18



XBL 863-893

Fig. 19. Time-of-flight distributions for the angles shown with the mass spectrometer set to m/q=16 at a collision energy of 18 kcal/mole with the ionizer emission off.



XBL 863-891

Fig. 20. Time-of-flight distributions for the angles shown with the mass spectrometer set to m/q=23 at a collision energy of 18 kcal/mole with the ionizer emission off.

Fig. 21. Time-of-flight distribution for 40° at a collision energy of 18 kcal/mole with the mass spectrometer set to m/q=39 and with the ionizer emission off.



13 kcal/mole in figures 22a, b, and c, respectively. These low collision energies were studied with a fairly broad Na beam velocity distribution (S =  $v/\Delta v<4$ ) and the very narrow  $0_2$  beam velocity distribution described in chapter I (S>10). The peak intensity increases strongly for these distributions, but does not change position in the laboratory. Also, a low angle peak seems to come into the distribution with increasing collision energy.

In an effort to understand why the observed process is stateselective, polarization dependences of the differential cross sections were measured in order to try to determine the symmetry of the transition state. These measurements were made as explained in chapter I by rotating the laser polarization in the scattering plane to align the Na excited orbitals, while the signals at fixed detector angles are recorded. These dependences were measured for 5 angles at a collision energy of 18 kcal/mole with the mass spectrometer set to m/q=23 (Na<sup>+</sup>). The important quantities deduced are the favored polarization angle for reaction at a given detector angle, and the amplitude of this polarization effect. The measured dependences are shown in figures 23a-e, and the best fits to the curves (as determined by the method described in chapter I) are sùmmarized in table II.



Fig. 22. Product angular distributions recorded with the ionizer emission off for Na(4D) + O<sub>2</sub> at collision energies of a) 11 kcal/mole, b) 12 kcal/mole, and c) 13 kcal/mole.

XBL 863-843

Fig. 23. Laser polarization dependences of the signal due to the Na(4D) +  $O_2$  reaction at a collision energy of 18 kcal/mole for the detector angles shown in each frame. The solid lines are the best fits to the data as discussed in chapter I.



XBL 863-892

Fig. 23

Laboratory Detector Angle	Laboratory Peak of Polarization Dependence (øLAB)	Center-of-Mass Peak of Polarization Dependence (ØCM)	Amplitude (2A)
30 <b>°</b>	-18°	9°	.130
36 <b>°</b>	-64 °	-37°	.080
40°	-88°	-61°	.106
50°	83°	-70°	.074
60°	29° .	56°	.084

Table II. The measured polarization dependences for Na(4D) + O<sub>2</sub> at a collision energy of 18 kcal/mole.

C. Analysis of Experimental Results

First, it is clear from the data that only the Na(4D) state undergoes whatever process produces the high laboratory angle peak. Since the Na(5S) state radiatively populates the Na(4S,4P) states, the Na(3S,3P,4S,4P,5S) states do not react at a collision energy of 18 kcal/mole.

In transforming the Na(4D) +  $0_2$  data to center-of-mass frame distributions, it is assumed that the reactive process (1), and thus product NaO was responsible for the data at m/q=23 (Na<sup>+</sup>). In fact, since the data is fitted to a center-of-mass angular distribution and a center-of-mass velocity distribution P(u), it does not matter what the product is in fitting the data. The fitted center-of-mass velocity distribution is reported as a function of translational energy, P(E), which is a more meaningful quantity for direct comparison with other experiments, which <u>is</u> dependent on the product masses, however. As shown below, there is a simple scaling factor allowing the interconversion of P(E) distributions due to different product masses, but the same center-of-mass velocity distributions. The possibility that other processes are responsible for the signal recorded is also discussed below.

The laboratory angular and velocity data at a collision energy of 18 kcal/mole were transformed to center-of-mass angular and velocity information using program GM as described in chapter I. The best fit

to the angular distribution is shown by the solid line in figure 24. and to the time-of-flight measurements in figure 25. This yielded the translational energy distribution  $(P(E_T))$  in figure 26, and the center-of-mass NaO product angular distribution  $(T(\theta))$  in figure 27. These remarkable distributions show the extremely limited conditions under which NaO is produced. As expected from examining the laboratory angular distributions, the center-of-mass scattering angular distributions are peaked sharply backwards with respect to the initial Na atom direction. The peak is clearly directly backwards (at  $\theta=180^{\circ}$ ), with 50% of the NaO product scattering to center-of-mass angles greater than 150° with respect to the incoming Na relative (center-of-mass frame) velocity (this is weighted NaO intensity, i.e.  $T(\theta)$  sin $\theta$ ). The very low translational energy of the products shows that the products must be very internally excited. The peak of the product recoil energy distribution is at 0.6 kcal/mole, and the mean recoil energy is 0.7 kcal/mole. From conservation of energy considerations, an average of 58 kcal/mole (2.5 eV) must be in the internal excitation of the products.

By weighting the reaction cross section according to collision energy using a step function, the reaction threshold was determined to be 15±1 kcal/mole. It is also evident from the constant peak position in the laboratory angular distributions for 11-13 kcal/mole mean collision energies that the threshold to reaction is dominating the shapes of these distributions. The assumption of a step function is clearly not valid, but is meant to be a zeroth order approximation useful in determining the threshold. Also, an implicit assumption is







XBL 863-793

Fig. 25. The best fit to the time-of-flight distributions for Na(4D) +  $O_2 \Rightarrow \underline{NaO} + O$  at a collision energy of 18 kcal/mole for the angles shown. The solid lines show the fit to the data.



Fig. 26. Center-of-mass product translational energy distribution for Na(4D) +  $0_2 \Rightarrow Na0$  + 0 at a collision energy of 18 kcal/mole.



Fig. 27. Center-of-mass scattering angular distribution for Na(4D) +  $0_2 \Rightarrow \underline{Na0} + 0$  at a collision energy of 18 kcal/mole.

that the center-of-mass product scattering distributions do not change as a function of collision energy. In the threshold region, these might be expected to change dramatically.

By determining the time at which the subsequent peaks occur in the time-of-flight distributions, it can be deduced at what point in the detector the species is ionized. In fact only the final peak, the one for which only the doorknob (ion target) and photomultiplier need be on, is readily interpretable. The others have unknown ion flight times, whereas the ions responsible for the final peak must be produced in the region of the doorknob, after which their ion flight time should be negligible. By calculating the peak velocities from the main peaks in each of the time-of-flight distributions for the 18 kcal/mole collision energy (figure 14), and from the positions of the final peaks, it is possible to calculate the flight length of the neutral. From the modulation point, the interaction region, to the detection point is a distance of 54±3 cm. This is derived from the peak separations of 19-21 channels of 10 usec duration, the peak of the detected product velocities of  $1.4-1.7 \times 10^5$  cm/sec, and the ion flight time of (electron bombardment ionization produced) Na<sup>+</sup>, ~14 µsec. This corresponds to the distance from the collision region to the front half of the region between the doorknob and the photomultiplier. Due to the high voltage of the doorknob (-30 kV), the electric field in this region is approximately 10 kV/cm. This would have a minimal effect on the NaO energy, Stark shifting the energy levels less than  $2 \text{ cm}^{-1}$ .

The back of the detector is approximately 60 cm. from the interaction region, and the ionization does not appear to be occurring there.

Since only signal at m/q=23 (Na<sup>+</sup>) has been shown to be due to electron bombardment ionization, it is necessary to address the question of whether some species other than NaO could be responsible for the signal observed. Since Na<sup>+</sup> is observed after ionization, the signal-producing species must contain Na. The non-electron bombardment ionization signal does not come from ground state Na atoms, since Na atoms at comparable velocities are produced in these and other experiments, and no similar signal is observed. Na atoms in low-lying excited levels have such short radiative lifetimes (tens of nsec, see chapter I) as to all radiatively decay before reaching the detector. Since the observed products do not occur at the center-of-mass, they are not a NaO<sub>2</sub> complex that lives long enough to make it into the detector. The placement of the products and the intentional destruction of dimers and clusters in the expansions negates the possibility that these could lead to the signal observed. Could metastable Na atoms be responsible for the signal? High Rydberg states of Na do have lifetimes sufficient to survive the transit through the detector to be field ionized. This possibility could be substantiated or eliminated by observing the complementary product 0 or  $0_2$ , however all attempts to do this were frustrated by the large detector background at these masses (m/q=16 and 32), and the fact that the  $0_2$  beam puts a constant load on the detector at these masses.

If the detected product is Na and not NaO, the fit of the center-of-mass angular and velocity distributions described above is still applicable, however the product translational energy distribution must be scaled by a factor related to the product reduced masses. The product translational energy in terms of the velocity of the detected product can be derived as:

$$E_{T} = \frac{1}{2} m_{det} u_{det}^{2} \qquad \frac{m_{det} + m_{undet}}{m_{undet}}, \qquad (9)$$

where  $m_{det}$  and  $m_{undet}$  are the detected and undetected products, respectively, and  $u_{det}$  is the center-of-mass velocity of the detected product. This is derived from conservation of momentum; the momenta of the two products are equal and opposite in the center-of-mass so that not only the masses change, the velocities of the undetected product do also. The ratio of the energies for a given  $u_{det}$  is then:

$$\frac{E(\det '=Na)}{E(\det =Na0)} = \frac{\overset{m'det mundet}{}}{\overset{m'det mundet}{}}, \qquad (10)$$

giving a conversion factor of 0.30 for the  $E_T$  distribution fitted for product NaO. Thus, instead of peaking at 0.6 kcal/mole for product NaO, assuming that the product is Na, the P(E) peaks at 0.2 kcal/mole. From conservation of energy, this gives a mean total product internal energy of 116.7 kcal/mole, or 40820 cm<sup>-1</sup>, for the 18.1 kcal/mole <u>mean</u> collision energy. Including the beam velocity spreads, enough energy exists to remove the Na electron. It should also be noted that all the collision energies for which product was observed are well above the energy required for electron transfer, process (3).

The signal due to ionization near the doorknob is one-tenth the signal due to electron bombardment ionization (see figures 15 and 18). By the time the atoms reach the detector ionizer, no matter in what state they are produced, the majority of the Na atoms will have radiatively decayed into the ground state. Their detection efficiency can be determined from the filament emission current  $\rm I_{p}{=}10~mA$ =6.2x10<sup>16</sup> e<sup>-</sup>/cm<sup>2</sup>-sec, and the ionization cross section of Na by 200 eV electrons,  $\sigma=2.46 \text{ Å}^2.19$  The ionization rate is then  $I\sigma = 13$  $sec^{-1}$ . The atomic velocity is very nearly the velocity of the centerof-mass,  $v_{cm}=1.5\times10^5$  cm/sec. The residence time in the 1 cm. long ionizer is  $6.7 \times 10^{-4}$  sec, giving the detection efficiency via electron bombardment as  $8.7 \times 10^{-5}$ . It will be assumed that if metastable Rydberg atoms are the source of this signal, they are detected with unit efficiency via field ionization. Whether a significant fraction of Rydberg atoms can be field ionized depends upon their electronic state. Classically, the field required to ionize a Rydberg atom in state nL is:<sup>22</sup>

$$E_c = 1/(16n^4)$$
 a.u. (11)

Thus a field of 10 kV/cm is required to ionize atoms with n=15.<sup>22</sup> The threshold field strength varies dramatically with n. Since the intensity of the signal at the final peak is one-tenth that due to electron bombardment, the fraction of the metastable atoms that survive in their excited state to the region of the doorknob must be  $9x10^{-6}$ . This means that  $-\ln(9\times10^{-6})=11.6$  or fewer lifetimes of the metastable have passed during the transit from the collision region to the back of the detector. The measured flight times to the final peaks were typically 340 µsec. The lifetime is then at least 340 µsec/11.6 = 29 µsec. The assumption of unit detection efficiency is not correct. The detector geometry is optimized to detect ions that have been accelerated through the quadrupole mass spectrometer, and designed to hit the center of the ion target. If an atom is field ionized it will not necessarily go to the center of the doorknob, and the secondary electrons will not necessarily go to the relatively small sensitive area of the scintillator/photomultiplier. It is not possible to estimate the detection efficiency. This requires that the lifetime be longer than the estimate given above.

If essentially all the available product internal energy goes into excitation of the sodium atoms, the first three levels listed in table III are accessible at the <u>mean</u> collision energy. Of these, only the nP states have lifetimes close to those predicted by the above arguments. Whatever fraction of the atoms might be produced in states other than nP (at this energy) will not be detected by field ionization, since they would not survive the transit to the detector, and this will lower the observed fraction of signal that could be due to field ionization. However, all states up to the ionization limit are accessible to the high velocity components of the beams. Since the Rydberg state lifetimes vary as  $n^{3}$ ,<sup>22</sup> nS, nD, and nF Rydberg levels have lifetimes sufficient to survive the transit to the doorknob for  $n^{-22}$ 

## Table III. Relevant Na atomic energy levels<sup>20</sup> of Na and their lifetimes.<sup>21</sup>

<u>State</u>	Energy (cm <sup>-1</sup> )	<u>Lifetime (µsec)</u>
14 <sup>2</sup> S	40769.5	2.8
14 <sup>2</sup> P	40814.5	22, 27 (J=1/2, 3/2)
13 <sup>2</sup> D	40798.8	2.1
15 <sup>2</sup> S	~40870	3.5
15 <sup>2</sup> P	40901.1	28, 34 (J=1/2, 3/2)
14 <sup>2</sup> D	40890.0	2.6

Table IV. Relevant Na atomic energy levels  $^{20}$  and their lifetimes. $^{21}$ 

State	Energy (cm <sup>-1</sup> )	Lifetime (µsec)
10 <sup>2</sup> S	39983.0	0.88
9 <sup>2</sup> P	39795.	5.1
9 <sup>2</sup> D	40091.	0.68
9 <sup>2</sup> F	40093.	0.77

and above. Additionally, peaks are observed in the TOF spectra too early for field ionization in the doorknob region, and are unrelated to the electron bombardment. The strongest field before the doorknob is that of the quadrupole mass spectrometer, where the electric field is at most 2 kV/cm for standard operating conditions. This also requires that higher Na levels are produced if they are to explain the detector off signal. For field ionization by 2 kV/cm, n must be greater than about 22 (see (11)). All these levels are possible to excite, but require higher beam velocities.

A number of observations contradict the possibility that the observed signal is due to Rydberg atoms. First, the threshold is inconsistent with Rydberg atom formation, since the detector off signal occurs even at threshold collision energies. At 16 kcal/mole collision energy, the highest excitable Rydberg levels are those shown in table IV. With a 5.1  $\mu$ sec lifetime, the Na(9<sup>2</sup>P) states could not possibly survive the transit to the detector, not to mention to the back of the detector. With a shorter than expected flight time to the ionizer of 130 µsec, 130/5.1=25 lifetimes would pass, and only  $e^{-25}=10^{-11}$  of the Rydberg atoms produced would retain their excitation, an undetectably small fraction. Secondly, a detector off signal due to Rydberg atoms would be extremely sensitive to product flight time, and thus to laboratory velocity. The detector on and off distributions would be very different with a strong weighting for observing high laboratory velocities. This is not seen. Thirdly, at a collision energy of 18 kcal/mole, it is necessary to assume that all the internal energy goes

into the Na atoms. If this were true at a collision energy of 22 kcal/mole, the Na atoms would be ionized, whereas similar center-of-mass distributions are observed. It is then necessary to invoke an increasing  $O_2$  internal excitation to exactly compensate for the increasing collision energy.

Thus, while the detector off signal cannot be adequately explained, Rydberg atom production is not consistent with the experimental observations. One possible process that could give such a signal is surface ionization of the very internally excited NaO. It is possible that each surface on which the NaO impinges gives some ion counts. There is no surface 54 cm. from the modulation position, however, for which the final peak is observed. This peak is thus left unexplained. The back of the detector which presumably would also lead to surface ionization is 60 cm. away from the modulation point.

As discussed in chapter II, the absolute reactive cross section can be determined by calibrating the reactive scattering with the small angle elastic scattering. This is done by estimating the absolute differential elastic cross sections due to the van der Waals attraction, and estimating the relative detection efficiencies of the reactively scattered molecules and the elastically scattered atoms.

The Slater-Kirkwood approximation can be used to estimate the van der Waals attraction,<sup>23</sup> as shown in chapter II. The polarizability of  $0_2$  in the gas phase has been calculated to be  $\alpha(0_2)=1.57$  Å<sup>3</sup>.<sup>24</sup> The effective number of electrons of  $0_2$  is  $N(0_2)=12$ . All the

values for Na remain the same as those used in chapter II. Using the Slater-Kirkwood approximation -- equation (10) of chapter II -- yields  $C_{6,disp} = 2620 \text{ kcal/mole } \mathbb{A}^6$ . Since  $O_2$  is a homonuclear diatomic molecule, its electric dipole moment is zero, and there is no inductive contribution to the  $C_6$  constant. Equation (8) of chapter II gives the approximate classical small angle elastic scattering of a van der Waals potential. This is used to calculate the absolute differential cross sections for small angle elastic scattering for comparison with the small laboratory angle signal due to the elastic scattering shown in figure 8. As in chapter II, the center-of-mass angular distribution is determined for only the peak of the beam velocities while distribution from the center-of-mass frame to the laboratory frame distributions and for the elastic recoil energy.

The ionization cross section of Na atoms by 200 eV electrons is known,  $\sigma_{ion}(Na)=2.46 \ \text{Å}^2.19$  The ionization cross section of NaO can be estimated from the ionization cross section of  $O_2$ , and the ratio of the polarizablilities of  $O_2$  and NaO.<sup>25</sup> The molecular polarizability of  $O_2$  is  $\alpha(O_2)=1.57 \ \text{Å}^3.24$  The molecular polarizability of NaO is approximately the sum of the polarizabilities of Na<sup>+</sup> and  $O^-.4,26$  These are  $\alpha(Na^+)=0.155 \ \text{Å}^3$ , and  $\alpha(O^-)=2.85 \ \text{Å}^3,27$  giving the polarizability of NaO as  $\alpha(NaO) = 3.0 \ \text{Å}^3$ . Since the ionization cross section of  $O_2$  by 200 eV electrons is 2.7  $\ \text{Å}^2,28$  the ionization cross section of NaO should be approximately  $\sigma_{ion}(NaO) = 5.2 \ \text{Å}^2$ . Since no signal was observed at  $m/q=39 \ (NaO^+)$ , it is assumed that all NaO ionized fragments to Na<sup>+</sup> + O + e<sup>-</sup>, and thus that the ionization cross section of NaO is the same as the ionization cross section of NaO to Na<sup>+</sup>. Neglecting those molecules not ionized by electron bombardment (which were found to contribute only a small fraction of the signal provided the detector was on), the relative detection efficiency (see equation (14) of chapter II) of NaO/Na is 2.1. Using the small angle elastic scattering, and assuming that the product is indeed NaO, gives the total reactive scattering cross section as  $\sigma_R=0.9$  Å<sup>2</sup> at a collision energy of 18 kcal/mole. If Rydberg atoms are formed, doing the same calibration as discussed earlier in this section gives the total cross section for their production as  $\sigma=2.2$  Å<sup>2</sup>.

## D. Discussion

Of course, the most spectacular feature of these measurements is the state-selectivity observed for the production of NaO + O. Because only the Na( $4^{2}D$ ) state reacts, and the Na( $5^{2}S$ ) state does not, and since both states are known to radiatively populate the Na(3S,3P,4S,4P) states, <sup>29</sup> it is clear that these states also do not react to produce NaO.

The strong back scattering of the NaO product with respect to the incoming Na atoms is evidence of a direct (i.e. no long-lived collision complex is formed), collinear reaction. The especially narrow NaO center-of-mass angular range points up the exceptionally restrictive

constraints on the impact parameter and relative orientation required for reaction to proceed. This requirement of very small impact parameters leading to reactive products puts an upper limit of a few  $\mathbb{A}^2$  on the reactive cross section, which is borne out in the estimated reactive cross section.

The polarization dependences show that the reaction mechanism is completely different than that observed for Na(4D) + HCl in chapter The favored polarization for reaction ( $\phi_{CM}$ ) decreases with II. increasing laboratory detector angle. Since the reaction product is so strongly backwards peaked, the reaction must proceed directly, so that there is a correspondence between laboratory and center-of-mass scattering angles. The smaller the laboratory scattering angle, the smaller the average center-of-mass scattering angle. The favored polarization for reaction appears to stay nearly perpendicular to the Na-O-O axis as shown in figure 28. This could have one or both of the following origins. First, the symmetry chosen by this favored orbital alignment is the least favorable for long range electron transfer (i.e. the opposite of that observed in chapter II). Secondly, this could be the required transition state symmetry for reaction. An interaction between the Na(4d) orbital and the half empty  $\pi^*$  antibonding orbitals could lead to a collinear  $(C_{mv})$  transition state of  $\Delta$  symmetry which would be inaccessible to Na(nS,nP) states (see table I), and would be capable of producing a NaO molecule with  $\pi$  or  $\Delta$  symmetry.





XBL 8511-11568Fig. 28. Higher impact parameters lead to scattering to smaller center-of-mass angles. The required geometry for reaction has the Na-O-O collinear, so that with increasing impact parameter, the Na-O-O axis must be tilted with respect to the relative velocity vector. The Na(4d<sub>2</sub>2) orbital remains perpendicular to the Na-O-O axis as shown.

- 196

The back scattering also shows that the reaction does not proceed via long range electron transfer. The outermost neutral-ionic potential energy curve crossing, using values of 0.854 eV for the ionization potential of the Na(4D) level (this is just the difference between the ionization potential of sodium, 5.138 eV, and the excitation energy of the Na(4D) level, 4.284 eV), and 0.44 eV for the adiabatic electron affinity of  $0_2$ , occurs at  $r_c=35$  Å. This is too far out for any significant interaction of the potential curves, and no electron transfer is expected at this long range. A second reason that this mechanism is not responsible for reaction is that in long range electron transfer there would be no discrimination between excited sodium levels of different symmetries (see table I). Thus, the Na(5S) level would have nearly the same cross section as the Na(4D) level. This is clearly not the case. A third reason is that long range electron transfer would likely lead to formation of a long-lived collision complex which is not observed for the reactively scattered NaO. It is not surprising that a long-lived complex does not lead to NaO formation, since a complex would statistically be most likely to decay to low energy products, in this case quenched Na atoms.

The very narrow product translational energy distribution in figure 26 shows that very little of the excess energy of this reaction goes into translation. How then is the energy distributed among the various forms of internal energy of the products? Since low impact parameter collisions lead to reaction directly, it is unlikely that a significant amount of energy is carried away in rotation. This leaves

electronic excitation of NaO and/or O, and vibrational excitation of the NaO molecule. Figure 29 is a diabatic correlation diagram which shows the relative energies of the reactants and products in the outermost columns. The only accessible electronically excited O atom level is  $O({}^{1}D)$ , 1.98 eV above the ground state. Only the NaO $(X^{2}\pi, A^{2}\Sigma^{+})$ levels are shown since these are the only <u>known</u> levels, although many other levels must exist in the first 3 eV above the ground state. Thus, reaction to electronically and/or vibrationally excited NaO is likely to be responsible for the signal observed, but not enough of the spectroscopy is known to determine what state is produced. It is unlikely that very specific, very high vibrational levels of NaO could be produced to give rise to such a narrow P(E), so that one might expect that a large fraction of the internal energy is in the electronic excitation of the NaO molecule.

The lack of reaction at low collision energy implies a barrier in the entrance channel, which the translational energy is required to surmount.<sup>30</sup> The height of the barrier must be at least 15±1 kcal/mole, as this was the observed threshold for reaction. That a barrier occurs on the particular surface leading to reaction is not terribly surprising considering the myriad of crossing potential surfaces in this region.

In fact, this lack of reaction at low collision energy also gives further evidence that long range electron transfer does not lead to reaction, as there would be no barrier for this process. An alternative explanation for the lack of reaction at low collision energy Fig. 29. Diabatic correlation diagram for Na +  $O_2$  in a collinear approach geometry for states of spin multiplicity 2. No attempt has been made to accurately portray the energies of any of the transition states with the exception of the lowest energy NaO<sub>2</sub> state calculated by M. Alexander in reference 4.



 $Na + O_2 \rightarrow NaO + O$ 

Fig. 29

XBL 863-826

<sup>200</sup> 

would be that the slower the excited atom and the molecule come together, the more likely they are to proceed adiabatically, along the lowest potential energy surface. This would imply long range electron transfer occurs with higher efficiency at lower collision energy, and this does not lead to reaction. At higher collision energy, however, the system is more likely to proceed through the crossing region remaining on the same diabatic (covalent) surface.

Looking once again at the alternative explanation, it is conceivable that Rydberg atom formation takes place. If electron transfer has taken place, and the Na<sup>+</sup> +  $0_2^-$  are separating after the collision, at large separation the electron on the  $0_2^-$  is in approximately the correct position for the Rydberg atom. That is, there is some overlap of the Rydberg orbital with the  $0_2^-$  orbitals. There are many crossings of the potential surfaces at these large separations, and it is conceivable that these could lead to some electron transfer back to the Na Rydberg levels. The very slow recoil (known from the P(E)) would help the Na- $0_2$  system follow adiabatic potentials back to neutral products. If this should occur with the cross section measured to be over 2 Å<sup>2</sup>, it would be highly unusual to say the least.

## E. Conclusions

There is strong circumstantial evidence pointing to the production of NaO, however no direct measurement was successful in definitively establishing this. The NaO is very internally excited, most likely to an excited electronic state. The production of NaO leaves some observations unexplained. An alternative explanation is that if Rydberg atoms are produced by the collisions, they are produced in very specific levels corresponding to a crossing far out in the exit channel. The level requirement comes from the need for a fraction of the excitation to survive to the back of the detector and to be field ionized by 2-10 kV/cm electric fields.

The Na +  $0_2$  reaction most likely produces NaO when the Na(4D) level is excited, but not when the Na(3S,3P,4S,4P,5S) levels are. The translational threshold to reaction is 15±1 kcal/mole, implying an entrance channel barrier of at least this magnitude. The reaction does not proceed via long range electron transfer, nor via complex formation, despite the favorable energetics of each. The NaO produced is very internally excited with less than 2% of the total available energy going into relative product translation. The neutral chemical reaction is only a minor product channel, with a reactive cross section at a collision energy of 18 kcal/mole of  $\sigma_R=0.9$  Å<sup>2</sup>. Clearly, further experiments are required to sort out the products of this reaction. Of particular interest would be to spectroscopically determine the neutral products and product states. It would also be interesting, and relatively simple, to see what the ionic products of the collisions are.
# F. References

- 1. Calculated from  $D_0^0$  values in reference 2.
- K. P. Huber and G. Herzberg, <u>Molecular Spectra and Molecular</u> <u>Structure, IV. Constants of Diatomic Molecules</u> (Van Nostrand Reinhold, Co., New York, 1979).
- 3. see figure 9 of chapter I.
- 4. M. H. Alexander, <u>J. Chem. Phys. 69</u>, 3502 (1978).
- 5. L. Andrews, J. Phys. Chem. 73, 3922 (1969).
- 6. R. R. Smardzewski and L. Andrews, J. Chem. Phys. 57, 1327, (1972).
- F. J. Adrian, E. L. Cochran, and V. A. Bowers, <u>J. Chem. Phys. 59</u>, 56 (1973).
- 8. P. A. G. O'Hare and A. C. Wahl, J. Chem. Phys. 56, 4516 (1972).
- 9. J. Pfeiffer and J. L. Gole, <u>J. Chem. Phys.</u> 80, 565 (1980).
- 10. J. L. Gole and P. A. Schulz, private communication (1985).
- V. Kempter, W. Mecklenbrauck, M. Menzinger, Ch. Schlier, <u>Chem.</u>
   Phys. Lett. 11, 353 (1971).
- 12. V. Kempter, Adv. Chem. Phys. 30, 417 (1975).
- R. H. Neynaber, B. F. Myers, and S. M. Trujillo, <u>Phys. Rev. 180</u>, 139 (1969).
- 14. A. W. Kleyn, Ph.D. Thesis, Stichting voor Fundamenteel Onderzoek der Materie, Amsterdam, The Netherlands (1982).
- 15. A. W. Kleyn, M. M. Hubers, and J. Los, <u>Chem. Phys. 34</u>, 55 (1978).

- 16. A. P. M. Baede, <u>Adv. Chem. Phys. 30</u>, 463 (1975), and references therein.
- 17. G. J. Schulz, <u>Rev. Mod. Phys. 45</u>, 423 (1973), and H. H. Michels and F. E. Harris, in <u>Seventh International Conference on the</u> <u>Physics of Electronic and Atomic Collisions: Abstracts of Papers,</u> <u>Volume II</u>, 1170 (North-Holland, Amsterdam, 1971).
- E. W. Rothe, B. P. Mather, and G. P. Reck, <u>Chem. Phys. Lett. 51</u>, 71 (1977).
- 19. R. H. McFarland and J. D. Kinney, Phys. Rev. 137, A1058 (1965).
- 20. C. E. Moore, National Bureau of Standards (U. S.) Circ. 467 (1949).
- 21. C. E. Theodosiou, Phys. Rev. A 30, 2881 (1984).
- S. A. Edelstein and T. F. Gallagher, <u>Adv. At. Mol. Phys. 14</u>, 365 (1978).
- J. H. Birely, R. R. Herm, K. R. Wilson, and D. R. Herschbach, <u>J.</u> Chem. Phys. 47, 993 (1967).
- 24. S. Nir, S. Adams, and R. Rein, J. Chem. Phys. 59, 3341 (1973).
- 25. C. H. Becker, P. Casavecchia, P. W. Tiedemann, J. J. Valentini, and Y. T. Lee, J. Chem. Phys. 73, 2833 (1980).
- 26. J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, <u>Molecular</u> <u>Theory of Gases and Liquids</u> (John Wiley and Sons, Inc., New York, 1964).
- 27. J. Thorhallsson, C. Fisk, and S. Fraga, <u>Theoret. Chim. Acta</u> (Berl.) 10, 388 (1968).
- S. C. Brown, <u>Basic Data of Plasma Physics</u>, 2<sup>nd</sup> Ed. (MIT Press, Cambridge, Massachusetts, 1967).

- 29. see Chapter I of this thesis, or G. Jamieson, W. Reiland, C. P. Schulz, H.-U. Tittes, and I. V. Hertel, <u>J. Chem. Phys. 81</u>, 5805 (1984).
- 30. A. M. G. Ding, L. J. Kirsch, D. S. Perry, and J. C. Polanyi, <u>Faraday Disc. Chem. Soc. 55</u>, 252 (1973).

# IV. THE REACTIONS OF GROUND AND EXCITED STATE ALKALI ATOMS WITH METHYL HALIDE MOLECULES

## A. Introduction

The reactions of alkali atoms (M) with methyl halide molecules  $(CH_3X)$  have been extensively studied for more than 50 years.<sup>1-12</sup> The first crossed molecular beam experiment on K +  $CH_3I$  revealed the rebound mechanism found in all subsequent experiments.<sup>2</sup> The elegant work of Bernstein and coworkers,<sup>4</sup> and Brooks and coworkers<sup>5</sup> in orienting methyl halide molecules in a molecular beam prior to reaction confirmed the steric effect earlier ascribed to the reaction. That is, for small impact parameter collisions if the alkali atom comes in at the halogen end of the molecule, then the reaction probability is unity, but if the alkali atom comes in at the methyl end of the molecule, then the reaction probability is zero.<sup>2</sup> Bernstein and coworkers have correlated this steric effect with a cardioid dependence of reaction probability on molecular orientation, as shown for Rb +  $CH_3I$  in the polar plot in figure 1 (reproduced with permission from reference 4).

The result of this steric effect is a backwards peaked angular distribution (with respect to the incoming alkali atoms' direction). The farthest backwards scattered MX product (highest laboratory angle) corresponds to the most collinear C-X-M geometry. Lower C-X-M approach



XBL 852-1140

Fig. 1. The dependence of reaction probability on methyl iodide orientation in the reaction  $Rb + CH_3I$  as measured by Parker et al. Reproduced with permission from reference 4.

angles ( $\alpha$ ) correspond to smaller center-of-mass scattering angles, and thus smaller laboratory scattering angles.

The mechanism for this reaction has an electron transferring from the alkali atom to the methyl halide molecule. This electron goes into the lowest unoccupied molecular orbital which is an antibonding orbital of the C-X bond ( $\sigma^*$  in diatomic terms) which causes the rupture of this bond. It was for the reaction of alkali atoms with halogen and alkali halide molecules that Kuntz, Mok, and Polanyi originally proposed the "Direct Interaction with Product Repulsion" (DIPR) model.<sup>13</sup> In this model only two interactions are taken into account; these are:  $[CH_3-X]^-$ , and  $M^+-X^-$ . Furthermore, the DIPR model uses a fixed repulsive energy release  $R_D$  to determine the translational energy release in the reaction.

The ideas in this model were extended by Harris and Herschbach to explain the dynamics of these alkali + methyl halide reactions more quantitatively.<sup>3,14</sup> Their work assumes a distribution of product energies as in photodissociation of the halogen containing molecule, thus the acronym DIPR-DIP -- "Direct Interaction with Product Repulsion - Distributed as In Photodissociation." This assumption is based on the fact that the molecular orbital configuration of  $CH_3X$  in the united atom (diatomic-like) approximation is:

$$\dots (\sigma_{g})^{2} (\pi_{u})^{4} (\pi_{g}^{*})^{4} (\sigma_{u}^{*})^{0}.$$
(1)

The  $\pi_g^*$  orbital is weakly antibonding, and the  $\sigma_u^*$  orbital is strongly antibonding with a node between the C and X atoms. The first

excited molecular orbital configuration, and the one that leads to photodissociation is then:

... 
$$(\sigma_g)^2(\pi_u)^4(\pi_g^*)^3(\sigma_u^*)^1.$$
 (2)

This is very similar to the ground state  $CH_3X$  molecular orbital configuration:

$$\dots (\sigma_{g})^{2} (\pi_{u})^{4} (\pi_{g}^{*})^{4} (\sigma_{u}^{*})^{1}, \qquad (3)$$

the only difference being the extra electron in the  $\pi_g^*$  orbital which makes little contribution to  $CH_3-X$  potential. Harris and Herschbach have concluded that the potential curves for  $CH_3X^*$  and  $CH_3X^-$  have very similar shapes and in particular similar slopes on their repulsive walls.

The DIPR-DIP model predicts very little variation of product angular and momentum distributions with changing mass combinations and alkali ionization potential. This was borne out in the crossed molecular beams experimental data on the reactions of various alkali atoms with a number of alkyl monoiodides.<sup>3,12</sup> A relatively large change in product angular distributions is predicted with varying geometric reaction probability. That is, as different collision geometries lead to reaction, the product angular distributions change substantially. This is expected from such an impulsive model.

Of the methyl monohalides, only the reactions of alkali atoms with  $CH_3I$  and  $CH_3Br$  have been studied in crossed molecular beams.<sup>12</sup> The reactions of all of the alkalis with  $CH_3I$  have been studied.<sup>12</sup> Many reactions of alkali atoms with alkyl polyhalides have also been studied in crossed beams, but these follow stripping mechanisms of the type described in the following chapter, and will not be discussed here. All the reactions of alkali atoms with alkyl monohalides studied follow the rebound mechanism described above.<sup>3,12</sup>

Bernstein and coworkers have measured the dependence of the reactive cross sections  $(\sigma_p)$  of:

$$K + CH_3 I \rightarrow KI + CH_3$$
(4)

$$Rb + CH_{3}I \rightarrow RbI + CH_{3}$$
(5)

as a function of collision energy.<sup>6,8</sup> They estimate that their values of the absolute cross sections are good to within a factor of two, but their relative values are good to 15%.<sup>6,8</sup> The cross section of process (4) peaks at  $\sigma_R=35$  Å<sup>2</sup> at 4 kcal/mole and drops to 1/3 of its peak value by 20 kcal/mole collision energy.<sup>6</sup> The cross section of process (5) drops steadily from  $\sigma_R=23$  Å<sup>2</sup> at 3 kcal/mole, falling to 40% of its low energy value at 20 kcal/mole collision energy.<sup>8</sup> Goldbaum and Martin have measured the reactive cross section for:

$$K + CH_3Br \rightarrow KBr + CH_3$$
 (6)

to be  $\sigma_R=3$  Å<sup>2</sup> by comparing it to the 35 Å<sup>2</sup> cross section for process (4).<sup>10</sup>

In process (5), with increasing collision energy, the RbI product is scattered less backwards (relative to the motion of the incoming alkali atom) in the center-of-mass frame, and actually peaks at angles lower than 180° for collision energies of 22 and 36 kcal/mole.<sup>9</sup> The product recoil energies increases approximately linearly with collision energy.<sup>9</sup> This behavior of the product recoil energy is also observed for K +  $CH_{3I}$ .<sup>7,11</sup> The increased product translational energy with increasing reactant translational energy is explained in terms of being able to run up higher on the repulsive walls of a potential surface, acquiring more potential energy than possible for lower collision energies, and then having a release of this energy into product translation.<sup>15</sup> This can also explain the center-of-mass angular distributions to some extent. The release of the more repulsive geometries accessible only to the higher collision energies allows scattering to wider angles than those seen at lower collision energies.

Polanyi's diffusion flame studies showed that for Na +  $CH_3X$  the reaction cross sections increase by several orders of magnitude as X is changed through the sequence F, Cl, Br, I.<sup>1</sup> Recent measurements of thermal energy rate constants for Na +  $CH_3F$ ,  $CH_3Cl$ ,  $CH_3Br$  concur with this strong trend in the reactivity.<sup>16</sup>

#### B. Results

The reaction:

 $Na + CH_3Br \Rightarrow NaBr + CH_3$   $\Delta H_0^0 = -20.7 \text{ kcal/mole}^{17}$  (7) was studied at a collision energy of 21 kcal/mole in the apparatus described in chapter I. This corresponds to seeding the sodium atoms in helium, and running a neat supersonic molecular beam of  $CH_3Br$ . A Newton velocity vector diagram for these conditions is shown in figure 2. The angular distributions recorded for  $Na(3S, 3P, 4D) + CH_{3}Br$  are shown in figure 3. The statistics for this data are very poor, but these distributions took 72 hours to record, so that a brute force attempt to improve their quality would not be fruitful. The data shows qualitatively similar angular distributions for the reactions of each of the three Na levels. The only difference between the distributions due to the Na(3S) and Na(3P) states appears to be on the low angle side of the reactive peak. This difference is more pronounced for the reactive scattering of the Na(4D) atoms.

Before the second laser became available, this region of the angular distribution for  $Na(3S, 3P) + CH_3Br$  was expanded over the useful angular range of the detector by seeding the  $CH_3Br$  in helium, as shown in the Newton diagram in figure 4 (note that the collision energy is increased to 25 kcal/mole). The NaBr product angular distribution (or more accurately -- partial angular distribution) obtained is shown in figure 5. The same behavior is exhibited in that the distribution of the  $Na(3P) + CH_3Br$  remains higher than the distribution for  $Na(3S) + CH_3Br$  in this region.

The statistics were too poor for this experiment to allow the measurement of time-of-flight or polarization dependences. These would certainly be of interest.

Against unfavorable kinematics and the reduced reaction probability,  $^{1,10,16}$  an attempt was made to measure the differential cross section of:

Na + CH<sub>3</sub>Cl  $\rightarrow$  NaCl + CH<sub>3</sub>.  $\Delta H_0^0 = -13.5 \text{ kcal/mole}^{20}$  (8)

E<sub>coll</sub>=21 kcal/mole



----- 5.0x10<sup>4</sup> cm/sec

XBL 863-852Fig. 2. Newton velocity vector diagram for Na(3S,3P,4D) +  $CH_3Br \Rightarrow$ NaBr +  $CH_3$  at a collision energy of 21 kcal/mole.



XBL 863-874

Fig. 3. NaBr product angular distributions for Na(3S,3P,4D) +  $CH_3Br$ at a collision energy of 21 kcal/mole.

20 0 10 30 40 / 1 ,50 <u>\_\_60</u> <u>6</u>5 V Na -70 3.0x10<sup>5</sup> cm/sec - 80 ЗS ЗP -90 V MeBr 1.5x10<sup>5</sup> cm/sec

\_\_\_\_ 5.0x10<sup>4</sup> cm/sec

XBL 863-854



216

E<sub>coll</sub>=25 kcal/mole



XBL 863-873

# Fig. 5. NaBr product angular distributions for $Na(3S, 3P) + CH_3Br$ at a collision energy of 25 kcal/mole.

Although some reactive scattering signal was observed, it was not possible to measure a laboratory angular distribution. A further complication was that the signal at m/q=23 (Na<sup>+</sup>) appeared to be contaminated by elastically scattered CH<sub>3</sub>Cl and modulated background at m/q=23.5-25 (CH<sub>0-3</sub>Cl<sup>++</sup>) due to fragmentation of CH<sub>3</sub>Cl in the electron bombardment ionizer. The m/q=58 (NaCl<sup>+</sup>) signal was too small to measure.

## C. Analysis of Experimental Results

As in the previous chapters the data was fitted using the program GM, a derivative of CMLAB,<sup>22</sup> as described in chapter I. Independent product center-of-mass angular and recoil energy distributions were assumed, and a laboratory angular distribution was calculated. Ordinarily this would be extremely difficult to do in the absence of product velocity (e.g. time-of-flight) measurements, however due to sharp back scattering the transformation from the center-of-mass frame to the laboratory frame is nearly one-to-one.

Due to the poor statistics in the angular distributions, and to the similarities between the angular distributions for the different states, rather than try to subtract the already noisy distributions, the raw data shown in figure 3 was used to represent the data for each level excited. This has the effect of underestimating the differences in the scattering of the different states, but should still be able to

<u>Na Level</u>	FWHM of T( <del>o</del> )	% Backward Scattered	Peak Recoil Energy	$\sigma(nL)/\sigma(3S)$
3S	73°	84%	33 kcal/mole	1
3P	74°	84%	33 kcal/mole	1.6
4D	83°	79%	33 kcal/mole	2.8

Table I. A summary of features of the center-of-mass distributions fitted to the laboratory angular distributions for Na + CH3Br at a collision energy of 21 kcal/mole.

point out the trends in the reactivity. Elastic scattering was estimated using the small angle elastic scattering formula given in chapter II, and the  $C_6$  constants estimated below. The signal subtracted away is shown as the dashed lines in figure 3.

The best fits to the angular distributions are shown as the solid curves in figures 6a, b, and c. The center-of-mass angular distributions used are shown in figure 7, and the recoil energy distributions used are shown in figure 8 for Na(3S,3P,4D) scattering, respectively. Some features of the fits are shown in table I. The relative cross sections (right hand column) have been corrected for the excitation fractions given in chapter I. It is assumed that 20% of the Na atoms are in the Na(3P) state when the  $D_2$  transition is excited. It is also assumed that this does not change on subsequent excitation for which it is assumed that 16% of the atoms are in the Na(4D) state. Fig. 6. Best fits to the laboratory angular distributions for Na + CH<sub>3</sub>Br at a collision energy of 21 kcal/mole for a) Na(3S),
b) Na(3P), and c) Na(4D). The solid lines are the fit to the data.



XBL 862-664



XBL 862-665









XBL 862-663

Fig. 7. Center-of-mass angular distributions used for the fits shown in figure 6 of Na +  $CH_3Br \Rightarrow NaBr + CH_3$  at a collision energy of 21 kcal/mole for Na(3S,3P,4D).



Fig. 8. Product recoil energy distributions for the fits shown in figure 6 of Na + CH<sub>3</sub>Br at a collision energy of 21 kcal/mole for Na(3S,3P,4D).

As in chapter II, the reactive cross sections can be calibrated using the low angle elastic scattering.<sup>23,24</sup> Briefly, the low angle elastic scattering is estimated from the small angle approximation to the classical scattering due to a van der Waals potential,<sup>24,25</sup> formula (8) of chapter II. This scattering depends upon the collision energy and the C<sub>6</sub> constant of the van der Waals potential. The C<sub>6</sub> constant is estimated using the Slater-Kirkwood formula for the dispersive portion,<sup>26,27</sup> formula (10) of chapter II, and the Debye formula for the dipole-induced-dipole portion,<sup>28,29</sup> formula (11) of chapter II. The two contributions are then added to get the C<sub>6</sub> constant,<sup>23,29</sup> as in formula (9) of chapter II.

The Slater-Kirkwood formula estimates the dispersive portion of the van der Waals potential from the polarizabilities ( $\alpha$ ) and effective number of electrons (N) of the collision partners. For Na, these are  $\alpha(Na(3S))=24.5 \text{ Å}^3, {}^{30}\alpha(Na(3P))=53.6 \text{ Å}^3, {}^{31} \text{ and } N(Na)=1$ . The polarizability of CH<sub>3</sub>Br can be estimated from the sum of the polarizabilities of its bonds:<sup>28</sup>

$$\alpha(CH_{3}Br) = \sum_{\substack{a11\\bonds}} \alpha_{bond} = \sum_{\substack{a11\\bonds}} \frac{1}{3} (\alpha_{//} + 2\alpha_{\perp}), \qquad (9)$$

where  $\alpha_{//}$  and  $\alpha_{\perp}$  are the bond polarizabilities parallel and perpendicular to the bond axis. The bond polarizabilities are  $\alpha_{//}(C-Br)=5.04 \text{ Å}^3$ ,  $\alpha_{\perp}(C-Br)=2.88 \text{ Å}^3$ ,  $\alpha_{//}(C-H)=0.79 \text{ Å}^3$ , and  $\alpha_{-}(C-H)=0.58 \text{ Å}^3$ .<sup>28</sup> From (8), the polarizability of methyl bromide is  $\alpha(CH_3Br)=5.55 \text{ Å}^3$ . The effective number of electrons is

N(CH<sub>3</sub>Br)=14. This gives the dispersive portion of the C<sub>6</sub> constant for Na(3S) scattering to be C<sub>6,disp</sub>=8830 kcal/mole Å<sup>6</sup>. For Na(3P) scattering, C<sub>6,disp</sub>=13560 kcal/mole Å<sup>6</sup>.

The dipole-induced-dipole contribution is determined from the Na atomic polarizability and the electric dipole moment of  $CH_3Br$ , using the Debye formula.<sup>28,29</sup> The electric dipole moment of methyl bromide is  $\mu(CH_3Br)=1.81 \text{ D.}^{32}$  This gives the inductive portion of the  $C_6$  constant for ground state scattering as  $C_{6, \text{ind}}=1160 \text{ kcal/mole } \text{Å}^6$ , and the total ground state  $C_6$  constant as 9990 kcal/mole  $\text{Å}^6$ . For Na(3P) scattering,  $C_{6, \text{ind}}=2530 \text{ kcal/mole } \text{Å}^6$ , and the total  $C_6=16090 \text{ kcal/mole } \text{Å}^6$ .

The absolute differential elastic scattering cross sections are. then estimated as described in chapter II, and compared to the low laboratory angle scattering intensity. A least squares fit is made to the Towest laboratory angles measured, and a scaling factor is determined between the absolute differential cross sections and the relative numbers measured in the scattering experiment. In order to convert the reactive cross sections into absolute cross sections, the relative detection efficiency of reactively and elastically scattered products must be known. The relative detection efficiency in this case is the ratio of the ionization cross sections for Na and NaBr to Na<sup>+</sup>. The ionization cross section of Na by 200 eV electrons has been measured as  $\sigma(Na)=2.1$  Å<sup>2</sup>.<sup>33</sup> The ionization cross section of NaBr can be estimated by comparing its polarizability to that of a species for which the ionization cross section is known.<sup>23,34</sup> The

Na Level	From 3S Data	From 3P Data	Average
35	5.6 Å <sup>2</sup>	5.8 Å <sup>2</sup>	5.7 Å <sup>2</sup>
3P	8.7 Å <sup>2</sup>	9.0 Å <sup>2</sup>	8.9 Å <sup>2</sup>
4D	15.7 Å <sup>2</sup>	16.2 Å <sup>2</sup>	16.0 Å <sup>2</sup>

Table II. Reactive cross sections of  $Na(nL) + CH_3Br \Rightarrow NaBr + CH_3$ .

polarizability of NaBr is approximately the sum of the polarizabilities of Na<sup>+</sup> and Br<sup>-</sup>.<sup>23</sup> These are  $\alpha(Na^+)=0.227$  Å<sup>3</sup>, and  $\alpha(Br^-)=3.585$  Å<sup>3</sup> in the NaBr molecule,<sup>35</sup> giving  $\alpha(NaBr)=3.812$  Å<sup>3</sup>. By comparing this to the polarizability of isoelectronic Kr,  $\alpha(Kr)=2.86$  Å<sup>3</sup>,<sup>36</sup> for which the cross section for ionization by 200 eV electrons is  $\sigma(Kr)=2.8$  Å<sup>2</sup>,<sup>37</sup> a total ionization cross section of  $\sigma(NaBr)=3.7$  Å<sup>2</sup> is derived. If it is assumed as for vibrationally excited NaCl that 90% fragmentation to Na<sup>+</sup> occurs, then the cross section for ionization to Na<sup>+</sup> is  $\sigma'(NaBr)=3.3$  Å<sup>2</sup> This gives a relative detection efficiency of Eff(NaBr)/Eff(Na)=1.6.

The cross sections derived for the reactions of Na(3S,3P,4D) atoms are shown in table II. As pointed out in chapter II, the absolute cross sections are probably only good to within a factor of 2, and are strongly dependent on the excitation fraction assumed. Since the reactive geometries are exactly those that would lead to large angle elastic scattering, and reactive signal appears at the highest laboratory angles studied, the hard sphere approximation cannot be used to estimate the reactive cross sections.

#### D. Discussion

1. Observations on the Experimental Results

The only differences with increasing electronic energy appear to be the increasing reactive cross sections and the slowly widening center-of-mass angular distributions. The rebound mechanism is seen for the reactions of the excited states as well as the ground state of Na with  $CH_3Br$ .

The slightly widening center-of-mass angular distributions (see figure 7 and table I) support the idea that an opening of the steric angle for reaction occurs with increasing electronic energy and thus sodium orbital size. The less strict geometrical requirements on the approach geometry lead to the broader center-of-mass angular distributions. This is visible in the raw data (figure 3) as the increase in signal on the low side of the reactive peak.

The fact that the recoil energy distributions do not shift with changing electronic energy implies that the idea of ions released at a distance and coming together to form a vibrationally excited ionic molecule as brought forth in the DIPR-DIP model appears to be correct. That is that as the electron jump distance moves further out for successively higher electronic states of sodium (with successively lower ionization potentials) the distance at which the  $M^+-X^$ attraction takes over becomes larger, and the resultant  $M^+X^$ molecule is produced in successively higher vibrational levels. A nearly vertical  $CH_3Br \rightarrow CH_3Br^-$  transition takes place so that the translational energy released is essentially the same for each electronic state.

