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1
2 **Contrasting impacts of the South Pacific Split Jet and the Southern Annular Mode**
3 **modulation on Southern Ocean circulation and biogeochemistry**
4

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15 **Key Points:**

- 16 • Ocean changes associated with wintertime South Pacific Split Jet is contrasted against
17 those associated with Southern Annular Mode
- 18 • Pronounced but different changes are seen to the South Pacific subtropical gyre, Drake
19 Passage throughflow, and SAMW/AAIW formation
- 20 • Changes to ventilation of deep waters occur, but impact on atmospheric CO₂ is limited by
21 compensating effects of biology and solubility
22

23 **Abstract**

24 A recent hypothesis postulated that paleoclimate changes to the Southern Hemisphere westerlies
25 were characterized by the modulation of the wintertime South Pacific Split Jet. We explore this
26 hypothesis further through simulating changes to the ocean circulation from Split Jet modulation,
27 contrasting them against changes associated with the wintertime Southern Annular Mode (SAM).
28 Three responses distinguish the Split Jet from the SAM impact on ocean circulation. (i) A
29 weaker Split Jet strengthens the South Pacific subtropical gyre, leading to stronger western
30 boundary currents and warming of the SSTs surrounding New Zealand. (ii) A positive SAM
31 leads to an increase in the Antarctic Circumpolar current and specifically Drake Passage
32 throughflow; and (iii) A weaker Split Jet leads to increased formation of Subantarctic Mode Water
33 whereas a positive SAM leads to increased Antarctic Intermediate Water. Both a weaker Split Jet
34 and positive SAM lead to increased Southern Ocean meridional overturning circulation, though
35 it is more pronounced for the latter. However, enhanced ventilation of deep water in both cases
36 increases atmospheric pCO₂ by only 1-3 ppm, because the associated cooling and efficient
37 nutrient utilization in the model effectively negates the venting of deep ocean carbon. Both a
38 weaker Split Jet and positive SAM enhance oxygenation of the deep ocean and intermediate
39 waters but diminish oxygenation of the eastern equatorial Pacific. Our results provide guidance
40 to distinguish SAM-like changes from Split Jet-like changes in paleoceanographic records, and
41 we discuss the case of early deglacial transition to Heinrich 1.

42 **Index terms:** 1635 Oceans, 4901 Abrupt/rapid climate change, 4912 Biogeochemical cycles,
43 processes, and modeling, 4223 Descriptive and regional oceanography, 4576 Western boundary
44 currents

45

46 1. **Introduction**

47 Changes to the southern hemisphere westerlies and their resulting impact on the ocean
48 circulation feature prominently in current conceptualization and debates on paleoclimate
49 changes. The pioneering studies are attributed to Toggweiler [*Toggweiler and Samuels, 1998*],
50 who showed that in the limit of no vertical mixing, a large-scale ocean overturning circulation
51 could be maintained by a wind-driven Antarctic circumpolar current. The wind-driven
52 upwelling brings up waters from the deep at the latitudes of the Drake Passage, and in so doing
53 ventilate the deep ocean. This ventilation brings CO₂-rich deep ocean waters to the surface and
54 exerts strong control on atmospheric CO₂.

55 Virtually all studies linking southern hemisphere westerlies to ocean circulation and
56 paleoclimate change assume (implicitly or explicitly) a zonally symmetric change to the
57 westerlies, usually a meridional shift (e.g. *Toggweiler 2009*) but also increase in the intensity or
58 variation in its meridional width (e.g. *Moreno et al., 2010*). This conceptual model follows the
59 Southern Annular Mode (SAM) variation in the westerlies that dominates current literature in
60 dynamical meteorology and climate (e.g. *Cai et al., 2005; Marshall, 2003; Thompson and*
61 *Wallace, 2000*). However, as paleoproxies occupy a limited number of discrete locations, few
62 studies question this assumption because of the inherent difficulty in distinguishing between
63 zonally symmetric from other kinds of westerly wind changes (there are a few notable
64 exceptions, including *Fletcher and Moreno [2011]* and *Bostock et al. [2015]*).

65 In a recent study [*Chiang et al., 2014*], we auditioned another type of change for southern
66 hemisphere westerly changes in paleoclimate: the modulation of the South Pacific Split Jet.
67 Southern hemisphere westerlies are largely zonally symmetric in the austral summer, but in
68 winter exhibit a pronounced zonal asymmetry (Figure 1a). In this season, there is a single

69 westerly jet over the Indian sector, but then splits in to a northern subtropical and a southern
70 subpolar branch at the longitudes of Australia. This split continues across the South Pacific,
71 until the subpolar branch weakens, and a single jet again emerges east of South America. This is
72 the so-called Split Jet [*Bals-Elsholz et al.*, 2001; *Taljaard*, 1972].

73 Current dynamical explanations link the zonal asymmetry in the Southern Hemisphere
74 wintertime westerlies to the presence of the strong summer Asian monsoon [*Nakamura and*
75 *Shimpo*, 2004]. The northern-positioned convection induces a Hadley circulation with a strong
76 southern branch, that in turn drives a strong southern subtropical jet over the Australia/South
77 Pacific sector through the meridional advection of angular momentum by the upper branch of the
78 cell (Figure 1a). This strong subtropical jet acts as waveguide, such that atmospheric transient
79 eddies from the Indian Ocean sector entering the South Pacific preferentially get swept into this
80 jet [*Nakamura and Shimpo*, 2004]. Otherwise, these transient eddies would continue to
81 propagate zonally over the South Pacific midlatitudes, resulting in strengthened midlatitude
82 westerlies through wave-mean flow interactions. In general, the zonal asymmetry in the
83 southern hemisphere storm track has been attributed to the zonal asymmetry in the tropical SST
84 and its resulting effect on tropical convection and planetary wave propagation [*Inatsu and*
85 *Hoskins*, 2004].

86 Our hypothesis, as stated in Chiang et al. (2014), is that paleoclimate changes to the
87 southern hemisphere westerlies are characterized by a modulation of this Split Jet. Specifically,
88 during the cold North Atlantic stadial periods, the Split Jet weakens and the southern hemisphere
89 wintertime westerlies more closely resemble a single zonally-symmetric jet. This contrasts with
90 current notions that southern hemisphere westerlies shifted southward (i.e. a more positive SAM)
91 during North Atlantic stadials [e.g. *Toggweiler*, 2009]. The reason for a weaker Split Jet is that

92 North Atlantic stadials are associated with a weakening of the Northern hemisphere summer
93 monsoons and a southward shift of the ITCZ (e.g. *Chiang and Friedman, 2012*, and references
94 therein). The South Pacific subtropical jet thus weakens, and by the dynamical arguments above
95 this ultimately leads to a weakening of the Split Jet. [*Chiang et al., 2014*] lent support to the
96 hypothesis by noting that land-based proxy records of abrupt changes over New Zealand and
97 Patagonia appear to be consistent with a modulation of the Split Jet. *Bostock et al. [2015]* also
98 proposed a modulation of the Split Jet as a potential explanation for changes to the latitude of the
99 subtropical front south of New Zealand.

100 The most intriguing implication of the hypothesis is the change to the wind stress over
101 the ocean and potential ramifications for ocean circulation; this is the focus of this paper. As
102 shown in *Chiang et al. [2014]*, a weaker Split Jet leads to a change in the wind stress curl over
103 the South Pacific, specifically a marked southward shift in the zero wind stress curl line from the
104 subtropical latitudes towards the southern tip of South America (see Figure 5 of *Chiang et al.*
105 *2014*, and also Figure 2b). This is a remarkable change to the wind stress pattern, with strong
106 implications for the wind-driven ocean circulation in the South Pacific and in particular the
107 subtropical gyre.

108 The goal of this study is to explore the consequences of a weakened Split Jet to the
109 Southern Ocean circulation and biogeochemistry, specifically carbon and oxygen. We undertake
110 ocean model simulations to explore the ocean circulation response to the specified westerly
111 changes imposed on the model ocean. We contrast the simulated changes to those induced by a
112 wintertime modulation of the SAM, to contrast the impacts of the two competing hypotheses for
113 Southern Hemisphere westerlies on ocean circulation. We will then compare the changes to
114 ocean paleoproxy records to assess how each are consistent (or not) to records of Southern

115 Ocean response the early deglacial period, specifically climate changes associated with the
116 transition to Heinrich 1.

117 2. **Materials and Methods**

118 2.1 Model

119 We undertake ocean-only simulations with the Parallel Ocean Model version 2 that comes as part
120 of the Community Earth System Model (CESM) 1.2.2 [Danabasoglu *et al.*, 2011]. The
121 component set used is a scientifically validated and supported compset of CESM
122 “2000_DATM%NYF_SLND_CICE_POP2_DROF%NYF_SGLC_SWAV”. It includes
123 prognostic ocean (POP2) and sea ice (CICE) modules, but with a prescribed ‘data’ atmosphere
124 (DATM) and ‘data’ runoff (DROF). We use the standard gx1v6 ‘displaced pole’ grid. The
125 CESM has been used extensively, including studies examining the ocean response to southern
126 hemisphere westerlies (e.g. *Gent and Danabasoglu, 2011*)

127 We include a simple biogeochemistry (BGC) module developed specifically for POP2
128 and made available to us by Keith Lindsay at NCAR. The BGC module follows the “BIOTIC”
129 formulation of OCMIP and features the so-called “nutrient-restoring” or diagnostic export
130 production [Najjar *et al.*, 2007]. Essentially, phosphate concentration in the surface ocean of the
131 model is nudged toward, or restored to, observed phosphate climatology on a time scale of
132 biological production of 30 days. For example, in upwelling regions, where phosphate supply is
133 large, the nudging and thus the diagnosed export production are large. Export production is
134 partitioned into particulate and dissolved phases. State variables in this BGC module include
135 phosphorus (PO₄, DOP, POP), carbon (DIC, DOC, POC), oxygen, and alkalinity. The
136 stoichiometry that relates the elements in organic matter is fixed: P:C:O₂=1:117:-170. Biological
137 production of organic matter near the surface removes these elements in exactly these ratios.

138 Respiration of organic matter in the ocean interior releases these elements in the same ratios.
139 Wind-driven gas exchange formulation allows CO₂ and O₂ to be exchanged between the
140 atmosphere and surface ocean. Atmospheric pCO₂ is thus a prognostic variable; it is 284.7 ppm
141 in the equilibrium state. Atmospheric O₂ is fixed at the modern mixing ratio. The OCMIP
142 formulation continues to respire organic matter even after O₂ runs out, assuming that nitrate,
143 which is not a model variable, would serve as the electron acceptor in respiration. This BGC
144 module has been thoroughly examined in the intercomparison project OCMIP Phase 2 [Najjar et
145 al., 2007]. It is appropriate for this study, because export production and associated net air-sea
146 CO₂ flux would respond to changes in wind-induced upwelling.

147 2.2 Forcing Data

148 The data atmosphere of CESM uses the ‘Co-ordinated ocean-ice reference experiments’ version
149 2 (CORE v2) normal year forcing [*Large and Yeager, 2009*] for driving the ocean model. They
150 are 6-hourly fields derived from several observational/reanalysis datasets, and intended as a
151 common ocean driver to compare different ocean model simulations. In addition to the ‘normal
152 year’ forcing, CORE v2 also includes forcings for individual years from 1948 through 2009.
153 For our purposes, we restrict ourselves to forcing data from 1979 onwards to avoid a potential
154 discontinuity resulting from the assimilation of satellite data starting in 1979.

155 To apply perpetual weak or strong Split Jet scenarios, we first derive surface spatial
156 patterns associated with Split Jet variation. We derive an interannual index of the Split Jet from
157 1979-2015 (see section 3), then use it to obtain a spatial pattern of the forcing field every 6-
158 hours. Using the zonal wind stress (τ_{zx}) for 00H of June 1 as an example, first a 21-day
159 average of τ_{zx} centered on 00H is formed (so taking all 00H time points 10 days before and 10
160 days after, and averaging), for each individual year in the CORE v2 τ_{zx} fields. The 21-day

161 average is applied to remove weather noise and obtain a reasonable estimate of the climate state
162 around that time. We then regress the normalized interannual Split Jet index against each
163 gridpoint of tau_x for 00H over 1979-2009 (the CORE v2 interannual data is only available until
164 2009) to obtain a spatial field corresponding to the change in the tau_x for 1 standard deviation
165 change in the Split Jet index. This spatial pattern is then multiplied by 2 and then added to the
166 CORE v2 normal year tau_x to form a representative strong Split Jet tau_x field for 00H on June 1;
167 to get the weak Split Jet field, we multiply by -2 and then add to the CORE v2 normal year; the
168 magnitude of our forcing anomalies is thus set to ± 2 standard deviations of the interannual
169 variability. Our study differs from many others investigating the ocean circulation changes to
170 southern hemisphere westerly changes in that the wind perturbations we apply- while large - are
171 spatially realistic and physically realizable.

