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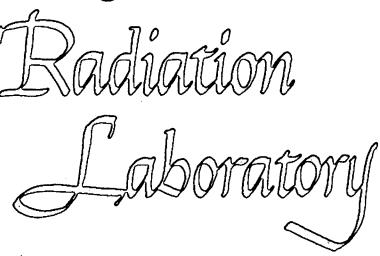
Glendenning, N.K. Sawicki, J.

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INTRINSIC QUADRUPOLE MOMENTS OF DEFORMED NUCLEI

N. K. Glendenning and J. Sawicki

August 1960

INTERNATIONAL CONFERENCE ON NUCLEAR STRUCTURE

AUGUST 29 - SEPTEMBER 3

1960

Secretary

Dr. L. G. Elliott

Atomic Energy of Canada Limited Chalk River, Ontario, Canada

RESEARCH CONTRIBUTION

Proposed Session No. No. 6 (Pairing Force and Nuclear Structure)

Title: Intrinsic Quadrupole Moments of Deformed Nuclei

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Abstract:

Intrinsic quadrupole moments of deformed nuclei are calculated on a particle picture as a function of Z and the deformation 6. Effects of the pairing force are examined. The single-particle levels are those of the Nilsson Model, but configuration mixing between major shells N' = N ± 2 must be included. Observed moments are accounted for by slightly smaller deformations than implied by the uniform-charge model.

INTRINSIC QUADRUPOLE MOMENTS OF DEFORMED NUCLEI

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July 12, 1960

Quadrupole moments of deformed nuclei have always been discussed in terms of the uniformly charged spheroid which establishes a connection between the quadrupole moment and the nuclear deformation. Since our knowledge of nuclear deformations rests principally on this connection, it is of interest to investigate the effects of the particle structure of nuclei on their quadrupole moments.

In this work we adopt a particle picture of the nucleus in calculating the quadrupole moments. The Z protons of the nucleus are assumed to move in a spheroidal well of fixed deformation and the interparticle interactions are approximated by the pairing force, introduced first to account for superfluidity by Bardeen, Cooper, and Schrieffer (BCS).²

Assume that a certain single-particle Hamiltonian representing the independent-particle motion in nuclei can be solved. Represent the quantum numbers of the eigenstates by $^{\nu}$ and the eigenvalues by E_{ν}^{0} . Then in terms of these eigenvalues, Migdal has shown how to compute the single-particle density matrix $P_{\nu\nu}$, corresponding to particles moving in the single-particle potential and interacting through the pairing force. If there is some additional perturbation V on the single-particle Hamiltonian, then the density matrix, up to first order in V, is

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$$P_{\nu\nu} = P_{\nu\nu}^{0} + P_{\nu\nu}^{1} \tag{1}$$

$$\rho_{\nu\nu}^{\circ} = \frac{1}{2} \left(1 - \frac{\epsilon_{\nu}}{E_{\nu}} \right) \delta_{\nu\nu}, \tag{2}$$

$$\rho_{\nu\nu}^{1} = \frac{\epsilon_{\nu}\epsilon_{\nu}^{1} - E_{\nu}E_{\nu} - \Delta^{2}}{2E_{\nu}E_{\nu}, (E_{\nu} + E_{\nu})} \quad \forall_{\nu\nu} = b_{\nu\nu}, \forall_{\nu\nu}, \qquad (3)$$

where

$$\epsilon_{\nu} = E_{\nu}^{\circ} - \lambda$$
, $E_{\nu} = \sqrt{\epsilon_{\nu}^{2} + \Delta^{2}}$ (4)

and A and A are the chemical potential and energy gap of the BCS theory.

We have computed the intrinsic quadrupole moment of a certain model nucleus comprising 2 protons moving in a permanently deformed potential well, and interacting through a pairing force. Because the deformation is assumed time-independent the results can have quantitative meaning for real nuclei only in the strongly deformed regions. Nevertheless the model may have at least qualitative meaning for weakly deformed nuclei near closed shells. For the basic states in Eqs. (1) - (4) we take the eigenvalues and eigenfunctions of Nilsson and Mottelson⁵ which result from an exact diagonalization of their Hamiltonian, except that only matrix elements that are diagonal in the oscillator quantum number N are included. That is, the deformation-dependent part of their Hamiltonian is

