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### Title

Custom-designed heat treatment simultaneously resolves multiple challenges facing 3D-printed single-crystal superalloys

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16 **Abstract**

17 Single-crystal Ni-based superalloys are currently the material of choice for turbine blade  
18 applications, especially with the emerging additive manufacturing (AM) that facilitates  
19 the manufacture/repair of these single crystals. This promising AM route, however,  
20 comes with a dilemma: in the fusion and heat affected zones after e-beam or laser  
21 induced melting, one needs a solutionizing annealing at a sufficiently high temperature to  
22 relieve the residual stresses and homogenize the chemical/microstructure. The  
23 solutionizing temperature is usually adopted from the protocol for the cast superalloys,  
24 but this heat treatment almost always causes recrystallization and stray grain growth,  
25 resulting in a polycrystalline microstructure that degrades the high-temperature  
26 mechanical performance. Here we demonstrate a post-printing sub-solvus solutionizing  
27 treatment to replace the conventional super-solvus one. The recovery and relatively low  
28 temperature diminish the driving force for recrystallization and the movement of stray  
29 grain boundaries, without suffocating the chemical/microstructural homogenization  
30 thanks to the narrow dendrite width and short element segregation distance. The duration  
31 of the sub-solvus solutionizing treatment is optimized to achieve atomic-diffusion  
32 mediated chemical homogenization while limiting  $\gamma'$ -particle coarsening in the  
33 interdendritic regions. Our solution therefore removes a seemingly formidable obstacle to  
34 effective 3D-printing of superalloy single crystal products.

35

36 **Keywords:**

37 3D-printing manufacture/repair; Ni-based superalloy single crystals; heat treatment;  
38 recrystallization and stray grain growth; chemical homogenization

## 39        **1. Introduction**

40        Ni-based superalloy single crystals are now widely used for turbine blades and vanes  
41 in modern aerospace industries <sup>1</sup>. However, making complicated shapes and internal  
42 cooling passages <sup>2</sup> in these single crystals has turned out to be much more difficult and  
43 expensive, when compared with conventional precision investment casting <sup>3</sup>. In recent  
44 years, 3D printing, also known as additive manufacturing, has emerged as a powerful  
45 solution to this problem, not only shortening the processing chain and minimizing the  
46 waste, but also providing the possibility to repair damaged and/or worn single crystal  
47 superalloy parts to extend their service life <sup>4,5</sup>. Rejuvenation (restoration/repair) of  
48 microstructures is effective in extending the life of superalloy blades <sup>6</sup>, whereas 3D-  
49 printing enables precise shape control. However, while crack free superalloy single  
50 crystals have been successfully printed <sup>7-9</sup>, their microstructure often fails to meet the  
51 homogeneity <sup>10</sup> and stability <sup>11</sup> requirements. As a result, it is always necessary to devise  
52 a suitable post-printing heat treatment, under either ambient or high pressure <sup>12-14</sup>.

53        To achieve excellent high-temperature performance, Ni-based superalloys  
54 incorporate tens of alloying elements, including refractory elements like Re, W and Mo,  
55 to form Ni<sub>3</sub>(Al,Ti)  $\gamma'$ -precipitates with L1<sub>2</sub> structure that are coherent with the solid-  
56 solution-strengthened austenitic  $\gamma$ -matrix <sup>15,16</sup>. When produced with the 3D printing  
57 approach, the superalloy is deposited layer-by-layer through local melting of the powder  
58 feedstock with designed chemical constituents using either a laser or an electron beam  
59 heat source <sup>17-19</sup>. Superalloy single crystals are produced under a steep temperature  
60 gradient <sup>20</sup>, which outperform their polycrystalline counterparts in terms of resistance to  
61 creep and fatigue <sup>7,11</sup>. Another advantage brought by the high cooling rate is the refined

62 dendrite structure (several micrometers in width <sup>21</sup>) that provides further strengthening  
63 compared to cast superalloys in which the dendrites are usually hundreds of microns wide  
64 <sup>22</sup>. Stray grains tend to form on metal surface where the temperature gradient is no longer  
65 parallel to the building direction <sup>23,24</sup>, thus they are inevitable in 3D-printed superalloy  
66 single crystals on both the outer surface of the blades and the inner surfaces of the  
67 internal cooling structures <sup>7</sup>. The outer surface can be machined and/or milled to regain  
68 the single crystalline structure, but not the inner surfaces of the cooling structures.  
69 Fortunately the local service temperature near the inner surfaces is lower than that at the  
70 outer surface by several hundred degrees Celsius; the stray grains there are bearable as  
71 long as they do not grow much larger into the interior <sup>25,26</sup>. During solidification, the  $\gamma$ -  
72 matrix forms first, and ejects the  $\gamma'$ -forming elements into the remaining liquid in-  
73 between the dendrites, lowering the liquidus temperature <sup>27</sup>. Once the interdendritic liquid  
74 solidifies, its solvus temperature is elevated due to enriched  $\gamma'$ -formers, and as a result the  
75  $\gamma'$ -precipitates in these regions form earlier and grow bigger than those in dendrite cores  
76 (DCs). Also, the  $\gamma'$ -precipitate volume fraction and mechanical properties become  
77 inhomogeneous across the dendrite width <sup>28-31</sup>. Another issue arising from the rapid  
78 cooling rate is the high thermal stress and plastic deformation <sup>32,33</sup>. The stored  
79 deformation energy provides the driving force for recrystallization in the bulk, and for  
80 stray grain boundary migration near the surface, both of which are detrimental or even  
81 disastrous, as they ruin the single-crystalline microstructure desired for retaining high-  
82 temperature mechanical properties.

83 From the discussion above, there appear to be four major challenges facing the post-  
84 printing heat treatment for 3D-printed superalloy single crystals. Specifically, a

85 successfully customized protocol must be able to (1) *r*elease most of the stored  
86 deformation energy, (2) *a*void recrystallization completely, (3) *s*uppress stray grain  
87 growth as much as possible, and (4) *h*omogenize the chemical/microstructure distribution  
88 to a level comparable with that in cast alloys. From here on these requirements are  
89 acronymized as the ***RASH*** challenges. In this paper, we evaluate previously reported heat  
90 treatment protocols against these four demands, and then design a novel and yet simple  
91 heat treatment, which will be demonstrated to accomplish all the ***RASH*** actions via a  
92 single-step solutionizing annealing at sub-solvus temperature prior to aging treatment.

93 This new strategy is conceived based on the following two considerations. First,  
94 since the diffusion distance, which is proportional to dendrite width, is greatly reduced in  
95 the 3D-printed superalloy due to the fast cooling rate, chemical homogenization at a level  
96 similar to the super-solvus solutionizing treatment in the traditional cast products can be  
97 achievable with a sub-solvus annealing treatment. Second, by setting the solutionizing  
98 temperature above the solvus of the  $\gamma'$ -precipitates in the DCs while below that in  
99 interdendritic regions (IRs), the dislocations moving freely in the DCs allow for a speedy  
100 recovery. In the meantime, the undissolved  $\gamma'$ -precipitates in IRs would be able to impede  
101 the massive motions and interactions of dislocations, nucleation of recrystallized grains,  
102 and migration of stray grain boundaries <sup>6</sup>. The success of such a single-step sub-solvus  
103 solutionizing treatment will be demonstrated in the following, in two types of superalloy  
104 single crystals with different chemical constituents produced using either a laser or an  
105 electron beam heat source.