172 Since the Split Jet and its variability only occur in the austral winter, we produce these
173 forcing fields every 6-hours from 00H of June 1 through 18H of September 30 only. This is
174 repeated for each forcing field in the CORE v2, including: zonal and meridional surface wind
175 stress, zonal and meridional winds at 10m, air density, temperature and specific humidity at 10m,
176 and sea level pressure. These fields are then used to drive the POP2 model for the strong and
177 weak Split Jet cases. The same is also done to the shortwave and longwave radiative forcing in
178 CORE v2, but they are at daily resolution; for those, we apply the same procedure but to daily
179 fields. Finally, the precipitation forcing in CORE v2 only has monthly fields, so we do the
180 regression for each month from June through September.

181 For forcings associated with opposite phases (positive and negative) of the SAM, we
182 follow the same procedure as for the Split Jet but using an interannual index for the annular
183 mode-like variation; section 3 describes how the SAM index is obtained.

184 2.3 Experiments

185 We first ran the POP under the standard CORE v2 normal year forcing for 500 years to obtain a
186 reasonable starting point. We then branch off the simulation into two runs, one with perpetual
187 weak Split Jet conditions and the other with perpetual strong Split Jet forcing. We run these two
188 cases for 600 years each so that the wind-driven surface circulation can come into quasi-
189 equilibrium, and changes to the deep ocean circulation are captured; examination of the
190 timeseries of southern ocean meridional overturning circulation strength and Drake Passage
191 throughflow indicate that this is the case (not shown). Our assessment of the impact of the Split
192 Jet modulation on ocean circulation and biogeochemistry is done by comparing the difference in
193 the annual mean climatology over the last 20 years of each simulation. Prior to analysis, the
194 ocean model output was regridded from its native gx1v6 horizontal resolution to a uniform $1^\circ \times 1^\circ$
195 grid. Mean fields reported here (in figures 5,6, and 7) are taken as the average of the weak Split
196 Jet and strong Split Jet annual mean climatologies. Finally, we undertake an analogous set of
197 two runs for the positive and negative SAM cases.

198

199 **3. Interannual variability of the Southern Hemisphere wintertime westerlies**

200 We apply an Empirical Orthogonal Function (EOF) analysis on the June-September averaged
201 zonal wind anomalies from 1979-2015 using the NCEP reanalysis to extract the dominant
202 interannual variability for the austral winter westerlies. Since the expression of the Split Jet is
203 most pronounced in the upper troposphere, we use the 3-dimensional zonal winds in the
204 calculation; this contrasts with the standard derivation of the SAM that uses the first EOF of sea
205 level pressure over the southern extratropics. The latitude range was restricted to 15°S - 80°S , and
206 the vertical co-ordinate was interpolated in 100mb intervals from 150mb to 950mb. Prior to the

207 calculation, the anomaly data was multiplied by the square root of the cosine of latitude to
208 weight for area.

209 The first 3 modes are significant (using the ‘Rule N’ test [*Preisendorfer*, 1988]) and well
210 separated from each other (applying North’s Rule of Thumb [*North et al.*, 1982]). The first
211 mode, accounting for 22.2% of the variance, has the distinct signature of the modulation of the
212 Split Jet: the primary signature is a positive loading of winds over the midlatitude South Pacific
213 (35-55°S), and negative loading to the north and south of this feature (Figure 1b): this indicates a
214 weakening of the Split Jet with weaker subtropical and subpolar jets, and stronger midlatitude
215 jet. The second mode, accounting for 16.1% of the variance, has a more zonally-symmetric
216 structure with negative loadings in the subtropical latitudes (25-45°S) and positive loadings from
217 45-70°S (Figure 1c). The mode represents a poleward shift of the midlatitude jet, resembling an
218 annular-mode like modulation. In fact, the 2nd principal component (PC2) correlates with a
219 standard index for the SAM [*Mo*, 2000] averaged over June-Sep at $r = 0.88$, significant at the
220 0.1% level; thus, PC2 is essentially the SAM index for the austral winter, and we will refer to
221 mode 2 as the SAM mode from now on. On the other hand, PC1 and 3 are not significantly
222 correlated with the SAM index. The third mode, accounting for 10.3% of the variance, is a
223 teleconnection pattern spanning from the Maritime continent across Australia to the high latitude
224 South Pacific.

225 Our EOF results indicate that we have successfully extracted the Split Jet (mode 1) from
226 SAM (mode 2) variations from data, and that they are linearly independent. We use these two
227 modes to derive anomalous forcing fields for each influence, through regression on the
228 corresponding normalized PC onto the CORE v2 interannual fields (see section 2).

229

230 4. Ocean circulation changes

231 We focus on key features of the Southern Ocean circulation that exhibit different responses
232 between the Split Jet and SAM wind forcings. For the surface circulation, they include the South
233 Pacific subtropical gyre circulation and the Antarctic Circumpolar current (ACC). For the deep
234 circulation, they include the meridional overturning circulation and deep ocean ventilation, and
235 formation of intermediate waters. Recall that the simulated ocean fields we use are the last 20
236 years of 600-year simulations, and with the same atmospheric forcing applied each year; hence,
237 the ocean circulation changes reported in this section (and section 5) can be thought of as
238 equilibrated ocean changes in response to permanent Split Jet or SAM changes.

239 4.1 Surface ocean circulation

240 4.1.1 *Subtropical gyre and currents off Eastern Australia and New Zealand*

241 ***Split Jet response:*** As shown in Chiang et al. [2014], one of the most striking features of a weak
242 Split Jet, contrasted with a strong Split Jet, is a change in the sign of the wind stress curl over the
243 subtropical South Pacific (30°S-40°S) to positive values, resulting in a southward shift of the
244 subtropical zero wind-stress curl line in the South Pacific (Figures 2a and 2b). This induces a
245 northward Sverdrup interior flow in the subtropical South Pacific that must be compensated for
246 by a stronger western boundary current, implying a strengthened and southward shifted oceanic
247 subtropical gyre. This feature is seen in the simulated ocean circulation changes, most
248 noticeably by the change to the barotropic streamfunction that shows a stronger and more
249 southward South Pacific subtropical gyre circulation for the weak Split Jet (Figure 3b). With
250 this, the southward extent of the South Pacific subtropical gyre resembles that for the South
251 Indian in the mean (Figure 3a).

252 Predictions of a stronger and more poleward-penetrating western boundary current is
253 indeed seen for the East Australian Current (figure 4b). The East Australian Current splits from
254 the Australian continental shelf at the latitude of the northern tip of New Zealand, though some
255 of the current proceeds farther south (see figure 4a for a schematic of the various currents). The
256 flow then moves eastwards to the northern tip of New Zealand as the Tasman Front, then turns
257 due south at the northern edge of North Island to form the East Auckland current, and then along
258 the eastern coast of North Island as the East Cape Current; these are the continuations of the
259 western boundary current. This boundary current finally splits away from New Zealand into the
260 southeastern Pacific by the latitude of Chatham Rise. In the weak Split Jet, these flows are
261 significantly strengthened (figure 4b). Farther south, the Antarctic Circumpolar Current south of
262 Tasmania and New Zealand is also strengthened, but this anomalous flow heads northwards
263 along the eastern coast of South Island and veers eastwards at the latitude of Chatham Rise to
264 merge with the current originating from the north.

265 A consequence of the stronger and more poleward-penetrating subtropical gyre is warmer
266 sea surface temperatures off the eastern coasts of Australia and New Zealand (Figure 4b), by up
267 to 2°C for the difference between the weak and strong Split Jet cases. This is notable as New
268 Zealand glaciers are known to have rapidly retreated during the Younger Dryas, indicating
269 warmer conditions; this is consistent with the Chiang et al. [2014] hypothesis of a weaker Split
270 Jet during that time. A stronger subtropical gyre also leads to a stronger equatorward return flow
271 over the eastern South Pacific, leading to moderately cooler SSTs there (Figure 4b).

272 Finally, there are intriguing though smaller northward shifts to the southern extent of the
273 subtropical gyre over the South Indian and South Atlantic oceans, as indicated by the changes to
274 the barotropic streamfunction (Figure 3b). The most notable consequence of this is a strong

275 cooling in the South Atlantic around 42°S at the poleward end of the South Atlantic subtropical
276 gyre boundary, near the South American coastline (figure 4b).

277 ***SAM response:*** The response of the ocean circulation to variation in the SAM produces little
278 change to the strength of the subtropical gyre circulation but does affect the East Australian
279 boundary current. The positive SAM (southward-shifted westerlies) relative to the negative
280 SAM leads to a more southward extension of the East Australian Current (Figure 4c). The
281 poleward extension is consistent with the increased positive wind stress in the interior South
282 Pacific between 40-50°S (Figure 2c, d). By Sverdrup balance, the latter implies increased
283 interior northward flow at those latitudes. A consequence of this is a reduced eastward current
284 by the Tasman Front where the East Australian Current partly separates from the coast around
285 35°S, but a stronger eastward current around 50°S where the East Australian Current merges
286 with the ACC (figure 4c). As a result, surface temperatures warm over the southeastern coast of
287 Australia and the Tasman sea, though not as pronounced as for the weaker Split Jet case (Figure
288 4b). By contrast, the SSTs east of the North Island of New Zealand cools, presumably because
289 of the reduced East Auckland and East Cape Currents; this is opposite to that seen for the weaker
290 Split Jet. This pattern of SST changes surrounding New Zealand was also obtained in an ocean
291 model study by *Cai et al.* [2005] to positive SAM winds.

292 4.1.2 *Antarctic Circumpolar Current and Drake Passage throughflow*

293 ***Split Jet response:*** There is a modest increase to ACC and Drake Passage throughflow for a
294 weaker Split Jet relative to a strong Split Jet, as indicated by the barotropic streamfunction
295 anomaly (figure 3b) that shows increased meridional gradient in streamfunction values across the
296 ACC latitudes. Drake Passage throughflow (evaluated at 69.5°W) is 116Sv in the strong Split
297 Jet case whereas in the weak Split Jet case is 133Sv, an increase of 17Sv.

298 ***SAM response:*** On the other hand, the positive SAM leads to a pronounced increase in the
299 volume transport of the Antarctic Circumpolar Current (ACC). Transport across the Drake
300 Passage increases from 100 Sv in the negative SAM run, to 142Sv, an increase of 42Sv, mainly
301 due to increased flow in the southern portion of the Drake Passage (figure 4c). The reason for
302 this increase is that the positive phase of the SAM leads to increased wind stress over the Drake
303 Passage latitudes ($\sim 60^{\circ}\text{S}$; see figure 1c). In the absence of a zonal pressure gradient to counter
304 the increased wind stress, an increased form drag at the ocean floor is required to balance it,
305 implying a stronger ACC circulation.

306 4.2 Overturning circulation

307 ***Split Jet response:*** Changes to the wind stress over the South Pacific leads to a pronounced
308 alteration in the spatial pattern of surface upwelling and downwelling. For a weak Split Jet, the
309 zone of upwelling shifts northwards over the South Pacific sector to around 50°S , and
310 downwelling is increased to the north of it (Figure 5b). This has the effect of enhancing the
311 spatial contrast of mean upwelling and downwelling over the South Pacific sector, making it
312 resemble the spatial structure of upwelling over the other ocean basins (Figure 5a). The sense is
313 that a weak Split Jet reduces the zonal asymmetry in the spatial pattern of surface upwelling and
314 downwelling.

315 The increased downwelling north of 40°S in the South Pacific leads to increased
316 subduction rate of Subantarctic Mode Water (SAMW) in the South Pacific. The rate increase is
317 suggested by a decrease in the ideal age of waters in the subduction region (ideal age is an
318 idealized tracer indicating the elapsed time since a water parcel's last exposure to the surface),
319 centered around 800m and 30°S (Figure 6b). This increased formation of SAMW does not occur
320 over the other basins, where downwelling does not increase. The change to the Pacific MOC

321 (Figure 6e) shows a slight increase to the subduction of southern subtropical surface waters into
322 the thermocline, consistent with the increased rate of SAMW formation.

323 A weaker Split Jet also leads to increased rate of Antarctic Bottom Water (AABW)
324 formation. Significantly younger ideal age values at the bottom of the South Pacific (Fig 6b)
325 suggests an increased rate of deep water formation, and this is collaborated by changes to the
326 Pacific MOC (Fig 6e) showing increased northward flow at the bottom of the South Pacific.

327 ***SAM response:*** Significant changes to the surface upwelling and downwelling also occurs, with
328 increased upwelling south of $\sim 55^{\circ}\text{S}$ and increased downwelling to the north of it (Figure 5c).
329 This occurs over all three basins; the sense is that the mean spatial pattern of upwelling and
330 downwelling over the Southern ocean is enhanced. The increased midlatitude subduction over
331 all basins leads to increased formation of Antarctic Intermediate Water (AAIW); this
332 interpretation is supported by the decrease in ideal age of waters around 40°S and 1000m in
333 depth (Figure 6c), as well as changes to the Pacific MOC (Fig 6f) showing increased subduction
334 of surface waters around $\sim 50^{\circ}\text{S}$ into the AAIW depths.

