Lawrence Berkeley National Laboratory

Lawrence Berkeley National Laboratory

Title

SOME CALCULATOR PROGRAMS FOR PARTICLE PHYSICS

Permalink

https://escholarship.org/uc/item/2s10279q

Author

Wohl, C.G.

Publication Date

1982

Peer reviewed

LBL--13987

DE82 012362

_DISCLAMER -

This took an proper of an incomprise of a control to in agrees of the blade State Controlled. Name the blade State Controlled to the growth the land of the enclosed, making the blade State Controlled to the growth the land of the enclosed, and as manufactured to the state of the enclosed to the enclosed of the enclosed of the enclosed of the enclosed of the enclosed controlled to the enclosed controlled to the enclosed controlled to the enclosed of the enclosed controlled to the enclosed of the enclosed controlled to the enclosed of the enclosed controlled and strategy acres to the enclosed controlled to the enclosed controlled to the enclosed of the enclosed of

SOME CALCULATOR PROGRAMS FOR PARTICLE PHYSICS

Charles G. Wohl Lawrence Berkeley Laboratory University of California Berkeley, CA 94720

January 1982

This work was supported by the Director, Office of Energy Research, Office of High Energy and Nuclear Physics, Division of High Energy Physics of the U.S. Department of Energy under Contract Number W-7405-ENG-48.



The ELLIPSE, DALITZ RECTANGULAR, and CM programs work as given on the HP-19C and HP-29C, while the other programs need changes from the following list. In storage-register addresses and in RCL and STO statements:

Also:

In the LEGENDRE program, change step 32 to RCL $\cdot 0$, and make $R_{\cdot 0}$, not R_{0} , the storage register for a_{0} .

The write-up for each program tells what the program does and how to run it, discusses any limitations or special cases or pit-falls, and gives an example or two to test that the program is stored properly and that its operation is understood. The examples have all been rechecked from the final typescript on an HP-97 and, with the necessary changes, on an HP-29C.

LEGENDRE

The LEGENDRE program calculates the values of the Legendre polynomial series

$$A(x) = C \sum_{n=0}^{N} a_n P_n(x)$$

at $x=x_0$, $x_0+\Delta x$, $x_0+2\Delta x$, The input data are the values of x_0 , Δx , N (<10), the overall normalization constant C, and the expansion coefficients a_n .

The method used is of interest. The straightforward way to calculate A(x) is to start at the bottom of the ladder of Legendre polynomials with P_0 and P_1 , use the standard recursion relation to climb to higher rungs, and accumulate the products $a_n P_n$ along the way. Suppose, however, we define a new set of (x-dependent) coefficients c_n recursively by

$$c_n = a_n + (\frac{2n+1}{n+1}) \times c_{n+1} - (\frac{n+1}{n+2}) c_{n+2}$$

and work downward from c_N , with $c_{N+1}=c_{N+2}=0$. This leads to the remarkable result that

$$\sum_{n=0}^{N} a_n v_n \approx c_0 ,$$

and so the need to accumulate $a_{n}^{\ P}_n$ terms and keep track of the sum along the way is eliminated. This method, which was discovered by C.W. Clenshaw, is widely applicable to the summation of series of orthogonal polynomials; see, for example, p. 11 of F.S. Acton, Numerical Methods that Work (Harper and Row, New York, 1970). The recursion relation for the new coefficients is of course related to that for the orthogonal polynomials, and has to be worked out separately for each case. In general, the two lowest coefficients (here c_0 and c_1) are involved in the final sum, but for Legendre polynomials the result is particularly simple.

To run the program, store the data in the indicated registers and start with a GSB O command. After some seconds, the value of x

is displayed briefly, and then A(x) is displayed and the program stops. For easy access, x is at this point in the Y register. Thereafter, the stored value of x is incremented by Δx and A(x) is calculated for the new value each time R/S is pressed.

To test the program, <u>zero the registers</u> and then set $a_9 = 1$, N = 9, $x_0 = 1$, $\Delta x = -0.25$, and C = 1. This input generates $P_9(x)$ at x = 1.0, 0.75, 0.5, The first few results are:

x = 1.0	P ₉ = 1.000000001
0.75	0.310331851
0.5	-0.267898560
0.25	0.176824421
0.0	0.000000000

These results agree perfectly with 8-place tables of P_9 in chap. 8 of M. Abramowitz and I. Stegun, Handbook of Mathematical Functions (Dover, New York, 1972). Some error is to be expected in the ninth place, as in $P_9(1)$ above, but this far exceeds normal requirements of accuracy.

