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Changes in the Saltation Flux Following a Step-Change in Macro-Roughness

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Under sediment supply limited saltation conditions a step change in macro-roughness alters the saltation reducing the horizontal flux by an amount that scales with the roughness density. The mean diameter of the saltating particles decreases in the horizontal and vertical dimensions with increasing distance. The e-folding height as a function of particle diameter increases with increasing distance, suggesting the particles are increasing their speed, which is likely due to the removal of slower particles as they travel through the roughness.

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1 Changes in the Saltation Flux Following a Step-Change in Macro-Roughness

2 J.A. Gillies, V. Etyemezian, G. Nikolich, W.G. Nickling, J.F. Kok

3 ABSTRACT

The effect of a step change in macro-roughness on the saltation process under sediment supply limited conditions was examined in the atmospheric boundary layer. For an array of roughness elements of roughness density λ =0.045 (λ =total element frontal area/total surface area of the array) the horizontal saltation flux was reduced by 90% (\pm 7%) at a distance of \approx 150 roughness element heights into the array. This matches the value predicted using an empirical design model and provides confidence that it can be effectively used to engineer roughness arrays to meet sand flux reduction targets. Measurements of the saltation flux characteristics in the vertical dimension, including: saltation layer decay (e-folding) height and particle size, revealed that with increasing distance into the array, the rate of mass flux change with increasing height decreased notably, and (geometric) mean particle diameter decreased. The distribution of the saltation mass flux in the vertical remains exponential in form with increasing distance into the roughness array, and the e-folding height increases as well increasing at a greater rate as particle diameter diminishes. The increase in e-folding height suggests the height of saltating particles is increasing along with their mean speed. This apparent increase in mean speed is likely due to the preferential removal, or sequestration, of the slower moving particles, across the size spectrum, as they travel through the roughness array.

Keywords: saltation flux, macro-roughness, e-folding height, mean particle size changes

21 Introduction

Wind blowing over a surface with loose sand that exceeds a critical shear stress on the surface will entrain these particles and move them downwind as bed-load in three recognized modes of transport: creep, reptation, and saltation. The characterization of the saltation flux, in which particles move in a series of repeated ejections and travel in ballistic trajectories followed by

impact on the surface has garnered considerable research effort in the earth sciences. The saltation process plays critical roles in aeolian bedform development across multiple scales (e.g., Anderson, 1987; Claudin and Andreotti, 2006; Dúran et al., 2011; Gillies et al., 2012; Parteli et al., 2014; Schmerler et al., 2016) and is a dominant mechanism driving the emission of dust-sized particles, which represent the suspended load in aeolian transport (e.g., Shao et al., 1993; Shao, 2001; Kok et al., 2012). It is now generally acknowledged that the horizontal mass flux of wind-driven sand decreases exponentially with elevation above the surface (e.g., Bagnold, 1941; Chepil, 1945; Williams, 1964; Sørensen, 1985; Namikas, 2003; Farrell and Sherman, 2006, 2013; Martin and Kok, 2017a) and the general form of the relationship is: $Q(z) = Q_0 \exp^{(-z/z_q)}$ (1)where Q(z) is the stream-wise mass flux at elevation z above the bed, Q_o is the scaling parameter for the profile and z_q the characteristic saltation layer decay height (i.e., the e-folding height). The magnitude of z_q depends on the typical hop height z_{hop} of saltating particles (Bagnold, 1941; Owen, 1964; Ungar and Haff, 1987; Namikas, 2003). That is (Martin and Kok, 2017a), (2) $z_q = C_0 z_{hop}$ where C_Q is a dimensionless constant of order 1. For ballistic trajectories, the hop height z_{hop}

depends on the vertical speed, v_{zo} (m s⁻¹) with which saltating particles leave the surface upon rebound or ejection:

$$z_{hop} = \frac{v_{zo}^2}{2g}$$
(3)

47 where *g* is the acceleration due to gravity (m² s⁻¹). Since v_{zo} is closely linked to the mean 48 particle speed (Owen, 1964; Kok, 2010), Eqs. (1) – (3) imply that the e-folding height z_q is a

49 sensitive function of the mean particle speed. In other words, a measured increase (decrease) 50 in z_q denotes an increase (decrease) in mean particle speed.

Observations of the change in mean grain size as a function of height show complex patterns including a decrease in mean grain size with height (Williams, 1964; Li et al., 2008) while other available data reveal instances of an increase in mean grain size with height (Williams, 1964; Xing, 2007; Speirs et al., 2008; Gillies et al., 2013). Additionally, it has often been observed that at some distance above the surface the grain size trend can reverse (Williams, 1964; Xing et al., 2007; Farrell et al., 2012; Zhang et al., 2017), which may be indicative of a change in transport system dynamics (Andreotti et al., 2004). The available data on particle size relationships in the vertical and horizontal dimensions do not present a simple or consistent picture of grain motions and trajectories leading to a unified model for a bed of mixed particle sizes.

