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### Title

Four-Dimensional Paleomagnetic Dataset: Plio-Pleistocene Paleodirection and Paleointensity Results From the Erebus Volcanic Province, Antarctica

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### Authors

Asefaw, H  
Tauxe, L  
Koppers, AAP  
[et al.](#)

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1 **Four-Dimensional Paleomagnetic Dataset: Plio-Pleistocene**  
2 **Paleodirection and Paleointensity Data from the Erebus**  
3 **Volcanic Province, Antarctica**

4 **H. A. Asefaw,<sup>1</sup>L. Tauxe,<sup>1</sup>A.A.P. Koppers,<sup>2</sup>H. Staudigel,<sup>1</sup>,**

5 <sup>1</sup>Geosciences Research Division, Scripps Institution of Oceanography, University of California San Diego, La Jolla,  
6 California, USA

7 <sup>2</sup>College of Earth, Ocean and Atmospheric Sciences, Oregon State University, Corvallis, Oregon, USA,

8 **Key Points:**

- 9 • Eleven new <sup>40</sup>Ar/<sup>39</sup>Ar age determinations from the Erebus Volcanic Province, Antarc-  
10 tica (-77.84°, 166.69°)
- 11 • One hundred and twenty-six site mean directions recover a paleopole (176.24°, 86.89°)  
12 consistent with a GAD field over the Plio-Pleistocene
- 13 • Twenty-eight site intensity estimates pass a set of strict selection criteria and recover a  
14 35.75  $\mu T \pm 7.30 \mu T$  time averaged field over the Plio-Pleistocene

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Corresponding author: Hanna Asefaw, hasafaw@ucsd.edu

## Abstract

The primary structure of the modern geomagnetic field may be accounted for by a Geocentric Axial Dipole (GAD) field. A GAD field is a magnetic field produced by a dipole positioned in the center of the Earth and aligned with the spin axis. Paleogeographic reconstructions of plate motion are based on the assumption that this predominately GAD structure extends to the paleomagnetic field structure so it is crucial to determine whether globally distributed paleodirectional and paleointensity datasets recover a GAD field. Global paleodirectional compilations that span 0 - 5 Myr support a field structure dominated by a GAD with minor non-GAD contributions. However, paleointensity datasets over the same period lack the global intensity structure expected of a GAD derived field. A notable deviation is the depressed intensities observed at the high latitudes which should preserve the highest average intensity in a purely GAD field. To determine whether the low intensities reflect the structure of the field, low quality data or inadequate temporal sampling, we have conducted a robust study of the paleomagnetic field at the high southerly latitudes. This study focuses on the paleomagnetic field structure over the Plio-Pleistocene to avoid corrections for plate motion. We present the results from one hundred and twenty-six site mean directions that were thermally or AF demagnetized and then subjected to a set of strict selection criteria. Along with twenty-eight new paleointensity estimates from samples that underwent the IZZI modified Thellier-Thellier experiment and were subjected to the strict CCRIT set of criteria. The recovered paleopole (176.24 °, 86.89 °) and its corresponding  $\alpha_{95}$  (4.92 °), supports the GAD hypothesis. Our time averaged field estimate,  $35.75\mu T \pm 7.30\mu T$ , is consistent with the low intensities measured at the poles in the global compilations.

## 1 Introduction

The spatial structure of modern geomagnetic field intensity (Figure) reveals latitudinal variability, longitudinal features, and regions with anomalously low intensities (the South Atlantic Anomaly). However, in mathematical representations of the geomagnetic field structure [Thébaud *et al.*, 2015] the geocentric axial dipole (GAD) accounts for much of the field [McElhinny, 2007; Lowes, 1973]. A geocentric axial dipole field is the magnetic field generated by a dipole positioned in the center of the Earth and aligned along the spin axis. In an ideal GAD field, both the intensity of the geomagnetic field ( $B$ ) and the inclination ( $I$ ) would vary with latitude ( $\lambda$ ) by

