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Nanoscale Iron Redistribution during Thermochemical Decomposition of CaTi1-x Fe x O3-δ Alters the Electrical Transport Pathway: Implications for Oxygen-Transport Membranes, Electrocatalysis, and Photocatalysis.

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# **Nanoscale Iron Redistribution during Thermochemical** Decomposition of  $\text{CaTi}_{1-x}\text{Fe}_x\text{O}_{3-\delta}$  Alters the Electrical Transport **Pathway: Implications for Oxygen-Transport Membranes, Electrocatalysis, and Photocatalysis**

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CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* maintained a stable perovskite phase at low Fe contents while exhibiting a phase decomposition to Fe/Fe oxide nanoparticles as the Fe content increases. In CaTi<sub>0.7</sub>Fe<sub>0.3</sub>O<sub>3−*δ*</sub> and CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−*δ*</sub>, the phase evolution to Fe/Fe oxide was greatly influenced by the temperature: Only temperatures of 300 °C and greater facilitated phase evolution. Fully coherent Fe-rich and Fedepleted perovskite nanodomains were observed directly by atomic-resolution scanning transmission electron microscopy. Prior evidence for such nanodomain formation was not found, and it is thought to result from a near-surface Kirkendall-like phenomenon caused by Fe migration in the absence of Ca and Ti co-migration. Density functional theory simulations of Fe-doped bulk models reveal that Fe in an octahedral interstitial site is energetically more favorable than in a tetrahedral site. In addition to coherent nanodomains, agglomerated Fe/Fe oxide nanoparticles formed on the ceramic surface during decomposition, which altered the electrical transport mechanism. From temperature-dependent electrical conductivity measurements, it was found that heat treatment and phase decomposition change the transport mechanism from thermally activated p-type electronic conductivity through the perovskite to electronic conduction through the iron oxide formed by thermochemical decomposition. This understanding will be useful to those who are developing or employing this and similar earth-abundant functional perovskites for use under reducing conditions, at elevated temperatures, and when designing materials syntheses and processes.

KEYWORDS: *oxide, perovskite, ceramic, earth-abundant, thermochemical reduction, decomposition, scanning transmission electron microscopy, electrical conductivity*

### **1. INTRODUCTION**

The earth-abundant, low-cost, and non-critical perovskite  $CaTiO<sub>3</sub>$  has been studied for its mechanical, optical, electrical, magnetic, and photocatalytic properties and for applications in oxygen-transport membranes (OTMs).[1](#page-10-0)−[9](#page-10-0) OTM materials are mixed oxygen ion- and electron-conducting solids that enable technologies for water splitting, value-added chemical production, oxyfuel combustion, and pre-combustion  $CO<sub>2</sub>$  separation, but functional membrane materials decompose under high oxygen chemical potential gradients.[10](#page-10-0)−[13](#page-10-0) Remarkably, even though the maximum oxygen permeation performance is inferior to state-of-the-art iron-cobaltite oxides, the oxygen flux of  $CaTi<sub>0.9</sub>Fe<sub>0.1</sub>O<sub>3−δ</sub>$  was shown to be more stable during a

heat treatments up to 600 °C. The results show that

semi-permeation experiment in a simulated service atmosphere of CO,  $CO_2$ ,  $H_2$ , and  $CH_4$ , where no evidence of membrane decomposition or reaction byproducts was found after 1600 h.<sup>[8](#page-10-0)</sup> This demonstrated the potential of Fe-doped  $CaTiO<sub>3</sub>$ 

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<span id="page-2-0"></span>

Figure 1. (a) Solid-state synthesis of porous ceramics was facilitated by the expansion of CaO upon reaction with H<sub>2</sub>O and/or CO<sub>2</sub> in air prior to sintering. (b) XRD confirmed that the orthorhombic perovskite-type structure was obtained for all CaTi1<sup>−</sup>*x*Fe*x*O3 samples where 0 ≤ *x* ≤ 0.4 (blue triangles). (c,d) Thermochemical reduction of the samples *x* = 0.3 and *x* = 0.4 caused secondary Fe/Fe oxide phasesto emerge at elevated temperatures (orange circles), as quantified in [Table](#page-4-0) 1 and detailed in [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf) S1.

(CaTi1<sup>−</sup>*x*Fe*x*O3−*δ*) as a more sustainable and non-critical material for OTM applications.

Considering the potential of OTM technologies to contribute to CO<sub>2</sub> utilization, CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> has attracted attention due to its durability in the presence of  $CO<sub>2</sub>$ . Introducing Fe and other transition metals as solutes in  $CaTiO<sub>3</sub>$  improves ionic and electronic conductivity $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  $\sqrt[4,7,14-17}$  and long-term stability due to its chemical and mechanical properties.[8,9](#page-10-0) In addition to optimizing the chemical composition, ceramic morphology control has proven useful: Adding porous CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> layers atop a dense membrane increases the active surface area and facilitates the impregnation of metal nanoparticles which increases oxygen permeability.[14](#page-10-0) However, CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* has also been shown to undergo thermal decomposition in  $H_2$  at elevated temperatures, forming metallic iron; such decomposition of OTM device components could affect surface phase stability and the oxygen surface exchange process that limits the overall oxygen flux.<sup>[14](#page-10-0),[18](#page-10-0)</sup> Thus, progress needs to be made to understand CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* phase stability in thermochemical applications, $13$  particularly at the unexplored atomic- and nanoscales and in porous ceramics.

