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Franck, Jack Vernon, M.S. Thesis

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UNIVERSITY OF CALIFORNIA

Radiation Laboratory

A CAVITY-STABILIZED OSCILLATOR WITH TWO FEEDBACK CIRCUITS

BERKELEY, CALIFORNIA

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A CAVITY -STABILIZED OSCILLATOR WITH TWO FEEDBACK CIRCUITS

Jack Vernon Franck

(M.S. Thesis)

September 28, 1955

A CAVITY -STABILIZED OSCILLATOR WITH TWO FEEDBACK CIRCUITS

Contents

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-2-

A STUDY OF A CAVITY-STABILIZED OSCILLATOR WITH TWO FEEDBACK CIRCUITS

Jack Vernon Franck

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Radiation Laboratory University of California Berkeley, California

September 28, 1955

ABSTRACT

The pre -exciter oscillator used on the Berkeley 32 -Mev proton linear accelerator is studied analytically and experimentally. The oscillator operates on 202.55 megacycles, using a single power tetrode, and is equivalent to a two -stage amplifier with two feedback circuits. Feedback circuit number one is a broadly tuned circuit around the first stage. Feedback circuit number two is a sharply tuned circuit around both stages. Feedback circuit No. 2 includes the cavity resonator. The object is a reduction in the number of high-frequency power tubes required in a conventional system. The frequency bandwidth over which the oscillator will "pull" to the resonator frequency is calculated and is found to be in good agreement with the measured value. Graphical calculators have been devised and used to reduce the computation to a minimum. The analysis of a self-oscillating system with two feedback circuits is believed to be new.

I. INTRODUCTION ..

A Cavity-Stabilized Oscillator

This paper is a study of a cavity-stabilized oscillator with two feedback circuits. The oscillator studied is used to supply the radiofrequency pre-excitation to the high-Q cavity of the Berkeley 32-Mev^{*} proton linear accelerator. ³

Because this particular oscillator is used to supply pre -excitation to the "Linac" cavity it has been defined as the "pre -exciter oscillator," and will be so designated throughout this paper.

Purpose of the Pre-Exciter Oscillator

The pre -exciter oscillator as used on the Linac fulfills several requirements. It drives the radiofrequency load cavity through the multipactor region. 3 it selects the correct mode of oscillation, and it supplies the low level of radiofrequency voltage necessary.to start the main power oscillators oscillating.

Requirements of the Pre -Exciter Oscillator

The pre -exciter oscillator should be frequency-stable and should deliver at least several percent of the normal output of the main power oscillators.

Design Specifications for the Pre -Exciter Oscillator

The pre -exciter oscillator under study was designed to conform to the following specifications:

In addition, during the "on" time of the main power oscillators the pre -exciter plate -circuit radiofrequency voltage is approximately

million electron volts

five times the voltage existing during the pre -excitation period. This is shown by the following relations. (For definition of symbols see table: $p. 74$):

$$
P_{\phi} = \frac{V_{ps}^2}{2 R_{ps}} , \quad V_{ps} = \sqrt{2 P_0 R_{ps}^2}
$$

$$
\frac{V_{ps}(\text{final level})}{V_{ps}(\text{pre-exc. level})} = \sqrt{\frac{2400}{100}} = \sqrt{24} = 5.
$$

The power output of the power oscillators is approximately 2400 kilowatts 3 at a pulse length of 600 microseconds and a repetition rate of 15 pulses per second.

The pre-exciter plate circuit has to be able to withstand this high voltage if no "transmit -receive" switch is to be used. From Appendix B, page 56 , this voltage is approximately

> 5 V _{ps}(pre-exc. level) = 5×4120 $= 20,600$ volts (crest).

Description of Pre -Exciter Oscillator

This pre-exciter oscillator is unique in that it has two feedback circuits. That is, it is a self-contained oscillator (cathode-gridscreen circuit) to which has been added a second feedback circuit. The line drawing, Fig. 1, shows the two feedback circuits.

In the particular pre -exciter under study, the complete function is filled by a single tetrode tube $(4W20,000A)$ operating as a self-oscillator (cathode -grid-screen circuit) electron-coupled to the plate output circuit. A tetrode tube was selected for this particular installation since electrically and mechanically this results in the most economical and compact unit. In other applications separate tubes for each amplifier stage might be desirable. The resonant load is the cavity resonator of the Berkeley proton linear accelerator. The complete pre -exciter and its internal parts are shown photographically in Figs. 2 and 3. The U-shaped transmission line protruding from the pre-exciter oscillator

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Fig. 1. Pre-exciter circuit.

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ZN-1396

Fig. 2. Pre -exciter oscillator.

ZN-1395

Fig. 3. Pre -exciter internal parts.

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is used solely for neutralizing the plate to grid capacity.

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Method of Operation

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., ~-.- In actual operation, the pre -exciter oscillator is turned on about 100 microseconds prior to turning on the main power oscillators. During this time the pre -exciter builds the rf cavity voltage up to the 100 kilowatt level.

When the pre-exciter is first turned on there is no energy in the rf cavity. However, the cathode -grid-screen circuit starts oscillating vigorously in a relatively few cycles. This seli -oscillating portion of the circuit is represented in the line drawing, Fig. l_r by amplifier No. 1 and feedback circuit No. l. Since the tube is neutralized (see Fig. 4) the oscillating grid circuit is unaffected by voltage in the plate output circuit. Feedback circuit No. 2 is very loosely coupled to the grid circuit and therefore affects the level of oscillation very little. The self -oscillating cathode -grid -screen circuit delivers a fully modulated space current to the output plate circuit. This is independent of whether or not the oscillating frequency is the same as the natural resonant frequency of the plate load circuit. If the grid circuit is oscillating near the natural resonant frequency of the load, some voltage, however small, will be developed as a result of the flow of plate current. Feedback circuit No.: 2. couples a portion of this load voltage back to the oscillating circuit. This feedback acts on the grid circuit as a frequencycorrecting voltage in such a manner that the frequencyof oscillation tends to be "pulled"·toward the natural resonant frequency of the cavityresonator load circuit. The amount of pulling and the manner in which this pulling occurs is the subject of study of this paper.

 \mathcal{L}^{max}

 \sim 100 \pm

Fig. 4. Schematic diagram of pre-exciter oscillator.

 $\mathcal{L}_{\mathcal{A}}$

 $\mathcal{L}_{\mathcal{A}}$ is a simple polynomial of the set of the set of the set of the set of \mathcal{A} \sim \sim **Contract** $\mathcal{L}^{\mathcal{L}}$ is a space of the space $\mathcal{P}_{\rm{max}}$ $\sim 10^{11}$ km $^{-1}$

II. ANALYSIS OF PRE-EXCITER OSCILLATOR

The following analysis uses equivalent lumped constants to repre sent the circuitry of the pre-exciter oscillator. The self-oscillating grid circuit and the feedback circuit are each represented by a simple generator with the appropriate internal impedance. Thevenin's theorem⁷ is the basis of this representation.

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The equivalent circuit of the oscillator is defined by the geometry of the experimental equipment. The geometry is, in turn, defined by the conditions necessary for proper operation of the tube. A typical set of operating conditions for the tube is evaluated in Appendix B. The schematic diagram of the oscillator is showh in Fig. 4. A schematic diagram using lumped constants is shown in Fig. 5. The equivalent lumped-constant circuit is given in Fig. 6.

Some of the equivalent circuit voltages and impedances will be expressed in terms of the tube currents and plate -load impedance. The dependence of the plate -load impedance on frequency has been deter·· mined in Appendix D.

The network representing the feed back circuit is considered first.

The voltage V_{fb} fed back to the grid circuit can be expressed in terms of the plate current and plate-load impedance of the tube:

$$
V_{ps} = I_{ps} Z_{ps}
$$
 (1)

Transmission line No. 1 is constructed of two quarter-wave sections of different impedance and is connected to the plate line at a voltage h V_{ps} , where h ≤ 1.0 . The voltage on loop No. 1 is then

$$
V_{\ell l} = \frac{Z_{0l1}}{Z_{0l}}
$$
 h V_{ps}. (2)

Since the magnetic flux density is a constant along the side of a cylindrical resonator excited in the electric 010 mode, the voltages across all loops along the side are proportional to their areas:

-11-

MU-10208

$$
V_{\ell 1} / V_{\ell 2} = A_{\ell 1} / A_{\ell 2}
$$

$$
V_{l1} = (A_{l1} / A_{l2}) V_{l2}.
$$

''

It is now convenient to define a coupling coefficient, c, which is determined by the adjustment of the network containing: transmission line No. 2:

$$
V_{fb} = c V_{l2}
$$
 (4)

 (3)

Writing Eq. (4) in terms of'Eqs. (3), (2), and (1) gives

$$
V_{\text{fb}} = h \circ \frac{A_{\ell 2}}{A_{\ell 1}} \frac{Z_{011}}{Z_{01}} \left[I_{\text{ps}} Z_{\text{ps}} \right]
$$
 (5)

The internal impedance Z_{fb} of feedback circuit No. 2 can now be evaluated. A power equality may be used to determine Z_{fb} in terms of Z_{ps} . At the natural resonant frequency of the load, the power generated by the tube is equal to the power delivered to the resonator via loop No. 1:

$$
P_{o} = \frac{V_{ps}^{2}}{2 R_{ps}} = \frac{V_{l1}^{2}}{2 Z_{l1} (f_{0})}.
$$
 (6)

