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### Effects of Magnetic Fields on HPGe Tracking Detectors

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### Abstract

We present a study of magnetic fields effects on the position resolution and energy response of hyper-pure germanium detectors. Our results provide realistic estimates of the potential impact on the resolving power of tracking-arrays from

are located close to large spectrometers where the magnetic fringe field may not be negligible, in some cases up to  $\approx 0.1$  T. For large crystal volumes, the deflection of the charge carriers in the magnetic field may result in deviations of the derived interaction point positions, a larger rise times of the signals due to longer drift paths and/or enhanced trapping and de-trapping, and changes in their

$$\frac{d\boldsymbol{r}(t)}{dt} = q\left(\mu\boldsymbol{E} + \mu_H \frac{d\boldsymbol{r}(t)}{dt} \times \boldsymbol{B}\right)$$
(1)

where  $\mu$  is the mobility and  $\mu_H$  is the Hall mobility. From a common initial position, the separated trajectories of

<sup>&</sup>lt;sup>1</sup>In a relevant work [13], the performance of silicon drift detectors in a magnetic field was studied.

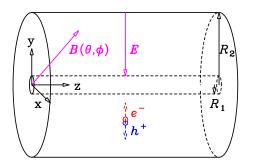


Figure 1: (Color online) Schematic illustration of the coaxial detector, the adopted axes system, and the  $\boldsymbol{E}$  and  $\boldsymbol{B}$  fields under which the electrons  $(e^-)$  and holes  $(h^+)$  drift in the Ge volume.  $\boldsymbol{E}$  is in the radial direction, while the orientation of  $\boldsymbol{B}$  is defined by the angles  $(\theta, \phi)$  with respect to the (x,y,z) reference frame.

the electron and the hole were solved numerically using their proper values of charge q and mobilities [18]. It is assumed that the electric field direction causes the holes to drift to the outer electrode and the electrons to the inner, as it happens in a typical n-type Ge detector. From the trajectories, a number of physical quantities can be derived. The key parameters, central to the understanding of the magnetic field effects on a tracking detector, are the position deviations and the change in drift times. The distance between the end positions of the trajectory with and without the magnetic filed gives the position deviation, and the difference in the drift time is similarly obtained. The pulse shape, resulting from the combined contributions from electrons and holes, is calculated using the weighting potential V(r) and the positions determined from their respective trajectories [19, 20]. For the outer electrode we have:

$$V(r) = \frac{\ln\left(\frac{r}{R_1}\right)}{\ln\left(\frac{R_2}{R_1}\right)} \tag{2}$$

Typical values of the input parameters we used are based on the average properties of GRETA detectors [8] and are given in Table 1.

To estimate the effects on the energy resolution and the shift of the full energy peak, charge trapping distances were introduced for both electrons and holes ( $\lambda_e$  and  $\lambda_h$ ). The charge attenuation is modeled after Refs. [21, 22]:

$$q_e = q_{0e} exp\left(-\frac{\Delta r_e}{\lambda_e}\right) \tag{3}$$

and

$$q_h = q_{0h} exp\left(-\frac{\Delta r_h}{\lambda_h}\right) \tag{4}$$

As the length  $\Delta r$  of the trajectory increase with the magnetic field, charge loss due to trapping increases and cause larger peak shift and worse energy resolution.

Table 1: Input parameters used in our calculations

Parameter	Value
$R_1$	$0.5~{\rm cm}$
$R_2$	$4 \mathrm{cm}$
E	1.46  kV/cm
$\mu_e$	$0.77 \text{ m}^2/\text{V/s}$
$\mu_h$	$0.62 \text{ m}^2/\text{V/s}$
$\mu_{He}$	$3\pi/8\mu_e$
$\mu_{Hh}$	$3\pi/8\mu_h$

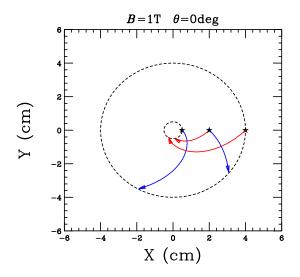


Figure 2: (Color online) Example of  $e^-$  (red) and  $h^+$  (blue) drift trajectories calculated for a 1 T longitudinal magnetic field. Three initial interaction points (stars) are shown.

