Title
Angular analysis of $B_d \rightarrow K^{*}\mu^{+}\mu^{-}$ decays in pp collisions at $s = 8 \sqrt{s} \text{ TeV}$ with the ATLAS detector

Permalink
https://escholarship.org/uc/item/7d5252zx

Journal
Journal of High Energy Physics, 2018(10)

ISSN
1029-8479

Authors
Aaboud, M
Aad, G
Abbott, B
et al.

Publication Date
2018-10-01

DOI
10.1007/jhep10(2018)047

Peer reviewed
Angular analysis of $B^0_d \rightarrow K^* \mu^+ \mu^-$ decays in $pp$ collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector

The ATLAS Collaboration

An angular analysis of the decay $B^0_d \rightarrow K^* \mu^+ \mu^-$ is presented, based on proton–proton collision data recorded by the ATLAS experiment at the LHC. The study is using 20.3 fb$^{-1}$ of integrated luminosity collected during 2012 at centre-of-mass energy of $\sqrt{s} = 8$ TeV. Measurements of the $K^*$ longitudinal polarisation fraction and a set of angular parameters obtained for this decay are presented. The results are compatible with the Standard Model predictions.
1 Introduction

Flavour-changing neutral currents (FCNC) have played a significant role in the construction of the Standard Model of particle physics (SM). These processes are forbidden at tree level and can proceed only via loops, hence are rare. An important set of FCNC processes involve the transition of a $b$-quark to an $s\mu^+\mu^-$ final state mediated by electroweak box and penguin diagrams. If heavy new particles exist, they may contribute to FCNC decay amplitudes, affecting the measurement of observables related to the decay under study. Hence FCNC processes allow searches for contributions from sources of physics beyond the SM (hereafter referred to as new physics). This analysis focuses on the decay $B^0_d \to K^{*}(892)\mu^+\mu^-$, where $K^{*}(892)$ is referred to as $K^*$ and charge conjugation is implied throughout, unless stated otherwise. In addition to angular observables such as the forward-backward asymmetry $A_{FB}$, there is considerable interest in measurements of the charge asymmetry, differential branching fraction, isospin asymmetry, and ratio of rates of decay into dimuon and dielectron final states, all as a function of the invariant mass squared of the dilepton system $q^2$. All of these observable sets can be sensitive to different types of new physics that allow for FCNCs at tree or loop level. The BaBar, Belle, CDF, CMS, and LHCb collaborations have published the results of studies of the angular distributions for $B^0_d \to K^*\mu^+\mu^-$ [1–8]. The LHCb Collaboration has reported a potential hint, at the level of 3.4 standard deviations, of a deviation from SM calculations [3, 4] in this decay mode when using a parameterization of the angular distribution designed to minimise uncertainties from hadronic form factors. Measurements using this approach were also reported by the Belle and CMS Collaborations [6, 8] and they are consistent with the LHCb experiment’s results and with the SM calculations. This paper presents results following the methodology outlined in Ref. [3] and the convention adopted by the LHCb Collaboration for the definition of angular observables described in Ref. [9]. The results obtained here are compared with theoretical predictions that use the form factors computed in Ref. [10].

This article presents the results of an angular analysis of the decay $B^0_d \to K^*\mu^+\mu^-$ with the ATLAS detector, using 20.3 fb$^{-1}$ of $pp$ collision data at a centre-of-mass energy $\sqrt{s} = 8$ TeV delivered by the Large Hadron Collider (LHC) [11] during 2012. Results are presented in six different bins of $q^2$ in the range 0.04 to 6.0 GeV$^2$, where three of these bins overlap. Backgrounds, including a radiative tail from $B^0_d \to K^* J/\psi$ events, increase for $q^2$ above 6.0 GeV$^2$, and for this reason, data above this value are not studied.

The operator product expansion used to describe the decay $B^0_d \to K^*\mu^+\mu^-$ encodes short-distance contributions in terms of Wilson coefficients and long-distance contributions in terms of operators [12]. Global fits for Wilson coefficients have been performed using measurements of $B^0_d \to K^*\mu^+\mu^-$ and other rare processes. Such studies aim to connect deviations from the SM predictions in several processes to identify a consistent pattern hinting at the structure of a potential underlying new-physics Lagrangian, see Refs. [13–15]. The parameters presented in this article can be used as inputs to these global fits.

2 Analysis method

Three angular variables describing the decay are defined according to convention described by the LHCb Collaboration in Ref. [9]: the angle between the $K^+$ and the direction opposite to the $B^0_d$ in the $K^*$ 

\[ \text{The forward-backward asymmetry is given by the normalised difference between the number of positive muons going in the forward and backward directions with respect to the direction opposite to } B^0_d \text{ momentum in the dimuon rest frame.} \]
centre-of-mass frame ($\theta_K$); the angle between the $\mu^+$ and the direction opposite to the $B^0_d$ in the dimuon centre-of-mass frame ($\theta_L$); and the angle between the two decay planes formed by the $K\pi$ and the dimuon systems in the $B^0_d$ rest frame ($\phi$). For $B^0_d$ mesons the definitions are given with respect to the negatively charged particles. Figure 1 illustrates the angles used.

![Diagram of the $B^0_d \rightarrow K^* \mu^+ \mu^-$ decay showing the angles $\theta_K$, $\theta_L$ and $\phi$ defined in the text. Angles are computed in the rest frame of the $K^*$, dimuon system and $B^0_d$ meson, respectively.](image.png)

The angular differential decay rate for $B^0_d \rightarrow K^* \mu^+ \mu^-$ is a function of $q^2$, $\cos \theta_K$, $\cos \theta_L$ and $\phi$, and can be written in several ways [16]. The form to express the differential decay amplitude as a function of the angular parameters uses coefficients that may be represented by the helicity or transversity amplitudes [17] and is written as

$$
\frac{1}{\Gamma} \frac{d^4 \Gamma}{d q^2 d \cos \theta_L d \cos \theta_K d \phi} = \frac{9}{32 \pi} \left[ \frac{3(1 - F_L)}{4} \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1 - F_L}{4} \sin^2 \theta_K \cos 2\theta_L - F_L \cos^2 \theta_K \cos 2\theta_L + S_3 \sin^2 \theta_K \sin^2 \theta_L \cos 2\phi + S_4 \sin 2\theta_K \sin 2\theta_L \cos \phi + S_5 \sin 2\theta_K \sin \theta_L \cos \phi + S_6 \sin^2 \theta_K \cos \theta_L + S_7 \sin 2\theta_K \sin \theta_L \sin \phi + S_8 \sin 2\theta_K \sin 2\theta_L \sin \phi + S_9 \sin^2 \theta_K \sin^2 \theta_L \sin 2\phi \right].
$$

(1)

Here $F_L$ is the fraction of longitudinally polarised $K^*$ mesons and the $S_i$ are angular coefficients. These angular parameters are functions of the real and imaginary parts of the transversity amplitudes of $B^0_d$ decays into $K^* \mu^+ \mu^-$. The forward-backward asymmetry is given by $A_{FB} = 3S_6/4$. The predictions for the $S$ parameters depend on hadronic form factors which have significant uncertainties at leading order. It is possible to reduce the theoretical uncertainty in these predictions by transforming the $S_i$ using ratios constructed to cancel form factor uncertainties at leading order. These ratios are given by Refs. [17, 18]

---

2 This equation neglects possible $K\pi$ S-wave contributions. The effect of an S-wave contribution is considered following the method used by LHCb in Ref. [3].
as

\[ P_1 = \frac{2S_3}{1 - F_L} \]  \hspace{1cm} (2)

\[ P_2 = \frac{2A_{FB}}{3(1 - F_L)} \]  \hspace{1cm} (3)

\[ P_3 = -\frac{S_0}{1 - F_L} \]  \hspace{1cm} (4)

\[ P'_{j=4,5,6,8} = \frac{S_{j=4,5,7,8}}{\sqrt{F_L(1 - F_L)}}. \]  \hspace{1cm} (5)

All of the parameters introduced, \( F_L \), \( S_i \) and \( P'_{j} \), may vary with \( q^2 \) and the data are analysed in \( q^2 \) bins to obtain an average value for a given parameter in that bin.

### 3 The ATLAS detector, data, and Monte Carlo samples

The ATLAS experiment at the LHC is a general-purpose detector with a cylindrical geometry and nearly \( 4\pi \) coverage in solid angle [19]. It consists of an inner detector (ID) for tracking, a calorimeter system and a muon spectrometer (MS). The ID consists of silicon pixel and strip detectors, with a straw-tube transition radiation tracker providing additional information for tracks passing through the central region of the detector.\(^3\) The ID has a coverage of \(|\eta| < 2.5\), and is immersed in a 2T axial magnetic field generated by a superconducting solenoid. The calorimeter system, consisting of liquid argon and scintillator-tile sampling calorimeter subsystems, surrounds the ID. The outermost part of the detector is the MS, which employs several detector technologies in order to provide muon identification and a muon trigger. A toroidal magnet system is embedded in the MS. The ID, calorimeter system and MS have full azimuthal coverage.

The data analysed here were recorded in 2012 during Run 1 of the LHC. The centre-of-mass energy of the \( pp \) system was \( \sqrt{s} = 8 \) TeV. After applying data-quality criteria, the data sample analysed corresponds to an integrated luminosity of 20.3 fb\(^{-1}\). A number of Monte Carlo (MC) signal and background event samples were generated, with \( b \)-hadron production in \( pp \) collisions simulated with Pythia 8.186 [20, 21]. The AU2 set of tuned parameters [22] is used together with the CTEQ6L1 PDF set [23]. The EvtGen 1.2.0 program [24] is used for the properties of \( b \)- and \( c \)-hadron decays. The simulation included modelling of multiple interactions per \( pp \) bunch crossing in the LHC with Pythia soft QCD processes. The simulated events were then passed through the full ATLAS detector simulation program based on Geant 4 [25, 26] and reconstructed in the same way as data. The samples of MC generated events are described further in Section 5.