From Eq. (6), by using the universal resonance curves, 7 we obtain (see Appendix D.)

$$
Z_{\ell 1} = \left[\frac{V_{\ell 1}}{V_{ps}}\right]^2 Z_{ps} \t\t(7)
$$

Using Eq. (2) , we have

/

$$
Z_{\ell 1} = \left[\frac{Z_{011} \quad h}{Z_{01}}\right]^2 \quad Z_{\text{ps}} \tag{8}
$$

Since the voltage of loop No. 2 is the same as if the power were being put into the tank through loop No. 2, the relations between the voltages and the impedances can be again determined as in Eqs. (6) and (7):

$$
v_{\ell 1}^{2} / z_{\ell 1} = v_{\ell 2}^{2} / z_{\ell 2} ,
$$

\n
$$
z_{\ell 2} = (v_{\ell 2} / v_{\ell 1})^{2} z_{\ell 1} .
$$
\n(9)

Using Eqs. (3) and (8), we get

 \blacktriangledown

$$
Z_{\ell 2} = (A_{\ell 2} / A_{\ell 2})^2 (Z_{011} h / Z_{01})^2 Z_{ps}
$$
 (10)

and from the coupling system schematic diagram (Fig. 7) it is evident that the feedback impedance has the form

$$
Z_{\text{fb}} = F(Z_1, Z_2, Z_3, Z_{\ell 2}, f). \tag{11}
$$

The reactance Z_1 , is adjusted in magnitude, by making trans-
on line No. 2 slightly longer or shorter than $\frac{n\lambda}{Z}$. Similar mission line No. 2 slightly longer or shorter than $\frac{n\lambda}{Z}$. Similarly, Z_3 is a reactance the magnitude of which is determined by adjusting the quarter -wave line (associated with loop No. 2) to be slightly longer or shorter than $\frac{11}{2}$. Z_2 is the equivalent shunt resistance of feedback transmission $\frac{4}{1}$ line No. 2.

This network has been interpreted from the experimental equipment. Once these impedances have been determined, Z_{fb} is found by solving the parallel-series network:

$$
Z_{\text{fb}} = \frac{\begin{bmatrix} Z_{12} + Z_3 \end{bmatrix} \begin{bmatrix} Z_1 & Z_2 \\ Z_1 + Z_2 \end{bmatrix}}{\begin{bmatrix} Z_{12} + Z_3 + \frac{Z_1}{Z_1 + Z_2} \end{bmatrix}}
$$
(12)

The voltage V_{fb} (Eq. 5) and the impedance Z_{fb} (Eq. 12) are the open-circuit voltage and internal impedance respectively of a simple generator which by Thévenin's theorem can be used to represent

MU-10209

the feedback circuit. The voltage V'_{sg} and the resistance r'_{sg} represent the cathode-grid-screen circuit in the constant-voltage form of equivalent circuit for a vacuum tube. These are shown in Fig. 6.

The shunt resistance $R_{\rm s}^{\rm t}$ and the capacity C' of the equivalent grid circuit are referred to the point at which the feedback line is coupled. The voltage at this point is V_{02} (see Fig. 22).

The value of R^{\dagger}_{s} can be determined by solving for the dissipated power in the relations used to calculate the Q_t of the grid circuit. Using the relation for Q_t , which includes the grid drive power, gives for W_d (Appendix A)

$$
W_d = \frac{\lambda r_s}{32 \pi} \left[8.46 \right] \quad V_{02}^2 = \frac{V_{02}^2}{2 R_s'}
$$

Expressing this in terms of a shunt resistance, we get

$$
R_{s}^{\dagger} = \frac{32 \pi}{2 \lambda r_{s}} \left[\frac{1}{8.46} \right] = 563 \text{ ohms}
$$
 (13)

The value of C^{\dagger} may be determined from the Q_t (Appendix A) and the value determined for R_S' (Eq. (13)):

$$
Q_{t} = \frac{\omega W_{s}}{W_{d}} = \frac{\omega \frac{C' V_{02}^{2}}{2}}{\frac{V_{02}^{2}}{2 R_{s}^{t}}} = \omega C' R_{s}^{t}.
$$

Solving for C', we obtain

$$
C' = \frac{Q_t}{\omega R_s'} = \frac{49.2}{2 \pi 202.55 \times 10^6 \times 563} ,
$$

$$
C' = 68.5 \times 10^{-12} \text{ farads.}
$$
 (14)

The value of r'_{sg} (see Fig. 6) is determined from the operating point for the tube (Appendix B) by transforming $r_{\rm sc}$ to the feedback point (V_{02}) :

$$
\mathbf{r}_{\text{sg}} = \frac{\mathbf{e}_{\text{sm}}}{\mathbf{I}_{\text{sg}}} = \frac{2540}{70.1} = 36.6 \text{ ohms},
$$
\n
$$
\mathbf{r}_{\text{sg}}' = \mathbf{r}_{\text{sc}} \left[\frac{V_{05}}{V_{\text{sc}}} \right]^2 \left[\frac{V_{02}}{V_{05}} \right]^2 \tag{15}
$$

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The grid-to-screen voltage, $V_{gg} = V_{05}$, is the sum of the gridcathode and screen-cathode voltages, since the screen voltage is 180° out of phase with the grid voltage:

 $-17-$

$$
V_{sg} = V_{05} = V_{g} + V_{sc}.
$$

From the tube operating point,

$$
V_{\text{sc}} = 460 \text{ volts},
$$

\n
$$
V_{\text{g}} = E_{\text{g}} + e_{\text{gm}} = 1800 \text{ volts},
$$
 (16)

therefore,

...

$$
V_{sg} = V_{05} = 460 + 1800 = 2260 \text{ volts}, \qquad (17)
$$

$$
\frac{V_{sg}}{V_{sc}} = \frac{V_{05}}{V_{sc}} = \frac{460 + 1800}{460} = \frac{2260}{460} = 4.92
$$
 (18)

From Appendix A, $V_{05} = 0.667 V_{02}$, so

$$
r'_{sg} = r_{sc} \left[\frac{V_{05}}{V_{sc}} \right]^2 \left[\frac{V_{02}}{V_{05}} \right]^2 = 36.6 \times 4.97^2 \times \frac{1}{0.667^2},
$$

$$
r'_{sg} = 1975 \text{ ohms.}
$$
 (19)

Since this circuit is self-oscillating, the fundamental frequency component of screen current is in phase with the grid-screen voltage. . When it is normalized to voltage V_{02} this means that

$$
V_{sg}^{\dagger} = I_{sg}^{\dagger} (R_{s}^{\dagger} + r_{sg}^{\dagger}). \qquad (20)
$$

The equivalent generator current $\mathbf{I}_{\mathbf{sg}}^{t\rightarrow\infty}$ can be expressed in terms of the space current in the tube by normalizing to the tube voltage. These voltage ratios were calculated in evaluating $r'_{\rm sc}$ (Eqs. (16), (17), (18), (19)):

$$
I'_{sg} = I_{sg} \begin{bmatrix} V_{sc} \\ V_{05} \end{bmatrix} \begin{bmatrix} V_{05} \\ V_{02} \end{bmatrix}.
$$
\n
$$
= I \begin{bmatrix} 1 \\ 667 \end{bmatrix} \begin{bmatrix} 667 \\ 50 \end{bmatrix} = 0.135 I
$$
\n(21)

$$
= 1_{sg} \left[4.92 \right] \left[\cdot ^{00} \right] = 0.133 \left[1_{sg} \right]
$$

This gives for the equivalent generator voltage

 \mathbf{T}

$$
V'_{sg} = 0.135 I_{sg} (478 + 1975),
$$

\n $V'_{sg} = 331 I_{sg}.$ (22)

The equivalent generator voltage can be expressed in terms of $\, {\rm I}_{\rm ps} \cdot$ Using the following ratio qbtained from the calculation of tube operation, Appendix B, we get

$$
\frac{I_{sg}}{I_{ps}} = \frac{70.1 \text{ amperes}}{55.0 \text{ amperes}} = 1.275;
$$
 (23)

the equivalent generator voltage is

$$
V'_{sg} = 331 \times 1.275 I_{ps}
$$

\n
$$
V'_{sg} = 422 I_{ps}
$$
 (24)

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All parts of the equivalent circuit (Fig. 6) have now been evaluated for Case I. The values are shown on the equivalent circuit diagram, Fig. 8. The frequency-dependent values are tabulated in Table I.

The superposition theorem $\begin{bmatrix} 7 & 2 & 1 \end{bmatrix}$ can be used here to solve for the network currents i_1 , i_2 , and i_3 , from which the effect of the feedback network can be determined. This would have to be subjected to the restriction that $i_1 = I_{sg}^{\dagger}$ must be in phase with V_{sg}^{\dagger} , since this part of the circuit is a self -oscillator. It is to be noted, however, that the impedance of the equivalent cathode -grid-screen circuit is at most a few percent of the combined feedback and coupling impedances $(Z_{fb} + Z_c)$. In this case $I_{gg'}^i$, to a good approximation, flows entirely through R_s^i .

