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**Structure of the Mn Complex in Photosystem II:  
Insights from X-ray Spectroscopy**

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## Abstract

We have used Mn K-edge absorption and K $\beta$  emission spectroscopies to determine the oxidation states of the Mn complex in the various S-states. We have started exploring the new technique of resonant inelastic X-ray scattering spectroscopy (RIXS); this technique can be characterized as a Raman process that uses K-edge energies (1s to 4p, ~6550 eV) to obtain L-edge-like spectra (2p to 3d, ~650 eV). The relevance of these data to the oxidation states and structure of the Mn complex is presented. We have obtained EXAFS data from the S<sub>0</sub> and S<sub>3</sub> states and observed heterogeneity in the Mn-Mn distances, leading us to conclude that there may be three rather than two di- $\mu$ -oxo bridged units present per tetranuclear Mn cluster. In addition, we have obtained data using Ca/Sr X-ray spectroscopy that provide evidence for a heteronuclear Mn/Ca cluster. The possibility of three di- $\mu$ -oxo-bridged Mn-Mn moieties and the proximity of Ca is incorporated into developing structural models for the Mn cluster. The involvement of bridging and terminal O ligands of Mn in the mechanism of oxygen evolution is discussed in the context of our X-ray spectroscopy results.

## 1. Introduction

Critical questions related to the process of photosynthetic water oxidation catalyzed by a Mn cluster in the oxygen-evolving complex (OEC) of photosystem II (PS II) are: 1) What are the oxidation state(s) and structural changes in the Mn complex as the OEC proceeds through the S-state cycle? and (2) What is the mechanism by which four electrons are removed from two water molecules by the Mn complex to produce an O<sub>2</sub> molecule? We have addressed these questions principally by the use of X-ray absorption near edge structure (XANES), K $\beta$  X-ray emission spectroscopy (XES) and extended X-ray absorption fine structure (EXAFS) along with electron paramagnetic resonance spectroscopy (Robblee et al. 2001; Yachandra et al. 1996).

X-ray spectroscopy is element specific and can monitor Mn (or Ca/Sr) directly in the membrane; hence, it is a good method for studying the structure of the Mn-OEC without interference from pigment molecules, the lipid and protein matrix, or other metals such as Ca, Mg, Cu, or Fe which are present in active PS II preparations. Mn K-edge XANES and K $\beta$  XES provide information about oxidation states and the site symmetry of the Mn complex, and EXAFS at the Mn/Sr/Ca K-edge furnishes information about the number, type and distances to neighboring ligand atoms in the S<sub>0</sub>, S<sub>1</sub>, S<sub>2</sub> and S<sub>3</sub> states of the OEC (Yachandra 1995; Yachandra & Klein 1996). X-ray spectroscopy does not require long-range order; the structural studies can be performed on frozen solutions. Several of the intermediate states mentioned above can be stabilized as frozen solutions. Few other spectroscopic techniques provide such specificity for studying the structure of Mn in the OEC.

This presentation focuses on the application of X-ray spectroscopic methods to resolve structural questions regarding the Mn cluster in the OEC and we present our view of the structure

of the Mn cluster which is refined based on new data, and a mechanism for water oxidation that emphasizes the recent results from our laboratory.

## 2. Oxidation States of the Mn cluster

A key question for the understanding of photosynthetic water oxidation is whether the four oxidizing equivalents generated by the reaction center are accumulated on the four Mn ions of the OEC during S-state turnover, or whether a ligand-centered oxidation takes place, especially, before the formation and release of molecular oxygen during the  $S_3$  to ( $S_4$ ) to  $S_0$  transition. We have addressed these questions by using Mn K-edge XANES (1s-4p absorption),  $K\beta$  XES (3p-1s emission) and the recently introduced resonant inelastic X-ray scattering spectroscopy (RIXS) (1s to 3d/4p absorption followed by 2p-1s  $K\alpha$  emission) to obtain L-edge-like spectra (2p-3d absorption).

A promising approach to study the Mn oxidation states in the native S-states is to step samples through the S-state cycle by the application of saturating single-turnover flashes and to characterize these samples by X-ray spectroscopy. We reported earlier XANES data from samples in the  $S_0$  through  $S_3$  states produced under these physiologically relevant conditions (Roelofs et al. 1996). The flash-advanced PS II samples prepared from spinach exhibited XANES spectra that showed Mn oxidation from  $S_0$  to  $S_1$  and from  $S_1$  to  $S_2$ , but no further oxidation during the  $S_2$  to  $S_3$  transition. These results are consistent with our earlier studies by Guiles *et al.* that were prepared by illumination at low temperature, or by chemical treatment (Guiles et al. 1990a; Guiles et al. 1990b).

Other spectroscopic techniques like NMR-PRE (Sharp 1992), EPR (Ahrling et al. 1997; Dismukes & Siderer 1981; Hansson & Andréasson 1982; Messinger et al. 1997b), UV

absorption (Dekker 1992; Van Leeuwen et al. 1993) and  $Y_D^{\text{ox}}$  EPR/spin echo (Styring & Rutherford 1988) studies show that, in PS II, the  $S_0$  to  $S_1$  and  $S_1$  to  $S_2$  transitions are Mn-centered oxidations, and most groups agree that Mn redox states in the  $S_1$  are (III<sub>2</sub>, IV<sub>2</sub>) which are oxidized to (III, IV<sub>3</sub>) in the  $S_2$  state. There is, however, a long standing controversy as to whether a Mn- or ligand-centered oxidation occurs in the  $S_2$  to  $S_3$  transition (Iuzzolino et al. 1998; Ono et al. 1992).

To resolve the controversy we used high-resolution Mn K $\beta$  X-ray fluorescence (XES) studies in collaboration with Dr. Steve Cramer, who has constructed an unique high-resolution emission spectrometer that operates at the Mn K $\beta$  fluorescence energy and has demonstrated the feasibility of determining the oxidation states of Mn (Bergmann & Cramer 1998). We have completed a comprehensive Mn K $\beta$  XES study in conjunction with XANES on sets of samples prepared in a similar manner and characterized by EPR (Messinger et al. 2001).

Because the actual S-state composition is of critical importance to the interpretation of the results, we characterized the samples by EPR spectroscopy, using the  $S_2$  multiline EPR signal (MLS) as a direct measure for the amount of  $S_2$  in our samples (Fig. 1). The relative S-state populations in samples given 0, 1, 2, 3, 4, 5 or 6 flashes were determined from fitting the flash-induced EPR multiline signal oscillation pattern to the Kok model (Kok et al. 1970) for each of the samples used in the X-ray spectroscopy experiments. It is essential to obtain high concentrations of PS II in the S-states to obtain EPR, XANES and K $\beta$  spectroscopy data with a good S/N ratio. A Nd-YAG laser system (Spectra-Physics PRO 230-10, 800 mJ/pulse at 532 nm, 9 ns pulse width) was used to illuminate PS II samples from both sides simultaneously. Because the pulse widths are narrow compared to those of a flash lamp, double hits become negligible. The damping of the oscillation pattern is minimized, thus making it possible to derive an unique

set of spectra by deconvolution. When the system becomes scrambled by frequent misses and double hits, the deconvolved spectra are not unique.

a) Mn K-edge XANES Our edge data are of sufficient S/N quality to allow for not only a determination of the absorption edge position (using the inflection point energy, defined as the zero-crossing in the second derivative of the edge), but also an analysis of the structure on the edge as a function of the number of flashes.

The edge spectra of samples given 0, 1, 2, or 3 flashes are combined with EPR information to calculate the pure S-state edge spectra. The shifts in inflection point energy positions and changes in shape are determined by second derivatives of the K-edge spectra and are shown in Fig. 2. In addition to the shift in edge position, the  $S_0$  to  $S_1$  (2.1 eV) and  $S_1$  to  $S_2$  (1.1 eV) transitions are accompanied by characteristic changes in the shape of the edge, also indicative of Mn-oxidation. The edge position shifts very little (0.3 eV) for the  $S_2$  to  $S_3$  transition, and the edge shape shows only minor changes (Messinger et al. 2001).

b) K $\beta$  X-ray Emission Spectroscopy (XES) High-resolution Mn X-ray emission spectroscopy was performed on  $S_0$ ,  $S_1$ ,  $S_2$  and  $S_3$  samples prepared by using high-power laser flashes as described above. In this technique, the energy of the Mn K $\beta$  emission ( $3p \rightarrow 1s$ ) is measured with a high-resolution dispersive spectrometer. The shape and energy of the K $\beta_{1,3}$  emission reflect the oxidation state(s) of the emitting Mn atom(s). The emission occurs from a 3p level that is mainly influenced by the number of unpaired 3d electrons and is less sensitive to the symmetry and bonding than the K-edge absorption, which involves transitions to the 4p level (Peng et al. 1994). The spin of the unpaired 3d valence electrons can be either parallel (K $\beta'$ ) or antiparallel (K $\beta_{1,3}$ ) to the hole in the 3p level. The splitting between the K $\beta'$  and K $\beta_{1,3}$  peaks becomes smaller for higher oxidation states because fewer 3d electrons interact with the 3p hole.

