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THE IBM SHARE PROGRAM D2 NU SCHR 1072 FOR SOLUTION OF THE
SCHRÖDINGER RADIAL EQUATION, BY J. W. COOLEY:
NECESSARY AND USEFUL MODIFICATIONS FOR ITS USE ON AN IBM 7090

R. N. Zare^v and J. K. Cashion^v

July 1963

The IBM SHARE Program D2 NU SCHR 1072 for Solution of the
Schrodinger Radial Equation, by J. W. Cooley; Necessary
and Useful Modifications for Its Use on an IBM 7090.

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Introduction

An efficient method for solving the radial Schrodinger equation numerically has been proposed recently by Cooley.¹ He has contributed a Fortran program embodying his procedure to the IBM SHARE project, D2 NU SCHR 1072. The authors of the present report have found his program extremely useful as a means of generating eigenvalues and eigenfunctions for diatomic molecule calculations. The program was written for an IBM 704 and must be modified slightly in order to be compatible with an IBM 7090. (The same modifications are necessary for its use on the 709 or 7094.) The principal purpose of this report is to make this information available to 7090 users.

A few modifications and additions which the authors have found useful also will be described. These in no way derogate from Cooley's skill as a programmer. On the contrary his program can be recommended not only for its utility but also as an excellent example of the efficient use of Fortran coupled with a clarity of presentation which makes it easy for others to adapt his work to their own particular needs.

Tests of the accuracy of this procedure when up to 200 intervals are used in the integration have been given by Cooley.¹ Additional information on its use with up to 2000 intervals has been given elsewhere by one of (J.K.C.²). Other applications in which this program has been very useful to us include a study of vibration-rotation interaction (J.K.C.³) and a calculation

of the intensity distribution in the iodine fluorescence spectrum based on Rydberg-Klein-Rees potentials (R.N.Z.⁴).

The Numerical Method

The following is taken largely from Cooley's summary of the method as given in the SHARE distribution description of this program. To facilitate comparison with the listing of the symbolic deck given in Appendix B, variable names employed in the program are also used in the following equations.

The program calculates the eigenvalue E and the normalized eigenfunction S of the second-order differential equation

$$\frac{d^2 S}{dr^2} = (V - E)S \quad (1)$$

where V is a given numerical potential function.

It is customary to separate the Schrodinger equation for a diatomic molecule into its radial and angular parts and express its solution in the form

$$\psi(r, \theta, \varphi) = R(r) \Theta(\theta) \Phi(\varphi). \quad (2)$$

Eq. (1) is equivalent to the radial Schrodinger equation when length and energy are expressed in dimensionless units. Its solution $S(r)$ is equal to r times $R(r)$. $S(r)$ provides a more convenient computational form since the product of any two functions S will include the volume element of integration. Hence, the expectation value for any quantity F is given by

$$[\bar{F}] = \int R F_{op} R \cdot r^2 dr = \int S F_{op} S dr \quad (3)$$

If length is measured in units of Bohr radii, $a_0 = 0.529172 \text{ \AA}$, the energies V and E must be expressed in a dimensionless unit equivalent to $h N_0 / 8\pi^2 c a_0^2 \mu_A$ wave numbers, where N_0 is Avagadro's number (physical scale) and μ_A is the reduced mass in Aston units. The numerical value of this factor is $60.2198/\mu_A$. Hence to convert eigenvalues generated by this program to the more familiar units of cm^{-1} , multiply them by this factor. Note that the Aston unit of reduced mass is based on the physical scale of atomic weight, i.e., $\mu_A = m_1 m_2 / (m_1 + m_2)$ where m_1 and m_2 are the atomic weights of atoms relative to $O^{16} = 16$. Herzberg's⁵ table 39 is a convenient source for μ_A values.

Regarding the potential $V(r)$ two things should be noted. First, its zero is taken at the dissociation limit. Therefore a Morse potential, for example, must be generated from the relation

$$V(r) = D_e [1 - e^{-\beta(r-r_e)}]^2 - D_e \quad (4)$$

where the dissociation energy D_e is a positive quantity. Secondly, $V(r)$ is an effective potential which may include a rotational term, $[J(J+1) - \Lambda]/r^2$, where J is the rotational quantum number and Λ is the quantum number for the z-component of electronic angular momentum. Since this term is inherently positive its addition to the potential for the rotationless state always raises the potential.

