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Detailed seismic imaging of the Mw 7.1 Ridgecrest earthquake rupture zone from data recorded by dense linear arrays

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1	Seismic imaging of the Mw 7.1 Ridgecrest earthquake rupture zone from data
2	recorded by dense linear arrays
3	
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21	Key points:
22 23	• Several 1- to 2-km-wide low-velocity zones with more intensely damaged inner cores (0.5-1.5 km wide) are identified beneath each array
24 25	• An automated detector, based on peak ground velocities and durations of high amplitude <i>S</i> waves, identifies fault-zone trapped waves
26 27	• The <i>P</i> wave delay time and <i>S</i> wave amplification patterns indicate consistent locations and widths of fault damage zones

# 28 Abstract

29 We analyze seismograms recorded by four arrays (B1-B4) with 100-m station spacing 30 and apertures of 4-8 km that cross the surface rupture of the 2019 Mw 7.1 Ridgecrest 31 earthquake. The arrays extend from B1 in the northwest to B4 in the southeast of the 32 surface rupture. Delay times between P-wave arrivals associated with ~1200 local 33 earthquakes and four teleseismic events are used to estimate local velocity variations 34 beneath the arrays. Both teleseismic and local P waves travel faster on the northeast than 35 the southwest side of the fault beneath arrays B1 and B4, but the velocity contrast is less 36 reliably resolved at arrays B2 and B3. We identify several 1- to 2-km-wide low-velocity 37 zones with much slower inner cores that amplify S waveforms, inferred as damage zones, 38 beneath each array. The damage zones at arrays B2 and B4 also generate fault-zone head 39 and trapped waves. An automated detector, based on peak ground velocities and 40 durations of high-amplitude waves, identifies candidate fault-zone trapped waves 41 (FZTWs) in a localized zone for ~600 earthquakes at array B4. Synthetic waveform 42 modeling of averaged FZTWs, generated by ~30 events with high-quality signals, indicates that the trapping structure at array B4 has a width of ~300 m, depth of 3-5 km, 43 44 S-wave velocity reduction of  $\sim 20\%$  with respect to the surrounding rock, Q-value of 45  $\sim$ 30, and S-wave velocity contrast of  $\sim$ 4% across the fault (faster on the northeast side). 46 The results show complex fault-zone internal structures (velocity contrasts and low-47 velocity zones) that vary along fault strike.

48

# 49 Plain Language Summary

50 The 2019 Mw 7.1 Ridgecrest earthquake in the Eastern California Shear Zone generated 51 a vigorous aftershock sequence that provided a wealth of seismic data. We derive 52 subsurface structural properties within and across the Ridgecrest rupture zone from 53 seismic waveforms generated by the earthquake sequence. The data are recorded by four 54 dense nodal arrays that were deployed across the Ridgecrest rupture zone with ~100 m 55 spacing and aperture of a few kilometers. Delay times of P wave and amplification of S 56 waves are used to infer on several 1-2-km-wide low-velocity zones with more intensely 57 damaged inner cores (0.5-1.5 km wide) beneath each array. Waveform modeling of fault

58 zone trapped waves well-recorded by one array provides geometrical and seismic 59 properties of a coherent waveguide in the damage fault zone structure at that location. 60 The results are complementary to tomographic models that provide a regional context but 61 do not resolve internal structural elements of the Ridgecrest rupture zone.

62

# 63 **1. Introduction**

64 The Mw 7.1 Ridgecrest earthquake of July 5, 2019 and the earlier Mw 6.4 event on 65 July 4 in the southern part of the Walker Lane shear zone (Fig. 1) were felt throughout 66 southern California and produced a vigorous aftershock sequence. These events led to 67 rapid deployments of seismic arrays across and around the Ridgecrest earthquake 68 sequence (Catchings et al., 2020). Kinematic rupture processes of the Mw 6.4 and Mw 69 7.1 events, surface deformation, and properties of the aftershocks show complex patterns, 70 with strong variations both along strike of the rupture zones and in depth (e.g., Chen et 71 al., 2020; Cheng & Ben-Zion, 2020; Jia et al., 2020; Ross et al., 2019; Xu et al., 2020). 72 Data recorded by several dense arrays crossing the rupture zone of the Mw 7.1 73 earthquake can be used to derive high-resolution seismic information on the internal 74 structure of the rupture zone. Detailed imaging of the structure associated with the 75 rupture zone can provide important information on various topics, including initiation and 76 arrest of ruptures (e.g., Aki, 1979; King, 1986), amplification of seismic waves (e.g., 77 Kurzon et al., 2014; Rovelli et al. 2002; Spudich & Olsen, 2001), interactions of ruptures 78 with fault zone properties (e.g., Ben-Zion & Huang 2002; Brietzke & Ben-Zion 2006; 79 Huang et al., 2014), and properties of earthquake sequences (e.g., Thakur et al., 2020).

80 Several velocity models for the Ridgecrest area provide information for volumes with 81 spatial resolutions ranging from several km (e.g., Lee et al., 2014; Shaw et al., 2015) to 82 about 500 m (White et al. 2020). Structures in the top 1-2 km are poorly resolved in these 83 velocity models due to limitations of the input data. Analyses of seismic data recorded by 84 arrays across faults and rupture zones have proven highly effective in complementing 85 regional velocity models and imaging sharp bimaterial interfaces and damage zones with 86 width of a few tens of meters (e.g., Ben-Zion et al., 2003; Cochran et al., 2009; Li et al., 87 1994; Peng et al., 2003; Qin et al., 2018; Qiu et al., 2017; Share et al., 2019).

88 In this study, we investigate the seismic and geometrical properties of the damage 89 structure associated with the 2019 Mw 7.1 Ridgecrest earthquake, based on the data 90 obtained from four dense linear seismic arrays (B1-B4; triangles in Figs. 1 and 2) located 91 across segments of the rupture. Analyses of the arrival patterns of P waves from both 92 teleseismic and local seismic events across each array helps to detect and constrain 93 properties of velocity contrasts across fault sections and low-velocity zones reflecting at 94 least partially damaged rock. We identified fault-zone trapped waves, i.e., amplified 95 motions of S waves associated with core damage zones that are sufficiently coherent to 96 act as a waveguide, at some locations and inverted for average geometrical and seismic 97 properties of the fault-zone waveguide.

In the following sections, we describe the deployment and data processing in section 2 and present the methodology and results on various aspects of the fault-zone structures from different types of observations in section 3. The imaging results from different phases and analyses are summarized and discussed in section 4. The results show overall complex fault-zone structures that vary along the rupture strike, in general agreement with fault surface traces (Xu et al., 2020), seismic catalog (Ross et al., 2019) and potency of aftershocks (Cheng & Ben-Zion, 2020) in the Ridgecrest area.

105

# 106 2. Data & basic processing

107 Four linear arrays, with about 100-m station spacing and apertures of 4-8 km, were 108 deployed across the surface rupture of the 2019 Mw 7.1 Ridgecrest earthquake (Fig. 1). 109 The arrays extended from B1 in the northwest to B4 in the southeast of the surface 110 rupture (Fig. 2). In total, the B-arrays consisted of 248 Fairfield and SmartSolo sensors 111 that recorded continuously at 500 Hz for about a one-month period (7/12/2019-8/8/2019). 112 For teleseismic delay time analysis (Section 3.1), we use the Taup toolkit (Crotwell et al., 1999) and velocity model IASP91 (Kennett & Engdahl, 1991) for predictions of P-113 114 arrival time at each station. The teleseismic earthquakes have epicentral distances between 30-90°, depth > 50 km, and Mw > 6.0. For analysis of local P waves (Section 115 3.2), we first extract the seismic waveforms generated by ~1200 local events (red box in 116 Fig. 1) at each station and use the catalog of Hauksson et al. (2012, extended to 2019) for 117

118 locations. The mean and linear trend are removed from the waveforms, and a bandpass 119 filter between 0.5 Hz and 20 Hz is applied. In the study of fault zone trapped waves 120 (Section 3.3), the north-south and east-west components are rotated to a coordinate 121 system parallel and perpendicular to the fault strike.