106

## 107 **2. Results**

## 2.1. Necessity for sub-solvus solutionizing heat treatment

Compared to the traditional casting technique, the solidification and cooling processes during 3D printing are faster by several orders of magnitude. In the resultant 3D-printed superalloy, the dendrite width is narrower. The  $\gamma'$ -precipitates are more irregular in morphology, smaller in size, lower in volume fraction, and less stable against temperature excursion<sup>9,34-36</sup>. Despite these differences between casting and 3D-printing, the micro-segregation in the resulting superalloy single crystals is similar. Thus in many cases standard heat treatment protocols of the cast superalloys have been applied to their 3D-printed counterparts, with no modification or just shortening the homogenization treatment duration<sup>13</sup>. As illustrated in the upper panel of Figure 1a, the standard heat treatment of cast superalloys consists of hours of super-solvus solutionizing treatment to homogenize the chemical distribution, followed by a long period of aging treatment to precipitate out, ripen and stabilize the  $\gamma'$ -precipitates<sup>37-39</sup>. However, with such a heat treatment, recrystallization sets in readily from the heat affected zone (HAZ), and the newly formed recrystallization grains as well as the existing stray grains grow big quickly as the grain boundary mobility is high once the specimen is heated above the solvus temperature (Figure 1b). Such observations have been reported previously by many<sup>13,40</sup>. This recrystallization, rendering the single crystal polycrystalline, wastes all the efforts that have been made to achieve the single-crystalline microstructure desired for high temperature mechanical properties. Clearly, the standard heat treatment cannot meet the **RASH** challenges for 3D-printed superalloy single crystals.

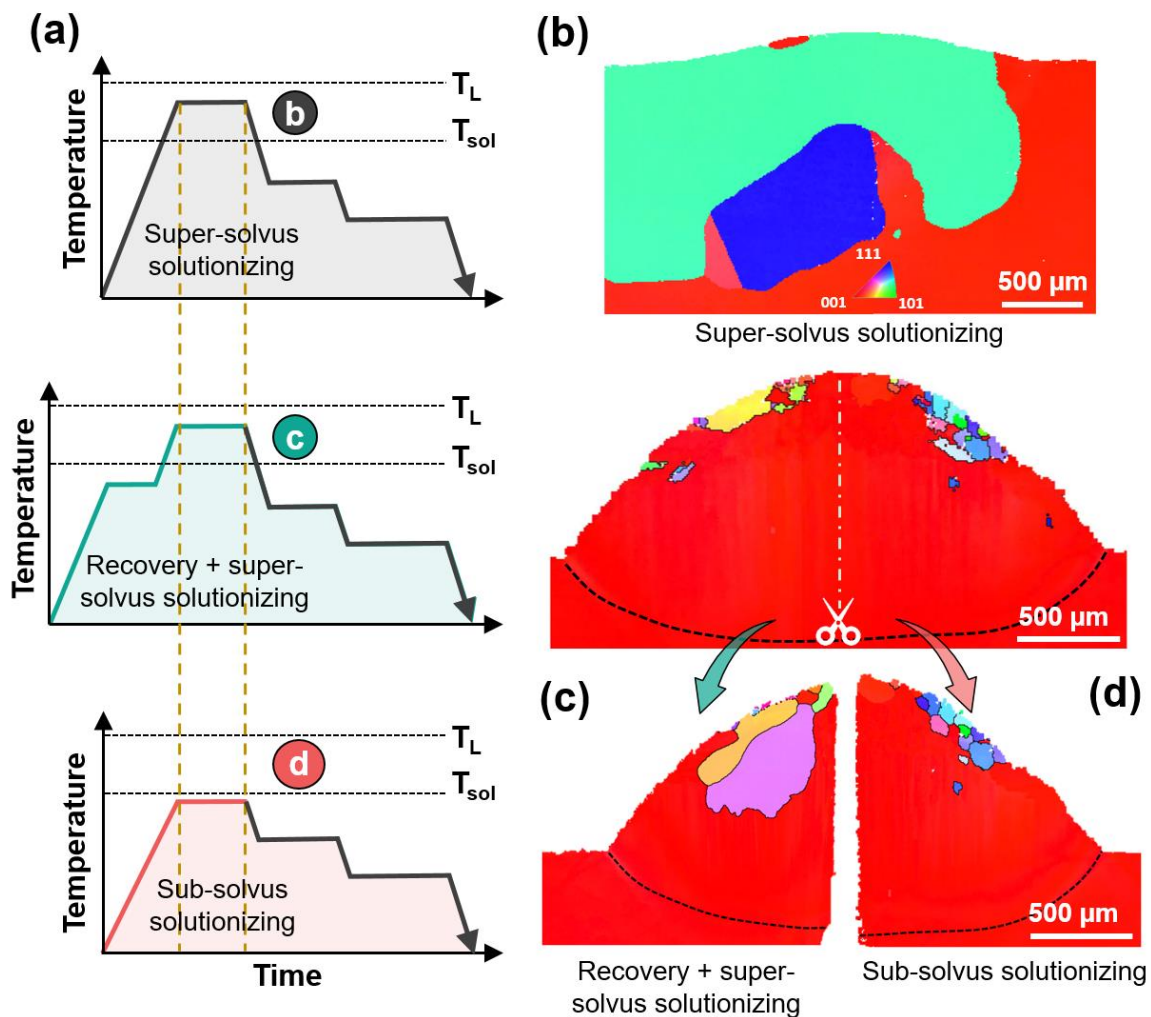
Efforts have been made before, to solve the recrystallization issue by applying a pre-solutionizing annealing step to the 3D-printed superalloy single crystals at sub-solvus

131 temperature, as shown in the middle panel of Figure 1a. With this step, a large fraction of  
132 the stored deformation energy is released prior to solutionizing treatment, taking away the  
133 driving force for recrystallization <sup>10</sup>. Afterwards, the standard ageing process is  
134 employed, which produces cuboidal  $\gamma'$ -precipitation microstructure indistinguishable  
135 from the cast base metal. However, stray grain growth is not precluded during the super-  
136 solvus solutionizing treatment, because the migration of high-angle grain boundaries is  
137 activated, as indicated by the crystal orientation distribution (inverse pole figure) maps  
138 obtained from electron backscatter diffraction (EBSD) scans before and after (Figure 1c)  
139 heat treatment of a laser 3D-printed single-crystalline superalloy AM3 specimen. Similar  
140 observations have also been recorded in electron beam melted superalloy single crystals  
141 (Figure S1). This necessitates a surface subtractive machining process before super-  
142 solvus solutionizing heat treatment, to get rid of the seeds of stray grains. Such an extra  
143 machining step is not feasible for internal cooling structures, and is in any case time-  
144 consuming and costly.

