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BEAM LOSS AND SHIELDING EXPERIMENTS RELEVANT TO HIGH ENERGY PROTON **SYNCHROTRONS**

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Author

Gilbert, William S.

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University of California **Ernest O. Lawrence Radiation Laboratory**

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Introduction

Our primary concern has been with the radiation problems associated with the contemplated 200-GeV strong-focusing proton synchrotron. We have been involved in various experiments at the existing proton synchrotrons in the 30-GeV range at Brookhaven and CERN, since these can be considered as models for the higher energy machine.

There are several distinct problems that are associated with radiation that is generated through the cascade process after the primary beam is lost. The magnitude and distribution of this primary beam loss is therefore an input to each of the radiation problems listed below in Table I.

Table I. Radiation Problems

* Work done under the auspices of the U. S. Atomic Energy Commission.

Problems 1-4 relate to the radiation levels in buildings on site and to the off site general population. Problems $5-8$ influence the m aintainance of the accelerator and it is in this m aintenance that the operators and other radiation workers get the bulk of their radiation exposures.

In 1965-1966 it was agreed by those involved in shielding calculations at CERN, the Lawrence Radiation Laboratory, and the Rutherford High Energy Laboratory that a need existed for a largescale shielding experiment at the CERN-PS in which the radiation field in the earth shield and the primary beam loss distribution would be measured simultaneously. The specific accelerator projects for which these new data were desired are: the U. S. 200-GeV proton synchrotron, the improved CERN-PS, the CERN intersecting storage rings, and the European 300-GeV proton synchrotron.

The above three laboratories** collaborated in a series of shielding experiments at the CERN-PS from September 28 to November 28, 1966. We had exclusive use of the accelerator for eight 12-hour periods.

1966 Shielding Experiment at the CERN-PS

The use of activation detectors enabled us to determine the radiation field at hundreds of locations inside the machine tunnel and within the'earth shield. The response of these detectors is well understood, and spectral information is obtainable. We were able to cover a dynamic range of flux from ≤ 1 to $> 10^8$ neutrons cm⁻² sec⁻¹. Several counting facilities, with a correspondingly large staff, were required to count the many samples within the times dictated by the induced activities and associated decay lives.

The geometry of the experiment can be seen in Figs. 1-3. Some preliminary data have been presented $^{(1)}$ and a formal report is in preparation.

Beam Loss Distribution

In Fig. 4 is shown the beam loss distribution as measured by aLuminum activation detectors placed on the vacuum tank. There is considerable variation in this 1088 in the quiet region upstream of the target as well as subsequent minima downstream of the target.

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These minima are associated with the gaps between magnets. The pattern is surprisingly yell approximated by a constant term plus an exponential that decreases with distance downstream of the target. All the target associated term is accounted for in some half betatron wavelength downstream of the target and some 70 to 80% of the total beam loss is found under this peak.

In addition to a thin foil target, a cleanup collimator, or clipper, was used at a point several betatron wavelengths around the machine. Hereward, Ranft, and Richter (2) have calculated what might be the effects of such a clipper. Figures 5 and 6 show new Monte-Carlo calculations by Ranft (3) as to what the beam loss distribution might be expected to be for the machine operating conditions for our experiment, together with the experimental data. The agreement is fairly good and once again one can see the influence of the magnet-gap geometry in the experimental data. Ranft discusses how these calculations could be carried out for the 200 to 300 GeV accelerators once the various lattice and aperture parameters have been decided upon. For the CERN-PS, the clipper did not clean up the beam and reduce the random losses around the quiet sections as much as the calculations predicted. The reason for this is not clear.

Fitting of Experimental Data

A model based on the modified Moyer method has been programmed for the CDC-6600 to yield a best fit to the experimental data.^{40} In a given exposure, for a given·type detector, there way be from 10 to 80 data points at different locations. At each of these locations a flux is calculated, these fluxes are integrals of the following type.