The feedback current i_3 , which by design is very small compared to I'_{ps} , will in general have an inphase and a quadrature component with respect to I'_{ps} . The quadrature component, as stated before, must flow only through the reactances of the equivalent cathode -grid-screen, resonant circuit. The inphase component of i_3 must divide between R'_s and r'_{sg} as a consequence of the use of the superposition theorem, subject to the approximation that $i_1 = I'_{sg}$ flows only through R'_s .

The current i_3 can be evaluated as a function of frequency as follows:

$$
i_3 = \frac{-V_{fb} + V'_{rs}}{Z_{fb} + Z_c} = \frac{-V_{fb} + \left[\frac{R'_s}{R'_s + r'_{sg}}\right]V'_{sg}}{Z_{fb} + Z_c}
$$
 (25)

All values needed are constant or have previously been evaluated. V_{fh} has been calculated as a function of frequency and is tabulated in Table I for Case I. We have, therefore,

$$
i_3 = \frac{-V_{fb} + 93.6 I_{ps}}{2500 + j4000}
$$

= $\left[-V_{fb} + 93.6 I_{ps} \right] \left[1.12 - j1.80 \right] 10^{-4}$ (26)

The current i₃ is tabulated as a function of frequency in Table II. The equivalent shunt capacity C'', representing the quadrature component of the current i₃, is also listed. This is obtained from the following

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Fig. 8. Case I: equivalent circuit.

Table II and the second of the second of

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 $\frac{1}{2}$

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 $-24-$

relation (the inphase component of i_3 is completely negligible in comparison with the inphase current I'_{sg} :

C^{II} =
$$
\frac{i_{3} (quadrature part)}{W\left[\frac{R'_{s}}{R'_{s} + \frac{r}{s_{g}}}\right] 422 I_{ps}}
$$

$$
= \frac{i_{3} (quadrature part)}{W 93.6 I_{ps}}
$$

The capacity C^{111} is the effective value of the capacity resulting from the feedback network after correcting the value of this capacity so that the oscillating circuit receives zero frequency correction when it is at the natural resonant frequency (f_0) of the load.

 (27)

I

The frequency shift that occurs because the capacity C^{111} is in parallel with the grid circuit is evaluated as follows. Let

$$
f = \frac{1}{2\pi \sqrt{LC}}
$$
 $\frac{1}{2\pi \sqrt{LC}}$ $\frac{1}{2\pi \sqrt{LC (C + \Delta C)}}$ $\frac{1}{2\pi \sqrt{LC \sqrt{1 + \frac{\Delta C}{C}}}}$

Since $\Delta f \ll f$, the following approximation can be made,

$$
\frac{1}{\sqrt{1+\frac{\Delta C}{C}}} \leq \frac{1}{1+\frac{\Delta C}{2C}} \leq 1 - \frac{\Delta C}{2C}
$$

and therefore

$$
f + \Delta f = f \left[1 - \frac{\Delta C}{2 C}\right],
$$

or, for $f = f_0$, $\Delta C = C^{(1)}$, and $C = C^1$ (see Fig. 8),

$$
\Delta f = \left[-\frac{C^{111}}{2C^1} \quad f_0 \right]. \tag{28}
$$

 $-25-$

The values of Δf are tabulated in Table II. The results of Case I are displayed in Fig. 9 along with the experimental values from Appendix F (Figs. 28 and 29). Figures 10 and 11 are merely the graphical solution of the relation +f = Δf_0 - Δf . Figure 9 is a plot of Δf_0 vs. f.

The power output of the oscillator as a function of frequency can be obtained by use of the relation

 $\ddot{}$

$$
P_o(f) = \frac{I_{ps}^2 R_{ps}(f)}{2}
$$
 (29)

In this equation, I_{ps} is a constant (see p. 9) and $R_{ps}(f)$ has been evaluated in Appendix D. The power output as a function of frequency is tabulated in Table III and is plotted in Fig. 9 along with the experimental results of Appendix F.

Table III

 $-26-$

Fig. 9. Case I: Frequency of oscillation and power output as functions of the frequency to which the oscillator is tuned.

Fig. 10. shift. Case I: Relation between frequency of oscillation and frequency (lower frequencies)

..

 $\mathcal{I}_1 \rightarrow \mathcal{I}_2$.

Fig. **li.** shift Case I: Relation between frequency of oscillation and frequency (higher frequencies).

'

B. Case II

..

The conditions in Case II are identically the same as for Case I, except that much of the tedious vector computation is bypassed by using several graphical calculators. Illustrations of these calculators are found on pages 78, 79, and 80. Their use was suggested and made possible by the fact that the internal impedance of the equivalent generator representing the feedback circuit is very large compared to the shunt impedance of the cathode -grid -screen circuit. The circular calculator is an expression of the fact that the locus of the impedance vector for a parallel resonant circuit is a circle. This is demonstrated in Appendix E.

To solve Case I with the graphical calculator, the grid circuit voltage V_{rs}^{i} is drawn to unit length as indicated on graphical calculator No. 1. Calculator No. 1 is then rotated so that the feedback voltage has the desired phase angle with respect to the grid voltage $\left. \mathrm{V}_{\mathrm{rs}}^{\mathrm{r}}\right\vert$ at the natural resonant frequency of the load circuit (f_0) . Calculator No. 2, on which has been drawn the impedance line representing the total impedance in the feedback circuit $(Z_{fb} + Z_c)$ in arbitrary units, is placed so that the origin is coincident with the head of the vector representing the feedback voltage V_{fb} . The magnitude of V_{fb} is relative to the unit length chosen for V'_{rs} . The impedance line is made to pass over the head of the vector representing the grid circuit voltage V_{rs}^i . The relative magnitude of the feedback current (i_3) , and its phase, can be read from calculator No.2 by use of a protractor or calculator No. 3. Calculator No. 3 can be used to read the quadrature component of the feedback current directly from calculator No. 2 without hothering with the relative magnitude.

The feedback current for other frequencies is obtained in the same manner as above. The magnitude and phase of the feedback voltage as a function of frequency are read directly on calculator No. 1. The frequency is given in terms of the Q of the load. The Δf can be readily determined by referring to Appendix D.

The relative magnitude of the feedback voltage can be determined by referring to Table I, and Eq. (25). Then we have

$$
\left|\frac{V_{\text{fb}}}{V_{\text{rs}}^i}\right| = \frac{104.}{93.6} = 1.11
$$
 (30)

A single numerical computation is required to find the normalizing constant for the current. This is done by evaluating Eqs. (25} and (26} for the natural resonant frequency of the load impedance (cavity resonator),

 $-31-$

$$
i_3 = \frac{V'_{rs} - V_{fb}}{Z_{fb} + Z_c} = \frac{93.6 I_{ps} - V_{fb}}{2500 + j4000}
$$

= $\left[93.6 I_{ps} - 104 \angle 30^{\circ} I_{ps}\right] \left[1.12 - j1.80\right] 10^{-4}$
= $\left[97.4 + j52.2\right] 10^{-4} I_{ps}$ (31)

The normalizing constant is

$$
N = \frac{52.2 \times 10^{-4}}{21} = 248. \times 10^{-4}
$$
 (32)

The equivalent shunt capacity C'' is by Eq. (27),

$$
C^{11} = \frac{i_3 (quadrature part)}{w 93.6 I_{ps}}
$$

=
$$
\frac{i_3 (quadrature part)}{I_{ps}} \times 0.839 \times 10^{-11}
$$
 (33)

The value of C'" is found as in Case I, page 26, and the Δf is again obtained by use of Eq. (28).

The results as obtained by use of the graphical calculators are tabulated in Table IV, and are plotted in Fig. 12, as obtained from Figs. 13 and 14. For comparison, the analytical results of Case I are also plotted in Fig. 12_.

