Lawrence Berkeley National Laboratory

LBL Publications

Title

Polarization in Proton-Proton Scattering Using a Polarized Target Part I. 0.330 to 0.740 GeV Part II. 1.70 to 6.15 GeV

Permalink

https://escholarship.org/uc/item/26b2q5d4

Authors

Betz, Frederick W Arens, John F Dost, Helmut E <u>et al.</u>

Publication Date

1964-06-01

University of California Ernest O. Lawrence Radiation Laboratory

UCRE-11440

TWO-WEEK LOAN COPY

This is a Library Circulating Copy which may be borrowed for two weeks. For a personal retention copy, call Tech. Info. Division, Ext. 5545

Berkeley, California

DISCLAIMER

This document was prepared as an account of work sponsored by the United States Government. While this document is believed to contain correct information, neither the United States Government nor any agency thereof, nor the Regents of the University of California, nor any of their employees, makes any warranty, express or implied, or assumes any legal responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by its trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof, or the Regents of the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof or the Regents of the University of California. Proceedings of the International Conference on High Energy Physics-Dubna, USSR, 1964.

UCRL-11440

UNIVERSITY OF CALIFORNIA

Lawrence Radiation Laboratory Berkeley, California

AEC Contract No. W-7405-eng-48

POLARIZATION IN PROTON-PROTON SCATTERING USING A POLARIZED TARGET

Part I. 0.330 to 0.740 GeV

Frederick W. Betz, John F. Arens, Helmut E. Dost, Michel J. Hansroul, Leland E. Holloway, Claude H. Schultz, Gilbert Shapiro, Wladyslaw K. Troka

Part II. 1.70 to 6.15 GeV

Herbert M. Steiner, Frederick W. Betz, Owen Chamberlain, Byron D. Dieterle, Paul D. Grannis, Claude H. Schultz, Gilbert Shapiro, Ludwig Van Rossum, David M. Weldon

June 24, 1964

POLARIZATION IN PROTON-PROTON SCATTERING USING A POLARIZED TARGET

Part I. 0.330 to 0.740 GeV^{\dagger}

Frederick W. Betz, John F. Arens, Helmut E. Dost, Michel J. Hansroul, Leland E. Holloway, Claude H. Schultz, Gilbert Shapiro, Wladyslaw K. Troka

Part II. 1.70 to 6.15 GeV[‡]

Herbert M. Steiner, Frederick W. Betz, Owen Chamberlain, Byron D. Dieterle, Paul D. Grannis, Claude H. Schultz, Gilbert Shapiro, Ludwig Van Rossum,[§] David M. Weldon

> Lawrence Radiation Laboratory University of California Berkeley, California

(Presented by Herbert M. Steiner)

June 24, 1964

Using the Berkeley polarized-proton target, we have measured the polarization parameter $P(\theta)$ for proton-proton (p-p) scattering. The measurements were obtained at beam kinetic energies of 0.330, 0.680, and . 0.740 GeV at the 184-in. synchrocyclotron and 1.70, 2.85, 3.50, 4.00, 5.05, and 6.15 GeV at the Bevatron. The angular regions measured were from 20° to 100° center of mass; the square of the four-momentum transfer ranged from 0.1 to 0.8 $(\text{GeV/c})^2$.

By means of copper absorber, the external cyclotron beam was degraded from the maximum energy of 0.74 GeV to the minimum of 0.33; the Bevatron external proton beam was spilled at various times during the acceleration cycle. The manner in which these beams were formed makes it unlikely that they contained any significant degree of polarization, and the symmetry of the arrangement was such that no component of beam polarization normal to the scattering plane would be expected.

I. EXPERIMENTAL METHOD

-2-

For these experiments, the target consisted of a 1-inch cube of $La_2Mg_3(NO_3)_{12} \cdot 24 H_2O$, with approximately 1 percent of the La replaced by Nd^{142} . The hydrogen content was about 3% by weight and the hydrogen thickness was 0.15 g/cm². The free protons in the crystal were polarized by the dynamic-nuclear-orientation¹ technique, which for this experiment involved immersion of the target in a 1.2 °K liquid helium bath inside a constant magnetic field of 18.75 kilogauss (kG). The appropriate "forbidden" transitions were excited by microwave radiation at about 71 kMc. A small variation of the microwave frequency made it possible to reverse the direction of the proton spins.