Calculated flux =
$$
\Phi_i = c \int_{-\infty}^{+\infty} \frac{S(Z)\Theta(\theta)}{r^2} e^{-\ell_E/\lambda_{\text{Fe}}} e^{-\ell_E/\lambda_{\text{E}}}
$$
 dZ

Experimental flux $\equiv \phi_i$

 $\frac{1}{\epsilon}$, "

Variance = $\sum_{i} \left(\frac{\Phi_i - \Phi_i}{\Phi_i}\right)^2$

This is minimized by changing the variables below.

$$
S(Z) = 1 + a_1 e^{-a_3 (Z - Z_{target})}
$$

 $\Theta(\theta) = a_2 e^{-a_1 \theta}$

In the range of 90° \pm 30[°], this is $\frac{1}{2}$ in the lange of $\frac{1}{20}$, $\frac{1}{20}$, $\frac{1}{20}$ in $\frac{1}{21}$ and Bi fission detectors.

 $r =$ distance from elemental source to field point

 $\lambda_{\text{Fe}}, \lambda_{\text{Earth}}, \text{C}$ = other adjustable constants

The fit, as shown in Fig. 7, agrees with the experimental data to within about 20%, over 5 orders of magnitude in flux. The $\lambda_{Earth} = 117g$ cm⁻² for this particular set of experimental data.

Shield Estimates for 200 to 300 GeV and Improved CERN-PS Accelerators

Using the assumed circulating proton beams, loss distributions and allowable surface radiation levels of the original design reports (1) and the constants for the FLUX program derived from an analysis of the CERN shielding data, Table II represents an interim estimate of the required shielding above the magnets, if one assumes that the 200 and 300 GeV accelerator magnets are the same thickness and geometry as the CERN-PSmagnets. The FLUX program allows one to insert the actual magnet configuration and one can do so at a later time.

Table II. Shielding required above main ring, in g cm^{-2}

These currents represent the total beam lost in the accelerator, not necessarily the total Circulating beam.

References

-5-

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- 2. H. G. Hereward, J. Ranft and W. Richter, Efficiency of Multi-Traversal Targets, CERN 65-1, January 6, 1965.
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Figure Captions

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- CERN Shielding Experiment Plan View. Fig. $1.$ Fig. $2.$ CERN Shielding Experiment - Cross-Section View. CERN Shielding Experiment - Elevation View Geometry for $Fig. 3.$ FLUX Program. Fig. $4.$ Beam Loss Distribution - Experimental Data - Analytic Fit. Beam Loss Distribution Target - Ranft Monte-Carlo Calculations Fig. $5.$ Experimental Data. Beam Loss Distribution Clipper - Ranft Monte-Carlo Calculations $Fig. 6.$ Experimental Data.
- Fig. 7. Orbit Hole Aluminum Detector Data FLUX Program Fit.

CERN Shielding Experiment-Elevation

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Fig. 4

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Fig. 5

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Fig. 7

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1994 - Paul Barristo, político estatubator $\label{eq:2.1} \omega_{\rm{max}} = \frac{1}{2} \left(\frac{1}{2} \left(\frac{1}{2} \left(\frac{1}{2} \right) - \frac{1}{2} \left(\frac{1}{2} \right) \right) - \frac{1}{2} \left(\frac{1}{2} \left(\frac{1}{2} \right) - \frac{1}{2} \left(\frac{1}{2} \right) \right) \right)$ $\mathcal{L}(\mathcal{A})$ and $\mathcal{L}(\mathcal{A})$ and $\label{eq:2.1} \frac{1}{2} \left(\frac{1}{2} \right) \left(\frac{1}{2} \right) \left(\frac{1}{2} \right) \left(\frac{1}{2} \right) \left(\frac{1}{2} \right)$ $\mathcal{A}=\mathcal{A}^{\mathcal{A}}$ $\label{eq:2.1} \mathcal{L}_{\text{max}}(\mathcal{L}_{\text{max}}) = \mathcal{L}_{\text{max}}(\mathcal{L}_{\text{max}}) = \mathcal{L}_{\text{max}}(\mathcal{L}_{\text{max}})$ $\hat{p}(\hat{x}) = \hat{p}^2(\hat{x})$ $\mu_{\rm{eff}}$ and α and β ~ 10 $\hat{\mathcal{L}}_{\text{max}}$ and $\hat{\mathcal{L}}_{\text{max}}$

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 $\sim 10^{-11}$

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