

Lawrence Berkeley National Laboratory

LBL Publications

Title

Combining theory and experiment for X-ray absorption spectroscopy and resonant X-ray scattering characterization of polymers

Permalink

<https://escholarship.org/uc/item/2pf7c7dk>

Authors

Su, Gregory M
Cordova, Isvar A
Brady, Michael A
[et al.](#)

Publication Date

2016-09-01

DOI

10.1016/j.polymer.2016.06.068

Peer reviewed

Combining theory and experiment for X-ray absorption spectroscopy and resonant X-ray scattering characterization of polymers

Gregory M. Su^a, Isvar A. Cordova^a, Michael A. Brady^{a,b}, David Prendergast^b, Cheng Wang^{a,*}

^aAdvanced Light Source, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA

^bThe Molecular Foundry, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA

Abstract

An improved understanding of fundamental chemistry, electronic structure, morphology, and dynamics in polymers and soft materials requires advanced characterization techniques that are amenable to *in situ* and *in operando* studies. Soft X-ray methods are especially useful in their ability to non-destructively provide material or chemical moiety specific information. Analysis of these experiments, which can be very dependent on X-ray energy and polarization, can quickly become complex. Complementary modeling and predictive capabilities are required to properly probe these critical features. Here, we present relevant background on this emerging suite of techniques. We focus on how the combination of theory and experiment has been applied and can be further developed to drive our understanding of how these methods probe relevant chemistry, structure, and dynamics in soft materials.

Keywords: polymer, resonant scattering, resonant reflectivity, NEXAFS, XPCS, simulations, *in situ*

1. Introduction

The intricate connections among chemical structure, local intermolecular and global morphology, and kinetics in polymeric materials determine the fundamental properties of many polymer-based applications. These include flexible electronics [1–4], gas separations [5–7], polymer electrolytes for batteries, fuel cells and water desalination [8–10], and lithographic patterning [11–13] just to name a few. It has remained challenging to fully understand the relationships between chemistry and structure, and the growing demand to characterize polymers under *in situ* or *in operando* conditions relevant to a specific application and as a function of time creates additional challenges. Probing detailed information such as this requires advanced characterization methods that need to be complemented by theory. Energy tunable soft X-rays are the basis for a unique set of tools that are sensitive to molecular and electronic structure, spatial and orientation information, and time-resolved dynamics.

Soft X-rays span an energy range that includes the core-level 1s electronic transitions (K edges) of some of the most common elements found in polymers, for example, carbon, nitrogen, and oxygen. Soft X-ray spectroscopy has an inherent ability to not only selectively probe for the presence of these elements, but also be sensitive to various bonding environments, functional groups, and the orientation of chemical moieties. In addition to the chemistry, soft X-rays can be used for energy-dependent scattering experiments that probe spatial information and characteristic length scales that range from a few to hundreds of nanometers. The interdependence of spectroscopy and scattering allows for unique capabilities to understand, for example, the length scales associated with certain components, moieties or molecular orientations. However, understanding the relationships that connect molecular structure, spectroscopy, and scattering can be very complex, and complementary simulations are needed to build fundamental knowledge and help interpret experiments. An improved ability to leverage theoretical predictions will enable soft X-ray methods to move forward into the realm of *in situ* characterization and bridge the gaps among chemistry, morphology, and dynamics in polymers and soft materials in general.

*Corresponding author

Email addresses: gsu@lbl.gov (Gregory M. Su), iacordova@lbl.gov (Isvar A. Cordova), mabrady@lbl.gov (Michael A. Brady), dgprendergast@lbl.gov (David Prendergast), cwang2@lbl.gov (Cheng Wang)

2. X-ray absorption spectroscopy and resonant scattering probe chemistry and structure

2.1. Connections between X-ray absorption and scattering contrast

Near edge X-ray absorption fine structure (NEXAFS) spectroscopy and resonant soft X-ray scattering (RSoXS) and reflectivity (RSoXR) are naturally complementary techniques, and information regarding a materials absorption properties is needed to understand RSoXS and RSoXR. Unlike hard X-ray scattering or electron microscopy where contrast arises due to differences in electron density, at energies near an absorption edge the complex index of refraction, $n(E) = 1 - \delta(E) + i\beta(E)$, and hence scattering contrast, varies significantly as a function of energy. NEXAFS data provides information related to β , and this can be extended in energy range and scaled based on a material's composition using the Henke atomic scattering factor database [14]. The real part of the index of refraction, δ , can be calculated from β through a Kramers-Kronig integral relation and the resulting contrast between different phases, $\Delta\delta^2 + \Delta\beta^2$, determined as a function of energy. This procedure is often done for organic systems, and an example is shown in Figure 1. It is clear that at energies near an absorption edge, scattering contrast even between materials with similar electron densities can be greatly enhanced. This is one of the main advantages that soft X-rays have for characterizing polymers. However, the detailed connections between a materials absorption properties and resonant scattering can be challenging to discern, for example, contributions from specific chemical moieties or orientation effects. This is an area where progress in effective combination of theoretical predictions and experiments are needed to improve advanced characterization.

2.2. Core-level X-ray absorption spectroscopy

Fundamental knowledge of NEXAFS is a prerequisite to also understanding RSoXS and RSoXR. The subtleties in scattering contrast between polymers stems from the unique absorption profiles of each material. Since NEXAFS is sensitive to specific functional groups and bonds, materials that are composed of the same elements, for example polystyrene and polyethylene which both contain just carbon and hydrogen, can have very distinct NEXAFS profiles. NEXAFS is itself a useful tool in polymer science, and has been widely used to examine thin film properties such as surface composition [15–18] and molecular orientation in ordered polymers [19–23]. Furthermore, absorption differences between polymers can be utilized for chemical sensitivity in real

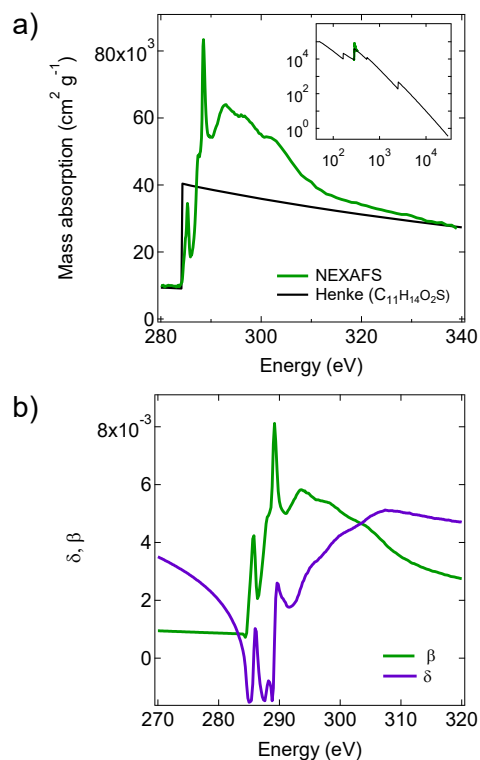


Figure 1: Measured NEXAFS data can be scaled and extended with the Henke database, as shown in (a). The inset shows a log-log plot to demonstrate how the Henke database is used to approximate values at energies outside of where NEXAFS data was collected. The real part of the complex index of refraction, δ , determined through a Kramers-Kronig relation, is plotted with the imaginary part, β , in (b). The example here is for a polythiophene, poly[3-(ethyl-5-pentanoate)thiophene-2,5-diyl].

space imaging of thin films via scanning transmission X-ray microscopy (STXM) [24–31].

NEXAFS spectra can be difficult to understand in detail, and complementary simulations are especially useful to help discern features such as contributions from specific atoms to certain peaks or intensity changes due to X-ray polarization and molecular orientation. Many recent studies utilize density functional theory (DFT) [32, 33] for predicting core-level transitions. Although DFT is limited to ground state properties and reliable predictions of excited states are more difficult [34], it has proven successful for simulations of X-ray absorption spectroscopy [35, 36]. The advantages of combining NEXAFS calculations with experiment have already shown success for polymer systems. Simulations based on a single polymer repeat unit can already be helpful in deconvoluting overall NEXAFS spectra into contributions from atoms at unique sites allowing for a better understanding of features such as relative peak intensi-

ties and origins of peak splitting [37, 38]. Calculations can aid in determining the overall and intrachain orientation in polymers. For instance chains in rub-aligned poly(tetrafluoroethylene) (PTFE) on gold tend to adopt a helical over a zig-zag conformation, and this is determined by comparing experimental results to simulations of model oligomers in both geometries [39]. However, the chain conformation of PTFE may vary depending on specific sample preparation conditions. Calculations that take into account different electric field polarization directions can reveal overall chain axis orientation, for example, along the rubbing direction [40]. NEXAFS is sensitive not only to structure and orientation, but also changes in chemistry. Chemical reactions, such as site specific bond scission and ionic fragmentation, can occur as a result of core excitation [41]. This allows NEXAFS to inherently be an *in situ* probe of soft X-ray induced chemical reactions and has been demonstrated for poly(methyl methacrylate) (PMMA) and poly-(isopropenyl acetate) (PiPAc) where a combination of experiment, DFT calculations and molecular dynamics help understand scission of the C-O bond and its relation to potential energy surfaces and bond elongation upon X-ray absorption [42].

Complementary NEXAFS predictions are also beneficial for functional polymers such as semiconducting [44] and ion conducting polymers. In semiconducting polymers electronic structure and molecular orientation are especially important for charge transport and NEXAFS can provide information on these properties. Additional complications need to be considered since the electronic wavefunctions in conjugated polymers are often delocalized over many repeat units [45]. Interpreting the NEXAFS profiles of complex molecular species is often based on the building block model, where a compound is broken down into local subunits (diatomic at the simplest level) to interpret overall spectra [46, 15, 47]. However, NEXAFS spectra of complex molecular species and polymers are often more complicated than the building block model predicts [48], especially in materials like semiconducting polymers with their delocalized π electrons [49, 50]. This highlights the need for X-ray absorption spectroscopy simulations for these materials. Only limited work has been done to predict NEXAFS of semiconducting polymers, but these studies demonstrate its potential and utility. For example, simulations have been able to help assign core-level transitions needed for orientation analysis [43, 20] (Figure 2), and help make connections between excited state delocalization and the formation of free charges versus excitons [51]. The combination of theory and experiment for X-ray absorption spectroscopy

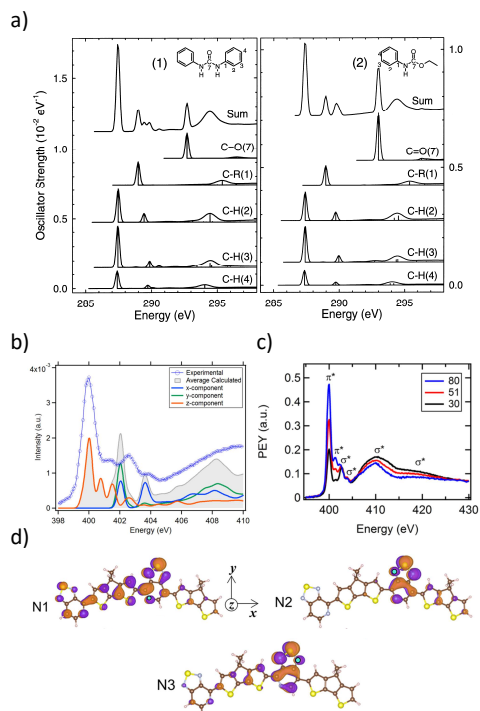


Figure 2: Simulated carbon K edge spectra of *N,N'*-diphenyl urea (1) and ethyl *N*-phenyl carbamate (2) are shown in (a). These calculations based on monomer units help understand transitions in the experimental data of the polymer. These were done using the GSCF3 package which is based on the improved virtual orbital approximation that takes into account the core-hole in the Hartree-Fock approximation. The overall simulated spectra is shown (Sum) as well as the contributions from each distinct carbon atom. Reprinted from Ref [38], with permission from Elsevier. An example of nitrogen K edge experiment and simulations of a pyridal[2,1,3]thiadiazole-containing conjugated polymer, PCDTPT, is shown in (b). Simulations were performed using the eXcited electron and Core Hole (XCH) approach [35]. Contributions from different electric field directions indicate that the first three peaks are π^* in character. This is supported by angle-dependent experiments of aligned PCDTPT shown in (c). The lowest energy electronic final state orbitals for the three distinct nitrogen atoms in PCDTPT are shown in (d), revealing the varying degree of delocalization possible in semiconducting polymers. Adapted with permission from Ref [43]. Copyright 2015 American Chemical Society.

in functional polymers is still in its early stages, and opportunities exist to improve fundamental understanding of how factors such as delocalized wavefunctions, charge transfer, and backbone geometry influence NEXAFS spectra. Examples of polymer-based applications that can take advantage of soft X-ray probes are presented in section 4.