Although it may challenge the implementation of CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* in OTMs, chemical instability presents an exciting opportunity in other contexts. For instance, the process of cation exsolution-which is akin to oxide phase decompositionprovides a finely tunable in situ synthesis route to durable metal nanoparticle catalysts embedded in the surface of various oxide supports.<sup>[19](#page-10-0)−[25](#page-10-0)</sup> Perovskite oxides readily incorporate transition



<span id="page-3-0"></span>

Figure 2. (a) As-sintered ceramics' high porosity was confirmed by SEM SE imaging. (b−d) Porosity was maintained after thermochemical treatments of  $x = 0.4$  after 300 °C (b),  $x = 0.3$  after 300 °C (c), and  $x = 0.3$  after 600 °C (d). Circles in (d) indicate the location of Fe/Fe-oxide phases formed on the ceramic surface, according to SEM-EDX and STEM analyses presented below.

metals into the lattice during material synthesis under oxidizing conditions; when subsequently exposed to reducing conditions, the oxide undergoes a controllable phase decomposition, and transition metals can be selectively exsolved as dispersed catalytically active nanoparticles. The exsolution method is increasingly popular as it can overcome typical drawbacksincluding nanoparticle agglomeration and deactivation—which are typical of traditional nanoparticle catalyst synthesis methods such as impregnation or vapor deposition. Very recently, it was found that Fe exsolution from  $SrTi_{0.65}Fe_{0.35}O_3$  thin films results in an Fe-depleted surface layer of ∼2 nm thickness.[26](#page-11-0) To further develop CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* as an earth-abundant catalyst support, exsolution of Fe (and/or other metals) could enable better catalytic activity and stability. It is thus useful to determine the critical temperatures and phenomena which occur during decomposition of the parent perovskite.

The tendency and mechanisms of CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−</sub>*δ* perovskites to decompose in reducing atmospheres at elevated temperatures is not well understood down to the atomic and nanoscale. Moreover, improved understanding of this process offers the potential to design advanced exsolution syntheses in this promising earth-abundant and non-critical perovskite. Defect chemistry studies have been performed on dense CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* ceramics down to very low oxygen partial pressures ( $pO_2$ ) and at temperatures around 1000 °C, assuming chemical phase stability. However, there are indications that this compound may not be stable in the absence of oxygen permeation flux, at which point Fe metal has been observedpresumably resulting from oxide decomposition. It is thus important to understand the decomposition mechanism and relevant processing conditions.

Here, we have done this in an unprecedented way by combining multiscale characterization down to the atomic scale, coupled with density functional theory (DFT) calculations, and electrical conductivity measurements over a range of chemical compositions and thermochemical heat-treatment conditions. We fabricated porous perovskite CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> ceramics with varied Fe B-site substitution and investigated the effects of thermochemical reduction by  $H_2$  diluted in Ar. We demonstrated variable phase evolution and decomposition of CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* over a range of temperatures and Fe doping amounts. Fully coherent Fe-rich and Fe-depleted perovskite nanodomains were, to our knowledge, observed directly for the first time by atomic-resolution scanning transmission electron microscopy (STEM) imaging and spectroscopy. Surprisingly, atomic-resolution elemental mapping by energy-dispersive Xray spectroscopy (EDX) revealed that Fe occupies interstitial sites in the Fe-rich nanodomains. DFT simulations were employed to identify the energetically favorable sites for these Fe interstitials. Along with the nanodomains, reduction yielded agglomerated Fe/Fe oxide nanoparticles on the ceramic surface, which proved consequential to the electrical transport mechanism. Electrical conductivity measurements by electrochemical impedance spectroscopy (EIS) demonstrated the effect of both Fe doping and thermochemical decomposition on the overall transport mechanism of CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub>, which was altered by nanoscale redistribution of Fe into Fe/Fe-oxide phases at the ceramic surface. Ultimately, correlating these experimental and computational methods allowed us to identify the critical Fe contents and temperatures of phase decomposition, the influence of decomposition on electrical transport, and to propose a multi-length-scale decomposition mechanism down to the atomic level.

#### **2. METHODS**

Porous CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> was synthesized by solid-state reaction of CaO, TiO<sub>2</sub>, and Fe<sub>2</sub>O<sub>3</sub> at 1500 °C for 2 h with a heating rate of 5 °C ([Figure](#page-2-0) [1a](#page-2-0)) following ref. $^{27}$  $^{27}$  $^{27}$  Stoichiometric amounts of precursor powders with  $0 \le x \le 0.4$  were mixed and ground by mortar and pestle and then pressed into pellets. Before sintering, the CaO in the pellets was allowed to react with air, likely forming a combination of  $Ca(OH)_2$  and  $CaCO_3$ , which caused the green body to increase in volume [\(Figure](#page-2-0) 1a) and ultimately introduced extensive porosity in the sintered ceramic (Figure 2).

## <span id="page-4-0"></span>Table 1. Evolution of Fe and CTO Phases after Heating in  ${{\rm H}_2}^a$



 $^a$ Tabulated are the ratios of XRD peak intensity for Fe{200}/CTO{200}. By quantifying the XRD peak intensity ratios of Fe{200}/CaTiO<sub>3</sub>{200}, it is apparent that the secondary phases emerged above 200 °C in  $x = 0.3$  and above 300 °C in  $x = 0.4$  (see also [Figures](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf) S1d and S2d).

Thermochemical reduction of samples  $0.1 \le x \le 0.4$  was performed in a quartz tube under flowing 5%  $H_2$  in Ar at 200–600 °C for 30 min within a Lindberg tube furnace. The samples were heated to each treatment temperature in air, then at temperature the reducing gas was flowed for 30 min before the samples were cooled in air to room temperature.