Frequency shift as a function of the frequency of oscillation Case II.						
Δf_0 (cps)	$i_3 \times 1/N$ (am)	$i_{\alpha} \times 1/N$ $(\texttt{quad}, \texttt{part})$ (amp)	(quad.png) (am)	C^{11} $(\mu\mu f)$	C^{111} $(\mu\mu f)$	Δf (kc)
$\mathbf{0}$	$0.47/$ +27.0 ^o I ps	$0.21/190^{\circ}$ DS	$+52.2 \times 10^{-4} / +90^{\circ}$ рs	$+0.044$	$\mathbf{0}$	Ω
+470	$0.71/\pm11.0^{\circ}$ DS.	$0.14/190^{\circ}$ DS	$+34.8 \times 10^{-4} / +90^{\circ}$	$+0.029$	-0.015	$+21.8$
-470	$0.18/+39.5^{\circ}$ ps	$0.12/190^{\circ}$ рs	$+90^{\circ}$ $+29.8 \times 10^{-4}$ ps	$+0.025$	-0.019	$+27.5$
+705	$0.80 \angle 13.5^{\circ}$ ps	$0.05/\pm90^{\circ}$ DS	$+12.4 \times 10^{-4}$ / $+90^{\circ}$ рs	+0.010	-0.033	$+49.3$
-705	$0.05/$ +12.7 ^o ps	$0.01/190^{\circ}$ ps	$+2.5 \times 10^{-4} / +90^{o}$ рs	$+0.002$	-0.042	$+62.3$
$+1410$	$0.93/ - 14.5^{\circ}$ DS.	$0.23/ -90^{\circ}$ рs	$-57.1 \times 10^{-4} / +90^{\circ}$ DS	-0.048	-0.091	$+135.$
-1410	$0.27/ -98.0$ ^o ps	$0.26/ -90^{\circ}$ рs	$1 + 90^{\circ}$ -64.5×10^{-4} рs	-0.050	-0.099	$+145.$
+2820	$0.97/ - 32.5^{\circ}$ DS	$0.51/\underline{-90}^{\circ}$ рs	$\frac{490}{9}$ $-127. x 10^{-7}$ рs	-0.106	-0.149	+220.
-2820	$0.55 / -82.5^{\circ}$ p s	$0.54/ -90^{\circ}$ ps	120° $-137. \times 10^{-7}$ ps	-0.113	-0.157	$+232.$
+5640	$0.93/ -44.5^\circ$ ps	$0.66 / -90^{\circ}$ ps	$-164. \times 10^{-4}$ 490° ps	-0.138	-0.181	$+268.$
-5640	$0.71/-71.0^{\circ}$ ps	$0.68 / -90^{\circ}$ p_{s}	490° -169×10^{-4} рs	-0.142	-0.186	$+276.$
+11280	$0.89/ -51.0$ ^o DS	$0.69/ -90^{\circ}$ ps^{γ}	$\frac{490^{\circ}}{2}$ $-171. \times 10^{-4}$ ps	-0.144	-0.187	$+276.$
-11280	$0.78 / -64.0^{\circ}$ рs	$0.70 / -90^{\circ}$ PS	490° $-174. \times 10^{-7}$ ps	-0.146	-0.191	$+283.$
22560	$0.87/ -54.5^{\circ}$ ps	$0.70 / -90^{\circ}$ ps	$-174. \times 10^{-4} / +90^{\circ}$ рs	-0.146	-0.191	$+283.$
-22560	$0.82 / -61.0^{\circ}$ DЗ.	$0.72 / -90^{\circ}$ рs	$-179. \times 10^{-4}$ /+90 ^o рs	-0.149	-0.193	$+283.$
45120	$0.86/ -56.5^{\circ}$ ps	$0.71/ -90^{\circ}$ DS	$-176. \times 10^{-4}$ /+90 ^o ps	-0.148	-0.191	$+283.$
45120	$0.83/ -59.5^{\circ}$ p s	$0.72 / -90^{\circ}$ p s	$-179. x 10^{-7}$ $+90$ ps	-0.149	-0.193	$+283.$

Table IV

 $-33-$

Fig. 12. Case II: Frequency of oscillation and power output as functions of the frequency to which the oscillator is tuned.

Fig. 13. shift Case II: Relation between frequency of oscillation and frequency (lower frequencies).

.•

Fig. 14. shift Case II: Reiation between frequency of oscillation and frequency (higher frequencies).

"
C. Case III

Case III is solved by using graphical calculators No. 1, No. 2, and No. 3. It differs from Cases I and II only by the nature of the feedback network and the phase of the feedback voltage, V_{fb} . The feedback-coupling impedance Z is a pure resistance instead of a capacitor, as in Cases I and II. The feedback voltage is in phase with the grid voltage, and as a result the feedback-coupling system is greatly simplified. The circuit values below can be identified by referring to Figs. 5, 6, and 7.

> $Z_c = 5000$ ohms (resistance) $Z_1 = \infty$. $Z_2 = \infty$ $Z_3 = 0$
 $h = 0$ $= 0.707$ = 1.0 $/0^{\circ}$ $\frac{1}{2}$ 4.4 z_{01} n_{l2} n_{11} $= 0.364$

The equivalent circuit for Case III is shown in Fig. 15.

The power output is dependent only on the characteristics of the load impedance and the frequency of oscillation, since the tube acts essentially as a constant-current generator (see p. 9°). The operating point of the tube and the load impedance are the same as in Cases I and II and therefore the power output as a function of frequency is again given by Table III.

The relative magnitudes of the feedback and grid voltages are the same as in the two previous cases and are given by Eq. (30).

$$
\left| \frac{\mathbf{V}_{\text{fb}}}{\mathbf{V'}_{\text{rs}}} \right| = \frac{104}{93.6} = 1.11 \quad . \tag{34}
$$

The normalizing constant can be evaluated by first calculating the feedback current at the natural resonant frequency of the load:

 $Z_{fb} + Z_c = 5000 + j0$

$$
i_{3} = \frac{V'_{rs} - V_{fb}}{Z_{fb} + Z_{c}} = \frac{93.6 \text{ I}_{ps} - V_{fb}}{5000}
$$

= $\left[93.6 \text{ I}_{ps} - 104. \left(\frac{0^{o}}{s}\right) \text{I}_{ps}\right] \left[2.0 \times 10^{-4}\right]$
= 20.8 x 10⁻⁴ I_{ps} (35)

The normalizing constant is

$$
N = \frac{20.8 \times 10^{-4}}{0.11} = 189. \times 10^{-4}
$$
 (36)

The equivalent shunt capacity C'' is again found by using Eq. (33). The value of C''' is found as in Case I, and the Δf is again obtained by use of Eq. (28) . (see p. 26).

The results as obtained by use of the graphical calculators are tabulated in Table V, and are plotted in Fig. 16, as obtained with the aid of Fig. 17.

Case IV D.

Case IV differs from Case I only in that the feedback voltage V_{fb} leads the grid voltage V_{rs} by the same phase angle as that by which in Case I it lagged the grid voltage. The feedback-coupling impedance Z_{ζ} is an inductance. The circuit values below can be identified by referring to Figs. 5, 6, and 7.

$$
Z_{c} = j330. (0.259 \times 10^{-6} \text{ henry})
$$

\n
$$
Z_{1} = j37,300 (29.3 \times 10^{-6} \text{ henry})
$$

\n
$$
Z_{2} = 10,000 \text{ ohms resistance}
$$

\n
$$
Z_{3} = -j5000 (0.157 \times 10^{-12} \text{ farad})
$$

\n
$$
h = 0.707
$$

\n
$$
C = 1.0/+30^{\circ}
$$

\n
$$
\frac{Z_{011}}{2} = 4.4
$$

\n
$$
\frac{A_{12}}{2} = 0.364
$$

 $\mathbf{A}_{\ell\,1}$

Case III. Frequency shift as a function of the frequency of oscillation										
		$i_3 \times 1/N$								
Δf_0	$i_3 \times 1/N$	(quad. part)	(quad. part)	C^{11}	C^{111}	Δf				
(cps)	(am)	(am)	\langle amp \rangle	$(\mu\mu f)$	$(\mu\mu f)$	(kc)				
$\mathbf{0}$	$0.11/ - 180.$ ^o I ps	0	Ω	$\mathbf{0}$	$C^{111} = C^{11}$	$\mathbf{0}$				
$+470$	$0.33/+88.5^{\circ}$ рs	$0.33/+90^{\circ}$ I ps	$+62.4 \times 10^{-4} / +90^{0}$ I рs	$+0.052$	$C^{111} = C^{11}$	-77.0				
-470	$0.33/-88.5^\circ$ рs	$0.33/-90^{0}$ ps	$-62.4 \times 10^{-4} / +90^{\circ}$ ps	-0.052	$C^{111} = C^{11}$	$+77.0$				
$+705$	$0.46/$ +74.5 ^o рs	$0.44/190^{\overline{O}}$ рs	$+83.2 \times 10^{-4} / +90^{0}$ рs	+0.070	$C^{111} = C^{11}$	$-104.$				
-705	$0.46 / -74.5^{\circ}$ рs	$0.44/-90^{\circ}$ рs	$-83.2 \times 10^{-4} / +90^{\circ}$ ps	-0.070	$C^{111} = C^{11}$	$+104.$				
$+1410$	$0,71/+50.5^{\circ}$ рs	$0.55/$ +90 ^o ps	490° $+104. \times 10^{-4}$ ps	+0.087	$C^{111} = C^{11}$	$-129.$				
-1410	$0.71/-50.5^{\circ}$ ps	$0.55/ -90^{\circ}$ DS.	$-104. \times 10^{-4}/+90^{\circ}$ ps	-0.087	$C_{111} = C_{11}$	$+129.$				
$+2820$	$0.89/+29.0^{\circ}$ $\mathbf{p}\mathbf{s} \neq$	$0.43/+90^{\circ}$ DS.	$+81.3 \times 10^{-4} / +90^{\circ}$ ps	$+0.068$	$C^{(1)} = C^{11}$	$-100.$				
-2820	$0.89/$ -29.0 ^o DS	$0.43/ -90^{\circ}$	$-81.3 \times 10^{-4} / +90^{o}$ ps	-0.068	$C^{111} = C^{11}$	$+100.$				
+5640	$0.97/+14.8^{o}$ I p s	$0.25/100$ ^o	$+46.3 \times 10^{-4} / +90^{\circ}$ ps :	$+0.039$	$C^{111} = C^{11}$	-57.7				
-5640	$0.97/ -14.8^{\circ}$ ps	$0.25/-90^{\circ}$ ps	$-46.3 \times 10^{-4} / +90^{0}$ ps	-0.039	$C^{111} = C^{11}$	$+57.7$				
+11280	$0.99/ + 7.5^{\circ}$ рs	$0.13/190^{\circ}$ DS.	$+23.6 \times 10^{-4} / +90^{\circ}$ ps	$+0.020$	$C^{111} = C^{11}$	-29.6				
-11280	$0.99/- 7.5^{\circ}$ ps	$0.13/-90^{\circ}$ ps	$-23.6 \times 10^{-4} / +90^{\circ}$ ps	-0.020	$C^{111} = C^{11}$	$+29.6$				
+22560	$0.99/ + 4.0^{\circ}$ DS.	$0.03/190^{\circ}$ ps	490° $+ 5.9 \times 10^{-4}$ ps	$+0.005$	$C^{111} = C^{11}$	-7.4				
-22560	$0.99/- 4.0^{\circ}$ ps	$0.03/190^{\circ}$ DS -	$-5.9 \times 10^{-4} / +90^{\circ}$ ps.	-0.005 =	$C^{111} = C^{11}$	$+ 7.4$				
+45120	$0.99/ + 2.0^{\circ}$ `ps	$0.02/190$ ^o ps	490° $+2.8 \times 10^{-4}$ `ps	$+0.002$	$C^{111} = C^{11}$	-3.0				
-45120	$0.99/ - 2.0^{\circ}$ 'ps	$0.02/-90^{\circ}$ `ps	-2.8×10^{-4} 490 ⁰ `ps	-0.002	$\mathbf{C}^{\dagger\,\mathbf{i}\,\mathbf{i}} = \mathbf{C}^{\,\mathbf{i}\,\mathbf{i}}$	$+3.0$				