Program: LEGENDRE

PAUSE

30

1109	Tem. Leasing			
Step	<u>Keys</u>	Step	Keys	Registers
*01	LBL O	31	x ‡ y	† ^R 0 ^a 0
02	0	†3 <i>2</i>	RCL 0	R_1 a_1
03	ENTER	33	+	R_2 a_2
04	ENTER	34	RCL A	R_{3} a_{3}
05	RCL B	35	x	R_4 a_4
06	1	36	R/S	R_5 a_5
07	+	37	RCL D	^R 6 ^a 6
08	STO I	38	RCL C	R ₇ a ₇
*09	LBL 7	39	+	R ₈ a ₈
10	1/x	40	STO C	R_{g} a_{g}
11	1	41	GTO O	R _A C
12	+	*42	LBL 2	R _B N
13	÷	43	RCL (i	$R_C \times_0 (+x)$
14	CHS	44	+	R_{D} Δx
15	x≠y	45	xŻy	$R_{\overline{I}}$ used
16	2	46	RCL I	
17	RCL I	47	GTO 1	
18	1/x			
19	-			
20	RCL C			
21	x			
22	$x \stackrel{\rightarrow}{+} y$			†
23	ENTER			† HP-19C and -29C users see the introduction.
24	R↓			
25	x			
	+			
27	DSZ I			
28	GTO 2			
29	RCL C			

ASSOCIATED LEGENDRE

The ASSOCIATED LEGENDRE program calculates the values of the first-associated Legendre polynomial series

$$B(x) = C \sum_{n=1}^{N} b_{n} P_{n}^{1}(x) = -C \sqrt{1-x^{2}} \sum_{n=1}^{N} b_{n} dP_{n}/dx$$

at $x = x_0$, $x_0 + \Delta x$, $x_0 + 2\Delta x$, \cdots ($|x| \le 1$). The input data are the values of x_0 , Δx , N (<10), the overall normalization constant C, and the expansion coefficients b_n . Note the minus sign in the definition of P_n^1 . Angular distributions are sometimes expanded with a sign convention opposite to the above, the remedy for which is to make C negative.

The method used here (see the LEGENDRE write-up) is to define a new set of coefficients \mathbf{d}_n recursively by

$$d_n = b_n + (\frac{2n+1}{n}) \times d_{n+1} - (\frac{n+2}{n+1}) d_{n+2}$$

and work downward from d_{N} , with $d_{N+1} = d_{N+2} = 0$. This leads to

$$\sum_{n=1}^{N} b_{n} P_{n}^{1} = -d_{1} \sqrt{1-x^{2}} ,$$

and so again the need to keep track of a partial sum along the way is eliminated.

To run the program, store the data in the indicated registers and start with a GSB 0 command. After some seconds, the value of x is displayed briefly, and then B(x) is displayed and the program stops. For easy access, x is at this point in the Y register. Thereafter, the stored value of x is incremented by Δx and B(x) is calculated for the new value each time R/S is pressed. If x is outside the range -1 to +1, Error is displayed.

To test the program, <u>zero</u> the <u>registers</u> and then set $b_9 = 1$, N = 9, $x_0 = 1$, $\Delta x = -0.25$, and C = 1. This input generates $P_9^1(x)$ at x = 1.0, 0.75, 0.5, The first few results are:

x = 1.0	$p_9^1 = 0.0000000000$
0.75	0.478134503
0.5	-0.626763685
0.25	1.827987321
0.0	-2.460937500

These results agree perfectly with those obtained using 7-place tables of $dP_{\rm g}/dx$ in chap. 8 of M. Abramowitz and I. Stegun, Handbook of Mathematical Functions (Dover, New York, 1972). In fact, they are almost certainly accurate to eight places, but some error is to be expected in the ninth place.

Program: ASSOCIATED LEGENDRE

Step	Keys	<u>S</u> :	<u>tep</u>	Keys
*01	LBL 0		31	RCL C
02	0		32	cos-1
03	ENTER	:	33	sin
94	ENTER		34	x
05	RCL B	3	5	RCL C
06	STO I	3	36	PAUSE
*07	LBL 1	3	37	x ‡ y
08	1	3	88	RCL A
09	+	3	19	x
10	1/x	4	10	R/S
11	1	4	1	RCL D
12	+	4	2	RCL C
13	x	4	3	+
14	CHS	4	4	STO C
15	x ‡ y	4	5	GTO 0
16	RCL I	*4	6	LBL 2
17	1/x	4	7	x≠y
18	2	4	8	RCL I
19	+	4	9	GTO ï
20	RCL C			
21	x			
22	x			
23	ENTER			
24	R↓			
25	x			
26	+			
2 7	RCL (i)			
28	-			
29	DSZ I			
30	GTO 2			

Registers ь₁ R_1 R_2 b_2 R₃ b₃ R₄ ь₄ R₅ b₅ R₆ b₆ R₇ b₇ R₈ b₈ R₉ b₉ R_A С N R_B R_C x₀ (→ x) R_D Δx used R_{T}

CONFIDENCE and EVEN N and POISSON

The CONFIDENCE program calculates confidence levels for the χ^2 probability distribution with N degrees of freedom. The calculation is much shorter when N is even than when it is odd, so there is a short program, EVEN N, for that case (CONFIDENCE handles any N). CONFIDENCE may also be used to get confidence levels for the Gaussian (or normal) and the Poisson probability distributions, but for the Poisson case there is a much shorter program, POISSON.

