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DEC 1965

THE VAPOR PRESSURE, THE EVAPORATION COEFFICIENT AND THE HEAT
OF SUBLIMATION OF BARIUM FLUORIDE

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ABSTRACT

The vapor pressure of barium fluoride was measured in the temperature range 1261° to 1548°K by the torsion-effusion method. The free-surface sublimation pressure from the [111] faces of single crystals was measured by the torsion-Langmuir method in the range 1315° to 1492°K. Comparison of the torsion-effusion and torsion-Langmuir data indicates a sublimation coefficient of 1.0 ± 0.1 . Extrapolation of the data in the Langmuir experiments yields a heat of sublimation at 298°K of 90.41 ± 3 kcal/mole and for the Knudsen experiments yields 90.25 ± 3 kcal/mole by the third-law method and 88.9 ± 3 kcal/mole and 88.8 ± 3 kcal/mole by the second-law method.

INTRODUCTION

Measurements of the vapor pressure of barium fluoride have been made by Ruff and LeBoucher¹ between 1960° and 2206°K by use of a dynamic method, by Green et al.² by means of a Knudsen effusion study with a mass spectrometer between 1232° and 1505°K, by Hildenbrand et al.³ by means of a torsion-effusion study between 1300° and 1550°K, and by Bautista and Margrave⁴ by means of a weight loss study on single crystals between 1167° and 1250°K.

From a comparison of various effusion and free-surface evaporation studies, Bautista and Margrave⁴ suggest that the evaporation coefficient for alkaline earth fluorides and chlorides are in the range 0.1 to 0.3.

However, the free-surface and effusion experiments compared were in no case performed in both the same apparatus and in the same temperature range, so that the apparent differences in measured pressures could arise from systematic experimental errors. It therefore seemed desirable to measure the vapor pressure of some alkaline earth difluoride by an effusion and by a Langmuir method in a single apparatus. Since one of us⁵ had already obtained some torsion-effusion data for barium fluoride and since single crystals of barium fluoride were available, barium fluoride was chosen as the substance for study by the torsion-effusion^{6,7} and torsion-Langmuir⁸ methods.

EXPERIMENTAL

Barium fluoride in the form of powder of optical quality and of single crystals was obtained from Semi Elements, Incorporated.

The torsion apparatus that was used in this work is described by

Hammer and Pask.⁹ The apparatus makes use of the interaction of an electric field with a small permanent magnet to balance the torque produced by effusion of vapor from orifices or surfaces which are offset from the center of support. Calibration of the apparatus was effectuated in the manner described by Hammer and Pask. For the present study, the torsion cells were made of National Carbon ZT-101 graphite in box shape. Pairs of orifices with three different areas¹⁰ were used in the Knudsen experiments. Corrections for the finite orifice lengths¹¹ were made in calculations from the data. Before effusion holes were machined in the cell walls, each cell was heated with barium fluoride inside. No observable deflection was found when the cells and samples were heated to 1450°K. Therefore, it was assumed that measured deflections during the experiment were due only to effusion of vapor through the orifices.

In the Langmuir free-surface evaporation experiments, cleaved [111] faces of barium fluoride single crystals were placed behind graphite washers mounted in the verticle cell walls. Two different washer orifice areas¹² were used to prove that leakage through or around the washers and deflections of evaporating molecules from the beveled washer walls did not constitute significant sources of error. The exterior face of each washer was beveled to zero thickness at the center hole. Microscopic examination at 250X of single crystals heated at 1450°K for 45 min showed no roughening of the [111] surface near the center of the region exposed for evaporation. Toward the edges of the exposed surfaces, successive edges of [111] surfaces rose from the flat central region toward the original crystal surface. The correction for molecules

reflected or condensed near the edges of the exposed surface are demonstrated to be negligible by the fact that data obtained from surfaces whose areas differed by a factor of 4 were in excellent agreement.

The furnace heating elements were of hairpin-shaped tungsten wires connected in parallel. Electromagnetic repulsion of the torsion assembly which was noted in other torsion-effusion work,¹³ was not observed because of the self-cancelling of the induced-field in the hairpin elements. Temperatures were measured by a thermocouple placed in a reference graphite cell located below the effusion cell in the furnace hot zone. To calibrate the effusion cell temperature against the reference cell temperature, a thermocouple was mounted in an effusion cell. This thermocouple was calibrated by a gold point determination when both cells were in the vapor pressure apparatus. The relative temperatures of thermocouples in the reference and effusion cells were measured at a series of power settings.