145       Considering that both recrystallization and stray grain growth occur under the  
146 precondition of the complete solid-solutioning of  $\gamma'$ -precipitates, we have conceived a  
147 different heat treatment protocol. It also consists of two steps of solutionizing and aging  
148 treatment, but the solutionizing treatment is carried out at sub-solvus temperature (1270  
149 °C for 30 min in this case) instead of the conventional super-solvus one, as displayed in  
150 the bottom panel of Figure 1a. The specimen is exactly the same as the one used for  
151 recovery pre-annealing plus super-solvus solutionizing heat treatment, but a pronounced  
152 difference in stray grain size is seen in Figure 1d in laser 3D-printed superalloy single  
153 crystal AM3. As shown in Figure S1 in the Supporting Information, when the additive



154 manufacturing heat source is changed to electron beam, this method still works well,  
 155 although here no feedstock is supplied. In other words, combining sub-solvus recovery  
 156 annealing and super-solvus solutionizing into a single sub-solvus solutionizing treatment  
 157 meets the “avoid recrystallization” and “suppress stray grain growth” requirements. In the  
 158 next section, we will further tune the temperature and duration of sub-solvus annealing  
 159 treatment to homogenize the chemical and microstructure distribution.



160

161 **Figure 1** Crystal orientation distributions resulted from various heat treatment protocols. (a) Three different  
 162 heat treatment protocols are considered. (b) Direct super-solvus treatment results in recrystallized  
 163 microstructure. (c) Recovery (pre-annealing) before super-solvus treatment helps to prevent  
 164 recrystallization, but stray grains grow bigger. (d) Sub-solvus solutionizing at 1270 °C eliminates both  
 165 recrystallization and stray grain growth.  $T_L$  and  $T_{sol}$  in (a) stand for liquidus and solvus temperatures,  
 166 respectively. (b-d) are the inverse pole figures obtained from EBSD mapping scanned with 10 μm, 3 μm,  
 167 and 3 μm step size, respectively.

168

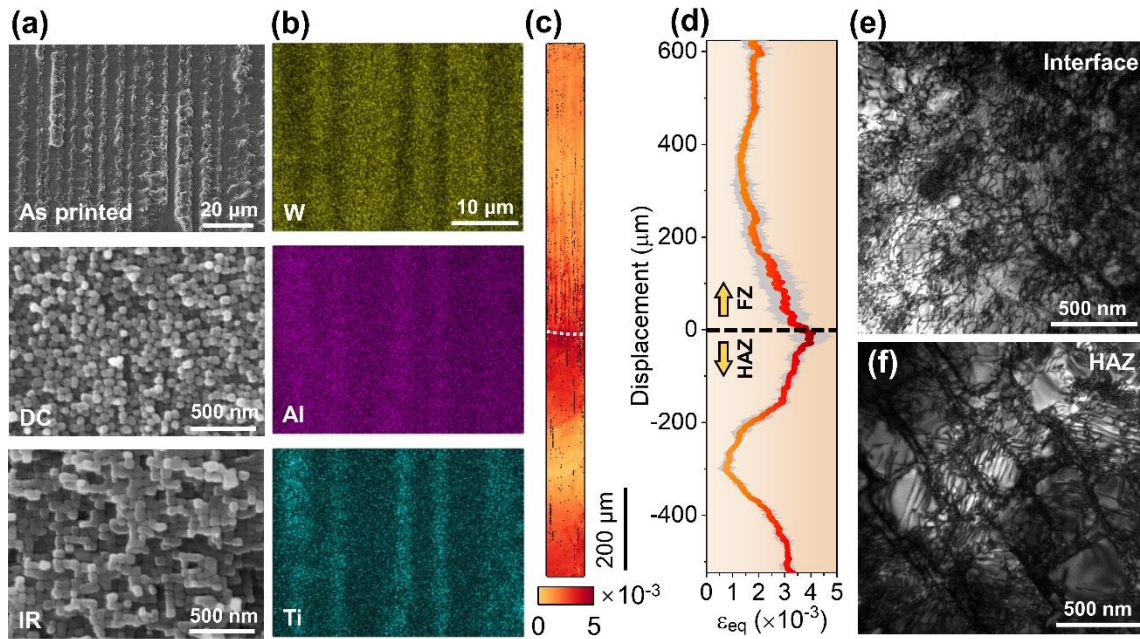
## 169 **2.2. Optimization of annealing temperature and duration**

170 In order to meet all the *RASH* requirements, the annealing temperature and duration  
171 need to be carefully tailored and optimized. Firstly, the microstructure and residual strains  
172 of the as-electron-beam-printed superalloy single crystal are demonstrated in Figure 2.  
173 The  $\gamma'$ -particles are tiny (30 - 50 nm in diameter) and irregular in the DCs, whereas in the  
174 IRs they appear in almost cuboidal shape with rounded corners and edge length of 70 -  
175 100 nm. Based on the contrastive size and shape distribution, the dendrite widths are  
176 measured to span from 3 to 10  $\mu\text{m}$  in the fusion zone (FZ). Micro-segregation is clearly  
177 seen from the X-ray wavelength dispersive spectra (WDS) maps, showing enriched W in  
178 DCs, and Ti and Al in IRs. After electron beam melting, the crystal is riddled with  
179 dislocations<sup>34</sup>. The elastic strain associated with these defects is measured using  
180 synchrotron based X-ray microdiffraction ( $\mu\text{XRD}$ ). An area of 80  $\mu\text{m}$  (horizontal)  $\times$  1150  
181  $\mu\text{m}$  (vertical) across the FZ and the heat affected zone (HAZ) is scanned using micro-  
182 focused polychromatic X-ray beam with 2  $\mu\text{m}$  spatial resolution, as marked in Figure  
183 S2a, and lattice strain tensor at each scanning position is measured from the Laue  
184 diffraction pattern. The equivalent strain  $\varepsilon_{eq}$  is calculated from the strain tensor, as a  
185 representative of the magnitude of the local strain,

186 
$$\varepsilon_{eq} = \frac{\sqrt{2}}{3} \sqrt{(\varepsilon_{xx} - \varepsilon_{yy})^2 + (\varepsilon_{yy} - \varepsilon_{zz})^2 + (\varepsilon_{zz} - \varepsilon_{xx})^2 + 6\varepsilon_{xy}^2 + 6\varepsilon_{xz}^2 + 6\varepsilon_{yz}^2}$$
, where  $\varepsilon_{ij}$  is one of

187 the six strain tensor components. As shown in Figure 2c, the equivalent strain varies  
188 pronouncedly in the building direction, but is quite uniform in the horizontal direction; in  
189 Figure 2d the average equivalent strain magnitude is therefore plotted as a function of the

190 distance to the melting line. The highest lattice strain,  $\sim 4 \times 10^{-3}$ , appears near the  
191 interface between FZ and HAZ. It decreases gradually into FZ, reaches the minimum  
192 value of  $1.5 \times 10^{-3}$  at the position about 350  $\mu\text{m}$  away from the interface, and increases  
193 slightly to an almost steady  $2 \times 10^{-3}$ . In the HAZ, the strain distribution is even more non-  
194 uniform. The equivalent strain drops to almost zero in a range of 300  $\mu\text{m}$  and then  
195 increases back to  $\sim 3 \times 10^{-3}$ . Although the exact elastic stiffness tensor for AM3 single  
196 crystal was not obtained, the peak tensile residual stress is estimated (assuming a  
197 modulus of the order  $\sim 200$  GPa) to be on the order of 0.8 GPa, which is an unacceptably  
198 high driving force from the recrystallization and high-angle grain boundary migration  
199 points of view. The dislocation density and structure, which is strongly influenced by not  
200 only the stress but also the  $\gamma/\gamma'$  structure, are studied under a transmission electron  
201 microscope (TEM). As seen in the bright-field image in Figure 2e, curly dislocations near  
202 the melting line, where only tiny  $\gamma'$ -particles exist, are heavily tangled, with a density of  
203 approximately  $6 \times 10^{14} \text{ m}^{-2}$  measured using the line-intercept method<sup>41</sup>. In the HAZ, the  
204 density of dislocations is similar to the interfacial area, while the structure is quite  
205 different. Because the primary  $\gamma'$ -particles are only partially dissolved,  $\gamma$ -channels are  
206 narrower and dislocations are straight and parallel, with pile-ups at the  $\gamma/\gamma'$  interfaces<sup>42</sup>.



