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Potential Control of Oxygen Non-Stoichiometry in Cerium Oxide and Phase Transition Away from Equilibrium

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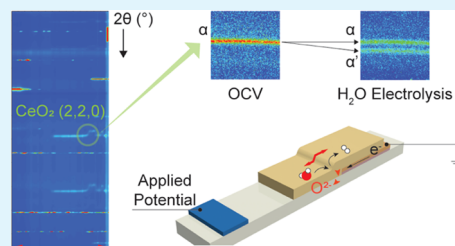
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Supporting Information

ABSTRACT: Cerium oxide (ceria, CeO_2) is a technologically important material for energy conversion applications. Its activities strongly depend on redox states and oxygen vacancy concentration. Understanding the functionality of chemical active species and behavior of oxygen vacancy during operation, especially in high-temperature solid-state electrochemical cells, is the key to advance future material design. Herein, the structure evolution of ceria is spatially resolved using bulk-sensitive operando X-ray diffraction and spectroscopy techniques. During water electrolysis, ceria undergoes reduction, and its oxygen non-stoichiometry shows a dependence on the electrochemical current. Cerium local bonding environments vary concurrently to accommodate oxygen vacancy formation, resulting in changes in Ce–O coordination number and $\text{Ce}^{3+}/\text{Ce}^{4+}$ redox couple. When reduced enough, a crystallographic phase transition occurs from α to an α' phase with more oxygen vacancies. Nevertheless, the transition behavior is intriguingly different from the one predicted in the standard phase diagram of ceria. This paper demonstrates a feasible means to control oxygen non-stoichiometry in ceria via electrochemical potential. It also sheds light on the mechanism of phase transitions induced by electrochemical potential. For electrochemical systems, effects from a large-scale electrical environment should be taken into consideration, besides effective oxygen partial pressure and temperature.

KEYWORDS: ceria, oxygen vacancy, phase transition, operando, X-ray diffraction



INTRODUCTION

Cerium oxide (ceria, CeO_2) is a technologically important and fundamentally interesting material because of its highly tunable redox chemical properties. It has been widely used in three-way catalysts,^{1,2} energy conversion devices,^{3–6} oxygen permeation membranes,⁷ and also in biotechnology and medicine.^{8,9} These applications are mostly related to the rapid formation/annihilation of oxygen vacancy ($\text{V}_\text{O}^{\bullet\bullet}$) and the strong impact of oxygen vacancy on the chemical and physical properties of ceria.¹⁰ Therefore, the ability to monitor, modify, and precisely control oxygen vacancy is a key prerequisite for being able to control ceria properties and design ceria-based materials in a rational manner.

Ceria has a robust fluorite structure that can accommodate a large oxygen non-stoichiometry over a wide temperature and oxygen partial pressure range as CeO_{2-x} .^{10,11} This feature enables the storage and supply of oxygen in ceria and ceria-based materials, which is important for catalytic applications.^{2,12} Upon oxygen release, the reduction of Ce^{4+} to Ce^{3+} and the incorporation of oxygen vacancies in its fluorite-type structure result in a volume expansion of the material.¹³ This is well known as chemical expansion,¹³ with a linear relationship between the cell parameter and oxygen non-stoichiometry x . This phenomenon presents a problem in mechanical engineering because of induced material strain and structural deflection. However, on the other hand, it is of particular interest for

mechanical sensing and actuation applications.¹⁴ For instance, one could monitor oxygen partial pressure from expansion/contraction of the structure; one could also control such oxide expansion/contraction for material actuation either by electrical tuning of oxygen vacancy concentration or by adjusting the gas atmosphere. Recently, Swallow et al.¹⁴ reported “breathing” response of a model material $\text{Pr}_x\text{Ce}_{1-x}\text{O}_{2-\delta}$ with strain $>0.1\%$ through electrochemically pumping of oxygen with electrical potential <0.1 V. This work noted a mechanism that couples electrical stimulus to mechanical properties via material defect chemistry, which leads to a new design path for electromechanical actuation with non-stoichiometric oxide films. In a recent review, Li and Chueh¹⁵ discussed an ion insertion process by chemical and electrochemical stimuli in many solid oxides. In addition to this, change in oxygen non-stoichiometry results also in distinct electrical properties^{15–18} and more fundamentally in phase transition, leading to a complex-associated phase diagram.

