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Measurement of the $D^*(2010)^+$-D^+ Mass Difference

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0031-9007/17/119(20)/202003(7)

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We measure the mass difference, Δm_+, between the $D^*(2010)^+$ and the D^+ using the decay chain $D^*(2010)^+ \to D^+\pi^0$ with $D^+ \to K^-\pi^+\pi^+$. The data were recorded with the *BABAR* detector at center-of-mass energies at and near the $\Upsilon(4S)$ resonance, and correspond to an integrated luminosity of approximately 468 fb^{-1}. We measure $\Delta m_+ = (140\,601.0 \pm 6.8[\mathrm{stat}] \pm 12.9[\mathrm{syst}])$ keV. We combine this result with a previous *BABAR* measurement of $\Delta m_0 \equiv m(D^*(2010)^+) - m(D^0)$ to obtain $\Delta m_D = m(D^+) - m(D^0) = (4824.9 \pm 6.8[\mathrm{stat}] \pm 12.9[\mathrm{syst}])$ keV. These results are compatible with and approximately five times more precise than the Particle Data Group averages.

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The difference between the masses of the D^0 and D^+ mesons [1], $\Delta m_D \equiv m(D^+) - m(D^0)$, is a key ingredient constraining calculations of symmetry breaking due to differing u and d quark masses and electromagnetic interactions in the frameworks of chiral perturbation theory [2] and lattice QCD [3]. Its value is reported by the Particle Data Group (PDG) [4] to be $\Delta m_D = (4.77 \pm 0.08)$ MeV. The most precise direct measurement, reported by the LHCb Collaboration, is $\Delta m_D = (4.76 \pm 0.12 \pm 0.07)$ MeV [5]. This was found by comparing the invariant mass distributions of $D^0 \to K^-K^+\pi^-\pi^+$ and $D^+ \to K^-K^+\pi^+$ decays. A more powerful constraint comes from the difference of measured $D^{*+} \to D^+\pi^0$ and $D^{*+} \to D^0\pi^+$ mass difference distributions. CLEO has previously reported $\Delta m_+ \equiv m(D^*(2010)^+) - m(D^+) = (140.64 \pm 0.08 \pm 0.06)$ MeV using the decay chain $D^{*+} \to D^+\pi^0$ with $D^+ \to K^-\pi^+\pi^+$ [6]. In the present Letter, we report a new measurement of Δm_+ and combine it with our previously measured $D^{*+} \to D^0\pi^+$ mass difference [7,8], $\Delta m_0 \equiv m(D^*(2010)^+) - m(D^0)$, using two decay modes $D^0 \to K^-\pi^+$ and $D^0 \to K^-\pi^+\pi^-\pi^+$, to determine $\Delta m_D \equiv \Delta m_0 - \Delta m_+$ with very high precision.

This analysis is based on a data set corresponding to an integrated luminosity of approximately 468 fb^{-1} recorded at, and 40 MeV below, the $\Upsilon(4S)$ resonance [9]. The data were collected with the *BABAR* detector at the PEP-II2 asymmetric energy e^+e^- collider, located at the SLAC National Accelerator Laboratory. The *BABAR* detector is described in detail elsewhere [10,11]. The momenta of charged particles are measured with a combination of a cylindrical drift chamber (DCH) and a five-layer silicon vertex tracker (SVT), both operating within the 1.5 T magnetic field of a superconducting solenoid. Information from a ring-imaging Cherenkov detector is combined with specific ionization (dE/dx) measurements from the SVT and DCH to identify charged kaon and pion candidates. Electrons are identified, and photons from π^0 decays are measured, with a CsI(Tl) electromagnetic calorimeter (EMC). The return yoke of the superconducting coil is instrumented with tracking chambers for the identification of muons.

We study the $D^{*+} \to D^+\pi^0$ transition, using the $D^+ \to K^-\pi^+\pi^+$ decay mode, to determine the difference between the D^{*+} and D^+ masses Δm_+. To extract Δm_+, we fit the distribution of the difference between the reconstructed D^{*+} and D^+ masses, Δm. The signal component in the Δm fit is a resolution function determined from our Monte Carlo (MC) simulation of the detector response, while the contaminations from the background are accounted for by a threshold function.

