

UC Berkeley

UC Berkeley Previously Published Works

Title

Co-Design of Multijunction Photoelectrochemical Devices for Unassisted CO₂ Reduction to Multicarbon Products

Permalink

<https://escholarship.org/uc/item/4rf1p76g>

Journal

Journal of The Electrochemical Society, 170(12)

ISSN

0013-4651

Authors

Wei, William J

King, Alex J

Bui, Justin C

et al.

Publication Date

2023-12-01

DOI

10.1149/1945-7111/ad10e7

Copyright Information

This work is made available under the terms of a Creative Commons Attribution License, available at <https://creativecommons.org/licenses/by/4.0/>

Peer reviewed

OPEN ACCESS

Co-Design of Multijunction Photoelectrochemical Devices for Unassisted CO₂ Reduction to Multicarbon Products

To cite this article: William J. Wei *et al* 2023 *J. Electrochem. Soc.* **170** 126502

View the [article online](#) for updates and enhancements.

You may also like

- [Design and on-Sun Testing of Tandem III-V Photoelectrochemical Water-Splitting Systems](#)
Micha Ben-Naim, Chase Aldridge, Myles A Steiner et al.
- [A Simplified Successive Ionic Layer Adsorption and Reaction \(s-SILAR\) Method for Growth of Porous BiVO₄ Thin Films for Photoelectrochemical Water Oxidation](#)
Wenlong Guo, Ding Tang, Oluwaniyi Mabayoje et al.
- [Dissolution Stability of Photoanodes and Co-Catalysts in Photoelectrochemical Water Splitting](#)
Julius Knöppel, Attila Kormanyos, Britta Mayerhöfer et al.



Your Lab in a Box!

The PAT-Tester-i-16: All you need for Battery Material Testing.

- ✓ All-in-One Solution with integrated Temperature Chamber!
- ✓ Cableless Connection for Battery Test Cells!
- ✓ Fully featured Multichannel Potentiostat / Galvanostat / EIS!

www.el-cell.com +49 40 79012-734 sales@el-cell.com

EL-CELL[®]
electrochemical test equipment





Co-Design of Multijunction Photoelectrochemical Devices for Unassisted CO₂ Reduction to Multicarbon Products

William J. Wei,^{1,2,*} Alex J. King,^{1,2,*} Justin C. Bui,^{1,2,*} Adam Z. Weber,^{2,**} and Alexis T. Bell^{1,2,z}

¹Department of Chemical and Biomolecular Engineering, University of California Berkeley, Berkeley, California 94720, United States of America

²Liquid Sunlight Alliance, Lawrence Berkeley National Laboratory, Berkeley, California 94720, United States of America

Photoelectrochemical (PEC) CO₂ reduction (PEC CO₂R) is a prospective approach for utilizing solar energy to synthesize a variety of carbon-containing chemicals and fuels, the most valuable of which are multicarbon (C₂₊) products, such as ethylene and ethanol. While these products can be produced with high faradaic efficiency using Cu, this occurs over a relatively narrow potential range, which, in turn, imposes constraints on the design of a device for PEC CO₂R. Herein, we used continuum-scale modeling to simulate the solar-to-C₂₊ (STC₂₊) efficiency of PEC CO₂R devices fed with CO₂-saturated, 0.1 M CsHCO₃. We then explored how cell architecture and the use of single or dual photoelectrode(s) alters the optimal combination of photoelectrode bandgaps for high STC₂₊ efficiency. Ultimately, this work provides guidance for the co-design of the device architecture and photoelectrode bandgaps required to achieve high STC₂₊ efficiency. The insights gained are then used to identify systems that yield the highest amount of C₂₊ products throughout the day and year.

© 2023 The Author(s). Published on behalf of The Electrochemical Society by IOP Publishing Limited. This is an open access article distributed under the terms of the Creative Commons Attribution 4.0 License (CC BY, <http://creativecommons.org/licenses/by/4.0/>), which permits unrestricted reuse of the work in any medium, provided the original work is properly cited. [DOI: 10.1149/1945-7111/ad10e7]



Manuscript submitted May 6, 2023; revised manuscript received October 24, 2023. Published December 8, 2023. *This paper is part of the JES Focus Issue on Multiscale Modeling, Simulation and Design: In Honor of Ralph E. White.*

Supplementary material for this article is available [online](#)

Solar-driven CO₂ reduction (CO₂R) holds great promise for the sustainable generation of energy-dense fuels and chemicals.^{1,2} Multicarbon (C₂₊) products (e.g., ethylene, ethanol, propanol) are particularly attractive because they have a large market size and can be further converted to higher molecular-weight hydrocarbon fuels that have high volumetric and mass energy densities.^{1,3} The majority of the effort on CO₂R has been directed towards the coupling of a photovoltaic (PV) energy source to an electrochemical (EC) cell, what is referred to as a PV+EC system.^{4–7} However, there is now increasing interest in developing systems for photoelectrochemical CO₂ reduction (PEC CO₂R), in which single or multiple photo-absorbers are coupled directly to catalysts at the cathode and/or anode of the device.^{8–10}

A metric for the performance of a solar-driven CO₂R system is its solar-to-C₂₊ (STC₂₊) efficiency, which gives the percentage of incident solar power used to transform CO₂ into C₂₊ products. State of the art PV+EC systems have demonstrated STC₂₊ efficiencies of ~3.8%.¹¹ For PEC CO₂R, the current state of the art is ~1% STC₂₊ efficiency, but this is achieved with a PEC device assisted externally by a PV stack.⁸ To the best of our knowledge, C₂₊ product generation in a monolithic, unassisted PEC device has not been demonstrated. The principal challenge to achieving high STC₂₊ efficiency in monolithic, unassisted PEC devices is the large potential required to generate C₂₊ products.¹² There are also additional challenges in the development of monolithic, unassisted PEC CO₂R devices, such as multijunction photoelectrode fabrication,¹³ catalyst and photo-absorber stability,^{14,15} device design,¹⁶ and others.⁷ But, to overcome these challenges, it is important to first understand how to design PEC systems that operate at potentials large enough for C₂₊ product formation.

