UC Riverside
UC Riverside Previously Published Works

Title

Spin Seebeck Effect from Antiferromagnetic Magnons and Critical Spin Fluctuations in Epitaxial FeF2 Films

Permalink

https://escholarship.org/uc/item/6gr8901z

Journal

Physical Review Letters, 122(21)

ISSN

0031-9007

Authors

Li, Junxue
Shi, Zhong
Ortiz, Victor H
et al.

Publication Date

2019-05-31

DOI

10.1103/physrevlett.122.217204

Peer reviewed

Spin Seebeck Effect from Antiferromagnetic Magnons and Critical Spin Fluctuations in Epitaxial FeF$_2$ Films

Junxue Li,[1] Zhong Shi,[2] Victor H. Ortiz,[1] Mohammed Aldosary,[1,3] Cliff Chen,[1] Vivek Aji,[1] Peng Wei,[1] and Jing Shi[1]

[1]*Department of Physics and Astronomy, University of California, Riverside, California 92521, USA*
[2]*School of Physics Science and Engineering, Tongji University, Shanghai 200092, China*
[3]*Department of Physics and Astronomy, King Saud University, Riyadh 11451, Saudi Arabia*

We report a longitudinal spin Seebeck effect (SSE) study in epitaxially grown FeF$_2$(110) antiferromagnetic (AFM) thin films with strong uniaxial anisotropy over the temperature range of 3.8–250 K. Both the magnetic-field and temperature-dependent SSE signals below the Néel temperature ($T_N = 70$ K) of the FeF$_2$ films are consistent with a theoretical model based on the excitations of AFM magnons without any net induced static magnetic moment. In addition to the characteristic low-temperature SSE peak associated with the AFM magnons, there is another SSE peak at T_N which extends well into the paramagnetic phase. All the SSE data taken at different magnetic fields up to 12 T near and above the critical point T_N follow the critical scaling law very well with the critical exponents for magnetic susceptibility of 3D Ising systems, which suggests that the AFM spin correlation is responsible for the observed SSE near T_N.

Antiferromagnetic insulators (AFMI) have recently attracted a great deal of interest in the emerging field of antiferromagnetic spintronics due to their unique properties such as robustness against magnetic-field perturbation and ultrafast spin dynamics [1–4]. In early reports, thin AFMI layers are found to enhance spin current transmission when they are inserted between a ferrimagnetic insulator (FMI) and a heavy metal (HM), such as the NiO and CoO AFMI layers in Y$_3$Fe$_5$O$_{12}$ (YIG)/NiO/Pt [5,6] and YIG/CoO/Pt [6], where an increased spin Seebeck effect (SSE) signal is attributed to the enhanced spin conductance in the AFMI spacer around its phase transition temperature [6,7]. In addition, AFMI themselves are reported as sources of pure spin current under an applied magnetic field which generate SSE signals in AFMI-HM heterostructures such as MnF$_2$/Pt [8], Cr$_2$O$_3$/Pt [9], and γ-Fe$_2$O$_3$/Pt [10].

The origin of SSE in AFMI-based heterostructures has been under active investigation. On one hand, the SSE signal in Cr$_2$O$_3$/Pt under a strong magnetic field is found to be proportional to the net equilibrium magnetization of Cr$_2$O$_3$ [9]; i.e., it is negligibly small until the magnetic field exceeds the spin-flop transition field (H_{SF}) and produces a net induced magnetic moment. On the other hand, a finite SSE signal is reported in MnF$_2$/Pt when the magnetic field is less than H_{SF}, indicating that the field-induced magnetic moment is not the only cause of the SSE in AFMI [8]. In ferromagnetic insulator-HM heterostructures, SSE below the Curie temperature is usually attributed to the spin current due to magnon flow driven by a temperature gradient [11–15]. In a uniaxial antiferromagnet, there are two branches of magnons, namely α- and β-mode magnons [Fig. 1(a)], which have opposite chiralities carrying

