## Title

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New structures on embedded contact homology and applications to low-dimensional topology

by<br>Luya Wang

A dissertation submitted in partial satisfaction of the requirements for the degree of Doctor of Philosophy
in

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Committee in charge:
Professor Michael Hutchings, Chair
Professor Ian Agol
Professor David Nadler

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New structures on embedded contact homology and applications to low-dimensional topology

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Abstract<br>New structures on embedded contact homology and applications to low-dimensional topology<br>by<br>Luya Wang<br>Doctor of Philosophy in Mathematics<br>University of California, Berkeley<br>Professor Michael Hutchings, Chair

Embedded contact homology ( ECH ) is a powerful three-manifold invariant that encodes information about the underlying contact manifold. It has been shown that embedded contact homology is isomorphic to certain versions of Seiberg-Witten Floer homology and Heegaard Floer homology. This thesis highlights the relations of ECH to other three-dimensional Floer theories. In Chapter 2, we use ideas from periodic Floer homology, a cousin of ECH, and techniques from pseudo-Anosov maps to answer a knot detection question in knot Floer homology. This is joint work with Ethan Farber and Braeden Reinoso. In Chapter 3, we show a connected sum formula for ECH by studying pseudo-holomorphic curves in the symplectization of the contact connected sum, without going through the tremendous isomorphisms to SeibergWitten Floer homology or Heegaard Floer homology. Our chain level description of the connected sum complex is useful for studying studying similar formulas for other contact homologies and ECH spectral invariants in the future. In Chapter 3, we give a relative version of ECH for contact three-manifolds with convex sutured boundaries. This generalizes both sutured ECH and Lipshitz's cylindrical reformulation of Heegaard Floer homology, and provides a potential framework for bordered ECH. This is joint work with Julian Chaidez, Oliver Edtmair, Yuan Yao and Ziwen Zhao.

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## Chapter 1

## Introduction

This thesis contains a collection of three papers that the author completed during her PhD:

1. Fixed-point-free pseudo-Anosovs, knot Floer homology and the cinquefoil (joint work with Ethan Farber and Braeden Reinoso)
2. A connected sum formula for embedded contact homology
3. Legendrian embedded contact homology (joint work with Julian Chaidez, Oliver Edtmair, Yuan Yao and Ziwen Zhao)

Each paper occupies a chapter. Each paper is self-contained, but we give some backgrounds and motivations here.

Chapter 2 resolves a question on whether knot Floer homology detects the cinquefoil knot $T(2,5)$. Knot Floer homology (HFK) defined by Ozsváth-Szabó and Rasmussen $[98,105]$ is a powerful knot invariant coming from the package of Heegaard Floer homology. Previously, HFK was known to detect the unknot [97], the trefoil knot and the figure-eight knot [40]. However, establishing such a result for the genus-two torus knot $T(2,5)$ has been very difficult, for example, because there are infinitely many genus-two hyperbolic fibered knot with the same Alexander polynomial as that of $T(2,5)$. Our approach is inspired by an isomorphism between a certain grading of HFK and the degree $d=1$ part of the periodic Floer homology, which is a cousin of ECH [41]. The result we need from [41] can also be recovered from [3]. The essential strategy of ours is to analyze dynamical, in particular pseudo-Anosov, properties of a hyperbolic genus-two fibered knot. We show that any genus-two fibered knot in $S^{3}$ satisfying a mild condition must contain a fixed point in the interior of its fibering surface.

Chapter 3 and 4 discuss foundational aspects of embedded contact homology (ECH). Chapter 3 gives a connected sum formula for embedded contact homology by studying how pseudo-holomorphic curves behave in the symplectization of a contact connected sum. Even though the similar formulas are known in Seiberg-Witten Floer and Heegaard Floer, our proof does not go through the isomorphisms between any of ECH, Seiberg-Witten Floer and Heegaard Floer, except for a direct limit argument when passing to homology $[109,110,111,112,113,19,24,26,22]$. See also $[76,77,78$, $75,79]$. Our techniques can also be used to study gluing formulas along higher-genus surfaces that admit certain pseudo-holomorphic foliations in their symplectization, e.g. mixed tori [16], which will be the subjects for future works. The chain level description of the connected sum ECH complex also gives information regarding ECH spectral invariants, which will be studied in future projects. Chapter 4 constructs a relative version of embedded contact homology for convex sutured contact manifolds that we call Legendrian ECH. We develop intersection theory for pseudo-holomorphic curves with boundaries on Lagrangians that are symplectizations of the boundary Legendrians chosen a priori. This allows us to define a relative ECH index and obtain a relative adjunction formula. We further show analogous properties of Legendrian ECH to the ordinary ECH. In particular, we have that $\partial^{2}=0$. The Legendrian ECH generalizes both sutured ECH and Lipshitz's reformulation of Heegaard Floer homology and provides a potential framework for bordered ECH.

## Chapter 2

# Fixed-point-free pseudo-Anosovs, knot Floer homology and the cinquefoil 

Ethan Farber, Braeden Reinoso and Luya Wang

### 2.1 Introduction

Recent developments in Heegaard Floer homology have highlighted intimate connections between link homology theories and the dynamics of surface diffeomorphisms (see e.g. [3], [95], [96], [41]). The aim of this paper is to use tools which may be familiar to some dynamicists, in order to study a particular open question in Heegaard Floer homology: whether knot Floer homology detects the torus knot $T(2,5)$. To that end, we answer this question in the affirmative:

Theorem A. If $\widehat{H F K}(K ; \mathbb{Q}) \cong \widehat{H F K}(T(2,5) ; \mathbb{Q})$ as bi-graded vector spaces, then $K=T(2,5)$. In particular, $T(2,5)$ is the only genus-two L-space knot in $S^{3}$.

This is the first knot Floer detection result for a knot of genus two. Prior detection results rely on a classification of fibered knots with genus at most one [97, 40]. But, there are infinitely many genus two, hyperbolic, fibered knots with the same Alexander polynomial as that of $T(2,5)$ [92], indicating that close attention must be paid to the structure of the fibration. In this vein, our proof completes a strategy outlined by Baldwin-Hu-Sivek in [3], which uses connections between knot Floer homology and symplectic Floer homology, to reduce the question to a problem about fixed points of pseudo-Anosov maps. Our first key result is a solution to this problem.

Definition 2.1.1. Let $K \subset Y$ be a hyperbolic, fibered knot in a 3-manifold $Y$. The knot $K$ specifies an open book decomposition $(S, h)$ for $Y$, where $h: S \rightarrow S$ is freely isotopic to a pseudo-Anosov $\psi$ on $S$. We say that $K$ is fixed point free if $\psi$ has no fixed points in the interior of $S$.

Theorem B. Let $K$ be a hyperbolic, genus-two, fibered knot in $S^{3}$. If the fractional Dehn twist coefficient $c(K) \neq 0$, then $K$ is not fixed point free.

Pseudo-Anosov maps and fractional Dehn twist coefficients are defined in Section 2.2. Theorem A follows immediately from Theorem B and the work of Baldwin- $\mathrm{Hu}-$ Sivek ([3], Theorem 3.5). We would like to emphasize that the proof of Theorem B is completely geometric in nature, and after the introduction we will only make passing references to Floer homology theories throughout the paper. An unfamiliar reader need not have expertise in any link homology theory to understand the proof of Theorem B.

The proof of Theorem B is broken down into two smaller theorems (Theorems B1 and B2), based on cases for the singularity type of $\psi$. An outline for the proof is given in subsection 2.1. We will now discuss various applications of our techniques and results.

## Applications to train tracks and the dilatation spectrum

One of the central tools in the proof of Theorem B is the theory of train tracks for pseudo-Anosov braids, including a theory of "tight splitting" developed in Section 2.5. We believe the techniques we use are broadly applicable elsewhere within surface dynamics. For example:

Theorem C (cf. Theorem 2.4.2). Let $\psi$ be a pseudo-Anosov on the genus-two surface with one boundary component, with singularity type $\left(4 ; \emptyset ; 3^{2}\right)$. Then, $\psi$ is conjugate to a pseudo-Anosov carried by the train track depicted on the bottom left in Figure 2.6. A similar statement holds for the closed genus-two surface.

This result suggests a strategy for systematically studying the set of dilatations of pseudo-Anosovs in genus two. Indeed, Theorem C reduces the study of dilatations in the stratum $\left(4 ; \emptyset ; 3^{2}\right)$ to the study of a special collection of maps on a single graph. One would hope for the development of a kneading theory generalizing the classical theory of Milnor-Thurston for interval maps (cf. [MT]). Applying such a theory to a small list of tracks in each of the other strata in genus two would provide a much clearer understanding of the full dilatation spectrum. This line of study was suggested to us by Chenxi Wu. Another step in this direction is the following result:

Theorem D. Let $\psi$ be a pseudo-Anosov on the punctured disk with at least one $k$-pronged singularity away from the boundary with $k \geq 2$. Then $\psi$ is carried by a standardly embedded train track $\tau$ with no joints.

See Definition 2.5.1 for the definition of a standardly embedded train track. A loop switch of a standardly embedded train track $\tau$ is a switch at a loop surrounding a 1-prong singularity (cf. Definition 2.5.16), and a joint is a loop switch that is incident to more than one expanding edge (cf. Definition 2.5.24). The track in Theorem C is the lift of a joint-less track on the punctured disk, so Theorem C may be seen as a specific case of Theorem D. See subsection 2.5 for the proofs of Theorems C and D.

## Applications to the Floer homology of branched covers

For a knot $K \subset S^{3}$, let $\Sigma_{n}(K)$ denote the $n$-fold cyclic cover of $S^{3}$ branched along $K$. There has been much interest recently in the Floer homology of $\Sigma_{n}(K)$ in terms of $K$. For example, Boileau-Boyer-Gordon have studied extensively in [10] and [11] the set of all integers $n \geq 2$ such that $\Sigma_{n}(K)$ is an L-space (see also e.g. [68] and [103]). One question that has persisted in this area is the following:

Question 2.1.2 (Boileau-Boyer-Gordon, Moore). Can $\Sigma_{n}(K)$ be an L-space for $K$ a hyperbolic L-space knot?

Combining Theorem A with ([10], Corollary 1.4) yields the following complete answer to this question for $n>2$ :

Corollary 2.1.3. If $K$ is an L-space knot and $\Sigma_{n}(K)$ is an L-space for some $n>2$, then $K$ is either $T(2,3)$ or $T(2,5)$. In particular, $K$ is not hyperbolic.

## Applications to instanton Floer homology and Dehn surgery

For a 3-manifold $Y$, let $R(Y)=\operatorname{Hom}\left(\pi_{1}(Y), S U(2)\right)$ denote the $S U(2)$-representation variety. We say that a 3-manifold $Y$ is $S U(2)$-abelian if $R(Y)$ contains no irreducibles. The name is motivated by the fact that $Y$ is $S U(2)$-abelian if and only if every $\rho \in R(Y)$ has abelian image.

Following work initiated by Kronheimer-Mrowka in their proof of the Property P conjecture [72], Baldwin-Li-Sivek-Ye [6], Baldwin-Sivek [4], and KronheimerMrowka [KM2] proved that $r$-surgery $S_{r}^{3}(K)$ on a nontrivial knot $K \subset S^{3}$ is not $S U(2)$-abelian for all slopes $r \in[0,3] \cup[4,5)$ with prime power numerator, and for some additional slopes $r \in[3,4)$.

The key theory which facilitates most of these results is the instanton Floer homology of the surgered manifold $S_{r}^{3}(K)$ (and related techniques arising from this theory, as in [KM2]). Combining Theorem B with ([6], Proposition 2.4) allows us to prove an analogue of Theorem A for instanton Floer homology:

Corollary 2.1.4. The cinquefoil $T(2,5)$ is the only genus-two instanton L-space knot, i.e. the only genus-two knot $K$ for which $\operatorname{dim} I^{\#}\left(S_{r}^{3}(K)\right)=\left|H_{1}\left(S_{r}^{3}(K)\right)\right|$ for some $r>0$.

Now, as described in ([6], Section 1.3), Corollary 2.1.4 completes the set of slopes $r$ for which $S_{r}^{3}(K)$ is not $S U(2)$-abelian, to all rational numbers $r \in[0,5)$ with prime power numerator:

Corollary 2.1.5. Let $K \subset S^{3}$ be a nontrivial knot, and $r \in[0,5)$ a rational number with prime power numerator. Then, $S_{r}^{3}(K)$ is not $S U(2)$-abelian.

Remark 2.1.6. Note that $S_{r}^{3}(K)$ may in general be $S U(2)$-abelian for $r \geq 5$, as $S_{5}^{3}(T(2,3))$ is the lens space $L(5,1)$, which has abelian fundamental group. However, Baldwin-Li-Sivek-Ye in [6] have extended the slopes for which $S_{r}^{3}(K)$ is not $S U(2)$ abelian to some additional $r \in(5,7)$. It is an open question whether $S_{r}^{3}(K)$ is $S U(2)$-abelian for all rational numbers $r \in[0,5)$, though it is known to be true for $r \in[0,2]$ by work of Kronheimer-Mrowka in [KM2].

## Applications to Khovanov homology

In [3], Baldwin-Hu-Sivek proved that Khovanov homology (with coefficients in $\mathbb{Z} / 2 \mathbb{Z}$ ) detects the cinquefoil $T(2,5)$. Combining Theorem A with previous work of Baldwin-Dowlin-Levine-Lidman-Sazdanovic ([5], Corollary 2), we can improve Baldwin-HuSivek's result from $\mathbb{Z} / 2 \mathbb{Z}$-coefficients to $\mathbb{Q}$-coefficients:

Corollary 2.1.7. If $K h(K ; \mathbb{Q}) \cong K h(T(2,5) ; \mathbb{Q})$ as bi-graded $\mathbb{Q}$-vector spaces, then $K=T(2,5)$.

We also obtain detection results in annular Khovanov homology. One may think of $T(2,5)$ as the lift of the braid axis for the 5 -braid $B=\sigma_{1} \sigma_{2} \sigma_{3} \sigma_{4}$ in $S^{3}$ seen as the double-branched cover over $\widehat{B}$ under the Birman-Hilden correspondence (see subsection 2.2 for background). From this perspective, we can adapt techniques of Binns-Martin in ([8], Theorems 10.2, 10.4, 10.7) to prove that annular Khovanov homology detects the aforementioned braid closure:

Corollary 2.1.8. Let $L \subset A \times I$ be an annular link with $A K h(L ; \mathbb{Q}) \cong A K h(\widehat{B} ; \mathbb{Q})$. Then, $L$ is isotopic to $\widehat{B}$ in $A \times I$.

The proof of this corollary is almost identical to those of the analogous results proved by Binns-Martin. So, we omit the proof in the present work, and instead refer the reader to [8].

## Outline of the proof of Theorem B

Let $K$ be a hyperbolic, genus-two, fibered knot in $S^{3}$ with associated open book decomposition $(S, h)$ and pseudo-Anosov representative $\psi_{h}$. If $c(K) \neq 0$ and $K$ is fixed point free, it follows from work of Baldwin-Hu-Sivek (see Theorem 2.2.5) that:

- $\psi_{h}$ has singularity type either:

Case 1: $(6 ; \emptyset ; \emptyset)$, or
Case 2: $\left(4 ; \emptyset ; 3^{2}\right)$

- $h$ is the lift of a 5 -braid $\beta$ with unknotted braid closure $\widehat{\beta}$
- The pseudo-Anosov representative $\psi_{\beta}$ of the braid $\beta$ has singularity type either:

Case 1: $\left(3 ; 1^{5} ; \emptyset\right)$
Case 2: $\left(2 ; 1^{5} ; 3\right)$
For our conventions of the singularity types, see Section 2.2. Cases 1 and 2 are mutually exclusive, and we will use notation $K, h, \beta, \psi_{h}, \psi_{\beta}$ as above for the rest of this outline, in either case.

Case 1 is dealt with in Section 2.3. In this case, Masur-Smillie proved in [89] that the foliations preserved by $\psi_{h}$ are orientable, so that the dilatation of $\psi_{h}$ is a root of the Alexander polynomial $\Delta_{K}(t)$. This is a special fact about the stratum $(6 ; \emptyset ; \emptyset)$ in genus two. Using the Lefschetz fixed point theorem and basic facts about Alexander polynomials of fibered knots in $S^{3}$, we completely determine the Alexander polynomial of $K$ and conclude that the dilatation of $\psi_{h}$ coincides with the minimal dilatation $\lambda_{2}$ for genus-two pseudo-Anosovs. Work of Lanneau-Thifeault in [81] further implies that $\psi_{h}$ is the almost unique genus-two pseudo-Anosov realizing $\lambda_{2}$ as its dilatation. It follows that $\beta$ is (up to inverse and composing with the hyperelliptic involution) conjugate to the dilatation-minimizing 5 -braid $\alpha$ from [44] within the mapping class group of the punctured sphere. We then show that no braid which is conjugate to $\alpha$ within the spherical mapping class group has unknotted closure.

Case 2 is harder: we perform our analysis using a splitting argument and a careful combinatorial analysis of train track maps. In this case, we focus on the braid $\beta$ and its pseudo-Anosov representative $\psi_{\beta}$. We show (Theorem 2.4.2) that any pseudoAnosov $\psi_{\beta}$ on the five-punctured disk in the stratum $\left(2 ; 1^{5} ; 3\right)$ is carried by a single

## CHAPTER 2. FIXED-POINT-FREE PSEUDO-ANOSOVS, KNOT FLOER

 HOMOLOGY AND THE CINQUEFOILcanonical train track $\tau$ (cf. also Theorems C and D). To prove this result, we develop a theory of "tight splitting" in Section 2.5, which allows us to study the splitting of all train track maps on any given track within the stratum.

In subsection 2.4, we find a collection of braids $\beta_{n}$ inducing special train track maps $f_{n}: \tau \rightarrow \tau$ on the distinguished track $\tau$ from the previous paragraph. We show, in subsection 2.4 , that $f_{n}$ are the only maps on $\tau$ which could lift to train track maps for fixed-point-free pseudo-Anosovs $\psi_{h}$ in the cover. So, if $\beta$ is any braid which lifts to a map $h$ with fixed-point-free pseudo-Anosov representative $\psi_{h}$, then $\beta$ is conjugate (in the mapping class group of the punctured sphere) to one of $\beta_{n}$. A similar argument as in Case 1 shows that no braid conjugate to $\beta_{n}$ has unknotted closure.

### 2.2 Background

## Mapping classes and fractional Dehn twists

Let $S=S_{g, n}^{r}$ be a compact surface of genus $g$ with $n$ marked points and $r$ boundary components. The mapping class group of $S$ is the group $\operatorname{Mod}(S)$ of isotopy classes of homeomorphisms $h: S \rightarrow S$ which fix $\partial S$ point-wise, and permute the marked points of $S$, where the isotopies fix all boundary components and marked points. The symmetric mapping class group of $S$ is the analogous group $\operatorname{SMod}(S)$ obtained by additionally requiring that the homeomorphisms commute with the hyperelliptic involution $\iota: S \rightarrow S$, i.e. $h \circ \iota=\iota \circ h$.

Definition 2.2.1. A pseudo-Anosov is a homeomorphism $\psi: S \rightarrow S$ preserving a pair of transverse singular measured foliations $\left(\mathcal{F}^{u}, \mu^{u}\right)$ and $\left(\mathcal{F}^{s}, \mu^{s}\right)$ such that

$$
\psi \cdot\left(\mathcal{F}^{u}, \mu^{u}\right)=\left(\mathcal{F}^{u}, \lambda \mu^{u}\right), \text { and } \psi \cdot\left(\mathcal{F}^{s}, \mu^{s}\right)=\left(\mathcal{F}^{s}, \lambda^{-1} \mu^{s}\right)
$$

for some fixed real number $\lambda>1$, called the dilatation of $\psi$.
We further require that each singularity $p$ of $\mathcal{F}^{u}$ or $\mathcal{F}^{s}$ is a " $k$-pronged saddle," as in Figure 2.1, where $k \geq 3$ for $p$ in the interior of $S$, or $k \geq 1$ for $p$ a marked point or puncture. Along $\partial S$, the singular points must all have a neighborhood of the form shown on the bottom left of Figure 2.1.

By a $k$-prong boundary singularity or a $k$-prong singularity on the boundary, we mean that $\psi$ has $k$ singular points on a particular boundary component of $S$. The singularity type of $\psi$ is the tuple $\left(b_{1}, \ldots, b_{r} ; m_{1}, \ldots, m_{n} ; k_{1}, \ldots, k_{s}\right)$ where the $i^{\text {th }}$ boundary component has $b_{i}$-prongs, the $i^{\text {th }}$ puncture or marked point has $m_{i}$-prongs, and the $i^{\text {th }}$ interior singularity has $k_{i}$ prongs. We will use $\emptyset$ if $\psi$ has no boundary


Figure 2.1: Top (left to right): a 3 -pronged saddle, a 4 -pronged saddle, and a 1pronged saddle at a marked point. Bottom (left to right): the neighborhood of a boundary singularity, and a 3-pronged boundary which permutes the first prong to the second.
(or if $\psi$ has no marked points or punctures, or interior singularities), and we will use an exponent to denote a repeated number of prongs. For example, the tuple $\left(3 ; 1^{5} ; \emptyset\right)$ indicates that $\psi$ has a 3 -pronged boundary, five 1-pronged marked points or punctures, and no interior singularities. A stratum on $S$ is the collection of all pseudo-Anosovs on $S$ with a given singularity type.

Theorem 2.2.2 (Nielsen-Thurston classification). Any element $h \in \operatorname{Mod}(S)$ is freely-isotopic rel. punctures (i.e. isotopic through homeomorphisms which fix the punctures, but may rotate a boundary component) to a representative $\psi$ with at least one of the following properties:
(1) $\psi^{n}=$ id for some $n$,
(2) $\psi$ preserves the isotopy class of a multicurve $C$ on $S$, or
(3) $\psi$ is pseudo-Anosov. This case is disjoint from the previous two.

We say that such a $\psi$ satisfying one of the properties (1)-(3) is geometric, and in case (3), we refer to $\psi$ as the pseudo-Anosov representative of $h$. The representative $\psi$ is unique for any such $h$, although when $S$ has non-empty boundary, $\psi$ will never be isotopic rel. boundary to an element of $\operatorname{Mod}(S)$, as $\psi$ may rotate $\partial S$.

The fractional Dehn twist coefficient $c(h)=n+m / k$ is a rational number which measures the rotation of $\psi$ along $\partial S$. Here, $n$ is an integer measuring the number of full-rotations of $h$ along $\partial S ; k$ is the number of prongs of $\psi$ on $\partial S$; and $\psi$ cyclically permutes the endpoint of the first prong to that of the $(m+1)^{t h}$ prong. For example, in the bottom right of Figure 2.1, the "fractional part" of $c(h)$ is $1 / 3$, where $h \in$ $\operatorname{Mod}(S)$ has pseudo-Anosov representative $\psi$. In particular $c(h) \in \mathbb{Z}$ if and only if the rotation number of $\psi$ along the boundary is zero, and when $\psi$ has a single boundary prong this is always the case.

The fractional Dehn twist coefficient can be extended analogously to braids, thought of as elements of $\operatorname{Mod}\left(S_{0, n}^{1}\right)$ for some $n>1$. In this case, we denote it by $c(\beta)$, where $\beta$ is a braid. When $K$ is a fibered knot with open book decomposition $(S, h)$, we define the fractional Dehn twist coefficient $c(K)$ to be that of $h$, i.e. $c(K):=c(h)$. It is crucial to note that, in general, $c(K)$ is not the same as $c(\beta)$ for $\beta$ a braid representative of $K$. The following theorem details a few well-known properties which we will make use of in this paper:

Theorem 2.2.3 ([49],[104],[69]). Let $h: S \rightarrow S$ be a mapping class with $\partial S$ connected, and let $\beta: S_{0, n}^{1} \rightarrow S_{0, n}^{1}$ be a braid.

- $c(h)$ and $c(\beta)$ are preserved under conjugation.
- $c\left(D_{\partial S}^{m} \circ h^{k}\right)=m+k c(h)$ for any $k, m \in \mathbb{Z}$, where $D_{\partial S}$ is a Dehn twist along $\partial S$.
- $c\left(\Delta^{2 m} \beta^{k}\right)=m+k c(\beta)$ for any $k, m \in \mathbb{Z}$, where $\Delta^{2}=\left(\sigma_{1} \ldots \sigma_{n-1}\right)^{n}$.
- If $\beta$ is $\sigma_{1}$-positive (i.e. $\beta$ can be written with only positive powers of $\sigma_{1}$ ) then $c(\beta) \geq 0$.
- If $\beta$ is a positive pseudo-Anosov braid, then $c(\beta)>0$.

See [115] for more details regarding the Nielsen-Thurston classification, and [49] or [70] for more details regarding fractional Dehn twist coefficients.

## Knots, braids, and the Birman-Hilden correspondence

We may use branched coverings to understand relationships between mapping class groups of different surfaces. The Birman-Hilden correspondence is a key tool for this study. For our purposes, the correspondence will help us study the mapping class groups of $S_{2}^{1}$ and $S_{2}$ seen as two-fold branched covers over the disk $S_{0,5}^{1}$ and sphere $S_{0,6}$, respectively, via the hyperelliptic involution $\iota$. Specifically, there is a diagram:


Here, the map $\Theta_{2}^{1}$ is an isomorphism, and the map $\Theta_{2}$ is surjective with $\operatorname{ker}\left(\Theta_{2}\right)=\langle\iota\rangle$. The cap-off maps are both given by setting a full-twist about $\partial S$ to 0 (geometrically, one may think about capping $\partial S_{2}^{1}$ with a marked disk), and the "forget" map forgets about the marked point on $S_{2,1}$. The cap-off map forgets about the "integer part" of the fractional Dehn twist coefficient but preserves the "fractional part" in each case (i.e. the rotation number of $\psi$ along $\partial S$ is the rotation number of the capped-off map $\widehat{\psi}$ at the marked point in the capping disk). See [38] for more details on the maps involved in this diagram.

Note that the braid $\Delta^{2}=\left(\sigma_{1} \sigma_{2} \sigma_{3} \sigma_{4}\right)^{5} \in \operatorname{Mod}\left(S_{0,5}^{1}\right)$ is isotopic to a full-twist about $\partial S_{0,5}^{1}$. It follows that, given a spherical mapping class $f \in \operatorname{Mod}\left(S_{0,6}\right)$, there are $\mathbb{Z}$-many lifts of $f$ to braids

$$
\ldots \Delta^{-4} \beta, \Delta^{-2} \beta, \beta, \Delta^{2} \beta, \Delta^{4} \beta \ldots \in \operatorname{Mod}\left(S_{0,5}^{1}\right)
$$

which are all related by powers of $\Delta^{2}$, and are distinguished by their fractional Dehn twist coefficient $c(\beta)$, as in Theorem 2.2 .3 . For any such $f$, only finitely many such $\beta$ may have braid closure $\widehat{\beta}$ an unknot. This may be seen, for example, from the following theorem of Ito-Kawamuro, which will be a key tool in this paper:
Theorem 2.2.4 (Ito-Kawamuro). If the braid closure $\widehat{\beta}$ is an unknot, then $|c(\beta)|<$ 1.

Moreover, one may check that $\Delta^{2}$ lifts to an element of $\operatorname{SMod}\left(S_{2}^{1}\right)$ which squares to a full twist about $\partial S_{2}^{1}$. This implies a very useful fact: the lift of $\Delta^{2}$ is freely isotopic to the hyperelliptic involution $\iota$. One may see this in a number of different ways- for example, by noting that the full twist about $\partial S_{2}^{1}$ is freely isotopic to the identity, and that the lift of $\Delta^{2}$ acts on $H_{1}\left(S_{2}^{1}\right)$ by -id.

A genus-two, hyperbolic, fibered knot $K \subset Y$ yields an open book decomposition $(S, h)$ for $Y$, where $S=S_{2}^{1}$, and $h \in \operatorname{Mod}\left(S_{2}^{1}\right)$ is freely isotopic to a pseudo-Anosov $\psi_{h}: S \rightarrow S$. If $\psi_{h}$ is fixed point free in the interior of $S$, then $h$ is symmetric (i.e. represents an element of $\operatorname{SMod}\left(S_{2}^{1}\right)$ ) by [3], so we may think of $h$ as the lift of a 5 -braid $\beta \in \operatorname{Mod}\left(S_{0,5}^{1}\right)$. From the perspective of 3 -manifolds, this means that $Y$ is the double cover of $S^{3}$ branched along the braid closure $\widehat{\beta}$. This implies, for example, that if $Y=S^{3}$ then $\widehat{\beta}$ is the unknot. From the perspective of knots, the original fibered knot $K=\partial S_{2}^{1} \subset Y$ is the lift of the braid axis $\partial S_{0,5}^{1} \subset S^{3}$ in the cover.

Baldwin-Hu-Sivek in [3] additionally computed the singularity type of $\psi_{h}$ (and originally observed several of the facts mentioned in the previous paragraph). These observations are recorded as the following result, which is the starting point for many of the ideas in this paper:

Theorem 2.2.5 (Baldwin-Hu-Sivek). Let $K \subset S^{3}$ be a hyperbolic, genus-two, fibered knot with associated open book decomposition $(S, h)$ satisfying $c(h) \neq 0$. If the pseudo-Anosov representative $\psi_{h}$ is fixed point free in the interior of $S$, then:

- $\psi_{h}$ has singularity type either:

$$
\text { Case 1: }(6 ; \emptyset ; \emptyset) \text {, or }
$$

Case 2: $\left(4 ; \emptyset ; 3^{2}\right)$

- $h$ is the lift of a 5 -braid $\beta$ under the Birman-Hilden correspondence, and
- as a knot in $S^{3}$, the braid closure $\widehat{\beta}$ is the unknot.

Theorem 2.2.5 yields strong constraints on the braid $\beta$. For example, because $h$ has pseudo-Anosov mapping class, we know that $\beta$ does, too, and we may determine the singularity type of its pseudo-Anosov representative $\psi_{\beta}$ from that of $\psi_{h}$, by appealing to Lemma 3.7 of [3]. If $\psi_{h}$ has singularity type $(6 ; \emptyset ; \emptyset)$ then $\psi_{\beta}$ has singularity type $\left(3 ; 1^{5} ; \emptyset\right)$. And, if $\psi_{h}$ has singularity type $\left(4 ; \emptyset ; 3^{2}\right)$ then $\psi_{\beta}$ has singularity type $\left(2 ; 1^{5} ; 3\right)$.

## Fibered surfaces and train tracks

For the remainder of the paper, we will denote by $S^{\prime}$ the surface $S$ with its marked points deleted, and by $\widehat{S}$ the closed surface obtained by capping-off the boundary components of $S$ with disks and marking a point in the interior of each disk. We will also assume that the surface $S^{\prime}$ has negative Euler characteristic.

In [7], Bestvina and Handel prove that one may associate to any geometric $\psi$ a fibered surface $F \subseteq S^{\prime}$. This fibered surface is decomposed into strips and junctions, where the strips are foliated by intervals, i.e. leaves. See Figure 2.2. Together, the leaves and junctions of $F$ are called decomposition elements, and $\psi(F) \subseteq F$, sending decomposition elements into decomposition elements and, in particular, junctions into junctions. Collapsing each decomposition element to a point produces a graph $G$ with a graph map $g: G \rightarrow G$. The vertices of $G$ correspond to the junctions of $F$, and the edges of $G$ correspond to the strips of $F$.

Roughly speaking, a graph map $g$ is efficient if the image of no edge backtracks under any power of $g$. After adjusting $F$ so that $g$ is efficient, Bestvina and Handel


Figure 2.2: Left: the fibered surface for some geometric $\psi$ on $S_{0,5}^{1}$. The shaded regions are the junctions, and the striped bands connecting them are the strips. Right: following Bestvina-Handel, one inserts additional, "infinitesimal" edges into the junctions. These will also be inserted into the graph $G$ that one obtains by collapsing all of the decomposition elements. Their inclusion will produce the smooth analog $\tau$ of $G$, which is a train track. See Figure 2.3 below.
construct a "smoothed" version of $G$ as follows. Within each junction $J \subseteq F$, one inserts additional edges that smoothly connect the strips of $F$ and encode how images of strips under $\psi$ pass through $J$.

In this way, we obtain a new graph $\tau$ smoothly embedded in the punctured surface $S^{\prime}$, called a train track. At each vertex $s$ of $\tau$, called a switch, there is a well-defined tangent line. Two arcs $a, b$ of $\tau$ are tangent at $s$ if $a(0)=b(0)=s$ and $a^{\prime}(0)=b^{\prime}(0)$. A cusp is the data of a pair $(a, b)$ of adjacent arcs tangent at $s$. See Figure 2.3 for an example.

The following proposition appears as Proposition 3.3.5 in [7].
Proposition 2.2.6. Suppose $\psi$ is pseudo-Anosov. Then in the capped surface, each component of $\widehat{S} \backslash \tau$ is either:

1. a disk with $k \geq 3$ cusps on its boundary, or
2. a disk with a single marked point in its interior and $k \geq 1$ cusps on its boundary.


Figure 2.3: A train track $\tau$ on the five-punctured disk $D_{5}$. The components of the complement of $\tau$ consist of: five once-punctured monogons, i.e. disks with a single boundary cusp; a trigon, i.e. a disk with three boundary cusps; and an exterior oncepunctured bigon. Pseudo-Anosovs carried by this track lie in the stratum $\left(2 ; 1^{5} ; 3\right)$.

Remark 2.2.7. From this perspective, the cusps of a component $C \subset \widehat{S} \backslash \tau$ correspond precisely to the prongs of a singularity $p \in C$ of the invariant foliations of $\psi$.

Definition 2.2.8. An edge path in $\tau$ is a map $e: I \rightarrow \tau$ such that $e(0)$ and $e(1)$ are switches. A train path is an edge path that is also a smooth immersion. The length of a train path $e$ is defined to be the number of edges traversed by $e(I)$, counting with multiplicity. Let $e(I)=e_{1} \cdots e_{k}$ denote a train path whose directed image traverses first $e_{1}$, then $e_{2}$, etc. See Figure 2.4 for examples.

Definition 2.2.9. A train track map is a map $f: \tau \rightarrow \tau$ such that for any train path $g: I \rightarrow \tau$ the composition $f \circ g: I \rightarrow \tau$ is a train path.

Remark 2.2.10. Note that if $f: \tau \rightarrow \tau$ is a train track map, then $f(e)$ is a train path for each edge $e$ of $\tau$. Indeed, from this it follows that $f^{k}(e)$ is a train path for each $k \geq 1$, and hence $f^{k}$ is a train track map for all $k \geq 1$.

The map $\psi: F \rightarrow F$, or equivalently the graph map $g: G \rightarrow G$ corresponding to $\psi$ and $F$, defines a map $f: \tau \rightarrow \tau$. The fact that $g$ is efficient implies that $f$ is a train track map. In this case, we say that the train track $\tau$ carries the map $\psi$, and the map $\psi$ induces the train track map $f$. The data of a geometric map will then be a triple $(\tau, \psi, f)$ in a commutative diagram:


Figure 2.4: Three edge-paths on the train track $\tau$ from Figure 2.3. Left: an edge path of length 7 which is not a train path, since it makes several sharp turns. Middle: a train path of length 7 which can be "pushed off" of $\tau$ into a small neighborhood so that it does not intersect itself. Right: a train path of length 6 which cannot be "pushed off" of $\tau$ so that it becomes injective.


Here, one should imagine $\tau$ being mapped forward by $\psi$ into $F$, meeting the leaves of $F$ transversely. The map $f$ is then defined by collapsing each leaf of $F$ to a point, while inside each junction the arcs of $\psi(\tau)$ are collapsed onto the appropriate edges of $\tau$. See Figure 2.5.

The edges of $G$ (other than those loops peripheral to marked points/punctures of $S)$ are in bijection with a subset of the edges of $\tau$, which we call the real edges. All other edges of $\tau$ are infinitesimal. In particular, all edges of $\tau$ contained in a junction of $F$ are infinitesimal. Enumerate the edges of $\tau$ so that $e_{1}, \ldots, e_{k}$ are the real edges and $e_{k+1}, \ldots, e_{n}$ are the infinitesimal edges. For each pair $(i, j)$ with $1 \leq i, j, \leq n$ define the integer

$$
m_{i, j}=\text { the number of times the train path } f\left(e_{j}\right) \text { traverses } e_{i} .
$$

The extended transition matrix of $f$ is the matrix $\widetilde{M}$ whose $(i, j)$-entry is the integer $m_{i, j}$. The transition matrix of $f$ is the submatrix $M \subset \widetilde{M}$ recording the transitions between real edges of $\tau$ : in other words,

$$
M=\left(m_{i, j}\right) \text { where } 1 \leq i, j \leq k
$$

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Figure 2.5: Top row: the train track $\tau$ from Figure 2.3 and the action of a pseudoAnosov $\psi$ that it carries. The real edges of $\tau$ are labeled $e_{1}, \ldots, e_{5}$. The shaded regions on the right denote the neighborhoods that deformation retract onto these edges. Bottom three rows: The action of $f=($ collapse $\circ \psi)$ on each edge of $\tau$, depicted separately.

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 HOMOLOGY AND THE CINQUEFOILDefinition 2.2.11. Let $M$ be a square matrix whose entries are non-negative integers. We say that $M$ is Perron-Frobenius if the entries of $M^{N}$ are strictly positive, for some power $N$. In this case, the Perron-Frobenius theorem states that the eigenvalue of $M$ of largest absolute value is in fact real, simple, and has an eigenvector all of whose entries are positive. We call this eigenvalue the Perron-Frobenius eigenvalue, and we say that such an eigenvector is positive.

The next theorem follows from work of Bestvina-Handel in [7].
Theorem 2.2.12. Let $(\tau, \psi, f)$ be the data of a geometric map, where $\tau$ satisfies the conclusion of Proposition 2.2.6. Let $M$ be the transition matrix of $f$. Then

$$
\psi \text { is pseudo-Anosov } \Longleftrightarrow M \text { is Perron-Frobenius. }
$$

The Perron-Frobenius eigenvalue $\lambda$ of $M$ is called the dilatation of $\psi$. There is a unique right $\lambda$-eigenvector $w$ of $M$, up to scale, and its entries $w_{i}$ for $i=1 \ldots, k$ define transverse weights on the real edges $e_{i}$ of $M$.

Remark 2.2.13. If a train track map $F$ is induced by some pseudo-Anosov $\psi$, then such a $\psi$ is unique up to conjugacy in $\operatorname{Mod}(\widehat{S})$. Indeed, Bestvina-Handel in [7] provide an algorithm to determine the measured foliations preserved by $\psi$. This will be a crucial idea in Section 2.4.

## Lifted train track maps and fixed points

Let $(\tau, \psi, f)$ be the data of a pseudo-Anosov on $S=S_{2}^{1}$. One of the key ideas in this paper is to use the transition matrix $M$ of $f$ to study fixed points of $\psi$ combinatorially. Our main tool to carry out this approach is the following theorem, which follows from work of Los in [88] and independently Cotton-Clay in [29]:

Theorem 2.2.14 (Los, Cotton-Clay). If $\psi$ is fixed point free in the interior of $S$, then $\operatorname{tr}(M)=0$.

Our goal will be to use Theorem 2.2 .14 to restrict the possible train tracks $\tau$ and maps $f: \tau \rightarrow \tau$ for a fixed-point-free $\psi$. In the case of genus-two, hyperbolic, fibered knots in $S^{3}$, the possible types of train tracks $\tau$ are already highly restricted by Theorem 2.2.5 and Remark 2.2.7.

Suppose we start with a genus-two, hyperbolic, fibered knot $K$ with open book decomposition $(S, h)$, where $h$ is the lift of a 5 -braid $\beta$ with pseudo-Anosov data $\left(\psi_{\beta}, \tau, f\right)$ on the disk $S_{0,5}^{1}$. We may lift $\tau$ to a train track $\widetilde{\tau}$ on $S$ which carries the pseudo-Anosov representative $\psi_{h}$ of $h$. As a graph, $\widetilde{\tau}$ is constructed by gluing


Figure 2.6: Lifting $\left(\psi_{\beta}, \tau, f\right)$ to $\left(\psi_{h}, \widetilde{\tau}, \widetilde{f}\right)$. The dotted lines indicate the "back" of the surface. In this example, $\tau$ is the Peacock track from Figure 2.7 and $\beta$ is conjugate to $\beta_{0}^{-1}$, from Proposition 2.4.4. Note here the 1-gons on the disk lift to 2 -gons upstairs, which are smoothed out to regular points.

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two copies of $\tau$ along the punctures and lifting the infinitesimal $k$-gons around these punctures to infinitesimal $2 k$-gons upstairs, as in Figure 2.6.

For this choice of train track $\widetilde{\tau}$, we may determine the train track map $\widetilde{f}: \widetilde{\tau} \rightarrow \widetilde{\tau}$ induced by $\psi_{h}$ (so that $h$ has pseudo-Anosov data $\left(\psi_{h}, \widetilde{\tau}, \widetilde{f}\right)$ ) from $f: \tau \rightarrow \tau$, up to a binary choice, as follows. For each edge $e \in \tau$, denote by $e^{1}$ or $e^{2}$ its two lifts in $\widetilde{\tau}$, and write $f(e)=e_{1} \ldots e_{n}$ where each $e_{i} \in \tau$ is an edge. Note that if two edges $e, e^{\prime} \in \tau$ form a train path $e e^{\prime}$ then exactly one of the paths $e^{i} e^{11}$ or $e^{i} e^{\prime 2}$ is a train path for each $i=1,2$. And, if $e^{1} e^{\prime i}$ is a train path, then so is $e^{2} e^{\prime j}$, where $i \neq j$. Set $\widetilde{f_{i_{1}}}\left(e^{1}\right)=e_{1}^{i_{1}} \ldots e_{n}^{i_{n}}$, where each $i_{j} \in\{1,2\}$ and each $e_{j}^{i_{j}} e_{j+1}^{i_{j+1}}$ is a train path. Choosing an image for $e^{1}$ also immediately determines an image for $e^{2}$, so the maps $\widetilde{f}_{1}$ and $\widetilde{f}_{2}$ are both defined on all of $\widetilde{\tau}$. See Figure 2.6.

Note that if $\widetilde{f}_{i}$ is induced by $\psi_{i}$ for $i=1,2$ then we have $\psi_{1}=\iota \circ \psi_{2}$. In particular, at most one of $\tilde{f}_{1}$ or $\widetilde{f}_{2}$ is induced by $\psi_{h}$, but this choice may be easily settled by examining $\beta$ as a braid, rather than a mapping class on the punctured sphere. So, we will denote simply by $\widetilde{f}$ the well-defined choice of $\tilde{f}_{i}$ induced by $\psi_{h}$.

### 2.3 The case with singularity type $(6 ; \emptyset ; \emptyset)$

The main goal of this section is to prove the following:
Theorem B1. Let $K$ be a genus-two, fibered, hyperbolic knot in $S^{3}$ with associated open book decomposition $\left(S_{2}^{1}, h\right)$. If $c(h) \neq 0$ and the pseudo-Anosov representative $\psi$ of $h$ has singularity type $(6 ; \emptyset ; \emptyset)$, then $K$ is not fixed point free.

This resolves Case 1 from the outline in subsection 2.1. Together with Theorem B2 in Section 2.4, this will complete the proof of Theorem B. Before turning to the proof of Theorem B1, it will be helpful to recall the Lefschetz fixed point theorem, which will be a key ingredient in our proof:

Theorem 2.3.1 (Lefschetz fixed point theorem). Let $S$ be a compact surface and $f: S \rightarrow S$ a homeomorphism. Let $f_{*}: H_{1}(S ; \mathbb{Z}) \rightarrow H_{1}(S ; \mathbb{Z})$ denote the induced map on first homology. Then

$$
2-\operatorname{tr}\left(f_{*}\right)=\sum_{p=f(p)} \operatorname{Ind}(f, p) .
$$

We will apply the Lefschetz fixed point theorem to read off information about the action of a pseudo-Anosov on homology, from its dynamical properties. The relevant result in this vein is an index calculation due to Lanneau-Thiffeault in [81]:

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 HOMOLOGY AND THE CINQUEFOILProposition 2.3.2 (Lanneau-Thiffeault). Let $\psi: S \rightarrow S$ be pseudo-Anosov with orientable invariant foliations, and let $p$ be a fixed $k$-prong singularity of $\psi$. Denote by $\psi_{*}: H_{1}(S) \rightarrow H_{1}(S)$ the action on homology, and denote by $\rho\left(\psi_{*}\right)$ the leading eigenvalue of this action, i.e. the eigenvalue with greatest absolute value.

1. If $\rho\left(\psi_{*}\right)<0$ then $\operatorname{Ind}(\psi, p)=1$; that is, every fixed point of $\psi$ has index 1 .
2. If $\rho\left(\psi_{*}\right)>0$ then either:
a) $\psi$ fixes each prong and $\operatorname{Ind}(\psi, p)=1-k<0$, or
b) $\psi$ cyclically permutes the prongs and $\operatorname{Ind}(\psi, p)=1$.

We can now use Lanneau-Thiffeault's calculation to restrict the dilatation of the pseudo-Anosov representative of a potential fixed-point-free knot $K$ in Theorem B1:

Proposition 2.3.3. Let $K \subset Y$ be a genus-two fixed-point-free knot with $c(K) \notin \mathbb{Z}$, and suppose that $Y$ is an integer homology sphere. If the pseudo-Anosov representative $\psi$ of $K$ has singularity type $(6 ; \emptyset ; \emptyset)$ then $\psi$ achieves the minimal dilatation $\lambda_{2}$ among pseudo-Anosovs in genus two.

Proof. Let $\left(S_{2}^{1}, h\right)$ be the open book decomposition of $Y$ associated to $K$, and suppose $\psi$ has no fixed points in the interior of $S$. Because $\psi$ has singularity type $(6 ; \emptyset ; \emptyset)$ by assumption, we may cap-off $\psi$ to a pseudo-Anosov $\widehat{\psi}$ on $S_{2}$ and extend the foliations preserved by $\psi$ over the capping disk. For this stratum on $S_{2}$, Masur-Smillie ([89]) prove that the foliations preserved by $\widehat{\psi}$, and therefore by $\psi$, are necessarily orientable. We will use this fact to apply the Lefschetz fixed point theorem and determine completely the Alexander polynomial of $K$.

Because $K$ is fibered, the Alexander polynomial $\Delta_{K}$ is equal to the characteristic polynomial $\chi\left(\psi_{*}\right)$ of the action of $\psi$ on homology: $\Delta_{K}=\chi\left(\psi_{*}\right)$. Because $K$ is a genus-two fibered knot in an integer homology sphere, $\Delta_{K}$ is a monic, degree-four, palindromic polynomial satisfying $\Delta_{K}(1)= \pm 1$. Moreover, because the fractional Dehn twist coefficient $c(h) \notin \mathbb{Z}$ by assumption, we know $\widehat{\psi}$ rotates the separatrices of the 6 -prong singularity $p$, so that $\operatorname{Ind}(\widehat{\psi}, p)=1$ regardless of the sign of $\rho\left(\psi_{*}\right)$. And, because $p$ is the unique fixed point of $\widehat{\psi}$ by assumption, it follows from the Lefschetz fixed point theorem that $\operatorname{tr}\left(\psi_{*}\right)=\operatorname{tr}\left(\widehat{\psi}_{*}\right)=1$.

From the discussion above, we conclude that the coefficients of $t^{4}$ and $t^{0}$ in $\Delta_{K}(t)$ are 1 , while the coefficients of $t^{3}$ and $t$ are $-\operatorname{tr}\left(\psi_{*}\right)=-1$. Now, using the fact that $\Delta_{K}(1)= \pm 1$, we see:

$$
\Delta_{K}(t)=t^{4}-t^{3} \pm t^{2}-t+1
$$

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 HOMOLOGY AND THE CINQUEFOILBecause the foliations preserved by $\psi$ are orientable, the dilatation $\lambda(\psi)$ is a root of $\Delta_{K}$. The polynomial $t^{4}-t^{3}+t^{2}-t+1$, however, has no real roots. We deduce:

$$
\Delta_{K}(t)=t^{4}-t^{3}-t^{2}-t+1
$$

Finally, note that this polynomial has a single root $\lambda_{2}$ greater than 1 , which is the minimal dilatation achieved by any pseudo-Anosov on the genus-two surface, see e.g. [81].

We will need the following lemma to finish the proof of Proposition B1:
Lemma 2.3.4. Let $h, h^{\prime} \in \operatorname{SMod}\left(S_{2}^{1}\right)$ be the lifts of braids $\beta, \beta^{\prime}$. Suppose that the capped-off maps $\widehat{h}$ and $\widehat{h^{\prime}}$ on $\widehat{S}$ are conjugate in $\operatorname{Mod}\left(S_{2}\right)$. Then, $\beta$ is conjugate to $\Delta^{2 k} \beta^{\prime}$ for some $k \in \mathbb{Z}$.

Proof. Because $\widehat{h}$ and $\widehat{h^{\prime}}$ are conjugate in $\operatorname{Mod}\left(S_{2}\right)$, and the hyperelliptic involution $\iota$ on $S_{2}$ is in the center of $\operatorname{Mod}\left(S_{2}\right)$, the conjugating mapping class in $\operatorname{Mod}\left(S_{2}\right)$ descends to the spherical mapping class group $\operatorname{Mod}\left(S_{0,6}\right)$. It follows that $\beta$ and $\beta^{\prime}$ are conjugate after capping-off to $\operatorname{Mod}\left(S_{0,6}\right)$. In particular, $\beta$ is conjugate to $\Delta^{2 k} \beta^{\prime}$ for some $k \in \mathbb{Z}$.

Though we will not need the following corollary for our purposes, it follows quickly from Lemma 2.3.4 and we believe it to be helpful in many other contexts, as well.

Corollary 2.3.5. Let $h, h^{\prime} \in \operatorname{SMod}\left(S_{g}^{r}\right)$ be the lifts of braids $\beta$, $\beta^{\prime}$, for $g, r \in\{1,2\}$. Then, $h$ and $h^{\prime}$ are conjugate in $\operatorname{Mod}\left(S_{g}^{r}\right)$ if and only if $\beta$ and $\beta^{\prime}$ are conjugate as braids.

Proof. For simplicity, suppose $g=2$ and $r=1$, though the same proof works for the other cases, with minor adjustments. If $\beta$ and $\beta^{\prime}$ are conjugate, it is clear that $h$ and $h^{\prime}$ are conjugate, too: we may simply lift the conjugating map to $S_{2}^{1}$. On the other hand, suppose $h$ and $h^{\prime}$ are conjugate in $\operatorname{Mod}\left(S_{2}^{1}\right)$. It follows that the capped-off maps $\widehat{h}$ and $\widehat{h^{\prime}}$ are conjugate in $\operatorname{Mod}\left(S_{2}\right)$. Lemma 2.3.4 now implies that $\beta$ is conjugate to $\Delta^{2 k} \beta^{\prime}$ for some $k \in \mathbb{Z}$. Because $h$ and $h^{\prime}$ are conjugate in $\operatorname{Mod}\left(S_{2}^{1}\right)$, we know that $c(h)=c\left(h^{\prime}\right)$ (see Theorem 2.2.3). It follows that $c(\beta)=2 c(h)=2 c\left(h^{\prime}\right)=c\left(\beta^{\prime}\right)$, whereas $c\left(\Delta^{2 k} \beta^{\prime}\right)=c\left(\beta^{\prime}\right)+k$, so we must have that $\beta$ and $\beta^{\prime}$ are conjugate as braids.

We need one last result before turning to the proof of Theorem B1.

Proposition 2.3.6. Let $\left(S_{2}^{1}, h\right)$ be an open book decomposition with $c(h) \notin \mathbb{Z}$, such that $h$ is symmetric and freely isotopic to a pseudo-Anosov $\psi$ with singularity type $(6 ; \emptyset ; \emptyset)$ and dilatation $\lambda(\psi)=\lambda_{2}$. Then, $\left(S_{2}^{1}, h\right)$ is not an open book decomposition for $S^{3}$.

Proof. Because $\psi$ has singularity type $(6 ; \emptyset ; \emptyset)$, we may cap-off $\psi$ to a pseudo-Anosov on $S_{2}$ with singularity type $(\emptyset ; \emptyset ; 6)$. Lanneau and Thiffeault [81] show that the pseudo-Anosov on $S_{2}$ with foliation type $(\emptyset ; \emptyset ; 6)$ and dilatation $\lambda_{2}$ is unique, up to conjugacy in $\operatorname{Mod}\left(S_{2}\right)$, inverse, and composition with the hyperelliptic involution $\iota$ on $S_{2}$. Note that the pseudo-Anosov representative $\psi_{\alpha}$ of the 5 -braid $\alpha=\sigma_{1} \sigma_{2} \sigma_{3} \sigma_{4} \sigma_{1} \sigma_{2}$ studied by Ham-Song in [44] achieves dilatation $\lambda\left(\psi_{\alpha}\right)=\lambda_{2}$. In particular, the pseudo-Anosov representative $\psi_{A}$ of the lift $A$ of $\alpha$ to $S_{2}^{1}$ achieves minimal dilatation $\lambda\left(\psi_{A}\right)=\lambda_{2}$ with the proper singularity type. It follows that $\widehat{\psi}$ is conjugate in $\operatorname{Mod}\left(S_{2}\right)$ to one of $\widehat{\psi}_{A}^{ \pm 1}$ or $\widehat{\psi}_{A}^{ \pm 1} \circ \iota$. This further implies that $\widehat{h}$ is conjugate in $\operatorname{Mod}\left(S_{2}\right)$ to one of $\widehat{A}^{ \pm 1}$ or $\widehat{A}^{ \pm 1} \circ \iota$.

Because $\iota$ is freely isotopic to the lift of the 5 -braid $\Delta^{2}$ (as described in subsection 2.2), $A^{ \pm 1} \circ \iota$ is freely isotopic to the lift of $\Delta^{2} \alpha^{ \pm 1}$. Since $h$ is symmetric by assumption, it is the lift of a braid $\beta$. $A$ is symmetric by construction, so Lemma 2.3.4 implies that $\beta$ is conjugate as a braid to $\Delta^{2 k} \alpha^{ \pm 1}$ for some $k \in \mathbb{Z}$. In particular, if $\left(S_{2}^{1}, h\right)$ is an open book decomposition for $S^{3}$, we can see that $\Delta^{2 k} \alpha^{ \pm 1}$ has unknotted closure for some $k \in \mathbb{Z}$. Moreover, note that the closure of $\Delta^{2 k} \alpha$ is unknotted if and only if the closure of $\Delta^{-2 k} \alpha^{-1}$ is also unknotted, because the unknot is amphicheiral.

By Theorem 2.2.4, if $\Delta^{2 k} \alpha^{ \pm 1}$ has unknotted closure, we must have $\left|c\left(\Delta^{2 k} \alpha^{ \pm 1}\right)\right|<$ 1. We may deduce that $0<c(\alpha)<1$ because $\alpha$ is a positive pseudo-Anosov braid, and $\Delta^{-2} \alpha$ is a negative pseudo-Anosov braid (see Theorem 2.2.3). In particular, $k<c\left(\Delta^{2 k} \alpha\right)<k+1$. So, it suffices to simply check that $\alpha$ and $\Delta^{2} \alpha^{-1}$ do not have unknotted closure. One may see this in a number of ways - for example by appealing to the self-linking number: the maximal self-linking number of the unknot is -1 , but the self-linking numbers $s l(\alpha)=1$ and $\operatorname{sl}\left(\Delta^{2} \alpha^{-1}\right)=9$ are both positive.

Proof of Theorem B1. Let $K, h$, and $\psi$ be as in the statement of the theorem. Recall that since $K$ is fixed point free by assumption, $h$ is symmetric (see Theorem 2.2.5). Since $\left(S_{2}^{1}, h\right)$ is an open book decomposition for $S^{3},|c(h)| \leq 1 / 2$ (see e.g. [70]). So if $c(h) \neq 0$ then $c(h) \notin \mathbb{Z}$. Proposition 2.3.3 then implies that the dilatation of $\psi$ is $\lambda(\psi)=\lambda_{2}$, but this contradicts Proposition 2.3.6.

Remark 2.3.7. Note that our argument does not show that no hyperbolic, genustwo, fibered knot in $S^{3}$ has the Alexander polynomial $t^{4}-t^{3}-t^{2}-t+1$ from the proof of Proposition 2.3.3. Rather, any such knot cannot have singularity type $(6 ; \emptyset ; \emptyset)$. Indeed, the knots $11 n_{38}$ and $11 n_{102}$ on KnotInfo [87] are genus-two fibered
hyperbolic knots with the given Alexander polynomial. But, their singularity types are $\left(1 ; \emptyset ; 3^{5}\right)$ and $\left(2 ; \emptyset ; 3^{4}\right)$, respectively, so the foliations preserved by their pseudoAnosov representatives are not orientable. One may check that their dilatations are $1.916 \ldots$ and $2.751 \ldots$, respectively, both of which are larger than $\lambda_{2}=1.722 \ldots$, and neither of which are roots of the given Alexander polynomial.

In the stratum $(6 ; \emptyset ; \emptyset)$, we may additionally lift the assumption that $c(h) \neq 0$ :
Proposition 2.3.8. Let $K$ be a hyperbolic, genus-two, fibered knot in $S^{3}$, with associated open book decomposition $\left(S_{2}^{1}, h\right)$. If $c(h)=0$ and the pseudo-Anosov representative $\psi$ of $h$ has singularity-type $(6 ; \emptyset ; \emptyset)$, then $K$ is not fixed point free.

Proof. Suppose that $\psi$ has no fixed points in the interior of $S_{2}^{1}$. As in the proof of Theorem B1, we may cap-off $\psi$ to a pseudo-Anosov $\widehat{\psi}$ on $S_{2}$ and extend the foliations preserved by $\psi$. Again, we have that these foliations are orientable. Note that if $\rho\left(\psi_{*}\right)<0$, then an argument identical to that of Theorem B1 will apply.

So, assuming that $\rho\left(\psi_{*}\right)>0$, the Lefschetz fixed point theorem then yields $\operatorname{tr}\left(\widehat{\psi}_{*}\right)=2-(-5)=7$, because the unique fixed point $p$ given by the boundary 6 -prong singularity is unrotated (since $c(h) \in \mathbb{Z}$ ). Consider the Markov matrix $M$ for a train track representative of $\widehat{\psi}$. It follows from a theorem of Rykken [106] that any eigenvalue of $\widehat{\psi}_{*}$ is also an eigenvalue of $M$ (counting multiplicity) except for possibly eigenvalues of 0 or roots of unity. Note here that a train track representative of $\widehat{\psi}$ has 8 real edges, so that $M$ is an $8 \times 8$ matrix, while $\widehat{\psi}_{*}$ is $4 \times 4$. In particular, $M$ has at most four more eigenvalues than $\widehat{\psi}_{*}$, and each has absolute value at most one. Hence:

$$
\operatorname{tr}(M) \geq \operatorname{tr}\left(\widehat{\psi}_{*}\right)-4=7-4=3
$$

On the other hand, a well-chosen train track carrying $\widehat{\psi}$ also carries $\psi$. In particular, by Theorem 2.2.14, we can see that $\operatorname{tr}(M)=0$ because $\psi$ is assumed to be fixed point free in the interior of $S$, which is a contradiction.

### 2.4 The case with singularity type ( $4 ; \emptyset ; 3^{2}$ )

The goal of this section is to prove the following theorem, which will resolve Case 2 from the outline in subsection 2.1:

Theorem B2. Let $K$ be a genus-two, fibered, hyperbolic knot in $S^{3}$ with associated open book decomposition $(S, h)$. If the pseudo-Anosov representative $\psi$ of $h$ has singularity type $\left(4 ; \emptyset ; 3^{2}\right)$, then $K$ is not fixed point free.


The Peacock


Figure 2.7: The two train track classes in the stratum $\left(2 ; 1^{5} ; 3\right)$ with no joints.

Remark 2.4.1. Note that we do not need to require $c(h) \neq 0$ here. Between this remark and Proposition 2.3.8, one may wonder why the assumption $c(h) \neq 0$ is necessary in the statement of Theorem B: it is purely to avoid additional singularity cases. In particular, if $c(h)=0$ then the boundary singularity may be 1-pronged or 2-pronged in Theorem 2.2.5.

The first step in proving Theorem B2 is to observe the following consequence of Theorem D , which we prove at the end of Section 2.5:

Theorem 2.4.2. Let $\psi$ be a pseudo-Anosov on $S_{0,5}^{1}$ with singularity type ( $2 ; 1^{5} ; 3$ ). Then $\psi$ is conjugate to a pseudo-Anosov carried by the Peacock train track shown in Figure 2.7.

Now, to prove Theorem B2, it suffices by Theorem 2.4.2 to look at pseudoAnosovs carried by the lift of the Peacock track. See Figure 2.6 for an image of the lifted track. We will perform a careful analysis of train track maps on this track, together with topological arguments to study a family of braids $\beta_{n}$ inducing a special collection of train track maps. We present the relevant family of braids $\beta_{n}$ and their corresponding train track maps in subsection 2.4. Then, in subsection 2.4, we study train track maps on the Peacock.

## A family of braids lifting to fixed-point-free maps

For the remainder of this section, $\tau$ will be the Peacock train track depicted on the left in Figure 2.8, with edges and vertices labeled as in the figure (with edges oriented towards the punctures); $\beta$ will be an arbitrary 5 -braid with pseudo-Anosov data $\left(\psi_{\beta}, \tau, f\right)$; and $h$ will be the lift of $\beta$ to $S=S_{2}^{1}$, with pseudo-Anosov data


Figure 2.8: Left: The Peacock train track with the labels and orientations we will use in this section. Center: $f(a)$ passes $b$ on the right. Right: $f(o)=g^{+} r^{-} g^{-} b^{\circ}$.
$\left(\psi_{h}, \widetilde{\tau}, \widetilde{f}\right)$, where $\widetilde{\tau}$ and $\widetilde{f}$ are constructed as in subsection 2.2. An image of $\widetilde{\tau}$ in this case is shown in Figure 2.6.

Let $a, b$ be any real edges in $\tau$. Note that because each peripheral 1-gon in $\tau$ is adjacent to a unique real edge, the image $f(a)$ is naturally a word of the form $w_{1} w_{2} \ldots w_{n}$, where $w_{i} \in\{o \bar{o}, g \bar{g}, p \bar{p}, b \bar{b}, r \bar{r}\}$ and $w_{n} \in\{o, g, p, b, r\}$.

Definition 2.4.3. For real edges $a, b$ of $\tau$, we say that $f(a)$ passes $b$ on the right if, before collapsing down to $\tau, b$ is to the left of $\psi_{\beta}(a)$ as in the middle of Figure 2.8. In this case, we write the letter $b^{+}$in place of $b \bar{b}$ in the word $f(a)$. We define passing on the left analogously, and denote it by $b^{-}$. If $b$ is the last letter in $f(a)$, we write $b^{\circ}$. When we allow for multiple possible options, we will write e.g. $b^{ \pm 0}, b^{+o}$, etc. See the right of Figure 2.8 for an example.

Here is the family of braids which we will study:
Proposition 2.4.4. Set $\beta_{n}=\sigma_{1}^{n+2} \sigma_{2} \sigma_{3} \sigma_{4} \sigma_{1} \sigma_{2} \sigma_{3} \sigma_{4}^{2}$ for $n \geq 0$. Then, $\beta_{n}^{-1}$ is pseudoAnosov, and conjugate to a braid carried by $\tau$, which induces the train track map $f_{n}: \tau \rightarrow \tau$ defined by:

$$
\begin{aligned}
& f_{n}(o)=p^{\circ} \quad f_{n}(g)=b^{\circ} \\
& f_{n}(p)=\left\{\begin{array}{lll}
\left(r^{-} o^{-}\right)^{\left(\frac{n}{2}+1\right)} r^{\circ} & n \text { even } \\
\left(r^{-} o^{-}\right)^{\frac{n+1}{2}} r^{-} o^{\circ} & n \text { odd }
\end{array} f_{n}(b)= \begin{cases}\circ & \left(r^{-} o^{-}\right)^{\frac{n}{2}} r^{-} o^{\circ} \\
\left(r^{-} o^{-}\right)^{\frac{n+1}{2}} r^{\circ} & n \text { even }\end{cases} \right.
\end{aligned}
$$

Proof. Figure 2.9 verifies that $H \beta_{0}^{-1} H^{-1}$ is carried by $\tau$ and induces the train track map $f_{0}: \tau \rightarrow \tau$, where the orientation-preserving homeomorphism $H$ is given by swinging the real edge $r$ around the train track to the other side. For $n \geq 1$, note


Figure 2.9: The train track map induced by $H \beta_{n}^{-1} H^{-1}$, where $H$ is an orientationpreserving homeomorphism which swings $r$ around the rest of the track.
that $\beta_{n}=\sigma_{1} \beta_{n-1}$, and the additional $\sigma_{1}$ simply adds more twists between the leftmost edges before composing with $H^{-1}$ in the last step. This additional twisting adds words of the form $\left(r^{-} o^{-}\right)$to $f_{n}(p)$ and $f_{n}(b)$, and swaps which edges $p$ and $b$ end on, as in the map in the proposition statement.

It now remains to verify that $\beta_{n}^{-1}$ is pseudo-Anosov. This can be seen by checking that the transition matrix $M_{n}$ associated to the train track map $f_{n}$ is PerronFrobenius. When determining the matrix $M_{n}$ from the map $f_{n}$ given above, it may be helpful to recall that, for each real edge $a, b$ of $\tau$, each instance of $b^{ \pm}$in $f(a)$ records the word $b b$, and each instance of $b^{\circ}$ records just $b$. Regardless of the parity of $n$, the transition matrix is:

$$
M_{n}=\left(\begin{array}{ccccc}
0 & 0 & n+2 & n+1 & 0 \\
0 & 0 & 0 & 0 & 1 \\
1 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
0 & 0 & n+3 & n+2 & 0
\end{array}\right)
$$

It is straightforward to check that $M_{n}^{7}$ is strictly positive for all $n \geq 0$, so $M_{n}$ is Perron-Frobenius.

It follows from Proposition 2.4.4 that any braid inducing the train track map $f_{n}: \tau \rightarrow \tau$ must be conjugate to $\beta_{n}^{-1}$ within the spherical mapping class group (see Remark 2.2.13). The following proposition then implies that no braid inducing the map $f_{n}$ has unknotted closure:


Figure 2.10: Isotopies of $\widehat{\Delta^{2} \beta_{n}^{-1}}$ to $P(3,3-n,-2)$

Proposition 2.4.5. The braid closure $\widehat{\Delta^{2 k} \beta_{n}^{ \pm 1}}$ is not an unknot for any $k, n \in \mathbb{Z}$.
Proof. Because $\beta_{n}$ is a positive pseudo-Anosov braid, we must have $c\left(\beta_{n}\right)>0$. On the other hand, one can check that $\beta_{n}$ is braid-isotopic to $\beta_{n}^{\prime}=\left(\sigma_{1} \sigma_{2} \sigma_{3} \sigma_{4}\right)^{2} \sigma_{4}^{n+2}$. One can easily check that $c\left(\beta_{n}^{\prime}\right) \leq 1$ in a variety of ways- for example by verifying that $\Delta^{-2} \beta_{n}^{\prime}$ is $\sigma_{1}$-negative (see Theorem 2.2.3). And, because $\beta_{n}$ and $\beta_{n}^{\prime}$ are braid-isotopic, it follows that $c\left(\beta_{n}\right) \leq 1$, too. So, we have:

$$
k<c\left(\Delta^{2 k} \beta_{n}\right) \leq k+1
$$

Now, by Theorem 2.2.4 and the fact that the unknot is amphicheiral, it suffices to check that neither $\widehat{\beta_{n}}$ nor $\widehat{\Delta^{2} \beta_{n}^{-1}}$ is an unknot for any $n \in \mathbb{Z}$. The closure $\widehat{\beta_{n}}$ is easily seen to be the torus knot $T(2, n+6)$. Figure 2.10 verifies that the closure $\widehat{\Delta^{2} \beta_{n}^{-1}}$ is the 3 -stranded pretzel $\operatorname{knot} P(3,3-n,-2)$. These knots are all known to not be unknotted.

## Train track maps on the Peacock

In this subsection, we retain the notation from the previous subsection. Our remaining goal is to prove the following proposition, which, together with Propositions 2.4.4 and 2.4.5, will imply Theorem B2 and complete the proof of Theorem B:

Proposition 2.4.6. Let $\psi_{\beta}$ be a pseudo-Anosov carried by $\tau$, which lifts to a map $\psi_{h}$ in the cover. If $\psi_{h}$ is fixed point free, then $\beta$ is conjugate in the spherical mapping class group to $\beta_{n}$ or $\beta_{n}^{-1}$ for some $n \in \mathbb{Z}$. In particular, $\beta$ is conjugate as a braid to $\Delta^{2 k} \beta_{n}^{ \pm 1}$ for some $n, k \in \mathbb{Z}$.

We begin with some helpful lemmas to simplify the case analysis.
Lemma 2.4.7 (Trace Lemma). If $\psi_{h}$ is fixed point free, then for any real edge $a \in \tau$, we have that $a^{ \pm 0} \notin f(a)$.

Proof. First, if $a^{\circ} \in f(a)$, then the marked point at the end of $a$ is fixed by $\psi_{\beta}$, and the lift of this marked point is fixed by $\psi_{h}$. Next, suppose that $a^{ \pm} \in f(a)$, and recall that this means that $f(a)$ contains a word of the form $a \bar{a}$. In the lift, it follows that $\widetilde{f}\left(a^{1}\right)$ contains a word of the form $a^{i} \bar{a}^{j}$ for some $i, j \in\{1,2\}$ (see the construction in subsection 2.2). Because the edges $a^{i}$ and $a^{j}$ are not adjacent to infinitesimal loops in the cover (the infinitesimal loops in $\tau$ lift to regular points in the cover), we can see that $i \neq j$. So, $\underset{\sim}{f}\left(a^{1}\right)$ contains either $a^{1}$ or $\bar{a}^{1}$ as a letter. In either case, the transition matrix of $\widetilde{f}$ has non-zero trace, so $\psi_{h}$ is not fixed point free by Theorem 2.2.14.


Figure 2.11: $c$ will be absorbed into $f(v)$.

The Trace Lemma also holds in general, with the same proof, for any jointless train track with only 1-pronged punctures. Because we will use the Trace Lemma with great frequency in this section, when we invoke this lemma we will often use the shorthand "by trace."

Lemma 2.4.8. If $\psi_{h}$ is fixed point free, then $f\left(v_{i}\right) \neq v_{i}$ for $i=1,2,3$.
Proof. If $f\left(v_{i}\right)=v_{i}$ for some $i$, then $D f(r)=r^{ \pm 0}$, which is forbidden by the Trace Lemma.

The above lemma implies that $f\left(v_{1}\right) \in\left\{v_{2}, v_{3}\right\}$, and this choice also determines the images $f\left(v_{i}\right)$ for $i=2,3$. Note that there is a natural horizontal symmetry of $\tau$ induced by reversing the orientation of the disk. Composing with this symmetry takes a braid to its reverse inverse, and a braid lifts to a fixed-point-free map if and only if its reverse inverse does. Hence, it suffices to choose either one of the images $f\left(v_{1}\right)$ as above. Therefore, without loss of generality, $f\left(v_{1}\right)=v_{3}$.

Lemma 2.4.9. Let $a, b$ be any two real edges of $\tau$ with $a$ and $b$ adjacent at initial vertex $v$. Then, for any real edge $c$ of $\tau, \psi_{\beta}(c)$ does not intersect the convex cone determined by the initial segments of $\psi_{\beta}(a)$ and $\psi_{\beta}(b)$. See Figure 2.11 for reference.

Proof. By injectivity of $\psi_{\beta}$, we know that $\psi_{\beta}(c)$ may not cross $\psi_{\beta}(b)$ or $\psi_{\beta}(a)$. If $\psi_{\beta}(c)$ enters the convex cone $X$ on the initial segments of $\psi_{\beta}(a)$ and $\psi_{\beta}(b)$, then $\psi_{\beta}(c)$ must either have its endpoint inside $X$ or must leave $X$. The first case is not possible because $c$ ends at a vertex of an infinitesimal monogon by assumption. Therefore $\psi_{\beta}(c)$ must enter and leave $X$. Let $A$ and $B$ denote the strips of the fibered surface $F$ that collapse onto $a$ and $b$, respectively. The $\operatorname{arc} \psi_{\beta}(c)$ lies transverse to the fibers of $F$. Assume without loss of generality that $\psi_{\beta}(c)$ enters $X$ along $A$. Since it must exit $X, \psi_{\beta}(c)$ must subsequently traverse either $A$ or $B$. Neither case is possible, however, since $a$ and $b$ form a cusp: the arc $\psi_{\beta}(c)$ is forced to be non-smooth.


Figure 2.12: Cases for the proof of Lemma 2.4.11.

When a situation as in Lemma 2.4.9 arises, we say that the edge $c$ is absorbed into $f(v)$. In practice, an edge being absorbed into a vertex is much easier to spot visually than by formal definition: the typical picture is the one depicted in Figure 2.11.

