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Authors

Gilbert, W S Althaus, R F Barale, P J [et al.](https://escholarship.org/uc/item/5wx9k88p#author)

Publication Date

1990-09-01

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Lawrence Berkeley Laboratory

UNIVERSITY OF CALIFORNIA

Accelerator & Fusion Research Division

Presented at the 1990 Applied Superconductivity Conference, Snowmass, CO, September 24–28, 1990, and to be published in the Proceedings

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W.S. Gilbert, R.F. Althaus, P.J. Barale, R.W. Benjegerdes, M.A. Green, M.I. Green, and R.M. Scanlan

September 1990

Prepared for the U.S. Department of Energy under Contract Number DE-AC03-76SF00098.

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LBL-2891O SSC-MAG-286

NEW MEASUREMENTS OF MAGNETIC FIELD DECAY IN 1 METER SSC-TYPE DIPOLES*

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w. S. Gilbert, R. F. Althaus, P. 1. Barale, R. W. Benjegerdes, M. A. Green, M. 1. Green, and R. M. Scanlan Lawrence Berkeley Laboratory University of California Berkeley, CA 94720

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*This work was supported by the Director, Office of Energy Research, Office of High Energy and Nuclear Physics, High Energy Physics Division, U.S. Department of Energy under Contract No. DE-AC03-76SF00098.

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Lawrence Berkeley Laboratory

University of California
Berkeley, CA 94720

Introduction

Previous studies of magnetic field decay^{1,2} in model SSC dipoles due to changes in magnetization currents caused by flux creep have used the assumed SSC injection energy of 1TeV, or 0.33 tesla central dipole field, and an excitation to the storage field of 6.6 tesla. More recently, it has been decided to inject at 2 TeV, or 0.66 tesla and so more recent tests have been carried out at the new injection field, or at both the new and old fields. Additionally, the effect of temperature changes and excitation cycles on the field decay have been studied.

Temperature Effect

A sextupole decay curve, at an injection field of 0.33 tesla, for
dipole magnet D-15C-3 is shown in Fig. 1. The curve is roughly log
linear with a slope of 0.6 unit/decade. The temperature is estimated to be $4.3K$ within \pm 0.05K.

Figure 1. Sextupole Decay of Dipole D-15C-3 at 4.3K

One suggestion to prevent or reduce the field decay is to reduce the temperature after the excitation cycle has set up the magnetization sextupole field³. After setting up the injection field at 0.33 tesla 4.3K and allowing a half hour decay at 4.3K, we then cooled the magnet at the rate of 0.4K/hr. for the next four and a half hours. This is shown in Fig. 2. Fig. 3 is the bath temperature record. Not only did the decay stop, but the sextupole increased due to the increase in the critical current density at the lower temperature. We did not have the cryogenic control needed to quickly reduce the bath temperature a small amount and hold it constant at that lower temperature.

Figure 2. Sextupole Decay of Dipole D-15C-3 with Decreasing Temperature

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There are arguments⁴, based on tests of wire samples, that the flux creep is activated by temperature fluctuations. The experiment cited used a temperature rise of 0.5K after the magnetization currents have been set up and the authors observed a decay increase of a factor of three. Our temperature monitors in the helium bath and in the magnet body only indicate a temperature variation of ±0.05K over a
one hour period. Therefore, we doubt that this is a principal cause of the observed field decay in the dipole magnets. Fig. 4 shows the temperature variation of one of the carbon glass sensors in the helium bath. The excursions between 25 and 35 minutes are due to the initiation of the HeII cooling cycle.

We also have decay data on a SSC-model quadrupole. Fig. 5 shows the 12 pole decay with time. The decay is similar, though smaller in magnitude, to that seen in the dipoles. In this case, we have temperture data taken during the decay, both inside the magnet and in the helium bath. Both sensors are carbon glass. The helium bath is coupled to a helium refrigerator and fluctuations in the compressor return pressure result in changes in the bath temperature. This bath temperature record is shown as Fig. 6. The temperature sensor inside the magnet seems to lag the bath temperature by several minutes, and is shown as Fig. 7. Short term fluctuations are in the few millikelvin range, and over the longer time the maximum temperature change was less than 50 millikelvin. Fig. 8 shows the temperature inside the magnet for a portion of the set up cycles, to 6.6 tesla, for the first 35 minutes; and for the decay between 35 and 60 minutes.

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Effect of Maximum Magnet Field

It was observed, at DESY^{5,6}, that the observed field decay was
influenced by the maximum field the magnet was cycled to before the injection field was measured for field decay. In our dipole D-15B-2, we cycled at 4.3K to different peak fields before measuring the
decays at 0.33 tesla. The setup cycles are listed on Fig. 9. The maximum currents for the various cycles were 750, 1500, 3000, & 6000 amperes (6 tesla). The decay curves are shown as Fig. 9. The curves are not very good log linear, but if we assume straight line fits, we arrive at the decay rates of Table I.

Figure 9. Sextupole Decay of D-15B-2 at 320A

Table I

0.33 Tesla Injection Field

Max I	Decay rate units/decade
750	0.25
1500	0.30
3000	0.90
6000	1.25

Fig. 10 shows the same type behavior for the 2 TeV injection energy, or 0.66 tesla. The decay rates are shown in Table II.

Figure 10. Sextupole Decay of D-15B-2 at 640A

Table II

0.66 Tesla Injection Field

Clearly, we see that the decay rate is larger when the peak
current, or peak field is higher. Also the decay rate, and magnetization, are less when the higher injection field is used.

Fig. 11 shows the sextupole field decay for a larger 5cm aperture dipole, D-16B-1 with peak current values of 3000 and 6000 amperes and an injection current of 640 amperes. In this particular magnet, a passive superconducting corrector was installed in the bore tube so a large portion of the magnetization sextupole and decay were corrected out. Even so, we see that the decay for the 6000A setup cycle was almost double that of the 3000A sextupole, 0.27 unit/decade vs. 0.14 unit/decade.

COLD MEASUREMENTS OF D-16B-1 WITH PASSIVE CORR.
640 AMP, 1 HOURD DECATS - SKAMP SEAMP SETUPS
10N-THE-FLY, 10 S/C, 15A/S, 4.3K

Figure 11. Sextupole Decay of D-15B-1

Summary

The roughly log linear decay of magnetization current multipole fields has been demonstrated in several new dipoles. A suggested mechanism of thermally induced activation in the additional 0.5 Kelvin range has been effectively countered by temperature measurements that show temperature fluctuation ten to thirty times lower, over a ninety minute time period. Decay of dipole fields still do not appear to quantatively agree with magnetization decay in wire samples.

It was found that multipole field decay in dipole magnets depended on the maximum field reached on magnet excitation, with the decay faster for the greater field reached.

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LAWRENCE BERKELEY LABORATORY UNIVERSITY OF CALIFORNIA INFORMATION RESOURCES DEPARTMENT BERKELEY, CALIFORNIA 94720

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