opposite angular momenta [16–19]. At zero magnetic field, both modes are degenerate; therefore, there is no net magnon flow until a magnetic field is applied along the AFM spin direction [18,20,21], as shown in Fig. 1(b). In a recent theoretical model proposed by Rezende *et al.* [20,21], a field-induced AFM magnon imbalance can lead to a characteristic SSE peak in the absence of any net
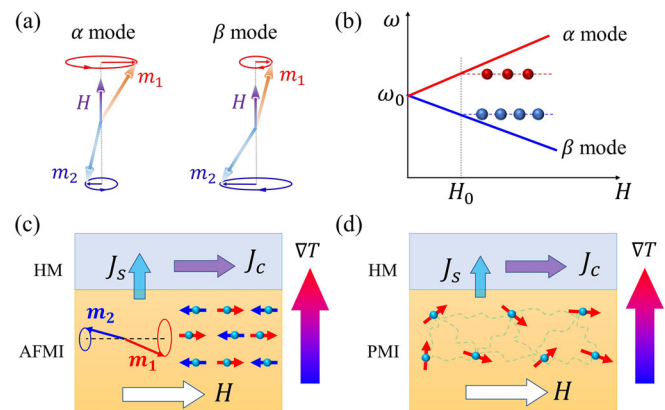

FIG. 1. (a) Illustration of α and β modes of magnons in a uniaxial AFM. m_1 and m_2 are magnetic moments on two spin sublattices, respectively. (b) Frequency ω of AFM magnons for the α and β modes as a function of the magnetic field H applied along the AFM easy axis. The red and blue symbols indicate the occupation of α and β AFM magnons, respectively. Schematic diagrams of SSE measurements for AFMI-HM (c) and PM insulator–HM (d) heterostructures. ∇T is the temperature gradient; J_S is the spin current density across the interface, and J_c is the charge current density generated by the inverse spin Hall effect.

equilibrium magnetization [Fig. 1(c)]. This AFM magnon picture explains the low-field SSE signal in MnF_2/Pt. Clearly, a unified picture is lacking regarding the low-field SSE responses. In order to clarify the physical origin of the SSE in AFMI, a uniaxial AFMI with an unusually high spin-flop field is desired so that no equilibrium magnetization is induced with any laboratory accessible magnetic field.

Besides the FMI and AFMI, paramagnetic insulators (PMIs) have also been reported as a source of pure spin currents. SSE signals were observed in $Gd_3Ga_5O_{12}$ (GGG)/Pt [22], $DyScO_3$ (DSO)/Pt [22], and $CoCr_2O_4$ (CCO)/Pt [23] over the temperature range where GGG, DSO, and CCO are PMIs. In the paramagnetic phase, the concept of magnons is no longer applicable; however, the short-range correlation of spin fluctuations is present [24–27]. Therefore, the observed SSE in PMI-HM heterostructure [as shown in Fig. 1(b)] must have a different origin, i.e., from spin fluctuations. A complete SSE picture in AFMI must contain ingredients of magnons and spin fluctuations in order to fully account for the data in both AFM and paramagnetic phases.

In this Letter, we report an experimental study of SSE in FeF_2/Pt heterostructures. There are two primary reasons why FeF_2 is chosen. First, compared with MnF_2 which has $H_{SF} \sim 9$ T, FeF_2 has stronger uniaxial magnetic anisotropy which gives rise to a larger SF field $H_{SF} \sim 42$ T [28], far greater than the maximum magnetic field used for this study. This ensures negligibly small induced magnetic moment at low temperatures for the magnetic fields applied along the easy axis direction. Second, since the AFM ordering temperature T_N of bulk FeF_2 is 78.4 K [29,30], we can systematically study the SSE response across the antiferromagnetic phase transition. Over the entire temperature range, SSE signals in FeF_2/Pt under different magnetic fields up to 12 T show very similar temperature-dependent behaviors. First, there is a SSE peak at ~ 11.6 K, which can be attributed to the effect of AFM magnons. Second, SSE shows a bump at the T_N of the FeF_2 thin films (70 K) and the finite SSE signal extends to 250 K. We can collapse all SSE data above T_N onto a single curve in a scaling plot by using the critical exponent for magnetic susceptibility, providing direct evidence of SSE probing the correlation of spin fluctuations in PMIs.

