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December 1990



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# Microstructure of the High-T<sub>c</sub> Phase (T<sub>c</sub> ~111 K) in the SB-Pb-Bi-Sr-Ca-Cu-O System

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# Microstructure of the High- $T_c$ Phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O System

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The microstructure of the high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system has been determined using transmission electron microscopy. Its crystal structure belongs to the superspace group  $N_{1\bar{1}1}^{Bbmb}$ ,  $N_{111}^{Bb2b}$ ,  $P_{111}^{Bbmb}$  or  $P_{111}^{Bb2b}$  with subcell lattice parameters a=5.411(1)Å, b=5.411(1)Å and c=37.22(6)Å. The high- $T_c$  phase has a modulated structure with b-axis wavelengths 26.9Å and 36.1Å. Stacking faults along the c-axis in the high- $T_c$  phase are much less numerous than in the Bi-Sr-Ca-Cu-O system, but comparable to the Pb-Bi-Sr-Ca-Cu-O system. Sb substitution for Ca may affect the internal strain of the crystal.

Since the discovery of the high- $T_c$  superconductor in the Bi-Sr-Ca-Cu-O system,<sup>1</sup> various dopants have been used to try to increase its critical temperature. It has been reported that the substitution of Sb and/or Pb for Bi in Bi<sub>2</sub>Sr<sub>2</sub>Ca<sub>2</sub>Cu<sub>3</sub>O<sub>x</sub> increases the endpoint critical temperature ( $T_{c, end}$ ) to 132K.<sup>2, 3</sup> Hongbao *et al.* have reported that this phase is so unstable that the  $T_{c, end}$  is decreased to 112K upon forming a stable phase after thermal cycling between 77K and room temperature.

In another manuscript,<sup>4</sup> the superconducting properties and chemical composition of the stable high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system were described; the Sb-bearing superconducting phase has a 4K higher critical temperature than that in the Pb-Bi-Sr-Ca-Cu-O system. The average chemical composition of the high- $T_c$  phase is 4.3%, 2.6%, 19.2%, 21.4%, 15.8% and 36.9% for Sb, Pb, Bi, Sr, Ca and Cu, respectively. The summation of the Sb concentration and the Ca concentration is approximately the same for all the samples of this phase, implying that Sb substitutes for Ca, and oxygen atoms are introduced to compensate the oxygen deficiency in the central Cu-O layer sandwiched by the two Ca layers in the crystal structure of the high- $T_c$  phase. The chemical composition corresponds to Sb<sub>0.39</sub>Pb<sub>0.23</sub>Bi<sub>1.72</sub>Sr<sub>1.92</sub>Ca<sub>1.42</sub>Cu<sub>3.31</sub>O<sub>x</sub>, , which has more Sb, Bi, Cu and less Pb, Sr, Ca than the starting nominal composition of  $Sb_{0.10}Pb_{0.50}Bi_{1.40}Sr_{200}Ca_{200}Cu_{300}O_x$ . Large amount of impurities with chemical compositions of  $Sr_{0.2}Ca_{1.9}Cu_{1.0}O_x$ ,  $Sr_{1.0}Ca_{0.8}Cu_{3.0}O_x$  and  $Sb_{0.5}Bi_{2.0}Sr_{1.6}Ca_{1.0}Cu_{0.4}O_x$  were observed with a small amount of the high- $T_c$  phase in the sample. The difference between the composition of the high- $T_c$  phase and the nominal composition may be due to the segregation of these impurities. In order to further clarify the role of the Sb atoms in the high- $T_c$  phase, this detailed study of its microstructure has been undertaken.

A sample with nominal composition  $Sb_{0.10}Pb_{0.50}Bi_{1.40}Sr_{2.00}Ca_{2.00}Cu_{3.00}O_x$  was prepared by a solid-state reaction in air as described in another report.<sup>4</sup> X-ray powder diffraction measurements were carried out using a Siemens D500 diffractometer with CuK $\alpha$  radiation. Electrical resistivity was measured by the standard dc four-terminal

method. Magnetic susceptibility was measured using a S.H.E. SQUID magnetometer in a magnetic field of 20 gauss. Transmission electron microscopy was carried out using a JEM-200CX high resolution electron microscope and a JEM-200CX analytical electron microscope, both operated at 200kV. All specimens for electron microscopy were prepared by ion milling.