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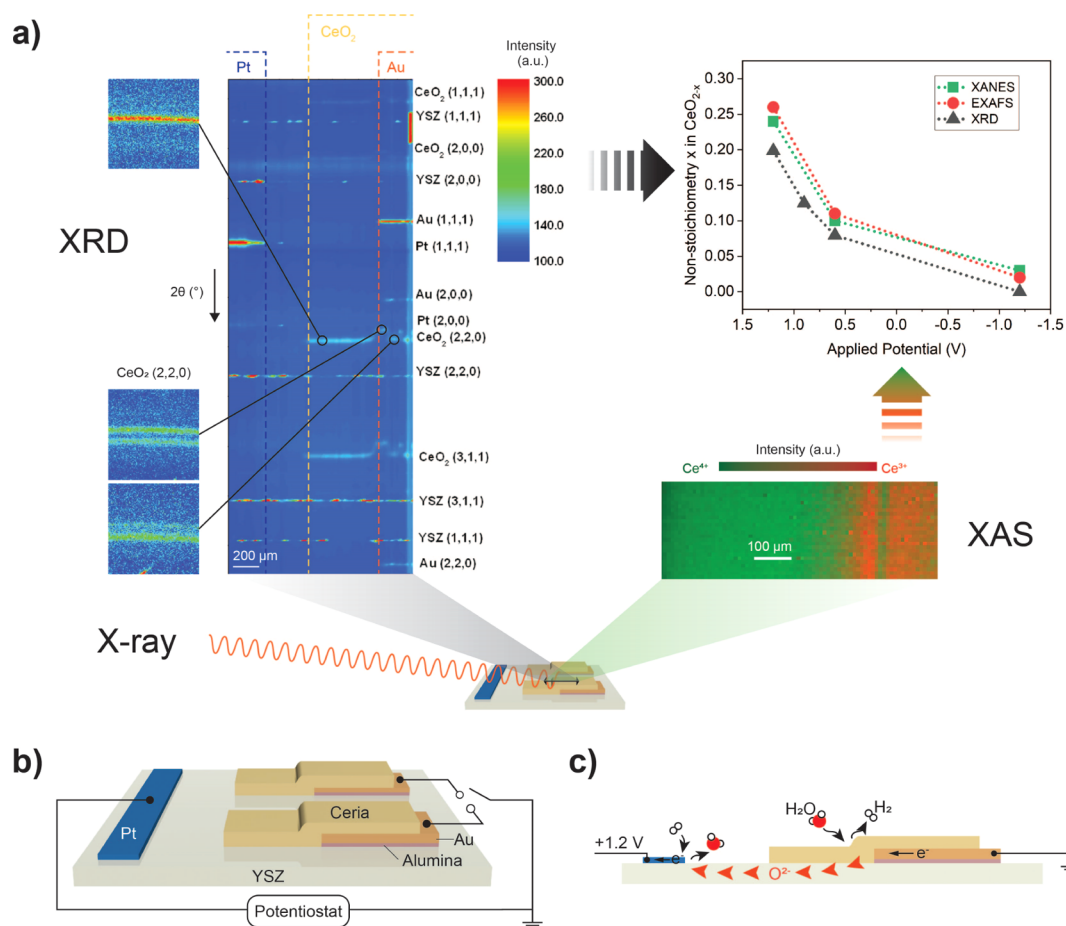


Figure 1. Operando behaviors of ceria are revealed by a suite of synchrotron-based X-ray techniques. (a) XRD and XAS are employed to study the bulk properties of a ceria electrode in a model SOC. XRD patterns are shown along a line crossing the SOC under the steady state of operation. Diffraction reflections of CeO_2 , Au, and YSZ are indicated. XAS mapping is obtained in fluorescence mode. Intensity of the contour plot represents the relative concentration of Ce^{3+} (red) and Ce^{4+} (green). Figure inset gives the extracted x values from XRD, X-ray absorption near edge structure, and extended X-ray absorption fine structure as a function of potential at 611 °C. (b) SOC consists of a 300 nm Pt CE (blue), 300 nm Au pad current collectors (orange) on top of 30 nm alumina films (red), and 1000 nm ceria WEs (yellow) patterned onto a polycrystalline YSZ substrate (light grey). During operation, the ceria WEs are grounded and negatively biased relative to the Pt CE. This schematic drawing is not to scale. (c) SOC is heated above 500 °C in a gas mixture of H_2 and H_2O . The negative potential on ceria drives H_2O electrolysis on the ceria surface, while the positive potential promotes H_2 electro-oxidation at the Pt side.

65 With a suite of synchrotron-based X-ray techniques, we aim
 66 to gain a more comprehensive picture of oxygen vacancy
 67 behavior from the surface through the bulk structure and verify
 68 the mechanism of potential-induced phase transition by using
 69 ceria as a model system. Such fundamental understanding is
 70 crucial for heterogeneous catalytic processes and especially for
 71 actuation applications where it is desirable to control the
 72 material properties by applied potentials. Our previous
 73 ambient pressure X-ray photoelectron spectroscopy (APXPS)
 74 studies have revealed an extended electrochemical surface
 75 active region beyond the three-phase boundary.^{19,20} In this
 76 paper, we probe the ceria electrode from the surface to the
 77 bulk. An external potential was applied to a solid oxide
 78 electrochemical cell (SOC) consisting of a ceria film as the
 79 working electrode (WE), an yttria-stabilized zirconia (YSZ)
 80 electrolyte, and a platinum counter electrode (CE). Our model
 81 SOC is designed such that the oxygen vacancy (and ion)
 82 transport is confined within a certain region, allowing one to
 83 discern changes in structure and vacancy concentration as a
 84 function of current and temperature.^{19–22} We used operando
 85 X-ray diffraction (XRD) and X-ray absorption spectroscopy
 86 (XAS) (Figure 1a) to examine the structural properties of

ceria, with gas pressure and temperature systematically 87
 controlled. Probing the vacancy distribution through ceria 88
 bulk, we obtained a direct correlation of oxygen vacancy 89
 concentration with the electrochemical current. In addition, 90
 our results clearly show that the non-stoichiometry x in 91
 CeO_{2-x} and associated phase transition in the relevant non- 92
 stoichiometry range can be tuned by controlling the electrical 93
 potential. 94