We suppress combinatorial backgrounds, and backgrounds with D^{*+} candidates from B decays, by requiring D^{*+} mesons produced in $e^+e^- \to c\bar{c}$ reactions to have momenta in the e^+e^- center-of-mass frame greater than 3.0 GeV. Decays $D^{*+} \to D^0\pi^+$ with $D^0 \to K^-\pi^+\pi^0$ create backgrounds when the π^+ daughter of the $D^{*+} \to D^0\pi^+$ decay replaces the π^0 in the D^0 decay by mistake and the two have similar momenta. To mitigate this problem, events are rejected if $m(K^-\pi^+\pi^+\pi^0) - m(K^-\pi^+\pi^0) < 160$ MeV for either of the two π^+. The value of 160 MeV is chosen to be very conservative in terms of removing $D^{*+} \to D^0\pi^+$ decays [7,8] and causes almost no loss of signal. The decay chain is fitted subject to geometric constraints at the D^{*+} production vertex and the D^+ decay vertex, and to a kinematic constraint that the D^+ laboratory momentum points back to the luminous region whose horizontal, vertical, and longitudinal rms dimensions are about 6, 9, and 120 μm, respectively [10]. The χ^2 p value from the fit is required to be greater than 0.1%.

The "slow pion" from the D^{*+} decay, denoted as π_s^0, has a typical laboratory momentum of 300 MeV. All photons from π_s^0 decays have energies below 500 MeV. Their energy resolution is $\sigma_E/E \sim 7\%$, and angular resolutions are σ_θ and $\sigma_\phi \sim 10$ mr where the resolutions are measured with large uncertainties. In the $\pi_s^0 \to \gamma\gamma$ reconstruction, we first require both photon energies to be above 60 MeV, the total energy to be greater than 200 MeV, and the diphoton invariant mass to be between 120 and 150 MeV (approximately $\pm 2.5\sigma$ around the nominal π^0 mass [4]). After the selection, each photon pair is kinematically fitted to the hypothesis of a π^0 originating from the event primary vertex, and with the diphoton mass constrained to the nominal π^0 mass. This greatly improves the reconstructed π^0 momentum resolution and, therefore, the Δm resolution. The π^0 relative momentum resolution after the kinematic fit is $\sigma_p/p \sim 3\%$; this is still considerably worse than the approximately 0.5% D^+ relative momentum resolution.

Our MC simulation attempts to track run-by-run variations in detector response. The standard MC energy calibration method that accounts for energy loss in the EMC differs from that used with real data. This results in a reconstructed π^0 mass ($m_{\gamma\gamma}$) peak in MC events that peaks about 0.5 MeV below the nominal mass for low energy π^0s. In contrast, the $m_{\gamma\gamma}$ peak value from the calibrated data events generally coincides with the nominal value. Therefore, we approximate the neutral energy correction algorithm used in data by rescaling the reconstructed photon energies in MC events by factors depending on photon energy and data-taking periods [11]. While this improves the data-MC agreement, the reconstructed π^0 momentum in MC events remains slightly biased when compared with its generated value. To account for this bias, we also rescale the π^0 momentum in each MC event by approximately 0.2%, depending on the diphoton opening angle. In addition to improving the data-MC agreement in peak positions and shapes of the background-subtracted $m_{\gamma\gamma}$ distributions, these MC corrections substantially improve the agreement in kinematic distributions, as described below.

Decay candidates $D^+ \to K^-\pi^+\pi^+$ are formed from well-measured tracks with kaon or pion particle identification and with a $K^-\pi^+\pi^+$ invariant mass $m_{K\pi\pi}$ within 1.86 and 1.88 GeV (approximately $\pm 2\sigma$ around the nominal D^+ mass [4]). This reduces background from random combinations of tracks, especially from $D^* \to D\pi_s^0$ decays with a correctly reconstructed π_s^0, which will also peak in the signal region of the Δm distribution. As in Ref. [7], we reject candidates with any D^+ daughter track for which the cosine of the polar angle measured in the laboratory frame $\cos\theta_t$ is above 0.89; this criterion reduces the final sample by approximately 10%. To further suppress peaking background events, we use a likelihood variable to select D^+ candidates, based on measured decay vertex separation from the primary vertex, and on Dalitz-plot positions. This likelihood criterion rejects about 70% of background events with incorrectly reconstructed D^+, while retaining about 77% of signal events. Figure 1 shows the $m_{K\pi\pi}$ distribution for data events passing all selection criteria except for the requirement on $m_{K\pi\pi}$. For illustrative purposes, we fit the $m_{K\pi\pi}$ distribution by modeling the D^+ signal with a sum of two Gaussian functions sharing a common mean and random background events with a linear function. After all selection criteria, the fraction of candidates with a correctly reconstructed D^+, as estimated from the $m_{K\pi\pi}$ fit, is about 95%.