\(^3\) ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates \((r, \Phi)\) are used in the transverse plane, \( \Phi \) being the azimuthal angle around the z-axis. The pseudorapidity is defined in terms of the polar angle \( \theta \) as \( \eta = -\ln \tan(\theta/2) \).
4 Event selection

Several trigger signatures constructed from the MS and ID inputs are selected based on availability during the data-taking period, prescale factor and efficiency for signal identification. Data are combined from 19 trigger chains where 21%, 89% or 5% of selected events pass one or more triggers with one, two, or at least three muons identified online in the MS, respectively. Of the events passing the requirement of at least two muons, the largest contribution comes from the chain requiring one muon with a transverse momentum $p_T > 4$ GeV and the other muon with $p_T > 6$ GeV. This combination of triggers ensures that the analysis remains sensitive to events down to the kinematic threshold of $q^2 = 4m^2_\mu$, where $m_\mu$ is the muon mass. The effective average trigger efficiency for selected signal events is about 29%, determined from signal MC simulation.

Muon track candidates are formed offline by combining information from both the ID and MS [27]. Tracks from signal MC simulation. Tracks from signal MC simulation.

The effective average trigger efficiency for selected signal events is about 29%, determined from signal MC simulation.

Muon track candidates are formed offline by combining information from both the ID and MS [27]. Tracks are required to satisfy $|\eta| < 2.5$. Candidate muon (kaon and pion) tracks in the ID are required to satisfy $p_T > 3.5$ (0.5) GeV. Pairs of oppositely charged muons are required to originate from a common vertex with a fit quality $\chi^2$/NDF < 10.

Candidate $K^*$ mesons are formed using pairs of oppositely charged kaon and pion candidates reconstructed from hits in the ID. Candidates are required to satisfy $p_T(K^*) > 3.0$ GeV. As the ATLAS detector does not have a dedicated charged-particle identification system, candidates are reconstructed with both possible $K\pi$ mass hypotheses. The selection implicitly relies on the kinematics of the reconstructed $K^*$ meson to determine which of the two tracks corresponds to the kaon. If both candidates in an event satisfy selection criteria, they are retained and one of them is selected in the next step following a procedure described below. The $K\pi$ invariant mass is required to lie in a window of twice the natural width around the nominal mass of 896 MeV, i.e. in the range [846, 946] MeV. The charge of the kaon candidate is used to assign the flavour of the reconstructed $B^0_d$ candidate.

The $B^0_d$ candidates are reconstructed from a $K^*$ candidate and a pair of oppositely charged muons. The four-track vertex is fitted and required to satisfy $\chi^2$/NDF < 2 to suppress background. A significant amount of combinatorial, $B^0_d$, $B^+$, $B^0_d$ and $\Lambda_b$ background contamination remains at this stage. Combinatorial background is suppressed by requiring a $B^0_d$ candidate lifetime significance $\tau/\sigma_\tau > 12.5$, where the decay time uncertainty $\sigma_\tau$ is calculated from the covariance matrices associated with the four-track vertex fit and with the primary vertex fit. Background from final states partially reconstructed as $B \to \mu^+\mu^-X$ accumulates at invariant mass below the $B^0_d$ mass and contributes to the signal region. It is suppressed by imposing an asymmetric mass cut around the nominal $B^0_d$ mass, $5150$ MeV < $m_{B^0_d\mu \mu} < 5700$ MeV. The high-mass sideband is retained, as the parameter values for the combinatorial background shapes are extracted from the fit to data described in Section 5. To further suppress background, it is required that the angle $\Theta$, defined between the vector from the primary vertex to the $B^0_d$ candidate decay vertex and the $B^0_d$ candidate momentum, satisfies $\cos \Theta > 0.999$. Resolution effects on $\cos \theta_K$, $\cos \theta_L$ and $\phi$ were found to have a negligible effect on the ATLAS $B^0_d \to J/\psi \phi$ analysis [28]. It is assumed to also be the case for $B^0_d \to K^+\mu^+\mu^-$. On average 12% of selected events in the data have more than one reconstructed $B^0_d$ candidate. The fraction is 17% for signal MC samples and 2–10% for exclusive background MC samples. A two-step selection process is used for such events. For 4% of these events it is possible to select a candidate with the smallest value of the $B^0_d$ vertex $\chi^2$/NDF. However, the majority, about 96%, of multiple candidates arise from four-track combinations where the kaon and pion assignments are ambiguous. As these candidates have degenerate values for the $B^0_d$ candidate vertex $\chi^2$/NDF, a second selection step is required. The $B^0_d$
candidate reconstructed with the smallest value of $|m_{K^*} - m_{K^0}|/\sigma(m_{K^0})$ is retained for analysis, where $m_{K^*}$ is the $K^*$ candidate mass, $\sigma(m_{K^0})$ is the per-event uncertainty in this quantity, and $m_{K^0}$ is the world average value of the $K^*$ mass.

The selection procedure results in an incorrect flavour tag (mistag) for some signal events. The mistag probability of a $B_d^0 (\bar{B}_d^0)$ meson is denoted by $\omega(\bar{\omega})$ and is determined from MC simulated events to be $0.1088 \pm 0.0005 (0.1086 \pm 0.0005)$. The mistag probability varies slightly with $q^2$ such that the difference $\omega - \bar{\omega}$ remains consistent with zero. Hence the average mistag rate $\langle \omega \rangle$ in a given $q^2$ bin is used to account for this effect. If a candidate is mistagged, the values of $\cos \theta_L$, $\cos \theta_K$ and $\phi$ change sign, while the latter two are also slightly shaped by the swapped hadron track mass hypothesis. Sign changes in these angles affect the overall sign of the terms multiplied by the coefficients for this effect. If a candidate is mistagged, the values of $\omega$ in a given $q^2$ bin is used to account for the effect. If a candidate is mistagged, the values of $\cos \theta_L$, $\cos \theta_K$ and $\phi$ change sign, while the latter two are also slightly shaped by the swapped hadron track mass hypothesis. Sign changes in these angles affect the overall sign of the terms multiplied by the coefficients $S_0$, $S_0$, $S_0$ and $S_0$ (similarly for the corresponding $P^{(r)}$ parameters) in Equation (1). The corollary is that mistagged events result in a dilution factor of $(1 - 2\langle \omega \rangle)$ for the affected coefficients.

The region $q^2 \in [0.98, 1.1] \text{ GeV}^2$ is vetoed to remove any potential contamination from the $\phi(1020)$ resonance. The remaining data with $q^2 \in [0.04, 6.0] \text{ GeV}^2$ are analysed in order to extract the signal parameters of interest. Two $K^*c\bar{c}$ control regions are defined for $B_d^0$ decays into $K^*J/\psi$ and $K^*\psi(2S)$, respectively as $q^2 \in [8, 11]$ and $[12, 15] \text{ GeV}^2$. The control samples are used to extract values for nuisance parameters describing the signal probability density function (pdf) from data as discussed in Section 5.3.

For $q^2 < 6 \text{ GeV}^2$ the selected data sample consists of 787 events and is composed of signal $B_d^0 \rightarrow K^*\mu^+\mu^-$ decay events as well as background that is dominated by a combinatorial component that does not peak in $m_{K\pi\mu\mu}$ and does not exhibit a resonant structure in $q^2$. Other background contributions are considered when estimating systematic uncertainties. Above 6 GeV$^2$ the background contribution increases significantly, including events coming from $B_d^0 \rightarrow K^*J/\psi$ with a radiative $J/\psi \rightarrow \mu^+\mu^-\gamma$ decay. Scalar $K\pi$ contributions are neglected in the nominal fit and considered only when addressing systematic uncertainties. The data are analysed in the $q^2$ bins $[0.04, 2.0]$, $[2.0, 4.0]$ and $[4.0, 6.0] \text{ GeV}^2$, where the bin width is chosen to provide a sample of signal events sufficient to perform an angular analysis. The width is much larger than the $q^2$ resolution obtained from MC simulated signal events and observed in data for $B_d^0$ decays into $K^*J/\psi$ and $K^*\psi(2S)$. Additional overlapping bins $[0.04, 4.0]$, $[1.1, 6.0]$ and $[0.04, 6.0] \text{ GeV}^2$ are analysed in order to facilitate comparison with results of other experiments and with theoretical predictions.

5 Maximum-likelihood fit

Extended unbinned maximum-likelihood fits of the angular distributions of the signal decay are performed on the data for each $q^2$ bin. The discriminating variables used in the fit are $m_{K\pi\mu\mu}$, the cosines of the helicity angles ($\cos \theta_K$ and $\cos \theta_L$), and $\phi$. The likelihood $L$ for a given $q^2$ bin is

$$L = \frac{e^{-n}}{N!} \prod_{k=1}^{N} \sum_{l} n_l P_{kl}(m_{K\pi\mu\mu}, \cos \theta_K, \cos \theta_L, \phi; \hat{p}, \hat{\theta}),$$

where $N$ is the total number of events, the sum runs over signal and background components, $n_l$ is the fitted yield for the $l$th component, $n$ is the sum over $n_l$, and $P_{kl}$ is the pdf evaluated for event $k$ and component $l$. In the nominal fit, $l$ iterates only over one signal and one background component. The $\hat{p}$ are parameters of interest ($F_L$, $S_i$) and $\hat{\theta}$ are nuisance parameters. The remainder of this section discusses the signal model.
(Section 5.1), treatment of background (Section 5.2), use of \( K^+ c\bar{c} \) decay control samples (Section 5.3), fitting procedure and validation (Section 5.4).

5.1 Signal model

The signal mass distribution is modelled by a Gaussian distribution with the width given by the per-event uncertainty in the \( K\pi\mu\mu \) mass, \( \sigma(m_{K\pi\mu\mu}) \), as estimated from the track fit, multiplied by a unit-less scale factor \( \xi \), i.e. the width given by \( \xi \cdot \sigma(m_{K\pi\mu\mu}) \). The mean values of the \( B^0_d \) candidate mass \( m_0 \) and \( \xi \) of the signal Gaussian pdf are determined from fits to data in the control regions as described in Section 5.3.