This work first demonstrates control of non-stoichiometry 95
 through electrochemical potential in functional oxides with a 96
 CeO_{2-x} model system. This electrochemical method is an 97
 easily accessible and practical means to tune electrical and 98
 mechanical properties of solid oxide materials. It also provides 99
 new possibility in heterogeneous catalysis, especially at 100
 elevated temperatures. Beyond changes in temperature, it 101
 adds an extra dimension to introduce/remove particular 102
 catalytic active sites by controlling the amount of oxygen 103
 vacancies. Second, this work also reveals a detailed 104
 mechanistic understanding of the phase change induced by 105
 electrochemical potential. It has been reported that the 106
 effective oxygen partial pressure as changed by electrical 107
 stimulus has the same effect on the phase transition induced by 108

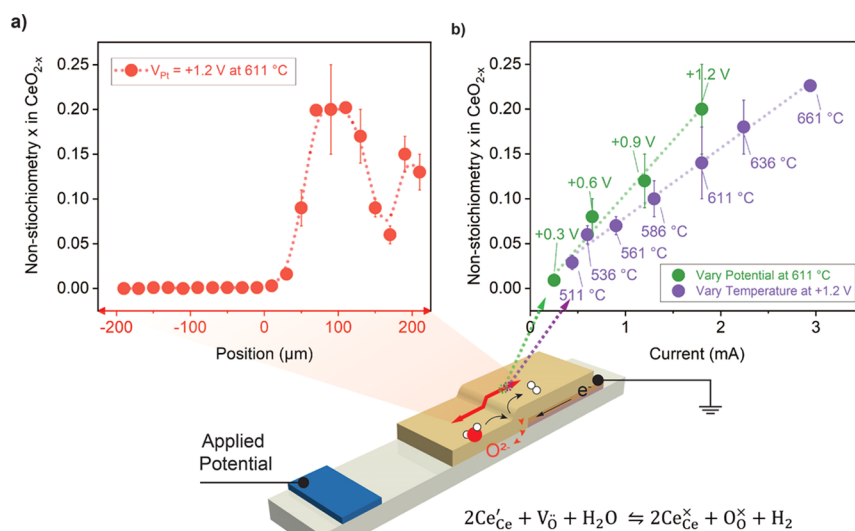


Figure 2. Ceria shows bulk reduction during operation, and the oxygen non-stoichiometry x has a linear dependence on cell current. (a) Oxygen non-stoichiometry x on the ceria electrode across the SOC at 611 °C WT and +1.2 V applied potential on Pt. These values are extracted from XRD measurement obtained along the red line. The WT values have been calibrated, and the +1.2 V applied potential denotes cathodic reaction on the ceria electrode. (b) Change in oxygen non-stoichiometry x as a function of applied cell potentials (denoted in green) and as a function of temperature (denoted in purple) on the ceria electrode. The linear fits are also displayed for reference.

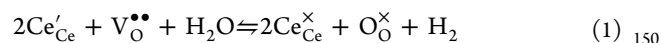
109 the actual change of oxygen partial pressure, following the
 110 Nernst equation.²³ However, a clear different phase transition
 111 behavior is observed in this study when compared with the
 112 standard ceria phase diagram.²⁴ New information is provided
 113 from this investigation to clarify the phase transition
 114 mechanism via potential control. Finally, there is a general
 115 implication on the relevance of ceria and similar fluorite-
 116 structure oxides for energy conversion and storage applica-
 117 tions. Under realistic non-open circuit voltage conditions, a
 118 large deviation in the structure and the chemical state can be
 119 induced by a finite change in electrical potential. The
 120 electrochemical activity, ionic and electronic transport proper-
 121 ties, and stability of these oxide materials can be largely altered
 122 as a consequence. Therefore, it is important to assess such
 123 changes under in situ and operando environments and to
 124 identify the threshold potential that leads to phase transitions.

125 ■ RESULTS AND DISCUSSION

126 **Operating the Ceria-Based SOC.** To perform simulta-
 127 neous electrochemical evaluation and X-ray powder diffrac-
 128 tion/absorption experiments, we employed a specially
 129 designed model SOC. The single-sided SOC consists of
 130 polycrystalline YSZ electrolyte, ceria WEs with Au pad current
 131 collectors, and Pt CE (Figure 1b). All the cell components are
 132 positioned on the same side to facilitate in situ X-ray
 133 measurements. Dense ceria films (1000 nm-thick) are
 134 deposited on top of Au pads in a way that only the ceria
 135 edge close to the Pt CE has direct contact with the YSZ
 136 electrolyte. Underneath the Au pads, insulating 30 nm thick
 137 alumina layers are deposited to block direct ionic transport
 138 between the YSZ electrolyte and the Au current collectors,
 139 thus allowing well-defined current flow regions. A detailed
 140 description of the cell geometry can be found in previous
 141 publications.^{19,20}