Figure 1 shows the temperature dependence of the electrical resistivity and the magnetic susceptibility of the material. The  $T_{c, end}$  and  $T_{c, mid}$  of the sample are 110K and 112K, respectively, while the critical temperature for the onset of diamagnetism is 111K. The sample therefore has a 4~5K higher critical temperature than the high- $T_c$  phase in the Pb-Bi-Sr-Ca-Cu-O system. The data also shows that the sample contains only the high- $T_c$  phase ( $T_c \sim 111$ K) and none of the lower- $T_c$  phase ( $T_c \sim 80$ K) of the Sb-Pb-Bi-Sr-Ca-Cu-O system. As mentioned before, however, the sample contains large amount of impurities which are not superconducting phases.

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Figure 2 shows four different zone axis electron diffraction patterns of the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system. Primary reflections with indices hkl (h+l=2n+1), 0kl (k=2n+1, l=2n+1) and hk0 (h=2n+1, k=2n+1) show extinctions, indicating that the space group of the average crystal structure of the high- $T_c$  phase is *Bbmb* or *Bb2b*, the same as that of the high- $T_c$  phase  $(T_c \sim 107K)$  in the Pb-Bi-Sr-Ca-Cu-O system.<sup>5</sup> Although Matsui *et al.*<sup>6</sup> have reported that the high- $T_c$  phase  $(T_c \sim 107K)$  in the Bi-Sr-Ca-Cu-O system has an ideal body-centered orthorhombic structure, a more accurate description is that this crystal structure belongs to either the *Bbmb* or *Bb2b* space groups because of the appearance of reflections with hkl indices satisfying (h+l=2n) in its diffraction patterns. The space group of the average crystal structure of the high- $T_c$  phase  $(T_c \sim 111K)$  in the Sb-Pb-Bi-Sr-Ca-Cu-O system is the same as the high- $T_c$  phase  $(T_c \sim 107K)$  in both the Pb-Bi-Sr-Ca-Cu-O and Bi-Sr-Ca-Cu-O systems.

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Satellite spots are clearly observed in the diffraction patterns of the  $a^*-b^*$  and  $b^*-c^*$  reciprocal planes. Figure 3 shows a magnified version of the electron diffraction pattern in

Fig.2(b) and the corresponding phase contrast image. Two types of satellite spots are observed in the electron diffraction pattern. One set of satellites is oblique to the  $b^*$ -axis, and is marked (A). The *b*-axis component of the corresponding modulation has a wavelength of 26.9Å. The other set of satellites occurs along the  $b^*$ -axis with a wavelength of 36.1Å, and is marked (B). The former modulated structure is the same as those observed in both the Bi-Sr-Ca-Cu-O system and Pb-Bi-Sr-Ca-Cu-O system, while the latter is closer to the one observed only in the Pb-Bi-Sr-Ca-Cu-O system. Many reports have been published on the modulated structure in the high- $T_c$  phase of the Pb-Bi-Sr-Ca-Cu-O system, citing wavelengths of 27Å and 45Å~54Å along the  $b^*$ -axis.<sup>79</sup> It is noted for comparison that although the direction of the modulation with the wavelength of 45Å~54Å in the Pb-Bi-Sr-Ca-Cu-O system is the same as the one (marked "B") in the Sb-Pb-Bi-Sr-Ca-Cu-O system, the wavelength is substantially smaller for the Sb-bearing material.

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Kabius *et al.*<sup>10</sup> have observed an endpoint critical temperature of 107K for their sample of the Sb-Pb-Bi-Sr-Ca-Cu-O superconductor, which is 4K lower than the temperature we have observed. Significantly, they also observed a modulated structure with a wavelength of 38Å along the  $b^*$ -axis in their specimen, which is slightly longer than that in our specimen. This may be due to differences in the preparation process of the samples. We used a long firing process at 865°C for 504 hours to obtain the high- $T_c$  phase with a critical temperature of 111K, while they employed a very short firing time (2 minutes) at higher temperature (950°C). Because we also observed a lower endpoint critical temperature for a sample prepared at 865°C for 24 hours than for one prepared at 865°C for 504 hours, we conclude that the modulated structure with the wavelength of 36.1Å is a characteristic feature of the high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system.

