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Note: Electrical and thermal characterization of a ferroelectric thin film with an electro-thermal nanoprobe

R. Jackson,¹ P. C. Fletcher,² K. Jambunathan,³ A. R. Damodaran,³ J. N. Emmerich,¹ H. Teng,¹ L. W. Martin,³ W. P. King,^{2,3} and Y. Wu^{1,a)}

¹College of Engineering, Mathematics, and Science, University of Wisconsin-Platteville, Platteville, Wisconsin 53818, USA

²Department of Mechanical Science and Engineering, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA

³Department of Materials Science and Engineering, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA

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The localized temperature-dependent piezoelectric response of ferroelectric barium strontium titanate (BST) thin films is studied using an electro-thermal (ET) nanoprobe. The ET probe provides independent electrical and thermal excitation to a nanometer-scale volume of the specimen and is capable of detecting the phase transition temperature of the BST thin films. The piezoresponse measured by the ET probe follows the temperature dependence of the piezoelectric constant, whereas with bulk heating the response follows the temperature dependence of the spontaneous polarization. The observed differences stem from the localized inhomogeneous electro-thermal field distribution at the specimen. © 2012 American Institute of Physics. [http://dx.doi.org/10.1063/1.4733730]

Ferroelectric materials are extensively used in nonvolatile data storage,¹ microscale actuators and sensors,² and microwave electronic components.³ The understanding of local polarization dynamics and domain structure in ferroelectric materials has been greatly facilitated by piezoresponse force microscopy (PFM).^{4–8} PFM is a scanning probe-based technique that detects the converse piezoelectric response of a material via bias-induced surface strains and the resulting cantilever deflection. Only a few articles consider ferroelectric response above room temperature^{9–12} and these studies rely on atomic force microscope (AFM) compatible heater stages to control temperature. This letter describes measurements of the piezoelectric response of ferroelectric barium strontium titanate (Ba_{0.6}Sr_{0.4}TiO₃, BST) thin films around the Curie temperature using an electro-thermal (ET) nanoprobe.¹

A few studies have used heated AFM probes for electrical measurements.^{13,14} Tip-based heating provides advantages over bulk heating because the integrated heater is capable of reaching temperatures in excess of 1200 °C with thermal time constants on the order of tens of microseconds.¹⁵ Tip-based heating also produces localized temperature gradients at the sample surface that can be utilized to separate local surface properties from bulk properties. The ET nanoprobe used here¹⁶ facilitates independent control of electrical and thermal tip excitations.

Thin films of ferroelectric BST and amorphous silicon dioxide (SiO_2) of 75 nm thickness were grown on an SrRuO₃ buffer layer on (001)-oriented SrTiO₃ single crystal substrates. The BST films were fabricated by pulsed-laser deposition. The SiO₂ films were deposited via plasma-enhanced chemical vapor deposition. The SiO₂ film was used as a control sample because it has a negligible piezoelectric response.

Figure 1(a) shows a scanning electron microscope micrograph of the ET probe. There are three basic elements of the probe: an integrated solid-state heater formed by two cantilever legs, a third high-doped n+-type silicon leg addressing the tip electrode, and an n-p-n semiconductor junction for electrical isolation between the heater region and the tip electrode. Figure 1(b) is a schematic diagram of the electrothermal contact between the heated tip and the specimen, showing the relevant temperatures and the electric impedance network model of the contact. We calibrated the heater temperature using micro-Raman spectroscopy.¹⁷ In the experiments, the cantilever heater temperature was varied over the range of 25 °C-553 °C. The relationship between tip-sample interface temperature, $T_{\text{interface}}$, and the heater temperature, T_{heater} , can be found using a previously developed thermal circuit model.¹⁸ For a heated tip contacting room temperature BST in air, a maximum heater temperature of 553 °C translates to a maximum interface temperature of 219 °C for an ambient temperature, T_{∞} , of 25 °C.

We simulated the electric and thermal field profile under the tip using the finite element program ANSYS. Simulation results show the temperature is less than 10% of its surface value and the electric field is less than 5% of its surface value beyond one tip radius from the tip into the film (see supplementary material for the details on the electrothermal modeling).¹⁹ Therefore, the piezoresponse depends on the film material within one tip radius from the tip. We refer to this volume as the signal generation volume. The signal generation volume scales with the radius of the tip, and for this study is around 2.6×10^5 nm³.

In addition to the ET probe, we also measured the piezoresponse of the BST thin film using a sample heating stage and commercial conductive AFM probes. In both local heating and bulk heating experiments, we measured the contact resonance amplitude of the cantilever under

a)Author to whom correspondence should be addressed. Electronic mail: wuy@uwplatt.edu.