High quality (110)-oriented FeF_2 thin films are grown using molecular beam epitaxy and characterized by reflection high energy electron diffraction and x-ray diffraction (see Supplemental Material [31] Sec. I). The (110) orientation is chosen so that the easy axis of the AFM spin, i.e., the [001] or the c-axis direction, lies in the film plane [39–41]. The uniaxial nature of AFM FeF_2 films are characterized by superconducting quantum interference device (SQUID) magnetometry (see Supplemental Material [31] Sec. II) and by performing field-cooling experiments with the applied magnetic field parallel and perpendicular to FeF_2[001]

during cooling and measuring the shifted loops of the anisotropic magnetoresistance in FeF_2/Co bilayers (see Supplemental Material [31] Sec. III). The blocking temperature of the 50 nm–thick FeF_2 thin film is determined to be ~ 70 K from the onset of the exchange bias field and the maximum coercive field in the same field-cooling experiments. The blocking temperature coincides with the temperature where the SSE peak appears in FeF_2/Pt; therefore, we believe the blocking temperature is very close to the Néel temperature T_N of the FeF_2 films. The deviation of T_N of thin film from the bulk value might be caused by the finite size effect.

To form FeF_2/Pt heterostructure for SSE measurements [Fig. 2(a)], 5 nm Pt is directly deposited on top of 50 nm FeF_2 with magnetron sputtering and patterned into a Hall bar with dimensions of $100~\mu m \times 630~\mu m$ perpendicular to the c axis. Then a 150 nm Al_2O_3 insulating layer is deposited by electron-beam evaporation, followed by a 50 nm Cr film covering the Hall bar channel area as a heater. In the SSE experiment, a dc current is applied to the Cr heater to generate a vertical temperature gradient across the interface. An external magnetic field is applied in the thin film plane at an azimuthal angle ϕ with respect to the AFM easy axis FeF_2[001], while the voltage response along
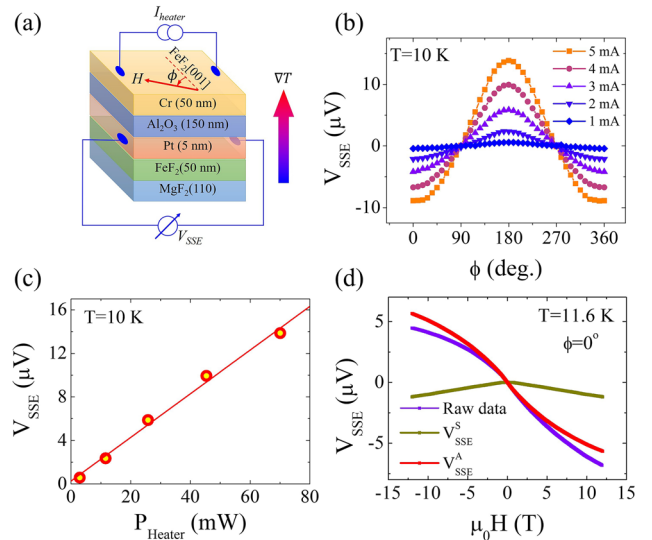