335 A more notable consequence of a positive SAM is the increased rate of bottom water
336 formation, significantly stronger than that for the weak Split Jet case. The increased rate of
337 formation of Antarctic Bottom Water (AABW) is indicated by its considerably younger ideal age
338 (Figure 6c), and the corresponding increase to the northward transport of bottom waters in the
339 Pacific as indicated by the Pacific MOC change (Figure 6f). This has consequences for the
340 oceanic carbon cycle, which we elaborate on in Section 5.

341

342 **5. Ocean biogeochemical changes**

343 The response of ocean circulation to the Split Jet and SAM wind forcings has important
344 implications for nutrient redistribution, water mass residence time, and deep water temperatures.
345 In turn, these have significant effects on biological production and air-sea gas exchange. Here we
346 focus on the effects of the ocean circulation changes to the biogeochemically important elements,
347 oxygen and carbon.

348 5.1 Oceanic Oxygen

349 Dissolved oxygen (O_2) is a sensitive indicator of ocean physics and biogeochemistry. Many areas
350 of the world surface ocean are close to O_2 saturation, because the residence time in those areas is
351 comparable to or exceeds the typical equilibration time scale of O_2 exchange with the
352 atmosphere. To first order, it is SST that determines the O_2 solubility and thus surface
353 concentration. Vertical mixing transports this surface O_2 signal into the ocean interior, where
354 respiration of organic matter consumes O_2 . Where dissolved O_2 becomes too low for oxic
355 respiration, nitrate replaces O_2 as the terminal electron acceptor in the process of denitrification.
356 This process changes the distribution and whole ocean inventory of nitrate, which limits
357 biological production in many parts of the world ocean. Thus, oceanic O_2 is not only an indicator
358 of ocean physics and biogeochemistry, but a driver of the latter at the same time.

359 In this study, increased deep ocean ventilation in both the weak Split Jet and positive
360 SAM simulations delivers dissolved O_2 in great abundance, as seen in the Pacific Ocean (Figure
361 7b,c). In the positive SAM case, the positive O_2 anomaly increases throughout the water column
362 in Southern Ocean (Figure 7c), where meridional overturning is strong and deep (Figure 6f).
363 Outside the Southern Ocean, the deep O_2 anomaly is large ($\sim 150 \mu\text{mol/kg}$) and extends to the far
364 northern extent of the Pacific basin. The same positive O_2 anomaly in the deep ocean is also
365 seen, but is weaker overall in the weak Split Jet simulations (Figure 7b). Higher in the water

366 column, there is a small positive O₂ anomaly associated with the SAMW waters in the weak
367 Split Jet case and with the AAIW waters in the positive SAM case; these are consistent with the
368 increased rate of SAMW/AAIW formation inferred from the ideal age and PMOC changes
369 (Figure 6). Spatially, the strong positive O₂ anomaly for the weak Split Jet case is restricted to
370 the South Pacific (Figure 8a). In contrast, the zonal symmetry of the forcing in the SAM runs
371 clearly translates to enhanced oxygenation occurring in all sectors of the circumpolar current
372 (Figure 8b).

373 Much of this positive O₂ anomaly can be understood in terms of apparent oxygen
374 utilization (AOU), which shows nearly the same spatial pattern as O₂ along the 150 °W section,
375 so that O₂ and AOU are tightly clustered about the 1:1 line in a scatter plot (figure not shown). In
376 a parcel of seawater, AOU is defined as the difference between the theoretical O₂ concentration
377 at saturation, calculated from the parcel's temperature, and the observed/simulated O₂
378 concentration. One assumes in AOU calculation that the parcel, when it was last at the surface,
379 was at saturation. It follows then that a parcel of interior water with AOU of zero is “fully
380 charged” with O₂, and a parcel with a large AOU has experienced significant O₂ depletion due to
381 respiration. In order to explain the positive O₂ anomalies in the weak Split Jet (Figure 7b) and
382 positive SAM (Figure 7c) simulations in terms of AOU, either the deep ventilation became more
383 vigorous, the interior respiration became smaller, or a combination of the two.

384 Our analysis shows that it is largely a more vigorous deep ventilation (Figure 6), because
385 respiration did not change much. In the OCMIP formulation, biological production largely keeps
386 up with ocean circulation changes and associated nutrient redistribution, so that global carbon
387 export in weak Split Jet (8.6 GtC yr⁻¹; GtC=10¹⁵ grams-C) is only slightly larger than strong Split
388 Jet (8.4 GtC yr⁻¹). The SAM positive simulation (8.8 GtC yr⁻¹) is somewhat larger than the

389 negative mode (8.3GtC yr^{-1}). If anything, changes in biological production in the weak Split Jet
390 and positive SAM cases should drive more respiration and thus O_2 consumption. The simulated
391 results are opposite (Figure 7).

392 Interestingly, even while the deep and intermediate waters are clearly more oxygenated in
393 the weak Split Jet and positive SAM runs (Figures 7, 8), the oxygen minimum zone in the
394 eastern equatorial Pacific Ocean is larger in those cases. The volume of water with O_2
395 concentrations below $2\ \mu\text{mol kg}^{-1}$ is larger in the weak Split Jet case ($5.36 \times 10^{15}\ \text{m}^3$) than in the
396 strong Split Jet case ($5.05 \times 10^{15}\ \text{m}^3$). It is also larger in the positive SAM case ($5.97 \times 10^{15}\ \text{m}^3$)
397 than in the negative SAM case ($4.77 \times 10^{15}\ \text{m}^3$). While the exact volumes would be overestimated
398 in the OCMIP formulation because denitrification is neglected, as already noted, the sense of
399 change should be correct. It seems puzzling then that the generally better oxygenated oceans
400 (weak Split Jet and positive SAM) have larger tropical Pacific OMZs. However, our results are
401 consistent with the theory that the southward expansion of the South Pacific subtropical gyres
402 prevents northward transport of well-ventilated thermocline waters originating from the
403 subantarctic [Getzlaff *et al.*, 2016].

404 In summary, both a weaker Split Jet and positive SAM lead to increased ventilation of
405 bottom waters, and with changes associated with positive SAM being larger. The positive SAM
406 also leads to increased ventilation of AAIW waters in all basins, whereas the weaker Split Jet
407 leads to increased ventilation of the SAMW in the South Pacific. In both cases, the oxygen
408 minimum zone region of the eastern equatorial ocean increases.

409 5.2 Atmospheric pCO_2 and oceanic DIC

410 Atmospheric pCO_2 is strongly controlled by deep ocean ventilation [Knox and McElroy, 1984;
411 Sarmiento and Toggweiler, 1984; Siegenthaler and Wenk, 1984]. Therefore, the enhanced deep

412 ventilation in the weak Split Jet (vs. strong Split Jet) and positive SAM (vs. negative SAM)
413 simulations would predict an increase in atmospheric pCO₂. The simulated response of
414 atmospheric pCO₂ in the weak Split Jet and positive SAM is indeed positive but unexpectedly
415 small. After 600 years of model integration, the positive pCO₂ anomaly is just over 1 ppm in the
416 weak Split Jet runs and about 3 ppm in the positive SAM runs.

417 In the Pacific, the zonally averaged the dissolved inorganic carbon (DIC) anomalies for
418 the two cases (Figure 7e,f) have spatial patterns very similar to those of ideal age (Figure 6b,c)
419 and O₂ (Figure 7b,c). The strong negative DIC anomaly in the deep ocean thus indicates that
420 carbon is lost by enhanced ventilation under both the weak-strong Split Jet and positive-negative
421 SAM simulations. Clearly this carbon lost from the deep ocean did not accumulate in the
422 atmosphere, since atmospheric pCO₂ changed only 1-3 ppm.

423 There are two processes that can put the carbon vented from the deep ocean back into the
424 ocean. The first is the efficient biological utilization of nutrients. This model's diagnostic
425 biological production, based on a 30-day restoring of surface PO₄, ensures that a significant
426 fraction of the nutrients and carbon upwelled in the Southern Ocean are sent back to the ocean
427 interior as organic carbon. To the extent that the restoring is not instantaneous and complete,
428 there will be some leakage of oceanic carbon as both ventilation and export production becomes
429 stronger.

430 The second process that prevents vented carbon from accumulating in the atmosphere is
431 cooling. The global ocean temperature is cooler by ~0.1 °C in the weak Split Jet case (mean
432 ocean temperature=3.58 °C) compared to the strong Split Jet case (3.67 °C). The cooling in the
433 positive SAM case (3.32 °C) compared to the negative SAM case (3.76 °C) is ~0.44 °C,
434 substantially larger. As gas solubility becomes larger with increased cooling, the weak Split Jet

435 and the positive SAM modes should increase oceanic uptake of carbon relative to their respective
436 opposite modes. In a rough estimate, we calculated the equilibrium DIC concentration with full
437 seawater carbon chemistry for the four cases using ocean mean properties of temperature (noted
438 above), alkalinity ($2418 \mu\text{eq kg}^{-1}$), salinity (34.73psu), and for equilibrium atmospheric pCO_2
439 (284.7 ppm). Multiplying the change in mean DIC with the ocean volume and converting the
440 oceanic carbon inventory to ppm, we determine the cooling effect to be ~ 8 ppm for the weak-
441 strong Split Jet case and ~ 38 ppm for the positive-negative SAM case. In other words, without
442 the cooling, atmospheric pCO_2 rise should have been much larger, 9 instead of 1 ppm in the Split
443 Jet case, and 41 instead of 3 ppm in the SAM case. This assumes the same degree of air-sea gas
444 equilibration, which is not strictly correct, because the simulated DIC anomaly does not follow
445 the change expected from PO_4 anomaly. If the air-sea CO_2 disequilibrium were constant, DIC
446 should be related to PO_4 by the organic matter C:P stoichiometry of 117, but the trend observed
447 in the model is larger. The simulated C:P ratio in the Southern Pacific is 136 in the Split Jet cases
448 and 190 in the SAM cases. The estimated cooling effect on atmospheric pCO_2 of 8 and 38 ppm
449 should thus be viewed as the upper bound.

450 Our overall interpretation is that dynamical response of atmospheric pCO_2 to the Split Jet
451 (weak-strong) and SAM (positive-negative) forcing is an increase in the atmospheric CO_2 from
452 increased ventilation of the deep oceans. This result, at least for the SAM response, is
453 qualitatively consistent with many previous studies on the link between ocean carbon and
454 Southern Hemisphere winds (e.g. *d'Orgeville et al.*, 2010; *Lovenduski et al.*, 2007; *Russell et al.*,
455 2006). The actual change would be much larger than 1-3 ppm, if biological and solubility effects
456 did not dampen the dynamical response. The combination of the deep ocean venting carbon in
457 the Southern Ocean, but the ocean largely maintaining its carbon inventory, results in a

458 redistribution of the DIC anomaly as seen in Figure 7. DIC is effectively stored outside the
459 Southern Ocean and deep ocean, where ventilation remains the same or is slightly less vigorous
460 (Figure 6). This is more clearly evident in the positive-negative SAM simulations (Figure 7f).

461

462 **6. Comparison to Paleoproxy records during the Early Deglacial Period**

463 To illustrate the applicability of the modeling results to interpreting Southern Hemisphere
464 westerly changes in past climates, we examine paleoceanographic changes during the early
465 deglacial period following the Last Glacial Maximum (LGM), linked to the Heinrich 1 event.
466 The period during Heinrich 1 is particularly interesting for examining changes to the Southern
467 Hemisphere westerlies and Southern Ocean circulation. This period marks the beginning of the
468 last glacial termination when CO₂ started rising and Northern Hemisphere ice sheets started
469 melting (*Denton et al. 2010*). While Heinrich 1 occurred in the North Atlantic with armadas of
470 icebergs entering the ocean there, studies have shown synchronous climate changes in distant
471 Southern Hemisphere locations timed to Heinrich 1, in particular changes in the Southern
472 Hemisphere westerlies and ocean circulation (e.g. *Lamy et al. 2007, Anderson et al., 2009*).
473 Increase to upwelling over the Southern Ocean and resulting flux of carbon from the deep ocean
474 to the atmosphere during Heinrich 1 are thought to provide the initial deglacial rise in CO₂ (*Lamy*
475 *et la. 2007, Anderson et al. 2009, Denton et al. 2010*).

476 We focus on specific circulation features as identified by the simulations, and which may
477 distinguish between a Split Jet response or SAM response; as it turns out, paleoproxy records
478 during Heinrich 1 suggest pronounced changes for all the features of interest. We note the
479 limitations to this comparison: while our simulated ocean changes are for the annual mean (and

480 using wintertime atmospheric forcing changes), proxy records can be records of annual means,
481 or of a particular season. This should be kept in mind in reading this section.