Figure 1 shows how the confidence levels will be defined here.

- (a) The confidence level $CL(\chi_0^2, N)$ for the χ^2 probability distribution $P_N(\chi^2)$ with N degrees of freedom is the probability that a χ^2 greater than χ_0^2 would be obtained.
- (b) The confidence level $\mathrm{CL}_g(\chi_0)$ for the Gaussian probability distribution $P(\chi)$ is the probability that a result more than χ_0 standard deviations from the mean would be obtained; this is related to $\mathrm{CL}(\chi_0^2,\ N)$ by

$$CL_{\alpha}(\chi_0) = CL(\chi_0^2, 1)$$
.

(c) The confidence level ${\rm CL}_p(n_0,\, \bar n)$ for the Poisson probability distribution $P_{\bar n}(n)$ with mean $\bar n$ is the probability that a value of n greater than n_0 would be obtained; this is related to ${\rm CL}(\chi_0^2,\, N)$ by

$$CL_{p}(n_{0}, \bar{n}) = 1.0 - CL(\chi_{0}^{2}, N)$$
,

with $\chi_0^2 = 2\overline{n}$ and $N = 2n_0 + 2$.

The equations used to calculate $CL(\chi_0^2,\,N)$ are as follows. When N is even,

CL(
$$\chi_0^2$$
, N) = $\sqrt{2\pi} \ \text{Z}(\chi_0) \left[1 + \sum_{n=1}^{N'} \frac{\chi_0^{2n}}{2 \cdot 4 \cdot 6 \cdots (2n)} \right]$,

where $Z(\chi_0) = (2\pi \exp \chi_0^2)^{-\frac{1}{2}}$ and N' = (N+2)/2. When N is odd,

$$CL(\chi_0^2, N) = CL_g(\chi_0) + 2Z(\chi_0) \sum_{n=1}^{N''} \frac{\chi_0^{2n-1}}{1 \cdot 3 \cdot 5 \cdot \cdot \cdot (2n-1)} ,$$

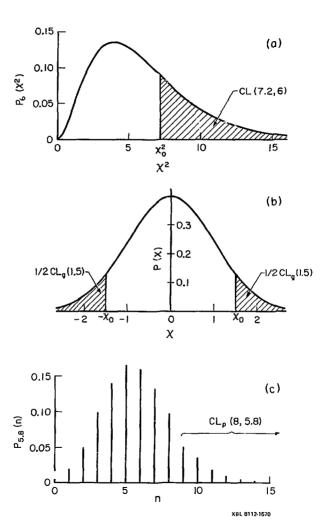


Figure 1

where N" = (N-1)/2. There is no closed expression for $CL_g(\chi_0)$. For $\chi_0\geqslant 2$, a truncated continued fraction has been used:

$$CL_g(\chi_0) = 2Z(\chi_0) \left[\frac{1}{\chi_0 + \chi_0 +$$

For χ_0 < 2, the equation used is

$$CL_g(\chi_0) = 1.0 - 2Z(\chi_0) \sum_{n=1}^{\infty} \frac{\chi_0^{2n-1}}{1 \cdot 3 \cdot 5 \cdot \dots (2n-1)}$$
,

where "infinity" is reached when the last term added to the series is smaller than 10^{-9} . For N > 1, the early terms of the series here are cancelling terms of the series in the equation at the bottom of p. 10. Account is taken of the fact that for very large N or very small χ_0^2 the "infinite" series can be shorter than the finite one.

To run CONFIDENCE or EVEN N, store χ_0^2 and N in the indicated registers and start the program with a GSB O command. When it stops, the confidence level is displayed. The range of χ_0^2 covered is 0.0 < χ_0^2 < 460.5; CL(0.0, N) is of course 1.0, and when χ_0^2 exceeds 460.5 there is an overflow.

To run POISSON, store n_0 and \bar{n} as indicated, press GSB 0, and see the confidence level when the program stops. It is easy with a few trials to home in on a value of \bar{n} that for a given n_0 produces a confidence level of, say, 0.95.

Following are some examples, with approximate running times.

The starred examples test all the parts of CONFIDENCE, and they at least should be tried.