$$V_{0} = \frac{8}{3} \text{ M }_{0}^{2} Q_{20}. \tag{5}$$

It is, however, essential to include the effects on the quadrupole moment of the nondiagonal matrix elements of V_8 . This can be done by treating the nondiagonal part,

$$V_{NE} = V_{R} \left(1 - \delta_{NE} \right), \tag{6}$$

as the perturbation V in Eq. (3). Then the quadrupole moment is

$$Q_{0} = Q_{D} + Q_{ND}, \qquad (7)$$

$$Q_{D} = \sum_{N \mid \Omega \mid \alpha} \left(1 - \frac{E_{N \mid \Omega \mid \alpha} - \lambda}{\left(E_{N \mid \Omega \mid \alpha} - \lambda \right)^{2} + \Delta^{2}} \right) \left\langle N \mid \Omega \mid \alpha \right| Q_{20} \mid N \mid \Omega \mid \alpha \right\rangle, \tag{8}$$

$$Q_{ND} = -\frac{1}{3}8 M \omega_0^2 \sum_{N|\Omega|\alpha \alpha'}^{N|\Omega|\alpha \alpha'} b(N \Omega \alpha; N + 2, \Omega \alpha')$$

$$\times |\langle \mathbf{N} \, \mathbf{n} \, \alpha | \mathbf{Q}_{20} | \mathbf{N} + \mathbf{z}, \, \mathbf{n} \, \alpha' \rangle|^2, \tag{9}$$

where b is defined in Eq. (3), Q_{20} is the single-particle quadrupole moment operator, and the basic states $V = N \Omega \alpha$ are defined by Nilsson.⁵

The results of our calculation are displayed in the figure and in the table. The quantity Q_0/ZR^2 , which for the uniform-charge model is a constant $\left(4/58(1+8/2)\right)$ is found to decrease monotonically with increasing Z through the rare earth region. In the presence of the pairing force Q_0/ZR^2 is a smooth function of Z. For small deformation $(5 \le 0.1)$ it shows shell-closure effects as the pairing force is reduced. However, at the deformations appropriate to the rare earth region, the pairing force has very little effect on the calculated quadrupole moment.

As indicated in the table, our values of the deformation 8 for the lower rare earths, deduced by a comparison of the experimental quadrupole moments with our calculated values, are slightly smaller than those deduced via the uniform-charge model.

Table I

Intrinsic quadrupole moments $Q_0 = Q_0 + Q_{00}$ (in 10^{-2k} cm²) for some strongly deformed nuclei, compared with the observed moments Q_{obs} . The deformation, 8, deduced from this comparison is compared with the deformation implied by analyzing the experimental values of Q_0 with Q = k/5 ZR^2 8 (1 + 8/2) for the uniform-charge model.