Recent research has documented the change in the mass flux of particles in saltation as the saltation cloud moves from relatively flat, unobstructed flow conditions to one in which the flow encounters and is perturbed by the presence of large roughness elements (e.g., Gillies et al., 2006, 2007, 2014, 2015; Gillies and Lancaster, 2013). Measurements of the saltation flux exterior and interior to a change of roughness show clearly that the mass flux decreases exponentially with increasing distance into the roughness, reaching a new equilibrium flux at >100 roughness element heights (Gillies et al., 2015). The rate of change of the mass flux with distance and the resultant roughness adjusted saltation flux scale with the roughness density (λ) (Gillies et al., 2015):

 $\lambda = n b h/S \tag{4}$

where *n* is the number of roughness elements occupying the area $S(m^2)$, *b* is element breadth (m), and *h* is element height (m). Both Gillies and Lancaster (2013) and Gillies et al. (2015)

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report this scaling relationship based on spatially-distributed measurements of saltation mass 73 74 flux at only one height above the surface, which allowed for a comparison of the spatial characteristics of the mass flux in response to its interaction with large roughness elements. 75 Although this recent research has established that the change in saltation mass flux at a specific 76 77 height scales as a function of λ as it responds to the presence of macro-roughness, it has not examined how the scaling changes over the vertical dimension. It has also not been 78 79 determined if the observed change in mass flux is accompanied by a change in the mean grain diameter. 80 Understanding how large roughness elements affect sand transport is important as there are 81 82 many natural surfaces where sand transport occurs among large roughness elements including 83 vegetation or sediment mounds covered in vegetation, for example nebkhas (Wolfe and Nickling, 1994; Al-Awadhi, 2014; Gillies et al., 2014), or varying amounts of solid roughness 84 85 elements composed of rocks (Gillies et al., 2010, 2013; Lancaster et al., 2010). Arrays of large 86 roughness elements that have been manufactured have been tested to evaluate their 87 effectiveness to serve as a method to reduce sand transport and the associated dust emissions, 88 for example, at the Keeler Dunes, Keeler, CA (Gillies et al., 2015). Increased understanding of 89 how roughness modulates saltation, which in turn affects the dust emission process, may result 90 in improved saltation and dust flux models that account for roughness effects and also ways to 91 better use roughness to reduce dust emissions, improve air quality, and meet air quality standards. 92 In this work, we present results from a field study where changes in vertical profiles of sand flux 93 94 and size, before and after encountering regularly spaced roughness elements, were measured

setting there were no constraints on the saltation system, the boundary layer flow, or the size of
the roughness elements that were used, as would be the case for wind tunnel based

to quantify the effect of the step change in roughness. By carrying out this study in a field

98 experiments (White and Mounla, 1991; Sherman and Farrell, 2008). The supply of sediment,
99 however, was not unlimited upwind of the array of roughness, nor were the elements resting on
100 a deformable bed of sand.

10 101 Materials and Methods

A new set of time-integrated, particle size distribution- and height-resolved mass saltation flux profiles were collected as saltating particles travelled through an array of macro-roughness elements in the atmospheric inertial sublayer in 2015 and 2016. This was carried out as part of an experiment to evaluate methods to reduce sand transport and dust emissions at Owens Lake, CA, which is a source of dust known to contribute to local and regional air quality degradation (Ono et al., 2011). Measurement of the vertical saltation flux profiles were made upwind of, and at defined distances through, a large array of large roughness elements to investigate how the mass flux is affected in the vertical and along-wind directions by a step-change in roughness.

The Roughness Elements and the Roughness Array

The roughness elements used in this experiment were commercially manufactured, rectangular plastic bins (model SNT300) with the dimensions of an individual element being: 0.725 m (length) × 0.38 m (height) × 0.45 m (width). The roughness elements were placed and organized into a staggered array pattern of fixed dimensions on an area of crusted playa surface at Owens Lake, CA, where sand transport was also known to occur (Fig. 1). The open fetch upwind of the northerly and southerly edges of the array well-exceeded 500 m. The roughness array was composed of 1620 elements with the length dimension of the elements set perpendicular to the predominantly bi-modal wind directions of 147° and 327°, determined from a 21 year record from a nearby 10 m high meteorological tower operated by the Great Basin Unified Air Pollution Control District (Bishop, CA). The center-to-center spacing between elements was 2.48 m as was the row-to-row distance (Figs. 1 and 2). The elements were

weighted down by adding sediment to prevent them from being shifted by strong winds. The length and width of the array was 100 m, creating an area of 1 hectare with λ =0.045. The normalized (total) sand flux (NSF_{τ} , the ratio of sand flux measured internal to the array to that measured upwind and external to the array) was expected to be 0.10 for sediment in transport deep into the roughness array (i.e., horizontal distance >100 element heights) based on Gillies et al.'s (2015) Eq. 4.