	CL(0.04,]	1) =	0.841480581	=	CL _a (0.2)	(7	sec)
	CL(3.61, 1	1) =	0.057433120	=	CL _q (1.9)	(19	sec)
			0.045500263			(17	sec)
	CL(25.0, 1	L) =	0.000000573	=	CL_(5.0)	(17	sec)
*	CL(3.61, 9	9) =	0.935159132		•	(24	sec)
*	CL{25.0, 9	9) =	0.002971180			(21	sec)
	CL(2.00, 2	2) =	0.367879441	=	1.0 - CL _D (0, 1.0)	(2	sec)
*	CL(5050)) =	0-473398469	=	1.0 - CL (24, 25.0)	(26	sec)

For N = 1, there are 15-place tables of (1-CL/2) in chap. 26 of M. Abramowitz and I. Stegun, Handbook of Mathematical Functions (Dover, New York, 1972). For all N, there are 9-place tables of (1-CL) in H.L. Harter, New Tables of the Incomplete Gamma-Function Ratio and of Percentage Points of the Chi-Square and Beta Distributions (Aerospace Research Laboratories, U.S. Air Force, 1964). Comparisons show that for N = 1 the absolute error from the program never exceeds 1×10^{-9} . For N < 100, the absolute error rarely exceeds 1×10^{-9} , but for N > 100 absolute errors as large as 3×10^{-9} have been found.

Program: CONFIDENCE

Step	<u>Keys</u>	<u>Step</u>	<u>Keys</u>	Step	Keys	Regist	:ers
*01	LBL 0	31	2	61	x‡y	R_{1}	x_0^2
02	RCL 2	32	0	62	x > y?	R_2	N
03	STO I	33	STO I	63	GTO 5	R_3	used
04	3	34	6	64	1	_	used
*05	LBL 1	*35	LBL 3	65	STO I	R_{T}^{T}	used
06	RCL I	36	RCL I	66	RCL 4	_	
07	2	37	x≠y	67	ABS		
08	-	38	÷	68	CHS		
09	x < 0?	39	RCL 3	*69	LBL 6		
10	GTO 2	40	+	70	2		
11	STO I	41	DSZ I	71	π		
12	x = 0?	42	GTO 3	7.3	÷		
13	GT0 7	43	1/x	73	√x		
14	RCL 1	44	RCL 4	74	X		
15	÷	45	+	75	ENTER		
16	÷	46	GTO 6	*76	LBL 7		
17	7	*47	LBL 4	77	R∔		
18	+	48	RCL 3	78	RCL 1		
19	GTO 1	49	STO - 4	79	2		
*20	LBL 2	*50	LBL 5	80	÷		
21	+	51	ISZ I	81	e ^x		
22	RCL .	5 <i>2</i>	ISZ I	82	÷		
23	√×	53	RCL 1	83	RCL I		
24	ST0 3	5 4	х	84	+		
25	÷	55	RCL I	85	RTN		
26	STO 4	56	÷				
27	RCL 1	57	ST0 - 4				
28	4	58	EEX				
29	x > y?	59	CHS				
30	GTO 4	60	9				

Program: EVEN N

Program: POISSON

Step	<u>Keys</u>	Registers	<u>Step</u>	<u>Keys</u>	Registers
*01	LBL 0	$R_1 \chi_0^2$	*01	LBL 0	R ₁ n ₀
02	RCL 2	R ₂ N	02	RCL 1	$R_2 \bar{n}$
03	STO I	R _I used	03	STO I	$R_{\overline{I}}$ used
04	1	1	04	1	-
*05	LBL 1		05	ENTER	
06	DSZ I		06	ISZ I	
07	DSZ I		*07	LBL 1	
08	GTO 2		08	DSZ I	
09	RCL 1		0 9	GTO 2	
10	2		10	RCL 2	
11	÷		11	ex	
12	e*		12	÷	
13	÷		13	-	
14	RTN		14	RTN	
*15	LBL 2		*15	LBL 2	
16	RCL 1		16	RCL 2	
17	RCL I		17	RCL I	
18	÷		18	÷	
19	X		19	x	
20	1		20	4	
21	+		21	GTO 1	
22	GTO 1				

TWO BODY and CM

In the 2-body reaction $a+b \to 1+2$, let m_i be the mass of particle i and let P_a be the momentum of particle a in the inertial frame in which particle b is at rest (the lab frame). The four masses and P_a are the input data for the TWO BODY program, which calculates the following quantities:

s = c.m. energy squared

E = c.m. energy

p = initial-state c.m. momentum

p' = final-state c.m. momentum

 $\Delta t = \Delta u = range of 4-momentum transfer squared$

 $t_0^{(t_{\pi})} = 4$ -momentum-transfer squared between particles a and 1 when 1 is produced at 0° (180°) in the c.m.

 $u_0^{\prime} (u_{\pi}) = 4$ -momentum-transfer squared between particles a and 2 when 2 is produced at 0° (180°) in the c.m.

CM is a shorter program that uses P_a , m_a , and m_b as input to calculate s, E, and p.

To run TWO BODY, store the four masses and P_a in the indicated recesters, and start the program with a GSB 0 command. When it stops, the calculated quantities are in the storage registers (u_{π} is displayed). If Error appears when the program is started, then P_a may be below the threshold for the reaction. In some cases, however, P_a can be too low without such notice being given.

To run CM, store P_a , m_a , and m_b in the indicated registers, and start the program with a GSB O command. When it stops, s, E, and p are in the storage registers (p is displayed).

