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# **Authors**

Chiang, M.H. Mahan, B.H. Tsao, C.W. <u>et al.</u>

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M. H. Chiang, B. H. Mahan, and C. W. Tsao,

and A. 🕱. Werner

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#### by

M. H. Chiang, B. H. Mahan, C. W. Tsao, and A. S. Werner Inorganic Materials Research Division of the Lawrence Radiation Laboratory and Department of Chemistry, University of California, Berkeley, California.

The effect of isotopic substitution on reaction rate constants has been a valuable aid in deducing both reaction mechanisms and the properties of transition states.<sup>1</sup> Recently, very large hydrogen-deuterium isotope effects have been found in ion beam studies of exothermic hydrogen atom transfer reactions carried out at relative energies of collision which were rather high on the usual chemical energy scale.<sup>2-4</sup> The most detailed of these investigations<sup>3</sup> showed that in the  $N_2^+$ -HD reaction, abstraction of hydrogen to form  $N_2H^+$  is as much as twenty times more probable as deuterium atom abstraction, in grazing collisions which lead to the prevalent forward scattering of the ion product. While the general shapes of the angular distributions of  $N_2H^+$  and  $N_2D^+$  are similar, the nearly head-on collisions which lead to large angle scattering preferentially produce  $N_2D^+$ .

We have also demonstrated<sup>5</sup> that a rather different isotope effect occurs in the  $O_2^+$ -HD reaction at low (< 5 eV) relative collision energies. In this instance, reaction occurs through

formation of a long-lived  $HDO_2^+$  complex which decays to give isotropic distributions of both  $HO_2^+$  and  $DO_2^+$ . The total intensity of the  $DO_2^+$  product greatly exceeds (~8×) that of the  $HO_2^+$ , just as one would expect from zero point energy and density of states considerations.

In this note we report the occurrence of a novel isotope effect in  $0_2^+$ -HD collisions at high (> 6 eV) relative energies. Figure 1 shows intensity contour maps of the distribution of H0<sup>+</sup><sub>2</sub>, D0<sup>+</sup><sub>2</sub>, and 0<sup>+</sup><sub>2</sub> from 0<sup>+</sup><sub>2</sub>-HD collisions at 8.59 eV relative energy. Both the  $HO_2^+$  and  $DO_2^+$  distributions are asymmetric about ±90° in the barycentric system, which is consistent with our observation<sup>5</sup> that these reactions proceed by a direct interaction mechanism at relative energies above 5 eV. The angular distributions of  $HO_2^+$  and  $DO_2^+$  are of totally different form, with the HO2 almost exclusively in the small angle, grazing collision region, and the  $DO_2^+$  predominantly in the large angle, rebound collision region. We have found that the  $OH^+$  and  $OD^+$  products also display asymmetric distributions of a much less extreme type, with OH<sup>+</sup> falling principally at angles smaller than 90°, and 0D<sup>+</sup> predominating at angles greater than 90°.

There are two simple explanations which may account for these angular dependent isotope effects. On the basis of geometrical factors alone, the  $HOOD^+$  collision complex will tend to be formed with  $O_2^+$  and HD roughly parallel to each other, and approximately perpendicular to the direction of flight of  $O_2^+$ .

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The center of mass of the resulting  $HOOD^+$  complex is closer to the D atom than the H atom, while the center of force is midway between these atoms. Therefore, the motion of the  $0_2^+$ projectile will, on the average, tend to cause the complex to start rotating with the OH end moving in the flight direction of the  $0_2^+$  projectile, and the OD end moving in the opposite direction. If the complex decomposes in less than one rotation, as it does in this high energy regime, any OH<sup>+</sup> or HO<sub>2</sub><sup>+</sup> would tend to be scattered into the forward or small angle region, while deuterated products would appear at large angles.

A second explanation for the  $HO_2^+-DO_2^+$  isotope effect involves the fact that  $HO_2^+$  formed by the spectator stripping process from HD at 8.6 eV relative energy is stable with respect to decomposition to H and  $O_2^+$ . In contrast,  $DO_2^+$  formed by stripping from HD at this relative energy has sufficient internal excitation to decompose to D and  $O_2^+$ . Thus,  $DO_2^+$  should appear only as a product of rebound collisions in which some of the internal energy of the molecule-ion can be dissipated as relative translation of H and  $DO_2^+$ . Some support for this picture is found in the nonreactive scattering of  $O_2^+$  from HD shown in Fig. 1C. The inelastically scattered  $O_2^+$  has an intensity peak close to the velocity which  $DO_2^+$  form by spectator stripping would have. This feature of the  $O_2^+$  distribution may, therefore, arise from the dissociation of the unstable forward scattered  $DO_2^+$ .

We feel that product stability may be the principal factor causing the very different angular distributions of  $HO_2^+$  and  $DO_2^+$  at these higher relative energies. In the formation of  $OH^+$  and

 $OD^+$ , the problem of product stability does not arise, and the smaller differences in the distributions of  $OH^+$  and  $OD^+$  may be caused by preferential tumbling of the short-lived collision complex to give forward scattered  $OH^+$  and back scattered  $OD^+$ .

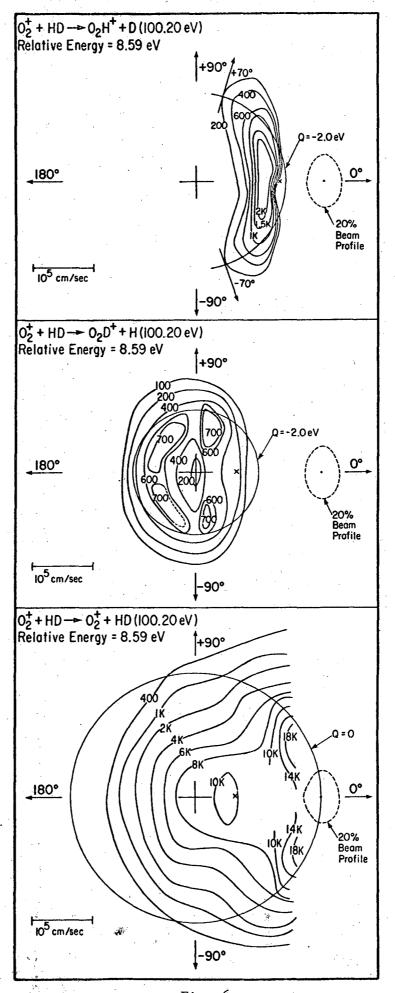
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### Figure Caption

Fig. 1. (a) Contour map of the specific intensity of  $HO_2^+$ from the  $O_2^+$ -HD reaction at 8.59 eV relative energy. The small cross marks the velocity of  $HO_2^+$  formed by the ideal stripping process, and the circle labelled Q = -2.0 eV is the approximate locus of product with no internal excitation. (b) The specific intensity of  $DO_2^+$ . Note the low intensity in the vicinity of the spectator stripping velocity of  $DO_2^+$ . (c) The specific intensity of  $O_2^+$  scattered from HD. The intensity maximum in the inelastic region falls very near to the velocity of  $DO_2^+$  formed by the spectator stripping process, which is marked by a cross.



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Fig. 6.

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