# 2. Other Possible Processes

Gersh and Bernstein have discussed energetically possible products in high energy (>1 eV) collisions of  $K + CH_3I$ , and several of these processes, <u>mutatis mutandis</u>, are energetically possible here:

Na +	$CH_3Br \Rightarrow NaCH_3 + Br$	∆H <sub>o</sub> =+26 kcal/mole <sup>38</sup>	(10)
Na +	CH <sub>3</sub> Br → NaH + CH <sub>2</sub> Br	∆H <sub>0</sub> =+59 kcal/mole <sup>39</sup>	(11)
Na +	CH <sub>3</sub> Br → Na + CH <sub>3</sub> + Br	$\Delta H_0^0 = +66.2 \text{ kcal/mole}^{17}$	(12)
Na +	CH <sub>3</sub> Br → Na <sup>+</sup> + CH <sub>3</sub> Br <sup>-</sup>	ΔH <sup>0</sup> =+103.5 kcal/mole <sup>40</sup>	(13)
Na +	$CH_{3}Br \rightarrow Na^{+} + CH_{3} + Br^{-}$	$\Delta H_0^0 = +107.3 \text{ kcal/mole}^{44}$	(14)

All these processes (10-14) are substantially endothermic, but with the electronic excitation to Na(4D), 98.8 kcal/mole, and the collision energy of over 20 kcal/mole, all are energetically possible.

It is unlikely that methyl sodium is formed via process (10) because this would require Na to approach the methyl end of the molecule. As shown for K +  $CH_3I$ , there is a substantial repulsion for the alkali approaching the methyl end of the methyl halide<sup>45</sup> in

the ground state potential, and this is certain to hold true for the excited state potentials as well.

Formation of NaH (process (11)) is also unlikely because it involves approach via a repulsive geometry and the breakage of the strong H-CH<sub>2</sub>Br bond. No evidence is seen for NaH, although it is also expected to be observed at m/q=23 (Na<sup>+</sup>).

It is possible that at high collision energy the Na could dissociate the  $CH_3Br$  as in process (12). This is most likely to happen where effectively NaBr is produced above its dissociation limit to atoms -- a neutral equivalent to the observations of Lacmann and Herschbach for K + HCl, <sup>46</sup> discussed further below.

The present experiments are not sensitive to the production of ions such as in processes (13-14). Futhermore,  $CH_3Br^-$  has such a small potential well (3.8 kcal/mole), offset significantly from the equilibrium geometry of the neutral,<sup>42</sup> that it is extremely unlikely that it could be produced. However, it is quite possible that the Na<sup>+</sup> + Br<sup>-</sup> ion pair could be produced via process (14). This is equivalent to the process seen by Lacmann and Herschbach for K + HCl.<sup>46</sup> As interpreted by Balint-Kurti and Yardley highly vibrationally excited KCl is produced above the dissociation limit to ions.<sup>47</sup> An experiment such as Mariella's described in chapter III would be able to observe such a process and measure its cross section.<sup>48</sup>

# E. Conclusions

The most striking feature of the reactions of the various states is their similarity. The rebound mechanism describes the reaction of the ground and excited Na states with CH<sub>3</sub>Br. Besides the increasing cross sections, the main difference in the reaction dynamics appears to be that the steric angle of acceptance opens somewhat for higher electronic states, and thus for larger electronic orbitals. Of course the lack of a change in the product translational energy with increasing electronic energy implies that the product internal energy increases substantially.

Further study, perhaps of the more reactive methyl iodide system, could provide insight into the symmetry aspects of the reactions of the excited states.

#### F. References

- M. Polanyi, <u>Atomic Reactions</u> (Williams and Northgate, London, 1932).
- D. R. Herschbach, G. H. Kwei, J. A. Norris, <u>J. Chem. Phys. 34</u>, 1842 (1961).
- 3. D. R. Herschbach, Faraday Disc. Chem. Soc. 55, 233 (1973).
- 4. D. H. Parker, K. K. Chakravorty, and R. B. Bernstein, <u>J. Phys.</u> Chem. 85, 466 (1981).
- 5. G. Marcelin and P. R. Brooks, J. Am. Chem. Soc. 97, 1710 (1975).
- 6. M. E. Gersch and R. B. Bernstein, J. Chem. Phys. 56, 6131 (1972).
- 7. A. M. Rulis and R. B. Bernstein, J. Chem. Phys. 57, 5497 (1972).
- H. E. Litvak, A. Gonzalez Urena, and R. B. Bernstein, <u>J. Chem.</u>
   Phys. 61, 4091 (1974).
- 9. A. Gonzalez Urena and R. B. Bernstein, ibid., 4101 (1974).
- R. H. Goldbaum and L. Robbin Martin, <u>J. Chem. Phys. 62</u>, 1181 (1975).
- 11. G. Rotzoll, R. Viard, and K. Schuegerl, <u>Chem. Phys. Lett. 35</u>, 353 (1975).
- 12. R. R. Herm, in <u>Alkali Halide Vapors: Structure, Spectra, and</u> <u>Reaction Dynamics</u>, 189 (Academic Press, New York, 1979), and references therein.
- P. J. Kuntz, M. H. Mok, and J. C. Polanyi, <u>J. Chem. Phys. 50</u>, 4623 (1969).

- R. M. Harris, Ph.D. Thesis, Harvard University, Cambridge, Massachusetts (1970).
- 15. A. M. G. Ding, L. J. Kirsch, D. S. Perry, and J. C. Polanyi, Faraday Disc. Chem. Soc. 55, 252 (1973).
- 16. D. Husain and P. Marshall, Int. J. Chem. Kin. 18, 83 (1986).
- 17. Calculated from the  $H_3C$ -Br dissociation energy (66.2 kcal/mole) in reference 18, and the  $D_0^0$  value for NaBr (86.9 kcal/mole) given in reference 19.
- G. N. A. van Veen, Ph. D. Thesis, Rijksuniversiteit te Utrecht, Utrecht, The Netherlands (1984).
- 19. K. P. Huber and G. Herzberg, <u>Molecular Spectra and Molecular</u> <u>Structure, IV. Constants of Diatomic Molecules</u> (Van Nostrand Reinhold, Co., New York, 1979).
- 20. Calculated from the  $H_3C-Cl$  bond strength (84.0 kcal/mole) in reference 21, and the  $D_0^0$  value for NaCl (97.5 kcal/mole) given in reference 19.
- 21. "CRC Handbook of Chemistry and Physics," 65th Ed., R. C. Weast, ed., (CRC Press, Boca Raton, Florida, 1984).
- 22. R. J. Buss, Ph.D. Thesis, University of California, Berkeley, California (1979).
- 23. See chapter II of this thesis.
- 24. J. H. Birely, R. R. Herm, K. R. Wilson, and D. R. Herschbach, <u>J.</u> Chem. Phys. 47, 993 (1967).
- 25. E. A. Mason, J. T. Vanderslice, and C. J. G. Raw, <u>J. Chem. Phys.</u> 40, 2153 (1964).

- 26. J. C. Slater and J. G. Kirkwood, <u>Phys. Rev. 37</u>, 682 (1931).
- 27. K. S. Pitzer, Adv. Chem. Phys. 2, 59 (1959).
- 28. J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, <u>Molecular</u> <u>Theory of Gases and Liquids</u> (John Wiley and Sons, Inc., New York, 1964).
- 29. H.-J. Werner and W. Meyer, Phys. Rev. A 13, 13 (1976).
- 30. P. Hannaford, W. R. MacGillivray, and M. C. Standage, <u>J. Phys. B:</u> Atom Molec. Phys. 12, 4033 (1979).
- 31. H. T. Duong and J.-L. Picque, J. Physique 33, 513 (1972).
- 32. R. D. Nelson, Jr., D. R. Lide, Jr., and A. A. Maryott, <u>National</u> <u>Standard Reference Data Series 10</u>, National Bureau of Standards (U. S.) (1967).
- 33. R. H. McFarland and J. D. Kinney, Phys. Rev. 137, A1058 (1965).
- 34. C. H. Becker, P. Casavecchia, P. W. Tiedemann, J. J. Valentini, and Y. T. Lee, J. Chem. Phys. 73, 2833 (1980).
- 35. P. Brumer and M. Karplus, J. Chem. Phys. 58, 3903 (1973).
- 36. J. Thorhallsson, C. Fisk, and S. Fraga, <u>Theoret. Chim. Acta</u> (Berl.) 10, 388 (1968).
- 37. S. C. Brown, <u>Basic Data of Plasma Physics</u>, 2<sup>nd</sup> Ed. (MIT Press, Cambridge, Massachusetts, 1967).
- 38. Calculated from the H<sub>3</sub>C-Br dissociation energy (66.2 kcal/mole) in reference 18, and an estimated Na-CH<sub>3</sub> bond strength of 40 kcal/mole.

- 39. Calculated from the H-CH<sub>2</sub>Br dissociation energy (102 kcal/mole) in reference 21, and the  $D_0^0$  value for NaH (43 kcal/mole) given in reference 19.
- 40. Calculated from the ionization potential of Na (118.5 kcal/mole),<sup>41</sup> and the H<sub>3</sub>C-Br<sup>-</sup> bond strength of 3.8 kcal/mole from reference 42, and the electron affinity of Br (77.4 kcal/mole) from reference 43.
- 41. W. L. Wiese, M. W. Smith, and B. M. Miles, <u>Atomic Transition</u> <u>Probabilities Volume II: Sodium Through Calcium</u>, Natl. Stand. Ref. Data Ser., Nat. Bur. Stand. (U. S.), Washington, DC (1969).
- 42. W. E. Wentworth, R. George, and H. Keith, <u>J. Chem. Phys. 51</u>, 1791 (1969).
- 43. R. S. Berry and C. W. Riemann, J. Chem. Phys. 38, 1540 (1963).
- 44. Calculated from the H<sub>3</sub>C-Br dissociation energy (66.2 kcal/mole) in reference 18, the ionization potential of Na (118.5 kcal/mole),<sup>41</sup> and the electron affinity of Br (77.4 kcal/mole) from reference 42.
- 45. L. M. Raff and M. Karplus, <u>J. Chem. Phys. 44</u>, 1212 (1966), and references therein.
- 46. K. Lacmann and D. R. Herschbach, Chem. Phys. Lett. 6, 106 (1970).
- 47. G. G. Balint-Kurti and R. N. Yardley, <u>Faraday Disc. Chem. Soc. 62</u>, 77 (1977).
- 48. R. P. Mariella, Jr., <u>J. Chem. Phys. 76</u>, 2965 (1981).

# V. THE REACTIONS OF GROUND AND EXCITED STATE ALKALI ATOMS WITH HALOGEN MOLECULES

## A. Introduction

The reactions of alkali atoms (M) with diatomic halogen molecules  $(X_2)$  were shown to have exceptionally large cross sections in Michael Polanyi's classic diffusion flame experiments in the early 1930's.<sup>1</sup> These reactions have been studied in detail in crossed molecular beams for over 20 years.<sup>2-16</sup> The alkali halide product is in all cases found to be: sharply peaked forwards, produced with large cross section, and with little translational energy. In analogy to nuclear stripping reactions, these reactions have been regarded as "spectator stripping reactions," in which the alkali atom and one halogen atom interact but leave the other halogen atom (the spectator) essentially unperturbed.<sup>2,3</sup> It was for these reactions that the "harpoon" mechanism was put forth.<sup>1,17</sup> The basic reaction mechanism has an electron transfering to (harpooning) the halogen molecule at long range. The alkali ion then "reels in" the halide ion. This leads to larger than gas kinetic cross sections as Polanyi originally observed.<sup>1</sup>

The spectator stripping features of these reactions can be quantified.<sup>18</sup> Since the spectator halogen atom is expected to continue undisturbed by the reaction, the product alkali halide molecule must remain confined to the center-of-mass angle and velocity defined by the alkali atom and the halogen <u>atom</u> involved in the reaction. As illustrated in figure 1, this means that the alkali halide product must appear at the laboratory angle

$$\boldsymbol{\Theta}_{s} = \tan^{-1}(\boldsymbol{m}_{X}\boldsymbol{v}_{X_{2}}/\boldsymbol{m}_{M}\boldsymbol{v}_{M}), \qquad (1)$$

where  $m_M$  and  $m_\chi$  are the masses of the alkali and halogen atoms, respectively, and  $v_M$  and  $v_{\chi_2}$  are the incident laboratory frame velocities of the alkali atom and the halogen molecule, respectively.<sup>5</sup> Furthermore, since the final velocity of the spectator halogen atom must be the same as the initial velocity of the halogen molecule in both the laboratory (v) and the center-of-mass (u) frames ( $v_{\chi}=v_{\chi_2}$ , and  $u_{\chi}=u_{\chi_2}$ ), the product recoil energy ( $E_{recoil}$ ) is related to the collision energy ( $E_{collision}$ ), and is simply derived to be:

$$E_{recoil} = \frac{m_{M}}{2m_{MX}} E_{collision}$$
(2)

in the absence of reactant internal energy.<sup>11</sup> If the reactant halogen molecules' internal energy ( $E_{internal}$ ) is included, an additional term is added, and the recoil energy becomes:

$$E_{\text{recoil}} = \frac{m_{M}}{2m_{MX}} E_{\text{collision}} + \frac{m_{M} + m_{X_{2}}}{2m_{MX}} E_{\text{internal}}$$
(3)

and the product MX is scattered to laboratory angles smaller than those given by (1).

Clearly, in describing nearly collinear, small impact parameter collisions the spectator stripping model is physically unrealistic, because due to geometric constraints alone, the alkali halide molecule Fig. 1.

Newton velocity vector diagram for  $M + X_2 \Rightarrow MX + X$  showing the spectator stripping angle  $\Theta_s$ , and the product scattering velocity vectors in the spectator stripping limit. The masses are taken to be those of Na (for M) and  $Cl_2$  (for  $X_2$ ). This figure is adapted from figure 1 of Botscwina, Meyer, Hertel, and Reiland, <u>J. Chem. Phys. 75</u>, 5438 (1981) with permission.


Fig. 1

XBL 863-898

must be backwards scattered. However, the electron transfer cross sections are expected to be so large that the fraction of collisions leading to backwards scattering in this way is small.

As shown in several experiments, the "spectator" aspect of these reactions must be regarded only as a limit, and the more general, phenomenological term "stripping mechanism" is sometimes adopted for a reaction in which scattering is predominantly forwards, but cannot be quantitatively described as in (1-3) above.<sup>19,20</sup>

As discussed in the previous chapter, it was for these reactions, as well as the reactions of alkali atoms with alkyl halide molecules, that the DIPR (Direct Interaction with Product Repulsion) model was developed.<sup>20,21</sup> The difference in the treatment of these two reactive systems is that for halogen molecules charge is allowed to migrate from one atom to the other, whereas for alkyl halides this is obviously not possible. Also, the repulsions for the negative ions of the diatomic halogen molecules are much lower than those for the (dissociative) negative ions of the alkyl halide molecules. Birely, Herschbach, and coworkers have pointed out that the electric field of the alkali ion is strong enough to dissociate the diatomic halogen ion in less than one vibrational period. $^{6,7}$  The electric field at 10 Å separation is approximately  $10^7$  V/cm,  $^{7,22}$  and diatomic ions such as vibrationally excited  $H_2^+$  have been shown to be dissociated by externally applied electric fields of  $10^6$  V/cm.<sup>6,23</sup> Additionally, if a vertical transition is assumed, the  $Cl_2^-$  is formed on the repulsive wall, above its dissociation limit as shown in the Cl<sub>2</sub> and

 $Cl_2^-$  potential energy curves in figure 2, adapted from reference 24 with permission.

The charge migration and lower repulsion lead to some possible problems in the use of the DIPR model. Up to 1/3 of the trajectories. for some mass combinations and neutral-ionic crossing points have "secondary collisions" in which the alkali atom interacts with the departing halogen atom.<sup>21</sup> As pointed out in reference 21, this clearly violates two of the conditions of the model -- first, that there be no interaction between the alkali atom and the departing halogen atom, and secondly, that the repulsive potential for the departing halogen atom be monotonically decreasing. In the calculations of Kuntz, Mok, and Polanyi (KMP), the effect that these secondary collisions have on the product distributions is that there is a large increase in forward scattering.<sup>21</sup> In fact, KMP dispute the conclusion that spectator stripping is the source of the forward scattering. Their calculations suggest that the predominance of forward scattering is due only to secondary encounters. However, the DIPR model reduces to spectator stripping in the limit when the collision energy greatly exceeds the strength of the halogen-halogen interatomic interaction.<sup>21</sup>

In the reactions of excited alkali atoms, where due to the decreased effective alkali ionization potential the crossing point of the neutral and ionic curves has moved out to larger distances, these problems should be less important because the halogen atom has more time to depart, and there should thus be fewer secondary encounters.





A naive view would predict very similar recoil distributions of product since the  $X_2^-$  repulsion remains the same. The only differences would be in the larger cross sections and product internal energies with increasing electronic energy. The product internal energies would be expected to increase simply because the coulombic  $M^+-X^-$  attraction begins at larger distance, giving more vibrational and rotational energy to the MX molecule (much like the reaction of Na(4D,5S) + HCl<sup>25</sup>). According to the conclusions of KMP, however, electronic excitation should lead to less forward scattering and more backwards scattering, because of the smaller number of secondary encounters.

Anderson has discussed the effect of increasing the distance at which the neutral-ionic curves cross,  $r_c$ .<sup>27</sup> The interaction of the ground state neutral and ionic curves is strongest for a collinear  $C_{ov}$  (M-X-X) geometry in which both the neutral M-X<sub>2</sub> (covalent) and the ionic  $M^+-X_2^-$  have  ${}^2\Sigma^+$  symmetry. The interaction has the effect of reducing the energy of the covalent potential energy. In the  $C_{2v}$  geometry (an isoceles triangle with the alkali atom at the apex), the covalent M-X<sub>2</sub> has  ${}^2A_1$  symmetry, but the ionic  $M^+-X_2^-$  has  ${}^2B_1$  symmetry. Thus in  $C_{2v}$  geometry the interaction between the neutral and ionic curves is small`(zero in the Born-Oppenheimer approximation), and the cross section for electron transfer remains small. In the intermediate case,  $C_s$  symmetry, both the neutral and ionic symmetries are  ${}^2A^+$ , and there is an interaction which will smoothly vary as a function of the M-X<sub>2</sub> orientation between the  $C_{ov}$  and  $C_{2v}$  limits. The anisotropic potential valid outside the crossing radius (r>r<sub>c</sub>) can be approximated:

$$V(r, \beta) = -V_{disp}(r) - Ae^{-\alpha r} \cos^2 \beta$$

where r is the  $M-X_2$  separation,  $\beta$  is the angle between the  $X_2$  molecular axis and the  $M-X_2$  axis (the line connecting the M atom with the  $X_2$  center-of-mass), and A and  $\alpha$  are parameters derivable from the properties of the separated atom and molecule.<sup>28,29</sup> The more collinear the configuration, the more attractive it is, and the more probable electron transfer is. Low energy collisions will allow some reorientation into these configurations -- this is termed "tracking"<sup>27</sup> -- that are more favorable for electron transfer. The higher energy the collision, the smaller the effects of tracking, and the more likely are noncollinear collisions to proceed diabatically, without electron transfer or reaction. If the crossing radius is increased, the interaction of the covalent and ionic curves is decreased, and not only is electron transfer at large r less likely for this reason, but also because of the reduced anisotropy. Excessive rotation of the halogen (at the flywheel limit) will reduce the effects of this anisotropy by averaging over it, however in a supersonic expansion the molecular rotation is efficiently cooled.

Another consequence of the mixing of the potential curves is that the harpoon can be thought of as slow. This helps to explain why the harpoon cross sections calculated from the vertical electron affinities of the halogens are often too low. $^{30,31}$  The longer interaction time allows some pre-stretching of the halogen bond, thus lowering the effective electron affinity. This has been shown by the measurements of Los and coworkers which show a decrease in electron transfer cross

245

(4)

sections with increasing collision energies (up to 120 eV) in alkalihalogen collisions. $^{32,33}$  This implies that the electron transfer cross section goes down as the collision is shortened so as not to allow this pre-stretching.

In what is called the orbiting model, the maximum distance at which the reaction can occur,  $r_0$ , is larger than  $r_c$  because it is determined only by the position of the centrifugal barrier on the potential given in (4) above, any collision energies above the centrifugal barrier are captured in the attractive potential, lead to electron transfer, and subsequently to reaction. The centrifugal energy is given by:

$$E_{\text{centrifugal}} = + \frac{b^2}{r_0^2} E_{\text{collision}}$$
(5)

where b is the impact parameter of the collision and  $r_0$  is the maximum in the potential, the sum of (4) and (5).<sup>27</sup> At low collision energies this leads to an increase in the cross section above that predicted by the harpoon mechanism. The orbiting model predicts a decreasing reactive cross section with increasing collision energy since the centrifugal barrier increases. It makes no predictions, however, as to the shape of the reactive product angular distributions.

The symmetry considerations discussed above apply to the scattering of the ground state. This discussion can be extended to the excited state, and the relevant symmetries for ground and excited state scattering are given in table I. Figures 3a-c show the three possible















XBL 862-630



Separated	C <sub>wV</sub>	C <sub>2v</sub>	C <sub>s</sub>	Figure
$M(n^2S) + X_2(X^1\Sigma_g^+)$ .	2 <sub>Σ</sub> +	<sup>2</sup> A1	2 <sub>A</sub> ,	
$M(n^{2}P) + X_{2}(X^{1}\Sigma_{q}^{+})$	2 <sub>Σ</sub> +	<sup>2</sup> A <sub>1</sub>	2 <sub>A</sub> ,	3a
$M(n^2P) + X_2(X^1\Sigma_q^+)$	2 <sub>π</sub> +	<sup>2</sup> 82	2 <sub>A</sub> ,	3b
$M(n^{2}P) + X_{2}(X^{1}\Sigma_{g}^{+})$	2 <b>"</b> –	2 <sub>81</sub>	2 <sub>A</sub> .,	3c
$M^{+}(^{1}S) + X_{2}^{-}(X^{2}\Sigma_{u}^{+})$	$2\Sigma^+$	° <sub>81</sub>	2 <sub>A</sub> ,	

alignments of the alkali np orbital with respect to the  $M-X_2$  axis for  $C_{\infty V}$  and  $C_{2V}$  geometries. It can be seen from table I that with these three alignments of the np orbital, different interactions take place. For the p orbital aligned along the  $M-X_2$  axis, as in figure 3a, the symmetries and interactions will be the same as for the ground state alkali atom (s orbital). For the p orbital in the  $MX_2$  plane, but perpendicular to the  $M-X_2$  axis, as in figure 3b, the interaction is weak for the  $C_s$  symmetry, and essentially zero for both the  $C_{\infty V}$  and  $C_{2V}$  geometries. Finally, for the p orbital perpendicular to the  $MX_2$  plane, as in figure 3c, the situation is the reverse of the ground state case. The symmetries of the covalent and ionic curves match for  $C_{2V}$  geometry, and mismatch for  $C_{\infty V}$  geometry. Therefore,

the maximum attraction and electron transfer probability will occur for  $C_{2v}$  geometry, and at low collision energies the M-X<sub>2</sub> system will be drawn toward this configuration. Because of the nature of the reaction mechanism, the reaction geometry would likely have little effect on the product scattering distributions, and it would be difficult to verify these predictions experimentally.

Several previous experimental measurements of ground state scattering have been performed with well defined collision energies -using velocity-selected effusive<sup>4-6,10,13</sup> or sputtered,<sup>13</sup> or supersonic<sup>11,12</sup> alkali atomic beams. Except for the K sputtering source of van der Meulen et al.,<sup>13</sup> this has given collision energies of up to only 5.5 kcal/mole. With relatively intense seeded supersonic beams, and a judicious choice of collision partners, it is possible to reach collision energies over 200 kcal/mole, and to vary the collision energy.

It has been shown that at low collision energies (0.5-1.3 kcal/mole) the center-of-mass angular distributions are essentially identical for any alkali atom with a particular halogen molecule.<sup>6,7</sup> Grice and Empedòcles found a single center-of-mass angular scattering distribution that adequately describes the product scattering for the reactions:

$$M + Cl_2 \Rightarrow MCl + Cl \tag{6}$$

at collision energies under 1 kcal/mole, where M = K, Rb, and Cs.

The first study of the effect of collision energy on these reactions appears to be that of Gillen, Rulis, and Bernstein who used a velocity-selected effusive K atom source.<sup>10</sup> They studied the reaction K + I<sub>2</sub> at collision energies ranging from 1.9 to 3.6 kcal/mole and saw relatively small changes in reactive cross sections, and center-of-mass angular and recoil energy distributions. Work on this and other systems was continued by Grice and coworkers using a neat supersonic K atom source.<sup>11,12</sup> For K + I<sub>2</sub> at a collision energy of 5.5 kcal/mole they saw an increase in forward scattering as well as a decrease in the reactive cross section when the collision energy was increased from the thermal energies of reference 10.<sup>12</sup> Similar behavior was observed for the reaction K + Br<sub>2</sub> at a collision energy of ~5 kcal/mole<sup>11</sup> when compared to thermal collision energy data.<sup>7</sup>

Van der Meulen, Rulis, and de Vries have studied the reaction K + Br<sub>2</sub> as a function of translational energy for collision energies up to 58 kcal/mole using velocity-selected effusive and sputtered K atomic beam sources.<sup>13</sup> They find that the reactive cross section decreases with increasing collision energy. The reaction cross section drops to 70% of its thermal energy value (measured for 1.3 kcal/mole<sup>26</sup>) at a collision energy of 19.4 kcal/mole, to 50% at 39 kcal/mole, and to a value too small for the authors to measure (<10 Å<sup>2</sup>) at 58 kcal/mole. These values correspond well to what is expected from the orbiting model of Grice, Anderson, Herschbach, and coworkers,<sup>27,28</sup> except at 58 kcal/mole, and not at all well to what is expected for spectator stripping. They see the product recoil energy increase from 27% of the total available energy at 1.3 kcal/mole, to 81% at 19.4 kcal/mole, to 95% at 39 kcal/mole. These values are somewhat suspect as they do not correspond to other experiments on the alkali-halogen systems at low collision energy in which the recoil energies are typically less than 10% of the total available energies of 40 kcal/mole or more. It is possible that this discrepancy is due to the difficulty of accurately extracting the reactive product angular distributions from the sharply skewed Newton diagrams that result from the very fast K atom beam. Nevertheless, the authors show that this behavior correlates better with a simple harmonic model of Kuntz, Polanyi, and coworkers<sup>20</sup> than with spectator stripping. The center-of-mass angular distributions change somewhat, but show no smooth variation with collision energy. All the center-of-mass angular distributions have full widths at half maxima between 13° and 18°, and all except the highest collision energy for which product KBr was detected (39 kcal/mole) peak at 0°.

Although the reaction dynamics of excited alkali atoms with halogens have not previously been studied in crossed molecular beams, the quenching cross sections of such collision partners have been measured.<sup>34</sup> These measurements have been made by dissociating alkali halide vapors in a gas cell using ultraviolet light, and then measuring the decay of the fluorescing excited alkali atomic state produced. The alkali state and velocity can be varied by changing the wavelength of the UV light source. The quenching of Li(2P), Na(3P,4P), and K(4P,5P) by I<sub>2</sub>, and of Na(3P) by Br<sub>2</sub> have been studied in this way.<sup>34-38</sup> For both Na and K the total quenching cross sections, which include all

processes that remove the excited state alkali atom (i.e. reaction and deexcitation), for the first excited state are larger than the ground state reactive cross sections, but then decrease on further excitation to the next nP state.  $^{36,37}$  This is most dramatic for K(4P,5P) in which the quenching cross section decreases by a factor of 5 for K(5P)as compared to K(4P).<sup>37</sup> This decrease was attributed to the neutral  $M \star - X_2$  potential energy curve being at all (M-X<sub>2</sub>) distances higher than the ionic curve. This is not strictly true, because at various halogen interatomic distances the vertical electron affinity can be higher or lower than the equilibrium vertical electron affinity, as discussed above. The first excited states of Na and K have cross sections substantially larger than (approximately twice) the ground state reactive cross sections. All the first excited state quenching cross sections measured decrease with increasing collision energy, typically 30-40% from thermal energies to 10-15 kcal/mole. Once again this can be attributed to the surmounting of a centrifugal barrier as described by Anderson.<sup>27</sup>

Mariella has studied the chemiionization of excited Li by  $Cl_2$ and  $F_2$ .<sup>39</sup> He has measured the cross sections of the processes:

$$Li(2P) + Cl_{2} \rightarrow LiCl^{+} + Cl^{-}$$
 (7)

$$Li(2P) + F_2 \rightarrow Li^+ + F + F^-$$
 (8)

The cross section of process (7) was measured to be  $\sigma=0.2$  Å<sup>2</sup>. The cross section of process (8) was measured to be  $\sigma=2$  Å<sup>2</sup>. No other ionic channels were detected in either case.

Since it has previously been possible to change only reactant combinations, and thus change mass ratios and ionization potential and/or electron affinities simultaneously, no direct experimental determination has been made to show that one or another set of factors is responsible for the variations seen or predicted for these reactions. Experiments on the reactions of excited state atoms should be able to isolate the effects due to the position of the neutral-ionic potential energy curve crossing as the excitation process to first order changes only the ionization potential of the alkali atom. Such experiments should also determine whether the large quenching cross sections observed for halogens with excited alkali atoms are due to reaction or deexcitation processes.

## B. Results

Angular distributions of the NaCl product for the exothermic reaction

Na + Cl<sub>2</sub>  $\Rightarrow$  NaCl + Cl  $\Delta H_0^0 = -40.4 \text{ kcal/mole}^{40}$  (7) were recorded as described in chapter I for three collision energies: 3, 6, and 19 kcal/mole. The Newton velocity vector diagrams for these collision energies are shown in figures 4-6. The measured angular distributions for the reactions of the Na(3S, 3P) states are shown in figures 7-9. All the angular distributions are very narrow compared to the possible product NaCl angular ranges shown in the Newton diagrams as has been observed in all previous alkali-halogen reactions. The shapes of the distributions change very little in going to higher electronic energy, although as will be shown below, the peaks of the angular distributions move to lower laboratory angles for the scattering due to Na(3P) versus that due to Na(3S). All these angular distributions were recorded with the quadrupole mass spectrometer set to a mass to charge ratio (m/q) of 23  $(Na^+)$  as it is expected that the vibrationally excited NaCl produced will fragment under electron bombardment predominantly to Na<sup>+</sup>.<sup>25</sup> The rising signal at low laboratory angles is due to elastically scattered Na which also is detected at m/q=23. Thus the effect of varying ionization efficiencies will be smallest in the data recorded at m/q=23.<sup>25,42</sup> An angular distribution was recorded for Na(3S,3P) +  $Cl_2$  at m/q=58 (NaCl<sup>+</sup>) at



XBL 863-845

Fig. 4. Newton velocity vector diagram for  $Na(3S, 3P) + Cl_2 \Rightarrow \underline{MaCl} + Cl$  at a collision energy of 3 kcal/mole. The sodium beam is seeded in argon, and the chlorine beam is neat.



XBL 863-847

Fig. 5. Newton velocity vector diagram for  $Na(3S, 3P) + Cl_2 \Rightarrow NaCl + Cl$  at a collision energy of 6 kcal/mole. The sodium beam is seeded in neon, and the chlorine beam is neat.

E<sub>coll</sub>=19 kcal/mole



XBL 863-846

Fig. 6. Newton velocity vector diagram for  $Na(3S, 3P) + Cl_2 \Rightarrow NaCl + Cl at a collision energy of 19 kcal/mole. The sodium beam is seeded in helium, and the chlorine beam is neat.$ 



XBL 863-866

Fig. 7. NaCl product angular distributions for  $Na(3S,3P) + Cl_2$  at a collision energy of 3 kcal/mole. The sodium beam is seeded in argon, and the chlorine beam is neat. These distributions were recorded at m/q=23 (Na<sup>+</sup>).









NaCl product angular distributions for Na(3S,3P) +  $Cl_2$  at a collision energy of 19 kcal/mole. The sodium beam is seeded in helium, and the chlorine beam is neat. These distributions were recorded at m/q=23 (Na<sup>+</sup>).

a collision energy of 6 kcal/mole, and is shown in figure 10. Within the rather large error bars of this measurement, the distribution is superimposable on the distribution recorded at m/q=23 (Na<sup>+</sup>).

Due to the small difference in the angular distributions between ground and excited state, no time-of-flight product velocity measurements, nor any polarization dependences were measured.

C. Analysis of Experimental Results

As in the previous chapters, the laboratory data was fitted using the program GM which calculates a laboratory scattering distribution from independent assumed center-of-mass angular and recoil energy distributions. Although no laboratory velocity data was recorded, the predominance of forward scattering and the narrow recoil energy distributions give nearly single-valued laboratory to center-of-mass frame transformations such as those which allowed the early studies of these reactions to derive center-of-mass distributions from their relatively primitive laboratory angular distributions.<sup>4-9</sup>

The single center-of-mass angular distribution derived by Grice and Empedocles for K,Rb,Cs +  $Cl_2$  at collision energies under 1 kcal/mole was unable to fit the measured angular distributions for Na +  $Cl_2$ .

The best fits to the laboratory angular distributions for  $Na(3S,3P) + Cl_2$  at a collision energy of 6 kcal/mole are shown in figure 11. The center-of-mass angular distributions (T( $\Theta$ )) derived



XBL 863-872

Fig. 10. NaCl product angular distributions recorded at m/q=58
(NaCl<sup>+</sup>) for Na(3S,3P) + Cl<sub>2</sub> at a collision energy of 19
kcal/mole. The sodium beam is seeded in helium, and the
chlorine beam is neat.



XBL 862-559

Fig. 11. Best fits to the laboratory angular distributions for Na(3S,3P) +  $Cl_2$  at a collision energy of 6 kcal/mole. The solid lines are the fit to the data.

for these fits are shown in figure 12. Significantly different recoil energy distributions (P(E)) were required to fit each, and these are shown in figures 13a and b for  $Na(3S,3P) + Cl_2$ , respectively. With increasing electronic energy the peak recoil energy moves out to higher energy. The peak recoil energies are 0.6 and 1.2 kcal/mole for  $Na(3S,3P) + Cl_2$ , respectively, at this collision energy. The total cross section for the reaction of the Na(3P) atoms is 58% higher than that for Na(3S) at a collision energy of 6 kcal/mole.

The best fits to the laboratory angular distributions for  $Na(3S,3P) + Cl_2$  at a collision energy of 19 kcal/mole are shown in figure 14. The center-of-mass angular distributions are shown in figure 15. The recoil energy distributions are shown in figures 16a and b for  $Na(3S,3P) + Cl_2$ , respectively. With increasing electronic energy, the peak recoil energy once again moves out to higher energy, from 0.8 and 1.0 kcal/mole for  $Na(3S,3P) + Cl_2$ , respectively. The relative total reaction cross section for Na(3P) is 16% higher than for Na(3S) at this collision energy.

A summary of some of the features of the four fits is given in table II. It can be seen that the center-of-mass angular distributions are somewhat wider for the reaction of Na(3P) than those for Na(3S). Also, as mentioned above, the peak recoil energies increase in going from Na(3S) to Na(3P). It should be noted that the column labeled "Percent Forward Scattered" gives the fraction of the weighted intensity,  $T(\Theta)$ sino, that is in the forward hemisphere.



XBL 862-662









XBL 862-667





XBL 862-668

Fig. 15. Center-of-mass angular distributions used for the fits in figure 14 of  $Na(3S,3P) + Cl_2 \Rightarrow NaCl + Cl$  at a collision energy of 19 kcal/mole.





<u>Na Level</u>	Ecoll	FWHM of T(e)	% Forward Scattered	Peak Recoil Energy	$\sigma(3P)/\sigma(3S)$
3S	6.	76°	83	0.6	1.58
3P	6.	97°	76	1.2	
3S	19.	64°	85	0.8	1.16
3P	19.	68°	82	1.0	

Table	II.	A summary	of features of the center-of-mass distributions
		fitted to	the laboratory angular distributions for
		Na(3S,3P)	+ Cl <sub>2</sub> . All energies are in kcal/mole.

As in chapter II, the total reactive cross section can be determined by calibrating the reactive scattering to the small angle elastic scattering. This is done using the small angle approximation to the classical scattering from a van der Waals r<sup>-6</sup> potential,<sup>6,43</sup> equation (8) of chapter II. As in chapter II, the van der Waals  $C_6$  constant can be estimated using the Slater-Kirkwood formula, 44,45 equation (10) of chapter II. The polarizability of Cl<sub>2</sub> is  $\alpha$ =4.61 Å<sup>3</sup>,<sup>46</sup> and its number of valence electrons is  $N(C1_2)=14$ . This yields a  $C_6$ constant for ground state Na(3S) of 7410 kcal/mole  $Å^6$ , using the polarizability of ground state Na  $\alpha$ (Na)=24.5 Å<sup>3</sup> as given in chapter II.<sup>47</sup> This agrees well with the theoretical value derived by Anderson and Herschbach for Na-Cl<sub>2</sub>,  $C_6=7460$  kcal/mole  $A^{6}$ .<sup>48</sup> The  $C_6$  constant determined for elastic scattering of the  $Na(3^2P_{3/2})$  atoms by  $Cl_2$  is 11340 kcal/mole  $Å^6$  as derived from the polarizability of the excited atom  $\alpha(Na(3^{2}P_{3/2}))=53.6 \text{ Å}^{3}.49,50$  The ionization cross section of Na by 200 eV electrons is 2.1  $Å^2$ .<sup>51</sup> It was assumed that all NaCl

Sodium Level	Using 3S Scaling Factor	Using 3P Scaling Factor	Average of Two Values	
35	79 Å <sup>2</sup>	64 Å <sup>2</sup>	72 Å <sup>2</sup>	
3P -	124 Å <sup>2</sup>	101 Å <sup>2</sup>	113 Å <sup>2</sup>	

Table III. Total reactive cross sections for Na(nL) + Cl<sub>2</sub> → NaCl + Cl at a collision energy of 6 kcal/mole.

produced was vibrationally excited, so that the ionization cross section of the NaCl to Na<sup>+</sup> (m/q=23) was  $\sigma^{1}=3.8$  Å<sup>2</sup> as in chapter II.<sup>52</sup> This gives a relative detection efficiency of 1.8 for Eff(NaCl<sup>‡</sup>)/Eff(Na). The low angle elastic scattering is then fit and scaled to the distribution given by formula (8) of chapter II. The scaling factors derived are used to determine the total reactive cross sections given in table III. As described in chapter I, it was assumed that 20% of the Na atoms are excited to the 3P level. Note that it is not possible to use the hard sphere elastic scattering approximation to calculate the reactive cross section as the large angle elastic scattering is depleted by reaction.

## D. Discussion

All the laboratory angular distributions peak at angles higher than the spectator stripping angles. From (1), using the most probable beam velocities at each energy, the spectator stripping angles,  $\theta_{e}$ , are 41°, 31°, and 17° for collision energies of 3, 6, and 19 kcal/mole, respectively. As discussed in section A, if the finite vibrational temperature of Cl<sub>2</sub> is considered, these angles shift to still lower values. Since the neutral-ionic curve crossing moves to larger distances for electronically excited Na, it would be expected that the NaCl product would be scattered more in the forward direction. While the recoil energy distribution does increase, moving the peak in the angular distribution closer to the spectator stripping angle, the center-of-mass angular distributions are actually wider for the reaction of the Na(3P) atoms. This is consistent with the idea that if the transfer takes place at longer range, there is less interaction of the Na and the Cl<sub>2</sub> before the transfer. Then there is a greater energy release due to the dissociation of the  $Cl_2^-$  higher up on the repulsive wall of its potential surface. This is inconsistent with the DIPR model calculation's prediction that without secondary interactions (MX-X) no preference for forward scattering will be observed.

The shift of the center-of-mass angular distributions to more forward scattering with increasing collision energy was observed, corroborating the results of Grice and coworkers.<sup>11,12</sup> The

## Table IV. The calculated crossing radii and electron transfer cross sections for the Na(3S,3P) + Cl<sub>2</sub> interaction. The adiabatic (Adia) electron affinity of Cl<sub>2</sub> is 2.39 eV, $^{41}$ and the vertical (Vert) electron affinity of Cl<sub>2</sub> is taken to be 0.7 eV.<sup>24</sup>

Sodium Level	Ionization Potential	Crossing r <sub>c</sub> Adia	Crossing Radius rc Adia Vert		Estimated Electron Transfer Cross Section Adia Vert	
35	5.138 eV	5.22 Å	3.2 Å	86 Å <sup>2</sup>	32 Å <sup>2</sup>	
3P	3.033 eV	22.3 Å	6.2 Å	1560 Å <sup>2</sup>	121 Å <sup>2</sup>	

center-of-mass angular distributions are all wider than the distribution derived by Grice and Empedocles for K, Rb, Cs +  $Cl_2$ .<sup>8</sup> It is possible that at low collision energies this distribution is also inadequate for Na +  $Cl_2$ , but that was not determined here.

The reaction cross section does increase with electronic excitation at low collision energy, but not as much as the estimates of the crossing of the neutral and ionic potential curves would predict. The crossing radii ( $r_c$ ) and estimated reactive cross sections ( $\sigma_R$ ), calculated as discussed in chapter II, are given in table IV. The adiabatic electron affinity of Cl<sub>2</sub> was taken to be 2.39 eV.<sup>41</sup> The vertical electron affinity of Cl<sub>2</sub> from reference 24 is approximately 0.7 eV. In a study of the cross sections of the reactions of alkali

atoms with halogen molecules, Maya and Davidovitz foundathat the relater molecule effective (reactive) electron affinity of  $Cl_2$  was 1.6 eV, 26 - 1 - 16 electron electron however, as they point out, this is certainly an overestimate because the Cl<sub>2</sub> in their experiments was at 1015°K, allowing the thermal train esternation population of several vibrational levels. Excited vibrational levels will have significantly greater electron affinities than v=0 due to the steep slope of the  $Cl_2$  repulsive wall in this region, as seen in figure 2. This smaller than expected increase in reactivity is due to the terminant the small electron transfer probability at long range  $(r_c>10^{-A})$ . Recall from chapter III that many Cl<sub>2</sub> vibrations occur during a collision at these energies, so that the electron affinity of the  $Cl_2$ is effectively modulated, giving the possibility of electron transfer each time the covalent and ionic potential energies become comparable. At very long range the interaction of the states is so small that electron transfer does not occur. On closer approach, the effective electron affinity has decreased, and the interaction of the states is smaller, so that if electron transfer occurs, it reaches higher on the repulsive wall of the  $Cl_2^-$  potential curve, giving higher recoil energies. These higher recoil energies are observed, but are a relatively small effect. A further explanation for the cross section not rising up to the harpoon limit for the Na(3P) reaction is that other processes involving electron transfer, such as quenching via electron transfer back to the Na atom, could be taking place. More likely is that electron transfer is taking place, that is

 $Na(3P) + Cl_2 \Rightarrow Na^+ + Cl_2^-$ 

274

(9)

which becomes energetically accessible at a collision energy of 15 kcal/mole. Also, as discussed in section A, one third of the atoms have a symmetry for which there is little or no interaction with the ionic state necessary for electron transfer. The effects of tracking may be important in the smaller increase in the cross section at the higher collision energy. Additionally, since the relative velocities are higher, the collision times are shorter, and there are then fewer crossings of the "oscillating" potential curves (see above). There is also less time for the "prestretching" of the Cl<sub>2</sub> bond, which has the effect of increasing the Cl<sub>2</sub> electron affinity, and thus increasing the reactive cross section.

Maya and Davidovitz have performed careful measurements of the cross sections of the reactions of alkali atoms with halogen molecules at thermal energies in a gas cell.<sup>26,30</sup> The value that they derive for Na + Cl<sub>2</sub> at a mean collision energy of 1.6 kcal/mole is  $\sigma_R$ =124 Å<sup>2</sup>, with a reported absolute accuracy of ±15%.<sup>30</sup> By comparison, the value obtained here for the ground state, 72 Å<sup>2</sup>, is qualitatively consistent with the drop-off of cross section observed by van der Meulen et al.<sup>13</sup> The limited accuracy of this measurement precludes quantative comparisons, but at a collision energy of 6 kcal/mole, van der Meulen et al. saw a 20% decrease in the cross section of K + Br<sub>2</sub> from the thermal energy value.

The low collision energy increase in cross section is in line with the quenching cross section expected for Na(3P) from the experiments of Bersohn, Herm, and coworkers described in section A. The small increase
in cross section at higher collision energy implies that some other mechanism, such as charge transfer (9), is responsible for the quenching of the Na(3P) excitation.

The fact that the recoil energy distributions change little upon excitation to Na(3P) implies that only a small fraction of the available energy goes into product translational energy with increasing electronic energy. This is consistent with the naive expectations of the dynamics. That is, after electron transfer, the molecular halogen ion  $(X_2^-)$  immediately begins to dissociate and the alkali ion  $(M^+)$  and the halide ion  $(X^-)$  are drawn together. If the electron transfer occurs at longer range, the  $M^+$  and  $X^-$  ions start out further apart, and thus more energy goes into the vibrational and rotational energy of the MX molecule. The small differences in the distributions could possibly be attributed to more  $Cl_2^-$  repulsion due to electron transfer inside the first crossing of the covalent and ionic curves, and thus higher up on the  $Cl_2^-$  repulsive wall.

# E. Conclusions

The cross section of the Na +  $Cl_2$  reaction does not increase as much as the simple harpoon mechanism transfer predicts upon excitation of the Na atom. A larger increase of the cross section for Na(3P) reaction is observed for the reaction at 6 kcal/mole, than at 19 kcal/mole

collision energy. This relates to the small interaction of the covalent and ionic states at long range.

The product recoil energy distributions change little with reagent Na electronic energy, and nearly all electronic energy must be converted into the internal energy of the product NaCl. The release of the alkali and halide ions at larger separation leads to the deposition of nearly all the electronic energy into the rotational and vibrational energy of the product alkali halide molecule. The small broadening of the center-of-mass angular distributions and increase of the product recoil energy with electronic excitation observed are attributable to slightly more repulsive energy release in the  $Cl_2^-$  dissociation.

Finally, while these reactions studied do not reach the spectator stripping limit. their general behavior is well described by the more general stripping mechanism.

# F. References

- M. Polanyi, <u>Atomic Reactions</u> (Williams and Northgate, London, 1932).
- 2. S. Datz and R. E. Minturn, <u>J. Chem. Phys.</u> 41, 1153 (1964).
- 3. K. R. Wilson, G. H. Kwei, J. A. Norris, R. R. Herm, J. H. Birely, and D. R. Herschbach, <u>ibid.</u>, 1154 (1964).
- 4. A. E. Grosser and R. B. Bernstein, J. Chem. Phys. 43, 1140 (1965).
- R. E. Minturn, S. Datz, and R. L. Becker, <u>J. Chem. Phys. 44</u>, 1149 (1966).
- 6. J. H. Birely and D. R. Herschbach, ibid., 1690 (1966).
- J. H. Birely, R. R. Herm, K. R. Wilson, and D. R. Herschbach, <u>J.</u> <u>Chem. Phys. 47</u>, 993 (1967).
- 8. R. Grice and P. Empedocles, J. Chem. Phys. 48, 5352 (1968).
- 9. D. D. Parrish and R. R. Herm, J. Chem. Phys. 51, 5467 (1969).
- K. T. Gillen, A. M. Rulis, and R. B. Bernstein, <u>J. Chem. Phys. 54</u>, 2831 (1971).
- 11. J. C. Whitehead, D. R. Hardin, and R. Grice, <u>Mol. Phys. 23</u>, 787 (1972).
- 12. S.-M. Lin, D. J. Mascord, and R. Grice, Mol. Phys. 28, 975 (1974).
- A. van der Meulen, A. M. Rulis, and A. E. deVries, <u>Chem. Phys. 7</u>, 1 (1975).
- 14. C. M. Scholeen, L. A. Gundel, and R. R. Herm, <u>J. Chem. Phys. 65</u>, 3223 (1976).

- 15. R. R. Herm, in <u>Alkali Halide Vapors: Structure, Spectra, and</u> <u>Reaction Dynamics</u>, 189 (Academic Press, New York, 1979), and references therein.
- J. L. Kinsey, <u>MTP Int. Rev. Sci., Phys. Chem. Ser. One 9</u>, 173 (1972), and references therein.
- 17. J. L. Magee, J. Chem. Phys. 8, 687 (1940).
- 18. D. R. Herschbach, Appl. Opt. Suppl. 2, 128 (1965).
- 19. J. C. Polanyi, Disc. Faraday Soc. 44, 293 (1967).
- 20. P. J. Kuntz, E. M. Nemeth, and J. C. Polanyi, <u>J. Chem. Phys. 50</u>, 4607 (1969).
- 21. P. J. Kuntz, M. H. Mok, and J. C. Polanyi, ibid., 4623 (1969).
- 22. The strength of the electric field at 7 Å was incorrectly reported as  $3x10^9$  V/cm in reference 6, but was corrected to  $3x10^7$  V/cm in reference 7.
- 23. A. C. Riviere, D. R. Sweetman, Phys. Rev. Lett. 5, 560 (1960).
- 24. T. L. Gilbert and A. C. Wahl, J. Chem. Phys. 55, 5247 (1971).
- 25. See Chapter II of this thesis.
- 26. J. Maya and P. Davidovitz, J. Chem. Phys. 59, 3143 (1973).
- R. W. Anderson, Ph. D. Thesis, Harvard University, Cambridge, Massachusetts (1968).
- R. Grice, Ph. D. Thesis, Harvard University, Cambridge, Massachusetts (1967).
- 29. R. Grice and D. R. Herschbach, Mol. Phys. 27, 159 (1974).
- 30. J. Maya and P. Davidovitz, J. Chem. Phys. 61, 1082 (1974).

- 31. P. Davidovits, in <u>Alkali Halide Vapors:</u> Structure, Spectra, and Reaction Dynamics, 331 (Academic Press, New York, 1979).
- 32. J. A. Aten and J. Los, <u>Chem. Phys. 25</u>, 47 (1977).
- 33. J. Los and A. W. Kleyn, in <u>Alkali Halide Vapors: Structure</u>, <u>Spectra, and Reaction Dynamics</u>, 275 (Academic Press, New York, 1979).
- 34. R. Bersohn, ibid., 345.
- 35. L. E. Brus, J. Chem. Phys. 52, 1716 (1970).
- 36. R. Bersohn and H. Horwitz, J. Chem. Phys. 63, 48 (1975).
- 37. B. L. Earl and R. R. Herm, J. Chem. Phys. 60, 4568 (1974).
- 38. S.-M. Lim and R. E. Weston, J. Chem. Phys. 65, 1443 (1976).
- 39. R. P. Mariella, Jr., J. Chem. Phys. 76, 2965 (1982).
- 40. Calculated from diatomic molecular dissociation constants  $(D_0^0)$  given in reference 41.
- F. P. Huber and G. Herzberg, <u>Molecular Spectra and Molecular</u> <u>Structure, IV. Constants of Diatomic Molecules</u>, Van Nostrand Reinhold, Co., New York, New York (1979).
- 42. M. F. Vernon, H. Schmidt, P. S. Weiss, M. H. Covinsky, and Y. T. Lee, <u>J. Chem. Phys.</u>, to be published, and M. F. Vernon, Ph. D. Thesis, University of California, Berkeley, Berkeley, California (1983).
- 43. E. A. Mason, J. T. Vanderslice, and C. J. G. Raw, <u>J. Chem. Phys.</u>
  40, 2153 (1964).
- 44. J. C. Slater and J. G. Kirkwood, Phys. Rev. 37, 682 (1931).
- 45. K. S. Pitzer, Adv. Chem. Phys. 2, 59 (1959).