142 The SOC was mounted on top of an Anton Paar DHS900
 143 heater and heated to a particular working temperature (WT)
 144 (from 510 to 665 °C after calibration, see Supporting
 145 Information for details in temperature calibration) under a

40 Torr of 1:1 H_2 and H_2O mixture (Figure S1). At open 146
 circuit voltage (OCV), the ceria surface reaches thermal 147
 equilibrium with H_2 and H_2O in the gas phase, as described by 148
 the following equation 149



where Ce'_{Ce} , Ce^x_{Ce} , V_{O}'' , and O_{O}^x are Ce^{3+} , Ce^{4+} , doubly positive 151
 charged oxygen vacancy, and oxide ion on oxygen site, 152
 respectively. With no net current flowing through the SOC, 153
 this equilibrium is determined by temperature and gas partial 154
 pressures. 155

If a potential is applied between the ceria WE and the Pt CE, 156
 charged species move through the cell as a current flow, driving 157
 the surface redox states and coverage away from thermal 158
 equilibrium and promoting the surface chemistry in eq 1. With 159
 the Au current collector grounded, the cell potential is 160
 controlled by the applied potential on the Pt CE through a 161
 potentiostat: $V_{\text{cell}} = V_{\text{Pt}} - V_{\text{Au}}$. For example, when the Pt 162
 electrode is under positive potential (e.g., +1.2 V in Figure 1c), 163
 electrons move from the Au current collector to the ceria 164
 electrode, while oxide ions are driven from the ceria film to the 165
 YSZ electrolyte and are conducted to the Pt CE. Therefore, the 166
 forward reaction of eq 1, H_2O electrolysis, is favored on ceria. 167
 If Pt is negatively biased, the fluxes of O_2^- and polarons are 168
 switched, leading to H_2 electro-oxidation (eq 1, reverse 169
 reaction) on the ceria surface and H_2O electrolysis at the Pt 170
 electrode. The current–voltage profile that reflects the activity 171
 of the SOC is given in Figure S2. Note that during 172
 electrochemical reactions, the local region of each cell 173
 component is subject to different “real” potential depending 174
 on ohmic loss and reaction kinetics.^{19–21} Hence, only the cell 175
 applied potential of +1.2 and –1.2 V are given to denote 176
 cathodic and anodic reactions on the ceria electrode, 177
 respectively. 178

Bulk Reduction of Ceria during Water Electrolysis. 179
 The accepted view on ceria reduction/oxidation behavior 180
 suggests several differences between the surface and the bulk. 181
 In general, the ceria surface is more enhanced with reduced 182

183 states compared with the bulk over wide ranges of temperature
 184 and oxygen partial pressure.^{25–28} Its redox dependence on the
 185 environment could exhibit another stark contrast.²⁹ Ceria
 186 redox behavior can continually evolve when it is driven further
 187 away from equilibria. In our previous APXPS studies, we have
 188 shown that ceria is surface-reduced upon cathodic potential, as
 189 signaled by shift in $\text{Ce}^{3+}/\text{Ce}^{4+}$ ratio.^{19–21} Herein, we find that
 190 ceria is also reduced through the bulk by using XRD and XAS.
 191 In Figure 1a, the XRD peaks of ceria shift toward smaller 2θ
 192 values during operation, consistent with chemical expansion of
 193 reduced ceria.^{13,30} The same reduction regions are concomitantly
 194 revealed in the $\text{Ce}^{3+}/\text{Ce}^{4+}$ chemical map derived from
 195 XAS measurement. The oxygen non-stoichiometry values
 196 obtained from X-ray absorption near edge structure
 197 (XANES), extended X-ray absorption fine structure
 198 (EXAFS), and XRD (inset Figure 1a) are consistent, with a
 199 slightly higher degree of reduction revealed in XANES and
 200 EXAFS measurements. This small discrepancy indicates that
 201 ceria reduction occurs through the bulk structure mainly with
 202 good crystal quality but accompanied by a small amount in the
 203 amorphous or disordered state.

204 Oxygen non-stoichiometry x values extracted from XRD
 205 characterization across the ceria electrode bulk during
 206 operation are summarized in Figure 2. This is achieved by
 207 first extracting the ceria cell parameter using Le Bail refinement
 208 and then applying Kim's law³¹ (see Supporting Information for
 209 more details). Ceria is mostly oxidized under the OCV
 210 condition at 611 °C (Table S2). Upon applied potential of
 211 +1.2 V, varying degrees of reduction are observed from the
 212 bulk non-stoichiometry x map along the ceria electrode
 213 (Figure 2a and Table S6), with the highest x values obtained
 214 close to the edge of the Au current collector. A similar pattern
 215 is delineated in our previous reports on surface redox state
 216 mapping.^{19–21} Such similarity reveals that reduction in the
 217 bulk is following the one on the surface. More essentially, it
 218 captures the propagation of reduction from the surface region
 219 to the bulk structure, as a result of the rate-limiting kinetics of
 220 surface reaction.¹⁹ It was noted that this reduction was
 221 observed only during cathodic potential. Ceria was oxidized
 222 back to the original OCV condition soon after the removal of
 223 applied potential, as confirmed by both XRD and XAS
 224 measurements. This also excludes possible reduction effects
 225 from the X-ray beam or by elevated temperatures.