The simultaneous extraction of all coefficients using the full angular distribution of Equation (1) requires a certain minimum signal yield and signal purity to avoid a pathological fit behaviour. A significant fraction of fits to ensembles of simulated pseudo-experiments do not converge using the full distribution. This is mitigated using trigonometric transformations to fold certain angular distributions and thereby simplify Equation (1) such that only three parameters are extracted in one fit: \( F_L, S_3 \) and one of the other \( S \) parameters. For these folding schemes the angular parameters of interest, denoted by \( \vec{p} \) in Equation (6), are \((F_L, S_3, S_i)\) where \( i = 4, 5, 7, 8 \). These translate into \((F_L, P_1, P'_j)\), where \( j = 4, 5, 6, 8 \), using Equation (5).

Following Ref. [3], the transformations listed below are used:

\[
\begin{align*}
F_L, S_3, S_4, P'_4 & : \\
\begin{cases}
\phi \rightarrow -\phi & \text{for } \phi < 0 \\
\phi \rightarrow \pi - \phi & \text{for } \theta_L > \frac{\pi}{2} \\
\theta_L \rightarrow \pi - \theta_L & \text{for } \theta_L > \frac{\pi}{2},
\end{cases} \\
(7) \\
F_L, S_3, S_5, P'_5 & : \\
\begin{cases}
\phi \rightarrow -\phi & \text{for } \phi < 0 \\
\phi \rightarrow \pi - \phi & \text{for } \theta_L > \frac{\pi}{2} \\
\theta_L \rightarrow \pi - \theta_L & \text{for } \theta_L > \frac{\pi}{2},
\end{cases} \\
(8) \\
F_L, S_3, S_7, P'_6 & : \\
\begin{cases}
\phi \rightarrow -\phi & \text{for } \phi > \frac{\pi}{2} \\
\phi \rightarrow -\pi - \phi & \text{for } \phi < -\frac{\pi}{2} \\
\theta_L \rightarrow \pi - \theta_L & \text{for } \theta_L > \frac{\pi}{2},
\end{cases} \\
(9) \\
F_L, S_3, S_8, P'_8 & : \\
\begin{cases}
\phi \rightarrow -\phi & \text{for } \phi > \frac{\pi}{2} \\
\phi \rightarrow -\pi - \phi & \text{for } \phi < -\frac{\pi}{2} \\
\theta_L \rightarrow \pi - \theta_L & \text{for } \theta_L > \frac{\pi}{2} \\
\theta_K \rightarrow \pi - \theta_K & \text{for } \theta_L > \frac{\pi}{2}.
\end{cases} \\
(10)
\end{align*}
\]

On applying transformation (7), (8), (9), and (10), the angular variable ranges become

\[
\begin{align*}
cos \theta_L \in [0, 1], & \quad cos \theta_K \in [-1, 1] \quad \text{and} \quad \phi \in [0, \pi], \\
cos \theta_L \in [0, 1], & \quad cos \theta_K \in [-1, 1] \quad \text{and} \quad \phi \in [0, \pi], \\
cos \theta_L \in [0, 1], & \quad cos \theta_K \in [-1, 1] \quad \text{and} \quad \phi \in [-\pi/2, \pi/2], \\
cos \theta_L \in [0, 1], & \quad cos \theta_K \in [-1, 1] \quad \text{and} \quad \phi \in [-\pi/2, \pi/2],
\end{align*}
\]
respectively. A consequence of using the folding schemes is that \( S_0 (A_{FB}) \) and \( S_0 \) cannot be extracted from the data. The values and uncertainties of \( F_L \) and \( S_1 \) obtained from the four fits are consistent with each other and the results reported are those found to have the smallest systematic uncertainty.

Three MC samples are used to study the signal reconstruction and acceptance. Two of them follow the SM prediction for the decay angle distributions taken from Ref. [29], with separate samples generated for \( B^0_d \) and \( \overline{B}^0_d \) decays. The third MC sample has \( F_L = 1/3 \) and the angular distributions are generated uniformly in \( \cos \theta_K, \cos \theta_L \) and \( \phi \). The samples are used to study the effect of potential mistagging and reconstruction differences between particle and antiparticle decays and for determination of the acceptance. The acceptance function is defined as the ratio of reconstructed and generated distributions of \( \cos \theta_K, \cos \theta_L, \phi \), i.e. it is compensating for the bias in the angular distributions resulting from triggering, reconstruction and selection of events. It is described by sixth-order (second-order) polynomial distributions for \( \cos \theta_K \) and \( \cos \theta_L (\phi) \) and is assumed to factorise for each angular distribution, i.e. using \( \varepsilon(\cos \theta_K, \cos \theta_L, \phi) = \varepsilon(\cos \theta_K)\varepsilon(\cos \theta_L)\varepsilon(\phi) \). A systematic uncertainty is assessed in order to account for this assumption. The acceptance function multiplies the angular distribution in the fit, i.e. the signal pdf is

\[
P_{kl} = \varepsilon(\cos \theta_K)\varepsilon(\cos \theta_L)\varepsilon(\phi) g(\cos \theta_K, \cos \theta_L, \phi) \cdot G(m_{K\pi\mu\mu}).
\]

where \( g(\cos \theta_K, \cos \theta_L, \phi) \) is an angular differential decay rate resulting from one of the four folding schemes applied to Equation (1) and \( G(m_{K\pi\mu\mu}) \) is the signal mass distribution. The MC sample generated with uniform \( \cos \theta_K, \cos \theta_L \) and \( \phi \) distributions is used to determine the nominal acceptance functions for each of the transformed variables defined in Equations (7)–(10). The other samples are used to estimate the related systematic uncertainty. Among the angular variables the \( \cos \theta_L \) distribution is the most affected by the acceptance. This is a result of the minimum transverse momentum requirements on the muons in the trigger and the larger inefficiency to reconstruct low-momentum muons, such that large values of \( |\cos \theta_L| \) are inaccessible at low \( q^2 \). As \( q^2 \) increases, the acceptance effects become less severe. The \( \cos \theta_K \) distribution is affected by the ability to reconstruct the \( K\pi \) system, but that effect shows no significant variation with \( q^2 \). There is no significant acceptance effect for \( \phi \). Figure 2 shows the acceptance functions used for \( \cos \theta_K \) and \( \cos \theta_L \) for two different \( q^2 \) ranges for the nominal angular distribution given in Equation (1).

![Figure 2: The acceptance functions for (left) \( \cos \theta_K \) and (right) \( \cos \theta_L \) for (solid) \( q^2 \in [0.04, 2.0] \) GeV\(^2\) and (dashed) \( q^2 \in [4.0, 6.0] \) GeV\(^2\), that shape the angular decay rate of Equation (1).](image-url)
5.2 Background modes

The fit to data includes a combinatorial background component that does not peak in the \( m_{K\pi\mu\mu} \) distribution. It is assumed that the background pdf factorises into a product of one-dimensional terms. The mass distribution of this component is described by an exponential function and second-order Chebychev polynomials are used to model the \( \cos \theta_K \), \( \cos \theta_L \) and \( \phi \) distributions. The values of the nuisance parameters describing these shapes are obtained from fits to the data independently for each \( q^2 \) bin.

Inclusive samples of \( b\bar{b} \rightarrow \mu^+\mu^-X \) and \( c\bar{c} \rightarrow \mu^+\mu^-X \) decays and eleven exclusive \( B_d^0, B_s^0, B^+ \) and \( \Lambda_b \) background samples are studied in order to identify contributions of interest to be included in the fit model, or to be considered when estimating systematic uncertainties. The relevant exclusive modes found to be of interest are discussed below. Events with \( B_c \) decays are suppressed by excluding the \( q^2 \) range containing the \( J/\psi \) and \( \psi(2S) \), and by charm meson vetoes discussed in Section 7. The exclusive background decays considered for the signal mode are \( \Lambda_b \rightarrow \Lambda(1520)\mu^+\mu^- \), \( \Lambda_b \rightarrow pK^-\mu^+\mu^- \), \( B^+ \rightarrow K^{(c+)\mu^+\mu^-} \) and \( B_s^0 \rightarrow \phi\mu^+\mu^- \). These background contributions are accounted for as systematic uncertainties estimated as described in Section 7.

Two distinct background contributions not considered above are observed in the \( \cos \theta_K \) and \( \cos \theta_L \) distributions. They are not accounted for in the nominal fit to data, and are treated as systematic effects. A peak is found in the \( \cos \theta_K \) distribution near 1.0 and appears to have contributions from at least two distinct sources. One of these arises from misreconstructed \( B^+ \) decays, such as \( B^+ \rightarrow K^+\mu\mu \) and \( B^+ \rightarrow \pi^+\mu\mu \). These decays can be reconstructed as signal if another track is combined with the hadron to form a \( K^+ \) candidate in such a way that the event passes the reconstruction and selection. The second contribution comes from combinations of two charged tracks that pass the selection and are reconstructed as a \( K^+ \) candidate. These fake \( K^+ \) candidates accumulate around \( \cos \theta_K \) of 1.0 and are observed in the \( K\pi \) mass sidebands away from the \( K^+ \) meson. They are distinct from the structure of expected \( S-, P- \) and \( D- \) wave \( K\pi \) decays resulting from a signal \( B_d^0 \rightarrow K\pi\mu\mu \) transition. The origin of this source of background is not fully understood. The observed excess may arise from a statistical fluctuation, an unknown background process, or a combination of both. Systematic uncertainties are assigned to evaluate the effect of these two background contributions, as described in Section 7.

Another peak is found in the \( \cos \theta_L \) distribution near values of \( \pm0.7 \). It is associated with partially reconstructed \( B \) decays into final states with a charm meson. This is studied using Monte Carlo simulated events for the decays \( D^0 \rightarrow K^-\pi^+ \), \( D^+ \rightarrow K^-\pi^+\pi^+ \) and \( D^+_s \rightarrow K^+K^-\pi^+ \). Events with a \( B \) meson decaying via an intermediate charm meson \( D^0, D^+ \) or \( D^+_s \) are found to pass the selection and are reconstructed in such a way that they accumulate around 0.7 in \( |\cos \theta_L| \). These are removed from the data sample when estimating systematic uncertainties, as described in Section 7.