### 3. Results and Discussions

Under the inference of both electric and magnetic fields and following from the solution of Eq. 1, a typical drift trajectory of a charge carrier in a HPGe detector is not a straight line. We now present some examples of our calculations at a magnetic field of 1 T.

Figure 2 shows the electron (red) and hole (blue) trajectories, corresponding to three initial positions. In this case, the magnetic field is parallel to the z-axis,  $\theta = 0 \deg$ , and thus perpendicular to the radial electric field. The trajectories, having spiral shapes, are confined in the x-y plane.

The example in Figure 3 corresponds to the case when the magnetic field is in the y-direction, perpendicular to both the z- and x-axis ( $\theta = 90 \deg, \phi = 90 \deg$ ). The position vector of the starting point of the charge carrier is on the x-direction. The trajectories are in the x-z plane and this is the only situation for a coaxial detector where the they are straight lines. In general, given the positioning of the array with respect to the spectrometer, the magnetic field will have an arbitrary direction ( $\theta, \phi$ ) relative to the

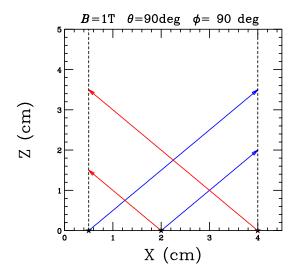


Figure 3: (Color online) Similar to Fig. 2 but for a 1 T magnetic field along the y-axis, with the three initial interaction points (stars) indicated. In these cases the carriers drift in the xz plane.

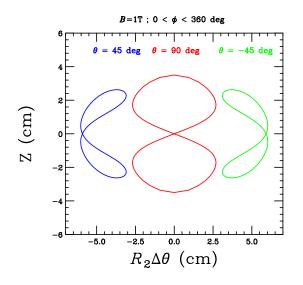


Figure 4: (Color online) Examples of the curves traced by the end points of the  $h^+$  trajectories, projected on the cylindrical outer surface of the detector. The three conditions are indicated in the plot. At zero field the end point is at (0,0).

axis of a particular detector and the trajectories will be more complicated.

### 3.1. Position deviations

Having calculated the carriers trajectories, we can now analyze the distribution of the end point positions. In Figure 4 we show the example of a hole trajectory on the detector outer surface, for various three magnetic field directions. The end points trace a 2-dimensional curve in the axial (z) and tangential (x-y) direction. The deviation of the end point position under magnetic field from that without magnetic field causes a two-dimensional position error in the tracking detector. Figure 5 shows the distri-

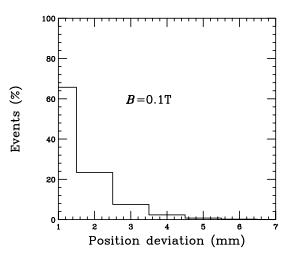


Figure 5: Calculated distribution of the position deviations for  $h^+$  at a 0.1 T field.

Table 2: Average position deviation and increase in drift time as functions of  ${\cal B}$ 

$B(\mathbf{T})$	Deviation (mm)	Drif time increase $(\%)$
0	0	0
0.1	0.9	0.3
0.2	1.8	1.9
0.3	2.8	3.5

bution of position deviations on the outer surface for the hole trajectories for all the magnetic field directions and all starting positions with a magnetic field strength of 0.1 T. In this case the deviation averaged over all events is  $\approx$  1 mm, and about 65% of the deviation is less than 1 mm. The average position deviation increases linearly with the magnetic field strength as seen in Table 2.

In addition to the two-dimensional positional sensitivity from the segmentation, tracking detectors use the drift time information for position determination in the drift direction. As the magnetic field increases the drift path length this introduces a position error in the drift direction. The example in Figure 6 shows the drift path length as functions of the magnetic field directions for B = 1T. The drift time deviations, averaged over the field orientations, are also given in Table 2 for several magnetic field strengths. The longer drift time will result in a longer raise time and change the pulse shape. The effect is relatively small at low field and increases approximately quadratically at high field

### 3.2. Energy resolution and peak shift

As it follows from Eqs. 3 and 4, charge losses due to trapping will increase as the path length increases, resulting in the reduction of the pulse height and thus peak en-

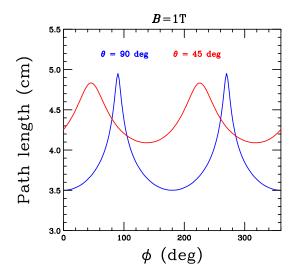


Figure 6: (Color online) Path lengths as a function of the magnetic field azimuthal orientation for two cases of the polar angle as indicated.

ergy shift. In addition, due to the variation of path length depending on the starting position and the relative angle between the drift direction and the field direction, the peak width will increase and degrade the energy resolution.