Initially, a non-normalized solution $P(r)$ for Eq. (1) is calculated from the integration formula

$$-Y_{i-1} + 2Y_i - Y_{i+1} + h^2(V_i - E)P_i = 0 \quad (5)$$

where

$$P_i = P(r_i)$$

$$Y_i = [1 - (h^2/12)(V_i - E)]P_i \quad (6)$$

$$V_i = V(r_i)$$

$$h = r_{i+1} - r_i.$$

The error associated with the use of Eq. (5) as a predictor is $(h^6/240)P_i^{(vi)}$.

Starting with the boundary values:

$$P_n = 10^{-30} \quad (7)$$

$$P_{n-1} = P_n \cdot \exp[r_n(V_n - E)^{1/2} - r_{n-1}(V_{n-1} - E)^{1/2}] \quad (8)$$

and a trial value of E , Eq. (5) is used to integrate inward, giving P_{n-2} , P_{n-3} , ..., P_m , where m is selected as the first point for which $|P_m| \leq |P_{m+1}|$. [Eq. (8) is based on the WKB approximation.] Then, P_i is replaced by P_i/P_m for $i = m, m+1, \dots, n$.

Starting with the boundary values

$$P_0 = 0, P_1 = 10^{-20}$$

Eq. (5) is used to integrate outwards, giving P_2, \dots, P_m , after which P_i is replaced by P_i/P_m for $i = 1, 2, \dots, m$. This yields a trial solution satisfying Eq. (5) at all points except r_m , the crossing-point for the outward and inward curves. A correction to E is calculated by applying the Newton-Raphson method to the calculation of the zeros of the function

$$F(E) = h^{-2}(-Y_{m-1} + 2Y_m - Y_{m+1}) + (V_m - E)P_m \quad (9)$$

which is a measure of the amount by which the m^{th} equation of Eq. (2) is not satisfied. The derivative of Eq. (9) is

$$F'(E) = - \sum_{i=1}^n P_i^2 \quad (10)$$

The correction to E, by the Newton-Raphson method, is

$$DE = -F(E)/F'(E). \quad (11)$$

After adding this correction to E to obtain a new trial eigenvalue the process is repeated until $DE \leq \epsilon$. When convergence is achieved, the normalized solution is calculated,

$$S_i = P_i / \left(h \sum_{j=1}^n P_j^2 \right)^{1/2}, \quad i=1, 2, \dots, n. \quad (12)$$

Necessary Modifications

The SHARE distribution includes the main program NU SCHR either in the form of 177 Fortran source program cards or as 43 relocatable binary cards. In addition binary cards labelled NU EMFT, 1-4 and NU LRT, 1-6 are also supplied. For 7090 use, discard all ten of these binary cards. Replace them with the FAP deck for subroutine EMFT which is listed in Appendix A. When compiled this subroutine will be contained on four binary cards.

Useful Modifications

Appendix B contains a listing of one Fortran source deck used by the authors. Frequently a particular application will

require recompilation with an altered DIMENSION statement, probably the inclusion of a COMMON card and perhaps some modification of output control or formats. The changes to be described below were either found to be generally useful or were made to remedy failures of the subroutine in particular applications. All symbolic cards which have been modified are labelled MOD 1 to 20; those labelled SCHRO004 to 176 are unaltered from the original SHARE distribution.

MOD 2 - The Call Statement

The variables which must be supplied are

- NI,NS: Output control parameters whose function is explained in the comment cards SCHRO004-007.
- RMAX,RMIN: Specify the range of r . $RMIN \leq r \leq RMAX$.
- V,S: Singly-dimensional arrays containing respectively the numerical potential and the solution, on exit. The source program must contain these in a DIMENSION statement.
- N: The number of equally-spaced intervals used in the integration. The length of each interval is $h=(RMAX-RMIN)/N$.
- KV: the number of nodes in the solution. See comments on this below.
- EO: Upon entry, EO is the first trial eigenvalue provided to SCHR by the user's program. EO must be a negative number. At exit time it is replaced by the E calculated on the last iteration. (Note: the second symbol in this variable name is a zero, not the letter O.)

EPS: The convergence criterion is $E' - E \leq \text{EPS}$ where $E' = E + \Delta E$. One may use $\text{EPS} = 0.$, except in rare circumstances where the magnitude of the solution E is extremely small relative to the V_1 's.