207

208 Figure 2 Microstructure of the electron-beam-melted AM3 superalloy single crystal. Distributions of (a)  $\gamma'$   
 209 particles, (b) constituent elemental species, (c-d) equivalent strain, and (e-f) dislocation structures are all  
 210 inhomogeneous.  
 211

212 Secondly, the solvi of the  $\gamma'$ -precipitates in DCs ( $T_1$ ) and IRs ( $T_2$ ) are measured by  
 213 annealing the specimen at various temperatures for a constant period of time (15 min in  
 214 this study) and then monitoring the evolution of the SEM image contrast. Once the  
 215 electron-beam-printed AM3 superalloy was annealed above  $T_1$  (1260 °C, as displayed in  
 216 Figure 3a), the  $\gamma'$ -precipitates in the DCs are fully dissolved and then re-precipitate as  
 217 bigger and almost uniform cuboids; in the meanwhile, the  $\gamma'$ -precipitates in the IRs are  
 218 partially dissolved, resulting in the disappearance of the contrast between DC and IR in  
 219 some regions. As the annealing temperature goes above  $T_2$  (1280 °C in this case), the  
 220 contrast between DC and IR becomes almost homogeneous. These observations suggest  
 221 that the optimized temperature for the sub-solvus homogenization heat treatment lies  
 222 between the solvi from 1260 °C to 1280 °C.

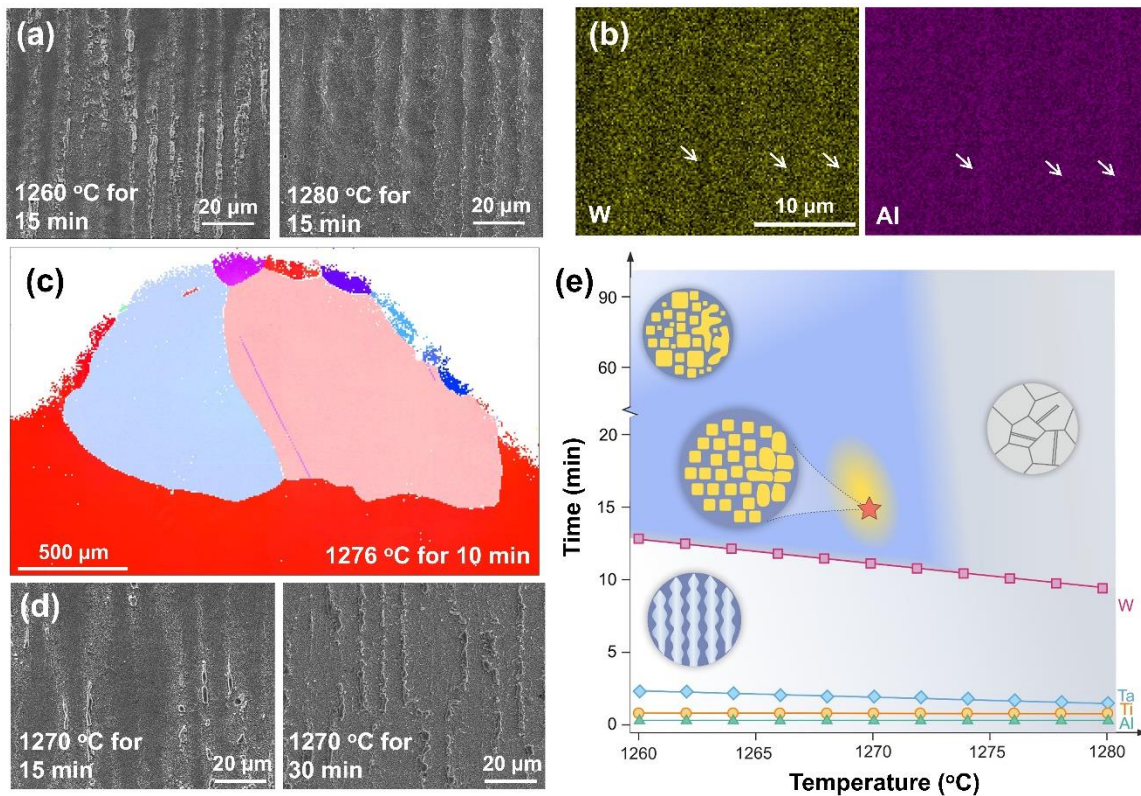
223 To guide the selection of the duration time, the classic diffusion equation  $d = \sqrt{4Dt}$  is

224 employed. By taking the reported diffusion coefficient  $D$  of alloying elements Al, Ti, Ta,  
225 and W<sup>43</sup> and setting the diffusion distance  $d$  as half of the dendrite width (5  $\mu\text{m}$ ), the  
226 diffusion time  $t$  is calculated for each temperature between 1260 °C and 1280 °C. The  
227 pace setter, W, migrates the most slowly and the annealing time needs to be longer than  
228 13 min when the specimen is heat treated at 1260 °C, and 10 min for 1280 °C. For  
229 experimental verification of the calculation, the 3D-printed specimen is annealed at 1270  
230 °C for 10 min, ~15% shorter than the calculated time. The chemical composition  
231 distribution is found still inhomogeneous, as seen in the WDS map in Figure 3b.  
232 Therefore, in the following the annealing time of all experiments is no shorter than the  
233 calculated values.

234 The resulted microstructure is sensitive to the annealing temperature. As displayed in  
235 Figure 3c, after annealing at 1276 °C for 10 min, stray grains grow rapidly and  
236 overwhelm the whole electron-beam-printed volume. Consequently, the desired annealing  
237 temperature should be set below 1276 °C to fulfill the “suppress stray grain growth” and  
238 “avoid recrystallization” requirements.

239 In order to understand how temperature and time influence the homogeneity of the  $\gamma'$   
240 microstructure, identical specimens were sub-solvus solutionizing heat treated at 1270 °C  
241 for 15 min and 30 min, respectively (Figure 3d). For the same period of time (15 min),  
242 1270 °C annealing results in the  $\gamma'$  microstructure with less contrast than that annealed at  
243 1260 °C; while extending the annealing time (from 15 min to 30 min, or even longer as  
244 shown in Figure S3) results in more pronounced contrast and thus microstructural  
245 inhomogeneity. Overall, using a shorter time at a higher temperature generates a more  
246 homogeneous microstructure.