For the following arguments, recall that we assume without loss of generality that $f\left(v_{1}\right)=v_{3}$. The figures provided in each proof below are single scenarios appearing in each main case, not exhaustive images of every possibility. We strongly encourage the reader to draw by hand the train track maps which are written in words, as they read through each argument.

Lemma 2.4.10. If $\psi_{h}$ is fixed point free, then $D f(r) \notin\left\{o^{+}, g^{+}\right\}$.
Proof. If $D f(r)=o^{+}$, then we must have $f(r)=o^{+} r^{ \pm 0} \ldots$, which is forbidden by the Trace Lemma. The same argument applies to $g^{+}$.

Lemma 2.4.11. If $\psi_{h}$ is fixed point free, then $D f(p)=r^{-}$.
Proof. First suppose that $D f(p)=r^{+}$. Then, the second letter in $f(p)$ is either $p$ or $b$. The former is not allowed by trace. In the latter case, note that then $f(b)=r^{+} b^{ \pm 0} \ldots$ since $p$ is to the left of $b$, which is again ruled out by trace.

Now, suppose that $f(p)=r^{\circ}$. Note that then $D f(b)=r^{+}$, and this further implies by trace that $f(b)=r^{+} p^{ \pm \circ} \ldots$ It then follows that $D f(o)=p^{ \pm \circ}$, so we will check these three possible cases individually. See Figure 2.12.

Case 1: $f(o)=p^{\circ}$. In this case, $f(b)=r^{+} p^{+} \ldots$ and $D f(g)=p^{+}$. By trace, it follows that $f(g)=p^{+} o^{ \pm \circ} \ldots$, and this in turn forces $D f(r)=o^{ \pm \circ}$. By Lemma 2.4.10, we know $D f(r) \neq o^{+}$, so $D f(r)=o^{-\circ}$. If $D f(r)=o^{\circ}$, then the real edges $o$, $p$, and


Figure 2.13: Cases for the proof of Lemma 2.4.12.
$r$ are permuted, implying that the transition matrix of $f$ is not Perron-Frobenius. Finally, if $D f(r)=o^{-}$, we can see that $f(r)=o^{-} p^{-} r^{ \pm \circ} \ldots$, which is ruled out by trace.

Case 2: $D f(o)=p^{-}$. In this case, we must have $f(o)=p^{-} r^{-} g^{ \pm 0} \ldots$ by trace (or else $f(o)$ is absorbed into either $f\left(v_{3}\right)$ or $f\left(v_{1}\right)$, if it goes "inside" $b$ or $g$, respectively). This further implies that $f(g)=p^{-} r^{-} g^{ \pm 0} \ldots$, which is ruled out by trace.

Case 3: $D f(o)=p^{+}$. Here, we have $f(o)=p^{+} g^{ \pm \circ} \ldots$, and therefore $f(b)=$ $r^{+} p^{+} g^{ \pm \circ} \ldots$ In particular, this forces $f(g)=p^{+} g^{ \pm 0} \ldots$, which is ruled out by trace.

Lemma 2.4.12. If $\psi_{h}$ is fixed point free, then $D f(b)=r^{-}$.
Proof. By Lemma 2.4.11, we may assume $D f(p)=r^{-}$. Note that $D f(b)=r^{+}$is not possible, because then one of $b$ or $p$ will be absorbed into $f\left(v_{3}\right)$.

So, suppose $f(b)=r^{\circ}$. We branch along cases for the second letter in $f(p)$. Note that if $f(p)=r^{-} g^{+0} \ldots$, then $D f(r)=g^{+}$, which contradicts Lemma 2.4.10. So, there

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are four cases left to consider, shown in Figure 2.13.
Case 1: $f(p)=r^{-} g^{-} \ldots$ In this case, we must have $f(p)=r^{-} g^{-} b^{ \pm 0} \ldots$, because otherwise $p$ is absorbed into $f\left(v_{3}\right)$ (if it follows $r$ next) or contributes trace (if it follows $p$ next). Now note that $D f(o)=b^{ \pm \circ}$. If $D f(o)=b^{+\circ}$ then $f(g)=b^{+} g^{ \pm \circ} \ldots$, which is ruled out by trace. So, $D f(o)=b^{-}$and we must have $f(o)=b^{-} r^{-} g^{ \pm \circ} \ldots$ by trace. Note that $D f(g) \neq b^{+}$by Lemma 2.4.9, so $D f(g)=b^{-\circ}$. If $D f(g)=b^{-}$, then we have $f(g)=b^{-} r^{-} g^{ \pm \circ} \ldots$, which is ruled out by trace. So, $D f(g)=b^{\circ}$. Finally, consider cases for $D f(r)$. We know $D f(r)=g^{ \pm \circ}$, and by Lemma 2.4.10, we must have $D f(r)=g^{-\circ}$. If $D f(r)=g^{\circ}$, then the transition matrix is not Perron-Frobenius. If $D f(r)=g^{-}$, then $f(r)=g^{-} b^{-} r^{ \pm 0} \ldots$, which is ruled out by trace.

Case 2: $f(p)=r^{-} o^{+} \ldots$ In this case, $f(p)=r^{-} o^{+} r^{+} b^{ \pm 0} \ldots$ by trace and, by Lemma 2.4.10, it follows that $D f(r)=g^{-\circ}$. Now, look at cases for $D f(o)$.

If $D f(o)=b^{+o}$ then $f(g)=b^{+} g^{ \pm 0} \ldots$, which is ruled out by trace. And, if $D f(o)=$ $b^{-}$then $f(o)=b^{-} r^{-} o^{ \pm \circ} \ldots$, also ruled out by trace. If $D f(o)=p^{-}$then similarly $f(o)=p^{-} r^{-} o^{ \pm 0} \ldots$, which is again ruled out by trace. There are two remaining subcases to consider:

Subcase 2A: $D f(o)=p^{\circ}$. Here, consider cases for $D f(g)$ : either $D f(g)=p^{+}$or $D f(g)=b^{ \pm \circ}$. If $D f(g)=p^{+}$then $f(g)=p^{+} g^{ \pm \circ} \ldots$ which is ruled out by trace. We cannot have $D f(g)=b^{+}$because then $g$ is absorbed into $f\left(v_{1}\right)$. Finally, we cannot have $D f(g)=b^{-\circ}$ because then either $p$ is absorbed into $f\left(v_{1}\right)$ (if $D f(g)=b^{\circ}$ or $D f(g)=b^{-}$with $p$ outside $g$ ) or $g$ is absorbed into $f\left(v_{3}\right)$ (if $D f(g)=b^{-}$and $g$ is outside $p$ ).

Subcase 2B: $D f(o)=p^{+}$. Here, consider cases for $D f(r)$ : either $D f(r)=g^{\circ}$ or $D f(r)=g^{-}$, by Lemma 2.4.10. If $D f(r)=g^{\circ}$ then $o$ will be absorbed into $f\left(v_{3}\right)$. If $D f(r)=g^{-}$, then either $o$ is inside $r$, in which case $f(r)=g^{-} p^{-} r^{ \pm 0} \ldots$ (which is ruled out by trace); or, $r$ is inside $o$, in which case $o$ will be absorbed into $f\left(v_{3}\right)$.

Case 3: $f(p)=r^{-} o^{-} \ldots$ In this case, we must have $f(p)=r^{-} o^{-} b^{ \pm 0} \ldots$ (otherwise $p$ will be absorbed into $f\left(v_{3}\right)$ or contribute trace), which forces $D f(o)=b^{ \pm \circ}$ and $D f(g)=b^{ \pm 0}$, as well. If $D f(o)=b^{+}$, then either $f(o)=b^{+} o^{ \pm 0} \ldots$ (which is ruled out by trace), or $f(o)=b^{+} g^{ \pm \circ} \ldots$ in which case $f(g)=b^{+} g^{ \pm \circ} \ldots$, too (which is again ruled out by trace). And, if $D f(o)=b^{-}$, then $f(o)=b^{-} r^{-} o^{ \pm 0} \ldots$, which is ruled out by trace.

So, we must have $f(o)=b^{\circ}$. Note here that we must have $D f(r)=o^{-0}$, by Lemma 2.4.10, and $D f(g)=b^{+}$. Now, look at $r$ : if $D f(r)=o^{-}$, then $f(r)=$ $o^{-} b^{-} r^{ \pm \circ} \ldots$ which is ruled out by trace. Finally, if $D f(r)=o^{\circ}$, then either $g$ is inside $p$, in which case $g$ is absorbed into $f\left(v_{3}\right)$, or $p$ is inside $g$, in which case $p$ will be absorbed into $f\left(v_{1}\right)$.

Case 4: $f(p)=r^{-} o^{\circ}$ In this case, note that $D f(r)=g^{-\circ}$ by Lemma 2.4.10, so consider subcases for $D f(r)$.

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 HOMOLOGY AND THE CINQUEFOILSubcase 4A: $D f(r)=g^{-}$. Here, consider cases for $D f(o)$. If $D f(o)=b^{+o}$ then $f(g)=b^{+} g^{ \pm \circ} \ldots$, which is ruled out by trace. If $D f(o)=b^{-}$then $f(o)=b^{-} r^{-} o^{-} \ldots$, again ruled out by trace. If $D f(o)=p^{-o}$ then $f(r)=g^{-} p^{-} r^{ \pm 0} \ldots$, ruled out by trace. And finally, if $D f(o)=p^{+}$then either $o$ is inside $r$ in which case $f(r)=g^{-} p^{-} r^{ \pm o} \ldots$ (which is ruled out trace), or $o$ is outside $r$ in which case it will be absorbed into $f\left(v_{3}\right)$.

Subcase 4B: $D f(r)=g^{\circ}$. Look first at $D f(o)$. If $D f(o)=p^{+}$or $D f(o)=b^{+}$ then $o$ will be absorbed into $f\left(v_{3}\right)$. If $D f(o)=p^{-}$or $D f(o)=b^{-}$then $f(o)=$ $p^{-} r^{-} o^{ \pm \circ} \ldots$ or $f(o)=b^{-} r^{-} o^{ \pm 0} \ldots$, both of which are ruled out trace. So, we must have either $D f(o)=p^{\circ}$ or $D f(o)=b^{\circ}$ and then it follows that $D f(g)=b^{\circ}$ or $D f(g)=p^{\circ}$, respectively, as well, after some simple analysis on $D f(g)$. But, note that $D f(o)=b^{\circ}$ and $D f(g)=p^{\circ}$ is not possible, since $\psi_{\beta}$ is orientation-preserving. And, $D f(o)=p^{\circ}$ and $D f(g)=b^{\circ}$ is not possible, because then the transition matrix is not Perron-Frobenius.

Lemma 2.4.13. If $\psi_{h}$ is fixed point free, then $D f(r)=g^{-\circ}$.
Proof. By Lemma 2.4.10, we know that $D f(r) \notin\left\{g^{+}, r^{+}\right\}$, so we just need to show that $D f(r) \neq o^{-\circ}$. Suppose otherwise, i.e. $D f(r)=o^{-\circ}$. By Lemmas 2.4.11 and 2.4.12, we know that $D f(p)=D f(b)=r^{-}$, and then because $D f(r)=o^{-\circ}$ by assumption, it follows that $f(p)=r^{-} o^{-} \ldots$ and $f(b)=r^{-} o^{-} \ldots$ From here, we must have $f(b)=r^{-} o^{-} p^{ \pm 0} \ldots$ by trace, but then $f(p)=r^{-} o^{-} p^{ \pm 0}$, which is ruled out by trace.

We are finally ready to prove Proposition 2.4.6.
Proof of Proposition 2.4.6. Note by the discussion after Lemma 2.4.8, it suffices to assume $f\left(v_{1}\right)=v_{3}$. By Lemmas 2.4.11 and 2.4.12, we know $D f(p)=D f(b)=r^{-}$, and by Lemma 2.4.13 we know $D f(r)=g^{-\circ}$. From here, we branch along cases for $D f(o)$. The cases $D f(o)=b^{ \pm \circ}$ can be ruled out quickly as follows.

If $D f(o)=b^{+}$, then $f(o)=b^{+} g^{ \pm \circ} \ldots$ by trace. But, then $f(g)=b^{+} g^{ \pm \circ} \ldots$, which is ruled out by trace. If $D f(o)=b^{\circ}$, then $f(g)=b^{+} g^{ \pm \circ} \ldots$ because $D f(r)=g^{-\circ}$, and this is ruled out by trace. Finally, if $D f(o)=b^{-}$, then $f(o)=b^{-} r^{-} g^{ \pm 0} \ldots$ by trace. Because $D f(r)=g^{-\circ}$, it follows that $f(p)=r^{-} g^{-} \ldots$ and $f(b)=r^{-} g^{-}$.. From here, we must have $f(b)=r^{-} g^{-} p^{ \pm 0} \ldots$ by trace. But, then $f(p)=r^{-} g^{-} p^{ \pm 0}$, too, which is ruled out by trace.

And, note that if $D f(o)=p^{-}$, then $f(o)=p^{-} r^{-} g^{ \pm 0} \ldots$ by trace. Here, we must have $f(p)=r^{-} g^{-} p^{ \pm \circ} \ldots$ because $D f(r)=g^{-\circ}$, which is ruled out by trace. So, we have two cases left to consider, shown in Figure 2.14.


Figure 2.14: Cases for the proof of Proposition 2.4.6. In this figure, we have chosen to omit the shaded collapsing regions for readability.

Case 1: $D f(o)=p^{+}$. In this case, we have $f(o)=p^{+} g^{ \pm 0} \ldots$ by trace. If $f(o)=$ $p^{+} g^{-\circ} \ldots$ then $f(r)$ is either absorbed into $f\left(v_{1}\right)$ or passes over $r$. So $f(o)=p^{+} g^{+} r^{ \pm 0} \ldots$ Next, consider $D f(g)$. If $D f(g)=b^{+}$then $g$ is absorbed into $f\left(v_{1}\right)$, and the case $D f(g)=p^{+}$is ruled out quickly by trace. So, we have $D f(g)=b^{-\circ}$. In either case, note that for both of $f(p)$ and $f(b)$, the second letter is in the set $\left\{o^{ \pm 0}, g^{-}\right\}$because $D f(r)=g^{-\circ}$. If $f(p)=r^{-} g^{-} \ldots$ then $f(p)=r^{-} g^{-} p^{ \pm \circ} \ldots$, which is ruled out by trace. So, $f(p)=r^{-} o^{ \pm 0} \ldots$

Similarly, if $f(b)=r^{-} g^{-} \ldots$ then either $b$ is outside $o$, in which case $b$ will be absorbed into $f\left(v_{1}\right)$, or $o$ is outside $b$, in which case $o$ will be absorbed into $f\left(v_{3}\right)$. So, we must have $f(b)=r^{-} o^{ \pm 0} \ldots$, too.

Now, if $f(p)=r^{-} o^{+o} \ldots$ then $f(b)=r^{-} o^{+} r^{+} b^{ \pm 0} \ldots$ which is ruled out by trace. So, we must have $f(p)=r^{-} o^{-} r^{ \pm 0} \ldots$ And, if $f(p)=r^{-} o^{-} r^{-0} \ldots$ then $f(o)=b^{+} g^{+} r^{-} o^{ \pm 0} \ldots$, which is ruled out by trace. So, we have instead $f(p)=r^{-} o^{-} r^{+} g^{-} p^{ \pm \circ}$... because $D f(r)=g^{-\circ}$, which is again ruled out by trace.

Case 2: $D f(o)=p^{\circ}$. Here, we branch along subcases for $D f(g)$. Note that if $D f(g)=b^{+}$then $g$ is absorbed into $f\left(v_{1}\right)$, and if $D f(g)=p^{+}$then $f(g)=p^{+} g^{ \pm 0} \ldots$, which is ruled out by trace. So we have two remaining subcases to consider:

Subcase 2A: $D f(g)=b^{-}$. Because $D f(r)=g^{-\circ}$, note that if $f(p)=r^{-} g^{-} \ldots$ then $p$ is absorbed into $f\left(v_{1}\right)$. A similar argument applies to $f(b)$, so we must have both $f(p)=r^{-} o^{ \pm 0} \ldots$ and $f(b)=r^{-} o^{ \pm 0} \ldots$ Now, if $f(b)=r^{-} o^{+} \ldots$ then either $b$ is absorbed into $f\left(v_{3}\right)$, or $f(b)=r^{-} o^{+} r^{+} b^{ \pm \circ}$, which is ruled out by trace. So, either $f(b)=r^{-} o^{-} r^{ \pm 0} \ldots$ or $f(b)=r^{-} o^{\circ}$.

In the first case, note that if $f(b)=r^{-} o^{-} r^{+} \ldots$, then either $b$ is absorbed into $f\left(v_{3}\right)$, absorbed into $f\left(v_{1}\right)$, or $f(b)=r^{-} o^{-} r^{+} o^{+} r^{+} b^{ \pm \circ} \ldots$, which is ruled out by trace.

So, we must have either $f(b)=r^{-} o^{-} r^{-} \ldots$ or $f(b)=r^{-} o^{-} r^{\circ}$.
We can then see that $f(b)=\left(r^{-} o^{-}\right)^{k} r^{-} o^{\circ}$ or $f(b)=\left(r^{-} o^{-}\right)^{k+1} r^{\circ}$ for some $k \geq 0$. The argument that follows will not depend on $k$ (with large $k$, all remaining strands will just turn more times along $r$ and $o$ ), so for simplicity suppose either $f(b)=r^{-} o^{\circ}$ or $f(b)=r^{-} o^{-} r^{\circ}$.

First, suppose $f(b)=r^{-} o^{\circ}$. Then, we must have $f(p)=r^{-} o^{-} r^{ \pm 0} \ldots$ and $f(g)=$ $b^{-} r^{-} o^{-} \ldots$ Note here that if $f(p)=r^{-} o^{-} r^{+\circ} \ldots$ then $g$ will be absorbed into $f\left(v_{3}\right)$. But, if $f(p)=r^{-} o^{-} r^{-} \ldots$, then $p$ will be absorbed into $f\left(v_{3}\right)$. This same argument will work for arbitrary $k$ after several twists around $r$ and $o$.

Next, suppose $f(b)=r^{-} o^{-} r^{\circ}$. The argument is very similar in this case. We must have $f(p)=r^{-} o^{-} r^{-} o^{ \pm \circ} \ldots$ and $f(g)=b^{-} r^{-} o^{-} r^{-} o^{ \pm \circ} \ldots$ If $f(p)=r^{-} o^{-} r^{-} o^{-\circ} \ldots$ then $g$ will be absorbed into $f\left(v_{3}\right)$. And, if $f(p)=r^{-} o^{-} r^{-} o^{+}$then either $g$ will be absorbed into $f\left(v_{3}\right)$ or $p$ will be absorbed into $f\left(v_{1}\right)$. As before, the same argument will work for arbitrary $k$ after several additional twists around $r$ and $o$.

Subcase 2B: $D f(g)=b^{\circ}$. We cannot have $D f(r)=g^{-}$since then $r$ will be absorbed into $f\left(v_{1}\right)$. So, we must have $D f(r)=g^{\circ}$. From here, note that $f(b)=r^{-} o^{-\circ} \ldots$ because otherwise $b$ will be absorbed into either $f\left(v_{1}\right)$ or $f\left(v_{3}\right)$. In either case, it follows that $f(p)=r^{-} o^{-} r^{-\circ} \ldots$ because otherwise $p$ will be absorbed into either $f\left(v_{1}\right)$ or $f\left(v_{3}\right)$. Iterating the same argument, it is now easy to see that $f(b)=\left(r^{-} o^{-}\right)^{\frac{n}{2}} r^{-} o^{\circ}$ or $f(b)=\left(r^{-} o^{-}\right)^{\frac{n+1}{2}} r^{\circ}$, and $f(p)=\left(r^{-} o^{-}\right)^{\left(\frac{n}{2}+1\right)} r^{\circ}$ or $f(p)=$ $\left(r^{-} o^{-}\right)^{\frac{n+1}{2}} r^{-} o^{\circ}$ for some $n \geq 0$.

One may observe that these final train track maps match identically with the ones given in Proposition 2.4.4. Proposition 2.4.4 and the subsequent discussion then implies that the braid $\beta$ is conjugate in the spherical mapping class group to $\beta_{n}^{-1}$ for some $n$.

### 2.5 The tight splitting

This section is devoted to developing a tool which will be integral to the proof of Theorem 2.4.2: a specialized form of "splitting," which will allow us to restrict our attention to pseudo-Anosovs carried by a single train track.

## Standardly embedded tracks

We first describe a particular class of train tracks on the punctured disk, called standardly embedded tracks, which will aid in the description of our splitting procedure. Standardly embedded train tracks have previously appeared in the work of

## CHAPTER 2. FIXED-POINT-FREE PSEUDO-ANOSOVS, KNOT FLOER

 HOMOLOGY AND THE CINQUEFOILKo-Los-Song ([71]), Cho-Ham ([15]), and Ham-Song ([44]), who used them to study pseudo-Anosovs on $S_{0, n}^{1}$ for small $n$.

Definition 2.5.1. An infinitesimal polygon of a train $\operatorname{track} \tau$ is a connected component of $S_{0, n}^{1} \backslash \tau$ whose boundary consists of finitely many infinitesimal edges of $\tau$. A train track $\tau$ on $S_{0, n}^{1}$ is standardly embedded if the following conditions hold:

1. Every component of $S_{0, n}^{1} \backslash \tau$ is an infinitesimal polygon, except for the one containing $\partial S_{0, n}^{1}$.
2. If two edges of $\tau$ are tangent at a switch, then either both are real or both are infinitesimal.
3. Cusps only occur at vertices of infinitesimal polygons.

Figure 2.3 is an example of a standardly embedded track, and Figure 2.5 shows a pseudo-Anosov carried by this track, as well as the induced train track map. Every train track may be adjusted to a standarly embedded one, and this adjustment does not affect which pseudo-Anosovs the track carries. So, we have:

Proposition 2.5.2. Every pseudo-Anosov on $S_{0, n}^{1}$ is carried by a standardly embedded train track.

We adapt the following definition from Ham-Song's notion of an elementary folding map [44].

Definition 2.5.3. Let $\tau, \tau_{1} \hookrightarrow S_{0, n}^{1}$ be standardly embedded train tracks. A Markov map is a graph map $p: \tau_{1} \rightarrow \tau$ that maps vertices to vertices, and is locally injective away from the preimages of vertices. An elementary folding map is a smooth Markov map such that for exactly one real edge $\alpha$, the image $p(\alpha)$ has word length 2 , while the images of all other edges have word length 1 . We require that the distinguished edge $\alpha$ belong to a cusp $(\alpha, \beta)$ of $\tau_{1}$, and that $p(\alpha)$ be of the form

$$
p(\alpha)=p(\beta) \cdot a
$$

where $a$ is a real edge joined to $p(\beta)$ by an infinitesimal edge.
For the purposes of this paper, an elementary folding map $p: \tau_{1} \rightarrow \tau$ will be the identity map away from the distinguished real edge $\alpha$. See Figure 2.15.

Remark 2.5.4. An elementary folding map in our terminology is the composition of two elementary moves in Ham-Song's terminology [44].


Figure 2.15: An example of an elementary folding map. The map $p$ is the identity except at the edge $\alpha$, which is mapped as a directed path to $\beta \cdot e \cdot a$.

## The tight splitting

We are ready to introduce a specialized variant of a classical operation on train tracks, known as splitting (cf. Section 2.1 of [HP]). Our variant, which we call tight splitting, takes as input the data $(\tau, \psi, f)$ of a pseudo-Anosov $\psi$ with an invariant standardly embedded train track $\tau$ and outputs another train track $\tau_{1}$ that is also invariant under $\psi$.

More precisely, suppose that $(\tau, \psi, f)$ is the data of a pseudo-Anosov $\psi$ on $S_{0, n}^{1}$ carried by the standardly embedded train track $\tau$ :


Suppose further that $\tau_{1} \hookrightarrow S_{0, n}^{1}$ is another standardly embedded train track such that there exists an elementary folding map $p: \tau_{1} \rightarrow \tau$. Then there is a welldefined elementary folding $\operatorname{map} p_{\psi}: \psi\left(\tau_{1}\right) \rightarrow \psi(\tau)$ such that the following diagram commutes:


If $\tau_{1}$ were invariant under $\psi$, we would then be able to complete the above commutative diagram as follows:


Unfortunately, $\tau_{1}$ need not be invariant under $\psi$, as Example 2.5.5 shows. We introduce tight splitting to deal with this problem by taking into account the train track map $f: \tau \rightarrow \tau$. In particular, we will show that if $(\tau, \psi, f)$ is the data of a pseudo-Anosov acting on a standardly embedded train track $\tau$, then there is always a split $p: \tau_{1} \rightarrow \tau$ such that $\tau_{1}$ is still invariant under $\psi$. See Proposition 2.5.10.

Example 2.5.5. Consider the pseudo-Anosov $\psi: S_{0,5}^{1} \rightarrow S_{0,5}^{1}$ represented in Figure 2.16. By computing a right eigenvector for the Perron-Frobenius eigenvalue of the transition matrix $M$, we see that edge $e_{2}$ has smaller transverse measure than edge $e_{3}$. Following Harer-Penner, we can perform a split by peeling $e_{2}$ away from $e_{3}$ (truthfully, this is two splits since we need to also peel $e_{2}$ away from the non-expanding edge connecting it to $e_{3}$ ). The resulting train track is shown on the left of Figure 2.17. Unhappily, $\tau_{1}$ is not invariant under $\psi$.

After some thought, one might notice that a different split does produce an invariant train track for $\psi$. Indeed, we can see from Figure 2.16 that $e_{4}$ has smaller transverse measure than $e_{5}$, and performing the corresponding pair of splits (first over the black edge connecting $e_{4}$ to $e_{5}$, then over $e_{5}$ ) produces an invariant train track $\tau_{2}$, as desired. See Figure 2.18. This second split differs from the first in that it is compatible with the action of $\psi$ on $\tau$ : all paths in the image $\psi(\tau)$ that collapse onto $e_{4}$ also collapse onto $e_{5}$. It is this property that allows us to easily isotope strands in $\psi(\tau)$ to lie transverse to the leaves of the fibered neighborhood of $\tau_{2}$. See Definition 2.5.8 and Proposition 2.5.10 for more.

Let $\tau \hookrightarrow S_{0, n}^{1}$ be standardly embedded, and let $v \in \tau$ be a switch. The link of $v$ is the collection $\operatorname{Lk}(v)$ of edges of $\tau$ incident to $v$. The elements of $\operatorname{Lk}(v)$ inherit a natural counterclockwise cyclic order $e_{1}, \ldots, e_{k}$. A subset $C \subseteq \operatorname{Lk}(v)$ is connected if whenever $e_{i}, e_{j} \in C$ and $i<j$, then either


Figure 2.16: The action of a particular pseudo-Anosov $\psi: S_{0,5}^{1} \rightarrow S_{0,5}^{1}$ on an invariant train track. Each of the loop edges contains a puncture.


Figure 2.17: The action of $\psi$ after naïvely splitting $e_{2}$ over $e_{3}$. On the left is the new track $\tau_{1}$, and on the right is the image $\psi\left(\tau_{1}\right)$, up to isotopy. The only difference between the righthand images of this figure and Figure 2.16 is that the image of of $e_{2}$ has been peeled back along that of $e_{3}$ and now starts at the leftmost loop. As we can see, $\tau_{1}$ is not invariant under $\psi$.


Figure 2.18: The action of $\psi$ after more carefully splitting $\tau$. On the left is the train track $\tau_{2}$, and on the right is $\psi\left(\tau_{2}\right)$. We obtain a new train track that is still invariant under $\psi$. This is an example of a "tight split."

1. $e_{i+1}, \ldots, e_{j-1} \in C$, or
2. $e_{j+1}, \ldots, e_{k}, e_{1}, \ldots, e_{i-1} \in C$.

The collections

$$
R(v)=\{\text { real edges in } \operatorname{Lk}(v)\}, \quad I(v)=\{\text { infinitesimal edges in } \operatorname{Lk}(v)\}
$$

are connected. We index the elements of $\operatorname{Lk}(v)$ so that the real edges are $e_{1}, \ldots, e_{m}$ under the cyclic order. In other words, from the perspective of $v$ facing its real edges, $e_{1}$ is the real edge furthest to the right and $e_{m}$ is the edge furthest to the left.

Definition 2.5.6. The right extremal edge of $v$ is $r(v)=e_{1}$, and the left extremal edge is $l(v)=e_{m}$. If $R(v)=\{e\}$ is a singleton, then we set $e=l(v)=r(v)$.

If $v$ is a switch at an infinitesimal loop of $\tau$, we treat each end of the loop as a distinct element of $\operatorname{Lk}(v)$. Hence $I(v)$ always consists of two elements, $i_{l}$ and $i_{r}$. These are defined so that, under the cyclic order, we have

$$
l(v)<i_{l}<i_{r}<r(v)
$$

Definition 2.5.7. We denote by $v_{l}$ the switch of $\tau$ at the other end of $i_{l}$ from $v$. Similarly, we denote by $v_{r}$ the switch of $\tau$ at the other end of $i_{r}$ from $v$. In the case that $v$ is at a loop of $\tau$, we set $v_{l}=v_{r}=v$.

From now on, we set the convention that, for a given switch $v$ of $\tau$, all edges in $R(v)$ are oriented into $v$ as paths.


Figure 2.19: Left: part of a train track $\tau$ and the image of a pseudo-Anosov $\psi$ carried by $\tau$. Here $\psi$ induces a train track map $f: \tau \rightarrow \tau$ for which $v$ splits tightly to the right. Right: The train track $\tau_{1}$ after $r$-splitting $v$, and the action of $\psi$ on $\tau_{1}$. Note in particular that $\psi$ has not changed, only $\tau$ and its fibered neighborhood $N(\tau)$. In each subfigure, the highlighted regions are collapsed by a deformation retraction onto the corresponding edges.

Definition 2.5.8. Let $\tau \hookrightarrow S_{0, n}^{1}$ be a standardly embedded train track. Let $v$ be a switch of $\tau$. Fix a train track map $f: \tau \rightarrow \tau$. We say that $v$ splits tightly to the left or l-splits if for every real edge $x \subseteq \tau$ the following two conditions hold:

1. Whenever $l(v)$ appears in the train path $f(x)$, it is followed by $\overline{r\left(v_{l}\right)}$, and
2. whenever $\overline{l(v)}$ appears in the train path $f(x)$, it is preceded by $r\left(v_{l}\right)$.

Similarly, we say that $v$ splits tightly to the right or $r$-splits if for every real edge $x \subseteq \tau$ the following two conditions hold:

1. Whenever $r(v)$ appears in the train path $f(x)$, it is followed by $\overline{l\left(v_{r}\right)}$, and
2. whenever $\overline{r(v)}$ appears in the train path $f(x)$, it is preceded by $l\left(v_{r}\right)$.

In either case, we say that $v$ splits tightly. See Figures 2.19 and 2.20.
If $v$ splits tightly, we define a new train track that maps to $\tau$ by an elementary folding map. In this way, we view splitting as an inverse operation to folding. In what follows we will restrict our attention to the case that $v$ tightly splits to the left: all definitions are analogous if $v$ splits tightly to the right. To obtain these analogous statements and proofs, one need only replace all l's with $r$ 's and vice versa.


Figure 2.20: Left: another train track $\tau$ and map $f: \tau \rightarrow \tau$ for which $v$ splits tightly to the right. Right: the train track $\tau_{1}$ after $r$-splitting $v$.

Suppose $v l$-splits. Define $\tau_{v}^{l}$ to be the standardly embedded train track obtained by deleting $l(v)$ and replacing it with a real edge $\alpha$ such that

1. As a directed edge, $\alpha(0)=l(v)(0)$ and $\alpha(1)=r\left(v_{l}\right)(1)$.
2. The edge $\alpha$ forms a bigon (i.e. a two-cusped disk) with the train path $l(v) \cdot \overline{r\left(v_{l}\right)}$, and there is an isotopy rel the punctures of $S_{0, n}^{1}$ so that $\alpha$ lies transverse to the leaves of the fibered neighborhood of $\tau$.

The standardly embedded track $\tau_{v}^{l}$ comes equipped with a natural elementary folding map $p: \tau_{v}^{l} \rightarrow \tau$, defined by

$$
p(x)= \begin{cases}x & \text { if } x \neq \alpha \\ l(v) \cdot \overline{r\left(v_{l}\right)} & \text { if } x=\alpha\end{cases}
$$

Definition 2.5.9. If $v$ splits tightly to the left, then the map $p: \tau_{v}^{l} \rightarrow \tau$ is called a tight left split or an l-split of $\tau$. We analogously define the tight right split or $r$-split $p: \tau_{v}^{r} \rightarrow \tau$.

Proposition 2.5.10. Suppose $(\tau, \psi, f)$ is the data of a pseudo-Anosov carried by the standardly embedded train track $\tau$ :


If $v l$-splits, then $\tau_{v}^{l}$ carries $\psi$. Hence the above diagram may be completed to the commutative diagram

where $f_{v}^{l}$ is a train track map.
Proof. Let $F \subseteq S_{0, n}^{1}$ be a fibered surface for $\psi$ from which the Bestvina-Handel algorithm produces $\tau$. Let $L, I$, and $R$ denote the strips of $F$ collapsing to the (unoriented) edges $l(v), i_{l}$, and $r\left(v_{l}\right)$ of $\tau$, respectively. Deleting $L$ and replacing it with a strip $A$ collapsing to $\alpha$ produces a new fibered surface $F^{\prime}$ from which the algorithm produces $\tau_{v}^{l}$. The fact that $F^{\prime}$ is a fibered surface for $\psi$ follows from the fact that $v l$-splits: any strip of $\psi(F)$ passing through $L$ in fact passes through all three of $L, I$, and $R$ in order, and hence after an isotopy we may arrange for the strip to pass through $A$ instead. Furthermore, since $\alpha$ is isotopic to $l(v) \cdot i_{l} \cdot \overline{r\left(l_{v}\right)}$ and $\psi(L), \psi(I)$, and $\psi(R)$ may be isotoped into $F^{\prime}$, it follows that $\psi(A)$ may be isotoped into $F^{\prime}$ as well.

Proposition 2.5.11. Suppose that $v l$-splits and let $M$ and $M_{v}$ be the transition matrices of $f: \tau \rightarrow \tau$ and $f_{v}^{l}: \tau_{v}^{l} \rightarrow \tau_{v}^{l}$, respectively. Then

$$
M_{v}=P^{-1} M P,
$$

where $P$ is the transition matrix of the elementary folding map $p: \tau_{v}^{l} \rightarrow \tau$ : that is, if $l(v)$ is the $j$ th edge and $r\left(v_{l}\right)$ is the $i$ th edge, then we have

$$
P=I_{n}+D_{i, j},
$$

where $\tau$ has $n$ real edges, $I_{n}$ is the identity, and $D_{i, j}$ is the square matrix with a 1 in the ( $i, j$ )-entry and 0 's elsewhere.

Proof. We will argue that we have the following commutative diagram:


From this the claim will follow, since each of the arrows is a Markov map, and so upon passing to transition matrices we obtain the relation

$$
P M_{v}=M P
$$

Suppose $x$ is an edge of $\tau_{v}^{l}$. By the definition of $p$ we have

$$
(f \circ p)(x)= \begin{cases}f(x) & \text { if } x \neq \alpha \\ f(l(v)) \cdot f\left(\overline{r\left(v_{l}\right)}\right) & \text { if } x=\alpha\end{cases}
$$

On the other hand, we must understand the map $f_{v}^{l}: \tau_{v}^{l} \rightarrow \tau_{v}^{l}$ in order to analyze the composition $p \circ f_{v}^{l}$. For any edge $y \in \tau$, define $f^{\prime}(y)$ to be the word obtained from the train path $f(y)$ by replacing each instance of $l(v) \cdot \overline{r\left(v_{l}\right)}$ with $\alpha$ and each instance of $r\left(v_{l}\right) \cdot l(v)$ with $\bar{\alpha}$. In other words, $f^{\prime}(x)$ is the unique word such that

$$
p\left(f^{\prime}(x)\right)=f(x)
$$

If $x \neq \alpha$ is an edge of $\tau_{v}^{l}$, then $f_{v}^{l}(x)=f^{\prime}(x)$. If $x=\alpha$, then $f_{v}^{l}(x)=f_{v}^{l}(\alpha)=$ $f^{\prime}(l(v)) \cdot f^{\prime}\left(\overline{r\left(v_{l}\right)}\right)$. In either case, we obtain the formula

$$
\left(p \circ f_{v}^{l}\right)(x)= \begin{cases}f(x) & \text { if } x \neq \alpha \\ f(l(v)) \cdot f\left(\overline{r\left(v_{l}\right)}\right) & \text { if } x=\alpha\end{cases}
$$

This agrees with the formula for $f \circ p$, so the proof is complete.
Recall that by the Perron-Frobenius theorem, the dilatation of $\psi$ is a simple eigenvalue of the transition matrix $M$, and there exists a positive right $\lambda$-eigenvector $\mu$ of $M$. For a fixed choice of $\mu$ we will denote by $\mu(x)$ the entry of $\mu$ corresponding to the real edge $x$.

Corollary 2.5.12. Let $(\tau, \psi, f)$ be the data of a pseudo-Anosov carried by a standardly embedded train track. Let $M$ be the transition matrix for $f: \tau \rightarrow \tau$, and let $\lambda$ be the dilatation of $f$. Fix a positive right $\lambda$-eigenvector $\mu$ of $M$. If $v l$-splits then $\mu_{v}=P^{-1} \mu$ is a positive right $\lambda$-eigenvector of $M_{v}$. Consequently,

$$
\mu(l(v))<\mu\left(r\left(v_{l}\right)\right) .
$$

Proof. Since $M_{v}=P^{-1} M P$, it immediately follows that $\mu_{v}=P^{-1} \mu$ is a right $\lambda$ eigenvector of $M_{v}$. At least one entry of $\mu_{v}$ is positive, since $\mu_{v}(\alpha)=\mu(l(v))>0$. Therefore $\mu_{v}$ is positive, since the Perron-Frobenius theorem states that $\lambda$ is a simple eigenvalue of $M_{v}$ and has a positive eigenvector.

To see that $\mu(l(v))<\mu\left(r\left(v_{l}\right)\right)$, observe that

$$
0<\mu_{v}\left(r\left(v_{l}\right)\right)=\mu\left(r\left(v_{l}\right)\right)-\mu(l(v)) .
$$

Example 2.5.13. Here is an extended example of a sequence of tight splittings. The maps appearing in this example are closely related to the maps studied in Section 2.4. Let $(\tau, \psi, f)$ be the data of the pseudo-Anosov represented in Figure 2.21. The transition matrix for $f: \tau \rightarrow \tau$ is

$$
M_{1}=\left(\begin{array}{lllll}
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 & 1 \\
1 & 2 & 0 & 0 & 0 \\
1 & 1 & 0 & 0 & 0
\end{array}\right) .
$$

The characteristic polynomial of $M_{1}$ is $\chi(t)=(t+1)\left(t^{4}-t^{3}-t^{2}-t+1\right)$. The dilatation of $\psi$ is the root $\lambda$ of this polynomial with largest absolute value, as in subsection 2.2. A positive right $\lambda$-eigenvector for $M_{1}$ is

$$
\mu_{1}=\left(\begin{array}{l}
\mu_{1}\left(e_{1}\right) \\
\mu_{1}\left(e_{2}\right) \\
\mu_{1}\left(e_{3}\right) \\
\mu_{1}\left(e_{4}\right) \\
\mu_{1}\left(e_{5}\right)
\end{array}\right)=\left(\begin{array}{c}
2+5 \lambda-\lambda^{2}-\lambda^{3} \\
-2-2 \lambda+\lambda^{2}+\lambda^{3} \\
1+\lambda+4 \lambda^{2}-2 \lambda^{3} \\
-1-\lambda-\lambda^{2}+2 \lambda^{3} \\
3
\end{array}\right)=\left(\begin{array}{c}
2.537 \ldots \\
2.628 \ldots \\
4.370 \ldots \\
4.526 \ldots \\
3
\end{array}\right)
$$

One can see that the vertex at loop 5 splits tightly to the left. Performing this $l$-split produces the track $\tau_{2}$, which also carries $\psi$. See Figure 2.21 . The transition matrix of the $l$-split $p_{1}: \tau_{2} \rightarrow \tau_{1}$ is

$$
P_{1}=\left(\begin{array}{lllll}
1 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 1 \\
0 & 0 & 0 & 0 & 1
\end{array}\right)=I_{5}+D_{4,5}
$$

and the transition matrix for $f_{2}: \tau_{2} \rightarrow \tau_{2}$ is


Figure 2.21: The track $\tau_{1}, \tau_{2}$ carries $\psi$. The track $\tau_{2}^{\prime}=\sigma_{4}^{-1}\left(\tau_{2}\right)$ carries $\sigma_{4}^{-1} \circ \psi \circ \sigma_{4}$. The track $\tau_{3}$ carries $\sigma_{4}^{-1} \circ \psi \circ \sigma_{4}$.

$$
M_{2}=P_{1}^{-1} M_{1} P_{1}=\left(\begin{array}{ccccc}
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 1 \\
0 & 0 & 0 & 1 & 2 \\
0 & 1 & 0 & 0 & 0 \\
1 & 1 & 0 & 0 & 0
\end{array}\right)
$$

which has right $\lambda$-eigenvector

$$
\mu_{2}=P_{1}^{-1} \mu_{1}=\left(\begin{array}{l}
\mu_{2}\left(e_{1}\right) \\
\mu_{2}\left(e_{2}\right) \\
\mu_{2}\left(e_{3}\right) \\
\mu_{2}\left(e_{4}\right) \\
\mu_{2}\left(e_{5}\right)
\end{array}\right)=\left(\begin{array}{c}
\mu_{1}\left(e_{1}\right) \\
\mu_{1}\left(e_{2}\right) \\
\mu_{1}\left(e_{3}\right) \\
\mu_{1}\left(e_{4}\right)-\mu_{1}\left(e_{5}\right) \\
\mu_{1}\left(e_{5}\right)
\end{array}\right)=\left(\begin{array}{c}
2.537 \ldots \\
2.628 \ldots \\
4.370 \ldots \\
1.526 \ldots \\
3
\end{array}\right)
$$

We may conjugate by $\sigma_{4}^{-1}$ to obtain the track $\tau_{2}^{\prime}$, which is slightly easier to read. See Figure 2.21. This move is a standardizing braid move in the language of [71]. It is not a tight splitting and is purely cosmetic. It does not alter the transition matrix or any other relevant dynamical information.

We can now see that the switch at loop 4 splits tightly to the right. Performing this $r$-split produces the track $\tau_{3}$, which also carries $\sigma_{4}^{-1} \circ \psi \circ \sigma_{4}$. See Figure 2.21. The transition matrix of the $r$-split $p_{2}: \tau_{3} \rightarrow \tau_{2}$ is

$$
P_{2}=\left(\begin{array}{lllll}
1 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 & 1
\end{array}\right)=I_{5}+D_{5,4}
$$

and the transition matrix for $f_{3}: \tau_{3} \rightarrow \tau_{3}$ is

$$
M_{3}=P_{2}^{-1} M_{2} P_{2}=\left(\begin{array}{ccccc}
0 & 0 & 1 & 0 & 0 \\
0 & 0 & 0 & 2 & 1 \\
0 & 0 & 0 & 3 & 2 \\
0 & 1 & 0 & 0 & 0 \\
1 & 0 & 0 & 0 & 0
\end{array}\right)
$$

which has right $\lambda$-eigenvector

$$
\mu_{3}=P_{2}^{-1} \mu_{2}=\left(\begin{array}{c}
\mu_{2}\left(e_{1}\right) \\
\mu_{2}\left(e_{2}\right) \\
\mu_{2}\left(e_{3}\right) \\
\mu_{2}\left(e_{4}\right) \\
\mu_{2}\left(e_{5}\right)-\mu_{2}\left(e_{4}\right)
\end{array}\right)=\left(\begin{array}{c}
\mu_{1}\left(e_{1}\right) \\
\mu_{1}\left(e_{2}\right) \\
\mu_{1}\left(e_{3}\right) \\
\mu_{1}\left(e_{4}\right)-\mu_{1}\left(e_{5}\right) \\
2 \mu_{1}\left(e_{4}\right)-\mu_{1}\left(e_{4}\right)
\end{array}\right)=\left(\begin{array}{c}
2.537 \ldots \\
2.638 \ldots \\
4.370 \ldots \\
1.526 \ldots \\
1.473 \ldots
\end{array}\right)
$$



Figure 2.22: An example of a rigid switch. On the left is the switch, on the right the image of the map near the switch.

## Switch rigidity

In this section we investigate when a tight splitting is possible at a given switch, identifying the essential obstruction. We call this obstruction switch rigidity and show that it is uncommon. Indeed, the orbit of every switch contains a switch that is tightly splittable (cf. Proposition 2.5.20).

Let $v$ be a switch of the train track $\tau$. Recall that $\operatorname{Lk}(v)$ is the set of edges of $\tau$ incident to $v$. A Markov map $f: \tau \rightarrow \tau$ induces a map $D f: \operatorname{Lk}(v) \rightarrow \operatorname{Lk}(f(v))$ as follows. Orient all edges in $\operatorname{Lk}(v)$ and $\operatorname{Lk}(f(v))$ away from $v$ and $f(v)$, respectively. Then

$$
D f(a)=b \text { if } f(a) \text { begins with } b
$$

As a consequence of the Bestvina-Handel algorithm, all elements of $R(v)$ belong to the same gate: that is, there exists an integer $k \geq 1$ such that $(D f)^{k}=D\left(f^{k}\right)$ is constant on $R(v)$.

Definition 2.5.14. Let $\tau \hookrightarrow S_{0, n}^{1}$ be standardly embedded, and let $f: \tau \rightarrow \tau$ be a train track map. Let $v$ be a switch of $\tau$ such that $R(v)$ is not a singleton, and set $R(v)=\left\{e_{1}, \ldots, e_{k}\right\}$. Let $w$ be the switch of $\tau$ such that $f(w)=v$. We say that $v$ is rigid if there exist $x_{1}, \ldots, x_{k} \in R(w)$ such that

$$
D f\left(x_{i}\right)=e_{i} \text { for all } i .
$$

Lemma 2.5.15. Let $(\tau, \psi, f)$ be the data of the pseudo-Anosov $\psi$ on $S_{0, n}^{1}$ carried by the standardly embedded $\tau$. Let $w$ be a switch of $\tau$. Write $\alpha=r(w), \beta=l(w)$, and $v=f(w)$. For any $c \in R(v)$ between $D f(\alpha)$ and $D f(\beta)$, there exists $\gamma \in R(w)$ such that $D f(\gamma)=c$. In other words, the set $D f(\operatorname{Lk}(w)) \subseteq \operatorname{Lk}(v)$ is connected.

Proof. Suppose $c \in R(v)$ is between $D f(\alpha)$ and $D f(\beta)$. Since $\psi$ is pseudo-Anosov, $f$ is surjective. Hence there exists a real edge $\gamma$ such that $f(\gamma)$ collapses onto $c$. But

## CHAPTER 2. FIXED-POINT-FREE PSEUDO-ANOSOVS, KNOT FLOER

 HOMOLOGY AND THE CINQUEFOILsince $\psi$ is a homeomorphism, $\psi(\gamma)$ cannot intersect $\psi(\alpha \cup \beta)$, so $\gamma$ must be incident to $w$. In other words, $c=D f(\gamma)$.

Definition 2.5.16. We say a switch $v$ of $\tau$ is a loop switch if it is incident to an infinitesimal loop.

The next lemma says that switch rigidity is the only barrier to the existence of a tight splitting at a loop switch. Note that if $v$ is a loop switch, then $v_{l}=v_{r}=v$.

Lemma 2.5.17. Let $(\tau, \psi, f)$ be the data of a pseudo-Anosov $\psi$ on $S_{0, n}^{1}$ carried by the standardly embedded $\tau$. Let $v$ be a loop switch, and suppose that $R(v)$ is not a singleton. Then exactly one of the following three possibilities is true.

1. The switch $v$ splits tightly to the left.
2. The switch $v$ splits tightly to the right.
3. The switch $v$ is rigid.

Proof. Let $w$ be the loop switch of $\tau$ such that $f(w)=v$. If either (1) or (2) holds then $v$ cannot be rigid: for example, if $v l$-splits then there does not exist $x \in R(w)$ such that $D f(x)=l(v)$. On the other hand, if $v$ is not rigid then Lemma 2.5.15 implies that at least one of $l(v), r(v)$ is not in the image $D f(\operatorname{Lk}(w))$.

Assume without loss of generality that $l(v) \notin D f(\operatorname{Lk}(w))$. Then any appearance of $l(v)$ in an image train path is in fact an appearance of $l(v) \cdot \bar{x}$, up to orientation. Here $x$ is some edge in $R(v)$ that might vary. If $x$ is always $r(v)$ then $v l$-splits. Otherwise, we claim that $v r$-splits.

Indeed, suppose that there exists a real edge $y \subseteq \tau$ such that $f(y)$ contains $l(v) \cdot \bar{x}$, up to orientation, for some real edge $x \neq r(v)$. Lemma 2.5.15 implies that $D f(\operatorname{Lk}(w))$ is a subset of the real edges between $l(v)$ and $x$. In particular, $r(v) \notin D f(\operatorname{Lk}(w))$. Let $z$ be a real edge such that $f(z)$ contains $r(v)$, up to orientation. Since $\psi$ is a homeomorphism and $f(z)$ is a train path, the appearance of $r(v)$ in $f(z)$ must be followed by $\overline{l(v)}$, due to the existence of $\psi(y)$. In other words, $v r$-splits.

Thus we have established that (1) or (2) holds if and only if (3) does not hold. It remains to show that (1) and (2) are mutually exclusive. Corollary 2.5 .12 says that if $v l$-splits then $\mu(l(v))<\mu(r(v))$. It follows that if (1) holds then (2) cannot. The proof is complete.

The same argument gives the following proposition for a switch not at a loop.

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 HOMOLOGY AND THE CINQUEFOILProposition 2.5.18. Let $(\tau, \psi, f)$ be the data of a pseudo-Anosov $\psi$ on $S_{0, n}^{1}$ carried by the standardly embedded $\tau$. Let $v$ be a switch of $\tau$, and suppose that $R(v)$ is not a singleton. Suppose additionally that $R\left(v_{l}\right)$ and $R\left(v_{r}\right)$ are singletons. Then at least one of the following three possibilities is true.

1. The switch $v$ splits tightly to the left.
2. The switch $v$ splits tightly to the right.
3. The switch $v$ is rigid.

Moreover, case (3) is disjoint from cases (1) and (2).
Lemma 2.5.17 says that if we cannot split at a particular switch $v$, then it is rigid. The natural next step is to consider the preimage switch $v_{1}$ causing $v$ to be rigid. If $v_{1}$ is also rigid, we look at its preimage $v_{2}$. It might happen that we never find a splittable switch. In this case, the periodic orbit of $v$ consists of a cycle of rigid switches.

Definition 2.5.19. A rigid cycle of length $k$ is a collection of rigid switches $v_{1}, \ldots, v_{k} \in$ $\tau$ such that $f\left(v_{j}\right)=v_{j-1}$ for all $j$, where the indices are taken modulo $k$.

Proposition 2.5.20. Rigid cycles do not exist.
Proof. Let $v \in \tau$ be a switch. Since $\tau$ is standardly embedded, every element of $R(v)$ belongs to the same gate of $v$, hence there exists $k \geq 1$ such that $(D f)^{k}$ is constant on $R(v)$. In fact, for all $n \geq k$ we have that $(D f)^{n}$ is constant on $R(v)$. On the other hand, if $v$ belonged to a rigid cycle of length $n$ then $(D f)^{n}: R(v) \rightarrow R(v)$ would be the identity map, a contradiction.

Corollary 2.5.21. Let $v \in \tau$ be a switch such that $R(v)$ is not a singleton. Then some iterated preimage switch $w$ of $v$ is not rigid.

It is well-known that if $(\tau, \psi, f)$ is the data of a pseudo-Anosov, then $f$ permutes the infinitesimal $k$-gons for each $k$ (cf. [7]). We obtain the following corollary, which will be of central importance in the following section. The real valence of a switch $v$ is the cardinality of $R(v)$.

Corollary 2.5.22. Let $n_{k}$ denote the maximal real valence of a switch at an infinitesimal $k$-gon of $\tau$, where $k \geq 1$. If $n_{k}>1$ then there exists a switch of valence $n_{k}$ at such a $k$-gon which is not rigid.

## CHAPTER 2. FIXED-POINT-FREE PSEUDO-ANOSOVS, KNOT FLOER

 HOMOLOGY AND THE CINQUEFOILProof. The infinitesimal $k$-gons are permuted by $f$. If every such maximal valence switch is rigid, then they must form a rigid cycle, since real valence cannot decrease when passing to the preimage of a rigid switch. This is impossible, since rigid cycles do not exist.

## The proofs of Theorems 2.4.2, C, and D

In this subsection, we will use the theory of tight splitting developed above to prove Theorems D, and see 2.4.2 as a consequence. Though Theorem D itself is more general than necessary to prove Theorem 2.4.2, we believe it has wider-reaching applications to surface dynamics.

Definition 2.5.23. Let $\tau \hookrightarrow S_{0, n}^{1}$ be a standardly embedded train track. We say a real edge $e$ of $\tau$ is a stem if at least one end of $e$ is incident to an infinitesimal $k$-gon, where $k \geq 2$.

Definition 2.5.24. Let $\tau \hookrightarrow S_{0, n}^{1}$ be a standardly embedded train track. We say a loop switch $v \in \tau$ is a joint if $|R(v)| \geq 2$.

Theorem D. Let $\psi$ be a pseudo-Anosov on $S_{0, n}^{1}$ with at least one $k$-pronged singularity away from the boundary with $k \geq 2$. Then $\psi$ is carried by a train track $\tau$ with no joints.

The central argument in the proof of Theorem D hinges on finding a maximalvalence vertex $v$ near a puncture, and then using Corollary 2.5.22 to tightly split at $v$. Before diving into the proof, we observe one crucial lemma. Although wellknown to experts, the authors could not find a complete proof of Lemma 2.5.25 in the literature. For the sake of completeness, we have included a proof which arose from a helpful conversation with Karl Winsor.

Lemma 2.5.25. For any fixed $n$ and $B>0$, there is a finite number of PerronFrobenius matrices of size $n$ and spectral radius at most $B$. In particular, there is a finite number of Perron-Frobenius matrices of a given size with a particular Perron-Frobenius eigenvalue.

Proof. Fix $n \geq 2$, and let $M$ be an $n \times n$ Perron-Frobenius matrix. Write $M_{i, j}$ for the $(i, j)$ th entry of $M$, and $C_{j}(M)$ for the $j$ th column of $M$. An exercise in matrix algebra shows that for each integer $k \geq 1$,

$$
C_{j}\left(M^{k}\right)=\sum_{i=1}^{n}\left(M^{k-1}\right)_{i, j} \cdot C_{i}(M)
$$

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It is well-known (cf. [124]) that $M^{n^{2}-2 n+2}$ has all entries positive. Hence the smallest column sum of $M^{n^{2}-2 n+3}$ is at least the sum $\|M\|_{1}$ of all the entries of $M$. It is not hard to see that the smallest column sum of a Perron-Frobenius matrix is a lower bound on its spectral radius $\rho(M)$. We now have

$$
\rho(M)^{n^{2}-2 n+3}=\rho\left(M^{n^{2}-2 n+3}\right) \geq\|M\|_{1} .
$$

In particular, $\rho(M) \geq\|M\|_{1}^{\frac{1}{n^{2}-2 n+3}}$. Since there are only finitely many integervalued matrices $M$ with $\|M\|_{1}$ below a given bound, the result follows.

Proof of Theorem D. Let $\tau_{0} \hookrightarrow S_{0, n}^{1}$ be a standardly embedded train track carrying $\psi$. We will algorithmically perform a finite sequence of tight splittings on $\tau_{0}$ to produce the desired track $\tau$ with no joints.

Let $J$ denote the number of cusps at the loop switches of $\tau$, i.e. $J=\sum_{v}(|R(v)|-$ 1), where $v$ ranges over the loop switches of $\tau$. If $J=0$ then there is nothing to prove, so assume $J \geq 1$. By Corollary 2.5 .22 there exists a loop switch of $\tau_{0}$ of maximal valence that can be tightly split. Therefore, we introduce the following simple algorithm.

1. Initialize $\tau=\tau_{0}$ and $\mathcal{M}=\left\{M_{0}\right\}$, where $M_{0}$ is the transition matrix associated to the data $\left(\tau_{0}, \psi, f_{0}\right)$.
2. Find a loop switch of $\tau$ of maximal valence that is not rigid, and split it, obtaining the data $\left(\tau_{1}, \psi, f_{1}\right)$ with transition matrix $M_{1}$. Set $\tau=\tau_{1}$.
3. If $J$ has decreased by one, return the data $\left(\tau_{1}, \psi, f_{1}\right)$.
4. If $J$ has not decreased, add $M_{1}$ to $\mathcal{M}$ and repeat Steps 2 and 3 with $\left(\tau_{1}, \psi, f_{1}\right)$.

We claim that this algorithm terminates in finitely many steps, and returns a train track $\tau$ with one fewer joint than $\tau_{0}$. First, note that splitting at a loop switch $v_{0}$ either preserves $J$ or decreases it by one. Indeed, let $b$ be the real edge that is split over, i.e. the edge whose transverse weight is reduced (cf. Corollary 2.5.12). Let $v_{b}$ denote the switch at the other end of $b$. The tight splitting transfers a cusp from the splitting switch $v_{0}$ to $v_{b}$. Thus, in the formula

$$
J=\sum_{v \text { a loop switch }}(|R(v)|-1),
$$

the contribution from $v_{0}$ decreases by one, whereas the contribution from $v_{b}$ either (1) increases by one, if $v_{b}$ is itself a loop switch; or (2) does not change, if $v_{b}$ is not a

## CHAPTER 2. FIXED-POINT-FREE PSEUDO-ANOSOVS, KNOT FLOER HOMOLOGY AND THE CINQUEFOIL

loop switch. In particular, if $b$ is a stem, then splitting over $b$ at the loop switch $v_{0}$ will always reduce $J$ by one.

It remains to show that, by repeatedly applying the above algorithm, we will eventually split over a stem. Indeed, by Lemma 2.5.25 there are only finitely many possible transition matrices that can appear, hence we will eventually produce a matrix $M_{j}=M_{i} \in \mathcal{M}$. Since this matrix is Perron-Frobenius, the dilatation $\lambda$ of $\psi$ is an eigenvalue with strictly positive eigenvectors $\mu_{i}$ and $\mu_{j}$. Moreover, $\lambda$ is simple, so in fact $\mu_{j}$ is a scalar multiple of $\mu_{i}$. According to Corollary 2.5.12, each tight splitting reduces one of the entries of this eigenvector, so recurring to a matrix in $\mathcal{M}$ implies that every entry of $\mu$ has been reduced, i.e. that every real edge of $\tau_{0}$ has been split over. In particular, the stems of $\tau_{0}$ have been split over. The preceding paragraph now implies that the algorithm must terminate in finite time.

Repeating this algorithm sufficiently many times will eventually reduce $J$ to 0 , proving the theorem.
proof of Theorem 2.4.2. Note that in the stratum $\left(2 ; 1^{5} ; 3\right)$, there are only two classes of standardly-embedded train tracks without joints: those shown in Figure 2.7. By Theorem D, any pseudo-Anosov in this stratum is conjugate to one carried by either the Peacock or the Snail. We will argue that any pseudo-Anosov $\psi$ carried by the Snail tightly splits to one carried by the Peacock.

First, observe that $\psi$ must split at the unique valence- 3 switch of the infinitesimal triangle in the Snail, by Corollary 2.5.22. Either a left or right split at this vertex yields a pseudo-Anosov $\psi^{\prime}$ conjugate to $\psi$, and carried by a track $\tau^{\prime}$ with a unique two-valent vertex $v$ at a puncture. This vertex $v$ is again splittable by Corollary 2.5.22. At $v$, note that $\psi^{\prime}$ splits either to another map carried by $\tau^{\prime}$, with strictly smaller edge weight on the edge running between two punctures, or to a map carried by the Peacock. In particular, after sufficiently many splits, $\psi^{\prime}$ splits to a pseudoAnosov carried by the Peacock.
proof of Theorem $C$. Note that if $\psi: S \rightarrow S$ has the given singularity type, we may cap-off $\psi$ to a pseudo-Anosov $\widehat{\psi}$ on the closed genus-two surface $\widehat{S}$ and extend the foliations preserved by $\psi$ along the capping disk. In this case, the 4-prong singularity $p$ in the capping disk is the unique 4 -prong singularity of $\widehat{\psi}$. In particular, $\widehat{\psi}$ commutes with the hyperelliptic involution $\iota$ on $\widehat{S}$ and $p$ is fixed by $\iota$, as in e.g. Lemma 3.7 of [3]. And, because $p$ is fixed by $\iota$, we see that $\psi$ commutes with the hyperelliptic involution on $S$, as well. We may then quotient $\psi$ to a pseudo-Anosov 5 -braid $\beta$, and from here the techniques of Section 2.5 apply. Theorem 2.4.2 implies that $\beta$ is carried by the "Peacock" train track depicted in Figure 2.7, and we can then lift this track to $S$ as described in subsection 2.2.

## Chapter 3

## A connected sum formula for embedded contact homology

Luya Wang

### 3.1 Introduction

## Embedded contact homology

Embedded contact homology (ECH) is a three-manifold invariant introduced by Hutchings in $[51,53]$. We assume for simplicity that our three-manifolds are connected. Given a contact three-manifold $(Y, \xi)$, fix a contact form $\lambda$ such that $\operatorname{ker}(\lambda)=\xi$ and all Reeb orbits associated to $\lambda$ are nondegenerate. Given $\Gamma \in H_{1}(Y)$, the ECH chain complex is a free $\mathbb{F}$-module generated by certain finite sets of Reeb orbits in the homology class $\Gamma$, where $\mathbb{F}:=\mathbb{Z} / 2 \mathbb{Z}$. The ECH differential counts pseudo-holomorphic curves with asymptotic ends at Reeb orbits and ECH index one in the symplectization of the contact three-manifold, with respect to a certain choice of $\mathbb{R}$-invariant almost complex structure $J$. It is shown in [59, 61] that $\partial^{2}=0$ for the ECH differential. Let $E C C_{*}(Y, \lambda, \Gamma, J)$ denote the ECH chain complex and let $E C H_{*}(Y, \lambda, \Gamma, J)$ denote the homology of the ECH chain complex. We often suppress the notation $\Gamma$ and understand that ECH decomposes as

$$
E C H_{*}(Y, \lambda, J)=\oplus_{\Gamma \in H_{1}(Y)} E C H_{*}(Y, \lambda, \Gamma, J) .
$$

ECH is a priori dependent on the choices of contact form $\lambda$ and almost complex structure $J$. In [109, 110, 111, 112, 113], Taubes showed an isomorphism between

## CHAPTER 3. A CONNECTED SUM FORMULA FOR EMBEDDED

 CONTACT HOMOLOGYECH and a certain version of Seiberg-Witten Floer cohomology:

$$
\begin{equation*}
E C H_{*}(Y, \lambda, \Gamma, J) \cong \widehat{H M}^{-*}(Y, \mathfrak{s}(\xi)+\operatorname{PD}(\Gamma)) \tag{3.1.1}
\end{equation*}
$$

as relatively graded $\mathbb{F}$-modules, where $\mathfrak{s}(\xi)$ is the spin-c structure determined by the 2-plane field $\xi$. For the definition of Seiberg-Witten Floer homology, see [73]. The isomorphism (3.1.1) establishes the well-definedness of ECH. In particular, ECH is independent of the choices of almost complex structures and contact forms and is sometimes denoted as $E C H_{*}(Y, \xi, \Gamma)$.

ECH is also equipped with a degree -2 chain map called the $U$ map. The chain level $U$ map is defined by counting ECH index 2 curves passing through a generic base point, which depends on the choice of the base point. The induced $U$ map on homology does not depend on such a choice and endows ECH with an $\mathbb{F}[U]$-module structure.

## The main theorem

Given two contact three-manifolds $\left(Y_{1}, \xi_{1}\right)$ and $\left(Y_{2}, \xi_{2}\right)$, one can form their contact connected sum $\left(Y_{1} \# Y_{2}, \xi_{1} \# \xi_{2}\right)$ by the Weinstein one-handle attachment [119]. Up to contactomorphism, this is a well-defined operation.

Define the derived tensor product of two chain complexes $C_{1}$ and $C_{2}$ over $\mathbb{F}[U]$ to be

$$
C_{1} \widetilde{\otimes}_{\mathbb{F}[U]} C_{2}:=H_{*}\left(C_{1} \otimes_{\mathbb{F}} C_{2} \xrightarrow{U_{1} \otimes i d+i d \otimes U_{2}} C_{1} \otimes_{\mathbb{F}} C_{2}[-1]\right),
$$

where the right hand side denotes the homology of the mapping cone of the chain map $U_{1} \otimes i d+i d \otimes U_{2}$ from $C_{1} \otimes_{\mathbb{F}} C_{2}$ to $C_{1} \otimes_{\mathbb{F}} C_{2}[-1]$, also sometimes denoted as $C o n e\left(U_{1} \otimes i d+i d \otimes U_{2}\right)$ in our paper.

We now state the main theorem:
Theorem 1. Let $\left(Y_{1}, \lambda_{1}\right)$ and $\left(Y_{2}, \lambda_{2}\right)$ be two closed connected contact three-manifolds with given nondegenerate contact forms $\lambda_{1}$ and $\lambda_{2}$. Then,

$$
\begin{equation*}
E C H\left(Y_{1} \# Y_{2}, \xi_{1} \# \xi_{2}, \Gamma_{1}+\Gamma_{2}\right) \cong\left(E C H\left(Y_{1}, \xi_{1}, \Gamma_{1}\right) \widetilde{\otimes}_{\mathbb{F}[U]} E C H\left(Y_{2}, \xi_{2}, \Gamma_{2}\right)\right. \tag{3.1.2}
\end{equation*}
$$

as $\mathbb{F}$-modules.
Remark 2. Similar theorems to our main theorem are known for Seiberg-Witten Floer homology [9, 79, 84] and Heegaard Floer homology [100, 99]. See [109, 110, 111, $112,113,19,24,26,22]$ for the isomorphisms between ECH and Seiberg-Witten Floer and between ECH and Heegaard Floer. See also [76, 77, 78, 75, 79]. Our motivations to have a connected sum formula proven in ECH include potential generalizations to

## CHAPTER 3. A CONNECTED SUM FORMULA FOR EMBEDDED

 CONTACT HOMOLOGYother contact homologies and obtaining more information on quantitative invariants in ECH. More details will be discussed in Section 3.1.

Theorem 1 can also be used to compute ECH of subcritical surgery. Note that the subcritical surgery that is attaching a 1-handle to the three-manifold $Y$ itself is simply a self connected sum. Topologically, this is the same as performing a connected sum with $S^{1} \times S^{2}$. The ECH of $S^{1} \times S^{2}$ has been computed in Section 12.2.1 in [58]. One can also understand the $U$ maps on $E C H\left(S^{1} \times S^{2}\right)$ from Seiberg-Witten theory [73, Section 36]. Given a contact structure $\xi_{0}$ on $S^{1} \times S^{2}$, let $\Gamma_{0} \in H_{1}\left(S^{1} \times S^{2}\right)$ be the homology class such that $c_{1}(\xi)+2 P D\left(\Gamma_{0}\right)=0$. Then,

$$
E C H\left(S^{1} \times S^{2}, \xi_{0}, \Gamma\right)=\left\{\begin{array}{l}
\left(\mathbb{F}\left[U, U^{-1}\right] / U \mathbb{F}[U]\right) \otimes H_{*}\left(S^{1} ; \mathbb{Z}\right), \Gamma=\Gamma_{0} \\
0, \text { otherwise }
\end{array}\right.
$$

as $\mathbb{F}[U]$-modules. Now by using (3.1.2) and Künneth formula, we have:
Corollary 3. Given a closed contact three-manifold $(Y, \xi)$,

$$
E C H\left(Y \#\left(S^{1} \times S^{2}\right), \xi \# \xi_{0}, \Gamma+\Gamma_{0}\right) \cong E C H(Y, \xi, \Gamma) \otimes_{\mathbb{F}} H_{*}\left(S^{1} ; \mathbb{Z}\right)
$$

as $\mathbb{F}$-modules.

## Ideas of the proof of the main theorem

To understand ECH chain complex of the contact connected sum $\left(Y_{1} \# Y_{2}, \xi_{1} \# \xi_{2}\right)$, one needs to understand not only the contact structure but also the contact form on the connected sum. Let $\lambda_{i}$ be a nondegenerate contact form such that ker $\lambda_{i}=\xi_{i}$ for $i=1,2$. Given two Darboux charts in $Y_{1}$ and $Y_{2}$, a particular model of the contact connected sum $\left(Y_{1} \# Y_{2}, \lambda_{1} \# \lambda_{2}\right)$ is carefully described in [39] and depends on various choices. In particular, the connected sum sphere $S_{+}$which is contained in the ascending manifold of the Weinstein one-handle contains a hyperbolic Reeb orbit which is the equator of $S_{+}$. This hyperbolic is denoted as the special hyperbolic orbit $h$ throughout our paper. In addition, one may adjust the size of the Weinstein one-handle in order to control the radius of $S_{+}$. Let $\left(Y_{1} \#_{R} Y_{2}, \lambda_{1} \#_{R} \lambda_{2}\right)$ denote the resulting contact connected sum with a connected sum sphere of radius $R$. More details about the contact connected sum operation will be given in Section 3.3.

The proof of Theorem 1 goes through a chain level statement on the filtered ECH chain complex. Recall that there is a filtration on ECH chain complex by the symplectic action functional integrating the contact form over Reeb orbits. The filtered ECH chain complex $E C C^{L}(Y, \lambda, \Gamma, J)$ is generated by orbit sets up to symplectic

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action $L$. To ease notations, we suppress $\Gamma$ and $J$ from input of ECH chain complex when it is clear.

Let

$$
C_{o}:=E C C\left(Y_{1}, \lambda_{1}\right) \otimes_{\mathbb{F}} E C C\left(Y_{2}, \lambda_{2}\right)
$$

denote all orbit sets in $Y_{1} \sqcup Y_{2}$ counted as ECH generators. Now note that up to action $L$, we may ignore orbits that cross the connected sum region by shrinking the connected sum sphere to be of a small enough radius $R(L)$, since these orbits would have actions greater or equal to $L$ by a compactness argument as in Lemma 41. Therefore, there is an obvious identification on the vector space level between the filtered ECH complex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \lambda_{1} \#_{R(L)} \lambda_{2}\right)$ and the filtered mapping cone complex

$$
\operatorname{Cone}^{L}\left(U_{1} \otimes i d+i d \otimes U_{2}\right):=C_{o}^{L} \oplus C_{h}^{L},
$$

where $C_{h} \cong C_{o}$ by appending the special hyperbolic orbit $h$ and the superscript $L$ denotes the filtration. In fact, one may find a chain homotopy equivalence that is triangular with respect to this obvious identification of the vector spaces:

Proposition 4. Given two closed connected contact three-manifolds ( $Y_{1}, \lambda_{1}$ ) and $\left(Y_{2}, \lambda_{2}\right)$ with nondegenerate contact forms $\lambda_{i}$, there exists a strictly decreasing function $R: \mathbb{R} \rightarrow \mathbb{R}$ with

$$
\lim _{L \rightarrow \infty} R(L)=0
$$

such that there is a chain homotopy equivalence

$$
f: E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \lambda_{1} \#_{R(L)} \lambda_{2}\right) \longrightarrow \operatorname{Cone}^{L}\left(U_{1} \otimes i d+i d \otimes U_{2}\right)
$$

More details about the identification on the vector space level will be provided in Section 3.4. The chain homotopy equivalence in Proposition 4 will be constructed in Section 3.6.

## Further directions

From the perspective of Symplectic Field Theory (SFT) constructed by Eliashberg, Givental and Hofer [35], one could also try to prove a connected sum formula in linearized contact homology. This is studied by Bourgeois and van-Koert [13]. We hope that our connected sum formula for embedded contact homology could give ideas to such formulas for other contact homologies that involve higher genus curves and more asymptotic ends.

It would also be interesting to study the ECH formulas and cobordism maps under additional contact surgeries. One natural candidate that is suitable for the

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Theorem 1 computes ECH of a connected sum as an $\mathbb{F}$-module. However, the model of the derived tensor product on the right hand side of (3.1.2) has a natural $U$ action, $U_{1} \otimes i d$, which is homotopic to $i d \otimes U_{2}$. It is natural to ask whether Theorem 1 holds over $\mathbb{F}[U]$-modules. This amounts to studying ECH $U$ maps on the chain level on the connected sum, which is also important for studying more refined invariants in ECH, such as the ECH spectral invariants. This will be discussed in future works, but we present here a conjecture.

Suppose $\lambda$ is nondegenerate. Let $0 \neq \sigma \in \operatorname{ECH}(Y, \lambda, \Gamma)$. Define $c_{\sigma}(Y, \lambda)$ to be the infimum over $L \in \mathbb{R}$ such that $\sigma$ is in the image of the inclusion-induced map

$$
E C H^{L}(Y, \lambda, \Gamma) \longrightarrow E C H(Y, \lambda, \Gamma)
$$

Recall that there is a canonical element $c(\xi):=[\emptyset] \in E C H(Y, \lambda, 0)$ called the $E C H$ contact invariant.