482 6.1 Changes to currents in the East Australian-New Zealand sector

483 There is paleoproxy evidence for changes to the ocean circulation and climate over the East
484 Australian – New Zealand sector during the early deglaciation consistent with the impacts of a
485 Split Jet weakening. First, *Felis et al.* [2014] showed that SST over the southeastern coral sea
486 (around 22.6°S, 161.7°E) as well as SSTs over the Great Barrier Reef (17.1°S, 146.6°E and,
487 19.7°S, 150.3°E) warmed during this period. Second, *Schiraldi et al.* [2014] argue, based on
488 proxy records in the Bay of Plenty (north of North Island of New Zealand, around 37°S, 177°E),
489 that the water source changed there from an Antarctic source to a tropical one; and that a proxy
490 farther south in Hawke Bay (MD97-2121 40.38°S, 178°E; *Carter et al.*, 2008) also shows more
491 influence from tropical sources. Both proxies are consistent with the impact of a weaker Split Jet
492 but not with a more positive SAM (Figure 4b, c; the green triangles in those figures show the
493 location of the proxies). In particular, strengthening of the South Pacific subtropical gyre in the
494 weak Split Jet leads to a stronger Tasman front that crosses over to northern New Zealand and
495 feeds into the East Auckland and then the East Cape Current off the eastern coast of North
496 Island; this would explain the increased prevalence of tropically-sourced waters. On the other
497 hand, a positive SAM leads to a *weaker* East Auckland and East Cape current (Figure 4c).

498 Furthermore, there is robust proxy evidence that argues for a southward shift of the
499 Subtropical Front during the early deglaciation that implicates changes to the Southern
500 Hemisphere westerlies, though both a weaker Split Jet and positive phase of the SAM may
501 produce this change. *Bostock et al.* [2015] show from a transect of cores (see figure 4, green line
502 south of New Zealand, for the location) that the southern Subtropical Front by the Solander

503 Trough rapidly shifted southwards during the early deglaciation (18-16ka) with respect to its
504 modern position. *Sikes et al.* [2009] show that the Subtropical Front south of Tasmania (figure 4,
505 green line south of Tasmania) also shifted southwards during deglaciation. Warming of SST
506 south of Tasmania and New Zealand indicates a southward shift of the Subtropical Front, which
507 is consistent with the East Australian Current penetrating farther south. This could be the result
508 of either a weaker Split Jet or positive SAM, or both.

509 6.2 Changes to Drake Passage Throughflow and southeastern Pacific

510 There is evidence for a substantial increase to the Drake Passage throughflow during the early
511 deglacial period that argues more for a positive-SAM like change. *Lamy et al.* [2015] used grain
512 size distributions in cores along the Drake Passage to argue for a 40% decrease in Drake Passage
513 northernmost flow during glacial times compared to present, coupled with a stronger northward
514 current along Chile. During Heinrich 1, Drake Passage transport – mainly over the Subantarctic
515 region – reinvigorated, though the magnitude was not reported. In our simulations, both a
516 weaker Split Jet and a positive SAM lead to a stronger Drake Passage throughflow, though the
517 change is much stronger with a positive SAM. This observation might argue that a positive
518 SAM was a more likely scenario for the Heinrich 1 period.

519 We note that sediment proxies along the coast of Southern Chile have shown a strong
520 warming in SST in the southeastern Pacific off the coast of southern Chile timed to Heinrich 1
521 (in particular site ODP1233 at 41°S, 74.2°W) (*Lamy et al.*, 2007; *Mohtadi et al.*, 2008). *Lamy et*
522 *al.* (2015) partly attribute this warming to an increased Drake Passage throughflow, that then
523 lessened the northward advection of subantarctic waters northwards towards southern Chile. We
524 note that our simulations show no evidence of significant warming for a weak Split Jet or
525 positive SAM (figure 4b,c), and in fact for a weaker Split Jet shows a moderate cooling around

526 41°S in the southeastern Pacific. We speculate that the ocean model may not have sufficient
527 horizontal resolution to resolve key processes that determine SST and its changes at the
528 ODP1233 location. This includes resolving the sharp meridional SST front over that region, as
529 well as coastal upwelling and coastal current processes.

530 6.3 Changes to Mode/Intermediate Water formation

531 Proxy evidence generally indicates an increase in either Antarctic Intermediate Water (AAIW) or
532 the Subantarctic Mode Water (SAMW) formation during early deglaciation (figure 9 shows the
533 location of the proxies discussed). *Pahnke and Zahn* [2005] found evidence for increased AAIW
534 formation during H stadials as seen in the record by Chatham Rise to the east of New Zealand;
535 this is corroborated by a record in the middle depths of the Brazil Margin [*Pahnke et al.*, 2008]
536 that showed similar changes during H1 and the Younger Dryas. *Pena et al.* [2013] argue, from
537 an Nd proxy record in the southeastern equatorial Pacific, for a stronger intrusion of southern
538 intermediate waters into the equatorial undercurrent during North Atlantic cold periods,
539 including H1.

540 Comparison of these records with the ideal age changes in our simulations – with a
541 reduced ideal age interpreted to be increased intermediate water formation - is equivocal. The
542 *Pahnke and Zahn (2005)* record is consistent with both a weak Split Jet or positive SAM change;
543 the simulated ocean ideal age decreases in both instances at the Pahnke and Zahn (2005) site off
544 Chatham rise (figure 9a & b, left panels), though the weak Split Jet response appears to have a
545 stronger influence. However, neither the Split Jet nor the SAM modulations produce the
546 necessary changes seen in the Brazil Margin (figure 9a & b, center panels), though we point out
547 that changes to the Atlantic Meridional Overturning Circulation (AMOC) during H1 – and not
548 represented in our simulations – likely have a larger influence on intermediate waters in this

549 region. There is intriguing evidence that a weaker Split Jet can explain the *Pena et al. (2013)*
550 equatorial undercurrent record. A weaker Split Jet leads to increased rate of SAMW formation
551 over the entire tropical Pacific (cf Fig 8a), and waters at intermediate depths over the equatorial
552 region are younger for a weaker Split Jet (figure 9a, right panel); on the other hand, they are
553 older for a positive SAM (Figure 9b, right panel).

554 *Muratli et al. (2010)* inferred *decreasing* AAIW ventilation during the early deglacial
555 warming, based on decreasing oxygen inferred from three ODP cores (1233, 1234, and 1235)
556 along the Chilean coast and intermediate depths record. This observation appears to be contrary
557 to the increased intermediate water formation simulated from the weaker Split Jet or positive
558 SAM phase. However, examination of the simulated dissolved O₂ along 75.5°W (the gridpoint
559 closest to the coastline) actually shows *reduced* oxygen concentration for the weak Split Jet, and
560 to a lesser extent a positive SAM (figure 9c,d), at least for the two equatorward cores (1234 and
561 1235). The decreased oxygen appears not to be from intermediate waters from the south (which
562 shows *increased* oxygen levels, consistent with the increased ventilation), but rather from the
563 surface undercurrent and Pacific Deep Water originating from the north. We thus suggest that
564 our inference of increased intermediate formation from the weaker Split Jet or positive SAM
565 might not be inconsistent with the *Muratli et al. (2010)* result.

566 Overall, the evidence from intermediate water formation somewhat favors a weakened
567 Split Jet over the positive SAM during the early deglacial period, as it more readily explains the
568 *Pahnke and Zahn (2005)*, *Pena et al. (2013)*, and *Muratli et al. (2010)* observations.

569 6.4 Changes to the Meridional Overturning Circulation

570 It is by now well-established that increased ventilation of Antarctic circumpolar deep water
571 occurred during Heinrich 1 (e.g. *Anderson et al.*, 2009; *Skinner et al.*, 2010; *Siani et al.*, 2013).
572 Since both a weaker Split Jet and a positive SAM lead to increased ventilation of bottom waters,
573 this observation does not distinguish between the two wind forcings (though the magnitude is
574 stronger in the positive SAM case). *Anderson et al.* [2009] however found increased burial of
575 biogenic opal, interpreted as increased upwelling, at three core locations during H1: 53.2S 5.1E
576 (TN057-13-4PC); 59.62S 155.24E(E27-29); and 61.9S 170W (NBP9802-6PC). If we compare
577 the simulated upwelling changes at these three locations (Figure 5b,c - green triangles), there is
578 no compelling match with either forcing – indeed in two of the three core locations, the change
579 in upwelling is of one sign with the weaker SJ, and of the other in the positive SAM. Hence, the
580 *Anderson et al.* (2009) proxy record cannot be readily explained by either the weak Split Jet nor
581 positive SAM forcing. Note however that the spatial structure of upwelling changes has
582 relatively small-scale variations, making it difficult to do a convincing model-proxy comparison
583 given the inherent spatial biases in model simulations.

584 While there is controversy regarding direct ventilation of the North Pacific Ocean in the
585 subarctic during H1 [*Jaccard and Galbraith*, 2013; *Okazaki et al.*, 2010], we are not aware of
586 any paleoproxy evidence in support or against the idea of enhanced deep Pacific oxygenation
587 from the south (Figure 7). Recently published redox sensitive trace metal data indicate increased
588 ventilation and oxygenation of the deep Southern Ocean during northern stadials [*Jaccard et al.*,
589 2016], but these data are from the Atlantic sector.

590

591 7. Conclusion

592 We investigate ocean circulation and biogeochemical changes resulting from a modulation of the
593 South Pacific Split Jet, the wind changes being previously proposed by the authors as an
594 alternative hypothesis for Southern Hemisphere westerly wind changes in contrast to Southern
595 Annular Mode changes. We apply wind changes – derived from observed Southern Hemisphere
596 wintertime interannual variability – on an ocean general circulation model (POP2) with a simple
597 formulation for ocean biogeochemistry. The most notable features of a weakening of the Split
598 jet is a strengthening and southward displacement of the South Pacific subtropical gyre, that
599 leads to pronounced circulation and temperature changes over the East Australian/New Zealand
600 sector. A weaker Split Jet and a positive SAM result in some similar responses and some
601 distinctly different ones. Similar responses include (i) increased ventilation of bottom waters,
602 though the SAM appears significantly more effective; and (ii) a southward shift of the subpolar
603 front over the New Zealand sector. Distinct responses of a weaker Split Jet include enhanced
604 Subantarctic Mode Water formation in the South Pacific, and a stronger subtropical gyre that
605 leads to a stronger East Australian current that feeds into the Tasman Front and East
606 Auckland/East Cape current; as a result, SSTs are warmer east of northeastern Australia and
607 surrounding New Zealand. Distinct responses to a positive SAM are significantly increased
608 Drake Passage throughflow and ventilation of the Antarctic Intermediate Waters over all basins.

609 We thus identify key oceanic responses that could uniquely distinguish between the
610 modulation of the Split Jet and a modulation of the SAM. Our limited comparison of existing
611 published paleoceanographic records of the early deglacial period however finds evidence for
612 both a weaker Split Jet and positive SAM: notably, the observed SST warming off the Great
613 Barrier Reef [*Felis et al.*, 2014] and northern New Zealand [*Schiraldi et al.*, 2014] suggests a
614 weaker Split Jet, whereas the increased Drake Passage throughflow [*Lamy et al.*, 2015] favors a

615 positive SAM response. Intermediate water proxies suggest that both a weaker Split Jet and
616 positive SAM are possible, though more consistent with the former.

617 Note that the wind responses are not exclusive of each other: both can occur at the same
618 time, and in fact this is the interpretation we favor for the early deglacial period based on our
619 comparison. We have done an additional set of ocean simulations using forcing where we added
620 the weak Split Jet and positive SAM winds, and compared against a simulation adding a strong
621 Split Jet with negative SAM conditions; to first order, we find that the results appear to be a
622 linear combination of the two individual forcings, with no significant nonlinearities (figures not
623 shown). We note that in our previous modeling study that proposed an atmospheric
624 teleconnection from North Atlantic cooling to the southern hemisphere westerlies via a
625 southward ITCZ shift [Lee *et al.*, 2011], the southern hemisphere westerly response resembled a
626 mixture of a weaker Split Jet and positive SAM (see figure 5a of Lee *et al.* 2011).

627 The modulation of the Split Jet appears to be extremely effective at changing the
628 temperature climate over the New Zealand sector. *Chiang et al. (2014)* previously noted that a
629 weaker Split Jet is effective at warming New Zealand by allowing tropical airmasses to penetrate
630 further southwards. This study shows that the ocean response to a weaker Split Jet also warms
631 the ocean temperatures surrounding New Zealand (figure 4b). Previous work had documented
632 pronounced and rapid reduction to glacier length there during Heinrich 1 (Putnam *et al.*, 2013)
633 and Younger Dryas (Kaplan *et al.* 2010), and the opposite during the Antarctic cold reversal
634 (Putnam *et al.* 2010). We advance the Split Jet modulation as a plausible mechanism for
635 explaining the rapid climate changes occurring over New Zealand during deglaciation.

636 There have been many previous studies on the ocean response to changing southern
637 hemisphere westerlies. A distinct advantage of our study over most other such studies is that, by

638 deriving wind changes from observed interannual variability, we are using wind forcings that are
639 physically realizable, albeit large. The realism of our forcing and model setup allows us to make
640 specific regional predictions that can be contrasted against paleoproxy data.