2.3. Resonant reflectivity and scattering

RSoXS and RSoXR take advantage of the unique core-level absorption profiles of polymers to provide statistically averaged characteristic length scales related to many factors including but not limited to, phase separation or molecular orientation [52]. As discussed in section 2.1, there are important connections between the real (δ) and imaginary (β) parts of the complex index of refraction that need to be understood. At this fundamental level, RSoXR can be combined with NEXAFS, atomic scattering factors from the Henke database, and the Kramers-Kronig relation to accurately determine δ and the index of refraction of a polymer in a self-consistent manner without the need for separate mass thickness measurements [53]. The unique features in the index of refraction can be leveraged in RSoXR to gain chemical specificity by tuning the incident X-ray energy and probe orientation or composition through the depth of thin film. For example, RSoXR can characterize roughness or interfacial width of buried polymer interfaces in bilayer or block copolymer systems [54–56], which would be difficult to probe with hard X-ray reflectivity, and likely require chemical modification through deuteration for neutron based experiments. RSoXR is able to assist in characterizing phase separated block copolymers to, for example, help distinguish between vertical and parallel lamellae [57].

Many polymers, including conjugated or liquid crystalline polymers, have asymmetric properties and absorption peaks that correspond to transition dipole moments with well-defined directions that are localized to specific chemical moieties. The control of X-ray polarization that undulator insertion devices provide allows RSoXR an opportunity to probe molecular orientation through the depth of a film, and this has been shown for polymers [58] and organic ionic liquids [59]. However, these properties also make it more difficult to properly determine the optical constants from NEXAFS data, and X-ray reflectivity can prove useful as a reference and comparison for *ab initio* calculations of the dielectric tensor [58]. This has also been demonstrated for a self-assembled monolayer of 1,4-benzenedimethanethiol on gold, where the film’s anisotropic dielectric constants were determined based on *ab initio* DFT simulations

of the polarization dependent absorption cross section. Structural parameters such as the molecular tilt angle, monolayer thickness, and packing structure were determined by fitting to experimental RSoXR [60], as shown in Figure 3a and 3b.

Complementary to reflectivity, RSoXS has recently emerged as an effective characterization tool for a range of soft materials, including polymeric systems [61–64, 21, 52, 65–67]. RSoXS has seen popular use for thin films of conjugated polymer and complex organic semiconductor blends [68–72]. RSoXS has the ability to interrogate in-plane and out-of-plane structure since information from multiple scattering angles can be collected simultaneously on a two-dimensional detector and experiments can be done in transmission and grazing incidence geometries. Similar to RSoXR, RSoXS takes advantage of the energy and orientation dependent complex index of refraction near absorption edges to achieve high contrast and chemical sensitivity. Consequently, it is important to understand the complex index of refraction of the materials of interest, at least based on NEXAFS measurements. Even a basic understanding of the absorption differences among individual polymers or components in a blend is sufficient to distinguish differences in characteristic length scales, such as the core vs. outer layer in structured polymer nanoparticles [73] and distinct spacing of various phases in a cylinder forming triblock copolymer [64] or a mesoporous block copolymer film [74].

The transmission geometry and multiple scattering angles detectable in an RSoXS experiment enable unique aspects of molecular orientation to be probed. This allows RSoXS to go beyond RSoXR and investigate in-plane orientation with polarized incident X-rays. A unique feature in polarized scattering that is often seen, especially in semicrystalline polymers, is anisotropic two-dimensional scattering patterns (varying intensity as a function of azimuthal angle for a given Q), which have been attributed to molecular alignment at domain interfaces or fibrillar bulk ordering [75]. This anisotropic scattering occurs in samples that are globally isotropic in-plane and the direction of scattering intensity anisotropy changes with incident X-ray polarization direction, suggesting this feature arises due to core-level transitions with specific dipole moment directions and local molecular arrangement. Several cases of scattering anisotropy have been observed in blends of conjugated polymers and polymer-fullerene blends [68, 71, 70, 76, 72, 75], but it can also be observed in a single semicrystalline polymer, as shown in Figure 3. Depolarized diffuse scattering at large scattering vectors and energies near absorption edges reveals short-range

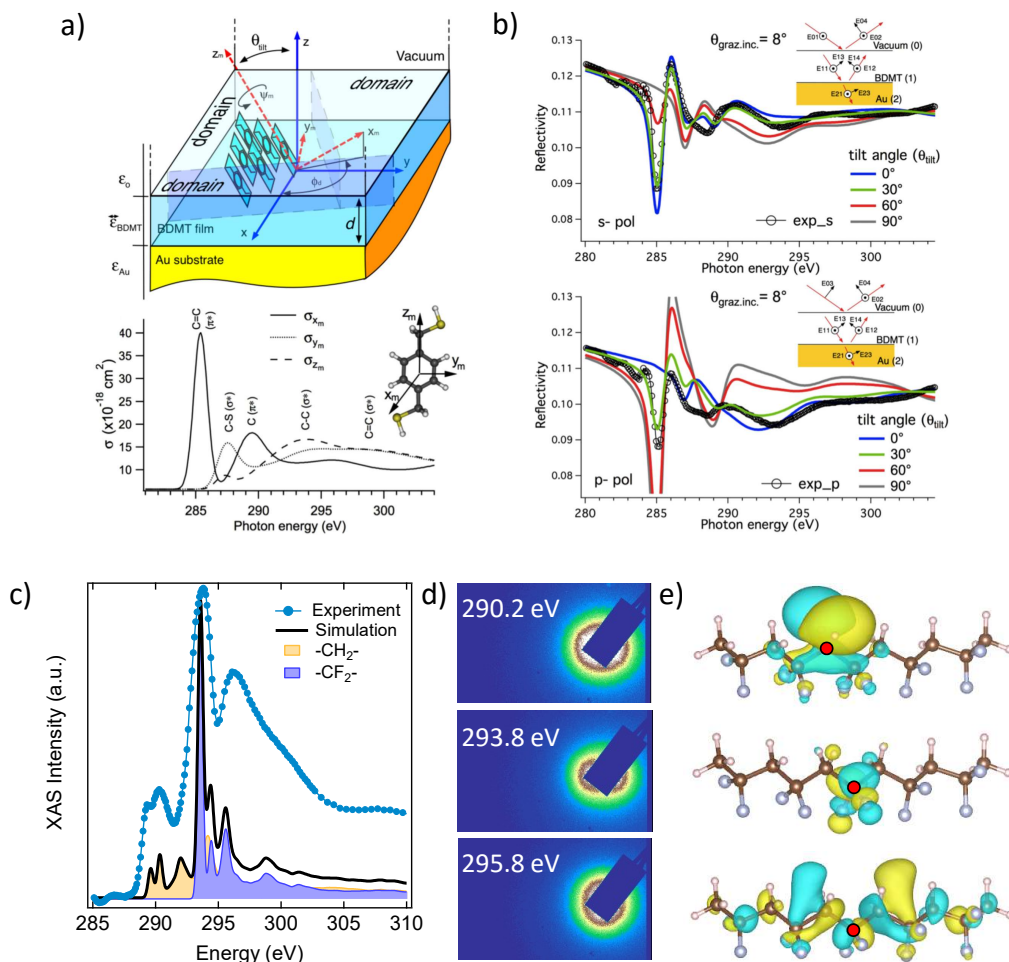


Figure 3: Polarized soft X-ray reflectivity and scattering are sensitive to incident X-ray energy, molecular orientation, and overall structure. A model of a self-assembled monolayer of 1,4-benzenedimethanethiol (BDMT) is shown in (a). Molecules were assumed to be standing up with planes parallel to each other and a tilt angle relative to the surface normal, θ_{tilt} . The anisotropic absorption cross section was calculated for a single molecule using DFT within the StoBe code [36]. In the lower panel of (a), the components of the absorption cross section, σ_{x_m} , σ_{y_m} , and σ_{z_m} , correspond to the electric field vector pointing along the x_m , y_m , and z_m axes of the molecule, respectively. Grazing incidence (8°) reflectivity as a function of energy is shown in (b). Experimental data for *s*-polarization (top panel) and *p*-polarization (bottom panel) are compared to simulations (solid traces) for a 1.5 nm thick BDMT film with varying molecular tilt angles. Reprinted with permission from Ref [60]. Copyright 2014 by the American Physical Society. Semicrystalline polymers often have anisotropic two-dimensional polarized scattering patterns. The experimental NEXAFS spectra of poly(vinylidene fluoride-co-trifluoroethylene) (PVDF-TrFE) is shown in (c). This is compared to XCH-calculated spectra of a PVDF oligomer, broken down into contributions from the two distinct carbon atoms. Experimental two-dimensional RSoXS patterns depicted in (d) reveal scattering anisotropy that changes with energy. Calculated final state orbitals at energies close to corresponding scattering patterns are shown in (e). This highlights how an understanding of NEXAFS transitions and transition dipole moments are connected to polarized RSoXS.

orientation fluctuations of functional groups, as shown for atactic polystyrene [77]. Anisotropic scattering patterns can reveal orientation in liquid crystalline materials, for example, the half-pitch in a bent core liquid crystal that forms a twisted smectic layer structure [78]. Furthermore, resonant scattering can uncover the spatial periodicity in a twist-bend liquid crystal phase that has no modulation in electron density [79].

A full understanding of the origins of this phenomenon is still underway [75], and challenges exist due to the many details present, ranging from an understanding of the core-level transitions and orientation dependence on the level of individual atoms and functional groups, to intermolecular interactions and local and global phase separation and structure. A complicated problem such as this highlights the need for theoretical developments to help understand subtle connections between spectroscopy and scattering, and push the capabilities of soft X-ray characterization forward.

3. Looking ahead: combining chemistry, morphology, and dynamics

The combination of X-ray absorption spectroscopy with resonant reflectivity and scattering undoubtedly can reveal critical information about chemistry and morphology in polymers and soft matter, not only in the bulk, but also in thin films and at interfaces. There is a need to understand this information *in situ* during relevant operating conditions and as a function of time to elucidate structural dynamics and chemical kinetics. This will require even more insights from theory in order to go beyond modeling results after data collection to predicting spectroscopy and scattering *a priori*.

3.1. Going beyond the carbon K edge

The majority of work done on XAS and RSoXS of polymers focuses on the carbon K edge since carbon is the most prevalent atomic species in these materials. There are often other elements present in polymers, for example, nitrogen, oxygen, fluorine, and sulfur, that can provide very useful complementary information. Moreover, since there tends to be fewer of these elements in distinct bonding environments compared to carbon, theoretical modeling and experimental data analysis can be more straightforward. For example, the nitrogen K edge can be used to quantify orientation in a planar conjugated polymer [43], as shown in Figure 2. In this case, nitrogen atoms are only present on the [1,2,5]thiadiazole[3,4-c]pyridine moiety in the backbone making NEXAFS simulations, peak assignment and orientation analysis less complicated.

In addition to elements with 1s absorption edges in the soft X-ray regime, many polymers contain elements with higher energy absorption K edges, for example phosphorus (2.145 keV) and sulfur (2.472 keV). Even though these K edges start to go beyond the energy range of many soft X-ray beamlines, there is much interest to increase capabilities in this “tender” X-ray (2 keV-7 keV) regime. Although typically less common due to the limited number of endstations that can perform XAS or RSoXS near the sulfur and phosphorus K edges, notable work has been done on sulfur edge XAS to probe molecular orientation of semiconducting oligothiophenes such as α -sexithiophene [80–82], polythiophenes [83–85, 20], and low-bandgap donor-acceptor polymers [84, 20, 86]. Assigning the transitions near the sulfur K edge is difficult not only because of the limited body of work published, but also due to challenges such as π^* and σ^* peaks being very close in energy. For example, thiophene based molecules exhibit $1s \rightarrow \pi^*$ and $1s \rightarrow \sigma^*(C-S)$ transitions that are very close in energy and hard to distinguish experimentally [80, 87, 88]. Complementary XAS simulations have seen only minimal use for the sulfur edge [86, 88], but the framework based on DFT used for other elemental edges, like carbon and nitrogen, works well at the sulfur K edge as well. Additionally, the L edge of sulfur (about 160 eV-190 eV) is in an energy range accessible at soft X-ray beamlines and offers further opportunities to understand electronic structure and molecular ordering [88, 89]. However, L edge calculations pose added challenges and is an area that would benefit from further developments. Predictions of L edge spectra require treatment via full relativistic spin-orbit coupling and accurate description of the relaxation of a core-hole electron, which is difficult for conventional plane-wave DFT methods [90].

Resonant scattering near the sulfur K edge has potential for providing complementary structural information of polymers and soft matter. It has proven useful in elucidating the internal structure of liquid crystals [91–93]. Resonant scattering at the sulfur K edge has rarely been used in polymers, and only proof of principles has been demonstrated [94]. There is room for growth in this area not only for experimental development, but also to expand our understanding of materials properties that may be revealed through theory and experiment in this regime of relevant X-ray energies.