- 46. J. O. Hirschfelder, C. F. Curtiss, and R. B. Bird, <u>Molecular</u> <u>Theory of Gases and Liquids</u> (John Wiley and Sons, Inc., New York, 1964).
- 47. H.-J. Werner and W. Meyer, Phys. Rev. A 13, 13 (1976).
- 48. R. W. Anderson and D. R. Herschbach, <u>J. Chem. Phys. 62</u>, 2666 (1975).
- 49. P. Hannaford, W. R. MacGillivray, and M. C. Standage, <u>J. Phys. B:</u> Atom Molec. Phys. 12, 4033 (1979).
- 50. H. T. Duong and J.-L. Picque, J. Physique 33, 513 (1972).
- 51. R. H. McFarland and J. D. Kinney, Phys. Rev. 137, A1058 (1965).
- 52. This is derived by assuming 90% fragmentation of NaCl to Na<sup>+</sup> in the electron bombardment ionizer, and by estimating the ionization cross section of NaCl from that of Ar using their relative polarizabilities.

#### VI. CONCLUSIONS

# A. Conclusions on the Reaction Dynamics of Electronically Excited Alkali Atoms

Representatives of four families of excited alkali atom reactions have been studied in crossed molecular beams. In the cases where electrons can effectively dissociate the molecular reaction partner, as inalkali atoms with methyl halide and halogen molecules, the simple ideas of the product repulsion qualitatively describe the reactions well. The very small changes in the reactivity and mechanism in these cases were surprising, but can be understood in terms of two competing effects. The lowered ionization potential of the excited states of Na has the tendency to make the upper states "super-reactive," but the shift of the crossing of the potential curves to larger separation reduces the interaction of the neutral and ionic states, and makes the excited states less reactive than they would otherwise be. This can be seen clearly in the case of Na(3P, 4D, 5S) + HC1. The first excited state does not proceed via long range electron transfer, because the crossing point of the neutral and ionic curves is small enough that the NaCl molecule is formed before the HCl<sup>-</sup> repulsion can play a role, and it is the NaCl-H repulsion that dominates the product scattering. When the ionization potential of the Na atoms is lowered further by subsequent excitation to Na(4D,5S), the reaction becomes more like the

electron transfer reactions, and the HCl<sup>-</sup> repulsion dominates the scattering. While these states are not "super-reactive," they are more reactive the ground and first excited states.

On the other hand, in Na + 0<sub>2</sub> the importance of orbital symmetry on the reaction dynamics is clearly apparent. Here the neutral chemical pathway is certainly a minor process. Nonetheless, it exhibits very unusual behavior, having a threshold to reaction, very small product recoil energies, and very specific angular scattering. As pointed out in chapter 3, measurements of the other processes taking place should certainly be performed, as well as further measurements of what was observed. Of particular interest would be a spectroscopic study of the products, first to prove their identity, and second to identify where the large amount of internal energy has been deposited. Also of interest would be the measurement of the ionic products, and the transition state symmetry required for those processes.

Future investigations should elucidate more of the important features of the transition states in these excited state reactions that are so complicated by the large number of potential energy surfaces involved.

## VII. APPENDICES

A. Data Acquisition for Angular Distributions, Polarization Dependences, and Machine Condition Optimization: The Data Acquisition System and Program SANG

1. Introduction to the Data Acquisition System

The data acquisition system is comprised of four types of electronics: NIM standard modules,<sup>1</sup> CAMAC standard modules,<sup>2</sup> and Q-Bus standard boards,<sup>3</sup> as well as some stand alone devices. As discussed in chapter I, data signals are received from two or three photomultipliers -- one photomultiplier tube (RCA  $8850^4$ ) that monitors the signal from the mass spectrometer, and one or two photomultipliers (RCA  $1P28^4$ ) on the fluorescence monitors. A further description of exactly what these measure can be found in chapter I. The gating and timing for the various electronics are derived from the 150 Hz tuning fork chopper used to modulate the secondary molecular beam.

A map of the electronics can be found in figure 1. The tuning fork chopper signal is input to a Timer-Gater module (LBL 13X-3050II). This uses a variable phase and a variable gate width to output two gate pulses "A" and "B," which correspond to the secondary beam on and off, respectively. The Timer-Gater module can be started and reset using a parallel line unit (DEC DRV- $11^5$ ) on the Q-Bus of the controlling minicomputer (DEC LSI- $11/23^5$ ). The gate width is set by matching



Fig. 1. Schematic of the data acquisition electronics.

XBL 861-10008

the width of the A and B pulses to the minimum of the measured beam on and off times. This measurement is made with the mass spectrometer and the time-of-flight program described in appendix B. The phase is coarsely set so as to maximize the mass spectrometer signal to noise ratio for beam on vs. beam off when looking directly into the secondary beam. The phase is later fine tuned on the reactive signal. Both of these optimizations are done using Test mode of program SANG, described below.

The A and B outputs go to a CAMAC powered Laser Gate Generator (Lawrence Berkeley Laboratory), built by Matt Vernon, Jim O'Brien, and Fred Vogelsberg. The A output is also sent to a CAMAC powered Motor Sync Scaler (LBL 13X3441-P1-A), which just divides down the A pulses by 25 giving pulses at 6 Hz. The Motor Sync Scaler on tput pulse goes to the Laser Gate Generator which has outputs to a stepper motor controller which changes the state of the laser beam flag, as described in chapter I. The Laser Gate Generator also outputs six gate levels corresponding to the detector signals described in chapter I. These go to two CAMAC Quad Scalers (Joerger S1-Ind<sup>6</sup>) which are read by the computer via the CAMAC dataway.

2. Introduction to Program SANG

Program SANG is designed to be used for machine condition optimization, and for recording angular distributions and polarization dependences. Because of this, there are several modes of data acquisition in which program SANG operates. These are designed to be

flexible enough to allow for any circumstances that could take place during an experiment, without the necessity of exiting and restarting the program. Along these lines, comments and machine conditions can be written directly to the data file and hard copy printer output, laboratory notebook style, data files can be closed, and new files begun. One function that does not exist is the deletion of data already recorded. Subsequent comments should be sufficient to mark bad points for later editing. The program has the facility (although not included in the listing below) of measuring the Doppler profile of the Na( $D_2$ ) transition at 45° and 90° laser crossings (as described in chapter I) automatically, and thus directly recording the velocity and speed ratio of the sodium beam.

The program is designed to be nearly independent of the hardware described in section 1 so that only minor changes would be needed to the MACRO subroutines in the file SANGLE.MAC to accommodate hardware changes. Everything else such as the number and definitions of data channels recorded can be modified while running SANG (using command DC). All of the gating of the data channels is done in hardware and SANG in no way depends upon this.

Commands are read from the console using "special screen mode" which is accessed by setting bit 12 in the job status word (see MACRO subroutines ENSPEC and EXSPEC).<sup>7</sup> All commands are one or two letters. There are help screens describing each set of commands. As the various functions of the program are described below, the commands relating to those functions are listed. Refer to the program listing

or the help screens for more complete discussions of the commands described.

#### Data Acquisition Modes

The various modes of data acquisition are described below. Of course there is a great deal of overlap between each of the modes. The only real differences between modes are in the outermost shells of control in which the data acquired at different points are linked together, and the points at which data are recorded are chosen according to the loop specific to that mode. A flag word (IFLAG in subroutines SDONE and SWRITE) is written to the data file on disk to specify under which mode each data point was recorded. The definition of this word is given in subroutine SDONE.

## Single Point Data Acquisition:

The simplest mode of data acquisition is the recording of a single point, with the user deciding where to proceed after each point. This is available through the commands CA, CI, and RD, which choose the angle to be measured, and the countdown interval, and start the countdown, respectively.

#### Normalization Mode:

After a certain number of points are recorded (often 10), data is recorded at a certain angle multiple times (often two or four) in order to allow for a correction for long term drifts in machine conditions. This is usually done in single point mode above, as well as loop mode, described below. The angle and interval at which data are recorded for normalization purposes are chosen with the commands NA and NI, respectively, as well as with NP and CN. The command NR determines how many times the data is acquired at the normalization point. After the requisite number of Timer-Gater countdowns, the user is queried as to whether the whole series of normalization countdowns should be repeated. This is especially useful where some optimization is done; a second set of normalization points provides a new starting point for normalizing data, and a direct comparison of data from before and after optimization.

# Loop Mode:

It is often useful to record data over an angular range, with a certain angular interval. The mass spectrometer is moved back and forth across the range, and the data recorded is normalized to a single value (see normalization mode above). All of the angular distributions reported in chapters II-V were recorded in this way using loop mode. The loop mode parameters can be set with the commands LA, LB, LD, LE, LI, LN, LP and CL. The loop can be started and continued with the commands LS and LC, respectively. Note that normalization mode can be called from within loop mode, and is automatically called at the beginning of a loop. The parameters can be changed during the loop to expand or contract the loop, etc.

Rotate Polarization Mode:

When lasers are used to optically prepare an excited state for reaction, as shown in chapters II and III it can be interesting to measure the polarization dependence of the various signals recorded at fixed detector angles. Rotate polarization mode can be used to measure the differences produced by left and right circularly polarized light or the dependences as the electric vector of linearly polarized light is rotated. The command RA changes the laboratory detector angle at which the polarization scan is recorded while the command RM changes the step size for the polarization rotator; the command RP changes both. The command RS starts a polarization scan. The data is stored in such a way that the detector and laser polarization angles can be easily determined. The angle written to disk is:

Angle = 
$$\Theta_{LAB}$$
 +  $\frac{\phi_{POLAR}}{10000}$ , (1)

where  $\Theta_{LAB}$  is the laboratory detector angle and  $\phi_{POLAR}$  is the laboratory polarization angle as described above. Thus if the detector is at 45° and the laser polarization angle is 90° (the electric vector of the laser/lasers is 90° from vertical in the laboratory), the angle written to disk would be 45.0090°. Left and right circularly polarized light have values of 0°, 90°, and 180°. The polarization is rotated by a stepping motor (Superior Electric Co. Slo-Syn Synchronous Stepping Motor M062-FC09E<sup>8</sup>) that rotates a double fresnel rhomb polarization rotator (Spectra-Physics 310A<sup>9</sup>) as described in chapter I. The stepping motor is sent pulses by a CAMAC Pulse Generator (Kinetic Systems

3360<sup>10</sup>). All of the polarization dependences reported in chapters II and III were recorded with rotate polarization mode. Like loop mode, the polarization is rotated back and forth from 0° to 180°. The motor is only driven in one direction however and is accurate to better than 1° after over 100 full rotations. A separate display mode is available for viewing polarization data, and recording the data (using the small angular offset related to the laser polarization) as discussed above facilitates setting up this special display window.

## Test Mode:

Machine optimization is accomplished by repeating the same measurement many times and optimizing the signal and/or the signal to noise ratio of a channel. To this end, test mode also has a special display window. In the same manner as the rotate polarization mode the data is recorded with a slight offset added to the detector angle. The data recorded in test mode is labeled sequentially and the angle recorded is then the detector angle plus an offset related to the number of the point recorded. This gives the test mode data point's angle as:

Angle = 
$$\Theta_{LAB}$$
 + Test Number, (2)

in the data file. Typically various machine parameters are varied while the reactive signal is monitored and optimized at a particular detector angle at which signal is expected or known to occur. The

display is used to histogram the effect of the variation of each machine parameter. When the optimization is complete, the test mode data is zeroed (this can only be done in the accumulated data arrays, not in the data already written to disk) and/or a new data file is opened and the full data acquisition is begun. The commands TA, TI, and TN change the detector angle, the countdown interval, and the number of measurements for test mode, respectively. The commands TS and TC start and continue test mode data acquisition, respectively, while the command TE terminates test mode after the current countdown is completed.

#### Other Functions of Program SANG

The machine conditions -- detector settings, pressures, temperatures, etc. -- are recorded by the user and can be output to various devices. The conditions recorded are meant to serve as a laboratory notebook, which is otherwise difficult information to integrate with data files stored on disk. The commands: MA, MB, MD, MF, MI, MM, MR, MS, and MT change these conditions. These conditions are purely cosmetic and have no bearing on the data or its acquisition. Whenever information is written to disk it is also written to the printer to serve as an easily available record of the experiment. The definitions and number of thermocouples and pressure regions can be changed with the commands DP and DT. The readings held by the computer can be written to console, the printer, or the disk data file with the commands: PM, PS, SM, SS, WM, and WS. The primary data acquisition parameters such as the countdown interval and the gate widths can be changed with the commands: CA, CC, CE, CG, CI, CN, and CP. The current data acquisition configuration including all the parameters of the data acquisition mode in use can be written to the console, the printer, and the disk data file with the commands SC, PC, and WC, respectively. In addition, the parameters of the data acquisition modes not in use can be written to these three devices with similar commands. The parameters for the various data acquisition modes can be changed with the commands discussed briefly above. Laboratory notebook style comments can be written to the disk data file and the printer or just to the printer with the commands WI and PI, respectively.

All of the settings for data acquisition, the definitions of pressure regions and thermocouples, as well as the actual pressures, temperatures, and detector settings have default values when SANG is begun. These are set in subroutine SINIT and are meant to approximately reproduce typical running conditions.

The size of program SANG far exceeds the memory limits of the LSI-11 computers. As such, many subroutines are "overlaid" -- moved in and out of memory as needed. This must be done explicitly in RT-11, and the LINKer command file lists the 65 segments in the three overlay regions. The size is still close enough to the upper limit that two other special steps are required. First, the FORTRAN files are all compiled with the /NOLINENUMBERS option as shown in the compilation

file CNSANG.COM. This has the effect that if a run time error should occur, the subroutine in which, but not the line at which the error occurred is displayed on the console. Secondly, the program is linked with the library file \$SHORT included in the LINKer command file LTUF.COM. This eliminates the text of the FORTRAN error messages, but not the error number, so that the text can be looked up in the DEC manuals.

All data files must be edited to remove the comments (and usually all bad data points) before further processing. The original unedited data file is of course archived for future reference. A number of routines are available on the LBL VAXes for the compilation and normalization of angular distributions and polarization data. Most frequently used is the program ANGBAB, which can do much of the data manipulation automatically.

3. Listing of Program SANG

A listing of program SANG follows. The order of the files is that done by backup/update/printout command file CPSANG.COM, which is also included.

#### PROGRAM SANG

File SANG.FOR begins on the previous line.

Version Date: November 11, 1985 Paul S. Weiss

Program SANG takes angular scan data on up to 16 scalers channels. Machine optimization is done and laser polarization measurements can be recorded at a number of angles. Control is maintained through subroutine SCOMM. All computer controlled hardware is driven through the MACRO routines found in SANGLE.MAC.

SANG must be linked with all the files found in LINKer file LSANG.COM.

INTEGER+4 JDATA COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /FNAME/ INAME(7), ISTORE COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB COMMON /DSPZER/ ITYZER, IBYZER COMMON /DSCR/ DSCR(20) COMMON /SCR/ SCR(20) COMMON /TGSMAX/ DWLMAX, MAXCHN COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /POLSET/ ANGPOL,POLSTP,POSIT, IPOSIT, IPLFLG COMMON /TSTSET/ TSTANG,NUMTST,TSTCUR,TSTFLG COMMON /VELO/ DELF, HALFH, DELPIP, VELB, SPDRIO COMMON /SDONE/ ISWFLG, JDATA(20) COMMON /CHADEF/ CHADEF(20,10) COMMON /PRES/ PRM(8), IPRE(8), PFORE(8), IDG(8), IPEM(8) 1.PBACK(2),PION(3) COMMON /TPRIM/ THMCPL(8) COMMON /MACDEF/ TCDEF(8,4),NTC,PRDEF(8,6),NPR,BACK(4,2) COMMON /DET/ AGRD, AFIL, VEB, VIE, VEXT 1, VEN, QPST, QPR, MAS, VEX, VDK, VPMT DEFINITION OF /COMMON/ VARIABLES AND ARRAYS C COMMON Block /CUR/ Contains the current data acquisition parameters:

С С C С C C

С

С

Շ Շ ANGLE The detector angle. 0000000000 INTERV The countdown interval on the TGS. The number of detector angles that have been recorded. NUMANG MANGLE The position of the current angle out of all the angles that have been recorded (Data is stored in order of increasing angle). The TGS running flag. If IRNFLG > 0, the TGS is running. IRNFLG Contains the accumulated data, and two work arrays: COMMON Block /DATA/ C С DATA(I, J)The accumulated data, stored in order of increasing С angle. Different detector angles are different С subscript I's. The J's are: CCCCCCCCC j=1 Detector Angle J=2 Time Recorded J=2\*n+1 Channel n Signal } n=1.8 J=2\*n+2 Channel n Background } n=1,8 WORK1(I) Work Buffer for upper display. Subscript I is the detector angle subscript, as in DATA(I,J). WORK2(I) The same as WORK1(I), but for the lower display channel. С COMMON Block /FNAME/ Contains the filename, and the number of data points C C C C C C C C C C C C C C C C C stored in it. INAME(I) The ASCII name of the data file, for example: DLO: ANG123 . ANG ISTORE The number of data points stored in the data file. C Č COMMON Block /DISP/ Contains the display array and modes. С С IDISP(I) The array for display by the Data Translation DT2771 С Direct Memory Access Display Driver. MODSPT The Display Mode of the upper (Top) display channel. MODSPB The Display Mode of the lower (Bottom) display channel. For the values of MODSPT/B, see Subroutines SCHDSP and SDCHOS. DSPB The initial detector angle for display. The angular step for the display (i.e. the display DSTEP resolution). NDSP The number of detector angles to be diplayed. Up to 256. INDEP Flag to determine whether or not to scale display channels independently. If INDEP > 0, display channels independently. С COMMON Block /DSPCHA/ Contains the data channels to be displayed. С

MDSPT, NDSPT The data channels to be used for the upper (Top) C C C C C C C C display channel. The data channels to be used for the lower (Bottom) MDSPB, NDSPB display channel. С С COMMON Block /DSPZER/ Contains the zero levels for the two display С channels. С С ITYZER The zero level for the top display channel. С IBYZER The zero level for the bottom display channel. С С C COMMON Block /DSCR/ Contains the scratch values of the current data С for the display loader Subroutine SDVAL. С C C DSCR(I) The scratch data for the display loader. I has the same meaning as the J subscript in DATA(I, J) above. Ċ С C COMMON Block /SCR/ Contains the scratch values of the data. С С С С С С The scratch data, kept here while determining if it is SCR(I) to be saved or not. I has the same meaning as the J subscript in DATA(I,J) above. C С COMMON Block /TGSMAX/ Contains the maximum values associated with the С Timer-Gater-Scaler system in use. These are set С in Subroutine SINIT, for easy manipulation. С C C DWLMAX No Meaning. MAXCHN The maximum number of data channels allowed. С С C COMMON Block /SETS/ Contains the TGS Settings. С С GATE(I) The gating times of each of up to 16 scalers. C The time each of up to 16 scalers are recorded. TIM(I) Č EXCFRC(I) The excited state fractions for each excited state. C C C CLOCK The TGS Clock Frequency. INTACQ The standard TGS countdown interval. NCHAN The number of data channels eing recorded. С С С COMMON Block /LASSET/ Contains the Velocity measurement settings. С C C LASDWL The dwell time for each point of a computer-driven laser scan. Not currently used. C FRQLAS The laser frequency in cm-1.

C The D/A step size for a computer-driven laser scan. LASSTP Not currently used. **FSRLAS** The Free Spectral Range of the Frequency Standard (i.e. Fabrey-Perot Etalon) in GHz. LASCHA The data channel on which the laser fluorescence and Frequency Standard signal come. COMMON Block /LOPSET/ Contains the Loop mode settings. ANLPBE The beginning angle of the Loop. ANLPFN The ending (finishing) angle of the Loop. ANLPIN The angular interval (step size) of the Loop. LPNRFQ The frequency with which normalizations are to be recorded for the Loop (after LPNRFQ angles are recorded, a normalization measurement is made). ANLPCU The current angle in the Loop. The flag to show that SANG is in Loop Mode. If LPFLAG LPFLAG < 0, Loop mode has just begun, and a normalization is to be done. If LPFLAG > 0, LPFLAG = the number of angles since the last normalization. COMMON Block /NRMSET/ Contains the Normalization Settings. С C ANGNRM The normalization detector angle. С INTNRM The countdown interval for normalization measurements. С NRMRPT The number of times the normalization measurement is С repeated. C C C C C C C C NRMFLG The Normalization Flag. If NRMFLG > 0, NRMFLG counts the number of times that the TGS has counted down (until NRMFLG = NRMRPT). С COMMON Block /POLSET/ Contains the Rotation of Polarization Settings. ANGPOL The detector angle at which rotation of polarization measurements are to be made. POLSTP The step size of the polarization rotator in degrees. POSIT The ideal position of the rotator in steps if step resolution were infinite. IPOSIT The actual position of the rotator (0-125). The rotation of polarization flag. If IPLFLG>O, IPLFLG SANG is acquiring polarization rotation information. С C COMMON Block /TSTSET/ Contains the Test Mode Settings.

С The detector angle for Test Mode Measurements. C C C C C C C C C C C TSTANG NUMTST The number of TGS countdowns to do in Test Mode. The current angular offset to use for storing a data TSTCUR point. Each test mode point is stored in a different С С COMMON Block /VELO/ Contains the measured beam velocity. С С DELF The measured distance between the 45 and 90 degree laser С crossing in arbitrary units. C C HALFH The measured Full Width Half Maximum for the 45 degree peak in the units of DELF. С DELPIP The measured distance between Frequency Standard peaks С in the units of DELF. Ċ VELB The beam velocity in cm/sec. C C SPDRIO The beam speed ratio. С C COMMON Block /SDONE/ Contains the flag to show that the TGS has counted С down, and the raw data in I+4. С Ċ ISWFLG The flag to show that the TGS has counted down. If 00000000 ISWFLG > 0, the TGS has counted down, and has been serviced and read by an interrupt service routine, and needs to have Subroutine SDONE, capture the data for display and saving. If ISWFLG = 0, accumulation has begun, and the display needs to be updated by Subroutine SDISP, via Subroutine SDONE. In both cases, ISWFLG is cleared by Subroutine SDONE. C C C C C C C C C JDATA(I) The raw data in an INTEGER\*4 array, as read from the scalers through the CAMAC crate. The subscript I corresponds to the scaler read: I=2\*n-1 Channel n Signal I=2∗n Channel n Background С C COMMON Block /CHADEF/ Contains the definitions of the data channels. С Շ Շ CHADEF(I,J) The definition of channel J, up to 80 ASCII characters, stored in I=1,20 С С C COMMON Block /PRES/ Contains the measured machine pressures. Note that С the pressure regions can be defined by the user. С С PRM(I) The mantissa of the ionization gauge reading in pressure C region I.

C۰ IPRE(I) The exponent of the ionization gauge reading in pressure C reaion I. 0000000000 The foreline pressure in region I. PFORE(I) IDG(I) The flag for whether the Degas is on in pressure region I. If IDG=1, the Ion Gauge is on; if IDG=2, the Ion Gauge is off. IPEM(I) The emission current of the ionization gauge reading in pressure region I. PBACK(I) The backing pressure of the sources. I=1 for the Primary Source; I=2 for the Secondary Source. PION(I) The ion pump currents in detector region I. C C COMMON Block /TPRIM/ Contains the measured thermocouple voltages. C C C TPRIM(I) The measured thermocouple voltages on thermocouple I. С C COMMON Block /MACDEF/ Contains the definitions and numbers of thermocouples CCCCCCCCCCCCCC pressure regions, and backing gases. TCDEF(I,J) The definition of thermocouple I, in up to 16 ASCII characters, in J=1-4. NTC The number of thermocouples defined. PRDEF(I,J) The definition of pressure region 1, in up to 24 ASCII characters, in J=1-6. NPR The number of pressure regions defined. BACK(I,J) The definition backing gas J, in up to 16 ASCII characters, in I=1-4. Ċ COMMON Block /DET/ Contains the Detector settings. CCCCCCCCCCCCC AGRD The grid (emission) current, in mA. AFIL The filament current, in A. VEB The electron beam energy, in V. VIE The ion energy, in V. VEXT The ion extractor voltage, in V. VEN The entrance lens voltage, in V. **QPST** The Quadrupole Mass Spectrometer Mass Potentiometer Setting. **QPRS** The Quadrupole Mass Spectrometer Resolution Potentiometer Settina. MAS The mass measured, in amu. VEX The exit lens voltage, in V. С VDK The doorknob voltage, in kV. С VPMT The photomultiplier tube supply voltage, in V. С Č С C\*

INITIALIZE ALL COMMON BLOCK DATA AND FLAGS

CALL SINIT

OPEN A DATA FILE FOR OUTPUT

CALL SFILE(1)

TRANSFER CONTROL TO SCOMM FOR USER COMMANDS AND CONTROL

CALL SCOMM END

SUBROUTINE SCOMM

Second routine in file SANG.FOR.

Version Date: August 15, 1985 Paul S. Weiss

Subroutine SCOMM maintains control of operations for the program.

It accepts 1 or 2 character commands from the user, then calls the appropriate subroutines.

INTEGER COM1, COM2, COMT INTEGER+4 JDATA COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /FNAME/ INAME(7), ISTORE COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB COMMON /DSPZER/ ITYZER, IBYZER COMMON /DSCR/ DSCR(20) COMMON /SCR/ SCR(20) COMMON /TGSMAX/ DWLMAX, MAXCHN COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG COMMON /VELO/ DELF, HALFH, DELPIP, VELB, SPDRIO COMMON /SDONE/ ISWFLG, JDATA(20) COMMON /CHADEF/ CHADEF(20,10) COMMON /PRES/ PRM(8), IPRE(8), PFORE(8), IDG(8), IPEM(8) 1, PBACK(2), PION(3) COMMON /TPRIM/ THMCPL(8)

С

С

Ç

Ċ

С

С

~	COMMON /MACDEF/ TCDEF(8,4),NTC,PRDEF(8,6),NPR,BACK(4,2) COMMON /DET/ AGRD,AFIL,VEB,VIE,VEXT 1,VEN,QPST,QPR,MAS,VEX,VDK,VPMT
	THE FOLLOWING IS A TABLE OF THE ASCII CHARACTERS USED AS COMMANDS
•	DATA IA, IB, IC, ID, IE /65,66,67,68,69/ DATA IF, IG, IH, II, IL, IM, IN /70,71,72,73,76,77,78/ DATA IO, IP, IR, IS, IT /79,80,82,83,84/ DATA IU, IV, IW, IX, IZ /85,86,87,88,90/ TYPE 2100
1 95 C	TYPE 95 FORMAT(' Command?',/)
C C	RECORD COMMAND ENTRY USING SPECIAL SCREEN MODE
200	CALL ENSPEC IF(ISWFLG.GE.O) CALL SDONE COM1 = ITTINR(1) IF(COM1 + T.O) CO TO 200
С	IGNORE CARRIAGE RETURNS AND LINE FEEDS IF((COM1.EQ.13).OR.(COM1.EQ.10)) GO TO 200 COM2 = 32
C C C	ECHO CHARACTERS
250 255 300	TYPE 255,COM1,COM2 FORMAT('+',2A1,25X) IF(ISWFLG.LT.O) GO TO 320 CALL SDONE GO TO 250
320	COMT = ITTINR(1) IF(COMT.LT.0) GO TO 300 IE(COMT.ED 13) GO TO 500
380 390 400	COM2 = COMT TYPE 255,COM1,COM2 IF(ISWFLG.LT.O) GO TO 420 CALL SDONE GO TO 390
C C	LOOK FOR CARRIAGE RETURN, OR ANOTHER CHARACTER
420	COMT = ITTINR(1) IF(COMT.LT.O) GO TO 400 IF(COMT.EQ.13) GO TO 500
	IF MORE THAN TWO CHARACTERS ARE TYPED IN, SHIFT COM2 TO COM1
v	COM1 = COM2 GD TO 380

READ SECOND HALF OF CARRIAGE RETURN (LINE FEED) 500 COMT = ITTINR(1)IF(COMT.LT.O) G0 T0 500 TYPE 510, COM1, COM2 510 FORMAT('+',2A1,25X,/) CALL EXSPEC EXIT SPECIAL SCREEN MODE BEFORE CHECKING COMMAND LIST GO TO 1300 MATCH UP COMMANDS AND EXECUTE DESIRED FUNCTION THE BASIC COMMANDS ARE: C\* Change Commands D\* Define and Display Commands EX Exit G\* Graph Commands Н Help Loop Commands L+ Machine Condition Commands M\* N+ Normalization Commands P+ Print Commands Record Commands R\* S\* Show Commands T+ Test Commands Velocity Measurement Commands V. Write Commands W+ Z\* Zero Commands C\* The Change Commands are: CA Change Angle CC Change Clock Frequency CE Change Excited State Fractions Change File CF CG Change Gate Widths CH C Help CI Change Interval for Acquisition Countdown CL Change Loop Parameters CN Change Normalization Parameters CP Change Data Acquisition Parameters CU Change Number of Channels in Use C۷ Change Velocity Measurement Parameters

C С

С

С С

C

С С С

С

С

С

С

С

С

С

С

Ċ

С

č

C C

C

CCCCCC

Ĉ

С

С

C C

С

C

С

С

C

С

C

C C

С

	С		
•	1300	IF(COM1.NE.IC) IF(COM2.NE.IA) CALL SCANG	GO TO 1401 GO TO 1311
	1311	GO TO 1 IF(COM2.NE.IC) CALL SCHPAR(2)	GO TO 1321
	1321	GU TU 1 IF(COM2.NE.IE) CALL SCHPAR(5)	GO TO 1326
	1326	IF(COM2.NE.IF) CALL SFILE(0)	GO TO 1331
	1331	IF (COM2.NE.IG) CALL SCHPAR (4)	GO TO 1336
	1336	IF (COM2.NE.II) CALL SCHPAR (3)	GO TO 1341
	1341	IF (COM2.NE.IL) CALL SCHLOP(O)	GO TO 1351
	1 <b>351</b>	IF (COM2.NE.IN) CALL SCHNRM(O)	GO TO 1356
	1356 <sub>.</sub>	IF (COM2.NE.IP) CALL SCHPAR(O)	GO TO 1376
	1376	IF(COM2.NE.IU) CALL SCHPAR(1)	GO TO 1381
•	1381	IF (COM2.NE.IV) CALL SCHVEL (O)	GO TO 1399
	1399	IF (COM2.NE.IH) CALL SCOMLC	TYPE 9100
	C		
	C	D*	
	C	The definition	and display commands are:
		UA DR	Display: Change All
	c		Define Channels
	Č	DH	D Help
	C	DI	Display: Scale Channels Independently
	C	DL	Display: Change Lower Channel
	C	DP	Define Pressure Regions
		DK	Display: Ketresh Display: Scale Changels to the Second Veture
	6	00	vispiay: scale channels to the same value

.

С.	DT	Define The	ermocouples
C	DU	Display:	Change Upper Channel
С	DW	Display:	Change Window
С			
1401	IF (COM1.NE.ID)	GO TO 1501	
	IF(COM2.NE.IA)	GO TO 1406	
	CALL SCHDSP(0)		
	GO TO 1		
1406	IF(COM2.NE.IB)	GO TO 1411	
	CALL SCHDEF(4)		
	GO TO 1		
1411	IF(COM2.NE.IC)	GO TO 1436	
	CALL SCHDEF(1)		
	GO TO 1		
1436	IF(COM2.NE.II)	GO TO 1441	
	INDEP = 1		
	GO TO 1		
1441	IF(COM2.NE.IL)	GO TO 1456	
	CALL SCHDSP(2)		
	GO TO 1		
1456	IF(COM2.NE.IP)	GO TO 1461	
	CALL SCHDEF (3)		
	GO TO 1		_
1461	IF(COM2.NE.IR)	GO TO 1471	
	CALL SDISP		
	GO TO 1		
1471	IF (COM2.NE.IS)	GU TU 1473	
	INDEP = 0		•
	GU IU 1		
1473	1F(CUM2.NE.11)	GU IU 1476	
	CALL SCHUEF (2)		
1 4 7 6		00 TO 1401	
14/0	IF (CUM2.NE.10)	GU IU 1481	
	CALL SCHUSP(I)		
1 4 0 1		00 TO 1400	
1481	IF (CUM2.NE.IW)	GU 10 1499	
	CALL SCHUWN(U)		
1400		TYPE 0100	
1499	CALL SCOULD	TIPE 9100	
	CALL SCUMED		
c	00 10 I		
c	F		
Ċ	Evit		
Ċ	LAIV		
1501	TE(COM1 NE TE)	GO TO 1701	
	CALL SLEAVE		
÷	GO TO 1		
С			
С			

,

C	0.
Č	u# The graph commands are:
Č	AP Creat Path Disalawa
	GD Graph Doth Displays
C	GH G Help
C	GL Graph Lower Display
С	GU Graph Upper Display
С	
1701	IF(COM1.NE.IG) GO TO 1801 IF(COM2.NE.IB) GO TO 1741 CALL SGRAPH(O)
	GO TO 1
1741	IF(COM2.NE.IL) GO TO 1776 CALL SGRAPH(-1) CO TO 1
1776	IF(COM2.NE.IU) GO TO 1799 CALL SGRAPH(1)
	GU IO 1
1799	IF(COM2.NE.IH) TYPE 9100 CALL SCOMLG
C I	
C	
C C	Help Print out command list.
1801	IF(COM1.NE.IH) GO TO 2201 CALL SCMLST GO TO 1
C	
č	
Ċ	L+ The loss commands and
	ine loop commands are:
C	LA Loop: Change Current Angle
C	LB ´ Loop: Change Beginning Angle
C	LC Loop: Continue
С	LD Loop: Change Direction
С	LE Loop: Change Ending Angle
Č	
č	
č	
	LN Loop: Change Normalization Frequency
	LP Unange All Loop Parameters
C A	LS Loop: Start
C	
2201	IF(COM1.NE.IL) GO TO 2301 IF(COM2.NE.IA) GO TO 2206 CALL SCHLOP(6) GO TO 1
2206	IF(COM2.NE.IB) GO TO 2211 CALL SCHLOP(1)
0011	UU IU I TE(COMO NE TC) OO TO OO1C
2211	IF(CUMZ.NE.IC) GU IU 2216

	CALL SLOOP(0)	
	GO TO 1	
2216	IF (COM2.NE.ID) CALL SCHLOP(5)	GU 10 2221
1	GO TO 1	
2221	IF(COM2.NE.IE)	GO TO 2236
•	CALL SCHLOP(2)	
	GO TO 1	
2236	IF(COM2.NE.II)	GO TO 2251
	CALL SCHLOP(3)	
	GO TO 1	
2251	IF(COM2.NE.IN)	GO TO 2256
	CALL SCHLOP(4)	
	GO TO 1	
2256	IF(COM2.NE.IP)	GD TD 2266
	CALL SCHLOP(0)	
	GO TO 1	· · · · · · · · · · · · · · · · · · ·
2266	IF(COM2.NE.IS)	GD TO 2299
	CALL SLOOP(1)	
	GO TO 1	
2299	IF (COM2.NE.IH)	IYPE 9100
	CALL SCUMEL	
~	GU TU I	, · · · ·
C	44	
	M# The Machine Co.	
		Naching Conditions, Change All Detector Setting
	MA	Machine Conditions: Change All Detector Settings
	MD	Machine Conditions: Change Dacking Pressures
	NE	Machine Conditions: Change Daily Detector (post-urms)
		and Topizon Settinge
c c	ML	Machine Conditions: Help
	MT	Machine Conditions: Change Ten Long Settings
c		Machine Conditions: Change Buadrupolo
c	141141	Mass Spectrometer Settings
c c	MR	Machine Conditions: Change Regional Pressures
c	MS	Machine Conditions: Change Source Conditions
ĉ	СМ ТМ	Machine Conditions: Change Source conditions
ĉ		waterine conditions. Change intermotouple headings
2301	TE(COM1 NE TM)	GD TD 2401
	TF(COM2 NE TA)	GD TD 2306
	CALL SCHMAD(5)	
	GO TO 1	
2306	IF (COM2.NE.IB)	GO TO 2316
	CALL SCHMAP(1)	
	GO TO 1	
2316	IF (COM2.NE.ID)	GO TO 2326
	CALL SCHMAD(4)	
	GO TO 1	
2326	IF(COM2.NE.IF)	GO TO 2336
	· · ·	

CALL SCHMAD(1) GO TO 1 2336 IF(COM2.NE.II) GO TO 2346 CALL SCHMAD(2) GO TO 1 2346 IF (COM2.NE.IM) GO TO 2361 CALL SCHMAD(3) GO TO 1 2361 IF(COM2.NE.IR) GO TO 2366 CALL SCHMAP(3) GO TO 1 IF(COM2.NE.IS) GO TO 2371 2366 CALL SCHMAP(4) GO TO 1 2371 IF(COM2.NE.IT) GO TO 2399 CALL SCHMAP(2) GO TO 1 2399 IF(COM2.NE.IH) TYPE 9100 CALL SCOMLM GO TO 1 CCCCCCCCC N\* The available Normalization Commands are: NA Change Normalization Augle NH N Help NI Change Normalization Countdown Interval NP Change All Normalization Parameters NR Change Normalization Repetitions С NS Normalization: Start č 2401 IF (COM1.NE.IN) GO TO 2601 IF (COM2.NE.IA) GO TO 2436 CALL SCHNRM(1) GO TO 1 2436 IF(COM2.NE.II) GO TO 2456 CALL SCHNRM(2) ~ GO TO 1 🕔 IF(COM2.NE.IP) GO TO 2461 2456 CALL SCHNRM(O) GO TO 1 2461 IF(COM2.NE.IR) GO TO 2466 CALL SCHNRM(3) GO TO 1 2466 IF(COM2.NE.IS) GO TO 2499 CALL SNORM GO TO 1 2499· IF (COM2.NE.IH) TYPE 9100 CALL SCOMLN GO TO 1

C C C C C P\* The P\*, S\*, and W\* commands are sorted out in subroutine SNUCH The Print Commands are as follows: C PA Print Accumulated Data С PC Print Current Configuration С PD Print Data Channel Definition С PG Print Curent Graphics Display С PH P Help C C C C C C C C PI Print Information PL Print Loop Mode Parameters PM Print Machine Condition PN Print Normalization Parameters PS Print Detector Settings С PT Print Test Mode Parameters С 2601 IF(COM1.NE.IP) GO TO 2801 NU = 6CALL SNUCH (NU, COM2) IF (NU.NE.O) GO TO 1 IF (COM2.NE.IH) TYPE 9100 2699 CALL SCOMLP GO TO 1 С С R С The available Record Data and Rotation of Polarization Commands are: С Rotation of Polarization: Change Detector Angle RA С RD Record Data С RH R Help Ĉ RI Reinitialize TGS System Ċ RM Rotation of Polarization: Change Motor Step Size C RP Rotation of Polarization: Change All Parameters C RS Rotation of Polarization: Start С 2801 IF(COM1.NE.IR) GO TO 2901 IF(COM2.NE.IA) GO TO 2816 CALL SCHPQL(1) GO TO 1 IF(COM2.NE.ID) GO TO 2836 2816 CALL SSTART(1) GO TO 1 2836 IF(COM2.NE.II) GO TO 2846 CALL INIT(ISWFLG, JDATA) GO TO 1 2846 IF(COM2.NE.IM) GO TO 2856 CALL SCHPOL(2) GO TO 1 2856 IF(COM2.NE.IP) GO TO 2866 CALL SCHPOL(0)

С

GO TO 1 2866 IF(COM2.NE.IS) GO TO 2899 CALL SPOLAR GO TO 1 IF(COM2.NE.IH) TYPE 9100 2899 CALL SCOMLR GO TO 1 S\* The P\*, S\*, and W\* commands are sorted out in subroutine SNUCH The available show commands are: SA Show Accumulated Data SC Show Current Configuration SD Show Data Channel Definition SG Show Curent Graphics Display SH S Help SI Show Time Information SL Show Loop Mode Parameters SM Show Machine Condition SN · Show Normalization Parameters SS Show Detector Settings ST Show Test Mode Parameters С 2901 IF(COM1.NE.IS) GO TO 3001 NU = 7CALL SNUCH (NU, COM2) IF (NU.NE.O) GO TO 1 2999 IF (COM2.NE.IH) TYPE 9100 CALL SCOMLS GO TO 1 С С T\* С The Test Commands are as follows: CCCCCCCCC TA Change Test Angle TC Test Mode: Continue TE Test Mode: Exit TH T Help TN Test Mode: Change Number of Countdowns TP Change All Test Mode Parameters TS Test Mode: Start 3001 IF(COM1.NE.IT) GO TO 3201 IF(COM2.NE.IA) GO TO 3011 CALL SCHTST(1) GO TO 1 3011 IF(COM2.NE.IC) GO TO 3021 CALL STEST(1) GO TO 1 3021 IF(COM2.NE.IE) GO TO 3051

	TSTFLG = - TST	FLG
	GO TO 1	
3051	IF (COM2.NE.IN)	GO TO 3056
	CALL SCHTST(2)	
	GO TO 1	
3056	IF(COM2.NE.IP)	GO TO 3066
	CALL SCHTST (0)	
	GO TO 1	
3066	IF (COM2.NE.IS)	GO TO 3099
	CALL STEST (0)	
	GO TO 1	
3099	TE(COM2 NE TH)	TYPE 9100
0000		
	CO TO 1	
r	60 10 I	
	V.	
	The Velseits W	annungen Commande and an fallower
	The velocity M	Channel Commands are as follows:
C C		Change Velocity Measurement Accumulation lime
		Change velocity Measurement Frequency Standard
۲ ۲	VH	
	VL.	Change Velocity Measurement Laser Frequency
	VM	Measure Velocity
C C	VP	Change All Velocity Measurement Parameters
C	VS	Change Velocity Measurement Step Size
C		
3201	IF (COM1.NE.IV)	GO TO 3301
	IF(COM2.NE.IA)	GD TD 3226
	CALL SCHVEL(4)	
	GO TO 1	
3226	IF(COM2.NE.IF)	GO TO 3241
	CALL SCHVEL(2)	
	GO TO 1	
3241	IF(COM2.NE.IL)	GO TO 3246
	CALL SCHVEL(1)	
	GO TO 1	
3246	IF(COM2.NE.IM)	GO TO 3256
	CALL SVELO	
	GO TO 1 🔪	
3256	IF(COM2.NE.IP)	GO TO 3266
	CALL SCHVEL(0)	
	GO TO 1	
3266	IF (COM2.NE.IS)	GO TO 3299
•	CALL SCHVEL (3)	
		*
3299	TE(COM2 NE TH)	TYPE 9100
0200		
	CO TO 1	
c		
č	Wat	·
č.		d We commande are corted out is subsoution CNUCH
6	ine r≠, o≠, an	U #+ COMMANUS ALE SOLVEY OUC IN SUDIOUCINE SNUCH
Ը Շ The Write Commands are as follows: WA Write Accumulated Data С WC Write Current Configuration С WD Write Data Channel Definition C C C C C C C C WG Write Curent Graphics Display WH W Help Write Information WI WL Write Loop Mode Parameters WM Write Machine Condition С WN Write Normalization Parameters Č WS Write Detector Settings С WT Write Test Mode Parameters С 3301 IF(COM1.NE.IW) GO TO 3601 NU = 9CALL SNUCH (NU, COM2) IF (NU.NE.O) GO TO 1 3399 IF(COM2.NE.IH) TYPE 9100 CALL SCOMLW GO TO 1 С С Ζ\* C The Zero Commands are as follows: C ZA Zero Arrays C C ZH Z Help ZS Zero Angular Segment С ΖT Zero Test Mode Data С 3601 IF(COM1.NE.IZ) GO TO 3700 IF (COM2.NE.IA) GO TO 3666 CALL SZERO(1) GO TO 1 3666 IF(COM2.NE.IS) GO TO 3671 CALL SZER0(2) GO TO 1 3671 IF(COM2.NE.IT) GO TO 3699 CALL SZERO(3) GO TO 1 3699 IF(COM2.NE.IH) TYPE 9100 CALL SCOMLZ GO TO 1 С С TYPE OUT THE FOLLOWING IF AN UNRECOGNIZED COMMAND WAS ISSUED C 3700 **TYPE 9100 TYPE 2100** FORMAT(' For a list of commands, type H') 2100 GO TO 1 9100 FORMAT(/,' This is not a command.') RETURN

END

C

C C

C C C C C C SUBROUTINE SACCUM(NU)

File SACCUM.FOR begins on the previous line.

Version Date: July 5, 1984 Paul S. Weiss

Subroutine SACCUM accumulates the data and outputs it to unit NU.

С COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN С С RETURN IF TGS IS RUNNING С IF(IRNFLG.EQ.O) GO TO 50 TYPE \*, 'Abort Data Accumulation -- TGS is running!' GO TO 9000 50 IF (NUMANG.GT.O) GO TO 100 TYPE \*, 'No data has been accumulated!' GO TO 9000 100 WRITE(NU,110) 12, NUMANG FORMAT(1X, A1, ' ACCUMULATED DATA', /, 1X, I3, ' angles') 110 DO 900 IN=1, NUMANG WRITE(NU, 300) DATA(IN, 1), DATA(IN, 2) FORMAT(/, ' Angle: ', F9.5, ' degrees, measured for: ', F8.3 1, ' sec.',/, ' Channel', 3X, 'A', 10X, 'B', 11X, 'Signal' 2,9X, 'Error', 6X, 'Signal/Noise') 300 DO 700 J=1,NCHAN IS = (2\*J) + 1IB = IS + 1SIG = (DATA(IN, IS) - DATA(IN, IB))/DATA(IN, 2)ERR = (SQRT(DATA(IN, IS) + DATA(IN, IB)))/DATA(IN, 2)SGTON = 0.IF(ERR.NE.O.) SGTON = SIG/ERR WRITE (NU, SOO) J, DATA (IN, IS), DATA (IN, IB), SIG, ERR, SGTON . 500 FORMAT(1X, I1, 2X, F9.0, 2X, F9.0, 5X, F11.2, '+/-', F9.2, 5X, F9.1) 700 CONTINUE CONTINUE 900 WRITE(NU,1100) 12 FORMAT(1X,A1,' RAW DATA',/) 1100 9000 RETURN END

SUBROUTINE SCANG

С

File SCANG.FOR begins on the previous line.

Version Date: June 28, 1984 Paul S. Weiss Subroutine SCANG changes the detector angle. COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG TYPE \*,'What is the new detector angle?' ACCEPT \*, ANGLE RETURN END SUBROUTINE SCHDEF (IOPT) File SCHDEF.FOR begins on the previous line. Version Date: July 18, 1984 Paul S. Weiss Subroutine SCHDEF changes the definitions of the: data channels, the thermocouples, the pressure regions, and the backing gases. The argument of SCHDEF determines the definitions to change, as follows: IOPT=1 Change Data Channel Definitions IOPT=2 Change Thermocouple Definitions IOPT=3 Change Pressure Region Definitions IOPT=4 Change Backing Gas Definitions COMMON /TGSMAX/ DWLMAX, MAXCHN COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /CHADEF/ CHADEF(20,10) COMMON /MACDEF/ TCDEF(8,4), NTC, PRDEF(8,6), NPR, BACK(4,2) GO TO(1100,2100,3100,4100),IOPT CHANGE DATA CHANNEL DEFINITIONS 1100 D0 1400 I=1,NCHAN TYPE 1210, I, (CHADEF (J, I), J=1, 20) 1210 FORMAT(' Change the definition of Channel #' 1,I1,'? The current definition is:',/,1X,20A4) ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE.'Y') GO TO 1400 **TYPE 1310** 1310 FORMAT(' Enter the new definition.') ENABLE LOWER CASE INPUT

С CCCCCCC

C

С С

С

С С Ċ CALL LCASE ACCEPT 1350, (CHADEF (J, I), J=1, 20) 1350 FORMAT(20A4) С С DISABLE LOWER CASE INPUT С CALL UCASE 1400 CONTINUE GO TO 9000 С С CHANGE THE NUMBER AND DEFINITIONS OF THERMOCOUPLES С 2100 TYPE 2110,NTC FORMAT(' There are currently ', I1,' thermocouples in use.' 2110 1,' Has this changed?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 2300 2200 TYPE \*, 'How many thermocouples are being used?' ACCEPT +,NTC MAXTC = 8IF((NTC.LE.O).OR.(NTC.GT.MAXTC)) GO TO 2200 2300 D0 2500 I=1,NTC TYPE 2410, I, (TCDEF (I, J), J=1, 4) FORMAT(' Thermocouple #', I1, ' is currently: ',4A4 2410 1,/,' Has this changed?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 2500 **TYPE 1310** С С ENABLE LOWER CASE INPUT С CALL LCASE ACCEPT 1350, (TCDEF (I, J), J=1, 4) С С DISABLE LOWER CASE INPUT С CALL UCASE 2500 CONTINUE GO TO 9000 С С CHANGE THE NUMBER AND DEFINITIONS OF THE PRESSURE REGIONS С 3100 TYPE 3110, NPR FORMAT(' There are ', I1,' pressure regions ' 3110 1, 'measured with ion gauges. Has this changed?') ACCEPT 1000, ANS IF (ANS NE. 'Y') GD TO 3300 3200 TYPE \*, 'How many pressure regions are being measured?' ACCEPT +, NPR

MAXPR = 8IF((NPR.LE.O).OR.(NPR.GT.MAXPR)) GO TO 3200 3300 D0 3500 I=1,NPR TYPE 3410, I, (PRDEF (I, J), J=1, 6) 3410 FORMAT(' Change the definition of Pressure Region #' 1, I1, '?', /, ' The current definition is: ',6A4) ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 3500 **TYPE 1310** С С ENABLE LOWER CASE INPUT С CALL LCASE ACCEPT 1350, (PRDEF (I, J), J=1,6) С С DISABLE LOWER CASE INPUT С CALL UCASE 3500 CONTINUE GO TO 9000 С С CHANGE BACKING GAS DEFINITIONS С 4100 TYPE 4210, (BACK(J,1), J=1, 4)4210 FORMAT(' The Primary backing gas is ',4A4 1,' Has this changed?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 4400 **TYPE 1310** С С ENABLE LOWER CASE INPUT С CALL LCASE ACCEPT 1350, (BACK (J, 1), J=1,4) С С DISABLE LOWER CASE INPUT С CALL UCASE 4400 TYPE 4410, (BACK(J, 2), J=1, 4)4410 FORMAT(' The Secondary backing gas is ',4A4 1,' Has this changed?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 9000 **TYPE 1310** С С ENABLE LOWER CASE INPUT C CALL LCASE ACCEPT 1350, (BACK (J,2), J=1,4) С

#### DISABLE LOWER CASE INPUT

С С

CCCCCCC

С

C

C

Ċ C C

Ċ

CALL UCASE 9000 RETURN END

SUBROUTINE	SCHDSP	(IUPL(	JW)
			<u> </u>

File SCHDSP.FOR begins on the previous line.

Version Date: July 5, 1984 Paul S. Weiss

IF(IUPLOW.EQ.2) GO TO 2100

Subroutine SCHDSP changes one or both of the display channels.