Assigning four indices H, K, L and m to both primary and satellite reflections in accordance with the procedure described by de Wolff *et al*,<sup>11</sup> we find that the allowed satellites marked (A) have indices HKLm (H+L+m=2n), HK0m (K=2n), and 0KLm

(K=2n) where H=h, K=k, and L=l+m. This corresponds to the superspace group  $N_{1\bar{1}1}^{Bbmb}$ or  $N_{111}^{Bb2b}$ , which is the same as that of the modulated structure observed in both the Bi-Sr-Ca-Cu-O system and Pb-Bi-Sr-Ca-Cu-O system. However, the allowed satellites marked (B) have indices HKLm (H+L=2n), HK0m (K=2n), and 0KLm (K=2n), where H=h, K=k, and L=l. This corresponds to the superspace group  $P_{1\bar{1}1}^{Bbmb}$  or  $P_{111}^{Bb2b}$ , which is the same as that of the modulated structure observed only in the Pb-Bi-Sr-Ca-Cu-O system.

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Lattice parameters of the average crystal structure of the high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system were determined by X-ray powder diffraction coupled with electron diffraction. The lattice parameters are a=5.411(1)Å, b=5.411(1)Å and c=37.22(6)Å, and are close to those of the high- $T_c$  phase ( $T_c \sim 107$ K) in the Pb-Bi-Sr-Ca-Cu-O system.<sup>5</sup> Although Sb<sup>3+</sup> has a smaller ionic radius than Ca<sup>2+</sup>, Sb substitution for about 20% of the Ca sites<sup>4</sup> does not cause significant shrinkage of the lattice parameter in the *c*-axis direction.

A stacking fault along the *c*-axis is observed in the high-resolution electron micrograph (marked "s" in Fig. 3(a)). It was found that the density of stacking faults in this crystal is much less than that in the Bi-Sr-Ca-Cu-O system but roughly the same as that in the Pb-Bi-Sr-Ca-Cu-O system. Previous studies suggest that Pb acts as a fluxing agent in the Bi-Sr-Ca-Cu-O system that decreases the density of the stacking faults and increases crystal size.<sup>12-14</sup> Given that Sb<sub>2</sub>O<sub>3</sub>, like PbO, is also a low-melting-point compound, it could be that Sb also acts as a fluxing agent, assisting the solid-state reaction that forms the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system with reduced stacking fault formation.

Table 1 shows a summary of the microstructures of the high- $T_c$  phases in the Bi-Sr-Ca-Cu-O system, the Pb-Bi-Sr-Ca-Cu-O system and the Sb-Pb-Bi-Sr-Ca-Cu-O system. The high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system has two types of structural modulations with wavelengths of 26.9Å and 36.1Å, similar to those of the high- $T_c$  phase in the Pb-Bi-Sr-Ca-Cu-O system. The high- $T_c$  phase ( $T_c \sim 111$ K) has a small number of stacking faults along the *c*-axis, equal in number to those in the high- $T_c$  phase of

the Pb-Bi-Sr-Ca-Cu-O system and much less than those in the high- $T_c$  phase of the Bi-Sr-Ca-Cu-O system. These results suggest that Sb not only acts as a flux in the solid state reaction but also substitutes directly into the lattice of the high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system, forming the characteristic modulated structures and a small number of stacking faults along the *c*-axis. As we mentioned in another report,<sup>4</sup> the chemical composition of the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system implies that Sb substitutes for Ca, and that oxygen atoms are introduced to compensate the resulting oxygen deficiency in the central Cu-O layer sandwiched by the two Ca layers in the crystal structure. Since the ionic radius of Sb<sup>3+</sup> is much smaller than that of Ca<sup>2+</sup>, Sb substitution for Ca may actually affect the internal strain of the crystal. More detailed observations of the crystal structure by high resolution electron microscopy are necessary to confirm the influence of the substitution of Sb for Ca on the crystal structure and the compensation of Sb for Ca on the crystal structure and the compensation of scores.