3.2. Simulating resonant X-ray scattering patterns

Scattering can be more challenging to interpret because information about characteristic length scales related to molecular packing, orientation, and morphol-

ogy are also encoded in scattering patterns. Simulations of scattering patterns in the hard X-ray regime, where scattering contrast is dominated by differences in electron density, have become prevalent [66, 95–97], but very little has been done to predict soft X-ray scattering for polymer and soft material systems. The origins of anisotropic scattering seen with polarized incident X-rays has been explored for a rigid-rod type conjugated polymer system [75]. Gann et al. simulated two-dimensional scattering patterns for various morphologies for a blend of a rigid-rod polymer, such as poly(3-hexylthiophene) (P3HT), and an isotropic fullerene-like molecule (PCBM). The morphologies included spherical domains of one component in a matrix of the other, and vice versa, as well as a fibrillar polymer structure in an amorphous matrix. The orientation of the polymer chains relative to the fullerene domain at the polymer-fullerene interface was also varied. The simulations alone can reproduce similar anisotropic scattering patterns (Figure 4) and help understand experimental observations [75]. This demonstrates how scattering patterns can be predicted based on a known morphology, but determining the precise morphology of a sample with unknown structure remains much more challenging.

Scattering experiments performed in transmission geometry typically probe in-plane structure, but variation of incident angle in combination with the contrast enhancement near resonance can allow for accurate determination of depth-dependent morphologies in ordered systems. This has been demonstrated for self-assembled block copolymer thin films, where resonant critical-dimension small angle X-ray scattering combined with thorough data fitting reveal the specific domain morphology of vertical lamella through the film thickness [98, 99].

3.3. *In situ and time-resolved studies for dynamics and chemical kinetics*

Advancements in polymer and soft matter based applications and technologies often depend on an understanding of how chemistry and morphology evolve in response to specific processing conditions, for example, thermal annealing, strain, solvent evaporation, or electrical biasing. Although examining structure pre- and post-processing is critical, a full understanding of structure development and chemical changes requires advanced *in situ* or *in operando* characterization that can probe dynamics and chemical kinetics. *In situ* and time-resolved investigations are more frequently done with hard X-rays [66, 100–103] since experiments do not need to be done under vacuum at these energies. The

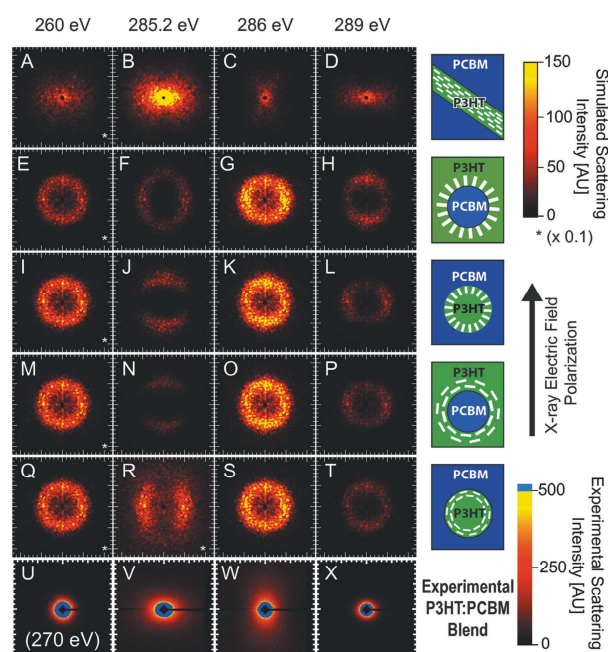


Figure 4: Simulated resonant scattering patterns at different energies (260 eV, 285.2 eV, 286 eV, and 289 eV) for various model morphologies. The morphologies shown are aligned P3HT fibrils in an amorphous PCBM matrix (A-D), P3HT molecules with cofacial alignment (face-on) (E-H) or chain axis perpendicular (edge-on) (M-P) to embedded PCBM domains, and embedded P3HT domains in a PCBM matrix with cofacial interfacial alignment (I-L) or perpendicular alignment (Q-T). Experimental scattering patterns of a spun-cast P3HT:PCBM blend are shown (U-X) at different X-ray energies (270 eV, 285.2 eV, 286 eV, and 289 eV). All of the two-dimensional scattering patterns shown are plotted as q_z (y-axis) vs. q_y (x-axis) ranging from -1 nm^{-1} to 1 nm^{-1} . Reprinted from Ref [75].

high vacuum chambers used for soft X-ray experiments makes *in situ* studies much more challenging, and further complications arise for scattering due to limitations on sample-detector distance.

The growing capability to quickly collect large data sets, especially for *in situ* experiments, creates a need for fast and even real-time data processing methods. This is actively being developed, especially for high throughput hard X-ray scattering beamlines. Packages such as *HipGISAXS* connect data collected at beamlines to supercomputers for on-the-fly analysis and modeling [104, 95, 105], opening new doors to increasing the efficiency of synchrotron experiments and understanding experimental results. Similar advancements are needed to complement soft X-ray techniques as high-throughput, *in situ*, and more complicated experiments become the norm.

3.3.1. *In situ* X-ray absorption spectroscopy

Simulations of core-level spectroscopy have already proven critical in elucidating changes in chemical bonding *in situ*. Metal-organic frameworks (MOFs) are coordination polymers with inherent pores from the cage-like structure that can be appended with certain chemistries, such as an amine functionality to selectively adsorb CO₂ to great capacity, and possibly serve as a catalytic center to derive useful chemicals from the greenhouse gas. MOFs have received immense research attention, and recent developments have focused on the *in situ* monitoring of gas sorption in this hybrid organic membranes from a structural and electronic point of view. Recent work on a diamine-appended MOF reveals a unique ‘phase-change’ that occurs upon adsorption of CO₂. This leads to large CO₂ separation capacities that can be attained with small temperature swings and low regeneration energies. This behavior is attributed to the formation of ammonium carbamate chains that form along the pores of the MOF due to adsorbed CO₂ molecules that insert into metal-nitrogen bonds [106]. This unique chemistry was determined based on X-ray spectroscopy measurements performed *in situ* at the nitrogen and oxygen K edges with varying amounts of CO₂ and compared to first-principles predictions on different adsorption structures. Theory verifies the specific bonding configuration needed to produce the observed changes in the NEXAFS spectra. In this case, the carbamate formed has a quasi-trigonal planar nitrogen and is bound to the metal through the oxygen atom [107], as shown in Figure 5. This underscores the need for predictive calculations to understand the subtle chemical and bonding environment changes that can be revealed by NEXAFS. The tunable energies in RSoXS make it

also intrinsically sensitive to chemistry in addition to characteristic length scales and *in situ* RSoXS studies are needed to reveal information not easily probed by hard X-ray scattering and for continued progress to understand the fundamental processes related to relevant applications.

3.3.2. Probing changes in interfacial structure and chemistry

Resonant reflectivity and scattering are useful tools for investigating interfacial phenomenon such as interfacial width in polymers and molecular orientation at domain interfaces as discussed in sections 2.3 and 3.2. These capabilities lend themselves to future work that understands interfacial processes under operating conditions or as a function of time, building off capabilities of *in situ* NEXAFS as discussed in section 3.3.1. For example, in a block copolymer, it is known that a neutral solvent that is equally soluble in both blocks reduces the non-favorable interactions between blocks A and B, lowering the Flory-Huggins interaction parameter, χ_{AB} . The interfacial width, a_I , is related to χ_{AB} by

$$a_I = \frac{2b}{\sqrt{6\chi_{AB}}}$$

where b is the segment length [108–110]. The energy tunability of RSoXS allows chemical sensitivity between distinct blocks, and this can be leveraged to selectively probe interfaces. The value of the invariant, Q , defined for an isotropic material as

$$Q = \frac{1}{2\pi^2} \int_0^\infty q^2 I(q) dq,$$

is reduced for diffuse phase boundaries relative to sharp interfaces [111]. Therefore, the invariant can be used to track changes in interfacial width as a function of various processing parameters. An example for cylinder forming polystyrene-*block*-poly(2-vinyl pyridine) (PS-*b*-P2VP) (Mn = 40-*b*-18 kg/mol) exposed to tetrahydrofuran (THF) (a nearly neutral solvent) to improve long-range order is shown in Figure 6, similar to systems studied previously [101]. This demonstrates the utility of using XAS and RSoXS to track not only interface localized changes in structure, but also changes in chemistry under *in situ* conditions.

3.3.3. Time-resolved dynamics with photon correlation spectroscopy

The time scales and dynamics of many polymers are dominated by complex rheological processes spanning length scales from nanometers to micrometers. In

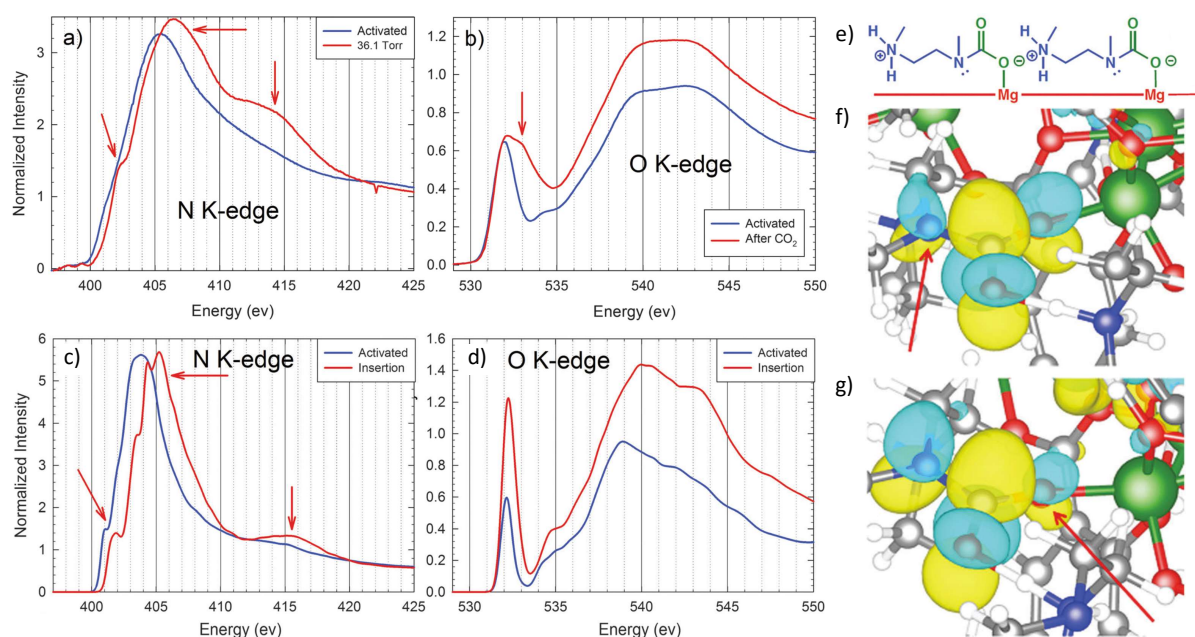


Figure 5: Experimental (a) nitrogen K edge and (b) oxygen K edge NEXAFS spectra of the MOF $\text{mmen-Mg}_2(\text{dobpdc})$ in vacuum (activated) and in the presence of 36.1 Torr CO_2 gas. Note the oxygen K edge spectra was taken after exposure to CO_2 . The main changes in the nitrogen K edge spectra (red arrows) upon CO_2 adsorption include a more pronounced pre-edge feature near 402.3 eV, an increase in intensity of the main peak near 405.4 eV and its blueshift by about 1 eV, and the appearance of a broad feature between 411 and 419 eV. The major change in the oxygen K edge spectra is an increase in overall intensity and a new π^* peak near 532.9 eV. All of these changes are captured in the DFT calculated spectra at 298 K (c,d) when the CO_2 insertion structure chemistry (e) is assumed. Electronic final state orbitals corresponding to the transition near (f) 402 eV and (g) the π^* peak in the oxygen K edge spectra is shown for the insertion structure. Red arrows indicate the excited atom. Similar electron density distributions in both cases suggest the electronic structure is similar for transitions from both nitrogen and oxygen excitations. Reproduced from Ref [107] with permission of The Royal Society of Chemistry.