FIG. 2. SSE signal V_{SSE} in FeF_2(50 nm)/Pt(5 nm) heterostructure. (a) Schematic diagram of the sample structure and longitudinal SSE measurement geometry. A current in Cr generates a temperature gradient ∇T which produces V_{SSE} in Pt. H is applied at an angle ϕ to the FeF_2[001] direction in the film plane. The patterned Pt strip is 100 μm wide and 630 μm long and the V_{SSE} is measured along the length of the strip. (b) ϕ dependence of V_{SSE} at 10 K for different heater currents with a constant magnetic field $H = 9$ T. (c) V_{SSE} vs heating power P at 10 K. (d) V_{SSE} as a function of H applied along FeF_2[001] at $T = 11.6$ K with a heater current of 3 mA. Purple, red, and brown curves represent the raw data, antisymmetric, and symmetric components, respectively.

the Pt Hall bar channel is recorded as the spin Seebeck signal V_{SSE}. Figure 2(b) plots the ϕ dependence of V_{SSE} at 10 K with different heater currents under a 9 T rotating magnetic field. V_{SSE} reaches maximum for the magnetic field along the AFM easy axis (i.e., $\phi = 0°$ and $180°$) and vanishes for the magnetic field perpendicular to the easy axis ($\phi = 90°$). The SSE signal shows opposite polarities at $\phi = 0°$ and $180°$, which is consistent with the expectation from the AFM magnon picture. As ϕ varies, the magnetic-field component projected to the FeF_2[001] direction oscillates; so do the Zeeman splitting of the two AFM magnon eigenmodes and the resulting spin current induced SSE signal. In addition, the magnitude of the SSE signal is directly proportional to the heating power [Fig. 2(c)], suggesting the thermoelectric nature of the signal similar to those reported in FMI materials such as YIG [42].

The magnetic-field dependence of the SSE signal at 11.6 K is shown in Fig. 2(d). To eliminate parasitic signals, we decompose the V_{SSE} into the symmetric and antisymmetric components: $V_{SSE} = V_{SSE}^S + V_{SSE}^A$, where $V_{SSE}^A = \frac{1}{2}[V_{SSE}(+H) - V_{SSE}(-H)]$ and $V_{SSE}^S = \frac{1}{2}[V_{SSE}(+H) + V_{SSE}(-H)]$, and plot V_{SSE}^S and V_{SSE}^A in Fig. 2(d). V_{SSE}^S could be due to the normal magneto-Seebeck signal produced by an incidental in-plane temperature gradient along the Hall bar channel. Such a background signal is inevitable and also observed in our control sample of MgF_2/Pt which does not produce any spin current (see Supplemental Material [31], Sec. IV). Only V_{SSE}^A is attributed to the SSE signal in the FeF_2/Pt heterostructure. When the magnetic field is swept along the easy axis of FeF_2, it does not induce any net magnetic moment below <20 K as the SQUID data indicate; however, it changes the energy splitting between the two AFM magnon eigenmodes. Furthermore, we note that the high-field characteristics of SSE in FeF_2/Pt heterostructure are distinctly different from those in YIG/Pt bilayers in the following aspects. First, the SSE voltage in FeF_2/Pt increases with the field and shows no sign of saturation up to 12 T. In contrast, the SSE in YIG/Pt decreases with strong magnetic fields due to suppression of thermal magnon population [43]. This effect is not seen in FeF_2 since the magnon energy is much higher due to the anisotropy gap. Second, the V_{SSE}^A in FeF_2/Pt vanishes at zero magnetic field because the two magnon eigenmodes are degenerate, which leads to zero net spin current; however, the SSE in YIG/Pt remains finite at zero magnetic field due to nonzero population of the sole right-hand magnon mode at finite temperatures [11–15]. We will refer V_{SSE}^A as V_{SSE} hereafter.