Thus, in contrast to the inflection points of the XANES edges, the  $K\beta_{1,3}$  peaks shift to lower energy with higher oxidation states.

Using a 1<sup>st</sup>-moment analysis, the position of the main  $K\beta_{1,3}$  peak has been calculated for each S-state. Based on the shifts, or lack thereof, of the 1<sup>st</sup> moments, we concluded that Mn is not oxidized during the  $S_2 \rightarrow S_3$  transition. Our results (shown as difference spectra in Fig. 2) show a shift in energy between the  $S_0$  and  $S_1$  state spectra and between the  $S_1$  and  $S_2$  state spectra and seen in the derivative-shaped difference spectra. However, there is very little change between the  $S_2$  and  $S_3$  states, as seen in the difference spectrum (Messinger et al. 2001).

We have used two independent spectroscopic techniques,  $K\beta$  and XANES spectroscopy, along with EPR, to examine the Mn redox states in PS II. Figure 3 summarizes the oscillation patterns of the 1<sup>st</sup> moments from the  $K\beta_{1,3}$  spectra and of the XANES spectra inflection point energies. Both patterns show large shifts between the 0F and 1F samples, which reflect the Mn-centered oxidation during the  $S_1 \rightarrow S_2$  transition. A much smaller shift occurs between 1F and 2F samples, which is inconsistent with a Mn-centered oxidation during the  $S_2 \rightarrow S_3$  transition. The 2<sup>nd</sup> derivatives of the S-state XANES spectra and the  $K\beta$  difference spectra (Fig. 2) confirm this conclusion. It is unlikely that both techniques would fail to detect a Mn-centered oxidation for the  $S_2 \rightarrow S_3$  transition if it did, in fact, occur. Support for this conclusion has also come from XANES and  $K\beta$  studies of two sets of structurally homologous Mn compounds in different oxidation states (Visser et al. 2001). We conclude that probably no direct Mn-oxidation is involved in this transition. The proposed Mn oxidation state assignments are as follows:  $S_0$  (II, III, IV, IV) or (III, III, III, IV);  $S_1$  (III, III, IV, IV);  $S_2$  (III, IV, IV, IV);  $S_3$  (III, IV, IV, IV) (Messinger et al. 2001).

c) Resonant Inelastic X-ray Scattering Spectroscopy - RIXS Mn L-edge spectra have several advantages over Mn K-edge spectra, because multiplet calculations can be applied to L-edge spectra which are sensitive to the number of holes in the 3d level, the oxidation state of Mn, and the spin state of Mn. This is a rich amount of information about oxidation states relative to what is available from K-edge XANES spectra. L-edge spectra have already been obtained for a number of Fe, Cu, and Ni metalloproteins, and have been simulated using multiplet calculations to provide quantitative conclusions about the oxidation states of the metals involved (Wang et al. 1997; Wang et al. 1996). However, the main disadvantage of using soft X-rays, where the Mn L-edge occurs (~650 eV), is that the samples are subjected to ultra high vacuum and are more prone to X-ray damage.

We have recently started exploring the technique of resonant inelastic scattering (RIXS) that circumvents the problems mentioned above in a very elegant manner. It is possible to combine the two types of K- and L-edge experiments by mapping the 1s2p resonant inelastic scattering (RIXS) landscape as shown in Fig. 4 in a simplified two-step process (Caliebe et al. 1998). RIXS involves exciting a pre-edge transition (1s to 3d) of the K shell using hard X-rays. Then, the high-resolution X-ray emission analyzer can be tuned to detect the  $K\alpha$  fluorescence of Mn, which is a 2p to 1s transition. The final state from this transition is a 2p hole, which is the same final state as the L-edge fluorescence (3d to 2p). An advantage of RIXS over both conventional techniques is that the resonances are separated according to the final states they decay into. As a result, the resonances appear as islands in the resonant landscape and their spectral shapes can be analyzed. A two-dimensional RIXS spectrum is obtained by plotting the scanned incident X-ray energy,  $\nu$  (1s to 3d), versus  $E_B$ , the binding energy, which is equal to  $(\nu - f)$ ,  $f$  is the energy of the scanned high resolution X-ray emission analyzer. Furthermore, this

technique offers an unprecedented opportunity to selectively excite into specific molecular orbitals with 3d character by tuning the incoming X-rays to look at a specific pre-edge transition; at least 2 – 3 resolvable pre-edge transitions are seen in the Mn K-edge spectrum of PS II. This means that it would be possible to obtain ‘nested’ RIXS spectra which correspond to fluorescence from different molecular orbitals with 3d character.

We have collected preliminary data from a PS II sample in the  $S_1$  and  $S_2$  states. The shape and ratio of the  $L_{III}/L_{II}$ -like peaks of the PS II RIXS spectrum are in between those for the Mn(III) and Mn(IV) complex, and very different from that for the Mn(II) complex. This indicates a mixed (III) and (IV) Mn oxidation state in the  $S_1$  state -  $(III_2,IV_2)$ ; a conclusion in accord with that derived from other studies. The cross sections of the 2D-RIXS spectra reveal the L-edge-like character, showing two well-resolved peaks due to spin-orbit coupled states  $J=3/2$  ( $L_{III}$ -like edge) and  $J=1/2$  ( $L_{II}$ -like edge). The ratio of the  $L_{III}/L_{II}$  peaks increases from (II) to (III) to (IV); the ratio for PS II changes from  $S_1$  to  $S_2$  in a manner consistent with an oxidation from  $(III)_2(IV)_2$  to  $(III)(IV)_3$ .

### 3. Structure of the Manganese Cluster

We have used Mn EXAFS to study the local structure of the Mn complex and have proposed topologically consistent models that can be reconciled with several structural arrangements of the four Mn atoms (DeRose et al. 1994). The proposed model consists of two or three di- $\mu$ -oxo bridged binuclear Mn units with Mn-Mn distances of  $\sim 2.7$  Å that are linked to each other by a mono- $\mu$ -oxo bridge with a Mn-Mn separation of  $\sim 3.3$  Å, and with Ca at a distance of 3.4 Å to Mn (DeRose et al. 1994; Yachandra et al. 1993). The Mn-Mn distances are largely invariant in the native  $S_1$  and  $S_2$  states (DeRose et al. 1994). EXAFS experiments on the

$S_3$  and  $S_0$  states have been difficult to perform on PS II samples using single flashes. The requirement of optically dilute samples to ensure saturation by single actinic flashes generally results in a low signal-to-noise ratio. We have recently been able to extend these studies to the  $S_3$  and  $S_0$  states prepared by flash illumination.

a)  $S_3$  State Earlier EXAFS data from the chemically treated  $S_3$  state samples produced by a double turnover method indicated that the two di- $\mu$ -oxo-bridged Mn-Mn dimer units may become nonequivalent (Guiles et al. 1990b). The EXAFS spectra from calcium-depleted  $S_3'$  samples ( $S_2Y_z'$ ) prepared by low pH treatment in a citrate buffer do not exhibit similar heterogeneity in 2.7 Å Mn-Mn distances (Latimer et al. 1998).

EXAFS data of  $S_3$  samples created under physiological conditions with saturating flash illumination show significant changes in the Mn-Mn distances in the  $S_3$  state compared to the  $S_1$  and the  $S_2$  states. The two 2.7 Å Mn-Mn distances that characterize the di- $\mu$ -oxo centers in the  $S_1$  and  $S_2$  states are lengthened to ~2.8 and 3.0 Å in the  $S_3$  state, respectively (Fig. 5). The 3.3 Å Mn-Mn and Mn-Ca distances also increase by 0.04-0.2 Å. These changes in Mn-Mn distances are interpreted as consequences of the onset of substrate water oxidation in the  $S_3$  state. Mn-centered oxidation is evident during the  $S_0$  to  $S_1$  and  $S_1$  to  $S_2$  transitions. During the  $S_2$  to  $S_3$  transition, we propose that the changes in Mn-Mn distances are the result of ligand or water oxidation, leading to the formation of an oxyl radical intermediate formed at a bridging or terminal position. We propose that substrate/water oxidation occurring at this transition provides the trigger for the formation of the O-O bond; the critical step in the water oxidation reaction. The reaction of the oxyl radical with  $\text{OH}^-$ ,  $\text{H}_2\text{O}$ , or an oxo group during the subsequent S-state conversion is proposed to lead to the formation of the O-O bond (Liang et al. 2000).

b) S<sub>0</sub> State Although many topological structures have been proposed, on the basis of EXAFS (Dau et al. 2001; Penner-Hahn 1998; Robblee et al. 2001) and EPR (Hasegawa et al. 1999a; Peloquin & Britt 2001) studies for the Mn cluster in PS II, the issue of whether there are two, or three di- $\mu$ -oxo bridged moieties in the Mn cluster has been an open question. However, a determination of the exact number of such interactions is important for narrowing the number of options from the many that have been proposed. We have obtained data from the S<sub>0</sub> state that shows that there may be three di- $\mu$ -oxo bridged motifs that are characterized by the  $\sim 2.7$  Å Mn-Mn distance.