MAXIT: Is the maximum number of the iterations to be performed. It might be noted that frequently the iteration procedure will not satisfy a zero epsilon test. For eigenfunctions with an odd number of nodes especially it can oscillate between two values differing by only 1 unit in the last figure. Hence, considerable time can be wasted by making MAXIT unnecessarily large. In most cases this could be prevented by making $\epsilon = 2 \times 10^{-8} E_{\text{trial}}$. A convenient way of doing this would be to insert the following statement near the beginning of the subroutine, following MOD 2, say,

```
EPS = 2.0 - 8 * EO
```

Cooley has written this subroutine as a FUNCTION with the following purpose in view. At exit time, SCHR = 0. if convergence has been achieved in MAXIT iterations or 1 less. Otherwise, SCHR = 1. and EO and S are the results for MAXIT iterations. Therefore, if one wants to test for convergence one should write in the source program

```
IF [SCHR(...)] N1, N2, N3.
```

When not testing for convergence, one may write

```
CALL SCHR (...).
```

MOD 3 - The Range of Integration

Cooley's published application of his program was to the H_2^+ molecule-ion. The values of its eigenfunctions at $r = 0$ are in the range of 10^{-25} . For most molecules the values at $r = 0$ are much smaller than 10^{-39} , which produces a machine zero on the 7090. Hence the use of $r = 0$ as a lower bound needlessly extends the range of the integration. Values of $RMIN = r_e - 2$ and $RMAX = r_e + 5$ (in atomic units) will give a satisfactory range for molecules having $\omega_e / \omega_e x_e \approx 50$. The larger this ratio, the more the range may be reduced. It is well to print the eigenvalues at the two extremes in order to be sure that they are sufficiently small.

MOD 4-6 - Trial Eigenvalue too High

It can happen that an iteration will result in a large correction to the trial eigenvalue and place the new trial eigenvalue beyond the range of bound states. Cooley's provision for this was to replace such a trial value by the third last potential value, $V(N-2)$, and then continue with the next iteration. Our experience was that when such a failure occurred resumption of the iteration through this artifice never led to the eigenvalue which was being sought originally. Hence we preferred to terminate the procedure with a print-out to indicate the source of the failure.

MOD 7-11 - Test for Crossing Point

Cooley used three separate tests for terminating the inward integration: a decrease in the eigenfunction, an increase in

the potential, and, if neither of these tests did it before the value r_2 was reached, it was made the crossing-point. We have removed the test on the potential for the following reason. In deriving some RKR-type potentials various interpolation schemes were used. In some cases it was found that small local irregularities could be introduced, especially near the classical turning points or near the extremities of the potential where various approximate forms were joined on to the RKR segment. If the inward integration is terminated very far away from the maximum of the eigenfunction, the efficiency of the iteration procedure is greatly reduced and may result in failure of the program. It is necessary to retain the IF(M-2) test, for if the eigenfunction test were never satisfied a tight loop would result. The number of cases in which termination was effected by this test were very few, and all were associated with potentials generated in part by interpolation. The detailed reasons for failure in these cases were never ascertained but were thought to be closely related to the small irregularities arising from unsatisfactory interpolation procedures. The print-out of the crossing-point (MOD 8-11) was made principally to investigate the difficulties encountered in these instances.

MOD 12-15 - Node-Count

Frequently the normalized solutions will have values smaller in magnitude than 10^{-39} near either end of the integration. On the 7090 underflow will occur producing a zero, but one which

The modifications made here prevent an oscillation in the signs of these zeros from being counted as a node. In connection with the node-count we might point out a danger in specifying by a number rather than by a variable any parameter in a CALL statement. If the parameter is changed by the subroutine its new value will be used in any subsequent Fortran statements employing that parameter. If one used the statement CALL SCHR (....,....,5,...) where 5 specified KV, errors could easily result if the node-count actually differed from 5. For instance, if 6 nodes were counted a later program statement $J = J+5$ would be executed as $J = J+6$. (A number used for EO in the CALL will be altered on exit.) In this program the input KV is never used so there is no point to specifying it numerically anyway.