The polarization was continuously monitored by measuring the strength of the proton magnetic resonance at the frequency v = 80 Mc. At approximately 12-hour intervals the spin system was allowed to come into thermal equilibrium with the liquid helium bath. Measurement of temperature and signal strength under these conditions gave the scale factor necessary for assigning the absolute polarization values.

The magnitude of target polarizations for these experiments ranged from 20% to 60%; the direction was reversed about every 45 minutes to minimize systematic error due to variations in beam geometry and detection efficiency.

Elastic p-p scattering events were separated kinematically from other events by counting protons in coincidence. Ten scintillation counters in an upper array were placed to catch the conjugate protons (Fig. 1). Acceptable events were required to satisfy the criteria: (a) coincidence in D_d , D_o , D_t , U_o ; (b) one and only one of the counters $a_0 - a_9$; and (c) one and only one of the counters $\beta_0 - \beta_9$. Each event detected caused a count to be stored in one of

UCRL-11440

100 channels of a pulse-height analyzer, the address being determined by the particular combination of a and β counters.

-3-

The background counting rate was continuously monitored by the coincidences between upper array counters and lower array counters for which elastic p-p events were kinematically impossible. In addition, data were taken with a dummy target that consisted of elements similar to those of the crystal but with no free protons. In this way the background caused by quasielastic scattering from the nonhydrogenous materials could be evaluated. Figure 2 shows the magnitude of the elastic-to-background counting ratio for one case (6 GeV).

The lower limit on momentum transfer for which polarization could be measured was determined by the range of the recoil protons that could reach the lower counter array. The upper limit was set both by the decrease in the differential cross section and space restrictions imposed by the magnet yoke.

The polarization parameter is related to the p-p differential cross section by

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{pol}} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{unpol}} [1 + P(\theta)P_T], \qquad (1)$$

where P_T is the target polarization. The data were analyzed by means of a least-squares fit to (1) after a proper background subtraction was made.

II. RESULTS

In Fig. 3, the measured values of $P(\theta)$ are shown as a function both of center-of-mass scattering angle θ and square of four-momentum transfer for each of the incident-proton energies. Our data at 0.330 GeV are compared with those of Chamberlain, Segrè, Tripp, Wiegand, and Ypsilantis² at 0.315 GeV, obtained by the usual double-scattering techniques. In addition to the random errors ascribed to counting statistics and to measurement of the strength of the proton magnetic resonance during the run, there were the following systematic effects. A systematic error common to all angles and energies is due to error in measuring the strength of the proton-magnetic-resonance signal at thermal equilibrium and in measuring the temperature at equilibrium.

-4-

The data must be corrected to account for the fact that the target crystal is nonuniformly polarized. This nonuniformity is due to the presence of temperature gradients within the crystal, radiation damage in the region of high beam intensity, and limited penetration of the microwave radiation to the interior of the crystal. In order to estimate the size of this correction, the variations in beam density and detection efficiency of the rf system across the crystal were folded in with an assumed distribution of target polarization. These results were compared with data taken with a 1/4-inch-diameter beam spot irradiating various portions of the crystal. For the runs at 1.70 to 6.15 GeV, a 15% positive correction was made to $P(\theta)$ and a 10% systematic error is due to this effect.

For the runs at 0.330 to 0.740 GeV, the thermal contact of the crystals to the helium bath was improved and the rf-detection-system sensitivity made more uniform, so that the correction was estimated to be less than 4% and was not applied to the data. The insert for each plot in Fig. 3 gives the total systematic error for that energy.

0.

III. DISCUSSION

-5-

Calculations based on Regge poles have given expressions for protonproton polarization in the limit of high energy and small momentum transfer. ^{3, 4} Consideration of only the Pomeranchuk pole and the nearest neighboring pole that communicates with the (pp) system gives $P \propto s \begin{pmatrix} a_n - a_p \end{pmatrix}$ [large s and small fixed (-t)]. Here t is the square of the four-momentum transfer given by $t = -2q^2(1 - \cos\theta)$, q being the c.m. momentum; s is the invariant mass squared, $s = 4M^2 + 4q^2$; a_p and a_n are the positions of the Pomeranchon and its nearest neighbor at low momentum transfer. Polarization is shown for (-t) = 0.28 (GeV/c)² in Fig. 4. The point at s = 20.3 (GeV)² is from reference 11. As shown in Fig. 4, there is not a good fit to any powerlaw behavior; however, it is not clear that these measurements can be considered asymptotic in energy.