641 There are clear limitations to our study. First, we are using simulations embedded within
642 a modern-day basic state, so the assumption is that the glacial basic state ocean circulation is
643 qualitatively similar. Second, our winds are applied only to the austral winter, so we omit the
644 influence of changes to the austral summer (though we could undertake a similar study for that
645 season). Third, the winds we apply are steady in time, and we do not allow changes to high-
646 frequency weather disturbances; in reality, the position of storm tracks will change as the mean
647 flow changes. Finally, we are using an ocean model that does not permit ocean eddies, which are
648 demonstrated to be important to understanding Southern Ocean circulation and in particular the
649 effect on the strength of the ACC and the Southern Ocean meridional overturning circulation
650 [*Gent*, 2016]; we are in effect relying on the model's parameterizations to provide estimates of
651 the influence of eddies on the ocean general circulation. However, we note that the ocean model
652 we use employs a variable coefficient in the Gent-McWilliams eddy parameterization, and *Gent*
653 *and Danabasoglu* [2011] showed that the response of this model to increased Southern
654 Hemisphere westerlies appear realistic as compared to eddy-permitting models.

655 Finally, we know that the early deglacial period was marked by a slowdown of the
656 AMOC, and there is a global ocean adjustment to this slowdown that is separate and distinct
657 from the changes in response to the southern hemisphere wind changes. It may well be that such
658 adjustments could also explain some of the paleoproxy changes discussed in section 6.
659 Contrasting the direct effect of AMOC changes on Southern Ocean circulation with those of the
660 southern hemisphere westerly changes will be the focus of a future study.

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672 index used to compare against the wind EOF2 timeseries in section 3 was obtained from the
673 NOAA CPC website:
674 http://www.cpc.ncep.noaa.gov/products/precip/CWlink/daily_ao_index/ao/ao.shtml.

675 References

- 676 Anderson, R. F., S. Ali, L. I. Bradtmiller, S. H. H. Nielsen, M. Q. Fleisher, B. E. Anderson, and
677 L. H. Burckle (2009), Wind-Driven Upwelling in the Southern Ocean and the Deglacial Rise in
678 Atmospheric CO₂, *Science*, 323(5920), 1443-1448, doi:Doi 10.1126/Science.1167441.
- 679 Bals-Elsholz, T. M., E. H. Atallah, L. F. Bosart, T. A. Wasula, M. J. Cempa, and A. R. Lupo
680 (2001), The wintertime Southern Hemisphere split jet: Structure, variability, and evolution,
681 *Journal of Climate*, 14(21), 4191-4215.
- 682 Bostock, H. C., B. W. Hayward, H. L. Neil, A. T. Sabaa, and G. H. Scott (2015), Changes in the
683 position of the Subtropical Front south of New Zealand since the last glacial period,
684 *Paleoceanography*, 30(7), 824-844.
- 685 Cai, W., G. Shi, T. Cowan, D. Bi, and J. Ribbe (2005), The response of the Southern Annular
686 Mode, the East Australian Current, and the southern mid-latitude ocean circulation to global
687 warming, *Geophysical Research Letters*, 32(23).
- 688 Carter, L., B. Manighetti, G. Ganssen, and L. Northcote (2008), Southwest Pacific modulation of
689 abrupt climate change during the Antarctic Cold Reversal–Younger Dryas, *Palaeogeography,*
690 *Palaeoclimatology, Palaeoecology*, 260(1), 284-298.

- 691 Chiang, J. C. H., and A. R. Friedman (2012), Extratropical Cooling, Interhemispheric Thermal
692 Gradients, and Tropical Climate Change, *Annual Review of Earth and Planetary Sciences*, 40(1),
693 383-412, doi:10.1146/annurev-earth-042711-105545.
- 694 Chiang, J. C. H., S.-Y. Lee, A. E. Putnam, and X. Wang (2014), South Pacific Split Jet, ITCZ
695 shifts, and atmospheric North–South linkages during abrupt climate changes of the last
696 glacial period, *Earth and Planetary Science Letters*, 406(0), 233-246,
697 doi:<http://dx.doi.org/10.1016/j.epsl.2014.09.012>.
- 698 Chiang, J. C. H., K. Tokos, S.-Y. Lee, and K. Matsumoto (2017), Forcing files and Model output
699 for "Contrasting impacts of the South Pacific Split Jet and the Southern Annular Mode
700 modulation on Southern Ocean circulation and biogeochemistry", v2, UC Berkeley Dash,
701 Dataset, <https://doi.org/10.6078/D1PM3T>
- 702 d'Orgeville, M., nbsp, W. Sijp, M. England, and K. Meissner (2010), On the control of glacial-
703 interglacial atmospheric CO2 variations by the Southern Hemisphere westerlies, *Geophysical*
704 *Research Letters*, 37(21).
- 705 Danabasoglu, G., S. C. Bates, B. P. Briegleb, S. R. Jayne, M. Jochum, W. G. Large, S. Peacock,
706 and S. G. Yeager (2011), The CCSM4 Ocean Component, *Journal of Climate*, 25(5), 1361-
707 1389, doi:10.1175/JCLI-D-11-00091.1.
- 708 Denton, G. H., Anderson, R. F., Toggweiler, J. R., Edwards, R. L., Schaefer, J. M., & Putnam,
709 A. E. (2010). The last glacial termination. *Science*, 328(5986), 1652-1656.
- 710 Felis, T., H. V. McGregor, B. K. Linsley, A. W. Tudhope, M. K. Gagan, A. Suzuki, M. Inoue, A.
711 L. Thomas, T. M. Esat, and W. G. Thompson (2014), Intensification of the meridional
712 temperature gradient in the Great Barrier Reef following the Last Glacial Maximum, *Nature*
713 *communications*, 5.
- 714 Fletcher, M. S., and P. I. Moreno (2011), Zonally symmetric changes in the strength and position
715 of the Southern Westerlies drove atmospheric CO2 variations over the past 14 k.y., *Geology*,
716 39(5), 419-422, doi:Doi 10.1130/G31807.1.
- 717 Gent, P. R. (2016), Effects of Southern Hemisphere wind changes on the meridional overturning
718 circulation in ocean models, *Annual review of marine science*, 8, 79-94.
- 719 Gent, P. R., and G. Danabasoglu (2011), Response to increasing Southern Hemisphere winds in
720 CCSM4, *Journal of climate*, 24(19), 4992-4998.

- 721 Getzlaff, J., H. Dietze, and A. Oschlies (2016), Simulated effects of southern hemispheric wind
722 changes on the Pacific oxygen minimum zone, *Geophysical Research Letters*, 43(2), 728-734.
- 723 Inatsu, M., and B. J. Hoskins (2004), The zonal asymmetry of the Southern Hemisphere winter
724 storm track, *Journal of Climate*, 17(24), 4882-4892.
- 725 Jaccard, S., and E. Galbraith (2013), Direct ventilation of the North Pacific did not reach the
726 deep ocean during the last deglaciation, *Geophysical Research Letters*, 40(1), 199-203.
- 727 Jaccard, S. L., E. D. Galbraith, A. Martinez-Garcia, and R. F. Anderson (2016), Covariation of
728 deep Southern Ocean oxygenation and atmospheric CO² through the last ice age, *Nature*,
729 530(7589), 207.
- 730 Kaplan, Michael R., Joerg M. Schaefer, George H. Denton, David JA Barrell, Trevor JH Chinn,
731 Aaron E. Putnam, Bjørn G. Andersen, Robert C. Finkel, Roseanne Schwartz, and Alice M.
732 Doughty. "Glacier retreat in New Zealand during the Younger Dryas stadial." *Nature* 467, no.
733 7312 (2010): 194-197.
- 734 Knox, F., and M. B. McElroy (1984), Changes in atmospheric CO₂: Influence of the marine
735 biota at high latitude, *Journal of Geophysical Research: Atmospheres*, 89(D3), 4629-4637.
- 736 Lamy, F., H. W. Arz, R. Kilian, C. B. Lange, L. Lembke-Jene, M. Wengler, J. Kaiser, O. Baeza-
737 Urrea, I. R. Hall, and N. Harada (2015), Glacial reduction and millennial-scale variations in
738 Drake Passage throughflow, *Proceedings of the National Academy of Sciences*, 112(44), 13496-
739 13501.
- 740 Lamy, F., J. Kaiser, H. W. Arz, D. Hebbeln, U. Ninnemann, O. Timm, A. Timmermann, and J.
741 R. Toggweiler (2007), Modulation of the bipolar seesaw in the southeast Pacific during
742 Termination 1, *Earth and Planetary Science Letters*, 259(3-4), 400-413, doi:Doi
743 10.1016/J.Epsl.2007.04.040.
- 744 Large, W., and S. Yeager (2009), The global climatology of an interannually varying air-sea flux
745 data set, *Climate Dynamics*, 33(2-3), 341-364.
- 746 Lee, S. Y., J. C. H. Chiang, K. Matsumoto, and K. S. Tokos (2011), Southern Ocean wind
747 response to North Atlantic cooling and the rise in atmospheric CO₂: Modeling perspective and
748 paleoceanographic implications, *Paleoceanography*, 26, doi:Artn Pa1214
749 Doi 10.1029/2010pa002004.

- 750 Lovenduski, N. S., N. Gruber, S. C. Doney, and I. D. Lima (2007), Enhanced CO₂ outgassing in
751 the Southern Ocean from a positive phase of the Southern Annular Mode, *Global Biogeochem*
752 *Cy*, 21(2).
- 753 Marshall, G. J. (2003), Trends in the Southern Annular Mode from observations and reanalyses,
754 *Journal of Climate*, 16(24), 4134-4143.
- 755 Mo, K. C. (2000), Relationships between low-frequency variability in the Southern Hemisphere
756 and sea surface temperature anomalies, *Journal of Climate*, 13(20), 3599-3610.
- 757 Mohtadi, M., P. Rossel, C. B. Lange, S. Pantoja, P. Böning, D. J. Repeta, M. Grunwald, F.
758 Lamy, D. Hebbeln, and H.-J. Brumsack (2008), Deglacial pattern of circulation and marine
759 productivity in the upwelling region off central-south Chile, *Earth and Planetary Science*
760 *Letters*, 272(1), 221-230.
- 761 Moreno, P. I., J. P. Francois, C. M. Moy, and R. Villa-Martinez (2010), Covariability of the
762 Southern Westerlies and atmospheric CO₂ during the Holocene, *Geology*, 38(8), 727-730,
763 doi:Doi 10.1130/G30962.1.
- 764 Muratli, J. M., Chase, Z., Mix, A. C., & McManus, J. (2010). Increased glacial-age ventilation of
765 the Chilean margin by Antarctic Intermediate Water. *Nature Geoscience*, 3(1), 23-26.
- 766 Najjar, R. G., X. Jin, F. Louanchi, O. Aumont, K. Caldeira, S. C. Doney, J. C. Dutay, M.
767 Follows, N. Gruber, and F. Joos (2007), Impact of circulation on export production, dissolved
768 organic matter, and dissolved oxygen in the ocean: Results from Phase II of the Ocean Carbon-
769 cycle Model Intercomparison Project (OCMIP-2), *Global Biogeochem Cy*, 21(3).
- 770 Nakamura, H., and A. Shimpo (2004), Seasonal variations in the Southern Hemisphere storm
771 tracks and jet streams as revealed in a reanalysis dataset, *Journal of Climate*, 17(9), 1828-1844.
- 772 North, G. R., T. L. Bell, R. F. Cahalan, and F. J. Moeng (1982), Sampling errors in the
773 estimation of empirical orthogonal functions, *MWR*, 110, 699-706.
- 774 Okazaki, Y., A. Timmermann, L. Menviel, N. Harada, A. Abe-Ouchi, M. O. Chikamoto, A.
775 Mouchet, and H. Asahi (2010), Deepwater Formation in the North Pacific During the Last
776 Glacial Termination, *Science*, 329(5988), 200-204, doi:Doi 10.1126/Science.1190612.
- 777 Pahnke, K., S. L. Goldstein, and S. R. Hemming (2008), Abrupt changes in Antarctic
778 Intermediate Water circulation over the past 25,000 years, *Nature Geoscience*, 1(12), 870-874.
- 779 Pahnke, K., and R. Zahn (2005), Southern Hemisphere water mass conversion linked with North
780 Atlantic climate variability, *Science*, 307(5716), 1741-1746.

781 Pena, L., S. L. Goldstein, S. R. Hemming, K. M. Jones, E. Calvo, C. Pelejero, and I. Cacho (2013),
782 Rapid changes in meridional advection of Southern Ocean intermediate waters to the
783 tropical Pacific during the last 30kyr, *Earth and Planetary Science Letters*, 368, 20-32.