Definition 5. If $(Y, \lambda)$ is a closed connected contact three-manifold with the contact invariant $c(\xi) \neq 0$ and if $k$ is a nonnegative integer, then define the $k$-th ECH spectral invariant to be

$$
c_{k}(Y, \lambda):=\inf \left\{c_{\sigma}(Y, \lambda) \mid \sigma \in E C H(Y, \lambda, 0), U^{k} \sigma=[\emptyset]\right\} .
$$

## Conjecture 6.

$$
\lim _{R \rightarrow 0} c_{k}\left(\left(Y_{1} \#_{R} Y_{2}, \lambda_{1} \#_{R} \lambda_{2}\right)\right)=\max \left\{c_{i}\left(Y_{1}, \lambda_{1}\right)+c_{j}\left(Y_{2}, \lambda_{2}\right) \mid i+j=k\right\} .
$$

Organization. Section 3.2 reviews basic definitions of embedded contact homology. Section 3.3 discusses the Reeb dynamics of the contact connected sum and the asymptotic behaviors of pseudo-holomorphic curves in its symplectization. Section 3.4 discusses how to ignore potential Reeb orbits that cross the connected sum region and identifies a filtered ECH chain complex of a connected sum with a filtered mapping cone complex associated to the ECH $U$ maps on the level of vector spaces. Section 3.5 relates some of the new ECH differentials in the connected sum to the ECH differentials in the original contact three-manifolds. Section 3.6 relates the remaining differentials to the ECH $U$ maps in the original contact three-manifolds and constructs a chain homotopy equivalence that proves Proposition 4. Section 3.7 uses a direct limit argument similar to that in [94] to prove the main result Theorem 1.

### 3.2 An overview of ECH

In this section we give a very quick overview of ECH . See $[51,55,53]$ for more details. Let $Y$ be a closed connected three-manifold with a contact form $\lambda$. Let $R$ be the Reeb vector field determined by $d \lambda(R, \cdot)=0$ and $\lambda(R)=1$. Over each closed Reeb orbit $\gamma: \mathbb{R} / T \mathbb{Z} \rightarrow Y$, where $T$ is the period of the orbit, the linearized Reeb flow determines a symplectic linear map

$$
d \psi_{T}: T_{\gamma(0)} Y \longrightarrow T_{\gamma(T)} Y
$$

If $d \psi_{T}$ does not have 1 as an eigenvalue, we say that $\gamma$ is nondegenerate. A contact form $\lambda$ is said to be nondegenerate if all Reeb orbits are nondegenerate. Note that this is a generic condition and from now on we fix $\lambda$ to be nondegenerate.

Let $\mathcal{P}$ be the set of embedded Reeb orbits of the Reeb vector field associated to $\lambda$. The ECH chain complex $\operatorname{ECC}(Y, \lambda, J)$ is generated as a $\mathbb{F}$ vector space by orbit sets $\gamma=\left\{\left(\gamma_{i}, m_{i}\right)\right\}$, such that:

- $\gamma_{i} \in \mathcal{P}$ are distinct Reeb orbits;
- $m_{i} \in \mathbb{Z}^{+}$is the covering multiplicity of the orbit $\gamma_{i}$;
- if $\gamma_{i}$ is hyperbolic, i.e. the linearized return map has real eigenvalues, then $m_{i}=1$.

We call orbit sets satisfying the above criteria admissible.
Let $\left(\mathbb{R}_{s} \times Y, d\left(e^{s} \lambda\right)\right)$ be the symplectization of $(Y, \lambda)$. Let $J$ be a $\lambda$-adapted almost complex structure. This means that $J$ is $\mathbb{R}$-invariant, $J\left(\partial_{s}\right)=R$ where $R$ is the Reeb vector field associated to $\lambda$, and $J$ sends the contact structure $\xi:=\operatorname{ker}(\lambda)$ to itself, positively rotating $\xi$ with respect to the orientation on $\xi$ given by $d \lambda$. We consider $J$-holomorphic curves $u:(\Sigma, j) \rightarrow(\mathbb{R} \times Y, J)$, where $(\Sigma, j)$ is a punctured Riemann surface, up to biholomorphisms. A positive asymptotic end of $u$ at a Reeb orbit $\gamma$ is a puncture with a neighborhood that can be given coordinates of a positive halfcylinder $(s, t) \in[0, \infty) \times(\mathbb{R} / T \mathbb{Z})$ such that $j(\partial s)=t, \lim _{s \rightarrow \infty} \pi_{\mathbb{R}}(u(s, t))=\infty$ and $\lim _{s \rightarrow \infty} \pi_{Y}(u(s, \cdot))=\gamma$. A negative asymptotic end is defined analogously, where the neighborhood of a negative puncture is identified with a negative half-cylinder $(-\infty, 0] \times(\mathbb{R} / T \mathbb{Z})$.

The ECH differential counts J-holomorphic currents of ECH index one. A Jholomorphic current is a finite set of pairs $\mathcal{C}=\left\{\left(C_{k}, d_{k}\right)\right\}$, where $C_{k}$ are distinct, connected, somewhere injective $J$-holomorphic curves in $\left(\mathbb{R} \times Y, d\left(e^{s} \lambda\right)\right)$ and $d_{k} \in \mathbb{Z}^{+}$. We say that $\mathcal{C}$ is positively asymptotic to an orbit set $\alpha=\left\{\left(\alpha_{i}, m_{i}\right)\right\}$, if $C_{k}$ has positive ends at covers of $\alpha_{i}$ with multiplicity $\operatorname{cov}\left(C_{k}\right)$ and $\sum_{k} \operatorname{cov}\left(C_{k}\right)=m_{i}$. Negative
asymptotic as currents is defined analogously. We now review ingredients of the ECH index and related properties.

## The Conley-Zehnder index

In this subsection we discuss a classical quantity that measures the "winding" of the linearized return map of a given Reeb orbit, with respect to a given trivialization.

Definition 7. Let $\tau$ be a trivialization of the contact structure $\xi$ over the Reeb orbit $\gamma$, parameterized as $\gamma: \mathbb{R} / T \mathbb{Z} \rightarrow Y$. The linearization of the Reeb flow $d \psi_{t}: T_{\gamma(0)} Y \rightarrow T_{\gamma(t)} Y$ induces a family of $2 \times 2$ symplectic matrices $\phi_{t}: \xi_{\gamma(0)} \rightarrow \xi_{\gamma(t)}$ with respect to $\tau$. The Conley-Zehnder index over $\gamma$ denoted as $C Z_{\tau}(\gamma) \in \mathbb{Z}$ is the Conley-Zehnder index of the family of symplectic matrices $\phi_{t \in[0, T]}$.

Since we assumed the contact form we started with is nondegenerate, the linearized return map $\phi_{T}(\gamma)$ does not have 1 as an eigenvalue ${ }^{1}$. If $\phi_{T}(\gamma)$ has eigenvalues in the unit circle, we call $\gamma$ elliptic. Otherwise, $\gamma$ is hyperbolic. If $\gamma$ is hyperbolic, we may choose $v \in \mathbb{R}^{2}$ to be an eigenvector of $\phi_{T}$ and the family $\left\{\phi_{t}(v)\right\}_{t \in[0, T]}$ rotates by $k \pi$, where $k \in \mathbb{Z}$. In this case,

$$
C Z_{\tau}(\gamma)=k
$$

If $\gamma$ is elliptic, we may adjust the trivialization $\tau$ so that $\phi_{t}$ is a rotation by angle $2 \pi \theta_{t}$, where $\theta_{t}$ is continuous in $t \in[0, T]$ and $\theta_{0}=0$. In this case,

$$
C Z_{\tau}(\gamma)=2\lfloor\theta\rfloor+1
$$

## The relative first Chern class

The relative first Chern class is a generalization of the usual first Chern class of a complex line bundle $\xi$ over a curve with boundaries.

Definition 8. Let $\alpha=\left\{\left(\alpha_{i}, m_{i}\right)\right\}$ and $\beta=\left\{\left(\beta_{j}, n_{j}\right)\right\}$ be orbit sets with $[\alpha]=[\beta] \in$ $H_{1}(Y)$, where $[\alpha]=\sum_{i} m_{i}\left[\alpha_{i}\right]$ and $[\beta]=\sum_{j} n_{j}\left[\beta_{j}\right]$. Then $H_{2}(Y, \alpha, \beta)$ denotes the set of relative homology classes of 2-chains $Z$ in $Y$ such that

$$
\partial Z=\sum_{i} m_{i} \alpha_{i}-\sum_{j} n_{j} \beta_{j}
$$

[^0]
## CHAPTER 3. A CONNECTED SUM FORMULA FOR EMBEDDED

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$$
c_{\tau}(Z):=c_{1}\left(\left.\xi\right|_{Z}, \tau\right) \in \mathbb{Z}
$$

counts algebraically the zeros of a generic section $\psi$ of the bundle $f^{*} \xi \rightarrow S$ transverse to the zero section, which is non-vanishing and has zero winding number with respect to $\tau$ on the boundary.

The relative first Chern class $c_{\tau}(Z)$ is well-defined and changes under the relative second homology class by

$$
\begin{equation*}
c_{\tau}(Z)-c_{\tau}\left(Z^{\prime}\right)=\left\langle c_{1}(\xi), Z-Z^{\prime}\right\rangle \tag{3.2.1}
\end{equation*}
$$

where $c_{1}(\xi) \in H^{2}(Y ; \mathbb{Z})$ is the ordinary first Chern class. Let $\tau^{\prime}:=\left(\left\{\tau_{i}^{+^{\prime}}\right\},\left\{\tau_{j}^{-^{\prime}}\right\}\right)$ be another trivialization over the positive orbit set $\alpha:=\left\{\left(a_{i}, m_{i}\right)\right\}$ and negative orbit set $\beta:=\left(\left\{b_{j}, n_{j}\right\}\right)$, then

$$
\begin{equation*}
c_{\tau}(Z)-c_{\tau^{\prime}}(Z)=\sum_{i} m_{i}\left(\tau_{i}^{+}-\tau_{i}^{+^{\prime}}\right)-\sum_{j} n_{j}\left(\tau_{j}^{-}-\tau_{j}^{-\prime^{\prime}}\right) . \tag{3.2.2}
\end{equation*}
$$

## The relative intersection pairing

The relative intersection pairing associates an intersection number to a pair of relative second homology classes. It is closely related to the behavior of the asymptotic linking number which will be defined in Section 3.3.

Definition 10. Let $\alpha=\left\{\left(\alpha_{i}, m_{i}\right)\right\}$ and $\beta=\left\{\left(\beta_{j}, n_{j}\right)\right\}$ be orbit sets with $[\alpha]=[\beta] \in$ $H_{1}(Y)$ and let $Z \in H_{2}(Y, \alpha, \beta)$. An admissible representative of $Z$ is a smooth map $f: S \rightarrow[-1,1] \times Y$, where $S$ is a compact oriented surface with boundary, such that:

- $\left.f\right|_{\partial S}$ consists of positively oriented covers of $\{1\} \times \alpha_{i}$ with total multiplicity $m_{i}$ and negatively oriented covers of $\{-1\} \times \beta_{j}$ with total multiplicity $n_{j}$;
- the composition of $f$ with the projection $[-1,1] \times Y \rightarrow Y$ represents the class Z;
- $\left.f\right|_{\dot{S}}$ is an embedding and $f$ is transverse to $\{-1,1\} \times Y$.

Now fix a trivialization $\tau$ of the contact structures $\xi$ over all Reeb orbits.

Definition 11. Let $Z \in H_{2}(Y, \alpha, \beta)$ and $Z^{\prime} \in H_{2}\left(Y, \alpha^{\prime}, \beta^{\prime}\right)$. The relative intersection pairing $Q_{\tau}\left(Z, Z^{\prime}\right): H_{2}(Y, \alpha, \beta) \times H_{2}\left(Y, \alpha^{\prime}, \beta^{\prime}\right)$ is defined by

$$
Q_{\tau}\left(Z, Z^{\prime}\right):=\#\left(\dot{S} \cap \dot{S}^{\prime}\right)-l_{\tau}\left(S, S^{\prime}\right)
$$

where $S$ and $S^{\prime}$ are admissible representatives of $Z$ and $Z^{\prime}$, and the interiors $\dot{S}$ and $\dot{S}^{\prime}$ are transverse and do not intersect near the boundary. The term $l_{\tau}\left(S, S^{\prime}\right)$ is the asymptotic linking number, which will be defined in Section 3.3.

An alternative definition for $Q_{\tau}\left(Z, Z^{\prime}\right)$ is to find representatives $S$ of $Z$ and $S^{\prime}$ of $Z^{\prime}$, called $\tau$-representatives, such that the projection of their intersections under the restricted projection maps $(1-\epsilon, 1] \times Y \rightarrow Y$ and $[-1,-1+\epsilon) \times Y \rightarrow Y$ are embeddings, and that near each orbit $\gamma_{i}$ at their ends, $S$ and $S^{\prime}$ are unions of rays emanating from points of $\gamma_{i}$, and the rays do not intersect and do not rotate with respect to the trivialization $\left.\tau\right|_{\gamma_{i}}$. One can check that these two definitions are the same (Lemma 8.5 in [51]). In fact, both definitions will be useful to us.

Both definitions of $Q_{\tau}\left(Z, Z^{\prime}\right)$ are well-defined by Lemma 2.5 and 8.5 in [51]. In particular, $Q_{\tau}$ is independent of the choice of admissible representatives of the relative second homology classes. Moreover, under change of second homology class,

$$
\begin{equation*}
Q_{\tau}\left(Z_{1}, Z^{\prime}\right)-Q_{\tau}\left(Z_{2}, Z^{\prime}\right)=\left(Z_{1}-Z_{2}\right) \cdot\left[\alpha^{\prime}\right] \tag{3.2.3}
\end{equation*}
$$

and under change of trivilizations,

$$
\begin{equation*}
Q_{\tau}\left(Z, Z^{\prime}\right)-Q_{\tau^{\prime}}\left(Z, Z^{\prime}\right)=\sum_{i} m_{i} m_{i}^{\prime}\left(\tau_{i}^{+}-\tau_{i}^{+\prime}\right)-\sum_{j} n_{j} n_{j}^{\prime}\left(\tau_{j}^{-}-\tau_{j}^{-\prime}\right) \tag{3.2.4}
\end{equation*}
$$

We call $Q_{\tau}(Z):=Q_{\tau}(Z, Z)$ the relative self-intersection pairing.
Notation 12. In the ECH literature, we often abuse notation and let $C$ denote the holomorphic curve $u:(C, j) \rightarrow(X, J)$. In addition, we often write $c_{\tau}(C):=c_{\tau}([C])$ and $Q_{\tau}(C):=Q_{\tau}([C])$.

## The Fredholm index and the ECH index

Let $\mathcal{M}\left(\gamma_{1}^{+}, \cdots, \gamma_{k}^{+} ; \gamma_{1}^{-}, \cdots, \gamma_{l}^{-} ; J\right)$ denote the moduli space of $J$-holomorphic curves with positive ends at $\gamma_{1}^{+}, \cdots, \gamma_{k}^{+}$and negative ends at $\gamma_{1}^{-}, \cdots, \gamma_{l}^{-}$.
Definition 13. The Fredholm index of a curve $C \in \mathcal{M}\left(\gamma_{1}{ }^{+}, \cdots, \gamma_{k}{ }^{+} ; \gamma_{1}{ }^{-}, \cdots, \gamma_{l}{ }^{-} ; J\right)$ is defined by

$$
\operatorname{ind}(C):=-\chi(C)+2 c_{\tau}(C)+\sum_{i=1}^{k} C Z_{\tau}\left(\gamma_{i}^{+}\right)-\sum_{j=1}^{l} C Z_{\tau}\left(\gamma_{j}^{-}\right)
$$

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Definition 14. Let $Z \in H_{2}(Y, \alpha, \beta)$, then the $E C H$ index is

$$
I(\alpha, \beta, Z):=c_{\tau}(Z)+Q_{\tau}(Z)+C Z_{\tau}^{I}(\alpha, \beta)
$$

where

$$
C Z_{\tau}^{I}(\alpha, \beta):=\sum_{i} \sum_{k=1}^{m_{i}} C Z_{\tau}\left(\alpha_{i}^{k}\right)-\sum_{j} \sum_{k=1}^{n_{j}} C Z_{\tau}\left(\beta_{j}^{k}\right) .
$$

Let $\mathcal{C} \in \mathcal{M}(\alpha, \beta)$. Then we define $I(\mathcal{C}):=I(\alpha, \beta,[\mathcal{C}])$.
ECH index has many useful properties. The following inequality relates it to the Fredholm index (See [53]).

Theorem 15. Let $\alpha$ and $\beta$ be orbit sets and $C \in \mathcal{M}(\alpha, \beta)$ be a somewhere injective curve, then

$$
\operatorname{ind}(C) \leq I(C)
$$

with equality if and only if $C$ is embedded.
The inequality in Theorem 15 has in particular the following applications.
Theorem 16 (Low-index currents). [53, Proposition 3.7] Suppose J is generic. Let $\alpha$ and $\beta$ be orbit sets and $\mathcal{C}=\left\{\left(C_{i}, d_{i}\right)\right\} \in \mathcal{M}(\alpha, \beta)$, be a not necessarily somewhere injective J-holomorphic current. Then:

- $I(\mathcal{C}) \geq 0$, with equality if and only if each $C_{i}$ is a trivial cylinder with multiplicity;
- If $I(\mathcal{C})=1$, then $\mathcal{C}=C_{0} \sqcup C_{1}$, where $I\left(C_{0}\right)=0$ and $C_{1}$ is an embedded curve with $I\left(C_{1}\right)=\operatorname{ind}\left(C_{1}\right)=1$;
- If $I(C)=2$, and $\alpha$ and $\beta$ are admissible orbit sets, then $\mathcal{C}=C_{0} \sqcup C_{2}$, where $I\left(C_{0}\right)=0$ and $C_{2}$ is an embedded curve with $I\left(C_{2}\right)=\operatorname{ind}\left(C_{2}\right)=2$.

In particular, whenever we are counting low-index currents, we may assume that the underlying curves, other than potentially branched covers of trivial cylinders, are embedded. Another useful result relating the different quantities introduced in the previous subsections is the relative adjunction formula, which often helps with ECH index computations.

Proposition 17 (Relative adjunction formula). [55, Proposition 4.9] Let $C \in \mathcal{M}(\alpha ; \beta ; J)$ be somewhere injective. Then $C$ has finitely many singularities and

$$
\begin{equation*}
c_{\tau}(C)=\chi(C)+Q_{\tau}(C)+w_{\tau}(C)-2 \delta(C) \tag{3.2.5}
\end{equation*}
$$

where $\delta(C)$ is an algebraic count of singularities and $w_{\tau}(C)$ is the asymptotic writhe to be discussed in Section 3.3.

## The ECH differential and the $U$ map

The ECH differential $\langle\partial \alpha, \beta\rangle$ is defined to be the $\bmod 2$ count of the ECH index 1 pseudo-holomorphic currents $u$ in $\mathbb{R} \times M$ from $\alpha$ to $\beta$. It is shown in [59, 61] that $\partial^{2}=0$, and in $[109,110,111,112,113]$ that the homology $\operatorname{ECH}(Y, \lambda, J)$ of this chain complex does not depend on the choice of $J$ or $\lambda$.

In addition, ECH is equipped with a degree -2 chain map given a base point $z \in Y$ not on a Reeb orbit. Given admissible orbit sets $\alpha$ and $\beta$, we define

$$
U_{z}: E C C_{*}(Y, \lambda, J) \longrightarrow E C C_{*-2}(Y, \lambda, J)
$$

where $\langle U \alpha, \beta\rangle$ is the $\bmod 2$ count of pseudo-holomorphic currents of ECH index $I=2$ from $\alpha$ to $\beta$ passing through $(0, z) \in \mathbb{R} \times Y$. Furthermore, one can show that $U_{z}$ is a chain map and that up to chain homotopy, $U_{z}$ does not depend on the choice of the base point $z$. More precisely, let $z^{\prime} \in Y$ be a different choice of the base point, then it is shown in [66] that there is a chain homotopy between $U_{z}$ and $U_{z^{\prime}}$. We will describe this further in Section 3.6.

## The filtered ECH

We can associate to a Reeb orbit $\gamma$ a symplectic action

$$
\mathcal{A}(\gamma):=\int_{\gamma} \lambda,
$$

and to each ECH generator $\alpha=\left\{\left(\alpha_{i}, m_{i}\right)\right\}$ the action

$$
\mathcal{A}(\alpha):=\sum_{i} m_{i} \mathcal{A}\left(\alpha_{i}\right)
$$

By Stokes' theorem and the definition of a symplectization-adapted almost complex structure $J$, we know that ECH differential decreases the symplectic action. Therefore, one may define the filtered version $E C C^{L}(Y, \lambda, J)$ to be the subcomplex generated by orbit sets $\alpha$ such that $\mathcal{A}(\alpha)<L$. For $L<L^{\prime}$, under the inclusion-induced map

$$
i_{J}^{L, L^{\prime}}: E C H_{*}^{L}(Y, \lambda, J) \longrightarrow E C H_{*}^{L^{\prime}}(Y, \lambda, J),
$$

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we may recover

$$
\begin{equation*}
E C H_{*}(Y, \lambda, J)=\lim _{L \rightarrow \infty} E C H_{*}^{L}(Y, \lambda, J) \tag{3.2.6}
\end{equation*}
$$

In the following theorem, we review definitions from [64] and summarize their results regarding the filtered ECH that we need.

Definition 18. Given an action $L$, a contact form $\lambda$ is L-nondegenerate if all orbits of action less than $L$ are nondegenerate, and there are no orbit sets of action exactly $L$.

Definition 19. An admissible deformation is a smooth family $\rho=\left\{\left(\lambda_{t}, L, J_{t}, r_{t}\right) \mid t \in\right.$ $[0,1]\}$ such that for each $t \in[0,1]$ :

- $\lambda_{t}$ is an $L_{t}$-nondegenerate contact form on $Y$;
- $J_{t}$ is a symplectization-admissible almost complex structure for $\lambda_{t}$;
- $r_{t}$ is a positive real number.

Definition 20. Given an action $L$, an almost complex structure is $E C H^{L}$-generic if the ECH differential is well-defined on admissible orbit sets with actions less than $L$ and satisfies $\partial^{2}=0$.

Theorem 21. [64, Theorem 1.3] Let $Y$ be a closed connected oriented three-manifold, then:

1. If $\lambda$ is an $L$-nondegenerate contact form on $Y$, then $E C H_{*}^{L}(Y, \lambda, J)$ does not depend on the choice of $E C H^{L}$-generic $J$;
2. If $L<L^{\prime}$ and $\lambda$ is $L^{\prime}$-nondegenerate, then $i_{J}^{L, L^{\prime}}$ induces a well-defined map

$$
i^{L, L^{\prime}}: E C H_{*}^{L}(Y, \lambda) \longrightarrow E C H_{*}^{L^{\prime}}(Y, \lambda)
$$

3. If $\lambda$ is a nondegenerate contact form on $Y$, then $E C H_{*}(Y, \lambda, J)$ does not depend on the choice of ECH-generic $J$, so it may be denoted by $E C H_{*}(Y, \lambda)$.

### 3.3 The connected sum for contact three-manifolds

Let $\left(Y_{1}, \lambda_{1}\right)$ and $\left(Y_{2}, \lambda_{2}\right)$ be two closed connected contact three-manifolds with specified contact forms $\lambda_{i}$. In this paper, we consider their connected sum via the 4 dimensional Weinstein 1-handle attachment as in [119]. We follow [119] and [39] for

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## Weinstein 1-handle model and the Reeb flow

Consider

$$
\mathbb{C}^{2}=\mathbb{R}^{4}=\{(x, y, z, w)\}
$$

and the standard symplectic form

$$
\omega:=d x \wedge d y+d z \wedge d w
$$

on it. Consider the Liouville vector field

$$
X=\frac{1}{2} x \partial x+\frac{1}{2} y \partial y+2 z \partial z-w \partial w
$$

Now consider the function

$$
f(x, y, z, w)=\frac{1}{4} x^{2}+\frac{1}{4} y^{2}+z^{2}-\frac{1}{2} w^{2} .
$$

The hypersurface $\{f=1\}$ is of contact-type. This is because

$$
d f(X)=\frac{1}{4} x^{2}+\frac{1}{4} y^{2}+4 z^{2}+w^{2}>0
$$

so $X$ is transverse to $\{f=1\}$. This is the contact manifold that contains the ascending sphere

$$
S_{+}:=\{w=0\} \cap\{f=1\}
$$

as in [119]. We call $S_{+}$the connected sum sphere. Similarly, we have the contact-type hypersurface $\{f=-1\}$ and the descending sphere

$$
S_{-}=\{x=y=z=0\} \cap\{f=-1\} .
$$

We will see in Lemma 24 that our resulting connected sum $Y_{1} \# Y_{2}$ contains the ascending sphere $S_{+}$. Therefore, in order to understand the Reeb dynamics of $Y_{1} \# Y_{2}$, we focus on the contact-type hypersurface $\{f=1\}$.

We may compute that the contact form

$$
\alpha=\left.i_{X} \omega\right|_{f=1}=\frac{1}{2} x d y-\frac{1}{2} y d x+2 z d w+\left.w d z\right|_{f=1}
$$

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and that $d \alpha=\omega$. We also compute the Reeb vector field

$$
R=\kappa\left(\frac{1}{2} x \partial y-\frac{1}{2} y \partial x+2 z \partial w+w \partial z\right)
$$

where $\kappa:=1 /\left(1+3 z^{2}+\frac{3}{2} w^{2}\right)$. Therefore, one can see that the equator

$$
h:=\left\{x^{2}+y^{2}=4, z=w=0\right\}
$$

of $S_{+}$is the only orbit in the ascending manifold of the 1-handle. To compute the Reeb flow at the orbit $h$, we need to solve the linear ODE

$$
\dot{\eta(t)}=A \eta(t)
$$

where

$$
A=\left(\begin{array}{cccc}
0 & \frac{1}{2} & 0 & 0 \\
-\frac{1}{2} & 0 & 0 & 0 \\
0 & 0 & 0 & 1 \\
0 & 0 & 2 & 0
\end{array}\right)
$$

and $\eta(t):=\left(\begin{array}{c}x(t) \\ y(t) \\ z(t) \\ w(t)\end{array}\right)$. We have the solution $\eta(t)=e^{A t} \eta(0)$, where

$$
\Phi(t):=e^{A t}=\left(\begin{array}{cccc}
-\cos (t / 2) & \sin (t / 2) & 0 & 0  \tag{3.3.1}\\
\sin (t / 2) & \cos (t / 2) & 0 & 0 \\
0 & 0 & \cosh (\sqrt{2} t) & \sinh (\sqrt{2} t) \\
0 & 0 & \sqrt{2} \sinh (\sqrt{2} t) & \sqrt{2} \cosh (\sqrt{2} t)
\end{array}\right)
$$

Note that the contact structure restricted to $h$ is

$$
\left.\xi\right|_{h}=\langle\partial z, \partial w\rangle .
$$

Now consider the trivialization $\tau_{0}$ over $h$ given by the second factor of $\mathbb{R}^{4}=\mathbb{R}^{2} \oplus \mathbb{R}^{2}$. Under this trivialization,

$$
C Z_{\tau_{0}}(h)=0 .
$$

This is because the flow does not rotate the contact plane $\langle\partial z, \partial w\rangle$ by looking at the lower $2 \times 2$ matrix of $\Phi(t)$.

Once we have the model for the connected sum "tube", the following proposition by Weinstein tells us how to glue two isotropic setups together via the tube that is the Weinstein 1-handle in our case. See [119, 17] for more details.

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 CONTACT HOMOLOGYDefinition 22. An isotropic setup is a quintuple ( $V, \omega, x, Y, \Lambda$ ) where $(V, \omega)$ is a symplectic manifold with Liouville vector field $x$ and $\omega=d \lambda ; Y \subset V$ is a codimension-1 hypersurface transverse to $x$; and $\Lambda \subset Y$ is a closed isotropic submanifold for the contact structure $\operatorname{ker}\left(\left.\lambda\right|_{Y}\right)$.

Proposition 23. [119, Proposition 4.2] Let $\left(V_{i}, \omega_{i}, x_{i}, Y_{i}, \Lambda_{i}\right), i \in\{1,2\}$ be two isotropic setups. Given a diffeomorphism $f: \Lambda_{1} \rightarrow \Lambda_{2}$ covered by an isomorphism $\Phi$ of their symplectic subnormal bundles $\left(T \Lambda_{i}\right)^{\omega} / T \Lambda_{i} \subset \xi_{i}$, there exists an isomorphism of isotropic setups

$$
F:\left(W_{1}, \omega_{1}, x_{1}, Y_{1}, \Lambda_{1}\right) \longrightarrow\left(W_{2}, \omega_{2}, x_{2}, Y_{2}, \Lambda_{2}\right)
$$

between neighborhoods $W_{i}$ of $\Lambda_{i}$ in $V_{i}$ inducing $f$ and $\Phi$.
Then, one may glue the descending sphere in the standard 1-handle as above to the product $[0,1] \times\left(Y_{1} \sqcup Y_{2}\right)$ along neighborhoods of the isotropic spheres in $\{1\} \times\left(Y_{1} \sqcup Y_{2}\right)$. This 1-handle attachment yields $\left(Y_{1} \# Y_{2}, \xi_{1} \# \xi_{2}\right)$ where the contact structure $\xi_{1} \# \xi_{2}$ is the same as $\xi_{1}$ and $\xi_{2}$ when restricted to the complement of the neighborhoods of the isotropic spheres (Theorem 5.1 in [119]). In fact, the proof of Theorem 5.1 in [119] shows the following.

Lemma 24. Let ( $Y_{1} \# Y_{2}, \lambda_{1} \# \lambda_{2}$ ) be obtained by attaching the Weinstein one-handle defined above along the isotropic sphere in the contact manifold ( $Y_{1} \sqcup Y_{2}, \lambda_{1} \sqcup \lambda_{2}$ ). Then, the contact form $\left(Y_{1} \# Y_{2}, \lambda_{1} \# \lambda_{2}\right)$ differs from that on $\left(Y_{1} \sqcup Y_{2}, \lambda_{1} \sqcup \lambda_{2}\right)$ only on the neighborhood of the isotropic sphere where the surgery takes place. In addition, $\lambda_{1} \# \lambda_{2}$ can be perturbed to be non-degenerate if $\lambda_{1} \sqcup \lambda_{2}$ is non-degenerate.

Proof. Consider the following isotropic setup
$(V, \omega, x, Y, \Lambda)=\left([0,1] \times\left(Y_{1} \sqcup Y_{2}\right), d\left(e^{t}\left(\lambda_{1} \sqcup \lambda_{2}\right)\right), \partial t,\{1\} \times\left(Y_{1} \sqcup Y_{2}\right), S^{0}=\{p\} \sqcup\{q\}\right)$.
where $p \in Y_{1}$ and $q \in Y_{2}$ are the two points away from orbits on which we are performing the connected sum operation. Consider the second isotropic setup given by the standard Weinstein handle described in Section 3.3

$$
\left(V^{\prime}, \omega^{\prime}, x^{\prime}, Y^{\prime}, \Lambda^{\prime}\right)=\left(\mathbb{R}^{4}, \omega, X,\{f=-1\}, S_{-}\right)
$$

Now applying Proposition 23 gives an isomorphism of isotropic setups and results in the contact manifold $\left(Y_{1} \# Y_{2}, \lambda_{1} \# \lambda_{2}\right)$. The strict contactomorphism away from the surgery region is given by flowing along $\partial t$. This is because $\mathcal{L}_{\partial t} \omega=\omega$ since $\partial t$ is a Liouville vector field, so $i_{\partial t} \omega$ is preserved under flowing along $\partial t$. Now, by construction $\lambda_{1} \# \lambda_{2}$ is nondegenerate on the complement of the surgery region.

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The linearized flow in (3.3.1) shows that $h$ is also nondegenerate. For other orbits potentially created during the connected sum operation, they must intersect the boundary spheres $S_{p}^{2} \sqcup S_{q}^{2}$ of a neighborhood of $\Lambda=\{p\} \sqcup\{q\}$. Therefore, we may find a $C^{\infty}$-small perturbation supported in an arbitrarily small neighborhood of the spheres $S_{p}^{2} \sqcup S_{q}^{2}$ as in [39] so that $\lambda_{1} \# \lambda_{2}$ is nondegenerate while maintaining the results in Theorem 25 in the next section.

## Almost complex structures

In this subsection, we discuss the pseudo-holomorphic behaviors in the symplectization of the connected sum region $\mathbb{R} \times S^{2} \subset \mathbb{R} \times\left(Y_{1} \# Y_{2}\right)$. This has been studied in [102, 39]. Fix a generic $\lambda_{i}$-adapted almost complex structure $J_{i}$ on $\mathbb{R} \times\left(Y_{i}, \lambda_{i}\right)$ and let

$$
\lambda^{\prime}:=\lambda_{1} \# \lambda_{2}
$$

be the connected sum contact form described in Lemma 24. Fish-Siefring found a $\lambda^{\prime}$-adapted almost complex structure such that one may see a pair of pseudoholomorphic planes asymptotic to $h$ in $\mathbb{R} \times S^{2}$. We continue to follow the paper [39], now for information regarding the pseudo-holomorphic curves.

Theorem 25. [39, Theorem 5.1] Let $Y:=Y_{1} \sqcup Y_{2}$ be a three-manifold equipped with a nondegenerate contact form $\lambda$. Let $Y^{\prime}$ be the connected sum manifold equipped with a nondegenerate contact form $\lambda^{\prime}$ given Lemma 24. Then, given any $\lambda^{\prime}$-adapted $J$, there exists a $\lambda^{\prime}$-adapted $J^{\prime}$ agreeing with $J$ outside the surgery region such that:

1. There exists a pair of embedded, disjoint $J^{\prime}$-holomorphic planes $P_{N}$ and $P_{S}$ both asymptotic to $h$.
2. The planes $P_{N}$ and $P_{S}$ approach $h$ in opposite directions, that is, their leading asymptotic coefficients are of opposite signs (see Section 3.3).
3. The planes $P_{N}$ and $P_{S}$ project to the northern and southern hemispheres of the connected sum sphere $S_{+}$under $\pi: \mathbb{R} \times Y^{\prime} \rightarrow Y^{\prime}$. Together with $\mathbb{R} \times h$, their $\mathbb{R}$-translations foliate $\mathbb{R} \times S_{+}$.

Proof. Consider the 1-form

$$
\alpha=\left.i_{X} \omega\right|_{f=1}=\frac{1}{2} x d y-\frac{1}{2} y d x+2 z d w+\left.w d z\right|_{f=1}
$$

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on $\mathbb{R} \times S_{+}$considered in Section 3.3. The coordinates in [39] are related to ours by the following coordinate change $\Phi$ :

$$
\begin{aligned}
& x(\rho, \phi, \theta)=\sin \theta \cos \phi \\
& y(\rho, \phi, \theta)=\sin \theta \sin \phi \\
& z(\rho, \phi, \theta)=4 \cos \theta \\
& w(\rho, \phi, \theta)=\frac{3}{8} \rho
\end{aligned}
$$

Now, at the sphere $w=\rho=0$,

$$
\begin{equation*}
\Phi^{*} \alpha=\frac{1}{2} \sin ^{2} \theta d \phi+3 \cos \theta d \phi \tag{3.3.2}
\end{equation*}
$$

which agrees with the contact form given in [39] up to a constant. Therefore, the proof of Lemma 5.7 in [39] goes through.

Now we compute the ECH index of the planes in Theorem 25.
Lemma 26. $I\left(P_{S}\right)=I\left(P_{N}\right)=\operatorname{ind}\left(P_{S}\right)=\operatorname{ind}\left(P_{N}\right)=1$.
Proof. Fix trivialization $\tau_{0}$ as in Section 3.3, then

$$
c_{\tau_{0}}\left(P_{S}\right)=1
$$

by considering the vector field

$$
z x \partial y-z y \partial x+\left(1-z^{2}\right) \partial w
$$

which has one positive zero at $z=-1$. The zero is positive by the Poincaré-Hopf theorem. This vector field is in the contact structure $\operatorname{ker} \alpha$ restricted to $P_{S}$ and tangent bundle of $\left.P_{S}\right|_{h}$ while being normal to the boundary, and under the Reeb flow this vector field does not rotate with respect to the contact structure restricted to $h$. Therefore,

$$
\operatorname{ind}\left(P_{S}\right)=-\chi(C)+2 c_{\tau_{0}}(C)+C Z_{\tau_{0}}(h)=-1+2+0=1
$$

Now, to compute the ECH index, we first need to compute $Q_{\tau_{0}}\left(P_{S}\right)$. Since our trivialization $\tau_{0}$ is the same trivialization associated to the constant section given by $\partial z$, one can see that

$$
Q_{\tau_{0}}\left(P_{S}\right)=0
$$

Therefore,

$$
I\left(P_{S}\right)=c_{\tau_{0}}\left(P_{S}\right)+Q_{\tau_{0}}\left(P_{S}\right)+C Z_{\tau_{0}}^{I}\left(P_{S}\right)=1+0+0=1
$$

Note that we did not need to compute the Fredholm index above, since if $I=1$, then ind $=1$ by Theorem 16 .

To see that all the curves are transversely cut out with the almost complex structure obtained in Theorem 25, we first observe that in the open set $W$ defined by $\mathbb{R} \times\left(Y_{1} \# Y_{2}\right)$ minus the symplectization of the connected sum sphere $\mathbb{R} \times S^{2}$, we can perturb our complex structure to be generic. Therefore, all the curves except for the two ind $=1$ planes bounding the hyperbolic orbit $h$ are transversely cut out, because they all intersect $W$. Then, we observe that for $P_{N}$ and $P_{S}$ which are embedded, we have automatic transversality (e.g. Proposition A. 1 in [123]). This is because

$$
\operatorname{ind}\left(P_{S}\right)>c_{N}\left(P_{S}\right)=\frac{\operatorname{ind}\left(P_{S}\right)-2+2 g+\# \Gamma_{0}}{2}=0
$$

where $\# \Gamma_{0}$ denotes the number of asymptotic ends with even Conley-Zehnder indices. Note that the parity of the Conley-Zehnder index is independent of the choice of trivialization. The calculation is the same for $P_{N}$.

Remark 27. Another useful way to obtain the contact connected sum is the following. First, drill two convex sutured balls $\left(D^{2} \times I, \Gamma\right)$, where the suture is $\Gamma=$ $\partial D^{2} \times\{1 / 2\}$, from the two contact three-manifolds $Y_{1}$ and $Y_{2}$. Turn one of the concave sutured manifolds, for example, $Y_{1} \backslash\left(D^{2} \times I, \Gamma\right)$, into a convex sutured manifold $Y_{1}(1)$ as described in [27]. This concave-to-convex process introduces a positive hyperbolic orbit. Finally, glue $Y_{1}(1)$ and $Y_{2} \backslash\left(D^{2} \times I, \Gamma\right)$ together to obtain a contact connected sum $Y_{1} \# Y_{2}$. Although this construction also results in exactly one extra positive hyperbolic orbit up to large action and yields isotopic contact structures as the Weinstein construction above, it is not a priori clear how the almost complex structure on its symplectization corresponds to that coming from the Weinstein 1-handle attachment.

## Asymptotic neighborhood of the special hyperbolic orbit $h$

The asymptotic neighborhood of a Reeb orbit encodes a lot of information about the potential pseudo-holomorphic curves that could asymptote to it. In order to understand the asymptotic neighborhood of $h$, we first review the asymptotic operator associated to a Reeb orbit. For more details, see for example [107] and [59].

Let $C$ be a somewhere injective $J$-holomorphic curve with an asymptotic end at an embedded Reeb orbit $\gamma$. By rescaling the $s$ and $t$ coordinates on $\mathbb{R} \times Y$ near the Reeb orbit $\gamma$, we may assume that $\gamma$ has period 1. The almost complex structure $J$ on $\left.\xi\right|_{\gamma}$ defines a family of $2 \times 2$ matrices $J_{t}$ such that $J_{t}^{2}=-1$ where $t \in \mathbb{R} / \mathbb{Z}$. The linearized Reeb flow $\Psi(t)$ along $\gamma$ gives a symplectic connection

$$
\nabla_{t}^{R}=\partial t+S_{t}
$$

on the $\left.\xi\right|_{\gamma}$, where $S_{t}$ is a symmetric matrix for each $t \in \mathbb{R} / \mathbb{Z}$.

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Definition 28. The asymptotic operator ${ }^{2}$ is defined as

$$
L_{\gamma}:=J_{t} \nabla_{t}^{R}: C^{\infty}\left(S^{1}, \gamma^{*} \xi\right) \rightarrow C^{\infty}\left(S^{1}, \gamma^{*} \xi\right)
$$

More specifically, let $\tau$ be a complex linear, symplectic trivialization of $\left.\xi\right|_{\gamma}$. For each $t \in S^{1}$, define

$$
S_{t}:=-J_{0} \frac{d \Psi(t)}{d t} \Psi^{-1}(t)
$$

Then we may write

$$
L_{\gamma}=J_{0} \frac{d}{d t}+S_{t} .
$$

Note $L_{\gamma}$ is self-adjoint.
Let $\eta(t)$ be an eigenfunction of $L_{\gamma}$ with eigenvalue $\lambda$. Then $\eta$ solves the ODE

$$
\frac{d \eta(t)}{d t}=J_{0}\left(S_{t}-\lambda\right) \eta(t)
$$

and hence $\eta$ is nonvanishing, if it's nonzero. Therefore, we may define $\operatorname{wind}_{\tau}(\eta)$ to be the winding number of the loop $\eta: \mathbb{R} / 2 \pi \mathbb{Z} \rightarrow \mathbb{C}$ around zero. The above discussion generalizes to multiple covers $\gamma^{d}$ of a simple orbit $\gamma$. Since we will focus on the hyperbolic orbit $h$ which does not allow multiple covers as ECH generators, we refer the readers to [107, 59] for further discussions.

We recall the following useful properties of eigenvalues and winding numbers associated to an asymptotic operator.

Lemma 29. [51, Lemma 6.4] and [45, Chapter 3]

1. If $\eta, \eta^{\prime}$ are eigenvalues of the asymptotic operator $L$ corresponding to eigenvalues $\lambda \leq \lambda^{\prime}$, then $\operatorname{wind}(\eta) \geq \operatorname{wind}\left(\eta^{\prime}\right)$.
2. For each winding number $w$, the space of eigenfunctions with winding number $w$ is 2-dimensional.
3. If $\gamma$ is nondegenerate, then $\operatorname{wind}_{\tau}(\phi) \leq\left\lfloor C Z_{\tau}\left(\gamma^{d}\right) / 2\right\rfloor$ for $\lambda>0$ i.e. $\lambda$ is an eigenvalue associated to an asymptotic operator at a positive end, and $\operatorname{wind}_{\tau}(\phi) \geq\left\lceil C Z_{\tau}\left(\gamma^{d}\right) / 2\right\rceil$ for $\lambda<0$ i.e. $\lambda$ is an eigenvalue associated to an asymptotic operator at a negative end.
[^1]
## CHAPTER 3. A CONNECTED SUM FORMULA FOR EMBEDDED

 CONTACT HOMOLOGYNow we review the asymptotic expansion of an asymptotically cylindrical pseudoholomorphic curve at a given nondegenenerate orbit. From now on, we focus on the case when a curve is positively asymptotic to a Reeb orbit, i.e. $s \gg 0$. The case for the negative asymptotic end is completely analogous. The following definition is from [107].

Definition 30. Let $\gamma$ be a Reeb orbit. Let $u$ be a $J$-holomorphic curve positively asymptotic to $\gamma$. Let $U:[R, \infty) \times S^{1} \rightarrow \gamma^{*} \xi$ be a smooth map satisfying $U(s, t) \in$ $\left.\xi\right|_{\gamma(t)}$ for all $(s, t) \in[R, \infty) \times S^{1}$, where $R$ is a large real number. Then, $U$ is an asymptotic representative of $u$ if there exists a proper embedding $\psi:[R, \infty) \times S^{1} \rightarrow$ $\mathbb{R} \times S^{1}$ asymptotic to the identity, so that

$$
u(\psi(s, t))=\left(s, \exp _{\gamma(t)} U(s, t)\right)=\widetilde{\exp }_{(s, \gamma(t))}(0, U(s, t))
$$

for all $(s, t) \in[R, \infty) \times S^{1}$, where exp and $\widetilde{\exp }$ are the exponential maps associated to a Riemannian metric on $Y$ and its $\mathbb{R}$-invariant lift to $\mathbb{R} \times Y$.

The asymptotic representative encodes the isotopy class of a braid $\zeta$ that is the intersection of the pseudo-holomorphic curve asymptotic to $\gamma$ and a tubular neighborhood $\mu$ of $\gamma$ at $s \gg 0$. Hofer-Wysocki-Zehnder [46] showed that the map $U$ can be written in the form

$$
\begin{equation*}
U(s, t)=e^{-\lambda s}(\eta(t)+r(s, t)) \tag{3.3.3}
\end{equation*}
$$

for an eigenvalue $\lambda>0$ and eigenfunction $\eta(t)$ of $L_{\gamma}$ and an error term $r(s, t)$ converging exponentially to zero as $s \rightarrow \infty$. The thesis of Siefring analyzed the difference of two pseudo-holomorphic half-cylinders asymptotic to the same Reeb orbit and as a result generalized the asymptotic form in Equation (3.3.3) to "higher orders" [107].

Theorem 31. [107, Theorem 2.2] Let $u, v$ be J-holomorphic curves with a positive asymptotic end at $\gamma$ and the maps $U, V:[R, \infty) \times S^{1} \rightarrow C^{\infty}\left(\gamma^{*} \xi\right)$ be their asymptotic representatives. Assume $U-V$ does not vanish identically. Then there exists a positive eigenvalue $\lambda$ of the asymptotic operator $L_{\gamma}$ with an eigenfunction $\eta(t)$ such that

$$
\begin{equation*}
U(s, t)-V(s, t)=e^{-\lambda s}(\eta(t)+r(s, t)), \tag{3.3.4}
\end{equation*}
$$

where the map $r(s, t)$ satisfies the decay estimate

$$
\begin{equation*}
\left|\nabla_{s}^{i} \nabla_{t}^{j} r(s, t)\right| \leq M_{i j} e^{-d s} \tag{3.3.5}
\end{equation*}
$$

for every $(i, j) \in \mathbb{N}^{2}$, where $M_{i j}$ and d consist positive constants.

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A consequence of the above theorem of Siefring is that, for any positive integer $N$, we may expand a positive end of a pseudoholomorphic curve by

$$
\begin{equation*}
u(s, t)=\sum_{i}^{N} e^{-\lambda_{i} s} a_{i} \eta_{i}(t)+o_{\infty}\left(\lambda_{N}\right) \tag{3.3.6}
\end{equation*}
$$

where $a_{i} \in \mathbb{R}$ keeps track of the scalar of the eigenfunction $\eta_{i}(t)$ and $o_{\infty}$ is a function $f:[R, \infty) \times S^{1} \rightarrow \mathbb{R}^{2}$ satisfying the decay estimate (3.3.5), c.f. Theorem 2.3 in [107]. We call $a_{i}$ the asymptotic coefficients of the asymptotic expansion of $u$ at $\gamma$. An immediate corollary is the following.

Corollary 32. Let $u, v$ be distinct pseudo-holomorphic curves with a common asymptotic end at $\gamma$. Then the asymptotic coefficients of the asymptotic expansions of $u$ and $v$ at $\gamma$ differ at some term.

To fix notations, from now on, let $\left(\lambda_{i}, \eta_{i}\right)$ be the corresponding eigenvalues and eigenvectors of the asymptotic operator $L_{h}$ associated to the hyperbolic orbit $h$. Fix some integer $N \gg 0$. Let the asymptotic expansion of $P_{N}$ be

$$
\sum_{i}^{N} e^{-\lambda_{i} s} a_{i} \eta_{i}(t)+o_{\infty}\left(\lambda_{N}\right)
$$

and the asymptotic expansion of $P_{S}$ be

$$
\sum_{i}^{N} e^{-\lambda_{i} s} b_{i} \eta_{i}(t)+o_{\infty}\left(\lambda_{N}\right)
$$

Notice that we use the same index $i$, and some of $a_{i}$ and $b_{i}$ could be zero. In addition, $a_{0}$ and $b_{0}$ are of opposite signs by Theorem 25(2).

By Lemma 29(3), we see that wind $\tau_{\tau_{0}}(h) \leq 0$ when the special hyperbolic orbit $h$ is at a positive end. In Lemma 37, we show that this bound is achieved. However, before that, we need to introduce a few more definitions in preparation of the proof.

Specifically, let $\zeta$ be a braid in a tubular neighborhood $\mu$ of $\gamma$ determined by a pseudo-holomorphic curve $u$ asymptotic to $\gamma$ as above. Identify

$$
\mu \cong S^{1} \times \mathbb{R}^{2} \cong A \times(0,1) \subset \mathbb{R}^{3}
$$

with respect to a trivialization $\tau$, where $A$ denotes an annulus. Let $\gamma$ be an embedded Reeb orbit such that $C$ has positive ends of multiplicities $q_{1}, \ldots, q_{n}$ at $\gamma$ with total multiplicity $m$. Each end of a pseudo-holomorphic curve determines a braid component $\zeta_{i}$ of the braid $\zeta$ with $q_{i}$ strands, where $\zeta=i m(u) \cap\{s\} \times \mu$ for $s \gg 0$ as before.

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 CONTACT HOMOLOGYDefinition 33. The writhe $w_{\tau}(\zeta) \in \mathbb{Z}$ is defined to be one half the signed count of crossings under the projection $A \times(0,1) \rightarrow A$.

Definition 34. Let $\zeta_{1}$ and $\zeta_{2}$ be two disjoint braids in the neighborhood of $\gamma$. Then the linking number $l_{\tau}\left(\zeta_{1}, \zeta_{2}\right) \in \mathbb{Z}$ is defined to be one half the signed count of crossings between $\zeta_{1}$ and $\zeta_{2}$ under the projection $A \times(0,1) \rightarrow A$.

In light of the braid picture, the winding number of an eigenfunction $\eta_{\zeta}$ associated to the braid $\zeta$ around the underlying simple orbit $\gamma$ is exactly

$$
\operatorname{wind}_{\tau}\left(\eta_{\zeta}\right)=l_{\tau}(\zeta, \gamma)
$$

Definition 35. Let $S$ and $S^{\prime}$ be admissible representatives of $H_{2}(Y, \alpha, \beta)$ and $H_{2}\left(Y, \alpha^{\prime}, \beta^{\prime}\right)$. Then,

$$
\begin{aligned}
w_{\tau}(S) & :=\sum_{i} w_{\tau}\left(\zeta_{i}^{+}\right)-\sum_{j} w_{\tau}\left(\zeta_{j}^{-}\right) \\
\operatorname{wind}_{\tau}(S) & :=\sum_{i} \operatorname{wind}_{\tau}\left(\zeta_{i}^{+}\right)-\sum_{j} \operatorname{wind}_{\tau}\left(\zeta_{j}^{-}\right), \\
l_{\tau}\left(S, S^{\prime}\right) & :=\sum_{i} l_{\tau}\left(\zeta_{i}^{+}, \zeta_{i}^{+\prime}\right)-\sum_{j} l_{\tau}\left(\zeta_{j}^{-}, \zeta_{j}^{-\prime}\right)
\end{aligned}
$$

where $\zeta_{i}^{+}$and $\zeta_{i}^{+\prime}$ are braids defined by the intersections of $S$ and $S^{\prime}$ with $(\{1-\epsilon\} \times Y)$, and $\zeta_{j}^{-}$and $\zeta_{j}^{-\prime}$ similarly with $\{-1+\epsilon\} \times Y$.

We have the following "linking bound" lemma that first appeared in [51].
Lemma 36. [53, Lemma 5.5(b)] Let

$$
\rho_{i}:=\left\lfloor\frac{C Z_{\tau}\left(\gamma^{q_{i}}\right)}{2}\right\rfloor .
$$

If $i \neq j$, then $\zeta_{i}$ and $\zeta_{j}$ are disjoint and

$$
l_{\tau}\left(\zeta_{i}, \zeta_{j}\right) \leq \max \left(\rho_{i} q_{j}, \rho_{j} q_{i}\right)
$$

In particular, any two braids $\zeta_{i}$ and $\zeta_{j}$ associated to the asymptotic neighborhood at $h$ have bounded linking number $l_{\tau_{0}}\left(\zeta_{i}, \zeta_{j}\right) \leq 0$ since $C Z_{\tau_{0}}(h)=0$.

Lemma 37. $\operatorname{wind}_{\tau_{0}}\left(P_{S}\right)=\operatorname{wind}_{\tau_{0}}\left(P_{N}\right)=0$.

Proof. We prove the lemma for $P_{S}$. The proof for $P_{N}$ is exactly the same. Applying the relative adjunction formula in Proposition 17 to $C:=P_{S} \sqcup P_{S}^{\prime}$, where $P_{S}^{\prime}$ is an $s$-translation of $P_{S}$, we get

$$
c_{\tau_{0}}(C)=\chi(C)+Q_{\tau_{0}}(C)+w_{\tau_{0}}(C)-2 \delta(C)
$$

By the construction of almost complex structure in Theorem 25 that admits the foliation of $\mathbb{R} \times S^{2}$ where $P_{S}$ is a leaf, we have that $\delta(C)=0$. We also know that $Q_{\tau_{0}}(C)=0$ since $Q_{\tau}$ is quadratic and $Q_{\tau_{0}}\left(P_{S}\right)=0$. In addition, $w_{\tau_{0}}\left(P_{S}\right)=0$. This is because the multiplicity of $h$ is one and therefore $w_{\tau_{0}}\left(\zeta_{h}\right)=0$ where $\zeta_{h}$ is the braid at the intersection of $P_{S}$ and $\{1-\epsilon\} \cap Y$. Now we have that

$$
2 c_{\tau_{0}}\left(P_{S}\right)=2+2 \operatorname{wind}_{\tau_{0}}\left(P_{S}\right)
$$

since

$$
w_{\tau_{0}}\left(P_{S} \sqcup P_{S}^{\prime}\right)=4 w_{\tau_{0}}\left(P_{S}\right)+2 \operatorname{wind}_{\tau_{0}}\left(P_{S}\right)=2 \operatorname{wind}_{\tau_{0}}\left(P_{S}\right),
$$

where the first equality is justified in the proof of Proposition 8.4 in [51]. Therefore,

$$
\operatorname{wind}_{\tau_{0}}\left(P_{S}\right)=0
$$

The above lemma is needed for our specific case when the almost complex structure on $\mathbb{R} \times S^{2}$ is not generic as constructed in Theorem 25 . When we have generic almost complex structure, the following proposition says that the extremal bound of the winding number is always achieved.

Proposition 38. [61, Proposition 3.2] If the symplectization-adapted almost complex structure $J$ on $\mathbb{R} \times Y$ is generic, then for any Fredholm index 1, connected, non-multiply-covered $J$-holomorphic curve $C$ having a positive end at $\gamma$, the winding number of the leading eigenfunction of the asymptotic expansion of $C$ achieves the equality in Lemma 29(3).

### 3.4 Correspondence of Reeb orbits and filtered ECH complex

In this section, we discuss how to correspond the Reeb orbits in the closed manifolds $Y_{1} \# Y_{2}$ to the ones in $Y_{1}$ and $Y_{2}$. We show that up to a sufficiently large action $L$, we may ignore orbits crossing the connected sum sphere. This allows an identification on the vector space level of a filtered ECH complex of the connected sum and a filtered mapping cone complex.

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## The mapping cone

Given a chain map between two chain complexes, one can form its mapping cone. In the following, we review this construction in our setting. We define

$$
C_{o}:=E C C\left(Y_{1} \sqcup Y_{2}, \lambda_{1} \sqcup \lambda_{2}\right)=E C C\left(Y_{1}, \lambda_{1}\right) \otimes_{\mathbb{F}} E C C\left(Y_{2}, \lambda_{2}\right) .
$$

Definition 39. Given the chain map $\varphi: C_{o} \rightarrow C_{o}$ defined by

$$
\varphi:=U_{1} \otimes i d+i d \otimes U_{2}
$$

one can form the mapping cone Cone $(\varphi)$ by the following. As a vector space,

$$
\operatorname{Cone}(\varphi):=C_{o} \oplus C_{o}[-1],
$$

where $C_{o}[-1]^{*}:=C_{o}^{*-1}$. The differential of $\operatorname{Cone}(\varphi)$ is given by

$$
\partial_{\text {cone }}:=\left(\begin{array}{cc}
\partial_{1} \otimes i d+i d \otimes \partial_{2} & 0 \\
\varphi & \partial_{1} \otimes i d+i d \otimes \partial_{2}
\end{array}\right)
$$

Let $C_{h}:=C_{o} \otimes h$ denote the collection of all orbit sets in $Y_{1} \sqcup Y_{2}$ concatenating with the special hyperbolic orbit $h$. By Theorem 25 and Lemma 26, we know that $h$ bounds $I=1$ planes $P_{N}$ and $P_{S}$. Therefore, one can think of $h$ as an element that increases the relative grading by one. We can then identify $C_{o}[-1]$ with $C_{h}$ and obtain that

$$
\operatorname{Cone}(\varphi)=C_{0} \oplus C_{h}
$$

as vector spaces. Furthermore, we define a filtered mapping cone complex.
Definition 40. Given a real number $L$, we define

$$
\operatorname{Cone}^{L}(\varphi):=C_{o}^{L} \oplus C_{h}^{L}=C_{o}^{L} \oplus C_{o}{ }^{L-\mathcal{A}(h)} .
$$

One may check that $\partial_{\text {cone }}$ preserves $\operatorname{Cone}^{L}(\varphi)$, since both the differential and the U map decrease the symplectic action by Stokes' theorem and the definition of a $\lambda$-adapted almost complex structure. Therefore, $\operatorname{Cone}^{L}(\varphi)$ is a subcomplex of Cone $(\varphi)$. In addition, we observe that

$$
\lim _{L \rightarrow \infty} \operatorname{Cone}^{L}(\varphi)=\operatorname{Cone}(\varphi)
$$

given by the obvious inclusion maps.

## Filtered ECH of a connected sum

We use filtered ECH to ignore potential orbits of large symplectic actions formed during the connected sum procedure, which cross the connected sum $S^{2}$. The following lemma is folklore.

Lemma 41. Let $L>0$. Let $p$ be a point which is not on a Reeb orbit of action $<L$. Then there exists a radius $R(L)$ such that any Reeb trajectory which starts on the ball $B$ of radius $R(L)$, leaves $B$ and returns to $B$, has symplectic action greater or equal $L$.

Proof. Suppose not. We take a sequence of balls $B_{n}$ around $p$ of radius converging to 0 , and Reeb trajectories $\gamma_{n}$ of action less than $L$ which start on $B_{n}$, leave $B_{n}$, and return to $B_{n}$. We can then pass to a subsequence so that the Reeb trajectories $\gamma_{n}$ converge to a Reeb orbit $\gamma$ of action less than $L$ which passes through $p$. Contradiction.

The above lemma tells us that up to an action $L$, we may ignore orbits that pass through the connected sum region other than the special hyperbolic orbit $h$. More precisely, using Lemma 41, we may find a strictly decreasing smooth function $R: \mathbb{R} \rightarrow \mathbb{R}$ such that $\lim _{L \rightarrow \infty} R(L)=0$ and

$$
\begin{equation*}
E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right) \cong C_{o}^{L} \oplus C_{h}^{L}=\text { Cone }^{L} \tag{3.4.1}
\end{equation*}
$$

as vector spaces, where $\alpha_{R(L)}:=\lambda_{1} \#_{R(L)} \lambda_{2}$ denotes the contact form on $Y_{1} \#_{R(L)} Y_{2}$ where the function $f$ in Section 3.3 is scaled by $R(L)$. Note that any orbit intersecting a connected sum sphere of radius $r$ would also intersect a connected sum sphere of radius $r^{\prime}>r$, so we may assume $R(L)$ in Lemma 41 is non-increasing. Furthermore, we can modify $R(L)$ to be smooth and strictly decreasing. Therefore, by (3.4.1), we obtain an isomorphism on the level of vector spaces of the following theorem.

Proposition 42 (Proposition 4). Given two closed connected contact three-manifolds $\left(Y_{1}, \lambda_{1}\right)$ and $\left(Y_{2}, \lambda_{2}\right)$ with nondegenerate contact forms $\lambda_{i}$. Let $R$ be the decreasing function of $L$ defined above. Then there is a chain homotopy equivalence

$$
f: E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right) \rightarrow \operatorname{Cone}^{L}\left(U_{1} \otimes i d+i d \otimes U_{2}\right)
$$

The main goal of Section 3.5 and Section 3.6 is to prove Proposition 42 by studying the connected sum differentials. Then, in Section 3.7, we will pass the filtered ECH chain complexes $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$ to a direct limit, where we consider smaller and smaller connected sum spheres as we let the symplectic action $L$ go to infinity, to obtain the main theorem.

### 3.5 The differentials for ECH of a connected sum

There are four types of curves to consider for the differential of $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$ :

- $I=1$ curves in $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive ends in $C_{o}^{L}$ and negative ends in $C_{o}^{L}$, denoted as $\partial_{o o}$;
- $I=1$ curves in $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive ends in $C_{h}^{L}$ and negative ends in $C_{o}^{L}$, denoted as $\partial_{o h}$;
- $I=1$ curves in $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive ends in $C_{o}^{L}$ and negative ends in $C_{h}^{L}$, denoted as $\partial_{h o}$;
- $I=1$ curves in $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive ends in $C_{h}^{L}$ and negative ends in $C_{h}^{L}$, denoted as $\partial_{h h}$.

We organize this as

$$
\partial_{\#}=\left(\begin{array}{ll}
\partial_{o o} & \partial_{o h} \\
\partial_{h o} & \partial_{h h}
\end{array}\right) .
$$

In this section, we describe how the differentials $\partial_{o o}, \partial_{o h}$ and $\partial_{h h}$ behave in terms of the original differentials in $E C C^{L}\left(Y_{1}, \lambda_{1}\right)$ and $E C C^{L}\left(Y_{2}, \lambda_{2}\right)$. In particular, we show that the connected sum differentials $\partial_{o o}$ and $\partial_{h h}$ are identified with the original differentials on $Y_{1}$ and $Y_{2}$, considering the additional trivial cylinder over $h$. The differential $\partial_{o h}$ is reminiscent of the linearized contact homology case [13], where the special hyperbolic orbit can be assumed to have small action and hence cannot admit a curve positively asymptotic to it while having negative asymptotics. In our case where a curve may have multiple positive asymptotic ends, we need to study the asymptotic behavior near the neighborhood of the hyperbolic orbit $h$ more closely. The differential $\partial_{h o}$ involves the original $U$ maps in $E C C^{L}\left(Y_{1}, \lambda_{1}\right)$ and $E C C^{L}\left(Y_{2}, \lambda_{2}\right)$ and will be discussed in Section 3.6.

## No-crossing lemma

This lemma appears in the draft of the connected sum formula for linearized contact homology [13] and we give essentially the same proof here using the relative intersection number $Q_{\tau}$ and the asymptotic linking number $l_{\tau}$.

Lemma 43 (No-crossing). Let $C$ be a connected, embedded $J$-holomorphic curve in the symplectization $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$ with ends that do not intersect the connected sum $S^{2}$. Then the image of $C$ cannot have ends in both $Y_{1}$ and $Y_{2}$ away from the connected sum region.

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 CONTACT HOMOLOGYProof. Suppose we have a curve $C$ with positive ends at $\alpha$ and negative ends at $\beta$, such that $\alpha \cup \beta$ lie in both $Y_{1}$ and $Y_{2}$ away from the Darboux balls used for the connected sum operation. Then $C$ must intersect $\mathbb{R} \times S_{+}$. In particular, $C$ must intersect one of the holomorphic planes that are the translations of $P_{N}$ the $P_{S}$ given in Theorem 25 or the trivial cylinder over $h$. We may assume without loss of generality that $C$ intersects $P_{N}$ or $C$ intersects the trivial cylinder over $h$. Therefore, we have $\#\left(\dot{C} \cap \dot{P_{N}}\right) \geq 1$ or $\#(\dot{C} \cap(\mathbb{R} \times h)) \geq 1$.

Suppose $\#\left(\dot{C} \cap \dot{P_{N}}\right) \geq 1$. Recall that

$$
Q_{\tau}\left(C, P_{N}\right)=\#\left(\dot{S} \cap \dot{S}^{\prime}\right)-l_{\tau}\left(S, S^{\prime}\right)
$$

where $S$ and $S^{\prime}$ are any admissible representatives of $[C]$ and $\left[P_{N}\right]$, and the interiors $\dot{S}$ and $\dot{S}^{\prime}$ are transverse and do not intersect near the boundary. We have that

$$
l_{\tau}\left(S, S^{\prime}\right)=0
$$

since by assumptions the ends of $C$ do not overlap with the end $h$ of $P_{N}$.
Now it suffices to show that

$$
Q_{\tau}\left(C, P_{N}\right)=0
$$

By a usual transversality argument, one can find another representative of $[C] \in$ $H_{2}(Y, \alpha, \beta)$, which is still denoted as $S$, such that the intersections of $S$ and $S^{\prime}$ miss a neighborhood of the north pole on $P_{N}$. Therefore, we may define a homotopy $S_{t \in[0,1]}$ moving all possible interior intersections of $S$ and $S^{\prime}$ to be on $P_{S}$. Therefore,

$$
\#\left(\dot{S} \cap \dot{S}^{\prime}\right)=\#\left(\dot{S}_{1} \cap \dot{S}^{\prime}\right)=0,
$$

so $Q_{\tau}\left(C, P_{N}\right)=0$. The arguments for $Q_{\tau}(C, \mathbb{R} \times h)=0$ is entirely analogous.

## $\partial_{o o}$ : differential from orbit sets without $\boldsymbol{h}$ to orbit sets without $h$

The discussions in Section 3.4 shows that, by using filtered ECH up to filtration $L$, we may focus only on orbits of actions less than $L$ in $\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$ and all such orbits are away from the connected sum $S^{2}$ except for the special hyperbolic orbit $h$. We then use the no-crossing lemma to show that the part of the connected sum differential not approaching the hyperbolic orbit $h$ at all is the same as the sum of the original differential $\partial_{i}$ on $E C C^{L}\left(Y_{i}, \lambda_{1}\right)$ :

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 CONTACT HOMOLOGYLemma 44. We have that the ECH differential

$$
\partial_{o o}=\partial_{1} \otimes i d+i d \otimes \partial_{2}
$$

when restricted to the subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.
Proof. By Lemma 41, orbits that intersect the connected sum $S^{2}$ have energy $\geq$ $L$, except for the special hyperbolic orbit $h$. To study $\partial_{o o}$, we may focus on the connected, embedded, $I=1$ curves given by Theorem 16. Then by the Lemma 43, the curves counted by $\partial_{o o}$ are exactly those that do not pass through $\mathbb{R} \times S^{2}$. These are exactly the $I=1$ curves either completely in $\mathbb{R} \times Y_{1}$ or completely in $\mathbb{R} \times Y_{2}$, i.e. those counted by $\partial_{1} \otimes i d+i d \otimes \partial_{2}$.

## $\partial_{o h}$ : differential from orbit sets with $\boldsymbol{h}$ to orbit sets without $\boldsymbol{h}$

The heuristics come from "local energy" of curves near the asymptotic end at the special hyperbolic orbit $h$. In this subsection, we eliminate all curves that have a positive asymptotic end at $h$ and have no negative asymptotic end at $h$, other than the two planes $P_{S}$ and $P_{N}$ and their translations bounding $h$ as described in Section 3.3.

Definition 45. A $J$-holomorphic curve $C$ approaches $h$ from the south (resp. north) if its leading asymptotic coefficient has the same sign as that of $P_{S}$ (resp. $P_{N}$ ).

Proposition 46. Let $J$ be generic in the open set $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}$. Any embedded, Fredholm index 1, connected curve with positive ends in $C_{h}$ and negative ends in $C_{o}$ is either $P_{S}$ or $P_{N}$, or a translation of them in the $\partial s$ direction.

Proof. Suppose $C$ is an embedded, Fredholm index 1, connected $J$-holomorphic curve with a positive asymptotic end at $h$. In particular, $C$ is not a trivial cylinder by Theorem 16. Suppose $C$ is not a translation of $P_{S}$ or $P_{N}$. In particular,

$$
C \cap\left(\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}\right) \neq \emptyset,
$$

so $J$ is generic for $C$. By Lemma 29(3) and Proposition 38, we know that the leading eigenfunction $\phi$ of $C$ at $h$ has

$$
\operatorname{wind}_{\tau_{0}}(\phi)=\left\lfloor C Z_{\tau_{0}}(h) / 2\right\rfloor=0
$$

Notice that this is a calculation that holds for any $C$ satisfying the assumptions.
Without loss of generality, assume that $C$ approaches $h$ from the south. In particular, the leading asymptotic coefficient of $C$ has the same sign as that of

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$$
\begin{equation*}
Q_{\tau_{0}}\left(C, P_{S}\right)+l_{\tau_{0}}\left(C, P_{S}\right)=\#\left(\dot{C} \cap \dot{P}_{S}\right) \geq 0 \tag{3.5.1}
\end{equation*}
$$

since both $C$ and $P_{S}$ are admissible representatives of their respective relative second homology classes.

Note that the leading eigenvalues of $C$ and $P_{S}$ both achieve the minimal possible eigenvalue of $L_{h}$ by Lemma 37 for curves completely inside the $\mathbb{R} \times S^{2}$ region, i.e. $P_{S}$ and $P_{N}$, and by Lemma 29(3) and Proposition 38 for curves that intersect the $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}$ region hence they are generic so they also achieve minimal winding numbers. Let the asymptotic expansions of $C$ and $P_{S}$ at $h$ be:

$$
\begin{align*}
C & =\sum_{i}^{N} e^{-\lambda_{i} s} c_{i} \eta_{i}+o_{\infty}\left(\lambda_{N}\right)  \tag{3.5.2}\\
P_{S} & =\sum_{i}^{N} e^{-\lambda_{i} s} b_{i} \eta_{i}+o_{\infty}\left(\lambda_{N}\right) \tag{3.5.3}
\end{align*}
$$

for some integer $N \gg 0$, where $\left(\lambda_{i}, \eta_{i}\right)$ are the corresponding eigenvalues and eigenfunctions of the asymptotic operator $L_{h}$ as discussed in Section 3.3. From the above discussions on minimal winding numbers, we have that $c_{0} \neq 0$ and $b_{0} \neq 0$. In addition, by the assumption that $C$ approaches $h$ from the south, we know that $c_{0}$ and $b_{0}$ have the same sign. Now we have two possibilities.

If $c_{0}=b_{0}$, then we may find the first $i>0$ such that $c_{i} \neq b_{i}$ by Corollary 32 and the assumption that $C \not \equiv P_{S}$. Now by Lemma 29(3), we know that the winding number that is 0 is achieved by the corresponding eigenfunctions of both the maximal negative and minimal positive eigenvalues. Since $\lambda_{0}<\lambda_{i}$, by Lemma 29(1), we have that $\operatorname{wind}\left(\eta_{i}\right) \leq \operatorname{wind}\left(\eta_{0}\right)$. Therefore, now by Lemma $29(2), \operatorname{wind}\left(\eta_{i}\right)<\operatorname{wind}\left(\eta_{0}\right)=0$. Therefore, we have that

$$
l_{\tau_{0}}\left(C, P_{S}\right) \leq \operatorname{wind}\left(\eta_{i}\right) \leq-1
$$

as in the proof of Lemma 6.9 in [51].
If $c_{0} \neq b_{0}$, then we can shift $P_{S}$ in the $\partial s$ direction to some $s^{\prime} \neq s$ corresponding to the curve $P_{S^{\prime}}$ so that $c_{0} e^{-\lambda_{0} s}=b_{0} e^{-\lambda_{0} s^{\prime}}$, since $c_{0}$ and $b_{0}$ have the same sign. Then some higher order $i>1$ term will contribute to at least an additional -1 to wind $\left(\eta_{0}\right)$ by using the same arguments as in the above paragraph. Therefore, in this case we also have

$$
l_{\tau_{0}}\left(C, P_{S^{\prime}}\right) \leq-1
$$

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 CONTACT HOMOLOGYRecall that $Q_{\tau_{0}}\left(C, P_{S}\right)$ is well-defined under relative second homology class. Now by the inequality in (3.5.1), it suffices to show that $Q_{\tau_{0}}\left(C, P_{S}\right)=0$. Combined with the fact that the connected sum $S^{2}$ is null-homologous, it suffices to show $Q_{\tau_{0}}\left(C, P_{N}\right)=0$ by (3.2.3).

