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Resilience Through **Community Engagement** and **Technology**

The High Frequency Attenuation Parameter in NGA-West3 EAS: Modeling κ_0 with Varying Levels of Complexity

G.A. Parker¹, G.M. Atkinson², A. Baltay³, D.M. Boore⁴, T.E. Buckreis⁵ and J.P. Stewart⁶

ABSTRACT

We are developing an earthquake ground-motion model (GMM) for effective amplitude spectra (EAS, the smoothed quadratic mean of two horizontal-component Fourier amplitude spectra) as part of the NGA-West3 project. κ_0 is a model parameter that determines the shape of the high frequency EAS spectrum by controlling how far it deviates from an ω^{-2} source model. We develop κ_0 models within an EAS GMM with increasing levels of complexity: regional constants, regional relationships with the time-averaged shear wave velocity in the upper 30 m (V_{S30}), and modular models based on the depth to the 1 km/s shear wave velocity isosurface (z_I). We estimate regional values of κ_0 between 0.023 and 0.065 s. Mexico, southern California, and Japan have lower than average κ_0 (less attenuation) whereas Taiwan and northern California have higher than average values (more attenuation). We find that sites with low V_{S30} have larger κ_0 , whereas κ_0 does not scale with $V_{S30} > 300$ m/s. Variations in z_I have a stronger control on κ_0 than V_{S30} , up to a factor of three in some regions; we parameterize a modular z_I model that can be combined with our $\kappa_0(V_{S30})$ model.

Modeling Workflow

The NGA-West3 project is a large multi-institution effort to develop a global earthquake ground motion database for active crustal regions [1], as well as GMMs for 5%-damped RotD50 pseudo-spectral acceleration and effective amplitude spectra (EAS, the smoothed quadratic mean of two horizontal-component Fourier amplitude spectra). We develop a model for EAS with the primary input parameters of moment magnitude (M), Joyner-Boore distance (R_{JB}) and the time-averaged shear-wave velocity in the upper 30 m (V_{S30}), with modeled adjustments to account for regional variations in stress drop, source depth, anelastic attenuation ($1/Q$), and the high frequency attenuation parameter kappa (κ_0). Our model utilizes seismological theory to constrain terms for geometrical spreading, $1/Q$, κ_0 , and the spectral shape of the seismic source, and combines them with empirical model terms for the near-source saturation of ground motions, linear site amplification, and basin response [2]. We consider over 91,000 EAS records from 1091 $M \geq 3$ earthquakes in ten global regions (Fig. 1a). To avoid tradeoffs between components, we use the following multi-phased model fitting approach:

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1. Adjust the EAS amplitude from each record for nonlinear site effects using an existing model [3].
2. Fix the geometrical spreading to decay as $R^{-0.5}$ beyond 50 km based on the decay of surface waves in a half-space [4] using a bilinear function [5]. Use records with $R_{JB} \geq 150$ km to fit regional anelastic attenuation coefficients, parameterized as the inverse of the seismic quality factor (Q), and an initial V_{S30} scaling model using a trilinear form [6], allowing for regional variations in both components. Both $1/Q$ and V_{S30} -scaling models are fit frequency-by-frequency, and then coefficients are smoothed.
3. Fit near-source saturation [7] and the geometrical spreading using records with $R_{JB} \leq 50$ km. These terms trade off against each other, resulting in non-uniqueness. We find that a constant geometrical spreading of R^{-1} works well to fit the data when combined with a near-source saturation model that depends on frequency (f), magnitude, and earthquake source depth [8].
4. Adjust records from step 1 using the path and site models from steps 2-3, and invert adjusted records over all frequencies to estimate event-specific earthquake source parameters (e.g., earthquake corner frequency, f_c [9]), regional amplitude constants, and site-specific κ_0 values (Eq. 1; Fig. 1b) [10] using a mixed-effects regression [11] where earthquake, region, and station are used as grouping factors.

