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Explorations of Computational Approaches to Chemical Reactions

by

Luke W. Bertels

A dissertation submitted in partial satisfaction of the

requirements for the degree of

Doctor of Philosophy

 $\mathrm{in}$ 

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in the

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of the

University of California, Berkeley

Committee in charge:

Professor Martin P. Head-Gordon, Chair Professor Teresa L. Head-Gordon Professor Jeffrey B. Neaton

Summer 2020

## Explorations of Computational Approaches to Chemical Reactions

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#### Abstract

## Explorations of Computational Approaches to Chemical Reactions

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Luke W. Bertels

#### Doctor of Philosophy in Chemistry

University of California, Berkeley

Professor Martin P. Head-Gordon, Chair

This thesis is concerned with the application, development, and assessment of computational approaches to treat chemical reactions from an electronic structure theory perspective. In Chapter 2, the recently developed density functional,  $\omega$ B97X-V, and coupled cluster with single, double and perturbative triple excitations [CCSD(T)] are applied to study the reaction networks of excited state silicon atoms with allene and methylacetylene in the circumstellar envelope of carbon stars. In Chapter 3, another recently developed density functional,  $\omega$ B97M-V, and CCSD(T) are used to benchmark the performance of the ReaxFF reactive force field as applied to hydrogen combustion systems. In Chapter 4, an extension of third-order Møller-Plesset perturbation theory (MP3) and regularized perturbation theory ( $\kappa$ -OOMP2) is developed and used to calculate thermochemistry, barrier heights, and noncovalent interaction energies. In Chapter 5, the use of  $\kappa$ -OOMP2 optimized orbitals is extended to CCSD(T) and used to investigate the vibrational frequencies of diatomic molecules.

For Mom. You were always there cheering, even when it was painful, and I heard you every time.

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Electronic structure theory is a subfield of chemistry (though at times its practice can resemble physics, computer science, and/or mathematics) that is concerned with modeling the motions of electrons and nuclei in molecules on the quantum scale and how these translate into observable chemical phenomena: reactivity, spectra, physical properties, etc. Though as theoretical chemists we are in possession of the exact laws that govern all chemical phenomena, however application of these laws to nontrivial systems requires the use of approximations. As chemistry occurs over broad domains in both space and time, one would desire a set of approximate methods that exist on a sliding scale from highly efficient treatment of large systems to highly accurate treatment of small model systems. In contrast to empirical science, which attempts to divine the physical laws through carefully planned observations, theoretical science begins from these laws to both elucidate observations and make predictions to be evaluated. In this sense, theorists and experimentalists complement each other: theoretical models give explanation to the observations of experiments while experiments demonstrate the reality (or lack thereof) of chemical model predictions.

In this thesis, we elucidate the reactivity of a chemical system using existing models to explain experimental observations, evaluate the performance of approximate models relative to more exact methods, and develop new modifications to existing models that improve upon their accuracy without altering the overall cost.

## 1.1 The Schrödinger Equation

In 1926 the physicist Erwin Schrödinger proposed an equation to describe the dynamics of a system of quantum particles:

$$i\hbar\frac{\partial}{\partial t}\left|\Psi\right\rangle = \hat{H}\left|\Psi\right\rangle,\tag{1.1}$$

where  $\hat{H}$  is the Hamiltonian operator and  $|\Psi\rangle$  is the wavefunction representing the system. The time-independent form of this equation,

$$\hat{H} |\Psi_k\rangle = E_k |\Psi_k\rangle$$
 (1.2)

is an eigenvalue equation where the wavefunction  $|\Psi_k\rangle$  is an eigenfunction of  $\hat{H}$  with a corresponding energy eigenvalue  $E_k$ . While  $\hat{H}$  is related to the Hamiltonian in classical mechanics, in quantum mechanics the particles are described by a wavefunction that belongs to a multidimensional Hilbert space rather than points in phase space. For a system of n electrons and M nuclei, the molecular Hamiltonian, in atomic units (assumed throughout), is

$$\hat{H} = \left(-\sum_{A}^{M} \frac{1}{2m_{A}}\hat{\nabla}_{A}^{2}\right) + \left(-\sum_{i}^{n} \frac{1}{2}\hat{\nabla}_{i}^{2}\right) + \left(\sum_{A}^{M} \sum_{B>A}^{M} \frac{Z_{A}Z_{B}}{\left|\hat{\mathbf{R}}_{A} - \hat{\mathbf{R}}_{B}\right|}\right) + \left(-\sum_{i}^{n} \sum_{A}^{M} \frac{Z_{A}}{\left|\hat{\mathbf{r}}_{i} - \hat{\mathbf{R}}_{A}\right|}\right) + \left(\sum_{i}^{n} \sum_{j>i}^{n} \frac{1}{\left|\hat{\mathbf{r}}_{i} - \hat{\mathbf{r}}_{j}\right|}\right),$$

$$(1.3)$$

where  $m_A$  is the mass of nucleus A,  $Z_A$  and  $Z_B$  are the charges of nuclei A and B, respectively,  $\hat{\mathbf{R}}_A$  and  $\hat{\mathbf{R}}_B$  are the spatial coordinates of nuclei A and B, respectively, and  $\hat{\mathbf{r}}_i$  and  $\hat{\mathbf{r}}_j$  are the spatial coordinates of electrons i and j, respectively. Here the terms in the Hamiltonian are the nuclear kinetic energy operator  $(\hat{T}_N)$ , the electronic kinetic energy operator  $(\hat{T}_e)$ , the nuclei-nuclei Coulomb potential operator  $(\hat{V}_{NN})$ , the nuclei-electron Coulomb potential operator  $(\hat{V}_{Ne})$ , and the electron-electron Coulomb potential operator  $(\hat{V}_{ee})$ , respectively. The eigenfunctions of  $\hat{H}$  in position-space take the form  $\Psi(\mathbf{r}, \mathbf{R})$  where  $\mathbf{r} = {\mathbf{r}_1, \mathbf{r}_2, ..., \mathbf{r}_n}$ and  $\mathbf{R} = {\mathbf{R}_1, \mathbf{R}_2, ..., \mathbf{R}_M}$ . The squared modulus  $|\Psi(\mathbf{r}, \mathbf{R})|^2$  gives the probability density of finding electron 1 at  $\mathbf{r}_1,...,$  electron n at  $\mathbf{r}_n$ , nucleus 1 at  $\mathbf{R}_1,...,$  and nucleus M at  $\mathbf{R}_M$ . Finding these wavefunctions therefore involves solving a 3n + 3M dimensional partial differential equation. While closed-form solutions exist for hydrogen-like atoms (n = 1, M = 1), the field of electronic structure theory relies on the use of careful approximations to simplify the system in question.

In order to reduce the dimensionality of the molecular Schrödinger equation, a popular approach is to apply the Born-Oppenheimer (BO) approximation where the nuclei are treated as fixed with respect to the electrons. This is rationalized by the observation that the nuclear motion, due to greater mass of nuclei  $(m_p > 1800m_e)$ , occurs on a much longer timescale than electronic motion. Applying this approximation to Eq. 1.3, we see that the first term will go to zero and the third term will become a positive constant. This allows for the factorization of the wavefunction into an electronic term  $\Psi_e(\mathbf{r}; \mathbf{R})$  that depends parametrically on the coordinates of the nuclei and a purely nuclear term  $\Phi_N(\mathbf{R})$ :

$$\Psi(\mathbf{r}, \mathbf{R}) = \Psi_e(\mathbf{r}; \mathbf{R}) \Phi_N(\mathbf{R}). \tag{1.4}$$

With this reduction, the problem of interest becomes the electronic Schrödinger equation,

$$\hat{H}_{e}(\mathbf{R}) |\Psi_{e,k}(\mathbf{R})\rangle = E_{e,k}(\mathbf{R}) |\Psi_{e,k}(\mathbf{R})\rangle$$
(1.5)

$$\left(-\sum_{i}^{n}\frac{1}{2}\nabla_{i}^{2}-\sum_{i}^{n}\sum_{A}^{M}\frac{Z_{A}}{|\mathbf{r}_{i}-\mathbf{R}_{A}|}+\sum_{i}^{n}\sum_{j>i}^{n}\frac{1}{|\mathbf{r}_{i}-\mathbf{r}_{j}|}\right)\Psi_{e,k}(\mathbf{r};\mathbf{R})=E_{e,k}(\mathbf{R})\Psi_{e,k}(\mathbf{r};\mathbf{R}),\ (1.6)$$

where  $|\Psi_{e,k}(\mathbf{R})\rangle$  is an electronic (eigen)state of the electronic Hamiltonian  $\hat{H}_e$  with electronic energy  $E_{e,k}(\mathbf{R})$ . The field of electronic structure theory is focused on finding accurate approximation solutions to this problem. Provided a solution to the electronic problem, one can define the corresponding BO potential energy surfaces (PES) on which the nuclei move:

$$V_k(\mathbf{R}) = E_{e,k}(\mathbf{R}) + V_{NN}(\mathbf{R}). \tag{1.7}$$

Accurate description of the ground state PES (k = 0) is a prominent aspect of many chapters in this thesis. It allows for a depiction of the energy change as systems move along reaction coordinates. The gradients of this potential energy are used to determine the forces for molecular dynamics simulations. One can also substitute Eqs. 1.4 and 1.7 into Eq. 1.2 to yield the BO nuclear Schrödinger equation:

$$[T_N + V_k(\mathbf{R})] \Phi_{K,k}(\mathbf{R}) = E_{K,k} \Phi_{K,k}(\mathbf{R}).$$
(1.8)

This equation has incorporated the BO approximation by neglecting the coupling of different electronic states by  $\hat{T}_N$ . This is valid when the electronic states are well-separated in energy for a given nuclear configuration, but can break down when the energies of two states become close. While one should be diligent about this issue, it is an appropriate approximation in the context of this thesis.

The nuclear Schrödinger equation is often neglected in electronic structure theory. A common approximation is to perform a second-order Taylor expansion of the PES at its extrema in order to model the nuclei as a set of coupled quantum harmonic oscillators. Normal mode analysis allows for computation of zero-point vibrational energies which are often used to correct zero-temperature electronic energy differences. In the context of molecular dynamics (MD) simulations, nuclei are often treated classically but approaches based on Feynman path integrals[1–5] and explicit quantum treatment of relevant light nuclei[6, 7] also are studied.

From here, we will focus on the electronic Schrödinger equation, leaving implicit the e subscript and  $\mathbf{R}$  dependence of the electronic Hamiltonian, wavefunction, and energy.

## **1.2** Hartree-Fock Theory

Despite much simplification, we are still left with a partial differential equation in 3n dimensions. A reasonable approximation to the ground state solution would be to minimize the expectation value of the Hamiltonian with a trial wavefunction  $|\tilde{\Psi}\rangle$  parameterized by some coefficients **C**. As the electronic Hamiltonian is a Hermitian operator, there exists an variational principle which guarantees this expectation value bounds the true ground state energy  $E_0$  from above:

$$E_{0} \leq \min_{\mathbf{C}} \frac{\left\langle \tilde{\Psi}(\mathbf{C}) \middle| \hat{H} \middle| \tilde{\Psi}(\mathbf{C}) \right\rangle}{\left\langle \tilde{\Psi}(\mathbf{C}) \middle| \tilde{\Psi}(\mathbf{C}) \right\rangle} = \tilde{E}_{0}$$
(1.9)

In order to determine a proper guess for the form of the wavefunction, one should consider the role of the permutation operator  $\hat{P}_{ij}$ , which exchanges electrons *i* and *j*. As electrons are indistinguishable,  $\hat{P}_{ij}$  commutes with  $\hat{H}$  and therefore the eigenfunctions of  $\hat{H}$  are also eigenvalues of  $\hat{P}_{ij}$  with eigenvalue *p*. Furthermore, applying the permutation operator twice will return the exact same wavefunction:

$$\hat{P}_{ij}\hat{P}_{ij}|\Psi_k\rangle = p^2|\Psi_k\rangle = |\Psi_k\rangle.$$
(1.10)

This eigenvalue p must be either 1, corresponding to bosonic particles with integer spin, or -1, corresponding to fermionic particles with half-integer spin. As electrons are spin-1/2 particles and therefore fermions, the trial electronic wavefunction must be antisymmetric (p = -1) to particle exchange. The determinant of a square matrix is a mathematical object which obeys the property of antisymmetry with respect to exchange of two columns. Therefore a reasonable trial wavefunction is a Slater determinant, the determinant of a matrix where the rows index one-electron spin orbitals  $\chi_i$  and columns index particle coordinates  $\mathbf{x}_j$ :

$$\Phi_{S}(\mathbf{x}_{1}, \mathbf{x}_{2}, ..., \mathbf{x}_{n}) = \frac{1}{\sqrt{n!}} \det \begin{vmatrix} \chi_{1}(\mathbf{x}_{1}) & \chi_{1}(\mathbf{x}_{2}) & \cdots & \chi_{1}(\mathbf{x}_{n}) \\ \chi_{2}(\mathbf{x}_{1}) & \chi_{2}(\mathbf{x}_{2}) & \cdots & \chi_{2}(\mathbf{x}_{n}) \\ \vdots & \vdots & \ddots & \vdots \\ \chi_{n}(\mathbf{x}_{1}) & \chi_{n}(\mathbf{x}_{2}) & \cdots & \chi_{n}(\mathbf{x}_{n}) \end{vmatrix}.$$
(1.11)

We require the set of one-electron spin orbitals to be orthonormal to each other. These oneparticle spin orbitals  $\chi_i$  can be decomposed into a spatial orbital  $\psi_i$  and a spin component  $\sigma_i$ :

$$\chi_i(\mathbf{x}_j) = \psi_i(\mathbf{r}_j)\sigma_i(\omega_j). \tag{1.12}$$

Here  $\mathbf{r}_j$  and  $\omega_j$  are the spatial coordinates and spin coordinate of electron j, respectively. Typically  $\sigma_i$  is chosen to be either  $\alpha$  or  $\beta$ , the orthonormal eigenvectors of the one-particle spin operator projected along the z-axis  $\hat{s}_z$  with eigenvalues of 1/2 and -1/2, respectively.

$$\langle \sigma(\omega_j) | \sigma'(\omega_j) \rangle = \delta_{\sigma,\sigma'} \tag{1.13}$$

$$\hat{s}_{z,j} |\alpha(\omega_j)\rangle = \frac{1}{2} |\alpha(\omega_j)\rangle \tag{1.14}$$

$$\hat{s}_{z,j} \left| \beta(\omega_j) \right\rangle = -\frac{1}{2} \left| \beta(\omega_j) \right\rangle \tag{1.15}$$

In practical calculations, the spatial orbitals are expanded in a finite basis  $\{|\phi_{\mu}\rangle\}$  of N atom-centered orbitals (AOs).

$$|\psi_i((r)\rangle = \sum_{\mu}^{N} C_{\mu i} |\phi_{\mu}(r)\rangle$$
(1.16)

As  $\sigma_i$  can be either  $\alpha$  or  $\beta$ , this yields a set of n occupied spin orbitals and 2N - n virtual (unoccupied) spin orbitals.

By taking this Slater determinant  $|\Phi_S\rangle$  as a trial wavefunction and performing a variational optimization with respect to the wavefunction coefficients **C** we have arrived at the Hartree-Fock (HF) approximation:

$$E_{\rm HF} = \min_{\mathbf{C}} \left\langle \Phi_S \right| \hat{H} \left| \Phi_S \right\rangle. \tag{1.17}$$

The electronic Hamiltonian in Eq. 1.6 can be further split into one- and two-electron operators:

$$\hat{O}_{1} = -\sum_{i}^{n} \frac{1}{2} \hat{\nabla}_{i}^{2} - \sum_{i}^{n} \sum_{A}^{M} \frac{Z_{A}}{|\hat{\mathbf{r}}_{i} - \mathbf{R}_{A}|} = \sum_{i}^{n} \hat{h}(\mathbf{r}_{i})$$
(1.18)

$$\hat{O}_2 = \sum_{i}^{n} \sum_{j>i}^{n} \frac{1}{|\hat{\mathbf{r}}_i - \hat{\mathbf{r}}_j|} = \sum_{i}^{n} \sum_{j>i}^{n} \frac{1}{\hat{r}_{ij}}$$
(1.19)

The expectation value in Eq. 1.17 can be evaluated using the Slater-Condon rules[8, 9] to yield the HF energy for a given Slater determinant:

$$E_{\rm HF}[\Phi_S] = \sum_{i}^{n} \left\langle i \left| \hat{h} \left| i \right\rangle + \frac{1}{2} \sum_{i}^{n} \sum_{j}^{n} \left( \left\langle ij \right| ij \right\rangle - \left\langle ij \right| ji \right\rangle \right)$$
(1.20)

where

$$\left\langle i \middle| \hat{h} \middle| i \right\rangle = \int \chi_i^*(\mathbf{x}_1) h(\mathbf{r}_1) \chi_i(\mathbf{x}_1) d\mathbf{x}_1$$
(1.21)

$$\langle ij|ij\rangle = \iint \chi_i^*(\mathbf{x}_1)\chi_j^*(\mathbf{x}_2)\frac{1}{r_{12}}\chi_i(\mathbf{x}_1)\chi_j(\mathbf{x}_2)d\mathbf{x}_1d\mathbf{x}_2$$
(1.22)

$$\langle ij|ji\rangle = \iint \chi_i^*(\mathbf{x}_1)\chi_j^*(\mathbf{x}_2)\frac{1}{r_{12}}\chi_j(\mathbf{x}_1)\chi_i(\mathbf{x}_2)d\mathbf{x}_1d\mathbf{x}_2.$$
(1.23)

Eq. 1.22 represents the classical Coulomb interaction between two charge densities  $|\chi_i(\mathbf{x}_1)|^2$ and  $|\chi_j(\mathbf{x}_2)|^2$  while Eq. 1.23, the exchange integral, represents a purely quantum mechanical effect due to the antisymmetry constraint on the wavefunction.

Minimization of Eq. 1.20 with respect to the spin orbitals with the constraint that the spin orbitals are orthogonal can be achieved using the method of Lagrange multipliers.

$$\mathcal{L}[\{\chi_i\}] = E_{\rm HF}[\{\chi_i\}] - \sum_{i}^{n} \sum_{j}^{n} \epsilon_{ij} \left(\langle \chi_i | \chi_j \rangle - \delta_{ij}\right)$$
(1.24)

Setting the variation of  $\mathcal{L}$  to zero leads to

$$\hat{f} |\chi_i\rangle = \sum_j^n \epsilon_{ij} |\chi_j\rangle.$$
(1.25)

The choice of the canonical orbitals which zero out the off-diagonal elements of  $\epsilon_{ij}$  yields the Hartree-Fock equations,

$$\hat{f} |\chi_i\rangle = \epsilon_i |\chi_i\rangle, \quad i = 1, 2, ..., n,$$
(1.26)

where the Fock operator,

$$\hat{f}(\mathbf{x}_1) = \hat{h}(\mathbf{r}_1) + \sum_{j=1}^{n} \int \chi_j^*(\mathbf{x}_2) \frac{1}{r_{12}} \left(1 - \hat{P}_{12}\right) \chi_j(\mathbf{x}_2) d\mathbf{x}_2, \qquad (1.27)$$

include the mean-field potential that electron 1 experiences due to the other electrons. The Hartree-Fock wavefunction can then be considered a minimally-correlated approximation, as the only correlation between electrons comes from the antisymmetry constraint on the wavefunction.

Looking at Eq. 1.27, the Fock operator depends on the spin orbitals and therefore must be solved self-consistently within a basis. We begin by rewriting Eq. 1.26 in terms of spatial orbitals  $\{\psi_{i,\sigma}\}$ :

$$\hat{f}_{\sigma} |\psi_{i,\sigma}\rangle = \epsilon_{i,\sigma} |\psi_{i,\sigma}\rangle \tag{1.28}$$

where

$$\hat{f}_{\sigma}(\mathbf{r}_{1}) = \hat{h}(\mathbf{r}_{1}) + \sum_{\sigma' \in \{\alpha,\beta\}} \sum_{j}^{n_{\sigma'}} \int |\psi_{j,\sigma'}|^{2} \frac{1}{r_{12}} d\mathbf{r}_{2} - \sum_{j}^{n_{\sigma}} \int \psi_{j,\sigma}^{*} \frac{1}{r_{12}} \hat{P}_{12} \psi_{j,\sigma}(\mathbf{r}_{2}) d\mathbf{r}_{2}$$
(1.29)

and  $\sigma \in \{\alpha, \beta\}$  and  $n_{\sigma}$  is the number of electrons with spin component  $|\sigma\rangle$ . Here the Fock operator and spatial orbital are tagged with a spin label due to the action of the exchange operator on different spins. Revisiting Eq. 1.16, each spatial orbital is expanded in terms of the AO basis  $\{\phi_{\nu}\}$ :

$$|\psi_{i,\sigma}\rangle = \sum_{\nu}^{N} C_{\nu i,\sigma} |\phi_{\nu}\rangle. \qquad (1.30)$$

Here the matrices  $\mathbf{C}_{\sigma}$  are the molecular orbital (MO) coefficient matrices. The basis functions  $\{\phi_{\nu}\}$  are generally not orthogonal, with an overlap integral given by

$$S_{\mu\nu} = \langle \phi_{\mu} | \phi_{\nu} \rangle \,. \tag{1.31}$$

Substituting Eq. 1.30 into Eq. 1.28 and projecting with  $\langle \phi_{\mu} |$  yields the Roothaan equation[10], a generalized matrix eigenvalue problem:

$$\left\langle \phi_{\mu} \right| \hat{f}_{\sigma} \sum_{\nu}^{N} \left| \phi_{\nu} \right\rangle C_{\nu i,\sigma} = \left\langle \phi_{\mu} \right| \sum_{\nu}^{N} \left| \phi_{\nu} \right\rangle C_{\nu i,\sigma} \epsilon_{i,\sigma}$$
(1.32)

$$\sum_{\nu}^{N} F_{\mu\nu,\sigma} C_{\nu i,\sigma} = \sum_{\nu}^{N} S_{\mu\nu} C_{\nu i,\sigma} \epsilon_{i}$$
(1.33)

$$\mathbf{F}_{\sigma}\mathbf{C}_{\sigma} = \mathbf{S}\mathbf{C}_{\sigma}\epsilon_{\sigma}.\tag{1.34}$$

where the Fock matrix elements in the AO basis are

$$F_{\mu\nu,\sigma} = h_{\mu\nu} + \sum_{\sigma' \in \{\alpha,\beta\}} \sum_{j}^{n_{\sigma'}} \langle \mu j_{\sigma'} | \nu j_{\sigma'} \rangle - \sum_{j}^{n_{\sigma}} \langle \mu j_{\sigma} | j_{\sigma} \mu \rangle.$$
(1.35)

The AO Fock matrix elements in Eq. 1.35 can be simplified via introduction of the oneparticle density matrix (1-PDM). In HF theory, the real-space electron density  $\rho(\mathbf{r})$  can be found by integrating the square of the wavefunction over all but one set of spatial coordinates:

$$\rho(\mathbf{r}_{1}) = n \iint \cdots \int |\Phi_{s}(\mathbf{x}_{1}, \mathbf{x}_{2}, ..., \mathbf{x}_{n})|^{2} d\omega_{1} d\mathbf{x}_{2} ... d\mathbf{x}_{n}$$
$$= \sum_{\sigma \in \{\alpha, \beta\}} \sum_{i}^{n_{\sigma}} |\psi_{i,\sigma}(\mathbf{r}_{1})|^{2}.$$
(1.36)

The spin-densities  $\rho_{\sigma}(\mathbf{r})$  can be defined similarly. Transforming this into the AO basis using Eq.1.30 yields

$$\rho_{\sigma}(\mathbf{r}) = \sum_{i}^{n_{\sigma}} |\psi_{i,\sigma}(\mathbf{r})|^2 \tag{1.37}$$

$$=\sum_{\mu}^{N}\sum_{\nu}^{N}\sum_{i}^{n_{\sigma}}C_{\mu i,\sigma}C_{\nu i,\sigma}^{*}\phi_{\mu}(\mathbf{r})\phi_{\nu}^{*}(\mathbf{r})$$
(1.38)

$$=\sum_{\mu}^{N}\sum_{\nu}^{N}P_{\mu\nu,\sigma}\phi_{\mu}(\mathbf{r})\phi_{\nu}^{*}(\mathbf{r}).$$
(1.39)

Here the 1-PDM for spin  $\sigma$  is given by

$$P_{\mu\nu,\sigma} = \sum_{i}^{n_{\sigma}} C_{\mu i,\sigma} C^*_{\nu i,\sigma}.$$
(1.40)

Incorporating Eq. 1.40 into Eq. 1.35 yields

$$F_{\mu\nu,\sigma} = h_{\mu\nu} + \sum_{\lambda}^{N} \sum_{\varsigma}^{N} \langle \mu\lambda | \nu\varsigma \rangle \left( P_{\lambda\varsigma,\alpha} + P_{\lambda\varsigma,\beta} \right) - \sum_{\lambda}^{N} \sum_{\varsigma}^{N} \langle \mu\lambda | \varsigma\nu \rangle P_{\lambda\varsigma,\sigma}$$
(1.41)

where one can define the Coulomb  $(\mathbf{J})$  and exchange  $(\mathbf{K})$  matrices as

$$J_{\mu\nu} = \sum_{\lambda}^{N} \sum_{\varsigma}^{N} \langle \mu \lambda | \nu\varsigma \rangle \left( P_{\lambda\varsigma,\alpha} + P_{\lambda\varsigma,\beta} \right)$$
(1.42)

$$K_{\mu\nu,\sigma} = \sum_{\lambda}^{N} \sum_{\varsigma}^{N} \langle \mu \lambda | \varsigma \nu \rangle P_{\lambda\varsigma,\sigma}$$
(1.43)

which are in turn contractions of the spin density matrices with the electron repulsion integrals (ERIs) in the AO basis. The Hartree-Fock energy, in terms of the spin density matrices is

$$E_{\rm HF} = \frac{1}{2} \left( \text{Tr}[\mathbf{P}_{\alpha}(\mathbf{h} + \mathbf{F}_{\alpha})] + \text{Tr}[\mathbf{P}_{\beta}(\mathbf{h} + \mathbf{F}_{\beta})] \right).$$
(1.44)

The derivative of the Hartree-Fock energy with respect to a spin density matrix is the corresponding Fock matrix.

Eq. 1.41 and 1.44 provide the working equations for constructing a Fock matrix and evaluating the Hartree-Fock energy. In order to self-consistently solve the Roothaan equations, the generalized eigenvalue problem is transformed into a eigenvalue problem via orthogonalization of the AO basis. This is accomplished using the symmetric orthogonalization scheme of Löwdin[11]. Beginning with the Roothaan equation, we have

$$\mathbf{F}_{\sigma}\mathbf{C}_{\sigma} = \mathbf{S}\mathbf{C}_{\sigma}\epsilon_{\sigma}.\tag{1.45}$$

Let  $\mathbf{X}$  be the symmetric orthogonalizer of  $\mathbf{S}$  such that

$$\mathbf{X}^{\dagger}\mathbf{S}\mathbf{X} = \mathbf{I},\tag{1.46}$$

$$\mathbf{X} = \mathbf{S}^{-1/2}.\tag{1.47}$$

Defining a transformed MO coefficient matrix  $\mathbf{C}'_{\sigma}$  and substituting into Eq. 1.45 yields

$$\mathbf{C}_{\sigma} = \mathbf{X}\mathbf{C}_{\sigma}',\tag{1.48}$$

$$\mathbf{F}_{\sigma}\mathbf{X}\mathbf{C}_{\sigma}' = \mathbf{S}\mathbf{X}\mathbf{C}_{\sigma}'\epsilon_{\sigma}.$$
(1.49)

(1.50)

Finally, both sides of Eq. 1.49 are left-multiplied by the hermitian conjugate of the orthogonalizer and terms are grouped:

$$\mathbf{X}^{\dagger} \mathbf{F}_{\sigma} \mathbf{X} \mathbf{C}_{\sigma}' = \mathbf{X}^{\dagger} \mathbf{S} \mathbf{X} \mathbf{C}_{\sigma}' \epsilon_{\sigma}$$
(1.51)

$$\mathbf{X}^{\dagger}\mathbf{F}_{\sigma}\mathbf{X} = \mathbf{F}_{\sigma}^{\prime} \tag{1.52}$$

$$\mathbf{F}'_{\sigma}\mathbf{C}'_{\sigma} = \mathbf{C}'_{\sigma}\epsilon_{\sigma} \tag{1.53}$$

Eq. 1.53 is now an eigenvalue problem that can be solved by diagonalizing the transformed Fock matrix  $\mathbf{F}'_{\sigma}$ . The eigenvectors  $\mathbf{C}'_{\sigma}$  can be back-transformed to  $\mathbf{C}_{\sigma}$  via Eq. 1.48, which can in turn be used to update the density matrix  $\mathbf{P}_{\sigma}$  via Eq. 1.40. As the Fock matrix depends explicitly on the MO coefficients, these equations must be solved iteratively, a process known as the self-consistent field method. Fig. 1.1 presents the Roothaan-Hall algorithm for SCF. Several methods[12–16], including direct inversion of the iterative subspace (DIIS)[17, 18] and geometric direct minimization (GDM)[19], have been proposed to accelerate convergence of the SCF procedure and are utilized throughout this thesis. The computational scaling of solving the HF equations is formally  $\mathcal{O}(N^4)$ , though algorithms which take advantage of the



Figure 1.1: Schematic representation of the self-consistent field (SCF) procedure used to solve the Hartree-Fock equations.

sparsity of the ERIs reduce the asymptotic scaling to  $\mathcal{O}(N^3)$  due to the diagonalization of the Fock matrix.

Throughout the derivations above, objects have been tagged with a spin index  $\sigma$  referring to the eigenstates of the one-particle operator  $\hat{s}_z$ ,  $|\alpha\rangle$  and  $|\beta\rangle$ . In doing so, we have assumed a common spin axis for all of the electrons in the system, making the wavefunction an eigenstate of the total spin operator projected along the z-axis,  $\hat{S}_z$ , where

$$\hat{S}_z = \sum_{i}^{n} \hat{s}_{z,i}.$$
(1.54)

As the molecular Hamiltonian commutes with  $\hat{S}_z$ , the exact wavefunction can be chosen to be an eigenstate of both operators with eigenvalue  $(n_{\alpha} - n_{\beta})/2$ . Similarly, the square of the total spin operator,  $\hat{S}^2$ , commutes with both  $\hat{H}$  and  $\hat{S}_z$  and therefore the exact wavefunction is expressible as an eigenstate of all three operators. Here

$$\hat{S}^2 = \hat{S}_x \hat{S}_x + \hat{S}_y \hat{S}_y + \hat{S}_z \hat{S}_z \tag{1.55}$$

$$= \hat{S}_{+}\hat{S}_{-} - \hat{S}_{z} + \hat{S}_{z}\hat{S}_{z} \tag{1.56}$$

where

$$\hat{S}_{+} = \sum_{i}^{n} \hat{s}_{+,i} \tag{1.57}$$

$$\hat{S}_{-} = \sum_{i}^{n} \hat{s}_{-,i} \tag{1.58}$$

and

$$\hat{s}_{+,i} \left| \alpha(\omega_i) \right\rangle = 0 \tag{1.59}$$

$$\hat{s}_{+,i} \left| \beta(\omega_i) \right\rangle = \left| \alpha(\omega_i) \right\rangle \tag{1.60}$$

$$\hat{s}_{-,i} |\alpha(\omega_i)\rangle = |\beta(\omega_i)\rangle$$
 (1.61)

$$\hat{s}_{-,i} \left| \beta(\omega_i) \right\rangle = 0. \tag{1.62}$$

In this mutual eigenbasis, the eigenvalues of  $\hat{S}^2$  are

$$\hat{S}^2 |\Psi\rangle = \left(\frac{n_{\alpha} - n_{\beta}}{2}\right) \left(\frac{n_{\alpha} - n_{\beta}}{2} + 1\right) |\Psi\rangle.$$
(1.63)

For a Slater determinant wavefunction, the expectation value of  $\hat{S}^2$  can be determined evaluated using the Slater-Condon rules[8, 9]:

$$\left\langle \Phi_s \middle| \hat{S}^2 \middle| \Phi_s \right\rangle = \left\langle \hat{S}^2 \right\rangle = \left( \frac{n_\alpha - n_\beta}{2} \right) \left( \frac{n_\alpha - n_\beta}{2} + 1 \right) + n_\beta - \sum_i^{n_\alpha} \sum_j^{n_\beta} \left| S_{ij}^{\alpha\beta} \right|^2 \tag{1.64}$$

where

$$S_{ij}^{\alpha\beta} = \langle \psi_{i,\alpha} | \psi_{j,\beta} \rangle \tag{1.65}$$

and  $n_{\alpha} \geq n_{\beta}$ . From Eq. 1.65, we see that restricting each  $\alpha$  spin orbital to have a corresponding  $\beta$  spin orbital with the same spatial component leads the second and third terms to cancel. Therefore this approximation, referred to as restricted HF (RHF), leads to a wavefunction that is an eigenvalue of  $\hat{S}^2$ . Conversely, unrestricted HF (UHF) allows the spatial components differ between  $\alpha$  and  $\beta$  spin orbitals and therefore yields a wavefunction that is not a eigenstate of  $\hat{S}^2$ . However, the variational nature of the HF energy functional states that additional constraints on an approximate wavefunction can only increase the energy. Therefore, there exists a tension between using an RHF ansatz in order to obtain an approximate wavefunction which shares the symmetries of the true wavefunction or a UHF ansatz in order to obtain a more accurate approximate energy, known in the field as the symmetry dilemma. While the works in this thesis give preference to the unrestricted approach, this symmetry dilemma plays an important role in Ch. 4 and 5 where orbital choice is examined in the context of methods built on top of HF. Additionally, one can consider other ansatze where the  $\hat{S}_z$  and/or complex-conjugation symmetries are broken by the approximate wavefunction to increase variational flexibility (general HF (GHF) and complex HF, respectively). These different classifications are categorized by Fukutome<sup>[20]</sup> and while interesting are not further explored in the context of this thesis.

## **1.3** Correlated Wavefunction Methods

In the previous section, the HF equations were derived by assuming a single Slater determinant wavefunction (Eq. 1.17). Pursuing this derivation led to the Roothaan equations (Eq. 1.34), a set of single-particle equations where the potential felt by each electron is the averaged field of the other electrons. In this sense, the electrons under the HF approximation are minimally correlated, with the only correlation arising from the antisymmetry conditions on the wavefunction. The energy contribution arising from the correlation between electrons in the exact ground state wavefunction is called the correlation energy and is given by

$$E_{\rm corr} = E_0 - E_{HF} \tag{1.66}$$

The correlation energy for a given post-HF method is defined similarly as the difference between that method's total energy and  $E_{\rm HF}$ .

For a system with n electrons and and AO basis of dimension N, the solution to the HF equations yields a set of o = n occupied spin orbitals  $\{\chi_i\}$  used to form the ground state wavefunction  $|\Phi_0\rangle$  and v = 2N - n virtual spin orbitals  $\{\chi_a\}$  (As a point of notation, from here on occupied spin orbitals are indexed with  $\{i, j, ...\}$ , virtual spin orbitals with  $\{a, b, ...\}$ , and unspecified spin orbitals with  $\{p, q, ...\}$  and sums are condensed). With these pieces one can define a basis of orthonormal Slater determinants, the Fock space, created by excitations (replacing one or more occupied orbitals with virtual orbitals) from  $|\Phi_0\rangle$ . Within this Fock space, the creation operator,  $\hat{a}_p^{\dagger}$ , and annihilation operator,  $\hat{a}_q$ , creates an electron in  $\chi_p$  and destroys an electron in  $\chi_q$ , respectively. Excited determinants are formed via the action of excitation operators  $\hat{T}_x$ , which are grouped by level of excitation x, e.g.

$$\hat{T}_1 \left| \Phi_0 \right\rangle = \sum_{ia} t_i^a \hat{a}_a^{\dagger} \hat{a}_i \left| \Phi_0 \right\rangle = \sum_{ia} t_i^a \left| \Phi_i^a \right\rangle, \qquad (1.67)$$

$$\hat{T}_2 \left| \Phi_0 \right\rangle = \sum_{\substack{i < j \\ a < b}} t_{ij}^{ab} \hat{a}_a^{\dagger} \hat{a}_b^{\dagger} \hat{a}_j \hat{a}_i \left| \Phi_0 \right\rangle = \sum_{\substack{i < j \\ a < b}} t_{ij}^{ab} \left| \Phi_{ij}^{ab} \right\rangle.$$
(1.68)

In the above example,  $|\Phi_i^a\rangle$  is a singly-excited determinant with corresponding amplitude  $t_i^a$  and  $|\Phi_{ij}^{ab}\rangle$  is a doubly-excited determinant with corresponding amplitude  $t_{ij}^{ab}$ . The Hamiltonian operator in this Fock space is given by

$$\hat{H} = \sum_{pq} h_{pq} \hat{a}_p^{\dagger} \hat{a}_q + \frac{1}{4} \sum_{pqrs} \langle pq | | rs \rangle \hat{a}_p^{\dagger} \hat{a}_q^{\dagger} \hat{a}_s \hat{a}_r$$
(1.69)

where

$$h_{pq} = \int \chi_p^*(\mathbf{x}_1) h(\mathbf{r}_1) \chi_q(\mathbf{x}_1) d\mathbf{x}_1$$
(1.70)

$$\langle pq||rs\rangle = \langle pq|rs\rangle - \langle pq|sr\rangle \tag{1.71}$$

$$\langle pq|rs\rangle = \iint \chi_p^*(\mathbf{x}_1)\chi_q^*(\mathbf{x}_2)\frac{1}{r_{12}}\chi_r(\mathbf{x}_1)\chi_s(\mathbf{x}_2)d\mathbf{x}_1d\mathbf{x}_2 \tag{1.72}$$

are the MO one- and two-electron integrals.

## **Configuration Interaction Theory**

Given this new basis of determinants, one might consider taking as a new trial wavefunction a linear combination of all possible Slater determinants, or configurations, in the Fock space:

$$|\Psi_{\rm FCI}\rangle = \left(\hat{T}_0 + \hat{T}_1 + \hat{T}_2 + ... + \hat{T}_n\right)|\Phi_0\rangle$$
 (1.73)

$$= t_{0} |\Phi_{0}\rangle + \sum_{ia} t_{i}^{a} |\Phi_{i}^{a}\rangle + \sum_{\substack{i,j > i \\ a,b > a}} t_{ij}^{ab} |\Phi_{ij}^{ab}\rangle + \dots$$
(1.74)

This approach, known as full configuration interaction (FCI), considers all possible excitation of all electrons and therefore completely spans the space given by the number of electrons and the choice of basis. Because of this completeness, the energies obtained by this method are exact within the corresponding basis set. Additionally, consideration of all possible configurations renders FCI invariant to changes in the reference state  $|\Phi_0\rangle$ . Solving for these *t*-amplitude coefficients is achieved by constructing the Hamiltonian matrix **H** in the basis of configurations and solving for its eigenvalues and corresponding eigenvectors. While this in principle could be achieved through diagonalization, in practice approximate methods are often used to obtain the few lowest energy eigenstates due to the size of **H**[21]. While this may seem like an ideal approach to incorporate correlation into a trial wavefunction, the total number of configurations (equivalent to the rank of **H**) is  $\binom{2N}{n}$  and therefore grows factorially with the size of the system, rendering the method intractable for all but the smallest systems.

Given this limitation, one could consider truncating the sum in Eq. 1.73, as each level of excitation x adds  $\binom{o}{x}\binom{v}{x}$  configurations. For example, truncation after  $\hat{T}_2$  would yield configuration interactions with single and double excitations (CISD) and inclusion of triply excited determinants through  $T_3$  would yield configuration interaction with single, double, and triple excitation (CISDT). As the full space of configurations these truncated configuration interaction (CI) methods are no longer independent of the choice of the reference  $|\Phi_0\rangle$ . This aspect is acceptable under the assumption that the reference determinant has a large overlap with the true wavefunction. A major weakness of truncated CI approaches, however, is the lack of size consistency. A method is size consistent if a calculation on a supersystem of non-interacting systems returns the same result as individual calculations on the subsystems. For example, if one were to imagine two non-interacting molecules, a doubly excited determinant on each molecule is accurately described by CISD in the subsystems whereas in the supersystem this effective quadruply excited determinant is excluded. Size consistency is an important property for methods used to study chemical reactivity, as reactants, intermediates, and products require equal treatment for accurate assessment of reaction thermodynamics.

## Coupled Cluster Theory

Coupled cluster (CC) theory utilizes an exponential ansatz to parameterize the wavefunction:

$$|\Psi_{\rm CC}\rangle = e^{\left(\hat{T}_1 + \hat{T}_2 + \dots\right)} |\Phi_0\rangle = e^{\hat{T}} |\Phi_0\rangle.$$
(1.75)

This wavefunction is equivalent to the FCI wavefunction when all excitations are included in the exponent. Unlike truncated CI, however, coupled cluster theory is size consistent with a truncated cluster operator  $\hat{T}$ . Considering coupled cluster with single and double excitation (CCSD)[22, 23], we see that expansion of the exponential leads to infinite-order excitations as products of single and double excitations,

$$|\Psi_{\rm CC}\rangle = \left[1 + \hat{T}_1 + \hat{T}_2 + \frac{1}{2!} \left(\hat{T}_1 \hat{T}_1 + \hat{T}_1 \hat{T}_2 + \hat{T}_2 \hat{T}_1 + \hat{T}_2 \hat{T}_2\right) + \dots\right] |\Phi_0\rangle, \qquad (1.76)$$

allowing for treatment of the effective quadruple and higher excitations lacking in previous example with CISD. This exponential form allows for separability of Hamiltonian in the limit of non-interacting systems.

While this new form of the wavefunction has provided a solution to the issue of size consistency, it comes at the cost of a variational approach to finding the t-amplitudes. Plugging Eq. 1.75 into Eq. 1.9 leads to a set of equations which do not truncate at finite order. In order to solve for the t-amplitudes one begins with the similar transformed Hamiltonian,

$$\bar{H} = e^{-\hat{T}} \hat{H} e^{\hat{T}}.$$
(1.77)

From here, the *t*-amplitudes and energy are determined via projected equations:

$$\left\langle \Phi_0 \middle| \bar{H} \middle| \Phi_0 \right\rangle = E_{\rm CC} \tag{1.78}$$

$$\left\langle \Phi_{i}^{a} \middle| \bar{H} \middle| \Phi_{0} \right\rangle = 0 \tag{1.79}$$

$$\left\langle \Phi_{ij}^{ab} \middle| \bar{H} \middle| \Phi_0 \right\rangle = 0 \tag{1.80}$$

$$\left\langle \Phi_{ij\ldots}^{ab\ldots} \middle| \bar{H} \middle| \Phi_0 \right\rangle = 0. \tag{1.81}$$

In these equations  $\hat{H}$  may be simplified using the Baker-Campell-Hausdorff expansion as a linear combination of nested commutators of  $\hat{T}$  and  $\hat{H}$ ,

$$\bar{H} = \hat{H} + \left[\hat{H}, \hat{T}\right] + \frac{1}{2!} \left[ \left[\hat{H}, \hat{T}\right], \hat{T} \right] + \frac{1}{3!} \left[ \left[ \left[\hat{H}, \hat{T}\right], \hat{T} \right], \hat{T} \right] + \dots,$$
(1.82)

which terminates after a finite number of terms. This termination can be rationalized by the observation that the Hamiltonian operator in the Fock space representation is given by Eq. 1.69 contains at most two de-excitation. These projected equations are nonlinear and must be iterated over to solve for the *t*-amplitudes and energy. The computational cost of CCSD is  $\mathcal{O}(o^2v^4)$  per iteration and the inclusion of explicit triple excitations, CCSDT[24] has a cost of  $\mathcal{O}(o^3v^5)$  per iteration. This high cost for the latter has inspired many approximations to treat the effects of the triples, most ubiquitous among them being CCSD(T)[25]. CCSD(T) is a non-iterative  $\mathcal{O}(o^3v^4)$  method that gives a perturbative approximation to the triples amplitudes using the CCSD  $t_1$ - and  $t_2$ -amplitudes. For many computational chemists CCSD(T) is considered the "gold standard" in terms of the accuracy given the cost, though for practical applications with a reasonable basis set it is limited to a few tens of heavy (non-hydrogen) atoms (architecture dependent). This method is used extensively throughout this thesis to predict energies and barrier heights for reactions (Ch. 2 and 3) and benchmark more approximate levels of theory (Ch. 3 and 4). In Ch. 5, we investigate the role of the reference  $|\Phi_0\rangle$  in predicting spectroscopic properties with CCSD(T).

More in-depth derivations and discussions of CC theory, including diagrammatic approaches to derive and represent the CC equations can be found in Refs. [26] and [27].

## Perturbation Theory

Møller-Plesset (MP) theory offers another avenue to incorporate electron correlation based on Rayleigh-Schrödinger perturbation theory[28]. The Fock operator is taken as the zero-order Hamiltonian and the zero-order wavefunction is the HF ground state Slater determinant  $|\Phi_0\rangle$ . Expanding the  $\hat{H}$ ,  $|\Psi_{MP}\rangle$ , and  $E_{MP}$ 

$$\hat{H} = \hat{F} + \lambda \hat{V} \tag{1.83}$$

$$\left|\Psi_{\rm MP}\right\rangle = \left|\Phi_{0}\right\rangle + \lambda \left|\Psi^{(1)}\right\rangle + \lambda^{2} \left|\Psi^{(2)}\right\rangle + \lambda^{3} \left|\Psi^{(3)}\right\rangle + \dots$$
(1.84)

$$E = E_{\rm MP0} + \lambda E_{\rm MP1} + \lambda^2 E_{\rm MP2} + \lambda^3 E_{\rm MP3} + \dots$$
(1.85)

Plugging Eq. 1.83–1.85 into the time-independent Schrödinger equation and separating by order in the perturbation  $\lambda$  yields

$$F \left| \Phi_0 \right\rangle = E_{\text{MP0}} \left| \Phi_0 \right\rangle \tag{1.86}$$

$$\hat{F} \left| \Psi^{(1)} \right\rangle + \hat{V} \left| \Phi_0 \right\rangle = E_{\text{MP0}} \left| \Psi^{(1)} \right\rangle + E_{\text{MP1}} \left| \Phi_0 \right\rangle \tag{1.87}$$

$$\hat{F} \left| \Psi^{(2)} \right\rangle + \hat{V} \left| \Psi^{(1)} \right\rangle = E_{\text{MP0}} \left| \Psi^{(2)} \right\rangle + E_{\text{MP1}} \left| \Psi^{(1)} \right\rangle + E_{\text{MP2}} \left| \Phi_0 \right\rangle \tag{1.88}$$

$$\hat{F} |\Psi^{(3)}\rangle + \hat{V} |\Psi^{(2)}\rangle = E_{\rm MP0} |\Psi^{(3)}\rangle + E_{\rm MP1} |\Psi^{(2)}\rangle + E_{\rm MP2} |\Psi^{(1)}\rangle + E_{\rm MP3} |\Phi_0\rangle$$
(1.89)

(1.90)

By employing the convention of intermediate normalization:

÷

$$\langle \Phi_0 | \Psi_{\rm MP} \rangle = 1, \tag{1.91}$$

these equations can be solved via projection with  $\langle \Phi_0 |$  to obtain energy contributions or with excited determinants to obtain higher-order wavefunction contributions.

The zeroth-order energy,  $E_{\rm MP0}$ , is

$$E_{\rm MP0} = \left\langle \Phi_0 \right| \hat{F} \left| \Phi_0 \right\rangle \tag{1.92}$$

$$=\sum_{i}\epsilon_{i}.$$
(1.93)

Here the occupied-occupied and virtual-virtual blocks of  $\mathbf{F}$  are assumed to be diagonal [(psuedo)canonical spin orbitals]. The first-order energy,  $E_{\text{MP1}}$ , is

$$E_{\rm MP1} = \left\langle \Phi_0 \middle| \hat{V} \middle| \Phi_0 \right\rangle \tag{1.94}$$

$$= \left\langle \Phi_0 \middle| \hat{H} - \hat{F} \middle| \Phi_0 \right\rangle \tag{1.95}$$

$$= E_{\rm HF} - \sum_{i} \epsilon_i. \tag{1.96}$$

We see that the total energy to first-order is  $E_{\rm HF}$ . Turning to the first-order wavefunction contributions, we begin by expanding it in the basis of zeroth-order states:

$$\left|\Psi^{(1)}\right\rangle = \sum_{ia} t_i^a \left|\Phi_i^a\right\rangle + \sum_{\substack{i,j>i\\a,b>a}} t_{ij}^{ab} \left|\Phi_{ij}^{ab}\right\rangle + \dots$$
(1.97)

Rearranging Eq. 1.87 and substituting in Eq. 1.97 yields

$$\left(\hat{F} - E_{\rm MP0}\right) \left( \sum_{ia} t_i^a \left| \Phi_i^a \right\rangle + \sum_{\substack{i,j > i \\ a,b > a}} t_{ij}^{ab} \left| \Phi_{ij}^{ab} \right\rangle + \dots \right) = \left( E_{\rm MP1} - \hat{V} \right) \left| \Phi_0 \right\rangle.$$
(1.98)

At this point, the *t*-amplitudes can be solved for by left-projection with excited determinants. Looking at the right hand side, however, we notice that the  $E_{\rm MP1} |\Phi_0\rangle$  term becomes zero under this projection due to the orthogonality of the configuration space. Additionally, if HF orbitals are assumed, the occupied-virtual blocks of **F** are zero, meaning that  $\hat{V}$  only couples states which are separated by two excitations. This means that only the  $t_2$ -amplitudes are nonzero. Incorporating this into Eq. 1.98 and rearranging yields an expression for  $|\Psi^{(1)}\rangle$ :

$$\left\langle \Phi_{kl}^{cd} \right| \left( \hat{F} - E_{\rm MP0} \right) \sum_{\substack{i < j \\ a < b}} t_{ij}^{ab} \left| \Phi_{ij}^{ab} \right\rangle = - \left\langle \Phi_{ij}^{ab} \right| \hat{V} \left| \Phi_0 \right\rangle \tag{1.99}$$

$$\left(\left\langle \Phi_{ij}^{ab} \middle| \hat{F} \middle| \Phi_{ij}^{ab} \right\rangle - E_{\rm MP0}\right) t_{ij}^{ab} = -\left\langle \Phi_{ij}^{ab} \middle| \hat{V} \middle| \Phi_0 \right\rangle \tag{1.100}$$

$$t_{ij}^{ab} = -\frac{\langle ij||ab\rangle}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j} \tag{1.101}$$

$$\left|\Psi^{(1)}\right\rangle = -\sum_{\substack{i,j>i\\a,b>a}} \frac{\langle ij||ab\rangle}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j} \left|\Phi^{ab}_{ij}\right\rangle.$$
(1.102)

The second-order energy,  $E_{\rm MP2}$ , is found by left projecting  $\langle \Phi_0 |$  on Eq. 1.88 and substituting

in Eq. 1.102:

$$E_{\rm MP2} = \left\langle \Phi_0 \middle| \hat{F} \middle| \Psi^{(2)} \right\rangle + \left\langle \Phi_0 \middle| \hat{V} \middle| \Psi^{(1)} \right\rangle$$
(1.103)

$$= -\sum_{\substack{i < j \\ a < b}} \left\langle \Phi_0 \right| \hat{V} \left| \Psi^{(1)} \right\rangle \frac{\langle ij||ab \rangle}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j}$$
(1.104)

$$= -\sum_{\substack{i < j \\ a < b}} \frac{\langle ab || ij \rangle \langle ij || ab \rangle}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j}$$
(1.105)

$$= -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}}.$$
(1.106)

In this last expression, the prefactor of 1/4 arises to account for unrestriction of the bounds of the sum and  $\Delta_{ij}^{ab}$  is the orbital energy difference  $\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j$ . The performance of the MP2 method lies between that of HF and CCSD and is generally accurate for closed-shell energies and properties. Computation of the energy has a cost of  $\mathcal{O}(o^2v^2)$ , however the computational bottleneck is the transfomation of the ERIs from AO to MO, leading to an asymptotic cost of  $\mathcal{O}(N^5)$ . Use of the "resolution of the identity" (RI) decomposition for the two-electron integrals helps save on the memory requirements and allows for MP2 to treat systems of around 100 atoms with a moderate basis set[29, 30].

For open-shell systems, symmetry breaking in the reference can lead to poor accuracy for MP2[31–34]. A possible solution to this issue is orbital-optimized MP2 (OOMP2), where the MP2 total energy is optimized with respect to occupied-virtual rotations in the reference  $|\Phi_0\rangle[35-37]$ . This is achieved through use of the Hylleraas functional,  $J_{\rm H}$  [27, 38]:

$$J_{\rm H}[\Psi^{(1)}] = \left\langle \Psi^{(1)} \middle| \hat{F} - E_{\rm MP0} \middle| \Psi^{(1)} \right\rangle + \left\langle \Phi_0 \middle| \hat{V} \middle| \Psi^{(1)} \right\rangle + \left\langle \Psi^{(1)} \middle| \hat{V} \middle| \Phi_0 \right\rangle.$$
(1.107)

Writing this as a functional of the *t*-amplitudes and combining it with the total first-order energy leads to the MP2 Lagrangian which depends on a matrix  $\Theta$  that parameterizes occupied-virtual rotation of the MO coefficients:

$$\mathcal{L}_{MP2}[\mathbf{t}, \boldsymbol{\Theta}] = E_{MP0}[\boldsymbol{\Theta}] + E_{MP1}[\boldsymbol{\Theta}] + J_{H}[\mathbf{t}, \boldsymbol{\Theta}].$$
(1.108)

Setting the variation of Eq. 1.108 with respect to  $\Theta$  to zero optimizes the orbitals in the reference state  $|\Phi_0\rangle$  with respect to the overall MP2 energy. A thorough derivation of the working equations for OOMP2 is given in Ref. [39]. From the derivation of  $|\Psi^{(1)}\rangle$ , we observe that orbital optimization, which leads to nonzero  $F_{ia}$ , in principle results in a non-Brillouin singles contribution to the energy. In practice, however, this contribution is neglected for orbital-optimized methods as justified by [36]. While OOMP2 is able to address the limitations of MP2 with respect to open-shell systems it is not without its own limitations. From Eq. 1.106, we see that in the limit of  $\Delta_{ij}^{ab} = 0$  the MP2 correlation energy diverges.

In order to remedy this, several methods have been developed to regularize the OOMP2 energy such that in the limit of  $\Delta_{ij}^{ab} = 0$  does not diverge [39–41]. On particular method,  $\kappa$ -OOMP2[39], achieves this through an exponential regularization of the two-electron integrals:

$$E_{\kappa\text{-OOMP2}}(\kappa) = -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}} \left(1 - e^{-\kappa \Delta_{ij}^{ab}}\right)^2.$$
(1.109)

The ability of this method to remedy the failings of traditional OOMP2 has inspired its use in extensions seen in Ch. 4 and 5.

The third-order energy,  $E_{\rm MP3}$ , can be found by projecting  $\langle \Phi_0 |$  on Eq. 1.89:

$$E_{\rm MP3} = \left\langle \Phi_0 \right| \hat{F} \left| \Psi^{(3)} \right\rangle + \left\langle \Phi_0 \right| \hat{V} \left| \Psi^{(2)} \right\rangle$$
(1.110)

The use of canonical HF orbitals eliminates the first term on the right hand side of Eq. 1.110, yielding

$$E_{\rm MP3} = \left\langle \Phi_0 \middle| \hat{V} \middle| \Psi^{(2)} \right\rangle. \tag{1.111}$$

While this expression seems to require an expression for  $|\Psi^{(2)}\rangle$  to evaluate, it can be rewritten using Wigner's 2n + 1 theorem, which states that corrections to the wavefunction to the *n*th order can be used to calculate the corrections to the energy up to the (2n + 1)th order. Therefore, we may rewrite

$$E_{\rm MP3} = \left\langle \Psi^{(1)} \middle| \hat{V} \middle| \Psi^{(1)} \right\rangle \tag{1.112}$$

$$=\sum_{\substack{i
(1.113)$$

Recalling that  $\hat{V}$  couples states that differ by two spin orbitals, we are left with terms where  $\langle \Phi_{ij}^{ab} |$  and  $| \Phi_{kl}^{cd} \rangle$  share two occupied indices, two virtual indices, or one occupied and one virtual index:

$$E_{\text{MP3}} = \frac{1}{8} \sum_{ijabcd} (t_{ij}^{ab})^* \langle ab || cd \rangle t_{ij}^{cd}$$
  
+  $\frac{1}{8} \sum_{ijklcd} (t_{ij}^{ab})^* \langle kl || ij \rangle t_{kl}^{ab}$   
-  $\sum_{ijkabc} (t_{ij}^{ab})^* \langle kb || ic \rangle t_{kj}^{ac}.$  (1.114)

Here the prefactors again arise to account for the redundancy after unrestricting the sums. This expression for  $E_{\text{MP3}}$  is therefore shown to be contractions of two *t*-amplitudes with two
electron integrals. Expanding Eq. 1.114 with Eq. 1.101 yields

$$E_{\rm MP3} = \frac{1}{8} \sum_{ijabcd} \frac{\langle ij||ab\rangle \langle ab||cd\rangle \langle cd||ij\rangle}{\Delta_{ij}^{ab} \Delta_{ij}^{cd}} + \frac{1}{8} \sum_{ijklcd} \frac{\langle ij||ab\rangle \langle kl||ij\rangle \langle ab||kl\rangle}{\Delta_{ij}^{ab} \Delta_{kl}^{ab}} - \sum_{ijkabc} \frac{\langle ij||ab\rangle \langle kb||ic\rangle \langle ac||kj\rangle}{\Delta_{ij}^{ab} \Delta_{kj}^{ac}}.$$
(1.115)

The cost to compute the terms in  $E_{\text{MP3}}$  are  $\mathcal{O}(o^2v^4)$ ,  $\mathcal{O}(o^4v^2)$ , and  $\mathcal{O}(o^3v^3)$ , respectively. The convergence of the MPx series for energies is slow in practice and not theoretically guaranteed and therefore higher-order methods than MP2 are traditionally disfavored relative to methods which scale similarly ( $\mathcal{O}(o^2v^4)$  for MP3 vs iterative  $\mathcal{O}(o^2v^4)$  for CCSD, for example). The development of a modified version of MP3 to improve upon this cost-benefit trade off is presented in Ch. 4.

# **1.4 Density Functional Theory**

In the previous section we discussed the use of several correlated wavefunctions to improve upon the accuracy of HF for energies and properties at the cost of increased computational scaling and therefore limited system size. A method with similar computational scaling to HF that accounts for electron correlation is therefore an attractive proposition. Density functional theory (DFT) provides an answer to this proposition and has therefore become the most widely utilized electronic structure method in chemistry, physics, and material science.

#### **Basic Principles**

Returning to the electronic Schrödinger equation, we see that the exact ground state energy as a functional of the wavefunction is given by

$$E_0[\Psi_0] = \min_{\Psi} \left\langle \Psi \middle| \hat{H} \middle| \Psi \right\rangle.$$
(1.116)

The electronic Hamiltonian  $\hat{H}$  uniquely determines the ground state energy and properties of the system.  $\hat{H}$ , in turn, is uniquely determined by the total number of electrons n and the external potential  $v_{\text{ext}}(\mathbf{r})$ , which for a molecule is the field of the nuclei given by

$$v_{\text{ext}}(\mathbf{r}) = -\sum_{A}^{M} \frac{Z_A}{|\mathbf{r} - \mathbf{R}_A|}.$$
(1.117)

As the name density functional theory suggests, the premise of DFT is the ground state energy of a system is a functional of the real-space electron density  $\rho(\mathbf{r})$ ,

$$E_0 = E_0[\rho(\mathbf{r})], \tag{1.118}$$

where

$$\rho(\mathbf{r}) = n \iint \cdots \int |\Psi(\mathbf{x}_1, \mathbf{x}_2, ..., \mathbf{x}_n)|^2 d\omega_1 d\mathbf{x}_2 ... d\mathbf{x}_n.$$
(1.119)

Given that the ground state energy  $E_0$  is uniquely determined by n and  $v_{\text{ext}}(\mathbf{r})$ , Eq. 1.118 holds if  $\rho(\mathbf{r})$  uniquely defines n and  $v_{\text{ext}}(\mathbf{r})$ . The fact  $\rho(\mathbf{r})$  uniquely defines n is straightforward to demonstrate:

$$n = \int \rho(\mathbf{r}) d\mathbf{r} \tag{1.120}$$

Pierre Hohenberg and Walter Kohn published two important theorems in a 1964 article proving  $\rho(\mathbf{r})$  uniquely defines  $v_{\text{ext}}(\mathbf{r})$ [42]. The first of these (HKI) proves a one-to-one mapping between the ground state density and the external potential by contradiction. For two distinct Hamiltonians

$$\hat{H}^{A} = \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{A}$$

$$\hat{H}^{B} = \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{B}$$
(1.121)

with corresponding ground state wavefunctions  $|\Psi_0^A\rangle$  and  $|\Psi_0^B\rangle$ , suppose that the wavefunction of each gives the same density  $\rho(\mathbf{r})$ . Using Eq. 1.116 we may construct a statement of the variational principle for each Hamiltonian by measuring its expectation value with the wavefunction of the other:

$$\left\langle \Psi_{0}^{A} \middle| \hat{H}^{A} \middle| \Psi_{0}^{A} \right\rangle < \left\langle \Psi_{0}^{B} \middle| \hat{H}^{A} \middle| \Psi_{0}^{B} \right\rangle$$

$$\left\langle \Psi_{0}^{B} \middle| \hat{H}^{B} \middle| \Psi_{0}^{B} \right\rangle < \left\langle \Psi_{0}^{A} \middle| \hat{H}^{B} \middle| \Psi_{0}^{A} \right\rangle.$$

$$(1.122)$$

Expanding  $\hat{H}^A$  and  $\hat{H}^B$  yields

$$\left\langle \Psi_{0}^{A} \middle| \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{A} \middle| \Psi_{0}^{A} \right\rangle < \left\langle \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{A} \middle| \Psi_{0}^{B} \right\rangle$$

$$\left\langle \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{B} \middle| \Psi_{0}^{B} \right\rangle < \left\langle \Psi_{0}^{A} \middle| \hat{T}_{e} + \hat{V}_{ee} + \hat{V}_{Ne}^{B} \middle| \Psi_{0}^{A} \right\rangle,$$

$$\left\{ \Psi_{0}^{A} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{A} \right\rangle + \int v_{\text{ext}}^{A}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r} < \left\langle \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{B} \right\rangle + \int v_{\text{ext}}^{A}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r} < \left\langle \Psi_{0}^{A} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{A} \right\rangle + \int v_{\text{ext}}^{B}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r},$$

$$\left\{ \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{A} \right\rangle < \left\langle \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{B} \right\rangle,$$

$$\left\{ \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{B} \right\rangle < \left\langle \Psi_{0}^{B} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{B} \right\rangle.$$

$$(1.125)$$

Comparing the two inequalities in 1.125 demonstrates the contradiction, therefore there exists a one-to-one mapping between  $\rho(\mathbf{r})$  and  $v_{\text{ext}}(\mathbf{r})$ . HKI therefore allows for the expression of  $E_0$  as a functional of  $\rho(\mathbf{r})$ ,

$$E_0[\rho(\mathbf{r})] = \left\langle \Psi_0 \middle| \hat{T}_e + \hat{V}_{ee} \middle| \Psi_0 \right\rangle + \int v_{\text{ext}}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r}$$
(1.126)

$$= F_{\rm HK}[\rho(\mathbf{r})] + \int v_{\rm ext}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r}$$
(1.127)

where  $F_{HK}[\rho(\mathbf{r})]$  is the Hohenberg-Kohn functional and is independent of the external potential. The second Hohenberg-Kohn theorem (HKII) provides a variational principle with respect to a trial density  $\tilde{\rho}(\mathbf{r})$ :

$$E_0[\rho(\mathbf{r})] = E_{v_{\text{ext}}}[\rho(\mathbf{r})] < E_{v_{\text{ext}}}[\tilde{\rho}(\mathbf{r})]$$
(1.128)

To prove HKII, we begin by with a wavefunction  $|\tilde{\Psi}_0\rangle$  that corresponds to the ground state wavefunction of a Hamiltonian  $\tilde{\hat{H}}$  with external potential  $\tilde{v}_{\text{ext}}(\mathbf{r})$  and corresponds to a density  $\tilde{\rho}(\mathbf{r})$ . Beginning from the Eq. 1.116, we take the expectation value of the Hamiltonian of interest with  $|\tilde{\Psi}_0\rangle$ :

$$\left\langle \Psi_{0} \middle| \hat{H} \middle| \Psi_{0} \right\rangle < \left\langle \tilde{\Psi}_{0} \middle| \hat{H} \middle| \tilde{\Psi}_{0} \right\rangle$$
(1.129)

$$F_{\rm HK}[\rho(\mathbf{r})] + \int v_{\rm ext}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r} < F_{\rm HK}[\tilde{\rho}(\mathbf{r})] + \int v_{\rm ext}(\mathbf{r})\tilde{\rho}(\mathbf{r})d\mathbf{r}$$
(1.130)

$$E_0[\rho(\mathbf{r})] = E_{v_{\text{ext}}}[\rho(\mathbf{r})] < E_{v_{\text{ext}}}[\tilde{\rho}(\mathbf{r})]$$
(1.131)

HKII therefore states that the ground state energy is found via minimization with respect to trial densities  $\tilde{\rho}(\mathbf{r})$ . However, in the derivation of HKI and HKII, the density has been assumed to come from the ground state wavefunction corresponding to a Hamiltonian with a valid external potential; a condition known as *v*-representability. Unfortunately, the conditions that guarantee *v*-representability of a trial density  $\tilde{\rho}(\mathbf{r})$  are unknown and therefore a search for a density that minimizes the energy functional may result in a non-*v*-representable density, invalidating the assumption of the variational principle.

The Levy constrained search formula, however, provides a loosening of the *v*-representability constraint via conversion to an *n*-representability constraint [43]. The *n*-representability conditions for a real-space density require the density to be continuous, positive semidefinite, integrate to *n*, and come from an antisymmetrized wavefunction. While the mapping of wavefunction to density is one-to-one via integration of the modulus squared, the inverse mapping from a *n*-representable density to an antisymmetrized wavefunction is (infinitely) many-to-one, that is many antisymmetrized wavefunctions integrate to the same density. For a given *n*-representable ground state density  $\rho_n(\mathbf{r})$ , the set of antisymmetrized wavefunctions which yield  $\rho_n(\mathbf{r})$  by integration is  $\{|\Psi_k^{\rho_n}\rangle\}$ . Using one of these wavefunctions as a trial wavefunction yields

$$\left\langle \Psi_{0}^{\rho_{n}} \middle| \hat{H} \middle| \Psi_{0}^{\rho_{n}} \right\rangle \leq \left\langle \Psi_{k}^{\rho_{n}} \middle| \hat{H} \middle| \Psi_{k}^{\rho_{n}} \right\rangle \tag{1.132}$$

$$\left\langle \Psi_{0}^{\rho_{n}} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{\rho_{n}} \right\rangle + \int v_{\text{ext}}(\mathbf{r})\rho_{n}(\mathbf{r})d\mathbf{r} \leq$$

$$(1.133)$$

$$\left\langle \Psi_{k}^{\rho_{n}} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{k}^{\rho_{n}} \right\rangle + \int v_{\text{ext}}(\mathbf{r})\rho_{n}(\mathbf{r})d\mathbf{r}$$

$$\left\langle \Psi_{0}^{\rho_{n}} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{0}^{\rho_{n}} \right\rangle \leq \left\langle \Psi_{k}^{\rho_{n}} \middle| \hat{T}_{e} + \hat{V}_{ee} \middle| \Psi_{k}^{\rho_{n}} \right\rangle.$$
(1.134)

For a given  $\rho_n(\mathbf{r})$  we can then define a functional  $F[\rho_n(\mathbf{r})]$  where

$$F[\rho_n(\mathbf{r})] = \min_{\{\left|\Psi_k^{\rho_n}\right\rangle\}} \left\langle \Psi_k^{\rho_n} \right| \hat{T}_e + \hat{V}_{ee} \left|\Psi_k^{\rho_n}\right\rangle.$$
(1.135)

When  $\rho_n(\mathbf{r}) = \rho(\mathbf{r})$ , this functional is identical to  $F_{\text{HK}}(\rho(\mathbf{r})$ . Having loosened the *v*-representability constraint, the energy functional for an *n*-representable density is given by

$$E[\rho_n(\mathbf{r})] = F[\rho_n(\mathbf{r})] + \int v_{\text{ext}}(\mathbf{r})\rho_n(\mathbf{r})d\mathbf{r}$$
(1.136)

and the ground state energy

$$E_0[\rho(\mathbf{r})] = \min_{\rho_n} E[\rho_n(\mathbf{r})]$$
(1.137)

can then be thought of as a double minimization where the inner minimization gives the ground state for a given n-representable density and the outer minimization minimizes over all n-representable densities.

Having established that the ground state energy can be written as a functional of the density, the next step is finding a representation for  $E[\rho(\mathbf{r})]$ . Decomposing the  $F[\rho(\mathbf{r})]$  into its components yields

$$E[\rho(\mathbf{r})] = T_e[\rho(\mathbf{r})] + V_{ee}[\rho(\mathbf{r})] + \int v_{ext}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r}.$$
(1.138)

The electron-electron interaction functional can further be broken down into classical  $(J[\rho(\mathbf{r})])$ and non-classical  $(K[\rho(\mathbf{r})])$  components:

$$J[\rho(\mathbf{r})] = \frac{1}{2} \iint \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r} d\mathbf{r}'$$
(1.139)

$$K[\rho(\mathbf{r})] = V_{ee}[\rho(\mathbf{r})] - J[\rho(\mathbf{r})].$$
(1.140)

Plugging these definitions into Eq. 1.138 yields

$$E[\rho(\mathbf{r})] = T_e[\rho(\mathbf{r})] + \frac{1}{2} \iint \frac{\rho(\mathbf{r})\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r} d\mathbf{r}' + K[\rho(\mathbf{r})] + \int v_{\text{ext}}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r}.$$
 (1.141)

At this point  $T_e[\rho(\mathbf{r})]$  and  $K[\rho(\mathbf{r})]$  are the only terms without a known exact form and must therefore be approximated. An accurate approximate representation of  $T_e[\rho(\mathbf{r})]$  is especially important for chemical applications as it's contribution is of the same magnitude as the energy. The Thomas-Fermi model[44, 45], which is exact for the uniform electron gas (UEG), has the following form for  $T_e[\rho(\mathbf{r})]$ :

$$T_e[\rho(\mathbf{r})] = \frac{3}{10} \left(3\pi^2\right)^{2/3} \int \rho(\mathbf{r})^{5/3} d\mathbf{r}.$$
 (1.142)

Given the non-uniformity of the electron density within molecules, the Thomas-Fermi model understandably breaks down and is unable to properly describe chemical bonding. Accurate description of the kinetic energy functional is a yet unaccomplished task in the application of orbital-free DFT to the chemistry of atoms and molecules.

In 1965 Kohn and Sham[46] developed a method to circumvent the kinetic energy functional through the introduction of a Slater determinant of n spin orbitals. This Kohn-Sham (KS) determinant describes a fictitious non-interaction ( $V_{ee} = 0$ ) system of electrons and integrates to give the same density  $\rho(\mathbf{r})$  as the exact wavefunction. The kinetic energy  $T_{e,\text{KS}}$ , associated with the KS determinant is computed in a nearly identical fashion to the Hartree-Fock:

$$T_{e,\mathrm{KS}} = \sum_{i}^{n} \left\langle \chi_{i} \right| - \frac{1}{2} \hat{\nabla}_{i}^{2} \left| \chi_{i} \right\rangle.$$
(1.143)

 $T_{e,\text{KS}}$  is a functional of  $\rho(\mathbf{r})$  because the KS determinant is the solution to a modified Schrödinger equation whose external potential  $v_{\text{ext}}(\mathbf{r})$  maps to  $\rho(\mathbf{r})$ . This allows us to write the exchange-correlation functional  $E_{xc}[\rho(\mathbf{r})]$  as the sum of non-classical electron-electron interaction  $K[\rho(\mathbf{r})]$  and the difference between the exact and KS kinetic energies:

$$E_{xc}[\rho(\mathbf{r})] = K[\rho(\mathbf{r})] + T_e[\rho(\mathbf{r})] - T_{e,\mathrm{KS}}[\rho(\mathbf{r})].$$
(1.144)

The overall KS energy functional is given by

$$E_{\rm KS}[\rho(\mathbf{r})] = T_{e,\rm KS}[\rho(\mathbf{r})] + J[\rho(\mathbf{r})] + \int v_{\rm ext}(\mathbf{r})\rho(\mathbf{r})d\mathbf{r} + E_{xc}[\rho(\mathbf{r})].$$
(1.145)

While Eq. 1.145 is formally exact, in practice approximations are made to the form of the exchange-correlation functional  $E_{xc}[\rho(\mathbf{r})]$ , which is further broken down into exchange  $(E_x[\rho(\mathbf{r})])$  and correlation  $(E_c[\rho(\mathbf{r})])$  functionals.

In analogy to the Hartree-Fock equations (Eq. 1.28), the Kohn-Sham equations can be recast as a set of effective one-electron equations

$$\hat{f}_{\sigma}^{\text{KS}} |\psi_{i,\sigma}\rangle = \epsilon_{i,\sigma} |\psi_{i,\sigma}\rangle \tag{1.146}$$

where the Fock operator is of the form

$$f_{\sigma}^{\rm KS}(\mathbf{r}) = h(\mathbf{r}) + \sum_{\sigma' \in \{\alpha,\beta\}} \int \frac{\rho_{\sigma'}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}' + v_{xc,\sigma}(\mathbf{r})$$
(1.147)

and the exchange-correlation potential

$$v_{xc,\sigma}(\mathbf{r}) = \frac{\delta E_{xc}[\rho_{\alpha}(\mathbf{r}), \rho_{\beta}(\mathbf{r})]}{\delta \rho_{\sigma}(\mathbf{r})}$$
(1.148)

is the functional derivative of the exchange-correlation energy with respect to the density. Expansion of Eq. 1.146 in a basis of AOs yields a generalized eigenvalue problem that is the KS equivalent of the Roothaan equations in HF theory (Eq. 1.34):

$$\mathbf{F}_{\sigma}^{\mathrm{KS}}\mathbf{C}_{\sigma} = \mathbf{S}\mathbf{C}_{\sigma}\epsilon_{\sigma} \tag{1.149}$$

where

$$\mathbf{F}_{\sigma} = \mathbf{h} + \mathbf{J} + \mathbf{V}_{xc,\sigma} \tag{1.150}$$

and

$$V_{xc,\mu\nu,\sigma} = \langle \phi_{\mu} | \, \hat{v}_{xc,\sigma} \, | \phi_{\nu} \rangle \,. \tag{1.151}$$

The 1PDMs,  $\mathbf{P}_{\sigma}$ , for KS-DFT are defined identically to those for HF (Eq. 1.40) and bear an obvious connection to the spin densities  $\rho_{\sigma}(\mathbf{r})$  (Eq. 1.39. The KS-DFT energy in terms of these density matrices is

$$E_{\rm KS} = {\rm Tr}\left[\left(\mathbf{P}_{\alpha} + \mathbf{P}_{\beta}\right)\left(\mathbf{h} + \mathbf{J}\right)\right] + E_{xc}\left[\mathbf{P}_{\alpha}, \mathbf{P}_{\beta}\right]. \tag{1.152}$$

From this expression one notices that the Fock matrices  $\mathbf{F}_{\sigma}$  are the matrix derivative of  $E_{\text{KS}}$  with respect to the corresponding 1PDMs,  $\mathbf{P}_{\sigma}$ . The same SCF approach is used to optimize the KS orbitals and densities as are used in optimizing the HF orbitals, albeit with a modified Fock matrix build and energy computation. Generally speaking the quantities  $E_{xc}[\mathbf{P}_{\alpha}, \mathbf{P}_{\beta}]$  and  $\mathbf{V}_{xc}$  are expressed as integrals over real-space evaluated via quadrature on a grid. The asymptotic computational scaling of KS-DFT is  $\mathcal{O}(N^3)$  due to the diagonalization of  $\mathbf{F}_{xc}$  and can in principle be applied to systems of hundreds of atoms.

A more in-depth derivation of both Hohenberg-Kohn and Kohn-Sham DFT can be found in Ref. [47].

#### Modern Density Functional Theory

Over the last 30 years Kohn-Sham DFT (hereafter referred to as DFT unless specified) has become the most widely used tool in electronic structure theory. This has been accompanied by a proliferation of approximate exchange-correlation functionals, both by *ab initio* approaches, which design functionals to meet criteria of the exact exchange-correlation functional, and semi-empirical approaches, which fit a parameterized functional to higher-level computational or experimental data. The levels of approximation used in defining these functional are arranged in a hierarchy, originally referred as "Jacob's Ladder" by Perdew[48], where ascending the ladder adds ingredients to the exchange-correlation functional and brings one closer to the "heaven" of chemical accuracy (typically thought of as errors less than 1

kcal mol<sup>-1</sup>). Exchange-correlation functionals are split into an exchange functional  $E_x$  and a correlation functional  $E_c$  which assume the general forms

$$E_x = \sum_{\sigma \in \{\alpha,\beta\}} \int e_{x,\sigma}(\mathbf{r}) d\mathbf{r}$$
(1.153)

$$E_c = \int \left(\rho_\alpha(\mathbf{r}) + \rho_\beta(\mathbf{r})\right) \epsilon_c(\mathbf{r}) d\mathbf{r}$$
(1.154)

where  $e_{x,\sigma}(\mathbf{r})$  is the exchange energy density for electrons of spin  $\sigma$  and  $\epsilon_c(\mathbf{r})$  is the correlation energy per electron. From these definitions we observe that exchange interactions only occur between same-spin electrons while correlation interaction occur between both same-spin and opposite-spin electrons. From this point, in our notation we suppress the dependence of the spin-densities on the real-space coordinates  $\mathbf{r}$ .

Local spin-density approximation (LSDA) functionals occupy the first rung of Jacob's Ladder and are functionals of the spin-densities  $\rho_{\sigma}(\mathbf{r})$ 

$$E_x^{\text{LSDA}}[\rho_{\alpha}, \rho_{\beta}] = \sum_{\sigma \in \{\alpha, \beta\}} \int e_{x,\sigma}(\rho_{\sigma}) \, d\mathbf{r}$$
(1.155)

$$E_{c}^{\text{LSDA}}[\rho_{\alpha},\rho_{\beta}] = \int \left(\rho_{\alpha} + \rho_{\beta}\right) \epsilon_{c} \left(\rho_{\alpha},\rho_{\beta}\right) d\mathbf{r}.$$
(1.156)

The LSDA exchange functional has an analytical form, known as the Slater-Dirac exchange functional  $E_x^{\rm S}[49]$  that is exact for the UEG:

$$E_x^{\rm S}[\rho_\alpha, \rho_\beta] = -C_x \sum_{\sigma \in \{\alpha, \beta\}} \int \rho_\sigma^{4/3} d\mathbf{r}.$$
 (1.157)

The correlation functional for LSDA does not have an exact functional form and therefore approximations are made by fitting to quantum Monte-Carlo data[50], yielding among others the VWN5[51], PZ81[52], and PW92[53] correlation functionals. While able to predict binding behavior in molecules not captured by Thomas-Fermi DFT, LSDA has been shown to systematically overbind molecules relative to atoms. The inaccuracy of LSDA for describing molecular systems should come as little surprise, as the method is parameterized for the UEG and electrons in molecules are not uniformly distributed.

A method to improve upon the purely local nature of LSDA involves including information about the gradient of the spin densities  $\nabla \rho_{\sigma}$  in order to better describe semi-local exchange and correlation effects. This generalized gradient approximation (GGA) is the second rung of Jacob's Ladder and has exchange and correlation functionals of the form

$$E_x^{\text{GGA}}[\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^2,|\nabla\rho_{\beta}|^2] = \sum_{\sigma\in\{\alpha,\beta\}} \int e_{x,\sigma} \left(\rho_{\sigma},|\nabla\rho_{\sigma}|^2\right) d\mathbf{r}$$
(1.158)

$$E_{c}^{\text{GGA}}[\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^{2},|\nabla\rho_{\beta}|^{2},\nabla\rho_{\alpha}\cdot\nabla\rho_{\beta}] = \int (\rho_{\alpha}+\rho_{\beta}) \epsilon_{c} (\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^{2},|\nabla\rho_{\beta}|^{2},\nabla\rho_{\alpha}\cdot\nabla\rho_{\beta}) d\mathbf{r}.$$
(1.159)

Popular GGA exchange-correlation functionals include the *ab initio* PBE (PBE exchange, PBE correlation)[54] and the semi-empirical BLYP (B88 exchange, LYP correlation)[55, 56].

Further improvement upon GGA exchange-correlation is achieved through the metageneralized gradient approximation (meta-GGA), the third rung on Jacob's ladder. Increased semi-local information is included by incorporating either the kinetic energy spin-densities,

$$\tau_{\sigma} = \sum_{i}^{n_{\sigma}} |\nabla \psi_{\sigma,i}|^2 \tag{1.160}$$

or the Laplacians of the spin-densities  $\nabla^2 \rho_{\sigma}$  in the exchange and correlation functionals:

$$E_{x}^{\text{meta-GGA}}[\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^{2},|\nabla\rho_{\beta}|^{2},\tau_{\alpha},\tau_{\beta}] = \sum_{\sigma\in\{\alpha,\beta\}} \int e_{x,\sigma} \left(\rho_{\sigma},|\nabla\rho_{\sigma}|^{2},\tau_{\sigma}\right) d\mathbf{r} \qquad (1.161)$$
$$E_{c}^{\text{meta-GGA}}[\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^{2},|\nabla\rho_{\beta}|^{2},\nabla\rho_{\alpha}\cdot\nabla\rho_{\beta},\tau_{\alpha},\tau_{\beta}] = \int \left(\rho_{\alpha}+\rho_{\beta}\right) \epsilon_{c} \left(\rho_{\alpha},\rho_{\beta},|\nabla\rho_{\alpha}|^{2},|\nabla\rho_{\beta}|^{2},\nabla\rho_{\alpha}\cdot\nabla\rho_{\beta},\tau_{\alpha},\tau_{\beta}\right) d\mathbf{r}. \qquad (1.162)$$

Popular meta-GGA exchange-correlation functionals include the *abinitio* TPSS[57] and SCAN[58] functionals and the semi-empirical M06-L[59] and MN15-L[60] functionals.

While the addition of terms into the exchange-correlation energy densities allows for more accurate treatment of chemical systems, semi-local density functionals exhibit three weakness that limit their universal applicability and accuracy: self-interaction error (SIE), poor treatment of long-range correlations, and failure to treat strongly correlated systems.

A simple to understand manifestation of SIE in DFT is seen when considering the neutral hydrogen atom. For this one electron system the exact electron-electron interaction term in the Hamiltonian,  $\hat{V}_{ee} = \hat{J} + \hat{K}$ , should be identically zero. Under the HF approximation, the Coulomb and exchange terms are equal in magnitude and opposite in sign and therefore cancel to give zere electron-electron interaction. For DFT, however, the non-classical part of  $\hat{V}_{ee}$  is folded into the exchange-correlation term and  $J[\rho] + E_{xc}[\rho]$  is no longer identically zero. This means that the electrons in DFT "feel" themselves, contributing to the charge-delocalization error observed in DFT.

A possible solution to mediate this self-interaction error is the inclusion of some fraction of HF exact exchange,  $E_x^{\text{HF}}$ , into the exchange functional.

$$E_x^{\rm GH} = c_x E_x^{\rm HF} + (1 - c_x) E_x^{\rm DFT}$$
(1.163)

These global hybrid (GH) density functional theory approaches belong to the fourth rung of Jacob's ladder and include explicit dependence on the occupied KS spin orbitals. Popular GH GGA exchange-correlation functionals include the B3LYP[61], B97[62], and PBE0[63] functionals, while popular GH meta-GGA exchange-correlation functionals include TPSSh[64], M06-2X[65], and MN15[66].

A more rigorous attempt at including HF exact exchange into the exchange functional is through the use of range-separated hybrid (RSH) functionals. In this approach, the Coulomb

interaction is separated into a long-range (lr) component and a short-range (sr) component through use of the error function and its complement, respectively:

$$\frac{1}{r_{12}} = \frac{\operatorname{erf}(\omega r_{12})}{r_{12}} + \frac{\operatorname{erfc}(\omega r_{12})}{r_{12}}$$
(1.164)

where the  $\omega$  parameter determines how quickly the HF exact exchange ramps up. Within the RSH formalism the DFT exchange is similarly separated into long-range and shortrange components; the long-range component usually being discarded such that the in the long-range SIE cancels. The semi-local nature of the DFT exchange makes it better suited for short range treatment, while a fraction of short-range HF exact exchange may also be included:

$$E_x^{\text{RSH}} = c_{x,\text{sr}} E_{x,\text{sr}}^{\text{HF}} + E_{x,\text{lr}}^{\text{HF}} + (1 - c_{x,\text{sr}}) E_{x,\text{sr}}^{\text{DFT}}$$
(1.165)

Popular examples of RSH exchange-correlation functionals include the CAM-B3LYP[67],  $\omega$ B97X-D[68], and M11[69] functionals.

Addressing the lack of long-range correlations in DFT is necessary for application to weakly-bound systems. One proposed remedy for this issue is the use of dispersion-corrections that add an analytic potential function of the internuclear distance to the DFT energy. Grimme's empirical DFT-D methods[70–73] employ pre-computed coefficients for pairs of nuclei trained on higher levels of theory. The potential for the recent DFT-D4 dispersion correction[73] takes the form

$$E_{\rm D4} = -\sum_{B>A} \sum_{k=6,8} s_k \frac{C_{AB}^{(k)}}{R_{AB}^k} f_{\rm damp,k}(R_{AB})$$
(1.166)

where  $s_k$  is a functional dependent scaling factor,  $C_{AB}^{(k)}$  are the pairwise coefficients, and  $f_{\text{damp},k}(R_{AB})$  is a damping function that smoothly switches of the correction when the internuclear distance  $R_{AB}$  is small. Many forms of this damping function have been investigated[72–77]. The exchange-hole dipole moment (XDM) [78, 79] and Tkatchenko-Scheffler (TS) [80] compute these  $C^{(k)}$  coefficients during a calculation from the electronic structure.

A less empirical method of introducing dispersion effects into DFT is the use of explicitly non-local correlation (NLC) functionals. These functionals couple the densities at points  $\mathbf{r}$  and  $\mathbf{r}'$  via the use of a NLC kernel  $\Phi(\mathbf{r}, \mathbf{r}')$ :

$$E_{\rm NLC}[\rho] = \iint \rho(\mathbf{r}) \Phi(\mathbf{r}, \mathbf{r}') \rho(\mathbf{r}') d\mathbf{r} d\mathbf{r}'.$$
(1.167)

While  $E_{\rm NLC}$  is expensive to evaluate due to the double quadrature, NLC functionals have seen increasing use, especially the two-parameter Vydrov-Van Voorhis kernel(VV10)[81]. Mardirossian and Head-Gordon have recently designed three minimally-parameterized semiempirical functionals that are trained with this NLC functional:  $\omega$ B97X-V[82], a RSH GGA, B97M-V[83], a semi-local meta-GGA, and  $\omega$ B97M-V[84], a RSH meta-GGA. Extensive assessment of these and other functionals on an enormous composite database has demonstrated these functionals to be the best performing in their corresponding categories[85]. The accuracy of these functionals has inspired their employment in this thesis to characterize the critical points along reactive pathways (Ch. 2 and 3), to benchmark the performance of reactive molecular dynamics models (Ch. 3), and as generators of reference orbitals for correlated calculations (Ch. 5).

Double-hybrid DFT methods, the fifth rung of Jacob's Ladder, provide a third way to incorporate missing correlation in semi-local DFT. The functionals incorporate a scaled secondorder perturbative correction on top of the KS Slater determinant and therefore depend on the virtual KS spin orbitals { $\chi_a$ }. Popular functionals in this rung include B2PLYP[86], XYG3[87], and  $\omega$ B97M(2) [88].

Despite corrections to the SIE and dispersion, existing density functionals still struggle to describe strongly correlated systems. This is by virtue of the single-determinant nature of both HF and DFT to describe systems which are inherently multireference. In passing, we mention the use of Hubbard model-like correction[89, 90] and multiconfigurational[91, 92] attempts to solve these problems, though this is still an area where there is much to be improved.

# 1.5 Outline

#### Chapter 2

Mechanisms of formation of complex organosilicon molecules have remained elusive until recently, despite the ubiquity of such species in the circumstellar environments of carbon rich stars. These organosilicon molecules, among them precursors to silicon carbide grains, account for up to 80% of the carbon and silicon budgets in the interstellar medium as ejecta from the circumstellar shells. In this work we provide a combined computational, experimental, and modeling approach to study the synthesis of cyclic silicon tricarbide species  $(c-SiC_3)$ , which has been observed in the circumstellar envelope of the asymptotic giant branch (AGB) star IRC+10216. Both experimental and computational approaches suggest reactions of hydrocarbons methylacetylene (CH<sub>3</sub>CCH) and allene (H<sub>2</sub>CCCH<sub>2</sub>) with excited state silicon atoms  $(Si(^1D))$ , all species present in the inner envelope of IRC+10216, lead to the formation of  $SiC_3H_2$  molecules via molecular hydrogen loss. Further computational studies suggest that these SiC<sub>3</sub>H<sub>2</sub> molecules can undergo a photodissociation-mediated dehydrogenation to form bare  $c-SiC_3$ . This mechanism of the formation of silicon tricarbide via a two-step neutral-neutral reaction followed by photodissociation offers a competing explanation for the origin of organosilicon molecules in the circumstellar medium. This work has been published in The Proceedings of the National Academy of Sciences of the United States of America [93].

#### Chapter 3

A thorough understanding of the kinetics and dynamics of combusting mixtures is of considerable interest, especially in regimes beyond the reach of current experimental validation. The ReaxFF reactive force field method has provided a way to simulate large-scale systems of hydrogen combustion via a parameterized potential that can simulate bond breaking. This modeling approach has been applied to hydrogen combustion, as well as myriad other reactive chemical systems. In this work, we benchmark the performance of several common parameterizations of this potential against higher-level quantum mechanical (QM) approaches. We demonstrate instances where these parameterizations of the ReaxFF potential fail both quantitatively and qualitatively to describe reactive events relevant for hydrogen combustion systems. This work has been published in *The Journal of Physical Chemistry A*[94].

