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# Precise measurement of $2\nu 2\beta$ decay of $^{100}\text{Mo}$ with $\text{Li}_2\text{MoO}_4$ low temperature detectors: preliminary results

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**Abstract.** The half-life of  $^{100}\text{Mo}$  relatively to the  $2\nu 2\beta$  decay to the ground state of  $^{100}\text{Ru}$  was measured as  $T_{1/2} = (6.99 \pm 0.15) \times 10^{18}$  yr with the help of enriched in  $^{100}\text{Mo}$  lithium molybdate scintillating bolometers in the EDELWEISS-III low background set-up at the Modane underground laboratory. This is the most accurate value of the  $2\nu 2\beta$  half-life of  $^{100}\text{Mo}$ .

## INTRODUCTION

Two-neutrino double-beta ( $2\nu 2\beta$ ) decay is detected in 11 nuclei with the half-lives in the range  $T_{1/2} \sim 10^{18} - 10^{24}$  yr [1, 2, 3]. The  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  was first observed by the ELEGANT V counting experiment [4]. After several measurements with the help of different detection techniques [5, 6, 7, 8, 9, 10], the most accurate study of the  $2\nu 2\beta$  decay have been performed in a calorimetric experiment using zinc molybdate ( $\text{ZnMoO}_4$ ) low temperature bolometers (with uncertainty at the level of  $\sim 11\%$ ) [11], and by the NEMO-3 collaboration by detecting the two electrons with a combination of tracking and calorimeter information ( $\sim 6\%$ ) [12]. In the present experiment the  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  was measured by enriched in  $^{100}\text{Mo}$  lithium-molybdate ( $\text{Li}_2^{100}\text{MoO}_4$ ) crystal scintillators as low temperature scintillating bolometers. The preliminary results of the measurements have been reported in [13, 14].