784 Preisendorfer, R. W. (1988), *Principal Component Analysis in Meteorology*, 425 pp., Elsevier,
785 Amsterdam.

786 Putnam, Aaron E., George H. Denton, Joerg M. Schaefer, David JA Barrell, Bjørn G. Andersen,
787 Robert C. Finkel, Roseanne Schwartz, Alice M. Doughty, Michael R. Kaplan, and Christian
788 Schlüchter. "Glacier advance in southern middle-latitudes during the Antarctic Cold
789 Reversal." *Nature Geoscience* 3, no. 10 (2010): 700-704.

790 Putnam, Aaron E., Joerg M. Schaefer, George H. Denton, David JA Barrell, Bjørn G. Andersen,
791 Tobias NB Koffman, Ann V. Rowan et al. "Warming and glacier recession in the Rakaia valley,
792 Southern Alps of New Zealand, during Heinrich Stadial 1." *Earth and Planetary Science
793 Letters* 382 (2013): 98-110.

794 Russell, J. L., K. W. Dixon, A. Gnanadesikan, R. J. Stouffer, and J. Toggweiler (2006), The
795 Southern Hemisphere westerlies in a warming world: Propping open the door to the deep ocean,
796 *Journal of Climate*, 19(24), 6382-6390.

797 Sarmiento, J., and J. Toggweiler (1984), A new model for the role of the oceans in determining
798 atmospheric pCO₂, *Nature*, 308(5960), 621-624.

799 Schiraldi, B., E. L. Sikes, A. C. Elmore, M. S. Cook, and K. A. Rose (2014), Southwest Pacific
800 subtropics responded to last deglacial warming with changes in shallow water sources,
801 *Paleoceanography*, 29(6), 595-611.

802 Siani, G., E. Michel, R. De Pol-Holz, T. DeVries, F. Lamy, M. Carel, G. Isguder, F. Dewilde,
803 and A. Lourantou (2013), Carbon isotope records reveal precise timing of enhanced Southern
804 Ocean upwelling during the last deglaciation, *Nature Communications*, 4.

805 Siegenthaler, U., and T. Wenk (1984), Rapid atmospheric CO₂ variations and ocean circulation,
806 *Nature*, 308(5960), 624-626.

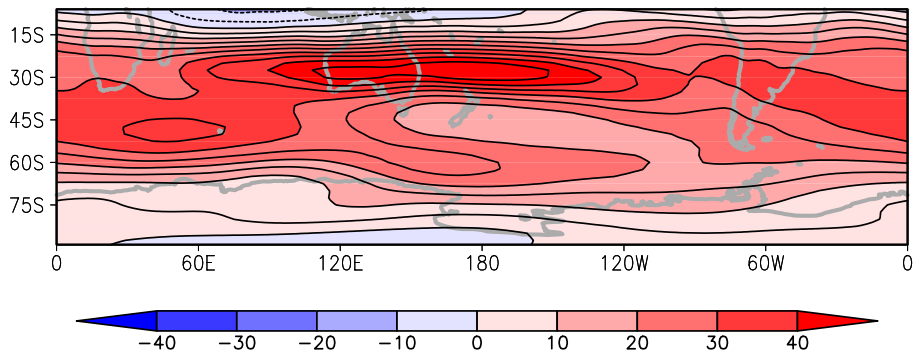
807 Sikes, E., W. Howard, C. Samson, T. Mahan, L. Robertson, and J. Volkman (2009), Southern
808 Ocean seasonal temperature and Subtropical Front movement on the South Tasman Rise in the
809 late Quaternary, *Paleoceanography*, 24(2).

810 Skinner, L., S. Fallon, C. Waelbroeck, E. Michel, and S. Barker (2010), Ventilation of the deep
811 Southern Ocean and deglacial CO₂ rise, *Science*, 328(5982), 1147-1151.

- 812 Taljaard, J. (1972), Synoptic meteorology of the Southern Hemisphere, in *Meteorology of the*
813 *Southern Hemisphere*, edited, pp. 139-213, Springer.
- 814 Thompson, D. W. J., and J. M. Wallace (2000), Annular modes in the extratropical circulation.
815 Part I: Month-to-month variability, *Journal of Climate*, 13(5), 1000-1016.
- 816 Toggweiler, J. R. (2009), Shifting Westerlies, *Science*, 323(5920), 1434-1435,
817 doi:10.1126/science.1169823.
- 818 Toggweiler, J. R., and B. Samuels (1998), On the ocean's large-scale circulation near the limit of
819 no vertical mixing, *Journal of Physical Oceanography*, 28(9), 1832-1852.
- 820

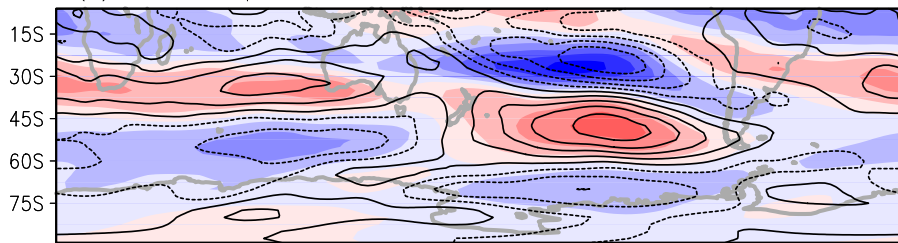
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(a) U 250mb JJAS climatology

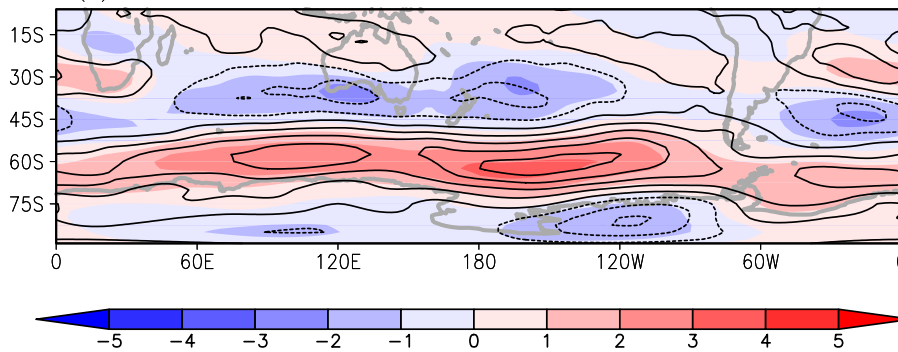


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(b) EOF1 – Split Jet



(c) EOF2 – Southern Annular Mode



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825 **Figure 1.** (a) June–September (JJA) 250 mb zonal wind climatology averaged over 1979–2009,

826 showing the presence of the Split Jet in the South Pacific. Contour interval is 5 m/s. (b and c)

827 Mode 1 and mode 2 (respectively) of the 3-D EOF of June-September mean zonal winds from

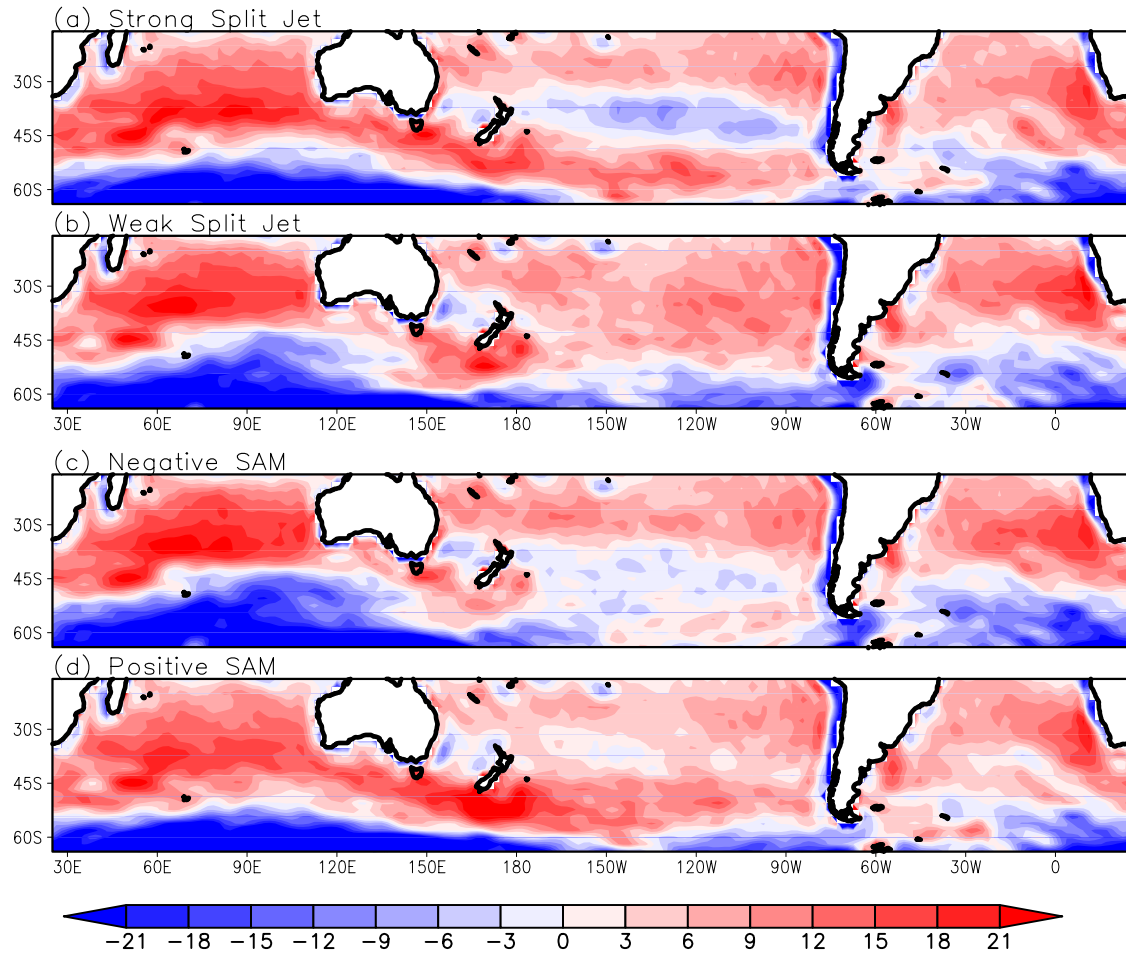
828 100-1000mb and 15°S-80°S. The spatial patterns are shown as regression of the normalized

829 principal component onto the June-Sep 250mb zonal winds (shaded) and June-Sep 700mb zonal

830 winds (contours). Units are m/s, and contour interval is 0.5m/s

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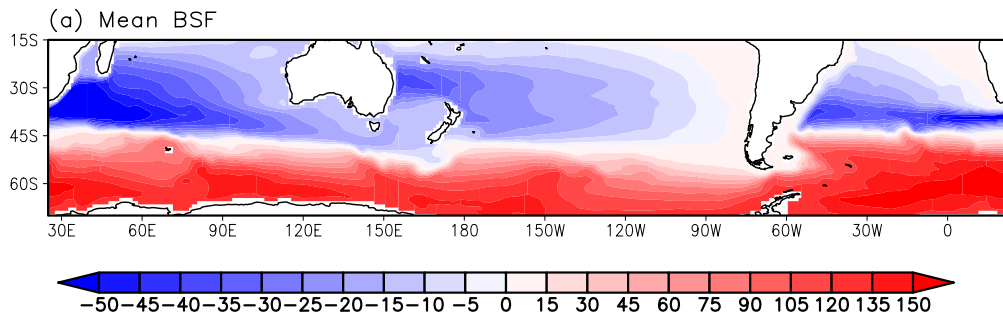
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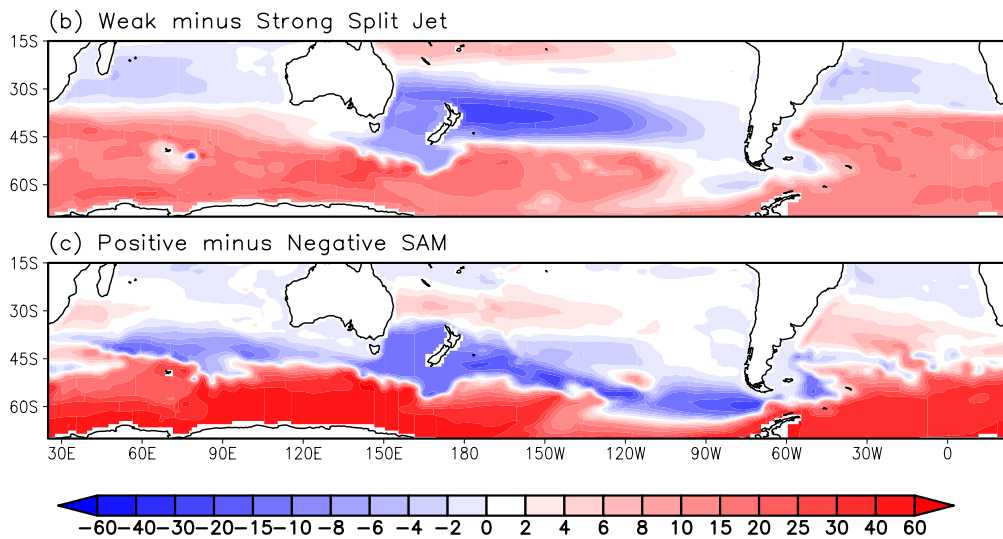
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Figure 2. June-September averaged wind stress curl imposed on the ocean model for (a) Strong Split Jet, (b) Weak Split Jet, (c) Negative SAM, and (d) Positive SAM. The fields correspond to a 2σ change to the interannual variation associated with each EOF mode, in both the positive and negative directions (i.e. they are sizable variations, but not unrealistically so). Units are $\times 10^{-8}$ Pa/m