$$B = M(1 + 3\cos^2\theta)^{\frac{1}{2}} \quad (1)$$

$$\tan(I) = 2\tan(\lambda) \quad (2)$$

where  $M$  is intensity of the field at the equator (nT) and  $\theta$  is co-latitude (°). The paleomagnetic field structure is preserved in the geological record and various techniques [Thellier and Thellier, 1959; Shaw, 1974; Coe, 1967; Yu *et al.*, 2004; Walton and Shaw, 1922; Hoffman and Biggin, 2005] allow us to recover paleodirections and paleointensities. Independent studies of the paleofield are then compiled in paleodirectional and paleointensity databases [Cromwell *et al.*, 2018; Brown *et al.*, 2015; Biggin *et al.*, 2009] which we can use to characterize the behavior of the time averaged field (TAF). The structure of the geomagnetic field reflects the dynamics and motion occurring in the fluid outer core so it is important to characterize the TAF. The GAD hypothesis is also central to paleogeographic reconstructions of plate motion.

Numerous studies [Opdyke and Henry, 1969; Cromwell *et al.*, 2018; Behar *et al.*, 2019] have recovered paleodirections from the Plio-Pleistocene that are consistent with a GAD field. However, a GAD structure does not emerge in the PINT (absolute paleointensity) database over the same time period [Biggin *et al.*, 2009]. The latitudinal variation of intensity, expected of a GAD field, is not evident in the PINT database for estimates that span the last 5 Myr. The field intensity at the high latitudes appears depressed, which may

64 either reflect a feature of the paleomagnetic field or the quality of the data underlying the  
 65 database. Recovering paleointensity is challenging due to the complex magnetization ac-  
 66 quisition behavior of non-ideal magnetic grains [Dunlop *et al.*, 2004; Dunlop and Özdemir,  
 67 2001] and the tendency for magnetomineralogical alteration during paleointensity experi-  
 68 ments [Smirnov and Tarduno, 2003]. To determine whether the low intensities measured at  
 69 the high southerly latitudes accurately represent the structure of the paleomagnetic field or if  
 70 they are an artifact of non-ideal magnetic recorders, we conducted an extensive study of the  
 71 paleomagnetic field in the Erebus Volcanic Province, Antarctica ( $-77.84^\circ$ ,  $166.69^\circ$ ).

## 72 2 Methods

### 73 2.1 Sample Collection

74 Our study examines 141 independent sites around the Erebus Volcanic Province, Antarc-  
 75 tica (Figure). Samples were collected during the 2016/2017 Antarctic Summer field season.  
 76 Site selection was based on the work of Mankinen and Cox [1988], Tauxe *et al.* [2004], and  
 77 Lawrence *et al.* [2009] who collected samples from the interior of lava flows primarily for  
 78 directional analysis. A compilation of all the paleodirectional and paleointensity experiments  
 79 were summarized in Lawrence *et al.* [2009]. Only a dozen of those sites yield paleointen-  
 80 sity data that pass modern, strict selection criteria (i.e. CCRIT of Cromwell *et al.* [2015]),  
 81 therefore, we re-sampled nearly all of the original sites for this study. Both Mankinen and  
 82 Cox [1988] and Lawrence *et al.* [2009] collected between four and seven cores with a gas-  
 83 powered drill. We used the 1-inch drill holes remaining in the outcrop to identify many of  
 84 the original sites (Figure ). The remainder were located by GPS coordinates from Lawrence  
 85 *et al.* [2009] and approximated from the maps and descriptions in Mankinen and Cox [1988].  
 86 Once we identified the sites, we re-sampled the finest-grained, glassy material from the lava  
 87 flow top or flow bottom. We collected hand samples using hammers and chisels. The out-  
 88 crops included lava flows, pillow lavas, and hyaloclastite cones that formed over the Plio-  
 89 pleistocene (Figure).