MOD 16-19 - Print Control

This modification is included merely as an example of an additional output option controlled by the parameter NS. With $NS = 10$ and $NI > 9$ only the node-count and eigenvalues are printed. It should be noted that MOD 17 and MOD 4 are not under the control of NI, nor have we included the option of on-line printing for these output statements. While it is most unlikely that a 7090 user would ever be printing results on line, the option provided by Cooley has not been removed from this deck. It is useful since some monitor systems interpret the PRINT command as WRITE OUTPUT TAPE 3. Where this is the case the user would have to replace cards MOD 4 and 17 by corresponding PRINT

statements, but could use Cooley's output formats by setting the appropriate control parameter equal to zero.

A Numerical Example

For the convenience of anyone wishing to check out a program which uses SCHR, we provide the following test case.

Using the input parameters $N = 1000$, $RMAX = 7.5$ and $RMIN = 0.5$ generate an array of r values in which

$$R(1) = RMIN + H$$

where

$$H = (RMAX - RMIN)/N$$

and

$$R(J) = R(J - 1) + H, \quad J = 2, N$$

Use these values to generate the Morse potential, Eq. (4), with the parameters $D_e = 605.559$, $r_e = 2.40873$ and $\beta = 0.988879$.

[Note that the value of $V(RMIN)$ is not used. $V(1) = V(RMIN + H)$.]

Call SCHR after specifying the trial eigenvalue $E0 = -581.46902$.

On exit print out the values of $R(300)$, $V(300)$, $S(300)$ and $E0$.

A program intended for generating vibrational-rotational eigenfunctions will include provision for altering the rotationless potential to the one appropriate for any J state. Use it to provide the V array for $J = 20$, change $E0$ to -508.62023 , call SCHR and then make the same print-outs as before. The values which should be obtained are given in Table I.

The first trial eigenvalues given in the text above were calculated from equations which may be found in reference 2. For purposes of this test there should be no need to specify them beyond two figures unless the value of MAXIT is very small.

Table 1. Morse Eigenvalues for the J=0 and J=20 rotational states of HCl. Values of the potentials and the eigenfunctions at $r=2.5999998$ atomic units are also given.

	J=0	J=20
E	-581.46913	-511.65467
V(r)	-587.57499	-525.44473
S(r)	1.1936730	1.4065177

For the sake of completeness the potentials and eigenfunctions obtained in the test are shown in Figures 1 and 2.

Finally, it might be noted that appendices A and B are photographic reproductions of the machine listings of the two decks. These decks were compiled and employed in a test program which used every statement with the exception of the on-line print commands. No errors were detected.

References

1. J. W. Cooley, Math. of Computation, XV, 363 (1961).
2. J. K. Cashion, "The Testing of Diatomic Potential Energy Functions by Numerical Methods," UCRL 10643, May 1963 and J. Chem. Phys., in press.
3. J. K. Cashion, "Vibration-Rotation Interaction Factors for Diatomic Molecules Calculated by Numerical Methods," UCRL-10644.
4. R. N. Zare, to be published.
5. G. Herzberg, "Spectra of Diatomic Molecules," 2nd Ed., D. Van Nostrand, New York, 1950.

Captions for Figures

1. Morse potentials for HCl with $J = 0$ and $J = 20$.
2. Ground state eigenfunctions for the potentials shown in Fig. 1.