The crystal structure and phase compositions were measured before and after thermochemical processing using an Ultima III X-ray diffractometer with Cu K*<sup>α</sup>* radiation under 40 kV and 30 mA. The scanned range was  $2\theta = 20 - 90^\circ$  with a step size of  $2\theta = 0.05^\circ$ . The morphology and elemental distribution were investigated using a Magellan SEM at 25 kV and 25 pA equipped with an EDX spectrometer.

STEM was used to study the micro- and nano-scale morphology, as well as the atomic structure and elemental distribution in CaTi0.6Fe0.4O3−*<sup>δ</sup>* after thermochemical processing at 600 °C, as this sample showed severe decomposition during heat treatment. To prepare the samples for STEM, focused ion beam (FIB) lift-out was performed on an FEI Quanta 3d dual-beam SEM system. An ∼200 nm thick Pt protection layer was first deposited on the region of interest with the electron beam at 5 kV, followed by a ∼4 *μ*m ion beam deposition of Pt at 30 kV. To reduce the damage induced by the Ga ion beam, decreasing ion beam currents from 1 nA to 300 pA were used for the thinning steps. A final cleaning step was performed at 5 kV and 16 pA for 5 min on each side of the specimen. STEM characterization was carried out on a double aberration-corrected JEOL-ARM300F Grand ARM operated at 300 kV, equipped with double X-ray detectors for EDX elemental mapping.

Temperature-dependent EIS was used to measure the resistance of selected samples in air, which was then used to calculate the conductivity. Silver ink (Fuel Cell Materials) was coated on opposite sides of samples as porous electrodes. After coating, the samples were calcined at 100 and 800 °C in air for 1 h for drying and firing the silver paste, respectively (Carbolite Gero CC-T1). EIS tests were conducted using a potentiostatic mode with a frequency range of 0.1 Hz to 1 MHz (Gamry 1010E) connected to an electrochemical workstation (Huber Scientific, Plug and Probe) inserted into a tube furnace (Carbolite Gero CC-T1). During the EIS test, samples were heated from 25 to 425 °C in increments of approximately 50 °C for Arrhenius analysis. These temperatures were recorded with a thermocouple and used in Arrhenius fitting. About 10−15 EIS curves were measured at each temperature and averaged before being fitted to a resistor-constant phase element equivalent circuit model ("1RQ").<sup>[28](#page-11-0)</sup> Selected samples were fitted to equivalent circuit models comprising series combinations of 1RQ subcircuits ("2RQ"), or 1RQ in series with a resistor, to improve the fitting if there were multiple arcs visible. The highest frequency arc was interpreted as the sample's impedance contribution.

To interpret the atomic-resolution STEM-EDX of the Fe cation distribution observed in CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−</sub>*δ*, DFT calculations were performed. All spin-polarized calculations were performed using the Vienna ab initio simulation package, VASP 5.4.[29](#page-11-0)−[32](#page-11-0) The strongly constrained and appropriately normed (SCAN) meta-generalized gradient approximation (meta-GGA) was employed.<sup>[33](#page-11-0)</sup> The Atomic Simulation Environment (ASE) was utilized for model construction.<sup>3</sup> All atoms were described using the default projector augmented wave (PAW) potentials available in VASP.<sup>[35,36](#page-11-0)</sup> Non-spherical contributions related to density gradient in the PAW spheres were included in the calculations, with d-orbitals included in kinetic energy density mixing. Bulk calculations were carried out using Γ-centered (9 × 9 × 9) *k*-point

density. Gaussian smearing of 0.05 eV was used along with 400 eV cutoff energy. Lattice optimizations were carried out to determine bulk lattice parameters using a single cubic unit cell for  $CaTiO<sub>3</sub>$ . To approximately represent the Fe-doped system while keeping computations tractable, the unit cell was doubled and Fe was placed at either the tetrahedral or the octahedral interstitial site, along similar lines of a previous study with  $BaZrO_3$ .<sup>[37](#page-11-0)</sup>

#### **3. RESULTS AND DISCUSSION**

**3.1. Crystal Structure, Microstructure, and Morphology Analyses by XRD and SEM.** XRD patterns of all assintered CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> samples were indexed to the singlephase CaTiO<sub>3</sub> orthorhombic perovskite according to the ICDD database no. 00-042-0423 [\(Figures](#page-2-0) 1b and [S1\)](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf). The absence of additional Fe-containing phases confirms the compound's hightemperature phase stability in air when substituting Ti with Fe, consistent with past work on this solid solution $\delta$  and the fact that  $Fe<sup>4+</sup>$  and Ti<sup>4+</sup> have similar ionic radii.<sup>[7](#page-10-0)</sup> Thermochemical heat treatments under flowing 5%  $H_2/Ar$  resulted in the formation of additional Fe-containing crystalline phases at 200 and 300 °C in  $x = 0.3$  and  $x = 0.4$ , respectively, which were found to coexist with the initial  $CaTiO<sub>3</sub>$  perovskite-type structure by XRD ([Figures](#page-2-0) 1c,d and [S1](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf)). No Fe-containing crystalline phases were detected in the  $x = 0-0.2$  samples, which remained entirely the  $CaTiO<sub>3</sub>$  perovskite-type structure and were thus not further investigated here. The phase percentage of the perovskite  $CaTiO<sub>3</sub>$  was calculated using the XRD peak intensity for the peaks  $Fe\{200\}/CTO\{200\}$ , Table 1. The estimation assumes that the peak intensity ratio equals the phase concentration ratio, and the phase concentrations sum to  $1: \frac{1}{1}$ *C C*  $FeO_x\{200\}$   $\qquad \qquad$   $FeO_x$ CTO 200 FeO  $\frac{C_{\text{FeO}_x}}{C_{\text{CTO}}}$  and