### 90 2.2 Paleointensity

#### 91 2.2.1 Recovering paleointensity

92 Magnetic grains in igneous rocks acquire a Thermal Remanent Magnetization (TRM)  
 93 by cooling from temperatures well above their curie temperature through their blocking tem-  
 94 peratures ( $T_b$ ). The resulting TRM captures an instantaneous record of the geomagnetic field  
 95 that remains stable over long timescales. The degree of alignment, between the magnetic  
 96 grains and the ambient field, depends on the strength of the field ( $B$ ) at the time of cooling  
 97 [Néel, 1955]

$$98 M_{TRM} = M_{rs} \tanh \frac{\nu M_s(T_b) B}{k T_b} \quad (3)$$

99 where  $T_b$  is blocking temperature,  $k$  is the Boltzmann constant,  $\nu$  is volume, and  $M_s(T_b)$   
 100 is spontaneous magnetization at  $T_b$ . In a weak magnetic field, on the order of the modern ge-  
 101 omagnetic field, TRM acquisition is linearly proportional to the strength of the ambient field.  
 102 This proportionality allows us to recover the intensity of the geomagnetic field when the rock  
 103 formed. The Natural Remanent Magnetization (NRM) may be removed by heating the rock  
 104 and a new partial Thermal Remanent Magnetization (pTRM) may overwrite the NRM by  
 105 cooling the rock in a controlled applied field. The ratio of the TRM acquired to the field ap-  
 106 plied is proportional to the ratio of the NRM to the paleomagnetic field [Néel, 1955]. We  
 107 then estimate the intensity of the paleomagnetic field by

$$108 B_{anc} = \frac{M_{NRM}}{M_{pTRM}} B_{lab}, \quad (4)$$

109 where  $M_{NRM}$  is the natural remanent magnetization,  $M_{pTRM}$  is the partial thermal re-  
 110 manent magnetization imparted by heating the sample in an applied field,  $B_{lab}$  is the field

111 applied in the lab, and  $B_{anc}$  is the strength of the paleomagnetic field. A rock contains an  
 112 assemblage of magnetic grains and each grain traps its magnetization at a different tempera-  
 113 ture, therefore incrementally demagnetizing and remagnetizing a rock sample results in sev-  
 114 eral independent estimates of the paleofield.

### 115 **2.2.2 Sample preparation**

116 We conducted a preliminary IZZI-modified Thellier Thellier experiment on 144 spec-  
 117 imens with, at minimum, two specimens from each site. The results from this preliminary  
 118 experiment refined our sites to the most promising from which we selected up to six addi-  
 119 tional specimens. Samples were crushed into 100 – 500 mg fragments. The fragments were  
 120 then examined under a binocular microscope to select the individual specimen that appeared  
 121 glassy. These glassy specimens may contain the uniaxial single domain grains of magnetite  
 122 needed to recover paleointensity. Each individual specimen was swaddled in glass microfiber  
 123 filter paper and affixed inside a borosilicate glass vial with  $K_2SiO_3$ . The specimen were then  
 124 placed in a transformer steel shielded room at the Paleomagnetic Laboratory at Scripps Insti-  
 125 tution of Oceanography for the duration of the experiment.

### 126 **2.2.3 IZZI modified Thellier-Thellier Experiment**

127 We conducted the IZZI-modified Thellier-Thellier protocol [Yu *et al.*, 2004], whereby  
 128 specimens are incrementally heated and cooled either in the absence of a magnetic field to  
 129 demagnetize the NRM (a zero-field step) or in the presence of an applied lab field to impart  
 130 a pTRM (an in-field step). Specimen were subjected to both an In-field (I) and Zero-field  
 131 (Z) treatments at each temperature step. Temperature steps were conducted at 100°C inter-  
 132 vals from 0°C to 400°C, then 25°C intervals to 500°C, and finally at 10°C intervals until  
 133 each specimen completely demagnetized. Specimens were heated in custom-built furnaces  
 134 with thermocouples in non-inductively wound heating elements to control the temperature  
 135 to within a few degrees. Specimen were rapidly air-cooled following treatment. During in-  
 136 field treatment steps, specimen were cooled in a 30  $\mu$ T field. The order of the treatment, IZ  
 137 or ZI [Coe, 1967], alternated with each temperature step in order to detect tails and zero-field  
 138 memory effects [Aitken *et al.*, 1988] in the ZI sequence. We applied a partial Thermal Re-  
 139 manent (pTRM) check, an additional in-field treatment at a previously measured temperature  
 140 step, between IZ-ZI sequences in order to monitor mineral neof ormation and magnetomin-  
 141 eral alteration. Immediately following treatment, we measured the magnetic remanence with  
 142 a 2G Cryogenic SQUID (Super Conducting Quantum Interference Device) magnetometer.