FIG. 1. (a) SEM micrograph of an electro-thermal probe showing the integrated heater, the independent electrode, and the n-p-n junction for electrical isolation. (b) Simple equivalent electric circuit model of the tip-sample contact and relevant temperatures.

fixed excitation voltage amplitude, V_{tip} , while controlling the temperature. The relation between V_{tip} and the effective voltage at the tip-sample interface, V_{eff} , is²⁰

$$\frac{V_{\rm eff}}{V_{\rm tip}} = \frac{1}{\sqrt{(1 + C_s/C_{cont})^2 + R_{cont}^2 C_s^2 \omega^2}},$$
(1)

where R_{cont} is the contact resistance between the tip and the sample, C_{cont} is the contact capacitance, C_s is the total capacitance of the sample, and ω is the frequency of the oscillating electric field. Both C_{cont} and C_s are on the order of tens of attofarads at room temperature.^{21,22}

Electrostatic coupling between the cantilever and sample can result in a spurious contribution to the measured signal in PFM.²³ The ET probe conducts an ac electric field to the tip while driving the integrated heater using dc. The electrostatic coupling from the dc driving voltage for different heater temperatures gives rise to flexural and torsional contact resonance modes, as shown in Figure 2. The response on the SiO₂ sample is due to electrostatic coupling and it is negligible at the fundamental mode (166 kHz near room temperature) compared to the piezoresponse on the BST sample. The electrostatic coupling is dominant at the second vibration mode (185 kHz near room temperature) because electrostatic force at the geometrically off-centered heater promotes torsional modes of the three-legged cantilever.²⁴

Figure 3(a) shows the local temperature-dependent piezoresponse at the fundamental contact resonance frequency of the ET probe on the BST and SiO₂ samples. The BST piezoresponse increases, then decreases after a transition point corresponding with an interface temperature of 97 °C, whereas the SiO₂ response only increases slightly with



FIG. 2. Cantilever response amplitude under electric excitation for various ET probe heater temperatures. The solid red line is the response on the BST sample and the dashed blue line is the response on the SiO_2 sample.

temperature due to the electrostatic effect. Figure 3(b) shows the temperature-dependent response for bulk heating. The BST response decreases monotonically until ~ 100 °C, above which the response shows little change with temperature, similar to the SiO₂ response for the entire temperature spectrum.



FIG. 3. (a) Contact resonance response on the BST and SiO₂ samples under localized thermal and electrical excitation. The BST transition temperature corresponds to ~97 °C. (b) Contact resonance response on the BST and SiO₂ samples as a function of bulk heating temperature. The BST transition temperature corresponds to ~100 °C.

For a homogeneously polarized, stress-free, ferroelectric material with polarization in the z direction, theoretical analysis⁵ shows the piezoresponse signal is proportional to the electric field strength within the signal generation volume and the piezoelectric constant d_{33} . The piezoelectric coefficient d_{33} increases with temperature.^{25–27} For bulk heating, the electric field strength decreases with increasing temperature due the increasing dielectric constant of the BST film and thus the ratio of C_s/C_{cont} in Eq. (1). Previous published results suggested that the temperature-dependent piezoelectric response follows the spontaneous polarization when C_s is much greater than C_{cont} .¹¹ In addition, the unscreened polarizationbound charges under the cantilever also contribute to the temperature dependence of the spontaneous polarization via electrostatic coupling.⁹ Our results using the heating stage are in agreement with results from other studies in literature using bulk heating methods.^{9,11} The piezoresponse is observed to continuously decrease with increasing temperature, consistent with a smearing of the phase transition, as has been reported in thin film samples of ferroelectrics.²⁸ The results from the ET probe differ from those using bulk heating methods^{9,11} and reflect the complex interplay between thermal and piezoelectric responses in the system. Our technique only heats a nanometer scale volume of material directly under the tip. The specimen away from the tip remains at room temperature, thus reducing the temperature dependent response due to the unscreened polarization-bound charges via stray capacitive coupling. The measured piezoresponse follows the temperature-dependence of d_{33} , which agrees with a recent study of the temperature-dependent piezoresponse of lead zirconate titanate using a micro-heater with minimization of the cantilever-microheater overlap region.¹² Furthermore, the non-uniform temperature distribution effectively limits the increase of capacitance C_s with increasing $T_{\text{interface}}$. With a steady electric field, the piezoresponse increases with $T_{\text{interface}}$ until the region directly under the heated tip reaches the Curie temperature of ~ 100 °C. Above the Curie temperature, the overall response decreases as a larger percentage of the signal generation volume is heated above the Curie temperature and becomes paraelectric with no piezoelectric response.

In summary, we measured the temperature-dependent piezoresponse of a ferroelectric thin film with nanometerscale electrical and thermal excitation using an electrothermal nanoprobe. The measured piezoresponse under localized heating differs from results using a bulk heating approach because of the localized non-uniform temperature distribution in the specimen. This study demonstrates an application of the ET nanoprobe in nanoscale electro-thermal characterization. This work was supported by the National Science Foundation (NSF) under Grant No. 0960232. J.K. and L.W.M. acknowledge support from the Office of Naval Research (ONR) under Grant No. N00014-10-10525. A.R.D. and L.W.M. acknowledge support from the (U.S.) Army Research Office (USARO) under Grant No. W911NF-10-1-0482. The authors thank Dr. Scott MacLaren at FSMRL Central Facilities for his assistance with part of the AFM work.

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