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Authors

Krishnagopal, S.
Siemann, R.H.

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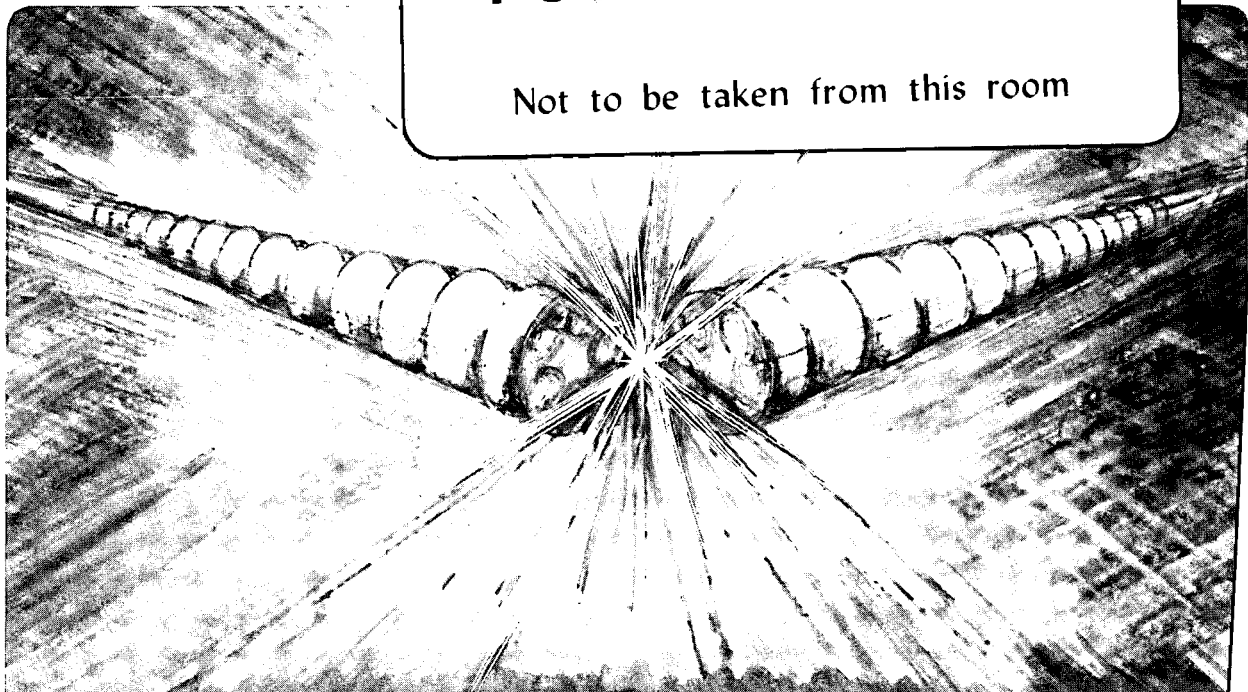
S. Krishnagopal and R. Siemann

May 1991

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The Coherent Beam-Beam Interaction

S. Krishnagopal*

Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720

R. Siemann**

Stanford Linear Accelerator Center, Stanford, CA 94305

May 1991

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The Coherent Beam-Beam Interaction

S. Krishnagopal (*)
Lawrence Berkeley Laboratory,
Berkeley, CA 94720

R. Siemann (**)
Stanford Linear Accelerator Center,
Stanford, CA 94309

Abstract

We present the results of a new beam-beam simulation program that allows for a self-consistent calculation of the electromagnetic fields of the beams by treating general (non-Gaussian) beam-distributions. We find that a new class of coherent instabilities, appearing at certain operating points, dominate the dynamics.

I. INTRODUCTION

One of the factors limiting the performance of e^+e^- storage ring colliders is the beam-beam interaction. There has been much speculation on the role of coherent (or collective) beam-beam effects as a mechanism for limiting the tune-shift, but no consensus has been reached on this issue[1]. Centroid (or dipole) motion, where the centroids of the two beams oscillate relative to each other, is routinely observed in operating storage rings, but there is no evidence that it affects luminosity. Centroid motion is easily detected and could be removed with feedback.