### 143 **2.2.4 Cooling Rate**

144 The TRM acquired by each specimen is affected by its rate of cooling [Dodson and  
 145 McClelland-Brown, 1980; Halgedahl and Fuller, 1980; Fox and Aitken, 1980]. After each  
 146 treatment, specimens were rapidly air-cooled to match the rate at which we suspect they ini-  
 147 tially cooled. To assess the impact of cooling rate on TRM acquisition in our specimens,  
 148 we conducted a cooling rate experiment. We heated the specimens to 620° in a 50  $\mu$ T field,  
 149 air-cooled them in under an hour, and then measured their TRM. We then re-heated the spec-  
 150 imens to 620° in a 50  $\mu$ T field, naturally cooled them over 12 hours, and then measured the  
 151 resulting TRM.

### 152 **2.2.5 Non-linear TRM Acquisition**

153 Néel theory is based on SD non-interacting grains of magnetite that acquire a TRM in  
 154 proportion to the ambient field, yet several studies have detected non-linear TRM acquisition  
 155 [Selkin *et al.*, 2007; Dunlop *et al.*, 2004]. Therefore after we completed the IZZI-experiment,  
 156 we selected the successful specimens to perform an additional set of steps that detect non-  
 157 linear TRM acquisition behavior. We subjected these specimens to a total TRM, at 630° C,

158 in several treatment fields. We treated the specimens in a  $0 \mu T$  ,  $15 \mu T$  ,  $20 \mu T$  ,  $30 \mu T$  ,  $40$   
 159  $\mu T$  ,  $50 \mu T$  , and  $60 \mu T$  field.

## 160 2.3 Paleodirection

### 161 2.3.1 Alternating Field Demagnetization and Thermal Demagnetization

162 We recovered paleodirection by stepwise thermal demagnetization and Alternating  
 163 Field (AF) demagnetization. Each oriented drill core was cut into one-inch specimens for AF  
 164 demagnetization or thermal demagnetization. At least five specimens per site were stepwise  
 165 demagnetized. 461 specimens were AF demagnetized in a Sapphire Instruments SI-4 uni-  
 166 axial AF demagnetizer. Specimens were treated in 5 mT steps from 5 mT – 20 mT, 10 mT  
 167 steps from 20 mT – 100 mT, and then at 120 mT, 150 mT, and 180 mT or until the NRM was  
 168 removed. An additional 323 specimens were thermally demagnetized by stepwise heating in  
 169  $50^\circ\text{C}$  intervals from  $0^\circ\text{C}$  –  $500^\circ\text{C}$ , in  $25^\circ\text{C}$  intervals from  $520^\circ\text{C}$  to  $560^\circ\text{C}$  and in  $5^\circ\text{C}$ - $10^\circ\text{C}$   
 170 intervals until the samples were entirely demagnetized. After each treatment, the remaining  
 171 NRM was measured. The demagnetization path, calculated from the resultant vector of the  
 172 NRM between treatment steps, monitors the stability and behavior of the magnetization.

## 173 2.4 Hysteresis and FORCs

174 We conducted paleointensity experiments on samples that were drilled from the inter-  
 175 rior of the lava flows [*Mankinen and Cox*, 1988; *Tauxe and Staudigel*, 2004] and samples that  
 176 were hand collected from the surface or base of the lava flow. At six sites, we recovered in-  
 177 tensity estimates from specimen collected from both the interior and the surface. We selected  
 178 sister specimens from these sites and measured hysteresis loops and FORC diagrams with  
 179 a Princeton Measurements Corporation Micromag Alternating Gradient Magnetometer to  
 180 diagnose domain state.

## 181 2.5 Ar-Ar

## 182 3 Results

### 183 3.1 Paleointensity

184 We present the results of our IZZI experiment as Arai diagrams [*Nagata and Arai*,  
 185 1963], in order to compare the pTRM acquired and the NRM removed at each temperature  
 186 step and to monitor any changes in this ratio between different temperature intervals. We  
 187 present the magnetization directions as zijderveld diagrams [*Zijderveld*, 1967] and calculate  
 188 the best fitting direction, or plane, through the vectors using Principal Component Analy-  
 189 sis [*Kirschvink*, 1980]. Our specimens rarely behave like the non-interacting uniaxial single  
 190 domain grains of magnetite assumed by Neel theory. Instead, many specimens exhibit non-  
 191 ideal behavior (i.e. zig-zagging, failed pTRM checks, or multiple components of magnetiza-  
 192 tion) resulting in unreliable paleointensity estimates.