5.3 \( K^*c\bar{c} \) control sample fits

The mass distribution obtained from the simulated samples for \( K^*c\bar{c} \) decays, respectively as \( q^2 \in [8, 11] \) and \( [12, 15] \) GeV\(^2\), and the signal mode, in different bins of \( q^2 \), are found to be consistent with each other. Values of \( m_0 \) and \( \xi \) for \( B_d^0 \rightarrow K^*J/\psi \) and \( B_d^0 \rightarrow K^*\psi(2S) \) events are used for the signal pdf and extracted from fits to the data. An extended unbinned maximum-likelihood fit is performed in the two \( K^*c\bar{c} \) control region samples. There are three exclusive backgrounds included: \( \Lambda_b \rightarrow \Lambda c\bar{c}, B^+ \rightarrow K^+c\bar{c} \) and \( B_s^0 \rightarrow K^*c\bar{c} \). The \( K^*c\bar{c} \) pdf has the same form as the signal model, combinatorial background is described by an exponential distribution, and double and triple Gaussian pdfs determined from MC
simulated events are used to describe the exclusive background contributions. A systematic uncertainty is evaluated by allowing for 0, 1, 2 and 3 exclusive background components. The control sample fit projections for the variant of the fit including all three exclusive backgrounds can be found in Figure 3. The impact of the used exclusive background model on the peak position and scale factor of the signal pdf is negligible. From these fits the statistical and systematic uncertainties in the values of \( m_0 \) and \( \xi \) are extracted for the \( B_d^0 \) component in order to be used in the \( B_d^0 \rightarrow K^\ast \mu^+\mu^- \) fits. From the \( J/\psi \) control data it is determined that the values for the nuisance parameters describing the signal model pdf in the \( K\pi\mu\mu \) mass are \( m_0 = 5276.6 \pm 0.3 \pm 0.4 \) MeV and \( \xi = 1.210 \pm 0.004 \pm 0.002 \), where the uncertainties are statistical and systematic, respectively. The \( \psi(2S) \) sample yields compatible results albeit with larger uncertainties. These results are similar to those obtained from the MC simulated samples, and the numbers derived from the \( K^\ast J/\psi \) data are used for the signal region fits.

Figure 3: Fits to the \( K\pi\mu\mu \) invariant mass distributions for the (left) \( K^\ast J/\psi \) and (right) \( K^\ast \psi(2S) \) control region samples. The data are shown as points and the total fit model as the solid lines. The dashed lines represent (black) signal, (red) combinatorial background, (green) \( \Lambda_b \) background, (blue) \( B^+ \) background and (magenta) \( B_s^0 \) background components.

### 5.4 Fitting procedure and validation

A two-step fit process is performed for the different signal bins in \( q^2 \). The first step is a fit to the \( K\pi\mu^+\mu^- \) invariant mass distribution, using the event-by-event uncertainty in the reconstructed mass as a conditional variable. For this fit, the parameters \( m_0 \) and \( \xi \) are fixed to the values obtained from fits to data control samples as described in Section 5.3. A second step adds the (transformed) \( \cos \theta_K \), \( \cos \theta_L \) and \( \phi \) variables to the likelihood in order to extract \( F_L \) and the \( S \) parameters along with the values for the nuisance parameters related to the combinatorial background shapes. Some nuisance parameters, namely \( m_0 \), \( \xi \), signal and background yields, and the exponential shape parameter for the background mass pdf, are fixed to the results obtained from the first step.

The fit procedure is validated using ensembles of simulated pseudo-experiments generated with the \( F_L \) and \( S \) parameters corresponding to those obtained from the data. The purpose of these experiments is to measure the intrinsic fit bias resulting from the likelihood estimator used to extract signal parameters. These ensembles are also used to check that the uncertainties extracted from the fit are consistent with expectations. Ensembles of simulated pseudo-experiments are performed in which signal MC events are
injected into samples of background events generated from the likelihood. The signal yield determined from the first step in the fit process is found to be unbiased. The angular parameters extracted from the nominal fits have biases with magnitudes ranging between 0.01 and 0.04, depending on the fit variation and $q^2$ bin. A similar procedure is used to estimate the effect of neglecting $S$-wave contamination in the data sample. Neglecting the $S$-wave component in the fit model results in a bias between 0.00 and 0.02 in the angular parameters. Similarly, neglecting exclusive background contributions from $\Lambda_B$, $B^+$ and $B^0_s$ decays that peak in $m_{K\pi\mu\mu}$ near the $B^0_d$ mass results in a bias of less than 0.01 on the angular parameters. All these effects are included in the systematic uncertainties described in section 7. The $P^{(r)}$ parameters are obtained using the fit results and covariance matrices from the second fit along with Equations (2)–(5).

6 Results

The event yields obtained from the fits are summarised in Table 1 where only statistical uncertainties are reported. Figures 4 through 9 show for the different $q^2$ bins the distributions of the variables used in the fit for the $S_5$ folding scheme (corresponding to the transformation of Equation (8)) with the total, signal and background fitted pdfs superimposed. Similar sets of distributions are obtained for the three other folding schemes: $S_4$, $S_7$ and $S_8$. The results of the angular fits to the data in terms of the $S_i$ and $P^{(r)}_j$ can be found in Tables 2 and 3. Statistical and systematic uncertainties are quoted in the tables. The distributions of $F_L$ and the $S_i$ parameters as a function of $q^2$ are shown in Figure 10 and those for $P^{(r)}_j$ are shown in Figure 11. The correlations between $F_L$ and the $S_i$ parameters and between $F_L$ and the $P^{(r)}_j$ are given in Appendix A.

Table 1: The values of fitted signal, $n_{\text{signal}}$, and background, $n_{\text{background}}$, yields obtained for different bins in $q^2$. The uncertainties indicated are statistical.

<table>
<thead>
<tr>
<th>$q^2$ [GeV$^2$]</th>
<th>$n_{\text{signal}}$</th>
<th>$n_{\text{background}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>[0.04, 2.0]</td>
<td>128 ± 22</td>
<td>122 ± 22</td>
</tr>
<tr>
<td>[2.0, 4.0]</td>
<td>106 ± 23</td>
<td>113 ± 23</td>
</tr>
<tr>
<td>[4.0, 6.0]</td>
<td>114 ± 24</td>
<td>204 ± 26</td>
</tr>
<tr>
<td>[0.04, 4.0]</td>
<td>236 ± 31</td>
<td>233 ± 32</td>
</tr>
<tr>
<td>[1.1, 6.0]</td>
<td>275 ± 35</td>
<td>363 ± 36</td>
</tr>
<tr>
<td>[0.04, 6.0]</td>
<td>342 ± 39</td>
<td>445 ± 40</td>
</tr>
</tbody>
</table>
Figure 4: The distributions of (top left) $m_{\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos \theta_K$, and (bottom right) $\cos \theta_L$ obtained for $q^2 \in [0.04, 2.0]$ GeV$^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S_5$ folding scheme.
Figure 5: The distributions of (top left) $m_{K\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos\theta_K$, and (bottom right) $\cos\theta_L$ obtained for $q^2 \in [2.0, 4.0]$ GeV$^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S_5$ folding scheme.
Figure 6: The distributions of (top left) $m_{K\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos \theta_K$, and (bottom right) $\cos \theta_L$ obtained for $q^2 \in [4.0, 6.0]$ GeV$^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S_5$ folding scheme.
Figure 7: The distributions of (top left) $m_{K\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos\theta_K$, and (bottom right) $\cos\theta_L$ obtained for $q^2 \in [0.04, 4.0] \text{ GeV}^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S_5$ folding scheme.
Figure 8: The distributions of (top left) $m_{K\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos \theta_K$, and (bottom right) $\cos \theta_L$ obtained for $q^2 \in [1.1, 6.0]$ GeV$^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S^5$ folding scheme.
Figure 9: The distributions of (top left) $m_{K\pi\mu\mu}$, (top right) $\phi$, (bottom left) $\cos \theta_K$, and (bottom right) $\cos \theta_L$ obtained for $q^2 \in [0.04, 6.0]$ GeV$^2$. The (blue) solid line is a projection of the total pdf, the (red) dot-dashed line represents the background, and the (black) dashed line represents the signal component. These plots are obtained from a fit using the $S_5$ folding scheme.
Table 2: The values of $F$, $S$, $S_1$, $S_2$, and $S_3$ parameters obtained for different bins in $q^2$. The uncertainties indicated are statistical and systematic, respectively.

<table>
<thead>
<tr>
<th>$q^2$ (GeV$^2$)</th>
<th>$F_1$</th>
<th>$S_1$</th>
<th>$S_2$</th>
<th>$S_3$</th>
</tr>
</thead>
<tbody>
<tr>
<td>[0.04, 0.20]</td>
<td>0.44 ± 0.08 ± 0.07</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
</tr>
<tr>
<td>[0.20, 0.40]</td>
<td>0.64 ± 0.11 ± 0.05</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
</tr>
<tr>
<td>[0.40, 0.60]</td>
<td>0.42 ± 0.13 ± 0.06</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
</tr>
<tr>
<td>[0.60, 0.80]</td>
<td>0.56 ± 0.07 ± 0.06</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
</tr>
<tr>
<td>[0.80, 1.00]</td>
<td>0.50 ± 0.06 ± 0.04</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
<td>0.33 ± 0.13 ± 0.08</td>
</tr>
</tbody>
</table>

Table 3: The values of $P_1$, $P_2$, $P_3$, $P_4$, and $P_5$ parameters obtained for different bins in $q^2$. The uncertainties indicated are statistical and systematic, respectively.