Table V

Fig. 16. Case III: Frequency of oscillation and power output as functions of the frequency to which the oscillator is tuned.

Fig. 17. Case III: Relation between frequency of oscillation and frequency shift.

The equivalent circuit for Case IV is shown in Fig. 18.

The power output, as in all previous cases (see p. 37), is obtained by the use of Table IH.

The relative magnitudes of the feedback and grid voltages are the same as in all previous cases and are given by Eqs. (30) and (34):

$$
\left.\frac{V_{fb}}{V_{rs}^i}\right| = \frac{104}{93.6} = 1.11
$$
 (37)

The normalizing constant can be evaluated by first calculating the feedback current at the natural resonant frequency of the load, as in Eqs. (25) and (26) :

$$
i_3 = \frac{V_{rs}^{\dagger} - V_{fb}}{Z_{fb} + Z_c} = \frac{93.6 I_{ps} - V_{fb}}{2500 - j4000}
$$

= $\left[93.6 I_{ps} - 104. / +30^{\circ} I_{ps}\right] \left[1.12 - j1.80\right] 10^{-4}$
= $\left[97.4 - j52.2\right] 10^{-4} I_{ps}$ (38)

In the above calculation the value of Z_{fh} is found to be the conjugate of the Z_{fb} of Case I, pgs. 21 and 23.

The normalizing constant is

$$
N = \frac{52.2 \times 10^{-4}}{0.21} = 248 \times 10^{-4}
$$
 (39)

The equivalent shunt capacity $C^{\dagger \dagger}$ is found by using Eq. (33).

The value of C^{111} is found as in Case I page 26, and the Δf is again obtained by use of Eq. (28).

The results as obtained by use of the graphical calculators are tabulated in Table VI, and are plotted in Fig. 19, as obtained with the aid of Figs. 20 and 21.

MU-10220

Case IV. Frequency shift as a function of the frequency of oscillation									
Δf_{0}	$i_3 \times 1/N$	$i_3 \times 1/\overline{N}$ (quad, part)	(quad. part)	C ¹¹	C^{III}	Δf			
(cps)	(am)	(am)	${\rm (amp)}$	$(\mu\mu f)$	$(\mu\mu f)$	(kc)			
$\mathbf{0}$	$0.47 / -27.0$ ^O I рs	$0.21/-90o$ I ps	$-52.2 \times 10^{-4} / +90^{0}$ I ps	-0.044	$\mathbf{0}$	$\mathbf{0}$			
+470	$0.18/+39.5^{\circ}$ ps	$0.12/-90^{\circ}$ $p_{\rm S}$	$-29.8 \times 10^{-4} / +90^{6}$ ps	-0.025	$+0.019$	-27.5			
-470	$0.71/111.0^{\circ}$ ps	$0.14/-90^{\circ}$ рs	$-34.8 \times 10^{-4} / +90^{o}$ рs	-0.029	$+0.015$	-21.8			
$+705$	$0.05/-12.7^{\circ}$ p _s	$0.01/-90^{\circ}$ ps	- 2.5 x $10^{-4}/+90^{\circ}$ ps	-0.002	$+0.042$	-62.3			
-709	$0.80/-3.5^{\circ}$ ps	$0.05/ -90^{\circ}$ DS.	$-12.4 \times 10^{-4} / +90^{\circ}$ рs	-0.010	$+0.033$	-49.3			
$+1410$	$0.27/ +98.0^{\circ}$ ps	$0.26/ +90^{\circ}$ ps	$+64.5 \times 10^{-4} / +90^{\circ}$ ps	$+0.050$	+0.099	$-145.$			
-1410	$0.93/114.5^{\circ}$ `ps	$0.23/+90^{\circ}$ ps	$+57.1 \times 10^{-4} / +90^{\circ}$ ps	$+0.048$	$+0.091$	$-135.$			
+2820	$0.55/ + 82.5^{\circ}$ рs	$0.54/+90^{\circ}$ рs	$+134. \times 10^{-4}/+90^{\circ}$ ps	$+0.113$	$+0.157$	$-232.$			
$= 2820$	$0.97/132.5^{\circ}$ ps	$0.51/+90^{\circ}$ $p_{\rm S}$	$+127. \times 10^{-4} / +90^{\circ}$ ps	$+0.106$	$+0.149$	$-220.$			
+5640	$0.71/+71.0^{\circ}$ ps	$0.68/+90^{\circ}$ ps	$+169. \times 10^{-4}/+90^{\circ}$ рs	$+0.142$	$+0.186$	$-276.$			
-5640	$0.93/$ +44.5 ^o DS	$0.66/+90^{\circ}$ \cdot ps	$+164. \times 10^{-4} / 90^{\circ}$ ps	$+0.138$	$+0.181$	$-268.$			
+11280	$0.78/ +64.0^{\circ}$ ps	$0.70/+90^{\circ}$ I ΰS	$+174. \times 10^{-4} / +90^{\circ}$ рs	$+0.146$	$+0.191$	$-283.$			
-11280	$0.89/+51.0^{\circ}$ ps	$0.69/+90^{\circ}$ ps	$+171. \times 10^{-4} / +90^{\circ}$ p s	$+0.144$	$+0.187$	$-276.$			
+22560	$0.82/+61.0^{\circ}$ ps	$0.72/+90^{\circ}$ ps	$+179. \times 10^{-4} / +90^{o}$ 1 рs	$+0.149$	$+0.193$	$-283.$			
-22560	$0.87/ +54.5^{\circ}$ ps	$0.70/+90^{\circ}$ ps	$+174. \times 10^{-4}/+90^{\circ}$ ps	$+0.146$	$+0.191$	$-283.$			
+45120	$0.83/+59.5^{0}$ ps	$0.72/+90^{\circ}$ `ps	$+179. \times 10^{-4}/+90^{\circ}$ ps	$+0.149$	$+0.193$	$-283.$			
-45120	$0.86/+56.5^{\circ}$ ps	$0.71/+90^{\circ}$ ps	$+176. \times 10^{-4} / +90^{\circ}$ ps	-0.148	$+0.191$	$-283.$			

Table VI

Fig. 19. Case IV: Frequency of oscillation and power output as functions of the frequency to which the oscillator is tuned.

Fig. 21. Relation between frequency of oscillation and frequency shift (higher frequencies).

IV. CONCLUSIONS

This oscillator is a stabilized oscillator and not a synchronized oscillator such as has been studied by Adler, $\frac{1}{1}$ Appleton, $\frac{4}{1}$ Huntoon, $\frac{5}{1}$ and others. However, it does exhibit characteristics strikingly similar to the "locking" of a synchronized oscillator, particularly in the resistance feedback-coupling impedance of Case III (p. 37).

The "pulling bandwidth" with the resistance feedback-coupling impedance was narrower than with the reactive coupling, although approximately the same voltage and circuit impedances were used in all cases. The oscillating frequency, in a compensating way, was very much closer to the natural frequency of the load and the power output was very constant over the narrower bandwidth.

The jump in the experimental curve B, (see Fig. 27 in Appendix F) was identified as an accidental resonance in the grid circuit at the time the experimental data were taken. This effect can be eliminated by properly adjusting the circuit parameters in the grid circuit.

The unstable regions indicated in Figs. 11, 14, 16, and 19 follow reasonably from the fact that the feedback changes the frequency during the buildup of oscillations in the resonant load, and the final equilibrium frequency will be the first stable situation found by the oscillator. This frequency modulation has been observed but not measured.

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¢.

V. APPENDICES

A. Calculation for the Q of the Cathode-Grid-Screen Circuit

The following calculations are based on physical measurements of the oscillator and the operating conditions of the tube. The voltages and the dimensions of the cathode-grid-screen circuit are both shown in Fig. 22.