Sand Flux and Wind Instrumentation

The test area was instrumented to measure sand flux external and internal to the array. Table 1 defines the positions of the instruments relative to the roughness array. Sand flux was measured using modified BSNE-style traps (Fryrear, 1986) as described by Gillies et al. (2012, 2013), which they used to measure sand flux in the McMurdo Dry Valleys in Antarctica. The version used here did not make use of the automated opening and closing of the trap inlets. Briefly, the traps have four wedge-shaped compartments that self-orient into the wind (Fig. 3). Each trap had four catchers spaced logarithmically with heights between ≈ 0.17 m and ≈ 1.3 m. The opening of each catcher is $0.02 \text{ m} \times 0.05 \text{ m}$. The total collection volume per receptacle is approximately 0.0025 m³. The four catchers are connected through the tail fin assembly, which keeps them all facing the same direction and aligned into the wind. By convention, the lowest trap compartment is 1, the highest 4. Eight traps were used in this experiment. By convention, the trap furthest to the north is designated as 1, the furthest south as 8. The traps interior to the array were set \approx 2.5 m from the centerline with the offset staggered to the west and then the east for successive traps through the array. In addition, 30 Cox Sand Catchers, a single measurement height trap (Gillies et al., 2015), were placed into the roughness array with their annular opening at 15 cm above the surface. Eight were collocated with the BSNE traps and the remaining were located at the same positions as the BSNE traps, with one along the center-

line, and the other two half way to the edge on either side of the center-line. This allowed for a more spatially diverse measurement of saltation flux through the roughness array. Wind speed at five heights above ground level (1.25 m, 2.20 m, 3.87 m, 6.82 m, and 10 m) and wind direction at 10 m were measured on four 10 m towers. Their positions exterior and interior

to the array are provide in Table 1. To characterize the airflow among the roughness elements 20 anemometers and four wind vanes were placed external and internal to the array at a height of 0.25 m above ground level (AGL). This is 0.13 m below the top of the roughness elements. One anemometer and one wind vane were placed 3 m upwind of the center-line of the array also at 0.25 m AGL on the north end of the array. The other 19 anemometers and four wind vanes were placed inside the array from its northern boundary extending \approx 50 m into the array. The positions of the anemometers and wind vanes relative to the roughness elements are shown in Fig. 2 and this pattern is repeated four times between the northern edge of the roughness array and halfway through (\approx 50 m).

Particle Size Distribution Measurements

Sediment samples retrieved from the trap compartments were returned to the lab for gravimetric and particle size analysis. The sediment samples were weighed on an electronic balance (Model MS1602S, Mettler-Toledo, Columbus OH) to a precision of 0.01 g. Following the weighing, sediment samples were submitted for particle size analysis to the DRI Soils Characterization Lab (Desert Research Institute, Reno NV USA) using a Malvern Mastersizer 3000 (Malvern Instruments Inc., Westborough MA). The particle size range covered was 0.75 µm to 1500 µm with 34 particle diameter bins reported.

Results

Wind Conditions

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170 Although saltation threshold was not monitored at the field site, as reported for a nearby sand-171 dominated surface (Gillies et al., 2015) transport conditions are likely to occur when wind speed measured at 4 m AGL exceeds 5 m s⁻¹. Wind roses for the met towers exterior to the edge of 172 the roughness array for winds $\geq 6 \text{ m s}^{-1}$ measured at 10 m shows the likely sand transport 173 174 directions are dominated by winds from the south and south-south-east and the north-northwest (Fig. 4). For the case of wind approaching $\pm 10^{\circ}$ from perpendicular to the rows, the near 175 surface spatially-averaged wind speed at the normalized distance ND≈137 (ND=horizontal 176 distance/element height) was reduced by approximately 49% (±13%) compared to the 177 measurement of wind at the same height upwind of the array (Fig. 5). 178 Although a small percentage of wind speeds exceeded 6 m s⁻¹ at 10 m with approach angles to 179 the array between 45° and 90° from perpendicular to the front of the array during the different 180 181 collection periods, their effect on the sand collected in the BSNE traps is likely not significant 182 enough to alter the sand flux reduction pattern with increasing distance into the roughness that 183 the array orientation was designed to characterize. The reasons are two-fold: 1) the above threshold winds for this direction range for the east and west sides of the array is between 6% 184 (east) and 12% (west) of the entire time winds were measured to be $\ge 6 \text{ m s}^{-1}$ at 10 m, and 2) as 185 wind direction approach angle to the front of the array changes between 45° and 90°, λ changes 186 187 from 0.052 to 0.033 due to the changing frontal area of the elements, which would reduce the 188 sand flux by 92% (at 45°) to 81% (at 90°) according to Gillies et al. (2015, Eq. 4) for BSNEs along the center-line of the array. 189 **Total and Vertically Distributed Mass Saltation Flux** 190

Four collections of the BSNE traps were made during the duration of the study: 2-10-2015, 29-192 12-2015, 17-3-2016, and 16-5-2016, with each collection representing multiple transport events. This results in mixing of different transport directions in this bi-modal sand transport wind