Metallic copper (Cu) has the unique ability to catalyze CO₂ to C₂₊ products with high faradaic efficiency (FE),^{12,17} which is the reason that it has been used in photocathodes developed for CO₂R.^{8,10,18–20} However, the product distribution of CO₂R on Cu is potential-dependent.^{8,12} The C₂₊ current density increases with cathode

potential and reaches a peak value, beyond which it decreases due to the low CO₂ concentration at the Cu surface and a shift to CH₄ and H₂ formation because of their high transfer coefficients.²¹ To achieve selective PEC CO₂R to C₂₊ products, the potential at which C₂₊ product formation is maximized, referred to here as V_{id} (ideal operating potential), needs to be attained. The operating potential (V_{op}) is determined by the intersection between the electrochemical (EC) and PV current-voltage curves (also called the electrochemical load and solar power curves). Operation of the overall device at V_{id} can be achieved with a photoelectrode that has a high photocurrent and a photovoltage near V_{id} (e.g., this photoelectrode performance enables intersection of the EC load and solar power curves near V_{id}; see Fig. S1).²¹

The photocurrent and photovoltage are related to the bandgaps of the semiconductors (V_g).²² A photoelectrode with two or three semiconductors of different bandgaps in a multijunction stack is required to obtain a sufficiently high photovoltage for operation near V_{id}. The photovoltage and photocurrent depend on the exact combination of bandgaps in the multijunction stack.²³ Moreover, V_{id} will depend on the overall device architecture because the overpotentials due to kinetics, species transport, and ohmic losses shift the potential at which C₂₊ products are maximized. Thus, there is a need to co-design the photoelectrode bandgaps and device architecture in order to ensure the photovoltage and photocurrent from the photoelectrode enables operation at the device-specific V_{id}. Understanding this relationship is critical for the design of monolithic, unassisted PEC CO₂R systems which yields high STC₂₊ efficiency. Previous efforts aimed at modeling PEC CO₂R systems have focused on engineering either the photoelectrode bandgap or the device architecture for generating C₂₊ products.^{24–27} Singh et al.²⁴ and Kalamaras et al.²⁵ investigated the theoretical maximum STC₂₊ efficiency as a function of bandgap, assuming that there are no potential losses and that the reactions occur at their thermodynamic equilibrium potentials. Chen et al.²⁶ have incorporated kinetic, ohmic, and transport potential losses, but only considered a single device architecture. Gutierrez et al.²⁷ considered alternative cell designs but did not account for reactant transport, a critical consideration for CO₂R,²⁸ and did not investigate the effect of photoelectrode bandgaps on STC₂₊ efficiency.

In this study, we report a model for PEC CO₂R that accounts for photo-absorber performance, anodic and cathodic catalyst kinetics,

^zThese authors contributed equally to this work.

*Electrochemical Society Student Member.

**Electrochemical Society Fellow.

^zE-mail: akingmi@lbl.gov; azweber@lbl.gov; alexbell@berkeley.edu

diffusion and migration through the boundary layers and anionic membrane, and ionic conduction through the bulk electrolyte. The simulations carried out using this model provide guidelines for the co-design of photoelectrode bandgaps and overall device design for achieving high STC_{2+} efficiency. The model is then used to understand how the optimal bandgaps for generating C_{2+} products change between single versus dual photoelectrode designs, and which bandgap combinations lead to high STC_{2+} efficiency throughout the day and year. These efforts enable identification of photoelectrode materials and device architectures for achieving high C_{2+} product generation in a monolithic, unassisted PEC CO_2R device.

Figure 1a shows a schematic of the modeled system. This cell design is similar to that used for electrochemical studies of CO_2R ^{1,12,29,30} and is currently being explored for PEC CO_2R .^{4,8,10} A liquid-based system is chosen because PEC devices are typically limited to current densities of $\sim 10 \text{ mA cm}^{-2}$ by the generated photocurrent, assuming no concentration of solar radiation. We note that the general findings (e.g., requirement of co-design between device architecture and photoelectrode bandgaps) are expected to be applicable to both liquid-based and gaseous-fed PEC devices. A photocathode (performing CO_2R) and a dark anode (performing the oxygen evolution reaction (OER)) are separated by a Sustainion[®] anion-exchange membrane. The electrolyte is chosen to be 0.1 M CsHCO_3 , because it promotes the formation of C_{2+} products,^{28,31–33} and is assumed to be saturated with CO_2 (34 mM). Ion transport through the bulk electrolyte is modeled using Ohm's law with an assumed constant concentration of 0.1 M CsHCO_3 since it is continuously replenished by feeding fresh electrolyte. CO_2 diffuses from the bulk electrolyte to the photocathode surface through a mass-transport boundary layer and reacts on the Cu foil to form C_1 (HCOOH , CO , and CH_4) and C_{2+} (C_2H_4 , EtOH , PrOH , and AllylOH) products. Water reduction to form H_2 is included as a competing cathodic reaction. The rate at which these products are formed is represented by a Tafel equation (kinetic parameters are listed in Table S3). The mass-transport boundary-layer thickness was assumed to be $100 \mu\text{m}$, which corresponds to a Reynolds number of 1800. The impact of electrolyte flowrate (e.g., Reynolds number) on STC_{2+} rates has been reported by King et al.²¹ Mass-transport boundary layers also exist near the membrane surfaces and the anode surface due to ion concentration gradients between the bulk electrolyte and the electrolyte layer adjacent to the membrane (a result from Donnan exclusion) and due to the generation of protons at the anode, respectively. The PV solar power curve is simulated by the ideal-diode equation in the Shockley–Queisser radiative efficiency limit under AM 1.5 G solar illumination.²² This PV model is coupled to the electrochemical model by determining the intersection point of the electrochemical load and solar power curves. MATLAB is used to identify the operating point and to model the PV component, and COMSOL 6.0 software is used to solve the electrochemical governing equations (details provided in SI sections S2 and S3).

The photo-absorber is assumed to operate at its theoretical maximum (Shockley–Queisser radiative efficiency limit under AM1.5 G solar illumination). This assumption enables the prediction of the STC_{2+} efficiency from a PEC CO_2R device based on the photoelectrode(s) bandgap combination; however, it neglects detailed semiconductor and multijunction properties (i.e., lattice matching, surface defects, semiconductor doping) and their effects on the photo-absorber performance. Moreover, the model excludes the impact of a protective metal-oxide interlayer between the semiconductor and the metal catalyst.³⁴ In general, the STC_{2+} efficiency will decrease whenever non-idealities in the photo-absorber are present. While the effects of these properties on the performance of photoelectrodes are important, they are outside the scope of the current study and their qualitative impact has been previously discussed.²¹ Instead, we assume that the multijunction stack operates at its theoretical maximum and focus on identifying

candidate semiconducting materials for achieving high STC_{2+} efficiency, as well as how the semiconductor bandgaps change with device architecture. The exclusion of the aforementioned properties is not expected to limit the applicability of the results.