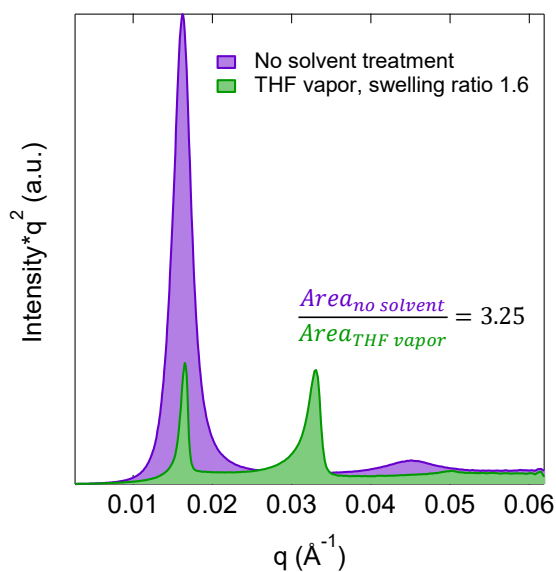


Figure 6: Circularly averaged transmission resonant X-ray scattering data of a PS-*b*-P2VP block copolymer taken at 285.2 eV, where good contrast exists due to the dominant absorption from PS. Data is plotted as Iq^2 vs. q and the area under the curve is related to the invariant, Q . The green trace represents a sample that was exposed to THF solvent vapor until a swelling ratio (thickness of swollen film divided by original thickness) of 1.6. The THF vapor leads to a reduction in χ_{AB} since THF is a nearly neutral solvent for PS and P2VP, and better long-range order is indicated by the sharper, multiple order peaks. However, the reduced χ_{AB} leads to more diffuse interfaces, revealed through resonant scattering by the reduction the area under the curve.

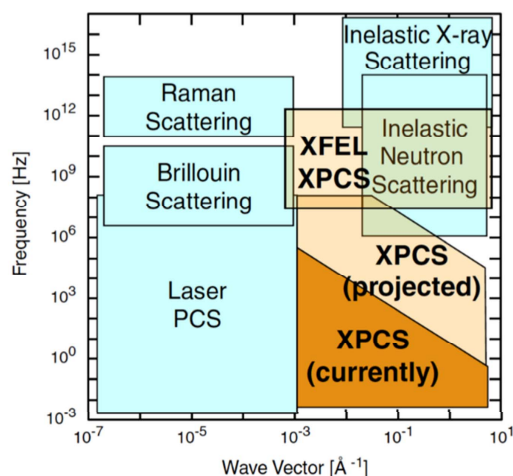


Figure 7: Plot showing the frequency-scattering wave vector domains accessible by various spectroscopic dynamic techniques. Special emphasis is placed on differentiating the current status of XPCS from the projected improvements upon completion of upgrades to synchrotron facilities. Taken from [112].

addition, the hierarchical internal structure inherent to many polymeric materials causes them to flow over long timescales, but respond elastically on much shorter timescales [112]. This behavioral crossover can span several orders of magnitude, thus calling for experimental techniques that are equally adept at accessing dynamic information over many scales. Recently, a coherent scattering technique known as photon correlation spectroscopy (PCS) has been increasingly applied to the X-ray photon regime and is an exceptionally well-suited tool to interrogate polymer dynamics over such scales. In short, PCS provides information on the dynamics of a given system by measuring intensity fluctuations within the scattering patterns that are produced by a coherent source. It is often subdivided into two categories based on the photon energy: dynamic light scattering (DLS), which uses visible light generated by a laser to monitor motion and measure particle sizes down to the sub micrometer-range; and X-ray photon correlation spectroscopy (XPCS), which uses X-rays to probe the dynamics of nanoscale inhomogeneities over a wide range of timescales (about 10^{-3} to 10^4 seconds) [113–116].

However, in order to understand the role that modeling and theory play in realizing the potential of XPCS to study polymer dynamics, it is important to review certain fundamental concepts such as the differences between coherent and incoherent scattering; the essential theory behind the autocorrelation analysis used to dissect the subtle fluctuations in the time-dependent scattering data; and the experimental conditions required to

produce a useful scattering plot. First, the scattering patterns from incoherent X-rays represents a statistical average over many incoherent electron density regions within the sample, which gives rise to an isotropically-broadened scattering peak. On the other hand, coherent X-ray experiments produce a scattering pattern with peaks composed of specific subregions with increased intensity, or speckles. These speckles are formed as a result of the mutual constructive interference of two X-ray wavefronts, which originate from a given scattering point. Thus, their inverse Fourier transform can give an electron density distribution with nanometer precision.

Snapshots of speckle patterns are taken at various time intervals, where their sequence reveals subtle intensity, I , fluctuations that can be analyzed by a second order autocorrelation function, g_2 , to extract correlation times, τ , characteristic to each wave vector, Q , according the following relation,

$$g_2(t) = \frac{I(\tau)I(\tau + t)}{I(\tau)^2} = 1 + A|F(Q, t)|^2,$$

where the time spacing of the speckle sequence, τ , determines the timescale of the dynamics probed. Finally, assuming a system with stationary dynamics around an equilibrium (*e.g.* Brownian motion), a plot can be obtained as shown in Figure 8 where the intermediate scattering function, $A |F(Q, t)|^2$, can be modeled to extract dynamic parameters.

Experimentally, the well-established autocorrelation theory described above can only be applied if these scattering events occur only once (limiting the thickness/density of the sample) and if the ergodic principle, which assumes that no information is lost within each snapshot time, is satisfied (limiting the short timescale resolution). Multiple scattering can be avoided by using a small micrometer sized sample (such as a single crystal). On the one hand, this can be very useful because it circumvents the requirement of a small slit to induce coherence since the sample itself will be the source. On the other hand, this approach severely limits the range of samples that can be used.

This divergence between the samples and length scales that are accessible are what distinguish the applicability of XPCS from DLS. Currently, DLS is used on a routine basis for the analysis of particle sizes in the sub micrometer range; it provides an estimation of the average size and its distribution within a measuring time of a few minutes. However, improvements in the ability to model dynamic behavior and relate it to scattering data has continued to make it a useful tool to study slow effects such as swelling mechanisms and phase transitions [117–127]. Still, XPCS is considered superior to

DLS in certain situations because 1) it enables high- q dynamics and 2) it can probe samples that are opaque or subject to multiple light scattering phenomena under DLS conditions. For example, the high q -range of XPCS can be used to extract elusive surface dynamics such as polymer chains and the relaxation rates of capillary waves on surfaces by directly probing a material in a grazing incidence geometry [128–130, 112, 131, 132].

There are also many samples exhibiting long-term relaxations of interest, like “jamming systems,” where the only way to distinguish slow sample movements from random fluctuations is by analyzing the slow drift of the speckle pattern. Such slow drifts cannot be observed in a DLS experiment due to the beam size, while the usual μm -scale beam used in XPCS can probe such flows. Therefore, several methods developed for light scattering experiments are now used with X-rays. One example is a method called heterodyning that can be used to obtain interferences between a reference and the sample. This provides correlation functions where heterodyning a sample’s diffracted amplitude against that of a reference can be used for phases retrieval, separate fluctuations from long-term flows [133], and mechanical relaxation from aggregate diffusion [134].

Another challenge is studying polymer dynamics in the μs to ms regime, which will require sources with greater coherent flux in order to obtain a useful signal to noise ratio. The recent push to improve coherence of many synchrotrons is an exciting development in this regard [135]. Access to this temporal regime will enable a deeper understanding of polymer systems away from equilibrium or near the onset of more complex dynamical heterogeneities. The detailed quantification of such phenomena is critical to elucidating the arrested states of systems. However, the autocorrelation analysis of the data produced by such transient states will require use of a more involved two-time normalized correlation function $C(q, t_1, t_2)$ with normalized variance and fourth-order (as opposed to traditional second order) autocorrelation functions. These complex autocorrelation quantities will result from transient speckle scattering patterns whose origin is likely to be chemical in nature. Therefore, the development of XPCS capabilities in this temporal regime will also rely heavily on complementary theoretical modeling of the absorptive and scattering properties of the polymer materials under investigation.

Future work must explore the dynamics of specific chemistries. This will help achieve a better understanding of structure-property relationships and selectively probe the dynamics of buried interfaces. Once again, this can be achieved by combining XPCS with spec-

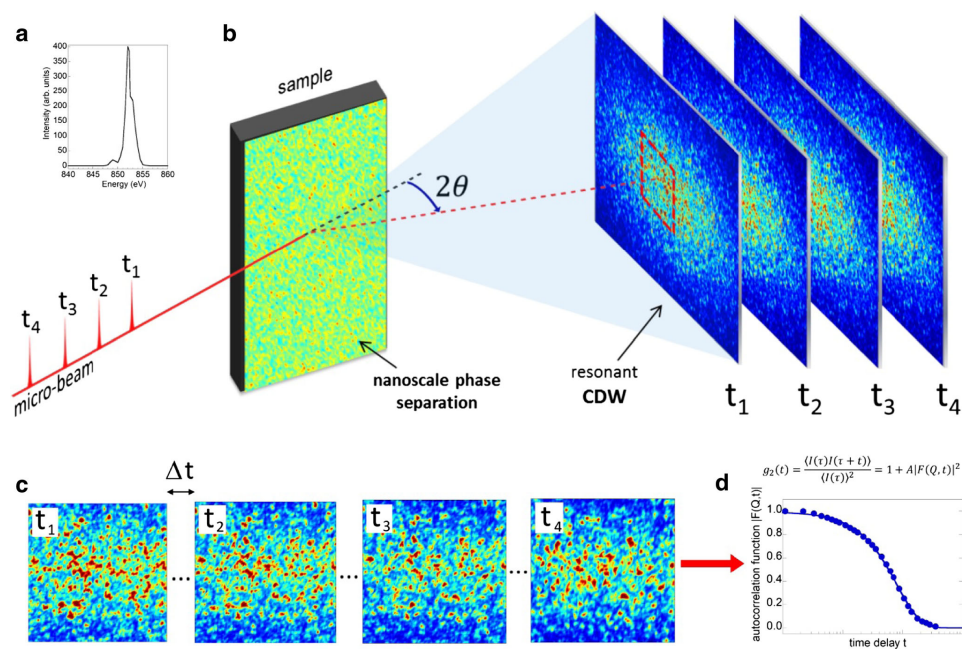


Figure 8: Resonant X-ray photon correlation spectroscopy experiment. a) The integrated intensity of resonant CDW peak in $\text{La}_{1.72}\text{Sr}_{0.28}\text{NiO}_4$ as a function of energy. The scattering intensity is strongly enhanced setting the energy to the Ni L_3 -edge. b) A sample showing nanoscale phase separation is illuminated by a coherent X-ray micro-beam. A time series of diffraction patterns is collected keeping external condition as stable as possible. c) The time-dependent coherent diffraction patterns are characterized by the presence of speckles that represent the granular local structure of the material showing an inhomogeneous landscape of charge-ordered domains in the illuminated spot. d) In order to determine the correlation time τ that characterizes the CDW domains dynamic, the intensity autocorrelation function can be calculated. This would allow to probe dynamics on timescales as short as the time spacing τ of the collected time series. Reproduced with permission from Springer [136].

troscopy. However, most of the resonance-enhanced XPCS research up to this point has been done using soft X-rays on solid inorganic materials (Fig 8) [137–141, 136]. However, despite the thorough background shown in the previous sections, resonance-enhanced XPCS still has not seen use in polymers.

Since scientific content in XPCS often derives from the deviation of the intensity correlation function from simple exponential line shapes, reliable and precise line shape theoretical analysis is fundamental to facilitating the use of theory to rationalize XPCS results. As synchrotron sources improve their coherent flux, the improved signal to noise ratios will enable researchers to reach even shorter timescales and probe more weakly scattering materials. However, even though this will enable unique insights into a new realm of material dynamics and *in-situ/operando* experiments, understanding of such studies will hinge even more upon the presence of a solid foundation linking scattering and spectroscopy theory with experiments.

4. Applications and outlook

The co-development of theoretical and experimental tools for soft X-ray absorption spectroscopy and scattering is focused on characterizing soft materials undergoing dynamic changes in relevant device and material application environments. In this section, two key areas of state-of-the-art polymer science that can greatly benefit from this combined theoretical and experimental approach of *in-situ* soft X-ray studies are highlighted: polymer materials for gas storage, ion conduction, and water filtration membranes and doped conjugated polymers for use in organic electronics.

