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Observations of coastal cliff base waves, sand levels, and cliff top shaking

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ABSTRACT: Concurrent observations of waves at the base of a southern California coastal cliff and seismic cliff motion were used to explore wave–cliff interaction and test proxies for wave forcing on coastal cliffs. Time series of waves and sand levels at the cliff base were extracted from pressure sensor observations programmatically and used to compute various wave impact metrics (e.g. significant cliff base wave height). Wave–cliff interaction was controlled by tide, incident waves, and beach sand levels, and varied from low tides with no wave–cliff impacts, to high tides with continuous wave–cliff interaction. Observed cliff base wave heights differed from standard Normal and Rayleigh distributions. Cliff base wave spectra levels were elevated at sea swell and infragravity frequencies. Coastal cliff top response to wave impacts was characterized using microseismic shaking in a frequency band (20–45 Hz) sensitive to wave breaking and cliff impacts. Response in the 20–45 Hz band was well correlated with wave–cliff impact metrics including cliff base significant wave height and hourly maximum water depth at the cliff base ($r^2 = 0.75$). With site-specific calibration relating wave impacts and shaking, and acceptable anthropogenic (traffic) noise levels, cliff top seismic observations are a viable proxy for cliff base wave conditions. The methods presented here are applicable to other coastal settings and can provide coastal managers with real time coastal conditions. Copyright © 2016 John Wiley & Sons, Ltd.

KEYWORDS: coastal cliff; wave–cliff interaction; wave impact; cliff shaking; microseismic

Introduction

Wave erosion is a fundamental process of coastal cliff morphology yet field studies of wave–cliff interaction are rare. Observations of cliff base waves (i.e. Stephenson and Kirk, 2000; Young et al., 2011a; Dickson and Pentney, 2012) and sand levels (Robinson, 1977), and the corresponding cliff erosion rate, are lacking (Naylor et al., 2010; Sunamura, 2015). Unvalidated parameterizations of marine forcing and cliff response are therefore necessarily used in model simulations and studies of past and future cliff erosion (e.g. Sunamura, 1982; Budetta et al., 2000; Young et al., 2009; Revell et al., 2011; Castedo et al., 2012; Hackney et al., 2013). Here, concurrent field observations of cliff top shaking, and waves and sand levels at a cliff base, are used to compare metrics of wave–cliff interaction, and to explore recent suggestions that seismic cliff top observations (at particular frequencies) provide a useful proxy for wave impacts on a cliff. Automated methods to extract time series of beach elevation and wave heights from a buried pressure sensor are also presented.

Wave impact forces

Wave-impact pressures vary widely depending on the type of wave–cliff interaction, generally classified as broken (bore) impacts, breaking impacts (waves breaking into the cliff), and reflective. Oumeraci et al. (1993) further divided breaking impact types into three categories based on air entrapment. Impulsive and pulsating wave loads are both observed on coastal structures. Many existing empirical and theoretical equations of wave loads on cliffs and seawalls depend on the wave height at the cliff base ($H_{\text{Cliff}}$) and bore speed ($C$). Example force ($F$) estimates are:

$$F = A \rho g H_{\text{Cliff}}$$

(1)

where $A$, $\rho$, and $g$ are a calibration constant, water density, and gravitational acceleration, respectively (Sunamura, 1977).

$$F = 4.5 \rho g H_{\text{Cliff}}^2$$

(2)


$$F = 0.5 \rho C^2 H_{\text{Cliff}}$$

(3)

(Larson et al. [2004] swash bore force).

Background

Methods to characterize the potential of waves to cause coastal erosion incorporate a variety of measured and estimated wave characteristics.
\[ F = 4.76 \rho g H_{\text{cliff-sig}}^2 \]  
(Cuomo et al. [2010] quasi-static load).

\[ F = F^* \rho g H_{\text{cliff-sig}}^2 \]  
(Oumeraci et al. [2001] maximum impulse force) where \( F^* \) is the relative maximum wave force, assumed to follow a Generalized Extreme Value distribution. Blackmore and Hewson (1984) found average pressures from broken waves of:

\[ P = 4 \rho T C^2 \]

where \( \lambda \) is aeration factor, and \( T \) is wave period.