#### Chapter 4

We develop and test methods that include second and third-order perturbation theory (MP3) using orbitals obtained from regularized orbital-optimized second-order perturbation theory,  $\kappa$ -OOMP2, denoted as MP3: $\kappa$ -OOMP2. Testing MP3: $\kappa$ -OOMP2 shows RMS errors that are 1.7 to 5 times smaller than MP3 across 7 data sets. To do still better, empirical training of the scaling factors for the second- and third-order correlation energies and the regularization parameter on one of those data sets led to an unregularized scaled ( $c_2 = 1.0$ ;  $c_3 = 0.8$ ) denoted as MP2.8: $\kappa$ -OOMP2. MP2.8: $\kappa$ -OOMP2 yields significant additional improvement over MP3: $\kappa$ -OOMP2 in 4 of 6 test data sets on thermochemistry, kinetics, and noncovalent interactions. Remarkably, these two methods outperform coupled cluster with singles and doubles in 5 of the 7 data sets considered, at greatly reduced cost (no  $\mathcal{O}(N^6)$  iterations). This work has been in *The Journal of Physical Chemistry Letters*[95].

#### Chapter5

While CCSD(T) with spin-restricted Hartree-Fock (RHF) orbitals has long been lauded for its ability to accurately describe closed-shell interactions, the performance of CCSD(T) on open-shell species is much more erratic, especially when using a spin-unrestricted HF (UHF) reference. Previous studies have shown improved treatment of open-shell systems when a non-HF set of molecular orbitals, like Brueckner or Kohn-Sham density functional theory (DFT) orbitals, is used as a reference. Inspired by the success of regularized orbital-optimized second-order Møller-Plesset perturbation theory ( $\kappa$ -OOMP2) orbitals as reference orbitals for MP3, we investigate the use of  $\kappa$ -OOMP2 orbitals and various DFT orbitals as reference orbitals for CCSD(T) calculations of the corrected ground-state harmonic vibrational frequencies of a set of 36 closed-shell (29 neutrals, 6 cations, 1 anion) and 59 open-shell diatomic species (38 neutrals, 15 cations, 6 anions). The aug-cc-pwCVTZ basis set is used for all calculations. The use of  $\kappa$ -OOMP2 orbitals in this context alleviates difficult cases observed for both UHF orbitals and OOMP2 orbitals. Removing 2 multireference systems and 12 systems with ambiguous experimental data leaves a pruned data set. Overall performance on the pruned data set highlights CCSD(T) with a B97 orbital reference (CCSD(T):B97), CCSD(T) with a  $\kappa$ -OOMP2 orbital reference (CCSD(T): $\kappa$ -OOMP2), and CCSD(T) with a B97M-rV orbital reference (CCSD(T):B97M-rV) with RMSDs of 8.48 cm<sup>-1</sup> and 8.50 cm<sup>-1</sup>, and 8.75 cm<sup>-1</sup> respectively, outperforming CCSD(T):UHF by nearly a factor of 5. Moreover, the performance on the closed- and open-shell subsets show these methods are able to treat open-shell and closed-shell systems with comparable accuracy and robustness. The use of  $\kappa$ -OOMP2 orbitals has also proven useful in diagnosing multireference character that can hinder the reliability of CCSD(T).

# Chapter 2

# Modeling of Gas Phase Formation of $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

This chapter was adapted with permission from Yang T.; Bertels, L.; Dangi, B. B.; Li, X.; Head-Gordon, M.; Kaiser, R. I. Gas phase formation of c-SiC<sub>3</sub> molecules in the circumstellar envelope of carbon stars. *Proc. Natl. Acad. Sci. U.S.A.* **2019**, 116, 14471-14478.

# 2.1 Introduction

Astronomical observations of the carbon-rich asymptotic giant branch (AGB) star IRC+10216 (CW Leo), coupled with astrochemical modeling and gas-phase reactions networks of ion-molecule[96–101] and neutral-neutral[99, 100, 102–106] reactions has furthered the understanding of the dynamic chemical environment of carbon-rich circumstellar envelopes. However for many silicon carbide species, including cyclic silicion dicarbide (c-SiC<sub>2</sub>)[107], and bicyclic silicon tricarbide (c-SiC<sub>3</sub>), and linear silicion tetracarbide (SiC<sub>4</sub>)[108], these modeling efforts underestimate their fractional abundance relative to observation by up to two orders of magnitude[109, 110]. From this discrepancy, it is clear that the mechanisms of organosilicon formation in the circumstellar medium are incomplete[110–112], despite the roles these molecules have in furnishing silicon and carbon in the interstellar medium[113, 114] and in the formation of silicon carbide grains.

Current astrochemical models propose that the initial carbon silicon coupling occurs in the inner envelope of carbon stars. In this region, shocks initiated by pulsations of the central star give rise to temperatures in excess of 3500 K[114] and generate metastable species such as silvlidyne radicals and excited state silicon atoms[115, 116]. These highly reactive fragments can then react with hydrocarbon molecules to form parent organosilicon species[115, 117]. From the inner envelope, stellar winds are proposed to drive these organosilicon species to the outer envelope. In this region, these parent organosilicon molecules can be photolyzed Chapter 2. Modeling of Gas Phase Formation of  $c-SiC_3$  Molecules in the Circumstellar Medium of Carbon Stars 31

to silicon carbides by the interstellar UV field, with carbon-hydrogen bonds (roughly 400 kJ  $\text{mol}^{-1}$ ) being preferentially cleaved over silicon-carbon double bonds (450-550 kJ  $\text{mol}^{-1}$ ). As these mechanisms have not been confirmed to date, deeper understanding of the molecular dynamics of circumstellar envelopes is required to concretely explain the genesis of interstellar silicon carbides.

In this work, we combine high-level *ab initio* calculations with crossed molecular beam experiments and astrochemical modeling to provide a mechanism for the formation of silicon tricarbide (c-SiC<sub>3</sub>) observed in the circumstellar medium of the carbon star IRC+10216. The reaction of excited state silicon atoms (Si(<sup>1</sup>D)) with methylacetylene (CH<sub>3</sub>CCH) and allene (H<sub>2</sub>CCCH<sub>2</sub> to form SiC<sub>3</sub>H<sub>2</sub> isomers (reaction I) was explored computationally and experimentally under single collision conditions, mimicking the environment of the inner envelope of IRC+10216. Further photolysis of these parent SiC<sub>3</sub>H<sub>2</sub> molecules in the outer envelope to form bicyclic silicon tricarbide (c-SiC<sub>3</sub>) via dehydrogenation (reaction II) was explored computationally.

$$\operatorname{Si}(^{1}\mathrm{D}) + \mathrm{C}_{3}\mathrm{H}_{4} \longrightarrow \operatorname{Si}\mathrm{C}_{3}\mathrm{H}_{2} + \mathrm{H}_{2}$$
 (I)

$$\operatorname{SiC}_{3}\operatorname{H}_{2} + h\nu \longrightarrow \operatorname{c-SiC}_{3} + \operatorname{H}_{2}$$
 (II)

# 2.2 Methods

A detailed accounting of the materials and methods used in the experimental study can be found in Ref. 93.

Structures for the reactants, intermediates, and products were obtained via geometry optimizations and frequency calculations using the  $\omega B97X-V$  [118] density functional and the cc-pVTZ basis set [119, 120]. This functional is known to be among the most accurate density functionals for thermochemistry and reaction barrier heights [85]. Transition state structures were calculated using the freezing string method [121] to generate an initial structure and Hessian which were then refined by a transition state search using the partitioned-rational function optimization eigenvector following method [122] and followed by a frequency calculation. These calculations were also carried out at the  $\omega B97X-V/cc-pVTZ$  level of theory. Vibrational analysis was used to confirm that the minima have no imaginary frequencies and the transition states have only one imaginary frequency each, as well as to calculate harmonic zero-point energy corrections for the structures. Density functional theory calculations were all carried out using an ultra-fine integration grid consisting of 99 radial points and 590 angular points. To further improve the accuracy of the results for relative energies and barrier heights, single point energy calculations were performed at the minima and transition states. The objective was to approach the complete basis set (CBS) limit using coupled cluster with single, double, and perturbative triple excitations (CCSD(T))[25]. To this end, CCSD(T)/cc-pVTZ calculations using a frozen core approximation were combined with second-order Møller-Plesset perturbation theory using the resolution of the identity approximation (RI-MP2) [29, 30], in larger basis sets. The working expression used for the energy of a given structure was:

$$E(\text{CCSD}(T)/\text{CBS}) = E(\text{HF/cc-pV5Z}) + E^{\text{corr}}(\text{RI-MP2/CBS}_{3,4,5}) + E^{\text{corr}}(\text{CCSD}(T)/\text{cc-pVTZ}) - E^{\text{corr}}(\text{RI-MP2/cc-pVTZ}) + ZPE(\omega\text{B97X-V/cc-pVTZ}).$$

 $E^{\rm corr}({\rm RI-MP2/CBS}_{3,4,5})$  is the extrapolated RI-MP2 correlation energy using the cc-pVTZ, cc-pVQZ, and cc-pV5Z basis sets to fit

$$E^{\text{corr}}(\text{RI-MP2/cc-pVNZ}) = E^{\text{corr}}(\text{RI-MP2/CBS}_{3,4,5}) + AN^{-3},$$

where N denotes the cardinal number for the cc-pVNZ basis sets [123]. These CCSD(T)/CBS energies are typically accurate to about 3-4 kJ mol<sup>-1</sup>. The energy of the silicon atom in the <sup>1</sup>D state was calculated by computing the energy of the atom on its triplet ground state and corrected with an experimentally derived triplet-singlet gap (75.4 kJ mol<sup>-1</sup>) taken from the National Institute of Standards and Technology Atomic Spectra Database. Time-dependent density functional theory[124] calculations were also carried out on select SiC<sub>3</sub>H<sub>2</sub> species to search for electronic excited states corresponding to the absorption of a Lyman- $\alpha$  photon. All calculations were performed using the Q-Chem suite of electronic structure programs [125]. For further analysis, the RRKM [126] rate constants for select dehydrogenation pathways of SiC<sub>3</sub>H<sub>4</sub> isomers were approximated using a Beyer-Swinehart direct state counting algorithm [127] modified to include vibrational tunneling as suggested by Miller [128].

## 2.3 Results and discussion

## Experimental Results: Summary of Crossed Molecular Beam Studies

#### Laboratory Frame

The reactions of excited-state silicon atoms (Si(<sup>1</sup>D)) with methylacetylene (CH<sub>3</sub>CCH) and allene (H<sub>2</sub>CCCH<sub>2</sub>) were explored via molecular crossed beam experiments under single collision conditions at a collision energy of  $30 \pm 2$  kJ mol<sup>-1</sup>. The silicon source was generated via photolysis of 0.5% disilane seeded in helium. In addition to excited state silicon, this primary molecular beam also included ground state silicon (Si(<sup>3</sup>P)) and the silylidyne radical (X<sup>2</sup>II). Ground state silicon atoms do not react with the hydrocarbons explored here, while the reactions of the silylidyne radical with allene and methylacetylene have been studied previously[129, 130]. These prior studies allow us to distinguish between signal arising from reactions of the silylidyne radical and the excited-state silicon atom with methylacetylene and allene in the present study. Products of these reactions were ionized via electron impact (80 eV) and detected by a rotatable quadrupole mass-spectrometer, with time of flight (TOF) spectra gathered from analysis of the mass-to-charge ratios (m/z) and velocities of the ionized fragments[131].

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For both the methylacetylene and allene systems reactive signals were observed at m/zof 68, 67, and 66, corresponding to  $SiC_3H_4^+$ ,  $SiC_3H_3^+$ , and  $SiC_3H_2^+$ , respectively. Previous studies have shown that the signal at m/z 68 originated from the reaction of methylacetylene and allene with the silvlidyne radical  $(SiH(X^2\Pi))$  to form 2-methyl-1-silacycloprop-2envlidence (c-SiC<sub>3</sub> $H_4$ ) via an indirect scattering mechanism and atomic hydrogen ejection [129, 130. For both reactions, scaling of the signal at m/z 67 reveals it to be indistinguishable from the signal at m/z 68, demonstrating that this signal arises from dissociative ionization of the SiC<sub>3</sub>H<sub>4</sub> fragments. The reactive signal at m/z 66, however, is distinct from the signals at m/z 68 and 67, suggesting that this fragment does not arise solely from the dissociative ionization of  $SiC_3H_4$  but also comes from the product(s) of indirect scattering reactions between excited state silicon atom and methylacetylene and allene, yielding  $SiC_{3}H_{2}$  isomer(s) via molecular hydrogen elimination (reaction I). Fig. 2.4, taken from Ref. 93, presents the laboratory angle distribution and select TOF spectra for m/z 66 corresponding to SiC<sub>3</sub>H<sub>2</sub><sup>+</sup> for both the allene and methylacetylene reactions. The wide distributions and symmetries about the center-of-mass angles in both the methylacetylene and allene cases suggests indirect mechanisms involving SiC<sub>3</sub>H<sub>4</sub> intermediates. The TOF spectra at m/z 66 can be fit to two channels in each case: a major channel comprising approximately 90% of the scattering signal accounting for the reactive scattering of  $Si(^{1}D)$  and  $C_{3}H_{4}$  and a minor channel corresponding to the dissociative ionization of the  $SiC_3H_4$  products of the  $SiH-C_3H_4$  scattering dynamics explored previously [129, 130]. The laboratory angular distributions of the TOF spectra presented in Fig. 2.1 show a spread over  $45^{\circ}$  and a center-of-mass (CM) angle of 33° for both reactions. The near forward-backward symmetry of these distributions suggest these reactions occur via indirect scattering mechanisms involving SiC<sub>3</sub>H<sub>4</sub> intermediates.

In order to further investigate the mechanism for molecular hydrogen loss in the Si(<sup>1</sup>D)-C<sub>3</sub>H<sub>4</sub> systems (reaction I), molecular cross beam experiments were carried out using partially deuterated D3-methylacetylene [CD<sub>3</sub>CCH,(X<sup>1</sup>A<sub>1</sub>)]. These reactions can proceed either via HD loss to form SiC<sub>3</sub>D<sub>2</sub> (reaction III) or via D<sub>2</sub> loss to form SiC<sub>3</sub>HD (reaction IV), leading to TOF signals at m/z 68 and 67, respectively. With the exception of other isotopically substituted products (5.1% <sup>29</sup>Si, 3.3% <sup>13</sup>C), these signals correspond exclusively to SiC<sub>3</sub>D<sub>2</sub><sup>+</sup> and SiC<sub>3</sub>DH<sup>+</sup>, respectively. Fig. 2.2 presents the TOF spectra for both m/z 68 (blue) and 67 (red) ionized products of the reaction of Si(<sup>1</sup>D) with D3-methylacetylene collected at the CM angle. These measurements predict a branching ratio of 60 ± 15% HD loss products and 40 ± 15% D<sub>2</sub> loss products. Therefore the reaction of Si(<sup>1</sup>D) with D3-methylacetylene is seen to exhibit both HD and D<sub>2</sub> loss channels leading to distinct isomers. Similar analysis could not be performed with allene due to the chemical equivalency of all the hydrogen atoms in allene.

$$\operatorname{Si}(^{1}\mathrm{D}) + \operatorname{CD}_{3}\mathrm{CCH} \longrightarrow \operatorname{SiC}_{3}\mathrm{D}_{2} + \mathrm{HD}$$
(III)  
$$\operatorname{Si}(^{1}\mathrm{D}) + \operatorname{CD}_{3}\mathrm{CCH} \longrightarrow \operatorname{SiC}_{3}\mathrm{D}_{2} + \mathrm{HD}$$
(III)

$$Si(^{1}D) + CD_{3}CCH \longrightarrow SiC_{3}HD + D_{2}$$
 (IV)



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Time-of-flight (us) Figure 2.1: A and B present laboratory angular distributions for the reactions of  $Si(^{1}D)$ (blue) and SiH( $X^2\Pi$ ) (green) with allene and methylacetly energy respectively. The red curves represent the overall fit, solid circles represent experimental data points with one  $\sigma$  error bars, and CM denotes the center of mass angle. C and D present TOF spectra at select angles for the allene and methylacetylene systems, respectively. Here the blue curves represent reactions with  $Si(^{1}D)$ , the green curves reactions with  $SiH(X^{2}\Pi)$ , the red curves the overall

fit, and the open circles experimental data points. Atom colors: silicon (green), carbon

500 1000

(black), and hydrogen (light grey). Taken with permission from Ref. 93.

Ó

500

1000

0

500 1000

0

500 1000

0

500

1000 1500



Figure 2.2: TOF spectra at m/z 68 (HD loss,  $SiC_3D_2^+$ , blue) and 67 (D<sub>2</sub> loss,  $SiC_3HD$ , red) taken at the respective CM angles for the reaction of  $Si(^{1}D)$  with D3-methylacetylene. Taken with permission from Ref. 93.

#### Center-of-Mass Frame

Interpretation of the raw TOF data suggests that the reactive scattering of excited state silicon atoms with allene and methylacetylene yields SiC<sub>3</sub>H<sub>2</sub> molecules and molecular hydrogen. Deuterium-labeled experiments further show that in the case of methylacetylene there are multiple channels, including channels where the ejected hydrogen molecule is composed of hydrogen atoms from both the acetylenic and methyl groups. In order to probe the identity of these possible  $SiC_3H_2$  products, the reaction dynamics of these systems was converted to the CM reference frame. CM translational energy distributions  $P(E_{\rm T})$  and CM angular distributions  $T(\theta)$  for the SiC<sub>3</sub>H<sub>2</sub> plus molecular hydrogen exit channels are presented in Fig. 2.3. The maximum product translational energies for the Si(<sup>1</sup>D)-allene and Si(<sup>1</sup>D)methylacetylene systems were found to be  $218 \pm 28$  and  $216 \pm 26$  kJ mol<sup>-1</sup>, respectively. These maximum product translation energies are a sum of the initial collision energies and the reaction energies for products in their ground ro-vibrational states. This suggests then that the products of these Si(<sup>1</sup>D)-C<sub>3</sub>H<sub>4</sub> should be excergic by 187  $\pm$  31 kJ mol<sup>-1</sup>. The broad peaks in the CM translational energy distributions centered at 50-80 kJ mol<sup>-1</sup> suggest molecular hydrogen loss from the  $SiC_3H_4$  complexes occurs via a tight exit transition states and involves significant electron density rearrangement [132]. The wide distributions in the CM angular distributions suggests intermediate scattering dynamics where  $SiC_3H_4$ intermediates persist with lifetimes comparable to or greater than the rotational period of the complex [133]. Both of these angular distributions are centered around 90°, suggesting a "sideways scattering" mechanism where the exit vector of the molecular hydrogen is nearly perpendicular to the rotation plane of the  $SiC_3H_4$  transition state structures [133].

#### **Electronic Structure Calculations for the Bimolecular Reactions**

#### Molecular hydrogen loss channels

Having experimentally identified  $\text{SiC}_3\text{H}_2$  isomers as the products of reactive scattering of excited state silicon with allene and methylacetylene, a computational approach was used to further elucidate the underlying reaction mechanisms and chemical dynamics. Reactants, products, intermediates, and transition state structures for species relevant to the reactions of electronically excited silicon atoms with allene, methylacetylene, and D3-methylacetylene are characterized in Tab. A.1-A.3 and summarized in Fig. 2.4, which presents zero-point corrected energies for species accessible within the experimental collision energy of  $30 \pm 2$  kJ mol<sup>-1</sup> (a full version is presented in Fig. A.1). For both allene and methylacetylene, products **p1-p4** fall within the experimentally determined reaction energy window of  $187 \pm 31$  kJ mol<sup>-1</sup>.

The calculations predict the barrierless addition of Si(<sup>1</sup>D) into the  $\pi$  electron density of the hydrocarbon. In the case of methylacetylene, this addition occurs via addition to the triple bond to yield [i1], the most stable of the SiC<sub>3</sub>H<sub>4</sub> intermediate species. From [i1], the system can isomerize to [i2], [i4], [i5], and [i6] with barriers of 305, 188, 239, and 256 kJ mol<sup>-1</sup>, respectively. Despite the lower barrier to [i4], rapid unimolecular dissociation to **p2** 

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Figure 2.3: CM translational energy distributions (A and B) and CM angular distributions (C and D) for the reactions of  $Si(^{1}D)$  with allene (A and C) and methylacetylene (B and **D**) to yield  $SiC_{3}H_{2}$  species via molecular hydrogen elimination. The shaded areas denote acceptable fits for one  $\sigma$  errors in the laboratory data. Taken with permission from Ref. 93.

via  $[i4] \rightarrow [i13] \rightarrow p2$  is stifled by a barrier of 260 kJ mol<sup>-1</sup> for the two-step process. Isomer [i5] can readily isomerize to [i2] with a barrier of 19 kJ mol<sup>-1</sup> or to [i9] with a barrier of 99 kJ mol<sup>-1</sup>. From [i6], there is fascile isomerization to [i9] and [i12] with barriers of 42 and 60 kJ mol<sup>-1</sup>, respectively, and access to [i10] through a barrier of 143 kJ mol<sup>-1</sup>. Most of the intermediates are thereby capable of interconverting, provided products are not too rapidly formed.

Qualitatively, this is confirmed by examining the exit barriers for molecular hydrogen loss to form  $SiC_3H_2$ . The most favorable exit transition states are the  $[i6] \rightarrow p4$  transition state at -114 kJ mol<sup>-1</sup>, i.e., a barrier of 188 kJ mol<sup>-1</sup>, the  $[i9] \rightarrow p1$  transition state at -89 kJ  $mol^{-1}$  with barrier of 190 kJ  $mol^{-1}$ , and two  $[i12] \rightarrow p3$  transition states at -71 and -76 kJ  $mol^{-1}$  with barriers of 177 and 172 kJ  $mol^{-1}$ , but from a higher energy intermediate. Exit transition states to  $\mathbf{p2}$  are much higher in energy, with energies of 7,-46, and -46 kJ mol<sup>-1</sup> for the  $[i1] \rightarrow p2$ ,  $[i6] \rightarrow p2$ , and  $[i13] \rightarrow p2$  transition states, respectively. Detailed prediction of branching ratios requires solving the coupled kinetic equations for all of these pathways, after applying Rice-Ramsperger-Kassel–Marcus (RRKM) theory to obtain the individual rate constants (Tab. A.9). The outcome for this network is a predicted product distribution of 38.2% **p1**,0.6% **p2**, 11.1% **p3**, and 45.8% **p4**. The experimental results from the silicon–D3methylacetylene system, which reveal a branching ratio of  $60\pm15\%$  HD loss and  $40\pm15\%$  D<sub>2</sub> loss, allows for validation of the computed reaction network. Comparing the moderate exit barriers to the comparatively much smaller barriers to isomerization, significant isomerization is expected. Including all possible constitutional isomers in the reaction network, our RRKM



Figure 2.4: Simplified potential energy surface for the reactions of excited state silicon with methylacetylene and allene. Energies relative to separated reactants are given in units of kJ mol<sup>-1</sup>. Pathways with barriers above the 30 kJ mol<sup>-1</sup> collision energy are excluded. Atom colors: silicon (green), carbon (black), and hydrogen (light grey). Taken from Ref. 93.

calculations predict a product distribution of 56.8% HD loss and 43.2%  $D_2$  loss. This result is in good agreement with experimental results and suggests that the energy landscape shown in Fig. 2.4 is faithfully capturing the critical aspects of the experimental dynamics observed for Si<sup>(1</sup>D) with methylacetylene.

Turning next to the computational results for the reaction of Si(<sup>1</sup>D) with allene, it is evident from Fig. 2.4 that the silicon atom can add barrierlessly to a C=C bond, yielding a 3-membered ring structure, [i7], or attack the central carbon atom to form the ring-opened isomer [i10]. There is virtually no barrier to ring-opening [i7] to yield [i10]. Intermediate [i10] provides access to the main intermediates already seen in silicon–methylacetylene system. For instance, the sequences  $[i10]\rightarrow[i2]\rightarrow[i1]$  and  $[i10]\rightarrow[i6]\rightarrow[i1]$  both access intermediate [i1], the initial encounter product for Si(<sup>1</sup>D)-methylacetylene. Therefore, the initial collision complexes of the silicon–methylacetylene ([i1]) and silicon–allene ([i7]/[i10]) surfaces are effectively coupled via intermediates [i2] and [i6] with reaction pathways eventually leading to  $\mathbf{p1}-\mathbf{p4}$  via H<sub>2</sub> loss involving tight exit transition states located 45-167 kJ mol<sup>-1</sup> above the energy of the separated products. The predicted product distribution for silicon–allene system: 38.7% **p1**, 0.5% **p2**, 11.3% **p3**, and 46.5% **p4**.

#### Atomic hydrogen and methyl loss channels

Having established that the molecular hydrogen loss channel leads predominantly to the formation of the 1-sila-1,2,3-butatrienylidene molecule( $H_2C=C=C=Si; X^1A_1; p1$ ) and the less stable 4-membered ring structure (c-SiCH=C=CH;X^1A\_1;p4), the possibility of atomic hydrogen (H) and methyl (CH<sub>3</sub>)loss channels were also investigated, though these channels

were not observed experimentally. For the atomic hydrogen loss channel, we located 22  $\operatorname{SiC}_{3}H_{3}$  isomers (Tab. A.4). Six isomers–one acyclic (**n1**) and five cyclic molecules (**n2-n6**)–are energetically accessible under our experimental conditions at collision energies of  $30\pm 2$  kJ mol<sup>-1</sup> holding overall reaction energies from -10 to 24 kJ mol<sup>-1</sup>. For the methyl loss channels, four SiC<sub>2</sub>H isomers could be identified (Tab. A.5), two of which were permitted under experimental conditions with overall reaction energies of -30 and -5 kJ mol<sup>-1</sup> for **m1** and **m2**, respectively. Statistical (RRKM) calculations were conducted to examine the branching ratios for the energetically accessible hydrogen atom and methyl loss channels (Tab. A.9) versus molecular hydrogen loss. These radical product channels were found to contribute only about 4% in total to the overall product stream for the silicon–methylacetylene system, and only 3% of the product stream for the silicon–allene system. These results bolster our experimental findings that the molecular hydrogen loss is the primary dissociation channel and also support the experimental non-detection of atomic hydrogen and methyl group loss channels.

# Electronic Structure Calculations: Photodissociation of $SiC_3H_2$ to $SiC_3$

We have thus far concluded that the reactions of excited state silicon atom with allene and methylacetylene lead primarily to the formation of **p1** and **p4** under the experimental conditions. While these experiments are designed to mimic the conditions in the inner envelope of carbon stars, we also consider computationally the ejection of these  $\text{SiC}_3\text{H}_2$ product species to the outer envelope, where subsequent photodissociation may occur to form silicon tricarbide  $\text{SiC}_3$  isomers including cyclic silicon tricarbide (c-SiC<sub>3</sub>), which has been astronomically observed. Reactants, products, and transition state structures for the isomerization and dehydrogenation reactions of  $\text{SiC}_3\text{H}_2$  are characterized in Tab. A.6 and summarized in Fig. 2.5, where the zero-point corrected energies are presented relative to the most stable  $\text{SiC}_3\text{H}_2$  isomer **p1**.

Under experimental conditions the total system energy is 220 kJ mol<sup>-1</sup> above **p1**, allowing for isomerization from **p1** to **p2** via ring closure, **p2** to **p5** via concerted ring opening and ring closure, **p4** to **p3** via ring opening, and **p3** to **p6** via two concerted ring closures, while leaving other the other isomerization channels and all of the dehydrogenation channels closed. Further activation of these species is required for dehydrogenation channels to be open, such as the absorption of a Lyman- $\alpha$  photon (10.2 eV or 984 kJ mol<sup>-1</sup> for hydrogen) from the interstellar UV field. As shown in Tab. A.7, **p1-p6** all have excited states with nonzero oscillator strength in the region of 10.2 eV, the species exhibiting the strongest absorption being **p2** (10.2 eV; f = 0.20) and **p4** (10.2 eV; f = 0.25). Following absorption and non-radiative relaxation to a vibrationally hot electronic ground state, these SiC<sub>3</sub>H<sub>2</sub> species possess enough internal energy to surmount the barriers to isomerization and dehydrogenation via molecular hydrogen loss to SiC<sub>3</sub> species [**1**], [**2**], and [**3**].

In agreement with our results, previous computational studies predict that the bicyclic

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Figure 2.5: Potential energy surface for the isomerization of low-energy  $SiC_3H_2$  isomers and dehydrogenation to yield  $SiC_3$ . Energies relative to the lowest energy  $SiC_3H_2$  isomer **p1** are given in units of kJ mol<sup>-1</sup>. The total internal energy on this scale for products of reaction **I** is 220 kJ mol<sup>-1</sup>. Atom colors: silicon (green), carbon (black), and hydrogen (light grey). Taken from Ref. 93.

SiC<sub>3</sub> isomer [1] represents the thermodynamically most stable isomer on the singlet electronic surface, with the other bicyclic isomer [2] being 8-52 kJ mol<sup>-1</sup> higher in energy[134–143]. This lowest energy isomer [1] has been observed in the circumstellar envelope of IRC+10216. Previous studies of the linear isomers of SiC<sub>3</sub> have also predicted the carbon-terminal CSiCC to be approximately 318 kJ mol<sup>-1</sup> higher in energy than the silicon-terminal SiCCC on the triplet electronic surface  $(X^3\Sigma^-)[144]$ .

We have also computationally investigated atomic hydrogen loss pathways from **p1**, **p2**, and **p5** as alternative photodissociation channels (Tab. A.8). These channels are endoergic by 387-410 kJ mol<sup>-1</sup> relative to **p1**. While these reaction energies are comparable to the molecular hydrogen loss barriers heights, the molecular hydrogen loss products are thermodynamically favored and the channels leading to them are enhanced by tunneling. We therefore suggest that molecular hydrogen loss to yield c-SiC<sub>3</sub> is the main dissociation channel.

#### Summary of Astrochemical Modeling

Following the combined computational and experimental conclusions that  $SiC_3H_2$  species including **p1** and **p4** are formed via bimolecular reactions of excited state silicon atoms with allene and methylacetylene and further computational findings that these  $SiC_3H_2$  species can form cyclic silicon tricarbide (c-SiC<sub>3</sub>) via molecular hydrogen-elimination photodissociation, Chapter 2. Modeling of Gas Phase Formation of  $c-SiC_3$  Molecules in the Circumstellar Medium of Carbon Stars 40

our co-authors used astrochemical modeling to relate these results to the circumstellar envelope of IRC+10216. This modeling seeks to account for the presence of other reactants and the overall reaction network not considered in the crossed beam experiments and electronic structure calculations. The simulations employed the chemical kinetic data from the RATE12 release of UMIST Database for Astrochemistry[110] with the addition of SiC<sub>3</sub>H<sub>2</sub> reactivity proposed in this work and with work physical parameters taken from Li et. al[145]. Fractional abundances (measured relative to molecular hydrogen) of species are also taken from refs. 110 and 145.

Firstly, these modeling studies approximated the abundance of excited state silicon atoms in the inner envelope. Sources of Si<sup>(1</sup>D) include photodissociation of silane (SiH<sub>4</sub>)[146], silicon monoxide (SO), and silicon monosulfide (SiS)[114, 147], which have detected fractional abundances of  $2.2 \times 10^{-7}$ ,  $1.8 \times 10^{-7}$ , and  $1.3 \times 10^{-6}$ , respectively, in the inner envelope of IRC+10218[110, 145]. For SiH<sub>4</sub>, quantitative studies by Suto and Lee revealed that excited state silicon atom contributes 1.2% of the photodissociation yield for 106-160 nm photons[146]. These photodissociation channels could lead to up to 1% of the total budget of silicon, which peaks at a fractional abundance of  $1.0 \times 10^{-6}$ , to be excited state silicon atoms. These finding confirm that excited state silicon atoms required for our mechanisms are present in the circumstellar envelope of IRC+10216.

Results of the modeling study including the channels presented in this work are summarized in Fig. 2.6. The present-day fractional abundances of allene, methylacetylene, silicon tricarbide, and silicon tetracarbide all peak at a radius of  $4 \times 10^{16}$  cm with fractional abundances of  $8 \times 10^{-9}$ ,  $8 \times 10^{-9}$ ,  $7 \times 10^{-9}$ , and  $2 \times 10^{-10}$ , respectively[148]. From these fractional abundances it is clear that dissociation of silicon tetracarbide cannot be the dominant precursor of silicon tricarbide. The inclusion of the two-step mechanism proposed in this work has increased the modeled column density of silicon tricarbide by 80% versus previous modeling efforts. Since only Lyman- $\alpha$  photons were considered to photodissociate SiC<sub>3</sub>H<sub>2</sub>, this modeled fractional abundance of silicon tricarbide should be seen as a lower bound.

Finally, the modeling study predicts a fractional abundance of  $SiC_3H_2$  to peak at a fractional abundance  $10^{-10}$  relative to molecular hydrogen. At this abundance signal from these species should be observable in future searches to further confirm the proposed mechanisms.

# 2.4 Conclusions

Our combined computational and experimental approach for the bimolecular reactions of excitated state silicon atom with allene and methylacetlyene has revealed the indirect scattering mechanisms to form  $\operatorname{SiC}_3H_2$  products, including 1-sila-1,2,3-butatrienylidene ( $H_2C=C=C=Si$ ;  $X^1A_1$ )(**p1**) among other isomers. These reactions begin with the barrierless addition of silicon into the  $\pi$  electron density of the carbon-carbon double bond(s) for allene and the carbon-carbon triple bond for methylacetlyene, yielding distinct  $\operatorname{SiC}_3H_4$  species. The internal energy available to these intermediates allows for interconversion between several  $\operatorname{SiC}_3H_4$ 

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Figure 2.6: Modeled present-day fractional abundances relative to molecular hydrogen are plotted versus the radius of the circumstellar envelope of the carbon star IRC+10216. Species color: silicon tricarbide  $(SiC_3)(blue)$ , silicon tetracarbide  $(SiC_4)(red)$ , allene  $(H_2CCCH_2)(magenta), methylacetylene (CH_3CCH)(cyan), SiC_3H_2$  (green). Taken from Ref. 93.

intermediate species before unimolecular decomposition via molecular hydrogen elimination. The computed reaction network was used to predict the product isomer distributions and provided excellent agreement with the the experimentally observed HD/D2 loss ratios observed in the experimental study of the silicon-D3-methylacetlyene system. The experimental collision energy of  $30 \pm 2$  kJ mol<sup>-1</sup> reflects the conditions in the circumstellar envelopes of carbon-rich stars undergoing shocks and pulsations of the stellar ejecta [114]. These  $SiC_3H_2$ products of reaction 1 can then be transported from the inner envelope to the outer envelope via solar winds. We computationally studied the photolysis of  $SiC_{3}H_{2}$  by the interstellar UV field to form cyclic silicon tricarbide  $(c-SiC_3)$  via molecular hydrogen elimination (reaction II). These computational mechanisms were then integrated into an astrochemical kinetic model to draw conclusions about the observation and abundance of  $c-SiC_3$  in the circumstellar envelope of IRC+10216.

These systems demonstrate excited-state neutral-neutral reactions in the inner circumstellar envelope to form SiC<sub>3</sub>H<sub>2</sub> species which are then dehydrogenated by the interstellar UV field in the outer envelope to form cyclic silicon tricarbide  $(c-SiC_3)$ . While previous models have suggested photodissociation of silicon tetracarbide  $(SiC_4)$  as the source of circumstellar silicon tricarbide[99, 108, 149], the modeled fractional abundance of silicon tricarbide in the circumstellar envelope of IRC+10216 exceeds that of silicon tetracarbide by two orders of magnitude [110, 145]. The computations and astrochemical modeling in this work propose  $SiC_3H_2$  species as important precursors of c-SiC<sub>3</sub>. While the computational photodissociation study only considered Lyman- $\alpha$  photons as opposed to the wider interstellar UV field. this study serves as a proof-of-concept for the two-step mechanisms taking place in the inner and outer circumstellar envelopes. These mechanism for silicon carbide formation via neutral-neutral bimolecular reactions and subsequent dehydrogenation provide an alternative to the conventional belief that silicon carbides could only be formed via high-temperature ion-molecule reactions or via photodissociation of larger silicon carbide molecules.

With the commission of the Atacama Large Millimeter/Submillimeter Array, the detection of novel organosilcon species will undoubtedly grow (including the 1-sila-1,2,3-butatrienylidene, which has not yet been searched for in the envelope of IRC+10216). An understanding of the chemistry of these species will rely on combined computational, experimental, and astrochemical modeling efforts, helping to close the long-standing gap between observational and laboratory data.

# Chapter 3

# Benchmarking the Performance of the ReaxFF Reactive Force Field on Hydrogen Combustion Systems

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# 3.1 Introduction

Hydrogen combustion is a promising energy source for both space exploration[150] and cleaner terrestrial transportation[151–154]. Hydrogen is an attractive fuel in part because of its high energy density and lack of carbon-containing combustion products[155–157]. Such applications require understanding of how the chemistry of hydrogen and hydrogen combustion plays a role in storage, efficiency, and safety[150, 153, 158–162]. Furthermore, many of the elementary reactions relevant to hydrogen combustion are present in the combustion of other hydrogen-containing fuels[163].

Extensive effort has been made to model the reaction kinetics of hydrogen combustion under a variety of conditions with ample feedback from experimental studies. Mueller et al.[164] developed a model based on their experimental findings to accurately describe the kinetics below the second explosion limit. Li et al.[165], beginning from the mechanisms proposed by Mueller et al., updated the model by modifying reaction rates to better predict high-pressure flame speed experiments. O'Connaire et. al[166] also developed a new mechanism validated on flames speeds, ignition time delays, and species composition, also modifying reaction rates to better account for higher pressure experiments. Konnov[167], in updating the  $H_2/O_2$  combustion mechanism, provided better analysis of the uncertainties in reaction rates and suggested the importance of  $HO_2 + H$  reactions in accurately predicting ignition delays. Hong et. al.[168] proposed updates to reaction rates based on their UV and

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IR measurements on high-temperature mixtures in shock tubes. Burke et al. [169] developed a model based on the Li mechanism that can push to more intermediate pressures towards the second extended explosion limit. To better treat high pressures, Shimuzi et al.[170], building of the work of Kitano et al.[171], further investigated the importance of  $HO_2 + H$ reactions and called for more accurate rate constats to be determined for several reactions. The development of the ELTE mechanism by Varga et. al[172] involved a global optimization of parameters and considered all available direct and indirect measurements for evaluation. The authors also benchmarked the performance of a variety of previous models on their measurement set. An experimental study of laminar flames across a wide parameter range was complemented by quantum-chemical calculations of rates constants and shock tube data to develop the HP-Mech model and extract pressure-dependent reaction orders [173] Burke and Klippenstein [174], investigating the elementary reactions of hydrogen combustion via a ab *initio*-master equation, proposed inclusion of new three-body radical propogation and termination reactions. Recent work by Konnov[175], however, has suggested that new transport parameters suggested by Jasper et al. [176, 177] provided a better match to flame speeds than the inculsion of Burke and Klippenstein's new reactions. Experimental studies on the roles of impurities in shock tube data have also been recently performed [178]. Sabia et al. [179], in an experimental work to better quantify the role that different bath gases, has found existing models lacking for bath gases other than  $N_2$ . In this work, the authors also highlight several existing challenges in the modeling of hydrogen combustion, including uncertainties in rate constants, unconsidered or excluded reaction mechanisms, and uncertainties relating to three-body effects. A more complete review of the history, current standings, and challenges of modeling  $H_2$  combustion is given in the discussion of their work.[179]

Three of us[180–185] have recently developed and applied a method of reactive symbol sequences to the study of hydrogen combustion as an alternative route to analyze chemical kinetics. This framework avoids the ideas of elementary reactions and rate coefficients, both of which can be strongly dependent on modeling assumptions [186, 187]. Instead, reactive symbol sequences can be seen as a way to probe the emergent chemical processes present in combustion chemistry. Applied to reactive molecular dynamics (MD), this approach naturally includes important non-equilibrium fluctuations. Beginning with an all-atom potential function that depends only on the atomic positions, in principle also allows for a method to be predictive from the outset instead of relying on parameterization which is only relevant near the particular conditions on which it is trained.[188]

To generate the trajectories for this approach, one could, in principle, consider using *ab initio* MD simulations where energies and forces are determined using electronic structure theory methods. Examples of this include the nanoreactor studies in which pressure and temperature fluctuations are used to induce reactive events [189, 190]. Despite their accuracy for barrier heights and reaction energies, the high computational cost of these methods limits both the size and timescale of simulations.

Alternatively, empirical methods including reactive force fields [188, 191–196] and tight binding [197–199] offer a computationally more feasible approach to larger and longer simulations, though potentially with a loss of accuracy. The former allow for the nanosecond

scale simulations on thousands of atoms. Within the ReaxFF reactive force field[188], energies are determined via a combination of connection dependent terms based on bond orders and a polarizable charge description. These potentials depend on system-specific parameters which are trained against QM data, including reaction energies and barrier heights. ReaxFF potentials have been developed and applied to combustion reactions[200–216], including hydrocarbons, hydrogen[203, 212], and syngas[216]. The ReaxFF method has also been extensively applied to material science[217–232], catalysis[233–237], and other chemical systems[188, 238–246].

The reactive symbol sequence studies for hydrogen combustion cited above have utilized ReaxFF potentials in order to generate the trajectories required for their analysis. While ReaxFF potential are derived from fits to QM data, there are no reports available on assessment of ReaxFF parameterizations against independent test sets of QM data at combustion relevant atomic configurations. In this work, we benchmark the performance of several ReaxFF parameter sets on a set of trajectories and reactions relevant to the combustion of hydrogen. The predictive power of the ReaxFF method hinges on its fidelity to QM data, especially in conditions where experimental validation is difficult or impossible.

# **3.2** Computational Methods

#### **ReaxFF** Reactive Force Field Method

ReaxFF is an atomistic potential capable of describing reactive events. Bond breaking and formation are captured in ReaxFF via computed bond orders between pairs of atoms. Intramolecular bonded energy contributions are captured using potential functions that depend on these bond orders, which are determined by the interatomic distances. Nonbonded intermolecular interactions including Coulomb, dispersion, and polarization are computed between all pairs of atoms using geometry-dependent charge equilibration schemes and are damped at short range[247, 248]. Analytic forms for these potential functions can be found in the the supporting information of Chenoweth et al.[200]. The total potential is:

$$E_{\text{system}} = E_{\text{bond}} + E_{\text{lp}} + E_{\text{over}} + E_{\text{under}} + E_{\text{val}} + E_{\text{tor}} + E_{\text{coa}} + E_{\text{conj}} + E_{\text{H-bond}} + E_{\text{vdW}} + E_{\text{coul}} + E_{\text{pol}}$$
(3.1)

where  $E_{\text{bond}}$ ,  $E_{\text{lp}}$ ,  $E_{\text{over}}$ ,  $E_{\text{under}}$ ,  $E_{\text{val}}$ ,  $E_{\text{tor}}$ ,  $E_{\text{conj}}$ ,  $E_{\text{H-bond}}$ ,  $E_{\text{vdW}}$ ,  $E_{\text{coul}}$ , and  $E_{\text{pol}}$ , are the bond, lone pair, overcoordination, undercoordination, valence angle, torsional angle, threebody conjugation, four-body conjugation, hydrogen bond, van der Waals, Coulomb, and polarization energy contributions, respectively.

For a given class of chemical systems, parameters for these energy contributions were trained on a data set comprised of electronic structure theory calculations and experimental data. For the purposes of this work, we have selected four combustion-specific parameter sets of the ReaxFF force field to benchmark against high-level QM calculations.

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- 1. The CHO2008 parameter set was developed in 2008 by Chenoweth, van Duin, and Goddard was a first parameterization of a ReaxFF potential for hydrocarbon oxidation[200].
- 2. The HO2011 parameter set was developed in 2011 by Agrawalla and van Duin specifically to simulate hydrogen combustion at high temperature and pressure[203].
- 3. The HO2014 parameter set was developed in 2014 by Cheng et al. as a reoptimization of the CHO2008 set for use in accelerated reactive molecular dynamics simulations[212].
- 4. The CHO2016 parameter set was developed in 2016 by Ashraf and van Duin as a systematic improvement of the CHO2008 to better capture small hydrocarbon oxidation and combustion initiation[216].

Our single-point calculations were performed using the reax/c package[249] as implemented in LAMMPS (http://lammps.sandia.gov)[250].

#### Molecular Dynamics Simulations

Molecular dynamics trajectories were calculated using the HO2011 parameter set[203] with the ReaxFF potential[200]. All trajectories simulations were carried out with constant particle number, volume, and energy (NVE) and utilized a time step of 0.1 fs and periodic boundary conditions.

We first considered a high-pressure simulations initialized with 66 H<sub>2</sub> molecules, 33 O<sub>2</sub> molecules, and an OH radical initiator at an initial pressure of 950 bar and an initial temperature of 1000 K (initial density  $\rho_0 = 250$  kg m<sup>-3</sup>, stoicheometric ratio  $\phi = 1$ ). These trajectories were allowed to evolve for 3 ns, at which point more than 80% of the reactants were converted to water[181]. This set of trajectories was calculated using the PuReMD-GPU simulation package[251]. This set of trajectories has already been reported and analyzed using reactive symbol sequences[181].

A larger, lower pressure simulation was also performed to better isolate bimolecular reactive events. This simulation was initialized with 320 H<sub>2</sub> molecules, 160 O<sub>2</sub> molecules, and 16 OH radical initiators at an initial pressure of 50 bar and an initial temperature of 1250 K ( $\rho_0 = 5.9$  kg m<sup>-3</sup>,  $\phi = 1$ ). This trajectory was allowed to evolve for approximately 70 ns, at which point the temperature had equilibrated and the reactants were nearly completely converted to water. The OH radicals were added to initiate the ignition of the reaction mixture. Figure 3.1 presents the kinetic temperature of this reacting mixture as a function of simulation time. This trajectory was calculated using the reax/c package[249] in LAMMPS[250].

From the simulated MD trajectories under these high and low pressure conditions, "reactive pathways" were extracted during the ignition events that characterize the combustion reaction at fixed energy [181]. These paths correspond to isolated elementary bi- and termolecular reaction events. Because they occur naturally during the course of ignition, the



Figure 3.1: Time series of the kinetic temperature for a hydrogen-oxygen mixture that is initially stoichiometric at a temperature of 1250 K and pressure of 50 bar.

paths of these reactions through the transition state regions include oscillations about the spectator modes. These reactive paths are a point of comparison to minimum energy paths found via quantum-mechanical calculations.

### Quantum-Mechanical Calculations

In order to benchmark the performance of the ReaxFF potential, QM calculations were performed using coupled cluster theory with single, double, and perturbative triple excitations [CCSD(T)][252] in the spin-unrestricted wave function formalism utilizing the frozen-core approximation and the cc-pVTZ basis set[119] and extrapolated to the complete basis set limit. The correlation energy was extrapolated with second-order Møller-Plesset perturbation theory using the resolution of the identity approximation (RI-MP2)[29, 30]. The expression for this overall extrapolated energy is

$$E(\text{CCSD}(T)/\text{CBS}) = E(\text{HF/cc-pV5Z}) + E^{\text{corr}}(\text{RI-MP2/CBS}_{3,4,5}) + E^{\text{corr}}(\text{CCSD}(T)/\text{cc-pVTZ}) - E^{\text{corr}}(\text{RI-MP2/cc-pVTZ})$$
(3.2)

where  $E^{\text{corr}}(\text{RI-MP2/CBS}_{3,4,5})$  is the extrapolated RI-MP2 correlation energy using the ccpVTZ, cc-pVQZ, and cc-pV5Z basis sets[119] to fit

$$E^{\text{corr}}(\text{RI-MP2/cc-pVNZ}) = E^{\text{corr}}(\text{RI-MP2/CBS}_{3,4,5}) + AN^{-3}$$
(3.3)

and N is the cardinality of the cc-pVNZ basis set [123]. From comparisons between extrapolated CCSD(T) and composite Weizmann methods as well as evaluation of CCSD(T) for bond stretching, we predict errors in our benchmark calculations on the order of 1–4 kcal  $mol^{-1}[253-255]$ .

As another point of comparison, QM calculations were also performed using the  $\omega$ B97M-V[256] density functional in the spin-unrestricted wave function formalism with the cc-pVTZ basis set[119]. This functional has been ranked as one of the most accurate in a recent assessment of 200 density functionals on systems including thermochemistry and barrier heights.[85] All DFT calculations were carried out using an ultra-fine integration grid using 99 radial points and 590 angular points.

In addition to the reactive pathways harvested from MD simulations, intrinsic reaction coordinate scans for a subset of bimolecular reactions relevant to hydrogen combustion from Li et al.[165] were used for benchmarking systems. Structures for the products and reactants were computed via geometry optimization at the  $\omega$ B97M-V/cc-pVTZ level of theory. Transition structures were computed using the freezing string method[121] to generate an approximate structure and Hessian, then refined by the partitioned-rational function optimization eigenvector following method (P-RFO)[122], also at the  $\omega$ B97M-V/cc-pVTZ level of theory. Vibrational analysis was used to confirm that the minima have no imaginary frequencies and that the transition state structures have only one imaginary frequency. The reaction path on either side of the transition state structure was computed using the intrinsic reaction coordinate method[257–259]. All QM calculations were performed using the Q-Chem[125] package of electronic structure programs.

# 3.3 Results and discussion

#### **Reactive Pathways**

A collection of seven reactive pathways were harvested from larger MD simulations and are presented in order to benchmark the performance of ReaxFF against electronic structure theory calculations. By this, we mean that the coordinates of all spectator species are removed, leaving a relatively small subset of atoms undergoing chemical changes, whose relative energies can be tractably benchmarked. As the overall MD simulations were performed with the HO2011 parameter set, energies of the configurations along these pathways will be relative to energies of the isolated reactants at geometries optimized with ReaxFF using the HO2011 parameter set.

Two termolecular reactions were extracted from the high-pressure simulation ( $N_{\rm H_2,0} = 66$ ,  $N_{\rm O_2,0} = 33$ ,  $N_{\rm OH,0} = 1$ ,  $P_0 = 950$  bar,  $T_0 = 1000$  K,  $\rho_0 = 250$  kg m<sup>-3</sup>,  $\phi = 1$ ). The first of these reactions is the near simultaneous cleaving of a hydrogen molecule by an oxygen atom and an oxygen molecule to yield a hydroxyl radical and a hydroperoxyl radical (O + H<sub>2</sub> + O<sub>2</sub>  $\longrightarrow$  OH + HO<sub>2</sub>). For the *ab initio* calculations the ground state is a singlet. The second of these termolecular reactions is the cleaving of a hydrogen molecule by an

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oxygen molecule and a hydroperoxyl radical to form a hydroperoxyl radical and hydrogen peroxide  $(O_2 + H_2 + HO_2 \longrightarrow HO_2 + H_2O_2)$ . The ground state is a doublet for this reaction. Plots of the potential energy along the pathways for the methods surveyed are presented in Figures B.1–B.2.

In addition to these termolecular reactive pathways, five bimolecular reactive pathways were extracted from the 50 bar simulation ( $N_{H_2,0} = 320$ ,  $N_{O_2,0} = 160$ ,  $N_{OH,0} = 16$ ,  $P_0 = 50$ bar,  $T_0 = 1250$  K,  $\rho_0 = 5.9$  kg m<sup>-3</sup>,  $\phi = 1$ ). The first of these is the initiation reaction between a hydrogen molecule and an oxygen molecule forming a hydroperoxyl radical and a hydrogen atom ( $H_2+O_2 \longrightarrow HO_2+H$ ), which has a triplet ground state. The second is the reaction of a hydrogen atom and an oxygen molecule to form a hydroperoxyl radical ( $H+O_2 \longrightarrow HO_2$ ), which has a doublet ground state. Third is the the reaction of a hydroperoxyl radical and a hydrogen atom to form two hydroxyl radicals ( $HO_2+H \longrightarrow 2OH$ ), occurring on the singlet surface. Fourth and fifth are both reactions of hydrogen molecules and hydroxyl radicals to form water and a hydrogen atom ( $H_2 + OH \longrightarrow H_2O + H$ ), which have doublet ground states. Plots of the potential energy along the trajectories for the methods surveyed are presented in Figures 3.2, B.3–B.7.

0					
Reaction	Method	RMSD	MSD	MAX	NPE
$O + H_2 + O_2 \longrightarrow OH + HO_2$	$\omega B97M-V$	3.00	-2.51	-4.81	5.28
	CHO2008	59.54	-57.06	-77.33	53.70
	HO2011	41.55	-38.49	-58.41	57.93
	HO2014	28.18	-26.18	-45.20	32.14
	CHO2016	18.93	-10.78	-41.07	55.93
$O_2 + H_2 + HO_2 \longrightarrow HO_2 + H_2O_2$	$\omega B97M-V$	2.02	0.74	4.38	7.33
	CHO2008	35.49	-33.70	-61.14	43.02
	HO2011	28.98	-24.46	-55.46	62.44
	HO2014	32.58	-29.67	-55.45	60.57
	CHO2016	16.86	-6.94	-37.15	58.38
$H_2 + O_2 \longrightarrow HO_2 + H$	$\omega B97M-V$	1.21	-0.43	-2.77	3.91
	CHO2008	22.34	-19.36	-37.54	32.87
	HO2011	16.35	-13.51	-32.23	35.36
	HO2014	11.31	3.30	17.02	30.45
	CHO2016	8.40	-0.94	-18.31	34.08
$H + O_2 \longrightarrow HO_2$	$\omega B97M-V$	2.17	1.57	4.79	5.10
	CHO2008	11.37	-10.63	-21.88	17.73
	HO2011	10.03	-1.21	15.43	30.85
	HO2014	12.60	-8.27	17.90	34.42
	CHO2016	8.63	-3.92	-14.17	26.69
$HO_2 + H \longrightarrow 2 OH$	$\omega B97M-V$	4.11	3.68	6.00	6.43

Table 3.1: Root mean square deviation, mean signed deviation, maximum absolute deviation, and non-parallelity error, in kcal  $mol^{-1}$ , are presented for a selection of reactive pathways extracted from larger MD simulations. CCSD(T)/CBS values were used as reference.

Table 5.1: (continued)						
Reaction	Method	RMSD	MSD	MAX	NPE	
	CHO2008	18.06	5.64	26.88	53.42	
	HO2011	8.55	2.43	-17.35	33.95	
	HO2014	21.08	-4.21	-59.19	90.63	
	CHO2016	14.80	-1.80	-38.68	58.01	
$\rm H_2 + OH \longrightarrow H_2O + H$	$\omega B97M-V$	1.69	-0.50	-3.42	5.79	
	CHO2008	14.66	-10.58	-28.76	33.98	
	HO2011	8.29	-3.46	-20.16	32.29	
	HO2014	8.27	-4.02	-15.89	24.25	
	CHO2016	10.48	-4.62	-19.28	32.09	
$H_2 + OH \longrightarrow H_2O + H$	$\omega B97M-V$	1.94	0.03	3.11	5.47	
	CHO2008	9.64	-4.75	-20.49	28.76	
	HO2011	8.05	-2.24	-16.22	26.00	
	HO2014	5.99	-1.60	-11.26	22.42	
	CHO2016	8.72	7.61	16.81	28.94	
All Pathways	Method	RMSD	MSD	MAX	NPE	
	$\omega B97M-V$	2.62	1.07	6.00	7.33	
	CHO2008	18.32	-8.30	-77.33	53.70	
	HO2011	12.46	-3.56	-58.41	62.44	
	HO2014	15.10	-5.25	-59.19	90.63	
	CHO2016	11.44	-2.59	-41.07	58.38	

Table 3.1: (continued)

Table 3.1 presents the root mean square deviations (RMSD), mean signed deviation (MSD), maximum absolute deviation (MAX), and non-parallelity error (NPE), all in kcal mol<sup>-1</sup>, for pathways extracted from MD simulations. Here, the non-parallelity error is defined as the difference between the maximum and minimum error along a reaction pathway. For the overall metrics, the modern density functional,  $\omega$ B97M-V, unsurprisingly outperforms all of the surveyed ReaxFF methods by a factor of 4 for RMSDs, 2 for MSDs, 7 for MAXs, and 7 for NPEs. Of the ReaxFF methods, the CHO2016 parameter set outperforms the other parameter sets in terms of RMSDs, MSDs, and MAXs with values of 11.44 kcal mol<sup>-1</sup>, -2.59 kcal mol<sup>-1</sup>, and 41.07 kcal mol<sup>-1</sup>, respectively. The next-best performing parameter set for these properties was HO2011, followed by HO2014 and CHO2008, respectively. The CHO2008 parameter set performs the best for NPEs with a value of 53.70 kcal mol<sup>-1</sup>. The negative MSDs for all methods demonstrates a systematic overstabilization of configurations relative to isolated reactants. Furthermore, the NPEs in excess of 15 kcal mol<sup>-1</sup> and up to 90 kcal mol<sup>-1</sup> demonstrate that ReaxFF does not capture the shape of the potential energy surfaces along the reactive pathways.

As an example, Figure 3.2 presents the potential energy along the first  $H_2 + OH \longrightarrow H_2O+H$  reactive pathway for the CCSD(T) reference,  $\omega B97M-V$ , and four ReaxFF parameter sets. Like the others, this reactive pathway was sampled from a larger mixture undergoing



Figure 3.2: Potential energy along the first  $H_2 + OH \longrightarrow H_2O + H$  reactive pathway, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to the energies of isolated  $H_2$  and OH computed at the ReaxFF/HO2011 optimized geometries. The total time to complete this reaction path is 30 fs with a time step of 0.1 fs.

ignition. This particular reactive path begins with the rotation of a  $H_2$  molecule relative to an OH radical, followed attack by the OH radical and oscillation of the transferring H atom between the OH and H fragments until a  $H_2O$  molecule and H atom are left. Of the ReaxFF potentials, the HO2014 parameter set yields the best RMSD, MAX, and NPE of 8.27 kcal mol<sup>-1</sup>, -15.89 kcal mol<sup>-1</sup>, and 24.25 kcal mol<sup>-1</sup>, respectively. The HO2011 parameter set yields the best MSD of -3.46 kcal mol<sup>-1</sup>. Qualitatively, the positions of local minima and maxima for the HO2014 and CHO2016 parameter sets agree with the reference calculations, though the actual values at these extrema often differ by over 10 kcal mol<sup>-1</sup>. The same can be said for the CHO2008 and HO2011 parameter sets in the first 100 time steps as well, however the CHO2008 parameter set significantly underestimates the barrier for hydrogen transfer. The HO2011 parameter set gives the most jarring performance qualitatively, with sharp features arising in the region of oscillation of H between OH and H.

In order to probe the origin of the sharp features in the HO2011 potential energy surface, we investigated the component-wise contributions to the potential energy. Figure 3.3 shows the overall potential energy, atom energy, and the difference, all in kcal mol<sup>-1</sup>. Here the atom energy is defined as the sum of the overcoordination energy  $E_{over}$  and undercoordination energy  $E_{under}$ . For the HO2011 parameter set, the sharp features in the overall potential energy can be traced to the atom energy, whereas for all other parameter sets this contribution changes much more smoothly. In this term, positive contributions are seen to arise mainly from overcoordination of the transferring hydrogen atom whereas negative



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Figure 3.3: Potential energy, atom energy, and the sum of all other energy contributions, in kcal mol<sup>-1</sup>, for the first  $H_2 + OH \longrightarrow H_2O + H$  reactive pathway, extracted from a larger MD simulation, are presented for the four ReaxFF parameter sets studied. The atom energy is the sum of the overcoordination and undercoordination energy contributions. The total time to complete this reaction path is 30 fs with a time step of 0.1 fs.

contributions arise from undercoordination of the OH radical in the entrance channel and H in the exit channel.

#### Intrinsic Reaction Coordinate Scans

We have calculated intrinsic reaction coordinate (IRC) scans for a subset of bimolecular reactions taken from Li et al.[165] as an additional set of systems on which to evaluate the performance of ReaxFF. These IRC scans were calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. Energies of isolated reactant and product species were also calculated at the  $\omega$ B97M-V/cc-pVTZ optimized geometries to present relative energies and compute reaction energies and barrier heights. Eleven scans are included (see Table 3.2 for the full list):

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the radical attack of a hydrogen atom and an oxygen molecule to form a hydroxyl radical and an oxygen atom  $(H + O_2 \longrightarrow OH + O)$  on the quartet surface, the abstraction of hydrogen from a hydrogen molecule by a hydroxyl radical to form a water molecule and a hydrogen atom  $(H_2 + OH \longrightarrow H_2O + H)$  on the doublet surface, the abstraction of hydrogen from a water molecule by an oxygen atom to form two hydroxyl radicals ( $H_2O + O \longrightarrow$ 2 OH) on the triplet surface, the formation of a hydroperoxyl radical from a hydrogen atom and an oxygen molecule  $(H + O_2 \longrightarrow HO_2)$  on the doublet surface, the abstraction of hydrogen from a hydroperoxyl radical by a hydrogen atom to form a hydrogen molecule and an oxygen molecule  $(HO_2 + H \longrightarrow H_2 + O_2)$  on the triplet surface, the radical attack of a hydrogen atom on a hydroperoxyl radical to form two hydroxyl radicals (HO<sub>2</sub> + H  $\longrightarrow$ 2 OH) on the triplet surface, the abstraction of hydrogen from a hydroperoxyl radical by a second hydroperoxyl radical to form a hydrogen peroxide molecule and an oxygen molecule  $(2 \text{HO}_2 \longrightarrow \text{H}_2\text{O}_2 + \text{O}_2)$  on the triplet surface, the radical attack of a hydrogen atom on a hydrogen peroxide molecule to form a water molecule and a hydroxyl radical  $(H_2O_2 + H \longrightarrow$  $H_2O + OH$ ) on the doublet surface, the abstraction of hydrogen from a hydrogen peroxide molecule by a hydrogen atom to form a hydrogen molecule and a hydroperoxyl radical  $(H_2O_2 + H \longrightarrow H_2 + HO_2)$  on the doublet surface, the abstraction of hydrogen from a hydrogen peroxide molecule by an oxygen atom to form a hydroperoxyl radical and a hydroxyl radical  $(H_2O_2 + O \longrightarrow HO_2 + OH)$  on the triplet surface, and the abstraction of hydrogen from a hydrogen peroxide molecule by a hydroxyl radical to form a water molecule and a hydroperoxyl radical  $(H_2O_2 + OH \longrightarrow H_2O + HO_2)$  on the doublet surface. Plots of the potential energy along the IRCs for the methods surveyed are presented in Figures 3.4,3.5,B.8-B.18.

Table 3.2: Root mean square deviation, mean signed deviation, maximum absolute deviation, and non-parallelity error, in kcal/mol, are presented for a selection of intrinsic reaction coordinate scans of elementary reactions relevant to hydrogen combustion. Geometries along the reaction coordinate were calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. CCSD(T)/CBS values were used as reference.

Reaction	Method	RMSD	MSD	MAX	NPE
$\overline{\mathrm{H} + \mathrm{O}_2 \longrightarrow \mathrm{OH} + \mathrm{O}}$	$\omega B97M-V$	1.52	1.36	2.85	2.33
	CHO2008	12.40	-9.37	-37.57	34.64
	HO2011	16.62	-12.14	-42.34	39.11
	HO2014	14.38	-0.96	33.47	45.50
	CHO2016	17.33	-12.77	-48.72	43.71
$\rm H_2 + OH \longrightarrow H_2O + H$	$\omega B97M-V$	2.20	0.95	2.91	5.22
	CHO2008	6.40	2.22	-12.38	21.04
	HO2011	15.50	9.56	44.54	48.60
	HO2014	9.35	3.56	23.35	30.88
	CHO2016	9.24	5.54	16.31	23.16
$H_2O + O \longrightarrow 2 OH$	$\omega B97M-V$	4.03	-3.51	-6.72	5.73
	CHO2008	28.88	-25.50	-53.15	42.86
Reaction	Method	RMSD	MSD	MAX	NPE
---	-----------------	-------	--------	--------	-------
	HO2011	11.65	-4.05	23.24	39.46
	HO2014	17.63	-13.75	-31.44	36.48
	CHO2016	12.91	-10.61	-24.57	23.96
$H + O_2 \longrightarrow HO_2$	$\omega B97M-V$	1.15	0.97	1.92	1.80
	CHO2008	10.25	-10.12	-14.06	6.76
	HO2011	9.40	-5.59	-14.17	21.20
	HO2014	14.92	-7.46	-21.49	36.32
	CHO2016	7.69	-2.68	-10.47	19.55
$HO_2 + H \longrightarrow H_2 + O_2$	$\omega B97M-V$	1.08	-0.40	-3.23	3.98
	CHO2008	10.03	4.98	19.05	31.95
	HO2011	5.18	-2.59	-11.57	15.35
	HO2014	21.36	-17.28	-32.82	33.26
	CHO2016	6.18	-1.29	12.06	19.38
$\rm HO_2 + H \longrightarrow 2  OH$	$\omega B97M-V$	2.38	1.87	4.04	3.68
	CHO2008	11.84	-9.93	-28.12	26.36
	HO2011	16.17	-6.87	-36.94	47.49
	HO2014	25.91	-21.95	-38.11	34.65
	CHO2016	17.39	-11.83	-49.89	47.07
$2 \operatorname{HO}_2 \longrightarrow \operatorname{H}_2\operatorname{O}_2 + \operatorname{O}_2$	$\omega B97M-V$	1.97	-0.69	-4.27	5.93
	CHO2008	5.31	-2.69	-15.33	22.31
	HO2011	12.00	10.01	24.65	23.63
	HO2014	25.32	-9.90	-41.68	59.91
	CHO2016	11.75	4.01	20.44	31.80
$H_2O_2 + H \longrightarrow H_2O + OH$	$\omega B97M-V$	2.12	1.81	3.12	3.42
	CHO2008	3.81	1.83	11.49	12.53
	HO2011	12.96	9.26	19.18	20.57
	HO2014	16.97	5.62	41.82	50.24
	CHO2016	11.84	7.00	17.84	23.45
$\mathrm{H_2O_2} + \mathrm{H} \longrightarrow \mathrm{H_2} + \mathrm{HO_2}$	$\omega B97M-V$	2.65	-1.59	-4.87	5.84
	CHO2008	6.49	-1.64	-17.13	22.76
	HO2011	9.05	5.78	30.34	37.39
	HO2014	9.61	5.88	16.51	22.95
	CHO2016	4.90	1.15	-11.22	17.91
$H_2O_2 + O \longrightarrow HO_2 + OH$	$\omega B97M-V$	3.86	-3.37	-7.71	6.77
	CHO2008	19.71	-17.65	-38.11	30.37
	HO2011	9.66	2.29	26.85	32.55
	HO2014	7.82	2.69	20.47	25.41
	CHO2016	4.50	1.09	7.78	13.85
$H_2O_2 + OH \longrightarrow H_2O + HO_2$	$\omega B97M-V$	2.68	-2.49	-6.79	4.88

Table 3.2: (continued)

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Reaction	Method	RMSD	MSD	MAX	NPE
	CHO2008	8.92	-3.27	-28.16	34.93
	HO2011	10.77	7.74	32.74	28.52
	HO2014	21.47	15.92	44.76	43.01
	CHO2016	13.27	8.13	22.78	25.31
All IRCs	Method	RMSD	MSD	MAX	NPE
	$\omega B97M-V$	2.49	-0.42	-7.71	6.77
	CHO2008	12.47	-5.54	-53.15	42.86
	HO2011	12.36	1.78	44.54	48.60
	HO2014	18.30	-2.46	44.76	59.91
	CHO2016	11.79	-0.40	-49.89	47.07

Table 3.2: (continued)

Table 3.2 presents the RMSD, MSD, MAX, and NPE, all in kcal mol<sup>-1</sup>, for the IRC scans. For the overall metrics,  $\omega$ B97M-V again outperforms all of the surveyed ReaxFF methods by at least a factor of 4 for RMSDs, 5 for MAXs, and 6 for NPEs. Of the ReaxFF methods, the CHO2016 parameter set is shown to outperform all other parameter sets in terms of RMSDs with a value of 11.79 kcal mol<sup>-1</sup>, followed closely by the HO2011 set the CHO2008 set, and finally the HO2014 set. For NPEs, the worst performances are 42.86 kcal mol<sup>-1</sup>, 47.07 kcal mol<sup>-1</sup>, 48.60 kcal mol<sup>-1</sup>, and 59.91 kcal mol<sup>-1</sup> for the CHO2008 set, the CHO2011 set, and the HO2014 set, respectively. The lowest NPE for any of the ReaxFF parameter sets is 12.53 kcal mol<sup>-1</sup>. This demonstrates the failure of ReaxFF to quantitatively replicate the shape of the CCSD(T) potential energy surface.

Table 3.3 presents the reaction energies, forward reaction barrier heights, and reverse reaction barrier heights for the IRC scans at the CCSD(T)/CBS level of theory and errors in these quantities for all other methods considered. For these calculations, forward and reverse barrier heights are measured as the difference between the maximum along the IRC and the isolated reactants or products at the  $\omega B97M-V/cc-pVTZ$  optimized geometries, respectively. Reaction energies are computed as the difference between these isolated products and reactants. For both reaction energies and barrier heights  $\omega B97M$ -V outperforms the best ReaxFF methods by a factor of 2 for reaction energy errors and a factor of 3 for barrier height errors. Of the ReaxFF parameter sets, the HO2011 set shows the best performance for reaction energies with an RMSD of 7.15 kcal  $mol^{-1}$ , followed by the CHO2016 set, the CHO2008 set, and the HO2014 set with RMSDs of 12.40 kcal  $mol^{-1}$ , 14.67 kcal  $mol^{-1}$ , and 25.09 kcal  $mol^{-1}$ , respectively. Turning to barrier heights, the performance of the ReaxFF methods is split, with the CHO2008 and CHO2016 parameter sets outperforming the HO2014 and HO2011 parameter sets by more than 8 kcal mol<sup>-1</sup>. The CHO2008 parameter set is generally seen to underestimate the forward and reverse barriers and in many cases stabilizes configurations in the region of the  $\omega B97M$ -V transition state, leading to a double-barrier potential energy surface. The CHO2016 parameter set shows a similar performance, both underestimating forward and reverse barriers and exhibiting multiple barriers on the potential energy surface.

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For the HO2011 parameter set, forward and reverse barriers are seen to be overestimated on average. In particular, reaction barriers are strongly overestimated for hydrogen transfer reactions, with the exception of the  $HO_2 + H \longrightarrow H_2 + O_2$  IRC scan. The HO2014 parameter set also overestimates both forward and reverse reaction barriers on average, with the performance being exacerbated by the poor performance on reaction energies.

Table 3.3: Reaction energy errors, forward reaction barrier height errors, and reverse reaction barrier height errors, in kcal mol<sup>-1</sup>, are presented for a selection of intrinsic reaction coordinate scans of elementary reactions relevant to hydrogen combustion. Geometries along the reaction coordinate were calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. CCSD(T)/CBS values were used as a reference and the reaction energies, forward reaction barriers, and reverse reaction barriers are reported, in kcal mol<sup>-1</sup>. Infinitely separated reactants and products optimized at the  $\omega$ B97M-V/cc-pVTZ level of theory are taken as the endpoints of the reaction paths.

Reaction	Method	$\Delta E$	$\Delta E_f^{\ddagger}$	$\Delta E_b^{\ddagger}$
$H + O_2 \longrightarrow OH + O$	CCSD(T)/CBS	11.73	41.05	29.32
	$\omega B97 M-V$	5.51	2.69	-2.82
	CHO2008	-17.22	-20.15	-2.92
	HO2011	6.43	-35.80	-42.23
	HO2014	-13.00	13.59	26.58
	CHO2016	-5.99	-33.38	-27.39
$\rm H_2 + OH \longrightarrow H_2O + H$	CCSD(T)/CBS	-16.63	5.61	22.24
	$\omega B97 M-V$	2.05	-2.31	-4.36
	CHO2008	8.55	3.27	-5.28
	HO2011	8.64	34.06	25.42
	HO2014	9.15	10.74	1.58
	CHO2016	13.62	2.72	-10.90
$H_2O + O \longrightarrow 2 OH$	CCSD(T)/CBS	19.00	22.05	3.06
	$\omega B97M-V$	-1.92	-6.65	-4.73
	CHO2008	-25.83	-32.91	-7.08
	HO2011	-13.82	21.44	35.26
	HO2014	-28.01	1.69	29.70
	CHO2016	-16.88	-13.76	3.12
$H + O_2 \longrightarrow HO_2$	CCSD(T)/CBS	-55.81	1.56	57.37
	$\omega B97M-V$	1.48	0.23	-1.25
	CHO2008	-14.25	-7.54	6.70
	HO2011	7.00	-7.98	-14.98
	HO2014	15.24	-16.32	-31.56
	CHO2016	5.61	-7.30	-12.91
$\rm HO_2 + H \longrightarrow H_2 + O_2$	CCSD(T)/CBS	-54.12	2.62	56.74
	$\omega B97M-V$	-0.17	-0.95	-0.78
	CHO2008	18.37	-2.85	-21.21

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Table 5.5. (commund)					
Reaction	Method	$\Delta E$	$\Delta E_f^{\ddagger}$	$\Delta E_b^{\ddagger}$	
	HO2011	0.02	1.15	1.13	
	HO2014	-13.89	-0.94	12.95	
	CHO2016	-0.09	7.84	7.94	
$\mathrm{HO}_2 + \mathrm{H} \longrightarrow 2 \mathrm{OH}$	CCSD(T)/CBS	-40.02	14.42	54.43	
	$\omega B97 M-V$	5.46	0.95	-4.52	
	CHO2008	-16.15	-15.79	0.35	
	HO2011	1.26	-15.54	-16.79	
	HO2014	-45.75	-8.68	37.07	
	CHO2016	-9.34	-16.70	-7.36	
$2 \operatorname{HO}_2 \longrightarrow \operatorname{H}_2\operatorname{O}_2 + \operatorname{O}_2$	CCSD(T)/CBS	-39.20	0.29	39.49	
	$\omega B97 M-V$	2.75	-4.00	-6.75	
	CHO2008	10.76	-11.09	-21.84	
	HO2011	-0.19	20.58	20.77	
	HO2014	-38.29	17.53	55.81	
	CHO2016	-10.53	17.04	27.57	
$\mathrm{H_2O_2} + \mathrm{H} \longrightarrow \mathrm{H_2O} + \mathrm{OH}$	CCSD(T)/CBS	-71.56	7.19	78.76	
	$\omega B97 M-V$	4.59	0.08	-4.51	
	CHO2008	0.01	10.06	10.05	
	HO2011	10.10	2.00	-8.10	
	HO2014	-12.20	29.67	41.87	
	CHO2016	14.72	-5.41	-20.14	
$\mathrm{H_2O_2} + \mathrm{H} \longrightarrow \mathrm{H_2} + \mathrm{HO_2}$	CCSD(T)/CBS	-14.92	10.24	25.16	
	$\omega B97 M-V$	-2.92	-3.23	-0.31	
	CHO2008	7.61	-9.55	-17.16	
	HO2011	0.20	28.29	28.09	
	HO2014	24.40	3.89	-20.51	
	CHO2016	10.44	-2.17	-12.62	
$\mathrm{H_2O_2} + \mathrm{O} \longrightarrow \mathrm{HO_2} + \mathrm{OH}$	CCSD(T)/CBS	-12.55	11.86	24.42	
	$\omega B97 M-V$	-2.79	-7.29	-4.50	
	CHO2008	-9.67	-21.31	-11.64	
	HO2011	-4.98	24.35	29.33	
	HO2014	5.54	10.19	4.65	
	CHO2016	7.18	-2.10	-9.29	
$H_2O_2 + OH \longrightarrow H_2O + HO_2$	CCSD(T)/CBS	-31.55	3.33	34.88	
	$\omega B97 M-V$	-0.87	-5.69	-4.81	
	CHO2008	16.16	-10.66	-26.82	
	HO2011	8.84	26.29	17.46	
	HO2014	33.55	25.20	-8.35	
	CHO2016	24.06	1.77	-22.29	

Table 3.3: (continued)

	X	/	
Reaction	Method	$\Delta E$	$\Delta E_f^{\ddagger}  \Delta E_b^{\ddagger}$
RMSD	Method	$\Delta E$	$\Delta E^{\ddagger}$
	$\omega B97M-V$	3.25	4.00
	CHO2008	14.67	15.08
	HO2011	7.15	23.69
	HO2014	25.09	23.45
	CHO2016	12.40	15.24

Table 3.3: (continued)

Figure 3.4 presents the potential energy along the  $H_2 + OH \longrightarrow H_2O + H$  IRC for the CCSD(T) reference,  $\omega$ B97M-V, and the four ReaxFF parameter sets. For the CHO2008 parameter set, we see a slightly overestimated barrier followed by a overstabilization of the configurations in the center of the scan. The HO2011 parameter set is seen to overestimate the forward and reverse barriers by 34 kcal mol<sup>-1</sup> and 25 kcal mol<sup>-1</sup>, respectively. The HO2014 parameter set exhibits an exaggerated forward barrier and underestimates the stabilization in the product channel. The CHO2016 parameter set yields a comparatively good forward barrier but clearly overstabilizes configurations near the transition state leading to a second barrier.

Figure 3.5 presents the potential energy along the  $H_2O_2 + OH \longrightarrow H_2O + HO_2$  IRC for the CCSD(T) reference,  $\omega$ B97M-V, and the four ReaxFF parameter sets. Here the CHO2008 parameter set fails to predict a forward barrier to reaction and significantly overstabilizes configuration in the transition state regime. The HO2011 parameter set yields an exaggerated barrier that overestimates the forward barrier height by 26 kcal mol<sup>-1</sup> and the reverse barrier height by 17 kcal mol<sup>-1</sup>. For the HO2014 parameter set the forward barrier is exaggerated nearly to the extent of the HO2011 parameter set and severely understabilizes the product channel, leading to a reaction energy error of 34 kcal mol<sup>-1</sup>. The CHO2016 parameter set again gives a second barrier for this reaction and poorly reproduces the reaction barrier.

For both of the reactions presented above, Agrawalla and van Duin[203] show that the forward reactions should have no forward barrier with the HO2011 parameter set. In order to better elucidate the nature of the potential energy surfaces, Figures 3.6 and 3.7 present unrelaxed two-dimensional potential energy surfaces for the reactions  $H_2 + OH \longrightarrow H_2O + H$  and  $H_2O_2 + OH \longrightarrow H_2O + HO_2$ , respectively. For these surfaces, the bonds which are breaking and forming are linearized while all other internal coordinates are held fixed in the configuration of the transition state calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. The interatomic distance along the IRC for the bonds breaking and forming are superimposed in each plot to approximate the one-dimensional slices given in Figures 3.4 and 3.5.

Beginning with  $H_2 + OH \longrightarrow H_2O + H$ , the transition state structure as optimized with  $\omega$ B97M-V has an HO-H bond length of 1.35 Å and a H-H bond length of 0.82 Å. The CHO2008 surface displays an entrance channel with a longer optimal H-H bond length and predicts a transition state structure to occur with a longer HO-H bond length and a longer H-H bond length. This surface also displays a well that extends to much shorter



Figure 3.4: Potential energies along the  $H_2 + OH \longrightarrow H_2O + H$  intrinsic reaction coordinate are presented for four ReaxFF parameter sets,  $\omega B97M$ -V, and CCSD(T). Geometries along the reaction coordinate were calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2$  and OH computed at the  $\omega B97M$ -V/cc-pVTZ optimized geometries.



Figure 3.5: Potential energies along the  $H_2O_2 + OH \longrightarrow H_2O + HO_2$  intrinsic reaction coordinate are presented for four ReaxFF parameter sets,  $\omega B97M$ -V, and CCSD(T). Geometries along the reaction coordinate were calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2O_2$  and OH computed at the  $\omega B97M$ -V/cc-pVTZ optimized geometries.



Figure 3.6: Unrelaxed potential energy surfaces for the  $H_2 + OH \longrightarrow H_2O + H$  reaction, in kcal mol<sup>-1</sup>, are presented for  $\omega B97M$ -V and four ReaxFF parameter sets. Here the HO···H···H angle is linearized while all other geometric parameters are held fixed in the transitions state configuration as optimized at the  $\omega B97M$ -V/cc-pVTZ level of theory. The interatomic distances for points along the intrinsic reaction coordinate calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory are overlaid on each surface. The energies presented are relative to the energies of the isolated atoms.

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H–H bond lengths than the  $\omega$ B97M-V surface. The HO2011 surface predicts a transition state structure with a longer H–H bond length. Following the IRC on the HO2011 surface elucidates the origin of the large errors in the barrier heights for this parameter set. The HO2014 surface predicts a transition state structure with a longer H–H bond length while the IRC skews toward a longer HO–H bond length, leading to an exaggerated barrier. The CHO2016 surface demonstrates the double barrier seen in Figure 3.4, with one transition state structure with a longer HO–H bond length and another with a longer H–H bond length. This surface also exhibits an exit channel with a slightly longer HO–H bond length than the  $\omega$ B97M-V surface.

For  $H_2O_2+OH \longrightarrow H_2O+HO_2$ , the transition state structure as optimized with  $\omega B97M-V$  has an HOO-H bond length of 1.04 Å and a HO-H bond length of 1.37 Å. The CHO2008 surface exhibits a transition state structure with a longer HO-H bond length and a well in the region immediately after the  $\omega B97M-V$  transition state. The HO2011 surface predicts a transition state structure with a longer HOO-H bond length, leading to an exaggeration of the barrier along the  $\omega B97M-V$  IRC. This behavior is also seen in the HO2014 surface, though the HO2014 surface predicts a longer optimum HOO-H bond length in the entrance channel. The CHO2016 surface exhibits very little energy variation across the IRC, which misses a narrow well in the product channel. The IRC is also destabilized by a short HOO-H bond length in the region of the  $\omega B97M-V$  transition state.

These two dimensional potential energy surfaces, coupled with the IRC scans, demonstrate failures of the ReaxFF methods to accurately model the QM methods. The ReaxFF potential show a general trend towards looser transition state structures. These differences can be expected to manifest themselves in MD simulations, where exaggerated barriers will prevent reactivity that would otherwise occur on the QM surfaces and under-estimated barriers will overrepresent other reactions.

#### **3.4** Conclusions

We have benchmarked the performance of several ReaxFF methods against state-of-the-art DFT calculations and CCSD(T). The data sets included both reactive pathways extracted from MD simulations and IRC scans, reaction energies, and barrier heights computed with  $\omega$ B97M-V. For the reactive pathways, the CHO2016 and HO2011 parameter sets had the best RMSDs (11 and 12 kcal mol<sup>-1</sup>, respectively) while the average NPEs were similar across all methods surveyed (38 kcal mol<sup>-1</sup>, 40 kcal mol<sup>-1</sup>, 42 kcal mol<sup>1</sup>, and 42 kcal mol<sup>-1</sup> for CHO2008, HO2011, HO2014, and CHO2016, respectively). Along the reaction coordinate scans, the CHO2016, HO2011, and CHO2008 parameter sets exhibit similar performance on RMSDs (12 kcal mol<sup>-1</sup>), while the average NPEs are favor CHO2008 and CHO2016 (26 kcal mol<sup>-1</sup> for CHO2008 and CHO2016; 32 kcal mol<sup>-1</sup> and 38 kcal mol<sup>-1</sup> for HO2011 and HO2014, respectively). While in principle the RMSDs are subject to systematic error, the NPEs serve as an independent metric.

While these results appear to be a condemnation of ReaxFF methods for these systems,



Figure 3.7: Unrelaxed potential energy surfaces for the  $H_2O_2 + OH \longrightarrow H_2O + HO_2$  reaction, in kcal mol<sup>-1</sup>, are presented for  $\omega B97M$ -V and four ReaxFF parameter sets. Here the HOO···H···OH angle is linearized while all other geometric parameters are held fixed in the transitions state configuration as optimized at the  $\omega B97M$ -V/cc-pVTZ level of theory. The interatomic distances for points along the intrinsic reaction coordinate calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory are overlaid on each surface. The energies presented are relative to the energies of the isolated atoms.

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studies by both the developers [188, 200–204, 206, 207, 210, 212, 216–221, 223, 225–230, 232–246] and users [180–183, 205, 208, 209, 211, 213–215, 222, 224, 231] have demonstrated the utility of ReaxFF as an exploratory tool for atomistic simulation. We also acknowledge the difficulty of parameter training for the ReaxFF system. Our results here point to ways to improve the treatment of hydrogen combustion.

- Amount of training data: Training data for the hydrogen combustion potentials with ReaxFF has largely centered around bond stretches, angle distortions, heats of formation, and reaction barriers. As seen in the IRC scans, the transition state structures computed with DFT are often higher-energy structures on the ReaxFF potential energy surfaces. More geometric properties of transition state structures in the ReaxFF training set should in principle help to relieve this discrepancy. The variable performance of the various ReaxFF potentials on the reactive pathways suggest that inclusion of data points far from the minimum energy pathways would improve the overall fidelity of the method.
- Quality of training data: References 200, 203, and 216 all specify that the QM training data was generated at the B3LYP/6-311G<sup>\*\*</sup> level of theory [56, 61, 260]. A recent review of DFT benchmarked the performance of over 200 functionals across diverse chemical systems, including hydrogen transfer and non-hydrogen transfer barrier heights. For these two test sets, B3LYP gave RMSDs of 4.89 kcal mol<sup>-1</sup> and 5.52 kcal mol<sup>-1</sup>, respectively, while  $\omega$ B97M-V gave RMSDs of 1.72 kcal mol<sup>-1</sup> and 1.98 kcal mol<sup>-1</sup>, respectively.
- Optimization methods for parameter training: The ReaxFF parameter sets employed in this study were generated using the successive one-parameter parabolic extrapolation (SOPPE) approach[261], which is both tedious and subject to converging to local minima. As a more global optimization scheme, several studies beginning with Pahari and Chaturvedi in 2012[262] and Larrson, van Duin, and Hartke[263, 264] in 2013 have used genetic algorithms to optimize parameters for different ReaxFF potentials. More recently, a study of ReaxFF parameter optimization with Monte-Carlo and evolutionary algorithms has shown that care needs to be used when employing such methods and gives recommendations for initial guesses[265]. Machine learning approaches have also been applied recently to a parameter optimization scheme for ReaxFF[266]
- Overall form of the potential: Furman and Wales[267] recently published a letter in which a Taper corrections applied to bond order and bond length terms in ReaxFF are able to greatly reduce energy drifts seen in MD simulations on long timescales. Such corrections could in principle improve the sharp changes in the atom energy term seen in Figure 3.3. Currently work is also being done to replace the charge equilibration scheme in ReaxFF with a more advanced C-GeM model[268].

Overall, the benchmarking of ReaxFF potential for hydrogen combustion systems shows that there is room for improvement in order to more closely match QM potential energies, Chapter 3. Benchmarking the Performance of the ReaxFF Reactive Force Field on Hydrogen Combustion Systems 64

energy differences important for thermochemistry, and energy barriers controlling chemical kinetics.

### Chapter 4

# Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

Reprinted with permission from Bertels, L. W.; Lee, J.; Head-Gordon, M. Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy for Thermochemistry, Kinetics, and Intermolecular Interactions. J. Phys. Chem. Lett. **2019**, 10, 4170-4176.

#### 4.1 Introduction

Single-reference second-order Møller-Plesset perturbation theory (MP2) is among the most popular correlated wavefunction methods in electronic structure theory, in part due to its economical  $\mathcal{O}(N^5)$  scaling, where N is the basis set size.

$$E_{MP2} = -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}}$$

$$\tag{4.1}$$

Equation (4.1) gives the correlation energy for MP2, where  $i, j, \ldots$  represent occupied molecular orbitals,  $a, b, \ldots$  represent virtual molecular orbitals, and  $\Delta_{ij}^{ab} = \varepsilon_a + \varepsilon_b - \varepsilon_i - \varepsilon_j$  is the (non-negative) energy denominator. The resolution-of-identity (RI) technique applied to MP2 has allowed for a much more widespread use due to the reduction of the prefactor in the overall computational cost of the algorithm.[29, 30]

Orbital-optimized MP2 (OOMP2) methods were developed to improve the performance of MP2 for energies and other properties.[35–37]. For systems where the unrestricted Hartree-Fock (UHF) reference exhibits spin-contamination (artificial spin-symmetry breaking), the Chapter 4. Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy 66

use of these reference orbitals can lead to catastrophic performance of MP2.[31–34] OOMP2 can also be thought of as a relatively inexpensive way to approximate Brückner orbitals.[35] Orbital optimization at the MP2 level often reduces the level of spin-contamination and improves energetics.[35, 36, 269]

Despite the benefits of OOMP2 described above, there are several unsatisfying characteristics of the method that limit its applicability. Orbital optimization at the MP2 level can produce divergent energy contributions due to small energy denominators. This is often observed when stretching bonds and leads to significant underestimation of harmonic vibrational frequencies.[40] OOMP2 also often fails to continuously transition from spin-restricted (R) to spin-unrestricted (U) solutions even when the U solution is lower in energy.[270] A continuous R to U transition requires a Coulson-Fischer point where the lowest eigenvalue of the R to U stability Hessian becomes zero.[271] Resolution of this issue is necessary to reach the proper dissociation limit for bond-breaking curves.

Our group has attempted to remedy these issues of OOMP2 through use of regularization to prevent divergence of the energy due to small energy denominators. The first of these approaches was to shift the energy denominator by a constant factor,  $\delta$ , so that  $\Delta_{ij}^{ab} \leftarrow \Delta_{ij}^{ab} + \delta$ .[40] This simple form was able to partially resolve the two issues with OOMP2 described above. The regularization parameter,  $\delta$ , both prevents the energy expression from diverging and damps the correlation energy contribution from MP2, leading the method to more closely resemble the continuous R to U transition seen in the HF reference. Unfortunately, in the case of scaled opposite spin OOMP2 (SOS-OOMP2), Razban et al.[41] found the values of  $\delta$  that could restore Coulson-Fischer points were very large and consequently led to poor performance on problems that are normally well-treated by MP2.