The S<sub>0</sub> samples were generated by three flash illumination followed by incubation with FCCP (Messinger et al. 1997a). The multiline EPR spectra of the S<sub>0</sub> and S<sub>2</sub> states were used to characterize all the samples.

Figure 6 shows the Fourier transforms of the  $k^3$ -weighted EXAFS data from the S<sub>0</sub> and S<sub>1</sub> states. The Fourier peak II was isolated and fit to the EXAFS equation with one or two distances for the S<sub>0</sub> and S<sub>1</sub> states. Since there has to be an integral number of Mn-Mn interactions, the fits were constrained to a ratio of 1:1 and 2:1 for the two Mn-Mn interactions. The fits for the S<sub>1</sub> state showed that the two distances are separated by significantly less than what can be resolved using the data set at present. However, the situation is very different for the fits for the S<sub>0</sub> state. The quality of fit parameter,  $\epsilon^2$  value, improved considerably on going from a one-shell to a two-shell fit, especially, for the case of the 2:1 ratio of Mn-Mn interactions. The result of the fitting provides evidence that there may be three short ( $\sim 2.7$  to  $2.8$  Å) Mn-Mn interactions in the Mn cluster.

Figure 7 shows the best fits to two Mn-Mn distances for the S<sub>0</sub> state as a function of the number of each of the Mn-Mn interactions. The contour plot graphically illustrates that the

global minimum at two Mn-Mn interactions at 2.7 Å and one Mn-Mn interaction at 2.85 Å is very well defined, and hence gives us confidence in its reliability. We have re-examined the data from the  $S_0^*$ , the  $g=4.1 S_2$ , the  $\text{NH}_3^-$  or  $\text{F}^-$ -inhibited  $S_2$ , and the  $S_3$  states, where we have shown that there is distance heterogeneity in the Mn-Mn vectors. In each case, the ratio of the number of the two Mn-Mn interactions was  $\sim 2:1$ . This leads us to think that there may be three Mn interactions at  $\sim 2.7$  Å in the Mn cluster in its native state, one of which is perturbed in the  $S_0$ ,  $\text{F}^-$ - and  $\text{NH}_3^-$ -inhibited  $S_2$  and the  $g=4.1 S_2$  states, and all three are perturbed in the  $S_3$  state.

We showed earlier that there are several topological models that are compatible with our EXAFS data. These models have been described in detail (DeRose et al. 1994; Robblee et al. 2001). However, only one of those options (A) is widely used as a working model, although options E, F and G (the nomenclature is from DeRose *et al.* 1994) have been shown to be preferred on the basis of EPR and ENDOR simulations [Peloquin, 2001 #2099; (Hasegawa et al. 1999b). On the basis of the present results from the  $S_0$  state, we think it is important to consider the options that include three di- $\mu$ -oxo bridged moieties in the Mn cluster. Figure 8 shows several such options, G, I, and J, that were proposed earlier, and two newer options, L and M, among several others that can be conceived. Options J and I are less likely structures because J lacks a Mn-Mn interaction at 3.3 Å, and I has two such interactions. The EXAFS data from an inorganic compound (Auger et al. 1990) with the motif in option J is very different from that obtained from a PS II sample (unpublished data). Options G, L and M all have three di- $\mu$ -oxo bridges and one mono- $\mu$ -oxo bridge; that is three 2.7 Å and one 3.3 Å Mn-Mn interaction and are qualitatively in agreement with the asymmetry seen in the electron density of the Mn cluster (Zouni et al. 2001). Options L and M are also similar to the structure proposed on the basis of density functional theory calculations (Siegbahn 2000). We are in the process of testing the

compatibility of various Mn models that satisfy the criteria set by the distance data and orientation of the various Mn-Mn/Ca vectors obtained from EXAFS with the electron density data obtained from X-ray crystallography.

### **3. The Role of the Ca Cofactor**

Calcium and chloride are necessary cofactors for the proper function of the OEC of PS II (Debus 1992). While the Mn complex has been extensively studied by X-ray absorption techniques (XANES and EXAFS) comparatively less is known about the Ca<sup>2+</sup> cofactor. The fewer number of studies on the Ca<sup>2+</sup> cofactor have sometimes relied on substituting the native cofactor with strontium or other metals (Boussac & Rutherford 1988; Ghanotakis et al. 1985) and have stirred some debate about the structure of the binding site. Past efforts using Mn EXAFS on Sr-substituted PS II are suggestive of a close link between the Mn cluster and Sr, within 3.5 Å (Latimer et al. 1995). The most recent published study using Sr EXAFS on similar samples confirms this finding of a 3.5 Å distance between Mn and Sr (Cinco et al. 1998). This finding was based on a second Fourier peak in the Sr EXAFS from functional samples, but is absent from inactive, hydroxylamine-treated PS II. This Fourier peak II was found to fit best to two Mn at 3.5 Å rather than lighter atoms (carbon). Nevertheless, other experiments have given contrary results (RiggsGelasco et al. 1996).

a) Oriented Sr EXAFS We have extended the technique by using polarized Sr EXAFS on layered Sr-substituted samples, to provide important angle information. Polarized EXAFS involves collecting spectra for different incident angles ( $\theta$ ) between the membrane normal of the layered sample and the X-ray electric field vector. Dichroism in the EXAFS can occur, depending on how the particular absorber-backscatterer (A-B) vector is aligned with the electric

field. Through analysis of the dichroism, we extract the average orientation ( $\phi$ ) of this A–B vector relative to the membrane normal, and the average number of scatterers per absorbing atom ( $N_{iso}$ ). Constraints on the structural model are then imposed by these parameters.

Sr-substituted PS II samples were made by a process of Ca depletion, Sr<sup>2+</sup> reactivation, and Chelex treatment to remove excess Sr (Cinco et al. 1998). Oriented samples were made by layering onto flat Mylar films (Dau et al. 1995; Mukerji et al. 1994). The Fourier transforms from the polarized Sr EXAFS showed extreme dichroism in Fourier peak II (Figure 9). Nonlinear least-squares regression analysis produced the solid curve shown in Figure 9 as the best fit of the 15 data points (angles from six separate samples) and the result translates to 1–2 Sr–Mn vectors with an average angle of  $23 \pm 4^\circ$  from the membrane normal.

The orientation data from the Sr EXAFS experiments can be combined with the dichroism data from Mn EXAFS data to calculate the orientation of the 3.3 Å Mn–Mn vector. Fourier Peak III in the Mn EXAFS which contains Mn–Mn (3.3 Å) and Mn–Ca (3.4 Å) contributions, is dichroic, with an average angle of  $43 \pm 10^\circ$  with respect to the membrane normal (Mukerji et al. 1994). By including the Mn–Ca vector at  $23^\circ$ , we calculate an angle of  $\sim 62^\circ$  for the 3.3 Å Mn–Mn vector. Previous polarized Mn EXAFS experiments on PS II have shown angles of  $55^\circ$  and  $67^\circ$  for the 2.7 Å Mn–Mn vectors (Dau et al. 1995; Mukerji et al. 1994). All Mn–Mn vectors then lie at roughly the same angle ( $\sim 60^\circ$ ) with respect to the membrane plane, but are not restricted to being collinear, because the PS II membranes are ordered in one dimension only.

b) Ca EXAFS In a complementary and definitive experiment, we have used Ca K-edge EXAFS studies to probe the binding site of the native cofactor for any nearby Mn, within  $\sim 4$  Å. This is analogous to the Sr EXAFS studies already published (Cinco et al. 1998), but it focuses

on the native cofactor and avoids the treatments involving Ca depletion and Sr substitution. This technique is a more sensitive and direct probe of the Ca binding site in PS II than Sr EXAFS.

