Title
Morphology–Conductivity Relationship of Single-Ion-Conducting Block Copolymer Electrolytes for Lithium Batteries

Permalink
https://escholarship.org/uc/item/5z25671v

Journal
ACS Macro Letters, 3(6)

ISSN
2161-1653

Authors
Inceoglu, Sebnem
Rojas, Adriana A
Devaux, Didier
et al.

Publication Date
2014-06-17

DOI
10.1021/mz5001948

Peer reviewed
Morphology–Conductivity Relationship of Single-Ion-Conducting Block Copolymer Electrolytes for Lithium Batteries

Sebnem Inceoglu, Adriana A. Rojas, Didier Devaux, X. Chelsea Chen, Greg M. Stone, and Nitash P. Balsara

Department of Chemical and Biomolecular Engineering, University of California, Berkeley, California 94720, United States
Environmental Energy Technologies Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, United States
Materials Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, United States
Joint Center for Energy Storage Research (JCESR), Lawrence Berkeley National Laboratory, Berkeley, California 94720, United States
Malvern Instruments Inc., 117 Flanders Road, Westborough, Massachusetts 01581, United States

ABSTRACT: A significant limitation of rechargeable lithium-ion batteries arises because most of the ionic current is carried by the anion, the ion that does not participate in energy-producing reactions. Single-ion-conducting block copolymer electrolytes, wherein all of the current is carried by the lithium cations, have the potential to dramatically improve battery performance. The relationship between ionic conductivity and morphology of single-ion-conducting poly(ethylene oxide)-b-polystyrenesulfonylethyllithium (trifluoromethylsulfonylethyl)imide (PEO−PSLiTFSI) diblock copolymers was studied by small-angle X-ray scattering and ac impedance spectroscopy. At low temperatures, an ordered lamellar phase is obtained, and the "mobile" lithium ions are trapped in the form of ionic clusters in the glassy polystyrene-rich microphase. An increase in temperature results in a thermodynamic transition to a disordered phase. Above this transition temperature, the lithium ions are released from the clusters, and ionic conductivity increases by several orders of magnitude. This morphology–conductivity relationship is very different from all previously published data on electrolytes. The ability to design electrolytes wherein most of the current is carried by the lithium ions, to sequester them in nonconducting domains and release them when necessary, has the potential to enable new strategies for controlling the charge–discharge characteristics of rechargeable lithium batteries.

Single-ion-conducting polymer electrolytes, with ionic groups covalently bonded to the polymer and free counterions, are fundamentally different from conventional electrolytes wherein salts comprising positive and negative ions are dissolved in either liquid or polymeric solvents. A popular example of a single-ion-conducting electrolyte is hydrated NaFion wherein negatively charged sulfonate groups are fixed onto a fluorinated polymer backbone while the associated hydronium counterions are mobile. Transport of counterions in such systems occurs in the absence of concentration gradients provided that the bound ion concentration is uniform. This follows from the electroneutrality constraint. The performance of batteries and fuel cells with single-ion-conducting electrolytes is predicted to be superior to that obtained using conventional electrolytes, due to the absence of concentration polarization effects. Commercial lithium-ion batteries contain electrolytes where a majority of the current is carried by the ion that does not participate in energy-producing reactions. It has long been recognized that dramatic improvements in batteries can be realized if single-ion conductors with high ionic conductivity can be designed.

In this paper, we focus on single-ion conductors containing "mobile" lithium ions. Most of the studies on such materials are limited to single-phase systems wherein ionic groups are covalently attached to homopolymers. Dramatically different single-ion conductors based on block copolymers were first proposed by Ryu et al. and more recently by Bouchet et al. The work in ref 9 is based on polystyrene-b-poly(ethylene oxide)-b-polystyrene triblock copolymers, wherein lithium bis(trifluoromethane)sulfonamide (LiTFSI) was covalently linked to styrene units in the polystyrene blocks. Polystyrene (PS) and poly(ethylene oxide) (PEO) chains are highly incompatible. Extensive works on the thermodynamics of polystyrene-b-poly(ethylene oxide) (SEO) copolymers with...
added lithium salts, particularly LiTFSI, indicate that the tendency for microphase separation is enhanced by the presence of ions.\textsuperscript{11,12} One thus expects microphase separation in the triblock copolymers, but no evidence for the presence or absence of such effects was reported in ref \textsuperscript{9}. The ionic conductivities of these single-ion-conducting triblock copolymers\textsuperscript{3} were similar to that of mixtures of SEO and LiTFSI\textsuperscript{13} In contrast, the cation transference number of these triblock copolymers was in the vicinity of unity, well above typical values obtained in electrolytes containing lithium salts (between 0.1 and 0.5).\textsuperscript{14–18}