The electrochemical behavior of the H-cell device is shown in Fig. 1b. The total current density (i_{tot}) increases monotonically with cell potential. We note the flat region in i_{tot} , which has been observed experimentally,^{4,5,8,9} is due to the appreciable OER rate that rapidly increases the proton concentration at the anode surface (see Fig. S2) and shifts the anodic equilibrium potential due to the Nernstian pH-shift. Importantly, the C_{2+} partial current density (PCD) does not increase monotonically but exhibits a peak at $\sim 4.25 \text{ V}$ (e.g., $V_{\text{id}} \approx 4.25 \text{ V}$), after which the C_{2+} PCD decreases (see Fig. S3 for all products). We note that this maximum in C_{2+} PCD can be increased by altering the Cu morphology,³⁵ tuning the electrolyte microenvironment (e.g., pH and CO_2 concentration) adjacent to the Cu surface,²⁸ and increasing the rate of CO_2 transport (e.g., increasing electrolyte flowrate).²¹ Employing these strategies along with solar-concentrators to increase the photocurrent density could enable STC_{2+} PCDs that are beyond the maximum C_{2+} PCD in this study ($\sim 10 \text{ mA cm}^{-2}$). To operate at the point of maximum C_{2+} generation, the electrochemical load and solar power curves must intersect at V_{id} . For a given device design, the operating point can be controlled by the bandgap combination of the multijunction photoelectrode, as shown in Fig. 1c. $V_{\text{g}}^{\text{top}}$, $V_{\text{g}}^{\text{mid}}$, and $V_{\text{g}}^{\text{bot}}$ correspond to the bandgap of the semiconductor at the top, middle, and bottom of the triple-junction stack, respectively. Shown here is the operating point for systems with a high (pink) and a low (teal) C_{2+} PCD. The teal system has a low C_{2+} PCD because it operates at a potential much lower than V_{id} , despite its photovoltage being greater than V_{id} . This low operating potential is due to the relatively low photocurrent, which causes the electrochemical load and solar power curves to intersect at a potential much lower than V_{id} . The pink system, on the other hand, has a relatively high photocurrent and is able to operate close to V_{id} , even though the photovoltage is less than V_{id} . Thus, the generation of C_{2+} products strongly depend on the photocurrent and photovoltage, both of which are controlled by the photoelectrode bandgaps.

By sweeping through an array of bandgap combinations for a triple-junction photoelectrode, a PEC operating point can be determined for each bandgap combination. The corresponding STC_{2+} efficiency for each photoelectrode can be visualized with slices of a 3D surface plot as shown in Fig. 1d. The STC_{2+} efficiency for the two operating points in Fig. 1c is labeled by circles in Fig. 1d. The z -axis slices in Fig. 1d correspond approximately to Si (1.11 eV), InP (1.27 eV), GaAs (1.43 eV), and CdSe (1.74 eV). This analysis reveals that there exists an optimal bandgap selection around $V_{\text{g,top}} = 2.10$, $V_{\text{g,mid}} = 1.65$, and $V_{\text{g,bot}} = 1.25 \text{ eV}$ for generating C_{2+} products (the pink line in Fig. 1c) in which the photocurrent is high and the photovoltage is near V_{id} . This optimal combination correlates to an $\text{AlAs}/\text{In}_{0.6}\text{Ga}_{0.4}\text{P}/\text{In}_{0.2}\text{Ga}_{0.8}\text{As}$ multijunction photocathode, which uses semiconductors that have approximately the same lattice constant.³⁶ These materials have been used in previous studies on multijunction PVs but, to our knowledge, have not been used together.^{36,37} A discussion on the use of more economical semiconducting materials (e.g., Si) is provided in SI section S4. For this device architecture, no bandgap combination leads to operation at V_{id} because the device is either limited by the photovoltage or photocurrent. Therefore, modification to the cell design is needed to reduce V_{id} and operate at the peak C_{2+} product generation, further highlighting that co-design of both the photoelectrode and the device architecture will be required to achieve optimal unassisted PEC CO_2R to C_{2+} products.

A voltage breakdown analysis, shown in Fig. 1e,²⁸ identifies the EC potential losses. The catholyte and anolyte chambers exhibit high ohmic potential losses ($< 1 \text{ V}$) due to their appreciable length (1.5 cm) and the low conductivity of the fed electrolyte. An effective approach for decreasing this potential loss is to reduce the distance