Chelex-treated PS II (2Ca/4Mn/PS II) was accumulated on the flat surface of solid Plexiglas (Ca-free) while not affecting oxygen-evolving activity or the S<sub>2</sub> EPR multiline signal. To prepare the parallel control (inactive) sample, 40  $\mu$ L of hydroxylamine (NH<sub>2</sub>OH, 100 mM) was added to the surface of the layered PS II, was allowed to soak and dry without loss of PS II material.

The Fourier transform of the Ca EXAFS is presented in Figure 10. The spectra are remarkably similar to the Fourier transforms of the earlier Sr EXAFS study with Sr substituted for Ca. The first (largest) Fourier peak corresponds to the coordinating oxygen atoms closest to calcium. In contrast to the control (NH<sub>2</sub>OH-treated) sample, the Chelex-treated PS II shows a second Fourier peak. When this peak II is isolated and simulated with possible scattering atoms, it corresponded best to Mn at 3.4 Å, rather than to light atom (C, O or Cl) neighbors. These results were consistent with the earlier Sr EXAFS studies (Cinco et al. 1998).

The Ca EXAFS protocol has directly probed the Ca cofactor in its native state, avoiding potentially disruptive treatments such as the low pH exposure involved in metal substitution. This result provides proof of the Mn/Ca heteronuclear cluster as the catalytic site of oxygen evolution in PS II.

Because significant angle information about Mn–Mn and Mn–Ca vectors is now available as described above, other topological models previously discussed (DeRose et al. 1994) can be refined to include the presence of Ca and account for the dichroism data. Three possible refined models are presented in Figure 11. Recent EXAFS data from the S<sub>0</sub> state has indicated

that there maybe three di- $\mu$ -oxo bridged Mn-Mn units (see above). However, such motifs were not included in Figure 11, but considering them is a subject for future work.

#### 4. Mechanism of Water Oxidation

The many mechanisms of water oxidation that have been proposed can be broadly divided into four groups, in which the kind of oxygen atom (Mn-terminal oxygen ligand, Mn-bridging oxygen ligand or exogenous oxygen from water or hydroxide) involved in O-O bond formation is different. In these four groups the oxygen atoms can be derived from two terminal oxygens bound to Mn atoms (Haumann & Junge 1999; Hillier & Wydrzynski 2001; Hoganson & Babcock 2000; Limburg et al. 2000; Messinger et al. 2001; Renger 2001; Robblee et al. 2001), one bridging or terminal oxygen and one exogenous oxygen (Dau et al. 2001; Hillier & Wydrzynski 2001; Kuzek & Pace 2001; Messinger et al. 2001; Robblee et al. 2001; Siegbahn & Crabtree 1999; Vrettos et al. 2001), one terminal and one bridging oxygen (Limburg et al. 2000; Nugent et al. 2001), or two bridging oxygens (Brudvig & Crabtree 1986; Christou & Vincent 1987; Messinger et al. 2001; Robblee et al. 2001). The results that have a direct bearing on this subject are the H<sub>2</sub><sup>16</sup>O/H<sub>2</sub><sup>18</sup>O exchange studies by Messinger and coworkers that support the presence of two non-equivalent exchangeable sites in the S<sub>3</sub> state (Hillier et al. 1998; Messinger et al. 1995).

We have proposed earlier a model for the S<sub>3</sub> state in which an oxyl radical is generated on one of the  $\mu$ -oxo-bridges, which results in an increase of one of the Mn–Mn distances to  $\sim 2.95$  Å (Figure 5) (Liang et al. 2000; Yachandra et al. 1996). Consistent with our XANES results is the implication in this structure that the oxidative equivalent is not stored on the Mn atoms *per se* during the S<sub>2</sub>  $\rightarrow$  S<sub>3</sub> transition but is delocalized with significant charge and spin density on the

bridging oxo ligand. These data are reinforced by our Mn K $\beta$  emission studies of the various S-states (Messinger et al. 2001). Second, the Mn–Mn distance in both of the di- $\mu$ -oxo-bridged units increases from 2.72 Å to 2.82 Å and 2.95 Å, upon the formation of the S<sub>3</sub> state. These changes imply a significant structural change in the Mn cluster as it proceeds to the S<sub>3</sub> state. It is difficult to rationalize such changes in Mn–Mn distance as arising purely from Mn oxidation, or involving Mn terminal ligands. It is difficult to understand how changes in terminal ligation can generate such a profound change on the Mn–Mn distances in the S<sub>3</sub> state. Replacement of terminal ligands in di- $\mu$ -oxo bridged model compounds has a minimal effect on the Mn–Mn distance of 2.7 Å that is characteristic of such di- $\mu$ -oxo bridged Mn compounds (Pecoraro 1992; Wieghardt 1989). However, it is easier to rationalize increases in Mn–Mn distances as being due to changes in the bridging structure. We propose that substrate/water oxidation chemistry is occurring at the S<sub>2</sub> to S<sub>3</sub> transition as evidenced by the lack of Mn oxidation, and leading to the significant structural changes as seen in the increase in the Mn–Mn distances. Our proposed mechanism, where the O–O bond is formed between one critical bridging oxygen and another oxygen atom derived from a bridging/terminal oxygen ligand or an exogenous oxygen also avoids the formation of the O–O bond until the most oxidized state (S<sub>4</sub>) is reached. This precludes the formation and release of peroxide or other oxidation products of water in the earlier S-states, thus preventing the system from ‘short circuiting’ and avoiding the risk of damaging the polypeptides of PS II.

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## Figure Captions

**Figure 1.** The  $S_2$  multiline EPR signal amplitudes are shown as a function of flash number (solid line). The best fit to the  $S_2$  multiline oscillation pattern is shown as a dashed line. The S-state compositions of the samples are derived from the fit. Individual S-State spectra can be uniquely determined from these data, if the oscillation pattern is as pronounced as shown in the Figure. It is important that a deep oscillation pattern is obtained, because all subsequent steps that are used in the deconvolution to generate pure S-state spectra are critically dependent on this data.

**Figure 2.** A): Second derivatives of the normalized pure S-state Mn K-edge spectra of the Mn-cluster of PS II. For clarity a vertical dashed line has been drawn at the inflection-point energy (IPE) of the  $S_1$ -state. The change in IPE ( $\Delta$  in eV) for each S-state advance is given at the right. For reference, the K-edge spectrum of the  $S_1$ -state is plotted as well (dashed line). B): The  $S_1$ - $S_0$ ,  $S_2$ - $S_1$ ,  $S_3$ - $S_2$ , and  $S_0$ - $S_3$  difference high resolution Mn X-ray  $K\beta$ -emission spectra. The derivative shapes indicate shifts in the energy of the spectra between the  $S_0$  and  $S_1$ , and  $S_1$  and  $S_2$  spectra, indicating oxidations of Mn. The  $S_3$ - $S_2$  difference spectra demonstrates the similarity of the spectra, suggesting that this advance does not involve a Mn centered oxidation.

**Figure 3.** A) Inflection-point energy (IPE, in eV) of the Mn K-edge of PS II membranes as a function of the number of applied flashes. B) The first moments of the Mn  $K\beta$  fluorescence spectra ( $\langle E \rangle$ ) as a function of flash number. The Mn K-edge shifts to higher energy as Mn is oxidized. The first moment of the Mn  $K\beta$  fluorescence shifts to lower energy as Mn is oxidized.

The complementary and mutually reinforcing results provide a strong case for Mn oxidation during the  $S_0$  to  $S_1$  and  $S_1$  and  $S_2$  advances but no Mn-centered oxidation at the  $S_2$  to  $S_3$  advance.

**Figure 4.** The electronic energy diagram of the transitions that occur in Mn K-edge, L-edge and RIXS spectroscopy. L-edge like information is extracted from RIXS by subtracting the emission energy,  $f$ , from the excitation energy,  $\nu$  (energy transfer,  $(\nu-f)$  vs the incident energy,  $\nu$ ). However, more final states are present in the RIXS spectra than in the L-edge spectra. Therefore, RIXS gives a wealth of information about the Mn oxidation state, ligand environment and local symmetry.

**Figure 5.** A): Fourier-filtered  $k$ -space EXAFS data from  $S_2$  (dashed line) and  $S_3$  (solid line) state of PS II samples. The differences in phase, frequency and amplitudes between the raw  $S_2$  and  $S_3$  state EXAFS spectra are very obvious in these spectra. B): Fourier transform power spectra of  $S_2$  (dashed line) and  $S_3$  (solid line) states of PS II. The major Fourier peaks are labeled I, II and III. The spectra are clearly different between the  $S_2$  (dashed line) and  $S_3$  states (solid line). There is a reduction in amplitude in all three peaks in the  $S_3$  state compared to the  $S_2$  state. More importantly, peaks II and III are at a greater apparent distance  $R^1$  for the  $S_3$  state compared to the  $S_2$  state as shown.

