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Solid Electrolyte Bimodal Grain Structures for Improved Cycling Performance

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Abstract 20

The application of solid-state electrolytes in Li batteries has been hampered by the occurrence of Li-dendrite-caused short circuits. To avoid cell failure, the electrolytes can only be stressed with rather low current densities, severely restricting their performance. Since grain size and pore distributions significantly affect dendrite growth in ceramic electrolytes like $Li₇La₃Zr₂O₁₂$ and its variants, here we have proposed a "detour and buffer" strategy to bring the superiority of both coarse and fine grains into play. To validate the mechanism, a coarse/fine bimodal grain microstructure was obtained by seeding unpulverized large particles in the green body. The rearrangement of coarse grains and fine pores was fine-tuned through changing the ratio of pulverized and unpulverized powders. The optimized bimodal microstructure, obtained when the two powders are equally mixed, allowed, without extra interface decoration, cycling for over 2000 h as the current density was increased from 1.0 mA·cm⁻² gradually up to 2.0 mA·cm⁻². The "detour and buffer" effects was confirmed from postmortem analysis. The complex grain boundaries formed by fine grains discourage the direct infiltration of Li. Simultaneously, the coarse grains further increase the tortuosity of the Li path. Our study sheds light on the microstructure optimization for the polycrystalline solid-state electrolytes. 21 22 23 24 25 26 27 28 29 30 31 32 33 34 35 36 37

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Keywords: 39

LLZO solid-state electrolyte; bimodal microstructure; Li dendrite suppression; cycling 40

performance; detour and buffer effects 41

Introduction 42

The deployment of solid-state electrolytes (SSEs) in all-solid-state Li batteries presents a compelling opportunity due to their nonflammability and compatibility with Li metal, and thus can potentially meet increasing demands for higher energy density and better safety^{1, 2}. Among the various SSEs, $Li₇La₃Zr₂O₁₂$ (LLZO), which is doped with Al, Ga, Ta and so on to stabilize the cubic structure, exhibits high ionic conductivities $(10^{-4} \text{ to } 10^{-3} \text{ S} \cdot \text{cm}^{-1})$, a wide operating voltage window, and excellent elastic and shear moduli³. These make LLZO attractive for use in batteries, but a formidable obstacle still exists: Li dendrites (filaments) form easily and eventually cause short-circuit failure to the cells during electrochemical cycling when a certain critical current density (CCD) is exceeded. The CCD values of SSEs are on the order of hundreds μA at room temperature and much lower than those of liquid electrolytes⁴⁻ ⁶, leading all-solid-state Li batteries to operate at low rates. Therefore, circumventing Li dendrites at increased current densities is essential to obtain satisfactory performance on commercially applicable all-solid-state Li batteries. 43 44 45 46 47 48 49 50 51 52 53 54 55 56

The exact mechanism of dendrite growth in SSEs remains elusive^{7, 8}, prompting various strategies to enhance all-solid-state Li battery performance, such as improving Li wettability on the interface⁹⁻¹², boosting the Coble creep of $Li^{13, 14}$, and adjusting the ratio of ionic and electrical conductivity in SSEs^{15, 16}. Besides interface and transport modification, the intrinsic polycrystalline ceramic nature of LLZO SSEs suggests that microstructures such as grain size, grain boundaries, and pores are closely relevant to Li dendrite growth and $CCD¹⁷⁻²⁴$. Consequently, designing and optimizing these microstructures are crucial to improving the cycling properties of SSEs. 57 58 59 60 61 62 63 64

Recent studies have shown that fine-grained solid electrolytes, featuring higher grain boundary densities and isolated pores²⁵, effectively disperse currents and enhance interfacial contacts, allowing higher current densities^{26, 27}. However, such almost straight grain boundary microstructures offer insufficient barriers to Li dendrite growth^{28, 29} (Figure 1a), limiting the CCD of fine-grained SSEs to below 1 mA·cm^{-2 20,} ³⁰⁻³². Additionally, the processing window for SSEs with normally distributed fine grains is narrow^{33, 34}. The low energy barrier for sintering fine powders means that grains grow easily. Randomly distributed abnormal grain growth and loose grain boundaries are often observed due to the different mass diffusion rate between grains 65 66 67 68 69 70 71 72 73

(In general the larger the grains, the faster they grow). A highly densified structure with Ostwald ripening and grain coarsening can occur when the sintering time is prolonged. Even though coarse grains can provide higher ionic conductivity, the loss of grain boundary complexity, usually accompanied with large interconnected pores encourage Li-dendrite growth^{22, 35} (Figure 1b) Consequently, neither fine nor coarse grains solely can effectively promote the electrochemical cycling performance of LLZO SSEs, since both have pros and cons. 74 75 76 77 78 79 80

To solve this dilemma, we adopted a "bimodal" strategy to the fabrication of dense LLZO samples. As shown in Figure 1c, by seeding unpulverized powders (mean diameter of approximately $20 - 30 \mu m$) into the green body with submicron-sized pulverized powders, we achieved inhomogeneous grain growth during the sintering process. The result was a microstructure with bimodally distributed grain sizes. Fine grains filled in the gaps between coarse ones, so that large interconnected pores are effectively avoided. The microstructure with a bimodal grain size distribution, together with finely dispersed pores, increased the tortuosity of Li dendrite paths resulting in a "buffer and detour" effect. As a result, without any additional interface modifications, the LLZO electrolytes with the bimodal microstructure operated stably at current densities higher than $1 \text{ mA} \cdot \text{cm}^{-2}$ for more than 2000 hours, and successfully cycled more than 100 hours at a current density of 2 mA·cm-2. Our results shed light on the rational design of the SSE microstructures and advancement of cycling performance of the all-solid-state Li batteries. 81 82 83 84 85 86 87 88 89 90 91 92 93 94