The field dependence of the SSE is studied over the temperature range of 3.8–250 K. As the temperature varies, the heater power P also varies due to the heater resistance change, so does the vertical heat current through the sample. To compare the SSE responses at different temperatures, we normalize the SSE voltage signals by the heater
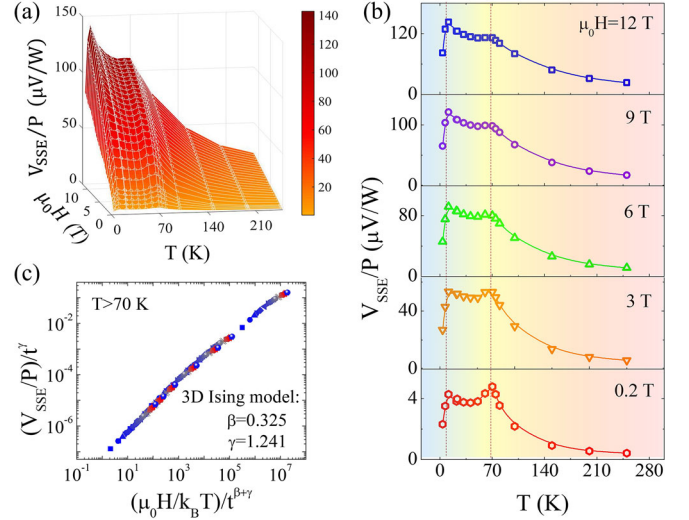

FIG. 3. Temperature dependence of V_{SSE} in FeF_2/Pt hetero-structure. (a) Three-dimensional plot of V_{SSE}/P (only the antisymmetric component) as a function of H and T. (b) Temperature dependence of V_{SSE}/P for $H = 12$, 9, 6, 3, and 0.2 T from top to bottom indicated by different colors. Two vertical dashed lines indicate the peak positions. Solid curves are guides to the eye. The lowest temperature in our experiments is 3.8 K. (c) Scaling plot of all the SSE data above T_N (70 K). Different colors represent data for different magnetic fields as shown in (b). We adopt $\beta = 0.325$ and $\gamma = 1.241$, which are the renormalization group critical exponents for three-dimension Ising systems.

power, which is proportional to the spin current density. We plot V_{SSE}/P against both temperature and magnetic field as displayed in Fig. 3(a). V_{SSE}/P curves under different magnetic fields show very similar characteristic temperature dependence as shown in Fig. 3(b). As the temperature increases from 3.8 K, V_{SSE}/P first increases, reaches a maximum at ~11.6 K, and then starts to decrease. A similar low-temperature peak was also observed in MnF_2/Pt, but at a lower temperature (~5 K at 8 T), which was later interpreted by Rezende *et al.* [20,21] in the AFM magnon model. Here, we attribute the low-temperature FeF_2 peak to the same AFM magnon mechanism. This model qualitatively accounts for the difference in the peak position between the two AFM films. Since the left-hand magnon branch (β mode) lies below the right-hand branch (α mode) under an external field, only the left-hand magnon states are predominately occupied at very low temperatures. As the temperature increases, the left-hand magnon population continues increasing until the right-hand magnon states start to be occupied. Because of the opposite angular momentum of the right-hand magnons, the rising temperature results in a peak in the net spin current and consequently a peak in SSE. In uniaxial AFM materials, the lower the zero-field magnon frequency is, the lower the peak temperature is. In comparison, the zero-field magnon frequencies are 1.6 and 0.27 THz in FeF_2 and MnF_2, respectively [20,21]; therefore, the AFM magnon peak

occurs at a higher temperature in FeF_2 than in MnF_2. Furthermore, a similar SSE signal is observed in FeF_2/Pt for $H \perp [001]$, which has also been reported in the MnF_2/Pt bilayer (see experimental details and discussions in Sec. V of Supplemental Material [31]).