#### $C_{FeO<sub>x</sub>} + C_{CTO} = 1$ .

A close inspection of the perovskite matrix  $\{200\}$  peaks in  $x =$ 0.3 and  $x = 0.4$  revealed a slight peak shift toward smaller Bragg angles with increasing heat-treatment temperature ([Figures](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf) S1b and [S2b](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf)). We estimated that this lattice expansion increased the unit cell volume of  $x = 0.3$  and  $x = 0.4$  from 0.161 to 0.166 nm<sup>3</sup>  $(+3%)$  and  $0.167$  nm<sup>3</sup>  $(+3.5%)$ , respectively ([Table](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf) S1). Expansion of the perovskite matrix can be attributed to an increased amount of Fe interstitial defects-for which we present direct experimental evidence below. Additionally, the ionic radii for 6-fold coordinated high-spin  $Fe^{4+}$ ,  $Fe^{3+}$ , and  $Fe^{2+}$ are 0.585, 0.645, and 0.78 Å, respectively.<sup>[38](#page-11-0)</sup> Although we do not have oxidation state experimental data for Fe, this too could contribute to the observed lattice expansion. It is not understood why the  $x = 0.4$  sample shows a peak in the amount of  $Fe/Fe$ oxide after thermochemical reduction at 400 °C.

The as-sintered porous ceramics had similar morphology according to SEM secondary electron (SE) imaging, [Figure](#page-3-0) 2a. The as-sintered  $x = 0.4$  porous ceramic surface is covered by small, ∼1 *μ*m, grains, which are assumed to be perovskite based on phase-pure XRD data, [Figure](#page-2-0) 1b,d. For this and all samples, we attempted SEM-EDX to map coarse surface agglomerates

<span id="page-5-0"></span>

Morphology and chemical mapping of  $x=0.4$  CaTi<sub>4</sub>  $x_F e_x O_{3,5}$  after 600 °C treatment

Morphology and chemical mapping of an Fe nanoparticle in porous Fe-FeO<sub>x</sub> phase



Figure 3. Nanoscale morphology and chemical analysis of CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−*δ*</sub> after thermochemical treatment at 600 °C revealed three distinct phases. (a) STEM-HAADF image near the ceramic surface (the FIB Pt protection layer is at the far left of the image and appears bright); the dashed lines indicate the borders between (1) the CaTi1<sup>−</sup>*x*Fe*x*O3 grain and the transition layer that contains Fe-rich and Fe-poor nanodomains and (2) the transition layer and the porous phase containing Fe metal and Fe oxide nanoparticles. (b−e) EDX mapsfor Ca, Fe, Ti, and O, respectively, of the region marked by a green box in (a). (f−h) STEM-HAADF image with Fe and O EDX maps of the border (dashed line) between the porous Fe/Fe-oxide nanoparticles and the transition layer. (i–m) STEM-HAADF image and EDX maps of an Fe nanoparticle in the boxed region in (f); see also [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf) S3 for the line profiles corresponding to the black arrow in (i).  $(n-p)$  Proposed decomposition mechanism of  $x = 0.3$  and  $x = 0.4$  CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3</sub> forms a coherent transition layer of Fe-rich and Fe-poor nanodomains and a porous phase of Fe metal and Fe oxide nanoparticles.

suspected to be Fe-rich but could not differentiate between Fe in surface agglomerates and Fe in the sub-surface perovskite matrix. It is unclear why these small grains formed on the *x* = 0.4 surface, though we speculate that the relatively high Fe content could contribute to solute drag effects which limit the incorporation of small grains into the porous superstructure. The morphology of the continuous solid shows negligible difference between the  $x = 0.3$  and  $x = 0.4$  at the initial stage of decomposition (after reduction at 300 °C) from SEM SE imaging [\(Figure](#page-3-0) 2b,c). In line with our XRD result, at later stages of decomposition, there was visible evidence in SEM SE images and STEM-EDX mapping (shown below) of morphology changes as a function of treatment temperature, such as the increased formation of surface features attributed to Fe metal and Fe oxide crystalline phases (e.g., *x* = 0.4 after 600 °C, [Figure](#page-3-0) [2](#page-3-0)d).

While our results on porous  $x = 0.3$  and  $x = 0.4$  ceramics are qualitatively consistent with findings by Salles et al.-who found that a dense  $CaTi<sub>0.9</sub>Fe<sub>0.1</sub>O<sub>3−δ</sub>$  ceramic underwent minor

decomposition to Fe metal during thermochemical reduction at [8](#page-10-0)50  $^{\circ} \text{C}^{\text{8}}$ —our samples underwent severe secondary phase formation at just 200−300 °C. Salles et al. observed partial reduction of Ti-substituting Fe<sup>4+</sup> to Fe<sup>3+</sup> (Fe<sub>Ti</sub>') under H<sub>2</sub> at 350  $\rm{^{\circ}C}$  via H<sub>2</sub> oxidation by lattice oxygen, followed by the formation of metallic Fe crystallites (visible by XRD) at 850 °C. Here, we hypothesize that the lower decomposition temperature is a combined effect of (i) higher Fe concentration and (ii) more surface area in our porous samples: the greater surface area facilitates oxygen vacancy  $(V_O^{\bullet\bullet})$  formation at the solid surface even at the low temperatures used in our study, which then electrostatically attract reduced Fe cations (Fe $'_{\rm Ti}$  and Fe $''_{\rm Ti}$ ) to the surface. Fe migration to regions of higher oxygen vacancy concentration is likely driven by their negative binding energy with oxygen vacancies  $\mathrm{Fe}_{Ti}^{\prime}-\mathrm{V}_{\mathrm{O}}^{\bullet\bullet}$  in dimer and trimer defect clusters.[7](#page-10-0) That we observe Fe metal and oxides via decomposition at significantly lower temperatures than Salles et al.'s Ca $Ti_{0.9}Fe_{0.1}O_{3-\delta}$  suggests that there is a balance between