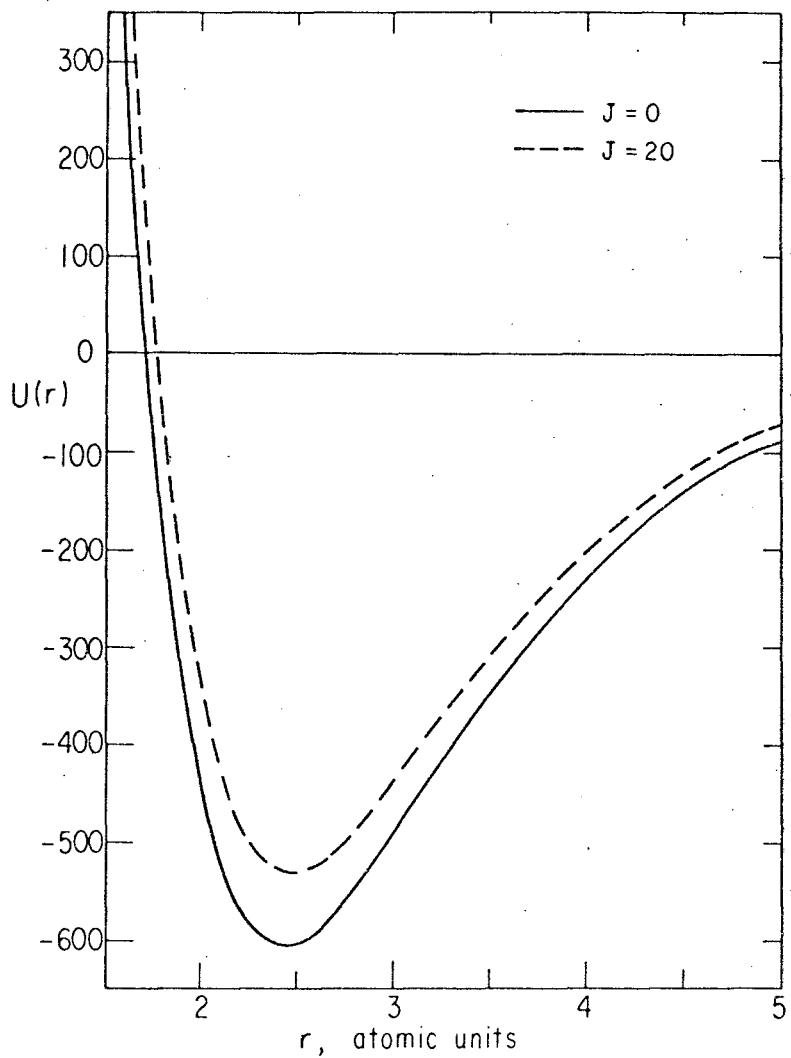


Fig. 1.

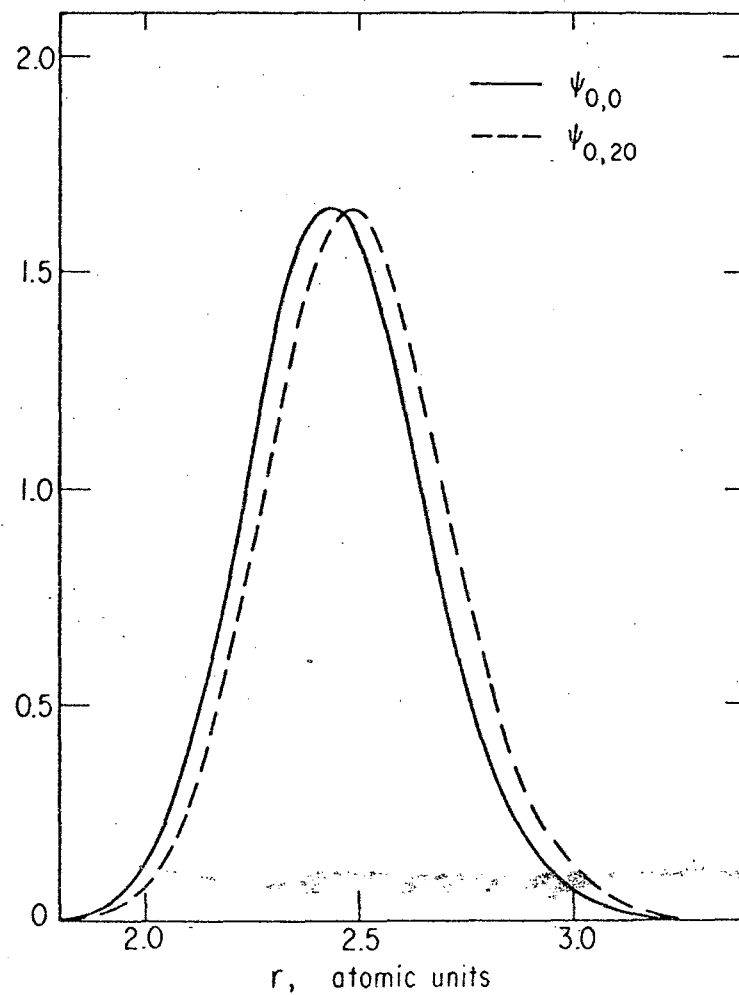


Fig. 2.