The equilibrium vapor pressure of tin was measured by the torsion-effusion method to check for any systematic errors that would affect pressure readings. The third-law heat of sublimation of tin at 298°K was determined as $\Delta H_{298}^{\circ} = 71.79 \pm 0.27$ kcal/mole. In an exhaustive study of tin by the same method, Schulz¹⁴ found $\Delta H_{298}^{\circ} = 71.8$ kcal/mole compared with $\Delta H_{298}^{\circ} = 72.2$ kcal, the value suggested by Hultgren et al.¹⁵ from earlier studies.

Prior to a run, the cell and sample were heated some 50° higher than the highest temperature to be reached during the run and until the ambient pressure in the furnace was no higher than 5×10^{-5} torr. In

any particular run from 2 to 6 points were taken with at least 15 min being allowed after each temperature change in order to insure that equilibrium was reached.

RESULTS AND DISCUSSION

The results of this study are shown in Fig. 1 and in Tables I and II. The lines of the graph are the least squares fit to the data collected.

Molecular flow equations for low-pressure gas flow through an orifice have been shown to apply only when the mean free path of the vapor species is greater than some critical value that depends on the dimensions of the effusion orifice.^{16,17} Mean free path to diameter ratios were calculated using the hard spheres approximation¹⁸ with the molecular diameter of barium fluoride taken as 8.14\AA , the sum of ionic diameters at room temperature, assuming a linear molecule. This value of λ/d can be used only for relative comparison. Recent experiments¹⁹ indicate barium fluoride may be a bent molecule and thus would have a slightly smaller effective molecular diameter than assumed. Knudsen pressures for which calculated λ/d values are less than 1.0 are excluded in the final calculations of the heats of sublimation and are not shown in Fig. 1 because systematic deviations in calculated values, presumably arising from failure of the molecular flow approximation, are apparent for measurements at lower λ/d 's. For orifice 1 when λ/d was less than 1.0, ΔH_{298}° values of 1.54 to 1.28 kcal/mole below the average were found. Schulz and Searcy¹³ also found substantial deviation in ΔH_{298}° values when λ/d became less than a critical value near unity.

Least-square calculations from Knudsen data where $\lambda/d > 1.0$ yield

for the vapor pressure in atmospheres $\log P = -\left(\frac{79.73 \pm 0.77}{4.576}\right) \times \frac{10^3}{T} + 7.15 \pm 0.12$ between 1261° and 1548°K and $\log P = -\left(\frac{80.08 \pm 1.2}{4.576}\right) \times \frac{10^3}{T} + 7.18 \pm 0.28$ between 1315° and 1492°K for the Langmuir experiment. The errors are the standard deviations from the least-squares fits.

The heats of sublimation of barium fluoride at 298°K for the Knudsen and Langmuir data were calculated by both the second-law and third-law methods. In all calculations, BaF_2 was considered to be the major vapor species² because thermodynamic calculations indicate that the pressures produced by the reactions $\text{BaF}_2(\text{s}) + \text{C}(\text{s}) = \text{BaF}(\text{g}) + \text{CF}(\text{g})$, $\text{BaF}_2(\text{s}) + \text{C}(\text{s}) = \text{Ba}(\text{g}) + \text{CF}_2(\text{g})$ and $\text{BaF}_2(\text{c}) + 1/2\text{C}(\text{s}) = \text{BaF}(\text{g}) + 1/2\text{CF}_2(\text{g})$, would be several orders of magnitude below the observed pressures.

The same heat capacity equations and free energy functions for solid barium fluoride²⁰ and the same estimated heat capacity and free energy functions for barium fluoride gas²¹ were used in calculating heats of sublimation as were used by Green et al.² and by Bautista and Margrave.⁴ The free energy functions used are expected to be as reliable as can presently be estimated²¹ whether or not the $\text{BaF}_2(\text{g})$ molecule is linear or bent as suggested by recent electric dipole studies.¹⁹ The second-law treatment yielded $\Delta H_{298}^\circ = 88.78 \pm 1.26$ kcal/mole from the Knudsen experiments and $\Delta H_{298}^\circ = 88.94 \pm 1.75$ kcal/mole from the Langmuir experiments. The third-law calculations for the 46 Knudsen measurements for which $\lambda/d > 1.0$ yielded $\Delta H_{298}^\circ = 90.25 \pm 0.28$ kcal/mole and for the 25 Langmuir measurements yielded $\Delta H_{298}^\circ = 90.41 \pm 0.35$. The indicated errors are the standard deviations. The free energy functions could be in error by enough to contribute a 4 kcal error to the calculated heats at 298°K; however, the

heat of sublimation calculated by the third-law method is probably correct to within ± 3 kcal/mole. Assuming a $+5^\circ$ error at the high end of measurements and a -5° error at the lower end, (correspondingly $\pm 10\%$ errors in pressures) yields an estimated maximum error of 2.8 kcal for the second law calculation.