The parameter κ_0 controls the decay rate of the acceleration spectrum at high frequencies (Eq. 1, Figure 1b) [10]. Here we focus on using site-specific estimates of κ_0 from step 4 above, which falls within the class of broad-band fitting methods described by [12]. We use these estimates to develop models for median κ_0 that depend on region, V_{S30} , and z_1 , here using northern California (NorCal) as an example region.

$$EAS \propto -\pi f \kappa_0 \quad (1)$$

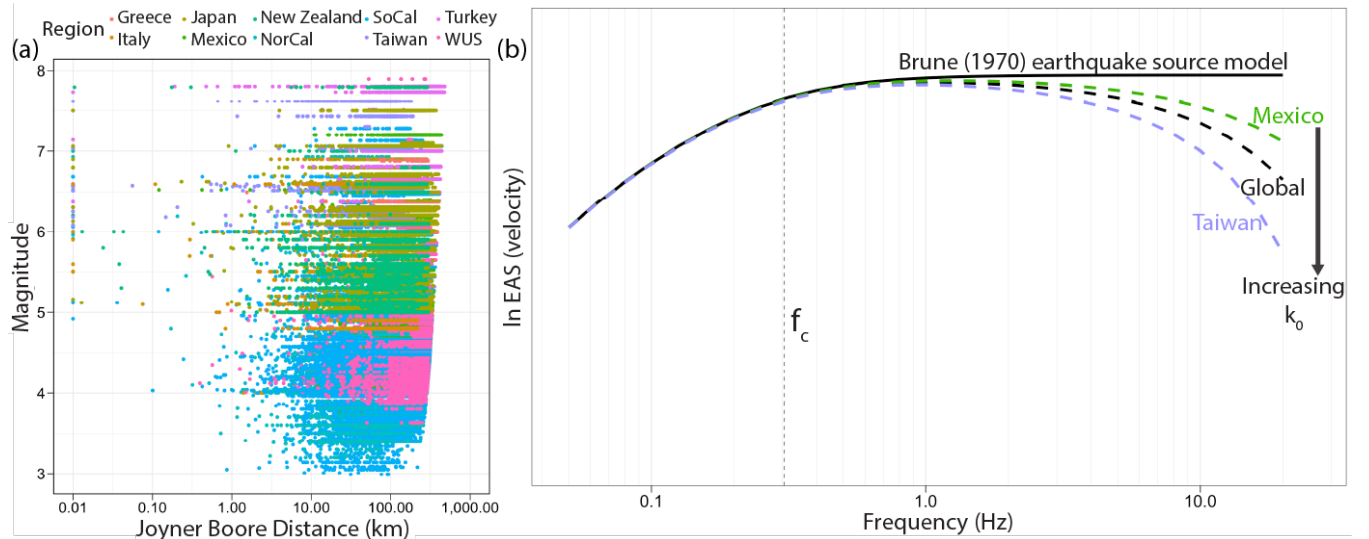


Figure 1. (a) The magnitude-distance distribution of EAS records used in model development, color-coded by earthquake region. SoCal = Southern California; WUS = western United States. (b) Schematic EAS spectra for an **M6** earthquake showing deviations from an ω^{-2} source model [9] for varying κ_0 values, including the global average and regional endmembers from this study (Fig 2a).

Models for the High Frequency Attenuation Parameter κ_0

We use site-specific estimates of κ_0 to develop several models with varying complexity. The simplest model is a constant value of κ_0 for each region. We find that κ_0 varies from 0.023s in Mexico to 0.065s in Taiwan, with a global average of 0.043s (Fig. 2a). In general, Japan, southern California and Mexico have more high-frequency energy (smaller κ_0 and hence less attenuation), whereas Taiwan and NorCal are more depleted in high-frequency energy with more attenuation (larger κ_0).

V_{S30} - Scaling

We find that there is some correlation between κ_0 and V_{S30} for the softest sites ($V_{S30} \leq 300$ m/s), with softer sites having larger κ_0 (more high frequency attenuation); there is no correlation for sites with $V_{S30} > 300$ m/s (Fig. 2b). We develop region-specific models using a sigmoid function following [5] that increases κ_0 for sites with low V_{S30} (Fig. 2b), fitting coefficients for each region separately. This is consistent with some previous studies that have found no or limited relationships between κ_0 and bedrock velocity for hard-rock sites [12–15] or with V_{S30} [13, 16]. However, there are also studies that have shown a relationship between κ_0 and V_{S30} in the 300–1000 m/s range that we do not see in our data [17, 18, 19]. There are several possible reasons for this difference, including dataset difference (including regional variations) and data sparsity; accuracy of input V_{S30} values; assumptions about site amplification and Q attenuation and their corresponding tradeoffs with κ_0 ; modeling choices (e.g., we could have fit a linear model across the range of V_{S30}); and differences in κ_0 estimation methods. For example, [17] and [19] use a high-frequency linear fitting approach, whereas this study uses a broadband spectral fitting approach.