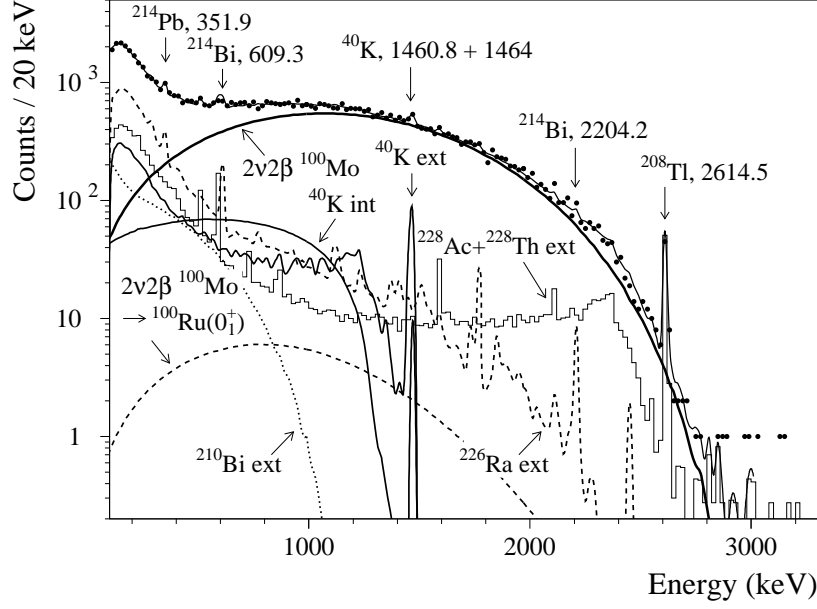
## EXPERIMENT

Four  $\text{Li}_2^{100}\text{MoO}_4$  crystal scintillators produced from molybdenum enriched in the isotope  $^{100}\text{Mo}$  to  $(96.9 \pm 0.2)\%$  with sizes  $\varnothing 44 \times (40 - 46)$  mm and the total mass 808.87 g were utilized in the experiment. Each crystal was equipped with a neutron transmutation doped (NTD) germanium temperature sensor glued on their surface, and with a heavily-doped silicone heater to control the detector thermal response. Germanium discs  $\varnothing 44 \times 0.17$  mm, also equipped with NTD sensors, were used as photo-detectors. Simultaneous detection of the heat and scintillation signals allows discrimination between  $\beta/\gamma$  and  $\alpha$  events to reduce  $\alpha$  background. The R&D of  $\text{Li}_2\text{MoO}_4$  based scintillating bolometers is described in [13, 14, 15]. The detector modules were operated in the low-background cryostat of the EDELWEISS-III dark-matter experiment [16] at the Modane underground laboratory (France). The energy scale and energy resolution of the detectors were calibrated with  $^{40}\text{K}$ ,  $^{133}\text{Ba}$ , and  $^{232}\text{Th}$  gamma sources. E.g., the energy resolution (full width at half of maximum) of the detectors was measured as  $\sim 6$  keV for  $\gamma$  quanta with energy 2614.5 keV of  $^{208}\text{Tl}$ .

## RESULTS AND DISCUSSION

The energy spectrum accumulated with  $\text{Li}_2^{100}\text{MoO}_4$  detectors over the exposure  $42.235 \text{ kg} \times \text{d}$  ( $3.797 \times 10^{23}$  nuclei of  $^{100}\text{Mo} \times \text{yr}$ ) is shown in Fig. 1.  $\alpha$  events have been eliminated from the data by using a light-assisted particle identification with at least  $9\sigma$   $\alpha/\beta$  selection efficiency [13, 14]. In addition, a pulse-shape discrimination cut was then applied to the signals. A total exposure-weighted average  $\beta$  events selection efficiency is  $(96.46 \pm 0.60)\%$ .

Several weak peaks in the spectrum can be ascribed to radioactive contamination by K, Ra and Th of the set-up, while the counting rate above  $\sim 1$  MeV is mainly caused by the  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$ . The observed  $^{40}\text{K}$  peak swelling is consistent with part of the population being due to EC decays of potassium inside the detector plus the atomic shell relaxation following this decay. A model of the experimental spectrum was built from the following components:  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  to the ground state of  $^{100}\text{Ru}$ ;  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  to the first  $0^+$  1130.3 keV excited level of  $^{100}\text{Ru}$  with the half-life  $T_{1/2} = [7.5 \pm 0.6(\text{stat.}) \pm 0.6(\text{syst.})] \times 10^{20}$  yr [17];  $\gamma$  quanta of  $^{40}\text{K}$  from the detector parts;  $\gamma$  quanta of  $^{214}\text{Pb}$  and  $^{214}\text{Bi}$  (contamination of the set-up by radium);  $\beta$  particles and bremsstrahlung  $\gamma$  quanta from  $^{210}\text{Bi}$  (daughter of  $^{210}\text{Pb}$ ) in the materials close to the detectors;  $\gamma$  quanta of  $^{228}\text{Ac}$ ,  $^{212}\text{Pb}$ ,  $^{212}\text{Bi}$  and  $^{208}\text{Tl}$  (contamination of the set-up by thorium; activity of  $^{228}\text{Ac}$  was taken as a free parameter,  $^{212}\text{Pb}$ ,  $^{212}\text{Bi}$  and  $^{208}\text{Tl}$  were assumed in equilibrium with  $^{228}\text{Th}$ ); internal contamination of the scintillators by  $^{40}\text{K}$ ,  $^{87}\text{Rb}$ ,  $^{90}\text{Sr}$  and  $^{90}\text{Y}$  (in



**FIGURE 1.** The energy spectrum of  $\beta/\gamma$  events accumulated with  $\text{Li}_2^{100}\text{MoO}_4$  scintillating bolometers (exposure is  $42.235 \text{ kg} \times \text{d}$ ) and its fit by the components of background in the energy interval 100 – 3000 keV. Energies of  $\gamma$  peaks are in keV.