The potential for performance limitations comes from effects that distort the beam shape. Such effects have been analyzed with two different types of models. In the first, of Furman *et al.*[2] and of Hirata[3], nonlinear maps for the colliding beam system are developed in the moments. In earlier work Hirata uses a nonlinear beam-beam kick calculated from a Gaussian beam, which is not consistent with the assumption that the beams remain Gaussian. He finds flip-flop solutions, where both beams maintain unequal sizes. Later attempts at including higher moments were in substantial agreement with the Gaussian calculation. Furman *et al.* maintain self-consistency at the cost of a simplified, linear model of the beam-beam force. They find that equal-size higher-period solutions coexist, and are responsible for performance limitations.

In the second type of model, of Chao and Ruth[4] and of Dikansky and Pestrikov[5], the phase space distributions of the two beams influence each other and modes develop in phase space. The stability of these modes is analyzed with the linearized Vlasov equation, assuming small perturbations from equilibrium. While these calculations indicate the potential importance of coherent beam-beam effects, there are open questions about the approximations used

in the calculations, Landau and radiation damping, and the relative importance compared to other effects.

Strong-strong computer simulations are an important tool in the study of the coherent beam-beam interaction. In the usual 'Gaussian simulations' the positions and rms sizes of the beams are first calculated from the coordinates of the test-particles just before the collision. They are then used in an expression for the beam-beam force that assumes the particle distribution to be Gaussian[6]. Unfortunately, this procedure does not allow for a self-consistent calculation of the fields; it restricts the fields directly and the distributions implicitly.

The simulation program discussed in this paper calculates the fields from the coordinates of the test-particles, and does not impose any restrictions on the beam-distributions or the beam-profiles. We find new coherent instabilities that, at certain operating points, dominate the dynamics.

II. FIELD CALCULATION

We worked with beams that were nominally round, *i.e.* the β^* 's and nominal emittances were equal in the two transverse dimensions. The beams were not restricted to remain round[7], but the field calculation is most efficient when the transverse sizes are comparable.

The electromagnetic fields of a bunch are calculated by Lorentz transforming to its rest frame and then solving Poisson's equation. Test particles are cast onto a circular mesh with radial step size Δr and azimuthal bin size $\Delta\phi$. The array that results, $N(r, \phi)$, must be smoothed, otherwise the simulation results are affected by statistical fluctuations due to the finite number of test particles. The average charge is treated first. It is smoothed by least-squares fitting using Forsythe's method[8], and the radial electric field it produces is calculated using Gauss' Law. The azimuthal variation of $N(r, \phi)$ is treated by Fourier analysing $N(r, \phi)$ and checking each Fourier coefficient for consistency with zero. Statistically significant coefficients are fitted and smoothed with Forsythe's method, and the Green's function for Poisson's equation in polar coordinates[9] is used to calculate the potential and the electric field.

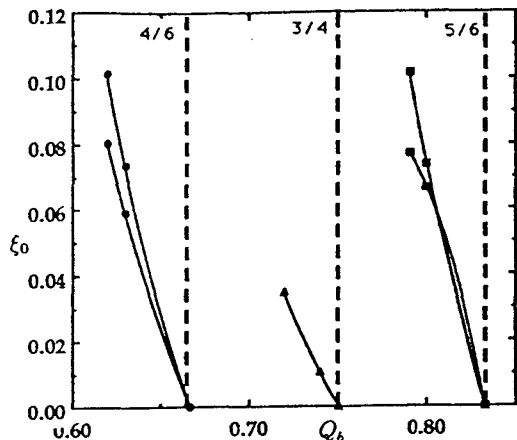


Figure 1: Onset and offset ξ_0 for the various coherent resonances, as a function of Q_b . At the nominal tune of the resonance it is assumed that the resonance ‘tube’ vanishes. The lines are drawn by interpolating from the calculated data points down to the zero- ξ_0 point.