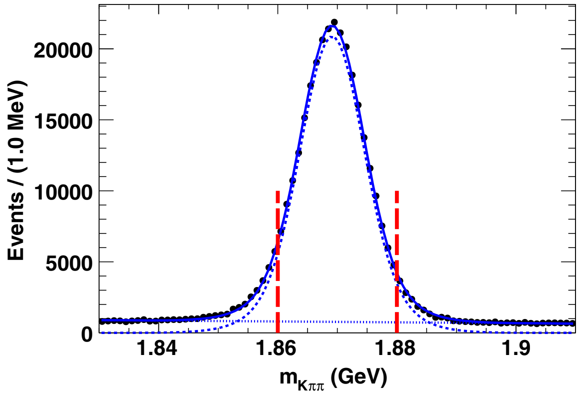

FIG. 1. The reconstructed D^+ mass distribution of real data, after all D^{*+} selection criteria except for the D^+ mass requirement, which is marked by the two vertical dashed lines. The result of the fit described in the text is superimposed (solid line), together with the background (dotted line) and signal (dashed line) components.

The value of Δm_+ is obtained from a fit to the Δm distribution in a two-step procedure as illustrated in Figs. 2(a) and 2(b). First, we model the Δm resolution function by fitting the Δm distribution for correctly reconstructed signal MC events using an empirically motivated sum of three Gaussian or Gaussian-like probability density functions (PDFs)

$$\begin{aligned}\mathcal{S}(\Delta m) = {} & f_1 G(\Delta m; \Delta m_+ + \delta_{\Delta m_+}, \sigma_1) \\ & + (1-f_1)[f_2 \mathrm{CB}(\Delta m; \Delta m_+ + \delta_{\Delta m_+}, \sigma_2, \alpha, n) \\ & + (1-f_2) \mathrm{BfG}(\Delta m; \Delta m_+ + \delta_{\Delta m_+}, \sigma_3^L, \sigma_3^R)], \qquad (1)\end{aligned}$$

where f_1 and f_2 give the fractions for the composite PDFs of Gaussian (G), crystal ball (CB [12], with α and n as two parameters to model the high mass tail), and BfG [a two-piece normal distribution with widths σ_3^L and σ_3^R on the left and right of $(\Delta m_+ + \delta_{\Delta m_+})$, respectively]. The sum $(\Delta m_+ + \delta_{\Delta m_+})$ is, therefore, the common peak position of the three PDFs. In the fit to the high-statistics MC sample [Fig. 2(a)], Δm_+ is fixed at the generated value of 140.636 MeV, and $\delta_{\Delta m_+}$ is a measure of the possible bias induced by our event selection procedure or the chosen form for the resolution function. The fitted functional distribution provides a reasonably good description of the data (with $\chi^2/\nu = 605/491$ for a sample more than seven times larger than the data). The fit gives $\delta_{\Delta m_+} = (+16.6 \pm 2.5)$ keV, with the uncertainty from the limited size of our MC sample. The fit results for the shape parameters are shown in Fig. 2(a); and the full-width at