**Figure 6.** Fourier Transforms of the  $S_0$  (dashed line) and  $S_1$  state (solid line) of the  $k^3$ -weighted Mn-EXAFS spectra (3.5 to 11.5  $\text{\AA}^{-1}$ ). The main Fourier peaks are labeled I, II and III. There are clear differences in amplitude and position of these peaks between the  $S_0$  and  $S_1$  states. The peak labeled II is from Mn-Mn interactions at  $\sim 2.7 \text{\AA}$ .

**Figure 7.** The contour plot shows the fit-quality-parameter  $\epsilon^2$  for the  $S_0$  state plotted versus the number of the two Mn-Mn interactions. The dashed lines shows the distinct minimum at one Mn-Mn vector at 2.85 Å and two Mn-Mn vectors at 2.7 Å.

**Figure 8.** Structural models for the Mn cluster with three di- $\mu$ -oxo bridged Mn moieties. Options G, I, and J are from DeRose et al., 1994, and L and M are two other viable options. One or two mono- $\mu$ -oxo bridged motifs are also present in Options G, I, L, and M. No such motif exists in option J.

**Figure 9.** Left panel. Fourier transforms of  $k^3$ -weighted Sr-EXAFS from oriented Sr-PS II samples at two angles ( $\theta$ ). The dichroism is most readily apparent in Fourier peak II ( $R' = 3.0$  Å) which is assigned to backscattering from Mn. Right Panel. The polar plot shows  $N_{\text{exp}}$  (filled circles) plotted vs the angle of detection,  $\theta$ . The solid line is the best fit from which we obtain  $N_{\text{iso}}$  (number of Mn neighbors to Sr) and  $\phi$ , the angle the Mn-Sr vector makes with the membrane normal.

**Figure 10.** Fourier Transform of Ca EXAFS from Chelex-treated, layered samples with 2 Ca/PS II ( $k^3$ -weighted,  $k = 2.5$ – $10.5$  Å<sup>-1</sup>). The FTs show the presence of a second Fourier peak in the 2Ca/PS II sample that fits to Ca-Mn that is absent in the control sample.

**Figure 11.** Refined models for the active site of the OEC in Photosystem II. These models combine the finding from oriented Sr-substituted PS II samples with previous results from Mn

EXAFS on oriented PS II samples. These are derived from core structures that have been described in earlier studies (Fig. 11A-B from option A and Fig. 11C from option F in DeRose et al., 1994).