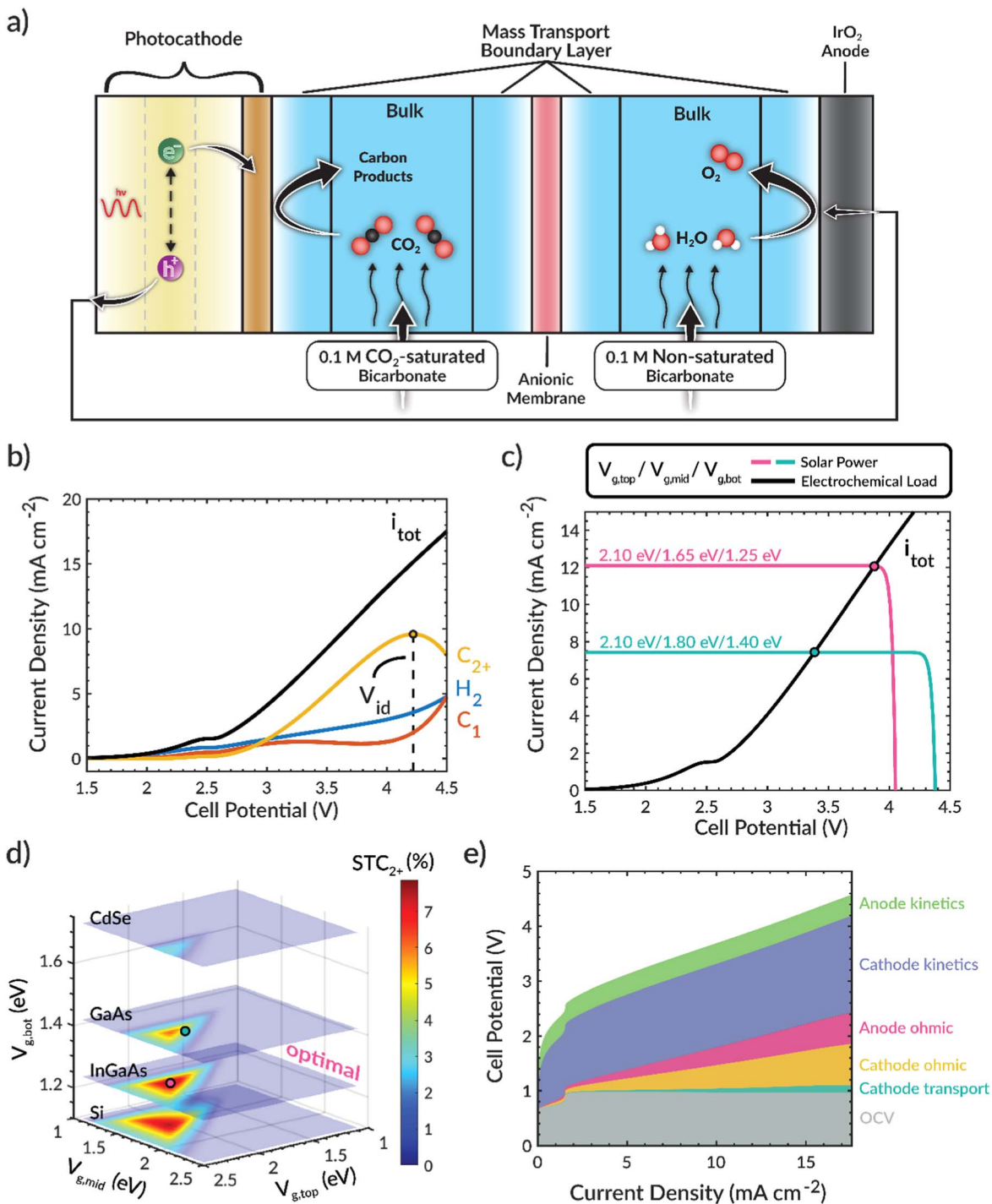


Figure 1. (a) Schematic of model system. (b) Total and partial current densities as a function of cell potential. (c) PEC operating points for two different photoelectrode bandgap combinations. (d) Solar-to-C₂₊ (STC₂₊) efficiency (color map) as a function of the top, middle, and bottom bandgap. (e) Electrochemical voltage breakdown versus total current density.

between the electrodes, as has been demonstrated for PEC water splitting devices.³⁸ Accordingly, we showed that the anodic ohmic loss can be reduced by eliminating the anodic chamber entirely (Fig. 2a), resulting in setup similar to that of a zero-gap membrane-electrode-assembly (MEA).^{39,40} We refer to this assembly as a half-MEA design because there is no electrolyte gap between the anode and the membrane. In this case, water from the catholyte must transport through the anion-exchange membrane to the anode catalyst layer where it is oxidized to O₂, which then diffuses out the back of the porous anode. We note that the model neglects O₂ transport because it is assumed to exit the system rapidly and that the

porous anode is approximated as a planar surface because of the low operating current densities (<20 mA cm⁻²). Moreover, water transport does not limit the performance of this device because the activity of liquid water is high (~1.0).³⁹ Significantly, this modification removes ~450 mV of potential loss from the cell (Fig. S4) and decreases the V_{id} by the same amount (Fig. 2b). The decrease in V_{id} causes a concomitant decrease in the optimal bandgap selection ($V_{g,top} = 2.00$, $V_{g,mid} = 1.55$, and $V_{g,bot} = 1.15$ eV), as seen in Fig. 2c, in order for the photovoltage to remain near V_{id} . This optimal bandgap combination corresponds approximately to an In_{0.4}Al_{0.6}As/In_{0.8}Ga_{0.2}P/Si multijunction photocathode^{36,41} and leads

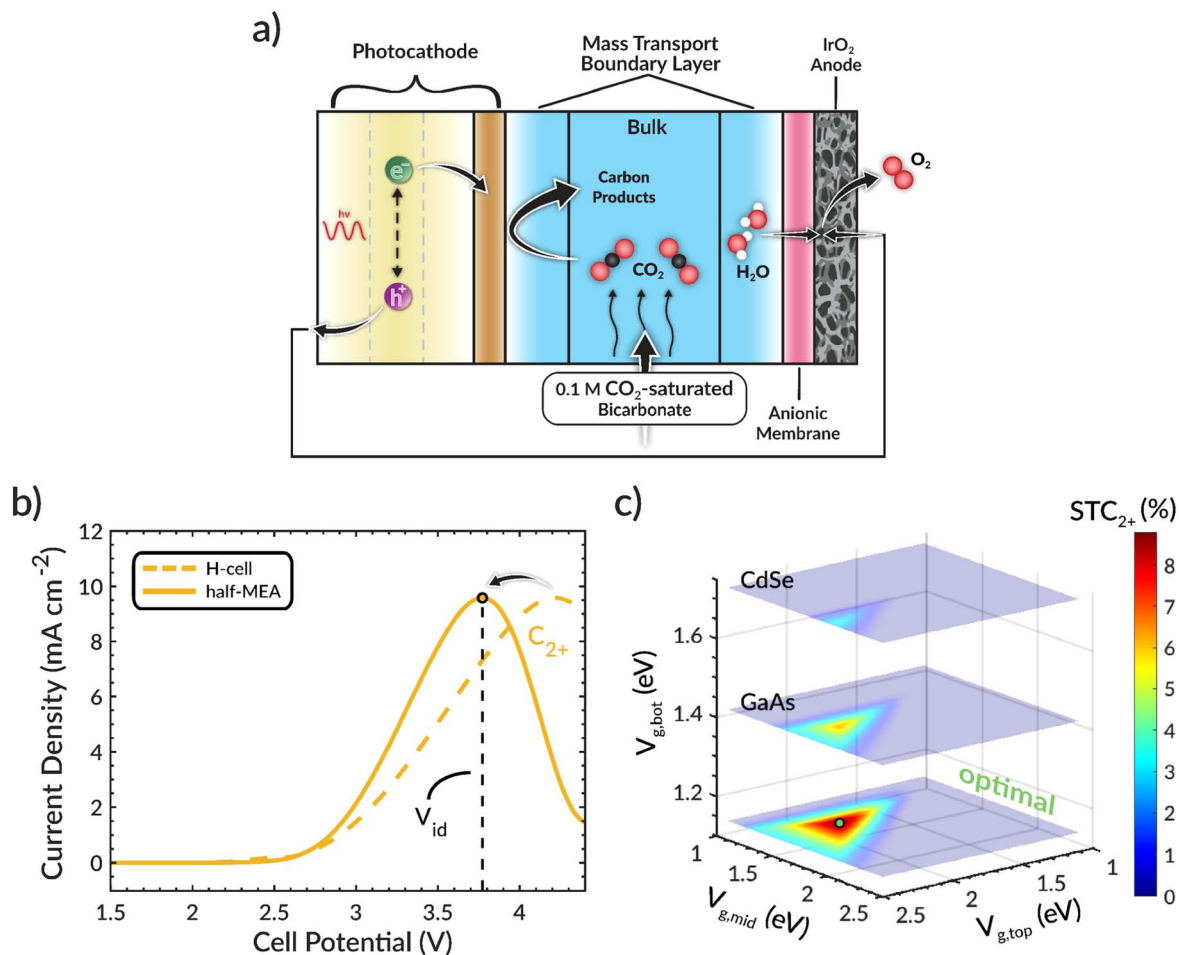


Figure 2. (a) Schematic of the half-MEA model system without anodic chamber. (b) Total and C_{2+} partial current densities as a function of cell potential. (c) Solar-to- C_{2+} (STC_{2+}) efficiency (color map) as a function of the top, middle, and bottom bandgap.