Recently, two of us[39] developed two new classes of orbital energy dependent regularizers for OOMP2, of which the most promising is denoted as  $\kappa$ -OOMP2. In  $\kappa$ -OOMP2, the matrix elements associated with small denominators are damped such that:

$$E_{MP2}(\kappa) = -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}} \left(1 - e^{-\kappa \left(\Delta_{ij}^{ab}\right)}\right)^2 \tag{4.2}$$

Unlike the case of  $\delta$ -OOMP2, for  $\kappa$ -OOMP2 the unregularized energy expression is recovered for large energy denominators, and in the limit of small energy denominators, the correlation energy contributions are zero. Regularization parameters of  $\kappa \leq 1.5 E_h^{-1}$  were found to restore Coulson-Fischer points for hydrogen, ethane, ethene, and ethyne bond-breaking curves.  $\kappa$ was trained on the W4-11 set to suggest a value for general application.[272] The result,  $\kappa =$  $1.45 E_h^{-1}$ , proved robust to further testing on the RSE43[273] and TA13[274] sets, and defines  $\kappa$ -OOMP2 as a replacement for OOMP2. Complex restricted (cR) and complex general (cG) orbital extensions of  $\kappa$ -OOMP2 have also been used to interrogate singlet biradicaloids[275] and the nature of symmetry breaking in fullerenes[276], respectively.

The success and ubiquity of MP2 and OOMP2 have led several research groups to develop modified second-order methods aimed at improving energetics. Notable examples are spin-component-scaled MP2 (SCS-MP2)[277–282] and orbital optimized SCS-MP2 (SCS-OOMP2)[35, 36] methods, which weight correlation contributions coming from same-spin Chapter 4. Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy 67

and opposite-spin pairs of electrons differently. These techniques have also been applied to the second-order correlation contribution in several double-hybrid density functionals.[86, 88, 283–285] A subset of these methods, scaled-opposite-spin MP2 (SOS-MP2)[278, 280, 282] and SOS-OOMP2[35], are notable in that they can be implemented via an overall  $\mathcal{O}(N^4)$ computational cost. Another example are the attenuated MP2 methods that partially cancel basis set superposition errors with errors in MP2 itself to yield improved intermolecular interaction energies in finite basis sets.[286–289] Density-fitting and Cholesky-decomposed variants of OOMP2 have also been developed to improve the computational efficiency of the method.[290]

Inclusion of higher-order terms in the perturbative expansion provides another approach to improve energetics from MP2.

$$E_{MP3} = \frac{1}{8} \sum_{ijabcd} (t_{ij}^{ab})^* \langle ab || cd \rangle t_{ij}^{cd}$$
  
+  $\frac{1}{8} \sum_{ijklab} (t_{ij}^{ab})^* \langle kl || ij \rangle t_{kl}^{ab}$   
-  $\sum_{ijkabc} (t_{ij}^{ab})^* \langle kb || ic \rangle t_{kj}^{ac}$  (4.3)

Equation (4.3) gives the third-order Møller-Plesset (MP3) contribution to the correlation energy in the spin-orbital basis. MP3 formally scales as  $\mathcal{O}(N^6)$  with basis set size, and describes the leading interaction of first-order pair-correlations,  $t_{ij}^{ab}$ , with each other. However, despite the higher compute cost, MP3 only modestly improves MP2 results (e.g. see data presented later). In passing we note that it is possible to utilize separable density fitting techniques such as tensor hypercontraction to achieve quartic scaling ( $\mathcal{O}(N^4)$ ) MP3 (and also MP2) energy evaluation.[291]

Grimme<sup>[277]</sup> developed a spin-component scaled MP3 (SCS-MP3) method that improved ground state energies over SCS-MP2 for reaction energies, atomization energies, ionization energies, and stretched geometries. This method applied an overall third-order correlation energy scaling factor of 0.25 in addition to the scaling factors for same-spin and opposite-spin components. For weak noncovalent interactions, application of MP3 has failed to substantially improve binding energies as compared to MP2.[277, 281, 292–296] Hobza and coworkers [294–296] proposed scaling the third-order correlation energy to interpolate between the MP2 and MP3, leading to the development of MP2.5 and MP2.X, in order to improve binding energies for noncovalent interactions. Following these successes, Bozkaya and coworkers [269, 290, 297–299] developed OOMP3 and OOMP2.5 and evaluated the performance of these methods on thermochemistry, kinetics, and noncovalent interactions. OOMP2.5 was shown to outperform coupled cluster theory with single and double excitations (CCSD)[22, 23] on reaction energies and barrier heights [269] and perform comparably to coupled cluster with single, double, and perturbative triple excitations [CCSD(T)][25] for noncovalent interactions [290]. These are very promising results. Analytic gradients for OOMP3, OOMP2.5, and their density-fitting variants have also been introduced. [300, 301]

Following the recent success of regularized OOMP2 in treating inherent problems in OOMP2, we decided to explore the use of  $\kappa$ -OOMP2 orbitals at the MP3 level. At the same time, we wanted to see if  $\kappa$  regularization in MP3 could improve the overall energetics. Beginning from  $\kappa$ -OOMP2 would allow this method to avoid energy divergences caused by small energy denominators.[39]

### 4.2 Theory

In  $\kappa$ -OOMP2, damping of the two-electron integrals leads to the following expression for the *t*-amplitudes:

$$t_{ij}^{ab}(\kappa) = -\frac{\langle ab||ij\rangle}{\Delta_{ij}^{ab}} \left(1 - e^{-\kappa \Delta_{ij}^{ab}}\right).$$

$$(4.4)$$

Inserting Equation (4.4) into Equation (4.3) we arrive at a regularized third-order correlation energy expression,  $E_{MP3}(\kappa)$ . Our first candidate ansatz involved calculating the  $\kappa$ -OOMP2 energy and applying a scaled single-shot  $E_{MP3}(\kappa)$  correction.

$$E_c(\kappa, \kappa_2, c_3) = E_{\kappa-OOMP2}(\kappa) + c_3 E_{MP3:\kappa-OOMP2}(\kappa_2)$$
(4.5)

As a second, alternative form, we considered using  $\kappa$ -OOMP2 (with  $\kappa = 1.45E_h^{-1}$ ) as a method to generate molecular orbitals for use in correlated calculations containing second and third order energies which could then be independently regularized and/or scaled:

$$E_c(\kappa, \kappa_2, c_2, c_3) = c_2 E_{MP2:\kappa-OOMP2}(\kappa_2) + c_3 E_{MP3:\kappa-OOMP2}(\kappa_2)$$

$$(4.6)$$

In this second case the non-Brillouin singles contribution,  $-\sum_{ia} f_{ia}^2 / \Delta_i^a$ , is included at secondorder as  $\kappa$ -OOMP2 does not obey the Brillouin theorem. For simplicity (and ease of implementation), we do not include a non-Brillouin singles contribution at third-order.

We trained both energy functionals on the non-multireference (non-MR) subset of the W4-11 thermochemistry data set.[272] We excluded the MR points in the set from the training data because the single reference methods we are investigating should not be able to describe MR systems adequately. Both reference and training calculations were performed using the aug-cc-pVTZ (aVTZ) basis set [119, 120, 302] and the corresponding RI basis[303, 304] without the frozen core approximation. Reference data were computed using CCSD(T)[25]. All calculations were performed in a development version of Q-Chem[125].

### 4.3 Results

Figure 4.1 presents the root mean square deviations (RMSD) for scans of the  $\kappa_2$  and  $c_3$  parameters in the first model, as given by Equation 4.5. Overall, we see that stronger regularization at third-order (smaller  $\kappa_2$ ) serves to lower the error on the training set. For  $\kappa_2 = 1.00E_h^{-1}$ , the optimal scaling parameter for the third-order regularized correlation

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Figure 4.1: Scans of the root mean square deviation on the non-MR subset of the W4-11 thermochemistry data set, in kcal mol<sup>-1</sup>, for the scaled, two regularization parameter correlation energy functional given in Equation 4.5. We fix  $\kappa = 1.45E_h^{-1}$ . All calculations use the aVTZ basis; CCSD(T) is used for the reference values.

energy is 0.5 for a RMSD of 4.85 kcal mol<sup>-1</sup>. If instead one applies the same regularization parameter ( $\kappa = 1.45E_h^{-1}$ ) at second- and third-order, we find an optimal scaling parameter for the third order correlation energy of  $c_3 = 0.4$  with a RMSD of 4.91 kcal mol<sup>-1</sup>. We note that  $c_3 = 0.0$  corresponds to  $\kappa$ -OOMP2 with a RMSD of 7.58 kcal mol<sup>-1</sup>. Inclusion of scaled, regularized third-order correlation energy contributions reduces the RMSD of  $\kappa$ -OOMP2 by more than 2.5 kcal mol<sup>-1</sup>, which is useful but not dramatic.

Turning to the second form we considered, Figure 4.2 presents the RMSDs for scans of the  $\kappa_2$  and  $c_3$  parameters in Equation 4.6 with  $c_2 = 1.0$  (the optimal value for all  $\kappa_2$  values plotted). For this ansatz, we see that the error relative to CCSD(T)/aVTZ is driven down quite dramatically by weakening the regularization (increasing  $\kappa$ ). Indeed, perhaps surprisingly, we find that computing energies with unregularized MP2 and scaled unregularized MP3 provides the lowest error. In this case, the optimal  $c_3$  parameter was found to be 0.8 and yields a RMSD of only 1.59 kcal mol<sup>-1</sup>, which is close to chemical accuracy. We also observe that increasing the regularization strength decreases the optimal fraction of regularized third-order correlation energy. Impressed with the performance on this set, we chose this method, which we denote as MP2.8: $\kappa$ -OOMP2, as our candidate for further evaluation.

In order to assess the transferability of MP2.8: $\kappa$ -OOMP2, we tested its performance on a series of benchmarks sets meant to encompass a variety of main group bonded and non-bonded interactions: the non-MR subset of the W4-11 set[272] (the training set), the BH76RC set[253, 305, 306], the RSE43 set[273, 307], the HTBH38 set[306], the NHTBH38 set[253], the TA13 set[274], and the A24 set[308]. We compare the performance of MP2.8: $\kappa$ -



Figure 4.2: Scans of the root mean square deviation on the non-MR subset of the W4-11 thermochemical data set, in kcal mol<sup>-1</sup>, for the regularized, second- and third-order correlation energy functional given in Equation 4.6. The optimal value of  $c_2$  was found to be 1.0 for all  $\kappa_2$  values plotted. SCF references were generated via  $\kappa$ -OOMP2 orbital optimization. The basis set used was aVTZ. Reference values are calculated at the CCSD(T)/aVTZ level of theory.

OOMP2 against an unscaled version of the method (MP3: $\kappa$ -OOMP2), CCSD[22, 23], MP2.5[294], MP3,  $\kappa$ -OOMP2[39], OOMP2, and MP2. Details of the computations (aVTZ basis, CCSDT(T) reference, no frozen core) are the same as given previously.

Table 4.1 presents the the RMSDs, mean signed deviations (MSD), minimum deviations (MIN), and maximum deviations (MAX), in kcal mol<sup>-1</sup>, for the non-MR subset of the W4-11 set (the training set). This set includes atomization energies (TAE140), bond dissociation energies (BDE99), heavy atom transfer energies (HAT707), nucleophilic substitution reaction energies (SN13), and isomerization energies (ISOMERIZATION20).[272] We see that CCSD has a RMSD of 4.94 kcal mol<sup>-1</sup> and a MSD of 1.49 kcal mol<sup>-1</sup>. MP2, MP2.5 and MP3 on top of UHF orbitals yields RMSDs of 11.99, 8.97, and 9.24 kcal mol<sup>-1</sup>, respectively. The use of  $\kappa$ -OOMP2 optimized orbitals for the computation of the MP3 energy reduces the error over the use of UHF orbitals by a remarkably large factor of 3. MP2.8/ $\kappa$ -OOMP2 yields a RMSD of 1.59 kcal mol<sup>-1</sup> and a MSD of -0.45 kcal mol<sup>-1</sup>, which is 6 times smaller than MP3. This is also an improvement on  $\kappa$ -OOMP2 optimized orbitals outperform CCSD on this data set.

To validate the performance for thermochemistry outside of the training set, we tested MP2.8: $\kappa$ -OOMP2 on the BH76RC[253, 305, 306] and RSE43[273, 307] sets. The BH76RC set contains reaction energies for 30 reactions involved in the HTBH38 and NHTBH38 sets.[253, 305, 306] On this set MP2.8: $\kappa$ -OOMP2 outperforms all other methods surveyed with an

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Method	RMSD	MSD	MIN	MAX
CCSD	4.94	1.49	-8.60	20.34
MP2.8: $\kappa$ -OOMP2	1.59	-0.45	-5.83	5.24
MP3: $\kappa$ -OOMP2	3.22	0.33	-8.12	14.63
MP2.5	8.97	-2.85	-40.24	24.87
MP3	9.24	-0.80	-38.57	31.66
$\kappa$ -OOMP2	7.58	-3.14	-38.94	13.56
OOMP2	10.82	-3.50	-48.81	17.87
MP2	11.99	-4.90	-51.14	27.28

Table 4.1: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal mol<sup>-1</sup> for the non-MR subset of the W4-11 set. All calculations use the aVTZ basis; CCSD(T) is used for the reference values.

Method	RMSD	MSD	MIN	MAX
CCSD	1.905	-0.645	-7.175	1.977
MP2.8: $\kappa$ -OOMP2	0.835	-0.143	-1.465	1.534
MP3: $\kappa$ -OOMP2	1.574	-0.437	-6.213	1.228
MP2.5	4.625	-0.536	-21.407	9.904
MP3	4.511	-0.975	-21.604	4.232
$\kappa$ -OOMP2	4.220	-0.276	-9.763	11.856
OOMP2	5.524	0.836	-10.010	20.496
MP2	6.341	-0.098	-21.211	15.577

Table 4.2: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the BH76RC set.

RMSD of 0.84 kcal mol<sup>-1</sup> and a MSD of -0.14 kcal mol<sup>-1</sup>. MP3: $\kappa$ -OOMP2 also performs very well. Of the  $\mathcal{O}(N^5)$  methods,  $\kappa$ -OOMP2 provides the lowest RMSD while MP2 provides the lowest overall MSD. The largest absolute error using the canonical MP methods is for the H + F<sub>2</sub>  $\longrightarrow$  HF + H reaction energy. This can be traced back to spin-contamination at the UHF level in the case of F<sub>2</sub> with a  $\langle S^2 \rangle$  of 0.293. Both  $\kappa$ -OOMP2 and MP2.8: $\kappa$ -OOMP2 show significant improvement on this case with errors of 1.425 kcal mol<sup>-1</sup> and 0.536 kcal mol<sup>-1</sup>, respectively. MP2.8: $\kappa$ -OOMP2 improves upon  $\kappa$ -OOMP2 in all but two cases in this set.

Table 4.3 contains benchmark results for the RSE43 set. The RSE43 set contains reaction energies for hydrogen abstraction from hydrocarbons by a methyl radical.[273, 307] For this set we see MP2.8: $\kappa$ -OOMP2, with an RMSD of 0.63 kcal mol<sup>-1</sup> and a MSD of -0.54 kcal mol<sup>-1</sup>, performs slightly worse than CCSD and MP3: $\kappa$ -OOMP2. However the RMSD is still almost 4 times smaller than MP3. MP2.8: $\kappa$ -OOMP2, MP3: $\kappa$ -OOMP2, and OOMP2 all underestimate the reaction energies on average. For the  $\mathcal{O}(N^5)$  methods,  $\kappa$ -OOMP2 outperforms OOMP2 and MP2 both in terms of RMSD and MSD. Several species in the set exhibit spin contamination, leading to failures of the canonical MP methods.

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Method	RMSD	MSD	MIN	MAX
CCSD	0.446	0.316	-0.815	0.973
$\mathrm{MP2.8}{:}\kappa\text{-}\mathrm{OOMP2}$	0.634	-0.538	-1.726	0.050
MP3: $\kappa$ -OOMP2	0.521	-0.416	-1.550	-0.002
MP2.5	3.234	1.907	0.061	12.899
MP3	2.433	1.563	0.109	9.361
$\kappa$ -OOMP2	0.476	0.119	-0.964	1.020
OOMP2	0.922	-0.607	-2.261	0.478
MP2	4.099	2.252	-0.028	16.445

Table 4.3: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the RSE43 set.

Method	RMSD	MSD	MIN	MAX
CCSD	2.206	1.877	-0.782	4.146
$\mathrm{MP2.8}{:}\kappa\text{-}\mathrm{OOMP2}$	0.711	-0.120	-1.424	1.301
MP3: $\kappa$ -OOMP2	0.730	0.346	-1.411	1.755
MP2.5	3.686	3.246	-0.273	7.323
MP3	3.883	3.506	0.695	7.214
$\kappa$ -OOMP2	2.918	1.658	-1.434	9.558
OOMP2	3.479	-0.952	-7.152	8.566
MP2	4.044	2.986	-1.487	12.142

Table 4.4: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the HTBH38 set.

To evaluate the performance of MP2.8/ $\kappa$ -OOMP2 on kinetics, we tested it on the HTBH38[306] and NHTBH38[253] data sets. The HTBH38 set contains forward and reverse barrier heights for 19 hydrogen transfer reactions.[306] Results for this set are presented in Table 4.4. On this set MP2.8: $\kappa$ -OOMP2 (and MP3: $\kappa$ -OOMP2) outperform the other methods surveyed with a RMSD of 0.71 kcal mol<sup>-1</sup> (and 0.73 kcal mol<sup>-1</sup>), corresponding to chemical accuracy. The RMSDs are around 3 times smaller than that for CCSD. MP2, MP2.5, and MP3 overestimate the barrier heights in nearly all cases in the test set, with worst performances for the HF + H  $\longrightarrow$  H<sub>2</sub> + F, HF + H  $\longrightarrow$  H<sub>2</sub> + F, and OH + NH<sub>3</sub>  $\longrightarrow$  H<sub>2</sub>O + NH<sub>2</sub> forward barriers, respectively. MP2.8: $\kappa$ -OOMP2 improves significantly on these cases with barrier height errors of 0.24 kcal mol<sup>-1</sup> and -0.51 kcal mol<sup>-1</sup>, respectively.

Assessment data for the NHTBH38[253] set are presented in Table 4.5. The NHTBH38 set contains forward and reverse barrier heights for 19 non-hydrogen transfer reactions. On this set, MP2.8: $\kappa$ -OOMP2 outperforms all other methods surveyed (RMSD of 0.76 kcal mol<sup>-1</sup>), with MP3: $\kappa$ -OOMP2 performing second best. The reduction in RMSD relative to MP3 is more than a factor of 8 for MP2.8: $\kappa$ -OOMP2. Remarkably, both methods improve substantially upon CCSD, with the improvement being more than a factor of 3 for MP2.8: $\kappa$ -OOMP2 MP2, MP2.5, and MP3 all exhibit large errors in the barrier heights for the reactions

Method	RMSD	MSD	MIN	MAX
CCSD	2.534	2.067	0.132	7.646
MP2.8: $\kappa$ -OOMP2	0.758	0.268	-0.949	1.579
MP3: $\kappa$ -OOMP2	1.668	1.076	-0.718	7.175
MP2.5	6.590	4.763	-0.328	24.455
MP3	6.651	5.158	1.099	23.283
$\kappa$ -OOMP2	2.766	1.553	-7.610	5.222
OOMP2	3.901	-1.650	-18.495	2.315
MP2	7.035	4.368	-2.676	25.627

Table 4.5: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the NHTBH38 set.

H + N<sub>2</sub>O → OH + N<sub>2</sub>, H + F<sub>2</sub> → HF + F, and CH<sub>3</sub> + ClF → CH<sub>3</sub>F + Cl. For H + N<sub>2</sub>O → OH + N<sub>2</sub> and CH<sub>3</sub> + ClF → CH<sub>3</sub>F + Cl, both forward and reverse barriers are overestimated due to spin-contamination of the UHF reference for the transition state structures. The UHF reference  $\langle S^2 \rangle$  values of 1.011 and 1.026, respectively, are corrected to mean-field  $\langle S^2 \rangle$  values of 0.765 and 0.775, respectively, via the  $\kappa$ -OOMP2 orbital optimization procedure. For H + F<sub>2</sub> → HF + F, the reverse barriers are overestimated by more than 20 kcal mol<sup>-1</sup> with MP2, MP2.5, and MP3 while errors in the forward barriers are of similar magnitude to other systems in the data set. Significant spin-contamination is present in the UHF reference for both F<sub>2</sub> and the transition state structure, leading to a cancellation of errors in the forward barrier that is not seen in the reverse barrier. Orbital optimization with  $\kappa$ -OOMP2 helps to mitigate this spin-contamination, reducing the mean-field  $\langle S^2 \rangle$  values of 0.293 and 1.212, respectively, to 0.000 and 0.767, respectively. For all three of these reactions MP2.8: $\kappa$ -OOMP2 gives errors that are reduced by a factor of 5-10 relative to MP2, MP2.5, and MP3.

To round out the test suite we assessed the performance of MP2.8: $\kappa$ -OOMP2 on two noncovalent interaction sets: the TA13 and A24 sets. The TA13 set contains 13 nonbonded interaction energies for radical closed-shell complexes.[274] We apply a counterpoise correction to these interaction energies to mitigate basis set superposition error (BSSE). Assessment data for the TA13 set is presented in Table 4.6. On this test set we see MP2.8: $\kappa$ -OOMP2 performs almost as well as CCSD. MP2.8: $\kappa$ -OOMP2 overbinds each interaction in the set, especially the H<sub>2</sub>O-Al interaction which is overbound by 2.05 kcal mol<sup>-1</sup>. Remarkably,  $\kappa$ -OOMP2 outperforms all methods surveyed on this set, with an RMSD of 0.35 kcal mol<sup>-1</sup>. Table 4.7 presents the counterpoise-corrected results for the A24 set. The A24 set contains noncovalent interaction energies for 24 closed-shell small molecule complexes.[308] MP2.8: $\kappa$ -OOMP2 outperforms all other methods with a RMSD of 0.08 kcal mol<sup>-1</sup> and a MSD of 0.01 kcal mol<sup>-1</sup>. For MP2, MP2.5, and MP3, artifactual spin-contamination at the UHF level causes underbinding for the ethene-ethene and ethene-ethyne dimers. For the ethene dimer, the MP2, MP2.5, and MP3 errors are in excess of 2 kcal mol<sup>-1</sup> while MP2.8: $\kappa$ -OOMP2 reduces this error to 0.04 kcal mol<sup>-1</sup>.

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Method	BMSD	MSD	MIN	MAX
	TUNDD	MDD	IVIII V	1111111
$\operatorname{CCSD}$	0.722	0.539	-0.259	1.470
MP2.8: $\kappa$ -OOMP2	0.823	-0.459	-2.054	-0.011
MP3: $\kappa$ -OOMP2	0.808	-0.442	-2.463	0.086
MP2.5	1.559	0.276	-3.888	3.708
MP3	1.435	0.391	-2.612	3.997
$\kappa$ -OOMP2	0.350	-0.019	-0.589	0.650
OOMP2	0.789	-0.149	-1.938	1.370
MP2	1.791	0.160	-5.164	3.419

Table 4.6: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the TA13 set.

Method	RMSD	MSD	MIN	MAX
CCSD	0.247	0.226	0.093	0.429
MP2.8: <i>κ</i> -OOMP2	0.075	0.007	-0.169	0.233
MP3: $\kappa$ -OOMP2	0.106	0.043	-0.113	0.373
MP2.5	0.492	0.132	-0.113	2.303
MP3	0.488	0.187	-0.010	2.203
$\kappa$ -OOMP2	0.184	-0.045	-0.631	0.199
OOMP2	0.193	-0.131	-0.475	0.063
MP2	0.515	0.078	-0.441	2.403

Table 4.7: Root mean square deviation, mean signed deviation, minimum deviation, and maximum deviation, in kcal  $mol^{-1}$ , for the A24 set.

### 4.4 Conclusion

Considering all the data presented, let us summarize the main conclusions obtained from this work.

- 1. At the MP2 level the choice of orbitals matters considerably, as is well known. In our work, for all 7 data sets considered, orbital optimized MP2 (OO-MP2) yields lower RMSD than MP2 (using unrestricted orbitals when necessary). Regularization via  $\kappa = 1.45$  has formal benefits in restoring Coulson-Fisher points. It also has practical benefits:  $\kappa$ -OOMP2 yields lower RMSD than OO-MP2 for all 7 data sets tested.
- 2. Use of  $\kappa$ -OOMP2 orbitals improves MP3 results to a surprising extent. MP3: $\kappa$ -OOMP2 has lower RMSD than MP3 by factors ranging from 1.7 to more than 5 across the 7 data sets reported here. MP3: $\kappa$ -OOMP2 is thus a far more robust method than MP3 itself, due to the reduced spin-contamination in  $\kappa$ -OOMP2 orbitals relative to HF orbitals.
- 3. Developing a semi-empirical variant based on scaling the MP2 and MP3 contributions yielded  $c_2 = 1.0$  and  $c_3 = 0.8$  based on the non-MR part of the W4-11 data set (no

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regularization is preferred). In transferability tests, this MP2.8: $\kappa$ -OOMP2 method improves over MP3: $\kappa$ -OOMP2 in 4 of our 6 test sets, with the other two being very similar.

- 4. Remarkably, the results obtained with MP3:κ-OOMP2 and MP2.8:κ-OOMP2 produce lower RMSD than CCSD itself in 5 of the 7 data sets (the remaining two show no large failures). This indicates a level of performance that is beyond the physical content of MP3 theory, and involves some rather systematic cancellation of the effects due to connected triples.
- 5. These improved MP3 methods are single reference of course, and the data sets considered here are suitable for single reference methods. Much poorer performance must be expected for systems where strong correlations are in play (perhaps with the exception of biradicaloids[275])

The improved performance granted by the use of  $\kappa$ -OOMP2 optimized orbitals suggests future developments in electronic structure theory. It will be interesting to assess results across additional data sets, and explore the use of larger basis sets. We intend to explore scaled fourth-order perturbation approaches (MP4) and coupled cluster methods with  $\kappa$ -OOMP2 reference orbitals. The latter would be especially interesting in the context of nonvariational failures of CCSD(T). In a different direction, perhaps MP3 should be considered as an independent descriptor of electron correlation in double hybrid density functional theory, where MP2 is at present most widely used.[71, 88, 283–285] With the advances in integral compression techniques such as tensor hypercontraction, both MP2 and MP3 energy evaluations scale quartically with system size.[291] Incorporating this into the development of new double hybrid density functionals will be a promising future direction.

### Chapter 5

## Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

### 5.1 Introduction

Coupled cluster theory with single, double, and perturbative triple excitations [CCSD(T)][25] with spin-restricted Hartree-Fock (RHF) orbitals is considered the "gold standard" by many quantum chemists for its ability to routinely provide results approaching chemical accuracy for energies and properties of closed-shell species at a reasonable computational cost [309– 311]. With the ability to achieve sub-kcal mol<sup>-1</sup> errors at a computational cost of  $\mathcal{O}(N^7)$  and a memory cost of  $\mathcal{O}(N^4)$ , where N is the size of the basis, CCSD(T) strikes an advantageous accuracy-to-cost balance between coupled cluster theory with single and double excitations (CCSD)[22] ( $\mathcal{O}(N^6)$  computational cost,  $\mathcal{O}(N^4)$  memory cost) and explicit treatment of the triple excitations (CCSDT)[24] ( $\mathcal{O}(N^8)$  computational cost,  $\mathcal{O}(N^6)$  memory cost). For openshell species, however, the performance of CCSD(T) (especially on top of spin-unrestricted HF (UHF) orbitals) is less clear. A study of bond lengths and frequencies of 33 small radical species by Byrd et. al [312] reported that CCSD(T) with a UHF reference (CCSD(T):UHF)shows little to no statistical improvement over CCSD[22] for geometries and frequencies. Beran et al.[313], in a study of harmonic vibrational frequencies of diatomic radicals, reported poor behaviour of CCSD(T):UHF for several species in the set including,  $CO^+$  and NO. Subsequent studies of these problem systems by Szalay et. al [314] investigated the source of this mixed performance. Stanton and Gauss[315] reviewed several potential factors leading to this discrepancy in describing open-shell species, including multireference (MR) character, spin-contamination, symmetry breaking, and instabilities and near-instabilities in the reference wavefunction. The use of restricted open-shell HF (ROHF) orbitals has been found to somewhat improve the performance on vibrational frequencies over UHF, though the former are prone to errors due to spatial symmetry breaking[315, 316]. While in the limit of full

configuration interaction the energy and properties are invariant to the choice of reference orbitals, any truncated approximate method will incur some level of orbital dependence.

Several strategies have been proposed as alternative references to UHF for CC calculations. The use of Brueckner orbitals in CC theory (BCC), which by definition have no singles contribution to the coupled cluster wavefunction, attempts to incorporate the most important electron correlations at the level of the reference [317–321]. BCC approaches have been known to preserve wavefunction symmetries as well, yielding more accurate properties [322–325]. In a similar vein, the molecular orbitals (MOs) can be directly optimized in the presence of CCD correlation energies as in the orbital-optimized coupled cluster doubles (OD) and orbital optimized coupled-cluster doubles and perturbative triples [OD(T)]approaches [326–328]. The use of OD(T) by Beran et al. [313] was shown to significantly suppress errors in the computed harmonic frequencies of CN, CO<sup>+</sup>, and NO compared to CCSD(T):UHF. Brueckner and/or optimized coupled cluster orbitals are costly to obtain; BCCD(T) and OD(T) typically are far more computationally expensive than CCSD(T) as orbital optimization is often more challenging than varying the singles amplitude. Given this steep computational cost, methods to approximate Brueckner or otherwise optimized orbitals at a lower cost are highly desirable.

One such approximate approach to incorporate correlation into the reference orbitals for a correlated calculation is the use of Kohn-Sham density functional theory (DFT) orbitals as a reference. The use of BLYP[55, 56] orbitals by Beran et al.[313] was shown to substantially improve the computed vibrational frequencies of radical diatomic species over UHF orbitals. More recently, Fang et al.[329] and Fang et al.[330] have demonstrated the efficacy of CCSD(T) with DFT orbitals in the prediction of thermodynamic properties of UCl<sub>6</sub> and several diatomic transition metal compounds, respectively. Aside from their use in CC calculations, DFT orbital references have been used successfully for excited state configuration interaction[331], as guiding functions in quantum Monte Carlo [332], and for second-order perturbation theory in the context of double-hybrid DFT[86]. In addition to their inexpensive computational cost ( $\mathcal{O}(N^3)$ ), DFT orbitals offer improved stability against symmetry breaking compared to HF orbitals[333]. A connection between DFT orbitals and Brueckner orbitals has been proposed by several researchers as well[334–336].

Orbital-optimized second-order Møller-Plesset perturbation theory (OOMP2) and its variants offer another way to approximate higher-order orbital optimized-methods at a cost of  $\mathcal{O}(N^5)$  per iteration[35, 36, 39]. Orbital optimization at the MP2 level, in addition to improving energetics, is often seen to reduce spin-contamination in the optimized reference[35, 36, 269]. Recently Haggag et al.[337] utilized OOMP2 reference orbitals for CC calculations on the triplet state of permanganate to combat spin contamination seen at the UHF level. Despite these benefits, OOMP2 exhibits three unsatisfying characteristics that limit its application: divergence in the cases of small orbital energy gaps[40], the loss of Coulson-Fischer points[271], and "artificial" symmetry restoration[276, 338]. The correlation energy

functional for MP2 is

$$E_{\rm MP2} = -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}}$$
(5.1)

where  $\Delta_{ij}^{ab} = \epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j$  is the non-negative orbital energy denominator. This energy is seen to diverge in cases where the denominator becomes small, as can occur when stretching bonds. This behavior leads to poor performance of OOMP2 when predicting harmonic vibrational frequencies[40]. Secondly, OOMP2 often fails to continuously transition from restricted to unrestricted solutions even when the unrestricted solution is lower in energy[270]. Thirdly, OOMP2 has been shown in some cases to "artificially" restore spin-symmetry to systems where the spin symmetry breaking is an "essential" feature of the system due to MR character of the system[276, 338].

In order to address the problematic aspects of OOMP2, two of us developed  $\kappa$ -OOMP2, a regularized variant of OOMP2[39]. The  $\kappa$ -OOMP2 energy functional is given by

$$E_{\kappa\text{-OOMP2}}(\kappa) = -\frac{1}{4} \sum_{ijab} \frac{|\langle ij||ab\rangle|^2}{\Delta_{ij}^{ab}} \left(1 - e^{-\kappa \Delta_{ij}^{ab}}\right)^2, \tag{5.2}$$

where  $\kappa$  is a regularization parameter that damps contributions to the correlation energy when the orbital energy denominator becomes small. Regularization parameter values  $\kappa \leq$  $1.5 E_h^{-1}$  were shown to restore Coulson-Fischer points for a series of bond-breaking curves[39]. Training of the regularization parameter on the W4-11 thermochemistry data set[272] led to an optimal  $\kappa$  value of 1.45  $E_h^{-1}$ [39]. With this parameter value,  $\kappa$ -OOMP2 was able to outperform OOMP2 on the TA13 data set[274] of radical–closed-shell interaction energies[39]. Further application to symmetry breaking in fullerenes revealed the ability of  $\kappa$ -OOMP2 to distinguish between essential and artificial symmetry breaking[276, 338].

We recently developed a scaled variant of third-order MP theory (MP3) that utilizes  $\kappa$ -OOMP2 orbitals as a reference which we will denote as MP2.8: $\kappa$ -OOMP2[95]. MP2.8: $\kappa$ -OOMP2 and its unscaled version, MP3: $\kappa$ -OOMP2, outperformed CCSD on five of the seven data sets investigated at the cost of a single  $\mathcal{O}(N^6)$  iteration. The use of  $\kappa$ -OOMP2 orbitals strongly improves upon the performance of MP3 with UHF orbitals as well, especially in cases of spin-symmetry breaking.

Inspired by MP2.8: $\kappa$ -OOMP2[95], the work of Beran et al.[313], and the success of  $\kappa$ -OOMP2 in treating radical species in the TA13 set[39] and in producing minimally spincontaminated references for biradicaloid systems[275], in this work we explore the use of  $\kappa$ -OOMP2 orbitals as a reference for CCSD(T) computation of vibrational frequencies. Errors are calculated with respect to experimental values and compared against CCSD(T) with UHF orbitals, OOMP2 orbitals, and several flavors of DFT orbitals. This use of  $\kappa$ -OOMP2 orbitals as a reference for CCSD(T) was previously explored in the computation of spin-gaps in an iron porphyrin complex and showed an improvement over conventional CCSD(T) [339].

#### 5.2 Computational Methods

We consider eight methods as generators of MOs for use as references: UHF, two OOMP2 methods (OOMP2[35–37] and  $\kappa$ -OOMP2[39]), and five density functionals (BLYP[55, 56], B97M-rV[340], B97[62],  $\omega$ B97X-V[82], and  $\omega$ B97M-V[256]). A regularization parameter value of  $\kappa = 1.45 E_h^{-1}$  was chosen for  $\kappa$ -OOMP2, as suggested by Lee and Head-Gordon[39]. Both the OOMP2 and  $\kappa$ -OOMP2 calculations were carried out using the resolution-of-theidentity (RI) approximation [29, 30]. The functionals B97M-rV, B97,  $\omega$ B97X-V, and  $\omega$ B97M-V were chosen on the basis of their performance in a recent benchmark of over 200 density functionals in which they were found to be the best performing meta-GGA, global hybrid GGA, range-separated hybrid GGA, and range-separated hybrid meta-GGA functionals, respectively[85]. The DFT calculations were performed using an ultra-fine integration grid of 99 radial points and 590 angular points per atom. Wavefunction stability analysis 270, 341] was performed on the UHF and DFT solutions to ensure that all orbitals used properly minimize their corresponding SCF energies. All unstable solutions (saddle points) were displaced and reoptimized to local minima. OOMP2 and  $\kappa$ -OOMP2 calculations were performed starting from a locally stable UHF solution. Both the reference calculations and CCSD(T) calculations were performed using unrestricted wavefunctions.

All calculations were performed with the aug-cc-pwCVTZ basis[119, 120, 302, 342, 343] set to capture the effects of core correlation. A core-valence basis set was utilized to account for the role of core-valence correlations in molecular property calculations. The use of the weighted, triple- $\zeta$  variant is justified by the faster convergence of properties to the complete basis set limit seen with cc-pwCVnZ over cc-pCVnZ[342]. Augmented functions were chosen to better treat the anions in the data set. The corresponding auxiliary basis set was utilized for the OOMP2 and  $\kappa$ -OOMP2 calculations (with cc-pwCVQZ-RI utilized for Li, Be, Na, and Mg)[303, 304, 344]. Atoms H-F have all electrons correlated and atoms Na-Cl utilize a frozen He core. All electronic structure calculations were performed using the Q-Chem package of electronic structure programs[125]. For a given reference, the corrected harmonic vibrational frequency was determined by fitting a quartic polynomial to seven equally-spaced points (0.005 Å between adjacent points) distributed near the minimum of the potential and applying Dunham analysis to the fitting coefficients to account for the effects of rotation[345]. Equilibrium bond lengths from the fitting procedure are reported in Tab. D.2 and D.3.

### 5.3 Experimental Data Selection

For simplicity and due to the lack of analytic first derivatives of the CCSD(T) energy with respect to nuclear displacements for non-HF references, we restricted the systems of study in this work to diatomic species for which the ground state potential energy surface can be determined via fitting to single point calculations. Beginning from all diatomics of row 2 and row 3 species (and hydrogen) for which Huber and Herzberg[346] report ground state frequencies, we include all species for which we were able to compute a smooth potential

energy surface about the corresponding equilibrium bond length for each method (excluded species are listed in Tab. D.7). The overall data set contains 36 closed-shell species (29 neutrals, 6 cations, 1 anion) and 59 open-shell species (38 neutrals, 15 cations, 6 anions; 46 doublets, 13 triplets). Among other species in this data set, we include all 12 species from Beran et al.[313] as well as several isovalent analogues of these species containing row 3 elements. Other notable inclusions are B<sub>2</sub>  $(X^3\Sigma_g^-)$  and C<sub>2</sub>  $(X^1\Sigma_g^+)$ , both known to exhibit MR behavior in their ground states[272, 347], and F<sub>2</sub>  $(X^1\Sigma_g^+)$ , a biradicaloid diamagnetic system know that is unbound at the UHF level of theory[348]. Where available, experimental frequencies were updated with data from Irikura[349]. Frequencies for OH<sup>-</sup>, F<sub>2</sub><sup>+</sup>, and SO<sup>+</sup> were updated with more recent experimental data from Hotop et al.[350], Cormack et al.[351], and Milkman et al.[352], respectively.

### 5.4 Results and Discussion

### 5.5 Vibrational frequencies

Fig. 5.1 presents the errors in the corrected vibrational frequencies on the closed- and openshell subsets for each method as box plots. In these plots, the boxed region represents data from the first to the third quartile of the distribution, a red line marks the median of the data, whiskers enclose all data within 1.5 times the inter-quartile distance of the upper and lower box edges, and points mark data lying outside of these regions. Tab. 5.1 presents the root mean square deviations (RMSD), mean signed deviations (MSD), most negative deviations (MIN), and most positive deviations (MAX), all in cm<sup>-1</sup>, of the corrected vibrational frequencies from the experimental frequencies for the overall data set.  $CCSD(T):\kappa$ -OOMP2 is seen to give the best overall performance in terms of RMSDs with a value of 17.66 cm<sup>-1</sup>, reducing the RMSD for CCSD(T):UHF by more than a factor of two. CCSD(T):OOMP2 is seen to perform slightly worse but still improves on the CCSD(T):UHF RMSD by a factor of 2. The performance of the DFT orbital approaches is hindered by the presence of C<sub>2</sub>, representing the MIN value for all functionals tested. The MAX value, corresponding to PH<sup>+</sup>, is shared among all non-HF methods tested. On average the frequencies are slightly blue-shifted for CCSD(T):UHF and slightly red-shifted for all other methods.

#### Closed-shell subset

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Figure 5.1: Box plots (overall, left, and enhanced, right) of the errors in corrected vibrational frequencies (in  $cm^{-1}$ ) are presented. Red lines mark the median errors, boxes bound the central 50% of the data, whiskers enclose all data points within 1.5 times the inter-quartile range of the box edges, and points denote outlying data.



Table 5.1: Root mean square deviations, mean signed deviations, most negative deviations, and most positive deviations (all in  $\text{cm}^{-1}$ ) for the predicted corrected vibrational frequencies of all species are summarized for the CCSD(T) methods utilizing different molecular orbital references.

	$\Delta(\text{CCSD}(\text{T})$	$\Delta(\text{CCSD}(T))$	: $\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(T)$	$: \Delta(\text{CCSD}(T))$	$: \Delta(\text{CCSD}(T))$	: $\Delta(\text{CCSD}(\text{T}))$ :	$\Delta(\text{CCSD}(\text{T})$
	UHF)	κ-	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M$ -
		OOMP2)						V)
RMSD	41.05	17.66	18.82	24.75	26.75	26.35	31.19	29.77
MSD	4.35	-2.19	-1.80	-4.99	-4.36	-4.05	-4.71	-5.45
MIN	-240.71	-64.06	-64.26	-142.01	-186.68	-187.50	-246.37	-184.40
MAX	177.65	90.68	88.43	87.87	86.24	88.25	89.26	89.52

for the	36 closed	-shell spe	ecies are pre	sented in	for the CC	SD(T) me	thods uti	lizing diffe	cent molec	ular orbit	als. Root
mean so	quare dev	iations, 1	mean signed	deviation	s, most neg	sative devi	ations, an	id most pos	itive devi	ations (all	in $\mathrm{cm}^{-1}$ )
for the	set of spe	cies and	subsets are	also presei	nted.						
Row 2– Row 2	Dimer	State	Expt.	$\Delta(\text{CCSD}(T))$ UHF)	$\Delta(CCSD(T))$ $\kappa$ -OOMP2)	$\Delta(CCSD(T)):$ 00MP2)	$\Delta(\text{CCSD}(T)$ BLYP)	$: \Delta(CCSD(T)): B97M-rV)$	$\Delta(CCSD(T))$ B97)	$\Delta(CCSD(T); \omega B97X-V)$	$\Delta(\text{CCSD}(T))$ $\omega B97M-V)$
	LiH	$X^{1}\Sigma^{+}$	$1405.49805^{a}$	-7.12	-8.81	-2.59	-8.41	-9.83	-2.50	-8.68	-6.13
	$\mathrm{Li}_2$	$X^1 \Sigma_a^+$	$351.4066^{a}$	-11.40	-16.69	-16.52	-15.37	-13.54	-1.76	-16.24	-11.40
	LiF	$X^1\Sigma^+$	$910.57272^{a}$	-5.10	-5.60	-6.23	-6.96	-6.37	-6.23	-6.09	-6.07
	${ m BeH^+}$	$X^1\Sigma^+$	$2221.7^{ m b}$	-11.81	-13.69	-12.35	-16.41	-13.93	-16.22	-11.43	-10.61
	BeO	$X^1\Sigma^+$	$1487.32^{\rm b}$	-20.79	-17.36	-21.83	-20.41	-19.22	-20.00	-18.16	-18.40
	BH	$X^1\Sigma^+$	$2366.7296^{a}$	-2.51	0.49	-1.09	-1.00	-1.63	-1.02	-0.97	2.51
	$\mathrm{BF}$	$X^1\Sigma^+$	$1402.15865^{a}$	-10.53	-12.64	-13.60	-15.41	-13.50	-13.60	-12.88	-13.02
	$C_2$	$X^1 \Sigma_g^+$	$1855.0663^{a}$	52.18	-29.19	-15.21	-142.01	-186.68	-187.50	-246.37	-184.40
	<u>c</u> o	$X^1\Sigma^+$	$2169.75589^{a}$	-16.15	-18.53	-19.25	-20.34	-18.52	-18.76	-18.10	-18.13
	$ m N_2$	$X^1\Sigma_g^+$	$2358.57^{a}$	-4.81	-5.92	-6.09	-6.74	-5.86	-6.13	-5.77	-5.85
	$^{+}ON$	$X^1\Sigma^+$	$2376.72^{a}$	-7.36	-9.86	-9.73	-11.48	-10.19	-9.78	-9.70	-9.60
	-HO	$X^1\Sigma^+$	$3735.2^{c}$	-0.73	-2.05	-6.68	-0.99	0.86	-0.98	0.55	-0.33
	НF	$X^1\Sigma^+$	$4138.385^{a}$	-7.04	-9.66	-8.96	-8.06	-7.40	-8.24	-7.48	-7.40
	$\mathrm{F}_2$	$X^1 \Sigma_g^+$	$916.929^{a}$	60.05	3.10	3.57	3.20	3.06	3.00	2.93	2.73
	RMSD			23.32	13.29	11.97	39.74	51.03	51.14	66.67	50.25
	MSD			0.49	-10.46	-9.75	-19.31	-21.62	-20.69	-25.60	-20.44
	MIN			-20.79	-29.19	-21.83	-142.01	-186.68	-187.50	-246.37	-184.40
	MAX			60.05	3.10	3.57	3.20	3.06	3.00	2.93	2.73
Row $2-$	Dimer	State	Expt.	$\Delta(CCSD(T))$	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(\text{CCSD}(T))$	$: \Delta(CCSD(T))$	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :
Row $3$				UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	$\operatorname{NaH}$	$X^1\Sigma^+$	$1171.968^{a}$	-8.85	-9.09	-9.02	7.76	-9.10	7.98	7.86	8.02
	NaLi	$X^1\Sigma^+$	$256.5412^{a}$	-1.17	-0.85	-3.97	-29.30	-1.75	-12.05	-1.76	-1.76
	$\operatorname{NaF}$	$X^1\Sigma^+$	$535.65805^{a}$	-5.05	-5.43	-5.87	-6.35	-5.99	-5.62	-5.61	-5.57
	$MgH^+$	$X^1\Sigma^+$	$1699.1^{\rm b}$	-10.18	-12.48	-12.48	-12.96	-11.67	-10.62	-11.98	-12.33
	AlH	$X^1\Sigma^+$	$1682.37474^{a}$	-15.37	-13.07	-12.95	-15.55	-21.21	-12.58	-16.38	-12.04
	AIF	$X^1\Sigma^+$	$802.32447^{a}$	-8.87	-9.64	-10.18	-11.23	-10.23	-10.19	-9.86	-9.95
	$SiH^+$	$X^1\Sigma^+$	$2157.17^{a}$	-4.24	-1.20	-4.25	-1.28	-4.20	-2.62	-6.51	-2.32
	SiO	$X^1\Sigma^+$	$1241.54388^{a}$	-10.51	-11.60	-15.61	-14.30	-12.50	-12.45	-11.07	-11.14
	PN	$X^1\Sigma^+$	$1336.948^{a}$	33.01	-5.90	-7.60	-6.56	-5.54	-5.84	-5.45	-5.57
	$\operatorname{BeS}$	$X^1\Sigma^+$	$997.94^{a}$	-240.71	-8.22	-8.10	-7.73	-8.00	-7.74	-7.96	-7.20
	CS	$X^1\Sigma^+$	$1285.08^{\mathrm{b}}$	-5.98	-7.54	-9.54	-9.31	-8.00	-7.98	-6.94	-7.13
	$NS^+$	$X^1\Sigma^+$	$1415^{b}$	52.10	3.37	3.04	2.40	3.51	3.15	3.47	3.24
	HCI	$X^1\Sigma^+$	$2990.9248^{a}$	16.30	14.11	15.27	16.82	16.86	15.88	16.93	16.10
	LiCl	$X^1\Sigma^+$	$642.95453^{a}$	-7.45	-7.55	-7.57	-7.65	-7.58	-7.73	-7.49	-7.82
	BCI	$X^1\Sigma^+$	$840.29472^{a}$	-2.91	-2.44	-3.45	-4.81	-4.04	-3.51	-2.66	-2.84
	CCI+	$X^1\Sigma^+$	$1175^{\mathrm{b}}$	2.19	-4.74	-11.35	-9.69	-6.52	-6.24	-3.41	-3.94
	CIF	$X^1\Sigma^+$	$783.4534^{a}$	-7.29	-7.02	-6.99	-6.08	-5.87	-6.00	-6.31	-6.33
	RMSD			60.79	8.31	9.47	11.86	9.68	8.89	8.83	8.23
	USN ,			-13.23	-5.25	-6.51	18.9-	-5.99	-4.95	-4.42	-4.03
	MIN			-240.71	-13.07	-15.61	-29.30	-21.21	-12.58	-16.38	-12.33

Table 5.2: Experimental vibrational frequencies (in  $\text{cm}^{-1}$ ) and errors (in  $\text{cm}^{-1}$ ) in the corrected vibrational frequencies

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16.10	$\Delta(\text{CCSD}(T \ \omega \text{B97M-V}))$	-5.29	-7.47	-8.43	-4.52	-11.63	7.88	-7.47	-11.63	-4.52	$\Delta(\text{CCSD}(T))$	$\omega B97M-V$	31.98	-10.89	-184.40	16.10
16.93	∆(CCSD(T): µB97X-V)	-5.27	-7.44	-8.42	-4.47	-11.66	7.87	-7.45	-11.66	-4.47	$\Delta(CCSD(T))$ :	$\mu B97X-V$	42.12	-13.08	-246.37	16.93
15.88	Δ(CCSD(T): . B97)	-5.31	-7.60	-8.96	-4.48	-11.59	8.01	-7.59	-11.59	-4.48	$\Delta(CCSD(T))$ :	B97)	32.61	-11.44	-187.50	15.88
16.86	$\Delta(CCSD(T): Z$ B97M-rV)	-5.41	-7.70	-9.05	-4.38	-11.63	8.06	-7.64	-11.63	-4.38	$\Delta(CCSD(T): Z$	B97M-rV)	32.65	-12.30	-186.68	16.86
16.82	$\Delta(\text{CCSD}(T); Z$	-5.31	-7.88	-9.62	-4.68	-11.64	8.25	-7.83	-11.64	-4.68	$\Delta(CCSD(T): 2$	BLYP)	26.27	-11.82	-142.01	16.82
15.27	$\Delta(\text{CCSD}(T); \lambda)$	-2	-7.50	-9.41	-4.67	-11.69	8.14	-7.71	-11.69	-4.67	$\Delta(\text{CCSD}(T))$ : $_{-}$	OOMP2)	10.36	-7.94	-21.83	15.27
14.11	$\Delta(CCSD(T))$ : -00MP2)	- 5.27	-7.38	-8.35	-4.41	-11.64	7.84	-7.41	-11.64	-4.41	$\Delta(CCSD(T))$ :	v-OOMP2)	10.48	-7.58	-29.19	14.11
52.10	$\Delta(\text{CCSD}(T); \lambda)$	-5.25	-7.43	-8.84	32.77	-11.83	16.58	-0.11	-11.83	32.77	$\Delta(\text{CCSD}(T))$ : $_{-}$	UHF)	44.66	-6.07	-240.71	60.05
	Expt.	$364.6842^{a}$	$481.77466^{a}$	$749.64559^{a}$	$780.77^{a}$	$559.751^{a}$										
	State	$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma_g^+$	$X^1 \Sigma_g^+$	0									
MAX	Dimer	NaCl	AICI	$\operatorname{SiS}$	$\mathbf{P}_2$	$\overline{\mathrm{Cl}}_2$	RMSD	MSD	MIN	MAX			RMSD	MSD	MIN	MAX
	Row 3– Row 3										Closed-	$_{ m shell}$				

 $\begin{array}{c} \mbox{Chapter 5. Polishing the gold standard: The role of orbital choice in $CCSD(T)$ frequency prediction $83$ } \\ \end{array}$ 

<sup>a</sup> From Ref. 349. <sup>b</sup> From Ref. 346. <sup>c</sup> From Ref. 350.

Tab. 5.2 presents the experimental vibrational frequencies, in cm<sup>-1</sup>, and errors in the corrected vibrational frequencies, in cm<sup>-1</sup>, for the 36 species in the closed-shell subset. We see that the use of OOMP2 or  $\kappa$ -OOMP2 orbitals is able to reduce the RMSD of UHF orbitals by a factor of 4 from 44.66 cm<sup>-1</sup> to 10.36 and 10.48 cm<sup>-1</sup>, respectively. The use of DFT orbitals is also seen to reduce the RMSD, with BLYP outperforming the other functionals and  $\omega$ B97X-V improving the RMSD by only 2.5 cm<sup>-1</sup>. On average, all methods are seen to red-shift the closed-shell frequencies by 5-12 cm<sup>-1</sup>.

The performance of CCSD(T) with DFT orbitals is most negatively impacted for the  $C_2$  system, with errors in excess of -100 cm<sup>-1</sup> for each. This data point also represents a negative outlier for CCSD(T) with  $\kappa$ -OOMP2 orbitals and a positive outlier for CCSD(T)with UHF orbitals. Its significant MR character renders  $C_2$  outside of the scope of the single-reference methods evaluated in this work [272, 347]. Judging from the mean-field  $\langle S^2 \rangle$ values (Tab. D.8) for  $\kappa$ -OOMP2 and OOMP2 of 0.89140 and 0, respectively, we observe that OOMP2 is "artificially" restoring symmetry in this case while the broken spin-symmetry of the  $\kappa$ -OOMP2 reference orbitals is diagnostic of a MR problem, consistent with the literature [272, 347]. Spin-symmetry breaking in  $\kappa$ -OOMP2 as a signal of MR character has been demonstrated earlier in a study of fullerenes by Lee and Head-Gordon [276]. Considering the other dimers in the closed-shell subset, CCSD(T) with UHF orbitals is seen to exhibit poor performance for  $F_2$ , PN, BeS, NS<sup>+</sup>, and  $P_2$ , especially given that all other methods have absolute errors below  $10 \text{ cm}^{-1}$  for these species. These errors can be attributed to spin contamination of the reference orbital set. For UHF, the mean field  $\langle S^2 \rangle$  values for F<sub>2</sub>, PN, BeS, NS<sup>+</sup>, and P<sub>2</sub> are 0.31922, 0.70716, 1.01647, 0.69154, and 0.67604, respectively while the mean-field  $\langle S^2 \rangle$  values of  $\kappa$ -OOMP2, OOMP2, and the five DFT functionals for these species are all zero. The large error for BeS at the CCSD(T):UHF level can be traced to the character of the UHF wavefunction, which localizes nearly an entire electron's spin density on each atomic center. The  $\langle S^2 \rangle$  value of 1.01647 shows significant triplet contamination. In the absence of these symmetry-broken species (and MR  $C_2$ ), the performance across the methods is largely equalized with RMSDs ranging from 9.5 to  $12 \text{ cm}^{-1}$ . The maximum error outlier for all non-HF references is attributed to HCl. NaLi is also an outlying case for  $CCSD(T):BLYP (-29.30 \text{ cm}^{-1}).$ 

#### **Open-shell subset**

mean sq	luare dev.	iations, n	nean signed	deviation	s, most neg	gative devi	ations, an	d most pos	itive devi	ations (all	in $\mathrm{cm}^{-1}$ )
for the $i$	set of spe	cies and s	subsets are	presented.							
Row 2– Row 2	Dimer	State	Expt.	$\Delta(\text{CCSD}(T))$ : UHF)	: $\Delta(CCSD(T))$ : $\kappa$ -OOMP2)	$\Delta(\text{CCSD}(T))$ 00MP2)	$\Delta(CCSD(T))$ BLYP)	$\Delta(CCSD(T))$ : B97M-rV)	$\Delta(CCSD(T))$ B97)	: $\Delta(CCSD(T))$ : $\omega B97X-V$ )	$\Delta(CCSD(T))$ : $\omega B97M-V)$
	LiO	$X^2 \Pi_i$	$814.62^{a}$	-6.94	-7.10	-7.22	-8.16	-7.02	-7.04	-6.97	-6.94
	$\operatorname{BeH}$	$X^2\Sigma^+$	$2061.235^{a}$	-5.15	-10.86	-3.14	-3.30	-7.48	-5.31	-4.56	-11.46
	$\mathrm{BeF}$	$X^2\Sigma^+$	$1247.36^{\mathrm{b}}$	9.50	8.85	8.47	7.42	8.23	8.30	8.61	8.42
	В,	$X^{3}\Sigma_{a}^{-}$	$1051.3^{\mathrm{b}}$	-16.10	-29.18	-1.90	-90.13	-72.72	-61.76	-45.34	-75.26
	ΒÑ	$X^3\Pi^{ m c}$	$1514.6^{\mathrm{b}}$	6.82	-8.55	-7.84	-7.27	-7.70	-7.12	-7.95	-8.55
	BO	$X^2\Sigma^+$	$1885.286^{a}$	-11.61	-21.02	-23.17	-24.32	-21.41	-22.38	-21.98	-21.92
	CH	$X^2 \Pi_r$	$2860.7508^{a}$	-10.43	1.87	-1.66	-1.56	-0.88	-0.74	-0.33	-2.12
	$C_{3}^{-}$	$X^2 \Sigma_a^+$	$1781.189^{a}$	-116.46	-6.89	-8.56	-7.90	-6.57	-4.75	-46.93	-133.20
	0N ON	$X^2\Sigma^+$	$2068.648^{a}$	54.67	-7.85	-10.32	-8.92	-7.36	-7.65	-7.05	-7.37
	$CO^+$	$X^2\Sigma^+$	$2214.127^{a}$	67.00	-16.90	-20.40	-19.77	-13.55	-16.29	-16.28	-16.00
	CF	$X^2 \Pi_r$	$1307.93^{a}$	-3.10	-6.70	-8.59	-10.31	-7.35	-7.36	-6.36	-6.52
	HN	$X^3\Sigma^-$	$3282.72^{a}$	0.19	-1.75	-0.73	-1.66	-1.85	-0.93	-0.68	-0.93
	×°+	$X^2\Sigma_q^+$	$2207.0115^{a}$	70.77	-6.08	-7.18	-6.92	-4.96	-5.51	-4.79	-5.09
	NŐ	$X^2 \Pi_r^j$	$1904.1346^{a}$	177.65	-6.77	-6.31	-7.80	-6.83	-7.09	-6.89	-6.90
	NF	$X^3\Sigma^-$	$1141.37^{a}$	-7.98	-10.86	-13.45	-14.29	-9.97	-10.40	-9.82	-9.56
	HO	$X^2 \Pi_i$	$3737.761^{a}$	-10.03	-9.36	-10.91	-11.21	-9.13	-9.16	-8.56	-10.23
	$^{+}\mathrm{HO}$	$X^3\Sigma^-$	$3113.37^{ m b}$	6.88	4.44	4.52	6.84	6.76	4.34	4.44	4.32
	$O_2$	$X^3 \Sigma_g^-$	$1580.161^{\rm a}$	-4.57	-3.09	-3.11	-3.25	-3.05	-3.23	-3.22	-3.24
	02+	$X^2 \Pi_g$	$1905.892^{a}$	28.13	4.68	4.77	3.87	4.66	4.40	4.54	4.45
	$0_2^-$	$X^2 \Pi_{g,i}$	$1090^{\mathrm{b}}$	24.81	27.20	27.23	27.20	27.12	26.89	26.82	26.85
	OF	$X^2\Pi$	$1053.0138^{a}$	29.52	2.56	0.62	1.23	5.31	4.65	5.21	5.54
	$HF^+$	$X^2 \Pi_i$	$3090.5^{\rm b}$	26.49	27.97	29.45	27.38	28.07	26.44	26.44	26.78
	$F_{2}^{+}$	$X^2 \Pi_{g,i}$	$1104^{d}$	-34.13	20.51	23.81	19.99	19.30	18.63	17.94	17.71
	$\mathbf{F_2}^-$	$X^2 \Sigma_u^+$	$510^{b*}$	-40.09	-64.06	-64.26	-63.25	-64.22	-63.46	-63.22	-63.09
	RMSD			51.77	18.77	18.44	25.78	23.00	21.46	21.76	35.74
	MSD			9.83	-4.96	-4.16	-8.17	-6.36	-6.11	-6.96	-12.26
	MIN			-116.46	-64.06	-64.26	-90.13	-72.72	-63.46	-63.22	-133.20
	MAX			177.65	27.97	29.45	27.38	28.07	26.89	26.82	26.85
Row $2^{-}$	Dimer	State	Expt.	$\Delta(\text{CCSD}(T))$ :	$: \Delta(CCSD(T)):$	$\Delta(CCSD(T))$	$\Delta(CCSD(T))$	$: \Delta(CCSD(T)):$	$\Delta(\text{CCSD}(T))$	$: \Delta(CCSD(T))$	$\Delta(CCSD(T))$ :
Row $3$				UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	NaO	$X^2 \Pi$	$526^{b*}$	-38.60	-38.52	-38.69	-38.45	-38.46	-38.58	-38.44	-38.59
	MgH	$X^2\Sigma^+$	$1492.7763^{a}$	-4.27	-2.49	-4.24	-5.69	-2.78	-0.51	-3.21	-4.82
	MgF	$X^2\Sigma^+$	$711.69^{b}$	0.08	-0.30	-0.66	-1.34	-0.93	-0.78	-0.50	-0.48
	$AIH^+$	$X^2\Sigma^+$	$1620^{\rm b}$	24.02	29.86	21.80	34.31	28.10	28.30	27.77	19.42
	$\operatorname{SiH}$	$X^2 \Pi_r$	$2042.5229^{a}$	-12.95	-9.14	-9.12	-9.46	-7.06	-7.05	-12.62	-12.63
	$\operatorname{SiF}$	$X^2_{ m , II}$	$837.32507^{a}$	10.84	9.61	8.85	7.61	8.75	8.85	9.39	9.29
	НЧ	$X^3 \Sigma^-$	$2363.774^{a}$	1.72	0.54	-1.71	1.43	-0.52	1.70	1.47	1.82
	- HH	$X^{z}\Pi_{r}$	$2299.6^{\mathrm{p}}$	90.81	90.68	88.43	87.87	86.24	88.25	89.26	89.52
	- HJ	$X^{\pm}\Pi_{i}$	2230~	30.09	31.71	33.09	27.12	33.24	24.68	31.71	31.07

Table 5.3: Experimental vibrational frequencies (in  $\rm cm^{-1}$ ) and errors (in  $\rm cm^{-1}$ ) in the corrected vibrational frequencies for the 59 open-shell species are presented in for the CCSD(T) methods utilizing different molecular orbitals. Root

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1.97	-5.33	34.19	-8.48	-7.83	6.39	-3.21	-4.66	-2.36	0.82	2.76	23.78	27.04	-3.02	-7.96	-10.81	-4.87	46.72	24.52	6.66	-38.59	89.52	$\Delta(\text{CCSD}(T))$ :	$\omega B97M-V)$	2.31	-13.33	0.34	-6.91	8.27	-5.71	-7.13	11.07	7.95	-1.39	-13.33	11.07	$\Delta(CCSD(T))$	$\omega B97M-V)$	28.34	-2.13	-133.20	89.52
6.27	-5.24	34.36	-8.38	-7.70	8.57	-3.09	-2.25	-1.69	0.86	2.85	24.97	27.15	-2.98	-7.69	-10.95	-4.87	46.76	24.88	7.47	-38.44	89.26	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	2.24	-13.45	0.32	-6.83	8.25	-5.81	-7.17	11.06	7.97	-1.42	-13.45	11.06	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	22.01	0.40	-63.22	89.26
0.64	-5.73	34.26	-8.78	-8.66	8.58	-3.49	-5.36	-2.67	0.91	2.61	23.34	26.71	-3.10	-9.01	-11.82	-5.15	47.24	24.38	6.87	-38.58	88.25	$\Delta(CCSD(T))$ :	B97)	2.27	-14.91	0.52	-7.19	8.56	-5.64	-7.02	11.25	8.37	-1.52	-14.91	11.25	$\Delta(CCSD(T))$ :	B97)	21.65	0.45	-63.46	88.25
1.77	-5.69	34.26	-8.75	-8.82	8.56	-3.32	-2.47	-2.91	1.28	2.81	24.90	27.11	-3.30	-9.15	-11.79	-5.24	47.73	24.58	7.17	-38.46	86.24	$\Delta(CCSD(T))$ :	B97M-rV)	2.19	-15.34	0.46	-7.34	8.68	-5.73	-7.05	11.32	8.52	-1.60	-15.34	11.32	$\Delta(CCSD(T))$ :	B97M-rV)	22.39	0.48	-72.72	86.24
-4.50	-6.74	33.46	-10.09	-11.04	8.59	-4.42	-10.72	-4.01	0.40	1.81	23.45	27.08	-3.26	-10.92	-13.73	-4.86	47.50	24.95	5.98	-38.45	87.87	$\Delta(CCSD(T))$ :	BLYP)	2.14	-15.43	0.50	-7.63	8.56	-5.94	-7.06	11.14	8.54	-1.71	-15.43	11.14	$\Delta(CCSD(T))$ :	BLYP)	23.77	-0.82	-90.13	87.87
-7.69	-8.11	30.68	-9.11	-9.70	8.60	-4.02	-13.40	-3.59	-2.19	-0.43	23.84	26.71	-3.03	-10.06	-12.85	-4.10	45.39	24.52	5.36	-38.69	88.43	$\Delta(CCSD(T))$ :	OOMP2)	2.27	-16.97	33.72	-7.09	60.95	-5.93	-7.44	10.91	25.99	8.80	-16.97	60.95	$\Delta(CCSD(T))$ :	OOMP2)	22.47	1.95	-64.26	88.43
4.78	-6.44	32.77	-8.36	-7.91	8.60	-2.78	-3.84	-3.17	-1.02	0.55	23.76	27.15	-2.97	-7.55	-10.50	-4.64	45.33	24.88	7.25	-38.52	90.68	$\Delta(CCSD(T))$ :	€-OOMP2)	2.27	-14.00	0.26	-6.88	8.27	-5.46	-7.18	11.06	8.07	-1.46	-14.00	11.06	$\Delta(CCSD(T))$ :	6-00MP2)	20.87	1.10	-64.06	90.68
35.49	-3.47	34.68	-7.52	-5.85	7.92	1.85	30.28	24.33	2.02	21.21	26.52	26.79	-2.93	-6.99	-11.14	-5.50	50.44	27.42	11.85	-38.60	90.81	$\Delta(\text{CCSD}(T))$ : 2	UHF) /	2.24	-8.83	-0.03	-6.92	33.88	19.14	-7.69	44.58	21.48	9.55	-8.83	44.58	$\Delta(\text{CCSD}(T); 2$	UHF) ,	38.69	10.71	-116.46	177.65
$1239.79924^{a}$	$1233.34^{a}$	$1000^{\mathrm{b}}$	$846.75^{a}$	$1053.25^{ m b}$	$2696.2475^{a}$	$1179.91^{a}$	$1384^{\mathrm{b}}$	$1218.7^{ m b}$	$1150.7913^{a}$	$1306.778^{e}$	$2673.69^{a}$	$480^{\mathrm{b}}$	$846.7^{\rm b}$	$876.89749^{a}$	$827.95767^{a}$	$853.64268^{a}$	$870^{b}$					Expt.		$462.12^{\rm b}$	$617.1169^{a}$	$510.98^{a}$	$535.59^{a}$	$672.2^{a}$	$739.1^{b}$	$725.7102^{a}$	4062										
$X^2\Sigma^+$	$X^{4}\Pi_{r}$	$X^3\Sigma^-$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^{2}\Pi$	$X^3\Sigma^-$	$X^2 \Pi_i$	$X^{2}\Pi$					State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^3 \Sigma_g^-$	$X^2\Pi_r^-$	$X^2 \Pi_u$	$X^2 \Pi_r$	$X^3 \Sigma_g^-$	$X^2 \Pi_{g,r}$										
CP	Ю	$PO^{-}$	ΡF	$PF^+$	HS	BS	$CS^+$	NS	SO	$^{+}Os$	$HCl^+$	LiCl <sup>-</sup>	BeCl	CCI	NCI	OCI	$CIF^+$	RMSD	MSD	MIN	MAX	Dimer		MgCl	AIS	$Si_2$	SiCI	$P_{2}^{+}$	PS	$^{ m S}_2$	$^{\mathrm{S}^+}_{\mathrm{S}^+}$	RMSD	MSD	MIN	MAX			RMSD	MSD	MIN	MAX
																						Row	3-Row 3													Open-	$_{ m shell}$				

<sup>a</sup> From Ref. 349. <sup>b</sup> From Ref. 346. <sup>d</sup> From Ref. 351. <sup>e</sup> From Ref. 352. \* Theoretical results.

Tab. 5.3 presents the experimental vibrational frequencies, in cm<sup>-1</sup>, and errors in the corrected vibrational frequencies, in cm<sup>-1</sup>, for the 59 species in the open-shell subset. Turning to Fig. 5.1, the open-shell non-HF methods exhibit an increase in the number of outliers as compared to the closed-shell cases. Overall, CCSD(T) with  $\kappa$ -OOMP2 orbitals provides the best RMSD of all references (20.87 cm<sup>-1</sup>), improving on the performance of CCSD(T) with UHF orbitals by nearly a factor of two (38.69 cm<sup>-1</sup>). The use of OOMP2, BLYP, B97M-rV, B97, and  $\omega$ B97X-V orbitals yields comparable performance to  $\kappa$ -OOMP2 orbitals (22-24 cm<sup>-1</sup>), while the overall performance of CCSD(T) with  $\omega$ B97M-V orbitals falls between the other non-UHF references and the UHF reference (28.34 cm<sup>-1</sup>). In terms of MSDs, CCSD(T) with UHF orbitals is seen to blue shift the open-shell frequencies by 11 cm<sup>-1</sup> while the non-UHF methods yield little-to-no systemic shift in frequencies ( $\pm 2$  cm<sup>-1</sup>). The performance of CCSD(T) with  $\kappa$ -OOMP2 or OOMP2 or OOMP2 orbitals on the open-shell systems is a factor of two worse than for the closed-shell systems.

Looking at individual cases, errors in corrected vibrational frequencies for the triplet ground state of B<sub>2</sub> are seen to range from -1.90 cm<sup>-1</sup> using OOMP2 orbitals to -90.13 cm<sup>-1</sup> using BLYP orbitals. B<sub>2</sub> is another system which is know to exhibit MR character and therefore the varied and often poor performance of these single reference methods is to be expected[272]. Evidence of MR character is seen in the mean-field  $\langle S^2 \rangle$  values, where the  $\kappa$ -OOMP2 and UHF orbitals both significantly break spin-symmetry ( $\langle S^2 \rangle$  of 2.81206 and 2.90778, respectively; Tab. D.8). The DFT orbital references are also seen to significantly break spin-symmetry while the OOMP2 reference artificially restores spin-symmetry for this system.

In agreement with the work of Beran et al. [313] and Tentscher and Arey [316] the predictions of the frequencies of CN, NO, OF, and their isoelectronic and isovalent counterparts with CCSD(T) using a UHF reference yield sizeable errors. CN is isoelectronic to  $C_2^-$ ,  $CO^+$ , and  $N_2^+$  and is isovalent to CP,  $CS^+$ , and  $P_2^+$ . CCSD(T) with UHF orbitals yields errors in the corrected frequencies of these species of 54.67 cm<sup>-1</sup>, -116.46 cm<sup>-1</sup>, 67.00 cm<sup>-1</sup>, 70.77 cm<sup>-1</sup>, 35.49 cm<sup>-1</sup>, 30.28 cm<sup>-1</sup>, and 33.88 cm<sup>-1</sup>, respectively, while CCSD(T) with  $\kappa$ -OOMP2 orbitals improves these errors to  $-7.85 \text{ cm}^{-1}$ ,  $-6.89 \text{ cm}^{-1}$ ,  $-16.90 \text{ cm}^{-1}$ ,  $-6.08 \text{ cm}^{-1}$ ,  $4.78 \text{ cm}^{-1}$ ,  $-3.84 \text{ cm}^{-1}$ , and  $8.27 \text{ cm}^{-1}$ , respectively. In all of these cases but  $C_2^{-1}$  these errors with UHF orbital references are accompanied by spin-symmetry breaking at the level of the reference with respective  $\langle S^2 \rangle$  values of 1.15755, 0.75627, 0.96970, 1.23889, 1.61297, 1.47259, 1.16866 for UHF and 0.76257, 0.75579, 0.76730, 0.75334, 0.82118, 0.80770, 0.75476 for the  $\kappa$ -OOMP2 reference (Tab. D.8). The spin-symmetry restoration from the  $\kappa$ -OOMP2 (or OOMP2/DFT) orbitals is seen to dramatically improve the predicted frequencies for these systems with significant symmetry-breaking occurring at the UHF level. For  $C_2^{-}$ , the error in predicted frequency for CCSD(T) with UHF orbitals has the opposite sign of the other frequency errors in this isoelectronic/isovalent family of dimers and the UHF  $\langle S^2 \rangle$  suggests little spin-contamination. Instead, UHF,  $\omega$ B97X-V, and  $\omega$ B97M-V are seen to favor brokenspatial-symmetry solutions for  $C_2^{-}$ , contributing to large errors in the predicted frequencies while  $\kappa$ -OOMP2, OOMP2, and the other density functionals preserve the spatial symmetry and yield much more reliable frequencies. For  $P_2^+$ , CCSD(T):OOMP2 yields an error of

60.95 cm<sup>-1</sup>, almost twice that of CCSD(T):UHF. The orbital optimization of the ground state at the OOMP2 level is shown to give preference to a higher symmetry orbital occupation where the  $\pi_{3p_x}$  and  $\pi_{3p_y}$  MOs are doubly-occupied and the  $\sigma_{3p_z}$  orbital is singly-occupied. The ground state for all other methods doubly occupies the  $\sigma_{3p_z}$  and singly-occupies one of the two  $\pi_{3p}$  orbitals, breaking the  $D_{\infty h}$  symmetry of the molecule. The latter occupation, however, is the filling predicted by MO theory and yields reasonable frequencies in comparison to the experimental benchmark. This represents another example of essential symmetry breaking that is quelled by OOMP2.

NO is isoelectronic to  $O_2^+$  and isovalent to PO, NS, SO<sup>+</sup>, PS, and  $S_2^+$ . CCSD(T) with UHF orbitals yields errors in the corrected frequencies of these species of  $177.65 \text{ cm}^{-1}$ , 28.13  $cm^{-1}$ , -3.47  $cm^{-1}$ , 24.33  $cm^{-1}$ , 21.21  $cm^{-1}$ , 19.14  $cm^{-1}$ , and 44.58  $cm^{-1}$ , respectively, while CCSD(T) with  $\kappa$ -OOMP2 orbitals yields error of -6.77 cm<sup>-1</sup>, 4.68 cm<sup>-1</sup>, -6.44 cm<sup>-1</sup>, -3.17  $cm^{-1}$ , 0.55  $cm^{-1}$ , -5.46  $cm^{-1}$ , and 11.06  $cm^{-1}$ , respectively. NO represents the most positive outlier for CCSD(T) with UHF orbitals while the error for PO is the smallest of all the orbital references considered. Comparing the mean-field  $\langle S^2 \rangle$  values between UHF and  $\kappa$ -OOMP2 for these species,  $O_2^+$  (1.12597 vs. 0.75272), NS (1.18951 vs. 0.75758), SO<sup>+</sup> (1.18709 vs. 0.75583), PS (1.05380 vs. 0.76119), and S<sub>2</sub><sup>+</sup> (1.23007 vs 0.75862) show significant spincontamination at the UHF level while cases exhibiting the largest and smallest errors, NO (0.79621 vs. 0.75407) and PO (0.77298 vs. 0.75758), respectively, show little evidence of significant spin-contamination (Tab. D.8). These findings are consistent with those of Szalay et al.[314], which find that in these errors arise from instabilities in the doublet wavefuction. This discrepancy between NO and PO arises due to the instability in the PO UHF wavefunction occurring at a larger internuclear separation than the equilibrium bond length around which data was collected.

OF is isoelectronic to  $F_2^+$  and isovalent to OCl and ClF<sup>+</sup>. CCSD(T) with UHF orbitals yields errors in the corrected frequencies of these species of 29.52 cm<sup>-1</sup>, -21.63 cm<sup>-1</sup>, -5.50 cm<sup>-1</sup>, and 50.44 cm<sup>-1</sup>, respectively, while CCSD(T) with  $\kappa$ -OOMP2 orbitals yields error of 2.56 cm<sup>-1</sup>, 33.01 cm<sup>-1</sup>, -4.64 cm<sup>-1</sup>, and 45.33 cm<sup>-1</sup>, respectively. The errors for UHF and non-UHF orbitals for  $F_2^+$  are seen to differ by over 50 cm<sup>-1</sup> while for OCl and ClF the errors in the frequencies are in good agreement. Similarly to the cases of NO and PO, the UHF wavefunctions for OF, OCl, and ClF<sup>+</sup> do not demonstrate signs of significant spin contamination ( $\langle S^2 \rangle$  values of 0.77257, 0.77010, 0.76463, respectively; Tab. D.8).

#### Potential issues with experimental data

Table 5.4: Experimental vibrational frequencies (in  $\text{cm}^{-1}$ ), mean errors for the non-HF CCSD(T) methods (in  $\text{cm}^{-1}$ ), ranges of the errors for the non-HF CCSD(T) methods (in  $\text{cm}^{-1}$ ), and alternative experimental reported frequencies (in  $\text{cm}^{-1}$ ) are presented for species where the experimental results are in question.

Dimer	State	Current Expt.	Mean Non-HF	Non-HF Error	Alternative
		Frequency	Error	Range	Expt.
					Frequency

$O_2^{-}$	$X^2 \Pi_{g,i}$	1090[346]	27.04	0.41	1090[353, 354]
					$1108 \pm 20[355]$
					1140[356]
					1145[357]
$\rm HF^+$	$X^2 \Pi_i$	3090.5[346, 358,	27.51	3.01	3061.8[360]
		359]			
		-			$3118[313]^*$
					3119[316]*
$F_{a}^{+}$	$X^2 \prod_{\alpha i}$	1104[351]	19.70	6.09	1091.5[361]
F -	$Y^2\Sigma^+$	510[346_362]*	63.65	1 17	100110[001]
1 <sup>2</sup> 2	$\Lambda \  \  \  \  \  \  \  \  \  \  \  \  \ $	510[540, 502]	-03.05	1.17	* o ([o oo]*
NaO	X 211	526[346]	-38.53	0.25	504[363]
					547[364]
$AlH^+$	$X^2\Sigma^+$	1620[346, 365]	27.08	14.89	
$PH^+$	$X^2 \Pi_r$	2299.6[346, 366]	88.61	4.44	$2382.75[367]^{*}$
$\rm PH^-$	$X^2 \Pi_i$	2230[346]	30.37	8.56	$2230 \pm 100[368]$
$PO^{-}$	$X^3\Sigma^-$	1000[346]	33.43	3.68	$1000 \pm 70[368]$
$HCl^+$	$X^2 \Pi_i$	2673.69[349]	24.00	1.63	$2702.6[369]^*$
$LiCl^{-}$	$X^2\Sigma^+$	480[346]	26.99	0.44	$480 \pm 80[370]$
$ClF^+$	$X^2\Pi$	870[346]	46.67	2.40	$870 \pm 30[371]$
		[]		-	$912 \pm 30[372]$
					[ ]

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<sup>\*</sup> Theoretical results.

In the cases of  $ClF^+$  and  $PH^+$  the predicted frequencies for CCSD(T) with a UHF orbital reference and CCSD(T) with non-UHF orbital references differ significantly (more than 25)  $cm^{-1}$ ) from the reported benchmark value while the CCSD(T) with non-UHF orbitals all vield predicted frequencies in agreement with each other. This observation is true of other molecules and ions in the open-shell set as well, as summarized in Tab. 5.4. In all of these cases except  $F_2^+$  and  $F_2^-$  the errors for CCSD(T) with UHF orbitals are consistent with the errors for CCSD(T) with non-HF orbitals. The agreement of all methods suggests that potentially inaccurate or imprecise experimental reference values should be revisited. For NaO[363, 364] and  $F_2^{-}$ [362], the reference data given by Huber and Herzberg[346] are sourced from calculations performed at the HF level of theory; any of the CCSD(T) methods surveyed should be seen as a more accurate result for these systems. For another subset of these systems the error bars on the experimental values contain the frequencies calculated from this work (PH<sup>-</sup>: 2230 ± 100 cm<sup>-1</sup>[368], PO<sup>-</sup>: 1000 ± 70 cm<sup>-1</sup>[368], LiCl<sup>-</sup>: 480 ± 80 cm<sup>-1</sup>[370]). For ClF<sup>+</sup>, Huber and Herzberg[346] reference the DeKock et al.[371] who report a vibrational frequency of  $870 \pm 30 \text{ cm}^{-1}$  while Anderson et al. [372] report a frequency of 912  $\pm$  30 cm<sup>-1</sup>. Our calculations, which predict a ClF<sup>+</sup> frequency of 917 cm<sup>-1</sup>, more closely agree with the work of Anderson et al. Similarly, for  $O_2^-$  Huber and Herzberg[346] cite a value of  $1090 \text{ cm}^{-1}$  based on the works of Boness and Schulz[353] and Linder and Schmidt[354] while also noting that Gray et al. [356] and Creighton and Lippincott[357] give values of 1140  $\rm cm^{-1}$  and 1145  $\rm cm^{-1}$ , respectively. In a more recent study, Ervin et al.[355] give a value of  $1108 \pm 20$  cm<sup>-1</sup>, which is in good agreement with our theoretical predictions of  $1117 \text{ cm}^{-1}$ . The most positive error point for many of the CCSD(T) method with non-UHF references, PH<sup>+</sup>, does not have much experimental data on its spectroscopic constants in the literature, with the value of 2299.6  $\text{cm}^{-1}$  tracing back to a study by Narasimham[366]. A recent modeling study by Reddy et al. [367] predicts a ground-state frequency of 2382.75  $\rm cm^{-1}$ , in good agreement with our calculated values of 2386–2390 cm<sup>-1</sup>, suggesting that
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further experimental study of this system is worthwhile. The experimental reference for  $F_2^+$  was taken from Cormack et al.[351] while Yang et al.[361], consistent with Tentscher and Arey[316], suggest a value of 1091.5 cm<sup>-1</sup>, further from our theoretical predictions. For HF<sup>+</sup>, the large-basis results from Beran et al.[313] and Tentscher and Arey[316] predict an error that agree with our error prediction of 28 cm<sup>-1</sup>. These calculated frequencies are closer to the value originally reported by Gewurtz et al.[358] and Hovde et al.[359] of 3090.5 cm<sup>-1</sup> adopted by Huber and Herzberg[346] than to the value of 3061.8 cm<sup>-1</sup> proposed by Yencha et al.[360]. Similarly, our predictions of the corrected vibrational frequency of HCl<sup>+</sup> are blue-shifted by approximately 24 cm<sup>-1</sup> compared to the experimental reference value of 2673.69 cm<sup>-1</sup> from Irikura[349]. A recent joint experimental and theoretical study of this system by Patanen et al.[369] has suggested a computed value of 2702.6 cm<sup>-1</sup>, in much better agreement with our results. For AlH<sup>+</sup>, Huber and Herzberg[346] cite a 1934 study by Hoslt[365] to approximate the vibrational frequency, though a more precise value is desired for the point of comparison to our computed frequencies to better assess the error between the different reference methods.

#### Pruned subset

In order to draw more meaningful conclusions about the performance of CCSD(T) with  $\kappa$ -OOMP2 orbitals, we consider a subset of data points where the MR species (C<sub>2</sub> and B<sub>2</sub>) and the species with ambiguous experimental values (Tab. 5.4) discussed above are excluded, leaving 35 closed-shell species (28 neutrals, 6 cations, 1 anion) and 46 open-shell species (36 neutrals, 9 cations, 1 anion; 35 doublets, 11 triplets). Tab. 5.5 presents the RMSDs, MSDs, MINs, and MAXs for the closed-shell species, open-shell species, and the overall pruned set. These data are presented graphically in Fig. 5.2. For the pruned data set CCSD(T):B97 and CCSD(T): $\kappa$ -OOMP2 are seen to yield the best performance with RMSDs of 8.48 cm<sup>-1</sup> and 8.50 cm<sup>-1</sup>, respectively. The performances of CCSD(T): $\omega$ B97X-V and CCSD(T): $\omega$ B97M-V are hindered by C<sub>2</sub><sup>-</sup>; excluding this point brings the RMSDs for these methods to the same level as the other DFT-based methods.

Comparing Fig. 5.1 and Fig. 5.2, we see that the pruning procedure removed many of the outlier cases that were shared between all or nearly all of the methods surveyed. For CCSD(T):UHF, the remaining outlying points are the spin-contaminated points from the closed-shell subset (F<sub>2</sub>, PN, BeS, NS<sup>+</sup>, and P<sub>2</sub>), and HCl, while the remaining open-shell outliers are NO, N<sub>2</sub><sup>+</sup>, CO<sup>+</sup>, and C<sub>2</sub><sup>-</sup>. The range covered by the CCSD(T):UHF whiskers is larger by a factor of two than nearly all the other methods tested. For CCSD(T): $\kappa$ -OOMP2, the pruning procedure leaves only one closed-shell and one open-shell outlier, HCl and BO, respectively. For CCSD(T):OOMP2, the remaining outliers in the closed-shell subset are HCl and BeO while the open-shell set has outliers P<sub>2</sub><sup>+</sup>, which has been previously discussed, Si<sub>2</sub>, which breaks spatial symmetry at the OOMP2 level, and BO. CCSD(T):BLYP has closedshell outliers in NaLi, NaH, and HCl and an open-shell outlier in BO. CCSD(T):B97M-rV and CCSD(T): $\omega$ B97X-V has closed-shell outliers in HCl and NaH and open-shell outliers in C<sub>2</sub><sup>-</sup>

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Figure 5.2: Box plots (overall, left, and enhanced, right) of the errors in corrected vibrational frequencies (in  $\text{cm}^{-1}$ ) are presented for the pruned subset of species. Red lines mark the median errors, boxes bound the central 50% of the data, whiskers enclose all data points within 1.5 times the inter-quartile range of the box edges, and points denote outlying data.



Table 5.5: Root mean square deviations, mean signed deviations, most negative deviations, and most positive deviations (all in  $\text{cm}^{-1}$ ) for the cropped subset of species are summarized for the CCSD(T) methods utilizing different molecular orbitals.

Closed-	$\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(\text{T}))$ :	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(T)$	$\Delta(CCSD(T))$	$\Delta(\text{CCSD}(T))$	$\Delta(\text{CCSD}(T))$ :
shell	UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega$ B97M-V)
pruned								
RMSD	44.43	9.41	10.18	11.56	10.05	9.45	9.50	8.97
MSD	-7.74	-6.96	-7.73	-8.10	-7.32	-6.41	-6.41	-5.93
MIN	-240.71	-18.53	-21.83	-29.30	-21.21	-20.00	-18.16	-18.40
MAX	60.05	14.11	15.27	16.82	16.86	15.88	16.93	16.10
Open-	$\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(\text{T}))$ :	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(T)$	: $\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(T))$	$\Delta(\text{CCSD}(T))$ :
shell	UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
pruned								
RMSD	38.33	7.75	13.45	8.80	7.61	7.67	10.26	21.11
MSD	9.27	-3.48	-2.83	-4.61	-3.32	-3.35	-3.97	-6.34
MIN	-116.46	-21.02	-23.17	-24.32	-21.41	-22.38	-46.93	-133.20
MAX	177.65	11.06	60.95	11.14	11.32	11.25	11.06	11.07
Total	$\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(\text{T}))$ :	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(T)$	: $\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(T))$	$\Delta(\text{CCSD}(T))$ :
pruned	UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
RMSD	41.07	8.50	12.15	10.08	8.75	8.48	9.94	16.96
MSD	1.92	-4.98	-4.95	-6.11	-5.04	-4.67	-5.03	-6.17
MIN	-240.71	-21.02	-23.17	-29.30	-21.41	-22.38	-46.93	-133.20
MAX	177.65	14.11	60.95	16.82	16.86	15.88	16.93	16.10

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and BO.  $CCSD(T):\omega B97M-V$  also has HCl and NaH as closed-shell outliers and  $C_2^-$  as an open-shell outlier. The non-HF methods are seen to have their mean values red-shifted compared to experiments by 5-6 cm<sup>-1</sup>.

### 5.6 Conclusions

We have evaluated the performance of CCSD(T) with different MO references to predict the vibrational frequencies of both closed-shell and open-shell diatomic molecules and ions. The at times problematic use of a UHF reference was compared against the use of  $\kappa$ -OOMP2, OOMP2, BLYP, B97M-rV, B97,  $\omega$ B97X-V, and  $\omega$ B97M-V molecular orbital references. Overall, CCSD(T):B97, CCSD(T): $\kappa$ -OOMP2, and CCSD(T):B97M-rV yield RMSDs on the pruned overall data set of 8.48 cm<sup>-1</sup>, 8.50 cm<sup>-1</sup>, and 8.75 cm<sup>-1</sup>, respectively, reducing the RMSD for CCSD(T):UHF by nearly a factor of 5. For the pruned closed- and open-shell subsets the associated RMSDs are 9.45 cm<sup>-1</sup> and 7.67 cm<sup>-1</sup>, respectively, 9.41 cm<sup>-1</sup> and 7.75 cm<sup>-1</sup>, respectively, and 10.05 and 7.61, respectively, for CCSD(T):B97, CCSD(T): $\kappa$ -OOMP2, and CCSD(T):B97M-rV. The slightly degraded performance of the other non-HF MO references are skewed by one or two data points and otherwise rank competitively with the B97,  $\kappa$ -OOMP2, and B97M-rV. These outlying data points are seen to arise from spatial or spin-symmetry breaking or erroneous symmetry restoration. The effect of regularization in  $\kappa$ -OOMP2 is seen to prevent the symmetry issues seen in OOMP2.

A major practical limitation of these very promising non-HF CCSD(T) methods is the present lack of implemented analytic gradients, relegating the current application of these to systems with only a few atoms. One may develop an approach like that of Taube and Bartlett's FNO-CCSD(T)[373] for  $\kappa$ -OOMP2 and/or DFT orbitals. The success of the non-HF CCSD(T) methods in treating closed-shell and open-shell systems with the same accuracy speaks to the usefulness of such an implementation. Such approaches do not affect the overall asymptotic scaling and would extend the "black box" utility of CCSD(T) currently seen for closed-shell systems. Furthermore, the use of  $\kappa$ -OOMP2 as the generator of MOs also provides a diagnostic tool in the mean-field  $\langle S^2 \rangle$  of the multireference character of the target system, informing the expected accuracy of the subsequent CCSD(T) results.