122

# 123 **3. Analysis**

We conduct three types of studies involving different signals and spatial scales to image several components of the fault-zone structure associated with the 2019 Mw 7.1 Ridgecrest earthquake beneath the four linear arrays (Fig. 2). We describe the analyses below, starting with large-scale structural features (e.g., overall velocity variations across the fault) and progressing to inner fault-zone components (e.g., geometry and velocity of the damage zone). The results are obtained using teleseismic delay-time analyses (DTA), local *P*-wave DTA, and analysis associated with FZTWs following the *S*-wave arrival.

131

#### 132 3.1 Teleseismic delay time analysis

During the one-month deployment, teleseismic *P* waves with sufficient signal to noise ratios (SNR > 5) between 0.5 Hz and 2 Hz were recorded for three events at array B1 and four earthquakes at arrays B2-B4 (Figs. S1- S4). We do not investigate teleseismic *S* waves since they have SNR < 5.

137

# 138 *3.1.1 Methodology*

139 As shown in previous studies (e.g., Ozakin et al., 2012; Qiu et al., 2017), there are 140 three contributing factors to travel-time delays observed on a linear array for a 141 teleseismic arrival: the geometry between the incoming plane wave and the array, 142 topography, and the crustal structure beneath the array. To obtain the travel-time delays 143 due to local crustal structures, we first predict the arrival time of the teleseismic *P*-wave 144 for each station and event pair using the IASP91 model and assume the station is at sea level. Then, teleseismic P waveforms are bandpass filtered between 0.5 Hz and 2 Hz. By 145 146 aligning the teleseismic P waves with respect to the corresponding predicted arrival time 147 (Fig. 3b) at each station, we remove the delay times associated with the non-vertical

incidence angle of incoming waves. This assumes a flat Moho interface beneath the
array, which is likely given the small area involved and receiver function results (e.g.,
Zhu & Kanamori, 2000).

151 In Qiu et al. (2017), the P waveform is first stacked over the entire array for a specific 152 teleseismic event and used as the reference. Then, the arrival-time pattern of P waves is 153 extracted from cross-correlations of the P waveform recorded at each station and the 154 reference waveform. However, this method is accurate only for arrays with short aperture (e.g., ~500 m in Qiu et al., 2017) when P-wave delay times are small (e.g., ~0.01 s). For 155 156 arrays with long aperture, the P waveform recorded at a specific station may be used as 157 the reference. Considering the long aperture of the B-arrays (4-8 km; Fig. 2), we estimate 158 the P-wave delay time pattern by cross-correlating waveforms within a narrow P-wave 159 window (Fig. 3b) for every pair of stations i and j. The center of the narrow P-wave window is determined based on the array-mean envelope function (Fig. 3b), and the peak 160 161 frequency of the array-mean P-wave amplitude spectrum (Fig. 3c) is used to set the 162 window width to be twice the dominant period. To further enhance the P-wave signals, we apply another filter with a narrower frequency band (black dashed lines in Fig. 3c) to 163 164 the teleseismic data prior to the cross-correlation.

165 Let  $\tilde{t}_{ij}$  be the time delay corresponding to where the cross-correlation function 166 between *P* waveforms recorded at the *i*- and *j*-th stations reaches the maximum,

$$\tilde{t}_{ij} = \tilde{T}_i - \tilde{T}_j, \tag{1a}$$

167 where  $\tilde{T}_i$  and  $\tilde{T}_j$  are the teleseismic *P*-wave travel times at the *i*- and *j*-th stations, 168 respectively. Since the mean of the arrival time pattern  $\tilde{T}_i$  has no significance for our 169 imaging, we can ignore the constant  $\tilde{T}_j$  in equation (1a) and the delay time pattern is 170 given by  $\tilde{T}_{i_j} = \tilde{t}_{ij}$  when the *j*-th station is fixed and set to be the reference. We note that 171 the delay time patterns obtained by using different reference stations should be 172 consistent, i.e.,  $\tilde{T}_{i_j} - \tilde{T}_{i_j}$  is a constant. Thus, we can minimize the measurement error 173 by averaging equation (1a) over *j*,

$$\tilde{T}_{i} = \sum_{j=1}^{N} \tilde{t}_{ij} / N + \sum_{j=1}^{N} \tilde{T}_{j} / N = \sum_{j=1}^{N} \tilde{t}_{ij} / N + C,$$
(1b)

174 where N is the number of stations and C is a constant. We again ignore the constant C175 and remove the effect of un-modeled topography from the teleseismic *P*-wave delay time

$$T_i = \sum_{j=1}^{N} \tilde{t}_{ij} / N - \Delta h_i / v_{\text{corr}},$$
(2)

176 where  $v_{corr}$  is the *P*-wave velocity (Vp) and  $\Delta h_i = h_i - \sum_{j=1}^N h_j/N$  represents the 177 relative topography, with  $h_i$  indicating the elevation at the *i*-th station (Figs. 2b-e). Here, 178 for the topographic correction, we assume a vertical-incidence angle for the incoming *P* 179 wave and a constant Vp that is likely representative of velocity structures averaged over 180 the top 1-2 km (e.g., Park et al., 2019). Since the velocity structure at the shallow section 181 is poorly constrained by existing velocity models, we use two constant Vp values, 2 km/s 182 and 4 km/s, to estimate the lower and upper limits of  $v_{corr}$ .

183

#### 184 *3.1.2 Results*

185 Coherent P arrivals, with different peak frequencies, are observed crossing the array 186 for the four events in Figure S1. Although the frequency content of the P waveforms is 187 different between events, the obtained arrival patterns prior to topographic correction are in general consistent (e.g., higher velocity in the northeast beneath array B4 in Fig. 4d). 188 189 The black curves in Figure 4 depict the teleseismic *P*-arrival patterns for each array 190 averaged over all events with the standard deviation of the mean giving the uncertainty, 191 and the delay times after correcting the array topography are illustrated as colored dashed 192 curves. Features of delay-time patterns associated with a velocity contrast across the fault 193 (i.e., a smoothed step function as in Fig. 6d of Qiu et al., 2017) and a low-velocity zone 194 (i.e., a mountain-shaped function as in Fig. 6b of Qiu et al., 2017) are both observed in 195 the results after the topographic correction.

Delay-time patterns resolved at arrays B1 and B4 yield clear velocity contrasts across the fault, with the southwest block having later arrivals indicating slower velocity ( $\sim 0.15$ s and  $\sim 0.25$  s in *P*-wave arrival time; Figs. 4a and 4d). Topographic corrections have minor effects on the resolved arrival-time patterns at both arrays, as differences between the two dashed curves are negligible in Figures 4a and 4d. Compared to the pattern dominated by the velocity contrast, the time delay associated with low-velocity zones is less obvious and comparable to the level of uncertainties (Figs. 4a and 4d). In contrast,

the dominant feature in arrival patterns resolved at arrays B2 and B3 yields several  $\sim 1$ km-wide low-velocity zones that generate a maximum time delay of  $\sim 0.04$  s (Figs. 4b-c): two centered at about 2.5 km southwest and 0.5 km northeast of the center of array B2; one centered at  $\sim 0.5$  km southwest of the B3 array center.