Wave–cliff interaction proxies

Simple proxies for marine forcing at coastal cliffs include incident wave height, energy (Davidson-Arnott and Oliverhead, 1995), and power (Kamphuis, 1987; Brown et al., 2005). These proxies neglect local influences such as shoaling, dissipation, runup, and beach elevation (although they are sometimes site-calibrated to incorporate local influences). Some studies include site conditions using a total water level (TWL) method, where TWL is the sum of tides and the vertical height of wave setup and runup (Shih et al., 1994; Kirk et al., 2000; Ruggiero et al., 2001). When TWL exceeds the cliff base beach elevation, wave–cliff impact occurs. The duration of wave attack provides an indicator of marine forcing (Ruggiero et al., 2001; Sallenger et al., 2002) and has been used in previous cliff studies (i.e. Young et al., 2009, 2011b; Young, 2015). However, detailed beach elevation data are often unavailable and if lacking, Collins and Sitar (2008) suggested using maximum daily TWL relative to seasonal elevation averages.

Methods neglecting cliff base beach elevations as an important factor are clearly insufficient (when beaches are present). At the other extreme, models that only consider wave impact duration neglect the relative amount of energy available for potential marine-driven erosion. Brown et al. (2005) and Swenson et al. (2006) derived marine-forcing proxies that include wave impact height using the cliff base water depth defined as the difference between TWL and cliff base beach elevation. TWL methods are based on maximum runup and neglect the distribution of wave impact heights. Hughes et al. (2010) explored wave runup distributions on a sandy coast, but there have been no similar studies on clifffaced coasts. Lacking cliff base wave observations, Wilcock et al. (1998) developed a frequency distribution of impact wave height from synthetic wave records transformed across the nearshore to estimate wave–cliff impact pressures and a proxy for wave–cliff interaction.

Other studies have included the influence of site conditions indirectly. Amin and Davidson-Arnott (1997) used incident wave energy as the marine-forcing proxy but also included sediment availability, cliff height, and potential longshore transport using a multi-variable regression. Benumof et al. (2000) defined a proxy of ‘wave power at the cliff’ as wave power multiplied by vertical runup and suggested that because slope is included in the runup calculation, local site conditions are inherently addressed. Kline et al. (2014) developed a numerical cliff retreat model using wave forcing established from wave impact height modified by beach elevation and water level. Trenhaile (2000) included local site effects and developed a proxy from breaker height, surfzone width, type of breaking waves, and roughness of the rock surface. Sunamura (1982), Wilcock et al. (1998), and Trenhaile (2000) also suggest using a threshold-based metric where a minimum site specific value is necessary to cause erosion.

Microseismic cliff motions

Seismic observations at coastal cliffs suggest ground motions generated by local ocean waves may provide a convenient proxy for wave–cliff impacts (long-term logistics for a cliff top seismometer are often easier than for a cliff base pressure sensor). Cliff ground motions are driven by different local mechanisms in two general frequency bands. Low-frequency cliff motion (0.01–0.1 Hz) or ‘flexing’ is generated by individual sea, swell (Adams et al., 2005), or infragravity (Young et al., 2011a, 2012, 2013) waves that load the foreshore with pressure fluctuations and exert gravitational attraction (Agnew and Berger, 1978). Double frequency (DF) motions are non-locally generated and propagate trans-continental distances (e.g. Bromirski, 2001). Higher frequency cliff motion (HF, > 0.3 Hz) or ‘shaking’ is generated from ocean waves directly impacting the cliff (Adams et al., 2002) or fronting shore platforms (Dickson and Pentney, 2012), wind (Norman et al., 2013), waves breaking in the nearshore (Young et al., 2013; Poppeliers and Mallinson, 2015), and anthropogenic sources such as vehicle traffic. These various contributions to HF shaking frequencies complicate the isolation of seismic signals from wave–cliff impacts. Peaks at various frequencies in HF energies have been observed at several sites (Dickson and Pentney, 2012; Young et al., 2013) and are possibly related to site conditions, geometry, and geologic composition.