<table>
<thead>
<tr>
<th>$q^2$ (GeV$^2$)</th>
<th>$P_1$</th>
<th>$P_2$</th>
<th>$P_3$</th>
<th>$P_4$</th>
<th>$P_5$</th>
</tr>
</thead>
<tbody>
<tr>
<td>[0.04, 0.20]</td>
<td>0.69 ± 0.26 ± 0.16</td>
<td>0.67 ± 0.26 ± 0.16</td>
<td>0.67 ± 0.26 ± 0.16</td>
<td>0.67 ± 0.26 ± 0.16</td>
<td>0.67 ± 0.26 ± 0.16</td>
</tr>
<tr>
<td>[0.20, 0.40]</td>
<td>0.31 ± 0.40 ± 0.20</td>
<td>0.31 ± 0.40 ± 0.20</td>
<td>0.31 ± 0.40 ± 0.20</td>
<td>0.31 ± 0.40 ± 0.20</td>
<td>0.31 ± 0.40 ± 0.20</td>
</tr>
<tr>
<td>[0.40, 0.60]</td>
<td>0.64 ± 0.33 ± 0.18</td>
<td>0.64 ± 0.33 ± 0.18</td>
<td>0.64 ± 0.33 ± 0.18</td>
<td>0.64 ± 0.33 ± 0.18</td>
<td>0.64 ± 0.33 ± 0.18</td>
</tr>
<tr>
<td>[0.60, 0.80]</td>
<td>0.32 ± 0.24 ± 0.17</td>
<td>0.32 ± 0.24 ± 0.17</td>
<td>0.32 ± 0.24 ± 0.17</td>
<td>0.32 ± 0.24 ± 0.17</td>
<td>0.32 ± 0.24 ± 0.17</td>
</tr>
<tr>
<td>[0.80, 1.00]</td>
<td>0.30 ± 0.21 ± 0.13</td>
<td>0.30 ± 0.21 ± 0.13</td>
<td>0.30 ± 0.21 ± 0.13</td>
<td>0.30 ± 0.21 ± 0.13</td>
<td>0.30 ± 0.21 ± 0.13</td>
</tr>
</tbody>
</table>
7 Systematic uncertainties

Systematic uncertainties in the parameter values obtained from the angular analysis come from several sources. The methods for determining these uncertainties are based either on a comparison of nominal and modified fit results, or on observed fit biases in modified pseudo-experiments. The systematic uncertainties are symmetrised. The most significant ones are described in the following, in decreasing order of importance.

- A systematic uncertainty is assigned for the combinatorial $K\pi$ (fake $K^*$) background peaking at $\cos \theta_K$ values around 1.0 obtained by comparing results of the nominal fit to that where data above $\cos \theta_K = 0.9$ are excluded from the fit.

- A systematic uncertainty is derived to account for background arising from partially reconstructed $B \rightarrow D^0 / D^+ / D_s^+ X$ decays, that manifest in an accumulation of events at $|\cos \theta_L|$ values around 0.7. Two-track or three-track combinations are formed from the signal candidate tracks, and are reconstructed assuming the pion or kaon mass hypothesis. A veto is then applied for events in which a track combination has a mass in a window of 30 MeV around the $D^0$, $D^+$ or $D_s^+$ meson mass. Similarly, a veto is implemented to reject $B^+ \rightarrow K^+ \mu^+ \mu^-$ and $B^+ \rightarrow \pi^+ \mu^+ \mu^-$ events that pass the event selection. Here $B^+$ candidates are reconstructed from one of the hadrons from the $K^*$ candidate and the muons in the signal candidate. Signal candidates that have a three-track mass within 50 MeV of the $B^+$ mass are excluded from the fit. A few percent of signal events are removed when applying these vetoes, with a corresponding effect on the acceptance distributions. The fit results obtained from the data samples with vetoes applied are compared to those obtained from the nominal fit and the change in each result is taken as the systematic uncertainty from these backgrounds. This systematic uncertainty dominates the measurement of $F_L$ at higher values of $q^2$.

- The combinatorial background pdf shape has an uncertainty arising from the choice of the model. For the mass distribution it is assumed that an exponential function model is adequate; however, for the angular variables the data are re-fitted using third-order Chebychev polynomials. The change from the nominal result is taken as the uncertainty from this source.

- The acceptance function is assumed to factorise into three separate components, for $\cos \theta_K$, $\cos \theta_L$ and $\phi$. To validate this assumption, the signal simulated events are fitted with the acceptance function obtained from that same MC sample. Differences in the fit results from expectation are small and taken as the uncertainty resulting from this assumption.

- A systematic uncertainty is assigned for the angular pdf model for the background by comparing the nominal result to that with a reduced fit range of $m_{K\pi\mu\mu} \in [5200, 5700]$ MeV, in particular to account for possible residues of the partially reconstructed $B$-decays.

- A correction is applied to the data by shifting the track $p_T$ according to the uncertainties arising from biases in rapidity and momentum scale. The change in results obtained is ascribed to the uncertainty in the ID alignment and knowledge of the magnetic field.

- The maximum-likelihood estimator used is intrinsically biased. Ensembles of MC simulated events are used in order to ascertain the bias in the extracted values of the parameters of interest. The bias is assigned as a systematic uncertainty.
• The $p_T$ spectrum of $B_d^0$ candidates observed in data is not accurately reproduced by the MC simulation. This difference in the kinematics results in a slight modification of the acceptance functions. This is accounted for by reweighting signal MC simulated events to resemble the $p_T$ spectrum found in data. The change in fitted parameter values obtained due to the reweighting is taken as the systematic uncertainty resulting from this difference.

• The signal decay mode is resonant $K^* \to K \pi$ decay, but scalar contributions from non-resonant $K \pi$ transitions may also exist. The LHCb Collaboration reported an $S$-wave contribution at the level of 5% of the signal [4, 30]. Ensembles of MC simulated events are fitted with 5% of the signal being drawn from an $S$-wave sample of events and the remaining 95% from signal. The observed change in fit bias is assigned as the systematic uncertainty from this source. Any variation in $S$-wave content as a function of $q^2$ would not significantly affect the results reported here.

• The values of the nuisance parameters of the fit model obtained from MC control samples and fits to the data mass distribution have associated uncertainties. These parameters include $m_0$, $\xi$, the signal and background yields, the shape parameter of the combinatorial background mass distribution, and the parameters of the signal acceptance functions. The uncertainty in the value of each of these parameters is varied independently in order to assess the effect on parameters of interest. This source of uncertainty has a small effect on the measurements reported here.

• Background from exclusive modes peaking in $m_{K\pi\mu\mu}$ is neglected in the nominal fit. This may affect the fitted results and is accounted for by computing the fit bias obtained when embedding MC simulated samples of $\Lambda_b \to \Lambda(1520)\mu^+\mu^-$. $\Lambda_b \to pK^-\mu^+\mu^-$, $B^+ \to K^{(*)+}\mu^+\mu^-$ and $B_d^0 \to \phi\mu^+\mu^-$ into ensembles of pseudo-data generated from the fit model containing only combinatorial background and signal components. The change in fit bias observed when adding exclusive backgrounds is taken as the systematic error arising from neglecting those modes in the fit.

• The difference from nominal results obtained when fitting the $B_d^0$ signal MC events with the acceptance function for $\overline{B_d^0}$ is taken as an upper limit of the systematic error resulting from event migration due to mistagging the $B_d^0$ flavour.

• The parameters $S_5$ and $S_8$, as well as the respective $P_j^{(*)}$ parameters are affected by dilution and thus have a multiplicative scaling applied to them. This dilution factor depends on the kinematics of the $K^*$ decay and has a systematic uncertainty associated with it. The effect of data/MC differences in the $p_T$ spectrum of $B_d^0$ candidates on the mistag probability was studied and found to be negligible. The uncertainty due to the limited number of MC events is used to compute the statistical uncertainty of $\omega$ and $\overline{\omega}$. Studies of MC simulated events indicate that there is no significant difference between the mistag probability for $B_d^0$ and $\overline{B_d^0}$ events and the analysis assumes that the average mistag probability provides an adequate description of this effect. The magnitude of the mistag probability difference, $|\omega - \overline{\omega}|$, is included as a systematic uncertainty resulting from this assumption.

The total systematic uncertainties of the fitted $S_i$ and $P_j^{(*)}$ parameter values are presented in Tables 2 and 3, where the dominant contributions for $F_L$ come from the modelling of the angular distributions of the combinatorial background and the partially reconstructed decays peaking in $\cos \theta_K$ and $\cos \theta_L$. These contributions and in addition also ID alignment and magnetic field calibration affect $S_3 (P_3)$. The largest systematic uncertainty contribution to $S_3 (P_3)$ comes from partially reconstructed decays entering the signal region. This also affects the measurement of $S_5 (P_4)$ and $S_7 (P_5)$. The partially reconstructed decays peaking in $\cos \theta_L$ affect the measurement of $S_4 (P_4)$ and $S_8 (P_8)$, whereas the fake $K^*$ background in $\cos \theta_K$ affects $S_4 (P_4)$, $S_5 (P_5)$, and $S_8 (P_8)$. The parameterization of the signal acceptance is another
significant systematic uncertainty source for $S_4 (P_4')$. The systematic uncertainties are smaller than the statistical uncertainties for all parameters measured.

8 Comparison with theoretical computations

The results of theoretical approaches of Ciuchini et al. (CFFMPSV) [31], Descotes-Genon et al. (DHMV) [32], and Jäger and Camalich (JC) [33, 34] are shown in Figure 10 for the $S$ parameters, and in Figure 11 for the $P^{(i)}$ parameters, along with the results presented here.\footnote{This result uses the experimental convention of Equations (2)–(5) following the LHCb Collaboration’s notation in Ref. [3]. In the DHMV calculation, a different convention is used as explained by Equation (16) in Ref. [15].}

QCD factorisation is used by DHMV and JC, where the latter focus on the impact of long-distance corrections using a helicity amplitude approach. The CFFMPSV group takes a different approach, using the QCD factorisation framework to perform compatibility checks of the LHCb data with theoretical predictions. This approach also allows information from a given experimentally measured parameter of interest to be excluded in order to make a fit-based prediction of the expected value of that parameter from the rest of the data.