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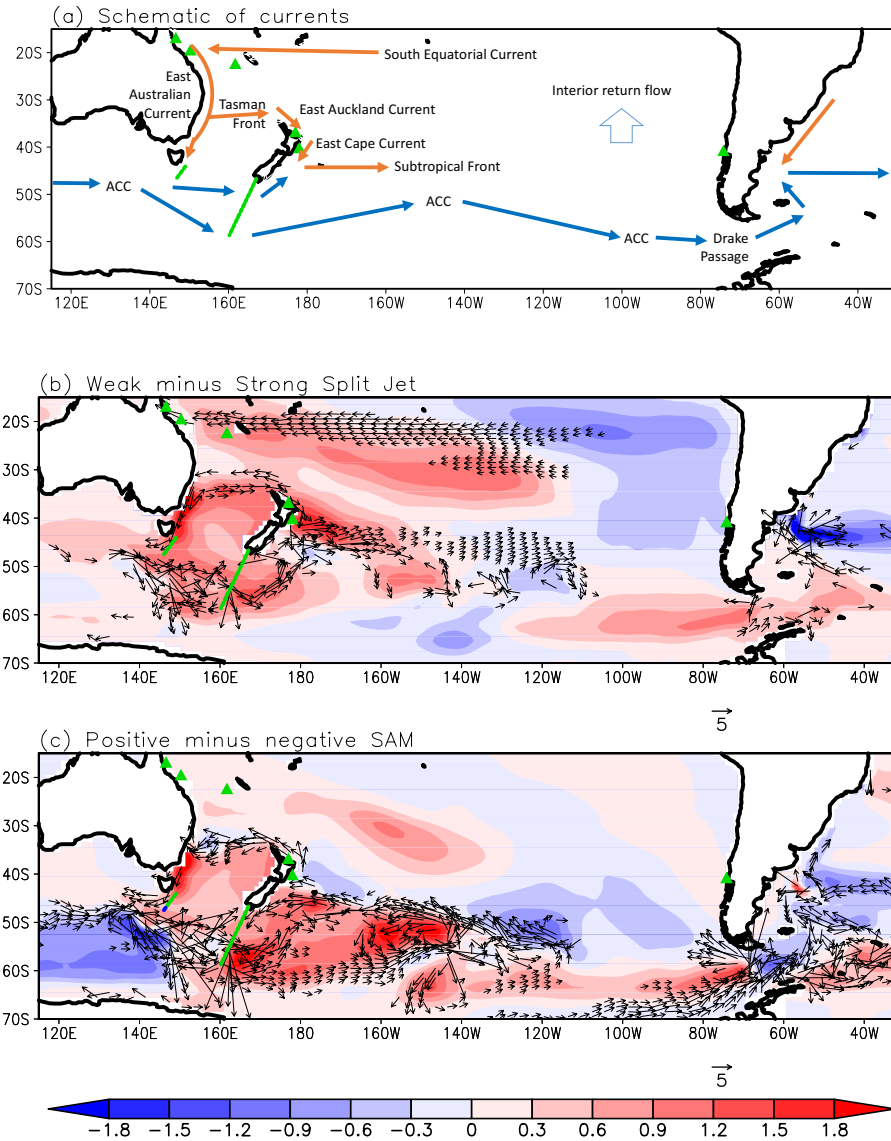


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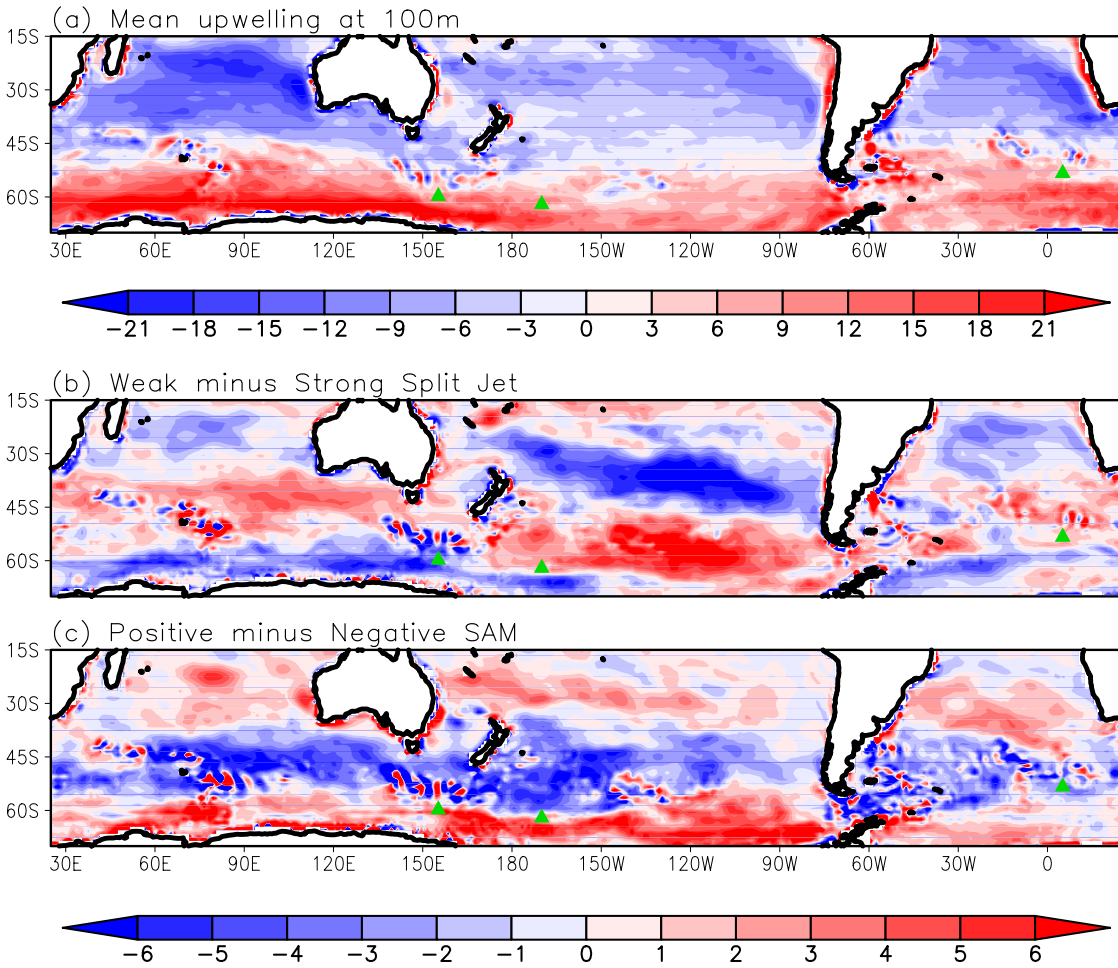
845 **Figure 3.** Barotropic streamfunction in the model simulation. (a) Annual mean streamfunction
 846 (negative values imply counterclockwise flow). (b) Difference between the weak and strong Split
 847 Jet cases. (c) Difference between the positive and negative SAM cases. Units are in Sv. Note
 848 the nonlinear scale – in particular for (a), positive intervals are 3x larger than negative intervals.
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854 **Figure 4.** (a) Schematic of ocean currents and geographic features discussed in the text. Warm
 855 currents are represented in orange, and cold currents in blue. We only include those features
 856 highlighted in our analyses. ACC = Antarctic Circumpolar Current. (b and c) Sea surface
 857 temperature anomaly (color) and surface current anomalies (vectors) for (b) Weak minus Strong
 858 Split Jet; and (c) Positive minus Negative SAM. Temperature is in K, and the reference vector is
 859 5 cm/s. Current anomalies below 2 cm/s are not shown. Anomalies are averaged over the upper
 860 100m. Green triangles denote locations of sediment core records for the following (from north to
 861 south): Great Barrier reef SST records at 17.1°S, 146.6°E and 19.7°S, 150.3°E (Felis et al. 2014);
 862 Southeastern coral sea 22.6°S, 161.7°E; Bay of Plenty 37°S, 177°E (Schiraldi Jr et al. 2014);
 863 Hawke Bay 40.38°S, 178°E (Carter et al. 2008); and ODP-1233 41°S, 74.2°W (Lamy et al.
 864 2007). The green lines denote the approximate locations of the surface ocean transect
 865 investigated in Bostok et al. (2015) south of New Zealand, and in Sikes et al. (2009) south of
 866 Tasmania.



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870 **Figure 5.** Upwelling at 100m (negative values are downwelling). (a) Annual mean upwelling;

871 (b) Weak minus Strong Split Jet; and (c) Positive minus negative SAM. Units are $\times 10^5$ cm/s.

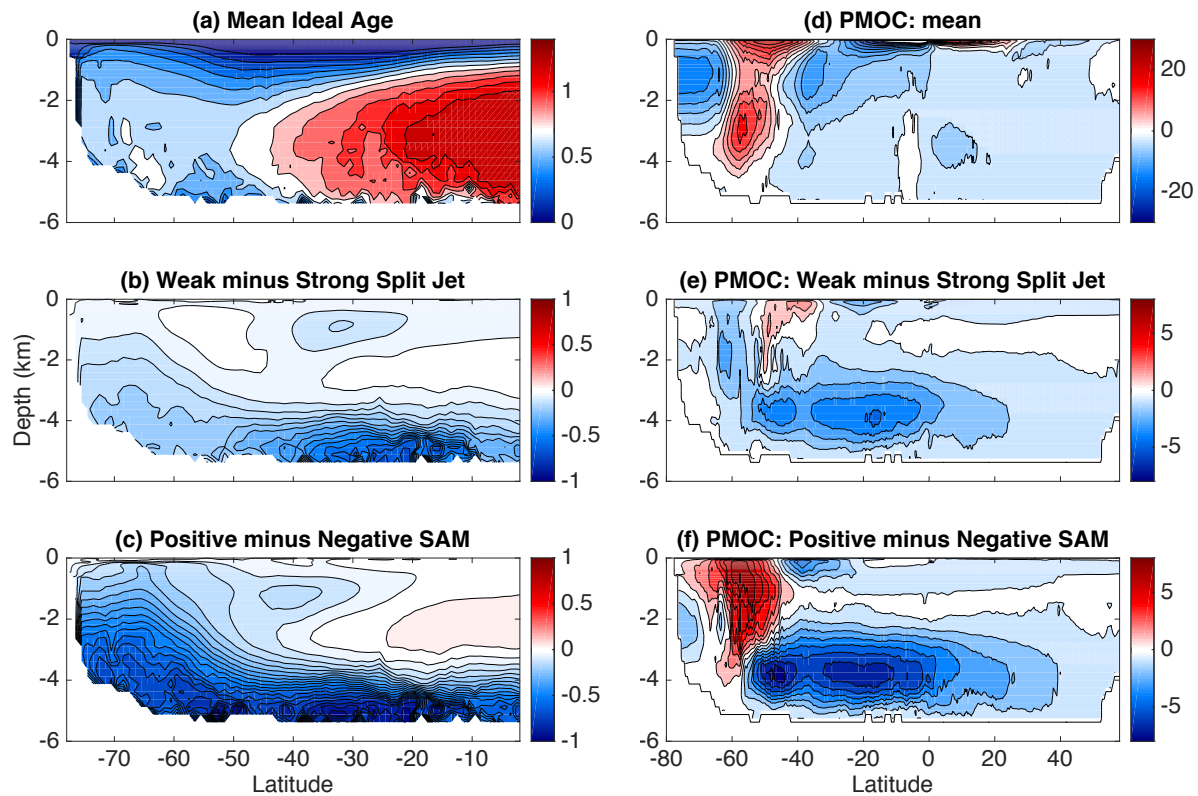
872 Green triangles are the locations to the three core locations (left to right - E27-23: 59.62°S

873 155.24°E; NBP9802-6PC: 61.9°S 170°W; and TN057-13PC: 53.2°S 5.1°E) in Anderson et al.