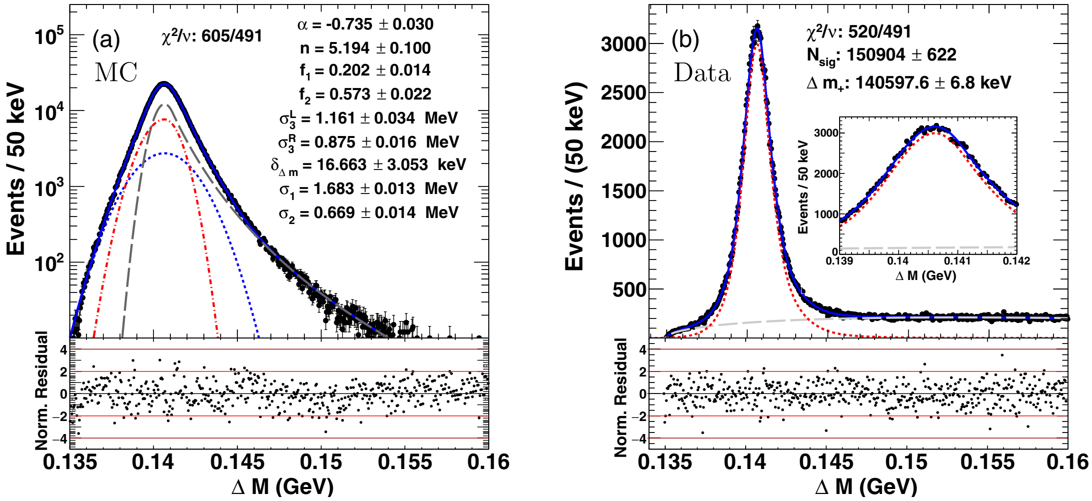

FIG. 2. (a) Δm fit to correctly reconstructed signal MC events. Shown are the total fit (blue solid line), Crystal Ball function (gray long-dashed line), Gaussian (blue short-dashed line), and two-piece normal distribution function (red dashed-dotted line). The fitted signal shape parameters defined in Eq. (1) are also shown in the text box. (b) Δm fit to real data. Shown are the total fit (blue solid line), signal PDF (magenta short-dashed line), and background PDF (gray long-dashed line). The inset shows the fit around the peak region. The Δm_+ central value from the fit is later corrected by the estimated fit bias. Normalized residuals shown underneath both fit plots are defined as $(N_{\rm observed} - N_{\rm predicted})/\sqrt{N_{\rm predicted}}$.

half maximum (FWHM) of the resolution function is found to be about 2.1 MeV, which is mainly due to the resolution of the π^0_s.

The second step [Fig. 2(b)] is an unbinned maximum-likelihood fit to real data using the PDF from the first step to model signal and a threshold function to model the combinatorial background [13]

$$T(\Delta m;\kappa) = \Delta m\sqrt{u}\exp(\kappa u), \tag{2}$$

where $u = (\Delta m/m_{\rm endpt})^2 - 1$, and κ is the slope parameter which is allowed to vary in the fit. We fix the end point $m_{\rm endpt}$ at the nominal π^0 mass [4] as the physical limit of Δm. In the data fit, we fix the bias $\delta_{\Delta m_+}$, fractions $f_{1,2}$, and CB tail parameters to the MC values from the first step, while allowing the widths $\sigma_{1,2,3}$ to be free in the fit to allow for differences between MC simulation and data. Figure 2(b) presents the data and the fit, with the normalized residuals showing good data and fit agreement. There are $150\,904 \pm 622$ signal events, the observed FWHM of the signal shape is about 2.0 MeV, and we determine $\Delta m_+ = (140597.6 \pm 6.8)$ keV, where the uncertainty is statistical only ($\sigma_{\rm stat}$). A bias correction to this result will be discussed later.

We estimate systematic uncertainties on Δm_+ from a variety of sources. Separately, we study the Δm_+ dependence on the D^{*+} laboratory momentum $p_{\rm lab}$, on the cosine of D^{*+} laboratory polar angle $\cos\theta$, on the D^{*+} laboratory azimuthal angle ϕ, on $m_{K\pi\pi}$, and on the diphoton opening angle $\theta_{\gamma\gamma}$ from $\pi^0 \to \gamma\gamma$, by collecting fit results for Δm_+ in ten subsets of data with roughly equal statistics for each parameter. Furthermore, we divide our data into four disjoint subsets of data-taking periods. For the data fit in each subset, the value of $\delta_{\Delta m_+}$ is determined separately from signal MC events with the same event selection criteria as for that subset. This is meant to expose possible detector response effects that have not been modeled in the simulation. We search for variations larger than those expected from statistical fluctuations based on a method similar to the PDG scale factor [4,8]. If the fit results from a given dependence study are compatible with a constant value, in the sense that $\chi^2/\nu < 1$, where ν is the number of degrees of freedom, we assign no systematic uncertainty. In the case that $\chi^2/\nu > 1$, we ascribe an uncertainty of $\sigma_{\rm sys} = \sigma_{\rm stat}\sqrt{\chi^2/\nu - 1}$ to account for unidentified detector effects. We observe $\chi^2/\nu > 1$ in the cases of $p_{\rm lab}$, $\cos\theta$, and $\theta_{\gamma\gamma}$ (shown in [14]). Systematic uncertainties of 5.0, 6.9, and 6.1 keV are assigned for the $D^{*+}p_{\rm lab}$, $D^{*+}\cos\theta$, and $\theta_{\gamma\gamma}$ dependences, respectively, for which the p values for the null hypotheses are 0.12, 0.03, and 0.06. The p values for the variations with D^{*+} azimuthal angle and D^+ mass are 0.99 and 0.47, and no systematic uncertainties are assigned for these observations.

The five signal shape parameters α, n, $f_{1,2}$, and $\delta_{\Delta_{m_+}}$, determined from the fit to signal MC events [Fig. 2(a)], possess statistical uncertainties that are highly correlated. We account for their uncertainties and correlations by producing 100 sets of correlated random numbers of signal shape parameters based on the central values and the covariance matrix from the fit to signal MC events. Then, for each set, we rerun the data fit by fixing α, n, $f_{1,2}$, and $\delta_{\Delta_{m_+}}$ to the corresponding random numbers in the set. The distribution of the 100 fit values for Δ_{m_+} has a root mean square of 2.1 keV which is taken as systematic uncertainty for the signal shape parameters.

