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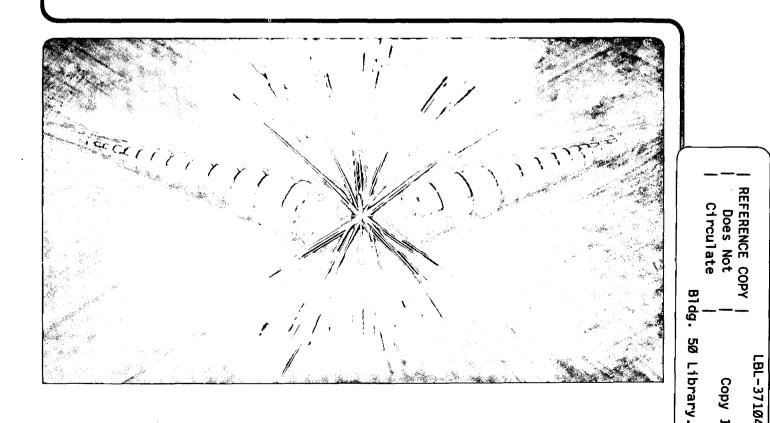
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RESONANT INELASTIC SCATTERING IN LOCALIZED SOLID SYSTEMS BY SOFT X-RAY FLUORESCENCE SPECTROSCOPY*

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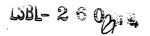
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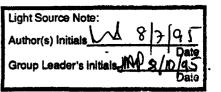
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Resonant Inelastic Scattering in Localized Solid Systems by Soft X-ray Fluorescence Spectroscopy

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An insertion device beamline, BL 8.0 at the ALS, utilizes the radiation from a 5.0 cm period undulator to cover the 100 - 1500 eV photon energy range. In a number of solid systems with localized conduction band states and excitonic states, by tuning the energy of the monochromatic x-rays near the core level absorption edges of cation atoms, a resonant inelastic scattering feature is observed. We have carried out careful studies in various systems. As an example, we present resonant X-ray inelastic scattering spectra from CaF₂, CaSi and CaSi₂ with distinctive localized empty 3d states, to give a general overview of resonant inelastic scattering observations in these systems. When a core electron is excited into a localized state, it stays there long enough eventually to recombine with the core hole, resulting in resonant elastic peaks. For the inelastic process, the energy loss observed indicates an electronic transition of the valence electron to the localized states in the conduction band accompanying the excitation. The valence band originating from the cation Ca is very narrow, yet the strong effects of inelastic scattering at the Ca 2p threshold are unambiguously observed.

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1. INTRODUCTION

Soft x-ray fluorescence (SXF) emission is especially powerful because the shallow core levels have a natural width that is about a factor of 10 less than deeply bound levels¹. Thus, the valence-band spectrum is not unduly broadened by the width of the core level. Furthermore, because the fluorescence photons are not affected by electric and magnetic fields, insulating samples can be studied. Tunable synchrotron radiation has revived interests in x-ray fluorescence spectroscopy². The major barrier to widespread use of soft x-ray fluorescence spectroscopy, the low intensity of available sources, has been overcome with the advent of high brightness third generation synchrotron radiation (SR) sources like the Advanced Light Source (ALS) in Berkeley. Such SR sources are optimized for high spectral brightness in the soft x-ray and vacuum ultra-violet wavelengths.

Resonant emission phenomenon has been observed before, using electron beam excitation³ (by varying the electron beam voltage, C. Bonnelle and co-workers observed resonant features in the emission spectra of rare earth metals and oxides, which coincide in energy with features in the absorption spectra), and laboratory x-ray sources⁴ resulting from the fortuitous location of the Ru L₃ absorption edge between the Pd L α_1 and L α_2 emission lines. The process involves the excitation of core electrons to a localized, empty state, and if the core hole lives long enough, the electron may recombine with the core hole and emit a photon subsequently. A resonant feature is observed only if a localized conduction band state (either an excitonic state or a discrete state hybridized with an extended state) exists. Thus selective excitations to the localized states provide an opportunity to study the localized states in a new angle⁵⁻⁷.