Figure 1. Illustration of polycrystalline SSE with different grain size distributions. 96

(a) Li dendrites propagate along relatively straight pathways inside the fine-grained SSE. 97

- (b) Interconnected pores encourage Li dendrite growth in the SSEs with coarse grains. (c) 98
- Expected "detour and buffer" effects to suppress Li dendrite growth in SSE with the 99

5

bimodal grain microstructure. 100

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Experimental 102

Synthesis of LLZO powder 103

The Al-doped LLZO powders were synthesized using a sol-gel method, employing LiNO₃ (AR, 99%; Aladdin), Al(NO₃)₃·9H₂O (AR, 99.0%; Aladdin), La(NO₃)₃·6H₂O (AR, 99.0%; Aladdin), and ZrO(NO₃)₂·xH₂O (99%; Sigma-Aldrich) as precursors. An extra $14wt\%$ LiNO₃ was added to offset any Li loss during the sintering process. These constituents were dissolved in a mixture of deionized water and citric acid. The resultant solution was boiled at 150 °C and then heated to 400 °C to form a uniform brown foam. This foam was crushed in an agate mortar and calcined at 750 °C for 2 h in an Al crucible to produce Al-doped LLZO powder with the cubic structure. 104 105 106 107 108 109 110 111 112

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Fabrication of bimodal LLZO pellets 114

The synthesized LLZO was sieved into two aliquots of fine and coarse powders. The fine powder was ball-milled at 400 rpm for 30 min and then air-dried. The particle size was measured using an Omic laser particle analyzer for unpulverized coarse powders and a Malvern Zetasizer Nano ZSE particle analyzer for pulverized fine powders. The size distributions of both powders are shown in Figure S1, having median diameters of 0.95 µm and 26.8 µm, respectively. Subsequently, the pulverized and unpulverized powders were mixed in specific weight ratios and pressed into pellets of 13 mm diameter. The weight percent of the unpulverized coarse powder in the green bodies were adjusted to 0%, 25%, 50%, 75%, and 100% for comparative analysis. These green bodies were sintered at 1180 °C for 6 h in Ar atmosphere to achieve densification. The densified pellets, denoted CP0, CP25, CP50, CP75, and CP100, respectively, were polished with sandpaper to clean the surface and eliminate potential surface impurities. The final LLZO pellets had a thickness of 850 µm and a diameter of approximately 10 mm. 115 116 117 118 119 120 121 122 123 124 125 126 127 128

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Microstructure characterizations 130

The phase purity of the LLZO was ascertained using a Shimadzu XRD-6100 131

diffractometer equipped with Cu Kα radiation. The synthesized LLZO powders and densified pellets were examined using scanning electron microscope (SEM) equipped with X-ray energy dispersive spectroscopy (EDS). A lab-based X-ray computed tomography (CT, Zeiss Xradia 610 Versa) was employed to analyze the pores in the densified LLZO pellets. The X-ray energy was set to 140 keV to ensure sample penetration. The voxel dimension, based on the setup, was 1.4-1.5 μm. More than 990 slices were collected for each sample, and tomographic reconstruction was performed using Avizo software from FEI. 132 133 134 135 136 137 138 139

Synchrotron X-ray nano-Laue-diffraction (nano-Laue) experiments were conducted on Beamline 21A at Taiwan Photon Source (TPS). A polychromatic X-ray beam (5 - 30 keV) was focused to 100×100 nm². The LLZO samples were mounted on a scan stage 45° oblique to the incident X-ray. The built-in SEM was utilized to locate the areas of interest. An area of $900 \times 900 \mu m^2$ was scanned with a step size of 5 μm on each sample. At each scanning position, a Laue diffraction pattern was recorded in reflection mode with a 2D detector situated 90° to the X-ray beam, and the fluorescence signals of Zr, Al and La were collected using a Si-drift detector concurrently. The Laue patterns were analyzed using a custom-developed algorithm^{36,} ³⁷ to map the crystal grain distribution on the sintered LLZO. 140 141 142 143 144 145 146 147 148 149

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Electrochemical tests 151

Ionic conductivity was measured via electrochemical impedance spectroscopy (EIS) tests. 10 mV AC voltage (0.1 Hz to 1 MHz) was applied at room temperature using a PARSTAT MC multichannel workstation. The top and bottom surfaces of the LLZO pellets were sputter coated with Pt as blocking electrodes. 152 153 154 155

Li symmetric cells were fabricated by affixing a Li foil on both sides of the LLZO pellets and assembling 2032-type coin cells in a glovebox. The Li foils were 8 mm in diameter and 600 μm thick. To secure the contact with Li metal, the assembled coin cells were annealed at 150 °C for 10 min. Galvanostatic and CCD tests were conducted using a Neware BTS 4000 at room temperature. 156 157 158 159 160