to a STC_{2+} efficiency of $\sim 9\%$, whereas $\sim 8\%$ STC_{2+} efficiency is obtained if the original optimal bandgap combination is used because it does not operate as close to V_{id} . Therefore, to obtain the maximum STC_{2+} efficiency, it is important to first optimize the device design and then identify the V_{id} for that design so that it can be paired with the right combination of semiconductors. In other words, to generate high yields of C_{2+} products, the selection of the photocathode materials and the design of the reactor are intimately coupled.

Another important design consideration is the use of single or dual photoelectrode(s). Dual photoelectrode devices can generate higher photovoltages and photocurrents for a given bandgap combination than a device with only a single photoelectrode.^{42,43} However, this requires illumination of both photoelectrodes.^{16,26} A simple design in which mirrors are used to reflect Sunlight toward the photoanode is depicted in Figs. 3a, 3b. This architecture can be used to evaluate how the optimal bandgap combination changes between single vs dual photoelectrode designs. Figures 3a, 3b also presents the STC_{2+} efficiency as a function of bandgaps, where the x -axis is the top photocathode bandgap ($V_{g,top}^c$), the y -axis is the photoanode bandgap (V_g^a), and the z -axis is the bottom photocathode bandgap ($V_{g,bot}^c$). The z -axis slices shown here correspond approximately to Si (1.11 eV), GaAs (1.43 eV), and CdSe (1.74 eV) (see Fig. S5). 2D surface plots at z -axis bandgaps corresponding to Si, GaAs, and CdSe are shown in Fig. S6.

It is evident that a dual photoelectrode device significantly increases the number of bandgaps capable of achieving high STC_{2+} efficiency compared to that for a single photoelectrode

device. This is because, for both the H-cell and half-MEA cell designs (Figs. 1d and 2c, a higher photovoltage and photocurrent for a given bandgap combination can be achieved using a dual photoelectrode system. We note that this finding is consistent with those previously seen in PEC water-splitting systems.^{44,45} Thus, the use of dual photoelectrodes enables greater flexibility in the choice of semiconductors and their photo-absorption characteristics than a single photoelectrode design. We note that the half-MEA dual photoelectrode design exhibits a region of lower STC_{2+} at the optimal bottom photocathode bandgap (1.15 eV), seen in Fig. 3b, because the operating potential exceeds V_{id} (see Fig. S7), leading to lower STC_{2+} efficiency. Although dual photoelectrodes broaden the selection of bandgaps that lead to high STC_{2+} efficiency, there still exists an optimal bandgap combination that maximizes STC_{2+} efficiency. This occurs for approximately $V_{g,top}^c = 2.00$ eV, $V_{g,bot}^c = 1.40$ eV, and $V^a = 2.00$ eV for the H-cell cell design (Fig. 3a) and $V_{g,top}^c = 1.80$ eV, $V_{g,bot}^c = 1.15$ eV, and $V^a = 1.75$ eV for the half-MEA cell design (Fig. 3b). Interestingly, the decrease in optimal bandgaps for each semiconductor upon removing potential losses is approximately the same (~ 0.25 eV). This is also seen in the single photoelectrode design, in which each bandgap in the multijunction stack decreases by ~ 0.1 eV (note that this value is less than 0.25 eV because of the smaller shift in operating potential, as discussed in SI section S11). The decrease in optimal bandgaps, for the single and dual photoelectrode systems, corresponds to a decrease of ~ 0.55 eV per volt of change in operating potential (e.g., a 1 V decrease in operating potential is achieved by decreasing each bandgap by ~ 0.55 eV). This analysis shows that there is a relationship between the optimal operating potential and the optimal bandgaps that is