By assumption, $C$ and $P_{N}$ have their leading asymptotic eigenfunctions of opposite signs. Therefore, $l_{\tau_{0}}\left(C, P_{N}\right)=0$. We then have that

$$
Q_{\tau_{0}}\left(C, P_{N}\right)=\#\left(\dot{C} \cap \dot{P_{N}}\right)-l_{\tau_{0}}\left(C, P_{N}\right)=\#\left(\dot{C} \cap \dot{P_{N}}\right) .
$$

In particular, $\#\left(\dot{C} \cap \dot{P_{N}}\right)$ is independent of the admissible representative $\Sigma$ of $[C]$, since $Q_{\tau_{0}}$ is. Now, by a usual transversality argument, we may find an admissible representative $\Sigma$ of $[C]$ such that $\Sigma$ does not pass through a neighborhood of the north pole of the connected sum $S^{2}$. Therefore, we may define a homotopy $\Sigma_{t \in[0,1]}$ moving all possible interior intersections of $\Sigma$ and $P_{N}$ to be on $P_{S}$. Therefore,

$$
\#\left(\dot{C} \cap \dot{P_{N}}\right)=\#\left(\dot{\Sigma_{1}} \cap \dot{P_{N}}\right)=0,
$$

so $Q_{\tau_{0}}\left(C, P_{N}\right)=0$. This gives a contradiction to (3.5.1).
Corollary 47. We have that the ECH differential $\partial_{o h}=0$ when restricted to the subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.

Proof. By Lemma 43, we may restrict our attention to orbits that lie in $C_{o}^{L} \oplus C_{h}^{L}$. By Proposition 46, the only ECH index 1 curves that have a positive asymptotic end at $h$ and no negative asymptotic end at $h$ are $P_{S}$ and $P_{N}$ up to $\mathbb{R}$-translation. This immediately means that $\partial_{o h}=0$ on $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$, since we are working over $\mathbb{F}$.

## $\partial_{h h}$ : differential from orbit sets with $\boldsymbol{h}$ to orbit sets with $\boldsymbol{h}$

In this subsection, we show that when restricted to $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$, any ECH index 1 current that have both a positive asymptotic end and a negative asymptotic end at $h$ must separate out a trivial cylinder over $h$.

Proposition 48. Let $J$ be generic in the open set $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}$. Any embedded, Fredholm index 1 curve with positive ends in $C_{h}$ and negative ends in $C_{h}$ must contain a trivial cylinder over $h$.

Proof. The proof is exactly the same as that of Proposition 46. Consider an embedded, Fredholm index 1, connected holomorphic curve $C$ with a positive end at $h$ and
a negative end at $h$. In particular, Proposition 38 still holds. Now consider again the equation

$$
Q_{\tau_{0}}\left(C, P_{S}\right)+l_{\tau_{0}}\left(C, P_{S}\right)=\#\left(\dot{C} \cap \dot{P}_{S}\right)
$$

As in the proof of Proposition 46, the left hand side is strictly negative while by positivity of intersections, the right hand side is greater or equal to zero. Therefore, the index 1 curve $C$ was disconnected. In order for $C$ to have a positive end at $h$, we know the only possibility is for $C$ to contain a trivial cylinder over $h$ or one of the $P_{S}$ and $P_{N}$ planes by Proposition 46. However, the latter case can be ruled out since both $P_{S}$ and $P_{N}$ are of Fredholm index 1 and the Fredholm index is additive.

Corollary 49. We have that the ECH differential

$$
\partial_{h h}=h \partial_{o o} \frac{1}{h}
$$

when restricted to the subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.
Proof. By Lemma 43, we may restrict our attention to orbits that lie in $C_{o}^{L} \oplus C_{h}^{L}$. By Proposition 48, we know that any current $C$ counted by $\partial_{h h}$ must contain a trivial cylinder over $h$, which is of ECH index 0 . The remaining components of $C$ contribute exactly to currents counted by $\partial_{o o}$.

### 3.6 The chain homotopy equivalence

In this section, we discuss the last piece of the differential in the connected sum, i.e. $\partial_{h o}$, which is the differential from orbit sets without $h$ to orbit sets with $h$. This uses a similar chain homotopy defined in the proof to show that the ECH $U$ map is independent of the change of base points (Section 2.5 in [66]). Roughly speaking, the differential $\partial_{h o}$ in the connected sum corresponds to $U$ maps in the original contact three-manifolds $Y_{1}$ and $Y_{2}$, up to a "chain homotopy" $K$ discussed in Section 3.6. This map $K$ will in turn help us construct the chain homotopy equivalence between the filtered ECH chain complex of the connected sum and the filtered mapping cone in Section 3.6.

## A "chain homotopy" induced by the change of base point

First, we need a lemma analogous to Lemma 44 for the $U$ maps. Let $U_{\#, z_{i}}$ denote the U map in $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with base point $z_{i}$ contained in $Y_{i}$ away from the connected sum region. A similar proof to that of Lemma 44 gives the following lemma, since
$I=2$ curves with asymptotic ends on ECH generators are also embedded by Theorem 16 and satisfy the no-crossing lemma.

Lemma 50. The moduli space of curves counted by $U_{\#, z_{i}}$ with asymptotic ends in $C_{o}$ is exactly the same as the moduli space of curves counted by $I=2$ curves in $\mathbb{R} \times Y_{i}$ passing through the point $z_{i}$, for $i \in\{1,2\}$. In particular,

$$
U_{1} \otimes i d+i d \otimes U_{2}=\left(U_{\#, z_{1}}\right)_{o o}+\left(U_{\#, z_{2}}\right)_{o o}
$$

when restricted to the filtered subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.
There are very strong restrictions on which curves could have a positive end at $h$, as shown in the previous section. When we require such curves to pass through a point $p$ in $P_{N}$, this naturally gives more restrictions to the curves. In fact, in the following, we obtain that the only curve with a positive end at $h$ that passes through $p$ is $P_{N}$.

Lemma 51. Let $J$ be generic in the open set $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}$. Let $C$ be an $I=1$ curve that has positive ends in $C_{h}$, negative ends in $C_{o} \oplus C_{h}$, and passes through a given point $p \in P_{N}$. Then $C \equiv P_{N}$.

Proof. This is a combination of Proposition 46 and Proposition 48.
Remark 52. We can also pick $p$ on $P_{S}$. Note that $p$ is not generic in the sense that it lies on a Fredholm index 1 curve.

Definition 53. Let $\mathcal{M}_{1}(\alpha ; \beta ; p)$ denote the moduli space of $I=1$ curves in $\mathbb{R} \times$ $\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive asymptotic ends at the orbit set $\alpha$ and negative asymptotic ends at orbit set $\beta$, passing through $\{0\} \times A$, where $A:=\left[z_{1}, p\right) \cup\left(p, z_{2}\right]$ is required to not pass through $P_{S}$ or any Reeb orbit, which is possible since there are only countably many Reeb orbits. See Figure 3.1.

Lemma 54. Let $J$ be generic in the open set $\mathbb{R} \times\left(Y_{1} \#_{R(L)} Y_{2}\right) \backslash \mathbb{R} \times S^{2}$. Then $\mathcal{M}_{1}(\alpha ; \beta ; p)$ is a compact 0-dimensional manifold.

Proof. Consider the moduli space $\widehat{\mathcal{M}}_{1}(\alpha ; \beta ; p)$ consisting of $I=1$ curves in $\mathbb{R} \times$ $\left(Y_{1} \#_{R(L)} Y_{2}\right)$ with positive asymptotic ends at the orbit set $\alpha$ and negative asymptotic ends at the orbit set $\beta$, passing through $\{0\} \times(A \cup\{p\})$. Observe that the moduli space $\mathcal{M}_{1}(\alpha ; \beta ; p)$ is the same as $\widehat{\mathcal{M}}_{1}(\alpha ; \beta ; p)$, except we are subtracting the holomorphic plane $P_{N}$ by Lemma 51. Since $\widehat{\mathcal{M}}_{1}(\alpha ; \beta ; p)$ is a compact 0-dimensional manifold (Section 2.5 of [66]), we have that $\mathcal{M}_{1}(\alpha ; \beta ; p)$ is also a compact 0-dimensional manifold.

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Figure 3.1: The base point $p$ on the pseudo-holomorphic plane $P_{N}$

In fact, by relating to the change of base point homotopy in the setting of [66], we obtain more information. First, we define the following related map.

Definition 55. Let $p$ be a point on $P_{N}$. Given an orbit set $\alpha$, the linear map $K$ is defined as

$$
K \alpha:=\sum_{\beta} \sum_{C \in \mathcal{M}_{1}(\alpha ; \beta ; p)} \beta
$$

where $\beta$ ranges over all orbit sets.
Lemma 56. We have that $K_{\text {oh }}=0$ when restricted to the subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.
Proof. Recall again that we may restrict our attention to $C_{o}^{L} \oplus C_{h}^{L}$ by the no-crossing Lemma 43. Then we apply Proposition 46 and the definition of the map $K$. In particular, $P_{S}$ also does not intersect the arc $A$, so there is no curve counted by $K_{o h}$.

Lemma 57. When restricted to the subcomplex $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$, we have the equation

$$
\begin{equation*}
U_{1} \otimes i d+i d \otimes U_{2}+\frac{1}{h} \partial_{h o}=\partial_{o o} K_{o o}+K_{o o} \partial_{o o} \tag{3.6.1}
\end{equation*}
$$

Proof. The chain homotopy $\widehat{K}$ induced by the change of base points in [66] gives us that

$$
\begin{equation*}
\left(U_{\#, z_{1}}\right)_{o o}+\left(U_{\#, z_{2}}\right)_{o o}=\partial_{o o} \widehat{K}_{o o}+\widehat{K}_{o o} \partial_{o o}+\partial_{o h} \widehat{K}_{h o}+\widehat{K}_{o h} \partial_{h o}, \tag{3.6.2}
\end{equation*}
$$

where $\widehat{K}$ is the degree -1 map defined by counting $I=1$ curves passing through the path $\{0\} \times(A \cup\{p\})$ as above. Restricting (3.6.2) to appropriate ends up with action $<L$ and applying Lemma 50 gives us that

$$
\begin{equation*}
U_{1} \otimes i d+i d \otimes U_{2}=\partial_{o o} \widehat{K}_{o o}+\widehat{K}_{o o} \partial_{o o}+\partial_{o h} \widehat{K}_{h o}+\widehat{K}_{o h} \partial_{h o} \tag{3.6.3}
\end{equation*}
$$

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 CONTACT HOMOLOGYRecall by Lemma 41, we have that $E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right) \cong C_{o}^{L} \oplus C_{h}^{L}=C o n e^{L}$ as vector spaces as in (3.4.1). By the definition of $K$ and Lemma 43, we immediately have that $K_{o o}=\widehat{K}_{o o}$ on the filtered complex Cone ${ }^{L}$. By Lemma 51, we have that

$$
K_{o h}=\widehat{K}_{o h}-\frac{1}{h}
$$

on $C o n e^{L}$. Moreover, by Lemma 56, we have that $K_{o h}=0$ hence

$$
\widehat{K}_{o h}=\frac{1}{h}
$$

on Cone ${ }^{L}$. Recall $\partial_{o h}=0$ by Corollary 47 on Cone ${ }^{L}$. Therefore, we obtain Equation (3.6.1) on Cone $^{L}=E C C^{L}\left(Y_{1} \#_{R(L)} Y_{2}, \alpha_{R(L)}\right)$.

Remark 58. There is another proof of Lemma 57 by directly analyzing the possible buildings in the breaking when one considers SFT compactness as following. Consider
$\mathcal{M}_{2}^{L}:=\left\{I=2\right.$ curves passing through $\left[z_{1}, p\right) \cup\left(p, z_{2}\right]$ with ends on orbit sets of actions $\left.<L\right\}$,
which is compact because the corresponding moduli space with curves passing through $p$ is compact (see Section 2.5 in [66]). We consider the boundary of $\mathcal{M}_{2}^{L}$. One may show that gluing the two $I=1$ levels gives back a curve passing through the same side of the connected sum $S^{2}$. Then, the boundary of $\mathcal{M}_{2}^{L}$ constitutes buildings $\partial K$, $K \partial, I=2$ curves passing through basepoint $z_{1}$ or $z_{2}$, and $I=2$ buildings that pass through $p$. Now the $I=2$ buildings that pass through $p$ are exactly the buildings that have two $I=1$ levels, and the bottom level constitutes of trivial cylinders and $P_{N}$, by Lemma 51. See Figure 3.2. Now since $\mathcal{M}_{2}^{L}$ is a compact 1-dimensional manifold, we have that

$$
\partial \mathcal{M}_{2}^{L}=(\partial K)_{o o}+(K \partial)_{o o}+\left(U_{\#, z_{1}}\right)_{o o}+\left(U_{\#, z_{2}}\right)_{o o}+\frac{1}{h} \partial_{h o}=0
$$

which is equivalent to (3.6.1).

## Proof of the main theorem

Now we construct a chain homotopy equivalence $F:\left(C_{\text {cone }}^{L}, \partial_{\text {cone }}\right) \rightarrow\left(C_{\#}^{L}, \partial_{\#}\right)$ in order to conclude the proof of Proposition 42.

Lemma 59. The map $F=\left(\begin{array}{cc}i d & 0 \\ h K_{o o} & h\end{array}\right)$ is a chain homotopy equivalence.

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Figure 3.2: An example of an $I=2$ building that passes through $p$.

Proof. Since the diagonal constitutes of the identity maps on vector spaces and $F$ is lower triangular, it suffices to show that $F$ is a chain map, i.e. $F \partial_{\text {cone }}=\partial_{\#} F$.

This amounts to checking the following equation:

$$
\left(\begin{array}{cc}
i d & 0 \\
h K_{o o} & h
\end{array}\right)\left(\begin{array}{cc}
\partial_{o o} & 0 \\
U_{1} \otimes i d+i d \otimes U_{2} & \partial_{o o}
\end{array}\right)=\left(\begin{array}{cc}
\partial_{o o} & \partial_{o h} \\
\partial_{h o} & \partial_{h h}
\end{array}\right)\left(\begin{array}{cc}
i d & 0 \\
h K_{o o} & h
\end{array}\right)
$$

Or, equivalently:

1. $i d \circ \partial_{o o}+0=\partial_{o o} \circ i d+\partial_{o h} h K_{o o}$. This is true by Corollary 47 .
2. $0+0=0+\partial_{\text {oh }} h$. This is true by Corollary 47 .
3. $h K_{o o} \partial_{o o}+h\left(U_{1} \otimes i d+i d \otimes U_{2}\right)=\partial_{h o} \circ i d+\partial_{h h} h K_{o o}$. This is true by Corollary 47, Corollary 49 and Lemma 57.
4. $0+h \partial_{o o}=0+\partial_{h h} h$. This is true by Corollary 49 .

Proof of Proposition 42. The inverse of $F$ in Lemma 59 gives the chain homotopy equivalence $f$.

### 3.7 The direct limit

We have shown the filtered version Proposition 42 of the main theorem. Now to obtain Theorem 1, we use a direct limit argument similar to that in [94] which involves Seiberg-Witten theory. First, we construct an isomorphism on the cobordism that

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 CONTACT HOMOLOGYinterpolates between connected sums with necks of different radii up to a certain threshold. The following lemma is useful for constructing the cobordism.

Lemma 60. Let $\alpha$ and $\alpha^{\prime}$ be two contact forms of the same contact structure $\xi$ on a contact manifold $Y$, then the convex combination $\lambda:=(1-s) \alpha+s \alpha^{\prime}$, where $s \in[0,1]$, is a contact form.

Proof. We have that

$$
d \lambda=-d s \wedge \alpha+(1-s) d \alpha+d s \wedge \alpha^{\prime}+s d \alpha^{\prime}
$$

and
$\lambda \wedge d \lambda=(1-s)^{2} \alpha \wedge d \alpha+\alpha \wedge d s \wedge \alpha^{\prime}+(1-s) s \alpha \wedge d \alpha^{\prime}+s(1-s) \alpha^{\prime} \wedge d \alpha+s^{2} \alpha^{\prime} \wedge d \alpha^{\prime}$.
Now since $\operatorname{ker} \alpha=\operatorname{ker} \alpha^{\prime}$, we have that $\alpha=\varphi \alpha^{\prime}$ for some positive function $\varphi: Y \rightarrow$ $\mathbb{R}_{>0}$. Therefore,

$$
\lambda \wedge d \lambda=(1-s)\left(1+\frac{s}{\varphi}\right) \alpha \wedge d \alpha+s(s+(1-s) \varphi) \alpha^{\prime} \wedge d \alpha^{\prime}>0
$$

since $\alpha$ and $\alpha^{\prime}$ are contact forms.
Proposition 61. Fix $L$ such that no orbit sets in $Y_{1}$ and $Y_{2}$ have action $L$. Then for

$$
r^{\prime}<r \leq R(L)
$$

there is a cobordism map

$$
\begin{equation*}
g_{r, r^{\prime}}^{L}: E C H_{*}^{L}\left(Y_{1} \#_{r} Y_{2}, \alpha_{r}\right) \rightarrow E C H_{*}^{L}\left(Y_{1} \#_{r^{\prime}} Y_{2}, \alpha_{r^{\prime}}\right) \tag{3.7.1}
\end{equation*}
$$

which is an isomorphism.
Proof. Consider the admissible deformation

$$
\begin{equation*}
\rho:=\left\{\left(\lambda_{t}:=(1-t) \alpha_{r^{\prime}}+t \alpha_{r}, L, J_{t}, \mu\right) \mid t \in[0,1]\right\}, \tag{3.7.2}
\end{equation*}
$$

where $\mu$ is sufficiently large. Note that $\lambda_{t}$ is a contact form since a convex combination of contact forms of the same contact structure is still a contact form by Lemma 60.

Now, observe by Lemma 41 that for $r^{\prime}<r \leq R(L)$, we have that any orbit intersecting the connected sum $S^{2}$ has to have action $\geq L$, except for the special hyperbolic orbit whose action stays $\ll L$. Since the symplectic cobordism ( $[0,1] \times$ $\left.\left(Y_{1} \# Y_{2}\right), d\left(e^{t} \lambda_{t}\right)\right)$ does not change orbits away from the connected sum region by the proof of Lemma 24, no orbit set of action equal to $L$ appears for any $\lambda_{t}$.

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 CONTACT HOMOLOGYBy Lemma 3.7 and Section 3.5 in [64], since $\lambda_{t}$ has no orbit sets of action $L$ and we may perturb $J$ to be $E C H^{L}$-generic, we have that for each $\Gamma \in H_{1}(Y)$, a well-defined isomorphism

$$
\begin{equation*}
E C H_{*}^{L}\left(Y, \lambda_{t}, \Gamma\right) \cong \widehat{H M}_{L}^{-*}\left(Y, \lambda_{t}, \mathfrak{s}_{\xi, \Gamma}\right) \tag{3.7.3}
\end{equation*}
$$

where $\mathfrak{s}_{\xi, \Gamma}$ is the spin-c structure $\mathfrak{s}(\xi)+\mathrm{PD}(\Gamma)$. For the definition of $\widehat{H M}_{L}^{-*}\left(Y, \lambda_{t}, \mathfrak{s}_{\xi, \Gamma}\right)$, see [64], whose exact definition we do not need in this proof.

By Lemma 3.4 in [64], the admissible deformation $\rho$ gives an isomorphism

$$
\begin{equation*}
\Phi: \widehat{H M}_{L}^{-*}\left(Y, \mathfrak{s}_{\xi, \Gamma} ; \lambda_{0}, J_{0}, \mu\right) \xrightarrow{\simeq} \widehat{H M}_{L}^{-*}\left(Y, \mathfrak{s}_{\xi, \Gamma} ; \lambda_{1}, J_{1}, \mu\right) . \tag{3.7.4}
\end{equation*}
$$

Composing the above isomorphisms (3.7.3) and (3.7.4) gives the desired map $g_{r, r^{\prime}}^{L}$. It is shown in Section 3.3 in [64] that $g_{r, r^{\prime}}^{L}$ is well-defined.

Given $L$, when $r>R(L)$, we no longer have the nice control of $E C C^{L}\left(Y_{1} \#_{r} Y_{2}\right)$, since there might be orbits of large action (despite being $<L$ ) that cross the connected sum $S^{2}$. However, we still have the following cobordism map induced by an exact symplectic cobordism given in [64].

Theorem 62 (Corollary 5.3(a) in [64]). Let $(X, \lambda)$ be an exact symplectic cobordism from $\left(Y_{+}, \lambda_{+}\right)$to $\left(Y_{-}, \lambda_{-}\right)$, where $\lambda_{ \pm}$is L-nondegenerate. Let $J_{ \pm}$be a symplectizationadapted almost complex structure for $\lambda_{ \pm}$. Suppose $\rho$ is sufficiently large. Fix appropriate 2 -forms $\mu_{ \pm}$and perturbations needed to define the chain complex $\widehat{C M}^{*}\left(Y_{ \pm} ; \lambda_{ \pm}, J_{ \pm}, \rho\right)$. Then there is a well-defined map

$$
\begin{equation*}
\widehat{H M}_{L}^{*}(X, \lambda): \widehat{H M}_{L}^{*}\left(Y_{+} ; \lambda_{+}, J_{+}, \rho\right) \longrightarrow \widehat{H M}_{L}^{*}\left(Y_{-} ; \lambda_{-}, J_{-}, \rho\right) \tag{3.7.5}
\end{equation*}
$$

depending only on $X, \lambda, L, \rho, J_{ \pm}, \mu_{ \pm}$and the perturbations, such that

$$
\begin{gather*}
\widehat{H M}_{L}^{*}\left(Y_{+} ; \lambda_{+}, J_{+}, \rho\right) \xrightarrow{i^{L, L^{\prime}}} \widehat{H M}_{L^{\prime}}^{*}\left(Y_{+} ; \lambda_{+}, J_{+}, \rho\right) \\
\widehat{H M}_{L}^{*}(X, \lambda) \downarrow  \tag{3.7.6}\\
\widehat{H M}_{L}^{*}\left(Y_{-} ; \lambda_{-}, J_{-}, \rho\right) \xrightarrow{{ }^{i^{L, L^{\prime}}}{ }^{\prime}(X, \lambda)} \widehat{H M}_{L^{\prime}}^{*}\left(Y_{-} ; \lambda_{-}, J_{-}, \rho\right)
\end{gather*}
$$

commutes for $L<L^{\prime}$ and $\lambda_{ \pm}$which are $L$-nondegenerate, where $i^{L, L^{\prime}}$ are induced by inclusions of chain complexes.

Therefore, we may consider the product exact symplectic cobordism

$$
\left(X:=[0,1]_{t} \times\left(Y_{1} \# Y_{2}\right), \lambda_{t}\right),
$$

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 CONTACT HOMOLOGYwhere $\lambda_{t}$ is defined in (3.7.2), from $Y_{1} \#_{r^{\prime}} Y_{2}$ to $Y_{1} \#_{r} Y_{2}$ where $r^{\prime}<r$. After passing through the isomorphism to Seiberg-Witten as in (3.7.3), Theorem 62 gives the following cobordism map:

$$
\begin{equation*}
E C H_{*}^{L}(X, \lambda): E C H_{*}^{L}\left(Y_{1} \#_{r} Y_{2}, \alpha_{r}\right) \longrightarrow E C H_{*}^{L}\left(Y_{1} \#_{r^{\prime}} Y_{2}, \alpha_{r^{\prime}}\right) \tag{3.7.7}
\end{equation*}
$$

Going back to the case when $r \leq R(L)$, the cobordism map induced by admissible deformation in Proposition 61, together with the cobordism map induced by inclusion, gives the following commutative diagram (as in proof of Lemma 3.7 in [64]) for $L_{1}<L_{2}$ and $R\left(L_{1}\right)>R\left(L_{2}\right)$ :

$$
\begin{align*}
& E C H_{*}^{L_{1}}\left(Y_{1} \#_{R\left(L_{1}\right)} Y_{2}, \alpha_{R\left(L_{1}\right)}\right) \xrightarrow{i^{L_{1}, L_{2}}} E C H_{*}^{L_{2}}\left(Y_{1} \#_{R\left(L_{1}\right)} Y_{2}, \alpha_{R\left(L_{1}\right)}\right) \\
& \quad g_{R\left(L_{1}\right), R\left(L_{2}\right) \downarrow}^{L_{1}} \downarrow  \tag{3.7.8}\\
& E C H_{*}^{L_{1}}\left(Y_{1} \#_{R\left(L_{2}\right)} Y_{2}, \alpha_{R\left(L_{2}\right)}\right) \xrightarrow{g_{R\left(L_{1}\right), R\left(L_{2}\right)}^{L_{1}, L_{2}}} E C H_{*}^{L_{2}}\left(Y_{1} \#_{R\left(L_{2}\right)} Y_{2}, \alpha_{R\left(L_{2}\right)}\right),
\end{align*}
$$

where $i^{L_{1}, L_{2}}$ are inclusion induced cobordism maps defined in Theorem 21. Furthermore, by Lemma 5.6 in [64], the map $g_{R\left(L_{1}\right), R\left(L_{2}\right)}^{L_{i}}$ identifies with the exact cobordism map $E C H_{*}^{L}(X, \lambda)$ in (3.7.7) when we consider $L=L_{i}, r=R\left(L_{1}\right)$ and $r^{\prime}=R\left(L_{2}\right)$.

Now we check that (3.7.8) gives a directed system, similar to that in [94]. We need to check that we have a well-defined composition for the cobordism maps

$$
\begin{equation*}
\Phi^{L_{1}, L_{2}}\left(R\left(L_{1}\right), R\left(L_{2}\right)\right): E C H_{*}^{L_{1}}\left(Y_{1} \#_{R\left(L_{1}\right)} Y_{2}, \alpha_{R\left(L_{1}\right)}\right) \rightarrow E C H_{*}^{L_{2}}\left(Y_{1} \#_{R\left(L_{2}\right)} Y_{2}, \alpha_{R\left(L_{2}\right)}\right) \tag{3.7.9}
\end{equation*}
$$

defined by either path in (3.7.8). For $L_{1}<L_{2}<L_{3}$, we define $r:=R\left(L_{1}\right), r^{\prime}:=$ $R\left(L_{2}\right)$ and $r^{\prime \prime}:=R\left(L_{3}\right)$. Then the composition of (3.7.9) is given by the following four-fold commutative diagram.


Now, we may pass the filtered ECH complexes to the direct limit with respect to the above maps. This requires some algebraic manipulations as in proof of Theorem 7.1 in [94], with which we will conclude the proof of our main theorem.

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 CONTACT HOMOLOGYProof of Theorem 1. In this proof we suppress the notation of the connected sum radius in $Y_{1} \# Y_{2}$ when it is encoded in the contact form, in order to simplify notations. Let $L(r)$ denote the value of $L$ such that $R(L)=r$, where $R(L)$ is the strictly decreasing function defined in Section 3.4. Therefore,

$$
E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{R(L)}\right)=E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right)
$$

Now,

$$
\begin{align*}
E C H\left(Y_{1} \# Y_{2}, \xi_{1} \# \xi_{2}\right) & =\lim _{r \rightarrow 0} E C H\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \\
& =\lim _{r \rightarrow 0} \lim _{L \rightarrow \infty} E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \\
& =\lim _{r \rightarrow 0} E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right)  \tag{3.7.10}\\
& =\lim _{L \rightarrow \infty} E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{R(L)}\right) \\
& =H_{*}\left(\operatorname{Cone}\left(U_{1} \otimes i d+i d \otimes U_{2}\right)\right) .
\end{align*}
$$

The first equation is because ECH is independent of the choice of contact forms [109]. The second equation is given by (3.2.6) and Theorem 21(1). The third equation comes from consideration of the following map:

$$
\begin{equation*}
\Psi: \lim _{r \rightarrow 0} E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \longrightarrow \lim _{r \rightarrow 0} \lim _{L \rightarrow \infty} E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \tag{3.7.11}
\end{equation*}
$$

by sending the equivalence class of an element $d_{r} \in E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right)$ under $\lim _{r \rightarrow 0}$ to the equivalence class of $d_{r}$ under $\lim _{r \rightarrow 0} \lim _{L \rightarrow \infty}$. We need to show that $\Psi$ is well-defined and a bijection. To establish well-definedness, consider $d_{r}$ and $d_{r^{\prime}}$ in $\lim _{r \rightarrow 0} E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right)$ where $r>r^{\prime}$ and $d_{r} \sim d_{r^{\prime}}$. This means that there is a common element $d_{r^{\prime \prime}} \in E C H^{L\left(r^{\prime \prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r^{\prime \prime}}\right)$ such that both $d_{r}$ and $d_{r^{\prime}}$ are mapped to under the direct limit. Then $\Psi\left(d_{r}\right) \sim \Psi\left(d_{r^{\prime}}\right)$ by composing the maps

$$
\begin{equation*}
E C H^{L(r)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \longrightarrow E C H^{L\left(r^{\prime \prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r}\right) \longrightarrow E C H^{L\left(r^{\prime \prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r^{\prime \prime}}\right) \tag{3.7.12}
\end{equation*}
$$

and

$$
\begin{equation*}
E C H^{L\left(r^{\prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r^{\prime}}\right) \longrightarrow E C H^{L\left(r^{\prime \prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r^{\prime}}\right) \longrightarrow E C H^{L\left(r^{\prime \prime}\right)}\left(Y_{1} \# Y_{2}, \alpha_{r^{\prime \prime}}\right) \tag{3.7.13}
\end{equation*}
$$

given by the right followed by down composition of the commutative diagram (3.7.8). Now we show that $\Psi$ is injective. Suppose $\Psi(d)=0$. That means there exists $L_{0}$ such that a representative $\widetilde{\Psi(d)} \in E C H^{L_{0}}\left(Y_{1} \# Y_{2}, \alpha_{r_{0}}\right)$, where $[\widetilde{\Psi(d)}]=\Psi(d)$, is zero for some $r_{0}$. See Figure 3.3. If $r_{0} \leq R\left(L_{0}\right)$, then we are done, since for $r \leq R(L)$,

$$
E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{r}\right)=E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{R(L))}\right)
$$

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Figure 3.3: Schematic illustration of the proof involving direct limits in Theorem 1. In the region above $R(L)$, i.e. when $r \leq R(L)$, we have isomorphisms of $E C H_{*}^{L}\left(Y_{1} \#_{r} Y_{2}, \alpha_{r}\right)$ given a fixed value of $L$ by Proposition 61.
by Proposition 61, so (3.7.11) is a bijection. Suppose $r_{0}>R\left(L_{0}\right)$. Then $\widetilde{\Psi(d)}$ is mapped to zero in $E C H^{L_{0}}\left(Y_{1} \# Y_{2}, \alpha_{R\left(L_{0}\right)}\right)$ under the map in (3.7.7). Therefore, $d \sim 0$. To show that $\Psi$ is surjective, let

$$
y \in \lim _{L \rightarrow \infty} E C H^{L}\left(Y_{1} \# Y_{2}, \alpha_{r_{0}}\right)
$$

for some $r_{0}$. Similar to the injective case, if $r_{0} \leq R\left(L_{0}\right)$, then we are done. Suppose $r_{0}>R\left(L_{0}\right)$. Now let $x \in E C H^{L_{0}}\left(Y_{1} \# Y_{2}, \alpha_{r_{0}}\right)$ be such that $[x]=y$. Let

$$
x^{\prime} \in E C H^{L_{0}}\left(Y_{1} \# Y_{2}, \alpha_{R(L)}\right)
$$

be the image of $x$ when taking the limit as $r \rightarrow 0$ defined by the exact cobordism map as in (3.7.7). Then $[y]=\Psi\left(\left[x^{\prime}\right]\right)$.

The fourth equation is by the definition of $L(r)$. The fifth equation is by Proposition 42 and the fact that taking direct limit commutes with taking homology.

Finally, note that the mapping cone in (3.7.10) is of the map
$U_{1} \otimes i d+i d \otimes U_{2}: E C C\left(Y_{1}, \lambda_{1}\right) \otimes_{\mathbb{F}} E C C\left(Y_{2}, \lambda_{2}\right) \longrightarrow E C C\left(Y_{1}, \lambda_{1}\right) \otimes_{\mathbb{F}} E C C\left(Y_{2}, \lambda_{2}\right)[-1]$.
Over $\mathbb{F}$, the homology of $\operatorname{Cone}\left(U_{1} \otimes i d+i d \otimes U_{2}\right)$ is isomorphic to the homology of the mapping cone of the induced map $\left(U_{1} \otimes i d+i d \otimes U_{2}\right)_{*}$ on homology, by observing that Künneth formula is natural over $\mathbb{F}$. This concludes the proof of the main theorem.

## Chapter 4

## Legendrian embedded contact homology

Julian Chaidez, Oliver Edtmair, Luya Wang, Yuan Yao, Ziwen Zhao

### 4.1 Introduction

The embedded contact homology of a closed contact 3-manifold $(Y, \xi)$ is a Floer theoretic invariant originally introduced by Hutchings [52]. These homology groups, denoted by

$$
E C H(Y, \xi)
$$

are, roughly speaking, computed as the homology of a chain complex freely generated by certain finite sets of closed simple Reeb orbits (with multiplicity). The differential counts certain embedded, possibly disconnected $J$-holomorphic curves in the symplectization of the contact manifold that can have arbitrary genus.

Since its inception, embedded contact homology and its associated invariants have been applied to many of the fundamental problems in 3-dimensional Reeb dynamics, to dramatic effect. These applications include a simple formal proof of the 3-dimensional Weinstein conjecture by [114], and the Arnold chord conjecture by Hutchings-Taubes [63, 65], and the existence of broken book decompositions [18]. The spectral invariants derived from the ECH groups (and related PFH groups), called the ECH capacities, have been used to prove the smooth closing lemma for Reeb flows [67] and area preserving surface maps [31, 33], and many results on symplectic embeddings [90, 54, 30].

Embedded contact homology can morally be regarded as a flavor of symplectic field theory (SFT), as formulated by Eliashberg-Givental-Hofer [36]. SFT is, broadly speaking, a framework for constructing invariants of contact manifolds in any dimension, acquired from chain complexes generated by Reeb orbits. Variants of SFT include cylindrical contact homology (cf. [hn2014, 93]), linearized contact homology (cf. [12]) and the contact homology algebra (cf. [101]).

Many flavors of SFT can be extended to invariants of a pair $(Y, \Lambda)$ of a closed contact manifold $Y$ and a closed Legendrian sub-manifold $\Lambda \subset Y$. The goal of this paper is to initiate the study of a corresponding Legendrian version of embedded contact homology.

## Standard ECH

We begin this introduction with a review of standard ECH, largely drawing on [52]. We discuss many of the key ideas that allow one to define ECH and to show that the differential in ECH, which counts certain $J$-holomorphic curves, defines a chain complex. Besides serving as a review, this will also provide a road map of the results required to construct Legendrian ECH.

## Holomorphic Currents

Let $(Y, \lambda)$ be a closed contact 3-manifold with a non-degenerate contact form $\lambda$. We start by considering the symplectization of $Y$, i.e. the cylindrical manifold

$$
\mathbb{R}_{s} \times Y
$$

There is a natural class of translation invariant complex structures $J$ on $\mathbb{R} \times Y$, which sends $\partial_{s}$ to the Reeb vector-field $R$ of $\lambda$ and positively preserves the contact structure $\operatorname{ker} \lambda$. We call such almost complex structure $\lambda$-adapted.

Recall that a $J$-holomorphic curve $C$ from a punctured Riemann surface (without boundary) is an equivalence class of map

$$
u:(\Sigma, j) \rightarrow(\mathbb{R} \times Y, J) \quad \text { satisfying } \quad d u \circ j=J \circ d u
$$

modulo holomorphic reparametrization of the domain $(\Sigma, j)$. We say that $C$ is proper if the map $u$ is proper and finite energy if the integral of $u^{*} d \lambda$ is finite. If $u$ is proper and finite energy, then $u$ converges (in an appropriate sense) to the cylinder over a collection of Reeb orbits

$$
\Gamma_{+}=\left(\gamma_{1}^{+}, \ldots, \gamma_{k}^{+}\right) \quad \text { and } \quad \Gamma_{-}=\left(\gamma_{1}^{-}, \ldots, \gamma_{k}^{-}\right)
$$

near $+\infty \times Y$ and $-\infty \times Y$, respectively. The Fredholm index $\operatorname{ind}(u)$ of $u$ is the Fredholm index of a certain linearized Cauchy-Riemann operator associated to $u$. It is given by the formula
$\operatorname{ind}(u)=-\chi(\Sigma)+2 c_{\tau}(u)+\mathrm{CZ}_{\tau}^{\text {ind }}\left(\Gamma_{+}\right)-\mathrm{CZ}_{\tau}^{\text {ind }}\left(\Gamma_{-}\right) \quad$ where $\quad \mathrm{CZ}_{\tau}^{\text {ind }}\left(\Gamma_{ \pm}\right):=\sum_{i} C Z_{\tau}\left(\gamma_{i}^{ \pm}\right)$
Here $c_{\tau}(u)$ is the relative first Chern number of $u^{*} \xi$ with respect to a trivialization $\tau$ of $\xi$ along $\Gamma_{+}$and $\Gamma_{-}$, and $C Z_{\tau}(\gamma)$ is the Conley-Zehnder index of the linearized Reeb flow around $\gamma$ in the trivialization $\tau$. Standard transversality result states that there is a generic class of regular $J$ such that, if $u$ is somewhere injective, the moduli space of proper, finite energy, somewhere injective $J$-holomorphic maps $v$ near $u$ asymptotic to $\Gamma_{ \pm}$at $\pm \infty$ is a manifold of dimension ind $(u)$.

A $J$-holomorphic current $\mathcal{C}=\left\{\left(C_{i}, m_{i}\right)\right\}$ in $Y$ is a finite collection of connected, proper, finite energy, somewhere injective $J$-holomorphic curves $C_{i}$ and positive integer multiplicities $m_{i}$.


Figure 4.1: An depiction of a $J$-holomorphic current.

An orbit set $\Theta=\left\{\left(\gamma_{i}, n_{i}\right)\right\}$ is, similarly, a collection of simple closed Reeb orbits $\gamma_{i}$ at $\pm \infty$. In analogy to curves, every $J$-holomorphic current is asymptotic orbit sets at $\pm \infty$. We note here we must count with multiplicities, e.g. if the $J$-holomorphic current $(C, 1)$ is positively asymptotic to the simple Reeb orbit $\gamma$, then $(C, 2)$ is positively asymptotic to $(\gamma, 2)$. We denote the moduli space of $J$-holomorphic currents asymptotic to $\alpha$ at $+\infty$ and $\beta$ at $-\infty$ by

$$
\mathcal{M}(\alpha, \beta)
$$

Note that this moduli space admits an $\mathbb{R}$-action given by $\mathbb{R}$-translation in $\mathbb{R} \times Y$. We denote the quotient by $\mathcal{M}(\alpha, \beta) / \mathbb{R}$.

## ECH Index

The key ingredient of ECH that differentiates it from other versions of symplectic field theory is the ECH index. It may be viewed as a map

$$
I: \mathcal{M}(\alpha, \beta) \rightarrow \mathbb{Z}
$$

On a $J$-holomorphic current $\mathcal{C}$ asymptotic to $\alpha$ at $+\infty$ and $\beta$ at $-\infty$, the ECH index is given by

$$
\begin{equation*}
I(\mathcal{C})=c_{\tau}(\mathcal{C})+Q_{\tau}(\mathcal{C})+\mathrm{CZ}_{\tau}^{E C H}(\alpha)-\mathrm{CZ}_{\tau}^{E C H}(\beta) . \tag{4.1.1}
\end{equation*}
$$

Here $Q_{\tau}$ is the relative self-intersection number [50, §2.4] that counts intersections between $\mathcal{C}$ and a push-off of $\mathcal{C}$ determined by $\tau$, and $\mathrm{CZ}_{\tau}^{E C H}$ is a Conley-Zehnder index term given by

$$
\begin{equation*}
\mathrm{CZ}_{\tau}^{E C H}(\Theta)=\sum_{i=1}^{k} \sum_{j=1}^{m_{i}} \mathrm{CZ}\left(\gamma_{i}^{j}\right) \quad \text { for } \quad \Theta=\left\{\left(\gamma_{i}, m_{i}\right)\right\} \tag{4.1.2}
\end{equation*}
$$

The fundamental property that the ECH index satisfies is the following index inequality.

Theorem 1 (Index Inequality). [50,55] Let $C$ is a somewhere injective $J$-holomorphic curve in $\mathbb{R} \times Y$, for a compatible $J$ and let $\delta(C)$ denote the count of singularities (with multiplicity) of $C$. Then

$$
\operatorname{ind}(C) \leq I(C)-2 \delta(C)
$$

This inequality places stringent constraints on curves and currents of low ECH index. Let us discuss the main ingredients of the proof of this inequality, as their Legendrian analogues will be the main topic of this paper.

The first ingredient of Theorem 1 is the writhe bound. Recall that the writhe $w(\zeta)$ of a braid $\zeta$ in $S^{1} \times D^{2}$ is an isotopy invariant that can be computed as a signed count of the self-intersections of the image $\pi(\zeta)$ under the projection $\pi: S^{1} \times D^{2} \rightarrow$ $S^{1} \times[-1,1]$. If $C$ is a somewhere injective J-holomorphic curve asymptotic to (covers of) a simple orbit $\gamma$ at some subset of its punctures, then $\mathcal{C}$ determines a braid in a tubular neighborhood of $\gamma$ (due to $J$-holomorphic curve asymptotics established by Siefring [108]).


Figure 4.2: The braid near a simple orbit $\gamma$ determined by a curve (or current) $C$.

Given a choice of trivialization $\tau$ of $\xi$, any braid near $\gamma$ is identified with a braid in $S^{1} \times D^{2}$. The writhe of this braid is denoted by $w_{\tau}(\mathcal{C} ; \gamma)$, and we define the writhe of $\mathcal{C}$ as

$$
w_{\tau}(\mathcal{C})=\sum_{(\gamma, m) \in \alpha} w_{\tau}(\mathcal{C} ; \gamma)-\sum_{(\gamma, m) \in \beta} w_{\tau}(\mathcal{C} ; \gamma)
$$

The writhe bound estimates the writhe of $\mathcal{C}$ in terms of the difference between the Conley-Zehnder terms in the ECH index and Fredhom index.

Theorem 2 (Writhe Bound). [55] Let $C$ be a somewhere injective $J$-holomorphic curve in $\mathbb{R} \times Y$. Then

$$
w_{\tau}(C) \leq C Z_{\tau}^{E C H}(C)-C Z_{\tau}^{\text {ind }}(C)
$$

The second ingredient to Theorem 1 is the adjunction inequality. This is a version of the classical adjunction inequality in complex geometry, tailored to the setting of ECH.

Theorem 3 (Adjunction). Let $C$ be a somewhere injective $J$-holomorphic curve in $\mathbb{R} \times Y$. Then

$$
c_{\tau}(C)=\chi(C)+Q_{\tau}(C)+w_{\tau}(C)-2 \delta(C)
$$

The index inequality is a short calculation using these two bounds.

## ECH Complex

We are now ready to explain the construction of the ECH chain complex, using the properties of the ECH index discussed above.

Remark 63 (Orientations). In this paper, we work without orientation for simplicity. In particular, we use $\mathbb{Z} / 2$-coefficient to define all ECH groups.

An ECH generator is an orbit set $\Theta$ such that any hyperbolic orbit $\gamma_{i}$ in $\Theta$ has multiplicity 1 . The ECH chain complex is simply the free vector-space over $\mathbb{Z} / 2$ generated by ECH generators.

$$
E C C(Y, \lambda):=\mathbb{Z} / 2\langle\mathrm{ECH} \text { generators } \Theta \text { of }(Y, \lambda)\rangle
$$

The differential on the ECH complex is defined by counting currents. To formalize this, we require the following classification of low ECH index currents deducible from the index bound.

Theorem 4 (Low-Index Currents). Let $J$ be a regular, compatible almost complex structure on $\mathbb{R} \times Y$ and let $\mathcal{C}$ be a $J$-holomorphic current of ECH index $I(\mathcal{C}) \leq 2$. Then

- $I(\mathcal{C}) \geq 0$ with equality only if $\mathcal{C}$ is a union of trivial cylinders.
- If $I(\mathcal{C})=1$ then $\mathcal{C}=C \sqcup \mathcal{T}$ where $C$ is Fredholm index 1 and embedded, and $\mathcal{T}$ is a union of trivial cylinders with multiplicity.
- If $I(\mathcal{C})=2$ and is asymptotic to ECH generators at $\pm \infty$, then $\mathcal{C}=C \sqcup \mathcal{T}$ where $C$ is Fredholm index 2 and embedded, and $\mathcal{T}$ is a union of trivial cylinders with multiplicity.

This classification can, in turn, be used to deduce compactness properties of low ECH index moduli spaces.

Theorem 5. [52, §5.3-5.4] Let $J$ be a regular, compatible almost complex structure on $\mathbb{R} \times Y$ and let $\mathcal{M}_{k}(\Theta, \Xi)$ be the space of ECH index $k J$-holomorphic currents from $\Theta$ to $\Xi$. Then

- the space $\mathcal{M}_{1}(\Theta, \Xi) / \mathbb{R}$ is 0 -dimensional and compact
- the space $\mathcal{M}_{2}(\Theta, \Xi) / \mathbb{R}$ is a 1-manifold with a compact truncation ${ }^{1} \mathcal{M}_{2}^{\prime}(\Theta, \Xi) / \mathbb{R}$ with a map

$$
\Pi: \partial \mathcal{M}_{2}^{\prime}(\Theta, \Xi) / \mathbb{R} \rightarrow \bigsqcup_{\Theta^{\prime}} \mathcal{M}_{1}\left(\Theta, \Theta^{\prime}\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta^{\prime}, \Xi\right) / \mathbb{R}
$$

[^2]- The inverse image $\Pi^{-1}\left(\mathcal{C}, \mathcal{C}^{\prime}\right)$ of a pair of currents in $\mathcal{M}_{1}\left(\Theta, \Theta^{\prime}\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta^{\prime}, \Xi\right) / \mathbb{R}$ has an odd number of points if and only if the orbit set $\Theta^{\prime}$ is an ECH generator.

Most of Theorem 5 follows from Theorem 4, a type of Gromov compactness due to Taubes and a bound on the topological complexity of low ECH index curves [52]. However, the last point require a delicate obstruction bundle gluing argument that is well beyond the scope of this introduction. However, we will remark that this analysis requires certain partition conditions obeyed by low ECH index currents, which restrict the braids that can appear at their ends. For a more detailed explanation of partition conditions, see [52].

By applying Theorem 5, one can simply define the ECH differential as the count of ECH index 1 curves, modulo reparametrization, for a regular choice of $J$.
$\partial: E C C(Y, \lambda) \rightarrow E C C(Y, \lambda)$ be given by $\partial \Theta=\sum_{\Xi} \#\left(\mathcal{M}_{1}(\Theta, \Xi) / \mathbb{R}\right) \cdot \Xi \bmod 2$
It is a simple consequence of Theorem 5 that $\partial$ is well-defined and that $\partial \circ \partial=0$.

## Legendrian ECH And Main Results

We now move to the main topic of this paper, providing an overview of the construction of Legendrian ECH. We shall see that all of the constructions in ordinary ECH have generalizations to the Legendrian setting.

## Holomorphic Currents With Boundary

Let $Y$ be a contact 3-manifold with convex sutured boundary $\partial Y$ and a non-degenerate, adapted contact form $\lambda$. We refer the reader to [28] for a detailed treatment of sutured contact manifolds. The boundary of $Y$ divides as

$$
\partial_{-} Y \quad \partial_{\sigma} Y \quad \text { and } \quad \partial_{+} Y
$$

where the Reeb vector-field is inward normal, tangent and outward normal respectively. Also fix a closed (possibly disconnected) Legendrian $\Lambda \subset \partial Y$ decomposing as

$$
\Lambda_{+} \subset \partial_{+} Y \quad \text { and } \quad \Lambda_{-} \subset \partial_{-} Y
$$

We assume that that these are exact Lagrangians in the Liouville domains $\left(\partial_{ \pm} Y, \lambda \mid \partial_{ \pm} Y\right)$, so that

$$
\left.\lambda\right|_{\partial_{ \pm} Y} \text { vanish along } \Lambda_{ \pm}
$$



Figure 4.3: A contact manifold $Y$ with convex boundary and Legendrians $\Lambda_{ \pm} \subset \partial_{ \pm} Y$.

We have included Figure 4.3 below to better illustrate this setup for the reader.
There is a natural class of tailored almost complex structures on $\mathbb{R} \times Y$, which we write as $\mathcal{J}_{T}(\mathbb{R} \times Y)$, that are compatible and also satisfy a set of assumptions near $\partial Y$ that gurantee that $J$-holomorphic curves do not cross the boundary. As defined in [28] tailoered almost complex structures exist in great abundance (see Section 4.3 for a description).

We consider finite energy, proper $J$-holomorphic maps from a punctured Riemann surface $(\Sigma, j)$ with boundary to the symplectization of $Y$, with boundary on the symplectization of $\Lambda$.

$$
\begin{equation*}
u:(\Sigma, \partial \Sigma) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda) \quad \text { with } \quad d u \circ j=J \circ d u \tag{4.1.3}
\end{equation*}
$$

A $J$-holomorphic current $\mathcal{C}$ in $(Y, \Lambda)$ is defined exactly as in the closed case, using $J$-holomorphic maps with boundary (4.1.3). These curves are now asymptotic at $\pm \infty$ to orbit-chord sets

$$
\Xi=\left\{\left(\gamma_{i}, m_{i}\right)\right\}
$$

Here $\gamma_{i}$ is either a simple closed Reeb orbit or a Reeb chord of $\Lambda$, and $m_{i}$ is a multiplicity. We continue to denote set of $J$-holomorphic currents asymptotic to orbit-chord sets $\Theta$ at $+\infty$ and $\Xi$ at $-\infty$ by

$$
\mathcal{M}(\Theta, \Xi)
$$

We will discuss the foundations of $J$-holomorphic currents with boundary in our setting more fully in Section 4.3.

## Legendrian ECH Index

We are now ready to introduce the Legendrian ECH index, generalizing the ECH index from the closed case.

Definition 6 (Definition 189). The (Legendrian) ECH index of a $J$-holomorphic current $\mathcal{C}$ in the pair $(Y, \Lambda)$ from the orbit-chord set $\Theta$ to the orbit-chord set $\Xi$ is

$$
I(\mathcal{C})=\frac{1}{2} \cdot \mu_{\tau}(\mathcal{C})+Q_{\tau}(\mathcal{C})+C Z_{\tau}^{E C H}(\Theta)-C Z_{\tau}^{E C H}(\Xi)
$$

where the terms in the index $I$ are as follows.

- $\tau$ is a trivialization of the bundle pair $(\xi, T \Lambda)$ over $\Theta$ and $\Xi$ (see Definition 92)
- $\mu_{\tau}$ is the (relative) Maslov number (see Definition 98)
- $Q_{\tau}$ is the relative self intersection number with respect to $\tau$ (see Definition 114)
- $\mathrm{CZ}_{\tau}^{E C H}$ is a Conley-Zehnder term associated to the orbit-chord set (see Definition 175).

All of the terms in Definition 6 directly correspond to (and generalize) the terms in (4.1.1). The Legendrian ECH index satisfies a direct generalization of the ECH index inequality.

Theorem 7 (Legendrian Index Inequality). Let $C$ be a somewhere injective, $J$ holomorphic curve with boundary in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ for tailored $J$. Then

$$
\operatorname{ind}(C) \leq I(C)-2 \delta(C)-\epsilon(C)
$$

Here $\delta(C)$ and $\epsilon(C)$ are the counts of the singularities of $C$ in its interior and its boundary, respectively.

This result is a consequence of Legendrian generalizations of adjunction and the writhe bound.

Theorem 8 (Legendrian Adjunction, §4.2). Let $C$ be a proper, finite energy, somewhere injective curve in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$. Then

$$
\frac{1}{2} \cdot \mu_{\tau}(C)=\bar{\chi}(C)+Q_{\tau}(C)+w_{\tau}(C)-2 \delta(C)-\epsilon(C)
$$

Note that $\bar{\chi}$ denotes a corrected Euler characteristic that only counts boundary punctures with a factor of $\frac{1}{2}$.

To study the writhe, we further adapt Siefring's [108] asymptotic formula for $J$-holomorphic curves asymptotic to chords.

Theorem 9. Let $u_{i}:[0, \infty) \times[0,1] \rightarrow \mathbb{R} \times Y$ be a collection of $J$-holomorphic strips asymptotic to a Reeb chord $\gamma(t)$. Here $\gamma$ connects between Legendrians $L_{1}$ and $L_{2}$, and $u_{i}$ maps the boundary of the strip to $\mathbb{R} \times L_{i}$. Then there exists a neighborhood $U$ of $\gamma$, a smooth embedding $\Phi: \mathbb{R} \times U \rightarrow \mathbb{R} \times Y$, proper reparametrizations $\phi_{i}$ : $\mathbb{R} \times[0,1] \rightarrow \mathbb{R} \times[0,1]$ exponentially asymptotic to the identity, and positive integers $N_{i}$ so that for large values of $s$,

$$
\Phi \circ u_{i} \circ \psi_{i}(s, t)=\left(s, t, \sum_{k=1}^{N_{i}} e^{\lambda_{i, j} s} e_{i, j}(t)\right)
$$

where $\lambda_{i, j}$ are negative eigenvalues of the asymptotic operator associated to $\gamma$, and the $e_{i, f}(t)$ are the corresponding eigenfunctions.

This theorem is discussed in detail in the Appendix. We will delay the precise statement of the writhe bound to Section 4.6.

## Legendrian ECH Complex

Finally, we present the construction of the Legendrian ECH chain complex of $(Y, \Lambda)$, mirroring the construction in the closed setting. These claims will be revisited and proven in Section 4.7.

Definition 64. An $E C H$ generator of $(Y, \Lambda)$ is an orbit-chord set $\Theta=\left\{\left(\gamma_{i}, m_{i}\right)\right\} \cup$ $\left\{\left(c_{i}, n_{i}\right)\right\}$ where

- Every hyperbolic orbit $\gamma_{i}$ has multiplicity 1.
- Every chord $c_{i}$ is multiplicity 1.
- There is at most one Reeb chord incident to $L$ in $\Theta$ for each connected component $L$ of $\Lambda$

As with the closed case, we define the differential by counting ECH index 1 curves. We again need a classification of low ECH index curves and an accompanying compactness statement.

Theorem 10 (Low-Index Currents With Boundary). Let $J$ be a regular, tailored almost complex structure on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ and let $\mathcal{C}$ be a $J$-holomorphic current of ECH index $I(\mathcal{C}) \leq 2$. Then

- $I(\mathcal{C}) \geq 0$ with equality only if $\mathcal{C}$ is a union of trivial cylinders and strips with multiplicity.
- If $I(\mathcal{C})=1$ then $\mathcal{C}=C \sqcup \mathcal{T}$ where $C$ is Fredholm index 1 and embedded, and $\mathcal{T}$ is a union of trivial cylinders and strips with multiplicity.
- If $I(\mathcal{C})=2$ and is asymptotic to ECH generators at $\pm \infty$, then $\mathcal{C}=C \sqcup \mathcal{T}$ where $C$ is Fredholm index 2 and embedded, and $\mathcal{T}$ is a union of trivial cylinders and strips with multiplicity.

Theorem 11. Let $J$ be a regular, tailored almost complex structure on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ and let $\mathcal{M}_{k}(\Theta, \Xi)$ be the space of ECH index $k J$-holomorphic currents in $(Y, \Lambda)$ from $\Theta$ to $\Xi$. Then

- the space $\mathcal{M}_{1}(\Theta, \Xi) / \mathbb{R}$ is 0 -dimensional and compact
- the space $\mathcal{M}_{2}(\Theta, \Xi) / \mathbb{R}$ is a 1-manifold with a compact truncation ${ }^{2} \mathcal{M}_{2}^{\prime}(\Theta, \Xi) / \mathbb{R}$ with a map

$$
\Pi: \partial \mathcal{M}_{2}^{\prime}(\Theta, \Xi) / \mathbb{R} \rightarrow \bigsqcup_{\Theta^{\prime}} \mathcal{M}_{1}\left(\Theta, \Theta^{\prime}\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta^{\prime}, \Xi\right) / \mathbb{R}
$$

- The inverse image $\Pi^{-1}\left(\mathcal{C}, \mathcal{C}^{\prime}\right)$ of a pair of currents in $\mathcal{M}_{1}\left(\Theta, \Theta^{\prime}\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta^{\prime}, \Xi\right) / \mathbb{R}$ has an odd number of points if and only if the orbit-chord set $\Theta^{\prime}$ is an ECH generator.

We will give detailed proofs of these claims in Section 4.7. They are analogous to the closed case.

Remark 65. Note that, due to the lack of multiply covered Reeb chords (in contrast to orbits), there is no need to carry out a new obstruction bundle gluing strategy. Instead, the work in [60] and [62] will yield the last claim in Theorem 11 after some minor modifications. We will discuss this in 4.7.

[^3]Finally, we can state the definition of Legendrian embedded contact homology.
Definition 66 (Legendrian ECH). The ECH chain complex ECC $(Y, \Lambda, \lambda)$ of $(Y, \Lambda, \lambda)$ is the free $\mathbb{Z} / 2$-module generated by ECH generators.

$$
E C C(Y, \Lambda, \lambda)=\mathbb{Z} / 2\langle\mathrm{ECH} \text { generators } \Theta\rangle
$$

The differential $\partial_{J}$ with respect to a regular, tailored almost complex structure $J$ on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ is given by the count of ECH index $1 J$-holomorphic currents.

$$
\partial_{J} \Theta=\sum_{\Xi} \#_{2} \mathcal{M}_{1}(\Theta, \Xi) / \mathbb{R} \cdot \Xi
$$

The Legendrian embedded contact homology $\operatorname{ECH}(Y, \Lambda ; \lambda, J)$ is the homology $E C C(Y, \Lambda, \lambda)$.
We prove that the differential defined above gives rise to a chain complex (i.e. $\partial^{2}=0$ ) in Section 4.7.

Remark 67 (Invariance). If $Y$ is closed, then $\operatorname{ECH}(Y, \lambda, J)$ is independent of $\lambda$ and $J$ up to canonical isomorphism. Thus we may write

$$
E C H(Y):=E C H(Y, \lambda, J) \quad \text { for any choice of } \lambda, J
$$

This is due to a chain level correspondence with the hat flavor $\widehat{H M}(Y)$ of MrowkaKronheimer's monopole Floer homology. An analogous proof of invariance in the Legendrian setting is beyond the scope of this paper and is an interesting topic for future work.

Remark 68 (Component Grading). Let $\Lambda$ and $K$ be Legendrians in $\partial Y$ satisfying the hypotheses of our construction, and suppose that $\Lambda \subset K$. Then this induces an injective map

$$
\iota_{\Lambda}^{K}: E C H(Y, \Lambda, \lambda, J) \rightarrow E C H(Y, K, \lambda, J)
$$

If $\Lambda$ is the empty set, then $\operatorname{ECH}(Y, \emptyset, \lambda, J)$ is simply the sutured ECH $E C H(Y)$ of Colin-Ghiggini-Honda-Hutchings [28]. In particular, there is always an inclusion

$$
E C H(Y) \rightarrow E C H(Y, \Lambda, \lambda, J)
$$

Remark 69. Given a homology class $A \in H_{1}(Y, \Lambda ; \mathbb{Z})$, there is a sub-group

$$
E C H_{A}(Y, \Lambda, \lambda, J) \subset E C H(Y, \Lambda, \lambda, J)
$$

generated by orbit sets $\Theta$ in the homology class $A$. This induces a direct sum decomposition

$$
E C H(Y, \Lambda, \lambda, J)=\bigoplus_{A \in H_{1}(Y, \Lambda)} E C H_{A}(Y, \Lambda, \lambda, J)
$$

Remark 70 (Reeb Chord Filtration). Let $\mathfrak{z}$ be a Reeb chord connecting $\partial_{-} Y$ to $\partial_{+} Y$ that is disjoint from $\Lambda$. Then there is an associated holomorphic sub-manifold $Z \subset \mathbb{R} \times Y$ given by

$$
Z=\mathbb{R} \times \mathfrak{z}
$$

Given this choice, we can define an extended ECH complex

$$
\overline{E C C}(Y, \Lambda, \lambda, J)=\mathbb{F}_{2}[t] \otimes E C C(Y, \Lambda, \lambda, J)
$$

and a differential on $\overline{E C H}(Y, \Lambda, \lambda, J)$ (as a module over $\mathbb{F}_{2}[t]$ ) determined by $J$ and $\mathfrak{z}$.

$$
\partial_{J, \mathfrak{z}}(\Theta)=\sum_{\Xi}\left(\sum_{\mathcal{C} \in \mathcal{M}_{1}(\Theta, \Xi)} t^{\mathcal{C} \cdot Z} \cdot \Xi\right)
$$

Here $\mathcal{C} \cdot Z$ denotes the count (with multiplicity) of interior intersections between $\mathcal{C}$ and $Z$. By intersection positivity, this must always be positive. This defines an extended homology

$$
\overline{E C H}(Y, \Lambda, \lambda, J):=H(\overline{E C C}(Y, \Lambda, \lambda, J))
$$

We can extract further homology groups by studying the associated graded to the $t$-filtrartion. As we will discuss in $\S 4.1$, our construction is related to a construction of Heegaard-Floer homology.

Remark 71 (Previous Work). The Legendrian ECH index has appeared in the works of Colin-Ghiggini-Honda [23, 25] in a more limited context, in the process of establishing an isomorphism between ECH and Heegaard Floer homology.

This work is a natural elaboration on [23]. In particular, we fully develop a general theory of holomorphic currents with boundary, adjunction, writhe bound and the Legendrian ECH index that goes beyond the specialized context of [23, 25]. In our version of the index inequality, our holomorphic currents, for instance, have no restrictions on their asymptotic braids and are permitted to have boundary singularities. ${ }^{3}$

[^4]
## Legendrian ECH In The Closed Case

The Legendrian ECH setup is, at first glance, highly constrained. We now explain how to use our framework to associate ECH groups to any pair

$$
(Y, \Lambda)
$$

of a closed contact 3-manifold $Y$ with contact form $\alpha$ and a closed Legendrian $\Lambda \subset Y$. Crucially, these ECH groups essentially count Reeb chords of $\Lambda$ with respect to the initial contact form.

To start, choose an arbitrary metric $g$ on $\Lambda$ and choose a small $\epsilon>0$. Let $D^{*} \Lambda$ denote the unit codisk bundle. The Weinstein neighborhood theorem for Legendrians states that, for small $\epsilon$ and after scaling the metric to shrink the codisk bundle, there is an embedding

$$
\iota: B:=[-\epsilon, \epsilon]_{t} \times D^{*} \Lambda \rightarrow Y \quad \text { such that } \quad \iota(0, \Lambda)=\Lambda \text { and } \iota^{*} \alpha=d t+\lambda_{\text {std }}
$$

The complement $M=Y \backslash \iota(\operatorname{int}(B))$ of the interior of $B$ in $Y$ is a concave sutured contact manifold (see [28, Def. 4.2]) with a decomposition of the boundary into three pieces

$$
\partial_{-} M=\epsilon \times D^{*} \Lambda \quad \partial_{+} M=-\epsilon \times D^{*} \Lambda \quad \text { and } \quad \partial_{\circ} M=[-\epsilon, \epsilon] \times S^{*} \Lambda
$$

Note that the Reeb vector-field of $\alpha$ points out of $\partial_{+} M$ and into $\partial_{-} M$. This is the reason for the sign reversal in the notation. There are two copies of $\Lambda$ on $\partial M$ given by

$$
\Lambda_{-}:=\epsilon \times \Lambda \subset \partial_{-} M \quad \text { and } \quad \Lambda_{+}:=\epsilon \times \Lambda \subset \partial_{+} M
$$

The next step is to apply the concave-to-convex operation described in [28, §4.2]. In particular, there is a plug $U$ that one can attach to a neighborhood of $\partial_{\circ} M$ to acquire a convex sutured contact manifold

$$
\check{Y}:=M \cup U \quad \text { with contact form } \check{\alpha}
$$

The plug only modifies $\partial M$ near $\partial_{\circ} M$, so that $\Lambda_{+} \subset \partial_{+} \check{Y}$ and $\Lambda_{-} \subset \partial_{-} \check{Y}$. Moreover, every Reeb chord $c$ from $\Lambda_{-}$to $\Lambda_{+}$arises as a sub-chord of a self Reeb chords of $\Lambda$, and the lengths differ by an error of $2 \epsilon$. We can now make the following definition.

Definition 72. The Legendrian embedded contact homology ECH ${ }^{L}(Y, \Lambda)$ of $(Y, \alpha)$ and $\Lambda$ is the Legendrian ECH of $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ of action filtration below $L$.

$$
E C H^{L}(Y, \Lambda ; \delta):=E C H^{L}(\check{Y}, \check{\Lambda} ; \check{\alpha}, J)
$$

Here $\delta$ denotes the set of all choices made during the construction: the metric $g$, the parameter $\epsilon$, the embedding $\iota$ and the tailored almost complex structure $J$.

The groups in Definition 72 certainly depend on the specific choices, since Reeb chords can appear and disappear depending on the size of $\iota(B)$ in $Y$. Addressing this issue is far beyond the scope of this paper. However, let us state an optimistic conjecture in this direction.

Definition 73. A choice of data $\delta=(g, \epsilon, \iota, J)$ is L-admissible if

- Every Reeb orbit of length less than or equal to $L$ is in $\check{Y}$.
- The Reeb chords $\Lambda_{+} \rightarrow \Lambda_{-}$in $\check{Y}$ of length less than or equal to $L$ are in bijection with the Reeb chords of $\Lambda \rightarrow \Lambda$ in $Y$ of length less than or equal to $L$

These criteria are always achievable by shrinking $\iota$ and scaling the metric to reduce the size of $B$.

Conjecture 74. Let $(Y, \Lambda)$ be a pair of a closed contact 3-manifold $Y$, a closed Legendrian link $\Lambda \subset Y$ and a contact form $\alpha$ that is non-degenerate for $(Y, \Lambda)$. Then

- (Well-Defined) If $\delta$ and $\delta^{\prime}$ are two choices of $L$-admissible data, then there is a natural isomorphism

$$
E C H^{L}(Y, \Lambda ; \delta) \simeq E C H^{L}\left(Y, \Lambda ; \delta^{\prime}\right)
$$

The resulting group $E C H^{L}(Y, \Lambda)$ is the filtered ECH of $(Y, \Lambda)$.

- (Filtration Map) For any $K>L$, there is a map

$$
\iota_{L}^{K}: E C H(Y, \Lambda) \rightarrow E C H^{K}(Y, \Lambda) \quad \text { such that } \quad \iota_{L}^{M}=\iota_{K}^{M} \circ \iota_{L}^{K}
$$

- (Colimit) The colimit of the filtered ECH groups of $(Y, \Lambda)$ over $L$

$$
E C H(Y, \Lambda):=\underset{L}{\operatorname{colim}} E C H^{L}(Y, \Lambda)
$$

depends only on $(Y, \Lambda)$ up to contactomorphism of the pair.
Remark 75. It is likely that the maps in Conjecture 74 would be built from cobordism maps arising in a Seiberg-Witten based model for Legendrian ECH. In particular, Conjecture 74 is of a similar difficulty to the invariance proof discussed in Remark 67.

Remark 76. In [28] they use sutured ECH to define an (conjectured) invariant of Legendrian in closed 3-manifolds by considering sutured ECH of the same convex sutured manifold. However, we expect our invariant to be different from theirs because we allow Reeb chords in our chain complex.

## Motivation And Future Directions

This paper lays the groundwork for several future projects on the structure of ECH, each of which provides ample motivation for the development of our theory. We conclude this introduction by giving an overview of these motivating projects.

## Circle-Valued Gradient Flows

Let $M$ be a closed 3-manifold equipped with a circle-valued Morse function. That is, a smooth function

$$
f: M \rightarrow S^{1}
$$

with isolated critical points $p$ that each have non-degenerate Hessian. Assume also that $f$ has no index 0 or index 3 critical points. In [57, 56], Hutchings and Lee defined a 3 -manifold invariant

$$
I_{3}: \operatorname{Spin}^{c}(M) \rightarrow \mathbb{Z}
$$

from the set of spin-c structures on $M$ to the integers, via counts of configurations of closed orbits and flow lines of the gradient vector field of $f$. This invariant was motivated by and related to a number of other previously known invariants.

First, in a series of papers $[118,117]$, Turaev introduced a form of Reidemeister torsion, later dubbed Turaev torsion, which is also a map

$$
\tau_{M}: \operatorname{Spin}^{c}(M) \rightarrow \mathbb{Z}
$$

Hutchings-Lee proved in [56] that $\tau_{M}=I_{3}$. On the other hand, rapid developments in low-dimensional topology and gauge theory contemporaneous to [118, 117] lead to the introduction of the Seiberg-Witten invariant

$$
\operatorname{SW}_{M}: \operatorname{Spin}^{c}(M) \rightarrow \mathbb{Z}
$$

This invariant is defined using a signed and weighted count of solutions to the 3dimensional Seiberg-Witten equation (cf. [83]). Turaev established in [116] that $\mathrm{SW}_{M}=\tau_{M}$ (up to sign), proving through indirect means that

$$
\begin{equation*}
I_{3}=\mathrm{SW}_{M} \tag{4.1.4}
\end{equation*}
$$

Through the equality (4.1.4), one is lead to the following question.
Question 77. Is there a direct proof of the equality $I_{3}=\mathrm{SW}_{M}$ that does not use Turaev torsion?

Moreover, the Seiberg-Witten invariant was categorified by Kronheimer-Mrowka's monopole Floer homology $H M_{\bullet}[74]$ and other variants of Seiberg-Witten-Floer theory. This suggests the following (roughly formulated) question, which is related to Question 77.

Question 78. Is there a Floer homology theory $F H_{\bullet}(Y, f)$ of 3-manifolds $Y$ with a circle valued Morse function $f: Y \rightarrow S^{1}$ as above such that
(a) $\mathrm{FH} \bullet$ categorifes $I_{3}$, i.e. it is computed as the homology of a complex generated by counts of configurations of gradient flow lines and closed orbits as in $I_{3}$.
(b) $\mathrm{FH}_{\bullet}$ is isomorphic to (the appropriate flavor of) monopole Floer homology.

Embedded contact homology, and its sister theory periodic Floer homology (PFH), answer both of these questions when the Morse function $f$ has no critical points. In this case, $M$ may be viewed as a mapping torus of a map

$$
\phi: \Sigma \rightarrow \Sigma \quad \text { where } \quad \Sigma=f^{-1}(0)
$$

and the generators of the hypothetical Floer homology groups $F H_{\bullet}$ must be configurations of periodic points of $\phi$. The PFH groups PFH• provides just such a theory and a result of Lee-Taubes [82] states that PFH• and $H M_{\bullet}$ are isomorphic ${ }^{4}$.

In the general case where $f$ can have critical points, Questions 77 and 78 are still open. However, there is a potential approach based on a slight generalization of the constructions in this paper. Choose a metric $g$ on $M$ such that $f$ is harmonic (this is possible if $f$ has no index 0 or 3 critical points). Assume that each index 2 critical point $p$ has a Morse chart where

$$
f(x, y, z)=2 x^{2}-y^{2}-z^{2}
$$

In this chart, we can remove a standard neighborhood $U$ and introduce a boundary component to $M$. The new boundary has corners, and on the smooth components the gradient vector-field is either orthogonal to or tangent to the boundary. This neighborhood is depicted in Figure 4.4.

We can perform this neighborhood removal around each critical point of $f$ (using an analogous local model near the index 1 points) to acquire a new space $Y \subset M$. This space is equipped with a stable Hamiltonian structure with 2-form $\omega_{f}$ given by the Hodge dual of $d f$ and stabilizing 1-form $\theta_{f}=d f$. The boundary of $Y$ contains a

[^5]

Figure 4.4: A standard neighborhood $U$ of an index 2 critical point. The critical point is in red and the descending sphere is in blue. Note the regions where the gradient points in and out of the boundary are level surfaces of $f$.


Figure 4.5: The space acquired by removing standard neighborhoods around each critical point from $Y$.
natural 1-dimensional sub-manifold $\Lambda$ given as the union of the ascending sphere of the index 1 critical points and the descending spheres of the index 2 critical points. This sub-manifold satisfies

$$
T \Lambda \subset \operatorname{ker}\left(\theta_{f}\right) \quad \text { and }\left.\quad \omega_{f}\right|_{T \Lambda}=0
$$

In other words, $\Lambda$ is the analogue of a Legendrian in the stable Hamiltonian manifold $\left(Y, \omega_{f}, \theta_{f}\right)$. Moreover, any gradient flow line $\eta$ from an index 1 critical point $p$ to an index 2 critical point $q$ becomes a chord $c$ connecting the corresponding components of $\Lambda$.

The pair $(Y, \Lambda)$ very much resembles a stable Hamiltonian analogue of the setup used in this paper to define Legendrian ECH. Thus, one approach to addressing Question 78 is to use the methods of this paper to develop a PFH version of our Legendrian ECH theory for this setting. There are some significant technical challenges to carrying out this program. For example, the boundary of $Y$ is naturally concave sutured, rather than convex, and thus we do not have immediate access to an appropriate maximum principle. This could be addressed by adapting the
concave-to-convex operation in [28] to our stable Hamiltonian setting and establish a maximum principle for $J$-holomorphic curves near the boundary. We hope to pursue this in future work.