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Results 162

X-ray diffraction analyses of the synthesized LLZO powder are consistent with a 163

cubic phase in the space group Ia-3d (Figure S2a). Minor impurities, primarily $LiAlO₂$ and $LaAlO₃$, were detected in the LLZO powder, mainly due to the locally inhomogeneous chemical reaction. These impurities, as shown in Figure S3, will be eliminated during the sintering process, and thus their effects are ignored in this study. The morphology and chemical distributions of both types of powders are displayed in Figure S2b to S2f. 164 165 166 167 168 169

After sintering, all the samples, CP0 to CP100, are well-densified and retain the cubic structure (Figure S3). Their relative densities are listed in Table S1. The microstructural characteristics of the post-sintered samples, from CP0 to CP100, were investigated using nano-Laue. Combining the diffraction peak intensity and grain maps plotted in Figure 2a and b, which depict the microscopic defect density and crystal orientation distributions^{38, 39}, CP0 exhibits a normally distributed grain size across the scanned area and does not have a bimodal microstructure, with limited number of pores among the grains. CP25, CP50, and CP75 display increasing area of fine grains (indicated by the mosaic regions) adjacent to the large ones. Considering that each pixel in the map represents an area of $5 \times 5 \mu m^2$, most of the fine grains within the mosaic regions are less than 10 μm. CP100 exhibits the largest grains and pores. 170 171 172 173 174 175 176 177 178 179 180 181

Subsequently, the sizes of coarse grains were measured quantitatively after filtering fine grains and pores. Pixels within an individual coarse grain were counted to calculate the area and equivalent diameter. As illustrated in Figure 2c, the area fraction of coarse grains initially decreased when the ratio of unpulverized powders in the green body rose to 50%. However, it increased in CP75 and dropped again in CP100. In the statistical analysis (Figure 2d, with total number of measured coarse grains exceeding 100 for each sample), the grain size of all tested samples was predominantly below 120 μm. Specifically, in CP50, grain growth appears inhibited, resulting in the lowest median grain size among the five samples. The number of grains sized larger than 120 μm and smaller than 200 μm reduced as the ratio of unpulverized powder increased from 0 to 50%; however, in CP75 and CP100, there was more abnormal grain growth with sizes exceeding 200 μm. 182 183 184 185 186 187 188 189 190 191 192 193

The SEM images of the fracture surfaces of CP25, CP50, and CP75 (Figure 2e) indicate that the coarse grains are surrounded by fine ones, with CP50 displaying the 194 195

densest structure. As shown in Figure 2f, slight element segregation was observed between grains. The grey color in the maps are the pores. Areas enriched with La have 196 197

less Zr , an effect attributed to Li excess⁴⁰. 198

Figure 2. Grain study of LLZO with bimodal arrangement. Nano-Laue analyzed grain distribution as shown in (a) intensity map and (b) color-labeled orientation map. (c) Area fraction of coarse grains in all five samples. (d) Box chart of coarse grain size distribution, red line and black dots inside the boxes indicate the mean and median, respectively. (e) SEM-BSE images of the cross-sectional morphology. (f) X-ray fluorescence maps on the same nano-Laue scanned area of CP50. 200 201 202 203 204 205 206

The influence of the bimodal microstructure on pore morphology and distribution within the densified LLZO was investigated using CT. A $431 \times 453 \times 293$ μm³ volume was selected from a similar position in each fabricated sample for comparison. In Figure 3a, by labelling adjacent pores with distinct colors, it is clear that CP0 features large and interconnected pores, while the pores are mostly isolated in the samples from CP25 to CP100. In Figure 3b, where the translucent-scarlet shades are employed to represent pore volume distribution, the scarlet color is hardly 207 208 209 210 211 212 213

observed in CP50, proving that pore growth and coalescence was suppressed. Furthermore, CP50 possesses the smallest total pore volume $(1.78 \times 10^6 \,\mu\text{m}^3)$ among all five samples, whereas the total number of pores in CP25, CP50 and CP75 are similar (Figure 3c). Although CP0 has the least number of pores, CP50 is the densest structure (95.1% density, Table S1). In the violin plot (Figure 3d), pore size distribution is better visualized. CP50 demonstrates the narrowest pore size distribution, averaging at 7.2 μm. Even though CP100 has the lowest average pore size (4.3 μm), a number of large pores ($> 40 \mu$ m) still exist, resulting in the lowest density among all five samples. 214 215 216 217 218 219 220 221 222

Figure 3. Pore rearrangement and morphology study of LLZO samples with bimodal structures. (a) Morphology and distribution density of the pores are observed using labbased CT in sequence, with sampling volume of $431 \times 453 \times 293 \mu m^3$. (b) Translucentscarlet shades are employed to show pore volume distribution and connectivity in each sample. The evolution of (c) total volume, population and (d) equivalent size are statistically studied in each sample. 224 225 226 227 228 229