207 Although the topographic correction has negligible effects on the low-velocity zones 208 resolved at arrays B2 and B3, the time delays associated with velocity contrasts across 209 the fault beneath the two arrays are much weaker and vary significantly with the Vp used 210 in the correction (Figs. 4b-c). Different from the other three arrays, the polarity of the 211 velocity contrast at array B3 is flipped for arrival patterns resolved using Vp of 2 km/s 212 and 4 km/s. This is likely due to the combination of the larger difference in topography 213 and smaller velocity contrast at arrays B2 and B3. Therefore, we do not discuss the 214 velocity contrasts resolved from teleseismic P waves recorded at arrays B2 and B3 in 215 later sections.

216

# 217 3.2 Local *P*-wave delay time analysis

P waves from local earthquakes recorded by the B-arrays are observed at higher frequencies (peaks at  $\sim$ 8 Hz; e.g., Fig. 5a) compared to those of teleseismic events (between 0.5-2 Hz; e.g., Fig. 3). Thus, higher resolution images of local fault zone structures can be achieved by analyzing arrival times of direct *P* waves from local earthquakes across each array.

223

# 224 3.2.1 Methodology

225 Compared with teleseismic arrivals, the effect of source-receiver geometry on P226 waves for local earthquakes recorded by an array requires additional processing than the 227 plane wave correction used in section 3.1.1. To extract the variations in *P*-wave arrival 228 times associated with local fault-zone structures, we first suppress the contributions from 229 source-receiver geometry and topographic variations by dividing the time axis of the P230 waveform recorded at the *i*-th station for event *j* with its corresponding propagation 231 distance  $H_{ii}$  (e.g., from Fig. 5a to 5b). Here, we use the hypocentral distance to approximate  $H_{ij}$ . P-wave picks,  $s_{ij} = t_{ij}/H_{ij}$  in units of slowness (e.g., Figs. 5b and S5b-232 233 S7b), are then picked via the short-term-average/long-term-average (STA/LTA) algorithm (Allen, 1978) using waveforms within the slowness range of 0.15-0.25 s/km (to exclude the effect of *S* waves). The *P*-wave SNR is calculated as the ratio between the maximum and root mean square amplitudes of waveforms in slowness windows from  $s_{ij}$ to 0.25 s/km and 0.15 s/km to  $s_{ij}$ , respectively. *P*-wave picks with SNR less than 10 are not used, and events are excluded if less than 80% of the array shows good quality *P*wave picks (SNR > 10).

240 The array-mean slowness  $\bar{s}_i = \sum_{i=1}^M s_{ij}/M$ , with M being the number of stations, can 241 vary significantly with focal depth and epicenter location (due to 3-D velocity structures). 242 Therefore, we use relative slowness,  $\hat{s}_{ij} = s_{ij}/\bar{s}_j$  (e.g., Qiu et al., 2017; Share et al., 243 2017) to characterize statistical features of the local structure-related P-wave arrival 244 pattern using all available events. We can also estimate the local structure-related *P*-wave arrival pattern in delay time,  $\delta t_{ij}$  for station *i* and event *j*, as  $\delta t_{ij} = (s_{ij} - \bar{s}_j) \cdot H_{ij}$ . Since 245 246 the existing velocity models indicate that structures in the Ridgecrest area are highly 247 heterogeneous, we only analyze P waves from events with depth > 5 km and close to 248 each array (red box in Figs. 5c and S5c-S7c) to ensure that the hypocentral distance is a 249 good approximation of the propagation distance  $H_{ij}$  and the resolved delay time pattern is 250 representative of local structures beneath the array.

Based on previous fault zone studies (e.g., Qiu et al., 2017; Share et al., 2017), the observed *P*-wave travel time at station i for a near-fault event j mainly consists of two components:

$$t_{ij} = s_{ij} \cdot H_{ij} \approx \bar{s}_j \cdot H_{ij} \cdot (1 + \eta_i) + \Delta t_{ij}.$$
(3)

254 The first term of equation (3) indicates the time delay associated with the cross-fault 255 velocity contrast beneath the array, and, ideally,  $\eta_i$  is a step function, i.e., equals  $-\eta/2$ 256 and  $\eta/2$  ( $\eta < 1$ ) for stations on the faster and slower crustal blocks, respectively. The 257 second term  $\Delta t_{ii}$  represents the contribution of local structures beneath the array at 258 shallow depth (e.g., fault damage zone, sedimentary basin). Let  $\theta_i$  be the average incident 259 angle of P waves from event j to the array,  $\Delta t_{ij}$  is inversely proportional to  $\cos \theta_i$ , i.e., 260  $\Delta t_{ij} \approx \tau_i / \cos \theta_j$ , where  $\tau_i$  is the delay time of a vertically incident P wave associated 261 with shallow structures beneath the array at station *i*.

Following the derivation in Text S1, the arrival patterns averaged over all near-fault and close-to-array events are approximately given by:

$$\hat{S}_{i} = \sum_{j=1}^{N} \hat{s}_{ij} / N \approx 1 + \eta_{i} + \tau_{i} / (\tilde{t}_{i} \cos \tilde{\theta}) \approx 1 + \eta_{i} + \tau_{i} / (\tilde{t} \cos \tilde{\theta}),$$
(4a)

in relative slowness, and

$$\delta T_{i} = \sum_{j=1}^{N} \delta t_{ij} / N \approx \eta_{i} \cdot \bar{t}_{i} + \tau_{i} / \cos \tilde{\theta} \approx \eta_{i} \cdot \bar{t} + \tau_{i} / \cos \tilde{\theta}, \qquad (4b)$$

265 in delay time, after dropping the higher order terms.  $\tilde{t}_i$  and  $\bar{t}_i$  are the harmonic and 266 arithmetic means of *P*-wave travel time  $t_{ij}$ , respectively, over all events. *N* is the number 267 of events and  $\tilde{\theta}$  is the mean incidence angle averaged over all events.  $\tilde{t}$  and  $\bar{t}$  denote the 268 array-mean travel times of  $\tilde{t}_i$  and  $\bar{t}_i$ , respectively. It is interesting to note that the shape of 269 arrival patterns  $\hat{S}_i$  and  $\delta T_i$  is the same, i.e.,  $\delta T_i / (\hat{S}_i - 1)$  is a constant, when  $\tilde{t} \approx \bar{t}$ .