Brewer et al.²¹ concluded from the early work of Ruff and LeBoucher¹ that the heat of sublimation of barium fluoride is about 88 kcal/mole at 298°K. Green et al.² in the mass spectrometer study found $\Delta H_{298}^\circ = 92.3$ kcal/mole by the third-law method and $\Delta H_{298}^\circ = 93.8$ kcal/mole by the second-law method. Torsion-effusion studies of Hildenbrand et al.³ yield $\Delta H_{298}^\circ = 88.6$ kcal when recalculated with the same free energy functions used by the other investigators. The Langmuir measurements of Bautista and Margrave⁴ yield $\Delta H_{298}^\circ = 92.3$ kcal/mole from third law analysis and $\Delta H_{298}^\circ = 94.7$ kcal/mole by second law analysis.

Hildenbrand and coworkers³ believe their experimental results indicate that barium fluoride is bent. We will not comment further on this point because the work of Hildenbrand et al. is not yet published except in company reports. The second and third law measurements of the present work agree well when Brewer's²¹ free energy functions are used.

Since the same free energy functions were used by Green et al.,² by Bautista and Margrave⁴ and by us and since the measurements were made in approximately the same temperature ranges, the differences in calculated heats of sublimation at 298°K reflect the systematic differences of approximately a factor of 2 in measured pressures and not possible errors in the free-energy functions. Mass spectrometer pressure measure-

ments are usually expected to be uncertain by a factor of two or more. It was primarily comparison of the Langmuir data with effusion measurements of other investigators that led Bautista and Margrave⁴ to hypothesis that evaporation coefficients for alkaline earth halides generally lie in the range 0.1 to 0.3. Recently, however, Loehman, Kent and Margrave²² have measured Knudsen data of strontium chloride above its melting point and Langmuir data below. The Langmuir experiments were performed on single crystals suspended in the furnace. Extrapolation of their Knudsen data into the range below the melting point clearly indicates an evaporation coefficient of 0.3 for the solid. For barium fluoride, comparison of Bautista's results to our effusion results leads to the conclusion that the value of α is about 0.5.

Comparison of our own free surface and effusion data, however, yields $\alpha = 1.0 \pm 0.1$. The relative error in our measurements of free surface and effusion data should probably not exceed 10% because most sources of systematic error are common to the free surface and effusion measurements. The most probable source of serious error in determination of α is a discrepancy between measured temperature and true temperature of the evaporating surface of a Langmuir sample. In this research the Langmuir specimens were mounted in the same position as the orifices of the effusion cells in a chamber of demonstrated uniform temperature. If the true value of α were 0.5 our free surface temperatures would have to be approximately 35° higher than the temperature inside an effusion cell placed in the same positions in order to account for the results presented in this paper. An actual temperature difference of more than

five degrees is very unlikely.

The 111 surfaces of the barium fluoride surfaces remained smooth when examined at 240 magnifications; surface roughening did not, therefore, significantly increase the effective rate of evaporation. The evaporation coefficient for this face of a barium fluoride crystal is unity within a small probable experimental error, and the evaporation coefficients for other alkaline earth fluorides are probably also unity.

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diameter 1.05 mm, length 1.30 mm
Orifice Pair 2, diameter 2.36 mm, length 1.30 mm
diameter 2.34 mm, length 1.30 mm
Orifice Pair 3, diameter 0.46 mm, length 0.75 mm
diameter 0.46 mm, length 0.77 mm.
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diameter 1.98 mm
Orifice Pair 5, diameter 0.99 mm
diameter 0.98 mm.
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Table I. Torsion-effusion sublimation pressures for barium fluoride.[†]