#### <span id="page-6-0"></span>Atomic structure and chemical mapping of transition layer in x=0.4 CaTi<sub>1-x</sub>Fe<sub>x</sub>O<sub>3-5</sub>

Figure 4. Transition layer in CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−*δ*</sub> is coherent with the perovskite grain, with Fe<sup>2+</sup> interstitials as octahedral sites. (a,b) STEM-HAADF images showing the Fe-rich and Fe-poor nanodomains in the transition layer with average intervals of around 20 nm. The dashed line is the border between the CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−*δ*</sub> grain and the transition layer. (c) Atomic-resolution STEM-HAADF image taken at the region marked by a white box in (b) along the [011] direction. (d−g) STEM-EDX maps for Ca, Ti, Fe, and O, respectively, of the region marked by a green box in (c), at the interface between Fe-poor [left side of (f)] and Fe-rich [right side of (f)] nanodomains. (h−j) Composite EDX cation maps (d−g), EDX signal line scan along the white arrow depicted in (h) and averaged over the width of the box, and perovskite crystal model indicating Fe occupying sites on the Fe−O plane.

increased ceramic surface area, Fe concentration, and thermochemical stability.

**3.2. Nanoscale Morphology and Elemental Distribution after Decomposition of CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−***δ***</sub>. To further** understand the decomposition mechanism, STEM high-angle annular dark field (HAADF) imaging and EDX spectroscopy were used to analyze the nanoscale morphology and elemental distribution in the  $x = 0.4$  sample after reduction by  $H_2$  at 600 °C. This sample contained three noteworthy phases: Fedeficient CaTi0.6Fe0.4O3−*<sup>δ</sup>* perovskite grains, a so-called "transition layer" containing Fe-rich and Fe-poor nanodomains which are coherent with the perovskite, and a nanoporous agglomerate of intermixed Fe-metal and Fe-oxide nanoparticles ([Figure](#page-5-0) 3). STEM-EDX elemental maps indicate that the remaining grains of the perovskite ceramic contain primarily Ca, Ti, and O and are depleted of Fe—which apparently migrated to the oxide's surface during decomposition [\(Figure](#page-5-0) 3a−e). Surprisingly, on the surface of the perovskite grain, there is a ∼200 nm thick transition layer (outlined by white dashed lines in [Figure](#page-5-0) 3a−e) with a decreasing concentration gradient of Ca and Ti and an increasing concentration gradient of Fe. Within the transition layer, there is a non-uniform distribution of Fe, with so-called "Fe-rich nanodomains" and "Fe-poor nanodomains" appearing bright and dark, respectively, in both Fe EDX maps ([Figure](#page-5-0) 3c,g) and HAADF images [\(Figure](#page-5-0) 3a,f). Opposite to the transition layer from the perovskite grains is the nanoporous phase comprising an intimate mixture of agglomerated Fe metal and Fe oxide nanoparticles that is strongly/ entirely absent of Ca and Ti ([Figure](#page-5-0) 3f−m). The Fe nanoparticles appear bright in the HAADF-STEM images (e.g., [Figure](#page-5-0) 3i) and EDX maps [\(Figure](#page-5-0) 3k) due to the local enrichment of the heavier Fe atoms, while nanopores appear as dark areas in the HAADF image and Fe and O EDX maps. The complexity of the Fe nanoparticles likely warrants further investigation beyond this study; for instance, EDX of an individual Fe nanoparticle also revealed trace amounts of Ca in the particle and Ti accumulation on the particle surface perhaps because Fe and Ti form the stable ilmenite perovskite  $FeTiO<sub>3</sub>$ .

The proposed nanoscale phase decomposition mechanism observed in the  $x = 0.3$  and  $x = 0.4$  samples is summarized as follows ([Figure](#page-5-0) 3n−p): (1) high-temperature annealing in air during porous ceramic synthesis yields a single-phase perovskite solid solution with a negligibly low concentration of oxygen vacancies compensating reduced Fe and Ti, which are randomly distributed.<sup>[39](#page-11-0)</sup> (2) The H<sub>2</sub> thermochemical reduction environment reduces the  $pO_2$  such that Fe metal is in thermodynamic equilibrium with the oxide.<sup>[8](#page-10-0),[40,41](#page-11-0)</sup> Additionally, reduction by H<sub>2</sub> creates lattice oxygen vacancies  $(\mathrm{V}^{\bullet\bullet}_\mathrm{O})$  near the ceramic's surface that are charge compensated by Fe reduction from 4+ to 3+ and/or 2+  $(\overline{Fe}_{Ti}^{\prime} \text{ and/or } Fe_{Ti}^{\prime\prime})^{\text{8}}$  $(\overline{Fe}_{Ti}^{\prime} \text{ and/or } Fe_{Ti}^{\prime\prime})^{\text{8}}$  $(\overline{Fe}_{Ti}^{\prime} \text{ and/or } Fe_{Ti}^{\prime\prime})^{\text{8}}$  like the atomistic exsolution process used to synthesize surface nanostructures in situ. The high concentration of near-surface oxygen vacancies contributes an electrostatic driving force for migration reduced Fe cations through the perovskite toward the surface where they are reduced to Fe metal,  $7,22$  albeit at unexpectedly low temperatures for cation diffusion. (3) The Fe diffusing out of the perovskite is believed to originate in the transition layer, a volume of perovskite lattice that is both coherent with the underlying grain and riddled with Fe-poor nanodomains after decomposition. Although cation substitutional diffusion would be the expected Fe transport mechanism toward the surface in single-phase perovskite, $\frac{7}{1}$  this is apparently not entirely the case here. As supported by direct atomic-resolution EDX maps and DFT calculations presented below, Fe apparently migrates out of the transition layer to the Fe metal/Fe oxide porous layer in part via interstitial diffusion in Fe-rich nanodomains observed in the transition layer.