270

#### 271 *3.2.2 Results*

272 Figure 6a shows the results of statistical analysis on the local structure-related *P*-wave arrival pattern in relative slowness ( $\hat{S}_i$  in eq. 4a) and delay time ( $\delta T_i$  in eq. 4b) for array 273 274 B1 using 189 near-fault events with depth > 5 km outlined by the red box in Figure 5c. 275 The relative slowness patterns are averaged over all analyzed events and the standard 276 deviation of the mean is used to estimate the uncertainty. The small error bars and 277 confined width of histograms suggest that the mean relative slowness curve is 278 representative of the patterns observed from all analyzed events. The mean pattern in 279 delay time  $(\delta T_i)$  with small uncertainties (Fig. 6a) is observed to have the same shape as that of the relative slowness, i.e.,  $\delta T_i/(\hat{S}_i - 1)$  is approximately a constant. This is 280 consistent with equation (4) when  $\tilde{t} \approx \bar{t}$ . 281

Good agreement between the mean patterns of relative slowness and delay time with small uncertainties is also observed at arrays B2-B4 in Figures 6c, 6e, and 6g. Consistent with the teleseismic *P*-wave arrival time pattern shown in Figure 4, we observe the features of delays in local *P*-wave arrival time associated with a step-function-like and several mountain-shaped components at each array (left panels of Fig. 6) that likely correspond to a velocity contrast across the fault and low-velocity zones, respectively. 288 Similar mean patterns of relative slowness are also observed for subsets of events 289 grouped according to different narrow ranges of depth (Figs. S8a-S11a). The polarity of 290 the velocity contrast is the same for arrays B1 and B4 (Figs. 6a and 6g), with the 291 southwest block being slower, consistent with results of the teleseismic delay time 292 analysis (Figs. 4a and 4d). However, delay time patterns at arrays B2 and B3 show the 293 opposite polarity (Figs. 6c and 6e), i.e., the northeast block having slower velocity. This 294 is consistent with results of the teleseismic delay time analysis before the topographic 295 correction (Figs. 4b-c).

296 We note that delay times associated with shallow materials ( $\tau_i$  in eq. 4a) could also 297 yield a step-function-like pattern (e.g., variations in sediment thickness across fault). 298 Thus, the amplitude of the step-function-like component,  $\tilde{\eta}$ , is an upper limit of the 299 magnitude of the cross-fault velocity contrast. We measure the amplitude of  $\tilde{\eta}$  as the 300 difference between relative slowness values averaged over stations on the southwest and 301 northeast edges of the array (left panels of Fig. 6), and the values are 5.3%, 0.8%, 3.9%, 302 and 8.8% at the sites of arrays B1-B4, respectively. Following equation (4a), if the 303 observed step-function-like pattern is due to variations in shallow materials, i.e.,  $\tau_i/(\tilde{t}\cos\tilde{\theta})$ , the amplitude of  $\tilde{\eta}$  should decrease with depth as the mean incidence angle 304  $\tilde{\theta}$  and travel time  $\tilde{t}$  increase with propagation distance. This contradicts the observations 305 306 that  $\tilde{\eta}$  increases with depth at array B2 (Fig. S9a) and only slightly varies with depth at 307 arrays B1 (Fig. S8a), B2 (Fig. S10a), and B4 (Fig. S11a), suggesting the observed step-308 function-like pattern is likely a good approximation of the cross-fault velocity contrast, if 309 the contribution due to elevation variation is negligible.

310 To further analyze the effect of array topography in the step-function-like component  $\tilde{\eta}$ , we first model the delay times of cross-fault velocity contrast (i.e.,  $\eta_i \cdot \bar{t}$  in eq. 4b) with 311 312 a smoothed step function (Text S2). The modeled delay times of cross-fault velocity 313 contrast consist of three components: a linear trend within a transition zone between two 314 groups of stations at the southwest and northeast edges with constant values (eq. S5 in 315 Text S2; Fig. 6). We then illustrate the location of the resolved transition zone beneath 316 each array in Figure 2. At array B2, the span of the transition zone covers the area with a 317 steep slope in the array topography (Fig. 2c). Considering the large variation in elevation (> 100 m) and the small amplitude of  $\tilde{\eta}$  (0.8%), we conclude that the modeled cross-fault 318

velocity contrast may be dominated by the residual in topographic correction rather than structure-related delay times. In contrast, for the other three arrays, the topography variation is less significant and the amplitude of  $\tilde{\eta}$  is much larger (> 4%), indicating the array topography likely has negligible contribution to the modeled cross-fault velocity contrast.

324 The width of these transition zones and their locations relative to fault surface traces 325 provide additional information on the internal structures of the fault zone. At array B1, 326 the southwest edge of the transition zone correlates well with the surface trace of the Mw 327 7.1 earthquake (Fig. 2b), which is indicative of asymmetric rock damage offset to the 328 northeast (faster side). At array B3, the northeast edge of the transition zone agrees well 329 with the surface trace of the Mw 6.4 earthquake. The data recorded by array B3, crossing 330 the surface ruptures of both the Mw 6.4 and Mw 7.1 events, likely detect the velocity 331 contrast across the fault that hosted the Mw 6.4 earthquake, with the northwest side being 332 higher in velocity. An asymmetric rock damage offset to the faster side (northwest) is 333 also observed beneath array B3 (Fig. 2d). At array B4, the transition zone is much wider 334 and almost covers the entire array. This is consistent with the fact that the array is at the 335 southeast end of the Mw 7.1 earthquake rupture, where the rupture zone has several 336 surface traces and is less localized (Fig. 2e).

Low-velocity zones that further delay the *P*-wave arrivals are also observed at each 337 338 array, in addition to the pattern associated with the cross-fault velocity contrast (left 339 panels of Fig. 6). To highlight contributions from these low-velocity zones (i.e.,  $(\tau_i/(\tilde{t}\cos\tilde{\theta}))$  in eq. 4a), the right panels of Figure 6 show the mean delay time pattern by 340 subtracting the modeled  $\eta_i \cdot \bar{t}$  from  $\delta T_i$  (eq. 4b). A heat map of delay times  $\delta T_i$  after 341 subtracting the modeled  $\eta_i \cdot \bar{t}$  is also shown. Similar delay patterns of  $\tau_i$  reflecting 342 343 shallow materials beneath the array are seen using events within different narrow ranges 344 of depth (Figs. S8b-S11b). The range of delay patterns related to the major low-velocity 345 zones underneath each array are outlined with green dashed lines in Figure 6. The slower 346 inner cores (with peak delay time > 0.03 s) of these low-velocity zones are marked by the 347 red dashed lines.

Consistent with the teleseismic *P*-wave arrival patterns obtained at arrays B2 and B3 (Figs. 4b-c), we observe low-velocity zones with comparable widths centered at the same

350 locations. In addition, we retrieve higher resolution images of the narrower and slower 351 inner cores that correlate well with locations of the local peaks identified in the 352 teleseismic P-wave delay time patterns (Figs. 4b-c). This is likely due to the shorter 353 wavelength of local seismic P waves that can provide high-resolution images of internal 354 fault-zone structures (e.g., Dahlen et al., 2000). Delay patterns related to low-velocity 355 zones are also observed in the results for arrays B1 and B4 (Figs. 6b and 6h), which are 356 hard to identify in the teleseismic delay time analyses (Figs. 4a and 4d) due to large 357 uncertainties (i.e., insufficient number of analyzed events).

358 We note that consistent patterns of the cross-fault velocity contrast and low-velocity 359 zones are observed (with different amplitudes) in results obtained using events with large 360 along-fault distances (Figs. S8c-d, S9c-d, S10c-d, and S11c-d). To better visualize the 361 locations of these major low-velocity zones with respect to the array configurations and 362 fault surface traces, we depict the core and entire range of these zones in Figure 2. In 363 general, locations of the identified low-velocity zones correlate well with surface traces 364 of the 2019 Ridgecrest earthquake sequence (Fig. 2) inferred from Xu et al. (2020). This suggests that these low-velocity zones are likely indicative of fault damage zones or 365 366 sedimentary basins around fault segments.