4.1. Polymer-based membranes

For membrane materials in applications of fuel cells, batteries, CO₂ capture with hybrid organic frameworks, and water filtration, it is critical to not only analyze the morphology of the material designed and fabricated to assess its suitability for diffusion of ions, gas, and water through its pores, but also to develop understanding for how the material behaves structurally *in operando*. The use of soft X-ray spectroscopy and scattering enable the contrast-enhanced probing of electronic structure and molecular to mesoscale material structure in complex blends of polymers and organic small molecules. The low energy of soft X-rays is critical to generating enhanced contrast between similarly low-Z, carbonaceous components that only differ slightly in local chemical bonding environment. In

the simplest case of polymers of intrinsic microporosity (PIMs), the membrane is composed of a neat polymer, and the contrast is derived between the polymer and vacuum (which comprises the pores when under vacuum conditions for measurement). These porous polymers and their very small (<2 nm) intrinsic pores are ideal structures for gas storage and separation [5, 7]. Essentially, the neat polymer comprises the matrix and provides mechanical integrity to the membrane, while the ‘unoccupied’ volume that arises from imperfect packing of rigid polymer backbones serves as the ‘highways’ for transport of ions, gas, or water, depending on the chemistry of the polymer. One time-dependent study using small angle (SAXS) and wide angle X-ray scattering (WAXS) and molecular dynamics (MD) simulations probes the dynamics of aging in microporous polymers with pore size <1 nm [142]. Physical aging is relevant because the closing of pores via this mechanism can drastically reduce molecular transport. Runt et al. found that there are two distinct aging mechanisms in PIMs: decrease of pore volume fraction followed by pore shrinkage. In another strategy for creating pores in polymer films for mesoporous conducting and storage membranes, inherently templating polyethylene-*b*-polystyrene-*b*-polyethylene (SES) films are impregnated with hPS and rinsed with selective solvent to generate pores within the PS domains that can be further treated by sulfonation to enhance hydrophilicity and ion conduction [74, 143, 144]. The structural analysis that is key to ion conduction will probe the location, size, and fraction of each block component relative to the generated voids, and RSoXS has been shown by Balsara et al. to be the ideal tool for characterizing these pores in heterogeneous soft matter membranes. Without the contrast-matching resonance of RSoXS, it would not be possible to distinguish between blocks and determine the relative locations of the pores and hydrophobic and hydrophilic domains. Similarly, the same group has mixed TiO₂ nanoparticles into templated thin films of PS-*b*-PEO battery membranes in an attempt to block dendrite runaway growth during electrochemical cycling [145]. Using RSoXS and mechanical testing, it was found that there exists an ideal nanoparticle loading near 20 wt.% that preserves the block copolymer ordering and optimizes the membrane toughness. This type of study could be extended to *in operando* RSoXS where membranes are cycled electrochemically, and the changing morphology near the block copolymer domains and nanoparticle interfaces can be probed.

Nafion is a well-studied material with a complex microstructure that is suitable for fuel cell, water filtration, and solar-fuels generation membranes. It forms a very

heterogeneous structure in which long, parallel, but randomly packed inverted micellar cylinders are comprised of the backbone on the outside and ionic side chains on the inside, and water and ions are able to be transported efficiently through this hydrophilic cylindrical micelle core [146]. *In situ* grazing incidence SAXS (GISAXS) has been used to study the existence of cylindrical domains of Nafion and how they swell during water treatment. Segalman et al. point out the importance of the substrate interface on directing either in-plane only or random orientation of cylinders; a random mixture is better to maximize water, oxygen, and ion transport. Thermal annealing during GISAXS revealed that water uptake is restricted following extensive thermal processing due to the formation of Nafion crystals that begin to take away from the compressibility of the Nafion matrix that allows substantial water uptake. The authors are able to monitor the swelling of the Nafion pores with water in real time with GISAXS [147]. Similarly, *in-situ* SAXS on Nafion membranes in the bulk form shows that there exists an enormous difference in time scale of water uptake and swelling for membranes treated with liquid versus water vapor. From these studies, the *d*-spacing of the ‘ionomer’ peak is observed, which represents the spacing between randomly packed inverse micellar cylinders of water transporting Nafion [148]. Some new preliminary RSoXS results reveal that the polymer chains are partially oriented inside ionomer domains (inverse micelles) near the interface with adjacent crystallites, and that this orientational ordering occurs even before the crystalline phase begins to form [65]. Future studies should continue to use RSoXS to probe the orientational and ionomer domain ordering, but in a time-resolved *in situ* study in aqueous environments with various salt and pH treatments. Using a solution environment cell that is currently under development, such studies will help to determine the molecular details of ionomer domain formation and to visualize the swelling of these membranes during operation.

4.2. Doped semiconducting polymers

The doping of conjugated polymers is a topic that has attracted much interest, but unanswered questions remain regarding the relationship of dopant-polymer microstructure (dopant incorporation), ionization efficiency and resulting device enhancement by doping. The combination of X-ray absorption spectroscopy and resonant scattering have potential to probe changes in electronic structure and corresponding spatial distributions that occur upon doping and help understand charge transfer.

The use of first-principles calculations is especially needed to understand details of electronic structure in conjugated polymers and small molecules. Organic semiconductors can have distinct NEXAFS spectra due to unique electronic structure characteristics. For example, a biradical hydrocarbon that exists as a stable radical with an unpaired electron in a singly occupied molecular orbital reveals a very low energy NEXAFS peak (283.4 eV), attributed to its low-lying lowest unoccupied molecular orbital [149]. Comparing experimental spectra to simulations can help determine where on the polymer backbone (or small molecule) an excited state exists, and how the dopant may change the local chemical bonding environment [51]. There have been many examples of doping in conjugated polymers [150], using varied dopants from organic small molecules to graphene oxide [151] and ranging in application from organic thin film transistors to sensors to devices for the tissue-electrode interface in biology.

Efforts seek to understand how the dopant-host structure determines ionization processes and efficiency. For example, the electronic structure of polyaniline (PANi) was investigated by XAS at various doping levels by organic acids. It was found with this soft X-ray strategy that doping occurs by protonation of the backbone nitrogen sites of PANi and changes the local chemical bonding configuration by breakup of the quinoid structure and disappearance of imine bonding, leaving only amine bonding characteristics [152]. Cu and Fe K edge NEXAFS were utilized in a different doping strategy to study the *in situ* electrochemical inclusion of Cu and Fe species in poly(3-methylthiophene) (P3MT). It was found that the Cu ion is an effective dopant for P3MT through its complexation with the sulfur atom within the backbone, following the reversible, time-dependent Cu^{II} - Cu^I - Cu^0 reduction process [153]. Analogously, NEXAFS was combined with DFT to probe local conformation and electronic structure of P3HT doped with F4TCNQ, a common p-type dopant for organic semiconductors. It was found that the dopant is inserted into some crystallites of P3HT in a planar fashion such that cofacial stacking of the two species occurs [154]. Another study confirmed that only some crystallites contain the dopant, and that weak and strong doping regimes occur [150]. In a similar polymer system, PBTTT- C_{14} , in which the thiophene monomers of P3HT are spaced farther by a thienothiophene unit, *in situ* GIWAXS monitored the incorporation of the F4TCNQ dopant [155]. A similar cofacial arrangement was determined that preserves planarity of the polymer backbone within crystallites, but this time the polymer crystallites all contain the dopant, and in nonrandom, correlated

locations. This correlated, nonrandom arrangement of dopants throughout the semiconductor crystallites and film has strong implications for transport models due to the nature of charge hopping in these 1D-like semiconductors. The dopants in this study were also found to cause crystallite contraction in the π -stacking direction (b axis) and slight expansion and paracrystalline disorder in the alkyl-stacking direction (a axis). This well-thought study should be extended to other polymers and dopants, for example using *in situ* ion gating where doping reversibility and effect of electrochemistry can be investigated.

Challenges exist since a theoretical basis to predict changes in X-ray absorption spectra for molecules that have donated or accepted an electron is still not well developed. This is needed to move forward in understanding the spatial distribution not only of the dopant [150], but also of electrically doped sites within an organic semiconductor, which could be probed through resonant scattering, even during device operating conditions.

5. Conclusions

X-ray absorption spectroscopy and resonant scattering and reflectivity are useful tools for combined chemical and spatial analysis of polymers and soft materials. Along with the benefits that these techniques provide, challenges also exist that must be overcome, including an improved understanding at the atomic level of electronic transitions, polarization and orientation dependent absorption and scattering properties, and an ability to predict resonant scattering patterns of intricate systems. As more complex material systems and advanced *in situ* and *in operando* experiments become essential to improve understanding related to various applications, addressing these challenges becomes more necessary to take full advantage of the capabilities of X-ray spectroscopy and resonant scattering. In the future, complementary calculations will be more closely integrated to experiments to predict and analyze data and accelerate materials development.

With this combination of theory and soft X-ray spectroscopy and scattering, soft X-rays will prove to be an essential part of the analytical toolkit for polymer scientists, especially with the ability to conduct time-resolved, *in situ* XPCS experiments that track distinct scattering centers under relevant sample conditions. Two key application areas of soft matter, doping of conjugated polymers and polymer membrane structure, were discussed both in terms of the morphological and electronic structure understanding developed thus

far, and the potential power soft X-ray spectroscopy and scattering have in determining a more complete picture for these materials. Many other systems exist where the combination of soft X-rays and electronic structure modeling will be essential in studying the details of dynamic structure of soft matter under operating conditions.

6. Acknowledgments

This work was supported by the Advanced Light Source (ALS) and the Molecular Foundry, user facilities located at the Lawrence Berkeley National Laboratory and supported by the Director, Office of Science, Office of Basic Energy Sciences, of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231. The authors thank Xiaodan Gu for helpful discussions.

References

- [1] Sirringhaus, H.. 25th anniversary article: Organic field-effect transistors: The path beyond amorphous silicon. *Adv Mater* 2014;26:1319–1335. doi:10.1002/adma.201304346.
- [2] Dou, L., You, J., Hong, Z., Xu, Z., Li, G., Street, R.A., et al. 25th anniversary article: A decade of organic/polymeric photovoltaic research. *Adv Mater* 2013;25(46):6642–6671. doi:10.1002/adma.201302563.
- [3] Holliday, S., Donaghey, J.E., McCulloch, I. Advances in charge carrier mobilities of semiconducting polymers used in organic transistors. *Chem Mater* 2014;26(1):647–663. doi:10.1021/cm402421p.
- [4] Huang, Y., Kramer, E.J., Heeger, A.J., Bazan, G.C.. Bulk heterojunction solar cells: Morphology and performance relationships. *Chem Rev* 2014;114(14):7006–7043. doi:10.1021/cr400353v.
- [5] Xu, C., Hedin, N.. Microporous adsorbents for {CO₂} capture - a case for microporous polymers? *Mater Today* 2014;17(8):397–403. doi:10.1016/j.mattod.2014.05.007.
- [6] Sreenivasulu, B., Sreedhar, I., Suresh, P., Raghavan, K.V.. Development trends in porous adsorbents for carbon capture. *Environ Sci Technol* 2015;49(21):12641–12661.
- [7] Du, N., Park, H.B., Robertson, G.P., Dal-Cin, M.M., Visser, T., Scoles, L., et al. Polymer nanosieve membranes for CO₂-capture applications. *Nature Mater* 2011;10:372–375. doi:10.1038/nmat2989.
- [8] Zhang, H., Shen, P.K.. Recent development of polymer electrolyte membranes for fuel cells. *Chem Rev* 2012;112(5):2780–2832. doi:10.1021/cr200035s.
- [9] Jr., D.T.H., Balsara, N.P. Polymer electrolytes. *Annu Rev Mater Res* 2013;43(1):503–525. doi:10.1146/annurev-matsci-071312-121705.
- [10] Young, W.S., Kuan, W.F., Epps, T.H.. Block copolymer electrolytes for rechargeable lithium batteries. *J Polym Sci B Polym Phys* 2014;52(1):1–16. doi:10.1002/polb.23404.
- [11] Nie, Z., Kumacheva, E.. Patterning surface with functional polymers. *Nat Mater* 2008;7:277–290. doi:10.1038/nmat2109.
- [12] Stoykovich, M.P., Müller, M., Kim, S.O., Solak, H.H., Edwards, E.W., de Pablo, J.J., et al. Directed assembly of block copolymer blends into nonregular device-oriented structures. *Science* 2005;308(5727):1442–1446. doi:10.1126/science.1111041.