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For room-temperature electrochemical tests, all LLZO pellets underwent similar polishing processes to minimize artificial effects. The pellets were then sandwiched between Li electrodes under consistent pressure and heated at 150 °C for 10 min to 231 232 233

fabricate symmetric cells. As revealed in the Nyquist plot (Figure 4a), the curves consist of one partial semicircle in the high-frequency range and flattened semicircles at lower frequencies. The high-frequency semicircles represent the total resistance of the LLZO pellets, while the lower-frequency impedances are attributed to the resistances at the Li/LLZO interfaces $41,42$, which is mainly due to the interface morphology caused constriction resistance and charge transfer resistance^{$43,44$}. The measured total conductivities of all samples are about the same order of magnitude, ranging from 2 to 4×10^{-4} S·cm⁻¹ (see Table S1). The impedance data are fitted and presented as continuous curves, and the fitted resistances are listed in Table S2. Generally, from CP0 to CP50, the interface resistance drops from 603 to 203 Ω ·cm². In CP50, the combined effects of more fine grains, fewer pores, and smaller pore size contribute to the decrease of both constriction resistance and charge transfer resistance. The decreased coarse grain fraction promotes charge transfer properties, which has been corroborated by previous reports²⁶. In contrast, in CP75, the presence of abnormally-grown grains and increased volume fraction of pores leads to increased constriction and charge transfer resistance. As a result, the interface impedance of CP75 is several times higher than that of the CP50. CP100 exhibits a relatively low constriction resistance but high charge transfer resistance, probably because most pores in it are small (Figure 3d), while the abnormal growth of grains is not prohibited. Additionally, the bulk resistances of CP25, CP75 and CP100 are high, because of their increased pore volume fractions that increase the tortuosity of Li transfer path⁴⁵. Apart from decreasing interfacial resistance, the optimal bimodal microstructure plays a significant role in enhancing the cycling behavior. Figure 4b shows DC stepped current cycling of the symmetrical cells containing CP50. The current density was raised to as much as 1.1 mA·cm⁻² at 25 °C, without causing any detectable hard short. When the CP50 sample was thinned to 500 μm, the critical current density was further elevated to $1.6 \text{ mA} \cdot \text{cm}^{-2}$ (see Figure S4, SEM, impedance and cycling test results of this sample are also presented). The voltage drops are not proportional to current densities (Figure S5), indicating the drop of cell resistance, probably because Li infiltrates into LLZO and fulfills the pores on/near the interfaces during the stripping and plating processes $46-48$. This is equivalent to the reduction of the LLZO thickness and the improvement of interface contact. Another possible 234 235 236 237 238 239 240 241 242 243 244 245 246 247 248 249 250 251 252 253 254 255 256 257 258 259 260 261 262 263 264 265

explanation to the lowered resistance might be the modification of solid electrolyte interphase via passage of current, making it less resistive⁴⁹. 266 267

To assess the stability of the interfaces, cells were subjected to galvanostatic charging and discharging at a consistent temperature of 25 °C. Each half-cycle lasted 30 min. For a holistic comparison, all symmetric cells from CP0 to CP100 were tested and exhibited significant differences in cycling stability under a low current density of ± 0.1 mA·cm⁻² (Figure S6). As depicted in Figure 4d, the galvanostatic profile of the Li/CP50/Li symmetric cell displayed an initial potential of approximately ± 90 mV, which subsequently dropped to around ± 50 mV, cycling steadily over 1000h, outperforming all others. The decrease in potential indicates that Li penetrates into LLZO, forming a stable interface post activation cycles⁴⁶. When further cycling the Li/ CP50/Li symmetric cell at a current density of ± 0.3 mA·cm⁻², the cell voltage remained at ~70 mV for an extended cycling period of up to 1000 h. Another sample experienced a quick increase in current density. Remarkably, stable cycling was successfully achieved on it at high current densities of $1 \text{ mA} \cdot \text{cm}^{-2}$ and $1.2 \text{ mA} \cdot \text{cm}^{-2}$ (Figure 4e) for approximately 2000 h. In a previously reported system that showed a similar voltage profile to those in Figure 4f and 4g, Li filaments were observed inside LLZO, but their growth was reversible and cyclable⁵⁰. Subsequently, as the current density was increased stepwise in units of ± 0.1 mA·cm⁻² from ± 1.2 mA·cm⁻² to ± 2.0 mA·cm-2 with each increment lasting 100 h (Figure 4h1-h8), the voltage profile remained similar to that observed at lower current densities. As the current density increased, the deviation of the voltage profile from the expected square shape became more severe, suggesting that Li was stripped faster at the interface than it could diffuse from the bulk to compensate the stripped Li. The cycling started to lose stability with an obvious voltage drop (Figure 4e, see voltage profile between h7 and h8) when the current density was approximately 1.8 mA \cdot cm⁻², due to soft shorting; however, the symmetric cell did not fail even at $2.0 \text{ mA} \cdot \text{cm}^{-2}$. This suggests that the bimodal structure of CP50 prevented the propagation of Li filaments/dendrites or healed them, 268 269 270 271 272 273 274 275 276 277 278 279 280 281 282 283 284 285 286 287 288 289 290 291 292 293

ensuring robust cycling. With another CP50 symmetrical cell, the evolution of the impedances as a function of cycles was investigated. As displayed in Figure S7, the cell (Cell #1) was cycled with the current densities ramping up in a stepwise manner from 0.3 to 1 mA·cm⁻² (Figure S7a). At 0.3 mA·cm⁻², the interface resistance first decreased and then stabilized in the following cycles, as shown in Figure S7b and c. As the current density was increased from 0.3 to 0.8 mA·cm⁻², the impedance became lower (Figure S7d), indicating that more Li had infiltrated into LLZO. The evolution of solid electrolyte interphase might also play a role. Soft shorting was detected from the impedance curve measured after 16 cycles at 1 mA \cdot cm⁻². Afterwards, the impedance became unstable and dropped down rapidly (Figure S7e). In another symmetrical cell (Cell #2), the cycling current density was increased from 0 to 1.0 $mA \cdot cm^{-2}$ in 10 cycles, kept constant at 1.0 mA $\cdot cm^{-2}$ for 50 cycles, ramped up in a stepwise manner to 1.8 mA·cm⁻² and finally kept at 1.8 mA·cm⁻² for 77 cycles (Figure S8a). The evolution of impedance during the 50 cycles at 1.0 mA·cm⁻² and 77 cycles at 1.8 mA·cm⁻² was quite similar to that in Cell #1 under 0.3 mA·cm⁻². The impedance dropped slightly at the beginning and then reached a stable state after about 40 cycles (Figure S8b). 294 295 296 297 298 299 300 301 302 303 304 305 306 307 308 309 310