367

368 3.3 Fault-zone trapped waves

369 A low-velocity fault-damage zone that is sufficiently uniform over a given distance 370 can act as a waveguide and generate, in addition to delay times and motion amplification, 371 trapped waves resulting from constructive interference of critically reflected phases 372 within the waveguide (e.g., Ben-Zion & Aki 1990; Igel et al., 1997; Jahnke et al., 2002). 373 Such waves have been observed at many locations, including the San Jacinto fault zone 374 (e.g., Lewis et al., 2005; Qin et al., 2018; Qiu et al., 2017; Share et al., 2019; Wang et al., 375 2019), the Parkfield section of the San Andreas fault (e.g., Li et al., 1990; Lewis and Ben-376 Zion, 2010; Ellsworth & Malin, 2011), and various other faults in California, Japan, Italy, 377 Turkey, and other places. Catchings et al. (2016) used peak ground velocities of P and S 378 waveforms recorded by cross-fault linear arrays to infer the location and width of the 379 West Napa-Franklin fault zone. Similarly, we find fault-damage-zone-related 380 amplification in data recorded by the B-arrays (e.g., Figs. 7 and S12) and use such

amplification to detect FZTW candidates. In this section, we first infer the location and
width of fault damage zones that produce FZTWs using waveforms for the fault-parallel
component, and then use waveforms of these candidates recorded by array B4 to invert

384 for properties (e.g., width, velocity, and attenuation) of the local fault-zone waveguide.

385

# 386 3.3.1 Methodology

387 Figure 7a shows S waveforms recorded on the fault-parallel component of array B4 for an example event (Fig. 1). Following Ben-Zion et al. (2003) and Qiu et al. (2017, Fig. 388 389 S6), several preprocessing steps are applied to the data prior to FZTW analyses. These 390 include instrument response removal, integration to displacement seismogram, tapering between 1 s before and 2 s after the S pick, and convolution with  $1/t^{1/2}$  to convert a point-391 392 source response to that of an equivalent SH line dislocation source (e.g., Igel et al., 2002; 393 Vidale et al., 1985). Compared to the raw recordings (Fig. S12), clearer Love-type 394 waveform sections with large amplitudes and lower frequency are found at a group of 395 stations (stations B416-423; Fig. 7a) in the southwest part of the array.

396 Figure 7b displays distributions of peak ground velocities (PGV) and root mean 397 square amplitudes (RMS) of the fault-parallel-component S waveforms, normalized by 398 the maximum value of the entire array. Large values of PGV and RMS are seen at 399 stations with FZTW, with considerably higher amplitudes than at the rest of the array. We 400 define the likelihood of a FZTW recorded by a station as the multiplication of PGV and 401 RMS (Fig. 7b). The likelihood curve measured for each event is normalized by the 402 maximum value of the entire array. After averaging the likelihood curves over all 403 analyzed events, we highlight the array sections that have high likelihood of recording 404 FZTWs (Fig. 8); these sections can be used to infer the location and width of fault-zone 405 waveguides.

Since FZTWs are observed in *S* waveforms recorded at stations B416-B423 of array B4 for the example event in Figure 7a, we can also identify candidate events with similar good-quality FZTWs through template matching, i.e., cross-correlating the fault-parallel component *S* waveforms recorded by stations B416-B423 for each event with those of the example event. The trapped waves of candidate events (Fig. 9a) that yield crosscorrelation coefficients higher than 0.85 (e.g., Fig. 9b) are averaged (Fig. 9c) and inverted for properties (e.g., width, shear velocities, and attenuation) of the average fault-zone
waveguide using a genetic inversion algorithm (e.g., Ben-Zion et al., 2003; Lewis et al.,
2005; Qiu et al., 2017).

415 We test a total of 10,000 models (50 generations and 200 models per generation) to 416 obtain a good estimate of the fault-zone parameters in the inversion. Parameters of the 417 best-fitting model and the 2,000 models (investigated in the last 10 generations) are 418 extracted from the inversion. Because there are strong trade-offs between model 419 parameters governing FZTWs (e.g., Ben-Zion, 1998; Jahnke et al., 2002; Peng et al. 420 2003), a successful inversion not only yields good waveform fits but also shows 421 consistency between parameters of the best-fitting model and peaks of the probability 422 density distributions of parameters developed in the last 10 generations. Additional 423 details on the method can be found in section 3.4 of Qiu et al. (2017) and Ben-Zion et al. 424 (2003).

425

# 426 *3.3.2 Results*

427 Figure 8 shows the distributions of FZTW likelihood values as a function of station 428 location, estimated at arrays B1-B4 for all events within the red box shown in Figure 1. 429 The likelihood values averaged over all analyzed events are representative of the in-situ 430 amplification in S-wave. A good correlation is found in locations between local maxima 431 of the likelihood curve and slower inner cores of the low-velocity zones identified from 432 the local *P*-wave delay time analysis (Section 3.2; Fig. 8). This suggests that the observed 433 inner cores of these low-velocity zones not only significantly delay the P arrivals but also 434 amplify the incoming S waves.

435 However, not all stations within the inner cores of low-velocity zones identified in 436 Figure 6 yield the same peak value of FZTW likelihood (Fig. 8). At array B4, the location 437 of the candidate fault-zone waveguide (B417-B422; Figs. 7a and 8d) is in good 438 agreement with the inner core of low-velocity zones that yields the highest likelihood 439 value. We do not detect candidate FZTW at stations within the other low-velocity zones 440 with much lower peak likelihood values. This is consistent with detailed fault-zone 441 studies at Parkfield (Lewis & Ben-Zion, 2010), the rupture zone of the 1992 Landers 442 earthquake (Peng et al., 2003), and fault zones in Japan (Mamada et al. 2004; Mizuno et

al. 2008). These studies showed that various sections of fault zones produce delay times
and other signals of damaged rocks but are either too heterogeneous or have significant
segmentation between sources and receivers to generate trapped waves (e.g., Igel et al.
1997, 2002; Jahnke et al. 2002).

447 Not all analyzed events show S-wave amplification patterns that are consistent with 448 the averaged curve. We first detect candidate events that show high FZTW likelihood 449 values at stations B417-B422 (Fig. 8d) by cross-correlating the likelihood pattern 450 measured for each event (e.g., black curve in Fig. 7b) with that of the mean (e.g., red 451 curve in Fig. 8d). More than 600 events with cross-correlation coefficients higher than 452 0.95 are identified as FZTW candidates. We further select 33 events (Fig. 9a) that 453 produce high-quality FZTWs (e.g., Fig. 9b) from the candidate events, through template 454 matching of fault-parallel component S waveforms at stations B416-B423 (using FZTWs 455 shown in Fig. 7a as the template; Section 3.3.1). These high-quality FZTW candidates 456 show a consistent source-receiver path (top inset of Fig. 9a), indicating the depth of the 457 fault-zone waveguide is likely shallower than 5 km, and there is an optimal range of 458 incidence angle for injecting seismic energy into the fault damage zone beneath array B4 459 (e.g., Fohrmann et al. 2004).