In conclusion, the microstructure of the high- $T_c$  phase ( $T_c \sim 111$ K) in the Sb-Pb-Bi-Sr-Ca-Cu-O system has been determined using transmission electron microscopy. The crystal structure of the high- $T_c$  phase has been found to belong to the superspace group  $N_{1\bar{1}1}^{Bbmb}$ ,  $N_{1\bar{1}1}^{Bb2b}$ ,  $P_{1\bar{1}1}^{Bbmb}$  or  $P_{111}^{Bb2b}$  with subcell lattice parameters a=5.411(1)Å, b=5.411(1)Å and c=37.22(6)Å, which is very similar to the high- $T_c$  phases in the Bi-Sr-Ca-Cu-O and Pb-Bi-Sr-Ca-Cu-O systems. The high- $T_c$  phase ( $T_c \sim 111$ K) has structural modulations with wavelengths of 26.9Å and 36.1Å. The former modulated structure is the same as those observed in the high- $T_c$  phases of both the Bi-Sr-Ca-Cu-O system and Pb-Bi-Sr-Ca-Cu-O system, while the latter one is closer (but much smaller) than that observed only in the high- $T_c$  phase of the Pb-Bi-Sr-Ca-Cu-O system. A few stacking faults along the *c*-axis are observed in the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system; however, the density of the stacking faults is much smaller than that of the high- $T_c$  phase in the Bi-Sr-Ca-Cu-O system. Sb substitution for Ca in the superconducting crystal appears to affect

the internal strain of the crystal, forming both the characteristic modulated structures and a small number of stacking faults along the *c*-axis.

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#### References

- <sup>1</sup> H. Maeda, Y. Tanaka, M. Fukutomi, and T. Asano, Jpn. J. Appl. Phys. 27, L209 (1988).
- <sup>2</sup> L. Hongbao, C. Liezhao, Z. Ling, M. Zhigiang, L Xiaoxian, Y. Zhidong, X. Bai, M. Xianglei, Z. Guien, R. Yaozhong, C. Zhaojia, and Z. Yuheng, *Solid State Commun.* 69, 867 (1989).

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- <sup>3</sup> M. R. Chandrachood, I. S. Mulla, and A. P. B. Sinha, Appl. Phys. Lett. 55, 1472 (1989).
- <sup>4</sup> N. Kijima, R. Gronsky, H. Endo, Y. Oguri, S. K. McKernan, and A. Zettl, unpublished.
- <sup>5</sup> N. Kijima, H. Endo, J. Tsuchiya, A. Sumiyama, M. Mizuno, and Y. Oguri, Jpn. J. Appl. Phys. 28, L787 (1989).
- <sup>6</sup> Y. Matsui, S. Takekawa, H. Nozaki, A. Umezono, E. Takayama Muromachi, and S. Horiuchi, *Jpn. J. Appl. Phys.* 27, L1241 (1988).
- <sup>7</sup> R. Ramesh, B. van Tendeloo, G. Thomas, S. M. Green, and H. L. Luo, *Appl. Phys. Lett.* 53, 2220 (1988).
- <sup>8</sup> S. Ikeda, K. Aota, T. Hatano, and K. Ogawa, Jpn. J. Appl. Phys. 27, L2040 (1988).
- <sup>9</sup> Y. Ikeda, M. Takano, Z. Hiroi, K. Oda, H. Kitaguchi, J. Takada, Y. Miura, Y. Takeda,
  O. Yamamoto, and H. Mazaki, Jpn. J. Appl. Phys. 27, L2067 (1988).
- <sup>10</sup> B. Kabius, M. E. Gomez, P. Prieto, S. Thierfeldt, and R. R. Arons, *Physica C* 162-164, 635 (1989).
- <sup>11</sup> P. M. de Wolff, T. Janssen, and A. Janner, Acta Cryst. A37, 625 (1981).
- <sup>12</sup> R. Ramesh, G. Thomas, S. Green, C. Jiang, Y. Mei, M. L. Rudee, and H. L. Luo, *Phys. Rev. B* 38, 7070 (1988).
- <sup>13</sup> M. Mizuno, H. Endo, J. Tsuchiya, N. Kijima, A. Sumiyama, and Y Oguri, Jpn. J. Appl. Phys. 27, L1225 (1988).