194 regime. The further downwind from the leading edge where the saltating particles enter the

array, the lesser this mixing effect occurs as it is expected that the flux will diminish with increasing distance into the roughness array as established by Gillies et al. (2015). The heights of the sample collection, mass of the trap compartment samples, and the corresponding normalized height (NH) and normalized mass (NM) samples are provided in Table 2. The height of collection is defined as the height above the surface of the geometric mean position of the rectangular trap opening (Ellis et al., 2009). NH is defined as: H_n/H_1 , where H_n is the distance above the ground to the position of the geometric mean of a rectangular opening in a trap compartment and H_1 is the distance above the ground to the position of the geometric mean of the rectangular opening of the lowest trap compartment. NM is defined as: $mass_n/mass_1$, where mass_n is the mass in trap compartment n divided by the mass in the lowest trap compartment, i.e., mass₁ at H_1 . For each collection period the total trap mass (i.e., the sum of the four compartments) divided by the total mass of a trap exterior to the array defines NSF_{T} , which is plotted as a function of ND for each collection period in Fig.6. These data show a rapid decrease in NSF_{τ} with increasing distance into the array from either side. For all sample periods the plots show that sediment transport from south to north was dominant. Although as Fig. 6 shows there must have been some transport events from north to south as traps 1 and 2 show higher values than the interior traps. The strongest signal of change in NSF_{T} as a function of ND is observed in the sequence of traps: 8, 7, 6, 5, 4, which represent the transport condition of sand entering the array on the

southern edge and its subsequent movement towards the opposite edge. The relationship between NSF_T for the four collection periods and ND, and two trap types, is shown in Fig. 7, with the error bars representing the standard deviation of the mean. As observed in previous studies NSF_{τ} decreases exponentially with increasing ND up to a point, after which it stabilizes

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3 4	219	to a constant value (Gillies et al., 2015). For this array of roughness elements NSF_{T} stabilizes
5 6	220	at ND $pprox$ 160 - 212, with a mean value of 0.10 (±0.07). This relationship can be expressed as:
7 8 9	221	$NSF_T = (1 - NSF_{\infty}) e^{(-a ND)} + NSF_{\infty} $ (5)
10 11 12	222	where <i>a</i> is a dimensional constant and NSF_{∞} the flux fully adjusted to the step change in
13 14	223	roughness.
15 16	224	Equation 5 was not derived formally; rather, it represents a plausible general form to fit the data
17 18 10	225	from a step increase (possibly also step decrease) in roughness. The parameter NSF_{∞} is
19 20 21	226	straightforward to understand and can be estimated easily either from direct measurement or
21 22 23	227	based on Gillies et al.'s (2015) Eq. 4.
24 25	228	NM plotted as a function of NH represents the form of the normalized horizontal flux relationship
26 27	229	(NF) based on the vertical variation of mass and defines the rate of change of the particle mass
28 29 30 31 32 33 34	230	flux above the surface at a measurement position. This relationship for pairs of traps with the
	231	same ND from the edges of the roughness array and the mean values of NM representing the
	232	four collection periods are shown in Fig. 8. The error bars in each case represent the standard
35 36	233	deviation of the mean NM. For each trap the vertical flux is well described by an exponential
37 38	234	function. The two traps external to the array (i.e., T1 and T8) show a greater rate of change in
39 40	235	<i>NM</i> with <i>NH</i> than do the trap pairs interior to the array (Fig. 8).
41 42 43	236	The z_q parameter in Eq. 1 is the characteristic saltation layer decay height. The change in this
44 45	237	variable as a function of ND for the traps combinations: 1, 2, 3, 4, and 8, 7, 6, 5, which groups
45 46 47 48 49	238	them based on having the same ND values for north to south and south to north transport
	239	directions is shown in Fig. 9.
50 51 52	240	Particle Size Distribution Changes in the Vertical Saltation Flux
53 54	241	The effect of the roughness on the particle size distribution characteristics of the saltation flux in
55 56	242	the vertical dimension is revealed through the Malvern Mastersizer-generated data and the
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2 3 4	243	calculation of grain size distribution characteristics using the Gradistat particle size analysis
5 6	244	software (Blott and Pye, 2001; Kenneth Pye Associates Ltd., UK). Grain size statistics
7 8	245	generated by Gradistat are all based on the method of moments and reported as geometric
9 10	246	mean diameter (μ m). The relationship between mean diameter and NH for the collection
11 12	247	periods: 02-10-2015, 29-12-2015, and 17-03-2016 for the sequence of traps 8, 7, 6, 5 (i.e.,
13 14	248	south to north transport) is shown in Fig. 10. In these cases, the mean diameter decreased with
15 16	249	increasing NH above the surface, and the relationship is well-described, most consistently, by a
17 18	250	best-fit logarithmic relationship. Pooling the data by trap number for the three sampling intervals
19 20	251	and normalizing the mean diameter by dividing each mean particle diameter by the mean
21 22 22	252	diameter for the highest receptacle (i.e., mean diameter _{$Hn/mean$ diameter_{$H4) gave the result that$}}
23 24 25	253	for Trap 8, exterior to the roughness, between the lowest and highest collection heights the
26 27	254	mean diameter differed by a factor of $3.0 (\pm 1.1)$. For the three interior traps (7, 6, and 5) the
28 29	255	difference in mean diameter from lowest to highest is a consistent decrease by a factor of 2.0
30 31	256	(\pm 0.1). The mean diameter at similar collection heights also changes systematically with ND as
32 33	257	shown in Fig. 11. The decrease in mean diameter as a function of collection height and ND is
34 35	258	observed for collection heights 1 (\approx 0.2 m), 2 (\approx 0.5 m), and 3 (\approx 0.8 m), and is most consistently
36 37	259	described by a logarithmic relationship. For collection height four (≈ 1.3 m) in two of the
38 39	260	collection periods an increase in mean diameter was observed moving from outside the array to
40 41	260	ND=48.9 thereafter little variation of mean diameter between ND =48.9 and ND =159.9 was
42 43	262	observed at this height
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46 47	263	By partitioning the mass flux in each trap compartment by the mass fraction associated with the
48 49	264	eight particle size bins for the sand-sized components of the material in each trap compartment
50 51	265	(i.e., 750 $\mu m,$ 375 $\mu m,$ 225 $\mu m,$ 175 $\mu m,$ 137.5 $\mu m,$ 112.5 $\mu m,$ 87.5 $\mu m,$ and 68.7 $\mu m),$ z_q for
52 53	266	each of these particle diameters can be estimated at the four measurement positions: ND=0,
54 55 56	267	16.3, 48.9, and 101.2. The regression-derived value for z_q as a function of mean diameter for
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the available data (from both transport directions), is shown in Fig. 12. As this Fig. shows, z_q increases for each particle diameter with increasing distance into the array, and as particle diameter decreases, z_q increases at a faster rate. The rate of change of z_q with increasing *ND* (i.e., *m* term in $z_q = m$ (*ND*)+*b*, for each particle diameter in Fig. 11), as a function of particle diameter is shown in Fig. 13 and is well-defined by a power relationship.