Lim et al. (2011) attributed cliff shaking (<300 Hz) to wave–cliff impacts and found distinct water levels were associated with an elevated cliff response. Norman et al. (2013) observed elevated cliff shaking signals (1.1–50 Hz) only during combined elevated water levels and waves. Norman et al. (2013) and Vann Jones et al. (2015) selected shaking at 9.6 Hz to represent wave–cliff interaction because it experiences the highest amplitudes without overlapping other shaking signals (i.e. wind generated). Vann Jones et al. (2015) found cliff shaking (9.6 Hz) correlated (multiple regression, \( R^2 = 0.53 \)) with modeled cliff base wave height (includes the influence of tide) and wave setup, supporting the use of cliff motion as a proxy for wave–cliff impacts. Previous studies lack quantitative analysis of concurrent observations of cliff base waves and cliff shaking. Here, pressure sensor observations at the base of a coastal cliff are used to determine the heights and frequency of wave–cliff impacts, and to measure cliff base beach elevations. Observations of cliff ground motions are examined in the context of the wave observations and associated metrics of wave–cliff interaction.

Study Site

Cliff setting

The studied 24 m high cliff, located in northern Del Mar, California, USA, consists of three geologic units (Figure 1). The lower unit is the Del Mar Formation, an Eocene sedimentary deposit composed of sandy claystone interbedded with coarse-grained sandstone, overlain conformably by Torrey Sandstone, a massive coarse-grained and well cemented Eocene sandstone (Kennedy, 1975). Together, these two units form the lower near-vertical portion of the cliff up an elevation of about 8 m (navd88). The upper cliff section is about 16 m thick, slopes at 35°–50° and consists of weakly cemented, fine-grained sandy Pleistocene terrace deposits. The contact
between the Del Mar and Torrey Sandstone Formations decreases in elevation toward the north and terminates abruptly at a fault immediately north of the instrumentation setup. The lower Eocene units are harder and more resistant to erosion compared to the softer marine terrace deposits. Typical cohesion and friction angle values for undisturbed Del Mar Formation and Terrace Deposits are 31 kPa, 47°, and 14 kPa, 36°, respectively (Leighton and Associates, 2003). The cliff is fronted by a narrow dissipative sand (and occasionally cobble) beach, which is often submerged during high tides. The underlying shore platform composed of the Eocene units is gently sloping and relatively smooth near the shoreline, and more irregular offshore, forming nearshore reefs. The shore platform elevation at the cliff base is close to mean lower low water (about 1 m navd88). Sand thickness at the cliff base is typically several meters, but is occasionally reduced to zero when winter storm waves expose the shore platform. The shore platform extends alongshore 0.4 km southward and 2.4 km northward. Offshore the platform contains transient sand pockets and plants.

Oceanographic setting

The cliffs are exposed to waves generated by local winds and distant storms in both hemispheres. During winter, swell from the North Pacific and Gulf of Alaska is most energetic, whereas swell from the South Pacific dominates in summer. Waves reaching southern California cliffs undergo a complex transformation, and ‘shadows’ of the Channel Islands create strong alongshore variations in wave height (e.g. Pawka, 1983). The seasonal cycle in the Del Mar region has maximum wave energy in winter. The tide range is approximately 2 m (http://tidesandcurrents.noaa.gov).

Methods

Incident waves

A wave buoy network (CDIP, http://cdip.ucsd.edu) was used to estimate hourly significant wave height ($H_{10}$, Figure 2a) and peak period ($T_p$) at a virtual buoy seaward of the study area in 10 m water depth. The effects of complex bathymetry in the southern California Bight, and of varying beach orientation and wave exposure, were simulated with a spectral refraction wave model initialized with offshore buoy data (O’Reilly and Guza, 1991, 1993, 1998).

Tides (offshore water level)

Hourly water levels seaward of the surfzone ($h_{wl}$), including tides, atmospheric pressure and wind effects, were obtained from the La Jolla tide gauge #94101230 (http://tidesandcurrents.noaa.gov), located in about 7 m water depth 12 km south of the study site.

