## **Lawrence Berkeley National Laboratory**

#### **Recent Work**

#### **Title**

HIGH FIELD MAGNET DEVELOPMENT ANALYSIS ""MECHANICS OF THIN RINGS WITH DIPOLE-LIKE LOADS

#### **Permalink**

https://escholarship.org/uc/item/5bx5n8dj

#### **Author**

Meuser, Robert B.

#### **Publication Date**

1979-09-01

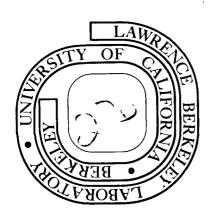
RECEIVED
AWRENCE
BERKELSY LABORATORY

UUI 3 1 1979

LIBRARY AND DOCUMENTS SECTION

# For Reference

Not to be taken from this room



#### **DISCLAIMER**

This document was prepared as an account of work sponsored by the United States Government. While this document is believed to contain correct information, neither the United States Government nor any agency thereof, nor the Regents of the University of California, nor any of their employees, makes any warranty, express or implied, or assumes any legal responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by its trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof, or the Regents of the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof or the Regents of the University of California.

LAWRENCE BERKELEY LABORATORY - UNIVERSITY OF CALIFORNIA

ENGINEERING NOTE

MD IIII M5399

1 of G

PROGRAM - PROJECT - JOB

HIGH-FIELD MANNET DEVELOPMENT

HIGH-FIELD MAGNET DEVELOPMENT ANALYSIS

TITLE

### MECHANICS OF THIN RINGS WITH DIPOLE-LIKE LOADS

This is a summary (believe it or not); no derivations are incl. Usual thin ring/shell assumptions apply: deflections due to extension of middle surface and shear are ignored.

We consider the following distributed unit loads on a circular thin ring.

$$Pr = Cr \cos^2\theta$$
 for unite,  
 $Pt = Ct \sin\theta \cos\theta$  see page 6  
 $Pm = C_m a \sin\theta \cos\theta$ 

These are the sorts of loads that can be produced by a circular-cross-section bending magnet;

the quantitative correspondence will be covered later.

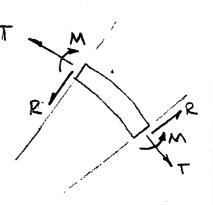
These loads result in the following internal forces in the ring:

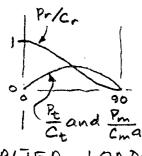
$$T = \frac{1}{3} \left[ C_r(1 + \sin^2 \theta) + C_t(1 - 2\sin^2 \theta) \right]$$

$$R = \frac{1}{3} \left( 2 C_r - C_t \right) \sin \theta \cos \theta$$

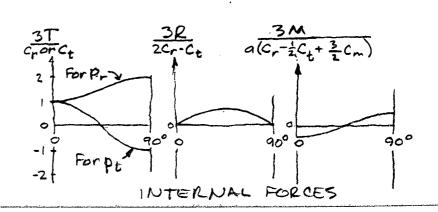
$$M = \frac{1}{3} \alpha \left( C_r - \frac{1}{2} C_t + \frac{3}{2} C_m \right) \left( \sin^2 \theta - \frac{1}{2} \right)$$











LAWRENCE BERKELEY LABORATORY - UNIVERSITY OF CALIFORNIA		CODE	SERIAL	PAGE
ENGINEERING NOTE		WDIII	M5399	2 OF 6
AUTHOR	DEPARTMENT	LOCATION	DATE	
R, MEUSER	MECH	BERLK	SEPT	7 1979

The maximum and minimum values are as follows:

$$T_{\theta=0} = \frac{1}{3} C_r$$
 for pronly

 $T_{\theta=0} = \frac{2}{3} C_r$  for pronly

 $T_{\theta=0} = \pm \frac{1}{3} C_t$  for pronly

 $R_{\theta=45} = \frac{1}{6} (2C_r - C_t)$ 
 $M_{\theta=0} = \mp \frac{1}{6} (C_r - \frac{1}{2} C_t + \frac{3}{2} C_m) a$ 

The body forces on an element of a thin winding having a cosme-0 distribution of currents

are  $dF_r = 2m \frac{Be}{Ba} \cos^2\theta d\theta$ 

$$dF_{t} = -2m\left(1 + \frac{B_{B}}{B_{A}}\right)\cos\theta\sin\theta d\theta$$

where 
$$m = \frac{B_a^2 \hat{a}}{4\pi}$$
 (cgs emunits)
$$= \frac{B_a^2 \hat{a}}{u_0}$$
 (S.I. mks units)

B<sub>A</sub> = contral field contributed by the winding.
B<sub>B</sub> = " " " iron yoke.