This algorithm was checked for sensitivity to the details of the mesh – the sizes of Δr and $\Delta\phi$ – and was found to be insensitive to them. On the other hand, the number of test particles and the smoothing provided by the least squares fitting were important. Without smoothing, statistical fluctuations dominated for any reasonable number of test particles. With smoothing, results were found to be insensitive to the number of test particles when that number exceeded 5,000; in this work we used 10,000 test particles.

III. SIMULATION RESULTS

Our results show that the coherent behaviour is sensitive to the radiation damping. In this work we concentrated on low-order resonances, up to sixth order, and on damping decrements of $\delta = 1 \times 10^{-3}$ and $\delta = 1 \times 10^{-4}$, corresponding to betatron damping times of 2,000 and 20,000 turns, respectively. Synchrotron motion was not included, and the beams were forced to collide head-on by a ‘feedback’ system that set the centroids of the beams to zero after each turn. The horizontal and vertical tunes were kept equal ($= Q_b$, say), and Q_b was restricted to the region $0.5 < Q_b < 1.0$. Results are presented only for the x -dimension; the y -dimension behaved analogously.

A Tune-Shift Scans

Figure 1 presents the results of scans over the nominal tune-shift ξ_0 , performed at different tunes, in order to determine the tune-shift region over which coherent motion broke out. The data-points in the figure correspond the lowest and highest ξ_0 ’s at which coherent motion was seen, at that particular tune.

For the fourth-order ($\frac{3}{4}$) resonance there seems to be no upper limit to the coherent motion, at least to the max-

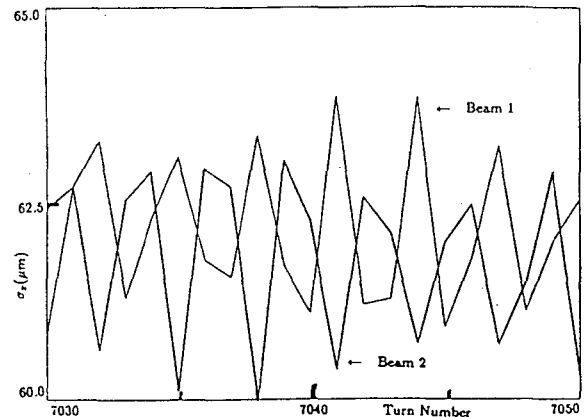


Figure 2: Horizontal beam-size variation over 20 turns. $Q_b = 0.79$, $\xi_0 = 0.08$.

imum ξ_0 of 0.14 investigated. A study of the beam-size variation indicates that the behaviour is period-2 and anti-correlated between the two beams; *i.e.* when one beam is tightly focussed, the other is blown-up. This is consistent with the results of Furman et al.[2] which show a strong fourth-order resonance. This is not a new feature of these simulations; the same behaviour is seen even in Gaussian simulations.

Two sixth-order resonances, the $\frac{4}{6}$ and the $\frac{5}{6}$, were identified and traced out in tune. In contrast to the $\frac{3}{4}$ resonance, they were found to have a *finite* width in ξ_0 . This is consistent with the predictions of the second class of models discussed above[4, 5]. The oscillations in the beam-sizes were found to be period-3 and anti-correlated (Fig. 2).

Just below $Q_b = 0.667$ both third ($\frac{2}{3}$) and sixth ($\frac{4}{6}$) order resonances could play a role in the dynamics. On the other hand below $Q_b = 0.833$ only the sixth order ($\frac{5}{6}$) resonance is possible. However, the similarity of the coherent dynamics in these two regions (Fig. 1) suggests that the third-order resonance does *not* contribute below $Q_b = 0.667$. More generally one may surmise that odd order coherent resonances do not occur at all. This agrees with the predictions of the second class of models[4, 5], and is significant from the practical point of view.