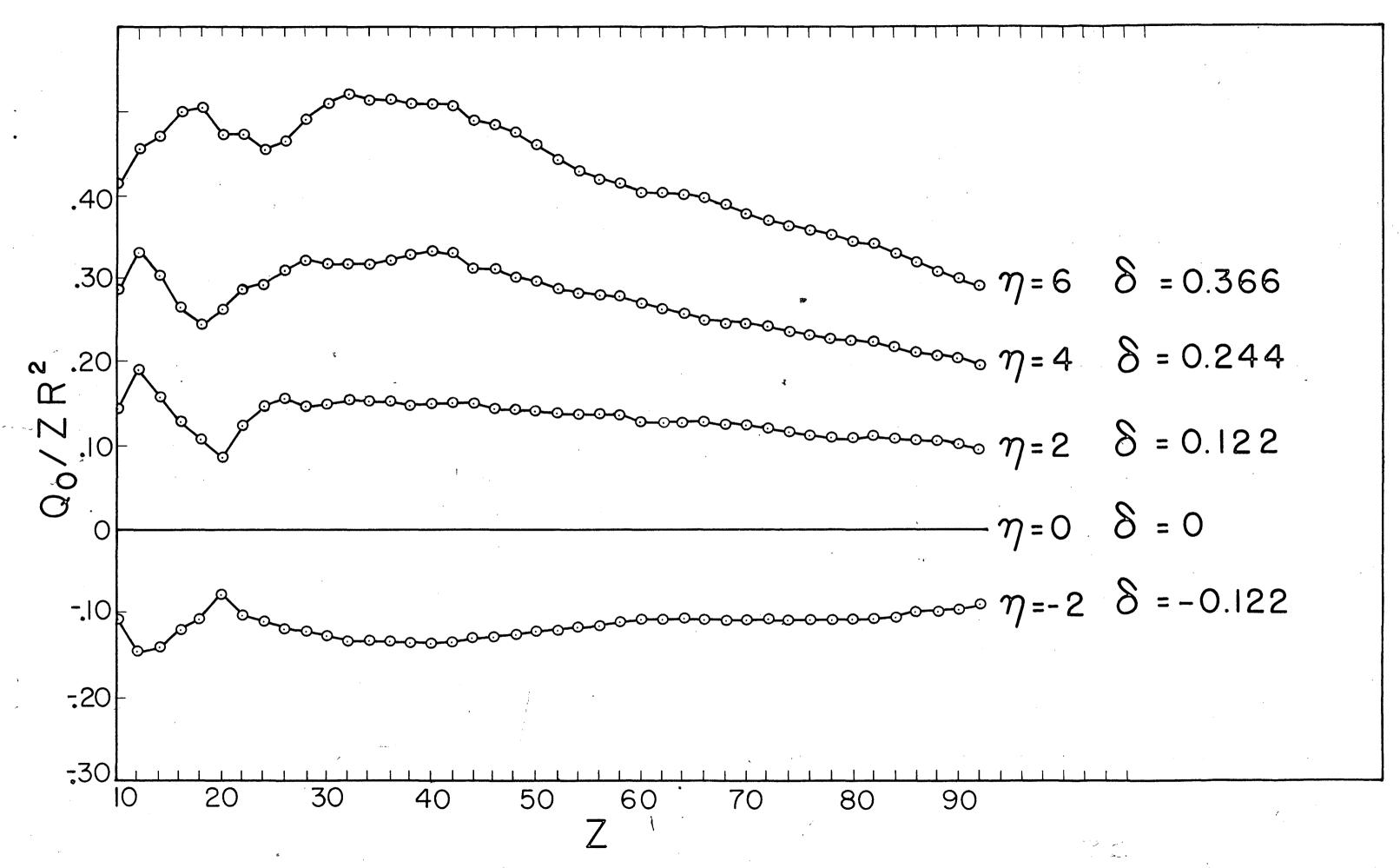
Z	A	Calculated Q			obs	8 (this calcu-	8 (uniform model)
		S = F	n = 4	η = 6		lation)	
78	24	0.35	0.63	0.94	0.45	0.24	0.33
62	150	3.25	6.60	10,11	3.62	0.13	0.17
62	152	3.29	6.71	10,26	5.82	0.20	0.26
64	154	3.41	6.85	10.58	5.75	0.20	0.28
64	156	3.45	6.90	10,72	6.75	0.23	0.28
64	158	3.44	6.91	10.67	7.35	0.25	0.30
64	160	3.45	6.94	10,71	7.59	0.25	0.30
66	160	3.60	7.10	11.07	6.65	0.22	0.26
66	162	3,63	7.13	11.17	7.14	0.23	0.28
66	164	3.62	7.16	11.16	7.48	0.24	0.29
68	164	3.76	7.35	11,41	7.07	0.23	0.26
68	166	3.78	7.39	11.46	7.52	0.24	0.28
68	168	3.78	7.41	11.50	7.55	0.24	0.28
68	170	3,89	7.44	11.55	7.34	0.23	0.27
70	172	3.91	7.72	11.79	7.42	0.23	0.26
72	176	3 .95	7.9 3	12.10	7.25	0.22	0.24
72	178	3.98	7.98	12.15	6.88	0.21	0.23
72	180	3.98	7.99	12.20	6.60	0.20	0.23
74	182	3.98	8.15	12.49	6.45	0.19	0.21
74	184	3.99	8.19	12.55	6.15	0.18	0.20

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FIGURE CAPTION

The quantity Q_0/ZR^2 , computed on the particle model of this article, is shown as a function of Z for several values of the deformation. For the uniform-charge model Q_0/ZR^2 has the constant values 0.346, 0.219, 0.103 for the deformations $\eta = 6$, 4, 2 respectively.



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