247 By plotting the observations above into a single “treasure map”, we close in on the  
 248 desired heat treatment protocol, as demonstrated in Figure 3e. In this map, the boundaries  
 249 are demarcated by the several considerations outlined in the preceding paragraphs.  
 250 Specifically, the grey areas in this map are not acceptable, as these conditions would  
 251 subject the alloy to either recrystallization/stray grain growth (increasing in degree when  
 252 moving to the right), or inhomogeneous chemical composition (increasingly  
 253 inhomogeneous moving towards the bottom). Only in the blue region, can  
 254 recrystallization and stray grain growth be successfully prevented, and the chemical  
 255 uniformity achieved. The bottom right corner of the blue region, indicated by the yellow  
 256 zone, is the most favorable, because lower temperature together with longer annealing  
 257 time would be harmful for the  $\gamma'$  particle size homogeneity<sup>44</sup>. Going in this direction  
 258 leads us to the best spot (1270 °C for 15 min, as marked by a red star).



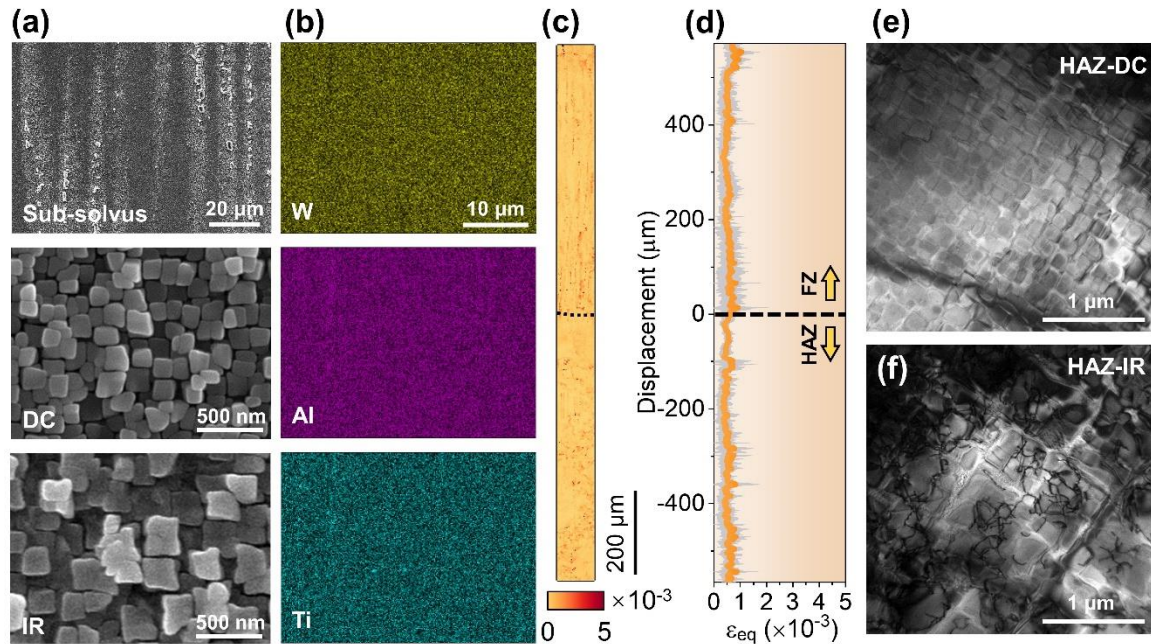
259



260 Figure 3 Optimization of heat treatment temperature and time. (a) The microstructure evolution of the  $\gamma'$ -  
261 particles in the electron-beam-printed AM3 superalloy, after annealing at 1260 and 1280 °C for 15 min. (c)  
262 is the inverse pole figures obtained from EBSD mapping scanned with 10  $\mu\text{m}$  step size. (d) The contrast of  
263  $\gamma'$ -particles becomes weaker after annealing at 1270 °C for 15 min. (e) shows a treasure map locating the  
264 optimized heat treatment parameters (yellow region with red star).  
265

### 266 2.3. Stress/strain and microstructural characterization

267 After our newly designed sub-solvus solutionizing heat treatment, the  $\gamma'$ -particles in  
268 the IRs are about 200 - 300 nm in edge length, bigger than those in the DCs (~100 nm).  
269 The  $\gamma$  channels in the IRs widens to ~30 nm, also wider than those in the DCs (only a few  
270 nm). No contrast remains between the DCs and IRs in the WDS maps of the FZ. The  
271 element maps of W, Al, and Ti displayed in Figure 4b indicate no detectable chemical  
272 inhomogeneity. The equivalent strain map in Figure 4c reveals directly that the lattice  
273 strain in the interfacial region, where the strain is high and inhomogeneous in the as-  
274 printed state, becomes low and uniform. Detailed analysis shows that the equivalent  
275 strain has a nearly constant magnitude at about  $0.5 \times 10^{-3}$ , which is roughly the lower-  
276 bound measurement limit of the  $\mu\text{XRD}$  technique. The dislocations still show a  
277 distribution that is moderately non-uniform, after sub-solvus heat treatment. The DCs  
278 become almost dislocation free, with only 2 dislocations in the observed area (Figure 4e).  
279 But in the IRs, the dislocation density is higher ( $5 \times 10^{13} \text{ m}^{-2}$ ) and more inhomogeneous,  
280 as observed in Figure 4f. From other TEM images taken in the same specimen, the  
281 dislocation density is in a range from  $3 \times 10^{13} \text{ m}^{-2}$  to  $1 \times 10^{14} \text{ m}^{-2}$ .



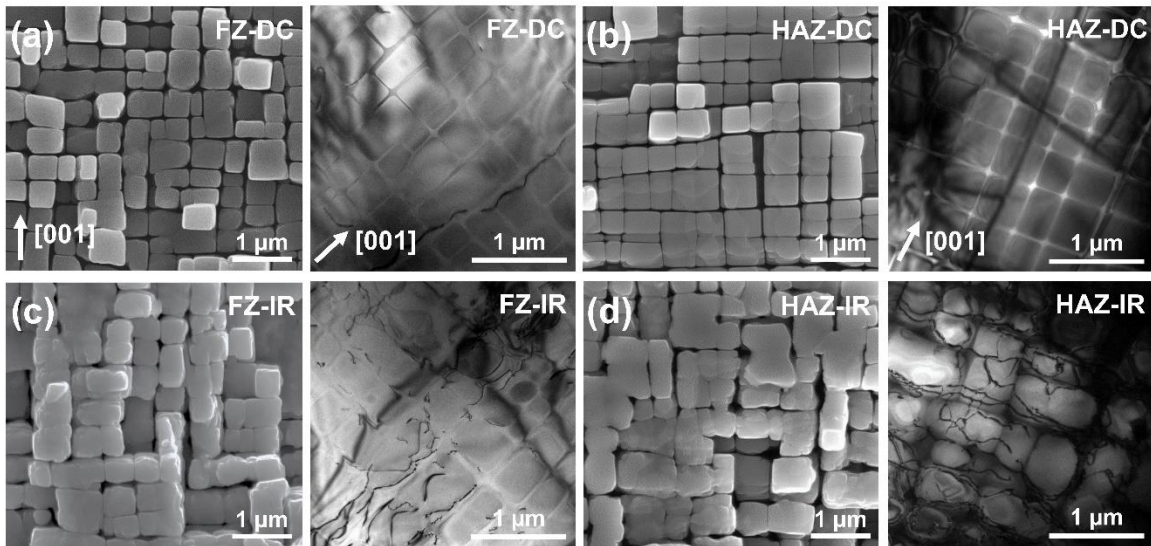
282

283 Figure 4 Microstructure after the sub-solvus solutionizing of the AM3 superalloy single crystal. The  
 284 distributions of (a) element and (b-c) equivalent strain become homogeneous, although (d) dislocation  
 285 densities in the DC and IR are different.  
 286