As a test example, consider $\pi^+p + \kappa^+\Sigma^+$ scattering at 4 GeV/c. Then, in units of GeV or GeV/c or GeV², the input numbers are $P_a=4.0$, $P_a=0.1396$, $P_b=0.9383$, $P_a=0.4937$, and $P_a=1.1894$, and the output numbers are $P_a=4.0$, $P_a=4.0$,

Program: TWO BODY Program: CM

Step	Keys	<u>Step</u>	<u>Keys</u>	<u>Step</u>	<u>Keys</u>	<u>Step</u>	<u>Keys</u>
*01	LBL 0	31	GSB 2	61	-	*01	LBL 0
02	RCL 0	32	-	62	GSB 2	02	RCL 0
03	RCL 2	33	\sqrt{x}	63	-	03	RCL 1
04	RCL 1	34	ST0 8	64	CHS	04	→ P
05	RCL 0	35	x	65	RTN	05	RCL 2
06	GSB 2	36	4	*66	LBL 2	06	+
07	+	37	x	67	x ²	07	x ²
08	√x	38	STO 9	68	x≠y	Oε	RCL 0
09	+	39	CHS	69	x ²	09	x^2
10	GSB 2	40	RCL 1	70	RTN	10	-
11	-	41	GSB 1			11	ST0 3
12	ST0 5	42	STO A			12	√x
13	\sqrt{x}	43	+			13	ST0 4
14	ST0 6	44	STO B			14	RCL 0
15	÷	45	R∔			15	x≠y
16	RCL 2	46	RCL 2	кедт	<u>sters</u>	16	÷
17	x	47	GSB 1	$R_{\mathcal{O}}$	P a	17	RCL 2
18	ST 0 7	48	STO C	R ₁	m _a	18	x
19	RCL 3	49	+	R ₂	m _b	19	ST0 5
20	RCL 4	50	STO D	R ₃	m ₁	20	RTN
21	RCL 3	51	RTN	R_{4}	m ₂		
22	GSB 2	*52	LBL 1	R ₅	s		
23	-	53	RCL 7	R ₆	E	Peni	****
24	RCL 5	54	GSB 2	R ₇	p	<u>keg i</u>	<u>sters</u>
25	+	55	+	R _g	P1	R _O	P _a
26	RCL 6	56	√x	$R_{m{g}}$	$\Delta t = \Delta u$	R_{1}	m _a
27	÷	57	RCL D	R_{A}	^t o	R_2	m _b
28	2	5 <i>8</i>	-	R_{B}	tπ	$R_{\mathcal{J}}$	s
29	÷	59	RCL 7	^{R}c	^u o	R_{4}	E
30	STO D	60	RCL B	R _D	u_	R_{E}	p

ESTIMATES OF RADIATION DOSES IN TISSUE AND ORGANS AND RISK OF EXCESS CANCER IN THE SINGLE-COURSE RADIOTHERAPY PATIENTS TREATED FOR ANKYLOSING SPONDYLITIS IN ENGLAND AND WALES

LBL--13999

DE82 008423

Jacob I. Fabrikant, M.D., Ph.D. 3,4 Lawrence Berkeley Laboratory, Donner Laboratory University of California, Berkeley Berkeley, California

and

John T. Lyman, Ph.D. Division of Biology and Medicine Lawrence Berkeley Laboratory University of California Berkeley, CA 94720

951 3026

This took was inspect on the root of the r

May

Presented at Scientific Session, Third International Symposium of Radiation Protection-Advances in Theory and Practice, Inverness, Scotland, June 6-11, 1982.

²Supported by the Office of Health and Environmental Research of the U.S. Department of Energy under contract No. W-7405-ENG-48.

³Professor of Radiology, University of California School of Medicine, San Francisco, California.

⁴Mailing address: Donner Laboratory, University of California, Berkeley, California 94720.

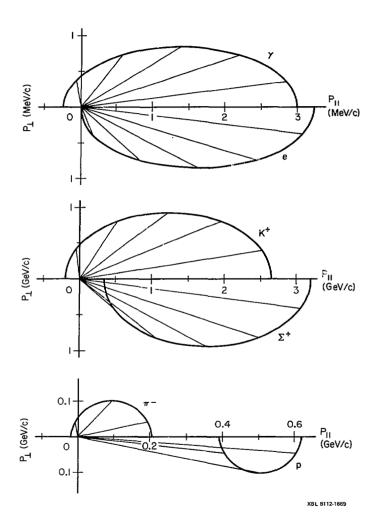


Figure 2

registers; for a decay a + 1+2, set $m_b = 0$. Start the program with a GSB 0 command. When it stops, the values of $\gamma\beta e_1$, $\gamma p'$, p', p', e (the total c.m. energy), and $\cos\theta$ will be in the storage registers, with $\cos\theta = +1$. The components $P_{1||}$ and $P_{1||}$ for this value of $\cos\theta$ will be in the X and Y registers. Each time R/S is pressed, the value of $\cos\theta$ is decreased by 0.1 and the program stops with the corresponding components of \vec{P}_1 in the X and Y registers:

$$(X, Y) = (P_{1||}, P_{11})$$
.