In Figure 1, we show a schematic of lithium ion migration in a microphase-separated poly(ethylene oxide)-b-poly(styrene)-polystyrenesulfonilimide diblock copolymer (PEO–PSLiTFSI) electrolyte. The most interesting aspect of these electrolytes is that “mobile” lithium counterions are ionically bound to the immobile and glassy polystyrene blocks. Transport of lithium ions can only occur if they migrate from the PS-rich microphase into a PEO-rich environment. The expected scenario is illustrated in Figure 1. The main purpose of this study is to describe the effect of lithium ion migration on morphology in single-ion-conducting block copolymer electrolytes.

![Figure 1. Schematic of lithium ion migration in a single-ion-conducting block copolymer electrolyte.](image)

Poly(ethylene oxide)-b-polystyrenesulfonilimide-(trifluoromethylsulfonyl)imide diblock copolymer (PEO–PSLiTFSI) were synthesized by following the previously reported procedures.\textsuperscript{9,10} In this study, we focus on a particular copolymer, PEO–PSLiTFSI(5.0–3.2) where 5.0 and 3.2 are the molecular weights in kg mol\textsuperscript{−1} of the PEO and PSLiTFSI blocks. The ratio of lithium ions to ethylene oxide monomers, $r$, is 0.088. We also present results obtained from a PEO–PSLiTFSI(5.0–2.0) diblock copolymer with $r = 0.056$.

Small-angle X-ray scattering (SAXS) intensities, $I$, as a function of the magnitude of the scattering vector, $q$, of PEO–PSLiTFSI(5.0–3.2), obtained during a heating run, are shown in Figure 2a. At room temperature, the scattering profile indicates the presence of a lamellar morphology, with a primary scattering peak at $q = q^* = 0.228$ nm\textsuperscript{−1}, and higher-order scattering peaks in the vicinity of $2q^*$, $3q^*$, and $4q^*$. The location of the higher-order peak in the vicinity of $2q^*$ differs substantially from the calculated value. Observed SAXS peak locations may be affected by the low-$q$ up-turn (Figure 2). Similar up-turns have been observed in other charged block copolymers.\textsuperscript{20} The domain spacing, $d$, defined as the center-to-center distance between adjacent PEO-rich lamellae is 28 nm ($d = 2\pi/q^*$). SEO copolymers with comparable molecular weights exhibit $d$ spacing in the vicinity of 15 nm.\textsuperscript{21} While the addition of charges on the chain is expected to increase $d$ spacing, the observed magnitude of this increase is surprising.

An additional SAXS peak is observed at $q = q_1 = 1.33$ nm\textsuperscript{−1} (Figure 2a). Similar peaks have been observed in a wide variety of ionomers including polystyrenesulfonate, NaNaf, and sulfonated polyethers.\textsuperscript{22,23} We thus conclude that the TFSI$^-$ and Li$^+$ ions are clustered at room temperature in PEO–PSLiTFSI(5.0–3.2). The average distance between clusters, $d_{\text{cluster}}$, is approximately 4.7 nm ($d_{\text{cluster}} = 2\pi/q_1$). Heating the sample results in a decrease of peak intensities as shown in Figure 2a. Scattering peaks indicative of the lamellar structure are absent at temperatures above 50 °C. The intensity of the ion cluster peak also decreases substantially above 50 °C (Figure 2b). Changes in the scattering intensity are reversible. The lamellar peaks are obtained during the cooling run at temperatures below 45 °C as shown in Figure 2c. The intensity of the ion cluster peak also increases as the sample is cooled.