The argument IUPLOW specifies which display channel will be changed, as follows:

IUPLOW=0	Change both channels
IUPLOW=1	Change Top channel
IUPLOW=2	Change Bottom Channel

COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN

С

50

TYPE 100, NCHAN, 'upper' 100 FORMAT(/,' The available diplay modes are (m,n=1-',I1,'):' Channel n Signal' 1,/,' Cn 2,/,' 3,/,' 4 / ' Dmn Difference of 2 Channels (m-n)' Emn Excited State Signal' 4,/, Fmn Signal/Noise of the Difference of' 'Channels (m-n)' 4, 5,/,' Sn Signal/Noise of Channel n' 6,/, Τn Time an Angle has been measured on Channel n', 7,/,' Please enter your choice for the ',A5,' display.') CALL SDCHOS(MODSPT,MDSPT,NDSPT) IF (MODSPT.EQ.O) GO TO 50 IF(IUPLOW.EQ.1) G0 T0 8000 TYPE 100, NCHAN, 'lower' 2100 CALL SDCHOS (MODSPB, MDSPB, NDSPB) IF (MODSPB.EQ.0) G0 T0 2100 TYPE \*, 'Scale the channels independently?' 8000 INDEP = 1ACCEPT 1000, ANS 1000 FORMAT(A1) IF(ANS.NE.'Y') INDEP = 0 RETURN

END · SUBROUTINE SDCHOS (MD, M, N) Second routine of file SCHDSP.FOR. Version Date: June 25, 1984 Paul S. Weiss Subroutine SDCHOS Sets the mode number for the various display modes. COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN DATA NANS /0/ M = 0N = 0ACCEPT 2000, NANS, M, N 2000 FORMAT(A1,2I1) MD = 0IF(NANS.EQ.'C') MD = 1IF(NANS.EQ.'D') MD = 2IF(NANS.EQ.'E') MD = 3IF(NANS.EQ.'F') MD = 4IF(NANS.EQ.'S') MD = 5IF(NANS, EQ. 'T') MD = 6C C C BE SURE THAT THERE IS A CHANNEL TO PLOT IF((M.EQ.O).OR.(M.GT.NCHAN)) MD = 0С IF TWO CHANNELS ARE REQUIRED, MAKE SURE BOTH ARE OK С IF ( (MD.GT.1) . AND. (MD.LT.5) . AND. ('(N.EQ.O) . OR. (N.GT. NCHAN) 1.0R.(M.EQ.N)) MD = 0IF (MD.EQ.0) TYPE 3100,7,7 FORMAT(/, 'Invalid Option!!',2A1) 3100 RETURN END SUBROUTINE SCHDWN (IOPT) С File SCHDWN.FOR begins on the previous line. Ċ С Version Date: September 27, 1984 C Paul S. Weiss С С С Subroutine SCHDWN changes the display window. С С The argument of SCHDWN is as follows: С

CCCCCCC

С

С IOPT=O Ask user to choose display window. С IOPT=1 Set window for test mode. С IOPT=2 Set window for polarization rotation mode. С COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG IF(IOPT.EQ.O) GO TO 1 GO TO (3200,2200),IOPT 1 **TYPE 100** FORMAT(/,' The display window choices are:' 100 1,/,' D Default setting (-10 to 100 degrees)' 2,/,' 2,/,' Ρ Polarization Rotation setting (single angle)' T Test Mode setting (single angle)' 3,/, W Window setting (User set)' ' Please enter choice') 4,/, ACCEPT 1000, ANS 1000 FORMAT(A1) IF(ANS.NE.'D') GO TO 2100 DSPB = -10. DSTEP = .5NDSP = 221GO TO 5000 IF(ANS.NE.'P') GO TO 3100 2100 2200 DSPB = ANGPOLDSTEP = 1.E-4NDSP = 181GO TO 5000 IF(ANS.NE.'T') GO TO 4100 3100 3200 DSPB = TSTANGDSTEP = 1.E-5NDSP = 256GO TO 5000 4100 IF(ANS.NE.'W') GD TO 9100 TYPE **\*,'Enter** beginning angle of display window'. ACCEPT +, DSPB TYPE \*, 'Enter angular step for display' ACCEPT +, DSTEP IF(I0PT.NE.0) G0 T0 9000 5000 TYPE 5010, DSPB, (DSPB + ((NDSP-1) \*DSTEP)), DSTEP FORMAT(/,' The display will be from ',F9.5,' to ',F9.5, 5010 1, ' degrees, ',/,' in ',F7.5, ' degree steps.' 2,/,' Is this correct?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 1 9000 RETURN TYPE \*, 'This is not an option' 9100 GO TO 1 END

SUBROUTINE SCHLOP(IOPT) File SCHLOP.FOR begins on the previous line. Version Date: July 19, 1984 Paul S. Weiss Subroutine SCHLOP changes one or all of the Loop Mode parameters in SANG. The argument IOPT specifies which parameters are to be changed as follows: IOPT=0 Change all Parameters IOPT=1 Change Beginning Angle (ANLPBE) IOPT=2 Change Ending Angle (ANLPFN) IOPT=3 Change Angular Interval (ANLPIN) IOPT=4 Change Normalization Frequency (LPNRFQ) Other changes: IOPT=5 Change Direction (Sign of ANLPIN) IOPT=6 Change Current angle of Loop COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG IF(IOPT.GT.O) GO TO (1100,2100,3100,4100,5100,6100),IOPT 1100 TYPE 1110, ANLPBE FORMAT(' The beginning angle of the Loop is ',F5.1 1110 1,' degrees. Change this?') ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE.'Y') GO TO 2000 TYPE \*, 'Enter beginning angle of Loop' ACCEPT +, ANLPBE 2000 IF(IOPT.NE.O) GO TO 9000 2100 TYPE 2110, ANLPFN FORMAT(' The ending angle of the Loop is ',F5.1 2110 1,' degrees. Change this?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 3000 TYPE \*, 'Enter ending angle of Loop' ACCEPT +, ANLPFN IF(I0PT.NE.O) G0 T0 9000 3000 3100 TYPE 3110, ANLPIN 3110 FORMAT(' The angular interval of the Loop is ',F5.1 1,' degrees. Change this?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 4000 TYPE \*, 'Enter angular interval of Loop' ACCEPT +, ANLPIN С MAKE SURE THAT SIGN OF INTERVAL IS CORRECT

CCCCCC

Ċ

Ċ

С

Ċ

C C

C C C C C C C C

С

С

С

С	
4000	IF(((ANLPFN-ANLPBE)*ANLPIN).LT.O.) ANLPIN = -ANLPIN IF(IOPT.NE.O) GO TO 9000
	CHANGE NORMALIZATION FREQUENCY
4100 4110	TYPE 4110,LPNRFQ FORMAT(' The data is Normalized every ',I5 1,' measurements. Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GD TO 9000
4300	TYPE *, 'Normalize after how many measurements?' ACCEPT *,LPNRFQ IF(LPNRFQ.LE.O) GO TO 4300 GO TO 9000
	CHANGE DIRECTION OF LOOP
5100	A = ANLPBE ANLPBE = ANLPFN ANLPFN = A ANLPIN = - ANLPIN GO TO 9000
C C C	CHANGE CURRENT ANGLE
6100 6110	TYPE 6110,ANLPCU FORMAT(' The current angle is ',F5.1,' degrees. Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GO TO 9000 TYPE *,'Enter angle to measure next in Loop' ACCEPT *,ANLPCU
9000	RETURN END
C	SUBROUTINE SCHMAD(IDPT) File SCHMAD.FOR begins on the previous line.
C C C	Version Date: July 18, 1984 Paul S. Weiss
	Subroutine SCHMAD changes the detector settings.
,C C	Argument IOPT determines which settings are to be changed, as follows:
С	IOPT=1 Change Filament Settings

С IOPT=2 Change Ion Lenses' Settings C IOPT=3 Change Quadrupole Mass Spectrometer Settings Č C IOPT=4 Change Daly Detector Settings Change All of the Above IOPT=5 С COMMON /DET/ AGRD, AFIL, VEB, VIE, VEXT 1, VEN, QPST, QPR, MAS, VEX, VDK, VPMT GO TO (1100,2100,3100,4100,1100),IOPT 1100 TYPE 1110, AGRD 1110 FORMAT(' The Grid Current is ',F4.1,' mA.') **TYPE 1120** 1120 FORMAT(' Has this changed?') ACCEPT 1000.ANS 1000 FORMAT(A1) IF (ANS.NE. 'Y') GO TO 1200 TYPE \*, 'Enter Grid Current (in mA).' ACCEPT +, AGRD TYPE 1210, AFIL 1200 1210 FORMAT(' The Filament Current is ', F4.2, ' A.') **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 1300 TYPE \*, 'Enter Filament Current (in A).' ACCEPT \*, AFIL 1300 TYPE 1310, VEB FORMAT(' The Electron Beam Energy is ',F5.0,' V.') 1310 **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 1400 TYPE \*, 'Enter Electron Beam Energy (in V).' ACCEPT \*, VEB 1400 TYPE 1410, VIE 1410 FORMAT(' The Ion Energy is ', F7.2, ' V.') TYPE 1120 ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 2000 TYPE \*, 'Enter Ion Energy (in V).' ACCEPT \*,VIE IF(IOPT.NE.5) GO TO 9000 2000 2100 TYPE 2110, VEXT 2110 FORMAT(' The Ion Extraction Voltage is ',F5.0,' V.') **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 2200 TYPE \*,'Enter Ion Extraction Voltage (in V).' ACCEPT \*,VEXT 2200 TYPE 2210, VEN 2210 FORMAT(' The Entrance Lens Voltage is ',F5.0,' V.') **TYPE 1120** ACCEPT 1000, ANS

IF (ANS.NE.'Y') GO TO 3000 TYPE \*, 'Enter Entrance Lens Voltage (in V).' ACCEPT +, VEN IF(IOPT.NE.5) G0 T0 9000 3000 3100 TYPE 3110, QPST FORMAT(' The Quadrupole Mass Setting is ',F5.2) 3110 **TYPE 1120** ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 3200 TYPE \*, 'Enter the Mass Setting.' ACCEPT \*, QPST 3200 TYPE 3210, QPR FORMAT(' The Quadrupole Resolution is ',F5.2) 3210 **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 3300 TYPE \*, 'Enter Quadrupole Resolution.' ACCEPT \*, QPR 3300 TYPE 3310, MAS 3310 FORMAT(' The Mass is ', I5) **TYPE 1120** ACCEPT 1000, ANS IF(ANS.NE.'Y') GD TD 4000 TYPE +,'Enter mass.' ACCEPT +, MAS IF(IOPT.NE.5) G0 T0 9000 4000 TYPE 4110, VEX 4100 4110 FORMAT(' The Exit Lens Voltage is ', F6.0,' V.') **TYPE 1120** ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 4200 TYPE \*, 'Enter Exit Lens Voltage (in V).' ACCEPT +, VEX 4200 TYPE 4210, VDK 4210 FORMAT(' The Doorknob Voltage is ',F5.1,' kV.') **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE, 'Y') GO TO 4300 TYPE \*, 'Enter the Doorknob Voltage (in kV).' ACCEPT +, VDK 4300 TYPE 4310, VPMT FORMAT(' The Photomultiplier Voltage is ', F6.0' V.') 4310 **TYPE 1120** ACCEPT 1000, ANS IF (ANS.NE.'Y') GD TO 9000 TYPE \*, 'Enter Photomultiplier Voltage (in V).' ACCEPT +, VPMT 9000 RETURN END

SUBROUTINE SCHMAP(IOPT) File SCHMAP.FOR begins on the previous line. Version Date: September 4. 1984 Paul S. Weiss Subroutine SCHMAP changes the Machine pressures and temperatures. DIMENSION NONOFF(2) COMMON /PRES/ PRM(8), IPRE(8), PFORE(8), IDG(8), IPEM(8) 1, PBACK(2), PION(3) COMMON /TPRIM/ THMCPL(8) COMMON /MACDEF/ TCDEF(8,4),NTC,PRDEF(8,6),NPR,BACK(4,2) DATA NONOFF /'n ','ff'/ GO TO (1100,2100,3100,1100) IOPT TYPE 1110, PBACK (1), (BACK (J,1), J=1,4) 1100 FORMAT(/,' The Primary Source backing pressure is', F8.1 1110 1,' torr with ',4A4,/,' Has this changed?') ACCEPT 1000, ANS FORMAT(A1) 1000 IF (ANS.NE. 'Y') GD TO 1300 TYPE 1210, (BACK(J, 1), J=1, 4)1210 FORMAT(' What is the pressure of ',4A4,'?') ACCEPT +, PBACK(1) 1300 TYPE 1310, PBACK (2), (BACK (J, 2), J=1, 4) 1310 FORMAT(/,' The Secondary Source backing pressure is', F8.1 1,' torr with ',4A4,/,' Has this changed?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 1900 TYPE 1210, (BACK (J, 2), J=1, 4) ACCEPT \*, PBACK (2) IF(I0PT.NE.4) G0 T0 9000 1900 CHANGE THE MEASURED THERMOCOUPLE VOLTAGES 2100 DO 2700 I=1,NTC TYPE 2210, (TCDEF (I, J), J=1, 4), THMCPL (I)FORMAT(' The ',4A4,' thermocouple voltage is ',F5.2 1,' mV.',/,' Has this changed?') 2210 ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 2700 TYPE 2310, (TCDEF (I, J), J=1, 4) 2310 FORMAT(' What is the ',4A4,' thermocouple voltage ' 1,'(in mV)?') ACCEPT +, THMCPL(I) 2700 CONTINUE

С 000000

С

С

С С

С

GO TO 9000

	CHANGE REGIONAL PRESSURES
3100	D0 3500 I=1,NPR TYPE 3210,(PRDEF(I,J),J=1,6),PRM(I),IPRE(I),PFORE(I)
2010	1, NUNUFF(IDG(I)), IPEM(I) FORMAT(2) The 2 SAA 2 Processor and 2
3210	TURMAI (* INE ', 0A4, ' Pressures are:'
	$2^{\prime}$ u Decas $0^{\prime}$ $A2^{\prime}$ , $12^{\prime}$ mA? / $2^{\prime}$ Have these chanced?
	ACCEPT 1000 ANS
	$IE(ANS, NE, 2Y^2)$ G0 T0 3500
	TYPE *.'Has the ion gauge reading changed?'
	ACCEPT 1000, ANS
	IF(ANS.NE.'Ý') GO TO 3300
	TYPE *,'Enter the ion gauge mantissa.'
	ACCEPT +, PRM(I)
	TYPE *, 'Enter the ion gauge exponent.'
	ACCEPT *, IPRE(I)
3300	IYPE *, 'Has the foreline pressure changed?'
	ACCEPT 1000, AND TE (ANS NE 222) ON TO 2250
	TYPE + 'Enter the foreline procesure in storm (4) '
	ACCEPT + PEARE(I)
3350	TYPE 3360 NONOFF( $3-IDG(I)$ )
3360	FORMAT(' Has the degas been turned o'.A2.'?')
	ACCEPT 1000, ANS
	IF(ANS.EQ.'Y') IDG(I) = 3 - IDG(I)
	TYPE *, 'Has the emission current changed?'
	ACCEPT 1000, ANS
	IF (ANS.NE.'Y') GO TO 3500
	ITE *, 'Enter the emission current in mA.'
2500	
C	
c	CHECK ION PUMP CURRENTS
č	
	DO 3700 I=1,3
	TYPE 3610, I, PION(I)
3610	FORMAT(' The Region ',I1,' Ion Pump Current is ',F6.1
	1,'uA.',/,'Has this changed?')
	ACCEPT 1000, ANS
	TYPE 2000 T
3660	FIPE 3000,1 FIPEMAT(? What is the Posion ? If ? Ton Pump (unreast (in uA)??)
5000	ACCEPT * PINALTS the region , II, ION Fullip Current (In UA): )
3700	CONTINUE
9000	RETURN
	END

SUBROUTINE SCHNRM(IOPT) File SCHNRM.FOR begins on the previous line. С 000000000000 Version Date: July 19, 1984 Paul S. Weiss Subroutine SCHNRM changes one or all of the Normalization parameters in SANG. The argument IOPT specifies which parameters are to be changed as follows: IOPT=0 Change all Parameters IOPT=1 Change Normalization Angle (ANGNRM) Ĉ Change Normalization Countdown Interval (INTNRM) IOPT=2 С IOPT=3 Change Normalization Repetitions (NRMRPT) С COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG С IF(IOPT.GT.O) GO TO (1100,2100,3100), IOPT 1100 TYPE 1110, ANGNRM FORMAT(' The Normalization angle is ',F5.1,' degrees.' 1110 1,' Change this?') ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE. 'Y') GO TO 2000 TYPE +, 'Enter Normalization angle' ACCEPT +, ANGNRM 2000 IF(IOPT.NE.O) GO TO 9000 2100 TYPE 2110, INTNRM FORMAT(' The Normalization countdown interval is ', I5 2110 1,' Change this?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GD TO 3000 TYPE \*, 'Enter countdown interval for Normalization' 2300 ACCEPT +, INTNRM IF (INTNRM.LE.O) GO TO 2300 3000 IF(I0PT.NE.O) G0 T0 9000 TYPE 3110, NRMRPT 3100 3110 FORMAT(' The Normalization measurement is repeated ', I5 1,' times. Change this?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 9000 TYPE \*, 'Enter number of repetitions at each Normalization' 3300 ACCEPT +, NRMRPT IF (NRMRPT.LE.O) GO TO 3300 9000 RETURN END

## SUBROUTINE SCHPAR (IOPT)

File SCHPAR.FOR begins on the previous line.

Version Date: July 19, 1984 Paul S. Weiss

Subroutine SCHPAR changes one or all of the data acquisistion parameters in SANG.

The argument IOPT specifies which parameters are to be changed as follows:

IOPT=0 Change all Parameters IOPT=1 Change the Number of Channels in Use (NCHAN) IOPT=2 Change Clock Frequency (CLOCK) IOPT=3 Change Interval for Data Acquisition (INTACQ) IOPT=4 Change Gate Widths (GATE)

IOPT=5 Change Excited State Fractions (EXCFRC)

COMMON /TGSMAX/ DWLMAX,MAXCHN COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /CHADEF/ CHADEF(20,10)

IF (IOPT.GT.O) GO TO (1100,2100,3100,4100,5100), IOPT

CHANGE THE NUMBER OF DATA CHANNELS IN USE

C 1100 TYPE 1110, NCHAN

1110 FORMAT(1X,I1,' Data Channels are in use. Change this?') ACCEPT 1000,ANS 1000 FORMAT(A1)

IF (ANS.NE. 'Y') GO TO 2000

1200 TYPE 1210

1210 FORMAT(' How many Data Channels are to be used' 1,' (A channel consists of 2 scalers)?') ACCEPT + NCHAN

IF((NCHAN.LE.O).OR.(NCHAN.GT.MAXCHN)) GO TO 1200

2000 IF (IDPT.NE.0) GD TD 9000

C CHANGE CLOCK FREQUENCY

2100 TYPE 2110, CLOCK

2110 FORMAT(' The clock frequency is ',F3.1,' Hz. Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GO TO 3000

2300 TYPE \*, 'Enter clock frequency'

- ACCEPT +, CLOCK
- IF((CLOCK.NE.1.0).AND.(CLOCK.NE.1.5)) GO TO 2300

3000 IF(IOPT.NE.O) GO TO 9000

C

С

C

С

С

C C C C C C

Ċ

C C

C

С

Ը Շ

Ċ

С

С

С

C

C C

C	CHANGE ACQUISTION INTERVAL
3100 3110	TYPE 3110,INTACQ FORMAT(' The countdown interval is ',I5,' Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GD TO 4000
3300	TYPE *, 'Enter countdown interval for data acquisition' ACCEPT *, INTACQ IF(INTACQ.LE.0) GD TO 3300
4000 C	IF(IOPT.NE.O) GO TO 9000
C C	CHANGE GATE WIDTHS
4100 4110	TYPE 4110,GATE(1) FORMAT(' The gate widths are set at ',F5.2 1,' msec. Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GD TO 5000 TYPE *,'Enter the Gate Width for the A and B Channels' ACCEPT *,GATE(1)
5000 C	IF(I0PT.NE.O) GÓ TO 9000
C C	CHANGE EXCITED STATE FRACTIONS
5100	D0 5400 I=1,2 TYPE 5210,I,EXCFRC(I)
5210	FORMAT(' The fraction in Excited State #',I1,' is ',F6.4 1,' Change this?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GO TO 5400
5300 5310	TYPE 5310,I FORMAT(' Enter Excited State fraction #',I1)
5400	ACCEPT *,EXCFRC(I) IF((EXCFRC(I).LT.O.).OR.(EXCFRC(I).GT.1.)) G0 T0 5300 CONTINUE
9000	RETURN END
C C	SUBROUTINE SCHPOL(IOPT) File SCHPOL.FOR begins on the previous line.
C C C	Version Date: September 5, 1984 Paul S. Weiss
Č C C	Subroutine SCHPOL changes one or all of the Polarization Rotation Mode parameters in SANG.
c	The argument IOPT specifies which parameters are to be changed

С as follows: С Change all Parameters IOPT=0 С Change Polarization Rotation Detector Angle (ANGPOL) IOPT=1 С IOPT=2 Change Polarization Rotation Motor Step Size (POLSTP) С COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG С С С IF(IOPT.GT.O) GO TO (1100,2100), IOPT TYPE 1110, ANGPOL 1100 FORMAT(' The detector angle for the Polarization Rotation ' 1110 1, 'Loop is ',F5.1,' degrees.' 2,/,' Change this?') ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE.'Y') GO TO 2000 TYPE \*, 'Enter detector angle.' ACCEPT \*, ANGPOL 2000 IF(IOPT.NE.O) GO TO 9000 TYPE 2110, POLSTP, INT((180./POLSTP) + 0.5) 2100 FORMAT(' The motor step size is ', F6.3,' degrees.' 2110 1,/,' Rotating 180 degrees in ',I3,' steps.' 2,/,' Change this?') ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 9000 **TYPE 2210** 2200 2210 FORMAT(' Enter the number of steps in which to rotate the' 1, 'polarization 180 degrees.') ACCEPT +, NSTEP IF((NSTEP.LT.2).OR.(NSTEP.GT.200)) GO TO 2200 POLSTP = 180. / (1. \* NSTEP)9000 RETURN END SUBROUTINE SCHTST(IOPT) С File SCHTST.FOR begins on the previous line. С C Version Date: June 28, 1984 C Paul S. Weiss Č С Subroutine SCHTST changes one or all of the Test Mode С parameters in SANG. С С The argument IOPT specifies which parameters are to be changed С as follows: С IOPT=0 Change all Parameters С IOPT=1 Change Angle (TSTANG)

C	IOPT=2 Change Number of Countdowns (NUMTST)
ć	COMMON /TGSMAX/ DWLMAX,MAXCHN COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /TSTSET/ TSTANG,NUMTST,TSTCUR,TSTFLG
1100	IF(IOPT.GT.O) GO TO (1100,2100),IOPT TYPE *,'Enter angle for Test' ACCEPT *,TSTANG IE(IOPT NE O) CO TO 2000
2100 9000	TYPE *,'Enter number of countdowns for Test' ACCEPT *,NUMTST RETURN END
C	SUBROUTINE SCHVEL(IOPT) File SCHVEL.FOR begins on the previous line.
	Version Date: July 19, 1984 Paul S. Weiss
	Subroutine SCHVEL changes one or all of the Velocity Measurement parameters in SANG.
	The argument IOPT specifies which parameters are to be changed as follows: IOPT=0 Change all Parameters IOPT=1 Change Laser Frequency (FRQLAS) IOPT=2 Change Relative Frequency Standard (FSRLAS) IOPT=3 Change Velocity Measurement Step Size (LASSTP) IOPT=4 Change Velocity Measurement Accumulation Time (LASDWL)
C	COMMON /TGSMAX/ DWLMAX,MAXCHN COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /LASSET/ LASDWL,FRQLAS,LASSTP,FSRLAS,LASCHA
1100 1110	IF(IOPT.GT.O) GO TO (1100,2100,3100,4100),IOPT TYPE 1110,FRQLAS FORMAT(' The laser frequency is ',F11.4 1,' cm-1. Has this changed?')
1000	FORMAT(A1) TE(ANS NE 'Y') CO TO 2000
1300	TYPE *,'Enter Laser Frequency (in cm-1).' ACCEPT *,FRQLAS IF(FRQLAS_LE_0_) G0 T0 1300
2000	IF(IDPT.NE.O) GD TD 9000
2100 2110	FORMAT(' The Fabrey-Perot Free Spectral Range is ',F5.2 1,' GHz. Has this changed?')

	ACCEPT 1000, ANS
	IF(ANS.NE.'Y') GO TO 3000
2300	TYPE 2310
2310	FORMAT(' Enter the Relative Frequency Standard '
	1.'(FSR of the Fabrey-Perot etalon, in GHz).')
	ACCEPT * FSRLAS
	TE (ESRIAS LE 0) GO TO 2300
3000	
2100	TYPE + 'Enton Valacity Massurement Stan Size (1-4006)'
3100	ACCEDT - LASSTD
	$\frac{1}{10000000000000000000000000000000000$
	IF(LASSIF, LE, U)  GU  IU  SIUU
4000	IF(IUPT.NE.O) GU TU 9000
4100	IYPE *, 'Enter Velocity Measurement Accumulation lime (1-30000)'
	ACCEPT +, LASDWL
	IF(LASDWL.LE.O) GO TO 4100
9000	RETURN
	END
	SUBROUTINE SCMLST
С	File SCMLST.FOR begins on the previous line.
Č	· · · · · · · · · · · · · · · · · · ·
č	Version Date: August 15, 1985
č	Paul S. Weiss
č	
ĉ	Subroutine SCMLST lists the command groups for SANG. It provides
C C	the general help screep for SANC
	the general help screen for SANG.
C	TYPE 100
* ^ ^	TOPMAT(// ) The solution of the second second for the solution
100	FURMAI(//, The available commands are as follows:
	1,/,' C* Change Commands'
	2,/,' D* Define and Display Commands'
	3,/,' EX Exit'
	4,/,' G* Graph Commands'
	5,/,'H Help'
	6,/,' L* Loop Commands'
	7,/,' M* Machine Condition Commands'
	7,/,' N* Normalization Commands')
	TYPE 200 V
200	FORMAT(' P* Print Commands'
	1./.' R* Record Data, and Rotate Polarization Commands'
	2 / ' S* Show Commands'
	3 / ' T* Test Commands'
	A / ' V+ Velocity Measurement Commands'
	5/' W+ Write Commands'
	$5/7$ , $\pi^{+}$ $\pi^{-}$ (compands)
200	TOPWAT(/ ) To not a lighter of the vertices of the
300	rukmai(/, lo get a listing of the various options of
	1, a starred command, ',/,' enter an "H" in place of the "*"'
	2, (i.e. "PH" for a list of the Print commands).')
	RETURN

SUBROUTINE SCOMLC File SCOMLC.FOR begins on the previous line. Version Date: July 7, 1984 Paul S. Weiss Subroutine SCOMLC lists all the C\* commands relating to changing the fundamental parameters of the data acquisition. **TYPE 100** FORMAT(/,' The available Change Commands are as follows:' 1,/,' CA Change Angle' 2,/,; СС Change Clock Frequency' CD Change Definitions of Channels' 3,/, CE Change Excited State Fractions' 4,/, CF Change File' 5,/, 3 CG 6,/, Change Gate Widths') TYPE 200 FORMAT(' CH C Help' 1,/,' CI Change Interval for Acquisition Countdown' > CL Change Loop Parameters' 2,/, , 3,/, CN Change Normalization Parameters' ) CP 4,/, Change Data Acquisition Parameters' , CU Change Number of Channels Used' 5,/, C۷ Change Velocity Measurement Parameters') 6,/, RETURN END SUBROUTINE SCOMLD File SCOMLD.FOR begins on the previous line. Version Date: July 18, 1984 Paul S. Weiss Subroutine SCOMLD lists all the D\* commands relating to display and the cosmetic definitions of channels and machine hardware. **TYPE 100** 100 FORMAT(/,' The available Definition and ' 1, 'Display Commands are as follows:' 1,/,' 2,/,' 3./.' DA Display: Change All' DB Define Beams' DC Define Channels' 3,/, DH D Help' 4,/,

END

C

CCCCCC

С

100

200

С

CCCCCC

C

	5,/,' DI Display: Scale Channels Independently' 6,/,' DL Display: Change Lower Channel' 7,/,' DP Define Pressure Regions') TYPE 200
200	FORMAT(' DR Display: Refresh' 9./.' DS Display: Scale Channels to the Same Value'
	1,/,' DT Define Thermocouples'
	2,/,' DU Display: Change Upper Channel'
	3,/,' DW Display: Change Window')
	END
·	SUBROUTINE SCOMLG
C	File SCOMLG.FOR begins on the previous line.
C	Version Date: September 17, 1984
C	Paul S. Weiss
C	Subroutine SCOMIG lists all the G* commands relating to plotting
č	Subroutine Scomed 11303 and the dw commands relating to projering.
	TYPE 100
100	TURMAI(/,' The available Graph Commands are as follows:'
	2,/,' GH G Help'
	3,/,' GL Graph Lower Display'
	4,/,' GP Graph Upper Display') RETURN
	END
с	File SCOMLL.FOR begins on the previous line.
č	
C	Version Date: June 23, 1984
c	raul J. Weiss
Č · C	Subroutine SCOMLL lists all the L $\star$ commands relating to loop mode.
C	TYPE 100
100	FORMAT(/,' The available Loop Commands are as follows:'
	1,7,7 LA Loop: Change Current Angle?
	3./.' LC Loop: Continue.'
	4,/, 'LD Loop: Change Direction'
	5,/,' LE Loop: Change Ending Angle')
200	FORMAT(' IH Help'
<b>F</b> AA	1,/,' LI Loop: Change Angular Interval'
	2// IN Loop: Change Normalization Erequency'

Change All Loop Parameters' 3,/,' LP 4,/,' Loop: Start') LS RETURN END SUBROUTINE SCOMLM File SCOMLM.FOR begins on the previous line. Version Date: July 18, 1984 Paul S. Weiss Subroutine SCOMLM lists all the M\* commands having to do with the current machine conditions. **TYPE 100** FORMAT(/,' The available Machine Condition' 100 1,' Commands are as follows:' 1,/,' MA Machine Conditions: 1, 'Change All Detector Settings' 2,/,' MB Machine Conditions: 2, 'Change Backing Pressures' 3,/, <sup>\*</sup> MD Machine Conditions: 3, 'Change Daly Detector (post-QPMS)') **TYPE 200** 200 FORMAT(' MF Machine Conditions: 3, 'Change Filament/Ionizer Settings' 4,/,' MH Machine Conditions: Help' 5,/,' MI Machine Conditions: Change Ion Lens Settings' MM Machine Conditions: 6, 'Change Mass Spectrometer Settings' 7,/,' 8,/,' MR Machine Conditions: Change Regional Pressures' Machine Conditions: MS 9,'Change Source Conditions' 1,/,' MT Machine Conditions: 1, 'Change Thermocouple Readings') RETURN END SUBROUTINE SCOMLN File SCOMLN.FOR begins on the previous line. Version Date: June 23, 1984 Paul S. Weiss Subroutine SCOMLN lists all the N\* commands relating to normalization mode of SANG.

CCCCCCCCCC

CCCCCCC

С

1,/,'	(/,'The NA NH	available Normalization Commands are as follows:' Change Normalization Angle' N Help'
3,/,'	NP	Change Normalization Countdown Interval' Change All Normalization Parameters'
5,/,'	NR	Change Normalization Repetitions'
6,/,'	NS	Normalization: Start')
RETURN		· · ·
SUBRUUT	INE SCU	MLP File SCOMLP.FOR begins on the previous I
	•	
Version Paul S	1 Date: Weiss	July 24, 1984
Subrout	ine SCO ation on	VLP lists all the P* commands relating to printing the line printer.
TYPE 10	00	
FORMAT	(/,' The	available Print Commands are as follows:'
1,/,'	PA	Print Accumulated Data'
2,/,'	PC	Print Current Configuration'
3,/,'	PD	Print Data Channel Definitions'
5,/,'	PG	Print Current Graphics Display'
6,/,'	PH	P Help'
7,/,'	PI	Print Information on the Line Printer')
TYPE 20	)0	
FORMAT	('	PL Print Loop Parameters'
9,/,'	PM	Print Machine Conditions'
1 / )	PN	Print Normalization Parameters'
±,/,	PR	Print Rotation of Polarization Parameters'
2,/,'	PS	Print Detector Settings'
1,/, 2,/,' 3,/,'		Print Test Mode Parameters'
1,/, 2,/,' 3,/,' 4,/,'	PT	
1,/, 2,/,' 3,/,' 4,/,' 5,/,'	PT PV	Print Velocity Parameters and Measurement.')
1,/, 2,/,' 3,/,' 4,/,' 5,/,' RETURN	PT PV	Print Velocity Parameters and Measurement.')
2,/,' 3,/,' 4,/,' 5,/,' RETURN END	PT PV	Print Velocity Parameters and Measurement.')
1,/, 2,/, 3,/, 4,/, 5,/, RETURN END SUBROUT	PT PV	Print Velocity Parameters and Measurement.') MLR
2,/,' 3,/,' 4,/,' 5,/,' RETURN END SUBROUT	PT PV	Print Velocity Parameters and Measurement.') MLR File SCOMLR.FOR begins on the previous !
2,/,' 3,/,' 4,/,' 5,/,' RETURN END SUBROUT	PT PV FINE SCO	Print Velocity Parameters and Measurement.') MLR File SCOMLR.FOR begins on the previous 1 July 24, 1984
2,/,' 3,/,' 4,/,' 5,/,' RETURN END SUBROUT Version Paul S	PT PV FINE SCO Date: Weiss	Print Velocity Parameters and Measurement.') MLR File SCOMLR.FOR begins on the previous I July 24, 1984
2,/,' 3,/,' 4,/,' 5,/,' RETURN END SUBROUT Version Paul S. Subrout	PT PV FINE SCO n Date: . Weiss tine SCO	Print Velocity Parameters and Measurement.') MLR File SCOMLR.FOR begins on the previous 1 July 24, 1984 MLR lists all the R* commands relating to polariza

**TYPE** 100

100

FORMAT(/,' The available Record Data and Rotation of' 1,' Polarization Commands are as follows:' RA Rotation of Polarization: Change Angle' 1,/, 2,/, RD Record Data' RH .R Help' 3,/, RI Reinitialize TGS System' 4,/, , RM 5,/, Rotation of Polarization: Change Motor' 5,' Step Size' 5,/,' RP Rotation of Polarization: Change All ' 5,'Parameters' 6,/,' RS Rotation of Polarization: Start') RETURN END

SUBROUTINE SCOMLS

ŚA

SC

SD

SG

File SCOMLS begins on the previous line.

July 24, 1984 Version Date: Paul S. Weiss

Subroutine SCOMLS lists all the S\* commands relating to typing information on the console terminal.

FORMAT(/,' The available Show Commands are as follows:'

Show Current Configuration'

Show Data Channel Definitions'

Show Current Graphics Display')

Show Accumulated Data'

**TYPE 100** 

1,/,'

2,/,;

3,/,

END

200

С

C C

C

100

С

CCCCCCC

5,/,' TYPE 200 FORMAT(' SH S Help' 1,/,' Show Time Information' SI 2,/,, SL Show Loop Mode Parameters' 9,/, SM Show Machine Conditions' , SN Show Normalization Parameters' 3,/, Show Rotation of Polarization Parameters' SR SS 3,/, Show Detector Settings' 5,/, ST Show Test Mode Parameters' S٧ Show Velocity Parameters and Measurement') 6,/, RETURN

SUBROUTINE SCOMLT

File SCOMLT.FOR begins on the previous line.

Version Date: June 23, 1984 Paul S. Weiss

Subroutine SCOMLT lists all the T\* commands relating to test mode of SANG.

Č C

С С

100

**TYPE 100** 

FORMAT (	/,'	The available Test Mode Commands are as follows:'
1,/,'	TA	Change Test Angle'
2,/,'	TC	Test Mode: Continue'
3,/,	ΤE	Test Mode: Exit'
4,/,	TH	T Help'
5,/,	TN	Test Mode: Change Number of Countdowns'
6,/,	TP	Change All Test Mode Parameters'
7,/,	TS	Test Mode: Start')
RETURN		
END		

SUBROUTINE SCOMLV

File SCOMLV.FOR begins on the previous line.

Version Date: June 23, 1984 Paul S. Weiss

Subroutine SCOMLV lists all the V\* commands relating to the measurement of a beam velocity by measuring the Doppler shift of a transition.

100

С

C C

С

C C

С

С

Ċ

С

000000

Ċ

**TYPE 100** FORMAT(/,' The available Velocity Measurement Commands' 1,'are as follows:' 2,/,'VA Chan Change Velocity Measurement Accumulation Time' 3,/,; ٧F Change Velocity Measurement Frequency Standard' 4,/, V Help' VH , ٧L Change Velocity Measurement Laser Frequency' 5,/, VM Measure Velocity' 6,/, VP 7,/, Change All Velocity Measurement Parameters' ) ٧S Change Velocity Measurement Step Size')

8,/,' VS Change Velocity Measurement S RETURN

END

SUBROUTINE SCOMLW

File SCOMLW.FOR begins on the previous line.

Version Date: September 3, 1984 Paul S. Weiss

Subroutine SCOMLW lists all the W\* commands relating to writing information to the data file.

**TYPE 100** 100 FORMAT(/,' The available Write to Disk Commands' 1,' are as follows:' ) 1,/, WA Write Accumulated Data' > WC 2,/, Write Current Configuration' `, 3,/, ₩D Write Data Channel Definitions' , ₩G Write Current Graphics Display' 5,/, ) WH 6,/, W Help' , 7,/, WI Write and Print Information') TYPE 200 FORMAT(' WL Write Loop Mode Parameters' 9,/,' WM Write Machine Conditions' 1,/,, WN Write Normalization Parameters' 2,/, WR Write Rotation of Polarization Parameters' ) **WS** 3,/, Write Detector Settings' WT Write Test Mode Parameters' 4,/, • 5,/, WV Write Velocity Parameters and Measurement.') RETURN END SUBROUTINE SCOMLZ File SCOMLZ.FOR begins on the previous line. Version Date: June 23, 1984 Paul S. Weiss Subroutine SCOMLZ lists all the Z\* commands relating to zeroing all or part of the accumulated data. С **TYPE 100** FORMAT(/,' The available Zero Commands are as follows:' 100 1,/,' ZA Zero Arrays' , 2,/, ZH Z Help' ` > ZS 3,/, Zero Angular Segment' `> ZT Zero Test Mode Data') 4,/, RETURN

SUBROUTINE SDISP

END

File SDISP.FOR begins on the previous line.

Version Date: September 25, 1984 Paul S. Weiss

Subroutine SDISP loads the display array for the DT2771 DMA display driver. The current point is brightened by

С

200

С

С

С

C۰ a factor of 10X. C C C C C WARNING: Subroutines SDISP and SDVAL make up the largest segment of overlay region 1. Enlarging either of them will directly increase the total size of SANG. С COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB COMMON /DSPZER/ ITYZER, IBYZER COMMON /SCR/ SCR(20) COMMON /DSCR/ DSCR(20) COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG С С FIRST STOP DISPLAY С CALL STOPD NBRITE = 0DO 100 I=1, NDSP IT = 2 \* I - 1IB = (2 \* NDSP) + ITIDSP(IB) = 16 \* I - 8IDSP(IT) = 16 \* I - 8IDSP(IB + 1) = 0IDSP(IT + 1) = 2040100 CONTINUE D0 200 I=1,256 WORK1(I) = 0.WORK2(I) = 0.200 CONTINUE С C FIRST DETERMINE THE FIRST ANGLE TO BE DISPLAYED С DO 1300 IBE=1, NUMANG IF(DATA(IBE,1).GE.DSPB) GO TO 1500 1300 CONTINUE C С INITIALIZE MAXIMA AND MINIMA С 1500 TMAX = 0.TMIN = 0.BMAX = 0.BMIN = 0.С С LOAD NEW DATA INTO SCRATCH ARRAY FOR DISPLAY LOADING С DO 1600 I=2,20 DSCR(I) = SCR(I)1600 CONTINUE C.

#### DETERMINE VALUES OF NEW DATA С С CALL SDVAL (MODSPT, MDSPT, NDSPT, TVAL, TMAX, TMIN) CALL SDVAL (MODSPB, MDSPB, NDSPB, BVAL, BMAX, BMIN) INTIALIZE THE MINIMUM AND MAXIMUM OF DISPLAY Y VALUES TO THE NEW DATA VALUES THE SUBROUTINE SOVAL WILL CONTINUE TO CHECK FOR MINIMA AND MAXIMA AS IT LOADS THE VALUES OF EACH DATA POINT DETERMINE WHICH ANGLES ARE DISPLAYED WHERE, NOTE THAT IF MORE THAN ONE DATA POINT WOULD BE DISPLAYED AT A DISPLAY POINT, THESE ARE COMBINED, AS IF TAKEN AT EXACTLY THE SAME ANGLE. THE DISPLAY RESOLUTION IS DETERMINED BY THE VARIABLE DSTEP D0 2500 I=1,NDSP IF (IBE-NUMANG.GT.O) GO TO 3000 NDANG = 0DO 2000 J=IBE, NUMANG IF(DATA(J,1).GT.(DSPB+(I\*DSTEP))) G0 T0 2100 NDANG = NDANG + 1CONTINUE 2000 С С IF THERE IS NO DATA FOR THIS DISPLAY POINT, GO ON TO NEXT С 2100 IF (NDANG.EQ.O) GO TO 2500 С С LOAD TOTAL DATA FOR THIS DISPLAY POINT INTO DSCR D0 2300 K=2,20 DSCR(K) = 0.D0 2200 J=1,NDANG DSCR(K) = DSCR(K) + DATA(IBE+J-1,K)2200 CONTINUE 2300 CONTINUE С С SDVAL DETERMINES THE VALUE OF THE DISPLAY POINT С CALL SDVAL (MODSPT, MDSPT, NDSPT, WORK1 (I), TMAX, TMIN) CALL SDVAL (MODSPB, MDSPB, NDSPB, WORK2(I), BMAX, BMIN) С C ADVANCE THE POINT AT WHICH THE DATA IS FIRST С CHECKED TO SEE IF A POINT IS TO BE DISPLAYED С IBE = IBE + NDANG2500 CONTINUE С

IF CHANNELS ARE TO SCALED TO THE SAME VALUE, FIND MAXIMUM AND MINIMUM

С С

С

Ċ С

С

С С

С

С

С С

С

С

С

С 3000 IF(INDEP.NE.O) GO TO 3500 TMAX = AMAX1(TMAX, BMAX)TMIN = AMIN1(TMIN, BMIN)BMAX = TMAXBMIN = TMIN С С DETERMINE THE Y RANGE OF THE UPPER AND LOWER DISPLAYS С 3500 TRANGE = (TMAX - TMIN) / 2000.BRANGE = (BMAX - BMIN) / 2000.С Ċ CONVERT TO 12 BIT INTEGERS FOR DISPLAY D0 5000 I=1,NDSP IF (TRANGE.LE.O.) GD TO 5000 IDSP(2\*I) = INT((WORK1(I) - TMIN) / TRANGE) + 2040IF (BRANGE.LE.O.) GD TO 5000 IDSP(2\*(I+NDSP)) = INT((WORK2(I) - BMIN) / BRANGE) 5000 CONTINUE С Ĉ LOAD BRIGHTENED POINT С ITYSCR = 2040IBYSCR = 0XSCR = (SCR(1) - DSPB)/DSTEPС C C SEE IF THE CURRENT DATA FITS ON THE DISPLAY WINDOW IF IT DOES NOT, ONLY DISPLAY ARRAYS (I.E. LEAVE NBRITE = 0) С IF((XSCR.LT.O.).OR.(XSCR.GT.256.)) GD TO 8500 NBRITE = 20IXSCR = INT(16.\*XSCR) - 8IF (TRANGE.GT.O.) ITYSCR = INT ((TVAL - TMIN)/TRANGE) + 2040 IF (BRANGE.GT.O.) IBYSCR = INT((BVAL - BMIN)/BRANGE) D0 6000 I=1,10 NDY = (4 \* NDSP) + (2 \* I)IDSP(NDY-1) = IXSCRIDSP(NDY) = ITYSCRIDSP(NDY + 19) = IXSCRIDSP(NDY + 20) = IBYSCR6000 CONTINUE С С CALL MACRO TO DO DISPLAY Ċ 8500 CALL DISPL(IDSP,((2\*NDSP)+NBRITE)) С С FIND ZERO LEVELS FOR TOP AND BOTTOM IN ORDER TO PLOT DISPLAYS C IF CALLED FOR С

ITYZER = 2040IBYZER = 0IF (TRANGE.GT.O.) ITYZER = INT ( - TMIN / TRANGE) + 2040 IF (BRANGE.GT.O.) IBYZER = INT( - BMIN / BRANGE) 9000 RETURN END SUBROUTINE SDVAL (MODD, MDS, NDS, VAL, VMAX, VMIN) С Second routine in file SDISP.FOR. CCCCCCCCC Version Date: July 27, 1984 Paul S. Weiss Subroutine SDVAL determines the correct display value for the cuurent point given the display. It then determines the new minimum and maximum. COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /DSCR/ DSCR(20) MS = (2 \* MDS) + 1MB = MS + 1NS = (2\*NDS) + 1NB = NS + 1VAL = 0.IF(DSCR(2).LT.1.E-5) GO TO 2000 GO TO (1100,1200,1300,1400,1500,1600),MODD VAL = (DSCR(MS) - DSCR(MB)) / DSCR(2)1100 GO TO 2000 VAL = (DSCR(MS) - DSCR(MB) - DSCR(NS) + DSCR(NB)) / DSCR(2)1200 GO TO 2000 1300 X = 1.IF(EXCFRC(1).NE.0.) X = EXCFRC(1)VAL = (DSCR(MS) - DSCR(MB) - ((1. - X))1 + (DSCR(NS) - DSCR(NB))) / (DSCR(2) + X)GO TO 2000 1400 VD = SQRT(DSCR(MS) + DSCR(MB) + DSCR(NS) + DSCR(NB))IF(VD.EQ.O.) GO TO 2000 VAL = (DSCR(MS) - DSCR(MB) - DSCR(NS) + DSCR(NB)) / VDGO TO 2000 1500 VD = SQRT(DSCR(MS) + DSCR(MB))IF(VD.EQ.0.) G0 T0 2000 VAL = (DSCR(MS) - DSCR(MB)) / VDGO TO 2000 1600 VAL = DSCR(2)IF(VAL.GT.VMAX) VMAX = VAL2000 IF (VAL.LT.VMIN) VMIN = VAL RETURN

SUBROUTINE SDLONO (ISTART) С File SDLOND.FOR begins on the previous line. С С September 27, 1984 Version Date: С Paul S. Weiss С С Subroutine SDLONO is called during Loop mode after a normalization. С С COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SCR/ SCR(20) COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG C IF (NRMFLG.NE.O) GO TO 9000 ISTART = 1GO TO (1100,2100),-LPFLAG C C CONTINUE LOOP C 1100 TYPE \*, 'End data taking loop?' ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS .EQ. 'Y') GO TO 8900 ANLPCU = ANLPCU + ANLPINС С HAS THE END OF A SCAN BEEN REACHED? С IF ( (ANLPCU.GT.AMAX1 (ANLPBE, ANLPFN) ) . OR . (ANLPCU 1.LT.AMIN1(ANLPBE,ANLPFN))) GO TO 7000 GO TO 6500 С С IF LOOP HAS JUST BEGUN, START AT BEGINNING OF LOOP С (AFTER NORMALIZATION) С 2100 ANLPCU = ANLPBEС С SET UP ANGLE FOR ACQUISITION C 6500 LPFLAG = 1ANGLE = ANLPCUINTERV = INTACQGD TD 9000 С C REVERSE DIRECTION IF THE END OF A SCAN HAS BEEN REACHED

END

C	
7000	A = ANLPFN
	ANLPFN = ANLPBE
	ANI PRE = A
	ANDERCO = ANDERDE
	TTPE 7010
7010	FURMAI(/,' The angular scan has been completed.'
	1,/,' Rescan in the opposite direction from ',F5.1,' to '
	2,F5.1,' degrees,',/,' in ',F5.1,' degree steps?')
	ACCEPT 1000, ANS
	IF (ANS, EQ, 'Y') GO TO 6500
C	
č	NO NOT RESTART TE LOOP IS NOT TO BE CONTINUED
8900	
	LPFLAG = O
9000	RETURN
	END
	•
	SUBROUTINE SDLOP(NKEEP.ISTART)
C	File SDLOP FOR begins on the previous line
č	
Ċ	Varian Data: Santarban 27 1084
Č	Paul C. Waine
	raul J. Weiss
C	
C	Subroutine SULUP decides whether or not to keep a data point,
C	and continue the data taking loop.
C	
C	
	COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG
	COMMON /SCR/ SCR(20)
	COMMON /SETS/ GATE(16).TIM(16).EXCFRC(2).CLOCK.INTACO.NCHAN
	COMMON /LOPSET/ ANI PRE ANI PEN ANI PIN I PNRED ANI PCU I PELAG
	COMMON /NRMSET/ ANGNEM INTNEM NEMEPT NEMELG
C	
L	NKEED - 1
	INCLE - I
100	I = I
100	117E 110, 1, 2
110	FURMAI(' Save this point', Al,:
	1,A1,'nd continue data taking loop?')
	ACCEPT 1000, ANS
1000	FORMAT(A1)
	IF(ANS.EQ.'Y') GO TO 5000
	TYPE 110, '?'
•	ACCEPT 1000 ANS
	TE (ANS EQ. 'Y') GO TO 3000
	$\mathbf{M}$

TYPE \*, 'Repeat this point?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') GD TO 9000 TYPE +, 'Change detector angle?' ACCEPT 1000, ANS IF(ANS.NE.'Y') GO TO 3000 TYPE \*, 'Enter correct angle.' ACCEPT \*, SCR(1) NKEEP = 13000 TYPE \*, 'End data taking loop?' ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 100 GO TO 8900 C С CONTINUE LOOP С С IS IT TIME FOR A NORMALIZATION? С IF SO, SET DETECTOR AND INTERVAL ACCORDINGLY С 5000 . IF (LPFLAG.LT.LPNRFQ) GO TO 6000 NRMFLG = 1ANGLE = ANGNRMINTERV = INTNRM LPFLAG = -1**TYPE 5210** FORMAT(/, ' Normalization!',/) 5210 GO TO 9000 С С INCREMENT COUNTER AND ANGLE FOR NEXT ANGLE IN LOOP С 6000 LPFLAG = LPFLAG + 1ANLPCU = ANLPCU + ANLPIN С С HAS THE END OF A SCAN BEEN REACHED? C IF ( (ANLPCU.GT. AMAX1 (ANLPBE, ANLPFN) ) . OR. (ANLPCU 1.LT.AMIN1 (ANLPBE, ANLPFN))) GO TO 7000 С С SET UP NEXT ANGLE С 6500 ANGLE = ANLPCUINTERV = INTACQGO TO 9000 С С REVERSE DIRECTION IF THE END OF A SCAN HAS BEEN REACHED С 7000 A = ANLPFNANLPFN = ANLPBEANLPBE = AANLPIN = - ANLPIN