460 Compared to FZTW observed from each candidate event (e.g., Fig. 9b), the stacked 461 recordings (Fig. 9c) yield much higher SNRs and can thus provide more reliable and 462 robust estimations of the average fault-zone waveguide properties. Figure 10 presents the 463 inversion results from modeling the stacked waveforms shown in Figure 9c (in red). The 464 best-fitting model yields good waveform fits (Fig. 10a) and suggests a simplified (i.e., 465 vertical rectangular-shaped) fault-zone waveguide with a width of ~280 m, Q value of 466  $\sim$ 30, and S-wave velocity  $\sim$ 80% of the surrounding host rocks (Fig. 10b). The estimated 467 propagation distance inside the waveguide is ~5.4 km (Fig. 10b). Because this includes a 468 propagation component along-strike (i.e., non-vertical incidence angle), it suggests a waveguide depth of ~4 km (=  $5.4/\sqrt{2}$ ). The estimated average S-wave velocity in the 469 470 host rock is ~4.1 km/s, with the northeast block being ~4.2% faster, consistent with 471 results from the *P*-wave delay time analysis at array B4 (~3% velocity contrast across 472 LVZ#5 with higher velocity in the northeast; Fig. 6g). The parameters of fault-zone 473 models from the last 10 generations (2,000 models) are shown in Figure 10b, along with

the corresponding probability density functions computed as the frequency of each parameter value weighted by the fitness values (Ben-Zion et al., 2003). The good waveform fits, combined with the consistency between best fitting parameters and peaks of the probability density functions (Fig. 10b), suggest that the best-fitting model parameters provide robust estimates of the average properties of the fault-zone waveguide.

480 In addition to FZTWs, we detect in data of stations B420 and B422 (Fig. S13) clear 481 fault-zone head waves (FZHWs) that refract along a bimaterial fault interface (e.g., Ben-482 Zion, 1990) and arrive ~0.1 s earlier than the direct *P*-wave. The early-arriving FZHWs 483 are inferred from horizontal particle motion analysis modified from the method of Bulut 484 et al. (2012). It is important to note that in the polarization analysis we only focus on the 485 rotation of the horizontal particle motion between FZHW and direct P-wave. This is 486 because the polarization direction of horizontal particle motion for the direct P wave at stations inside a damage zone (Fig. S13) can deviate significantly from the source-487 488 receiver back azimuth. Since the differential time between the FZHW and direct *P*-wave 489 decreases significantly from the northeast (B422) to the southwest (B420) over a short 490 distance (~0.1 km), the observed FZHW is likely traveling along a local interface that is associated with the edge of the damage zone (e.g., Qiu et al., 2017) on the northeast side 491 492 between stations B422 and B423. Similar FZTW and FZHW signals are also clearly 493 observed in the data of array B2 (e.g., Fig. S14) but not for arrays B1 and B3.

494

### 495 **4. Discussion**

496 We analyze systematically time delays of P arrivals from teleseismic and local 497 earthquakes, S-wave amplification, and fault zone trapped waves in data recorded by four 498 long-aperture (4-8 km) arrays at different locations along the rupture zone of the 2019 499 Mw 7.1 Ridgecrest earthquake (Fig. 2). The analyses allow us to derive information on 500 velocity contrast interfaces and low velocity damage zones at scales of ~100 m associated 501 with the Ridgecrest rupture zone. The results are complementary to tomographic velocity 502 models that provide a regional context but do not resolve the internal structural elements 503 of the Ridgecrest rupture zone imaged in this work.

504 The delay times of *P*-waves from both teleseismic and local earthquakes, after proper 505 corrections for propagation and topography effects, show clear and consistent velocity 506 contrasts across the rupture zone, with a higher velocity in the northeast at arrays B1 507 (Figs. 4a and 6a) and B4 (Figs. 4d and 6g). Combining the delay times inferred from 508 teleseismic P waves in Figure 4 and the amplitudes of cross-fault velocity contrast in the 509 left panels of Figure 6, we can estimate the depth extent of the contrast interface beneath 510 each array. For a constant velocity contrast across the fault with an amplitude of  $\eta$ 511 extending to a depth of  $H_c$ , the time delay between teleseismic P arrivals, assuming a 512 vertical angle of incidence, at stations on different sides of the fault is given by

$$\Delta t \approx \eta \cdot H_c / \overline{V_p}, \tag{5}$$

where  $\overline{V_p}$  is the average *P*-wave velocity of the upper crust. We set  $\overline{V_p} = 5.6$  km/s as inferred from the mean of the average slowness  $\overline{s_j}$  (eq. 3) over all event-array pairs. The *H*<sub>c</sub> beneath arrays B1 ( $\Delta t = 0.15$  s from Fig. 4a and  $\eta = 5.3\%$  from Fig. 6a) and B4 ( $\Delta t = 0.25$  s from Fig. 4d and  $\eta = 8.8\%$  from Fig. 6g), computed using equation (5), are both ~16 km, suggesting a deep fault interface that likely extends to the seismogenic depth (10-12 km) after taking measurement uncertainties (Figs. 4a and 4d) and nonvertical angle of incidence into account.

520 We do not estimate the depth  $H_c$  for arrays B2 and B3 as the teleseismic *P*-wave 521 delay times  $\Delta t$  are highly dependent on the topographic correction (Section 3.1.2; Figs. 522 4b-c). The cross-fault velocity contrast observed in local *P*-wave delay time patterns 523 resolved at arrays B2 and B3 (Figs. 6c and 6e) suggests a higher velocity in the 524 southwest, the opposite of the polarity obtained at arrays B1 and B4 (Figs. 6a and 6g). 525 Further analysis of the transition zones (Figs. 2b-e) resolved from the modeled cross-fault 526 velocity contrasts (left panels of Fig. 6) indicates that: the observed velocity contrast 527 beneath array B2 is likely an artifact due to insufficient topographic correction; array B3 528 may detect the velocity contrast across the rupture zone of the Mw 6.4 earthquake; and 529 the broader transition zone beneath array B4 suggests the fault zone at the southeast edge 530 of the Mw 7.1 earthquake rupture is likely not as well localized (e.g., several rupture 531 surface traces; Fig. 2) as those beneath the other three arrays.

532 Symmetry properties of fault damage zones with respect to the main slip surface can 533 provide information on the direction of earthquake ruptures (e.g., Ben-Zion & Shi, 2005; 534 Dor et al., 2006; Mitchell et al., 2011; Xu et al., 2012). Statistically preferred rupture 535 direction is expected for prominent bimaterial faults (e.g., Ampuero & Ben-Zion, 2008; 536 Andrews and Ben-Zion, 1997; Shlomai & Fineberg, 2016; Weertman, 1980). However, a 537 prominent large-scale velocity contrast is not expected across the Ridgecrest rupture zone 538 based on the regional velocity models (Lee et al., 2014; Shaw et al., 2015; White et al., 539 2020) and total offset across the structure. The transition zones of the cross-fault velocity 540 contrast resolved at arrays B1 (Fig. 2b) and B3 (Fig. 2d) are asymmetrically distributed to 541 the faster side of the fault surface trace, which is likely indicative of asymmetric rock 542 damage offset to the faster side of the rupture zone for the Mw 7.1 and Mw 6.4 543 earthquakes, respectively. Since the transition zones resolved at arrays B2 and B4 are not 544 well resolved, we do not discuss results obtained at these two locations (Figs. 2c and 2e).

545 Good agreement between locations of slower inner cores of the low-velocity zones 546 identified from delay time analyses and the group of stations with amplified S waveforms 547 (Fig. 8) is found beneath all four arrays. Combined with the narrow width (300-400 m; 548 right panels of Fig. 6) of the inner cores, these zones that both significantly delay the P 549 arrivals and amplify the incoming S waveforms are likely representative of localized 550 damage zones, rather than variations in sedimentary thickness that usually yield much 551 wider patterns of S-wave amplification. Waveform modeling of FZTWs detected at array 552 B4 yields good waveform fits and an average waveguide with fault-zone parameters 553 comparable to those inferred from previous studies in the SJFZ (Oin et al., 2018; Oiu et 554 al., 2017; Share et al., 2017, 2019): width of ~300 m, Q of ~30, S-wave velocity 555 reduction of ~20% inside the damage zone, and depth extent of 3-5 km (Fig. 10).