T, °K	P, atm	$-\left(\frac{\Delta F_T^\circ - \Delta H_{298}^\circ}{T}\right)$ cal/deg	ΔH_{298}° kcal/mole
1491	3.32×10^{-5}	39.77	89.81
1459	1.69×10^{-5}	39.96	90.17
1440	1.03×10^{-5}	40.07	90.57
1414	5.83×10^{-6}	40.22	90.74
1388	3.56×10^{-6}	40.35	90.61
1548	1.24×10^{-4}	39.42	88.71*
1520	6.90×10^{-5}	39.57	89.09*
1358	1.59×10^{-6}	40.46	90.99
1378	2.56×10^{-6}	40.39	90.92
1395	3.95×10^{-6}	40.32	90.70
1425	8.77×10^{-6}	40.10	90.21
1358	2.18×10^{-6}	40.47	90.55
1448	1.44×10^{-5}	40.03	90.03
1526	7.55×10^{-5}	39.57	89.17*
1401	5.27×10^{-6}	40.28	90.27
1305	6.10×10^{-7}	40.69	90.22
1430	9.90×10^{-4}	40.13	90.13
1388	3.80×10^{-6}	40.37	90.46
1402	5.22×10^{-6}	40.27	90.33
1293	4.68×10^{-7}	40.74	90.13
1308	6.76×10^{-7}	40.67	90.13
1318	8.70×10^{-7}	40.64	90.12
1375	3.08×10^{-6}	40.41	90.25
1359	2.13×10^{-6}	40.48	90.28
1277	3.47×10^{-7}	40.82	89.87
1261	2.52×10^{-7}	40.89	89.64
1375	3.10×10^{-6}	40.41	90.22
1384	4.20×10^{-6}	40.37	89.92
1370	2.53×10^{-6}	40.43	90.48
1350	1.73×10^{-6}	40.52	90.29
1336	1.19×10^{-6}	40.58	90.43
1455	1.44×10^{-5}	39.98	90.40
1390	4.04×10^{-6}	40.33	90.36
1431	9.22×10^{-6}	40.12	90.39
1486	2.70×10^{-5}	39.80	90.21
1486	2.95×10^{-5}	39.80	90.05
1506	3.73×10^{-5}	39.69	90.29

Table I. Continued

T, °K	P, atm	$-\left(\frac{\Delta F_T^\circ - \Delta H_{298}^\circ}{T}\right)$ cal/deg.	ΔH_{298}° kcal/mole
1535	5.86×10^{-5}	39.53	90.41
1394	4.76×10^{-6}	40.32	90.24
1413	7.52×10^{-6}	40.22	89.96
1440	1.13×10^{-5}	40.07	90.30
1464	1.99×10^{-5}	39.93	89.95
1473	2.34×10^{-5}	39.88	89.95
1548	7.11×10^{-6}	39.45	90.45
1360	2.02×10^{-5}	40.48	89.49
1523	4.93×10^{-5}	39.58	90.30
1395	4.51×10^{-6}	40.32	90.37
1501	3.66×10^{-5}	39.72	90.09
1529	5.29×10^{-5}	39.55	90.39
			ave = 90.25 ± 0.28

* Indicates where λ/d was calculated to be less than 1.0

† First 14 points taken with orifice set 1
 Middle 18 points with orifice set 2
 Last 18 points with orifice set 3

Table II. Torsion-Langmuir Sublimation Pressure for Barium Fluoride[†]