7010 C C 8900 9000	ANLPCU = ANLPBE TYPE 7010, ANLPBE, ANLPFN, ANLPIN, 7, 7, 7, 7, 7 FORMAT(/,' The angular scan has been completed.' 1,' Rescan in the opposite direction',/,' from ',F5.1,' to ' 2,F5.1,' degrees, in ',F5.1,' degree steps?',5A1) ACCEPT 1000, ANS IF(ANS.EQ.'Y') GD TO 6500 D0 NOT RESTART IF LOOP IS NOT TO BE CONTINUED ISTART = 0 LPFLAG = 0 RETURN END
	SUBROUTINE SDNRM(NKEEP,ISTART) File SDNRM.FOR begins on the previous line. Version Date: July 19, 1984
	Faul 5. Weiss Subrouting SDNRM decides whether or not to keep a data point. It is called during a normalization.
c	COMMON /CUR/ ANGLE,INTERV,NUMANG,MANGLE,IRNFLG COMMON /NRMSET/ ANGNRM,INTNRM,NRMRPT,NRMFLG
1000	ISTART = -1 NKEEP = 1 TYPE *,'Save this normalization point?' ACCEPT 1000,ANS FORMAT(A1)
1210	IF(ANS.EQ.'Y') GO TO 2000 NKEEP = O TYPE 1210 FORMAT(/,' Abort this normalization point.' 1,/,' Measure it again?')
C C C	ACCEPT 1000,ANS IF POINT NEEDS TO BE RETAKEN, DO NOT INCREMENT NORMALIZATION FLAG, AND QUERY THAT DETECTOR IS READY
C	IF(ANS.NE.'Y') GO TO 2000 ISTART = 1 GO TO 9000
C	TO CONTINUE NORMALIZATION, INCREMENT REPETITION COUNTER, AND

С SEE IF THE CORRECT NUMBER OF REPETITIONS HAVE BEEN DONE. С 2000 IF (NRMFLG.GE.NRMRPT) G0 T0 7000 NRMFLG = NRMFLG + 1GO TO 9000 С С IF THE NUMBER OF REPETITIONS REQUESTED HAVE BEEN COMPLETED, С STOP NORMALIZING, AND SEE IF ANOTHER NORMALIZATION IS DESIRED С 7000 NRMFLG = 0ISTART = 08000 TYPE +,'End of Normalization. Normalize again?' ACCEPT 1000, ANS IF(ANS.NE.'Y') G0 T0 9000 ISTART = 1NRMFLG = 19000 RETURN END SUBROUTINE SDONE С File SDONE.FOR begins on the previous line. C С Version Date: October 30, 1984 С Paul S. Weiss С С Subroutine SDONE outputs the scratch data from the scalers С (using SOUT), displays the data (using SDISP), and calls the С appropriate routines depending upon how the data was taken С (SDLOP for Loop Mode, SDNRM for Normalization, SDPOL for Rotation) С of Polarization, SDTST for Test Mode, and SDSIN for single С angle acquisition). С С INTEGER+4 JDATA COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG

COMMON /DATA/ DATA(256,20),WORK1(256),WORK2(256) COMMON /SCR/ SCR(20) COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /LASSET/ LASDWL,FRQLAS,LASSTP,FSRLAS,LASCHA COMMON /LOPSET/ ANLPBE,ANLPFN,ANLPIN,LPNRFQ,ANLPCU,LPFLAG COMMON /NRMSET/ ANGNRM,INTNRM,NRMRPT,NRMFLG COMMON /POLSET/ ANGPOL,POLSTP,POSIT,IPOSIT,IPLFLG COMMON /TSTSET/ TSTANG,NUMTST,TSTCUR,TSTFLG COMMON /SDONE/ ISWFLG,JDATA(20)

LEAVE SPECIAL SCREEN MODE

CALL EXSPEC

с • с с
C C IF AN OVERFLOW HAS OCCURRED, HANDLE IT IN ROUTINE SOVER Ċ IF(ISWFLG.LT.2) GO TO 200 CALL SOVER (ISTART) GO TO 8000 C Č IF THERE IS NEW DATA, GET IT С 200 IF(ISWFLG.EQ.1) CALL SGTSCR С С DISPLAY NEW DATA С CALL SDISP С C RESET CALLING FLAG, AND RETURN IF ONLY CALLED FOR DISPLAY PURPOSES С I = ISWFLGISWFLG = -1IF(I.EQ.0) GO TO 8900 С С OUTPUT NEW DATA ON TERMINAL С CALL SOUT (7) С С DETERMINE WHAT MODE THE DATA WAS TAKEN IN AND CALL THE APPROPRIATE С ROUTINES C NKEEP = 0ISTART = 0C C C ACCUMULATE FLAGS' STATUS SO THAT THIS CAN BE WRITTEN TO DISK AS ONE WORD, IFLAG С IFLAG = 0IF(NRMFLG.NE.O) IFLAG = 1 $IF(LPFLAG_NE.0)$  IFLAG = IFLAG + 10IF(IPLFLG.NE.O) IFLAG = IFLAG + 100 IF(TSTFLG.NE.O.) IFLAG = IFLAG + 1000 С С TEST MODE? ' С IF(TSTFLG.EQ.O.) GO TO 2100 CALL SDTST (NKEEP, ISTART) GO TO 7000 С С NORMALIZING? С IF(NRMFLG.EQ.O) GO TO 3100 2100 CALL SDNRM (NKEEP, ISTART)

Ը Շ IF NORMALIZATION IS IN LOOP, ALSO CALL SDLONO С IF (LPFLAG.NE.O) CALL SDLONO (ISTART) GO TO 7000 С С LOOP MODE, BUT NOT CURRENTLY NORMALIZING? С 3100 IF(LPFLAG.EQ.O) GO TO 4100 CALL SDLOP (NKEEP, ISTART) GO TO 7000 С С ROTATION OF POLARIZATION? С 4100 IF(IPLFLG.EQ.0) GO TO 5100 CALL SDPOL (NKEEP, ISTART) GO TO 7000 С С SINGLE POINT ACQUISITION? С 5100 CALL SDSIN(NKEEP) 7000 IF (NKEEP.EQ.0) GO TO 8000 С С IF DATA IS SATISFACTORY, WRITE TO LINE PRINTER AND DATA FILE С CALL SSAVE CALL SOUT(6) CALL SWRITE(IFLAG) 8000 IF(ISTART.NE.O) CALL SSTART(ISTART) **TYPE 8010** 8010 FORMAT(' Command?',/) С С REENTER SPECIAL SCREEN MODE С 8900 CALL ENSPEC RETURN END SUBROUTINE SDPOL (NKEEP, ISTART) С File SDPOL.FOR begins on the previous line. С С September 25, 1984 Version Date: С Paul S. Weiss С С Subroutine SDPOL decides whether or not to keep a data point, С and continue the polarization rotation data taking loop. C С

COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG С POLCUR = POSIT + 180./200.NKEEP = 1ISTART = 1INTERV = INTACQ100 TYPE 110,' ','a' FORMAT(' Save this point', A1, : 110 1,A1, 'nd continue polarization rotation loop?') ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.EQ.'Y') GO TO 6000 TYPE 110, '?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') GO TO 3000 NKEEP = 0TYPE \*, 'Repeat this point?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') GO TO 8900 3000 TYPE +, 'End polarization rotation loop?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') GO TO 8800 GO TO 100 С С CONTINUE POLARIZATION ROTATION LOOP С С INCREMENT COUNTER FOR NEXT POLARIZATION IN LOOP С 6000 STEP = POLSTP + 200. / 180.POSIT = POSIT + STEP С С HAS THE END OF A SCAN BEEN REACHED? C IF(POSIT.LT.O.) GO TO 7000 C С HAS 180 DEGREES BEEN REACHED? С IF (POSIT.GT.200.) GO TO 8000 С С IF IN THE MIDDLE OF A SCAN, MOVE ROTATOR, AND INCREMENT/DECREMENT С POLARIZATION ANGLE С MOVE = INT(POSIT + 0.5) - IPOSITIF(MOVE.LT.O) MOVE = MOVE + 800IPOSIT = INT(POSIT + 0.5)С С MACRO MOT MOVES THE POLARIZATION ROTATOR С

CALL MOT(-MOVE) ANGLE = ANGLE + (POLSTP + 1.E-4)GO TO 8900 С С IF THE SCAN HAS REACHED O DEGREES, QUERY USER AS TO WHETHER OR NOT С TO CONTINUE С POSIT = 0.7000 IPOSIT = 0ANGLE = ANGPOLIF (POLSTP.LT.O.) POLSTP = - POLSTP TYPE 7110,7,7,7,7,7 7110 FORMAT(/,' The polarization rotation scan has been completed.' 1,/,' Rescan from 0 to 180 degrees and back?', 5A1) ACCEPT 1000, ANS С С DO NOT RESTART IF LOOP IS NOT TO BE CONTINUED C IF(ANS.NE.'Y') GO TO 8800 GO TO 8900 С С IF THE SCAN HAS REACHED 180 DEGREES, REPEAT 180, BUT REVERSE STEP С DIRECTION С 8000 POSIT = 200.IPOSIT = 200POLSTP = - POLSTPGO TO 8900 С С IF THE LOOP IS TO BE CANCELLED, SET POL FLAG OFF, С AND RESTART FLAG OFF С 8800 TYPE \*, 'Abort further Rotation of Polarization Measurements' С С MOVE ROTATOR BACK TO VERTICAL С CALL MOT (IPOSIT-800) POSIT = 0.IPOSIT = 0IPLFLG = 0ISTART = 0IF(ISTART.NE.0) TYPE 8910, (POSIT+180./200.) 8900 8910 FORMAT(' The laser polarization is ', F7.3 1,' degrees from vertical.') RETURN END SUBROUTINE SDSIN(NKEEP) C C File SDSIN.FOR begins on the previous line.

CCCCCCC

Version Date:

Paul S. Weiss

July 1, 1984

Subroutine SDSIN decides whether or not to keep a data point. It is called when in single point acquisition mode. COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /SCR/ SCR(20) COMMON /SETS/ GATE (16), TIM (16), EXCFRC (2), CLOCK, INTACQ, NCHAN COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG NKEEP = 1TYPE \*, 'Save this data?' ACCEPT 1000,ANS FORMAT(A1) IF(ANS.NE.'Y') NKEEP = 0RETURN END SUBROUTINE SDTST(NKEEP, ISTART) File SDTST.FOR begins on the previous line. Version Date: July 24, 1984 Paul S. Weiss Subroutine SDTST continues test mode unless TSTFLG has been made 0. COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /SCR/ SCR(20) COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG ISTART = -1NKEEP = 1IRNFLG = 0INTERV = INTACQ

1000

С

0000000000

С

С

C C C 2000 2010	CONTINUE TEST MODE?
	IF(TSTFLG.LE.O.) GO TO 2000
	DONE WITH CYCLES?
	IF(TSTFLG.LT.NUMTST) GO TO 4000 TSTFLG = 0. ISTART = 0 TYPE 2010,7,7,7,7,7 FORMAT(/,' End of Test',5A1) GO TO 9000
C	ADD .00001 TO ANGLE AND CONTINUE TEST
4000 9000	TSTFLG = TSTFLG + 1. TSTCUR = TSTCUR + 00001 ANGLE = TSTANG + TSTCUR RETURN END
C C	SUBROUTINE SFILE(IOPT) File SFILE.FOR begins on the previous line.
C C	Version Date: July 7, 1984 Paul S. Weiss
	Subroutine SFILE opens a data file on unit 9 for output of data in ASCII form.
c	LOGICAL*1 BNAME(14) COMMON /FNAME/ INAME(7),ISTORE EQUIVALENCE (BNAME(1),INAME(1)) DATA INAME /'DL','O:','AN','GD','AT','.A','NG'/
C C C 100 110 500	IF(IOPT.EQ.O) CLOSE(UNIT=9,DISPOSE='SAVE')
	RESET THE NUMBER OF DATA POINTS STORED
	<pre>ISTORE = 0 TYPE 110 FORMAT(' Enter the six letter data file name:') ACCEPT 500, (INAME(I), I=3,5) FORMAT(3A2) INAME(6) = '.A' INAME(7) = 'NG'</pre>

C C C C C C IF THE FIRST LETTER IS A SPACE, OR ANY CHARACTER OTHER THAN A LETTER GO BACK AND GET THE CORRECT FILENAME IF(BNAME(5).LT.64) GO TO 100 С С IF THERE ARE FEWER THAN SIX LETTERS, SHIFT THE EXTENSION DOWN С LNAME = 14DO 600 I=1,6 B = BNAME(11-I)C C C C IF THE CHARACTER IN QUESTION IS A LETTER OR NUMBER, SAVE IT OTHERWISE ELIMINATE IF((B.GT.64).OR. ((B.GT.47).AND.(B.LT.58))) GO TO 700 С С. KEEP TRACK OF NAME LENGTH С LNAME = 14 - IDO 550 J=11-I,14-I BNAME(J) = BNAME(J+1)550 CONTINUE BNAME(15-I) = 0600 CONTINUE С С CHECK AND SEE IF FILE ALREADY EXISTS С 700 OPEN(UNIT=9, ERR=2000, NAME=BNAME, TYPE='OLD') С С IF IT EXISTS SEE IF IT SHOULD BE WRITTEN OVER С CLOSE(UNIT=9,DISPOSE='SAVE') С С FILL IN NULLS IN DATA FILE NAME WITH BLANKS FOR OUTPUT Č IF(LNAME.EQ.14) GO TO 800 D0 750 I=LNAME+1,14 BNAME(I) = 32750 CONTINUE TYPE 810, (BNAME(I), I=1, LNAME), 7, 7, 7, 7, 7, 0 800 FORMAT(1X,14A1,' already exists! 810 1,'Would you like to write over it?',6A1) ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE. 'Y') GO TO 100 2000 CALL ASSIGN (9, INAME, LNAME, 'NEW', 'CC') С С FILL IN NULLS IN DATA FILE NAME WITH BLANKS FOR OUTPUT

С IF (LNAME. EQ. 14) GO TO 2200 DO 2150 I=LNAME+1,14 BNAME(I) = 322150 CONTINUE 2200 FORMAT(' The data file: ',14A1 2210 1,' has been opened successfully!',11A1) RETURN END SUBROUTINE SGRAPH(IOPT) С File SGRAPH.FOR begins on the previous line. С С Version Date: November 11, 1985 С Paul S. Weiss С С Subroutine SGRAPH plots the angular distribution on the С HP 7470A or HP 7475A plotter which is defined to be on unit 8. С С The argument is: С Ċ IOPT=-1 Plot lower display. C C IOPT=0 Plot both displays. IOPT=1 Plot upper display. С С COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPZER/ ITYZER, IBYZER DATA NESC /27/ С С IOP = IOPTС С SET UP THE HP 7470A (7475A) PLOTTER С WRITE(8,10) NESC,'.I11;;17:',NESC,'.N;19:' 10 FORMAT('\$',A1,A9,A1,A6) С C XON AND XOFF HAVE NOW BEEN ESTABLISHED, ESC I11;;17 SETS THE XON THRESHOLD TO 11 REMAINING BYTES AND XON CHARACTER TO DC1, ESC .N;19: С C SETS THE XOF CHARACTER TO DC3. С С INITIALIZE THE PLOTTER. С C WRITE(8,\*) 'IN;SC 0,4095,0,2048;' С

С GIVE RANGE OF PLOT C WRITE(8,\*) 'IP 500,600,9500,7000;' С C C C C C C GET OFFSET OF ARRAY ARGUMENT FOR POINTS TO BE PLOTTED IF LOWER DISPLAY IS TO BE PLOTTED IAROFF = 0100 IF(IOP.LE.O) IAROFF = 2 \* NDSP С C C SET ZERO OFFSET OF DISPLAY TO BE PLOTTED TO ONLY PLOT NON-ZERO POINTS IYOFF = IBYZERIF(IOP.GT.O) IYOFF = ITYZER С С SEND SYMBOL FOR PLOT (\* FOR UPPER, O FOR LOWER) Č CHOOSE PEN FOR PLOT (LEFT FOR UPPER, RIGHT FOR LOWER) С IPEN = 1SYMBOL = '\*'IF(I0P.GT.0) G0 T0 1400 SYMBOL = '0'IPEN = 21400 WRITE(8,1410) SYMBOL, IPEN FORMAT(' SM',A1,';SP',I1,';') 1410 C Č PLOT ALL NON-ZERO POINTS С D0 2000 I=1,NDSP IARG = (2\*I) + IAROFFIF(IDSP(IARG).NE.IYOFF) WRITE(8,\*) 'PU;PA ' 1, ÌDSP(ÌARG-1), ', ', (IDSP(IARG)-IYOFF), '; PD; ' 2000 CONTINUE WRITE(8,\*) 'SP;PU;PA 4000,2000;' IF(IOP.NE.O) G0 T0 9000 С C C IF BOTH ARE TO BE PLOTTED, GO BACK AND DO UPPER IOP = IOP + 1GO TO 100 **REWIND 8** 9000 RETURN END

SUBROUTINE SGTSCR

С

С

File SGTSCR.FOR begins on the previous line.

С Version Date: September 25, 1984 С Paul S. Weiss č C C C Subroutine SGTSCR converts scaler contents, currently stored in JDATA, into real numbers stored in SCR, and resets the run flag. INTEGER+4 JDATA COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SCR/ SCR(20) COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /SDONE/ ISWFLG, JDATA (20) С IRNFLG = 0SCR(1) = ANGLEС С DETERMINE THE TIME COUNTED С SCR(2) = .150 + GATE(1) + INTERV / (CLOCK + 2.)DO 100 I=1, ((2\*NCHAN) - 2)SCR(I+2) = AJFLT(JDATA(I))100 CONTINUE С С IN LAST CHANNEL, SHOW THE DIFFERENCE OF THE FIRST TWO CHANNELS Ċ WITH THE CORRECT SIGNAL TO NOISE BY: С nA = 1A + 2BС nB = 1B + 2AС SCR((2\*NCHAN) + 1) = SCR(3) + SCR(6)SCR((2\*NCHAN) + 2) = SCR(4) + SCR(5)С С MAKE THE REMAINDER OF THE CHANNELS O С IF (NCHAN.GE.9) G0 T0 9000 D0 200 I=(2\*NCHAN)+3,20 SCR(I) = 0.200 CONTINUE 9000 RETURN END SUBROUTINE SINIT С File SINIT.FOR begins on the previous line. C C C C C Version Date: November 11, 1985 Paul S. Weiss C C Subroutine SINIT initializes the parameters of angular Ē scan program SANG to the default values given below. C С

## INTEGER+4 JDATA

С

С

С

С

C C C

С

С

1

```
COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG
 COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256)
 COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP
 COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB
 COMMON /TGSMAX/ DWLMAX, MAXCHN
 COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN
 COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA
 COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG
 COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG
 COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG
 COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG
 COMMON /VELO/ DELF, HALFH, DELPIP, VELB, SPDRIO
 COMMON /SDONE/ ISWFLG, JDATA
 COMMON /CHADEF/ CHADEF(20,10)
 COMMON /PRES/ PRM(8), IPRE(8), PFORE(8), IDG(8), IPEM(8)
 1,PBACK(2),PION(3)
 COMMON /TPRIM/ THMCPL(8)
 COMMON /MACDEF/ TCDEF (8,4), NTC, PRDEF (8,6), NPR, BACK (4,2)
COMMON /DET/ AGRD, AFIL, VEB, VIE, VEXT
 1, VEN, QPST, QPR, MAS, VEX, VDK, VPMT
 DATA CHADEF /'Lase','r On',' Det','ecto','r Si','gnal',14*' '
                                                          ',13*<sup>',</sup> '
 1, 'Lase', 'r Of', 'f De', 'tect', 'or S', 'igna', 'l
 2,'Lase','r Fl','uore','scen','ce ',15*''
3,'Lase','r On',' - L','aser',' Off',' Det','ecto','r Si'
3,'gnal',11*''
```

```
4,120*' '/
```

```
DATA TCDEF /'Na T', 'Na B', 'Na N', 'Na p', 'Na S', 'Cold', 'Seco'
1, 'None', 'op O', 'otto', 'ozzl', 're-S', 'kimm', 'Shi', 'ndar', '
2, 'ven ', 'm Ov', 'e ', 'kimm', 'er ', 'eld ', 'y No', 2*'
3, 'en ', ', 'er ', 2*' ', 'zzle', ''/
```

```
DATA PRDEF /2*'Prim', 'Main',2*'Seco', 'Regi', 'Betw', 'None'
1,2*'ary ', 'Cha',2*'ndar', 'on 3', 'een ',' '
2, 'Sour', 'Diff', 'mber', 'y So', 'y Di', ' ', 'Turb', ' '
4, 'ce ', 'eren', ' ', 'urce', 'ffer', ' ', 'os', '
5, ' ', 'tial',2*' ', 'enti',7*' ', 'al ',3*' '/
```

```
DATA BACK /'Heli','um ',2*' ','HCl ',3*' '/
```

RELAY ADDRESSES TO MACROS

CALL INIT (ISWFLG, JDATA)

TYPE 1 FORMAT(//, ' Angular Scan Program' 1,/,' Version Date: November 11, 1985',/)

## SET DATA ACQUISITION PARAMETER DEFAULTS

MAXCHN = 8ISWFLG = 0CLOCK = 1.5INTACQ = 90NCHAN = 4EXCFRC(1) = .3EXCFRC(2) = 0.DO 2200 I=1,16 GATE(I) = 3.22200 CONTINUE SET CURRENT ACQUISITION DEFAULTS ANGLE = 40. INTERV = 90NUMANG = 0MANGLE = 1IRNFLG = 0SET NORMALIZATION DEFAULTS ANGNRM = 40. INTNRM = 90NRMRPT = 2NRMFLG = 0SET DISPLAY DEFAULTS MDSPT = 1MDSPB = 2NDSPT = 1NDSPB = 2DSPB = -10. NDSP = 221DSTEP = .5MODSPT = 1MODSPB = 1INDEP = 0SET LOOP MODE DEFAULTS ANLPBE = 25. ANLPFN = 75.ANLPIN = 5.ANLPCU = 25. LPFLAG = 0

С С С

C C C C

C C C

C C C

C C C

LPNRFQ = 10

SET ROTATION OF POLARIZATION DEFAULTS ANGPOL = 45. POLSTP = 22.5POSIT = 0.IPOSIT = 0IPLFLG = 0SET TEST MODE DEFAULTS TSTANG = 0.NUMTST = 200TSTCUR = 0.TSTFLG = 0.SET MACHINE CONDITION DEFAULTS FIRST THERMOCOUPLE DEFAULTS NTC' = 7 $\mathsf{THMCPL}(1) = 23.$  $\mathsf{THMCPL}(2) = 21.3$ THMCPL(3) = 25.  $\mathsf{THMCPL}(4) = 19.$ THMCPL(5) = 17.5THMCPL( $\delta$ ) = -6.2  $\mathsf{THMCPL}(7) = 7.33$ THMCPL(8) = 0.SET PRESSURE DEFAULTS NPR = 7DO 2500 I=1,8 PRM(I) = 1.0IPRE(I) = -4PFORE(I) = 10.IDG(I) = 1IPEM(I) = 1CONTINUE IPRE(2) = -5IPRE(3) = -7IPRE(5) = -5IPRE(6) = -10IPRE(7) = -7IPEM(3) = 10IPEM(6) = 10IPEM(7) = 10PFORE(6) = 0.PFORE(7) = 0.PION(1) = 35.

С С С

с с с с с

С С С

PION(2) = 6.PION(3) = 0.PBACK(1) = 500. $\mathsf{PBACK}(2) = 500.$ IDG(6) = 2IDG(7) = 2C C C SET DETECTOR SETTING DEFAULTS AGRD = 10. AFIL = 5.5VEB = 200. VIE = -75. VEXT = -150. VEN = -250. QPST = 4.68QPR = 3.75MAS = 23VEX = -600. VDK = -30. VPMT = -2350. C C C SET VELOCITY MEASUREMENT DEFAULTS FRQLAS = 16973.FSRLAS = 1.5LASDWL = 100LASSTP = 1LASCHA = 4VELB = 0.SPDRIO = -1. C C SET CHANNEL DEFINITION DEFAULTS С D0 3200 I=5,10 CHADEF(1, I) = 'None'3200 CONTINUE 🔨 С Č ZERO DATA FILES С DO 4200 I=1,256 DATA(I,1) = -100.DO 4100 J=2,20 DATA(I, J) = 0.4100 CONTINUE 4200 CONTINUE RETURN END

C	SUBROUTINE SLEAVE File SLEAVE.FOR begins on the previous line.
	Version Date: July 4, 1984 Paul S. Weiss
	Subroutine SLEAVE is the termination routine for SANG, the angular scan program.
1000	COMMON /FNAME/ INAME(7),ISTORE TYPE *,'Exit? Are you sure?' ACCEPT 1000, ANS FORMAT(A1) IF(ANS.NE.'Y') RETURN
C C	TURN OFF INTERRUPTS AND DISPLAY WITH MACRO SUBROUTINE SEXIT
9000 9100 9110	CALL SEXIT CLOSE(UNIT=9,DISPOSE='SAVE',ERR=9100) STOP TYPE 9110,(INAME(I),I=1,7),7,7,7,7 FORMAT(/,'Error closing file ',7A2,4A1) GO TO 9000 END
C	SUBROUTINE SLOOP(IOPT) File SLOOP.FOR begins on the previous line.
	Version Date: July 4, 1984 Paul S. Weiss
	Subroutine SLOOP starts or continues the data taking loop. It is called from SCOMM.
	Its argument IOPT is used as follows: IOPT=O Continue IOPT=1 Start Loop
C	COMMON /CUR/ ANGLE,INTERV,NUMANG,MANGLE,IRNFLG COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /LOPSET/ ANLPBE,ANLPFN,ANLPIN,LPNRFQ,ANLPCU,LPFLAG COMMON /NRMSET/ ANGNRM,INTNRM,NRMRPT,NRMFLG COMMON /SDONE/ ISWFLG,JDATA(20)
	TE (TONEL ONE O) ON TO ROOO

IF(I0PT.NE.0) G0 T0 2100 С С CONTINUE LOOP FROM WHERE IT LEFT OFF С LPFLAG = 1ANGLE = ANLPCUINTERV = INTACQGO TO 8000 С С IF STARTING A LOOP, SET ANGLE FOR A NORMALIZATION С 2100 TYPE \*, 'Start Loop with a Normalization!' ANGLE = ANGNRMINTERV = INTNRM LPFLAG = -2NRMFLG = 1С С SET FLAG FOR LOOP MODE С CALL SSTART(1) 8000 GO TO 9000 8900 TYPE \*, 'Abort Acquisition -- TGS is Running!' 9000 RETURN END SUBROUTINE SNORM С File SNORM FOR begins on the previous line. C Ċ Version Date: July 2, 1984 C C C C C C C C Paul S. Weiss Subroutine SNORM starts the Normalization. С COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /SDONE/ ISWFLG, JDATA (20) С С TYPE \*, 'Normalization? Are you sure?' ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE.'Y') GO TO 9000 NRMFLG = 1ANGLE = ANGNRMINTERV = INTNRM CALL SSTART(1) 9000 RETURN

SUBROUTINE SNUCH (NU, ICOM) File SNUCH.FOR begins on the previous line. Version Date: July 24, 1984 Paul S. Weiss Subroutine SNUCH chooses which subroutine to call for the P\*, S\*, and W\* commands. It was created to streamline the program by removing some of the command burden from subroutine SCOMM. The arguments of SNUCH are as follows: NU Logical Unit for output. ICOM Letter of command typed in. COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG COMMON /TSTSET/. TSTANG, NUMTST, TSTCUR, TSTFLG COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG DATA IA, IC, ID, IG /65, 67, 68, 71/ DATA II, IL, IM, IN, IR, IS, IT, IV /73, 76, 77, 78, 82, 83, 84, 86/ The commands are: \*A Write Accumulated Data \*C Write Current Configuration \*D Write Data Channel Definition ∗G Write Curent Graphics Display \*I Write Information \*L Write Loop Mode Parameters \*M Write Machine Condition Write Normalization Parameters \*N \*R Write Rotation of Polarization Parameters \*S Write Detector Settings \*T Write Test Mode Parameters \*¥ Write Velocity Measurement and Parameters IF(ICOM.NE.IA) GO TO 11 CALL SACCUM (NU) GO TO 100 IF(ICOM.NE.IC) GO TO 16 CALL SSHCUR(NU) IF LOOP MODE, TEST MODE, OR NORMALIZATION IS ACTIVE, ALSO SHOW THOSE SETTINGS IF (LPFLAG.NE.O) CALL SSHLOP (NU)

END

С

CCCCCCCCCCCC

С

C

11

C C

C

IF (TSTFLG.NE.O.) CALL SSHTST (NU) IF (NRMFLG.NE.O) CALL SSHNRM (NU) IF(IPLFLG.NE.O) CALL SSHPOL(NU) GO TO 100 IF(ICOM.NE.ID) GO TO 31 16 CALL SSHDEF (NU) GO TO 100 IF(ICOM.NE.IG) GO TO 36 31 CALL SSHDSP(NU) GO TO 100 36 IF(ICOM.NE.II) GO TO 41 CALL SPRII(NU) GO TO 100 IF(ICOM.NE.IL) GO TO 46 41 CALL SSHLOP (NU) GO TO 100 IF(ICOM.NE.IM) GO TO 51 46 CALL SSHMP (NU) GO TO 100 51 IF (ICOM.NE.IN) GO TO 61 CALL SSHNRM(NU) GO TO 100 61 IF(ICOM.NE.IR) GO TO 66 CALL SSHPOL (NU) GO TO 100 66 IF(ICOM.NE.IS) GO TO 71 CALL SSHMD (NU) GO TO 100 IF(ICOM.NE.IT) GO TO 81 71 CALL SSHTST(NU) GO TO 100 IF(ICOM.NE.IV) GO TO 99 81 CALL SSHVEL (NU) GO TO 100 NU = 099 100 RETURN END SUBROUTINE SOUT (NU) С File SOUT.FOR begins on the previous line. С С Version Date: September 5, 1984 С Paul S. Weiss С С

С

С

С

Subroutine SOUT outputs the current data to logical unit NU.

COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /CUR/ ANGLE,INTERV,NUMANG,MANGLE,IRNFLG

COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG COMMON /SCR/ SCR(20) С LOGICAL\*1 ITIME(8) CALL TIME(ITIME) WRITE(NU, 100) SCR(1), SCR(2), (ITIME(I), I=1,8), 7 100 FORMAT(/,' Angle: ',F5.1,' degrees, measured for: ',F8.3,' sec.' 1,6X,'Time: ',9A1) С С IS THIS A POLARIZATION ROTATION SCAN? IF SO, PRINT POLARIZATION С ANGLE С IF(IPLFLG.NE.0) WRITE(NU,120) (10000.\*(SCR(1)-(INT(SCR(1))))) 120 FORMAT(' The laser polarization is ',F5.1 1,' degrees from vertical.') WRITE (NU, 140) FORMAT(' Channel', 3X, 'A', 10X, 'B', 11X, 'Signal' 140 3,9X, 'Error',6X, 'Signal/Noise') DO 250 I=1,NCHAN IS = (2\*I) + 1IB = IS + 1SIG = (SCR(IS) - SCR(IB))/SCR(2)ERR = SCR(IS) + SCR(IB)ERR = SQRT(ERR)/SCR(2)SGTON = 0.IF(ERR.NE.O.) SGTON = SIG/ERR WRITE(NU, 200) I, SCR(IS), SCR(IB), SIG, ERR, SGTON 200 FORMAT(1X, I1, 2X, F9.0, 2X, F9.0, 5X, F11.2, '+/-', F9.2, 5X, F9.1) 250 CONTINUE RETURN END SUBROUTINE SOVER (ISTART) С File SOVER.FOR begins on the previous line. С С С С С October 30, 1984 Version Date: Paul S. Weiss C C Subroutine SOVER "handles" overflows. It tells the user of the condition, and asks whether or not to restart the acquisition. C TYPE 100,7,7,7,7,7,7,7,7,7,7,7 100 FORMAT(' \*\*\*\*\*\*\*\*\*\*\*\*\*\* 1,/, ' An overflow has occurred!' 2,/,' \*\*\* 3,///,' Try retaking this point?') ISTART = 0ACCEPT 1000, ANS

1000 FORMAT(A1) IF(ANS.NE.'Y') GO TO 9000 ISTART = 1TYPE \*,'Be sure to lower the gain in some way before restarting!' 9000 RETURN END SUBROUTINE SPOLAR С File SPOLAR.FOR begins on the previous line. C С Version Date: September 27, 1984 С Paul S. Weiss C С Subroutine SPOLAR starts the polarization rotation data taking loop. С It is called from SCOMM. С С COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG С IF(IRNFLG.NE.O) GO TO 8800 TYPE \*, 'Polarization Rotation Loop? Are you sure?' ACCEPT 1000, ANS 1000 FORMAT(A1) IF(ANS.NE.'Y') GO TO 8900 TYPE \*, 'Change display window for polarization rotation scan?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') CALL SCHDWN (2) С С START A LOOP, SET ANGLE FOR POLARIZATION С ANGLE = ANGPOLC С SET COUNTDOWN INTERVAL С INTERV = INTACQС С SET FLAG FOR POLARIZATION ROTATION LOOP MODE С IPLFLG = 1С С SET STEP INDICATORS, AND OTHER PARAMETERS OF MEASUREMENT C POSIT = 0.IPOSIT = 0IF (POLSTP.LT.O.) POLSTP = - POLSTP TYPE \*,'Set the laser polarization vertical at the Main Door.'

8000 8800 8900 9000	CALL SSTART(1) G0 T0 9000 TYPE *,'TGS is Running!' TYPE *,'Abort Acquisition' RETURN END
C	SUBROUTINE SPRII(NU) File SPRII.FOR begins on the previous line.
	Version Date: September 3, 1984 Paul S. Weiss
	Subroutine SPRII prints information on the line printer or into a disk file to be associated with a data file.
	The arguments of SPRII are as follows: NU Unit number to which information is to be written. Note that if information is written to disk, the same information is also written to the printer.
	LOGICAL+1 CHAR(80),ITIME(8) INTEGER JDATE(3) IF(NU.EQ.7) GO TO 150 TYPE 110
110	FORMAT(/, ' Include date and time?')
1000	FORMAT (A1)
150	CALL IDATE(JDATE(1), JDATE(2), JDATE(3)) CALL TIME(ITIME)
200	WRITE(NU,200) JDATE(1), JDATE(2), (JDATE(3)+1900), (ITIME(I), I=1,8) FORMAT(/,' Date: ',2(I2,'/'),I4,/,' Time: ',8A1,/) IF(NU.EQ.9) WRITE(6,200) JDATE(1), JDATE(2) 1,(JDATE(3)+1900), (ITIME(I),I=1,8)
C	
ç	IF UNLT WRITING TO THE TERMINAL, EXIT
с	IF (NU.EQ.7) GU IU 9000
C C	ENABLE LOWER CASE INPUT
1100 1110	CALL LCASE TYPE 1110 FORMAT(/,' Type information in. To stop, begin a line with' 1.' a ='./)
1200	D0 1205 $I=1,80$ CHAR(I) = 0

1205 CONTINUE ACCEPT 1210, (CHAR(I), I=1,80) FORMAT(80A1) 1210 IF(CHAR(1).EQ.61) GO TO 8000 WRITE(NU,1310) (CHAR(I), I=1,80) 1310 FORMAT(1X,80A1) С С IF WRITING TO DISK, ALSO WRITE TO THE PRINTER С IF(NU.EQ.9) WRITE(6,1310) (CHAR(I), I=1,80) GO TO 1200 С С DISABLE LOWER CASE INPUT С 8000 CALL UCASE 9000 RETURN END SUBROUTINE SSAVE С File SSAVE.FOR begins on the previous line. С С Version Date: July 4, 1984 С Paul S. Weiss С С Subroutine SSAVE puts new data in the accumulated data C array DATA, and adjusts the order of the array so that Ċ there is room for the new data. С C C COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256) COMMON /SCR/ SCR(20) COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN С IF (NUMANG.EQ.0) GO TO 1200 D0 100 I=1, NUMANG IF(SCR(1) - DATA(I,1)) 2000,5000,100100 CONTINUE С C C IF THE DO LOOP HAS RUN OUT, THEN INCREASE NUMANG, AND PUT NEW DATA AT THE LAST POINT С IF (NUMANG.LT.256) GO TO 1200 TYPE 1110,7,7 FORMAT(/, ' Accumulated data array full!', 2A1 1110 1,/,' Abort Storage!') GO TO 9000 1200 NUMANG = NUMANG + 1MANGLE = NUMANG

C	GŨ TŨ 3000
	IF THE NEW DATA ANGLE FALLS IN THE MIDDLE OF THE ARRAY DATA (ORGANIZED BY ANGLE), THEN MOVE ALL SUBSEQUENT DATA UP BY ONE AND PUT THE NEW DATA IN THE CORRECT PLACE
2000	IF(NUMANG.LT.256) GO TO 2200 TYPE 1110,7,7 GO TO 9000
2200	NUMANG = NUMANG + 1 MANGLE = I
	SHIFT OLD DATA TO MAKE ROOM FOR NEW DATA
	DO 2400 I=1,NUMANG-MANGLE IOLD = NUMANG - I INEW = IOLD + 1 DO 2300 J=1,20 DATA(INEW I) = DATA(IOLD I)
2300 2400 C	CONTINUE CONTINUE
C C	INSERT NEW DATA
3000	D0 3200 J=1,20 DATA(MANGLE,J) = SCR(J)
3200	CONTINUE IF(NUMANG.EQ.256) TYPE 3310,7,7,7
3310	FORMAT(/,' Warning Accumulated data array is now full!' 1,/,' Do not take data at any new angles!',3A1) GD TO 9000
5000	MANGLE = I D0 5200 J=2,20 DATA(MANGLE,J) = DATA(MANGLE,J) + SCR(J)
5200 9000	CONTINUE RETURN END
C	SUBROUTINE SSHCUR(NU) File SSHCUR.FOR begins on the previous line.
	Version Date: October 17, 1984 Paul S. Weiss
č c c	Subroutine SSHCUR prints the current settings on logical unit NU.

	COMMON /COR/ ANGLE, INTERV, NOMANG, MANGLE, IRNELG COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /LASSET/ LASDWL EDDLAS LASSTR ESPLAS LASCHA
	COMMON /LOPSET/ ANLPBE, ANLPFN, ANLPIN, LPNRFQ, ANLPCU, LPFLAG
•	COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG
	COMMON /FULSET/ ANGFUL,FULSTF,FUSIT,IFUSIT,IFLFLG
	COMMON /FNAME/ INAME(7), ISTORE
C	
100	NUNIT = NU WRITE (NUNIT 110) NCHAN (LOCK INTACO CATE(1) CATE(2) EXCERC(1)
100	1.EXCFRC(2)
110	FORMAT(/,' The Quad Scalers and Laser Gate Generator are in use.'
	1,/,1X,I1,' data channels are being used.'
	2,/,' the clock frequency is ',F3.1,' Hz.'
	4/. The date widths are: A = '.F5.3.' msec. B = '.F5.3.' msec.'
	5,/,' The Excited State Fractions are: ',F6.4,' and ',F6.4,/)
	WRITE(NUNIT,2100) (INAME(I),I=1,7),ISTORE
2100	FURMAT(' The file ',7A2,' has recorded ',15,' IGS countdowns.')
3100	FORMAT(' The detector angle is currently ' F5.1.' degrees'
	1,/,' The countdown interval is ',I5,'.'
	2,/,1X,I3,' detector angles have been measured.')
	IF (LPFLAG.NE.O) WRITE (NUNIT, 3200) 'in Loop Mode.'
	IF (IPLEIG.NE.O) WRITE (NUNIT. 3200) 'rotating the '
	1, 'laser polarization.'
	IF(TSTFLG.GT.O.) WRITE(NUNIT,3200) 'in Test Mode.'
3200	FORMAT(' The acquisition is currently ',A13,:,A19)
3300	FORMAT(' The TGS is counting down ')
C	
C	IF CONFIGURATION WAS WRITTEN TO DISK, WRITE IT TO THE PRINTER ALSO
C	TE (NUNTTINE 0) CO TO 0000
	$NUNIT = 6^{1}$
•	GO TO 100
9000	RETURN
· .	END
c	SUBROUTINE SSHDEF (NU)
c	File SSHUEF.FUK begins on the previous line.
č	Version Date: October 17, 1984
C	Paul S. Weiss
C	
L L	

ŝ

C Subroutine SSHDEF prints the array definitions to logical unit С number NU. C COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /CHADEF/ CHADEF(20,10) С NUNIT = NU 100 D0 2000 I=1,NCHAN WRITE (NUNIT, 1200) I, (CHADEF (J, I), J=1, 20) FORMAT(' Channel #', I1, ' is:',/, 1X, 20A4) 1200 2000 CONTINUE С С IF THIS IS TO BE WRITTEN TO DISK, PRINT ALSO С IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6GO TO 100 9000 RETURN END -SUBROUTINE SSHDSP(NU) С File SSHDSP.FOR begins on the previous line. С Ċ Version Date: October 23, 1984 C C C C C C C C Paul S. Weiss Subroutine SSHDSP shows the current display mode on logical unit NU. C COMMON /DISP/ IDSP(1064), MODSPT, MODSPB, DSPB, DSTEP, NDSP, INDEP COMMON /DSPCHA/ MDSPT, NDSPT, MDSPB, NDSPB DATA TOP1, TOP2, BOT1, BOT2 /'uppe', 'r', 'lowe', 'r'/ С NUNIT = NUMD = MODSPT100 TB1 = TOP1TB2 = TOP2M = MDSPTN = NDSPTGO TO (1100,1200,1300,1400,1500,1600),MD 1050 WRITE(NUNIT, 1110) TB1, TB2, M, 'Signa', 'I.' 1100

1110 FORMAT(' The ',A4,A1,' display is Channel ',I1,1X,A5,A2,A6) G0 T0 2000

1200 WRITE(NUNIT, 1210) TB1, TB2, M, N, 'Difference.'

1210 FORMAT(' The ',A4,A1,' display is Channels ',I1,',',I1 1,1X,A11,A10,A3) GO TO 2000

1300	WRITE(NUNIT,1210) TB1,TB2,M,N,'Excited Sta','te Signal.' GO TO 2000
1400	WRITE(NUNIT,1210) TB1,TB2,M,N,'Difference ','Signal/Noi','se.' G0 TO 2000
1500	WRITE(NUNIT,1110) TB1,TB2,M,'Signa','//','Noise.'
1600	WRITE (NUNIT 1110) TR1 TR2 M 'Time '
2000	IF (TB1.NE.TOP1) G0 T0 8000 $MD = M0DSPB$ $TB1 = B0T1$ $TB2 = B0T2$ $M = MDSPB$ $N = NDSPB$ $G0 T0 1050$
8000	IF(INDEP.EQ.0) WRITE(NUNIT,8010) 'to the same scale. ' IF(INDEP.NE.0) WRITE(NUNIT,8010) 'scaled independently.'
8010	FORMAT(' The upper and lower display channels are ',A21) WRITE(NUNIT.8110) DSPB.(DSPB + ((NDSP-1.)*DSTEP)) DSTEP
8110	FORMAT(' The display window is from ',F9.5,' to ',F9.5
c	
C C	IF THIS IS TO BE WRITTEN TO DISK, PRINT ALSO
9000	IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6 GO TO 100 RETURN END
C	SUBROUTINE SSHLOP(NU) File SSHLOP.FOR begins on the previous line.
	Version Date: October 23, 1984 Paul S. Weiss
C C C	Subroutine SSHLOP prints the current Loop Mode parameters on logical unit NU.
° ·	COMMON /LOPSET/ ANLPBE,ANLPFN,ANLPIN,LPNRFQ,ANLPCU,LPFLAG
100 110	NUNIT = NU WRITE(NUNIT,110) ANLPBE,ANLPFN,ANLPIN,LPNRFQ,ANLPCU FORMAT(/,' The Loop goes from ',F5.1,' to ',F5.1,' degrees in ' 1,F5.2,' degree steps.' 2,/,' Normalization occurs every',I5,' angles.' 3,/,' The current angle in the Loop is ',F5.1,' degrees.')
C C C	IF THIS IS TO BE WRITTEN TO DISK, PRINT IT ALSO
*	

## IF(NUNIT.NE.9) GO TO 9000 NUNIT = 6 GO TO 100 RETURN END

SUBROUTINE SSHMP (NU)

C C

9000

SUBROUTINE SSHMD (NU) С File SSHMD.FOR begins on the previous line. С C Version Date: October 17, 1984 C C C Paul S. Weiss C C Subroutine SSHMD writes out the detector settings to logical unit NU. С COMMON /DET/ AGRD, AFIL, VEB, VIE, VEXT 1, VEN, QPST, QPR, MAS, VEX, VDK, VPMT С NUNIT = NU100 WRITE(NUNIT, 110) AGRD, AFIL, VEB, VIE, VEXT 1, VEN, QPST, QPR, MAS, VEX, VDK, VPMT FORMAT(/, ' The detector settings are: ' 110 1,/,' Grid Current: ',F4.1,' mA' 2,/,' Filament Current: ',F4.2,' A' 3,/,' Electron Beam Energy: ',F5.0,' V' 4,/,' Ion Energy: ',F7.2,' V' 5,/,' Ion Extraction Voltage: ',F5.0,' V' ,'Entrance Lens Voltage: ',É5.0,'V' 6,/, 7,/,' Quadrupole Mass Setting = ',F5.2,' Resolution = ',F5.2 8, Mass = ', I59,/,' Exit Lens Voltage = ',F6.0,' V' 1,/,' Doorknob Voltage = ',F5.1,' kV' 2,/,' Photomultiplier Voltage = ',F6.0,' V') С С IF WRITING TO DISK, PRINT ALSO С IF(NUNIT.NE.9) GO TO 9000 NUNIT = 6GO TO 100 9000 RETURN END

File SSHMP.FOR begins on the previous line.

C Version Date: October 23, 1984 С Paul S. Weiss C Ċ C C Subroutine SSHMP outputs the Machine pressures and temperatures to logical unit NU. Ċ DIMENSION NONOFF(2) COMMON /PRES/ PRM(8), IPRE(8), PFORE(8), IDG(8), IPEM(8) 1, PBACK(2), PION(3) COMMON /TPRIM/ THMCPL(8) COMMON /MACDEF/ TCDEF(8,4), NTC, PRDEF(8,6), NPR, BACK(4,2) DATA NONOFF /'n ','ff'/ С NUNIT = NU.100 WRITE (NUNIT, 110) (PBACK (I), (BACK (J, I), J=1, 4), I=1, 2) FORMAT(/,' The source backing pressures are: 110 1,/,' Primary Source: ',F8.1,' torr with ',4A4 2,/,' Secondary Source: ',F8.1,' torr with ',4A4) WRITE(NUNIT,200) ((PRDEF(I,J),J=1,6),PRM(I),IPRE(I),PFORE(I) 1,NONOFF(IDG(I)), ÌPEM(I), I=1,NPR) FORMAT(/,' The Regional Pressures are:' 200 1,8(/,1X,6A4,1X,F4.2,'x10',I3,' torr, foreline: ',F4.0 2,' u, Degas 0',A2,', ',I2,' mA',:)) WRITE(NUNIT,300) (I,PION(I),I=1,3) FORMAT(/,' The ion pump currents are:',/ 300 1,3(' Region ',I1,': ',F6.1,' uA ')) WRITE(NUNIT, 400) ((TCDEF(I, J), J=1, 4), THMCPL(I), I=1, NTC) 400 FORMAT(/,' The measured Thermocouple voltages are:' 1,4(/,2(1X,4A4,1X,F5.2,' mV ',:))) С С IF THIS IS TO BE WRITTEN TO DISK, PRINT ALSO С IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6GO TO 100 9000 RETURN END SUBROUTINE SSHNRM (NU) С File SSHNRM.FOR begins on the previous line. С С Version Date: October 23, 1984 Č Paul S. Weiss С С C C Subroutine SCHNRM prints the current Normalization parameters on logical unit NU. С COMMON /NRMSET/ ANGNRM, INTNRM, NRMRPT, NRMFLG

С NUNIT = NUWRITE (NUNIT, 110) ANGNRM, INTNRM, NRMRPT 100 FORMAT(/, ' Normalization occurs at ', F5.1, ' degrees.' 110 1,/,' The countdown interval is ', I5,', and is ' 2, 'repeated ', I3, ' times.',/) С С IF THIS IS TO BE WRITTEN TO DISK, PRINT ALSO С IF (NUNIT.NE.9) GD TO 9000 NUNIT = 6GO TO 100 RETURN 9000 END SUBROUTINE SSHPOL(NU) С File SSHPOL.FOR begins on the previous line. C C Version Date: October 23, 1984 C C C C C Paul S. Weiss Subroutine SSHPOL outputs the Polarization Rotation Mode parameters to logical unit NU. С С COMMON /POLSET/ ANGPOL, POLSTP, POSIT, IPOSIT, IPLFLG Ը Շ С NUNIT = NU100 WRITE(NUNIT,110) ANGPOL,POLSTP,INT((180./POLSTP)+0.5) 1, (POSIT\*180./200.) FORMAT(/,' The detector angle for the Polarization Rotation ' 110 1,'Loop is ',F5.1,' degrees.' 2,/,' The motor step size is ',F6.2,' degrees,' 3,/,' which scans through 180 degrees in ',I3,' steps.' 3,/, 4,/,' The polarization position is ',F7.3,' degrees from' 4, 'vertical.') C C C C IF THIS IS TO BE WRITTEN TO DISK, PRINT ALSO IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6GO TO 100 9000 RETURN END SUBROUTINE SSHTST (NU)

File SSHTST.FOR begins on the previous line.

Version Date: October 23, 1984 Paul S. Weiss

С

С С

C C C C C C C C

Č

100

110

C С

С

9000

С

C C

Ċ

C C C

С

С

Subroutine SSHTST prints the current Test mode parameters on logical unit NU.

COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG NUNIT = NU

WRITE(NUNIT, 110) TSTANG, NUMTST, INT(1.E5+TSTCUR) FORMAT(/,' The current Test mode angle is ',F5.1,' degrees. 1,/,1X,I3,' countdowns can be taken at this angle.' 2,/,' The current countdown will be #',I3,/)

IF THIS IS TO BE WRITTEN TO DISK PRINT IT ALSO

IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6GO TO 100 RETURN END

SUBROUTINE SSHVEL (NU)

File SSHVEL.FOR begins on the previous line.

Version Date: October 23, 1984 Paul S. Weiss

Subroutine SSHVEL outputs the Velocity measurement parameters and measurement on logical unit NU.

COMMON /LASSET/ LASDWL, FRQ	AS, LASSTP, FSRLAS, LASCHA
COMMON /VELO/ DELF, HALFH, DI	ELPIP, VELB, SPDRIO
NUNIT = NU	
WRITE(NUNIT, 110) FRQLAS, FS	RLAS

100 FORMAT(/,' The Laser Frequency is ',F10.3,' cm-1.' 110 1,/,' The Frequency Standard is ',F5.2,' GHz.') IF (VELB.LE.O.) GO TO 8000 IF(SPDRIO.LE.O.) GO TO 7000 WRITE (NUNIT, 2100) VELB, SPDRID, DELF, HALFH, DELPIP 2100

FORMAT(' The beam velocity is ',F10.0,' cm/sec.'