When $\cos\theta$ eventually goes below -1, Error is displayed. Then, to obtain the half ellipse for \dot{P}_2 , interchange m_1 and m_2 in the registers (i.e., now consider particle 2 to be particle 1), and start over with a GSB 0 command. If Error appears when the program is started, then P_a may be below the threshold for the reaction. Error also appears if $P_a = 0$.

The $\rightarrow P$ key changes the rectangular coordinates P_{11} and P_{11} in the X and Y registers to the polar coordinates P_{1} and Θ_{1} , where Θ_{1} is the lab angle at which particle 1 is produced. After the first pass through the program has been made, the lab angle Θ_{1} corresponding to any c.m. angle Θ_{1} may be found by storing the value of $(\cos \theta + 0.1)$ in the R_{9} register, then pressing R/S, then $\rightarrow P: \Theta_{1}$ is in the Y register.

As a test example, consider $\pi^+p + K^+\Sigma^+$ scattering at 3 GeV/c. The input numbers, in GeV/c or GeV, are $P_a=3.0$, $m_a=0.1396$, $m_b=0.9383$, $m_1=0.4937$, and $m_2=1.1894$. Start the program with a GSB 0 command, and it stops with $P_{1\parallel}=2.659$ GeV/c, $P_{1\perp}=0.0$ GeV/c, $\gamma e_1=1.231$ GeV, $\gamma p'=1.427$ GeV/c, p'=0.926 GeV/c, E=2.557 GeV, and $\cos\theta=1.0$. Press R/S and the program stops with $P_{1\parallel}=2.516$ GeV/c, $P_{1\perp}=0.404$ GeV/c ($P_1=2.548$ GeV/c, $P_1=9.113^\circ$), and $\cos\theta=0.9$.

Program: ELLIPSE

Step	Keys	Step	<u>Keys</u>
			
	LBL 0		RCL 3
	RCL 0	32	GSB 2
03	RCL 2	33	-
04	RCL 1	34	CHS
	RCL 0	35	\sqrt{x}
06	GSB 2	36	ST0 x 6
07	+	<i>37</i>	ST0 7
08	\sqrt{x}	3 8	1
09	+	39	ST0 9
10	ST0 6	*40	LBL 1
11	GSB 2	41	RCL 7
12	-	42	RCL 9
13	ST0 8	43	cos ⁻¹
14	÷	44	sin
15	RCL 4	45	х
16	RCL 3	46	RCL 6
17	GSB 2	47	RCL 9
18	-	48	x
19	x	49	RCL 5
20	+	50	+
21	2	51	R/S
22	*	5 2	•
23	ST0 5	53	1
24	x	54	STO - 9
25	÷	55	GTO 1
26	RCL 8	*56	LBL 2
27	\sqrt{x}	57	x ²
28	ST0 8	58	x≠y
29	ST0 ÷ 6	59	x ²
30	x	60	RTN

Registers

R_{O}	Pa
R_1	m _a
R_2	m _b
R_3	m ₁
R_4	m ₂
₽ ₅	γβe l
R_6	γp'
^R 7	p'
^R 8	E
R_g	cos θ

DALITZ RECTANGULAR

The DALITZ RECTANGULAR program calculates coordinates of points on the boundary of the rectangular Dalitz plot for the decay of a particle or system of mass M into three particles having masses $\mathbf{m_1}$, $\mathbf{m_2}$, and $\mathbf{m_3}$. The x and y coordinates of the plot are $\mathbf{m_{12}}^2$ and $\mathbf{m_{13}}^2$, where $\mathbf{m_{11}}$ is the invariant mass of the system of particles i and j.

Let m_{12}^2 + and m_{12}^2 + be the smallest and largest values that m_{12}^2 attains anywhere on the boundary. The values of m_{12}^2 for which boundary coordinates (m_{12}^2, m_{12}^2) are going to be calculated are

$$m_{12}^2$$
, m_{12}^2 + Δ , m_{12}^2 + 2Δ , ..., m_{12}^2 ,

where Δ is a step size in m_{12}^2 . (If, however, $m_1 = m_2 = 0$, see below.) This sequence repeats over and over, first for points along the upper boundary (the boundary on which m_{13}^2 is larger), then for points along the lower boundary, then back to the upper boundary, and so on.

Store the four masses and the step size Δ in the indicated registers (Δ in fact need not be chosen until after the first pass is made through the program, when m_{12}^2+ and m_{12}^2+ are available). Start the program with a GSB 0 command. When it stops, the values of m_{12}^2+ and m_{12}^2+ will be in the indicated registers, and the coordinates $\begin{pmatrix} 2 \\ m_{12} \end{pmatrix}$, m_{13}) of the leftmost point of the boundary will be in the X and Y registers. From then on a new boundary point is obtained each time R/S is pressed:

$$(x, y) = (m_{12}^2, m_{13}^2)$$
.