00101829334

APPENDIX A

REM NU EFM,EFM AND LFM FOR FORTRAN 2		EFMT	02
REM		EFMT	03
ENTRY EFM			
ENTRY EFMT			
LFM	OCT 476000000004	EFMT	16
	TTR 1,4	EFMT	17
EFMT	CLA 1,4	EFMT	18
	STA SETN+1	EFMT	19
	CLA TTR	EFMT	20
	TXI EFM+1,4,-1	EFMT	21
EFM	CLA HPR	EFMT	22
	STO OVER+2	EFMT	23
	OCT 476000000002	EFMT	24
SET	CLA X	TTR ANALY	EFMT 25
	STO 8		EFMT 26
	TTR 1,4		EFMT 27
X	TTR ANALY		EFMT 28
ANALY.	SXD SAVEX,1		EFMT 29
	LXD 0,1		EFMT 30
	TTR *+14,1		EFMT 31
	TTR MQO	MQ OVER	EFMT 32
SAVEX	PZE	SAVE XA IN DEC.,THEN ACC.	EFMT 33
	TTR ACMQ	11	EFMT 34
	TTR AC	10	EFMT 35
	TTR MQ	9	EFMT 36
ZERO	PZE		EFMT 37
	TTR ACMQO	AC AND MQ OVER	EFMT 38
	TTR ACO		EFMT 39
AC	CLA ZERO		EFMT 40
	TTR RETRN		EFMT 41
ACMQ	CLA ZERO	3	EFMT 42
	TTR MQ		EFMT 43
MQ	LDQ ZERO	1, MQ UNDER	EFMT 44
RETRN	LXD SAVEX,1		EFMT 45
	STO SAVEX		EFMT 46
	CLA 0		EFMT 47
	STA *+2		EFMT 48
	CLA SAVEX		EFMT 49
	TTR **		EFMT 50
ACMQO	ORA MAX		EFMT 51
MQO	STO SAVEA		EFMT 52
	LLS 0		EFMT 53
	LDQ MAX		EFMT 54
	LRS 0		EFMT 55
	TTR OVER+1		EFMT 56
ACO	ORA MAX		EFMT 57
OVER	STO SAVEA		EFMT 58
	CLA 0		EFMT 59
	HPR 63	PAUSE OVERFLOW	EFMT 60
	CLA SAVEA	PRESS START TO CONTINUE	EFMT 61
	TTR RETRN		EFMT 62
SAVEA	PZE		EFMT 63
THREE	PZE 0,0,3		EFMT 64
TTR	TTR SETN		EFMT 65
HPR	HPR 63		EFMT 66

SETN ANA THREE
STO -
TTR OVER+3
MAX OCT 3777777777
END

EFMT 67
EFMT 68
EFMT 69
EFMT 70
EFMT 71

APPENDIX B

C	MODIFIED SYMBOLIC DECK FOR COOLEYS D2 NU SCHR 1072	MOD 1
	FUNCTION SCHR(NI,NS,RMAX,RMIN,V,N,KV,E0,EPS,MAXIT,S)	MOD 2
C	NI=0 PRINT ITER. ON-LINE	SCHR0004
C	NI=1,...,9 PRINT OFF-LINE TAPE NI	SCHR0005
C	NI=OTHERWISE, DONT PRINT	SCHR0006
C	NS, SAME EXCEPT FOR SOLUTIONS	SCHR0007
C	DIMENSION S(500),V(500),P(500),Y(3)	SCHR0008
C		SCHR0009
	IF(NI)6,2,1	SCHR0010
1	IF(NI-10)4,6,6	SCHR0011
2	PRINT 3	SCHR0012
	GO TO 6	SCHR0013
3	FORMAT(70H0 ITER E F(E) DF(E)	SCHR0014
X	D(E))	SCHR0015
4	WRITE OUTPUT TAPE NI,3	SCHR0016
6	CALL EFMT(K)	SCHR0017
	H = (RMAX-RMIN)/FLOATF(N)	MOD 3
	H2=H**2	SCHR0019
8	HV=H2/12.	SCHR0020
	E=E0	SCHR0021
	TEST = -1.	SCHR0022
	DE=0.	SCHR0023
C		SCHR0024
C	START ITER LOOP	SCHR0025
C		SCHR0026
12	DO 171 IT=1,MAXIT	SCHR0027
CSTART INWARD INTEGRATION	SCHR0028
30	P(N)=1.E-30	SCHR0029
32	GN=V(N)-E	SCHR0030
34	GI=V(N-1)-E	SCHR0031
CTEST IF E TOO HIGH	SCHR0032
	IF(GI) 35, 36, 36	SCHR0033
35	WRITE OUTPUT TAPE 3, 899	MOD 4
899	FORMAT(30H DIFFERENCE EQUATION TECHNIQUE FAILS	MOD 5
	RETURN	MOD 6
		<i>MOD -5 SCHR=1-108 P 5.5</i>
36	P(N-1)=P(N)*EXPF(RMAX*SQRTF(GN)-(RMAX-H)*SQRTF(GI))	SCHR0036
38	Y=(1.-HV*GN)*P(N)	SCHR0037
40	Y(2)=(1.-HV*GI)*P(N-1)	SCHR0038
CINTEGRATE	SCHR0039
	K=0	SCHR0040
44	M=N-2	SCHR0041
46	Y(3)=Y(2)+((Y(2)-Y)+H2*GI*P(M+1))	SCHR0042
48	GI=V(M)-E	SCHR0043
50	P(M)=Y(3)/(1.-HV*GI)	SCHR0044
CTEST FOR OVERFLOW	SCHR0045
52	IF(K)54,70,54	SCHR0046
COVERFLOW	SCHR0047
54	K=0	SCHR0048
	M1=M+1	SCHR0049
	PM=P(M1)	SCHR0050
55	DO 56 J=M1,N	SCHR0051
56	P(J)=P(J)/PM	SCHR0052
58	Y=Y/PM	SCHR0053
60	Y(2)=Y(2)/PM	SCHR0054
62	Y(3)=Y(3)/PM	SCHR0055
	GI= V(M+1) - E	SCHR0056