**3.3. Atomic Structure and Elemental Distribution of the Transition Layer and Nanodomains in CaTi0.6Fe0.4O3**−*δ***.** Given our unexpected discovery of the transition layer, atomic-resolution STEM imaging and EDX coupled with DFT calculations were used to conclude that Fe in the Fe-rich nanodomains exists as  $\text{Fe}^{2+}$  interstitials, which are energetically most favorable in octahedral sites [\(Figure](#page-6-0) 4). The transition layer borders the perovskite CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−δ</sub> matrix and consists of atomically coherent Fe-rich and Fe-poor nanodomains 20−100 nm in size [\(Figure](#page-6-0) 4a−c). Fe-rich and -poor nanodomains show brighter and darker HAADF contrast, respectively, caused by different Fe contents, as confirmed by EDX ([Figure](#page-6-0) 4f). Atomic-resolution STEM-EDX confirms that Ca, Ti, and O ions exist in both types of nanodomains. To our knowledge, this is the first observation of such structurally coherent nanodomains formed during perovskite oxide decomposition. Nanodomain formation was not observed in the perovskite bulk and so is assumed to be associated with the near-surface decomposition process. A possible nanodomain formation mechanism is related to the nanoscale Kirkendall effect, where in this case Fe migration from the perovskite to the surface proceeds faster than the reverse process, thereby depleting the transition layer of Fe and creating the Fe-poor nanodomains. However, rather than creating Kirkendall voids (e.g., as seen in  $NiO<sup>42</sup>$ ), Fe-poor nanodomains form in a stable  $CaTiO<sub>3</sub>$ -like perovskite because Fe is the only cation diffusing from the transition layer to the surface.

Surprisingly, a closer inspection of the composite cation EDX maps indicated that Fe is located at cation interstitial sites in the Fe-rich nanodomains, suggesting local interstitial Fe diffusion

([Figures](#page-6-0) 4d−j and [S5\)](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf). The possibility that this observation is caused by projection through the perovskite and an additional Fe-rich oxide particle is precluded by the fact that the Ca/O and Ti/O EDX signal intensity ratios remain constant across the Ferich/Fe-poor phase boundary in the analyzed area. If the STEM specimen contained an additional Fe-rich oxide particle overlapping the perovskite, instead of Fe interstitials, the Ca/ O and Ti/O would decrease in the area occupied by the Fe-rich oxide particle because the contribution of Ca and Ti to the total cation signal would necessarily decrease, while the anion signal would remain approximately uniform. This direct spectroscopic evidence of Fe between two {0-11} planes of alternating Ca and Ti atoms indicates that Fe is located on the same perovskite lattice planes as the O sites within the Fe-rich nanodomain (red arrows in [Figure](#page-6-0) 4j). This means that Fe occupies either octahedral or tetrahedral interstitial sites in the perovskite lattice; we calculate the relative energies of these Fe interstitial defects using DFT below. This is quite surprising, given that cation interstitials and Frenkel disorder in the close-packed perovskite system are expected to be energetically unfavor-able<sup>[43](#page-11-0),[44](#page-11-0)</sup> at equilibrium (e.g., >3.6 eV formation energy<sup>[7](#page-10-0)</sup>). However, first-principles calculations by Polfus et al. have shown that because of their reduced ionic radius, divalent transitionmetal cations (e.g.,  $Ni^{2+}$ ) can dissolve into zirconate perovskite oxides with large lattice volumes by occupying the octahedral interstitial site that is square planar coordinated to four oxide ions on the cell edge face center. $37$  (Occupation of the tetrahedral interstitial site�which in our case would be coordinated to Ca and three oxide ions-was found to be unfavorable in ref.<sup>[37](#page-11-0)</sup>) Thus, ours could be an observation of Fe<sup>2+</sup> migration via octahedral interstitials or the system in a metastable state amenable to cation interstitials.

**3.4. Determination of the Energetically Favorable Fe Interstitial Location by DFT.** The surprising existence of interstitial Fe cations in the Fe-rich nanodomain is interpreted initially following the work of Polfus et al., who showed that  $Ni<sup>2+</sup>$ interstitial solutes in  $BaZrO<sub>3</sub>$  are more energetically favorable at octahedral sites than tetrahedral sites according to  $DFT$ .<sup>[37](#page-11-0)</sup> Following this approach and considering our direct experimental evidence of interstitial Fe, we employed DFT calculations to suggest which interstitial site-octahedral or tetrahedral—is the most energetically favorable for  $Fe^{2+}$  to occupy. We only considered Fe interstitials based on direct evidence from STEM-EDX mapping [\(Figure](#page-6-0) 4i,j), which are effective positive charges (e.g., Fext). We assumed charge compensation by Ti reduction from  $Ti^{4+}$  to  $Ti^{3+}$   $(Ti_{Ti}^2)$  or  $Ti^{2+}$  $(Ti_{\text{Ti}}^{7/7})$ , and that interstitial Fe was Fe<sup>2+</sup> as this would require less total charge compensation than interstitial  $Fe<sup>3+</sup>$ .