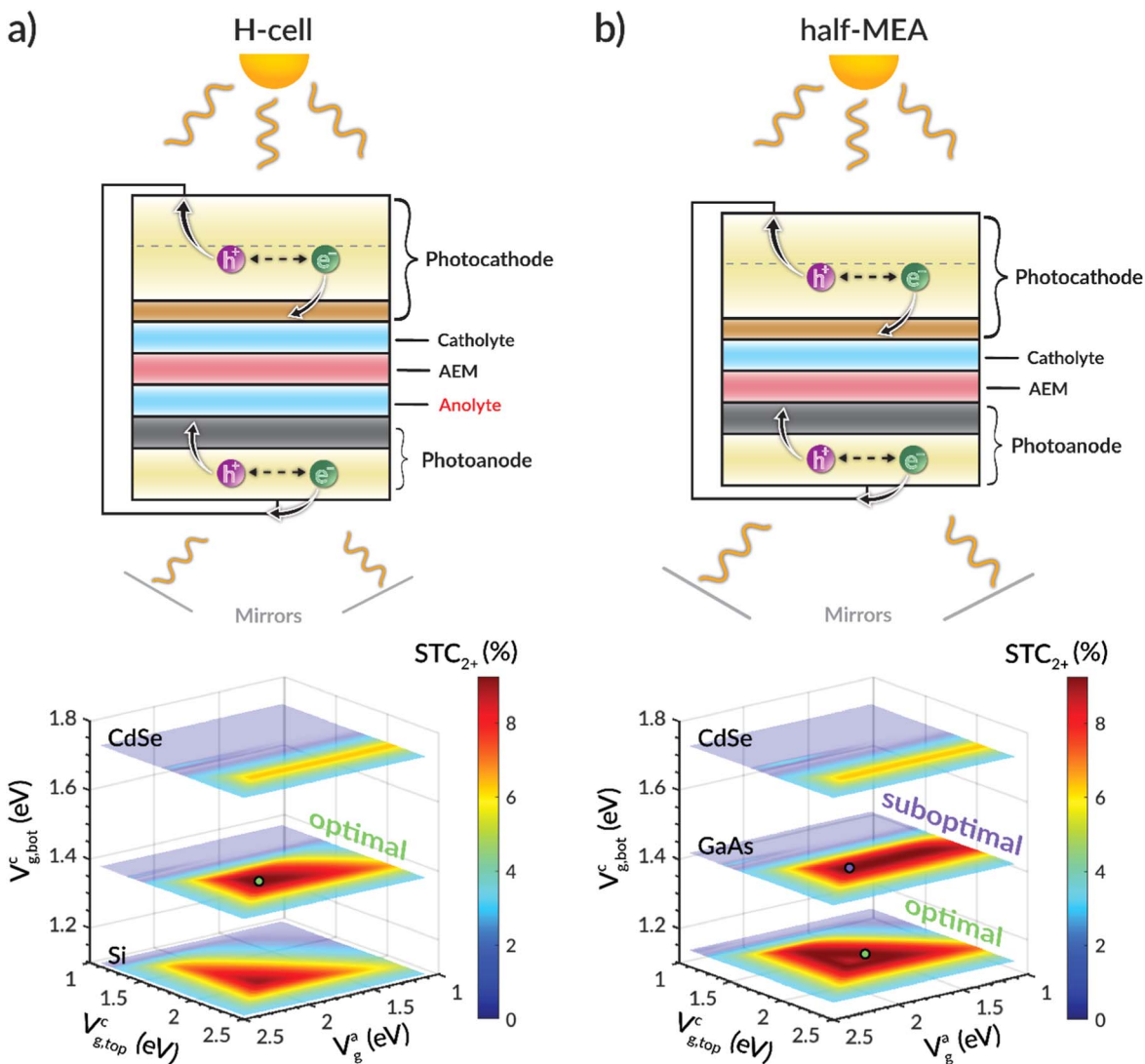


Figure 3. (a) Schematic of the H-cell model system with dual photoelectrodes and corresponding solar-to- C_{2+} (STC_{2+}) efficiency (color map) as a function of the top cathode, bottom cathode, and anode bandgap, where the top, middle and bottom z-slices are CdSe, the optimal bandgap, and Si respectively. (b) Schematic of the half-MEA model system with dual photoelectrodes and corresponding solar-to- C_{2+} (STC_{2+}) efficiency (color map) as a function of the top cathode, bottom cathode, and anode bandgap, where the top, middle and bottom z-slices are CdSe, GaAs, and the optimal bandgap respectively.

dictated by the device architecture and choice of photoelectrode materials, respectively. We note this decrease in optimal bandgap with reduced potential losses is also seen in PEC water splitting,⁴³ but this ~ 0.55 eV per volt for each bandgap is unique to CO_2 -saturated liquid-flow cells for which the operating potential is ≥ 2.5 V; otherwise, a peak C_{2+} PCD versus potential would not exist and only two semiconductors would be needed to achieve sufficient photovoltage to drive electrolysis. This observed relationship applies to commonly used device architectures and operating conditions, and it provides a useful guide for choosing the bandgaps of the photoelectrode(s) required to achieve operation at V_{id} and the maximum STC_{2+} efficiency. These results also underscore the need to co-design both the device architecture and photoelectrode (s) in order to achieve optimal PEC CO_2R to C_{2+} products.

Another useful guideline drawn from this analysis is the choice of bandgaps for the photoanode and the semiconductor at the top of the photocathode. For both the H-cell and half-MEA cell designs, the optimal bandgap combination for a dual photoelectrode system occurs when the photoanode and top photocathode have approximately the same bandgap. This combination maximizes the photocurrent and generates a large photovoltage, and then the bottom photocathode bandgap is tuned so that the photovoltage is near V_{id} .

Figure S8 shows how the photocurrent and photovoltage change as a function of each bandgap. These results suggest that, for CO_2 -saturated liquid-flow cells in which the operating potential is ≥ 2.5 V, the photoanode and top photocathode should be the same semiconducting material in order to achieve high STC_{2+} efficiency.

The importance of maximizing photocurrent can be easily seen from an analysis of the STC_{2+} efficiency throughout the diurnal cycle. Using AM 1.5 G solar illumination and approximating the hourly solar concentration from the NREL's TMY3 dataset for the summer and winter solstices in 2020 at Barstow, CA (see Fig. S9), we estimate the hourly STC_{2+} efficiency, photocurrent, and operating potential for the optimal ($V_{g,top}^c = 1.80$, $V_{g,bot}^c = 1.15$, and $V^a = 1.75$ eV) and a suboptimal ($V_{g,top}^c = 1.90$, $V_{g,bot}^c = 1.45$, and $V^a = 1.90$ eV) combination of semiconductor bandgaps for the half-MEA dual photoelectrode design (these bandgap combinations are shown as circles in Fig. 3b). We note this analysis assumes that the PEC device operates at ambient temperature. The impact of temperature needs to be investigated in the future, but such work needs information about temperature-dependent properties (recombination current, species transport coefficients, material properties, kinetic rate coefficients, etc), which in many cases are not currently known. Figure 4a shows that the hourly STC_{2+} efficiency between

the optimal and suboptimal bandgap selections follows the same trend and that they are similar during the midday summer solstice. Near Sunrise and Sunset and during the winter solstice, however, the optimal bandgap combination has higher STC_{2+} efficiency than the suboptimal combination. This behavior is also evident from 3D surface plots presented in Fig. S10. The higher STC_{2+} efficiency when the solar intensity is less than one Sun is due to the high photocurrent of the optimal bandgap combination (Fig. 4b). Because the optimal bandgap selection is able to absorb more of the solar spectrum and generate a higher photocurrent than the suboptimal combination, its solar power curve can intersect the electrochemical load curve at a potential that is closer to V_{id} throughout the diurnal and annual solar cycles (Fig. 4c). In summary, the optimal bandgap combination generates $1.13\times$ and $1.39\times$ more C_{2+} products throughout the day during the summer and winter solstices, respectively, compared to the suboptimal bandgap combination. Throughout the entire year of 2020, the optimal bandgap combination generates $1.21\times$ more C_{2+} products than the suboptimal combination. While there are various combinations of bandgaps leading to high STC_{2+} efficiency at full Sun (Fig. 3), only the bandgaps that lead to a high photocurrent and photovoltage near V_{id} result in the highest C_{2+} generation throughout the day and year.