- [13] Ruiz, R., Kang, H., Detcheverry, F.A., Dobisz, E., Kercher, D.S., Albrecht, T.R., et al. Density multiplication and improved lithography by directed block copolymer assembly. *Science* 2008;321(5891):936–939.
- [14] Henke, B., Gullikson, E., Davis, J. X-ray interactions: Photoabsorption, scattering, transmission, and reflection at $E = 50\text{--}30,000$ eV, $Z = 1\text{--}92$. *At Data Nucl Data Tables* 1993;54(2):181 – 342. doi:10.1006/adnd.1993.1013.
- [15] Outka, D.A., Stöhr, J., Rabe, J.P., Swalen, J.D.. The orientation of langmuir-blodgett monolayers using nexafs. *J Chem Phys* 1988;88(6):4076–4087. doi:10.1063/1.453862.
- [16] Krishnan, S., Ayothi, R., Hexemer, A., Finlay, J.A., Sohn, K.E., Perry, R., et al. Anti-biofouling properties of comb-like block copolymers with amphiphilic side chains. *Langmuir* 2006;22(11):5075–5086. doi:10.1021/la052978l.
- [17] Dimitriou, M.D., Zhou, Z., Yoo, H.S., Killops, K.L., Finlay, J.A., Cone, G., et al. A general approach to controlling the surface composition of poly(ethylene oxide)-based block copolymers for antifouling coatings. *Langmuir* 2011;27(22):13762–13772. doi:10.1021/la202509m.
- [18] van Zoelen, W., Zuckermann, R.N., Segalman, R.A.. Tunable surface properties from sequence-specific polypeptoid polystyrene block copolymer thin films. *Macromolecules* 2012;45(17):7072–7082. doi:10.1021/ma3009806.
- [19] DeLongchamp, D.M., Kline, R.J., Fischer, D.A., Richter, L.J., Toney, M.F. Molecular characterization of organic electronic films. *Adv Mater* 2011;23(3):319–337. doi:10.1002/adma.201001760.
- [20] Aygül, U., Peisert, H., Batchelor, D., Dettinger, U., Ivanovic, M., Tournebise, A., et al. Molecular orientation in polymer/fullerene blend films and the influence of annealing. *Solar Energy Materials & Solar Cells* 2014;128:119 – 125. doi:10.1016/j.solmat.2014.05.017.
- [21] McNeill, C.R., Ade, H.. Soft x-ray characterisation of organic semiconductor films. *J Mater Chem C* 2013;1:187–201. doi:10.1039/C2TC00001F.
- [22] Schuettfort, T., Thomsen, L., McNeill, C.R.. Observation of a distinct surface molecular orientation in films of a high mobility conjugated polymer. *J Am Chem Soc* 2013;135(3):1092–1101. doi:10.1021/ja310240q.
- [23] Gann, E., McNeill, C.R., Szumilo, M., Siringhaus, H., Sommer, M., Maniam, S., et al. Near-edge x-ray absorption fine-structure spectroscopy of naphthalene diimide-thiophene co-polymers. *J Chem Phys* 2014;140(16):164710. doi:10.1063/1.4871463.
- [24] McNeill, C.R., Watts, B., Thomsen, L., Belcher, W.J., Greenhamand, N.C., Dastoor, P.C.. Nanoscale quantitative chemical mapping of conjugated polymer blends. *Nano Lett* 2006;6(6):1202–1206. doi:10.1021/nl060583w.
- [25] McNeill, C.R., Asadi, K., Watts, B., Blom, P.W.M., de Leeuw, D.M.. Structure of phase-separated ferroelectric/semiconducting polymer blends for organic non-volatile memories. *Small* 2010;6(4):508–512. doi:10.1002/sml.200901719.
- [26] Watts, B., Ade, H.. NEXAFS imaging of synthetic organic materials. *Materials Today* 2012;15(4):148 – 157. doi:10.1016/S1369-7021(12)70068-8.
- [27] Ade, H., Zhang, X., Cameron, S., Costello, C., Kirz, J., Williams, S.. Chemical contrast in x-ray microscopy and spatially resolved XANES spectroscopy of organic specimens. *Science* 1992;258:972–975. doi:10.1126/science.1439809.
- [28] Ade, H., Smith, A., Cameron, S., Cieslinski, R., Mitchell, G., Hsiao, B., et al. X-ray microscopy in polymer science: prospects of a new imaging technique. *Polymer* 1995;36(9):1843 – 1848. doi:10.1016/0032-3861(95)90930-Z.
- [29] Ade, H., Hsiao, B.. X-ray linear dichroism microscopy. *Science* 1993;262:1427 – 1429. doi:10.1126/science.262.5138.1427.
- [30] Watts, B., McNeill, C.R.. Simultaneous surface and bulk imaging of polymer blends with x-ray spectromicroscopy. *Macromol Rapid Commun* 2010;31(19):1706–1712. doi:10.1002/marc.201000269.
- [31] Ade, H., Stoll, H.. Near-edge X-ray absorption fine-structure microscopy of organic and magnetic materials. *Nature Mater* 2009;8:281–290. doi:10.1038/nmat2399.
- [32] Hohenberg, P., Kohn, W.. Inhomogeneous electron gas. *Phys Rev* 1964;136:B864. doi:http://dx.doi.org/10.1103/PhysRev.136.B864.
- [33] Kohn, W., Sham, L.J.. Self-consistent equations including exchange and correlation effects. *Phys Rev* 1965;140:A1133. doi:10.1103/PhysRev.140.A1133.
- [34] Elliott, P., Furche, F., Burke, K.. *Excited States from Time-Dependent Density Functional Theory*; chap. 3. John Wiley & Sons, Inc. ISBN 9780470399545; 2009, p. 91–165. doi:10.1002/9780470399545.ch3.
- [35] Prendergast, D., Galli, G.. X-ray absorption spectra of water from first principles calculations. *Phys Rev Lett* 2006;96:215502. doi:10.1103/PhysRevLett.96.215502.
- [36] Hermann, K., Pettersson, L.G.M., Casida, M.E., Daul, C., Goursot, A., Koester, A., et al. Stobe-demon version 3.3. 2014.
- [37] Magnuson, M., Yang, L., Guo, J.H., Sâthe, C., Agui, A., Nordgren, J., et al. The electronic structure of poly(pyridine-2,5-diyl) investigated by soft X-ray absorption and emission spectroscopies. *Chem Phys* 1998;237(3):295 – 304. doi:10.1016/S0301-0104(98)00262-6.
- [38] Urquhart, S.G., Hitchcock, A.P., Smith, A.P., Ade, H.W., Lidy, W., Rightor, E.G., et al. NEXAFS spectromicroscopy of polymers: overview and quantitative analysis of polyurethane polymers. *J Electron Spectrosc Relat Phenom* 1999;100:119 – 135. doi:10.1016/S0368-2048(99)00043-2.
- [39] Gamble, L.J., Ravel, B., Fischer, D.A., Castner, D.G.. Surface structure and orientation of PTFE films determined by experimental and FEFF8-calculated NEXAFS spectra. *Langmuir* 2002;18(6):2183–2189. doi:10.1021/la011258l.
- [40] Ågren, H., Carravetta, V., Vahtas, O., Pettersson, L.G.M.. Orientational probing of polymeric thin films by nexafs: Calculations on polytetrafluoroethylene. *Phys Rev B* 1995;51:17848–17855. doi:http://dx.doi.org/10.1103/PhysRevB.51.17848.
- [41] Eberhardt, W., Sham, T.K., Carr, R., Krummacher, S., Strongin, M., Weng, S.L., et al. Site-specific fragmentation of small molecules following soft-x-ray excitation. *Phys Rev Lett* 1983;50:1038–1041. doi:10.1103/PhysRevLett.50.1038.
- [42] Takahashi, O., Tabayashi, K., Wada, S.i., Sumii, R., Tanaka, K., Odelius, M., et al. Theoretical study of ion desorption from poly-(methyl methacrylate) and poly-(isopropenyl acetate) thin films through core excitation. *J Chem Phys* 2006;124(12):124901. doi:10.1063/1.2176605.
- [43] Patel, S.N., Su, G.M., Luo, C., Wang, M., Perez, L.A., Fischer, D.A., et al. NEXAFS spectroscopy reveals the molecular orientation in blade-coated pyridal[2,1,3]thiadiazole-containing conjugated polymer thin films. *Macromolecules* 2015;48(18):6606–6616. doi:10.1021/acs.macromol.5b01647.
- [44] Nahid, M.M., Gann, E., Thomsen, L., McNeill, C.R.. NEXAFS spectroscopy of conjugated polymers. *Eur Polym J* 2016;doi:10.1016/j.eurpolymj.2016.01.017.
- [45] Brédas, J.L., Beljonne, D., Cornil, V.C.J.. Charge-transfer and energy-transfer processes in π -conjugated oligomers and polymers: A molecular picture. *Chem Rev*

- 2004;104(11):4971–5004. doi:10.1021/cr040084k.
- [46] Stöhr, J.. NEXAFS Spectroscopy. Springer Series in Surface Sciences; Springer; 1992. ISBN 9783540544227.
- [47] Outka, D.A., Stöhr, J., Rabe, J., Swalen, J.D., Rotermund, H.H.. Orientation of arachidate chains in langmuir-blodgett monolayers on si(111). *Phys Rev Lett* 1987;59:1321–1324. doi:10.1103/PhysRevLett.59.1321.
- [48] Pettersson, L.G.M., Ågren, H., Schürmann, B.L., Lippitz, A., Unger, W.E.S.. Assembly and decomposition of building blocks to analyze polymer nexafs spectra. *Int J Quantum Chem* 1997;63(3):749–765.
- [49] Carravetta, V., Ågren, H., Pettersson, L.G.M., Vahtras, O.. Near-edge core photoabsorption in polyenes. *J Chem Phys* 1995;102(14):5589–5597. doi:10.1063/1.469290.
- [50] Watts, B., Swaraj, S., Nordlund, D., Lüning, J., Ade, H.. Calibrated nexafs spectra of common conjugated polymers. *J Chem Phys* 2011;134(2):024702. doi:10.1063/1.3506636.
- [51] Gliboff, M., Sulas, D., Nordlund, D., deQuilettes, D.W., Nguyen, P.D., Seidler, G.T., et al. Direct measurement of acceptor group localization on donor-acceptor polymers using resonant auger spectroscopy. *J Phys Chem C* 2014;118(10):5570–5578. doi:10.1021/jp412150j.
- [52] Ade, H.. Characterization of organic thin films with resonant soft x-ray scattering and reflectivity near the carbon and fluorine absorption edges. *Eur Phys J Special Topics* 2012;208(1):305–318. doi:10.1140/epjst/e2012-01626-y.
- [53] Yan, H., Wang, C., McCarn, A.R., Ade, H.. Accurate and facile determination of the index of refraction of organic thin films near the carbon absorption edge. *Phys Rev Lett* 2013;110:177401. doi:10.1103/PhysRevLett.110.177401.
- [54] Wang, C., Araki, T., Watts, B., Harton, S., Koga, T., Basu, S., et al. Resonant soft x-ray reflectivity of organic thin films. *J Vac Sci Technol A* 2007;25:575–586. doi:10.1116/1.2731352.
- [55] Yan, H., Wang, C., Garcia, A., Swaraj, S., Gu, Z., McNeill, C.R., et al. Interfaces in organic devices studied with resonant soft x-ray reflectivity. *J Appl Phys* 2011;110(10):102220. doi:10.1063/1.3661991.
- [56] Sunday, D.F., Kline, R.J.. Reducing block copolymer interfacial widths through polymer additives. *Macromolecules* 2015;48(3):679–686. doi:10.1021/ma502015u.
- [57] Welch, C.F., Hjelm, R.P., Mang, J.T., Hawley, M.E., Wroblewski, D.A., Bruce Orlor, E., et al. Resonant soft x-ray scattering and reflectivity study of the phase-separated structure of thin poly(styrene-*b*-methyl methacrylate) films. *J Polym Sci B Polym Phys* 2013;51(2):149–157. doi:10.1002/polb.23190.
- [58] Mezger, M., Jérôme, B., Kortright, J.B., Valvidares, M., Gullikson, E.M., Giglia, A., et al. Molecular orientation in soft matter thin films studied by resonant soft x-ray reflectivity. *Phys Rev B* 2011;83:155406. doi:10.1103/PhysRevB.83.155406.
- [59] Mezger, M., Ocko, B.M., Reichert, H., Deutsch, M.. Surface layering and melting in an ionic liquid studied by resonant soft x-ray reflectivity. *Proc Natl Acad Sci* 2013;110(10):3733–3737. doi:10.1073/pnas.1211749110.
- [60] Pasquali, L., Mukherjee, S., Terzi, F., Giglia, A., Mahne, N., Koshmak, K., et al. Structural and electronic properties of anisotropic ultrathin organic films from dichroic resonant soft x-ray reflectivity. *Phys Rev B* 2014;89:045401. doi:10.1103/PhysRevB.89.045401.
- [61] Ade, H., Hitchcock, A.P.. Nexafs microscopy and resonant scattering: Composition and orientation probed in real and reciprocal space. *Polymer* 2008;49(3):643 – 675. doi:10.1016/j.polymer.2007.10.030.
- [62] Fink, J., Schierle, E., Weschke, E., Geck, J.. Resonant elastic soft x-ray scattering. *Rep Prog Phys* 2013;76(5):056502. doi:10.1088/0034-4885/76/5/056502.
- [63] Virgili, J.M., Tao, Y., Kortright, J.B., Balsara, N.P., Segalman, R.A.. Analysis of order formation in block copolymer thin films using resonant soft x-ray scattering. *Macromolecules* 2007;40(6):2092–2099. doi:10.1021/ma061734k.
- [64] Wang, C., Lee, D.H., Hexemer, A., Kim, M.I., Zhao, W., Hasegawa, H., et al. Defining the nanostructured morphology of triblock copolymers using resonant soft x-ray scattering. *Nano Lett* 2011;11(9):3906–3911.
- [65] Liu, F., Brady, M.A., Wang, C.. Resonant soft x-ray scattering for polymer materials. *Eur Polym J* 2016;doi:10.1016/j.eurpolymj.2016.04.014.
- [66] Hexemer, A., Müller-Buschbaum, P.. Advanced grazing-incidence techniques for modern soft-matter materials analysis. *IUCrJ* 2015;2:106–125. doi:10.1107/S2052252514024178.
- [67] Gleeson, H.F., Hirst, L.S.. Resonant x-ray scattering: A tool for structure elucidation in liquid crystals. *ChemPhysChem* 2006;7:321–328. doi:10.1002/cphc.200500426.
- [68] Collins, B.A., Cochran, J.E., Yan, H., Gann, E., Hub, C., Fink, R., et al. Polarized x-ray scattering reveals non-crystalline orientational ordering in organic films. *Nat Mater* 2012;11:536–543. doi:10.1038/nmat3310.
- [69] Collins, B.A., Li, Z., Tumbleston, J.R., Gann, E., McNeill, C.R., Ade, H.. Absolute measurement of domain composition and nanoscale size distribution explains performance in PTB7:PC71BM solar cells. *Adv Energy Mater* 2013;3(1):65–74. doi:10.1002/aenm.201200377.
- [70] Ma, W., Tumbleston, J.R., Wang, M., Gann, E., Huang, F., Ade, H.. Domain purity, miscibility, and molecular orientation at donor/acceptor interfaces in high performance organic solar cells: Paths to further improvement. *Adv Energy Mater* 2013;3(7):864–872. doi:10.1002/aenm.201200912.
- [71] Tumbleston, J.R., Collins, B.A., Yang, L., Stuart, A.C., Gann, E., Ma, W., et al. The influence of molecular orientation on organic bulk heterojunction solar cells. *Nat Photonics* 2014;8:385–391. doi:10.1038/nphoton.2014.55.
- [72] Diao, Y., Zhou, Y., Kurosawa, T., Shaw, L., Wang, C., Park, S., et al. Flow-enhanced solution printing of all-polymer solar cells. *Nature Commun* 2015;6(7955). doi:10.1038/ncomms8955.
- [73] Araki, T., Ade, H., Stubbs, J.M., Sundberg, D.C., Mitchell, G.E., Kortright, J.B., et al. Resonant soft x-ray scattering from structured polymer nanoparticles. *Appl Phys Lett* 2006;89(12):124106. doi:10.1063/1.2356306.
- [74] Wong, D.T., Wang, C., Beers, K.M., Kortright, J.B., Balsara, N.P.. Mesoporous block copolymer morphology studied by contrast-matched resonant soft x-ray scattering. *Macromolecules* 2012;45(22):9188–9195. doi:10.1021/ma3019206.
- [75] Gann, E., Collins, B.A., Tang, M., Tumbleston, J.R., Mukherjee, S., Ade, H.. Origins of polarization-dependent anisotropic x-ray scattering from organic thin films. *J Synchrotron Rad* 2016;23(1):219–227. doi:10.1107/S1600577515019074.
- [76] Liu, F., Wang, C., Baral, J.K., Zhang, L., Watkins, J.J., Briseno, A.L., et al. Relating chemical structure to device performance via morphology control in diketopyrrolopyrrole-based low band gap polymers. *J Am Chem Soc* 2013;135(51):19248–19259. doi:10.1021/ja408923y.
- [77] Stone, K.H., Kortright, J.B.. Molecular anisotropy effects in carbon *k*-edge scattering: Depolarized diffuse scattering and optical anisotropy. *Phys Rev B* 2014;90:104201. doi:10.1103/PhysRevB.90.104201.
- [78] Zhu, C., Wang, C., Young, A., Liu, F., Gunkel, I., Chen, D., et al. Probing and controlling liquid crystal he-