Figure 5 shows the variation of maximum polarization over a range of 1.5 decades in incident beam energy. The values shown below 0.30 GeV are representative of several measurements made in this region.

FOOTNOTES AND REFERENCES

Work supported in part by the U. S. Atomic Energy Commission.

⁴Abstract for Part II submitted as UCRL-11439 Abstract.

³Now at Centre d'Etude, Nucleaire, Saclay.

1. For a discussion of the dynamic nuclear-orientation technique as applied to polarized targets and nuclear scattering experiments, see Owen Chamberlain, Claude Schultz, and Gilbert Shapiro, Use of a Polarized Proton Target in High Energy Scattering Experiments, Lawrence Radiation Laboratory Report UCRL-11438 (unpublished); also in <u>Proceedings of the 1964</u> <u>International Conference on High Energy Physics</u>, Dubna, U.S.S.R.,

August 5-15, 1964.

6

ີ ງ

2. O. Chamberlain, E. Segrè, R. D. Tripp, C. Wiegand, and T. Ypsilantis, Phys. Rev. 105, 288 (1957).

3. Ivan J. Muzinich, Phys. Rev. Letters 9, 475 (1962).

4. Y. Hara, Phys. Letters 2, 246 (1962).

- 5. J. N. Palmieri, A. M. Cormack, N. F. Ramsey, and R. Wilson Ann. Phys. (N.Y.) 5, 299 (1958).
- 6. J. Tinlot and R. E. Warner, Phys. Rev. 124, 890 (1964).
- 7. J. A. Kane, R. A. Stallwood, R. B. Sutton, T. M. Fields, J. G. Fox, Phys. Rev. 95, 1694 (1954).
- M. G. Mescheryakov, S. B. Nurashev, G. D. Stoletov, Zh. Eksperim. i Teor. Fiz. 33, 37 (1957).
- 9. R. J. Homer, W. K. MacFarlane, A. W. O'Dell, E. J. Sacharidas, and G. H. Eaton, Nuovo Cimento 23, 690 (1962).

10. P. Bareyre, J. F. Detoef, L. W. Smith, R. D. Tripp, and L. Van Rossum, Nuovo Cimento 20, 1049 (1961).

-7-

V. P. Kanavets, I. I. Levintov, B. V. Morozov, M. D. Shafranov,
Zh. Eksperim. i. Teor. Fiz. 45, 1272 (1963).

FIGURE CAPTIONS

-8-

Fig. 1. Diagram of experimental arrangement.

Fig. 2. Example of elastic-to-background counting ratio at 6 GeV.

Fig. 3. Polarization as a function of center-of-mass angle θ and four momentum transfer squared (-t) for various beam energies $T_{\rm p}$. RSE means relative systematic error. (a) Δ , Data from

Ref. 2; O, data from this experiment.

Fig. 4. Log P vs log s; The errors include systematic error. (a) $P \propto s^{-0.32}$ (b) $P \propto s^{-1.27}$

Fig. 5. Maximum polarization as a function of beam energy T_p. Values from this experiment include systematic error. O, data from this experiment; ∇, data from Ref. 5; Q, data from Ref. 6; Δ, data from Ref. 2; □, data from Ref. 7; Q, data from Ref. 8; □, data from Ref. 9; Δ, data from Ref. 10; ∇, data from Ref. 14.









Fig. 3. Polarization as a function of center-of-mass angle θ and four momentum transfer squared (-t) for various beam energies T p. RSE means relative systematic error. (a) \triangle , Data from Ref. 2; O, data from this experiment.



Fig. 4



()

¥

MUB-3318

Fig. 5

This report was prepared as an account of Government sponsored work. Neither the United States, nor the Commission, nor any person acting on behalf of the Commission:

- A. Makes any warranty or representation, expressed or implied, with respect to the accuracy, completeness, or usefulness of the information contained in this report, or that the use of any information, apparatus, method, or process disclosed in this report may not infringe privately owned rights; or
- B. Assumes any liabilities with respect to the use of, or for damages resulting from the use of any information, apparatus, method, or process disclosed in this report.

As used in the above, "person acting on behalf of the Commission" includes any employee or contractor of the Commission, or employee of such contractor, to the extent that such employee or contractor of the Commission, or employee of such contractor prepares, disseminates, or provides access to, any information pursuant to his employment or contract with the Commission, or his employment with such contractor. and the second and and a second se Second الم المراجع المراجع المراجع ومن المراجع المراجع