273 Discussion

17 274 Effect of Roughness on the Mass Flux

As previously observed by Gillies et al. (2006, 2015), NSF_T scales as a function of λ . Gillies et al.'s (2015) Eq. 4 predicts that for λ =0.045, the equilibrium NSF₇ for an array of large roughness elements should be 0.10. The mean NSF_{T} for the interior of the array for this experiment for ND=159.9 was 0.10 (±0.07), which matches the model prediction very closely. The roughness adjusted NSF_{τ} value was based on using measurements at multiple heights above the surface and as can be seen from Fig. 7, matches closely the saltation flux reduction pattern measured at the same height low to the ground and over a wider area.

The rate of change of NSF_{T} with increasing ND for that portion of the curve shown in Fig. 7 before the fully-adjusted flux is attained (i.e., ND=0 to ND=101) for this roughness array can be compared with other experiments that have measured the effect of roughness on NSF_{T} . Defining this portion of the curve similarly to Gillies et al. (2006, 2015) and Lancaster and Gillies, 2013) as $NSF_T = a e^{(-B ND)}$, the value of B for this experiment for the BSNE trap data is -0.17 and -0.021 for the CSC trap data. Given the stochastic and variable nature of aeolian transport (e.g., Gares et al., 1996; Stout and Zobeck, 1997; Jackson et al., 2006) and the variability in wind direction, transport and surface conditions among the available studies, the values of B for this experiment fit closely with the predicted value of B=-0.016 based on the relationship presented by Gillies et al. (2015, their Fig. 11) that uses data from three

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experiments. In general, the new Owens Lake NSF_{τ} relationships fit quite well with previous measurements, which indicates that the established scaling relationship between sand flux and λ (Gillies et al., 2015, Eq. 4) is very robust at least for roughness that is relatively closely spaced and of all the same height. Chappell and Webb (2016) argue that λ is less successful in predicting shear stress partitioning and shelter effects for superposed roughness that covers a range of element heights and widths, which may also reduce the applicability of the Gillies et al., (2015) equation for estimating the effect of non-homogeneous roughness on sand transport. Their equation does, however, appear very useful for designing arrays of large roughness elements to control saltation flux and the associated dust emissions. Effect of Roughness on the Particle Size Distribution of the Saltators In general, for all trap samples the mean particle diameter decreases as a function of height above the surface (for H > 0.17 m), exterior and interior to the roughness. This matches the observations of Li et al. (2008) for wind tunnel tests using sand of mixed size and this pattern of decreasing grain diameter with height has also been observed by others in wind tunnel experiments (e.g., Williams, 1964) and in field measurements (Gillies et al., 2013), but as noted earlier this pattern is not universal (e.g., Williams, 1964; Xing, 2007; Speirs et al., 2008; Gillies et al., 2013). The reduction of mean particle diameter with passage of the saltators through the roughness array suggests that there is preferential removal of coarser grains through the height of the saltation layer. The removal is likely due to the preferential trapping of grains in zones of wind shear in the lee of the roughness elements that are below the threshold to keep particles in motion as well as long enough for particle inertia to be insufficient to carry them through to a zone of above threshold conditions. Deposits of sand were observed in the lee of roughness elements, most notably for the rows nearest the edges (i.e., rows 1 to 8 and 32 to 40). Effect of Roughness on the Mass Flux in the Vertical Dimension

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316 This new data set, for the first time, allowed for an examination of the effect of roughness on the 317 mass flux based on integration in the vertical dimension using the multi-height BSNE traps. As 318 Figs. 9, 10 and 11 show, the saltation flux is altered by its interaction with the roughness array in 319 the horizontal and vertical dimensions. With increasing distance into the array the sand in 320 saltation shows an increase in z_a (Fig. 9), which can, in part, be explained by the diminishing 321 total mass flux. In addition, the mean particle diameter at all but the top measurement position 322 decreases with increasing distance into the array (Fig. 11). The particle size distribution data 323 offers some explanatory power for why these patterns occur.