Modeled total water level (TWL)

The modeled hourly TWL is the sum of offshore water levels and vertical runup

$$TWL = h_{wl} + R_{2\%}$$

where $R_{2\%}$ is the level exceeded by 2% of wave uprushes,

$$R_{2\%} = 1.1 \left\{ 0.35\beta_f (H_o L_o)^{0.5} 
+ \left[ (H_o L_o (0.563\beta_f^2 + 0.004))^{0.5} \right] / 2 \right\}$$

where $H_o$ and $L_o$ characterize the incident deep water wave height and wavelength, respectively (Stockdon et al., 2006). The value of $H_o$ was calculated by backing out $H_{10}$ to deep water by reverse shoaling using linear wave theory, while $L_o$ was calculated using the deep water linear dispersion relationship and the peak period. The beach slope ($\beta_f$) was estimated as the mean slope of a 20 m horizontal swath centered on the intersection of hourly mean sea level and interpolated beach profile from four ATV-GPS (all terrain vehicle equipped with differential global positioning system) surveys (conducted on October 11, 2010, November 8, 2010, December 17, 2010, and January 4, 2011).

Cliff base pressure sensor

A Paroscientific pressure sensor (model# 245A-102), sampling at 8 Hz from November 4, 2010 to December 20, 2010, was located on the shore platform (1.02 m elevation, datum-NAVD88) approximately 4 m shoreward of the cliff base. Atmospheric pressure was removed from the record using linearly interpolated six-minute observations collected approximately 12 km south of the study site on a pier.

Figure 1. (left) Cliff study site, location map and (right) cliff profile, geology, and instrumentation.
Pressure time series were corrected for clock drift and converted to hydrostatic elevation relative to NAVD88.

Automated beach elevation extraction
After a wave recedes, the fully saturated subaerial beach can be exposed, and the water level is relatively constant for several seconds (Figure 3b, Raubenheimer et al., 1995). Beach elevations were detected by locating points in the time series with slope (running mean over a 10 second window) near zero, 25–35 seconds after a wave crest passes (red dots in Figures 3b and 3c). Multiple points identified between successive waves were reduced to a single representative beach elevation, yielding 3269 values, generally at the beginning and end of the high tide cycles. Beach elevations from the pressure sensor were consistent with ATV-GPS beach surveys on November 8, 2010 and December 17, 2010.

Automated wave height estimation
Discontinuous wave time series at the cliff base precluded using typical zero-crossing wave detection methods. Instead, wave peaks and troughs were detected programmatically using consecutive local minima and maxima (similar to the methods of Hughes et al., 2010) and a minimum wave height threshold of 1 cm. When the interval between waves was long enough, drainage from the beach (Figure 3c) generated local water level minima unrelated to waves at the cliff base. These subsurface minima were programmatically excluded from wave height estimates.

Wave–cliff interaction metrics
Cliff base wave heights ($H_{C_{\text{Cliff}}}$) and water levels from the pressure sensor observations (water levels include tide, wave height, setup, etc.) were used to calculate the following observed hourly wave–cliff interaction metrics (Figure 2c):

- **M1**: Variance of cliff base water level (about the mean, water table fluctuations removed)
- **M2**: Number of cliff base waves ($N$)
- **M3**: Sum of cliff base wave heights ($\Sigma H_{C_{\text{Cliff}}}$)
- **M4**: Sum of squared cliff base wave heights ($\Sigma H_{C_{\text{Cliff}}}^2$)
- **M5**: Duration of wave attack (fraction during each hour)
- **M6**: Significant cliff base wave height ($H_{C_{\text{Cliff-Sig}}}$, average of 1/3 highest $H_{C_{\text{Cliff}}}$)
- **M7**: Significant cliff base wave height squared ($H_{C_{\text{Cliff-Sig}}}^2$)
- **M8**: $H_{C_{\text{Cliff-Sig}}}^2 \times N$ ($M8 = M7^* M2$)
- **M9**: Swash wave force (Equation 3, Larson et al., 2004) where wave speed is:
  \[ C = (g^*(\text{wave peak elevation} – \text{beach elevation}))^{0.5} \]  

- **M10**: Maximum instantaneous water depth (maximum water elevation – beach elevation)
Cliff top ground motions

Ground motions were measured at 100 Hz with a Nanometrics Compact Trillium broadband seismometer near the cliff top edge (elevation 23.5 m, NAVD88), 26 m shoreward of the pressure sensor. The raw vertical velocity data were phase and magnitude corrected in the frequency domain according to the instrument response curve. Seismic and cliff base water levels, divided into one hour records, were processed with standard Fourier spectral and cross-spectral methods (Jenkins and Watts, 1968). Hours containing significant ground motion from earthquakes, post-installation settlement, or local noise were removed manually. Spectral seismic data was used to identify both useful wave–cliff interaction metrics and the specific seismic frequency range (20–45 Hz) suited for impact detection. All correlations presented are statistically significant \((p < 0.05)\).