To further investigate the structures that produce the observed low-velocity zones, we map the inner core and entire span of these zones together with locations of fault surface traces (or their extrapolations) and the surface geology compiled by the USGS in Figure 2. The low-velocity zone (LVZ#1 in Fig. 6b) identified at array B1, in contrast to the resolved transition zone (Fig. 2b), is on the southwest (slower) side of the surface trace of the main rupture zone. There are two well-separated inner cores (Fig. 2b) that yield both large delay times (Fig. 6b) and *S*-wave amplitudes (Fig. 8a). The inner core in the

southwest (Fig. 2b) is likely associated with the damage zone of past earthquake ruptures, since it is centered on a surface displacement mapped before the 2019 Ridgecrest earthquake. The northeast inner core (Fig. 2b), however, is likely representative of the damage zone of the Mw 7.1 Ridgecrest earthquake ruptures, as it is centered on the main rupture zone.

568 At array B2, two low-velocity zones (LVZ #2 and #3) are identified from Figures 6d. 569 Different from the other three arrays, the surface geology beneath array B2 consists of 570 two different types of rocks (alternating playa and igneous rocks; Fig. 2c), which likely 571 explains the weak cross-fault velocity contrast resolved at this location (Figs. 4b and 6c). 572 Both low-velocity zones are centered on the surface traces (or their extrapolations) of the 573 main rupture zone, with the slower inner cores close to the fault traces. Good agreement 574 between locations of low-velocity zones and surface traces of the 2019 Ridgecrest 575 earthquake ruptures is also observed beneath arrays B3 (Fig. 2d) and B4 (Fig. 2e). Since 576 the locations of these low-velocity zones correlate well with the fault surface traces rather 577 than the surface geology, we conclude that: these low-velocity zones with slower inner 578 cores that considerably amplify S waveforms depict the fault damage zones of the Mw 579 7.1 earthquake rupture at arrays B2 and B4, and ruptures of both the Mw 6.4 and Mw 7.1 580 earthquakes at array B3.

581 It is interesting to note that we find a clear and coherent secondary phase between P 582 and S arrivals in waveforms recorded by part of array B4 for more than 10 events located 583 beneath the array. Figure S15b shows such signals for an example event (circle in Fig. 584 S15a). The strong phases are visible at stations B423-B457 (blue curve in Fig. S15b) and 585 correlate well with the shape of the direct *P*-wave but with the opposite first-motion 586 polarity (e.g., Fig. S15d) suggesting the signals are likely reflected or converted waves. 587 Since the array-mean travel time of the direct P-wave is between 1-1.5 s, based on the 588  $\sim 0.5$  s differential time between the direct P and the strong secondary waves (e.g., black 589 and blue dashed lines in Fig. S15d), the reflection or conversion interface should be a few 590 kilometers away from the array and is thus likely associated with the Garlock fault in the 591 south rather than the rupture zone of the Mw 7.1 Ridgecrest earthquake (Fig. S15a). 592 Detailed analysis of these candidate reflected/converted phases may require more 593 sophisticated techniques, such as seismic migration of the arrival times and waveform

594 simulation with accurate source information (e.g., focal mechanism) for the 595 understanding of the wave amplitudes. As these secondary signals detected in Figure S15 596 are outside the scope of this paper, additional analysis of these phases will be the subject 597 of a follow-up study.

598

# 599 **5.** Conclusions

The rupture zone of the Mw 7.1 Ridgecrest earthquake is shown to have seismic velocity contrast interfaces and damage structures with significant along-strike variations of local damage zones. The results are in overall agreement with and complementary to the mapped fault surface traces in the region (e.g., Xu et al., 2020) and previous analyses of aftershocks in the Ridgecrest sequence (e.g., Ross et al., 2019; Cheng & Ben-Zion, 2020).

606Delay times of P waves from teleseismic and local earthquakes recorded by arrays B1607and B4 resolve velocity contrasts ranging from 5%-9% in Vp across the Ridgecrest608rupture zone, which extend to a depth of ~16 km, with the northeast being locally faster.609Data recorded by array B3, crossing the surface ruptures of both the Mw 6.4 and Mw 7.1610events, likely detects a ~4% velocity contrast in Vp across the fault that hosted the Mw6116.4 earthquake, with the northwest side having higher velocity. Asymmetric rock damage612concentrated on the stiffer side of the fault is seen beneath arrays B1 and B3.

613 Low-velocity zones that further delay the P waves of local seismic events are 614 centered on mapped surface traces of faults associated with both past and the 2019 615 Ridgecrest earthquake ruptures. Considerable amplification is seen consistently in S 616 waveforms recorded at stations within the slower inner core of the identified LVZ. Clear 617 FZTWs are identified at arrays B2 and B4, and inversion of high-quality FZTWs at array 618 B4 indicates a waveguide with average properties comparable to those found in previous 619 studies at sites of the San Jacinto fault zone. Phases identified as FZHWs, associated with 620 the northeast boundary of the fault zone waveguide, are observed at arrays B2 and B4. 621 Coherent candidate reflected or converted waves are observed in data recorded at array 622 B4. Additional analysis of these phases can provide additional constraints on prominent 623 impedance interfaces associated with the Garlock fault and/or the Ridgecrest rupture 624 zone.

625

# 626 Data Availability

627 The digital data are available in mseed day volume format, with each component in a 628 separate volume. The data samples are 4 byte floats and consistently sampled at 500 629 samples/second. Data described in this report are available from the IRIS Data 630 Management Center (https://ds.iris.edu/ds/nodes/dmc/data/). An accompanying report for 631 data acquisition is available from Catchings et al. (2020). Surface traces of the 2019 632 Ridgecrest earthquake sequence shown in this study were determined by Xu et al. (2020). 633 The geologic data depicted as the background of Figure 2 were obtained from the USGS 634 website (https://mrdata.usgs.gov/geology/state/state.php?state=CA).

635

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852 Figure 1. Location map for the 2019 Ridgecrest earthquake sequence (colored 853 circles, square, and stars) and four linear arrays (B1, B2, B3, and B4 as red, green, 854 blue, and purple triangles, respectively) analyzed in this study. The catalog of 855 Hauksson et al. (2012, extended to 2019) is used for earthquake locations, with color 856 representing the focal depth (colorbar). Epicenters of the 2019 Mw 6.4 and Mw 7.1 857 Ridgecrest earthquakes are marked as stars. Bandpass filtered waveforms, fault zone 858 head waves, and fault zone trapped waves of an example event (orange square) 859 recorded at array B4 are shown in Figs. S12, S13a, and 7a. Fault surface traces are 860 depicted as black lines with ruptures of the 2019 Ridgecrest earthquake sequence 861 being thicker. Near-fault seismic events outlined by the red box are analyzed in 862 sections 3.2 and 3.3. The background gray colors indicate the local topography. 863 WLSZ - Walker Lane Shear Zone; ECSZ - Eastern California Shear Zone; EF -Elsinore Fault; GF – Garlock Fault; SAF – San Andreas Fault; SJF – San Jacinto 864 865 Fault.

866 Figure 2. Zoomed-in maps of all the Ridgecrest linear arrays (triangles) in (a) and 867 B1-B4 in (b)-(e), respectively. Color of the map background depicts the surface 868 geology (i.e., rock types; legend in (a)) extracted from the USGS website (see Data 869 Availability). The station elevation is illustrated by the color of the triangles in (b)-870 (e). The green circle and black lines denote the center of the array and surface traces 871 of faults associated with the 2019 Mw 7.1 and Mw 6.4 Ridgecrest earthquakes (red 872 and blue stars), respectively. The green bars outline the span of the low-velocity 873 zones identified in the local P-wave delay time analysis (Section 3.2), whereas the red 874 bars illustrate the region of the slower inner cores (Fig. 6). The blue arrows depict the 875 extent of the transition zone pointing from the slower to faster crustal blocks (black 876 solid curves in Fig. 6).