<sup>14</sup> T. Hatano, K. Aota, S. Ikeda, K. Nakamura, and K. Ogawa, Jpn. J.Appl. Phys. 27, L2055 (1988).

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#### **Figure Captions**

- Fig. 1. Temperature dependence of the electrical resistivity and the magnetic susceptibility of the superconductors in the Sb-Pb-Bi-Sr-Ca-Cu-O system. The  $T_{c, end}$  and  $T_{c, mid}$  of the sample are 110K and 112K, respectively. The critical temperature for the onset of diamagnetism is 111K.
- Fig. 2. Selected area diffraction patterns of the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system. (a)  $a^*-b^*$  plane, (b)  $b^*-c^*$  plane, (c)  $a^*-c^*$  plane, (d). $(a^*+b^*)-c^*$  plane. The superspace group of the high- $T_c$  phase is  $N_{1\bar{1}1}^{Bbmb}$ ,  $N_{1\bar{1}1}^{Bbmb}$ ,  $P_{1\bar{1}1}^{Bbmb}$  or  $P_{111}^{Bb2b}$ .
- Fig. 3. Phase contrast high-resolution transmission electron microscope image (a) and corresponding selected area diffraction pattern (b) of the high- $T_c$  phase in the Sb-Pb-Bi-Sr-Ca-Cu-O system with the incident beam along the *a*-axis. Note that there are two types of satellite spots marked (A) and (B) in the diffraction pattern. The *b*-axis components of wave vectors of modulated structures corresponding to the satellite spots (A) and (B) have magnitudes of 26.9Å and 36.1Å, respectively. The modulated structure with a wavelength of 26.9Å is observed in the image. Note also that only one *c*-plane stacking fault is apparent in the image and is marked (s).

System	$T_c$	Superspace	Subcell lattic			Wavelength of	Stacking
,	(11)	group	pu pu			modulated structure	iduits
			a	<u>b</u>	С	(A)	
Bi-Sr-Ca-Cu-O <sup>a, b</sup>	107	$N^{Bbmb}_{1\overline{1}1}$ or $N^{Bb2b}_{1\overline{1}1}$	5.4	5.4	36.8	27	many
Pb-Bi-Sr-Ca-Cu-O <sup>c</sup>	107	$N^{Bbmb}_{1\overline{1}1}$ or $N^{Bb2b}_{1\overline{1}1}$	5.407(6)	5.407(6)	37.051(7)	27	few
	,	$P^{Bbmb}_{1\bar{1}1}$ or $P^{Bb2b}_{111}$				45 ~ 54	
Sb-Pb-Bi-Sr-Ca-Cu-O	111	$N {B \atop 1 \ \overline{1} \ 1}^{B \atop 1 \ \overline{1} \ 1}$ or $N {B \atop 1 \ \overline{1} \ 1}^{B \atop 1 \ \overline{1} \ 1}$	5.411(1)	5.411(1)	37.22(6)	26.9	few
		$P \stackrel{B \ bm \ b}{1 \ \overline{1} \ 1} $ or $P \stackrel{B \ b2 \ b}{1 \ 1 \ 1}$				36.1	

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# Table 1. Summary of the microstructures of the Bi-based high- $T_c$ phases.

<sup>a</sup>Reference 6.

<sup>b</sup>N. Kijima, H. Endo, J. Tsuchiya, A. Sumiyama, M. Mizuno, and Y. Oguri, Jpn. J. Appl. Phys. 27, L821 (1988).

<sup>c</sup>Reference 5, 7-9.

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Fig. 1.

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XBB 912-919

Fig. 2.



XBB 912-920

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Fig. 3

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Fig. 4.



XBB 912-918

Fig. 5.

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