Now we turn to the SSE signal near and above 70 K, i.e., T_N of the FeF_2 thin film. As shown in Fig. 3(b), the SSE signal shows a second peak at $\sim$70 K, which becomes less sharp as the magnetic field increases but its position remains unchanged. The SSE signal decays above 70 K but remains finite up to 250 K. This is reminiscent of the critical behavior of continuous phase transitions. Since there is no long-range AFM order above T_N, the high-temperature SSE signal cannot be interpreted by AFM magnons. Although there is the absence of a spontaneous magnetic moment above T_N, the magnetic field can induce a finite magnetic moment which is proportional to the magnetic susceptibility χ. This can lead to an anomaly in SSE at T_N [44]. Following the critical theory [44], $\chi = t^\gamma f(h/t^{\beta+\gamma})$, where β and γ are the critical exponents for spontaneous staggered magnetization and magnetic susceptibility χ for $H = 0$, respectively, h is $H/k_B T$, $t = (T - T_N)/T_N$ is the reduced temperature for $T > T_N$, and f is a scaling function. Because of the overlap with the tail of the low-temperature AFM magnon peak below T_N, we analyze only the SSE data above T_N. In Fig. 3(c), we show that the heating power-normalized SSE data for all magnetic fields above 70 K in a scaling plot: $(V_{SSE}/P)/t^\gamma$ vs $h/t^{\beta+\gamma}$ using $\beta = 0.325$, $\gamma = 1.241$, the critical exponents from the renormalization group calculations for 3D Ising systems [45]. All data collapse onto a single curve in the scaling plot. The excellent agreement indicates that the SSE signal in the PM region measures the magnetic susceptibility which is governed by critical spin fluctuations near T_N.

Moreover, we find that the scaling analysis does not work for $(V_{SSE}/P)/t^\beta$ vs $h/t^{\beta+\gamma}$ using critical exponent β for the staggered magnetization instead of γ for the magnetic susceptibility to scale the SSE data with all possible combinations of β and γ (see Supplemental Material [31], Sec. VI). Clearly, the SSE signal scales as the magnetic susceptibility χ rather than field-induced sublattice magnetization of the AFMI. Even in a ferro- or ferrimagnet, it is not obvious whether the SSE signal near the Curie temperature scales as the magnetization with critical exponent β. In fact, SSE measurements in YIG were carried out near the Curie temperature T_c and a $(T_c - T)^3$ power-law behavior was found below T_c. This behavior was attributed to the critical behavior of YIG [46]; however, the exponent 3 is not related to the mean-field or critical exponent for the magnetization, nor for the magnetic susceptibility. No satisfactory explanation was given besides acknowledging possible complications due to YIG being a ferrimagnet. Theoretically, the exponent for SSE in ferromagnets was linked to the spontaneous

magnetization [47]. Recently, both SSE [22,23] and spin pumping [48] results in PMI-HM heterostructure were reported, but no quantitative relationship between spin correlation and SSE signal was discussed. Our experimental data and the analyses have established a clear connection between the SSE signal and the magnetic susceptibility near the critical point; therefore, we concluded that SSE is capable of probing correlations of spin fluctuations in magnetic systems.

In summary, we have demonstrated the epitaxial growth of the $FeF_2(110)$ thin films on $MgF_2(110)$ substrate with controlled AFM anisotropy. The exchange bias data confirm that the uniaxial AFM spin axis is along the $FeF_2[001]$ direction and the AFM ordering temperature is 70 K. The temperature dependence of the SSE signal shows two peaks located at 11.6 and 70 K, respectively. The former is attributed to the AFM magnons, and the latter to the enhanced correlation of critical fluctuations near the AFM ordering temperature. Our results suggest that both magnons in magnetically ordered phases and the correlated spin fluctuations near phase transitions can act as pure spin current sources. This picture provides an alternative interpretation of the enhanced SSE and spin pumping signals near the AFM ordering temperature previously observed in FMI-AFM-NM heterostructures [5,6,49].

[1] V. Baltz, A. Manchon, M. Tsoi, T. Moriyama, T. Ono, and Y. Tserkovnyak, Rev. Mod. Phys. **90**, 015005 (2018).

[2] P. Němec, M. Fiebig, T. Kampfrath, and A. V. Kimel, Nat. Phys. **14**, 229 (2018).

