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ANGULAR CORRELATIONS IN THE PRODUCTION AND DECAY OF SPIN- $\frac{3}{2}$ HYPERONS

Richard Spitzer and Henry P. Stapp

October 4, 1957

Printed for the U.S. Atomic Energy Commission

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> Radiation Laboratory University of California Berkeley, California

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ABSTRACT

Simple relationships between the angular distribution in the decay of a spin- $\frac{3}{2}$ hyperon and the directions defined by the production mechanism are obtained for the case in which only final S and P waves contribute to the production process.

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The general expressions given by Spitzer and Stapp in UCRL-3796 (Rev)¹ (hereafter referred to as SS) relating the angular distribution of the decay products of a $\operatorname{spin} \frac{3}{2}$ hyperon to the parameters of the production process may be reduced to a rather simple form if only S and P waves are included.

Let the production cross section be written

 $I(\theta) = A + B \cos \theta + C \cos^2 \theta$.

Equations (2.13), (2.18), and (2.19) and Tables I and II of SS show that neglecting the \underline{N} dependent terms (i.e., averaging over up-down), one may write the decay angular distribution in the form

 $W(\theta, (\widehat{H}), (\widehat{H})') = \frac{1}{4\pi} \left[1 + I^{-1}(\theta)(\alpha + \beta \cos \theta + \gamma \cos^2 \theta)(3 \cos^2 (\widehat{H})' - 1) + I^{-1}(\theta)(\delta + \epsilon \cos \theta)(3 \cos (\widehat{H}) \cos (\widehat{H})' - \cos \theta) \right]$

+ $I^{-1}(\theta)\rho(3\cos^2(\hat{H} - 1))$

Richard Spitzer and Henry P. Stapp, Polarization and Angular Correlation in the Production and Decay of Particles of Spin $\frac{1}{2}$ and Spin $\frac{3}{2}$, UCRL-3796 (Rev.), July 1957. The coefficients α , β , γ , etc. are constants, and the angles are defined (see SS)

= X (initial nucleon, hyperon),

 $(\widehat{\mathbb{H}}) = \Delta$ (hyperon, final nucleon),

 $(\mathbb{H})' = \mathcal{A}$ (initial nucleon, final nucleon).

This may be reduced to the form

$$\pi I(\theta)W(\theta, (\mathbb{H}^{'}; \Phi) = I(\theta) + \left\{ \alpha + (\beta + \delta)\cos\theta + (\gamma + \epsilon + \rho)\cos^{2}\theta \right\}$$

$$\times (3 \cos^{2} (\mathbb{H}^{'} - 1) + (\delta + \epsilon \cos\theta)(3 \cos (\mathbb{H}^{'}) \sin (\mathbb{H}^{'}) \sin\theta \cos\Phi)$$

$$+ \rho(6 \cos (\mathbb{H}^{'}) \cos\theta \sin (\mathbb{H}^{'}) \sin\theta \cos\Phi)$$

+ $\rho \sin^2 \theta (3 \sin^2 (H) \cos^2 \Phi - 1)$,

where Φ is the azimuthal angle of the final nucleon with respect to a polar axis along the velocity of the initial nucleon ($\Phi = 0$ if $\widehat{H} = 0$). From Eqs. (2.12), (2.13), and (2.18) and Tables I and II of SS, one verifies

> $\beta + \delta = \frac{1}{2} B ,$ $\gamma + \epsilon + \rho = \frac{1}{2} (A + C) - \alpha .$

Substituting these and averaging over Φ , one obtains

 $4\pi I(\theta)W(\theta, \widehat{\mathbb{H}}') = I(\theta)(\frac{3}{2}\cos^2(\widehat{\mathbb{H}}' + \frac{1}{2}) - (A + \rho - 2\alpha)\sin^2\theta(\frac{3}{2}\cos^2(\widehat{\mathbb{H}}' - \frac{1}{2}))$

If the production angular distribution $I(\theta)$ is known, then the Φ -averaged decay angular distribution is completely determined by the single additional (constant) parameter (A + ρ - 2α). This parameter is positive and is

smaller than $\left[A - C + \sqrt{(A + C)^2 - B^2}\right] \leq 2A$. This expression provides a convenient test of the hypothesis that the hyperon has $\operatorname{spin}_2 \frac{3}{2}$. In the limiting case $\sin \Theta \to 0$ one obtains the form given by Adair.²

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R. K. Adair, Phys. Rev. 100, 19540 (1955).

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