In order to assess errors relating to the computational treatment presented in this work we consider errors in the basis set and errors related to the approximate treatment of triple excitations. Peterson and Dunning[342] demonstrated a blue shift of computed harmonic frequencies for row 3 diatomic species of 2-15 cm<sup>-1</sup> when going from the cc-pwCVTZ to cc-pwCV5Z. Provided this trend holds with the aug-cc-pwCVnZ basis sets the use of the aug-cc-pwCVQZ basis should help to correct the systematic red-shift seen for the non-HF CCSD(T) methods. For exact treatment of triples Tentscher and Arey[316] found CCSDT tends to red-shift in computed frequencies compared to CCSD(T). This effect is exacerbated for species where CCSD(T) with UHF orbitals is particularly poor (N<sub>2</sub><sup>+</sup>, CN) while for other other species the shift was on the order of 5 cm<sup>1</sup>. We also note that CCSDT, free of the perturbative nature of CCSD(T), should be less sensitive to the choice of reference orbitals. Chapter 5. Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction 93

Further extension of this study approach to cases where DFT traditionally struggles would be of significant interest. In particular, extension to transition metals species and systems with significant charge separation is highly desirable to discriminate if CCSD(T):DFT is able to remedy the traditional failures of DFT. Additionally, the exponential regularization schemes of Lee and Head-Gordon[39] could be extended to the perturbative triples calculation in CCSD(T) as an attempt to handle cases of nonvariational failure in traditional CCSD(T).

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## Appendix A

# Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

### A.1 Branching Ratio Calculations

We performed RRKM calculations [126] via direct counting of vibrational states [127] to determine the unimolecular reaction rate constants for the silicon-methylacetylene, silicon-allene, and silicon–D3-methylacetylene systems. We also applied a tunneling correction to the reaction rates by approximating the barriers as Eckart potentials [128]. These rate constants were used to form the coupled systems of equations for each network, which was then solved and evaluated in the longtime limit to extract branching ratios. For the silicon–methylacetylene system, the trimmed reaction network of reactions considered in the RRKM calculations is given in Fig. S3. For this system our network began with a unit population in isomer [i1] and no population in any other isomers, corresponding to the entrance channel for this reaction. Using the tunneling-corrected rates, we predicted a branching ratio of 38.2% **p1**, 0.6% **p2**, 11.1% **p3**, and 45.8% **p4**. For the silicon–allene system, the same reaction network and rates were used as in the silicon-methylacetylene case. For this system, however, our network began with unit population in isomer [i7]. While we observed that the silicon-allene system has entrance channels leading to both isomers [i7] and [i10], we note that all flux out of isomer [i7] into the reaction network must proceed through [i10] according to our network. Using the tunneling-corrected rates, we predicted a branching ratio of 38.7% p1, 0.5% p2, 11.3% **p3**, and 46.5% **p4**.

Comparison of the silicon–methylacetylene and silicon–allene branching ratios shows similar branching ratios for each product. While the reaction rates show the fastest isomerization from [i1] to be  $[i1]\rightarrow[i4]$ , exit from this channel is stifled by the high barrier to form product

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**p2** going through intermediate [i13]. This allows for the other channels out of [i1] to be expressed, particularly [i1] $\rightarrow$ [i5] and [i1] $\rightarrow$ [i6]. Both of these channels can eventually lead to isomers [i6], [i9], and [i12], from which we see much lower overall barriers to for products **p4**, **p1**, and **p3**, respectively. Starting from isomer [i7], we predict fast isomerization to [i10], which then preferentially isomerizes to [i6] over [i2]. From [i6] we have facile isomerization to isomers [i9] and [i12] via ring opening. These computational results corroborate the experimental results for these systems, where the product translational energy distributions of the silicon– methylacetylene and silicon–allene systems are nearly identical.

We also examined the branching ratios for hydrogen atom and methyl loss in the silicon-methylacetylene and silicon-allene systems. These channels were found to contribute negligibly (4.2% for methylacetylene, 2.9% for allene) to the overall product stream, corroborating the experimental findings. Specifically, when considering the silicon-methylacetylene and silicon-allene networks, we included exits to atomic hydrogen loss products **n1**, **n2**, **n3**, **n4**, and **n6** and methyl loss product **m2**. For the silicon-methylacetylene system, we predicted branching ratio of 3.9% **n1**, 0.06% **n2**, 0.02% **n3**, 0.01% **n4**, 0.01% **n6**, and 0.2% **m2**. For the silicon-allene system, we predicted branching ratio of 2.7% **n1**, 0.03% **n2**, 0.02% **n3**, 0.01% **n4**, 0.004% **n6**, and 0.1% **m2**. These lower branching ratios are due to the higher barriers for hydrogen atom or methyl loss when compared to the more favorable molecular hydrogen loss channels.

For the silicon-D3-methylacetylene system, we begin from the same reaction network as in the non-deuterated case. From here, we expanded the reaction network to include all possible constitutional isomers for each corresponding non-deuterated isomer to account for the potential of hydrogen-deuterium scrambling. Hydrogen-deuterium scrambling reaction can occur for the  $[i2] \leftrightarrow [i3]$ ,  $[i2] \leftrightarrow [i10]$ ,  $[i6] \leftrightarrow [i10]$ , and  $[i6] \rightarrow [i9] \rightarrow [i8] \rightarrow [i12] \rightarrow [i6]$  reactions. We further included the  $[i2] \rightarrow [i2]$  and  $[i2] \rightarrow [i3]$  reactions in our reaction network. Using the tunneling-corrected rates, we predicted a branching ratio of 21.7% **p1a** + HD, 11.0% **p1b** + D<sub>2</sub>, 0.3% **p2a** + HD, 0.01% **p2b** + D<sub>2</sub>, 0.01% **p2c** + D<sub>2</sub>, 2.6% **p3a** + D<sub>2</sub>, 1.8% **p3b** + D<sub>2</sub>, 5.2% **p3c** + HD, 27.7% **p4a** + D<sub>2</sub>, and 29.5% **p4b** + HD. Therefore, we predicted a total branching ratio of 56.8% HD loss products and 43.2% D<sub>2</sub> loss products. These values are in good agreement with the experimental findings of 60 ± 15% HD loss and 40 ± 15% D<sub>2</sub> loss. Furthermore, this corroborates the non-deuterated results, which utilize the same general reaction network framework. Reaction rate constants, imaginary transition state frequencies, and reaction path degeneracies are reported in Table A.9.

### A.2 Tables

# Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC3Molecules in the Circumstellar Medium of Carbon Stars124

Table A.1: Reactants, products, intermediates, and transition state structures relevant to the reaction of electronically excited atomic silicon with methylacetylene and allene. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

Reactants									
CH <sub>3</sub> CCH				$CH_2CCH_2$					
	0.0	$0 \text{ kJ mol}^{-1}$		$3.9 \text{ kJ mol}^{-1}$					
		$C_{3v}$ -1 $A_1$				$D_{2d}$ – $^1A_1$			
C	-1.919449	1.519750	0.000000	C	-7.512217	1.836332	0.000036		
C	-0.721221	1.524736	0.000000	C	-6.210183	1.823559	0.000153		
C	-3.382029	1.513693	0.000000	C	-4.908232	1.810798	0.000267		
H	0.343441	1.529181	0.000000	H	-8.068002	1.984399	-0.919492		
H	-3.768514	1.986640	-0.905294	H	-8.070971	1.698436	0.919347		
H	-3.768812	2.058826	0.863618	H	-4.360016	0.885485	-0.140011		
H	-3.762327	0.490833	0.041676	H	-4.342009	2.725331	0.139709		
		Si							
<sup>1</sup> D									
Si	0.000000	0.000000	0.000000	1					

Appendix A.	Additional Information:	Modeling of Gas	Phase Formation	of $c$ -SiC <sub>3</sub>	
Molecules in th	e Circumstellar Mediun	n of Carbon Stars			12

Products									
H <sub>2</sub>					p1				
					-189	$9.6 \text{ kJ mol}^{-1}$			
		$D_{\infty h}^{-1}\Sigma$				$C_{2v}^{-1}A_1$			
H	0.003216	0.000000	0.000000	Si	-2.667415	2.191276	0.000000		
H	0.746784	0.000000	0.000000	C	-0.985018	2.220382	0.000000		
				C	0.298405	2.242563	0.000000		
					1.607796	2.265203	0.000000		
				H	2.142021	3.207971	0.000000		
				H	2.174299	1.341470	0.000000		
		$\mathbf{p2}$				p3			
-173.9 kJ mol <sup>-1</sup>				-163.4 kJ mol <sup>-1</sup>					
$C_s^{-1}A'$						$C_s - A'$			
Si	-3.530077	0.607099	0.333008	Si	-1.440031	2.135965	1.571949		
	-2.115964	1.751749	-0.167586	C	-1.258585	1.493357	-0.509032		
C	-1.762971	0.449585	-0.634584	C	-1.918764	2.648067	0.002796		
C	-2.980050	2.474570	0.443216	C	-0.632564	0.456668	-0.644013		
H	-0.954113	-0.065250	-0.122807	H	-2.429665	3.372178	-0.610959		
H	-1.792019	0.273905	-1.706863	H	-0.091486	-0.444116	-0.815669		
		p4		p5					
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	-159	$9.3 \text{ kJ mol}^{-1}$			-138	$8.6 \text{ kJ mol}^{-1}$			
		$C_2$ -1 $A_1$				$C_{2v}$ - <sup>1</sup> $A_1$			
Si	0.015583	1.421694	0.441321	Si	-0.153925	2.791130	0.597108		
C	-1.710909	1.657474	1.228824	C	-0.701267	1.901637	-0.854814		
C	-1.833196	1.452865	-0.080188	C	-1.470675	0.643621	-0.899626		
C	-1.051473	1.232997	-1.134185	C	-1.079103	1.277031	0.374194		
H	-2.213088	1.188068	2.060067	H	-0.958764	-0.265564	-1.204780		
H	-1.030405	1.710504	-2.101157	H	-2.515112	0.686694	-1.198739		
		p6				$\mathbf{p7}$			
					0-				
	-99	$.8 \text{ kJ mol}^{-1}$			-83	$.5 \text{ kJ mol}^{-1}$			
	0.000000	$C_1^{-1}A$	0.0000000		0.000000	$C_{2v}^{-1}A_1$	0.000000		
Si	0.000000	0.000000	0.000000	Si	0.000000	0.000000	0.000000		
	C 0.000000 0.000000 1.806628				0.000000	0.000000	1.719212		
	C -1.286808 0.627089 1.815653				0.639430	0.000000	3.039625		
	C -0.765193 1.502612 0.924959				-0.652963	-0.004736	3.033189		
	0.697327	0.000000	2.637474		1.582650	0.005577	3.557070		
H	-2.217623	0.494213	2.353974	H	-1.601967	-0.006606	3.539526		

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		p8		p9				
$-70.6 \text{ kJ mol}^{-1}$					-62	$.1 \text{ kJ mol}^{-1}$		
$C_s^{-1}A'$						$C_s$ – $^1A'$		
Si	0.000000	0.000000	0.000000	Si	0.000000	0.000000	0.000000	
C	0.000000	0.000000	1.843287	C	1.937037	0.000000	-0.025559	
C	-1.401563	-0.160952	2.078228	C	2.673224	1.126057	-0.140369	
C	-0.973371	-1.112994	1.101537	C	1.353436	1.386523	-0.023475	
H	-1.159475	-2.178800	1.178760	H	0.000000	0.000000	1.517171	
H	0.745816	0.000000	2.630932	H	3.661304	1.544024	-0.241081	
		p10				p11		
	-6.	$\frac{0 \text{ kJ mol}^{-1}}{2}$			26.	$.5 \text{ kJ mol}^{-1}$		
		$C_1^{-1}A$				$C_{2v}^{-1}A_1$		
Si	0.000000	0.000000	0.000000	Si	0.000000	0.000000	0.000000	
	1.834756	0.000000	-0.012712		-1.406915	-0.732528	-0.876099	
	3.087703	0.481614	-0.103968		-2.436321	0.000674	-1.517362	
	2.999800	-0.889315	-0.084047		-1.406701	0.733908	-0.875881	
H	0.000000	0.000000	1.519674	H H	0.000000	0.000000	1.472112	
H	3.675556	1.385395	-0.152491	H	1.321012	0.000000	-0.649666	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

		p12		p13				
	58.	$.7 \text{ kJ mol}^{-1}$			82.	$7 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_{2v}$ -1 $A_1$		
Si	0.000000	0.000000	0.000000	Si	0.000000	0.000000	0.000000	
C	1.906751	0.000000	-0.088904	C	-1.604147	-0.000599	-0.965813	
C	1.719817	-1.323251	-0.158049	C	-1.321163	1.286223	-0.796186	
C	0.828888	-2.195636	-0.498761	C	-1.319392	-1.286995	-0.795253	
H	0.000000	0.000000	1.522410	H	0.000000	0.000000	1.471313	
H	2.826815	0.569597	-0.151117	H	1.299958	0.000000	-0.689114	
		p14				p15		
	127	$\frac{7.8 \text{ kJ mol}^{-1}}{\Omega}$			147	$\frac{.4 \text{ kJ mol}^{-1}}{\Omega}$		
	0.000000	$\frac{\mathrm{U}_{\mathrm{s}}^{-1}\mathrm{A}'}{0.0000000000000000000000000000000000$	0.000000	<u>a.</u>	0.000000	$\frac{U_{2v}^{-1}A_1}{O_{2v}^{-1}O_{2v}^{-1}O_{2v}^{-1}A_1}$	0.000000	
	0.000000	0.000000	0.000000		0.000000	0.000000	0.000000	
	C = 1.264017 = 0.0000000 = -1.312498				-1.482108	0.004467	-0.835093	
	0.030420	-1.228041	-1.404983		-2.002384	0.000108	-1.40(455)	
	-0.859350	-0.920370	-1.31317U		-3.121018	0.003497	-2.101804	
	0.000000	0.000000	1.4/0/80			0.000000	1.409804	
H	0.950336	-2.176678	-1.791926	H	1.257655	0.000000	-0.760691	

Intermediates										
		[i1]		[i2]						
	-381	$1.9 \text{ kJ mol}^{-1}$			-366	$\frac{5.5 \text{ kJ mol}^{-1}}{2}$				
		$C_s^{-1}A'$				$C_1^{-1}A$				
Si	-1.656605	1.891479	-0.619303	Si	1.217161	2.092109	-0.478041			
C	-3.043502	0.820878	-0.149821	C	-0.718085	2.105551	-0.106733			
C	-1.887126	0.137774	-0.160470		0.896350	0.503330	0.378512			
	-1.441882	-1.256744	0.108866		-0.384938	0.659690	-0.076621			
	-2.267948	-1.914363	0.393239		-1.435881	2.452625	-0.849645			
	-0.951711	-1.667938	-0.777728		-0.927003	2.520049	0.875808			
	-0.695679	-1.260617	0.907968	H	-1.014639	-0.115450	-0.520096			
H	-4.076299	0.569772	0.067786	H	1.393630	-0.440136	0.568491			
		[i3]				[i4]				
					0					
	-307	$7.0 \text{ kJ mol}^{-1}$			-306	$3.5 \text{ kJ mol}^{-1}$				
		$C_s^{-1}A'$				$C_1 - A$				
Si	-2.584650	2.826243	0.102164	Si	2.269227	0.314672	0.012013			
C	-1.277255	1.533566	0.133824	C	0.446698	0.519161	0.002884			
	-2.614644	1.183903	0.928846		-0.767059	0.515702	-0.005318			
	-1.976747	0.221766	-0.020795		-2.225868	0.529673	-0.008062			
	-1.488341	-0.647336	0.417639		2.557659	1.810564	-0.006117			
	-2.696767	0.944785	1.983522		-2.625510	-0.364956	-0.490733			
	-2.476241	-0.001272	-0.958483		-2.607844	1.412410	-0.525637			
H	-0.273455	1.579469	0.542808	H	-2.593502	0.553573	1.020967			

		[i5]		[i6]				
$-305.2 \text{ kJ mol}^{-1}$					-301	$1.9 \text{ kJ mol}^{-1}$		
$C_s^{-1}A'$						$C_1-^1A$		
Si	-2.791074	1.331229	0.575045	Si	1.271276	2.401422	0.597200	
C	0.041883	1.957213	0.098223	C	-0.676445	2.076496	-0.045615	
C	1.276133	1.530665	-0.176961	C	1.201685	0.439514	-0.022520	
C	-1.134334	1.101465	0.199981	C	0.071228	0.921046	-0.285832	
H	-0.109806	3.022058	0.263830	H	1.712703	2.984364	-0.742591	
H	2.111404	2.219175	-0.232581	H	-0.965219	2.689118	-0.894417	
H	1.481749	0.478866	-0.350967	H	-1.396336	2.031078	0.765002	
H	-0.944290	0.035021	0.000428	H	1.843949	-0.418828	-0.047808	
		[i7]				[i8]		
					0			
	-292	$2.0 \text{ kJ mol}^{-1}$			-280	$0.5 \text{ kJ mol}^{-1}$		
		$C_s^{-1}A'$				$C_s^{-1}A'$		
Si	-0.002314	3.177910	-0.169760	Si	-2.123430	2.732956	0.435983	
C	-0.649663	1.369072	-0.142375	C	-0.578881	1.995631	0.420032	
	0.826543	1.531957	-0.066304		-0.360326	0.594266	0.333811	
	C 1.909821 0.771693 0.020207				-0.179590	-0.593207	0.258429	
	-1.073424	0.953192	-1.057167		-2.265622	4.199834	0.528920	
	-1.173338	1.016264	0.746817		-3.379454	1.961598	0.360797	
H	1.853963	-0.315790	0.055069	H	0.306230	2.623143	0.479257	
H	2.902621	1.211905	0.059823	H	-0.015456	-1.642789	0.191810	

	[i9]				[i10]			
					<b>99-0-0</b> 0			
$-279.4 \text{ kJ mol}^{-1}$					-277	$7.3 \text{ kJ mol}^{-1}$		
$C_{s}^{-1}A'$						$C_s-A'$		
Si	-0.948677	2.863727	0.654535	Si	-0.333163	1.486867	-0.671031	
C	-1.217505	1.006333	0.704275	C	-0.144820	0.334917	0.683483	
C	-0.284681	0.271592	0.142353	C	-1.419478	-0.081122	0.274337	
C	0.670871	-0.400535	-0.428995	C	1.062694	0.035706	0.040856	
H	-2.235815	3.206243	1.405402	Н	-1.610019	-1.027212	-0.240233	
H	-2.068992	0.501898	1.155645	Н	-2.282879	0.366410	0.756404	
H	1.531568	-0.731501	0.140907	Н	1.241688	-0.893852	-0.506956	
H	0.619686	-0.645957	-1.483673	Н	1.954699	0.571689	0.349342	
		[i11]				[i12]		
					0.0			
	-273	$\frac{3.3 \text{ kJ mol}^{-1}}{2}$			-247	$7.9 \text{ kJ mol}^{-1}$		
	1.100700	$C_{s}^{-1}A'$			0.0000000	$C_s^{-1}A'$	1.0.10.170	
Si	1.409563	2.464346	0.895067	Si	-0.980308	2.921623	1.046459	
	-0.674602	1.916613	-0.147441		-1.407690	1.314359	0.084885	
	0.532533	1.098713	0.523749		-0.359350	0.302722	0.082071	
	-0.611275	0.405511	-0.024987		0.531367	-0.500595	0.106582	
	-0.504797	2.315714	-1.149586		-2.330725	3.551344	0.727266	
	-1.403810	2.412592	0.481019		-1.0/138/	1.007227	-0.940700	
	-0.4///05	-0.14/52/	-0.951107		-2.3300/1	0.913777	0.513390	
Н	-1.3128/3	-0.051284	0.008184	н	1.315062	-1.220401	0.122189	



Appendix A. Additional Informat	tion: Modeling of Gas Phase Formation of $c$ -Si $C_3$	3
Molecules in the Circumstellar Me	edium of Carbon Stars	13

		[i1]→[i5]		[i1]→[i6]				
$-142.6 \text{ kJ mol}^{-1}$					-125	$5.8 \text{ kJ mol}^{-1}$		
$C_1-^1A$						$C_1-^1A$		
Si	0.485782	2.153579	-0.469008	Si	0.778035	1.356341	0.013393	
C	0.332198	-1.132452	-0.095064	C	-1.041349	0.972089	0.089159	
C	0.161878	0.240941	-0.155304	C	-0.726348	-0.144886	-0.571048	
C	-0.869881	1.190140	0.041691	C	0.465469	-0.839573	-0.180070	
H	1.058802	-0.010869	0.587659	H	0.248932	-1.572742	0.597463	
H	-0.474833	-1.810179	0.181935	H	1.180352	-1.204148	-0.916275	
H	1.295799	-1.572078	-0.331053	H	1.081984	-0.069503	0.844891	
H	-1.919433	1.016042	0.235748	H	-1.980789	1.508295	0.137363	
		[i2]→[i3]				[i2]→[i5]		
	-188	$8.1 \text{ kJ mol}^{-1}$			-294	$4.1 \text{ kJ mol}^{-1}$		
		$C_1 - A$				$C_1-A$		
Si	-2.150053	2.178667	0.625370	Si	0.342188	2.735051	-0.130811	
C	-0.578457	1.585315	-0.456669	C	0.391241	-1.205372	-0.299825	
C	-1.232877	0.355972	-0.039025	C	0.039692	-0.188254	0.478770	
C	-0.432102	1.821019	0.963102	C	-0.308252	1.168082	-0.001663	
H	0.445188	1.756078	1.604662	H	0.000480	-0.344895	1.555941	
H	0.020532	1.864829	-1.311772	H	0.632717	-2.174234	0.123862	
H	-0.716789	-0.405155	0.544939	H	0.451087	-1.098752	-1.378354	
H	-2.145988	-0.008539	-0.518207	H	-1.322618	1.262237	-0.434147	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

	$[i2] \rightarrow [i10]$				[i2]→[i11]			
$-27.6 \text{ kJ mol}^{-1}$					8.5	$3 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.728857	1.050351	-1.035226	Si	0.711581	1.120945	-1.147132	
C	0.745086	-0.738470	-0.340304	C	0.472146	-0.659767	-0.443763	
C	0.053676	-0.073521	0.854344	C	-0.426581	0.126484	0.507436	
C	-0.693567	0.957443	0.165361	C	0.185285	1.423246	0.742216	
H	1.677882	-1.260348	-0.131790	Н	-1.131751	-0.376092	1.184455	
H	0.063165	-1.398234	-0.873462	Н	1.329286	-1.008742	0.132221	
H	-1.158678	-0.176597	0.686967	Н	0.056519	-1.475193	-1.037839	
H	-1.402374	1.604957	0.677393	Н	-1.169658	0.875873	-0.033549	
		$[i3] \rightarrow [i11]$			[	i4]→[i13]		
						•		
	-25	$1.7 \text{ kJ mol}^{-1}$			-224	$\frac{1.8 \text{ kJ mol}^{-1}}{1.1 \text{ kJ}}$		
		$C_1^{-1}A$				$C_s^{-1}A'$		
Si	-2.117498	2.386481	0.331886	Si	-1.742300	2.701762	0.558436	
C	-1.109699	0.390124	0.271138	C	-1.453027	0.995062	-0.232032	
	-0.453814	1.674840	-0.208502		-0.226982	1.212635	-0.272076	
	-1.026002	1.305652	1.486479		1.233250	1.262100	-0.377354	
	-2.106608	0.234772	-0.198696		-0.321943	3.146011	0.987967	
	0.243508	1.734296	-1.034771		1.549669	2.031502	-1.083304	
H	-0.560203	-0.548099	0.270135	H	1.675869	1.477113	0.596678	
H	0.211278	1.873915	0.737564	H	1.597521	0.293445	-0.727392	

		[i5]→[i9]		[i6]→[i9]				
	-206	$3.4 \text{ kJ mol}^{-1}$			-259	$0.9 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-1.992862	2.228890	0.527181	Si	1.820646	2.497868	0.585175	
C	-1.168691	0.701841	0.088179	C	1.254585	0.742483	0.074294	
C	0.014708	1.354943	0.010861	C	0.032876	1.141174	-0.073976	
C	1.299058	1.583736	-0.198471	C	-1.054450	1.887589	-0.083800	
H	-0.660165	2.892930	0.016175	H	2.460246	2.819309	-0.774176	
H	-1.325198	-0.327819	0.375544	H	1.690241	-0.197006	-0.235942	
H	1.802093	2.476996	0.142678	H	-1.294651	2.500444	-0.946357	
H	1.862043	0.843052	-0.756205	H	-1.774537	1.839849	0.726578	
		$i6] \rightarrow [i10]$				$i6] \rightarrow [i12]$		
					0-0			
	-198	$\frac{0.1 \text{ kJ mol}^{-1}}{0.1 \text{ kJ mol}^{-1}}$			-24	$\frac{1.9 \text{ kJ mol}^{-1}}{0.14}$		
C:	1 100009	$\frac{O_1^{}A}{2 42^{\circ \circ 44}}$	0 806755	C;	0.852600	$\frac{\cup_1^{}A}{1.625694}$	0 591065	
	1.100090 1.954157	2.420044 0.804199	0.090700 0.100514		0.002009	1.050084 0.493009	0.061000	
	-0 0/7300	0.004122	0.190314		-0.000000	-0.910555	0.009091	
	-0.677368	2 066019	0.000403		0.544992	-1 989473	-0.037835	
Н	2.054356	0.150676	-0 497980	Н	0.514552 0.555684	2 682867	-0 493235	
H	1 608219	2 139909	-0.662094	H	-1 048992	0.683920	-0.883382	
H	-0.645862	2.587746	-0.883451	H	-1.329382	0.463880	0.848535	

Η

-1.601289

2.227878

0.621898

Η

1.001958

-2.950630

-0.071481

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	136

		[i7]→[i10]		[i8]→[i9]				
$-276.7 \text{ kJ mol}^{-1}$					-56	$.0 \text{ kJ mol}^{-1}$		
$C_1-^1A$						$C_1-^1A$		
Si	-1.965407	2.054224	0.949023	Si	1.041061	0.917892	0.436916	
C	-2.941031	0.741834	-0.252682	C	0.760892	-0.738193	-0.161392	
C	-1.529983	0.959789	-0.397630	C	-0.609128	-0.750344	-0.306514	
C	-0.519325	0.330323	0.274811	C	-1.670823	-0.094074	-0.212414	
H	-3.585322	1.284613	-0.938932	H	2.501462	1.265763	0.648959	
H	-3.363470	-0.192898	0.123519	H	-0.818889	0.943557	0.334551	
H	-0.582867	-0.703830	0.623728	H	1.411713	-1.596634	-0.075037	
H	0.457528	0.798760	0.345200	H	-2.618347	0.185964	-0.624789	
		[i8]→[i12]				$[i5] \rightarrow p1$		
	-115	$5.2 \text{ kJ mol}^{-1}$			36	$\frac{.0 \text{ kJ mol}^{-1}}{2}$		
		$C_1$ - <sup>1</sup> A				$C_s^{-1}A'$		
Si	-2.521663	2.463165	1.136938	Si	-2.989127	-0.912311	-0.067647	
C	-0.857503	2.002660	0.565118		-1.436574	-0.366039	-0.033811	
	-0.462648	0.671903	0.277186		-0.145854	0.096600	-0.001158	
	-0.127979	-0.463490	0.061318		1.148919	-0.335722	0.029154	
	-0.053439	2.725894	0.410733	H	1.968519	0.370590	0.051314	
H	0.166755	-1.470185	-0.125512	H	1.364749	-1.396295	0.033802	
H	-2.334378	3.957638	1.028947	H	-0.052376	1.252833	0.469567	
H	-2.405675	1.983846	-0.345689	H	-0.028713	1.251089	-0.470006	

Appendix A. Additional Inform	nation: Modeling of Gas Phase Formation of c-SiC	ч /3
Molecules in the Circumstellar	Medium of Carbon Stars	13

[i9]→ <b>p1</b>				$[i1] \rightarrow p2$			
					•		
	-88	$.7 \text{ kJ mol}^{-1}$			7.	$0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_s-^1A'$	
Si	-0.744770	2.157892	1.133843	Si	0.528665	2.400596	-0.746415
C	-0.117447	0.739329	0.301978	C	0.638969	-0.741790	0.047537
C	0.508868	-0.265289	-0.184605	C	0.334217	0.578629	-0.236986
C	1.137912	-1.291212	-0.706981	C	-0.782728	1.293996	-0.211915
H	2.192710	-1.451493	-0.511820	H	1.142678	-0.981263	0.980498
H	0.620937	-2.000203	-1.344227	H	0.810919	-1.435189	-0.771673
H	-1.537155	0.742156	0.712436	H	-1.559747	-0.175915	0.315690
H	-2.215299	1.527879	0.692674	H	-1.065292	-0.900941	0.410427
		$[i6] \rightarrow p2$				$[i11] \rightarrow p2$	
					8.		
	-46	$.2 \text{ kJ mol}^{-1}$			98	$.6 \text{ kJ mol}^{-1}$	
		$C_1 - A$				$C_1 - A$	
Si	1.178923	2.205371	0.803390	Si	0.042413	-0.714013	0.848100
	-0.641206	2.137118	-0.050763		-0.717668	0.642379	-0.366158
	-0.130702	0.814244	0.053579		0.748846	0.531954	-0.607018
	0.888991	0.086553	0.200705		1.444960	-0.634113	-0.453494
H	-0.835280	2.504167	-1.054512	H	-1.128405	1.560574	0.057494
H	-1.380706	2.457777	0.683302	H	-1.303184	0.282882	-1.207928
H	1.725043	1.088245	-0.475804	H	1.486648	1.537155	-0.404833
H	1.923344	2.086626	-0.687100	H	1.243355	1.496133	-1.295948

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	138

[i13]→ <b>p2</b>				$[i12] \rightarrow \mathbf{p3}$			
						yed	
	-45	$.6 \text{ kJ mol}^{-1}$			-70	$.6 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-0.830773	0.883817	0.953003	Si	-1.151997	1.718947	0.728726
C	-1.811287	-0.696687	0.301549	C	-0.793552	-0.036891	0.485201
C	-0.590983	-0.923847	0.102221	C	0.356363	-0.744575	0.018188
C	0.749568	-0.461799	-0.040845	C	1.337204	-1.326892	-0.364743
H	1.185475	-0.546441	-1.038520	H	-1.597641	-0.704860	0.800559
H	1.455315	-0.696271	0.749183	H	2.205962	-1.845464	-0.696218
H	0.271987	0.942873	-0.437740	H	-0.463884	1.036474	-0.582223
H	-0.545547	1.525644	-0.576927	H	0.061553	1.931814	-0.358027
	[	i12]→ <b>p3</b> ′				$[i6] \rightarrow p4$	
	-76	$3 \text{ kJ mol}^{-1}$			-114	$4.2 \text{ kJ mol}^{-1}$	
		$C_1 - A$				$C_1 - A$	
Si	-1.555112	0.642765	0.911988	Si	0.831125	0.640403	0.777373
	0.049390	0.418870	0.114001		-1.185225	0.266693	0.157410
	0.581034	-0.893905	-0.073288		-0.415827	-0.797532	0.102736
	1.006592	-2.011224	-0.207659		0.854625	-1.135711	0.195070
H	0.726733	1.208649	-0.209397	H	-2.051331	0.445181	0.777573
H	1.387736	-2.997804	-0.330653	H	1.482402	-1.878498	-0.264555
H	-1.435106	2.141508	0.235326	H	-0.203090	1.162708	-0.807825
H	-1.002771	1.445970	-0.426216	H	0.676107	1.295027	-0.933420

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars

	$[i11] \rightarrow p5$						
	66.	$.7 \text{ kJ mol}^{-1}$					
		$C_1-^1A$					
Si	-0.987538	1.955490	-0.108967				
C	0.391227	0.544056	-0.099353				
C	-0.222777	-0.796271	-0.455253				
C	-1.149070	0.247028	-0.031728				
H	-0.029086	-1.630912	0.222284				
H	-0.176856	-1.062271	-1.507519				
H	1.344670	0.560033	0.664829				
H	0.776750	0.207977	1.273359				

## Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC3Molecules in the Circumstellar Medium of Carbon Stars140

Table A.2: Reactants, products, intermediates, and transition state structures relevant to the reaction of electronically excited atomic silicon with D3-methylacetylene. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

	Reactants								
	(	CD <sub>3</sub> CCH		Si					
Store Conce									
	0.0	) kJ mol <sup><math>-1</math></sup>							
$C_{3v}$ $-1A_1$					$^{1}\mathrm{D}$				
С	-1.919449	1.519750	0.000000	Si	0.000000	0.000000	0.000000		
C	-0.721221	1.524736	0.000000						
C	-3.382029	1.513693	0.000000						
H	0.343441	1.529181	0.000000						
D	-3.768514	1.986640	-0.905294						
D	-3.768812	2.058826	0.863618						
D	-3.762327	0.490833	0.041676						
			Proc	lucts					
		HD				$D_2$			
		$C_{\infty v}^{-1}\Sigma$				$D_{\infty h}^{-1}\Sigma$	_		
D	0.003216	0.000000	0.000000	D	0.003216	0.000000	0.000000		
H	0.746784	0.000000	0.000000	D	0.746784	0.000000	0.000000		

Appendix A. Addi	itional Information: N	Modeling of Gas Phase	Formation of $c$ -SiC <sub>3</sub>	
Molecules in the C	ircumstellar Medium	of Carbon Stars		141

p1a				p1b			
0=0=0=0							
	-183	$3.1 \text{ kJ mol}^{-1}$			-179	$0.8 \text{ kJ mol}^{-1}$	
		$C_{2v}$ - <sup>1</sup> $A_1$				$C_s-A'$	
Si	-2.667415	2.191276	0.000000	Si	-2.667415	2.191276	0.000000
C	-0.985018	2.220382	0.000000	C	-0.985018	2.220382	0.000000
C	0.298405	2.242563	0.000000	C	0.298405	2.242563	0.000000
C	1.607796	2.265203	0.000000	C	1.607796	2.265203	0.000000
D	2.142021	3.207971	0.000000	H	2.142021	3.207971	0.000000
D	2.174299	1.341470	0.000000	D	2.174299	1.341470	0.000000
		p2a				$\mathbf{p2b}$	
	-169	$\frac{9.2 \text{ kJ mol}^{-1}}{\Omega}$			-165	$\frac{5.1 \text{ kJ mol}^{-1}}{2}$	
	2 500055	$\frac{\mathrm{C_{s}}^{-1}\mathrm{A}'}{2}$	0.000000			$\frac{C_1 - A}{C_1 - A}$	0.000000
	-3.530077	0.607099	0.333008		-3.530077	0.607099	0.333008
	-2.115964	1.751749	-0.167586		-2.115964	1.751749	-0.167586
	-1.762971	0.449585	-0.634584		-1.762971	0.449585	-0.634584
	-2.980050	2.474570	0.443216		-2.980050	2.474570	0.443216
D	-0.954113	-0.065250	-0.122807		-0.954113	-0.065250	-0.122807
D	-1.792019	0.273905	-1.706863	D	-1.792019	0.273905	-1.706863

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	142

p2c				p3a			
	-165	$5.1 \text{ kJ mol}^{-1}$			-154	$4.0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_s-^1A'$	
Si	-3.530077	0.607099	0.333008	Si	-1.440031	2.135965	1.571949
C	-2.115964	1.751749	-0.167586	C	-1.258585	1.493357	-0.509032
C	-1.762971	0.449585	-0.634584	C	-1.918764	2.648067	0.002796
C	-2.980050	2.474570	0.443216	C	-0.632564	0.456668	-0.644013
D	-0.954113	-0.065250	-0.122807	D	-2.429665	3.372178	-0.610959
H	-1.792019	0.273905	-1.706863	H	-0.091486	-0.444116	-0.815669
		$\mathbf{p3b}$				p3c	
	-153	$\frac{3.5 \text{ kJ mol}^{-1}}{2}$			-157	$7.0 \text{ kJ mol}^{-1}$	
	1 4 40001	$\frac{\mathrm{C_{s}}^{-1}\mathrm{A}'}{2.125065}$	1 551040		1 4 40001	$\frac{\mathrm{C_{s}}^{-1}\mathrm{A}'}{2.125065}$	1 5 5 1 0 4 0
S <sub>1</sub>	-1.440031	2.135965	1.571949		-1.440031	2.135965	1.571949
	-1.258585	1.493357	-0.509032		-1.258585	1.493357	-0.509032
	-1.918764	2.048067	0.002796		-1.918764	2.048067	0.002796
	-0.032504	0.450008	-0.644013		-0.032504	0.450008	-0.644013
	-2.429665	3.372178	-0.610959		-2.429665	3.372178	-0.610959
$\mid D$	-0.091486	-0.444116	-0.815669	$\mid D$	-0.091486	-0.444116	-0.815669

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

p4a				p4b			
	-149	$9.7 \text{ kJ mol}^{-1}$			-15:	$3.1 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_2 - A_1$	
Si	0.015583	1.421694	0.441321	Si	0.015583	1.421694	0.441321
C	-1.710909	1.657474	1.228824	C	-1.710909	1.657474	1.228824
C	-1.833196	1.452865	-0.080188	C	-1.833196	1.452865	-0.080188
C	-1.051473	1.232997	-1.134185	C	-1.051473	1.232997	-1.134185
H	-2.213088	1.188068	2.060067	D	-2.213088	1.188068	2.060067
D	-1.030405	1.710504	-2.101157	D	-1.030405	1.710504	-2.101157
Intermediates							
	[i1a]					[i1b]	
	-381	$1.9 \text{ kJ mol}^{-1}$			-38	$1.2 \text{ kJ mol}^{-1}$	
		$C_s - A'$				$C_1 - A$	
Si	-1.656605	1.891479	-0.619303	Si	-1.656605	1.891479	-0.619303
C	-3.043502	0.820878	-0.149821	C	-3.043502	0.820878	-0.149821
	-1.887126	0.137774	-0.160470		-1.887126	0.137774	-0.160470
	-1.441882	-1.256744	0.108866		-1.441882	-1.256744	0.108866
D	-2.267948	-1.914363	0.393239	D	-2.267948	-1.914363	0.393239
D	-0.951711	-1.667938	-0.777728	D	-0.951711	-1.667938	-0.777728
D	-0.695679	-1.260617	0.907968	H	-0.695679	-1.260617	0.907968
H	-4.076299	0.569772	0.067786	D	-4.076299	0.569772	0.067786

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	144

		[i2a]		[i2b]			
	-360	$3.5 \text{ kJ mol}^{-1}$			-360	$3.5 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	1.217161	2.092109	-0.478041	Si	1.217161	2.092109	-0.478041
C	-0.718085	2.105551	-0.106733	C	-0.718085	2.105551	-0.106733
C	0.896350	0.503330	0.378512	C	0.896350	0.503330	0.378512
C	-0.384938	0.659690	-0.076621	C	-0.384938	0.659690	-0.076621
D	-1.435881	2.452625	-0.849645	D	-1.435881	2.452625	-0.849645
D	-0.927003	2.520049	0.875808	D	-0.927003	2.520049	0.875808
D	-1.014639	-0.115450	-0.520096	H	-1.014639	-0.115450	-0.520096
H	1.393630	-0.440136	0.568491	D	1.393630	-0.440136	0.568491
		[i2c]				i2d	
	-360	$\frac{5.2 \text{ kJ mol}^{-1}}{2}$			-360	$\frac{5.2 \text{ kJ mol}^{-1}}{2}$	
	1.0151.01	$C_1^{-1}A$	0.4500.41			$C_1^{-1}A$	
Si	1.217161	2.092109	-0.478041	Si	1.217161	2.092109	-0.478041
	-0.718085	2.105551	-0.106733		-0.718085	2.105551	-0.106733
	0.896350	0.503330	0.378512		0.896350	0.503330	0.378512
	-0.384938	0.659690	-0.076621		-0.384938	0.659690	-0.076621
	-1.435881	2.452625	-0.849645		-1.435881	2.452625	-0.849645
	-0.927003	2.520049	0.875808	H	-0.927003	2.520049	0.875808
D	-1.014639	-0.115450	-0.520096	D	-1.014639	-0.115450	-0.520096
D	1.393630	-0.440136	0.568491	D	1.393630	-0.440136	0.568491

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i2e]					[i2f]				
	-360	$3.5 \text{ kJ mol}^{-1}$			-360	$3.5 \text{ kJ mol}^{-1}$			
		$C_1-^1A$				$C_1-^1A$			
Si	-1.217161	2.092109	-0.478041	Si	-1.217161	2.092109	-0.478041		
C	0.718085	2.105551	-0.106733	C	0.718085	2.105551	-0.106733		
C	-0.896350	0.503330	0.378512	C	-0.896350	0.503330	0.378512		
C	0.384938	0.659690	-0.076621	C	0.384938	0.659690	-0.076621		
D	1.435881	2.452625	-0.849645	D	1.435881	2.452625	-0.849645		
D	0.927003	2.520049	0.875808	D	0.927003	2.520049	0.875808		
D	1.014639	-0.115450	-0.520096	Н	1.014639	-0.115450	-0.520096		
H	-1.393630	-0.440136	0.568491	D	-1.393630	-0.440136	0.568491		
		[i2g]				[i2h]			
[i2g]									
	-360	$5.2 \text{ kJ mol}^{-1}$			-360	$3.2 \text{ kJ mol}^{-1}$			
		$C_1$ - <sup>1</sup> A				$C_1 - A$			
Si	-1.217161	2.092109	-0.478041	Si	-1.217161	2.092109	-0.478041		
	0.718085	2.105551	-0.106733	C	0.718085	2.105551	-0.106733		
	-0.896350	0.503330	0.378512		-0.896350	0.503330	0.378512		
	0.384938	0.659690	-0.076621		0.384938	0.659690	-0.076621		
H	1.435881	2.452625	-0.849645	D	1.435881	2.452625	-0.849645		
D	0.927003	2.520049	0.875808	H	0.927003	2.520049	0.875808		
D	1.014639	-0.115450	-0.520096	D	1.014639	-0.115450	-0.520096		
D	-1.393630	-0.440136	0.568491	D	-1.393630	-0.440136	0.568491		

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 146

		[i3a]		[i3b]			
	-307	$7.7 \text{ kJ mol}^{-1}$			-307	$7.3 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_s-^1A'$	
Si	-2.584650	2.826243	0.102164	Si	-2.584650	2.826243	0.102164
C	-1.277255	1.533566	0.133824	C	-1.277255	1.533566	0.133824
C	-2.614644	1.183903	0.928846	C	-2.614644	1.183903	0.928846
C	-1.976747	0.221766	-0.020795	C	-1.976747	0.221766	-0.020795
D	-1.488341	-0.647336	0.417639	H	-1.488341	-0.647336	0.417639
H	-2.696767	0.944785	1.983522	D	-2.696767	0.944785	1.983522
D	-2.476241	-0.001272	-0.958483	D	-2.476241	-0.001272	-0.958483
D	-0.273455	1.579469	0.542808	D	-0.273455	1.579469	0.542808
		[i3c]				[i3d]	
	-307.5	$2 \text{ kJ kJ mol}^{-1}$			-307	$7.7 \text{ kJ mol}^{-1}$	
		$C_s - A'$				$C_1 - A$	
Si	-2.584650	2.826243	0.102164	Si	-2.584650	2.826243	0.102164
C	-1.277255	1.533566	0.133824	C	-1.277255	1.533566	0.133824
	-2.614644	1.183903	0.928846		-2.614644	1.183903	0.928846
	-1.976747	0.221766	-0.020795		-1.976747	0.221766	-0.020795
D	-1.488341	-0.647336	0.417639	D	-1.488341	-0.647336	0.417639
D	-2.696767	0.944785	1.983522	D	-2.696767	0.944785	1.983522
H	-2.476241	-0.001272	-0.958483	D	-2.476241	-0.001272	-0.958483
D	-0.273455	1.579469	0.542808	H	-0.273455	1.579469	0.542808

Appendix A. A	dditional Information:	Modeling of Gas Phase	Formation of $c$ -SiC <sub>3</sub>	
Molecules in the	e Circumstellar Medium	n of Carbon Stars	14	17

		[i4a]				[i4b]	
		0-0-0					
	-306	$3.4 \text{ kJ mol}^{-1}$			-302	$2.8 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	2.269227	0.314672	0.012013	Si	2.269227	0.314672	0.012013
C	0.446698	0.519161	0.002884	C	0.446698	0.519161	0.002884
C	-0.767059	0.515702	-0.005318	C	-0.767059	0.515702	-0.005318
C	-2.225868	0.529673	-0.008062	C	-2.225868	0.529673	-0.008062
H	2.557659	1.810564	-0.006117	D	2.557659	1.810564	-0.006117
D	-2.625510	-0.364956	-0.490733	D	-2.625510	-0.364956	-0.490733
D	-2.607844	1.412410	-0.525637	D	-2.607844	1.412410	-0.525637
D	-2.593502	0.553573	1.020967	H	-2.593502	0.553573	1.020967
		[i5a]				[i5b]	
	-308	$\frac{5.2 \text{ kJ mol}^{-1}}{\Omega \text{ s}^{-1} \text{ A}}$			-304	$\frac{4.1 \text{ kJ mol}^{-1}}{\Omega \sqrt{1}}$	
		$C_{\rm s}^{-1}A'$				$C_{s}^{-1}A'$	
	-2.791074	1.331229	0.575045		-2.791074	1.331229	0.575045
	0.041883	1.957213	0.098223		0.041883	1.957213	0.098223
	1.276133	1.530665	-0.176961		1.276133	1.530665	-0.176961
	-1.134334	1.101465	0.199981		-1.134334	1.101465	0.199981
	-0.109806	3.022058	0.203830		-0.109806	3.022058	0.203830
	2.111404	2.219175	-0.232581		2.111404	2.219175	-0.232581
	1.481749	0.478866	-0.350967		1.481749	0.478866	-0.350967
H	-0.944290	0.035021	0.000428	$\mid D$	-0.944290	0.035021	0.000428

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i5c] [i5d	[i5d]			
$-304.3 \text{ kJ mol}^{-1}$ $-304.4 \text{ kJ}$	$mol^{-1}$			
$C_s^{-1}A'$ $C_s^{-1}$	A'			
Si -2.791074 1.331229 0.575045 Si -2.791074 1.33	31229 $0.575045$			
C 0.041883 1.957213 0.098223 C 0.041883 1.95	57213 0.098223			
C 1.276133 1.530665 -0.176961 C 1.276133 1.55	30665 -0.176961			
C -1.134334 1.101465 0.199981 C -1.134334 1.10	01465 0.199981			
D -0.109806 3.022058 0.263830 D -0.109806 3.02	0.263830			
H 2.111404 2.219175 -0.232581 D 2.111404 2.2	19175 -0.232581			
D 1.481749 0.478866 -0.350967 H 1.481749 0.4'	78866 -0.350967			
D -0.944290 0.035021 0.000428 D -0.944290 0.03	35021 0.000428			
[i6a] [i6b	)]			
$-299.6 \text{ kJ mol}^{-1}$ $-301.5 \text{ kJ}$	mol <sup>-1</sup>			
$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	A 01400 0 505000			
$\begin{bmatrix} 51 & 1.271276 & 2.401422 & 0.597200 & 51 & 1.271276 & 2.40 \\ C & 0.676445 & 2.076406 & 0.045615 & C & 0.676445 & 2.00 \end{bmatrix}$	01422  0.597200			
$\begin{bmatrix} 0 & -0.070440 & 2.070490 & -0.040010 \\ C & 1.201685 & 0.420514 & 0.022520 \\ \end{bmatrix} \begin{bmatrix} 0 & -0.070440 & 2.0 \\ C & 1.201685 & 0.420514 \\ \end{bmatrix}$	10490 -0.049010			
$\begin{bmatrix} 0 & 1.201005 & 0.459514 & -0.022520 & 0 & 1.201085 & 0.46\\ C & 0.071528 & 0.021046 & 0.2858222 & C & 0.071528 & 0.06\\ \end{bmatrix}$	09014 -0.022020 01046 0.005020			
$\begin{bmatrix} 0 & 0.071220 & 0.921040 & -0.200052 & 0 & 0.071220 & 0.9. \\ D & 1.719703 & 2.084364 & 0.749501 & U & 1.719702 & 2.09 \end{bmatrix}$	21040 -0.20002 84364 0.749501			
$\begin{array}{ c c c c c c c c c c c c c c c c c c c$	04304 -0.742391			
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	QQ11Q <u>0Q</u> 0//17			
$(1)$ _1 306336 2 031078 0 765002 $(1)$ D 1 306336 2 05	89118 -0.894417 31078 0.765002			

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	149

		[i6c]		[i6d]			
	-298	$8.9 \text{ kJ mol}^{-1}$			-299	$0.0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	1.271276	2.401422	0.597200	Si	1.271276	2.401422	0.597200
C	-0.676445	2.076496	-0.045615	C	-0.676445	2.076496	-0.045615
C	1.201685	0.439514	-0.022520	C	1.201685	0.439514	-0.022520
C	0.071228	0.921046	-0.285832	C	0.071228	0.921046	-0.285832
D	1.712703	2.984364	-0.742591	D	1.712703	2.984364	-0.742591
D	-0.965219	2.689118	-0.894417	H	-0.965219	2.689118	-0.894417
H	-1.396336	2.031078	0.765002	D	-1.396336	2.031078	0.765002
D	1.843949	-0.418828	-0.047808	D	1.843949	-0.418828	-0.047808
		[i6e]				[i6f]	
	-299	$\frac{9.6 \text{ kJ mol}^{-1}}{2}$			-301	$1.5 \text{ kJ mol}^{-1}$	
	1.001000	$\frac{C_1 - A}{2 + 401 + 202}$	0 505200	C'	1.051050	$\frac{C_1 - A}{2 + 401 + 222}$	0.505000
$\begin{vmatrix} S_1 \\ C \end{vmatrix}$	-1.271276	2.401422	0.597200		-1.271276	2.401422	0.597200
	0.070443 1.201695	2.070490	-0.043013		0.070443 1.001695	2.070490	-0.043013
	-1.201080	0.439314	-0.022020		-1.201085	0.439314	-0.022020
	-0.071228 1 719702	0.921040 0.094264	-0.200832 0.749501		-0.071228 1.719702	0.921040 0.094264	-U.200002 0.749501
	-1.112100	2.304304 9.600110	-0.742091 0.804417		-1./12/03	2.304304 9.600110	-0.742091 0.804417
	0.900219	2.009110	-0.094417		1 206226	2.009110 2.021079	-0.094417
ц Ц	1.090000	2.031070 0.419999	0.700002		1 843040	2.031070 0.419999	0.705002
	-1.040949	-0.410020	-0.041000		-1.040949	-0.410020	-0.041000

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

		[i6g]		[i6h]			
	-298	$8.9 \text{ kJ mol}^{-1}$			-299	$0.0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-1.271276	2.401422	0.597200	Si	-1.271276	2.401422	0.597200
C	0.676445	2.076496	-0.045615	C	0.676445	2.076496	-0.045615
C	-1.201685	0.439514	-0.022520	C	-1.201685	0.439514	-0.022520
C	-0.071228	0.921046	-0.285832	C	-0.071228	0.921046	-0.285832
D	-1.712703	2.984364	-0.742591	D	-1.712703	2.984364	-0.742591
D	0.965219	2.689118	-0.894417	H	0.965219	2.689118	-0.894417
H	1.396336	2.031078	0.765002	D	1.396336	2.031078	0.765002
D	-1.843949	-0.418828	-0.047808	D	-1.843949	-0.418828	-0.047808
		[i7a]				[i7b]	
[i7a]							
	-291	$1.2 \text{ kJ mol}^{-1}$			-291	$1.2 \text{ kJ mol}^{-1}$	
		C <sub>1</sub> - <sup>1</sup> A				C <sub>1</sub> - <sup>1</sup> A	
Si	-0.002314	3.177910	-0.169760	Si	-0.002314	3.177910	-0.169760
	-0.649663	1.369072	-0.142375		-0.649663	1.369072	-0.142375
	0.826543	1.531957	-0.066304		0.826543	1.531957	-0.066304
	1.909821	0.771693	0.020207		1.909821	0.771693	0.020207
D	-1.073424	0.953192	-1.057167	H	-1.073424	0.953192	-1.057167
	-1.173338	1.016264	0.746817		-1.173338	1.016264	0.746817
D	1.853963	-0.315790	0.055069	D	1.853963	-0.315790	0.055069
D	2.902621	1.211905	0.059823	D	2.902621	1.211905	0.059823

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 151

		[i7c]				[i7d]	
	-290	$0.9 \text{ kJ mol}^{-1}$			-291	$1.1 \text{ kJ mol}^{-1}$	
		$C_s-^1A'$				$C_s-A'$	
Si	-0.002314	3.177910	-0.169760	Si	-0.002314	3.177910	-0.169760
C	-0.649663	1.369072	-0.142375	C	-0.649663	1.369072	-0.142375
C	0.826543	1.531957	-0.066304	C	0.826543	1.531957	-0.066304
C	1.909821	0.771693	0.020207	C	1.909821	0.771693	0.020207
D	-1.073424	0.953192	-1.057167	D	-1.073424	0.953192	-1.057167
D	-1.173338	1.016264	0.746817	D	-1.173338	1.016264	0.746817
D	1.853963	-0.315790	0.055069	H	1.853963	-0.315790	0.055069
H	2.902621	1.211905	0.059823	D	2.902621	1.211905	0.059823
		[i8a]				[i8b]	
	-275	$5.2 \text{ kJ mol}^{-1}$			-274	$\frac{1.7 \text{ kJ mol}^{-1}}{2}$	
		$C_s^{-1}A'$				$C_s^{-1}A'$	
Si	-2.123430	2.732956	0.435983	Si	-2.123430	2.732956	0.435983
C	-0.578881	1.995631	0.420032	C	-0.578881	1.995631	0.420032
C	-0.360326	0.594266	0.333811	C	-0.360326	0.594266	0.333811
	-0.179590	-0.593207	0.258429	C	-0.179590	-0.593207	0.258429
D	-2.265622	4.199834	0.528920	D	-2.265622	4.199834	0.528920
D	-3.379454	1.961598	0.360797	D	-3.379454	1.961598	0.360797
D	0.306230	2.623143	0.479257	H	0.306230	2.623143	0.479257
H	-0.015456	-1.642789	0.191810	D	-0.015456	-1.642789	0.191810

		[i8c]		[i8d]			
	-276	$5.8 \text{ kJ mol}^{-1}$			-276	$3.9 \text{ kJ mol}^{-1}$	
		$C_{s}-^{1}A'$				$C_{s}-^{1}A'$	
Si	-2.123430	2.732956	0.435983	Si	-2.123430	2.732956	0.435983
C	-0.578881	1.995631	0.420032	C	-0.578881	1.995631	0.420032
C	-0.360326	0.594266	0.333811	C	-0.360326	0.594266	0.333811
C	-0.179590	-0.593207	0.258429	C	-0.179590	-0.593207	0.258429
H	-2.265622	4.199834	0.528920	D	-2.265622	4.199834	0.528920
D	-3.379454	1.961598	0.360797	H	-3.379454	1.961598	0.360797
D	0.306230	2.623143	0.479257	D	0.306230	2.623143	0.479257
D	-0.015456	-1.642789	0.191810	D	-0.015456	-1.642789	0.191810
		[i9a]				[i9b]	
	6						
	-278	$5.3 \text{ kJ mol}^{-1}$			-278	$3.3 \text{ kJ mol}^{-1}$	
	0.010055	$C_{s}^{-1}A'$	0.00.00		0.040055	$C_s^{-1}A'$	0.05.1505
Si	-0.948677	2.863727	0.654535	Si	-0.948677	2.863727	0.654535
	-1.217505	1.006333	0.704275		-1.217505	1.006333	0.704275
	-0.284681	0.271592	0.142353		-0.284681	0.271592	0.142353
	0.070871	-0.400535	-0.428995		0.070871	-0.400535	-0.428995
	-2.235815	3.200243	1.405402		-2.235815	3.200243	1.405402
	-2.008992	0.301898	1.100040		-2.008992	0.301898	1.133043 0.140007
	1.001008	-0.731301	0.140907 1 402672		1.331308	-0.731301	0.140907 1 402672
D	0.019090	-0.043937	-1.403073	$D$	0.019080	-0.043937	-1.403073

		[i9c]		[i9d]			
	<u>و</u>						
	-275	$5.4 \text{ kJ mol}^{-1}$			-275	$5.4 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-0.948677	2.863727	0.654535	Si	-0.948677	2.863727	0.654535
C	-1.217505	1.006333	0.704275	C	-1.217505	1.006333	0.704275
C	-0.284681	0.271592	0.142353	C	-0.284681	0.271592	0.142353
C	0.670871	-0.400535	-0.428995	C	0.670871	-0.400535	-0.428995
D	-2.235815	3.206243	1.405402	D	-2.235815	3.206243	1.405402
D	-2.068992	0.501898	1.155645	D	-2.068992	0.501898	1.155645
D	1.531568	-0.731501	0.140907	H	1.531568	-0.731501	0.140907
H	0.619686	-0.645957	-1.483673	D	0.619686	-0.645957	-1.483673
		[i10a]				[i10b]	
	-277	$7.4 \text{ kJ mol}^{-1}$			-277	$7.1 \text{ kJ mol}^{-1}$	
		$C_1$ -'A				$C_1^{-1}A$	
Si	-0.333163	1.486867	-0.671031	Si	-0.333163	1.486867	-0.671031
	-0.144820	0.334917	0.683483		-0.144820	0.334917	0.683483
	-1.419478	-0.081122	0.274337		-1.419478	-0.081122	0.274337
	1.062694	0.035706	0.040856		1.062694	0.035706	0.040856
	-1.010019	-1.02/212	-0.240233		-1.010019	-1.027212	-0.240233
	-2.282879	0.300410	0.700404		-2.282879	0.300410	0.750404
	1.241088	-0.893852	-0.506956		1.241088	-0.893852	-0.200920
	1.954699	0.571689	0.349342	D	1.954699	0.571689	0.349342

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars



Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i12c]					[i12d]			
	-243	$3.7 \text{ kJ mol}^{-1}$			-243	$3.7 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-0.980308	2.921623	1.046459	Si	-0.980308	2.921623	1.046459	
C	-1.407690	1.314359	0.084885	C	-1.407690	1.314359	0.084885	
C	-0.359350	0.302722	0.082071	C	-0.359350	0.302722	0.082071	
C	0.531367	-0.500595	0.106582	C	0.531367	-0.500595	0.106582	
D	-2.336725	3.551344	0.727266	D	-2.336725	3.551344	0.727266	
H	-1.671387	1.607227	-0.940700	D	-1.671387	1.607227	-0.940700	
D	-2.336671	0.913777	0.513390	H	-2.336671	0.913777	0.513390	
D	1.315662	-1.220461	0.122189	D	1.315662	-1.220461	0.122189	
		[i13a]		[i13b]				
	-224	$4.2 \text{ kJ mol}^{-1}$		-220.8 kJ mol <sup>-1</sup>				
		$C_s^{-1}A'$		~		$C_s^{-1}A'$		
Si	-1.688737	2.624682	0.616789	Si	-1.688737	2.624682	0.616789	
	-0.244907	1.218317	-0.220340		-0.244907	1.218317	-0.220340	
	1.218639	1.272651	-0.343790		1.218639	1.272651	-0.343790	
	-1.465384	0.927557	-0.237325		-1.465384	0.927557	-0.237325	
	-0.260660	3.204713	0.754427		-0.260660	3.204713	0.754427	
	1.524821	2.129620	-0.945098		1.524821	2.129620	-0.945098	
	1.687125	1.351436	0.638071		1.687125	1.351436	0.638071	
D	1.564998	0.357273	-0.829223	H	1.564998	0.357273	-0.829223	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

Transition State Structures								
	[j	$i1a] \rightarrow [i2a]$			[]	$i1a] \rightarrow [i2e]$		
	-75	$0.0 \text{ kJ mol}^{-1}$			-75	$.0 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.387656	1.918057	-0.868165	Si	-0.387656	1.918057	-0.868165	
C	0.507204	-1.109474	-0.273582	C	-0.507204	-1.109474	-0.273582	
C	-0.001009	0.171130	-0.093798	C	0.001009	0.171130	-0.093798	
C	-0.892669	1.297958	0.108223		0.892669	1.297958	0.108223	
D	1.347937	-1.296852	-0.939211	D	-1.347937	-1.296852	-0.939211	
D	0.177136	-0.265142	0.972772	D	-0.177136	-0.265142	0.972772	
D	0.170559	-1.954642	0.325094	D	-0.170559	-1.954642	0.325094	
H	-1.765240	1.299326	0.746709	H	1.765240	1.299326	0.746709	
	li	i1b]→[i2b]			li	$i2b] \rightarrow [i2c]$		
	-70	$\frac{0.5 \text{ KJ mol}^{-1}}{C^{-1}}$			- ( 4	$\frac{1}{C}$		
C;	0.387656	$\frac{0.1^{-} \text{ A}}{1.018057}$	0.869165	Q;	0.387656	$\frac{O_1 - A}{1.018057}$	0.869165	
	0.387030	-1 100/7/	-0.000100		0.307030	-1 100/7/	-0.000100	
		0 171120	-0.273382			0 171120	-0.273382	
	-0.802660	1 207058	0.108223		-0.892660	1 297958	0.108223	
	1 347937	-1 296852	-0.939211	D D	1.347937	-1.296852	-0.939211	
H	0.177136	-0.265142	0.972772		0.177136	-0.265142	0.972772	
	0.170559	-1 954642	0.325094	Н	0 170559	-1 954642	0.325094	
D	-1.765240	1.299326	0.746709	D	-1.765240	1.299326	0.746709	

[i1b]→[i2d]					[i1b]→[i2f]			
	-74	$.5 \text{ kJ mol}^{-1}$			-76	$.5 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.387656	1.918057	-0.868165	Si	-0.387656	1.918057	-0.868165	
C	0.507204	-1.109474	-0.273582	C	-0.507204	-1.109474	-0.273582	
C	-0.001009	0.171130	-0.093798	C	0.001009	0.171130	-0.093798	
C	-0.892669	1.297958	0.108223	C	0.892669	1.297958	0.108223	
H	1.347937	-1.296852	-0.939211	D	-1.347937	-1.296852	-0.939211	
D	0.177136	-0.265142	0.972772	Н	-0.177136	-0.265142	0.972772	
D	0.170559	-1.954642	0.325094	D	-0.170559	-1.954642	0.325094	
D	-1.765240	1.299326	0.746709	D	1.765240	1.299326	0.746709	
	[i	$1b] \rightarrow [i2g]$		[i1b]→[i2h]				
					G			
	-74	$\frac{1.7 \text{ kJ mol}^{-1}}{2}$			-74	$\frac{.5 \text{ kJ mol}^{-1}}{2}$		
	0.005050	$C_1 - A$	0.0001.07			$C_1 - A$	0.0001.07	
Si	-0.387656	1.918057	-0.868165	Si	-0.387656	1.918057	-0.868165	
	-0.507204	-1.109474	-0.273582		-0.507204	-1.109474	-0.273582	
	0.001009	0.171130	-0.093798		0.001009	0.171130	-0.093798	
	0.892669	1.297958	0.108223		0.892669	1.297958	0.108223	
	-1.347937	-1.296852	-0.939211		-1.347937	-1.296852	-0.939211	
	-0.177136	-0.265142	0.972772		-0.177136	-0.265142	0.972772	
	-0.170559	-1.954642	0.325094		-0.170559	-1.954642	0.325094	
$\mid D$	1.765240	1.299326	0.746709	$\mid D$	1.705240	1.299326	0.746709	

[i1a]→[i4a]					[i1b]→[i4b]			
	•							
	-194	$1.4 \text{ kJ mol}^{-1}$			-190	$0.8 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-2.122577	1.493357	-0.465731	Si	-2.122577	1.493357	-0.465731	
C	-3.577778	0.142016	0.174868	C	-3.577778	0.142016	0.174868	
C	-2.426216	-0.379803	0.065365	C	-2.426216	-0.379803	0.065365	
C	-1.341337	-1.393100	0.060616	C	-1.341337	-1.393100	0.060616	
D	-1.737882	-2.337896	-0.315266	D	-1.737882	-2.337896	-0.315266	
D	-0.507637	-1.078652	-0.566987	D	-0.507637	-1.078652	-0.566987	
D	-0.973850	-1.547956	1.076886	Н	-0.973850	-1.547956	1.076886	
H	-3.309526	2.412152	-0.373543	D	-3.309526	2.412152	-0.373543	
	[i	$1a] \rightarrow [i5a]$			[i	1b]→[i5b]		
	•							
	-14(	$0.2 \text{ kJ mol}^{-1}$		-141.9 kJ mol <sup>-1</sup>				
		C <sub>1</sub> - <sup>1</sup> A				C <sub>1</sub> - <sup>1</sup> A		
Si	0.485782	2.153579	-0.469008	Si	0.485782	2.153579	-0.469008	
	0.332198	-1.132452	-0.095064		0.332198	-1.132452	-0.095064	
	0.161878	0.240941	-0.155304		0.161878	0.240941	-0.155304	
	-0.869881	1.190140	0.041691		-0.869881	1.190140	0.041691	
	1.058802	-0.010869	0.587659		1.058802	-0.010869	0.587659	
	-0.474833	-1.810179	0.181935		-0.474833	-1.810179	0.181935	
	1.295799	-1.572078	-0.331053		1.295799	-1.572078	-0.331053	
H	-1.919433	1.016042	0.235748	$\mid D$	-1.919433	1.016042	0.235748	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i1b]→[i5c]					[i1b]→[i5d]			
	-139	$9.5 \text{ kJ mol}^{-1}$			-139	$9.8 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.485782	2.153579	-0.469008	Si	0.485782	2.153579	-0.469008	
C	0.332198	-1.132452	-0.095064	C	0.332198	-1.132452	-0.095064	
C	0.161878	0.240941	-0.155304	C	0.161878	0.240941	-0.155304	
C	-0.869881	1.190140	0.041691	C	-0.869881	1.190140	0.041691	
D	1.058802	-0.010869	0.587659	D	1.058802	-0.010869	0.587659	
D	-0.474833	-1.810179	0.181935	H	-0.474833	-1.810179	0.181935	
H	1.295799	-1.572078	-0.331053	D	1.295799	-1.572078	-0.331053	
D	-1.919433	1.016042	0.235748	D	-1.919433	1.016042	0.235748	
	[j	$(1a] \rightarrow [i6a]$			[]	i1a]→[i6e]		
	-122	$\frac{2.9 \text{ kJ mol}^{-1}}{\text{G}^{-1}}$		-122.9 kJ mol <sup>-1</sup>				
	0.550005	$\frac{C_1 - A}{1 - 25 + 20 + 1}$	0.010000	a.	0.550005	$\frac{C_1 - A}{1.2562.41}$	0.010000	
	0.778035	1.356341	0.013393		-0.778035	1.356341	0.013393	
	-1.041349	0.972089	0.089159		1.041349	0.972089	0.089159	
	-0.726348	-0.144886	-0.5/1048		0.726348	-0.144886	-0.5/1048	
	0.405469	-0.8395/3	-0.180070		-0.405469	-0.8395/3	-0.180070	
	0.248932	-1.3(2(42)	0.59/403		-0.248932	-1.3(2(42)	0.016075	
	1.100302	-1.204148	-0.910273		-1.100302	-1.204148	-0.910279	
	1.081984	-0.009503 1.509905	0.127262		-1.081984	-0.009503 1 509905	0.844891	
П	-1.980789	1.308295	0.137303	П	1.980789	1.308293	0.137303	

Appendix A. Add	itional Information: 1	Modeling of Gas	Phase Formation	of $c$ -SiC <sub>3</sub>	
Molecules in the $C$	'ircumstellar Medium	of Carbon Stars			160

[i1b]→[i6b]					[i1b]→[i6c]			
	-125	$5.8 \text{ kJ mol}^{-1}$			-122	$2.4 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.778035	1.356341	0.013393	Si	0.778035	1.356341	0.013393	
C	-1.041349	0.972089	0.089159	C	-1.041349	0.972089	0.089159	
C	-0.726348	-0.144886	-0.571048	C	-0.726348	-0.144886	-0.571048	
C	0.465469	-0.839573	-0.180070	C	0.465469	-0.839573	-0.180070	
D	0.248932	-1.572742	0.597463	D	0.248932	-1.572742	0.597463	
D	1.180352	-1.204148	-0.916275	H	1.180352	-1.204148	-0.916275	
H	1.081984	-0.069503	0.844891	D	1.081984	-0.069503	0.844891	
D	-1.980789	1.508295	0.137363	D	-1.980789	1.508295	0.137363	
	[i	.1b]→[i6d]			[	i1b]→[i6f]		
	-122	$\frac{2.4 \text{ kJ mol}^{-1}}{2.4 \text{ kJ mol}^{-1}}$		$-125.8 \text{ kJ mol}^{-1}$				
	0.880005	$\frac{C_1 - A}{1 - 25 + 25 + 4}$	0.010002		0.880005	$C_1^{-1}A$	0.010002	
$  S_1   C$	0.778035	1.356341	0.013393		-0.778035	1.356341	0.013393	
	-1.041349	0.972089	0.089159		1.041349	0.972089	0.089159	
	-0.720348	-0.144880	-0.5(1048)		0.720348	-0.144880	-0.371048	
	0.400409	-0.8395/3	-0.180070		-0.403409	-0.839373	-0.180070	
	0.248932	-1.3(2(42)	0.397403		-0.248932	-1.3(2(42)	0.09/403	
	1.180392	-1.204148	-0.910275		-1.180302	-1.204148	-0.910270	
	1.081984	-0.009503	0.127262		-1.081984	-0.009503	0.127262	
	-1.980789	1.508295	0.137363	D	1.980789	1.508295	0.137363	

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC	3
Molecules in the Circumstellar Medium of Carbon Stars	161

[i1b]→[i6g]					[i1b]→[i6h]			
	-122	$2.4 \text{ kJ mol}^{-1}$			-122	$2.4 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-0.778035	1.356341	0.013393	Si	-0.778035	1.356341	0.013393	
C	1.041349	0.972089	0.089159	C	1.041349	0.972089	0.089159	
C	0.726348	-0.144886	-0.571048	C	0.726348	-0.144886	-0.571048	
C	-0.465469	-0.839573	-0.180070	C	-0.465469	-0.839573	-0.180070	
D	-0.248932	-1.572742	0.597463	H	-0.248932	-1.572742	0.597463	
H	-1.180352	-1.204148	-0.916275	D	-1.180352	-1.204148	-0.916275	
D	-1.081984	-0.069503	0.844891	D	-1.081984	-0.069503	0.844891	
D	1.980789	1.508295	0.137363	D	1.980789	1.508295	0.137363	
	[]	$i2a] \rightarrow [i2e]$			[	$i2b] \rightarrow [i2f]$		
	-343	$\frac{3.8 \text{ kJ mol}^{-1}}{2}$		-344.1 kJ mol <sup>-1</sup>				
		$C_s^{-1}A'$				$C_s^{-1}A'$		
Si	1.278722	2.298406	-0.055446	Si	1.278722	2.298406	-0.055446	
C	-0.670175	2.114670	-0.122843		-0.670175	2.114670	-0.122843	
	0.931362	0.457675	0.076046		0.931362	0.457675	0.076046	
	-0.403313	0.629572	0.008121		-0.403313	0.629572	0.008121	
	-1.103674	2.406338	-1.054805		-1.103074	2.406338	-1.054805	
	-1.238437	2.548998	0.705255		-1.238437	2.548998	0.705255	
	-1.174293	-0.147915	0.038856		-1.174293	-0.147915	0.038856	
H	1.448755	-0.490942	0.172710	D	1.448755	-0.490942	0.172710	
Appendix A. Additiona	al Information: Modeling of C	Gas Phase Formation of $c$ -Si $C_3$						
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Molecules in the Circun	nstellar Medium of Carbon S	tars 1	6					

$[i2c] \rightarrow [i2h]$				$[i2d] \rightarrow [i2g]$			
-343.7 kJ mol <sup>-1</sup>					-343	$3.7 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	1.278722	2.298406	-0.055446	Si	1.278722	2.298406	-0.055446
C	-0.670175	2.114670	-0.122843	C	-0.670175	2.114670	-0.122843
C	0.931362	0.457675	0.076046	C	0.931362	0.457675	0.076046
C	-0.403313	0.629572	0.008121	C	-0.403313	0.629572	0.008121
H	-1.163674	2.406338	-1.054805	D	-1.163674	2.406338	-1.054805
D	-1.238437	2.548998	0.705255	H	-1.238437	2.548998	0.705255
D	-1.174293	-0.147915	0.038856	D	-1.174293	-0.147915	0.038856
D	1.448755	-0.490942	0.172710	D	1.448755	-0.490942	0.172710
	[j	$i2a] \rightarrow [i3a]$			[i	$2b] \rightarrow [i3d]$	
	-187	$7.7 \text{ kJ mol}^{-1}$			-187	$7.3 \text{ kJ mol}^{-1}$	
	0.150050	$C_1^{-1}A$	0.005050		0.150050	$C_1^{-1}A$	0.005250
Si	-2.150053	2.178667	0.625370	Si	-2.150053	2.178667	0.625370
	-0.578457	1.585315	-0.456669		-0.578457	1.585315	-0.456669
	-1.232877	0.355972	-0.039025		-1.232877	0.355972	-0.039025
	-0.432102	1.821019	0.963102		-0.432102	1.821019	0.963102
	0.445188	1.756078	1.604662		0.445188	1.756078	1.604662
	0.020532	1.864829	-1.311772		0.020532	1.864829	-1.311772
D	-0.716789	-0.405155	0.544939	D	-0.716789	-0.405155	0.544939
D	-2.145988	-0.008539	-0.518207	D	-2.145988	-0.008539	-0.518207

Appendix A. Ad	ditional Information:	Modeling of Gas	Phase Formation	of $c$ -Si $C_3$	
Molecules in the	Circumstellar Mediun	n of Carbon Stars			16

[i2c]→[i3c]				[i2d]→[i3b]			
	-18	$7.5 \text{ kJ mol}^{-1}$			-187	$7.0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-2.150053	2.178667	0.625370	Si	-2.150053	2.178667	0.625370
C	-0.578457	1.585315	-0.456669	C	-0.578457	1.585315	-0.456669
C	-1.232877	0.355972	-0.039025	C	-1.232877	0.355972	-0.039025
C	-0.432102	1.821019	0.963102	C	-0.432102	1.821019	0.963102
D	0.445188	1.756078	1.604662	D	0.445188	1.756078	1.604662
D	0.020532	1.864829	-1.311772	D	0.020532	1.864829	-1.311772
D	-0.716789	-0.405155	0.544939	H	-0.716789	-0.405155	0.544939
H	-2.145988	-0.008539	-0.518207	D	-2.145988	-0.008539	-0.518207
	[j	$i2e] \rightarrow [i3d]$			[	i2f]→[i3a]	
	-187	$7.7 \text{ kJ mol}^{-1}$			-187	$7.3 \text{ kJ mol}^{-1}$	
		$C_1 - A$	0.005050			$C_1 - A$	0.005050
	2.150053	2.178667	0.625370		2.150053	2.178667	0.625370
	0.578457	1.585315	-0.456669		0.578457	1.585315	-0.456669
	1.232877	0.355972	-0.039025		1.232877	0.355972	-0.039025
	0.432102	1.821019	0.963102		0.432102	1.821019	0.963102
	-0.445188	1.756078	1.004662		-0.445188	1.756078	1.004062
	-0.020532	1.804829	-1.311772		-0.020532	1.804829	-1.311772
	0.716789	-0.405155	0.544939		0.716789	-0.405155	0.544939
$\mid D$	2.145988	-0.008539	-0.518207	$\mid D$	2.145988	-0.008539	-0.518207

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i2g]→[i3c]				$[i2h] \rightarrow [i3b]$			
$-187.5 \text{ kJ mol}^{-1}$					-18	$7.0 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	2.150053	2.178667	0.625370	Si	2.150053	2.178667	0.625370
C	0.578457	1.585315	-0.456669	C	0.578457	1.585315	-0.456669
C	1.232877	0.355972	-0.039025	C	1.232877	0.355972	-0.039025
C	0.432102	1.821019	0.963102	C	0.432102	1.821019	0.963102
D	-0.445188	1.756078	1.604662	D	-0.445188	1.756078	1.604662
D	-0.020532	1.864829	-1.311772	D	-0.020532	1.864829	-1.311772
D	0.716789	-0.405155	0.544939	H	0.716789	-0.405155	0.544939
H	2.145988	-0.008539	-0.518207	D	2.145988	-0.008539	-0.518207
	[]	$i2a] \rightarrow [i5a]$			[i	2b]→[i5b]	
	-294	$\frac{4.0 \text{ kJ mol}^{-1}}{1000000000000000000000000000000000$			-292	$2.4 \text{ kJ mol}^{-1}$	
		$C_1$ - <sup>1</sup> A				$C_1$ -1A	
Si	0.342188	2.735051	-0.130811	Si	0.342188	2.735051	-0.130811
	0.391241	-1.205372	-0.299825		0.391241	-1.205372	-0.299825
C	0.039692	-0.188254	0.478770	C	0.039692	-0.188254	0.478770
	-0.308252	1.168082	-0.001663		-0.308252	1.168082	-0.001663
	0.000480	-0.344895	1.555941		0.000480	-0.344895	1.555941
	0.632717	-2.174234	0.123862		0.632717	-2.174234	0.123862
	0.451087	-1.098752	-1.378354	D	0.451087	-1.098752	-1.378354
H	-1.322618	1.262237	-0.434147	$\mid D$	-1.322618	1.262237	-0.434147

Appendix A. A	dditional Information:	Modeling of Gas	Phase Formation	of $c$ -Si $C_3$
Molecules in the	e Circumstellar Medium	n of Carbon Stars		16!