![Figure 2. Temperature dependence of SAXS for the PEO–PSLiTFSI electrolyte. (a) SAXS intensity versus scattering vector, $q$, during a heating scan. (b) Scattering in the vicinity of the ion cluster peak during a heating scan. (c) SAXS intensity versus scattering vector, $q$, during a cooling scan. (d) Scattering in the vicinity of the ion cluster peak during the cooling scan of PEO–PSLiTFSI(5.0–3.2), $r = 0.088$.](image)
The lamellar morphology of PEO–PSLiTFSI(5.0–3.2) at room temperature was confirmed by dark field scanning transmission electron microscopy (STEM) as shown in Figure 3. The $d$ spacing determined by STEM is 23 nm which is in reasonable agreement with SAXS measurements. It is evident that PEO crystallization is confined within the microphase-separated structure.

The temperature dependence of the conductivity of PEO–PSLiTFSI(5.0–3.2) is shown in Figure 4. At room temperature, the conductivity is $3.0 \times 10^{-8}$ S cm$^{-1}$. The conductivity increases from $7.4 \times 10^{-8}$ S cm$^{-1}$ at 45 °C to $2.7 \times 10^{-7}$ S cm$^{-1}$ at 60 °C. A further increase in temperature results in a slower increase in conductivity, reaching a value of $3.8 \times 10^{-7}$ S cm$^{-1}$ at 90 °C. Also shown in Figure 4 is the temperature dependence of the SAXS intensity at $q = 0.228$ nm$^{-1}$ in the vicinity of the primary peak. There is a correlation between morphology and conductivity: high conductivity values are obtained when the sample is disordered and the normalized $I(q = 0.228$ nm$^{-1}$) is nearly zero. This suggests that the migration of the lithium ions into a PEO-rich environment coincides with homogenization of the block copolymer microstructure. Further work is needed to establish the quantitative relationship between SAXS profiles and conductivity. To verify the morphology–conductivity relationship described above, we performed in situ SAXS experiments wherein the conductivity was measured concurrently with the SAXS experiment during a heating run. The conductivity data obtained during this run are included along with three separate ex situ conductivity measurements to obtain the values and error bars shown in Figure 4. The SAXS experiments were also repeated several times.

The PEO–PSLiTFSI(5.0–3.2) electrolyte was placed between two lithium foils, and the current needed to sustain a potential drop of 80 mV ($\Delta V$) across the electrodes was measured as a function of time at 90 °C. The results of this experiment are shown in Figure 5. The current density values obtained when the experiment was started and after 1 h were $I^0 = 6.99 \times 10^{-2}$ and $I^\infty = 6.90 \times 10^{-2}$ mA cm$^{-2}$, respectively. The results of ac impedance measurements before starting the experiment and after its completion are shown in the inset in Figure 5. The Nyquist plot of the impedance data contains two semicircles, a small semicircle representing electrolyte resistance, $R_{\text{elect}}$, and a large semicircle representing interfacial resistances, $R_{\text{int}}$. The resistance values obtained before and after the cell polarization are $R_{\text{elect}}^0 = 127$ $\Omega$ cm$^2$, $R_{\text{elect}}^\infty = 119$ $\Omega$ cm$^2$, $R_{\text{int}}^0 = 1018$ $\Omega$ cm$^2$, and $R_{\text{int}}^\infty = 1028$ $\Omega$ cm$^2$, respectively.

The ionic conductivity values obtained from symmetrical cells using either lithium or aluminum electrodes are similar. In principle, determination of cation transfer number, $t^+$, requires knowledge of salt activity coefficients. In typical lithium battery electrolytes, the ratio $I^\infty/I^0$ obtained during the...
polarization experiment described in the preceding paragraph is about 0.3. In contrast, the value of \( \frac{I_\infty}{I_0} \) obtained from PEO−PSTLiTFSI(5.0−3.2) is 0.99. This implies that most of the current in our electrolyte is carried by the cation. In this case, the cation transference number, \( t^+ \), is approximately given by eq 1.

\[
 t^+ = \frac{I_\infty(\Delta V - I_0^0R_{\text{int}}^0)}{I_0(\Delta V - I_\infty R_{\text{int}}^\infty)} 
\]  

This gives \( t^+ = 0.95 \). To a very good approximation, our PEO−PSTLiTFSI(5.0−3.2) electrolyte is a single-ion-conductor.