226 To find possible correlation between the non-stoichiometry
 227 x and other parameters, the cell current was systematically
 228 varied, either by changing the applied potential at a constant
 229 calibrated WT (611 °C) or by changing the WT at a constant
 230 applied potential (+1.2 V on Pt). A linear dependence of bulk
 231 oxygen non-stoichiometry x on the cell current was observed
 232 (Figure 2b). This agrees with our previous finding of rate-
 233 limiting surface kinetics on ceria.²¹ If surface kinetics of water
 234 electrolysis is fast, we should expect the oxygen vacancies to be
 235 rapidly filled, and therefore, the linear increase of oxygen non-
 236 stoichiometry x with the cell current should not be observed.
 237 The oxygen non-stoichiometry determined at +1.2 V (on Pt)
 238 while scanning the potential (denoted in green) is larger than,
 239 yet within the error of measurement, the one obtained at 611
 240 °C during temperature scan (denoted in purple) for a given
 241 current. This is due to the difference in the measurement spot
 242 of the two scans. The EXAFS measurements also show similar
 243 dependence of bulk oxygen vacancy (i.e., non-stoichiometry)
 244 on electrical potential (Figure 1a inset and Figure S12). The
 245 linear dependence (at a given cell location) reveals that the

ceria redox states can be effectively controlled by tuning 246
 electrical stimulus and/or thermal stimulus. A more 247
 fundamental correlation of ceria non-stoichiometry and the 248
 electrochemical potential is currently under investigation with 249
 different cell geometries and will be included in a future work. 250

Non-Equilibrium α - α' Phase Couple. In the temper- 251
 ature range of 450–650 °C and for oxygen non-stoichiometry 252
 x of 0.01–0.16, there is a region in the ceria phase diagram 253
 where two cubic phases of ceria coexist, denoted as α and α' 254
 (higher oxygen-deficient phase). This is the so-called 255
 miscibility gap as depicted in the reproduced phase diagram 256
 in Figure 3a. At our WT of 611 °C, any deficient ceria with an 257

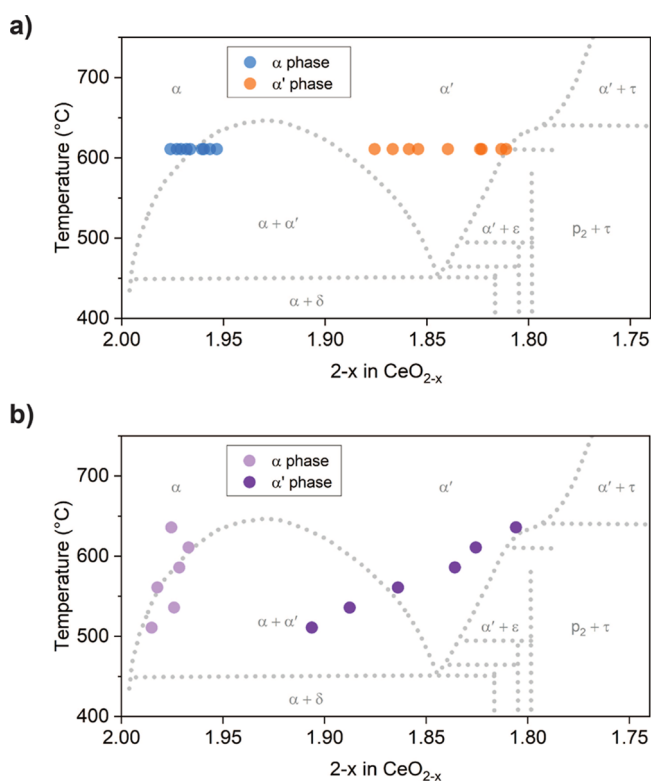


Figure 3. Ceria α - α' phase couple during water electrolysis. (a) Oxygen non-stoichiometry reported in the ceria phase diagram. The $2-x$ values are extracted from operando XRD measurements on the ceria electrode at 611 °C during H_2O electrolysis. (b) XRD-derived oxygen non-stoichiometry at +1.2 V under different WT-s. Data are collected at a fixed position where α and α' phases co-exist and are reported in the reproduced ceria phase diagram.³² The equilibrium phase diagram is reproduced with permission from ref 32. Copyright 1985 Elsevier.

oxygen non-stoichiometry x of 0.036–0.106 would separate 258
 into α and α' phases, with a composition of $\text{CeO}_{1.964}$ and 259
 $\text{CeO}_{1.894}$, respectively. As predicted by the lever rule, the only 260
 difference in phase separation between two deficient ceria 261
 inside the miscibility gap (e.g., $\text{CeO}_{1.95}$ and $\text{CeO}_{1.92}$ at 611 °C) 262
 lies in the relative amounts (relative intensities) of the two 263
 phases. This is illustrated by our simulation of powder 264
 diffraction patterns of ceria for several oxygen non- 265
 stoichiometries (Figure S3). 266

However, when the α - α' phase couples derived from our 267
 operando XRD measurements at 611 °C are plotted (Figure 268
 3a), such expected behaviors of the α - α' phase couples are 269
 not observed. If data points corresponding to the α phase 270

271 mostly gather at the left boundary of the miscibility gap, those
272 corresponding to the α' phase exhibit an unexpected behavior
273 as they do not cluster at the right boundary. The lowest $2 - x$
274 value of the α' phase in a α - α' phase couple is found at 1.811.
275 Again, these α - α' phase couples were observed only during
276 H_2O electrolysis. Ceria regained its oxidized state in the α
277 phase once the cell was returned to the OCV condition. Beam
278 effect and temperature effect are thus excluded as well.