The step size Δ , or the current value of m_{12}^2 (in R_5), or the boundary one is on (±1 in R_g) may be changed at any time (Δ should not be made negative). No error message is given if M is less than $(m_1 + m_2 + m_3)$.

There is a special case. If $m_1 = m_2 = 0$, then m_{12}^2 is zero, and to avoid a division by zero the program skips the first in the sequence of m_{12}^2 values given above (i.e., the sequence becomes Δ , 2Δ , ..., m_{12}^2). In this case, a nonzero value of Δ does need to be stored before the first pass is made. Now the upper and lower

boundaries end at the points $(0, M^2)$ and $(0, m_3^2)$, and the segment of the m_{12}^2 axis between these points is part of the boundary.

As a test example, consider the decay of a 3-GeV system into $\tilde{K}^0\pi^-p$. The input masses, in GeV, are M = 3.0, $m_1=0.4977$, $m_2=0.1396$, and $m_3=0.9383$. After the GSB 0 command, the program stops with the coordinates (0.406, 7.152) — here and below the units are GeV². The values of m_{12}^2 † and m_{12}^2 † are 0.406 and 4.251, which suggests that a reasonable value of Δ , at least to start with, might be 0.2. Storing this and pressing R/S gives the coordinates (0.606, 8.179). Pressing R/S again gives the coordinates (0.806, 8.119). Another test is to set M = 1.0, $m_1=m_2=m_3=0$, and $\Delta=0.2$. The boundary is the triangle with vertices at (0, 0), (1, 0), and (0, 1).

Program: DALITZ RECTANGULAR

<u>Step</u>	<u>Keys</u>	Step Keys	Step	Keys
*01	LBL 0	<i>31</i> X	61	2
02	RCL 1	32 -	62	+
03	RCL 2	33 RCL 9	63	ST0 + 9
04	+	34 GSB 6	64	GSB 6
05	x ²	35 RCL 5	65	ABS
06	S70 6	<i>36</i> R/S	66	√x
07	RCL 0	*37 LBL 3	67	RTN
08	RCL 3	38 RCL 7	*68	LBL 6
09	-	39 RCL 5	69	x ²
10	x ²	40 x = y?	70	x⊋y
11	ST0 7	41 GTO 1	71	x ²
12	1	42 RCL 4	72	-
13	STO 8	43 +	73	RTN
*14	LBL 1	44 x ≠y		
15	RCL 6	45 x > y?		
*16	LBL 2	46 GTO 4		
27	ST0 5	47]		
18	$x \approx 0$?	48 CHS		
19	GTO 3	49 ST0 x 8		
20	0	*50 LBL 4	Regis	ters
21	ST 0 9	51 R ↓	R_{O}	M
22	RCL 3	52 GTO 2	R ₁	m ₁
23	RCL 0	*53 LBL 5	R_2	m ₂
24	RCL 3	54 GSB 6	R_{3}	m ₃
25	GSB 5	55 RCL 5	R ₄	Δ
26	RCL 1	56 +	R ₅	m ₁₂
27	RCL 2	57 ABS	R ₆	m _{1.2} +
28	RCL 1	58 RCL 5	R ₇	m12/2
29	GSB 5	59 √ x	R ₈	±1
30	RCL 8	60 ÷	R_g	used

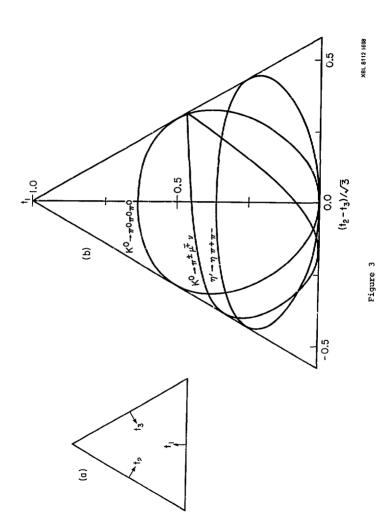
DALITZ TRIANGULAR

The DALITZ TRIANGULAR program calculates coordinates of points on the boundary of the normalized triangular Dalitz plot for the decay of a particle or system of mass M into three particles having masses $\mathbf{m_1}$, $\mathbf{m_2}$, and $\mathbf{m_3}$. Let $\mathbf{T_i}$ be the kinetic energy of particle i in the rest frame of M, $\mathbf{Q} = (\mathbf{T_1} + \mathbf{T_2} + \mathbf{T_3})$ be the total kinetic energy released by the decay, and $\mathbf{t_i} = \mathbf{T_i}/\mathbf{Q}$ be the fraction of \mathbf{Q} taken up by particle i. Then the fractions $\mathbf{t_1}$, $\mathbf{t_2}$, and $\mathbf{t_3}$ are the distances measured inward from the base, left-hand side, and right-hand side of an equilateral triangle whose altitude is unity. Figure 3 (a) shows a schematic of the plot, and fig. 3 (b) shows some boundaries drawn using the program. Each boundary touches each side of the circumscribing triangle at one point. A boundary has a sharp corner where it touches a side if the particle whose kinetic energy is measured from that side has a mass of zero.