### 3.3. Comparison with experimental data

In order to benchmark our calculations, we contrast our results with selected data from Refs. [14, 15, 17]. First we consider a typical example of the effect of the magnetic field on the pulse shape [17], shown in Fig. 7. While a fully quantitative comparison may not be possible given that we do not know the full details of the HPGe detector used in that investigation, such as the depletion voltage and bias voltage, it is clear that the effects are well reproduced. The time response of the system was adjusted to reproduce the zero field pulse.

Similarly, when comparing the measured rise-time distributions in Ref. [14] with our calculations in Fig. 8, the impact of the magnetic field appears to be captured in the model assumptions. In this case, the distributions were aligned to a common centroid.

Finally, we show in Fig. 9 the calculated dependence of the peak shift (negative) and the energy resolution (FWHM) for a 1.33 MeV <sup>60</sup>Co line as a function of magnetic field strength together with data from Refs. [14, 15, 17]. Here, the initial FWHM was adjusted to be consistent with the measurements. Charge trapping distances (see Eqs. 3 and 4),  $\lambda_e = 6.0$  m and  $\lambda_h = 2.1$  m, were used to reproduce the results from Ref. [17], while 4.5 m and 1.6 m respectively were used for the results from Refs. [14, 15]. Once again our results are in good agreement with the empirical observations.

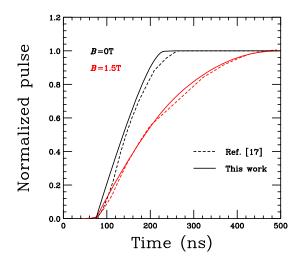


Figure 7: (Color online) Comparison of calculated pulse shapes (solid lines) to the measurements in Ref. [17] (dashed lines) for B = 0 (black) and B = 1.5 T (red). The calculations were normalized to qualitatively reproduce the zero field data.

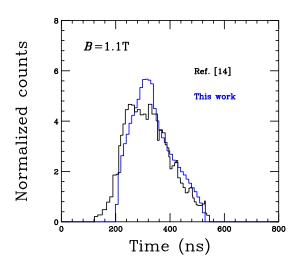


Figure 8: (Color online) Similar to Fig. 7, comparing the calculated rise-time distribution (blue) with the data in Ref. [14] (black). The centroids of the distributions were aligned for display purposes.

### 4. Conclusions

In this work we have studied the effects of magnetic field on the position resolution and energy response of HPGe tracking detectors, such as those used in GRETA. We believe our study provides a simple and robust framework to estimate the potential impact of (fringe) magnetic fields on the resolving power of tracking arrays, when operating together with large spectrometers. This situation is likely to be encountered in a large number of experiments running at rare isotope facilities.

The major effect is in the deviation of the transverse (to the electric field direction) position at the end point of the trajectories due to the deflection of the electron and hole

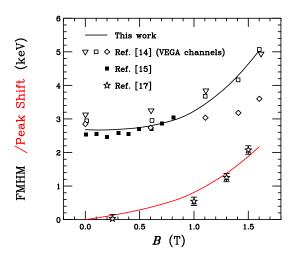


Figure 9: (Color online) Energy response, FWHM (black line) and peak shift (red) as a function of B compared to results of Refs. [14, 15, 17]. Note that in our calculations the intrinsic FWHM at B=0 was adjusted to agree with the data.

paths. In a field of  $\approx 0.1$  T this deviation, averaged over all magnet field orientations and interaction points positions, is  $\approx 1$ mm with 65% of the cases less than 1 mm. In the longitudinal direction the averaged drift time increases by  $\approx 0.3\%$ , which equals to a 0.1 mm deviation. Thus, if no corrective action is taken in the signal-decomposition and tracking data analysis procedures, to limit the deviation below the intrinsic position resolution of 2 mm (RMS) [8, 9] the maximum allowed field strength is  $\approx 0.1$  T. At field strengths higher than 1 T, the degradation of the energy response (resolution and peak shift) becomes observable. This effect is caused by the increase in charge trapping due to the lengthening of the trajectory.

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