To test whether our fit procedure introduces a bias on Δm_+, we generate an ensemble of data sets with signal and background events generated from appropriately normalized PDFs based on our nominal data fit. The data sets are then fitted with exactly the same fit model as for real data ("pure pseudoexperiment"). By performing 500 pseudoexperiments, we collect Δm_+ pulls, defined as the differences of fitted and input values normalized by the fitted errors. The mean of the pulls is $-(50 \pm 4)\%$, while the root mean square is consistent with being unity. Thus, we correct for the bias in our fit model by adding $50\% \times \sigma_{\rm stat} = 3.4$ keV to the fit value of Δm_+ from the data, and assign a systematic uncertainty equal to half this bias correction (1.7 keV). We perform another type of pseudoexperiment by fitting to ensembles of data sets where signal and background events are produced by randomly sampling the corresponding MC events. Background events from decays such as $D^{*+} \to D^+\pi^0$ with $D^+ \to \pi^-\pi^+\pi^+\pi^0$ misreconstructed as $K^-\pi^+\pi^+$ produce small peaks in the signal region, but the fit does not account for them explicitly. The collected pulls show a mean fit bias consistent with that found in our pure pseudoexperiments, and we assign no additional systematic uncertainty related to peaking backgrounds.

To account for the systematic uncertainty due to imperfect photon energy simulation and calibration in the MC simulation, we rescale photon energies in signal MC events by $+0.3\%$ and -0.3%, and take the larger of the two variations in the Δm peak position, 7.0 keV, as the corresponding systematic uncertainty. The values $\pm 0.3\%$ correspond to the difference between MC and data π^0 mass peak positions after the nominal MC neutral energy corrections are applied. Because the MC and data $m_{\gamma\gamma}$ distribution shapes differ, aligning the peak positions does not produce equal mean values. We also account for the associated uncertainties on the π^0 momentum rescaling factors due to the limited size of our MC sample and find the related systematic uncertainty to be 0.5 keV.

Besides the systematic studies, we also perform a series of consistency checks that are not used to assess systematics but, rather, to reassure us that the experimental approach and fitting technique behave reasonably. We vary the upper limit of the Δm fit range from its default position of 0.160 GeV to a series of values between 0.158 and 0.168 GeV. Also, we vary the selection criteria on the invariant masses $m_{K\pi\pi}$ and $m_{\gamma\gamma}$, as well as the Dalitz-plot based likelihood. The resulting fit values of Δm_+ from all these checks are consistent.

All systematic uncertainties of Δ_{m_+} are summarized in Table I; adding them in quadrature leads to a total of 12.9 keV. After adding the fit bias of 3.4 keV, our final result is $\Delta m_+ \equiv m(D^{*+}) - m(D^+) = [140601.0 \pm 6.8(\text{stat}) \pm 12.9(\text{syst})]$ keV. This result is consistent with the current world average of (140.66 ± 0.08) MeV, and about five times more precise. Combining with the *BABAR* measurement of $\Delta m_0 = [145\,425.9 \pm 0.5(\text{stat}) \pm 1.8(\text{syst})]$ keV based on the same data set, we obtain the D meson mass difference of $\Delta m_D = [4824.9 \pm 6.8(\text{stat}) \pm 12.9(\text{syst})]$ keV. This result is, as for Δm_+, about a factor of 5 more precise than the current world average, (4.77 ± 0.08) MeV. Adding the statistical and systematic uncertainties in quadrature, $\Delta m_D = (4824.9 \pm 14.6)$ keV. This can be compared with the corresponding values for the pion and kaon systems, $\Delta m_\pi = (4539.6 \pm 0.5)$ keV and $\Delta m_K = (-3934 \pm 20)$ keV [4].

TABLE I. Assigned systematic errors from all considered sources.

Source	Δm_+ systematic [keV]
Fit bias	1.7
D^{*+} $p_{\rm lab}$ dependence	5.0
D^{*+} $\cos\theta$ dependence	6.9
D^{*+} ϕ dependence	0.0
$m(D^+_{\rm reco})$ dependence	0.0
Diphoton opening angle dependence	6.1
Run period dependence	0.0
Signal model parametrization	2.1
EMC calibration	7.0
MC π^0 momentum rescaling	0.5
Total	12.9

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