Observations

General conditions

Incident wave heights \((H_{10})\) ranged from 0.36 to 1.84 m (average 0.83 m, Figure 2a). The observed water level exceeded the cliff base beach elevation during 46% of the study period (513 of 1123 hours). Continuous wave–cliff interaction occurred for 89 hours, during the highest tides. Hours with discontinuous wave–cliff interaction occurred more frequently (424 hours), usually before and after high tides, but also occasionally during entire high tide cycles. Waves did not reach the cliff during low tide. Beach elevations at the cliff base fluctuated about 0.59 m over the study period (1.03–1.62 m, NAVD88), with a maximum change of 26 cm over a single high tide cycle. Negative pressures, yielding observed water levels as much as 20 cm below the sensor (Figure 2b) are likely associated with the capillary fringe (Cartwright et al., 2006).

Cliff seismic motions generally increased with increasing tide and incident wave energy, consistent with previous studies. Cliff shaking energies, wave impact metrics, and pressure sensor energies are all tidally modulated (Figures 2c–2e). Spectra of cliff top vertical ground velocity are most energetic at infragravity and single frequencies. The HF shaking band has peaks at about 1 and 7 Hz (Figure 4b).

Total water level (TWL)

Modeled TWL ranged from about 1.5 m above to 1.5 m below the cliff base sand level (green relative to red curves in Figure 2b). Modeled TWL generally agreed well with observed maximum water elevation at the cliff base (M10) for hours with observed wave–cliff interaction \((r^2 = 0.81, \text{Figure 2b})\). Some offset (about 10 cm in Figure 2b) is expected, given the substantial scatter in the underlying Stockdon et al. (2006) parameterization (Equation 8). Overall, TWL provides reasonable estimates of the observed maximum cliff base water level when compared to the extracted beach elevations, correctly determining hours with wave–cliff impacts 90% of the time. However without consideration of beach elevation, TWL alone is not a direct measure of wave–cliff interaction because sand levels influence whether wave–cliff interaction occurs.

Figure 3. Observed water level versus time (blue curve) with detected wave peaks, troughs and beach elevations for time periods of (a) continuous cliff base waves, (b) discontinuous waves and fully saturated beach between waves, and (c) discontinuous waves and water levels draining below beach elevation. In (a), much of the variance is at infragravity frequencies (green curve).
Observed cliff base wave heights

Cliff base wave heights ranged up to 1.17 m with mean and standard deviation both equal to 0.13 m. Binned (0.05 m bin size) hourly wave height distributions (hours with at least 100 waves) were consistently skewed with relatively more small waves than a Rayleigh distribution and often peaked in the minimum bin (0–0.05 m, Figure 5a). Wave heights increase with hourly maximum cliff base water depth (M10), causing an increase in the cumulative hourly wave height distribution and a wider range of wave heights (Figure 6). The ensemble-average of the measured probability distribution differs from both standard Rayleigh and Normal distributions often used to describe offshore wave heights (Figure 5).

Observed cliff base wave spectra

Hourly cliff base wave energy spectra generally increased with tide and cliff base water depth (Figure 2e). Spectra levels increased with wave period and exhibited high infragravity energies (Figure 4a), consistent with recent research highlighting the importance of infragravity waves in rock coast processes.
Observed wave-impact metrics

All impact metrics (Figure 2c) were tidally modulated and usually mutually correlated (Table I; mean $r^2 = 0.68$, median $r^2 = 0.71$). M5 was the least correlated overall ($r^2$ between 0.21–0.8, mean 0.50), while M1, M3, M4, and M8 were best correlated (mean $r^2 = 0.75$). For a Rayleigh wave height distribution M4 and M8 are linearly correlated ($r^2 = 1$, $M8 = 2^*M4$). Deviations from the Rayleigh distribution alter the observed proportionality to $M8 \sim 2.2^*M4$. During periods of intermittent swash (when the wave trough was equivalent to beach elevation) M9 and M4 differ only by a constant because squared bore speed is proportional to wave height ($H_{Cliff}$), and $r^2 = 1$. When waves are continuous, bore wave speed (Equation 9) included the variable depth below the wave trough and is not proportional to wave height ($H_{Cliff}$), causing a deviation between M9 and M4. However the deviation only causes a slight decrease in correlation and correlations between M9, M4, and similarly M8 are high ($r^2$ ranges from 0.92 to 1.00).