#### 193 3.1.1 Non-ideal behavior: Zig-zagging

194 Zig-zagging in the Arai diagram, (Figure a). occurs when the ratio of NRM removed  
 195 to pTRM acquired varies between different temperature intervals based on the sequence of  
 196 treatment steps (IZ or ZI). The IZZI modified Thellier-Thellier experiment alternates the or-  
 197 der in which the treatments are applied for each temperature step [*Yu et al.*, 2004]. The al-  
 198 ternating sequence, In-field then Zero-field or Zero-field then In-field, is used to detect the  
 199 presence of tails and zero-field memory effects [*Aitken et al.*, 1988]. Tails occur when the  
 200 pTRM acquired by heating to temperature T in a field is not entirely removed when the spec-  
 201 imen is reheated to temperature T and cooled in a zero-field. This may indicate the presence  
 202 of MD grains.

### 3.1.2 Non-ideal behavior: Failed pTRM checks

A pTRM check, where a previously measured in-field treatment is repeated, is inserted into every IZ-ZI sequence [Shaw, 1974]. Any deviation in the remanence ( $\beta$ ) indicates magneto-mineral alteration or changes in the blocking and unblocking temperature spectra due to multidomain grains

## 3.2 Ideal behavior and Selection Criteria

To filter out the specimen that exhibit non-ideal behavior, (??) we apply a set of selection criteria at the specimen and site level. A wide range of selection criteria [Leonhardt *et al.*, 2004; Kissel and Laj, 2004] and paleointensity statistics [Paterson *et al.*, 2014] exist to separate low and high quality paleointensity data. We modeled our criteria (Table) after those of Cromwell *et al.* [2015], where they successfully recovered the paleointensity of the 1960 Hawaiian lava flow. CCRIT applies two directional statistics, Deviation ANGLE (DANG [Tanaka and Kobayashi, 2003]) and Maximum Angle of Deviation (MAD [Kirschvink, 1980]) to determine the variability in the direction of the NRM. MAD (maximum angle of deviation) quantifies the amount of scatter in the directions while DANG (deviation angle) calculates the angle between the center of the demagnetization direction and the origin. Three additional parameters, SCAT, Frac [Shaar and Tauxe, 2013], and  $k$  [Paterson, 2011], are applied to test the assumption of linearity. SCAT constrains the amount of scatter permitted between the best fit proportionality constant and the demagnetization data and pTRM checks; frac ensures the majority of the remanence is used to calculate paleointensity;  $k$  detects deviations from linearity by fitting a circle to the data to determine the amount of curvature. CCRIT also tests for consistency between estimates at the site level by setting thresholds on the percentage of scatter,  $\beta_{\%}$  and intensity of scatter,  $\beta_{\sigma}$  permitted at a site. Twenty-eight of our original 135 sites passed these selection criteria (Table)

## 3.3 Paleodirection

The results of the demagnetization experiment vary (Figure) from a single stable direction to multiple unstable directions. Multiple directions with distinct coercivity and blocking temperature spectra decay along one direction at low field and temperature treatments then abruptly shift and decay along a different direction for the final characteristic remanent magnetization (ChRM) (Figure a). The low temperature and low coercivity component may result from a viscous remanent magnetization or a partial overprint that is typically removed after the first or second treatment. Multiple components with overlapping coercivity or blocking temperature spectra appear as zig-zagging or gradual shifts in the demagnetization curve. The zig-zagging may result from tails, if the thermal demagnetization data was derived from an IZZI experiment. If the directional components are removed in different proportions between each treatment step, then we would observe gradual changes in the magnetization direction. We applied a set of criteria (Table) to determine the final stable component of the demagnetization vector, the ChRM (Figure b). We set  $n$  - the minimum number of consecutive demagnetization steps - to 4 and constrained the direction with MAD and DANG. To ensure consistency within a site, we set  $N$ , the minimum number of samples per site, to 5, set a maximum threshold for  $\alpha_{95}$  and a minimum threshold for  $\kappa$  [Fisher, 1953] a precision parameter to quantify the dispersion in the directions. One-hundred and twenty-six sites yield reliable paleodirections (Table)