Figure 4. Electrochemical performance of the symmetric cells containing bimodal microstructured LLZO. (a) The ionic conductivity of Li symmetric cells containing samples with various types of bimodal microstructures was studied using EIS. (b) The CCD of the symmetric cell containing CP50 was measured by increasing the current stepwise and (c) a comparison of the CCD of cells containing the different samples. (d, e) Voltage profiles of the CP50 symmetric cells cycled under increased current densities, with zoom-in details shown in (f-h).

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- The cells containing CP50 outperform in every aspect, be it CCD (Figure 4c), overpotential, or cycling lifetime. For reference, the cycling stability and CCD test results for symmetric cells containing other bimodal microstructures are displayed in Figure S6 and Figure S9. Importantly, the interfaces in the tested cells are unadorned, i.e., free of any additional coatings. As demonstrated in Figure 5, the outcomes from

our CP50 specimens not only outperform the previously reported symmetric cells without interface and/or grain boundary decoration (in the blueish-grey area), but also surpass or at least run on a par with those incorporating interface modifications (in the yellow region). 326 327 328 329

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Figure 5. Galvanostatic cycling performance of Li symmetric cells with bulky garnet-type solid electrolytes (references are listed in Table S3). 331 332

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Discussion 334

The formation of bimodal microstructure 335

The bimodal microstructure was obtained by sintering a mixture of pulverized and unpulverized powders. In the mixed state, small particles fill the interstices between the large ones. During sintering, the inhomogeneous powder distribution results in non-uniform mass diffusion. The intra-agglomerate densification and grain growth occur in unpulverized powders more easily and rapidly, which consume the mass that is necessary for the pulverized fine powders to grow⁵¹. Moreover, sintering also influences the size, morphology, and spatial distributions of pores. As small particles coalesce, small pores are easily formed between them. On the other hand, when there are more large particles, more intragranular pores are trapped and cannot escape from the interior to the surface of the powder. According to experimental observations, grain growth takes place under the regular coarsening process in CP0, forming normally distributed coarse grains with large interconnected pores. CP25 has 336 337 338 339 340 341 342 343 344 345 346 347

more isolated pores but still suffers adverse effects from the interconnected ones. In CP75 and CP100, the unpulverized coarse powders become the principal component in green bodies, and the fast grain growth and fine grain consumption lead to more abnormal large grains. In CP50, a synergy is achieved, as the competing grain growths of unpulverized and pulverized powders are balanced. The coarse grains grow in a constricted manner, sparing the fine grains opportunities to sinter well. Additionally, the presence of excess Li can cause more pores to form $40, 52$, because gas is entrapped between particles during the rapid densification and volatilization of the Li-containing compounds employed to replenish Li loss. It should be noticed that the adding of Licontaining compounds which melt at low temperature, such as $LiNO₃$ and $Li₂CO₃$, can boost densification due to liquid-phase sintering. In summary, the initial powder size, grain growth rate, as well as Li excess affects the size and spatial distributions of pores. 348 349 350 351 352 353 354 355 356 357 358 359 360

With these observations, we surmise that a balanced ratio of pulverized and unpulverized powders, as in CP50, facilitates a bimodal microstructure of uniformly distributed coarse grains encapsulated by fine grains. Small and isolated pores, as well as a dense structure, are also obtained. Three guiding principles can be summarized for microstructure control concerning particle size effects in the green body: (1) adjacent large grains and rapid grain growth in polycrystalline SSE tend to create interconnected pores and intragranular pores; (2) introducing a certain size of large particles into the green body can significantly reduce interconnected pores via differential densification; (3) by controlling the ratio of two different types of powders, the size and population of pores can be adjusted, regardless of whether they are between or inside grains. 361 362 363 364 365 366 367 368 369 370 371

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Buffer and detour effects in bimodal microstructure 373

After cycling, the cells were disassembled to examine Li deposition behavior within the bimodal microstructure. Figure 6a shows the interface that distinguishes between areas with and without Li foil contact. As the Li foil diameter is 2 mm smaller than LLZO, the periphery not in contact with Li is clean, while the dark spots signify the existence of Li under back-scattered electron mode of the SEM image^{35, 53}. Magnified images reveal prominent Li deposition within grain boundaries and pores. 374 375 376 377 378 379

In Figure 6b, widespread dark color is readily detected in CP25, indicating conspicuous Li propagation due to the detour effect primarily caused by coarse grains. In contrast, dark spots are barely seen in CP50 due to the considerable area fraction of fine grains, leading to a finer and more complex Li propagation path. However, upon closer examination (Figure 6c), it can be seen that Li propagation in the fine grain area forms a net-like distribution. The fine grain areas, as well as fine pores, even those that are submicron size, act as Li metal reservoirs, thereby enhancing cycling stability under high current densities due to its buffering effect. 380 381 382 383 384 385 386 387