à = coil mean radius

(For a thick winding, see MS256, MEUSER, NOV 14, 1978

Ba and BB are expressed in terms of the current density and geometry as follows

where 1 = 10 coso, to = peak current/circumference

The state of the s		CODE	SERIAL	PAGE
ENGINEERING NOTE		WOIII	M5399	3 OF 6
AUTHOR	DEPARTMENT	LOCATION	DATE	
R.MEUSER	MECH	BEIRK	SEP 1	7 1979

The resulting loads on the surrounding rings are as follows:

For a coil that sticks to the rings, with no force transmitted tangentially within the coil, the body forces are transmitted directly to the rings:

$$P_{t} = \frac{dF_{t}}{d\theta} = 2m \frac{B_{B}}{B_{A}} \cos^{2}\theta$$

$$P_{t} = \frac{dF_{t}}{d\theta} = -2m \left(1 + \frac{B_{B}}{B_{A}}\right) \cos\theta \sin\theta$$

$$P_{m} = -\frac{dF_{t}}{d\theta} (a-\hat{a}) = 2m \left(1 + \frac{B_{B}}{B_{A}}\right) (a-\hat{a}) \cos\theta \sin\theta$$

$$C_r = 2m \frac{B_B}{B_A}$$

$$C_t = -2m \left(1 + \frac{B_B}{B_A}\right)$$

$$C_m = 2m \left(1 + \frac{B_B}{B_A}\right) \frac{a - \hat{a}}{a}$$

(For applied bending moments, pm, the philosophy is that the ring has a small but finite thickness, so shear stresses applied at the inside surface result in applied moments:  $P_m = T(a-\hat{a})$ .)

LAWRENCE BERKELEY LABORATORY - UNIVERSITY OF CALIFORNIA		CODE	SERIAL	PAGE
ENGINEERING NOTE		MDIII	M5399	4 OF 6
AUTHOR	DEPARTMENT	LOCATION	DATE	
R.MEUSER	MECH	BERK	SEP 17	1979

For a "slippery" coil, no shear is transmitted to the ring, but tangential forces are transmitted from element to element within the coil. The resulting radial loads are

$$P_r = m(1 + 3\frac{B_B}{B_A}) \cos^2\theta$$
or
 $C_r = m(1 + 3\frac{B_B}{B_A})$ ,  $C_t = 0$ ,  $C_m = 0$ 

The internal ring forces expressed in terms of the aporture field components for a thin cosine-0 winding are:

For a sticky coil;
$$T_{\theta=0}=-\frac{2}{3}m$$

$$T_{\theta=90} = \frac{2}{3}m\left(1+3\frac{B_B}{B_A}\right)$$

for a slippery coil:

LAWRENCE BERKELEY LABORATORY - UNIVERSITY OF CALIFORNIA

ENGINEERING NOTE

MD IIII M5399 5 of 6

AUTHOR

R. MEUSER MECH BEYLK SEP 17 19 79

The local deflections are:

$$\delta_{x} = \frac{1}{3} \frac{M_{0}a^{2}}{EI} \cos^{3}\theta$$

$$\delta_{y} = -\frac{1}{3} \frac{M_{0}a^{2}}{EI} \sin^{3}\theta$$

$$\delta_{y} = \frac{1}{3} \frac{M_{0}a^{2}}{EI} (\cos^{3}\theta + \sin^{3}\theta) = \frac{1}{3} \frac{M_{0}a^{2}}{EI} \cos^{3}2\theta$$

$$\delta_{t} = -\frac{1}{3} \frac{M_{0}a^{2}}{EI} \sin^{3}\theta \cos^{3}\theta = -\frac{1}{3} \frac{M_{0}a^{2}}{EI} \cdot \frac{1}{2} \sin^{3}2\theta$$

$$S_{m} = -\frac{M_{0}Q}{ET} \sin \theta \cos \theta = -\frac{M_{0}Q}{ET} \cdot \frac{1}{2} \sin 2\theta$$

where -Mo is the value of the bending moment at 0=0, or +Mo is its value at 0=900.

The maximum and minimum values are as follows

 $S_{m_{\Theta}=45^{\circ}} = -\frac{1}{2} \frac{M_{\circ} a}{ET}$ 

$$\delta_{\star \theta=0} = \delta_{\tau \theta=0} = -\delta_{\tau \theta=90} = \frac{1}{3} \frac{M_0 a^2}{EI}$$

$$\delta_{t \theta=45^{\circ}} = -\frac{1}{6} \frac{M_0 a^2}{EI}$$

LAWRENCE BERKELEY LABORATORY - UNIVERSITY OF CALIFORNIA		CODE	SERIAL	PAGE
ENGINEERING NOTE		MDIIII	M5399	6 OF 6
AUTHOR	DEPARTMENT	LOCATION	DATE	
R. MEUSER	MECH	BERK	SEP 17	1979

Momendature

Pr, Pt, Pm Applied loads, see sketch, p. 1

Cr, Ct, Cm See equations, top of p.1

a Radius of ring

à Coil aug radius

T, R, M Internal forces, see sketch, p. 1

dFr, dFt Element body forces, radial and tangential

Ba Field in aperture for no iron

BB Field in aperture from iron

b Iron inside radius

m see lower-middle, p. 2

8x, Sy, Sr, St, Sm See sketch, p. 5

E Elastic modulus

I Section moment of inertia

Units, dimensions: Any consistent units are OK except in is defined only in terms of SI and cgs-emu units

The loads, internal torces, and section modulus can be either for the whole ring or for unit axial length In English units the dimensions are as follows:

This report was done with support from the Department of Energy. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the Department of Energy.

TECHNICAL INFORMATION DEPARTMENT LAWRENCE BERKELEY LABORATORY UNIVERSITY OF CALIFORNIA BERKELEY, CALIFORNIA 94720