- lical nanofilaments. *Nano Lett* 2015;15(5):3420–3424. doi: 10.1021/acs.nanolett.5b00760.
- [79] Zhu, C., Tuchband, M.R., Young, A., Shuai, M., Scarbrough, A., Walba, D.M., et al. Resonant carbon *k*-edge soft x-ray scattering from lattice-free heliconical molecular ordering: Soft dilative elasticity of the twist-bend liquid crystal phase. *Phys Rev Lett* 2016;116:147803. doi: 10.1103/PhysRevLett.116.147803.
- [80] Kiguchi, M., Yoshikawa, G., Saiki, K.. Temperature and thickness dependence of molecular orientation of α -sexithienyl on Cu(111). *J Appl Phys* 2003;94(8):4866–4870. doi: 10.1063/1.1609637.
- [81] Kiguchi, M., Entani, S., Saiki, K., Yoshikawa, G.. One-dimensional ordered structure of α -sexithienyl on Cu(110). *Appl Phys Lett* 2004;84(18):3444–3446. doi: 10.1063/1.1736315.
- [82] Yoshikawa, G., Kiguchi, M., Ikeda, S., Saiki, K.. Molecular orientations and adsorption structures of α -sexithienyl thin films grown on Ag(1 1 0) and Ag(1 1 1) surfaces. *Surf Sci* 2004;559:77–84. doi:10.1016/j.susc.2004.04.045.
- [83] Onoki, R., Yoshikawa, G., Tsuruma, Y., Ikeda, S., Saiki, K., Ueno, K.. Nanotransfer of the Polythiophene Molecular Alignment onto the Step-Bunched Vicinal Si(111) Substrate. *Langmuir* 2008;24(20):11605–11610. doi:10.1021/la8016722.
- [84] Aygül, U., Batchelor, D., Dettinger, U., Yilmaz, S., Allard, S., Scherf, U., et al. Molecular orientation in polymer films for organic solar cells studied by nexafs. *J Phys Chem C* 2012;116(7):4870–4874. doi:10.1021/jp205653n.
- [85] Ikeura-Sekiguchi, H., Sekiguchi, T.. Molecular ordering effect of regioregular poly(3-hexylthiophene) using sulfur *k*-edge x-ray absorption spectroscopy. *Jpn J Appl Phys* 2014;53:02BB07. doi:10.7567/JJAP.53.02BB07.
- [86] Batchelor, D., Aygül, U., Dettinger, U., Ivanovic, M., Tournebise, A., Mangold, S., et al. Insight into the orientation of LBG polymer films by XANES experiment and calculation. *Eur Polym J* 2016;doi:10.1016/j.eurpolymj.2016.04.005.
- [87] Arantes, C., Borges, B.G.A.L., Beck, B., Araújo, G., Roman, L.S., Rocco, M.L.M.. Femtosecond Electron Delocalization in Poly(thiophene) Probed by Resonant Auger Spectroscopy. *J Phys Chem C* 2013;117(16):8208–8213. doi: 10.1021/jp312660d.
- [88] Hitchcock, A.P., Horsley, J.A., Stöhr, J.. Inner shell excitation of thiophene and thiolane: Gas, solid, and monolayer states. *J Chem Phys* 1986;85(9):4835–4848. doi:10.1063/1.451718.
- [89] Skotheim, T., Yang, X., Chen, J., Hale, P., Inagaki, T., Samuelson, L., et al. Highly ordered thin films of polyheterocycles: A synchrotron radiation study of polypyrrrole and polythiophene Langmuir-Blodgett films. *Synt Met* 1989;28(1):229–236. doi:10.1016/0379-6779(89)90526-2.
- [90] Fransson, T., Burdakova, D., Norman, P. *K*- and *l*-edge x-ray absorption spectrum calculations of closed-shell carbon, silicon, germanium, and sulfur compounds using damped four-component density functional response theory. *Phys Chem Chem Phys* 2016;18:13591–13603. doi:10.1039/C6CP00561F.
- [91] Mach, P., Pindak, R., Levelut, A.M., Barois, P., Nguyen, H.T., Baltes, H., et al. Structures of chiral smectic-*C* mesophases revealed by polarization-analyzed resonant x-ray scattering. *Phys Rev E* 1999;60:6793–6802. doi: 10.1103/PhysRevE.60.6793.
- [92] Cady, A., Pitney, J.A., Pindak, R., Matkin, L.S., Watson, S.J., Gleeson, H.F., et al. Orientational ordering in the chiral smectic- c_{n2}^* liquid crystal phase determined by resonant polarized x-ray diffraction. *Phys Rev E* 2001;64:050702. doi: 10.1103/PhysRevE.64.050702.
- [93] Folcia, C.L., Ortega, J., Etxebarria, J., Rodriguez-Conde, S., Sanz-Enguita, G., Geese, K., et al. Spontaneous and field-induced mesomorphism of a silyl-terminated bent-core liquid crystal as determined from second-harmonic generation and resonant x-ray scattering. *Soft Matter* 2014;10:196–205. doi: 10.1039/C3SM51277K.
- [94] Mårdalen, J., Riekel, C., Müller, H.. Anomalous X-ray scattering at the sulfur edge of poly(3-octylthiophene). *J Appl Cryst* 1994;27(2):192–195. doi:10.1107/S0021889893009835.
- [95] Chourou, S.T., Sarje, A., Li, X.S., Chan, E.R., Hexemer, A.. HipGISAXS: a high-performance computing code for simulating grazing-incidence x-ray scattering data. *J Appl Cryst* 2013;46(6):1781–1795. doi:10.1107/S0021889813025843.
- [96] Renaud, G., Lazzari, R., Leroy, F.. Probing surface and interface morphology with grazing incidence small angle x-ray scattering. *Surf Sci Rep* 2009;64(8):255–380. doi: 10.1016/j.surfrep.2009.07.002.
- [97] Jiang, Z., Lee, D.R., Narayanan, S., Wang, J., Sinha, S.K.. Waveguide-enhanced grazing-incidence small-angle x-ray scattering of buried nanostructures in thin films. *Phys Rev B* 2011;84:075440. doi:10.1103/PhysRevB.84.075440.
- [98] Sunday, D.F., Hammond, M.R., Wang, C., li Wu, W., DeLongchamp, D.M., Tjio, M., et al. Determination of the internal morphology of nanostructures patterned by directed self assembly. *ACS Nano* 2014;8(8):8426–8437.
- [99] Sunday, D.F., List, S., Chawla, J.S., Kline, R.J.. Evaluation of the effect of data quality on the profile uncertainty of critical dimension small angle x-ray scattering. *J Micro/Nanolith* 2016;15:014001. doi:10.1117/1.JMM.15.1.014001.
- [100] Chou, K.W., Yan, B., Li, R., Li, E.Q., Zhao, K., Anjum, D.H., et al. Spin-cast bulk heterojunction solar cells: A dynamical investigation. *Adv Mater* 2013;25(13):1923–1929. doi: 10.1002/adma.201203440.
- [101] Gu, X., Gunkel, I., Hexemer, A., Gu, W., Russell, T.P. An in situ grazing incidence x-ray scattering study of block copolymer thin films during solvent vapor annealing. *Adv Mater* 2014;26(2):273–281. doi:10.1002/adma.201302562.
- [102] Richter, L.J., DeLongchamp, D.M., Bokel, F.A., Engmann, S., Chou, K.W., Amassian, A., et al. In situ morphology studies of the mechanism for solution additive effects on the formation of bulk heterojunction films. *Adv Energy Mater* 2015;5(3):1400975. doi:10.1002/aenm.201400975.
- [103] Pröllner, S., Liu, F., Zhu, C., Wang, C., Russell, T.P., Hexemer, A., et al. Following the morphology formation in situ in printed active layers for organic solar cells. *Adv Energy Mater* 2016;6(1):1501580. doi:10.1002/aenm.201501580.
- [104] Sarje, A., Li, X.S., Chourou, S., Chan, E.R., Hexemer, A.. Massively parallel x-ray scattering simulations. In: *High Performance Computing, Networking, Storage and Analysis (SC)*, 2012 International Conference for. 2012, p. 1–11. doi: 10.1109/SC.2012.76.
- [105] Sarje, A., Li, X.S., Hexemer, A.. *High Performance Computing Systems. Performance Modeling, Benchmarking and Simulation*. Springer International Publishing; 2014.
- [106] McDonald, T.M., Mason, J.A., Kong, X., Bloch, E.D., Gygi, D., Dani, A., et al. Cooperative insertion of CO₂ in diamine-appended metal-organic frameworks. *Nature* 2015;519:303–308. doi:10.1038/nature14327.
- [107] Drisdell, W.S., Poloni, R., McDonald, T.M., Pascal, T.A., Wan, L.F., Pemmaraju, C.D., et al. Probing the mechanism of CO₂ capture in diamine-appended metal-organic frameworks using measured and simulated x-ray spectroscopy. *Phys Chem Chem Phys* 2015;17:21448–21457. doi:10.1039/c5cp02951a.
- [108] Helfand, E., Tagami, Y.. Theory of the interface between immiscible polymers. *J Polym Sci B Polym Phys* 1971;9(10):741–746. doi:10.1002/pol.1971.110091006.