In recent studies, the presence of pores has been identified as one of the key factors for Li dendrite initiation and propagation^{17, 54, 55}. The size and distance between the pores and interfaces affect the current density that the SSE is able to withstand. Reducing the size and quantity of pores at the subsurface, and preventing the build-up stress caused by Li-filled pores from exceeding the local fracture toughness has been shown to be beneficial. Our bimodal microstructure design is consistent with this strategy. The introduction of coarse grains reduces subsurface pores and the chances of Li dendrite contact. Coarse LLZO grains with lower defect densities provide higher fracture toughness and thus can withstand higher current density. Furthermore, the size of subsurface pores is effectively reduced by the bimodal microstructure in which fine grains surround coarse ones, bringing the Li-pore-filling induced pressure to a magnitude far below the fracture toughness. Additionally, the pores along LLZO grain boundaries have the chance to form interconnected channels, which further lowers the pressure during plating. The benefit of the microstructure is that crack initiation is circumvented; even if cracking occurs, it consumes the energy for crack propagation and prevents hard shorts. 388 389 390 391 392 393 394 395 396 397 398 399 400 401 402 403

Generally, the reasons for the enhanced cycling stability in bimodal SSE can be divided into two aspects: (1) the fine grains surround the coarse ones, eliminating the interconnected pores and providing a zigzagging propagation pathway to hamper the direct and rapid infiltration of Li. In the meantime, the coarse grains lying on the infiltration route of Li detours its migration and delays the failure of SSEs; (2) pores filled with Li aid in moving the Li/LLZO interface forward, significantly increasing the Li contact area, thereby dispersing the interfacial current and facilitating uniform and reversible stripping and plating cycles (Figure 6d). This mitigates the void 404 405 406 407 408 409 410 411

formation and contact loss at the interface caused by high current density. Such improvements contribute to all-solid-state Li battery cycles without necessitating additional mechanical pressure for accelerating Li creep, which is deleterious to longterm cycling¹⁷. 412 413 414 415

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Figure 6. Postmortem analysis of Li deposition after electrochemical cycling. SEM-BSE images of (a) LLZO surface morphology and (b) cross-section after short circuit. (c) Li network propagation between the fine grains. (d) Schematic illustration of Li network infiltration between fine grains and in pores pushing the interface forward. 417 418 419 420 421

On the LLZO surface, we also notice spallation cracks (Figure S10), predominantly around the pores in the coarse grain areas and along the grain boundaries where large wedges are not entirely filled with fine grains. These cracks are caused by stress buildup during the pre-filling procedure of viscoelastic Li metal 54 , ⁵⁵. This suggests two aspects of future microstructure control. The intra-granular pores 422 423 424 425 426

should be minimized in both quantity and size, which can be achieved by reducing the excess Li in the green body and curtailing the growth rate of coarse grains. Fine grains 427 428

- should also be more tightly packed around the coarse grains. 429
- 430

Conclusion 431

In conclusion, our study focuses on the microstructure design and optimization of polycrystalline LLZO SSEs to enhance cycling performance, incorporating both CCD and long-term cycling stability. By mixing pulverized and unpulverized powders into the green body, a bimodal microstructure was obtained. Tuning the ratio of these two types of powders allows for the alteration of the distribution of grains and pores in the sintered samples, and thus influences the electrochemical properties significantly. When the pulverized and unpulverized powders are equally mixed, a microstructure in which coarse grains are surrounded by fine-grains and isolated fine pores are obtained, which results in efficient Li plating/stripping at a high current density. Long-term cycling stability experiments underscore the exceptional stability of symmetric cells containing the LLZO SSEs with the optimal bimodal microstructure, over 2000 h at current densities of 1.0 mA·cm⁻² and 1.2 mA·cm⁻². It even exhibits the capacity to endure over 100 h under a current density of $2 \text{ mA} \cdot \text{cm}^{-2}$. 432 433 434 435 436 437 438 439 440 441 442 443 444

Postmortem analysis reveals that intricate grain boundaries, formed by the fine grains in the bimodal structure, make the Li penetration path more complex, thereby reducing direct Li infiltration. Simultaneously, the coarse grains enhance the tortuosity of the Li path. With the "buffer and detour" effects, the LLZO electrolyte gains the potential to operate at elevated current densities. The propensity of Li to permeate between fine grains indirectly propels the interface forward, promoting reversible Li stripping and plating, and showcasing excellent cycling stability. Spallation cracks, primarily surrounding the pores in coarse grains and the edges of coarse grains without full contact with fine grains, are detected at the interface. These observations hint at future opportunities for microstructure optimization. In the realm of battery 445 446 447 448 449 450 451 452 453 454

applications, it is also essential to investigate the compatibility of the solid-state electrolytes with the cathode electrodes. This involves a thorough consideration of cathode materials, additives, interfacial decorations, and cathodic composite structure design. Our findings illuminate the path to rational and optimal design of bimodal structures, offering new insights on mitigating dendrite growth-induced failure in polycrystalline SSEs. 455 456 457 458 459 460