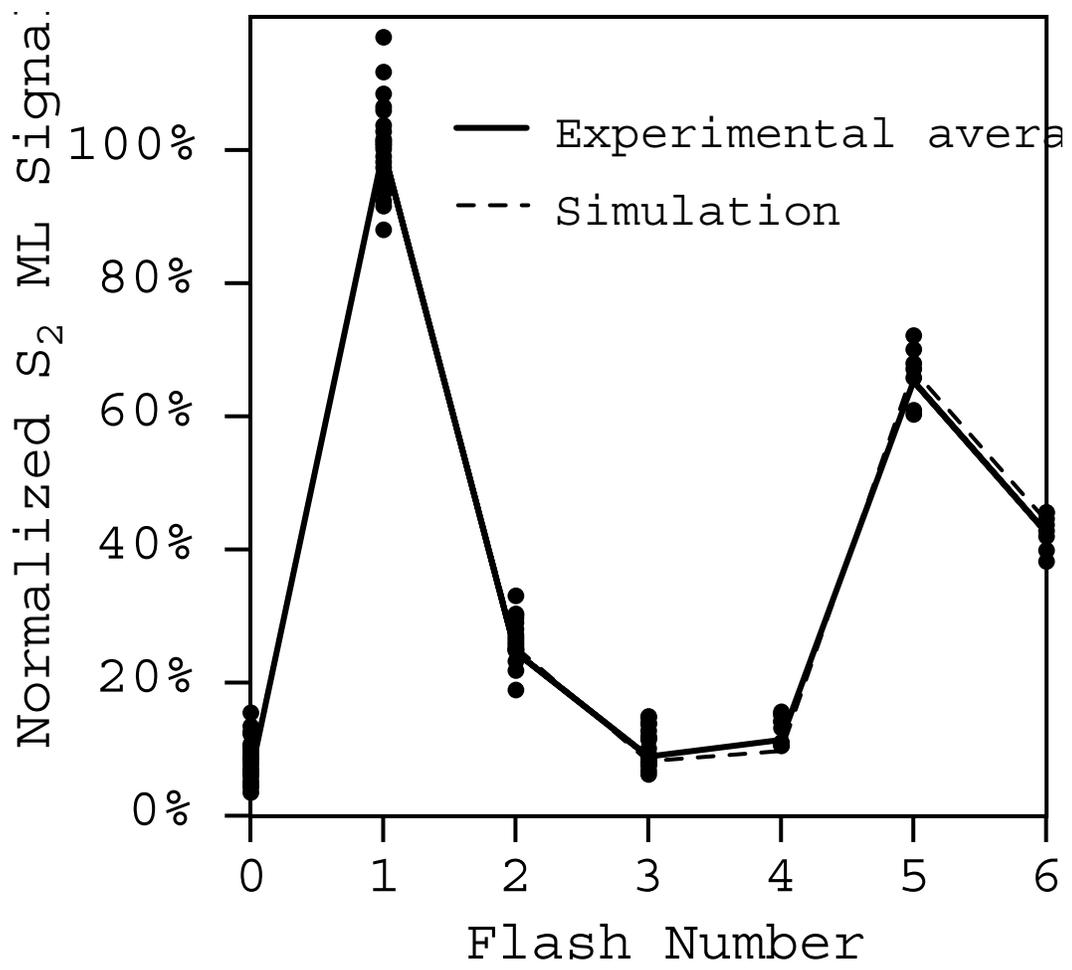


Figure 1

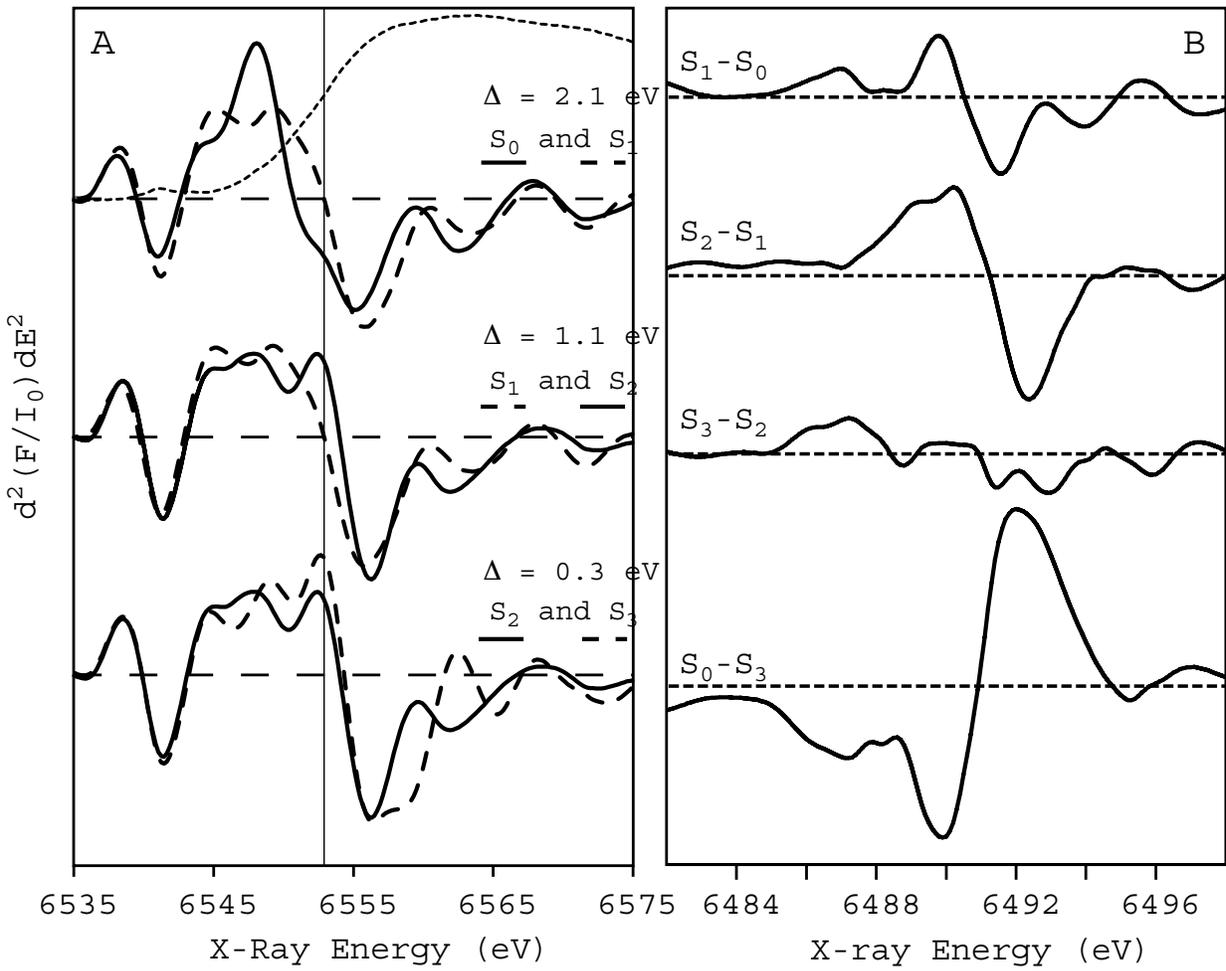


Figure 2

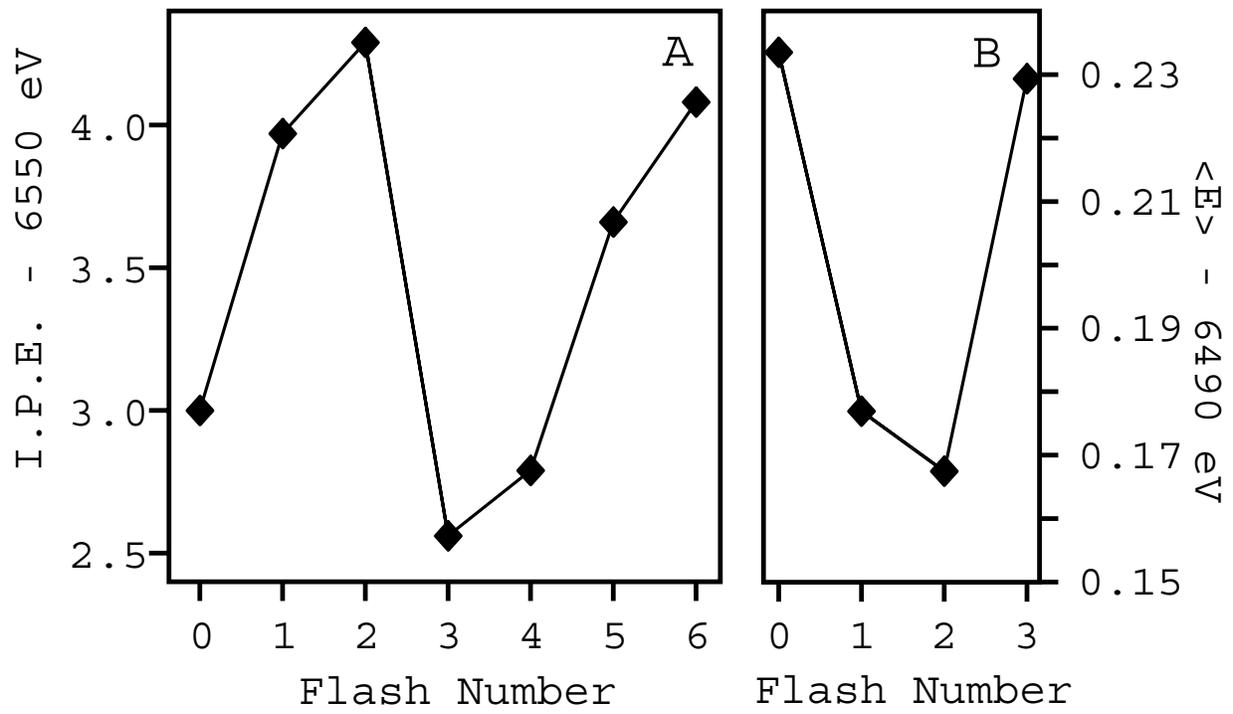


Figure 3

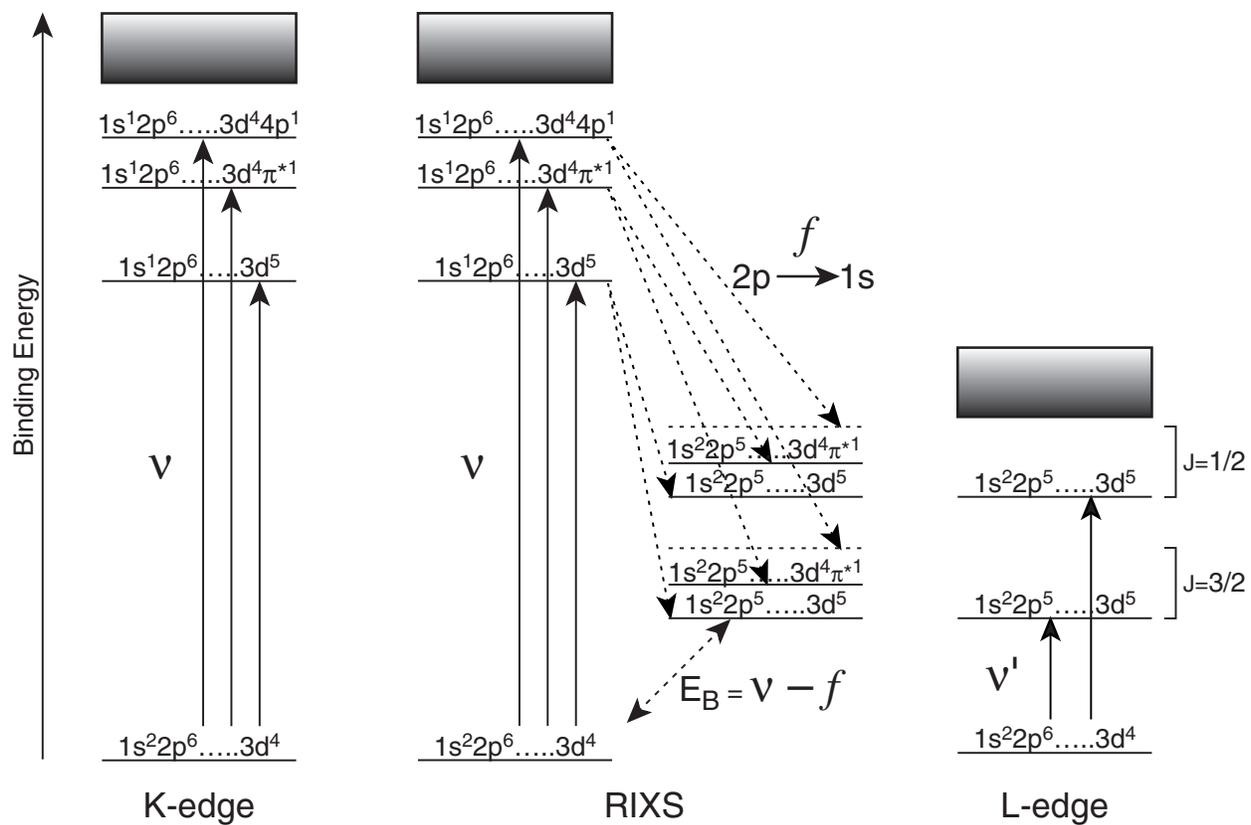