With the exception of the $P_4'$ and $P_5'$ measurements in $q^2 \in [4.0, 6.0] \text{ GeV}^2$ and $P_8'$ in $q^2 \in [2.0, 4.0] \text{ GeV}^2$ there is good agreement between theory and measurement. The $P_4'$ and $P_5'$ parameters have statistical correlation of 0.37 in the $q^2 \in [4.0, 6.0] \text{ GeV}^2$ bin. The observed deviation from the SM prediction of $P_4'$ and $P_5'$ is for both parameters approximately 2.7 standard deviations (local) away from the calculation of DHMV for this bin. The deviations are less significant for the other calculation and the fit approach. All measurements are found to be within three standard deviations of the range covered by the different predictions. Hence, including experimental and theoretical uncertainties, the measurements presented here are found to agree with the predicted SM contributions to this decay.
Figure 10: The measured values of $F_L$, $S_3$, $S_4$, $S_5$, $S_7$, $S_8$ compared with predictions from the theoretical calculations discussed in the text (Section 8). Statistical and total uncertainties are shown for the data, i.e. the inner mark indicates the statistical uncertainty and the total error bar the total uncertainty.
Figure 11: The measured values of $P_1$, $P'_4$, $P'_5$, $P'_6$, $P'_8$ compared with predictions from the theoretical calculations discussed in the text (Section 8). Statistical and total uncertainties are shown for the data, i.e. the inner mark indicates the statistical uncertainty and the total error bar the total uncertainty.
9 Conclusion

The results of an angular analysis of the rare decay $B_d^0 \rightarrow K^* \mu^+ \mu^-$ are presented. This flavour-changing neutral current process is sensitive to potential new-physics contributions. The $B_d^0 \rightarrow K^* \mu^+ \mu^-$ analysis presented here uses a total of 20.3 fb$^{-1}$ of $\sqrt{s} = 8$ TeV pp collision data collected by the ATLAS experiment at the LHC in 2012. An extended unbinned maximum-likelihood fit of the angular distribution of the signal decay is performed in order to extract the parameters $F_L, S_i$ and $P(j')$ in six bins of $q^2$. Three of these bins overlap in order to report results in ranges compatible with other experiments and phenomenology studies. All measurements are found to be within three standard deviation of the range covered by the different predictions. The results are also compatible with the results of the LHCb, CMS and Belle collaborations.

Acknowledgements

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently.

We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF and DNSRC, Denmark; IN2P3-CNRS, CEA-DRF/IRFU, France; SRNSFG, Georgia; BMBF, HGF, and MPG, Germany; GSRT, Greece; RGC, Hong Kong SAR, China; ISF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; NWO, Netherlands; RCN, Norway; MNiSW and NCN, Poland; FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TAEK, Turkey; STFC, United Kingdom; DOE and NSF, United States of America. In addition, individual groups and members have received support from BCKDF, the Canada Council, CANARIE, CRC, Compute Canada, FQRNT, and the Ontario Innovation Trust, Canada; EPLANET, ERC, ERDF, FP7, Horizon 2020 and Marie Skłodowska-Curie Actions, European Union; Investissements d’Avenir Labex and Idex, ANR, Région Auvergne and Fondation Partager le Savoir, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF; BSF, GIF and Minerva, Israel; BRF, Norway; CERCA Programme Generalitat de Catalunya, Generalitat Valenciana, Spain; the Royal Society and Leverhulme Trust, United Kingdom.

The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN, the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA), the Tier-2 facilities worldwide and large non-WLCG resource providers. Major contributors of computing resources are listed in Ref. [35].
Appendix

A Correlation Matrices

Four folding schemes are applied to the data in order to extract \( F_L, S_3, S_4, S_5, S_7 \) and \( S_8 \) from four separate fits. The \( P^{(i)} \) parameters are subsequently derived from the fit results using Equations (2)–(5). It is not possible to extract a full correlation matrix between fitted parameters obtained from different fits. In order to reconstruct the correlation matrix, ensembles of pseudo-experiments are simulated using the pdf corresponding to the nominal angular distributions. Each simulated ensemble has the four folding schemes applied to it and four fits are performed on the resulting samples. The distributions obtained for pairs of parameters obtained from fits to these ensembles are used to compute Pearson correlation coefficients for those pairs. Correlation matrices for \( F_L \) and the \( S \) parameters are reconstructed from all possible pairings for a given \( q^2 \) bin. A similar method is used to extract the correlation matrices for the \( P^{(i)} \) parameters. This procedure is repeated for each \( q^2 \) bin studied in order to obtain correlation matrices given in the remainder of this appendix. The correlation matrices are statistical only. Contributions from systematic uncertainties are not included, since the measurement precision is statistically limited.

- Table 4 (5) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([0.04, 2.0] \) GeV\(^2\).
- Table 6 (7) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([2.0, 4.0] \) GeV\(^2\).
- Table 8 (9) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([4.0, 6.0] \) GeV\(^2\).
- Table 10 (11) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([0.04, 4.0] \) GeV\(^2\).
- Table 12 (13) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([1.1, 6.0] \) GeV\(^2\).
- Table 14 (15) shows the statistical correlation matrix for \( F_L \) and \( S \) \( (P^{(i)}) \) parameters for the \( q^2 \) bin \([0.04, 6.0] \) GeV\(^2\).

Table 4: Statistical correlation matrix for the \( F_L \) and \( S \) parameters obtained for \( q^2 \in [0.04, 2.0] \) GeV\(^2\).