[3] M. B. Jungfleisch, W. Zhang, and A. Hoffmann, Phys. Lett. A **382**, 865 (2018).

[4] M. W. Daniels, W. Guo, G. M. Stocks, D. Xiao, and J. Xiao, New J. Phys. **17**, 103039 (2015).

[5] A. Prakash, J. Brangham, F. Yang, and J. P. Heremans, Phys. Rev. B **94**, 014427 (2016).

[6] W. W. Lin, K. Chen, S. F. Zhang, and C.-L. Chien, Phys. Rev. Lett. **116**, 186601 (2016).

[7] K. Chen, W. Lin, C. L. Chien, and S. Zhang, Phys. Rev. B **94**, 054413 (2016).

[8] S. M. Wu, W. Zhang, A. KC, P. Borisov, J. E. Pearson, J. S. Jiang, D. Lederman, A. Hoffmann, and A. Bhattacharya, Phys. Rev. Lett. **116**, 097204 (2016).

[9] S. Seki, T. Ideue, M. Kubota, Y. Kozuka, R. Takagi, M. Nakamura, Y. Kaneko, M. Kawasaki, and Y. Tokura, Phys. Rev. Lett. **115**, 266601 (2015).

[10] P. Jiménez-Cavero, I. Lucas, A. Anadón, R. Ramos, T. Niizeki, M. H. Aguirre, P. A. Algarabel, K. Uchida,

M. R. Ibarra, E. Saitoh, and L. Morellón, APL Mater. **5,** 026103 (2017).

[11] J. Xiao, G. E. W. Bauer, K. C. Uchida, E. Saitoh, and S. Maekawa, Phys. Rev. B **81,** 214418 (2010).

[12] H. Adachi, K. Uchida, E. Saitoh, and S. Maekawa, Rep. Prog. Phys. **76,** 036501 (2013).

[13] S. M. Rezende, R. L. Rodriguez-Suarez, R. O. Cunha, A. R. Rodrigues, F. L. A. Machado, G. A. Fonseca Guerra, J. C. Lopez Ortiz, and A. Azevedo, Phys. Rev. B **89,** 014416 (2014).

[14] A. Kehlberger *et al.*, Phys. Rev. Lett. **115,** 096602 (2015).

[15] K. Uchida, J. Xiao, H. Adachi, J. Ohe, S. Takahashi, J. Ieda, T. Ota, Y. Kajiwara, H. Umezawa, H. Kawai, G. E. W. Bauer, S. Maekawa, and E. Saitoh, Nat. Mater. **9,** 894 (2010).

[16] C. Kittel, Phys. Rev. **82,** 565 (1951).

[17] F. Keffer and C. Kittel, Phys. Rev. **85,** 329 (1952).

[18] R. Cheng, J. Xiao, Q. Niu, and A. Brataas, Phys. Rev. Lett. **113,** 057601 (2014).

[19] Ø. Johansen and A. Brataas, Phys. Rev. B **95,** 220408 (2017).

[20] S. M. Rezende, R. L. Rodríguez-Suárez, and A. Azevedo, Phys. Rev. B **93,** 014425 (2016).

[21] S. M. Rezende, A. Azevedo, and R. L. Rodríguez-Suárez, J. Phys. D **51,** 174004 (2018).

[22] S. M. Wu, J. E. Pearson, and A. Bhattacharya, Phys. Rev. Lett. **114,** 186602 (2015).

[23] A. Aqeel, N. Vlietstra, J. A. Heuver, G. E. W. Bauer, B. Noheda, B. J. van Wees, and T. T. M. Palstra, Phys. Rev. B **92,** 224410 (2015).

[24] C. J. Liu, S. M. Wu, J. E. Pearson, J. S. Jiang, N. d'Ambrumenil, and A. Bhattacharya, Phys. Rev. B **98,** 060415(R) (2018).

[25] T. Chatterji, G. J. McIntyre, and P.-A. Lindgard, Phys. Rev. B **79,** 172403 (2009).