B Gaussian vs General Simulations

To show that the sixth-order resonances are a new feature of these simulations, we performed a set of comparison runs with a Gaussian simulation. For the $\frac{5}{6}$ resonance, at a tune of $Q_b = 0.80$, the region over which coherent motion was seen was scanned using a Gaussian simulation.

The results are presented in Fig. 3 in terms of the beam-size and its variation as a function of ξ_0 for both, the general and the Gaussian simulations. For the former, with 10,000 test particles, the statistical fluctuation in the beam-size is less than a micron. For the latter, with 1,000 particles, it is just over a micron. Larger size-variations are indicative of coherent motion.

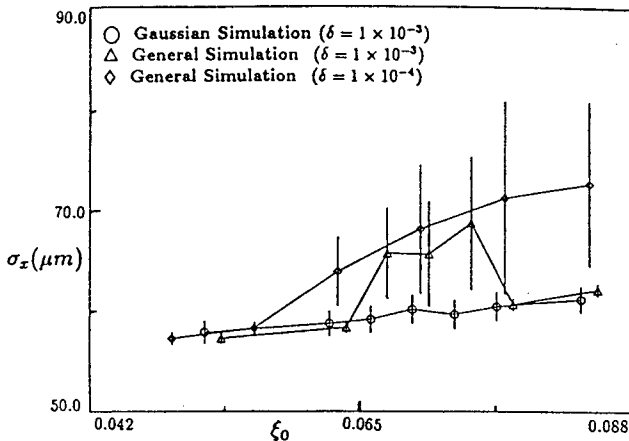


Figure 3: Horizontal beam-size as a function of ξ_0 at $Q_b = 0.80$, for three different cases.

From the figure one sees that for the sixth-order resonance coherent oscillations are present *only* in the more general simulation; they are absent in the Gaussian simulation.

IV. DISCUSSION

The last observation above suggests that it is critical to use general field calculations in the study of nonlinear coherent resonances. The reason is that this allows for a *self-consistent* calculation of the electromagnetic fields from the positions of the test particles. It results in higher-order fixed-point coherent motion.

It should be noted that centroid motion has been removed in these simulations, allowing us to concentrate on quadrupolar coherent oscillations. This is also reasonable from the experimental point of view, because such motion could be removed using feedback systems. If it is allowed to remain, we find that the beam-size variations persist, although they are reduced in magnitude. There is now a large degree of centroid motion between the beams. Thus both centroid and shape oscillations take place, and the situation is complex.

As mentioned earlier, the coherent behaviour is sensitive to the damping decrement δ . With lesser damping the width of the resonance becomes considerably larger, as can be seen from Fig. 3. Further, this could potentially result in higher-order resonances entering the picture; some preliminary results with $\delta = 1 \times 10^{-4}$ indicate that this is indeed the case. Thus, in machines with low radiation damping these coherent resonances could play an important role.

It should be emphasized that the signature of these coherent resonances is a swift, turn-to-turn, variation in the beam-sizes. Since existing detectors are not sensitive to such rapid variations, these resonances could not, hitherto, have been detected.

V. FUTURE WORK AND CONCLUSIONS

Much work remains to be done, and is in progress. Resonance structure with lower damping is being investigated. The impact of removing feedback has to be studied. Sensitivity to initial conditions, consequences of asymmetric parameters between the two beams, and methods of handling flat beams; all these avenues can be explored.

We conclude by observing that a self-consistent calculation of the electromagnetic fields is essential in a study of the coherent beam-beam interaction. It leads to a new class of higher-period coherent instabilities that dominate the dynamics at certain operating points.

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LAWRENCE BERKELEY LABORATORY
UNIVERSITY OF CALIFORNIA
INFORMATION RESOURCES DEPARTMENT
BERKELEY, CALIFORNIA 94720