Figure 3. (a) Locations of array B4 (triangle) and four teleseismic events analyzed in this study (stars), with the red star indicating the event analyzed in (b)-(c). (b) Pwaves recorded on vertical-component sensors of array B4 for the target teleseismic event are illustrated in colors with red and blue indicating positive and negative values. The *P*-wave arrival time predicted from the model IASP91 is used to align the 882 P waveforms and is set to be zero in the time axis. The P waveforms are bandpass 883 filtered twice. After applying a bandpass filter between 0.5 Hz and 2 Hz, the array-884 mean envelope function and an array-mean P-wave pick are computed and depicted 885 as the curve and the vertical solid line in black, respectively. (c) Amplitude spectrum 886 averaged over the entire array with the red star and horizontal dashed lines indicating 887 the peak frequency and median of the amplitude spectrum between 0.5 Hz and 2 Hz, 888 respectively. Then, a second bandpass filter between the frequency range outlined by 889 the vertical dashed lines in (c) is applied. The red dashed curve in (b) depicts the 890 teleseismic P-wave delay times, measured using the P waveforms tapered between 891 the vertical dashed lines ( $\pm$  one dominant period relative to the array-mean *P*-wave 892 pick).

893 Figure 4. Teleseismic P-wave delay times for arrays (a) B1, (b) B2, (c) B3, and 894 (d) B4. The colored stars indicate P-wave delay times measured from different 895 teleseismic events and are labeled in the legend by the corresponding peak frequency 896 of the array-mean P-wave amplitude spectrum. The black dots depict the delay-time 897 pattern averaged over all teleseismic events, with error bars representing the standard 898 deviation of the mean. The blue and red dashed curves illustrate the delay times after 899 a topographic correction, assuming constant *P*-wave velocities of 2 km/s and 4 km/s. 900 The green bars outline the span of the low-velocity zones identified in local *P*-wave 901 delay time analysis (Section 3.2), while the red bars illustrate the region of the slower 902 inner cores (Fig. 6).

903 Figure 5. (a) P waveforms of an example local seismic event, shown as the blue 904 star in (c), recorded on vertical-component sensors of array B1. Waveform at each 905 station is normalized by its corresponding maximum amplitude and bandpass filtered 906 between 0.5 Hz and 20 Hz. Red stars denote the automatic *P* picks. (b) Same as (a) 907 but shown in slowness domain, i.e., the time axis of each station is divided by the 908 corresponding hypocentral distance. (c) Distribution of seismic events that pass the 909 selection criteria (Section 3.2.1) for array B1 (red triangles) are shown as circles with 910 color representing the focal depth. Events used in the delay time analysis (Section 911 3.2) for array B1 are outlined by the red box (Area1) and the results are shown in

Figs. 6a-b. Data from events outside the red box (Area2 and Area3) are also analyzed
(Figs. S8c-d). The black lines represent fault surface traces. Earthquakes that are
excluded from the delay-time analysis at array B1 are shown as grey dots.

915 Figure 6. Statistical analysis of local *P*-wave arrival patterns for all four arrays. 916 (a) Red dots illustrate the *P*-wave relative slowness variation within array B1, 917 averaged over 189 local seismic events with depth > 5 km and along-fault distance < 918 5 km relative to the array. The error bars represent a range of three standard 919 deviations of the mean about each respective mean value. The solid black lines depict 920 the contribution associated with a *P*-wave velocity contrast (~5.3% faster in the NE) 921 across the fault beneath the array. The corresponding mean P-wave delay time 922 variation is illustrated as the blue curve with dashed blue curves indicating the 923 uncertainties. (b) The P-wave delay time variation after removing the modeled cross-924 fault velocity contrast (the solid black curve in (a)) is shown in blue for array B1. The 925 green dashed vertical lines outline the identified low velocity zone (LVZ), where P 926 waves are significantly delayed with respect to the background level (black dashed 927 line). The red dashed vertical lines denote the inner core of the LVZ, i.e., the 928 amplitude of the blue curve is large whereas its gradient is relatively small. The gray-929 scale heat map of values for relative slowness and delay times after the cross-fault-930 velocity-contrast correction obtained at each station for all analyzed events is 931 illustrated in (a) and (b), respectively. PDF – Probability Density Function. Results 932 for arrays B2, B3, and B4 are shown in (c)-(d), (e)-(f), and (g)-(h), respectively.

933 Figure 7. (a) Fault zone trapped waves (FZTWs) following the S-wave arrivals for 934 an example event (square in Fig. 1) observed at the fault-parallel component of array 935 B4. The waveforms are preprocessed following the steps of Figure S6 of Qiu et al. 936 (2017). The blue bar outlines the stations with FZTWs. (b) Red dots and blue stars 937 denote the distributions of normalized peak ground velocities (PGV) and root mean 938 square amplitudes (RMS) of the S waveforms shown in (a). The black curve 939 represents the likelihood of FZTW, i.e., the normalized multiplication of PGV and 940 RMS values, and is used to identify FZTW candidates.

941 Figure 8. S-wave amplification analysis for arrays B1-B4 in (a)-(d). Gray-scale 942 heat maps of FZTW-likelihood values (e.g., black curve in Fig. 7b) computed at each 943 station in the array for all analyzed events are shown. Red dots depict the likelihood 944 curve averaged over all analyzed events. Error bars represent a range of two standard 945 deviations of the mean about each respective mean value. The bottom green bar 946 marks the low velocity zone (LVZ) identified from the local P-wave delay time 947 analysis (green dashed lines in Fig. 6), whereas the red bar illustrates the slower inner 948 core of the LVZ (red dashed lines in Fig. 6). The blue bar in (d) outlines the stations 949 with FZTWs at array B4 (blue bar in Fig. 7a).

950 Figure 9. (a) Locations of earthquakes (gray dots) analyzed in section 3.3. FZTW candidates identified through template matching, with correlation coefficient greater 951 952 than 0.85 marked as stars and color representing the focal depth. Red triangles denote 953 the location of array B4. The along-fault cross section of seismicity (dots and stars) 954 and array B4 (triangle) are shown in the top inset. (b) FZTWs recorded at stations 955 B416-B423 for nine high-quality candidate events (red) with the highest correlation 956 coefficients. The template waveforms (Fig. 7a) are shown in black. The array-mean S 957 pick and correlation coefficient of each candidate event are labeled in the top left. (c) 958 Comparison between FZTWs of the reference event (in black) and those averaged 959 over all the high-quality candidate events (in red) observed between stations B416-960 B423.

961 Figure 10. Inversion results for FZTWs observed between stations B416-423, 962 averaged over candidates shown in Figure 9a. (a) Comparison between synthetic 963 waveforms (red) computed using the best-fitting model parameters (black dots in (b)) 964 and the observed FZTWs (in black). (b) Fitness values of fault-zone model 965 parameters from the last 10 generations of the inversion (green dots). The best-fitting parameters (black circles) are displayed in each panel and used to generate the 966 967 synthetic waveforms shown in (a). Black curve indicates the probability density function of model parameters shown as green dots. 968

969

Figure 1.



Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.



Figure 9.



Figure 10.