## 3.4 Hysteresis and FORC

Several sites- mc1030, mc1115, mc1147, and mc1157 – passed CCRIT and included estimates from samples that were collected from the interior [Mankinen and Cox, 1988; Tauxe and Staudigel, 2004; Lawrence *et al.*, 2009] and from the surface of the same lava flow (Figure). At each of these sites, the estimates from the interior are  $2^{\circ} - 8^{\circ}$  lower than the paleointensity estimates from the lava flow tops. We selected sister specimen for hysteresis



loops and FORCs to examine the micromagnetic components - domain state and interaction-  
that may explain the difference.

Although the sites passed CCRIT, the specimen exhibit a mixture of magnetic components (Figure). We interpret the ridge in the FORC diagram at  $B_u = 0$  (Figure Xb) as the contribution from single domain grains after *Roberts and Verosub* [2000] and *Pike et al.* [2001]. The distribution of coercivities ( $B_c$ ) ranges from 0 and 50 mT and peaks between 0 and 20 mT. The contours spread vertically from this ridge which reflects the level of interaction fields between the single domain grains. In multi-domain grains this peak is offset from the  $B_u = 0$  mT axis and the contours follow a steep gradient that extends beyond 30 mT (Figure a). Each specimen contains some degree of superparamagnetic behavior as inferred from the vertical ridge at  $B_c = 0$  mT that peaks around  $B_u = 0$  mT (Figure c).

### 3.5 Ar-Ar

## 4 Discussion

### 4.1 Examining the GAD structure

#### 4.1.1 Paleodirections

Our new, robust dataset consists of 126 site-mean directions that pass our selection criteria (**table 3**). It includes 54 reverse polarity and 79 normal polarity site-mean directions (**table 5**). The paleomagnetic site-mean directions were separated by polarity then transformed to their corresponding Virtual Geomagnetic Poles (VGPs) (Figure ). VGP is the position of the geomagnetic dipole that would generate the direction measured at a particular latitude. We calculated the paleomagnetic pole and  $\alpha_{95}$  by taking the average of the VGPs for the normal polarity sites (declination  $208.3^\circ$ , inclination  $86.2^\circ$ , and  $\alpha_{95}$   $5.66^\circ$ ) and the reverse polarity sites (declination  $308.7^\circ$ , inclination  $-85.6^\circ$ , and  $\alpha_{95}$   $8.88^\circ$ ). We applied a bootstrap reversal test ([*Tauxe et al.*, 1991]) on the reverse and normal directions. The directions pass the reversal test, so the two sets are indistinguishable (see the supplementary material) and we can calculate the paleopole from the successful VGPs of the entire dataset (declination  $176.24^\circ$ , inclination  $86.89^\circ$ , and  $\alpha_{95}$   $4.92^\circ$ ). The 95% confidence bounds of the paleopole includes the spin axis (Figure ), so the paleodirections from our study are consistent with a GAD.

#### 4.1.2 Paleointensities

Our new paleointensity dataset consists of twenty-eight sites that pass CCRIT (**table 1**) and include both normal and reverse directions (**table 6**). We converted the paleointensities to their corresponding Virtual Axial Dipole Moment (VADM) to compare intensity estimates across latitudes (Figure). VADM is the strength of the axial dipole moment that would generate the intensity observed at a given latitude. Our twenty-eight sites yield a  $35.75 \pm 7.30$  median intensity and a  $44.02 \text{ ZAm}^2 \pm 3.05 \text{ ZAm}^2$  median VADM. Our median intensity estimate is consistent with *Lawrence et al.* [2009] and half of the modern intensity measured in the Erebus Volcanic Province ( $\sim 62 \mu\text{T}$ ).

To assess the structure of the paleomagnetic field over the Plio-Pleistocene, we compare our results to globally distributed paleointensity data stored in the PINT database. We do not observe the latitudinal dependence of intensity expected of a GAD generated field. The paleointensity measured at the high southerly latitudes still appears depressed when compared to the global paleointensity dataset over the Plio-Pleistocene. Before we conclude this depressed intensity near the pole reflects the structure of the paleomagnetic field, we must repeat this same robust study of the 0 – 5 Myr paleomagnetic field at several latitudes [*Dossing et al.*, 2016; *Wang et al.*, 2015].