[i2c]→[i5c]			[i2d]→[i5d]				
$-292.4 \text{ kJ mol}^{-1}$					-292	$2.5 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	0.342188	2.735051	-0.130811	Si	0.342188	2.735051	-0.130811
C	0.391241	-1.205372	-0.299825	C	0.391241	-1.205372	-0.299825
C	0.039692	-0.188254	0.478770	C	0.039692	-0.188254	0.478770
C	-0.308252	1.168082	-0.001663	C	-0.308252	1.168082	-0.001663
D	0.000480	-0.344895	1.555941	D	0.000480	-0.344895	1.555941
H	0.632717	-2.174234	0.123862	D	0.632717	-2.174234	0.123862
D	0.451087	-1.098752	-1.378354	H	0.451087	-1.098752	-1.378354
D	-1.322618	1.262237	-0.434147	D	-1.322618	1.262237	-0.434147
	[j	$i2e] \rightarrow [i5a]$			[-	$i2f] \rightarrow [i5b]$	
	-294	$4.0 \text{ kJ mol}^{-1}$			-292	$2.4 \text{ kJ mol}^{-1}$	
		$C_1$ - <sup>1</sup> A				$C_1$ - <sup>1</sup> A	
Si	-0.342188	2.735051	-0.130811	Si	-0.342188	2.735051	-0.130811
C	-0.391241	-1.205372	-0.299825		-0.391241	-1.205372	-0.299825
	-0.039692	-0.188254	0.478770		-0.039692	-0.188254	0.478770
	0.308252	1.168082	-0.001663		0.308252	1.168082	-0.001663
	-0.000480	-0.344895	1.555941		-0.000480	-0.344895	1.555941
	-0.632717	-2.1(4234	0.123862		-0.632717	-2.1(4234	0.123862
	-0.451087	-1.098752	-1.378354		-0.451087	-1.098752	-1.378354
H	1.322618	1.262237	-0.434147	D	1.322618	1.262237	-0.434147

	[j	[2g]→[i5c]		[i2h]→[i5d]				
							6	
	-292	$2.4 \text{ kJ mol}^{-1}$		$-292.5 \text{ kJ mol}^{-1}$				
		$C_1 - A$				$C_1 - A$		
Si	-0.342188	2.735051	-0.130811	Si	-0.342188	2.735051	-0.130811	
C	-0.391241	-1.205372	-0.299825	C	-0.391241	-1.205372	-0.299825	
C	-0.039692	-0.188254	0.478770	C	-0.039692	-0.188254	0.478770	
	0.308252	1.168082	-0.001663	C	0.308252	1.168082	-0.001663	

D

D

Η

D

-0.000480

-0.632717

-0.451087

1.322618

-0.344895

-2.174234

-1.098752

1.262237

 $[i2b] \rightarrow [i10c]$ 

1.555941

0.123862

-1.378354

-0.434147

1.555941

0.123862

-1.378354

-0.434147

D

Η

D

D

-0.000480

-0.632717

-0.451087

1.322618

-0.344895

-2.174234

-1.098752

1.262237

 $[i2a] \rightarrow [i10d]$ 

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

$-25.2 \text{ kL mol}^{-1}$								
$-25.2 \text{ kJ mol}^{-1}$					$-27.6 \text{ kJ mol}^{-1}$			
		$C_1 - A$		$C_1$ $^1A$				
Si	0.728857	1.050351	-1.035226	Si	0.728857	1.050351	-1.035226	
С	0.745086	-0.738470	-0.340304	C	0.745086	-0.738470	-0.340304	
С	0.053676	-0.073521	0.854344	C	0.053676	-0.073521	0.854344	
$\mathbf{C}$	-0.693567	0.957443	0.165361	C	-0.693567	0.957443	0.165361	
D	1.677882	-1.260348	-0.131790	D	1.677882	-1.260348	-0.131790	
D	0.063165	-1.398234	-0.873462	D	0.063165	-1.398234	-0.873462	
D	-1.158678	-0.176597	0.686967	H	-1.158678	-0.176597	0.686967	
Η	-1.402374	1.604957	0.677393	D	-1.402374	1.604957	0.677393	



Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

[i2g]→[i10d]				$[i2h] \rightarrow [i10c]$			
-25.0 kJ mol <sup>-1</sup>					-25	$.3 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-0.728857	1.050351	-1.035226	Si	-0.728857	1.050351	-1.035226
C	-0.745086	-0.738470	-0.340304	C	-0.745086	-0.738470	-0.340304
C	-0.053676	-0.073521	0.854344	C	-0.053676	-0.073521	0.854344
C	0.693567	0.957443	0.165361	C	0.693567	0.957443	0.165361
H	-1.677882	-1.260348	-0.131790	D	-1.677882	-1.260348	-0.131790
D	-0.063165	-1.398234	-0.873462	Н	-0.063165	-1.398234	-0.873462
D	1.158678	-0.176597	0.686967	D	1.158678	-0.176597	0.686967
D	1.402374	1.604957	0.677393	D	1.402374	1.604957	0.677393
	[i4	4a]→[i13a]			[i4	$4b] \rightarrow [i13b]$	
	-224	$\frac{1.9 \text{ kJ mol}^{-1}}{2.0 \text{ kJ mol}^{-1}}$			-221	$1.4 \text{ kJ mol}^{-1}$	
	1	$C_s^{-1}A'$				$C_s^{-1}A'$	
Si	-1.742300	2.701762	0.558436	Si	-1.742300	2.701762	0.558436
	-1.453027	0.995062	-0.232032		-1.453027	0.995062	-0.232032
	-0.226982	1.212635	-0.272076		-0.226982	1.212635	-0.272076
	1.233250	1.262100	-0.377354		1.233250	1.202100	-0.377354
	-0.321943	3.140011	0.987967		-0.321943	3.140011	0.987967
	1.349009	2.051502	-1.083304		1.349009	2.031302	-1.083304
	1.075809	1.4//113	0.590078		1.075809	1.4//113	0.596678
	1.597521	0.293445	-0.727392	μН	1.597521	0.293445	-0.727392

[i5a]→[i9a]			$[i5b] \rightarrow [i9b]$				
	-202	$2.5 \text{ kJ mol}^{-1}$			-205	$5.3 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-1.992862	2.228890	0.527181	Si	-1.992862	2.228890	0.527181
C	-1.168691	0.701841	0.088179	C	-1.168691	0.701841	0.088179
C	0.014708	1.354943	0.010861	C	0.014708	1.354943	0.010861
C	1.299058	1.583736	-0.198471	C	1.299058	1.583736	-0.198471
D	-0.660165	2.892930	0.016175	H	-0.660165	2.892930	0.016175
H	-1.325198	-0.327819	0.375544	D	-1.325198	-0.327819	0.375544
D	1.802093	2.476996	0.142678	D	1.802093	2.476996	0.142678
D	1.862043	0.843052	-0.756205	D	1.862043	0.843052	-0.756205
	ĺ	$i5c] \rightarrow [i9d]$			ĺi	5d]→[i9c]	
	-202	$\frac{2.3 \text{ kJ mol}^{-1}}{2.3 \text{ kJ mol}^{-1}}$			-202	$\frac{2.5 \text{ kJ mol}^{-1}}{2.5 \text{ kJ mol}^{-1}}$	
	1.000000	$\frac{\text{U}_1^{-1}\text{A}}{2,222222}$	0 505101	<u> </u>	1.000000	$\frac{\text{U}_1 - \text{A}}{\text{A}}$	0 505101
$\begin{vmatrix} S_1 \\ C \end{vmatrix}$	-1.992862	2.228890	0.527181		-1.992862	2.228890	0.527181
	-1.108091	U. (U1841 1 254042	0.008179		-1.108091	0.701841	0.088179
	0.014/08	1.334943	0.010801		0.014/08	1.334943	0.010801
	1.299038	1.003/30	-0.1984/1		1.299008	1.383/30	-0.1984/1
	-0.000105	2.892930	0.010175		-0.000105	2.892930	0.0101/5
	-1.325198	-0.32/819	0.375544		-1.325198	-0.327819	0.375544
	1.802093	2.476996	0.142678		1.802093	2.476996	0.142678
$\mid D$	1.862043	0.843052	-0.756205	H	1.862043	0.843052	-0.756205

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$		
Molecules in the Circumstellar Medium of Carbon Stars	170	

	[j	i6a]→[i9a]		[i6b]→[i9b]				
	-256	$3.0 \text{ kJ mol}^{-1}$			-258	$8.8 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	1.820646	2.497868	0.585175	Si	1.820646	2.497868	0.585175	
C	1.254585	0.742483	0.074294	C	1.254585	0.742483	0.074294	
C	0.032876	1.141174	-0.073976	C	0.032876	1.141174	-0.073976	
C	-1.054450	1.887589	-0.083800	C	-1.054450	1.887589	-0.083800	
D	2.460246	2.819309	-0.774176	H	2.460246	2.819309	-0.774176	
H	1.690241	-0.197006	-0.235942	D	1.690241	-0.197006	-0.235942	
D	-1.294651	2.500444	-0.946357	D	-1.294651	2.500444	-0.946357	
D	-1.774537	1.839849	0.726578	D	-1.774537	1.839849	0.726578	
	[	i6c]→[i9c]			[i	6d]→[i9d]		
	-258	$\frac{5.8 \text{ kJ mol}^{-1}}{\text{C}}$		$-255.7 \text{ kJ mol}^{-1}$				
	1.000040	$\frac{C_1 - A}{2 407000}$	0 505175	<u> </u>	1.000040	$\frac{C_1 - A}{2 + 4070}$	0 505155	
	1.820646	2.497868	0.585175		1.820646	2.497808	0.585175	
	1.204080	0.742465 1 141174	0.074294 0.072076		1.204000	0.742400 1 1 4 1 1 7 4	0.074294 0.072076	
	0.032870	1.1411/4 1.007500	-U.U/39/0 0 000000		0.032870	1.1411/4 1.997590	0.000000	
	-1.004400 0.460046	1.00/009	-0.00000000000000000000000000000000000		-1.004400 0.460946	1.00/009 9.810200	-0.00000000000000000000000000000000000	
ם	2.400240 1.600941	2.019309 0.107006	-0.774170	ם   ת	2.400240 1.600941	2.019309 0.107006	-0.774170	
	1.030241	2 500444	-0.233942	и Ц	1.030241	2 500444	-0.200942	
ч Н	-1.294001 _1.774537	2.000444 1 830840	-0.940337 0 796578		-1.294001 _1.774537	2.000444	-0.940337 0 796578	
1 11	-1.114001	1.003043	0.120010		-1.114001	1.003043	0.120010	

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC <sub>3</sub>	
Molecules in the Circumstellar Medium of Carbon Stars	171

	[i	$6e] \rightarrow [i9a]$			[i	$[i6f] \rightarrow [i9b]$		
					C			
	-256	$3.0 \text{ kJ mol}^{-1}$			-258	$3.8 \text{ kJ mol}^{-1}$		
		$C_1 - A$				$C_1-^1A$		
Si	-1.820646	2.497868	0.585175	Si	-1.820646	2.497868	0.585175	
C	-1.254585	0.742483	0.074294	C	-1.254585	0.742483	0.074294	
C	-0.032876	1.141174	-0.073976	C	-0.032876	1.141174	-0.073976	
C	1.054450	1.887589	-0.083800	C	1.054450	1.887589	-0.083800	
D	-2.460246	2.819309	-0.774176	H	-2.460246	2.819309	-0.774176	
H	-1.690241	-0.197006	-0.235942	D	-1.690241	-0.197006	-0.235942	
D	1.294651	2.500444	-0.946357	D	1.294651	2.500444	-0.946357	
D	1.774537	1.839849	0.726578	D	1.774537	1.839849	0.726578	
	[i	6g]→[i9d]		[i6h]→[i9c]				
	-255	$5.8 \text{ kJ mol}^{-1}$			-255	$5.7 \text{ kJ mol}^{-1}$		
		$C_1 - A$				$C_1^{-1}A$		
Si	-1.820646	2.497868	0.585175	Si	-1.820646	2.497868	0.585175	
	-1.254585	0.742483	0.074294		-1.254585	0.742483	0.074294	
C	-0.032876	1.141174	-0.073976	C	-0.032876	1.141174	-0.073976	
	1.054450	1.887589	-0.083800		1.054450	1.887589	-0.083800	
D	-2.460246	2.819309	-0.774176	D	-2.460246	2.819309	-0.774176	
D	-1.690241	-0.197006	-0.235942		-1.690241	-0.197006	-0.235942	
D	1.294651	2.500444	-0.946357	H	1.294651	2.500444	-0.946357	
$\mid$ H	1.774537	1.839849	0.726578	D	1.774537	1.839849	0.726578	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 172

	[i6	$ba] \rightarrow [i10b]$			[i6	$bb] \rightarrow [i10a]$		
	-155	$5.7 \text{ kJ mol}^{-1}$			-158	$3.9 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	1.108893	2.428844	0.896755	Si	1.108893	2.428844	0.896755	
C	1.254157	0.804122	-0.190514	C	1.254157	0.804122	-0.190514	
C	-0.047399	0.766331	0.066469	C	-0.047399	0.766331	0.066469	
C	-0.677368	2.066019	0.073662	C	-0.677368	2.066019	0.073662	
H	2.054356	0.150676	-0.497980	D	2.054356	0.150676	-0.497980	
D	1.608219	2.139909	-0.662094	Н	1.608219	2.139909	-0.662094	
D	-0.645862	2.587746	-0.883451	D	-0.645862	2.587746	-0.883451	
D	-1.601289	2.227878	0.621898	D	-1.601289	2.227878	0.621898	
	[i6	$bc] \rightarrow [i10d]$		$[i6d] \rightarrow [i10c]$				
	-155	$6.4 \text{ kJ mol}^{-1}$			-155	$\frac{1}{1}$ $\frac{1}{1}$ $\frac{1}{1}$		
	1 100000	$C_1^{-1}A$			1 100000	$C_1^{-1}A$		
Si	1.108893	2.428844	0.896755	Si	1.108893	2.428844	0.896755	
	1.254157	0.804122	-0.190514		1.254157	0.804122	-0.190514	
	-0.047399	0.700331	0.000469		-0.047399	0.700331	0.000469	
	-0.077308	2.000019	0.073662		-0.077308	2.000019	0.073062	
	2.034330	0.100070	-0.497980	ם	2.034330 1.609210	0.100070	-0.497980	
	1.000219	2.139909	-0.002094		1.000219	2.139909	-0.002094	
	-0.045802	2.38/(40	-0.883431		-0.045802	2.38/(40	-0.883431	
П	-1.001289	2.221010	0.021898	ע ן	-1.001289	2.221010	0.021898	

Appendix A. A.	dditional Information:	Modeling of Gas	Phase Formation	of $c$ -SiC <sub>3</sub>	
Molecules in the	e Circumstellar Mediun	n of Carbon Stars		1'	73

	[i6	$\delta e] \rightarrow [i10d]$		[i6f]→[i10c]				
	-155	$5.7 \text{ kJ} \text{ mol}^{-1}$			-158	$3.9 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-1.108893	2.428844	0.896755	Si	-1.108893	2.428844	0.896755	
C	-1.254157	0.804122	-0.190514	C	-1.254157	0.804122	-0.190514	
C	0.047399	0.766331	0.066469	C	0.047399	0.766331	0.066469	
C	0.677368	2.066019	0.073662	C	0.677368	2.066019	0.073662	
H	-2.054356	0.150676	-0.497980	D	-2.054356	0.150676	-0.497980	
D	-1.608219	2.139909	-0.662094	H	-1.608219	2.139909	-0.662094	
D	0.645862	2.587746	-0.883451	D	0.645862	2.587746	-0.883451	
D	1.601289	2.227878	0.621898	D	1.601289	2.227878	0.621898	
	[i6	$\delta g] \rightarrow [i10b]$		$[i6h] \rightarrow [i10a]$				
	-155	$5.4 \text{ kJ mol}^{-1}$		$-155.7 \text{ kJ mol}^{-1}$				
		$C_1 - A$		-		$C_1 - A$		
Si	-1.108893	2.428844	0.896755	Si	-1.108893	2.428844	0.896755	
C	-1.254157	0.804122	-0.190514	C	-1.254157	0.804122	-0.190514	
	0.047399	0.766331	0.066469		0.047399	0.766331	0.066469	
	0.677368	2.066019	0.073662		0.677368	2.066019	0.073662	
D	-2.054356	0.150676	-0.497980	D	-2.054356	0.150676	-0.497980	
D	-1.608219	2.139909	-0.662094	D	-1.608219	2.139909	-0.662094	
D	0.645862	2.587746	-0.883451	H	0.645862	2.587746	-0.883451	
H	1.601289	2.227878	0.621898	D	1.601289	2.227878	0.621898	

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	174

[i6a]→[i12a]					[i6b]→[i12b]			
	-238	$3.3 \text{ kJ mol}^{-1}$			-241	$1.1 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.852609	1.635684	0.581065	Si	0.852609	1.635684	0.581065	
C	-0.565860	0.423002	0.059391	C	-0.565860	0.423002	0.059391	
C	0.001882	-0.919555	0.000234	C	0.001882	-0.919555	0.000234	
C	0.544992	-1.989473	-0.037835	C	0.544992	-1.989473	-0.037835	
D	0.555684	2.682867	-0.493235	H	0.555684	2.682867	-0.493235	
D	-1.048992	0.683920	-0.883382	D	-1.048992	0.683920	-0.883382	
D	-1.329382	0.463880	0.848535	D	-1.329382	0.463880	0.848535	
H	1.001958	-2.950630	-0.071481	D	1.001958	-2.950630	-0.071481	
	[i	6c]→[i12d]		[i6d]→[i12c]				
	-237	$7.7 \text{ kJ mol}^{-1}$			-237	$7.4 \text{ kJ mol}^{-1}$		
	0.050000	$\frac{C_1 - A}{1 \cos 2 \cos 4}$	0 501005		0.050000	$\frac{C_1 - A}{1 \cos 2 \cos 4}$	0 501005	
	0.852609	1.635684	0.581065		0.852609	1.635684	0.581065	
	-0.505800	0.423002	0.059391		-0.303800	0.423002	0.059391	
	0.001882	-0.919555	0.000234		0.001882	-0.919555	0.000234	
	0.544992	-1.9894/3	-0.037833		0.544992	-1.9894/3	-U.U37833	
	0.000084	2.082807	-0.493233		0.000084	2.082807	-0.493233	
	-1.048992	0.083920	-U.883382		-1.048992	0.083920	-U.883382	
	-1.329382	0.403880	0.848535		-1.329382	0.403880	0.848535	
D	1.001958	-2.950630	-0.071481	D	1.001958	-2.950630	-0.071481	

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC	1 '3
Molecules in the Circumstellar Medium of Carbon Stars	175

	[i	6e]→[i12a]		[i6f]→[i12b]				
			D					
	-238	$3.3 \text{ kJ mol}^{-1}$			-241	$1.1 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-0.852609	1.635684	0.581065	Si	-0.852609	1.635684	0.581065	
C	0.565860	0.423002	0.059391	C	0.565860	0.423002	0.059391	
C	-0.001882	-0.919555	0.000234	C	-0.001882	-0.919555	0.000234	
C	-0.544992	-1.989473	-0.037835	C	-0.544992	-1.989473	-0.037835	
D	-0.555684	2.682867	-0.493235	H	-0.555684	2.682867	-0.493235	
D	1.048992	0.683920	-0.883382	D	1.048992	0.683920	-0.883382	
D	1.329382	0.463880	0.848535	D	1.329382	0.463880	0.848535	
H	-1.001958	-2.950630	-0.071481	D	-1.001958	-2.950630	-0.071481	
	[i	$6g] \rightarrow [i12c]$		[i6h]→[i12d]				
	-237	$7.7 \text{ kJ mol}^{-1}$			-237	7.4 kJ mol <sup>-1</sup>		
		$C_1 - A$				$C_1^{-1}A$		
Si	-0.852609	1.635684	0.581065	Si	-0.852609	1.635684	0.581065	
C	0.565860	0.423002	0.059391	C	0.565860	0.423002	0.059391	
C	-0.001882	-0.919555	0.000234	C	-0.001882	-0.919555	0.000234	
	-0.544992	-1.989473	-0.037835		-0.544992	-1.989473	-0.037835	
D	-0.555684	2.682867	-0.493235	D	-0.555684	2.682867	-0.493235	
D	1.048992	0.683920	-0.883382	H	1.048992	0.683920	-0.883382	
H	1.329382	0.463880	0.848535	D	1.329382	0.463880	0.848535	
D	-1.001958	-2.950630	-0.071481	D	-1.001958	-2.950630	-0.071481	

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	176

	[i'	7a]→[i10a]		[i7b]→[i10b]				
	-276	$3.6 \text{ kJ mol}^{-1}$			-276	$3.4 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-1.965407	2.054224	0.949023	Si	-1.965407	2.054224	0.949023	
C	-2.941031	0.741834	-0.252682	C	-2.941031	0.741834	-0.252682	
C	-1.529983	0.959789	-0.397630	C	-1.529983	0.959789	-0.397630	
C	-0.519325	0.330323	0.274811	C	-0.519325	0.330323	0.274811	
D	-3.585322	1.284613	-0.938932	H	-3.585322	1.284613	-0.938932	
H	-3.363470	-0.192898	0.123519	D	-3.363470	-0.192898	0.123519	
D	-0.582867	-0.703830	0.623728	D	-0.582867	-0.703830	0.623728	
D	0.457528	0.798760	0.345200	D	0.457528	0.798760	0.345200	
	[i'	7c]→[i10d]			[i'	7d]→[i10c]		
	-276	$\frac{5.4 \text{ kJ mol}^{-1}}{2}$			-276	$5.7 \text{ kJ mol}^{-1}$		
	1.005.105	$C_1^{-1}A$	0.040000		1.005.05	$C_1^{-1}A$	0.040000	
	-1.965407	2.054224	0.949023		-1.965407	2.054224	0.949023	
	-2.941031	0.741834	-0.252682		-2.941031	0.741834	-0.252682	
	-1.529983	0.959789	-0.397030		-1.529983	0.959789	-0.397030	
	-U.319323	0.330323	0.274811		-0.319323 2 505399	0.330323	0.274811	
	-3.303322 2.262470	1.204013	-U.YJ8YJ2 0 192510		-0.0000 <u>7</u> 2 2.262470	1.204013	-U.YJ8YJZ 0 192510	
	-3.303470 0 509067	-0.192090 0 709990	0.120019 0.609700	ע   נו	-3.303470 0 589967	-0.192090 0 702020	0.120019 0.609700	
и 10	-0.002007	-0.703030 0.709760	0.023728		-0.002007	-U.1U303U 0.709760	0.020720	
11	0.407020	0.190100	0.040200		0.40/020	0.190100	0.040200	

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	177

	[j	[i9d] →[i9d]		[i8b]→[i9a]			
	-47	$1.9 \text{ kJ mol}^{-1}$			-47	$1.6 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	1.041061	0.917892	0.436916	Si	1.041061	0.917892	0.436916
C	0.760892	-0.738193	-0.161392	C	0.760892	-0.738193	-0.161392
C	-0.609128	-0.750344	-0.306514	C	-0.609128	-0.750344	-0.306514
C	-1.670823	-0.094074	-0.212414	C	-1.670823	-0.094074	-0.212414
D	2.501462	1.265763	0.648959	D	2.501462	1.265763	0.648959
D	-0.818889	0.943557	0.334551	D	-0.818889	0.943557	0.334551
D	1.411713	-1.596634	-0.075037	H	1.411713	-1.596634	-0.075037
H	-2.618347	0.185964	-0.624789	D	-2.618347	0.185964	-0.624789
	[j	i8c]→[i9b]			[]	i8d]→[i9c]	
	-50	$.0 \text{ kJ mol}^{-1}$			-51	$.4 \text{ kJ mol}^{-1}$	
		$C_1^{-1}A$				$C_1-A$	
Si	1.041061	0.917892	0.436916	Si	1.041061	0.917892	0.436916
C	0.760892	-0.738193	-0.161392	C	0.760892	-0.738193	-0.161392
	-0.609128	-0.750344	-0.306514		-0.609128	-0.750344	-0.306514
	-1.670823	-0.094074	-0.212414		-1.670823	-0.094074	-0.212414
H	2.501462	1.265763	0.648959	D	2.501462	1.265763	0.648959
D	-0.818889	0.943557	0.334551	H	-0.818889	0.943557	0.334551
D	1.411713	-1.596634	-0.075037	D	1.411713	-1.596634	-0.075037
D	-2.618347	0.185964	-0.624789	D	-2.618347	0.185964	-0.624789

Appendix A. Addi	tional Information: M	Iodeling of Gas I	Phase Formation	of $c$ -SiC $_3$	
Molecules in the Ci	rcumstellar Medium o	of Carbon Stars		1	78

	[i8	$a] \rightarrow [i12a]$			[i8	$b \rightarrow [i12c]$	
	-108	$3.7 \text{ kJ mol}^{-1}$			-108	$3.2 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-2.521663	2.463165	1.136938	Si	-2.521663	2.463165	1.136938
C	-0.857503	2.002660	0.565118	C	-0.857503	2.002660	0.565118
C	-0.462648	0.671903	0.277186	C	-0.462648	0.671903	0.277186
C	-0.127979	-0.463490	0.061318	C	-0.127979	-0.463490	0.061318
D	-0.053439	2.725894	0.410733	H	-0.053439	2.725894	0.410733
H	0.166755	-1.470185	-0.125512	D	0.166755	-1.470185	-0.125512
D	-2.334378	3.957638	1.028947	D	-2.334378	3.957638	1.028947
D	-2.405675	1.983846	-0.345689	D	-2.405675	1.983846	-0.345689
	[i8	$Bc] \rightarrow [i12b]$			[i8	8d]→[i12d]	
	-11(	$\frac{1.5 \text{ kJ mol}^{-1}}{2}$			-111	$1.6 \text{ kJ mol}^{-1}$	
	2 221 222	$C_1$ - <sup>1</sup> A	1 1 2 6 2 2 2		2 501000	$C_1 - A$	1 1 2 2 2 2 2
	-2.521663	2.463165	1.136938		-2.521663	2.463165	1.136938
	-0.857503	2.002660	0.505118		-0.857503	2.002660	0.565118
	-0.462648	0.671903	0.277186		-0.462648	0.671903	0.277186
	-0.127979	-0.463490	0.061318		-0.127979	-0.463490	0.061318
	-0.053439	2.725894	0.410733		-0.053439	2.725894	0.410733
	0.100755	-1.4/0185	-0.125512		0.100755	-1.4/0185	-0.125512
	-2.334378	3.957638	1.028947		-2.334378	3.957638	1.028947
$\mid D$	-2.405675	1.983846	-0.345689	H	-2.405675	1.983846	-0.345689

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 179

[i9a]→ <b>p1a</b>			[i9b]→p1a				
							)
	-84	$.8 \text{ kJ mol}^{-1}$			-83	$.8 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-0.744770	2.157892	1.133843	Si	-0.744770	2.157892	1.133843
C	-0.117447	0.739329	0.301978	C	-0.117447	0.739329	0.301978
C	0.508868	-0.265289	-0.184605	C	0.508868	-0.265289	-0.184605
C	1.137912	-1.291212	-0.706981	C	1.137912	-1.291212	-0.706981
D	2.192710	-1.451493	-0.511820	D	2.192710	-1.451493	-0.511820
D	0.620937	-2.000203	-1.344227	D	0.620937	-2.000203	-1.344227
H	-1.537155	0.742156	0.712436	D	-1.537155	0.742156	0.712436
D	-2.215299	1.527879	0.692674	H	-2.215299	1.527879	0.692674
	[j	$[9c] \rightarrow p1b$			[j	9d]→ <b>p1b</b>	
	-81	$.1 \text{ kJ mol}^{-1}$			-81	$.1 \text{ kJ mol}^{-1}$	
		$C_1 - A$				$C_1 - A$	
Si	-0.744770	2.157892	1.133843	Si	-0.744770	2.157892	1.133843
	-0.117447	0.739329	0.301978		-0.117447	0.739329	0.301978
C	0.508868	-0.265289	-0.184605	C	0.508868	-0.265289	-0.184605
C	1.137912	-1.291212	-0.706981	$\mid C$	1.137912	-1.291212	-0.706981
D	2.192710	-1.451493	-0.511820	H	2.192710	-1.451493	-0.511820
H	0.620937	-2.000203	-1.344227	D	0.620937	-2.000203	-1.344227
D	-1.537155	0.742156	0.712436	D	-1.537155	0.742156	0.712436
D	-2.215299	1.527879	0.692674	D	-2.215299	1.527879	0.692674

Appendix A. A	Additional Information:	Modeling of Gas Phase	Formation of $c$ -SiC <sub>3</sub>	
Molecules in th	e Circumstellar Medium	n of Carbon Stars		180

	[	i1a] $\rightarrow$ <b>p2a</b>			[j	$i1b] \rightarrow p2a$	
	10	$.8 \text{ kJ mol}^{-1}$			11	$.0 \text{ kJ mol}^{-1}$	
		$C_s-^1A'$				$C_s-^1A'$	
Si	0.528665	2.400596	-0.746415	Si	0.528665	2.400596	-0.746415
C	0.638969	-0.741790	0.047537	C	0.638969	-0.741790	0.047537
C	0.334217	0.578629	-0.236986	C	0.334217	0.578629	-0.236986
C	-0.782728	1.293996	-0.211915	C	-0.782728	1.293996	-0.211915
D	1.142678	-0.981263	0.980498	D	1.142678	-0.981263	0.980498
D	0.810919	-1.435189	-0.771673	D	0.810919	-1.435189	-0.771673
H	-1.559747	-0.175915	0.315690	D	-1.559747	-0.175915	0.315690
D	-1.065292	-0.900941	0.410427	H	-1.065292	-0.900941	0.410427
	[j	1b]→ <b>p2b</b>			[1	$i1b] \rightarrow p2c$	
	13	$\frac{.9 \text{ kJ mol}^{-1}}{2}$			13	$\frac{.9 \text{ kJ mol}^{-1}}{2}$	
		$C_1$ - <sup>1</sup> A				$C_1$ - <sup>1</sup> A	
Si	0.528665	2.400596	-0.746415	Si	0.528665	2.400596	-0.746415
	0.638969	-0.741790	0.047537		0.638969	-0.741790	0.047537
	0.334217	0.578629	-0.236986		0.334217	0.578629	-0.236986
	-0.782728	1.293996	-0.211915		-0.782728	1.293996	-0.211915
	1.142678	-0.981263	0.980498		1.142678	-0.981263	0.980498
	0.810919	-1.435189	-0.771673		0.810919	-1.435189	-0.771673
D	-1.559747	-0.175915	0.315690	D	-1.559747	-0.175915	0.315690
D	-1.065292	-0.900941	0.410427	D	-1.065292	-0.900941	0.410427

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$		
Molecules in the Circumstellar Medium of Carbon Stars	181	

	[i	$[i6a] \rightarrow p2a$			[i6b]→ <b>p2a</b>			
	-43	$.5 \text{ kJ mol}^{-1}$			-41	$.9 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	1.178923	2.205371	0.803390	Si	1.178923	2.205371	0.803390	
C	-0.641206	2.137118	-0.050763	C	-0.641206	2.137118	-0.050763	
C	-0.130702	0.814244	0.053579	C	-0.130702	0.814244	0.053579	
C	0.888991	0.086553	0.200705	C	0.888991	0.086553	0.200705	
D	-0.835280	2.504167	-1.054512	D	-0.835280	2.504167	-1.054512	
D	-1.380706	2.457777	0.683302	D	-1.380706	2.457777	0.683302	
H	1.725043	1.088245	-0.475804	D	1.725043	1.088245	-0.475804	
D	1.923344	2.086626	-0.687100	H	1.923344	2.086626	-0.687100	
	[i	$6c] \rightarrow p2b$			[i	$6d] \rightarrow p2c$		
	-39	$\frac{.4 \text{ kJ mol}^{-1}}{2}$			-39	$\frac{.4 \text{ kJ mol}^{-1}}{2}$		
	4 4 10 0 0 0 0 0	C <sub>1</sub> - <sup>1</sup> A	0.0000000		4 4 10 0 0 0 0	C <sub>1</sub> - <sup>1</sup> A	0.000000	
Si	1.178923	2.205371	0.803390	Si	1.178923	2.205371	0.803390	
	-0.641206	2.137118	-0.050763		-0.641206	2.137118	-0.050763	
	-0.130702	0.814244	0.053579		-0.130702	0.814244	0.053579	
	0.888991	0.086553	0.200705		0.888991	0.086553	0.200705	
	-0.835280	2.504167	-1.054512		-0.835280	2.504167	-1.054512	
	-1.380706	2.457777	0.683302		-1.380706	2.457777	0.683302	
D	1.725043	1.088245	-0.475804	D	1.725043	1.088245	-0.475804	
D	1.923344	2.086626	-0.687100	D	1.923344	2.086626	-0.687100	

Appendix A. Additional	Information: Modeling	of Gas Phase	Formation	of $c$ -SiC <sub>3</sub>	
Molecules in the Circumst	tellar Medium of Carbo	on Stars			182

	[j	$[6e] \rightarrow p2a$			[i	$[6f] \rightarrow p2a$	
	-43	$.5 \text{ kJ mol}^{-1}$			-41	$.9 \text{ kJ mol}^{-1}$	
		$C_1-^1A$				$C_1-^1A$	
Si	-1.178923	2.205371	0.803390	Si	-1.178923	2.205371	0.803390
C	0.641206	2.137118	-0.050763	C	0.641206	2.137118	-0.050763
C	0.130702	0.814244	0.053579	C	0.130702	0.814244	0.053579
C	-0.888991	0.086553	0.200705	C	-0.888991	0.086553	0.200705
D	0.835280	2.504167	-1.054512	D	0.835280	2.504167	-1.054512
D	1.380706	2.457777	0.683302	D	1.380706	2.457777	0.683302
H	-1.725043	1.088245	-0.475804	D	-1.725043	1.088245	-0.475804
D	-1.923344	2.086626	-0.687100	H	-1.923344	2.086626	-0.687100
	[j	$[6g] \rightarrow p2c$			[i	$6h] \rightarrow p2b$	
	-39	$.4 \text{ kJ mol}^{-1}$			-39	$.4 \text{ kJ mol}^{-1}$	
		$C_1 - A$				$C_1^{-1}A$	
Si	-1.178923	2.205371	0.803390	Si	-1.178923	2.205371	0.803390
C	0.641206	2.137118	-0.050763	C	0.641206	2.137118	-0.050763
C	0.130702	0.814244	0.053579	C	0.130702	0.814244	0.053579
	-0.888991	0.086553	0.200705		-0.888991	0.086553	0.200705
D	0.835280	2.504167	-1.054512	H	0.835280	2.504167	-1.054512
H	1.380706	2.457777	0.683302	D	1.380706	2.457777	0.683302
D	-1.725043	1.088245	-0.475804	D	-1.725043	1.088245	-0.475804
D	-1.923344	2.086626	-0.687100	D	-1.923344	2.086626	-0.687100

Appendix A. Additional Information: Modeling	of Gas Phase Formation of $c$ -SiC <sub>3</sub>
Molecules in the Circumstellar Medium of Carbo	n Stars 18

[i13a]→ <b>p2a</b>				$[i13b] \rightarrow p2a$				
	-42	$.3 \text{ kJ mol}^{-1}$			-43	$.1 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-0.830773	0.883817	0.953003	Si	-0.830773	0.883817	0.953003	
C	-1.811287	-0.696687	0.301549	C	-1.811287	-0.696687	0.301549	
C	-0.590983	-0.923847	0.102221	C	-0.590983	-0.923847	0.102221	
C	0.749568	-0.461799	-0.040845	C	0.749568	-0.461799	-0.040845	
D	1.185475	-0.546441	-1.038520	D	1.185475	-0.546441	-1.038520	
D	1.455315	-0.696271	0.749183	D	1.455315	-0.696271	0.749183	
D	0.271987	0.942873	-0.437740	H	0.271987	0.942873	-0.437740	
H	-0.545547	1.525644	-0.576927	D	-0.545547	1.525644	-0.576927	
	[1]	13b]→ <b>p2b</b>			[i]	$13b] \rightarrow p2c$		
	-39	$.9 \text{ kJ mol}^{-1}$			-39	$.7 \text{ kJ mol}^{-1}$		
		$C_1 - A$				$C_1^{-1}A$		
Si	-0.830773	0.883817	0.953003	Si	-0.830773	0.883817	0.953003	
C	-1.811287	-0.696687	0.301549		-1.811287	-0.696687	0.301549	
	-0.590983	-0.923847	0.102221		-0.590983	-0.923847	0.102221	
	0.749568	-0.461799	-0.040845		0.749568	-0.461799	-0.040845	
	1.185475	-0.546441	-1.038520		1.185475	-0.546441	-1.038520	
	1.455315	-0.696271	0.749183		1.455315	-0.696271	0.749183	
D	0.271987	0.942873	-0.437740	D	0.271987	0.942873	-0.437740	
D	-0.545547	1.525644	-0.576927	D	-0.545547	1.525644	-0.576927	

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	184

[i12a]→ <b>p3a</b>				$[i12b] \rightarrow p3c$				
	-64	$.2 \text{ kJ mol}^{-1}$			-66	$3 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-1.151997	1.718947	0.728726	Si	-1.151997	1.718947	0.728726	
C	-0.793552	-0.036891	0.485201	C	-0.793552	-0.036891	0.485201	
C	0.356363	-0.744575	0.018188	C	0.356363	-0.744575	0.018188	
C	1.337204	-1.326892	-0.364743	C	1.337204	-1.326892	-0.364743	
D	-1.597641	-0.704860	0.800559	D	-1.597641	-0.704860	0.800559	
H	2.205962	-1.845464	-0.696218	D	2.205962	-1.845464	-0.696218	
D	-0.463884	1.036474	-0.582223	D	-0.463884	1.036474	-0.582223	
D	0.061553	1.931814	-0.358027	H	0.061553	1.931814	-0.358027	
	[i	12c]→ <b>p3c</b>			[1]	12d]→ <b>p3b</b>		
	-66	$\frac{1}{1000}$ kJ mol <sup>-1</sup>			-63	$\frac{1}{100}$ kJ mol <sup>-1</sup>		
		$C_1-A$				$C_1-A$		
Si	-1.151997	1.718947	0.728726	Si	-1.151997	1.718947	0.728726	
C	-0.793552	-0.036891	0.485201	C	-0.793552	-0.036891	0.485201	
C	0.356363	-0.744575	0.018188	C	0.356363	-0.744575	0.018188	
	1.337204	-1.326892	-0.364743		1.337204	-1.326892	-0.364743	
D	-1.597641	-0.704860	0.800559		-1.597641	-0.704860	0.800559	
	2.205962	-1.845464	-0.696218		2.205962	-1.845464	-0.696218	
H	-0.463884	1.036474	-0.582223	D	-0.463884	1.036474	-0.582223	
D	0.061553	1.931814	-0.358027	D	0.061553	1.931814	-0.358027	

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC <sub>3</sub>	
Molecules in the Circumstellar Medium of Carbon Stars	

[i12a]→ <b>p3a</b> ′					[i12b]→ <b>p3c</b> ′			
	-69	$.9 \text{ kJ mol}^{-1}$			-71	$.9 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-1.555112	0.642765	0.911988	Si	-1.555112	0.642765	0.911988	
C	0.049390	0.418870	0.114001	C	0.049390	0.418870	0.114001	
C	0.581034	-0.893905	-0.073288	C	0.581034	-0.893905	-0.073288	
C	1.006592	-2.011224	-0.207659	C	1.006592	-2.011224	-0.207659	
D	0.726733	1.208649	-0.209397	D	0.726733	1.208649	-0.209397	
H	1.387736	-2.997804	-0.330653	D	1.387736	-2.997804	-0.330653	
D	-1.435106	2.141508	0.235326	H	-1.435106	2.141508	0.235326	
D	-1.002771	1.445970	-0.426216	D	-1.002771	1.445970	-0.426216	
	[i]	$12c] \rightarrow p3b'$			[i]	$12d] \rightarrow p3c'$		
	-69	$2 \text{ kJ mol}^{-1}$		$-72.6 \text{ kJ mol}^{-1}$				
		$C_1 - A$				$C_1 - A$		
Si	-1.555112	0.642765	0.911988	Si	-1.555112	0.642765	0.911988	
C	0.049390	0.418870	0.114001	C	0.049390	0.418870	0.114001	
C	0.581034	-0.893905	-0.073288	C	0.581034	-0.893905	-0.073288	
	1.006592	-2.011224	-0.207659		1.006592	-2.011224	-0.207659	
H	0.726733	1.208649	-0.209397	D	0.726733	1.208649	-0.209397	
D	1.387736	-2.997804	-0.330653	D	1.387736	-2.997804	-0.330653	
D	-1.435106	2.141508	0.235326	D	-1.435106	2.141508	0.235326	
D	-1.002771	1.445970	-0.426216	H	-1.002771	1.445970	-0.426216	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 186

[i6a]→ <b>p4a</b>				$[i6b] \rightarrow p4b$				
	-107	$7.7 \text{ kJ mol}^{-1}$			-109	$9.7 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	0.831125	0.640403	0.777373	Si	0.831125	0.640403	0.777373	
C	-1.185225	0.266693	0.157410	C	-1.185225	0.266693	0.157410	
C	-0.415827	-0.797532	0.102736	C	-0.415827	-0.797532	0.102736	
C	0.854625	-1.135711	0.195070	C	0.854625	-1.135711	0.195070	
D	-2.051331	0.445181	0.777573	D	-2.051331	0.445181	0.777573	
H	1.482402	-1.878498	-0.264555	D	1.482402	-1.878498	-0.264555	
D	-0.203090	1.162708	-0.807825	D	-0.203090	1.162708	-0.807825	
D	0.676107	1.295027	-0.933420	Н	0.676107	1.295027	-0.933420	
	[j	$i6c] \rightarrow p4a$		$[i6d] \rightarrow p4b$				
-107.7 kJ mol <sup>-1</sup>					-110.6 kJ mol <sup>-1</sup>			
		$C_1 - A$				$C_1 - A$		
Si	0.831125	0.640403	0.777373	Si	0.831125	0.640403	0.777373	
C	-1.185225	0.266693	0.157410	C	-1.185225	0.266693	0.157410	
	-0.415827	-0.797532	0.102736		-0.415827	-0.797532	0.102736	
	0.854625	-1.135711	0.195070		0.854625	-1.135711	0.195070	
H	-2.051331	0.445181	0.777573	D	-2.051331	0.445181	0.777573	
D	1.482402	-1.878498	-0.264555	D	1.482402	-1.878498	-0.264555	
D	-0.203090	1.162708	-0.807825	H	-0.203090	1.162708	-0.807825	
D	0.676107	1.295027	-0.933420	D	0.676107	1.295027	-0.933420	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

$[i6e] \rightarrow p4a$				$[i6f] \rightarrow p4b$				
	-107	$7.7 \text{ kJ mol}^{-1}$			-109	$9.7 \text{ kJ mol}^{-1}$		
		$C_1-^1A$				$C_1-^1A$		
Si	-0.831125	0.640403	0.777373	Si	-0.831125	0.640403	0.777373	
C	1.185225	0.266693	0.157410	C	1.185225	0.266693	0.157410	
C	0.415827	-0.797532	0.102736	C	0.415827	-0.797532	0.102736	
C	-0.854625	-1.135711	0.195070	C	-0.854625	-1.135711	0.195070	
D	2.051331	0.445181	0.777573	D	2.051331	0.445181	0.777573	
H	-1.482402	-1.878498	-0.264555	D	-1.482402	-1.878498	-0.264555	
D	0.203090	1.162708	-0.807825	D	0.203090	1.162708	-0.807825	
D	-0.676107	1.295027	-0.933420	Н	-0.676107	1.295027	-0.933420	
	[j	$[6g] \rightarrow p4a$		$[i6h] \rightarrow p4b$				
	-107	$7.7 \text{ kJ mol}^{-1}$			-11(	$\frac{0.6 \text{ kJ mol}^{-1}}{2}$		
		$C_1$ - <sup>1</sup> A		<u> </u>		$C_1$ - <sup>1</sup> A		
Si	-0.831125	0.640403	0.777373	Si	-0.831125	0.640403	0.777373	
	1.185225	0.266693	0.157410		1.185225	0.266693	0.157410	
	0.415827	-0.797532	0.102736		0.415827	-0.797532	0.102736	
	-0.854625	-1.135/11	0.195070		-0.854625	-1.135/11	0.195070	
	2.051331	0.445181	0.777573		2.051331	0.445181	0.111513	
	-1.482402	-1.878498	-0.204000		-1.482402	-1.878498	-0.204333	
	0.203090	1.102708	-0.807825		0.203090	1.102708	-0.807825	
D	-0.070107	1.295027	-0.933420	ע ן	-0.070107	1.295027	-0.933420	

## Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC3Molecules in the Circumstellar Medium of Carbon Stars188

Table A.3: Exit transition state structures for the emission of  $H_2/HD/D_2$ , principle moments of inertia (AMU  $a_0^2$ ), and angles of emission relative to the principal axes.



Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars



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Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars



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Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars







Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars



## Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC3Molecules in the Circumstellar Medium of Carbon Stars195

Table A.4: Products relevant to the reaction of electronically excited atomic silicon with methylacetylene and allene to form atomic hydrogen and SiC<sub>3</sub>H<sub>3</sub>. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

n1					n2			
Sec.								
	-10	$.0 \text{ kJ mol}^{-1}$			3.	$9 \text{ kJ mol}^{-1}$		
		$C_{3v}$ -2 $A_1$				$C_s$ – $^2A'$		
Si	2.256938	0.386814	-0.000054	Si	-1.648670	1.864824	-0.639004	
C	0.452597	0.445576	0.001238	C	-3.059962	0.865119	-0.096199	
C	-0.767469	0.480303	0.004880	C	-1.956341	0.083084	-0.283303	
C	-2.224034	0.526369	0.004637	C	-1.616436	-1.271962	-0.143675	
H	-2.643608	-0.348009	-0.498856	H	-2.264110	-1.960263	0.392886	
H	-2.561104	1.418906	-0.528946	H	-0.692145	-1.662759	-0.553724	
H	-2.617180	0.570276	1.023216	H	-4.087409	0.662816	0.185587	
		n3				n4		
					8			
$7.0 \text{ kJ mol}^{-1}$					$11.4 \text{ kJ mol}^{-1}$			
		$C_s$ – <sup>2</sup> A′				$C_1$ -2A		
Si	1.265159	2.162075	-0.444465	Si	1.208012	2.089231	-0.448538	
C	-0.607874	2.024234	-0.328499	C	-0.745936	2.116845	-0.093209	
	0.896654	0.500285	0.357941		0.910530	0.446506	0.329165	
C	-0.395785	0.642838	-0.176604	$\mid C$	-0.351043	0.718990	-0.018821	
H	-1.518838	2.509970	-0.653582	H	-1.431207	2.427755	-0.878519	
H	-1.060126	-0.160934	-0.504981	H	-0.968549	2.572982	0.868603	
H	1.374406	-0.420749	0.666057	H	1.419426	-0.479092	0.553088	

n5				n6			
					9	e ge	
	11.	5 kJ mol <sup>-1</sup>			24.	4 kJ mol <sup>-1</sup>	
		$C_s - A'$				$C_s - A'$	
Si	-0.357120	1.246352	-0.941397	Si	-1.573453	1.860263	-0.624254
C	-0.224659	-0.413295	0.358340		-3.053427	0.877832	-0.162194
C	-1.552969	-0.035456	0.272406	C	-2.014789	0.087378	-0.128936
	0.911547	-0.089182	-0.083811	C	-1.441504	-1.261433	0.109938
H	-2.226882	-0.652993	-0.312865	H	-2.222907	-1.968179	0.399486
H	-1.995601	0.484661	1.115990	H	-0.947918	-1.627458	-0.792862
H	1.959706	-0.318373	-0.071805	H	-0.690455	-1.217933	0.901573
$33.1 \text{ kJ mol}^{-1}$				$45.1 \text{ kJ mol}^{-1}$			
		$C_s - A'$				$C_s - A'$	
Si	-2.187581	1.634287	1.295559	Si	-2.654104	0.750695	0.461772
	0.051611	1.182614	-0.115557		0.054200	2.167105	0.127665
	1.314261	1.492673	-0.038176		1.204947	1.500096	-0.166927
	-1.232532	0.906849	-0.142409		-1.174983	1.525933	0.279596
	1.718281	2.306210	-0.630132		0.101347	3.249863	0.248999
	1.982839	0.942873	0.614082		2.148416	2.023473	-0.279824
	1 COF 400	0.188197	-0.870198	ΙĤ	1.196132	0.423506	-0 294711

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars

n9				n10				
	62.	$1 \text{ kJ mol}^{-1}$			72.	$1 \text{ kJ mol}^{-1}$		
		$C_1-^2A$				$C_s - A'$		
Si	-2.291802	2.588971	1.004979	Si	1.258456	2.565463	1.003571	
C	0.397110	1.552862	0.026512	C	-0.623501	1.956099	-0.193860	
C	1.610713	1.062007	-0.294224	C	0.616569	1.034679	0.382244	
C	-0.724745	2.021162	0.301926	C	-0.361129	0.485074	-0.277162	
H	-2.437580	3.782829	0.074313	H	-0.481565	2.545990	-1.097167	
H	2.447989	1.727500	-0.468446	H	-1.405388	2.291595	0.484508	
H	1.769210	-0.006354	-0.382326	H	-0.795541	-0.412938	-0.675421	
		n11				n12		
80.2 kJ mol <sup>-1</sup>				92.2 kJ mol <sup>-1</sup>				
<u> </u>	0 574000	$C_{\rm s}$ -"A'	0 100050	<u> </u>	1 000440	$\frac{U_1 - A}{2.007470}$	0 590517	
$\begin{vmatrix} 51 \\ C \end{vmatrix}$	-2.374003	2.832038	0.123250 0.101799	$\begin{vmatrix} 51 \\ C \end{vmatrix}$	1.209442	2.09/4/0	-0.030017	
	-1.238200	1.027070 1.156447	0.101728 0.042780		-0.098444	2.114228 0.510741	-0.082080	
	-2.004002 1 802585	0.264440	0.343709		0.011210	0.519741	0.400490 0.077667	
Ч	-1.092000 1 517944	0.204449	0.139070 0.331673	Ч	-0.300720	0.042040 9 116272	-0.011001 0.846667	
	-1.011244 2823004	-0.752291 0.055161	2 000084	11   Ц	-1.404080	2.440373	-0.040007	
	-2.020994 0.911910	0.900101	2.000904		-0.940030 1.011954	2.019009 0.191045	0.039044	
11	-0.211210	1.000921	0.441000	11	-1.011204	-0.121940	-0.042210	
Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

		n13		n14				
	۹ ۱							
	95.	$3 \text{ kJ mol}^{-1}$			96.	$2 \text{ kJ mol}^{-1}$		
		$C_1 - ^2A$				$C_s - A'$		
Si	-0.985382	2.986144	1.021451	Si	1.266892	2.417929	0.545124	
C	-1.112970	1.510143	-0.086940	C	-0.513422	2.108452	-0.345251	
C	-0.249537	0.428763	0.000560	C	1.220128	0.475897	0.012334	
C	0.510002	-0.511322	0.124227	C	0.113298	0.965102	-0.496432	
H	-2.442307	3.405788	0.874610	H	2.016603	2.947139	-0.672214	
H	-1.803839	1.493576	-0.929501	H	-1.434786	2.608535	-0.574972	
H	1.175602	-1.336872	0.224343	H	1.790465	-0.429923	0.089828	
		n15				n16		
	8							
	100	$1.5 \text{ kJ mol}^{-1}$			103	$.4 \text{ kJ mol}^{-1}$		
	1.10000.1	$C_s$ - <sup>2</sup> A'	0.0001.11		1 100000	<u>C1-4A</u>	0.0010.10	
Si	-1.190224	2.801061	0.833141	Si	1.193399	2.298845	0.391243	
	-1.168840	1.085298	0.732883		-0.704840	2.116248	-0.023680	
	-0.098546	0.202454	0.301977		1.220452	0.404007	-0.037734	
	U.81U3Z3 9.421EGE	-0.430003	-0.003099		0.009520	0.893913	-0.211514	
	-2.431505	5.480292 0.599491	1.319383		1.934239	J.120040	-0.013101	
	-2.086989	0.582431	1.037028		-1.095714	2.589399	-0.921264	
μН	1.012011	-1.03/1/4	-0.386591	μН	-1.404109	2.114579	0.807277	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

		n17		n18				
			D		<b>و</b>			
	108	$3.7 \text{ kJ} \text{ mol}^{-1}$			117	$1.5 \text{ kJ} \text{ mol}^{-1}$		
		$C_1 - A$				$C_{2v}$ – $^2A_1$		
Si	-2.464141	2.945399	0.274357	Si	-2.311951	2.797319	0.436930	
C	-1.402738	1.539986	0.017239	C	-1.347058	1.427980	0.366860	
C	-2.656658	1.177833	0.870404	C	-0.579950	0.339258	0.311177	
C	-1.964943	0.185393	-0.030971	C	0.125545	-0.661917	0.259982	
H	-1.409247	-0.612469	0.460416	H	-1.741994	4.155135	0.537058	
H	-2.744651	0.954547	1.933667	H	-3.785613	2.721926	0.402298	
H	-2.472267	-0.129035	-0.938397	H	0.738262	-1.531412	0.215476	
		n19				n20		
					162			
	142	$\frac{1.5 \text{ kJ mol}^2}{C^2 \Lambda'}$			102	$\frac{1.9 \text{ kJ mol}}{C}$		
Si	0.043742	$\frac{\sim_{\rm s}}{3154114}$	-0 166368	Si	-2.830058	$\frac{\sim_{\rm s}}{1.396545}$	0 563817	
	-0.724823	1.416847	-0.148213		0.020435	1.929632	0.104027	
Č	0.785745	1.448685	-0.065657	Ċ	1.226258	1.487302	-0.182331	
C	1.844977	0.723600	0.018307	C	-1.178048	1.088902	0.213717	
H	-1.148970	1.007794	-1.064990	H	-0.107872	2.998903	0.273833	
H	-1.249067	1.071020	0.742447	H	1.697174	0.537596	-0.400024	
H	2.139985	-0.317762	0.070963	H	-1.007629	0.017638	0.036540	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars 200

		n21		n22				
	221	$.1 \text{ kJ mol}^{-1}$		$303.6 \text{ kJ mol}^{-1}$				
		$C_s$ -2A'				$C_s$ – $^2A'$		
Si	-2.111813	2.725963	0.436207	Si	-0.952202	2.895186	1.046884	
C	-0.502023	2.028434	0.423706	C	-1.442045	1.273467	0.049434	
C	-0.381293	0.665379	0.337252	C	-0.404343	0.306776	0.063506	
C	-0.337798	-0.605402	0.255140	C	0.600189	-0.477548	0.149868	
H	-2.333611	4.183554	0.526784	H	-2.310437	3.508799	0.720181	
H	-3.302425	1.864084	0.355458	H	-1.683793	1.658231	-0.950735	
H	0.387891	2.652209	0.482682	H	-2.368134	0.945547	0.540814	

# Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC3Molecules in the Circumstellar Medium of Carbon Stars201

Table A.5: Products relevant to the reaction of electronically excited atomic silicon with methylacetylene and allene to form methyl radical and SiC<sub>2</sub>H. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

		$CH_3$		m1			
					-30	$\frac{10 \text{ kJ mol}^{-1}}{2 \text{ mol}^{-2}}$	
		$D_{3h}-A_{2}''$				$C_{\infty v}$ – <sup>2</sup> $\Pi$	
C	-0.951849	0.641383	0.030228	Si	-2.442919	0.826756	0.000000
H	0.066970	0.283748	-0.012939	C	-0.664513	1.191684	0.000000
H	-1.461269	0.932347	-0.877290	C	0.527499	1.436282	0.000000
H	-1.461264	0.707991	0.980921	H	1.572824	1.650730	0.000000
		m2				m3	
-5.0 kJ mol <sup>-1</sup>					144	$.5 \text{ kJ mol}^{-1}$	
		$C_s$ – $^2A'$				$C_s$ -2A'	
Si	-0.263369	1.916480	0.123746	Si	-0.098748	1.527158	-0.929747
C	1.306779	2.912456	-0.144538	C	-0.713212	0.293539	0.287855
C	1.559669	1.640124	-0.007267	C	0.539217	0.390597	0.368126
H	1.692598	3.909459	-0.305242	H	-0.233597	3.003472	-0.658751

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars



# Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$ 203Molecules in the Circumstellar Medium of Carbon Stars203

Table A.6: Reactants, products, and transition state structures relevant to the singlet SiC<sub>2</sub>H<sub>2</sub> system. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

	Reactants									
		p1		p2						
	0-0									
	0	kJ mol <sup>-1</sup>			15.	$7 \text{ kJ mol}^{-1}$				
		$C_{2v}$ - <sup>1</sup> $A_1$				$C_s^{-1}A'$				
Si	-2.667415	2.191276	0.000000	Si	-3.530077	0.607099	0.333008			
C	-0.985018	2.220382	0.000000	C	-2.115964	1.751749	-0.167586			
C	0.298405	2.242563	0.000000	C	-1.762971	0.449585	-0.634584			
C	1.607796	2.265203	0.000000	C	-2.980050	2.474570	0.443216			
H	2.142021	3.207971	0.000000	H	-0.954113	-0.065250	-0.122807			
H	2.174299	1.341470	0.000000	H	-1.792019	0.273905	-1.706863			
		p3				p4				
					00					
	26.	$2 \text{ kJ mol}^{-1}$			30.	$3 \text{ kJ mol}^{-1}$				
		$C_s^{-1}A'$				$C_s - A'$				
Si	-1.497455	1.581760	-1.469801	Si	0.015583	1.421694	0.441321			
$\mid C$	-1.339958	2.070906	0.170228	C	-1.710909	1.657474	1.228824			
C	0.014251	1.630794	0.108585	C	-1.833196	1.452865	-0.080188			
C	1.093044	1.182725	-0.237854	C	-1.051473	1.232997	-1.134185			
H	-1.794955	2.499942	1.048186	H	-2.213088	1.188068	2.060067			
H	2.054975	0.801432	-0.489589	H	-1.030405	1.710504	-2.101157			

Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars



		[3]		]			
_		<b>D-O-(</b>	D				
	394	$1.5 \text{ kJ mol}^{-1}$					
	0.145000	$\frac{\mathrm{C}_{\infty \mathrm{v}}^{-1}\mathrm{II}}{0.440617}$	0.00050				
	-2.145893	0.448617	-0.269059				
	-0.515622	0.067722	0.114369				
	0.708771 1.042400	-0.218171	0.402492				
	1.942409	-0.000000 T	0.093099	to S	tructures		
		 n1→n?			inuctures	<u>n1⊸n3</u>	
	73.	$7 \text{ kJ mol}^{-1}$			292	$1.3 \text{ kJ mol}^{-1}$	
		$C_1 - A$				$C_s$ - <sup>1</sup> A'	
Si	-1.658581	1.905372	0.779310	Si	-2.507882	1.510190	-1.237408
C	-1.148281	2.429951	-0.785493	C	-1.349116	0.507922	-0.612484
	0.093463	2.368595	-0.368703		-0.344930	-0.152986	0.049360
	1.352227	2.282628	0.047174		0.577768	-0.599243	0.746176
H	2.023595	3.126519	-0.048111	H	-0.353452	0.864522	0.641182
		1 940940	0 449559	I Ц	1 360507	0 008868	1 200550

		$p2 \rightarrow p5$		$p2 \rightarrow p6$				
_			R					
	216	$5.6 \text{ kJ mol}^{-1}$		$326.5 \text{ kJ mol}^{-1}$				
$C_s^{-1}A'$					$C_1$ –1A			
Si	-1.747415	$2.19\overline{2117}$	$0.63\overline{1384}$	Si	$-0.05\overline{5032}$	$1.78\overline{3261}$	0.075704	
C	-0.726407	2.290059	-0.882772	C	-1.168897	0.125688	-0.408579	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

	216	$.6 \text{ kJ mol}^{-1}$		$326.5 \text{ kJ mol}^{-1}$				
		$C_s-^1A'$				$C_1$ -1A		
Si	-1.747415	2.192117	0.631384	Si	-0.055032	1.783261	0.075704	
C	-0.726407	2.290059	-0.882772	C	-1.168897	0.125688	-0.408579	
C	0.016254	2.219441	0.409350	C	-1.787386	1.063113	0.218017	
C	1.348831	2.236480	0.302214	C	-1.427514	2.444367	1.067918	
H	1.899910	3.173627	0.295983	H	-2.035607	3.266860	0.681532	
H	1.914855	1.315559	0.186033	H	-2.549883	1.381730	1.028390	
		$p3 \rightarrow p4$				$p3 \rightarrow p6$		
							Ð	
	34.	$\frac{8 \text{ kJ mol}^{-1}}{2 \text{ c}^{-1}}$			190	$1.5 \text{ kJ mol}^{-1}$		
		$C_1$ -'A				$C_1$ -'A		
Si	-1.069242	1.492753	-1.390925	Si	-1.913929	1.504303	-1.242424	
C	-1.320148	1.984111	0.282526	C	-0.837758	2.188007	-0.048120	
C	0.040325	1.733392	0.253700	C	-0.271626	0.758029	-0.566567	
C	1.042599	1.415660	-0.447969	C	0.707158	1.171608	0.093043	
H	-1.929010	2.506285	0.997954	H	-0.902898	2.598930	0.949658	
H	1.756851	0.623631	-0.570277	H	1.716663	1.490465	0.251411	

Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$	
Molecules in the Circumstellar Medium of Carbon Stars	207

		$\mathrm{p5}{ ightarrow}[1]$		$\mathbf{p2}{ ightarrow}[2]$				
	20						8	
	417	$7.8 \text{ kJ} \text{ mol}^{-1}$			400	$0.4 \text{ kJ} \text{ mol}^{-1}$		
		$C_s$ – $^1A'$				$C_1-^1A$		
Si	-0.467866	3.174918	0.595550	Si	-3.654008	0.659771	0.202068	
C	-0.477089	2.021651	-0.792545	C	-3.156407	2.579324	0.092405	
C	-0.673171	0.590493	-0.708054	C	-2.151736	1.771052	-0.080262	
C	-0.877286	1.422661	0.457992	C	-1.646692	0.497172	0.138953	
H	-1.559940	-0.070334	-1.308318	H	-1.557132	-0.024346	-1.630107	
H	-2.249902	0.491914	-1.259838	H	-1.044179	-0.140306	-1.013032	
	0-1			-				
	412	$2.5 \text{ kJ mol}^{-1}$		-				
$C_{\rm s}^{-1}{\rm A}'$				1				
Si	-1.284546	0.614725	0.341830	1				
C	0.280930	-0.071001	0.309093					
C	1.457679	-0.594085	0.274620					
C	2.575093	-1.242370	0.051990					
H	3.725873	-0.775049	1.241686					
H	3.408591	-0.262386	1.718808					

# Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$ 208Molecules in the Circumstellar Medium of Carbon Stars208

	p1								
State	Excitation	Oscillator	Sym	Major Amplitudes	Attachment Density	Detachment Density			
	Energy (eV)	Strength							
42	9.865	$5.61 \times 10^{-2}$	$^{1}A_{1}$	$\begin{array}{l} 2b_2 \rightarrow 4b_2 = 0.300 \\ 2b_1 \rightarrow 5b_1 = -0.199 \\ 11a_1 \rightarrow 12a_1 = -0.440 \\ 11a_1 \rightarrow 13a_1 = -0.179 \\ 3b_2 \rightarrow 5b_2 = -0.277 \\ 3b_1 \rightarrow 5b_1 = -0.196 \end{array}$	ysoc				
46	10.074	$1.99 \times 10^{-4}$	$^{1}\mathrm{B}_{2}$	$9a_1 \rightarrow 4b_2 = 0.225$ $10a_1 \rightarrow 4b_2 = 0.617$	yo	<b>B</b>			
47	10.175	$6.77 \times 10^{-2}$	$^{1}\mathrm{B}_{2}$	$3b_2 \rightarrow 13a_1 = 0.675$	Sourc'	<b>*</b>			
50	10.430	0.177	$^{1}A_{1}$	$\begin{array}{l} 2b_2 \rightarrow 4b_2 = 0.546 \\ 11a_1 \rightarrow 12a_1 = 0.225 \\ 11a_1 \rightarrow 14a_1 = -0.190 \\ 3b_2 \rightarrow 5b_2 = 0.165 \\ 3b_1 \rightarrow 5b_1 = 0.161 \end{array}$	yo	e con			
52	10.544	$2.45 \times 10^{-2}$	<sup>1</sup> B <sub>1</sub>	$\begin{array}{l} 10a_1 \to 4b_1 = 0.454 \\ 11a_1 \to 5b_1 = -0.440 \\ 3b_1 \to 6b_1 = -0.257 \end{array}$	3 Andrews	Concella S			

Table A.7: Excited states of six  $SiC_3H_2$  isomers **p1-p6** with excitation energies in the region of a Lyman- $\alpha$  photon.

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

p2										
State	Excitation Energy (eV)	Oscillator Strength	Sym	Major Amplitudes	Attachment Density	Detachment Density				
46	10.029	0.538	<sup>1</sup> A′	$12a' \rightarrow 17a' = 0.375$ $3a'' \rightarrow 5a'' = 0.313$ $13a' \rightarrow 17a' = -0.369$						
49	10.207	0.203	1A'	$13a' \rightarrow 18a' = 0.633$	<b>\$</b> **					
50	10.320	$1.04 \times 10^{-3}$	1A″	$\begin{array}{l} 11a' \to 4a'' = -0.365 \\ 3a' \to 17a' = -0.159 \\ 3a'' \to 18a' = 0.553 \end{array}$						
51	10.436	0.235	<sup>1</sup> A′	$2a^{\prime\prime} \rightarrow 4a^{\prime\prime} = 0.637$						
52	10.588	0.182	<sup>1</sup> A''	$11a' \rightarrow 4a'' = 0.566$ $3a'' \rightarrow 18a' = 0.371$						

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

				p3		
State	Excitation	Oscillator	Sym	Major Amplitudes	Attachment Density	Detachment Density
48	9.879	1.30×10 <sup>-2</sup>	<sup>1</sup> A'	$13a' \rightarrow 18a' = 0.574$ $14a' \rightarrow 18a' = -0.343$		
49	9.957	$1.08 \times 10^{-2}$	<sup>1</sup> A″	$12a' \rightarrow 4a'' = 0.612$ $3a'' \rightarrow 20a' = -0.279$		
50	10.007	0.178	<sup>1</sup> A′	$12a' \rightarrow 15a' = 0.620$		
51	10.220	$1.81 \times 10^{-2}$	<sup>1</sup> A″	$3a^{\prime\prime} \rightarrow 21a^{\prime} = 0.650$		
55	10.678	0.120	<sup>1</sup> A′	$ \begin{array}{r} 14a' \to 18a' = 0.195 \\ 14a' \to 19a' = 0.580 \end{array} $	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars2

				p4		
State	Excitation	Oscillator	Sym	Major Amplitudes	Attachment Density	Detachment Density
43	9.826	2.49×10 <sup>-2</sup>	<sup>1</sup> A	$5b \rightarrow 8b = 0.534$ $6b \rightarrow 9b = -0.237$ $6b \rightarrow 10b = -0.348$		
44	9.889	0.895	<sup>1</sup> B	$7b \rightarrow 11a = 0.163$ $7b \rightarrow 12a = 0.560$ $9a \rightarrow 9b = 0.249$ $10a \rightarrow 10b = -0.177$		
47	10.207	0.248	<sup>1</sup> A	$8a \rightarrow 11a = 0.169$ $5b \rightarrow 8b = 0.384$ $6b \rightarrow 9b = 0.231$ $6b \rightarrow 10b = 0.426$ $9a \rightarrow 11a = 0.164$		
50	10.422	$5.62 \times 10^{-2}$	<sup>1</sup> A	$6b \rightarrow 9b = 0.575$ $6b \rightarrow 10b = -0.320$		
52	10.507	0.147	<sup>1</sup> B	$8a \rightarrow 8b = -0.168$ $5b \rightarrow 11a = -0.251$ $10a \rightarrow 12b = 0.602$		
53	10.598	$5.95 \times 10^{-3}$	<sup>1</sup> B	$5b \rightarrow 11a = 0.618$ $10a \rightarrow 12b = 0.259$		

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

				$\mathbf{p5}$		
State	Excitation	Oscillator	Sym	Major Amplitudes	Attachment Density	Detachment Density
40	Energy (eV) 9.778	Strength 1.17×10 <sup>-2</sup>	$^{1}B_{2}$	$4b_2 \rightarrow 11a_1 = 0.621$ $4b_2 \rightarrow 12a_1 = -0.264$	a contraction of the second	~ <b>~</b>
41	9.848	$5.06 \times 10^{-7}$	$^{1}A_{2}$	$3b_1 \rightarrow 6b_2 = 0.678$		
42	9.909	$1.32 \times 10^{-2}$	$^{1}B_{2}$	$4b_2 \rightarrow 11a_1 = 0.259$ $4b_2 \rightarrow 12a_1 = 0.636$	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	~ <b>*</b>
44	10.025	0.499	$^{1}A_{1}$	$9a_1 \rightarrow 11a_1 = -0.155$ $9a_1 \rightarrow 12a_1 = 0.659$	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~
46	10.156	$1.17 \times 10^{-2}$	<sup>1</sup> B <sub>1</sub>	$9a_1 \rightarrow 5b_1 = 0.208$ $10a_1 \rightarrow 5b_1 = 0.666$		~
50	10.468	0.436	$^{1}A_{1}$	$\begin{array}{c} 9a_1 \rightarrow 11a_1 = 0.589 \\ 9a_1 \rightarrow 13a_1 = 0.154 \\ 3b_1 \rightarrow 5b_1 = 0.201 \\ 10a_1 \rightarrow 11a_1 = -0.155 \\ 10a_1 \rightarrow 13a_1 = -0.186 \end{array}$	~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	<b>A</b>

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars

				p6		
State	Excitation	Oscillator	Sym	Major Amplitudes	Attachment Density	Detachment Density
45	9.801	0.158	<sup>1</sup> A	$15a \rightarrow 22a = 0.554$ $16a \rightarrow 23a = 0.291$ $16a \rightarrow 24a = 0.195$		
47	9.857	$5.58 \times 10^{-2}$	<sup>1</sup> A	$17a \rightarrow 25a = 0.673$		
49	9.991	$4.45 \times 10^{-2}$	<sup>1</sup> A	$13a \rightarrow 18a = 0.167$ $15a \rightarrow 22a = -0.286$ $16a \rightarrow 22a = 0.162$ $16a \rightarrow 23a = 0.572$		
51	10.154	$4.37 \times 10^{-2}$	<sup>1</sup> A	$13a \rightarrow 18a = 0.410$ $15a \rightarrow 23a = 0.228$ $16a \rightarrow 22a = 0.160$ $16a \rightarrow 23a = -0.201$ $16a \rightarrow 24a = 0.382$		
53	10.236	0.187	<sup>1</sup> A	$13a \rightarrow 18a = -0.451$ $15a \rightarrow 23a = 0.367$ $16a \rightarrow 24a = 0.259$		
54	10.292	$2.07 \times 10^{-2}$	<sup>1</sup> A	$13a \rightarrow 19a = 0.573$ $15a \rightarrow 23a = 0.321$	Ŷ	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ 214Molecules in the Circumstellar Medium of Carbon Stars214

57	10.338	$5.44 \times 10^{-2}$	<sup>1</sup> A	$13a \rightarrow 19a = 0.317$ $15a \rightarrow 23a = -0.403$ $16a \rightarrow 24a = 0.381$	•	
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#### Appendix A. Additional Information: Modeling of Gas Phase Formation of c-SiC<sub>3</sub> Molecules in the Circumstellar Medium of Carbon Stars

Table A.8: Products relevant to the photodissociation of  $SiC_3H_2$  to form atomic hydrogen and SiC<sub>3</sub>H. Geometries were calculated at the  $\omega$ B97X-V//cc-pVTZ level of theory; relative energies were calculated at the CCSD(T)//CBS level of theory.

		1				2	
							•
	380	$3.9 \text{ kJ mol}^{-1}$			399	$0.7 \text{ kJ mol}^{-1}$	
		$C_{\infty v}^{-2}\Pi$				$C_s$ – $^2A'$	
Si	-0.151031	2.995567	0.000000	Si	0.160190	2.314648	0.000000
C	-0.284810	1.318339	0.000000	C	1.981786	2.002412	0.000000
C	-0.391670	-0.020999	0.000000	C	1.341367	0.789581	0.000000
C	-0.488903	-1.239734	0.000000	C	0.164322	0.230650	0.000000
H	-0.573593	-2.301281	0.000000	H	3.012843	2.311956	0.000000
	•••			-			
	409	$\frac{9.7 \text{ kJ mol}^{-1}}{\text{Co}^{-2}\text{A}_{1}}$		-			
Si	0.375572	$\frac{2}{2}$ 2.261711	0.000000	-			
	-0.388292	0.542499	0.000000				
Ċ	1.102760	0.526660	0.000000				
Ċ	0.345057	-0.611559	0.000000				
H	0.333627	-1.686713	0.000000				

### Appendix A. Additional Information: Modeling of Gas Phase Formation of $c-SiC_3$ 216Molecules in the Circumstellar Medium of Carbon Stars216

Table A.9: Tunneling-corrected RRKM reaction rate constants, transition state imaginary frequencies, and reaction path degeneracies for the silicon–methylacetylene, silicon–allene, and silicon–D3-methylacetlyene systems.

Reaction	$k_{\rm f}  ({\rm s}^{-1})$	$k_{\rm b}~({\rm s}^{-1})$	$\omega_{\rm b}~({\rm cm}^{-1})$	Forward	Backwards
				Reaction Path	Reaction Path
				Degeneracy	Degeneracy
$[i1] \rightarrow [i2]$	$3.74 \times 10^{6}$	$5.64 \times 10^{7}$	1235	2	1
$[i1] \rightarrow [i4]$	$1.37 \times 10^{11}$	$5.67 \times 10^{10}$	168	2	2
[i1]→[i5]	$3.44 \times 10^{8}$	$2.88 \times 10^{9}$	1399	2	2
[i1]→[i6]	$6.03 \times 10^{7}$	$1.39 \times 10^{9}$	1018	2	1
$[i2] \rightarrow [i5]$	$8.55 \times 10^{12}$	$4.76 \times 10^{12}$	190	1	2
[i2]→[i10]	$2.78 \times 10^5$	$1.52 \times 10^{6}$	1207	1	2
[i4]→[i13]	$1.45 \times 10^{11}$	$6.54 \times 10^{12}$	203	2	2
$[i5] \rightarrow [i9]$	$1.04 \times 10^{11}$	$2.38 \times 10^{11}$	883	2	2
[i6]→[i9]	$4.62 \times 10^{12}$	$3.83 \times 10^{12}$	240	1	2
[i6]→[i10]	$1.16 \times 10^{10}$	$4.13 \times 10^{10}$	989	1	2
[i6]→[i12]	$3.73 \times 10^{12}$	$3.22 \times 10^{12}$	217	1	2
[i7]→[i10]	$7.49 \times 10^{11}$	$3.28 \times 10^{12}$	124	2	2
$[i8] \rightarrow [i9]$	$2.05 \times 10^{7}$	$5.54 \times 10^{7}$	1532	2	2
[i8]→[i12]	$2.11 \times 10^9$	$5.95 \times 10^{9}$	1082	2	2
$[i9] \rightarrow p1$	$1.36 \times 10^{9}$		1470	2	
$[i1] \rightarrow p2$	$4.81 \times 10^2$		1225	1	
$[i6] \rightarrow p2$	$8.41 \times 10^{6}$		1429	1	
[i13]→ <b>p2</b>	$1.09 \times 10^{7}$		1399	2	
[i12]→ <b>p3</b>	$1.65 \times 10^{8}$		1475	2	
[i12]→ <b>p3</b> ′	$2.49 \times 10^{8}$		1424	2	
$[i6] \rightarrow p4$	$1.97 \times 10^{9}$		908	1	
[i1]→ <b>n2</b>	$1.47 \times 10^5$			1	
[i13]→ <b>n6</b>	$1.39 \times 10^5$			1	
[i1]→ <b>n6</b>	$6.58 \times 10^{3}$			1	
[i2]→ <b>n3</b>	$4.94 \times 10^{5}$			1	
$[i2] \rightarrow n4$	$7.54 \times 10^4$			1	
[i4]→ <b>n1</b>	$4.11 \times 10^{6}$			1	
$[i6] \rightarrow n4$	$2.32 \times 10^5$			1	
$[i1] \rightarrow m2$	$5.40 \times 10^5$			1	

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ 217Molecules in the Circumstellar Medium of Carbon Stars217

Reaction	$k_{\rm f}  ({\rm s}^{-1})$	$k_{\rm b}~({\rm s}^{-1})$	$\omega_{\rm b}~({\rm cm}^{-1})$	Forward	Backwards
				Reaction Path	Reaction Path
				Degeneracy	Degeneracy
$[i1a] \rightarrow [i2a]$	$1.28 \times 10^{6}$	$1.93 \times 10^{7}$	977	1	1
$[i1b] \rightarrow [i2b]$	$1.56 \times 10^{6}$	$2.33 \times 10^{7}$	1184	1	1
$[i1b] \rightarrow [i2c]$	$1.34 \times 10^{6}$	$1.99 \times 10^{7}$	1006	1	1
$[i1b] \rightarrow [i2d]$	$1.35 \times 10^{6}$	$2.01 \times 10^{7}$	1000	1	1
$[i1a] \rightarrow [i2e]$	$1.28 \times 10^{6}$	$1.93 \times 10^{7}$	977	1	1
$[i1b] \rightarrow [i2f]$	$1.56 \times 10^{6}$	$2.33 \times 10^{7}$	1184	1	1
$[i1b] \rightarrow [i2g]$	$1.34 \times 10^{6}$	$1.99 \times 10^{7}$	1006	1	1
$[i1b] \rightarrow [i2h]$	$1.35 \times 10^{6}$	$2.01 \times 10^{7}$	1000	1	1
$[i1a] \rightarrow [i4a]$	$9.67 \times 10^{10}$	$4.35 \times 10^{10}$	166	2	2
$[i1b] \rightarrow [i4b]$	$9.27 \times 10^{10}$	$5.44 \times 10^{10}$	152	2	2
$[i1a] \rightarrow [i5a]$	$1.49 \times 10^{8}$	$1.48 \times 10^{9}$	1063	2	2
[i1b]→[i5b]	$2.00 \times 10^{8}$	$1.93 \times 10^{9}$	1371	2	2
$[i1b] \rightarrow [i5c]$	$1.57 \times 10^{8}$	$1.58 \times 10^{9}$	1075	2	2
[i1b]→[i5d]	$1.53 \times 10^{8}$	$1.51 \times 10^{9}$	1078	2	2
[i1a]→[i6a]	$2.58 \times 10^{7}$	$6.60 \times 10^{8}$	801	1	1
[i1b]→[i6b]	$3.44 \times 10^{7}$	$8.92 \times 10^{8}$	993	1	1
[i1b]→[i6c]	$2.50 \times 10^{7}$	$6.49 \times 10^8$	803	1	1
[i1b]→[i6d]	$2.53 \times 10^{7}$	$6.55 \times 10^{8}$	816	1	1
[i1a]→[i6e]	$2.58 \times 10^{7}$	$6.60 \times 10^{8}$	801	1	1
[i1b]→[i6f]	$3.44 \times 10^{7}$	$8.92 \times 10^{8}$	993	1	1
[i1b]→[i6g]	$2.50 \times 10^{7}$	$6.49 \times 10^8$	803	1	1
[i1b]→[i6h]	$2.53 \times 10^{7}$	$6.55 \times 10^{8}$	816	1	1
$[i2a] \rightarrow [i2e]$	$3.64 \times 10^{12}$	$3.64 \times 10^{12}$	227	1	1
[i2b]→[i2f]	$3.98 \times 10^{12}$	$3.98 \times 10^{12}$	249	1	1
[i2c]→[i2h]	$3.65 \times 10^{12}$	$3.68 \times 10^{12}$	229	1	1
[i2d]→[i2g]	$3.68 \times 10^{12}$	$3.65 \times 10^{12}$	229	1	1
[i2a]→[i3a]	$9.68 \times 10^{9}$	$1.04 \times 10^{11}$	587	1	1
[i2b]→[i3d]	$8.81 \times 10^9$	$9.40 \times 10^{10}$	551	1	1
[i2c]→[i3c]	$9.22 \times 10^{9}$	$1.02 \times 10^{11}$	589	1	1
$[i2d] \rightarrow [i3b]$	$9.12 \times 10^9$	$1.01 \times 10^{11}$	566	1	1
$[i2e] \rightarrow [i3d]$	$9.68 \times 10^{9}$	$1.04 \times 10^{11}$	587	1	1
[i2f]→[i3a]	$8.81 \times 10^9$	$9.40 \times 10^{10}$	551	1	1
[i2g]→[i3c]	$9.22 \times 10^9$	$1.02 \times 10^{11}$	589	1	1
[i2d]→[i3b]	$9.12 \times 10^9$	$1.01 \times 10^{11}$	566	1	1
[i2a]→[i5a]	$6.86 \times 10^{12}$	$4.52 \times 10^{12}$	177	1	1
[i2b]→[i5b]	$6.17 \times 10^{12}$	$3.98 \times 10^{12}$	161	1	1

Table A.9: continued

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ 218Molecules in the Circumstellar Medium of Carbon Stars218

Reaction	$k_{\rm f} \; ({\rm s}^{-1})$	$k_{\rm b}~({\rm s}^{-1})$	$\omega_{\rm b}~({\rm cm}^{-1})$	Forward	Backwards
				Reaction Path	Reaction Path
				Degeneracy	Degeneracy
$[i2c] \rightarrow [i5c]$	$5.79 \times 10^{12}$	$3.93 \times 10^{12}$	155	1	1
$[i2d] \rightarrow [i5d]$	$6.08 \times 10^{12}$	$4.00 \times 10^{12}$	160	1	1
$[i2e] \rightarrow [i5a]$	$6.86 \times 10^{12}$	$4.52 \times 10^{12}$	177	1	1
$[i2f] \rightarrow [i5b]$	$6.17 \times 10^{12}$	$3.98 \times 10^{12}$	161	1	1
$[i2g] \rightarrow [i5c]$	$5.79 \times 10^{12}$	$3.93 \times 10^{12}$	155	1	1
$[i2h] \rightarrow [i5d]$	$6.08 \times 10^{12}$	$4.00 \times 10^{12}$	160	1	1
$[i2a] \rightarrow [i10d]$	$5.41 \times 10^4$	$3.34 \times 10^{5}$	912	1	1
$[i2b] \rightarrow [i10c]$	$9.27 \times 10^{4}$	$5.70 \times 10^{5}$	1195	1	1
$[i2c] \rightarrow [i10b]$	$5.41 \times 10^4$	$3.34 \times 10^5$	903	1	1
$[i2d] \rightarrow [i10a]$	$5.28 \times 10^{4}$	$3.26 \times 10^{5}$	903	1	1
$[i2e] \rightarrow [i10b]$	$5.41 \times 10^4$	$3.34 \times 10^5$	912	1	1
$[i2f] \rightarrow [i10a]$	$9.27 \times 10^{4}$	$5.70 \times 10^{5}$	1195	1	1
[i2g]→[i10d]	$5.41 \times 10^4$	$3.34 \times 10^{5}$	903	1	1
$[i2h] \rightarrow [i10c]$	$5.28 \times 10^4$	$3.26 \times 10^5$	903	1	1
[i4a]→[i13a]	$1.16 \times 10^{11}$	$6.46 \times 10^{12}$	202	2	2
[i4b]→[i13b]	$1.28 \times 10^{11}$	$5.47 \times 10^{12}$	171	2	2
[i5a]→[i9a]	$7.40 \times 10^{10}$	$7.21 \times 10^{10}$	706	2	2
[i5b]→[i9b]	$8.26 \times 10^{10}$	$8.19 \times 10^{10}$	832	2	2
$[i5c] \rightarrow [i9d]$	$7.22 \times 10^{10}$	$6.91 \times 10^{10}$	706	2	2
$[i5d] \rightarrow [i9c]$	$7.01 \times 10^{10}$	$6.89 \times 10^{10}$	716	2	2
[i6a]→[i9a]	$3.40 \times 10^{12}$	$1.29 \times 10^{12}$	199	1	1
[i6b]→[i9b]	$4.04 \times 10^{12}$	$1.49 \times 10^{12}$	224	1	1
[i6c]→[i9c]	$3.61 \times 10^{12}$	$1.34 \times 10^{12}$	205	1	1
[i6d]→[i9d]	$3.47 \times 10^{12}$	$1.29 \times 10^{12}$	197	1	1
$[i6e] \rightarrow [i9a]$	$3.40 \times 10^{12}$	$1.29 \times 10^{12}$	199	1	1
[i6f]→[i9b]	$4.04 \times 10^{12}$	$1.49 \times 10^{12}$	224	1	1
[i6g]→[i9d]	$3.61 \times 10^{12}$	$1.34 \times 10^{12}$	205	1	1
[i6h]→[i9c]	$3.47 \times 10^{12}$	$1.29 \times 10^{12}$	197	1	1
[i6a]→[i10b]	$6.54 \times 10^{9}$	$2.39 \times 10^{10}$	795	1	1
[i6b]→[i10a]	$8.57 \times 10^{9}$	$3.05 \times 10^{10}$	980	1	1
[i6c]→[i10d]	$6.42 \times 10^9$	$2.26 \times 10^{10}$	794	1	1
$[i6d] \rightarrow [i10c]$	$6.65 \times 10^9$	$2.36 \times 10^{10}$	794	1	1
[i6e]→[i10d]	$6.54 \times 10^9$	$2.39 \times 10^{10}$	795	1	1
[i6f]→[i10c]	$8.57 \times 10^{9}$	$3.05 \times 10^{10}$	980	1	1
[i6g]→[i10b]	$6.42 \times 10^9$	$2.26 \times 10^{10}$	794	1	1
[i6h]→[i10a]	$6.65 \times 10^9$	$2.36 \times 10^{10}$	794	1	1

Table A.9: continued

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ 219Molecules in the Circumstellar Medium of Carbon Stars219

Reaction	$k_{\rm f}  ({\rm s}^{-1})$	$k_{\rm b}~({\rm s}^{-1})$	$\omega_{\rm b}~({\rm cm}^{-1})$	Forward	Backwards
				Reaction Path	Reaction Path
				Degeneracy	Degeneracy
$[i6a] \rightarrow [i12a]$	$2.39 \times 10^{12}$	$2.39 \times 10^{12}$	163	1	1
[i6b]→[i12b]	$3.32 \times 10^{12}$	$3.10 \times 10^{12}$	213	1	1
$[i6c] \rightarrow [i12d]$	$2.52 \times 10^{12}$	$2.39 \times 10^{12}$	163	1	1
$[i6d] \rightarrow [i12c]$	$2.42 \times 10^{12}$	$2.29 \times 10^{12}$	163	1	1
$[i6e] \rightarrow [i12a]$	$2.39 \times 10^{12}$	$2.39 \times 10^{12}$	163	1	1
[i6f]→[i12b]	$3.32 \times 10^{12}$	$3.10 \times 10^{12}$	213	1	1
$[i6g] \rightarrow [i12c]$	$2.52 \times 10^{12}$	$2.39 \times 10^{12}$	163	1	1
$[i6h] \rightarrow [i12d]$	$2.42 \times 10^{12}$	$2.29 \times 10^{12}$	163	1	1
$[i7a] \rightarrow [i10a]$	$6.96 \times 10^{11}$	$3.07 \times 10^{12}$	116	2	2
[i7b]→[i10b]	$7.04 \times 10^{11}$	$3.09 \times 10^{12}$	116	2	2
[i7c]→[i10d]	$6.86 \times 10^{11}$	$3.08 \times 10^{12}$	117	2	2
[i7d]→[i10c]	$7.29 \times 10^{11}$	$3.21 \times 10^{12}$	123	2	2
$[i8a] \rightarrow [i9d]$	$4.65 \times 10^{6}$	$5.06 \times 10^{6}$	1134	2	2
[i8b]→[i9c]	$7.22 \times 10^{6}$	$8.24 \times 10^{6}$	1498	2	2
[i8c]→[i9b]	$5.32 \times 10^{6}$	$5.90 \times 10^{6}$	1148	2	2
$[i8d] \rightarrow [i9a]$	$4.54 \times 10^{6}$	$5.21 \times 10^{6}$	1134	2	2
[i8a]→[i12a]	$8.84 \times 10^{8}$	$2.51 \times 10^{9}$	789	2	2
[i8b]→[i12c]	$8.48 \times 10^8$	$2.46 \times 10^9$	811	2	2
[i8c]→[i12b]	$8.72 \times 10^{8}$	$2.45 \times 10^9$	794	2	2
[i8d]→[i12d]	$1.23 \times 10^{9}$	$3.63 \times 10^{9}$	1058	2	2
[i9a]→ <b>p1a</b>	$2.29 \times 10^{8}$		1403	2	
[i9b]→ <b>p1a</b>	$1.84 \times 10^{8}$		1118	2	
[i9c]→ <b>p1b</b>	$1.43 \times 10^{8}$		1064	2	
[i9d]→p1b	$1.44 \times 10^{8}$		1064	2	
[i1a]→p2a	$4.83 \times 10^{1}$		1104	1	
[i1b]→p2a	$4.93 \times 10^{1}$		1105	1	
[i1b]→ <b>p2b</b>	$1.42 \times 10^{1}$		980	1	
[i1b]→p2c	$1.42 \times 10^{1}$		980	1	
$[i6a] \rightarrow p2a$	$2.67 \times 10^{6}$		1372	1	
$[i6b] \rightarrow p2a$	$2.12 \times 10^{6}$		1080	1	
$[i6c] \rightarrow p2b$	$1.57 \times 10^{6}$		1032	1	
$[i6d] \rightarrow p2c$	$1.58 \times 10^{6}$		1032	1	
$[i6e] \rightarrow p2a$	$2.67 \times 10^{6}$		1372	1	
$[i6f] \rightarrow p2a$	$2.12 \times 10^{6}$		1080	1	
$[i6g] \rightarrow p2c$	$1.57 \times 10^{6}$		1032	1	
$[i6h] \rightarrow p2b$	$1.58 \times 10^{6}$		1032	1	

Table A.9: continued

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ 220Molecules in the Circumstellar Medium of Carbon Stars220

Reaction	$k_{\rm f}~({\rm s}^{-1})$	$k_{\rm b}~({\rm s}^{-1})$	$\omega_{\rm b} \ ({\rm cm}^{-1})$	Forward	Backwards
				Reaction Path	Reaction Path
				Degeneracy	Degeneracy
[i13a]→ <b>p2a</b>	$3.10 \times 10^{6}$		1067	2	
$[i13b] \rightarrow p2a$	$4.78 \times 10^{6}$		1327	2	
[i13b]→ <b>p2b</b>	$2.65 \times 10^{6}$		1030	2	
$[i13b] \rightarrow p2c$	$2.65 \times 10^{6}$		1024	2	
$[i12a] \rightarrow p3a$	$4.26 \times 10^{7}$		1060	2	
[i12b]→ <b>p3c</b>	$5.00 \times 10^{7}$		1111	2	
$[i12c] \rightarrow p3c$	$6.76 \times 10^{7}$		1409	2	
$[i12d] \rightarrow p3b$	$4.61 \times 10^{7}$		1076	2	
[i12a]→ <b>p3a</b> ′	$7.47 \times 10^{7}$		1023	2	
[i12b]→ <b>p3c</b> ′	$8.62 \times 10^{7}$		1071	2	
$[i12c] \rightarrow p3b'$	$7.30 \times 10^{7}$		1034	2	
[i12d]→ <b>p3c</b> ′	$1.19 \times 10^{8}$		1364	2	
$[i6a] \rightarrow p4a$	$7.46 \times 10^{8}$		675	1	
$[i6b] \rightarrow p4b$	$8.81 \times 10^{8}$		720	1	
$[i6c] \rightarrow p4a$	$7.71 \times 10^{8}$		679	1	
$[i6d] \rightarrow p4b$	$1.05 \times 10^{9}$		837	1	
$[i6e] \rightarrow p4a$	$7.46 \times 10^{8}$		675	1	
$[i6f] \rightarrow p4b$	$8.81 \times 10^{8}$		720	1	
$[i6g] \rightarrow p4a$	$7.71 \times 10^{8}$		679	1	
$[i6h] \rightarrow p4b$	$1.05 \times 10^{9}$		837	1	

Table A.9: continued

### A.3 Figures

Appendix A. Additional Information: Modeling of Gas Phase Formation of  $c-SiC_3$ Molecules in the Circumstellar Medium of Carbon Stars



Figure A.1: Complete potential energy surfaces involved in the reactions of electronically excited silicon atoms  $(Si(^{1}D))$  with allene and methylacetylene. Relative energies are given in units of kJ mol<sup>-1</sup>. Colors of the atoms: silicon (green), carbon (black) and hydrogen(light grey). Taken with permission from Ref. 93.



Figure A.2: Pathways exploited in the RRKM calculations to determine the branching ratios in the reactions of electronically excited silicon atoms  $(Si(^{1}D))$  with methylacetylene, allene, and D3-methylacetylene. Relative energies are given in units of kJ mol<sup>-1</sup>. Colors of the atoms: silicon (green), carbon (black) and hydrogen (light grey). Taken with permission from Ref. 93.

Appendix B



#### **B.1** Reactive Pathways

Figure B.1: Potential energies, in kcal mol<sup>-1</sup>, along the  $O + H_2 + O_2 \longrightarrow OH + HO_2$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated O, H<sub>2</sub>, and O<sub>2</sub> at the ReaxFF/HO2011 optimized geometries.



Figure B.2: Potential energies, in kcal mol<sup>-1</sup>, along the  $O_2 + H_2 + HO_2 \longrightarrow HO_2 + H_2O_2$ reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated  $O_2$ ,  $H_2$ , and  $HO_2$  at the ReaxFF/HO2011 optimized geometries.



Figure B.3: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2 + O_2 \longrightarrow HO_2 + H$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated  $H_2$  and  $O_2$  at the ReaxFF/HO2011 optimized geometries.



Figure B.4: Potential energies, in kcal mol<sup>-1</sup>, along the  $H + O_2 \longrightarrow HO_2$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated H and  $O_2$  at the ReaxFF/HO2011 optimized geometries.



Figure B.5: Potential energies, in kcal mol<sup>-1</sup>, along the  $HO_2 + H \longrightarrow 2 \text{ OH}$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated  $HO_2$  and H at the ReaxFF/HO2011 optimized geometries.



Figure B.6: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2 + OH \longrightarrow H_2O + H$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated  $H_2$  and OH at the ReaxFF/HO2011 optimized geometries.



Figure B.7: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2 + OH \longrightarrow H_2O + H$  reactive trajectory, extracted from a larger MD simulation. The parent simulation was run using ReaxFF with the HO2011 parameter set. The energies presented are relative to isolated  $H_2$  and OH at the ReaxFF/HO2011 optimized geometries.

#### **B.2** Intrinsic Reaction Coordinate Scans



Figure B.8: Potential energies, in kcal mol<sup>-1</sup>, along the  $H+O_2 \longrightarrow OH+O$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M-V/cc-pVTZ$  level of theory. The energies presented are relative to the energies of isolated H and O<sub>2</sub> computed at the  $\omega B97M-V/cc-pVTZ$  optimized geometries.

Appendix B. Additional Information: Benchmarking the Performance of the ReaxFF Reactive Force Field on Hydrogen Combustion Systems



Figure B.9: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2 + OH \longrightarrow H_2O + H$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega$ B97M-V/ccpVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2$  and OH computed at the  $\omega$ B97M-V/cc-pVTZ optimized geometries.



Figure B.10: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2O + O \longrightarrow 2OH$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega$ B97M-V/ccpVTZ level of theory. The energies presented are relative to the energies of isolated H<sub>2</sub>O and O computed at the  $\omega$ B97M-V/cc-pVTZ optimized geometries.



Figure B.11: Potential energies, in kcal mol<sup>-1</sup>, along the  $H + O_2 \longrightarrow HO_2$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated H and  $O_2$  computed at the  $\omega$ B97M-V/cc-pVTZ optimized geometries.



Figure B.12: Potential energies, in kcal mol<sup>-1</sup>, along the  $HO_2 + H \longrightarrow H_2 + O_2$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega$ B97M-V/ccpVTZ level of theory. The energies presented are relative to the energies of isolated  $HO_2$ and H computed at the  $\omega$ B97M-V/cc-pVTZ optimized geometries.



Figure B.13: Potential energies, in kcal mol<sup>-1</sup>, along the HO<sub>2</sub> + H  $\longrightarrow$  2 OH intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega$ B97M-V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated HO<sub>2</sub> and H computed at the  $\omega$ B97M-V/cc-pVTZ optimized geometries.



Figure B.14: Potential energies, in kcal mol<sup>-1</sup>, along the  $2 \text{HO}_2 \longrightarrow \text{H}_2\text{O}_2 + \text{O}_2$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M\text{-V/cc-}$  pVTZ level of theory. The energies presented are relative to the energies of two isolated HO<sub>2</sub> computed at the  $\omega B97M\text{-V/cc-}$  pVTZ optimized geometries.



Figure B.15: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2O_2 + H \longrightarrow H_2O + OH$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2O_2$  and H computed at the  $\omega B97M$ -V/cc-pVTZ optimized geometries.



Figure B.16: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2O_2 + H \longrightarrow H_2 + HO_2$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2O_2$  and H computed at the  $\omega B97M$ -V/cc-pVTZ optimized geometries.



Figure B.17: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2O_2 + O \longrightarrow HO_2 + OH$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M$ -V/cc-pVTZ level of theory. The energies presented are relative to the energies of isolated  $H_2O_2$  and O computed at the  $\omega B97M$ -V/cc-pVTZ optimized geometries.



Figure B.18: Potential energies, in kcal mol<sup>-1</sup>, along the  $H_2O_2+OH \longrightarrow H_2O+HO_2$  intrinsic reaction coordinate. The intrinsic reaction coordinate was calculated at the  $\omega B97M-V/cc-pVTZ$  level of theory. The energies presented are relative to the energies of isolated  $H_2O_2$  and OH computed at the  $\omega B97M-V/cc-pVTZ$  optimized geometries.

### Appendix C

### Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

#### C.1 W4-11

Table C.1 presents CCSD(T) reaction energies for the non-multireference (non-MR) subset of the W4-11 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for the MP2, OOMP2, and  $\kappa$ -OOMP2 calculations. For this set, we did not use the RI approximation for the CCSD(T), CCSD, or MP3 contributions.