## Heegaard-Floer Homology And Embedded Contact Homology

It has been shown that embedded contact homology is isomorphic to (the appropriate flavors of) several other Floer homologies, notably Heegaard-Floer homology and monopole Floer homology. The isomorphisms are, however, highly non-trivial to construct. For instance, the construction of the isomorphism relating ECH to Heegaard-Floer theory (through open book decompositions) occupies four long papers due to Colin-Ghiggini-Honda [20, 23, 25, 21].

The work of Colin-Ghiggini-Honda utilizes a cylindrical reformulation of HeegaardFloer homology due to Lipshitz [85] that we now describe in broad terms. The construction begins with a pointed Heegaard diagram that we write as follows.

$$
(\Sigma, \alpha, \beta, \mathfrak{z})
$$

This consists of a closed orientable surface $\Sigma$ of genus $g$, a distinguished point $\mathfrak{z} \in \Sigma$ and two collections of $g$ simple, non-separating, closed curves

$$
\alpha=\alpha_{1} \cup \cdots \cup \alpha_{g} \quad \text { and } \quad \beta=\beta_{1} \cup \cdots \cup \beta_{g}
$$

Recall that ( $\Sigma, \alpha, \beta, \mathfrak{z}$ ) determines a 3 -manifold $M$ (uniquely, up to diffeomorphism). More precisely, we take the 3-manifold

$$
Y=[0,1] \times \Sigma \quad \text { with curves } \quad \Lambda_{\alpha}=(0 \times \alpha) \quad \text { and } \quad \Lambda_{\beta}=(1 \times \beta)
$$

and attach 2-handles to each curve in $\Lambda_{\alpha}$ and $\Lambda_{\beta}$. The resulting manifold has a boundary consisting of two 2 -spheres, and after attaching 3 -handles to these areas, we acquire $M$.

The space $Y=[0,1]_{t} \times \Sigma$ may be viewed as a stable Hamiltonian manifold with the two-form $\omega$ equal to an area form on $\Sigma$ and stabilizing 1-form $d t$. The analogue of the Reeb vector-field is $R=\partial_{t}$. Moreover, $\Lambda_{\alpha}$ and $\Lambda_{\beta}$ are Legendrians in the sense that

$$
\left.\omega\right|_{\Lambda_{ \pm}}=0 \quad \text { and }\left.\quad d t\right|_{\Lambda_{ \pm}}=0
$$

In this picture, Reeb chords between $\Lambda_{\alpha}$ and $\Lambda_{\beta}$ are equivalent to intersection points in $\alpha \cap \beta$. The symplectization of $Y$ is given the symplectic manifold

$$
W=\mathbb{R}_{s} \times Y \quad \text { with symplectic form } \quad \Omega=d s \wedge d t+\omega_{\Sigma}
$$



Figure 4.6: A picture of the setup of Lipshitz' cylindrical formulation of HeegaardFloer homology.

The cylindrical sub-manifolds $L_{\alpha}=\mathbb{R} \times \Lambda_{\alpha}$ and $L_{\beta}=\mathbb{R} \times \Lambda_{\beta}$ are both Lagrangians. There is a natural class of compatible almost complex structures $J$ on $W$ : those that are translation invariant and that satisfy $J\left(\partial_{s}\right)=\partial_{t}$ and $J(T \Sigma)=T \Sigma$. Finally, note that the marked point $\mathfrak{z}$ determines a $J$-holomorphic strip

$$
Z=\mathbb{R} \times[0,1] \times \mathfrak{z}
$$

The formulation of Heegaard-Floer homology of Lipshitz [85] can now be described as follows. The chain complex $C F_{\text {Lip }}(\Sigma, \alpha, \beta, \mathfrak{z})$ is generated by sets of the form

$$
\Xi=\left\{c_{1}, \ldots, c_{g}\right\}
$$

where $\Xi$ consists of chords $c_{i}$ from $\alpha_{i}$ to $\beta_{\sigma(j)}$ for some permutation $\sigma$, or equivalently a set of intersection points $p_{i} \in \alpha_{i} \cap \beta_{\sigma(i)}$. The differential $\partial$ counts pseudo-holomorphic curves of the following form. Let $\Gamma$ and $\Xi$ be two generators. Let $S$ with a Riemann surface with $2 g$ boundary punctures, $g$ of which we label positive and $g$ of which we label negative. We consider $J$-holomorphic maps

$$
u:(S, \partial S) \rightarrow\left(W, L_{\alpha} \cup L_{\beta}\right)
$$

of Fredholm index 1 (modulo translation), where the positive punctures are asymptotic to the chords corresponding to $\Gamma$ at $+\infty$ of the $\mathbb{R}$ direction in $W$; the negative punctures are asymptotic to the chords corresponding to $\Xi$ at $-\infty$ of the $\mathbb{R}$ direction in $W$; the image of $S$ under $u$ is embedded; and $u(S)$ is disjoint from $Z$. For more technical formulations of these conditions see section 1 in [85]. For a depiction of this setup, see Figure 4.6.

Theorem 79. [85] The homology groups of $C F_{\text {Lip }}(\Sigma, \alpha, \beta, \mathfrak{z})$ are isomorphic to the hat version of the Heegaard-Floer homology groups.

$$
H\left(C F_{L i p}(\Sigma, \alpha, \beta, \mathfrak{z})\right) \simeq \widehat{H F}(Y)
$$

The setup for Lipshitz' construction is very similar to the one for Legendrian ECH. In fact, it may be viewed as a special case of our construction adapted to the stable Hamiltonian manifold

$$
(Y, \omega, \theta)=\left([0,1] \times \Sigma, \omega_{\Sigma}, d t\right) \quad \text { with "Legendrians" } \quad \Lambda_{\alpha} \cup \Lambda_{\beta} \subset \partial Y
$$

This is the perspective adopted by Colin-Ghiggini-Honda, and in [23, §4] they prove the following claim:

Claim 80. [23, §4] The differential $C F_{\text {Lip }}(\Sigma, \alpha, \beta, \mathfrak{z})$ precisely counts ECH index 1 holomorphic curves $C$ in $\mathbb{R} \times Y$ with boundary $\mathbb{R} \times\left(\Lambda_{\alpha} \cup \Lambda_{\beta}\right)$.

The theory articulated in this paper can therefore be viewed as providing a common framework for describing both Heegaard-Floer theory and embedded contact homology as special cases of a single construction, when one includes only the Reeb chords or Reeb orbits as generators.

## Bordered Embedded Contact Homology And Gluing Formulas

In the paper [86], Lipshitz-Oszvath-Thurston formulate a theory of Heegaard-Floer homology for 3-manifolds with boundary, called bordered Heegaard-Floer homology.

Roughly speaking, the bordered Heegaard-Floer homology groups depend on a 3manifold $Y$ and on a diffeomorphism $\phi: F \simeq \partial Y$ of $\partial Y$ with a surface $F$ determined by a datum called a matched circle $\mathcal{Z}$. The hat version of this theory comes in two flavors.

$$
\widehat{C F A}(Y, \phi) \quad \text { and } \quad \widehat{C F D}(Y, \phi)
$$

which are, respectively, an $A_{\infty}$-module and a dg-module over a certain free dg-algebra $\mathcal{A}$ associated to $\mathcal{Z}$. A fundamental application of this theory is the following gluing result.

Theorem 81. [86] If $Y=M \cup_{F} N$ is a closed union of two 3-manifolds with boundary via the identifications of $\phi: F \simeq \partial M$ and $\psi: F \simeq \partial N$ with $F$, then

$$
\widehat{H F}(Y) \simeq H\left(\widehat{C F A}(M, \phi) \otimes_{\mathcal{A}} \widehat{C F D}(N, \psi)\right)
$$

Here $\otimes_{\mathcal{A}}$ denotes the tensor product as $A_{\infty}$-modules over $\mathcal{A}$.

This gluing formula has since become a fundamental tool in computing HeegaardFloer invariants and their knot counterparts. In particular, they are useful in studying submanifolds e.g. knots and surfaces [48, 42].

Although many variants of symplectic field theory can be formulated for contact manifolds with boundary (e.g. [28]), a general gluing formula in the spirit of Theorem 81 has yet to appear. As a first step, one can try to prove Theorem 81 where the sutured Heegaard-Floer groups are replaced with variants of ECH for manifolds with boundary.

To further understand what bordered ECH might look like, we observe that the two types of hat bordered Heegaard-Floer complexes are constructed by extending Lipshitz' cylindrical formulation [85] in two ways.

First, the complexes are computed using a bordered Heegaard diagram $(\Sigma, \alpha, \beta)$ for $Y$, where $\Sigma$ has boundary and some $\alpha$ curves are permitted to have boundary in $\partial \Sigma$. Second, the bordered complexes incorporate Reeb chords between the Legendrians (i.e. points) $\alpha \cap \partial \Sigma$ in the matched circle $\mathcal{Z}$, which is identified with $\partial \Sigma$. The dgalgebra $\mathcal{A}$ associated to $\mathcal{Z}$ is, in fact, a variant of the Chekanov-Eliashberg dg-algebra of the Legendrians in $\mathcal{Z}$. The holomorphic curves counted in bordered HeegaardFloer theory are allowed to limit to chords of $\alpha \cap \partial \Sigma$ at boundary punctures, and the chain complexes themselves are (roughly speaking) augmented by changing the ordinary Heegaard-Floer complexes to have coefficients in $\mathcal{A}$.

In light of these observations, we expect that a bordered theory of ECH will require an analogous extension of Legendrian ECH that incorporates Legendrians $\Lambda \subset \partial_{ \pm} Y$ with boundary on $\partial \Lambda$ contained in $\partial\left(\partial_{ \pm} Y\right) \simeq \sigma$, the suture, and Reeb chords between the points $\partial \Lambda \subset \sigma$. The precise formulation of this theory will be the subject of future work based on this paper.

### 4.2 Intersection Theory

In this section, we describe an intersection theory for surfaces in symplectic cobordisms with boundary on a Lagrangian cobordism in dimension 4. This mirrors the theory for surfaces without Lagrangian boundary described by Hutchings in [55].

## Bundle Pairs

We begin by introducing the notion of a bundle pair with punctures over a punctured Riemann surface with boundary.

Definition 82. A symplectic bundle pair with punctures $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ consists of
(a) A compact, oriented surface with boundary and corners

$$
\Sigma \quad \text { with } \quad \partial \Sigma=\partial_{\star} \Sigma \cup \partial_{0} \Sigma
$$

where $\partial_{\star} \Sigma$ and $\partial_{\circ} \Sigma$ are each unions of smooth strata of the boundary.
(b) A bundle pair of a symplectic vector-bundle and a Lagrangian sub-bundle

$$
(E, \omega) \rightarrow \Sigma \quad \text { and } \quad F \rightarrow \partial_{\circ} \Sigma \quad \text { with }\left.\quad F \subset E\right|_{\partial_{0} \Sigma}
$$

We refer to $\partial_{\circ} \Sigma$ as the Lagrangian boundary and to $\partial_{\star} \Sigma$ as the puncture boundary. We note by our previous notation $\partial_{\star} \Sigma$ is the same as $\partial_{ \pm} \Sigma$. We equip both $\partial_{\circ} \Sigma$ and $\partial_{\star} \Sigma$ with the induced boundary orientation.

Definition 83. A puncture trivialization $\tau$ of $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ consists of the following. First a symplectic trivialization of $E$ over $\partial_{\star} \Sigma$, which we write as

$$
\left.E\right|_{\partial_{\star} \Sigma} \simeq \mathbb{C}^{n}
$$

so that

$$
\left.F\right|_{\partial_{\star} \Sigma \cap \partial_{0} \Sigma} \simeq \mathbb{R}^{n} \subset \mathbb{C}^{n}
$$

We fix the convention of thinking of $\tau$ as a map from $E \rightarrow \partial_{\star} \Sigma$ to $\mathbb{C}^{n} \rightarrow \partial_{\star} \Sigma$.
Bundle pairs with punctures admit an integer invariant analogous to the Chern class, generalizing the Maslov number of an ordinary bundle pair.

Proposition 84 (Maslov Number). For any symplectic bundle pair with punctures $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ and puncture trivialization $\tau$, there is a well-defined Maslov number

$$
\mu(E, F ; \tau) \in \mathbb{Z}
$$

Furthermore, the Maslov number is uniquely determined by the following axioms.
(a) (Isomorphism) The Maslov number is invariant under isomorphism of the pair and trivialization.

$$
\mu(E, F ; \tau)=\mu\left(E^{\prime}, F^{\prime} ; \tau^{\prime}\right) \quad \text { if } \quad(E, F ; \tau) \simeq\left(E^{\prime}, F^{\prime} ; \tau^{\prime}\right)
$$

(b) (Direct Sum) The Maslov number is additive with respect to direct sum.

$$
\mu\left(E \oplus E^{\prime}, F \oplus F^{\prime} ; \tau \oplus \tau^{\prime}\right)=\mu(E, F ; \tau)+\mu\left(E^{\prime}, F^{\prime} ; \tau^{\prime}\right)
$$

(c) (Disjoint Union) The Maslov number is additive under disjoint union.

$$
\mu\left(E \cup E^{\prime}, F \cup F^{\prime} ; \tau \cup \tau^{\prime}\right)=\mu(E, F ; \tau)+\mu\left(E^{\prime}, F^{\prime} ; \tau^{\prime}\right)
$$

(d) (Puncture Gluing) Let $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ and $\left(E^{\prime}, F^{\prime}\right) \rightarrow\left(\Sigma^{\prime}, \partial \Sigma^{\prime}\right)$ be two bundle pairs with punctures. Let $R \subset \partial_{\star} \Sigma$ and $R^{\prime} \subset \partial_{\star} \Sigma^{\prime}$ be components of the puncture boundary. Let $-R^{\prime}$ denote $R^{\prime}$ with reverse orientation. Assume we have a bundle pair isomorphism

$$
\psi:\left(\left.E\right|_{R}\right) \simeq\left(\left.E^{\prime}\right|_{-R^{\prime}}\right) \quad \text { covering a diffeomorphism } \quad R \simeq-R^{\prime}
$$

Further, composed with trivializations $\tau$ and $\tau^{\prime}$, the map $\psi$ takes the form

$$
\tau^{\prime-1} \circ \psi \circ \tau:\left(\mathbb{C}^{n}, \mathbb{R}^{n}\right) \rightarrow\left(\mathbb{C}^{n}, \mathbb{R}^{n}\right)
$$

over $R$ and $-R^{\prime}$. Then

$$
\mu\left(E \cup_{\psi} E^{\prime}, F \cup_{\psi} F^{\prime}, \tau \cup \tau^{\prime}\right)=\mu(E, F, \tau)+\mu\left(E^{\prime}, F^{\prime}, \tau^{\prime}\right)
$$

Here $\left(E \cup_{\psi} E^{\prime}, F \cup_{\psi} F^{\prime}\right)$ is the bundle pair over ( $\Sigma, \partial \Sigma$ ) acquired by gluing of $(E, F)$ and $\left(E^{\prime}, F^{\prime}\right)$, and $\tau \cup_{\psi} \tau^{\prime}$ is the glued trivialization.
(e) (Normalization) If $\Sigma=D$ is the 2-disk with $\partial \Sigma=\partial_{0} \Sigma=S^{1}$, and $(E, F)$ is the bundle pair

$$
E=D \times \mathbb{C} \quad \text { and } \quad F=\left\{e^{i \theta} \mathbb{R}: \theta \in \partial D \simeq \mathbb{R} / \pi \mathbb{Z}\right\}
$$

then the Maslov number is given by $\mu(E, F ; \tau)=1$.
Proof. To prove existence, let $\bar{F} \rightarrow \partial \Sigma$ denote the Lagrangian sub-bundle of $\left.E\right|_{\partial \Sigma}$ given by

$$
\left.\bar{F}\right|_{\partial_{0} \Sigma}=F \quad \text { and }\left.\quad \bar{F}\right|_{\partial_{\star} \Sigma}=\tau^{-1}\left(\mathbb{R}^{n}\right)
$$

Then $\mu(E, F ; \tau)$ is simply the standard Maslov number $\mu(E, \bar{F})$ of the pair $(E, \bar{F})$ (see [91, §C.3]).

$$
\mu(E, F ; \tau)=\mu(E, \bar{F})
$$

The properties (a)-(e) except (d) follow immediately from the analogous properties of the ordinary Maslov number (see [91, Theorem C.3.5(a)-(d)]).

To see (d), consider for $(E, \bar{F}) \rightarrow(\Sigma, \partial \Sigma)$, there is already a trivialization of $E$ over the boundary punctures. Extend this to a trivialization of $E$ over the entire boundary. This is possible since all complex bundles over $S^{1}$ are trivial. In the interior of $\Sigma$, we cut out a disk with boundary $(D, \partial D)$, then the trivialization of $E$ can be extended to $E \backslash D$. We further choose a family of Lagrangians over $\left.E\right|_{\partial D}$, which we denote by $\widetilde{F}$. Then by the additivity property of the Maslov index we have

$$
\mu(E, F ; \tau)=\mu\left(\left.E\right|_{\Sigma \backslash D},\left.\bar{F} \cup \widetilde{F}\right|_{\partial \Sigma \cup-\partial D}\right)+\mu\left(\left.E\right|_{D}, \widetilde{F}\right)
$$

Further more, we observe $\left.E\right|_{\Sigma \backslash D} \simeq \Sigma \backslash D \times \mathbb{C}^{n}$, hence $\mu\left(\left.E\right|_{\Sigma \backslash D},\left.\bar{F} \cup \widetilde{F}\right|_{\partial \Sigma \cup-\partial D}\right)$ is the sum of the Maslov indices of the Lagrangians $\bar{F} \cup \widetilde{F} \subset \mathbb{C}^{n}$ along $\partial \Sigma \cup-\partial D$. An entirely analogous construction holds for $\left(E^{\prime} \bar{F}^{\prime}\right) \rightarrow\left(\Sigma^{\prime}, \partial \Sigma^{\prime}\right)$ - we just add prime to all of our previous symbols. Then the act of gluing induces a gluing between $(\partial \Sigma \backslash D, \partial \Sigma \cup-\partial D)$ and $\left(\partial \Sigma^{\prime} \backslash D^{\prime}, \partial \Sigma^{\prime} \cup-\partial D^{\prime}\right)$, then this implies that $\mu\left(\left.E\right|_{\Sigma \backslash D},\left.\bar{F} \cup \widetilde{F}\right|_{\partial \Sigma \cup-\partial D}\right)+\mu\left(\left.E^{\prime}\right|_{\Sigma^{\prime} \backslash D^{\prime}},\left.\bar{F}^{\prime} \cup \widetilde{F}^{\prime}\right|_{\partial \Sigma \cup-\partial D}\right)=\mu\left(\left.E \cup E^{\prime}\right|_{(\Sigma \backslash D) \cup_{\psi}\left(\Sigma^{\prime} \backslash D^{\prime}\right)}, F \cup_{\psi} F^{\prime} \cup \widetilde{F} \cup \widetilde{F}^{\prime}\right)$
because the contributions for the Lagrangians $F$ and $F^{\prime}$ cancel due to our assumptions. The additivity formula then follows from the composition formula in the regular Maslov index case.

To prove uniqueness, we argue as follows. First, consider the case where $\Sigma$ is the half-disk

$$
\begin{align*}
& \Sigma=\mathbb{D} \cap \mathbb{H} \quad \text { with } \quad \partial_{\star} \Sigma=\partial \mathbb{D} \cap \mathbb{H} \text { and } \partial_{0} \Sigma=\mathbb{D} \cap \mathbb{R}  \tag{4.2.1}\\
& E=\mathbb{C}^{n} \quad F=\mathbb{R}^{n} \quad \text { and } \quad \tau \text { is the tautological trivialization }
\end{align*}
$$

Then by gluing $(E, F)$ to itself along $\partial_{\star} \Sigma$ and applying axiom (d), we find that

$$
\mu(E, F ; \tau)=\frac{1}{2} \cdot \mu\left(\mathbb{C}^{n}, \mathbb{R}^{n}\right)=0
$$

Second, the uniqueness of the ordinary Maslov index [91, Theorem C.3.5] implies that for any bundle pair $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ with $\partial_{\star} \Sigma=\emptyset$, we have

$$
\mu(E, F)=\mu(E, F ; \tau)
$$

Finally, any bundle pair with punctures $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ transforms into a bundle pair with no puncture boundary by gluing on trivial bundles over the half-disk along $\partial_{\star} \Sigma$. Thus the gluing axiom (d) determines $\mu(E, F ; \tau)$. Uniqueness then follows from the uniqueness of the Maslov index for surfaces without boundary punctures.

Remark 85. In the case where $\partial_{\star} \Sigma=\emptyset$, we omit the (empty) trivialization from the notation and denote the Maslov number by $\mu(E, F)$.

Example 86 (Special Cases). The following integer invariants can be viewed as special cases of the Maslov number.
(a) The 1st Chern number of a bundle $E \rightarrow \Sigma$ of a surface with boundary with respect to a trivialization $\tau$ along $\partial \Sigma$ is

$$
c_{1}(E, \tau)=\frac{1}{2} \cdot \mu(E, \emptyset ; \tau)
$$

In keeping with our previous conventions, we take $\partial_{*} \Sigma$ to be the entire boundary of $\Sigma$, and $\partial_{o} \Sigma=\emptyset$. In particular, if $\Sigma$ is closed then $c_{1}(E)=\frac{1}{2} \cdot \mu(E, \emptyset)$.
(b) [91, Thm C.3.6] Let $\Sigma$ be a compact surface with boundary and let $F: \partial \Sigma \rightarrow$ $\mathrm{LGr}(2 n)$ be a set of loops in the Lagrangian Grassmannian. Then

$$
\mu(F):=\mu\left(\Sigma \times \mathbb{C}^{n}, F\right)
$$

(c) The Maslov index of a loop $\Phi: S^{1} \rightarrow \operatorname{Sp}(2 n)$ is

$$
\mu(\Phi):=\frac{1}{2} \cdot \mu\left(D \times \mathbb{C}^{n}, F\right) \quad \text { where } \quad F_{\theta}:=\Phi_{\theta}\left(\mathbb{R}^{n}\right) \subset \mathbb{C}^{n}
$$

Example 87 (Maslov Class). Let $(X, L)$ be a pair of a symplectic manifold and a Legrangain sub-manifold $L \subset X$. The Maslov class

$$
\mu(X, L): H_{2}(X, L) \rightarrow \mathbb{Z}
$$

is the map defined so that the value on the class $A$ of a surface $\Sigma \subset X$ with boundary $\partial \Sigma \subset L$ is the Maslov number of $\left.(T X, T L)\right|_{\Sigma}$.

The Maslov class $\mu(\xi, \Lambda)$ of a contact manifold $(Y, \xi)$ with Legendrian $\Lambda \subset Y$ is defined similarly, and we have

$$
\mu(\xi, \Lambda)=\mu(T X, T L) \quad \text { if } \quad(X, L)=([0,1] \times Y,[0,1] \times \Lambda)
$$

We will require a few properties of the Maslov number. First, we must record how the Maslov number changes under change of end trivializations.

Lemma 88 (Trivial Bundle). Consider the a bundle pair with punctures

$$
\left(\mathbb{C}^{n}, F\right) \rightarrow\left(\partial \Sigma, \partial_{\circ} \Sigma\right)
$$

where $F: \partial_{0} \Sigma \rightarrow \operatorname{LGr}(2 n)$ is a collection of arcs in the Lagrangian Grassmanian with ends on $\mathbb{R}^{n}$. Then

$$
\mu(E, F ; \tau)=\mu(F)+\mu\left(\tau^{-1}\left(\mathbb{R}^{n}\right)\right)
$$

Here $\mu(F)$ and $\mu\left(\tau^{-1}\left(\mathbb{R}^{n}\right)\right)$ are the Maslov indices of the (collections of) loops of Lagrangians $F$ and $\tau^{-1}\left(\mathbb{R}^{n}\right)$.

Proof. As in Proposition 84, the Maslov index $\mu(E, F ; \tau)$ is given by

$$
\mu(E, F ; \tau)=\mu(E, \bar{F})
$$

where $\bar{F}$ is the Lagrangian sub-bundle acquired by taking the union of $F$ over $\partial_{\circ} \Sigma$ and $\tau^{-1}\left(\mathbb{R}^{n}\right)$ over $\partial_{\star} \Sigma$. The result thus follows immediately from Example 86(b).

As with the Chern number, the Maslov number can be interpreted as a signed count of zeros.

Lemma 89 (Zero Count). Let $(E, F) \rightarrow(\Sigma, \partial \Sigma)$ be a bundle pair with punctures equipped with boundary trivialization $\tau$. We further assume $E$ is two dimensional. Let $\psi:\left(\Sigma, \partial_{\circ} \Sigma\right) \rightarrow(E, F)$ be a section with

$$
\psi \text { is transverse to } \Sigma \quad \text { and } \quad \psi\left(\partial_{\star} \Sigma\right) \subset \tau(\mathbb{R} \backslash 0)
$$

Then the Maslov number $\mu(E, F ; \tau)$ is given by

$$
\begin{equation*}
\mu(E, F ; \tau):=2 \cdot \#(\psi \cap \operatorname{int}(\Sigma))+\#\left(\psi \cap \partial_{0} \Sigma\right) \tag{4.2.2}
\end{equation*}
$$

Proof. Let $-\Sigma$ and $-E$ denote the surface and bundle with reversed orientations, respectively. Recall a trivialization $\tau$ is achoice of a map

$$
\tau:\left.E\right|_{\partial_{*} \Sigma} \rightarrow \mathbb{C}
$$

With this trivialization, as with the case how we defined Maslov indices, we extend $F$ over $\partial_{*} \Sigma$ as

$$
\tau:\left.(E, F)\right|_{\partial_{*} \Sigma} \rightarrow(\mathbb{C}, \mathbb{R})
$$

Furthermore, let $\bar{\tau}$ denote the composition of $\tau$ with complex conjugation (which is anti-symplectic).

$$
\left.(E, F)\right|_{\partial_{\star} \Sigma} \simeq(\mathbb{C}, \mathbb{R}) \xrightarrow{z \rightarrow \bar{z}}(\mathbb{C}, \mathbb{R})
$$

This composition is a trivialization of $-E$ over $\partial_{\star} \Sigma$ in $-\Sigma$. Viewing $-E$ as a bundle over $-\Sigma$, we have

$$
\mu(-E, F ; \bar{\tau})=\mu(E, F ; \tau)
$$

Moreover, there is a natural (isotopy class of) symplectic bundle map

$$
\left.\left.(-E, F)\right|_{\partial_{\circ} \Sigma} \rightarrow(E, F)\right|_{\partial_{\circ} \Sigma} \quad \text { covering } \quad \text { Id }: \partial_{\circ} \Sigma \rightarrow \partial_{\circ} \Sigma
$$

given by complex conjugation with respect to the totally real subspace $F$. Thus, we can glue $E$ and $-E$ to a bundle $D E$ over the double $D \Sigma=\Sigma \cup_{\partial_{0} \Sigma}-\Sigma$ of $\Sigma$ along the boundary region $\partial_{0} \Sigma$. Using the composition property of Maslov class (Theorem C.3,5 in [91] and the fact we have fixed trivializations around punctures so that the totally bundle is trivial) gives us that

$$
\mu(E, F ; \tau)=\frac{1}{2} \mu(D E, \emptyset ; \tau)=c_{1}(D E, \emptyset ; \tau)
$$

Now take a section $\psi$ as in the lemma statement. We may assume (after a small isotopy leaving the zeros unchanged) that this section doubles to a section $\phi$ of $D E$ that is transverse to the zero section. The relative Chern class of $D E$ with respect to $\tau$ is computed by the zeros of $\phi$, and thus

$$
\mu(E, F ; \tau)=c_{1}(D E, \emptyset ; \tau)=\#(\phi \cap D \Sigma)=2 \cdot \#(\psi \cap \operatorname{int}(\Sigma))+\#\left(\psi \cap \partial_{\circ} \Sigma\right)
$$

## Euler Characteristic

Consider a compact surface $\Sigma$ with boundary and corners $\partial \Sigma=\partial_{0} \Sigma \cup \partial_{\star} \Sigma$. In this paper, we will use the following version of the Euler characteristic.

Definition 90. The orbifold Euler characteristic $\bar{\chi}(\Sigma)$ is the quantity

$$
\chi(\Sigma)-\frac{1}{2} \chi\left(\partial_{\circ} \Sigma\right) \quad \text { or equivalently } \quad \frac{1}{2} \cdot \chi\left(\Sigma \cup_{\partial_{0} \Sigma}-\Sigma\right)
$$

Like the ordinary Euler characteristic, this quantity is a special case of the Maslov number. Specifically, consider the tangent bundle of $\Sigma$. This naturally has the structure of a bundle pair with punctures.

$$
\left(T \Sigma, T\left(\partial_{0} \Sigma\right)\right) \rightarrow(\Sigma, \partial \Sigma)
$$

Moreover, this bundle pair comes with a canonical isotopy class of trivialization

$$
\tau_{\partial \Sigma}:\left.\left(T \Sigma, T\left(\partial_{0} \Sigma\right)\right)\right|_{\partial_{\star} \Sigma} \simeq(\mathbb{C}, \mathbb{R}) \quad \text { with } \quad \tau_{\partial \Sigma}^{-1}(\mathbb{R}) \pitchfork T\left(\partial_{\star} \Sigma\right)
$$

The Maslov number of the tangent pair in the canonical trivialization is precisely twice the orbifold Euler characteristic.

Lemma 91 (Euler Characteristic). Let $\Sigma$ be a surface with boundary $\partial \Sigma$. Then

$$
\frac{1}{2} \cdot \mu\left(T \Sigma, T\left(\partial_{o} \Sigma\right) ; \tau_{\partial \Sigma}\right)=\bar{\chi}(\Sigma)
$$

Proof. First, consider two special cases. If $\Sigma$ is closed, this is equivalent to the fact that $c_{1}(T \Sigma)=\chi(\Sigma)$. Likewise, if $\Sigma=\mathbb{D}$ with $\partial_{\star} \Sigma=\partial \mathbb{D}$, then $T(\partial \mathbb{D}) \subset T \mathbb{D} \simeq \mathbb{D} \times \mathbb{R}^{2}$ is a loop of Lagrangians with Maslov index 2. Thus by Example 86(c), we have

$$
\frac{1}{2} \cdot \mu(T \mathbb{D}, T(\partial \mathbb{D}))=1=\chi(\mathbb{D})
$$

This verifies these special cases. In the general case we can consider the double of $\Sigma$ along $\partial_{\circ} \Sigma$.

$$
S=\Sigma \cup_{\partial_{0} \Sigma} \bar{\Sigma}
$$

Applying the proof of Lemma 89 and the property of the Euler characteristic under gluing, we deduce

$$
\mu\left(T S, \emptyset ; \tau_{\partial S}\right)=2 \cdot \mu\left(T \Sigma, T\left(\partial_{0} \Sigma\right) ; \tau_{\partial \Sigma}\right) \quad \text { and } \quad \chi(S)=2 \cdot \bar{\chi}(\Sigma)
$$

On the other hand, $S$ can be acquired by taking a closed surface and removing a collection of disks to produce puncture boundary. Thus, by Proposition 84(d) and the special cases, we have

$$
\mu\left(T S, \emptyset ; \tau_{\partial S}\right)=\chi(S)
$$

## Setup And Trivializations

Before proceeding, it will be useful for us to fix a common setup for the remainder of this section. For each $\bullet \in\{+,-\}$, we fix
(a) a compact contact 3 -manifold $\left(Y_{\bullet}, \xi_{\bullet}\right)$ with boundary $\partial Y_{\bullet}$.
(b) a closed Legendrian sub-manifold $\Lambda_{\bullet} \subset \partial Y_{\bullet}$ in the boundary of $Y_{\bullet}$.
(c) a contact form $\alpha_{\bullet}$ with non-degenerate Reeb orbits and chords.
(d) an orbit-chord set $\Xi_{\bullet}=\Gamma_{\bullet} \cup C_{\bullet}$ consisting of a Reeb orbit set $\Gamma_{\bullet}=\left\{\left(\gamma_{i}^{\bullet}, m_{i}^{\bullet}\right)\right\}$ and Reeb chord set $C_{\bullet}=\left\{\left(c_{i}^{\bullet}, n_{i}^{\bullet}\right)\right\}$.
(e) a symplectic 4-dimensional cobordism $X: Y_{+} \rightarrow Y_{-}$with symplectic form $\omega$. This is a manifold with boundary and corners, where $\partial X=\partial_{+} X \cup \partial_{0} X \cup \partial_{-} X$ such that $\partial_{\bullet} X \simeq Y_{\bullet}$.
(d) a Lagrangian cobordism $L: \Lambda_{+} \rightarrow \Lambda_{-}$contained in the horizontal boundary $\partial_{0} X$ of $X$.

We will also need to refer to a symplectic trivialization of $\xi$ along the orbits and chords in $\Xi$. We fix the following terminology.

Definition 92 (Trivialization). A trivialization $\tau$ of $\xi$ over $\Xi$ is a symplectic bundle isomorphism

$$
\tau: \xi\left|\gamma_{i} \simeq \mathbb{C} \quad \tau:(\xi, T \Lambda)\right| c_{i} \simeq(\mathbb{C}, \mathbb{R})
$$

The difference $\sigma-\tau$ between two trivializations over $\Xi$ is defined by

$$
\sigma-\tau=\sum_{i} m_{i} \cdot \mu\left(\left.\sigma \circ \tau^{-1}\right|_{\gamma_{i}}(\mathbb{R})\right)+\sum_{i} n_{i} \cdot \mu\left(\left.\sigma \circ \tau^{-1}\right|_{c_{i}}(\mathbb{R})\right)
$$

We use $\mathcal{T}\left(\Xi_{+}, \Xi_{-}\right)$to denote the space of pairs $\tau=\left(\tau_{+}, \tau_{-}\right)$of trivializations $\tau_{+}$over $\Xi_{+}$and $\tau_{-}$over $\Xi_{-}$.

Note that a trivialization $\tau:\left.\xi\right|_{\gamma} \simeq \mathbb{C}$ of the contact structure over an orbit induces a trivialization over any iterate of $\gamma$. By abuse of notation, we also denote this trivialization by $\tau$.

## Surface Classes

We next discuss the set of surface classes associated to the pair $(X, L)$.
Definition 93. A proper surface $S=(\iota, S)$ in $(X, L)$ from $\Xi_{+}$to $\Xi_{-}$consists of a compact surface $S$ with boundary $\partial S=\partial_{+} S \cup \partial_{-} S \cup \partial_{0} S$ and a continuous map $\iota: S \rightarrow X$ such that
(a) The restriction of $\iota$ to $\partial_{ \pm} S$ is a map

$$
\iota: \partial_{ \pm} S \rightarrow Y_{ \pm}=\partial_{ \pm} X
$$

consisting of a collection of orbits and chords in $Y_{ \pm}$representing the orbit set $\Xi_{ \pm}$.
(b) The boundary region $\partial_{0} S$ maps to the Lagrangian $L$, i.e. $\iota\left(\partial_{0} S\right) \subset L$

A proper surface $(\iota, S)$ will be called well-immersed if it also satisfies the following criteria.
(a) The restriction of $\iota$ to $S \backslash\left(\partial_{+} S \cup \partial_{-} S\right)$ is an immersion with transverse double points. Note that these double points may occur on $\partial_{\circ} S$.
(b) The map $\iota$ is an embedding in a neighborhood of $\partial_{+} S \cup \partial_{-} S$ (except at $\partial_{+} S \cup \partial_{-} S$ itself) and is transverse to $\partial_{+} S \cup \partial_{-} S$.

Definition 94. The set $S\left(\Xi_{+}, \Xi_{-}\right)$of surface class $A: \Xi_{+} \rightarrow \Xi_{-}$is the set of proper maps $\iota: S \rightarrow X$ bounding $\Xi_{+}$and $\Xi_{-}$modulo the equivalence relation that

$$
\iota \sim \iota^{\prime} \quad \text { if } \quad\left[S \cup-S^{\prime}\right]=0 \in H_{2}(X, L)
$$

Here $S \cup-S^{\prime}$ is the 2-chain in $X$ with boundary on $L$ represented by the map $\iota \cup \iota^{\prime}: S \cup S^{\prime} \rightarrow X$.

The set of surface classes $S\left(\Xi_{+}, \Xi_{-}\right)$is (tautologically) a torsor over the relative homology group $H_{2}(X, L)$ and we use the notation

$$
A-B \in H_{2}(X, L) \quad \text { for the class such that } \quad A=B+(A-B)
$$

Example 95. Let $(Y, \Lambda)$ be a closed contact manifold and a Legendrian and fix an orbit-chord set $\Xi$. Consider the cobordism

$$
X:=Y \times[0,1] \quad L:=\Lambda \times[0,1]
$$

A trivial branched covers of $\Xi$ will refer to a map from a surface $C$ to $Y \times[0,1]$ of the form

$$
C \xrightarrow{\kappa} \Xi \times[0,1] \subset Y \times[0,1]
$$

Here $\kappa$ is a branched cover whose covering multiplicity at $\Xi_{+}$and $\Xi_{-}$satisfies Definition 94.

Any two collection of orbits and chords representing $\Xi$ (i.e. with the same underlying simple orbits and chords with multiplicity) can be connected by a trivial branched cover. Furthermore, the surface classes of all such maps all agree since $H_{2}(\Xi, \partial \Xi)=0$.

There are some important operations on proper surfaces and surface classes that we will use in later parts of this paper.

Definition 96 (Union). The union $S \cup T$ of two proper surfaces $(\iota . S)$ and $(\jmath, T)$ in $(X, L)$ is

$$
\iota \sqcup \jmath: S \sqcup T \rightarrow(X, L)
$$

This descends to a map on the space of surface classes of the form

$$
S\left(\Xi_{0}, \Theta_{0}\right) \times S\left(\Xi_{1}, \Theta_{1}\right) \rightarrow S\left(\Xi_{0} \cup \Xi_{1}, \Theta_{0} \cup \Theta_{1}\right) \quad \text { with } \quad(A, B) \mapsto A \cup B
$$

Definition 97. A composition $S \circ T$ of two proper surfaces $S: \Xi_{0} \rightarrow \Xi_{1}$ in $(X, L)$ and $T: \Xi_{1} \rightarrow \Xi_{2}$ in $\left(X^{\prime}, L^{\prime}\right)$ is proper surface of the form

$$
S \cup Z \cup T \subset X \circ X^{\prime}
$$

where $Z$ is a trivial branched cover of $\Xi$ connecting $\partial_{-} S$ and $\partial_{+} T$. This descends to a map of surface classes of the form

$$
S\left(\Xi_{0}, \Xi_{1}\right) \times S\left(\Xi_{1}, \Xi_{2}\right) \rightarrow S\left(\Xi_{0}, \Xi_{1}\right) \quad \text { given by } \quad(A, B) \rightarrow A \circ B
$$

## Maslov Number Of A Surface Class

Given any proper surface $S$ in $(X, L)$, there is a natural bundle pair with punctures over $S$ induced by pullback.

$$
\left(\iota^{*} T X, \iota^{*} T L\right) \rightarrow\left(S, \partial_{\circ} S\right)
$$

We extend the trivialization of $\xi_{c_{i}}$ to trivializations of $\left.\iota^{*} T X\right|_{c_{i}}$. And similarly we extend trivilizations of $\xi_{\gamma_{i}}$ to $\left.\iota^{*} T X\right|_{\gamma_{i}}$. We here describe our choice for such an extension (this is essentially the same choice has we have for Lemma 101.) We assume near $Y_{ \pm}$the cobordism $X$ has collar neighborhoods of the form $(1-\epsilon, 1] \times Y_{+}$ and $[0, \epsilon) \times Y_{-}$. We call the direction in the half open interval the symplectization direction. Then $\xi$ and the plane given by the symplectization direction and the Reeb direction on over chords and orbits span $T X$. We extend $\tau$ in this way to $T X$ over both chords and orbits. We further assume that with this choice of trivialization $\tau:\left.\left.\iota^{*}(T X)\right|_{c_{i}} \rightarrow \mathbb{C}^{2}\right|_{c_{i}}$, we have $\tau\left(\left.\iota^{*}(T X)\right|_{\partial_{ \pm} S \cap \partial_{o} S}=\mathbb{R}^{2} \subset \mathbb{C}^{2}\right.$.

Definition 98. The Maslov number $\mu_{\tau}(A, \tau)$ of a surface class $A$ with respect to a trivialization $\tau$ over $\Xi$ is given by

$$
\mu(A, \tau):=\mu\left(\iota^{*} T X, \iota^{*} T L ; \tau\right) \quad \text { for any representative } \quad \iota: S \rightarrow X
$$

Here $\mu\left(\iota^{*} T X, \iota^{*} T L ; \tau\right)$ is the Maslov index of $\left(\iota^{*} T X, \iota^{*} T L\right)$ as a bundle pair (see Proposition 84).

A priori, it is not clear that $\mu(A, \tau)$ is independent of the choice of representative. We now prove that this is the case.

Proposition 99. The Maslov number $\mu(A, \tau)$ is well-defined and has the following properties.

- (Trivialization) Let $\sigma$ and $\tau$ be two trivializations in $\mathcal{T}\left(\Xi_{+}, \Xi_{-}\right)$. Then

$$
\mu(A, \tau)-\mu(A, \sigma)=\left(\sigma_{+}-\tau_{+}\right)-\left(\sigma_{-}-\tau_{-}\right)
$$

- (Maslov Class) Let $A: \Xi \rightarrow \Theta$ and $B: \Xi \rightarrow \Theta$ be a pair of surface classes. Then

$$
\mu(A, \tau)-\mu(B, \tau)=\mu(X, L) \cdot(A-B)
$$

- (Union) Let $A: \Xi_{0} \rightarrow \Theta_{0}$ and $B: \Xi_{1} \rightarrow \Theta_{1}$ be a pair of surface classes and let $\tau$ be a trivialization that agrees on $\Xi_{0} \cap \Xi_{1}$ and $\Theta_{0} \cap \Theta_{1}$. Then

$$
\mu(A \cup B, \tau)=\mu(A, \tau)+\mu(B, \tau)
$$

- (Composition) Let $A: \Xi_{0} \rightarrow \Xi_{1}$ and $B: \Xi_{1} \rightarrow \Xi_{2}$ be composible surface classes, and let $\tau$ be a trivialization along $\Xi_{0} \cup \Xi_{1} \cup \Xi_{2}$. Then

$$
\mu(A \circ B, \tau)=\mu(A, \tau)+\mu(B, \tau)
$$

To proceed with the proof, we first consider the case of branched covers as in Example 95.

Lemma 100. Let $\Xi$ be an orbit-chord set in $Y$ and consider a trivial branched cover of $\Xi$, as in Example 95.

$$
\iota:\left(C, \partial_{\circ} C\right) \rightarrow(X, L)=([0,1] \times Y,[0,1] \times \Lambda)
$$

Choose a pair of trivializations $\sigma$ over $\Xi_{-}=\Xi \times\{0\}$ and $\tau$ over $\Xi_{+}=\Xi \times\{1\}$. Then

$$
\mu\left(\iota^{*} T X, \iota^{*} T L ; \iota^{*}(\sigma \cup \tau)\right)=\tau-\sigma
$$

Here $\iota^{*}(\sigma \cup \tau)$ is the trivialization of $\iota^{*} T X$ over $\partial_{\star} C=\partial_{+} C \cup \partial_{-} C$ induced by $\sigma$ and $\tau$.

Proof. The trivialization $\sigma$ extends to a trivialization

$$
(\xi, T \Lambda) \simeq(\mathbb{C}, \mathbb{R}) \quad \text { over } \quad \Xi \times[0,1]
$$

This pulls back to a trivialization of $\iota^{*} \xi$ over $C$. Then $\tau$ and $\sigma$ are identified, respectively, with

$$
\iota^{*}\left(\tau \circ \sigma^{-1}\right) \quad \text { along } \quad \partial_{+} C=\iota^{-1}\left(\partial_{+} X\right) \quad \text { and } \quad \text { Id } \quad \text { on } \quad \partial_{-} C=\iota^{-1}\left(\partial_{-} X\right)
$$

Now decompose $\partial_{+} C$ into regions $S_{i}=\iota^{-1}\left(\gamma_{i}\right)$ and $T_{i}=\iota^{-1}\left(c_{i}\right)$. Then as a simple application of Lemma 88, we find that

$$
\mu\left(\iota^{*} \xi, \iota^{*} T \Lambda ; \iota^{*}(\sigma \cup \tau)\right)=\sum_{i} \operatorname{deg}\left(\left.\iota\right|_{S_{i}}\right) \cdot \mu\left(\left.\tau \circ \sigma^{-1}\right|_{S_{i}}\right)+\operatorname{deg}\left(\left.\iota\right|_{T_{i}}\right) \cdot \mu\left(\left.\tau \circ \sigma^{-1}\right|_{T_{i}}\right)
$$

This is precisely the difference $\tau-\sigma$ of the trivializations over $\Xi$. The same formula follows for $\left(\iota^{*} T X, \iota^{*} T L\right)$ since it is isomorphic to the direct sum of $\left(\iota^{*} \xi, \iota^{*} T \Lambda\right)$ and a trivial bundle pair $(\mathbb{C}, \mathbb{R})$ (spanned by the $\mathbb{R}$-direction and Reeb direction in $Y \times$ $[0,1]$ ).

Proof. (Proposition 99) The union and composition properties follow immediately from the disjoint union and puncture gluing properties of the Maslov number (Proposition 84). Here we will argue the other two properties.

Well-Definedness and Maslov Class. We prove well-definedness and the Maslov class property together, since the argument is the same. Choose representatives of $A$ and $B$ in $S\left(\Xi_{+}, \Xi_{-}\right)$respectively.

$$
\iota: S \rightarrow X \quad \text { and } \quad \jmath: T \rightarrow X
$$

Since $S$ and $T$ bound the same orbit sets $\Xi_{+}$and $\Xi_{-}$, there exist trivial branched covers

$$
\kappa_{\bullet}: C_{\bullet} \rightarrow[0,1] \times \Xi_{\bullet} \subset[0,1] \times Y_{\bullet} \quad \text { for } \quad \bullet \in\{+,-\}
$$

such that $\kappa_{+}$connects the positive ends of $S$ and $T$ and $\kappa_{-}$connects the negative ends of $S$ and $T$. We can form a new surface

$$
\Sigma=S \cup C_{+} \cup C_{-} \cup-T
$$

by gluing $S$ and $T$ to $C_{+}$along their common orbit-chord ends, and likewise at $C_{-}$ (see Figure 4.7). This surface inherits a continuous map


Figure 4.7: The surface $\Sigma$

$$
\kappa:(\Sigma, \partial \Sigma) \rightarrow(X, L)
$$

restricting to $\iota$ on $S$, to $\jmath$ on $T$ and to the map $C_{ \pm} \rightarrow \Xi_{ \pm}$on $C_{ \pm}$. By the gluing property of the Maslov number in Proposition 84(d), we have
$\mu\left(\kappa^{*} T X, \kappa^{*} T L\right)=\mu\left(\iota^{*} T X, \iota^{*} T L ; \tau\right)-\mu\left(\jmath^{*} T X, \jmath^{*} T L ; \tau\right)+\mu\left(\left.T X\right|_{C_{+}}, F_{C_{+}}\right)+\mu\left(\left.T X\right|_{C_{-}}, F_{C_{-}}\right)$
The left hand side is simply $\mu(X, L) \cdot[\Sigma]$ where $\mu(X, L)$ is the Maslov class (see Example 87) and $[\Sigma]=A-B$. On the right hand side, the cylinder $C_{+}$is equipped with the trivialization $\left.\tau\right|_{\Xi_{+}}$on both ends and likewise for $C_{-}$, so by Lemma 100

$$
\mu\left(\left.T X\right|_{C_{+}}, F_{C_{+}}\right)=0 \quad \text { and } \quad \mu\left(\left.T X\right|_{C_{+}}, F_{C_{-}}\right)=0
$$

Therefore, (a) follows from (4.2.3).
Trivialization. Fix a representative $S$ of $A$ as in (a) and let $C_{ \pm}$be the trivial (unbranched cover) of $\Xi_{ \pm} \times[0,1]$, as in Example 95. Equip $C_{ \pm}$with the trivialization $\rho_{ \pm}$with

$$
\left.\rho_{ \pm}\right|_{\Xi_{ \pm} \times 0}=\left.\sigma\right|_{\Xi_{ \pm}} \quad \text { and }\left.\quad \rho_{ \pm}\right|_{\Xi_{ \pm} \times 1}=\left.\tau\right|_{\Xi_{ \pm}}
$$

We may form a glued surface $T=S \cup C_{+} \cup C_{-}$and a map $\jmath: T \rightarrow X$ representing the class $A$ by a similar process to (a). By applying the puncture gluing property Proposition 84(c), we have

$$
\mu(A, \tau)=\mu(A, \sigma)+\mu\left(C_{+}, \rho_{+}\right)-\mu\left(C_{-}, \rho_{-}\right)
$$

Therefore, the result follows from 100 since

$$
\mu\left(C_{+}, \rho_{+}\right)=\mu\left(\left.\sigma \circ \tau^{-1}\right|_{\Xi_{+}} \mathbb{R}\right) \quad \text { and } \quad \mu\left(C_{-}, \rho_{-}\right)=\mu\left(\left.\sigma \circ \tau^{-1}\right|_{\Xi-} \mathbb{R}\right)
$$

We conclude this section by considering the Maslov number of a symplectic, wellimmersed proper surface $(\iota, S)$ with class $A \in S\left(\Xi_{+}, \Xi_{-}\right)$. In this case, we have a decomposition

$$
\begin{equation*}
(T X, T L)=\left(T \Sigma, T\left(\partial_{\circ} \Sigma\right)\right) \oplus\left(\nu S, \nu\left(\partial_{\circ} S\right)\right) \tag{4.2.4}
\end{equation*}
$$

Here $\nu S$ is the symplectic perpendicular to $T \Sigma$ and $\nu(\partial S)$ is any choice of transverse sub-bundle to $T\left(\partial_{\circ} S\right)$ in $T L$.

Note that $\nu S$ is canonically identified with $\xi_{ \pm}$near the positive and negative ends of $(\iota, S)$. In particular, a choice of trivialization $\tau$ over $\Xi$ induces a puncture trivialization of $(T X, T L)$ (also denoted $\tau$ ) by direct sum with the canonical puncture trivialization of the tangent pair $\left(T \Sigma, T\left(\partial_{0} \Sigma\right)\right)$. By the direct sum property of the Maslov number, we thus acquire

Lemma 101. Let $(\iota, S)$ be a symplectic, well-immersed proper surface. Then

$$
\mu(S, \tau)=2 \bar{\chi}(S)+\mu\left(\nu S, \nu\left(\partial_{\circ} S\right) ; \tau\right)
$$

## Writhe and Linking

We next introduce the notion of the writhe and linking number of a braid, and of an admissible representative of a surface class.

To introduce these concepts, we must first clarify the notion of braid that we will use.

Definition 102. A braid $\zeta$ around a Reeb chord or orbit $\eta$ in $Y$ is a disjoint union of arcs (if $\eta$ is a chord) or of loops (if $\eta$ is an orbit) in a tubular neighborhood $\operatorname{Nbhd}(\eta)$ of $\eta$ such that

$$
\zeta \text { is transverse to } \xi \quad \text { and } \quad \partial \zeta \in \Lambda \text { if } \eta \text { is a chord }
$$

Definition 103. Let $S$ be a well immersed surface representative of a surface class $A \in S\left(\Xi_{+}, \Xi_{-}\right)$. Suppoose for each chord or orbit $\eta_{i}^{ \pm}$in the set $\Xi_{+}$, there is an associated (isotopy class of) braid

$$
\zeta_{i}^{ \pm} \quad \text { around } \quad \eta_{i}^{ \pm}
$$

constructed as follows. Choose a collar neighborhood of pairs

$$
\left(Y_{ \pm} \times[0,1), \Lambda \times[0,1)\right) \rightarrow(X, L) \quad \text { near } \quad \partial_{ \pm} X
$$

Then, under this identification, we define $\zeta_{i}^{ \pm}$to be the intersection

$$
\zeta_{i}^{ \pm}:=S \cap(\operatorname{Nbhd}(\eta) \times\{\epsilon\})
$$

Here $\operatorname{Nbhd}(\eta)$ is a small tubular neighborhood disjoint from other orbits and chords in $\Xi$ determined by a choice of trivialization of the contact structure along the orbit or chord:

$$
\phi_{\tau}: \operatorname{Nbhd}(\eta) \simeq[0,1] \times D^{2} \quad \text { or } \quad \phi_{\tau}: \operatorname{Nbhd}(\eta) \simeq S^{1} \times D^{2}
$$

In the chord case, $\phi_{\tau}$ maps the boundary $\partial \eta$ to the $x$-axis $(\mathbb{R} \times 0) \cap D^{2}$ in $0 \times D^{2}$ and $1 \times D^{2}$. We take $\epsilon>0$ to be small. Suppose the braids constructed as above does not depend on $\epsilon$ if it is small enough. Then we say the surface $S$ is an admissible representative. We collectively refer to these braids as the ends of $S$.

Given a trivialization $\tau$ along Reeb orbits or chords, as in the above definition, this gives rise to a choice of tubular neighborhood around the chord or Reeb orbit of the form $[0,1] \times D^{2}$ (resp. $S^{1} \times D^{2}$ ). Via the projection map $D^{2} \rightarrow \mathbb{R}$ to the $x$-axis, we acquire projection maps

$$
\pi:[0,1] \times D^{2} \rightarrow[0,1] \times \mathbb{R} \quad \text { and } \quad \pi: S^{1} \times D^{2} \rightarrow S^{1} \times \mathbb{R}
$$

Definition 104 (Writhe). The writhe $w_{\tau}(\zeta) \in \mathbb{Z}$ of a braid $\zeta$ around $\eta$ with respect to a trivialization $\tau$ is the signed count of self-intersections of the curve

$$
\pi \circ \phi_{\tau}(\widetilde{\zeta})
$$

Here $\widetilde{\zeta}$ is a perturbation of $\zeta$ relative to $\partial \zeta$ such that $\pi \circ \phi_{\tau}\left(\zeta^{\prime}\right)$ has only transverse double points. The sign convention in ECH is that anticlockwise rotations in the disk $D^{2}$ contributes a positive crossing. This is described, for instance after Definition 2.7 in [55]

Likewise, the writhe $w_{\tau}(S)$ of an admissible representative $S$ of a surface class $A$ with respect to a trivialization $\tau \in \mathcal{T}\left(\Xi_{+}, \Xi_{-}\right)$is the signed sum of writhes of the ends.

$$
\begin{equation*}
w_{\tau}(S):=\sum_{i} w_{\tau}\left(\zeta_{i}^{+}\right)-\sum_{j} w_{\tau}\left(\zeta_{j}^{-}\right) \tag{4.2.5}
\end{equation*}
$$

Definition 105. The linking number $l_{\tau}\left(\zeta, \zeta^{\prime}\right) \in \frac{1}{2} \mathbb{Z}$ of a pair of disjoint braids $\zeta$ and $\zeta^{\prime}$ around $\eta$ is half of the signed count of intersections between the pair of curves

$$
\pi \circ \phi_{\tau}(\zeta) \quad \text { and } \quad \pi \circ \phi_{\tau}\left(\zeta^{\prime}\right)
$$

The linking number $l_{\tau}\left(S, S^{\prime}\right)$ of a pair of disjoint, admissible representatives of two surface classes $A$ and $B$ is the analogous signed sum of linking numbers to (4.2.5).

We will require a few elementary properties of the linking number and the writhe.
Lemma 106. The linking number and the writhe satisfy the following properties.

- (Union) The linking number and writhe are related by the formula

$$
\begin{equation*}
w_{\tau}\left(\zeta \cup \zeta^{\prime}\right)=w_{\tau}(\zeta)+w_{\tau}\left(\zeta^{\prime}\right)+2 \cdot l_{\tau}\left(\zeta, \zeta^{\prime}\right) \tag{4.2.6}
\end{equation*}
$$

- (Trivialization) If $\sigma$ and $\tau$ are two trivializations of $\xi$ along $\eta$ then

$$
w_{\tau}(\zeta)-w_{\sigma}(\zeta)=m(m-1) \cdot(\tau-\sigma)
$$

where $m$ is the number of strands in the braid $\zeta$.

## Classical Intersection Pairing

We next consider a half-integer valued intersection pairing associated to a pair of a 4 -manifold with a 2 -manifold in its boundary. This is analogous to the intersection pairing of closed 4-manifolds.

Definition 107 (Intersection). Let $M$ be a compact 4-manifold with corners and let $N \subset \partial M$ be an oriented embedded 2-manifold with boundary. The intersection pairing

$$
Q: H_{2}(M, N) \otimes H_{2}(M, N) \rightarrow \frac{1}{2} \mathbb{Z}
$$

is defined as follows. Given classes $A, B \in H_{2}(M, N)$, choose immersed representatives

$$
\iota:(S, \partial S) \rightarrow(M, N) \quad \text { and } \quad \jmath:(T, \partial T) \rightarrow(M, N)
$$

of $A$ and $B$, respectively, that intersect transversely (including along $N$ ). Then each intersection point $p \in S \cap T$ is isolated and has a well-defined $\operatorname{sign} \iota(M, S, T, p)= \pm 1$ (cf. [121]). We let
$\#_{M}(\operatorname{int}(S) \cap \operatorname{int}(T))=\sum_{p \in \operatorname{int}(S) \cap \operatorname{int}(T)} \iota(M, S, T, p) \quad$ and $\quad \#_{M}(\partial S \cap \partial T)=\sum_{p \in \partial S \cap \partial T} \iota(M, S, T, p)$

The boundaries $\partial S$ and $\partial T$ are also transverse as sub-manifolds of $N$. Thus, we have well-defined intersect numbers in $N$ for each $p \in \partial S \cap \partial T$, denoted

$$
\iota(N, \partial S, \partial T, p)= \pm 1 \quad \text { and } \quad \#_{N}(\partial S \cap \partial T)=\sum_{p \in \partial S \cap \partial T} \iota(N, \partial S, \partial T, p)
$$

After an isotopy of $S$ (or $T$ ) that fixes $\partial S$, we may assume that

$$
\begin{equation*}
\iota(M, S, T, p)=\iota(N, \partial S, \partial T, p) \quad \text { for each } \quad p \in S \cap T \tag{4.2.7}
\end{equation*}
$$

Note that this isotopy may introduce intersections in the interior of $S$ and $T$, but we may perturb $S$ and $T$ to be transverse to each other. We now define the intersection pairing $Q$ as follows.

$$
\begin{equation*}
Q(A, B):=\#_{M}(\operatorname{int}(S) \cap \operatorname{int}(T))+\frac{1}{2} \cdot \#_{M}(\partial S \cap \partial T) \tag{4.2.8}
\end{equation*}
$$

In general, the intersection numbers $\#_{N}(\partial S \cap \partial T)$ and $\#_{M}(\partial S \cap \partial T)$ may not agree. However, we can always isotope $S$ so that this is the case.

Lemma 108. Let $p \in \partial S \cap \partial T$ be a boundary intersection and assume that

$$
\iota(M, S, T, p)=-\iota(N, \partial S, \partial T, p)
$$

Then there is a small neighborhood $U$ of $p$ and a surface $S^{\prime}$ isotopic to $S$ such that

$$
\partial S^{\prime}=\partial S \quad S=S^{\prime} \text { on } M \backslash U \quad S^{\prime} \cap T \cap U=\{p, q\}
$$

Here the boundary intersection point $p$ has intersection numbers

$$
\iota\left(N, \partial S^{\prime}, \partial T, p\right)=-\iota(N, \partial S, \partial T, p)
$$

and $q$ is an interior intersection point such that

$$
\iota\left(M, S^{\prime}, T, q\right)=\iota(M, S, T, p)
$$

Proof. By choosing coordinates and (possibly) reversing the orientation of $T$, we can reduce to the local picture where $M$ and $N$ are given by

$$
M:=\left\{\left(z_{1}, z_{2}\right): \operatorname{Im}\left(z_{1}+z_{2}\right) \geq 0\right\} \subset \mathbb{C}^{2} \quad \text { and } \quad N=\mathbb{R}^{2}
$$

Moreover, we may assume that $S$ and $T$ take the form

$$
S:=(\mathbb{C} \oplus 0) \cap M \quad \text { and } \quad T:=(0 \oplus \mathbb{C}) \cap M
$$

Now consider the sub-space

$$
S^{\prime \prime}:=\mathbb{R} \oplus i \mathbb{R} \cap M \quad \text { and } \quad T^{\prime \prime}:=i \mathbb{R} \oplus \mathbb{R} \cap M
$$

Note that $S^{\prime \prime} \cap T^{\prime \prime}=S \cap T=(0,0)$ and we have

$$
\#_{M}(S \cap T)=-\#_{M}\left(S^{\prime \prime} \cap T^{\prime \prime}\right)=1 \quad \text { and } \quad \#_{N}(S \cap T)=\#_{N}\left(S^{\prime \prime} \cap T^{\prime \prime}\right)=1
$$

Now note that there is a pair of braids in the hemisphere $S^{3} \cap M$ given by

$$
B=(S \cup T) \cap S^{3} \quad \text { and } \quad B^{\prime \prime}=\left(S^{\prime \prime} \cup T^{\prime \prime}\right) \cap S^{3} .
$$

We can define homotopies of the components $S_{t}$ from $S$ to $S^{\prime \prime}$ and $T_{t}$ from $T$ to $T^{\prime \prime}$ as follows.
$S_{t}=\operatorname{span}_{\mathbb{R}}\left(1 \oplus 0, v_{t}\right) \cap S^{3} \cap M \quad$ where $\quad v_{t}=\cos \left(\frac{\pi t}{2}\right) \cdot(i \oplus 0)+\sin \left(\frac{\pi t}{2}\right) \cdot(0 \oplus i)$
$T_{t}=\operatorname{span}_{\mathbb{R}}\left(0 \oplus 1, w_{t}\right) \cap S^{3} \cap M \quad$ where $\quad w_{t}=\cos \left(\frac{\pi t}{2}\right) \cdot(0 \oplus i)+\sin \left(\frac{\pi t}{2}\right) \cdot(i \oplus 0)$
Note that $S_{t}$ and $T_{t}$ are disjoint, except at $t=\frac{1}{2}$. We can use $S_{t}$ and $T_{t}$ to form a pair of surfaces

$$
\begin{aligned}
& \Sigma(S)=\left\{(2-t) \cdot z: 0 \leq t \leq 1 \text { and } z \in S_{t}\right\} \\
& \Sigma(T)=\left\{(2-t) \cdot z: 0 \leq t \leq 1 \text { and } z \in T_{t}\right\}
\end{aligned}
$$

These surfaces intersect at one point with sign -1 . We now let $S^{\prime}$ and $T^{\prime}$ be, respectively, smoothings of the $C^{0}$ embedded surfaces

$$
\left(S \cap B^{3}(1)\right) \cup \Sigma(S) \cup\left(S^{\prime \prime} \backslash B^{3}(2)\right) \quad \text { and } \quad\left(T \cap B^{3}(1)\right) \cup \Sigma(S) \cup\left(T^{\prime \prime} \backslash B^{3}(2)\right)
$$

These surfaces are smooth except along some curves that are disjoint from their intersections, so the intersections of $S^{\prime}$ and $T^{\prime}$ are given by

$$
\operatorname{int}\left(S^{\prime}\right) \cap \operatorname{int}\left(T^{\prime}\right)=S^{\prime \prime} \cap T^{\prime \prime}=\{0\} \quad \partial S^{\prime} \cap \partial T^{\prime}=\Sigma(S) \cap \Sigma(T)=\left\{\frac{3 i}{2} \cdot(1,1)\right\}
$$

This is precisely the local picture required in the lemma, so we are done.

Corollary 109. There are isotopic surfaces $S^{\prime}$ and $T^{\prime}$ to $S$ and $T$ respectively such that

$$
Q\left(S^{\prime}, T^{\prime}\right)=Q(S, T) \quad \text { and } \quad \#_{M}(\partial S \cap \partial T)=\#_{N}(\partial S \cap T \partial T)
$$

Proposition 110. The intersection pairing $Q$ in Definition 107 is well-defined.
Proof. Fix classes $A, B \in H_{2}(M, N)$. Choose immersed representatives

$$
\iota:(S, \partial S) \rightarrow(M, N) \quad \iota^{\prime}:\left(S^{\prime}, \partial S^{\prime}\right) \rightarrow(M, N) \quad \text { and } \quad \jmath:(T, \partial T) \rightarrow(M, N)
$$

of the classes $A, A$ and $B$, respectively, such that $S$ and $S^{\prime}$ are both transverse to $T$. Moreover, we may assume by Corollary 109 that

$$
\#_{M}(S \cap T)=\#_{N}(S \cap T) \quad \text { and } \quad \#_{M}\left(S^{\prime} \cap T\right)=\#_{N}\left(S^{\prime} \cap T\right)
$$

To prove well-definedness, it now suffices to show that

$$
\begin{equation*}
Q(S, T)-Q\left(S^{\prime}, T\right)=0 \tag{4.2.9}
\end{equation*}
$$

To prove (4.2.9), we must consider the standard short exact sequence of the pair $(M, N)$.

$$
\begin{equation*}
H_{2}(N) \rightarrow H_{2}(M) \rightarrow H_{2}(M, N) \rightarrow H_{1}(N) \tag{4.2.10}
\end{equation*}
$$

By hypothesis, we have an immersion of $S \cup-S^{\prime}$ in the homology class $\iota_{*}[S]-\jmath_{*}\left[S^{\prime}\right]=$ 0 in $H_{2}(M, N)$. Therefore, $\partial S \cup-\partial S^{\prime}$ is an immersed null-homologous 1-manifold in $N$, and

$$
\begin{equation*}
\#\left(\left(\partial S \cup-\partial S^{\prime}\right) \cap \partial T\right)=0 \tag{4.2.11}
\end{equation*}
$$

Furthermore, we may choose a map $\kappa: R \rightarrow N$ from a compact surface $R$ bounding $\partial S \cup-\partial S^{\prime}$, and acquire an immersion of a closed surface.

$$
f: \Sigma=S \cup R \cup-S^{\prime} \rightarrow X
$$

The class $[\Sigma] \in H_{2}(M)$ maps to $[S \cup-S]=0 \in H_{2}(M, N)$ under the map $H_{2}(M) \rightarrow$ $H_{2}(M, N)$ in (4.2.10). Thus by the exactness of (4.2.10), we know that $[\Sigma \cup Z]=0$ for some closed immersed surface $Z \rightarrow N$. We may assume that $[\Sigma]=0$ by absorbing $Z$ into the choice of bounding surface $R$.

Now choose a collar neighborhood $N \times(-1,1)_{t}$ of $N$ in $\partial M$ and extend this to a collar into $M$ as $U=N \times(-1,1)_{t} \times[0,1)_{s}$. By choosing this collar to be very small, we can assume that

$$
U \cap(f(\Sigma) \cap \jmath(T))=\jmath(T) \cap(R \times 0 \times 0) \subset N \times 0 \times 0
$$

We can choose a vector-field $v=\phi(s) \cdot \partial_{t}$ where $\phi:[0,1) \rightarrow[0,1)$ is a non-negative function with $\phi=0$ near $s=1$. By flowing $\Sigma$ along $v$ for a small amount of time, we perturb $f$ to a new map $f^{\prime}$ such that

$$
f^{\prime}(\Sigma) \cap T \cap U=\emptyset
$$

Finally, we can smooth $f^{\prime}$ to a smooth immersion $\phi: \Sigma \rightarrow \operatorname{int}(X)$ agreeing with $f^{\prime}$ away from a neighborhood of $\partial X$ that satisfies

$$
\#(\Sigma \cap T)=\#(S \cap T)-\#\left(S^{\prime} \cap T\right)=Q(S, T)-Q\left(S^{\prime}, T\right)
$$

We now simply observe that since $[\Sigma]=[S]-\left[S^{\prime}\right]=0$, we have

$$
Q(S, T)-Q\left(S^{\prime}, T\right)=\#(\Sigma \cap T)=\operatorname{PD}[T] \cdot[\Sigma]=0 \text { where }[T] \in H_{2}(X, \partial X)
$$

There is also an analogue of the intersection pairing for a 3-manifold with boundary, equipped with a 1-manifold in the boundary.

Definition 111. Let $Y$ be a compact 3 -manifold with boundary and corners, and let $Z \subset \partial Y$ be a closed 1-manifold in $\partial Y$. The intersection pairing

$$
Q: H_{1}(Y, Z) \otimes H_{2}(Y, Z) \rightarrow \frac{1}{2} \mathbb{Z}
$$

is defined as follows. Consider the 4-manifold

$$
M=[0,1] \times Y \quad \text { with sub-manifold } \quad N=[0,1] \times Z \subset \partial X
$$

Let $S=[0,1] \times \Gamma$ where $\Gamma \subset Y$ is a 1-manifold representing a class $A \in H_{1}(Y, Z)$ and let $T$ be a 2-manifold representing $B$ in $H_{2}(M, N) \simeq H_{2}(Y, Z)$ contained in $s \times Y$ for $s \in(0,1)$, and intersecting $\Gamma$ transversely.

$$
Q(A, B):=\#_{M}(\operatorname{int}(S) \cap \operatorname{int}(T))+\frac{1}{2} \cdot \#(\partial S \cap \partial T)
$$

The proof of well-definedness is analogous to Proposition 110.

## Relative Intersection Number

Next, we generalize the classical intersection number to an intersection number for surface classes.

Definition 112. Let $(X, L)$ be a pair of a symplectic cobordism with boundary and a Lagrangian cobordism $L \subset X$, and let $A \in S(\Xi, \Theta)$ be a surface class. The associated intersection map

$$
q_{A}: H_{2}(X, L) \rightarrow \frac{1}{2} \mathbb{Z} \quad \text { given by } \quad B \mapsto q_{A}(B)
$$

is defined as the same count of intersections in (4.2.8) where $S$ is a representative of the surface class $A$ and $T$ is a compact surface with boundary in the interior of $X$ representing $B$.

The proof of well-definedness is analogous to Proposition 110. Moreover, we can prove the following lemma.

Lemma 113. Consider a trivial symplectic cobordism pair ( $X, L$ ) of the form
$([0,1] \times Y,[0,1] \times \Lambda) \quad$ for a contact manifold $Y$ with Legendrian $\Lambda \subset \partial Y$
Let $A \in S(\Xi, \Theta)$ be any surface class between orbit-chord sets in 1-homology class $\Gamma=[\Xi]=[\Theta]$. Then

$$
q_{A}(B)=Q(\Gamma, B) \quad \text { for any } \quad B \in H_{2}(X, L) \simeq H_{2}(Y, \Lambda)
$$

Proof. Choose a representative surface $S \subset[0,1] \times Y$ of the surface class $A$ that is a cylinder $[0, \epsilon) \times \eta$ over a 1 -manifold $\eta$ representing $\Gamma$ in $[0, \epsilon) \times Y$. We may choose an immersed 2-manifold $T \subset Y$ with boundary $\partial T \subset \Lambda$ transverse to $\eta$. Then the corresponding count of intersections computes both $q_{A}(B)$ and $Q(\Gamma, B)$.

More generally, we define the intersection number between two surface classes as follows.

Definition 114 (Relative Intersection). Fix a pair of surface classes

$$
A: \Xi \rightarrow \Theta \quad \text { and } \quad B: \Xi^{\prime} \rightarrow \Theta^{\prime}
$$

and a trivialization $\tau$ of $\xi$ along $\Xi \cup \Xi^{\prime}$ and $\Theta \cup \Theta^{\prime}$. The relative intersection pairing of $A$ and $B$ with respect to $\tau$ is the half-integer

$$
Q_{\tau}(A, B) \in \frac{1}{2} \mathbb{Z}
$$

is defined as follows. Pick admissible surfaces $S$ and $T$ representing $A$ and $B$, respectively.

$$
\iota: S \rightarrow X \quad \text { and } \quad \jmath: T \rightarrow X
$$

Assume that $S$ and $T$ are disjoint near $\partial_{ \pm} X$ (except along $\partial_{ \pm} X$ ) and transversely intersecting away from $\partial_{+} X \cup \partial_{-} X$. Then we let

$$
\begin{equation*}
Q_{\tau}(A, B):=\#(\operatorname{int}(S) \cap \operatorname{int}(T))+\frac{1}{2} \cdot \#\left(\partial_{0} S \cap \partial_{\circ} T\right)-l_{\tau}(S, T) \tag{4.2.12}
\end{equation*}
$$

The relative intersection number satisfies a number of very useful axioms. We now prove these axioms in detail.

Proposition 115. The relative intersection pairing $Q_{\tau}$ is well-defined and has the following properties.