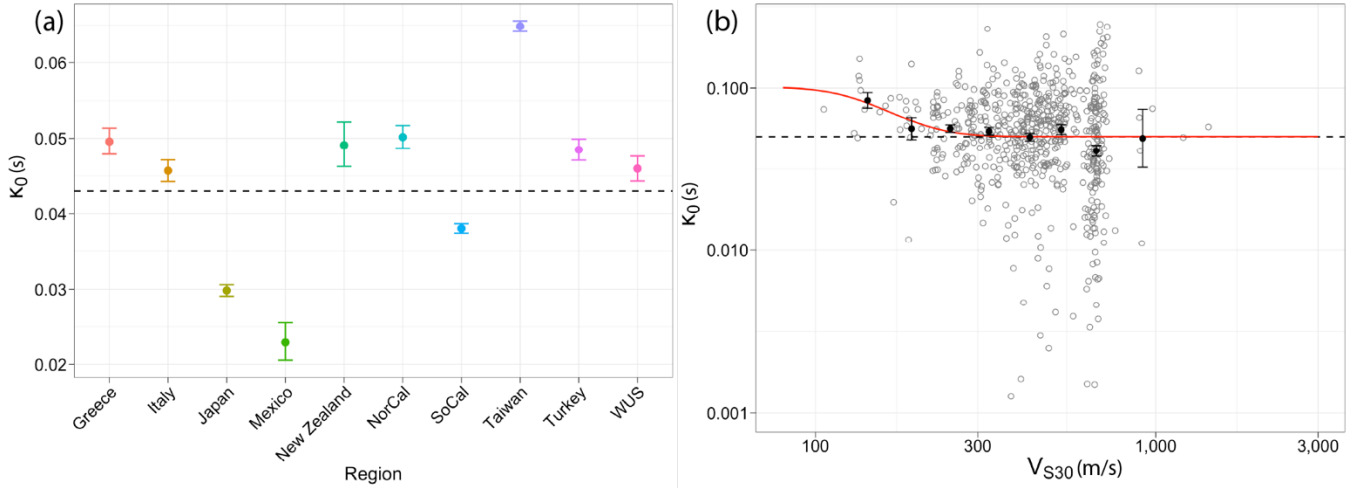


Figure 2. (a) Regional variations in κ_0 shown with ± 1 standard error bars. Black dashed line at the global average value. (b) The relationship between κ_0 and V_{S30} for NorCal. Black points show binned means with ± 1 standard errors, red line shows best-fit sigmoid function, and dashed line shows regional average κ_0 for NorCal (0.05 s).

Relationship With z_I

We find that the depth to the 1 km/s shear wave velocity isosurface (z_I) has a stronger control on κ_0 than V_{S30} . However, we acknowledge that V_{S30} is more readily available from proxy-based models, whereas estimating z_I requires a site investigation or an available seismic velocity model at the site of interest. Therefore, we parameterize a model for κ_0 using z_I that can be used in conjunction with the V_{S30} -scaling model if z_I is known. If z_I is not known, the $\kappa_0(V_{S30})$ model can be used independently. To do this, we develop a centering model between z_I and V_{S30} for each region using a sigmoid function following [6] (Fig. 3a). We do not recommend this model to predict z_I , but to account for correlation between the V_{S30} and z_I model terms.