324 According to Martin and Kok (2017a), for saltating sand over an essentially roughness-element-325 free and relatively flat surface with no sediment supply restrictions, mean particle speed and z_q remain constant with changing surface shear stress (τ , N m⁻²) above threshold. Martin and Kok 326 (2017a) note that z_q is not constant for all surfaces, but is likely controlled by the particle size 327 328 distribution of the source sand. In the case studied in our experiment, where the saltating sand 329 encounters and interacts with the large scale roughness, there are profound changes in the mass flux as well as z_q as a function of particle diameter. As z_q increases with distance into the 330 331 array it suggests that particle speeds are changing as the saltating particles move through the 332 array at least to ND=101.2. An increase in z_q represents an increase in the characteristic saltator hop height (Owen, 1964; Ungar and Haff, 1987; Namikas, 2003; Kok, 2010). Since hop 333 334 height is determined by the speed at which particles are launched from the surface (Eq. 3), which in turn depends on the mean particle speed (Owen, 1964; Ungar and Haff, 1987; Kok, 335 336 2010), our observation of increasing z_{d} into the roughness array implies that the mean particle speed also increases for all particle sizes and with increasing distance into the roughness. 337 338 The observed effect of z_q increasing with increasing distance into the array and with decreasing particle diameter in the presence of the roughness is different from that observed for saltation in 339 the absence of superposed roughness. Martin and Kok (2017b) report for unconstrained 340

saltation on sand sheets, z_q increases with particle diameter up to the modal diameter of the sand (that is of the source sand) and then decreases for particle diameters less than the modal diameter.

We suggest that the effects on z_q and mean diameter in the case of the superposed roughness are occurring due to the reduction in saltation flux caused by the presence of the roughness elements. Mechanisms that are driving the loss are the change in wind shear exerted on the saltation layer with increasing distance into the array due to shear stress partitioning as the wind adjusts to this step change in roughness (Raupach et al., 1993; Gillies et al., 2007), and the interaction of the saltation with the roughness elements (Gillies and Lancaster, 2013).

According to Gillies et al. (2007), the shear stress ratio (*R*, Raupach et al., 1993) for λ =0.045 at *ND* >127.5, should be \approx 0.62, which would reduce the shear stress on the surface (τ_s , N m⁻²) at that position by ≈38%. Applying the Shao (2005) scaling relationship between shear stress, expressed as the shear velocity u_* (m s⁻¹, where $\tau = \rho_a u_*^2 [\rho_a \text{ is air density, kg m}^3]$) and sand flux elie $(q, kq m^{-1} s^{-1})$:

 $q = C_{Shao} \frac{\rho_a u_*^3}{g} \left(1 - \left(\frac{u_{*t}}{u_*} \right)^2 \right)$

(6)

where C_{Shao} is a coefficient and u_{*t} is the threshold shear velocity, the reduction in q at *ND*=127.5 should be \approx 76% compared to the flux external to the array. At *ND* <127.5 the partitioning of shear stress between the roughness elements and the surface will be changing as a function of ND into the array at an approximately exponential rate to approximately ND=73 to 93 (Gillies et al., 2007). Gillies and Lancaster (2013) observed first that the reduction in saltation flux observed in an array of large roughness elements was greater than could be attributed to shear stress partitioning by itself and suggested the physical interaction of the saltating sand with the roughness and the environment around the elements accounted for the additional reduction in sand flux by the elements physically impeding the movement of the

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particles. The shear stress partitioning, the physical dimensions of the roughness, and their
distribution on the surface result in a reduction in mass flux through the roughness as defined by
Fig. 7.

The loss of particles from the saltating population explains why the particles appear to increase 368 their mean speed as indicated by the increase in z_q (Figs. 12 and 13). A plausible mechanism 369 370 for this is that slower particles are being preferentially removed as they have a reduced 371 probability of making it deeper into the roughness. Slower moving particles are less likely to 372 have the momentum necessary to carry them through zones of lower shear stress created by 373 the presence of the roughness. Hence as saltators of a given size range pass through the 374 array, their mean speed increases as the distribution of speed is increasingly truncated at the 375 lower end of the distribution, which is supported by the observed increase in z_a with increasing 376 distance into the array (Fig. 9). The wind close to the surface decreases with increasing distance into the array (Fig. 5), so the particles cannot be gaining speed through interaction with 377 higher speed winds. 378

379 The measurements of particle size resolved flux exterior to the roughness provide additional information on the scaling of sand flux with wind speed under these supply-limited conditions. 380 381 Martin and Kok (2017a) demonstrated with saltation flux data that the saltation layer height does 382 not change with u_{\star} implying that mean particle speed also remains constant with u_{\star} . When the 383 flux versus height relationships are binned by particle diameter, it is observed (Fig. 12 at ND=0) 384 that z_q increases with decreasing particle diameter, implying that for this sand transport condition the smaller diameter particles are travelling at a faster speed than the larger diameter 385 particles, suggesting in this case that particle launch velocity increased with decreasing particle 386 387 diameter. A likely reason why z_a increases with particle diameter exterior to the array, as opposed to the observation of Martin and Kok (2017b) showing an increase in z_{α} up to the 388 389 modal diameter, followed thereafter by a decrease, is due to the supply-limited sediment

conditions on the playa. Under supply limited conditions, particle speed is controlled by the restitution coefficient, determined by the soil surface condition, and the wind speed. Smaller diameter particles enter into the array with greater speed than larger ones so only the fastest particles in each particle size bin make it deep into the roughness array, and hence z_q increases for decreasing particle diameter.