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References 473

- 1. J. Janek, W. G. Zeier, *Nat. Energy.* **2016**, 1, 1. 474
- 2. Q. Zhao, S. Stalin, C.-Z. Zhao, L. A. Archer, *Nat. Rev. Mat.* **2020**, 5, 229. 475
- 3. A. J. Samson, K. Hofstetter, S. Bag, V. Thangadurai, *Energy Environ. Sci.* **2019**, 12, 2957. 476 477
- 4. T. Famprikis, P. Canepa, J. A. Dawson, M. S. Islam, C. Masquelier, *Nat. Mater.* **2019**, 18, 1278. 478 479
- 5. J. Janek, W. G. Zeier, *Nat. Energy.* **2023**, 8, 230. 480
- 6. M. J. Wang, R. Choudhury, J. Sakamoto, *Joule* **2019**, 3, 2165. 481
- 7. T. Krauskopf, F. H. Richter, W. G. Zeier, J. r. Janek, *Chem. Rev.* **2020**, 120, 7745. 482
- 8. H. Shen, K. Chen, J. Kou, Z. Jia, N. Tamura, W. Hua, W. Tang, H. Ehrenberg, M. Doeff, *Mater. Today.* **2022**, 57, 180. 483 484
- 9. X. Han, Y. Gong, K. Fu, X. He, G. T. Hitz, J. Dai, A. Pearse, B. Liu, H. Wang, G. Rubloff, *Nat. Mater.* **2017**, 16, 572. 485 486
- 10. R. Dubey, J. Sastre, C. Cancellieri, F. Okur, A. Forster, L. Pompizii, A. Priebe, Y. E. Romanyuk, L. P. Jeurgens, M. V. Kovalenko, *Adv. Energy Mater.* **2021**, 11, 487 488
- 2102086. 489
- 11. N. J. Taylor, S. Stangeland-Molo, C. G. Haslam, A. Sharafi, T. Thompson, M. Wang, R. Garcia-Mendez, J. Sakamoto, *J. Power Sources.* **2018**, 396, 314. 490 491
- 12. K. Fu, Y. Gong, Z. Fu, H. Xie, Y. Yao, B. Liu, M. Carter, E. Wachsman, L. Hu, *Angew. Chem., Int. Ed.* **2017**, 56, 14942. 492 493
- 13. Y. Chen, Z. Wang, X. Li, X. Yao, C. Wang, Y. Li, W. Xue, D. Yu, S. Y. Kim, F. Yang, *Nature.* **2020**, 578, 251. 494 495
- 14. J. Kasemchainan, S. Zekoll, D. Spencer Jolly, Z. Ning, G. O. Hartley, J. Marrow, P. G. Bruce, *Nat. Mater.* **2019**, 18, 1105. 496 497
- 15. W. Ping, C. Wang, Z. Lin, E. Hitz, C. Yang, H. Wang, L. Hu, *Adv. Energy Mater.* **2020**, 10, 2000702. 498 499
- 16. X. Wang, J. Chen, Z. Mao, D. Wang, *Chem. Eng. J.* **2022**, 427, 130899. 500
- 17. Z. Ning, G. Li, D. L. Melvin, Y. Chen, J. Bu, D. Spencer-Jolly, J. Liu, B. Hu, X. Gao, J. Perera, *Nature.* **2023**, 618, 287. 501 502
- 18. E. J. Cheng, A. Sharafi, J. Sakamoto, *Electrochim. Acta.* **2017**, 223, 85. 503
- 19. G. Li, C. W. Monroe, *Phys. Chem. Chem. Phys.* **2019**, 21, 20354. 504
- 20. A. Sharafi, C. G. Haslam, R. D. Kerns, J. Wolfenstine, J. Sakamoto, *J. Mater. Chem. A.* **2017**, 5, 21491. 505 506
- 21. M. B. Dixit, M. Regala, F. Shen, X. Xiao, K. B. Hatzell, *ACS Appl. Mater. Interfaces.* **2018**, 11, 2022. 507 508
- 22. F. Shen, M. B. Dixit, X. Xiao, K. B. Hatzell, *ACS Energy Lett.* **2018**, 3, 1056. 509
- 23. K. Tantratian, H. Yan, K. Ellwood, E. T. Harrison, L. Chen, *Adv. Energy Mater.* **2021**, 11, 2003417. 510 511
- 24. H.-K. Tian, Z. Liu, Y. Ji, L.-Q. Chen, Y. Qi, *Chem. Mater.* **2019**, 31, 7351. 512
- 25. F. Okur, H. Zhang, D. T. Karabay, K. Muench, A. Parrilli, A. Neels, W. Dachraoui, M. D. Rossell, C. Cancellieri, L. P. Jeurgens, *Adv. Energy Mater.* **2023**, 13, 2203509. 513 514 515
- 26. L. Cheng, W. Chen, M. Kunz, K. Persson, N. Tamura, G. Chen, M. Doeff, *ACS Appl. Mater. Interfaces.* **2015**, 7, 2073. 516 517
- 27. D. K. Singh, A. Henss, B. Mogwitz, A. Gautam, J. Horn, T. Krauskopf, S. Burkhardt, J. Sann, F. H. Richter, J. Janek, *Cell Rep. Phys. Sci.* **2022**, 3, 101043. 518 519
- 28. H. Huo, Y. Chen, R. Li, N. Zhao, J. Luo, J. G. P. da Silva, R. Mücke, P. Kaghazchi , X. Guo, X. Sun, *Energy Environ. Sci.* **2020**, 13, 127. 520 521
- 29. Q. Dai, J. Yao, C. Du, H. Ye, Z. Gao, J. Zhao, J. Chen, Y. Su, H. Li, X. Fu, J. Yan, D. Zhu, X. Zhang, M. Li, Z. Luo, H. Qiu, Q. Huang, L. Zhang, Y. Tang, J. Huang, *Adv. Funct. Mater.* **2022**, 32, 2208682. 522 523 524
- 30. F. M. Pesci, R. H. Brugge, A. O. Hekselman, A. Cavallaro, R. J. Chater, A. Aguadero, *J. Mater. Chem. A* **2018**, 6, 19817. 525 526
- 31. Z. Huang, L. Chen, B. Huang, B. Xu, G. Shao, H. Wang, Y. Li, C.-A. Wang, *ACS Appl. Mater. Interfaces.* **2020**, 12, 56118. 527 528
- 32. S. Lee, K.-s. Lee, S. Kim, K. Yoon, S. Han, M. H. Lee, Y. Ko, J. H. Noh, W. Kim, K. Kang, *Sci. Adv.* **2022**, 8, eabq0153. 529 530