$$\delta z_I = \ln(z_I) - \ln(\bar{z}_I) \quad (2)$$

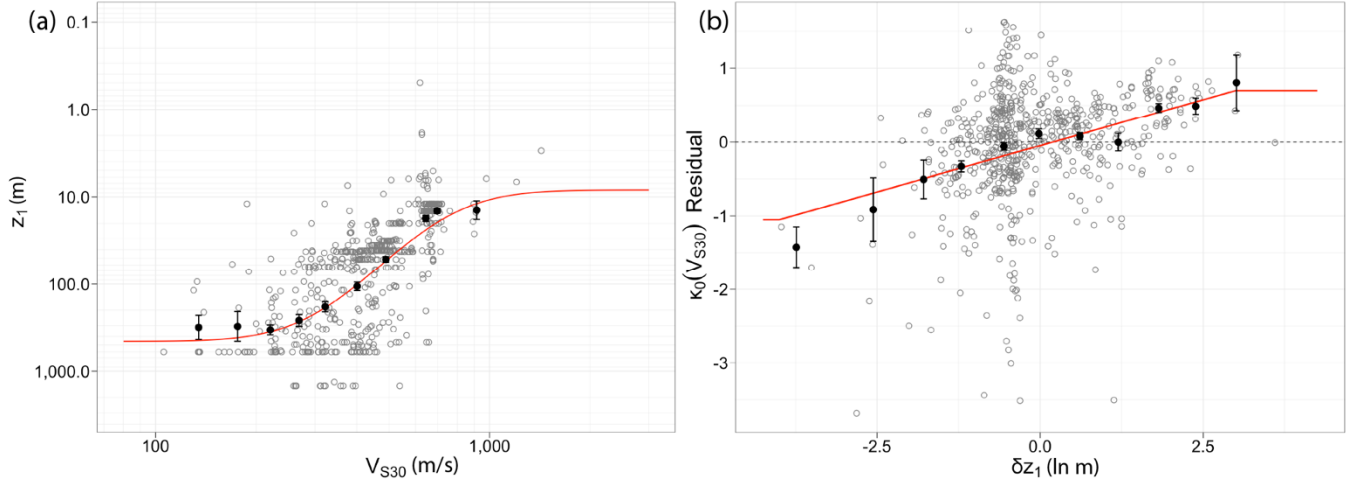


Figure 3. (a) The best-fit sigmoid centering model (red) to account for correlation between V_{S30} and z_1 in NorCal. (b) The trilinear $\kappa_0(\delta z_1)$ model for NorCal based on residuals of the $\kappa_0(V_{S30})$ model, in natural logarithm units. In both panels black points show binned means with ± 1 standard error of the mean. In (b) black dashed line shows 0 for reference.

We subtract the $\kappa_0(V_{S30})$ model from our site-specific κ_0 estimates, and fit a model to those values conditioned on δz_1 (the $\ln$ difference between the z_1 at a site and the average value $\bar{z}_1$ given the V_{S30} at the site, Eq. 2). We model $\kappa_0(\delta z_1)$ as a trilinear function [20] (Fig. 3b). We find that variations in z_1 can modify κ_0 by up to a factor of three in NorCal. This is consistent with past studies that have found κ_0 to be correlated with sediment thickness [18, 21] and depth to the 2.5 km/s shear-wave velocity isosurface [19], but here we offer a more sophisticated parameterization based on δz_1 that takes into account any correlation between z_1 and V_{S30} .

Conclusions

Based on estimates of site-specific κ_0 from the NGA-West3 database, we find that Japan, southern California, and Mexico have smaller than average κ_0 leading to larger high frequency ground motions, whereas Taiwan and NorCal have larger than average κ_0 leading to a relative depletion of high frequency energy. Our results show that κ_0 and V_{S30} are correlated for sites with $V_{S30} \leq 300$ m/s; for sites with larger V_{S30} we find no correlation. This is consistent with some previous studies investigating the relationship between κ_0 and V_{S30} [12-16], however, unlike some past studies [17-19], we do not see a trend with κ_0 across the entire range of V_{S30} considered. This may be due to dataset or methodology differences (e.g., high frequency linear fitting versus broadband spectral fitting [12]). Additionally, in estimating κ_0 we assume a V_{S30} - dependent linear site amplification that may already include implicit κ_0 effects (step 4, above), although we constrain the VS30-scalign to be 0 for frequencies ≥ 12.5 Hz. We also find that sites with deeper than average z_1 have larger κ_0 , and sites with shallow z_1 have smaller κ_0 , indicating that thicker sediment columns increase high frequency damping. This is consistent with previous findings that investigated the correlation between κ_0 and sediment depth [18, 19, 21], but here we parameterize a model for κ_0 on δz_1 (Eq. 2) such that it can be used in conjunction with our $\kappa_0(V_{S30})$ model, accounting for any correlation between the two site condition parameters. The κ_0 can then be estimated using the sum of either the regional constant or the $\kappa_0(V_{S30})$ with the $\kappa_0(\delta z_1)$ term if z_1 is available. Going forward, we will consider interpolating our site-specific κ_0 to generate spatially continuous maps. Combining all κ_0 models presented here with the linear site response model developed in step 2 above will ensure high frequency spectral shapes consistent with our other GMM components.

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