<table>
<thead>
<tr>
<th></th>
<th>( F_L )</th>
<th>( S_3 )</th>
<th>( S_4 )</th>
<th>( S_5 )</th>
<th>( S_7 )</th>
<th>( S_8 )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( F_L )</td>
<td>1.00</td>
<td>0.11</td>
<td>-0.13</td>
<td>0.03</td>
<td>0.16</td>
<td>0.24</td>
</tr>
<tr>
<td>( S_3 )</td>
<td>1.00</td>
<td>0.31</td>
<td>0.28</td>
<td>0.73</td>
<td>0.45</td>
<td></td>
</tr>
<tr>
<td>( S_4 )</td>
<td>1.00</td>
<td>0.58</td>
<td>0.14</td>
<td>0.28</td>
<td></td>
<td></td>
</tr>
<tr>
<td>( S_5 )</td>
<td>1.00</td>
<td>0.14</td>
<td>0.28</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>( S_7 )</td>
<td>1.00</td>
<td>0.59</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>( S_8 )</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>
Table 5: Statistical correlation matrix for the $P^{(\prime)}$ parameters obtained for $q^2 \in [0.04, 2.0] \text{GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P_4'$</th>
<th>$P_5'$</th>
<th>$P_6'$</th>
<th>$P_8'$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>0.04</td>
<td>0.05</td>
<td>0.62</td>
<td>0.32</td>
</tr>
<tr>
<td>$P_4'$</td>
<td>1.00</td>
<td>0.53</td>
<td>$-0.08$</td>
<td>$-0.06$</td>
<td></td>
</tr>
<tr>
<td>$P_5'$</td>
<td>1.00</td>
<td></td>
<td>0.00</td>
<td>0.22</td>
<td></td>
</tr>
<tr>
<td>$P_6'$</td>
<td>1.00</td>
<td></td>
<td></td>
<td>0.55</td>
<td></td>
</tr>
<tr>
<td>$P_8'$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 6: Statistical correlation matrix for the $F_L$ and $S$ parameters obtained for $q^2 \in [2.0, 4.0] \text{GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$F_L$</th>
<th>$S_3$</th>
<th>$S_4$</th>
<th>$S_5$</th>
<th>$S_7$</th>
<th>$S_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$F_L$</td>
<td>1.00</td>
<td>0.27</td>
<td>0.35</td>
<td>$-0.04$</td>
<td>$-0.15$</td>
<td>$-0.37$</td>
</tr>
<tr>
<td>$S_3$</td>
<td>1.00</td>
<td>$-0.08$</td>
<td>$-0.44$</td>
<td>$-0.09$</td>
<td>$-0.20$</td>
<td></td>
</tr>
<tr>
<td>$S_4$</td>
<td>1.00</td>
<td></td>
<td>0.60</td>
<td>$-0.02$</td>
<td>$-0.12$</td>
<td></td>
</tr>
<tr>
<td>$S_5$</td>
<td>1.00</td>
<td></td>
<td></td>
<td>$-0.11$</td>
<td>$-0.20$</td>
<td></td>
</tr>
<tr>
<td>$S_7$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
<td>0.63</td>
</tr>
<tr>
<td>$S_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 7: Statistical correlation matrix for the $P^{(\prime)}$ parameters obtained for $q^2 \in [2.0, 4.0] \text{GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P_4'$</th>
<th>$P_5'$</th>
<th>$P_6'$</th>
<th>$P_8'$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>$-0.12$</td>
<td>$-0.21$</td>
<td>0.05</td>
<td>0.05</td>
</tr>
<tr>
<td>$P_4'$</td>
<td>1.00</td>
<td></td>
<td>0.51</td>
<td>0.08</td>
<td>0.03</td>
</tr>
<tr>
<td>$P_5'$</td>
<td>1.00</td>
<td></td>
<td></td>
<td>$-0.23$</td>
<td>0.22</td>
</tr>
<tr>
<td>$P_6'$</td>
<td>1.00</td>
<td></td>
<td></td>
<td></td>
<td>0.66</td>
</tr>
<tr>
<td>$P_8'$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Table 8: Statistical correlation matrix for the $F_L$ and $S$ parameters obtained for $q^2 \in [4.0, 6.0] \text{GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$F_L$</th>
<th>$S_3$</th>
<th>$S_4$</th>
<th>$S_5$</th>
<th>$S_7$</th>
<th>$S_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$F_L$</td>
<td>1.00</td>
<td>0.33</td>
<td>$-0.18$</td>
<td>0.04</td>
<td>0.22</td>
<td>0.28</td>
</tr>
<tr>
<td>$S_3$</td>
<td>1.00</td>
<td>0.15</td>
<td>0.23</td>
<td>0.60</td>
<td>0.05</td>
<td></td>
</tr>
<tr>
<td>$S_4$</td>
<td>1.00</td>
<td>0.52</td>
<td>0.03</td>
<td>0.01</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_5$</td>
<td>1.00</td>
<td>0.28</td>
<td>0.27</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_7$</td>
<td>1.00</td>
<td>0.60</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_8$</td>
<td></td>
<td>1.00</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
Table 9: Statistical correlation matrix for the $P'(\prime)$ parameters obtained for $q^2 \in [4.0, 6.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P'_4$</th>
<th>$P'_5$</th>
<th>$P'_6$</th>
<th>$P'_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>0.11</td>
<td>0.34</td>
<td>0.41</td>
<td>0.16</td>
</tr>
<tr>
<td>$P'_4$</td>
<td>1.00</td>
<td>0.37</td>
<td>0.06</td>
<td>0.04</td>
<td></td>
</tr>
<tr>
<td>$P'_5$</td>
<td>1.00</td>
<td>0.39</td>
<td>0.33</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_6$</td>
<td>1.00</td>
<td>0.62</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 10: Statistical correlation matrix for the $F_L$ and $S$ parameters obtained for $q^2 \in [0.04, 4.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$F_L$</th>
<th>$S_3$</th>
<th>$S_4$</th>
<th>$S_5$</th>
<th>$S_7$</th>
<th>$S_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$F_L$</td>
<td>1.00</td>
<td>0.08</td>
<td>0.05</td>
<td>0.01</td>
<td>0.18</td>
<td>0.14</td>
</tr>
<tr>
<td>$S_3$</td>
<td>1.00</td>
<td>-0.04</td>
<td>0.03</td>
<td>0.29</td>
<td>-0.16</td>
<td></td>
</tr>
<tr>
<td>$S_4$</td>
<td>1.00</td>
<td>0.79</td>
<td>0.08</td>
<td>0.03</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_5$</td>
<td>1.00</td>
<td>0.03</td>
<td>-0.02</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_7$</td>
<td>1.00</td>
<td>0.60</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 11: Statistical correlation matrix for the $P'(\prime)$ parameters obtained for $q^2 \in [0.04, 4.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P'_4$</th>
<th>$P'_5$</th>
<th>$P'_6$</th>
<th>$P'_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>-0.07</td>
<td>0.00</td>
<td>0.21</td>
<td>0.12</td>
</tr>
<tr>
<td>$P'_4$</td>
<td>1.00</td>
<td>0.78</td>
<td>0.08</td>
<td>0.02</td>
<td></td>
</tr>
<tr>
<td>$P'_5$</td>
<td>1.00</td>
<td>0.03</td>
<td>-0.04</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_6$</td>
<td>1.00</td>
<td>0.59</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 12: Statistical correlation matrix for the $F_L$ and $S$ parameters obtained for $q^2 \in [1.1, 6.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$F_L$</th>
<th>$S_3$</th>
<th>$S_4$</th>
<th>$S_5$</th>
<th>$S_7$</th>
<th>$S_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$F_L$</td>
<td>1.00</td>
<td>0.05</td>
<td>0.00</td>
<td>0.07</td>
<td>0.18</td>
<td>-0.03</td>
</tr>
<tr>
<td>$S_3$</td>
<td>1.00</td>
<td>0.41</td>
<td>0.46</td>
<td>0.32</td>
<td>-0.01</td>
<td></td>
</tr>
<tr>
<td>$S_4$</td>
<td>1.00</td>
<td>0.60</td>
<td>0.09</td>
<td>0.03</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_5$</td>
<td>1.00</td>
<td>0.17</td>
<td>0.24</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_7$</td>
<td>1.00</td>
<td>0.67</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>
Table 13: Statistical correlation matrix for the $P^{(r)}$ parameters obtained for $q^2 \in [1.1, 6.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P'_4$</th>
<th>$P'_5$</th>
<th>$P'_6$</th>
<th>$P'_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>0.23</td>
<td>-0.09</td>
<td>0.08</td>
<td>-0.07</td>
</tr>
<tr>
<td>$P'_4$</td>
<td>1.00</td>
<td>0.53</td>
<td>0.15</td>
<td>0.08</td>
<td></td>
</tr>
<tr>
<td>$P'_5$</td>
<td>1.00</td>
<td>0.28</td>
<td>0.24</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_6$</td>
<td>1.00</td>
<td>0.67</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 14: Statistical correlation matrix for the $F_L$ and $S$ parameters obtained for $q^2 \in [0.04, 6.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$F_L$</th>
<th>$S_3$</th>
<th>$S_4$</th>
<th>$S_5$</th>
<th>$S_7$</th>
<th>$S_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$F_L$</td>
<td>1.00</td>
<td>0.03</td>
<td>0.01</td>
<td>-0.10</td>
<td>0.13</td>
<td>0.06</td>
</tr>
<tr>
<td>$S_3$</td>
<td>1.00</td>
<td>-0.02</td>
<td>-0.09</td>
<td>0.32</td>
<td>-0.01</td>
<td></td>
</tr>
<tr>
<td>$S_4$</td>
<td></td>
<td>1.00</td>
<td>0.68</td>
<td>0.00</td>
<td>0.04</td>
<td></td>
</tr>
<tr>
<td>$S_5$</td>
<td></td>
<td></td>
<td>1.00</td>
<td>-0.05</td>
<td>0.03</td>
<td></td>
</tr>
<tr>
<td>$S_7$</td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
<td>0.65</td>
<td></td>
</tr>
<tr>
<td>$S_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>

Table 15: Statistical correlation matrix for the $P^{(r)}$ parameters obtained for $q^2 \in [0.04, 6.0] \text{ GeV}^2$.

<table>
<thead>
<tr>
<th></th>
<th>$P_1$</th>
<th>$P'_4$</th>
<th>$P'_5$</th>
<th>$P'_6$</th>
<th>$P'_8$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$P_1$</td>
<td>1.00</td>
<td>-0.02</td>
<td>-0.14</td>
<td>0.17</td>
<td>0.04</td>
</tr>
<tr>
<td>$P'_4$</td>
<td>1.00</td>
<td>0.65</td>
<td>0.00</td>
<td>0.05</td>
<td></td>
</tr>
<tr>
<td>$P'_5$</td>
<td>1.00</td>
<td>-0.06</td>
<td>0.09</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_6$</td>
<td>1.00</td>
<td>0.61</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$P'_8$</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>1.00</td>
</tr>
</tbody>
</table>
References


of Physics, TOBB University of Economics and Technology, Ankara; Turkey.
5 LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy; France.
6 High Energy Physics Division, Argonne National Laboratory, Argonne IL; United States of America.
7 Department of Physics, University of Arizona, Tucson AZ; United States of America.
8 Department of Physics, University of Texas at Arlington, Arlington TX; United States of America.
9 Physics Department, National and Kapodistrian University of Athens, Athens; Greece.
10 Physics Department, National Technical University of Athens, Zografou; Greece.
11 Department of Physics, University of Texas at Austin, Austin TX; United States of America.
12 (a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; (b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; (c) Department of Physics, Bogazici University, Istanbul; (d) Department of Physics Engineering, Gaziantep University, Gaziantep; Turkey.
13 Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
14 Institut de Física d’Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain.
15 (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (b) Physics Department, Tsinghua University, Beijing; (c) Department of Physics, Nanjing University, Nanjing; (d) University of Chinese Academy of Science (UCAS), Beijing; China.
16 Institute of Physics, University of Belgrade, Belgrade; Serbia.
17 Department for Physics and Technology, University of Bergen, Bergen; Norway.
18 Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA; United States of America.
19 Institut für Physik, Humboldt Universität zu Berlin, Berlin; Germany.
20 Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern; Switzerland.
21 School of Physics and Astronomy, University of Birmingham, Birmingham; United Kingdom.
22 Centro de Investigaciónes, Universidad Antonio Nariño, Bogota; Colombia.
23 (a) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna; (b) INFN Sezione di Bologna; Italy.
24 Physikalisches Institut, Universität Bonn, Bonn; Germany.
25 Department of Physics, Boston University, Boston MA; United States of America.
26 Department of Physics, Brandeis University, Waltham MA; United States of America.
27 (a) Transilvania University of Brasov, Brasov; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; (d) National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; (e) University Politehnica Bucharest, Bucharest; (f) West University in Timisoara, Timisoara; Romania.
28 (a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice; Slovak Republic.
29 Physics Department, Brookhaven National Laboratory, Upton NY; United States of America.
30 Departamento de Física, Universidad de Buenos Aires, Buenos Aires; Argentina.
31 Cavendish Laboratory, University of Cambridge, Cambridge; United Kingdom.
32 (a) Department of Physics, University of Cape Town, Cape Town; (b) Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; (c) School of Physics, University of the Witwatersrand, Johannesburg; South Africa.
33 Department of Physics, Carleton University, Ottawa ON; Canada.
34 (a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université
Hassan II, Casablanca; (b) Centre National de l'Énergie des Sciences Techniques Nucléaires (CNESTEN), Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V, Rabat; Morocco.
35 CERN, Geneva; Switzerland.
36 Enrico Fermi Institute, University of Chicago, Chicago IL; United States of America.
37 LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand; France.
38 Nevis Laboratory, Columbia University, Irvington NY; United States of America.
39 Niels Bohr Institute, University of Copenhagen, Copenhagen; Denmark.
40 (a) Dipartimento di Fisica, Università della Calabria, Rende; (b) INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Italy.
41 Physics Department, Southern Methodist University, Dallas TX; United States of America.
42 Physics Department, University of Texas at Dallas, Richardson TX; United States of America.
43 (a) Department of Physics, Stockholm University; (b) Oskar Klein Centre, Stockholm; Sweden.
44 Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen; Germany.
45 Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund; Germany.
46 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden; Germany.
47 Department of Physics, Duke University, Durham NC; United States of America.
48 SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh; United Kingdom.
49 INFN e Laboratori Nazionali di Frascati, Frascati; Italy.
50 Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.
51 II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.
52 Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
53 (a) Dipartimento di Fisica, Università di Genova, Genova; (b) INFN Sezione di Genova; Italy.
54 SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow; United Kingdom.
55 LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble; France.
56 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA; United States of America.
57 (a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; (b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; (c) School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; (d) Tsung-Dao Lee Institute, Shanghai; China.
58 (a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; Germany.
59 (a) Department of Applied Information Science, Hiroshima Institute of Technology, Hiroshima; Japan.
60 Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; (b) Department of Physics, University of Hong Kong, Hong Kong; (c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong; China.
61 Department of Physics, National Tsing Hua University, Hsinchu; Taiwan.
62 Department of Physics, Indiana University, Bloomington IN; United States of America.
63 INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; (b) ICTP, Trieste; (c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine; Italy.
64 INFN Sezione di Lecce; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce; Italy.
65 INFN Sezione di Milano; (b) Dipartimento di Fisica, Università di Milano, Milano; Italy.
INFN Sezione di Napoli; Dipartimento di Fisica, Università di Napoli, Napoli; Italy.