Soft x-ray fluorescence spectroscopy using photon excitation has a number of unique properties. Photons penetrate many atomic layers into the material,. It is a true bulk probe of the electronic structure and is less affected by surface contamination than would occur if one used electron spectroscopy². Since a core hole of welldefined angular momentum is generated by a dipole transition in SXF, the method is site and angular momentum selective. We have studied molecular gasses², buried interfaces⁶, and solid systems⁷ using resonant x-ray emission spectroscopy. As an example of resonant x-ray inelastic scattering, spectra from CaF₂, CaSi, and CaSi₂ with distinctive localized empty 3d states will be presented in this paper. Detailed analyses of these systems are presented elsewhere⁸.

2. EXPERIMENT

An insertion device beamline BL 8.0 at ALS is utilizing radiation from a 5.0 cm period undulator (U5.0) to deliver photons in the 70-1200 eV photon range⁹. The energy resolution, $E/\Delta E$ is better than 10⁴ for the 100-200 eV region. Using refocussing optics, over 10¹² ph/s in a 0.1% bandwidth can be delivered to about 100 µm spot size¹⁰.

One of the four end stations to utilize this undulator radiation is a high resolution soft x-ray emission (SXE) spectrometer, a grazing incidence, Rowland Circle instrument in which the diameter of the Rowland Circle can vary from 5 to 10 m. This spectrometer is designed to match the small phase-space attributes of the high spectral brightness of the beamline optical system. The fluorescence spectrometer has a fixed entrance slit, four spherical diffraction gratings, and a position-sensitive area detector. The gratings are interchangeable from outside the vacuum. The detector is a true photon counting device utilizing micro channel plates and a resistive anode readout. Emission spectra are measured by positioning the detector along the Rowland circle to intercept the spectrum of interest, and utilizing the area detector to record the entire spectrum simultaneously. Scanning is accomplished by a precision X-Y-q table on which the detector chamber is mounted. Both scanning and data acquisitions are automated through an IEEE-488 bus interfaced to a personal computer. Detailed description of the spectrometer is presented elsewhere¹¹. All measurements reported in this paper were made with an 100 mm entrance slit for the spectrometer, 1000 lines/mm, 10 meter radius grating. At the Ca L_{2,3} edges, the resolution of the spectrometer is approximately 1.0 eV. The energy resolution of the monochromator was set to be 0.3 eV during these measurements.

3. RESULTS AND DISCUSSIONS

In Fig. 1, the Ca $L_{2,3}$ fluorescence spectra of CaF₂ for excitation energy varying from 346.0 to 353.3 eV are shown. The numbers labeling the curves are the excitation energies. Sharp elastic (recombination) peaks with varying intensity are obvious in these spectra for excitation energies

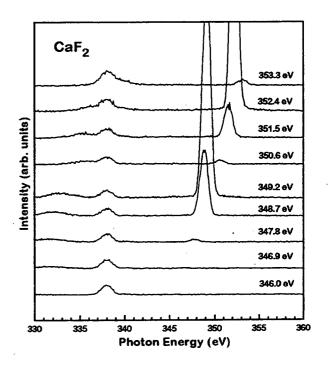


Fig. 1. The Ca $L_{2,3}$ fluorescence spectra of CaF₂ exited by monochromatic photons of various energies near the Ca L edge. The numbers labeling the curves are the energies of the incident photons.

above 347.8 eV. It is apparent from figure 1 that the recombination peaks show strong resonance at both the L₃ $(2p_3/2)$ and L₂ $(2p_1/2)$ edges. The normal fluorescence (Ca L) is the narrow peak present at about 338 eV. This feature persists at all excitation energies and remains unchanged until the onset of the second Ca 2p-3d resonance, when the intensity of the high energy side increases to give it an asymmetric look. Further checking indicates that the assymetry has significant contribution from the second order F K fluorescence excited by the higher-order radiation that is inherent for a SGM monochromator.

The most interesting part of these spectra in figure 1 appears at about 333 eV region. At an excitation energy of 346 eV, the spectrum is clean without the presence of an elastic peak in this spectral region. With an increase of the excitation energy by just one eV, a broad spectral feature at about 333 eV is observed. The intensity of this spectral feature increased as the excitation subsequently being raised. It reaches its maximum value relative to the normal Ca L fluorescence, when the excitation energy is 349.2 eV, coinciding with the excitation of the largest elastic recombination peak.