1,/,' The Speed Ratio is ',F5.2 2,/,' The 45 to 90 degree separation was ',F10.4 

377 .

7000 7010	GO TO 8500 TYPE 7010,VELB,DELF,DELPIP FORMAT(' The beam velocity is ',F10.3,' cm/sec.' 1,/,' The Speed Ratio was not measured.' 2,/,' The 45 to 90 degree separation was ',F10.4 4,/,' The frequency standard separation was ',F10.4,/) GO TO 8500
8000 C	TE WRITTING TO DISK ALSO PRINT
C 8500 9000	IF (NUNIT.NE.9) GO TO 9000 NUNIT = 6 GO TO 100 RETURN END
C	SUBRUUTINE SSTART(IUPT) File SSTART.FOR begins on the previous line
	Version Date: September 3, 1984 Paul S. Weiss
	Subroutine SSTART starts the TGS counting down using MACRO subroutine SGO. It sets the display correctly for the angle under study, so that the current point is brightened.
C	The argument of SSTART is:
	IOPT=-1 Skip all queries. IOPT=1 Query as to whether angle and settings are correct.
C	INTEGER*4 JDATA COMMON /CUR/ ANGLE,INTERV,NUMANG,MANGLE,IRNFLG COMMON /DATA/ DATA(256,20),WORK1(256),WORK2(256) COMMON /SETS/ GATE(16),TIM(16),EXCFRC(2),CLOCK,INTACQ,NCHAN COMMON /SDONE/ ISWFLG,JDATA(20) COMMON /SCR/ SCR(20)
	IF(IRNFLG.EQ.1) GO TO 8000 IF(IOPT.LE.O) GO TO 2000 TANG = ANGLE + 270. IF(TANG.GE.360.) TANG = TANG - 360.
90 100	TYPE 100,TANG,INTERV FORMAT(' The detector angle is ',F5.1,' degrees.' 1,/,' The Countdown interval is ',I5,'.' 2,/,' Are these correct and are you ready?')

	ACCEPT 1000, ANS
1000	FORMAT(A1)
	IF (ANS. NE. 'Y') GO TO 8500
2000	CALL SGO(ISWFLG, JDATA)
	TYPE 4100.7.0
4100	FORMAT(' ŤGŚ is running.'.2A1)
	IRNFLG = 1
	ISWFLG = O
	SCR(1) = ANGLE
	D0 5000 I=2.20
	SCR(I) = 0.
5000	CONTINUE
	RETURN
8000	TYPE 4100.7.7
8500	TYPE * 'Abort Requested Acquisition?'
	ACCEPT 1000 ANS
	TE (ANS NE 'Y') GO TO 90
	RETURN
	FND

SUBROUTINE STEST(IOPT)

File STEST.FOR begins on the previous line.

Version Date: September 27, 1984 Paul S. Weiss

Subroutine STEST starts or continues the Test Mode of TGS operation, running the TGS over continuously without interruption.

The argument IOPT determines whether Test Mode is to be continued or started over, as follows:

IOPT=0 Start Test Mode IOPT=1 Continue Test Mode

WARNING: Subroutine STEST is the largest segment of overlay region 2. Enlarging it will directly increase the size of SANG.

INTEGER\*4 JDATA COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG COMMON /SDONE/ ISWFLG, JDATA(20) COMMON /CHADEF/ CHADEF(20,10)

IF (IRNFLG.NE.O) GO TO 8800

C

С

С

**TYPE 100** FORMAT(' Test Mode? Are you sure?') 100 ACCEPT 1000, ANS FORMAT(A1) 1000 IF (ANS.NE.'Y') GO TO 8900 TYPE \*, 'Change display window for test mode?' ACCEPT 1000, ANS IF (ANS.EQ.'Y') CALL SCHDWN(1) IF(IOPT.NE.O) GO TO 3000 С Ը Շ IF STARTING TEST MODE, START FLAG COUNTER AT O, AND QUERY AS TO WHETHER OR NOT TO CLEAR OLD TEST MODE DATA, IF THERE IS ANY С IF (TSTCUR.EQ.O.) GO TO 3000 -TYPE \*,'Zero previous Test Mode Data?' ACCEPT 1000,ANS IF (ANS.EQ.'Y') CALL SZERO(3) TSTCUR = 0.С С START HERE FOR BOTH STARTING AND CONTINUING TEST MODE C 3000 TSTFLG = 1.TSTCUR = TSTCUR + .00001ANGLE = TSTANG + TSTCUR INTERV = INTACQС С QUERY AS TO DETECTOR POSITION, ETC. AND START TGS IN С SUBROUTINE SSTART С CALL SSTART(1) GO TO 9000 C С IF TGS IS ALREADY RUNNING, SHOW THIS AND ABORT С 8800 **TYPE 8810** 8810 FORMAT(' TGS is running') С С IF TEST MODE IS CANCELLED AUTOMATICALLY OR BY USER, SHOW THIS С 8900 TYPE 8910,7,7 8910 FORMAT(' Abort Test Mode',2A1) RETURN 9000 END SUBROUTINE SVELO С File SVELO.FOR begins on the previous line. С

Version Date: July 20, 1984

С

Paul S. Weiss C С С Subroutine SVELO determines the beam velocity from the Ċ 45 degree crossing of the laser. С C COMMON /LASSET/ LASDWL, FRQLAS, LASSTP, FSRLAS, LASCHA COMMON /VELO/ DELF, HALFH, DELPIP, VELB, SPDRID **TYPE 110** 100 FORMAT(' Enter the distance between the 45 and 90 ' 110 1,'degree peaks.') ACCEPT \*, DELF IF(DELF.LE.O.) G0 T0 100 200 TYPE 210 FORMAT(' Enter the distance between frequency standard' 210 1,' mark.') ACCEPT +, DELPIP IF(DELPIP.LE.O.) GO TO 200 300 **TYPE 310** FORMAT(' Enter the half width of the 45 degree peak.') ACCEPT \*,HALFH 310 IF(HALFH.EQ.O.) G0 T0 300 C. DETERMINE CONVERSION OF FREQUENCY TO VELOCITY С С FTOV = (3.00E+10 \* 1.41421) / (FRQLAS \* 30.)С С DETERMINE MEASURED DISTANCE TO VELOCITY С DTOV = FTOV \* FSRLAS / DELPIP С С DETERMINE VELOCITY С VELB = DELF + DTOV TYPE 2010, VELB 2010 FORMAT(' The beam velocity is ',F10.0,' cm/sec.') С DETERMINE SPEED RATIO С С SPDRIO = DELF / HALFH IF(SPDRI0.GT.O.) GO TO 5000 SPDRIO = 0.TYPE \*, 'The speed ratio was not measured.' GO TO 9000 TYPE 5010, SPDRID 5000 5010 FORMAT(' The speed ratio is ', F6.3) 9000 RETURN END

SUBROUTINE SWRITE(IFLAG) File SWRITE FOR begins on the previous line. Version Date: September 4, 1984 Paul S. Weiss Subroutine SWRITE writes the raw data to logical unit 9, opened in subroutine SFILE. The argument of SWRITE is IFLAG, and it is a combination of all of the flags as determined in SDONE. LOGICAL \*1 ITIME(8) COMMON /FNAME/ INAME(7), ISTORE COMMON /SETS/ GATE(16), TIM(16), EXCFRC(2), CLOCK, INTACQ, NCHAN COMMON /SCR/ SCR(20) CALL TIME(ITIME) WRITE(9,100, ERR=9200) (ITIME(J), J=1,8), IFLAG 1, (SCR(I), I=1, ((2\*NCHAN)+2))FORMAT(1X, 8A1, 1X, I4, 1X, F9.5, 1X, F7.2, 2(2(1X, F9.0), :),/ 1,2(4(2(1X,F9.0),:),/)) ISTORE = ISTORE + 19000 RETURN 9200 TYPE 9210, (INAME(I), I=1,7),7,7,7,7 9210 FORMAT(' Error writing to file ',7A2,4A1) GO TO 9000 END SUBROUTINE SZERO (IOPT) File SZERO.FCR begins on the previous line. Version Date: October 12, 1984 Paul S. Weiss Subroutine SZERO zeroes some or all of the accumulated data in array DATA, sets NUMANG and MANGLE accordingly. The argument of SZERO specifies what angles to zero as follows: IOPT=1 Zero all data. IOPT=2 Zero data over a specified angular range. IOPT=3 Zero all Test Mode data. COMMON /CUR/ ANGLE, INTERV, NUMANG, MANGLE, IRNFLG COMMON /TSTSET/ TSTANG, NUMTST, TSTCUR, TSTFLG COMMON /DATA/ DATA(256,20), WORK1(256), WORK2(256)

NUMZER = 0

CCCCCCCCCCC

100

С

C C

Ċ

CCCCCCCCC

С

	<b>C</b>	GO TO(1100,2100,3100),IOPT
	C	ZERO ALL DATA
	1100	TYPE *,'Zero Data Arrays? Are you sure?' ACCEPT 1000,ANS
	1000	FORMAT(A1) IF(ANS.NE.'Y') GO TO 9000 NUMZER = NUMANG GO TO 5000
		ZERO DATA OVER A SEGMENT
	2100 2110	TYPE 2110,'first' FORMAT(/,' Enter ',A5,' angle of segment to zero') ACCEPT *,ANBE TYPE 2110,'last' ACCEPT *,ANFN TYPE 2210 ANBE ANEN
	2210	FORMAT(' Zero data from ',F9.5,' to ',F9.5 1,' degrees? Are you sure?') ACCEPT 1000,ANS IF(ANS.NE.'Y') GD TO 9000
	c	D0 2400 I=1, NUMANG
	C C C	HAVE ALL THE ANGLES BEEN LOOKED AT? IF SO, LEAVE LOOP
	2230 C	IF(I+NUMZER.GT.NUMANG) GO TO 2500
	Č ·	IS DATA WITHIN SEGMENT?
		IF (DATA(I,1).LT.ANBE) GO TO 2400 IF (DATA(I,1).GT.ANFN) GO TO 2500 DO 2300 J=1,20 DO 2250 IK=I,NUMANG-1 DATA(IK,J) = DATA(IK+1,J)
	2250 2300	CONTINUE CONTINUE NUMZER = NUMZER + 1
×	C C C C	SINCE DATA HAS BEEN SHIFTED, STAY AT THIS POSITION TO CHECK NEXT POINT
	2400	GD TO 2230 CONTINUE GD TO 6000
	C	SHIFT LATER DATA INTO PLACE
С IF (NUMZER.EQ.O) GO TO 8000 2500 DO 2550 IK=NUMANG-NUMZER, NUMANG-1 DO 2540 J=1,20 DATA(IK, J) = DATA(IK+1, J)2540 CONTINUE 2550 CONTINUE GO TO 6000 С C ZERO DATA TAKEN IN TEST MODE С 3100 **TYPE 3110** FORMAT(' Zero all Test Mode Data? Are you sure?') 3110 ACCEPT 1000, ANS IF (ANS.NE. 'Y') GO TO 9000 С DO 3500 I=1, NUMANG С С HAVE ALL THE POINTS BEEN CHECKED? С IF SO, LEAVE LOOP С 3200 IF (I+NUMZER.GT.NUMANG) GO TO 5000 С С IS THIS DATA FROM TEST MODE? С IF (INT(((20.\*DATA(I,1))-FLOAT(INT(20.\*DATA(I,1)))) 1\*1.E5).LT.1) GO TO 3500 С Č SHIFT DATA INTO SPACE VACATED BY REMOVED TEST MODE ANGLE С IF(I.GT.NUMANG-1) GO TO 3450 D0 3400 IK=I, NUMANG-1 DO 3300 J=1,20 DATA(IK, J) = DATA(IK+1, J)3300 CONTINUE 3400 CONTINUE 3450 NUMZER = NUMZER + 1 GO TO 3200 3500 CONTINUE С С ZERO TEST MODE ANGLE COUNTERS FOR OPTIONS 1 AND 3 С 5000 T= TSTFLG TSTFLG = 0.IF(T.GT.O.) TSTFLG = 1.TSTCUR = 0.C C EXECUTE FOR EACH OPTION: C C ZERO LAST SECTION OF DATA THAT IS NO LONGER USED

6000	IF(NUMZER.EQ.O) GO TO 8000
	DO 6500 IK=NUMANG-NUMZER+1, NUMANG
	D0 6400 J=1,20
	DATA(IK, J) = 0.
6400	CONTINUÉ
6500	CONTINUE
8000	TYPE 8010, NUMZER, 7
8010 C	FORMAT(' Data at ',I3,' angles zeroed.',A1)
Ċ C	REDUCE THE NUMBER ANGLES TO THE NEW LOWER VALUE
	NUMANG = NUMANG - NUMZER
	MANGLE = MANGLE - NUMZER
9000	RETURN END

ANGULAR SCAN ROUTINES .TITLE File SANGLE.MAC begins on the previous line. ; . IDENT /103084/ Version Date: October 30, 1984 ;Paul S. Weiss ;Angular Scan Routines to run the Timer-Gater Unit with program SANG. ;FORTRAN GLOBALS: .GLOBL SCHDEF, SCOMM, SDISP, SDPOL, SINIT, SPRII, SSTART ;MACRO GLOBALS: . GLOBL CLSHTR, DISPL . GLOBL ENSPEC, EXSPEC, INIT . GLOBL LCASE, MOT, OPSHTR, SCASUB . GLOBL SCLEAR, SEXIT, SGO, STOPD, UCASE .MCALL .INTEN, .PRINT . ENABL LC ; ì .SBTTL ADDRESSES JS₩=44 DATAL0=166000 DATAHI=166002 SCAVEC=504 SCAPRI=506 LASVEC=510 LASPRI=512 M0TCSR=167000 SCAL0=167040 SCAL1=167042 SCAL2=167044 SCAL3=167046 SCALAM=167074 LASE0=167100 LASE1=167102 LASE2=167104 LASE3=167106 LASLAM=167134 STATU1=166010 STATU2=166012 TTRCSR=177560 LFCSR=167760 TGCSR=167770 TG0TBF=167772

TGINBF=167774
TGVECA=350
TGAPRI=352
TGVECB=354
TGBPRI=356
AI0DAB=170406
DSPCSR=170440
DSPXYB=170442
DSPPCR=170444
DSPCAR=170446
DSPINV=340
DSPPRI=342

;

## . PAGE

## .SBTTL MACRO SUBROUTINES

;Subroutine INIT loads the interrupt vectors with the addresses ;of the interrupt service routines, and enables the Request B (Timer ;countdown completed) and Scaler (count overflow) interrupts.

INIT: RESE

RESEI		
BIS	#100, <b>0</b> #TTRCSR	;RESET TERMINAL INPUT ENABLE
TST	(R5) +	SKIP ARGUMENT COUNTER
MOV	(R5) + , 0#SFLAG	LOAD END COUNTDOWN FLAG ADDRESS
MOV	(R5),0#JDATAD	LOAD DATA ARRAY ADDRESS
MOV	#TGINTÄ,0#TGVECA	
MOV	#O', <b>Q</b> #TGAPRI	;KEEP PRIORITY LOW TO ENABLE INPUT
MOV	#TĠIŇTB, <b>0</b> #TGVECB	
MOV	#200, <b>0</b> #TGBPRI	
MOV	#O, <b>O</b> #DŜPCSR	CLEAR DISPLAY CSR
MOV	#DŚPÏNT, <b>0</b> #DSPINV	LOAD DISPLAY INTERRUPT VECTOR
MOV	₩O, <b>Q</b> #DSPPRI	AND PRIORITY
JSR	PC, SCLEAR	CLEAR SCALERS
MOV	#0, <b>0</b> #STATU1	
MOV	#O, <b>O</b> #STATU2	
MOV	#SCAINT,0#SCAVEC	;LOAD SCALER OVERFLOW INTERRUPT
MOV	#0.0#SCAPRI	KEEP SCALER PRIORITY I OW
MOV	#SCAINT. 9#LASVEC	LOAD LASER FLUORESCENCE SCALER
		INTERRUPT VECTOR
MOV	#O.Q#LASPRI	KEEP THIS PRIORITY LOW TOO
BIC	#301.0#STATU1	DISABLE SCALER INTERRUPT
MOVB	#24. O#SCALAM	
MOVB	#24.,0#LASLAM	
MOV	₩O, <b>Q</b> #TĞFLAG	CLEAR INTERRUPT FLAG WORD
MOV	#OÍ <b>O</b> #TGCSR	CLEAR TIMER-GATER CSR
BIS	#40,0#TGCSR	ENABLE B INTERRUPT
TST	<b>Ö</b> #TĠIŇBF	SEND DATA TRANSMITTED PULSE

MOV #0,LFCSR ;DISABLE LASER FLUORESCENCE ; PARALLEL LINE UNIT, AND CLOSE ; SHUTTER PC RTS Subroutine SEXIT disables all the interrupts in the system to allow ;termination of the program. SEXIT: MOV #0,0#STATU1 ;TURN OFF QUAD SCALER #0,0#STATU2 MOV MOVB #24.,**0**#SCALAM MOVB #24.,0#LASLAM MOV #O,O#TGCSR ;TURN OFF TIMER-GATER #0,0#DSPCSR MOV ;TURN OFF DISPLAY RTS PC ;Subroutine ENSPEC puts the monitor in special screen mode ENSPEC: BIS #10000,0#44 ;ENTER SPECIAL SCREEN MODE RTS PC Subroutine EXSPEC takes the monitor out of special screen mode: EXSPEC: BIC #10000,0#44 ;LEAVE SPECIAL SCREEN MODE RTS PC Subroutine LCASE enables lower case input from terminal LCASE: BIS #40000,0#44 RTS PC · ;Subroutine UCASE disables lower case input from terminal UCASE: BIC #40000,0#44 ΡC RTS ;Subroutine STOPD stops the display. STOPD: ;CLEAR DISPLAY INTERRUPTS BIC #103,0#DSPCSR RTS PC ;Subrouine SGO starts the Timer-Gater Module counting down.

(R5) + SGO: TST SKIP ARGUMENT COUNTER MOV (R5) + . 0#SFLAG :LOAD END COUNTDOWN FLAG ADDRESS MOV (R5),Q#JDATAD ;LOAD DATA ARRAY ADDRESS PC, SCLEAR **JSR** MOV ; ENABLE REQUEST B INTERRUPT #40,0#TGCSR MOV #O, O#TGOTBF ;SEND NEW DATA READY PULSE RTS PC ;Subroutine MOT sends pulses from a CAMAC pulse generator ; to a stepping motor controller, which turns a polarization :rotator. ;The argument of MOT is an integer, positive for clockwise ;rotation, and negative for counterclockwise rotation. MOT: MOVB #1,0#MOTCSR ;READ MOTOR STATUS WORD #1,**0**#DATALO ; IS LAM SET? BIT BEQ MOT ;NO,WAIT TST (R5)+ SKIP ARGUMENT COUNT MOV Q(R5),R1 ;GET ARGUMENT, NUMBER OF STEPS TST ; IS R1 NEGATIVE? **R1** BGE 30\$ ; NO NEG **R1** ;YES-MAKE POSITIVE BIS #100000,R1 ;SIGN BIT = 130\$: BEQ 40\$ ;DO NOT LOAD A ZERO MOV R1,0#DATALO MOVB #16.,**0**#MOTCSR ;LOAD COUNTER 405: RTS PC Subroutine DISPL loads the display buffer for the DMA display unit. ;It must be called with two arguments: DSPAD: The address of a single precision integer array, prepared to be displayed as a point plot, with x and y coordinates stored alternately. All values range from O to 4095. The address of this array is stored in DSPAD. NUMDSP: The number of points to be displayed. DISPL: BIC #103,0#DSPCSR ; DISABLE DISPLAY INTERRUPTS TST ;SKIP ARGUMENT COUNT (R5)+ MOV (R5) + , 0#DSPAD STORE ADDRESS OF DISPLAY BUFFER MOV Q(R5), Q#DSPPCR ;LOAD POINT COUNT REGISTER MOV ;LOAD NUMBER OF POINTS REGISTER 0(R5),0#DSPNPT ;FOR FUTURE REFERENCE

	NEG MOV BIS RTS	0#DSPPCR 0#DSPAD,0#DSPCAF #103,0#DSPCSR PC	;TWO'S COMPLEMENT CODING OF PCR ;INITIALIZE CURRENT ADDRESS REGISTER ;ENABLE DISPLAY INTERRUPT, AND START	
; ; ;Subrout ;enable:	tine SCLI s Scaler	EAR clears the ei interrupts.	ight Quad Scaler channels, and	
ŚCLEAR:	BIC MOVB MOVB MOVB MOVB MOVB MOVB MOVB BIS MOVB BIS MOVB BIS	#301,0#STATU1 #24.,0#SCALAM #2,0#SCAL0 #2,0#SCAL0 #2,0#SCAL1 #2,0#SCAL2 #2,0#LASE0 #2,0#LASE0 #2,0#LASE1 #2,0#LASE2 #2,0#LASE3 #301,0#STATU1 #26.,0#LASLAM #26.,0#LASLAM	;DISABLE SCALER OVERFLOW INTERRUPTS ;REENABLE SCALER INTERRUPTS	
; ; ;Subrout ;them.	tine SCAS	SUB reads the six	Quad Scaler channels in use, then clears rned in a double precision integer	
;array:	JDATA.			
;SCASUB ;	is call	ed from TGINTB		
SCASUB:	MOV BIC MOVB MOVB	R5,-(SP) #301,0#STATU1 #24.,0#SCALAM #24.,0#LASLAM	;SAVE R5 ON STACK ;DISABLE SCALER INTERRUPTS	
, LOAD THE ADDRESS OF JDATA INTO R5 MOV 0#JDATAD,R5				
; FIRST, EACH CHANNEL IS READ IN THROUGH DATALO ;AND DATAHI, AND THEN CLEARED, THE RESULT IS STORED IN ;THE ARRAY JDATA(1-4) MOVB #2,0#SCALO JSR PC,VALUE MOVB #2,0#SCAL1 JSR PC,VALUE				

390

ì

MOVB	#2,0#SCAL2
JSR	PC, VALUE
MOVB	#2,0#SCAL3
JSR	PC, VALUE

;

HERE THE VALUE OF THE LASER FLUORESCENCE SCALER CONTENTS ARE READ INTO DOUBLE PRECISION INTEGER ARRAY JDATA

MOVB	#2,0#LASE0	
JSR	PÇ, VÄLUE	
MOVB	#2,0#LASE1	
JSR	PC,VALUE	
MOVB	#2,0#LASE2	
JSR	PC, VALUE	
MOVB	#2,0#LASE3	
JSR	PC, VÄLUE	
MOV	(SP)+,R5	
RTS	PC	

; ;Subroutine VALUE stores the contents of one channel into one of ;the elements of a double precision integer array.

;RETURN

;RESTORE R5 FROM STACK

VALUE is called from SCASUB

ALUE:	MOV MOV	DATALO,(R5)+ DATAHI,(R5)+
	RTS	PC

Subroutine WAIT maintains control while the Timer-Gater is reset. WAIT waits for an interrupt from the Timer-Gater Module, then

NAIT:	TST	O#TGFLAG	HAS THERE BEEN AN INTERRUPT?
	BNE	RÊQ	, IF SO, RETURN
	BR	WAIT	, IF NOT, LOOP
REQ :	MOV	#O,0#TGFLAG	REMOVE REQUEST A FLAG
-	RTS	PC .	;RETURN TO COMAND VIA REQIN & START

;Subroutine LASDAB sends a value to Digital to Analog converter ;channel B. This is used to scan the laser.

ASDAB:	TST	(R5)+	;SKIP	ARGUMENT C	OUNT
• •	MOV	Q(R5),Q#AIODAB	;LOAD	D/A VALUE	•
	RTS	PC			

Subroutine OPSHTR opens the shutter to the 45 degree laser crossing; of the molecular beam by setting the CSRO bit of the secondary DRV11; Parallel Line Unit.

OPSHTR: MOV #1,**0**#LFCSR ;OPEN SHUTTER RTS PC ; ; ; ;Subroutine CLSHTR closes the shutter to the 45 degree laser crossing

;of the molecular beam by clearing the CSRO bit of the secondary DRV11 ;Parallel Line Unit.

CLSHTR: MOV #0,0#LFCSR ;CLOSE SHUTTER RTS PC ; ; .PAGE .SBTTL INTERRUPT SERVICE ROUTINES ;

;Interrupt Service Routine DSPINT is for the DMA display unit. ;The point counter is reinitialized, as is the current address register, ;and then the unit is restarted.

DSPINT:	BIC	#103,0#DSPCSR	;DISABLE DISPLAY INTERRUPTS
	.INTEN	0,PIC	;RUN REMAINING CODE AT PRIORITY O
	MOV	0#DSPNPT,0#DSPPCR	;LOAD POINT COUNTER
	NEG MOV BIS RTS	0#DSPPCR 0#DSPAD,0#DSPCAR #103,0#DSPCSR PC	;LOAD CURRENT ADDRESS REGISTER ;REENABLE DISPLAY INTERRUPTS ;AND START

;Interrupt Service Routine SCAINT is for Scaler count overflows. ;The Quad Scaler is cleared, and the end of the countdown occurs.

<b>SCAINT</b> :	MOV	#O, <b>O</b> #TGCSR	;DISABLE REQUEST B INTERRUPT
	BIS	#4,0#TGFLAG	SET OVERFLOW FLAG
	BIS	₩40,0#TGCSR	REENABLE REQUEST B INTERRUPT
	. INTEN	Ö,PÍC	
	JSR	PC, SCLEAR	CLEAR SCALER INTERRUPTS
	BIC	#301,0#STATU1	•
	MOVB	#24.,0#SCALAM	
	MOVB	#24.,0#LASLAM	۰.
	MOV	₩O,Q#TĞFLAG	;CLEAR FLAG FOR REQ A INTERRUPT
	TST	<b>Ö#</b> ŤGÏNBF	SEND A DATA TRANSMITTED PULSE

	BIS JSR	#100, <b>0</b> #TGCSR PC,WAIT	;ENABLE A REQUEST A INTERRUPT ;WAIT FOR A REQUEST A INTERRUPT
	MOV MOV RTS RTS	<b>0#</b> SFLAG,R5 #2,(R5) PC PC	;SET ISWFLG FOR OVERFLOW CONDITION
; ; ; ;			
;Interro ;the Tin	upt Serv mer-Gate	ice Routine TGINTB is for r Module Interface Board.	the Request B line of
;It read ;SCASUB	ds the Q , then s	uad Scaler's contents usi ets a flag and returns to	ng the MACRO subroutine > WAIT
TGINTB:	MOV JSR INTEN	#0,0#TGCSR PC,SCASUB	;DISABLE REQUEST B INTERRUPT ;READ AND CLEAR QUAD SCALER
. 3	MOV TST BIS JSR	#0,0#TGFLAG 0#TGINBF #100,0#TGCSR PC,WAIT	CLEAR FLAG FOR REQ A INTERRUPT SEND A DATA TRANSMITTED PULSE ENABLE A REQUEST A INTERRUPT WAIT FOR A REQUEST A INTERRUPT I.E. TG HAS RESET
	MOV MOV RTS	<b>0</b> #SFLAG,R5 #1,(R5) PC	;SET ISWFLG
; ;Interro ;of the ;contro	upt Serv Timer-G I to WAI	ice Routine TGINTA is for ater Interface Board. It T, which then returns to	- the Request A line t sets a flag, and returns COMAND.
ŤGINTA:	MOV BIS RTI	#0,0#TGCSR #1,0#TGFLAG	;DISABLE REQUEST A INTERRUPT ;SET REQUEST A FLAG BIT
;			
, <b>,</b>	. PAGE . SBTTL	ARGUMENT BLOCKS, VARIABL	ES, ETC.
JDATAD: SFLAG: TGFLAG: DSPNPT: DSPAD:	. WORD . WORD . WORD . WORD . WORD . WORD		

IFile CNSANG.COM !Compilation File for Angular Scan Program SANG !Version Date: October 30, 1984 **!Paul S. Weiss** FORTRAN/NOLINENUMBERS SANG SANG contains the Main program SANG, and subroutine SCOMM FORTRAN/NOLINENUMBERS SACCUM FORTRAN/NOLINENUMBERS SCANG FORTRAN/NOLINENUMBERS SCHDEF FORTRAN/NOLINENUMBERS SCHOSP FORTRAN/NOLINENUMBERS SCHOWN FORTRAN/NOLINENUMBERS SCHLOP FORTRAN/NOLINENUMBERS SCHMAD FORTRAN/NOLINENUMBERS SCHMAP FORTRAN/NOLINENUMBERS SCHNRM FORTRAN/NOLINENUMBERS SCHPAR FORTRAN/NOLINENUMBERS SCHPOL FORTRAN/NOLINENUMBERS SCHTST FORTRAN/NOLINENUMBERS SCHVEL FORTRAN/NOLINENUMBERS SCMLST FORTRAN/NOLINENUMBERS SCOMLC FORTRAN/NOLINENUMBERS SCOMLD FORTRAN/NOLINENUMBERS SCOMLG FORTRAN/NOLINENUMBERS SCOMLL FORTRAN/NOLINENUMBERS SCOMLM FORTRAN/NOLINENUMBERS SCOMLN FORTRAN/NOLINENUMBERS SCOMLP FORTRAN/NOLINENUMBERS SCOMLR FORTRAN/NOLINENUMBERS SCOMLS FORTRAN/NOLINENUMBERS SCOMLT FORTRAN/NOLINENUMBERS SCOMLV FORTRAN/NOLINENUMBERS SCOMLW FORTRAN/NOLINENUMBERS SCOMLZ FORTRAN/NOLINENUMBERS SDISP FORTRAN/NOLINENUMBERS SDLONO FORTRAN/NOLINENUMBERS SDLOP FORTRAN/NOLINENUMBERS SDNRM FORTRAN/NOLINENUMBERS SDONE FORTRAN/NOLINENUMBERS SDPOL FORTRAN/NOLINENUMBERS SDSIN FORTRAN/NOLINENUMBERS SDTST FORTRAN/NOLINENUMBERS SFILE FORTRAN/NOLINENUMBERS SGRAPH FORTRAN/NOLINENUMBERS SGTSCR FORTRAN/NOLINENUMBERS SINIT FORTRAN/NOLINENUMBERS SLEAVE FORTRAN/NOLINENUMBERS SLOOP FORTRAN/NOLINENUMBERS SNUCH FORTRAN/NOLINENUMBERS SNORM

FORTRAN/NOLINENUMBERS SOUT FORTRAN/NOLINENUMBERS SOVER FORTRAN/NOLINENUMBERS SPOLAR FORTRAN/NOLINENUMBERS SPRII FORTRAN/NOLINENUMBERS SSAVE FORTRAN/NOLINENUMBERS SSHCUR FORTRAN/NOLINENUMBERS SSHDEF FORTRAN/NOLINENUMBERS SSHDSP FORTRAN/NOLINENUMBERS SSHLOP FORTRAN/NOLINENUMBERS SSHMD FORTRAN/NOLINENUMBERS SSHMP FORTRAN/NOLINENUMBERS SSHNRM FORTRAN/NOLINENUMBERS SSHPOL FORTRAN/NOLINENUMBERS SSHTST FORTRAN/NOLINENUMBERS SSHVEL FORTRAN/NOLINENUMBERS SSTART FORTRAN/NOLINENUMBERS STEST FORTRAN/NOLINENUMBERS SVELO FORTRAN/NOLINENUMBERS SWRITE FORTRAN/NOLINENUMBERS SZERO MAC SANGLE

!File CPSANG.COM !Copy command file for Angular Scan Program SANG ! !Version Date: October 30, 1984 !Paul S. Weiss !

COPY I:SANG.FOR 0:SANG.FOR COPY I:SACCUM.FOR D:SACCUM.FOR COPY I:SCANG.FOR 0:SCANG.FOR COPY I:SCHDEF.FOR 0:SCHDEF.FOR COPY I:SCHDSP.FOR 0:SCHDSP.FOR COPY I:SCHDWN.FOR O:SCHDWN.FOR COPY I:SCHLOP.FOR 0:SCHLOP.FOR COPY I: SCHMAD.FOR O: SCHMAD.FOR COPY I:SCHMAP.FOR D:SCHMAP.FOR COPY I: SCHNRM. FOR D: SCHNRM. FOR COPY I:SCHPAR.FOR 0:SCHPAR.FOR COPY I:SCHPOL.FOR 0:SCHPOL.FOR COPY I:SCHTST.FOR 0:SCHTST.FOR COPY I:SCHVEL.FOR 0:SCHVEL.FOR COPY I:SCMLST.FOR 0:SCMLST.FOR COPY I:SCOMLC.FOR 0:SCOMLC.FOR COPY I:SCOMLD.FOR 0:SCOMLD.FOR COPY I:SCOMLG.FOR O:SCOMLG.FOR COPY I:SCOMLL.FOR D:SCOMLL.FOR COPY I:SCOMLM.FOR O:SCOMLM.FOR

COPY I:SCOMLN.FOR O:SCOMLN.FOR COPY I:SCOMLP.FOR 0:SCOMLP.FOR COPY I:SCOMLR.FOR O:SCOMLR.FOR COPY I:SCOMLS.FOR 0:SCOMLS.FOR COPY I:SCOMLT.FOR O:SCOMLT.FOR COPY I:SCOMLV.FOR 0:SCOMLV.FOR COPY I:SCOMLW.FOR O:SCOMLW.FOR COPY I:SCOMLZ.FOR 0:SCOMLZ.FOR COPY I:SDISP.FOR 0:SDISP.FOR COPY I:SDLONO.FOR 0:SDLONO.FOR COPY I:SDLOP.FOR 0:SDLOP.FOR COPY I:SDNRM.FOR O:SDNRM.FOR COPY I:SDONE.FOR 0:SDONE.FOR COPY I:SDPOL.FOR 0:SDPOL.FOR COPY I:SDSIN.FOR 0:SDSIN.FOR COPY I:SDTST.FOR 0:SDTST.FOR COPY I:SFILE.FOR 0:SFILE.FOR COPY I:SGRAPH.FOR D:SGRAPH.FOR COPY I:SGTSCR.FOR O:SGTSCR.FOR COPY I:SINIT.FOR O:SINIT.FOR COPY I:SLEAVE.FOR 0:SLEAVE.FOR COPY I:SLOOP.FOR 0:SLOOP.FOR COPY I: SNORM. FOR O: SNORM. FOR COPY I: SNUCH, FOR O: SNUCH, FOR COPY I:SOUT.FOR 0:SOUT.FOR COPY I:SOVER.FOR 0:SOVER.FOR COPY I: SPOLAR.FOR 0: SPOLAR.FOR COPY I:SPRII.FOR 0:SPRII.FOR COPY I:SSAVE.FOR 0:SSAVE.FOR COPY I:SSHCUR.FOR O:SSHCUR.FOR COPY I:SSHDEF.FOR 0:SSHDEF.FOR COPY I:SSHDSP.FOR 0:SSHDSP.FOR COPY I:SSHLOP.FOR O:SSHLOP.FOR COPY I:SSHMD.FOR D:SSHMD.FOR COPY I:SSHMP.FOR D:SSHMP.FOR COPY I:SSHNRM.FOR 0:SSHNRM.FOR COPY I:SSHPOL.FOR 0:SSHPOL.FOR COPY I:SSHTST.FOR O:SSHTST.FOR COPY I:SSHVEL.FOR 0:SSHVEL.FOR COPY I:SSTART.FOR D:SSTART.FOR COPY I:STEST.FOR 0:STEST.FOR COPY I:SVELO.FOR 0:SVELO.FOR COPY I:SWRITE.FOR O:SWRITE.FOR COPY I:SZERO.FOR O:SZERO.FOR COPY I: SANGLE.MAC D: SANGLE.MAC COPY I: CNSANG.COM D: CNSANG.COM COPY I: CPSANG.COM 0: CPSANG.COM COPY I: LSANG.COM D: LSANG.COM

!File LSANG.COM !LINK command file for Angular Program !Version Date: October 30, 1984 **!Paul S. Weiss R** LINK SANG, SANG.MAP, =SANG, SANGLE/I// SACCUM/0:1 SCANG/0:1 SCHDEF/0:1 SCHDSP/0:1 SCHDWN/0:1 SCHLOP/0:1 SCHMAD/0:1 SCHMAP/0:1 SCHNRM/0:1 SCHPAR/0:1 SCHPOL/0:1 SCHTST/0:1 SCHVEL/0:1 SCMLST/0:1 SCOMLC/0:1 SCOMLD/0:1 SCOMLG/0:1 SCOMLL/0:1 SCOMLM/0:1 SCOMLN/0:1 SCOMLP/0:1 SCOMLR/0:1 SCOMLS/0:1 SCOMLT/0:1 SCOMLV/0:1 SCOMLW/0:1 SCOMLZ/0:1 SDISP/0:1 SDLONO/0:1 SDLOP/0:1 SDNRM/0:1 SDPOL/0:1 SDSIN/0:1 SDTST/0:1 SFILE/0:1 SGRAPH/0:1 SGTSCR/0:1 SINIT/0:1 SLEAVE/0:1 SOUT/0:1 SOVER/0:1 SPRII/0:1 SSAVE/0:1 SSHCUR/0:1

SSHDEF/0:1 SSHDSP/0:1 SSHLOP/0:1 SSHMD/0:1 SSHMP/0:1 SSHNRM/0:1 SSHPOL/0:1 SSHTST/0:1 SSHVEL/0:1 SSTART/0:1 SVEL0/0:1 SWRITE/0:1 SZER0/0:1 SDONE/0:2 SL00P/0:2 SNORM/0:2 SNUCH/0:2 SPOLAR/0:2 STEST/0:2 // \$SHORT

^C

1

398

# B. Data Acquisition for Time-of-Flight Measurements: The New Multichannel Scaler and Program TUF

1. The Purpose of the Multichannel Scaler

A multichannel scaler (MCS) is used for measuring the time-offlight of scattered atoms or molecules from a point of modulation (e.g. reaction or collision center, or modulating wheel, chopper, or laser) to a point of detection (e.g. ionizer and subsequent ion optics). This time is typically on the order of microseconds, as is the spread in the distribution of typical time-of-flight (TOF) spectra.

A multichannel scaler used for measuring TOF spectra counts a series of pulses in consecutive channels, where each channel corresponds to a well defined time interval after an initial trigger pulse. For example, after a trigger pulse arrives at the MCS, the first channel records data pulses from 0 to 1 usec, the second channel records data pulses from 1 to 2 usec, and the  $n\frac{th}{t}$  channel records data to have. When another trigger pulse arrives at the MCS, the first channel mediately begins again with the first channel adding the present data to that recorded after previous trigger pulses. The scan described above has a "dwell time" of 1 usec, that is, after each trigger pulse arrives at the MCS the channels are addressed sequentially with each channel recording data for 1 usec.

A 4096 channel scaler has been designed and built in collaboration with Jacques Millaud and Fred Vogelsberg of the LBL Department of Instrument Science and Engineering. The scaler is a double width CAMAC module, and can interface to any of the numerous computers which can be equipped with a CAMAC Dataway Interface (IEEE Standard 583-1975).<sup>2</sup> Only Digital Equipment Corporation LSI-11 series computers (namely 11, 11/2, 11/23, and 11/73)<sup>5</sup> have been programmed to run the MCS thus far. The specifications for the scaler are given below.

This multichannel scaler was built to replace a 256 channel scaler built by Randall Sparks and Vince Randolph in 1978.<sup>11</sup> Because of the advance in IC design, it was clear that a new scaler could be built with substantially better capabilities. The new scaler has 16 times the memory, 16 times the counting ability, can count over 3 times as long, has 6 times the time resolution, is half the size, and uses approximately one third of the power of the old scaler.

At the time of its design (1984), the wait for some of the most advanced chips was over 1 year. Thus, the choice of IC logic often was not at the cutting edge, but was sufficient for our space and time constraints. Even now, it is possible to see significant improvements that could be made in reducing the dead time, and increasing the time resolution and the maximum counting rate.

Table I. The CAMAC commands for the MCS.

Command Function

N.FO.AO Read MCS Control Status Register.

N.FO.A1 Read the number of triggers remaining.

N.FO.A3 Read the memory data.

N.F8.AO Test LAM.<sup>5</sup>

N.F18.AO Load MCS Control Status Register.

N.F18.A1 Load the number of triggers.

N.F18.A2 Load the dwell counts.

N.F18.A3 Load the memory address.

2. CAMAC Commands for the New Multichannel Scaler

Table I is a list of CAMAC commands which drive the multichannel scaler. In order to use these refer to the manual for the CAMAC crate controller in use. All CAMAC crate controllers have the ability to issue these commands, but each has its own protocol for doing so. All programs and examples below are written for an LSI-11 computer with a Kinetic Systems Model 3912-ZIG crate controller.<sup>10</sup> All load (F18) commands load values from the data register (low 16 bit word: "DATALO" in CAMAC parlance) into the appropriate MCS register. All read (F0) commands load values from the MCS into the data register (DATALO). A description of the various commands follows. Table II. Definitions of Control Status Register Bits. A "(1)" implies that the corresponding bit is set when the condition is true.

<u>Bit</u>	<u>Read/Write</u>	Meaning
0	R/W	LAM Enable (1).
1		Not used.
2	R/W	Read Data Enable (1).
3	R/W	Acquisition Enable (1).
(4	W	CLEAR, not implemented in hardware)
5 5	R W	Overflow Flag (1) Test mode. (1)
6	R	End Sweep (1).
7	R	LAM Flag (1).
8–15		Not Used.

Reading and Writing the Control Status Register:

The meaning of the bits of the Control Status Register (CSR) are shown in Table II. R refers to read only bits, W refers to write only bits, and R/W refers to bits which can be both read and written. If an attempt is made to write to a read only bit, the word being written is masked so that the bit is not written, and no error of any sort is generated. If a condition is true, the bit is set, as is shown by a "(1)" after the bit definition in table II, and below. If a bit is cleared a "(0)" is shown.

In order to start accumulating data for a TOF spectrum, bit 3 must be set, while bit 2 must be clear. This has the effect that data acquisition is enabled, and the MCS memories cannot be read. In program TUF, bit 0 is always set at this point to enable interrupts. An alternative to this method of detecting an end sweep is to repeatedly poll CSR bit 6 for an end of sweep. Bit 5 is used for a self test of the MCS unit. The 20 MHz clock is shunted in as the input to the scalers, and a constant signal is observed if the triggers are supplied. This is enabled in the Test command of program TUF. If connected, CSR bit 4 would provide a means to check the MCS with no data input (when 0) by gating off the data input, as well as a clear function (when 0). Normal operation and data acquisition would occur if bit 4 was set high (1). The CSR bits are set by first writing the appropriate word to the data word (DATALO), and then loading these into the MCS with the N.F18.AO command. This is shown in the MACRO subroutine BEGIN, in program TUF.

The interrupt-driven structure chosen leaves the program a freer hand to do more lengthy work between end sweeps without fear of significantly reducing the data acquisition rate. The disadvantage of the interrupt structure is that if the interrupts occur too often for the program to get through the necessary overhead, it will stop. This can be avoided by lengthening the time between interrupts, by increasing the number of triggers per sweep, and/or by reducing the trigger rate. This is particularly important in the case of cross-correlated data (discussed in chapter I), for which the preparation of the display involves a rather lengthy calculation (which can take several seconds). The control status register can be queried to determine whether a sweep has been completed, an overflow has occurred, or the MCS is set so that its memories can be read (bits 6 and 7, bit 6, or bit 2, respectively). This is done by issuing the N.FO.AO command and then checking the appropriate bits of the CAMAC data word (DATALO), or storing this value for future reference. This is done in the MACRO interrupt service routine MCSINT in program TUF.

Reading and Writing the Number of Triggers:

The number of triggers is loaded in the same way as the control status register, that is through the CAMAC data word (DATALO). The number of triggers must be converted to its two's complement before being loaded. The loading is done by the N.F18.A1 command as shown in MACRO subroutine LOADTR in program TUF. The maximum number of triggers is 65535, but the program TUF allows a maximum of only 32767 so that two byte integers can conveniently be used.

Without interfering with the data acquisition the number of triggers remaining in a sweep can be read out of the MCS with the N.FO.Al command. Issuing the command moves the number of triggers remaining into the CAMAC data word. In program TUF, this is only done in the case of a scaler overflow, as shown in the MACRO interrupt service routine MCSINT.

#### Write the Dwell Count:

The dwell count determines the time interval for which pulses are counted in each channel. The count is the number of 50 nsec intervals for which data collection is to occur for each channel. Thus the dwell count is simply the dwell time divided by 50 nsec. The dwell count is first loaded into the CAMAC data word, and then written to the MCS register using the N.F18.A2 command, as shown in the MACRO subroutine LOADDW in program TUF. There is no facility to read back the dwell count. The minimum dwell count is 3, and the maximum dwell count is 4095, giving minimum and maximum dwell times of 150 nsec, and 204.75 µsec, respectively. The maximum dwell time in program TUF is arbitrarily set at 200 µsec.

Load the Memory Address and Read the Memory Contents:

In order to read the memory contents, bit 2 of the control status register must be set as discussed above. Then, each memory address must be entered by loading the address of the channel to be read (0 to 4095, or 0 to 7777<sub>8</sub>) into the CAMAC data word (DATALO), and then issuing the N.F18.A3 command. Next, the contents of the memory are read out to the CAMAC data word by issuing the N.F0.A3 command. The CAMAC data word is then added to the appropriate array element stored in the controlling computer's memory. This is shown in MACRO subroutine READMS in program TUF. Note that there is no reason that the memory must be read out sequentially.

.E

#### Test LAM:

The interrupt (LAM -- for "look at me") flag can be tested by issuing the command N.F8.AO.<sup>12</sup> This causes the LAM to be asserted and if enabled causes an interrupt of the computer. This feature is not used in program TUF.

## 3. Multichannel Scaler Hardware Specifications

The above summary of CAMAC commands for use with the MCS has included many of its specifications, but a more complete list is given in table III. The scaler speed could be increased (by interchanging faster pin-compatible memories) and the dead time decreased<sup>13</sup> in a straightforward manner if this were found to be necessary, but at present counting rates in excess of 1 MHz are rare, and the current set-up is sufficient.

The schematics, circuit board layout, and a tape for the automatic wire wrap of the circuit board were generated by a CAD system developed by the Instrument Engineering group at LBL. A 2 A fuse and a 1 V drop diode (to change CAMAC +6V power to +5V) are soldered between the +6Vpower line from the CAMAC crate and the board's power in circuit. Otherwise, wire wraps interconnect all pins of the 80 IC sockets.

Table III. Multichannel Scaler Hardware Specifications.

Size:	Double width CAMAC module <sup>14</sup> with one wire-wrapped
	socketed circuit board.
Power:	+6V, 2 A supplied by CAMAC crate.
Dwell:	Minimum: 100 nsec per channel. <sup>15</sup>
	Maximum: 204.75 usec per channel. <sup>16</sup>
	Increment: 50 nsec per channel.

Triggers per sweep:

Minimum: 1.

Maximum: 65535.<sup>17</sup>

Increment: 1.

Sweeps: Under software control, no limit.

Dead Time: Between channels: <10 nsec.

Average: 25 nsec after previous data pulse.

Maximum: 50 nsec after previous data pulse.<sup>13</sup>

Maximum Counts per Channel:

Per Trigger: 32767.

Per Sweep: 32767.

Scaler Speed: 20 MHz (120 MHz internal scalers).

Memory Speed: 90 nsec.

Maximum Trigger Rate:

Nominally 900 kHz -- the time between trigger pulses must be at least 500 nsec + 4 dwell periods.

Table IV. Front Panel Connections and Indicators.

Inputs:	Data:	TTL (+5V), 50 nsec.				
	Trigger:	TTL (+5V), 500 nsec.				
	Trigger Gate:	TTL (+5V), Constant when trigger				
		is to be suppressed.				
Output:	RCO:	TTL (+5V), 50 nsec, on the				
		trailing edge of each dwell period.				

LED Indicators:

Power on:	0n	when	modu	ıle	is	powered.
LAM:	0n	when	LAM	is	set	<b>t.</b> .

Trigger Bits:

On when associated trigger word bit is 0. Off when associated trigger word bit is 1.

Inputs and Outputs

The front panel connections and LED indicators are summarized in table IV. The triggers supplied to the MCS should be TTL (+5V), and at least 500 nsec long. No data acquisition is done in the first 500 nsec as this time is used to set up the address counters, memories, and store data from a previous trigger if necessary. Data pulses should be TTL (+5V), and as close to 50 nsec long as possible. This ensures that the pulse is counted exactly once, as it will then straddle the rising edge of one 20 MHz clock pulse to which it will be synched. The triggers can be gated with a constant TTL voltage applied to the Trigger Gate input. The RCO output provides a 50 nsec TTL pulse out at the beginning of each dwell period.

Front Panel Indicators

The "power on" LED is on at all times that the CAMAC crate is powered. The LAM light is set when a LAM is set, and is cleared on reset. As triggers are supplied to the MCS, the front panel LEDs show the scaler counting down triggers to their preset value. After the preset trigger countdown, all LEDs are lit. When there is some number remaining to be counted, the lights summing up to the number remaining are <u>off</u>.

4. Multichannel Scaler Circuit Logic and Signal Descriptions

What follows is a description of the signals and IC chips on the circuit diagrams.<sup>18</sup> It is not necessary to understand this in order to use the MCS, but it would be helpful in modifying or repairing the MCS. Figures 2a and 2b schematically show the functions of the various IC groups as well as the common read and write lines they share. The "U" numbers below and in figure 2 refer to the chips used in the MCS. The acronymic names refer to signals that can be found on the circuit diagram. Table V lists the various signal names and their functions and/or derivations.

Fig. 2. (a) and (b) Schematic of the new LBL multichannel scaler. The "U" numbers refer to the IC numbers on the circuit diagram and layout drawings.<sup>18</sup>



XBL 863-787

Fig. 2(a)



XBL 863-786

Fig. 2(b)

Table V. The signals used in the circuit diagram of the  $MCS^{18}$  and their definitions and/or derivations.

ACNTCK [BCNTCK] Clocked data counts to Channel A [B] counters.

ACQEN Enables the start of data acquisition.

ADDR Address word (12 bit) for memory read out.

AGAT [BGAT] Square wave signals to control switching between A and B count channels -- period is twice the dwell time (AGAT = BGAT).

CLEAR When set data is allowed, when clear data is inhibited.

DATEN Delayed trigger signal -- when set enables data counters.

DL Delayed trigger signal which provides predata resets and configures an address latch.

DWELL Dwell count word.

- ENDSWP End of sweep -- set when trigger counters have completed count up from preset level to overflow.
- F Set after the first scan is completed, allows counters to be preloaded with memory data before counting the current data.
- FA, FB Levels set after trigger delay to provide phase of channel A [B] counter resets and write sequences.
- LAM Interrupt signal to CAMAC bus -- identical to CSR bit 7.

LAMEN LAM enable set by CAMAC -- identical to CSR bit 0.