DFT simulations optimizing the bulk  $2 \times 2 \times 2$  CaTiO<sub>3</sub> supercell showed that Fe occupying an octahedral interstitial site is 5.8 eV lower in energy than that in a tetrahedral interstitial site ([Figure](#page-8-0) 5), in agreement with the work of Polfus et al. Fe at the tetrahedral site leads to uniform lattice expansion along all directions (lattice parameter increases from 7.726 Å in  $CaTiO<sub>3</sub>$ to 7.822 Å upon Fe doping). Fe at the octahedral site distorts the cubic structure, leading to lattice contraction along two dimensions (7.673 Å) and expansion along the third (7.856 Å). Bader analysis is performed to assess charge compensation effects among cationic species. Fe is slightly more cationic at the tetrahedral site  $(+0.73 \text{ e})$  relative to the octahedral site  $(+0.58 \text{ e})$  $e)$ .<sup>[45,46](#page-11-0)</sup> However, it is difficult to draw meaningful conclusions regarding the mechanism of charge transfer. This is because changes in oxidation state are not always easy to discern from

<span id="page-8-0"></span>

Figure 5. DFT-optimized structures of Fe atom at the tetrahedral (left) and octahedral (right) positions of  $CaTiO<sub>3</sub>$ . Fe at the octahedral site is 5.8 eV lower in energy relative to Fe at the tetrahedral site.

Bader analysis. In this study, for instance, the reduction in charges on Ca (from +1.63 e in CaTiO<sub>3</sub> to +1.33 e for octahedral and  $+1.23$  e for tetrahedral) and Ti (from  $+1.89$  e in CaTiO<sub>3</sub> to +1.91 e for octahedral and +1.63 for tetrahedral) in species that are closest to the Fe cation are both small.

**3.5. Electrical Transport of CaTi<sub>1−***x***</sub>Fe<sub>***x***</sub>O<sub>3−***δ***</sub> and the Effect of Thermochemical Decomposition.** Because the perovskites CaTiO<sub>3</sub> and CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> are researched for applications relying on their electronic and mixed oxygen ionic and electronic transport (e.g., OTMs, photocatalysis, and electrocatalysis), we used temperature-dependent electrical measurements on select samples to both benchmark our materials and understand the effect of decomposition on transport properties (Figures 6 and [S4](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf)). EIS was performed at low and intermediate temperatures (<400 °C) on the asprepared CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−</sub> $_{\delta}$  samples (*x* = 0.1, 0.2, and 0.5) before thermochemical treatment and on the  $CaTi<sub>0.7</sub>Fe<sub>0.3</sub>O<sub>3−<sub>δ</sub>}</sub>$  sample after thermochemical treatments at 400, 500, and 600 °C (Figure 6a,b). This provided a convenient way to compare our as-fabricated materials' conductivity to that in the literature and to infer the effect of thermochemical treatment temperature and phase decomposition on the electrical properties. The EIS samples were prepared by applying liquid Ag paste to opposite sides of the ceramic pellets and then firing the sample at 800 °C in air to decompose all organics in the Ag ink solution, ultimately forming a porous Ag electrode. This firing procedure is also expected to oxidize any pure Fe metal particles (discussed above) in the heat-treated CaTi<sub>0.7</sub>Fe<sub>0.3</sub>O<sub>3−*δ*</sub>, yielding a porous ceramic composite of perovskite CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* and Fe oxide nanostructures.

As expected from the original publication by Iwahara et al., $^5$  $^5$ and more recent work by Salles et al.<sup>[17](#page-10-0)</sup> and others,<sup>[15,16](#page-10-0),[48](#page-11-0)</sup> the conductivity of our as-prepared CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* samples (before thermochemical treatment) is attributed to thermally activated electron hole conduction in the measurement temperature range and  $pO<sub>2</sub>$  of air. From Arrhenius plots of  $ln(\sigma T)$  versus  $1/T$ , we measured conductivity activation energies of 0.45−0.67 eV (Figure 6c), which is 0.20−0.33 eV less than the reported activation energy for oxygen ionic conductivity in similar compositions of 0.87−1 eV (e.g., refs [4](#page-10-0)



Figure 6. Electrical conductivity of select CaTi<sub>1−x</sub>Fe<sub>x</sub>O<sub>3−*δ*</sub> samples before and after thermochemical decomposition showed a distinct change in the charge-transport mechanism. Representative EIS data acquired in air at various temperatures for (a) CaTi<sub>0.7</sub>Fe<sub>0.3</sub>O<sub>3−*δ*</sub> following thermochemical treatment at 400 °C and for (b) CaTi<sub>0.8</sub>Fe<sub>0.2</sub>O<sub>3−*δ*</sub> before thermochemical treatment. The solid lines in (a,b) are equivalent circuit models fitted to the measured data. The resistance values extracted from equivalent circuit modeling were converted to conductivities and plotted against temperature (d), which were fitted in Arrhenius plots to calculate the activation energies for conductivity (c). For the thermochemically treated samples (CaTi<sub>0.7</sub>Fe<sub>0.3</sub>O<sub>3−*δ*</sub>), the conductivity below and above 150 °C exhibits the opposite temperature dependence consistent with electron transport through Fe oxide reported by Ito et al.; reproduced with permission from ref [47](#page-11-0).

<span id="page-9-0"></span>and [17](#page-10-0)). Also, like Iwahara et al.,  $x = 0.4$  showed the highest conductivity compared to both  $x = 0.1$  and  $x = 0.2$ .