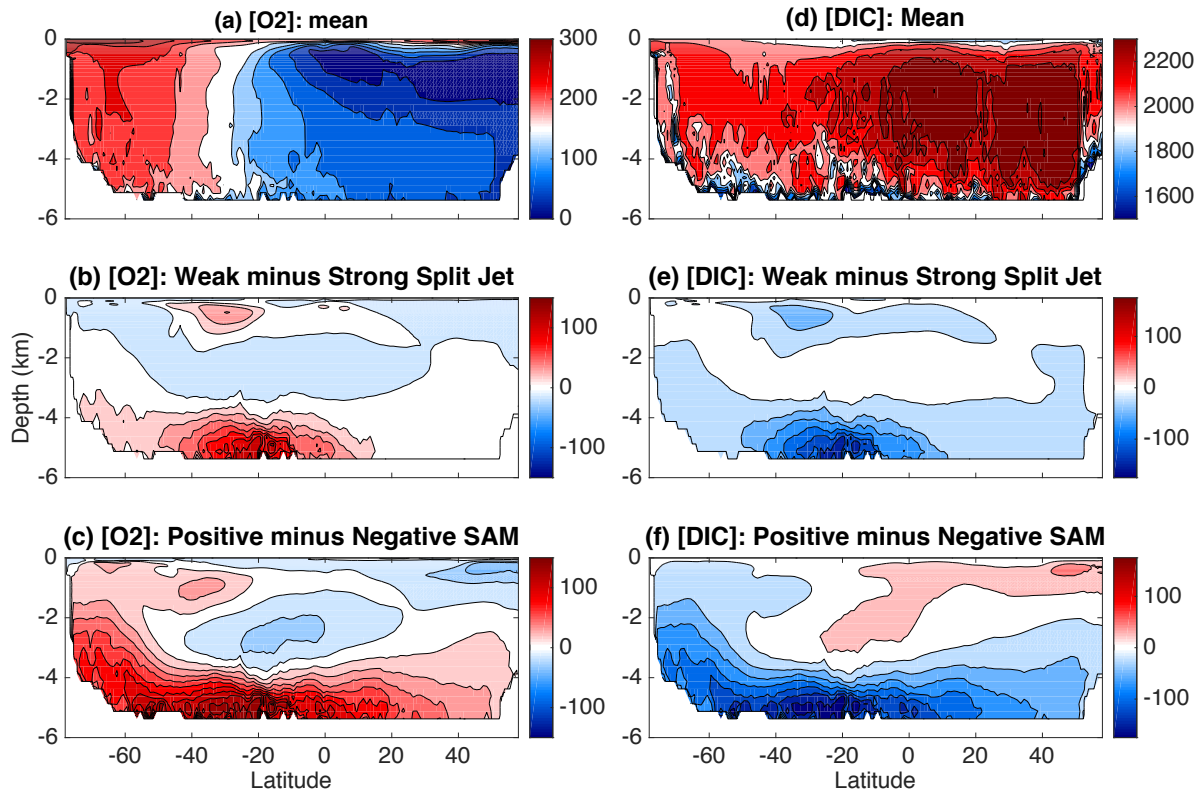
874 (2009).

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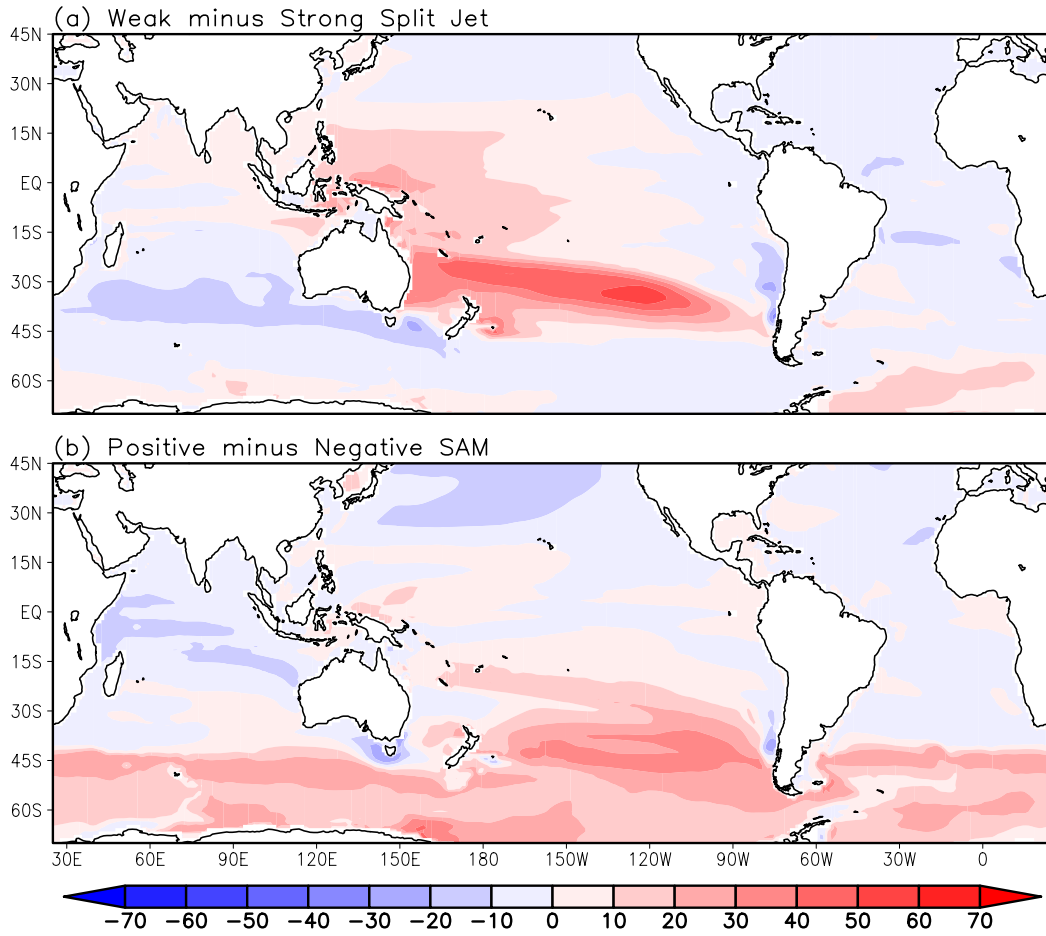


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 878 **Figure 6.** (a-c) Change to the ideal age zonally averaged over the South Pacific Sector (150°E-
 879 280°E). The ideal age is the amount of time elapsed since a water parcel left the surface. (a)
 880 Mean Ideal Age; (b) Weak minus Strong Split Jet; and (c) Positive minus Negative SAM. Units
 881 are kiloyears. Figure shows reduced ideal ages in the ocean interior for both cases indicating
 882 increased ventilation, but more pronounced for the SAM case. For a weak Split Jet, ideal age
 883 suggests increased ventilation of Subantarctic Mode Water (~30°S and 800m) and Antarctic
 884 Bottom Water. For a positive SAM, ideal age suggests increased ventilation of Antarctic
 885 Intermediate Water (~40°S and 1200m) and Antarctic Bottom Water. (d-f) Pacific MOC and its
 886 change. (d) Mean circulation; (e) Weak minus Strong Split Jet; and (f) Positive minus Negative
 887 SAM. Units are Sv ($10^6 \text{ m}^3 \text{ s}^{-1}$) and the contour interval is 1Sv. Positive values represent
 888 clockwise circulation.
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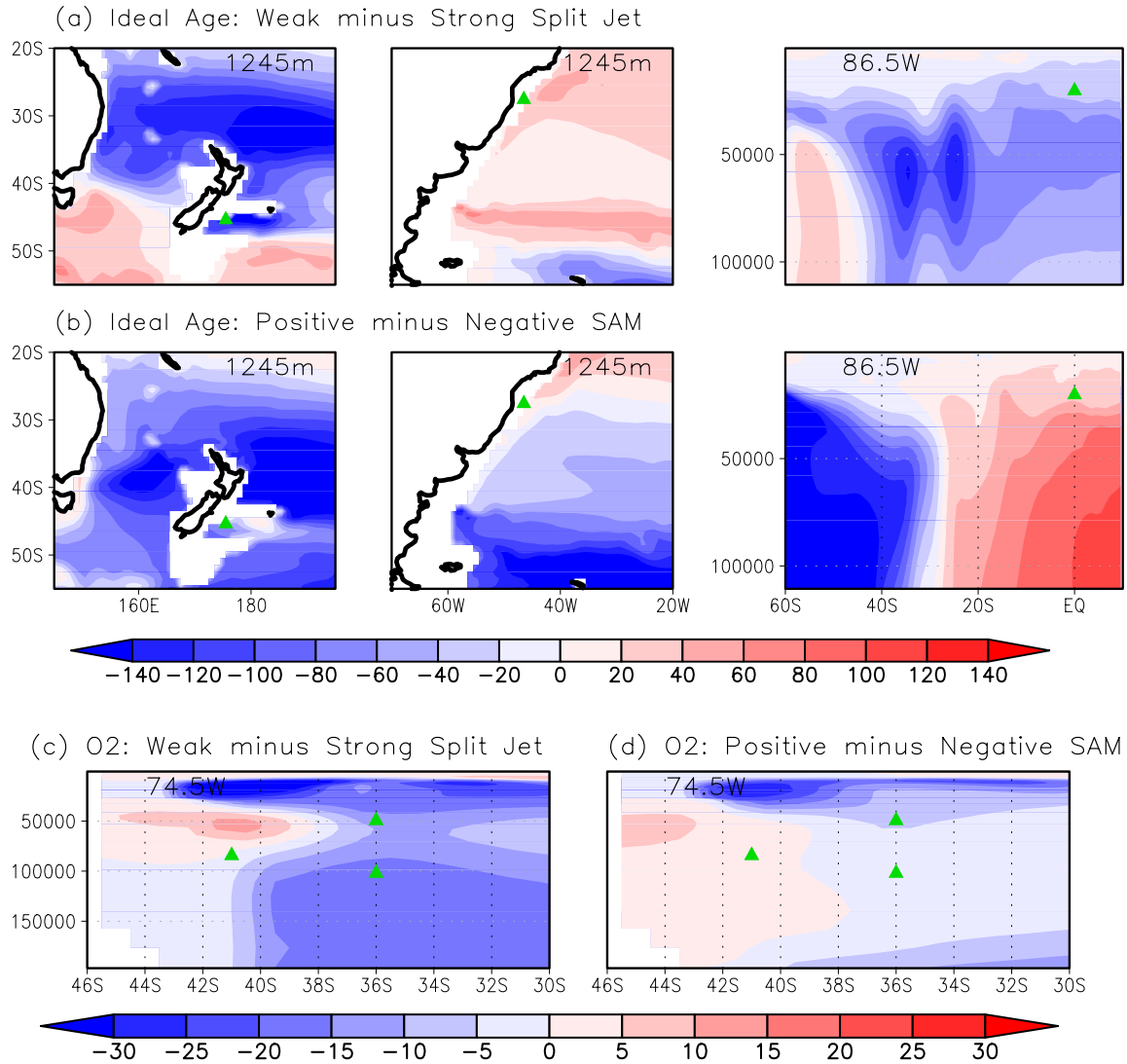
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Figure 7. (a-c) Mean field and change in the zonally averaged (150°E - 280°E) Pacific dissolved oxygen concentration. (a) Mean (contour interval $30 \mu\text{mol/kg}$); (b) Weak minus Strong Split Jet (contour interval $15 \mu\text{mol/kg}$); and (c) Positive minus Negative SAM (contour interval $15 \mu\text{mol/kg}$). (d-f). Mean and change in Pacific dissolved inorganic carbon averaged over 150°E - 280°E . (d) Mean (contour interval $100 \mu\text{mol/kg}$); (e) Weak minus Strong Split Jet (contour interval $25 \mu\text{mol/kg}$); and (f) Positive minus Negative SAM (contour interval $25 \mu\text{mol/kg}$).



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Figure 8. Change to the dissolved O₂ averaged over the upper 1000 m (a) Weak minus Strong Split Jet; and (b) Positive minus Negative SAM. Units are μmol/kg.



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 904 **Figure 9.** (a-b) Simulated difference in the ideal age of ocean water over the New Zealand
 905 sector (left panels), western south Atlantic (middle panels), and at the southeastern Pacific cross-
 906 section at 86.5°W. (a) Weak minus Strong Split Jet. (b) Positive minus Negative SAM. Units
 907 are years, and we interpret negative values to mean younger waters and hence more ventilation.
 908 For the New Zealand and western south Atlantic sectors, the depth of the water is as indicated in
 909 the figure, chosen as the level closest to the proxy record that it is being compared to. For the
 910 cross-section, the y-axis depth is in cm. Location of proxy records are as shown in the green
 911 triangles (MD27-2120 45.3°S, 175.5°E off New Zealand [Pahnke and Zahn 2005]; KNR159-5-
 912 36GGC 27.5°S, 46.5°W off the Brazil Margin [Pahnke et al. 2008]; and ODP 1240 0°N 86.3°W
 913 [Pena et al. 2013]). (c,d) Cross section of dissolved O₂ difference at 74.5°W off the coast of
 914 Chile; (c) is for weak minus strong Split Jet, and (d) is for positive minus negative SAM. Units
 915 are μmol/kg. The green triangles indicate the positions of ODP cores 1233 (41°S, 74°W, 838
 916 m), 1234 (36°S, 73°W, 1,015 m), and 1235 (36°S, 73°W, 489 m) reported in Muratli et al.
 917 (2010). Y-axis units of depth are in cm.