OV Overflow flag, set on counter overflow -- identical to CSR bit 5.

PCP, PTRIG, PTW Signals at end of scan to ensure that only data from a complete dwell period are transferred to memory.

PEA [PEB] Load Counter A [B] signal.

RCO Dwell Counter Overflow, the falling edge of which ends one dwell period and starts another --- identical to front panel RCO out.

#### Table V, cont.

RD DATA Set by CAMAC to allow data readout -- identical to CSR bit 2. RSTA [RSTB] Reset for channel A [B] during the first scan of each sweep.

RTRIG, TDL Delayed trigger signals to configure counters and address for the following data scan.

TRIG External pulse required to start a scan.TRIGGER Trigger word loaded through CAMAC, preset in trigger counters.WEA [WEB] Write enable to transfer Counter A [B] data to memory.

The clear function of the MCS is currently not implemented. A <u>de facto</u> clear is made when the previous data is not read from the memories into the counters on the first trigger of each sweep. Acquisition is started with the signals ACQ, F, FA, and FB. The flipflops U34, U35, and U36 control the conditions after the first trigger and for subsequent scans.

The dwell timing is managed by counters U19, U20, and U21. They generate the clock-synchronized RCO signal pulses which drive the address advance in the scalers and indirectly drive the load, latch, read, and output functions of the counters and memories.

The data acquisition is accomplished with minimal dead time by having two identical sets of counters and memories. The two sets are toggled back and forth, so that while data is accumulated on one channel, the other channel is storing its previously acquired data and preparing for its next acquisition. The two channels are referred to as A and B. Note that all odd channels are in A and all even channels are in B. The addressing is simplified as the final bit of the address determines whether that address is in channel A or B. A special timing sequence has been worked out to allow the loading, latching, etc. of the data to the correct channel at all times.<sup>19</sup> The address counters are U45, U46, and U57. The addressing is accomplished with an asymmetric set of channel latches, buffers, and registers. Each of these is associated primarily with one channel, but services both channels. These are: the A channel address latches U50 and U51, the B channel address buffers U48 and U49, and the B channel address registers U52 and U53. Signals AGAT and BGAT drive the switching between channels A and B.

Multiplexers U32 and U33 generate: RSTA and RSTB to carry out the required resets on the first trigger of each sweep, PEA and PEB to load the counters, ACNTCK and BCNTCK to provide clock pulses to count the incoming data, and WEA and WEB to enable memory writing. The multiplexers are gated by AGAT. These signals control: the channel A counters U54, U55, U56, and U57, the channel A latches U58 and U59, the channel A memories U60 and U61, the channel B counters U62, U63, U64, and U65, the channel B latches U66 and U67, and the channel B memories U68 and U69. The input data pulses are shifted and latched by U80, U39, and U73. This switching avoids either missing a data pulse or counting one pulse twice, since only one data pulse can occur per clock pulse. The memories are currently 6116 type with a read or write cycle of 90 nsec. If the memories were exchanged for faster ones, the minimum dwell time could be lowered to 100 nsec per channel. $^{15}$ 

Trigger counters U22, U23, U24, and U25 accept four bits each of the trigger word, count up to the maximum possible (64k), and drive the front panel LEDs. Their overflow signals the end of sweep. The ENDSWP and LAM signals are set by flipflops U16 and U44.

If a trigger pulse comes before the end of a scan through the channels, a new scan is begun after 500 nsec. U28 and U31 generate signals RTRIG, DL, TDL, and DATEN which permit address and data resets, and synchronize the start of the next scan, as well as PTRIG, PCP and PTW which ensure that only data from the final complete dwell period is written into memory.

Communication with the CAMAC crate is accomplished through the CAMAC read lines R1, R2, R4, R8, and R16, write lines W1, W2, W4, W8, and W16, function codes F1, F2, F4, F8, and F16, and address lines A1, and A2, which are decoded and carried into the MCS. The input buffers for writing are U6 and U7, for instructions are U1 and U2, and for the status word bits LAMEN, ACQEN, DATEN, and TEST is U5. The readouts buffers are: U74 and U75 for the trigger word, U8, U9, U10, and U11 for data, and U12 for the status word. The status word (control status register in software terms) signals LAMEN, RD DATA, ACQEN, OV (overflow), ENDSWP, and LAM are latched at U12. The selected function commands (e.g. write trigger word -- F18.A1) control the destination of the data words DWELL, TRIGGER, and ADDR through common write lines with U4. The selected function commands also control the readout of the

data words (stored) CSR, DATA, and TRIGGER. Data readout of all 4096 channels takes approximately 10-20 msec.

## 5. Introduction to Program TUF

The program TUF drives the CAMAC multichannel scaler described above. A listing follows in section 6. It has been designed to run on DEC LSI 11, 11/2, 11/23, and 11/73 computers<sup>5</sup> equipped with the MCS in a CAMAC crate driven by a Kinetic Systems 3912–Z1G crate controller.<sup>10</sup> A direct memory access display driver (Data Translation DT2771<sup>20</sup>) is required on the Q-Bus of the LSI-11. This drives an oscilloscope or equivalent, and is referred to only as "display" below. The oscilloscope is not necessary for program execution, but is extremely useful at all times. If plotting is to be done, a Hewlett Packard 7470A or 7475A plotter<sup>21</sup> is required. The operating system used is RT-11 and the version of TUF shown is in use with RT-11 version 5.0.

The command structure is very similar to that of program SANG described in Appendix A. Execution begins with the main routine TUF, but after initialization and setting defaults in subroutine TINIT, control is passed to subroutine TCOMM where it remains except for subsequent subroutine calls until program completion.

All commands from the keyboard are processed by TCOMM using DEC function subroutine ITTINR as in program SANG. Commands are one or two letters followed by a carriage return. The one exception to this is the use of arrows on VT52 and VT100 compatible terminals. The arrows are used to adjust the display window, and no carriage return is required. It is possible to avoid the use of the arrows by using the normal structured commands.

The program is heavily commented throughout. At the beginning of each routine the file in which it is found, a version date, and a brief resume of its function are given. The meanings of all of the COMMON variables are listed in the first two pages of the main routine TUF. The commands are listed and explained throughout subroutine TCOMM, and on the help screens provided by TCMLST and all nine of the TCOML\* subroutines. The use of the arrows for changing the display is explained in subroutines TCOMM and TCDISP, and on the help screen provided by subroutine TCOMLD.

All input to the program is from the console terminal. All output is in ASCII form and when written to a disk file can be easily manipulated with the normal DEC editors. In data files, a header section is terminated with a control-C, and followed by rows of 8 data points listed as real numbers. A program entitled PAN is currently used to process the data.

All hardware functions are carried out by MACRO subroutines located in the file XBL.MAC. These are all modular and are often called from more than one FORTRAN subroutine. This file also contains the hardware initialization routine INIT and the interrupt service routines MCSINT and DSPINT for the MCS and the display driver, respectively. All MACRO subroutines are declared Global to allow access from

the FORTRAN subroutines. Note that all calling subroutines must also be declared Global.

The program TUF is far too large for the LSI-11 to load all at once (as with program SANG) so that it is overlaid with a total of 32 segments using the LINKer command file LTUF.COM. Also, a number of concessions have been made to decrease the size of the program, such as separating all of the help screens into separate files TCMLST.FOR, and TCOML\*.FOR, and overlaying them into different segments. The listing shown is within a few words of the limit of the system, and great care should be taken if any section of it is expanded. As with program SANG, all of the FORTRAN routines are compiled with the /NOLINENUMBERS option as shown in the compilation command file CNLTUF.COM, and the LINKed with the \$SHORT error messages as shown in the LINKer command file LTUF.COM so as to save space. If a section absolutely must be enlarged and will affect the total program size, the command structure of TCOMM could be split up into overlaid subroutines to save some room. Alternatively, all processors running TUF could be converted to LSI-11/23 or LSI-11/73 CPUs with at least 128k bytes memory, and some of the data arrays could be declared VIRTUAL.

# 6. Listing of Program TUF

On the following pages program TUF is listed. The order of the files is exactly that done by the backup/update/printout command file COPTUF.COM, which is included as well.
с		PROGRAM	TUF File TUF.FOR begins on the previous line.
C C C C		Version Paul S.	Date: November 10, 1985 Weiss
		This ma two init control	in routine just establishes all the COMMON areas, calls tialization routines, starts the display, and transfers over to subroutine TCOMM.
		INTEGER COMMON COMMON COMMON COMMON COMMON COMMON COMMON	4 JTRIG /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT /SWFLAG/ ISWFLG,JTRIG /TERM/ ITERM /TITANG/ TITLE(10),ANGLE
	****	******	DEFINITIONS OF /COMMON/ VARIABLES AND ARRAYS
C	соммо	N Block ,	/DATA/ Contains the data and work arrays.
		ACCUM(I) XCORR(I) WORK(I) NEW(I) NSWEEP	) The accumulated data for channels I=1-4096. ) The cross-correlated data for I=1-256. The work array for display I=1-256. The new data array for channels I=1-4096. Not used.
	СОММО	N Block	/DISP/ Contains the Display array and parameters.
C		IDSP(I)	The array for display by the Direct Memory Access Display Driver (Data Translation DT2771 or equivalent)
Ĉ		INEW	The display new data flag. INEW=0 => Display accumulated
Č C		ICORR	The cross-correlation display flag. ICORR=0 => Display uncorrelated data. ICORR=1 => Display cross-correlated
C C C		IBACK	data. The display background subtraction flag. IBACK=0 => Subtract background from displayed data. IBACK=1 =>
C C C C C		NCORR IADV	Display data with background included. The number of channels per cross-correlation channel. The number of channels after ICHAN to begin displaying data
C		IRES	The display resolution in channels.

IQIK The quick display flag. If IQIK=0, the data is displayed С before restarting the MCS (this prevents the program from C C C C C C C C getting hung up if the number of triggers is small). If IQIK=1, the data is accumulated and displayed while the MCS is counting down (for maximum effeciency, and minimal overhead delay). С С С COMMON Block /MCSMAX/ Contains the minima and maxima for various С MCS settings. С С DWLMAX The maximum dwell time. С DWLMIN The minmum dwell time. С MAXSWP The maximum number of sweeps. С MAXTRG The maximum number of triggers per sweep. С С C COMMON Block /SETS/ Contains the MCS settings. C С DWELL The dwell time in usec. С The initial channel (for display and output only, all channels ICHAN С are recorded and kept in memory). С ISWEEP The number of sweeps requested (a sweep is of ITRIG triggers). С ITRIG The number of triggers per sweep. C EDELAY The external delay in usec. C The number of channels (for output only, see ICHAN above). NCHAN C IOFSET The cross-correlation offset, in channels. С IEXT The external dwell flag. Not currently used. С С С COMMON Block /SWFLAG/ Contains the sweep counter, and the number of triggers measured. С С С The sweep counter. ISWFLG<0 for no new data to be ISWFLG С examined => remain in Command Mode. ISWFLG>0 => С The number of sweeps remaining. ISWFLG=0 for the С acquisition has been completed. C C JTRIG The triggers measured counter in INTEGER\*4. If JTRIG(0, an overflow has occurred. Č C С COMMON Block /TTERM/ Contains the terminal type specification. С С ITERM The terminal type: ITERM=0 for unknown, ITERM=52 for С VT52 compatible, and ITERM=100 for VT100 compatible С С C COMMON Block /TITANG/ Contains the title and detector angle for output purposes only. С С

	TITLE(I) The title, for output up to 40 ASCII characters. ANGLE The detector angle.
C C C C	
C C	INIT SETS UP INTERRUPT VECTORS
с	MCS = INIT(ISWFLG, JTRIG)
	TINIT SETS UP THE INITIAL VALUES OF ALL: MCS PARAMETERS (STORED IN /SETS/) TERMINAL TYPE (STORED IN /TERM/) DISPLAY PARAMETERS (STORED IN /DISP/) ALLOWED RANGE OF MCS PARAMETERS (STORED IN /MCSMAX/)
	CALL TINIT
	TINIT HAS JUST ASKED IF DEFAULT VALUES ARE SATISFACTORY GET ANSWER
1000 1100 C	ACCEPT 1000,ANS FORMAT(A1) IF(ANS.EQ.'Y') GO TO 1100 CALL TCHAN(O) CALL TDEF(O) CALL TTERM(O) CALL TTERM(O) CALL TDISP TRANSFER CONTROL TO ROUTINE COMMAN FOR ACTUAL USE OF TOF CALL TCOMM END
C	SUBROUTINE TCOMM Second routine of file TUF.FOR.
	Version Date: November 8, 1985 Paul S. Weiss
	Subroutine TCOMM maintains control of operations for the program.
	It accepts 1 or 2 character commands from the user, then calls the appropriate subroutines.
C	The arrows of VT52 and VT100 compatible terminals are also used.

С These change the window data displayed at any time. Ċ INTEGER COM1, COM2, COMT INTEGER+4 JTRIG COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /MCSMAX/ DWLMAX, DWLMIN, MAXSWP, MAXTRG COMMON /SETS/ DWELL, ICHAN, ISWEEP. ITRIG. EDELAY.NCHAN. IDFSET. IEXT COMMON /SWFLAG/ ISWFLG, JTRIG COMMON /TERM/ ITERM COMMON /TITANG/ TITLE(10), ANGLE С č THE FOLLOWING IS A TABLE OF THE ASCII CHARACTERS USED AS COMMANDS С DATA IA, IB, IC, ID, IE /65, 66, 67, 68, 69/ DATA IF, IG, IH, II; IL, IM, IN /70, 71, 72, 73, 76, 77, 78/ DATA I0, IP, IQ, IR, IS, IT /79, 80, 81, 82, 83, 84/ DATA IU, IW, IZ /85,87,90/ **TYPE 2100** ISWFLG = -11 TYPE 95 95 FORMAT(' Command?',/) С RECORD COMMAND ENTRY USING SPECIAL SCREEN MODE С С CALL ENSPEC 200 IF(ISWFLG.GE.O) CALL TDONE COM1 = ITTINR(1)IF(COM1.LT.0) G0 T0 200 С IF AN ARROW WAS TYPED, DETERMINE WHICH IF(COM1.EQ.27) GO TO 700 С IGNORE CARRIAGE RETURNS AND LINE FEEDS IF((COM1.EQ.13).OR.(COM1.EQ.10)) GO TO 200 COM2 = 32С С ECHO CHARACTERS С 250 TYPE 255, COM1, COM2 FORMAT('+',2A1,25X) 255 300 IF(ISWFLG.LT.O) GO TO 320 CALL TDONE GO TO 250 COMT = ITTINR(1)320 IF(COMT.LT.0) G0 T0 300 IF(COMT.EQ.13) GO TO 500 380 COM2 = COMT390 TYPE 255, COM1, COM2 400 IF(ISWFLG.LT.O) GO TO 420 CALL TDONE GO TO 390

C C	LOOK FOR CARRIAGE RETURN, OR ANOTHER CHARACTER
420	COMT = ITTINR(1) IF(COMT.LT.O) GO TO 400 IF(COMT.EQ.13) GO TO 500
	IF MORE THAN TWO CHARACTERS ARE TYPED IN, SHIFT COM2 TO COM1
	COM1 = COM2 GO TO 380
	READ SECOND HALF OF CARRIAGE RETURN (LINE FEED)
500 510	COMT = ITTINR(1) IF(COMT.LT.O) GO TO 500 TYPE 510,COM1,COM2 FORMAT('+',2A1,25X,/) CALL EXSPEC
C C	EXIT SPECIAL SCREEN MODE BEFORE CHECKING COMMAND LIST
	GO TO 1100
	IF AN ARROW WAS TYPED IN, THIS SECTION DETERMINES WHICH ARROW IT WAS AND THEN CALLS SUBROUTINE DISPR TO CHANGE THE FORMAT OF THE DISPLAY
C C C	TO USE WITH A VT52 COMPATIBLE TERMINAL, THE SECOND CHARACTER OF THE ARROW DETERMINES THE DIRECTION OF THE ARROW TYPED
700 700	IF(ITERM.LT.100) G0 T0 710 COMT = ITTINR(1) IF(COMT.LT.0) G0 T0 700
	THE ARROWS ARE DEFINED AS FOLLOWS:
C C	DIRECTION VT100 VT52 EFFECT
C C C C C C C	UP\$ [A\$ARAISE DISPLAY STEP SIZEDOWN\$ [B\$BLOWER DISPLAY STEP SIZERIGHT\$ [C\$CINCREASE INITIAL DISPLAY CHANNEL BY STEP SIZELEFT\$ [D\$DDECREASE INITIAL DISPLAY CHANNEL BY STEP SIZE
τ 710	COMT = ITTINR(1) IF(COMT.LT.O) GO TO 710 IF(COMT.EQ.IA) CALL TCDISP(1) IF(COMT.EQ.IB) CALL TCDISP(2) IF(COMT.EQ.IC) CALL TCDISP(3)

C	IF(COMT.EQ.ID) CALL TCDISP(4) G0 T0 200
	MATCH UP COMMANDS AND EXECUTE DESIRED FUNCTION THE BASIC COMMANDS ARE: A* Analyze Commands C* Change Commands D* Display Commands E Exit G* Graph Commands H Help L* Load Commands P* Print Commands R* Read and Reset Commands S* Show and Start Commands T Test MCS W* Write Commands Z Zero Arrays
	A* The analyze commands are: AC Analyze Chopper for Gate Width AD Analyze Dip in TOF AH Analyze Help AP Analyze Peak in TOF
1100	IF(COM1.NE.IA) GO TO 1201 IF(COM2.NE.IC) GO TO 1121 CALL TANAL(1) GO TO 1
1121	IF(COM2.NE.ID) GO TO 1131 CALL TANAL(2) GO TO 1
1131	IF(COM2.NE.IP) GO TO 1141 CALL TANAL(3) GO TO 1
1141	IF(COM2.NE.IH) TYPE 9100 CALL TCOMLA GD TD 1
с с с с с с с с с с	C* The change commands are: CA Change All Settings CC Change Correlation Method CD Change Dwell Time CE Change External Delay CH C Help

с с с с с с с с с с с с с с с	CI CN CD CP CS CT CU CW	Change Initial Channel Change Number of Channels Change Offset for Cross-Correlation Change Detector Position Change Sweep Counts Change Trigger Counts Change User Terminal Change Display Window
1201	IF(COM1.NE.IC) IF(COM2.NE.IA) CALL TCHAN(O) CALL TDEF(O) GO TO 1	GO TO 1301 GO TO 1211
1211	IF(COM2.NE.IC) CALL TCHCOR(O) GO TO 1	GO TO 1216
1216	IF(COM2.NE.ID) CALL TCHAN(1) GO TO 1	GO TO 1221
1221	IF(COM2.NE.IE) CALL TCHAN(5) GO TO 1	GO TO 1236
1236	IF(COM2.NE.II) CALL TCHAN(4) GO TO 1	GO TO 1251
1251	IF(COM2.NE.IN) CALL TCHAN(6) GO TO 1	GO TO 1255
1255	IF(COM2.NE.IO) CALL TCHCOR(1) CALL TDSMOD GO TO 1	GO TO 1256
1256	IF(COM2.NE.IP) CALL TDEF(2) GO TO 1	GO TO 1266
1266	IF(COM2.NE.IS) CALL TCHAN(3) GO TO 1	GO TO 1271
1271	IF (COM2.NE.IT) CALL TCHAN(2) GO TO 1	GO TO 1276
1276	IF (COM2.NE.IU) CALL TTERM(1) GO TO 1	GO TO 1281
1281	IF(COM2.NE.IW) GO TO 1381	GO TO 1299
1299	IF (COM2.NE.IH) CALL TCOMLC GD TD 1	TYPE 9100

C C C ( С C 1 1 1

Č	D*	
С	The define and d	lisplay commands are:
C	DA	Display Accumulated Data
C	DB	Display Background
č	DC	Display Correlated Data
č	· DH	D Help
č	DI	Display Low Trigger Count
č	DN DN	Display No Background
c		Display Nuickly
		Display Quickly
		Display Refresh
		Display Single Sweeps
C		
C	DU	Display Uncorrelated Data
C .	Dw	Display: Change Window
С		
1301	IF(COM1.NE.ID) G	O TO 1401
	IF(COM2.NE.IA) G	0 TO 1321
	INEW = O	
	GO TO 1389	
1321	IF(COM2.NE.IB) G	0 TO 1331
	IBACK = 1	
	GO TO 1389	
1331	IF(COM2.NE.IC) G	0 TO 1336
1001	TE(TCORR ED 1) G	n Tn 1
	TCORR = 1	
	CO TO 1388	
1336	TE(COM2 NE TL) C	N TN 1341
1550	$\frac{11}{10} \left( \frac{10}{10} - 0 \right)$	0 10 1341
	1011 = 0	
1241	TE (CONO NE TN) O	0 T0 1246
1341	IF (CUM2.NE.IN) G	0 10 1346
	1BACK = 0	
	GU TU 1389	0 70 4054
1346	IF (CUM2.NE.IQ) G	0 10 1351
	IQIK = 1	
	GO TO 1389	
1351	IF(COM2.NE.IR) G	0 TO 1361
	CALL TDISP	
	GO TO 1	
1361	IF(COM2.NE.IS) G	0 TO 1366
	INEW = 1	
	GO TO 1389	
1366	IF(COM2.NE.IT) G	0 TO 1371
	CALL TDEF(1)	
	GO TO 1	
1371	IF(COM2.NE.IU) G	O TO 1381
-	IF (ICORR.ED.O) G	0 TO 1
	ICORR = 0	
	GD TD 1388	
1381	TE(COM2.NE TW) G	0 T0 1399
avvė		

	•
1388 1389 1399	CALL TCDISP(5) CALL TDISP CALL TDSMOD GO TO 1 IF(COM2.NE.IH) TYPE 9100 CALL TCOMLD GO TO 1
	E Exit
1401	IF(COM1.NE.IE) GO TO 1501 CALL TLEAVE GO TO 1
	G* The graph commands are: GH G Help GL Graph Line Plot GP Graph Point Plot
1501	IF(COM1.NE.IG) GO TO 1555 IF(COM2.NE.IL) GO TO 1521 CALL TPLOT(1) GO TO 1
1521	IF (COM2.NE.IP) GO TO 1531 CALL TPLOT(O) GO TO 1
1531	IF(COM2.NE.IH) TYPE 9100 CALL TCOMLG GO TO 1
	H Help Print out command list.
1555	IF(COM1.NE.IH) GO TO 1601 CALL TCMLST GO TO 1
	L* The load commands are: LA Load All LD Load Dwell Time LH L Help LI Load Initial Channel LN Load Number of Channels LS Load Sweep Counts LT Load Trigger Counts LT Load Trigger Counts
TOAT	TL(COMI''NE''I'') AN IN INDI

	IF (COM2.NE.IA) CALL TLOAD (O)	GO TO 1621
1621	IF(COM2.NE.ID) CALL TLOAD(1)	GO TO 1641
1641	IF (COM2.NE.IN) CALL TLOAD (2)	GO TO 1651
1651	IF(COM2.NE.II) CALL TLOAD(3)	GO TO 1661
1661	IF (COM2.NE.IS) CALL TLOAD (4) GO TO 1	GO TO 1671
1671	IF (COM2.NE.IT) CALL TLOAD (5) GO TO 1	GO TO 1681
1681	IF(COM2.NE.IH) CALL TCOMLL	TYPE 9100
C	00 10 I	
č	P*	
č	The available P	rint commands are
č	PC	Print Configuration
č	РН	P Heln
č	PŤ	Print Information on the Line Printer
c c	PI	Print Data on the Line Printer
c ·		Print Data on the Terminal with Pauses
c c	DD	Print Cursent Sector Contents on the Line Prints
c c	DT	Print Current Scaler Contents on the Line Frinte
c c	11	
1701	TE (CONT NE TO)	CO TO 2001
1701	IF (COM2.NE.IC) IF (COM2.NE.IC) CALL TSHOCO(6) GO TO 1	GO TO 1711
1 <b>711</b>	IF(COM2.NE.II) CALL TPRII(6) CD TO 1	GO TO 1721
1721	IF (COM2.NE.IL) CALL TPRDT(6)	GO TO 1731
1731	IF (COM2.NE.IP) CALL TPRDT (-7)	GO TO 1736
1736	IF (COM2.NE.IR) CALL READMS (NEW	GO TO 1741 /)
1737	FORMAT(16(/,8(2 CALL RSTART	((11),11=1CHAN, (1CHAN+NCHAN-1)) X,I6,:)))

GO TO 1 1741 IF(COM2.NE.IT) GO TO 1751 CALL TPRDT(7) GO TO 1 1751 IF (COM2.NE.IH) TYPE 9100 CALL TCOMLP GO TO 1 С С R\* С RA Read Ascii Data File C C RH Read and Reset Help Reset MCS (Halt Acquisition) RM С RS Reset System Hardware С 2801 IF(COM1.NE.IR) GO TO 3001 IF(COM2.NE.IA) GO TO 2846 CALL TREDAS GO TO 1 IF(COM2.NE.IM) GO TO 2866 2846 CALL STOPR GO TO 1 2866 IF(COM2.NE.IS) GO TO 2899 CALL RINIT GO TO 1 2899 IF (COM2.NE.IH) TYPE 9100 CALL TCOMLR GO TO 1 С С S\* С The available show and start commands are: С SC Show Configuration C SD Show Display Mode С SE Scan Extend Ĉ SH S Help C ST Start С 3001 IF(COM1.NE.IS) GO TO 3501 IF (COM2.NE.IC) GO TO 3011 CALL TSHOCO(7) GO TO 1 3011 IF(COM2.NE.ID) GO TO 3021 CALL TDSMOD GO TO 1 3021 IF(COM2.NE.IE) GO TO 3031 CALL TSTART(1) GO TO 1 3031 IF(COM2.NE.IT) GO TO 3041 CALL TSTART(0) GO TO 1 3041 IF (COM2.NE.IH) TYPE 9100

CALL TCOMLS GO TO 1 C \_ C \_ C \_ C Т Test MCS 3501 IF(COM1.NE.IT) GO TO 4001 CALL TSTMCS GO TO 1 С C C C ₩ Write Commands WF Write Data File with Full Header C C WH Write Help WI Write Data File with Informational Header C WN Write Data File with no Header С 4001 IF(COM1.NE.IW) GO TO 4101 IF (COM2.NE.IF) GO TO 4036 CALL TWRA(0) GO TO 1 4036 IF(COM2.NE.II) GO TO 4051 CALL TWRA(1) GO TO 1. 4051 IF(COM2-NE.IN) GO TO 4099 CALL TWRA(-1) GO TO 1 IF(COM2.NE.IH) TYPE 9100 4099 CALL TCOMLW GO TO 1 С С Z Command С Zero Data Arrays С 4101 IF(COM1.NE.IZ) G0 T0 4500 TYPE \*,'Zero Árrays? Are you sure?' ACCEPT 1000,ANS 1000 FORMAT(A1) IF (ANS.EQ.'Y') CALL TZERO GO TO 1 С C TYPE OUT THE FOLLOWING IF AN UNRECOGNIZED COMMAND WAS ISSUED C **TYPE 9100** 4500 **TYPE 2100** 2100 FORMAT(' For a list of commands, type H') GO TO 1 9100 FORMAT(/, 'This is not a command.') RETURN END

### SUBROUTINE TANAL (IOPT)

IF(I0PT.NE.1) G0 T0 3000

IBB = IBIBE = IEGO TO 3200

BACK = 0.

ACCEPT \*, IBB, IBE

LEN = IBE - IBB

File TANAL.FOR begins on the previous line.

Version Date: November 8, 1985 Paul S. Weiss

Subroutine TANAL does the analysis of the TOF scan to give peak or dip positions in a region of the TOF spectrum of choice. It also determines the proper gate width for a Timer-Gater module based upon the beam modulation done by a tuning fork chopper (option 1).

WARNING: TANAL is the limiting segment of overlay region 2. Enlarging it will increase the total size of TUF.

Argument IOPT determines the manner in which the data is interpreted, as follows:

IOPT = 1	The data is assumed to be the TOF of
	beam chopped by a tuning fork, and the
	appropriate Gate Width is determined.
IOPT = 2	The data is assumed to be from hole
	burning, and the depth and width of the
	hole are determined.
IOPT = 3	The data is assumed to be a normal TOF,
	•

and the height and width of the peak are determined.

COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /MCSMAX/ DWLMAX, DWLMIN, MAXSWP, MAXTRG COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT

TYPE 2000, 'nalyze.', ' FORMAT(/,' Enter beginning and ending channels of segment to a'

1,A7,A21) ACCEPT \*,IB,IE

3000

2000

С

C

C

3200

C C

IF SEGMENT ENDS WERE ENTERED IN THE WRONG ORDER IF(LEN.GE.O) GO TO 4000 I = IBE IBE = IBB IBB = I LEN = - LEN

TYPE 2000, 'verage ', 'for background level.'

D0 4100 I=IBB,IBE 4000 BACK = BACK + ACCUM(I)4100 CONTINUE BACK = BACK / (LEN + 1)LEN = IE - IBIF(LEN.GE.O) G0 T0 5000 I = IEIE = IBIB = ILEN = - LEN С С IOP DETERMINES WHETHER TO LOOK FOR A PEAK OR A DIP С 5000 IOP = 1IF(IOPT.EQ.2) IOP = -1С С INITIALIZE VALUE AND POSITION OF MAXIMUM (/MINIMUM) C 5050 TOP = ACCUM(IB)ITOP = IBD0 5100 I=IB,IE С С CHECK FOR NEW MAXIMUM (/MINIMUM) VALUE IF(((ACCUM(I)-TOP)\*IOP).LT.O.) GO TO 5100 TOP = ACCUM(I)ITOP = ICONTINUE 5100 UPDN = 'maxi' IF(IOP.EQ.-1) UPDN = 'mini' TYPE 5200, UPDN, ITOP FORMAT(/, 'The ', A4, 'mum of this segment occurs in channel ', I4) HALFH = (TOP + BACK) / 2 5200 D0 6100 I=ITOP.IE IF(((ACCUM(I) - HALFH) \* IOP).LT.O.) GO TO 6300 6100 CONTINUE IGS = 0TYPE 6200, UPDN, 'slow' 6200 FORMAT(' No half ',A4, 'mum point found on the ',A4, ' side.') GO TO 700Ò 6300 IGS = I - ITOPTYPE 6400, UPDN, 'slow', IGS, (IGS \* DWELL) FORMAT(' Half', A4, 'mum point on the', A4, ' side is ' 6400 1,I4, ' channels = ',F7.0, ' usec.') LEN = ITOP - IB + 17000 D0 7100 I=1,LEN IF(((ACCUM(ITOP-I+1) - HALFH) \* IOP).LT.O.) GO TO 7500 7100 CONTINUE IGF = 0TYPE 6200,UPDN,'fast' GO TO 8000

7500 IGF = I - 1TYPE 6400, UPDN, 'fast', IGF, (DWELL \* IGF) C С IF ANALYZING A TUNING FORK CHOPPER PRODUCED TOF. EITHER GO BACK C AND FIND DIP NOW, OR IF THAT HAS BEEN DONE, FIND PROPER GATE WIDTH С 8000 IF(IOPT.NE.1) GO TO 9000 IF(IOP.NE.1) GO TO 8100 GWON = (IGF + IGS) \* DWELLIOP = -1GO TO 5050 8100 GWOFF = (IGF + IGS) + DWELLTYPE 8110, 'on ', GWON, 'off', GWOFF, AMIN1 (GWON, GWOFF) FORMAT(/,2(/,' The beam is ',A3,' for ',F7.0,' usec.') 1,/,' Set the gate width at ',F7.0,' usec.',/) 8110 9000 RETURN END SUBROUTINE TCDISP(ICDISP) С File TCDISP.FOR begins on the previous line. С С Version Date: July 23, 1984 С Paul S. Weiss С С Subroutine TCDISP changes the display window either from the arrow С function keys or through the DW command. С С ARROWS ARE DEFINED AS FOLLOWS: С UP \$[A RAISE DISPLAY STEP SIZE C C \$ **B** DOWN LOWER DISPLAY STEP SIZE RIGHT **\$**[C INCREASE INITIAL DISPLAY CHANNEL BY STEP SIZE C C DECREASE INITIAL DISPLAY CHANNEL BY STEP SIZE LEFT **S**[D C COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT GO TO (2100,2200,2300,2400,2500),ICDISP 2100 IF(IRES.LT.16) IRES = IRES + 1GO TO 8000 2200 IF(IRES.GT.1) IRES = IRES - 1 GO TO 8000 2300 IF(IADV.LT.(4096 - IRES)) IADV = IADV + IRES GO TO 8000 IF(IADV.GE.IRES) IADV = IADV - IRES 2400 GO TO 8000 2500 TYPE \*, 'Enter initial channel to be displayed' ACCEPT +, J IADV = J - ICHANIF(IADV.LT.O) IADV = 0IF(J.LE.4096) GO TO 2900 TYPE \*, 'Out of Range.'

2900 8000	GO TO 2500 TYPE *,'Enter display step size in channels (1-16)' ACCEPT *,IRES IF((IRES.GT.16).OR.(IRES.LE.O)) GO TO 2900 RETURN END
C	SUBROUTINE TCHAN(IDEVIC) File TCHAN.FOR begins on the previous line.
	Version Date: October 15, 1984 Paul S. Weiss
	Subroutine TCHAN accepts new values of the various settings of the multichannel scaler registers. It then loads these new values into the proper registers of the computer and MCS unit.
	The argument IDEVIC determines which of these settings are to be entered by the user, as follows: IDEVIC = 0 Change all settings. IDEVIC = 1 Change the dwell time. IDEVIC = 2 Change the number of triggers per sweep. IDEVIC = 3 Change the number of sweeps. IDEVIC = 4 Change the initial channel. IDEVIC = 5 Change the external delay time. IDEVIC = 6 Change the number of channels.
C	INTEGER+4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG IF(IDEVIC.EQ.0) G0 TO 2100 GD TO (2100,2200,2300,2400,2500,2600),IDEVIC
2100 2101	TYPE 2101,DWELL FORMAT(' Enter the dwell time in microseconds (currently ',F7.2 1,' usec).') ACCEPT *,DWELL
· .	IF((DWELL.GI.(DWLMIN-0.05)).AND.(DWELL.LE.(DWLMAX+.05))) 1 GO TO 2150 TYPE 2110.DWLMIN.DWLMAX
2110	FORMAT(/,' The dwell time must be between ',F5.2,' and ' 1,F7.2,' microseconds!')
2150	CALL TLOAD(1) IF(IDEVIC.NE.O) RETURN
2200 2201	TYPE 2201,ITRIG FORMAT(' Enter the number of triggers per sweep (currently ',I5

1,').') ACCEPT \*, ITRIG IF((ITRIG.GT.O).AND.(ITRIG.LE.MAXTRG)) GO TO 2250 TYPE 2210, MAXTRG 2210 FORMAT(/,' The number of triggers must be between 1 and ', I5) GO TO 2200 2250 CALL TLOAD(5) IF(IDEVIC.NE.O) RETURN 2300 TYPE 2301, ISWEEP 2301 FORMAT(' Enter number of sweep counts (currently ', I5, ').') ACCEPT \*, ISWEEP IF((ISWEEP.GT.O).AND.(ISWEEP.LE.MAXSWP)) GD TO 2350 TYPE 2310, MAXSWP FORMAT(/,' The number sweeps must be between 1 and ', I5) 2310 GO TO 2300 2350 CALL TLOAD(4) IF(IDEVIC.NE.O) RETURN TYPE 2401, ICHAN 2400 2401 FORMAT(' Enter initial channel (currently ',I4,').') ACCEPT \* ICHAN IF((ICHAN.GT.O).AND.(ICHAN.LT.4096)) GD TO 2450 **TYPE 2420** 2420 FORMAT(/,' The initial channel must be between' 1,' 1 and 4096!') GO TO 2400 2450 CALL TLOAD(3) IF(IDEVIC.EQ.0) G0 T0 2600 RETURN 2500 TYPE 2501, EDELAY FORMAT(' Enter external delay in microseconds (currently ' 2501 1,F10.3,' usec).') ACCEPT \*,EDELAY IF (IDEVIC.NE.O) RETURN 2600 TYPE 2601, NCHAN FORMAT(' Enter the number of channels (currently ',I4,').') 2601 ACCEPT +, NCHAN IF((NCHAN.GT.O).AND.(NCHAN.LE.4096)) GO TO 2650 TYPE \*, 'The number of channels must be between 1 and 4096' GO TO 2600 2650 CALL TLOAD(2) 9000 RETURN END

SUBROUTINE TCHCOR (IOPT)

С

C C

С

File TCHCOR.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

С C C Subroutine TCHCOR changes the number of channels per TOF wheel slot and the cross-correlation offset С COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT IF(I0PT.GT.0) G0 T0 5000 **TYPE 100** 1 100 FORMAT(/, ' How many MCS channels per TOF wheel slot?') ACCEPT \*, NCORR IF((NCORR.GT.O).AND.(NCORR.LE.16)) GO TO 9000 TYPE 200,1,16 FORMAT(/, 'Value must be between ', I1, ' and ', I3, '.') 200 G0 T0 1 С С CHANGE OFFSET С 5000 **TYPE 5100** FORMAT(/,' Enter zero offset for cross-correlation') 5100 ACCEPT \*, IOFSET IF((IOFSET.GE.O).AND.(IOFSET.LE.255)) GO TO 9000 TYPE 200,0,255 GO TO 5000 9000 RETURN END SUBROUTINE TCMLST С File TCMLST.FOR begins on the previous line. С С Version Date: October 15, 1984 С Paul S. Weiss С С Subroutine TCMLST lists the main commands for TPUF. It serves as С the main help screen. С **TYPE 5100** FORMAT(//,' 5100 The available commands are as follows:' 1,/, , Ą\* Analyze Commands' 2,/, Change Commands' C\* , Define and Display Commands' 3,/, D\* , Ε 4,/, Exit' , Graph Commands' 5,/, G\* , 6,/, Η Help') **TYPE 5200** 5200 FORMAT (' Load Commands.' L\* 2,/,' P\* Print Commands' 3,/,; Read and Reset Commands' R\* 4,/, S\* Show and Start Commands' , T Test MCS' 5,/, ₩\* Write Data Commands' 6,/,

5300	7,/,' Z Zero Data Arrays') TYPE 5300 FORMAT(/,' To get a listing of the various options of' 1,' a starred command,',/,' enter an "H" in place of the "*"' 2,' (i.e. "PH" for a list of the Print commands).')
	RETURN END
C	SUBROUTINE TCOMLA File TCOMLA.FOR begins on the previous line.
	Version Date: July 23, 1984 Paul S. Weiss
	Subroutine TCOMLA lists all the A* commands.
100	TYPE 100 FORMAT(/,' The available Analyze options are as follows:' 1,/,' AC Analyze Chopper for Gate Width' 3,/,' AD Analyze Dip (for Hole Burning)' 5,/,' AH A Help' 2,/,' AP Analyze Peak',/) RETURN END
C	SUBROUTINE TCOMLC File TCOMLC.FOR begins on the previous line.
	Version Date: October 12, 1984 Paul S. Weiss
C C	Subroutine TCOMLC lists all the C* commands.
100	TYPE 100 FORMAT(/,' The available Change options are as follows:' 1,/,' CA Change All Settings' 2,/,' CC Change Correlation Method' 3,/,' CD Change Dwell Time' 4,/,' CE Change External Delay' 5,/,' CH C Help' 6,/,' CI Change Initial Channel')
200	FORMAT(' CN Change Number of Channels' 8,/,' CO Change Offset for Cross-Correlation' 9,/,' CP Change Detector Position' 9,/,' CS Change Sweep Counts' 1,/,' CT Change Trigger Counts' 2,/,' CU Change User Terminal Type'

	3,/,' CW Change Window for Display',/) RETURN END
C	SUBROUTINE TCOMLD File TCOMLD.FOR begins on the previous line.
	Version Date: October 12, 1984 Paul S. Weiss
	Subroutine TCOMLD types out the available D* commands and explains the use of the arrows in changing the display.
200	TYPE 100 FORMAT(/,' The available define and display options ' 1,'are as follows:' 1,/,' DA Display Accumulated Data' 2,/,' DB Display Background' 3,/,' DC Display Correlated Data' 4,/,' DH D Help' 5,/,' DL Display Speed for Low Trigger Counts' 5,/,' DN Display No Background' 6,/,' DQ Display Quickly' 6,/,' DR Display Refresh' 7,/,' DS Display Single Sweeps' 8,/,' DT Define Title' 8,/,' DU Display Uncorrelated Data' 9,/,' DW Display: Change Window',/) TYPE 200 FORMAT(/,' The arrows on your terminal can be used'
	<pre>1,' to change the',/,' display window as follows:' 2,/,' Up/Down Increase/Decrease the Display Step Size' 3,' by 1 channel' 4,/,' Left/Right Change the Initial Channel to be displayed by' 5,/,' -/+ the Display Step',/) RETURN END </pre>
C	SUBROUTINE TCOMLG File TCOMLG.FOR begins on the previous line.
	Version Date: July 23, 1984 Paul S. Weiss
C C	Subroutine TCOMLG lists all of the G* commands.
1535	TYPE 1535 FORMAT(/,' The available Graph options are as follows:'

1,/,' G Help' GH 2,/,' 3,/,' RETURN GL Graph Line Plot' GP Graph Point Plot',/) END SUBROUTINE TCOMLL File TCOMLL.FOR begins on the previous line. Version Date: July 23, 1984 Paul S. Weiss Subroutine TCOMLL lists all the L\* load commands. **TYPE 100** FORMAT(/, ' The available Load options are as follows' 1,/,' LA Load All' 2,/,, LD Load Dwell Time' LH 3,/, L Help' , LI 4,/, Load Initial Channel' LN Load Number of Channels' 5,/, LS 6,/, Load Sweep Counts' , LT 7,/, Load Trigger Counts') RETURN END SUBROUTINE TCOMLP File TCOMLP.FOR begins on the previous line. July 23, 1984. Version Date: Paul S. Weiss Subroutine TCOMLP type out the available P\* print commands. **TYPE 100** FORMAT(/, ' The available print options are as follows:' 1,/,' ÝĈ Print Configuration' , 2,/,' 3,/,' PH P Help' ΡI Print Information (from the User) on ' 'the Line Printer' 3, 4,/,' PL Print Data on the Line Printer' , 5,/,' PP 6,' Screen' PP Print Data on the Terminal Pausing for each' 7,/,' PR Print Current Scaler Contents on the Line Printer' 8,/,' PT Print Data on the Terminal') RETURN END

100

CCCCCCC

100

### SUBROUTINE TCOMLR

## File TCOMLR.FOR begins on the previous line.

Version Date: October 11, 1984 Paul S. Weiss

Subroutine TCOMLR lists all the R\* commands.

TYPE 100

C C C C C C C C

С

С

С

C C C

C C

С

С

C C C C C C C C C C C C

100

100

	-		
FORMAT (	//,?	The	Read and Reset Commands are: '
1,/,'	RA		Read Ascii Data File'
2,/,'	RH		Read and Reset Help'
3,/, '	RM		Reset MCS (Halt Acquisition)'
4,/,'	RS		Reset System Hardware',//)
RETURN			• • • • • • • • • • • • • • • • • • • •
END			

SUBROUTINE TCOMLS

**TYPE 100** 

File TCOMLS.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TCOMLS lists the available S\* show and start commands.

100

FORMAT(/,' The Show and Start options are as follows:' 1,/,' SC Show Configuration'

1,/,' SC Show Configuration 2,/,' SD Show Display Mode' 3,/,' SE Scan Extend' 4,/,' SH S Help' 5,/,' ST Start') RETURN END

SUBROUTINE TCOMLW

File TCOMLW.FOR begins on the previous line.

Version Date: October 11, 1984 Paul S. Weiss

Subroutine TCOMLW lists all the W\* commands

TYPE 100

FORMAT(//,' The Write Commands are:' 1,/,' WF Write Data File with a Full Header'

2,/,' WH Write Help' 3,/, WI Write Data File with Entered Information in' 3,' Header' 4,/,' WN Write Data File with No Header',//) RETURN END SUBROUTINE TDEF (IOPT) C C File TDEF.FOR begins on the previous line. C C C C C C C C Version Date: November 8, 1985 Paul S. Weiss Subroutine TDEF changes the title and detector angle for TUF, both these values are cosmetic and have no effect other than for output C C to console, printer, or data file. С The argument IOPT is defined as follows: Č C C IOPT=0 Change Title and Detector Angle. IOPT=1 Change Title. Ċ IOPT=2 Change Detector Angle. С С COMMON /TITANG/ TITLE(10), ANGLE IF(IOPT-1) 1100,1100,2100 1100 TYPE 1110, (TITLE(I), I=1,10) FORMAT(' The current title is: ',/,1X,10A4,/,' Change this?') 1110 ACCEPT 1000, ANS 1000 FORMAT(A1) IF (ANS.NE. 'Y') GO TO 2000 TYPE \*,'Enter title.' С С ENABLE LOWER CASE INPUT С CALL LCASE ACCEPT 1310, (TITLE(I), I=1,10) 1310 FORMAT(10A4) С С DISABLE LOWER CASE INPUT С CALL UCASE 2000 IF(I0PT.NE.O) G0 T0 9000 2100 TYPE 2110, ANGLE FORMAT(' The current detector angle is ', F6.2,' degrees.' 2110 1,/,' Change this?') ACCEPT 1000, ANS IF (ANS.NE.'Y') GO TO 9000 TYPE \*, 'Enter detector angle' ACCEPT +, ANGLE

C C C C C C C

С

С

C

С

Ċ

Ĉ

С

C C

С

C C

С

С

Ը Շ

C 2000

Ը Ը

С

1500

# RETURN

SUBROUTINE TDISP

File TDISP.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TDISP sets up the array to be displayed IDSP. If the data is cross-correlated, subroutine TCORR is called. Note that the display is stopped before the array is altered and restarted after it has been recalculated.

WARNING: TDISP is in the root segment of TUF. Enlarging it will enlarge the already precariously large size of TUF.

INTEGER+4 JTRIG

COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IDFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG

SET VALUES IF DATA IS TO BE CROSS-CORRELATED

IF(ICORR.NE.1) GO TO 1500 ND = 255 IBEG = 1 ISTEP = NCORR GO TO 2000

SET FINAL DISPLAY CHANNEL ND SO THAT IT IS NOT GREATER THAN THE LAST CHANNEL RECORDED

IBEG = ICHAN + IADV
ISTEP = IRES
ND = 256
IF((IBEG → (256\*ISTEP)).GT.4097)
1 ND = ((4097 - IBEG)/ISTEP)

SEE IF ONLY SINGLE SWEEP DATA IS TO BE DISPLAYED

IF(INEW.EQ.1) GO TO 3000

PREPARE ACCUMULATED DATA

D0 2500 I=1,ND KPLAT = IBEG + ((I-1) \* ISTEP) - 1 WORK(I) = 0.

D0 2400 J=1, ISTEP ISUB = KPLAT + JWORK(I) = WORK(I) + ACCUM(ISUB) 2400 CONTINUE CONTINUE 2500 GO TO 4000 С С PREPARE SINGLE SWEEP DATA С 3000 DO 3500 I=1,ND KPLAT = IBEG + ((I-1) \* ISTEP) - 1WORK(I) = 0.D0 3400 J=1, ISTEP ISUB = KPLAT + JCONTINUE 3400 WORK(I) = WORK(I) + FLOAT(NEW(ISUB)) 3500 CONTINUE С C IF DISPLAY IS TO BE CROSS-CORRELATED DATA, CALL TCORR TO DO THIS С 4000 IF(ICORR.NE.1) GO TO 4100 WORK(255) = (WORK(1) + WORK(254))/2.CALL TCORR С Ç FIND MINIMUM AND MAXIMUM OF DISPLAY Ĉ 4100 YMIN = WORK(1)YMAX = YMIN DO 5000 I=1,ND IF(WORK(I).GT.YMAX) YMAX = WORK(I) IF(WORK(I).LT.YMIN) YMIN = WORK(I) 5000 CONTINUE С С IF NO BACKGOUND SUBTRACTION IS TO BE DONE, SET YMIN TO O C 7000 IF((IBACK, EQ.1), AND.(YMIN.GT.0.)) YMIN = 0.С С SCALE DISPLAY ARRAY TO 12 BITS (0 TO 4095) Ç DIVR = YMAX - YMIN IF(DIVR.EQ.O.) DIVR = 1.DIVR = 4095./DIVRC С BEFORE ALTERING IDSP, STOP THE DISPLAY C CALL STOPD С С LOAD DISPLAY ARRAY IDSP С D0 7500 I=2,512,2

#### IDSP(I - 1) = (8 \* I) - 8IDSP(I) = INT((WORK(I/2) - YMIN) \* DIVR)CONTINUE LOAD NEW DISPLAY ARRAY FOR DISPLAY DRIVER, AND RESTART CALL DISPL(IDSP,ND) RETURN END SUBROUTINE TCORR Second routine of file TDISP.FOR. Version Date: July 23, 1984 Paul S. Weiss Subroutine TCORR deconvolutes the data in WORK that has been cross-correlated by chopping the product beams with the crosscorrelation sequence. The correlation sequence is stored in ISEQZ LOGICAL\*1 ISEQ(512), ISEQZ(256) COMMON /DATA/ ACCUM(4096), XCORR(256), WORK(256), NEW(4096), NSWEEP COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /TERM/ ITERM EQUIVALENCE (ISEQ, ISEQZ) ,T,F,F,F,F DATA ISEQZ/T,T,T,T,T,T,T,F,F,T,F,T,T,T,T,F,T,F,F,T,F 2,T,T,F,T,T,T,F,T,T,F,Ť,Ť,Ť,Ť,Ť,Ť,Ě,Ť,Ť,Ť,Ť,Ě,Ť ,F ,F . F .T.T.F 3, F, T, F, T, F, T, F, F, F, F, T, T, F, T, F, T, T, F, F ,F .F.T.T.F.F , T , F . F ,T,F,F ,T,T,F,T,F,F 4,T,F,T,T,F,T,T,F,T,F ,T,T,T,T,T,T,F,T,T,F, F, F,F,F,F,T,F,F • , F 5, T, T, F, F, T, T, T, T, F, T, T, F, F, T, ,F,F,F,T,T ,T,F ,F,T,T,F,F,F,T,F,F,T,T,T,F,T,F,T,F,T,T,F,T,F,T,F 6.F.F.T.F . F .F 7, F, T, F, T, F, F, F, F, F, T, T, T, T, F, F, F, F, F, F, F, F, T, T, T, T, F, F, F, T, T, T, T, T, F, F, F, T 8, T, T, F, T, T, T, T, F, F, F, T, F, T, T, F, F, T, T, F, T, T, F .F.F . F 9, T, T, T, F, F, F, F, T, F, T, F, O/ DO 100 I=1,256 ISEQ(I+255) = ISEQ(I)100 CONTINUE DO 2500 I=1,255 XCORR(I) = 0.KNTR = 256 - I + IOFSETDO 2000 J=1,255 KNTR = KNTR + 1WHEN BIT IS TRUE (ISEQ = T), ADD VALUE OF WORK

C C C C C С С Ĉ

> С С С

С C

С С С

С

7500 С С С

> ,F ,F ·,T .T.F.F T,T,T,T,F,F

C WHEN BIT IS FALSE (ISEQ = F), SUBTRACT VALUE OF WORK

6	
	IF(ISEQ(KNTR)) GO TO 1800
	XCORR(I) = XCORR(I) - WORK(J)
	GD TO 2000
1800	XCORR(I) = XCORR(I) + WORK(J)
2000	CONTINUÉ
2500	CONTINUE
	DO 4000 I=1,255
	WORK(I) = XCORR(I)
4000	CONTINUE
	RETURN
	END

SUBROUTINE TDONE

File TDONE.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TDONE collects the data from the MCS using the MACRO subroutine READMS. It determines whether a counter overflow has occurred by checking the flag/data word JTRIG set in interrupt service routine MCSINT. If JTRIG is negative, an overflow has occurred after -JTRIG triggers. If JTRIG is positive, a sweep of JTRIG triggers has been completed successfully. The new data is placed in integer array NEW by READMS, and is added to the accumulated data array ACCUM. If the display quickly mode has been chosen (i.e. if the DQ command has been executed) the MCS is restarted if sweeps remain to be accumulated, and then the display routine is called. This eliminates the overhead involved in calculating the display array. If the display speed for low trigger counts has been chosen (default, or by command DL), the current data is displayed, and then the MCS is restarted if sweeps remain.