395 Conclusions

The effect of a step change in macro-roughness on the saltation process was examined at the full scale in the atmospheric boundary layer. The saltation flux was profoundly altered in the horizontal and vertical dimensions as a result of its interactions with the roughness. The horizontal flux was reduced by 90% (\pm 7%) at a distance of \approx 150 roughness element heights into the array as predicted by the empirical model of Gillies et al. (2015), for an array of roughness with λ =0.045. These results provide further corroboration of the robust nature of this simple model to estimate sand flux reduction for roughness elements of height dimension of ≈ 0.4 m that are evenly distributed across space, providing increased confidence that it can be effectively used to engineer roughness arrays to meet sand flux reduction targets. Measurements of the saltation flux characteristics in the vertical dimension, including particle size characteristics, revealed that with increasing distance into the array mean particle diameter became smaller, reducing in size as a function of distance as well as a logarithmic function of height above the surface. The distribution of the saltation mass flux in the vertical as a function of increasing distance into the roughness array, trended towards a more uniform distribution, although it remained exponential in form. This is only possible if there is a diminishment of the particle flux, which was observed. This diminishment of flux appears to be associated with the loss of the slower speed particles so that with increasing distance into the roughness, the mean

413 particle speed for a given particle diameter is shifted to a higher value, but remains less than the

maximum speed of the initial distribution of speeds. It will require additional experimentation to

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3 4	415	resolve the behaviour of the particle speed distribution as a function of particle diameter through
5 6	416	measurement of actual particle speeds within superposed roughness for both supply and
7 8	417	transport limited conditions. Further information on the particle speed distribution in the
9 10	418	presence of the roughness would also be informative on how the roughness modulated saltation
11 12	419	flux affects the dust emission process.
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Table 1. The positions of the instruments relative to the roughness array.

Instrument	Between Row #s (from north edge)	Horizontal Distance (from north edge), m	Normalized Distance* (<i>ND</i> , from upwind [north] edge defined as zero to southern edge defined as 260.0)							
Trap 1	N/A (exterior)	-10.0	0							
Trap 2	3-4	6.2	16.3							
Trap 3	8-9	18.6	48.9							
Trap 4	16-17	38.5	101.2							
Trap 5	26-27	60.8	159.9							
Trap 6	33-34	80.6	212.1							
Trap 7	38-39	93.0	244.7							
Trap 8	N/A (exterior)	110.0	261.0							
Tower 1	N/A (exterior)	-10.0	0							
Tower 2	16-17	38.4	101.2							
Tower 3	25-26	60.8	159.9							
Tower 4	N/A (exterior)	110.0	261.0							

Table 2. The collected mass and its height of collection for the eight BSNE traps for the four

collection periods.