$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	action	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta$ (MP2.8: $\kappa$ -OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(\mathrm{MP2.5})$	$\Delta(MP3)$	$\overline{\Delta}(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\rightarrow \mathrm{H}_2$	-108.55	0.00	1.90	1.20	3.04	1.30	5.64	4.74	4.78
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 3 \mathrm{H} \rightarrow \mathrm{AlH}_3$	-214.45	0.18	1.73	0.74	3.15	0.65	4.43	5.74	5.64
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H \rightarrow AlH$	-73.97	0.29	1.42	0.93	4.97	3.96	4.65	3.41	5.98
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 4 \mathrm{H}  ightarrow \mathrm{SiH}_4$	-324.46	0.91	2.74	1.46	5.13	1.91	7.73	8.01	8.34
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 3 H \rightarrow B H_3$	-282.54	0.76	-0.97	-1.70	0.18	-1.69	0.84	2.04	2.04
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H \rightarrow BH$	-85.81	0.56	1.86	1.42	6.93	6.58	4.21	3.34	7.29
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 2 H \rightarrow {}^{3}CH_{2}$	-192.08	0.89	-1.01	-1.23	0.05	-0.53	0.20	0.14	0.63
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H \rightarrow SiH$	-73.58	0.44	1.61	1.15	3.41	2.32	4.20	3.55	4.50
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$ m i+6H ightarrow m Si_{2}H_{6}$	-536.29	2.95	3.57	1.83	7.43	2.96	11.80	10.85	11.90
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+3 \text{H} \rightarrow \text{CH}_3^{-1}$	-308.24	1.77	-0.15	-0.59	1.24	0.04	2.83	1.96	2.44
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$ m + 4H  ightarrow { m CH}_4^{ m CH}$	-420.43	2.76	0.42	-0.02	1.94	0.68	4.46	2.65	3.19
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 6 H \rightarrow B_2 H_6$	-611.16	4.26	-2.59	-3.61	0.14	-2.62	1.42	2.07	2.89
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$+ 3 \mathrm{H} + \mathrm{F} \rightarrow \mathrm{SiH}_3 \mathrm{F}$	-379.79	2.95	1.10	2.61	0.97	4.02	-4.79	-3.65	-2.09
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 3 \mathrm{H}  ightarrow \mathrm{PH}_3$	-240.35	2.05	3.19	2.18	6.03	3.43	10.37	7.59	8.63
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+6{ m H}  ightarrow { m C}_2{ m H}_6$	-714.06	6.09	-0.58	-0.45	1.04	0.96	2.85	-0.11	1.12
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 8 H \rightarrow propane$	-1010.13	9.65	-1.58	-0.77	0.10	1.42	1.04	-3.27	-1.21
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 2 \mathrm{H} \rightarrow {}^{1}\mathrm{CH}_{2}$	-181.88	1.23	0.04	-0.40	2.06	0.98	2.88	2.12	3.14
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H \rightarrow CH$	-84.78	0.94	1.15	0.68	5.91	5.17	4.19	3.20	6.65
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$(+ S \rightarrow H_2 S)$	-181.88	2.31	1.38	1.17	1.93	1.28	5.67	2.31	2.57
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ S \rightarrow HS$	-86.56	1.08	0.87	0.66	1.48	0.91	3.30	1.80	2.05
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+5 H + F \rightarrow C_2 H_5 F$	-722.71	8.49	-1.57	1.00	-1.36	3.92	-6.63	-9.48	-6.64
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ \mathrm{N} + 5 \mathrm{H} \rightarrow \mathrm{CH}_3 \mathrm{NH}_2$	-580.56	6.93	1.05	2.12	1.87	3.93	2.22	-1.93	-0.18
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+3 \mathrm{H} + \mathrm{F} \rightarrow \mathrm{CH}_3 \mathrm{F}$	-422.88	5.00	-0.39	1.55	-0.54	3.51	-4.79	-6.46	-4.58
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$+ 6 H \rightarrow propene$	-863.23	10.67	-2.29	-0.54	-1.38	2.21	-0.61	-7.50	-4.98
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 3 H \rightarrow NH_3$	-295.48	3.67	1.86	2.40	2.36	3.37	3.64	0.38	1.35
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ O + 6 H \rightarrow ethanol$	-811.22	10.01	-0.82	2.00	-1.15	4.70	-5.38	-9.91	-6.99
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ \text{ N} + 4 \text{ H} \rightarrow \text{ CH}_3 \text{ NH}$	-473.50	5.81	0.19	0.71	2.93	3.57	3.10	-0.29	2.30
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 4{ m H}  ightarrow { m C}_2{ m H}_4$	-564.36	7.09	-1.19	-0.18	-0.24	1.83	1.57	-3.99	-2.32
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ O + 4 H \rightarrow methanol$	-512.19	6.48	0.26	2.40	-0.25	4.23	-3.46	-6.73	-4.73
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ \text{Cl} \rightarrow \text{HCl}$	-106.39	1.52	0.37	0.60	-0.08	0.40	1.44	-0.54	-0.57
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$ m + 2  H  ightarrow  m NH_2$	-180.84	2.33	1.50	1.41	2.90	2.49	4.57	2.42	3.31
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ \mathrm{H} \rightarrow \mathrm{NH}$	-82.32	1.06	0.94	0.65	2.38	1.63	3.46	2.46	3.14
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ \text{ N} + 4 \text{ H} \rightarrow \text{ CH}_2\text{ NH}_2$	-480.36	6.48	-0.05	1.24	1.68	4.18	0.18	-3.69	-0.82
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H + 2 F \rightarrow BHF_2$	-411.02	5.67	-4.17	0.49	-6.89	3.28	-21.38	-20.21	-17.06
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 0 \rightarrow H_2 0$	-230.63	3.29	1.04	2.74	-0.20	3.48	-2.45	-4.98	-3.89
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\vdash F \rightarrow HF$	-140.37	1.98	0.43	2.03	-0.88	2.64	-4.12	-5.26	-4.40
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 3 H \rightarrow CH_2 CH$	-446.39	6.26	-2.31	-1.23	4.61	6.22	-0.03	-5.11	3.00
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ H \rightarrow OH$	-106.00	1.56	0.50	0.94	0.69	1.59	0.37	-0.84	-0.21
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+ 4 \mathrm{H} \rightarrow \mathrm{propyne}$	-706.71	11.43	-3.25	0.28	-4.44	3.27	-5.15	-15.19	-12.14
acetaldehyde       acetaldehyde       -706.21       11.76       -3.05       -0.32       -2.83       3.02 $3C + 4H \rightarrow$ allene       -706.21       11.76       -3.05       -0.32       -2.83       3.02 $Si + 4F \rightarrow$ SiF $_4$ -565.82       9.26       -3.98       5.27       -10.90       9.68 $B + 3F \rightarrow BF_3$ -470.35       7.65       -5.83       1.17       -10.73       4.79	$+ 0 + 4 H \rightarrow$	-677.49	11.06	-1.68	2.69	-2.35	6.73	-9.28	-16.01	-11.43
$3.0 \pm 4.11 \rightarrow aucuce$ -100.21 11.10 -3.09 -0.32 -2.03 -0.02 Si $\pm 4.7 \rightarrow SiF_4$ -565.82 9.26 -3.98 5.27 -10.90 9.68 B $\pm 3.7 \rightarrow BF_3$ -470.35 7.65 -5.83 1.17 -10.73 4.79	taldehyde	706 91	94 11	20 G	66 0	60 C	60.6	10 0	11 01	090
$D1 + 4\Gamma \rightarrow D1\Gamma_4$ -D05.62 9.20 -5.96 0.21 -10.90 9.05 $B + 3\Gamma \rightarrow B\Gamma_3$ -10.73 -10.73 4.79		17.001-	0.111	00.6-		00.4-	20.0	10.2-	10'TT-	10.00
$B + 3F \rightarrow BF_3$ -470.35 7.65 -5.83 1.17 -10.73 4.79	$+4.F \rightarrow SiF_4$	28.000-	9.20	-3.98	7.7.0	-10.90	9.08	-38.94	-30.05	-31.47
	$+ 3 F \rightarrow B F_3$	-470.35	7.65	-5.83	1.17	-10.73	4.79	-31.89	-30.34	-26.26

Appendix C.Additional Information: Third-Order Møller-Plesset Perturbation TheoryMade Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy233
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$8:\kappa - \Delta(MP3:\kappa - \infty) = 0.0000000000000000000000000000000000$	$\Delta(\mathrm{MP2.5})$	(GTM)	$\Delta(\kappa^{-2})$	$\Delta(OOMP2)$	$\Delta(MP2)$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	2) UUMP2)		1	UUMP2)	007	07
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	2.48 1.12	-2.88	4.73	-8.59	-14.09	-10.49
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.81 2.51	-3.26	5.91	-10.78	-16.23	-12.42
$\begin{array}{cccccc} \mathrm{AIF}_3 & -423.42 & 7.18 & -3.0 \\ \mathrm{BeF}_2 & -510.40 & 5.22 & -4.8 \\ \mathrm{CH}_2 \mathrm{NH} & -437.60 & 8.00 & 0.61 & -1.6 \\ \mathrm{H} \rightarrow \operatorname{acetic} & -373.85 & 7.96 & 5.4 \\ \mathrm{H} \rightarrow \operatorname{acetic} & -302.21 & 15.31 & -2.7 \\ \mathrm{H} \rightarrow \operatorname{acetic} & -302.21 & 15.31 & -2.7 \\ \mathrm{C} + 160.40 & 3.13 & 8.01 & -2.2 \\ \mathrm{C} - + 120 & -373.20 & 7.51 & -0.0 \\ \mathrm{C} - + 120 & -373.20 & 7.51 & -0.0 \\ \mathrm{C} - + 120 & -373.20 & 7.51 & -0.0 \\ \mathrm{C} - + 237.19 & -1.6 & -0.0 \\ \mathrm{C} - + 237.19 & -1.6 & -0.0 \\ \mathrm{C} - + 237.19 & -1.6 & -0.0 \\ \mathrm{C} - + 237.19 & -1.6 & -0.0 \\ \mathrm{C} - + 237.19 & -338.63 & -2.2 \\ \mathrm{C} - + 237.19 & -338.63 & -2.2 \\ \mathrm{C} - + 237.19 & -338.03 & -332.66 & -3.4 \\ \mathrm{C} - + 232.00 & -332.06 & -3.4 \\ \mathrm{C} - + 232.00 & -332.06 & -3.4 \\ \mathrm{C} - + 232.00 & -332.0 & -0.1 \\ \mathrm{C} - + 232.00 & -332.0 & -0.1 \\ \mathrm{C} - + 232.00 & -332.0 & -0.1 \\ \mathrm{C} - + 232.00 & -333.0 & -0.1 \\ \mathrm{C} - + 232.00 & -333.0 & -0.1 \\ \mathrm{C} - + 232.00 & -333.0 & -0.1 \\ \mathrm{C} - + 233.00 & -333.0 & -0.1 \\ \mathrm{C} - + 233.00 & -333.0 & -0.1 \\ \mathrm{C} + 233.00 & -0.1 \\ \mathrm{C} + 233.00 & -333.0 & -0.1 \\ \mathrm{C} + -33.0 & -333.0 & -0.1 \\ \mathrm{C}257.54 & -7.63 & -0.2 \\ \mathrm{C}265.74 & -1.0 \\ \mathrm{C}265.74 & -2.0 \\ \mathrm{C}265.24 & -2.0 \\ \mathrm{C}265.$	1.70 2.79	-3.23	6.30	-14.76	-16.45	-12.77
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	3.06    4.53	-9.34	7.43	-31.88	-30.12	-26.11
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.87 0.30	-11.65	-0.42	-26.88	-23.42	-22.88
	1.60 -0.40	3.10	5.28	1.82	-4.83	0.93
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2.81 4.69	3.12	7.00	2.58	-3.06	-0.75
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	0.63 2.78	1.38	5.84	0.70	-6.20	-3.08
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2 20	-2.43	3 47	-8.96	-10.15	-8.34
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	0 77 A 73	-5 U3	10 55	-10.08	-97.84	-90.61
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	61. <del>1</del>	en.e-	00.UI	06.61-	-21.04	10.02-
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	0.47 0.47	-3,25	2.86	-2.55	-11.67	-9.36
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	337 337	-133	6.40	-6 92	-12.51	-9.05
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	00 0 22	8 03	10.24	0.78	-5.49	5.82
$\begin{array}{cccccc} F_{12} & -2.7.19 & -2.6.7 & -2.4.0 & -2.7.19 & -2.6.4 & -0.0 & -2.2.119 & -2.6.4 & -0.0 & -0.1 & -0.0 $	1.63 1.08	-0.93	50 G	0.00	-0.40	1 0 0 1 0 0
$ \begin{array}{cccccc} & -227.13 & -4.00 & -2.3 \\ -AICl_3 & -220.67 & 6.94 & -0.0 \\ -AICl_3 & -315.91 & 7.07 & -0.0 \\ AICl & + & -232.03 & 6.58 & -4.0 \\ AICl & -232.93 & 12.44 & -3.1 \\ H \rightarrow ketene & -532.93 & 12.44 & -3.1 \\ H \rightarrow formic & -139.06 & 11.89 & -0.1 \\ H \rightarrow formic & -138.40 & 3.30 & -0.1 \\ H \rightarrow formic & -138.20 & 12.44 & -3.1 \\ 2H \rightarrow glyoxal & -632.90 & 16.23 & -2.2 \\ H \rightarrow HCOF & -244.42 & 6.53 & 11.89 & -0.2 \\ H \rightarrow HCOF & -244.42 & 6.53 & 11.89 & -0.1 \\ H \rightarrow HCOF & -244.12 & 6.53 & 11.8 \\ - \rightarrow HCOF & -241.13 & 10.54 & -3.6 \\ CCH & -265.01 & 6.92 & -0.2 \\ CCH & -265.01 & 6.92 & -0.2 \\ CCH & -265.01 & 6.92 & -0.2 \\ - +HCO & -265.01 & 6.92 & -0.2 \\ CCH & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\ - +HCO & -265.01 & 0.54 & -2.6 \\267.31 & 2.86.23 & 9.55 & -2.5 \\2.0 \\2.$	00'T 00'T	00.7-	20.2	00.6-	26 V I	-0.00- 16 10
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	07.00.20	07.11-	-0.40	-14.00	10.41-	- 10.40
$ \begin{array}{ccccccc} & \mbox{AICl}_3 & \mbox{-}306.63 & \mbox{-}6.58 & \mbox{-}4.0 \\ \mbox{AICl} & \mbox{-}2.32.93 & \mbox{-}1.22.00 & \mbox{-}2.32.93 & \mbox{-}1.244 & \mbox{-}1.16 \\ \mbox{H} \rightarrow \mbox{ketene} & \mbox{-}532.93 & \mbox{-}1.244 & \mbox{-}3.1 \\ \mbox{H} \rightarrow \mbox{ketene} & \mbox{-}532.93 & \mbox{-}1.244 & \mbox{-}3.1 \\ \mbox{-}1.2 & \mbox{-}1.232.93 & \mbox{-}1.244 & \mbox{-}3.1 \\ \mbox{-}1.2 & \mbox{-}1.232.90 & \mbox{-}1.2 \\ \mbox{-}1.2 & \mbox{-}2.32.90 & \mbox{-}1.2 \\ \mbox{-}1.2 & \mbox{-}2.32.90 & \mbox{-}1.2 \\ \mbox{-}1.4 & \mbox{-}2.32.90 & \mbox{-}1.6 \\ \mbox{-}2.1 & \mbox{-}2.32.90 & \mbox{-}1.6 \\ \mbox{-}2.1 & \mbox{-}2.12 & \mbox{-}2.12 \\ \mbox{-}2.12 & \mbox{-}2.12 & \mbox{-}2.12 & \mbox{-}2.12 \\ \mbox{-}2.12 & \mbox{-}2.12 & \mbox{-}2.12 & \mbox{-}2.12 & \mbox{-}2.2$	0.08 2.70	4.45	10.10	-4.29	-9.UJ	-1.21
$\begin{array}{cccc} -\mathrm{AICl}_{3} & -\mathrm{b.05} & -\mathrm{b.05}$	00	66.0	00 F	10 7 7	2 1 7	- - - -
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	F.UU	-9.33	-1.33	-14.15	-17.03	-17.34
$\begin{array}{ccccc} \mathrm{AICI} & -120.00 & 2.86 & -1.0 \\ \mathrm{H} \rightarrow \mathrm{ketene} & -532.93 & 12.44 & -3.1 \\ \mathrm{F} & -1.00 & -1.00 & 3.30 & -0.1 \\ \mathrm{H} \rightarrow \mathrm{formic} & -498.86 & 11.89 & -1.6 \\ \mathrm{H} \rightarrow \mathrm{HCNH} & -334.08 & 8.09 & -0.5 \\ \mathrm{H} \rightarrow \mathrm{HCOF} & -31.08 & 8.09 & -0.5 \\ \mathrm{H} \rightarrow \mathrm{HCOF} & -324.42 & 6.53 & -2.6 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -344.42 & 6.53 & -2.6 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -398.03 & 10.54 & -2.1 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -398.03 & 10.54 & -2.1 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -398.03 & 10.54 & -2.6 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -398.03 & 10.54 & -2.6 \\ \mathrm{H} \rightarrow \mathrm{HCC} & -296.11 & 8.43 & -0.2 \\ \mathrm{H} \rightarrow \mathrm{HC} & -257.54 & 7.63 & -0.2 \\ \mathrm{H} \rightarrow \mathrm{HOC} & -257.54 & 7.63 & -0.2 \\ \mathrm{H} \rightarrow \mathrm{HOC} & -257.54 & 7.63 & -0.2 \\ \mathrm{H} \rightarrow \mathrm{HOC} & -265.01 & 6.92 & -0.2 \\ \mathrm{H} \rightarrow \mathrm{HOC} & -265.74 & 8.43 & 1.7 \\ \mathrm{H} \rightarrow \mathrm{HNC} & -265.74 & 8.43 & 1.7 \\ \mathrm{H} \rightarrow \mathrm{HNC} & -265.73 & 12.92 & -0.2 \\ \mathrm{CC} & -413.49 & 12.92 & -0.2 \\ \mathrm{CC} \rightarrow \mathrm{HOC} & -265.73 & 8.73 & 9.55 & -2.5 \\ \mathrm{CO} & -257.31 & 8.67 & -2.0 \\ \mathrm{CO} & -257.31 & 8.67 & -2.2 \\ \mathrm{CO} & -257.31 & 8.67 & -2.2 \\ \mathrm{CO} & -2.6 & -2.6 & -2.6 \\ \mathrm{CO} & -2.$	0.02 $2.75$	5.15	10.78	-3.94	-8.89	-0.47
$\begin{array}{ccccccc} H \rightarrow \mathrm{ketene} & -532.93 & 12.44 & -3.1 \\ H \rightarrow \mathrm{formic} & -139.06 & 11.89 & -0.1 \\ H \rightarrow \mathrm{formic} & -138.86 & 11.89 & -0.1 \\ -1.6 & -1.6 & -3.34.08 & 8.09 & -0.5 \\ -1.6 & -1.6 & -3.41.82 & 16.23 & -2.6 \\ H \rightarrow \mathrm{HH}_2\mathrm{CI} & -2.44.42 & 6.53 & 1.5 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -3.6 & -3.6 & -3.6 \\ -2.44.42 & -2.65.01 & 0.54 & -3.6 \\ -2.67 & -2.65.01 & 0.54 & -3.6 & -3.6 \\ -2.67 & -2.65.01 & 0.54 & -3.6 & -3.6 \\ -2.67 & -2.67 & -2.63 & -3.6 & -3.6 \\ +1 \rightarrow \mathrm{MCO} & -2.67.13 & 12.75 & -3.2 \\ +1 \rightarrow \mathrm{HOOH} & -265.74 & 8.43 & -1.7 \\ +1 \rightarrow \mathrm{HOC} & -2.66.33 & -2.57.31 & 8.67 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.6 & -2.6 & -2.6 & -2.6 & -2.6 & -2.6 \\ -2.6 & -2.$	0.00 0.09	-2.69	-0.12	-2.03	-5.43	-5.25
$ \begin{array}{ccccccc} F & -139.06 & 3.30 & -0.1 \\ H \rightarrow formic & -198.86 & 11.89 & -1.6 \\ 2 H \rightarrow HCNH & -334.08 & 8.09 & -0.5 \\ 2 H \rightarrow HCOF & -314.08 & 8.09 & -0.5 \\ + H \rightarrow HCOF & -304.182 & 10.48 & -2.1 \\ -7 H \rightarrow HCOF & -401.82 & 10.48 & -2.1 \\ -7 H \rightarrow HCOF & -311.37 & 8.79 & -0.1 \\ -7 HCN & -211.37 & 8.79 & -0.1 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -296.11 & 8.43 & -0.2 \\ -7 HCN & -265.01 & 6.92 & -0.2 \\ -7 H \rightarrow HOCN & -265.01 & 6.92 & -0.2 \\ -7 H \rightarrow HOCN & -265.01 & 6.92 & -0.2 \\ -7 H \rightarrow HOCN & -265.74 & 8.43 & 1.7 \\ + H \rightarrow HOCN & -265.74 & 8.43 & 1.2 \\ + H \rightarrow HOCN & -265.74 & 8.43 & 1.7 \\ + H \rightarrow HOCN & -265.73 & 9.55 & -2.2 \\ -277.31 & 8.67 & -2.0 \\ -277.21 & -257.31 & 8.67 & -2.0 \\ -267 & -200 & -286.23 & 9.55 & -2.2 \\ -200 & -257.31 & 8.67 & -2.0 \\ -200 & -257.31 & 8.67 & -2.0 \\ -200 & -286.23 & 9.55 & -2.2 \\ -200 & -201 & -2.0 \\ -201 & -2$	3.17 2.86	-4.96	7.61	-14.11	-23.46	-17.53
$\begin{array}{c} \mathrm{H} \rightarrow \mathrm{formic} \qquad -498.86 \qquad 11.89 \qquad -1.6 \\ \mathrm{Z} \rightarrow \mathrm{HCNH} \qquad -334.08 \qquad 8.09 \qquad -0.5 \\ \mathrm{Z} \rightarrow \mathrm{HCOF} \qquad -334.08 \qquad 8.09 \qquad -0.5 \\ \mathrm{Z} \rightarrow \mathrm{HCOF} \qquad -334.08 \qquad 8.09 \qquad -0.5 \\ \mathrm{Z} \rightarrow \mathrm{HCOF} \qquad -308.03 \qquad 16.23 \qquad -2.1 \\ \mathrm{Z} \rightarrow \mathrm{HCCF} \qquad -398.03 \qquad 10.54 \qquad -3.6 \\ \mathrm{Z} \rightarrow \mathrm{HCC} \qquad -311.37 \qquad 8.79 \qquad -0.1 \\ \mathrm{Z} \rightarrow \mathrm{HCO} \qquad -311.37 \qquad 8.73 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HCO} \qquad -311.37 \qquad 8.73 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HCO} \qquad -256.11 \qquad 8.43 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HCO} \qquad -257.54 \qquad 7.74 \qquad -1.1 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -257.54 \qquad 7.63 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -257.54 \qquad 7.63 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -265.74 \qquad 8.43 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -265.74 \qquad 8.43 \qquad -0.2 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -265.74 \qquad 8.43 \qquad -1.7 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -265.74 \qquad 8.43 \qquad -2.2 \\ \mathrm{Z} \rightarrow \mathrm{HOOH} \qquad -265.73 \qquad 8.67 \qquad -2.2 \\ \mathrm{Z} \rightarrow \mathrm{Z} \rightarrow$	0.15 2.64	-0.94	5.01	-9.17	-9.67	-6.89
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.60 5.33	-4.06	10.37	-17.89	-24.97	-18.48
$ \begin{array}{c} \mbox{I} \rightarrow \mbox{HCNH} & -334.08 & 8.09 & -0.5 \\ \mbox{2} \mbox{H} \rightarrow \mbox{HCOF} & -344.42 & 6.53 & -1.5 \\ \mbox{H} \rightarrow \mbox{HCOF} & -244.42 & 6.53 & -1.5 \\ \mbox{CF} & -244.42 & 6.53 & -1.5 \\ \mbox{CF} & -398.03 & 10.54 & -3.6 \\ \mbox{H} \rightarrow \mbox{HCO} & -398.03 & 10.54 & -3.6 \\ \mbox{H} \rightarrow \mbox{HCN} & -398.03 & 10.54 & -3.6 \\ \mbox{H} \rightarrow \mbox{HCN} & -296.11 & 8.43 & -0.2 \\ \mbox{H} \rightarrow \mbox{HCN} & -296.11 & 8.43 & -0.2 \\ \mbox{H} \rightarrow \mbox{HCN} & -265.01 & 6.92 & -3.6 \\ \mbox{H} \rightarrow \mbox{HCO} & -257.54 & 7.63 & -0.8 \\ \mbox{H} \rightarrow \mbox{HOOH} & -257.54 & 7.63 & -0.8 \\ \mbox{H} \rightarrow \mbox{HOOH} & -257.54 & 7.63 & -0.8 \\ \mbox{H} \rightarrow \mbox{HOOH} & -265.74 & 8.43 & -1.7 \\ \mbox{H} \rightarrow \mbox{HOOH} & -265.74 & 8.43 & -2.2 \\ \mbox{H} \rightarrow \mbox{HOOH} & -265.74 & 8.43 & -1.7 \\ \mbox{H} \rightarrow \mbox{HOOH} & -265.74 & 8.43 & -2.2 \\ \mbox{H} \rightarrow \mbox{HOOH} & -265.74 & 8.43 & -2.2 \\ \mbox{H} \rightarrow \mbox{H} \rightarrow \mbox{HOOH} & -265.73 & 9.55 & -2.6 \\ \mbox{CO} & -387.26 & 13.19 & -2.5 \\ \mbox{CO} & -257.31 & 8.67 & -2.0 \\ \mbox{H} \rightarrow \mbox{H} $						
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	0.58 2.34	4.38	9.48	-1.94	-9.84	-0.72
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2.63 5.96	0.13	16.85	-20.40	-31.75	-16.60
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2.12 4.67	-4.34	9.71	-19.07	-24.66	-18.38
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.50 3.38	1.19	5.10	1.16	-4.79	-2.72
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.76 4.51	-9.17	10.84	-34.62	-35.95	-29.18
$\begin{array}{ccccc} \rightarrow \mbox{HCN} & \rightarrow \mbox{HCN} & -311.37 & 8.79 & -0.1 \\ \rightarrow \mbox{HNC} & -296.11 & 8.43 & -0.2 \\ \mbox{CCH} & -265.01 & 6.92 & -3.0 \\ \rightarrow \mbox{HCO} & -277.78 & 7.74 & -1.0 \\ \rightarrow \mbox{HCO} & -257.54 & 7.63 & -0.8 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{F}_2 & 0 & -1.1 \\ \mbox{I} & \rightarrow \mbox{I} & -2.1 \\ \mbox{I} & \rightarrow \mbox{I} & -2.1 \\ \mbox{I} & -2.1 & -2.1 \\ \mbox{I} & -2.1 & -2.1 \\ \mbox{I} & -2.1 & -2.1 \\ \mbox{CF} & -2.1 & -2.1 \\ \$	3.67 1.57	-5.99	5.45	-13.17	-21.55	-17.42
$\begin{array}{cccc} \rightarrow \mbox{HNC} & -296.11 & 8.43 & -0.2 \\ CCH & -265.01 & 6.92 & -3.0 \\ \rightarrow \mbox{HCO} & -277.78 & 7.74 & -1.0 \\ 0 & -257.54 & 7.63 & -0.8 \\ H \rightarrow \mbox{oxirene} & -454.13 & 13.57 & -2.2 \\ H \rightarrow \mbox{oxirene} & -418.18 & 12.75 & -3.8 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -418.18 & 12.75 & -3.8 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -418.18 & 12.75 & -3.8 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -406.69 & 12.92 & -0.9 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -406.69 & 12.92 & -0.9 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -265.74 & 8.43 & 1.7 \\ + \mbox{H} \rightarrow \mbox{HOCN} & -265.74 & 8.43 & 1.7 \\ + \mbox{H} \rightarrow \mbox{HNCO} & -311.49 & 13.93 & -1.7 \\ + \mbox{H} \rightarrow \mbox{HNCO} & -286.23 & 9.55 & 2.5 \\ CF_2 & -257.31 & 8.67 & -2.0 \\ CO_2 & -387.26 & 13.19 & -2.5 \\ \end{array}$	0.10 4.17	-1.98	7.26	-4.76	-14.78	-11.21
$\begin{array}{cccc} {\rm CCH} & -265.01 & 6.92 & -3.0 \\ \rightarrow {\rm HCO} & -277.78 & 7.74 & -1.0 \\ 0 & -257.54 & 7.63 & -0.8 \\ {\rm H} \rightarrow {\rm oxirene} & -454.13 & 13.57 & -2.2 \\ + {\rm H} \rightarrow {\rm HOCN} & -418.18 & 12.75 & -3.8 \\ + {\rm H} \rightarrow {\rm HOCN} & -406.69 & 12.92 & -0.9 \\ + {\rm H} \rightarrow {\rm HOCN} & -265.74 & 8.43 & 1.7 \\ trans-{\rm N}_2{\rm H}_2 & -291.95 & 9.36 & 2.5 \\ + {\rm H} \rightarrow {\rm HNCO} & -431.49 & 13.93 & -1.7 \\ + {\rm H} \rightarrow {\rm HNCO} & -286.23 & 9.55 & 2.5 \\ {\rm CF}_2 & -257.31 & 8.67 & -2.0 \\ {\rm CO}_2 & -387.26 & 13.19 & -2.5 \\ \end{array}$	0.27 3.15	-1.12	6.19	-3.79	-11.79	-8.43
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	3.00 -0.29	8.10	12.32	-3.50	-11.92	3.88
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1.03 3.44	-1.52	7.78	-9.65	-15.85	-10.82
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	).88 4.02	-3.06	7.35	-11.52	-17.57	-13.47
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	2.26 3.48	3.84	15.20	-11.04	-21.83	-7.52
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	3.84 $5.20$	-7.71	11.49	-29.53	-34.22	-26.90
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	0.08 0.14	-4.48	10.72	-14.46	-25.58	-19.68
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	1.72 5.81	0.52	9.06	-5.60	-11.96	-8.03
$\begin{array}{ccccc} + H \rightarrow HN\bar{C}O & -431.49 & 13.93 & -1.7 \\ & \epsilon  cis-N_2H_2 & -286.23 & 9.55 & 2.5 \\ \mathrm{CF}_2 & -257.31 & 8.67 & -2.0 \\ \mathrm{CO}_2 & -387.26 & 13.19 & -2.5 \\ \end{array}$	2.54 $5.60$	12.63	17.18	1.62	-7.26	8.08
$\begin{array}{cccc} & \epsilon  ise N_2 H_2 & -286.23 & 9.55 & 2.5 \\ {\rm CF}_2 & -257.31 & 8.67 & -2.0 \\ {\rm CO}_2 & -387.26 & 13.19 & -2.5 \\ \end{array}$	1.71 6.17	-4.44	11.89	-17.52	-28.44	-20.77
CF <sub>2</sub> -2.0 -257.31 8.67 -2.0 CO <sub>2</sub> -387.26 13.19 -2.5	2.56 5.62	13.28	17.75	2.02	-7.18	8.81
$CO_2$ -387.26 13.19 -2.5	2.01 3.75	0.86	12.40	-16.82	-20.35	-10.68
1	2.57 6.67	-6.71	12.52	-25.10	-34.00	-25.94
• FCCF -385.08 12.97 -4.8	1.89 2.70	-8.52	7.99	-23.24	-30.82	-25.04

			Table (	U.I: (conti	nued)				
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta$ (MP2.8: $\kappa$ -OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(\mathrm{MP2.5})$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$C + 2 O + 2 H \rightarrow dioxirane$	-407.73	13.86	-0.84	6.17	9.11	20.63	-13.46	-23.84	-2.41
$C + F \rightarrow CF$	-131.92	4.70	-1.04	2.01	-0.21	5.80	-8.25	-10.35	-6.21
$2\mathrm{S} + \mathrm{H}  ightarrow \mathrm{SSH}$	-158.96	5.81	0.08	1.45	1.54	4.09	3.33	-4.92	-1.01
$H + O + CI \rightarrow HOCI$	-162.68	6.44	0.92	4.03	-0.42	6.00	-4.08	-9.80	-6.84
$2 \text{ N} + 2 \text{ C} \rightarrow \text{NCCN}$	-498.43	19.60	-2.02	8.48	15.04	31.66	-16.57	-39.19	-1.58
$2\mathrm{N} ightarrow\mathrm{N}_{2}$	-223.58	9.02	1.78	7.07	-0.87	10.50	-6.87	-16.33	-12.24
$2 \mathrm{N} + \mathrm{H} \stackrel{\sim}{ ightarrow} \mathrm{N}_{2} \mathrm{H}$	-220.12	8.61	1.00	4.83	9.16	14.72	-1.84	-10.81	3.60
$O + C + S \rightarrow OCS$	-331.41	13.96	-1.89	5.86	-5.80	10.52	-14.40	-29.19	-22.11
$Si + O \rightarrow SiO$	-186.90	8.13	1.85	7.66	-1.56	10.30	-8.83	-19.22	-13.42
$CI + C + N \rightarrow CICN$	-282.06	11.93	-0.83	5.42	-4.65	8.91	-10.34	-23.05	-18.20
$H + 2.0 \rightarrow HOO$	-172.37	7.45	-0.48	3.95	4.07	11.18	-7.60	-13.76	-3.05
$C + O + N + H \rightarrow HCNO$	-362.44	16.54	0.18	9.90	-3.31	16.03	-15.04	-33.16	-22.65
$C + O + N + H \rightarrow HONC$	-346.90	13.37	-0.16	6.31	-1.97	11.47	-10.57	-22.10	-15.41
$ONH \leftarrow O + N + H$	-201.40	9.32	1.69	6.28	15.44	21.82	-4.51	-13.34	9.06
$H + O + F \rightarrow HOF$	-156.38	7.65	1.12	5.03	1.25	9.17	-5.46	-11.30	-6.68
$H + N + 2 O \rightarrow cis-HONO$	-306.61	15.69	1.44	10.95	4.78	21.04	-15.82	-29.57	-11.49
$H + N + 2O \rightarrow$	-307.09	15 75	1 50	11 07	6.63	22.48	-15.55	-29.82	-9.22
$t_{rans-HONO}$	00.100	0.00	00.1	0.11	0000		00.01		1
	-074 00	15 33	-0.88	5 64	-5 44	8 77	-9 53	-95.95	-10.64
	2014 - 17- 2014 R 1	17.61	-0.00 2.16	10.01	1010- 1010	10.08	20.01	20.67	10.01 
		10.11	01.0	16.91	01.0-	00.61	10-01-	10.00-	67.6T-
	60.701-	9.49	-0.23	3.01	-0.02	0.93	2.03	-13.54	0T.8-
$C + N \rightarrow CN$	-177.09	9.10	-0.58	5.69	17.10	23.79	-6.16	-21.93	10.42
$S + 3 O \rightarrow SO_3$	-320.43	19.29	-1.32	12.55	-7.25	20.91	-34.71	-48.81	-35.42
$C + 2 CI \rightarrow CCI_2$	-172.12	10.17	-2.13	1.82	4.31	11.64	-4.19	-15.86	-3.02
$B + N \rightarrow {}^{3}BN$	-103.84	5.94	-2.74	-0.50	0.08	4.55	-4.86	-10.26	-4.39
$S + 2 O \rightarrow SO_{2}$	-242.81	15.18	0.94	11.92	-2.91	18.10	-21.83	-37.29	-23.92
$N + O \rightarrow NO^{-1}$	-148.93	8.55	0.86	6.23	2.59	12.71	-8.88	-16.60	-7.53
$S + O \rightarrow SO$	-120.04	7.12	-0.42	4.49	-0.66	8.46	-9.36	-17.09	-9.79
$2\mathrm{N}+\mathrm{O} ightarrow\mathrm{N}_{2}\mathrm{O}$	-265.47	17.98	2.65	14.63	-2.15	21.08	-18.37	-38.24	-25.38
$2S \rightarrow S_{2}$	-97.71	6.59	-1.04	1.21	-1.05	3.58	1.05	-9.45	-5.68
$4\mathrm{P}  ightarrow \mathrm{P}_4$	-277.29	20.34	5.24	13.32	0.19	17.60	9.18	-26.41	-17.23
$2 \operatorname{Cl}  ightarrow \operatorname{Cl}_{2}$	-55.39	4.31	-0.13	1.35	-1.45	1.85	-0.57	-5.75	-4.75
$20  ightarrow 0_{2}^{-}$	-118.26	7.70	-0.97	4.71	-0.80	10.12	-12.43	-19.33	-11.73
$2F \rightarrow F_2^-$	-37.02	7.34	-0.20	3.51	23.99	27.18	-5.28	-11.51	20.81
$2S + O \rightarrow S_2O$	-191.58	15.50	-1.03	7.95	12.03	23.95	-9.96	-34.17	0.11
$2P  ightarrow P_2$	-108.52	9.03	1.19	3.77	19.02	21.08	12.47	-8.63	16.96
$CI + F \rightarrow CIF$	-60.01	5.01	0.32	3.23	-0.35	5.54	-5.47	-9.27	-6.24
$N + 2 O \rightarrow NO_2$	-222.79	17.98	1.37	13.53	0.16	22.37	-20.91	-39.11	-22.06
$CI + O \rightarrow CIO$	-60.64	5.83	-1.84	1.83	4.62	9.11	-4.98	-12.58	0.13
$H + CIO \rightarrow HOCI$	-102.04	0.61	2.76	2.20	-5.03	-3.10	0.00	2.78	-6.96
$OH + CI \rightarrow HOCI$	-56.69	4.88	0.42	3.09	-1.11	4.41	-4.45	-8.96	-6.63
$OH + F \rightarrow HOF$	-50.39	6.09	0.62	4.09	0.56	7.59	-5.83	-10.46	-6.47
$\rm NO + O \rightarrow NO_2$	-73.86	9.42	0.50	7.31	-2.43	9.67	-12.02	-22.51	-14.53
$NO + N \rightarrow N_2 \bar{O}$	-116.54	9.42	1.78	8.40	-4.74	8.37	-9.49	-21.64	-17.85
$N_2 + O \rightarrow N_2 O$	-41.89	8.96	0.87	7.56	-1.28	10.58	-11.50	-21.91	-13.14

$\begin{array}{c c} (CCSD) & \Delta(MP2.8:\kappa- \ \Delta(MP3:\kappa- \ \Delta(M)2)) \\ OOMP2) & OOMP2) \end{array}$
-2.22 0.114 -2.46
5.64 0.14 $3.91$
-0.68 0.32 -2.54
7.37 -1.21 2.50
5.24 -1.77 1.50 7.05 1.83 3.00
8.26 0.76 5.63
3.16 -0.09 2.05
4.44 -1.73 1.09
-0.81 0.95 -0.74
2.59 -1.06 0.95
5.84 -1.57 3.26
3.92 -1.29 0.07
0.75 -0.56 -0.92
1.55 -0.50 -0.15
4.08 -1.83 1.45
2.74 -1.08 1.23
-0.23 0.74 -0.07
6.62 $0.71$ $4.60$
0.11 -0.15 -0.58
6.80 -2.18 2.76
5.84 -0.65 2.03
5.56 -1.69 $2.65$
3.26 -0.54 1.07
-1.01 -0.00 -1.04 2.08 -0.61 0.65
-2.28 -0.44 -2.07
0.45 1.10 0.87
1.12 $0.86$ $1.40$
0.00 $1.21$ $0.44$
6.09 -2.67 1.01
-0.70 -0.48 -1.84
-0.31 0.48 -1.52
7.85 -1.25 3.50
1.39 -0.85 -2.91
1.23 $0.51$ $0.51$
8.97 -0.84 6.08
7.17 -0.85 4.38
0.98 2.20 1.87
5.32 0.73 3.94
4.73 -0.79 0.79
-0.78 1.12 0.23

$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta(MP2.8:\kappa-OOMP2)$	$\Delta$ (MP3: $\kappa$ -OOMP2)	$\Delta(\mathrm{MP2.5})$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + O \rightarrow HOO$	-66.38	5.89	-0.98	3.01	3.38	9.60	-7.97	-12.92	-2.84
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$0.0 + H \rightarrow HOO$	-54.12	-0.25	0.49	-0.76	4.87	1.06	4.83	5.58	8.68
	$0$ H $+$ H $\rightarrow$ H $_{\circ}$ O	-124.64	1.73	0.54	1.80	-0.89	1.90	-2.82	-4.15	-3.68
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$I_{3} + NH \rightarrow \tilde{H}NNN$	-19.61	7.74	0.44	5.19	-1.62	6.95	-6.66	-16.80	-10.15
$\begin{split} NH^+ \wedge NH_3^+ H_3 & 12731 & 725 & 0.06 & 4.29 & 7.87 & 13392 & 5.30 & 114 \\ H_3 + N + NH_3 & 13730 & 756 & 0.00 & 1197 & 1171 & 0.88 & 0.03 \\ H_3 + H + NH_3 & 14464 & 11341 & 0.36 & 0.173 & 0.054 & 0.08 & 0.173 & 0.04 & 1131 & 0.57 & 0.054 & 0.034 & 0.08 & 0.173 & 0.06 & 1131 & 0.57 & 0.054 & 0.08 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.173 & 0.06 & 0.17 & 0.08 & 0.017 & 0.126 & -7.04 & 0.06 & 0.17 & 0.08 & 0.017 & 0.036 & 0.037 & 0.036 & 0.017 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.037 & 0.036 & 0.017 & 0.036 & 0.057 & 0.036 & 0.017 & 0.026 & 0.036 & 0.013 & 0.064 & 0.044 & 0.086 & 0.013 & 0.036 & 0$	${   \check {   N}}{   H}_{ m o}  ightarrow {   N}_{ m o}{   H}_{ m o}$	-72.18	3.30	-0.18	1.87	-2.68	2.01	-6.57	-7.89	-7.37
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$NH^{-} \rightarrow trans-N_{2}H_{2}$	-127.31	7.25	0.66	4.29	7.87	13.92	-5.30	-12.18	1.81
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{HH} + \mathrm{N}  ightarrow \mathrm{N_{2}H}^{2}$	-137.80	7.56	0.06	4.18	6.77	13.08	-5.30	-13.27	0.46
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{IH}_{\mathrm{o}} + \mathrm{Cl}  ightarrow \mathrm{\tilde{NH}}_{\mathrm{o}}\mathrm{Cl}$	-63.58	4.20	0.00	1.97	-1.71	2.61	-3.42	-7.20	-6.03
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$(\mathrm{H}_{2}^{2}+\mathrm{H} ightarrow\mathrm{NH}_{3}^{2}$	-114.64	1.34	0.36	1.00	-0.54	0.88	-0.94	-2.04	-1.96
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{IH}^+ \mathrm{H} \to \mathrm{NH}_{\circ}^{\circ}$	-98.52	1.27	0.56	0.75	0.52	0.86	1.11	-0.04	0.17
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$(H, C + {}^{3}CH, \xrightarrow{f} allene$	-154.14	4.27	-0.44	1.31	-5.99	-1.73	-4.33	-7.22	-10.24
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$C\tilde{H} + CH_3 \xrightarrow{\sim} propyne$	-133.45	2.74	-0.10	1.16	-13.77	-9.09	-4.49	-5.23	-18.46
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H_3CH + CH_3 \rightarrow propene$	-108.59	2.65	0.17	1.29	-7.24	-4.06	-3.41	-4.35	-10.42
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$(\mathrm{H}_{2}^{-}\mathrm{CH} + \mathrm{F} \rightarrow \mathrm{C}_{2}\mathrm{H}_{3}\mathrm{F}$	-127.51	3.45	-0.17	2.35	-7.49	-1.49	-8.57	-8.98	-13.49
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$CH_3 \rightarrow C_2H_6$	-97.57	2.55	-0.29	0.73	-1.44	0.87	-2.82	-4.03	-3.75
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$^{3}\mathrm{CH}_{2}  ightarrow \mathrm{C}_{2}\mathrm{H}_{4}$	-180.21	5.31	0.82	2.28	-0.35	2.88	1.17	-4.27	-3.57
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	${ m H}_2{ m C}+{ m H} ightarrow{ m CH}_2{ m CH}$	-86.41	-0.35	-0.71	-0.83	1.51	0.95	-1.85	-0.29	2.07
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$_2{ m H}_2 + { m H}  ightarrow { m CH}_2{ m CH}$	-41.07	-1.75	-0.02	-1.70	7.86	3.37	2.52	6.55	12.30
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$CH_2 + C \rightarrow CH_2C$	-167.91	5.72	-0.60	0.82	3.05	5.80	1.62	-4.97	0.3(
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$CH + F \rightarrow HCCF$	-133.02	3.62	-0.66	1.86	-14.08	-6.87	-9.67	-9.63	-21.3(
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + CF \rightarrow HCCF$	-181.33	4.90	-3.78	-1.11	-11.69	-5.53	-9.12	-14.40	-17.8
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\rm CH + H \rightarrow C_2 H_2$	-140.31	1.08	0.71	0.75	-11.35	-9.46	0.94	0.26	-13.25
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\rm CH  ightarrow \rm C_2 H_2$	-235.77	6.12	-4.59	-0.89	-15.07	-7.48	-10.94	-18.07	-22.6(
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{H}_3 + \mathrm{F}  ightarrow \mathrm{CH}_3 \mathrm{F}$	-114.64	3.23	-0.24	2.14	-1.78	3.46	-7.62	-8.42	-7.0
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	${ m H}_3+{ m H} ightarrow{ m CH}_4$	-112.18	1.00	0.56	0.57	0.70	0.64	1.62	0.70	0.75
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{CH}_2 + \mathrm{H} \xrightarrow{\sim} \mathrm{CH}_3$	-116.17	0.88	0.86	0.64	1.19	0.57	2.63	1.82	1.8
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	${ m H}+{ m H} ightarrow { m ^3CH_2}$	-107.30	-0.05	-2.15	-1.90	-5.86	-5.70	-3.99	-3.06	-6.0
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{F} + \mathrm{F}  ightarrow \mathrm{CF}_2$	-125.39	3.98	-0.97	1.74	1.06	6.60	-8.57	-10.01	-4.4
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + CI \rightarrow CICN$	-104.97	2.83	-0.25	-0.27	-21.75	-14.89	-4.19	-1.12	-28.6
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	${ m BH}_3  ightarrow { m B}_2 { m H}_6$	-46.08	2.73	-0.65	-0.20	-0.22	0.76	-0.25	-2.01	-1.1
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$_2 + \mathrm{O}  ightarrow \mathrm{S}_2\mathrm{O}$	-93.87	8.91	0.01	6.73	13.08	20.37	-11.01	-24.71	5.78
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$+$ SO $\rightarrow$ S <sub>2</sub> O	-71.54	8.38	-0.61	3.45	12.69	15.49	-0.60	-17.08	9.9(
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	${\rm O}_2+{\rm O} ightarrow{ m SO}_3$	-77.62	4.12	-2.26	0.63	-4.35	2.81	-12.88	-11.52	-11.5(
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$0 + 0  ightarrow \mathrm{SO}_2$	-122.77	8.05	1.36	7.43	-2.25	9.64	-12.46	-20.21	-14.13
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{P}_2  ightarrow \mathrm{P}_4$	-60.26	2.27	2.86	5.78	-37.84	-24.55	-15.75	-9.15	$-51.1^{\circ}$
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$\mathbf{F} + \mathbf{CH}_3 \rightarrow \mathbf{H} + \mathbf{CH}_3 \mathbf{F}$	25.74	1.25	-0.68	0.11	-0.90	0.83	-3.51	-3.16	-2.6
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$H + HCI \rightarrow H + HOCI$	49.70	3.37	0.05	2.50	-1.02	4.02	-5.89	-8.42	-6.0
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$H + HF \rightarrow H + HOF$	89.99	4.11	0.19	2.07	1.44	4.95	-1.72	-5.21	-2.0′
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$CI + O \rightarrow H + CIO$	45.75	4.31	-2.21	1.24	4.70	8.71	-6.42	-12.03	0.69
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + CI \rightarrow H + CIO$	45.36	4.27	-2.34	0.89	3.93	7.52	-5.35	-11.74	0.33
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{H} + \mathrm{NO} \rightarrow \mathrm{H} + \mathrm{NO}_2$	32.13	7.87	0.01	6.37	-3.12	8.08	-12.40	-21.68	-14.32
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + O_2 \rightarrow H + NO_2$	-22.21	9.22	1.40	8.17	-1.42	10.62	-11.94	-22.24	-13.47
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\mathrm{OH} + \mathrm{N}_2 \rightarrow \mathrm{H} + \mathrm{N}_2\mathrm{O}$	64.11	7.40	0.37	6.62	-1.97	8.99	-11.87	-21.08	-12.94
$0.11 + HNO \rightarrow 0.31 + 8.8 -0.69 -3.85 -9.50 -0.93 -11.41 -17$	$(\mathrm{H}+\mathrm{NO} ightarrow\mathrm{H}+\mathrm{N}_2\mathrm{O}$	-34.22	8.37	0.85	7.75	-7.12	6.74	-12.95	-24.10	-20.98
	$\rightarrow ONH + HO$	0.31	4.88	-0.69	3.85	-9.50	-0.93	-11.41	-15.64	-18.07

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table (	C.1: (cont	inued)				
uc	CCSD(T)	$\Delta(CCSD)$	$\Delta$ (MP2.8: $\kappa$ - OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(\mathrm{MP2.5})$	$\Delta({ m MP3})$	$\Delta(\kappa - OOMP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$HOO \rightarrow H +$	-52.39	7.25	1.04	6.47	0.18	9.66	-11.42	-18.53	-9.31
$HCO \rightarrow H + HNCO$	-71.39	5.13	-1.61	2.07	-5.30	2.48	-11.33	-15.06	-13.09
$HO \rightarrow H + HNCO$	-145.31	3.67	-4.55	-0.79	-25.79	-15.10	-17.20	-18.31	-36.48
$HCN \rightarrow H + HNCO$	-14.12	3.59	-2.10	1.05	-3.15	3.05	-13.13	-12.83	-9.35
$OH \rightarrow H + HNO$	-13.08	6.71	0.26	4.69	12.37	18.60	-8.35	-14.96	6.13
$O \rightarrow H + NO$	-66.61	7.50	-0.07	5.57	0.20	11.08	-12.34	-19.06	-10.67
$N \rightarrow H + NO$	-42.93	7.00	0.37	5.29	1.90	11.12	-9.26	-15.76	-7.33
$C_2H_2 \rightarrow H + C_2H_2$	57.20	4.01	-0.47	2.07	6.40	10.76	-8.86	-9.33	2.04
${}^{}_{2}\mathrm{H}_{4}  ightarrow \mathrm{H} +$	18.90	2.53	-1.12	1.75	-3.70	2.50	-12.72	-11.41	-9.90
$H_A \rightarrow H +$	14.23	2.16	-0.65	1.48	-2.87	1.96	-8.29	-8.55	-7.71
$\mathrm{hol}^{2}$ $\mathrm{H}_{2}$ $\mathrm{O} \rightarrow \mathrm{H} +$	-5.26	2.70	-2.27	-0.71	-1.26	1.44	-6.03	-7.27	-3.97
IČOH	00 00	07.7	01.0		1	10 0	00 1	C C	60 F
$CH_2 \rightarrow H + COH$	-22.60	4.49	0.42	3.00	3.70	9.04	-4.86	-8.35	-1.63
$trans-HCOH \rightarrow$	-72.19	3.40	-2.01	1.69	-9.19	-1.32	-13.98	-15.09	-17.07
HCO $\rightarrow$ H +	9.55	0.86	-1.60	-0.85	-2.34	-0.90	-5.79	-4.14	-3.77
acid $HOOH \rightarrow H +$	-148.34	2.52	-4.47	-1.16	-10.49	-3.87	-16.49	-16.22	-17.10
acid $CH_2C \rightarrow H +$	-66.95	4.28	-2.07	2.32	-8.75	0.75	-16.30	-17.80	-18.26
$I_2 CO \rightarrow H +$	-74.95	3.99	-4.03	-1.19	-9.55	-3.96	-11.38	-14.16	-15.13
- xirene → H +	-72.78	111	-0.87	1.54	-4.40	0.07	-9.73	-9.08	-8.87
-									
formic acid $\rightarrow$ H +	-49.26	3.40	-2.18	-0.05	-1.73	1.32	-6.70	-9.98	-4.78
methanol $\rightarrow$ H +	-80.52	3.64	-3.09	-0.38	-8.01	-2.67	-10.01	-12.48	-13.35
$^{1}_{2}\mathrm{H}_{4} \rightarrow \mathrm{H} +$	-7.13	2.42	-0.99	1.94	-2.80	3.32	-11.22	-11.18	-8.91
$F_{3} \rightarrow H + F_{3}CO$	-54.88	2.53	-2.32	0.51	-9.26	-2.50	-13.08	-13.03	-16.02
$ICO \rightarrow H + \tilde{H}COF$	16.33	0.76	-1.52	-0.80	-1.94	-0.71	-5.30	-3.56	-3.17
$CH_2 \rightarrow H + H_2CO$	-75.13	5.07	0.21	3.66	-2.07	5.34	-7.49	-11.81	-9.47
$I_2O \rightarrow H + H_2O$	-57.79	3.28	-2.48	-0.04	-7.03	-2.25	-8.66	-10.73	-11.81
$\mathrm{H} \rightarrow \mathrm{H} + \mathrm{HCO}$	-87.00	5.24	-2.68	1.83	-8.12	1.03	-14.21	-18.21	-17.26
$S \rightarrow H + CS_2$	-19.77	4.76	-1.52	1.37	-6.30	0.93	-7.86	-13.51	-13.53
$2 \rightarrow \Pi + Co_2$ S $\downarrow \Pi + OCS$	-91.03	1.80	-0.99	3.73 1.31	-10.30	0.07 2.00	-1.18 -16.81	-19.00 -14.82	-20.02 -13.75
$O \rightarrow H + OCS$	-126.58	5.89	-2.62	0.69	-0.01	-3.12	-10.01	-15.31	-18.98
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Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

Table C.1: $CCSD(T)  \Delta(CCSD)  \Delta(MP2.8:\kappa - \Delta(0))$	Table C.1: $\Delta(CCSD)  \Delta(MP2.8:\kappa - \Delta(0))$	Table C.I: $\frac{\Delta(MP2.8:\kappa-\Delta(1))}{\Delta(MP2.8)}$		(conul <u>MP3:6-</u>	nued) $\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - \Delta)$	$\Delta(OOMP2)$	$\Delta(MP2)$
00MP2) 00 12.70 5.25 -1.88	$0.000 \pm 0.000$ $0.000 \pm 0.000$		Ŭ	1.18 1.18	-4.22	2.25	00MP2) -6.17	-13.42	-10.70
-23.72 $4.00$ $-2.19$	4.00 -2.19	-2.19		1.72	-4.34	3.58	-13.95	-15.60	-12.27
-184.23 4.55 -2.75	4.55 -2.75	-2.75		1.29	-11.82	-2.77	-16.86	-17.87	-20.86
-82.91 8.55 -1.38	8.55 -1.38	-1.38		2.93	-6.53	1.76	-2.16	-16.74	-14.81
-81.12 8.42 -1.11	8.42 -1.11	-1.11		2.95	-2.10	6.02	-1.26	-15.34	-10.22
-113.32 $7.69$ $0.24$	7.69 0.24	0.24		6.52	-4.97	7.98	-13.03	-22.77	-17.91
-80.91 $6.58$ $1.35$	6.58 1.35	1.35		6.73	-2.25	8.72	-9.20	-18.38	-13.21
-172.76 6.69 -2.03	6.69 -2.03	-2.03		3.34	-8.97	2.18	-15.71	-20.77	-20.12
-151.54 6.08 -1.38	6.08 -1.38	-1.38		3.08	-3.75	5.77	-11.90	-16.73	-13.26
-144.27 6.79 -2.19	6.79 -2.19	-2.19		2.84	-10.27	0.46	-12.41	-20.44	-21.00
-92.31 8.16 -1.73	8.16 -1.73	-1.73		5.01	11.19	18.62	-10.35	-25.13	3.76
-94.77 8.04 -1.52	8.04 -1.52	-1.52		5.04	14.72	22.16	-9.62	-24.39	7.28
13.23 1.23 -1.39	1.23 -1.39	-1.39		-0.69	-2.60	-0.93	-3.26	-4.24	-4.26
-46.02 2.57 -2.42	2.57 -2.42	-2.42		-0.84	-7.66	-4.08	-6.68	-8.83	-11.24
-13.61 1.34 -1.38	1.34 -1.38	-1.38		-0.50	-2.60	-0.52	-3.47	-4.88	-4.69
-86.43 2.84 -1.80	2.8480	-1.80		136	-7.41	-0.49	-13.61	-13.55	-14.33
-20.03 1.15 -0.90	1 15 -0 90	-0.90		0.26		0.84	-4.84	-4.89	-3.95
70.92 2.52 0.10	2 K2 0 10	0.10		0 1 0	0.03	100	2 53	613	3 03
	0.00 1 PA	01.10		FT:0	64 1	00.F	10.0-	01.0-	-0.00 F
	00.1- 00.0	-1.00		0.10	-1.40 2 2 2	7.71	-0.21	10.0-	-0.12
39.62 4.34 -1.47	4.34 -1.47	-1.47		2.07	2.69	8.01	-8.34	-12.08	-2.64
-16.08 1.73 -1.36	1.73 -1.36	-1.36		0.60	-3.94	0.60	-8.46	-8.89	-8.47
-28.79 $5.98$ $0.10$	5.98 0.10	0.10		3.54	7.35	13.06	-6.42	-12.14	1.64
-57.57 $3.40$ $-3.20$	3.40 -3.20	-3.20		-0.22	-10.10	-3.73	-10.91	-14.41	-16.48
						5			
9 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 GA 0 0 0	20 C		0 76	100	60.6	1 65	10 60	10.75
-04.03 0.04 -2.00	0.04	00.7-		01.0-	+0·0-	-0.34	00.1-	ec.01-	C1.21-
		0000		00 0	500		1		000
12.80 1.48 -0.00	1.48 -U.0U	-0.00		0.33	10.0-	-4.12	64.4-	-3.12	-9.09
							0		
-00.25 3.77 -3.24	3.77 -3.24	-3.24		-1.10	-8.25	-3.94	-8.00	-10.83	-12.50
-59.16 $3.38$ $-2.75$	3.38 -2.75	-2.75		-0.83	-8.09	-4.02	-7.08	-9.84	-12.16
-53.37 $3.55$ $-3.31$	3.55 -3.31	-3.31		-1.32	-2.54	1.01	-7.06	-10.27	-6.09
-42.99 3.35 -4.03	3.35 -4.03	-4.03		-1.73	-15.29	-9.58	-10.61	-13.67	-21.01
7.35 1.64 -1.10	1.64 $-1.10$	-1.10		-0.17	-13.20	-9.51	-5.56	-4.37	-16.90
-128.47 6.17 -2.43	6.17 -2.43	-2.43		1.02	-9.21	-1.79	-6.95	-15,01	-16.64
21 DK A D1 2 1K	A 01 2 15	2 15		16.0	110	1.69	68.7	13 39	601
		01.0-		17.0-	0T-0-	40.F	70	70.01-	40.H-
-135.51 $0.39$ $-2.96$	0.39 -2.96	-2.96		-0.44	-29.05	-19.95	-15.74	-12.05	-38.14
14.99 2.00 -1.41	2.00 -1.41	-1.41		-0.29	1.95	3.96	-4.46	-4.75	-0.07
1.25 -0.05 -4.05	-0.05 -4.05	-4.05		-3.10	-8.90	-6,99	-9.63	-7.81	-10.81
1 20 1 20 0 50	1 20 0 50	0 50		0.69	0.06	999	R OR		01.6
	00'0- 70'T	-000		20.0	00.0-	2.00	-0.0-	-4.41	-4.43
-65.47 2.86 -1.76	2.86 -1.76	-1.76		1.50	-4.35	2.69	-13.37	-13.22	-11.39
1.41 1.31 -0.62	1.31 -0.62	-0.62		-0.86	-21.66	-15.28	-5.63	-0.57	-28.05
8.45 2.72 -1.47	2.72 -1.47	-1.47		-0.02	0.68	3.17	-4.13	-5.09	-1.82
-47.15 3.75 -2.18	3.75 -2.18	-2.18		1.33	-6.12	0.63	-12.44	-13.55	-12.86
	2010 101	90 C		100	но - 1 О 1	100 0	Г И С И	ни 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	110
00.2- T0.1 C0.14-	00.2- 10.1	-2.00		-0.34	-1.40	0.43	00.0-	-0.04	-4.10

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta(MP2.8:\kappa-OOMP2)$	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - \Omega)$	$\Delta(00MP2)$	$\Delta(MP2)$
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	:H + F → H + BF	-96.30	3 23			97 0-	-3.65		-14.13	-15.87
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$		18.03	0.10 7.20	01.0 A 60	1 0.0	0 12 00 - 00 - 00	0.00	0.07	13.60	11.68
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	ити – ит – ли и – в – и – Зри	60.0T-	00 1	00.4-	7 - T - 2 - 7 - 1 - 1 - E	00.0-	0.7- 10 c	10.6-	00.01-	00.111-
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$		70.12-	4.09	-0.00	01.1-	10.2-	16.7 1 1 1 0	70.0-	-12.12	10.1
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$S + SO \rightarrow H + S_2O$	15.03	7.30	-1.48	2.79	11.21	14.58	-3.90	-18.88	( <b>8</b> . 7
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + S_2 \rightarrow H + S_2O$	12.12	7.36	-0.49	5.80	12.39	18.78	-11.39	-23.88	5.99
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + SO_2 \rightarrow H + SO_3$	28.38	2.56	-2.76	-0.31	-5.04	1.23	-13.26	-10.68	-11.30
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + SO \rightarrow H + SO_2$	-16.77	6.50	0.87	6.49	-2.94	8.05	-12.84	-19.37	-13.92
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$S + O_{2} \rightarrow H + SO_{2}$	-37.99	6.40	1.04	6.56	-3.59	7.07	-12.70	-19.76	-14.24
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$S + O \rightarrow H + SO$	-33.48	6.05	-1.29	3.83	-2.15	7.55	-12.66	-18.88	-11.84
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$H + S \rightarrow H + SO$	-14.05	5.57	-0.92	3.56	-1.35	6.88	-9.74	-16.25	-9.58
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$S + S \rightarrow H + S$	-11.15	5.51	-1.91	0.55	-2.53	2.67	-2.25	-11.25	-7.73
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$CI + CI \rightarrow H + CI_{o}$	51.00	2.80	-0.50	0.75	-1.37	1.46	-2.02	-5.20	-4.19
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$CI + F \rightarrow H + CIF$	46.37	3.50	-0.05	2.64	-0.26	5.14	-6.92	-8.73	-5.67
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$F + CI \rightarrow H + CIF$	80.36	3.03	-0.11	1.21	0.53	2.90	-1.36	-4.02	-1.84
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$F + F \rightarrow H + F$ ,	103.35	5.37	-0.64	1.48	24.87	24.54	-1.16	-6.25	25.21
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + O \rightarrow H + O_{2}$	-12.26	6.14	-1.47	3.77	-1.49	8.54	-12.80	-18.50	-11.52
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$H + N \rightarrow H + N_{2}$	-141.26	7.96	0.84	6.41	-3.25	8.87	-10.33	-18.79	-15.37
$\begin{array}{llllllllllllllllllllllllllllllllllll$	$\mathbf{F} + \mathbf{CH}_{s} \rightarrow \mathbf{C} + \mathbf{CH}_{s} \mathbf{F}$	17.29	-1.46	0.79	0.13	-1.57	-2.34	0.62	1.93	-0.81
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$O + HF \rightarrow C + HOF$	241.53	-1.97	1.57	-1.01	5.19	-0.82	10.18	11.53	11.19
$\begin{array}{llllllllllllllllllllllllllllllllllll$	$O + CI \rightarrow C + CIO$	196.90	-1.81	-0.96	-2.19	7.67	1.75	6.54	4.99	13.59
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$O + NO \rightarrow C + NO$ ,	183.68	1.79	1.38	3.29	0.63	2.31	-0.50	-4.95	-1.06
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + O_{2} \rightarrow C + NO_{2}$	72.56	1.17	2.92	3.14	-16.14	-11.54	-2.32	2.15	-20.75
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	$O + N_2 \rightarrow C + N_2 \tilde{O}$	215.65	1.32	1.75	3.54	1.78	3.23	0.02	-4.35	0.32
$\begin{array}{llllllllllllllllllllllllllllllllllll$	$N + NO \rightarrow C + N_{2}O$	60.55	0.32	2.37	2.71	-21.84	-15.42	-3.33	0.29	-28.26
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + HNO \rightarrow C + C$	151.85	-1.20	0.69	0.77	-5.75	-6.69	0.48	1.09	-4.81
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	ans-HONO									
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + HOO \rightarrow C +$	42.38	-0.80	2.57	1.43	-14.54	-12.50	-1.80	5.87	-16.59
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	ans-HONO									
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + HCO \rightarrow C + HNCO$	23.38	-2.91	-0.09	-2.96	-20.02	-19.68	-1.72	9.33	-20.37
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + HCN \rightarrow C + HNCO$	137.42	-2.49	-0.73	-2.03	0.60	-2.72	-1.24	3.90	3.91
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + OH \rightarrow C + HNO$	81.69	-1.34	1.78	-0.34	-2.35	-3.56	1.27	9.43	-1.15
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$N + O \rightarrow C + NO$	28.16	-0.55	1.45	0.54	-14.52	-11.08	-2.73	5.33	-17.95
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + N \rightarrow C + NO$	108.61	0.92	1.74	2.21	5.65	5.36	2.64	0.97	5.94
$ \begin{array}{c} \mbox{irene} \\ 0 + C_2 H_4 \rightarrow C + & 170.44 & -3.55 & 0.26 & -1.33 & 0.05 & -3.27 & -0.83 & 5.32 & 3.5 \\ 0 + C H_4 \rightarrow C + & 165.78 & -3.92 & 0.72 & -1.60 & 0.87 & -3.80 & 3.60 & 8.18 & 5. \\ 0 + chanol \\ 0 + ^3 CH_2 \rightarrow C + & 128.95 & -1.59 & 1.80 & -0.08 & 7.45 & 3.27 & 7.03 & 8.38 & 11. \\ 0 + trans-HCOH \rightarrow & 79.35 & -2.68 & -0.63 & -1.39 & -5.44 & -7.08 & -2.08 & 1.65 & -3. \\ 1.65 & 0.6 + trans-cid \\ 0 + C H_2 C \rightarrow C + ketene & 84.59 & -1.80 & -0.69 & -0.76 & -5.01 & -5.02 & -4.41 & -1.07 & -5.02 \\ 0 + C H_2 C \rightarrow C + ketene & 84.59 & -1.80 & -0.69 & -0.76 & -5.01 & -5.02 & -4.41 & -1.07 & -5.02 \\ 0 + c nire a cid \\ 0 + c nire a cid$	$O + C_2H_2 \rightarrow C + C$	208.74	-2.07	0.91	-1.01	10.15	4.99	3.03	7.40	15.30
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	irene									
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + C_2 H_4 \rightarrow C + $	170.44	-3.55	0.26	-1.33	0.05	-3.27	-0.83	5.32	3.36
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + CH_A \rightarrow C +$	165.78	-3.92	0.72	-1.60	0.87	-3.80	3.60	8.18	5.55
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	ethanol									
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$O + {}^{3}CH_{2} \rightarrow C +$	128.95	-1.59	1.80	-0.08	7.45	3.27	7.03	8.38	11.63
+ formic acid 0+CH <sub>2</sub> C $\rightarrow$ C+ ketene 84.59 -1.80 -0.69 -0.76 -5.01 -5.02 -4.41 -1.07 -5. 0+cme $\rightarrow$ C+ 78.76 -4.97 0.51 -1.54 -0.65 -5.70 2.17 7.65 4.	$O + trans-HCOH \rightarrow$	79.35	-2.68	-0.63	-1.39	-5.44	-7.08	-2.08	1.65	-3.80
$ \begin{array}{ccccc} 0 + \mathrm{CH}_2\mathrm{C} \rightarrow \mathrm{C} + \mathrm{ketene} & 84.59 & -1.80 & -0.69 & -0.76 & -5.01 & -5.02 & -4.41 & -1.07 & -5.0 \\ 0 + \mathrm{oxtrene} \rightarrow \mathrm{C} + & 78.76 & -4.97 & 0.51 & -1.54 & -0.65 & -5.70 & 2.17 & 7.65 & 4.00 \\ \end{array} $	+ formic acid	0 1 - 0		4			1	:	1	1
$O + 0$ oxirene $\rightarrow C + (8.76) -4.97 0.51 -1.54 -0.05 -5.70 2.17 7.05 4.$	$O + CH_2 C \rightarrow C + \text{ketene}$	84.59 70 70	-1.80	-0.69	-0.76	-5.01	-5.02	-4.41	-1.07	-5.00
	$O + \text{oxirene} \rightarrow C +$	78.76	-4.97	0.51	-1.54	-0.65	-5.70	2.17	7.65	4.35

Appendix C.Additional Information: Third-Order Møller-Plesset Perturbation TheoryMade Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy240

			Table (	C.1: (conti	nued)				
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta(MP2.8:\kappa-OOMP2)$	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(\mathrm{MP2.5})$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(OOMP2)$	$\Delta(MP2)$
$CO + C_2H_4 \rightarrow C + C_4$	144.42	-3.66	0.39	-1.14	0.95	-2.45	0.68	5.55	4.35
acetaldenyde CO + CF C + F CO	06.67	ນ ນ ຕ	-0 0K	-0 R7	ц Ц	8 96 8	1 10	3 70	-076
$CF + HCO \rightarrow C + F_2 CO$	78.0	-3.33	-0.05	10.2-	-2.62	-0.20	-1.17	0.10 1.53	-2.70
$CO + {}^{3}CH, \rightarrow C + H_{0}CO$	76.41	-1.01	1.59	0.58	1.68	-0.43	4.41	4.92	3.79
$CO + CH \rightarrow C + HCO$	64.54	-0.84	-1.30	-1.25	-4.37	-4.74	-2.32	-1.48	-4.00
$2 \text{ CS} \rightarrow \text{C} + \text{CS},$	61.35	-3.65	-0.41	-1.59	-4.20	-5.09	-6.60	1.83	-3.32
$CO + CS \rightarrow C + OCS$	93.82	-3.17	-0.78	-1.77	-2.12	-3.76	-4.91	1.92	-0.48
$2\mathrm{CO} \rightarrow \mathrm{C} + \mathrm{CO}_2$	127.82	-2.08	-0.81	-1.36	-0.60	-2.19	-2.05	1.14	1.00
$CO + Si \rightarrow C + \overline{SiO}$	70.63	0.50	2.73	3.65	1.50	2.95	2.69	-1.65	0.05
$CN + CH \rightarrow C + HCN$	-49.50	-1.25	-0.67	-2.19	-24.99	-21.70	-2.80	3.95	-28.28
$CS + H_2 \rightarrow C + H_2S$	94.36	-7.18	-0.28	-3.64	-0.50	-6.95	-2.00	11.10	5.95
$CF + AI \rightarrow C + AIF$	-28.48	-1.57	0.56	0.28	-2.23	-2.33	-0.71	0.20	-2.13
$CO + H_2O \rightarrow C + HOOH$	222.43	-2.50	1.56	-0.94	3.78	-1.77	8.38	10.60	9.33
$CS + HS \rightarrow C + SSH$	95.29	-4.76	-0.56	-2.82	0.67	-3.75	-2.01	6.82	5.10
$CO + OH \rightarrow C + HOO$	191.16	-1.74	-0.10	-1.01	6.44	2.25	3.55	4.65	10.63
$CO + H_2 \rightarrow C + H_2O$	135.46	-4.34	0.02	-2.48	-0.19	-5.17	3.44	7.84	4.79
$CN + N\tilde{H}_2 \rightarrow C + \tilde{c}$	65.98	-2.07	1.62	-1.50	-7.37	-9.10	3.20	12.25	-5.64
$trans$ -N $_2 H_2$									
$CF + CH_2CH \rightarrow C + C_2 + C_2 + C_2$	4.41	-1.24	0.87	0.35	-7.28	-7.29	-0.32	1.37	-7.28
$CF + CCH \rightarrow C + HCCF$	-1.10	-1.08	0.37	-0.15	-13.88	-12.67	-1.43	0.72	-15.09
$2CF \rightarrow C + CF$ ,	6.54	-0.72	0.06	-0.26	1.27	0.79	-0.33	0.34	1.75
$CF + Si \rightarrow C + SiF$	-7.13	-1.40	0.89	0.64	-0.73	-0.79	-0.92	0.68	-0.68
$CF + B \rightarrow C + BF$	-50.27	-0.91	-0.59	-0.92	-2.62	-2.88	-1.41	-0.45	-2.37
$CN + B \rightarrow C + {}^{3}BN$	73.25	-3.16	-2.15	-6.19	-17.03	-19.25	1.30	11.67	-14.80
$CS + SO \rightarrow C + S_2O$	96.15	-1.12	-0.38	-0.16	13.31	8.56	-2.63	-3.54	18.06
$CO + S_2 \rightarrow C + S_2 O$	163.67	1.28	0.89	2.72	16.13	13.02	0.51	-7.15	19.25
$CO + SO_3 \rightarrow C + SO_3$	179.92	-3.52	-1.38	-3.39	-1.29	-4.54	-1.36	6.05	1.96
$CO + SO \rightarrow C + SO_2$	134.77	0.42	2.24	3.41	0.81	2.28	-0.94	-2.64	-0.66
$CS + O_2 \rightarrow C + SO_2^-$	43.13	-2.02	2.15	3.60	-1.49	1.05	-11.43	-4.42	-4.03
$CS + O \rightarrow C + SO$	47.64	-2.37	-0.18	0.88	-0.05	1.54	-11.40	-3.55	-1.63
$CO + S \rightarrow C + SO$	137.49	-0.51	0.46	0.48	2.39	1.11	2.16	0.48	3.68
$CS + S \rightarrow C + S_2$	69.98	-2.90	-0.81	-2.40	-0.43	-3.35	-0.98	4.08	2.49
$CF + CI \rightarrow C + CIF$	71.91	0.32	1.36	1.23	-0.14	-0.26	2.77	1.07	-0.03
$CF + F \rightarrow C + F_2$	94.90	2.65	0.83	1.50	24.20	21.37	2.97	-1.16	27.02
$CO + O \rightarrow C + O_2$	139.28	0.07	-0.09	0.69	2.25	2.77	-0.91	-1.76	1.74
$CN + N \rightarrow C + N_2$	-46.49	-0.08	2.36	1.38	-17.97	-13.29	-0.72	5.60	-22.65
$NO + HCI \rightarrow N + HOCI$	92.63	-3.63	-0.32	-2.79	-2.92	-7.10	3.36	7.35	1.26
$NO + HF \rightarrow N + HOF$	132.92	-2.89	-0.18	-3.22	-0.46	-6.17	7.54	10.56	5.25
$NO + CI \rightarrow N + CIO$	88.29	-2.72	-2.70	-4.40	2.03	-3.60	3.90	4.02	7.66
$NO + NO \rightarrow N + NO_2$	75.06	0.87	-0.36	1.08	-5.02	-3.04	-3.14	-5.91	-7.00
$\rm N_2 + O_2 \rightarrow \rm N + \rm NO_2$	119.05	1.26	0.55	1.76	1.83	1.75	-1.61	-3.45	1.90
$\mathrm{NO} + \mathrm{N_2}  ightarrow \mathrm{N} + \mathrm{N_2O}$	107.04	0.40	0.00	1.33	-3.87	-2.13	-2.62	-5.31	-5.61
$NO + HNO \rightarrow N + ON$	43.24	-2.12	-1.05	-1.44	-11.40	-12.05	-2.16	0.12	-10.75
trans-HONO									

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

	2) $\Delta(MP2)$	5 6.06	3 2.29	3 -40.24	4 -2.02	3 21.50	7 4.70	4 9.37	3 -2.57	0 30	- 0.0- -	7 -7.73	1 5.70	0	-9.14	2 -20.87	4 -10.93	18.00	06.01- C	-1.55	3 -8.54		0 -17.12	-1.58	3 -8.69	5 -2.15	1 -15.58	5 -9.93	4 -24.38	5 -6.42	3 -22.74	7 -4.94	5 -24.63	9 -18.58	2 -5.88 3 -23.88
	$\Delta(OOMP2$	0.2(	3.75	6.8	$2.9_{4}$	3.8	-0.27	6.4	4.30	4 0, 4		17.87	7.47		0.00	8.92	-2.04	0001	T 0.30	6.6	15.10	0	12.0	4.58	2.75	3.95	14.4	-2.4	$6.1_{4}$	0.95	9.8	0.17	7.20	×.3	-2.62
	$\Delta(\kappa - OOMP2)$	-1.09	-1.00	-6.85	-3.88	1.98	-2.01	0.39	-3.47	0 O	00.0	4.32	4.39	110	-4.12	-6.14	-7.05	1 03	CU.1-	-0.47	3.65	0	0.34	-1.96	-3.83	1.77	1.69	-4.96	2.57	-7.55	1.12	-4.69	-6.52	8.19	0.05 -5.37
	$\Delta(MP3)$	0.79	-6.39	-33.72	-8.07	9.73	2.21	-0.37	-8.63	-0.16	01.0-	-17.18	-2.09	77 OF	-12.44	-22.49	-10.37	02 66	00.77-	-11.05	-17.31	00 10	-21.30	-7.81	-13.62	-5.78	-20.88	-10.10	-18.60	-9.12	-21.74	-7.54	-21.39	-16.86	-2.41 -16 44
ued)	$\Delta(\mathrm{MP2.5})$	3.43	-2.05	-36.98	-5.05	15.62	3.46	4.50	-5.60	22 V -	F	-12.46	1.81	00 11	60.11-	-21.68	-10.65	12.00	-20.14	-6.30	-12.92		12.61-	-4.69	-11.16	-3.97	-18.23	-10.02	-21.49	-7.77	-22.24	-6.24	-23.01	-17.72	-4.15 -20 16
1.1: (contin	$\Delta(MP3:\kappa-OOMP2)$	0.06	-4.34	-5.81	-4.24	-1.72	-0.84	-3.22	-3.54	-3.81	10.0-	-5.72	-2.30	09 E	00.6-	-6.18	-2.97	00 8	-0.40	-3.75	-5.06	л Сс	-0.39	-3.35	-4.78	-1.63	-5.06	-3.46	-1.27	-3.98	-4.33	-3.57	-3.72	-2.08	1.44 -1 67
Table C	$\Delta(MP2.8:\kappa-OOMP2)$	0.20	-2.46	-2.82	-2.47	-0.58	-0.92	-0.83	-1.48	-1.09	70.T_	-0.54	0.06	96.0	00.7-	-2.74	-2.44	00 0	67.7-	-1.24	-0.45	-	-1.30	-1.36	-2.69	-0.15	-0.75	-3.04	0.74	-2.52	-0.89	-2.56	-1.02	0.35	0.99 -0.30
	$\Delta(\text{CCSD})$	-0.71	-2.83	-4.49	-3.41	-1.26	-0.47	-2.99	-4.47	7 83	00-E	-5.46	-2.51	09 6	-0.00	-5.64	-2.72	717	-4.1/	-5.89	-4.76	7 0	-4.52	-4.58	-4.47	-1.93	-4.88	-1.75	-0.36	-4.09	-2.27	-3.00	-3.61	0.39	-0.42 -1 47
	CCSD(T)	88.87	69.87	-53.00	28.81	128.18	74.65	100.13	61.83	57 1 G	0110	87.05	20.33	00 00	07.67-	-56.03	-24.02	17 96	00.11	-29.85	43.05	00 - 1	67.11	35.80	-11.95	-32.20	34.52	-44.07	0.78	-14.79	-34.27	19.21	-91.91	9.40	-37.98 -80 45
	Reaction	$N_2 + HOO \rightarrow N + MONO$	$N_{2} + HCO \rightarrow N + HNCO$	$C\tilde{N} + HNO \rightarrow N + HNCO$	$NO + HCN \rightarrow N + HNCO$	$N_2 + OH \rightarrow N + HNO$	$N_2 + O \rightarrow N + NO$	$\rm NO + C_2H_2 \rightarrow N +$	oxirene NO + $C_2H_4 \rightarrow N +$	oxirane NO ± CH = ► N ±	methanol	$CN + H_2O \rightarrow N + $	$NO + {}^{3}CH_{2} \rightarrow N +$	trans-HCOH	$NO + vrans-nOOn \rightarrow N + formic acid$	$CN + HOOH \rightarrow N +$	formic acid NO + CH <sub>2</sub> C $\rightarrow$ N +	ketene	UN + 11200 -> IN + ketene	$NO + \text{ oxirene} \rightarrow N +$	glyoxal $CN + $ formic acid $\rightarrow N + $	glyoxal	$CN + methanol \rightarrow N +$ acetaldehvde	$NO + C_2 H_4 \rightarrow N +$	$NO + CF_2 \rightarrow N + F_2CO$	$NO + {}^{3}CH_{2} \rightarrow N + H_{2}CO$	$CN + H_2O \rightarrow N + H_2CO$	$NO + CH \rightarrow N + HCO$	$CN + S_2 \rightarrow N + CS_2$	$NO + CS \rightarrow N + OCS$	$CN + SO \rightarrow N + OCS$	$\rm NO + CO \rightarrow N + CO_2$	$CN + O_2 \rightarrow N + CO_2$	$CN + S \rightarrow N + CS$	$NO + Si \rightarrow N + SiO$ $CN + O \rightarrow N + CO$

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table (	C.1: (conti	inued)				
ction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta$ (MP2.8: $\kappa$ -OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta({ m MP3})$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta({ m MP2})$
$+ C \rightarrow N + CO$	-108.61	-0.92	-1.74	-2.21	-5.65	-5.36	-2.64	-0.97	-5.94
$- CH \rightarrow N + HCN$	-3.01	-1.17	-3.03	-3.57	-7.02	-8.42	-2.08	-1.65	-5.63
$- C \rightarrow N + CN$	46.49	0.08	-2.36	-1.38	17.97	13.29	0.72	-5.60	22.65
$+ H_2O \rightarrow N + HOOH$	113.82	-3.42	-0.18	-3.15	-1.87	-7.13	5.74	9.63	3.40
$+ OH \rightarrow N + HOO$	82.55	-2.66	-1.84	-3.22	0.79	-3.11	0.91	3.68	4.69
$+ H_2 \rightarrow N + H_2O$	26.85	-5.26	-1.72	-4.69	-5.83	-10.53	0.80	6.87	-1.14
$+ \operatorname{NH}_2 \rightarrow \operatorname{N} + $	112.47	-1.99	-0.74	-2.87	10.60	4.19	3.92	6.65	17.01
$^{s-\mathrm{N}_{2}\mathrm{H}_{2}}$ + $\mathrm{C}_{2}\mathrm{H}_{4}$ $\rightarrow$ N +	34.74	-4.76	-1.47	-5.23	-21.30	-22.35	-0.56	10.73	-20.24
yne $+ C_2 H_6 \rightarrow N +$	27.92	-4.52	-1.12	-5.78	-19.53	-22.55	2.70	14.54	-16.52
hene $+ \operatorname{CH}_{\operatorname{o}} F \to \operatorname{N} + \operatorname{C}_{\operatorname{o}} \operatorname{H}_{\operatorname{o}} F$	26.06	-4.39	-1.50	-6.12	-19.44	-22.56	2.35	14.30	-16.32
$+ CH, \rightarrow N + C, H$	33.15	-4.78	-1.02	-5.85	-19.28	-22.65	3.27	15.29	-15.92
$+ CH_{\circ} \rightarrow N + CH_{\circ}CH$	38.94	-4.61	-1.58	-6.33	-13.73	-17.62	3.29	14.86	-9.85
$+ CF_{o} \rightarrow N + FCCF$	49.32	-4.81	-2.30	-6.74	-26.49	-28.20	-0.27	11.47	-24.78
$+ {}^{3}C\tilde{H}_{2} \rightarrow N + C_{2}H_{2}$	-36.16	-1.99	-0.70	-4.00	-20.41	-20.41	3.40	10.13	-20.40
$+ \operatorname{Cl}_2 \xrightarrow{-}{} \mathrm{N} + \operatorname{CCl}_2 \xrightarrow{-}{}$	60.36	-3.25	-1.42	-5.22	-11.34	-14.00	2.53	11.82	-8.68
$+ F_2 \rightarrow N + CF_2$	-43.20	-7.77	-1.22	-5.45	-40.24	-38.57	-5.39	13.09	-41.91
$+ H_2 \rightarrow N + {}^3CH_2$	93.57	-8.21	-2.32	-8.12	-20.09	-25.62	0.72	17.33	-14.57
$- F \rightarrow N + CF$	45.16	-4.41	-0.45	-3.68	-17.31	-17.99	-2.09	11.59	-16.63
$+ F \rightarrow N + BF$	-78.36	-2.16	1.11	1.58	-2.91	-1.62	-4.79	-0.53	-4.19
$B \rightarrow N + ^{3}BN$	119.74	-3.08	-4.52	-7.56	0.95	-5.96	2.01	6.07	7.85
$+S_2 \rightarrow N+S_2O$	55.05	0.36	-0.85	0.50	10.49	7.66	-2.13	-8.11	13.32
$+ SO_2 \rightarrow N + SO_3$	71.31	-4.44	-3.13 0.70	-5.60	-6.93	-9.90	-4.00	5.08	-3.97
$+50 \rightarrow N + 50_2$	20.16	0.50-	06.0	1.20	-4.83 2 95	-3.07	-3.58	-3.60	-0.59
O + N + O	30.67	-1.40	-1.20	- 1.10	-3.20	-4.20	0.40 7.7.6	-0.40	07.7-
$HCI \rightarrow 0 + HOCI$	61.96	-2.78	1.52	-1.27	0.47	-4.52	6.91	10.08	5.46
$\text{HF} \rightarrow \text{O} + \text{HOF}$	102.25	-2.03	1.66	-1.70	2.93	-3.59	11.09	13.29	9.45
$CI \rightarrow O + CIO$	57.62	-1.87	-0.87	-2.87	5.42	-1.02	7.45	6.75	11.86
$\rm NO \rightarrow O + NO_2$	44.39	1.72	1.47	2.60	-1.63	-0.46	0.40	-3.18	-2.80
$HNO \rightarrow O +$	12.57	-1.27	0.78	0.08	-8.01	-9.46	1.39	2.85	-6.55
$-HOO \rightarrow O +$	14.21	-0.25	1.12	0.89	-0.03	-1.41	0.93	0.53	1.36
-ONOH-								0	
$+ HCO \rightarrow O + HNCO$	-4.79	-2.36	-1.54	-3.50	-5.51	-8.60	1.01	4.00	-2.42
- HNO $\rightarrow$ O + HNCO	27.45	-3.02	-2.52	-4.14	-16.82	-17.28	-1.49	2.46	-16.36
$HCN \rightarrow O + HNCO$	-1.86	-2.56	-0.64	-2.71	-1.66	-5.49	-0.33	5.67	2.17
OH + O + HO +	53.52	-0.79	0.33	-0.88	12.16	7.53	4.00	4.10	16.80
$N \rightarrow O + NO$	-30.67	0.85	1.83	1.52	3.39	2.59	3.55	2.73	4.20
$C_2H_2 \rightarrow 0 + \text{ oxirene}$	69.46	-2.14	1.00	-1.69	7.89	2.22	3.94	9.17	13.56
$C_2H_4 \rightarrow O + \text{ oxirane}$	31.16	-3.62	0.35	-2.02	-2.21	-6.04	0.08	7.09	1.62
$-\mathrm{CH}_4 \rightarrow \mathrm{O} +$	26.49	-3.98	0.81	-2.29	-1.38	-6.57	4.51	9.95	3.81
ithanol									

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table (	C.1: (conti	nued)				
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\begin{array}{c} \Delta(\mathrm{MP2.8:}\kappa-\\ \mathrm{OOMP2}) \end{array}$	$\Delta$ (MP3: $\kappa$ -OOMP2)	$\Delta(\mathrm{MP2.5})$	$\Delta({ m MP3})$	$\Delta(\kappa - 00MP2)$	$\Delta(OOMP2)$	$\Delta({ m MP2})$
$\frac{\text{CO} + \text{H}_2\text{O} \rightarrow \text{O} +}{tmns-\text{HCOH}}$	167.50	-3.99	-0.24	-4.05	7.71	-0.74	9.69	13.50	16.15
$\begin{array}{c} 0_2 + {}^{3}\mathrm{CH}_2 \rightarrow \mathrm{O} + \\ t_{muss} + \mathrm{HCOH} \end{array}$	-10.34	-1.66	1.89	-0.77	5.20	0.50	7.94	10.15	9.89
$O_2 + trans-HCOH \rightarrow O_2 + trans-HCOH \rightarrow O_2$	-59.93	-2.75	-0.54	-2.08	-7.70	-9.85	-1.17	3.41	-5.54
O + 1011110 actu $CO + HOOH \rightarrow O +$	24.42	-4.17	-2.44	-4.50	-1.52	-6.05	-0.77	4.55	3.02
$O_2 + CH_2C \rightarrow O + ketene$	-54.69	-1.86	-0.60	-1.45	-7.26	-7.78	-3.50	0.69	-6.74
$CO + H_2CO \rightarrow O + ketene$	97.81	-2.70	-2.00	-4.53	-0.58	-6.14	4.33	6.61	4.99
$O_2 + oxirene \rightarrow O + O_2$	-60.52	-5.04	0.60	-2.22	-2.91	-8.47	3.08	9.42	2.65
$\begin{array}{c} \text{glyoxal}\\ \text{CO} + \text{ formic acid} \rightarrow \text{O} + \end{array}$	123.50	-3.29	-0.15	-3.39	7.24	-0.87	9.02	10.79	15.34
glyoxal $CO + methanol \rightarrow O + O + O$	92.24	-3.06	-1.06	-3.72	0.96	-4.86	5.71	8.29	6.77
acetaldenyde $O_2 + C_2 H_4 \rightarrow O +$	5.13	-3.73	0.48	-1.83	-1.30	-5.22	1.58	7.31	2.62
acetaldehyde O + CF -> O + F CO	69 67-	-3 69	-0 86 -0	-3 96 -	7 76	-11 03	86 U-	74 A	07 70
$0_{0}^{2} + {}^{3}CH_{0} \rightarrow 0 + H_{0}^{2}CO$	-42.02 -62.87	-0.02 -1.08	1.68	-0.11	-0.57	-3.20	5.31	0.41 0.68	2.05
$CO + H_{2}O \rightarrow O + H_{3}CO$	114.97	-3.41	-0.45	-3.38	1.94	-4.44	7.06	10.04	8.31
$O_2 + CH \rightarrow O + HCO$	-74.74	-0.90	-1.21	-1.94	-6.62	-7.51	-1.41	0.28	-5.74
$SO + CS \rightarrow O + CS_2$	13.71	-1.28	-0.23	-2.47	-4.16	-6.62	4.80	5.37	-1.69
$CO + S_2 \rightarrow O + CS_2$	81.23	1.11	1.04	0.41	-1.33	-2.16	7.94	1.77	-0.50
$O_2 + CS \rightarrow O + OCS$	-45.46	-3.23	-0.69	-2.46	-4.38	-6.53	-4.00	3.68	-2.22
$sO + CO \rightarrow O + CO$	40.18 11 46	-0.80	-0.59	-2.65 2.05	-2.08 9.85	-5.3U A QK	6.49 1 15	5.46 2.00	1.14 0.74
$C_2 + C_2 \rightarrow C_2 + C_2$ $CO + S \rightarrow O + CS$	-11.40 89.85	$^{-2.14}$	-0.13	-0.41	-2.00	-4.30	13.56	2.30 4.03	5.31
$SO + C \rightarrow O + CS$	-47.64	2.37	0.18	-0.88	0.05	-1.54	11.40	3.55	1.63
$O_2 + Si \rightarrow O + SiO$	-68.65	0.43	2.82	2.96	-0.75	0.18	3.60	0.11	-1.69
$O_2 + C \rightarrow O + CO$	-139.28	-0.07	0.09	-0.69	-2.25	-2.77	0.91	1.76	-1.74
$NO + CH \rightarrow O + HCN$	-77.67	-0.71	-2.12	-2.73	-10.48	-10.62	-0.07	-1.38	-10.33
$CO + N \rightarrow O + CN$	80.45	1.47	0.30	1.67	20.16	16.44	5.37	-4.36	23.88
	01.02-	0.00	CF.1-	-0.04	14.02	0 10 0 10	2.13	-0.00 14 65	11.95 7 5 6
$CIO + AI \rightarrow O + AICI$	40.12	-2.97	-0.10	-174	-7.30	-0.43	9.40 2.95	14.00 7 15	-5.38
$O_0 + H_0 O \rightarrow O + HOOH$	83.14	-2.56	1.65	-1.63	1.53	-4.54	9.28	12.36	7.59
$SO + HS \rightarrow O + SSH$	47.64	-2.39	-0.37	-3.71	0.72	-5.29	9.39	10.37	6.72
$O_2 + OH \rightarrow O + HOO$	51.88	-1.81	-0.01	-1.69	4.18	-0.52	4.46	6.41	8.89
$ ilde{O_2} +  ext{H}_2  ightarrow  ext{O} +  ext{H}_2  ext{O}$	-3.82	-4.41	0.11	-3.17	-2.44	-7.94	4.34	9.60	3.06
$NO + NH_2 \rightarrow O + $	37.81	-1.52	0.18	-2.04	7.14	1.98	5.93	6.92	12.30
CO + C + C + O + O + O + O + O + O + O +	115 19	-3.30	-1 18	-356	-113	-5 91	4 80	6.36	3.64
propyne	27.07 F	22.2	)	22.22	DT-T_	13.5	20.4	22.2	F0.0