- (Trivialization) Let $A: \Xi_{+} \rightarrow \Xi_{-}$and $B: \Theta_{+} \rightarrow \Theta_{-}$be two surface classes and let $\sigma$ and $\tau$ be two trivializations that differ only along one orbit or chord

$$
\eta \text { of multiplicity } m \text { in } \Xi \text { and } n \text { in } \Theta
$$

Then the self-intersection numbers of $\sigma$ and $\tau$ differ as follows.

$$
Q_{\tau}(A, B)-Q_{\sigma}(A, B)=m \cdot n \cdot(\tau-\sigma)
$$

- (Difference) Let $A, A^{\prime}: \Xi_{+} \rightarrow \Xi_{-}$and $B \in S\left(\Psi_{+}, \Psi_{-}\right)$be surface classes between the same orbit-chord sets. Then

$$
Q_{\tau}(A, B)-Q_{\tau}\left(A^{\prime}, B\right)=q_{B}\left(A-A^{\prime}\right)
$$

- (Union) Let $A, A^{\prime}$ and $B$ be surface classes. Then

$$
Q_{\tau}\left(A \cup A^{\prime}, B\right)=Q_{\tau}(A, B)+Q_{\tau}\left(A^{\prime}, B\right)
$$

- (Composition) Let $A, A^{\prime}: \Xi_{0} \rightarrow \Xi_{1}$ and $B, B^{\prime}: \Xi_{1} \rightarrow \Xi_{2}$ be two pairs of composible surface classes in $(X, L)$ and $\left(X^{\prime}, L^{\prime}\right)$ respectively. Then

$$
Q_{\tau}\left(A \circ B, A^{\prime} \circ B^{\prime}\right)=Q_{\tau}\left(A, A^{\prime}\right)+Q_{\tau}\left(B, B^{\prime}\right)
$$

In order to prove these axioms, we will need the following lemma.

Lemma 116. Let $\eta$ be a Reeb chord or orbit. Let $\zeta^{+}, \zeta^{-}$and $\beta$ be disjoint braids such that $\zeta^{+}$and $\zeta^{-}$have the same degree. Then there exists a symplectically immersed cobordism $Z$ from $\zeta^{+}$to $\zeta^{-}$such that

$$
\#([0,1] \times \beta \cap Z)+\frac{1}{2} \#\left([0,1] \times \partial \beta \cap \partial_{\circ} Z\right)=l_{\tau}\left(\zeta^{+}, \beta\right)-l_{\tau}\left(\zeta^{-}, \beta\right)
$$

Proof. We prove the result for braids near a Reeb chord, as the Reeb orbit case can be treated identically. It suffices to consider the following setup. Let

$$
Y:=[0,1]_{t} \times D^{2} \quad \text { and } \quad \Lambda:=0 \times I \cup 1 \times I
$$

Here $I:=(\mathbb{R} \times 0) \cap D^{2}$ is the segment of the x -axis on the disk. We consider the braid diagram given by the projection onto the $(x, t)$-plane.

$$
\pi: Y \subset \mathbb{R}^{3} \rightarrow \mathbb{R}^{2}
$$

Recall that any two braids of the same degree can be related by a series of braid isotopies and the following crossing changes (and the reverse moves).


We call these Type 1 and Type 2 moves. A sequence of braid isotopies, Type 1 and Type 2 moves can be viewed as a regular homotopy and thus lifted to an immersed cobordism

$$
\iota: \Sigma \rightarrow[0,1]_{s} \times Y
$$

from the initial braid at $s=0$ to the final braid at $s=1$. This immersed cobordism is symplectic as in Lemma 2.4 of Golla-Etnyre [37]. A Type 1 move yields a transverse interior self-intersection of $\Sigma$ with sign -1 and a Type 2 move yields a transverse boundary self-intersection of $\Sigma$ with sign -1 . The reverse moves yield sign +1 .

Now choose a sequence of isotopies and Type $1 / 2$ moves from $\zeta^{+} \cup \beta$ to $\zeta^{-} \cup \beta$. We may assume that these moves leave $\beta$ fixed, so that the corresponding cobordism $\Sigma$ decomposes as

$$
\Sigma=Z \cup[0,1] \times \beta
$$

Here $Z$ is an immersed cobordism from $\zeta^{+}$to $\zeta^{-}$induced by a regular homotopy of braids $\zeta_{t}$ and intersections between $Z$ and $[0,1] \times \beta$ correspond to Type 1 and 2 moves between the braids $\zeta_{t}$ and $\beta$. A Type 1 move adds -1 intersections and changes the linking number by -1 . A Type 2 move adds $-1 / 2$ intersections and changes the linking number by $-1 / 2$. This yields the result.

Proof. (Proposition 115) We demonstrate each property individually. Note that the argument for the difference property suffices to prove well-definedness, since the latter follows from the same argument by taking $A=A^{\prime}$.

Trivialization. This follows immediately from the corresponding transformation law for the linking number.

Well-Definedness And Difference. Let $T$ be an admissible representative of a class $B$ in $S\left(\Psi_{+}, \Psi_{-}\right)$and let

$$
\iota: S \rightarrow X \quad \text { and } \quad \iota^{\prime}: S^{\prime} \rightarrow X
$$

be two admissible representatives of classes $A$ and $A^{\prime}$ in $S\left(\Xi_{+}, \Xi_{-}\right)$, respectively, that satisfy the transversality conditions in Definition 114 with respect to $T$. We construct an immersion $\phi:(\Sigma, \partial \Sigma) \rightarrow(X, L)$ transverse to $T$ such that

$$
\begin{equation*}
[\Sigma]=A-A^{\prime} \in H_{2}(X, L) \quad \text { and } \quad \#(\Sigma \cap T)+\frac{1}{2} \#(\partial \Sigma \cap \partial T)=Q_{\tau}(S, T)-Q_{\tau}\left(S^{\prime}, T\right) \tag{4.2.13}
\end{equation*}
$$

The result will then follow since $q_{B}([\Sigma])$ is precisely given by $\#(\Sigma \cap T)+\frac{1}{2} \#(\partial \Sigma \cap \partial T)$.
To begin the construction, fix the following notation for a collar near $\partial_{ \pm} X$.

$$
V_{ \pm}:=\partial_{ \pm} X \times[0, \mp \epsilon] \subset X \quad \text { and } \quad U=\overline{X \backslash\left(V_{+} \cup V_{-}\right)} \text {with } \partial_{ \pm} U:=\partial_{ \pm} X \times\{\mp \epsilon\}
$$

Here $\epsilon$ is an arbitrarily small parameter. Note that $\iota^{-1}(U) \subset S$ and $\left[\iota^{\prime}\right]^{-1}(U) \subset S^{\prime}$ are equal to $S$ and $S^{\prime}$ minus collars near $\partial_{ \pm} S$ and $\partial_{ \pm} S^{\prime}$, respectively. The underlying surface $\Sigma$ is of the form

$$
\Sigma=\iota^{-1}(U) \cup Z_{+} \cup Z_{-} \cup-\left[\iota^{\prime}\right]^{-1}(U)
$$

Here $Z_{ \pm}$is a surface with boundary and corners that we will specify shortly.
To define the immersion $\phi$ (and in the process, $Z_{ \pm}$), we proceed as follows. First, we let

$$
\begin{equation*}
\phi=\iota \text { on } \iota^{-1}(U) \quad \text { and } \quad \phi=\iota^{\prime} \text { on }\left[\iota^{\prime}\right]^{-1}(U) \tag{4.2.14}
\end{equation*}
$$

Note that the image of $\phi$ along $\partial_{ \pm} U$ consists of the braids at the positive and negative ends of $S$ and $S^{\prime}$, and the remaining boundary lies in $L$. Next, we assert that $\phi\left(Z_{ \pm}\right)$ is contained in a union of disjoint tubular neighborhoods of the orbits and chords in $\Xi_{ \pm}$.

$$
\begin{equation*}
\phi\left(Z_{ \pm}\right) \subset \bigsqcup_{i} \operatorname{Nbhd}\left(\eta_{i}^{ \pm}\right) \times[0, \mp \epsilon] \subset V_{ \pm} \tag{4.2.15}
\end{equation*}
$$

To describe $\phi$ in each neighborhood, fix an chord or orbit $\eta=\eta_{i}^{ \pm}$in $\Xi_{ \pm}$and a tubular neighborhood $N=\operatorname{Nbhd}\left(\eta_{i}^{ \pm}\right)$of $\eta$ in $Y_{ \pm}$. Denote the braids of $S, S^{\prime}$ and $T$ around $\eta_{i}^{ \pm}$by $\zeta, \zeta^{\prime}$ and $\beta$ respectively. Note that the braid $\beta$ will be empty if $\eta$ is not in the orbit set $\Psi$. For definiteness let's focus on what happens on $V_{+}$. By Lemma 116, we may choose a cobordism $Z$ in $N \times[0,-\epsilon)$ from $\zeta$ to $\zeta^{\prime}$ with

$$
\#(\operatorname{int}(Z) \cap \beta \times[0,-\epsilon])+\frac{1}{2} \cdot \#\left(\partial_{\circ} Z \cap \beta \times[0, \epsilon]\right)=l_{\tau}(\zeta, \beta)-l_{\tau}\left(\zeta^{\prime}, \beta\right)
$$

The definition for $V_{-}$is similar. We thus may associate a surface $Z_{i}^{ \pm}$to the end chord or orbit $\eta_{i}^{ \pm}$, namely

$$
Z_{i}^{ \pm}=Z \cup \zeta_{ \pm}^{\prime} \times[0, \epsilon]
$$

We may view $Z_{i}^{ \pm}$as a smooth surface with boundary and corners that is topologically embedded into $V_{ \pm}$and meets $\iota^{-1}(U)$ and $\left[\iota^{\prime}\right]^{-1}(U)$ along $\partial_{ \pm} U$. We then let

$$
\begin{equation*}
Z_{ \pm}:=\bigsqcup_{i} Z_{i}^{ \pm} \tag{4.2.16}
\end{equation*}
$$

and we let $\left.\phi\right|_{Z_{ \pm}}$be given by a smoothing of the tautological map $Z_{ \pm} \rightarrow V_{ \pm}$. Note that since the asymptotic braids $\beta$ and $\zeta^{\prime}$ are disjoint, we have

$$
\begin{equation*}
\#\left(\left(Z_{+} \cup Z_{-}\right) \cap T\right)=l_{\tau}(\zeta, \beta)-l_{\tau}\left(\zeta^{\prime}, \beta\right) \tag{4.2.17}
\end{equation*}
$$

The $\operatorname{map} \phi: \Sigma \rightarrow \mathbb{Z}$ is smoothly immersed and satisfies the criteria in (4.2.13) by construction. This concludes the proof of well-definedness and the difference property.

Union. We can choose well-immersed representatives $S, S^{\prime}$ and $T$ of $A, A^{\prime}$ and $B$ respectively so that $S, S^{\prime}$ and $S \cup S^{\prime}$ are transverse to $T$ away from $\partial_{ \pm} X$, disjoint from $T$ near (but not along) $\partial_{ \pm} X$ and transverse to $\partial_{ \pm} X$. Then

$$
Q_{\tau}\left(A \cup A^{\prime}, B\right)=\#\left(\operatorname{int}\left(S \cup S^{\prime}\right) \cap \operatorname{int}(T)\right)+\frac{1}{2} \cdot \#\left(\partial_{\circ}\left(S \cup S^{\prime}\right) \cap \partial_{\circ} T\right)-l_{\tau}\left(S \cup S^{\prime}, T\right)
$$

$$
=Q_{\tau}(A, B)+Q_{\tau}\left(A^{\prime}, B\right)
$$

Composition. Choose representatives $S, S^{\prime}, T$ and $T^{\prime}$ of $A, A^{\prime}, B$ and $B^{\prime}$, respectively. Via Lemma 116, we may assume that the braid of $S$ at the positive end agrees with the braid of $S^{\prime}$ at the negative end, and likewise for $T$ and $T^{\prime}$. Then

$$
S \cup T \quad \text { and } \quad S^{\prime} \cup T^{\prime}
$$

are representatives in $\left(X \circ X^{\prime}, L \circ L^{\prime}\right)$ of $A \circ B$ and $A^{\prime} \circ B^{\prime}$, respectively. We can smooth both to surfaces $U$ and $V$ so that near $\partial_{-} X=\partial_{+} X^{\prime}$, they agree with the negative braid of $S \cup S^{\prime}$ (or equivalently, the positive braid of $T \cup T^{\prime}$ ). Then we have

$$
\begin{gathered}
\#\left(\operatorname{int}(U) \cap \operatorname{int}(V)=\#\left(\operatorname{int}(S) \cap \operatorname{int}\left(S^{\prime}\right)\right)+\#\left(\operatorname{int}(T) \cap \operatorname{int}\left(T^{\prime}\right)\right)\right. \\
\#\left(\partial_{\circ}(U) \cap \partial_{\circ}(V)=\#\left(\partial_{\circ}(S) \cap \partial_{\circ}\left(S^{\prime}\right)\right)+\#\left(\partial_{\circ}(T) \cap \partial_{\circ}\left(T^{\prime}\right)\right)\right. \\
l_{\tau}(U, V)=l_{\tau}\left(S, S^{\prime}\right)+l_{\tau}\left(T, T^{\prime}\right)
\end{gathered}
$$

Note that the last formula above involves the cancelation of the linking of the negative ends of $S$ and $S^{\prime}$ with the linking at the positive ends of $T$ and $T^{\prime}$. This proves the result.

## Topological Adjunction

We are now ready to prove a topological version of Legendrian adjunction. The holomorphic curve version will be proven in Section 4.3 after the appropriate discussion of background.

Theorem 117 (Topological Adjunction). Let $S: \Xi \rightarrow \Theta$ be a symplectic, wellimmersed surface in a symplectic cobordism $(X, L)$ with a trivialization $\tau$ along $\partial S$. Then $S$ satisfies the adjunction formula

$$
\mu(S, \tau)=2\left(\bar{\chi}(A)+Q_{\tau}(S)+w_{\tau}(S)-2 \delta(S)-\epsilon(S)\right)
$$

If $S$ is simply smooth (and not necessarily symplectic) then the adjunction formula holds mod 2.

Proof. By Lemma 101, we simply need to show that

$$
\begin{equation*}
\mu_{\tau}\left(\nu S, \nu\left(\partial_{0} S\right)\right)=2\left(Q_{\tau}(S, S)+w_{\tau}(S)-2 \delta(S)-\epsilon(S)\right) \tag{4.2.18}
\end{equation*}
$$

where $\nu\left(\partial_{\circ} S\right)$ is the normal bundle of $\partial_{\circ} S \subset L$ as before. To prove (4.2.18), choose a section
$\phi:(S, \partial S) \rightarrow\left(\nu S, \nu\left(\partial_{\circ} S\right)\right) \quad$ satisfying $\left.\quad \phi\right|_{\partial_{ \pm} S} \neq 0 \quad \phi \pitchfork \partial_{\circ} \Sigma \quad$ and $\quad \phi \pitchfork \operatorname{int}(\Sigma)$
Also assume that, near the Reeb chords and orbits in $\partial S$, the section $\phi$ is constant with respect to our chosen trivialization $\tau$. By Lemma 89, the Maslov index is given by the following count of intersections

$$
\begin{equation*}
\mu_{\tau}\left(\nu S, \nu\left(\partial_{\circ} S\right)\right)=2 \cdot \#(\phi \cap \operatorname{int}(S))+\#\left(\phi \cap \partial_{\circ} S\right) . \tag{4.2.19}
\end{equation*}
$$

On the other hand, let $S^{\prime}$ be the perturbation of $S$ by the section $\phi$ in a neighborhood of $S$. In both $\operatorname{int}(S)$ and $\partial_{\circ} S$, there is a single transverse intersection in $S \cap S^{\prime}$ for each 0 of $\phi$ and two intersections in $S \cap S^{\prime}$ each transverse double point. Thus

$$
\begin{gather*}
\#(\phi \cap \operatorname{int}(S))=\#\left(\operatorname{int}\left(S^{\prime}\right) \cap \operatorname{int}(S)\right)-2 \delta(S)  \tag{4.2.20}\\
\#\left(\phi \cap \partial_{\circ} S\right)=\#\left(\partial_{\circ} S^{\prime} \cap \partial_{\circ} S\right)-2 \epsilon(S) \tag{4.2.21}
\end{gather*}
$$

Finally, we note that by Definition 114, we know that

$$
\begin{equation*}
Q_{\tau}(S, S)+w_{\tau}(S)=Q_{\tau}(S, S)+l_{\tau}\left(S, S^{\prime}\right)=\#(\phi \cap \operatorname{int}(S))+\frac{1}{2} \cdot \#\left(\partial_{\circ} S^{\prime} \cap \partial_{\circ} S\right) \tag{4.2.22}
\end{equation*}
$$

Combining (4.2.19)-(4.2.22) proves (4.2.18), and thus the proposition.
If $S$ is simply a smooth, well-immersed surface then we still have a decomposition of oriented, real bundle pairs

$$
(T X, T L)=\left(T S, T\left(\partial_{\circ} S\right)\right) \oplus\left(\nu S, \nu\left(\partial_{\circ} S\right)\right)
$$

After isotopy of $S$ near the ends, we may assume that this splitting respects the symplectic structure near $\partial S$. Modifying the symplectic stricture on $\left.T X\right|_{S}$ away from $\partial S$ is equivalent to direct summing with a symplectic bundle on $S^{2}$, and thus changes $\mu(T X, T L)$ by an even integer. Therefore

$$
\mu_{\tau}\left(\left.T X\right|_{S},\left.T L\right|_{\partial S}\right)=\mu_{\tau}\left(T S, T\left(\partial_{0} S\right)\right)+\mu\left(\nu S, \nu\left(\partial_{0} S\right)\right)
$$

The proof now reduces to (4.2.18), and proceeds by essentially the same argument.

### 4.3 J-Holomorphic Curves with Boundary

In this section, we examine holomorphic curves with Legendrian boundary conditions in convex sutured contact manifolds.

## Sutured Contact Manifolds

We stary by reviewing contact manifolds with sutured boundary, and the appropriate classes of contact forms and complex structures.

Remark 118. Our setting is most similar to [28]. We will only state our definitions for 3 - manifolds, but we note many of our definitions have higher dimensional versions, as found in [28].

Definition 119. A sutured 3-manifold $Y$ is a compact 3-manifold $Y$ with boundary and corners, a closed sub-manifold $\Gamma \subset \partial Y$ of codimension 2 and a neighborhood of $\Gamma$ of the form

$$
\begin{equation*}
V(\Gamma) \simeq[-1,1]_{t} \times(-\epsilon, 0]_{\tau} \times \Gamma \quad \text { with } \quad \Gamma \simeq 0 \times 0 \times \Gamma \tag{4.3.1}
\end{equation*}
$$

The boundary $\partial Y$ must divide into smooth strata $\partial_{-} Y, \partial_{\sigma} Y$ and $\partial_{+} Y$ such that, in the chart (4.3.1)

$$
\partial_{\sigma} Y \simeq[-1,1] \times 0 \times \Gamma \quad \text { and } \quad \partial_{ \pm} Y \cap V(\Gamma) \simeq \pm 1 \times(-\epsilon, 0] \times \Gamma
$$

Definition 120. [28, §2] A contact form $\alpha$ on a sutured manifold $Y$ is adapted to $Y$ if
(a) $\left.\alpha\right|_{\partial_{ \pm} Y}$ is a Liouville form on $\partial_{ \pm} Y$.
(b) In the neighborhood $V(\Gamma)$ in (4.3.1), $\alpha$ is given by

$$
\alpha=C \cdot d t+e^{\tau} \cdot \beta
$$

where $C>0$ is a constant and $\beta$ is a one-form on $\Gamma$ independent of $t$ and has no $d t$-term. Consequently in this neighborhood the Reeb vector field is given by $\frac{1}{C} \partial_{t}$.

Remark 121. When $Y$ is equipped with an adapted contact form, we may extend $t$ to a function

$$
t: \operatorname{Nbhd}(\partial Y) \rightarrow \mathbb{R}
$$

on a neighborhood of $\partial Y$ by setting $t\left(\partial_{ \pm} Y\right)= \pm 1$ and then extending $t$ to be Reeb invariant near $\partial_{ \pm} Y$. We will regularly use this extension without further comment.
Definition 122. A (convex) sutured contact manifold $(Y, \xi)$ is a sutured manifold $Y$ and a contact structure $\xi$ on $Y$ that is the kernel of an adapted contact form.

Definition 123. A collection of Legendrians $\Lambda$ in the horizontal boundary of the convex sutured convex contact manifold is called exact if the Liouville form $\beta_{ \pm}$ vanishes on $\Lambda$.

Consequently on the exact Legendrians the contact distribution $\xi$ is tangent to the horizontal boundary.

Remark 124. We give some examples where one can find exact Legendrians. It suffices to find a two dimensional Liouville manifold $(\Sigma, \beta)$ where $\beta$ is the Liouville form, and a collection of Lagrangians $\left\{L_{i}\right\}$ on which $\beta$ vanishes. The simplest possible example is $\left(D^{*} S^{1}, \beta\right)$, the codisk bundle of $S^{1}$; and $\beta$ is the canonical one form. The Lagrangian in question is then the zero section.

More generally speaking, let $\left\{L_{i}\right\}$ denote a collection of circles. Consider the Liouville manifolds $\left(D^{*} L_{i}, \beta_{i}\right)$ as above. We attach one handles between this collection of Liouville manifolds to form the Livioulle manifold $(S, \beta)$. Then the collection of curves $\left\{L_{i}\right\}$ can be taken to be the required collection of Lagrangians.

Definition 125. [28, §3.1] A complex structure $J$ on $\xi$ that is tailored to $(Y, \Lambda)$ if
(a) $J$ is Reeb invariant near $\partial Y$.
(b) Consider the completion of $\partial_{ \pm} Y$ near $V(\Gamma)$ given by $\partial_{ \pm} Y \cup[0, \infty) \times \Gamma$, with Liouville form $e^{\tau} \beta$ in a neighborhood of the form $(-\epsilon, \infty) \times\{ \pm 1\} \times \Gamma$. We require that in both $V(\Gamma)$ and small tubular neighborhoods $\partial_{+} Y \times[1,1-\epsilon]$ and $\partial_{-} Y \times[-1,-1+\epsilon]$, the almost complex structure $J$ is the pullback via the natural projection $\pi: \xi \rightarrow T\left(\partial_{ \pm} Y\right)$ of a complex structure $J_{0}$ that is $\beta$ compatible on $(-\epsilon, \infty) \times \Gamma$, and compatible with $\left.\alpha\right|_{\partial_{ \pm} Y}$ on $\partial_{ \pm} Y$.

We let $\widehat{J}$ denote the complex structure induced on $\mathbb{R} \times Y=\mathbb{R}_{s} \times Y$, given by

$$
\widehat{J}\left(\partial_{s}\right)=\left.R \quad \widehat{J}\right|_{\xi} \text { is a tailored complex structure } J
$$

Remark 126. Given a sutured contact manifold, it is also helpful to think about its completion as in Section 2.4 in [28]. First, "vertically" complete $V(\Gamma)$ by gluing $[1, \infty) \times \partial_{+} Y$ and $(-\infty,-1] \times \partial_{-} Y$ with the forms $C d t+\left.\alpha\right|_{\partial_{+} Y}$ and $C d t+\left.\alpha\right|_{\partial_{-} Y}$ respectively. Now the boundary is $\{0\} \times \mathbb{R} \times \Gamma$. Second, "horizontally" complete by gluing $[0, \infty) \times \mathbb{R} \times \Gamma$ with the form $C d t+e^{\tau} \beta$. We denote the completion as $\left(M^{*}, \alpha^{*}\right)$.

## Holomorphic Maps

We now recall the basic definitions regarding $J$-holomorphic maps in SFT, as required in this paper.

Let $(\Sigma, j)$ be a Riemann surface with boundary and corners, and assume that $\Sigma$ is equipped with a boundary decomposition into smooth components

$$
\partial \Sigma=\partial_{+} \Sigma \cup \partial_{0} \Sigma \cup \partial_{-} \Sigma
$$

We think of of $\partial_{0} \Sigma$ as boundaries of the surface. For $\partial_{ \pm} \Sigma$, we think of $\Sigma$ as surface with boundary/interior punctures, and each near boundary puncture we equip the puncture with a semi-infinite strip-like (resp. cylindrical) neighborhood $[0, \pm \infty] \times$ $[0,1]$ (resp. $[0, \infty] \times S^{1}$, and we think of $\partial_{ \pm} \Sigma$ as the components of the boundary at infinity, of the form $\pm \infty \times[0,1]$ (resp. $\pm \infty \times S^{1}$ ).

Let $(Y, \xi)$ be a convex sutured contact manifold with closed Legendrians $\Lambda \subset$ $\partial Y_{ \pm}$and fix an adapted contact form $\alpha$ and tailored complex structure $J$. A $J$ holomorphic map

$$
u:\left(\Sigma, \partial_{0} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)
$$

is a smooth map $u$ from $\Sigma \backslash\left(\partial_{+} \Sigma \cup \partial_{-} \Sigma\right)$ to the symplectization $\mathbb{R} \times Y=\mathbb{R} \times Y$ that maps $\partial_{\circ} \Sigma$ to $\mathbb{R} \times \Lambda=\mathbb{R} \times \Lambda$, and that satisfies the non-linear Cauchy-Riemann equations

$$
\begin{equation*}
d u \circ j=\widehat{J} \circ d u \tag{4.3.2}
\end{equation*}
$$

A connected component $C$ of $\partial_{+} \Sigma \cup \partial_{-} \Sigma$ is called a puncture of $u$. A puncture is positive if it is in $\partial_{+} \Sigma$ and negative if it is in $\partial_{-} \Sigma$. Likewise, a closed component is an interior puncture and a component with boundary is a boundary puncture. We require positive punctures under the map $u$ are at $+\infty$ of the symplectization direction, and negative punctures map to $-\infty$ in the symplectization direction.

The energy of a $J$-holomorphic curve $u$ is given by

$$
E(u)=\sup _{\phi}\left(\int_{\Sigma} u^{*} d(\phi(s) \alpha)\right) \in[0, \infty)
$$

Here the supremum is over all non-decreasing smooth maps $\phi: \mathbb{R} \rightarrow[0,1]$.

## Local Maximum Principles

Tailored complex structures satisfy two harmonicity results (which may be viewed as maximum principles) that will be used throughout this paper.

Lemma 127. [28, Lem. 5.6] Let $u: \Sigma \rightarrow \mathbb{R} \times Y$ be a $J$-holomorphic map (possibly with boundary) with respect to a tailored $J$. Then

$$
t \circ u: \Sigma \rightarrow \mathbb{R} \quad \text { defined on } \quad u^{-1}\left(\operatorname{Nbhd}\left(\partial_{+} Y \cup \partial_{-} Y\right)\right)
$$

is harmonic in a neighborhood of $\partial_{+} Y$ and $\partial_{-} Y$.
Proof. We follow the proof of Lemma 5.6 in [28]. Without loss of generality, we focus on a tubular neighborhood of $\partial_{+} Y$ of the form $\partial_{+} Y \times[1+\epsilon, 1-\epsilon$ ) (we imagine extending the $t$ coordinate slightly in the upwards direction). The key observation is that since $\partial_{+} Y$ is two dimensional, $J_{0}$ is Stein, which means there exists $f: \partial_{+} Y \rightarrow \mathbb{R}$ with the boundary of $\partial_{+} Y$ its level set. Further, the 1-form $\beta^{\prime}:=-d f \circ J_{0}$ gives $\partial_{+} Y$ the structure of a Liouville manifold, and the symplectic form $d \beta^{\prime}$ is $J_{0}$ compatible. If we take the contact form $\alpha^{\prime}=d t+\beta^{\prime}$ on $(1-\epsilon,+\infty)_{t} \times \partial_{+} Y$ (which we think of the upwards completion of $Y$, see [28]), then the same proof as Lemma 5.6 in [28] tells us that $t \circ u$ is harmonic in this region.

Lemma 128. [28] Let $J$ be a tailored almost complex structure as above. Restricted to the Liouville manifold

$$
(-\epsilon, 0]_{\tau} \times \Gamma \quad \text { with Liouville form }\left.\quad e^{\tau} \alpha\right|_{\Gamma},
$$

the almost complex structure $J$ is compatible with $\left.\alpha\right|_{\Gamma}$. Then the $\tau$ coordinate $\tau:(-\epsilon, 0]_{\tau} \times \Gamma \rightarrow(-\epsilon, 0]$ is pluri-subharmonic.

Proof. See lemma 5.5 in [28].

## Local Properties Of Boundary Singularities

We will also need a number of local results governing the singularities of holomorphic curves with boundary. To state these results, we fix the following notation. Let $\mathbb{U}$ denote the upper half-disk

$$
\mathbb{U}:=\mathbb{D} \cap \mathbb{H} \quad \text { with } \quad \partial_{\circ} \mathbb{U}:=\mathbb{U} \cap \mathbb{R}
$$

We adopt coordinates $z=s+i t$ on $\mathbb{U} \subset \mathbb{C}$. We also denote the upper half-ball in $\mathbb{C}^{n}$ as follows.

$$
U^{2 n}:=B^{2 n} \cap\left(\mathbb{H} \times \mathbb{C}^{n-1}\right) \quad \text { with } \quad \partial_{0} U^{2 n}=U^{2 n} \cap \mathbb{R} \times \mathbb{C}^{n-1}
$$

Finally, we let $R^{n}:=B^{2 n} \cap \mathbb{R}^{n}$ denote the Lagrangian in $\partial_{\circ} U^{2 n}$ given by the real unit disk.

Lemma 129. Let $J$ be an almost complex structure on $U^{4}$ such that $J\left(T R^{n}\right) \cap T R^{n}=$ 0 and consider a $J$-holomorpic map

$$
u:\left(\mathbb{U}, \partial_{0} \mathbb{U}\right) \rightarrow\left(U^{2 n}, R^{n}\right) \quad \text { with } \quad u(0)=0 \quad \text { and } \quad d u(0) \neq 0
$$

Then there is an open neighborhood $\Omega \subset U^{2 n}$ of 0 and a local diffeomorphism $\phi: \Omega \rightarrow U^{2 n}$ such that

$$
\phi \circ u(z)=(z, 0, \ldots, 0) \quad \text { and } \quad J=\phi^{*} J_{0} \quad \text { on } \quad \mathbb{U} \times 0 \times \cdots \times 0 .
$$

Here $J_{0}$ is the standard almost complex structure on $\mathbb{C}$.
Proof. Choose a complex bundle isomorphism

$$
\tau:\left(\mathbb{C}^{n}, J_{0}\right) \simeq u^{*}\left(\mathbb{C}^{n}, J\right) \quad \text { with } \quad \tau\left(\mathbb{R}^{n}\right)=\mathbb{R}^{n} \text { along } \partial_{\circ} \mathbb{U} .
$$

Also assume that $\tau$ satisfies the following constraint

$$
\tau_{u(z)}(1,0, \ldots, 0)=\frac{\partial u}{\partial s} \quad \text { in coordinates } z=s+i t \text { on } \mathbb{U} .
$$

Then we may define $\phi$ in terms of the exponential map as follows.

$$
\phi^{-1}\left(z, z_{2}, \ldots, z_{n}\right):=\exp _{u(z)}\left(\tau_{u(z)}\left(0, z_{2}, \ldots, z_{n}\right)\right)
$$

Verifying the claimed properties of $\phi$ is standard (see [91, Lem. 2.4.2]).
Next we prove that boundary intersections of (locally) distinct curves are isolated. The analogue for interior intersections is proven in [91, Lem. 2.4.3] and the proof is directly analogous.

Lemma 130. Fix an almost complex structure $J$ on $U^{2 n}$ with $J\left(T R^{n}\right) \cap T R^{n}=0$ and consider a pair of $J$-holomorphic maps

$$
u, v:\left(\mathbb{U}, \partial_{\circ} \mathbb{U}\right) \rightarrow\left(U^{2 n}, R^{n}\right) \quad \text { with } \quad u(0)=v(0) \quad \text { and } \quad d u(0) \neq 0
$$

Suppose that there are sequences $\left\{z_{k}\right\}$ and $\left\{w_{k}\right\} \in \mathbb{U}$ such that

$$
u\left(z_{k}\right)=v\left(w_{k}\right) \quad \lim _{k \rightarrow \infty} z_{k}=\lim _{k \rightarrow \infty} w_{k}=0 \quad \text { and } \quad z_{k} \neq 0
$$

Then there exists a neighborhood $\Omega \subset \mathbb{U}$ of 0 and a map $\phi: \Omega \rightarrow \mathbb{U}$ such that

$$
\phi(0)=0 \quad \text { and } \quad v=u \circ \phi
$$

Proof. By Lemma 129, we can assume that $u(z)=(z, 0, \ldots, 0)$ and that $J=J_{0}$ along $u$. Write

$$
v(z)=\left(v_{1}(z), \widetilde{v}(z)\right) \quad \text { for maps } \quad v_{1}: \mathbb{U} \rightarrow \mathbb{U} \quad \text { and } \quad \widetilde{v}: \mathbb{U} \rightarrow \mathbb{C}^{n-1}
$$

We claim that $\widetilde{v}$ vanishes to infinite order at $z=0$. Indeed, suppose that $\widetilde{v}$ is order $l \geq 1$, i.e. that $\widetilde{v}(z)=O\left(|z|^{l}\right)$ and $\widetilde{v}(z) \neq O\left(|z|^{l+1}\right)$. Since $J=J_{0}$ along $u$, we know that

$$
J(v(z))=J_{0}+O\left(|z|^{l}\right)
$$

and therefore that the order $l$ Taylor series of $v$ is holomorphic, i.e. that

$$
v_{1}(z)=p(z)+O\left(|z|^{l+1}\right) \quad \text { and } \quad \widetilde{v}(z)=\widetilde{a} z^{l}+O\left(|z|^{l+1}\right)
$$

for a polynomial $p$ of order $l$ and a non-zero constant $\widetilde{a}$. In particular, $\widetilde{v}(z) \neq 0$ in some neighborhood of 0 and the intersections of $v$ and $u$ are isolated away from 0 , contradicting the hypotheses. Finally, note that since $J=J_{0}$ along $u$, we have (for all $\left.w=\left(w_{1}, \widetilde{w}\right)\right)$

$$
\frac{\partial J\left(w_{1}, 0\right)}{\partial x_{1}}=\frac{\partial J\left(w_{1}, 0\right)}{\partial y_{1}}=0 \quad \text { and thus } \quad\left|\frac{\partial J(w)}{\partial x_{1}}\right|+\left|\frac{\partial J(w)}{\partial y_{1}}\right| \leq C|\widetilde{w}|
$$

Since $v$ is $J$-holomorphic, this implies that

$$
\begin{equation*}
|\Delta \widetilde{v}| \leq C\left(|\widetilde{v}|+\left|\partial_{s} \widetilde{v}\right|+\left|\partial_{t} \widetilde{v}\right|\right) \tag{4.3.3}
\end{equation*}
$$

and we can apply Aronszajn's theorem (cf. [91, Thm 2.3.4]). In particular we extend $\widetilde{v}$ smoothly past the boundary, then Equation 4.3.3 continues to hold in a small neighborhood of the origin, and the origin is still a zero of infinite order. Then we can apply Aronszajn's theorem in its original form to conclude that $\widetilde{v}=0$ identically. Since $\mathbb{C} \times 0 \times \cdots \times 0 \cap U^{2 n}$ is precisely the (embedded) image of $u$, this proves the result. See also the results in [80].

## Singularities

We now apply the local maximum principles of $\S 4.3$ to study the singularities of holomorphic curves with boundary in the symplectization of $Y$.

For the rest of this sub-section, fix a convex sutured contact 3-manifold $(Y, \xi)$ and a closed (possibly disconnected) Legendrian $\Lambda \subset \partial_{+} Y \cup \partial_{-} Y$. Also fix a nondegenerate, adapted contact form $\alpha$ and a tailored complex structure $J$.

Lemma 131 (Boundary Immersion). Let $u:\left(\Sigma, \partial_{\circ} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ be a finite energy $\widehat{J}$-holomorphic map with boundary. Then

$$
u^{-1}(\mathbb{R} \times \partial Y)=\partial_{0} \Sigma \quad \text { and } \quad u \text { is transverse to } \mathbb{R} \times \partial Y
$$

Proof. Since $u$ has finite energy, it is asymptotic to $\Xi_{ \pm}$at $\pm \infty$, in the sense that $u$ exponentially approaches the trivial cylinders and strips over $\Xi_{ \pm}$at $\pm \infty$. Recall by our notation $\Xi_{ \pm}$denotes a collection of orbits and chords at $\pm \infty$ in the symplectization direction. The specific manner in which they approach the orbits and chords are detailed in [108] and [2]. Since the Reeb chords and orbits are transverse to $\partial_{ \pm} Y$ and disjoint from $\partial_{\sigma} Y$ (we assume implicitly the Legendrians are in the interior of $\partial_{ \pm} Y$ ), our theorem holds on

$$
u^{-1}((-\infty,-T] \cup[T, \infty) \times Y) \subset \Sigma \quad \text { for any sufficiently large } \quad T>0
$$

To show that the result holds elsewhere, consider

$$
A:=u^{-1}\left([-T, T] \times \operatorname{Nbhd}\left(\partial_{\sigma} Y\right)\right) \subset \Sigma
$$

For sufficiently large $T$ and sufficiently small $\operatorname{Nbhd}\left(\partial_{\sigma} Y\right)$, this is a smooth surface, whose boundary is disjoint from $\mathbb{R} \times \partial_{\sigma} Y$. Thus by Lemma $128,\left.u\right|_{A}$ (and thus $u$ ) is disjoint from $\partial_{\sigma} Y$. Finally, consider the compact surface with boundary

$$
B:=u^{-1}\left([-T, T] \times[-1,-1+\epsilon] \times \partial_{-} Y\right) \subset \Sigma
$$

The restriction $\left.t \circ u\right|_{B}$ is harmonic by Lemma 127. Thus, the minimum -1 is achieved only on the boundary of $B$, which is the inverse image of

$$
\{-T,+T\}_{s} \times[-1,-1+\epsilon]_{t} \times \partial_{\circ} Y \cup[-T, T]_{s} \times\{-1,-1+\epsilon\}_{t} \times \partial_{\circ} Y \subset[-T, T] \times Y
$$

For large $T$ the minimum cannot be achieved on

$$
\{-T,+T\} \times(0, \epsilon] \times \partial_{\circ} Y \quad \text { or } \quad[-T, T] \times \epsilon \times \partial_{\circ} Y
$$

Thus $t \circ u$ only takes the value of -1 on $\partial \Sigma \cap B=u^{-1}\left([-T, T] \times\{-1\} \times \partial_{\circ} Y\right)$. A similar argument holds near $\partial_{+} Y$, and this proves that

$$
u^{-1}(\mathbb{R} \times \partial Y)=\partial \Sigma
$$

Finally, we argue that $u$ is an immersion near $\mathbb{R} \times \partial Y$. This is evident away from $[-T, T] \times Y$ by the normal forms of the curve near punctures, as in [108] and [2]. Thus it suffices to show that

$$
t \circ u: u^{-1}\left([-T, T] \times \operatorname{Nbhd}\left(\partial_{-} Y\right)\right) \rightarrow \mathbb{R}
$$

has no critical points along $\partial_{0} \Sigma$ (and similarly for $\partial_{+} Y$ ). Thus, pick any $p \in \partial_{0} \Sigma$ with $u(p) \in \partial_{-} Y$. In a neighborhood of $p$, since $t \circ u$ is harmonic, it is modelled on the real part of a map holomorphic map with a 0 of order $k$, i.e.

$$
z \mapsto z^{k} \quad \text { for } \quad k \geq 1
$$

In particular, the number of connected components of the set

$$
(t \circ u)^{-1}((0, \infty)) \cap \operatorname{Nbhd}(p)
$$

is equal to the order $k$ for small enough $\operatorname{Nbhd}(p)$. Since $t \circ u>0$ in $\operatorname{Nbhd}\left(\partial_{-} Y\right)$, we thus know that $k=1$. This concludes the proof.

Lemma 132 (Boundary Submersion). Let $u:\left(\Sigma, \partial_{0} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ be a finite energy $\widehat{J}$-holomorphic map with boundary and let $\Gamma$ be a connected component of $\partial_{0} \Sigma$. Then the map

$$
\pi_{\mathbb{R}} \circ u: \Gamma \xrightarrow{u} \mathbb{R} \times \Lambda \xrightarrow{\pi_{\mathbb{R}}} \mathbb{R}
$$

is a proper submersion (and therefore a diffeomorphism).
Proof. By Lemma 131, $u$ is transverse to $\mathbb{R} \times \partial Y$. On the otherhand, suppose that

$$
d\left(\left.\pi_{\mathbb{R}} \circ u\right|_{\Gamma}\right)=\left.d \pi_{\mathbb{R}} \circ d u\right|_{\Gamma}=0 \quad \text { at some } p \in \Gamma
$$

Then $d u$ maps $T \Gamma$ to $0 \oplus T Y$ at $p$, since the kernel of $d \pi_{\mathbb{R}}$ is precisely $T Y$. Since $u\left(\partial_{0} \Sigma\right) \subset \mathbb{R} \times \Lambda$, we must then have

$$
d u_{p}(T \Sigma) \subset(0 \oplus T \Lambda) \oplus(0 \oplus J(T \Lambda))=0 \oplus \xi
$$

Moreover, $\xi_{p}=T(\partial Y)_{p}$ at $p$ since $\Lambda$ is an exact Legendrian. Thus $u$ is tangent to $\mathbb{R} \times \partial Y$ at $p$. This contradicts Lemma 131 , so this proves that $\left.\pi_{R} \circ u\right|_{\Gamma}$ is a submersion. Since $\pi_{\mathbb{R}}$ and $u$ are proper, $\left.\pi_{\mathbb{R}} \circ u\right|_{\Gamma}$ is also proper.

As a consequence of the boundary immersion property, we can easily prove that there are a finite number of critical points.

Lemma 133 (Finite Critical Points). Let $u:\left(\Sigma, \partial_{0} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ be a finite energy, non-constant $J$-holomorphic curve. Then the set of critical points $\operatorname{Crit}(u) \subset \Sigma$ is finite.

Proof. By Lemma 131, the set $\operatorname{Crit}(u) \subset \Sigma$ is a closed (and thus compact) subset of the interior of $u^{-1}([-T, T] \times Y)$ for some large $T>0$. By [91, §2.4, p. 26], interior critical points of a non-constant holomorphic curve are isolated. Therefore, Crit(u) is finite.

We next state a key proposition that will be essential to the proof of $d^{2}=0$.
Proposition 134. (No closed boundary components) Let $u:\left(\Sigma, \partial_{\circ} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times$ $\Lambda$ ) be a finite energy $J$-holomorphic map, then no component of $\partial_{o} \Sigma$ is closed unless $u$ is constant.

Proof. Suppose not, let $S^{1}$ denote a closed component of $\partial_{o} \Sigma$ with coordinate $q \in S^{1}$. It is mapped by $u$ to $\mathbb{R} \times L$, where $L$ is an exact Legendrian. Let $\pi_{\mathbb{R}}$ denote the projection to the symplectization direction, and assume $\left.\pi_{\mathbb{R}} u\right|_{S^{1}}$ achieves its maximum at $p$. Then $\frac{d}{d q} \pi_{\mathbb{R}} u(p)=0$. Thus at $p$ the tangent space of $u$ coincides with the contact distribution $\xi$, which at $p$ is tangent to the horizontal boundary. But this implies the curve $u$ is not transverse to the boundary of $\mathbb{R} \times Y$ at $p$, which contradicts the local form in Lemma 131, 129.

## Simple And Somewhere Injective Curves

We can now prove an equivalence between the class of simple curves and somewhere injective curves, generalizing the analogous statement in the case of curves without boundary.

Definition 135. A point $p$ in the domain of a $J$-holomorphic map $u:(\Sigma, j) \rightarrow(X, J)$ is injective if

$$
\begin{equation*}
u^{-1}(u(p))=\{p\} \quad \text { and } \quad d u_{p} \neq 0 \tag{4.3.4}
\end{equation*}
$$

The $J$-holomorphic map $u$ is called somewhere injective if it has an injective point.

Definition 136. A $J$-holomorphic map $u:(\Sigma, j) \rightarrow(X, J)$ is simple if it does not factor as

$$
(\Sigma, j) \xrightarrow{\phi}(S, i) \xrightarrow{v}(X, J)
$$

where $\phi$ is a branched cover of degree 2 or more, and $v$ is $J$-holomorphic.
Lemma 137 (Factorization). Any finite energy, proper $J$-holomorphic map $u$ : $\left(\Sigma, \partial_{0} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ factors as

$$
(\Sigma, j) \xrightarrow{\phi}(S, i) \xrightarrow{v}(X, J)
$$

where $\phi$ is a branched cover and $v$ is a proper $J$-holomorphic map whose injective points form a dense and open set.

Proof. Let $C r$ denote the set of critical values of $u$ and let $B \subset u(\Sigma) \backslash C r$ be the set of non-critical values where multiple branches of $u$ meet. That is, $y \in B$ if and only if there are points $p, q$ such that

$$
y=u(p)=u(q) \quad p \neq q \quad \text { and } \quad u(\operatorname{Nbhd}(p)) \neq u(\operatorname{Nbhd}(q))
$$

for any neighborhoods $\operatorname{Nbhd}(p)$ and $\operatorname{Nbhd}(q)$. The set $C r$ is finite by Lemma 133 and the set $B$ is discrete in $u(\Sigma) \backslash C r$ by [91, Lem. 2.4.3] and Lemma 130. Therefore, $S^{\prime}=u(\Sigma) \backslash(B \cup C r) \subset \mathbb{R} \times Y$ is a Riemann surface with (interior and boundary) punctures, equipped with a tautological map $v: S^{\prime} \rightarrow \mathbb{R} \times Y$. The punctures of $S^{\prime}$ mapping to $B$ and $C$ are removable, and so we can choose extensions of $S^{\prime}$ and $v$

$$
S^{\prime} \subset S \quad \text { and } \quad v: S \rightarrow \mathbb{R} \times Y
$$

so that the resulting map is proper (the extension on the interior works the same was as Proposition 2.5.1 in [91], the extension near the boundary comes from the boundary immersion property in Lemma 131). The map $\Sigma \backslash u^{-1}(B \cup C r) \rightarrow S^{\prime}$ extends across $u^{-1}(B \cup C r)$ to a map $\phi: \Sigma \rightarrow S$ of Riemann surfaces with punctures. The maps $\phi: \Sigma \rightarrow S$ and $v: S \rightarrow \mathbb{R} \times Y$ are precisely the desired maps as we observe that $v$ is simple by construction, and its injective points form an open and dense set.

Remark 138. We emphasize that the results of this lemma do not hold for general holomorphic curves with Lagrangian (or totally real) boundary conditions. Indeed, simple curves with Lagrangian boundary need not have finite self-intersections and
may not be determined by their image. An example given in Remark 2.5.6 in [91]: a disc mapping in $S^{2}$ with boundary on the equator can wrap one and half times around the sphere. Here we are able to obtain nice results because we restricted the Lagrangian to lie at the boundary of our manifold, and $t \circ u$ is harmonic as in Lemma 131. This allows us to conclude that $B$ is finite near the boundary (which is not the case for the partially wrapped disk around $S^{2}$ in the example we just gave), and ensures the factorization Lemma 137 holds.

Lemma 139 (Simple Is Mostly Injective). Fix a non-constant, simple, finite energy $J$-holomorphic map

$$
u:(\Sigma, \partial \Sigma) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)
$$

Then the set of injective points is cofinite.
Proof. Let $S(u) \subset \Sigma \times \Sigma$ denote the set of pairs $(p, q)$ that satisfy $u(p)=u(q)$ and $p \neq q$. Then the set of non-injective points in $\Sigma$ is precisely

$$
\pi_{1}(S(u)) \cup \operatorname{Crit}(u) \quad \text { where } \quad \pi_{1}(p, q)=p
$$

Crit $(u)$ is finite by Lemma 133, and so it suffices to show that $S(u)$ is finite. Note that by Lemma 131, we can decompose $S(u)$ as

$$
S(u)=S_{\partial}(u) \cup S_{\mathrm{int}}(u)
$$

where $S_{\partial}(u)$ consists of pairs in $\partial \Sigma \times \partial \Sigma$ and $S_{\text {int }}(u)$ consists of pairs int $(\Sigma) \times \operatorname{int}(\Sigma)$.
Now, note that $u$ is locally injective, i.e. for every $z \in \Sigma$ there exists a neighborhood $U$ of $z$ such that $\left.u\right|_{U}$ is injective. If $z \in \partial \Sigma$, then this follows from the fact that $d u_{z} \neq 0$ by Lemma 131. If $z \in \operatorname{int}(\Sigma)$, see [91, Rmk. E.1.3]. This implies that $S(u)$ is separated from an open neighborhood of the diagonal.

Next, note that $S(u)$ is a discrete set. Indeed, Lemma 130 implies that any pair $(z, w) \in S_{\partial}(u)$ is isolated from other pairs in $S(u)$. Moreover, pairs in $S_{\text {int }}(u)$ are isolated by [91, Thm. E.1.2, Claim (ii)].

Finally, note that the image of pairs of points in $S(u)$ must be contained in a compact region $[-T, T] \times Y$ because near the punctures of $u$, we have asymptotic formulas (see [108] and the Appendix). Since $u$ is proper, this implies that $S(u)$ is bounded in $\Sigma \times \Sigma$.

Thus, $S(u)$ is a closed, bounded, isolated subset of $(\Sigma \times \Sigma) \backslash \Delta$ and is therefore finite. This proves the result.

Corollary 140. Fix a pair of connected, simple, finite energy $J$-holomorphic maps

$$
u:(\Sigma, \partial \Sigma) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda) \quad \text { and } \quad v:\left(\Sigma^{\prime}, \partial \Sigma^{\prime}\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)
$$

Suppose that $u(\Sigma)=v\left(\Sigma^{\prime}\right)$. Then there is a biholomorphism $\varphi: \Sigma \simeq \Sigma^{\prime}$ such that $u=v \circ \varphi$.

Proof. The same proof as Corollary 2.5.4 in [91] follows through.

## Counts Of Singularities And Legendrian Adjunction

We can now associate a count of singularities to the $J$-holomorphic maps of interest. We first must explain the local situation, starting with the interior case.

Proposition 141. (cf. [121, Lemma 2.6]) Let $J$ be an almost complex structure on $B^{4} \subset \mathbb{C}^{2}$ and let $\Omega=\mathbb{D}_{1} \cup \cdots \cup \mathbb{D}_{m}$ where $\mathbb{D}_{i}=\mathbb{D}$ is a copy of the disk with origin $z_{i}$. Consider a $J$-holomorphic map

$$
v:(\Omega, i) \rightarrow\left(B^{4}, J\right) \quad \text { with } \quad v\left(z_{i}\right)=0 \quad \text { and } \quad v(\partial \Omega) \subset \partial B^{4}
$$

Assume that the only non-injective points of $v$ are $z_{i}$. Then there is a unique positive integer

$$
\delta(v)=\delta(v ; 0) \in \mathbb{Z}_{+}
$$

with the following property: there exists a symplectic immersion

$$
\widetilde{v}: \Omega \rightarrow B^{4} \quad \text { with } \quad \widetilde{v}=v \text { on } \operatorname{Nbhd}(\partial \Omega)
$$

with precisely $\delta(v)$ positive, transverse double points .
We will use the following analogue of Proposition 141 for boundary singularities. The proof is significantly simpler than analogous results (e.g. [121, Lemma 2.6]) in the closed case because we may assume that there are no critical points near the boundary.

Proposition 142. Let $J$ be an almost complex structure on $U^{4}=B^{4} \cap \mathbb{H} \times \mathbb{C}$ and let $\Omega=\mathbb{U}_{1} \cup \cdots \cup \mathbb{U}_{m}$ where $\mathbb{U}_{i}=\mathbb{D} \cap \mathbb{H}$ is a copy of the upper half-disk with origin $z_{i}$. Consider a $J$-holomorphic immersion

$$
v:(\Omega, i) \rightarrow\left(U^{4}, J\right) \quad \text { with } \quad v\left(z_{i}\right)=0 \quad v\left(\partial_{\circ} \mathrm{U}_{i}\right) \subset R^{2} \quad \text { and } \quad v(\partial \Omega) \subset \partial U^{4}
$$

Assume that the only non-injective points of $v$ are $z_{i}$. Then there is a unique positive integer

$$
\epsilon(v)=\epsilon(v ; 0) \in \mathbb{Z}_{+}
$$

with the following property: there exists a symplectic immersion

$$
\widetilde{v}: \Omega \rightarrow U^{4} \quad \text { with } \quad \widetilde{v}=v \text { on } \operatorname{Nbhd}\left(\partial \Omega \backslash \partial_{0} \Omega\right)
$$

with precisely $\epsilon(v)$ positive, transverse boundary double points.
Proof. We first consider the simple case where $\left.v\right|_{\partial \mathbb{U}_{i}}$ and $\left.v\right|_{\partial \mathbb{U}_{j}}$ intersect transversely on the Lagrangian boundary $\mathbb{R} \times \mathbb{R} \subset \mathbb{C} \times \mathbb{C}$. Then the intersection viewed as an intersection between two surfaces in $\mathbb{C} \times \mathbb{C}$ is positive, since it is modelled on the intersection of two complex planes in $\mathbb{C}^{2}$.

In the next case consider a collection of $\left.v\right|_{\partial \mathbb{U}_{i}}, i=1, \cdots k$ intersecting at a single point $p \in \mathbb{R}^{2}$, but such that pairwise their intersections at $p$ are transverse. Then by perturbing the $\left.v\right|_{\mathbb{U}_{i}}$ slightly we recover the conclusion of the theorem.

The most difficult case is when $\left.v\right|_{\partial \mathbb{U}_{i}}$ and $\left.v\right|_{\partial \mathbb{U}_{j}}$ are tangent to each other at $p \in \mathbb{R}^{2}$. We turn to this situation next.

Let $U^{4}(\epsilon)$ be the half-ball of radius $\epsilon$. Note that the image of $v$ is disjoint from $0 \times \mathbb{C}$ except at 0 . Therefore, by Lemma 130, we may acquire a (transverse) braid given by

$$
L=\partial B^{4}(\epsilon) \cap v(\Omega) \quad \text { with components } \quad L_{i}=\partial B^{4}(\epsilon) \cap v\left(\mathbb{U}_{i}\right)
$$

To prove this, fix $i \neq j$ and consider the maps $\left.v\right|_{\mathbb{U}_{i}}$ and $\left.v\right|_{\mathbb{U}_{j}}$. By Lemma 129, we can (after a change of coordinates) assume that

$$
\left.v\right|_{\mathbb{U}_{i}}(z)=(z, 0) \quad \text { and }\left.\quad J\right|_{\mathbb{U} \times 0}=J_{0}
$$

In these coordinates, we may write $v$ restricted to $\mathbb{U}_{j}$ as follows.

$$
\left.v\right|_{\mathbb{U}_{j}}(z)=\left(z, a \cdot z^{l}\right)+O\left(|z|^{l+1}\right)
$$

Write the second factor of $\left.v\right|_{\mathbb{U}_{j}}$ as $w(z)=a \cdot z^{l}+O\left(|z|^{l+1}\right)$. After an orientation check, we find that the linking number of $\left.v\right|_{\mathbb{U}_{i}}$ and $\left.v\right|_{\mathbb{U}_{j}}$ is now precisely the (half-integer) winding number of the path

$$
[0,1] \rightarrow \mathbb{C} \backslash 0 \quad \text { given by } \quad \theta \mapsto w\left(e^{\pi i \theta}\right)
$$

Since $w(z)=a \cdot z^{l}+O\left(|z|^{l+1}\right)$, this winding number is simply $l / 2>0$. This concludes the proof.

The most important application of Proposition 142 for our purposes is the following resolution of singularities result. It states that we can replace any somewhere injective $J$-curve in our setting with a well-behaved surface in the same surface class.

Lemma 143. Fix a non-constant, simple, finite energy $J$-holomorphic map

$$
u:(\Sigma, \partial \Sigma) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)
$$

Then there exists an symplectic, well-immersed surface

$$
\iota:\left(S, \partial_{\circ} S\right) \rightarrow([0,1] \times Y,[0,1] \times \Lambda)
$$

representing the surface class $A$ corresponding to $u$ such that
(a) $S$ has $\delta(C)$ transverse, positive double points in the interior $\operatorname{int}(S)$.
(b) $S$ has $\epsilon(C)$ transverse, positive double points in the boundary $\partial_{\circ} \Sigma$.
(c) The domain $S$ is the same as the domain of $u$.
(d) $\iota$ agrees with $C$ in a neighborhood of $\partial_{ \pm} \Sigma$.

Proof. By Lemma 139, we know that $u$ has a finite number of isolated non-injective points $\operatorname{Sing}(u) \subset \Sigma$. To acquire a map $\iota: S \rightarrow \mathbb{R} \times Y$ satisfying $(c)$ and $(d)$, we simply modify $u$ in a neighborhood the points $\operatorname{Sing}(u)$. To ensure that this map satisfies (a) and (b), we perform this modification by applying Proposition 141 at the interior singularities and Proposition 142 to boundary singularities.

As a corollary, we can prove the Legendrian adjunction result stated in the introduction.

Corollary 144 (Theorem 8). Let $C$ be an finite energy, simple holomorphic curve in $\mathbb{R} \times Y$ with boundary on $\mathbb{R} \times \Lambda$. Then

$$
\mu(C, \tau)=2\left(\bar{\chi}(C)+Q_{\tau}(C)+w_{\tau}(C)-2 \delta(C)-\epsilon(C)\right)
$$

Proof. By Lemma We 143, we can find a symplectic, well-immersed surface $S$ representing the same surface class of $C$ that has the same domain, writhe and count of singularities, so this follows from topological adjunction (Theorem 117).

### 4.4 J-Holomorphic Currents With Boundary

We can now formally introduce $J$-holomorphic currents and describe their basic properties.

Definition 145. A J-holomorphic current with boundary $C$ in an almost complex manifold with boundary $(X, J)$ is a finite set of pairs

$$
\left(C_{i}, m_{i}\right) \quad \text { for } \quad i=1, \ldots, k
$$

where $C_{i}$ is a distinct, simple, somewhere injective $J$-holomorphic curves $C \subset X$ with boundary on $\partial X$ and $m_{i}$ is a positive integer called the multiplicity of $C_{i}$.

We will utilize a specific topology on the space of currents, introduced by Taubes [52, §2.4].

Definition 146. A sequence of $J$-holomorphic currents $\mathcal{C}^{\nu}=\left\{\left(C_{i}^{\nu}, m_{i}^{\nu}\right)\right\}$ for $\nu \in \mathbb{N}$ in an almost complex manifold $(X, J)$ converges to an $J$-holomorphic current $\mathcal{C}=$ $\left\{\left(C_{i}, m_{i}\right)\right\}$ if

- $\mathcal{C}^{\nu}$ converges to $\mathcal{C}$ as a point set. More precisely, the point sets $S^{\nu}=\cup_{i} C_{i}^{\nu}$ converge to the set $S=\cup_{i} C_{i}$ in the Hausdorff metric, on each compact subset of $X$.
- $\mathcal{C}^{\nu}$ converges to $\mathcal{C}$ as a current, i.e. for every compactly supported 2-form $\sigma$ on $X$, we have

$$
\lim _{\nu \rightarrow \infty} \sum_{i}\left(m_{i}^{\nu} \cdot \int_{C_{i}^{\nu}} \sigma\right) \rightarrow \sum_{i}\left(m_{i} \cdot \int_{C_{i}} \sigma\right)
$$

A key property of holomorphic currents, which does not hold for curves, is a general compactness property given a uniform action bound.

Theorem 147. (Taubes, cf. [52]) Let $(X, \omega)$ be a symplectic manifold with boundary and compatible almost complex structure $J$. Let $C_{i}$ be a sequence of $J$-holomorphic currents with boundary on $\partial X$ such that

$$
\int_{C_{i}} \omega<C \quad \text { for all } i \text { and some } C>0
$$

Then there is a subsequence that converges to a J-holomorphic current $\mathcal{C}$.

## Currents With Boundary In Symplectizations

We are interested in currents in the symplectization of a pair $(Y, \Lambda)$ of a convex sutured contact 3-manifold $Y$ and a pair of closed Legendrians $\Lambda_{ \pm} \subset \partial_{ \pm} Y$.

In this setting, $J$-holomorphic maps, curves and currents admit a natural $\mathbb{R}$ action, denoted by

$$
\mathcal{C} \mapsto \mathcal{C}+s \quad \text { for any } \quad s \in \mathbb{R}
$$

On curves, this action is given by composing the map $u: \Sigma \rightarrow \mathbb{R} \times Y$ with translation $\mathbb{R} \times Y \rightarrow \mathbb{R} \times Y$ by $s$. Moreover, the translation invariant $J$-holomorphic currents coincide with cylinders over sets of orbits and chords, as follows.

Example 148. The trivial current $\mathbb{R} \times \Xi$ over an orbit-chord set $\Xi$ consisting of simple orbits and chords $\gamma_{i}$ of multiplicity $m_{i}$ is the current of pairs

$$
\left(\mathbb{R} \times \gamma_{i}, m_{i}\right)
$$

Every tame enough current in $\mathbb{R} \times Y$ is asymptotic to an orbit-chord sets near infinity, as follows.

Lemma 149. Let $\mathcal{C}$ be a $J$-holomorphic current in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ such that each curve $C_{i}$ in $\mathcal{C}$ is proper and finite energy. Then there are orbit sets $\Xi_{+}$and $\Xi_{-}$such that

$$
\begin{equation*}
\mathcal{C}+s_{i} \xrightarrow{\text { Def } 146} \mathbb{R} \times \Xi_{ \pm} \quad \text { for any sequence } s_{i} \in \mathbb{R} \text { with } s_{i} \rightarrow \pm \infty \tag{4.4.1}
\end{equation*}
$$

Definition 150. A $J$-holomorphic current $\mathcal{C}$ from $\Xi_{+}$to $\Xi_{-}$is a current satisfying (4.4.1). We denote the space of currents from $\Xi_{+}$to $\Xi_{-}$by

$$
\mathcal{M}\left(\Xi_{+}, \Xi_{-}\right)
$$

We denote the quotient by the $\mathbb{R}$-action on currents by $\mathcal{M}\left(\Xi_{+}, \Xi_{-}\right) / \mathbb{R}$. More generally, a broken J-holomorphic current $\overline{\mathcal{C}}$ from $\Xi_{+}$to $\Xi_{-}$is a sequence

$$
\mathcal{C}_{i} \in \mathcal{M}\left(\Xi_{i}, \Xi_{i+1}\right) / \mathbb{R} \quad \text { for } i=1, \ldots, m \quad \text { where } \quad \Xi_{1}=\Xi_{+} \text {and } \Xi_{m+1}=\Xi_{-}
$$

The space of broken $J$-holomorphic currents from $\Xi_{+}$to $\Xi_{-}$is denoted by $\overline{\mathcal{M}}\left(\Xi_{+}, \Xi_{-}\right)$.
Remark 151. Any $J$-holomorphic current in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ appearing in the rest of this paper will be assumed to be finite energy and proper, so that

$$
\mathcal{C} \in \mathcal{M}(\Xi, \Theta) \quad \text { for some orbit-chord sets } \Xi \text { and } \Theta
$$

There is an tautological map sending a finite energy, proper $J$-holomorphic curve $C$ from orbit-chord set $\xi$ to orbit-chord set $\Theta$ to a corresponding surface class [ $C]$. This extends to a map

$$
\begin{equation*}
\overline{\mathcal{M}}(\Xi, \Theta) \rightarrow S(\Xi, \Theta) \quad \text { denoted by } \quad \overline{\mathcal{C}} \mapsto[\overline{\mathcal{C}}] \tag{4.4.2}
\end{equation*}
$$

that is compatible with union and composition in the sense that

$$
\begin{array}{ll}
{[\mathcal{C} \cup \mathcal{D}]=[\mathcal{C}] \cup[\mathcal{D}]} & \text { for any pair of currents } \mathcal{C} \text { and } \mathcal{D} \\
{[\overline{\mathcal{C}} \circ \overline{\mathcal{D}}]=[\overline{\mathcal{C}}] \circ[\overline{\mathcal{D}}]} & \text { for any pair of broken currents } \overline{\mathcal{C}} \text { and } \overline{\mathcal{D}}
\end{array}
$$

## Gromov Compactness

We next introduce the notion of Gromov compactness of broken currents. This compactness result resembles SFT compactness of Bourgeouis-Eliashberg-Hofer-WyosockiZehnder, but does not depend on it.

Definition 152. A sequence of (equivalence classes of) $J$-holomorphic currents

$$
\mathcal{C}^{\nu} \in \mathcal{M}\left(\Xi_{+}, \Xi_{-}\right) / \mathbb{R} \quad \text { for } \quad \nu \in \mathbb{N}
$$

converges in the Gromov topology to a broken $J$-holomorphic current $\overline{\mathcal{C}}=\left(\mathcal{C}_{1}, \ldots, \mathcal{C}_{m}\right)$ if there are representatives of $\mathcal{C}^{\nu}$ and $\mathcal{C}_{i}$, denoted respectively by

$$
\mathcal{S}^{\nu} \in \mathcal{M}\left(\Xi_{+}, \Xi_{-}\right) \quad \text { and } \quad \mathcal{S}_{i} \in \mathcal{M}\left(\Xi_{i}, \Xi_{i+1}\right)
$$

and a set of sequences $s_{i}^{\nu} \in \mathbb{R}$ for each $i=1, \ldots, m$ such that

$$
\mathcal{S}^{\nu}+s_{i}^{\nu} \xrightarrow{\text { Def } 146} \mathcal{S}_{i} \quad \text { as } \nu \rightarrow \infty
$$

The using the same argument as [52], we have the following compactness result for currents.

Proposition 153 (Gromov Compactness). Let $\mathcal{C}^{\nu} \in \mathcal{M}\left(\Xi_{+}, \Xi_{-}\right)$be a sequence of currents. Then, after passing to a subsequence, there is a broken $J$-holomorphic current $\overline{\mathcal{C}}$ from $\Xi_{+}$to $\Xi_{-}$such that

$$
\mathcal{C}^{\nu} \rightarrow \overline{\mathcal{C}} \quad \text { and } \quad[\overline{\mathcal{C}}]=\left[\mathcal{C}^{\nu}\right]
$$

## Intersection Numbers And Linking Of Currents

It will be convenient to extend the intersection number and linking number to currents under a certain disjointness hypothesis.

Definition 154. A pair of $J$-holomorphic currents $\mathcal{C}$ and $\mathcal{D}$ have distinct components if the component curves $\left(C_{i}, m_{i}\right)$ and $\left(D_{j}, n_{j}\right)$ satisfy

$$
C_{i} \neq D_{j} \quad \text { for all } i, j
$$

Currents with disjoint components have a well-defined geometric intersection number that counts actual intersections (with multiplicity).

Definition 155. The geometric intersection number of $J$-holomorphic currents $\mathcal{C}$ and $\mathcal{D}$ with distinct components is the non-negative half-integer given by

$$
\mathcal{C} \cdot \mathcal{D} \quad \text { is given by } \quad \mathcal{C} \cdot \mathcal{D}:=\sum_{i, j} m_{i} \cdot n_{j} \cdot\left(C_{i} \cdot C_{j}\right)
$$

where the geometric intersection $C \cdot D$ of two distinct, somewhere injective $J$ holomorphic curves $C$ and $D$ is defined as the following count of singularities.

$$
C \cdot D:=(\epsilon(C \cup D)-\epsilon(C)-\epsilon(D))+\frac{1}{2}(\delta(C \cup D)-\delta(C)-\delta(D))
$$

Remark 156. It is possible to generalize the geometric intersection number to a self-intersection number of curves and currents (see [55]). However, we will not carry this out in this paper.

The finiteness (and thus well-definedness) of the geometric intersection number follows from the analysis of singularities of $J$-holomorphic curves in our setting (see Section 4.3).

Definition 157. The linking number of two currents $\mathcal{C}$ and $\mathcal{D}$ with disjoint components with respect to a trivialization $\tau$, denoted by

$$
l_{\tau}(\mathcal{C}, \mathcal{D})
$$

is defined as follows. Let $\left(C_{i}, m_{i}\right)$ and $\left(D_{j}, n_{j}\right)$ denote the component curves of $\mathcal{C}$ and $\mathcal{D}$. Then

$$
l_{\tau}(\mathcal{C}, \mathcal{D})=\sum_{i, j} m_{i} \cdot n_{j} \cdot l_{\tau}\left(C_{i}, D_{j}\right)
$$

Here we adopt the following conventions for the linking number of curves.

- If $C$ is non-trivial curve with an end on $\eta$ and $D=\mathbb{R} \times \eta$ is a trivial curve, then

$$
l_{\tau}(C, D)=\operatorname{wind}_{\tau}\left(\zeta_{+}\right)-\operatorname{wind}_{\tau}\left(\zeta_{-}\right)
$$

where $\zeta_{ \pm}$re the asymptotic braids of $C$ at $\pm \infty$.

- If $C$ and $D$ are both non-trivial, then

$$
l_{\tau}(C, D)=l_{\tau}\left(\zeta_{+}, \xi_{+}\right)-l_{\tau}\left(\zeta_{-}, \xi_{-}\right)
$$

where $\zeta_{ \pm}$and $\xi_{ \pm}$are the asymtptotic braids of $C$ and $D$, respectively.
Note that the asymptotic braids can be empty. In this case, the winding number and linking number with any other braid are zero by convention.

This extended linking number transforms in precisely the same way as its topological counterpart under change of trivialization.

Lemma 158. Let $\mathcal{C} \in \mathcal{M}\left(\Xi_{+}, \Xi_{-}\right)$and $\mathcal{D} \in \mathcal{M}\left(\Theta_{+}, \Theta_{-}\right)$be $J$-holomorphic currents with distinct components.

$$
\eta \text { of multiplicity } m \text { in } \Xi_{ \pm} \text {and } n \text { in } \Theta
$$

Then the linking number of $\mathcal{C}$ and $\mathcal{D}$ differ as follows.

$$
l_{\tau}(\mathcal{C}, \mathcal{D})-l_{\sigma}(\mathcal{C}, \mathcal{D})=m \cdot n \cdot(\sigma-\tau)
$$

The proof is straightforward so we omit it. The linking number is also related to the geometric intersection number as follows.

Lemma 159. Let $\mathcal{C}$ and $\mathcal{D}$ be $J$-holomorphic currents with distinct components. Then

$$
\mathcal{C} \cdot \mathcal{D}=Q_{\tau}(\mathcal{C}, \mathcal{D})+l_{\tau}(\mathcal{C}, \mathcal{D})
$$

Proof. By bilinearity with respect to union, we assume that $\mathcal{C}$ and $\mathcal{D}$ are somewhere injective (and even connected) curves, and that $C$ is non-trivial. Then $C \cup D$ is somewhere injective and we can use Lemma 143 to replace $C$ and $D$ with wellimmersed symplectic surfaces $S$ and $T$ in the same surface class that satisfy

$$
C \cdot D=\#(\operatorname{int}(S) \cap \operatorname{int}(T))+\frac{1}{2} \cdot \#\left(\partial_{\circ} S \cap \partial_{\circ} T\right) \quad \text { and } \quad l_{\tau}(C, D)=l_{\tau}(S, T)
$$

Note that we are using the fact that, when $D$ is a trivial cylinder over an orbit or chord

$$
l_{\tau}(C, D)=\operatorname{wind}_{\tau}\left(\zeta_{+}\right)-\operatorname{wind}_{\tau}\left(\zeta_{-}\right)
$$

The result follows from Definition 114.

### 4.5 The Conley-Zehnder And Fredholm Indices

In this section, we review several versions of the Conley-Zehnder index. In particular, we introduce the Conley-Zehnder index of a Lagrangian path.

## Paths Of Lagrangians and the Conley-Zehnder Index

Recall that given a path of symplectic matrices, $\Phi_{t}, t \in[0,1]$ satisfying $\Phi_{0}=I d$ and $\Phi_{1}$ does not have 1 as an eigenvalue by nondegeneracy, as in [122] we can associate an integer valued Conley-Zehnder index to this path of symplectic matrices, which we call the Conley-Zehnder index. We refer the reader to lecture 3 of [122] for properties of this index.

We next provide a definition for the Conley-Zehnder index of a path of Lagrangian sub-spaces. We adopt the perspective that this is naturally a half-integer.

Definition 160. A path $L:[0,1] \rightarrow \operatorname{LGr}(2 n)$ is non-degenerate if the ends are transverse.

$$
L_{1} \pitchfork L_{0}
$$

Definition 161. The Conley-Zehnder index of a non-degenerate path of Lagrangians $L:[0,1] \rightarrow \operatorname{LGr}(2 n)$ is the half-integer

$$
\mathrm{CZ}(L):=\mu(\bar{L})+\frac{n}{2} \in \frac{1}{2} \mathbb{Z}
$$

Here $\bar{L}: S^{1} \rightarrow \operatorname{LGr}(2 n)$ is the loop of Lagrangians constructed as follows. Choose a complex structure $J$ on $\mathbb{C}^{n}$ such that

$$
J \text { is compatible with } \omega_{0} \quad \text { and } \quad J L_{1}=L_{0}
$$

Then $\bar{L}$ is the loop acquired by joining $L$ to the path $\exp \left(-J \cdot \frac{\pi t}{2}\right)$ for $t \in[0,1]$ from $L_{1}$ to $L_{0}$.