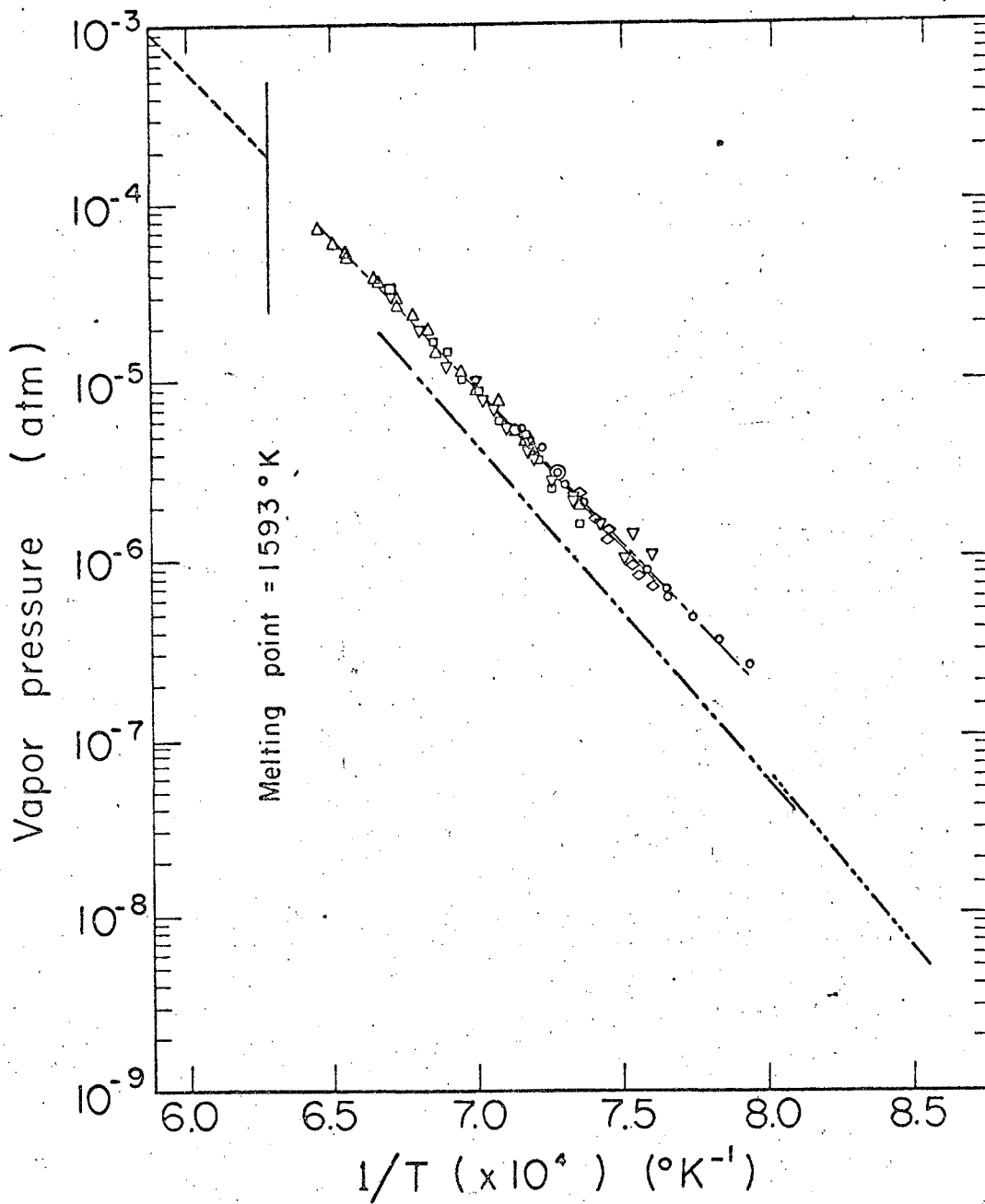
T, °K	P, atm	$-\left(\frac{\Delta F_T^\circ - \Delta H_{298}^\circ}{T}\right)$ cal/deg	ΔH_{298}° kcal/mole
1445	1.19×10^{-5}	40.05	90.42
1470	1.91×10^{-5}	39.90	90.70
1492	2.86×10^{-5}	39.75	90.33
1318	1.04×10^{-6}	40.63	89.63
1394	3.90×10^{-6}	40.31	90.69
1425	7.66×10^{-6}	40.16	90.59
1395	4.40×10^{-6}	40.31	90.42
1390	3.56×10^{-6}	40.34	90.73
1331	1.37×10^{-6}	40.59	89.73
1347	1.58×10^{-6}	40.53	90.34
1365	2.23×10^{-6}	40.45	90.51
1417	6.51×10^{-6}	40.20	90.52
1365	2.14×10^{-6}	40.45	90.63
1408	5.41×10^{-6}	40.24	90.60
1377	2.70×10^{-6}	40.40	90.71
1430	9.68×10^{-6}	40.12	90.18
1402	5.27×10^{-6}	40.28	90.33
1331	9.63×10^{-7}	40.59	90.68
1328	9.19×10^{-7}	40.60	90.61
1315	6.84×10^{-7}	40.65	90.55
1345	1.23×10^{-6}	40.53	90.88
1323	7.85×10^{-7}	40.62	90.70
1340	1.44×10^{-6}	40.55	90.16
1361	2.25×10^{-6}	40.48	90.27
1350	1.70×10^{-6}	40.52	89.34
			ave = 90.41 ± 0.35

[†] First 18 points with orifice set 4
Last 7 points with orifice set 5

Figure Caption

Figure 1. Vapor Pressure of Barium Fluoride

- — — — Knudsen experiment (this work) $\lambda/d > 1.0$
□ orifice 1, ○ orifice 2, △ orifice 3;
- Langmuir experiment (this work)
▽ orifice 4, ◇ orifice 5;
- - - - - Buff and LeBouche (extrapolated);
- — — — Green et al.;
- - - - - Bautista and Margrave.



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