33. J. Su, X. Huang, Z. Song, T. Xiu, M. E. Badding, J. Jin, Z. Wen, *Ceram. Int.* **2019**, 45, 14991. 34. Y. Tian, Y. Zhou, Y. Liu, C. Zhao, W. Wang, Y. Zhou, *Solid State Ionics.* **2020**, 354, 115407. 35. Y. Ren, Y. Shen, Y. Lin, C.-W. Nan, *ACS Appl. Mater. Interfaces.* **2019**, 11, 5928. 36. J. Kou, K. Chen, N. Tamura, *Scr. Mater.* **2018**, 143, 49. 37. G. Zhou, W. Zhu, H. Shen, Y. Li, A. Zhang, N. Tamura, K. Chen, *Sci. Rep.* **2016**, 6, 28144. 38. R. Barabash, G. Ice, B. Larson, G. Pharr, K.-S. Chung, W. Yang, *Appl. Phys. Lett.* **2001**, 79, 749. 39. G. Zhou, J. Kou, Y. Li, W. Zhu, K. Chen, N. Tamura, *Quantum Beam Sci.* **2018**, 2, 13. 40. C. Schwab, G. Häuschen, M. Mann, C. Roitzheim, O. Guillon, D. Fattakhova-Rohlfing, M. Finsterbusch, *J. Mater. Chem. A* **2023**, 11, 5670. 41. H. Buschmann, J. Dölle, S. Berendts, A. Kuhn, P. Bottke, M. Wilkening, P. Heitjans, A. Senyshyn, H. Ehrenberg, A. Lotnyk, *Phys. Chem. Chem. Phys.* **2011**, 13, 19378. 42. L. Zhang, L. Cheng, J. Cabana, G. Chen, M. M. Doeff, T. J. Richardson, *Solid State Ionics* **2013**, 231, 109. 43. S. J. Cooper, A. Bertei, D. P. Finegan, N. P. Brandon, *Electrochim. Acta.* **2017**, 251, 681. 44. J. K. Eckhardt, P. J. Klar, J. Janek, C. Heiliger, *ACS Appl. Mater. Interfaces.* 2022, 14, 35545. 45. M. B. Dixit, M. Regala, F. Shen, X. Xiao, K. B. Hatzell, *ACS Appl. Mater. Interfaces.* 2018, 11, 2022. 46. X. Ji, S. Hou, P. Wang, X. He, N. Piao, J. Chen, X. Fan, C. Wang, *Adv. Mater.* **2020**, 32, 2002741. 47. H. Huo, J. Liang, N. Zhao, X. Li, X. Lin, Y. Zhao, K. Adair, R. Li, X. Guo, X. Sun, *ACS Energy Lett.* **2020**, 5, 2156. 48. B. Kinzer, A. L. Davis, T. Krauskopf, H. Hartmann, W. S. LePage, E. Kazyak, J. Janek, N. P. Dasgupta, J. Sakamoto, *Matter.* **2021**, 4, 1947. 49 Y. Xu, K. Dong, Y. Jie, P. Adelhelm, Y. Chen, L. Xu, P. Yu, J. Kim, Z. Kochovski, Z. Yu, *Adv. Energy Mater.* 2022, 12, 2200398. 50. E. Kazyak, R. Garcia-Mendez, W. S. LePage, A. Sharafi, A. L. Davis, A. J. Sanchez, K.-H. Chen, C. Haslam, J. Sakamoto, N. P. Dasgupta, *Matter.* **2020**, 2, 1025. 51. R. M. German, *Metall. Trans. A* **1992**, 23, 1455. 52. M. N. Rahaman, *Sintering of Ceramics*, CRC press, **2007**. 53. F. Aguesse, W. Manalastas, L. Buannic, J. M. Lopez del Amo, G. Singh, A. Llordés , J. Kilner, *ACS Appl. Mater. Interfaces.* **2017**, 9, 3808. 54. F. Han, A. S. Westover, J. Yue, X. Fan, F. Wang, M. Chi, D. N. Leonard, N. J. Dudney, H. Wang, C. Wang, *Nat. Energy.* **2019**, 4, 187. 531 532 533 534 535 536 537 538 539 540 541 542 543 544 545 546 547 548 549 550 551 552 553 554 555 556 557 558 559 560 561 562 563 564 565 566 567 568 569 570 571 572