Discussion

Correlation of wave impact metrics with shaking

Impact metrics show higher correlations (Figure 7) with seismic energy at cliff impact frequencies (20–45 Hz). Cliff impact energy has the highest correlation $r^2 > 0.7$ with impact metrics M5, M6, and M10 (Figure 7, Table I). Metrics related to $\Sigma H_{Cliff}^2$ (M4, M8, and M9) are least correlated with cliff impact energy ($r^2 = 0.34$–0.48). The 9.6 Hz wave impact band of Vann Jones et al. (2015) is not correlated with any metrics (Figure 7), suggesting site conditions influence cliff impact seismic response consistent with previous research (Dickson and Pentney, 2012; Young et al., 2013). Correlations are elevated slightly ($r^2 \sim 0.3$) around 1 Hz and 7 Hz, corresponding to peaks in the cliff top HF energy spectra (Figure 4b). Cliff impact energy levels are much lower than these peak HF frequencies (Figure 4b), but the $r^2$ of cliff impact with impact metrics is much higher (Figure 7).

Effect of conditions on wave impacts

The ability of waves to interact with the cliff is largely controlled by the difference between tide (observed) and beach elevation (M10 is shown in Figure 8a). All metrics except M6 and M7 were maximum during two consecutive hours with

Table 1. The $r^2$ correlation matrix of wave-cliff impact metrics and cliff impact shaking (CI, 20–45 Hz)

<table>
<thead>
<tr>
<th></th>
<th>M1</th>
<th>M2</th>
<th>M3</th>
<th>M4</th>
<th>M5</th>
<th>M6</th>
<th>M7</th>
<th>M8</th>
<th>M9</th>
<th>CI</th>
</tr>
</thead>
<tbody>
<tr>
<td>M1</td>
<td>1.00</td>
<td>0.69</td>
<td>0.83</td>
<td>0.84</td>
<td>0.51</td>
<td>0.78</td>
<td>0.87</td>
<td>0.85</td>
<td>0.65</td>
<td>0.85</td>
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<tr>
<td>M2</td>
<td>1.00</td>
<td>0.91</td>
<td>0.75</td>
<td>0.57</td>
<td>0.53</td>
<td>0.43</td>
<td>0.74</td>
<td>0.71</td>
<td>0.72</td>
<td>0.56</td>
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<tr>
<td>M3</td>
<td>1.00</td>
<td>0.95</td>
<td>0.45</td>
<td>0.57</td>
<td>0.57</td>
<td>0.94</td>
<td>0.90</td>
<td>0.90</td>
<td>0.75</td>
<td>0.36</td>
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<tr>
<td>M4</td>
<td>1.00</td>
<td>0.34</td>
<td>0.54</td>
<td>0.63</td>
<td>1.00</td>
<td>0.93</td>
<td>0.69</td>
<td>0.48</td>
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<tr>
<td>M5</td>
<td>1.00</td>
<td>0.80</td>
<td>0.54</td>
<td>0.34</td>
<td>0.21</td>
<td>0.77</td>
<td>0.73</td>
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<tr>
<td>M6</td>
<td>1.00</td>
<td>0.87</td>
<td>0.54</td>
<td>0.34</td>
<td>0.93</td>
<td>0.75</td>
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<tr>
<td>M7</td>
<td>1.00</td>
<td>0.64</td>
<td>0.39</td>
<td>0.81</td>
<td>0.63</td>
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<tr>
<td>M8</td>
<td>1.00</td>
<td>0.92</td>
<td>0.69</td>
<td>0.48</td>
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<tr>
<td>M9</td>
<td>1.00</td>
<td>0.51</td>
<td>0.34</td>
<td></td>
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<tr>
<td>M10</td>
<td>1.00</td>
<td>0.75</td>
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elevated tide, moderate waves ($H_{10} = 1.17$ m), and near minimum beach sand elevations (1.08 m NAVD88, Figure 8b). At the same tide level, small waves with depleted beaches can cause more cliff shaking than larger waves with an elevated beach.