279 Furthermore, we investigated these two phases by their
280 relative intensities and coherently diffracting domain size as a
281 function of $2 - x$ (Figures S4 and S5). The α phase, in both its
282 relative intensity and coherent domain size, seems to randomly
283 scatter over a small $2 - x$ range without following any obvious
284 trend. The α' phase, on the other hand, reveals a general
285 increase in the relative intensity and domain size against the 2
286 $- x$ values. This observation is in relation with the splitting
287 mechanism when going through the miscibility gap: phase
288 transition occurs inside each nanocrystal of ceria, and both
289 intensity and domain size of the α' phase increase at the
290 expense of the α phase.

291 We conclude that each α - α' phase couple seems to denote
292 its own discrete phase diagram. Phase separation is a direct
293 consequence of the Gibbs free energy of mixing versus the
294 chemical composition at a given temperature. If the
295 homogeneous solution is unstable within a composition
296 range (x of 0.036–0.106 at 611 °C), the separation of the
297 mixture into two phases (α and α') would occur and effectively
298 reduce the Gibbs free energy. The common tangent line to the
299 Gibbs free energy curve gives the points of the two
300 compositions ($\text{CeO}_{1.964}$ and $\text{CeO}_{1.894}$), which hence mark the
301 limit of thermodynamic stability. Because the Gibbs free
302 energy curve is a function of temperature, the thermodynamic
303 stability limits should be fixed if the temperature is unchanged.
304 Therefore, our results indicate that, even at the same
305 temperature, there is a shift of free energy curves of the two
306 phases induced by electrochemical potential. Such relative shift
307 leads to the consequence that the chemical composition values
308 (i.e., $2 - x$ of α and α' phases) determined from the common
309 tangent of the free energy curves deviate from the equilibrium
310 stability limits, even though they are obtained at the same
311 temperature of 611 °C. As we have noted above, the “real”
312 potential differs for local regions; thus, only the cell applied
313 potential is given for qualitative discussion herein. Precise
314 determination of the electrochemical potential is beyond the
315 scope of this work and will be a future focus that requires
316 different cell geometry and experimental design.

317 Finally, we take a closer look at how this system is affected
318 by the temperature parameter. On the ceria electrode where
319 both α and α' phase co-exist at cathodic applied potential
320 (+1.2 V on Pt), the WT was systematically varied between 511
321 and 661 °C (as illustrated in Figure 2b). The $2 - x$ values
322 obtained at six different temperatures are reported in Figure 3b
323 in the reproduced ceria phase diagram. Similar to that in Figure
324 3a, the α phase composition stays approximately on the left
325 boundary of the miscibility gap, while the α' phase
326 composition deviates significantly from the right boundary.
327 In addition, when temperature is decreased, the α' phase
328 border shifts toward more oxidized stoichiometry and
329 eventually ends inside the classic miscibility gap. This observed
330 shrinking of the miscibility gap indicates a relative decrease of
331 the α' phase free energy curve, compared to that of the α
332 phase. Therefore, at lower temperatures, the α' phase becomes

energetically favorable and phase transition from α to α' 333
finishes at a more oxidized non-stoichiometry. 334

Phase Separation under Electrochemical Potentials. 335
It was reported in literature for perovskite oxides that structural 336
evolution and the associated phase change induced by applying 337
electrochemical potential are equivalent to that triggered by 338
changing the oxygen partial pressure in the environment.^{23,33,34} 339
However, our findings in this study suggest some difference. 340
On one hand, consistency is shared with previous reports as we 341
have also verified the control of oxygen non-stoichiometry in 342
ceria via electrochemical potential without changing the gas 343
atmosphere. But on the other hand, we did not observe the 344
phase change behavior typical for cases under thermal 345
equilibrium, that is, phase change does not follow the classic 346
phase diagram of deficient ceria. Upon electrochemical 347
potentials, the miscibility gap right boundary for CeO_{2-x} at a 348
given relatively high temperature is found to expand toward 349
the reduced non-stoichiometry, while at lower temperatures, 350
this boundary is observed to shrink toward the oxidized non- 351
stoichiometry. 352