INFN Sezione di Pavia; Dipartimento di Fisica, Università di Pavia, Pavia; Italy.

INFN Sezione di Pisa; Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.

INFN Sezione di Roma; Dipartimento di Fisica, Sapienza Università di Roma, Roma; Italy.

INFN Sezione di Roma Tor Vergata; Dipartimento di Fisica, Università di Roma Tor Vergata, Roma; Italy.

INFN Sezione di Roma Tre; Dipartimento di Matematica e Fisica, Università Roma Tre, Roma; Italy.

INFN-TIFPA; Università degli Studi di Trento, Trento; Italy.

Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck; Austria.

University of Iowa, Iowa City IA; United States of America.

Department of Physics and Astronomy, Iowa State University, Ames IA; United States of America.

Joint Institute for Nuclear Research, Dubna; Russia.

Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; Universidade Federal de São João del Rei (UFSJ), São João del Rei; Instituto de Física, Universidade de São Paulo, São Paulo; Brazil.

KEK, High Energy Accelerator Research Organization, Tsukuba; Japan.

Graduate School of Science, Kobe University, Kobe; Japan.

AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow; Poland.

Institute of Nuclear Physics Polish Academy of Sciences, Krakow; Poland.

Faculty of Science, Kyoto University, Kyoto; Japan.

Kyoto University of Education, Kyoto; Japan.

Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka; Japan.

Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata; Argentina.

Physics Department, Lancaster University, Lancaster; United Kingdom.

Oliver Lodge Laboratory, University of Liverpool, Liverpool; United Kingdom.

Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana; Slovenia.

School of Physics and Astronomy, Queen Mary University of London, London; United Kingdom.

Department of Physics, Royal Holloway University of London, Egham; United Kingdom.

Department of Physics and Astronomy, University College London, London; United Kingdom.

Louisiana Tech University, Ruston LA; United States of America.

Fysiska institutionen, Lunds universitet, Lund; Sweden.

Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne; France.

Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid; Spain.

Institut für Physik, Universität Mainz, Mainz; Germany.

School of Physics and Astronomy, University of Manchester, Manchester; United Kingdom.

CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.

Department of Physics, University of Massachusetts, Amherst MA; United States of America.

Department of Physics, McGill University, Montreal QC; Canada.

School of Physics, University of Melbourne, Victoria; Australia.

Department of Physics, University of Michigan, Ann Arbor MI; United States of America.

Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
America.

105 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk; Belarus.

106 Research Institute for Nuclear Problems of Byelorussian State University, Minsk; Belarus.

107 Group of Particle Physics, University of Montreal, Montreal QC; Canada.

108 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow; Russia.

109 Institute for Theoretical and Experimental Physics (ITEP), Moscow; Russia.

110 National Research Nuclear University MEPhI, Moscow; Russia.

111 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow; Russia.

112 Fakultät für Physik, Ludwig-Maximilians-Universität München, München; Germany.

113 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München; Germany.

114 Nagasaki Institute of Applied Science, Nagasaki; Japan.

115 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya; Japan.

116 Department of Physics and Astronomy, University of New Mexico, Albuquerque NM; United States of America.

117 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.

118 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam; Netherlands.

119 Department of Physics, Northern Illinois University, DeKalb IL; United States of America.

120 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk; Novosibirsk State University Novosibirsk; Russia.

121 Department of Physics, New York University, New York NY; United States of America.

122 Ohio State University, Columbus OH; United States of America.

123 Faculty of Science, Okayama University, Okayama; Japan.

124 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK; United States of America.

125 Department of Physics, Oklahoma State University, Stillwater OK; United States of America.

126 Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc; Czech Republic.

127 Center for High Energy Physics, University of Oregon, Eugene OR; United States of America.

128 LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.

129 Graduate School of Science, Osaka University, Osaka; Japan.

130 Department of Physics, University of Oslo, Oslo; Norway.

131 Department of Physics, Oxford University, Oxford; United Kingdom.

132 LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris; France.

133 Department of Physics, University of Pennsylvania, Philadelphia PA; United States of America.

134 Konstantinov Nuclear Physics Institute of National Research Centre "Kurchatov Institute", PNPI, St. Petersburg; Russia.

135 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA; United States of America.

136 Laboratório de Instrumentação e Física Experimental de Partículas - LIP; Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; Departamento de Física, Universidade de Coimbra, Coimbra; Centro de Física Nuclear da Universidade de Lisboa, Lisboa; Departamento de Física, Universidade do Minho, Braga; Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica; Portugal.

137 Institute of Physics, Academy of Sciences of the Czech Republic, Prague; Czech Republic.
138 Czech Technical University in Prague, Prague; Czech Republic.
139 Charles University, Faculty of Mathematics and Physics, Prague; Czech Republic.
140 State Research Center Institute for High Energy Physics, NRC KI, Protvino; Russia.
141 Particle Physics Department, Rutherford Appleton Laboratory, Didcot; United Kingdom.
142 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette; France.
143 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA; United States of America.
144(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso; Chile.
145 Department of Physics, University of Washington, Seattle WA; United States of America.
146 Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.
147 Department of Physics, Shinshu University, Nagano; Japan.
148 Department Physik, Universität Siegen, Siegen; Germany.
149 Department of Physics, Simon Fraser University, Burnaby BC; Canada.
150 SLAC National Accelerator Laboratory, Stanford CA; United States of America.
151 Physics Department, Royal Institute of Technology, Stockholm; Sweden.
152 Departments of Physics and Astronomy, Stony Brook University, Stony Brook NY; United States of America.
153 Department of Physics and Astronomy, University of Sussex, Brighton; United Kingdom.
154 School of Physics, University of Sydney, Sydney; Australia.
155 Institute of Physics, Academia Sinica, Taipei; Taiwan.
156 Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei; Taiwan.
157(a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (b) High Energy Physics Institute, Tbilisi State University, Tbilisi; Georgia.
158 Department of Physics, Technion, Israel Institute of Technology, Haifa; Israel.
159 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv; Israel.
160 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki; Greece.
161 International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo; Japan.
162 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo; Japan.
163 Department of Physics, Tokyo Institute of Technology, Tokyo; Japan.
164 Tomsk State University, Tomsk; Russia.
165 Department of Physics, University of Toronto, Toronto ON; Canada.
166(a) TRIUMF, Vancouver BC; (b) Department of Physics and Astronomy, York University, Toronto ON; Canada.
167 Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba; Japan.
168 Department of Physics and Astronomy, Tufts University, Medford MA; United States of America.
169 Department of Physics and Astronomy, University of California Irvine, Irvine CA; United States of America.
170 Department of Physics and Astronomy, University of Uppsala, Uppsala; Sweden.
171 Department of Physics, University of Illinois, Urbana IL; United States of America.
172 Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia; Spain.
173 Department of Physics, University of British Columbia, Vancouver BC; Canada.
174 Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada.
175 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg; Germany.
176 Department of Physics, University of Warwick, Coventry; United Kingdom.
Waseda University, Tokyo; Japan.

Department of Particle Physics, Weizmann Institute of Science, Rehovot; Israel.

Department of Physics, University of Wisconsin, Madison WI; United States of America.

Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany.

Department of Physics, Yale University, New Haven CT; United States of America.

Yerevan Physics Institute, Yerevan; Armenia.

a Also at Department of Physics, University of Malaya, Kuala Lumpur; Malaysia.

b Also at Borough of Manhattan Community College, City University of New York, NY; United States of America.

c Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town; South Africa.

d Also at CERN, Geneva; Switzerland.

e Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.

f Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.

g Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.

h Also at Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); Spain.

i Also at Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; Chile.

j Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah; United Arab Emirates.

k Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.

l Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.

m Also at Department of Physics, California State University, Fresno CA; United States of America.

n Also at Department of Physics, California State University, Sacramento CA; United States of America.

o Also at Department of Physics, King’s College London, London; United Kingdom.

p Also at Department of Physics, Nanjing University, Nanjing; China.

q Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.

r Also at Department of Physics, Stanford University; United States of America.

s Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.

t Also at Department of Physics, University of Michigan, Ann Arbor MI; United States of America.

u Also at Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.

v Also at Giresun University, Faculty of Engineering, Giresun; Turkey.

w Also at Graduate School of Science, Osaka University, Osaka; Japan.

x Also at Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; Romania.

y Also at II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.

z Also at Institut Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.

aa Also at Institut de Física d’Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain.

ab Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.

ae Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen; Netherlands.

ad Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.

ae Also at Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and
Particle Irradiation (MOE), Shandong University, Qingdao; China.

Also at Institute of Particle Physics (IPP); Canada.

Also at Institute of Physics, Academia Sinica, Taipei; Taiwan.

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.

Also at Institute of Theoretical Physics, Ilia State University, Tbilisi; Georgia.

Also at LAL, Université Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, Orsay; France.

Also at Louisiana Tech University, Ruston LA; United States of America.

Also at LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris; France.

Also at Manhattan College, New York NY; United States of America.

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

Also at National Research Nuclear University MEPhI, Moscow; Russia.

Also at Novosibirsk State University, Novosibirsk; Russia.

Also at Ochadai Academic Production, Ochanomizu University, Tokyo; Japan.

Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.

Also at School of Physics, Sun Yat-sen University, Guangzhou; China.

Also at The City College of New York, New York NY; United States of America.

Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China.

Also at Tomsk State University, Tomsk, and Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.

Also at TRIUMF, Vancouver BC; Canada.

Also at Universita di Napoli Parthenope, Napoli; Italy.

* Deceased