A fundamental definition of Q^7 is

$$
Q = \frac{2\pi \text{ energy stored at peak of cycle}}{\text{energy dissipated per cycle}}
$$

=
$$
\frac{\omega \text{ energy stored at peak of cycle}}{\text{power dissipated}}
$$

The Q of this grid circuit, which is composed of several transmission lines tightly coupled together, is

$$
Q = \frac{\omega \left[W_{s1} + W_{s2} + \dots W_{s.5}\right]}{\left[W_{d1} + W_{d2} + \dots W_{d5}\right]}
$$

The energy stored per quarter wavelength of coaxial transmission line is calculated first. Since the circuit losses are small when large brass conductors are used, the voltage and current distributions are nearly sinusoidal. A small error is introduced in the calculations if they are assumed to be exactly sinusoidal. We have

$$
\Delta W_{s} = \frac{\Delta C_{o}V^{2}}{2}, \quad \Delta C_{o} = C_{o}\Delta l,
$$

$$
\Delta W_{s} = \frac{C_{o}\Delta lV^{2}}{2},
$$

MU-10225

 $\sim 10^6$ $\gamma_{\rm{th}}$

 ~ 10 \sim

$$
\int_0^W \frac{s}{dW_s} = \frac{\lambda C_o V_0^2}{4\pi} \int_0^{\pi/2} \sin^2\left(\frac{2\pi \ell}{\lambda}\right) d\left(\frac{2\pi \ell}{\lambda}\right)
$$

$$
= \frac{\lambda C_o V_0^2}{4\pi} \left[\frac{1}{2} \left(\frac{2\pi \ell}{\lambda}\right) - \frac{1}{4} \sin 2\left(\frac{2\pi \ell}{\lambda}\right)\right]
$$

$$
W_s = \frac{\lambda C_o V_0^2}{16}.
$$

The energy dissipated per quarter wavelength of coaxial transmission line is as follows.

$$
I = I_0 \cos\left(\frac{2\pi l}{\lambda}\right),
$$

\n
$$
\Delta W_d = \frac{\Delta R I^2}{2}, \quad \Delta R = R \Delta l,
$$

\n
$$
\Delta W_d = \frac{R \Delta l I^2}{2},
$$

\n
$$
\int W_d = \frac{\lambda R I_0^2}{4\pi} \int_0^{\pi/2} \cos^2\left(\frac{2\pi l}{\lambda}\right) d\left(\frac{2\pi l}{\lambda}\right),
$$

\n
$$
W_d = \frac{\lambda R I_0^2}{4\pi} \left[\frac{1}{2}\left(\frac{2\pi l}{\lambda}\right) + \frac{1}{4}\sin 2\left(\frac{2\pi l}{\lambda}\right)\right],
$$

\n
$$
W_d = \frac{\lambda R I_0^2}{16}.
$$

For both conductors of a coaxial transmission line,

$$
W_{d} = \frac{\lambda I_0^2}{16} \left[R_i + R_0 \right] , R_i = \frac{r_{si}}{\pi d_i} , R_o = \frac{r_{so}}{\pi d_o}
$$

If the inner and outer conductors are of the same material,

$$
\begin{bmatrix} R_i + R_o \end{bmatrix} = \frac{r_s}{2\pi} \begin{bmatrix} \frac{2}{d_i} & + \frac{2}{d_o} \end{bmatrix},
$$
\n
$$
W_d = \frac{\lambda r_s I_o^2}{32\pi} \begin{bmatrix} \frac{2}{d_i} & + \frac{2}{d_o} \end{bmatrix}.
$$

The total Q of the composite system is

$$
Q_{t} = \frac{\left(\frac{\omega\lambda}{16}\right)\left[C_{o1}V_{01}^{2} + C_{o2}V_{02}^{2} + \dots C_{o5}V_{05}^{2}\right]}{\left(\frac{\lambda r_{s}}{32\pi}\right)\left[\frac{r_{01}^{2}}{r_{01}}\left(\frac{2}{d_{11}} + \frac{2}{d_{01}}\right) + \frac{r_{02}^{2}\left(\frac{2}{d_{12}} + \frac{2}{d_{02}}\right) + \dots + \frac{r_{05}^{2}\left(\frac{2}{d_{15}} + \frac{2}{d_{05}}\right)}{r_{02}^{2}}\right]}
$$

Assuming sinusoidal distributions of voltage and current in a quarter wavelength resonant transmission line, we find the relation

$$
\mathbf{V}_0 = \mathbf{I}_0 \mathbf{Z}_0
$$

The composite Q is then

$$
Q_{t} = \frac{2\pi\omega}{r_s} \frac{\begin{bmatrix} C_{o1} V_{01}^2 + C_{o2} V_{02}^2 + \dots C_{o5} V_{0.5}^2 \\ V_{01}^2 + C_{o2} V_{02}^2 + \dots C_{o5} V_{0.5}^2 \end{bmatrix}}{\begin{bmatrix} V_{01}^2 \\ Z_{01}^2 \end{bmatrix} \begin{bmatrix} \frac{2}{\alpha_1} + \frac{2}{\alpha_0} \\ 1 + \frac{2}{\alpha_1} \end{bmatrix} + \frac{V_{02}^2}{Z_{02}^2} \begin{bmatrix} \frac{2}{\alpha_1} + \frac{2}{\alpha_0} \\ 1 + \frac{2}{\alpha_0} \end{bmatrix} + \dots + \frac{V_{05}^2}{Z_{05}^2} \begin{bmatrix} \frac{2}{\alpha_1} + \frac{2}{\alpha_0} \\ 1 + \frac{2}{\alpha_0} \end{bmatrix}}
$$

The characteristic impedance (Z_0) of a coaxial transmission line is given by

$$
Z_0 = 138 \log_{10} \left(\frac{d_o}{d_i} \right) \text{ ohms} .
$$

With the aid of the above relation and the physical dimensions of the oscillator, the characteristic impedances of the transmission lines in the cathode -grid -screen circuit have been calculated. These are listed in Table VII. The relative voltages are based on the tube calculations in Appendix B (p. 56), and are given in Table VII^T in terms of V_{02} for the condition that $V_{04} = 0.796 V_{05}$.

Open-end voltage of the several transmission lines
in terms of the open-end voltage of Line 2

By using the relation⁷

 $C_0 = \frac{0.241 \epsilon_1}{4} 10^{-10}$ farad per meter, $\frac{1}{\log_{10} \left(\frac{d_o}{d_i}\right)}$

we get the final expression for Q_t :

$$
Q_{\frac{1}{t}} = \frac{0.48 \text{ m} \epsilon_1 10^{-10} \left[\frac{v_{01}^2}{\log \left(\frac{d_{01}}{d_{11}} \right)} + \frac{v_{02}^2}{\log \left(\frac{d_{02}}{d_{12}} \right)} + \dots + \frac{v_{05}^2}{\log \left(\frac{d_{05}}{d_{15}} \right)} \right]}{r_s}
$$

$$
r_s = \frac{v_{01}^2}{z_{01}^2} \left(\frac{2}{d_{11}} + \frac{2}{d_{01}} \right) + \frac{v_{02}^2}{z_{02}^2} \left(\frac{2}{d_{11}} + \frac{2}{d_{02}} \right) + \dots + \frac{v_{05}^2}{z_{05}^2} \left(\frac{2}{d_{11}} + \frac{2}{d_{05}} \right)
$$

$$
Q_{t} = \frac{0.482 \omega \pi \epsilon_1 10^{-10} [0.670 + 5.19 + 3.05 + 3.22 + 3.45 \cdot \cdot \cdot]}{r_{s} [0.003 + 0.041 + 0.074 + 0.152 + 0.088]}
$$

$$
Q_{\rm E} = \frac{0.482 \text{ m} \epsilon_1 10^{-10}}{r_{\rm s}} \left[\frac{15.6}{0.358} \right].
$$

The surface resistivity of brass⁶ is

 $r_s = 5.01 \times 10^{-7} \sqrt{f}$ ohm per square.

This gives, for the total Q,

$$
Q_{t} = \frac{0.482 \times 2 \pi 202.55 10^{6} \times \pi \times 10^{-10}}{5.01 \times 10^{-7} \sqrt{202.55 10^{6}} [43.6]},
$$

 $Q_t = 1170$ (measured value^{*} $Q_m = 570$).