In order to examine the miscibility gap behavior under 353
electrochemical potentials, we calculated the right boundary 354
position limits. From the linear dependence of α' relative 355
intensity on $2 - x$ values at 611 °C (Figure S4), we found that 356
phase transition occurs at the composition $\text{CeO}_{1.879}$, where the 357
 α' phase starts to exist and the nanocrystal turns completely 358
into the α' phase at $\text{CeO}_{1.798}$. In addition, the evolution of α' 359
phase relative intensity with temperature (shown in Figure 360
S6a) allows us to extract the temperature limits. A single α 361
phase is found for temperatures below 480 °C with a chemical 362
composition of $\text{CeO}_{1.933}$, while a single α' phase is obtained at 363
temperatures above 659 °C as $\text{CeO}_{1.780}$. The deduced limits 364
(Figure S7) suggest that, compared to the thermal equilibrium 365
condition shown in the classic phase diagram, a cathodic 366
electrochemical potential shifts the right boundary toward 367
more reduced stoichiometry (smaller $2 - x$ value) at high 368
temperatures and to a more oxidized state (larger $2 - x$ value) 369
at lower temperatures. Phenomenally, because only cathodic 370
potentials are introduced in this study, such shifts are unlikely 371
to be associated with direction of the current passage but most 372
probably with the magnitude of it. The higher the temperature, 373
the larger the cell current is delivered. Once it passes a 374
threshold, the miscibility gap is switched from shrinking to 375
expanding. More fundamentally, this observation can suggest 376
the presence of a threshold in electrochemical potential that 377
essentially shifts the relative free energy between α and α' 378
phases. An electrochemical system is special in that it requires 379
the same magnitude in electrochemical potential of the two 380
phases where phases separate under steady state. However, the 381
chemical potential and electrical potential of the two phases are 382
not necessarily the same because of their different properties. 383
Ultimately, the electrochemical potential is determined by the 384
coupling of oxygen vacancy concentration and the electrical 385
driving force in the two phases. Therefore, the observed 386
deviation in phase separation behavior from the conventional 387
phase diagram is understandable because the latter considers 388
free energy (i.e., chemical potential) only. Also, it is unlikely 389
that material structure feature²¹ could play a role in this 390
unexpected phase separation behavior because the ceria 391
crystallites are small in size (Figure S5) with no preferred 392
orientation. Note that our experimental data from the α phase 393
do not allow a certain exclusion of any possible evolution of 394
the left boundary position. Future work will focus on precis- 395

396 the exact shape of the continuum of miscibility gap and on
397 better understanding the mechanism behind it. We would like
398 to reiterate in this work that while considering phase transition
399 in a system subject to electrochemical potential, one should
400 consider *electrochemical free energy* instead of chemical free
401 energy, which includes effects from the large-scale electrical
402 environment.³⁵ Moreover, volume change due to chemical
403 expansion and entropy change due to oxygen vacancy
404 incorporation upon reduction bring additional contribution
405 to the relative shift in free energy curves of the two phases.¹⁵

406 ■ CONCLUSIONS

407 In conclusion, the oxygen vacancy evolution in ceria and the
408 associated phase transition were studied using operando
409 synchrotron XRD and XAS. We found that upon cathodic
410 potential, the ceria is reduced not only on the surface but also
411 through the bulk. This is consistent with our previous reports
412 of the rate limiting surface exchange reaction.^{19–21} Within the

413 electrochemical active region, ceria is bulk-reduced both in
414 crystalline structure (characterized by XRD) and in local
415 atomic order (by XAS). Therefore, from these bulk-sensitive
416 studies and previous surface-sensitive measurements, we obtain
417 a full picture of the oxygen vacancy behavior in ceria operando.
418 By using this model system of ceria, we aim to demonstrate the
419 capabilities of in situ/operando X-ray synchrotron-based
420 techniques in material characterization far from equilibrium.

421 Analysis of our operando XRD data allowed a quantitative
422 correlation of oxygen vacancy non-stoichiometry to cathodic
423 cell current. This demonstrates that by electrochemical
424 stimulus, we could feasibly modify the concentration of oxygen
425 vacancy without change in gas atmosphere and essentially
426 control the chemical composition CeO_{2-x} in a steady state.
427 Another implication lies on the electromechanical coupling.
428 For ceria and ceria-based energy conversion and storage
429 devices, the operation under non-equilibrium condition could
430 introduce substantial chemical expansion/contraction of the
431 material, to which one needs to pay attention for mechanical
432 durability concerns. The same applies to other fluorite-
433 structured materials.

434 When the electrochemical potential is large enough to drive
435 ceria into a miscibility gap, phase separation occurs and a
436 mixture of α and α' cubic phases is observed. Relative
437 intensities of the two phases suggest that the phase transition
438 occurs inside each nanocrystal of ceria. An interesting find is
439 that the phase transition behavior from α to α' induced by
440 electrochemical potentials differs from that under thermal
441 equilibrium condition. Depending on the magnitude of cell
442 current, phase transition completes at different stoichiometric
443 states even at the same temperature. Such difference indicates a
444 relative shift of the electrochemical free energy curves of the
445 two phases regardless of a retained temperature. This phase
446 transition achieved by electrochemical potential opens up new
447 possibilities in using potential to control the phase of a
448 functioning oxide materials. However, a more quantitative
449 understanding of the phase change mechanism at steady state
450 is required and will be the focus of our future work.

451 ■ EXPERIMENTAL SECTION

452 **In Situ Setup.** The cell was mounted on a specially designed
453 sample holder with an X-ray transparent dome to insure a constant
454 atmosphere during cell operation. The dome was air-cooled when the
455 temperature reached over 200 °C.