## 4.2 Examining the role of sampling material

Several of our sites that passed CCRIT include specimens from both the interior and the surface of the same lava flow. We assume a single lava flow cooled instantaneously, so the surface and interior of the flow should preserve identical intensities however, at several sites the interior yields systematically lower paleointensities (Figure ) than the samples from the surface. A slower cooling rate may result in a higher intensity of magnetization [Dodson and McClelland-Brown, 1980] so we tested the effect of cooling rate on the TRM of these samples by conducting a cooling rate experiment. Each specimen measured a higher remanence following rapid cooling than slow cooling, but the correction required for the slowly cooled specimens from the interiors is greater than the corrections required for slowly cooled flow tops at the same site (see the supplementary material). Therefore differences in the cooling history between the two sampling regions does not explain the lower paleointensities we measure in the interior.

Next, we tested whether differences in domain state and interaction could explain the behavior by measuring hysteresis loops and FORC diagrams. The magnetic moments in specimen from mc1115 (Figure ) and mc1147 (see the supplementary material) include a superparamagnetic component, a single domain component and some degree of interaction [Roberts and Verosub, 2000], but the domain structure of specimen from the interiors appears identical to those from the flow tops at the same site. Therefore, differences in domain states does not account for the higher paleointensities measured in the samples collected from the surface.

In addition to cooling rate and domain state, we investigated whether non-linear TRM acquisition could explain the bias in the intensity estimates from the interior. During the in-field step of the IZZI experiment, we applied a  $30 \mu T$  field to our specimens collected from the surface during the 2016/2017 field season. Lawrence *et al.* [2009] cooled some specimens from the interior in a  $25 \mu T$  field and others in a  $30 \mu T$  field. Therefore, we performed a non-linear TRM acquisition test to determine whether the lower intensities measured in the interiors resulted from the lower intensities applied during the IZZI experiment. Each successful specimen was subjected to a total TRM in a  $10 \mu T$ ,  $20 \mu T$ ,  $30 \mu T$ ,  $40 \mu T$ ,  $50 \mu T$ , and  $60 \mu T$  field and each specimen acquired a remanence in proportion to the strength of the applied field (see the see the supplementary material). Neither cooling rate, domain state, nor non-linear TRM acquisition accounts for the lower intensities recorded by the specimen sampled from the interior of the lava flows. Only six of our twenty-eight successful sites include paleointensity estimates from both the surface and the interior, so a full investigation on the role of sampling material on paleointensity estimates would require a larger sample size.

## 5 Conclusions

We present a robust study of the paleomagnetic field over the Plio-Pleistocene in the Erebus Volcanic Province, Antarctica ( $-77.84^\circ$ ,  $166.69^\circ$ ) and eleven new  $^{40}\text{Ar}/^{39}\text{Ar}$  results. We recovered a paleopole at  $176.24^\circ$ ,  $86.89^\circ$  from 126 independent sites that were subjected to both thermal and AF demagnetization and then filtered using a set of strict selection criteria. The  $\alpha_{95}$  of the paleopole is  $4.92^\circ$  and encompasses the spin axis so the paleodirections measured from the EVP during the Plio-Pleistocene are consistent with a GAD field. We also conducted an IZZI-modified Thellier-Thellier experiment and applied the CCRIT set of selection criteria to estimate paleointensity. Twenty-eight sites passed our criteria and recorded a  $35.75 \mu T \pm 7.30 \mu T$  median intensity and a  $44.02 \text{ ZAm}^2 \pm 3.05 \text{ ZAm}^2$  median VADM. Compared with global paleointensity estimates stored in the PINT database, our results from Antarctica are lower than expected for a purely GAD generated field. Before we conclude that this result is representative of the paleomagnetic field structure, we recommend that this extensive study is replicated at different latitudes to ensure high quality paleointensity estimates, appropriate temporal overlap, and adequate global coverage.

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The supporting information can be found — and the entire data set at —

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