	GO TO 46	SCHR0057
CTEST FOR CROSSING PT.	SCHR0058
70	IF(ABSF(P(M))-ABSF(P(M+1))) 90, 90, 72	SCHR0059
72	IF(M-2)90,90,81	MOD 7
81	Y= Y(2)	SCHR0062
82	Y(2)=Y(3)	SCHR0063
84	M=M-1	SCHR0064
86	GO TO 46	SCHR0065
C		SCHR0066
90	PM=P(M)	SCHR0067
	MSAVE = M	MOD 8
92	YIN=Y(2)/PM	SCHR0068
94	DO 96 J=M,N	SCHR0069
96	P(J)=P(J)/PM	SCHR0070
C		SCHR0071
CSTART OUTWARD INTEGRATION	SCHR0072
C		SCHR0073
100	P(1)=1.E-20	SCHR0074
102	Y=0.	SCHR0075
104	GI=V-E	SCHR0076
106	Y(2)=(1.-HV*GI)*P	SCHR0077
	K = 0	SCHR0078
108	DO 132 I=2,M	SCHR0079
110	Y(3)=Y(2)+((Y(2)-Y)+H2*GI*P(I-1))	SCHR0080
112	GI=V(I)-E	SCHR0081
114	P(I)=Y(3)/(1.-HV*GI)	SCHR0082
CTEST FOR OVERFLOW	SCHR0083
116	IF(K)118,130,118	SCHR0084
118	K=0	SCHR0085
	I1=I-1	SCHR0086
	PM=P(I1)	SCHR0087
	DO 120 J=1,I1	SCHR0088
120	P(J)=P(J)/PM	SCHR0089
122	Y=Y/PM	SCHR0090
124	Y(2)=Y(2)/PM	SCHR0091
126	Y(3)=Y(3)/PM	SCHR0092
	GI=V(I1)-E	SCHR0093
	GO TO 110	SCHR0094
C		SCHR0095
130	Y=Y(2)	SCHR0096
132	Y(2)=Y(3)	SCHR0097
C		SCHR0098
CFINISHED OUTWARD INTEGRATION	SCHR0099
134	PM=P(M)	SCHR0100
	IF(PM)135,149,135	SCHR0101
135	YOUT=Y/PM	SCHR0102
136	YM=Y(3)/PM	SCHR0103
138	DO 140 J=1,M	SCHR0104
140	P(J)=P(J)/PM	SCHR0105
C		SCHR0106
CCORRECTION	SCHR0107
C		SCHR0108
142	DF=0.	SCHR0109
144	DO 146 J=1,N	SCHR0110
146	DF=DF-P(J)**2	SCHR0111
148	F=(-YOUT-YIN+2.*YM)/H2+(V(M)-E)	SCHR0112
	DOLD=DE	SCHR0113
	IF(K)149,150,149	SCHR0114
149	F=9.999999E+29	SCHR0115
	DF=-F	SCHR0116

```
DE=ABSF(.0001*E)
GO TO 152
150 DE=-F/DF
152 IF(NI)164,158,154
154 IF(NI-10)162,164,164
156 FORMAT( 1H0 14,2X 1P4E16.7, 5X 31H THE CROSSING POINT OCCURS AT
1 I 14)
158 PRINT 156,IT,E,F,DF,DE
160 GO TO 164
162 WRITE OUTPUT TAPE NI, 156, IT,E,F,DF,DE,MSAVE
164 EOLD = E
E=E+DE
166 TEST=MAX1F(ABSF(DOLD)-ABSF(DE),TEST)
168 IF(TEST)171,170,170
170 IF( ABSF( E-EOLD) - EPS ) 172,172,171
171 CONTINUE
SCHR=1.
GO TO 173
C .....CONVERGED-COUNT NODES
172 SCHR=0.
173 KV=0
NL=N-2
174 DO 192 J=3,NL
176 IF(P(J))178,177,177
177 IF(P(J-1))180,192,192
178 IF(P(J-1))192,270,184
C POS. NODE
180 IF(P(J+1))192,182,182
182 IF(P(J-2))190,192,192
C NEG. NODE
184 IF(P(J+1))186,192,192
186 IF(P(J-2))192,190,190
C IGNORE FALSE NODE DUE TO UNDERFLOW
270 IF(P(J+1))280,192,192
280 IF(P(J-2))192,192,190
190 KV=KV+1
192 CONTINUE
C .....NORMALIZE
200 SN=SQRTF(-H*DF)
202 DO 204 J=1,N
204 S(J)=P(J)/SN
C .....PRINT SOLUTION
206 IF(NS)236,210,208
208 IF(NS-10)210,236,236
210 DO 234 JF=1,N,300
212 IF(NS)236,216,218
214 FORMAT(47H1SCHR- SOLUTION OF RADIAL SCHR. EQUATION FOR V= 13, 7H
X E= 1PE15.7 /20H0 I S(I) 5(20H I S(I) ))
216 PRINT 214,KV,E
GO TO 220
218 WRITE OUTPUT TAPE NS,214,KV,E
220 JL=XMINOF(JF+49,N)
222 DO 234 J=JF,JL
224 IL=XMINOF(J+250,N)
226 IF(NS)236,230,232
228 FORMAT(6(I5,1PE15.7))
230 PRINT 228,(I,S(I),I=J,IL,50)
GO TO 234
232 WRITE OUTPUT TAPE NS,228,(I,S(I),I=J,IL,50)
234 CONTINUE
```