This study demonstrates the use of multiphysics modeling to simulate the electrochemical and photovoltaic performance of PEC CO_2R systems involving a Cu catalyst for the unassisted generation of multicarbon (C_{2+}) products. The C_{2+} product selectivity is potential dependent, which necessitates simultaneous co-design of photoelectrode(s) and system architecture in order to ensure that the overall device operates at the potential for which C_{2+} production is maximized (V_{id}). The model is used to identify optimal device designs and photoelectrode bandgaps, as well as general relationships between the two, for achieving high solar-to- C_{2+} (STC_{2+}) efficiencies. The simulation shows that the optimal combination of bandgaps in a triple-junction stack for achieving high STC_{2+} efficiency strongly depends on the device architecture. Modifying the device design changes V_{id} , which necessitates a change in photoelectrode bandgaps in order to ensure the system operates at the new V_{id} and maintains high STC_{2+} efficiency. Specifically, the model predicts that the bandgap of each layer in the photoelectrode stack should change by approximately the same amount, the amount depending on how much V_{id} changes. This finding holds for both single photoelectrode and dual photoelectrode systems. In addition to these results, the model predicts that photoelectrodes with a high photocurrent and a photovoltage near V_{id} are required to achieve

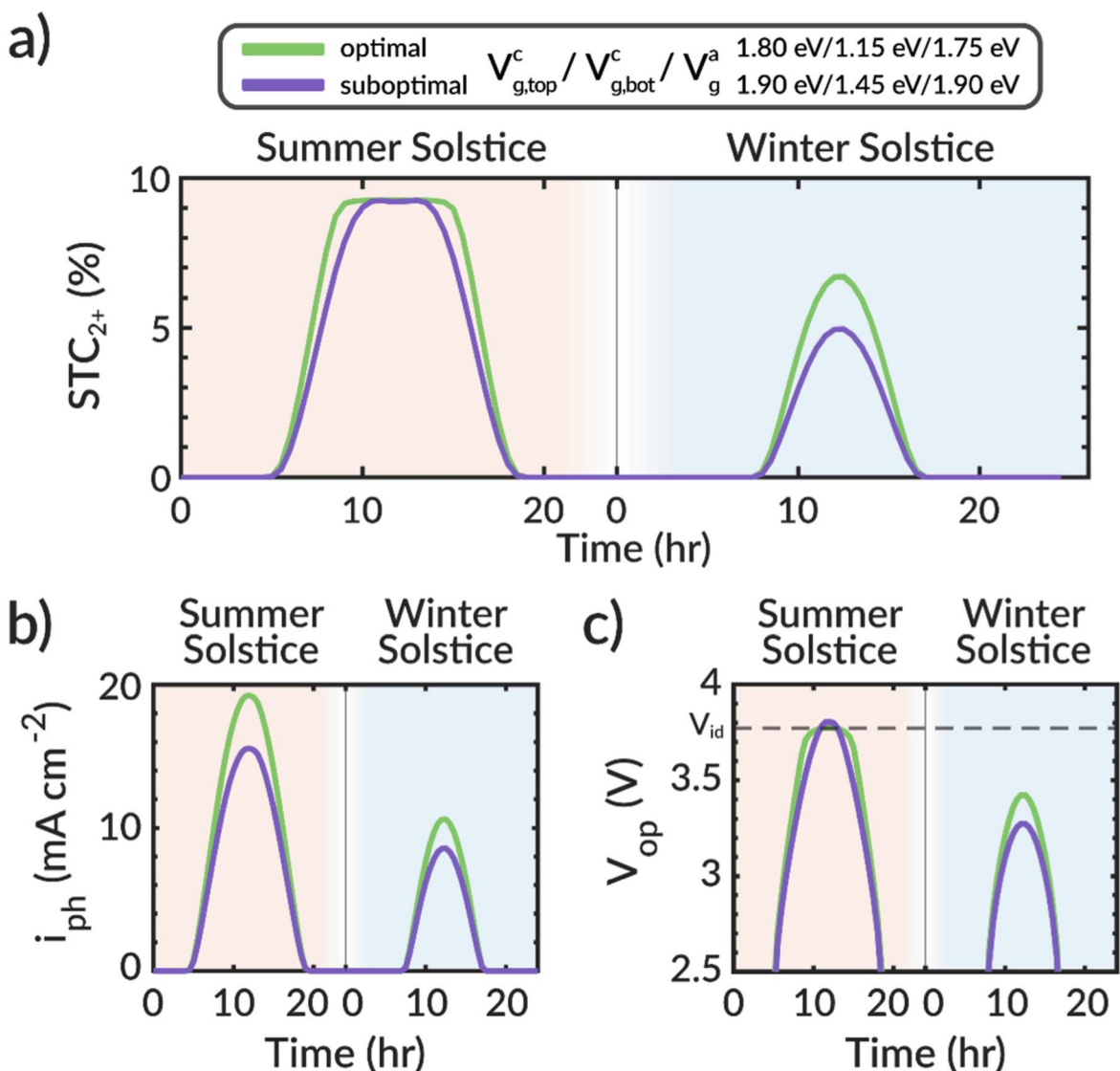


Figure 4. (a) Solar-to- C_{2+} (STC_{2+}) efficiencies of optimal and suboptimal bandgap combinations for Barstow, CA during the 2020 summer and winter solstices. (b) Photocurrent density (i_{ph}) of optimal and suboptimal bandgap combinations for Barstow, CA during the 2020 summer and winter solstices. (c) Operating potential (V_{op}) of optimal and suboptimal bandgap combinations during the 2020 summer and winter solstices in Barstow, CA.

high STC_{2+} efficiency throughout the diurnal and annual solar cycles. This is because high photocurrents enable the solar power curve to intersect the electrochemical load curve at potentials close to V_{id} even when the incident solar power is $<100 \text{ mW cm}^{-2}$. To achieve high photocurrents, it is desired to employ semiconductors with moderate bandgaps, but the exact combination of these bandgaps must generate a photovoltage near V_{id} to attain high STC_{2+} efficiency. To improve these predictions, future model developments should include non-idealities and series and shunt resistances in the diode equation to account for detailed semiconductor, multijunction, and solid-state interfacial properties because they will impact the ultimate performance of each photoelectrode. While the present analysis neglects these phenomena, it nevertheless reveals useful relationships between photoelectrode bandgaps and device design for achieving high STC_{2+} efficiency throughout the day, a key step in the development of monolithic, unassisted PEC CO_2R systems.