The ratio of the calculated Q to the measured Q gives the ratio of the actual conductor losses to the calculated losses. The stored energy is the same in either case. We find

$$
\frac{Q_c}{Q_m} = \frac{W_{dm}}{W_{dc}} = \frac{1170}{570} = 2.06.
$$

The effect of the grid drive power is to further lower the Q by adding an additional loss in the circuit. From Appendix B, the grid drive power is

$$
P_g = I_g V_g \text{ (crest value)}
$$

= 5.2 (600 + 1200)
= 9340 watts

$$
R_g = \frac{V^2}{2 P_g} = \frac{1800^2}{2 \times 9340} = 174 \text{ ohms}
$$

Since the grid drive power is dissipated at the open end of transmission line No. 4, it may be included in the expression for Q_t by writing the expression in the form

$$
W_g = \frac{V_{04}^2}{2 \text{ R g}} = \frac{\left[0.532 \text{ V}_{02}\right]^2}{2 \times 174} = 8.14 \times 10^{-4} \text{ V}_{02}^2
$$

see Appendix \mathbf{F} .

and normalizing,

$$
W_{g} = \frac{32\pi}{\lambda r_{s}} \times 8.14 \times 10^{-4} \text{ V}_{02}^{2} = 7.75 \text{ V}_{02}^{2}
$$

The corrected value of Q_t including the grid drive power is
\n
$$
Q_t = 1170 \times \frac{1}{43.6} \left[\frac{15.6}{0.358 \times 1.97 + 7.75} \right]
$$
\n
$$
= \left[\frac{1170}{43.6} \right] \left[\frac{15.6}{8.49} \right] \quad \text{or}
$$
\n
$$
Q_t = 49.2
$$

B. Calculation of Tube Operation

The function of the oscillating cathode-grid-screen Circuit is to supply the maximum fundamental component of space current to the output (anode) circuit. An operating point for the tube-must be chosen that insures strong self-oscillation. Since the maximum possible current is to flow to the anode, the efficiency must be as low as possible consistent with strong self-oscillation. As is shown in the analysis, the following operating point for the tube fits these requirements:

$$
E_p = 10,000
$$
 volts dc
\n
$$
E_s = 3,000
$$
 volts dc
\n
$$
E_g = 600
$$
 volts dc
\n
$$
e_{gm} = 1,200
$$
 volts (instantaneous)
\n
$$
e_{pm} = 5,420
$$
 volts (instantaneous)
\n
$$
e_{sm} = 2,540
$$
 volts (instantaneous)
\n
$$
V_{pc} = E_p - e_{pm} = 4,580
$$
 volts ac
\n
$$
V_{sc} = E_s - e_{sm} = 460
$$
 volts ac

Since it is assumed that e_{gm} , e_{pm} , and e_{sm} go through their crest values at the same time, the graphical Fourier analysis for the

- 55 -

case of a symmetrical wave shape can be used. This has been worked out in Appendix $C(pg, \sqrt{64})$. $\mathcal{F} = \frac{1}{\sqrt{2}}$

For the above operating point, the electrode voltages and currents corresponding to the graphical analysis points are tabulated in Table VIII. These points are also plotted on the load lines for the tube $(4W20,000A)$ in Fig.23. On examination of the values of the currents at each of the analysis points we note that only $F(0)$ through $F(6)$ contribute, and all others beyond F(6) are zero.

The dc value of the grid current is calculated by using the relation from Appendix C $(p. 64)$:

$$
I_g = \frac{1}{12} \left[0.5 F(0) + F(1) + F(2) + \dots + F(23) + 0.5 F(24) \right]
$$

= $\frac{1}{12} \left[0.5 x 23.5 + 22.5 + 17.4 + 10.2 + 0.38 + 0.0 \right]$
= $\frac{1}{12} \left[62.2 \right] = 5.2$ amperes

The screen currents are calculated by using the relations in Appendix C (pg. .64): \therefore ...

 $-56-$

Electrode voltages and currents corresponding to graphical analysis points

Table VIII

The plate currents are calculated by using the same relations:

 $I_{ps} = \frac{1}{13} \int 100 + 1.93 \times 101 + 1.73 \times 99.0 + 1.41 \times 91.0 + 61.0 + 0.518 \times 10.7$ $=\frac{1}{12}$ [661] = 55.0 amperes (crest value of the fundamental component). $I_{p} = \frac{1}{12} \left[0.5 \times 100 + 101 + 99.0 + 91.0 + 61.0 + 10.7 \right]$ $=\frac{1}{12}$ [413] = 34.4 amperes (dc value).

The power output is the product of the fundamental components of plate current and plate-to-screen voltage:

$$
P_o = \frac{I_{ps}}{\sqrt{2}} \frac{V_{ps}}{\sqrt{2}} = 55 \left(\frac{V_{pc} + V_{sc}}{2} \right) = 138 \text{ kilowatts.}
$$

The plate-load impedance at the natural resonant frequency (f_0) of the load is the fundamental component of plate-to-screen voltage divided by the plate current:

$$
Z_{\text{ps}}(f_0) = R_{\text{ps}} = \frac{5040}{55} = 91.7 \text{ ohms.}
$$

Power input = $I_p \times E_p = 34.4 \times 10,000 = 344$ kilowatts.

The plate dissipation is the difference between the plate power input and the plate power output:

Plate dissipation = $344 - 138 = 206$ kilowatts.

The grid drive power can be obtained with sufficient accuracy by using the following relation:⁷

$$
P_g = I_g V_g
$$
 (crest value)
= $I_g (E_g + e_{gm})$
= 5.2 (600 + 1200),
 $P_g = 9340$ watts.

The plate efficiency is the ratio of the plate power output to the plate power input:

Plate efficiency = $\frac{138}{111}$ = 40.1 percent. 344

The screen dissipation is the difference between the screen de power input and the power generated by the screen circuit:

Screen dissipation =
$$
(E_{s}I_{s}) - \left[\frac{V_{sc}}{\sqrt{2}} \frac{I_{sg}}{\sqrt{2}} \left(\frac{I_{s}}{I_{s} + I_{p}}\right)\right]
$$

\n
$$
= (3000 \times 8.4) - \left[\frac{460}{\sqrt{2}} \frac{70.1}{\sqrt{2}} \left(\frac{8.4}{8.4 + 34.4}\right)\right]
$$

= 22.1 kilowatts.

The grid dissipation can be approximated by using the following relation: $8,9$

Grid dissipation =
$$
I_g V_g - E_g I_g
$$

= (5.2 x 1800) - (600 x 5.2)

 $= 6.2$ kilowatts.

C. Graphical Fourier Analysis of Nonsinusoidal Wave Forms

The basis of the Fourier analysis is the assumption that the periodic function F(t) may be written in the form

$$
F(t) = \begin{bmatrix} a_0 + a_1 \cos{(wt)} + a_2 \cos{(2wt)} + \dots a_n \cos{(nw t)} \\ + b_1 \sin{(wt)} + b_2 \sin{(2wt)} + \dots + b_n \sin{(nw t)} \end{bmatrix}
$$

The coefficients in this series can be evaluated as a result of the orthogonality properties of sinusoids. These coefficients have the following forms: 6

$$
a_0 = \frac{1}{2 \pi} \int_0^{2\pi} F(t) d(wt),
$$

$$
a_n = \frac{1}{\pi} \int_0^{2\pi} F(t) \cos n(wt) d(wt),
$$

$$
b_n = \frac{1}{\pi} \int_0^{2\pi} F(t) \sin n(wt) d(wt).
$$

For use in a graphical analysis these integrals can be approximated by summations. The wave shape in Fig. 24 can be used as an aid in setting up these summations.

·•

a =
$$
\frac{1}{24}
$$
 $\frac{m=24}{m=0}$ = $\left[\frac{1}{24} \frac{F(0)}{2} + F(1) + F(2) + \dots + \frac{F(24)}{2}\right]$,
\na_n = $\frac{1}{12}$ $\frac{m=24}{m=0}$ cos n (m 15°)
\n= $\frac{1}{12}$ $\left[\frac{F(0) \cos (n 0^0)}{2} + F(1) \cos (n 15^0) + \frac{F(24) \cos (n 360^0)}{2}\right]$,
\nb_n = $\frac{1}{12}$ $\sum_{m=0}^{m=24} F(m) \sin n (m 15^0)$
\n= $\frac{1}{12}$ $\left[\frac{F(0) \sin (n.0^0)}{2} + F(2) \sin (n 15^0) + \frac{F(24) \sin (n 15^0)}{2}\right]$.
\n= $\frac{F(24) \sin (n.360^0)}{2}$.

For the purpose for which these calculations are to be used, only the fundamental frequency $(n = 1)$ is of interest. Summing the terms that have the same numerical coefficients gives the following expressions for a_0 , a_1 , and b_1 :

$$
a_0 = \frac{1}{24} \left[\frac{F(0)}{2} + F(1) + F(2) + \ldots + \frac{F(24)}{2} \right]
$$

$$
a_{1} = \frac{1}{12} \begin{cases} \begin{bmatrix} F(0) - F(12) \end{bmatrix} \cos 0^{\circ} \\ \begin{bmatrix} F(1) - F(11) + F(23) - F(13) \end{bmatrix} \cos 15^{\circ} + \\ \begin{bmatrix} F(2) - F(10) - F(14) + F(22) \end{bmatrix} \cos 30^{\circ} + \\ \begin{bmatrix} F(3) - F(9) - F(15) + F(21) \end{bmatrix} \cos 45^{\circ} + \\ \begin{bmatrix} F(4) - F(8) - F(16) + F(20) \end{bmatrix} \cos 60^{\circ} + \\ \begin{bmatrix} F(5) - F(7) - F(17) + F(19) \end{bmatrix} \cos 75^{\circ} + \\ \begin{bmatrix} F(6) - F(18) \end{bmatrix} \cos 90^{\circ} \\ \begin{bmatrix} F(1) + F(11) - F(13) - F(23) \end{bmatrix} \sin 15^{\circ} + \\ \begin{bmatrix} F(2) + F(10) - F(14) - F(22) \end{bmatrix} \sin 30^{\circ} + \\ \begin{bmatrix} F(2) + F(9) - F(15) - F(21) \end{bmatrix} \sin 45^{\circ} + \\ \begin{bmatrix} F(4) + F(8) - F(16) - F(20) \end{bmatrix} \sin 60^{\circ} + \\ \begin{bmatrix} F(5) + F(7) - F(17) - F(19) \end{bmatrix} \sin 75^{\circ} + \\ \begin{bmatrix} F(6) - F(18) \end{bmatrix} \sin 90^{\circ} \end{cases}
$$

For the special case of a wave shape that is symmetrical about F(0), m = 0 the sine term is zero $(b_1 = 0)$. Also,

$$
a_1 = \frac{1}{12} \begin{bmatrix} [F(0) - F(12) + 1.98 [F(1) - F(11)] \\ + 1.73 [F(2) - F(10)] + 1.41 [F(3) - F(9)] \\ [F(4) - F(8)] + 0.518 [F(5) - F(7)] \end{bmatrix}.
$$

..