- [109] Helfand, E., Tagami, Y. Theory of the Interface between Immiscible Polymers. II. *J Chem Phys* 1972;56(7):3592–3601. doi:10.1063/1.1677735.
- [110] Broseta, D., Fredrickson, G.H., Helfand, E., Leibler, L.. Molecular weight and polydispersity effects at polymer-polymer interfaces. *Macromolecules* 1990;23(1):132–139. doi:10.1021/ma00203a023.
- [111] Roe, R.J.. *Methods of X-ray and Neutron Scattering in Polymer Science*. Oxford University Press; 2000.
- [112] Leheny, R.L.. Xpcs: Nanoscale motion and rheology. *Curr Opin Colloid Interface Sci* 2012;17(1):3–12. doi:10.1016/j.cocis.2011.11.002.
- [113] Grubel, G., Zontone, F. Correlation spectroscopy with coherent x-rays. *J Alloys Compd* 2004;362(1-2):3–11. doi:10.1016/S0925-8388(03)00555-3.
- [114] Livet, F. Diffraction with a coherent x-ray beam: dynamics and imaging. *Acta Crystallogr, Sect A: Found Crystallogr* 2007;63:87–107. doi:10.1107/S010876730605570x.
- [115] Sutton, M.. A review of x-ray intensity fluctuation spectroscopy. *Comptes Rendus Physique* 2008;9(5-6):657–667. doi:10.1016/j.crhy.2007.04.008.
- [116] Cristofolini, L.. Synchrotron x-ray techniques for the investigation of structures and dynamics in interfacial systems. *Curr Opin Colloid Interface Sci* 2014;19(3):228–241. doi:10.1016/j.cocis.2014.03.006.
- [117] Gulari, E., Gulari, E., Tsunashima, Y., Chu, B.. Polymer diffusion in a dilute theta solution .I. polystyrene in cyclohexane. *Polymer* 1979;20(3):347–355. doi:10.1016/0032-3861(79)90099-5.
- [118] Jung, S.C., Bae, Y.C.. The effects of interaction energy on the volume phase transition of n-isopropylacrylamide-co-n-isopropylmethacrylamide nano-sized gel particles: Applicability of molecular simulation technique. *Polymer* 2009;50(20):4957–4963. doi:10.1016/j.polymer.2009.08.006.
- [119] Jung, S.C., Oh, S.Y., Bae, Y.C.. Reentrant swelling behavior of thermosensitive n-isopropylacrylamide nano-sized gel particles. *Polymer* 2009;50(14):3370–3377. doi:10.1016/j.polymer.2009.05.011.
- [120] Kim, S.M., Bae, Y.C.. Co-nonsolvency effect of thermosensitive n-isopropylacrylamide nanometer-sized gel particles in water-peg systems. *Polymer* 2013;54(8):2138–2145. doi:10.1016/j.polymer.2013.02.014.
- [121] Lee, S.M., Bae, Y.C.. Enhanced solvation effect of recollapsing behavior for cross-linked pmma particle gel in aqueous alcohol solutions. *Polymer* 2014;55(18):4684–4692. doi:10.1016/j.polymer.2014.07.033.
- [122] Oh, S.Y., Kim, H.J., Bae, Y.C.. Molecular thermodynamic analysis for phase transitions of linear and cross-linked poly(n-isopropylacrylamide) in water/2-propanol mixtures. *Polymer* 2013;54(25):6776–6784. doi:10.1016/j.polymer.2013.10.026.
- [123] Sedlak, M., Konak, C., Stepanek, P., Jakes, J.. Semidilute solutions of poly(methacrylic acid) in the absence of salt - dynamic light-scattering study. *Polymer* 1987;28(6):873–880. doi:10.1016/0032-3861(87)90156-X.
- [124] Tang, H.I., Johnson, P.L., Gulari, E.. Styrene polymerized in an oil-in-water microemulsion. *Polymer* 1984;25(9):1357–1362. doi:10.1016/0032-3861(84)90391-4.
- [125] Yi, F.P., Zheng, S.X.. Effect of hydrophobic polystyrene microphases on temperature-responsive behavior of poly(n-isopropylacrylamide) hydrogels. *Polymer* 2009;50(2):670–678. doi:10.1016/j.polymer.2008.11.038.
- [126] Gupta, P., Elkins, C., Long, T.E., Wilkes, G.L.. Electrospinning of linear homopolymers of poly(methyl methacrylate): exploring relationships between fiber formation, viscosity, molecular weight and concentration in a good solvent. *Polymer* 2005;46(13):4799–4810. doi:10.1016/j.polymer.2005.04.021.
- [127] Kita-Tokarczyk, K., Grumelard, J., Haeefe, T., Meier, W.. Block copolymer vesicles - using concepts from polymer chemistry to mimic biomembranes. *Polymer* 2005;46(11):3540–3563. doi:10.1016/j.polymer.2005.02.083.
- [128] Sinha, S.K., Jiang, Z., Lurio, L.B.. X-ray photon correlation spectroscopy studies of surfaces and thin films. *Adv Mater* 2014;26(46):7764–7785. doi:10.1002/adma.201401094.
- [129] Jackle, J.. The spectrum of surface waves on viscoelastic liquids of arbitrary depth. *J Phys Condens Matter* 1998;10(32):7121–7131. doi:10.1088/0953-8984/10/32/004.
- [130] Hoshino, T., Kikuchi, M., Murakami, D., Harada, Y., Mitamura, K., Ito, K., et al. X-ray photon correlation spectroscopy using a fast pixel array detector with a grid mask resolution enhancer. *J Synchrotron Rad* 2012;19:988–993. doi:10.1107/S0909049512038769.
- [131] Frieberg, B., Kim, J., Narayanan, S., Green, P.F. Surface layer dynamics in miscible polymer blends. *ACS Macro Lett* 2013;2(5):388–392. doi:10.1021/mz400104p.
- [132] Frieberg, B., Kim, J., Narayanan, S., Green, P.F. Surface dynamics of miscible polymer blend nanocomposites. *ACS Nano* 2014;8(1):607–613. doi:10.1021/nn405233a.
- [133] Livet, F., Bley, F., Ehrburger-Dolle, F., Morfin, I., Geissler, E., Sutton, M.. Homodyne and heterodyne x-ray photon correlation spectroscopy: latex particles and elastomers. *J Appl Crystallogr* 2007;40:S38–S42. doi:10.1107/S0021889807003561.
- [134] Livet, F., Bley, F., Ehrburger-Dolle, F., Morfin, I., Geissler, E., Sutton, M.. X-ray intensity fluctuation spectroscopy by heterodyne detection. *J of Synchrotron Rad* 2006;13:453–458. doi:10.1107/S0909049506030044.
- [135] Steier, C., et al. Proposal for a Soft X-ray Diffraction Limited Upgrade of the ALS. In: *Proceedings, 5th International Particle Accelerator Conference (IPAC 2014)*. 2014, p. MOPME084.
- [136] Ricci, A.. Nanoscale dynamics in complex materials by resonant x-ray photon correlation spectroscopy (rxpcs). *J Supercond Nov Magn* 2015;28(4):1295–1298. doi:10.1007/s10948-014-2907-3.
- [137] Konings, S., Schussler-Langeheine, C., Ott, H., Weschke, E., Schierle, E., Zabel, H., et al. Magnetic domain fluctuations in an antiferromagnetic film observed with coherent resonant soft x-ray scattering. *Phys Rev Lett* 2011;106(7). doi:10.1103/PhysRevLett.106.077402.
- [138] Bao, Z., Springell, R., Walker, H.C., Leiste, H., Kuebel, K., Prang, R., et al. Antiferromagnetism in uo_2 thin epitaxial films. *Phys Rev B* 2013;88(13). doi:10.1103/PhysRevB.88.134426.
- [139] Chen, S.W., Guo, H., Seu, K.A., Dumesnil, K., Roy, S., Sinha, S.K.. Jamming behavior of domains in a spiral antiferromagnetic system. *Phys Rev Lett* 2013;110(21). doi:10.1103/PhysRevLett.110.217201.
- [140] Matsumura, T., Nakao, H., Murakami, Y.. Resonant x-ray scattering experiments on the ordering of electronic degrees of freedom. *J Phys Soc Jpn* 2013;82(2). doi:10.7566/Jpsj.82.021007.
- [141] Ewerlin, M., Pfau, B., Gunther, C.M., Schaffert, S., Eisebitt, S., Abrudan, R., et al. Exploration of magnetic fluctuations in $pdfe$ films. *J Phys Condens Matter* 2013;25(26). doi:10.1088/0953-8984/25/26/266001.
- [142] McDermott, A.G., Budd, P.M., McKeown, N.B., Colina, C.M., Runt, J.. Physical aging of polymers of intrinsic microporosity: a saxs/waxs study. *J Mater Chem A* 2014;2:11742–11752. doi:10.1039/C4TA02165G.

- [143] Chen, X.C., Kortright, J.B., Balsara, N.P. Water uptake and proton conductivity in porous block copolymer electrolyte membranes. *Macromolecules* 2015;48(16):5648–5655. doi: 10.1021/acs.macromol.5b00950.
- [144] Wong, D.T., Wang, C., Pople, J.A., Balsara, N.P. Effect of nonsolvent exposure on morphology of mesoporous semicrystalline block copolymer films. *Macromolecules* 2013;46(11):4411–4417. doi:10.1021/ma400051x.
- [145] Gurevitch, I., Buonsanti, R., Teran, A.A., Gludovatz, B., Ritchie, R.O., Cabana, J., et al. Nanocomposites of titanium dioxide and polystyrene-poly(ethylene oxide) block copolymer as solid-state electrolytes for lithium metal batteries. *J Electrochem Soc* 2013;160(9):A1611–A1617. doi: 10.1149/2.117309jes.
- [146] Schmidt-Rohr, K., Chen, Q. Parallel cylindrical water nanochannels in nafion fuel-cell membranes. *Nat Mater* 2008;7:75–83. doi:10.1038/nmat2074.
- [147] Modestino, M.A., Kusoglu, A., Hexemer, A., Weber, A.Z., Segalman, R.A.. Controlling nafion structure and properties via wetting interactions. *Macromolecules* 2012;45(11):4681–4688. doi:10.1021/ma300212f.
- [148] Kusoglu, A., Modestino, M.A., Hexemer, A., Segalman, R.A., Weber, A.Z.. Subsecond morphological changes in nafion during water uptake detected by small-angle x-ray scattering. *ACS Macro Lett* 2012;1(1):33–36. doi: 10.1021/mz200015c.
- [149] Kanai, K., Noda, Y., Kato, K., Kubo, T., Iketaki, K., Shimizu, A., et al. Electronic structure of delocalized singlet biradical Ph₂-IDPL solid film. *Phys Chem Chem Phys* 2010;12:12570–12577. doi:10.1039/c0cp00178c.
- [150] Duong, D.T., Phan, H., Hanifi, D., Jo, P.S., Nguyen, T.Q., Salleo, A.. Direct observation of doping sites in temperature-controlled, p-doped P3HT thin films by conducting atomic force microscopy. *Adv Mater* 2014;26(35):6069–6073. doi: 10.1002/adma.201402015.
- [151] Tian, H.C., Liu, J.Q., Wei, D.X., Kang, X.Y., Zhang, C., Du, J.C., et al. Graphene oxide doped conducting polymer nanocomposite film for electrode-tissue interface. *Biomaterials* 2014;35(7):2120–2129. doi: 10.1016/j.biomaterials.2013.11.058.
- [152] Magnuson, M., Guo, J.H., Butorin, S.M., Agui, A., S athe, C., Nordgren, J., et al. The electronic structure of polyaniline and doped phases studied by soft x-ray absorption and emission spectroscopies. *J Chem Phys* 1999;111(10):4756–4761. doi:10.1063/1.479238.
- [153] Guay, D., Tourillon, G., Fontaine, A.. Electrochemical inclusion of copper and iron species in a conducting polymer observed in situ using time-resolved x-ray absorption spectroscopy. *Faraday Discuss Chem Soc* 1990;89:41–50. doi: 10.1039/DC9908900041.
- [154] Aziz, E., Vollmer, A., Eisebitt, S., Eberhardt, W., Pingel, P., Neher, D., et al. Localized charge transfer in a molecularly doped conducting polymer. *Adv Mater* 2007;19(20):3257–3260. doi:10.1002/adma.200700926.
- [155] Cochran, J.E., Junk, M.J.N., Gludell, A.M., Miller, P.L., Cowart, J.S., Toney, M.F., et al. Molecular interactions and ordering in electrically doped polymers: Blends of PBTTT and F4TCNQ. *Macromolecules* 2014;47(19):6836–6846. doi: 10.1021/ma501547h.