Cliff shaking as a wave impact proxy

Infragravity (IG) signals increase approximately linearly with cliff base water depth (TWL, Figure 9a), consistent with forcing by nearshore loading and gravitational attraction. Cliff impact ground motions trend consistently only during hours of wave–cliff interaction (cliff base depths > 0, Figure 9b). During hours without wave–cliff interaction, traffic noise may contribute to the scattered seismometer energy between $10^{-3.5}$ and $10^{-4}$, independent of cliff base water depth. Traffic noise may contaminate shaking band signals at this and other sites.

Elevated cliff impact signals during periods of wave–cliff interaction and good correlations between cliff impact energy and several cliff base wave metrics suggest that cliff motions probably provide a proxy for wave–cliff impacts at certain site-specific frequencies. However, bursts of elevated cliff top ground motion were not always coincident with peaks in cliff base water levels (Figure 10), possibly because the seismometer responds to wave breaking displaced alongshore from the seismometer. The spatial range of detectable wave–cliff impacts is unknown. Similarly, wave breaking on the beach fronting the cliff excites ground motion at the seismometer, as observed at other sites (e.g., Young et al., 2013). Detailed observations of wave breaking and cliff impacts in space and time are needed to link individual wave characteristics to cliff shaking. More studies of concurrent quantitative cliff base wave observations, video monitoring, and coastal and inland seismic motion are needed to reduce uncertainties associated with the generation of coastal seismic motions. Long-term observations are needed to link cliff erosion with wave impacts and/or shaking.

Modeling and monitoring applications

Lacking cliff base wave observations, maximum hourly water depth (M10) can be modeled using beach slope and elevation,
wave models and empirical formulae for runup (for example the TWL method used here). M10 correlates very well \(r^2=0.97\) with observed maximum hourly cliff base wave height and can be used to generate synthetic time series of cliff base waves from the normalized wave height distribution (Figure 6b). Modeled M10 is a convenient measure of marine forcing in cliff models for this site, in lieu of wave impact and seismic observations. M10 is similar to the wave impact metric (modeled cliff base wave height plus wave setup) that Vann Jones et al. (2015) found best correlated with cliff shaking, suggesting M10 could be applicable at other sites and used to improve rock coast models (e.g.Trenhaile, 2011; Limber et al., 2014).

This study describes two methods to monitor wave–cliff interaction and nearshore conditions. A cliff base pressure sensor was used to observe waves, water levels, and beach sand levels at one location. The sand level observations derived from the pressure sensor require less labor than repeated GPS field surveys, and provide high temporal resolution data. Seismic observations can provide a proxy for nearshore conditions and wave–cliff impacts, and can survive storm conditions when in situ instrumentation may fail. Both methods can provide coastal managers with real-time quantitative coastal conditions and contribute new approaches to the growing research on rock coast processes (Naylor et al., 2014).

Summary

Wave heights and beach elevations, programmatically extracted from cliff base pressure sensor observations, are used to explore wave–cliff interaction and develop wave–cliff impact metrics. Spectra of cliff base water levels were elevated at sea swell and infragravity energies. Cliff shaking signals between 20–45 Hz were above noise levels only during periods of wave–cliff interaction and correlate well with several wave–cliff impact metrics. Cliff motions at site specific frequencies appear to provide a proxy for wave–cliff interaction. Impact metrics of hourly cliff base significant wave height, maximum cliff base water depth, and duration of wave attack were best correlated with cliff shaking \(r^2=0.73–0.75\). Maximum hourly cliff base water depth provides a robust wave-based metric for marine forcing and can be modeled with runup equations and beach elevations. Cliff shaking frequencies most correlated with cliff base wave metrics differed from peak frequencies in the seismic signals that are generated by incident wave loading and distant ocean waves. The methods presented here are applicable to a wide range of settings beyond rock coasts, and can provide novel real-time time series of beach sand levels and hydrodynamic conditions.

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