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			TOT		mann				
eaction	CCSD(T)	$\Delta(CCSD)$	$\Delta$ (MP2.8: $\kappa$ - OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$O + C_2 H_6 \rightarrow O +$	108.37	-3.05	-0.83	-4.10	0.63	-6.11	8.06	10.18	7.37
$\begin{array}{l} \text{Opene} \\ \text{O} + \text{CH}_3 \text{F} \rightarrow \text{O} + \text{C}_3 \text{H}_3 \text{F} \end{array}$	106.51	-2.93	-1.21	-4.45	0.72	-6.12	7.72	9.94	7.56
$\mathrm{O}+\mathrm{CH}_4^{} ightarrow\mathrm{O}+\mathrm{C}_2\mathrm{H}_4^{}$	113.60	-3.31	-0.72	-4.17	0.88	-6.21	8.64	10.93	7.96
$O + CH_3 \rightarrow O + CH_3\dot{C}H$	119.39	-3.15	-1.28	-4.66	6.43	-1.18	8.66	10.50	14.03
$O + CF_2 \rightarrow O + FC\overline{CF}$	129.77	-3.34	-2.00	-5.07	-6.33	-11.76	5.10	7.10	-0.89
$0 + {}^{3}C\tilde{H}_{2} \rightarrow 0 + C_{2}H_{2}$	44.29	-0.52	-0.41	-2.32	-0.24	-3.97	8.77	5.76	3.48
$0 + Cl_2 \xrightarrow{\sim} 0 + CCl_2 \xrightarrow{\sim}$	140.81	-1.78	-1.13	-3.55	8.82	2.44	7.90	7.45	15.20
$O + F_2 \rightarrow O + CF_2$	37.25	-6.31	-0.93	-3.78	-20.08	-22.13	-0.02	8.72	-18.02
$\mathrm{O} + \mathrm{H}_2^{}  ightarrow \mathrm{O} + {}^3\mathrm{C}\mathrm{\tilde{H}}_2$	174.02	-6.74	-2.03	-6.44	0.07	-9.18	6.09	12.96	9.31
$O + F \rightarrow O + SiF$	47.85	-4.84	-2.00	-5.02	0.62	-5.29	-0.34	9.55	6.53
$O + CN \rightarrow O + CICN$	-44.33	-3.00	1.59	-2.10	-26.37	-23.99	0.79	11.46	-28.74
$O + F \rightarrow O + CF$	125.61	-2.94	-0.16	-2.01	2.85	-1.55	3.28	7.22	7.26
$O + B \rightarrow O + {}^{3}BN$	45.09	-2.61	-3.60	-6.73	-2.51	-8.16	4.02	6.34	3.14
$O \rightarrow O + S_2O$	48.51	1.25	-0.19	-1.04	13.36	7.03	8.76	0.00	19.69
$c + S_2 \rightarrow O + S_2 O$	24.38	1.21	0.98	2.03	13.88	10.25	1.41	-5.38	17.51
$(+SO_2 \rightarrow O + SO_3)$	40.64	-3.58	-1.30	-4.08	-3.54	-7.31	-0.46	7.81	0.23
$+$ SO $\rightarrow$ O $+$ SO <sub>2</sub>	-4.51	0.35	2.33	2.72	-1.44	-0.48	-0.03	-0.87	-2.40
$+ S \rightarrow O + SO$	-1.79	-0.58	0.55	-0.21	0.14	-1.66	3.06	2.25	1.94
$0 + S \rightarrow O + S_{2}$	22.33	-0.54	-0.62	-3.28	-0.38	-4.88	10.41	7.63	4.11
$0 + CI \rightarrow 0 + CI_2$	5.25	-1.52	1.71	-0.48	-6.07	-7.25	4.41	6.83	-4.88
$O + F \rightarrow O + CIF$	0.63	-0.82	2.16	1.40	-4.97	-3.57	-0.50	3.31	-6.36
$O + N \rightarrow O + N_2$	-74.65	0.47	0.92	0.84	-3.46	-2.21	2.01	0.27	-4.70
$+ CH_3 \rightarrow F + CH_3F$	-77.61	-4.11	-0.04	-1.37	-25.77	-23.71	-2.35	3.09	-27.83
$\mathbf{F} + \mathbf{O} \rightarrow \mathbf{F} + \mathbf{CIO}$	-0.63	0.82	-2.16	-1.40	4.97	3.57	0.50	-3.31	6.36
$h + HNO \rightarrow F + HNCO$	-98.17	-0.08	-2.37	-2.13	-19.67	-15.73	-4.76	-4.76	-23.61
$^{+}$ H <sub>2</sub> O $\rightarrow$ F +	41.88	-1.05	-0.09	-2.04	4.85	0.81	6.41	6.28	8.89
ns-hCOH		00 <del>-</del>	0000	0					
+ n∪On → r +	6T.TU1-	-1.23	97.7-	- 2.49	-4.37	-4.00	cu. <del>1</del> -	10.2-	-4.24
шис асци л н ПО , Б , 1t	10 70		10.1	6 2 6	67 6	A 60	1.05	0.61	70.0
TII200 -7 F + Kevelle	TO:17-	1.0 95	#0'T-	40.4- 00 1	06 4	-4-U3 0 69	ГО.Т Т	10.0-	19:7-
oxal	71.7-	-0.0-	00.00	OC'T-	4.03	0.00	1.0	10.0	0.09
$^{\mathrm{r}}$ + methanol $\rightarrow$ F +	-33.37	-0.12	-0.90	-1.71	-1.90	-3.31	2.43	1.07	-0.49
etaldehyde									
$+ \text{HCO} \rightarrow \text{F} + \text{HCOF}$	-87.01	-4.61	-0.88	-2.28	-26.81	-25.25	-4.14	2.69	-28.38
$^{\mathrm{r}} + \mathrm{H}_{2}\mathrm{O} \rightarrow \mathrm{F} + \mathrm{H}_{2}\mathrm{CO}$	-10.65	-0.47	-0.30	-1.37	-0.92	-2.88	3.78	2.82	1.05
$r + S_2 \rightarrow F + CS_2^-$	-44.38	4.05	1.20	2.42	-4.18	-0.61	4.66	-5.45	-7.76
$\mathbf{F} + \mathbf{SO} \rightarrow \mathbf{F} + \mathbf{OCS}$	-79.44	2.14	-0.44	-0.64	-4.93	-3.75	3.21	-1.76	-6.11
$^{\mathrm{r}}\mathrm{+O}_{2}\mathrm{ ightarrow}\mathrm{F}\mathrm{+CO}_{2}$	-137.08	0.80	-0.57	-0.04	-5.70	-3.40	-4.43	-4.32	-8.00
$\mathbf{T} + \mathbf{S} \rightarrow \mathbf{F} + \mathbf{CS}$	-35.76	4.80	0.80	1.60	-0.41	1.13	10.28	-3.19	-1.95
$\mathbf{f} + \mathbf{O} \rightarrow \mathbf{F} + \mathrm{SiO}$	-47.85	4.84	2.00	5.02	-0.62	5.29	0.34	-9.55	-6.53
$\mathbf{F} + \mathbf{O} \rightarrow \mathbf{F} + \mathbf{CO}$	-125.61	2.94	0.16	2.01	-2.85	1.55	-3.28	-7.22	-7.26
$r + N \rightarrow F + CN$	-45.16	4.41	0.45	3.68	17.31	17.99	2.09	-11.59	16.63
$F + CI \rightarrow F + AlCI$	40.41	-0.27	-0.53	-2.20	-0.25	-3.60	6.93	4.72	3.09

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table	C.1: (conti	nued)				
ttion	CCSD(T)	$\Delta(CCSD)$	$\frac{\Delta(\text{MP2.8:}\kappa-00\text{MP2})}{00\text{MP2}}$	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(OOMP2)$	$\Delta(MP2)$
$+ AI \rightarrow F + AICI$	-59.99	-2.16	-1.32	-3.14	-2.34	-5.66	3.44	3.85	0.99
$- AI \rightarrow F + AIF$	-123.38	-4.21	-0.27	-1.22	-26.43	-23.70	-3.68	1.36	-29.15
$+ C_2 H_4 \rightarrow F +$	-10.42	-0.36	-1.02	-1.55	-3.99	-4.36	1.52	-0.86	-3.62
$+ C_2 H_6 \rightarrow F +$	-17.24	-0.11	-0.67	-2.09	-2.22	-4.56	4.79	2.96	0.11
ene - CH <sub>3</sub> CH →	-90.49	-3.89	0.03	-1.16	-31.48	-28.66	-3.29	2.53	-34.30
$c_2 H_3 F$									
$- \operatorname{CH}_3 \operatorname{F} \to \operatorname{F} + \operatorname{C}_2 \operatorname{H}_3 \operatorname{F}$	-19.10	0.01	-1.05	-2.44	-2.13	-4.57	4.44	2.72	0.31
- $\mathrm{CH}_4 \to \mathrm{F} + \mathrm{C}_2 \mathrm{H}_4$	-12.01	-0.37	-0.57	-2.16	-1.98	-4.66	5.36	3.70	0.71
$- CH_3 \rightarrow F + CH_2 CH$	-6.22	-0.21	-1.13	-2.65	3.58	0.37	5.38	3.28	6.78
$- \operatorname{CF}_2 \to \operatorname{F} + \operatorname{FCCF}$	4.16	-0.40	-1.84	-3.06	-9.18	-10.21	1.82	-0.12	-8.15
$CCH \rightarrow F + HCCF$	-96.00	-3.73	-0.46	-1.65	-38.08	-34.05	-4.40	1.88	-42.11
$^{-3}CH_{2} \rightarrow F + C_{2}H_{2}$	-81.32	2.42	-0.25	-0.31	-3.10	-2.42	5.49	-1.46	-3.77
$- \operatorname{Cl}_{2} \xrightarrow{\circ} F + \operatorname{CCl}_{2} \xrightarrow{\circ}$	15.20	1.16	-0.97	-1.54	5.97	3.99	4.62	0.23	7.95
$CF \rightarrow F + CF_{2}$	-88.36	-3.37	-0.77	-1.77	-22.93	-20.58	-3.30	1.50	-25.28
$- H_{2} \rightarrow F + {}^{3}CH_{2}$	48.40	-3.81	-1.87	-4.43	-2.79	-7.63	2.81	5.74	2.06
$Si \rightarrow F + SiF$	-102.03	-4.05	0.05	-0.87	-24.93	-22.16	-3.89	1.84	-27.70
$+ CN \rightarrow F + CICN$	-44.96	-2.18	-0.57	-3.50	-21.40	-20.43	1.29	8.16	-22.37
$C \rightarrow F + CF$	-94.90	-2.65	-0.83	-1.50	-24.20	-21.37	-2.97	1.16	-27.02
$B \rightarrow F + BF$	-145.17	-3.56	-1.42	-2.43	-26.82	-24.25	-4.38	0.72	-29.39
$- N \rightarrow F + {}^{3}BN$	78.36	2.16	-1.11	-1.58	2.91	1.62	4.79	0.53	4.19
$+ \mathrm{Cl} \rightarrow \mathrm{F} + \mathrm{Cl}_{2}$	4.62	-0.70	-0.45	-1.88	-1.10	-3.69	4.90	3.52	1.48
$CI \rightarrow F + CIF$	-22.99	-2.33	0.53	-0.28	-24.34	-21.64	-0.20	2.24	-27.05
$HCI \rightarrow S + HOCI$	63.75	-2.20	0.97	-1.06	0.33	-2.86	3.84	7.83	3.52
$HF \rightarrow S + HOF$	104.04	-1.46	1.11	-1.49	2.79	-1.93	8.02	11.04	7.51
$\cdot \mathrm{Cl} \rightarrow \mathrm{S} + \mathrm{ClO}$	59.41	-1.30	-1.42	-2.66	5.28	0.64	4.39	4.51	9.92
$\rm NO  ightarrow S + NO_2$	46.18	2.30	0.92	2.81	-1.77	1.20	-2.66	-5.43	-4.74
$N_2  ightarrow { m S} { m S} + { m N}_2 { m O}$	78.16	1.83	1.28	3.07	-0.62	2.11	-2.14	-4.83	-3.35
$HNO \rightarrow S +$	14.36	-0.69	0.23	0.29	-8.15	-7.80	-1.67	0.60	-8.49
-HONO									
HNO $\rightarrow$ S + HNCO	-62.41	-4.88	-3.17	-3.73	-19.26	-16.86	-15.04	-1.56	-21.66
$+$ HCN $\rightarrow$ S $+$ HNCO	-0.07	-1.98	-1.19	-2.50	-1.80	-3.83	-3.40	3.42	0.23
$-N \rightarrow S + NO$	-28.88	1.43	1.28	1.73	3.25	4.25	0.48	0.48	2.26
$-C_2H_2 \rightarrow S + \text{ oxirene}$	71.25	-1.56	0.45	-1.48	7.75	3.88	0.87	6.92	11.62
$C_{0}H_{4} \rightarrow S + oxirane$	32.95	-3.04	-0.20	-1.80	-2.35	-4.38	-2.99	4.84	-0.32
$- \operatorname{CH}_4 \rightarrow \operatorname{S} +$	28.28	-3.41	0.26	-2.08	-1.52	-4.91	1.44	7.70	1.87
anol									
$H_2O \rightarrow S + UCOU$	77.64	-5.85	-0.89	-3.64	5.27	-0.31	-3.87	9.48	10.85
s-n∪∪u - <sup>3</sup> CH <sub>o</sub> → S +	-8.55	-1.08	1.34	-0.56	5.06	2.16	4.88	7.90	7.95
s-HCOH									
- trans-HCOH $\rightarrow$	-58.14	-2.17	-1.09	-1.87	-7.84	-8.20	-4.24	1.16	-7.48
formic acid									
$HOOH \rightarrow S + S$	-65.43	-6.03	-3.08	-4.10	-3.96	-5.63	-14.33	0.53	-2.29

P2) $\Delta(MP2)$	.55 -8.68	.58 -0.32	.17 0.71	.76 10.04		.26 1.46		.07 0.68	679 GU		.44 $0.11$	.01 $3.00$	.96 -7.68	.26 -5.81	.43 -4.16	.17 -2.97	.65 -2.68	.13 -6.05	.08 -2.49	.13 -3.63	.03 -5.31	.48 -3.68	.39 18.58	.01 3.47	.11 5.65	.36 1.12	.33 -1.67		.15 $2.06$	.91 2.26	.90 2.66	.47 8.73	.07 -6.20	.73 -1.82	.42 9.90	.69 -23.33	.93 4.01	.19  1.95	.63 15.57	.56 -1.71	.12 -4.34	-6.51	.63 -4.11
Δ(OOM	7 -1	2	1 7	4 6		5 4	L	Ω Ω	с и	0 ×	5	0	8 -1	2 -2	7 1	3 -2	1 0	1 -1	8 -4	4 -2	6 -4	9 -0	8- 6	1 7	2 10	8 7	5 2		9	4 5	2 6	0 6	6 3	9 1	6 3	8	7 8	8	5 -7	2	-3	5 8 1 8	1 -7
$\Delta(\kappa - 00MP2)$	-6.5	-9.2	0.0	-4.5		-7.8	Ţ	-1.4	0 0	-0.0 0.0	2.2	-6.5	-4.4	-5.6	-7.0	-3.9	-4.2	-14.7	0.9	0.5	-13.5	-2.1	-8.1	-1.0	6.2	1.2	-8.7		-5.4	-5.8	-4.9	-4.9	-8.4	-4.7	-5.6	-13.5	-7.4	-10.2	-1.6	-3.5	-3.1	-10.4	-10.4
$\Delta(MP3)$	-6.13	-5.72	-6.81	-0.44		-4.43	с С	-3.50	06 0	-4.00	-1.54	-4.01	-5.85	-1.74	-4.88	-0.42	-3.30	-4.53	3.35	1.84	0.42	-1.11	16.86	-3.60	-2.88	-6.28	-5.48		-5.68	-5.70	-5.78	-0.75	-11.33	-3.55	2.86	-21.71	-8.75	-1.13	11.91	-5.65	1.17	4.40	4.88
$\Delta({ m MP2.5})$	-7.40	-3.02	-3.05	4.80		-1.49		-1.44	00 2	-1.90	-0.71	-0.51	-6.76	-3.77	-4.52	-1.69	-2.99	-5.29	0.43	-0.89	-2.44	-2.39	17.72	-0.07	1.39	-2.58	-3.57		-1.81	-1.72	-1.56	3.99	-8.77	-2.69	6.38	-22.52	-2.37	0.41	13.74	-3.68	-1.58	-1.06	0.38
$\begin{array}{c} \Delta(\mathrm{MP3:}\kappa-\\ \mathrm{OOMP2}) \end{array}$	-1.23	-4.13	-2.01	-2.98		-3.31	00 F	-1.62	9 05	01.0-	0.10	-2.98	-1.73	0.81	-2.25	0.63	-1.84	-1.64	2.40	3.17	0.41	-0.48	2.08	-1.24	-1.42	-2.96	-3.15		-3.70	-4.04	-3.77	-4.25	-4.67	-1.92	-3.14	-3.37	-6.04	-1.60	2.24	-3.86	2.93	6.00	3.28
$\begin{array}{c} \Delta(\mathrm{MP2.8:}\kappa\text{-}\\ \mathrm{OOMP2}) \end{array}$	-1.15	-2.64	0.05	-0.80		-1.71		10.0-	LV L	-1.41	1.13	-1.10	-1.76	0.39	-1.24	0.03	-1.28	-1.37	0.81	2.27	-0.65	-0.46	-0.35	0.52	1.10	-0.44	-1.82		-1.47	-1.85	-1.37	-1.93	-2.65	-1.05	-1.77	-1.57	-2.67	-0.80	0.43	-1.85	1.78	2.96	0.62
$\Delta(\text{CCSD})$	-1.29	-4.56	-4.46	-5.15		-4.91	с 1	-3.15	V 0 6	-0.04	-0.50	-5.27	-0.33	-0.75	-2.66	-0.26	-1.57	-4.00	2.90	1.01	-1.86	0.51	-0.39	-4.28	-1.99	-3.83	-5.15		-4.91	-4.78	-5.17	-5.01	-5.20	-2.38	-3.64	-8.16	-8.60	-4.80	1.79	-3.01	0.93	0.89	0.54
CCSD(T)	-52.90	7.96	-58.73	33.65		2.39	00 0	0.92	00 01	-40.00	-61.08	25.11	-72.96	-8.62	-43.68	23.84	-9.67	-101.32	-69.98	-66.86	-89.85	-137.49	-9.40	24.38	84.93	-2.03	25.34		18.52	16.66	23.75	29.54	39.92	-45.56	50.96	-52.60	84.16	35.76	26.17	42.42	-2.72	-26.85	-22.33
Reaction	$SO + CH_2C \rightarrow S + ketene$	$CS + H_2CO \rightarrow S + ketene$	$SO + oxirene \rightarrow S + \sigma$	$CS + formic acid \rightarrow S +$	glyoxal	$CS + methanol \rightarrow S +$	acetaldehyde	$SO + C_2H_4 \rightarrow S + c_2 c_1 c_2 c_2 c_2 c_2 c_2 c_2 c_2 c_2 c_2 c_2$	acetaldellyde SO i CE , S i E CO	$3O + Cr_2 \rightarrow 3 + r_2 OO$	$SO + {}^{3}CH_{2} \rightarrow S + H_{2}CO$	$CS + H_2O \rightarrow S + H_2CO$	$\rm SO + CH \rightarrow S + HCO$	$CS + S_2 \rightarrow S + CS_2$	$SO + CS \rightarrow S + OCS$	$S_2 + CO \rightarrow S + OCS$	$SO + CO \rightarrow S + CO$ ,	$CS + O_3 \rightarrow S + CO_3^2$	$S_{2} + C \xrightarrow{\sim} S + CS$	$SO + Si \rightarrow S + SiO$	$CS + O \rightarrow S + CO$	$SO + C \rightarrow S + CO$	$CS + N \rightarrow S + CN$	$S_2 + H_2 \rightarrow S + H_2S$	$SO + H_{o}O \rightarrow S + HOOH$	$\rm SO + H_2^2 \rightarrow S + H_2O$	$CS + C_2 H_4 \rightarrow S + \tilde{C}$	propyne	$CS + C_2H_6 \rightarrow S +$ propene	$CS + CH_3F \rightarrow S + C_2H_3F$	$CS + CH_4 \rightarrow S + C_2 H_4$	$CS + CH_3 \rightarrow S + CH_2CH$	$CS + CF_3 \rightarrow S + FC\overline{CF}$	$CS + {}^{3}CH_{2} \rightarrow S + C_{2}H_{2}$	$CS + CI_2 \rightarrow S + CCI_2$	$CS + F_2 \rightarrow S + CF_2$	$CS + H_2 \rightarrow S + {}^3CH_2$	$CS + F \rightarrow S + CF$	$S_2 + SO \rightarrow S + S_2O$	$SO + SO_2 \rightarrow S + SO_3$	$2{ m SO} ightarrow{ m S+SO}_2$	$\mathbf{S}_2 + \mathbf{O}_2 \rightarrow \mathbf{S} + \mathbf{SO}_2$	$S_2 + O \rightarrow S + SO$

$\Delta(MP2)$		-1.94	-0.78	-6.40	-2.40	4.88	-14.66	-13.27	-18.41		-9.68	-7.66	1.71	10.93	07.01-	-8.05		-1.96	-17.40		-18.59		-9.21		-9.24	-16.35	-1.33	-9.81	-17.59	-14.08	-12.60	-13.54	-13.59	-0.50	-3.09	-2.10	-4.26	-2.97	-8.80	-7.25		-15.06	1 77
$\Delta(OOMP2)$	1	-2.25	0.85	3.33	0.54	-6.83	-9.94	-9.33	-3.90		-1.08	-4.02	2.41	0.37	F0.0	3.20		3.39	-3.34		-6.06		2.66		0.56	-1.29	0.46	-0.07	-6.47	-3.07	-3.85	-6.64	-4.99	0.32	-4.72	-0.87	5.61	-0.34	2.85	0.30		-0.35	0.12
$\Delta(\kappa - \lambda)$	OOMP2)	-3.06	-2.15	-0.54	3.64	-4.41	-7.04	-6.52	-6.06		-7.78	-3.90	-3.51	7.97	IO: 1-	-2.94		0.49	-8.62		-10.95		-4.37		-5.87	-7.73	-3.94	-2.14	-8.86	-11.45	-8.60	-3.85	-6.54	-1.46	-6.93	-3.48	1.83	-2.99	-3.11	-3.09		-4.20	01.0
$\Delta({ m MP3})$	4 4	1.66	-2.08	-3.50	-2.57	7.25	0.56	1.47	-8.45		-4.47	3.60	3.24	с0 л	70.0-	-5.56		1.52	-8.84		-6.77		-7.45		-4.20	-10.02	-3.61	-2.18	-6.49	-5.52	-3.94	1.20	-1.75	-1.98	3.60	-2.07	-3.52	0.49	-6.92	-7.03		-12.41	1 06
iuea) ∆(MP2.5)		-0.14	-1.43	-4.95	-2.49	6.07	-7.05	-5.90	-13.43		-7.08	-2.03	2.47	763	· · · ·	-6.80		-0.22	-13.12		-12.68		-8.33		-6.72	-13.18	-2.47	-5.99	-12.04	-9.80	-8.27	-6.17	-7.67	-1.24	0.25	-2.08	-3.89	-1.24	-7.86	-7.14		-13.73	1 1 1
$\Delta(MP3:\kappa-\Delta(MP3:\kappa-\lambda))$	OOMP2)	0.21	-1.10	1.60	7.17 7	0.48	5.48	5.73	2.96		0.16	4.40	1.18	98 0	00.0	0.59		2.10	0.79		1.43		0.65		1.04	-0.39	-2.00	2.77	0.93	0.42	0.82	5.83	2.19	-1.26	2.20	-0.94	1.25	1.18	-0.29	-0.88	1	-1.37	1 40
Lable $\Delta(MP2.8:\kappa)$	OOMP2)	-0.55	-0.57	2.39	2.53	-1.71	2.34	2.71	1.65		0.23	2.70	1.87	1 22	77.T	1.68		2.76	0.33		0.27		1.47		1.35	0.01	-1.41	2.55	-0.34	0.18	0.14	3.69	0.96	-0.88	0.53	-0.80	2.52	0.86	0.98	-0.49	0	-0.99	1 20
$\Delta(\text{CCSD})$	0 1 0	0.58	-1.78	-0.91	-0.16	1.52	3.59	3.13	0.60		-0.69	2.72	-0.27	1 77	F. I . T.	-2.11		0.22	-0.88		0.01		-3.16		-1.86	-1.75	-2.27	0.79	0.97	-1.36	-0.27	2.31	1.81	-1.46	0.27	-1.88	-0.69	0.06	-2.54	-1.56		-1.39	1 0.0
CCSD(T)		1.79	-54.62	4.34	44.63	-5.25	-13.22	18.75	-45.05		-59.48	-88.29	11.84	96 <i>1</i> 6	01.07-	-31.13		-67.96	-117.55		-112.31		-118.14		-52.48	-100.24	-64.03	-120.49	-132.36	-103.08	-69.08	-126.27	-196.90	-64.61	-40.41	-100.39	25.53	-5.74	-61.44	-67.50		-73.01	010
Reaction	5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	$SO + O \rightarrow S + O_2$	$\operatorname{CIF} + \operatorname{CH}_3 \to \operatorname{CI} + \operatorname{CH}_3 \operatorname{F}$	$CIO + HCI \rightarrow CI + HOCI$	$CIO + HF \rightarrow CI + HOF$	$CI_2 + O \rightarrow CI + CIO$	$CIO + NO \rightarrow CI + NO_2$	$CIO + N_2 \rightarrow CI + N_2O$	$CIO + HNO \rightarrow CI + $	trans-HONO	$CIO + HCN \rightarrow CI + HNCO$	$CIO + N \rightarrow CI + NO$	$ClO + C_2H_2 \rightarrow Cl + C_2H_2$	OXITEDE	$0.0 \pm 0.014 = 0.01 \pm 0.015$	$ClO + CH_4 \rightarrow Cl +$	methanol	$CIO + {}^{3}CH_{2} \rightarrow CI +$	$CIO + trans-HCOH \rightarrow$	Cl + formic acid	$CIO + CH_2C \rightarrow CI +$	ketene	$ClO + oxirene \rightarrow Cl +$	glyoxal	$CIO + C_2H_4 \rightarrow CI +$ acetaldehyde	$CIO + CF$ , $\rightarrow CI + F$ , $CO$	$CIF + HCO \rightarrow CI + HCOF$	$CIO + {}^{3}CH_{2} \rightarrow CI + H_{2}CO$	$CIO + CH \rightarrow CI + HCO$	$CIO + CS \rightarrow CI + OCS$	$CIO + CO \rightarrow CI + CO$ ,	$CIO + Si \rightarrow CI + SiO$	$CIO + C \rightarrow CI + CO$	$Cl_2 + Al \rightarrow Cl + AlCl$	$AI\tilde{C}I + F \rightarrow CI + AIF$	$CIF + AI \rightarrow CI + AIF$	$CIO + H_2O \rightarrow CI + HOOH$	$OH + CIO \rightarrow CI + HOO$	$CIO + H_2 \rightarrow CI + H_2O$	$CIF + CH_2CH \rightarrow $	$CI + C_2 H_3 F$	$CIF + CCH \rightarrow CI + HUCF$	בכינכי בכי בכ

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta(MP2.8:\kappa)$	$\Delta(MP3;\kappa)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa$ -	A(OOMP2)	$\Delta(MP2)$
			OOMP2)	OOMP2)			OOMP2)		
$F + Si \rightarrow Cl + SiF$	-79.04	-1.72	-0.47	-0.59	-0.59	-0.53	-3.69	-0.39	-0.65
$_2 + CN \rightarrow CI + CICN$	-49.58	-1.48	-0.12	-1.62	-20.30	-16.74	-3.61	4.63	-23.86
$F + C \rightarrow CI + CF$	-71.91	-0.32	-1.36	-1.23	0.14	0.26	-2.77	-1.07	0.03
$\mathrm{F} + \mathrm{B}  ightarrow \mathrm{Cl} + \mathrm{BF}$	-122.19	-1.23	-1.95	-2.15	-2.48	-2.62	-4.18	-1.52	-2.34
$O + S_2 \rightarrow CI + S_2O$	-33.23	3.08	1.85	4.90	8.46	11.27	-6.04	-12.14	5.66
$O + SO_2 \rightarrow CI + SO_3$	-16.98	-1.71	-0.42	-1.20	-8.96	-6.29	-7.90	1.06	-11.63
$O + SO \rightarrow CI + SO$	-62.13	2.22	3.20	5.60	-6.86	0.53	-7.48	-7.63	-14.25
$O + S \rightarrow CI + SO$	-59.41	1.30	1.42	2.66	-5.28	-0.64	-4.39	-4.51	-9.92
$c + F \rightarrow Cl + ClF$	-4.62	0.70	0.45	1.88	1.10	3.69	-4.90	-3.52	-1.48
$+ F \rightarrow CI + F$	22.99	2.33	-0.53	0.28	24.34	21.64	0.20	-2.24	27.05
$O + O \rightarrow CI + O_{O}$	-57.62	1.87	0.87	2.87	-5.42	1.02	-7.45	-6.75	-11.86
$F + CH_3 \rightarrow$	-64.25	-2.86	-0.87	-1.96	-2.34	-4.12	-1.79	2.04	-0.55
$I + CH_3F$									
00 + HCl → I + HOCl	10.08	-0.97	1.52	0.42	-3.71	-3.99	2.45	3.67	-3.43
$O + HF \rightarrow OH + HOF$	50.37	-0.22	1.66	-0.01	-1.25	-3.06	6.63	6.88	0.56
$CI + O \rightarrow OH + CIO$	-3.95	0.95	-2.26	-1.26	5.72	4.69	-0.53	-3.62	6.76
$O + CI \rightarrow OH + CIO$	5.74	-0.06	-0.86	-1.18	1.24	-0.49	2.99	0.34	2.97
$O + NO \rightarrow OH + NO_{2}$	-7.48	3.53	1.48	4.30	-5.81	0.07	-4.06	-9.59	-11.65
$O + O_{3} \rightarrow OH + NO_{3}$	-9.13	2.51	1.14	3.48	-13.79	-7.98	-3.59	-7.28	-19.60
$O + N_{s} \rightarrow OH + N_{s} O_{s}$	24.49	3.07	1.84	4.55	-4.66	0.98	-3.53	-8.99	-10.30
$O + NO \rightarrow OH + N_{2}O$	-21.14	1.66	0.59	3.05	-19.49	-11.86	-4.60	-9.14	-27.11
$0 + HO \rightarrow OH + $	-39.31	0.54	0.78	1.78	-12.19	-8.94	-3.07	-3.56	-15.44
ONOH-su									
$\rightarrow OH + O$	-58.31	-1.57	-1.87	-2.62	-17.67	-16.12	-2.99	-0.09	-19.22
+ HNCO									
$O + HCN \rightarrow$	-53.74	-0.75	-0.63	-1.02	-5.84	-4.96	-4.79	-0.74	-6.71
+ HNCO									
$O + O \rightarrow OH + NO$	-53.52	0.79	-0.33	0.88	-12.16	-7.53	-4.00	-4.10	-16.80
$O + N \rightarrow OH + NO$	-82.55	2.66	1.84	3.22	-0.79	3.11	-0.91	-3.68	-4.65
$O + C_2H_2 \rightarrow OH +$	17.58	-0.33	1.01	0.00	3.71	2.74	-0.52	2.75	4.68
rene O ± C H _ → OH ±	64.06-	-1 81	0.36	-0 39	-6 30	ת ה	85 V-	0.68	20.7
してし2114 -7 UII T rane	41.04-	TO'T-	00.0	70.0-	60.0-	T0'0-	00.1	00.0	4.
$O + CH_4 \rightarrow OH +$	-25.39	-2.17	0.82	-0.59	-5.56	-6.05	0.05	3.54	-5.08
thanol									
$O + H_2O \rightarrow OH +$ <i>ms</i> -HCOH	81.74	-2.54	0.41	-2.53	6.85	0.42	8.19	10.95	13.29
$O + {}^{3}CH_{2} \rightarrow OH +$	-62.22	0.15	1.90	0.92	1.02	1.03	3.48	3.73	1.01
$O + trans-HCOH \rightarrow$	-111.81	-0.94	-0.54	-0.39	-11.88	-9.33	-5.63	-3.00	-14.45
I + formic acid									
$+ HO \leftrightarrow HOOH \rightarrow OH + O$	-61.34	-2.72	-1.79	-2.99	-2.37	-4.90	-2.27	2.00	0.16
$\text{Interaction} (C \to OH + C)$	-106.57	-0.05	-0.60	0.25	-11 44	-7.26	-7 96 -7	-5.72	-15.65
ene	10.00T-	00.0-	00.0-	0770	F.F. T.T.	07.	06.1-		70.01-

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table (	U.1: (conti <u> </u>	nued)	A (MD3)	V (		(PULDA) V
5	(.T.)(TS)	$\Delta(\text{UCSD})$	$\Delta$ (MP2.8: $\kappa$ - 00MP2)	$\Delta(MP3:\kappa-00MP2)$	(6.2ЧМ)∆	$\Delta(MP3)$	$\Delta^{(\kappa-}$ 00MP2)	∆(UUMP2)	$\Delta(MP2)$
	12.05	-1.25	-1.35	-3.02	-1.43	-4.98	2.83	4.05	2.13
	-112.40	-3.23	0.60	-0.53	-7.09	-7.94	-1.38	3.00	-6.24
	37.74	-1.84	0.50	-1.87	6.39	0.29	7.52	8.23	12.49
	6.49	-1.61	-0.41	-2.20	0.10	-3.70	4.21	5.73	3.91
	-46.74	-1.92	0.48	-0.14	-5.48	-4.70	-2.88	06.0	-6.27
	-94.50	-1.81	-0.85	-1.57	-11.95	-10.51	-4.74	-0.95	-13.38
	-73.65	-3.35	-1.71	-2.86	-3.38	-5.66	-3.58	1.64	-1.10
	-114.75	0.73	1.68	1.59	-4.76	-2.67	0.85	0.27	-6.84
	29.21	-1.96	0.20	-1.87	1.08	-3.28	5.56	7.49	5.45
	-126.62	0.91	-1.21	-0.25	-10.80	-6.98	-5.87	-6.13	-14.63
	-4.53	2.56	1.69	1.92	-2.18	-1.01	6.44	-0.79	-3.36
	-91.34	-1.42	-0.08	-0.70	-8.00	10.0-	-8.40	-2.73	-11.11
	-39.58 -63.34	0.65	0.06	-1.14 -0.36	-2.93	-4.14 -7 73	4.99 -5.61	2.90	-1.71 -0.63
	-97.22	-0.69	-0.0-	-0.54	-3.70	-3.80	-2.65	0.34	-3.60
	4.10	3.31	1.30	1.11	1.59	0.73	12.06	1.47	2.45
	-120.53	2.24	2.83	4.65	-4.93	0.70	-0.86	-6.30	-10.57
	-85.76	1.45	0.65	1.52	-0.85	1.16	-1.50	-2.56	-2.86
	-191.16	1.74	0.10	1.01	-6.44	-2.25	-3.55	-4.65	-10.63
	-131.19	0.00	-2.45	-1.60	-22.04 10.21	01.01-	-4.07	-0.40 6 00	-21.13 01.09
	-81.69	1.34	-1 78	0.34	2.35	3.56	-1.27	-0.92	115
	-63.31	-2.02	-1.43	-3.00	-1.58	-4.54	2.42	3.54	1.38
	-110.02	-2.96	-1.09	-1.81	-2.99	-4.11	-3.13	0.32	-1.87
	31.27	-0.76	1.66	0.07	-2.66	-4.02	4.82	5.95	-1.29
	-55.70	-2.60	0.11	-1.47	-6.62	-7.41	-0.12	3.19	-5.83
	-15.71	-0.73	-0.16	-1.15	-5.02	-5.54	1.93	2.82	-4.49
	29.43	-1.85	-0.52	-2.04	-1.99	-4.75	3.30	3.80	0.78
	22.62	-1.60	-0.17	-2.59	-0.22	-4.95	6.56	7.62	4.51
	-77.13	-2.63	-0.79	-1.74	-8.05	-9.07	-2.73	1.48	-7.02
	20.76	-1.48	-0.56	-2.93	-0.13	-4.97	6.22	7.38	4.70

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table (	U.1: (conti	nued)				
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\frac{\Delta(\text{MP2.8:}\kappa)}{\text{OOMP2}}$	$\Delta$ (MP3: $\kappa$ -OOMP2)	$\Delta(\mathrm{MP2.5})$	$\Delta({ m MP3})$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$\frac{\text{HCO} + \text{CH}_4}{\text{OH} + \text{C}_2 \text{H}_2} \rightarrow$	27.85	-1.86	-0.07	-2.66	0.03	-5.05	7.14	8.37	5.10
$\begin{array}{c} \operatorname{OH}_{-1} \subset \operatorname{OH}_{3} \\ \operatorname{HCO}_{-1} \subset \operatorname{CH}_{3} \\ \operatorname{OH}_{-1} \subset \operatorname{CH}_{-1} \\ \operatorname{CH}_{-1} \\$	33.63	-1.70	-0.63	-3.14	5.58	-0.02	7.16	7.94	11.17
$\begin{array}{c} \operatorname{OH} + \operatorname{CH}_{2} \\ \operatorname{HCO} + \operatorname{CF}_{2} \\ \operatorname{OH} + \operatorname{FCCF} \end{array}$	44.01	-1.89	-1.35	-3.56	-7.18	-10.60	3.60	4.55	-3.75
$\begin{array}{c} \operatorname{OH} + \operatorname{ICCH} \rightarrow \\ \operatorname{OH} + \operatorname{HCCF} \end{array}$	-82.63	-2.47	-1.28	-2.23	-14.64	-14.46	-3.84	0.83	-14.83
$HCO + {}^{3}CH_{2} \rightarrow HCO + C_{0}H_{2}$	-41.47	0.93	0.25	-0.81	-1.10	-2.82	7.27	3.21	0.62
$HCO + Cl_2 \rightarrow OH + CCl_2$	55.05	-0.33	-0.47	-2.03	7.97	3.59	6.40	4.90	12.34
$HCO + F_2 \rightarrow OH + CF_2$	-48.50	-4.86	-0.28	-2.26	-20.93	-20.97	-1.52	6.16	-20.88
$HOF + CF \rightarrow OH + CF_2$	-75.00	-2.11	-1.60	-2.35	0.51	-0.99	-2.74	0.45	2.00
$\begin{array}{c} HCO + H_2 \rightarrow OH + {}^{3}CH_2 \\ HOF + Si \rightarrow OH + SiF \end{array}$	88.26 88.67	-5.30	-1.37	-4.93 1 45	-0.78	-8.02 2 57	4.59	10.40	6.45 0.42
	-00.01-	-20.05	-0.67	-1.40	06.1-	-10.30	0.06	0.00	-91.98
$HOF + C \rightarrow OH + CF$	-81.54	-1.39	-0.01	-2.09	-20.05-	-1.78	-2.41	0.12	0.26
$HCO + F \rightarrow OH + CF$	39.86	-1.49	0.49	-0.49	2.00	-0.40	1.78	4.66	4.40
$HOF + B \rightarrow OH + BF$	-131.81	-2.30	-2.25	-3.01	-3.39	-4.66	-3.82	-0.33	-2.11
$HNO + B \rightarrow OH + {}^{3}BN$	-8.43	-1.82	-3.93	-5.84	-14.67	-15.69	0.03	2.24	-13.65
$HOO + S_2 \rightarrow OH + S_2O$	-27.50	3.02	0.99	3.72	9.70	10.77	-3.05	-11.80	8.63
$HOO + SO_2 \rightarrow OH + SO_3$	-11.24	-1.77	-1.29	-2.38	-7.72	-6.79	-4.92	1.40	-8.66
$HOO + SO \rightarrow OH + SO_2$	-56.39	2.16	2.34	4.42	-5.62	0.04	-4.50	-7.29	-11.29
$HOO + S \rightarrow OH + SO$	-53.67	1.23	0.56	1.48	-4.04	-1.13	-1.40	-4.17	-6.95
$HOCI + CI \rightarrow OH + CI_2$	1.30	-0.57	-0.55	-1.74	-0.34	-2.56	3.88	3.21	1.88
$HOCI + F \rightarrow OH + CIF$	-3.33	0.13	-0.10	0.14	0.76	1.13	-1.02	-0.31	0.39
$HOF + CI \rightarrow OH + CIF$	-9.63	-1.08	-0.30	-0.86	-0.91	0.36	1.19	0.23	
$\mathrm{HOF} + \mathrm{F} \rightarrow \mathrm{OH} + \mathrm{F}_2$	13.36	1.25	-0.82	-0.58	23.44	19.59	0.56	-1.05	27.28
$HOO + O \rightarrow OH + O_2$	-51.88	1.81	0.01	1.69	-4.18	0.52	-4.46	-6.41	-8.89
$\text{HNO} + \text{N} \rightarrow \text{OH} + \text{N}_2$	-128.18	1.26	0.58	1.72	-15.62	-9.73	-1.98	-3.83	-21.50
$cis$ -HONO $\rightarrow$	-0.48	0.06	0.06	0.12	1.85	1.44	0.27	-0.26	2.27
	00 1	000	0 7 0	CC T		l T		00 0	
	-10.20	0.30	01.0	1.02	-0.80	1.07	-0.97	-2.99	-2.18
$HOCN \rightarrow HNCO$	-24.80	10.1	-0./3	0.03	0.04	1.17	-3.00	-2.80	-1.U9
HONC + HNCO	-84.59	0.56	-1.55	-0.15 -	-2.47	0.43	-6.96	-6.34	-5.36
$HCNO \rightarrow HNCO$	-69.05	-2.61	-1.90	-3.73	-1.13	-4.14	-2.48	4.72	1.88
$HONC \rightarrow HOCN$	-59.78	-0.45	-0.82	-0.18	-2.51	-0.75	-3.90	-3.48	-4.27
$HCNO \rightarrow HOCN$	-44.25	-3.62	-1.17	-3.76	-1.17	-5.31	0.58	7.58	2.97
$HONC \rightarrow HCNO$	-15.54	3.17	0.34	3.59	-1.34	4.56	-4.48	-11.06	-7.24
$trans$ -HCOH $\rightarrow$ H <sub>2</sub> CO	-52.53	0.58	-0.21	0.67	-5.77	-3.70	-2.63	-3.46	-7.84
$cis$ -HCOH $\rightarrow$	-4.77	-0.13	-0.06	-0.05	-0.71	-0.68	-0.35	-0.15	-0.74
trans-HCOH									
$CH_2C \rightarrow C_2H_2$	-45.34	1.40	-0.69	0.87	-6.35	-2.42	-4.38	-6.84	-10.29
$cis$ -HOOO $\rightarrow t_{mnn e}$ -HOOO	1.24	0.58	-0.42	0.25	6.12	4.84	0.00	-4.52	7.40

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

			Table	C.1: (contin	nued)				
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\begin{array}{c} \Delta(\mathrm{MP2.8:}\kappa-\\ \mathrm{OOMP2}) \end{array}$	$\Delta$ (MP3: $\kappa$ -OOMP2)	$\Delta(\mathrm{MP2.5})$	$\Delta({ m MP3})$	$\Delta(\kappa - 00MP2)$	$\Delta(OOMP2)$	$\Delta({ m MP2})$
$cis-N_2H_2 \rightarrow trans-N_2H_2$	-5.72	-0.18	-0.02	-0.02	-0.65	-0.57	-0.40	-0.09	-0.73
allene $\rightarrow$ propyne	-0.51	-0.34	-0.20	0.60	-1.61	0.25	-2.84	-3.29	-3.46
dioxirane $\rightarrow$ formic acid	-91.13	-1.97	-0.75	-0.84	-13.17	-10.26	-4.43	-1.13	-16.07
oxirene $\rightarrow$ ketene	-78.81	-1.12	-0.91	-0.62	-8.80	-7.58	-3.06	-1.63	-10.01
oxirane $\rightarrow$ acetaldehyde	-26.02	-0.11	0.13	0.18	0.91	0.82	1.51	0.22	0.99
$HCNH \rightarrow H_2CN$	-7.98	-1.28	-0.42	-1.57	3.65	0.76	2.72	4.35	6.54
$CH_3 NH \rightarrow CH_2 NH_2$	-6.86	0.67	-0.24	0.53	-1.26	0.61	-2.92	-3.40	-3.12
$F + CH_4 \rightarrow CH_3F + H$	-2.45	2.24	-0.81	1.57	-2.47	2.82	-9.25	-9.12	-7.77
$\mathrm{F} + \mathrm{CH}_3 \mathrm{F} \rightarrow \mathrm{CH}_3 \mathrm{F}_3 + \mathrm{H}$	-14.55	2.47	-1.31	1.24	-2.69	2.80	-9.97	-9.99	-8.18
$F + C_3 H_6 \rightarrow CH_3 F + CH_3$	-17.07	0.68	0.05	1.41	-0.34	2.59	-4.81	-4.39	-3.27
$F + C_2^- H_5^- F \rightarrow C_1^- F_5 + C_1^-$	-22.97	0.75	-0.27	1.20	-0.63	2.43	-5.29	-5.01	-3.69
$F + propane \rightarrow$	-20.82	0.60	-0.14	1.17	-0.22	2.54	-4.84	-4.25	-2.99
$C_2H_5F + CH_3$									
$F + CH_3 NH_2 \rightarrow CH_5 F + NH_5$	-23.16	0.40	0.06	0.84	0.49	2.06	-2.44	-2.12	-1.09
$F + CH_3NH \rightarrow CH_3F + NH$	-31.70	0.25	0.36	1.49	-1.09	1.57	-4.43	-3.72	-3.74
$F + methanol \rightarrow CH_3 F + OH$	-16.68	0.07	-0.15	0.09	0.40	0.86	-0.95	-0.57	-0.06
$F + {}_{3} acetaldehyde \rightarrow CH_{3}F + HCO$	-23.17	1.68	0.26	2.29	0.30	4.56	-5.16	-6.30	-3.97
$F + {}_{3} acetaldehyde \rightarrow CH_{3} + HCOF$	-32.58	1.19	-0.58	1.39	-0.74	3.03	-6.96	-6.70	-4.52
$F + H_2 CO \rightarrow HCOF + H$	-28.61	2.97	-1.82	1.30	-3.01	3.31	-12.15	-12.15	-9.33
$F + H\overline{C}OF \rightarrow F_2CO + H$	-16.37	2.27	-1.72	0.52	-3.37	1.77	-10.46	-9.56	-8.52
$F + glyoxal \rightarrow$ HCOF + HCO	-46.69	1.99	-0.52	2.15	-5.98	0.64	-8.32	-8.76	-12.60
RMSD		4.94	1.59	3.22	8.97	9.24	7.58	10.82	11.99
MSD		1.49	-0.45	0.33	-2.85	-0.80	-3.14	-3.50	-4.90
MIN		-8.60	-5.83	-8.12	-40.24	-38.57	-38.94	-48.81	-51.14
MAX		20.34	5.24	14.63	24.87	31.66	13.56	17.87	27.28

# C.2 RSE43

Table C.2 presents CCSD(T) reaction energies for the RSE43 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for the MP2, OOMP2, and  $\kappa$ -OOMP2 calculations. For this set, we did not use the RI approximation for the CCSD(T), CCSD, or MP3 contributions.

Table C.2: (	JUSD('I')/a	V'I'Z energ	gies (in kcal	/mol)for t	he RSE43	set and erre	ors for oth	er methods	
Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\frac{\Delta(\text{MP2.8:}\kappa)}{\text{OOMP2}}$	$\Delta$ (MP3: $\kappa$ -OOMP2)	$\Delta(\mathrm{MP2.5})$	$\Delta(MP3)$	$\Delta(\kappa - 00 MP2)$	$\Delta(00MP2)$	$\Delta({ m MP2})$
$\begin{array}{c} C_6H_5CH_3+CH_3 \rightarrow \\ C_H & CH_3 \rightarrow \\ C_H & CH_3 \rightarrow \end{array}$	-14.54	-0.82	-1.33	-1.55	2.81	0.96	0.75	0.14	4.67
$\begin{array}{c} C_{6} H_{5} CH_{2} + CH_{4} \\ CH_{2} CHCN + CH_{3} \rightarrow \\ CH_{2} CM + CH_{3} \rightarrow \end{array}$	2.96	0.52	-1.73	-1.40	12.90	9.36	0.81	-2.24	16.44
$CH_2 CUN + CH_4$ $CH_2 CHF + CH_3 \rightarrow$ $CH_2 CF + CH_3 \rightarrow$	7.31	0.44	-0.74	-0.47	5.36	5.17	-0.41	-1.06	5.55
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{CC} + \operatorname{CH}_{4} \\ \operatorname{CH}_{3}\operatorname{CCI}_{3} + \operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{CCI}_{3} + \operatorname{CH}_{3} \rightarrow \end{array}$	-0.26	0.31	-0.41	-0.18	0.87	0.57	0.93	0.11	1.17
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{CU}_{3}+\operatorname{CH}_{4}\\\operatorname{CH}_{3}\operatorname{CF}_{2}\operatorname{CH}_{3}+\operatorname{CH}_{3}\rightarrow\\\operatorname{CH}_{3}\operatorname{CF}_{3}\operatorname{CH}_{3}+\operatorname{CH}_{3}\rightarrow\end{array}$	0.32	0.14	-0.24	-0.46	0.42	0.40	0.16	-0.12	0.44
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{CF}_{2}\operatorname{CH}_{3}+\operatorname{CH}_{4}\\ \operatorname{CH}_{3}\operatorname{CF}_{3}+\operatorname{CH}_{3}\rightarrow\\ \operatorname{CH}_{3}\operatorname{CF}_{4}\operatorname{CH}_{3} \rightarrow\end{array}$	1.85	0.13	-0.21	-0.14	0.35	0.37	0.11	-0.16	0.33
$\begin{array}{c} \operatorname{CH}_{3}\operatorname{CH}_{3}+\operatorname{CH}_{4}\\ \operatorname{CH}_{3}\operatorname{CH}_{2}\operatorname{Cl}+\operatorname{CH}_{3}\rightarrow\\ \operatorname{CH}_{1}\operatorname{CH}_{1}\operatorname{CH}_{1}\operatorname{CH}_{2}\rightarrow\\ \operatorname{CH}_{1}\operatorname{CH}_{1}\operatorname{CH}_{1}\operatorname{CH}_{2}\end{array}$	-2.56	0.41	-0.31	-0.31	0.96	0.80	0.95	0.02	1.12
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{CH}_{2}\operatorname{CH}_{2}\operatorname{CH}_{2}+\operatorname{CH}_{3} \\ \operatorname{CH}_{3}\operatorname{CH}_{2}\operatorname{F}+\operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{CH}_{2}\operatorname{F}+\operatorname{CH}_{3} \end{array} \rightarrow \end{array}$	-1.19	0.18	-0.20	-0.13	0.39	0.43	0.13	-0.18	0.36
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{CH}_{2}\Gamma + \operatorname{CH}_{4}\\ \operatorname{CH}_{3}\operatorname{CH}_{2}\operatorname{OH} + \operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{CH}_{2}\operatorname{OH} + \operatorname{CH}_{3} \rightarrow \end{array}$	-1.76	0.08	-0.23	-0.27	0.31	0.18	0.24	0.09	0.45
$CH_{3}CHCH_{2} + CH_{4}$ $CH_{3}CHCH_{2} + CH_{3} \rightarrow$ $CH_{3}CHCH_{4} + CH_{3} \rightarrow$	-16.91	0.06	-0.80	-0.80	4.73	3.94	00.00	-0.12	5.52
$CH_{3}CHOH_{2} + CH_{4}$ $CH_{3}CHO + CH_{3} \rightarrow$ $CH_{3}CHO + CH_{3} \rightarrow$	-9.41	-0.10	-0.71	-0.73	6.51	4.10	0.69	-0.07	8.93
$CH_{3}CN + CH_{4}$ $CH_{3}CN + CH_{3} \rightarrow$ $CH_{3}CN + CH_{3} \rightarrow$	-8.09	0.18	-0.75	-0.70	7.59	5.02	1.02	-0.31	10.17
$CH_{3}CONH_{2} + CH_{4}$ $CH_{3}CONH_{2} + CH_{3} \rightarrow$ $CH_{3}CONH_{-} + CH_{3} \rightarrow$	-5.96	0.16	-0.37	-0.31	1.21	0.81	0.32	-0.31	1.62
$\begin{array}{c} CH_2 CONH_2 + CH_4 \\ CH_3 CONHCH_3 + CH_3 \rightarrow \\ CH_3 CONHCH_4 + CH_3 \rightarrow \end{array}$	-6.01	0.22	-0.40	-0.30	1.15	0.86	0.29	-0.47	1.45
$CH_{3}COOCH_{3} + CH_{4}$ $CH_{3}COOCH_{3} + CH_{3} \rightarrow$ $CH_{3}COOCH_{4} + CH_{3} \rightarrow$	-6.16	0.15	-0.42	-0.32	1.59	1.04	0.23	-0.44	2.14
$CH_3COOH + CH_3 \rightarrow CH_3 COOH + CH_3 \rightarrow CH_3 COOH + CH_3 \rightarrow CH_3 COOH + CH_3 \rightarrow CH$	-6.02	0.11	-0.38	-0.32	1.68	1.04	0.31	-0.28	2.31
$CH_3-c-C_3H_5+CH_3 \rightarrow CH_3-c-C_3H_5+CH_3 \rightarrow CH_3-c-U_3H_5$	-2.94	0.10	-0.27	-0.30	0.39	0.24	0.29	0.02	0.54
$\begin{array}{c} \operatorname{CH}_{3}^{-c-\operatorname{C3}\operatorname{H}_{5}} + \operatorname{CH}_{4} \\ \operatorname{CH}_{3}^{-F} + \operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{5}^{-F} + \operatorname{CH}_{3} \end{array}$	-3.32	0.38	-0.41	-0.17	0.33	0.68	-0.53	-0.90	-0.02
$\begin{array}{c} \operatorname{CH}_{3}\operatorname{CH}_{3} + \operatorname{CH}_{4} \\ \operatorname{CH}_{3}\operatorname{NH}_{2} + \operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{NH}_{-} + \operatorname{CH}_{3} \end{array}$	-12.01	0.56	-0.53	-0.30	0.47	0.87	-0.43	-1.08	0.08
$CH_3 NH_3^+ + CH_4$ $CH_3 NH_3^+ + CH_3 \rightarrow$ $CH nH_4 + CH_3 \rightarrow$	4.76	0.06	-0.12	-0.10	0.06	0.11	-0.05	-0.15	0.01
$\begin{array}{c} \operatorname{CH}_{2}\operatorname{NH}_{3}^{-1} + \operatorname{CH}_{4} \\ \operatorname{CH}_{3}\operatorname{NH}\operatorname{CH}_{3} + \operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{NH}\operatorname{CH}_{-1} + \operatorname{CH}_{3} \end{array}$	-12.63	0.63	-0.53	-0.27	0.72	1.12	-0.31	-1.14	0.33
$\begin{array}{c} \operatorname{CH}_{2}^{1}\operatorname{MICHO}_{3}+\operatorname{CH}_{4}\\ \operatorname{CH}_{3}\operatorname{NHCHO}+\operatorname{CH}_{3} \rightarrow \\ \operatorname{CH}_{3}\operatorname{MHCHO}+\operatorname{CH}_{3} \end{array}$	-10.80	0.63	-0.40	-0.24	0.98	1.11	0.21	-0.70	0.85
$\begin{array}{c} \operatorname{CH}_{2}^{1}\operatorname{NHOH}+\operatorname{CH}_{3}^{1} \\ \operatorname{CH}_{2}^{1}\operatorname{NHOH}+\operatorname{CH}_{3}^{1} \\ \operatorname{CH}_{2}^{1}\operatorname{NHOH}+\operatorname{CH}_{4}^{1} \end{array}$	-8.42	0.79	-0.69	-0.33	1.25	1.63	-0.27	-1.49	0.88

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation TheoryMade Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy254

$D(T) \Delta$	Tac (CCSD) ∆(MP2.8 OOMP2	$\begin{array}{c c} \text{JIC } \bigcup \\ 8:\kappa^{-} & \bigtriangleup(\\ 2) & 0 \end{array}$	. (сольни МРЗ:к- ОМР2)	Δ(MP2.5)	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
0.67	9	.51	-0.26	0.99	1.32	-0.12	-1.07	0.66
0.34	-0	.78	-0.64	2.17	1.38	0.08	-1.12	2.97
0.34	-0-	.41	-0.27	0.54	0.68	-0.14	-0.57	0.40
0.20	-0-	.34	-0.26	0.37	0.42	-0.12	-0.37	0.32
0.46	0	.37	-0.25	0.58	0.75	-0.03	-0.53	0.42
0.53	-0.	.41	-0.25	0.78	0.97	0.03	-0.65	0.60
0.19	-0.	.41	-0.30	0.16	0.28	-0.38	-0.53	0.04
-0.06	-0.	.03	-0.11	0.41	0.20	0.34	0.40	0.61
0.63	-0.	.70	-0.42	1.30	1.55	-0.16	-1.26	1.05
0.65	-0.	.61	-0.35	1.44	1.57	0.07	-1.13	1.32
0.06	-0	.24	-0.21	0.40	0.39	-0.09	-0.14	0.42
0.51	-0-	.63	-0.38	1.06	1.30	-0.27	-1.06	0.83
0.31	-0.	.40	-0.21	1.07	1.02	0.34	-0.61	1.13
0.29	-0.	.65	-0.58	1.19	1.02	0.21	-0.55	1.36
0.76	-1.	.25	-0.85	6.57	4.97	-0.42	-2.21	8.16
0.94	-0-	.85	-0.40	0.69	1.25	-0.77	-2.16	0.13
0.97	-0.	06.	-0.44	0.53	1.10	-0.96	-2.26	-0.03
0.30	-0.	88.	-0.94	2.09	5.25	0.92	-0.08	8.93
0.06	-0-	.60	-0.68	0.76	0.37	0.42	0.17	1.14
0.09	0.	.05	0.00	0.81	0.58	0.73	0.48	1.05
0.45	0.0	.63 7 1	0.52	3.23	2.43	0.48	0.92	4.10
0.02 0.80		-04 73	-0.42 -155	1.91 0.06	11.00 0.11	71.0 90 O-	10.0-	2.20 0.03
0.97	i C	.05	0.00	12.90	9.36	1.02	0.48	16.44

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

## C.3 BH76RC

Table C.3 presents CCSD(T) reaction energies for the BH76RC set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for all methods.

Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta(MP2.8:\kappa-$ 00MP2)	$\Delta(MP3;\kappa-00MP2)$	$\Delta(\text{MP2.5})$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta(MP2)$
$\rm H + N_{2}O \rightarrow OH + N_{2}$	-64.16	-7.18	-0.16	-6.21	2.00	-8.62	11.86	20.50	12.61
$\mathrm{H} + \mathrm{C} \tilde{\mathrm{H}}_{3} \mathrm{F} \rightarrow \mathrm{H} \mathrm{F} + \tilde{\mathrm{C}} \mathrm{H}_{3}$	-25.68	-1.26	0.68	-0.11	0.89	-0.84	3.51	3.16	2.62
$H + F_{2} \xrightarrow{\sim} HF + F$	-103.25	-5.01	0.54	-1.50	-21.41	-21.60	1.43	5.94	-21.21
$CH_3 + CIF \rightarrow CH_3F + CI$	-54.62	-1.78	-0.60	-1.12	-1.46	-2.09	-2.20	0.79	-0.82
$F^- + CH_3CI \rightarrow C$	-31.56	-1.67	-1.30	-3.02	-0.51	-3.71	2.45	4.67	2.69
$CI^- + CH_3F$ $CH^- CH^- E^-$	96 41	1 00	60 U	06 C	000	69 6	1.02	917	700
$CH_3 CI \cdots r \rightarrow CH_5 F \cdots CI^-$	-20.41	-1.00	00.U-	00.7-	-0.93	-0.03	T./U	4.10	0 <i>1</i> .1
$OH^{-} + CH_{3}F \rightarrow OH^{-}OH + F^{-}$	-19.59	-0.48	-0.70	-0.97	-0.50	-1.06	-2.10	0.68	0.07
$CH_3 CH^- \rightarrow CH_3 CH.OH^- \rightarrow CH_3 OHF^-$	-36.59	-0.05	-1.15	-1.54	-0.66	-1.34	-2.34	0.87	0.01
$H + N_{o} \rightarrow HN_{o}$	3.44	-0.35	-0.72	-2.13	9.90	4.23	5.03	5.41	15.58
$H + CO \rightarrow HCO$	-20.24	0.10	-0.16	-0.58	1.54	0.43	1.84	1.71	2.64
$\mathrm{H} + \mathrm{C}_{2}\mathrm{H}_{4}  ightarrow \mathrm{C}_{2}\mathrm{H}_{5}$	-40.86	-1.87	-0.17	-1.06	0.89	-1.27	-0.14	3.15	3.05
$CH_3 + C_2H_4 \rightarrow C_3H_7$	-28.20	-0.06	-1.03	-0.80	-1.19	-0.80	-4.68	-1.94	-1.58
$HNC \rightarrow HCN$	-15.25	0.34	0.14	0.98	-0.86	1.02	-1.00	-2.94	-2.74
$\rm H + HCl \rightarrow H_2 + Cl$	-2.18	-1.52	1.53	0.61	3.14	0.91	4.21	5.31	5.36
$OH + H_2 \rightarrow H_2O + H$	-16.07	1.73	-1.36	0.60	-3.94	0.60	-8.46	-8.89	-8.47
$CH_3 + H_2 \rightarrow CH_4 + H$	-3.62	1.00	-1.34	-0.63	-2.35	-0.66	-4.02	-4.05	-4.03
$OH + CH_4 \rightarrow H_2O + CH_3$	-12.45	0.74	-0.02	1.23	-1.59	1.26	-4.44	-4.84	-4.44
$OH + NH_3 \rightarrow H_2O + NH_2$	-9.99	0.40	0.18	0.81	-0.34	1.03	-1.86	-2.10	-1.71
$\mathrm{HCl} + \mathrm{CH}_3 \rightarrow \mathrm{Cl} + \mathrm{CH}_4$	-5.81	-0.52	0.20	-0.02	0.79	0.25	0.19	1.25	1.33
$OH + C_2H_6 \rightarrow$	-15.80	0.83	-0.25	0.99	-1.31	1.46	-4.31	-4.87	-4.08
$\mathrm{R_2O}+\mathrm{C_2H_5}$ F + Ho $ ightarrow$ HF + H	-31.81	1.98	-1.47	0.83	-3.92	1.34	-9.76	-10.01	-9.19
$\mathrm{O} + \mathrm{CH}_{A} \rightarrow \mathrm{OH} + \mathrm{CH}_{3}$	6.18	0.56	-0.07	0.37	-0.01	0.94	-1.25	-1.53	-0.96
$\mathrm{H} + \mathrm{PH}_3^{-} \rightarrow \mathrm{H}_{2} + \mathrm{PH}_{2}^{-}$	-20.93	-0.78	1.22	0.75	2.13	1.04	2.70	3.12	3.23
$H + OH \rightarrow H_2 + O$	-2.56	-1.55	1.40	0.26	2.36	-0.28	5.27	5.58	4.99
$\mathrm{H} + \mathrm{H}_2\mathrm{S} \rightarrow \mathrm{H}_2 + \mathrm{HS}$	-13.25	-1.23	1.39	0.69	2.60	0.94	3.28	4.24	4.27
$O + HCI \rightarrow OH + CI$	0.38	0.04	0.13	0.35	0.78	1.19	-1.06	-0.28	0.37
$\rm NH_2 + CH_3 \rightarrow NH + CH_4$	-13.65	-0.27	0.00	-0.18	0.18	-0.22	0.51	0.73	0.58
$NH_2 + C_2H_5 \rightarrow NH + C_2H_2$	-10.30	-0.37	0.23	0.06	-0.10	-0.42	0.38	0.76	0.22
$C_2H_6 + NH_2 \rightarrow C_1H_2 + NH_2$	-5.81	0.43	-0.44	0.18	-0.97	0.43	-2.45	-2.77	-2.37
$\begin{array}{c} \sim_{2^{115}}+1^{1113}\\ \mathrm{NH}_2+\mathrm{CH}_4 \rightarrow \mathrm{NH}_3+\mathrm{CH}_3 \end{array}$	-2.46	0.34	-0.21	0.42	-1.25	0.23	-2.58	-2.74	-2.72
RMSD		1.90	0.84	1.57	4.62	4.51	4.22	5.52	6.34
MSD		-0.64	-0.14	-0.44	-0.54	-0.97	-0.28	0.84	-0.10
MIN		-7.18	-1.47	-6.21	-21.41	-21.60	-9.76	-10.01	-21.21
MAX		1.98	1.53	1.23	9.90	4.23	11.86	20.50	15.58

# C.4 HTBH38

Table C.4 presents CCSD(T) barrier heights for the HTBH38 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for all methods.

		(1000) <b>7</b>	$\Delta(MP2.8:\kappa-OOMP2)$	$\Delta(MP3:\kappa-00MP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	Δ(υυμιτε)	
$\mathrm{H} + \mathrm{HCl} \rightarrow \mathrm{TS1}$	5.21	0.10	1.18	0.60	3.61	2.11	4.39	3.84	5.11
$H_2 + Cl \rightarrow TS1$	7.39	1.61	-0.36	-0.01	0.47	1.19	0.17	-1.47	-0.25
$OH + H_2 \rightarrow TS2$	4.99	1.82	-0.69	0.43	3.10	4.33	-0.99	-3.29	1.88
${ m H_2O}+{ m H} ightarrow{ m TS2}$	21.06	0.09	0.67	-0.17	7.04	3.73	7.48	5.60	10.34
$CH_3 + H_2 \rightarrow TS3$	11.29	1.37	-0.03	0.23	1.29	1.78	0.79	-0.57	0.80
$CH_4 + H \rightarrow TS3$	14.91	0.38	1.30	0.86	3.64	2.44	4.81	3.48	4.83
$OH^+ CH_4 \rightarrow TS4$	5.64	2.61	-0.53	0.97	3.21	5.03	-1.42	-4.56	1.39
$H_2O + CH_3 \rightarrow TS4$	18.09	1.87	-0.50	-0.26	4.80	3.77	3.02	0.28	5.83
$\mathrm{H}\mathrm{\bar{+}}\mathrm{H}_{2} ightarrow\mathrm{TS5}$	9.77	0.30	1.26	0.96	2.67	2.01	3.40	2.67	3.32
${ m H}_2+{ m  ilde{H}} ightarrow{ m TS5}$	9.77	0.30	1.26	0.96	2.67	2.01	3.40	2.67	3.32
$OH + NH_3 \rightarrow TS6$	3.17	3.54	-0.51	1.54	5.06	7.21	-1.43	-6.42	2.91
$H_2O + NH_2 \rightarrow TS6$	13.16	3.14	-0.70	0.74	5.40	6.18	0.43	-4.32	4.63
$HCI + CH_3 \rightarrow TS7$	0.09	1.79	-0.15	0.22	1.39	2.03	0.81	-0.99	0.76
$CH_4 + CI \rightarrow TS7$	5.89	2.31	-0.35	0.24	0.60	1.77	0.62	-2.24	-0.57
$OH + C_2 H_6 \rightarrow TS8$	2.68	2.69	-0.83	0.71	3.63	5.32	-1.41	-4.80	1.94
$H_2O + \overline{C}_2H_5 \rightarrow TS8$	18.48	1.85	-0.58	-0.28	4.94	3.87	2.90	0.06	6.02
$F + H_2 \rightarrow TS9$	1.14	1.20	-1.23	-0.58	3.40	3.84	-0.21	-1.44	2.95
$HF + H \rightarrow TS9$	32.95	-0.78	0.24	-1.41	7.32	2.50	9.56	8.57	12.14
$\mathrm{O}+\mathrm{CH}_4  ightarrow \mathrm{TS10}$	13.61	2.37	-0.51	0.14	3.27	4.11	0.75	-1.58	2.43
$OH + CH_3 \rightarrow TS10$	7.43	1.82	-0.45	-0.23	3.28	3.17	2.01	-0.05	3.39
$\mathrm{H} + \mathrm{PH}_3 \rightarrow \mathrm{TS11}$	2.29	0.59	0.97	0.80	2.16	1.73	2.90	1.84	2.59
$H_2 + PH_2 \rightarrow TS11$	23.21	1.37	-0.25	0.04	0.03	0.70	0.19	-1.28	-0.64
${ m H+OH}  ightarrow { m TS12}$	10.24	0.44	0.83	0.36	5.43	3.73	5.27	3.69	7.14
${ m H}_2 + { m O}  ightarrow { m TS12}$	12.80	1.99	-0.57	0.10	3.08	4.01	0.00	-1.89	2.14
${ m H} + { m H_2S}  ightarrow { m TS13}$	3.16	0.65	0.97	0.79	2.33	1.88	3.06	1.87	2.78
${ m H}_2 + { m HS}  ightarrow { m TS13}$	16.41	1.88	-0.43	0.10	-0.27	0.94	-0.22	-2.38	-1.49
$O + HCl \rightarrow TS14$	9.73	3.58	-1.29	-0.44	5.89	6.88	1.97	-3.54	4.90
$OH + CI \rightarrow TS14$	9.35	3.55	-1.42	-0.79	5.11	5.69	3.03	-3.27	4.53
$\mathrm{CH}_3 + \mathrm{NH}_2  ightarrow \mathrm{TS15}$	7.66	1.99	-0.27	0.05	2.62	2.94	1.81	-0.54	2.29
$\mathrm{CH}_4 + \mathrm{NH}  ightarrow \mathrm{TS15}$	21.31	2.26	-0.27	0.23	2.44	3.16	1.30	-1.27	1.72
$\mathrm{C_2H_5} + \mathrm{NH_2}  ightarrow \mathrm{TS16}$	8.21	2.09	-0.21	0.20	2.50	2.98	1.53	-0.88	2.02
${ m C}_2{ m H}_6+{ m NH} ightarrow{ m TS16}$	18.51	2.46	-0.44	0.15	2.60	3.40	1.16	-1.64	1.79
$\rm NH_2 + C_2 H_6 \rightarrow TS17$	10.39	2.73	-0.39	0.72	2.00	3.74	-0.35	-3.59	0.26
$\rm NH_3 + C_2H_5 \rightarrow TS17$	16.20	2.30	0.04	0.54	2.97	3.31	2.09	-0.83	2.63
$\rm NH_2^+ + CH_4^- \rightarrow TS18$	13.23	2.55	-0.23	0.77	1.81	3.46	-0.18	-3.17	0.16
$\rm NH_3 + CH_3 \rightarrow TS18$	15.68	2.21	-0.03	0.36	3.06	3.23	2.40	-0.43	2.89
<i>s-trans cis</i> - $C_5H_8 \rightarrow TS19$	40.42	4.15	-0.04	1.75	4.40	6.51	-1.02	-7.15	2.29
<i>s-trans cis</i> - $C_5H_8 \rightarrow TS19$	40.42	4.15	-0.04	1.75	4.40	6.51	-1.02	-7.15	2.29
RMSD		2.21	0.71	0.73	3.69	3.88	2.92	3.48	4.04
MSD		1.88	-0.12	0.35	3.25	3.51	1.66	-0.95	2.99
MIN		-0.78	-1.42	-1.41	-0.27	0.70	-1.43	-7.15	-1.49
MAX		4.15	1.30	1.75	7.32	7.21	9.56	8.57	12.14

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

#### **C.5** NHTBH38

Table C.5 presents CCSD(T) barrier heights for the NHTBH38 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for all methods.

Reaction	CCSD(T)	$\Delta(\text{CCSD})$	$\Delta$ (MP2.8: $\kappa$ -OOMP2)	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(00MP2)$	$\Delta({ m MP2})$
$H + N_2O \rightarrow TS1$	17.97	0.47	1.22	0.96	12.59	7.35	4.25	2.00	17.83
$OH + \tilde{N}_3 \rightarrow TS1$	82.13	7.65	1.38	7.18	10.59	15.96	-7.61	-18.49	5.22
$\mathrm{H} + \mathrm{HF} \rightarrow \mathrm{TS2}$	41.25	1.40	1.31	1.34	3.91	3.63	3.59	1.86	4.20
$\mathrm{HF} + \mathrm{H} \rightarrow \mathrm{TS2}$	41.25	1.40	1.31	1.34	3.91	3.63	3.59	1.86	4.20
$\rm H + HCl \rightarrow TS3$	18.23	1.19	1.16	1.00	4.06	3.56	4.47	2.32	4.55
$\rm H + HCl \rightarrow TS3$	18.23	1.19	1.16	1.00	4.06	3.56	4.47	2.32	4.55
$\mathrm{H} + \mathrm{CH}_{3}\mathrm{F} \rightarrow \mathrm{TS4}$	30.09	1.98	1.17	1.52	5.29	5.29	4.51	0.94	5.28
$\text{HF} + \text{CH}_3 \rightarrow \text{TS4}$	55.77	3.24	0.49	1.63	4.39	6.13	1.00	-2.22	2.66
$\mathrm{H} + \mathrm{F}_2  imes \mathrm{TS5}$	1.34	0.35	-0.01	0.35	3.05	1.68	3.35	0.36	4.42
$\mathrm{HF}+\mathrm{ ilde{F}} ightarrow\mathrm{TS5}$	104.60	5.36	-0.54	1.84	24.46	23.28	1.93	-5.58	25.63
$CH_3 + FCI \rightarrow TS6$	5.60	2.61	-0.31	0.45	11.64	10.43	1.97	-2.62	12.85
$CH_3F + CI \rightarrow TS6$	60.22	4.39	0.29	1.57	13.10	12.52	4.17	-3.41	13.67
$F^-$ + $CH_3F \rightarrow TS7$	-1.37	2.66	-0.47	0.69	2.78	4.27	-0.08	-2.53	1.29
$\mathrm{CH}_3\mathrm{F}+\mathrm{F}^-  ightarrow \mathrm{TS7}$	-1.37	2.66	-0.47	0.69	2.78	4.27	-0.08	-2.53	1.29
$F^ CH_3F \rightarrow TS7$	13.15	2.22	-0.24	0.96	2.30	3.98	-0.26	-2.84	0.61
$\mathrm{CH}_{3}\mathrm{F}\mathrm{F}^{-}  ightarrow \mathrm{TS7}$	13.15	2.22	-0.24	0.96	2.30	3.98	-0.26	-2.84	0.61
$Cl^{-}$ + $CH_3Cl$ $\rightarrow$ TS8	1.48	2.43	-0.53	0.02	2.35	3.03	2.49	-1.69	1.67
$CH_3CI + CI^- \rightarrow TS8$	1.48	2.43	-0.53	0.02	2.35	3.03	2.49	-1.69	1.67
$Cl^{-}CH_{3}Cl \rightarrow TS8$	13.02	1.92	-0.47	-0.01	2.17	2.70	2.34	-1.42	1.64
$CH_3CICI^- \rightarrow TS8$	13.02	1.92	-0.47	-0.01	2.17	2.70	2.34	-1.42	1.64
$F^- + CH_3CI \rightarrow TS9$	-13.19	1.78	-0.95	-0.72	2.23	2.03	1.99	-0.41	2.42
$CH_3F + Cl^- \rightarrow TS9$	18.37	3.45	0.35	2.30	2.74	5.74	-0.46	-5.07	-0.27
$F^- \cdots CH_3 CI \rightarrow TS9$	3.24	1.16	-0.52	-0.22	1.61	1.74	1.33	-0.68	1.49
$CH_3FCl^- \rightarrow TS9$	28.65	3.04	0.31	2.16	2.55	5.37	-0.37	-4.84	-0.28
$\rm OH^- + CH_3F \rightarrow TS10$	-3.22	2.79	-0.56	0.61	2.64	4.10	-0.70	-2.76	1.17
$\rm CH_3OH + F^- \rightarrow TS10$	16.37	3.27	0.14	1.58	3.13	5.16	1.40	-3.43	1.10
$OH^- \cdots CH_3F \rightarrow TS10$	11.06	2.38	-0.41	0.78	2.31	3.91	-0.60	-2.86	0.71
$\rm CH_3OHF^-  ightarrow TS10$	47.66	2.43	0.74	2.32	2.97	5.25	1.74	-3.73	0.69
$\mathrm{H} + \mathrm{N}_2  ightarrow \mathrm{TS11}$	14.47	0.53	0.86	0.73	9.58	6.25	3.59	1.99	12.90
$HN_2 \rightarrow TS11$	11.03	0.88	1.58	2.86	-0.33	2.02	-1.43	-3.42	-2.68
$\mathrm{H}+\mathrm{CO}  ightarrow \mathrm{TS12}$	3.11	0.52	0.84	0.78	2.09	1.80	1.92	1.33	2.37
$\mathrm{HCO} \rightarrow \mathrm{TS12}$	23.35	0.42	1.00	1.35	0.55	1.37	0.09	-0.37	-0.27
$\mathrm{H} + \mathrm{C}_{2}\mathrm{H}_{4} \to \mathrm{TS13}$	2.04	0.13	0.51	0.36	5.91	4.80	1.66	1.28	7.02
$\mathrm{C_2H_5}  ightarrow \mathrm{TS13}$	42.90	2.01	0.68	1.42	5.02	6.07	1.80	-1.87	3.97
$CH_3 + C_2H_4 \rightarrow TS14$	5.46	1.37	-0.28	0.01	6.01	5.69	0.54	-0.86	6.34
$C_3H_7 \rightarrow TS14$	33.66	1.43	0.75	0.80	7.20	6.48	5.22	1.07	7.92
$HCN \rightarrow TS15$	48.24	0.45	-0.10	-0.34	2.72	1.10	1.82	1.25	4.34
$\text{HNC} \rightarrow \text{TS15}$	32.99	0.79	0.04	0.64	1.85	2.12	0.82	-1.68	1.59
RMSD		2.53	0.76	1.67	6.59	6.65	2.77	3.90	7.03
MSD		2.07	0.27	1.08	4.76	5.16	1.55	-1.65	4.37
MIN		0.13	-0.95	-0.72	-0.33	1.10	-7.61	-18.49	-2.68
MAX		7.65	1.58	7.18	24.46	23.28	5.22	2.32	25.63

Appendix C. Additional Information: Third-Order Møller-Plesset Perturbation Theory Made Useful? Choice of Orbitals and Scaling Greatly Improves Accuracy

#### C.6 TA13

Table C.6 presents CCSD(T) interactions for the TA13 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for all methods.

Table C.6:	CCSD(T)/aVTZ	interaction	ı energies (	in kcal/mc	ol)for the T.	A13 set and	d errors fo	r other met	hods.
Jomplex	CCSD(T)	$\Delta(\text{CCSD})$	$\frac{\Delta(\text{MP2.8:}\kappa)}{\text{OOMP2}}$	$\Delta(MP3:\kappa-OOMP2)$	$\Delta(MP2.5)$	$\Delta(MP3)$	$\Delta(\kappa - 00MP2)$	$\Delta(OOMP2)$	$\Delta(MP2)$
I <sub>2</sub> 0…Al	-6.50	0.79	-2.05	-2.46	0.49	0.65	-0.09	-0.33	0.33
$H_2^-OBe^+$	-62.90	-0.26	-0.06	-0.47	0.29	-0.49	0.54	1.37	1.08
$H_2^{-}OBr$	-2.62	0.71	-0.12	0.09	0.51	0.82	0.05	-0.67	0.21
HŌH…CH₃	-1.64	0.24	-0.02	0.01	0.13	0.16	0.03	-0.06	0.09
$H_2 0CI$	-2.79	0.78	-0.27	-0.03	0.71	1.00	0.04	-0.77	0.41
$H_2^-OF$	-2.99	1.42	-2.01	-1.27	3.71	4.00	-0.34	-1.94	3.42
H <sub>2</sub> O…Li	-11.72	0.14	-0.08	-0.22	0.63	0.35	0.65	0.53	0.91
$H_2^-OHNH_2^+$	-24.67	0.29	-0.32	-0.36	-0.26	-0.31	-0.59	-0.14	-0.21
$H_2ONH_3^+$	-17.88	0.65	-0.52	-0.35	1.06	1.16	-0.04	-0.41	0.96
$\overline{\rm FH}$ $\overline{\rm BH}_2$	-4.03	0.25	-0.01	0.03	0.17	0.22	-0.01	-0.05	0.12
HFCO <sup>+</sup>	-29.07	1.47	-0.25	-0.52	-3.89	-2.61	0.25	0.87	-5.16
$\rm FHNH_2$	-9.84	0.29	-0.15	-0.10	-0.03	0.07	-0.54	-0.25	-0.13
FHOH	-5.73	0.23	-0.11	-0.09	0.06	0.08	-0.21	-0.06	0.05
RMSD		0.72	0.82	0.81	1.56	1.44	0.35	0.79	1.79
MSD		0.54	-0.46	-0.44	0.28	0.39	-0.02	-0.15	0.16
MIN		-0.26	-2.05	-2.46	-3.89	-2.61	-0.59	-1.94	-5.16
MAX		1.47	-0.01	0.09	3.71	4.00	0.65	1.37	3.42

## C.7 A24

Table C.7 presents CCSD(T) interaction energies for the A24 set in kcal mol<sup>-1</sup> and errors for other methods. All calculations were performed with the aVTZ basis set and the corresponding RI basis was used for the MP2, OOMP2, and  $\kappa$ -OOMP2 calculations. For this set, we did not use the RI approximation for the CCSD(T), CCSD, or MP3 contributions.

$\Delta(MP2)$	-0.07	0.03	-0.17	0.10	-0.01	0.05	0.06	0.04	-0.01	-0.18	-0.07	-0.13	-0.13	-0.17	-0.05	0.05	0.04	0.06	0.04	-0.02	-0.08	0.49	2.40	-0.44	0.52	0.08	-0.44	2.40
$\Delta(OOMP2)$	-0.21	-0.06	-0.18	0.06	-0.08	-0.03	0.01	0.01	-0.47	-0.22	-0.16	-0.12	-0.14	-0.20	-0.04	-0.06	-0.03	0.01	0.00	-0.04	-0.08	-0.35	-0.36	-0.40	0.19	-0.13	-0.47	0.06
$\Delta(\kappa - 00MP2)$	-0.30	-0.18	-0.37	-0.08	-0.03	-0.02	0.05	0.04	-0.63	-0.19	-0.07	-0.07	-0.04	0.08	0.05	0.20	0.11	0.11	0.09	0.02	00.0	0.06	0.06	0.04	0.18	-0.05	-0.63	0.20
$\Delta(MP3)$	0.12	0.05	0.05	-0.01	0.09	0.09	0.04	0.03	0.34	0.02	0.06	-0.01	0.01	0.02	-0.01	0.11	0.09	0.06	0.05	0.05	0.03	0.80	2.20	0.21	0.49	0.19	-0.01	2.20
$\Delta({ m MP2.5})$	0.02	0.04	-0.06	0.05	0.04	0.07	0.05	0.04	0.16	-0.08	0.00	-0.07	-0.06	-0.07	-0.03	0.08	0.06	0.06	0.05	0.02	-0.03	0.65	2.30	-0.11	0.49	0.13	-0.11	2.30
$\Delta(MP3:\kappa-$ OOMP2)	-0.02	-0.07	0.02	-0.11	0.03	-0.01	0.01	0.01	-0.06	0.03	0.03	0.05	0.05	0.08	0.03	0.02	0.05	0.03	0.03	0.04	0.06	0.26	0.13	0.37	0.11	0.04	-0.11	0.37
$\Delta(MP2.8:\kappa-$ OOMP2)	-0.07	-0.08	-0.02	-0.09	0.00	-0.02	0.01	0.01	-0.17	-0.02	-0.01	0.03	0.01	0.03	0.02	0.00	0.03	0.02	0.02	0.02	0.04	0.15	0.04	0.23	0.07	0.01	-0.17	0.23
$\Delta(\text{CCSD})$	0.33	0.22	0.12	0.14	0.25	0.21	0.12	0.10	0.43	0.28	0.25	0.18	0.21	0.31	0.13	0.35	0.21	0.16	0.13	0.09	0.10	0.36	0.41	0.33	0.25	0.23	0.09	0.43
CCSD(T)	-6.26	-4.78	-4.63	-4.32	-3.01	-1.56	-0.72	-0.62	-4.22	-2.44	-1.52	-1.45	-1.31	-1.03	-0.48	-1.37	-0.79	-0.57	-0.50	-0.34	-0.30	0.96	1.15	1.23				
Complex	$H_2ONH_3$	$H_{2}^{-}OH_{2}O$	HCNHCN	$HF \cdots HF$	$\mathrm{NH}_3\mathrm{NH}_3$	$HFCH_4$	$\mathrm{NH}_3\mathrm{CH}_4$	$H_2\bar{O}CH_4$	$CH_2OCH_2O$	$\mathrm{H_2}\mathrm{O}\mathrm{C_2}\mathrm{H_4^-}$	$\operatorname{CH}_2\operatorname{O}\cdots\operatorname{C}_2\operatorname{H}_4$	$\mathrm{C_2H_2C_2H_2}$	${ m NH_3}{}^{-}{ m C_2H_4}^{-}$	$\mathrm{C_2} \mathrm{\ddot{H}_4} \mathrm{\cdots} \mathrm{\ddot{C}_2} \mathrm{\ddot{H}_4}$	$\mathrm{CH}_4 ^{-}\mathrm{C}_2 \mathrm{H}_4$	$BH_{3}$ $CH_{4}$	$\mathrm{CH}_4^{}\mathrm{C}_2\mathrm{H}_6^{}$	$\mathrm{CH}_4 \cdots \mathrm{C}_2 \mathrm{H}_6$	$CH_4 \cdots CH_4$	$ArCH_4$	$ArC_2H_4$	$\mathrm{C_2H_4}\mathrm{C_2H_2}$	$\mathrm{C_{2}H_{4}^{-}C_{2}H_{4}^{-}}$	$C_2H_2^C_2H_2^-$	RMSD	MSD	MIN	MAX

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# Appendix D

# Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

D.1 Equilibrium bond lengths

Table D.1: Root mean square deviations, mean signed deviations, most negative deviation, and most positive deviations in the equilibrium bond lengths for all species are are presented for the CCSD(T) methods utilizing different molecular orbital references.