287 From the results above, we conclude that all the four *RASH* challenges are resolved  
 288 via the single step of sub-solvus solutionizing treatment. That is, the stored deformation  
 289 energy is released, recrystallization is avoided, stray grain growth is suppressed, and the  
 290 chemical distribution is homogenized. What is to be dealt with next, is the morphology of  
 291 the  $\gamma'$ -particles that subsequently evolve during ageing, and the remaining (left-over)  
 292 dislocation contents. After the sub-solvus solutionizing treatment at 1270 °C for 15 min,  
 293 the electron-beam-printed AM3 single crystal is aged following the standard heat  
 294 treatment protocol to evolve the  $\gamma'$ -particles. In the DCs, regular cuboidal  $\gamma'$ -particles with  
 295 the side length of 450 nm are obtained in both FZ and HAZ, and recovery reduces the  
 296 dislocation density (Figure 5a and b) to less than 1% of that in the as-printed single  
 297 crystals, and only individual dislocation can be detected occasionally. However, inside  
 298 the IRs the microstructure is not fully uniform. In the IRs in FZ,  $\gamma'$ -particles coarsen



299 moderately, and dislocations are observed to align parallel to the  $\gamma/\gamma'$  phase boundaries  
 300 (Figure 5c). The dislocation density in this region is  $5 \times 10^{13} \text{ m}^{-2}$ , about 10 times lower  
 301 than that in the as-printed state while 10 times higher than that in the DCs. In the IRs in  
 302 HAZ, directional coarsening, *i.e.* rafting, of  $\gamma'$ -precipitates occurs, similar to what has  
 303 been observed at the early stage of creep test when  $\gamma'$ -precipitates are not well connected  
 304 and thus  $\gamma$ - $\gamma'$  topological phase inversion has not yet set in. From the bright-field TEM  
 305 image in Figure 5d, the  $\gamma'$ -precipitates are embraced by dislocations; this is because the  
 306 majority of the dislocations trying to recover are blocked effectively by the strengthening  
 307 precipitates, while only a small proportion of dislocations penetrate into  $\gamma'$ -precipitates.  
 308 The dislocation density in this region is measured to be approximately  $10^{14} \text{ m}^{-2}$ , higher  
 309 than that in the IRs of FZ but still 5 times lower than that in the as-printed single crystals.  
 310 The microstructure distributions of the  $\gamma'$ -particles of the fully heat treated laser 3D-  
 311 printed superalloys (Figure S4 in Supporting Information) are quite similar to those in the  
 312 electron beam melted one.



313

314 Figure 5 Microstructure of fully heat treated AM3 superalloy single crystal.  $\gamma'$ -precipitates in the DCs of (a)  
 315 FZ and (b) HAZ are uniform and dislocations are completely annihilated. Coarsening of  $\gamma'$ -precipitates in  
 316 the IRs of (c) FZ and (d) HAZ occurs and dislocations tend to reside along the  $\gamma/\gamma'$  interface.

317

318 It must be noticed that the size and/or morphology of the  $\gamma'$ -particles in the fully  
319 conventionally heat-treated cast superalloys are not completely homogeneous either <sup>45</sup>,  
320 mainly because the dendrite widths of cast superalloys are on the order of hundreds of  
321 microns such that the chemical composition from the DCs to the IRs cannot be  
322 homogenized even when the solutionizing heat treatment is carried out above the solvus  
323 temperature. Therefore, the minor inhomogeneity resulted from our novel sub-solvus heat  
324 treatment is believed to be acceptable. More details about this inhomogeneity will be  
325 discussed further in the next section.

326

### 327 **3. Discussion**

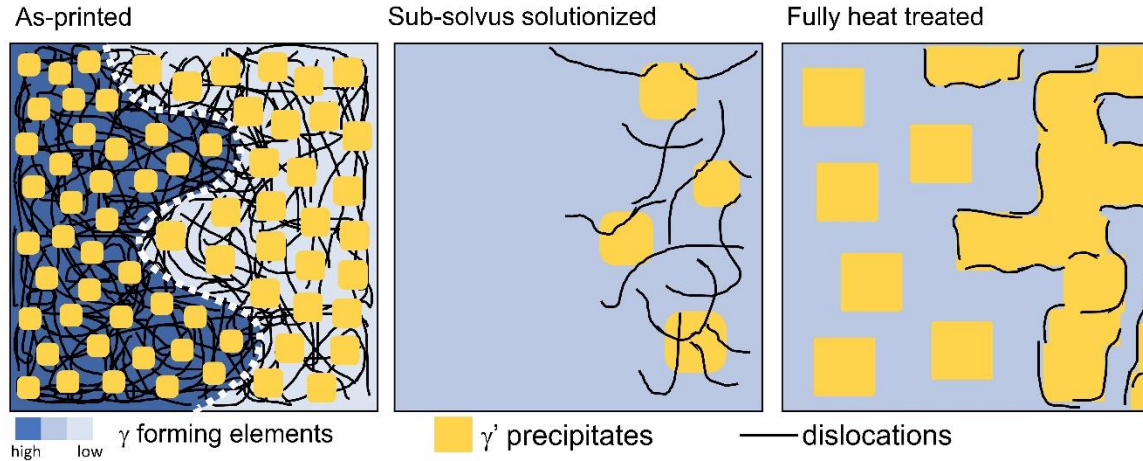
328 Cracking is known to be one of the difficulties facing 3D-printed Ni-based  
329 superalloys, it is also noticed that in most cases cracks occur along with the high-angle  
330 grain boundaries. In other words, cracking is more frequently observed in polycrystalline  
331 superalloys <sup>46,47</sup>. For superalloy single crystals, cracks are also observed when the  
332 orientation is not perfectly controlled <sup>33,48</sup>. As reported previously <sup>47</sup>, in the first few  
333 layers of 3D-printing deposition of superalloy on a polycrystalline stainless steel base  
334 metal, the as-printed superalloy inherits the crystal orientation of the base metal, and  
335 cracks are detected along with the high-angle grain boundaries. As more layers are  
336 cladded, [001] preferred orientation overwhelms, and eventually a single crystal is  
337 obtained. In the meanwhile, no more cracks exist in the single crystal parts. In our current  
338 study, a single crystal superalloy is employed as the base metal; therefore no cracks are  
339 observed in the printed and melted superalloys with both laser and electron beam as the

340 heat source. Minor porosity is observed in the 3D-printed superalloy single crystals,  
341 which leads to the concentration of strain/stress. But due to the small size of the pores in  
342 the 3D-printed superalloys observed in this study (Figure S5 in Supporting Information)  
343 and reported in previous literatures <sup>7</sup>, as well as the recovery effect of the sub-solvus  
344 treatment, no recrystallization is caused due to the porosity.