SCHR0117
SCHR0118
SCHR0119
SCHR0120
SCHR0121
MOD 9
MOD 10
SCHR0123
SCHR0124
MOD 11
SCHR0126
SCHR0127
SCHR0128
SCHR0129
SCHR0130
SCHR0131
SCHR0132
SCHR0133
SCHR0134
SCHR0135
SCHR0136
SCHR0137
SCHR0138
SCHR0139
SCHR0140
MOD 12
SCHR0142
SCHR0143
SCHR0144
SCHR0145
SCHR0146
SCHR0147
MOD 13
MOD 14
MOD 15
SCHR0148
SCHR0149
SCHR0150
SCHR0151
SCHR0152
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SCHR0168
SCHR0169
SCHR0170
SCHR0171
SCHR0172

```
236   E0=E                                     SCHR0173
      IF(NS-10)250,875,250                     MOD 16
      875   WRITE OUTPUT TAPE 3, 876, KV,E     MOD 17
      876   FORMAT(54H0 SOLUTION OF THE RADIAL SCHRODINGER EQUATION FOR V = MOD 18
1      I3,8H   E = 1PE15.7 )                   MOD 19
250   RETURN                                    SCHR0174
      FREQUENCY 52(0,1,0),70(0,0,1),72(0,0,1),94(100),55(50),108(100),11 SCHR0175
X      6(0,1,0),138(100),144(200),152(1,0,0),154(0,1,1),202(200) SCHR0176
      END                                       MOD 20
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