For the XRD and XAS experiments, a WT of 650 °C was set at the
456 heating stage. However, because of temperature gradient from the
457 bottom to the top of the cell, the probed region experienced a lower
458 WT of 610–615 °C (see Supporting Information for details). The cell
459 was exposed to a controlled atmosphere composing 20 Torr of H_2O
460 and 20 Torr of H_2 . Applied voltage varied from +1.2 to –1.2 V for
461 H_2O splitting and H_2 oxidation on the ceria electrode, respectively.
462

Micro-XRD. Powder X-ray microdiffraction experiments were
463 conducted at beamline 12.3.2 of the Advanced Light Source of the
464 Lawrence Berkeley National Laboratory. A high-brilliance X-ray beam
465 (6.5 keV) was focused down to about $4\ \mu\text{m} \times 10\ \mu\text{m}$ using a pair of
466 Kirkpatrick–Baez mirrors.^{36,37} Diffraction patterns were collected
467 using a two-dimensional DECTRIS Pilatus 1M X-ray detector,
468 positioned at 70° with respect to the incident beam. The sample was
469 tilted to an angle of 30° relative to the incident beam (reflection
470 geometry). The distance from the sample to the center of the detector
471 was ~138 mm. An exposure time of 160 s per pattern was used. The
472 sample–detector distance, the center channel of the detector, and the
473 tilt of the detector relative to the sample surface were calibrated using
474 a Laue pattern obtain from a strain-free Si single-crystal. Data
475 processing is detailed in Supporting Information.
476

Micro-XAS. Hard X-ray microprobe measurements were per-
477 formed at beamline 10.3.2 of the Advanced Light Source of the
478 Lawrence Berkeley National Laboratory.³⁸ This beamline includes a
479 pair of focusing Kirkpatrick–Baez mirrors and a Si 111 two-crystal
480 monochromator which forms the X-ray micro-beam. A Canberra 7-
481 element UltraLeGe solid-state detector records the fluorescence from
482 the sample.
483

A series of X-ray fluorescence maps were collected at the ceria/Au
484 electrode. The Ce^{3+} and Ce^{4+} chemical species were identified by
485 mapping the sample using several incident energies around the Ce
486 emission edge: 5710 eV (pre-edge), 5727 eV (mostly Ce^{3+}), 5738.4
487 eV (mostly Ce^{4+}), and 5747 eV (post-edge). Energy calibration was
488 such that the first peak of the XANES of CeO_2 was defined to be at
489 5730.39 eV. An additional map was collected above the Au L3 edge in
490 order to locate the Au region.
491

XAS data were acquired in fluorescence mode using the quick XAS
492 method, by slewing the monochromator continuously and at a
493 constant speed from low to high energies as the X-ray absorption data
494 got recorded. This mode allowed the collection of spectra in the
495 5620–5940 eV range with 2 eV step size (0.3 eV in the XANES
496 region) and ~5 s/point summed over 30 sweeps (0.81 s/point in the
497 XANES region). Each sweep from low to high energy took about 1
498 min; so a scan of 30 sweeps took half an hour. The energy scale of
499 each sweep was repeatable to ~0.1 eV. Ten scans were acquired at
500 each spot of interest in order to improve the signal-to-noise ratio. For
501 simplicity, a linear combination of two models were used in data
502 fitting. The first model is the fluorite-type cubic structure of ceria with
503 the Ce atoms surrounded by eight oxygen neighbors; while the
504 second one has the same structure consisting of four oxygen
505 neighbors and four vacancies surrounding the Ce atoms. Further
506 details can be found in Supporting Information.
507

508 ■ ASSOCIATED CONTENT

509 ⓘ Supporting Information

The Supporting Information is available free of charge at
510 <https://pubs.acs.org/doi/10.1021/acsami.0c08284>.
511

Anton-Paar DHS heater during operando measurement,
512 current–voltage profile, simulation of powder diffraction
513 patterns at 611 °C, relative intensity of α and α' phases
514 as a function of $2 - x$, coherent domain size of the two
515 phases as a function of respective intensity, evolution of
516 relative intensity and $2 - x$ value of α' as a function of
517 WT, position limits of α' phase, powder diffraction
518 patterns, calibration of WT, oxygen non-stoichiometry
519 calculation, refinement results for 611 °C at OCV, 520
521 applied potentials of +0.3, +0.6, +0.9, and +1.2 V,
522 refinement results for an applied potential of +1.2 V at

523 511–661 °C, data processing from XAS measurements,
524 k^2 -weighted Ce L₃ EXAFS spectra, and corresponding
525 Fourier transform results (PDF)

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572 Author Contributions

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574 Author Contributions

575 C.D., F.B., and Z.L. conceived the experiment. Y.Y. and C.Z.
576 fabricated the SOCs. C.D., N.T., M.K., F.B., M.A.M., and Z.L.
577 carried out the measurements. C.D., N.T., and M.K. performed
578 the XRD analysis and M.A.M., Y.L.H., and F.B. interpreted the
579 XAS and electrochemical results. B.W.E. and Z.L. supervised
580 the project. All authors contributed to writing the manuscript.

Notes

The authors declare no competing financial interest.

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