Acknowledgments

This material is based on work performed within the Liquid Sunlight Alliance, which is supported by the U.S. Department of Energy, Office of Science, Office of Basic Energy Sciences, Fuels from Sunlight Hub under Award Number DE-SC0021266. A.J.K. acknowledges funding from the National Science Foundation Graduate Research Fellowship under Grant No. DGE 2146752. J.C.B. would like to acknowledge support from the National Defense Science and Engineering Graduate Fellowship (NDSEG) supported by the Army Research Office (ARO). The authors acknowledge the National Institute of Health under Grant No. S10OD034382 for its support of the Molecular Graphics and Computation Facility in the College of Chemistry at the University of California, Berkeley. The authors also thank Aditya Prajapati for his help in resolving convergence issues and Kyra Yap for her constructive feedback on the paper.

ORCID

William J. Wei  <https://orcid.org/0000-0001-8187-5041>
 Alex J. King  <https://orcid.org/0000-0002-3156-1607>
 Justin C. Bui  <https://orcid.org/0000-0003-4525-957X>
 Adam Z. Weber  <https://orcid.org/0000-0002-7749-1624>

References

1. S. Nitopi et al., *Chem. Rev.*, **119**, 7610 (2019).
2. J. C. Bui et al., *Chem. Rev.*, **122**, 11022 (2022).
3. S. Verma, B. Kim, H. R. M. Jhong, S. Ma, and P. J. A. Kenis, *ChemSusChem*, **9**, 1972 (2016).
4. Y. Wang, J. Liu, Y. Wang, Y. Wang, and G. Zheng, *Nat. Commun.*, **9**, 1 (2018).
5. Y. Xiao et al., *J Mater Chem A Mater*, **8**, 18310 (2020).
6. A. Prajapati et al., *Cell Reports Physical Science*, **3**, 1 (2022).
7. C. E. Creissen and M. Fontecave, *Adv. Energy Mater.*, **11**, 2002652 (2020).
8. Gurudayal et al., *Energy Environ. Sci.*, **12**, 1068 (2019).
9. B. Liu et al., *Nat. Commun.*, **13**, 1 (2022).
10. C. Kim et al., *Energy Environ. Sci.*, **16**, 2968 (2023).
11. Z. Chen et al., *J. Am. Chem. Soc.*, **142**, 6878 (2020).
12. K. P. Kuhl, E. R. Cave, D. N. Abram, and T. F. Jaramillo, *Energy Environ. Sci.*, **5**, 7050 (2012).
13. E. A. Polyzoeva, *Tradeoffs of the Use of SiGe Buffer Layers in Tandem GaAsP/Si Solar Cells*, Massachusetts Institute of Technology (2004).
14. R. V. Mom et al., *ACS Appl. Mater. Interfaces*, **15**, 30052 (2023).
15. F. Nandjou and S. Haussener, *ACS Applied Materials & Interfaces*, **14**, A (2022).
16. S. Haussener et al., *Energy Environ. Sci.*, **5**, 9922 (2012).
17. D. Corral et al., *Chem Catalysis*, **2**, 3239 (2022).
18. K. Wang et al., *Green Chem.*, **23**, 3207 (2021).
19. I. Roh et al., *J. Am. Chem. Soc.*, **144**, 8002 (2022).
20. K. M. K. Yap et al., *Chem Catalysis*, **3**, 1 (2023).
21. A. J. King, J. C. Bui, A. T. Bell, and A. Z. Weber, *ACS Energy Lett.*, **7**, 2694 (2022).
22. W. Shockley and H. J. Queisser, *J. Appl. Phys.*, **32**, 510 (1961).
23. S. M. Sze and K. K. Ng, *Physics of Semiconductor Devices* (John Wiley, New York) 3rd ed., p. 1 (1981).
24. M. R. Singh, E. L. Clark, and A. T. Bell, *Proc Natl Acad Sci U S A*, **17**, 18924 (2015).
25. E. Kalamaras, H. Wang, M. Mercedes Maroto-Valer, J. M. Andresen, and J. Xuan, *ChemPhysChem*, **21**, 232 (2020).
26. Y. Chen, N. S. Lewis, and C. Xiang, *ACS Energy Lett.*, **1**, 273 (2016).
27. R. R. Gutierrez and S. Haussener, *J. Electrochem. Soc.*, **163**, H1008 (2016).
28. J. C. Bui et al., *Acc. Chem. Res.*, **55**, 484 (2021).
29. C. Kim, L. C. Weng, and A. T. Bell, *ACS Catal.*, **10**, 12403 (2020).
30. C. Kim et al., *Nat. Energy*, **6**, 1026 (2021).
31. S. Ringe et al., *Energy Environ. Sci.*, **12**, 3001 (2019).
32. M. R. Singh, Y. Kwon, Y. Lum, J. W. Ager, and A. T. Bell, *J. Am. Chem. Soc.*, **138**, 13006 (2016).
33. J. Resasco, Y. Lum, E. Clark, J. Z. Zeddon, and A. T. Bell, *Chem. Electro. Chem.*, **5**, 1064 (2018).
34. A. J. King, A. Z. Weber, and A. T. Bell, *ACS Appl. Mater. Interfaces*, **15**, 23024–23039 (2023).
35. J. A. Gauthier, J. H. Stenlid, F. Abild-Pedersen, M. Head-Gordon, and A. T. Bell, *ACS Energy Lett.*, **6**, 3252 (2021).
36. M. S. Leite et al., *Appl. Phys. Lett.*, **102**(3), 1 (2013).
37. A. C. Varonides, *Proceedings of the World Renewable Energy Congress—Sweden, 8–13 May, 2011, Linköping, Sweden*, **57**, 2767 (2011).
38. F. F. Abdi, R. R. Gutierrez Perez, and S. Haussener, *Sustain Energy Fuels*, **4**, 2734 (2020).
39. L. C. Weng, A. T. Bell, and A. Z. Weber, *Energy Environ. Sci.*, **12**, 1950 (2019).
40. L. Weng, A. T. Bell, and A. Z. Weber, *Energy Environ. Sci.*, **13**(10), 3592–3606 (2020).
41. R. Cariou et al., *Nat. Energy*, **3**, 326 (2018).
42. M. G. Walter et al., *Chem. Rev.*, **110**, 6446 (2010).
43. S. Hu, C. Xiang, S. Haussener, A. D. Berger, and N. S. Lewis, *Energy Environ. Sci.*, **6**, 2984 (2013).
44. C. Ding et al., *Phys. Chem. Chem. Phys.*, **16**, 15608 (2014).
45. B. Seger et al., *Energy Environ. Sci.*, **7**, 2397 (2014).