The DSC, conductivity, and transference number results presented here are qualitatively consistent with those reported in ref 9. To assess the generality of our conclusions, we also studied the morphology−conductivity relationship in PEO−PSTLiTFSI(5.0−2.0) diblock copolymer with \( r = 0.056 \). The data obtained from this sample are very similar to those obtained from PEO−PSTLiTFSI(5.0−3.2) and thus presented in the Supporting Information, Figure S3.

The morphology−conductivity relationship derived from the results is pictured in Figure 6. At low temperatures, below 50 °C, we have an ordered microphase-separated structure with crystalline PEO-rich domains and glassy PSTLiTFSI-rich domains (\( T_g, \text{PSTLiTFSI} = 160 °C \)). The Li⁺ and TFSI⁻ ions form clusters in the glassy domains. The concentration of lithium ions in the PEO-rich domains is negligible, and this results in very low conductivity values (below 10⁻⁷ S cm⁻¹). At high temperatures, above 50 °C, we have a disordered morphology wherein amorphous PEO and PSTLiTFSI blocks are intimately mixed, and most of the ions are no longer in clusters. A majority of the lithium ions in this state are “mobile”, and conductivity values as high as 3.8 × 10⁻⁴ S cm⁻¹ are obtained at 90 °C.

In conclusion, we have shown that the transfer of lithium ions from the glassy PSTLiTFSI-rich domains into a PEO-rich environment results in disordered and highly conductive single-ion-conducting block copolymer electrolytes. The ability to design electrolytes wherein most of the current is carried by the lithium ions, to sequester them in nonconducting domains, and release them when necessary has the potential to enable new strategies for controlling the charge−discharge characteristics of rechargeable lithium batteries.

**ASSOCIATED CONTENT**

* Supporting Information
Additional experimental details, SAXS, and conductivity results of PEO−PSTLiTFSI(5.0−2.0) single-ion-conducting block copolymer electrolyte. This material is available free of charge via the Internet at http://pubs.acs.org.

**AUTHOR INFORMATION**

*Corresponding Author
E-mail: nbalsara@berkeley.edu.

**Notes**
The authors declare no competing financial interest.

**ACKNOWLEDGMENTS**

This work was primarily supported as part of the Joint Center for Energy Storage Research, an Energy Innovation Hub funded by the U.S. Department of Energy (DOE), Office of...
Science, Basic Energy Sciences (BES). X-ray scattering research at the Advanced Light Source was supported by DOE, Office of Science, BES. The electrochemical testing equipment was supported by the Assistant Secretary for Energy Efficiency and Renewable Energy, Office of Vehicle Technologies of the U.S. Department of Energy under Contract DE-AC02-05CH11231 under the Batteries for Advanced Transportation Technologies (BATT) Program. Work at the Molecular Foundry, Lawrence Berkeley National Laboratory was supported by DOE, Office of Science, BES under Contract No. DE-AC02-11231. STEM work was provided by the Electron Microscopy of Soft Matter Program from the Office of Science, Office of Basic Energy Sciences, Materials Sciences and Engineering Division of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231. The STEM experiments were performed as user projects at the National Center for Electron Microscopy, Lawrence Berkeley National Laboratory, under the same contract. We thank Dr. Eric Schaible for his help at Advanced Light Source, Lawrence Berkeley National Laboratory. We also gratefully acknowledge Jacob L. Thelen for helpful discussions.

REFERENCES

(17) Evans, J.; Vincent, C. A.; Bruce, P. G. Polymer 1987, 28 (13), 2324–2328.
(22) Yarusso, D. J.; Cooper, S. L. Macromolecules 1983, 16, 1871–1880.