INTEGER+4 JTRIG COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /DATA/ ACCUM(4096), XCORR(256), WORK(256), NEW(4096), NSWEEP COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /SWFLAG/ ISWFLG, JTRIG

FIRST CALL READMS, THE MACRO ROUTINE TO READ THE SCALER

TRIGS = AJFLT(JTRIG) IF(TRIGS.LT.O) GO TO 6000 IF(ISWFLG.EQ.O) GO TO 3000 TYPE 1200,ISWFLG,7 FORMAT('+',I5,' sweeps remaining.',A1)

1200

С

Č

CCCCCCCCCC

Č C

Č

С

Č

C C

C C C

> Ը Շ

> С

GO TO 4000 3000 3100 FORMAT('+\* 1,/,' Scan has been completed' 2,/,' with ',F9.0,' triggers.' 2,/,' with ' 3,/,' \*\*\*\*\*\* 4,//, ' Command?',/) 4000 CALL READMS (NEW) С С RESTART MCS USING THE MACRO RSTART IF SWEEPS REMAIN TO BE DONE С AND IF IN "QUICK DISPLAY" MODE С IF((ISWFLG.GT.O).AND.(IQIK.GT.O)) CALL RSTART D0 4200 I=1,4096 IF(NEW(I).GE.O) GO TO 4150 ACCUM(I) = ACCUM(I) + 65536.4150 ACCUM(I) = ACCUM(I) + FLOAT(NEW(I))4200 CONTINUE GO TO 7000 TYPE 6100, (-TRIGS), ISWFLG, 7, 7, 7, 7 6000 FORMAT(/,' Count Overflow has occurred after ',F6.0,' triggers' 1,' with ',I5,' sweeps remaining.',4A1 6100 2,/,' Lower trigger count before restarting.' 3,//, ' Command?',/) GO TO 8900 7000 CALL TDISP С C RESTART MCS USING THE MACRO RSTART IF SWEEPS REMAIN TO BE DONE С AND IF IN "LOW TRIGGER COUNT DISPLAY" MODE С IF((ISWFLG.GT.O).AND.(IQIK.EQ.O)) CALL RSTART 8900 ISWFLG = -1RETURN END SUBROUTINE TDSMOD С File TDSMOD.FOR begins on the previous line. Version Date: July 24, 1984 Paul S. Weiss С Subroutine TDSMOD types out the current display parameters. С COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT IF(INEW.EQ.1) GO TO 2200 TYPE 2100, 'accumulated data, ' FORMAT(/,' The current display is of ',A18) 2100

2200 2500	GO TO 2500 TYPE 2100,'single sweep data,' IF(ICORR.EQ.1) GO TO 3000 IB = ICHAN + IADV IE = IB + (256*IRES) - 1 IF(IE.LE.4096) GO TO 2600 ND = (4097 - IB)/IRES
2600 2650	IE = IB + (ND*IRES) - 1 TYPE 2650,IB,IE,IRES FORMAT(' channels ',I4,' to ',I4,', in steps of ',I2,' channels,') GO TO 4000
3000 3100 4000	TYPE 3100, IOFSET FORMAT(' cross-correlated with an offset of ',I4) IF(IBACK.EQ.0) GO TO 4500 TYPE 4100 'included ' ''
4100	FORMAT(' with the background level ',A9,A2)
4500 5000	TYPE 4100, 'subtracte', 'd.' IF(IQIK.EQ.O) GD TO 5200 TYPE 5110, 'quickly.'
5110	FORMAT(' The display is refreshed ',A8,/)
5200 9000	TYPE 5110,'slowly. ' RETURN END
C	SUBROUTINE TINIT File TINIT.FOR begins on the previous line.
	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss
	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss Subroutine TINIT initializes most of the common variables for the multichannel scaler.
C C C C C C C C C C C C C C C C C C C	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss Subroutine TINIT initializes most of the common variables for the multichannel scaler. INTEGER+4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG COMMON /TERM/ ITERM COMMON /TITANG/ TITLE(10),ANGLE
C C C C	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss Subroutine TINIT initializes most of the common variables for the multichannel scaler. INTEGER*4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG COMMON /TERM/ ITERM COMMON /TITANG/ TITLE(10),ANGLE DATA TITLE /'Seco','ndar','y So','urce',' TOF',5*' '/
С С С С С	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss Subroutine TINIT initializes most of the common variables for the multichannel scaler. INTEGER*4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG COMMON /TITANG/ TITLE(10),ANGLE DATA TITLE /'Seco','ndar','y So','urce',' TOF',5*' '/ TYPE 5
C C C C C S	SUBROUTINE TINIT File TINIT.FOR begins on the previous line. Version Date: November 10, 1985 Paul S. Weiss Subroutine TINIT initializes most of the common variables for the multichannel scaler. INTEGER*4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /MCSMAX/ DWLMAX,DWLMIN,MAXSWP,MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /TERM/ ITERM COMMON /TITANG/ TITLE(10),ANGLE DATA TITLE /'Seco','ndar','y So','urce',' TOF',5*' '/ TYPE 5 FORMAT(//,' TOF Program' 1.(.' Version Date: November 10, 1985',///)

INITIALIZE ALLOWED VARIABLE RANGES DWLMIN = 0.15DWLMAX = 200.MAXSWP = 32767MAXTRG = 32767INITIALIZE MCS SETTINGS DWELL = 100.ICHAN = 1NCHAN = 256ISWEEP = 10ITRIG = 100EDELAY = 0IOFSET = 0IEXT = 0ANGLE = 0. INITIALIZE TERMINAL SETTING ITERM = 100INITIALIZE DISPLAY SETTINGS IADV = 0IRES = 1INEW = 0ICORR = 0IBACK = 1IOFSET = 0NCORR = 1IQIK = 0CALL TZERO SHOW CURRENT CONFIGURATION CALL TSHOCO(7) TYPE \*, 'Are these default values satisfactory?' ANSWER IS READ IN MAIN ROUTINE TPUF RETURN END

SUBROUTINE TLEAVE

С

С

C C

С

C C

č

C

С

C C

С

C C

С

С С File TLEAVE.FOR begins on the previous line.

С С С С С

C

C C

С

C C

С

С

CCCCCCC

C C

C C

С

CCCCCC

С

1000

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TLEAVE is the termination routine.

TYPE \*,'Exit? Are you sure?' ACCEPT 1000,ANS FORMAT(A1) IF(ANS.NE.'Y') RETURN

STOP THE MCS

CALL STOPR

STOP THE DISPLAY

CALL STOPD STOP END

SUBROUTINE TLOAD (IDEVIC)

File TLOAD.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TLOAD loads the already stored value of a specified setting into the appropriate register of the computer or multichannel scaler unit using the MACRO subroutines LOADDW, LOADSW, and LOADSW.

The value of the argument IDEVIC determines which settings are to be loaded into their registers, as follows:

IDEVIC	=	0	Load	all settings.
IDEVIC	=	1	Load	dwell counts.
IDEVIC		2	Load	number of channels
IDEVIC		Έ	Load	initial channel.
IDEVIC		4	Load	sweep counter.
IDEVIC		5	Load	trigger counter.
		•		,

COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT IF(IDEVIC.EQ.O) GD TO 2100 GO TO (2100,2200,2300,2400,2500),IDEVIC 2100 CALL LOADDW(INT(4096. - (DWELL\*20.))) IF(IDEVIC.NE.O) RETURN 2200 IF(IDEVIC.NE.O) RETURN 2300 IF(IDEVIC.NE.O) RETURN

2400 CALL LOADSW(ISWEEP) IF (IDEVIC.NE.O) RETURN CALL LOADTR (ITRIG) 2500 RETURN END

С

С С

C

С Ċ

C

Ċ

C С

C

C C C C C C C C

C C

С

C С

С

С С

С

С

C C C

C Č

С

С

# INTEGER FUNCTION TLRDY (BNAME, LNAME) File TLRDY.FOR begins on the previous line. Version Date: October 15, 1984 Paul S. Weiss Function TLRDY checks to see if a file already exists with the name in the arguments INAME and BNAME, and its length in characters is argument LNAME. WARNING: TLRDY is the limiting segment of overlay region 1. Enlarging it will expand the total size of TUF. The value of TLRDY returned is as follows: TLRDY=-1 Successfully opened file, therefore it already exists. TLRDY=1 Unsuccessfully tried to open file, therfore it does not already exist. LOGICAL\*1 BNAME(10) ASSUME FAILURE TO OPEN FILE, MEANING IT IS OK TO USE THIS NAME TLRDY = 1CHECK AND SEE IF FILE ALREADY EXISTS OPEN (UNIT=9, ERR=9000, NAME=BNAME, TYPE='OLD') 1700 TLRDY = -1RETURN 9000 END SUBROUTINE TPLOT(LINE) File TPLOT.FOR begins on the previous line. Version Date: November 8, 1985 Paul S. Weiss Subroutine TPLOT plots out the TOF scans on the HP 7470A or HP 7475A plotter.

SET UP THE HP 7470A OR HP 7475A PLOTTER CALL ASSIGN(9, 'PL:',3, 'NEW') WRITE(9,10) NESC, '.I11;;17:',NESC, '.N;19:' FORMAT('\$',A1,A9,A1,A6) 10 C XON AND XOFF HAVE NOW BEEN ESTABLISHED, ESC I11;;17 SETS THE XON C THRESHOLD TO 11 REMAINING BYTES AND XON CHARACTER TO DC1, ESC .N;19: C SETS THE XOF CHARACTER TO DC3. INITIALIZE THE PLOTTER. WRITE(9,\*) 'IN;SP 1;SC 0,4095,0,4095;' GIVE RANGE OF PLOT WRITE(9,\*) 'IP 500,600,9500,7000;' PLOT THE FIRST POINT WRITE(9,\*) 'PA ', IDSP(1), ', ', IDSP(2), ';PD;' DO 2000 I=2,256 IARG = (2\*I)IF(LINE.EQ.1) WRITE(9,\*) 'PA ', IDSP(IARG-1), ', ', IDSP(IARG), '; ' IF (LINE.EQ.O) WRITE (9,\*) 'PU; PA', IDSP (IARG-1) 1,',',IDSP(IARG),';PD;' CONTINUE 2000 PEN UP, RETURN PEN, AND MAKE PAPER EASY TO REMOVE WRITE(9,\*) 'PU;SP;PA 4000,4000;' MAKE SURE PLOT IS COMPLETED BY CLEARING OUT BUFFER **REWIND 9** CLOSE (UNIT=9) 9000 RETURN END SUBROUTINE TPRDT (IDOUT) File TPRDT.FOR begins on the previous line.

COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT

DATA NESC/27/

C C С

С

С

С С

С С

С С

С

С C Č

C

С С

С

С С

С

С С

Ը Շ October 11, 1984 Version Date: Paul S. Weiss C C C C C C C C Subroutine TPRDT outputs the data on the line printer or the terminal. The argument IDOUT determines to which device to write data, and С if to the terminal, whether or not to pause every 128 channels. C The available values of IDOUT are: С C IDOUT=-7 Write data to the terminal with pauses. C C C IDOUT=6 Write data to the line printer. IDOUT=7 Write data to the terminal. С INTEGER+4 JTRIG COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /MCSMAX/ DWLMAX, DWLMIN, MAXSWP, MAXTRG COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /SWFLAG/ ISWFLG, JTRIG С С NUNIT IS THE LOGICAL UNIT NUMBER TO WHICH DATA WILL BE WRITTEN С NUNIT = IABS(IDOUT) IST = ICHAN 900 IFIN = IST + 128IF(IFIN.GT.(ICHAN + NCHAN)) IFIN = ICHAN + NCHAN WRITE(NUNIT,1200) IST,(IFIN-1),(ACCUM(I),I=IST,(IFIN-1)) FORMAT(' Channels', I4, ' to ', I4, 16(/, 8(1X, F9.0, :))) 1200 IF(IFIN.EQ.(ICHAN + NCHAN)) GO TO 8000 IST = IFINС Ċ IF PRINTOUT IS TO TERMINAL WITH PAUSES EVERY 128 CHANNELS, QUERY C AS TO WHETHER OR NOT TO CONTINUE, WHILE PAUSING С IF(IDOUT.GT.O) GO TO 900 TYPE \*, 'Discontinue output?' ACCEPT 1000, ANS 1000 FORMAT (A1) IF(ANS.NE.'Y') GD TO 900 С С COMPLETE PRINTOUT IF TO LINE PRINTER С 8000 IF (NUNIT.EQ.6) REWIND 6 9000 RETURN END SUBROUTINE TPRII(NU) С File TPRII.FOR begins on the previous line. С

Version Date: November 8, 1985 C C C C C C C Paul S. Weiss Subroutine TPRII prints information on the line printer or into a disk file to be associated with a data file. С LOGICAL +1 CHAR(72), ITIME(8) INTEGER JDATE(3) IF (NU.EQ.9) GO TO 500 **TYPE 310** FORMAT(/, ' Include date and time?') 310 ACCEPT 1000, ANS FORMAT(A1) 1000 IF (ANS.NE.'Y') GO TO 1100 CALL IDATE(JDATE(1), JDATE(2), JDATE(3)) 500 CALL TIME (ITIME) WRITE(NU,610) JDATE(1), JDATE(2), (JDATE(3)+1900), (ITIME(I), I=1,8) 610 FORMAT(/, ' Date: ',2(I2,'/'),I4,/,' Time: ',8A1,/) С С IF INFORMATION IS TO BE WRITTEN TO DISK, ALSO PRINT IT С IF(NU.EQ.9) WRITE(6,610) JDATE(1), JDATE(2) 1, (JDATE(3)+1900), (ITIME(I), I=1,8) 1100 **TYPE 1110** FORMAT(/,' Type information in. To stop, begin a line with' 1110 1,' a = ',/)С С ENABLE LOWER CASE INPUT С CALL LCASE 1200 DO 1205 I=1,72 CHAR(I) = 01205 CONTINUE ACCEPT 1210, (CHAR(I), I=1, 72) FORMAT(72A1) 1210 IF (CHAR(1).EQ.61) GO TO 8500 WRITE(NU,1310) (CHAR(I), I=1,72) 1310 FORMAT(1X,72A1) С C IF INFORMATION IS TO BE WRITTEN TO DISK, ALSO PRINT IT C IF(NU.EQ.9) WRITE(6,1310) (CHAR(I), I=1,72) GO TO 1200 С С DISABLE LOWER CASE INPUT С 8500 CALL UCASE 9000 RETURN END

Version Date: November 8, 1985 Paul S. Weiss Subroutine TREDAS reads in a previously recorded data file. INTEGER\*4 JTRIG, JJTRIG LOGICAL\*1 CHAR(72) COMMON /DATA/ ACCUM(4096), XCORR(256), WORK(256), NEW(4096), NSWEEP COMMON /MCSMAX/ DWLMAX, DWLMIN, MAXSWP, MAXTRG COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /SWFLAG/ ISWFLG, JTRIG TYPE \*, 'Read Ascii data file into data array? Are you sure?' ACCEPT 1000, ANS 1000 FORMAT(A1) IF(ANS.NE.'Y') GO TO 9000 **TYPE 1110** FORMAT(/,' Enter the full data file name',/) 1110 CALL ASSIGN(9,,-10,'OLD') С С READ AND TYPE THROUGH THE CONTROL C THAT DESIGNATES С THE END OF THE HEADER C. TYPE \*, 'The header is as follows:' READ (9,1210, END=2000) (CHAR(I), I=1,72) 1200 1210 FORMAT(72A1) IF(CHAR(1).EQ.3) GO TO 2000 TYPE 1310, (CHAR(I), I=1,72) 1310 FORMAT(1X,72A1) GO TO 1200 TYPE 2010, ICHAN 2000 FORMAT(/,' What is the channel to which the first data point ' 2010 1,'in the stored array' 2,/,' should be added (the initial channel is now ',I4,')?') ACCEPT \*, IC IF((IC.LE.O).OR.(IC.GE.4097)) GD TD 2000 TYPE 2110, NCHAN 2100 FORMAT(/,' How many points should be read from the stored array' 1,' (the number of',/,' channels is now ',I4,').') 2110 ACCEPT +, NC IF(NC.LE.0) GO TO 2100 IF(MOD(NC,8).NE.O) NC = NC + 8С **TYPE 2210** 2210 FORMAT(' How many triggers were recorded in this file (please' 1, ' include a decimal point)?') ACCEPT +, TRIGN С
С CONVERT REAL TO INTEGER+4 С I = JAFIX(TRIGN, JJTRIG) С С ADD NUBER OF TRIGGERS TO THOSE RECORDED С I = JADD(JTRIG, JJTRIG, JTRIG)С READ AND ADD TO ACCUMULATED DATA С С OFFSET INITIAL CHANNEL TO TAKE INTO ACCOUNT THE INITIAL C C VALUES OF I AND J IN THE DO LOOPS C NOTE THAT IF THE NUMBER OF POINTS IN THE DATA FILE IS NOT DIVISIBLE C BY 8, THEN THE LAST FEW POINTS ARE LOST C IC = IC - 9D0 4000 I=1,NC/8 READ (9, \*, END=8000) (WORK (J), J=1,8) D0 3900 J=1.8 ACCUM(IC + (8\*I) + J) = ACCUM(IC + (8\*I) + J) + WORK(J)3900 CONTINUE 4000 CONTINUE CLOSE(UNIT=9) 8000 9000 RETURN END SUBROUTINE TSHOCO (NUNIT) C C C File TSHOCO.FOR begins on the previous line. Version Date: November 8, 1985 С Paul S. Weiss С C C Subroutine TSHOCO shows the present hardware and software configuration in the system. С С The argument NUNIT is the logical unit to which the С configuration will be written. С INTEGER+4 JTRIG COMMON /SWFLAG/ ISWFLG, JTRIG COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /TITANG/ TITLE(10), ANGLE С WRITE(NUNIT, 1510) (TITLE(I), I=1, 10), ANGLE 1510 FORMAT(/,1X,10A4,/,' The detector angle is ',F6.2,' degrees.') С WRITE (NUNIT, 2010) DWELL, ITRIG, ISWEEP, ICHAN, (ICHAN+NCHAN-1), EDELAY

2010 FORMAT(//,' The 4096 Channel Multichannel Scaler is in use.' 1,/,' The dwell time is ',F7.2,' microseconds per channel.' 2,/,' There are ',I5,' triggers per sweep, and ',I5,' sweeps.' 4,/,' Channels ',I4,' to ',I4,' are recorded.' 5,/,' The external delay time is ',F10.1,' microseconds.') С IF(IEXT.GT.O) WRITE(NUNIT, 2210) 2210 FORMAT(' An external dwell clock is in use.') С TRIGS = ABS(AJFLT(JTRIG))WRITE(NUNIT, 4110) TRIGS 4110 FORMAT(/,1X,F10.0,' triggers have been measured.',/) RETURN END SUBROUTINE TSTART(ILK) C File TSTART.FOR begins on the previous line. C C Version Date: July 23, 1984 С Paul S. Weiss С С Subroutine START loads all the appropriate registers, С then starts the TOF unit. С This is all done by calling the appropriate macros. С С START will clear the data arrays if ILK = 0. С INTEGER+4 JTRIG COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /DISP/ IDSP(512), INEW, ICORR, IBACK, NCORR, IADV, IRES, IQIK COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT COMMON /SWFLAG/ ISWFLG, JTRIG NSWEEP = 0IF(ILK.NE.O) GO TO 2000 CALL TZERO С С IF MCS IS RUNNING, WAIT FOR COMPLETION С 2000 ISWFLG = -1IF(IRUN(1).NE.1) GO TO 2500 С IF MCS IS RUNNING, SEND IT ONE LAST TRIGGER TO DO С С CALL LOADTR(1) 2100 IF(ISWFLG.GE.O) GO TO 2400 GO TO 2100 С С REMOVE THE EXTRA TRIGGERS COUNTED IN JTRIG, AS THE LAST SWEEP WAS С NOT SAVED С

2400	JTRIG = JTRIG - ITRIG
2500	IF (ILK.EQ.O) CALL CLEAR
	CALL TLOAD (O)
	CALL BEGIN (IEXT)
	TYPE 5000
5000	FORMAT(' MCS running')
	RETURN
	END

SUBROUTINE TSTMCS

С

C C

С

C C

C C

C C

C C

C C

С

C C

С

C

C

С

File TSTMCS.FOR begins on the previous line.

Version Date: July 23, 1984 Paul S. Weiss

Subroutine TSTMCS loads all the appropriate registers, then starts the TOF unit in a test mode which inputs the 20 MHz clock instead of Data.

This is all done by calling the appropriate macros.

INTEGER\*4 JTRIG COMMON /DATA/ ACCUM(4096),XCORR(256),WORK(256),NEW(4096),NSWEEP COMMON /DISP/ IDSP(512),INEW,ICORR,IBACK,NCORR,IADV,IRES,IQIK COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG NSWEEP = 0

ZERO DATA BUFFERS

CALL TZERO

IF MCS IS RUNNING, WAIT FOR COMPLETION

2000 ISWFLG = -1 IF(IRUN(1).NE.1) GO TO 2500

IF MCS IS RUNNING, SEND IT ONE LAST TRIGGER TO DO

CALL LOADTR(1) 2100 IF(ISWFLG.GE.O) GO TO 2500 GO TO 2100

C REMOVE THE EXTRA TRIGGERS COUNTED IN JTRIG, AS THE LAST SWEEP WAS C NOT SAVED

2500 CALL CLEAR

CALL TLOAD(0) 3000 **TYPE 3100** FORMAT(/, ' Test Mode? Are you sure?') 3100 ACCEPT 1000, ANS FORMAT(A1) 1000 IF(ANS.NE.'Y') GO TO 9000 С С START MCS WITH TEST BIT SET HIGH IN CONTROL STATUS WORD С CALL BEGIN(32) **TYPE 5000** 5000 FORMAT(' MCS running') 9000 RETURN END SUBROUTINE TTERM(N) 0000000 File TTERM.FOR begins on the previous line. Version Date: March 10, 1984 Paul S. Weiss Subroutine TTERM changes the terminal setting. COMMON /TERM/ ITERM TYPE 100, ITERM FORMAT(/,' The terminal you are using is currently ' 100 1, 'listed as a VT', I3,//) IF(N.GT.O) GO TO 2000 **TYPE 200** 200 FORMAT(' To change the terminal listing, type CU',/) RETURN **TYPE 2100** 2000 if your terminal is compatible with:' 2100 FORMAT(' Enter 3,/,' 52 VT52' 4,/,' 5,/,' 100 VT100' 0 none of the above',/) ACCEPT +, ITERM TYPE 100, ITERM RETURN END SUBROUTINE TWRA(IOPT) С File TWRA.FOR begins on the previous line. С Ĉ Version Date: October 16, 1984 С Paul S. Weiss

С С С Subroutine TWRA writes an ASCII file to disk. It includes a header С terminated with a <sup>°</sup>C. С CCCCCCCCC The argument IOPT determines whether or not to write a full header to the file, as follows: IOPT=-1 Write no header. IOPT=0. Write a full header, including the title, detector angle, configuration, and any other important information. IOPT=1 Write only information entered by the user. С С LOGICAL +1 BNAME (10), B, OTHER INTEGER INAME(6), TLRDY COMMON /DATA/ ACCUM (4096), XCORR (256), WORK (256), NEW (4096), NSWEEP COMMON /SETS/ DWELL, ICHAN, ISWEEP, ITRIG, EDELAY, NCHAN, IOFSET, IEXT EQUIVALENCE (INAME(1), BNAME(1)) С С GET NAME FOR FILE С 100 **TYPE 110** 110 FORMAT(/,' Enter the six letter data file name.',/) ACCEPT 210, (INAME(I), I=1,3) 210 FORMAT(3A2)  $INAME(\dot{4}) = \dot{,} T'$ INAME(5) = 'UF'INAME(6) = 0С С IF THE FIRST LETTER IS A SPACE, OR ANY CHARACTER OTHER THAN A LETTER С GO BACK AND GET THE CORRECT FILENAME С IF(BNAME(1).LT.64) G0 T0 100 С С IF THERE ARE FEWER THAN SIX LETTERS, SHIFT THE EXTENSION DOWN С LNAME = 10DO 600 I=1.6 B = BNAME(7-I)С Ը Շ IF THE CHARACTER IN QUESTION IS A LETTER OR NUMBER, SAVE IT OTHERWISE ELIMINATE С IF((B.GT.64).OR. ((B.GT.47).AND.(B.LT.58))) GO TO 700 С С KEEP TRACK OF NAME LENGTH

C LNAME = 10 - ID0 550 J=7-I,10-I BNAME(J) = BNAME(J+1)550 CONTINUE BNAME(11-I) = 0600 CONTINUE С С SEE IF FILE ALREADY EXISTS, IF SO, SEE IF IT CAN BE WRITTEN OVER, С IF IT MUST BE PRESERVED, GET A NEW FILE NAME С 700 I = TLRDY(INAME, BNAME, LNAME)С С IF FILE DOES NOT EXIST, GO AHEAD AND OPEN IT С IF(I.GT.0) G0 T0 1100 С С IF IT DOES EXIST, CLOSE IT AND SEE IF IT IS TO BE WRITTEN OVER С CLOSE(UNIT=9,DISPOSE='SAVE') С С IF FILE IS NOT TO BE WRITTEN OVER, GET A NEW NAME С TYPE 810, (BNAME(I), I=1, LNAME), 7, 7, 7, 7, 7 810 FORMAT(1X,10A1,' already exists! 1,'Would you like to write over it?',5A1) ACCEPT 1000, ANS 1000 FORMAT(A1) С С IF THE FILE ALREADY EXISTS, AND IS NOT TO BE WRITTEN OVER, С RETURN A VALUE OF O, SO THAT A NEW FILE NAME CAN BE ENTERED С IF(ANS.NE.'Y') GO TO 100 С С OTHERWISE, OPEN FILE С 1100 CALL ASSIGN(9, INAME, LNAME, 'NEW', 'CC') 2200 2210 FORMAT(' The file: ',10A1 1,' has been opened successfully!',8A1) С С GO TO THE TOP OF THE PRINTER PAGE, AND PRINT FILE NAME С PRINT 3010, 12, (BNAME(I), I=1, LNAME) 3010 FORMAT(1X,11A1) С С DECIDE WHAT TYPE OF HEADER TO WRITE С IF(I0PT) 3300,3100,3200 C

WRITE CONFIGURATION TO FILE AND PRINTER С С 3100 CALL TSHOCO(9) CALL TSHOCO(6) TYPE \*, 'Write any additional information to file?' 3200 ACCEPT 1000, ANS С Ċ CALL SUBROUTINE TPRII TO WRITE A HEADER TO THE DATA FILE С IF (ANS.EQ.'Y') CALL TPRII(9) С С WRITE A CONTROL C TO DESIGNATE THE END OF THE HEADER С 3300 WRITE(9,3301) 3 3301 FORMAT(1X,A1) С C C WRITE THE DATA TO THE DATA FILE FROM THE INITIAL CHANNEL, THROUGH THE NUMBER OF CHANNELS REQUESTED С NC = NCHAN/8IF(MOD(NCHAN, 8).NE.O) NC = NC + 1DO 4500 I=1,NC IBASE = ICHAN + (8 \* (I-1)) - 1WRITE(9,4010) (ACCUM(IBASE + J), J=1,8) FORMAT(8(1X, F9.0))4010 4500 CONTINUE С С CLOSE FILE С CLOSE(UNIT=9, DISPOSE='SAVE') С С ALSO, PRINT DATA С CALL TPRDT(6) RETURN END SUBROUTINE TZERO C C C C C File TZERO.FOR begins on the previous line. February 22, 1984 Version Date: Paul S. Weiss Ĉ С Subroutine ZERO zeroes the arrays used to store data. С INTEGER\*4 JTRIG

COMMON /DATA/ ACCUM(4096), XCORR(256), WORK(256), NEW(4096), NSWEEP COMMON /MCSMAX/ DWLMAX, DWLMIN, MAXSWP, MAXTRG COMMON /SETS/ DWELL,ICHAN,ISWEEP,ITRIG,EDELAY,NCHAN,IOFSET,IEXT COMMON /SWFLAG/ ISWFLG,JTRIG D0 500 I=1,256 XCORR(I) = 0. WORK(I) = 0. CONTINUE D0 600 I=1,4096 NEW(I) = 0 ACCUM(I) = 0. CONTINUE END

600

.TITLE NEW LBL TOF MACROS

File XBL.MAC begins on the previous line.

Macros for using the new LBL multichannel scaler

Version Date: November 5, 1985 Paul S. Weiss

.SBTTL GLOBALS

TLEAVE, TANAL, TCHAN, TLOAD, TDSMOD, TSTMCS . GLOBL . GLOBL TPLOT, TPRDT, TSHOCO, TSTART, TWRA, TCOMM . GLOBL TDISP, TDONE, TCORR, TCHCOR, TZERO, MCSINT, DSPINT TPRII, TDEF . GLOBL INIT, RINIT, LOADDW, LOADSW, LOADTR . GLOBL . GLOBL BEGIN, CLEAR, DISPL, ENSPEC, EXSPEC, STOPR, READMS . GLOBL RSTART, STOPD, IRUN, LCASE, UCASE .MCALL .INTEN, .PRINT

.PAGE

;

.SBTTL ADDRESSES DSPVEC=340 DSPPRI=342 MCSVEC=410 MCSPRI=412 DATAL0=166000 DATAHI=166002 LAML0=166004 LAMHI=166006 STATU1=166010 STATU2=166012 MCSA0=166100 MCSA1=166102 MCSA2=166104 MCSA3=166106 DSPCSR=170440 DSPXYB=170442 DSPPCR=170444 DSPCAR=170446 TTRCSR=177560 TTTCSR=177564 TTTBUF=177566 . PAGE

.SBTTL MACRO SUBROUTINES

;Subroutine INIT initializes the system ;and then loads all of the interrupt vectors and ;priorities. INIT is called from the main routine of the TOF ;program. INIT: TST ;SKIP ARGUMENT COUNT (R5)+ (R5) +, O#FLAGAD MOV ;LOAD SWEEP FLAG ADDRESS MOV (R5) + , **0**#TRGAD ;LOAD TRIGGER COUNTER ADDRESS MOV #O,0#DATALO ;CLEAR TRIGGERS MOVB #18.,0#MCSA1 MOV #O,**Q**#DATALO MOVB #18.,0#MCSAO ; CLEAR MCS CSR MOV #6,0#STATU1 ;CLEAR CAMAC STATUS REGISTERS MOV #0,0#STATU2 MOV #MCSINT, 0#MCSVEC ;LOAD MCS INTERRUPT VECTOR MOV #340,0#MCSPRI SET MAXIMUM PRIORITY FOR MCS : INTERRUPT BIC #103,**0**#DSPCSR ;CLEAR DISPLAY MOV ;LOAD DISPLAY INTERRUPT VECTOR #DSPINT, 0#DSPVEC MOV #340,0#DSPPRI :LOAD DISPLAY PRIORITY PC RTS ;Subroutine RINIT initializes all the hardware on the system. RINIT: RESET ; ISSUE HARDWARE RESET (B INITL) BIS #100,0#TTRCSR RESET TERMINAL INPUT ENABLE ;Subroutine LOADDW loads the dwell counts into the proper register ;of the MCS unit. LOADDW: TST ;SKIP ARGUMENT COUNT (R5)+ 0(R5),0#DATALO MOV ;GET DWELL COUNTS #18.,**0**#MCSA2 MOVB ;LOAD DWELL COUNTS RTS PC ;Subroutine LOADSW loads the number of sweeps into a word accesible to ;the MACRO routines, SWEEPS. LOADSW: TST (R5)+ ;SKIP ARGUMENT COUNT Q(R5), O#SWEEPS MOV ; GET SWEEP COUNTS RTS PC ;Subroutine LOADTR loads the number of triggers into the trigger counter. LOADTR: TST SKIP ARGUMENT COUNT (R5)+ MOV ;GET TRIGGER COUNTS **Q**(R5),R1 MOV R1, C#TRIGS ;SAVE NUMBER OF TRIGGERS

MOV #0,R0 SUB R1,R0 MOV R0,0#DATAL0 MOV R0,0#TRIGLD MOVB #18.,0#MCSA1 RTS PC

;GET TWO'S COMPLEMENT OF TRIGGER ;COUNTS ;STORE VALUE TO BE LOADED FOR TRIGGERS ;LOAD TRIGGER COUNTS

;Subroutine CLEAR clears the trigger counter

CLEAR :

MOV 9#TRGAD,R5 MOV #0,(R5)+ MOV #0,(R5) RTS PC ;CLEAR NUMBER OF TRIGGERS

## Subroutine BEGIN starts the TOF

;BEGIN is called with one argument, the argument contains the bits that need ;to be set in the control status register to enable the test frequency ;(these are 40, 100, and 200 octal) or external clock (this is 2) ;It is not necessary to send the Acquisition and LAM enable bits.

;If BEGIN is called as an integer function, the value loaded into the CSR ; is returned.

BEGIN:	BIC	#300,0#STATU1	;DISABLE CAMAC LAM
	BIS	#1,0#STATU1	SET DII BIT, SEE CAMAC CRATE
			CONTROLLER MANUAL
	MOV	#O, <b>G</b> #DATALO	;DISABLE MCS LAM AND DATA ACQUISITION
	MOVB	#18.,0#MCSAO	
	TST	(R5) +	SKIP ARGUMENT COUNT
	MOV	#11,R0	SET BITS FOR DATA ACQUISITION AND
			;LAM ENABLE
	BIS	<b>0</b> (R5),R0	EXTERNAL DWELL OR TEST FREQUENCY?
15:	MOV	RO, O#CSRLD	STORE PROPER VALUE OF CSR FOR MCSINT
	MOV	RO, G#DATALO	LOAD DATALO
	MOVB	#18.,0#MCSAO	;LOAD CSR INTO MCS
	BIS	#301, <b>0</b> #STATU1	;ENABLE CAMAC LAMS
	RTS	PC	· · · · · · · · · · · · · · · · · · ·

;Subroutine RSTART restarts the MCS

RSTART:	BIC	#300, <b>0</b> #statu1	
	MOV	<b>0</b> #TRÍGĽD,0#DATALO	

;RELOAD TRIGGERS

#18.,**0**#MCSA1 MOVB O#CSRLD, O#DATALO MOV #18.,**0**#MCSAO START MCS AND ENABLE LAM MOVB BIS #300,0#STATU1 ; ENABLE CAMAC LAM'S PC RTS ;Function subroutine IRUN returns a 1 as its value if the MCS acquisition ;enable bit is on. IRUN: ;RETURN O IF MCS NOT RUNNING MOV #0,R0 MOVB #O, Q#MCSAO ;READ MCS CSR BIT #1,0#DATALO BEQ 15 MOV #1,R0 ;RETURN 1 IF MCS IS RUNNING 15: PC RTS ;Subroutine DISPL loads the display buffer for the DMA display unit. ;It must be called with two arguments: ;DSPAD: The address of a single precision integer array, prepared to be displayed as a point plot, with x and y coordinates stored alternately. All values range from 0 to 4095. The address of this array is stored in DSPAD. NUMDSP: The number of points to be displayed. DISPL: BIC #103,**0**#DSPCSR ;DISABLE DISPLAY INTERRUPTS TST (R5) + ;SKIP ARGUMENT COUNT MOV STORE ADDRESS OF DISPLAY BUFFER (R5) + , **0**#DSPAD MOV Q(R5), Q#DSPPCR ;LOAD POINT COUNT REGISTER MOV O(R5),O#DSPNPT ;LOAD NUMBER OF POINTS REGISTER FOR FUTURE REFERENCE NEG 0#DSPPCR ;TWO'S COMPLEMENT CODING OF PCR MOV **Q**#DSPAD, **Q**#DSPCAR ; INITIALIZE CURRENT ADDRESS REGISTER BIS #103.0#DSPCSR ; ENABLE DISPLAY INTERRUPT, AND START PC RTS ;Subroutine ENSPEC puts the monitor in special screen mode ENSPEC: BIS ;ENTER SPECIAL SCREEN MODE #10000,0#44 RTS PC ;Subroutine EXSPEC takes the monitor out of special screen mode

EXSPEC:	BIC RTS	#10000, <b>0</b> #44 PC	;LEAVE SPECIAL SCREEN MODE
; ;Subrou	tine LCA	SE enables lower case i	nput from terminal
LCASE:	BIS RTS	#40000, <b>0</b> #44 PC	
; ;Subrou	tine UCA	SE disables lower case	input from terminal
UCASE:	BIC RTS	#40000, <b>0</b> #44 PC	
; ;Subrou	tine STO	PR turns off interrupts	<b>i</b>
; STOPR:	MOV MOVB RTS	#0, <b>0</b> #DATAL0 #18.,0#MCSAO PC	;CLEAR MCS CSR
; ; ;Súbrou	tine STO	PD stops the display	
; STOPD:	CLR RTS	9#DSPCSR PC	CLEAR DISPLAY INTERRUPTS
; ;Subrou ;the ar ; ;Subrou	tine REA ray spec tine REA	DMS reads the 4096 mult ified in NEW DMS must be called with	ichannel scaler channels into n one argument,
NEW	is the	array the new data is t	to be entered into.
ŔEADMS :	TST MOV MOV MOVB MOV MOV	(R5)+ (R5),R4 #4,0#DATALO #18.,0#MCSAO #4096.,R5 #0,R3	GET NEW DATA ADDRESS ENABLE READ MEMORY AND DISABLE DATA ACQUISITION AND LAM READ ALL 4096 POINTS INTO NEW REGISTER 3 CONTAINS MEMORY ADDRESS
2\$:	MOV MOVB MOVB	R3,0#DATAL0 #18.,0#MCSA3 #0.,0#MCSA3 0#DATAL0.R2	;LOAD MEMORY ADDRESS ,READ MEMORY

MOV R3,0#DATALO MOVB #18.,0#MCSA3 #0.,0#MCSA3 MOVB CMP 0#DATAL0,R2 BNE 2\$ MOV **0**#DATAL0, (R4)+ INC R3 SOB R5,2\$ MOV #O,O#DATALO MOVB #O, O#MCSAO PC RTS

;LOAD MEMORY ADDRESS ;READ MEMORY ;WAS VALUE THE SAME TWO TIMES? ;IF NOT, READ IT TWO MORE TIMES ;IF SO, MOVE IT TO ARRAY NEW(I)

;RESET CSR

.PAGE

.SBTTL INTERRUPT SERVICE ROUTINES

;Interrupt service routine MCSINT services the Multichannel Scaler. ;It checks to see whether an overflow or an end sweep has caused ;the interrupt. Then, it adds the number of triggers taken in ;that sweep, decrements the sweep counter, and in so doing, sets ;a flag for the FORTRAN routine COMMAN to see.

MCSINT: MOVB #0,0#MCSAO ;READ CSR OF MCS MOV O#DATALO, O#OVRFLG #O, Q#DATALO MOV ;DISABLE DATA ACQUISITION AND LAM MOVB #18.,0#MCSAO BIC #300,0#STATU1 ;DISABLE CAMAC LAM'S . INTEN 0,PIC ; DROP PRIORITY :GET ADDRESS OF TRIGGER COUNTER MOV O#TRGAD, R5 ; ADD THE NUMBER OF TRIGGERS IN THE ADD **0**#TRIGS, (R5)+ ADC ;LATEST SWEEP (R5) BIT #100,0#0VRFLG :CHECK FOR END SWEEP ; IF SWEEP END DETECTED CONTINUE WITH BNE 15 ;NORMAL INTERRUPT SERVICE ;OTHERWISE SERVICE OVERFLOW READ NUMBER OF TRIGGERS REMAINING MOVB #0,0#MCSA1 MOV O#TRGAD, R5 SUB Q#DATALO, (R5) ;SUBTRACT THE NUMBER OF :TRIGGERS REMAINING ON MCS SBC 2(R5) COM (R5)+ ;MAKE INTEGER+4 NEGATIVE COM (R5) MOV ;LOAD NUMBER OF SWEEPS REMAINING 0#FLAGAD,R5

> ;DECREMENT SWEEP COUNT ;LOAD NUMBER OF SWEEPS REMAINING

15:

MOV

RTS

DEC

MOV

MOV

RTS

O#SWEEPS, (R5)

**Q#SWEEPS** 

O#FLAGAD,R5

**0**#SWEEPS, (R5)

PC

PC

;

; ;Interrupt Service Routine DSPINT is for the DMA display unit. ;The point counter is reinitialized, as is the current address register, ;and then the unit is restarted. ;				
ÓSPINT:	MOV	DSPNPT, <b>0</b> #DSPPCR		;LOAD POINT COUNTER
	MOV	DSPAD,0#DSPCAR #103,0#DSPCSR		;LOAD CURRENT ADDRESS REGISTER ;REENABLE DISPLAY INTERRUPTS
;	KI1			; AND START
,	. PAGE . SBTTL	WORDS, ARGUMENT	BLOCKS,	DATA BLOCKS, ETC.
; FLAGAD : TRGAD : CSRLD :	. WORD . WORD . WORD			
SWEEPS: TRIGS:	. WORD . WORD			
TRIGLD: DSPNPT: DSPAD	. WORD . WORD WORD			
OVRFLG: SAVR:	.WORD	4		· · · ·
OUTARG:	.BLKW .END IN	4 IT		

١.

!File CNLTUF.COM !Compile command file for Time of Flight Programs !Version Date: November 13, 1984 **!Paul S. Weiss** FORT/NOLINENUMBERS TUF.FOR FORT/NOLINENUMBERS TANAL.FOR FORT/NOLINENUMBERS TCDISP.FOR FORT/NOLINENUMBERS TCHAN.FOR FORT/NOLINENUMBERS TCHCOR.FOR FORT/NOLINENUMBERS TCMLST.FOR FORT/NOLINENUMBERS TCOMLC.FOR FORT/NOLINENUMBERS TCOMLA.FOR FORT/NOLINENUMBERS TCOMLD.FOR FORT/NOLINENUMBERS TCOMLG.FOR FORT/NOLINENUMBERS TCOMLL.FOR FORT/NOLINENUMBERS TCOMLP.FOR FORT/NOLINENUMBERS TCOMLR.FOR FORT/NOLINENUMBERS TCOMLS.FOR FORT/NOLINENUMBERS TCOMLW.FOR FORT/NOLINENUMBERS TDEF.FOR FORT/NOLINENUMBERS TDISP.FOR FORT/NOLINENUMBERS TOONE.FOR FORT/NOLINENUMBERS TDSMOD.FOR FORT/NOLINENUMBERS TINIT.FOR FORT/NOLINENUMBERS TLEAVE.FOR FORT/NOLINENUMBERS TLOAD.FOR FORT/NOLINENUMBERS TLRDY.FOR FORT/NOLINENUMBERS TPLOT.FOR FORT/NOLINENUMBERS TPRDT.FOR FORT/NOLINENUMBERS TPRII.FOR FORT/NOLINENUMBERS TREDAS.FOR FORT/NOLINENUMBERS TSHOCO.FOR FORT/NOLINENUMBERS TSTART.FOR FORT/NOLINENUMBERS TSTMCS.FOR FORT/NOLINENUMBERS TTERM.FOR FORT/NOLINENUMBERS TWRA.FOR FORT/NOLINENUMBERS TZERO.FOR MAC XBL.MAC OLTUF

!File COPTUF.COM !Copy command file for Time of Flight Programs Version Date: October 15, 1984 **!Paul S. Weiss** COPY I:TUF.FOR 0:TUF.FOR COPY I: TANAL.FOR O: TANAL.FOR COPY I: TCDISP.FOR 0: TCDISP.FOR COPY I: TCHAN.FOR 0: TCHAN.FOR COPY I: TCHCOR.FOR 0: TCHCOR.FOR COPY I:TCMLST.FOR 0:TCMLST.FOR COPY I: TCOMLA.FOR 0: TCOMLA.FOR COPY I: TCOMLC.FOR 0: TCOMLC.FOR COPY I: TCOMLD.FOR 0: TCOMLD.FOR COPY I: TCOMLG.FOR 0: TCOMLG.FOR COPY I: TCOMLL.FOR 0: TCOMLL.FOR COPY I: TCOMLP.FOR 0: TCOMLP.FOR COPY I: TCOMLR.FOR 0: TCOMLR.FOR COPY I: TCOMLS.FOR 0: TCOMLS.FOR COPY I: TCOMLW.FOR 0: TCOMLW.FOR COPY I: TDEF.FOR 0: TDEF.FOR COPY I: TDISP.FOR 0: TDISP.FOR COPY I: TDONE.FOR 0: TDONE.FOR COPY I: TDSMOD.FOR 0: TDSMOD.FOR COPY I:TINIT.FOR 0:TINIT.FOR COPY I:TLEAVE.FOR 0:TLEAVE.FOR COPY I:TLOAD.FOR 0:TLOAD.FOR COPY I:TLRDY.FOR 0:TLRDY.FOR COPY I: TPLOT.FOR 0: TPLOT.FOR COPY I: TPRDT.FOR 0: TPRDT.FOR COPY I: TPRII.FOR 0: TPRII.FOR COPY I: TREDAS.FOR 0: TREDAS.FOR COPY I:TSHOCO.FOR 0:TSHOCO.FOR COPY I: TSTART.FOR 0: TSTART.FOR COPY I: TSTMCS.FOR 0: TSTMCS.FOR COPY I:TTERM.FOR O:TTERM.FOR COPY I:TWRA.FOR O:TWRA.FOR COPY I: TZERO.FOR 0: TZERO.FOR COPY I:XBL.MAC D:XBL.MAC COPY I: CNLTUF. COM 0: CNLTUF. COM COPY I:COPTUF.COM 0:COPTUF.COM COPY I:LTUF.COM 0:LTUF.COM

!File LTUF.COM

!LINK command file for Time of Flight Program !Version Date: October 15, 1984 R LINK TUF, TUF.MAP,=TUF, TDISP, XBL/I// !TUF contains the main routine TUF, and the subroutine COMMAN. !TDISP contains the display subrroutines: TDISP, and CORR !XBL contains the MACROs: INIT, LOADDW, LOADSW, LOADTR, CLEAR, BEGIN, RSTART, IRUN, DISPL, ENSPEC, EXSPEC, STOPR, STOPD, READMS, and the interrupt service routines: MCSINT, and DSPINT TCDISP/0:1 TCHCOR/0:1 TCMLST/0:1 TCOMLA/0:1 TCOMLC/0:1 TCOMLD/0:1 TCOMLG/0:1 TCOMLL/0:1 TCOMLP/0:1 TCOMLR/0:1 TCOMLS/0:1 TCOMLW/0:1 TDEF/0:1 TDSMOD/0:1 TLEAVE/0:1 TL0AD/0:1 TLRDY/0:1 TPLOT/0:1 TPRDT/0:1 TPRII/0:1 TSHOCO/0:1 TZER0/0:1 TANAL/0:2 TCHAN/0:2 TDONE/0:2 TINIT/0:2 TREDAS/0:2 TSTART/0:2 TSTMCS/0:2 TTERM/0:2 TWRA/0:2  $\boldsymbol{H}$ \$SHORT

°C

## C. References

- 1. NIM is a standard modular instrumentation bin which supplies power only on a bus at  $\pm 12$  V,  $\pm 24$  V, and sometimes  $\pm 6$  V.
- <u>IEEE Standard Modular Instrumentation and Digital Interface System</u> (CAMAC), (Institute of Electrical and Electronic Engineers, New York, 1975).
- 3. Q-Bus refers to a DEC standard bus for the LSI-11 series computers.<sup>5</sup>
- 4. RCA Electronic Components, Harrison, New Jersey.
- Digital Equipment Corporation, Maynard, Massachusetts. For a description of the various processors used, see the <u>Processor</u> <u>Handbook</u>, Digital Equipment Corporation, Maynard, Massachusetts (1984).
- 6. Joerger Enterprises, Inc., East Northport, New York.
- see the <u>RT-11 Programmer's Reference Manual</u>, Digital Equipment
   Co., Maynard, Massachusetts (1983).
- 8. Superior Electric Co., Bristol, Connecticut.
- 9. Spectra-Physics, Mountain View, California.
- 10. Kinetic Systems, Evansport, Illinois.
- R. K. Sparks, Ph.D. Thesis, University of California, Berkeley, California (1979).
- 12. Any F command in which the 8 bit (bit 3) is set (i.e. 8-15, 24-31) and any subaddress is used will test the LAM function.

- 13. The data pulses are tied to the next 20 MHz clock pulse, and no other pulses can be recorded until this clock pulse. Thus the maximum dead time is the time between clock pulses -- 50 nsec, and the average dead time is half this -- 25 nsec. By changing to a 100 MHz clock (divided down for its other timing purposes), it would be a relatively simple matter to reduce the maximum dead time to 10 nsec, and the average dead time to 5 nsec.
- 14. A prototype unit exists that is a triple width CAMAC module, but
  the front and back are offset from the standard positions, so that
  4 CAMAC slots are required.
- 15. The MCS has not been adequately tested with a 100 nsec dwell time as yet, so that TUF uses 150 nsec as the minimum dwell time.
- 16. The program TUF arbitrarily sets the maximum dwell time at 200  $\mu\text{sec.}$
- 17. The program TUF sets the maximum number of triggers at 32767 for convenience, as described in the text.
- LBL Electronics Drawings: A22952 and A22953, Lawrence Berkeley Laboratory, Berkeley, California (1985).
- 19. See the timing diagrams for the MCS, not included here due to space limitations, LBL Electronics Drawings: A153T-1, Lawrence Berkeley Laboratory, Berkeley, California (1985).
- 20. Data Translation, Inc., Natick, Massachusetts.
- 21. Hewlett Packard, Palo Alto, California.

This report was done with support from the Department of Energy. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the Department of Energy.

Reference to a company or product name does not imply approval or recommendation of the product by the University of California or the U.S. Department of Energy to the exclusion of others that may be suitable. LAWRENCE BERKELEY LABORATORY TECHNICAL INFORMATION DEPARTMENT UNIVERSITY OF CALIFORNIA BERKELEY, CALIFORNIA 94720

7