Figure 4

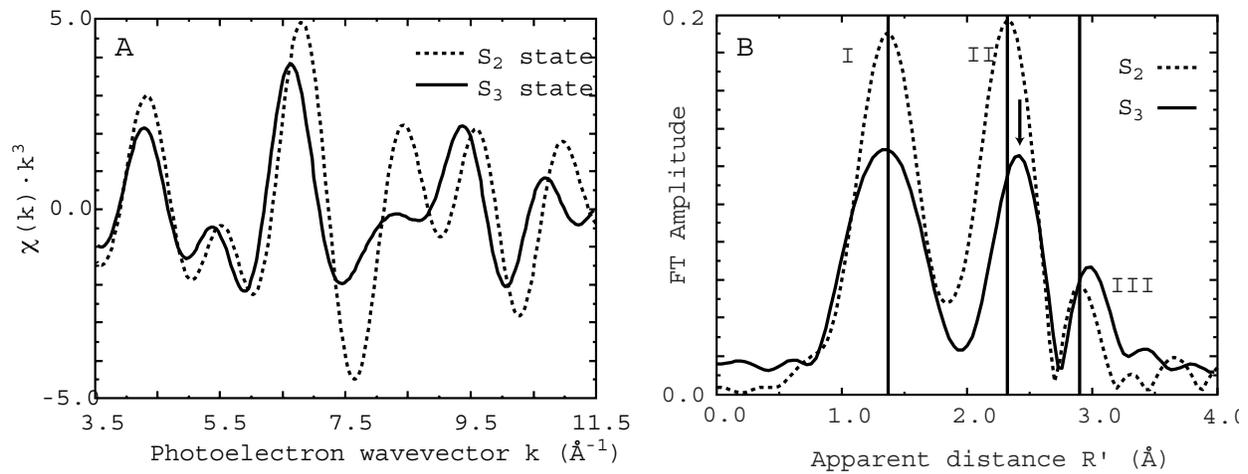


Figure 5

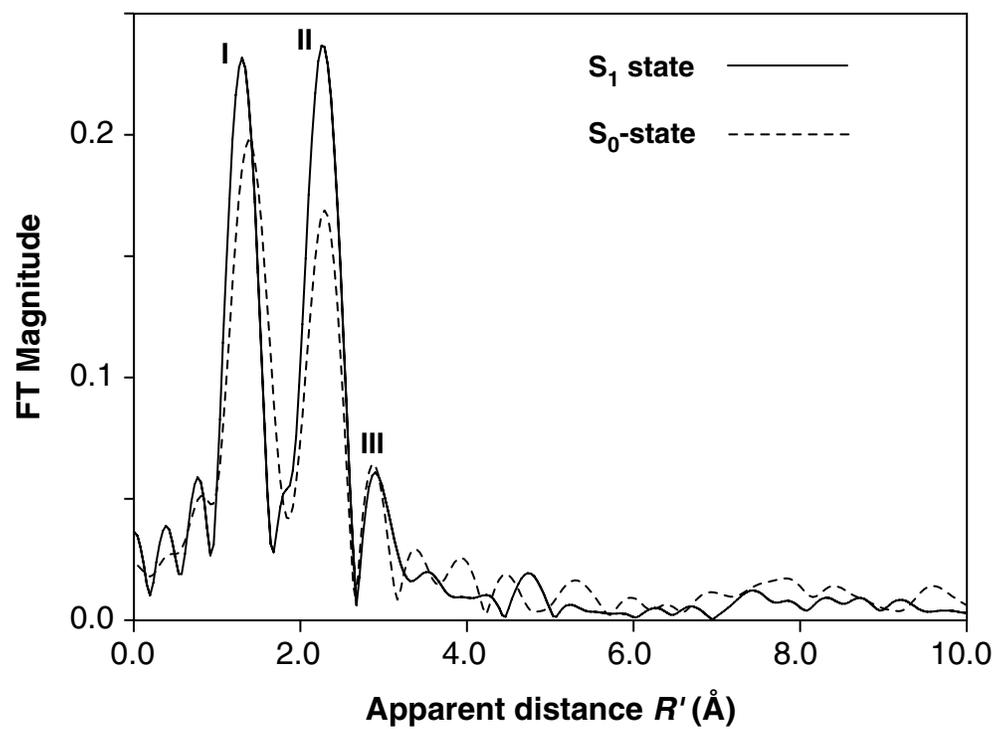


Figure 6

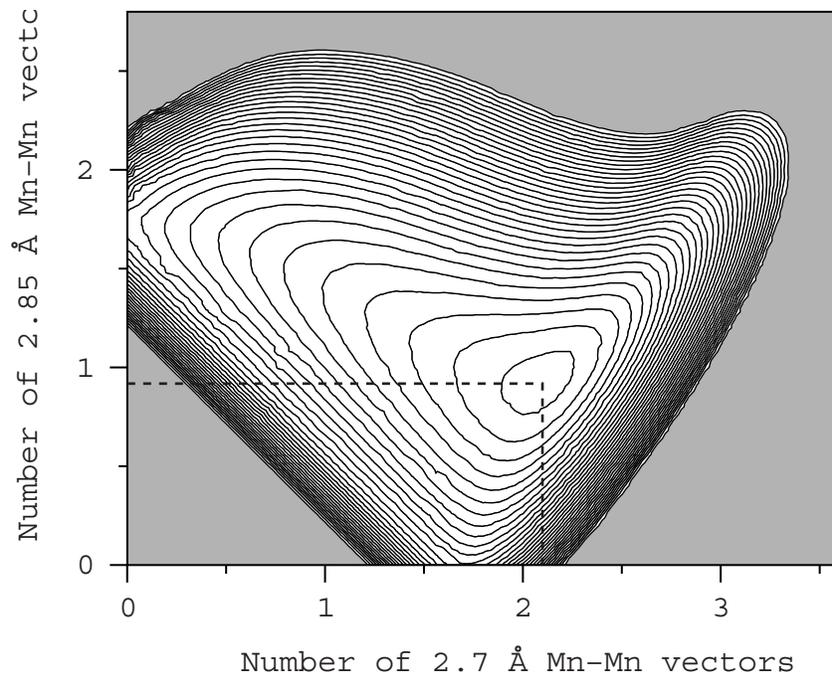
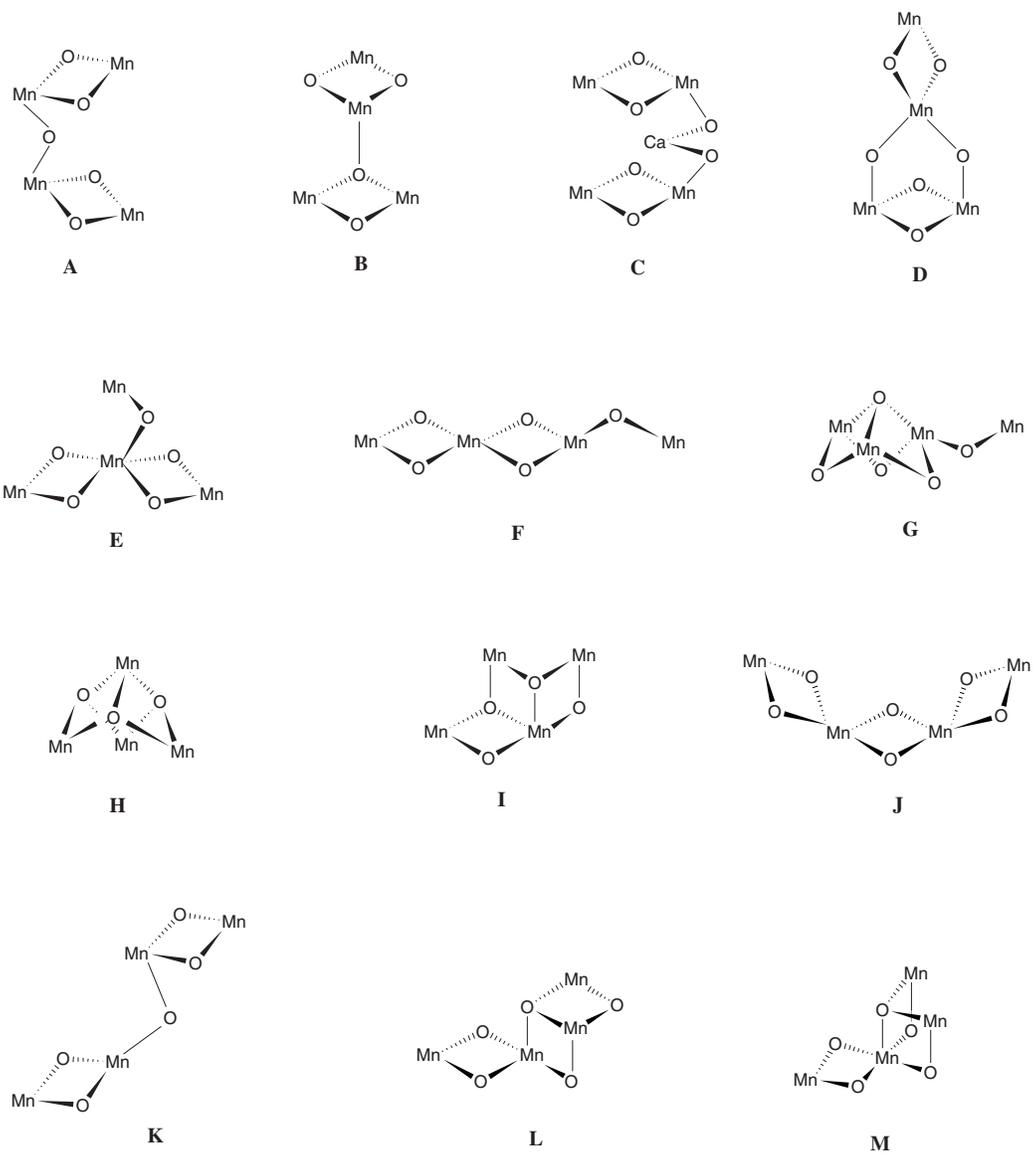