Remark 162. By abuse of terminology and notation, we will refer to $\mu(\bar{L})$ as the Maslov index of the path $L$ and denote it by $\mu(L)$.

We can also define the Conley-Zehnder index of a Lagrangian path as half of the ordinary Conley-Zehnder index of an associated path of symplectic matrices. See [43] for an exposition of how various flavours of Conley-Zehnder indices are related to each other.

## Bundles Over 1-Manifolds

Next, we introduce the notion of an asymptotic operators and discuss the ConleyZehnder index of an asymptotic operator.

Let $S$ be an oriented 1-manifold with boundary. Let $(E, F) \rightarrow(S, \partial S)$ be a symplectic bundle pair and fix a symplectic connection

$$
\nabla: C^{\infty}(S, \partial S ; E, F) \rightarrow C^{\infty}\left(S ; \Omega^{1} S \otimes E\right)
$$

Definition 163. The connection $\nabla$ is non-degenerate if the kernel of $\nabla$ is trivial.
Given a choice of trivialization $\tau:(E, F) \simeq\left(\mathbb{C}^{n}, \mathbb{R}^{n}\right)$, we can assign a ConleyZehnder index to the above data. To explain how, let $\eta \subset S$ be a component of $S$ and choose an oriented identification

$$
\phi: \mathbb{R} / \mathbb{Z} \simeq \eta \quad \text { or } \quad \phi:[0,1] \simeq \eta
$$

In either case, parallel transport via the connection $\nabla$ determines a family of symplectic maps

$$
\mathrm{PT}_{t}: E_{\phi(0)} \rightarrow E_{\phi(t)} \quad \text { for each } \quad t \in[0,1]
$$

Given the choice of trivialization $\tau$, this then determines a 1-parameter family of matrices

$$
\Phi^{\tau}:[0,1] \rightarrow \operatorname{Sp}(2 n) \quad \text { with } \quad \Phi_{t}^{\tau}:=\tau_{\phi(t)} \circ \mathrm{PT}_{t} \circ \tau_{\phi(0)}^{-1}
$$

and a family of Lagrangians $L^{\tau}:=\Phi^{\tau}\left(\mathbb{R}^{n}\right)$.
Definition 164. The Conley-Zehnder index $\mathrm{CZ}(\eta, \tau)$ of a component $\eta \subset S$ with respect to a trivialization $\tau$ of $(E, F)$ is

$$
\mathrm{CZ}\left(\Phi^{\tau}\right) \text { if } \eta \simeq \mathbb{R} / \mathbb{Z} \quad \text { and } \quad \mathrm{CZ}\left(L^{\tau}\right) \text { if } \eta \simeq[0,1]
$$

The Conley-Zehnder index $\mathrm{CZ}(S, \tau)$ of $(E, F, \nabla)$ with respect to $\tau$ is simply the sum

$$
\mathrm{CZ}(S, \tau):=\sum_{\eta} \mathrm{CZ}(\eta, \tau)
$$

Let $S$ be a compact 1-manifold diffeomorphic to the interval $[0,1]$ and let $(E, F) \rightarrow$ $(S, \partial S)$ be a symplectic bundle pair. From this point onward assume $E$ is twodimensional. Given a nowhere vanishing vector field $R$ on $S$, a compatible complex structure $J$ on $E$ and a symplectic connection $\nabla$, we define the associated asymptotic operator to be

$$
A: C^{\infty}(S, \partial S ; E, F) \rightarrow C^{\infty}(S, E) \quad A X:=-J \nabla_{R} X
$$

Lemma 165. The asymptotic operator $A$ satisfies the following basic properties:
(i) $A$ has trivial kernel if and only if $\nabla$ is non-degenerate.
(ii) All eigenvalues of $A$ are simple.
(iii) All eigenfunctions of $A$ are nowhere vanishing.

Proof. The first assertion is immediate because the operators $A$ and $\nabla$ have the same kernel. Let $\lambda$ be an eigenvalue of $A$. Since $A-\lambda$ is a linear first order ordinary differential operator, the initial value homomorphism

$$
\operatorname{ker}(A-\lambda) \rightarrow F_{p} \quad X \mapsto X(p)
$$

is injective. Here $p \in \partial S$ is a boundary point of $S$. Thus the eigenspace of $\lambda$ has dimension 1. This proves the second assertion. Let $X \in \operatorname{ker}(A-\lambda) \backslash\{0\}$ be an eigenvector. Again because $A-\lambda$ is a linear first order ordinary differential operator, $X$ vanishes at one single point $q \in S$ if and only if it is identically equal to zero. The third assertion follows.

In particular, given an asymptotic operator $A$, a symplectic trivialization $\tau$ of $(E, F)$ and an eigenvalue $\lambda$ of $A$, there is a well-defined half-integer-valued winding number $w(\tau, A ; \lambda) \in \frac{1}{2} \mathbb{Z}$ associated to the eigenfunction of $A$ with eigenvalue $\lambda$. To be specific, under our choice of trivialization $\tau$, if $e(t):[0,1] \rightarrow \mathbb{C}$ is such an eigenfunction that maps the end points $\{0,1\}$ of $[0,1]$ to $\mathbb{R} \subset \mathbb{C}$, using the fact that $e(t)$ is never vanishing, the winding number is $\frac{1}{2}$ times the number of half turns $e(t)$ makes about the origin.

Lemma 166. The function $w(\tau, A ; \cdot): \operatorname{Spec} A \rightarrow \frac{1}{2} \mathbb{Z}$ is monotonic and a bijection.
Proof. We solve this problem by deforming to a simplified operator $A_{0}$ with the same Conley-Zehnder index (see below for description). Recall the asymptotic operator $A$ is of the form $-J_{0} \partial_{t}-S(t): C^{\infty}\left([0,1],\{0,1\} ; \mathbb{R}^{2}, \mathbb{R} \times 0\right) \mapsto C^{\infty}\left([0,1], \mathbb{R}^{2}\right)$. Its Conley-Zehnder index is computed by considering the path of symplectic matrices $\Phi(t)$ satisfying

- $\Phi(0)=I d$
- $-J_{0} \partial_{t} \Phi(t)-S(t) \Phi(t)=0$
and considering the loop of Lagrangians $\Phi_{t}(\mathbb{R}) \subset \mathbb{C}$. In particular, for general operators of this form, the path of Lagrangians $\Phi_{t}(\mathbb{R})$ being nondegenerate implies that $A$ has trivial kernel.

Now suppose we have an asymptotic operator $A_{0}$ satisfying the conclusion of the lemma (we will specify what this operator is directly in a bit) such that $A$ and $A_{0}$ have the same Conley-Zehnder index as described above, then we can choose a path $\left(A_{s}\right)_{s \in[0,1]}$ of nondegenerate asymptotic operators connecting $A_{0}$ to $A_{1}=A$. Let $\lambda_{0}^{k / 2}$ for $k \in \mathbb{Z}$ denote the eigenvalues of $A_{0}$ labelled by the winding number $k / 2$. Since the spectrum of $A_{s}$ is simple for all $s \in[0,1]$, we may continuously extend $\lambda_{0}^{k / 2}$ to a family $\lambda_{s}^{k / 2}$ such that the spectrum of $A_{s}$ is given by $\lambda_{s}^{k / 2}$ for $k \in \mathbb{Z}$. By assumption, the function $k \mapsto \lambda_{0}^{k / 2}$ is monotonic. Thus the same must continue to hold for $s>0$. We may choose families $X_{s}^{k / 2}$ of associated eigenvectors. Since $X_{s}^{k / 2}$ is nowhere vanishing for all $s$, the winding number of $X_{1}^{k / 2}$ agrees with the winding number of $X_{0}^{k / 2}$, which is equal to $k / 2$. Thus the winding number of $\lambda_{1}^{k / 2}$ is actually equal to $k / 2$. This proves that the lemma holds for $A=A_{1}$.

Let us prove the lemma for the following explicit asymptotic operator $A_{0}$.

$$
A_{0}: C^{\infty}\left([0,1],\{0,1\} ; \mathbb{R}^{2}, \mathbb{R} \times 0\right) \mapsto C^{\infty}\left([0,1], \mathbb{R}^{2}\right) \quad X \mapsto-J_{0} \partial_{t} X-\pi / 2(1+2 l) X
$$

This can be solved explicitly as follows. Consider $(x(t), y(t)) \in \mathbb{R}^{2}$, this satisfies the equation

$$
-\left(\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right)\binom{x^{\prime}}{y^{\prime}}-\pi / 2(1+2 l)\binom{x}{y}=\lambda\binom{x}{y}
$$

This is the system of equations

$$
y^{\prime}=(\lambda+\pi / 2(1+2 l)) x, \quad x^{\prime}=-(\lambda+\pi / 2(1+2 l)) y
$$

differentiating twice to get

$$
y^{\prime \prime}=-(\lambda+\pi / 2(1+2 l))^{2} y, \quad x^{\prime \prime}=-(\lambda+\pi / 2(1+2 l))^{2} x
$$

with boundary conditions:

$$
y(0)=y(1)=0
$$

Hence the eigenvalues of the asymptotic operator are $\lambda=\pi / 2(2 n+1)(n \in \mathbb{Z})$, and the corresponding eigenvectors are

$$
y_{n}(t)=a \sin ((\lambda+\pi / 2(2 l+1)) t), \quad x_{n}(t)=a \cos ((\lambda+\pi / 2(2 l+1)) t)
$$

For each choice of eigenvalue $\lambda$ of $A_{0}$, we may compute the winding number $w$ from these explicit eigenvectors and see that the theorem is satisfied.

In addition, with this choice of $A_{0}$ we also compute the Conley-Zehnder index associated to this choice of trivialization. We consider the matrix satisfying

$$
\partial_{t} \Phi=\pi / 2(1+2 l) J \Phi
$$

Then identifying $\mathbb{R}^{2}$ with $\mathbb{C}$, the matrix $\Phi$ can be identified with $e^{i \pi / 2(1+2 l) t}$. We consider the path of Lagrangians $\Phi(t)(\mathbb{R})$ and append the path $e^{-i \frac{\pi}{2} t}(i \mathbb{R})$ to make this a loop of Lagrangians. We see this loop has Maslov index $l$, and hence the Conley-Zehnder index of the operator is $l+\frac{1}{2}$. The fact that $\left\{\left.l+\frac{1}{2} \right\rvert\, l \in \mathbb{Z}\right\}$ enumerates all possible Conley-Zehnder indices that an asymptotic operator $A$ can have, proves that the path $\left(A_{s}\right)_{s \in[0,1]}$ we chose exists.

Lemma 167. Suppose that $\nabla$ is non-degenerate. We have
$\mathrm{CZ}(S, \tau)=\min \left\{w(\tau, A ; \lambda) \mid \lambda \in \operatorname{Spec}(A) \cap \mathbb{R}_{>0}\right\}+\max \left\{w(\tau, A ; \lambda) \mid \lambda \in \operatorname{Spec}(A) \cap \mathbb{R}_{<0}\right\}$
Proof. We assume the Conley-Zehnder index of $A$ is $l+\frac{1}{2}(l \in \mathbb{Z})$. Same as what we did in the proof of Lemma 166, we begin by choosing a path of non-degenerate asymptotic operators $\left(A_{s}\right)_{s \in[0,1]}$ such that $A_{1}=A$ and $A_{0}=-J_{0} \partial_{t}-\pi / 2(1+2 l)$. As in the proof of Lemma 166, let $\lambda_{s}^{k / 2}$ denote the unique eigenvalue of $A_{s}$ with winding number $k / 2$. Since for every $s$, the asymptotic operator $A_{s}$ is non-degenerate, we observe that 0 is never an eigenvalue of $A_{s}$, and hence the paths of eigenvalues $\left(\lambda_{s}^{k / 2}\right)_{s \in[0,1]}$ do not cross 0 . It follows that

$$
\min \left\{w\left(\tau, A_{s} ; \lambda\right) \mid \lambda \in \operatorname{Spec}\left(A_{s}\right) \cap \mathbb{R}_{>0}\right\}+\max \left\{w\left(\tau, A_{s} ; \lambda\right) \mid \lambda \in \operatorname{Spec}\left(A_{s}\right) \cap \mathbb{R}_{<0}\right\}
$$

is constant for $s \in[0,1]$, and hence it suffices to verify the claim for $A_{0}$.
By our explicit calculation in the proof of Lemma 166, the smallest positive eigenvalue of $A_{0}$ is $\frac{\pi}{2}$ with winding number $\frac{l+1}{2}$, and the largest negative eigenvalue of $A_{0}$ is $-\frac{\pi}{2}$ with winding number $\frac{l}{2}$, so the claim holds true for $A_{0}$.

## Fredholm Index

We the previous section on Conley-Zehnder indices we are now prepared to discuss the Fredholm index for pseudo-holomorphic curves in our setting.

Theorem 168. Let $Y$ be a sutured contact manifold of dimension $2 n-1$. Let $u:\left(\Sigma, \partial_{0} \Sigma\right) \rightarrow(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ be a pseudoholomorphic curve. Suppose that $u$ is positively asymptotic to orbits and chords $\gamma_{1}^{+}, \ldots, \gamma_{k}^{+}$and negatively asymptotic to orbits and chords $\gamma_{1}^{-}, \ldots, \gamma_{\ell}^{-}$. Then the Fredholm index of $u$ is given by

$$
\operatorname{ind}(u)=(n-3) \bar{\chi}(\Sigma)+\mu(u, \tau)+\sum_{i=1}^{k} \mathrm{CZ}_{\tau}\left(\gamma_{i}^{+}\right)-\sum_{i=1}^{\ell} \mathrm{CZ}_{\tau}\left(\gamma_{i}^{-}\right)
$$

The way to interpret the above formula is as follows: if $u$ is somewhere injective, and $J$ is generic, then $u$ lives in a moduli space that has dimension given by the index formula.

Let $\Sigma$ be a compact connected Riemann surface, possibly with boundary. Let $\Gamma \subset \Sigma$ be the finite set of punctures (previously we denoted this by $\partial_{o} \Sigma$ ), possibly on the boundary. Let $\Gamma=\Gamma_{i} \cup \Gamma_{b}$ be the partition of $\Gamma$ into interior and boundary punctures. Moreover, let $\Gamma=\Gamma^{+} \cup \Gamma^{-}$be a partition of $\Gamma$ into positive and negative punctures. Let $\dot{\Sigma}=\Sigma \backslash \Gamma$ denote the punctured surface. Let $(E, F) \rightarrow \dot{\Sigma}$ be a bundle pair and let $\tau$ denote a trivialization near the punctures.

Lemma 169. Let

$$
D: W^{k, p}(E, F) \rightarrow W^{k-1, p}\left(\Lambda^{0,1} T^{*} \dot{\Sigma} \otimes E\right)
$$

be a Cauchy-Riemann type operator with non-degenerate asymptotic operators $A_{z}^{\tau}$ at the punctures $z \in \Gamma$. Then $D$ is Fredholm and its index is given by

$$
\operatorname{ind}(D)=n \cdot \bar{\chi}(\dot{\Sigma})+\mu^{\tau}(E, F)+\sum_{z \in \Gamma^{+}} \mathrm{CZ}\left(A_{z}^{\tau}\right)-\sum_{z \in \Gamma^{-}} \mathrm{CZ}\left(A_{z}^{\tau}\right) .
$$

Proof. The form of the asymptotic operators $A_{z}^{\tau}$ is described in [14], and using our Definition 164 we can associate to it a Conley-Zehnder index. The lemma is exactly the statement of the Fredholm index formula proved in [14]. Note that [14] uses different definitions for the Euler characteristic of a punctured surface with boundary and for the Conley-Zehnder index. The Euler characteristic $X\left(\Sigma, \Gamma^{ \pm}\right)$appearing in [14] is related to the orbifold Euler characteristc $\bar{\chi}(\dot{\Sigma})$ via

$$
X\left(\Sigma, \Gamma^{ \pm}\right)=\bar{\chi}(\dot{\Sigma})+\frac{1}{2} \# \Gamma_{b}^{+}-\frac{1}{2} \# \Gamma_{b}^{-} .
$$

At interior punctures the Conley-Zehnder index $\mu_{\mathrm{CZ}}\left(A_{z}^{\tau}\right)$ showing up in [14] agrees with our Conley-Zehnder index $\mathrm{CZ}\left(A_{z}^{\tau}\right)$. At boundary punctures, the relationship is given by

$$
\mu_{\mathrm{CZ}}\left(A_{z}^{\tau}\right)=\mathrm{CZ}\left(A_{z}^{\tau}\right)-\frac{n}{2}
$$

where $n$ denotes the complex rank of $E$. Plugging these identities into the index formula given in [14] readily yields the lemma.

Lemma 170 (Section 3, [120]). The real dimension of the moduli space of punctured Riemann surfaces with boundary of topological type $\left(\Sigma, \Gamma=\Gamma_{i} \cup \Gamma_{b}\right)$ is given by

$$
-3 \chi(\Sigma)+2 \# \Gamma_{i}+\# \Gamma_{b}=-3 \bar{\chi}(\dot{\Sigma})-\# \Gamma_{i}-\frac{1}{2} \# \Gamma_{b}
$$

Proof of Theorem 168. The moduli space containing $u$ is the zero set of a certain Fredholm section. For a general description of this setup see Section 3.3 of [120] and Section 5 of [34]. The linearization of this Fredholm section at $u$ is a Fredholm operator of the form

$$
D: W^{k, p, \delta}\left(u^{*} T \mathbb{R} \times Y, u^{*} T \mathbb{R} \times \Lambda\right) \oplus V \oplus T \rightarrow W^{k-1, p, \delta}\left(\Lambda^{0,1} T^{*} \dot{\Sigma} \otimes u^{*} T \mathbb{R} \times Y\right)
$$

Here $\delta$ is a small positive exponential weight. The space $W^{k, p, \delta}\left(u^{*} T \mathbb{R} \times Y, u^{*} T \mathbb{R} \times \Lambda\right)$ means we use a weighted Sobolev space with weight of the form $e^{\delta|s|}$ near each puncture of $\dot{\Sigma}$ around which we have already chosen cylindrical coordinates of the form $(s, t) \in[0, \pm \infty) \times S^{1}$ or $[0, \pm \infty) \times[0,1]$ The dimension of the tangent space $V$ of the space of asymptotic markers is given by $2 \# \Gamma_{i}+\# \Gamma_{b}$. By Lemma 170 the dimension of the Teichmüller slice $T$ is given by $-3 \bar{\chi}(\dot{\Sigma})-\# \Gamma_{i}-\frac{1}{2} \# \Gamma_{b}$. The restriction $D^{\prime}$ of $D$ to $W^{k, p, \delta}\left(u^{*} T \mathbb{R} \times Y, u^{*} T \mathbb{R} \times \Lambda\right)$ is a Cauchy-Riemann type operator. Let $z \in \Gamma$ be a puncture and let $\gamma$ denote the corresponding Reeb chord or orbit. Then the asymptotic operator of $D^{\prime}$ at $z$ takes the form $A_{z}^{\tau} \oplus-i \partial_{t}$ where $A_{z}^{\tau}$ acts on sections of $\gamma^{*} \xi$ and $-i \partial_{t}$ acts on sections of the trivial bundle spanned by $\partial_{s}$ and $R$. It follows from Lemma 169 that if the exponential weight $\delta$ is sufficiently small, the index of $D^{\prime}$ is given by
$\operatorname{ind}\left(D^{\prime}\right)=n \cdot \bar{\chi}(\dot{\Sigma})+\mu(u, \tau)+\sum_{z \in \Gamma^{+}} \mathrm{CZ}\left(A_{z}^{\tau} \oplus\left(-i \partial_{t}+\delta\right)\right)-\sum_{z \in \Gamma^{-}} \mathrm{CZ}\left(A_{z}^{\tau} \oplus\left(-i \partial_{t}-\delta\right)\right)$.
The Conley-Zehnder index is additive under direct sums, i.e.

$$
\mathrm{CZ}\left(A_{z}^{\tau} \oplus\left(-i \partial_{t} \pm \delta\right)\right)=\mathrm{CZ}\left(A_{z}^{\tau}\right)+\mathrm{CZ}\left(-i \partial_{t} \pm \delta\right)
$$

A direct computation shows that if we regard $-i \partial_{t} \pm \delta$ as an asymptotic operator over the circle, then

$$
\mathrm{CZ}\left(-i \partial_{t} \pm \delta\right)=\mp 1
$$

for $\delta>0$ sufficiently small. If we regard it as an asymptotic operator over the interval, then

$$
\mathrm{CZ}\left(-i \partial_{t} \pm \delta\right)=\mp 1 / 2
$$

Combining these identities we compute
$\operatorname{ind}(D)=\operatorname{ind}\left(D^{\prime}\right)+\operatorname{dim} V+\operatorname{dim} T=(n-3) \bar{\chi}(\Sigma)+\mu(u, \tau)+\sum_{i=1}^{k} \mathrm{CZ}_{\tau}\left(\gamma_{i}^{+}\right)-\sum_{i=1}^{\ell} \mathrm{CZ}_{\tau}\left(\gamma_{i}^{-}\right)$

## Regular Almost Complex Structures

We conclude this section by reviewing the notion of regular almost complex structures.

Definition 171. A tailored almost complex structure $J$ on $(Y, \Lambda)$ is regular if every somewhere injective, finite energy $J$-holomorphic curve in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ is transversely cut out.

A standard argument shows that regular, tailored almost complex structures are generic. In particular, we have the following proposition.

Proposition 172. The set $\mathcal{J}_{\text {reg }}(Y, \Lambda)$ of regular, tailored almost complex structures on $(Y, \Lambda)$ is comeager in the space of all tailored almost complex structures on $(Y, \Lambda)$.

The proof follows the analogous arguments in the case of $J$-holomorphic curves without boundary in symplectizations (cf. [122, §7, 8] and more specifically [122, Thm 7.2]).

Proposition 173. If $J$ is a regular, tailored almost complex structure on $(Y, \Lambda)$, then the set of simple, finite energy, $J$-holomorphic curves with boundary is a smooth manifold and the dimension at a curve $u$ is given by ind $(u)$.

This also follows from analogues of the arguments in $[122, \S 7,8]$ for somewhere injective curves in symplectizations, and the fact that simple implies somewhere injective in our setting (Lemma 139).

Remark 174. In general, for a curve $u$ with Lagrangian boundary condition, there is no guarantee that $u$ can be factored into a branched cover and a somewhere injective (and therefore regular) curve. It is key that we can use Lemma 139 to replace
somewhere injective with simple in the statement of Proposition 173. It implies that the underlying $J$-holomorphic curves in any $J$-holomorphic current are transversely cut out when $J$ is regular.

### 4.6 Legendrian ECH Index And Its Properties

In this section, we introduce the Legendrian ECH index and establish its basic properties.

## ECH Conley-Zehnder Term

We begin by introducing the Conley-Zehnder term that appears in the Legendrian ECH index.

Definition 175. The ECH Conley-Zehnder term of an orbit-chord set $\Xi$ with respect to a trivialization $\tau$ along $\Xi$ is the unique half-integer

$$
\mathrm{CZ}_{\tau}^{E C H}(\Xi) \in \frac{1}{2} \mathbb{Z}
$$

that is additive with respect to disjoint union, in that

$$
\mathrm{CZ}_{\tau}^{E C H}(\Xi \sqcup \Theta)=\mathrm{CZ}_{\tau}^{E C H}(\Xi)+\mathrm{CZ}_{\tau}^{E C H}(\Theta)
$$

and that is given by the following formula for a simple Reeb orbit $\gamma$ and Reeb chord c
$\mathrm{CZ}_{\tau}^{E C H}((\gamma, m))=\sum_{i=1}^{m} \mathrm{CZ}_{\tau}\left(\gamma^{i}\right) \quad$ and $\quad \mathrm{CZ}_{\tau}^{E C H}((c, m))=\frac{m}{2}+\frac{m(m+1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(c)-\frac{1}{2}\right)$
The ECH Conley-Zehnder term of a surface class $A \in S\left(\Xi_{+}, \Xi_{-}\right)$is defined to be

$$
\mathrm{CZ}_{\tau}^{E C H}(A):=\mathrm{CZ}_{\tau}^{E C H}\left(\Xi_{+}\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\Xi_{-}\right)
$$

We adopt the same definition for $J$-holomorphic curves and currents.
We will make heavy use of the following change of trivialization rule.
Lemma 176 (Trivialization). Let $\sigma, \tau$ be two trivializations of $\xi$ along an simple orbit or chord $\eta$. Then

$$
\begin{equation*}
\mathrm{CZ}_{\tau}^{E C H}((\eta, m))-\mathrm{CZ}_{\sigma}^{E C H}((\eta, m))=\frac{m(m+1)}{2} \cdot(\sigma-\tau) \in \frac{1}{2} \mathbb{Z} \tag{4.6.1}
\end{equation*}
$$

Proof. Recall that the Conley-Zehnder index itself transforms as

$$
\mathrm{CZ}_{\tau}(\eta)-\mathrm{CZ}_{\sigma}(\eta)=2(\sigma-\tau) \in \mathbb{Z} \quad \text { by Lemma } 176
$$

Note that this formula holds in the Reeb orbit and chord case (and in the former case, the difference $\sigma-\tau$ is a whole integer). In the case of a Reeb orbit, we thus have

$$
\mathrm{CZ}_{\tau}^{E C H}((\eta, m))-\mathrm{CZ}_{\sigma}^{E C H}((\eta, m))=\sum_{i=1}^{m} i \cdot(\sigma-\tau)=\frac{m(m+1)}{2} \cdot(\sigma-\tau)
$$

In the case of a chord, the difference formula follows directly from the definition of $\mathrm{CZ}{ }_{\tau}^{E C H}$.

## Partition Conditions

We next briefly review the ECH partition conditions. Let $P(m)$ denote the set of partitions of an integer $m \in \mathbb{N}$.

Definition 177. The positive partition $p_{\gamma}^{+}(m)$ and negative partition $p_{\gamma}^{-}(m)$ associated to a non-degenerate Reeb orbit $\gamma$ in $(Y, \alpha)$ are partitions in $P(m)$ defined as follows.

Fix a trivialization of $\xi$ along $\gamma$ and let $L$ be the period of $\gamma$. Consider the linearized flow

$$
\phi_{\tau}=\left.\tau_{\gamma(0)} \circ d \phi_{L}(\gamma(0))\right|_{\xi} \circ \tau_{\gamma(0)}^{-1} \in \operatorname{Sp}(2)
$$

- If $\gamma$ is positive hyperbolic (i.e. $\phi_{\tau}$ has positive real eigenvalues) then

$$
p_{\gamma}^{+}(m)=p_{\gamma}^{-}(m)=(1,1, \ldots, 1)
$$

- If $\gamma$ is negative hyperbolic (i.e. $\phi_{\tau}$ has negative real eigenvalues) then

$$
p_{\gamma}^{+}(m)=p_{\gamma}^{-}(m)=\left\{\begin{array}{cl}
(2,2, \ldots, 2) & \text { if } m \text { is even } \\
(2,2, \ldots, 2,1) & \text { if } m \text { is odd }
\end{array}\right.
$$

- If $\gamma$ is elliptic (i.e. $\phi_{\tau}$ has unit complex eigenvalues) then $\phi_{\tau}$ is conjugate to a rotation by angle $2 \pi \theta$ for $\theta \in \mathbb{R} / \mathbb{Z}$. Then

$$
p_{\gamma}^{+}(m)=p_{\theta}^{+}(m) \quad \text { and } \quad p_{\gamma}^{-}(m)=p_{\theta}^{-}(m)
$$

Here $p_{\theta}^{ \pm}(m)$ are partitions defined using only the rotation angle $\theta$ as follows. Let $\Lambda_{\theta}^{+}(m) \subset \mathbb{R}^{2}$ be the maximal concave polygonal graph with vertices at lattice points in $\mathbb{Z}^{2}$ that starts on $(0,0)$, ends on $(m,\lfloor m \theta\rfloor)$ and lies below the line $y=\theta x$ where $\theta \in(0,1)$. Then $p_{\theta}^{+}(m)$ is the sequence of horizontal displacements of the consecutive vertices of $\Lambda_{\theta}^{+}(m)$. We can define $p_{\theta}^{-}(m)=p_{-\theta}^{+}(m)$.

Definition 178. A somewhere injective $J$-holomorphic curve $C$ in $(Y, \Lambda)$ from orbit set $\Xi_{+}$to orbit set $\Xi_{-}$satisfies the ECH partition conditions if

- For each orbit $\gamma$ of multiplicity $m \geq 1$ in $\Gamma_{+}$, the partition of $\gamma$ determined by the positive end of $C$ at $\gamma$ is $p_{+}(m, \gamma)$.
- For each orbit $\gamma$ of multiplicity $m \geq 1$ in $\Gamma_{-}$, the partition of $\gamma$ determined by the negative end of $C$ at $\gamma$ is $p_{-}(m, \gamma)$.

Given a simple, elliptic Reeb orbit $\gamma$ of $Y$ and two partitions $p, q \in P(m)$, we write

$$
p \prec_{\gamma} q
$$

if there is a degree $m J$-holomorphic branched cover

$$
u: \Sigma \rightarrow \mathbb{R} \times \gamma \subset \mathbb{R} \times Y
$$

with Fredholm index 0 and whose ends yield the partition $q$ on the positive end and $p$ at the negative end.

Lemma 179. [52, Exercise 3.1 and A.1.4] The relation $\prec_{\gamma}$ is a partial order on $P(m)$. Moreover, $p_{\gamma}^{+}(m)$ and $p_{\gamma}^{-}(m)$ are respectively minimal and maximal with respect to this partial order.

A helpful consequence of Lemma 179 is the following lemma controlling the triviality of trivial cylinders over closed Reeb orbits.

Lemma 180. Let $\gamma$ be a simple closed Reeb orbit in $Y$. Fix a sequence of partitions $p_{1}, \ldots, p_{n}$ of $m$ and let $Z_{i}$ be a branched cover of $\mathbb{R} \times \gamma$ with positive partition $p_{i}$ and negative partition $p_{i+1}$. Assume that

$$
\begin{equation*}
p_{1}=p_{n}=p_{\gamma}^{+}(m) \quad \text { or } \quad p_{1}=p_{n}=p_{\gamma}^{-}(m) \tag{4.6.2}
\end{equation*}
$$

Then $Z_{i}$ is a trivial cover and $p_{i}=p_{\gamma}^{ \pm}(m)$ for each $i$.

Proof. We may view $\mathbf{Z}=\left(Z_{1}, \ldots, Z_{n}\right)$ as an SFT building (cf. [1]) with Fredholm index

$$
\operatorname{ind}(\mathbf{Z})=\sum_{i} \operatorname{ind}\left(Z_{i}\right)=0
$$

By [60, Lem. 1.7(a)], we know that $\operatorname{ind}\left(Z_{i}\right) \geq 0$ for each $i$. Therefore $\operatorname{ind}\left(Z_{i}\right)=0$ for each $i$. Moreover, [60, Lem. 1.7(b)] states that the cover is unbranched if $\gamma$ is hyperbolic. If $\gamma$ is elliptic, then

$$
p_{\gamma}^{ \pm}(m) \prec_{\gamma} p_{i} \prec_{\gamma} p_{\gamma}^{ \pm}(m)
$$

for each $i$. It follows from the fact that $\prec_{\gamma}$ is a partial order that $p_{i}=p_{\gamma}^{ \pm}(m)$ for each $i$. The Fredholm index is then given by

$$
\operatorname{ind}\left(Z_{i}\right)=-\chi\left(\Sigma_{i}\right)=0
$$

where $\Sigma_{i}$ is the domain of the cover $Z_{i}$. Riemann-Hurwitz then implies that $Z_{i}$ is unbranched.

There is also an analogous lemma for chords, which we include here for completeness.

Lemma 181. Let $c$ be a Reeb chord in $(Y, \Lambda)$ and let $Z$ be an $m$-fold branched cover of $\mathbb{R} \times c$. Then $\operatorname{ind}(Z) \geq 0$ with equality if and only if $Z$ is unbranched.

Proof. Let $\Sigma$ be the domain of $Z$. We may double $\mathbb{R} \times c$ to a cylinder $\mathbb{R} \times S^{1}$ and double $\Sigma$ to a surface $S$. Then the curve $Z$ is an equivalence class of branched cover

$$
u: S \rightarrow \mathbb{R} \times S^{1}
$$

By choosing a trivialization of $\xi$ along $c$ and extending it along $Z$, we find that

$$
\operatorname{ind}(Z)=-\bar{\chi}(\Sigma)=\frac{-\chi(S)}{2}
$$

Note that $S$ has at least 2 punctures, and thus has at non-negative Euler characteristic. Thus $\operatorname{ind}(Z)=0$. If $\operatorname{ind}(Z)=-\frac{\chi(S)}{2}=0$, then Riemann-Hurwitz implies that the cover $u$ is unbranched. This implies that $Z$ is unbranched.

## Writhe And Linking Bounds

We are now ready to prove the Legendrian version of the writhe bound, and a separate linking bound that will be used to prove the ECH index inequality.

Lemma 182 (Writhe-Linking Bound At Chord). Let $C$ be a somewhere injective, $J$-holomorphic curve with respect to a tailored $J$ on $(Y, \Lambda)$, asymptotic to the orbitchord sets $\Xi_{ \pm}$at $\pm \infty$. Then

- For any chord $\eta$ in the orbit-chord set $\Xi_{+}$, the asymptotic braid $\zeta$ of $C$ in Nbhd $(\eta)$ satisfies

$$
w_{\sigma}(\zeta) \leq 0 \quad \text { in the unique trivialization } \sigma \text { with } \mathrm{CZ}_{\sigma}(\eta)=\frac{1}{2}
$$

and the component braids $\zeta_{i}$ of $\zeta$ satisfy

$$
l_{\sigma}\left(\zeta_{i}, \zeta_{j}\right) \leq 0 \quad \text { and } \quad \operatorname{wind}_{\sigma}\left(\zeta_{i}\right) \leq 0
$$

- For any chord $\eta$ in the orbit-chord set $\Xi_{-}$, the asymptotic braid $\zeta$ of $C$ in $\operatorname{Nbhd}(\eta)$ satisfies

$$
w_{\sigma}(\zeta) \geq 0 \quad \text { in the unique trivialization } \sigma \text { with } \mathrm{CZ}_{\sigma}(\eta)=-\frac{1}{2}
$$

and the component braids $\zeta_{i}$ of $\zeta$ satisfy

$$
l_{\sigma}\left(\zeta_{i}, \zeta_{j}\right) \geq 0 \quad \text { and } \quad \operatorname{wind}_{\sigma}\left(\zeta_{i}\right) \geq 0
$$

Moreover, equality occurs if and only if $\zeta$ is the trivial braid in the trivialization $\sigma$.
Proof. We prove the result for a chord in $\Xi_{+}$. The proof is entirely analogous in the other case.

Consider a strip-like end of the holomorphic curve $C$ which is positively asymptotic to the chord $\eta$. For $s_{0} \gg 0$ sufficiently large, the intersection of $C$ with the half infinite cylinder $\left[s_{0}, \infty\right) \times Y$ is given by the graph of a function

$$
\left[s_{0}, \infty\right) \times[0,1] \rightarrow\left[s_{0}, \infty\right) \times Y \quad \text { given by } \quad(s, t) \mapsto\left(s, \exp _{\eta(t)} U(s, t)\right)
$$

Here exp denotes the exponential map of an auxiliary Riemannian metric on $Y$ with the property that the Legendrian $\Lambda$ consists of closed geodesics. Moreover, $U$ is a family of sections

$$
U:\left(\left[s_{0}, \infty\right) \times[0,1],[0, \infty) \times\{0,1\}\right) \rightarrow \eta^{*}(\xi, T \Lambda)
$$

of the bundle pair $\eta^{*}(\xi, T \Lambda)$. The following result of Abbas states that $U$ decays exponentially as $s \rightarrow \infty$.

Theorem 183. ([2]) There exists an eigenvector $e$ of the asymptotic operator of $\eta$ with negative eigenvalue $\lambda$ and a family $r(s, t)$ of sections of $\eta^{*} \xi$ exponentially decaying as s tends to $\infty$ such that

$$
U(s, t)=e^{\lambda s}(e(t)+r(s, t))
$$

We require two ingredients from other parts of this paper for this proof. First, we have the following refinement of Theorem 183. Consider two strip-like ends of $C$ positively asymptotic to $\eta$. Let $U$ and $V$ denote the associated families of sections of $\eta^{*} \xi$.

Theorem 184. Assume that $U$ and $V$ do not agree identically. Then there exist an eigenvector $e$ of the asymptotic operator of $\eta$ with negative eigenvalue $\lambda$ and a family $r(s, t)$ of sections of $\eta^{*} \xi$ exponentially decaying as $s$ tends to $\infty$ such that

$$
U(s, t)-V(s, t)=e^{\lambda s}(e(t)+r(s, t))
$$

The proof of Theorem 184 requires a rather long analytical digression and is deferred to the Appendix. Second, we require the following result on the winding number of $e$ in Theorem 184.

Lemma 185. Let $\sigma$ be a representative of the unique homotopy class of trivializations of $\eta^{*} \xi$ such that

$$
\mathrm{CZ}_{\sigma}(\eta)=\frac{1}{2}
$$

Then the winding number of $e$ with respect to $\sigma$ is strictly positive if $\lambda>0$ and non-positive if $\lambda<0$.

This is an immediate corollary of Lemma 166 and Lemma 167 in Section 4.5.
We now proceed with the proof of Lemma 182. Let $\zeta_{i}$ denote the components of the braid $\zeta$. By Theorem 184 and Lemma 185, the winding numbers and pairwise linking numbers satisfy

$$
\operatorname{wind}_{\sigma}\left(\zeta_{i}\right) \leq 0 \quad \text { and } \quad l_{\sigma}\left(\zeta_{i}, \zeta_{j}\right) \leq 0 \quad \text { for all } i \neq j
$$

Hence it follows from (4.2.6) that the writhe $w_{\sigma}(\zeta)$ is non-positive, proving the inequality.

Finally, we claim that the braid $\zeta$ must be trivial with respect to $\sigma$ if $w_{\sigma}(\zeta)=0$. The asymptotic operator of $\eta$ has a 1-dimensional eigenspace whose elements have zero winding number with respect to $\sigma$. Let $e$ be a non-zero vector in this eigenspace, $\lambda<0$ be the associated eigenvalue and $U_{i}(s, t)$ be the section of $\eta^{*} \xi$ associated to the braid $\zeta_{i}$. Since $l_{\sigma}\left(\zeta_{i}, \zeta_{j}\right)$ is non-positive for all $i \neq j$ and the writhe $w_{\sigma}(\zeta)$ vanishes by assumption, we can deduce that $l_{\sigma}\left(\zeta_{i}, \zeta_{j}\right)=0$ for all $i \neq j$. Thus, for every tuple $i \neq j$, there exist a non-zero real number $a_{i j}$ and an exponentially decaying family of sections $r_{i j}(s, t)$ such that

$$
\begin{equation*}
U_{i}(s, t)-U_{j}(s, t)=e^{\lambda s}\left(a_{i j} e(t)+r_{i j}(s, t)\right) \tag{4.6.3}
\end{equation*}
$$

We may replace the trivialization $\sigma$ by a homotopic one with the property that $e(t)$ is constant with respect to the new trivialization. Then it follows from (4.6.3) that the braid $\zeta$ is trivial.

We can now prove Proposition 186. It largely follows from the local version for chords.

Proposition 186 (Writhe Bound). Let $C$ be a somewhere injective $J$-holomorphic curve in $(Y, \Lambda)$ asymptotic to $\Gamma_{ \pm}=\left(\gamma_{i}^{ \pm}\right)$at $\pm \infty$. Then

$$
w_{\tau}(C) \leq \mathrm{CZ}_{\tau}^{E C H}(C)-\mathrm{CZ}_{\tau}\left(\Gamma_{+}\right)+\mathrm{CZ}_{\tau}\left(\Gamma_{-}\right)
$$

Moreover, there is equality only if following conditions are satisfied.

- The orbit parts of $\Gamma_{+}$and $\Gamma_{-}$satisfy the ECH partition conditions.
- The chords in $\Gamma_{-}$are multiplicity one.

Proof. The writhe bound is additive under disjoint union of simple orbits and chords. Therefore, it suffices to consider each orbit and chord in $\Xi_{ \pm}$independently. For Reeb orbits, the result is proven in [52, §5.1]. For Reeb chords, we may write

$$
\mathrm{CZ}_{\tau}^{E C H}(c, m)-m \cdot \mathrm{CZ}_{\tau}(\eta)=\frac{m(m-1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(c)-\frac{1}{2}\right)
$$

Thus, it is sufficicient to prove to the following pair of statements.

- If $\eta$ is a Reeb chord that appears in $\Xi_{+}$with multiplicity $m$ and $\zeta$ is the corresponding asymptotic braid of $C$, then

$$
\begin{equation*}
w_{\tau}(\zeta) \leq \frac{m(m-1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(\eta)-\frac{1}{2}\right) \tag{4.6.4}
\end{equation*}
$$

- If $\eta$ is a Reeb chord that appears in $\Xi_{-}$with multiplicity $n$ and $\zeta$ is the corresponding asymptotic braid of $C$, then

$$
\begin{equation*}
w_{\tau}(\zeta) \geq \frac{m(m-1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(\eta)-\frac{1}{2}\right) \tag{4.6.5}
\end{equation*}
$$

Moreover, equality holds if only if $m=1$.
Also, observe that the above inequalities are equivalent for different choices of the trivialization $\tau$ on $c$. If $\sigma$ and $\tau$ are two trivializations of $\xi$ along $c$ with difference $\tau-\sigma \in \frac{1}{2} \mathbb{Z}$, then

$$
\begin{gather*}
w_{\tau}(\zeta)-w_{\sigma}(\zeta)=m(m-1) \cdot(\sigma-\tau) \quad \text { by Lemma } 106  \tag{4.6.6}\\
\left.C Z_{\tau}(\eta)-C Z_{\sigma}(\eta)=2 \cdot(\sigma-\tau)\right) \quad \text { by Lemma } 176 \tag{4.6.7}
\end{gather*}
$$

So it suffices to prove the inequalities for any choice of trivialization.
For the first claim, we choose the trivialization $\tau$ so that $\mathrm{CZ}_{\tau}(\eta)=\frac{1}{2}$. Then Proposition 182 implies that

$$
\begin{equation*}
w_{\tau}(\zeta) \leq 0=\frac{m(m-1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(\eta)-\frac{1}{2}\right) \tag{4.6.8}
\end{equation*}
$$

For the second claim, we choose $\tau$ so that $\mathrm{CZ}_{\tau}(\eta)=-\frac{1}{2}$. Proposition 182 implies that

$$
\begin{equation*}
w_{\tau}(\zeta) \geq 0 \geq-\frac{m(m-1)}{2}=\frac{m(m-1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(\eta)-\frac{1}{2}\right) \tag{4.6.9}
\end{equation*}
$$

with equality holds only if the multiplicity $n=1$.
Remark 187 (Writhe Asymmetry). The observant reader will notice that there is an asymmetry in the chord condition required for the writhe bound to yield an equality. This is an artifact of our convention for the Legendrian ECH index. We could have alternatively defined it as

$$
-\frac{m}{2}+\frac{m(m+1)}{2} \cdot\left(\mathrm{CZ}_{\tau}(c)+\frac{1}{2}\right)
$$

Then the partition conditions imposed on the positive chords and negative chords would have been reversed. An asymmetry of this type is more or less unavoidable.

There is also an analogous statement to the writhe inequality about the linking number of holomorphic currents.

Proposition 188. Let $\mathcal{C}$ and $\mathcal{D}$ be a pair of $J$-holomorphic currents with boundary that have disjoint components. Then

$$
\mathrm{CZ}_{\tau}^{E C H}(\mathcal{C} \cup \mathcal{D}) \geq \mathrm{CZ}_{\tau}^{E C H}(\mathcal{C})+\mathrm{CZ}_{\tau}^{E C H}(\mathcal{D})+2 \cdot l_{\tau}(\mathcal{C}, \mathcal{D})
$$

Proof. Let $\left(C_{i}, m_{i}\right)$ and $\left(D_{j}, n_{j}\right)$ denote the components of $\mathcal{C}$ and $\mathcal{D}$. Fix an orbit or chord $\eta$ and let

$$
\zeta_{i}^{ \pm} \quad \text { and } \quad \xi_{j}^{ \pm}
$$

denote the braids of $C_{i}$ and $D_{j}$ asymptotic to $\eta$ at the positive and negative ends respectively. We fix the shorthand notation

$$
l_{\tau}\left(\zeta^{+}, \xi^{+}\right)=\sum_{i, j} m_{i} \cdot n_{j} \cdot l_{\tau}\left(\zeta_{i}^{+}, \xi_{j}^{+}\right) \quad \text { and } \quad l_{\tau}\left(\zeta^{-}, \xi^{-}\right)=\sum_{i, j} m_{i} \cdot n_{j} \cdot l_{\tau}\left(\zeta_{i}^{-}, \xi_{j}^{-}\right)
$$

Finally, let $m_{ \pm}$and $n_{ \pm}$denote the multiplicity of $\eta$ at the positive and negative ends of $\mathcal{C}$ and $\mathcal{D}$, respectively. It suffices to prove that

$$
\begin{align*}
& \mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, m_{+}+n_{+}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, m_{+}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, n_{+}\right)\right)-2 \cdot l_{\tau}\left(\zeta^{+}, \xi^{+}\right) \geq 0 \\
& \mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, m_{-}+n_{-}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, m_{-}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, n_{-}\right)\right)-2 \cdot l_{\tau}\left(\zeta^{-}, \xi^{-}\right) \leq 0 \tag{4.6.10}
\end{align*}
$$

for any choice of $\eta$. This is proven by Hutchings [55, Eq. 5.6 on p. 47] if $\eta$ is a Reeb orbit. Thus, we assume that $\eta$ is a Reeb chord. Finally, as in the proof of Proposition 186, Lemma 158 and Lemma 176 imply that the claim is independent of trivialization.

Next choose the trivialization $\tau$ so that $\mathrm{CZ}_{\tau}(\eta)=\frac{1}{2}$. Then the ECH ConleyZehnder term is

$$
\mathrm{CZ}_{\tau}^{E C H}((\eta, k))=0 \quad \text { for any integer } k
$$

Moreover, by Lemma 182, the linking numbers of $\zeta_{i}^{+}$and $\xi_{j}^{+}$satisfies

$$
l_{\tau}\left(\zeta_{i}^{+}, \xi_{j}^{+}\right) \leq 0 \quad \text { for all } i, j
$$

Note that the linking number is the winding number, by convention, if $C_{i}$ or $D_{j}$ is a trivial cylinder over $\eta$. These two inequalities imply (4.6.10). Finally, choose the trivialization $\sigma$ so that $\mathrm{CZ}_{\tau}(\eta)=-\frac{1}{2}$. Then

$$
\mathrm{CZ}_{\tau}^{E C H}((\eta, k))=-\frac{k(k-1)}{2} \quad \text { for any integer } k
$$

In particular, the ECH Conley-Zehnder term in (4.6.11) satisfies

$$
\mathrm{CZ} Z_{\tau}^{E C H}\left(\left(\eta, m_{-}+n_{-}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, m_{-}\right)\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\left(\eta, n_{-}\right)\right)<0
$$

Moreover, by Lemma 182, the linking numbers of $\zeta_{i}^{-}$and $\xi_{j}^{-}$satisfies

$$
l_{\tau}\left(\zeta_{i}^{-}, \xi_{j}^{-}\right) \geq 0 \quad \text { for all } i, j
$$

This proves (4.6.11) and completes the proof.

## ECH Index And Basic Properties

We are now ready to introduce the Legendrian ECH index in detail, and prove its basic properties.

Definition 189. The Legendrian ECH index $I(A)$ of a surface class $A \in S(\Theta, \Xi)$ between orbit-chord sets $\Theta$ and $\Xi$ is given by the following formula.

$$
\begin{equation*}
I(A):=Q_{\tau}(A)+\frac{1}{2} \mu_{\tau}(A)+\mathrm{CZ}_{\tau}^{E C H}(\Theta)-\mathrm{CZ}_{\tau}^{E C H}(\Xi) \tag{4.6.12}
\end{equation*}
$$

The terms in the formula above are as follows.

- $\tau$ is any trivialization of the bundle pair $(\xi, T \Lambda)$ over $\Theta$ and $\Xi$ (see Definition 92)
- $\mu_{\tau}$ is the (relative) Maslov number (see Definition 98)
- $Q_{\tau}$ is the relative self intersection number with respect to $\tau$ (see Definition 114)
- $\mathrm{CZ}_{\tau}^{E C H}$ is the ECH Conley-Zehnder term (see Definition 175).

The Legendrian ECH index of a $J$-holomorphic current $\mathcal{C} \in \mathcal{M}(\Theta, \Xi)$ is simply the index of its corresponding surface class $A \in S(\Theta, \Xi)$.

$$
I(\mathcal{C}):=I(A)
$$

Lemma 190. The Legendrian ECH index $I(A)$ of a surface class $A: \Theta \rightarrow \Xi$ is an integer that is independent of the trivialization in (4.6.12).

Proof. Since $I(A)$ is a sum of half-integer valued terms, it is a half integer. To show that $I(A)$ is an integer, fix a well-immersed surface $S$ representing $A$. Let $m(S)$ denote the number of open boundary components of $\partial_{+} S \cup \partial_{-} S$, and note that

$$
\bar{\chi}(S)=\frac{1}{2} \cdot m(S) \quad \bmod 1
$$

By topological adjunction (Theorem 117), we have

$$
\frac{1}{2} \mu_{\tau}(S)+Q_{\tau}(S)=\bar{\chi}(S)=\frac{1}{2} \cdot m(S) \quad \bmod 1
$$

On the other hand, the ECH Conley-Zehnder term is a sum of the (integer) CZindices of closed orbits and the CZ-indices of chords. The Conley-Zehnder index of a chord in dimension three is automatically a strict half-integer. Since the total number of chords in $\Theta$ and $\Xi$ (with multiplicity) is $m(S)$, we find that

$$
\mathrm{CZ}_{\tau}^{E C H}(\Xi)-\mathrm{CZ}_{\tau}^{E C H}(\Theta)=\frac{1}{2} \cdot m(S) \bmod 1
$$

Thus we find that

$$
I(A)=\frac{1}{2} \mu_{\tau}(S)+Q_{\tau}(S)+\mathrm{CZ}_{\tau}^{E C H}(\Xi)-\mathrm{CZ}_{\tau}^{E C H}(\Theta)=m(S)=0 \quad \bmod 1
$$

To see that $I(A)$ is independent of trivialization, fix two trivializations $\sigma$ and $\tau$ of $\xi$ along $\Xi_{ \pm}$. We write the orbit sets $\Theta$ and $\Xi$ as follows.

$$
\Xi=\left\{\left(\gamma_{i}, m_{i}\right)\right\} \quad \text { and } \quad \Theta=\left\{\left(\eta_{j}, n_{j}\right)\right\}
$$

Fix an orbit or chord $\eta \subset \Xi$ with multiplicity $m$. Let $\sigma$ and $\tau$ be two trivializations of $\xi$ along $\Xi$ and $\Theta$ that agree everywhere except along $\eta$. We compute that

$$
\begin{gathered}
Q_{\tau}(A)-Q_{\sigma}(A)=m^{2} \cdot(\tau-\sigma) \quad \text { by Lemma } 115 \\
\mu_{\tau}(A)-\mu_{\sigma}(B)=\sigma-\tau \quad \text { by Proposition } 99 \\
\mathrm{CZ}_{\tau}^{E C H}(\Xi)-\mathrm{CZ}_{\sigma}^{E C H}(\Xi)=m(m-1) \cdot(\sigma-\tau) \quad \text { by Lemma } 176
\end{gathered}
$$

The sum of these terms is zero, thus yielding the desired invariance property. An identical proof works if $\tau$ and $\sigma$ only differ along an orbit or chord of $\Theta$. Any two trivializations are related by a sequence of changes supported on one orbit or chord, so this proves the result.

The Legendrian ECH index has the following basic properties, generalizing the corresponding properties of the ordinary ECH index (cf. [50, Proposition 1.6]).

Proposition 191 (Basic Properties). The Legendrian ECH index satisfies the following axioms.

- (Index Ambiguity) If $A, B: \Xi_{+} \rightarrow \Xi_{-}$are two classes with difference $B-A \in$ $H_{2}(Y, \Lambda)$, and let $\Gamma \in H_{1}(Y, \Lambda)$ denote the homology class of $\Xi_{ \pm}$. Then

$$
I(B)-I(A)=\frac{1}{2}\langle\mu(\xi, \Lambda), B-A\rangle+2 \cdot Q(\Gamma, B-A)
$$

where $\mu(\xi, \Lambda)$ is the Maslov class of $(\xi, \Lambda)$ and $Q$ is the intersection pairing of $(Y, \Lambda)$.

- (Composition) If $A: \Xi_{0} \rightarrow \Xi_{1}$ and $B: \Xi_{1} \rightarrow \Xi_{2}$ are composible surface classe then

$$
I(A \circ B)=I(A)+I(B)
$$

Proof. We prove each of these properties separately.
Index Ambiguity. Let $A, B: \Theta \rightarrow \Xi$ be two surface classes with the same ends and fix a trivialization $\tau$ over $\Theta$ and $\Xi$. We compute the difference of each term in $I(A)$ and $I(B)$. Starting with the Maslov number, we have

$$
\mu_{\tau}(B)-\mu_{\tau}(A)=\mu(\xi, \Lambda) \cdot(B-A) \quad \text { by Proposition } 99
$$

To compute the difference between the self-intersection numbers, we note that
$Q_{\tau}(B)-Q_{\tau}(A)=Q_{\tau}(B, B)-Q_{\tau}(A, B)+Q_{\tau}(B, A)-Q_{\tau}(A, A)=q_{B}(B-A)+q_{A}(B-A)$
Here $q_{A}$ and $q_{B}$ are the homomorphisms $H_{2}(Y, \Lambda) \rightarrow \frac{1}{2} \mathbb{Z}$ in Definition 112. By Lemma 113

$$
q_{A}(B-A)+q_{B}(B-A)=2 \cdot Q(\Gamma, B-A)
$$

Finally, since $A$ and $B$ have the same ends, the ECH Conley-Zehnder terms coincide. This proves the desired formula.

Composition. This follows from the corresponding composition property for $Q_{\tau}$ and $\mu_{\tau}$.

## Index Inequality And Union Property

The Legendrian ECH index has a number of properties that hold for currents with boundary due to the writhe and linking inequalities. We are now ready to demonstrate these properties.

We begin with the the Legendrian ECH index inequality as stated in the introduction.

Theorem 192 (Theorem 7). Let $C$ be a somewhere injective, J-holomorphic curve with boundary in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ for tailored $J$. Then

$$
\begin{equation*}
\operatorname{ind}(C) \leq I(C)-2 \delta(C)-\epsilon(C) \tag{4.6.13}
\end{equation*}
$$

Proof. Let $C$ be a somewhere injective $J$-holomorphic curve asymptotic to Reeb chords and orbits $\Gamma_{ \pm}$at $\pm \infty$. Then the difference between ECH index and Fredholm index is
$I(C)-\operatorname{ind}(C)=Q_{\tau}(C)-\frac{1}{2} \mu_{\tau}(C)+\mathrm{CZ}_{\tau}^{E C H}\left(\Gamma_{+}\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\Gamma_{-}\right)+\bar{\chi}(C)-\mathrm{CZ}_{\tau}\left(\Gamma_{+}\right)+\mathrm{CZ}_{\tau}\left(\Gamma_{-}\right)$
Then by Legendrian adjunction, Theorem 144, the above formula becomes
$Q_{\tau}(C)+\mathrm{CZ}_{\tau}^{E C H}\left(\Gamma_{+}\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\Gamma_{-}\right)-\mathrm{CZ}_{\tau}\left(\Gamma_{+}\right)+\mathrm{CZ}_{\tau}\left(\Gamma_{-}\right)-Q_{\tau}(C)-w_{\tau}(C)+2 \delta(C)+\epsilon(C)$
Here the writhe $w_{\tau}(C)$ is the sum of the writhes of the positive asymptotic braid $\zeta_{+}$ and negative asymptotic braid $\zeta_{-}$of $C$. Thus we may write the above formula as
$=2 \delta(C)+\epsilon(C)+\left(\mathrm{CZ}_{\tau}^{E C H}\left(\zeta_{+}\right)-\mathrm{CZ}_{\tau}\left(\Gamma_{+}\right)-w_{\tau}\left(\zeta_{+}\right)\right)-\left(\mathrm{CZ}_{\tau}^{E C H}\left(\zeta_{-}\right)-\mathrm{CZ}_{\tau}\left(\Gamma_{-}\right)-w_{\tau}\left(\zeta_{-}\right)\right)$
The writhe bound, Proposition 186, implies that the middle and right terms are non-negative. Thus we have proven that

$$
I(C)-\operatorname{ind}(C) \geq 2 \delta(C)+\epsilon(C)
$$

Next, we prove a fundamental sub-additivity property of the ECH index under union. This generalizes [55, Thm. 5.1] to currents with boundary.

Theorem 193 (Union). Let $\mathcal{C}$ and $\mathcal{D}$ be J-holomorphic currents for a tailored J on $(Y, \Lambda)$, with distinct components. Then the ECH index satisfies

$$
I(\mathcal{C} \cup \mathcal{D}) \geq I(\mathcal{C})+I(\mathcal{D})+2(\mathcal{C} \cdot \mathcal{D})
$$

Proof. We compute that the difference of ECH indices is given by
$I(\mathcal{C} \cup \mathcal{D})-I(\mathcal{C})-I(\mathcal{D})=2 \cdot Q_{\tau}(\mathcal{C}, \mathcal{D})+\mathrm{CZ}_{\tau}^{E C H}(\mathcal{C} \cup \mathcal{D})-\mathrm{CZ}_{\tau}^{E C H}(\mathcal{C})-\mathrm{CZ}_{\tau}^{E C H}(\mathcal{D})$
By Lemma 159, this can be rewritten using the geometric intersection number as

$$
2(\mathcal{C} \cdot \mathcal{D})+\left(\mathrm{CZ}_{\tau}^{E C H}(\mathcal{C} \cup \mathcal{D})-\mathrm{CZ}_{\tau}^{E C H}(\mathcal{C})-\mathrm{CZ}_{\tau}^{E C H}(\mathcal{D})-2 \cdot l_{\tau}(\mathcal{C}, \mathcal{D})\right)
$$

The second term above is non-negative by the linking bound for currents (Proposition 188).

### 4.7 Legendrian ECH

We are now ready to construct Legendrian embedded contact homology in detail. Specifically, we prove the two key technical results Theorem 10 and Theorem 11 from the introduction.

Setup 194. To proceed with our construction, we fix the following setup for the rest of this section.
(a) $(Y, \xi)$ is a contact 3 -manifold with convex sutured boundary $\partial Y$.
(b) $\Lambda=\Lambda_{+} \cup \Lambda_{-}$is a union of exact Legendrians $\Lambda_{ \pm} \subset \partial_{ \pm} Y$.
(c) $\alpha$ is a non-degenerate, adapted contact form on $Y$.
(d) $J$ is a tailored complex structure on $\xi$

We start by recalling the definition of an ECH generator.
Definition 195. An ECH generator of $(Y, \Lambda)$ is an orbit-chord set $\Theta=\left\{\left(\gamma_{i}, m_{i}\right)\right\} \cup$ $\left\{\left(c_{i}, n_{i}\right)\right\}$ where

- Every hyperbolic orbit $\gamma_{i}$ has multiplicity 1.
- Every chord $c_{i}$ is multiplicity 1.
- There is at most one Reeb chord incident to $L$ in $\Theta$ for each connected component $L$ of $\Lambda$


## Classification Of Low ECH Index Currents

We begin by classifying low ECH index currents with boundary in Lemmas 196 and Proposition 197 and 198 below. These results, together, yield Theorem 10 in the introduction.

Lemma 196. Let $\mathcal{C}$ be an $J$-holomorphic current for a regular, tailored almost complex structure $J$ on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$. Then the following are equivalent.

- $\mathcal{C}$ is a trivial current $\mathbb{R} \times \Xi$ over an orbit-chord set $\Xi$.
- $\mathcal{C}$ has ECH index less than or equal to zero.

Proof. By the Legendrian ECH index inequality (Theorem 192) and the sub-additivity of the ECH index

$$
\operatorname{ind}\left(C_{i}\right) \leq I\left(C_{i}\right)-2 \delta\left(C_{i}\right)-\epsilon\left(C_{i}\right) \leq I(\mathcal{C}) \leq 0 \quad \text { where } \quad \mathcal{C}=\left\{\left(C_{i}, m_{i}\right)\right\}
$$

For any regular almost complex structure $J$, the Fredholm index of any somewhere injective curve is non-negative, and zero if and only if the curve is a trivial cylinder or strip. Therefore

$$
C_{i}=\mathbb{R} \times \gamma_{i} \quad \text { for an simple orbit or chord } \gamma_{i} \subset Y
$$

It follows that $\mathcal{C}$ is the trivial current $\mathbb{R} \times \Xi$ over the orbit-chord set $\Xi=\left\{\left(\gamma_{i}, m_{i}\right)\right\}$.

Proposition 197. Let $\mathcal{C}$ be a $J$-holomorphic current for a regular, tailored almost complex structure on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ of ECH index one. Then

$$
\mathcal{C}=C \sqcup \mathcal{T}
$$

where $\mathcal{T}$ is trivial and $C$ is an embedded, connected curve with $\operatorname{ind}(C)=I(C)=1$. Proof. We may write $\mathcal{C}=\mathcal{S} \cup \mathcal{T}$ where $\mathcal{T}$ is a trivial current $\mathbb{R} \times \Xi$ and $\mathcal{S}$ is a current with no trivial components. By Lemma 196 and the sub-additivity of the ECH index (Theorem 193), we have

$$
0<I(\mathcal{S}) \quad \text { and } \quad I(\mathcal{S}) \leq I(\mathcal{S})+I(\mathcal{T})+2(\mathcal{S} \cdot \mathcal{T}) \leq I(\mathcal{C}) \leq 1
$$

Thus, $\mathcal{S}$ is non-empty and $\mathcal{T}$ is disjoint from $\mathcal{S}$. We can then write

$$
\mathcal{S}=\left\{\left(C_{i}, m_{i}\right)\right\} \quad \text { for simple, non-trivial curves } C_{i} \text { and multiplicities } m_{i}
$$

If $\left(C_{i}, m_{i}\right)$ has multiplicity $m_{i}>1$, we may translate $m_{i}$ of copies of $C_{i}$ along the symplectization direction to obtain $m_{i}$ somewhere injective curves whose union forms a new $J$-holomorphic curve $C^{\prime}$.

We now apply the ECH index inequality (Theorem 192) to $C^{\prime}$ and obtain:

$$
\begin{equation*}
\sum_{i} m_{i} \operatorname{ind}\left(C_{i}\right) \leq \sum_{i} I(\mathcal{C})-2 \delta\left(C^{\prime}\right)-\epsilon\left(C^{\prime}\right) \tag{4.7.1}
\end{equation*}
$$

since the Fredholm index is additive under taking unions and ECH index is only dependent on the relative homology class. Since $J$ is regular, $\operatorname{ind}\left(C_{i}\right) \geq 0$ and in fact $\operatorname{ind}\left(C_{i}\right)>0$ by our assumption of $\mathcal{S}$ and Lemma 196. This contradicts (4.7.1). Furthermore,

$$
\mathcal{S}=\left\{\left(C_{1}, 1\right)\right\} \quad \text { for } C_{1} \text { an embedded, non-trivial curve }
$$

since $\delta\left(C_{1}\right)=\epsilon\left(C_{1}\right)=0$ by (4.7.1) again. Therefore, $\operatorname{ind}\left(C_{1}\right)=I\left(C_{1}\right)=1$.

Proposition 198. Let $\mathcal{C}$ be a $J$-holomorphic current for a regular, compatible almost complex structure on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ of ECH index two. Also assume that the ends $\Xi_{+}$and $\Xi_{-}$of $\mathcal{C}$ are ECH generators. Then

$$
\mathcal{C}=\mathcal{S} \sqcup \mathcal{T}
$$

where $\mathcal{T}$ is trivial and $\mathcal{S}$ is a current of one of the following types.

- A pair of disjoint, embedded curves $C_{1} \sqcup C_{2}$ with $I\left(C_{i}\right)=\operatorname{ind}\left(C_{i}\right)=1$. Note here that $C_{1}$ and $C_{2}$ are distinct in the sense they are not $\mathbb{R}$ translates of each other.
- A single embedded curve $C$ of multiplicity 1 with $I(C)=\operatorname{ind}(C)=2$.

To prove Proposition 198, the following lemma will be helpful. It severely limits the type of boundary singularities that can occur in currents connecting ECH generator.

Lemma 199 (Singularity Censorship). Let $\mathcal{C}$ be a (finite energy, proper) $J$-holomorphic current with boundary between ECH generators $\Xi_{+}$and $\Xi_{-}$. Then

- $\mathcal{S} \cdot \mathcal{T}$ is an integer for any pair of currents $\mathcal{S}$ and $\mathcal{T}$ with disjoint components such that $\mathcal{S} \cup \mathcal{T} \subset \mathcal{C}$.
- Every component curve $C$ of $\mathcal{C}$ is non-singular near the boundary. That is $\epsilon(C)=0$.

Proof. For the first claim, it suffices to show that any pair of component curves $C$ and $D$ satisfy

$$
C \cap D \cap(\mathbb{R} \times \Lambda)=\emptyset
$$

Thus suppose otherwise. Then there are components $\Gamma$ and $\Gamma^{\prime}$ of the boundary of $C$ and $D$, respectively, such that $\Gamma \cap \Gamma^{\prime}$ is non-empty and lie on $\mathbb{R} \times L$ where $L \subset \Lambda$ is a component.

By Lemma 131, both components $\Gamma$ and $\Gamma^{\prime}$ are diffeomorphic to $\mathbb{R}$. Thus $C$ and $D$ both have punctures asymptotic at $+\infty$ to a Reeb chord of $L$. Thus $\Xi_{+}$has either a Reeb chord of multiplicity two or two Reeb chords incident to a single component $L$ of $\Lambda$. This is a contradiction, since $\Xi_{ \pm}$is an ECH generator. The second claim follows immediately from Lemma 131.
proof of Proposition 198. We again write $\mathcal{C}=\mathcal{S} \cup \mathcal{T}$ where $\mathcal{T}$ is a trivial current $\mathbb{R} \times \Xi$ and $\mathcal{S}$ is a current with no trivial components. Applying Lemma 196 and Theorem 193 now yields

$$
0<I(\mathcal{S}) \quad \text { and } \quad I(\mathcal{S}) \leq I(\mathcal{S})+I(\mathcal{T})+2(\mathcal{S} \cdot \mathcal{T}) \leq I(\mathcal{C}) \leq 2
$$

Since $\mathcal{C}$ has ends on ECH generators, $\mathcal{S}$ and $\mathcal{T}$ are disjoint on $\mathbb{R} \times \Lambda$. By Lemma 199, we must have $\mathcal{S} \cdot \mathcal{T}=0$ or 1 . In the later case, $I(\mathcal{S})=0$. Thus we must have

$$
I(\mathcal{S})=2 \quad \text { and } \quad \mathcal{S} \cdot \mathcal{T}=\emptyset
$$

As in Proposition 197, we can assume that $\mathcal{S}$ consists of connected, somewhere injective components $C_{i}$ of multiplicity 1 . We apply the ECH index inequality (Theorem 192) and the sub-additivity of the ECH index (Theorem 193) to find that
$\sum_{i} \operatorname{ind}\left(C_{i}\right) \leq \sum_{i} I\left(C_{i}\right)-2 \delta\left(C_{i}\right)-\epsilon\left(C_{i}\right) \quad$ and $\quad \sum_{i} I\left(C_{i}\right)+2 \sum_{i<j} C_{i} \cdot C_{j} \leq I(\mathcal{S})=2$
Since each $C_{i}$ is non-trivial, we must have $\operatorname{ind}\left(C_{i}\right)>0$ for each $i$. Thus we infer from Lemma 199 that $C_{i} \cdot C_{j}=0$ for each $i$ and $j$ and $\delta\left(C_{i}\right)=\epsilon\left(C_{j}\right)=0$.

Finally we rule out the case of $\mathcal{C}$ containing a nontrivial curve $C$ with multiplicity two. By assumption on ECH generators, $C$ cannot contain chords or be asymptotic to hyperbolic Reeb orbits. Then the same proof as in regular ECH shows $I(C)-\operatorname{ind}(C)$ is even. Hence we cannot have nontrivial curves of multiplicity $\geq 2$.

## Compactness Of Low ECH Index Moduli Spaces

Next, we prove the requisite compactness properties for the low ECH index moduli spaces required in ECH.

We begin by proving a crude topological lower bound for the (corrrected) Euler characteristic of $J$-holomorphic curves between ECH generators.

Lemma 200 (Topological Index Bound). Let $C$ be a (proper, finite energy) $J$ holomorphic curve in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ between ECH generators $\Xi_{+}$and $\Xi_{-}$in the surface class $A$. Then

$$
-\bar{\chi}(C) \leq-\frac{1}{2} \cdot \mu_{\tau}(A)+Q_{\tau}(A)+\mathrm{CZ}_{\tau}^{E C H}\left(\check{\Xi}_{+}\right)-\mathrm{CZ}_{\tau}^{E C H}\left(\check{\Xi}_{-}\right)
$$

Here $\check{\Xi}_{+}$and $\check{\Xi}_{-}$are the orbit-chord sets made from $\Xi_{+}$and $\Xi_{-}$by reducing the multiplicity of each orbit and chord by one (and removing everything of multiplicity one).

Proof. We apply Legendrian adjunction (Corollary 144) to see that
$-\bar{\chi}(C)=-\frac{1}{2} \cdot \mu_{\tau}(C)+Q_{\tau}(C)+w_{\tau}(C)-2 \delta(C)-\epsilon(C) \leq-\frac{1}{2} \cdot \mu_{\tau}(C)+Q_{\tau}(C)+w_{\tau}(C)$
We may decompose the writhe of $C$ into a sum of writhes

$$
w_{\tau}(C)=\sum_{i} w_{\tau}\left(\zeta_{i}^{+}\right)-\sum_{j} w_{\tau}\left(\zeta_{j}^{-}\right)
$$

where the sum is over all braids over each orbit and chord appearing as a positive and negative end of $C$. The writhe is automatically 0 for any chord and orbit of multiplicity 1. Since $\Xi_{+}$and $\Xi_{-}$are ECH generators, this includes all chords and hyperbolic orbits. This reduces the desired result to the inequality

$$
\begin{equation*}
w_{\tau}\left(\zeta^{+}\right) \leq \mathrm{CZ}_{\tau}^{E C H}((\eta, m-1)) \tag{4.7.2}
\end{equation*}
$$

for any simple elliptic orbit of multiplicity $m>1$ in $\Xi_{+}$and the corresponding inequality

$$
\begin{equation*}
w_{\tau}\left(\zeta^{-}\right) \geq \mathrm{CZ}_{\tau}^{E C H}((\gamma, n-1)) \tag{4.7.3}
\end{equation*}
$$

for any simple elliptic orbit of multiplicity $n>1$ in $\Xi_{-}$. This is proven in [55, Prop. 6.9]. In particular, [55, Eq. 6.2] states that

$$
\mathrm{CZ}^{E C H}((\eta, m-1))-w_{\tau}\left(\zeta^{+}\right) \geq n-1
$$

where $\eta$ is elliptic and $n$ denotes the number of components of the braid $\zeta^{+}$. This implies (4.7.2), and (4.7.3) can be argued analogously.

Remark 201. In [55], Hutchings also develops a general theory of the $J_{0}$-index, a replacement for the ECH index that provides a topological filtration on any type of contact homology. Proposition 200 is essentially using this theory in the relative case. We will not develop a relative $J_{0}$-invariant in this paper.

Next, we demonstrate the compactness of the space of $J$-holomorphic currents of ECH index 1. This is the first part of Theorem 11 in the introduction.

Lemma 202. The moduli space of $J$-holomorphic currents $\mathcal{M}_{1}\left(\Xi_{+}, \Xi_{-}\right) / \mathbb{R}$ of ECH index 1 is finite.

This is a virtual repeat of the same argument in ECH using Gromov compactness and bounds on topological complexity of $J$ holomorphic curves. We sketch it here briefly for completeness.