			2/10/2015		29/12/2015		17/03/2016		16/05/2016			
	Height (m)	NH	mass (g)	NM	mass (g)	NM	mass (g)	NM	mass (g)	NM	Mean NM	Std. D. NN
T1H1 ¹	0.21	1	11.5	1	101.17	1.00	2.79	1	178.78	1.00	1	0.00
T1H2	0.50	2.38	13.43	1.17	30.03	0.30	0.82	0.29	46.47	0.26	0.50	0.44
T1H3	0.83	3.95	NS ²		15.70	0.16	0.46	0.16	16.93	0.09	0.14	0.04
T1H4	1.28	6.10	NS		7.72	0.08	0.04	0.01	3.93	0.02	0.04	0.03
T2H1	0.18	1	11.07	1	174.89	1.00	1.34	1.00	167.81	1.00	1	0.00
T2H2	0.48	2.62	8.95	0.81	56.17	0.32	0.70	0.52	51.03	0.30	0.49	0.23
T2H3	0.79	4.30	9.42	0.85	22.21	0.13	0.42	0.31	17.73	0.11	0.35	0.35
T2H4	1.26	6.90	9.6	0.87	7.61	0.04	0.13	0.10	4.06	0.02	0.26	0.41
T3H1	0.17	1	8.92	1	73.52	1.00	1.18	1.00	56.98	1.00	1	0.00
T3H2	0.48	2.79	8.34	0.93	25.18	0.34	0.36	0.31	20.28	0.36	0.48	0.30
T3H3	0.78	4.61	7.43	0.83	15.98	0.22	0.68	0.57	9.76	0.17	0.45	0.31
T3H4	1.26	7.39	NS		6.57	0.09	0.49	0.41	3.90	0.07	0.19	0.19
T4H1	0.17	1.00	7.86	1	35.95	1.00	2.05	1.00	5.10	1.00	1	0.00
T4H2	0.47	2.76	7.72	0.98	18.05	0.50	1.52	0.74	5.48	1.07	0.82	0.26
T4H3	0.77	4.55	NS		14.28	0.40	1.36	0.66	1.19	0.23	0.43	0.22
T4H4	1.26	7.41	NS		10.13	0.28	0.78	0.38	3.30	0.65	0.44	0.19
T5H1	0.18	1.00	21.49	1	69.52	1.00	2.35	1.00	67.44	1.00	1	0.00
T5H2	0.48	2.67	9.65	0.45	19.42	0.28	2.03	0.86	14.85	0.22	0.45	0.29
T5H3	0.78	4.34	8.43	0.39	13.44	0.19	0.97	0.41	6.82	0.10	0.27	0.15
T5H4	1.26	7.01	7.41	0.34	8.95	0.13	NS		2.42	0.04	0.17	0.16
T6H1	0.20	1.00	40.02	1	70.30	1.00	3.64	1.00	38.16	1.00	1	0.00
T6H2	0.54	2.70	19.37	0.48	29.98	0.43	3.15	0.87	15.05	0.39	0.54	0.22
T6H3	0.81	4.05	7.57	0.19	21.26	0.30	2.49	0.68	7.44	0.20	0.34	0.23
T6H4	1.29	6.43	9.97	0.25	9.34	0.13	1.34	0.37	2.64	0.07	0.20	0.13
T7H1	0.19	1.00	118.01	1	179.80	1.00	23.01	1.00	150.75	1.00	1	0.00
T7H2	0.49	2.58	31.4	0.27	68.81	0.38	9.13	0.40	42.41	0.28	0.33	0.07
T7H3	0.78	4.16	11.78	0.10	38.57	0.21	5.75	0.25	20.02	0.13	0.17	0.07
T7H4	1.27	6.73	8.76	0.07	11.60	0.06	2.41	0.10	4.27	0.03	0.07	0.03
T8H1	0.20	1.00	230.6	1	331.94	1.00	33.53	1.00	412.79	1.00	1	0.00
T8H2	0.46	2.32	44.86	0.19	72.88	0.22	6.69	0.20	44.07	0.11	0.18	0.05
T8H3	0.80	4.02	14.97	0.06	26.50	0.08	2.02	0.06	12.15	0.03	0.06	0.02
T8H4	1.28	6.44	7.82	0.03	7.76	0.02	1.41	0.04	3.35	0.01	0.03	0.01
¹ T1H1: trap	o 1, height 1											
² NS: no sar	nple, trap coi	mpartmen	nt was empty	of sand								



Figure 1. The roughness array on Owens Lake, CA. The view into the frame is towards the south and the dimensions of the roughness array are $\approx 100 \text{ m} \times \approx 100 \text{ m}$.



Figure 2. The spacing of the elements and the positions of the near surface anemometers and wind vane with respect to the roughness elements (gray rectangles). The pattern of near surface anemometers and wind vanes is repeated at four locations through to 50 m into the roughness array.



Figure 3. The modified BSNE style trap (left panel) and the Cox Sand Catcher (right panel) used to collect saltating particles.



Figure 4. Wind roses for the site based on 10 m AGL wind speed and direction measurements at the northern (T1) and southern (T2) edges of the roughness array.

Perez



Figure 5. The change in near-surface wind speed ratio, *WSR*, as a function of normalized distance, *ND*. The x-error bars represent the standard deviation of *ND* based on the *ND* for each individual instrument (Positions 1 through 4) in the clusters of instruments from *ND*=5.2 (\pm 2.9) through to *ND*=137 (\pm 0) (see Fig. 2). The y-error bars represent the standard deviation of the mean *WSR* based on the 5-minute mean *WSR* values.



Figure 6. Mean normalized total sand flux, NSF_{τ} (normalized to trap 8) for the four collection periods. The lines serve only to emphasize the groupings of the data points by collection period Perez.



Figure 7. Mean normalized total sand flux, NSF_T , as a function of normalized distance, *ND*, for traps 4 (*ND*=160), 5 (*ND*=101.2), 6 (*ND*=48.9), 7 (*ND*=16.3), and 8 (*ND*=0) representing the dominant transport of sand from south to north during the observation period for the BSNE traps (black circles, top panel) and the CSC traps (open circles, bottom panel). Error bars represent the standard deviation of the mean value for the multiple measurement periods. $NSF\infty$ was set to the calculated NSF_T value (0.10) for fitting the relationship defined by Eq. (5).



Figure 8. The relationship between normalized mass, *NM*, and normalized height, *NH*, as a function of shared trap positions at *ND*=0 (Traps 1 and 8, top panel), *ND*=16.3 (Traps 2 and 7, middle panel), *ND*=48.9 (Traps 3 and 6), and *ND*=101.2 (Traps 4 and 5, bottom panel). http://mc.manuscriptcentral.com/esp



Figure 9. The mean e-folding height (i.e., z_q in Eq. 1) for the sequence of four traps for sand transport direction from north to the south and south to north at equivalent *ND* positions. Error bars represent the standard deviation of the mean.







Figure 11. Observed change in the mean particle diameter as a function of normalized distance (*ND*) for trap grouping: 8, 7, 6, 5, characterizing the south to north transport direction for three monitoring periods. http://mc.manuscriptcentral.com/esp



Figure 12. z_q as a function of *ND* and particle diameter.



Figure 13. The rate of change in z_q as a function of particle diameter as defined by the *m* term in the relationship $z_q=m ND + b$ for each particle size bin.

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