D. Calculation of the Plate -Load Impedance as a Function of Frequency

The universal resonance curves for parallel resonant circuits⁷ have been used to determine the plate -load impedance as a function of frequency. The exact values are fixed by the Q of the resonant load (measured $Q = 72,000$) and the coupled impedance at resonance. By design, this value has been set at $Z_{ps}(f'_{0}) = R_{ps}(f_{0}) = 91.7$ ohms. This is calculated in Appendix B.

The plate-load impedance and the resistance component of the plate -load impedance are listed as functions of frequency in Table IX.

E. Proof that the Locus of the Impedance Vector for a Parallel Circuit is a Circle

In the following proof Q_0 and ω_0 are the Q and the angular frequency at resonance respectively. The proof is concerned with Figs. 25 and 26.

The parallel impedance is as follows:

$$
Z = \frac{\left(\frac{j\omega LR}{R + j\omega L}\right)\left(\frac{-j}{\omega C}\right)}{\left(\frac{j\omega LR}{R + j\omega L} - \frac{j}{\omega C}\right)} = \frac{\omega LR}{\omega L - j\left[R - \omega^{2} L RC\right]}
$$

$$
= \frac{\omega L}{\frac{\omega L}{R} - j [1 - \omega^{2} LC]} = \frac{\omega L [\frac{\omega L}{R} + j (1 - \omega^{2} LC)]}{\frac{\omega^{2} L^{2}}{R^{2}} + (1 - \omega^{2} LC)^{2}}
$$

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Now at resonance,

$$
LC = \frac{1}{\omega_0^2},
$$

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Table IX

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$$
Q_0 = \frac{\omega_0 W_s}{W_d} = \frac{\omega_0 \frac{1}{2} C V^2}{\frac{V^2}{2R}} = \omega_0 C R = \frac{R}{\omega_0 L}
$$

Therefore

 \sim

$$
Z = \frac{R + j \frac{R^{2}}{\omega L} (1 - \omega^{2} LC)}{1 + \frac{R^{2}}{\omega^{2} L^{2}} (1 - \omega^{2} LC)^{2}}
$$
\n
$$
= \frac{R + jR \left(\frac{\omega_{0}}{\omega}\right) Q_{0} \left[\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0}^{2}}\right]}{1 + \left(\frac{\omega_{0}}{\omega}\right)^{2} Q_{0}^{2} \left[\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0}^{2}}\right]^{2}}
$$
\n
$$
Z = R \left[\frac{1 + jQ_{0} \left(\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0} \omega}\right)}{1 + Q_{0}^{2} \left(\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0} \omega}\right)^{2}}\right]
$$
\n
$$
= R \frac{\sqrt{1 - Q_{0}^{2} \left(\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0} \omega}\right)^{2}}}{1 + Q_{0}^{2} \left(\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0} \omega}\right)^{2}}
$$
\n
$$
= R + R \left[\frac{1 + jQ_{0} \left(\frac{\omega_{0}^{2} - \omega^{2}}{\omega_{0} \omega}\right)}{\omega_{0} \omega}\right]
$$

 $\overline{}$

$$
\frac{R}{2} - \frac{R}{2} + R \left[\frac{3\omega_0}{1 + \omega_0^2} \left(\frac{\omega_0^2 - \omega^2}{\omega_0} \right)^2 \right]
$$

$$
= \frac{R}{2} + \frac{R}{2} \left[\frac{-1 - Q_0^2 (\frac{\omega_0^2 - \omega^2}{\omega_0 \omega}) + 2 + j2 Q_0 (\frac{\omega_0^2 - \omega^2}{\omega_0 \omega})}{1 + Q_0^2 (\frac{\omega_0^2 - \omega^2}{\omega \omega_0})^2} \right]
$$

 $\frac{\frac{2}{\omega_0 \omega}}{\frac{2}{\omega_0 \omega}}\left\{\frac{1}{1} + j Q_0 \left(\frac{\omega_0^2 - \omega^2}{\omega_0 \omega}\right)\right\}$ ω $rac{R}{2}$ $\frac{\mathbf{R}}{2}$ $^{\circ}$ 0 $\frac{\sqrt{2}}{2}$ $\overline{2}$

$$
= \frac{R}{2} + \frac{R}{2} \left[\frac{1}{\theta} \left(\theta - \tan^{-1} \frac{Q_0 \left(\frac{\omega^2 - \omega^2}{\omega_0 \omega} \right) - \tan^{-1} \left(- \right) \frac{Q_0 \left(\frac{\omega^2 - \omega^2}{\omega_0 \omega} \right)}{\omega_0 \omega} \right) \right].
$$

But the simple relation between the inverse tangets $\tan \frac{1}{2}(-\theta) = + \tan^{-1} \theta$ gives the final result:

$$
Z = \frac{R}{2} + \frac{R}{2} \left[1 / \theta = 2 \tan^{-1} Q_0 \left(\frac{\omega_0^2 - \omega^2}{\omega_0 \omega} \right) \right],
$$

which clearly shows the locus to be a circle.

 $-69-$

F. Experimental Data

Results are given in the following tables.

Measurement of oscillation characteristics - Cases I, II

* Clockwise

., ... -

t Counterclockwise

Fig. 27. Frequency of oscillation and power output as functions of the frequency to which the oscillator is tuned: experimental data corresponding to Cases I and II.

Fig. 28. Corrected experimental data corresponding to Cases I and II.

NOTE: The experimental data are plotted against a linear frequency line to correct for the fact that the tuning capacitor varies inversely as the spacing where the spacing is proportional to knob numbers. Also, the calculations do not admit to coupling back through feedback line No. 2 nor do they admit to accidental resonances in the oscillating grid circuit as indicated by the jump in Curve B.

DEFINITIONS OF 5 YMBOLS

 $A_{0,1}$ = area of coupling loop No. 1 (p. 13) A_{12} = area of coupling loop No. 2 (p. 13) $a =$ with subscripts, the coefficients of the Fourier series (p. 61) $b =$ with subscripts, the coefficients of the Fourier series $(p. 61)$ $C =$ capacity (p. 2 θ C^1 = equivalent lumped circuit capacity of the cathode-grid-screen circuit (p. 17) *C''* = equivalent shunt capacity appearing across the cathode -grid screen circuit as a result of the feedback network $(p. 23)$ $C^{(i)}$ = corrected value of the equivalent shunt capacity appearing across the cathode -grid -screen circuit as a result of the feedback network (p. 26) $\mathbf{C_{o}}$ $=$ with additional subscripts, the capacity per meter of transmission line (p, 50) = coupling coefficient (p. 13) c $\mathbf{d}_{\mathbf{i}}$ = diameter of the inner conductor of a coaxial transmission line (p.53) d 0 $=$ diameter of the outer conductor of a coaxial transmission line (p. 53) E_{g} $= d \cdot c$. control grid bias (p. 18). E_p^e $= d c \cdot$ plate voltage (p. 56) $E_{\bf g}$ $= dc$. screen voltage (p. 56) e g $=$ instantaneous value of the grid-to-cathode voltage (p. 58) e_{gm} = maximum instantaneous value of the grid-to-cathode voltage $(p. 18)$ e _p $=$ instantaneous value of the plate-to-cathode voltage (p.58) e_{pm} = minimum instantaneous value of the plate-to-cathode voltage (p. 56) $e_{\rm g}$ = instantaneous value of the soreen-to-cathode voltage (p.58) $e_{\bf sm}$ = minimum instantaneous value of the screen-to-cathode voltage (p. 18) $F(t)$ = arbitrary periodic function of time (p. 58, 61) \mathbf{f} = frequency in cycles per second f_{0} = natural frequency of oscillation of the resonant load in cycles per second (p.21.) h = fraction of, the plate-to-screen voltage across which transmission line No. 1 is coupled $(p. 11)$

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 $\label{eq:2.1} \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{$

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 $\mathcal{A}_{\mathcal{A}}$, $\mathcal{A}_{\mathcal{A}}$

 $\sim 10^6$

 $\label{eq:2.1} \Psi_{\alpha\beta} = \Psi_{\alpha\beta} + \Psi_{\alpha\beta} + \Psi_{\alpha\beta} + \Psi_{\alpha\beta} + \Psi_{\alpha\beta} + \Psi_{\alpha\beta}$ $\label{eq:2} \frac{1}{2}\sum_{i=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{1}{2}\sum_{j=1}^n\frac{$

 $\label{eq:2} \frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1}{\sqrt{2}}\sum_{i=1}^n\frac{1$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$

 $\frac{1}{2}$ \mathbb{R}^n

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