**TABLE I.** Estimated systematic uncertainties of the  $^{100}\text{Mo}$  half-life (%).

Number of $^{100}\text{Mo}$ nuclei	$\pm 0.4$
Live time	$\pm 0.22$
Pulse-shape discrimination cut to accept $\beta$ events	$\pm 0.60$
Interval of fit	$+0.80$ $-0.86$
Localization of radioactive sources in the set-up	$\pm 0.85$
Monte Carlo simulated models statistic	$\pm 0.30$
Energy scale instability	$\pm 0.46$
$2\nu 2\beta$ spectral shape	$\pm 1.0$
Mechanism of decay (HSD instead of SSD)	$+0.14$
Total systematic error	$+1.80$ $-1.83$

equilibrium). Bulk U/Th radioactivity of the crystals is omitted, taking into account that the activity of  $^{100}\text{Mo}$  in the crystals is three orders of magnitude higher than the limits on activities of U/Th daughters [13, 14]. The models were Monte Carlo simulated using GEANT4 package [18, 19, 20] with initial kinematics given by the DECAY0 event generator [21]. The  $2\nu 2\beta$  distribution was simulated using an assumption about the single-state dominance (SSD) hypothesis, taking into account that the data of the NEMO-3 experiment favors the SSD mechanism in  $^{100}\text{Mo}$  [12]. The model well describes the experimental data in a wide energy interval (see Fig. 1).

The best fit achieved in the energy interval 940 keV – 2860 keV provides the half-life  $T_{1/2} = [6.988 \pm 0.074(\text{stat.})] \times 10^{18}$  yr. The statistical error already does include correlations to the background models. The systematic error includes uncertainties of the number of  $^{100}\text{Mo}$  nuclei, live time of the experiment, pulse-shape discrimination cut to accept  $\beta$  events, interval of fit, localization of radioactive sources in the set-up, statistical fluctuations of the simulated background models, energy scale instability, theoretically calculated  $2\nu 2\beta$  spectral shape, mechanism of decay (high-state dominance instead of SSD). A summary of the systematic uncertainties is given in Table I.

Summing all the systematic uncertainties and the statistical error in quadrature, the half-life of  $^{100}\text{Mo}$  relative to the  $2\nu 2\beta$  decay to the ground state of  $^{100}\text{Ru}$  is:

$$T_{1/2} = (6.99 \pm 0.15) \times 10^{18} \text{ yr.}$$

The half-life value, being the most accurate one, is in an agreement with all the previous counting experiments [4, 5, 6, 7, 8, 10, 11, 12].

An effective nuclear matrix element for  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  to the ground state of  $^{100}\text{Ru}$  can be calculated as  $|M_{2\nu}^{eff}| = 0.186 \pm 0.002$  by using the phase-space factor  $4134 \times 10^{-21} \text{ yr}^{-1}$  from [22]. The effective nuclear matrix element can be written as product  $M_{2\nu}^{eff} = g_A^2 \times M_{2\nu}$ , where  $g_A$  is the axial vector coupling constant,  $M_{2\nu}$  is nuclear matrix element.

## CONCLUSION

The half-life of the  $2\nu 2\beta$  decay of  $^{100}\text{Mo}$  to the ground state of  $^{100}\text{Ru}$  was measured with the highest up-to-date accuracy as  $T_{1/2} = (6.99 \pm 0.15) \times 10^{18} \text{ yr}$  with enriched  $\text{Li}_2^{100}\text{MoO}_4$  scintillating bolometers at the Modane underground laboratory (France). The systematic error is mainly due to the uncertainty in the background model. The half-life and the spectral shape accuracy are expected to be further improved in the CUPID-Mo experiment running now in its first phase with 20 enriched  $\text{Li}_2^{100}\text{MoO}_4$  scintillating bolometers (with mass  $\approx 0.2 \text{ kg}$  each). Precise measurement of the  $2\nu 2\beta$  decay spectral shape can be realized by measurement also of four  $\text{Li}_2\text{MoO}_4$  detectors (already produced from molybdenum depleted in the isotope  $^{100}\text{Mo}$  to 0.007%) in the same conditions.

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