345 Post 3D-printing heat treatment, as one of the most effective approaches to tune the  
346 microstructure and thus mechanical properties of superalloys, has attracted a lot of  
347 attention. Because of the fine dendritic structures, shortening the heat treatment time has  
348 been proposed; however, the solutionizing temperatures in previous reports are still  
349 higher than the solvus of  $\gamma'$ -precipitates, under either ambient or high pressures <sup>13,14</sup>. This  
350 is because in many of these previous investigations the 3D-printed superalloys are  
351 polycrystalline, and thus recrystallization is acceptable. For 3D-printed single crystals,  
352 *RASH* issues pose major challenges – low temperature for short time leads to  
353 inhomogeneous chemical distribution, while high temperature (higher than 1320 °C for  
354 example) for even very short time may still trigger recrystallization. Exploiting the fact  
355 that the presence of  $\gamma'$ -particles can impede the motion of dislocations and high-angle  
356 grain boundaries, we have designed a novel heat treatment approach to achieve plastic  
357 deformation recovery and chemical homogenization with a single annealing step.  
358 Different from the previously employed super-solvus solutionizing heat treatment, which  
359 just adopts the solutionizing temperature from the standard heat treatment protocol  
360 established for the cast superalloys, the solutionizing temperature is optimized in a two-  
361 step manner: firstly measuring the solvi of  $\gamma'$ -particles in DCs and IRs, and then further  
362 specifying the critical temperature that does not induce recrystallization and stray grain

363 growth.

364 We now take a closer look at the mechanisms as to how the *RASH* issues are  
365 resolved. As illustrated in Figure 6, a high-density of dislocations, non-uniform chemical  
366 composition, and non-identical  $\gamma'$ -particle size/morphology are formed in the as-printed  
367 superalloys, although the non-uniformity is the most obvious at locations on the surface.  
368 When the temperature is elevated to above the solvus temperature (1280 °C for the AM3  
369 superalloy single crystal in this report) of the DCs but below the IRs, atomic diffusion  
370 easily covers half of the dendrite width (only a distance of 10  $\mu\text{m}$  or less in 3D-printed  
371 superalloys), to achieve compositional homogeneity. Since the  $\gamma'$ -particles in the DCs are  
372 totally dissolved, dislocations move readily without obstacles to mediate recovery. In the  
373 IRs, although the much wider  $\gamma$ -channels provide a spatial playground for dislocations to  
374 interact with one another,  $\gamma'$ -particles in these regions are not fully dissolved and the  
375 remnants impede the motion of dislocations and high-angle grain boundaries. This is why  
376 the dislocation density is brought down effectively and meanwhile recrystallization is  
377 largely avoided and stray grain growth is suppressed. These left-over dislocations and  
378 their associated deformation energy, during the aging heat treatment, provide the driving  
379 force for  $\gamma'$ -particle growth, which is similar to rafting during creep, except that the  
380 residual strain/stress caused by the dislocations is more inhomogeneous in both  
381 magnitude and direction on the microscopic scale, such that the  $\gamma'$ -particle coarsening is  
382 less directional than that in a conventional creep test.



383

384 Figure 6 Evolution of  $\gamma'$ -precipitate size/morphology and dislocation density/configuration during the sub-  
 385 solvus solutionizing and subsequent aging heat treatment. Light and dark blue colors in the as-printed state  
 386 indicate chemical inhomogeneity, which is homogenized after heat treatment.  
 387

388 An interesting question is whether it is possible, and how, to further enhance the  
 389 homogeneity. It is obvious that elevating the solutionizing heat treatment temperature  
 390 would promote the homogeneity of the size and morphology of  $\gamma'$ -precipitates, because  
 391 recovery would be more thorough, leaving less dislocations and lower residual  
 392 strain/stress after sub-solvus solutionizing treatment. However, higher solutionizing  
 393 temperature also increases the likelihood for recrystallization and stray grain growth,  
 394 especially the latter. As we observed, stray grains grow bigger significantly at 1276 °C  
 395 after annealing for 15 minutes (Figure 3c). Consequently, we conclude that while the  
 396 compositional homogeneity is relatively easy to achieve, the  $\gamma'$ -morphology homogeneity  
 397 is counter-balanced by the risk of forming high-angle grain boundaries that are prone to  
 398 migration. This trade-off could lead to an optimized solutionizing temperature. As for  
 399 treatment durations, systematic investigation shows that extending the annealing time to  
 400 30 min leads to even more severe  $\gamma'$ -coarsening (Figure S6). Reducing the annealing time  
 401 to 10 min or even shorter exposes the 3D-printed superalloys to the risk of inhomogeneity

402 in chemical composition. By balancing all these factors, we determined the heat treatment  
403 parameters to be 1270 °C for 15 min.

404 How much the non-uniform  $\gamma'$ -precipitates in the IRs would influence the high-  
405 temperature mechanical properties of the superalloy single crystal is of interest. To  
406 quantify the effects would require extensive future research but a rough estimate can be  
407 made here. It has been reported that the high temperature creep properties of crept  
408 superalloy single crystal René 5 could be restored by 82 – 85%, compared to the cast  
409 René 5 superalloy, by means of rejuvenation heat treatment with solutionizing at the  
410 temperature 28 °C below the solvus and then standard aging<sup>6</sup>. Although in this reference  
411 paper, no detailed  $\gamma'$  microstructure evolution during the creep test and rejuvenation was  
412 reported, it was believed that solutionizing at 28 °C below solvus must have resulted in an  
413 incomplete recovery, and thus rafted  $\gamma'$ -precipitates. In our study, the solutionizing  
414 temperature is only 10 °C below the solvus of the  $\gamma'$ -precipitates in IRs, so the coarsening  
415 is much more moderate. Therefore, we expect the creep properties of the superalloy  
416 single crystals after the novel heat treatment protocol to be at least better than 85% of  
417 their cast counterparts after full heat treatment.

418 We note that the one-step annealing strategy to tackle the *RASH* issues reported in  
419 this study can be applicable to a variety of 3D-printed superalloys. As demonstrated in  
420 Figure S6 in Supporting Information, our heat treatment is successful for both laser and  
421 electron beam 3D-printed (or melted) AM3 superalloy single crystals. Laser 3D-printed  
422 SRR99 single crystal has also been tested (Figure S7), with the same outcome when it  
423 comes to accomplishing *RASH*.

424 Finally, the simplicity and efficiency are appealing attributes of our new one-step

425 annealing approach. Compared to the previous “recovery annealing plus standard heat  
426 treatment”, both time and energy consumption are markedly reduced, shortening the  
427 processing chain and reducing the associated costs and wasted parts. More importantly,  
428 the existing stray grains, which are unavoidable due to the nature of 3D printing, do not  
429 coarsen with this sub-solvus solutionizing heat treatment. Since the service temperature  
430 of the single crystal blades will be lower than the solutionizing treatment temperature, we  
431 expect the stray grains to be stable during service. Our method also offers a double  
432 insurance in case there are occasionally some tiny stray grains leftover on the surface  
433 after post-printing machining or etching.