Interestingly, the combination of thermochemical treatment of CaTi0.7Fe0.3O3−*<sup>δ</sup>* and EIS electrode fabrication reversed the temperature dependence of conductivity below ∼150 °C ([Figure](#page-8-0) 6c,d) and increased the conductivity activation energy above ∼150 °C to 0.86−0.97 eV [\(Figure](#page-8-0) 6c). This is attributed to (1) the fact that thermochemical reduction heat treatment causes Fe and Fe oxide nanostructures to form on the perovskite surface and (2) that EIS electrode fabrication oxidizes the Fe. The conductivity that is therefore measured by EIS is electronic transport through the iron oxide formed during thermochemical treatment and/or during EIS electrode fabrication. Ito et al. have reported a very similar temperature dependence of the electrical conductivity of sintered  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub>, which showed a decrease in conductivity from 70 to 150 °C caused by phonon scattering that lowered the carrier mobility.<sup>[47](#page-11-0)</sup> At temperatures >150 °C, they measured the conductivity activation energy to be 1.05  $eV,47}$  $eV,47}$  $eV,47}$  which is only slightly greater than that measured in  $CaTi_{0.7}Fe_{0.3}O_{3-\delta}$  here. Thus, we attribute the conductivity after heat treatment to electron conductivity through iron oxide formed via decomposition by  $H_2$  and subsequent oxidation.

Alternative explanations for the transport behavior of  $x = 0.3$ after heat treatments were deemed unlikely. First, below 150 °C, conductivity decreases with increasing temperature, which could be facilitated by band conduction through Fe metal, though this is unlikely considering that exposing the sample to air at 800 °C during preparation of the Ag EIS electrode should oxidize any metallic iron created during thermochemical treatments. Second, above 150 °C, conductivity increases with temperature, as expected for a thermally activated conduction mechanism. The conductivity activation energy calculated in this temperature regime ranges from 0.86 to 0.97 eV for 400 and 600 °C treatments, respectively, which could be interpreted as oxygen anion conduction through the perovskite, for example, 0.91 to 1.19 eV from ref.[16](#page-10-0) However, the electrolytic domain of  $CaTi_{1-x}Fe_xO_{3-\delta}$  is typically reported to be below pO<sub>2</sub> = 10<sup>-7</sup> atm, which is several orders of magnitude lower than that of our air measurements and likely precludes oxygen conduction as a relevant transport mechanism in this study.

#### **4. CONCLUSIONS**

This study reveals unprecedented details about the nano- and atomic-level phase evolution of porous CaTi<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>3−*δ*</sub> ceramics caused by thermochemical reduction heat treatments and the implication of decomposition on electrical transport. This is a crucial property for applications including OTMs, electrocatalysis, and photocatalysis. By varying the amount of Fe in the system, we find that CaTi1<sup>−</sup>*x*Fe*x*O3−*<sup>δ</sup>* undergoes a complex phase decomposition to agglomerated Fe/Fe oxide nanoparticles accompanied by the creation of a non-uniform transition layer near the surface in higher Fe-containing systems. XRD analysis after thermochemical treatments show secondary phases which were indexed to Fe/Fe oxide. SEM, STEM, and EDX revealed the agglomeration of Fe nanoparticles due to phase decomposition, which is facilitated by iron cation migration. Interestingly, Fe-rich and Fe-poor nanodomains were created in a transition layer between the single-phase perovskite and the porous nanostructured surface layer comprising agglomerated Fe and Fe oxide nanoparticles. This has been probed in the exemplary system CaTi<sub>0.6</sub>Fe<sub>0.4</sub>O<sub>3−*δ*</sub> with atomic spatial resolution and DFT calculations, which together revealed that Fe is likely located at octahedral interstitials square

planar coordinated by oxide ions in the Fe-rich nanodomains. Prior evidence for such nanodomain formation was not found and is thought to result from a near-surface Kirkendall-like phenomena caused by Fe migration in the absence of Ca and Ti migration. The effect of decomposition on the electrical properties of the samples was determined using EIS, which revealed that the charge-transport mechanism changes from thermally activated p-type conductivity through the perovskite to, most likely, electronic conduction through the iron oxide formed by thermochemical decomposition. The knowledge produced by this study is useful to readers who are interested in using electrochemical and functional perovskites under reducing conditions and/or at elevated temperatures and designing nanoscale materials via phase decomposition such as by the exsolution mechanism. While there is no evidence of submerged or partially submerged nanostructures resulting from the exsolution mechanism, the formation of agglomerated Fe nanoparticles and the Fe-poor/-rich nanodomains observed in this study expands our understanding of key aspects of the temperature-dependent thermochemical phase decomposition of CaTi1<sup>−</sup>*x*Fe*x*O3−*δ*.

#### ■ **ASSOCIATED CONTENT**

#### $\bullet$  Supporting Information

The Supporting Information is available free of charge at [https://pubs.acs.org/doi/10.1021/acsanm.2c04537](https://pubs.acs.org/doi/10.1021/acsanm.2c04537?goto=supporting-info).

XRD patterns and unit cell volume change of CaTi0.7Fe0.3O3−*<sup>δ</sup>* and CaTi0.6Fe0.4O3−*<sup>δ</sup>* after thermochemical treatments of 200−600 °C; EDX elemental signal profile along the line in Figure 3i; and representative EIS measurements for all samples across three temperatures ([PDF](https://pubs.acs.org/doi/suppl/10.1021/acsanm.2c04537/suppl_file/an2c04537_si_001.pdf))

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#### **Author Contributions**

W.J.B. conceived the study. J.L. synthesized the materials, performed thermochemical experiments, and performed XRD, SEM, and SEM-EDX characterization. X.W. performed STEM experiments. H.G. and B.X.L. performed EIS experiments. A.T. and N.H. performed DFT simulations, supervised by SMS. All co-authors contributed to the final version. W.J.B. supervised the study.

#### **Notes**

The authors declare no competing financial interest.

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