[26] Z. Qiu, J. Li, D. Hou, E. Arenholz, A. T. N'Diaye, A. Tan, K.-i. Uchida, K. Sato, S. Okamoto, Y. Tserkovnyak, Z. Q. Qiu, and E. Saitoh, Nat. Commun. **7,** 12670 (2016).

[27] L. Frangou, S. Oyarzún, S. Auffret, L. Vila, S. Gambarelli, and V. Baltz, Phys. Rev. Lett. **116,** 077203 (2016).

[28] V. Jaccarino, A. R. King, M. Motokawa, T. Sakakibara, and M. Date, J. Magn. Magn. Mater. **31–34,** 1117 (1983).

[29] D. P. Belanger, P. Nordblad, A. R. King, and V. Jaccarino, J. Magn. Magn. Mater. **31–34,** 1095 (1983).

[30] J. Nogués, D. Lederman, T. J. Moran, and I. K. Schuller, Phys. Rev. Lett. **76,** 4624 (1996).

[31] See Supplemental Material at http://link.aps.org/supplemental/10.1103/PhysRevLett.122.217204 for details, which includes Refs. [32–38].

[32] R. L. Stamps, J. Phys. D **33,** R247 (2000).

[33] J. Wu, J. S. Park, W. Kim, E. Arenholz, M. Liberati, A. Scholl, Y. Z. Wu, C. Hwang, and Z. Q. Qiu, Phys. Rev. Lett. **104,** 217204 (2010).

[34] T. Ambrose and C. L. Chien, Phys. Rev. Lett. **76,** 1743 (1996).

[35] R. H. Kodama, S. . Makhlouf, and A. E. Berkowitz, Phys. Rev. Lett. **79,** 1393 (1997).

[36] K. Behnia and H. Aubin, Rep. Prog. Phys. **79,** 046502 (2016).

[37] D. Qu, S. Y. Huang, J. Hu, R. Q. Wu, and C. L. Chien, Phys. Rev. Lett. **110,** 067206 (2013).

[38] D. . Cahill, Rev. Sci. Instrum. **61,** 802 (1990).

[39] M. Hutchings, B. Rainford, and H. Guggenheim, J. Phys. C **3,** 307 (1970).

[40] R. Morales, Z. P. Li, O. Petracic, X. Batlle, I. K. Schuller, J. Olamit, and K. Liu, Appl. Phys. Lett. **89,** 072504 (2006).

[41] R. Morales, A. C. Basaran, J. E. Villegas, D. Navas, N. Soriano, B. Mora, C. Redondo, X. Batlle, and I. K. Schuller, Phys. Rev. Lett. **114,** 097202 (2015).

[42] K. Uchida, H. Adachi, T. Ota, H. Nakayama, S. Maekawa, and E. Saitoh, Appl. Phys. Lett. **97,** 172505 (2010).

[43] T. Kikkawa, K.-i. Uchida, S. Daimon, Z. Qiu, Y. Shiomi, and E. Saitoh, Phys. Rev. B **92,** 064413 (2015).

[44] N. Goldenfeld, *Lectures on Phase Transitions and the Renormalization Group*, 1st ed. (CRC Press, Boca Raton, FL, 2018).

[45] S. N. Kaul, J. Magn. Magn. Mater. **53,** 5 (1985).

[46] K. I. Uchida, T. Kikkawa, A. Miura, J. Shiomi, and E. Saitoh, Phys. Rev. X **4,** 041023 (2014).

[47] H. Adachi, Y. Yamamoto, and M. Ichioka, J. Phys. D **51,** 144001 (2018).

[48] Y. Shiomi and E. Saitoh, Phys. Rev. Lett. **113,** 266602 (2014).

[49] H. Wang, C. Du, P. C. Hammel, and F. Yang, Phys. Rev. Lett. **113,** 097202 (2014).