The Ca $L_{2,3}$ spectra of CaSi are shown in Fig. 2. Each spectrum is labeled with the photon energy used for its excitation. The normal fluorescence spectra, generated by the radiative recombination of Ca valence band states with the core hole, produce the spectral features that are centered near 343 eV and observed with offresonance excitation. The energy loss features at about 335 eV region reaches its maximum value relative to the normal Ca L fluorescence, when

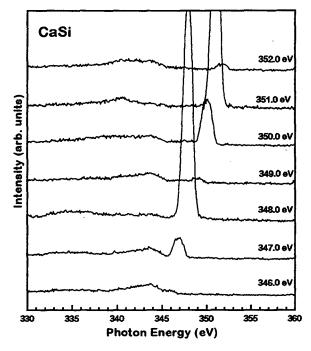


Fig. 2. The Ca $L_{2,3}$ fluorescence spectra of CaSi with the excitation energies indicated.

the excitation energy is 348.0 eV, corresponding to the excitation of the maximum elastic recombination peak.

The Ca $L_{2,3}$ fluorescence spectra of CaSi₂ are shown in Fig. 3. The normal Ca L fluorescence is observed at about 343 eV. Note that no prominent resonant inelastic peaks are resolved for CaSi₂, though a very weak displaced peak may be present centered at about 336 eV in the spectrum excited on the 348 eV resonance.

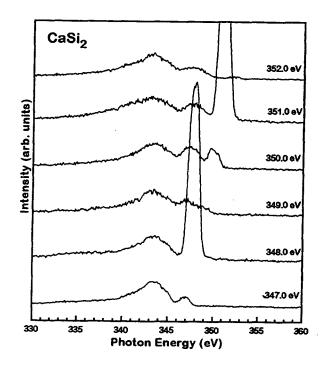


Fig. 3. The Ca $L_{2,3}$ fluorescence spectra of CaSi₂ with the excitation energies indicated.

The resonant process can be understood as following: when monochromatic photons are tuned to excite Ca 2p core electrons into 3d states, the electrons stay in the d states due to its semidiscrete nature; and eventually they recombine with the core hole and give such large resonances. Extended conduction band states in de localized solid systems do not usually appear in fluorescence spectra, because electrons in these states quickly diffuse away from the core hole within the core lifetime before they radiate. The 3d states in Ca compounds are an intermediate case between atomic and well-extended states in

solids. The conduction band states are essentially a hybridization between Ca 3d states and extended s-p states. Thus, selective excitation of Ca 2p electrons using photons with energies corresponding to the Ca 2p-3d transition may also excite the core electron to the extended states, in which case the valence band electrons will annihilate the core hole, resulting in normal valence band fluorescence. The intensity ratio between the resonance from the 3d states and valence band fluorescence may provide an indication of the degree of hybridization between the Ca 3d states and the extended s-p conduction states. From Figs. 1, 2 & 3 the intensity ratios of the elastic peak to the valence band peaks for CaF₂ and CaSi are about 20, whereas the corresponding intensity ratio for CaSi₂ is about 7.

The appearance of inelastic scattering feature in Ca L spectra seems to be closely related to the localization factor of Ca 3d states. In a highly ionic compound like CaF2, the valence band is derived largely from fluorine 2p atomic states. As a result, CaF₂ has very narrow valence band. The energy-loss feature is also present for the CaSi spectra, but not for CaSi2. The Ca 3d states are more localized in CaSi than in CaSi₂, as evidenced by the ratio of recombination peak and main emission feature from the valence band. The extra silicon concentration in CaSi2 obviously increased the interaction of the Ca d orbitals with de localized Si s- and p- orbitals. The resonant elastic recombination is quite strong in CaSi, as in CaF₂.

In conclusion, we have measured the Ca $L_{2,3}$ fluorescence of CaF₂, CaSi and CaSi₂. The localized excitation of Ca 2p - 3d produces resonant recombination peaks. Prominent inelastic scattering features were observed below the valence emission. A coherent Raman-like inelastic scattering process is believed to be responsible for the observation of such features. The absence of the inelastic peak in CaSi₂ is consistent with both the reduced magnitudes of the elastic resonance and the stronger Ca d-, Si sand p- bonding expected in the di-silicide.

4. ACKNOWLEDGMENTS

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