Proof. Fix a sequence of distinct currents with boundary of index 1 , denoted by

$$
\mathcal{C}^{\nu} \in \mathcal{M}\left(\Xi_{+}, \Xi_{-}\right) / \mathbb{R}
$$

By Proposition 197, we have $\mathcal{C}^{\nu}=C^{\nu} \sqcup \mathcal{T}^{\nu}$ where $C^{\nu}$ is a connected, embedded $J$-holomorphic curve with $I\left(C^{\nu}\right)=1$ and $\mathcal{T}^{\nu}$ is a trivial current over an orbit-chord subset $\Theta^{\nu}$ of $\Xi_{+}$and $\Xi_{-}$. There are only finitely many orbit-subsets of $\Xi_{+}$, so after passing to a subsequence we may assume that

$$
\mathcal{T}^{\nu}=\mathcal{T}=\mathbb{R} \times \Theta \quad \text { for all } \nu
$$

Moreover, let $\Theta_{+}$and $\Theta_{-}$be the unique orbit sets such that $\Xi_{ \pm}=\Theta_{ \pm} \cup \Theta$. Then

$$
C^{\nu} \in \mathcal{M}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}
$$

By Gromov compactness for currents (Theorem 153), we can choose a subsequence of $C^{\nu}$ that lie in the same surface class $A$ in $S\left(\Theta_{+}, \Theta_{-}\right)$. By Lemma 200, we have

$$
-\bar{\chi}(C) \leq-\frac{1}{2} \cdot \mu_{\tau}(A)+Q_{\tau}(A)+\mathrm{CZ}_{\tau}^{E C H}\left(\check{\Theta}_{+}\right)-\mathrm{CZ} Z_{\tau}^{E C H}\left(\check{\Theta}_{-}\right)
$$

Thus the genus and number of boundary components of $C$ is uniformly bounded. We can thus apply SFT compactness for curves with boundary (cf. Abbas [1, Thm. $3.6]$ ) to acquire a $J$-holomorphic building

$$
\mathbf{C}=\left(C_{1}, C_{2}, \ldots, C_{m}\right)
$$

Claim 203. The building $\mathbf{C}$ consists of a single level $C$ of ECH index 1 .
Proof. Since the ECH index is additive under composition and the surface class of C agrees with that of $C^{\nu}$ for large $\nu$, we know that

$$
\sum_{j} I\left(C_{j}\right)=I\left(C^{\nu}\right)=1
$$

Moreover, $I\left(C_{j}\right) \geq 0$ for each $j$ since $J$ is regular. Therefore, for some fixed $i$ we must have

$$
I\left(C_{i}\right)=\operatorname{ind}\left(C_{i}\right)=1
$$

and all other levels $C_{j}$ are ECH index 0 . By Lemma 196, we have the following equality of currents

$$
C_{j}=\mathbb{R} \times \Xi_{+} \quad \text { for } j<i \quad \text { and } \quad C_{j}=\mathbb{R} \times \Xi_{-} \quad \text { for } j>i
$$

and each $C_{j}$ is (as a curve) an explicit branched cover of the trivial cylinders over the simple orbits and chords of $\Xi_{ \pm}$.

Now fix an orbit or chord $\gamma$ of multiplicity $m$ in $\Xi_{+}$. Let $Z_{j}$ be the sub-curve $C_{j}$ that is a cover of $\mathbb{R} \times \gamma$ for $1 \leq j<i$ and consider the building $\mathbf{Z}=\left(Z_{1}, \ldots, Z_{i-1}, Z_{i}\right)$. Then Lemma 180 implies that $\mathbf{Z}$ consists of unbranched covers if $\gamma$ is an orbit, and Lemma 181 implies the same if $\gamma$ is a chord. Such components cannot appear in a building, so we must have $i=1$ and $C_{1}=C_{i}$ is non-trivial.

The same argument shows that there are no trivial levels $C_{j}$ for $j>i$.
Returning to the proof of Lemma 202, we now find that the embedded curves $C^{\nu}$ converge to the curve $C$ in the moduli space of somewhere injective Fredholm index 1 curves. Since $J$ is regular, this moduli space is discrete. Thus $C^{\nu}=C$ for sufficiently large $\nu$, and

$$
\mathcal{C}^{\nu}=C^{\nu} \sqcup \mathcal{T}=C \sqcup \mathcal{T} \quad \text { for large } \nu
$$

This contradicts the assumption that the currents $\mathcal{C}^{\nu}$ are distinct. This concludes the proof.

Next, we demonstrate the compactness of the space of $J$-holomorphic currents of ECH index 2, under the assumption that the positive and negative ends are on ECH generators.

Lemma 204. Let $\Xi_{1}, \Xi_{2}$ be ECH generators. Fix a regular, tailored $J$ on $(\mathbb{R} \times Y, \mathbb{R} \times$ $\Lambda$ ). Then the moduli space of ECH index two holomorphic currents

$$
\mathcal{M}_{2}\left(\Xi_{1}, \Xi_{2}\right) / \mathbb{R}
$$

is a topological 1-manifold that admits a (possibly singular) compactification

$$
\overline{\mathcal{M}}_{2}\left(\Xi_{1}, \Xi_{2}\right) / \mathbb{R}:=\mathcal{M}_{2}\left(\Xi_{1}, \Xi_{2}\right) / \mathbb{R} \cup\left(\bigcup_{\Theta} \mathcal{M}_{1}\left(\Xi_{1}, \Theta\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta, \Xi_{2}\right) / \mathbb{R}\right)
$$

Note that the above compactification is topologized with the Gromov topology on currents. It is not necessarily a manifold itself. There could be many connected components of $\mathcal{M}_{2}\left(\Xi_{1}, \Xi_{2}\right) / \mathbb{R}$ that have the same boundary point in the compactification. The number of such components is given by a obstruction bundle count that we describe in more detail below (Section 4.7).

Proof. By Gromov compactness for currents, we know that $C^{\nu}$ converges in the Gromov topology to a broken current

$$
\overline{\mathcal{C}}=\left(\mathcal{C}_{1}, \ldots, \mathcal{C}_{m}\right)
$$

The ECH index is additive (Proposition 191) and non-negative (by Lemma 196), so the broken current must have $m \leq 2$ component currents. It suffices to show that if $m=1$ and $\overline{\mathcal{C}}=\mathcal{C}$, then

$$
\mathcal{C}=C
$$

where $C$ is a connected embedded curve of ECH index two. This is the content of Proposition 198.

## Moduli Spaces Truncation

We now formally describe the truncated moduli space of ECH index two currents in Theorem 11.

In the case of standard ECH, the construction of this moduli space and the proof of its key properties is spread over several papers [52, 60, 62]. Thus, although the construction is exactly analogous to that case, we collect several details here for the reader.

We will require a version of gluing pairs in the sense of [60, Definition 1.9]. We adopt the following definition.

Definition 205. An ECH gluing pair in $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$ is a pair of $J$-holomorphic curves

$$
\left(u_{+}, u_{-}\right)
$$

satisfying the following properties.

- $u_{+}$and $u_{-}$have ECH index one.
- $u_{+}$and $u_{-}$are embedded, except for unbranched covers of cylinders of orbits and chords.
- The orbit-chord sets at the negative end of $u_{+}$and the positive end of $u_{-}$are the same.

Moreover, let $\gamma$ be a simple elliptic orbit of total multiplicity $m$ at the negative end of $u_{-}$(or equivalently, at the positive end of $u_{+}$).

- The partition $p_{-}\left(u_{+} ; \gamma\right)$ of $m$ determined by the negative end of $u_{+}$is the negative ECH partition $p_{-}(m)$.
- The partition $p_{+}\left(u_{-} ; \gamma\right)$ of $m$ determined by the positive end of $u_{-}$is the positive ECH partition $p_{+}(m)$

There is a space of ECH index two curves that are nearly broken at $\left(u_{+}, u_{-}\right)$, denoted by

$$
\mathcal{G}_{\delta}\left(u_{+}, u_{-}\right) \subset \mathcal{M}_{2}\left(\Xi_{+}, \Xi_{-}\right) / \mathbb{R}
$$

consisting of all connected, embedded curves $C$ in $\mathcal{M}_{2}\left(\Xi_{1}, \Xi_{2}\right)$ that admit constants $R_{+}, R_{-}$and a decomposition

$$
C=C_{+}^{\prime} \cup C_{0} \cup C_{-}^{\prime}
$$

that satisfy the following conditions.

- $C_{0}$ is contained in the radius $\delta$ neighborhood of $\mathbb{R} \times \Theta$, where $\Theta$ is the orbitchord set at the negative end of $u_{+}$.
- $R_{+}-R_{-}>\frac{2}{\delta}$.
- There is a section $\psi_{+}$of the normal bundle of $u_{+}$such that

$$
\left|\psi_{+}\right|<\delta \quad \text { and } \quad C_{+}^{\prime}=\left(\exp _{C_{+}}\left(\psi_{+}\right)+R_{+}\right) \cap[-1 / \delta, \infty) \times Y
$$

- There is a section $\psi_{-}$of the normal bundle of $u_{-}$such that

$$
\left|\psi_{-}\right|<\delta \quad \text { and } \quad C_{-}^{\prime}=\left(\exp _{C_{-}}\left(\psi_{-}\right)+R_{-}\right) \cap(-\infty, 1 / \delta] \times Y
$$

Lemma 206. Fix an ECH gluing pair of the form

$$
\left(u_{+}, u_{-}\right)
$$

Then there is an $\epsilon>0$ such that any sequence of curves $C_{i} \in \mathcal{G}_{\epsilon}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right) / \mathbb{R}$ has a subsequence that converges in the SFT topology (see [1, Thm 3.6]) to one of the following types of $J$-holomorphic buildings.

- A connected, embedded curve $C$ of ECH index two in $\mathcal{G}_{\delta}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)$
- A building $\left(u_{+}, v_{1}, \ldots, v_{k}, u_{-}\right)$where $v_{i}$ is a union of branched covers of trivial cylinders representing $\mathbb{R} \times \Theta$.

Proof. Let $C_{i} \in \mathcal{G}_{\epsilon}\left(u_{+}, u_{-}\right) / \mathbb{R}$ be a sequence. By SFT compactness in the relative case [1, Thm 3.6] and the topological bounds in Lemma 200, we know that there is a limit SFT building of the form

$$
\mathbf{C}=\left(w_{1}, \ldots, w_{m}\right)
$$

The sum of the ECH indices of the levels must be 2 by addivity of the ECH index. Therefore, there are two cases.

Case 1. In the first case, there is a single level of ECH index two and every other level is ECH index zero. In this case, we can argue that there is only one level $C=\mathbf{C}$, by an identical as in Claim 203. It follows that $C_{i} \rightarrow C$ in the SFT topology.

Case 2. In the second case, there are two levels $w_{a}$ and $w_{b}$ of $\mathbf{C}$ of ECH index one, and the remaining levels are ECH index zero. The levels $w_{a}$ and $w_{b}$ must be equal to $u_{+}$and $u_{-}$due to the construction of $\mathcal{G}_{\delta}$ and the definition of SFT convergence. Moreover, since the positive and negative ends of the curves $C_{i}$ satisfy the ECH partition conditions, by [60, Lemma 1.7] and [52, Exercise 3.14], the levels $w_{j}$ must be trivial for $j<a$ and $j>b$.

Definition 207. The count of gluings $\# G\left(u_{+}, u_{-}\right)$of an ECH gluing pair ( $u_{+}, u_{-}$) is defined as follows. Let $\epsilon$ be as in Lemma 206 and choose an open subset $U \subset$ $\mathcal{G}_{\epsilon}\left(u_{+}, u_{-}\right) / \mathbb{R}$ such that

- $\bar{U}$ has finitely many boundary points.
- $U$ contains $\mathcal{G}_{\delta}\left(u_{+}, u_{-}\right) / \mathbb{R}$ for some $\delta<\epsilon$.

Then we define $\# G\left(u_{+}, u_{-}\right)$to be the number of boundary points in the closure $\bar{U}$ of $U$.

We are now ready to construct the truncated moduli space of ECH index two curves. For each ECH gluing pair $\left(u_{+}, u_{-}\right)$between ECH generators $\Xi_{+}$and $\Xi_{-}$, fix for the remainder of the section an open set

$$
U\left(u_{+}, u_{-}\right) \subset \mathcal{G}_{\epsilon}\left(u_{+}, u_{-}\right)
$$

as in Definition 207. For any ECH generators $\Theta_{+}$and $\Theta_{-}$, we adopt the notation $W\left(\Theta_{+}, \Theta_{-}\right):=\left\{C \sqcup \mathcal{T}: C \in U\left(u_{+}, u_{-}\right)\right.$for some ECH gluing pair $\left(u_{+}, u_{-}\right)$and $\mathcal{T}$ is trivial $\}$

Definition 208 (Truncated Moduli Space). Let $\Theta_{+}$and $\Theta_{-}$be ECH generators and fix a regular, tailored $J$ on $(\mathbb{R} \times Y, \mathbb{R} \times \Lambda)$. The truncated moduli space

$$
\mathcal{M}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R} \subset \mathcal{M}_{2}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}
$$

is defined as follows. By Proposition 198, we may divide $\mathcal{M}_{2}\left(\Theta_{+}, \Theta_{-}\right)$into disjoint pieces

$$
\mathcal{A}_{2}\left(\Theta_{+}, \Theta_{-}\right)=\{C \sqcup \mathcal{T}: C \text { is connected with } I(C)=2 \text { and } \mathcal{T} \text { trivial. }\}
$$

$\mathcal{B}_{2}\left(\Theta_{+}, \Theta_{-}\right)=\{C \sqcup D \sqcup \mathcal{T}: C, D$ are connected with $I(C)=I(D)=1$ and $\mathcal{T}$ trivial. $\}$ We define $\mathcal{M}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}$ as a union of pieces

$$
\mathcal{A}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) \subset \mathcal{A}_{2}\left(\Theta_{+}, \Theta_{-}\right) \quad \text { and } \quad \mathcal{B}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) \subset \mathcal{B}_{2}\left(\Theta_{+}, \Theta_{-}\right)
$$

Truncation Of $\mathcal{A}$. To truncate $\mathcal{A}_{2}\left(\Theta_{+}, \Theta_{-}\right)$, we note that by Gromov compactness for currents (Theorem 153), any sequence of currents $\mathcal{C}_{i} \in \mathcal{A}_{2}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}$ has a subsequence of the form

$$
\mathcal{C}_{i}=C_{i} \sqcup \mathcal{T}
$$

where $C_{i}$ is a connected curve with $I\left(C_{i}\right)=2$ from $\Theta_{+}$to $\Theta_{-}$converging to

- a connected ECH index two curve $C \in \mathcal{M}_{2}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}$ or
- a building $\left(u_{+}, v_{1}, \ldots, v_{k}, u_{-}\right)$where $\left(u_{+}, u_{-}\right)$are an ECH gluing pair and $v_{i}$ are covers of trivial cylinders and chords.

In the latter case, it follows that $C_{i} \notin W\left(\Theta_{+}, \Theta_{-}\right)$for sufficiently large $i$. We thus let

$$
\mathcal{A}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right)=\mathcal{A}_{2}\left(\Theta_{+}, \Theta_{-}\right) \backslash W\left(\Theta_{+}, \Theta_{-}\right)
$$

There is a natural projection map

$$
\Pi: \partial \mathcal{A}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R} \rightarrow \bigsqcup_{\Theta^{\prime}} \mathcal{M}_{1}\left(\Theta_{+}, \Xi\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Xi, \Theta_{-}\right) / \mathbb{R}
$$

defined by mapping $C \sqcup \mathcal{T} \in \partial \mathcal{A}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}$ to the broken $J$-holomorphic current

$$
\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)=\left(u_{+} \sqcup \mathcal{T}, u_{-} \sqcup \mathcal{T}\right)
$$

where $\left(u_{+}, u_{-}\right)$is the unique ECH gluing pair such that $C \in \mathcal{G}_{\delta}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)$. Moreover, we have

$$
\# \Pi^{-1}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)=\# \bar{U}\left(u_{+}, u_{-}\right)=\# G\left(u_{+}, u_{-}\right)
$$

Truncation Of $\mathcal{B}$. To truncate $\mathcal{B}_{2}\left(\Theta_{+}, \Theta_{-}\right)$, note that a component $S \subset$ $\mathcal{B}_{2}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}$ is 1-dimensional, consisting of curves of the form

$$
\mathcal{C}_{s}=C \sqcup(D+s) \sqcup \mathcal{T} \quad \text { for } s \in \mathbb{R}
$$

Here $C$ and $D$ are connected, embedded curves of ECH index 1 and $\mathcal{T}$ is a trivial current. Let $\mathcal{S}_{+}$and $\mathcal{S}_{-}$be the trivial currents over the positive and negative ends of $C$, and let $\mathcal{T}_{+}$and $\mathcal{T}_{-}$be the analogous currents for $D$. Then

$$
\begin{aligned}
& \lim _{s \rightarrow-\infty} \mathcal{C}_{s}=\left(C \sqcup \mathcal{T}_{+} \sqcup \mathcal{T}, D \sqcup \mathcal{S}_{-} \sqcup \mathcal{T}\right) \\
& \lim _{s \rightarrow+\infty} \mathcal{C}_{s}=\left(D \sqcup \mathcal{S}_{+} \sqcup \mathcal{T}, C \sqcup \mathcal{T}_{+} \sqcup \mathcal{T}\right)
\end{aligned}
$$

where the limit is taken in the Gromov topology on broken currents. We now truncate $S$ by setting

$$
S^{\prime}=\left\{\mathcal{C}_{s}: s \in[-1,1]\right\} \subset S
$$

and let $\mathcal{B}_{2}^{\prime}$ be the union of these pieces over all components $S$. There is a natural projection map

$$
\Pi: \partial \mathcal{B}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R} \rightarrow \bigsqcup_{\Theta^{\prime}} \mathcal{M}_{1}\left(\Theta_{+}, \Xi\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Xi, \Theta_{-}\right) / \mathbb{R}
$$

sending $C_{-1}$ to $\lim _{s \rightarrow-\infty} \mathcal{C}_{s}$ and $C_{1}$ to the broken current $\lim _{s \rightarrow \infty} \mathcal{C}_{s}$. Note that

$$
\# \Pi^{-1}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)=1 \quad \text { for any } \quad\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right) \in \Pi\left(\partial \mathcal{B}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) / \mathbb{R}\right)
$$

## Gluing Counts

We conclude this section by explaining the applications of the obstruction bundle gluing results of Hutchings-Taubes [60,62] to our setting. In particular, we prove the last item of Theorem 11 as Corollary 211. This completes the proofs of all the main results of the paper.

Let $\gamma$ be a simple Reeb orbit and fix two partitions $p$ and $q$ of an integer $m$. In [60, §1.5, 1.6], Hutchings-Taubes define gluing coefficients

$$
c_{\gamma}(p, q) \in \mathbb{Z}
$$

that count (roughly) the number of ways to glue two curves asymptotic to $\gamma$ with multiplicity $m$, and corresponding partitions $p$ and $q$. For chords, we adopt the following definition.

Definition 209. The gluing coefficient of a Reeb chord $\eta$ of $(Y, \Lambda)$ and of multiplicity one is given by

$$
c_{\gamma}(\gamma, 1)=1
$$

Given a $J$-holomorphic curve $v$, we let $m_{ \pm}(v, \eta)$ be the multiplicity of the orbit or chord $\eta$ at the $\pm$-end of $v$, and $p_{ \pm}(v, \gamma)$ be the partition of $m_{ \pm}(v, \gamma)$ determined by the $\pm$-end of a finite energy curve $v$ at any orbit $\gamma$. Finally, if $\gamma$ is a simple orbit and $\eta$ is a chord, then

$$
c_{\gamma}\left(u_{+}, u_{-}\right):=c_{\gamma}\left(p_{-}\left(u_{+}, \gamma\right), p_{+}\left(u_{-}, \gamma\right)\right) \quad \text { and } \quad c_{\eta}\left(u_{+}, u_{-}\right):=1
$$

Proposition 210. Fix a tailored, regular almost complex structure $J$ on $(\mathbb{R} \times Y, \mathbb{R} \times$ $\Lambda)$ and let $\left(u_{+}, u_{-}\right)$be an ECH gluing pair. Then

$$
\# G\left(u_{+}, u_{-}\right)=\prod_{\gamma} c_{\gamma}\left(u_{+}, u_{-}\right)
$$

The product is over all simple closed orbits and chords appearing as a negative end of $u_{+}$.

Proof. The essential ingredients to this formula appear already in Hutchings-Taubes [60, 62]. Here we briefly sketch the argument.

Let $u_{1}$ and $u_{2}$ be a ECH gluing pair. Throughout we assume we have chosen a generic $J$. For simplicity of exposition we first assume none of the $u_{1}$ and $u_{2}$ have boundary, and that all negative ends of $u_{1}$ are asymptotic to one (simple) Reeb orbit
with multiplicities (therefore the same is true for all positive punctures of $u_{2}$ ). We illustrate the general methodology in this case and explain the more general case later.

Then recall from Section 5.2 [62] in order to glue $u_{1}$ and $u_{2}$ together, we must first preglue a branched cover of trivial cylinder $v$ of Fredholm index zero between them, so that the partition conditions at positive (resp. negative ends) of $v$ match the partition conditions of negative ends of $u_{1}$ (resp. positive ends of $u_{2}$ ). Then as in Section 5.4 that there is a gluing is translated into three equations as we describe below.

To borrow the notation of Section 5.3 in [62], let $N_{*}$ denote the normal bundle of $u_{i}, v$. Let $\mathcal{H}_{1}\left(N_{*}\right)$ denote a suitable completion of sections of $N_{i}$. Let $\psi_{i} \in \mathcal{H}_{1}\left(N_{i}\right)$ and $\phi \in \mathcal{H}_{2}\left(N_{v}\right)$ denote sections of the normal bundle. We consider deforming the preglued curve consisting of $u_{1}, u_{2}$ and $v$ using the sections $\psi_{i}$ and $\phi$ (patched together using cutoff functions, as in Section 5.3 in [62]). The condition that gluing exists is translated into a system of three equations

$$
\Theta_{1}\left(\psi_{1}, \phi\right)=0, \quad \Theta_{v}\left(\psi_{1}, \psi_{2}, \phi\right)=0, \quad \Theta_{2}\left(\psi_{2}, \phi\right)=0
$$

so that given fixed $\phi$, the equations $\Theta_{1}, \Theta_{2}$ can always be solved essentially uniquely (with $\psi_{i}$ expressed as functions of $\phi$ ). See Section 5.6 of [62].

The number of gluings is equivalent to the number of solutions of $\Theta_{v}$ with independent variable $\phi$; and $\psi_{1}$ and $\psi_{2}$ are functions of $\phi$. This is the content of Theorem 7.3 (b) in [62].

The number of solutions of $\Theta_{v}=0$ is counted by the number of zeroes of a section $\mathfrak{s}$ of an obstruction bundle $\mathcal{O}(\Sigma) \rightarrow \mathcal{M}_{\Sigma}$. The base $\mathcal{M}_{\Sigma}$ is the space of branched covers of the trivial cylinder satisfying the same multiplicity and parition conditions as $v$; and the fiber is the (dual of) cokernel of a linearized Cauchy Riemann operator $D_{\Sigma}$ acting on the normal bundle of the branched cover. This obstruction bundle is defined carefully in Definition 2.16 of [60]. See Section 2.1 of [60] for definitions and properties of the base $\mathcal{M}_{\Sigma}$ and Section 2.2 and 2.3 for the linearized Cauchy Riemann $D_{\Sigma}$. This obstruction section $\mathfrak{s}$ whose zero corresponds to the number of solutions of $\Theta_{v}=0$ is described carefully in Equation 5.43 of [62]. The count of zeroes can be viewed as computing some version of a relative Euler class of this obstruction bundle. The zeros of $\mathfrak{s}$ are transverse (proved in Section 10 of [62]), and the number of zeroes counted with sign is equal to $c_{\gamma}$. The definition of $c_{\gamma}$ is given in equation
1.7 of [60], and described more carefully in Sections 1.5 and 1.6 of [60]. That the number of zeros is equal to $c_{\gamma}$ is the content of Theorem 1.13 in [60].

From this way of gluing, we see readily that we can count the number of gluings as a local count of zeroes of obstructions sections over spaces of branched covers of trivial cylinders. If $u_{1}$ and $u_{2}$ meet along multiple Reeb orbits, the count of total number of gluings is the product of the count of zeroes of such obstruction sections over all the Reeb orbits. This is the content of Theorem 1.13 in [60].

In the case of ECH gluing pairs for Legendrian ECH, we can again reduce the count of gluing to count of zeroes of obstruction sections over branched covers of trivial cylinders and trivial strips. We here use the observation that by partition conditions all chords between $u_{+}$and $u_{-}$must have mulitplicity one. Over branched covers of trivial cylinders such counts over closed Reeb orbits is given by the obstruction bundle counts given in [60, Theorem 1.13], and the gluing over Reeb chords is just standard gluing.

Corollary 211 (Theorem 11, Truncation). The inverse image $\Pi^{-1}\left(\mathcal{C}_{+}, \mathcal{C}_{-}\right)$of a pair of currents in $\mathcal{M}_{1}\left(\Theta, \Theta^{\prime}\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Theta^{\prime}, \Xi\right) / \mathbb{R}$ under the map

$$
\Pi: \mathcal{M}_{2}^{\prime}\left(\Theta_{+}, \Theta_{-}\right) \rightarrow \bigsqcup_{\Theta^{\prime}} \mathcal{M}_{1}\left(\Theta_{+}, \Xi\right) / \mathbb{R} \times \mathcal{M}_{1}\left(\Xi, \Theta_{-}\right) / \mathbb{R}
$$

has an odd number of points if and only if the orbit set $\Theta^{\prime}$ is an ECH generator.
Proof. For any ECH gluing pair, the gluing coefficients $c_{\gamma}\left(u_{+}, u_{-}\right)$are always odd when $\gamma$ is is an elliptic orbit [60, Prop. 7.26]. They are odd if and only if the multipliciy of $\gamma$ in $\Xi_{+}$or $\Xi_{-}$is one when $\gamma$ is a hyperbolic orbit [60, Def. 1.14] or chord (see Definition 209). Thus this follows from Proposition 210.

### 4.8 Decay Estimates

In this appendix, we prove asymptotic formulas for ends of $J$ holomorphic curves converging to a Reeb chord in the same style of Siefring [108]. The asymptotic formula for a single end is already established in [2]. Our goal is to explain the general case, where multiple ends approach the same chord. This will be essential in our proof of the general writhe bound.

Remark 212. The general writhe bound is not strictly necessary for our formulation of Legendrian ECH in Section 4.1. This is due to our combination of maximum
principles and the definition of ECH generators prevents breaking along multiply covered chords in the ECH setting.

However, the general asymptotic formula and writhe bound may play a role in other theories that count curves asymptotic to chords with higher multiplicities.

The presented proof is an amalgamation of techniques and ideas found in [47], [108], and [2]. The proof has the following outline.
(i) First, we describe the local geometric setup as in [2]. In particular, we follow the choices of metrics and connections in [2].
(ii) Suppose we have two ends approaching the same chord. Following [108], we show the difference between the two ends satisfy a particular PDE (see equation 4.8.3).
(iii) We show that the bounded solutions to Equation 4.8.3 have an asymptotic representation formula. We do this using the setup of [2], and essentially the same techniques there that proved asymptotic formula holds for a single end.
(iv) Finally we translate the asymptotic formula of the difference between two ends into behaviour of two ends of the holomorphic curve, following [108].

## Local differential geometry

We start by recalling the geometric setup of [2]. By rescaling the contact form, we may assume that the Reeb chord has action 1. We choose a neighborhood of the Reeb chord and local coordinates $x, y, z$ such that the following holds.

- The Reeb chord $\eta$ is the line segment $\{(0,0, t): t \in[0,1]\}$
- The two Legendrians consist of the lines

$$
L_{1}:=\{(t, 0,0), t \in \mathbb{R}\} \quad \text { and } \quad L_{2}:=\{(0, t, 1), t \in \mathbb{R}\}
$$

- The contact form $\lambda$ is given by

$$
\lambda=f d z+c_{1} d x+c_{2} d y
$$

where $f=1$ along $\eta$, and $c_{1}$ and $c_{2}$ are smooth functions satisfying $c_{1}=c_{2}=0$ on the $z$-axis.

- The contact form $\lambda$ satisfies

$$
d \lambda(0,0, z)=a d x \wedge d y
$$

where $a(x, y, z)$ is positive in a neighborhood of the Reeb chord.

- The Reeb vector field is given by

$$
X(x, y, z)=f(x, y, z) \cdot(0,0,1)
$$

where $f(0,0, z)=1$.
Given any $\lambda$-adapted almost complex structure $J$ on the symplectization $\mathbb{R} \times \mathbb{R}^{3}$, we write $M(t)$ to denote the 2 by 2 matrix that is the restriction of $J$ to the contact plane $\operatorname{ker} \lambda_{0,0, z}$. It satisfies

$$
M^{T} J_{0} M=J_{0} \quad \text { and } \quad-J_{0} M>0
$$

where $J_{0}$ is the standard 2 by 2 almost complex structure.
We now consider two $J$-holomorphic curves in the symplectization $\mathbb{R} \times \mathbb{R}^{3}$ that have boundary punctures that asymptotic to the Reeb chord $\eta$.

$$
U, V:[0, \infty) \times[0,1] \rightarrow \mathbb{R} \times \mathbb{R}^{3} \quad \text { with } \quad \lim _{s \rightarrow \infty} U(s,-)=\lim _{s \rightarrow \infty} V(s,-)=\eta
$$

In particular we consider how boundary punctures approach Reeb chords. We consider the case where there are two strip-like ends approaching the same chord. Assume that in the local conformal coordinates $(s, t)$ of the domain, the first end is parametrized as

$$
U(s, t)=(b(s, t), x(s, t), y(s, t), z(s, t)) .
$$

Using this parametrization, we define the local embedding $E:[R, \infty) \times D_{\epsilon} \times[0,1] \rightarrow$ $\mathbb{R} \times \mathbb{R}^{3}$ to be

$$
\begin{equation*}
E\left(s, h_{1}, h_{2}, t\right)=\left(b(s, t), x(x, t)+h_{1}, y(s, t)+h_{2}, z(s, t)-c_{1} f h_{1}-c_{2} f h_{2}\right) \tag{4.8.1}
\end{equation*}
$$

where $R>0$ is sufficiently large, $\epsilon>0$ is sufficiently small, and $D_{\epsilon}$ is the $\epsilon$ neighborhood of the origin. With the embedding $E$ understood, any nearby end could be viewed as a graph $E\left(s, \eta_{1}(s, t), \eta_{2}(s, t), t\right)$. In particular, we now parametrize the second end as

$$
\begin{align*}
V(s, t) & =E\left(s, \eta_{1}(s, t), \eta_{2}(s, t), t\right) \\
& =\left(b(s, t), x(s, t)+\eta_{1}(s, t), y(s, t)+\eta_{2}(s, t), z(s, t)-f c_{1} \eta_{1}(s, t)-f c_{2} \eta_{2}(s, t)\right) . \tag{4.8.2}
\end{align*}
$$

Note that by abuse of notation we have used $(s, t)$ to denote the coordinates for the domains of both ends. We also note here that $u(s, t)=(x(s, t), y(s, t), z(s, t))$ is the projection of $U$ onto the last three coordinates, and we shorthand the functions $f(u(s, t)), c_{1}(u(s, t))$ and $c_{2}(u(s, t))$ as $f, c_{1}$ and $c_{2}$ respectively. Under these parametrizations, the map $\left(\eta_{1}(s, t), \eta_{2}(s, t)\right)$ satisfies the following PDE:

$$
\begin{equation*}
\partial_{s}\binom{\eta_{1}}{\eta_{2}}+M(u(s, t)) \partial_{t}\binom{\eta_{1}}{\eta_{2}}+\Delta(s, t)\binom{\eta_{1}}{\eta_{2}}=0 \tag{4.8.3}
\end{equation*}
$$

where the matrix $\Delta(s, t)$ decays exponentially, i.e. there are positive constants $d$ and $M_{\beta}$ such that for any multi-index $\beta,\left|\partial^{\beta} \Delta\right| \leq M_{\beta} e^{-d s}$.

## The asymptotic formula

We now use the estimates from [2] to derive an asymptotic formula for $\left(\eta_{1}, \eta_{2}\right)$. The main difference in Abbas' approach from that of Hofer or Siefring is choosing a 1parameter family of norms dependent on $s$ which are equivalent to the ordinary $L^{2}$ norm on $[0,1]$. We first review this part of his construction, then we will state and prove the analogous lemmas in [2] that will culminate in the asymptotic formula.

We first recall properties of the matrix $M(u(s, t))$, which we will often abbreviate $M(s, t)$ for convenience. It converges to a matrix $M_{\infty}(t)$ as $s \rightarrow \infty$ at an exponential rate. Further, since $M(s, t)$ comes from restricting the ambient $\lambda$-adapted almost complex structure on the contact distribution, it satisfies the following algebraic properties:
(i) $M^{2}=-1$
(ii) $M^{T} J_{0} M=J_{0}$
(iii) $-J_{0} M>0$.

We consider the following family of inner products on $L^{2}\left([0,1], \mathbb{R}^{2}\right)$, let $v_{1}, v_{2} \in$ $L^{2}\left([0,1], \mathbb{R}^{2}\right):$

$$
\begin{equation*}
\left(v_{1}, v_{2}\right)_{s}:=\int_{0}^{1}\left\langle v_{1},-J_{0} M(s, t) v_{2}\right\rangle d t \tag{4.8.4}
\end{equation*}
$$

and it follows from the fact that $M(s, t)$ exponentially decays to $M_{\infty}(t)$ that for large values of $s$ this norm is uniformly equivalent to the ordinary $L^{2}$ norm.

We define

$$
\begin{equation*}
W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right):=\left\{v \in W^{1,2}\left([0,1], \mathbb{R}^{2}\right) \mid v(s, 0) \in \mathbb{R} \cdot(1,0), v(s, 1) \in \mathbb{R} \cdot(0,1)\right\} \tag{4.8.5}
\end{equation*}
$$

hence we can further define an inner product

$$
\begin{equation*}
\left(v_{1}, v_{2}\right)_{s, 1,2}:=\left(v_{1}, v_{2}\right)_{s}+\left(v_{1}^{\prime}, v_{2}^{\prime}\right)_{s} \tag{4.8.6}
\end{equation*}
$$

where ' denotes the derivative with respect to $s$.
We consider the trivial vector bundle $E:=\left[s_{0}, \infty\right) \times[0,1] \times \mathbb{R}^{2}$ with the fiber $\mathbb{R}^{2}$. We view maps $\zeta(s, t):\left[s_{0}, \infty\right) \times[0,1] \rightarrow \mathbb{R}^{2}$ as sections of $E$. With the above choice of metric, we choose a connection on $E$ by defining the following covariant derivatives:

$$
\begin{aligned}
\nabla_{s} \zeta & :=\partial_{s} \zeta-\frac{1}{2} M(s, t) \partial_{s} M(s, t) \cdot \zeta(s, t) \\
\nabla_{t} \zeta & :=\partial_{t} \zeta-\frac{1}{2} M(s, t) \partial_{t} M(s, t) \cdot \zeta(s, t)
\end{aligned}
$$

We abbreviate:

$$
\begin{aligned}
\Gamma_{1} & :=-\frac{1}{2} M(s, t) \partial_{s} M(s, t) \\
\Gamma_{2} & :=-\frac{1}{2} M(s, t) \partial_{t} M(s, t) .
\end{aligned}
$$

The following properties of this connection are proved in [2]:
Proposition 213. Let $X=a_{1} \partial_{s}+a_{2} \partial_{t}$, then we define $\nabla_{X}:=a_{1} \nabla_{s}+a_{2} \nabla_{t}$.
(i) For $u_{1}, u_{2} \in W^{1,2}\left(\left[s_{0}, \infty\right) \times[0,1], \mathbb{R}^{2}\right)$, we have:

$$
\begin{equation*}
\frac{d}{d s}\left(u_{1}, u_{2}\right)_{s}=\left(\nabla_{s} u_{1}, u_{2}\right)_{s}+\left(u_{1}, \nabla_{s} u_{2}\right)_{s} . \tag{4.8.7}
\end{equation*}
$$

(ii) If $B$ is a section in the endomorphism bundle of $E$, i.e. $B \in \Gamma(\operatorname{End}(E))=$ $\Gamma\left(\left[s_{0}, \infty\right) \times[0,1] \times \operatorname{End}\left(\mathbb{R}^{2}\right)\right)$, then we define:

$$
\begin{equation*}
\nabla_{s} B:=\partial_{s} B+\frac{1}{2}\left[B \cdot M(s, t) \partial_{s} M(s, t)-M(s, t) \partial_{s} M(s, t) \cdot B\right] \tag{4.8.8}
\end{equation*}
$$

This satisfies:

$$
\begin{equation*}
\nabla_{s}(B \cdot \zeta)=\nabla_{s} B \cdot \zeta+B \cdot \nabla_{s} \zeta . \tag{4.8.9}
\end{equation*}
$$

In particular we have

$$
\begin{equation*}
\nabla_{s} M(s, t)=0 \tag{4.8.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\partial_{t} \nabla_{s} \zeta-\nabla_{s} \partial_{t} \zeta=\partial_{t} \Gamma_{1} \cdot \zeta \tag{4.8.11}
\end{equation*}
$$

We also define the following family of unbounded self-adjoint operators:

$$
\begin{equation*}
A(s): W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right) \subset L^{2}\left([0,1], \mathbb{R}^{2}\right) \longrightarrow L^{2}\left([0,1], \mathbb{R}^{2}\right) \tag{4.8.12}
\end{equation*}
$$

as

$$
\begin{equation*}
A(s):=-M(s, t) \frac{d}{d t} \tag{4.8.13}
\end{equation*}
$$

We also use $A_{\infty}$ to denote the operator $-M_{\infty}(t) \frac{d}{d t}$. The following properties are established in Proposition 3.4 in [2].

Proposition 214. (i) $A(s)$ is self-adjoint on $\left(L^{2}\left([0,1], \mathbb{R}^{2}\right),(\cdot, \cdot)_{s}\right)$.
(ii) $\operatorname{Ker} A(s)=\{0\}$.
(iii) There exists $\delta>0$, independent of $s$, so that for all $s \in\left[s_{0}, \infty\right)$ ( $s_{0}$ some fixed number), and for all $\gamma \in W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right)$ we have

$$
\|A(s) \gamma\|_{s} \geq \delta\|\gamma\|_{s}
$$

where here and in what follows, we use $\|\cdot\|_{s}$ to denote the norm defined by $(\cdot, \cdot)_{s}$.

We also need the following proposition on Hilbert spaces, also stated and proved in [2, Theorem 3.7].

Proposition 215. Let $T: D(T) \in H \rightarrow H$ be a self-adjoint operator in a Hilbert space $H$. Let $A_{0}: H \rightarrow H$ be a linear symmetric, bounded operator. Let $\sigma(\cdot)$ denote the spectrum of an operator. Then

$$
\begin{align*}
\operatorname{dist}\left(\sigma(T), \sigma\left(T+A_{0}\right)\right) & :=\max \left\{\sup _{\lambda \in \sigma(T)} \operatorname{dist}\left(\lambda, \sigma\left(T+A_{0}\right)\right), \sup _{\lambda \in \sigma\left(T+A_{0}\right)} \operatorname{dist}(\lambda, \sigma(T))\right\}  \tag{4.8.14}\\
& \leq\left\|A_{0}\right\|_{\mathcal{L}(H)} \tag{4.8.15}
\end{align*}
$$

Assume further that the resolvent $\left(T-\lambda_{0}\right)^{-1}$ of $T$ exists and is compact for some $\lambda_{0} \notin \sigma(T)$, then $(T-\lambda)^{-1}$ exists and is compact for every $\lambda \notin \sigma(T)$, and $\sigma(T)$ consists of isolated eigenvalues $\left\{\mu_{k}\right\}_{k \in \mathbb{Z}}$ with finite multiplicities $\left\{m_{k}\right\}_{k \in \mathbb{Z}}$. If we assume $\sup _{k \in \mathbb{Z}} m_{k} \leq M<\infty$, and that for each $L>0$ there is a number $m_{T}(L) \in \mathbb{N}$ so that every interval $I \subset \mathbb{R}$ of length $L$ contains at most $m_{T}(L)$ points of $\sigma(T)$ (counted with multiplicity), then for each $L>0$ there is also a number $m_{T+A_{0}}(L)$ so that every interval $I \in \mathbb{R}$ of length $L$ contains at most $m_{T+A_{0}}(L)$ points of $\sigma\left(T+A_{0}\right)$.

To lighten the exposition we also isolate the following lemma from [2], whose proof depends on the above proposition.

Lemma 216. Consider the operators $A_{\infty}: W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right) \rightarrow L^{2}\left([0,1], \mathbb{R}^{2}\right)$ and $A(s): W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right) \rightarrow L^{2}\left([0,1], \mathbb{R}^{2}\right)$. We have:
(i) Given any $L>0$, there is some positive integer $m$ so that any interval in $\mathbb{R}$ of length $L$ contains at most $m$ eigenvalues of $A_{\infty}$ (up to multiplicity).
(ii) $\operatorname{dist}\left(\sigma(A(s)), \sigma\left(A_{\infty}\right)\right) \rightarrow 0$ as $s \rightarrow \infty$.
(iii) Let $I_{n}:=[-(n+1) L,-n L]$, then each $I_{n}$ contains at most $m$ points of $\sigma\left(A_{\infty}\right)$. There is a closed interval $J_{n} \subset I_{n}$ of length $2 d>0$ so that for $s>s_{0}$, the interval $J_{n}$ does not contain any element of $\sigma(A(s))$. In the following, we will write $J_{n}$ to be of the form $\left[r_{n}-d, r_{n}+d\right]$.

Let $\zeta(s, t):[0,1] \times \mathbb{R} \rightarrow \mathbb{R}^{2}$ be a solution to equation 4.8.3 that exponentially decays to zero as $s \rightarrow \infty$. Then we have (see [2] Theorem 3.6):

Proposition 217. If $\zeta$ does not vanish identically, then we have the following asymptotic formula:

$$
\zeta(s, t)=e^{\int_{s_{0}}^{s} \alpha(\tau) d \tau}(e(s, t)+r(s, t)),
$$

where
(i) $e(s, t)$ is an eigenvector of $A_{\infty}$ in $W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right)$ with eigenvalue $\lambda<0$.
(ii) $\alpha(s):\left[s_{0}, \infty\right) \times \mathbb{R}$ is a smooth function satisfying $\alpha(s) \rightarrow \lambda$ as $s \rightarrow \infty$.
(iii) $r(s, t):\left[s_{0}, \infty\right) \times[0,1] \rightarrow \mathbb{R}$ is a smooth function with

$$
\left|\partial^{\beta} r(s, t)\right| \rightarrow 0
$$

as $s \rightarrow \infty$. Here $\beta \in \mathbb{N}^{2}$ is some multi-index.
We break down the proof into a series of lemmas, as was done analogously in [2].
Lemma 218 (Lemma 3.9 in [2]). If the assumptions of the above proposition are satisfied, then $\alpha(s) \rightarrow \lambda$, where $\lambda$ is a negative eigenvalue of $A_{\infty}$.

Proof. We first assume $\zeta(s, t) \neq 0$ for all $s \geq s_{0}$. The alternative case is handled later. We start with some preliminary manipulations.
Step 0 Define

$$
\begin{equation*}
\alpha(s):=\frac{d / d s\|\zeta(s, t)\|_{s}^{2}}{2\|\zeta(s, t)\|_{s}^{2}} \tag{4.8.16}
\end{equation*}
$$

Then trivially we can write

$$
\frac{d}{d s}\|\zeta(s, t)\|_{s}^{2}=2 \alpha\|\zeta(s, t)\|_{s}^{2}
$$

from which we deduce

$$
\|\zeta\|_{s}^{2}=e^{2 \int_{s_{0}}^{s} \alpha d s^{\prime}}\left\|\zeta\left(s_{0}\right)\right\|_{s}^{2}
$$

Hence we need to establish the asymptotics of $\alpha$. We define

$$
\xi(s, t):=\frac{\zeta(s, t)}{\|\zeta(s, t)\|_{s}}
$$

Consequently

$$
\partial_{t} \xi=\frac{\partial_{t} \zeta}{\|\zeta\|_{s}}
$$

and

$$
\partial_{s} \xi=\frac{\partial_{s} \zeta}{\|\zeta\|_{s}}-\frac{\zeta}{2\|\zeta\|_{s}^{3}} \frac{d}{d s}\|\zeta\|_{s}^{2}
$$

Recall that $\zeta$ satisfies the equation 4.8.3, we have

$$
\partial_{s} \xi+\alpha \xi+M(s, t) \partial_{t} \xi+\Delta(s, t) \xi=0
$$

which we can further rewrite as

$$
\nabla_{s} \xi-\Gamma_{1} \xi+\alpha \xi+M(s, t) \partial_{t} \xi+\Delta(s, t) \xi=0
$$

Taking the $(\cdot, \cdot)_{s}$ innter product of the above with $\xi$, we obtain

$$
\alpha(s)=\left(\Gamma_{1} \xi, \xi\right)_{s}-\left(M \partial_{t} \xi, \xi\right)_{s}-(\Delta \xi, \xi)_{s}
$$

where we used $\left(\nabla_{s} \xi, \xi\right)_{s}=0$. Taking the $s$ derivative of both sides, we obtain

$$
\begin{aligned}
\alpha^{\prime}(s)= & -\left(\nabla_{s}\left(M(s, t) \partial_{t} \xi\right), \xi\right)_{s}-\left(M \partial_{t} \xi, \nabla_{s} \xi\right)_{s}+\left(\nabla_{s}\left(\Gamma_{1} \xi\right), \xi\right)_{s}+\left(\Gamma_{1} \xi, \nabla_{s} \xi\right) \\
& -\left(\nabla_{s}(\Delta(s, t) \xi, \xi)_{s}-\left(\Delta(s, t) \xi, \nabla_{s} \xi\right)_{s}\right. \\
= & -\left(M \partial_{t} \nabla_{s} \xi, \xi\right)_{s}+\left(M \partial_{t} \Gamma_{1} \xi, \xi\right)_{s}+\left(A \xi, \nabla_{s} \xi\right)_{s}+\left(\left(\nabla_{s} \Gamma_{1}\right) \xi, \xi\right)_{s}+2\left(\Gamma_{1} \xi, \nabla_{s} \xi\right)_{s} \\
& -\left(\left(\nabla_{s} \Delta(s, t) \cdot \xi, \xi\right)_{s}+\left(\Delta(s, t) \cdot \nabla_{s} \xi, \xi\right)_{s}+\left(\Delta(s, t) \xi, \nabla_{s} \xi\right)_{s}\right) \\
= & 2\left(\nabla_{s} \xi, A \xi\right)+2\left(\Gamma_{1} \xi, \nabla_{s} \xi\right)_{s}+\left(M \partial_{t} \Gamma_{1} \xi, \xi\right)_{s}+\left(\left(\nabla_{s} \Gamma_{1}\right) \xi, \xi\right)_{s} \\
& -\left(\left(\nabla_{s} \Delta(s, t) \cdot \xi, \xi\right)_{s}+\left(\Delta(s, t) \cdot \nabla_{s} \xi, \xi\right)_{s}+\left(\Delta(s, t) \xi, \nabla_{s} \xi\right)_{s}\right)
\end{aligned}
$$

where in the passage from first line to second line we used properties of $\nabla$, as well as the fact that

$$
\left(\Gamma_{1} \nabla_{s} \xi, \xi\right)_{s}=\left(\nabla_{s} \xi, \Gamma_{1} \xi\right)_{s}
$$

which follows from the definition of $(\cdot, \cdot)_{s}$ as follows:

$$
\begin{aligned}
\left(\Gamma_{1} \nabla_{s} \xi, \xi\right)_{s} & =\int_{0}^{1}\left\langle-\frac{1}{2} M \partial_{s} M \nabla_{s} \xi,-J_{0} M \xi\right\rangle d t \\
& =-\frac{1}{2} \int_{0}^{1}\left\langle\partial_{s} M \nabla_{s} \xi,-J_{0} \xi\right\rangle d t \\
& =\frac{1}{2} \int_{0}^{1}\left\langle\nabla_{s} \xi, \partial_{s} M^{T} J_{0} \xi\right\rangle d t \\
& =\frac{1}{2} \int_{0}^{1}\left\langle\nabla_{s} \xi,-\partial_{s}\left(M^{T} J_{0} M M\right) \xi\right\rangle d t \\
& =-\frac{1}{2} \int_{0}^{1}\left\langle\nabla_{s} \xi, J_{0} \partial_{s} M \xi\right\rangle d t \\
& =\int_{0}^{1}\left\langle\nabla_{s} \xi,-J_{0} M\left(-\frac{1}{2} M \partial_{s} M \xi\right)\right\rangle d t
\end{aligned}
$$

Now inserting

$$
A \xi=\nabla_{s} \xi-\Gamma_{1} \xi+\alpha \xi+\Delta(s, t) \xi
$$

into the expression for $\alpha^{\prime}$, we obtain:

$$
\begin{aligned}
\alpha^{\prime}(s)= & 2\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}-2\left(\nabla_{s} \xi, \Gamma_{1} \xi\right)+2\left(\nabla_{s} \xi, \Delta(s, t) \xi\right)+2\left(\nabla_{s} \xi, \Gamma_{1} \xi\right)_{s} \\
& +\left(M \partial_{t} \Gamma_{1} \xi, \xi\right)_{s}+\left(\left(\nabla_{s} \Gamma_{1}\right) \xi, \xi\right)_{s} \\
& -\left(\left(\nabla_{s} \Delta(s, t) \cdot \xi, \xi\right)_{s}+\left(\Delta(s, t) \cdot \nabla_{s} \xi, \xi\right)_{s}+\left(\Delta(s, t) \xi, \nabla_{s} \xi\right)_{s}\right) \\
= & 2\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}+\left(\nabla_{s} \xi, \Delta(s, t) \xi\right)_{s}+\left(M \partial_{t} \Gamma_{1} \xi, \xi\right)_{s}+\left(\left(\nabla_{s} \Gamma_{1}\right) \xi, \xi\right)_{s} \\
& -\left(\left(\nabla_{s} \Delta(s, t) \cdot \xi, \xi\right)_{s}+\left(\Delta(s, t) \cdot \nabla_{s} \xi, \xi\right)_{s}\right) .
\end{aligned}
$$

We observe the following bounds

$$
\left(\nabla_{s} \xi, \Delta(s, t) \xi\right)_{s} \leq \epsilon(s)\left\{\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}+(\xi, \xi)_{s}\right\} .
$$

where $\epsilon(s)$ is a function that decays exponentially to zero. Likewise

$$
\left|\left(\Delta(s, t) \cdot \nabla_{s} \xi, \xi\right)_{s}\right| \leq \epsilon(s)\left\{\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}+(\xi, \xi)_{s}\right\} .
$$

Similarly the terms $\left(M \partial_{t} \Gamma_{1} \xi, \xi\right)_{s},\left(\left(\nabla_{s} \Gamma_{1}\right) \xi, \xi\right)_{s}$, and $\left(\left(\nabla_{s} \Delta(s, t) \cdot \xi, \xi\right)_{s}\right.$ have absolute value uniformly upper bounded by $\epsilon(s)$ as well. Hence we arrive at the following inequality for $\alpha^{\prime}(s)$ :

$$
\begin{equation*}
\alpha^{\prime}(s) \geq(2-\epsilon(s))\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}-\epsilon(s) \tag{4.8.17}
\end{equation*}
$$

We note our precise definition of $\epsilon(s)$ may change from line to line, but it will always denote a function that exponentially decays to zero. With the above inequality established, from this point onward we can then repeat the arguments in [2].

Step 1. We first show that $\alpha(s)$ is bounded above. Suppose not, then we can find a sequence

$$
\left\{s_{k}\right\}_{k} \rightarrow \infty \quad \text { satisfying } \quad \alpha\left(s_{k}\right) \rightarrow \infty
$$

If we had $\alpha(s) \geq \eta$ for some $\eta>0$ for all $s$ large enough, then by definition of $\alpha$ this implies:

$$
\frac{d}{d s}\|\zeta\|_{s}^{2} \geq 2 \eta\|\zeta\|_{s}^{2}
$$

which would imply (assuming $\|\zeta\|_{s}$ is nonzero for large enough $s$ ) $\|\zeta\|_{s} \rightarrow \infty$ as $s \rightarrow \infty$, a contradiction.

The above argument gives us that for any $\eta>0$, we can find $s_{k}^{\prime} \rightarrow \infty$ so that $\alpha\left(s_{k}^{\prime}\right)<\eta$. Now choose $\eta \in(0, \delta)$, where $\delta$ is given by Proposition 214. Let $\widehat{s}_{k}$ be the smallest number satisfying $\widehat{s}_{k}>s_{k}$ and $\alpha\left(\widehat{s}_{k}\right)=\eta$. Observe this $\eta$ cannot be an eigenvalue of any $A(s)$, because for any $\gamma \in W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right)$, we have

$$
\|A(s) \gamma-\eta \gamma\| \geq\|A(s) \gamma\|-\eta\|\gamma\| \geq(\delta-\eta)\|\gamma\|>0
$$

Thus

$$
\begin{aligned}
\left\|\nabla_{s} \xi\left(\widehat{s}_{k}\right)\right\|_{\widehat{s}_{k}} & \geq\left\|A\left(\widehat{s}_{k}\right) \xi\left(\widehat{s}_{k}\right)-\eta \xi\left(\widehat{s}_{k}\right)\right\|_{\widehat{s}_{k}}-\epsilon(s)\left\|\xi\left(\widehat{s}_{k}\right)\right\|_{\widehat{s}_{k}} \\
& \geq(\delta-\eta)-\epsilon(s) \\
& \geq \tau
\end{aligned}
$$

for some $\tau>0$. Consequently

$$
\alpha^{\prime}\left(\widehat{s}_{k}\right) \geq\left(2-\epsilon\left(\widehat{s}_{k}\right)\right)\left\|\nabla_{s} \xi\left(\widehat{s}_{k}\right)\right\|_{\widehat{s}_{k}}^{2}-\epsilon\left(\widehat{s}_{k}\right) \geq \tau^{2}
$$

for large enough $k$. This is a contradiction because we picked $\widehat{s}_{k}$ to be the smallest number greater than $s_{k}$ such that $\alpha\left(\widehat{s}_{k}\right)=\eta$, and that $\alpha\left(s_{k}\right)>\eta$. The fact $\alpha^{\prime}\left(\widehat{s}_{k}\right)>0$ implies there is some number in $\left[s_{k}, \widehat{s}_{k}\right)$ such that $\alpha<\eta$, which is a contradiction.

Step 2. We next proceed to show $\alpha$ is also bounded from below. Assume not, then we can find a sequence $s_{n} \rightarrow \infty$ so that $\alpha\left(s_{n}\right)=r_{n}$ and $\alpha^{\prime}\left(s_{n}\right) \leq 0$, where $r_{n}$ and $d$ are given in Lemma 216. Then we consider

$$
A\left(s_{n}\right) \xi\left(s_{n}\right)=\nabla_{s} \xi\left(s_{n}\right)-\Gamma_{1} \xi\left(s_{n}\right)+\alpha \xi\left(s_{n}\right)+\Delta\left(s_{n}, t\right) \xi\left(s_{n}\right)
$$

where both sides are functions of $t$. We have

$$
\begin{aligned}
\left\|\nabla_{s} \xi\left(s_{n}\right)\right\|_{s_{n}} & =\left\|A \xi\left(s_{n}\right)-r_{n} \xi\left(s_{n}\right)+\Gamma_{1} \xi\left(s_{n}\right)-\Delta\left(s_{n}, t\right) \xi\left(s_{n}\right)\right\|_{s_{n}} \\
& \geq d-\epsilon\left(s_{n}\right) \\
& \geq d / 2
\end{aligned}
$$

for large enough $n$. Combining this with

$$
\begin{equation*}
\alpha^{\prime}(s) \geq(2-\epsilon(s))\left(\nabla_{s} \xi, \nabla_{s} \xi\right)_{s}-\epsilon(s) \tag{4.8.18}
\end{equation*}
$$

we conclude $\alpha^{\prime}\left(s_{n}\right) \geq d^{2} / 4$, which is contrary to our assumptions. Hence $\alpha$ is bounded from below as well.

Step 3. We observe we can always find a sequence $\left\{s_{k}\right\}$ so that $\left\|\nabla_{s} \xi\left(s_{k}\right)\right\|_{s_{k}} \rightarrow 0$. Suppose not, then we can find $\eta>0$ so that

$$
\left\|\nabla_{s} \xi\right\|_{s} \geq \eta
$$

for large enough $s$, which implies

$$
\alpha^{\prime}(s) \geq \eta^{2}
$$

contradicting the claim $\alpha(s)$ is bounded.
Step 4. Because $\alpha$ is bounded, we can find a subsequence of $\left\{s_{k}\right\}$, which we also denote by $\left\{s_{k}\right\}$ so that

$$
\alpha\left(s_{k}\right) \rightarrow \lambda .
$$

As in [2], we first show $\lambda \in \sigma\left(A_{\infty}\right)$. Suppose not, let $\epsilon>0$ be defined by

$$
\epsilon:=\operatorname{dist}\left(\lambda, \sigma\left(A_{\infty}\right)\right)
$$

Then for also large enough $s$ we have:

$$
\operatorname{dist}(\lambda, \sigma(A(s))) \geq \epsilon / 2
$$

Then as before we can compute:

$$
\begin{aligned}
\left\|\nabla_{s} \xi\left(s_{k}\right)\right\|_{s_{k}} & \geq\left\|A\left(s_{k}\right) \xi\left(s_{k}\right)-\alpha\left(s_{k}\right) \xi\left(s_{k}\right)\right\|_{s_{k}}-\epsilon\left(s_{k}\right)\|\xi(k, \cdot)\|_{s_{k}} \\
& \geq \epsilon / 4-\epsilon\left(s_{k}\right) \\
& \geq \epsilon / 8
\end{aligned}
$$

for large enough $k$. This contradicts $\left\|\nabla_{s} \xi\left(s_{k}\right)\right\|_{s} \rightarrow 0$, and hence $\lambda \in \sigma\left(A_{\infty}\right)$.
Step 5. Finally we show $\lim _{s \rightarrow \infty} \alpha(s)=\lambda$. Suppose not, then we can find $\left\{s_{k}^{\prime}\right\}$ so that

$$
\lim _{k \rightarrow \infty} \alpha\left(s_{k}^{\prime}\right) \rightarrow \mu
$$

Without loss of generality, let us assume $\mu<\lambda$. Then we can find some $\nu \in(\mu, \lambda)$ and $d>0$ so that

$$
\operatorname{dist}(\nu, \sigma(A(s))) \geq d
$$

for all $s$ sufficiently large. Then for any $\widehat{s}$ so that $\alpha(\widehat{s})=\nu$, we have:

$$
\left\|\nabla_{s} \xi(\widehat{s})\right\|_{\widehat{s}} \geq d / 2
$$

as before, but this implies $\alpha^{\prime}(\widehat{s}) \geq d^{2} / 4>0$, which implies for large enough $s$ we have $\alpha(s) \geq \nu$, which is a contradiction. A very similar argument can be applied in case $\mu>\lambda$. We have therefore $\alpha(s) \rightarrow \lambda$. We claim $\lambda<0$ because the norm of $\zeta$ exponentially decays to zero as $s \rightarrow \infty$.

Step 6 Finally we return to the case where for some $s^{*}$ we have $\left\|\zeta\left(s^{*}, \cdot\right)\right\|_{s^{*}}=0$, this would then imply $\zeta\left(s^{*}, t^{*}\right)=0$ for all $t^{*}$, and the Carleman similarity principle then implies $\xi$ is zero everywhere.

Lemma 219. With the same notation as above. Let $\beta=\left(\beta_{1}, \beta_{2}\right) \in \mathbb{N}^{2}$, and $j \in \mathbb{N}$. We have

$$
\begin{gathered}
\sup _{(s, t) \in\left[s_{0}, \infty\right) \times[0,1]}\left|\partial^{\beta} \xi(s, t)\right|<\infty \\
\sup _{s \in\left[s_{0}, \infty\right)}\left|\frac{d^{j} \alpha}{d s^{j}}(s)\right|<\infty
\end{gathered}
$$

Proof. The proof is almost verbatim the proof of lemma 3.10 in [2]. The only difference in our case is the appearance of the term $\Delta(s, t) \xi$, but in the notation of lemma 3.10 of Abbas this can be absorbed into $\widehat{\alpha} \xi$, and using the fact $\Delta(s, t)$ and its derivatives decay exponentially, we can recover bounds of derivatives of $\xi$ and $\alpha$.

Proposition 220. Let $E$ denote the eigenspace of the operator $A_{\infty}$ in $W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right)$ corresponding to eigenvalue $\lambda$, then

$$
\operatorname{dist}(\xi(s), E) \rightarrow 0
$$

as we take $s \rightarrow \infty$. The notion of distance is defined with respect to $W_{\Gamma}^{1,2}\left([0,1], \mathbb{R}^{2}\right)$.
Proof. The same proof as in Lemma 3.6 in [47].
Lemma 221. There exists $e \in E$ so that $\xi(s) \rightarrow e$ as $s \rightarrow \infty$.
Proof. The proof is the same as the proof of Lemma 3.12 in [2], as such we will only sketch the notable differences between our proof and theirs. As in their case for every sequence $\left\{s_{n}\right\}$ that converges to infinity, we can extract a subsequence $\left\{s_{n}^{\prime}\right\}$ so that $\xi\left(s_{n}^{\prime}\right) \rightarrow e \in E$. Hence it remains to show if we have sequences $\left\{s_{n}\right\}$ and $\left\{\tau_{n}\right\}$ so that

- $\xi\left(s_{n}\right) \rightarrow e \in E$,
- $\xi\left(\tau_{n}\right) \rightarrow e^{\prime} \in E$
then we have $e=e^{\prime}$. To this end, consider the inner product on $L^{2}\left([0,1], \mathbb{R}^{2}\right)$ given by:

$$
\left(u_{1}, u_{2}\right):=\int_{0}^{1}\left\langle u_{1}(t),-J_{0} M_{\infty} u_{2}(t)\right\rangle d t
$$

Let $P$ denote the orthogonal projection to $E$ with respect to the above inner product, and we define

$$
\widehat{\xi}:=P \xi .
$$

Then it follows as in the proof in [2] that:

- $A_{\infty} P \xi=P A_{\infty} \xi$.
- If we define $\epsilon^{\prime}(s):=A(s)-A_{\infty}$, we have the equation

$$
\partial_{s} \widehat{\xi}=(\lambda-\alpha(s)) \widehat{\xi}+P \epsilon^{\prime}(s) \xi-P \Delta \xi(s)
$$

- We have that both $\left\|P \epsilon^{\prime}(s) \xi\right\|$ and $\|P \Delta(s, t) \xi(s)\|$ are bounded above by $C e^{-\rho s}$.
- We have the norm bounds:

$$
\begin{gathered}
\|\xi(s)\| \rightarrow 1 \\
\|\widehat{\xi}-\xi(s)\| \rightarrow 0
\end{gathered}
$$

which in particular implies for large enough $s$ we have

$$
\|\widehat{\xi}\| \geq 1 / 2
$$

With the above we define

$$
\eta:=\frac{\widehat{\xi}(s, t)}{\|\widehat{\xi}(s, t)\|}
$$

Then we have the equation

$$
\begin{aligned}
\partial_{s} \eta & =\frac{\partial_{s} \widehat{\xi}}{\|\widehat{\xi}\|}-\frac{\left(\widehat{\xi}, \partial_{s} \widehat{\xi(s)}\right)}{\|\widehat{\xi}\|^{3}} \widehat{\xi} \\
& =\frac{(\lambda-\alpha(s)) \widehat{\xi}+P \epsilon^{\prime} \xi-P \Delta \xi}{\|\widehat{\xi}\|}-\frac{\left(\widehat{\xi},(\lambda-\alpha(s)) \widehat{\xi}+P \epsilon^{\prime} \xi-P \Delta \xi\right) \widehat{\xi}}{\|\widehat{\xi}\|^{3}} \\
& =\frac{P \epsilon^{\prime} \xi-P \Delta \xi}{\|\widehat{\xi}\|}-\frac{\left(\widehat{\xi}, P \epsilon^{\prime} \xi-P \Delta \xi\right)}{\|\widehat{\xi}\|^{3}} \widehat{\xi}
\end{aligned}
$$

From the above norm estimates we obtain

$$
\left|\partial_{s} \eta\right| \leq e^{-\rho s}
$$

where the above is the genuine absolute value. Then we have

$$
\begin{aligned}
\left|\eta\left(s_{n}\right)-\eta\left(\tau_{n}\right)\right| & \leq\left|\int_{\tau_{n}}^{s_{n}} \partial_{s} \eta d s\right| \\
& \leq\left|\int_{\tau_{n}}^{s_{n}} e^{-\rho s} d s\right| \\
& \rightarrow 0
\end{aligned}
$$

as $n \rightarrow \infty$. This concludes the proof of our lemma.
Proof of Proposition 217. With the previous lemmas assembled, the proof of Proposition 217 follows verbatim to that of Proposition 3.6 in [2]. We use Sobolev embedding and induction to show that all derivatives of $r(s, t)$ converge uniformly to zero. We use Arzela-Ascoli to show that $\alpha(s)$ approaches $\lambda$ in the $C^{\infty}$ norm.

We are now ready to apply the above estimates to analyze the difference of two ends asymptotic to the same Reeb orbit. Namely, we prove the following:

Proposition 222. For two different ends asymptotic to the same Reeb orbit with local coordinates:

$$
U(s, t)=(b(s, t), x(s, t), y(s, t), z(s, t))
$$

and
$V(s, t)=\left(b(s, t), x(s, t)+\eta_{1}(s, t), y(s, t)+\eta_{2}(s, t), z(s, t)-f c_{1} \eta_{1}(s, t)-f c_{2} \eta_{2}(s, t)\right)$.
There are reparametrizations $\phi_{u}, \phi_{v}:[R, \infty) \times[0,1] \rightarrow[R, \infty) \times[0,1]$ for sufficiently large $R$, such that

$$
\begin{aligned}
U\left(\phi_{u}(s, t)\right) & =\left(s, U_{0}(s, t), t\right), \\
V\left(\phi_{v}(s, t)\right) & =\left(s, V_{0}(s, t), t\right),
\end{aligned}
$$

and

$$
V_{0}(s, t)-U_{0}(s, t)=e^{\lambda s}(e(t)+r(s, t))
$$

where $e(t)$ is an eigenvector of the asymptotic operator with eigenvalue $\lambda$, and $r(s, t)$ decays exponentially to zero as $s \rightarrow \infty$.

Proof of Proposition 222. Fix a sufficiently large positive number $R$, and $\epsilon>0$ sufficiently small. Let $D_{\epsilon}$ be the $\epsilon$-disk around the origin in $\mathbb{R}^{2}$, and $E:[R, \infty) \times D_{\epsilon} \times$ $[0,1] \rightarrow \mathbb{R} \times \mathbb{R}^{3}$ be the embedding:

$$
\begin{equation*}
E\left(s, h_{1}, h_{2}, t\right)=\left(b(s, t), x(x, t)+h_{1}, y(s, t)+h_{2}, z(s, t)-c_{1} f h_{1}-c_{2} f h_{2}\right) \tag{4.8.19}
\end{equation*}
$$

where, as before, we shorthand $c_{1}(u(s, t)), c_{2}(u(s, t))$ and $f(u(s, t))$ as $c_{1}, c_{2}$ and $f$ respectively. Now the first end $U$ is parameterized as:

$$
U(s, t)=E(s, 0,0, t)
$$

and since any different end that is asymptotic to the same Reeb chord can be viewed as an embedded surface in the image of $E, V$ can be parametrized as

$$
V(s, t)=E\left(s, \eta_{1}(s, t), \eta_{2}(s, t), t\right)
$$

Since both ends limit to the Reeb chord exponentially, there are reparametrizations $\phi_{u}, \phi_{v}:[R, \infty) \times[0,1] \rightarrow[R, \infty) \times S^{1}$ such that

$$
U\left(\phi_{u}(s, t)\right)=\left(s, U_{0}(s, t), t\right)
$$

$$
V\left(\phi_{v}(s, t)\right)=\left(s, V_{0}(s, t), t\right) .
$$

Let $\lambda$ be the negative eigenvalue that appears in the asymptotic expansion of

$$
\eta(s, t)=\left(\eta_{1}(s, t), \eta_{2}(s, t)\right)
$$

provided by Proposition 217, and $e(t)$ be the corresponding eigenvector. We have:

$$
\phi_{u}(s, t)-\phi_{v}(s, t)=o_{\infty}(\lambda)
$$

Here we use the same notation as in [108]. Namely, a function $f:[R, \infty) \times[0,1] \rightarrow$ $\mathbb{R}^{2}=o_{\infty}(\lambda)$ if and only if there are constants $\delta>0$ and $M_{\beta}>0$ for every multi-index $\beta$, such that

$$
\left|\partial^{\beta}\left(e^{\lambda s} f\right)\right| \leq M_{\beta} e^{-\delta s}
$$

The reason is that, by definition
$\phi_{u}^{-1}(s, t)-\phi_{v}^{-1}(s, t)=\left(0,-c_{1}(u(s, t)) f(u(s, t)) \eta_{1}(s, t)-c_{2}(u(s, t)) f(u(s, t)) \eta_{2}(s, t)\right)$,
therefore $\phi_{u}^{-1}(s, t)-\phi_{v}^{-1}(s, t)=o_{\infty}(\lambda)$, and hence $\phi_{u}(s, t)-\phi_{v}(s, t)=o_{\infty}(\lambda)$. Now let $\pi: \mathbb{R}^{4} \rightarrow \mathbb{R}^{2}$ be the projection to the second and third coordinates. We have:

$$
\begin{aligned}
V_{0}(s, t)-U_{0}(s, t) & =\pi\left(V\left(\phi_{v}(s, t)\right)-U\left(\phi_{v}(s, t)\right)\right) \\
& =\pi\left(E\left(\phi_{v}(s, t), \eta\left(\phi_{v}(s, t)\right)\right)-E\left(\phi_{u}(s, t), 0,0\right)\right) \\
& =\pi\left(E\left(\phi_{u}(s, t), \eta\left(\phi_{u}(s, t)\right)\right)-E\left(\phi_{u}(s, t), 0,0\right)\right)+o_{\infty}(\lambda) \\
& =\eta\left(\phi_{u}(s, t)\right)+o_{\infty}(\lambda) \\
& =e^{\lambda s} e(t)+o_{\infty}(\lambda)
\end{aligned}
$$

The above discussion finishes the proof of Theorem 184. To see this fact, choose an Riemannian metric on $Y$ so that near the Reeb chord $\eta$ with the chosen local coordinates, the metric agrees with the standard Riemannian metric on $\mathbb{R}^{3}$. Extend this metric to an $\mathbb{R}$-invariant metric on $\mathbb{R} \times Y$, and we see by definition that the functions $U_{0}$ and $V_{0}$ are asymptotic representatives of the two ends (see Definition 2.1 of [108]). The above Proposition shows that there is a section $r(s, t)$ of $\eta^{*} \xi$ that exponentially decays to zero as $s$ tends to $\infty$, so that under the chosen local coordinates the two asymptotic representatives satisfy

$$
U_{0}(s, t)-V_{0}(s, t)=e^{\lambda s}(e(t)+r(s, t))
$$

Having established the above proposition, the exact same argument as in Section 4 in [108] proves the following:

Proposition 223. (Compare Theorem 2.4 in [108]) Let $\gamma$ be a Reeb chord connecting two Legendrians $\Gamma_{0}$ and $\Gamma_{1}$, with action 1, and fix a trivialization $\tau$ of $\gamma^{*} \xi$. Let $\left\{u_{i}\right\}_{i=1}^{n}$ be a collection of pseudo-holomorphic curves with $\gamma$ as a positive (resp. negative) end. Then there exist a neighborhood $U$ of $\gamma$, a smooth embedding $\Phi: \mathbb{R} \times U \rightarrow$ $\mathbb{R} \times \mathbb{R}^{2} \times[0,1]$ with the property

$$
\Phi(s, \gamma(t))=(s, 0,0, t),
$$

proper reparametrizations $\psi_{i}:[R, \infty) \times[0,1] \rightarrow \mathbb{R} \times[0,1]$ asymptotic to the identity, and positive integers $N_{i}$, such that near the asymptotic ends

$$
\Phi \circ u_{i} \circ \psi_{i}(s, t)=\left(s, t, \sum_{k=1}^{N_{i}} e^{\lambda_{i, j} s} e_{i, j}(t)\right)
$$

where $\lambda_{i, j}$ are negative (resp. positive) eigenvalues of the asymptotic operator associated to $\gamma$ and $\tau$, and $e_{i, j}$ are eigenvectors with eigenvalue $\lambda_{i, j}$.

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[^0]:    ${ }^{1}$ Indeed, we may not have any roots of unity as eigenvalues of the linearized return map since we may consider covers of Reeb orbits.

[^1]:    ${ }^{2}$ Our asymptotic operator is defined with an opposite sign from others in the literature, e.g. [46].

[^2]:    ${ }^{1}$ It is a technical point that we cannot guarantee a priori that the moduli space of index 2 currents is compact. The truncation may be viewed as a replacement for the compactification.

[^3]:    ${ }^{2}$ It is a technical point that we cannot guarantee a priori that the moduli space of index 2 currents is compact. The truncation may be viewed as a replacement for the compactification.

[^4]:    ${ }^{3}$ However we must impose more restrictions if we want the differential to square to zero, see Section 4.7 for details.

[^5]:    ${ }^{4}$ Note that in the general case, the isomorphism between $H M_{\bullet}$ and $P F H_{\bullet}$ uses variants of both Floer groups with appropriate twisted Novikov coefficients.