434

#### 435 **4. Conclusion**

436 In summary, we have designed a new heat treatment protocol to release stored  
437 deformation energy, avoid recrystallization, suppress stray grain growth, and homogenize  
438 the chemical and microstructure distribution, all of which are mandated for 3D-printing  
439 manufacture and repair of Ni-based superalloy single crystals. It is remarkable that the  
440 multiple *RASH* requirements are satisfied all at once, via a single-step sub-solvus  
441 solutionizing treatment. Specifically, as ultra-fine dendrite width is generated by the steep  
442 temperature gradient of the 3D-printing process, the necessary distance for diffusion to  
443 cover is greatly reduced, making sub-solvus annealing adequate to accomplish chemical  
444 homogenization. Meanwhile, by setting the heat treatment temperature between the  
445 solvus points of DCs and IRs, the dislocations move freely in the DCs to annihilate fully,  
446 eliminating the stored energy that drives the nucleation of recrystallizing new grains,  
447 while the remaining precipitates in the IRs are able to hinder the motion and interactions

448 of dislocations, the nucleation of recrystallization and the migration of stray grain  
449 boundaries. Meanwhile, our experiments and diffusion analysis successfully singled out  
450 an annealing time that is sufficiently long to homogenize the chemical species  
451 distribution, while as short as possible to limit the  $\gamma'$ -particle coarsening in the  
452 undissolved IRs. Via the construction of a temperature-time “treasure map”, 1270 °C for  
453 15 min is found to be the “sweet spot” for optimal solutionizing to resolve the **RASH**  
454 issues. Such a tactfully crafted heat treatment thus provides a much-needed stepping  
455 stone, for making 3D printing practical to the manufacture and repair of single-crystal  
456 superalloy parts, as exemplified above by the AM3 superalloy single crystals 3D-printed  
457 using either electron or laser beams, including those with leftover surface stray grains.



## 458 **Materials and methods**

459 In this study, three types of superalloy single crystals were investigated, which were  
460 electron beam melted AM3, laser 3D-printed AM3, and laser 3D-printed SRR99. The  
461 nominal compositions of AM3 and SRR99 are Ni-7.82Cr-5.34Co-2.25Mo-4.88W-  
462 6.02Al-1.94Ti-3.49Ta and Ni-8.39Cr-5.01Co-9.47W-5.47Al-2.14Ti-2.92Ta in weight  
463 percentage, respectively. The [001] cast AM3/SRR99 single crystal base-metal boules  
464 were cut into cylinders ~4 mm in height. Electron beam melting with no feedstock was  
465 carried out using a DMAMS Zcomplex3<sup>TM</sup> electron-beam 3D-printing system operated in  
466 10<sup>-3</sup> mbar vacuum. Electron beam of 15 mA was accelerated to 60 keV and focused onto  
467 the base metal surface to form a melt pool. Line scanning was programmed with the  
468 velocity of 10 mm/s to ensure epitaxial dendrite growth in the melt pool. A FZ of about  
469 1500 μm in width and 800 μm in depth was generated. Laser 3D-printing was conducted  
470 on an in-house developed co-axial laser cladding apparatus equipped with a CO<sub>2</sub> laser  
471 with the beam size of 2 mm. The gas atomized superalloy powders with diameters  
472 ranging from 48 to 180 μm with similar composition to the base metal were coaxially  
473 injected at a 11 g/min feeding speed by high-purity Ar gas carrier into the molten pool  
474 formed by the laser beam with a power of 2000 W and 2 mm/s laser scanning speed. The  
475 interlayer spacing is 0.2 mm with a back-and-forth scan path. Therefore the molten  
476 powder solidified on top of the crystal and deposited layer by layer. More detailed  
477 information about the manufacturing process can be found elsewhere <sup>49</sup>.

478 In our novel heat treatment protocol with sub-solvus solutionizing, the electron beam  
479 melted and laser 3D-printed single crystal superalloys were first solutionized at sub-  
480 solvus temperature and then aged at 1100 °C and 870 °C for 6 h and 20 h, respectively.

481 The optimized solutionizing condition was optimized to be 1270 °C for 15 min.  
482 Comparisons were made with identical electron beam melted and laser 3D-printed  
483 samples, heat treated via two other protocols. The “standard heat treatment” was carried  
484 out in a similar manner as the sub-solvus heat treatment, except that the solutionizing  
485 treatment was at 1300 °C. Note that the duration of the so-called standard solutionizing  
486 heat treatment was significantly shorter than that in the superalloy handbook, because of  
487 the narrow dendrite width. The last heat treatment process involves a recovery annealing  
488 step at 1100 °C for 6 h, prior to the standard heat treatment. The heating rate of all  
489 specimens was set at 15 °C/min. All the heat treatment experiments were performed using  
490 a CARBOLITE® RHF 1500 muffle furnace, equipped with an R-type thermocouple  
491 installed at the center of the back inside wall of the furnace chamber to monitor the  
492 temperature. In order to verify the accuracy of the temperature, an AM3 cast superalloy  
493 single crystal was firstly solutionized at 1300 °C for 3 h and then annealed at  
494 temperatures from 1260 to 1280 °C for 30 min, respectively. Observed from the  
495 secondary electron images taken in the scanning electron microscopy, the solvus  
496 temperature of the  $\gamma'$ -precipitates in the DCs was measured to be between 1260 to 1280  
497 °C (Figure S8 in Supporting Information), which is in agreement with the literature<sup>39</sup>. All  
498 samples were air cooled at the rate of approximately 300 °C/min once the heat treatment  
499 was finished.

500 The microstructure was examined under secondary electron mode in a SEM after  
501 etching in 25% phosphoric acid water solution at the voltage of 5 V for 10 s. WDS for  
502 element distribution study was conducted in a SuperProbe JXA-8230 Electron Probe  
503 Microanalyzer at the accelerating voltage of 20 kV. EBSD was carried out after the

504 sample surface was polished electrochemically in 10% perchloric acid alcohol solution at  
505 the voltage of ~30 V for 60 s. All TEM images displayed in this manuscript were taken in  
506 a JEM-2100F field emission electron microscope at the accelerating voltage of 200 kV.  
507  $\mu$ XRD sample was electro-polished in the same way, and then scanned using micro-  
508 focused synchrotron polychromatic X-ray beam on Beamline 12.3.2 at the Advanced  
509 Light Source of Lawrence Berkeley National Laboratory <sup>50</sup>. The collected Laue  
510 diffraction data were processed using a custom-developed software based on the peak  
511 position comparison method to measure <sup>51</sup> and visualize <sup>52</sup> the strain distribution  
512 accurately. Diffraction peaks were searched based on a user-defined peak to background  
513 threshold and fitted using a 2D Lorentzian function to obtain the peak position precisely.  
514 By comparing the angles between experimentally measured diffraction peak positions  
515 with the theoretically calculated ones, the strain tensor and thus the equivalent strain are  
516 calculated. TEM specimens were prepared using the conventional twin-jet  
517 electropolishing.

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528

529 **Author contributions**

530 K.C. designed the project in consultation with E.M. S.L. conducted the experiments and  
531 electron microscopy characterization. W.H. provided the additive-manufactured and re-  
532 melted specimens. Y.L. and S.L. analyzed and interpreted the  $\mu$ XRD data under the  
533 supervision of K.C. and N.T. K.C. and E.M. wrote the paper based on the draft from S.L.  
534 All authors contributed to the discussions of the results.

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