	$\Delta(\text{CCSD}(\text{T}):$	$\Delta(\text{CCSD}(\text{T}))$ :	$\Delta(\text{CCSD}(\text{T}))$	$\Delta(\text{CCSD}(T))$	$: \Delta(\text{CCSD}(\text{T}))$	$\Delta(\text{CCSD}(T))$	: $\Delta(\text{CCSD}(T))$ :	$\Delta(\text{CCSD}(\text{T}):$
	UHF)	κ-	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M$ -
		OOMP2)						V)
RMSD	0.02258	0.01283	0.01772	0.01338	0.01328	0.01330	0.01328	0.01331
MSD	0.00676	0.00627	0.00468	0.00639	0.00617	0.00616	0.00604	0.00611
MIN	-0.01283	-0.01153	-0.08190	-0.02182	-0.02153	-0.02277	-0.02355	-0.02448
MAX	0.17325	0.08127	0.08457	0.08479	0.08477	0.08473	0.08472	0.08471

Table D	.2: Expe	rimental e	quilibrium	n bond len	gths (in Å	() and erre	ors in the	computed	bond len	gths (in Å	() for the
closed-sł	nell specie	es are pres	ented for	the CCSD	(T) methc	ds utilizin	g differen	t molecula	corbitals.	Root me	an square
deviatio	ns, mean	signed de	viations, 1	most nega	tive deviat	ions, and	most posi	tive deviat	ions (all i	n Å) for t	the set of
species <i>ɛ</i>	und subse	ts are pres	ented. Ex	rperimenta	l values we	ere compile	ed by Hub	er and Her	zberg[346]		
Row 2– Row 2	Dimer	State	Expt.	$\Delta(\text{CCSD}(T))$ UHF)	$\frac{\Delta(\text{CCSD}(T))}{\kappa \text{-OOMP2}}$	$\Delta(\text{CCSD}(T))$ 00MP2)	$\Delta(CCSD(T)$ BLYP)	$\Delta(CCSD(T))$ : $\Delta(CCSD(T))$ B97M-rV)	$\Delta(\text{CCSD}(T))$ : B97)	$\Delta(CCSD(T))$ : $\omega B97X-V)$	$\Delta(CCSD(T); \omega B97M-V)$
	LiH	$X^1\Sigma^+$	1.5957	0.00180	0.00181	0.00181	0.00182	0.00182	0.00182	0.00181	0.00181
	${\rm Li}_2$	$X^1\Sigma_g^+$	2.6729	0.00462	0.00463	0.00463	0.00465	0.00466	0.00463	0.00462	0.00462
	LiF	$X^1\Sigma^+$	1.56386	0.00439	0.00462	0.00477	0.00511	0.00480	0.00478	0.00472	0.00476
	$BeH^+$	$X^1\Sigma^+$	1.3122	0.00028	0.00028	0.00028	0.00029	0.00029	0.00029	0.00029	0.00029
	BeO	$X^1\Sigma^+$	1.3309	0.00544	0.00556	0.00735	0.00676	0.00591	0.00598	0.00546	0.00560
	BH	$X^1\Sigma^+$	1.2324	-0.00038	-0.00047	-0.00080	-0.00065	-0.00075	-0.00081	-0.00082	-0.00081
	ВF	$X^1\Sigma^+$	1.26259	0.00606	0.00642	0.00654	0.00687	0.00654	0.00658	0.00649	0.00651
	$C_2$	$X^1\Sigma_g^+$	1.24253	-0.00631	0.00372	0.00464	-0.02182	-0.02153	-0.02277	-0.02355	-0.02448
	co	$X^1\Sigma^+$	1.12832	0.00411	0.00429	0.00432	0.00445	0.00430	0.00433	0.00428	0.00428
	$\mathrm{N}_2$	$X^1\Sigma_q^+$	1.09769	0.00260	0.00268	0.00268	0.00275	0.00267	0.00270	0.00267	0.00268
	+0u	$X^1\Sigma^+$	1.06322	0.00245	0.00258	0.00257	0.00270	0.00258	0.00261	0.00258	0.00258
	-HO	$X^1\Sigma^+$	0.97	-0.00330	-0.00322	-0.00310	-0.00321	-0.00326	-0.00325	-0.00327	-0.00327
	HF	$X^1\Sigma^+$	0.91681	0.00349	0.00355	0.00356	0.00356	0.00351	0.00352	0.00352	0.00352
	$F_2$	$X^1\Sigma_g^+$	1.41193	-0.01148	0.00452	0.00446	0.00447	0.00449	0.00452	0.00454	0.00456
	RMSD	0		0.00489	0.00385	0.00415	0.00706	0.00689	0.00718	0.00732	0.00755
	MSD			0.00098	0.00293	0.00312	0.00127	0.00115	0.00107	0.00095	0.00090
	MIN			-0.01148	-0.00322	-0.00310	-0.02182	-0.02153	-0.02277	-0.02355	-0.02448
	MAX			0.00606	0.00642	0.00735	0.00687	0.00654	0.00658	0.00649	0.00651
Row $2^-$	Dimer	State	Expt.	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$	$: \Delta(CCSD(T))$	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :
Row $3$				UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	$_{ m NaH}$	$X^1\Sigma^+$	1.8874	0.00481	0.00481	0.00479	0.00498	0.00483	0.00490	0.00493	0.00487
	NaLi	$X^1\Sigma^+$	2.81	0.08249	0.08127	0.08457	0.08479	0.08477	0.08473	0.08472	0.08471
	NaF	$X^1\Sigma^+$	1.92595	0.00882	0.00904	0.00926	0.00968	0.00936	0.00930	0.00918	0.00919
	$MgH^+$	$X^1\Sigma^+$	1.6519	0.00387	0.00387	0.00384	0.00415	0.00386	0.00402	0.00400	0.00389
	AIH	$X^1\Sigma^+$	1.6478	0.00092	0.00039	0.00039	0.00043	0.00046	0.00041	0.00039	0.00040
	AIF	$X^1\Sigma^+$	1.65437	0.00577	0.00607	0.00623	0.00663	0.00625	0.00629	0.00617	0.00620
	$SiH^+$	$X^1\Sigma^+$	1.5041	-0.00065	-0.00089	-0.00089	-0.00085	-0.00083	-0.00086	-0.00087	-0.00086
	SiO	$X^1\Sigma^+$	1.50974	0.00527	0.00548	0.00600	0.00597	0.00563	0.00563	0.00541	0.00541
	PN	$X^1\Sigma^+$	1.49087	-0.00451	0.00504	0.00523	0.00515	0.00496	0.00502	0.00497	0.00499
	$\operatorname{BeS}$	$X^1\Sigma^+$	1.74153	0.17325	0.00814	0.00821	0.00832	0.00821	0.00815	0.00812	0.00800
	CS	$X^1\Sigma^+$	1.53494	0.00444	0.00471	0.00495	0.00504	0.00479	0.00481	0.00465	0.00468
	$NS^+$	$X^1\Sigma^+$	1.44	-0.00806	0.00120	0.00125	0.00135	0.00117	0.00125	0.00119	0.00122
	HCI	$X^1\Sigma^+$	1.27455	-0.00054	-0.00053	-0.00052	-0.00054	-0.00053	-0.00053	-0.00054	-0.00053
	LiCl	$X^1\Sigma^+$	2.02067	0.00775	0.00777	0.00779	0.00790	0.00784	0.00782	0.00778	0.00780
	BCI	$X^1\Sigma^+$	1.7159	0.00408	0.00429	0.00460	0.00511	0.00479	0.00470	0.00436	0.00445
	CIF	$X^1\Sigma^+$	1.62831	0.00677	0.00682	0.00686	0.00662	0.00647	0.00653	0.00658	0.00660
	RMSD			0.04824	0.02097	0.02180	0.02188	0.02183	0.02182	0.02180	0.02180
	MSD			0.01841	0.00922	0.00954	0.00967	0.00950	0.00951	0.00944	0.00944
	MIN			-0.00806	-0.00089	-0.00089	-0.00085	-0.00083	-0.00086	-0.00087	-0.00086

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

0.08471	$\Delta(CCSD(T))$ :	$\omega B97M-V$	0.01382	0.01505	0.01033	0.00874	0.01582	0.01305	0.01275	0.00874	0.01582	$\Delta(CCSD(T))$ :	$\omega B97M-V$	0.01626	0.00650	-0.02448	0.08471
0.08472	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	0.01386	0.01504	0.01031	0.00873	0.01582	0.01305	0.01275	0.00873	0.01582	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	0.01622	0.00652	-0.02355	0.08472
0.08473	$\Delta(CCSD(T))$ :	B97)	0.01387	0.01520	0.01051	0.00874	0.01575	0.01310	0.01281	0.00874	0.01575	$\Delta(CCSD(T))$ :	B97)	0.01621	0.00660	-0.02277	0.08473
0.08477	$\Delta(CCSD(T))$ :	B97M-rV)	0.01389	0.01521	0.01053	0.00871	0.01578	0.01312	0.01282	0.00871	0.01578	$\Delta(CCSD(T))$ :	B97M-rV)	0.01617	0.00663	-0.02153	0.08477
0.08479	$\Delta(CCSD(T))$ :	BLYP)	0.01392	0.01540	0.01076	0.00882	0.01578	0.01322	0.01294	0.00882	0.01578	$\Delta(CCSD(T))$ :	BLYP)	0.01624	0.00678	-0.02182	0.08479
0.08457	$\Delta(CCSD(T))$ :	OOMP2)	0.01382	0.01508	0.01057	0.00880	0.01586	0.01311	0.01283	0.00880	0.01586	$\Delta(CCSD(T))$ :	OOMP2)	0.01577	0.00744	-0.00310	0.08457
0.08127	$\Delta(CCSD(T))$ :	$\kappa$ -OOMP2)	0.01381	0.01491	0.01031	0.00874	0.01583	0.01301	0.01272	0.00874	0.01583	$\Delta(CCSD(T))$ :	$\kappa$ -OOMP2)	0.01520	0.00720	-0.00322	0.08127
0.17325	$\Delta(CCSD(T))$ :	UHF)	0.01382	0.01496	0.01042	-0.00649	0.01591	0.01280	0.00972	-0.00649	0.01591	$\Delta(CCSD(T))$ :	UHF)	0.03311	0.01020	-0.01148	0.17325
	Expt.		2.36079	2.13011	1.92932	1.8934	1.9879										
	State		$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma_g^+$	$X^1\Sigma_g^+$										
MAX	Dimer		NaCl	AlCI	$\operatorname{SiS}$	$\mathbf{P}_2$	$\mathrm{Cl}_2$	RMSD	MSD	MIN	MAX			RMSD	MSD	MIN	MAX
	Row 3–	Row $3$										Closed-	$_{ m shell}$				

Appendix D. Additional Information: Polishing the gold standard: The role of orbital<br/>choice in CCSD(T) frequency prediction2

open-sh	ell specie	s are pres	sented for 1	the CCSD	(T) method	ds utilizing	g different	molecular	· orbitals.	Root me	an square
deviatic	ns, mear	signed d	eviations,	most nega	tive deviat	ions, and	most posi	tive deviat	ions (all j	in Å) for 1	the set of
species	and subs	ets are pre	esented. Ex	sperimenta	l values we	ere compile	d by Hub	er and Her	zberg[346		
Closed- shell	Dimer	State	Expt.	$\Delta(\text{CCSD}(T))$ UHF)	$\frac{\Delta(\text{CCSD}(T))}{\kappa-\text{OOMP2}}$	$\Delta(\text{CCSD}(T))$ 00MP2)	$\Delta(\text{CCSD}(T)$ BLYP)	$\Delta(CCSD(T))$ : D97M-rV)	$\Delta(\text{CCSD}(T))$ B97)	$\Delta(CCSD(T))$ $\omega B97X-V)$	$\Delta(CCSD(T))$ $\omega B97M-V)$
	LiO	$X^2 \Pi_i$	1.695	-0.00222	-0.00218	-0.00209	-0.00173	-0.00199	-0.00203	-0.00211	-0.00207
	$\operatorname{BeH}$	$X^2\Sigma^+$	1.3426	0.00043	0.00043	0.00043	0.00047	0.00055	0.00046	0.00044	0.00048
	$\mathrm{BeF}$	$X^2\Sigma^+$	1.361	0.00378	0.00393	0.00400	0.00425	0.00404	0.00405	0.00399	0.00404
	$\mathrm{B}_2$	$X^3 \Sigma_g^-$	1.59	-0.00741	-0.00463	0.00106	0.00905	0.00398	0.00280	-0.00081	0.00335
	ΒÑ	$X^3\Pi$	1.281	0.04824	0.04909	0.04915	0.04917	0.04909	0.04908	0.04912	0.04914
	BO	$X^2\Sigma^+$	1.2045	0.00498	0.00561	0.00580	0.00600	0.00573	0.00586	0.00571	0.00571
	CH	$X^2 \Pi_r$	1.1199	0.00095	0.00016	0.00016	0.00015	0.00017	0.00015	0.00015	0.00016
	$C_2^{-}$	$X^2 \Sigma_q^+$	1.26821	0.04244	0.00393	0.00407	0.00407	0.00389	0.00394	0.00390	0.00591
	CN	$X^2\Sigma^+$	1.17182	-0.00157	0.00328	0.00340	0.00338	0.00321	0.00324	0.00322	0.00325
	$CO^+$	$X^2\Sigma^+$	1.11522	0.00023	0.00412	0.00441	0.00448	0.00399	0.00419	0.00413	0.00409
	CF	$X^2 \Pi_r$	1.2718	0.00518	0.00577	0.00598	0.00636	0.00588	0.00591	0.00578	0.00580
	HN	$X^3\Sigma^-$	1.03621	0.00182	0.00181	0.00181	0.00180	0.00180	0.00181	0.00180	0.00180
	$^{2}$	$X^2 \Sigma_q^+$	1.11642	-0.00207	0.00252	0.00254	0.00263	0.00248	0.00254	0.00249	0.00251
	N0 N	$X^2 \Pi_r^{J}$	1.15077	-0.00019	0.00288	0.00285	0.00298	0.00289	0.00291	0.00289	0.00289
	NF	$X^3\Sigma^-$	1.31698	0.00503	0.00572	0.00605	0.00644	0.00566	0.00576	0.00566	0.00563
	HO	$X^2 \Pi_i$	0.96966	0.00247	0.00247	0.00247	0.00246	0.00245	0.00246	0.00246	0.00246
	+OH	$X^3\Sigma^-$	1.0289	0.00134	0.00133	0.00133	0.00132	0.00133	0.00132	0.00131	0.00132
	$O_2$	$X^3 \Sigma_g^-$	1.20752	0.00340	0.00335	0.00336	0.00337	0.00333	0.00335	0.00333	0.00334
	$0_2^+$	$X^2 \Pi_g$	1.1164	-0.00158	0.00227	0.00226	0.00235	0.00227	0.00230	0.00228	0.00229
	$\mathrm{O_2}^-$	$X^2\Pi_{g,i}$	1.35	0.00396	0.00373	0.00374	0.00361	0.00365	0.00370	0.00371	0.00370
	OF	$X^2\Pi$	1.35411	0.00324	0.00492	0.00508	0.00516	0.00459	0.00472	0.00466	0.00463
	HF+	$X^2 \Pi_i$	1.00105	0.00302	0.00302	0.00302	0.00302	0.00299	0.00300	0.00300	0.00301
	$F_{2}^{+}$	$X^2 \Pi_{g,i}$	1.3119	-0.00197	-0.00024	-0.00058	-0.00024	-0.00014	-0.00001	0.0000	0.00011
	$\mathbf{F_2}^-$	$X^2 \Sigma_u^+$	1.88	0.03882	0.04568	0.04590	0.04528	0.04553	0.04555	0.04553	0.04544
	RMSD			0.01562	0.01410	0.01413	0.01419	0.01406	0.01406	0.01404	0.01408
	MSD			0.00635	0.00621	0.00651	0.00691	0.00656	0.00654	0.00636	0.00662
	MIN			-0.00741	-0.00463	-0.00209	-0.00173	-0.00199	-0.00203	-0.00211	-0.00207
	MAX			0.04824	0.04909	0.04915	0.04917	0.04909	0.04908	0.04912	0.04914
Row $2-$	Dimer	$\operatorname{State}$	Expt.	$\Delta(CCSD(T))$ :	$\Delta(CCSD(T))$ :	$\Delta(\text{CCSD}(T))$ :	$\Delta(CCSD(T))$	$: \Delta(CCSD(T))$	$\Delta(CCSD(T))$	$\Delta(CCSD(T))$ :	$\Delta(\text{CCSD}(T))$ :
Row 3				UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	MgH	$X^{4}\Sigma^{+}$	1.7297	0.00203	0.00201	0.00201	0.00242	0.00209	0.00229	0.00220	0.00211
	MgF	$X^2\Sigma^+$	1.75	0.00585	0.00605	0.00618	0.00658	0.00628	0.00629	0.00617	0.00616
	$AIH^+$	$X^2\Sigma^+$	1.6018	0.00589	0.00583	0.00582	0.00586	0.00581	0.00582	0.00583	0.00582
	SiH	$X^2 \Pi_r$	1.5201	0.00085	0.00030	0.00031	0.00035	0.00037	0.00033	0.00032	0.00035
	$\operatorname{SiF}$	$X^2 \Pi_r$	1.6011	0.00604	0.00638	0.00657	0.00700	0.00662	0.00662	0.00647	0.00650
	Н	$X^{3}\Sigma^{-}$	1.42234	-0.00007	-0.00006	-0.00006	-0.00004	-0.00004	-0.00005	-0.00006	-0.00005
	$^{+}\mathrm{Hd}$	$X^2 \Pi_r$	1.4352	-0.01149	-0.01153	-0.01153	-0.01149	-0.01149	-0.01151	-0.01151	-0.01150
	-Hd	$X^2 \Pi_i$	1.407	0.02776	0.02776	0.02777	0.02781	0.02777	0.02780	0.02778	0.02780
	CP CP	$X^4\Sigma^+$	1.5622	-0.00410	0.00333	0.00491	0.00447	0.00362	0.00385	0.00334	0.00368
	D'A	$X^{\pm}\Pi_r$	L 4759	0.00602	0.00626	0.00649	0.00639	0.00619	0.00018	0.00608	0.00609

Table D.3: Experimental equilibrium bond lengths (in Å) and errors in the computed bond lengths (in Å) for the

Appendix D. Additional Information: Polishing the gold standard: The role of orbital<br/>choice in CCSD(T) frequency prediction2
0.00133	0.00660	0.00691	-0.00043	0.00532	0.00067	0.00498	0.00522	0.00586	-0.00147	0.00610	0.00624	0.00487	0.00650	0.00771	0.00440	-0.01150	0.02780	$\Delta(CCSD(T))$ :	$\omega B97M-V)$	0.00790	0.01291	0.00839	0.01422	-0.00206	0.00597	0.01004	0.00684	0.00928	0.00803	-0.00206	0.01422	$\Delta(CCSD(T))$ :	$\omega B97M-V)$	0.01108	0.00629	-0.01150	0.04914
0.00126	0.00658	0.00688	-0.00043	0.00531	0.00052	0.00492	0.00520	0.00586	-0.00147	0.00607	0.00614	0.00487	0.00649	0.00769	0.00437	-0.01151	0.02778	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	0.00797	0.01305	0.00842	0.01417	-0.00207	0.00597	0.01004	0.00684	0.00931	0.00805	-0.00207	0.01417	$\Delta(CCSD(T))$ :	$\omega B97X-V)$	0.01106	0.00616	-0.01151	0.04912
0.00134	0.00670	0.00707	-0.00042	0.00542	0.00078	0.00500	0.00523	0.00591	-0.00147	0.00613	0.00656	0.00516	0.00652	0.00776	0.00448	-0.01151	0.02780	$\Delta(CCSD(T))$ :	B97)	0.00800	0.01342	0.00840	0.01437	-0.00216	0.00597	0.01004	0.00681	0.00941	0.00811	-0.00216	0.01437	$\Delta(\text{CCSD}(T))$ :	B97)	0.01110	0.00629	-0.01151	0.04908
0.00139	0.00667	0.00707	-0.00043	0.00536	0.00044	0.00500	0.00519	0.00588	-0.00146	0.00618	0.00655	0.00510	0.00654	0.00775	0.00445	-0.01149	0.02777	$\Delta(CCSD(T))$ :	B97M-rV)	0.00798	0.01368	0.00841	0.01444	-0.00220	0.00601	0.01004	0.00679	0.00947	0.00814	-0.00220	0.01444	$\Delta(\text{CCSD}(T))$ :	B97M-rV)	0.01110	0.00629	-0.01149	0.04909
0.00159	0.00711	0.00754	-0.00043	0.00559	0.00137	0.00515	0.00539	0.00605	-0.00147	0.00621	0.00708	0.00568	0.00646	0.00790	0.00469	-0.01149	0.02781	$\Delta(CCSD(T))$ :	BLYP)	0.00806	0.01355	0.00844	0.01467	-0.00214	0.00606	0.01007	0.00685	0.00952	0.00820	-0.00214	0.01467	$\Delta(CCSD(T))$ :	BLYP)	0.01122	0.00655	-0.01149	0.04917
0.00198	0.00671	0.00712	-0.00044	0.00540	0.00154	0.00512	0.00578	0.00629	-0.00147	0.00609	0.00672	0.00537	0.00647	0.00784	0.00463	-0.01153	0.02777	$\Delta(CCSD(T))$ :	OOMP2)	0.00793	0.01411	-0.07907	0.01432	-0.08190	0.00606	0.01019	0.00694	0.04125	-0.01268	-0.08190	0.01432	$\Delta(CCSD(T))$ :	OOMP2)	0.01884	0.00395	-0.08190	0.04915
0.00158	0.00652	0.00685	-0.00044	0.00519	0.00048	0.00501	0.00559	0.00617	-0.00147	0.00606	0.00609	0.00475	0.00644	0.00770	0.00438	-0.01153	0.02776	$\Delta(CCSD(T))$ :	$\kappa$ -OOMP2)	0.00792	0.01346	0.00846	0.01415	-0.00204	0.00596	0.01007	0.00687	0.00938	0.00811	-0.00204	0.01415	$\Delta(CCSD(T))$ :	$\kappa$ -OOMP2)	0.01110	0.00611	-0.01153	0.04909
0.00107	0.00616	0.00642	-0.00046	0.00431	-0.00616	-0.00140	0.00498	-0.00047	-0.00149	0.00604	0.00576	0.00466	0.00581	0.00752	0.00308	-0.01149	0.02776	$\Delta(CCSD(T))$ :	UHF)	0.00793	0.01279	0.00857	0.01415	-0.01283	-0.00514	0.01018	-0.00387	0.01006	0.00397	-0.01283	0.01415	$\Delta(CCSD(T))$ :	UHF)	0.01197	0.00506	-0.01283	0.04824
1.54	1.5897	1.5003	1.3409	1.6092	1.4954	1.49402	1.48109	1.424	1.31468	1.7971	1.645	1.6144	1.56963					Expt.		2.1991	2.029	2.246	2.058	1.9859	1.9009	1.8892	1.825					Expt.					
$X^3\Sigma^-$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Pi$	$X^3\Sigma^-$	$X^2 \Pi_i$					State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^3 \Sigma_g^-$	$X^2\Pi_r^-$	$X^2 \Pi_u$	$X^2 \Pi_r$	$X^3 \Sigma_g^-$	$X^2 \Pi_{g,r}$					State					
$PO^{-}$	ΡF	$PF^+$	HS	BS	$CS^+$	NS	SO	$^{\rm SO^+}$	$HCl^+$	BeCl	CCI	NCI	OCI	RMSD	MSD	MIN	MAX	Dimer		MgCl	AlS	$\mathrm{Si}_2$	SiCI	$P_2^+$	PS	$\mathbf{S}_2$	$s_2^+$	RMSD	MSD	MIN	MAX	Dimer		RMSD	MSD	MIN	MAX
																		Row 3-	Row 3													Open-	$_{\rm shell}$				

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

Table D.4: Root mean square deviations, mean signed deviations, most negative deviation, and most positive deviations in the equilibrium bond lengths for all species are are presented for the CCSD methods utilizing different molecular orbital references.

	$\Delta$ (CCSD:	$\Delta$ (CCSD:						
	UHF)	$\kappa$ -	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M$ -
		OOMP2)						V)
RMSD	0.02276	0.01352	0.020060	0.01465	0.01442	0.01445	0.01435	0.01437
MSD	-0.00021	-0.00259	-0.00454	-0.00245	-0.00202	-0.00205	-0.00177	-0.00182
MIN	-0.03572	-0.02706	-0.10195	-0.02699	-0.02629	-0.02621	-0.02596	-0.02589
MAX	0.17058	0.08265	0.08996	0.08987	0.08983	0.09008	0.09041	0.08989

Table D.	.5: Exper	imental eq	uilibrium	bond lengt	ths (in Å) a	and errors i	in the com	puted bond	l lengths (	in Å) for t	he closed-
shell spe	scies are l	presented f	for the CC	SD metho	ds utilizing	g different	molecular	orbitals.	Root mean	n square d	eviations,
mean sig	gned dev.	iations, me	ost negatir	ve deviatio	ons, and m	nost positir	ve deviatio	ons (all in	Å) for the	e set of sp	ecies and
subsets ;	are prese	nted. Expe	erimental <sup>-</sup>	values wer	e compiled	by Huber	and Herz	berg[346].			
Row 2– Row 2	Dimer	State	Expt.	$\Delta(CCSD: UHF)$	$\Delta$ (CCSD: $\kappa$ -OOMP2)	$\Delta(CCSD: OOMP2)$	$\Delta(CCSD: BLYP)$	$\Delta(CCSD: B97M-rV)$	$\Delta(CCSD: B97)$	$\Delta(CCSD: \omega B97X-V)$	$\Delta$ (CCSD: $\omega$ B97M-V)
	LiH	$X^1\Sigma^+$	1.5957	0.00203	0.00204	0.00203	0.00203	0.00201	0.00203	0.00204	0.00204
	${\rm Li}_2$	$X^1\Sigma_g^+$	2.6729	0.00946	0.00748	0.00746	0.00735	0.00711	0.00737	0.00747	0.00743
	LiF	$X^1\Sigma^+$	1.56386	0.00041	-0.00048	-0.00073	-0.00097	-0.00059	-0.00060	-0.00054	-0.00061
	${ m BeH^+}$	$X^1\Sigma^+$	1.3122	0.00001	0.0001	0.00000	0.00000	0.00002	0.00000	0.00000	0.00000
	BeO	$X^1\Sigma^+$	1.3309	-0.00998	-0.01298	-0.01468	-0.01405	-0.01310	-0.01331	-0.01256	-0.01281
	ВН	$X^1\Sigma^+$	1.2324	-0.00180	-0.00168	-0.00215	-0.00211	-0.00210	-0.00226	-0.00221	-0.00220
	BF	$X^1\Sigma^+$	1.26259	0.00188	0.00081	0.00063	0.00028	0.00073	0.00063	0.00074	0.00070
	$C_2$	$X^1\Sigma_a^+$	1.24253	-0.01475	-0.01391	-0.00579	0.03649	0.03668	0.03615	0.03636	0.03743
	CÕ	$X^1\Sigma^+$	1.12832	-0.00318	-0.00443	-0.00463	-0.00470	-0.00431	-0.00437	-0.00423	-0.00424
	$\mathrm{N}_2$	$X^1 \Sigma_q^+$	1.09769	-0.00473	-0.00565	-0.00580	-0.00584	-0.00551	-0.00559	-0.00544	-0.00545
	+0u	$X^1\Sigma^+$	1.06322	-0.00638	-0.00775	-0.00792	-0.00790	-0.00750	-0.00757	-0.00742	-0.00742
	-HO	$X^1\Sigma^+$	0.97	-0.00732	-0.00821	-0.00854	-0.00826	-0.00803	-0.00802	-0.00793	-0.00796
	HF	$X^1\Sigma^+$	0.91681	0.00062	0.00019	0.00014	0.00005	0.00027	0.00022	0.00022	0.00020
	Ъ2	$X^1 \Sigma_q^+$	1.41193	-0.03572	-0.02017	-0.02067	-0.02037	-0.01982	-0.01987	-0.01972	-0.01969
	RMSD	c		0.01140	0.00850	0.00816	0.01252	0.01236	0.01229	0.01225	0.01249
	MSD			-0.00496	-0.00462	-0.00433	-0.00129	-0.00101	-0.00109	-0.00094	-0.00090
	MIN			-0.03572	-0.02017	-0.02067	-0.02037	-0.01982	-0.01987	-0.01972	-0.01969
	MAX			0.00946	0.00748	0.00746	0.03649	0.03668	0.03615	0.03636	0.03743
Row $2^-$	Dimer	State	Expt.	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(\text{CCSD};$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$
Row 3				ÚHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	$_{\rm NaH}$	$X^1\Sigma^+$	1.8874	0.00651	0.00674	0.00671	0.00700	0.00677	0.00685	0.00698	0.00694
	NaLi	$X^1\Sigma^+$	2.81	0.08327	0.08265	0.08996	0.08987	0.08983	0.09008	0.09041	0.08989
	NaF	$X^1\Sigma^+$	1.92595	0.00475	0.00380	0.00346	0.00338	0.00364	0.00369	0.00382	0.00373
	${ m MgH^+}$	$X^1\Sigma^+$	1.6519	0.00236	0.00239	0.00233	0.00259	0.00236	0.00250	0.00254	0.00243
	AIH	$X^1\Sigma^+$	1.6478	-0.00022	-0.00130	-0.00136	-0.00150	-0.00131	-0.00145	-0.00132	-0.00133
	AIF	$X^1\Sigma^+$	1.65437	0.0003	-0.00018	-0.00044	-0.00083	-0.00029	-0.00041	-0.00027	-0.00032
	$SiH^+$	$X^1\Sigma^+$	1.5041	-0.00247	-0.00290	-0.00293	-0.00299	-0.00293	-0.00295	-0.00289	-0.00289
	SiO	$X^1\Sigma^+$	1.50974	-0.00615	-0.00802	-0.00892	-0.00862	-0.00816	-0.00816	-0.00775	-0.00778
	PN	$X^1\Sigma^+$	1.49087	-0.01339	-0.01050	-0.01150	-0.01088	-0.01047	-0.01042	-0.01002	-0.01001
	$\operatorname{BeS}$	$X^1\Sigma^+$	1.74153	0.17058	-0.00488	-0.00585	-0.00611	-0.00551	-0.00559	-0.00488	-0.00507
	CS	$X^1\Sigma^+$	1.53494	-0.00762	-0.00962	-0.01049	-0.01026	-0.00969	-0.00976	-0.00921	-0.00932
	$NS^+$	$X^1\Sigma^+$	1.44	-0.01821	-0.01577	-0.01669	-0.01611	-0.01558	-0.01560	-0.01519	-0.01523
	HCI	$X^1\Sigma^+$	1.27455	-0.00306	-0.00319	-0.00324	-0.00330	-0.00320	-0.00323	-0.00316	-0.00322
	LiCI	$X^1\Sigma^+$	2.02067	0.00558	0.00549	0.00538	0.00517	0.00536	0.00535	0.00542	0.00534
	BCI	$X^1\Sigma^+$	1.7159	0.00065	0.00024	-0.00024	-0.00090	-0.00046	-0.00041	0.00005	-0.00007
	CIF	$X^1\Sigma^+$	1.62831	-0.00640	-0.00772	-0.00807	-0.00799	-0.00748	-0.00756	-0.00739	-0.00742
	RMSD			0.04796	0.02172	0.02359	0.02352	0.02343	0.02350	0.02352	0.02340
	MSD			0.01357	0.00233	0.00238	0.00241	0.00268	0.00268	0.00295	0.00285
	MIN			-0.01821	-0.01577	-0.01669	-0.01611	-0.01558	-0.01560	-0.01519	-0.01523

Appendix D. Additional Information: Polishing the gold standard: The role of orbital<br/>choice in CCSD(T) frequency prediction2'

0.08989	$\Delta(CCSD:$	$\omega B97M-V)$	0.01272	0.01178	-0.00390	-0.00949	0.00563	0.00935	0.00335	-0.00949	0.01272	$\Delta(CCSD:$	$\omega B97M-V)$	0.01803	0.00142	-0.01969	0.08989
0.09041	$\Delta(CCSD:$	$\omega B97X-V)$	0.01281	0.01184	-0.00376	-0.00945	0.00567	0.00938	0.00342	-0.00945	0.01281	$\Delta(CCSD:$	$\omega B97X-V)$	0.01804	0.00146	-0.01972	0.09041
0.09008	$\Delta(CCSD:$	B97)	0.01269	0.01158	-0.00448	-0.01012	0.00556	0.00947	0.00305	-0.01012	0.01269	$\Delta(CCSD:$	B97)	0.01805	0.00123	-0.01987	0.09008
0.08983	$\Delta(CCSD:$	B97M-rV)	0.01265	0.01164	-0.00447	-0.01009	0.00556	0.00947	0.00306	-0.01009	0.01265	$\Delta(CCSD:$	B97M-rV)	0.01802	0.00126	-0.01982	0.08983
0.08987	$\Delta(CCSD:$	BLYP)	0.01259	0.01126	-0.00500	-0.01054	0.00526	0.00948	0.00271	-0.01054	0.01259	$\Delta(CCSD:$	BLYP)	0.01812	0.00097	-0.02037	0.08987
0.08996	$\Delta(CCSD:$	OOMP2)	0.01265	0.01174	-0.00481	-0.01062	0.00544	0.00963	0.00288	-0.01062	0.01265	$\Delta(CCSD:$	OOMP2)	0.01716	-0.00023	-0.02067	0.08996
0.08265	$\Delta(CCSD:$	$\kappa$ -OOMP2)	0.01276	0.01200	-0.00378	-0.00962	0.00566	0.00944	0.00340	-0.00962	0.01276	$\Delta(CCSD:$	$\kappa$ -OOMP2)	0.01604	-0.00030	-0.02017	0.08265
0.17058	$\Delta(CCSD:$	UHF)	0.01274	0.01222	-0.00253	-0.01803	0.00599	0.01165	0.00208	-0.01803	0.01274	$\Delta(CCSD:$	UHF)	0.03351	0.00452	-0.03572	0.17058
	Expt.		2.36079	2.13011	1.92932	1.8934	1.9879										
	State		$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1\Sigma_g^+$	$X^1\Sigma_g^+$										
MAX	Dimer		NaCl	AlCI	$\operatorname{SiS}$	$\mathbf{P}_2$	$\operatorname{Cl}_2$	RMSD	MSD	MIN	MAX			RMSD	MSD	MIN	MAX
	Row $3-$	Row $3$										Closed-	$_{ m shell}$				

shell spe mean si <sub>i</sub> subsets a	scies are gned dev are prese	iations, me nted. Expe	ost negati erimental	values wer	e compiled	by Huber	allu nerz				
Row 2– Row 2	Dimer	State	Expt.	$\Delta(CCSD: UHF)$	$\Delta$ (CCSD: $\kappa$ -OOMP2)	$\Delta(CCSD: OOMP2)$	$\Delta(CCSD: BLYP)$	$\Delta(CCSD: B97M-rV)$	$\Delta(CCSD: B97)$	$\Delta(CCSD: \omega B97X-V)$	$\Delta(CCSD: \omega B97M-V)$
	LiO	$X^2 \Pi_i$	1.695	-0.00729	-0.00808	-0.00846	-0.00889	-0.00844	-0.00844	-0.00828	-0.00842
	BeH	$X^2\Sigma^+$	1.3426	-0.00045	-0.00051	-0.00051	-0.00054	-0.00045	-0.00054	-0.00053	-0.00054
	$\mathrm{BeF}$	$X^2\Sigma^+$	1.361	0.00033	-0.00042	-0.00056	-0.00085	-0.00049	-0.00054	-0.00048	-0.00054
	$\mathrm{B}_2$	$X^3 \Sigma_g^-$	1.59	-0.02379	-0.02361	-0.00059	-0.02699	-0.02375	-0.02397	-0.02270	-0.02356
	BN	$X^3\Pi$	1.281	0.03714	0.03588	0.03526	0.03517	0.03575	0.03578	0.03600	0.03581
	BO	$X^2\Sigma^+$	1.2045	-0.00266	-0.00412	-0.00455	-0.00470	-0.00420	-0.00436	-0.00407	-0.00411
	CH	$X^2 \Pi_r$	1.1199	-0.00108	-0.00185	-0.00187	-0.00200	-0.00186	-0.00193	-0.00187	-0.00192
	$C_2^{-}$	$X^2 \Sigma^+_g$	1.26821	0.03456	-0.00645	-0.00720	-0.00663	-0.00613	-0.00624	-0.00608	-0.00579
	CN	$X^2\Sigma^+$	1.17182	-0.00895	-0.00851	-0.00865	-0.00830	-0.00805	-0.00816	-0.00805	-0.00825
	CO+	$X^2\Sigma^+$	1.11522	-0.00599	-0.00804	-0.00855	-0.00849	-0.00771	-0.00793	-0.00774	-0.00773
	CF	$X^2 \Pi_r$	1.2718	-0.00126	-0.00325	-0.00361	-0.00402	-0.00339	-0.00341	-0.00321	-0.00325
	HN	$X^3\Sigma^-$	1.03621	-0.00064	-0.00087	-0.00090	-0.00102	-0.00089	-0.00090	-0.00088	-0.00092
	$N_{2}^{+}$	$X^2\Sigma_q^+$	1.11642	-0.00980	-0.00828	-0.00899	-0.00819	-0.00764	-0.00773	-0.00745	-0.00752
	NÕ	$X^2\Pi_r$	1.15077	-0.00682	-0.00878	-0.00907	-0.00899	-0.00852	-0.00860	-0.00842	-0.00844
	$\rm NF$	$X^3\Sigma^-$	1.31698	-0.00409	-0.00726	-0.00797	-0.00844	-0.00722	-0.00735	-0.00700	-0.00696
	НО	$X^2 \Pi_i$	0.96966	-0.00029	-0.00061	-0.00066	-0.00075	-0.00061	-0.00061	-0.00060	-0.00063
	$^{+}OH$	$X^3\Sigma^-$	1.0289	-0.00101	-0.00117	-0.00118	-0.00133	-0.00126	-0.00122	-0.00124	-0.00126
	$O_2$	$X^3\Sigma_g^-$	1.20752	-0.00845	-0.01060	-0.01101	-0.01075	-0.01022	-0.01023	-0.01006	-0.01009
	$O_2^+$	$X^2 \Pi_g$	1.1164	-0.01172	-0.01209	-0.01243	-0.01216	-0.01167	-0.01172	-0.01152	-0.01154
	$0_2^-$	$X^2 \Pi_{g,i}$	1.35	-0.01316	-0.01649	-0.01747	-0.01637	-0.01565	-0.01566	-0.01552	-0.01552
	OF	$X^2\Pi$	1.35411	-0.00842	-0.01446	-0.01657	-0.01550	-0.01338	-0.01356	-0.01296	-0.01283
	HF+	$X^2 \Pi_i$	1.00105	0.00023	-0.00003	-0.00005	-0.00025	-0.00005	-0.00011	-0.00013	-0.00013
	$F_2^+$	$X^2 \Pi_{g,i}$	1.3119	-0.02617	-0.02706	-0.02809	-0.02695	-0.02629	-0.02621	-0.02596	-0.02589
	$F_2^-$	$X^2 \Sigma_u^+$	1.88	0.00653	0.03413	0.03232	0.03358	0.03458	0.03461	0.03464	0.03463
	RMSD			0.01393	0.01437	0.01365	0.01461	0.01421	0.01425	0.01409	0.01413
	MSD			-0.00264	-0.00427	-0.00381	-0.00472	-0.00406	-0.00413	-0.00392	-0.00398
	MIN			-0.02617	-0.02706	-0.02809	-0.02699	-0.02629	-0.02621	-0.02596	-0.02589
4	MAX	ł	1	0.03714	0.03588	0.03526	0.03517	0.03575	0.03578	0.03600	0.03581
Row 2–	Dimer	State	Expt.	$\Delta(\text{CCSD})$	$\Delta(\text{CCSD})$	$\Delta(\text{CCSD}; \Omega)$	$\Delta(\text{CCSD};$	$\Delta(\text{CCSD}; \mathbf{M})$	$\Delta(\text{CCSD};$	$\Delta(\text{CCSD};$	$\Delta(\text{CCSD};$
HOW 3	MaH	$V^2 \nabla^+$	1 7907		R-UUMPZ)	OUMP2)	DLYP)	D9/M-FV)	D 00089	$\frac{\omega B9(A-V)}{0.00076}$	$\frac{\omega B9/(M-V)}{0.00064}$
	MeF	$X^2\Sigma^+$	1.75	0.00142	0.00045	0.00019	-0.00003	0.00024	0.00028	0.00041	0.00031
	+HIH+	$X^2\Sigma^+$	1.6018	0.00274	0.00223	0.00210	0.00194	0.00200	0.00198	0.00206	0.00206
	SiH	$X^2 \Pi_r$	1.5201	-0.00073	-0.00200	-0.00206	-0.00217	-0.00209	-0.00211	-0.00200	-0.00203
	SiF	$X^2 \Pi_r$	1.6011	0.00007	-0.00129	-0.00160	-0.00204	-0.00150	-0.00155	-0.00138	-0.00141
	НЧ	$X^3\Sigma^-$	1.42234	-0.00307	-0.00315	-0.00321	-0.00329	-0.00321	-0.00324	-0.00314	-0.00319
	$^{\rm +Hd}$	$X_{\widetilde{I}}^{2}\Pi_{r}$	1.4352	-0.01396	-0.01406	-0.01409	-0.01414	-0.01409	-0.01412	-0.01405	-0.01409
	-Hd	$X^2 \Pi_i$	1.407	0.02464	0.02448	0.02432	0.02420	0.02434	0.02431	0.02452	0.02441
	PO Cr	$X^2 \Pi_r$	1.4759	-0.00649	-0.00840	-0.0017 71900.0	-0.00890	-0.00848 -0.00848	-0.00842 -0.00842	-0.00810	-0.00810 -0.00810

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

-0.01144	-0.00291	-0.00341	-0.00337	-0.00525	-0.01335	-0.00993	-0.00799	-0.01068	-0.00428	0.00372	-0.00064	-0.00125	-0.00394	0.00836	-0.00360	-0.01409	0.02441	$\Delta(CCSD:$	$\omega B97M-V)$	0.00535	0.00518	-0.00724	0.00988	-0.02196	-0.01009	-0.00350	-0.01171	0.01083	-0.00426	-0.02196	0.00988	$\Delta(CCSD:$	$\omega B97M-V$	0.01150	-0.00385	-0.02589	0.03581
-0.01136	-0.00291	-0.00340	-0.00330	-0.00507	-0.01313	-0.00982	-0.00794	-0.01068	-0.00422	0.00380	-0.00044	-0.00120	-0.00382	0.00832	-0.00351	-0.01405	0.02452	$\Delta(CCSD:$	$\omega B97X-V)$	0.00542	0.00521	-0.00715	0.00997	-0.02187	-0.00998	-0.00339	-0.01162	0.01079	-0.00418	-0.02187	0.00997	$\Delta(CCSD:$	$\omega B97X-V)$	0.01147	-0.00378	-0.02596	0.03600
-0.01167	-0.00311	-0.00365	-0.00341	-0.00558	-0.01370	-0.01031	-0.00821	-0.01099	-0.00425	0.00365	-0.00120	-0.00204	-0.00481	0.00851	-0.00383	-0.01412	0.02431	$\Delta(CCSD:$	$\mathbf{B97}$	0.00527	0.00404	-0.00788	0.00959	-0.02263	-0.01063	-0.00386	-0.01222	0.01111	-0.00479	-0.02263	0.00959	$\Delta(CCSD:$	B97)	0.01165	-0.00409	-0.02621	0.03578
-0.01182	-0.00307	-0.00363	-0.00339	-0.00547	-0.01346	-0.01033	-0.00831	-0.01105	-0.00427	0.00365	-0.00116	-0.00186	-0.00446	0.00850	-0.00381	-0.01409	0.02434	$\Delta(CCSD:$	B97M-rV)	0.00524	0.00344	-0.00771	0.00958	-0.02252	-0.01059	-0.00385	-0.01219	0.01103	-0.00483	-0.02252	0.00958	$\Delta(CCSD:$	B97M-rV)	0.01161	-0.00406	-0.02629	0.03575
-0.01222	-0.00372	-0.00423	-0.00346	-0.00604	-0.01452	-0.01096	-0.00873	-0.01149	-0.00430	0.00346	-0.00184	-0.00311	-0.00621	0.00883	-0.00427	-0.01452	0.02420	$\Delta(CCSD:$	BLYP)	0.00518	0.00352	-0.00848	0.00914	-0.02310	-0.01116	-0.00422	-0.01262	0.01135	-0.00522	-0.02310	0.00914	$\Delta(CCSD:$	BLYP)	0.01197	-0.00460	-0.02699	0.03517
-0.01253	-0.00319	-0.00386	-0.00339	-0.00578	-0.01497	-0.01193	-0.00929	-0.01204	-0.00427	0.00373	-0.00139	-0.00249	-0.00797	0.00910	-0.00438	-0.01497	0.02432	$\Delta(CCSD:$	OOMP2)	0.00530	0.00204	-0.09502	0.00973	-0.10195	-0.01114	-0.00423	-0.01283	0.04982	-0.02601	-0.10195	0.00973	$\Delta(CCSD:$	OOMP2)	0.02168	-0.00723	-0.10195	0.03526
-0.01166	-0.00281	-0.00341	-0.00333	-0.00513	-0.01351	-0.01066	-0.00856	-0.01122	-0.00426	0.00387	-0.00037	-0.00092	-0.00504	0.00852	-0.00369	-0.01406	0.02448	$\Delta(CCSD:$	$\kappa$ -OOMP2)	0.00544	0.00444	-0.00697	0.01015	-0.02187	-0.00996	-0.00351	-0.01185	0.01079	-0.00427	-0.02187	0.01015	$\Delta(CCSD:$	$\kappa$ -OOMP2)	0.01167	-0.00402	-0.02706	0.03588
-0.00968	-0.00123	-0.00140	-0.00320	-0.00411	-0.01279	-0.01009	-0.00657	-0.01295	-0.00421	0.00404	0.00099	0.00039	-0.00068	0.00821	-0.00279	-0.01396	0.02464	$\Delta(CCSD:$	ÙHF)	0.00554	0.00689	-0.00635	0.01053	-0.02823	-0.01566	-0.00251	-0.01718	0.01402	-0.00587	-0.02823	0.01053	$\Delta(CCSD:$	UHF)	0.01184	-0.00317	-0.02823	0.03714
1.54	1.5897	1.5003	1.3409	1.6092	1.4954	1.49402	1.48109	1.424	1.31468	1.7971	1.645	1.6144	1.56963					Expt.		2.1991	2.029	2.246	2.058	1.9859	1.9009	1.8892	1.825										
$X^3\Sigma^-$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2\Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Pi$	$X^3\Sigma^-$	$X^2 \Pi_i$					State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^3 \Sigma_q^-$	$X^2\Pi_r^{J}$	$X^2 \Pi_u$	$X^2 \Pi_r$	$X^3 \Sigma_g^-$	$X^2 \Pi_{g,r}$										
$^{-}\mathrm{Od}$	ΡF	$PF^+$	HS	BS	$CS^+$	NS	SO	$^{\rm SO^+}$	$HCl^+$	BeCl	CCI	NCI	OCI	RMSD	MSD	MIN	MAX	Dimer		MgCl	AlS	${\rm Si}_2$	SiČl	$P_{2}^{+}$ +	PS	$\mathbf{S}_2$	$s_{2}^{+}$	RMSD	MSD	MIN	MAX			RMSD	MSD	MIN	MAX
																		Row $3^-$	Row 3													Open-	$_{ m shell}$				

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

Table D.7: Species and experimental vibrational frequencies (in  $\text{cm}^{-1}$ ) for which we were unable to compute a continuous PES with one or more of the methods surveyed. Experimental values were compiled by Huber and Herzberg[346].

Dimer	State	Expt.
$\mathrm{CH}^-$	$X^{3}\Sigma^{-}$	979.23
$\mathrm{CN}^+$	$a^1\Sigma$	2033.05
$\mathrm{N_2}^-$	$X^2 \Pi_g$	1968
MgO	$X^1\Sigma^+$	785.0
AlN	$X^3 \Pi_i$	746.93
AlO	$X^2\Sigma^+$	979.23
SiN	$X^2\Sigma^+$	1151.3
$Na_2$	$X^1 \Sigma_q^+$	159.124
$Na_2^+$	$X^2 \Sigma_q^{-+}$	126
$Mg_2$	$X^1 \Sigma_q^+$	51.12
MgS	$X^1 \Sigma^+$	528.74
$Al_2$	$X^3 \Sigma_q^-$	350.01
$\operatorname{Cl}_2^+$	$X^2 \Pi_g^{g}$	645.61

## D.2 Species not included in the test set

Tab. D.7 presents the species for which we were not able to obtain a continuous potential energy surface. For some cases this breakdown occurred at the SCF level with the reference curve changing character near equilibrium, leading to a discontinuous change in the CCSD energy. In other cases the *t*-amplitude optimization of CCSD was found to converge to different roots for different points along the surface.

## **D.3** Mean-field $\langle S^2 \rangle$ values

Table D.	8: Mean-	field $\langle S^2 \rangle$	values fo	or the refere	nce method	ds at the C	CSD(T) I	ninimum e	mergy poin	nts used in	the fitting
procedur	e.										
Closed- shell	${ m Row}~2-{ m Row}~2$	Dimer	State	CCSD(T): UHF	$CCSD(T): \kappa-OOMP2$	CCSD(T): OOMP2	CCSD(T): BLYP	CCSD(T): B97M-rV	CCSD(T): B97	$CCSD(T): \omega B97X-V$	$CCSD(T): \omega B97M-V$
		LiH	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\mathrm{Li}_2$	$X^1\Sigma_g^+$	0.31860	0.0000	0.00000	0.0000	0.01065	0.00000	0.00000	0.00000
		LiF	$X^1\Sigma^+$	0.0000	0.0000	0.00000	0.0000	0.0000	0.00000	0.00000	0.00000
		$\mathrm{BeH}^+$	$X^1\Sigma^+$	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
		BeO	$X^1\Sigma^+$	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
		ВН	$X^1\Sigma^+$	0.47646	0.04871	0.00000	0.03506	0.00000	0.00000	0.00000	0.00000
		$\mathrm{BF}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$C_2$	$X^1\Sigma_g^+$	1.65493	0.89140	0.00000	0.85447	0.94029	0.93904	1.01320	0.97020
		QŌ	$X^1\Sigma^+$	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
		$ m N_2$	$X^1 \Sigma_a^+$	0.0000	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000
		+0u	$X^1\Sigma^+$	0.0000	0.0000	0.00000	0.0000	0.0000	0.00000	0.00000	0.00000
		-HO	$X^1\Sigma^+$	0.0000	0.0000	0.00000	0.0000	0.0000	0.00000	0.00000	0.00000
		HF	$X^1\Sigma^+$	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
		$\mathbf{F}_2$	$X^1 \Sigma_g^+$	0.31922	0.0000	0.00000	0.0000	0.0000	0.00000	0.00000	0.00000
	Row $2^-$	Dimer	State	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):
	Row 3			UHF	$\kappa$ -OOMP2	00MP2	BLYP	B97M-rV	B97	$\omega B97X-V$	$\omega B97M-V$
		NaH	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		NaLi	$X^1\Sigma^+$	0.32431	0.03641	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\operatorname{NaF}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$MgH^+$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		AIH	$X^1\Sigma^+$	0.17105	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		AIF	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$SiH^+$	$X^1\Sigma^+$	0.06997	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		SiO	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		PN	$X^1\Sigma^+$	0.70716	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\operatorname{BeS}$	$X^1\Sigma^+$	1.01647	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CS	$X^1\Sigma^+$	0.0000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
		$NS^+$	$X^1\Sigma^+$	0.69154	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		HCI	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		LiCI	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		BCI	$X^{1}\Sigma^{+}$	0.17749	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CCI+	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CIF	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	Row 3– Row 3	Dimer	State	CCSD(T): UHF	CCSD(T): k-OOMP2	CCSD(T): OOMP2	CCSD(T): BLVP	CCSD(T): B97M-rV	CCSD(T): R97	CCSD(T):	CCSD(T): <sup>(1)</sup> B97M-V
		NaCl	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.0000	0.0000	0.00000	0.00000
		AICI	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\operatorname{SiS}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000
		$\mathbf{P}_2$	$X^1\Sigma_a^+$	0.67604	0.00000	0.00000	0.0000	0.0000	0.00000	0.0000	0.00000
		$C_{1_2}$	$X^1 \Sigma_g^+$	0.00000	0.0000	0.00000	0.0000	0.00000	0.00000	0.00000	0.00000
Open-	Row 2–	Dimer	State	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):	CCSD(T):
snell	Kow2		V2H	UHF 0 75070	K-UUIMF'2	OUMP'2	BLYF 0 75971	D9/INFrV	D97 0 75 190	WB91A-V	wB9/IM-V
		DIJ DIJ	$X$ $^{-11i}$	U.139/U	cuaci.u	7.1007.0	U.733/1	0.70019	U.7548U	0.703/2	U.73439

Appendix D. Additional Information: Polishing the gold standard: The role of orbital<br/>choice in CCSD(T) frequency prediction2

0.75052	0.75050	2.57710	2.00340	0.75832	0.75210	0.75835	0.76049	0.76630	0.75220	2.00603	0.75330	0.75279	2.00626	0.75279	2.00583	2.00857	0.75221	0.75631	0.75376	0.75209	0.75292	0.75599	CCSD(T):	$\omega B97M-V$	0.75688	0.75388	0.75049	0.75678	0.75211	0.75174	2.00395	0.75231	0.75394	0.79607	0.75258	2.00607	2.00285	0.75208	0.75266	0.76416	0.79312	0.75736	2.00851	0.75304	0.75257	0.75018 0.75078
0.75107	0.75065	2.90525	2.00361	0.75704	0.75279	0.76265	0.76000	0.76466	0.75333	2.00562	0.75423	0.75324	2.00687	0.75252	2.00511	2.00954	0.75259	0.75710	0.75348	0.75191	0.75308	0.75578	CCSD(T):	$\omega B97 X-V$	0.75546	0.75516	0.75052	0.75903	0.75548	0.75597	2.00641	0.75437	0.75523	0.80446	0.75529	2.01013	2.00654	0.75478	0.75383	0.76101	0.79583	0.75990	2.01195	0.75513	0.75332	0.75016 0.75079
0.75068	0.75065	2.51071	2.00383	0.75698	0.75326	0.75544	0.75958	0.76532	0.75370	2.00706	0.75392	0.75347	2.00795	0.75299	2.00594	2.00981	0.75263	0.75739	0.75378	0.75219	0.75303	0.75597	CCSD(T):	B97	0.75888	0.75330	0.75062	0.75545	0.75485	0.75510	2.00709	0.75448	0.75612	0.78991	0.75438	2.00860	2.00662	0.75464	0.75420	0.76235	0.79111	0.75751	2.00982	0.75413	0.75373	0.75038 0.75078
0.75042	0.75212	2.47622	2.00644	0.76350	0.75428	0.75519	0.76131	0.77442	0.75435	2.01010	0.75417	0.75368	2.00948	0.75419	2.00758	2.01025	0.75292	0.75787	0.75436	0.75306	0.75314	0.75568	CCSD(T):	B97M-rV	0.76283	0.75128	0.75076	0.75413	0.75455	0.75480	2.00850	0.75465	0.75708	0.80531	0.75418	2.00812	2.00876	0.75552	0.75530	0.77086	0.80860	0.75636	2.00929	0.75398	0.75489	0.75219 0.75221
0.75100	0.75058	2.35395	2.00289	0.75571	0.75270	0.75332	0.75447	0.75964	0.75299	2.00481	0.75241	0.75215	2.00507	0.75224	2.00419	2.00526	0.75154	0.75412	0.75224	0.75169	0.75152	0.75283	CCSD(T):	BLYP	0.75946	0.75320	0.75025	0.75539	0.75404	0.75400	2.00492	0.75364	0.75478	0.76599	0.75301	2.00468	2.00410	0.75350	0.75337	0.75858	0.77031	0.75377	2.00507	0.75259	0.75301	0.75008 0.75094
0.75144	0.75103	2.00956	2.00510	0.75971	0.75551	0.75337	0.75580	0.76261	0.75584	2.00985	0.75207	0.75343	2.00989	0.75442	2.00858	2.00751	0.75229	0.75395	0.75390	0.75334	0.75193	0.75533	CCSD(T):	00MP2	0.75866	0.75575	0.75057	0.76425	0.75871	0.75908	2.01010	0.75781	0.75950	0.76087	0.75572	2.00729	2.00788	0.75772	0.75666	0.76520	0.76914	0.75365	2.00667	0.75422	0.75598	0.75019 0.75150
0.75136	0.75096	2.81206	2.00636	0.76145	0.75579	0.75579	0.76257	0.76730	0.75627	2.01005	0.75334	0.75407	2.01047	0.75446	2.00868	2.00962	0.75272	0.75617	0.75447	0.75336	0.75265	0.75746	CCSD(T):	$\kappa$ -OOMP2	0.75797	0.75613	0.75048	0.76616	0.76132	0.76186	2.01225	0.75888	0.76106	0.82118	0.75758	2.00989	2.01047	0.75909	0.75738	0.77318	0.80770	0.75758	2.01059	0.75583	0.75633	0.75016 0.75140
0.75204	0.75083	2.90778	2.05171	0.80067	1.10075	0.75627	1.15755	0.96970	0.76275	2.01660	1.23889	0.79621	2.02283	0.75710	2.01414	2.04908	1.12597	0.79527	0.77257	0.75515	0.82975	0.77276	CCSD(T):	UHF	0.76211	0.75985	0.75036	0.77723	0.79747	0.77677	2.02909	0.79438	0.77186	1.61297	0.77298	2.04447	2.03511	0.77166	0.76481	0.85482	1.47259	1.18951	2.05909	1.18709	0.76226	0.75014 0.75155
$X^{2}\Sigma^{+}$	$X^2\Sigma^+$	$X^{3}\Sigma^{-}_{g}$	$X^{3}\Pi$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^2 \Sigma_g^+$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^{3}\Sigma^{-}$	$X^2 \Sigma_g^+$	$X^2\Pi_r^j$	$X^{3}\Sigma^{-}$	$X^2 \Pi_i$	$X^{3}\Sigma^{-}$	$X^{3}\Sigma_{g}^{-}$	$X^2 \Pi_g$	$X^2 \Pi_{g,i}$	$X^{2}\Pi$	$X^2 \Pi_i$	$X^2\Pi_{g,i}$	$X^2\Sigma_u^+$	State		$X^{2}\Pi$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^2 \Pi_r$	$X^{3}\Sigma^{-}$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Pi_r$	$X^{3}\Sigma^{-}$	$X^{3}\Sigma^{-}$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^{3}\Sigma^{-}$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$
$\operatorname{BeH}$	$\mathrm{BeF}$	$^{ m B_2}$	BN	BO	CH	$C_2^{-}$	CN	$CO^+$	CF	HN	$N_{2}^{+}$	0N N	NF	НО	+OH	$O_2$	$O_2^+$	$0^{-2}_{2}$	OF	$\mathrm{HF}^+$	н Ч	Ъ	Dimer		NaO	MgH	MgF	$AIH^+$	HiS	$\mathrm{SiF}$	НЧ	$\rm PH^+$	$^{-}\mathrm{Hd}$	CP	РО	$PO^{-}$	РF	$PF^+$	HS	BS	$CS^+$	NS	SO	$^{+}$ OS	HCl+	LiCl - BeCl
																							닎	$\sim$																						

Row 2 Row 3

0.75405	2.01037	0.75521	0.75267	CCSD(T):	$\omega B97M-V$	0.75082	0.75552	2.00332	0.75208	0.75166	0.75445	2.00936	0.75392
0.75486	2.01189	0.75629	0.75381	CCSD(T):	$\omega B97X-V$	0.75094	0.75917	2.00422	0.75601	0.75237	0.75791	2.01559	0.75747
0.75493	2.01122	0.75536	0.75348	CCSD(T):	B97	0.75084	0.75498	2.00554	0.75517	0.75289	0.75557	2.00936	0.75457
0.75543	2.01180	0.75543	0.75387	CCSD(T):	B97M-rV	0.75059	0.75466	2.00890	0.75541	0.75510	0.75496	2.00775	0.75386
0.75355	2.00654	0.75289	0.75221	CCSD(T):	BLYP	0.75045	0.75355	2.00441	0.75402	0.75262	0.75340	2.00494	0.75285
0.75656	2.01178	0.75405	0.75456	CCSD(T):	00MP2	0.75093	0.75923	2.00560	0.75850	0.75651	0.75590	2.00876	0.75431
0.75772	2.01532	0.75664	0.75541	CCSD(T):	$\kappa$ -OOMP2	0.75085	0.76188	2.00946	0.76144	0.75476	0.76119	2.01837	0.75862
0.76579	2.03589	0.77010	0.76463	CCSD(T):	UHF	0.75092	0.77492	2.01897	0.77631	1.16866	1.05380	2.06163	1.23007
$X^2\Pi$	$X^{3}\Sigma^{-}$	$X^2 \Pi_i$	$X^2\Pi$	State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^{3}\Sigma_{g}^{-}$	$X^2\Pi_r^{-}$	$X^2 \Pi_u$	$X^2 \Pi_r$	$X^{3}\Sigma_{g}^{-}$	$X^2 \Pi_{g,r}$
CCI	NCI	0CI	FC1+	Dimer		MgCl	AIS	$\mathrm{Si}_2$	SiCI	$P_{2}^{+}$	$P\tilde{S}$	$^{ m S}_2$	$s_2^+$
				Row $3-$	Row 3								

Table D. procedur	9: Mean- e.	-field $\langle S^2 \rangle$	values fc	or the refer	ence metho	ods at the	CCSD m	inimum ene	ergy poin	ts used in	the fitting
Closed- shell	Row 2- Row 2	Dimer	State	CCSD: UHF	CCSD: k-OOMP2	CCSD: OOMP2	CCSD: BLYP	CCSD: B97M-rV	CCSD: B97	CCSD: (),B97X-V	CCSD: (1)B97M-V
TIONS	7 0001	LiH	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		Li,	$X^1 \Sigma_a^+$	0.31690	0.00000	0.00000	0.00000	0.01065	0.00000	0.00000	0.00000
		LiF	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.0000	0.00000	0.00000	0.00000
		$\mathrm{BeH^{+}}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		BeO	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		ВН	$X^1\Sigma^+$	0.47646	0.04871	0.00000	0.03506	0.00000	0.00000	0.00000	0.00000
		$\mathrm{BF}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$_{\rm C}^{\rm C}$	$X^1 \Sigma_g^+$	1.62372	0.90291	0.00000	0.92892	1.00478	1.00057	1.05327	1.02472
		CO	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\mathrm{N}_2$	$X^1\Sigma_g^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		+ON	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		-HO	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		HF	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\mathbf{F}_2$	$X^1\Sigma_g^+$	0.21545	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	Row $2^-$	Dimer	State	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:
	Row 3			UHF	$\kappa$ -OOMP2	00MP2	BLYP	B97M-rV	B97	$\omega B97X-V$	$\omega B97M-V$
		NaH	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		NaLi	$X^1\Sigma^+$	0.32702	0.03641	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\operatorname{NaF}$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$MgH^+$	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		AlH	$X^1\Sigma^+$	0.17105	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		AIF	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		SiH+	$X^{1}\Sigma^{+}$	0.06997	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		SiO	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		PN	$X^1\Sigma^+$	0.65868	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$\operatorname{BeS}$	$X^1\Sigma^+$	1.01646	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CS	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$NS^+$	$X^1\Sigma^+$	0.64038	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		HCI	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		LiCI	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		BCI	$X^{1}\Sigma^{+}$	0.17638	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CCI+	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		CIF	$X^1\Sigma^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	Row 3–	Dimer	State	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:
	Row $3$			UHF	k-OOMP2	00MP2	BLYP	B97M-rV	B97	$\omega B97X-V$	$\omega B97M-V$
		NaCl	$X^{1}\Sigma^{+}$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		AICI	$X^{1}\Sigma^{+}$	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
		212	$X^{-}X$	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000	0.0000.0	0.0000
		$\mathrm{P}_2$	$X^{1}\Sigma^{1}g$	0.60962	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
		$Cl_2$	$X^1\Sigma_g^+$	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
Open- shell	${ m Row}~2^-$	Dimer	State	CCSD: UHF	CCSD: <i>k</i> -OOMP2	CCSD: OOMP2	CCSD: BLYP	CCSD: B97M-rV	CCSD: B97	CCSD: ∞B97X-	$\omega B97M-V$
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Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

		G								
	LiO	$X^{2}\Pi_{i}$	0.75966	0.75599	0.75607	0.75366	0.75613	0.75475	0.75369	0.75435
	BeH	$X^2\Sigma^+$	0.75204	0.75136	0.75144	0.75100	0.75042	0.75068	0.75107	0.75052
	$\mathrm{BeF}$	$X^2\Sigma^+$	0.75084	0.75096	0.75103	0.75058	0.75212	0.75065	0.75065	0.75050
	В,	$X^{3}\Sigma_{a}^{-}$	2.91265	2.82216	2.00956	2.40515	2.50830	2.54165	2.69514	2.61548
	ΒÑ	$X^3\Pi^{J}$	2.04561	2.00633	2.00510	2.00286	2.00650	2.00377	2.00355	2.00337
	BO	$X^2\Sigma^+$	0.79782	0.76122	0.75955	0.75559	0.76313	0.75683	0.75689	0.75813
	CH	$X^2 \Pi_r$	1.10075	0.75579	0.75551	0.75270	0.75428	0.75326	0.75279	0.75210
	$C_2^{-}$	$X^2 \Sigma_g^+$	0.75628	0.75577	0.75338	0.75331	0.75516	0.75534	0.76049	0.75748
	CN	$X^2\Sigma^+$	1.11881	0.76253	0.75590	0.75441	0.76117	0.75942	0.75973	0.76030
	$CO^+$	$X^2\Sigma^+$	0.95712	0.76676	0.76253	0.75924	0.77298	0.76444	0.76379	0.76536
	CF	$X^2 \Pi_r$	0.76271	0.75624	0.75582	0.75297	0.75432	0.75368	0.75331	0.75218
	HN	$X^3\Sigma^-$	2.01660	2.01005	2.00985	2.00481	2.01010	2.00706	2.00562	2.00603
	$^+$ °N	$X^2 \Sigma_q^+$	1.19138	0.75333	0.75212	0.75232	0.75406	0.75377	0.75416	0.75316
	NÕ	$X^2\Pi_r^j$	0.78638	0.75402	0.75344	0.75209	0.75356	0.75333	0.75311	0.75265
	NF	$X^{3}\Sigma^{-}$	2.02281	2.01045	2.00989	2.00505	2.00944	2.00791	2.00683	2.00616
	HO	$X^2 \Pi_i$	0.75710	0.75446	0.75442	0.75224	0.75419	0.75299	0.75252	0.75279
	$^{+}OH$	$X^{3}\Sigma^{-}$	2.01414	2.00868	2.00858	2.00419	2.00758	2.00594	2.00511	2.00583
	0,	$X^3\Sigma_a^-$	2.04740	2.00967	2.00772	2.00516	2.01004	2.00958	2.00927	2.00834
	-'0	$X^2\Pi_a$	1.08640	0.75274	0.75233	0.75153	0.75289	0.75259	0.75255	0.75217
	0,- _	$X^2\Pi_{a,i}$	0.79295	0.75622	0.75417	0.75405	0.75767	0.75718	0.75684	0.75610
	OF	$X^2\Pi^2$	0.76696	0.75440	0.75391	0.75219	0.75425	0.75364	0.75335	0.75353
	$HF^+$	$X^2 \Pi_i$	0.75510	0.75334	0.75332	0.75168	0.75304	0.75218	0.75189	0.75207
	н° Н	$X^2 \Pi_{q,i}$	0.79656	0.75276	0.75213	0.75150	0.75312	0.75298	0.75302	0.75285
	بر 12	$X^2 \Sigma_u^+$	0.77581	0.75739	0.75531	0.75281	0.75565	0.75591	0.75571	0.75592
Row $2-$	Dimer	State	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:	CCSD:
Row $3$			UHF	$\kappa$ -OOMP2	00MP2	BLYP	B97M-rV	B97	$\omega B97X-V$	$\omega B97M-V$
	NaO	$X^2\Pi$	0.76205	0.75794	0.75863	0.75938	0.76271	0.75882	0.75544	0.75684
	MgH	$X^2\Sigma^+$	0.75985	0.75613	0.75575	0.75320	0.75128	0.75330	0.75516	0.75388
	MgF	$X^2\Sigma^+$	0.75036	0.75048	0.75057	0.75024	0.75076	0.75061	0.75051	0.75048
	$AIH^+$	$X^2\Sigma^+$	0.77723	0.76584	0.76400	0.75531	0.75403	0.75536	0.75887	0.75666
	SiH	$X^2 \Pi_r$	0.79747	0.76132	0.75871	0.75404	0.75455	0.75485	0.75548	0.75211
	$\operatorname{SiF}$	$X^2 \Pi_r$	0.77677	0.76186	0.75908	0.75400	0.75479	0.75510	0.75597	0.75174
	ΗЧ	$X^3\Sigma^-$	2.02913	2.01224	2.01007	2.00492	2.00848	2.00709	2.00641	2.00394
	$_{\rm HH^+}$	$X^2 \Pi_r$	0.79445	0.75887	0.75780	0.75364	0.75464	0.75447	0.75437	0.75231
	$^{-}\mathrm{Hd}$	$X^2 \Pi_i$	0.77189	0.76106	0.75949	0.75478	0.75707	0.75612	0.75523	0.75394
	CP	$X^2\Sigma^+$	1.59712	0.81614	0.76132	0.76525	0.80179	0.78701	0.79906	0.79240
	РО	$X^2 \Pi_r$	0.77267	0.75758	0.75576	0.75300	0.75413	0.75434	0.75522	0.75252
	$PO^{-}$	$X^3\Sigma^-$	2.04349	2.00978	2.00720	2.00460	2.00799	2.00841	2.00993	2.00591
	ΡF	$X^{3}\Sigma^{-}$	2.03511	2.01048	2.00788	2.00411	2.00875	2.00663	2.00655	2.00284
	$PF^+$	$X^2 \Pi_r$	0.77166	0.75910	0.75772	0.75350	0.75550	0.75464	0.75478	0.75207
	HS	$X^2 \Pi_i$	0.76481	0.75737	0.75664	0.75337	0.75530	0.75420	0.75383	0.75266
	BS	$X^2\Sigma^+$	0.85146	0.77281	0.76503	0.75844	0.77046	0.76215	0.76083	0.76394
	$CS^+$	$X^2\Sigma^+$	1.44109	0.80436	0.76927	0.76925	0.80472	0.78827	0.79187	0.78999
	$\mathbf{NS}$	$X^2 \Pi_r$	1.16463	0.75742	0.75373	0.75364	0.75599	0.75702	0.75917	0.75672
	SO	$X^{3}\Sigma^{-}$	2.05772	2.01052	2.00679	2.00498	2.00905	2.00958	2.01162	2.00823
	$^{+}$ Os	$X^2_{ m D}\Pi_r$	1.12991	0.75587	0.75431	0.75259	0.75396	0.75408	0.75502	0.75297
	HCl+	$X^2 \Pi_i$	0.76225	0.75631	0.75597	0.75300	0.75488	0.75372	0.75332	0.75257

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

0.75018	0.75078	0.75403	2.01031	0.75515	0.75265	CCSD:	$\omega B97M-V$	0.75081	0.75537	2.00335	0.75208	0.75167	0.75429	2.00924	0.75388
0.75016	0.75079	0.75486	2.01188	0.75626	0.75380	CCSD:	$\omega B97X-V$	0.75093	0.75886	2.00423	0.75602	0.75237	0.75780	2.01545	0.75743
0.75038	0.75078	0.75492	2.01119	0.75531	0.75348	CCSD:	B97	0.75083	0.75487	2.00555	0.75517	0.75289	0.75545	2.00931	0.75456
0.75219	0.75221	0.75541	2.01177	0.75540	0.75388	CCSD:	B97M-rV	0.75059	0.75457	2.00891	0.75540	0.75508	0.75484	2.00775	0.75388
0.75008	0.75094	0.75354	2.00653	0.75287	0.75222	CCSD:	BLYP	0.75045	0.75350	2.00442	0.75402	0.75262	0.75338	2.00493	0.75286
0.75019	0.75150	0.75655	2.01178	0.75412	0.75461	CCSD:	00MP2	0.75093	0.75911	2.00565	0.75850	0.75663	0.75592	2.00891	0.75444
0.75016	0.75140	0.75771	2.01531	0.75662	0.75542	CCSD:	$\kappa$ -OOMP2	0.75084	0.76156	2.00949	0.76144	0.75477	0.76093	2.01829	0.75864
0.75014	0.75155	0.76583	2.03586	0.77006	0.76481	CCSD:	UHF	0.75092	0.77348	2.01899	0.77634	1.13370	1.02029	2.06113	1.20499
$X^2\Sigma^+$	$X^2\Sigma^+$	$X^{2}\Pi$	$X^{3}\Sigma^{-}$	$X^2 \Pi_i$	$X^{2}\Pi$	State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^{3}\Sigma_{q}^{-}$	$X^2\Pi_r^2$	$X^2\Pi_u$	$X^2 \Pi_r$	$X^3 \Sigma_g^-$	$X^2\Pi_{g,r}$
LiCl <sup>-</sup>	BeCl	CCI	NCI	OCI	$FC1^+$	Dimer		MgCl	AlS	${\rm Si}_2$	SiCI	$P_{2}^{+}$	$P_{S}$	$\mathbf{S}_2$	$s_2^+$
						Row $3-$	Row 3								

Appendix D. Additional Information: Polishing the gold standard: The role of orbita	l
choice in $CCSD(T)$ frequency prediction	282

## D.4 CCSD data

Table D	0.10: Exp	erimenta	d vibrational	l frequenci	es and errc	ors in the c	orrected v	ibrational	frequencie	s for the cl	losed-shell
species	are prese	inted in	for the CCS	D method	ls utilizing	different	molecular	orbitals.	Root mea	n square d	leviations,
mean si	igned dev	riations,	minimum d	eviations,	and maxi	mum devi	ations for	the set o	f species a	und subset	s are also
presente	∋d.										
Row 2-Row 2	Dimer	State	Expt.	$\Delta(CCSD: UHF)$	$\Delta$ (CCSD: $\kappa$ -OOMP2)	$\Delta(CCSD: OOMP2)$	$\Delta(CCSD: BLYP)$	$\Delta(CCSD: B97M-rV)$	$\Delta(CCSD: B97)$	$\Delta(\text{CCSD}: \omega \text{B97X-V})$	$\Delta$ (CCSD: $\omega$ B97M-V)
	LiH	$X^1\Sigma^+$	$1405.49805^{a}$	-4.95	-5.93	-5.25	-3.96	-6.91	-7.45	-4.46	-5.49
	$\mathrm{Li}_{2}$	$X^1 \Sigma_a^+$	$351.4066^{a}$	-4.86	-10.16	-3.10	-4.03	-1.75	-3.36	-26.61	-4.36
	LiF	$X^1\Sigma^+$	$910.57272^{a}$	4.40	6.56	7.19	7.72	6.80	6.80	6.57	6.45
	${ m BeH^+}$	$X^1\Sigma^+$	$2221.7^{ m b}$	-10.75	-10.66	-8.80	-10.75	-16.01	-8.7694	-13.26	-8.97
	BeO	$X^1\Sigma^+$	$1487.32^{b}$	75.94	88.76	93.78	91.43	89.12	90.2704	86.98	87.69
	BH	$X^1\Sigma^+$	$2366.7296^{a}$	12.83	8.70	12.76	15.81	8.90	14.66	11.03	10.41
	$_{ m BF}$	$X^1\Sigma^+$	$1402.15865^{a}$	9.33	15.33	16.66	18.28	16.10	16.28	15.51	15.68
	C,	$X^1 \Sigma_a^+$	$1855.0663^{a}$	96.46	85.08	78.07	-385.19	-366.32	-362.08	-340.93	-356.99
	CO	$X^1\Sigma^+$	$2169.75589^{a}$	55.68	70.60	73.72	72.33	68.78	68.93	67.27	67.34
	$ m N_2$	$X^1\Sigma_g^+$	$2358.57^{a}$	76.36	88.42	91.17	89.58	86.00	86.76	84.80	84.93
	+ON	$X^1\Sigma^+$	$2376.72^{a}$	103.61	123.66	127.13	123.61	119.47	118.90	117.38	117.44
	-HO	$X^1\Sigma^+$	$3735.2^{\circ}$	52.95	69.59	76.05	69.57	65.42	65.63	64.56	64.61
	HF	$X^1\Sigma^+$	$4138.385^{a}$	39.17	48.26	49.86	49.91	46.21	48.53	46.91	46.61
	$\mathrm{F}_2$	$X^1\Sigma_g^+$	$916.929^{ m d}$	170.29	104.70	107.50	104.74	103.50	103.32	102.90	102.84
	RMSD	1		69.95	66.52	68.26	120.82	115.37	114.54	109.40	112.80
	MSD			48.32	48.78	51.20	17.07	15.67	17.03	15.62	16.30
	MIN			-10.75	-10.66	-8.80	-385.19	-366.32	-362.08	-340.93	-356.99
	MAX			170.29	123.66	127.13	123.61	119.47	118.90	117.38	117.44
Row $2^-$	Dimer	State	Expt.	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(\text{CCSD};$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$
Row $3$			,	ÙHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B_{0}^{0} 7 M-V)$
	$_{\rm NaH}$	$X^1\Sigma^+$	$1171.968^{a}$	-8.37	-8.55	-8.45	-10.03	-8.69	-9.60	-9.63	-9.82
	NaLi	$X^1\Sigma^+$	$256.5412^{a}$	2.01	2.11	-2.86	-0.47	-0.59	-15.03	-3.63	-2.49
	NaF	$X^1\Sigma^+$	$535.65805^{a}$	0.17	1.57	1.93	2.13	1.72	1.52	1.57	1.63
	${ m MgH^+}$	$X^1\Sigma^+$	$1699.1^{\mathrm{b}}$	-0.10	-0.17	1.38	-1.86	-3.39	0.23	-1.82	-3.41
	AIH	$X^1\Sigma^+$	$1682.37474^{a}$	-7.02	-2.63	-6.91	-3.94	-8.11	-1.94	-2.19	-4.07
	AIF	$X^1\Sigma^+$	$802.32447^{a}$	2.01	4.73	5.55	6.33	5.17	5.24	4.87	4.99
	$SiH^+$	$X^1\Sigma^+$	$2157.17^{a}$	8.25	12.28	12.44	14.13	14.30	13.86	8.73	12.63
	SiO	$X^1\Sigma^+$	$1241.54388^{a}$	43.96	53.88	59.60	56.14	54.35	53.99	51.85	51.92
	PN	$X^1\Sigma^+$	$1336.948^{a}$	73.08	72.09	79.35	73.46	71.43	71.25	69.02	69.04
	$\operatorname{BeS}$	$X^1\Sigma^+$	$997.94^{a}$	-234.21	27.17	30.29	30.00	28.94	28.73	26.82	26.97
	CS	$X^1\Sigma^+$	$1285.08^{ m b}$	51.17	62.22	68.15	65.08	62.36	62.60	59.60	60.19
	$^{+SX}$	$X^1\Sigma^+$	$1415^{\rm b}$	100.80	102.09	109.97	103.13	100.22	100.32	97.92	98.31
	HCI	$X^1\Sigma^+$	$2990.9248^{a}$	43.02	44.49	43.89	44.69	46.48	44.55	44.46	46.52
	LiCl	$X^1\Sigma^+$	$642.95453^{a}$	-3.92	-3.82	-3.91	-3.63	-3.61	-3.68	-4.01	-3.62
	BCI	$X^1\Sigma^+$	$840.29472^{a}$	5.09	7.21	8.77	10.30	9.28	8.85	7.72	7.96
	CCI+	$X^1\Sigma^+$	$1175^{b}$	44.47	58.67	66.11	63.50	60.07	60.05	56.25	57.15
	CIF	$X^1\Sigma^+$	$783.4534^{a}$	28.96	32.27	33.38	32.36	31.19	31.30	31.19	31.16
	RMSD			68.52	41.95	45.59	43.30	42.05	42.01	40.38	40.78
	MSD			8.79	27.39	29.34	28.31	27.13	26.60	25.81	26.18
	MIN			-234.21	-8.55	-8.45	-10.03	-8.69	-15.03	-9.63	-9.82

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

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98.31	$\Delta(\text{CCSD};$	$\omega B97M-V$	-3.76	-4.30	24.83	38.96	8.46	21.16	12.84	-4.30	38.96	$\Delta(CCSD:$	$\omega B97M-V)$	76.13	20.48	-356.99	117.44
97.92	$\Delta(\text{CCSD};$	$\omega B97X-V$	-3.77	-4.33	24.47	38.80	8.46	21.02	12.73	-4.33	38.80	$\Delta(CCSD:$	$\omega B97X-V)$	74.07	20.03	-340.93	117.38
100.32	$\Delta(CCSD:$	B97)	-3.67	-3.98	26.39	40.57	8.56	22.11	13.57	-3.98	40.57	$\Delta(CCSD:$	B97)	77.48	21.07	-362.08	118.90
100.22	$\Delta(\text{CCSD};$	B97M-rV)	-3.70	-4.00	26.51	40.53	8.74	22.15	13.62	-4.00	40.53	$\Delta(CCSD:$	B97M-rV)	77.97	20.79	-366.32	119.47
103.13	$\Delta(\text{CCSD};$	BLYP)	-3.55	-3.65	27.66	41.59	9.01	22.81	14.21	-3.65	41.59	$\Delta(CCSD:$	BLYP)	81.45	21.98	-385.19	123.61
109.97	$\Delta(\text{CCSD};$	OUMP2)	-3.69	-4.13	28.00	42.93	8.85	23.39	14.39	-4.13	42.93	$\Delta(CCSD:$	OOMP2)	53.57	35.76	-8.80	127.13
102.09	$\Delta(\text{CCSD:}$	$\kappa$ -OOMP2)	-3.88	-4.48	24.75	39.47	8.43	21.34	12.86	-4.48	39.47	$\Delta(CCSD:$	$\kappa$ -OOMP2)	51.14	33.69	-10.66	123.66
100.80	$\Delta(\text{CCSD};$	UHF)	-3.81	-4.67	21.34	45.20	7.82	22.79	13.18	-4.67	45.20	$\Delta(CCSD:$	ÚHF)	64.75	24.77	-234.21	170.29
	Expt.		$364.6842^{a}$	$481.77466^{a}$	$749.64559^{a}$	$780.77^{a}$	$559.751^{a}$					Expt.					
	State		$X^{1}\Sigma^{+}$	$X^1\Sigma^+$	$X^1\Sigma^+$	$X^1 \Sigma_g^+$	$X^1\Sigma_g^+$					State					
MAX	Dimer		NaCl	AlCI	$\operatorname{SiS}$	$\mathbf{P}_2$	$\mathrm{Cl}_2$	RMSD	MSD	MIN	MAX	Dimer		RMSD	MSD	MIN	MAX
	Row 3–	Row 3										Closed-	$_{\rm shell}$				

<sup>a</sup> From Ref. 349. <sup>b</sup> From Ref. 346. <sup>c</sup> From Ref. 350.

Appendix D. Additional Information: Polishing the gold standard: The role of orbita	al
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Table D	.11: Exp	erimental	l vibrational	frequenci	ies and erro	ors in the o	corrected v	ribrational	frequenci	es for the e	open-shell
species .	are prese	ented in f	or the CCS	D method	ls utilizing	different 1	molecular	orbitals.	Root mea	n square d	eviations,
mean si	gned dev.	iations, m	iinimum dev	riations, a	nd maximu	um deviatio	ons for the	set of spe	cies and su	ubsets are ]	presented.
Row 2– Row 2	Dimer	State	Expt.	$\Delta(CCSD: UHF)$	$\Delta$ (CCSD: $\kappa$ -OOMP2)	$\Delta(CCSD: OOMP2)$	$\Delta(CCSD: BLYP)$	$\Delta(CCSD: B97M-rV)$	$\Delta(CCSD: B97)$	$\Delta(CCSD: \omega B97X-V)$	$\Delta(CCSD: \omega B97M-V)$
	LiO	$X^2 \Pi_i$	$814.62^{a}$	3.94	5.61	6.08	7.31	6.48	6.51	5.95	$5.9\hat{1}$
	$\operatorname{BeH}$	$X^2\Sigma^+$	$2061.235^{a}$	1.68	1.23	4.44	2.76	0.82	1.41	2.78	1.21
	$\mathrm{BeF}$	$X^2\Sigma^+$	$1247.36^{\mathrm{b}}$	22.61	25.70	26.43	27.35	26.14	26.17	25.81	26.03
	$\mathrm{B}_2$	$X^3 \Sigma_q^-$	$1051.3^{ m b}$	44.59	40.38	3.78	35.49	33.05	31.27	32.03	32.73
	BN	$X^3\Pi^{ m c}$	$1514.6^{\mathrm{b}}$	64.98	68.38	71.04	68.20	66.68	65.12	64.42	66.43
	BO	$X^2\Sigma^+$	$1885.286^{a}$	43.31	55.40	59.77	58.97	54.94	55.58	54.35	54.42
	CH	$X^2 \Pi_r$	$2860.7508^{a}$	14.37	22.42	20.91	23.07	22.30	20.99	20.73	22.84
	C, -	$X^2\Sigma_a^+$	$1781.189^{a}$	-65.13	57.66	64.77	58.17	55.07	54.40	6.09	25.65
	oN^ C	$X^2\Sigma^{J}$	$2068.648^{a}$	103.60	78.37	81.29	76.63	75.27	74.70	73.98	75.70
	$CO^+$	$X^2\Sigma^+$	$2214.127^{a}$	91.85	112.67	119.04	113.90	105.88	107.95	106.78	107.09
	CF	$X^2 \Pi_r$	$1307.93^{a}$	28.27	39.55	42.36	43.55	40.34	40.11	38.60	38.82
	HN	$X^3\Sigma^-$	$3282.72^{a}$	28.52	34.79	32.86	36.15	33.80	34.03	30.57	32.40
	$^+$	$X^2 \Sigma_a^+$	$2207.0115^{a}$	116.78	101.46	112.23	97.51	92.54	92.63	90.12	90.71
	NON	$X^{2}\Pi_{r}^{3}$	$1904.1346^{a}$	89.82	104.51	108.95	105.22	100.93	101.24	99.29	99.51
	ЪF	$X^3\Sigma^-$	$1141.37^{a}$	28.52	42.43	47.39	47.22	41.43	41.84	39.87	39.49
	HO	$X^2\Pi_i$	$3737.761^{a}$	30.46	36.53	36.58	38.33	35.95	36.01	36.65	35.49
	OH <sup>+</sup>	$X^3\Sigma^-$	$3113.37^{\rm b}$	33.25	35.47	35.67	38.43	39.19	36.47	34.66	36.60
		$X^{3}\Sigma^{-}$	1580 161a	81 40	00.60	10.4 70	00 11	05.60	05.12	03.78	03.88
	2°4	$v_{2\Pi}^{2}$	TOOL DOOR	01-10	150.00	10 0 T	TT-66	20.05	407.00 77.00	01.05	00.06
	0 <sup>2</sup>	$X^{2}\Pi_{g}$	1905.892 <sup>a</sup>	133.11	150.051	106.05	148.73	144.25	144.04	142.11	142.24
	0 0 0	$X^{2}\Pi_{g,i}$	1090 <sup>n</sup>	102.15	119.79	127.46	118.74	116.16	115.49	114.22	114.17
	OF.	$X^{2}$ III	$1053.0138^{a}$	50.04	75.93	93.00	79.55	68.15	68.72	65.06	64.20
	HF+	$X^2 \Pi_i$	$3090.5^{\rm b}$	57.35	59.59	59.35	66.31	60.95	62.06	61.75	61.82
	$F_{2}^{+}$	$X^2 \Pi_{g,i}$	$1091.5^{d}$	153.52	152.75	161.40	150.43	148.16	147.68	146.35	145.97
	$\mathbf{F}_{2}^{-}$	$X^2\Sigma_u^+$	$510^{\mathrm{b}}$	-24.34	-49.85	-47.22	-49.21	-50.23	-50.31	-50.52	-50.67
	RMSD			71.63	77.25	81.44	77.39	74.18	74.08	72.10	72.54
	MSD			51.44	61.29	63.71	62.16	58.92	58.72	55.64	56.78
	MIN			-65.13	-49.85	-47.22	-49.21	-50.23	-50.31	-50.52	-50.67
	MAX			153.52	152.75	161.40	150.43	148.16	147.68	146.35	145.97
Row $2^-$	Dimer	State	Expt.	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$	$\Delta(CCSD:$
Row $3$				UHF)	$\kappa$ -OOMP2)	OOMP2)	BLYP)	B97M-rV)	B97)	$\omega B97X-V)$	$\omega B97M-V)$
	NaO	$X^{2}\Pi$	$526^{\rm b}$	-31.86	-30.77	-29.95	-30.34	-30.32	-30.43	-30.74	-30.43
	MgH	$X^2\Sigma^+$	$1492.7763^{a}$	5.71	5.65	2.85	6.38	5.07	4.69	6.68	4.41
	MgF	$X^2\Sigma^+$	$711.69^{b}$	8.65	10.70	11.32	11.61	11.24	11.14	10.76	10.96
	$AIH^+$	$X^2\Sigma^+$	$1620^{\mathrm{b}}$	35.64	39.98	49.31	53.80	55.42	53.32	56.24	56.34
	SiH	$X^2 \Pi_r$	$2042.5229^{a}$	3.70	10.43	9.01	8.12	9.23	8.72	5.85	8.80
	$\operatorname{SiF}$	$X^2 \Pi_r$	$837.32507^{a}$	25.35	29.11	30.25	31.16	29.92	29.82	29.22	29.28
	Ηd	$X^3\Sigma^-$	$2363.774^{a}$	22.73	20.77	24.52	28.98	25.38	25.99	26.01	25.42
	$^{\rm +Hd}$	$X^2 \Pi_r$	$2299.6^{\rm b}$	111.17	106.73	108.42	110.46	110.30	111.31	109.79	111.25
	$^{-}\mathrm{Hd}$	$X^2 \Pi_i$	$2230^{\mathrm{b}}$	51.52	53.38	52.00	57.70	55.47	55.64	51.35	55.97
	СР	$X^2\Sigma^+$	$1239.79924^{a}$	46.65	49.54	65.23	56.05	50.55	50.81	46.51	49.01
	РО	$X^2 \Pi_r$	$1233.34^{a}$	56.82	66.22	71.61	68.09	66.12	65.76	64.09	64.08

Appendix D. Additional Information: Polishing the gold standard: The role of orbital choice in CCSD(T) frequency prediction

81.48	15.46	30.60	34.04	35.84	58.66	63.65	56.34	91.58	50.43	29.65	2.63	9.10	100.43	1.87	16.02	51.28	39.36	-30.43	111.25	$\Delta(CCSD:$	$\omega B97M-V)$	5.23	-12.27	21.99	-1.57	49.28	32.74	24.48	61.02	32.57	22.61	-12.27	61.02	$\Delta(CCSD:$	$\omega B97M-V)$	59.06	44.18	-50.67	145.97
81.41	15.49	30.52	33.01	34.95	56.66	62.92	56.15	91.67	50.09	29.86	2.48	8.66	100.22	1.93	15.72	50.70	38.80	-30.74	109.79	$\Delta(CCSD:$	$\omega B97X-V)$	5.17	-12.30	21.82	-1.70	49.06	32.45	24.32	60.82	32.42	22.46	-12.30	60.82	$\Delta(CCSD:$	$\omega B97X-V)$	58.60	43.43	-50.52	146.35
82.61	16.25	31.73	35.14	37.24	60.42	65.95	57.53	93.44	47.95	29.54	2.83	10.78	102.89	4.14	20.11	52.00	40.20	-30.43	111.31	$\Delta(CCSD:$	B97)	5.35	-8.69	22.99	-1.13	50.83	34.18	25.24	62.37	33.38	23.89	-8.69	62.37	$\Delta(CCSD:$	B97)	60.17	45.52	-50.31	147.68
82.98	16.23	32.02	34.95	37.17	59.11	66.09	57.70	93.69	50.09	29.52	2.99	10.97	102.63	4.12	19.09	52.03	40.29	-30.32	110.30	$\Delta(CCSD:$	B97M-rV)	5.30	-6.96	22.94	-1.07	50.71	34.10	25.43	62.42	33.32	24.11	-6.96	62.42	$\Delta(CCSD:$	B97M-rV)	60.23	45.67	-50.23	148.16
84.47	17.82	34.16	37.34	39.23	65.99	69.90	59.59	95.92	50.83	30.19	3.28	12.79	106.86	7.27	25.96	53.93	42.36	-30.34	110.46	$\Delta(CCSD:$	BLYP)	5.42	-7.46	23.52	-0.52	51.80	35.72	26.26	63.46	34.12	24.78	-7.46	63.46	$\Delta(CCSD:$	BLYP)	62.65	48.03	-49.21	150.43
87.48	16.87	33.88	35.10	39.48	72.80	78.28	64.42	102.29	50.61	29.55	2.76	11.95	109.34	5.88	35.08	55.65	43.35	-29.95	109.34	$\Delta(CCSD:$	OOMP2)	5.31	-2.64	55.58	-1.29	106.91	36.60	26.94	65.35	51.10	36.59	-2.64	106.91	$\Delta(CCSD:$	OOMP2)	66.86	50.71	-47.22	161.40
82.86	15.47	31.21	35.89	35.92	60.63	68.59	59.71	95.50	48.92	29.89	2.41	8.46	101.77	1.02	19.08	51.27	39.23	-30.77	106.73	$\Delta(CCSD:$	$\kappa$ -OOMP2)	5.14	-11.12	21.66	-1.99	49.27	32.27	24.59	61.62	32.59	22.68	-11.12	61.62	$\Delta(CCSD:$	$\kappa$ -OOMP2)	61.44	45.96	-49.85	152.75
74.82	11.24	22.32	33.10	31.87	50.49	51.77	50.36	87.10	51.07	29.43	1.99	5.41	82.04	0.17	8.61	46.34	34.37	-31.86	111.17	$\Delta(CCSD:$	UHF)	5.04	-19.95	20.68	-2.40	49.78	40.49	22.26	60.19	33.71	22.01	-19.95	60.19	$\Delta(CCSD:$	UHF)	56.78	39.64	-65.13	153.52
$1000^{\mathrm{b}}$	$846.75^{a}$	$1053.25^{ m b}$	$2696.2475^{a}$	$1179.91^{a}$	$1384^{ m b}$	$1218.7^{ m b}$	$1150.7913^{a}$	$1306.778^{e}$	$2673.69^{a}$	$480^{ m b}$	$846.7^{ m b}$	$876.89749^{a}$	$870^{\mathrm{p}}$	$827.95767^{a}$	$853.64268^{a}$					Expt.		$462.12^{\rm b}$	$617.1169^{a}$	$510.98^{a}$	$535.59^{a}$	$672.2^{a}$	$739.1^{ m b}$	$725.7102^{a}$	$_{ m q}062$					Expt.					
$X^3\Sigma^-$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2 \Pi_r$	$X^3\Sigma^-$	$X^2 \Pi_r$	$X^2 \Pi_i$	$X^2\Sigma^+$	$X^2\Sigma^+$	$X^2\Pi$	$X^2\Pi$	$X^3\Sigma^-$	$X^2 \Pi_i$					State		$X^2\Sigma^+$	$X^2\Sigma^+$	$X^3\Sigma_q^-$	$X^2\Pi_r^{}$	$X^2 \Pi_u$	$X^2 \Pi_r$	$X^3 \Sigma_g^-$	$X^2\Pi_{g,r}$					State					
$^{-}\mathrm{Od}$	ЪF	$PF^+$	HS	$_{\mathrm{BS}}$	$CS^+$	NS	SO	$^{\rm SO^+}$	$HCl^+$	$LiCl^{-}$	BeCl	CCI	$CIF^+$	NCI	OCI	RMSD	MSD	MIN	MAX	Dimer		MgCl	AIS	$\mathrm{Si}_2$	$si\bar{c}l$	Р <sub>2</sub> +	$\tilde{PS}$	$\mathbf{S}_2$	$s_{+}^{2}$	RMSD	MSD	MIN	MAX	Dimer		RMSD	MSD	MIN	MAX
																				Row 3–	Row 3													Open-	$_{\rm shell}$				

<sup>a</sup> From Ref. 349. <sup>b</sup> From Ref. 346. <sup>d</sup> From Ref. 351. <sup>e</sup> From Ref. 352. \* Theoretical results.