Figure 7



**Figure 8**

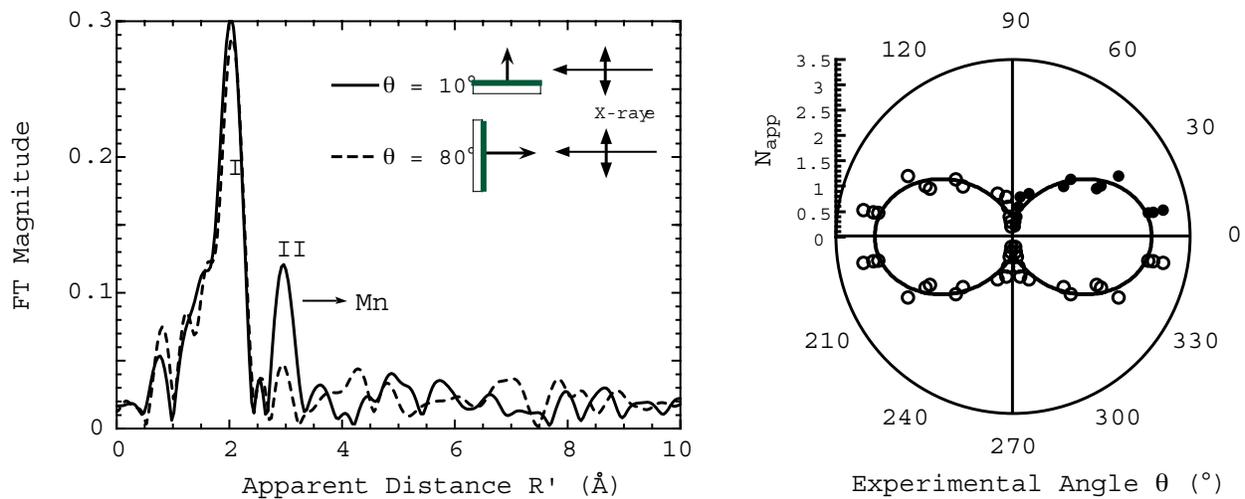
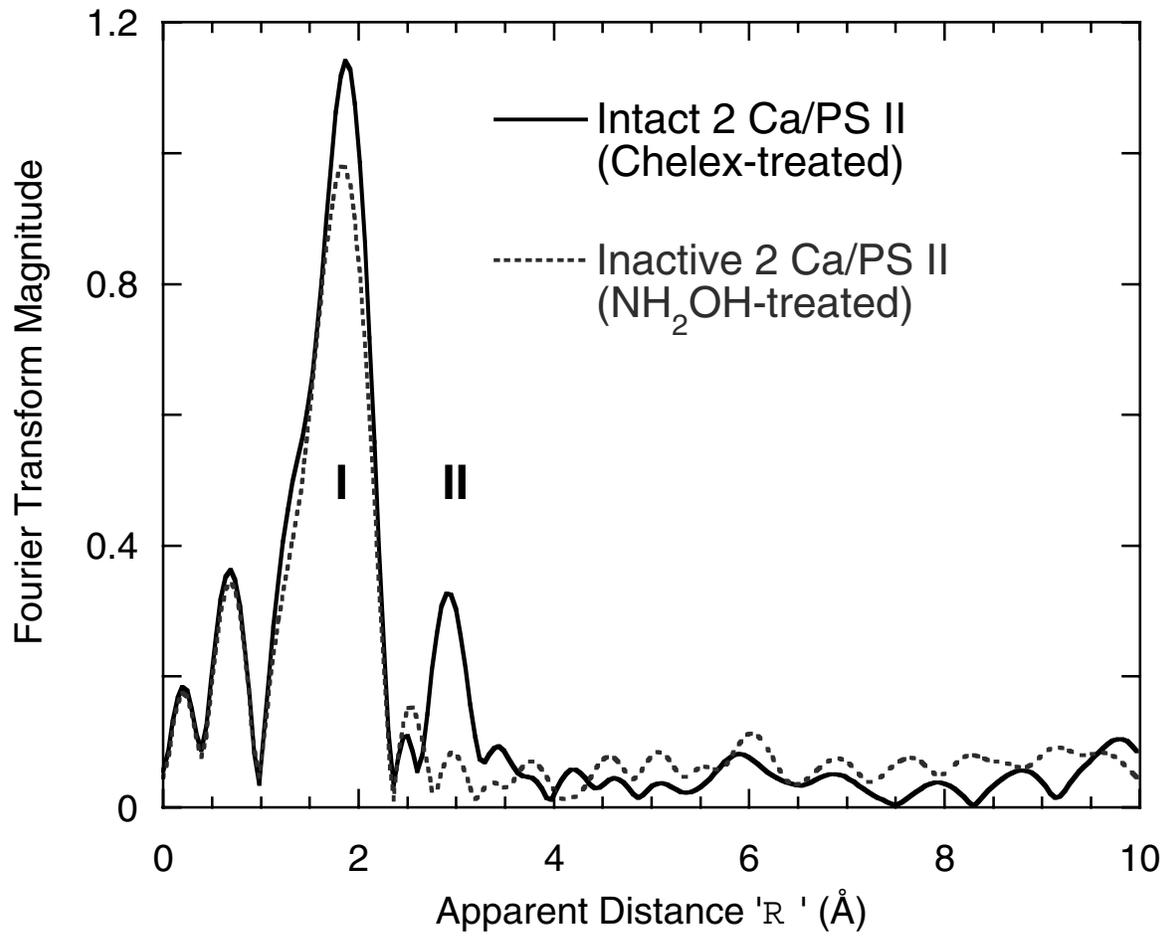
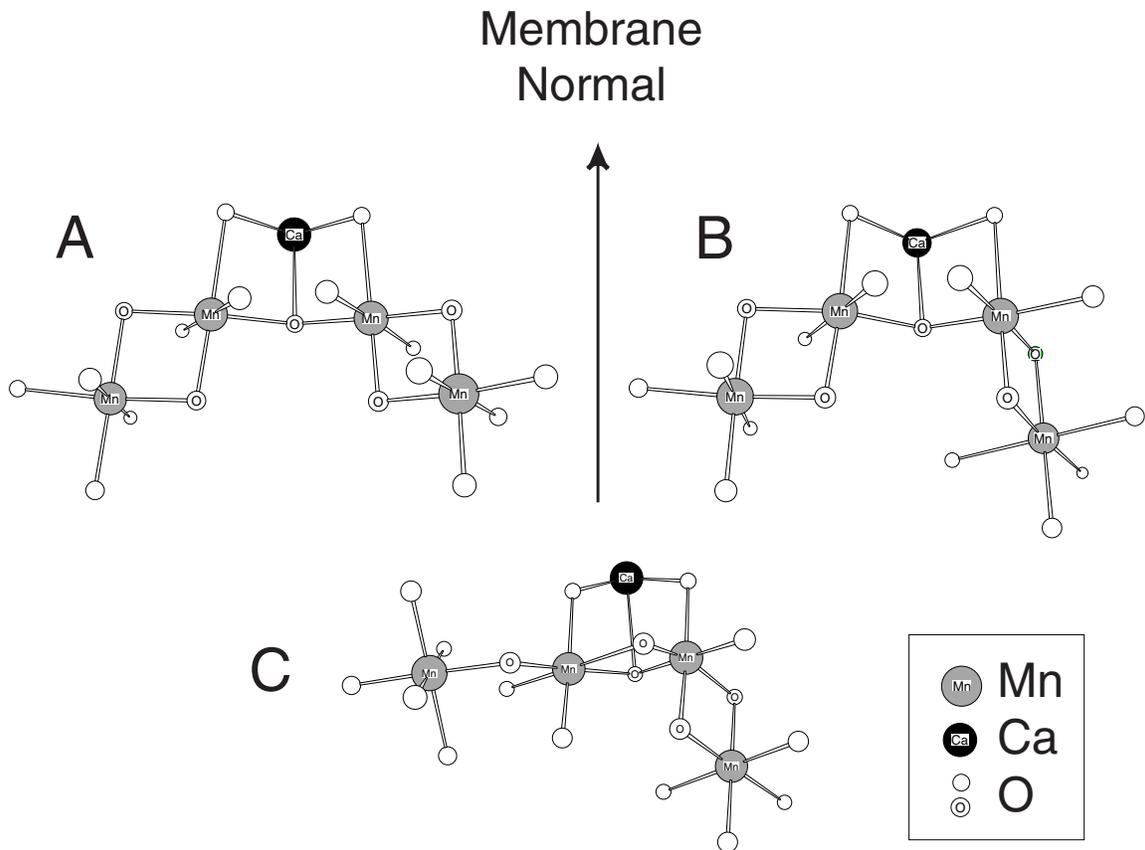


Figure 9



**Figure 10**



**Figure 11**