To avoid actually having to plot in triangular coordinates, the boundary points are calculated using a rectangular system having its origin at the center of the base: the <u>vertical</u> coordinate is t_1 and the <u>horizontal</u> coordinate is $(t_2-t_3)/\sqrt{3}$. Let t_1 † be the largest value that t_1 attains anywhere on the boundary. The values of t_1 for which boundary coordinates are going to be calculated are

where Δ is a step size in t_1 . (If, however, $m_2 = m_3 = 0$, see below.) This sequence repeats over and over, first for points along the right side of the boundary, then for points along the left side, then back to the right side, and so on.

Store the four masses and the step size Δ in the indicated registers (a value of 0.05 or 0.1 for Δ is about right), and start the program with a GSB 0 command. When it stops, the values of t_1^+ and Q will be in the indicated registers, and the coordinates $(t_1, (t_2-t_3)/\sqrt{3})$ of the lowermost point of the boundary will be in the X and Y registers. From then on a new boundary point is obtained



each time R/S is pressed:

$$(x, y) = (t_1, (t_2-t_3)/\sqrt{3})$$

The step size Δ , or the current value of t_1 (in R_5), or the boundary one is on (± 1 in R_8) may be changed at any time (Δ should not be made negative). No error message is given if M is less than $(m_1 + m_2 + m_3)$.

There is a special case. If $m_2=m_3=0$, then to avoid a division by zero the program skips the last in the sequence of t_1 values given above. In this case, the upper boundary of the plot is the horizontal line segment at $t_1=t_1^+$ that crosses from one side to the other of the circumscribing triangle.

As a test example, consider $K^0 \to \pi^\pm \mu^\mp \nu$ decay. The input masses, in GeV, are M = 0.4977, m_1 = 0.1396, m_2 = 0.1057, and m_3 = 0; and set Δ = 0.05. When, after the GSB 0 command, the program stops, the coordinates of the first point to be plotted are (0.0, -0.170); and t_1^+ = 0.466 and Q = 0.2524 GeV. Press R/S and the next coordinates are (0.05, -0.040). Press it again, and the coordinates are (0.10, +0.020). Another test is to set M \neq 0 and m_1 = m_2 = m_3 = 0. Then t_1^+ = 0.5, Q = M, and the boundary is the triangle with vertices at the centers of the sides of the circumscribing triangle.

Program: DALITZ TRIANGULAR

<u>Step</u>	<u>Keys</u>	<u>Step</u>	<u>Keys</u>	<u>Step</u>	<u>Keys</u>	<u>Step</u>	Keys
*01	LBL 0	31	RCL 1	61	-	91	ABS
02	RCL 0	32	GSB 5	62	-	92	√x
03	RCL 1	33	STO A	63	RCL 7	93	RTN
04	-	34	GSB 5	64	3		
05	ST0 7	35	x = 0?	65	\sqrt{x}		
06	RCL 2	36	GTO 3	66	x		
07	RCL 3	37	STO B	67	*		
08	+	38	RCL 2	6 0	RCL 5		
09	STO - 7	39	x ²	69	R/S		
10	+	40	RCL 3	*70	LBL 3		
11	RCL 0	41	x ²	71	RCL 6		
12	2	42	-	72	RCL 5		
13	x	43	RCL B	73	x = y?		
14	÷	44	÷	74	GTO 7		
15	ST0 6	45	ST0 x 9	75	RCL 4		
16	1	46	+	76	+		
17	ST0 8	47	RCL 2	77	x≵y		
*18	LBL 1	48	2	78	x > y?	Regis	ters
19	0	49	x	79	GTO 4	^R o	M
*20	LBL 2	50	GSB 5	80	1	R,	m ₁
21	ST0 5	51	RCL A	81	CHS	R_2	m ₂
22	RCL 7	52	x	82	ST0 x 8	R_3	m ₃
2 3	x	5 <i>3</i>	RCL 8	*83	LBL 4	R ₄	Δ
24	RCL 1	54	X	84	R∔	R_{5}	t ₁
25	+	55	RCL 9	85	GTO 2	R ₆	t ₁ †
26	RCL 0	56	+	*86	LBL 5	R ₇	Q
27	x≠y	<i>57</i>	RCL B	87	x ²	R _B	±1
28	<u>-</u>	58	÷	88	x≠y	R ₉	used
	ST0 9	5 9	RCL 2	89	x^2	$R_{\overline{A}}$	used
30	LAST x	60	RCL 3	90	-	R_{B}	used